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# Model realization and model reduction for quantum systems

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- **Model realization and system identification**

- Estimating unknown Hamiltonian parameters

[arXiv:1401.5780](#) [[pdf](#), [other](#)]

**Quantum Hamiltonian identification from measurement time traces**

[Jun Zhang](#), [Mohan Sarovar](#)

Comments: 6 pages, 2 figures

- **Model reduction**

- Reducing simulation cost for certain many-body quantum systems
  - Preliminary results

# System Identification

Identify system from input-output behavior



*e.g.* Process tomography: identify process (CP-map, unitary) at a particular time

Alternative: identify generator/Hamiltonian of system

# System Identification

How powerful are time traces?

Takens' theorem



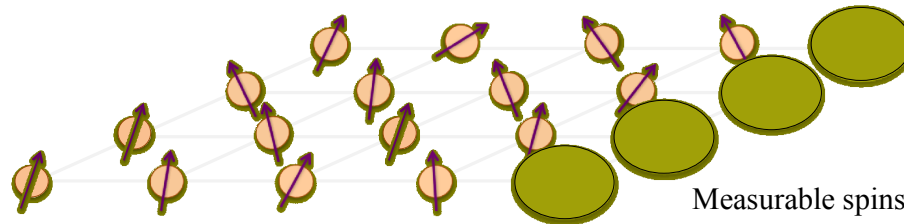
Additional desiderata

- Measurements could be restricted
- May have partial information about system

**Assumptions:**

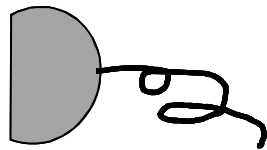
1. system is finite dimensional
2. Hamiltonian dynamics (closed system)

# An example



Parametric Hamiltonian

$$H(\theta_1, \theta_2, \dots, \theta_M)$$



$$\langle \sigma_z^1(t_0) \rangle, \langle \sigma_z^1(t_1) \rangle, \dots, \langle \sigma_z^1(t_n) \rangle$$

Time trace of some  
accessible observable

Can we back out the parameters in the Hamiltonian from just this?

# The setup

Choose an orthogonal operator basis for the linear operator space (e.g. generalized Paulis)

$$[iX_j, iX_k] = \sum_{l=1}^{N^2-1} C_{jkl}(iX_l), \quad j, k = 1, \dots, N^2 - 1,$$

Hamiltonian can be expanded in this basis

$$H = \sum_{m=1}^M a_m(\theta) X_m,$$

**Goal:** to identify  $a_m$

Leads to a linear, autonomous equation for state  $x(t)$

$$\frac{d}{dt} x_k = \sum_{l=1}^{N^2-1} \left( \sum_{m=1}^M C_{mkl} a_m \right) x_l.$$

$$\frac{d}{dt} \mathbf{x} = \mathbf{A} \mathbf{x}, \quad x_k(0) = \langle \psi(0) | X_k | \psi(0) \rangle,$$

$$|\psi\rangle \in \mathbb{C}^N$$

$$\dim H = N \times N$$

$$\mathbf{x} \in \mathbb{R}^{(N^2-1)}$$

$$\dim A = (N^2 - 1) \times (N^2 - 1)$$

# The setup

Similarly, each directly measured observable can be expanded in the same basis. Resulting in a AC system

$$\frac{d}{dt}\mathbf{x} = \mathbf{A}\mathbf{x}$$

$$\mathbf{y}(t) = \mathbf{C}\mathbf{x}(t)$$

But this may be too complex a description. E.g.

$$\mathbf{C} = \begin{bmatrix} 1 & 0 & 0 & 0 \end{bmatrix}$$

$$\mathbf{A} = \begin{bmatrix} a & b & 0 & 0 \\ c & d & 0 & 0 \\ 0 & 0 & a' & b' \\ 0 & 0 & c' & d' \end{bmatrix}$$

