

Boosting Photoelectric Conductivity in Porphyrin-Based MOFs Incorporating C₆₀

Saied Md Pratik, Laura Gagliardi, Christopher J. Cramer**

Department of Chemistry, Minnesota Supercomputing Institute, and Chemical Theory Center,
University of Minnesota, Minneapolis, Minnesota 55455, United States

Corresponding Authors

*E-mail: gagliardi@umn.edu and cramer@umn.edu.

Abstract:

Electronic structure calculations show that guest C₆₀ in the porphyrin-containing metal organic frameworks Zn₂(TCPB)(DA-ZnP) (DA-MOF; H₄TCPB = 1,2,4,5-tetrakis(4-carboxyphenyl)benzene, DA-ZnP = [5,15-bis[(4-pyridyl)ethynyl]-10,20-diphenylporphinato]zinc(II)) and Zn₂(TCPB)(F-ZnP) (F-MOF; F-ZnP = [5,15-di(4-pyridyl)-10,20-bis(pentafluorophenyl)porphinato]-zinc(II)) engenders high photoelectrical conductivity due to efficient donor-acceptor charge transfer (CT) interactions. Structural modifications at the meso position of the porphyrin influence the preferred positions of C₆₀ within the frameworks, giving rise to host-guest interactions with different anisotropic structural, electronic, and optoelectronic properties. A preferred slipped-parallel π -stacked interaction of C₆₀ that is predicted for NH₂-substituted DA-MOF and F-MOF fosters strong charge-transfer (CT) transitions and lowers band gaps by ~1.0 eV compared to the pristine DA-MOF and F-MOF. Hopping rates computed using Marcus theory are found to be anisotropic and accelerated by multiple orders of magnitude across π -stacked interfaces created by C₆₀ incorporation, a consequence of strong electronic coupling between initial and final diabatic states. Calculations indicate that photoinduced electron transfer (PET), as well as direct CT from porphyrin to C₆₀ upon irradiation, triggers a charge separation process that leads to the formation of what should be long-lived electron-trapped states at the heterojunctions. Design principles revealed here for the control of photophysical and electron transfer processes will be useful for constructing new MOF-derived visible- and infrared-based optoelectronics.

Introduction

Metal organic frameworks (MOFs) are porous crystalline materials that have attracted extensive attention due to their wide-range of potential applications encompassing gas storage and separation, catalysis, chemical sensing, photocatalysis, and drug delivery.¹⁻⁹ Most MOFs are electrically insulating or poor conductors because of the insulating character of the organic ligands and/or energy mismatch or otherwise ineffective interaction between the ligand and metal orbitals (e.g., π -d interactions).¹⁰⁻¹¹ Intrinsic electrical conductivity or fast charge delocalization within the framework is the key to electrochemical applications.¹² Therefore, their implementation for such applications is significantly limited and further improving their electronic properties remains challenging.¹³ While, structural regularity and synthetic flexibility of MOFs allow hierarchical and unambiguous assemblies of multiple chromophoric units and nodes, it remains a challenge to achieve a systematic structural-functional basis in MOFs that improves their electrical conductivity.¹⁴⁻¹⁹ A detailed understanding of the nature of charge transport mechanisms in MOFs is vital to design the next generation of MOF-based optoelectronic materials.

Until recently, several experimental and theoretical reports have addressed the strategies to unlock the low energy charge transfer (CT) pathways within MOFs to engineer electrical conductivity -for example by utilization of redox active linkers,^{15, 20} mixed valency,²¹⁻²³ formation of π -stacking with the MOFs,²⁴ and formulation of 2D π -conjugated MOFs.²⁵⁻²⁶ An alternative strategy to induce electrical conductivity is to infiltrate the pores with an appropriate electron acceptor or donor that can facilitate host-guest CT interaction within the framework and introduce new pathways for carrier migration.^{24, 27-29} In this context, Talin et al. reported that the electrical conductivity of benzenetricarboxylate-based MOFs (HKUST-1) increased by six orders of magnitude by infiltrating it with tetracyanoquinodimethane (TCNQ), as the nitrile of TCNQ bridges Cu(II) sites and induces a continuous CT pathway throughout the MOF unit cell.³⁰ Farha and coworkers introduced electron deficient Nickel(IV) bis(dicarbollide) as a guest molecule in the microporous channel of NU-1000 MOF to facilitate donor-acceptor, interactions and the resulting MOF became electrically conductive.³¹ Recently, covalent and non-covalent infiltrations of C₆₀ in the pores of MOFs have received attention as they offer unique pathways for efficient CT from linker-based donors to the C₆₀ acceptor and render the MOF electrically

conductive.³²⁻³⁵ Hupp and coworkers have reported that the conductivity of NU-901 MOF increased substantially upon C₆₀ incorporation due to CT from the pyrene-based linker to C₆₀.³² Shustova and co-workers have integrated porphyrin and fullerene-based linkers within MOF scaffolds and demonstrated emergent photophysical and energy-transfer processes.³³ Very recently, Souto et al. reported that the electrical conductivity in a TTF-containing MOF (MUV-2) increased by ~2 orders of magnitude after C₆₀ incorporation due to CT between C₆₀ and the TTF-based framework.³⁶

