

Exceptional service in the national interest



Active control of complex far infrared plasmonic structures



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Gregory C. Dyer
Sandia National Laboratories

Monday, October 28th, 2013



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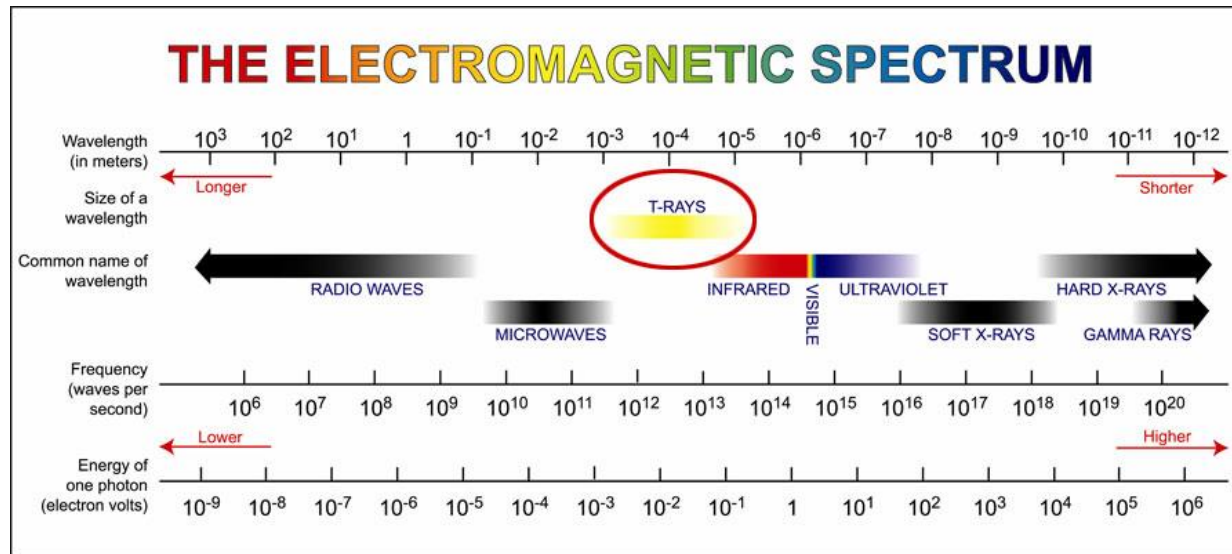
Outline

- Introduction
- Low dimensional far infrared plasmonic structures
 - Integrated 2D plasmonic devices
 - Strong coupling of tunable 2D plasmonic resonators
 - Generalized 2D plasmonic transmission line model
 - Active control of plasmonic crystal band structure
 - Electromagnetically induced transparency
- Conclusions & future directions

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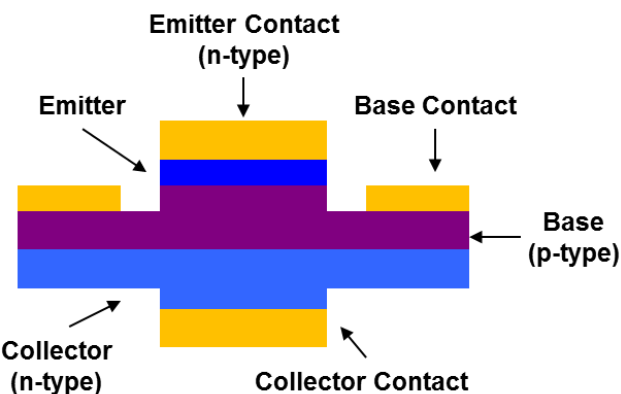
The 'THz Gap'



from
http://web.njit.edu/~barat/RBB_research.html

Electronics

RC limited



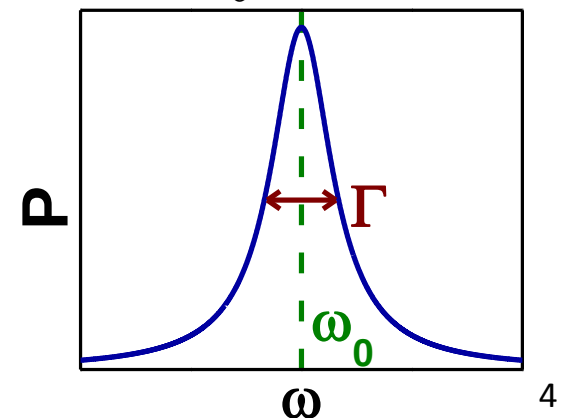
Photonics

kT limited

$$E_2 - E_1 = \hbar\omega \gg kT$$

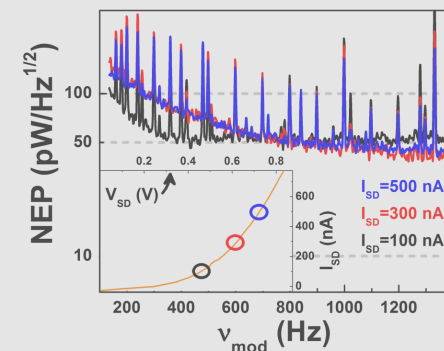
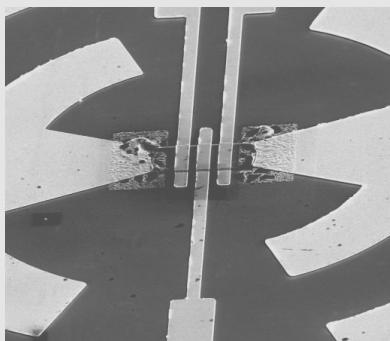
Plasmonics

$Q = \omega_0/\Gamma$ limited

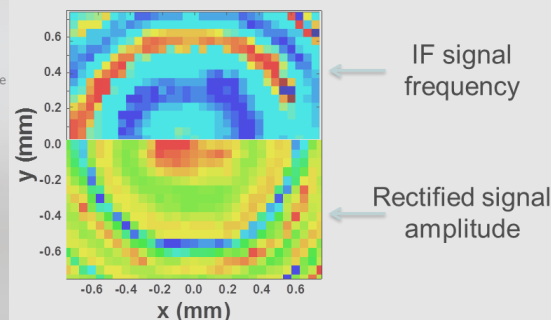
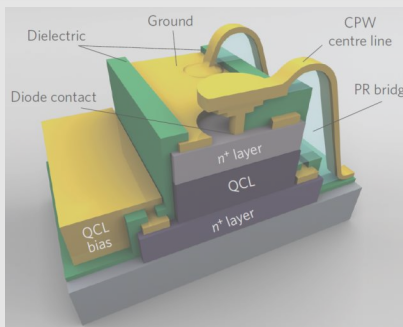


THz Electro-Optics at Sandia

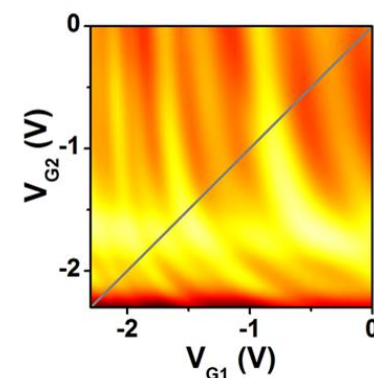
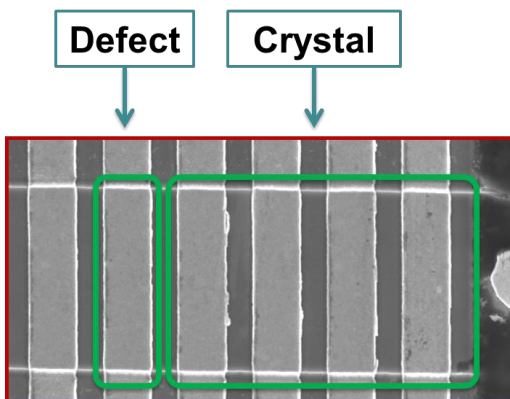
- Far infrared plasmonic detectors



- Quantum cascade laser transceivers

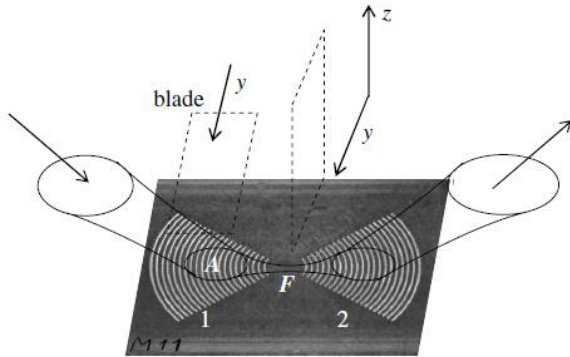


- Coherent low dimensional plasmonic phenomena

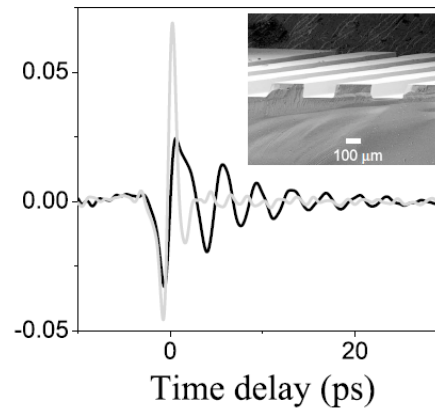


Far Infrared Plasmonics

Metal & Semiconductor Gratings

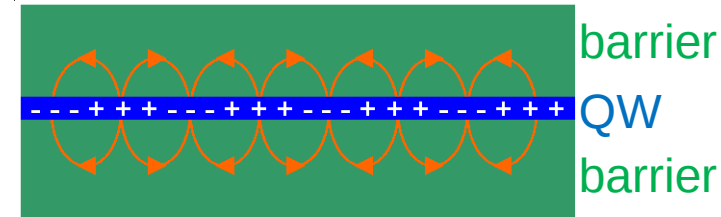


G. Gaborit *et al.*, *APL* **94**,
231108 (2009)

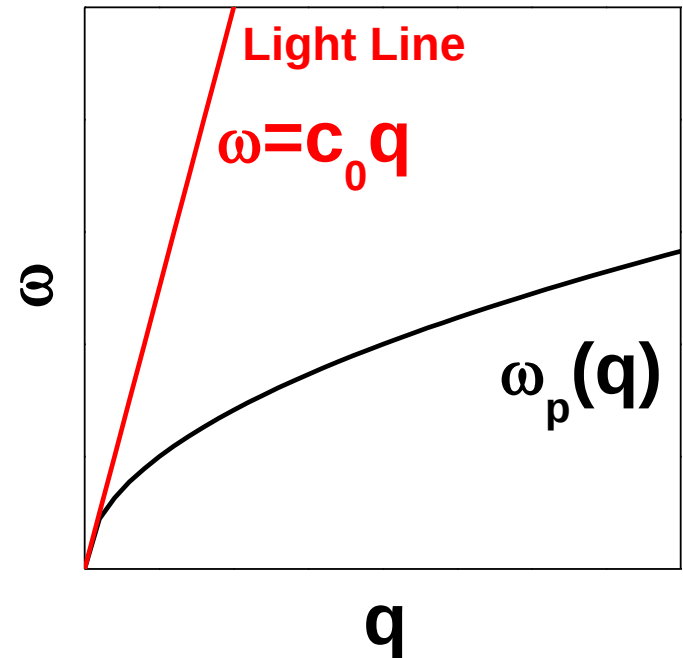


J. G. Rivas *et al.*, *PRL* **93**,
256804 (2004)

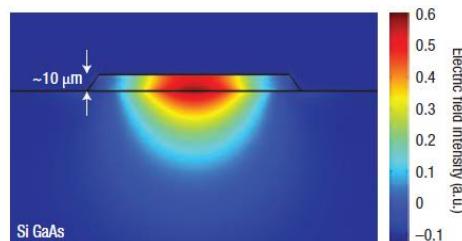
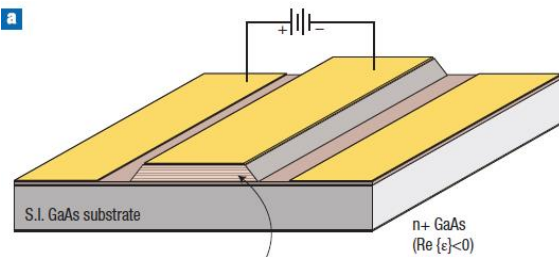
Semiconductor 2DEGs



$$\omega_p = \sqrt{n_{2D} e^2 q / 2 \epsilon m^*}$$

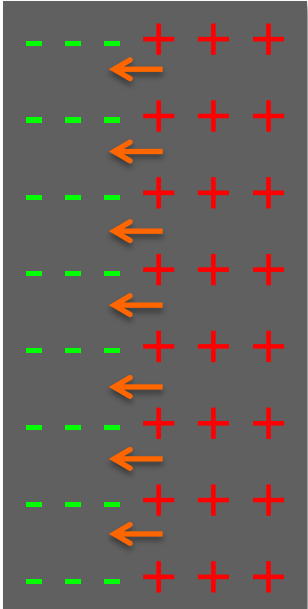
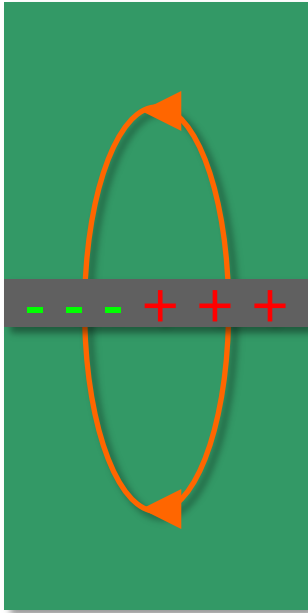
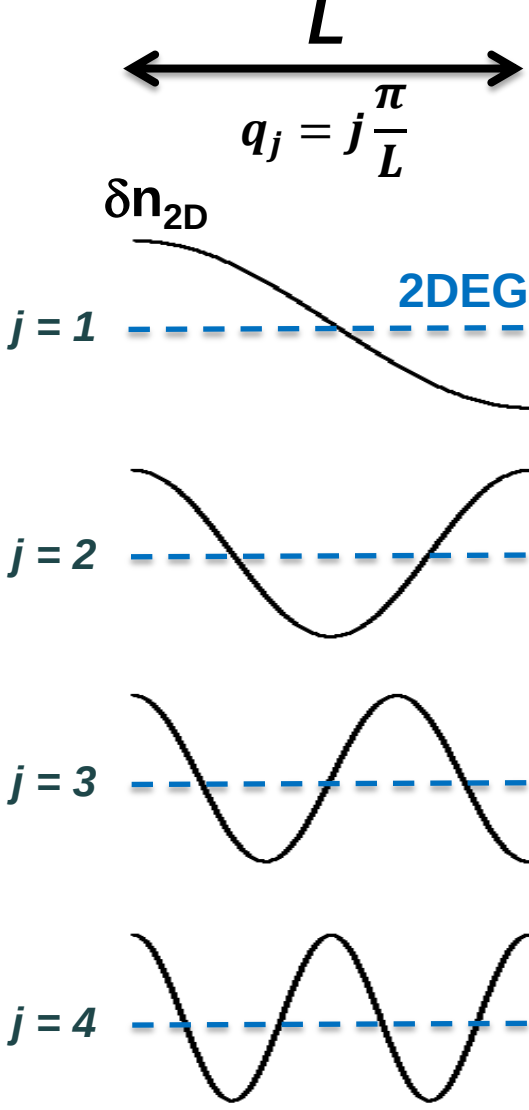


Quantum Cascade Laser Cavities



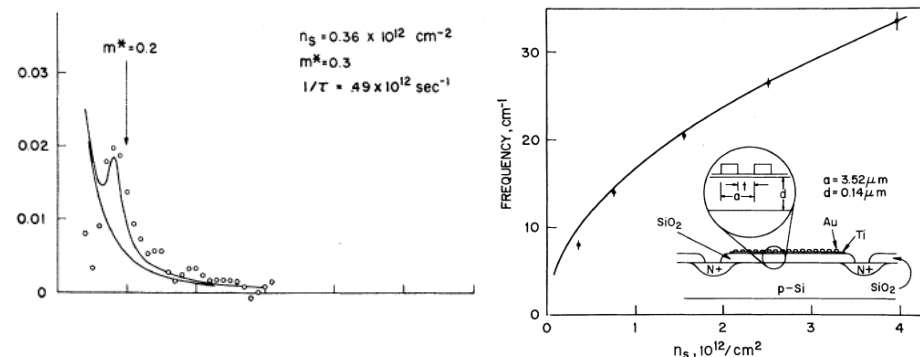
B. S. Williams, *Nat. Photonics* **1**, 517 (2007)

Bulk to 2D plasmons

	Bulk	2D	
Geometry			L $q_j = j \frac{\pi}{L}$ 
Units	$n_{3D} [m^{-3}]$	$n_{2D} [m^{-2}] \times q [m^{-1}]$	
Plasma Frequency	$\omega_p^{3D} = \sqrt{\frac{n_{3D} e^2}{\epsilon m^*}}$	$\omega_p^{2D} = \sqrt{\frac{n_{2D} e^2 q}{2 \epsilon m^*}}$	

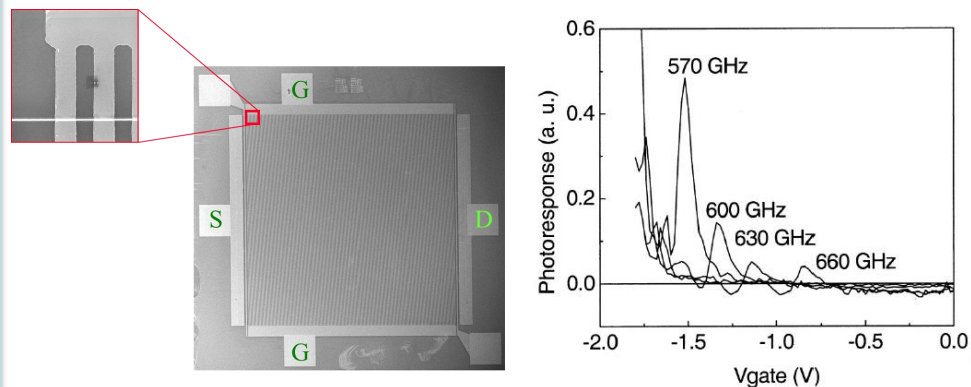
FIR 2D Plasmonic Systems

Grating Gated Si MOSFETs



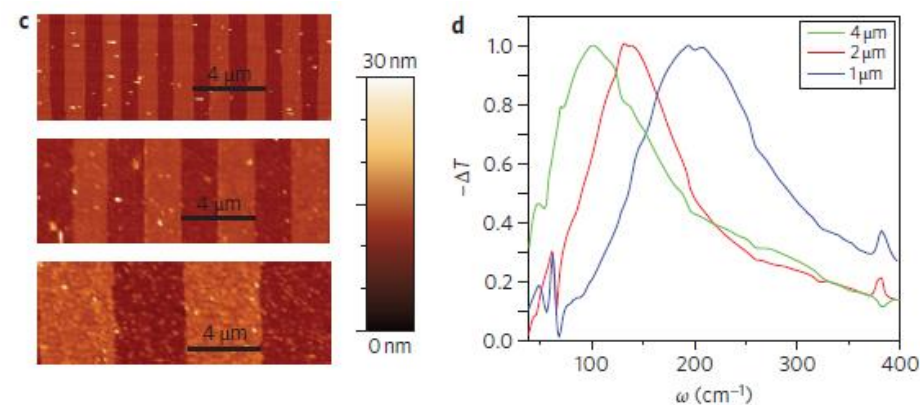
S.J. Allen *et al.*, *PRL* **38**, 980 (1977)

Grating Gated GaAs/AlGaAs HEMTs



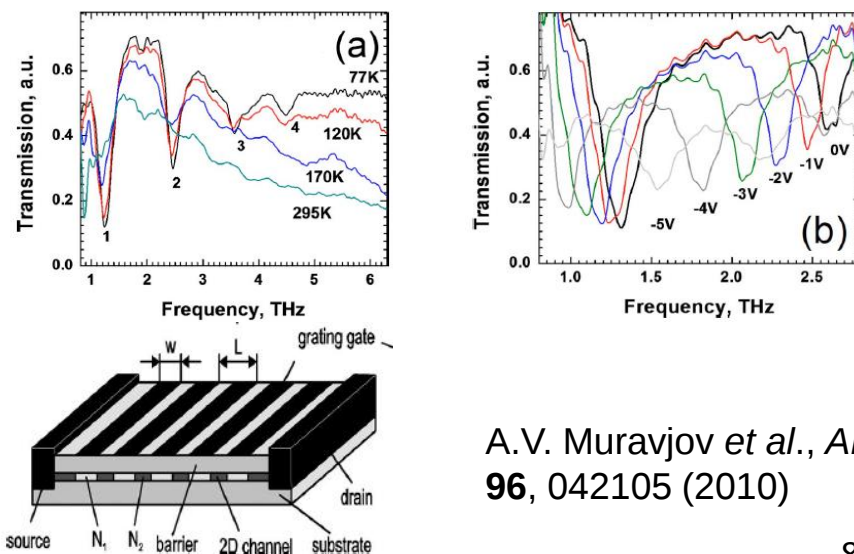
X.G. Peralta *et al.*, *APL* **81**, 1627 (2002)

Graphene Gratings



L. Ju *et al.*, *Nat. Nano.* **6**, 630 (2011)

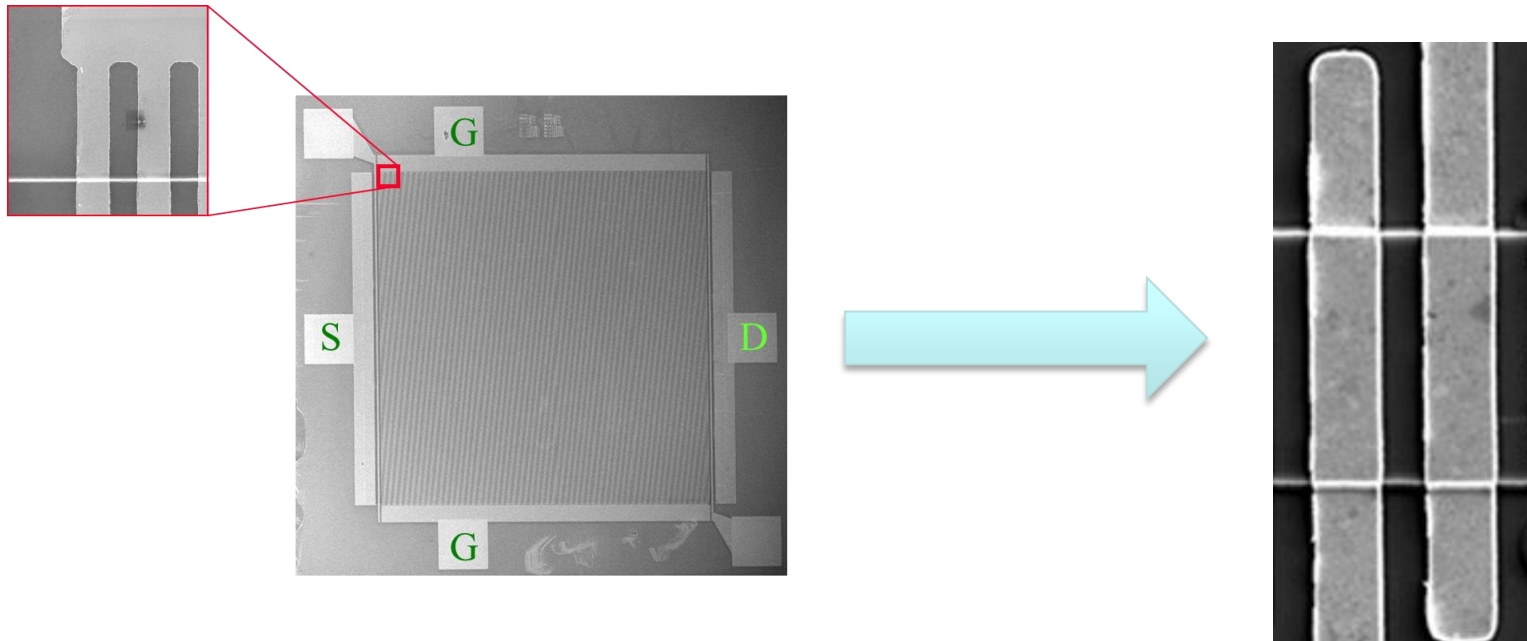
Grating Gated GaN/AlGaN HEMTs



A.V. Muravjov *et al.*, *APL* **96**, 042105 (2010)

From Plasmonic Arrays to Elements

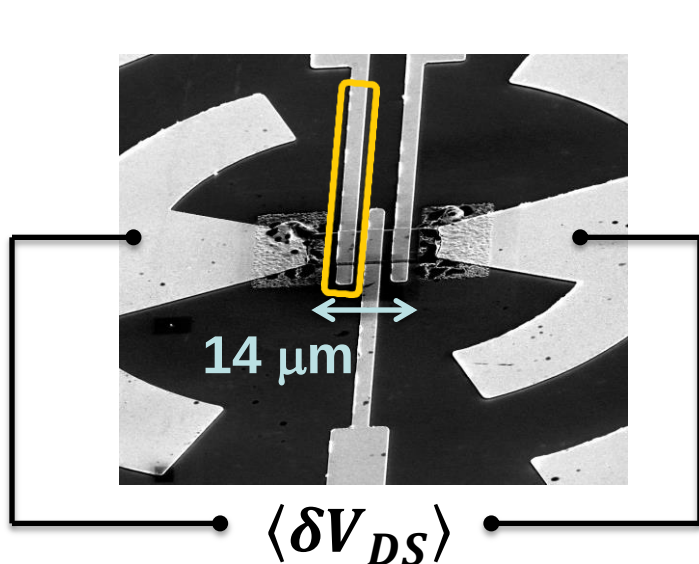
- Plasmon-plasmon coupling?
- Plasmonic dispersion in coupled structures?
- Plasmon damping & coherence length?



