

Viscous Added Mass of a Moving Solid Object in a Closed Liquid-Filled Container

John R. Torczynski¹ and Louis A. Romero²

¹ Engineering Sciences Center

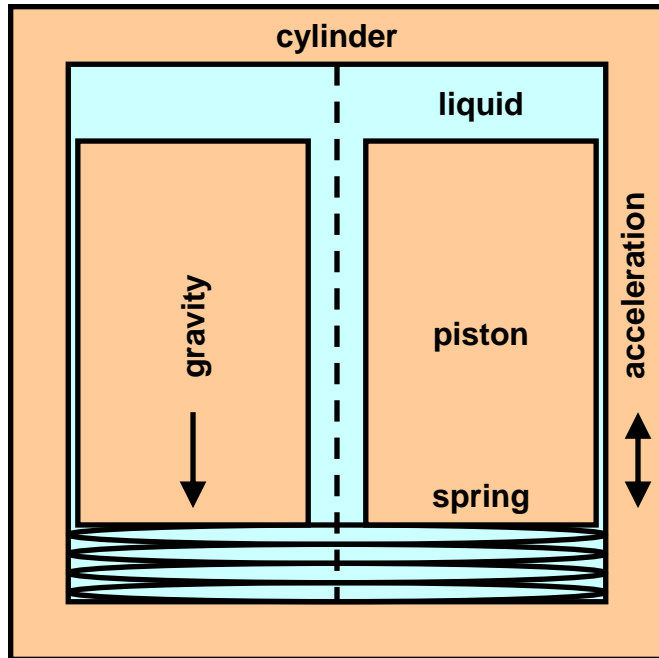
² Computing Research Center

Sandia National Laboratories

Albuquerque, New Mexico, USA

***DFD12, American Physical Society Division of Fluid Dynamics
65th Annual Meeting; San Diego, California; November 18-20, 2012***

Piston in Liquid-Filled Cylinder



$$\frac{dZ_p}{dt} = U_p$$

$$M_p \frac{dU_p}{dt} = ???$$

Motion of piston in closed cylinder

- Piston fills most of cylinder
- Spring supports it against gravity
- Viscous liquid fills open volume
- Cylinder is transiently accelerated

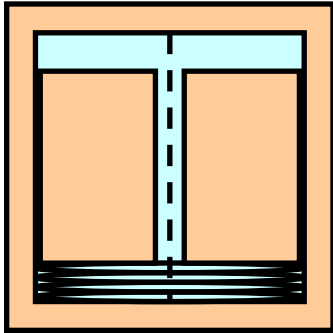
Develop ODE model for dynamics

- Consider cylinder reference frame
- Find liquid forces on piston
- Quasi-steady approximation

General approach is simplified here

- Axisymmetric geometry
- Motion only along axis

Rigorous Dynamics



$$\frac{dZ_P}{dt} = U_P, \quad M_P \frac{dU_P}{dt} = -K_S (Z_P - Z_S) - M_P (1 - (\rho_L / \rho_P)) (g + A_C) + F_L$$

spring

gravity/acceleration/buoyancy

liquid

Liquid force on piston in cylinder can be found rigorously

- Unsteady incompressible Navier-Stokes equations
- Cylinder reference frame: acceleration acts like gravity
- Subtract hydrostatic pressure: gravity/acceleration/buoyancy
- Piston position and velocity both vary in time

Liquid force is integral of stress tensor over piston surface

$$\frac{\partial}{\partial \mathbf{x}} \cdot \mathbf{u} = 0 \quad \rho_L \left\{ \frac{\partial \mathbf{u}}{\partial t} + \left(\mathbf{u} \cdot \frac{\partial}{\partial \mathbf{x}} \right) \mathbf{u} \right\} = \frac{\partial}{\partial \mathbf{x}} \cdot \boldsymbol{\sigma} \quad \boldsymbol{\sigma} = -p\mathbf{I} + \mu_L \left(\frac{\partial \mathbf{u}}{\partial \mathbf{x}} + \frac{\partial \mathbf{u}^T}{\partial \mathbf{x}} \right)$$

$$\mathbf{u} = U_P \hat{\mathbf{e}}_z \text{ on piston} \quad \mathbf{u} = \mathbf{0} \text{ on cylinder}$$

$$F_L = \int_{S_p} \hat{\mathbf{e}}_z \cdot \boldsymbol{\sigma} \cdot \hat{\mathbf{n}} dS$$

