

Polarization as a field variable from molecular dynamics simulations

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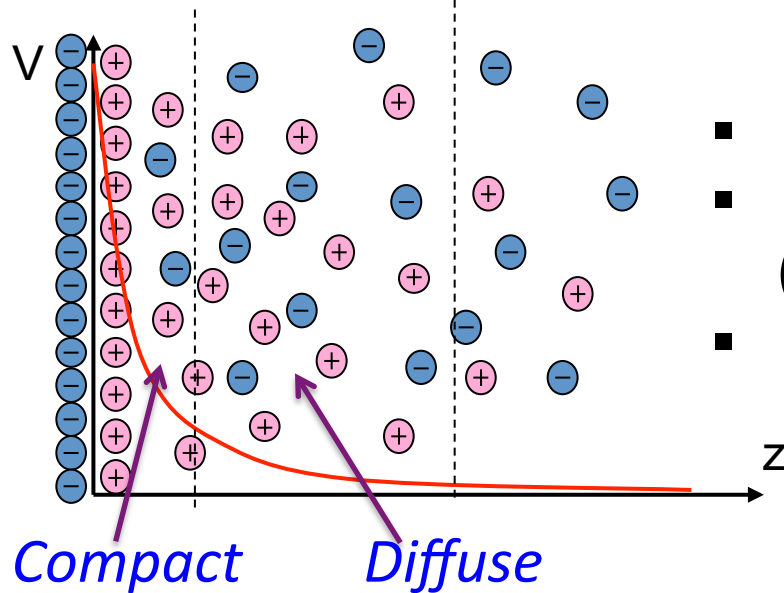
Sandia National Laboratories, Livermore, CA, USA

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Electrical Double Layers

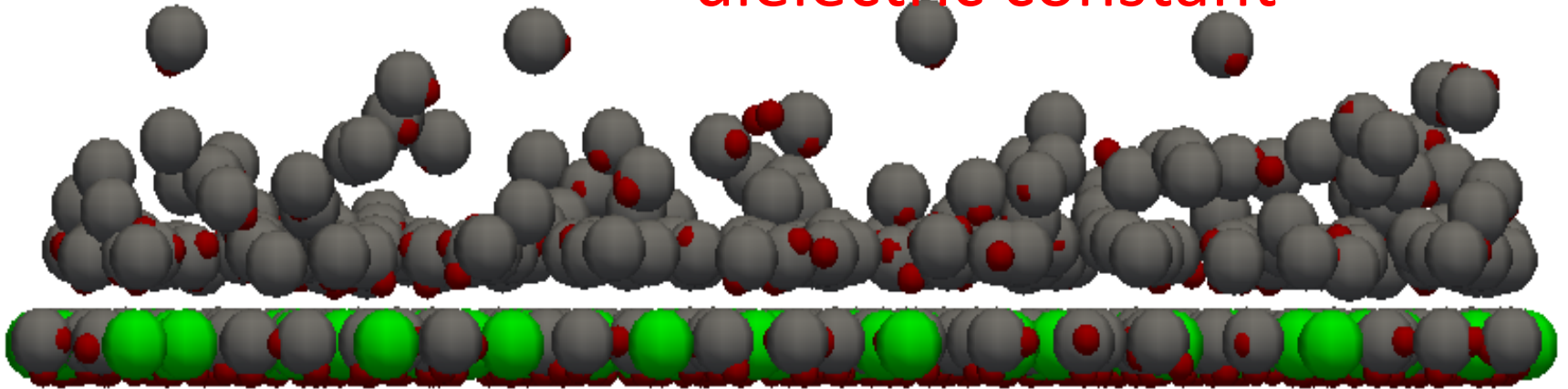
Models:

- Poisson-Boltzmann Theory
- Compact-Diffuse Layer Models (Kilic et. al. 2007)
- Mean-Field Models (Borukhov et. al. 2000)



$$D = \epsilon E = \epsilon_0 E + P$$

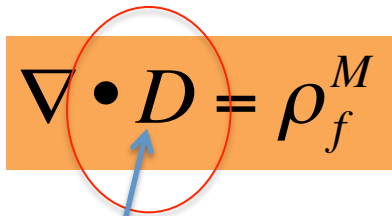
All layers modeled by bulk dielectric constant



We need better polarization models to depict reality.

Objective

Macroscopic Electrostatics:


$$\nabla \cdot \mathbf{D} = \rho_f^M$$

$\rho_f^M \rightarrow$ macroscopic free-charge density

$\mathbf{D} \rightarrow$ macroscopic electric displacement

$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} = \epsilon \mathbf{E}$$

$\mathbf{E} \rightarrow$ macroscopic electric field

$\epsilon \rightarrow$ dielectric constant

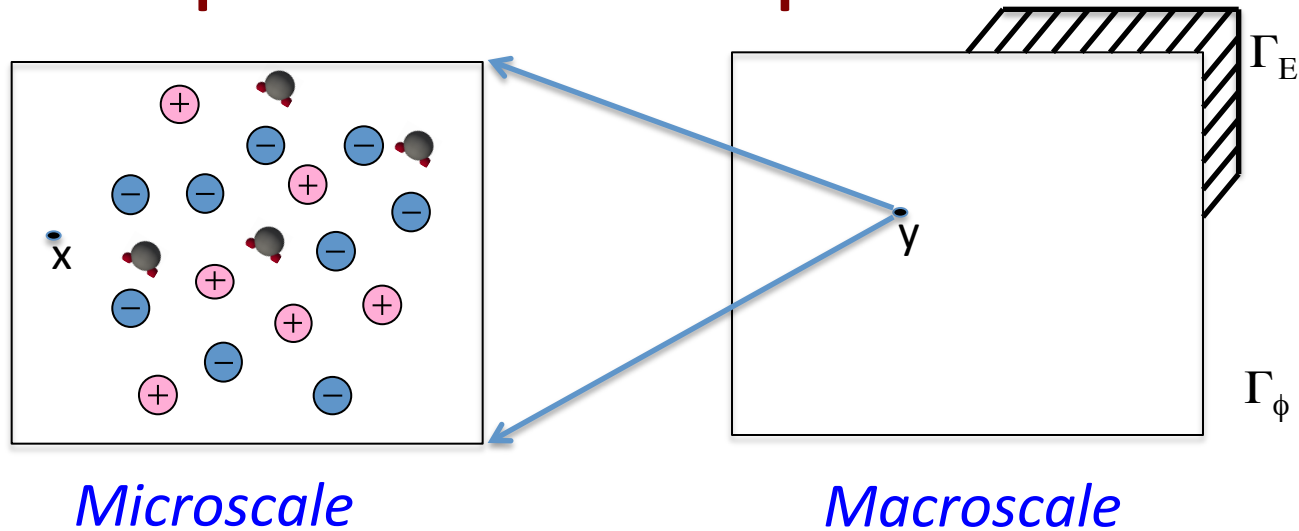
$\mathbf{P} \rightarrow$ macroscopic polarization

How do we calculate “P” using Molecular Information using Molecular Dynamics?

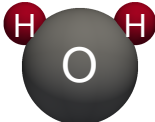
Today...

- Polarization in terms of molecular positions and charges.
- Dielectric constant of bulk water
- Application to electrical double layers
- At the interface
 - Water at the interface: Dipole and Quadrupole Moments

Microscopic-Macroscopic Connection



$$\nabla \cdot \mathbf{e}(\mathbf{x}, t) = \frac{1}{\epsilon_0} \rho^m(\mathbf{x}, t) \xrightarrow[\text{Coarse Grain}]{\text{Coarse Grain}} \nabla \cdot \mathbf{D}(\mathbf{y}, t) = \rho_f^M(\mathbf{y}, t)$$

$$\rho^m(\mathbf{x}, t) = \sum_{j=1}^{N_f} q_j \delta(\mathbf{x} - \mathbf{x}_j) + \sum_{n=1}^{N_m} \left(\sum_{j=1}^{N_n} q_j^n \delta(\mathbf{x} - \mathbf{x}_j^n) \right)$$


Can we derive macroscopic electrostatics equation from the microscopic electrostatics equation?

Mandadapu et. al. J. Chem. Phys. (2013);
Jackson, Classical Electrodynamics (1999)⁵

Microscopic-Macroscopic Connection

Coarse-graining Charge:

$$\rho^M(\mathbf{y}, t) = \int_{\mathcal{R}} \int_{\Gamma} \rho^m(\mathbf{x}, t) \Delta(\mathbf{y} - \mathbf{x}) f(\Gamma, t) d\Gamma d\mathbf{x}$$

*Phase-space
distribution*

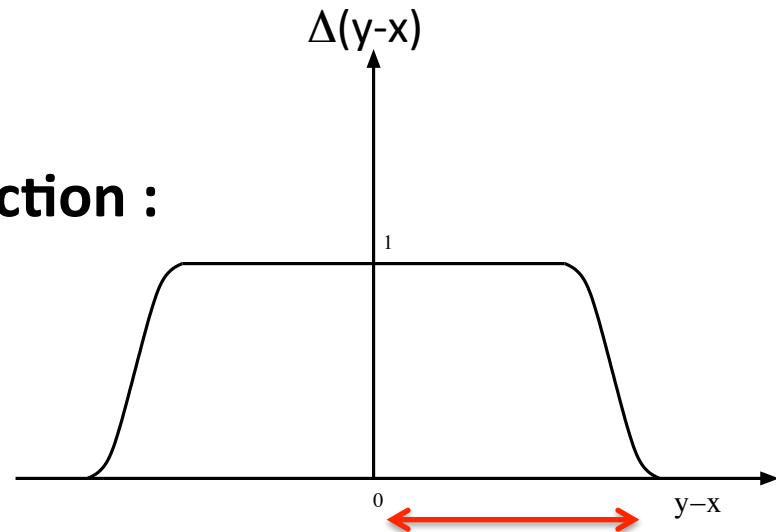
Coarse-graining function

Properties of coarse-graining function :

$$\int_R \Delta(y-x) dx = 1$$

$$\Delta(y-x) = 0, \text{ when } x \in \partial R$$

$$\text{supp } \Delta \ll R$$



Intermediate asymptotic length scale
Or
Coarse-graining length scale

Coarse-graining the microscopic electrostatic equation:

$$\int_{\mathcal{R}} \int_{\Gamma} \frac{\partial}{\partial \mathbf{x}'} \cdot \epsilon_0 \mathbf{e}(\mathbf{x}', t) \Delta(\mathbf{x} - \mathbf{x}') f(\Gamma, t) d\Gamma d\mathbf{x}' = \int_{\mathcal{R}} \int_{\Gamma} \rho^m(\mathbf{x}', t) \Delta(\mathbf{x} - \mathbf{x}') f(\Gamma, t) d\Gamma d\mathbf{x}'$$