Outline

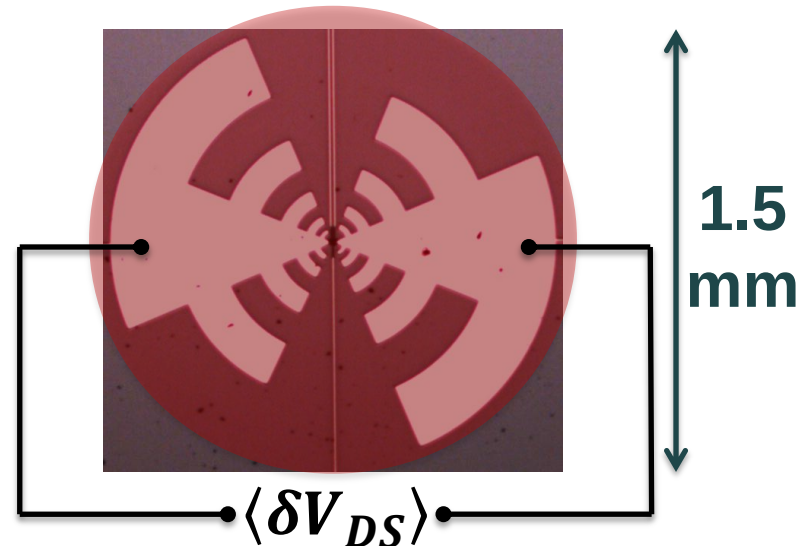
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Integrated 2D Plasmonic Devices

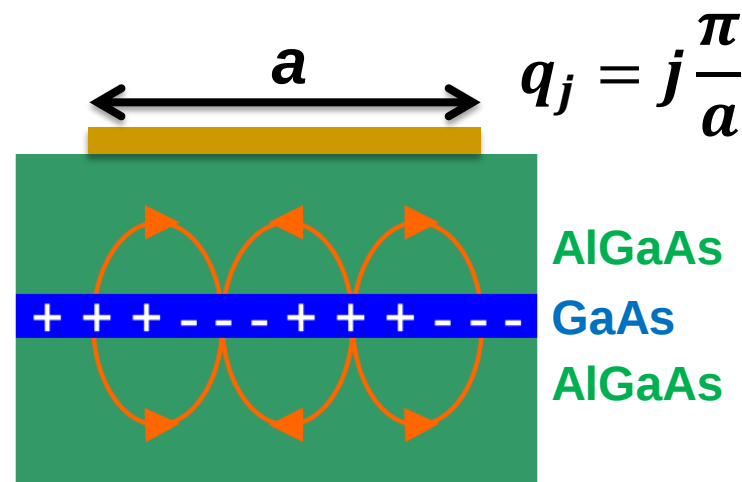


$\vec{E}$
 $\vec{S}$
 $\vec{B}$

0.1 - 0.5
THz

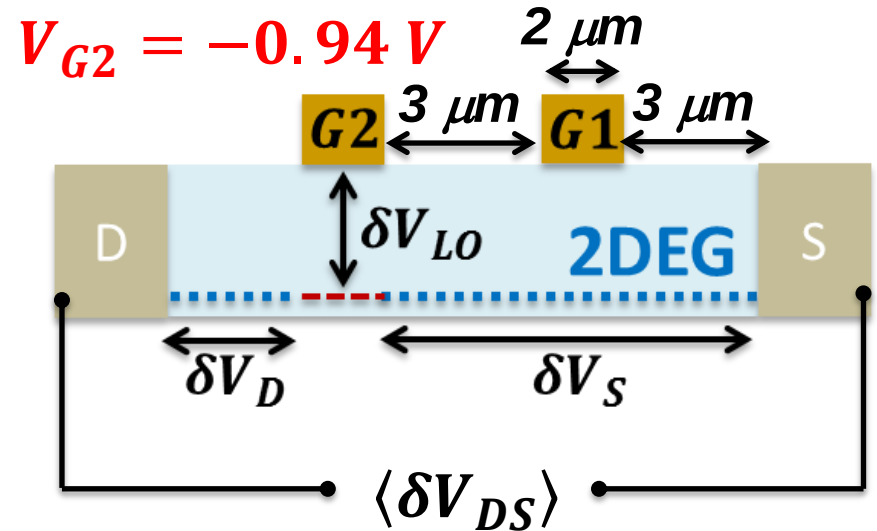
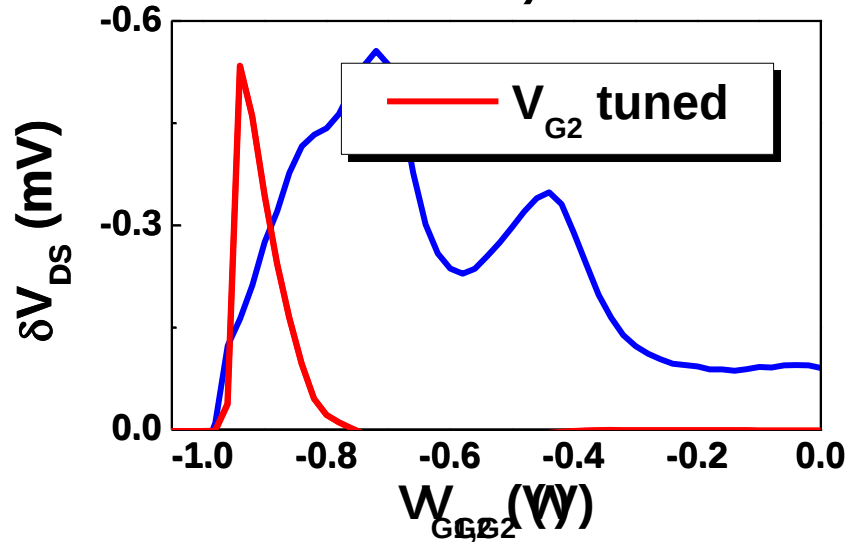


- **Detection:** depleted 2DEG
- **Coupling:** log-periodic antenna
- **Resonance:** 2D plasmon wavelength set by device geometry



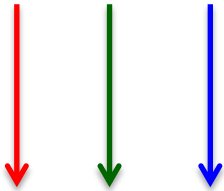
Single cavity & detector

270 GHz, 10 K

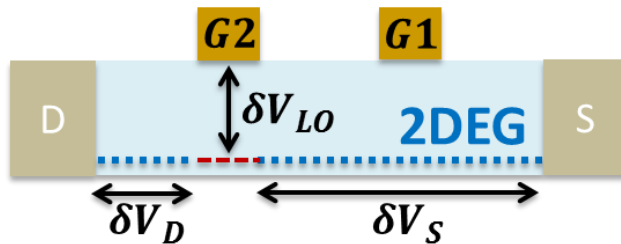
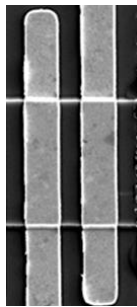


Plasmonic Spectroscopy

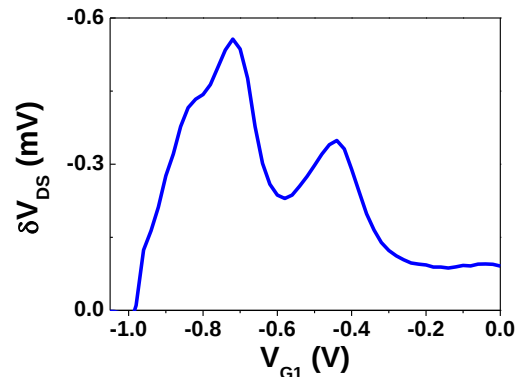
FIR Signal(s)



Detector + Plasmonic Cavity



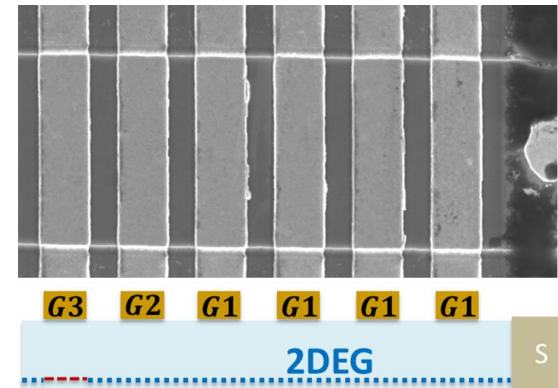
Spectrum of Incident Signal



Monochromatic FIR Source



Detector + Plasmonic Structure



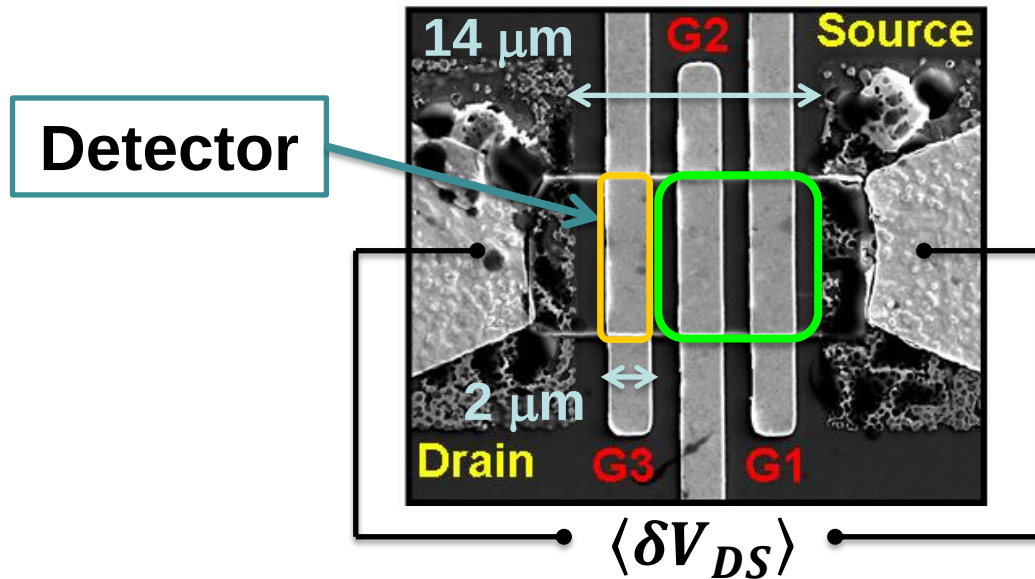
Spectrum of Plasmonic Structure

???

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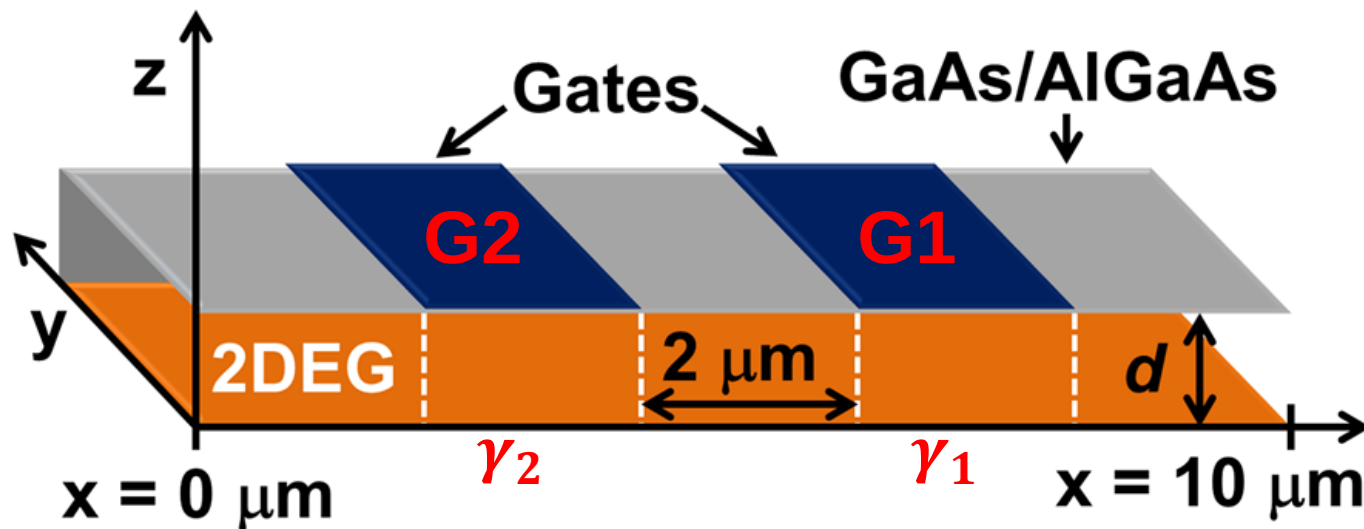
Inhomogeneous 2DEG Structures



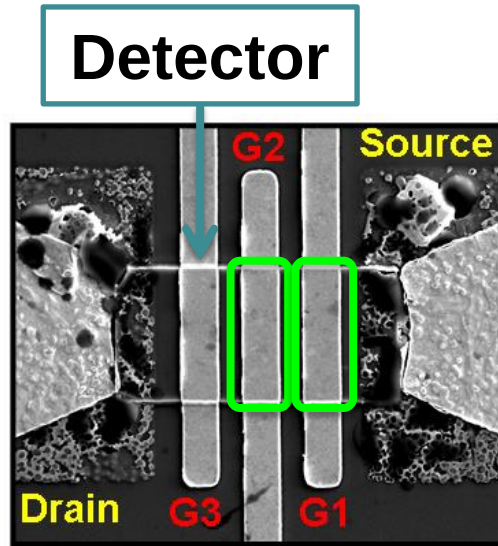
Normalized 2DEG Density

$$\gamma_1 = \frac{V_{th} - V_{G1}}{V_{th}}$$

$$\gamma_2 = \frac{V_{th} - V_{G2}}{V_{th}}$$

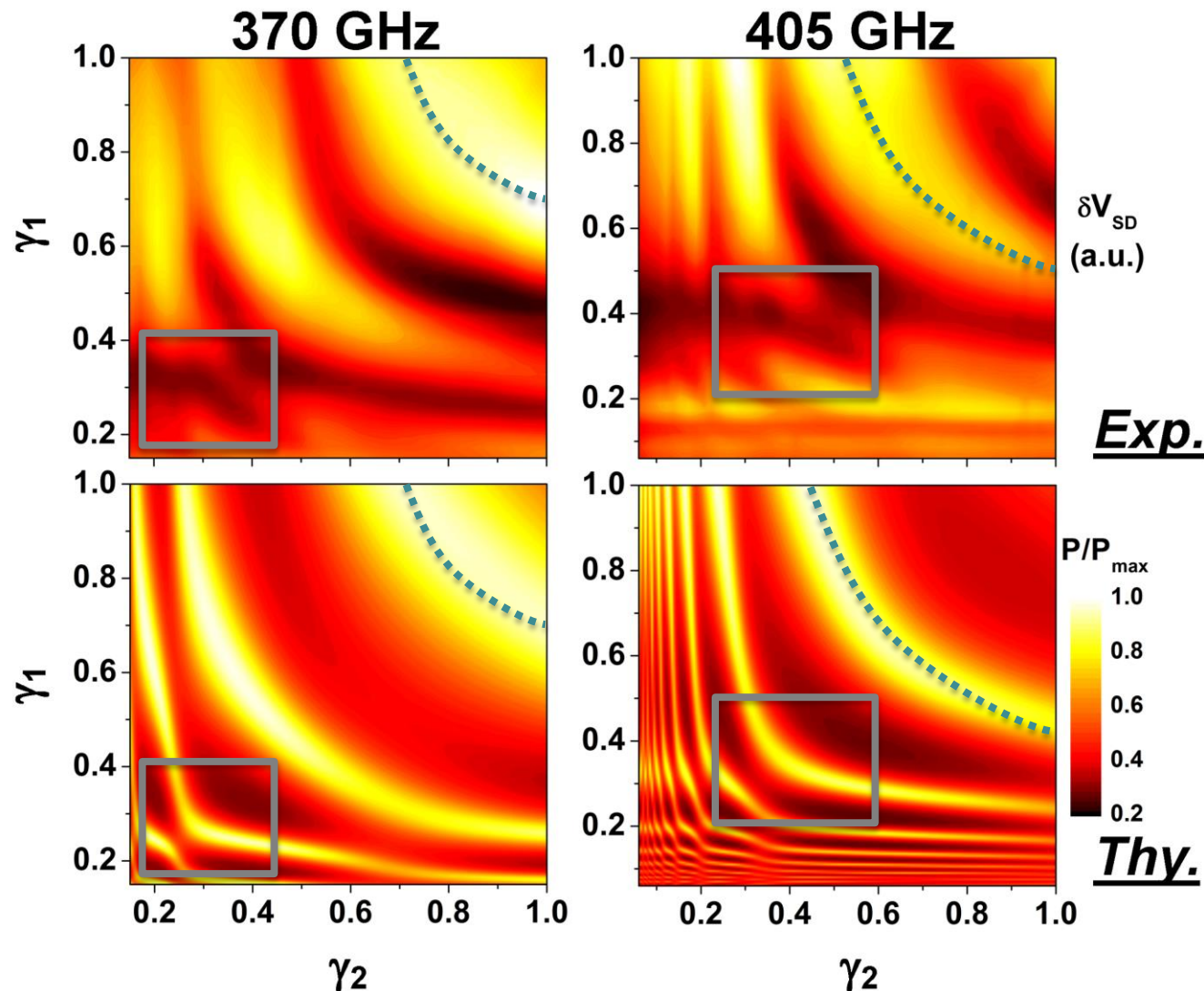


Coherent Plasmonic Coupling

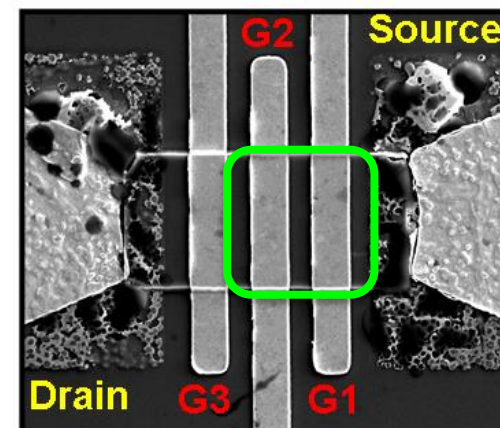
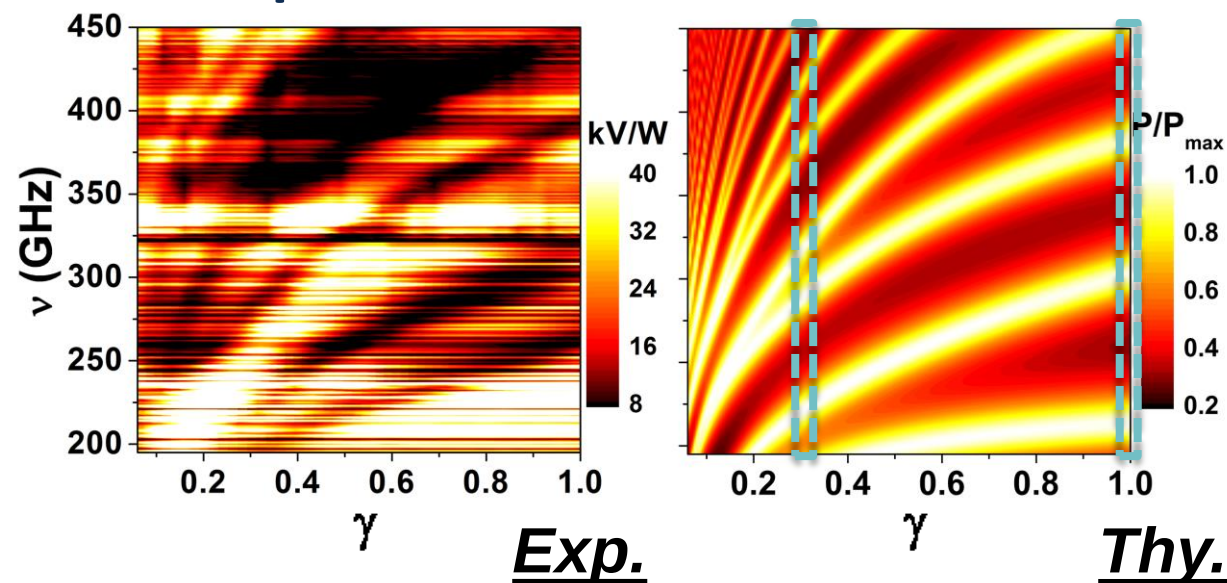


$$\gamma_{1,2} = \frac{V_{th} - V_{G1,G2}}{V_{th}}$$

Strongly coupled
2D plasmonic
resonators



Incipient Finite Plasmonic Crystal



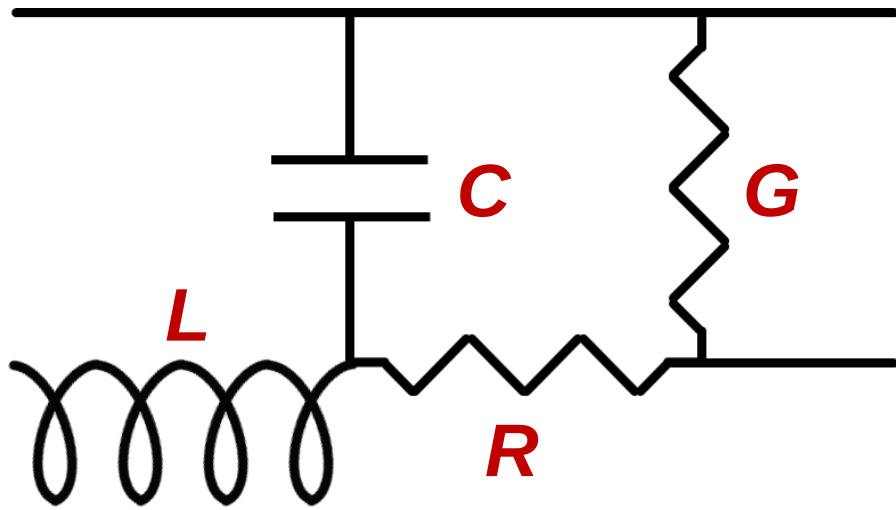
$$\gamma = \gamma_1 = \gamma_2$$

Two period
plasmonic crystal
embedded in 10 μm
plasmonic cavity

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Transmission Line (TL) Theory



Telegrapher's Equations

$$\frac{\partial V(x)}{\partial x} + \left(R + L \frac{\partial}{\partial t} \right) I(x) = 0$$

$$\frac{\partial I(x)}{\partial x} + \left(G + C \frac{\partial}{\partial t} \right) V(x) = 0$$

Wave Solutions

$$V(x) = V_+ e^{-iqx} + V_- e^{+iqx}$$

$$I(x) = \frac{1}{Z_0} (V_+ e^{-iqx} - V_- e^{+iqx})$$

Characteristic Z & Dispersion

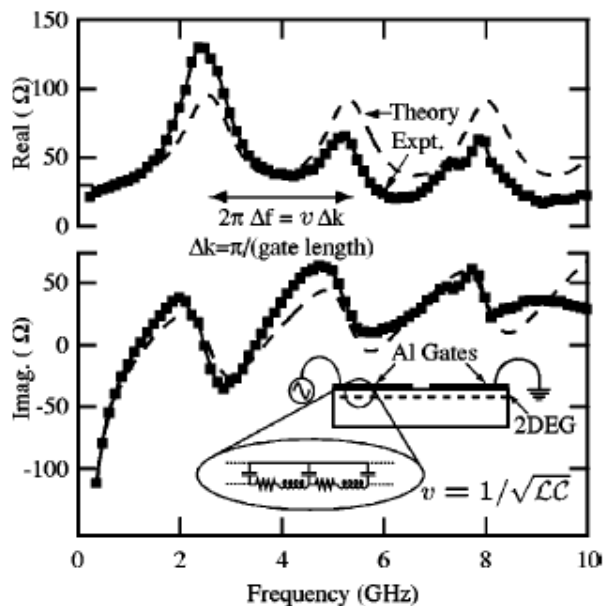
$$Z_0 = \sqrt{\frac{R + i\omega L}{G + i\omega C}}$$

$$iq = \sqrt{(R + i\omega L)(G + i\omega C)}$$

Plasmonic TL Model

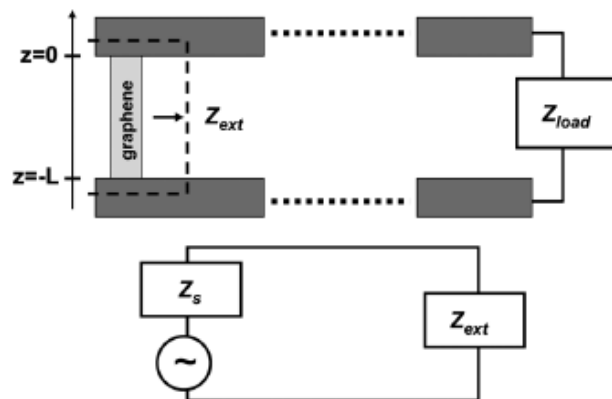
High frequency conductivity of the high-mobility two-dimensional electron gas

Graphene Terahertz Plasmon Oscillators



$$Z_0 = \sqrt{\frac{Z_{2DEG}}{i\omega C(\omega)}}$$

$$iq = \sqrt{i\omega C(\omega)Z_{2DEG}}$$



$$Z_{2DEG} = \frac{m^*}{e^2 n_{2D} \tau W} (1 + i\omega\tau)$$

from Drude model $\text{Im}[Z_{2DEG}]$
is kinetic inductance

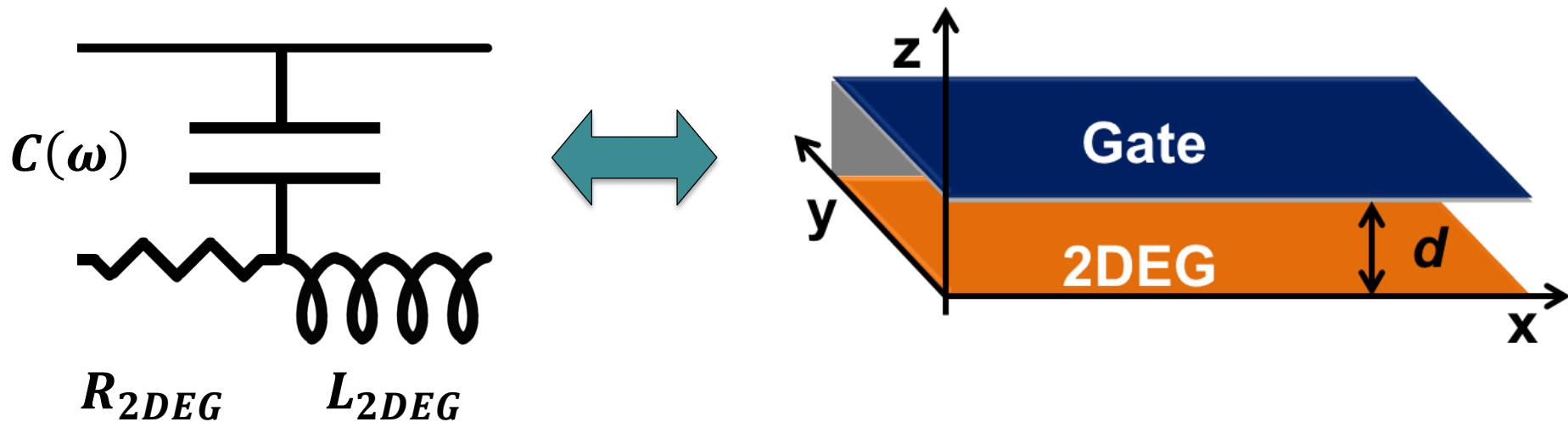
$$\lim_{qd \rightarrow \infty} C(\omega) = 2W\epsilon q$$

(unscreened limit)

$$\lim_{qd \rightarrow 0} C(\omega) = \frac{W\epsilon}{d}$$

(fully screened limit)

