

Surpassing the High-Temperature Thermoelectric Figure-of-Merit by Nano-Structured Phononic Crystals

Ihab El-Kady

Principal Member of Technical Staff

Sandia National Laboratories

PO Box 5800, MS 1082

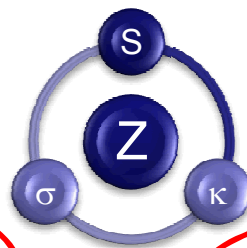
Albuquerque, NM 87185

Ph: (505) 284-4308

Fax: (505) 284-7690

Email: ielkady@sandia.gov





Objective High-ZT: ZT – a 3D dilemma!

The ZT-Problem:

The task is to min κ and max S & σ :

➤ The κ - σ Problem:

- » In semiconductors, “ κ ” is dominated by the phonon contribution
- » → to reduce “ κ ”, we need to increase phonon scattering!
- » → leads to an increase in electron scattering
- » → $\kappa \downarrow \Leftrightarrow \sigma \downarrow$

➤ The S - σ Problem:

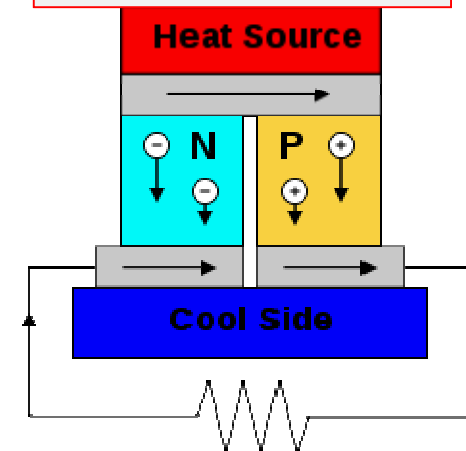
- » S : is like a **measure of the electronic entropy (disorder/scattering) in the system**
- » σ : is a measure of how fast we can propagate a **current** → **decreases with increasing disorder!**
- » **Mott Equation:**

$$S \sim \left. \frac{d \ln(\sigma(\varepsilon))}{d\varepsilon} \right|_{\varepsilon=\varepsilon_F} = \frac{1}{\sigma} \left. \frac{d \ln(\sigma(\varepsilon))}{d\varepsilon} \right|_{\varepsilon=\varepsilon_F}$$

- » → $S \uparrow \Leftrightarrow \sigma \downarrow$

TE Peltier Effect:

TE Generator



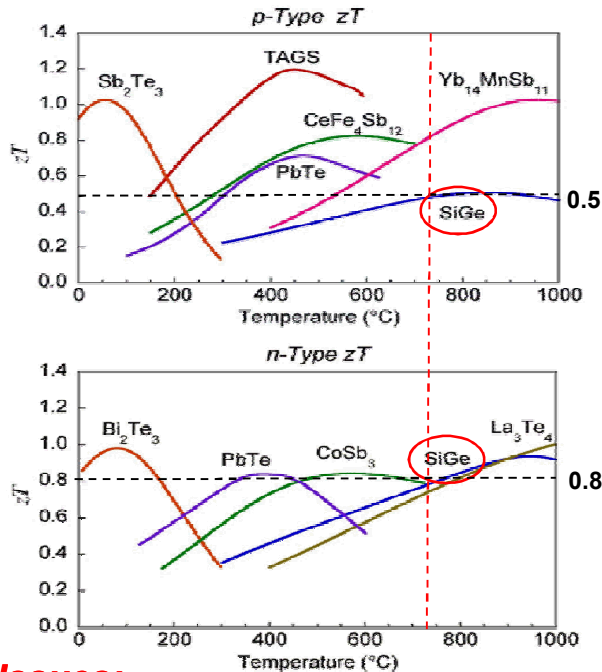
- **Cooler:** Apply DC voltage → heat moves from cold-side to hot-side
- **Generator:** Heat gradient → current flow
- **Efficiency:** Quantified by the dimensionless quantity ZT
- **Figure of merit :** $Z = S^2\sigma/\kappa$
- S : Seebeck coefficient ($\Delta V / \Delta T$)
- σ : Electrical conductivity
- κ : Thermal conductivity ($= \kappa_e + \kappa_{ph}$)

Z is a 3D problem in nature, meanwhile almost all literature attempts deal with one parameter at a time!