# Filtration to find minimal description

$$O_i = \sum_j o_j^{(i)} X_j \quad \mathcal{M} = \{X_{\nu_1}, X_{\nu_2}, \dots, X_{\nu_p}\}$$

$$H = \sum_{m=1}^M a_m(\theta) X_m, \quad \Delta = \{X_m\}_{m=1}^M$$

Filtration recursively constructed as:

$$G_0 = \mathcal{M}, \text{ and}$$

$$G_i = [G_{i-1}, \Delta] \cup G_{i-1}$$

where

$$[G_{i-1}, \Delta] \equiv \{\langle [g, h] \rangle : g \in G_{i-1}, h \in \Delta\}$$

Results in minimal description

$$\frac{d}{dt} \mathbf{x}_a = \tilde{\mathbf{A}} \mathbf{x}_a$$



# Discretization

$$\mathbf{x}_a(j+1) = \tilde{\mathbf{A}}_d \mathbf{x}_a(j)$$

$$\tilde{\mathbf{A}}_d = e^{\tilde{\mathbf{A}}\Delta t}$$

$$\mathbf{y}(j) = \tilde{\mathbf{C}} \mathbf{x}_a(j)$$

Explicit solution

$$\mathbf{y}(j) = \tilde{\mathbf{C}} \tilde{\mathbf{A}}_d^j \mathbf{x}_a(0)$$

**Goal:**

Use  $\{\mathbf{y}(j)\}_{j=0}^J$  to estimate  $\{a_m\}_{m=1}^M$

Strategy:

1. Find the minimal linear model that generates the collected data
2. Back out the unknown parameters from this model

# Eigenstate realization algorithm

Step 1: Form Hankel matrix from data

$$\mathbf{H}_{rs}(k) = \begin{bmatrix} \mathbf{y}(k) & \mathbf{y}(k+t_1) & \cdots & \mathbf{y}(k+t_{s-1}) \\ \mathbf{y}(j_1+k) & \mathbf{y}(j_1+k+t_1) & \cdots & \mathbf{y}(j_1+k+t_{s-1}) \\ \vdots & \vdots & & \vdots \\ \mathbf{y}(j_{r-1}+k) & \mathbf{y}(j_{r-1}+k+t_1) & \cdots & \mathbf{y}(j_{r-1}+k+t_{s-1}) \end{bmatrix}$$

Step 2: Take SVD of Hankel matrix at  $k=0$

$$\mathbf{H}_{rs}(0) = P \begin{bmatrix} \Sigma & 0 \\ 0 & 0 \end{bmatrix} Q^\top = [P_1 \ P_2] \begin{bmatrix} \Sigma & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} Q_1^\top \\ Q_2^\top \end{bmatrix},$$

Step 3: Form realizations of linear model from SVD components

$$\hat{\mathbf{A}}_d = \Sigma^{-\frac{1}{2}} P_1^\top \mathbf{H}_{rs}(1) Q_1 \Sigma^{-\frac{1}{2}}, \quad \hat{\mathbf{C}} = \mathbf{E}_p^\top P_1 \Sigma^{\frac{1}{2}},$$

$$\hat{\mathbf{x}}(0) \equiv \Sigma^{\frac{1}{2}} Q_1^\top \mathbf{e}_1,$$

# The Hankel matrix

Interpret Hankel matrix

# Realization to parameter estimation

The triple  $(\hat{\mathbf{A}}_d, \hat{\mathbf{C}}, \hat{\mathbf{x}}(0))$  is a realization of the triple  $(\tilde{\mathbf{A}}_d, \tilde{\mathbf{C}}, \mathbf{x}_a(0))$

$$\mathbf{y}(j) = \mathbf{C}\tilde{\mathbf{A}}_d^j \mathbf{x}_a(0) = \hat{\mathbf{C}}\hat{\mathbf{A}}_d^j \hat{\mathbf{x}}(0), \quad \text{for all } j \geq 0,$$