Quasi-Steady Approximation

Steady: Drag

$$\begin{aligned}\frac{\partial}{\partial \mathbf{x}} \cdot \mathbf{u}_0 &= 0 & \rho_L \frac{\partial}{\partial \mathbf{x}} \cdot (\mathbf{u}_0 \mathbf{u}_0) &= \frac{\partial}{\partial \mathbf{x}} \cdot \boldsymbol{\sigma}_0 & \boldsymbol{\sigma}_0 &= -p_0 \mathbf{I} + \mu_L \left(\frac{\partial \mathbf{u}_0}{\partial \mathbf{x}} + \frac{\partial \mathbf{u}_0}{\partial \mathbf{x}}^T \right) \\ \mathbf{u}_0 &= \hat{\mathbf{e}}_z U_P \text{ on piston} & \mathbf{u}_0 &= \mathbf{0} \text{ on cylinder} & F_{L0} &= \int_{S_p} \hat{\mathbf{e}}_z \cdot \boldsymbol{\sigma}_0 \cdot \hat{\mathbf{n}} dS \equiv -\beta_D U_P\end{aligned}$$

Unsteady: Added Mass

$$\begin{aligned}\frac{\partial}{\partial \mathbf{x}} \cdot \mathbf{u}_1 &= 0 & \rho_L \frac{\partial \mathbf{u}_0}{\partial U_P} \frac{dU_P}{dt} + \rho_L \frac{\partial}{\partial \mathbf{x}} \cdot (\mathbf{u}_0 \mathbf{u}_1 + \mathbf{u}_1 \mathbf{u}_0) &= \frac{\partial}{\partial \mathbf{x}} \cdot \boldsymbol{\sigma}_1 & \boldsymbol{\sigma}_1 &= -p_1 \mathbf{I} + \mu_L \left(\frac{\partial \mathbf{u}_1}{\partial \mathbf{x}} + \frac{\partial \mathbf{u}_1}{\partial \mathbf{x}}^T \right) \\ \mathbf{u}_1 &= \mathbf{0} \text{ on piston} & \mathbf{u}_1 &= \mathbf{0} \text{ on cylinder} & F_{L1} &= \int_{S_p} \hat{\mathbf{e}}_z \cdot \boldsymbol{\sigma}_1 \cdot \hat{\mathbf{n}} dS \equiv -M_A (dU_P/dt)\end{aligned}$$

Find liquid-force expression for quasi-steady flow

- Closed container has time scale to achieve steady flow
- Piston position & velocity change little over this time scale

Decompose flow field into sum of two contributions

- Steady part (“0”, large) yields **drag** force
- Unsteady part (“1”, small) yields **added-mass** force

Quasi-Steady Stokes Limit

Stokes Limit

$$\frac{\partial}{\partial \mathbf{x}} \cdot \mathbf{u} = 0$$

$$\mathbf{u} = U_P \hat{\mathbf{e}}_z \text{ on piston}$$

$$\mathbf{0} = \frac{\partial}{\partial \mathbf{x}} \cdot \boldsymbol{\sigma}$$

$$\mathbf{u} = \mathbf{0} \text{ on cylinder}$$

$$\boldsymbol{\sigma} = -p\mathbf{I} + 2\mu_L \mathbf{S}$$

$$\mathbf{S} = \frac{1}{2} \left(\frac{\partial \mathbf{u}}{\partial \mathbf{x}} + \frac{\partial \mathbf{u}^T}{\partial \mathbf{x}} \right)$$

$$F_L = -\beta_D U_P - M_A (dU_P/dt)$$

$$\frac{1}{2} \beta_D U_P^2 = \mu_L \int_{V_L} \mathbf{S} : \mathbf{S} dV$$

$$\frac{1}{2} M_A U_P^2 = \frac{1}{2} \rho_L \int_{V_L} \mathbf{u} \cdot \mathbf{u} dV$$