Objective High-Temperature High-ZT : “Traditional Approaches” and their Drawbacks

Bulk TE Materials:

G. Jeffrey et. al., Nature Materials, Feb 2008



Issues:

» Fabrication/Integration:

- » Compatibility with standard Fab.
- » Need p-n junction → SiGe

» Common Literature Mistake:

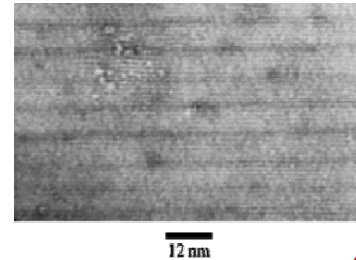
- » Mention only n-type or p-type ZT
- » $ZT_{SiGe-n} \sim 0.8$ & $ZT_{SiGe-p} \sim 0.5$
- » $ZT_{SiGe-pn} \sim 0.7$ (@1000K) low!

$$Z\bar{T} = \frac{(S_p - S_n)^2 T}{[(\rho_n \kappa_n)^{1/2} + (\rho_p \kappa_p)^{1/2}]^2}$$

» K-σ Problem: Super Lattices:

- » “Hard interfacial boundary” → Lots of electron scattering
- » Fabrication:
 - » “Small layer thicknesses” → ALD: Non scalable!
- » Integration:
 - » “Small separation between hot and cold sides” → low performance!

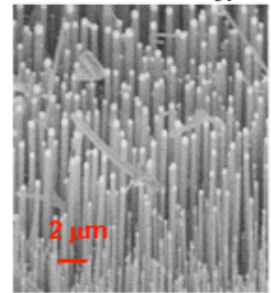
10Å/50Å
Bi₂Te₃/Sb₂Te₃
Structure



» K-σ Decoupled: Nano-Wires:

- » “Unimpeded transport of electrons”
- » Fabrication:
 - » “Too many wires” → lose quantum effects and approach thin films
- » Integration:
 - » “Requires vertical integration” → Practically impossible

Doped Si
nanowire array

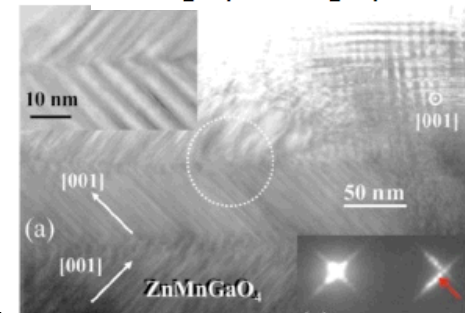


Ceramics (Oxides):

» K-σ Problem (low κ ~10w/mK & low σ):

- » “Incoherent scattering” → Affects both electrons and phonons
- » Fabrication:
 - » “Nanostructures self-assembly” → Nondeterministic outcome
- » Integration (vertical):
 - » “very thin films ~0.5 μm” → Questionable!

ZnMn₂O₄-ZnGa₂O₄



Reducing κ :

Phononic Crystal Approach A lateral (planar) Device

➤ The thermal conductivity is given by

$$\kappa = \frac{1}{6\pi} \sum_j \int_q \frac{\hbar^2 \omega_j^2(q)}{k_B T^2} \frac{\exp\left[\frac{\hbar \omega_j(q)}{k_B T}\right]}{\left(\exp\left[\frac{\hbar \omega_j(q)}{k_B T}\right] - 1\right)^2} v_j^2(q) \tau_j(q) q^2 dq$$

$$\frac{1}{\tau_j(q)} = \frac{1}{\tau_{Umklapp,j}} + \frac{1}{\tau_{impurity,j}} + \frac{1}{\tau_{boundary,j}}$$

$$\tau_{boundary,j} = \frac{L}{v_j(q)}$$

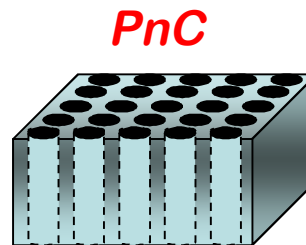
$\omega(q)$ is the phonon dispersion
 $v(q) = \partial \omega(q) / \partial q$ is the phonon group velocity
 $\tau(q)$ is the scattering time of the phonons
 q is the wavevector
 $j = 1, 2, 3$ (1 longitudinal, and 2 transverse)
 L is the minimum feature size