Macroscopic electric field

$$\frac{\partial}{\partial \mathbf{x}} \cdot \epsilon_0 \mathbf{E}(\mathbf{x}, t) = \int_{\mathcal{R}} \int_{\Gamma} \rho^m(\mathbf{x}', t) \Delta(\mathbf{x} - \mathbf{x}') f(\Gamma, t) d\Gamma d\mathbf{x}'$$

$$\rho^m(\mathbf{x}, t) = \sum_{j=1}^{N_f} q_j \delta(\mathbf{x} - \mathbf{x}_j) + \sum_{n=1}^{N_m} \left(\sum_{j=1}^{N_n} q_j^n \delta(\mathbf{x} - \mathbf{x}_j^n) \right)$$

*Macroscopic
free charge*

$$\sum_{n=1}^{N_m} \left(\sum_{j=1}^{N_n} q_j^n \Delta(\mathbf{x} - \mathbf{x}_n - \mathbf{x}_{jn}) \right)$$

Gives rise to dipole and quadrupole moments when expanded around the center of mass of the solvent molecule.

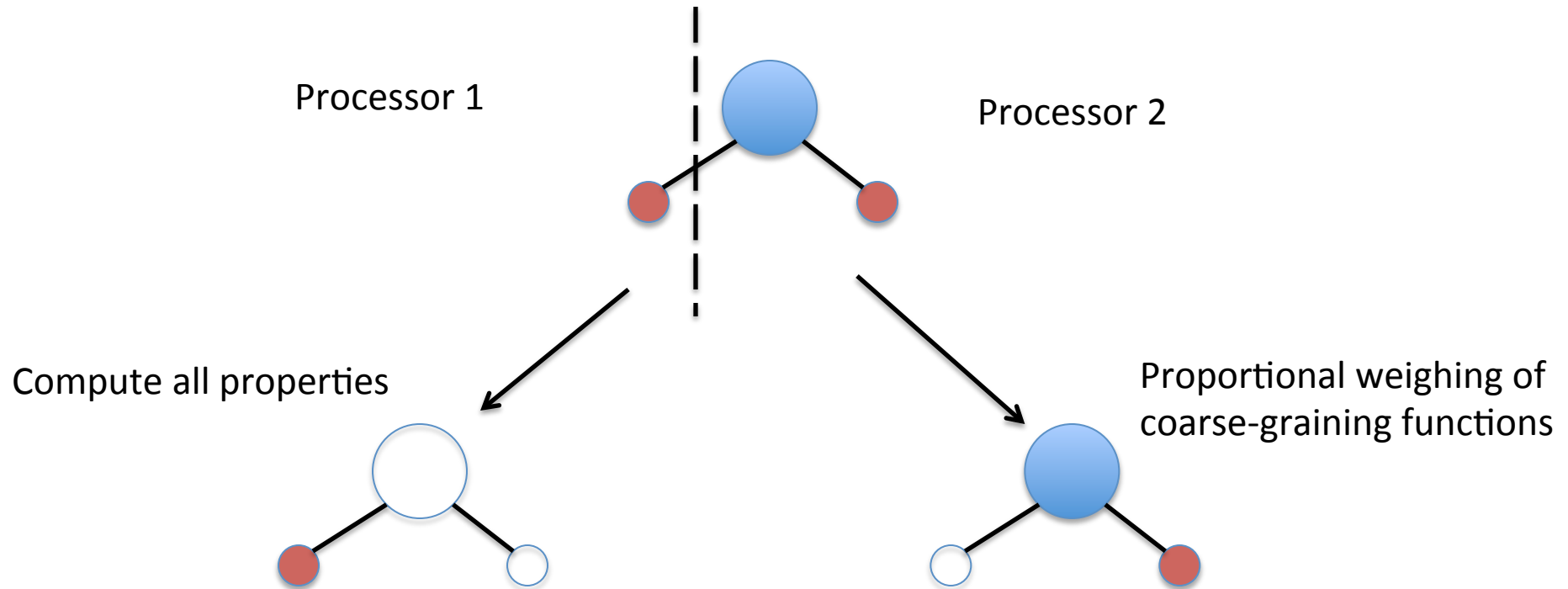
$$\frac{\partial}{\partial \mathbf{x}} \cdot \epsilon_0 \mathbf{E}(\mathbf{x}, t) = \int_{\Gamma} \left[\sum_{j=1}^{N_f} q_j \Delta(\mathbf{x} - \mathbf{x}_j) + \sum_{n=1}^{N_m} \left(\sum_{j=1}^{N_n} q_j^n \Delta(\mathbf{x} - \mathbf{x}_n) \right) \right] f(\Gamma, t) d\Gamma +$$

$$\frac{\partial}{\partial \mathbf{x}} \cdot \left[\underbrace{\int_{\Gamma} \sum_{n=1}^{N_m} \sum_{j=1}^{N_n} q_j^n \mathbf{x}_{jn} \Delta(\mathbf{x} - \mathbf{x}_n) f(\Gamma, t) d\Gamma}_{\text{Macroscopic Dipole Moment}} - \underbrace{\int_{\Gamma} \sum_{n=1}^{N_m} \sum_{j=1}^{N_n} q_j^n \frac{\mathbf{x}_{jn} \otimes \mathbf{x}_{jn}}{2!} \cdot \frac{\partial}{\partial \mathbf{x}} \Delta(\mathbf{x} - \mathbf{x}_n) f(\Gamma, t) d\Gamma + \dots}_{\text{Macroscopic Quadrupole Moment}} \right]$$

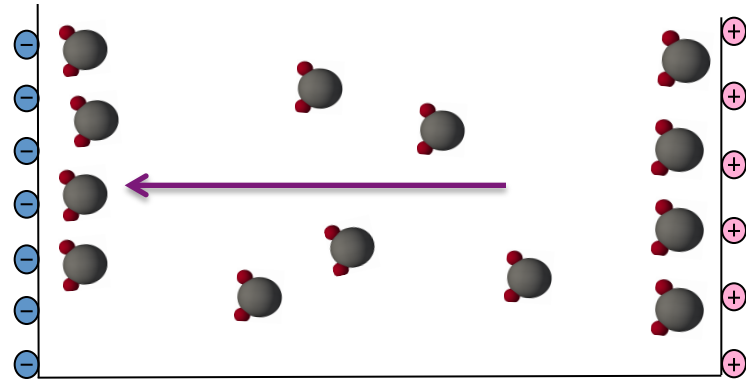
Order of Coarse Graining function Δ :

- **Constant:** gives only dipole moment
- **Linear:** gives up to quadrupole moment

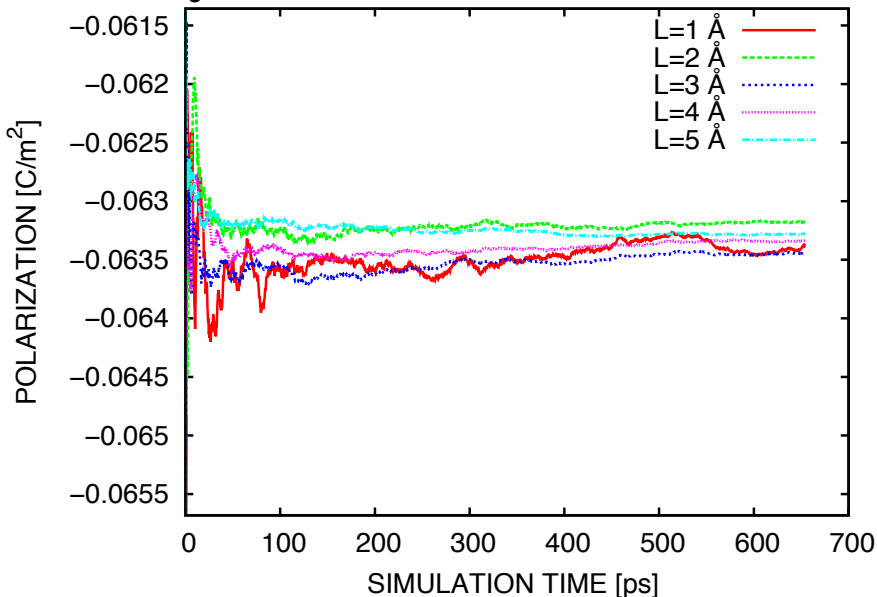
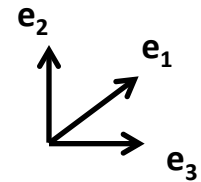
Scalable Implementation