Generalized Plasmonic TL Model

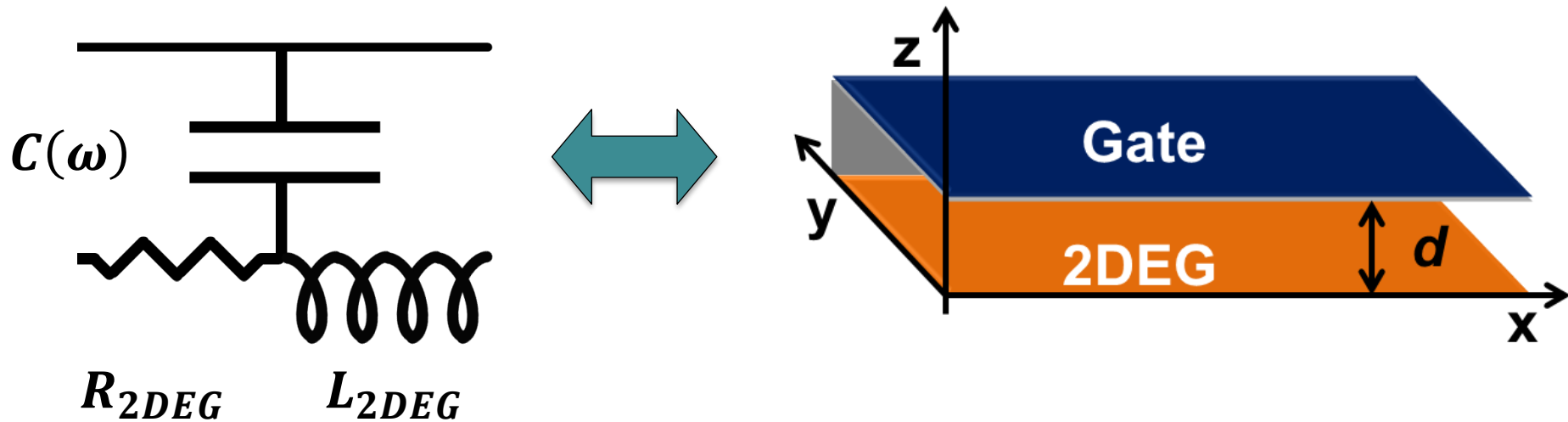


$$C(\omega) = \frac{\delta n(x)}{\delta \Phi(x, z = 0)} = W \epsilon q (1 + \coth[qd])$$

$$\frac{1}{2} \tilde{V}(x) \tilde{I}^*(x) = \frac{1}{2} \Phi(x, 0) I^*(x) \chi^*(\omega, d)$$

$$\begin{aligned} \tilde{V}(x) &\equiv \Phi(x, 0) \\ \tilde{I}(x) &\equiv I(x) \chi(\omega, d) \end{aligned}$$

Generalized Plasmonic TL Model



$$C(\omega) = W\epsilon q(1 + \coth[qd]) \frac{m^*}{e^2 n_{2D} W}$$

$$L_{2DEG} = \frac{m^*}{e^2 n_{2D} W}$$

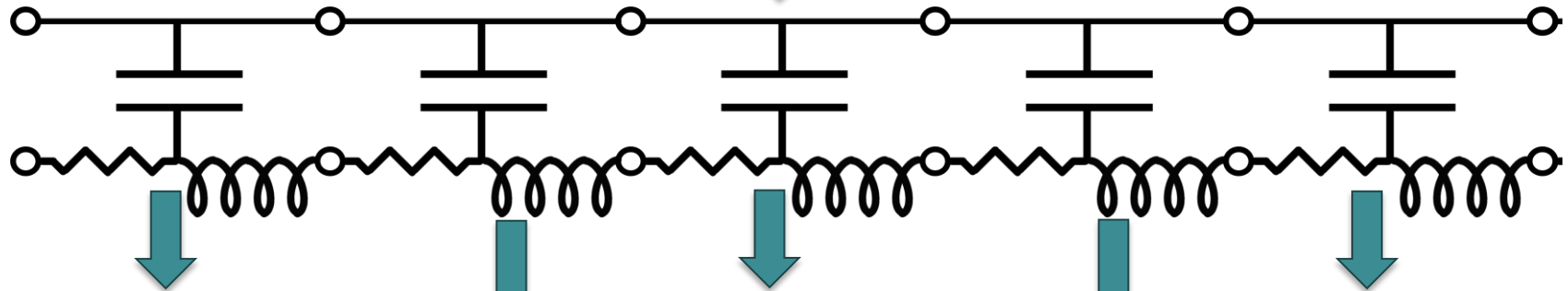
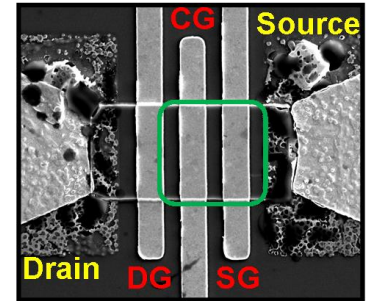
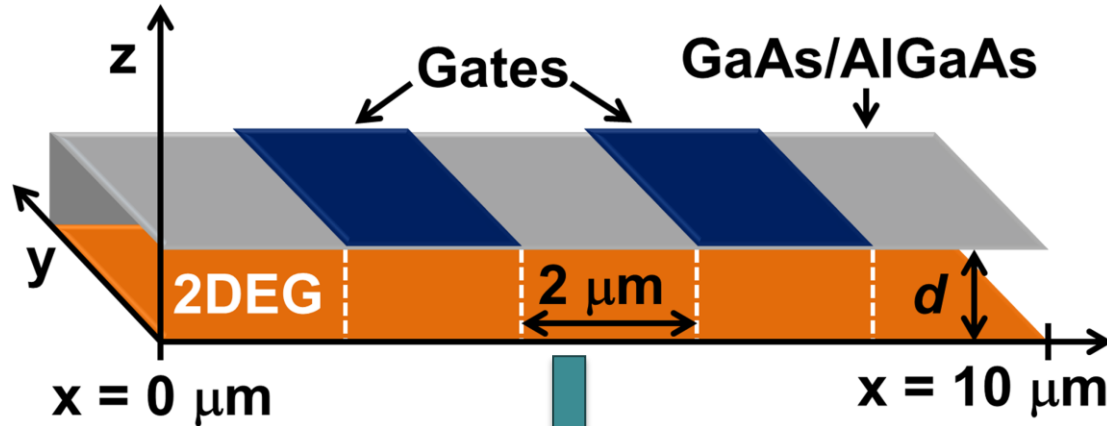
$$R_{2DEG} = L_{2DEG} / \tau$$

$$Z = \frac{1}{\chi(\omega, d)} \sqrt{\frac{R_{2DEG} + i\omega L_{2DEG}}{i\omega C(\omega)}}$$

$$q = -i\sqrt{i\omega C(\omega)(R_{2DEG} + i\omega L_{2DEG})}$$

where: $\chi(\omega, d) = 1 - \frac{q(1 - e^{-2q'd} \cos[2q''d])}{2q'(1 - e^{-2q^*d})} + \frac{q e^{-2q'd} \sin[2q''d]}{2q''(1 - e^{-2q^*d})}$

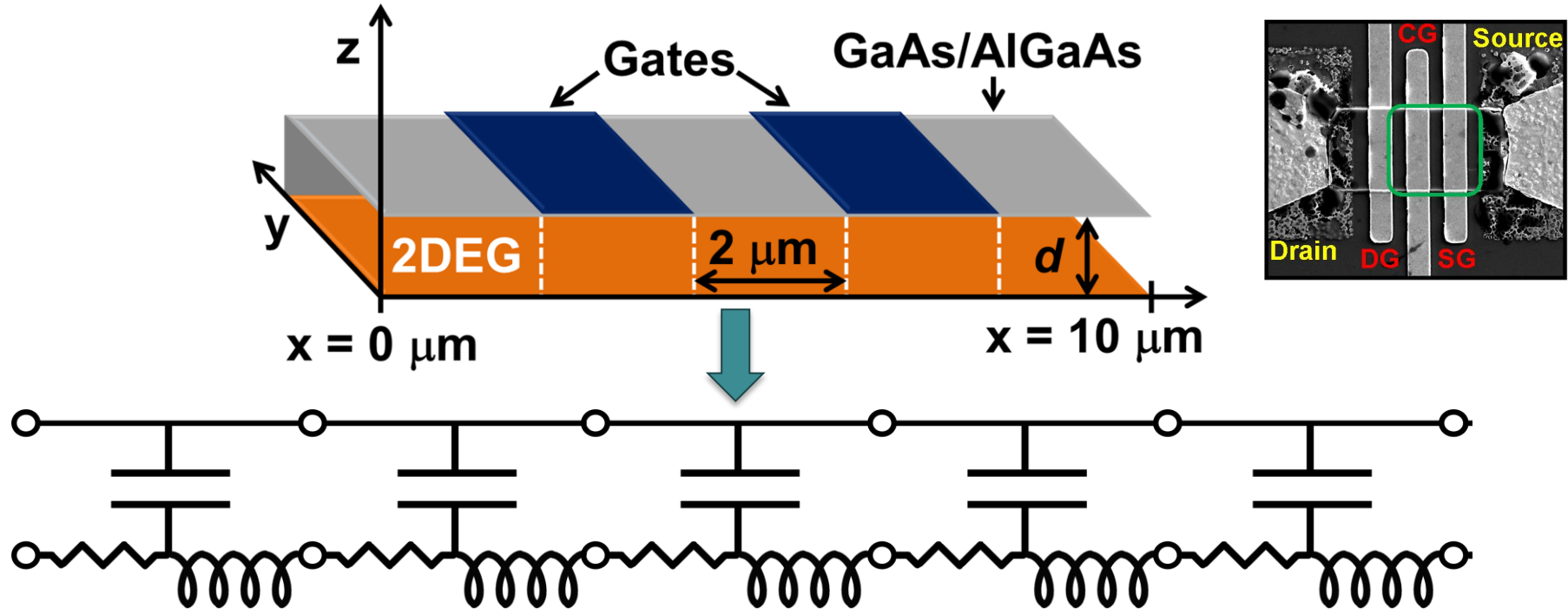
Discretized Plasmonic TL



$$\begin{pmatrix} \cos[q_1 a_1] & i\tilde{Z}_1 \sin[q_1 a_1] \\ i\sin[q_1 a_1]/\tilde{Z}_1 & \cos[q_1 a_1] \end{pmatrix} \quad \begin{pmatrix} \cos[q_3 a_3] & i\tilde{Z}_3 \sin[q_3 a_3] \\ i\sin[q_3 a_3]/\tilde{Z}_3 & \cos[q_3 a_3] \end{pmatrix} \quad \begin{pmatrix} \cos[q_5 a_5] & i\tilde{Z}_5 \sin[q_5 a_5] \\ i\sin[q_5 a_5]/\tilde{Z}_5 & \cos[q_5 a_5] \end{pmatrix}$$

$$\begin{pmatrix} \cos[q_2 a_2] & i\tilde{Z}_2 \sin[q_2 a_2] \\ i\sin[q_2 a_2]/\tilde{Z}_2 & \cos[q_2 a_2] \end{pmatrix} \quad \begin{pmatrix} \cos[q_4 a_4] & i\tilde{Z}_4 \sin[q_4 a_4] \\ i\sin[q_4 a_4]/\tilde{Z}_4 & \cos[q_4 a_4] \end{pmatrix}$$

Discretized Plasmonic TL



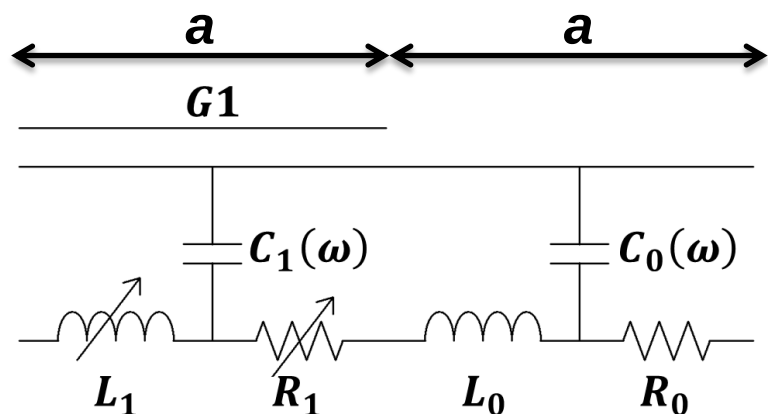
Solution for an inhomogeneous cavity with step-like boundaries between TL elements is reduced to linear algebra:

$$\begin{pmatrix} \tilde{V}_m \\ \tilde{I}_m \end{pmatrix} = \left[\prod_m \begin{pmatrix} \cos[q_m a_m] & i\tilde{Z}_m \sin[q_m a_m] \\ i\sin[q_m a_m] / \tilde{Z}_m & \cos[q_m a_m] \end{pmatrix} \right] \begin{pmatrix} \tilde{V}_{n+1} \\ \tilde{I}_{n+1} \end{pmatrix}$$

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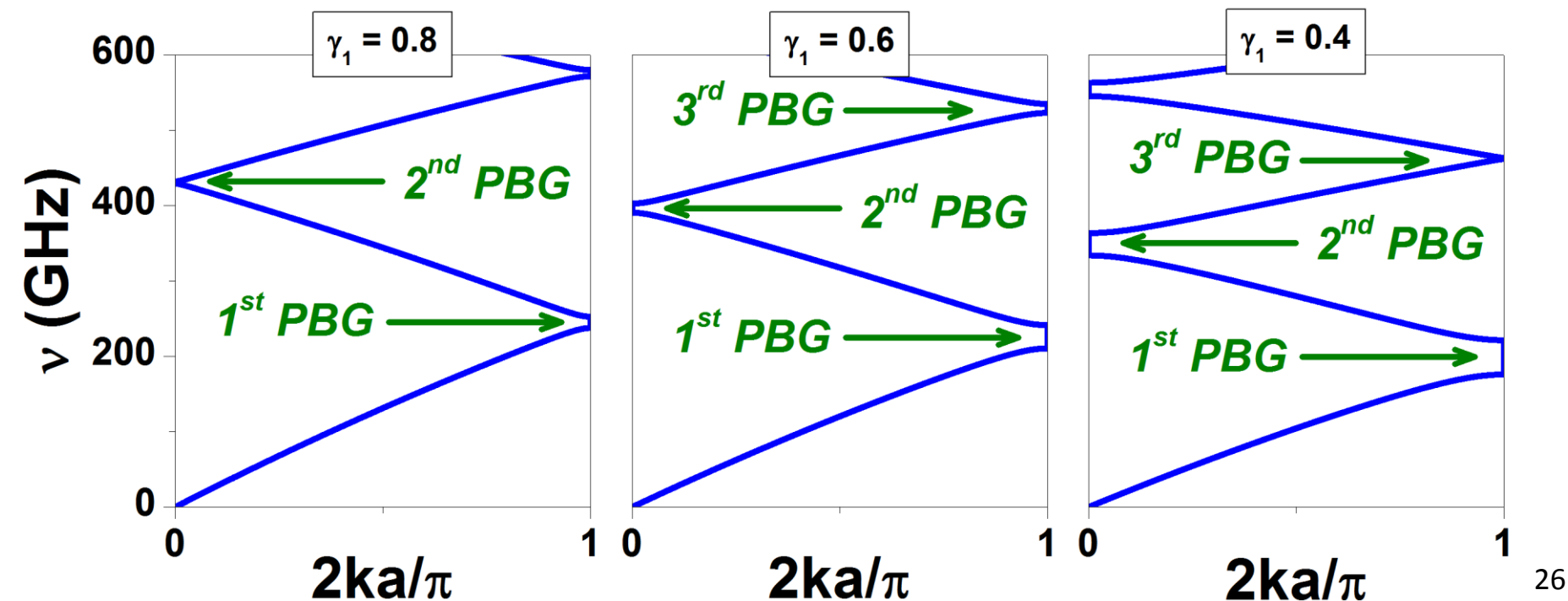
Plasmonic Crystal Band Structure



1D Kronig-Penney:

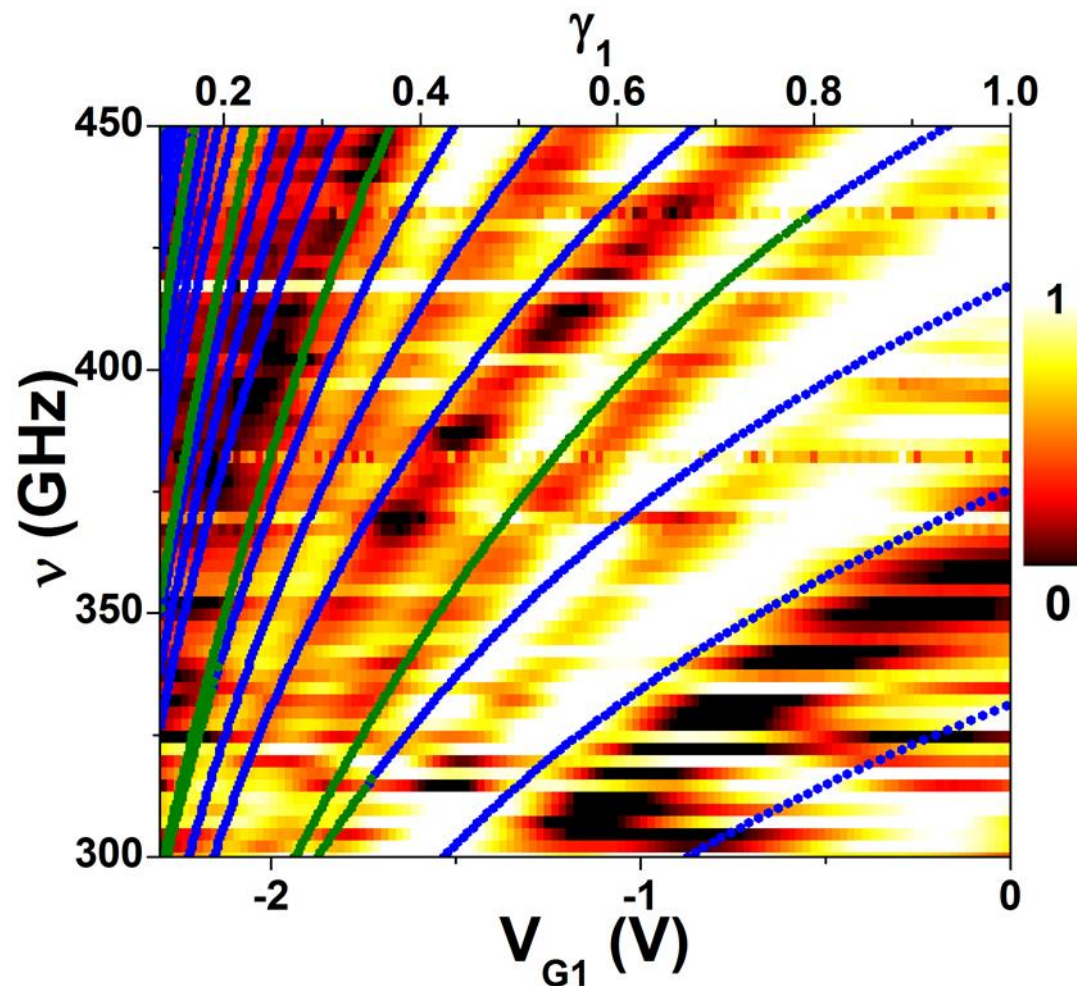
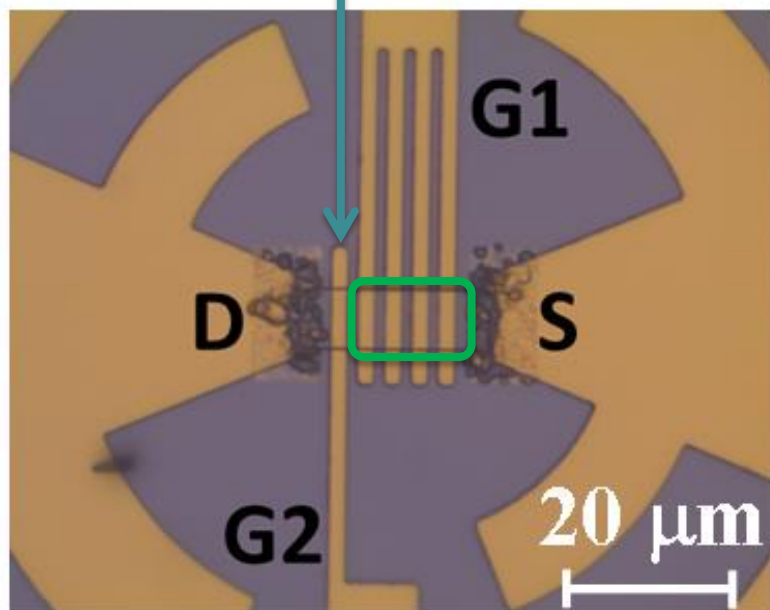
$$\cos 2k_B a =$$

$$\cos q_0 a \cos q_1 a - \frac{1}{2} \left(\frac{Z_0}{Z_1} + \frac{Z_1}{Z_0} \right) \sin q_0 a \sin q_1 a$$



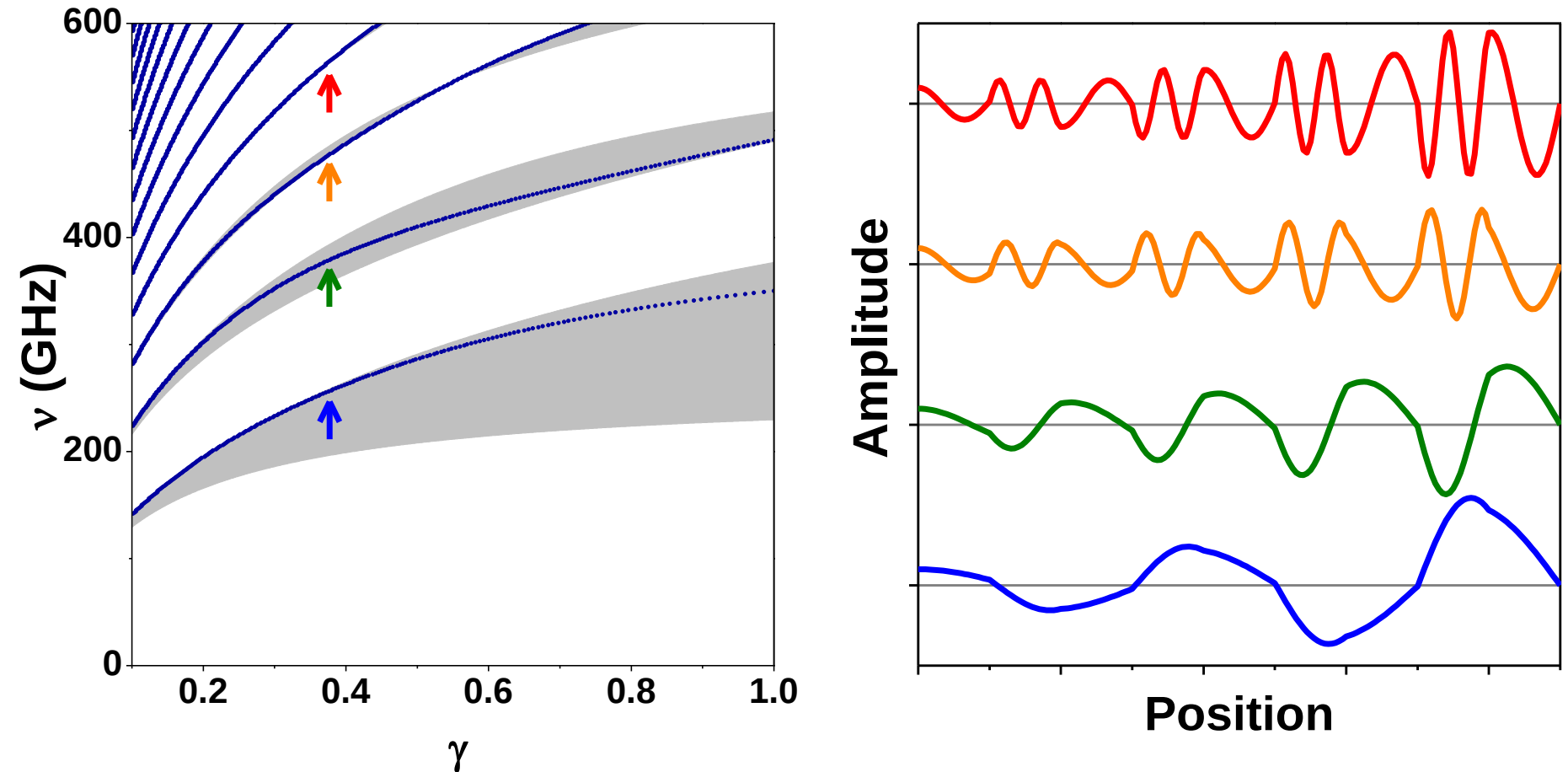
Four Period Plasmonic Crystal

Detector



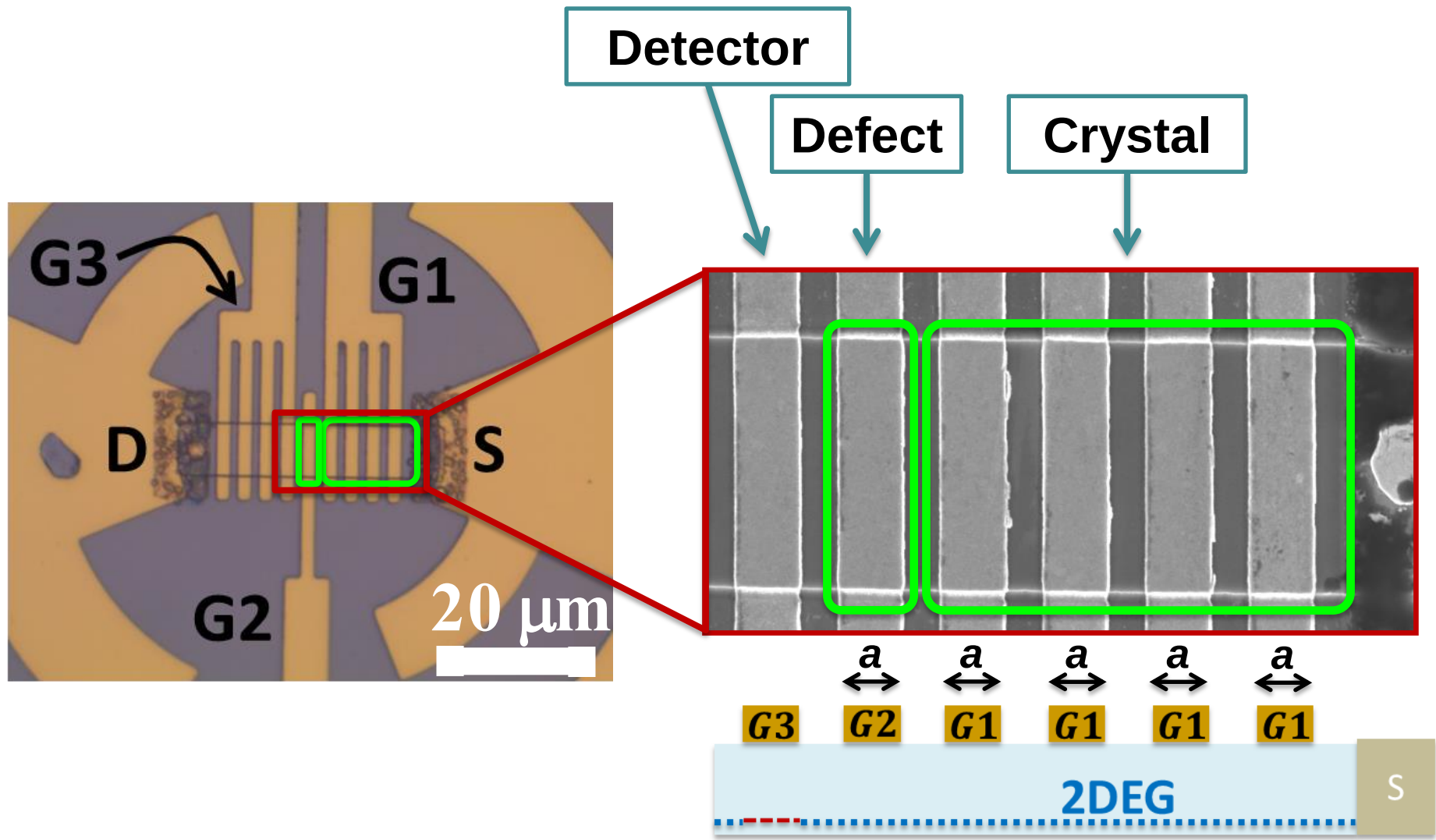
$$V_{G1} \rightarrow 0, L_1 \rightarrow \infty$$