Porphyrins and metalloporphyrins have seen widespread use in the design of molecular optoelectronics owing to their exceptional photophysical, electrochemical, and structural properties.³⁷⁻³⁹ Their installation as linkers in a rigid framework keeps them well separated and prevents self-quenching, while at the same time permitting utilization of the full range of porphyrin properties. Such architecture have the potential to unlock new pathways for solar light harvesting, energy transport and exciton migration.⁴⁰⁻⁴² Porphyrin derived MOFs are superior to their monomeric units and have been exploited for solar energy applications, in particular for the design of photovoltaics and photocatalysts.⁴³⁻⁴⁵ As two recent examples, Hupp and coworkers have synthesized two zinc porphyrin-derived MOFs (Figure 1, a and b), Zn₂(TCPB)(DA-ZnP) (DA-MOF; H₄TCPB = 1,2,4,5-tetrakis(4-carboxyphenyl)benzene, DA-ZnP = [5,15-bis[(4-pyridyl)ethynyl]-10,20-diphenylporphinato]zinc(II)) and Zn₂(TCPB)(F-ZnP) (F-MOF; F-ZnP = [5,15-di(4-pyridyl)-10,20-bis(pentafluorophenyl)porphinato]-zinc(II)).⁴⁶ We are particularly interested in these two MOFs as the highly ordered arrangement of porphyrin units within their frameworks should facilitate long-distance and directional migration of photogenerated excitons, making them efficient for energy-transport and related applications. In a follow up study, these authors reported the light harvesting efficiency of DA-MOF and F-MOF to be enhanced upon incorporation of CdSe/ZnS core/shell quantum dots (QDs) because of energy transfer from the QDs to the MOFs.⁴⁷

In this work, we systematically apply density functional theory (DFT) to study the structures of DA-MOF and F-MOF with incorporated C₆₀ (referred to hereafter as C60@DA-MOF and C60@F-MOF, respectively) as well as photophysical and electron-transfer processes that occur therein. We show that electrical conductivity is significantly enhanced through efficient CT interactions from the porphyrin-based donor to the C₆₀-based acceptor. To improve

the degree of CT between the donor-acceptor conjugates, we also explore structural modifications at the meso position of the porphyrin, which are found to influence the preferred positioning of C₆₀ within the crystallographic framework. We find that charge carrier mobility across π -stacked heterojunction interfaces increases by several orders of magnitude upon C₆₀ doping. Finally, we demonstrate that the donor-acceptor interface readily promotes photoinduced electron transfer from porphyrin to C₆₀, allowing the formation of charge separated states that should make these heterostructures excellent platforms for designing optoelectronics.

Computational details:

First-principles calculations for structural optimizations were performed within the framework projector augmented wave (PAW) approach as implemented in the Vienna ab initio simulation package (VASP 5.4.4).⁴⁸ Electron exchange-correlation was incorporated by the generalized gradient approximation (GGA) employing semi-local PBEsol functional.⁴⁹ The ionic core-electron interactions were taken into account using ultrasoft pseudopotentials. During geometry optimization, all atomic positions were fully relaxed unless Hellman–Feynman forces become less than 0.05 eV/Å. A planewave kinetic energy cut-off of 500 eV and 2×2×2 k-point mesh was used with an electronic energy convergence criterion of 10⁻⁶ eV. Lattice parameters of the unit cell varies ~1.0% (Table S1) upon optimization with respect to the experimental CIFs which thereby justifying the choice of our level of calculations. Subsequent single-point calculations for band energies and densities of states were performed using the screened hybrid functional of Heyd, Scuseria and Ernzerhof (HSE06).⁵⁰ Dispersion interactions were taken into account employing Grimme’s DFT-D3 empirical formalism.⁵¹ We have incorporated one C₆₀ per unit cell and the binding energy (ΔE) within the MOF scaffolds was computed as $\Delta E = E(\text{C}_{60} \text{ incorporated MOF}) - E(\text{MOF}) - E(\text{C}_{60})$.

Time-dependent DFT (TD-DFT) calculations using HSE06 were performed on cluster models as retrieved from each of the system containing bare porphyrin-based linkers and C₆₀. The 6-31+G(d) basis set was used for C, H, N and F, while for the Zn Stuttgart–Dresden–Dunning (SDD) pseudopotential was employed.⁵²⁻⁵³ The singlet–singlet excitations were modeled considering 100 excited states. These calculations were performed using Gaussian 16 suite of programs.⁵⁴

Charge transfer integral (or electronic coupling) calculations were carried out on cluster models, chosen on the basis of center to center distances between the neighboring molecules in each system. These calculations were performed using the fragment approach as implemented in the Amsterdam density functional (ADF) package⁵⁵⁻⁵⁶ with the M06-2X⁵⁷ functional and the TZP all-electron basis set.

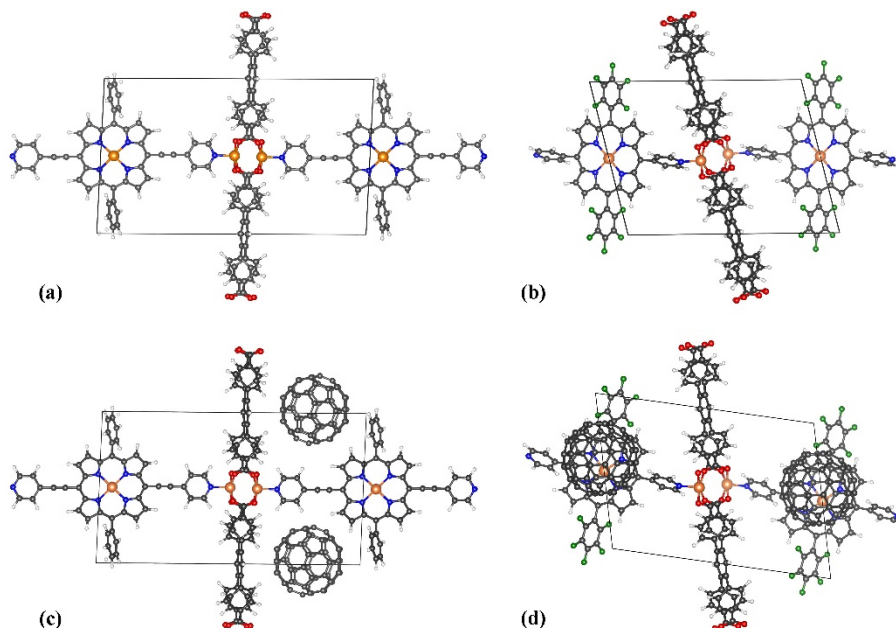


Figure 1. Optimized crystal structures of (a) DA-MOF and (b) F-MOF. (c) and (d) are their C₆₀ incorporated analogues. Atomic colors are H (white), C (gray), N (blue), O (red), F (blue-green), and Zn (orange).