➤ To change κ need to change :

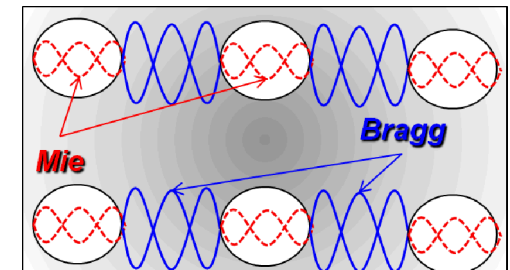
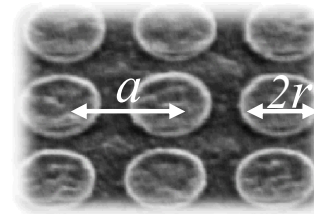
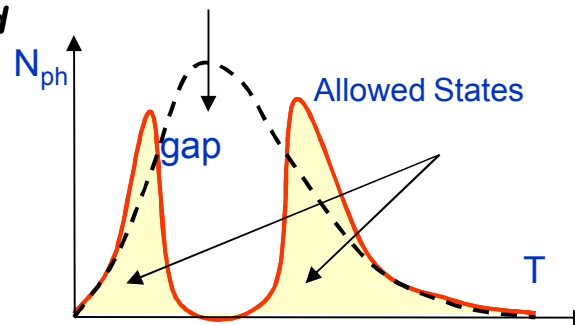
- Dispersion " $\omega(q)$ "
- Minimum feature size " L "

➤ What is a Photonic Crystal?

- Created by superposing Mie and Bragg resonant scattering by a periodic arrangement of scattering centers in a lattice.
- Requires sufficient acoustic impedance mismatch.

