



# Pulse shaping techniques with nested wire arrays (and applications of novel wire arrays)

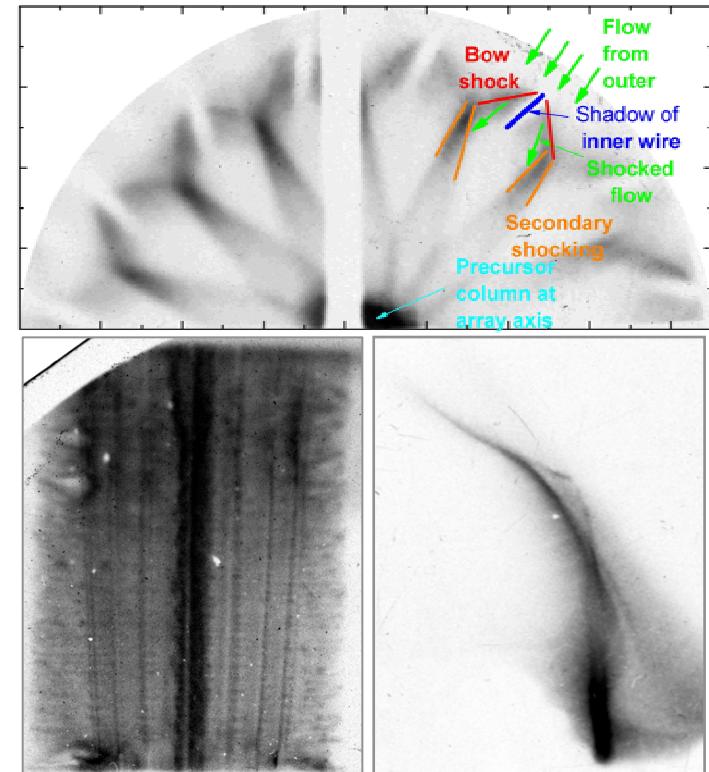
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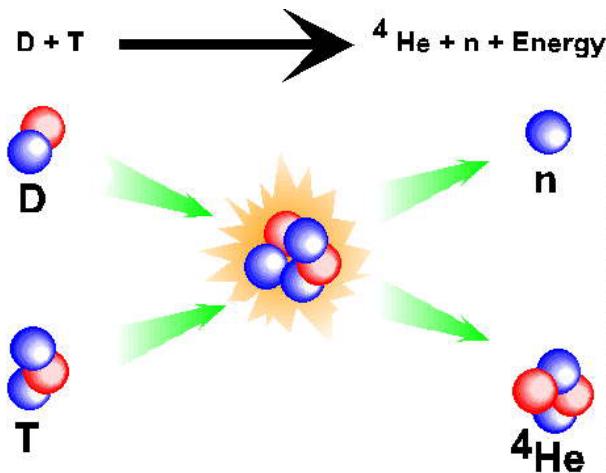
Research at Imperial College is sponsored by the NNSA under DOE Cooperative Agreement DE-F03-02NA00057.

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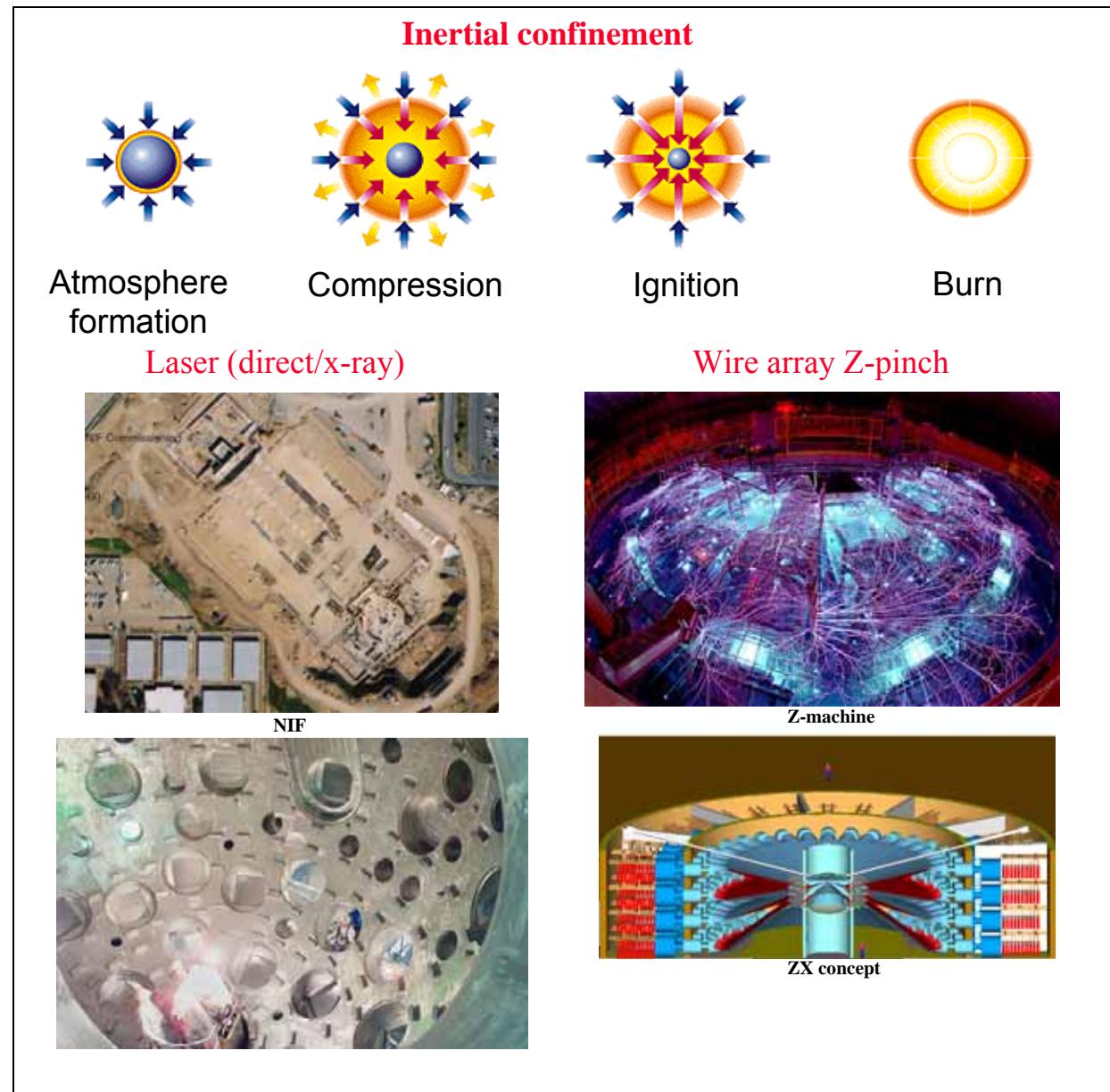
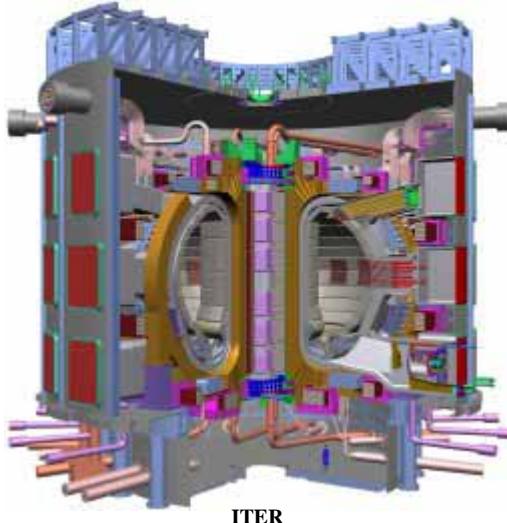


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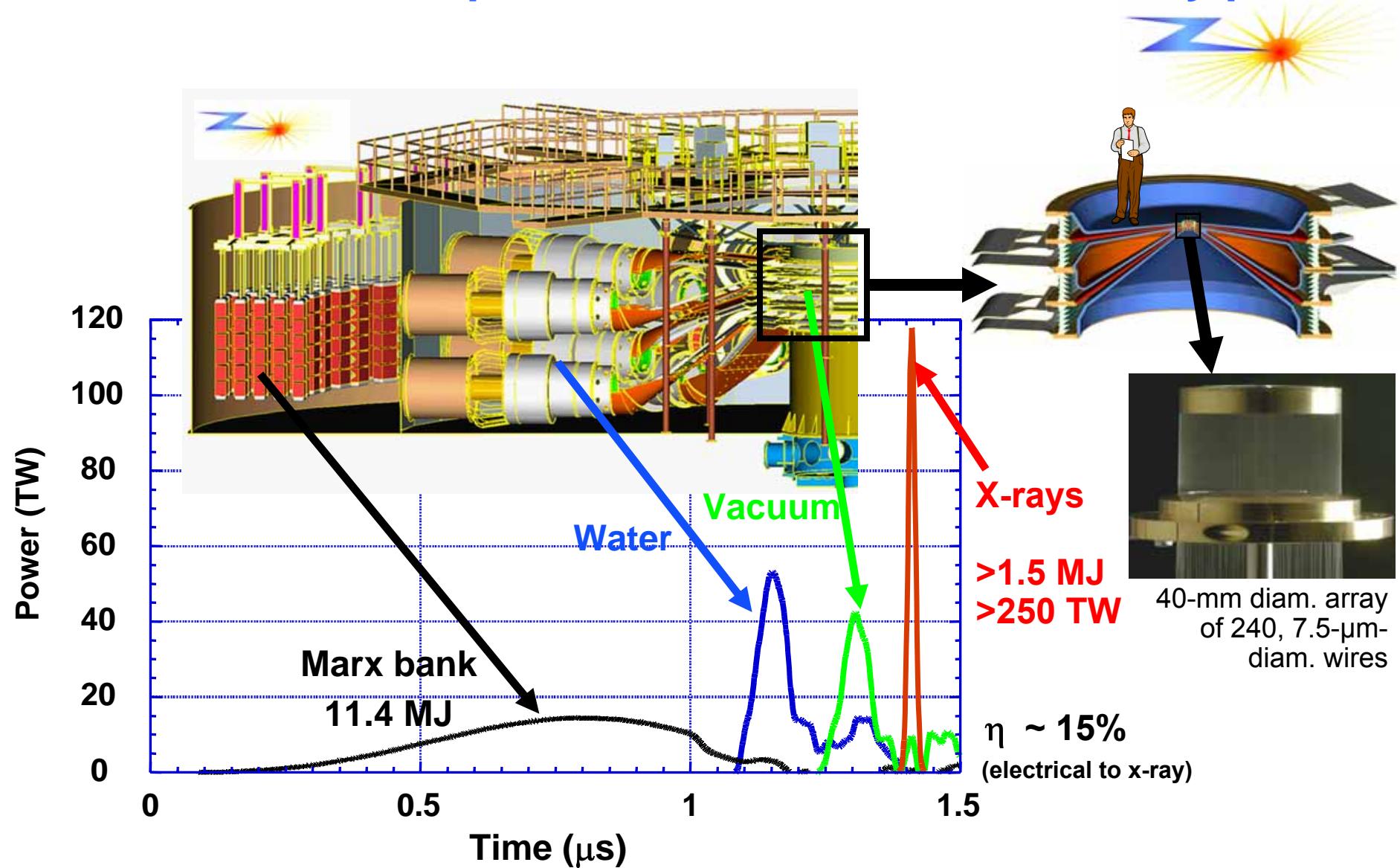
# Fusion needs a hot, confined plasma



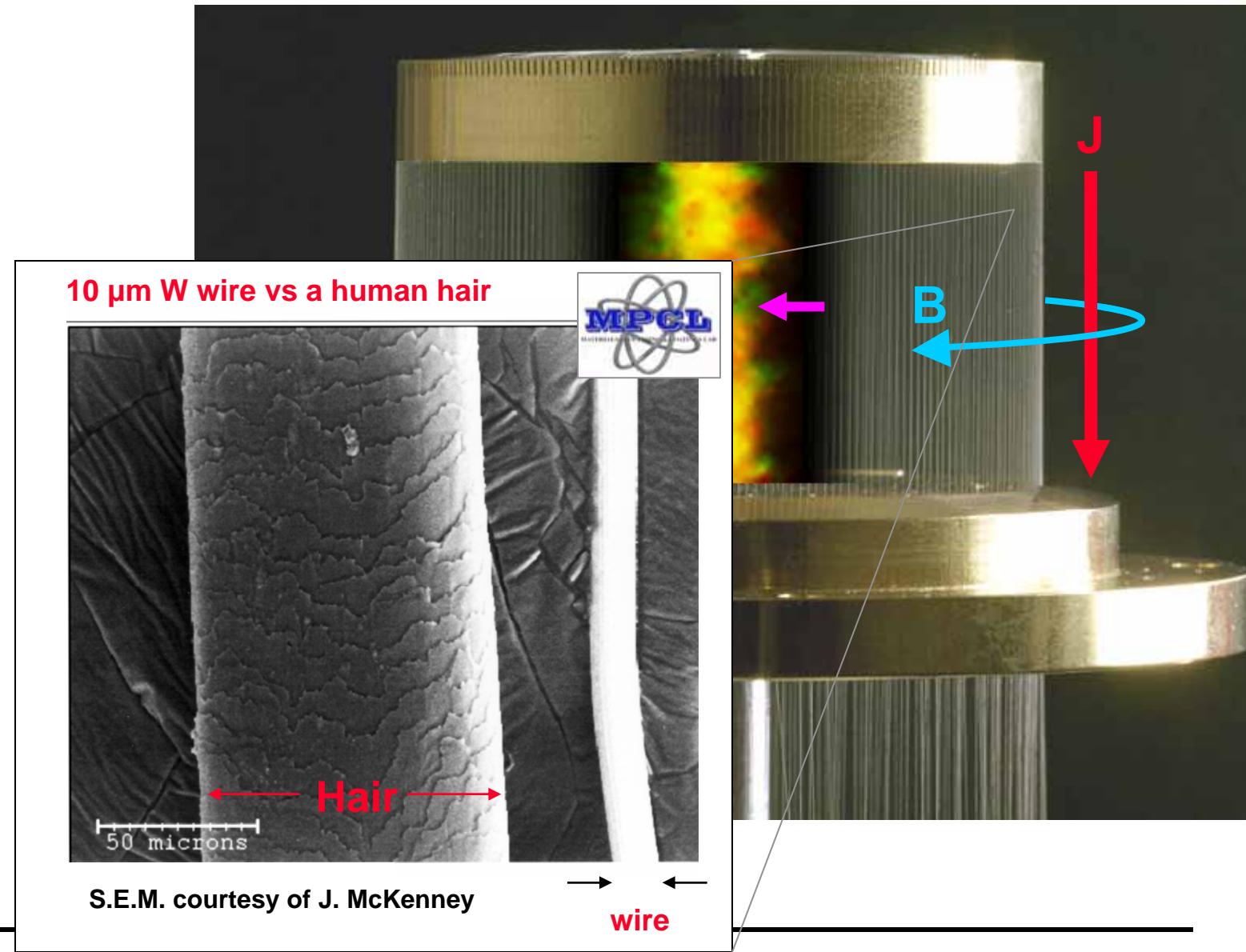
**Magnetic confinement:**  
Tokomak (and classical z-pinch)



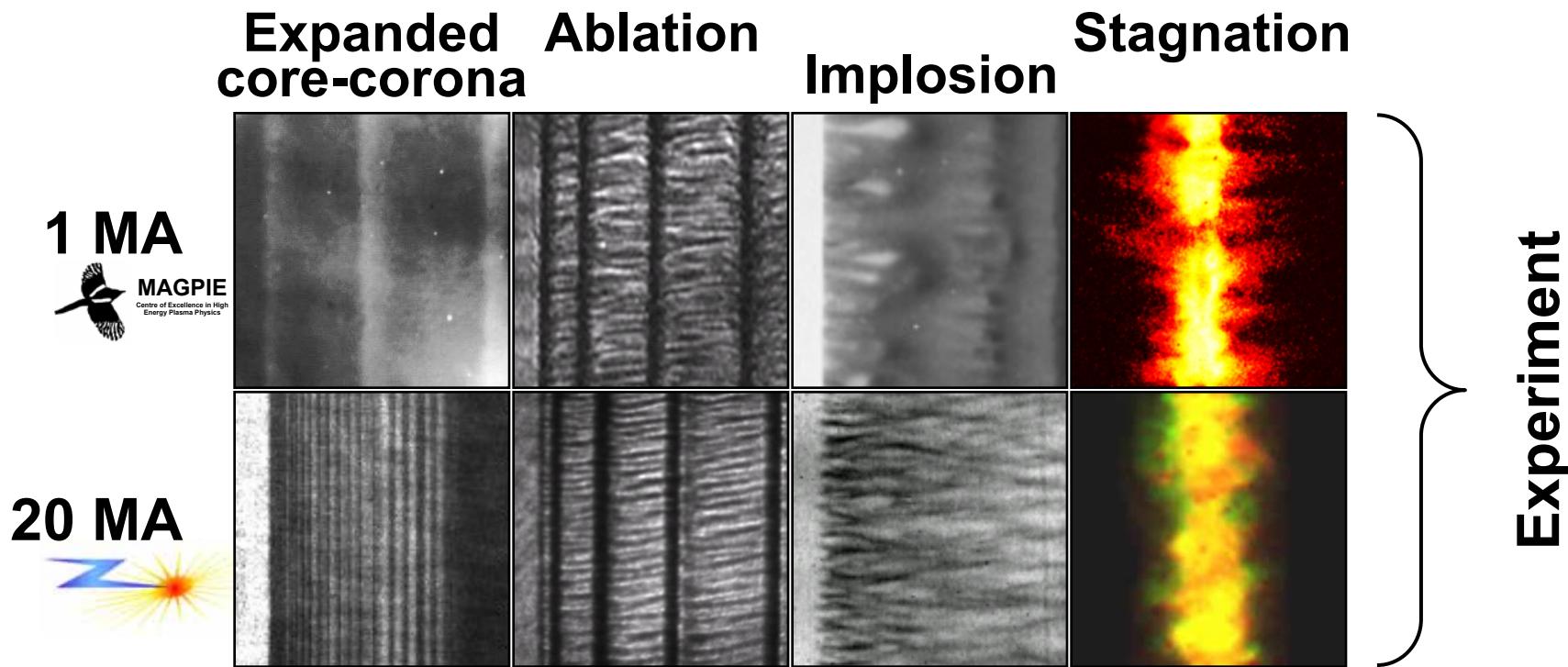
# Sandia's Z machine produces world-record soft x-ray powers