Plasmonic Tamm States

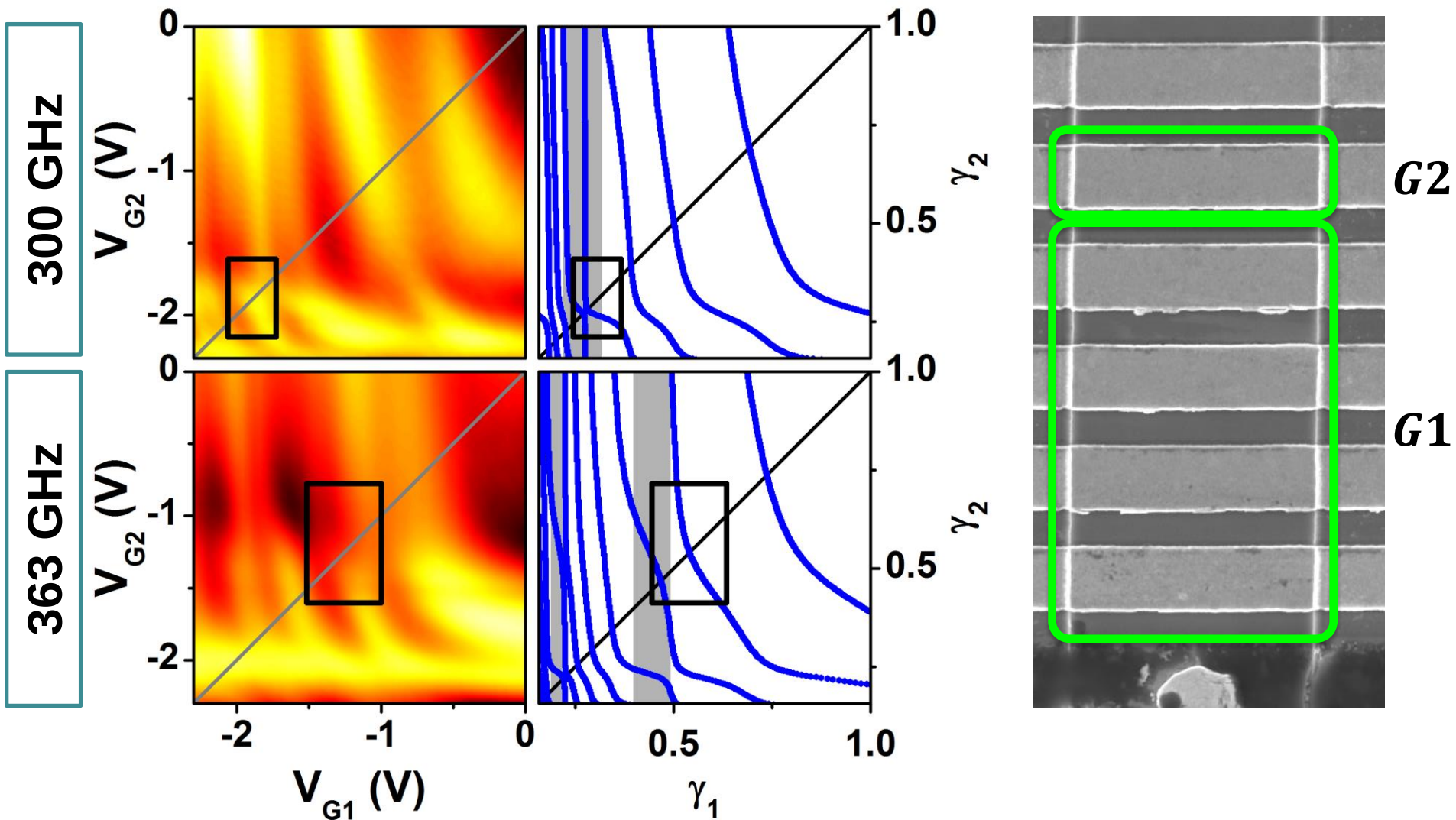


Crystal surface states typically found in crystal band gap that have complex Bloch wavevector ($Im[k_B] \neq 0$)

Defect Probing of Tamm States



Coupled Tamm & Defect States



Paired Tamm States

On the Surface States Associated with a Periodic Potential

WILLIAM SHOCKLEY

Bell Telephone Laboratories, New York, New York

(Received June 19, 1939)

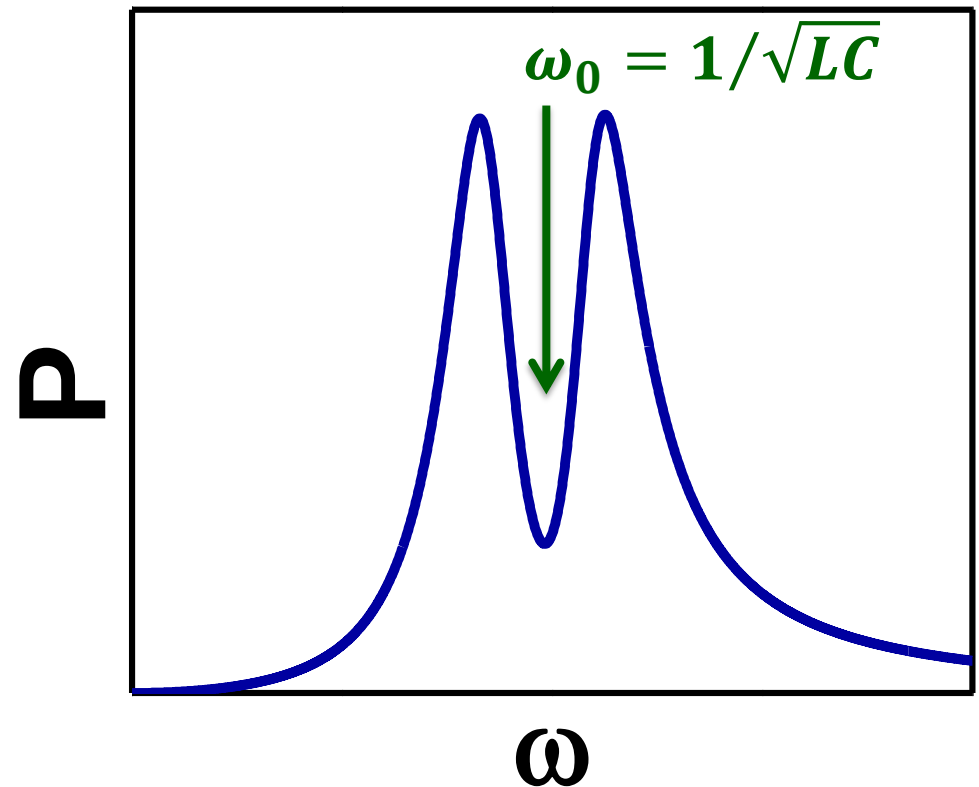
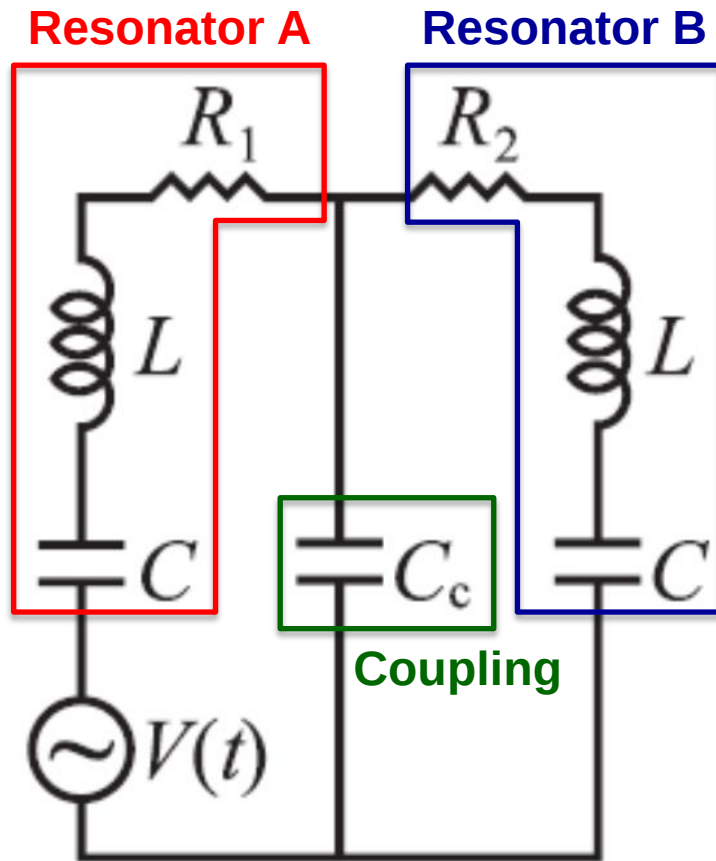
Tamm found in this case that it was possible to have energy levels whose wave functions were localized at the surface of the crystal. In this treatment a semi-infinite crystal was used, and Tamm found that one surface level was possible for each energy gap between the ordinary allowed bands of energies. In a later paper Fowler³ discussed Tamm's levels for a finite crystal and pointed out that the levels should occur in pairs because of the two faces of the crystal. They

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 - Electromagnetically induced transparency
- Conclusions & future directions

Classical Analogue of EIT

“Low-Loss Metamaterials Based on
Classical Electromagnetically
Induced Transparency”

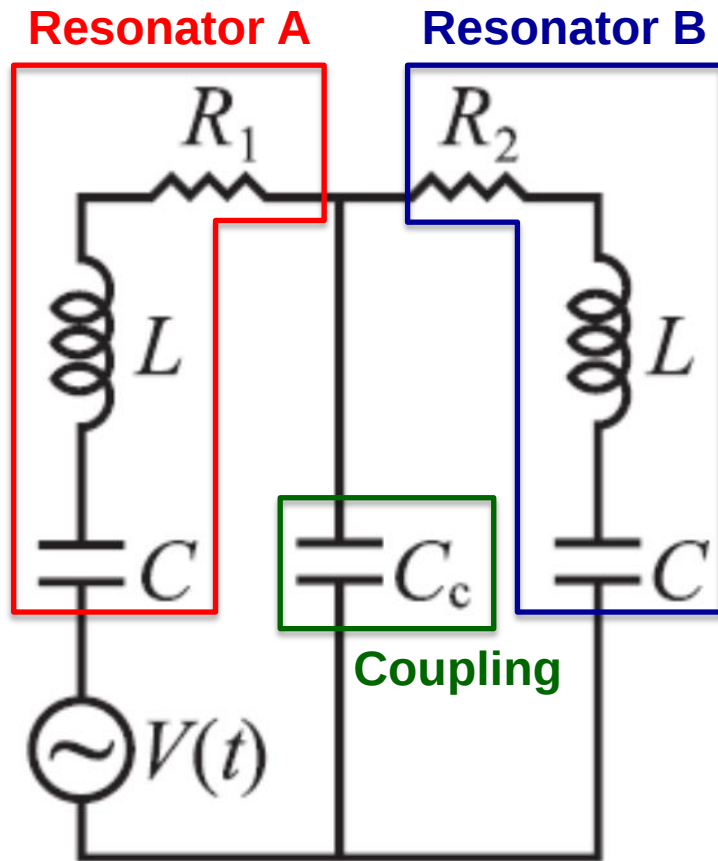


$$P = \text{Re} \frac{i\omega \frac{V^2}{2} (L\omega^2 - iR_2\omega - \frac{1}{C_{\parallel}})}{\frac{1}{C_c^2} - (L\omega^2 - iR_2\omega - \frac{1}{C_{\parallel}})(L\omega^2 - iR_1\omega - \frac{1}{C_{\parallel}})}$$

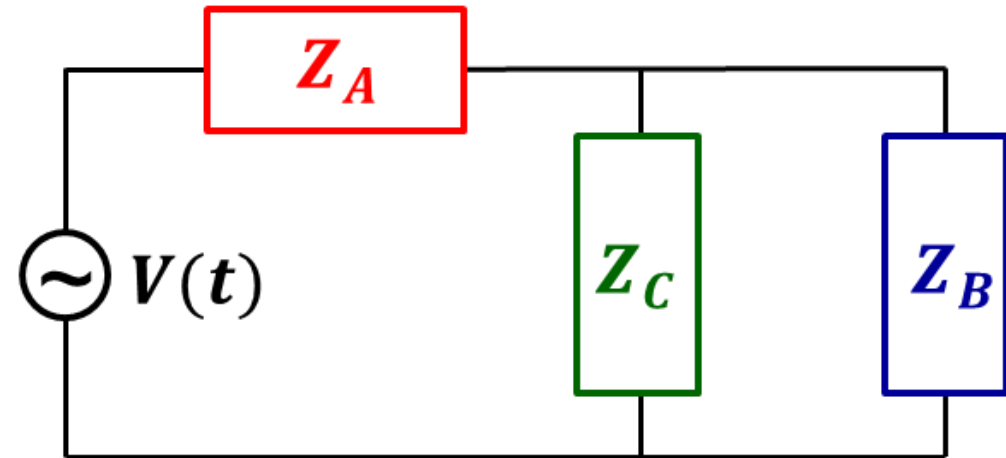
with $C_{\parallel} = (1/C + 1/C_c)^{-1}$

Classical Analogue of EIT

“Low-Loss Metamaterials Based on
Classical Electromagnetically
Induced Transparency”



$$Z_A = R_1 + \frac{1 - \omega^2 LC}{i\omega C}$$



$$Z_C = \frac{1}{i\omega C_c} \quad Z_B = R_2 + \frac{1 - \omega^2 LC}{i\omega C}$$

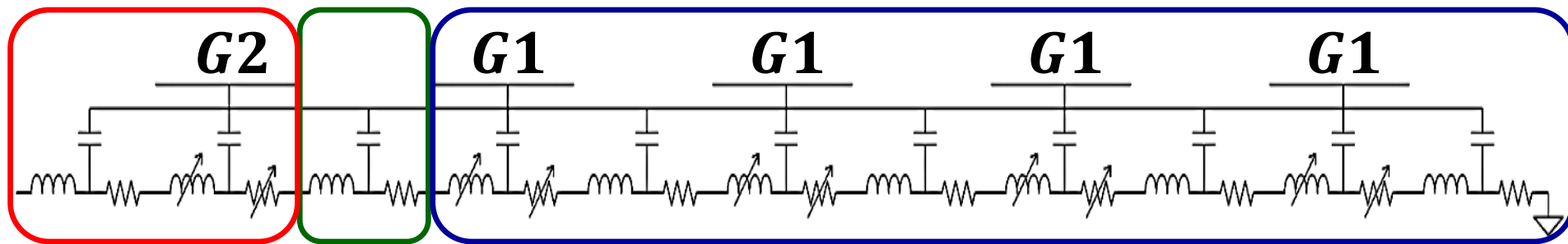
$$Z_{eq} = Z_A + \frac{Z_B Z_C}{Z_B + Z_C}$$

Plasmonic Crystal-Defect EIT



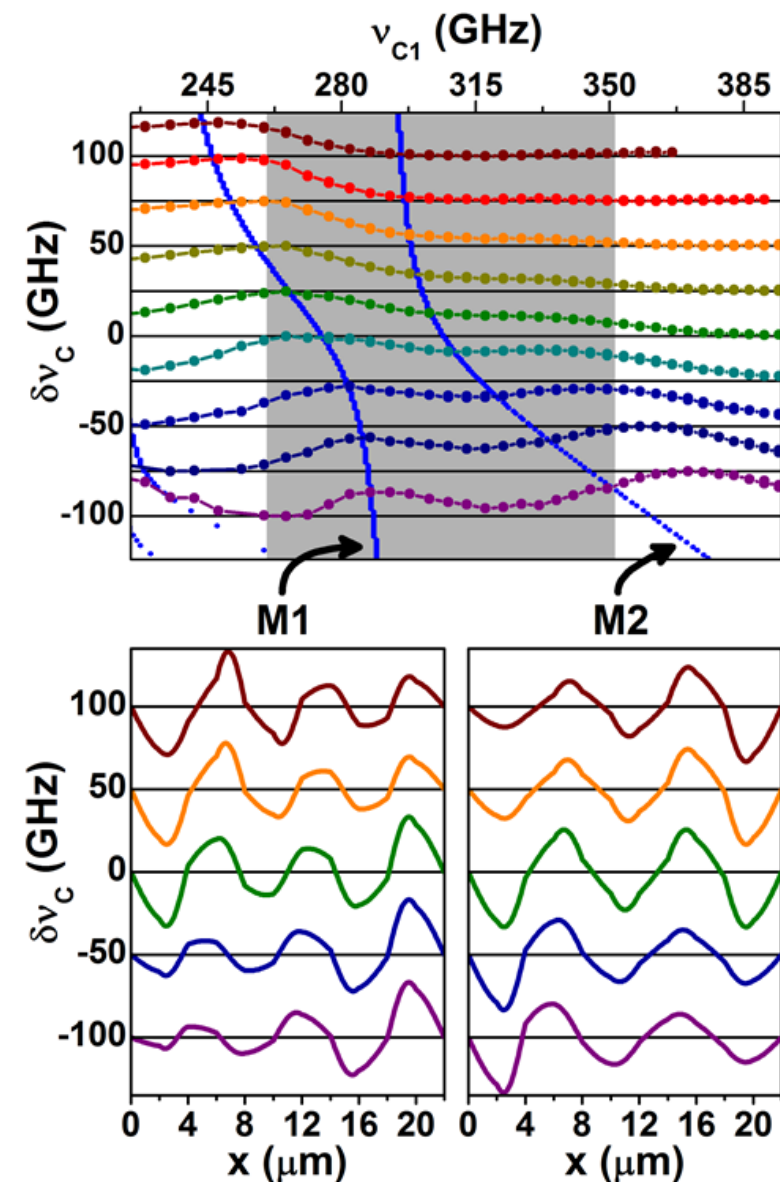
Resonator A:
Defect Mode

Resonator B:
Localized Crystal Tamm State



Coupling

Active Control of Plasmonic EIT

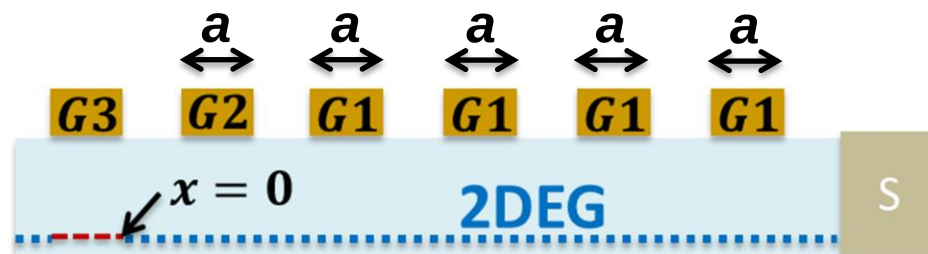


Crystal-defect eigenmodes in or near crystal band gap at **210 GHz**

$$\nu_{Cj} = \frac{\beta_C}{2\pi\sqrt{L_j C_j}}$$

$$\beta_C = \pi/a$$

$$\delta\nu_C = \nu_{C2} - \nu_{C1}$$

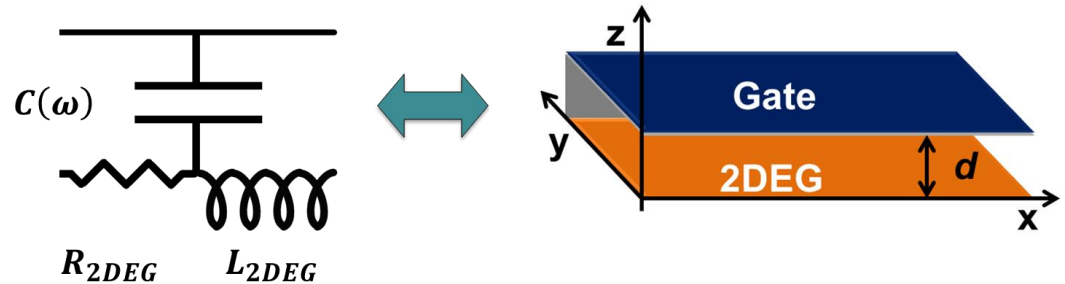


Outline

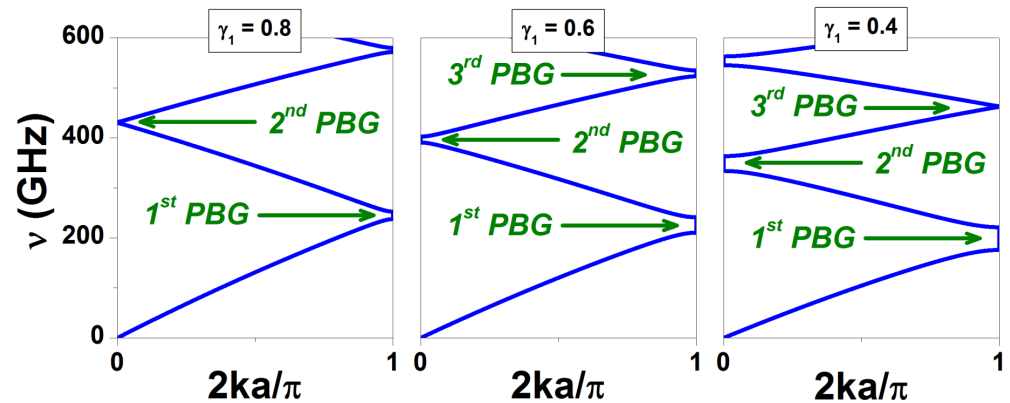
- Introduction
- Low dimensional far infrared plasmonic structures
 - Integrated 2D plasmonic devices
 - Strong coupling of tunable 2D plasmonic resonators
 - Generalized 2D plasmonic transmission line model
 - Active control of plasmonic crystal band structure
 - Electromagnetically induced transparency
- Conclusions & future directions

Conclusions

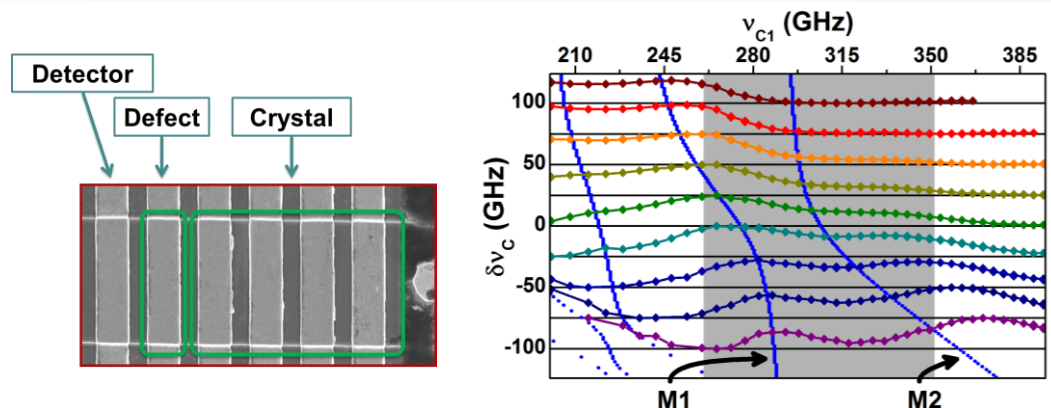
Plasmonic waveguide



Active control of plasmonic dispersion



Monolithically integrated & reconfigurable



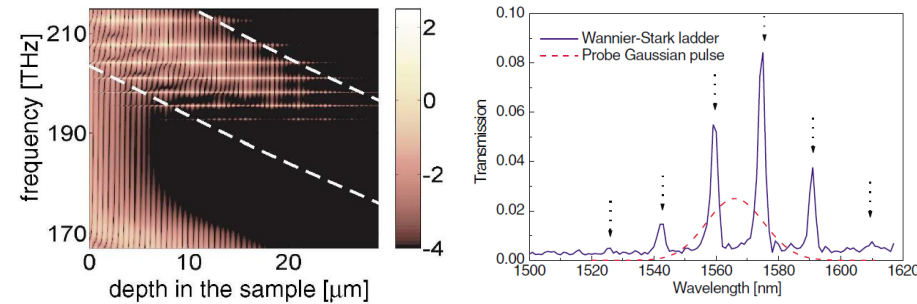
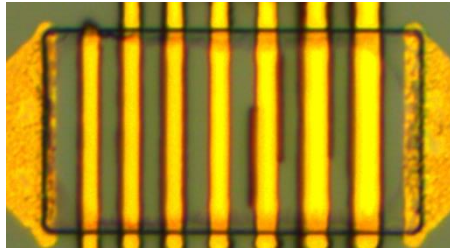
Credits

- Low dimensional plasmonics
 - **Grower:** John Reno
 - **Device Design/Mask Layout:** Greg Dyer, Eric Shaner
 - **Processing:** Eric Shaner, Don Bethke
 - **Characterization:** Greg Dyer, Albert Grine
 - **Data Analysis:** Greg Dyer
 - **Numerical Modeling:** Greg Dyer, Greg Aizin, John Mikalopas
 - **Theory:** Greg Aizin, Greg Dyer
 - **Funding:** Eric Shaner (STTR, BES), Joel Hensley (STTR)
- THz transceivers
 - **Grower:** John Reno
 - **Device Design/Mask Layout:** Mike Wanke, Chris Nordquist, Mike Cich, Mark Lee, Eric Young
 - **Processing:** Mike Cich, Melissa Cavaliere
 - **Characterization:** Greg Dyer, Mike Wanke, Albert Grine, Chuck Fuller, Mark Lee, Eric Young
 - **Data Analysis:** Greg Dyer, Mike Wanke
 - **Funding:** Mike Wanke (LDRD)

ADDITIONAL MATERIAL

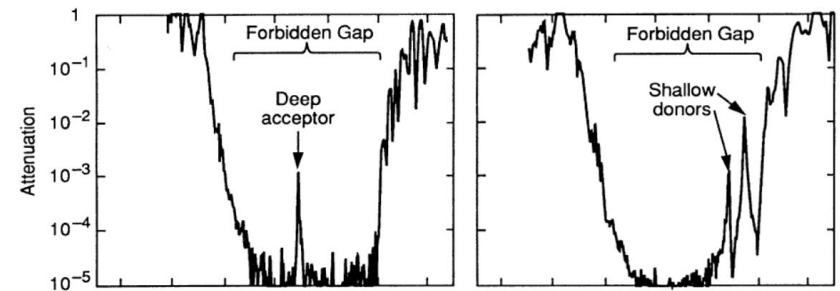
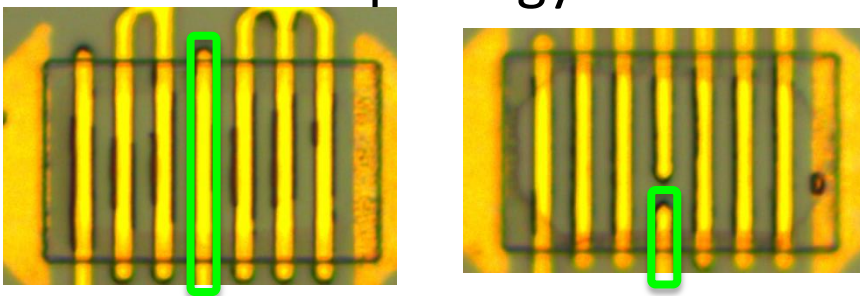
Quantum & photonic analogues