Define  $\hat{\mathbf{A}} = \log \hat{\mathbf{A}}_d / \Delta t$

Expanding the exponential in a power series and equating terms,

$$\mathbf{C}\tilde{\mathbf{A}}^j \mathbf{x}_a(0) = \hat{\mathbf{C}}\hat{\mathbf{A}}^j \hat{\mathbf{x}}(0), \quad \text{for all } j \geq 0$$

*Determined by the data*

*Polynomial equation in unknown parameters*

# Realization to parameter estimation

$$\mathbf{C}\tilde{\mathbf{A}}^j \mathbf{x}_a(0) = \hat{\mathbf{C}}\hat{\mathbf{A}}^j \hat{\mathbf{x}}(0), \quad \text{for all } j \geq 0$$

*Determined by the data*

*Polynomial equation in unknown parameters*

Solving these equations yields estimates of parameters

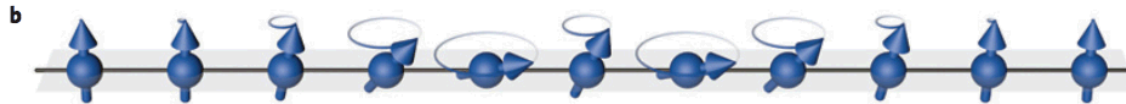
Notes:

1. Parameter estimates can be non-unique (gauge freedom/symmetries)
2.  $\Delta t$  must be small enough

# Example

XY spin chain

$$H = \sum_{k=1}^n \frac{\omega_k}{2} \sigma_z^k + \sum_{k=1}^{n-1} \delta_k (\sigma_+^k \sigma_-^{k+1} + \sigma_-^k \sigma_+^{k+1}).$$



Fukuhara et al. Nature  
Physics, **9** 235 (2013)

$$\tilde{\mathbf{A}} = \begin{bmatrix} 0 & \omega_1 & 0 & -\delta_1 & & & & \\ -\omega_1 & 0 & \delta_1 & 0 & 0 & & & \\ 0 & -\delta_1 & 0 & \omega_2 & 0 & \ddots & & \\ \delta_1 & 0 & -\omega_2 & 0 & \ddots & \ddots & 0 & \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots & \ddots & -\delta_{n-1} \\ \ddots & \ddots & \ddots & 0 & \delta_{n-1} & 0 & & \\ & 0 & 0 & -\delta_{n-1} & 0 & \omega_n & & \\ & & \delta_{n-1} & 0 & -\omega_n & 0 & & \end{bmatrix}$$

# Example

System of equations

$$\begin{aligned}\omega_1 &= \hat{\mathbf{C}}_2 \hat{\mathbf{A}}_2 \hat{\mathbf{x}}_2(0) \\ \omega_1^2 + \delta_1^2 &= -\hat{\mathbf{C}}_1 \hat{\mathbf{A}}_1^2 \hat{\mathbf{x}}_1(0) \\ \omega_1^3 + \delta_1^2(2\omega_1 + \omega_2) &= -\hat{\mathbf{C}}_2 \hat{\mathbf{A}}_2^3 \hat{\mathbf{x}}_2(0) \\ \omega_1^4 + \delta_1^2(3\omega_1^2 + 2\omega_1\omega_2 + \omega_2^2 + \delta_1^2 + \delta_2^2) &= \hat{\mathbf{C}}_1 \hat{\mathbf{A}}_1^4 \hat{\mathbf{x}}_1(0).\end{aligned}$$

Coupling parameters only occur up to even order (symmetry) => can only be determined up to sign

# Summary

System identification through model realization

Most useful when

- measurements are restricted
- Prior information about process is available

Future work:

- Noisy measurements
- Markovian open-system evolution



- **Model realization and system identification**
  - Estimating unknown Hamiltonian parameters
- **Model reduction**
  - Reducing simulation cost for certain many-body quantum systems
    - Preliminary results