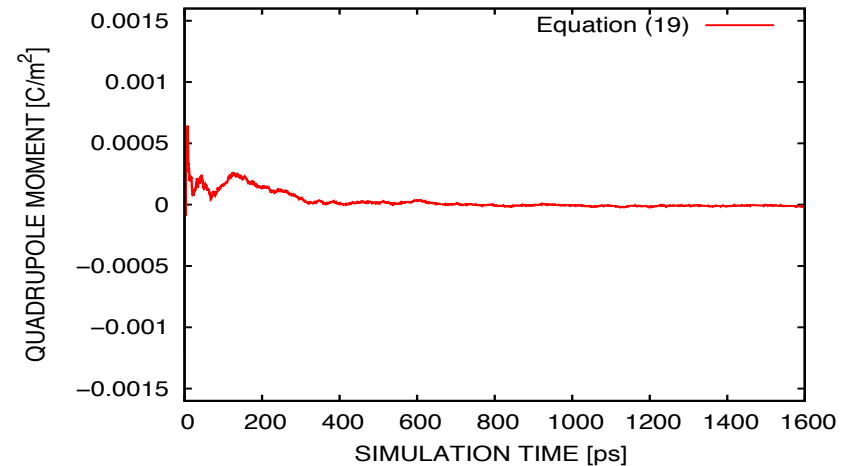
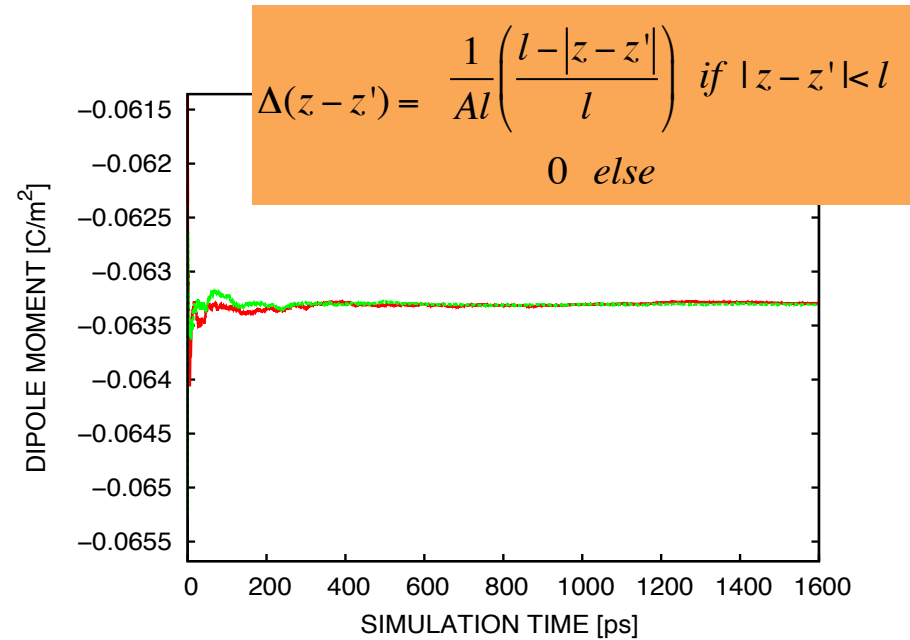
Bulk Water-TIP3P(Validation)



$$\Delta(z - z') = \begin{cases} \frac{1}{Al} & \text{if } |z - z'| < l \\ 0 & \text{else} \end{cases}$$

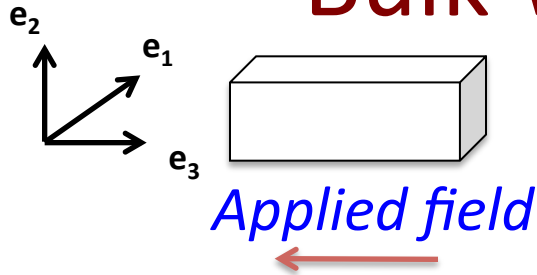


Polarization as a time average at the center of the simulation box for varying averaging length

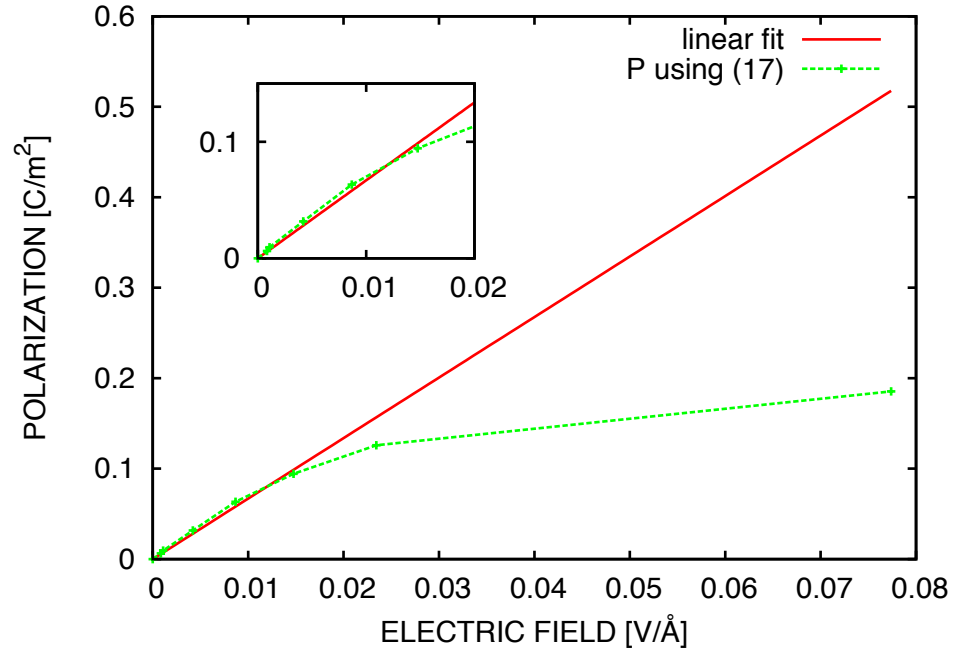


Quadrupole moments are negligible in the bulk. No surprises !

Bulk Water-TIP3P(Validation)



$$\Delta(z-z') = \begin{cases} \frac{1}{Al} & \text{if } |z-z'| < l \\ 0 & \text{else} \end{cases}$$

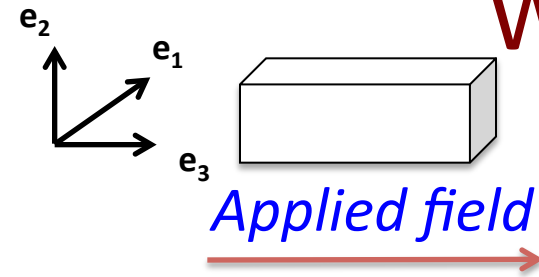


Polarization vs. Resultant electric field

$$\frac{\epsilon}{\epsilon_0} = 73$$

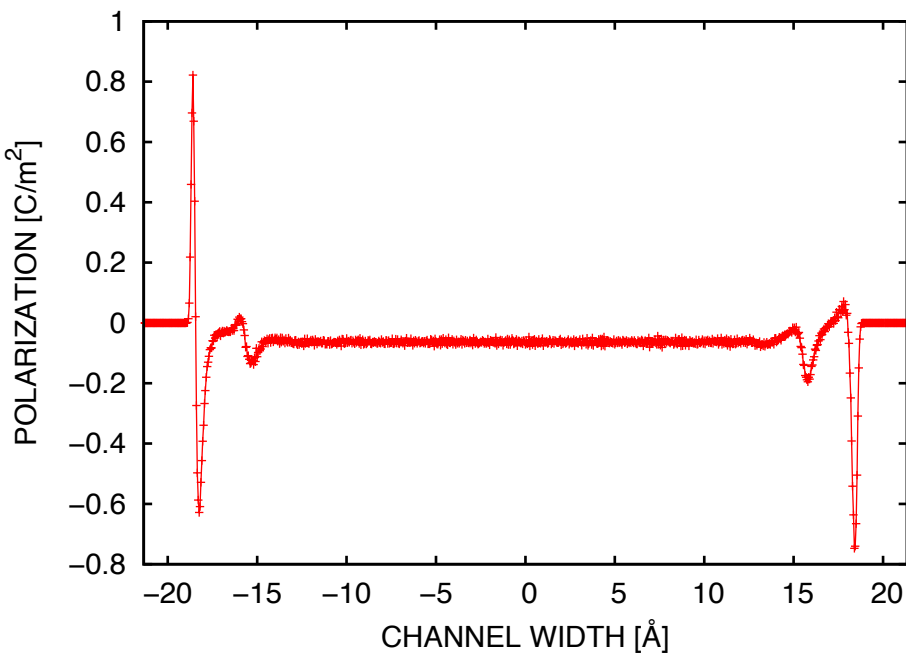
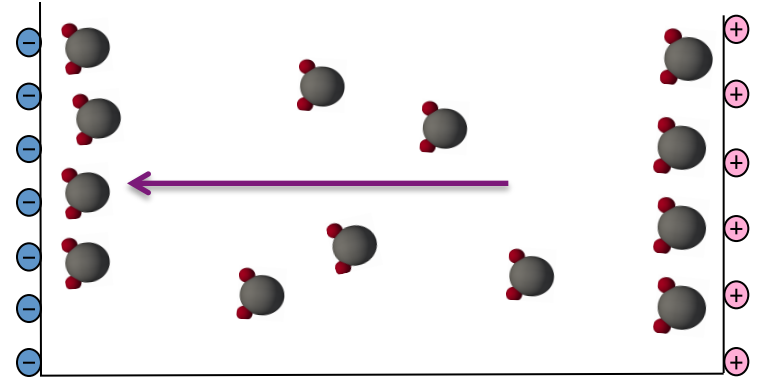
Note that the resultant electric fields are supposed to be averaged with the same averaging volume.