# J x B pinches wire array into a dense, radiating plasma



# Wire array dynamics are 3D

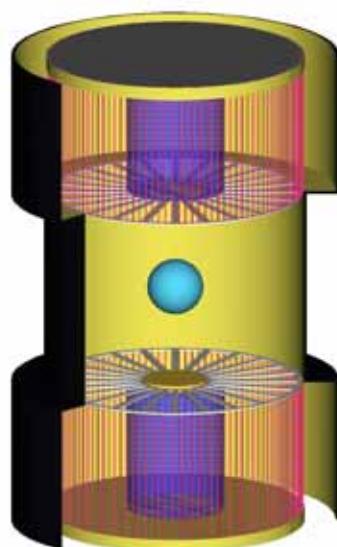
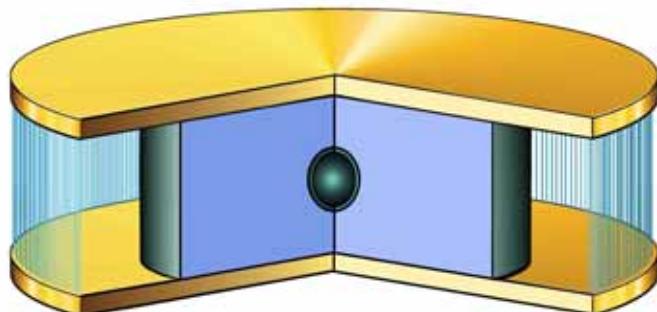


Images courtesy of: Brent Jones, Dan Sinars, Dave Bliss, Gennady Sarkisov, Christine Coverdale, Chris Garasi et al. (Sandia), Sergey Lebedev, Jerry Chittenden Malcolm Haines et al. (Imperial), Bob Clark, Jack Davis et al. (NRL), Darrell Peterson (LANL)

# High yield fusion is one of the major motivators for the z-pinch research at Sandia

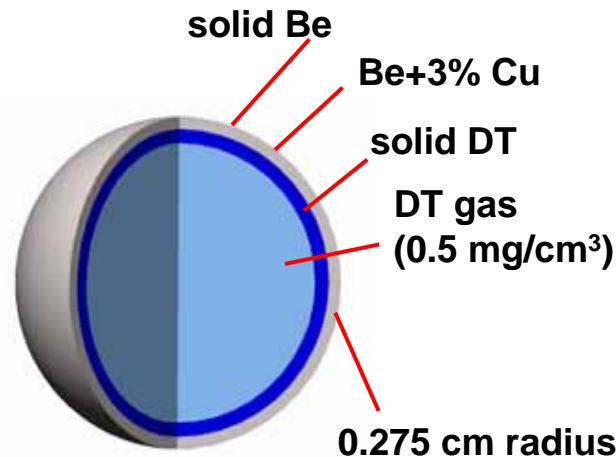


Dynamic Hohlraum



Z-pinch Driven Hohlraum

Estimates predict 2 drivers, each 60MA required for ignition  
Controlled radiation pulse-shape vital!



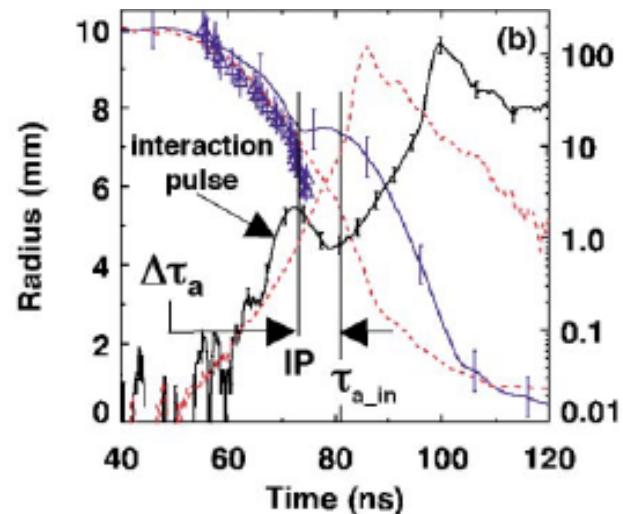
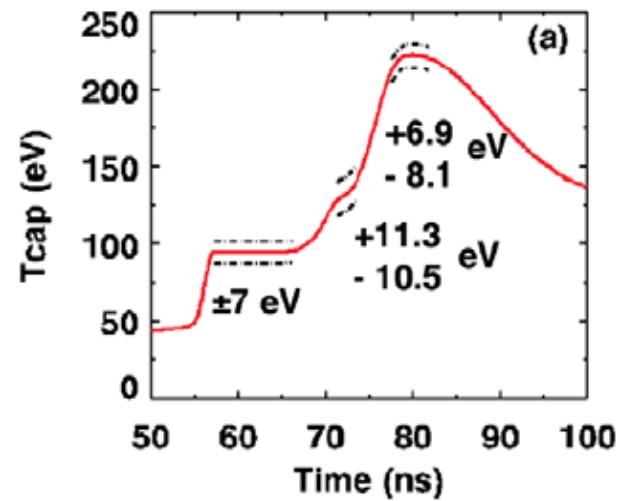
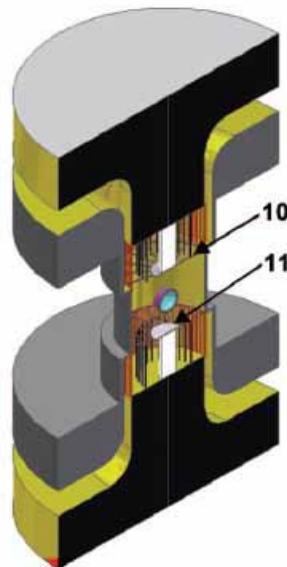
## Important considerations:

- Intensity of radiation incident
- Radiation symmetry
- Radiation pulse shape
- Mechanics of power station

Peak drive temperature	350 eV
In-flight aspect ratio	48
Implosion velocity	$3.3 \times 10^7$ cm/s
Convergence ratio	27
DT KE @ ignition	50%
Peak density	444 g/cm³
Total $\rho r$	$2.14 \text{ g/cm}^2$
Driver energy	12 MJ
Absorbed energy	2.3 MJ
Yield	527 MJ
Burnup fraction	34%



# Pulse shaping is vital to z-pinch ICF concepts



- Three or more controlled x-ray pulses are required in order to heat a fusion capsule
- One suitable pre-pulse is observed as the imploding outer array of two nested arrays interacts with the inner array, however detailed physical mechanism of the interaction pulse is not fully understood
- Necessary to broaden main pulse (and interaction pulse)
  - nested arrays on Z are *too good* at temporal compression
- Increasing energy density in the main pulse useful for ICF

Images reproduced from  
M.E. Cuneo et al. Phys Plas 13 056318, 2006

# This talk will use MAGPIE data to help understand / interpret / predicting Z data

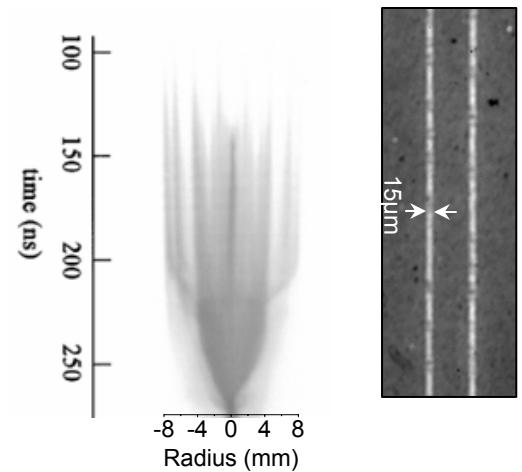
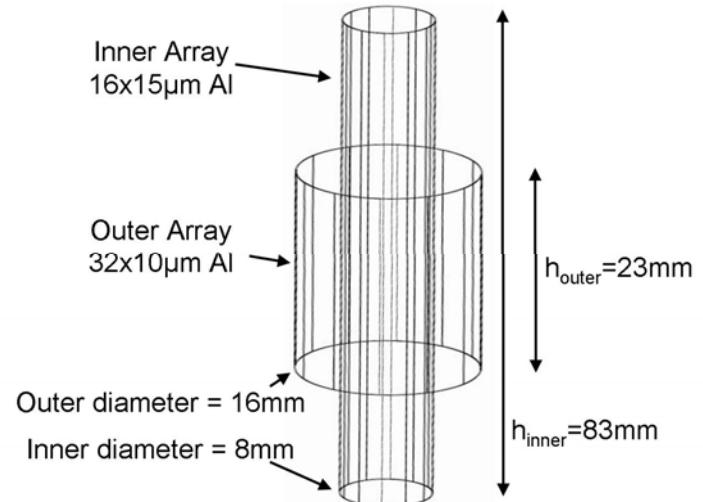
## Aims:

- Understand our present pulse shaping capability
  - New understanding of mechanism possibly responsible for interaction pulse
- New tools in the pulse-shaping toolbox
  - Look at what conical nested can bring to pulse shaping
- Novel arrays
  - Discuss experiments performed on Saturn to study elimination of cathode effects
  - Radial arrays will be discussed both for ICF and lab-astrophysics
  - Conical wire arrays will be discussed with regard to lab-astrophysics

1MA experiments can be used to understand 20MA experiments if careful consideration given to getting appropriate setup and being aware of differences

## Nested wire arrays on MAGPIE use high inductance inner to suppress current through the inner array to be similar to Z

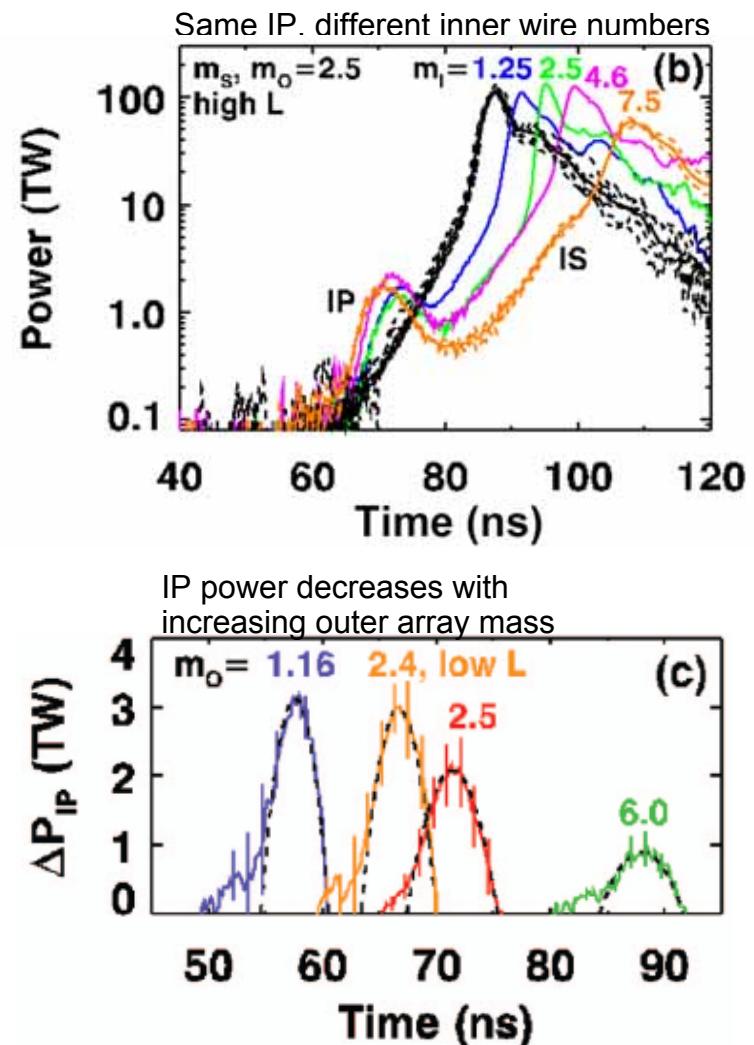
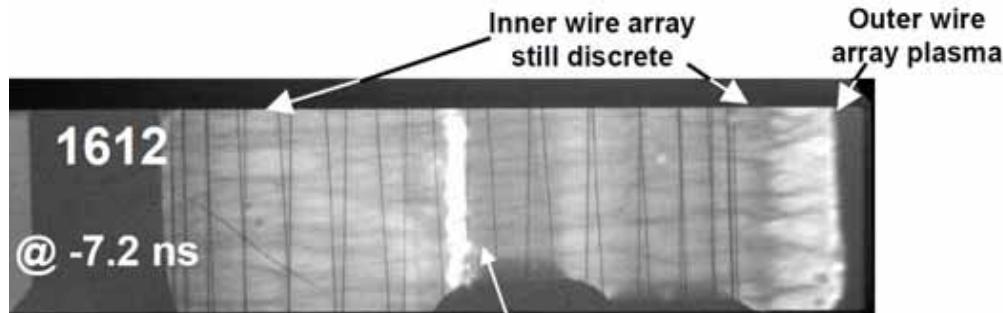
- High wire number in outer at 20MA leads to  
Inductive contrast:  $L_{\text{outer}} \ll L_{\text{inner}}$   
e.g. Cuneo et al PRL 94, 225003 (2005)
- High wire number not possible at  $\sim 1\text{MA}$
- Array design can give same inductive contrast  
(by lengthening inner)  
Lebedev et al. PRL 84, 1708 (2000)
- Negligible inner current confirmed by
  - Radial optical streak
  - X-pinch radiography
  - B-dot probes
- Present experiments use
  - Outer array 16-32 x  $10\mu\text{m}$  Al 5056 at 16mm
  - Inner array 16 wire Al, W or CH at 8mm





## Interaction pulse on Z is critical for pulse shaping, but remains a puzzle

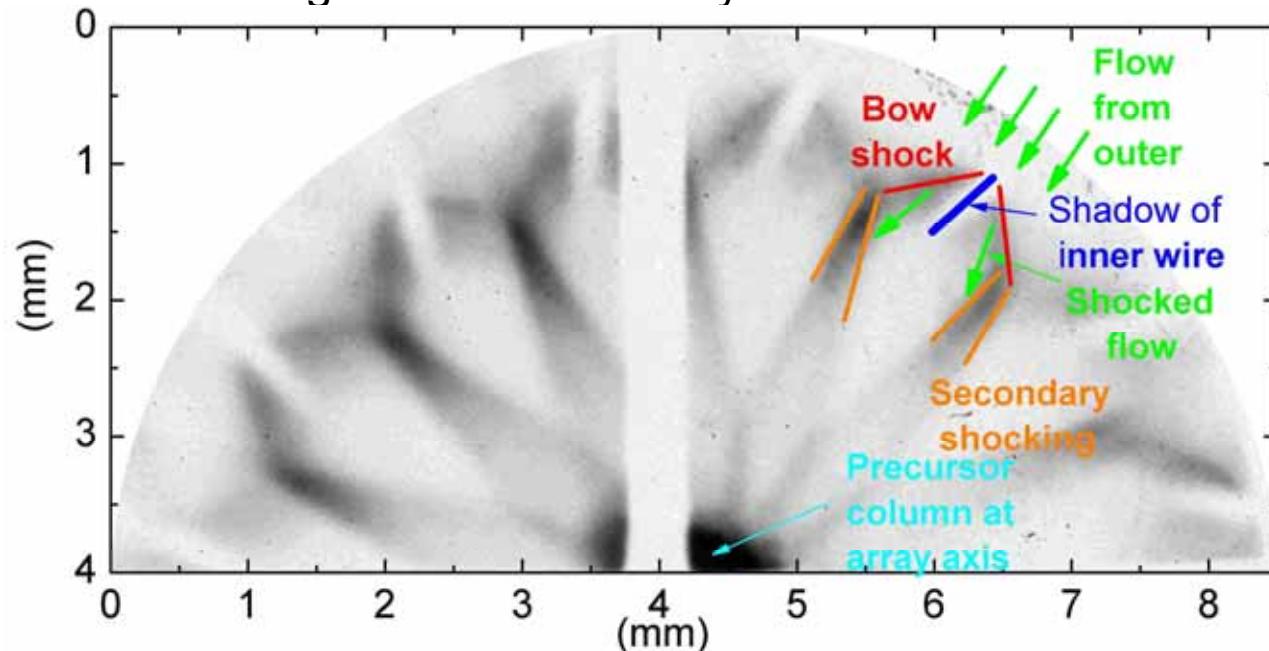
- Interaction energy measured is less than that predicted from hydrodynamic collision of outer and inner shells
- Nested arrays on Z now recognized to operate in a transparent mode
- Partial transparency cannot explain same power for different inner wire numbers
- Alternative models (e.g. ohmic heating, possibly by flux compression), do not recreate dependency on outer array mass, or explain small cores prior to interaction



Plots reproduced from  
M.E. Cuneo *et al.* Phys Plas 13, 056318, 2006

## Ablation streams from outer array are supersonic and will shock on the inner array

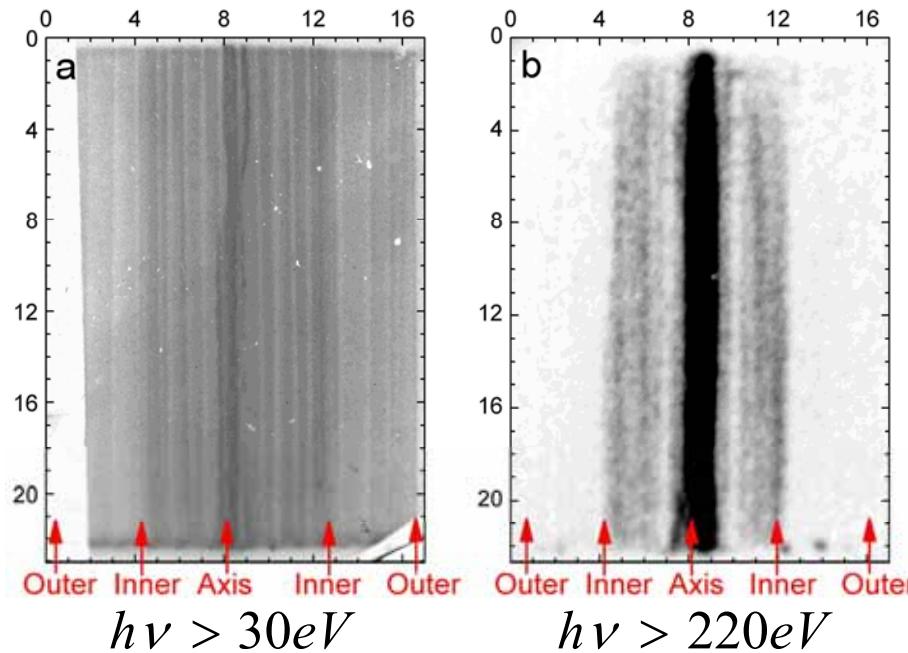
- Precursor plasma flows from outer array are supersonic at position of inner array:
  - MAGPIE (from spectra):  $T_e \sim 40\text{eV}$ ,  $Z \sim 6$      $c_s \sim 3\text{cm}/\mu\text{s}$      $M \sim 5$
  - Z (from MHD):  $T_e \sim 25\text{eV}$ ,  $Z \sim 11$      $c_s \sim 1.3\text{cm}/\mu\text{s}$      $M > 11$
- At reaching the inner array the precursor flow will be shocked
  - end-on XUV image inside inner array on MAGPIE