Results and Discussion

Both DA-MOF and F-MOF contain zinc porphyrin cores as photoactive building blocks and their meso positions (10- and 20- positions) can accommodate “out of plane” phenyl and pentafluorophenyl substituents, respectively (Figure S1 in SI). In contrast to F-MOF, DA-MOF possesses extended conjugation through the incorporation of two additional acetylene moieties at the 5- and 15- positions which makes the pyridine units coplanar with the porphyrin (Figure S1 in SI) and increases the pore size as well. Our calculations indicate that C₆₀ can be incorporated

in the pore channel of DA-MOF only through edge-on interaction (Figure 1c), while in F-MOF C_{60} can be incorporated in between two co-facial porphyrin moieties, leading to a complete π -stacked arrangement (Figure 1d). Optimized lattice parameters following C_{60} incorporation are almost unchanged for DA-MOF, but change significantly for F-MOF (Table S1 in SI). While this suggests that there may be some kinetic barrier to the incorporation of C_{60} in the latter, the C_{60} binding energies (ΔE) computed at the PBEsol-D3 level are -1.23 eV and -1.32 eV for DA-MOF and F-MOF, respectively, suggesting that incorporation is quite favorable in both instances. The higher binding energy in F-MOF is consistent with the better aligned π - π stacking interactions between the porphyrin and C_{60} in its case. Such strong C_{60} binding energies are consistent with dispersion type interactions reported for fullerene encapsulated porphyrin-based cage complexes.⁵⁸

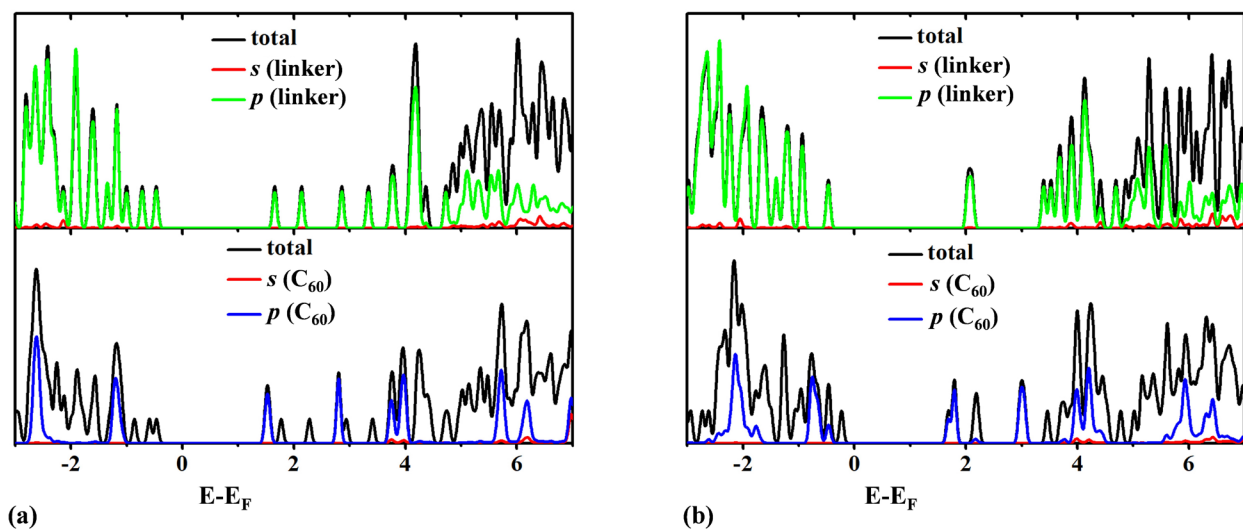


Figure 2. Calculated density of states for (a) DA-MOF and (b) F-MOF, before (upper panel) and after (lower panel) C_{60} incorporation. The energy is relative to the Fermi level and the y axis for the s - and p -orbitals is scaled by a factor 1.75.

To understand the influence of C_{60} incorporation, we have calculated the total density of states (TDOS) for DA-MOF and F-MOF before and after incorporation of C_{60} (Figure 2). Calculated optical band gaps (E_g) for DA-MOF and F-MOF are 1.66 eV and 2.04 eV,

respectively, which is in the visible region as expected for the porphyrin-based linkers. Partial density of states (PDOS) analysis reveals that both the highest occupied crystal orbitals (HOCO) and lowest unoccupied crystal orbitals (LUCO) predominantly comprise the p-orbitals of the linkers, and thus the absorption and emission properties of both MOFs are mainly governed by the electronic structure of the linkers (Figure 2, upper panels and Figure S1.b). The band gaps in $C_{60}@DA$ -MOF and $C_{60}@F$ -MOF (Figure 2, lower panels) are 1.50 eV and 1.68 eV, respectively, which values are 0.16 eV and 0.36 eV lower than those of the undoped DA-MOF and F-MOF. Though the band gap reduction is not large, the heterostructures possess delocalized electronic states, where the HOCO mainly comprises p-orbitals of the linkers and the LUCO predominantly comprises p-orbitals of C_{60} (lower panels of Figure 2). There is significant orbital overlap between them as is important for typical donor-acceptor based CT complexes.³² This indicates that photoelectrons and holes generated on the linker should readily transfer to C_{60} at the interfacial region (vide infra for cluster calculations that further explore this).