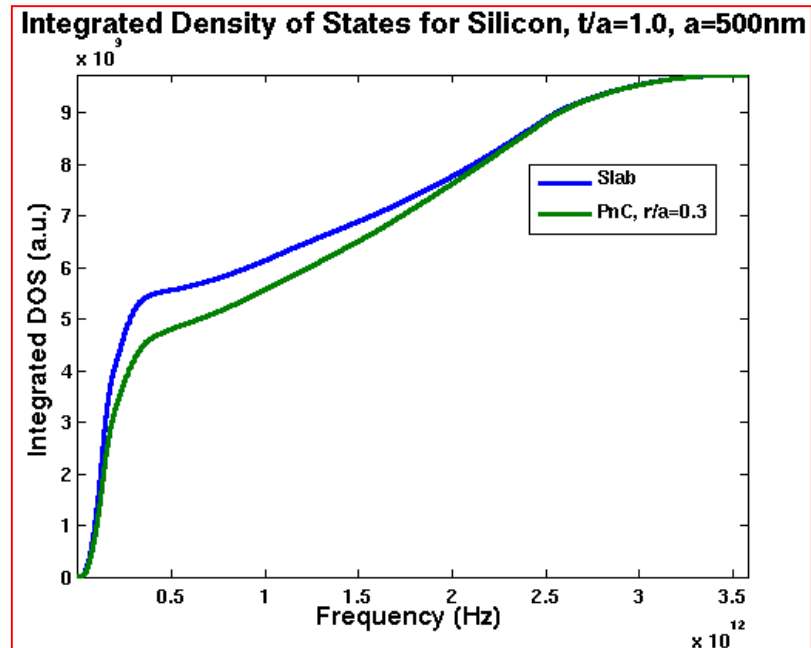
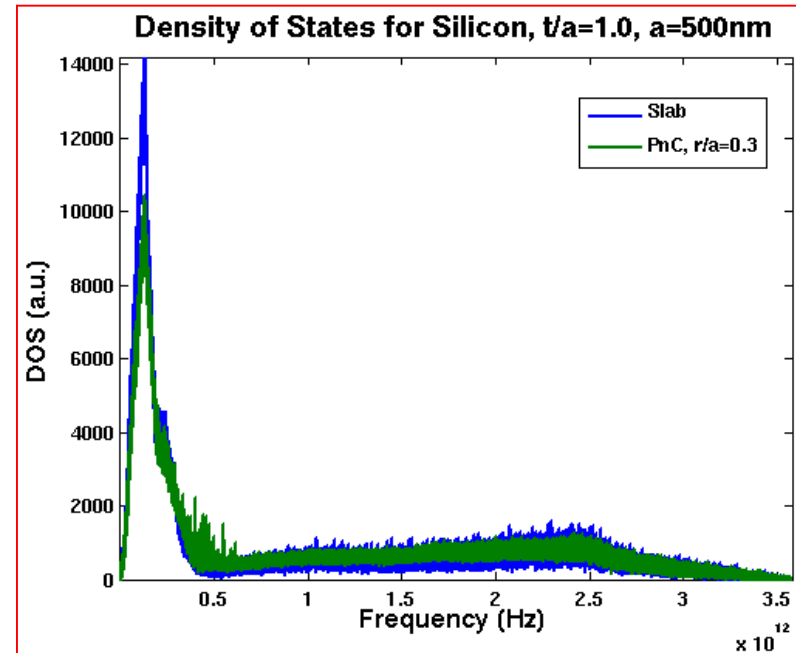
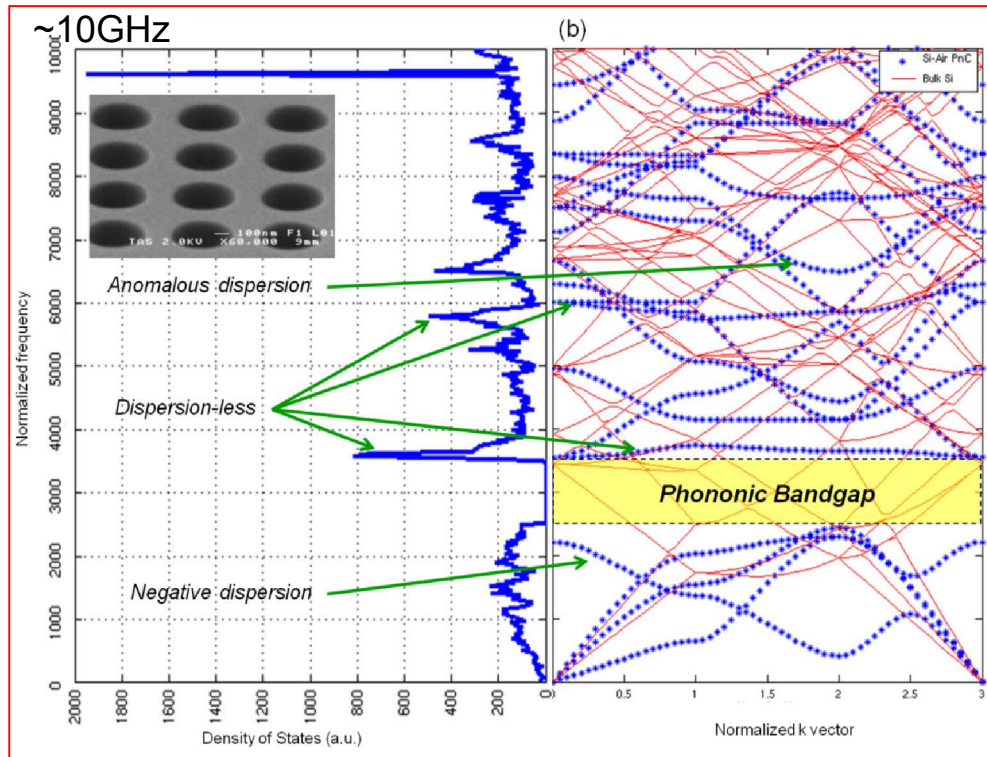