- Define angle  $\beta$  as angle between initial precursor flow and shock

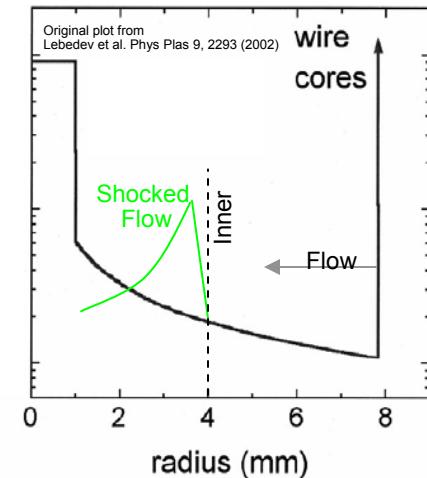
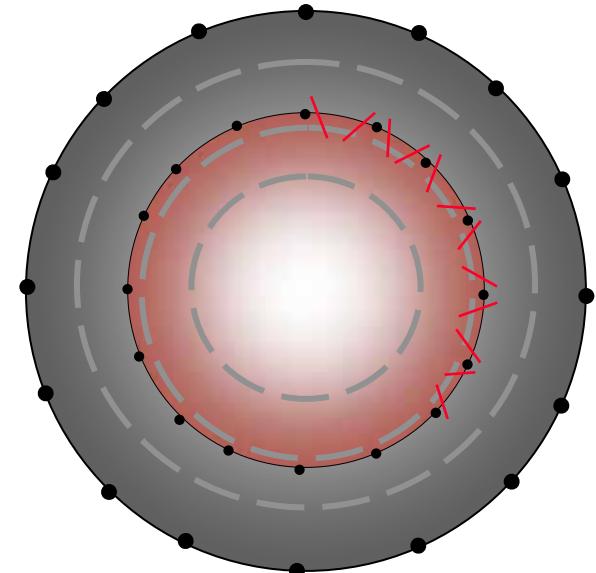
## Shock will perturb plasma conditions in streams as they pass inner

- Perpendicular component of stream velocity will be reduced across shock
$$\frac{v_{\perp sh}}{v_{\perp abl}} = \frac{1}{\eta} = \frac{M_{\perp}^2(\gamma - 1) + 2}{M_{\perp}^2(\gamma + 1)} \sim 0.14$$
- Density will be increased by the compression ratio  $\eta$
- Temperature of streams will also be increased
- Temperature and/or density jumps inferred from side-on emission imaging during ablation process
  - Definite change in plasma conditions near inner array, despite ‘transparency’ :



## Perturbing the pre-fill will affect the snowplow

- For single array snowplow of pre-fill by implosion results in emission
- Power radiated by snowplow emission is
  - $P_{SP} \propto \rho(r,t) (v_{piston} - v_{prefill})^3$
- Comparing nested with single,  $\rho$  and  $v_{prefill}$  both altered by jump conditions
- Modifications act to enhance snowplow emission.
- Can adapt a snowplow model to incorporate these jumps



# Snowplow model for perturbed system predicts enhanced emission above that of a single array

Variable	MAGPIE
$v_{\text{abl}}(\text{cm}/\mu\text{s})$ [1, 2]	15 (Ablation velocity)
$c_s(\text{cm}/\mu\text{s})$ [2–4]	3 (Sound speed)
$\beta$ (end-on image)	39° (Shock angle)
$\gamma$ [5]	1.1 (Adiabatic index)
$M_{\perp} = \frac{v_a \sin(\beta)}{c_s}$	3.4 (Mach numb perp)
$\eta = \frac{v_{\text{abl}}}{v_{\text{sh}}}$	7.7 (Compression)

[1] S. V. Lebedev et al., *Plas. Phys. Contr. Fus.* 47, A91 (2005).

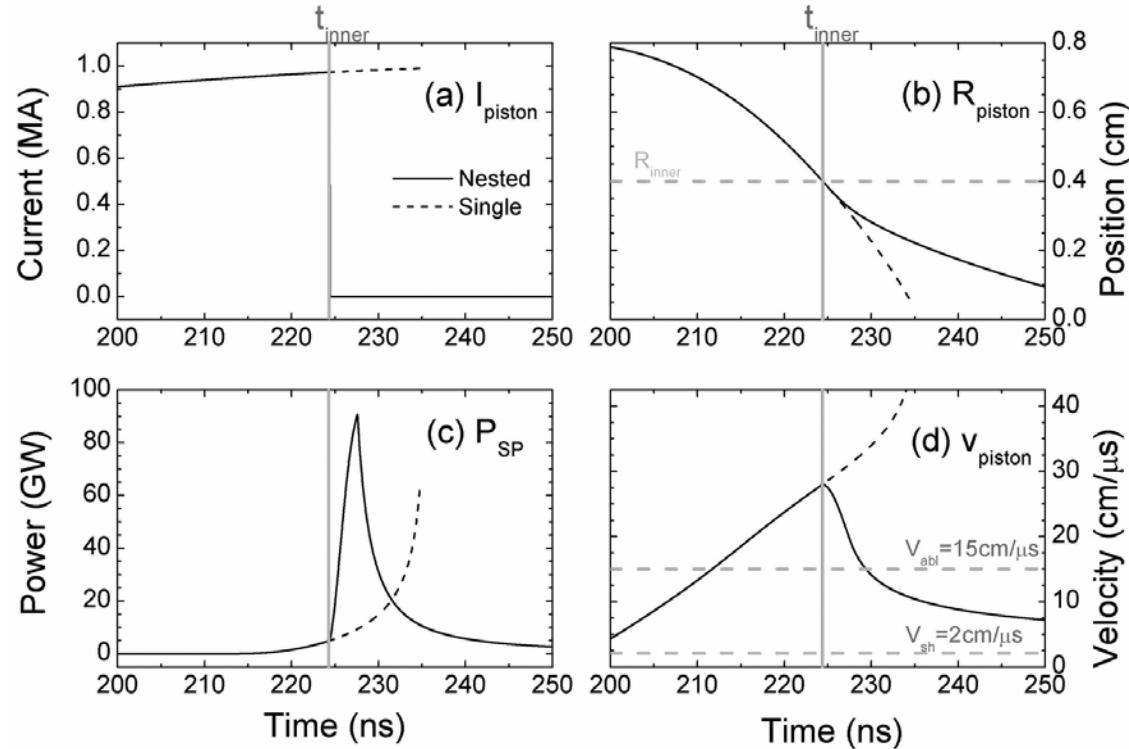
[3] S. V. Lebedev et al., *Laser Particle Beams* 19, 355 (2001).

[4] J. P. Chittenden et al., *Phys. Plasmas* 8, 675 (2001).

[5] R. P. Drake, *High Energy Density Physics* (Springer, 2006).

## Model setup

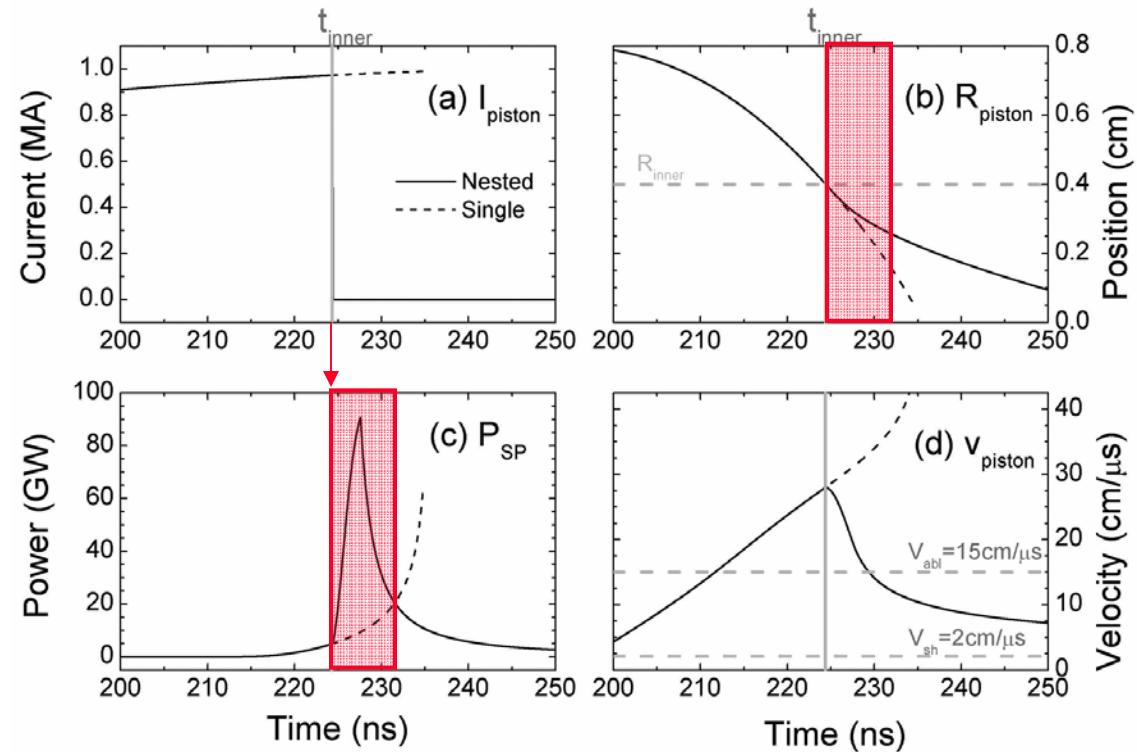
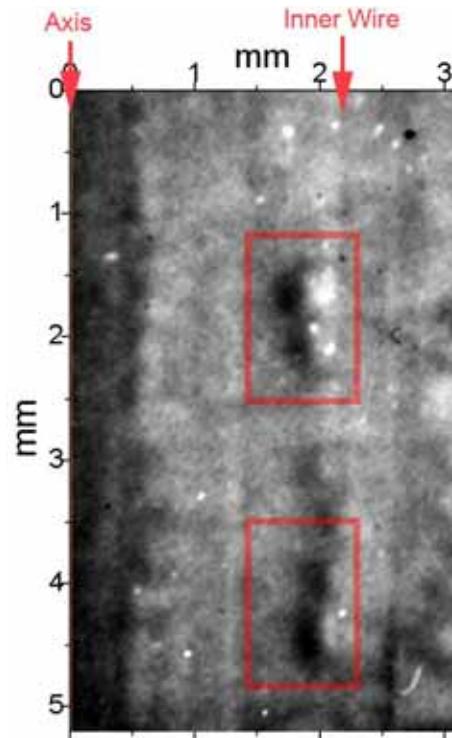
- Ablation model for fill
- Jump conditions at inner
- Snowplow model for implosion trajectory



## Results of model

- Snowplow model shows excess emission despite current being switched out of piston prior to experiencing perturbed density
- Excess emission is AFTER piston passes inner wires (in shocked region)
- Experiments show that piston slows below ablation velocity, despite 100% transparency
- Averaging emission over total MAGPIE array smoothes out interaction pulse due to azimuthal and axial non-uniformities

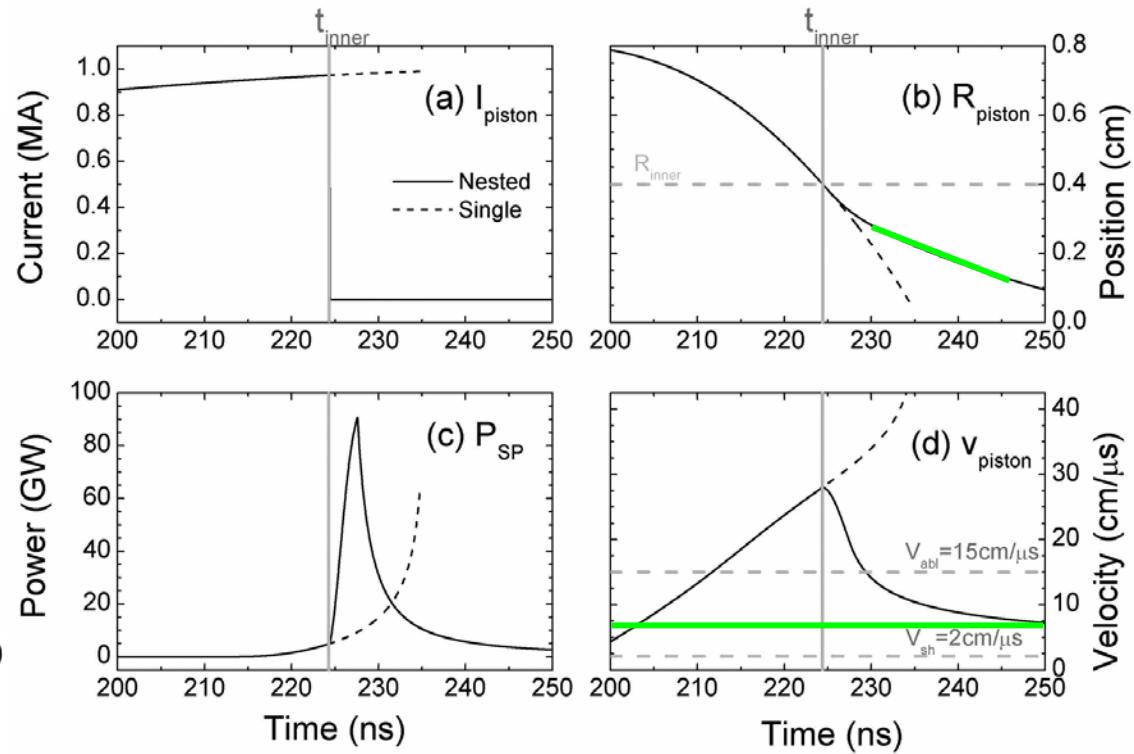
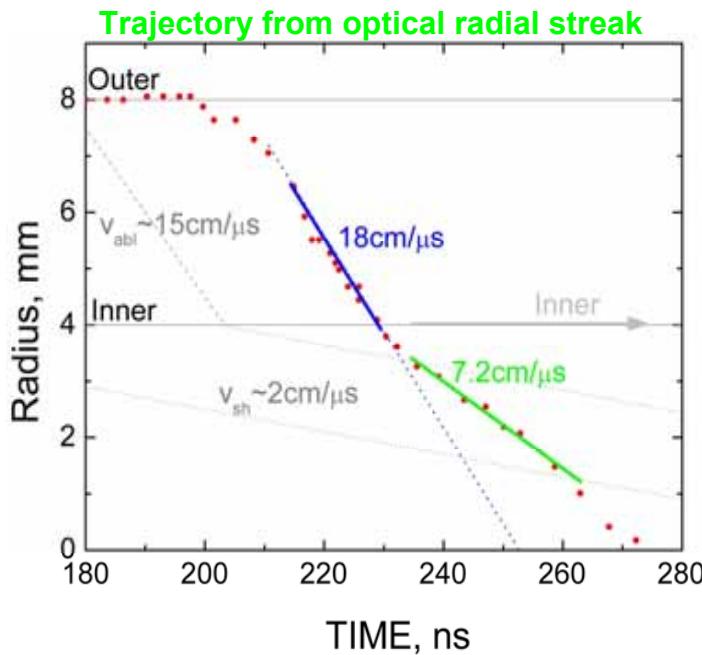
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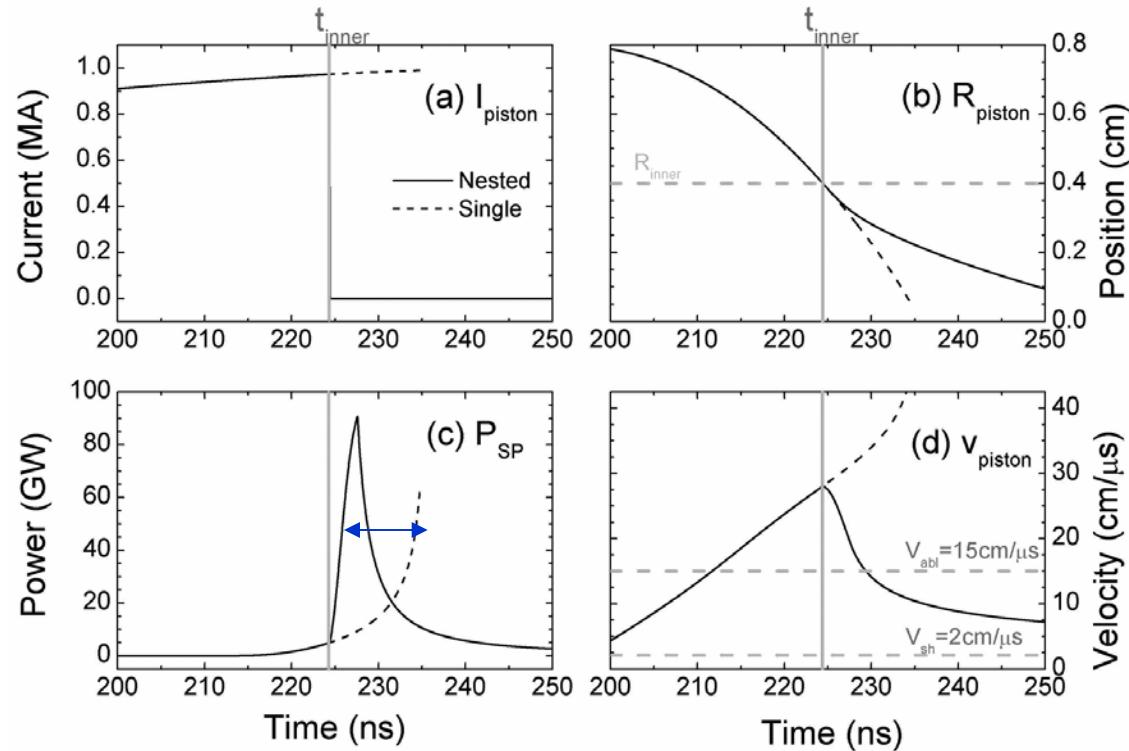
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$\gamma$ [5]	1.1
$M_{\perp} = \frac{v_a \sin(\beta)}{c_s}$	3.4
$\eta = \frac{v_{ahl}}{v_{ah}}$	7.7