To further engineer the physicochemical properties of these species, we removed the “out of plane” (perfluoro)phenyl rings from the meso positions of the porphyrins in (F)DA-MOF to assess the effect on donor-acceptor conjugate structures. The resulting MOFs are referred to as H_DA -MOF and H_F -MOF, respectively (Figure S2). Since this operation creates additional void space within the pore channel, different polymorphs of $C_{60}@H_DA$ -MOF and $C_{60}@H_F$ -MOF become possible. For $C_{60}@H_DA$ -MOF, we found three polymorphs, Type-I, Type-II, and Type-III, while for $C_{60}@H_F$ -MOF we found only two, Type-A and Type-B (Figure 3). Similar to the parent $C_{60}@DA$ -MOF, Type-I possesses an edge-on interaction between the porphyrin and C_{60} . In Type-II, C_{60} is incorporated between two co-facial porphyrin moieties to generate a complete π -stacking interaction and a double-decker complex. In Type-III, C_{60} is positioned in between two diagonally disposed porphyrin linkers, resulting in a slipped π -stacked complex. The lattice parameters for Type-I and Type-III $C_{60}@H_DA$ -MOF remain almost unchanged compared to H_DA -MOF, while they expand slightly along the stacking direction in Type-II $C_{60}@H_DA$ -MOF to accommodate C_{60} (Table S1). The computed binding energies (ΔE) of C_{60} in Type-I, Type-II, and Type-III are -0.91 eV, -1.16 eV and -0.89 eV, respectively. The larger ΔE in Type-II presumably results from the optimal π - π stacking interaction between the donor-acceptor conjugates.

Type-A for $C_{60}@H_F\text{-MOF}$ shows complete π - π stacking interactions as is found in parent $C_{60}@F\text{-MOF}$, while Type-B has slipped stacked interactions between the porphyrin and C_{60} . The binding energy of C_{60} in Type-A is larger than in Type-B (-1.37 eV vs -1.11 eV) again owing to better π - π interactions in the former.

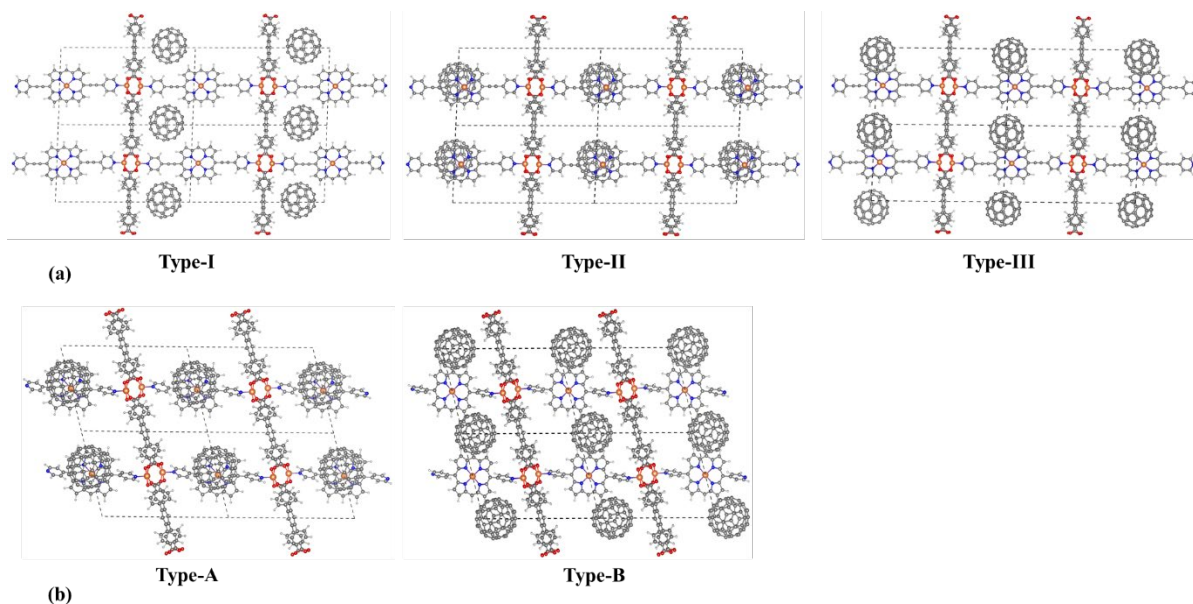


Figure 3. Different crystallographic packing arrangements of C_{60} within the pore channel of (a) $H_DA\text{-MOF}$ and (b) $H_F\text{-MOF}$ shown using a $2 \times 2 \times 2$ supercell. The dashed line represents the unit cell. See caption to Figure 1 for atomic colors.

$H_DA\text{-MOF}$ has a similar density of states and band gap (1.69 eV) as $DA\text{-MOF}$ (1.66 eV) (Table 1, Figure S3), indicating that the diphenyl rings in the meso position have a negligible effect on the band gap. By contrast, $H_F\text{-MOF}$ has an increased band gap (2.34 eV) compared to $F\text{-MOF}$ (2.04 eV) following the removal of the pentafluorophenyl substituents. The computed band gaps of Type-I, Type-II, and Type-III of $C_{60}@H_DA\text{-MOF}$ are lower by ~ 0.2 - 0.4 eV compared to $H_DA\text{-MOF}$, while $C_{60}@H_F\text{-MOF}$ has band gaps lower by 0.63 eV and 0.77 eV in Type-A and Type-B, respectively, compared to $H_F\text{-MOF}$. While the excitations in $H_DA\text{-}$ and $H_F\text{-MOFs}$ are localized in nature and assignable as intramolecular π - π^* transitions, their C_{60} encapsulated analogues access delocalized electronic states, where the HOCOs and LUCOs comprise p-orbitals of the linkers and C_{60} , respectively.

Table 1. Computed band gaps (E_g , eV) and binding energies (ΔE , eV) for DA-MOF, F-MOF, and their analogues after C_{60} incorporation.