■ Wannier-Stark ladder



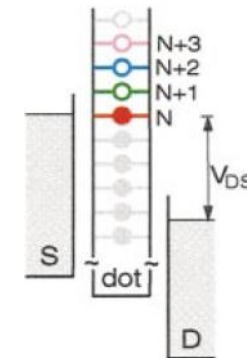
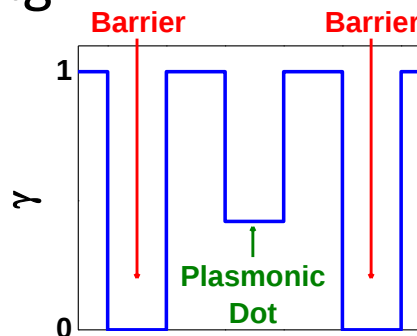
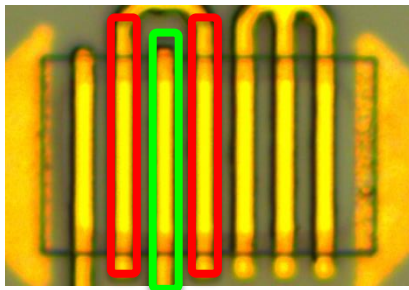
R. Sapienza *et al.* *PRL* **91**, 263902 (2003)

■ Defect morphology & location



E. Yablonovitch *et al.* *PRL* **67**, 3380-3383 (1991)

■ Plasmonic 'tunnel' coupling

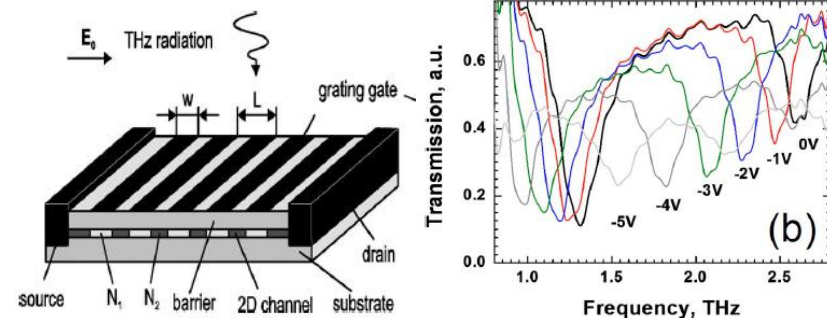


D.R. Stewart *et al.*
Science **278**, 1784-1788 (1997)

Potential material systems

Nitride 2DEGs

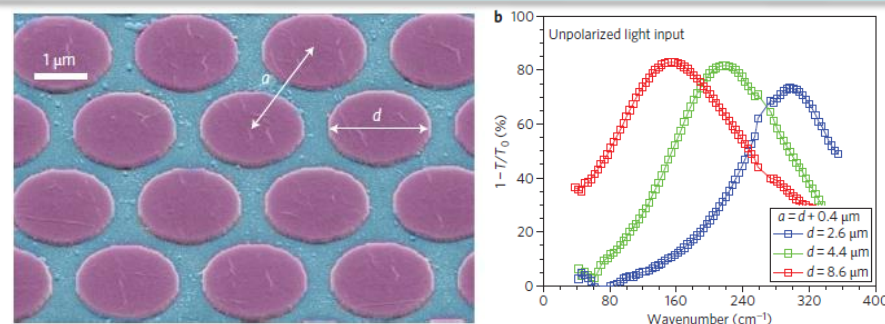
- >1 THz, > 77 K
- Damping: phonon energies, coupling



A. V. Muravjov, *et al.*, *APL* **96**, 042105 (2010)

Graphene

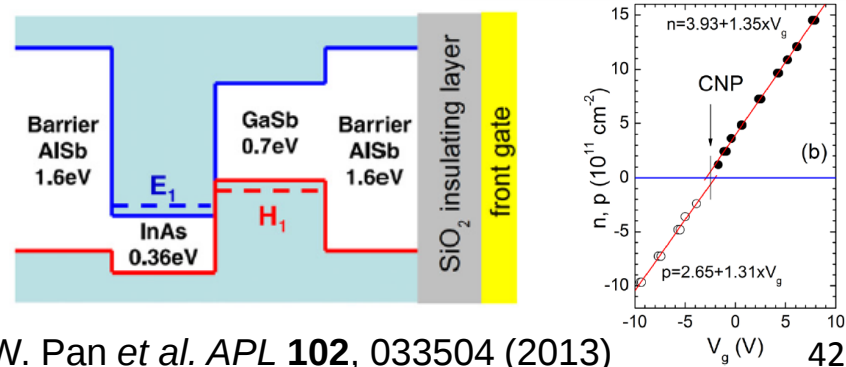
- FIR to MIR, room temperature
- Dirac plasmons



H. Yan *et al.*, *Nat. Nano.* **7**, 330 (2012)

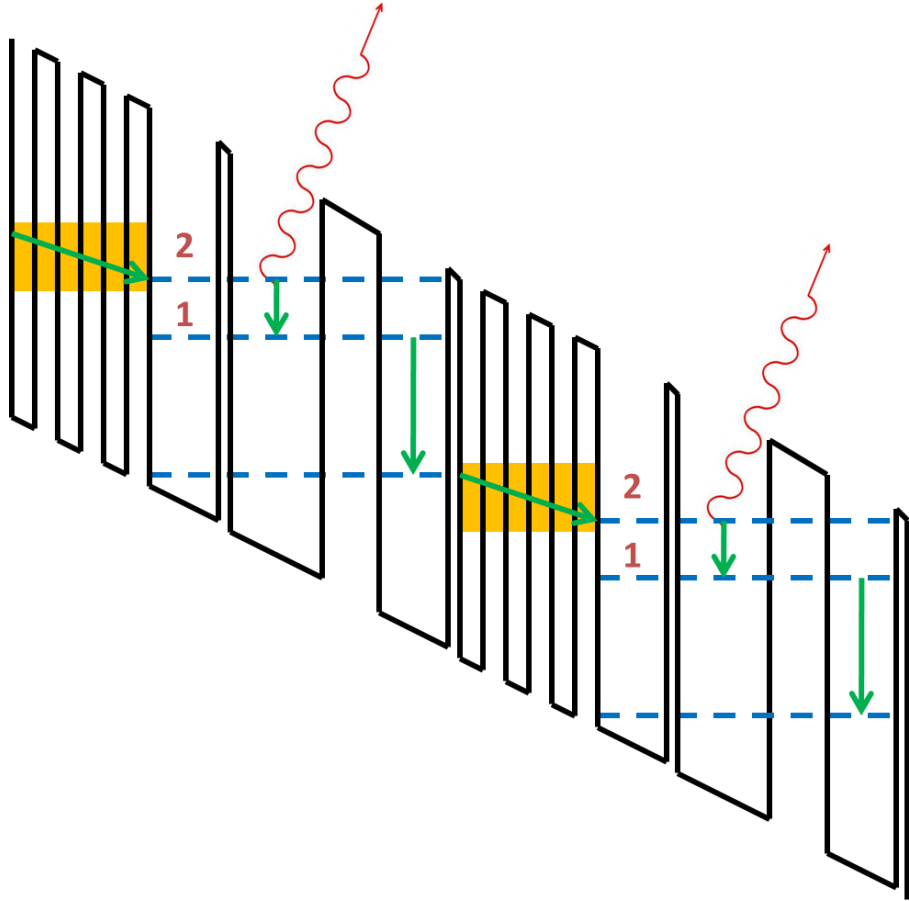
Type-II InAs/GaSb

- 2DEG, 2DHG & CNP
- p-n: $m_e = 0.03 m_0$, $m_h = 0.37 m_0$

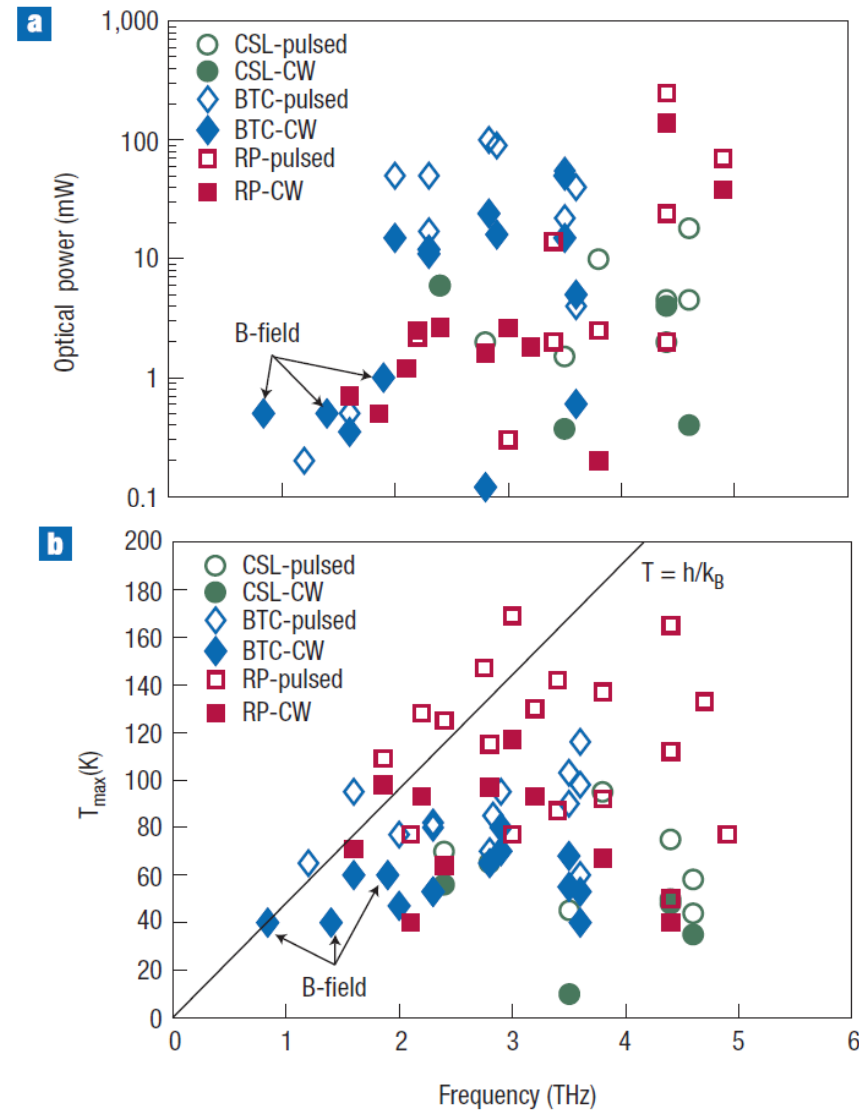


W. Pan *et al.* *APL* **102**, 033504 (2013)

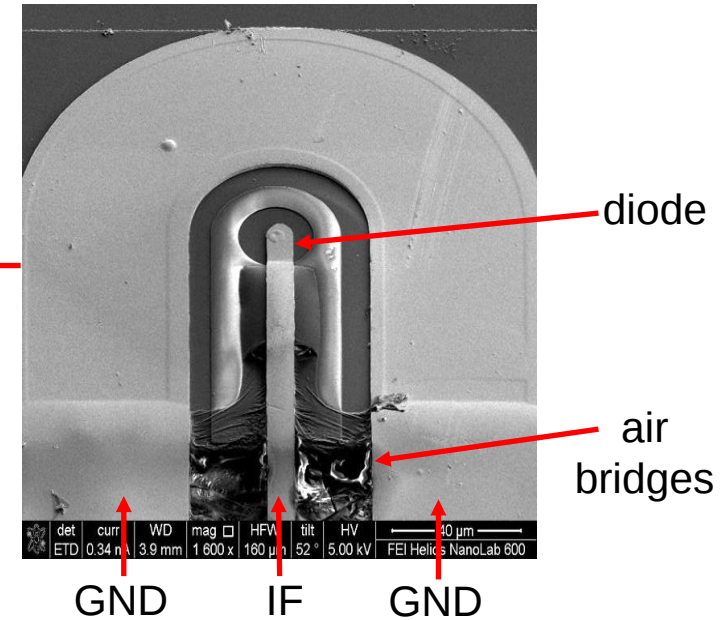
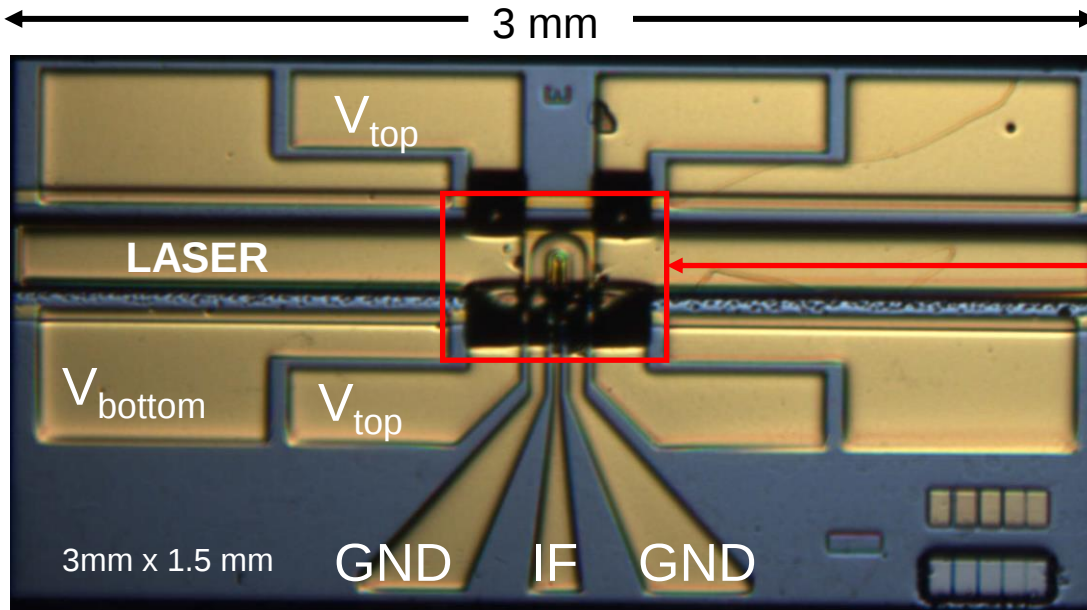
THz Quantum Cascade Lasers



Unipolar semiconductor laser
based on engineering of electronic
states in conduction band

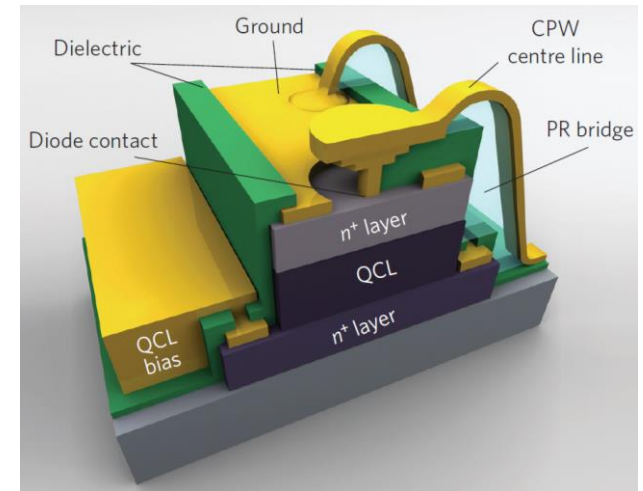


THz QCL Transceivers

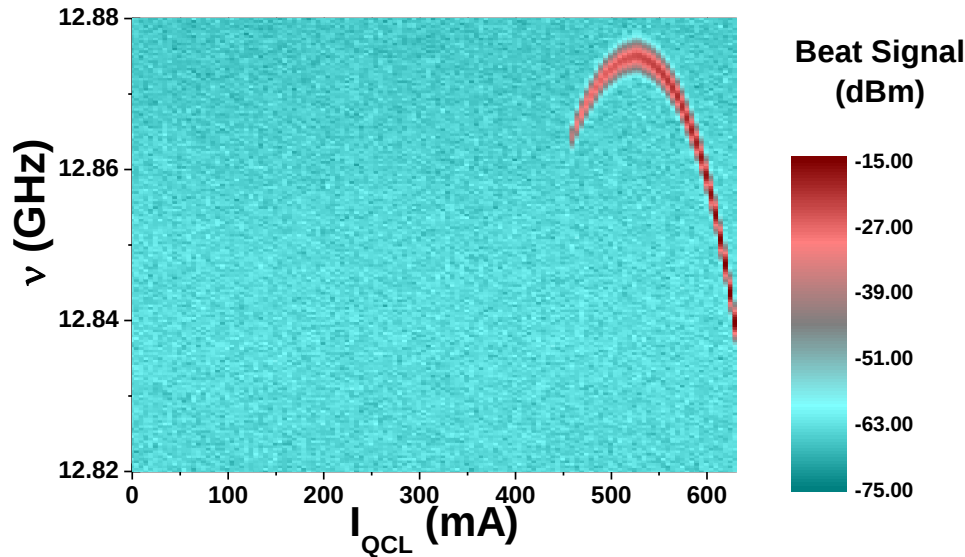
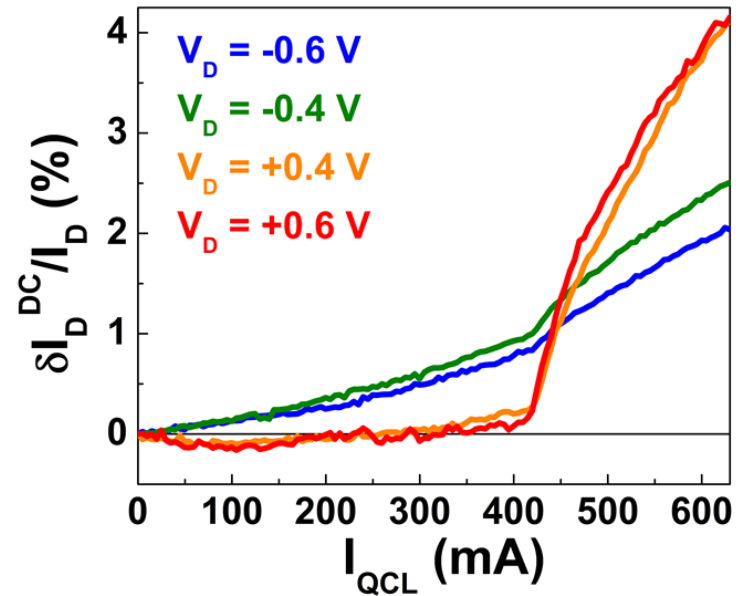
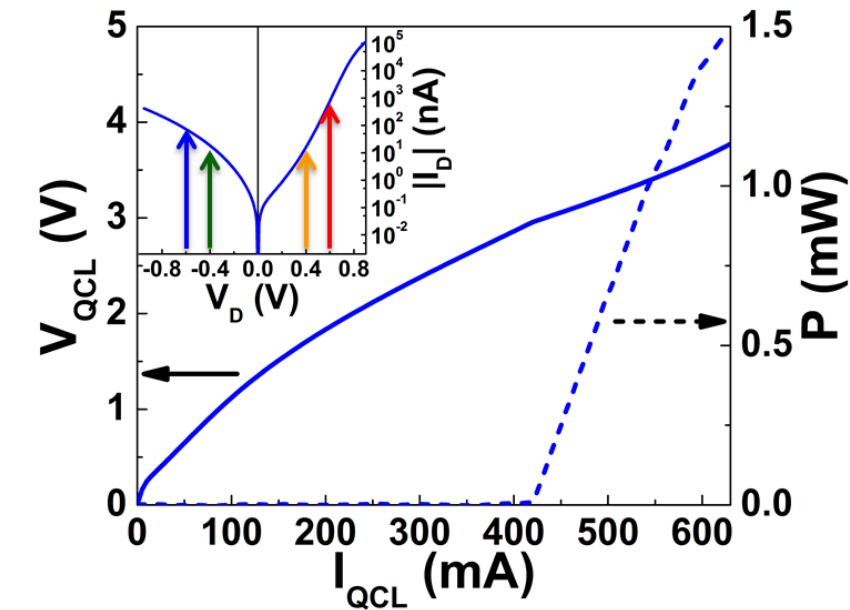


Schottky diode directly embedded into
2.8 THz QCL waveguide at cavity center

QCL functions as local oscillator (LO)



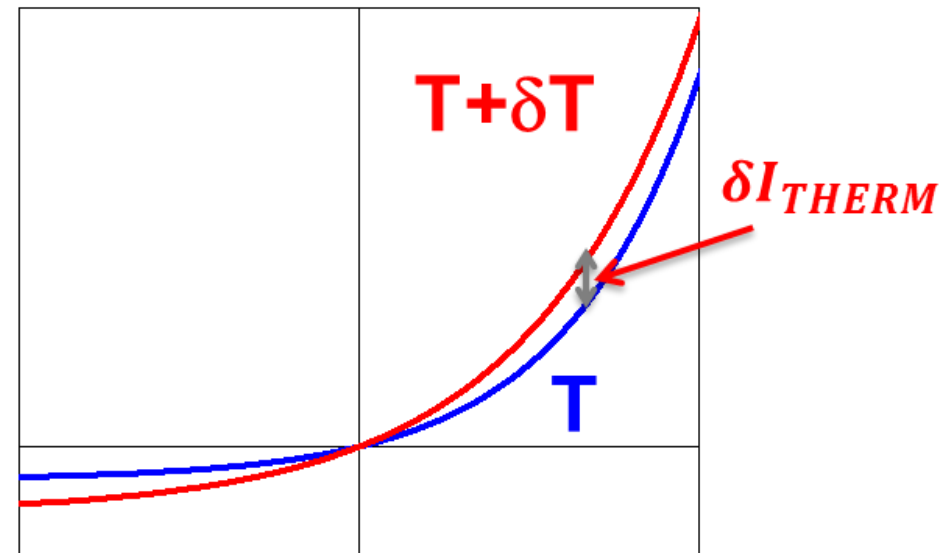
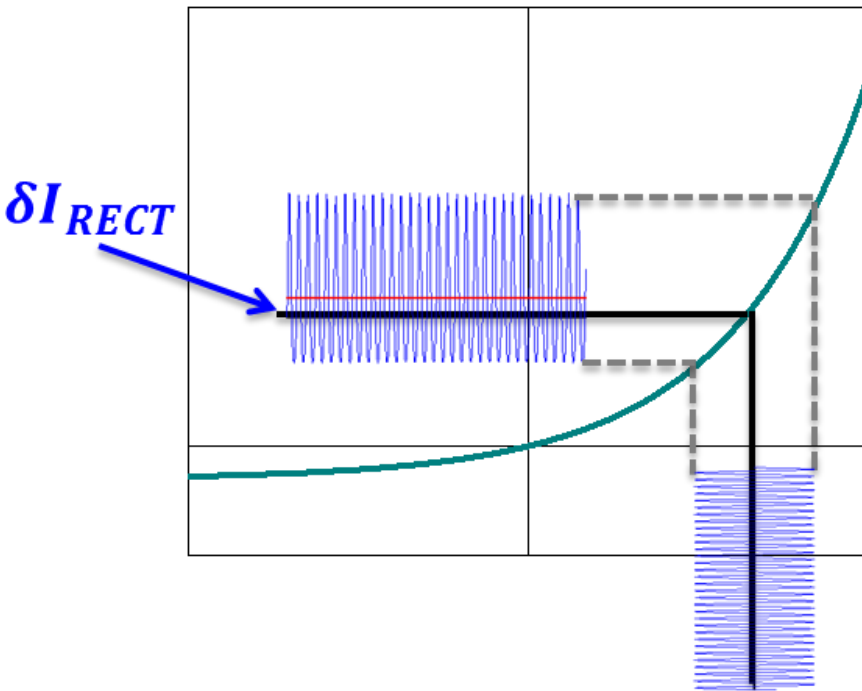
Rectified Diode Response



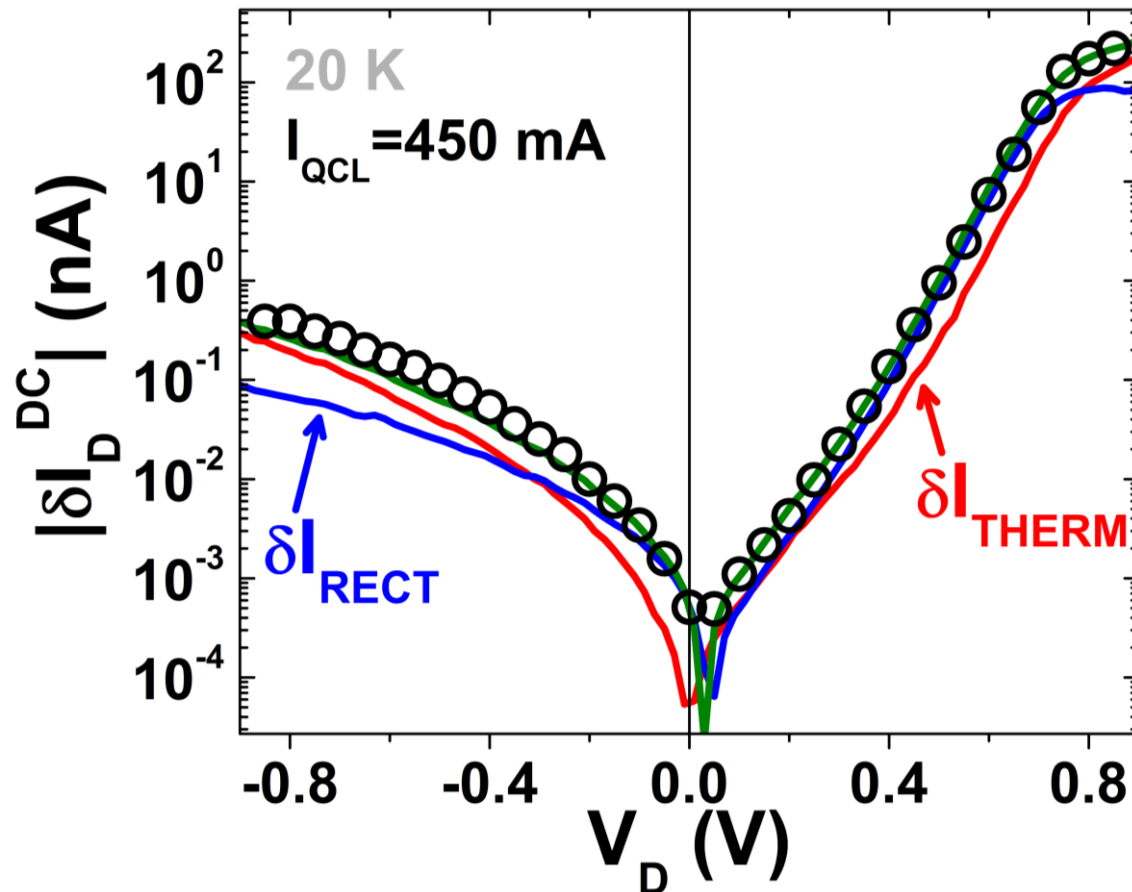
Rectified (DC) and beat (microwave) signals from QCL coupling to diode