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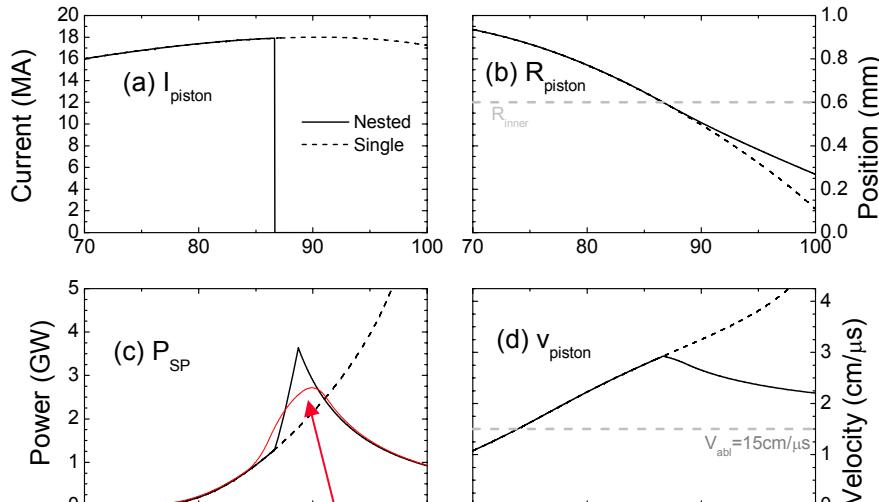
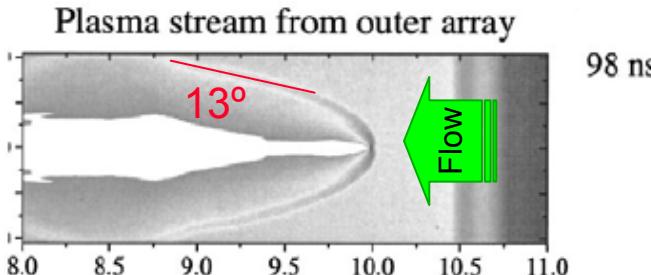
# Applying similar model to a (more uniform) Z implosion allows a comparison of powers



Variable	MAGPIE Z-6mg	
$v_{\text{abl}}(\text{cm}/\mu\text{s})$ [1, 2]	15	15
$c_s(\text{cm}/\mu\text{s})$ [2–4]	3	1.3
$\beta$ (end-on image)	39°	→ 5°
$\gamma$ [5]	1.1	1.1
$M_{\perp} = \frac{v_a \sin(\beta)}{c_s}$	3.4	1.25
$\eta = \frac{v_{\text{abl}}}{v_{\text{sh}}}$	7.7	1.5

Shock angle smaller in sims of Z than measured on MAGPIE

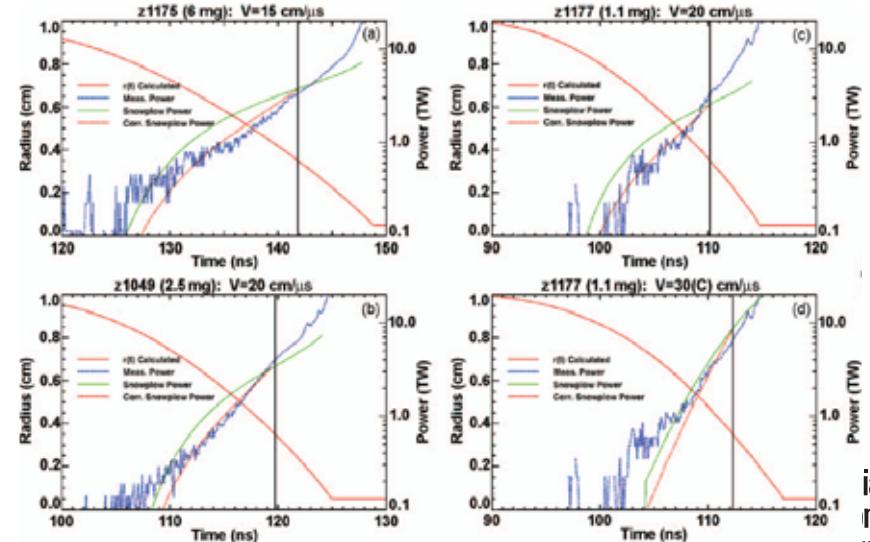
Chittenden *et al.* Phys Plas 8, 675 (2001)



Temporally spread to account for experimentally measured width of piston

Snowplow for 6mg outer on Z gives good fit for  $v_{\text{abl}}$

Sinars *et al.* Phys Plas 13, 042704 (2006)



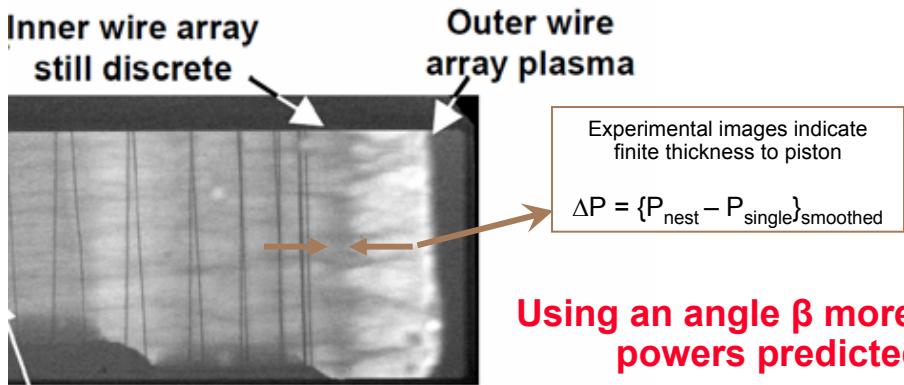
# Snowplow model recreates interaction pulse for Z outer array mass scan



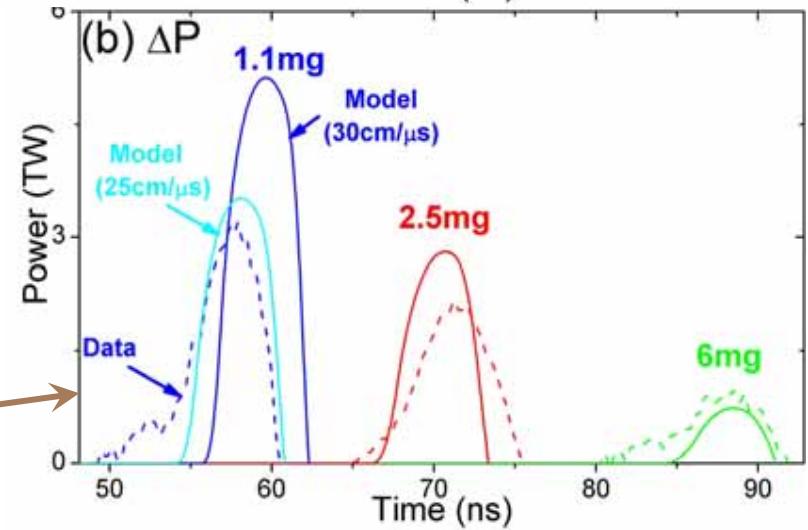
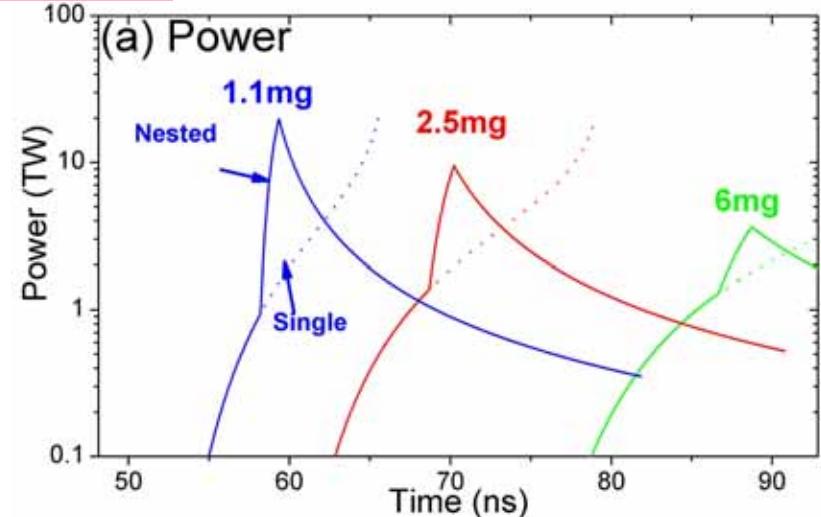
Variable	MAGPIE	Z-6mg	Z-2.5mg	Z-1.1mg	Z-1.1mg
$v_{abl}(cm/\mu s)$ [1, 2]	15	15	20	30	25
$c_s(cm/\mu s)$ [2-4]	3	1.3	1.3	1.3	1.3
$\beta$ (end-on image)	39°	5°	5°	5°	5°
$\gamma$ [5]	1.1	1.1	1.1	1.1	1.1
$M_{\perp} = \frac{v_a \sin(\beta)}{c_s}$	3.4	1.25	1.66	2.5	2.0
$\eta = \frac{v_{ahl}}{v_{ah}}$	7.7	1.5	2.5	5.0	3.7
$E_{Sp}(kJ)$	-	3.0	12.4	22.5	15.3
$E_{Exp}(kJ)$ [6]	-	5.2	12.1	15.0	15.0
$P_{Sp}(TW)$	-	0.7	2.8	5.1	3.5
$P_{Exp}(TW)$ [6]	-	1.2	2.2	3.2	3.2

Variable ablation velocity  
Sinars et al. Phys Plas 13, 042704 (2006)

- [1] S. V. Lebedev et al., Plas. Phys. Contr. Fus. 47, A91 (2005).
- [2] D. B. Sinars et al., Phys. Plasmas 13, 042704 (2006).
- [3] S. V. Lebedev et al., Laser Particle Beams 19, 355 (2001).
- [4] J. P. Chittenden et al., Phys. Plasmas 8, 675 (2001).
- [5] R. P. Drake, High Energy Density Physics (Springer, 2006).
- [6] M. E. Cuneo et al., Phys. Plasmas 13, 056318 (2006).

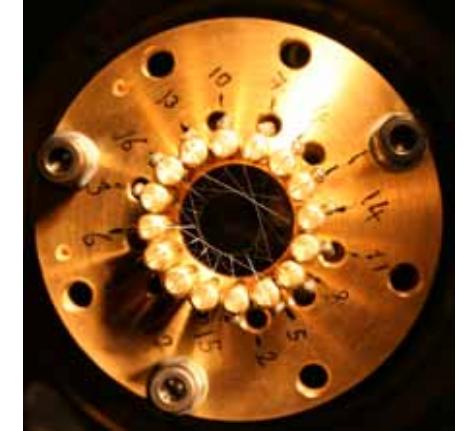
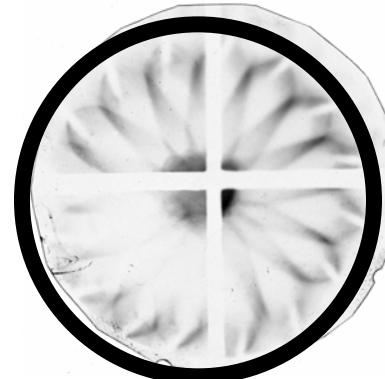


Using an angle  $\beta$  more like MAGPIE or MHD would increase powers predicted, but reduce mass dependence



## Upcoming experiments will aim to better diagnose shocks on MAGPIE, for both W & Al

- Normal nested array hardware prevents the end-on viewing of the streams prior to meeting the inner array
- We do not need to see the full evolution of the nested wire array to image shock structures and diagnose jump conditions
- New array configuration aims to allow imaging of streams before and after passing the inner array
- Potentially evaluate shock jump conditions using end-on diagnostics, and investigate differences between W & Al

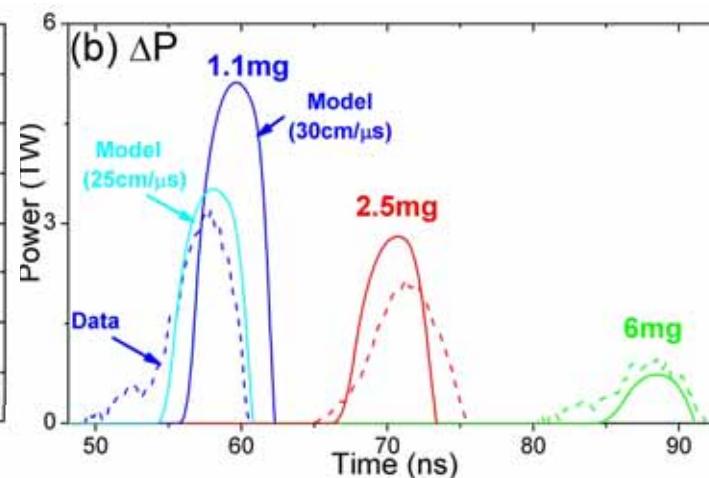
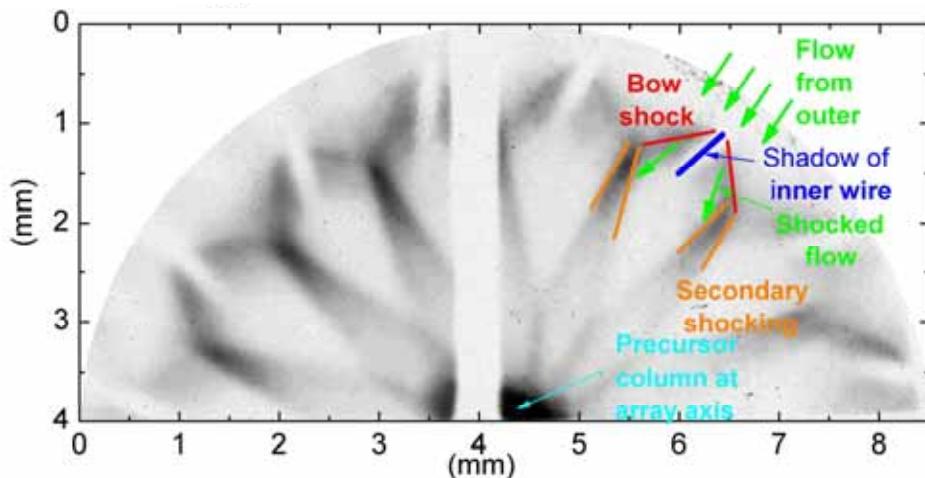


Hardware design (*Spider Array*) and photos courtesy of Dr G.N. Hall (Imperial)



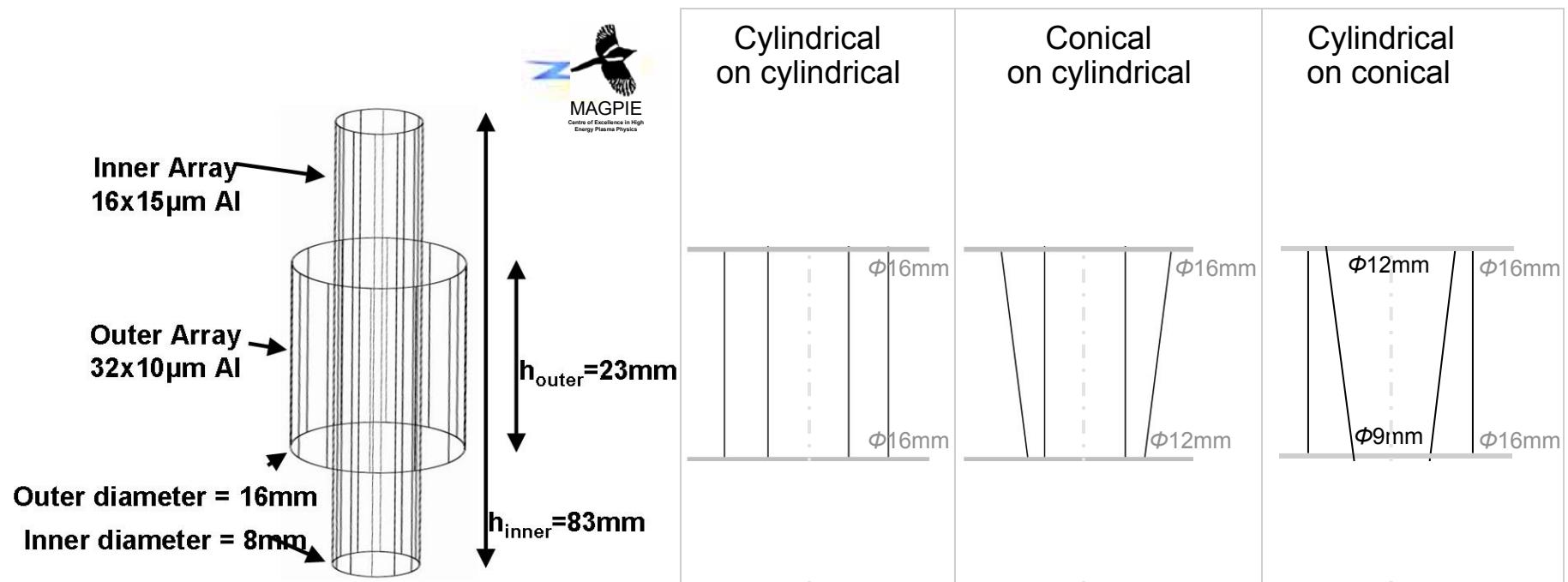
## Part 1 Summary : Perturbed snowplow can lead to observed Interaction Pulse

- Inner array shocks precursor plasma streams on MAGPIE
- Shock will alter  $\rho$ ,  $T$ ,  $v$  of the streams
- This jump is likely to alter the snowplow emission as the outer array implodes
- For Z conditions this change in snowplow radiation can be comparable to the observed interaction pulse
- Able to recreate correct outer mass dependency
- Future plans include analytics, simulations and experiments to better determine correct angle  $\beta$  for W arrays on Z



## Part 2: Conical nested arrays may be able to enhance control of x-ray pulse

- One suggested technique to lengthen the main pulse from nested arrays is to seed a zipper using conical arrays  
c.f. zipper observed in gas puff experiments
- Two modified setups tested on both MAGPIE and Z