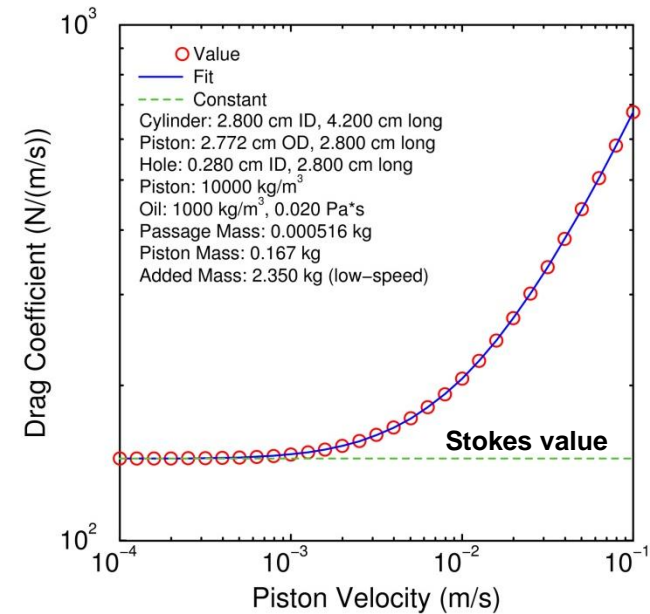
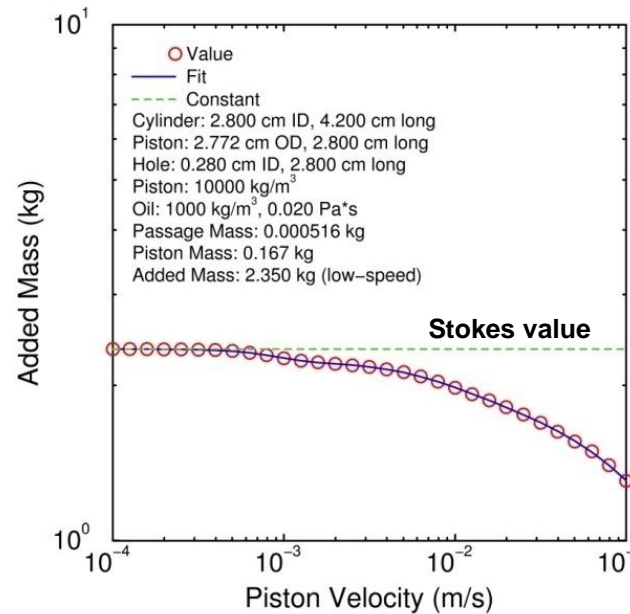
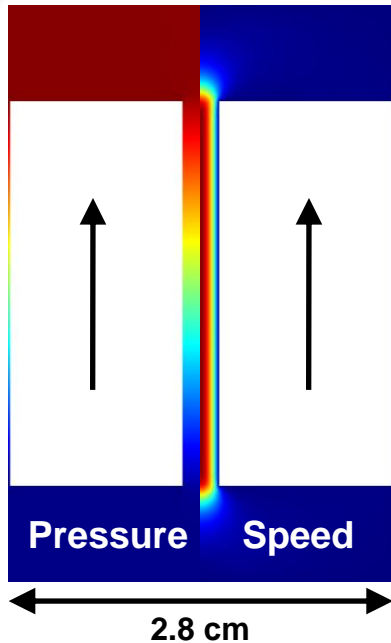
Expressions simplify considerably in Stokes limit

- Navier-Stokes: 3 CFD solutions for each piston velocity
 - Drag coefficient & added mass depend on piston velocity
- Stokes: only 1 CFD solution for all piston velocities
 - Drag coefficient & added mass are constants in this limit

Added mass caused by change in liquid kinetic energy

- Acceleration changes both piston & liquid kinetic energy
- Liquid acts like flywheel storing kinetic energy

Piston-Cylinder Example

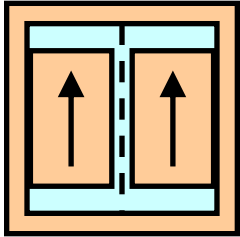


Added mass is very large (Stokes below)

- Liquid mass: 0.000516 kg (in passages)
- Piston mass: 0.167 kg
- Added mass: 2.350 kg (~14 × piston, ~4500 × liquid)

At large piston velocities, added mass & drag coefficient depart from corresponding Stokes values