Forbidden States



Reducing κ : Planar PnC Approach

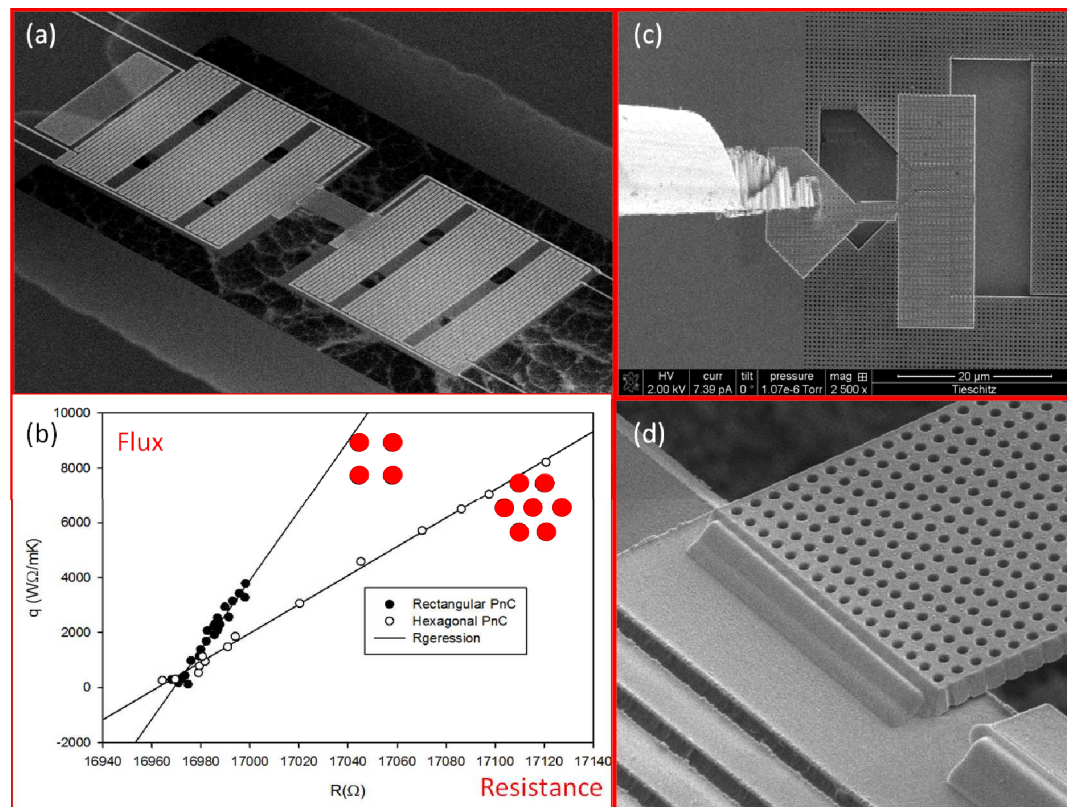
- Dispersion Engineering: How Coherent Scattering Works



- **Bragg Resonance** → Creation of gap (long wavelength acoustic phonons)
- **Flat Bands** → Less in number but higher in density
→ Reduced group velocity
- **Negative Bands** → Backward propagation (backward scattering)

Reducing κ : Planar PnC Approach

- **Additional Evidence of Coherent Scattering**
- **Suspended-island technique:**
 - The **only thermal path** between the heating stage and the detector stage is the PnC sample.
 - **Pick-and-place** technique using a Omniprobe needle tip used to place the PnC samples between the suspended islands.
 - **SEM** close-up showing the quality of the FIB welding of the PnC sample to the islands using platinum.
 - **Comparison** of the square PnC lattice with a triangular PnC lattice of the same thickness, hole diameter, critical dimension L (minimum feature size).



- **Since all other parameters are constant:**

- Same hole size
- Same slab thickness
- Same critical length
- Same material

➢ **Then:**

$$2.88 \leftarrow \frac{\kappa_{Sq}}{\kappa_{Hex}} = \frac{ff_{Hex}}{ff_{Sq}} \rightarrow 1.15$$

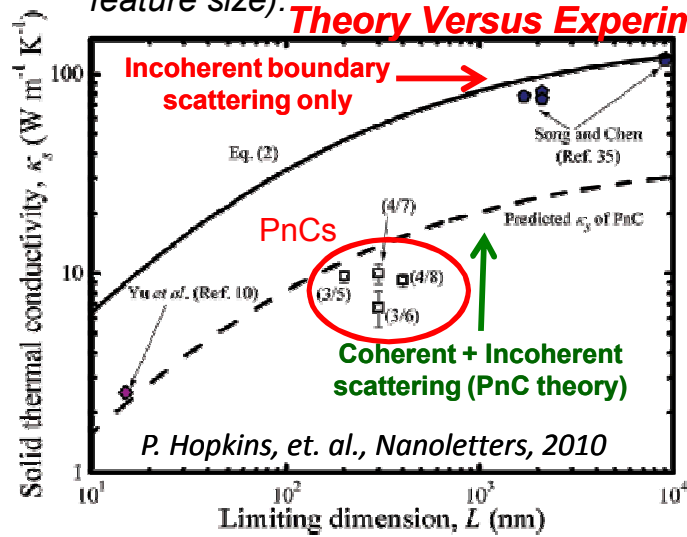
- **Off by a factor of 2.5!**
- **Must be topology!**
- **Dispersion !**
- **Coherent !**