# Use of a conical outer array will alter mass ablation rate

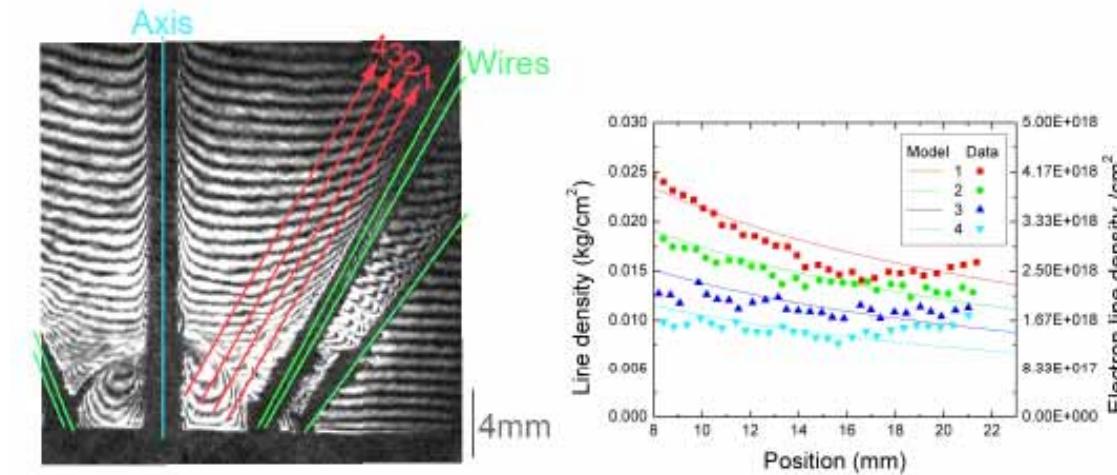
- With conical outer array the global field strength is varied along the length of the wire

$$B(z, t) = \frac{\mu_0 I(t)}{2\pi R(z)} = \frac{\mu_0 I(t)}{2\pi(R_0 + z \tan(\alpha))}$$

- Variation in global field alters mass ablation (due to fixed  $v_{abl}$ )

$$\dot{m} = \frac{\mu_0 I(t)}{2\pi V_{abl}(R_0 + z \tan(\alpha))}$$

- Mass density profiles in single conical arrays are consistent ( $v_{abl} \sim 15 \text{ cm}/\mu\text{s}$ )

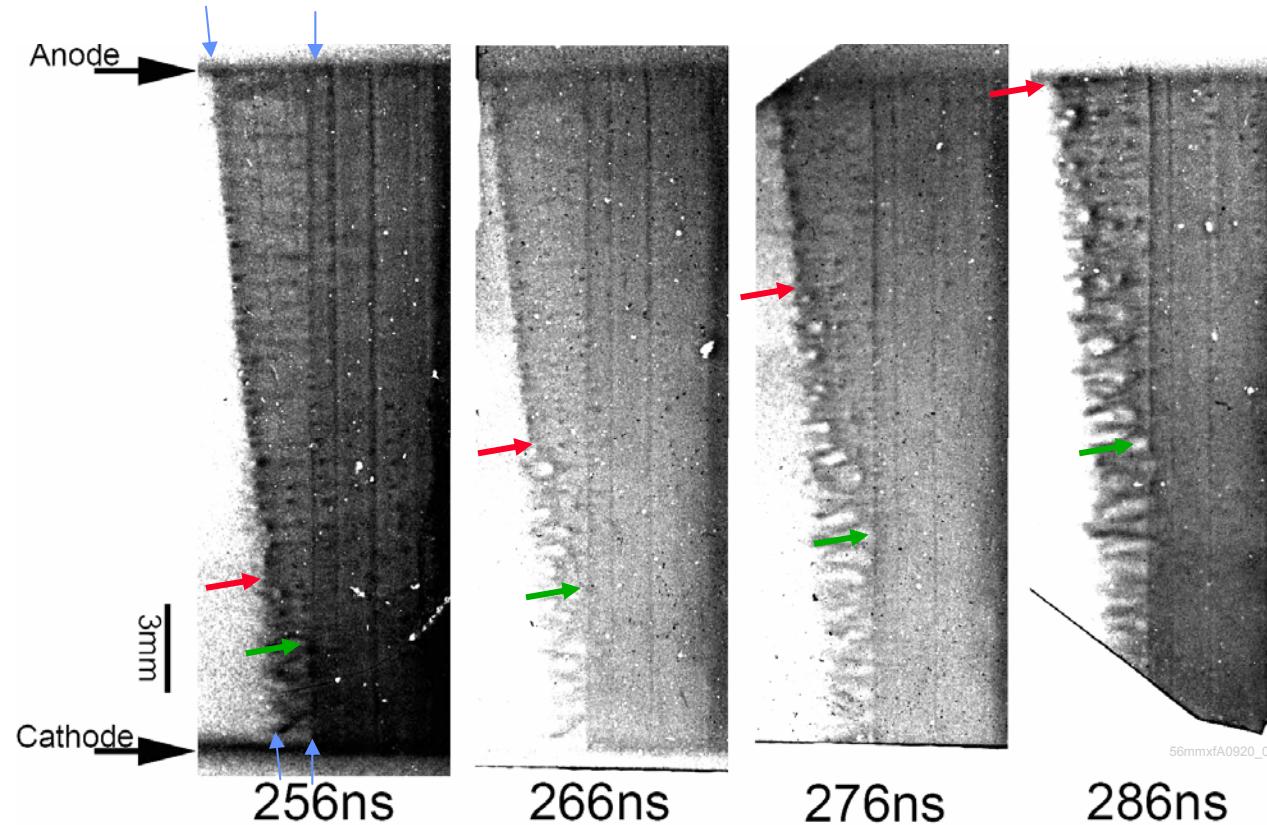


- Time of mass depletion (for single or nested) is expected to be a function of axial position

D.J. Ampleford, PhD Thesis, Imperial College London, 2005  
D.J. Ampleford et al., to be submitted

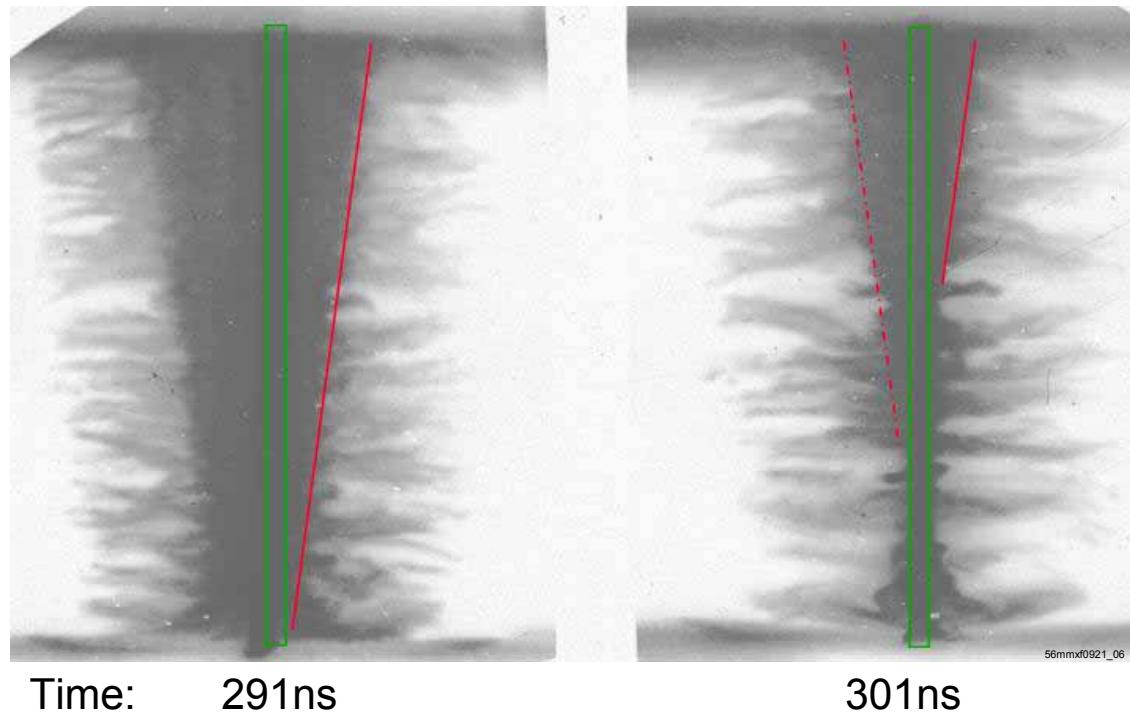
# A conical outer array can be used to seed an axial zipper

- Nested conical data agrees with predicted delay in time of wire breakage



- Also see that zipper of implosion gives an axial dependence time of interaction with the inner array

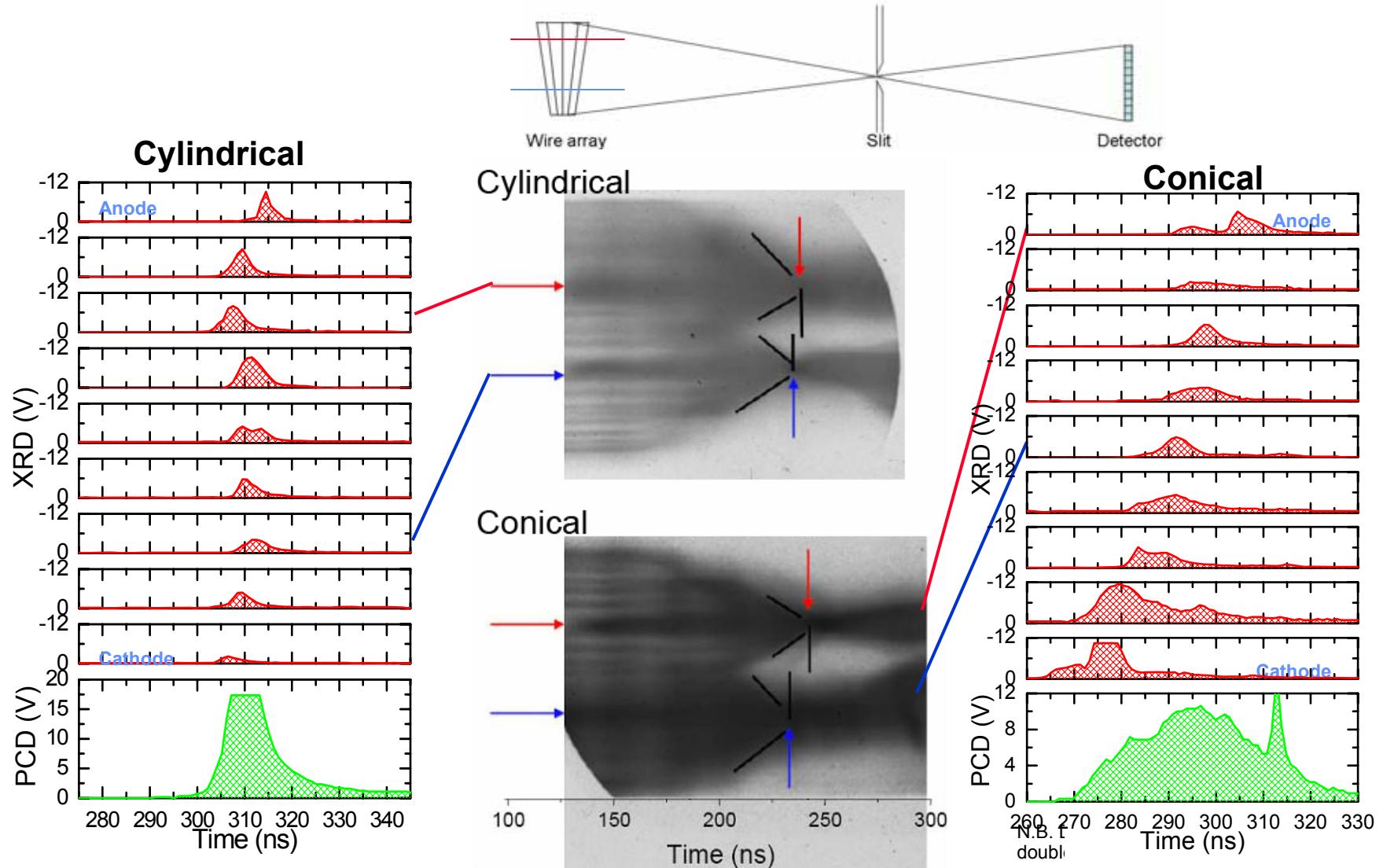
**Axial zipper survives through the interaction, leading to a zippered implosion of inner**



- Data indicates change in implosion position 13.9mm in 10ns
- Implies zipper velocity  $\sim 139\text{cm}/\mu\text{s}$
- Zipper can be extrapolated to estimate zipper along full axis  $\Delta t_{\text{zip}} \sim 16.5\text{ns}$
- On MAGPIE, left-right asymmetry also present due to concentricity issue,
  - will effect pulse, but temporal effect is less ( $\Delta t_{\text{L-R}} \sim 7\text{ns}$ )

# Seeded zipper translates into a zippered stagnation, and elongated x-ray pulse

- Twin radial optical streak and zipper array each indicate conical outer zippers stagnation
- Axial dependence of stagnation time leads to pulse lengthening

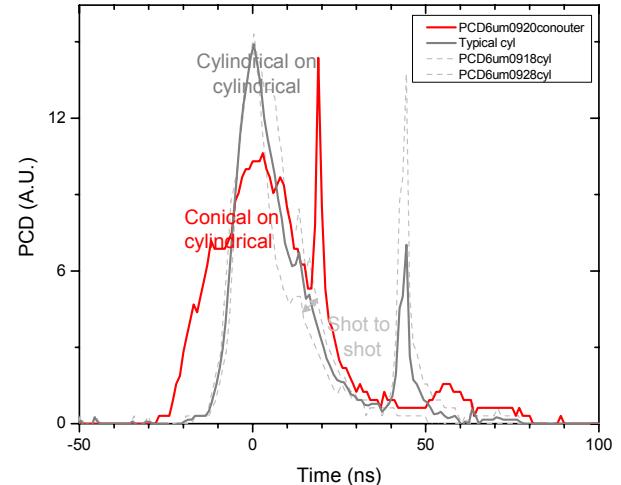


# Comparison of conical and cylindrical outers indicates success at widening main pulse



## For conical outer on MAGPIE:

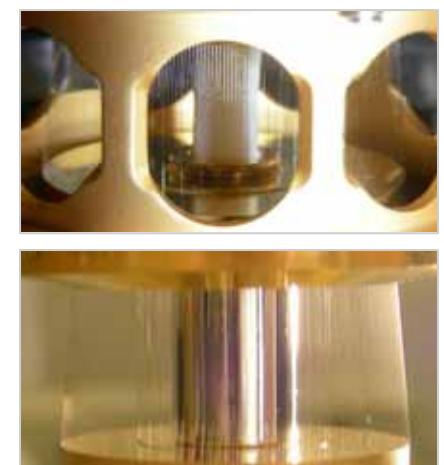
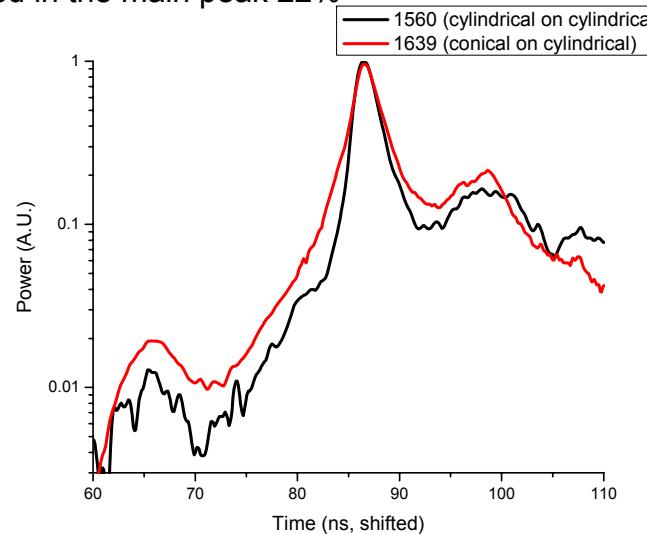
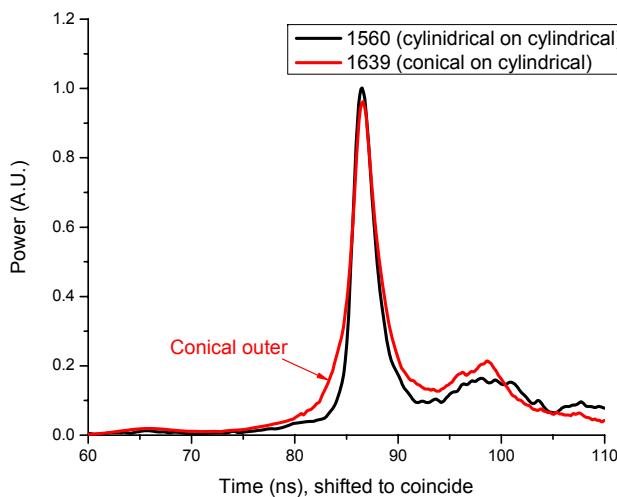
- See a longer rise in x-ray pulse with outer array
- Peak power down,
- Total energy similar (possibly 25% higher for conical outer)