## Quantum state space: exponential

- Full-scale simulation of quantum systems very difficult
  - “Hilbert space is a big place” – Carlton Caves
- Formal state is exponentially large in the number of particles

$$\rho_1 \in \mathcal{H}_1 \qquad \dim \mathcal{H}_1 = n_1$$

$$\rho_2 \in \mathcal{H}_2 \qquad \dim \mathcal{H}_2 = n_2$$

$$\rho_c \in \mathcal{H}_1 \otimes \mathcal{H}_2 \qquad \dim \mathcal{H}_1 \otimes \mathcal{H}_2 = n_1 n_2 \neq n_1 + n_2$$

## Quantum state space: NOT exponential

- However, for most practical systems, this exponential scaling is only formal

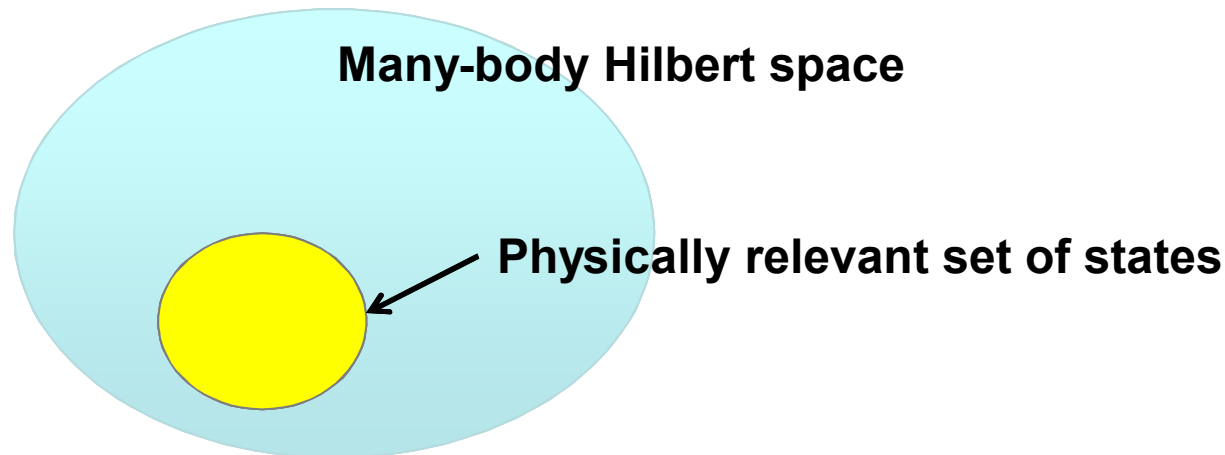
PRL **106**, 170501 (2011)

PHYSICAL REVIEW LETTERS

week ending  
29 APRIL 2011

### Quantum Simulation of Time-Dependent Hamiltonians and the Convenient Illusion of Hilbert Space

David Poulin,<sup>1</sup> Angie Qarry,<sup>2,3</sup> Rolando Somma,<sup>4</sup> and Frank Verstraete<sup>2</sup>

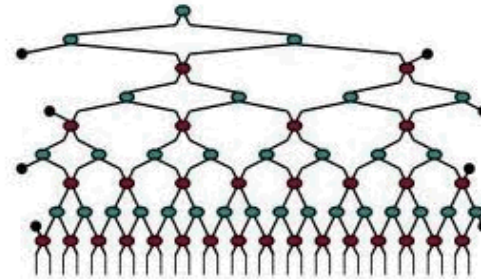
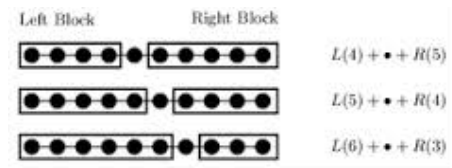


- Identifying this set of relevant states is difficult

# Identifying reduced order models

## Techniques in physics:

- static: DMRG, MPS, etc.