MOFs	E_g for Free MOF	E_g (and ΔE) for $C_{60}@MOF$		
		Type-I	Type-II	Type-III
DA-MOF	1.66	1.50 (-1.23)		
H_DA-MOF	1.69	1.49 (-0.91)	1.54 (-1.16)	1.29 (-0.89)
NH ₂ _DA-MOF	1.30	0.85 (-1.06)	1.10 (-1.31)	0.71 (-0.94)
		Type-A		Type-B
F_MOF	2.04	1.68 (-1.33)		
H_F-MOF	2.34	1.71 (-1.37)		1.57 (-1.11)
NH ₂ _F-MOF	1.66	1.03 (-1.45)		0.67 (-1.36)

The photophysical properties of porphyrin are known to be sensitive to substitution of the meso position with electron-donating or accepting groups.⁵⁹ We considered the strong electron donor -NH₂ as such a group and computed the structures of NH₂_DA-MOF and NH₂_F-MOF, respectively (Figure S4) as well as $C_{60}@NH_2_DA-MOF$ and $C_{60}@NH_2_F-MOF$, both of which were found to have the same number of polymorphs as their H-MOF analogs (Figure S5 and S6). The calculated C_{60} binding energies for NH₂_DA-MOF are -1.06 eV, -1.31 eV, and -0.94 eV for the Type-I, Type-II, and Type-III polymorphs, respectively, while for NH₂_F-MOF they are -1.45 eV and -1.36 eV for Type-A and Type-B, respectively. More complete π - π overlap in both Type-II and Type-A interfaces again leads to the highest binding energies. Compared to H_DA-MOF and H_F-MOF, the C_{60} binding energies are significantly higher in the NH₂-substituted cases (Table1), consistent with increased donor-acceptor interactions between the linkers and C_{60} that have the potential to impart more interesting optoelectronic properties.

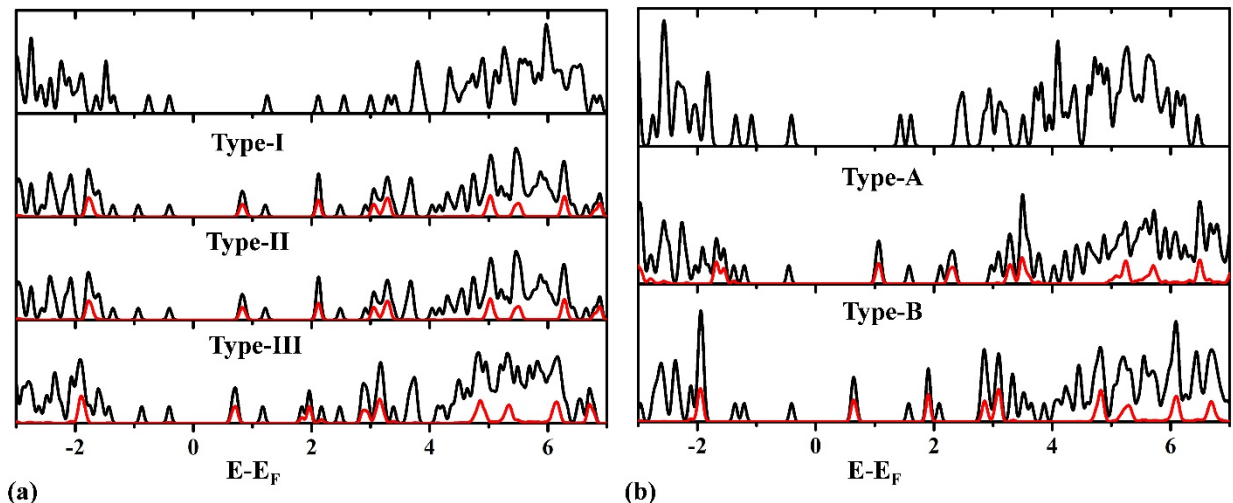
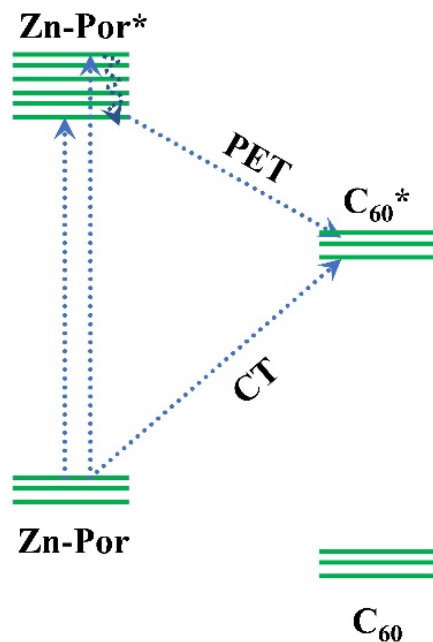


Figure 4. Calculated TDOS for (a) NH₂_DA and (b) NH₂_F-MOF are shown in upper panels (contribution from *s*- and *p*-orbitals and -NH₂ are given in Figure S7). In other panels TDOS (black line) and PDOS of C₆₀ (red line) for different polymorphs are shown. The energy is relative to the Fermi energy.

Amine functionalization affects the electronic properties of the NH₂_MOFs (Figures 4, S7). The band gaps decrease by ~ 0.39 eV and ~ 0.68 eV, respectively, compared to H_DA-MOF and H_F-MOF (Table 1). PDOS analysis indicates that low-energy optical transitions are intralinker $\pi-\pi^*$ absorptions (Figure S7). The band gaps of the Type-I, Type-II, and Type-III polymorphs of C₆₀@NH₂_DA-MOF are 0.85, 1.10 and 0.71 eV, respectively, which values are lower than that for NH₂_DA-MOF by 0.45, 0.21 and 0.59 eV, respectively. Similarly, the band gaps of Type-A and Type-B (1.03 eV and 0.67 eV, respectively) are lower (by 0.63 and 0.99 eV, respectively) than that of NH₂_F-MOF. The incorporation of C₆₀ in NH₂_DA-MOF and NH₂_F-MOF thus leads to infrared-bandgap tuned materials, again with delocalized electronic states that should permit efficient charge transfer across interfacial regions leading to improved electrical conductivity.



Scheme 1. Schematic representation for direct charge transfer (CT) and photoinduced electron transfer (PET) processes between porphyrin-based donor and C₆₀ based acceptor.