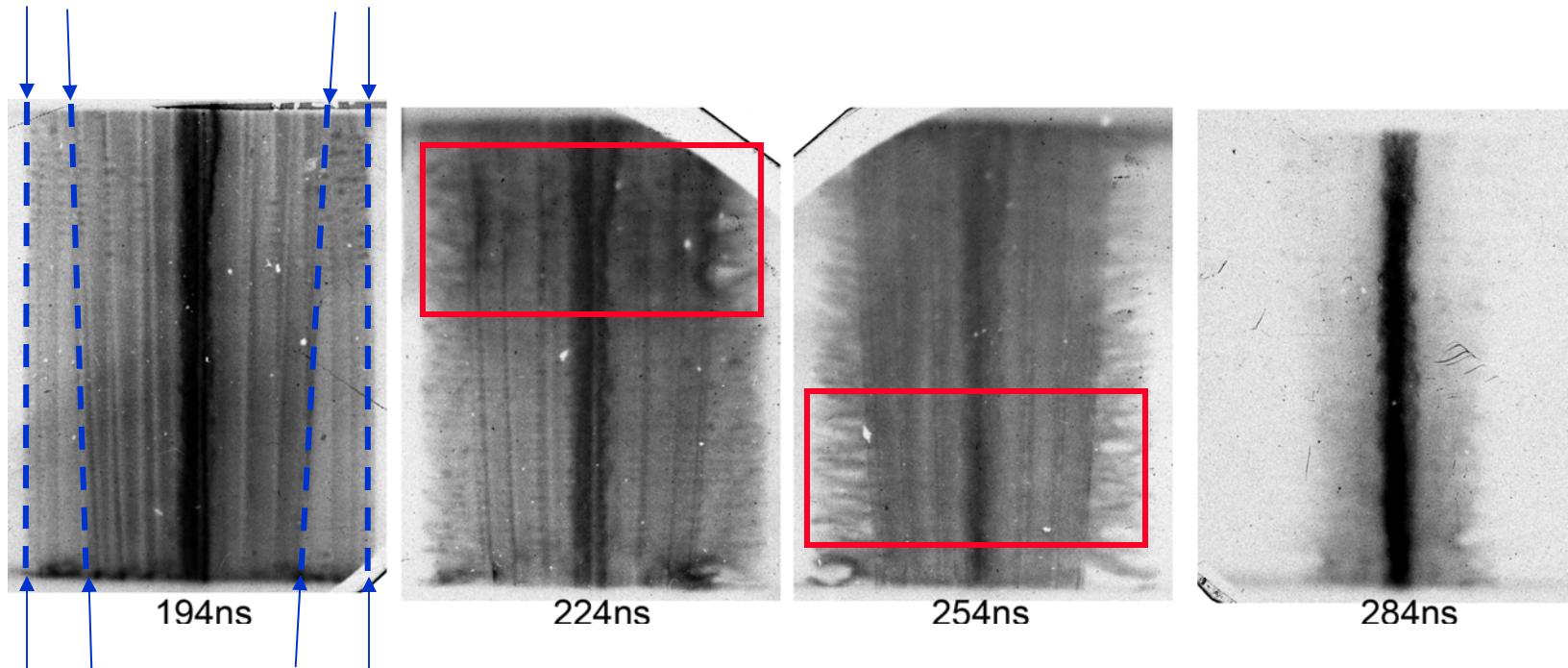
## For conical outer on Z:

(Z1639, Cuneo et al., 22mm Cathode, 20m anode, 12mm inner)

- Increased foot pulse power by 45% because of increase in outer velocity
- Increase first step by zippering implosion onto foam
  - power by a factor of ~4.6
  - energy by ~4.2
- Increase energy in the first step from 25 kJ to 104 kJ
- Energy radiated after the first step is unchanged
- Conical outer increases energy radiated in the main peak 22%

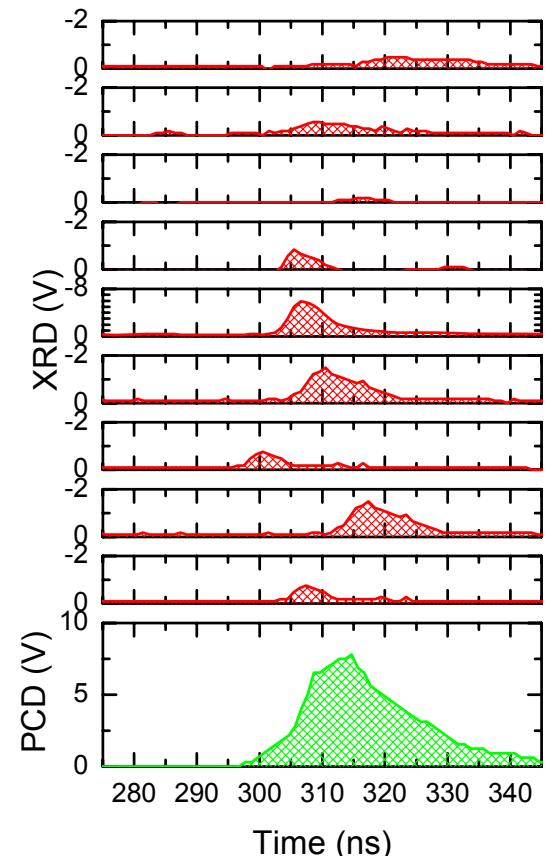
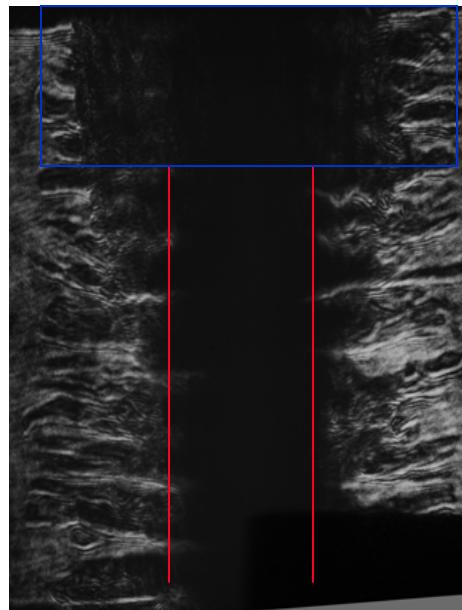
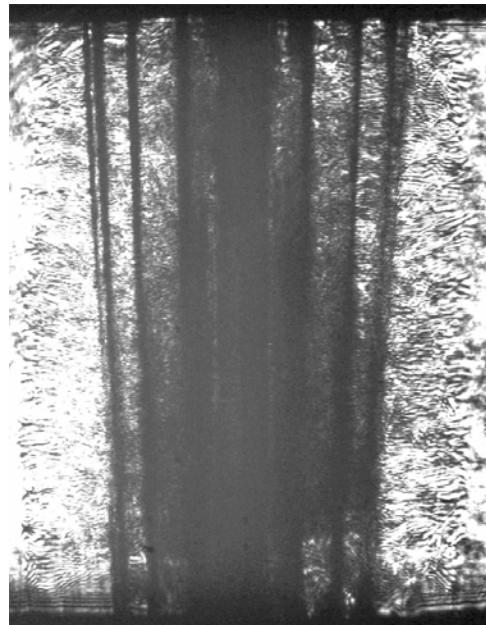


## Inclining inner leads to change of time of flight of outer to inner, and alters timing of Interaction



- Cylindrical wire arrays show variation in time of inner collision with inner array diameter
- Conical inner shows that time of interaction varies with  $z$
- Power pulse will be lengthened
- Zipper less significant after inner ablation

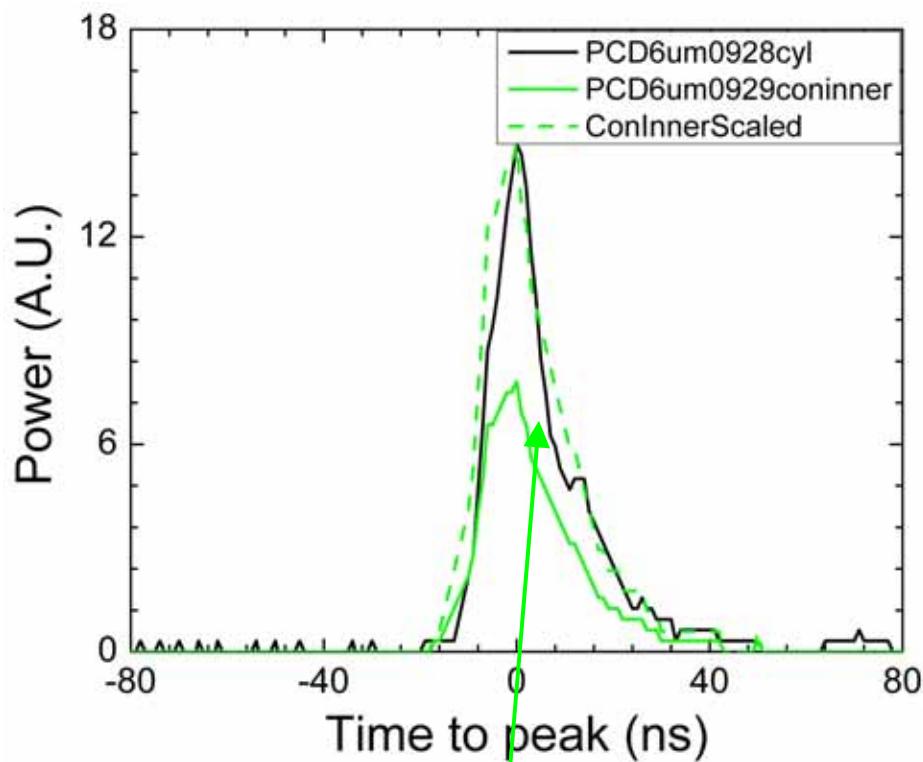
## Nested array with conical inner does not globally zipper stagnation on MAGPIE, but would alter Interaction Pulse



- Laser imaging after interaction indicates no substantial zipper
- Zipper array confirms no zipper in stagnation
- Laser imaging does show top section does not participate in implosion

## Conical inner has some effects, but does not alter width of main pulse on MAGPIE or Z

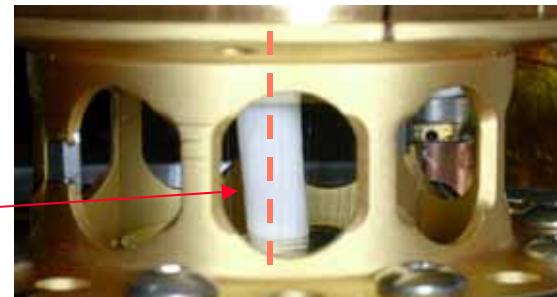
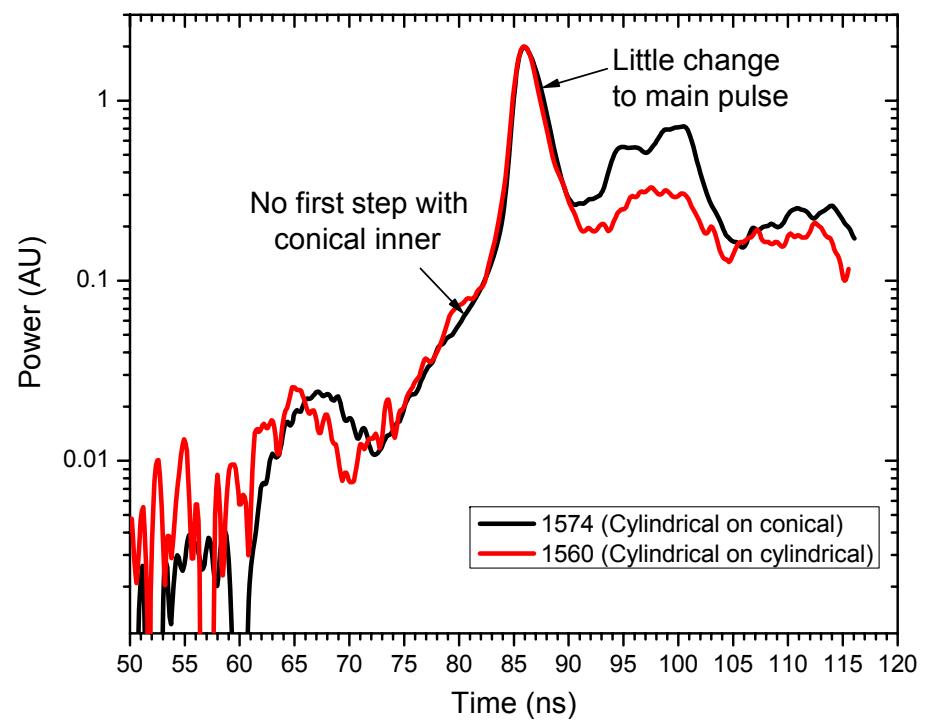
MAGPIE: PCD (hv>200eV)



MAGPIE power lower than cylindrical  
due to part of array not participating in implosion

Lack of first step on Z may be related to foam position

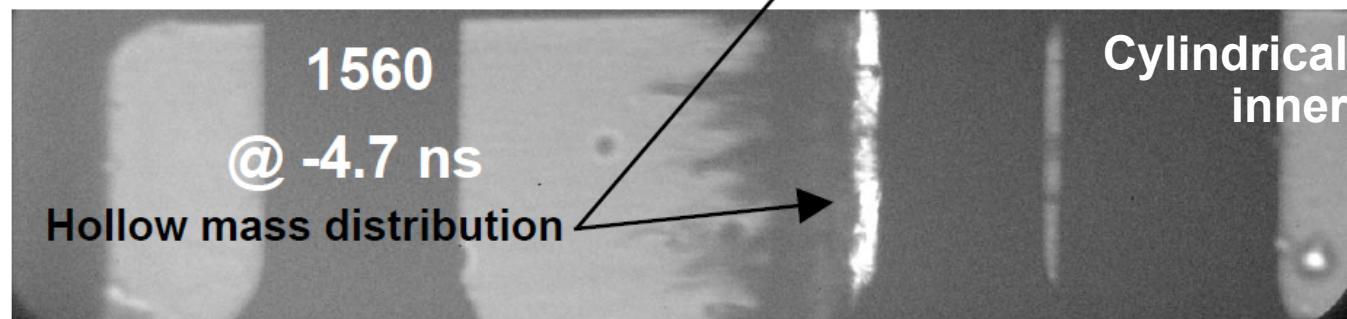
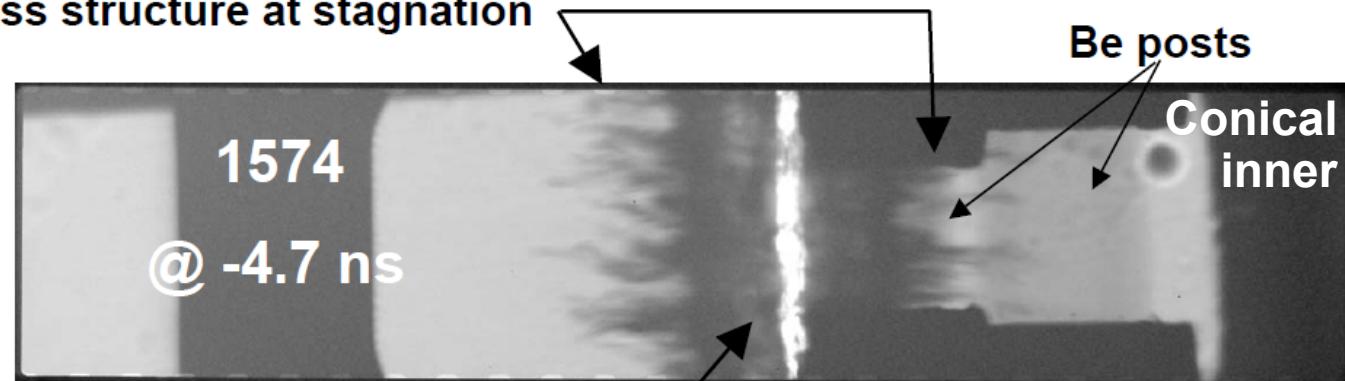
Z: XRD (Z1574, Cuneo et al.)





## Radiography on Z indicates no zipper present after interaction

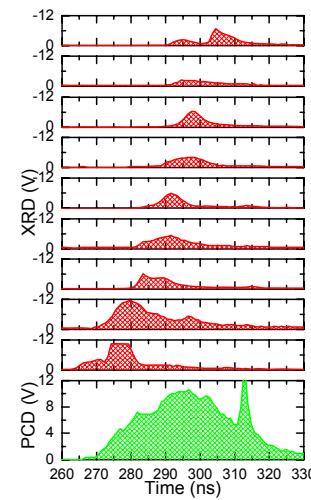
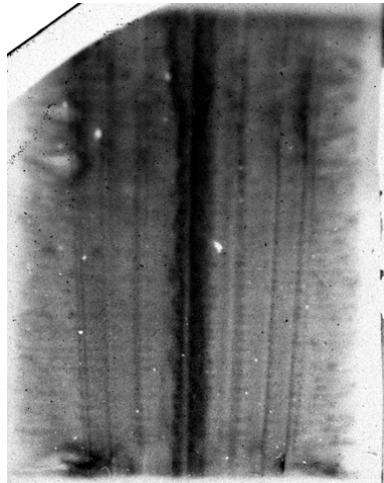
3D mass structure at stagnation



- For Z see no evidence of change to the mass distribution post-interaction



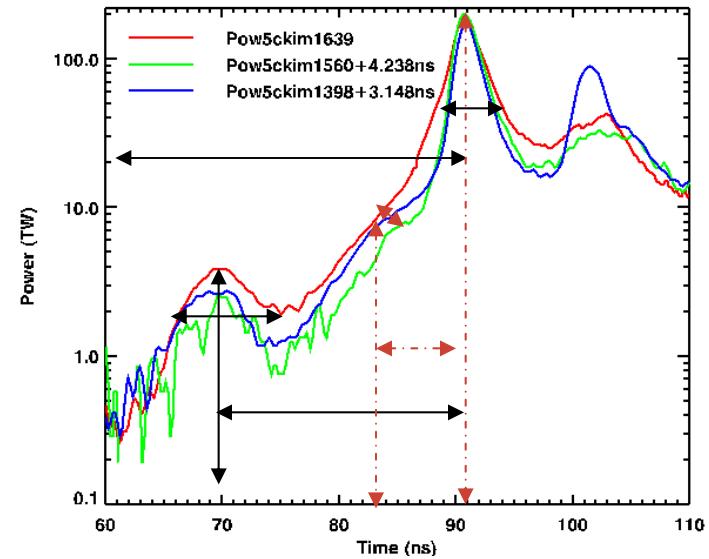
## Part 2 Summary: What can we now control