Water at the Interfaces

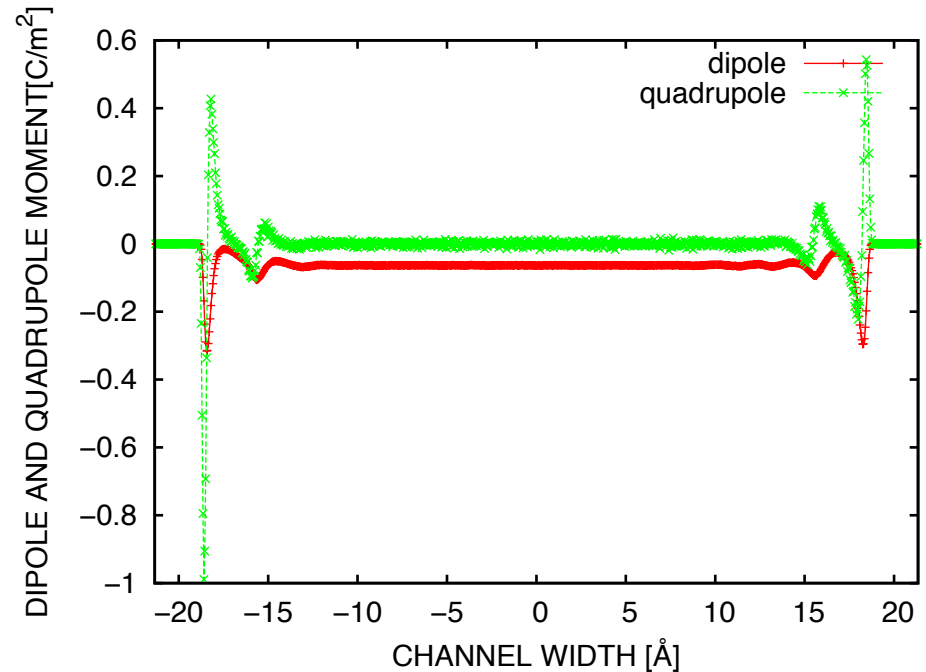


$$\Delta(z-z') = \frac{1}{Al} \left(\frac{l-|z-z'|}{l} \right) \text{ if } |z-z'| < l$$

$$0 \text{ else}$$

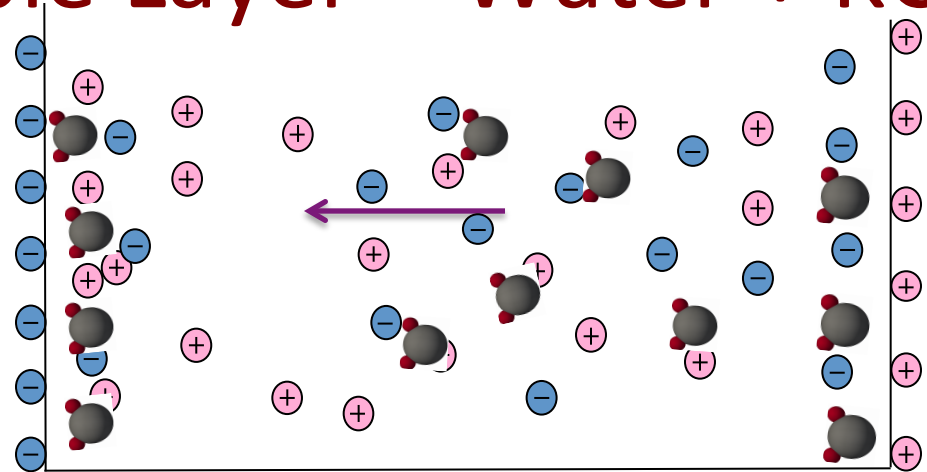
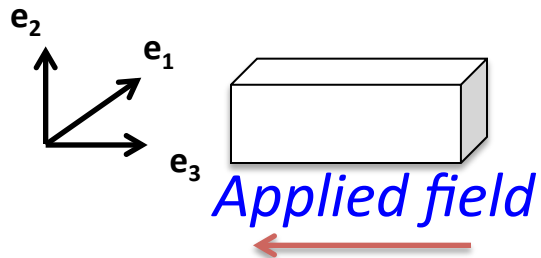


Quadrupole moments are NOT negligible at the interface !!!



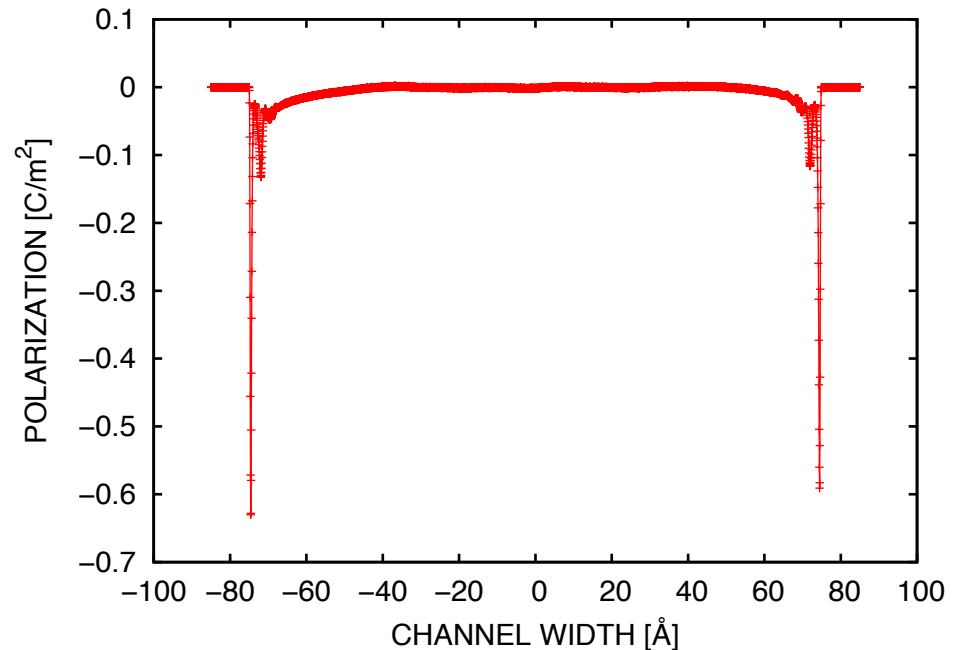
Dipole and Quadrupole Moments across the channel width for averaging length of 0.05 Angstroms.

Electrical Double Layer – Water + KCl

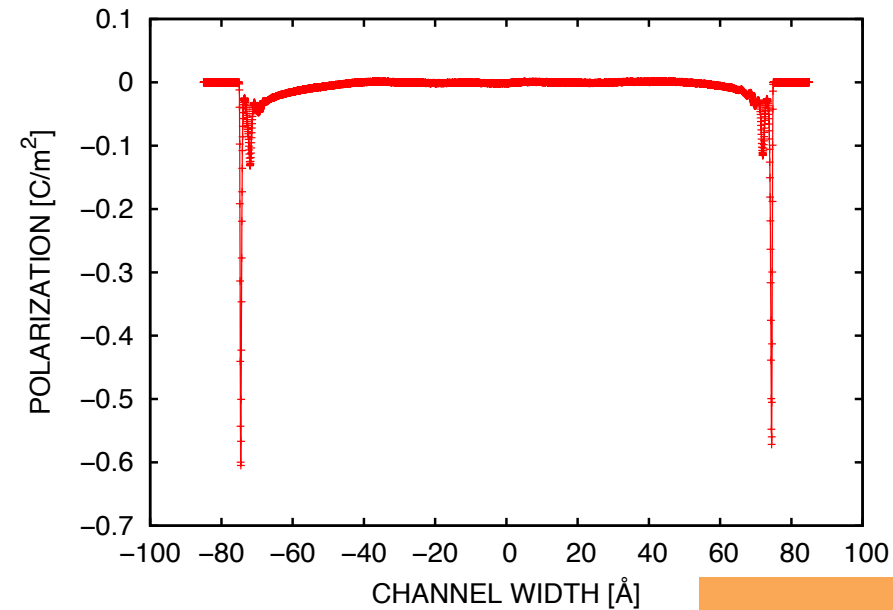


KCl concentration: 100mM
Debye length : 1nm

$$\Delta(z - z') = \begin{cases} \frac{1}{Al} & \text{if } |z - z'| < l \\ 0 & \text{else} \end{cases}$$

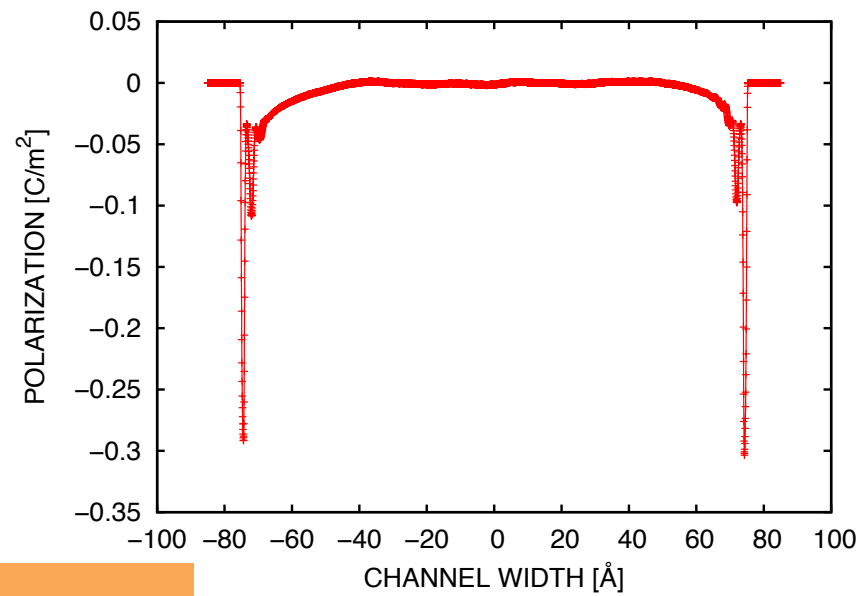


Polarization across the channel width for averaging length of 0.05 Angstroms.