Our solid-state electronic structures indicate that the empty states of C₆₀ are positioned in between the band extrema of the empty MOFs. In addition, several donor-acceptor conjugates have geometries consistent with strong π - π interactions, which should facilitate charge transfer (CT) and photoinduced electron transfer (PET) processes from porphyrin to C₆₀ (Scheme 1), which phenomena are well known for analogous molecular and supramolecular architectures.⁶⁰⁻⁶¹ To obtain further insight into the localized character of these electronic processes, we performed time dependent DFT (TD-DFT) calculations on cluster models of bare linkers and C₆₀ extracted from our systems. Each C₆₀ is in close proximity to four porphyrin units within the framework and we have decomposed the system into corresponding fragments accordingly (see Section S2 of SI). We restrict discussion here to results from C₆₀@NH₂_DA-MOF Type-III and C₆₀@NH₂_F-MOF Type-B, as minimum band gaps are predicted for these systems; details for other systems may be found in the SI.

The cluster model of C₆₀@NH₂_DA-MOF Type-III may be split into two fragments, namely, stacked, and diagonal, while for C₆₀@NH₂_F-MOF Type-B, three fragments are considered, namely, stacked, diagonal and parallel. Figure 5 shows that in the stacked fragment of C₆₀@NH₂_DA-MOF Type-III, the highest-energy occupied frontier orbitals are mainly

localized on the porphyrin while the lowest-energy unoccupied frontier orbitals are localized on both C₆₀ and the porphyrin. The nature of the excited states, their associated vertical excitation energies (E_{0-n}), oscillator strengths (f), and the molecular orbitals (MOs) computed to dominate the transition are reported in Table 2. In the stacked fragment, low energy near IR transitions such as S₀→S₁ ($E_{0-1}= 0.62$ eV; $f=0.0003$), S₀→S₂ ($E_{0-2}= 0.67$ eV; $f=0.0047$), S₀→S₃ ($E_{0-3}= 0.69$ eV; $f=0.0054$), and S₀→S₆ ($E_{0-6}= 0.83$ eV; $f=0.0312$) occur from HOMO and HOMO-1 localized over the porphyrin to LUMO, LUMO+1, and LUMO+2 localized on C₆₀, indicating the direct CT nature of these excitations. This is consistent with the periodic calculations that predict the band gap for C₆₀@NH₂_DA-MOF Type-III to involve a porphyrin-based HOCO and a C₆₀-based LUCO. In contrast, higher energy intense transitions, e.g., S₀→S₉ ($E_{0-9}= 1.59$ eV; $f=0.0524$), S₀→S₁₀ ($E_{0-10}= 1.62$ eV; $f=1.0257$), and S₀→S₉₂ ($E_{0-92}= 2.98$ eV; $f=0.3585$) involve visible-light excitation and correspond to intramolecular $\pi \rightarrow \pi^*$ transitions within the porphyrin moiety (Table 2 and Figure 5). These higher-energy states should rapidly internally convert to the lowest-lying porphyrin-based excited state (Kasha's rule, Scheme1), and subsequent electron transfer to the LUMO of C₆₀ would constitute a PET process facilitated by suitable spatial and geometric overlap in the donor-acceptor conjugate. A similar situation is found for the diagonal fragment of C₆₀@NH₂_DA-MOF Type-III, with low-lying absorptions ascribed as direct CT excitations and higher energy absorptions associated with $\pi \rightarrow \pi^*$ transitions in the porphyrin moieties (Table S11 and Figure S24).

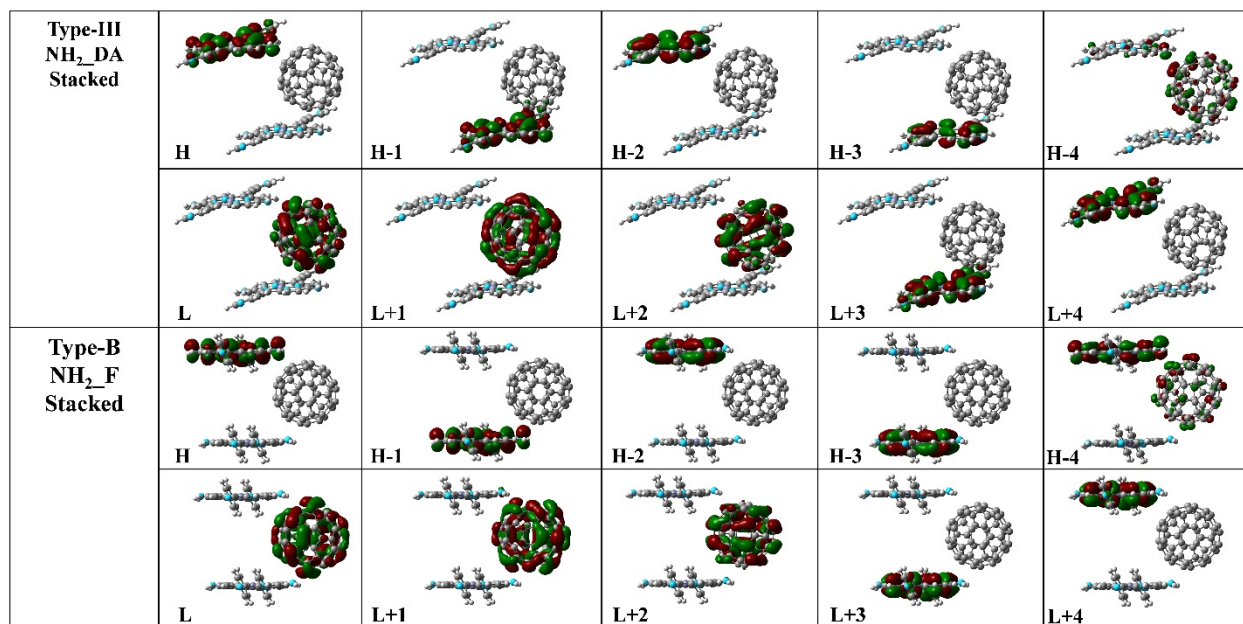


Figure 5. Frontier molecular orbitals for stacked fragments extracted from Type-III and Type-B polymorphs of C₆₀@NH₂_DA-MOF and C₆₀@NH₂_F-MOF, respectively (H = HOMO, L = LUMO). See caption to Figure 1 for atomic colors.