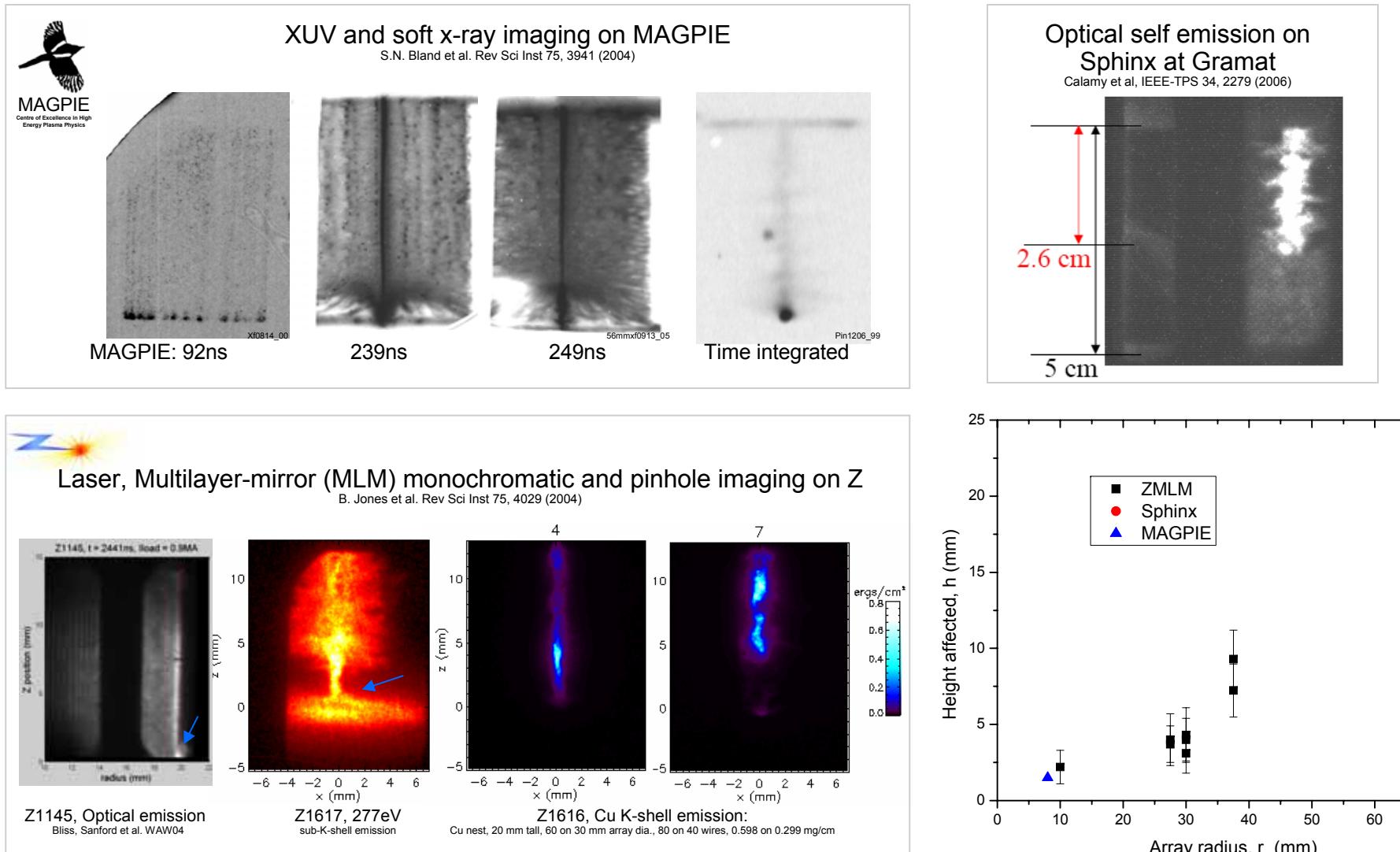
- Conical inner can control
  - Time scale of interaction
- Conical outer can control
  - Time scale of interaction (MAGPIE)
  - Time scale of main stagnation
- Need a more quantitative comparison

Combine with previous data (Cuneo et al.), now have control of:

- Time of peak (Outer mass, inner mass, outer diameter)
- Interaction to peak (inner diameter and mass)
- Pulse length of stagnation (outer angle)
- Pulse length of interaction (relative angle between outer and inner – needs verification)
- Understanding of interaction may lead to control of amplitude and length (need more experiments)



# End effects impact arrays at all current levels



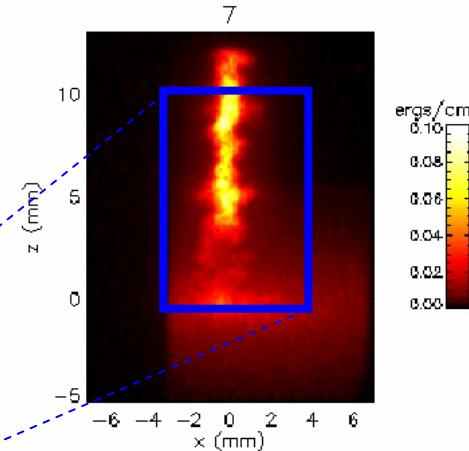
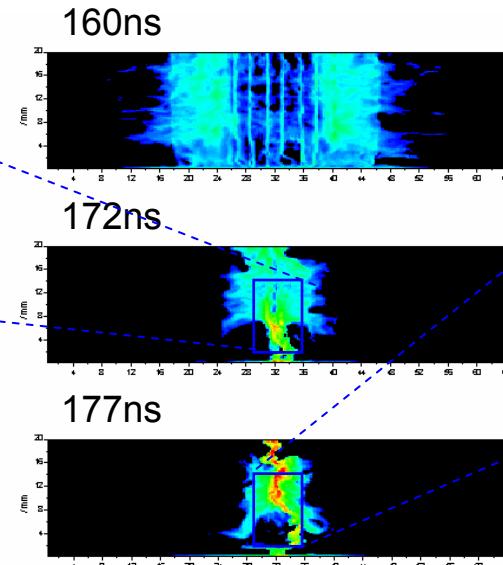
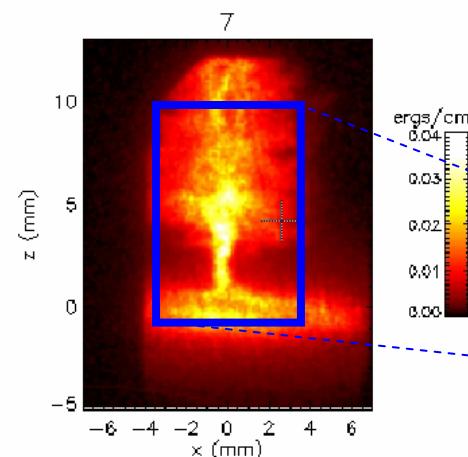
Data and diagnostics from:

B. Jones, C.A. Coverdale, D. Bliss, T.W.A. Sanford et al. (SNL),  
H. Calamy (CEA), S.N. Bland, S.V. Lebedev & J.P. Chittenden (IC)



# Simulations can reproduce effect and show possible mitigation with step

## Integrated emission profiles

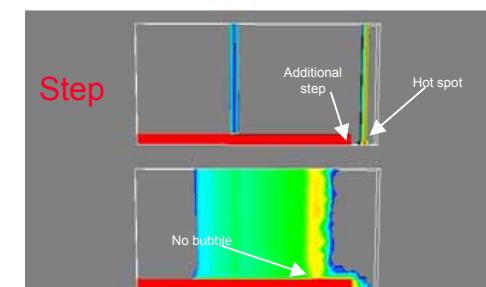
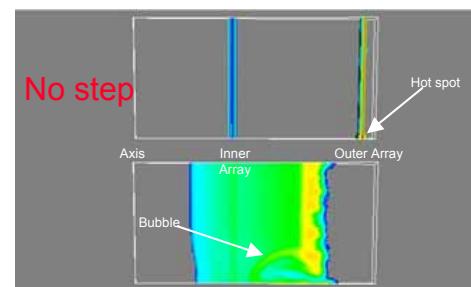
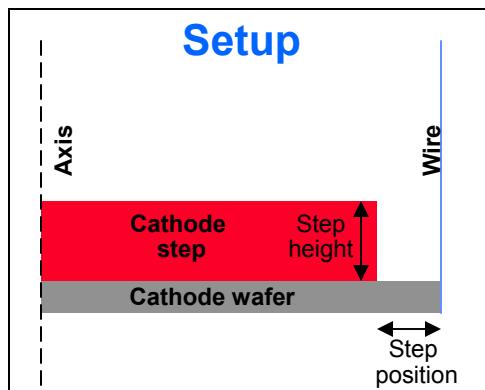


Bubble initiated with 0.5eV hotspot at the base of each wire

## Simulation of Z Shot 1617

- 60mm on 30mm array diameters
- 80 on 40 copper wires

## Presence of step on cathode step changes dynamics



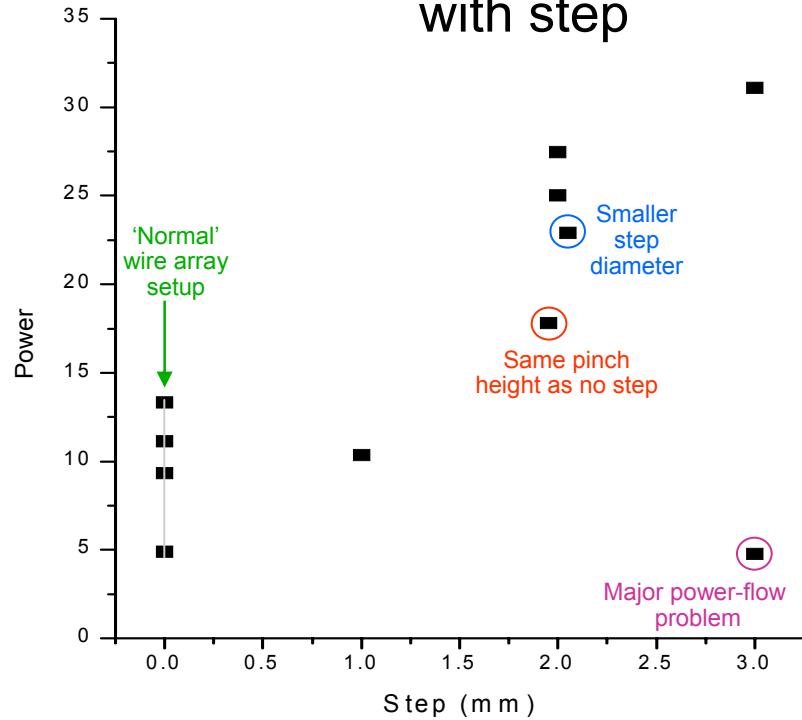
Simulations by C. Jennings using  
Gorgon 3D resistive MHD code  
J.P. Chittenden et al. PPCF 46, B457 2004

Data and diagnostics from B. Jones, C.A. Coverdale (SNL)  
Simulations by C.A. Jennings (SNL)

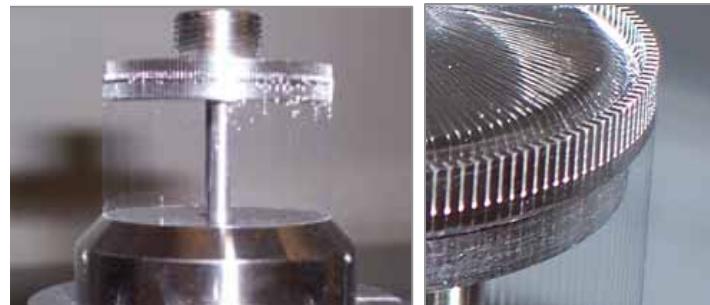
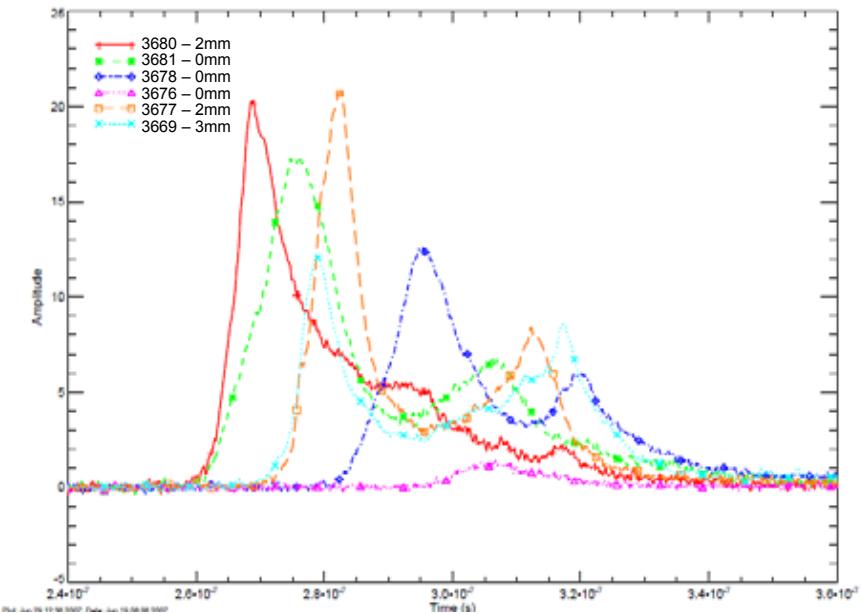
# Initial analysis of data from recent Saturn shots indicates zippering is eliminated



Peak total power increased with step



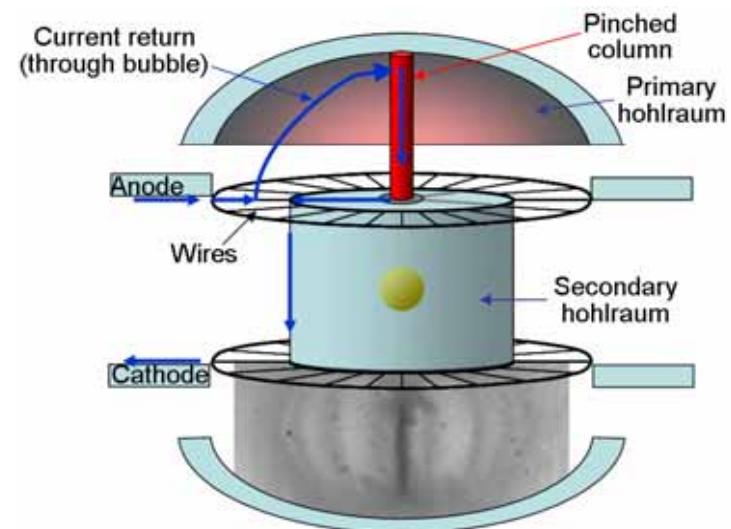
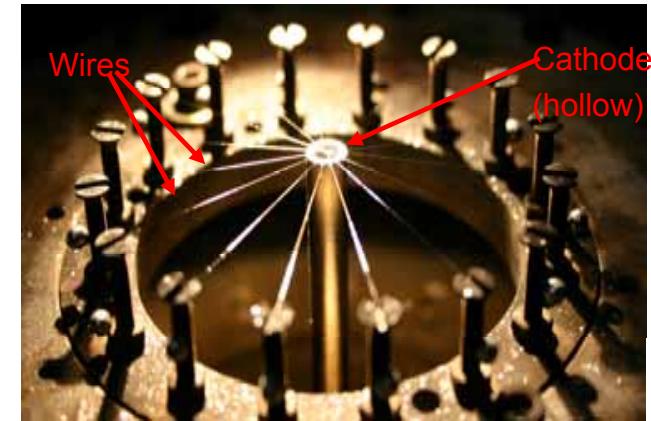
K-shell rise time shortened



- Data indicates higher powers with the step present
- K-shell diagnostics demonstrate a much faster rise with the step, however not higher K-shell yield
- Analysis still in progress!!

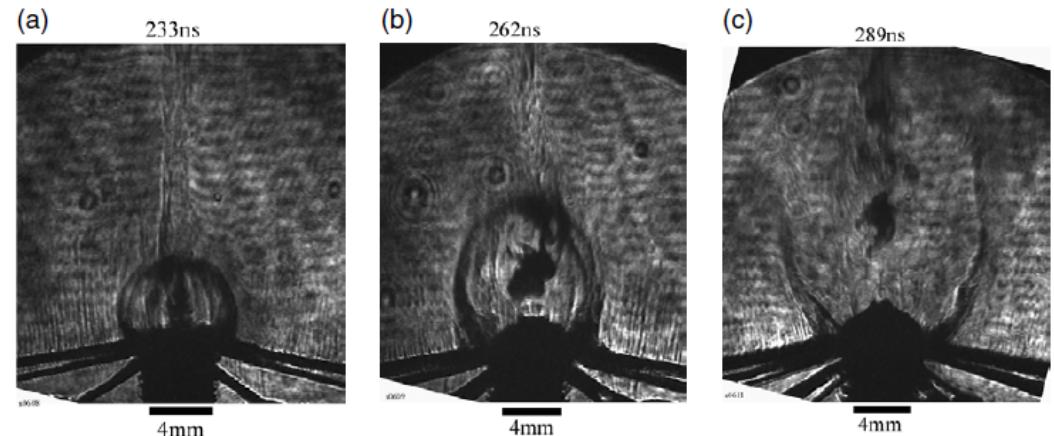
# Bubble formation not necessarily a bad thing! Radial wire arrays have applications as an ICF driver

- New idea for a compact x-ray source to drive ICF hohlraums
- MAGPIE data indicates significantly higher power densities than cylindrical arrays
- Experiments planned for Saturn to investigate scaling to higher current drives
- Potentially could get  $>4x$  higher power densities

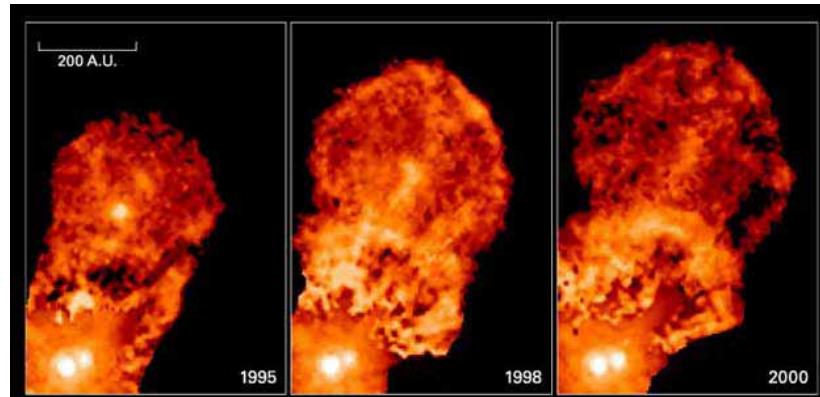


## And also to astrophysics

- The evolution of a radial wire array includes the formation of a magnetically driven jet
- For some configurations multiple jets can be formed
- Magnetic cavities are thought to be most likely formation mechanism for protostellar jets
- Simulations recreate features of lab-jet