Ideal Diode Picture

$$I_D = I_0 \left[\exp\left(\frac{V_D}{k_B T}\right) - 1 \right]$$



Model of Rectified Response



Circles: measured change in diode current with 'QCL on' vs. 'QCL off'

Derivatives: calculated from static diode I-V

δP_{QCL} calculated from static QCL I-V

Fit: $\Theta_T = 12.28 \text{ W/K}$ & $\delta V_{THz}^2 = 6.92 \times 10^{-5} \text{ V}^2$

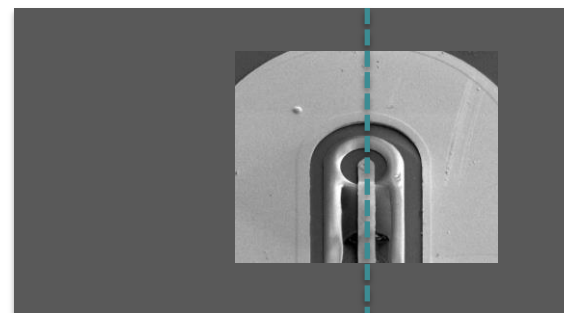
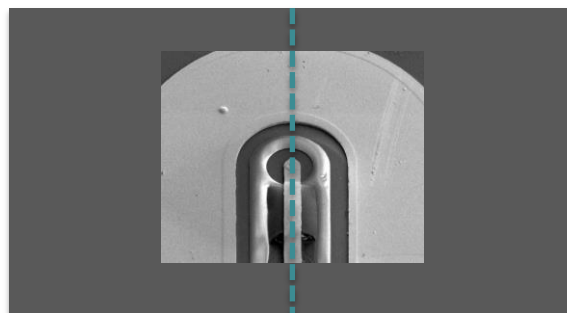
$$\delta I_D^{DC} = \frac{1}{4} \frac{\partial^2 I_D}{\partial V_D^2} \delta V_{THz}^2 + \frac{\partial I_D}{\partial T} \delta P_{QCL} \Theta_T$$

Relative Diode Positioning

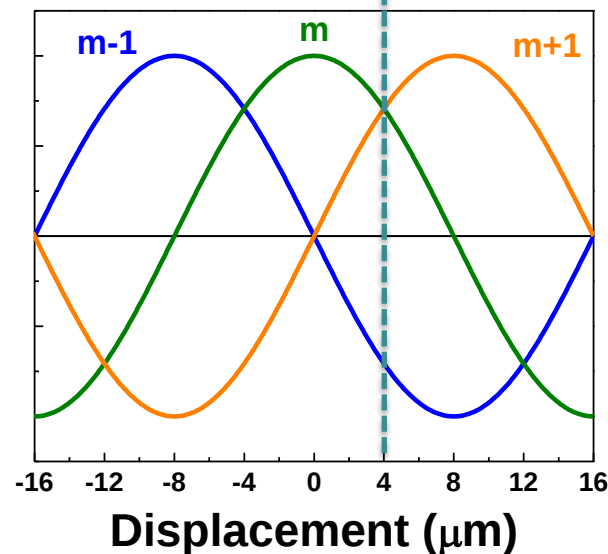
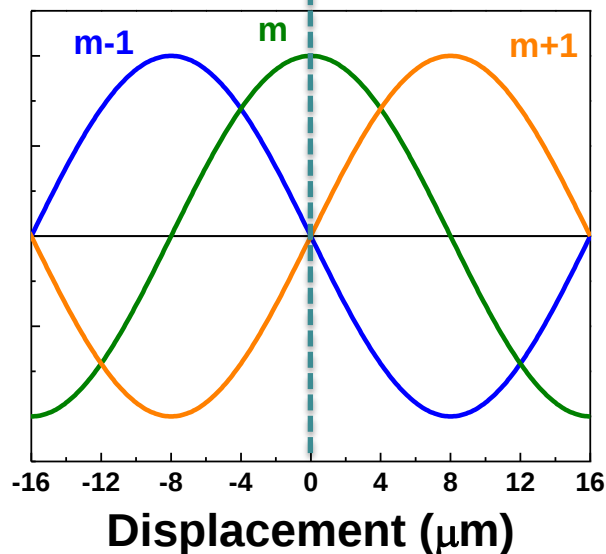
*diode at $x = 0 \mu\text{m}$
(mid-cavity)*

*diode at $x = +4 \mu\text{m}$
($\lambda/8$ offset)*

*Diode location at
center of QCL
cavity*

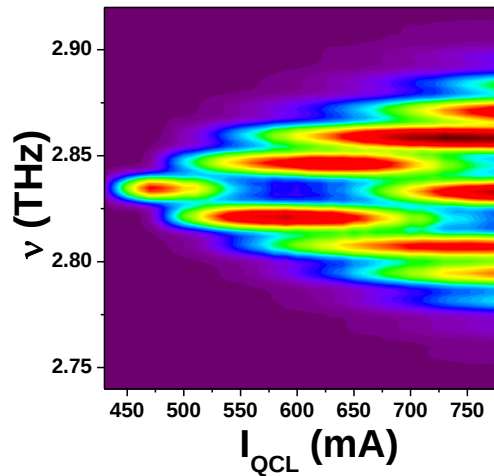


*Fabry-Perot mode
field distributions*

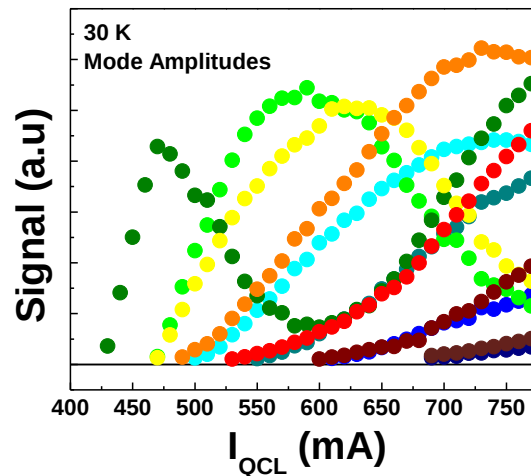


Lithographic Tuning

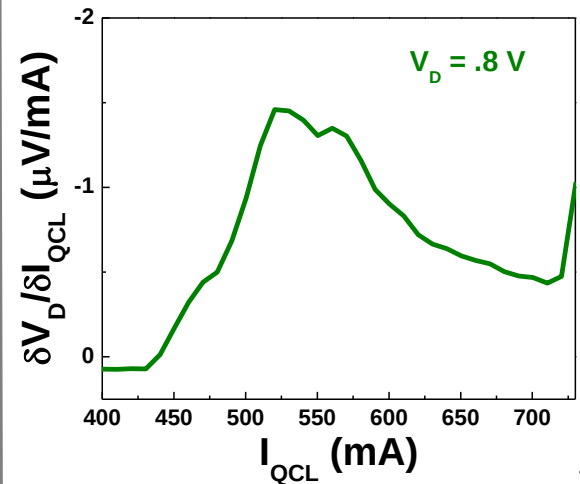
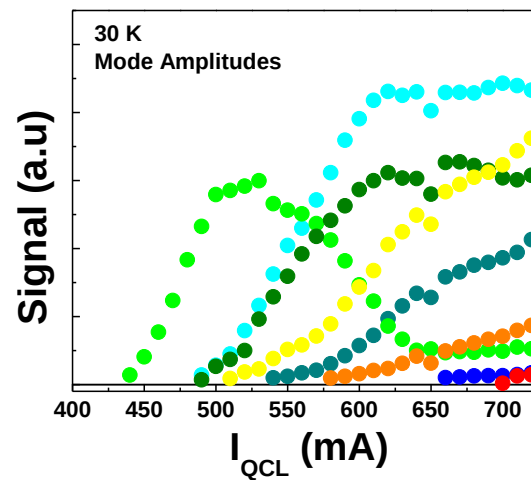
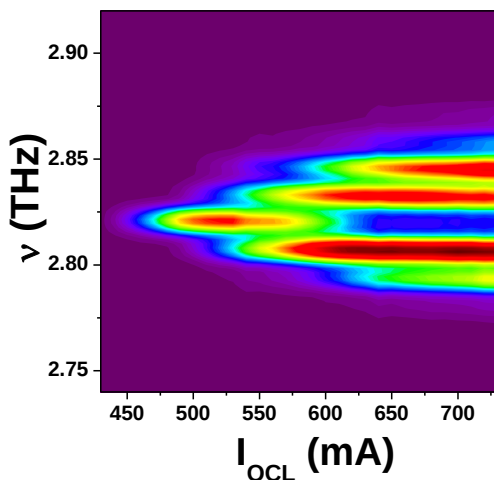
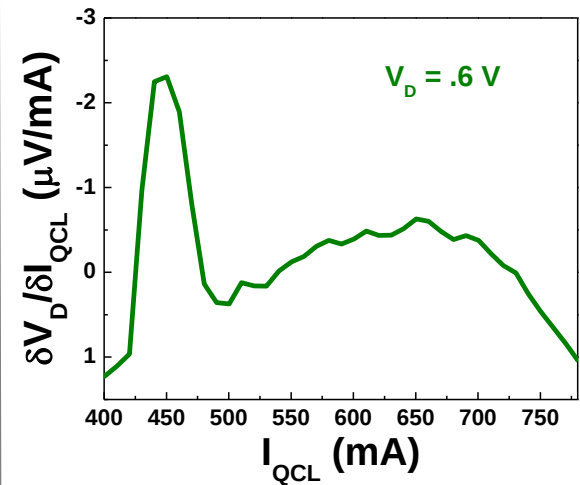
QCL Emission



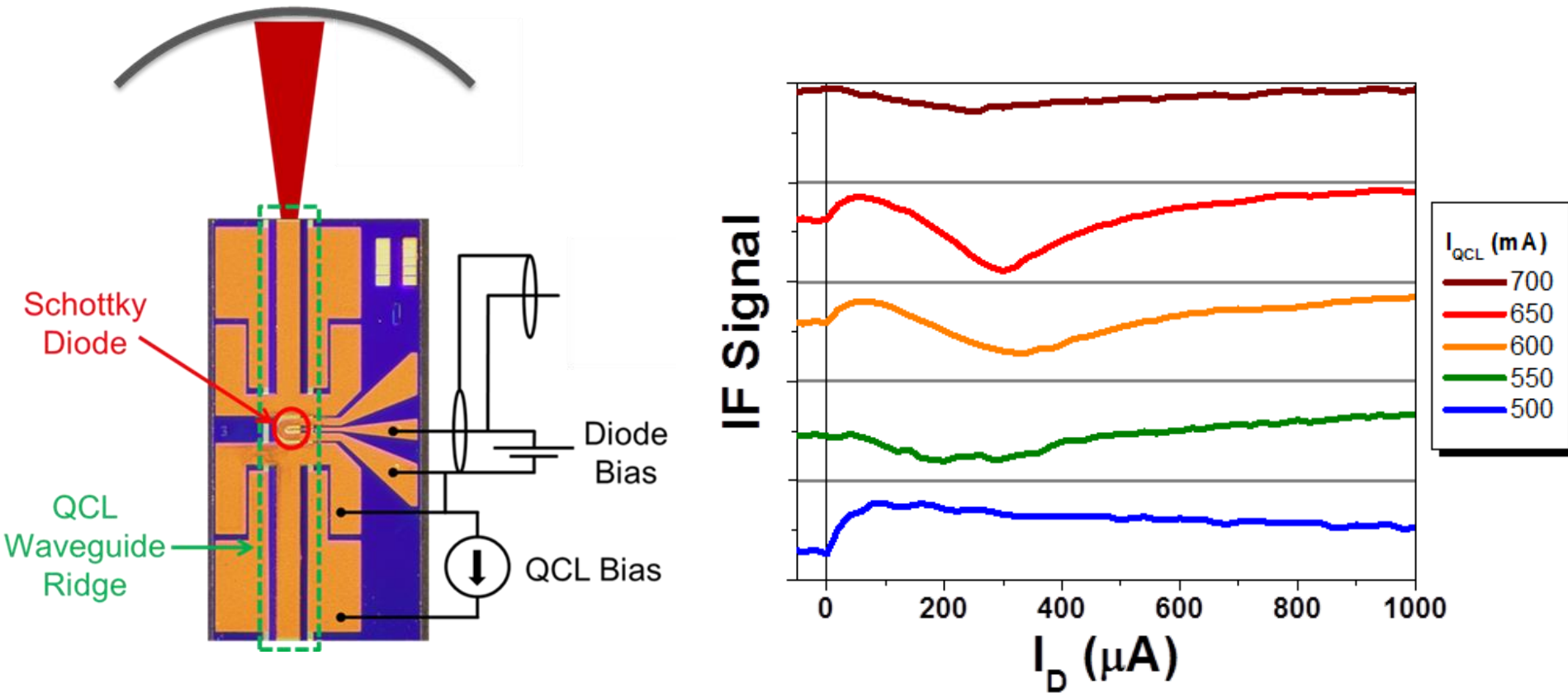
Mode Amplitudes



Rectified Signal

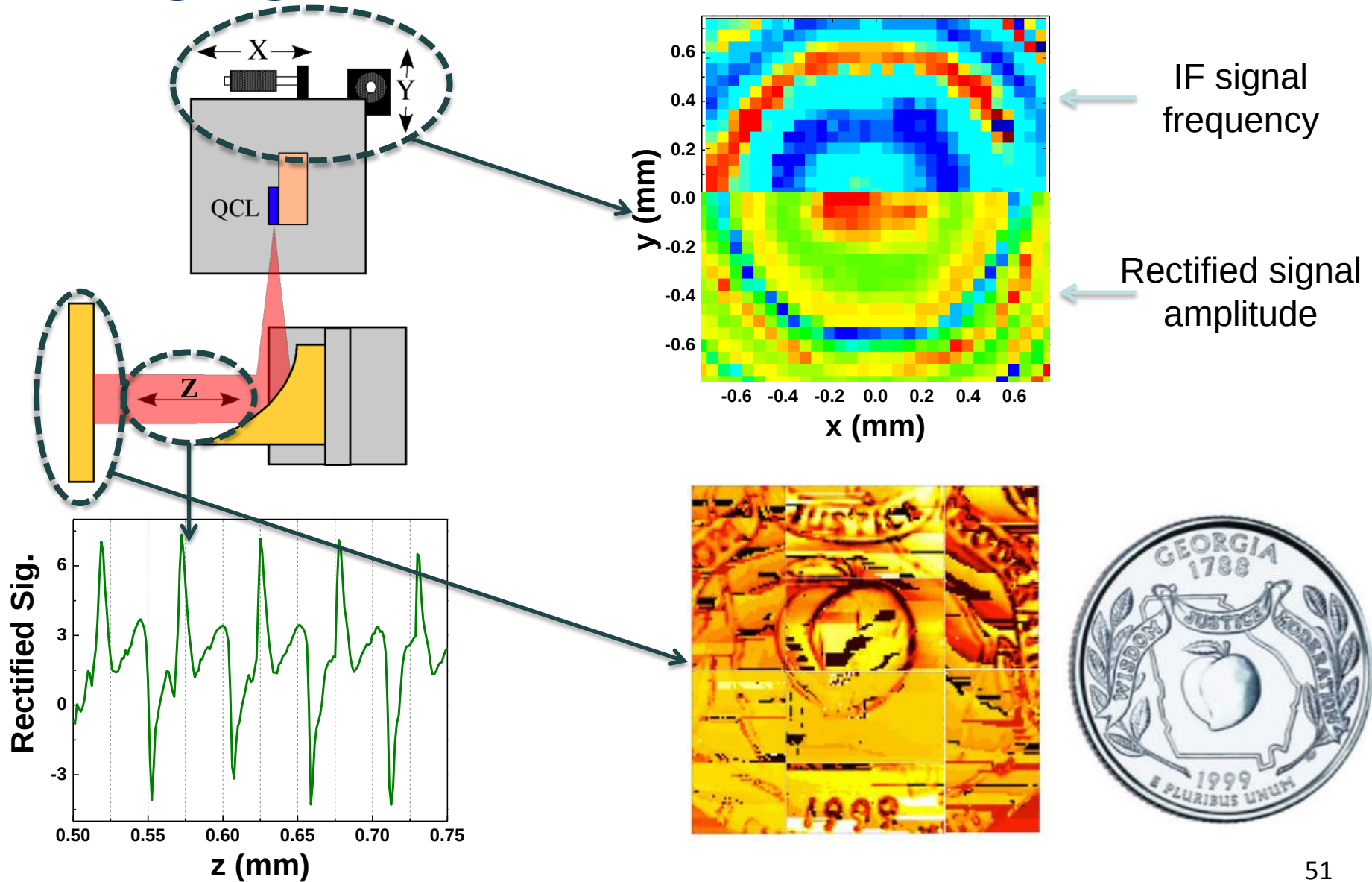


Diode-Laser Perturbation

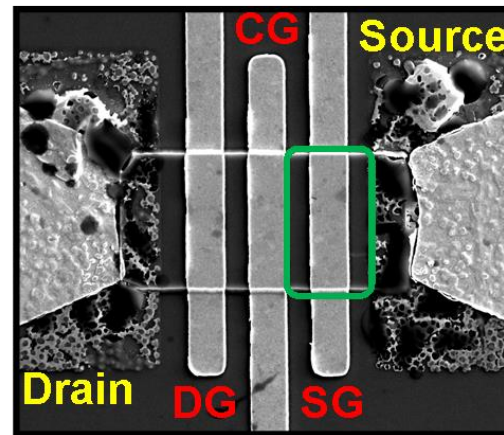
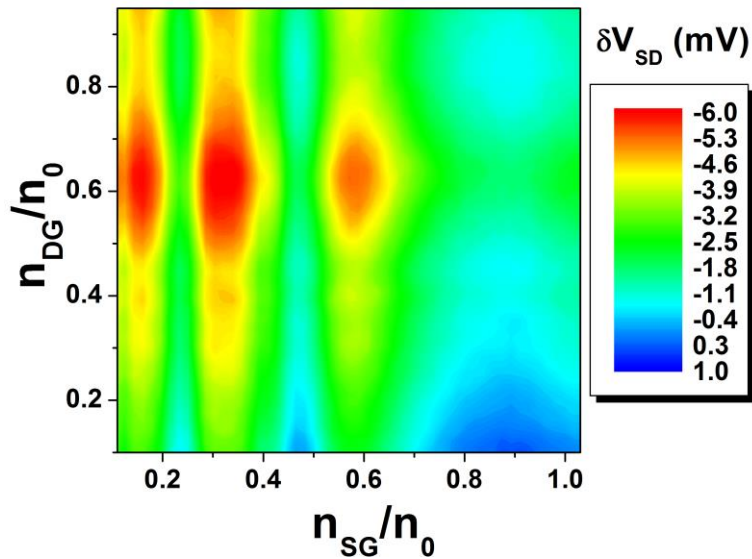


**Non-monotonic change in beat signal tuning diode:
perturbation of QCL modes**

Imaging via Feedback

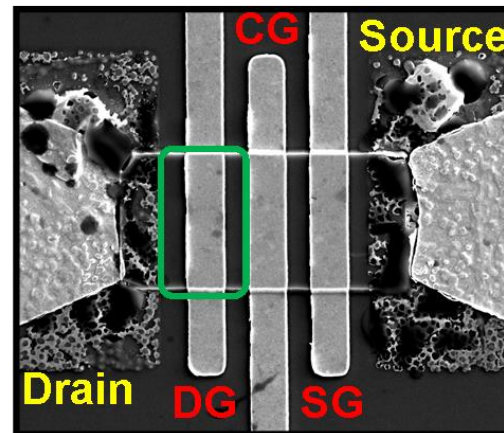
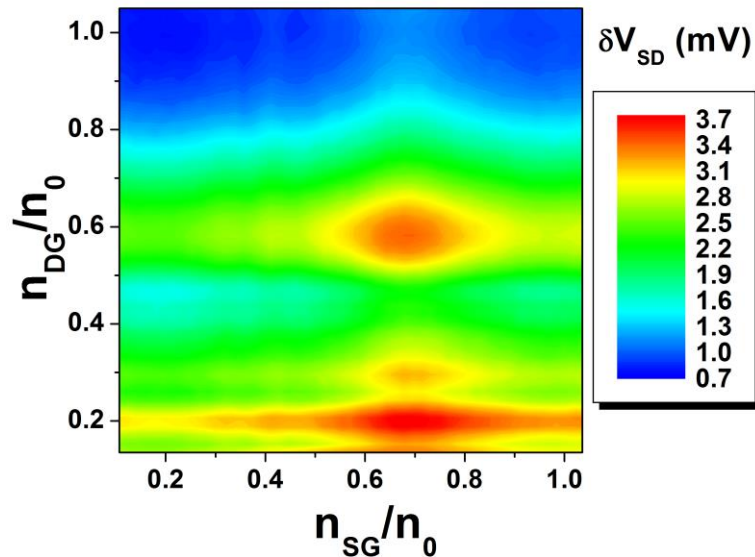


Source-Drain Bias Cavity Selection



370 GHz

$I = +500$ nA



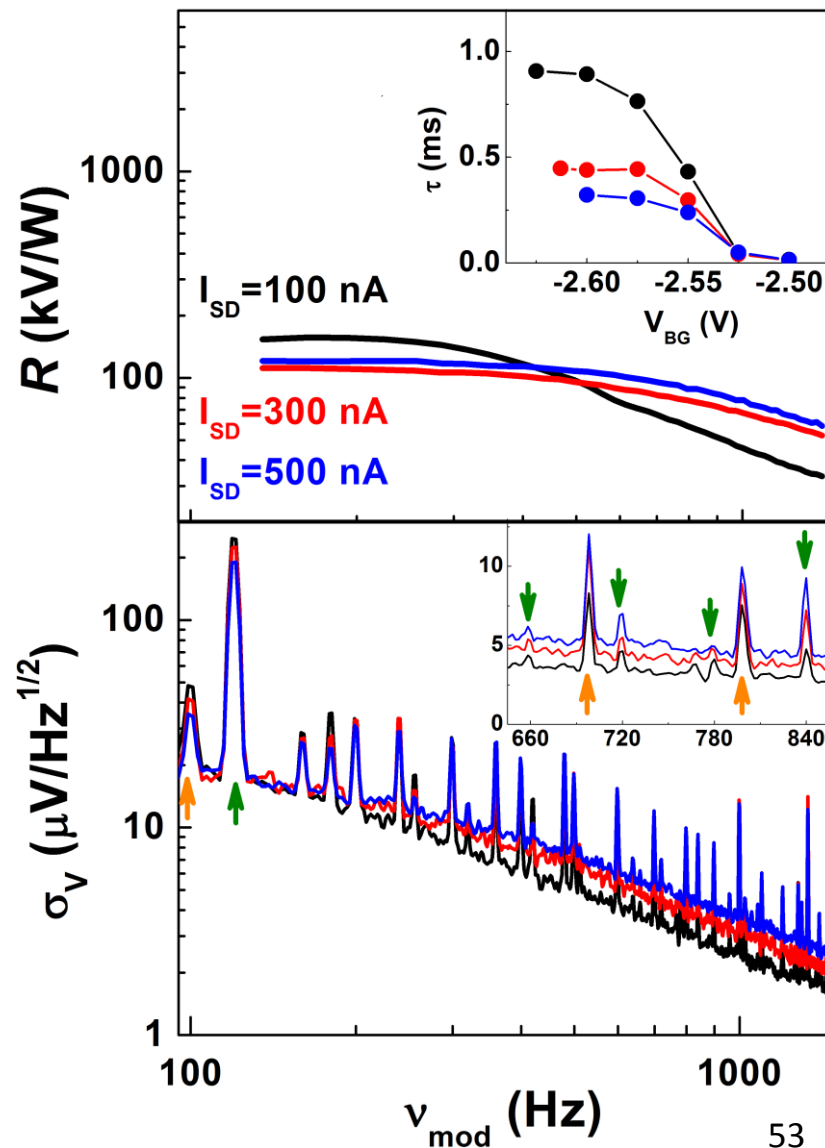
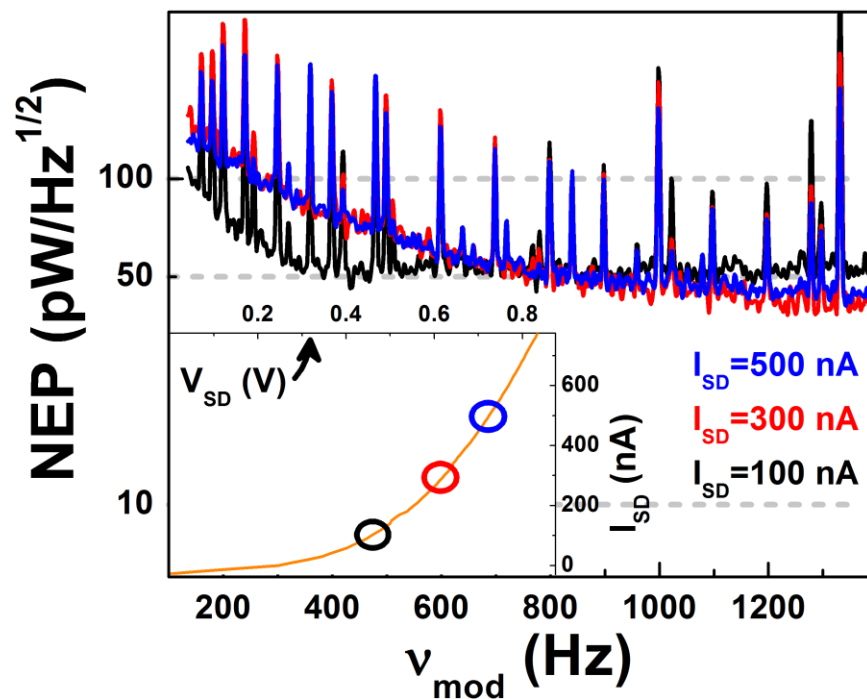
$I = -500$ nA

Bolometric Detection FOMs

11K Operation, 243 GHz

$$R > 100 \text{ kV/W}$$

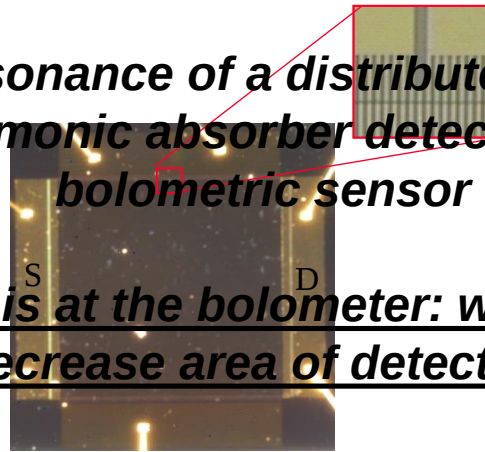
$$NEP < 50 \text{ pW/Hz}^{1/2}$$



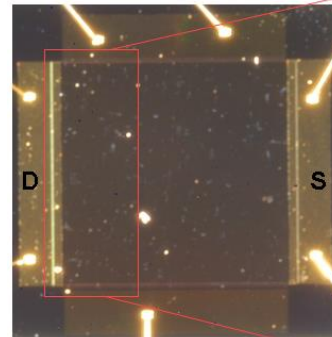
Resonant Terahertz Response

Resonance of a distributed 2D plasmonic absorber detected by bolometric sensor

Action is at the bolometer: what if we decrease area of detector?