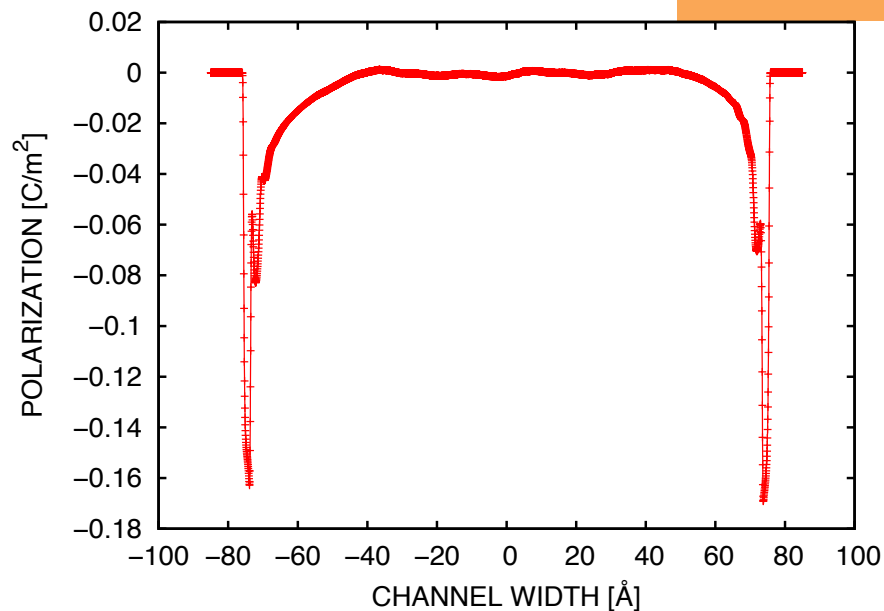


Averaging length = 0.1 Angs.

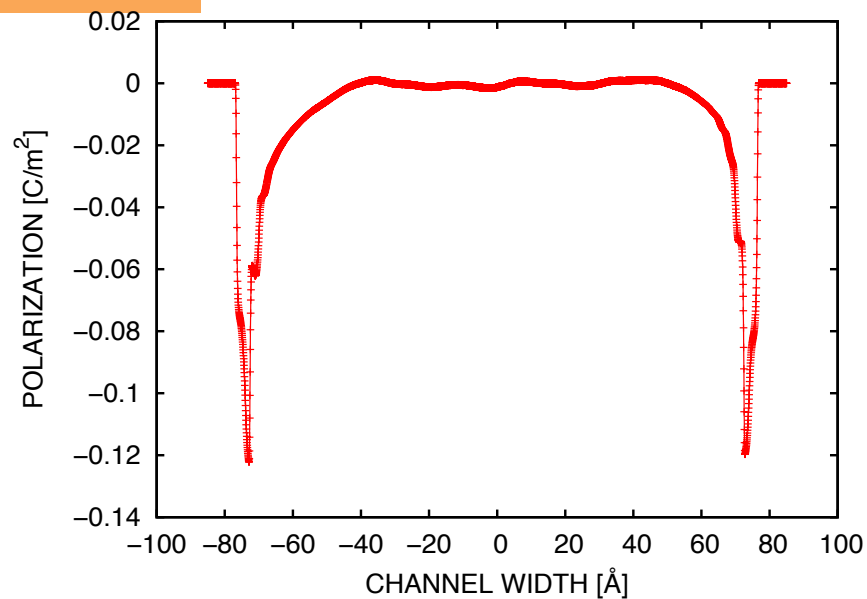
$$\Delta(z - z') = \begin{cases} \frac{1}{Al} & \text{if } |z - z'| < l \\ 0 & \text{else} \end{cases}$$



Averaging length = 0.5 Angs.

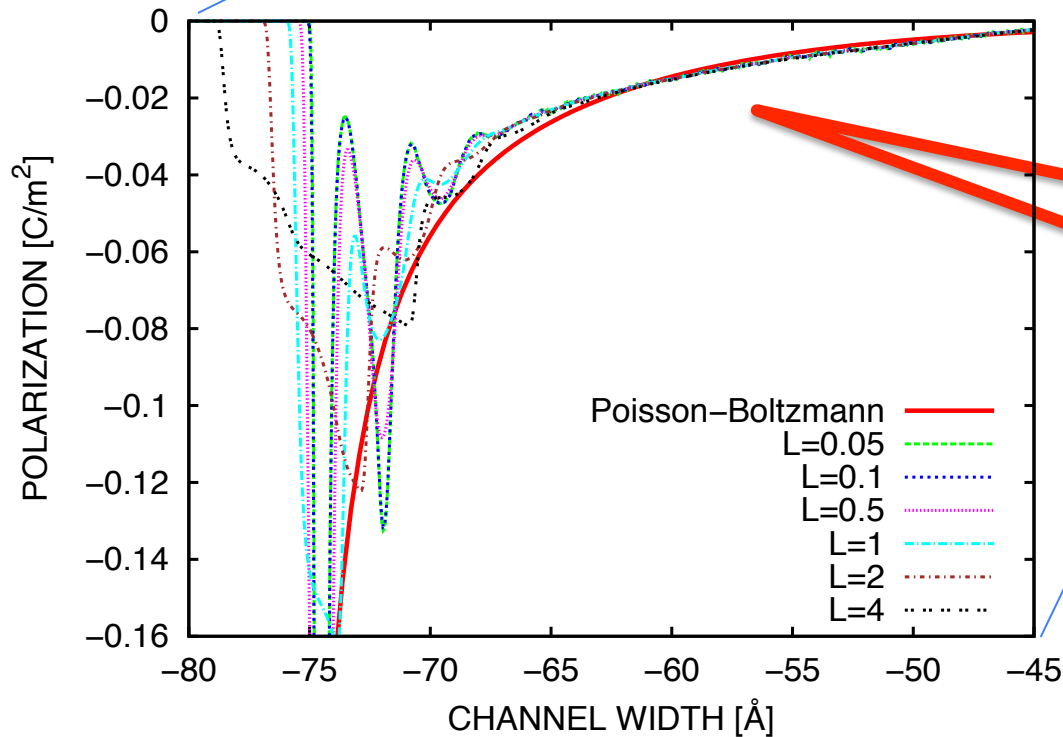
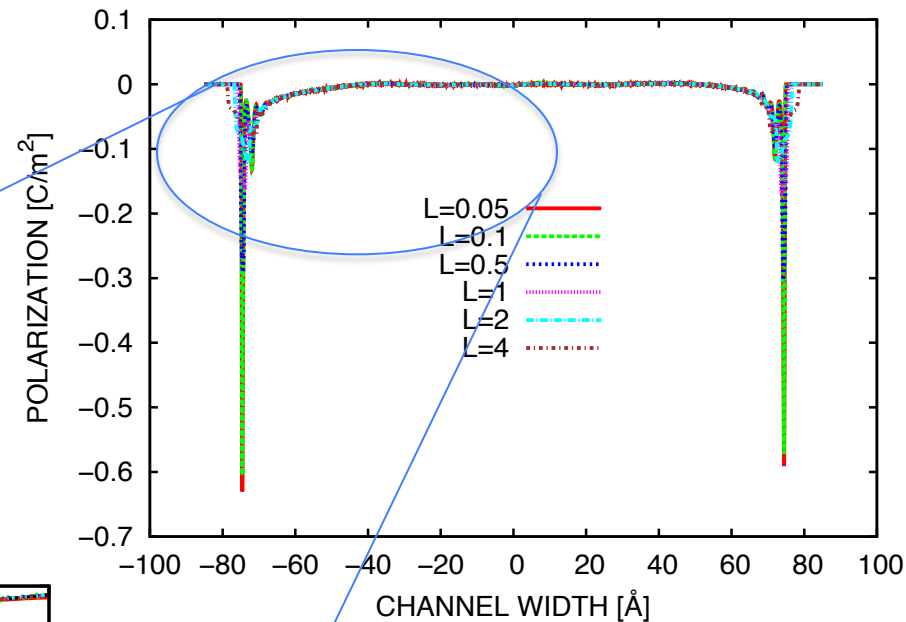


Averaging length = 1.0 Angs.



Averaging length = 2.0 Angs.