- **2D Air Fraction:**

$$ff_{sq} = \frac{\pi r^2}{a^2} \quad ff_{hex} = \frac{2\pi r^2}{\sqrt{3}a^2}$$

- **Ratio:**

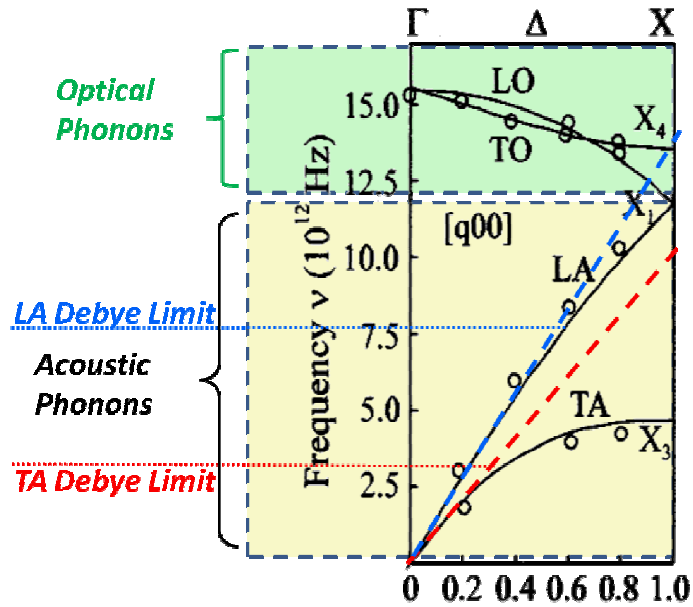
$$\frac{ff_{hex}}{ff_{sq}} = \frac{\frac{2\pi r^2}{\sqrt{3}a^2}}{\frac{\pi r^2}{a^2}} = \frac{2}{\sqrt{3}} = 1.155$$



Reducing κ :

How can such widely spaced holes affect high frequency phonons?

- When do our predictions breakdown?:
- Sensitive to the lattice periodic topology



Si Phonon Dispersion **LA**: Longitudinal Acoustic Mode; **TA**: Transverse Acoustic Mode.

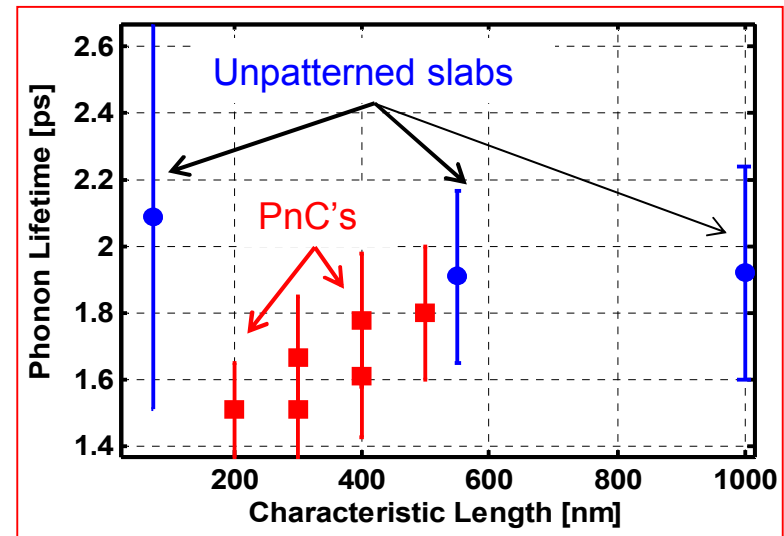
PWE-PnC theory predictions remain:

- Quantitatively valid :
 - In the Debye Limit portion of the dispersion:
 - ~ 7.5 THz for LA
 - ~ 3.5 THz for TA
- Qualitatively valid :
 - In the linear/slow portion of the material dispersion