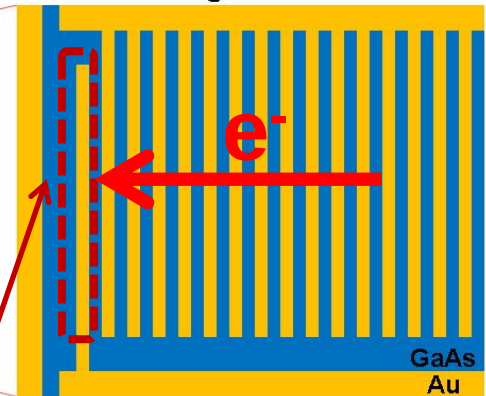


1 mm by 1 mm channel



Grating Period: 4 μm

Tuning Gate Contact

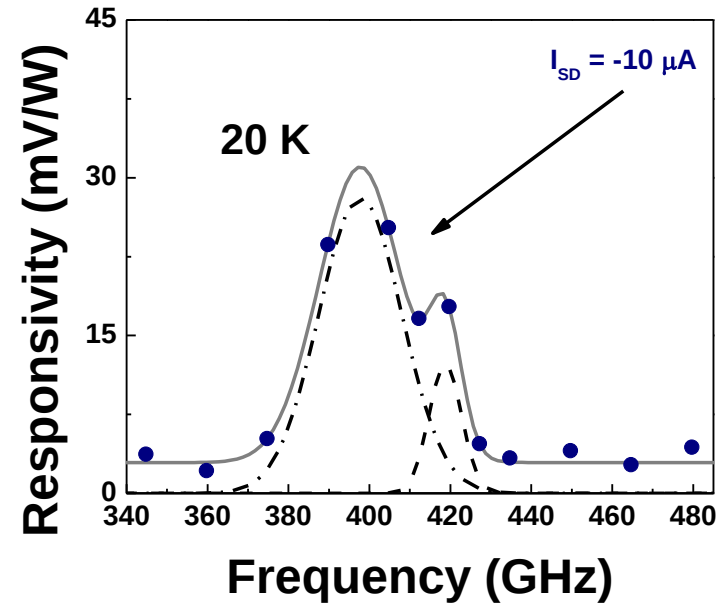
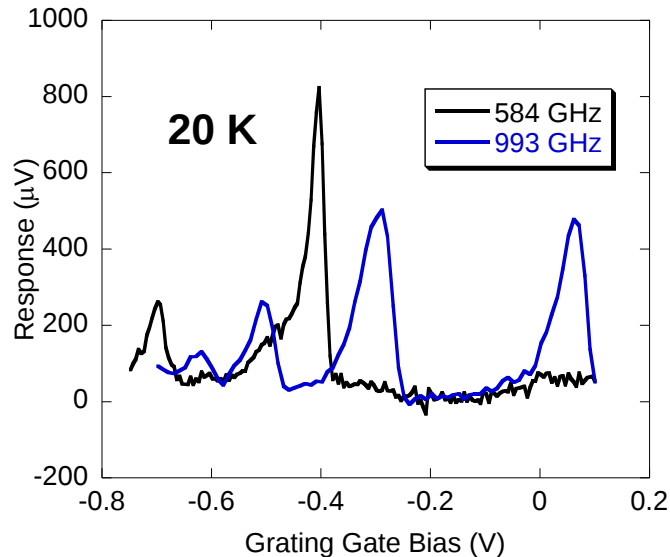


Barrier Gate Contact

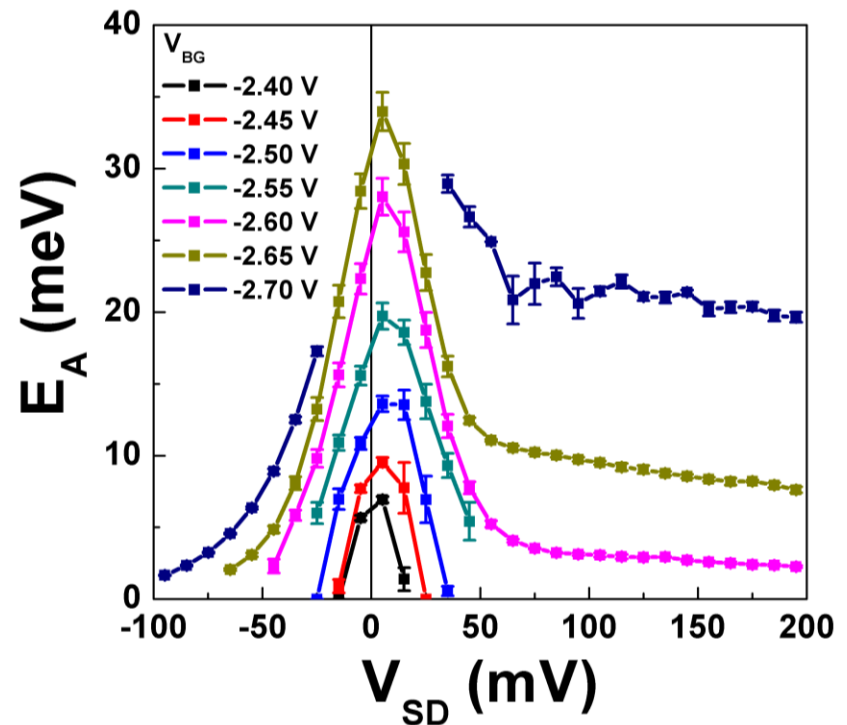
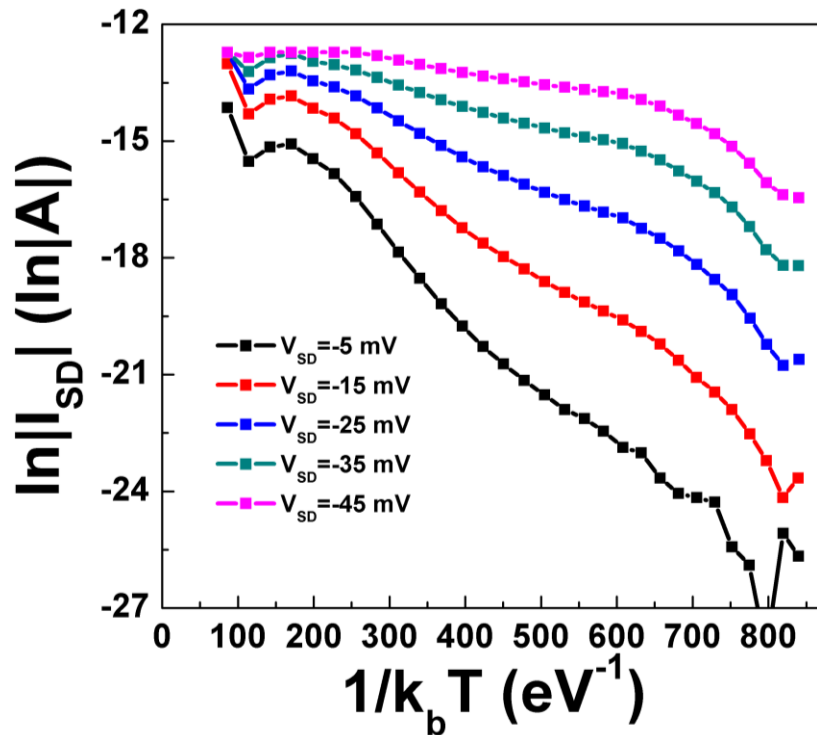
Bolometric Element

$$V_{BG} < V_{TH}$$

Split-Gate Detector Photoresponse

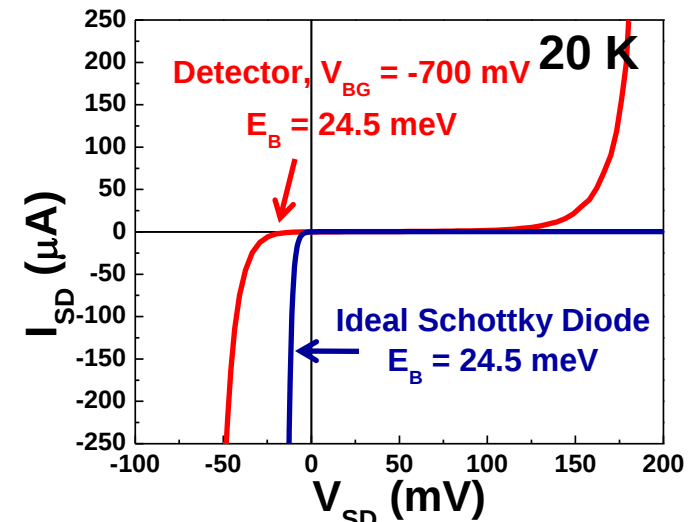
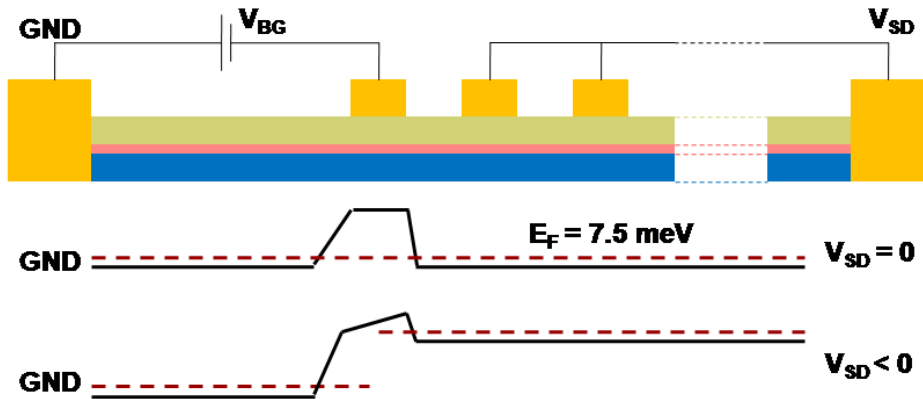


Activation Energy



- Double quantum well material, $V_{TH} \sim -2.40$ V
- Left: Arrhenius plots $V_{BG} = -2.6$ V, fits from 200-600 eV^{-1}
- Right: activation energy as a function of bias

Barrier Lowering ($V_{SD} < 0$)



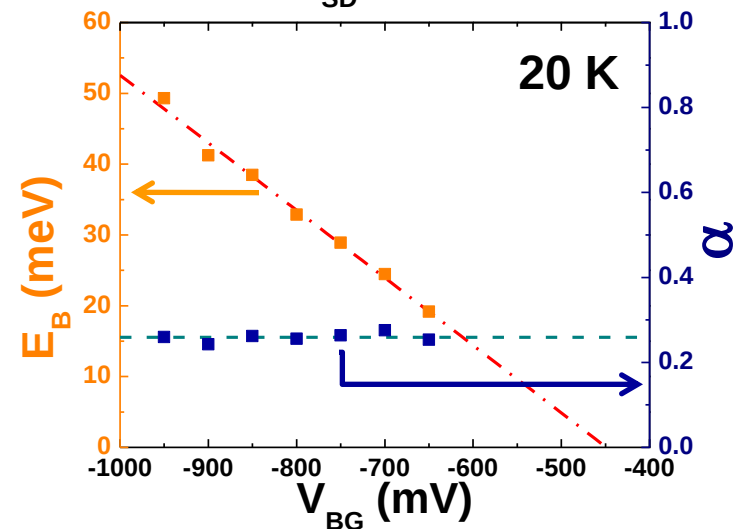
Non-Degenerate Limit, Negative Bias

IF $E_B - \mu + e\alpha V_{SD} \gg k_B T$:

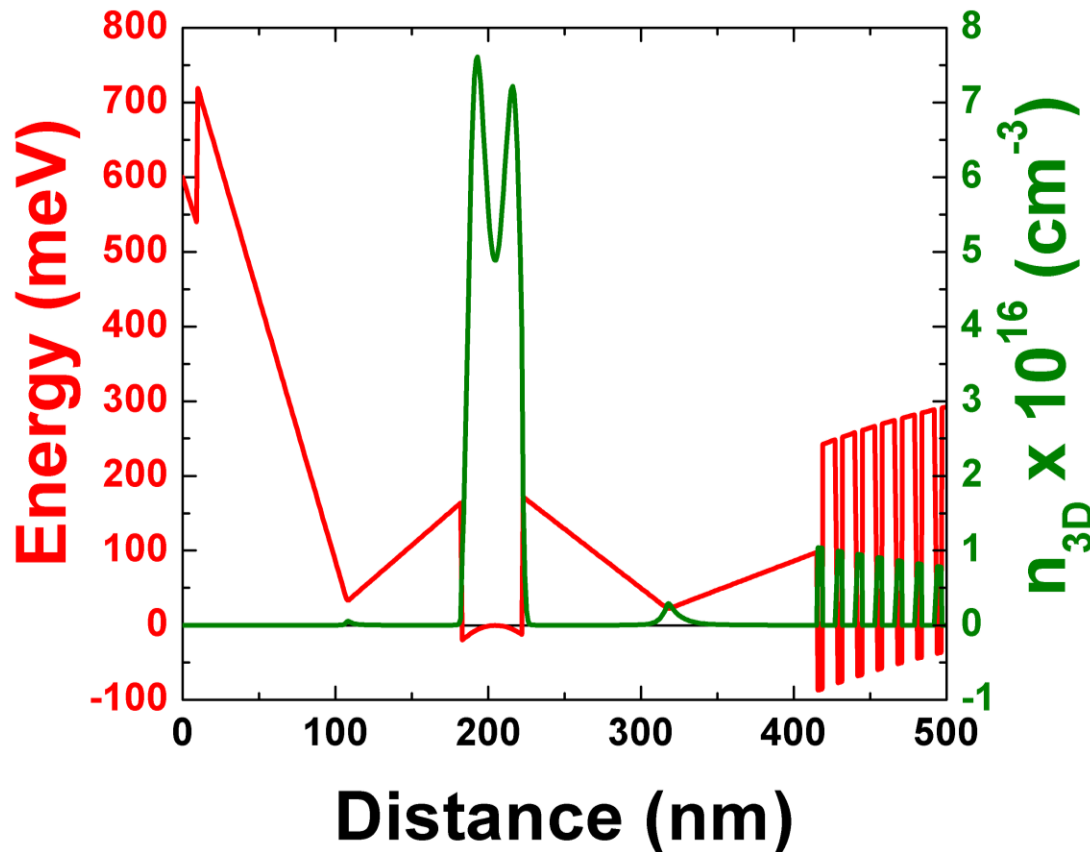
$$J_{2D} \propto -e^{\frac{-E_B}{k_B T}} e^{\frac{-e\alpha V_{SD}}{k_B T}}$$

Barrier Height E_B

Barrier Lowering Parameter α



Schrodinger Poisson Model

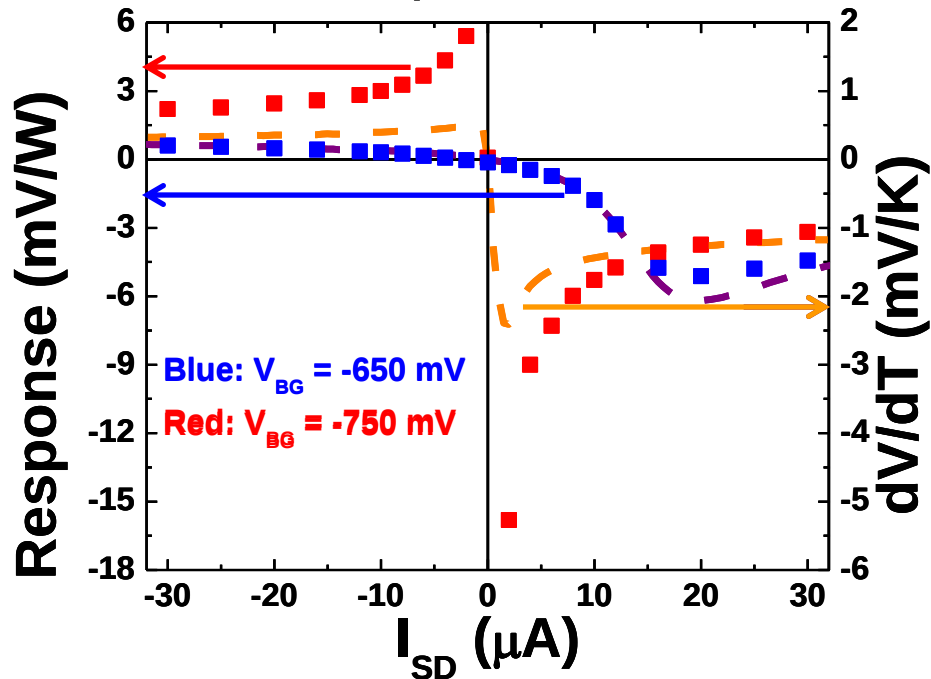


- Via Greg Snider's 1D Schrodinger-Poisson solver
- Carrier density in backside delta doping and superlattice
- Potential for parallel channel, 'lever arm'
- From lateral barrier activation model:

$$\left[\frac{\partial V_{SD}}{\partial T} \right]_{\max} \approx \mp 0.08625 \frac{E_A}{\alpha_{\pm} k_B T}$$

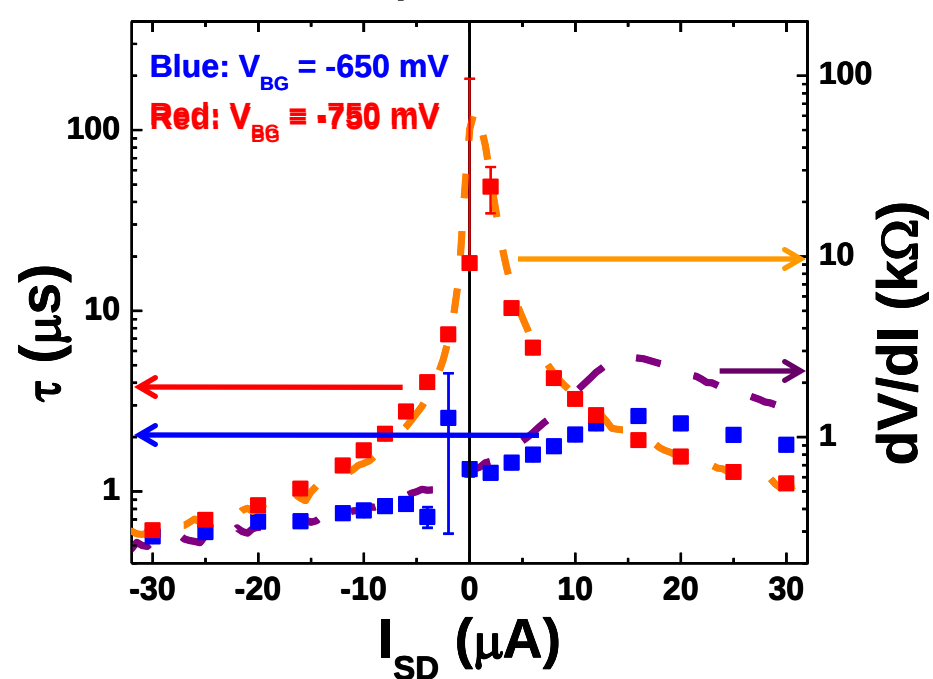
Hot Electron Bolometric Response

40 K, 420 GHz



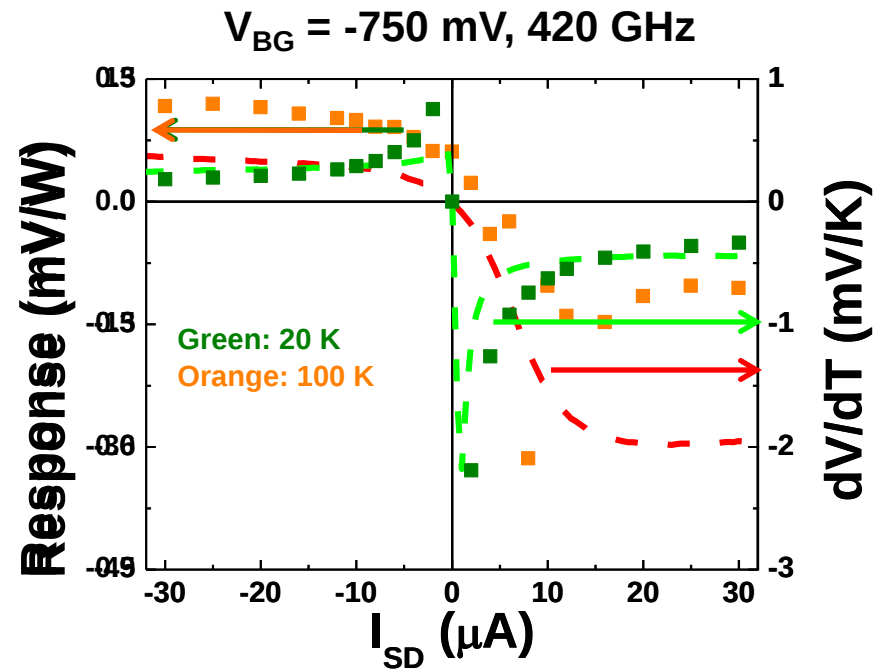
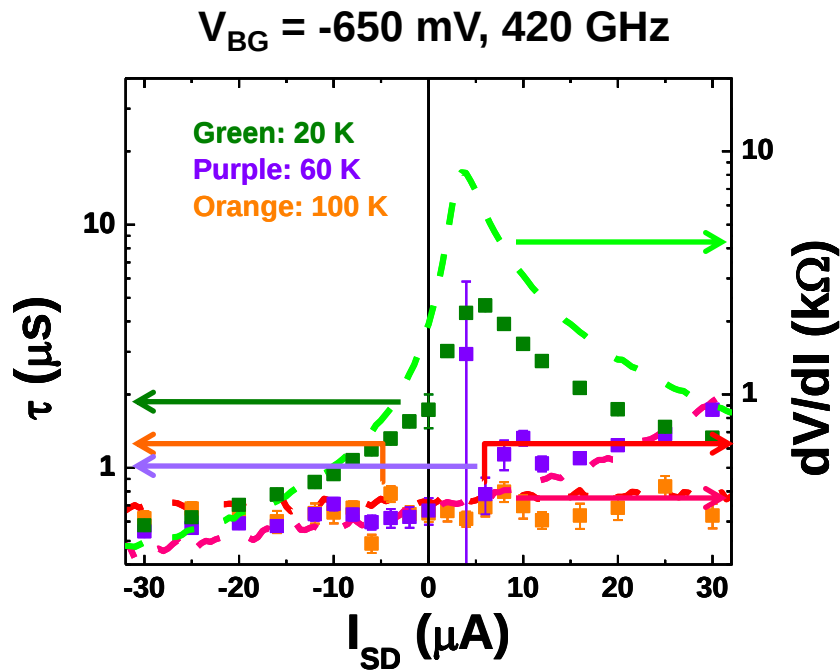
$$\text{Response} \propto \frac{\partial V}{\partial T}$$

40 K, 420 GHz



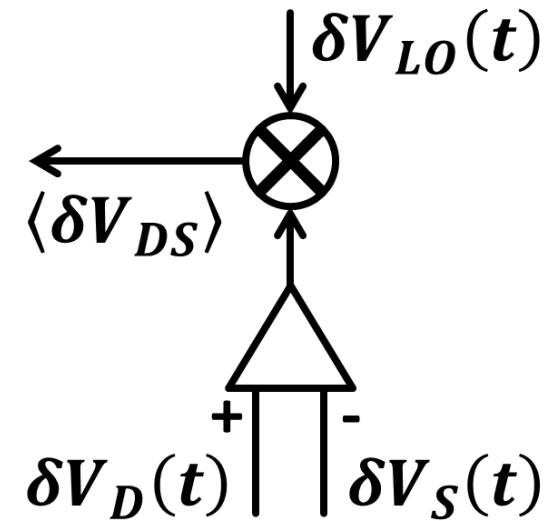
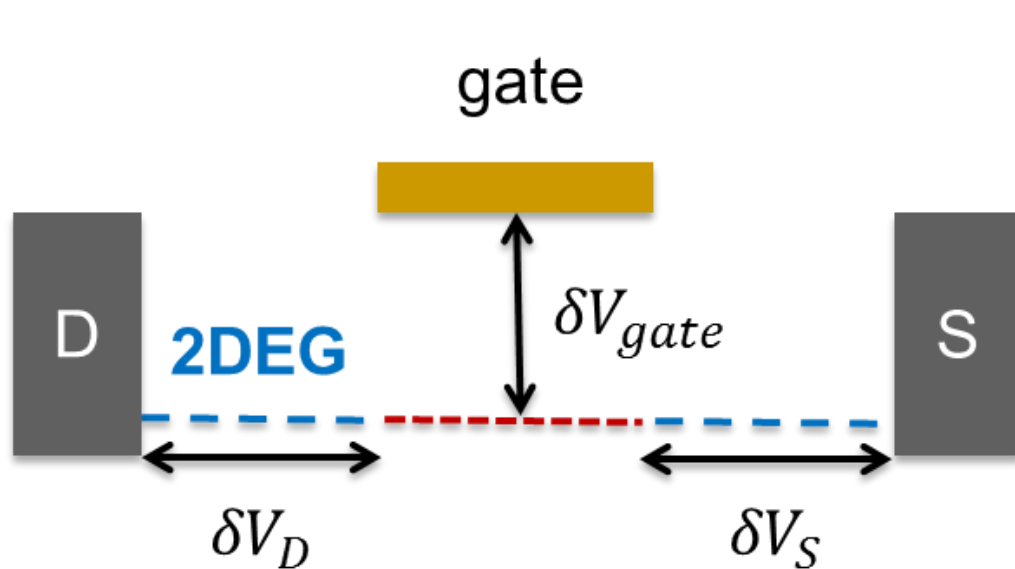
$$\tau = R_d C = \frac{\partial V}{\partial I} C$$

Temperature Dependence



- Response described by “RC” over broad range of T
- Response must be scaled to agree with dV/dT
- Scale factor: approx. .07 W/K at 20 K, 3 W/K at 100 K