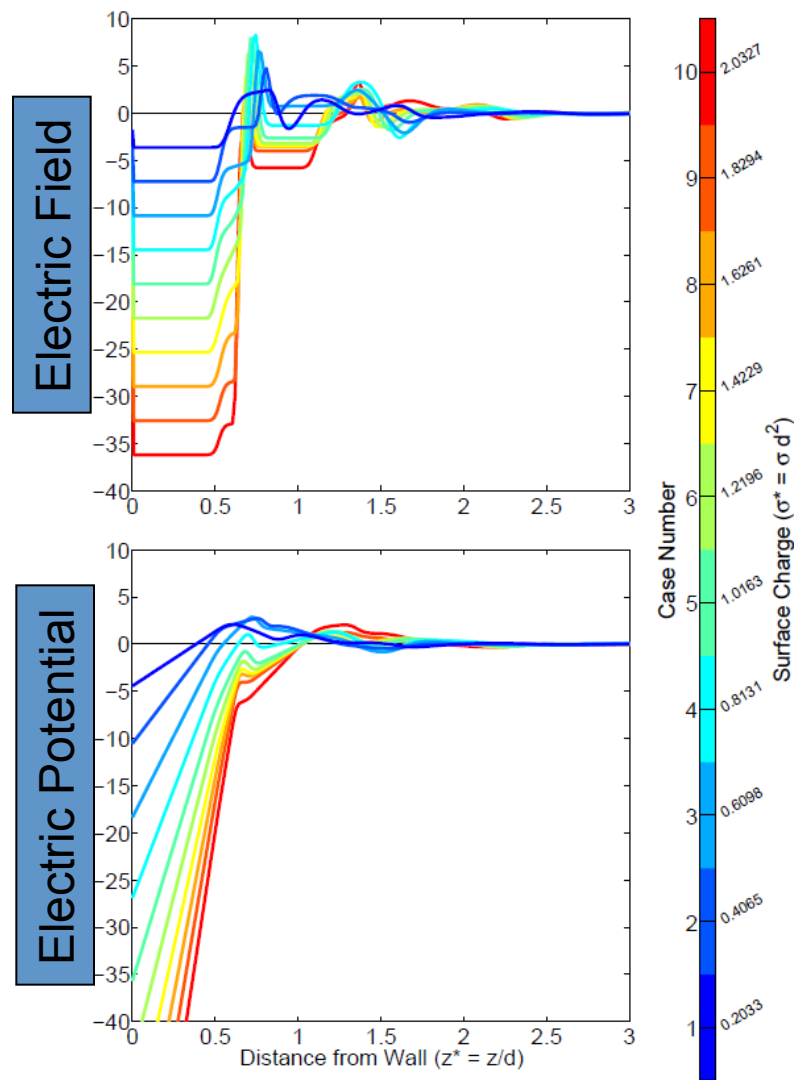
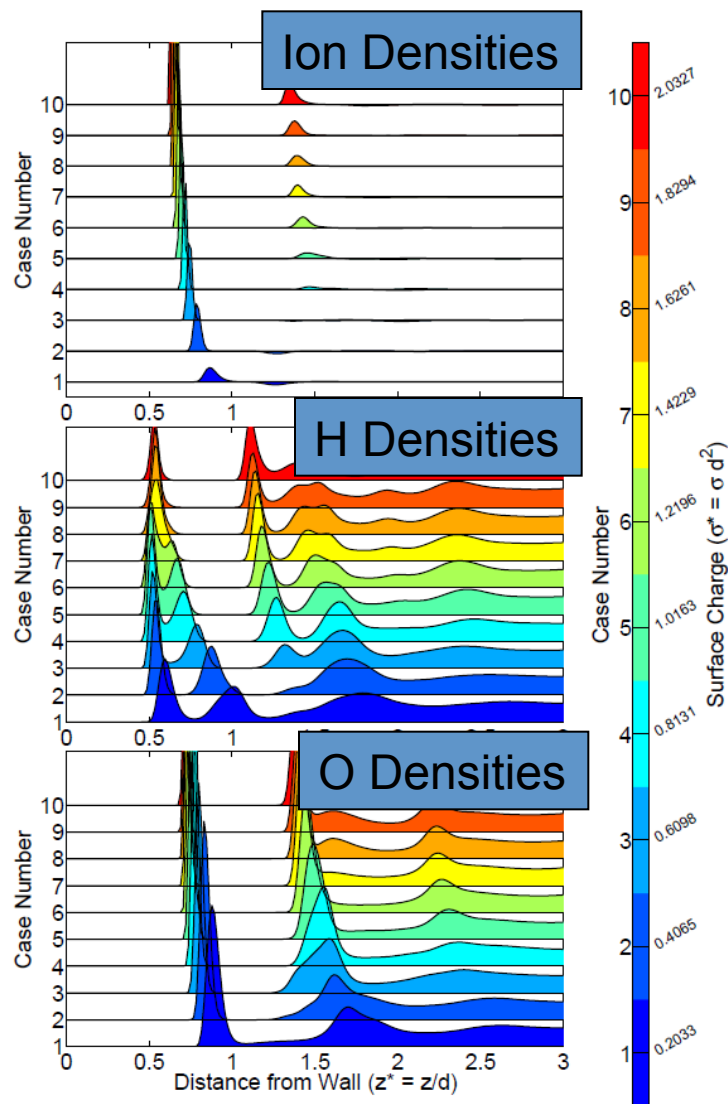
$$\Delta(z - z') = \begin{cases} \frac{1}{Al} & \text{if } |z - z'| < l \\ 0 & \text{else} \end{cases}$$



INTERMEDIATE ASYMPTOTIC
LENGTH SCALE !!!

Barenblatt, Scaling (2003)

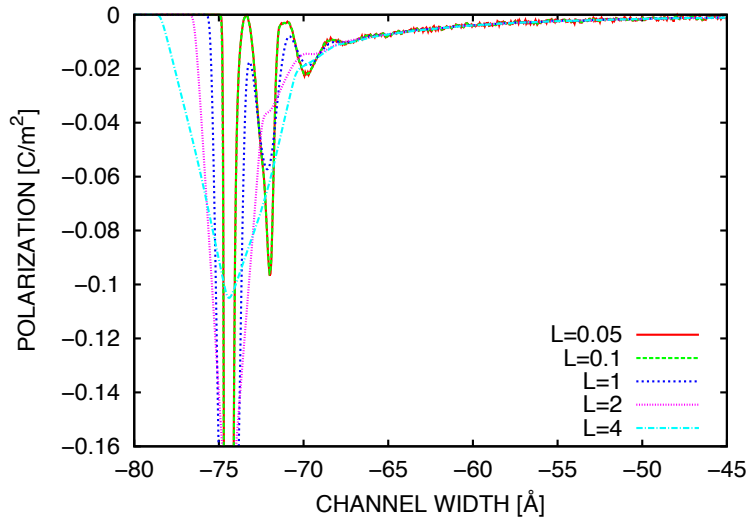
Electrical Double Layer – High Charge



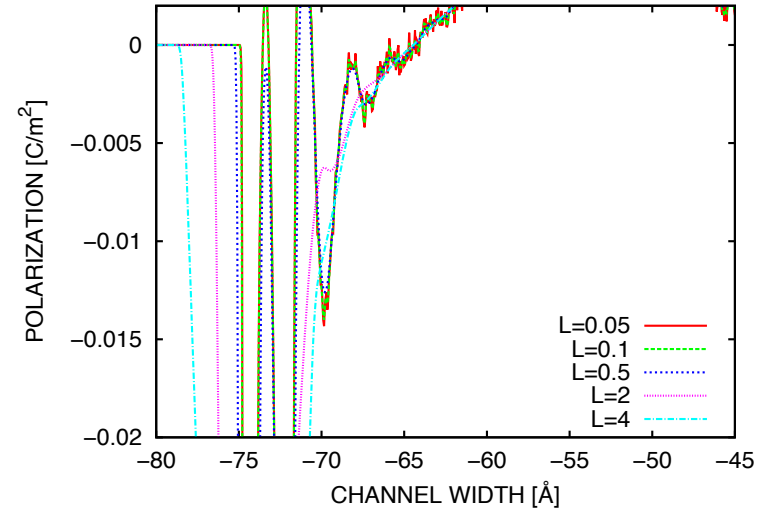
Lee, Templeton, Mandadapu, Zimmerman, J. Chem. Theor. Comp (2013)