- dynamic: Nakajima-Zwanzig (statistical), Bloch equations, Glauber dynamics

$$\begin{aligned} \frac{d}{dt} \mathbf{P}x(t) = & \mathbf{PAP}x(t) + \mathbf{PB}u(t) \\ & + \mathbf{PAG}(t, 0)\mathbf{Q}x(0) \\ & + \int_0^t \mathbf{PAG}(t, s)\mathbf{QAP}x(s)ds \\ & + \int_0^t \mathbf{PAG}(t, s)\mathbf{QB}u(s)ds. \end{aligned}$$

# Model reduction

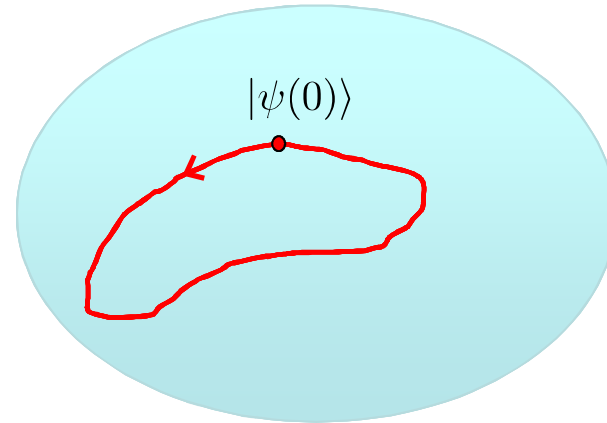
<div>Resources available</div> <div>Desired Output</div>	State snapshots	Input-output map	Dynamical model
Full state vector want to reproduce $ \psi(t)\rangle$	Proper orthogonal decomposition (POD)	-	Identify invariant subspaces
Input-output map want to reproduce $y(t)$	Empirical balanced truncation (BPOD)	Minimal model realization algorithms	Balanced truncation

# Compressible dynamics

$$|\psi(t)\rangle = e^{iHt} |\psi(0)\rangle$$

$$H = \sum_i \lambda_i h_i$$

e.g. adiabatic QC



## Problem:

Identify subspace of Hilbert space that contains  $|\psi(0)\rangle$   
and is invariant under Hamiltonian for all choices of  $\lambda$

Projective model reduction: columns of  $P$  are basis vectors in this invariant subspace

$$\frac{d}{dt} P |\psi(t)\rangle = P^\dagger H P |\psi(t)\rangle$$

$$\dim P = N \times q, \quad q \ll N$$

$q \times q$  compressed description

# E.g. Quench dynamics

## Quantum Ising model



Fukuhara et al. Nature Physics, **9** 235 (2013)

$$H = -B \sum_i \sigma_x^i - J \sum_{\langle i,j \rangle} \sigma_z^i \sigma_z^j$$

- Basic model for magnetism in crystalline material
- Competition between B and J results in phase transition behavior
- Can be emulated using cold atoms
- As a result: intense interest in dynamical phase transitions, quenching dynamics

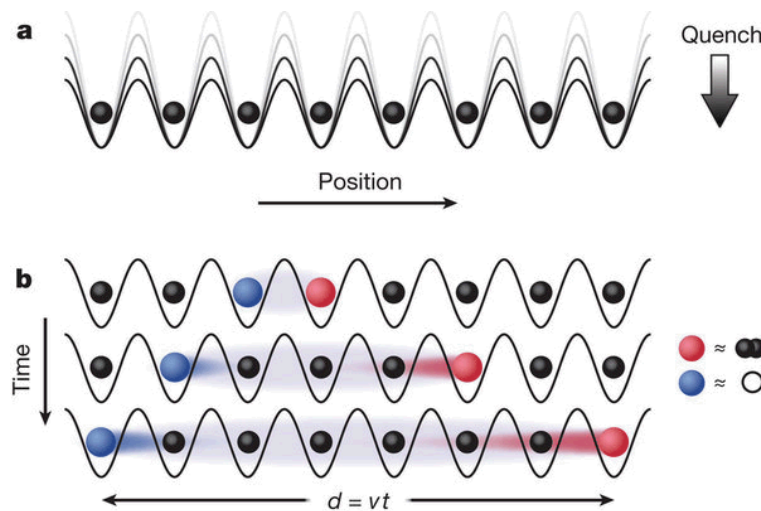
### Quenching dynamics:

1. Prepare ground state of  $H^0 = -B^0 \sum_i \sigma_x^i - J \sum_{\langle i,j \rangle} \sigma_z^i \sigma_z^j$
1. Rapidly change B and evolve system under  $H^1 = -B^1 \sum_i \sigma_x^i - J \sum_{\langle i,j \rangle} \sigma_z^i \sigma_z^j$
1. The resulting dynamics is very informative; e.g. contains information about static phases of system

# E.g. Quench dynamics

Can identify many important features of many-body model by looking at dynamics after quench

Especially relevant now with cold-atom quantum simulators that are capable of quenched dynamics



Cheneau et al. Nature,  
**481** 484 (2012)

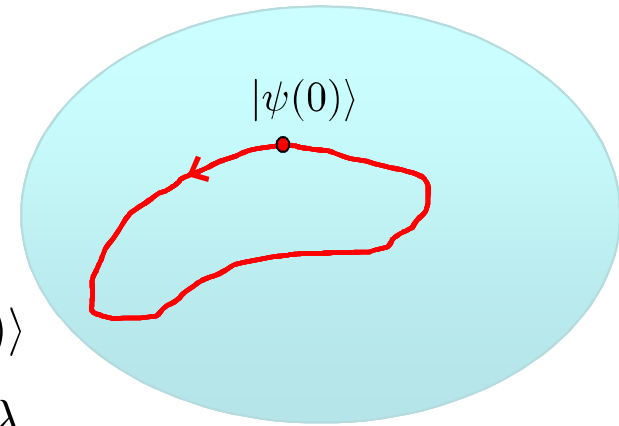


# Compressible dynamics

$$H = \sum_i \lambda_i h_i$$

## Problem:

Identify subspace of Hilbert space that contains  $|\psi(0)\rangle$  and is invariant under Hamiltonian for all choices of  $\lambda$



1. Certificates
  1. Is this dynamics compressible?
2. Computing reduced order models
  1. What is the invariant subspace and moreover, what is the compressed dynamical model?

$$H = \sum_i \lambda_i h_i$$

$$H \in L(\mathcal{H}) \quad \dim \mathcal{H} = N$$

$$\text{Coeff}(H) \equiv \{h_i\}$$

Theorem: (algebraic certificate)

The Hamiltonian acting on  $\mathcal{H}$  keeps invariant a non-trivial proper subspace iff the subalgebra generated by  $\text{Coeff}(H)$  is a proper subalgebra of  $L(\mathcal{H})$ .

Intuition:

$$|\psi(t)\rangle = \exp\{i(\lambda_1 h_1 + \lambda_2 h_2)t\} |\psi(0)\rangle$$

$$|\psi(t)\rangle = \sum_n \frac{(it)^n}{n!} (\lambda_1 h_1 + \lambda_2 h_2)^n |\psi(0)\rangle$$

Products of  $h_i$  generate an algebra. If the full operator algebra is not generated, there are directions not explored in state space

# Certificate

## Special case: Pauli Hamiltonian

$$H = \sum_i \lambda_i \sigma_i$$

$$H \in L(\mathcal{H}) \quad \mathcal{H} = \mathbb{C}^{2^n}$$

$$\sigma_i : \sigma_x^{(1)} \otimes \mathbf{1} \otimes \dots \otimes \sigma_y^{(n)}$$

$$\text{Coeff}(H) \equiv \{h_i\} \quad \dim \mathcal{H} = 2^n$$

Theorem: (Pauli algebraic certificate)

Any Pauli Hamiltonian acting on  $n$  qubits with fewer than  $2n$  terms has a non-trivial proper invariant subspace.