Response to Constant Force



applied force
but no spring

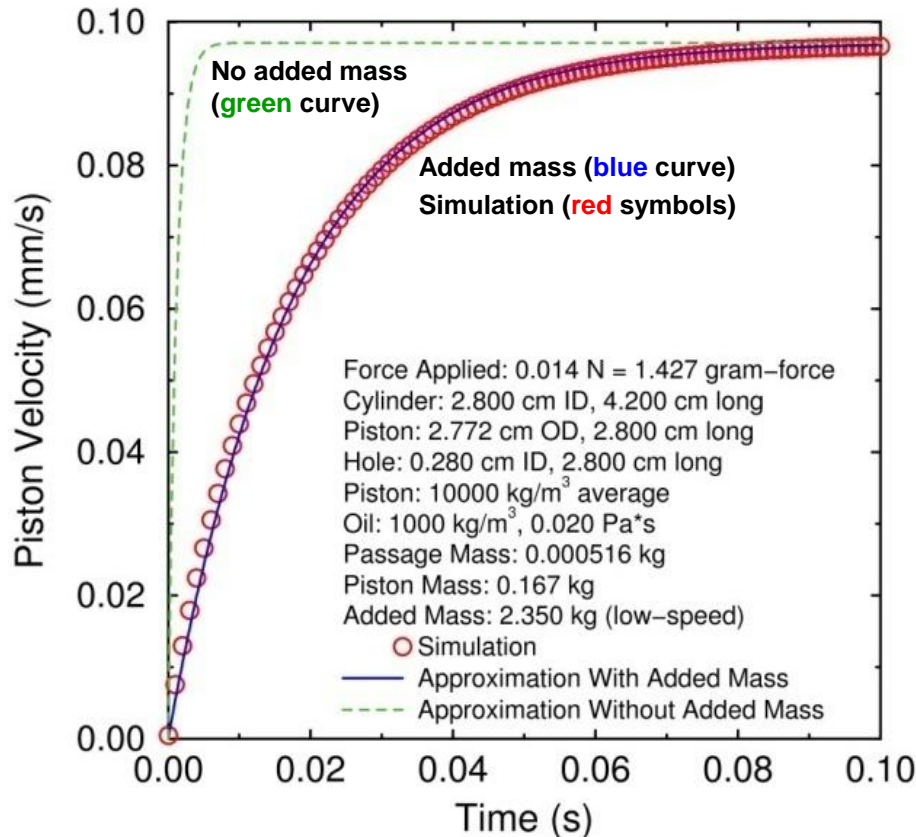


$$\frac{dZ_P}{dt} = U_P, \quad \left(M_P + \underset{\text{added mass}}{M_A[U_P]} \right) \frac{dU_P}{dt} = \underset{\text{applied force}}{F_{\text{applied}}} - \underset{\text{drag force}}{\beta_D[U_P]U_P}$$