Electrical Double Layer – Anode

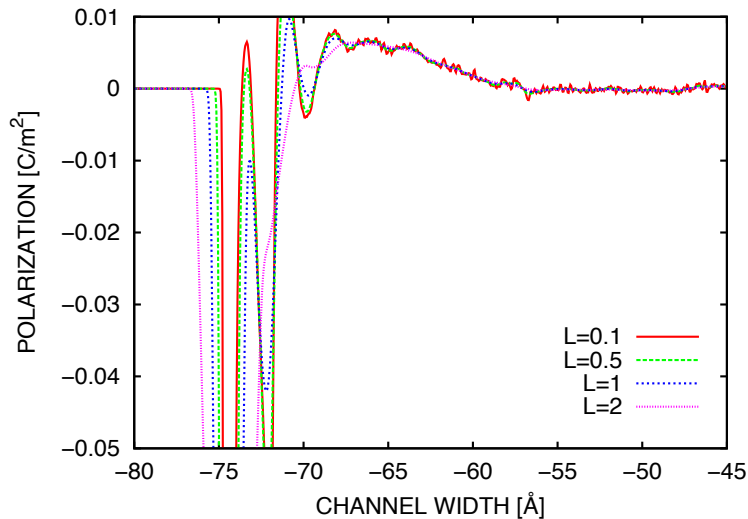
Case 1



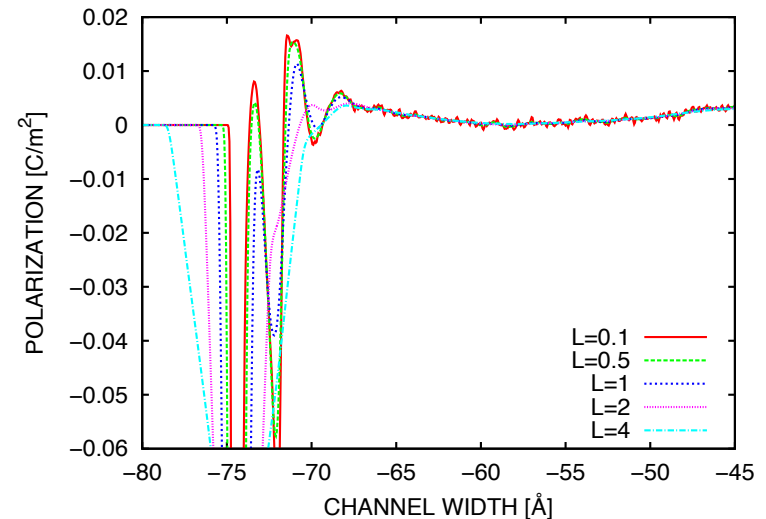
Case 3



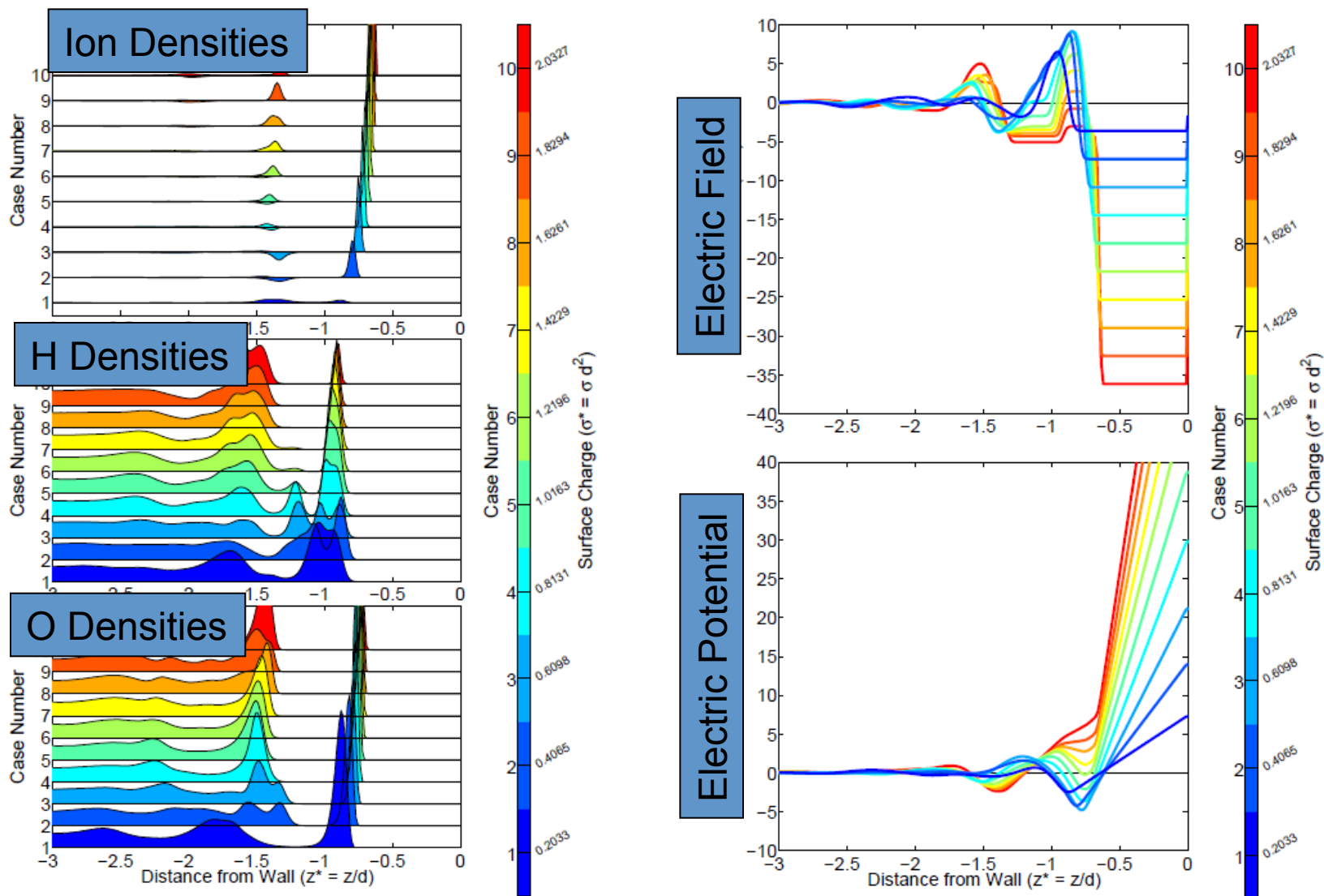
Case 5



Case 10

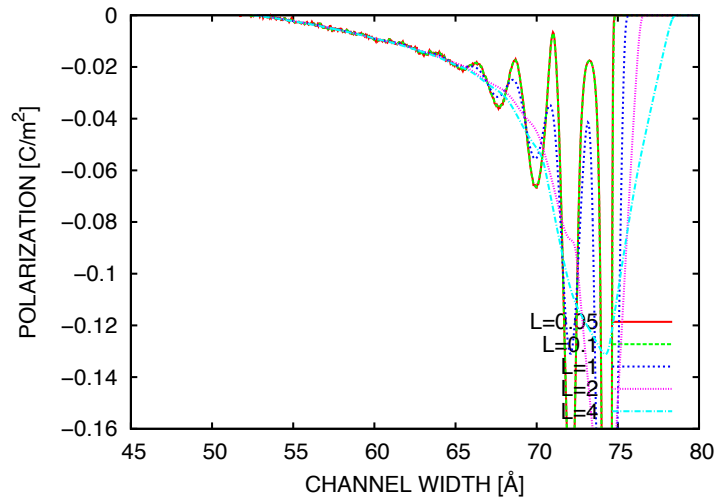


Electrical Double Layer – Cathode

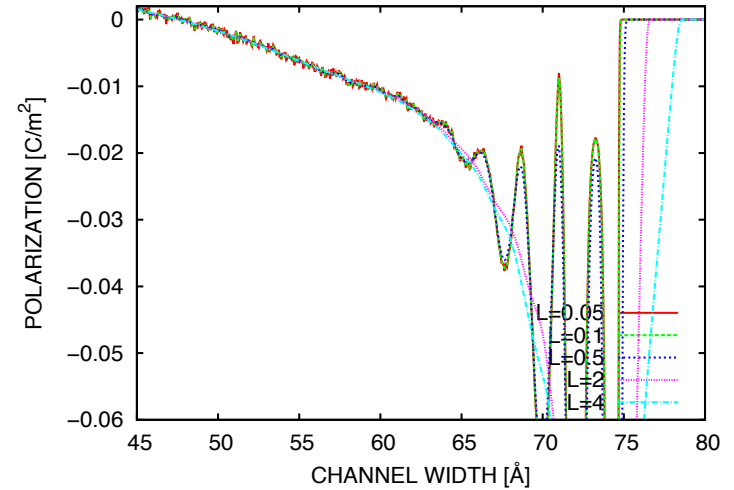


Electrical Double Layer – Cathode

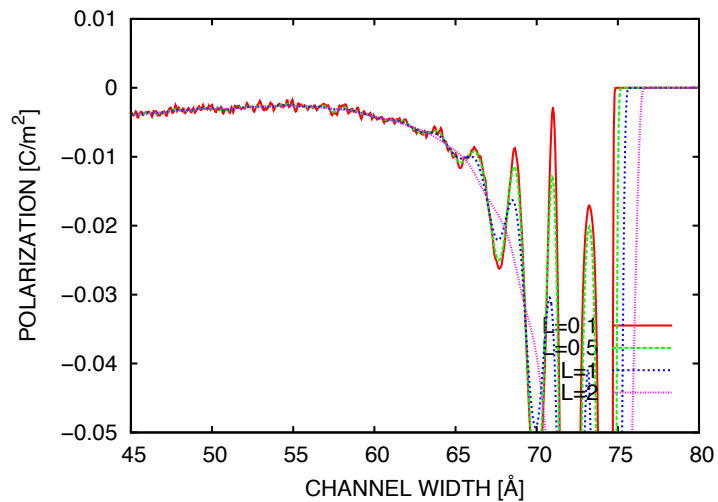
Case 1



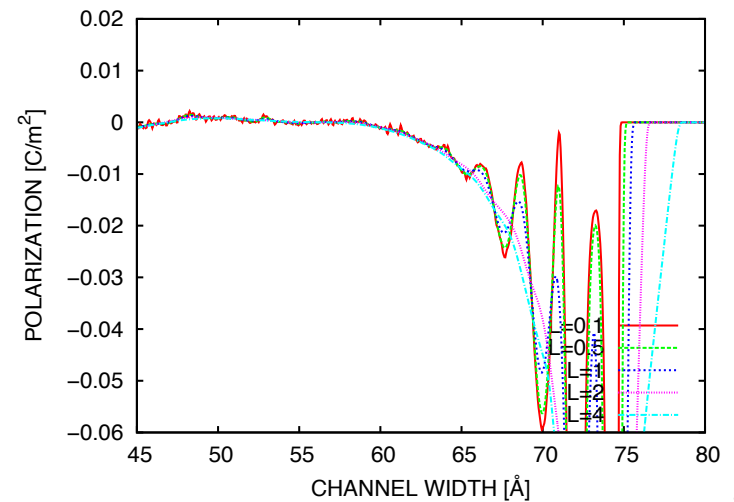
Case 3



Case 5

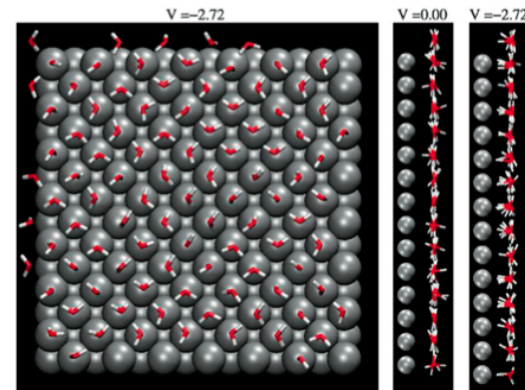
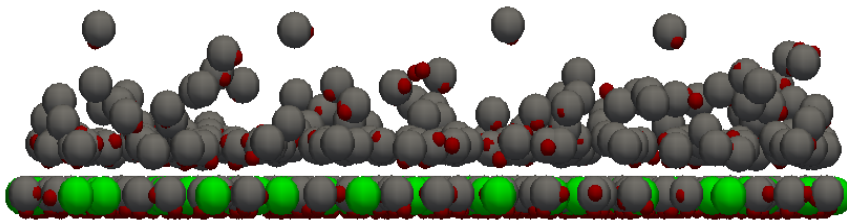


Case 10



Conclusions

- A theoretical and computational procedure in real space exists now for systematically calculating the moments that contribute to the total polarization vector.
- Intermediate asymptotic scales and therefore length scales of the diffuse layers can be found for various concentrations and surface charges.
- Quadrupole moments do not seem to be negligible at the interface.