e.g. Quantum (transverse field) Ising model with random couplings and energies

# Example: quench dynamics

## Quantum Ising model



Fukuhara et al. Nature Physics, **9** 235 (2013)

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### Quenching dynamics:

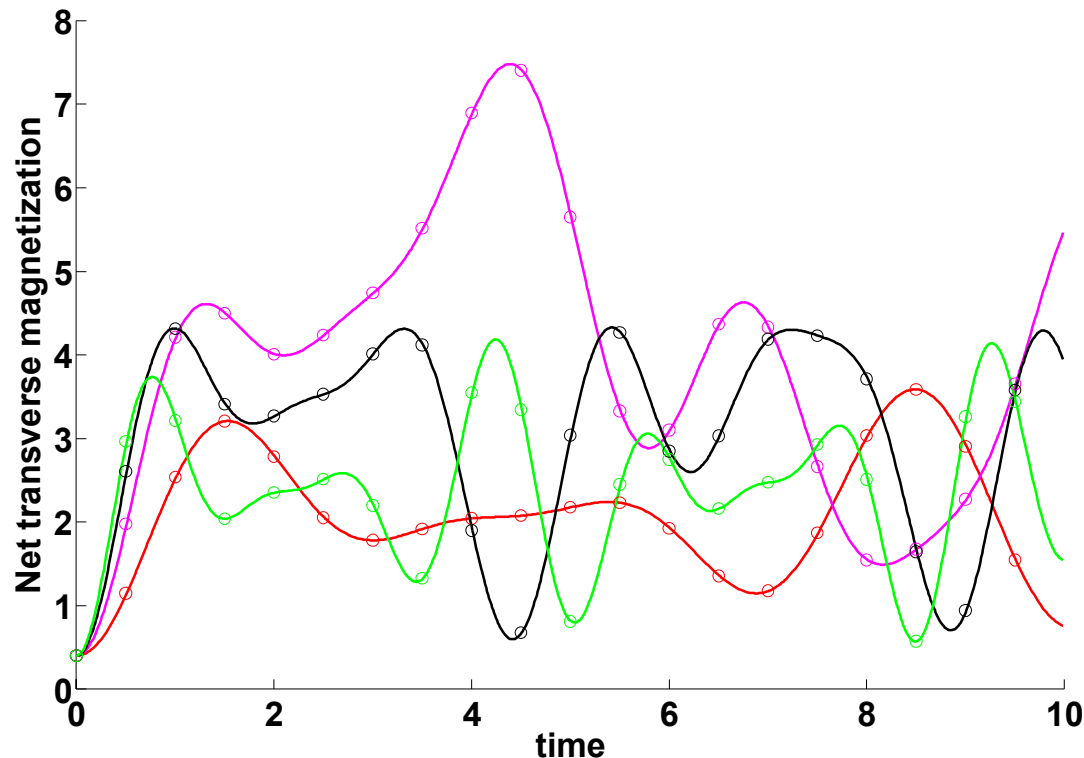
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1. The resulting dynamics is very informative; e.g. contains information about static phases of system

# Example

## POD-based model reduction to simulate quenching dynamics of the quantum Ising model

Simulation of quench dynamics in quantum Ising model  
(Circles: full model, lines: reduced order model)

Quenches to different parameters are indicated by different colors



Full order model:  
 $2^8 - 1 = 255$  complex  
numbers

Reduced order model:  
23 complex numbers

Order of magnitude  
improvement in simulation  
complexity

# Conclusions and continuing work

- The exponentially scaling of simulation complexity of many-body quantum dynamics may be an “illusion” in many cases
- Model reduction techniques can identify the relevant set of states
- In extreme cases, there may be invariant subspaces
  - Have developed certificates to identify such cases
- Continuing work
  - Develop practical methods for constructing reduced order models
  - Develop code to implement model reduction techniques

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## Model reduction

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