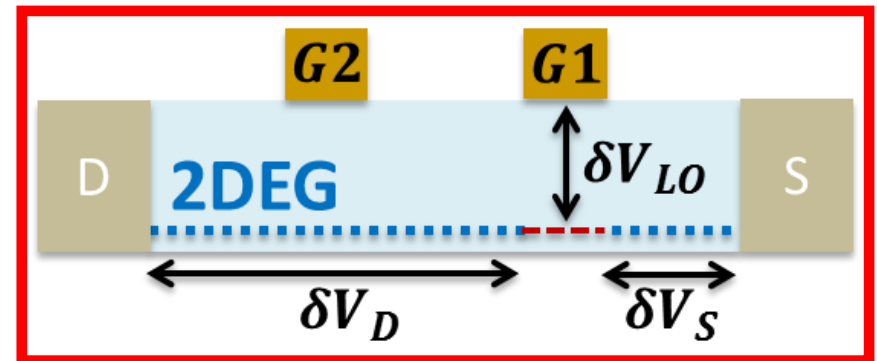
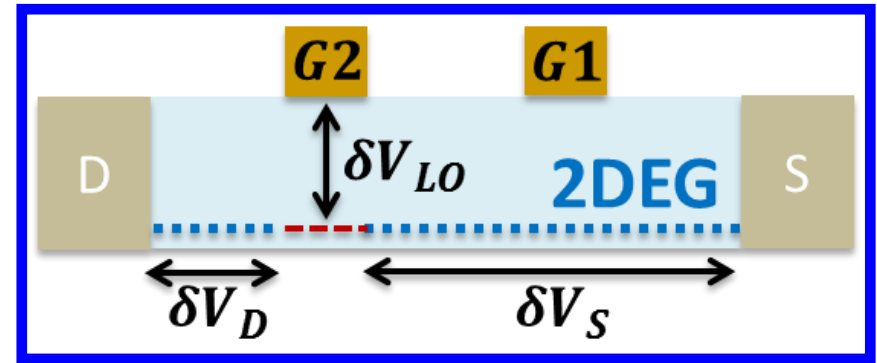
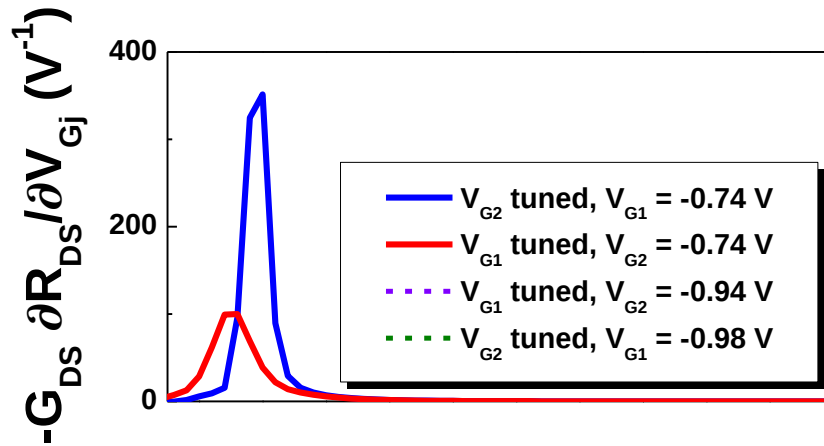
Plasmonic Homodyne Mixing



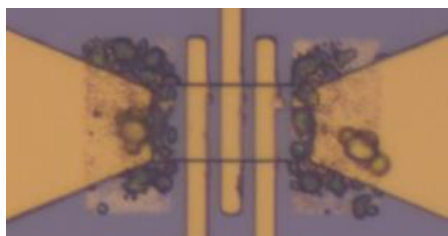
$$\begin{aligned} \langle \delta V_{DS} \rangle &= \left\langle \left[\frac{\partial R_{chan}}{\partial V_{gate}} \delta V_{gate} e^{-i(\omega t + \phi)} \right] [\delta V_D - \delta V_S] e^{-i\omega t} G_{chan} \right\rangle \\ &= \frac{1}{2} \left[\boxed{G_{chan} \frac{\partial R_{chan}}{\partial V_{gate}}} \right] \delta V_{gate} [\delta V_D - \delta V_S] \cos[\phi] \end{aligned}$$

ϕ is relative phase between gate & channel THz potentials

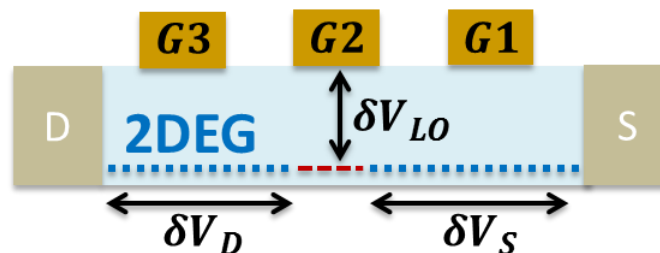
Homodyne Mixing Photovoltage



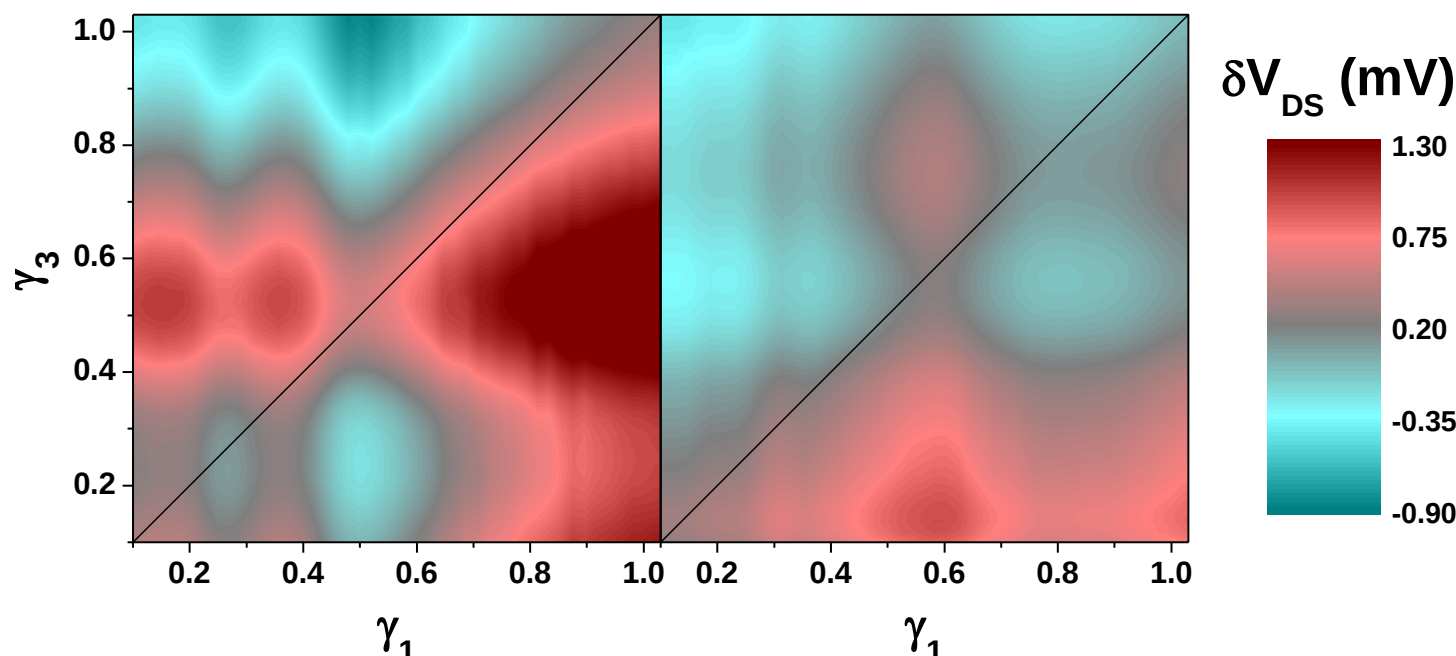
Two Path Plasmonic Mixing



327 GHz

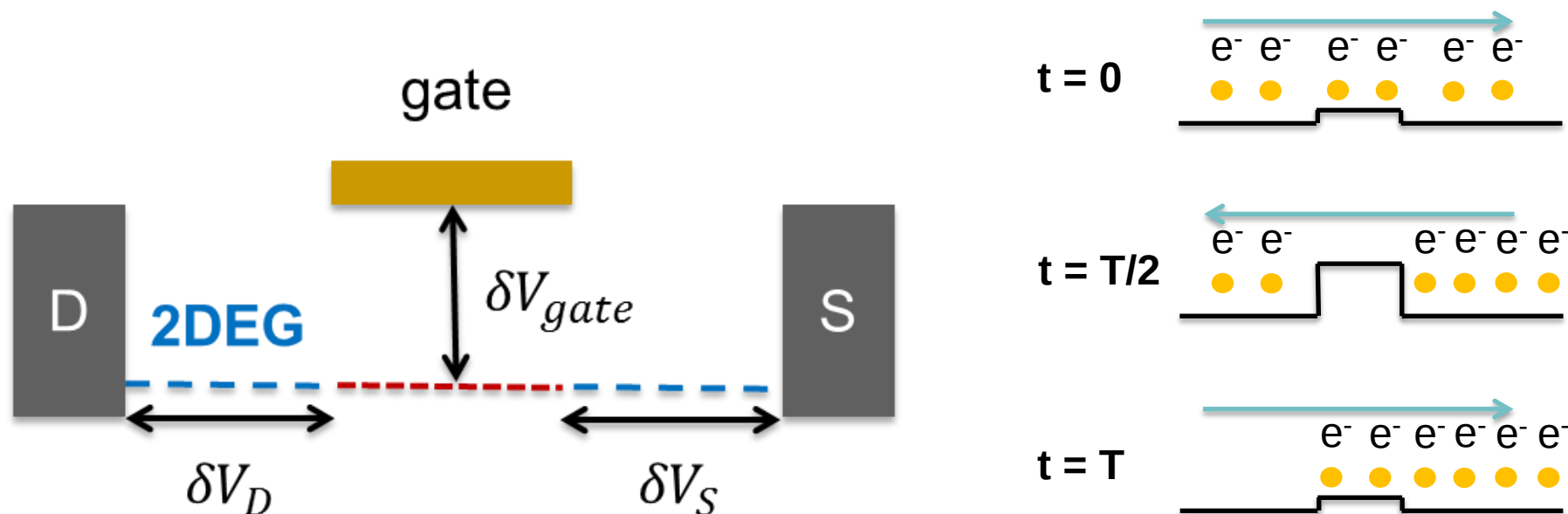


370 GHz



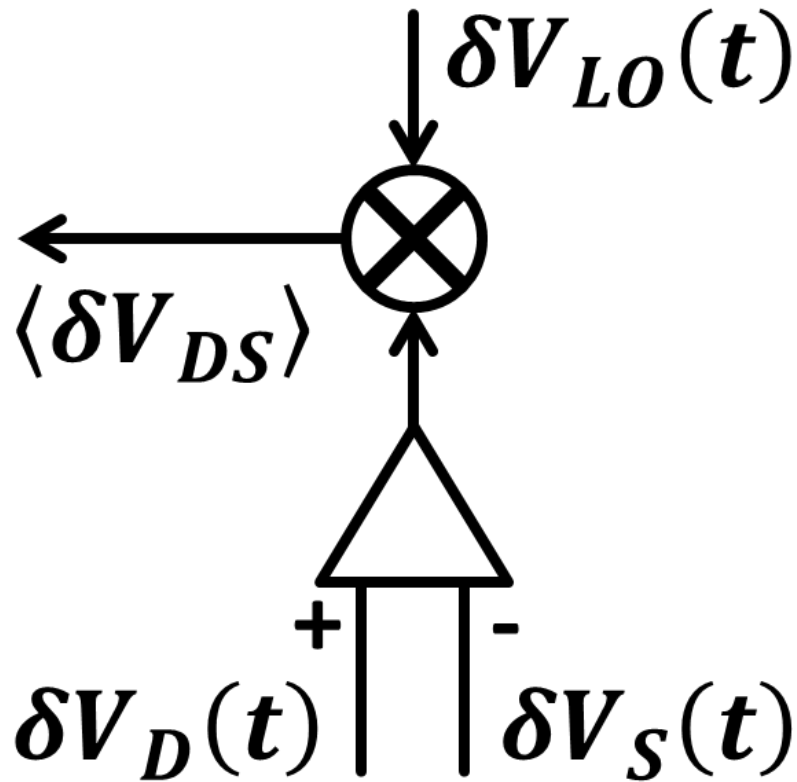
Physical path length L fixed, electrical length qL is tuned

Plasmonic Homodyne Mixing

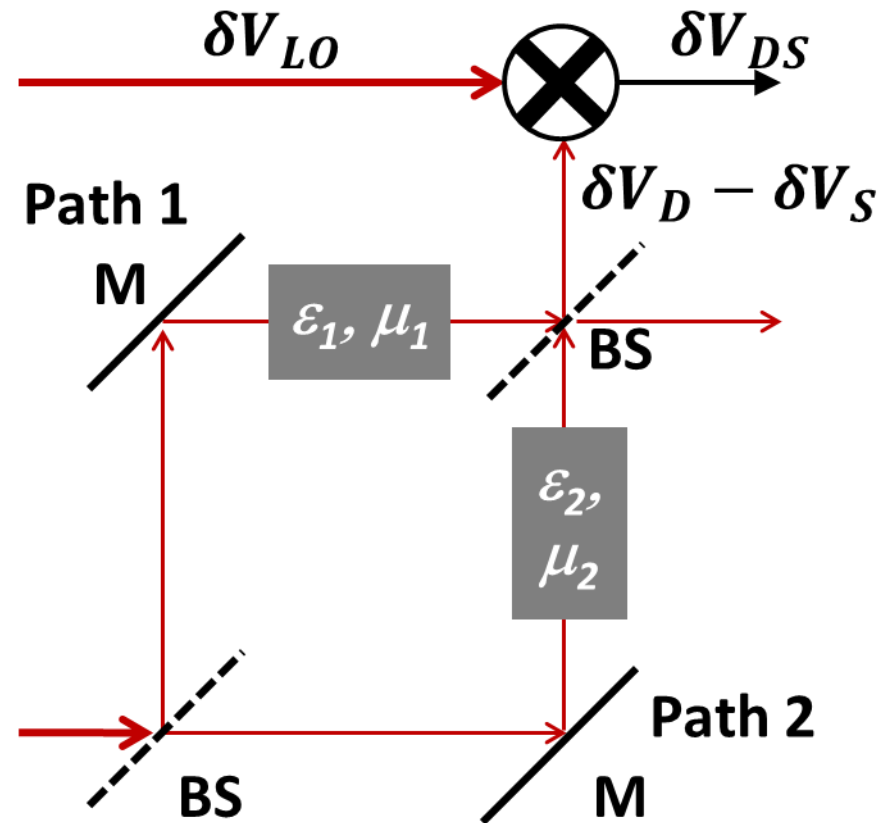


THz field coupled to channel ([blue](#)) drives current across depleted region ([red](#)) which is rectified by gate THz field

Mach-Zehnder Analogue



Equivalent Optical Interferometer



Inductance Comparison

Treat gated 2DEG like a single wire loop



Inductance per unit length
along x-direction

$$L_{kinetic} = \frac{m^*}{e^2 n_{2D} W} \sim 50 \mu_0$$

$$L_{geometric} = \frac{d}{W} \mu_0 \sim 0.05 \mu_0$$

Magnetic flux per unit length
through x-z plane

$$\Phi_m \sim \frac{\mu_0 I d}{W} = I L_{geometric}$$

Growth Sheet



EPI-A GROWTH SHEET

EA0741

SAMPLE ID: EA0741

OBJECTIVE: High Mobility Structure

DATE: 02-May-01

CUSTOMER: Lilly,M

GROWER: Reno, J.

MATERIAL: GaAs, AlGaAs

Structure						
#	Material	Thickness	Dopant	Density	Temp.	SL
1	GaAs	10.0 nm	undoped		635 °C	Cap
2	AlGaAs	98.0 nm	undoped		635 °C	24.1%
3	delta-dope	0.0 nm	Si (n)	1.00E+12	635 °C	
4	AlGaAs	75.0 nm	undoped		635 °C	Upper Setback 24.1%
5	GaAs	30.0 nm	undoped		635 °C	QW
6	AlGaAs	95.0 nm	undoped		635 °C	Lower Setback 24.1%
7	delta-dope	0.0 nm	Si (n)	1.00E+12	635 °C	
8	AlGaAs	98.0 nm	undoped		635 °C	24.1%
9	GaAs	3.0 nm	undoped		635 °C	[x300 Smoothing & EtchStop SL
10	AlGaAs	10.0 nm	undoped		635 °C	] 54.8%
11	GaAs	100.0 nm	undoped		635 °C	Buffer

Substrate Information			Source Information				
Ingot:	G079K3400	Source	Temp °C	Flux	G.R. ML/s	Rheed	UFG
Slice:		Al7	1155.0		0.242	Yes	Yes
Material:	GaAs	As2_Va	175.0		1.760	Yes	Yes
Orientation:	(100)	As2C	650.0		0.000	No	Yes
Thickness:	625 µm	As2X	390.0		0.000	No	Yes
Diameter:	3"	Ga3	1158.0		0.762	Yes	Yes
Type:	undoped	Ga9	1090.0		0.199	Yes	Yes
Vendor:	AXT	Rotate	-15.0		0.000	No	Yes
		Rotate	-5.0		0.000	No	Yes
		Si	1212.0		0.000	No	Yes

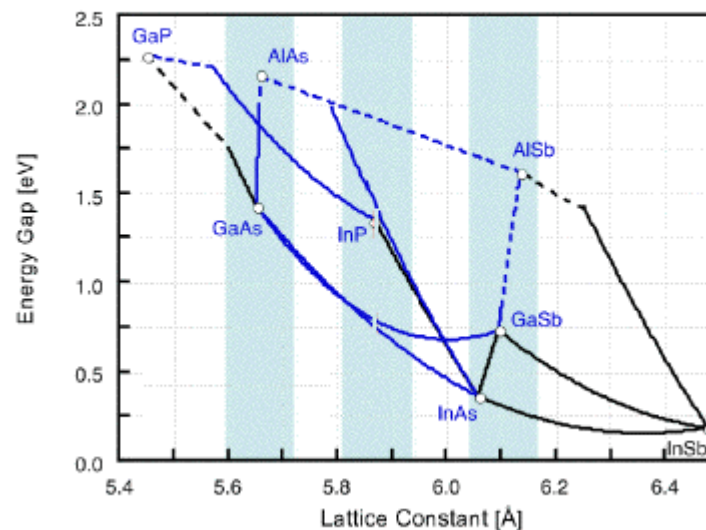
Notes and Comments:

Duplicate of EA0739 but grown with lower As flux.

Characterization Results:

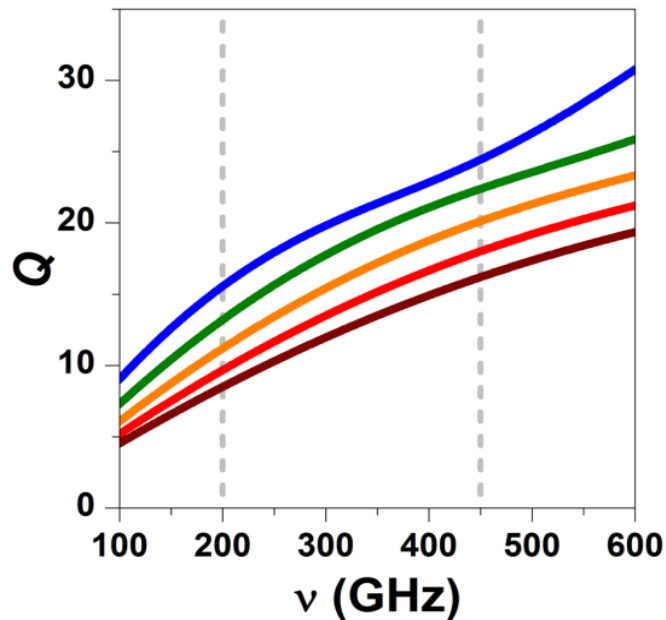
GaAs / Al_xGa_{1-x}As

x = .24



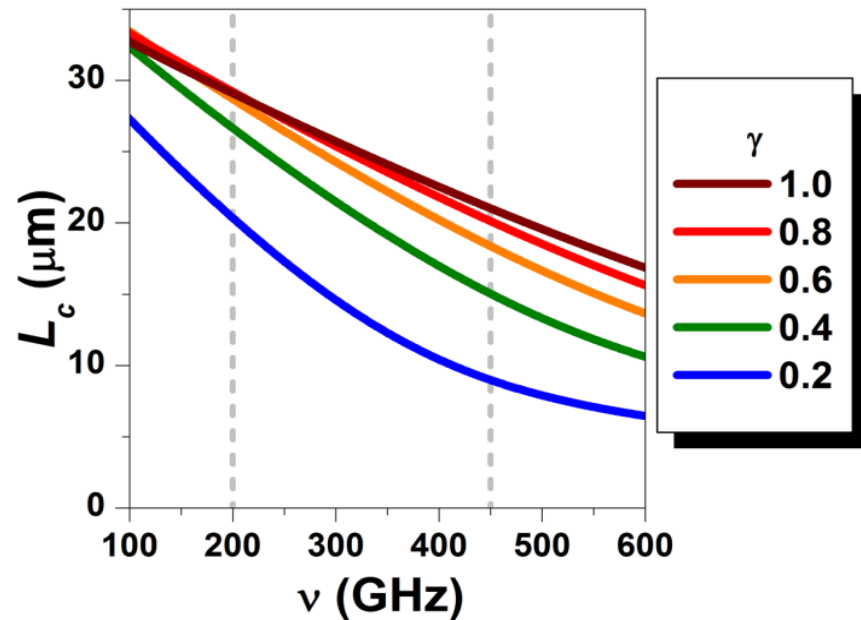
Q & Coherence Length

a



$$Q = \frac{\beta}{2\alpha}$$

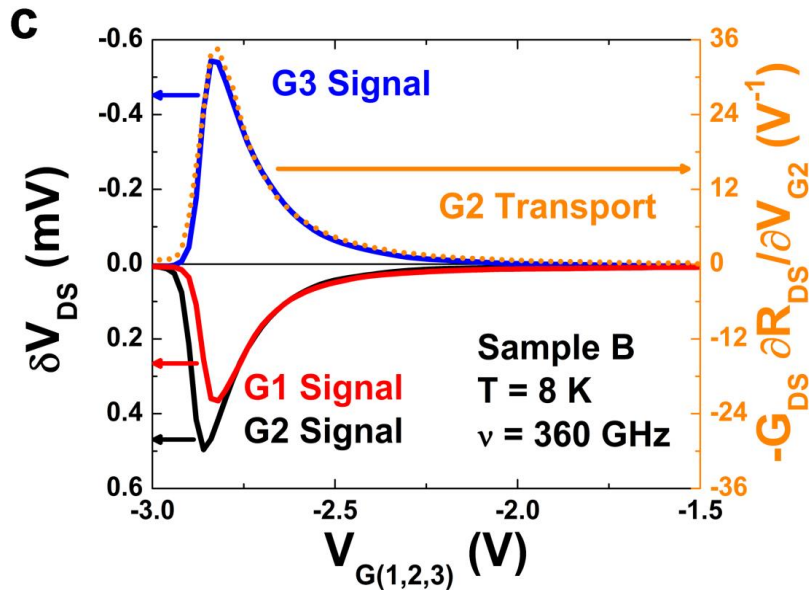
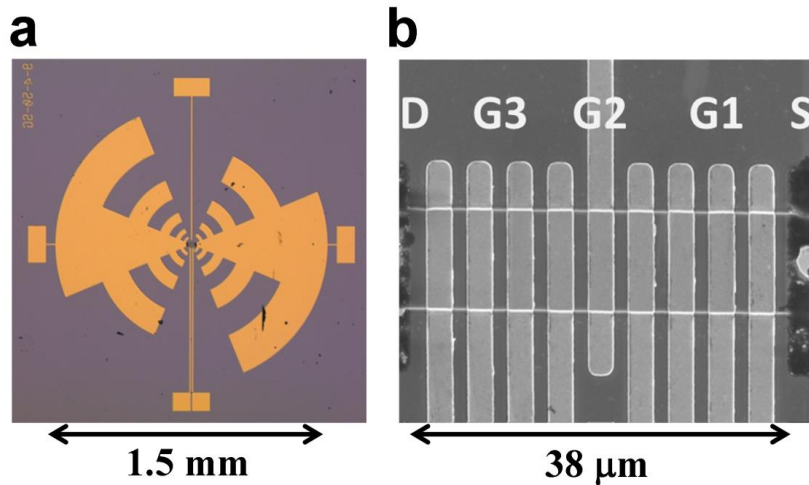
b



$$L_c = \frac{1}{\alpha}$$

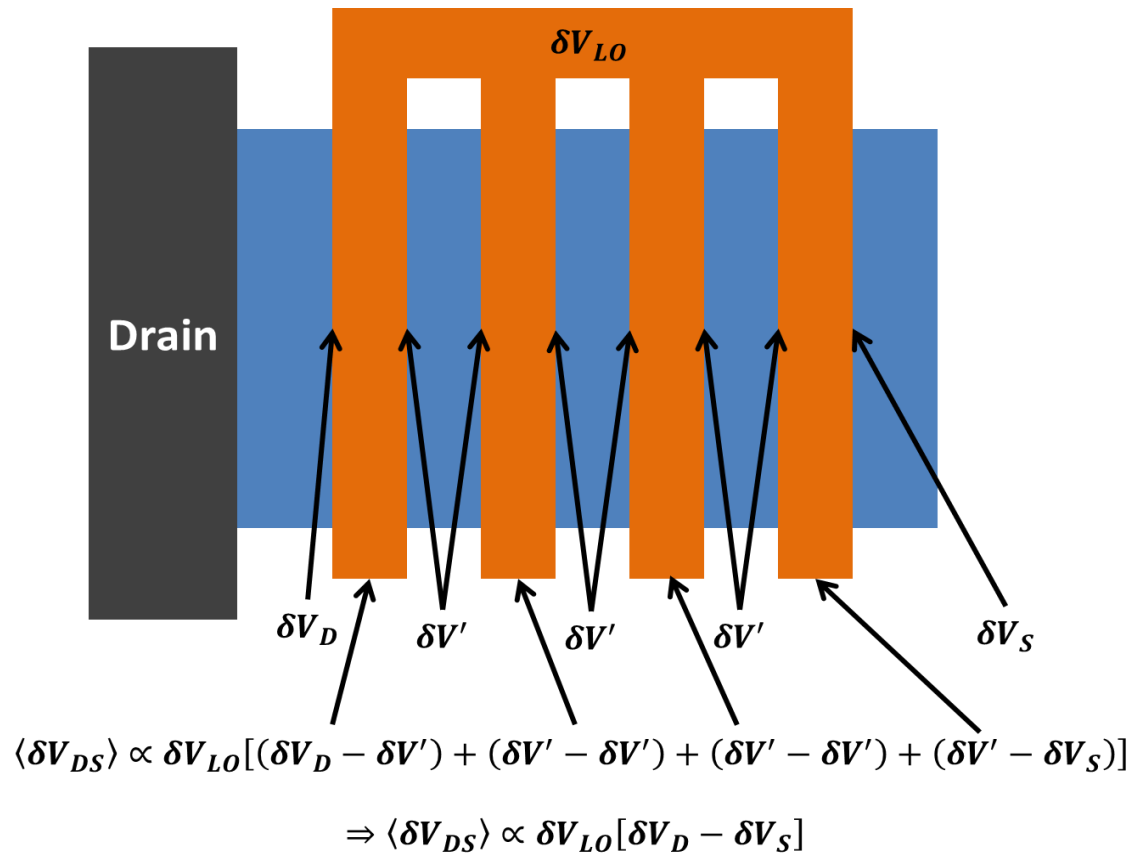
$$q \equiv \beta + i\alpha$$

3 gates, response vs. transport



$$\langle \delta V_{DS} \rangle = -G_{DS} \frac{\partial R_{DS}}{\partial V_{G2}} \langle \delta V_{LO}(t) [\delta V_D(t) - \delta V_S(t)] \rangle$$

Four stripe rectifier



Scalability