Willard, Chandler (2008)

- Study polarization in stern, compact and diffuse layers.
- Develop *new models* by understanding hydrogen bonding networks.

Acknowledgements

Funding: Sandia LDRD Program

K. K. Mandadapu, J. A. Templeton and J. W. Lee, Polarization as a field variable from molecular dynamics simulations.

The Journal of Chemical Physics (2013)

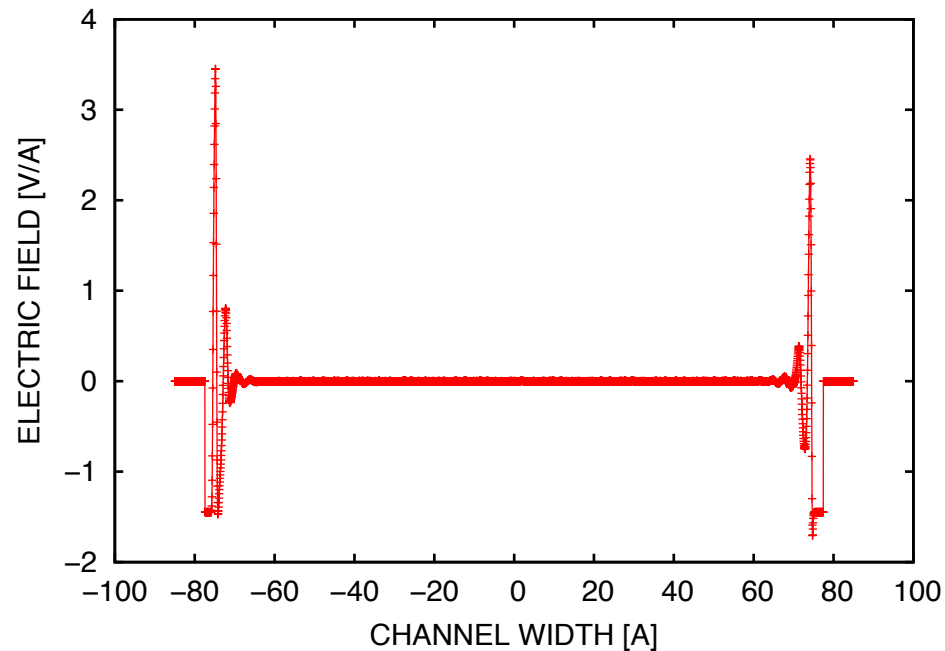
Computational software –
AtC package of LAMMPS.

Available to the public:

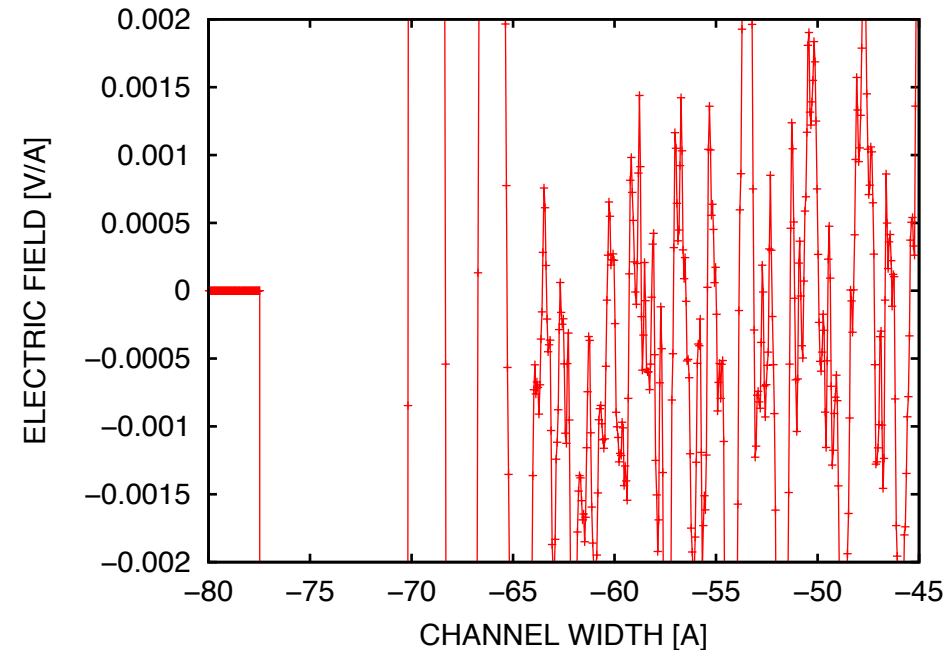
<http://lammps.sandia.gov>

Thank You

Advantages - Polarization



Electric field across the channel width using averages from 10 independent simulations using 10 million time steps.



Electric fields do not converge even for large number of timesteps !!!

Coarse-graining the microscopic electrostatic equation:

$$\int_{\mathcal{R}} \int_{\Gamma} \frac{\partial}{\partial \mathbf{x}'} \cdot \epsilon_0 \mathbf{e}(\mathbf{x}', t) \Delta(\mathbf{x} - \mathbf{x}') f(\Gamma, t) d\Gamma d\mathbf{x}' = \int_{\mathcal{R}} \int_{\Gamma} \rho^m(\mathbf{x}', t) \Delta(\mathbf{x} - \mathbf{x}') f(\Gamma, t) d\Gamma d\mathbf{x}'$$

Macroscopic electric field

$$\frac{\partial}{\partial \mathbf{x}} \cdot \epsilon_0 \mathbf{E}(\mathbf{x}, t) = \int_{\mathcal{R}} \int_{\Gamma} \rho^m(\mathbf{x}', t) \Delta(\mathbf{x} - \mathbf{x}') f(\Gamma, t) d\Gamma d\mathbf{x}'$$

$$\begin{aligned} \frac{\partial}{\partial \mathbf{x}} \cdot \epsilon_0 \mathbf{E}(\mathbf{x}, t) = & \int_{\Gamma} \left[\sum_{j=1}^{N_f} q_j \Delta(\mathbf{x} - \mathbf{x}_j) + \sum_{n=1}^{N_m} \left(\sum_{j=1}^{N_n} q_j^n \Delta(\mathbf{x} - \mathbf{x}_n) \right) \right] f(\Gamma, t) d\Gamma + \\ & \frac{\partial}{\partial \mathbf{x}} \cdot \left[\underbrace{\int_{\Gamma} \sum_{n=1}^{N_m} \sum_{j=1}^{N_n} q_j^n \mathbf{x}_{jn} \Delta(\mathbf{x} - \mathbf{x}_n) f(\Gamma, t) d\Gamma}_{\text{Dipole Moment}} - \underbrace{\int_{\Gamma} \sum_{n=1}^{N_m} \sum_{j=1}^{N_n} q_j^n \frac{\mathbf{x}_{jn} \otimes \mathbf{x}_{jn}}{2!} \cdot \frac{\partial}{\partial \mathbf{x}} \Delta(\mathbf{x} - \mathbf{x}_n) f(\Gamma, t) d\Gamma}_{\text{Quadrupole Moment}} + \dots \right] \end{aligned}$$

Order of Coarse Graining function Δ :

Macroscopic

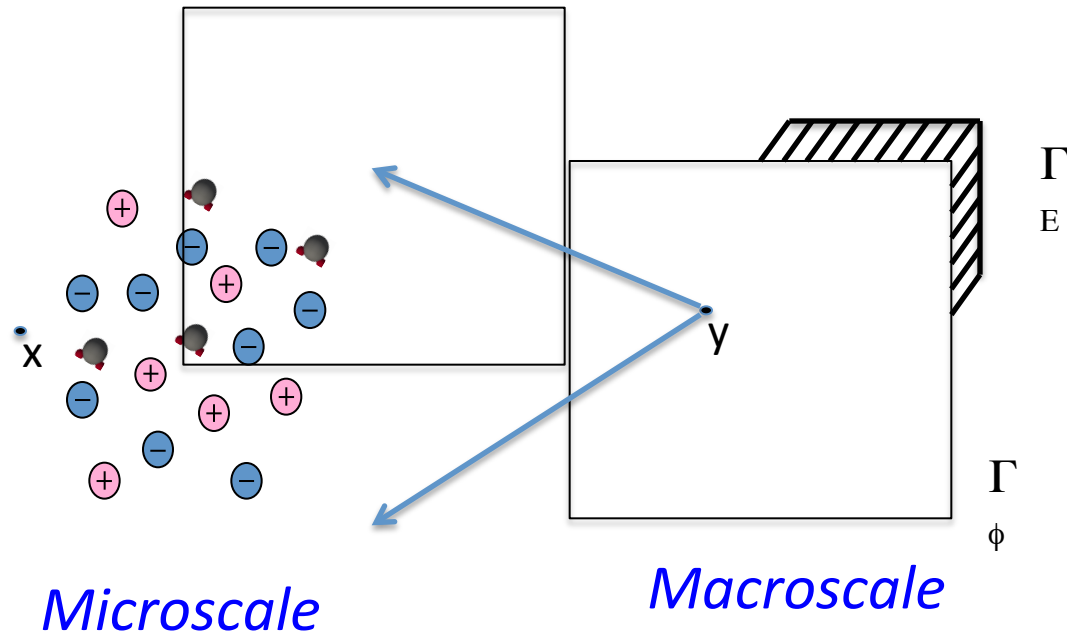
▪ **Constant:** gives only dipole moment

▪ **Linear:** gives up to quadrupole moment

Macroscopic

Quadrupole Moment

Descriptions of Scales



$$\nabla \cdot \mathbf{e}(\mathbf{x}, t) = \frac{1}{\epsilon_0} \rho^m(\mathbf{x}, t)$$

$$\nabla \cdot \mathbf{D}(\mathbf{y}, t) = \rho_f^M(\mathbf{y}, t)$$