For C₆₀@NH₂_F-MOF Type-B, excitations are found to be similar to those for C₆₀@NH₂_DA-MOF Type-III. In the stacked fragment, lower energy transitions S₀→S₂ ($E_{0-2}=0.55$ eV; $f=0.0034$), S₀→S₃ ($E_{0-3}=0.59$ eV; $f=0.0263$), S₀→S₄ ($E_{0-4}=0.60$ eV; $f=0.0152$), S₀→S₆ ($E_{0-6}=0.66$ eV; $f=0.0288$) occur from porphyrin localized HOMO, HOMO-1 and HOMO-2 to C₆₀ based LUMO, LUMO+1 and LUMO+2 (Table 2 and Figure 5), while higher energy transitions S₀→S₂₁ ($E_{0-21}=1.91$ eV; $f=0.0852$), S₀→S₂₂ ($E_{0-22}=1.93$ eV; $f=0.2893$), and S₀→S₃₄ ($E_{0-34}=2.19$ eV; $f=0.0288$) are molecular in nature on the porphyrin but should permit electron flow from $\pi \rightarrow \pi^* \rightarrow \text{C}_{60}\text{-LUMO}$. The diagonal and parallel fragments of C₆₀@NH₂_F-MOF Type-B are also similar (Table S13 and Figure S26). We investigated analogous electronic excitations for parent C₆₀@DA- and C₆₀@F-MOFs, as well as for all the other polymorphs of C₆₀@H_ and C₆₀@NH₂_DA- and F-MOFs (see Section S2 of SI).

Interestingly, edge-on interactions between porphyrins and C₆₀ in Type-I interfaces found in both H_DA- and NH₂_DA-MOF fail to permit direct CT, although PET is energetically allowed (Figures S17 and S22; Tables S4 and S9), which emphasizes the importance of the

crystallographic position of C₆₀ in the MOF pore with respect to controlling the electronic processes taking place at the host-guest interfaces. Geometry notwithstanding, molecular and supramolecular assemble of porphyrin and fullerene are well known to engender long-lived radical cations and radical anions following photolysis.^{60, 62} For instance, Chen et. al. recently reported that C₆₀ spatially confined via covalent anchoring within the nanochannels of a porphyrin based covalent organic framework ([C₆₀]y-ZnPc-COFs, where y = acceptor content) promotes photoinduced electron transfer and allows charge separation to form ZnPc^{•+} and C₆₀^{•-} species, and the resulting COF shows photoelectric conductivity.³⁴ Similarly, the C₆₀@DA-MOF and C₆₀@F-MOFs and their analogues reported here could allow the formation of ionized species that could contribute to photoelectric conductivity.

Table 2. Excited State, excitation energies (eV), wavelength (nm), oscillator strengths (*f*) and important orbitals involved in transition for Type-III and Type-B interfaces of NH₂_DA-MOF and NH₂_F-MOF, respectively, in their stacked orientation.

Systems with C ₆₀	Excited State (Ex. St.)	Excitation energies (E _{0-n} , eV)	Wavelength (nm)	Oscillator strengths (<i>f</i>)	Important orbitals involved
Type-III NH ₂ -DA Stack	1	0.62	1999.41	0.0003	H→L
	2	0.67	1849.33	0.0047	H→L+1; H→L+2
	3	0.69	1799.39	0.0054	H→L+1; H→L+2
	6	0.83	1490.67	0.0312	H-1→L; H-1→L+1; H-1→L+2
	9	1.59	781.84	0.0524	H-1→L+3; H→L+4
	10	1.62	763.43	1.0257	H-2→L+9; H-1→L+3; H→L+4
	20	1.93	642.35	0.0315	H-3→L+1; H-3→L+3; H-1→L+5; H-1→L+6
	66	2.75	451.12	0.0153	H-4→L+4
	81	2.93	423.77	0.3170	H-3→L+3; H-2→L+7; H-1→L+8

	92	2.98	416.36	0.3585	H-9→L+4; H-2→L+4; H→L+12; H→L+13
Type-B NH ₂ -F Stacked	1	0.51	2440.32	0.0001	H→L; H→L+1
	2	0.55	2274.70	0.0034	H→L; H→L+1; H→L+2
	3	0.59	2099.13	0.0263	H-1→L; H→L+1; H→L+2
	4	0.60	2081.01	0.0152	H-1→L; H-1→L+1; H→L+1; H→L+2
	6	0.66	1865.26	0.0288	H-1→L; H-1→L+1; H- 1→L+2
	11	1.58	784.30	0.0067	H-3→L+1
	21	1.91	648.46	0.0852	H-2→L+9; H-1→L+3; H→L+4
	22	1.93	642.15	0.2893	H-1→L+3; H-1→L+6; H→L+4
	34	2.19	565.10	0.0288	H-4→L+1; H-4→L+2; H→L+9
	99	3.00	412.66	0.0108	H→L+15; H→L+16

Charge carrier and exciton migrations play a crucial role in controlling electron conduction in optoelectronic devices.⁶³⁻⁶⁴ Carrier mobilities across molecular interfaces can be assessed from Marcus theory (Eq. S1 in SI) and depend on the electronic transfer integral (V) between the adjacent molecules as well as the reorganization energy (λ). Transfer integrals are extremely sensitive to the relative orientation of the associated molecules and the direction of transport. Therefore, we have considered different possible hopping pathways between adjacent molecular fragments and computed associated transfer integrals. Analyzing center-to-center distances between the nearest neighbor porphyrinic units in DA-MOF and F-MOF (and also in their analogues, but all without C₆₀ incorporation), we have considered four possible pathways

namely, AB, AC, AD and AE for exciton migrations (Figure S27 in SI). Computed hole (V_h) and electron (V_e) transfer integrals are negligible in all of the pathways for both MOFs (Table S14 in SI), suggesting the electronic transport between the porphyrinic units is ineffective for conduction in the absence of a dopant of some sort.