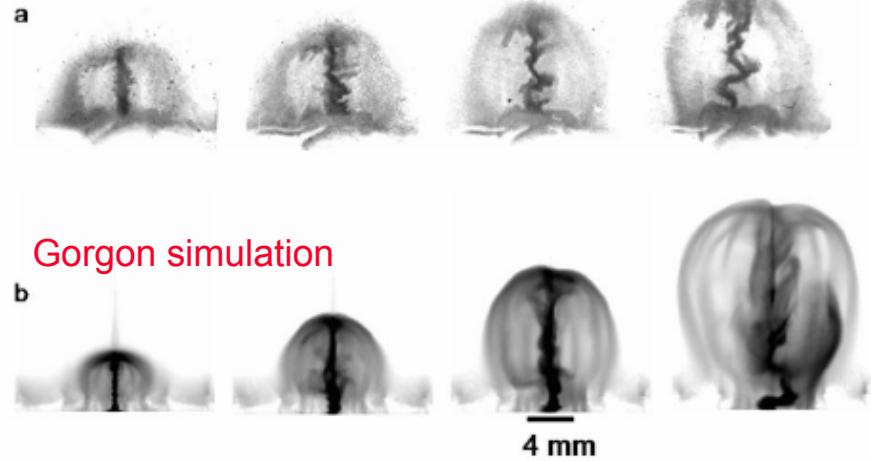


Astrophysical jet



XZ TAU. Credits: John Krist (STScI) et.al. WPFC2, HST, NASA

MAGPIE data

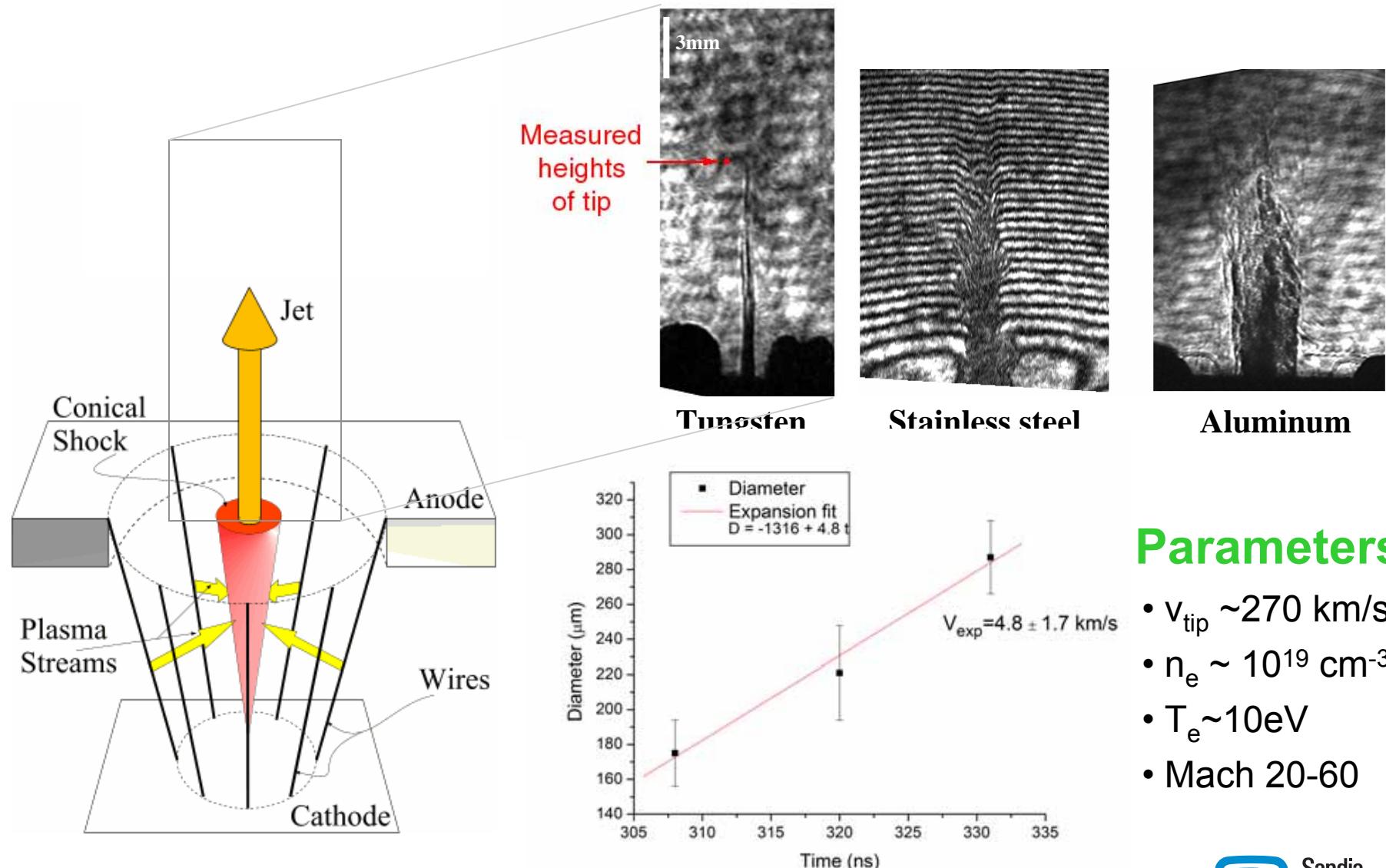


Gorgon simulation

- High Energy Density Laboratory Astrophysics is an emerging area of research, which provides understanding and benchmarking to the astrophysical community
- Need to carefully consider appropriate scaling requirements to compare systems/codes

S.V. Lebedev et al., Mon Not Royal Ast. Soc 361, 97 (2005).  
A. Ciardi et al., Phys Plasmas 14, 056501 (2007)

# Jets produced by conical wire array z-pinches can be used to model protostellar jets far from source

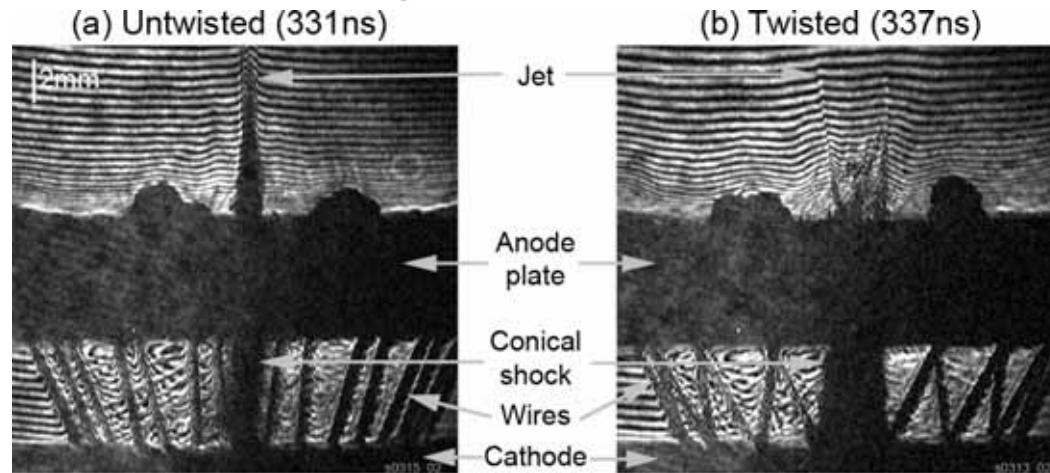


## Parameters

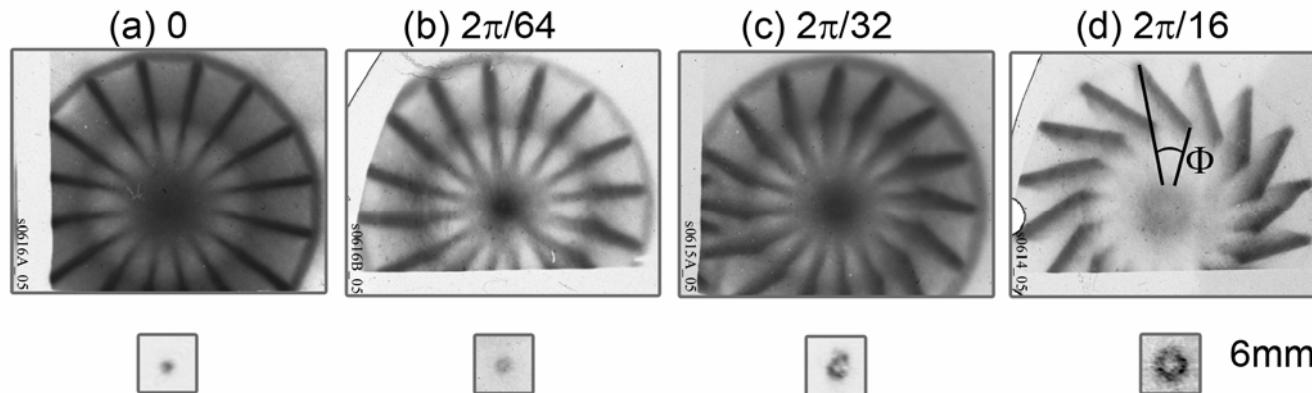
- $v_{\text{tip}} \sim 270 \text{ km/s}$
- $n_e \sim 10^{19} \text{ cm}^{-3}$
- $T_e \sim 10 \text{ eV}$
- Mach 20-60

S.V. Lebedev et al., *Astrophysical Journal*, 564:113, 2002.  
D.J. Ampleford, PhD Thesis, Imperial College London, 2005

# Conical wire arrays can be used to model protostellar jets further from the source object.

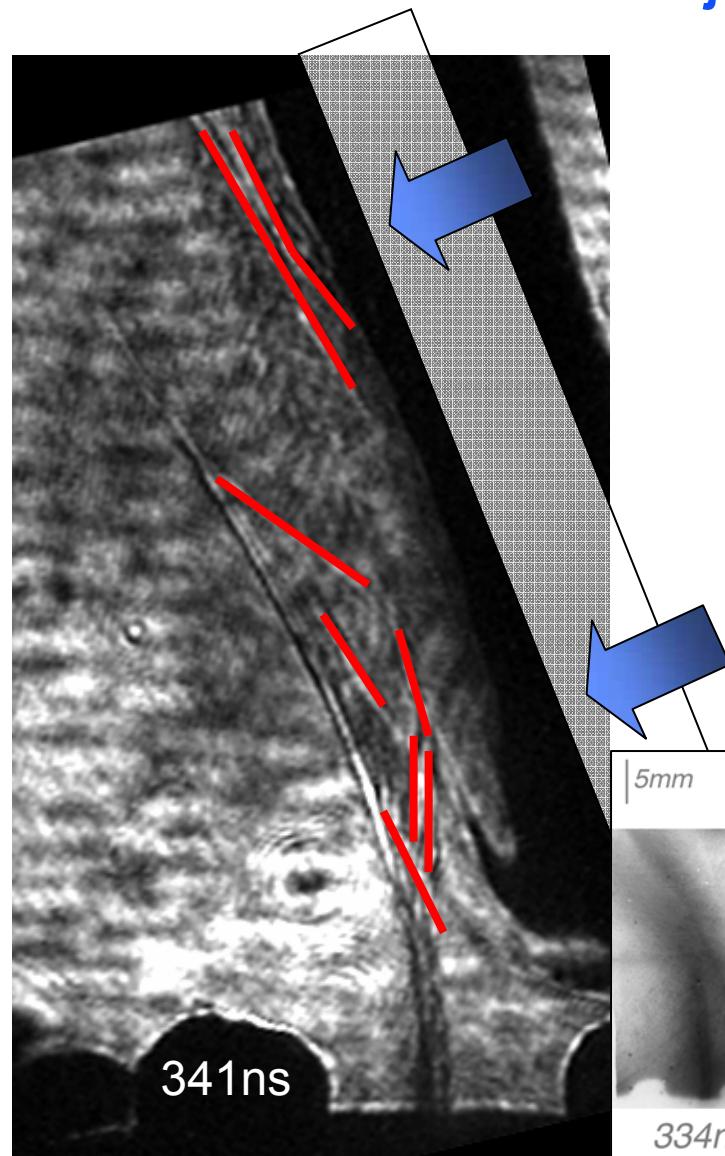


- By twisting array, generate  $B_z$ , hence introducing  $J_r \times B_z$  component to Lorentz force
- Rotation leads to hollow shock on axis
- Jet produced maintains angular momentum, and becomes divergent for larger rotation rates.



D.J. Ampleford et al., *Astrophys Space Sci* 307, 51 (2007)  
D.J. Ampleford et al. *submitted* (2007)

## Conical wire arrays can also be used to understand interaction of jets with ambient medium



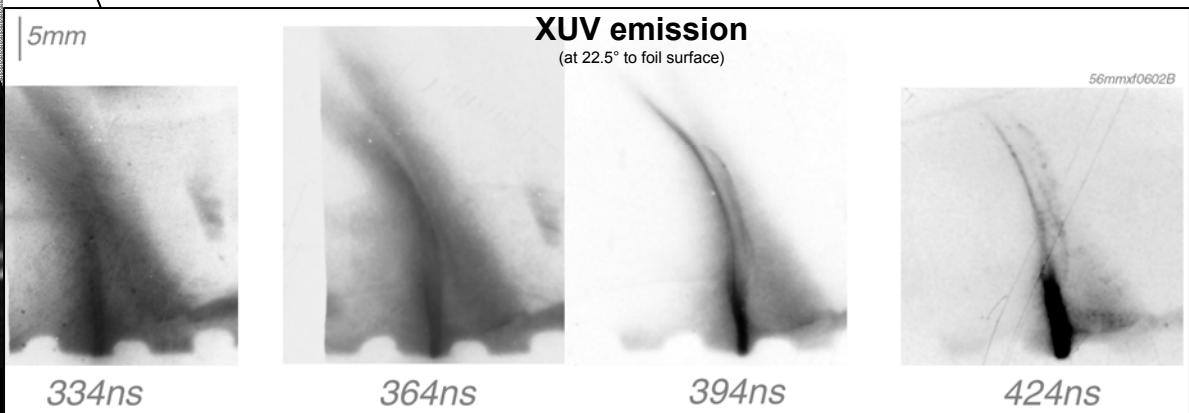
Use emission from the array to ablate a CH foil, producing a side-wind

Side-wind interacts with jet, causing deflection

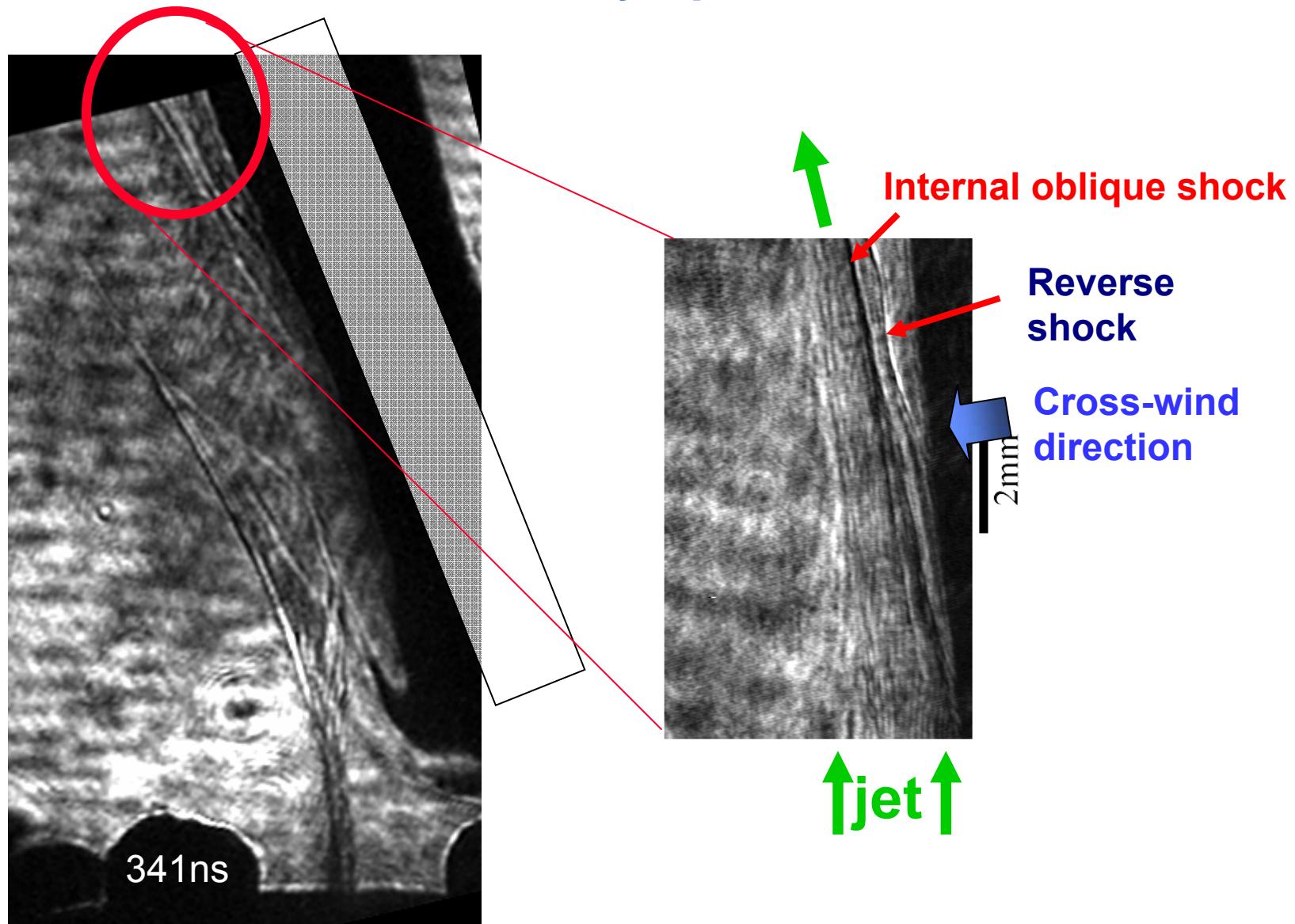
Various internal shocks seen in jet

Internal oblique shock causes thermalisation of kinetic energy

Formation of new bow-shock starting



# Low density tip interaction



M.A.G.P.I.E.

D.J. Ampleford, PhD Thesis, Imperial College London, 2005  
D.J. Ampleford et al., *Astrophys Space Sci* 307, 29 (2007)

# Summary

- Wire array Z-pinchers have many applications, especially when tailored to meet the needs
- For ICF, nested arrays can provide a tailored pulse shape
  - Interaction of implosion with a prefill that has been perturbed by the presence of the inner array leads to an interaction pulse
  - Altering array configuration (e.g. conical nested) can lead to additional pulse shaping controls
- Novel array configurations can be used for
  - Improving the performance of cylindrical arrays
  - As a source for compact ICF hohlraums
  - To understand physics of astrophysical jets
    - In MHD mode (formation and near the source)
    - In hydrodynamic mode (far from the source, including interaction with ambient medium)