added mass

applied force

drag force



Compare to rigorous Navier-Stokes simulation

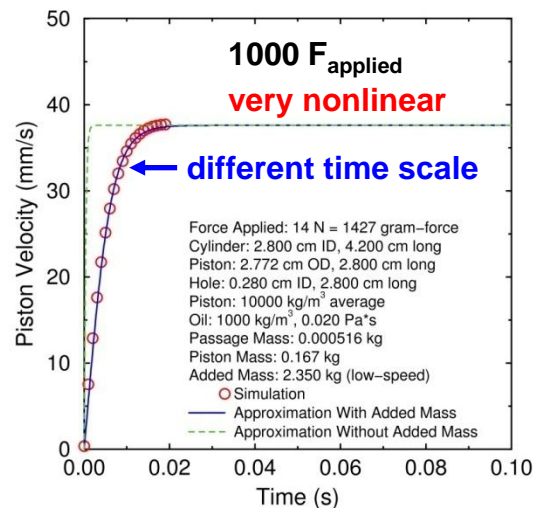
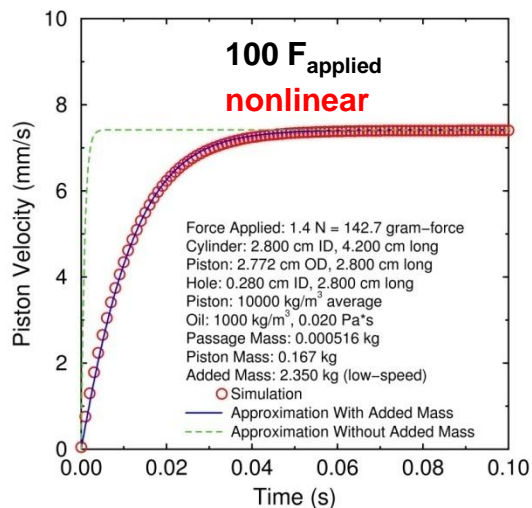
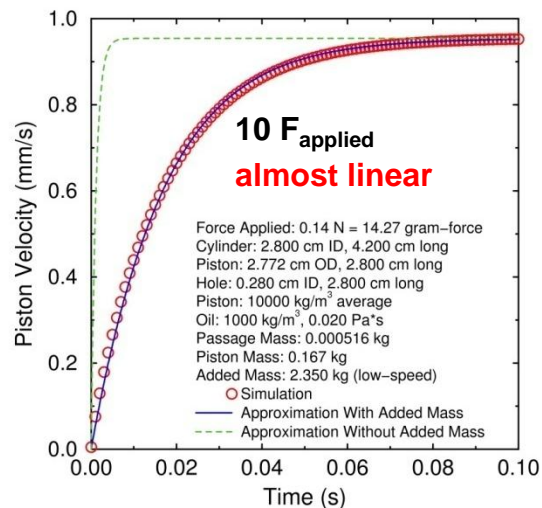
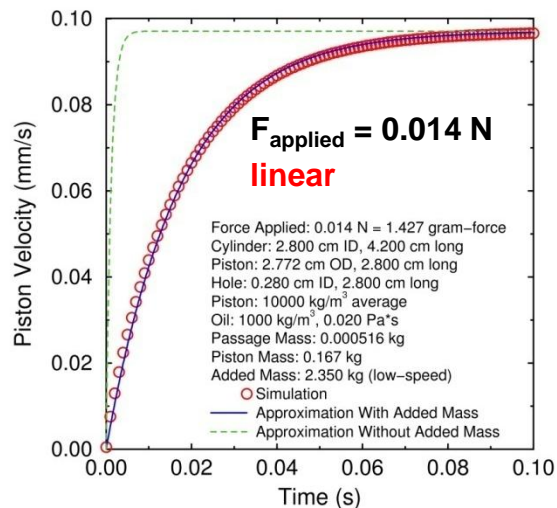
- Apply constant force to piston (no spring)
- Piston approaches terminal velocity
- Added mass determines time scale of response

Added mass is essential for accurate dynamics

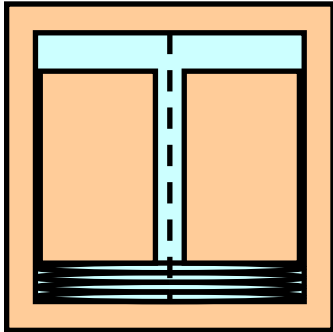
Linear to Nonlinear Conditions

Added mass produces accurate dynamics for all cases

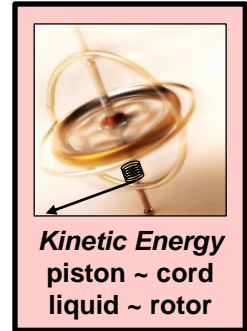
– From **linear** (Stokes) to **very nonlinear** (Navier-Stokes)



Viscous Added Mass



$$\frac{dZ_P}{dt} = U_P, \quad \left(M_P + \overset{\text{added mass}}{M_A[U_P]} \right) \frac{dU_P}{dt} = - \overset{\text{drag coefficient}}{\beta_D[U_P]} U_P - \underset{\text{gravity/acceleration/buoyancy}}{M_P(1 - (\rho_L/\rho_P))(g + A_C)} - \underset{\text{spring force}}{K_S(Z_P - Z_S)}$$



Kinetic Energy
piston ~ cord
liquid ~ rotor

Observed when piston accelerates in liquid-filled cylinder

- In low-frequency limit, how liquid kinetic energy changes

Can be accurately computed for complicated shapes

- General approach based on standard CFD simulations

Large for pistons with thin passages in closed cylinders

- Relative to piston mass & liquid mass in piston passages

Dynamically significant during transient acceleration

- Must include it to agree with rigorous dynamics

Enables rich dynamic behavior during vibration

- Position dependence, interaction with gas pockets