Table3. Computed electron-transfer (V_e , cm^{-1}) and hole-transfer (V_h , cm^{-1}) integrals between porphyrin-based linkers and C_{60} along different pathways for $\text{C}_{60}@\text{DA-MOF}$, $\text{C}_{60}@\text{F-MOF}$, and their analogues.

C_{60} incorporated Systems	Interface	Pathways			
		AF V_e (V_h)	BF V_e (V_h)	CF V_e (V_h)	DF V_e (V_h)
DA-MOF	Parent	0.32 (1.05)	21.78 (82.11)	8.71 (8.79)	9.52 (24.12)
H_DA-MOF	Type-I	0.65 (11.53)	2.42 (11.13)	2.98 (1.45)	13.79 (4.60)
	Type-II	142.92 (774.29)	695.17 (306.41)	0.00 (0.00)	0.00 (0.00)
	Type-III	392.79 (687.51)	173.97 (216.40)	1.29 (0.48)	111.39 (58.39)
$\text{NH}_2\text{-DA-MOF}$	Type-I	0.65 (11.53)	2.42 (11.13)	2.98 (1.45)	13.79 (4.60)
	Type-II	197.44 (992.95)	669.36 (460.62)	0.00 (0.08)	0.00 (0.00)
	Type-III	207.61 (610.72)	228.66 (241.64)	0.48 (0.73)	116.55 (627.10)
F-MOF	Parent	487.32 (371.26)	534.75 (383.60)	0.08 (0.00)	0.08 (0.00)
H_F-MOF	Type-A	0.48 (109.93)	591.29	0.00 (0.00)	0.00 (0.00)

			(297.94)		
	Type-B	40.33 (72.75)	32.10 (5.40)	212.53 (129.05)	116.71 (194.06)
NH ₂ _F-MOF	Type-A	113.24 (100.74)	406.50 (427.55)	0.00 (0.00)	0.00 (0.00)
	Type-B	17.10 (0.32)	13.79 (31.70)	84.12 (143.89)	190.18 (410.21)

We next computed transfer integrals between donor-acceptor conjugates in the various C₆₀@MOFs considered here (Figures S28-S32 in SI) addressing four pathways, namely, AF, BF, CF and DF. Computed coupling values (V) are anisotropic in nature and highly dependent on pathway (Table 3). For C₆₀@DA-MOF, the maximum V_h and V_e values are predicted to be 82.1 and 21.8 cm⁻¹, respectively, along the BF pathway. For Type-I interfaces, where C₆₀ has a similar orientation as in parent DA-MOF, we compute still lower V values, as the removal of the phenyl rings from the meso positions further decreases the $\pi\cdots\pi$ overlap between them (Figures S28, S29 and S31). By contrast, these values increase significantly for the Type-II interfaces of both H_DA-MOF (V_h= 774.3 cm⁻¹ and V_e= 695.2 cm⁻¹) and NH₂_DA-MOF (V_h= 993.0 cm⁻¹ and V_e= 669.4 cm⁻¹) due to optimal $\pi\cdots\pi$ stacking overlap between the porphyrin and C₆₀ (Table 3). Relative to these highest values, the predicted integrals in Type-III interfaces decrease somewhat (V_h= 687.5 cm⁻¹ and V_e= 392.8 cm⁻¹ for H_DA-MOF and V_h= 627.1 cm⁻¹ and V_e= 228.7 cm⁻¹ for NH₂_DA-MOF) as might be anticipated given the the slip-stacked arrangement of C₆₀ relative to the porphyrin (Figure S29 and S31). Interestingly, though, Type-III interfaces *do* provide an additional mechanism, along the DF pathway, for charge transport that is not available in the other cases.

Similarly, for C₆₀@F-MOF as well as for the Type-A polymorphs of its analogues, the largest transfer integrals are associated with ideal, eclipsed $\pi\cdots\pi$ stacking interaction (Table 3, Figure S28, S30 and S32). The coupling values decrease considerably for both C₆₀@H_F-MOF and C₆₀@NH₂_F-MOF Type-B polymorphs where the slip-stacked geometry is less optimal.

It is apparent, then, that computed transfer integrals increase significantly upon C₆₀ incorporation in all of the MOFs considered here (cf. Table S15 in SI). Importantly, C₆₀ is known to have very small reorganization energies in electron transfer reactions, which is essential for the generation of long-lived charge-separated states.⁶¹ Taking this together with our solid-state calculations, molecular TD-DFT calculations, and computed electronic coupling values, we may posit that electrical transport will be high in C₆₀ encapsulated porphyrin-based MOFs and that these will be promising materials for optoelectronics, particularly when linker modifications are chosen to enhance the likelihood of eclipsed π - π stacking interactions between the fullerene and the porphyrin.

Conclusions

The incorporation of C₆₀ into the pores of porphyrin-containing DA-MOF and F-MOF analogs shows great promise for engendering massive electrical conductivity through π -mediated donor-acceptor interactions, as confirmed through calculation of electronic transfer integrals for materials with and without C₆₀. Substitution at the meso positions of the porphyrin can lead to C₆₀ occupying different positions within the framework, and calculations indicate that these polymorphic heterostructures possess anisotropic structural, electronic, opto-electronic, and transport properties that depend on variable donor-acceptor interactions. TD-DFT calculations also indicate that electrons can be transferred from porphyrin to C₆₀ either directly through near-infrared charge-transfer transitions or as well by photoinduced electron transfer processes upon visible light excitation. This organization could prolong the lifetime of charge separated states at interfaces and substantially improve the performance of the doped MOFs for optoelectronic applications. Our detailed insights into the structural, electronic, and photophysical properties of porphyrin-based MOFs incorporating C₆₀ as guest will hopefully further stimulate the rational design of improved MOF-derived photoelectrochemical devices.

Supporting Information:

Lattice parameters, electronic structures, cluster modelling and excited states calculations results, hopping pathways and other related data, full references of Gaussian 16 and ADF.

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