Reduction Mechanism:

1. PnC-coherent scattering directly affects low frequency acoustic phonons
2. PnC-coherent scattering affects high frequency optical phonons in an indirect manner:
 - Optical Phonons are responsible for upto $\sim 27\%$ of cross-plane κ [Sellan, et. al, JAP 108, 113524 (2010)].
 - $\sim 30\%$ of all optical phonon scattering involves an acoustic phonon mode [Broido, PRB 81, 085205 (2010)].
 - PnC anomalous dispersion affects low laying acoustic modes
3. Incoherent PnC boundary scattering affects high frequency acoustic and optical phonon modes

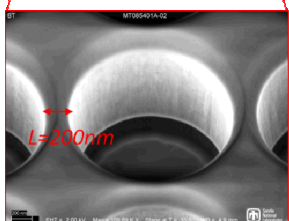
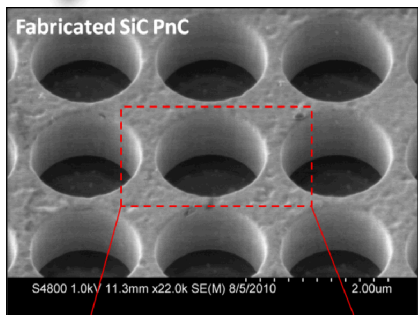
End result: Is a reduction in the overall thermal conductivity



Anticipated Performance

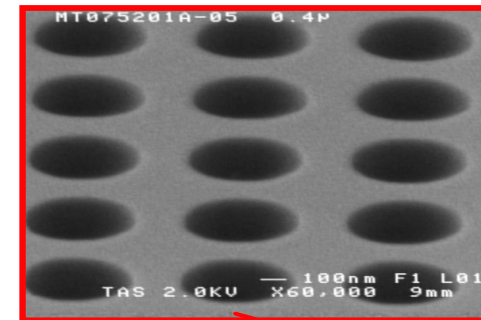
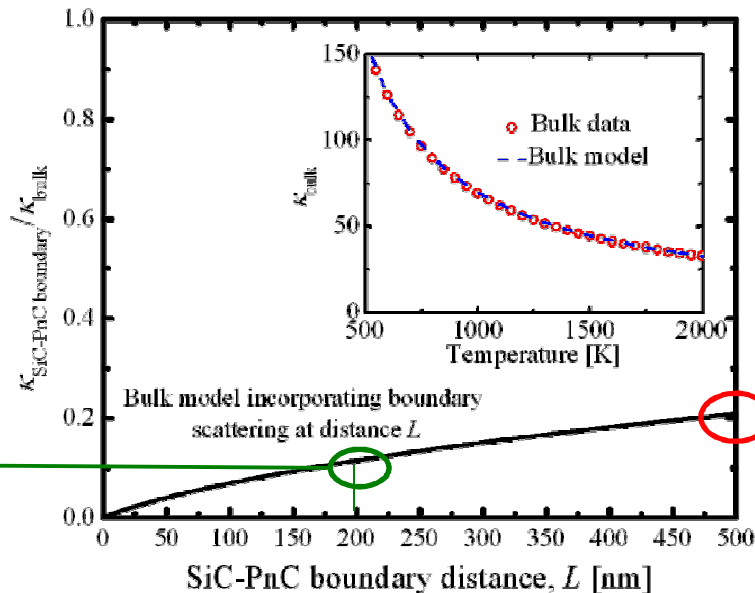
Semiconductor-Air PnC ZT Estimate

- Predicted κ Reduction: Holland-Callaway Equation

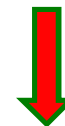


Controlled Oxidation

$a=1\mu\text{m}$, $r=250\text{nm}$, $L=500\text{nm}$



Standard Photolithography
 $a=1\mu\text{m}$, $r=250\text{nm}$, $L=500\text{nm}$



- Ok, but what happens to the electrons @ $L \sim 500\text{nm}$?
- Mean free path of the electrons is $\sim 10\text{nm}$ or less @ $1000\text{K} \rightarrow \sigma$ should be affected at most by the porosity factor

For $a = 1\mu\text{m}$ & $r = 250\text{nm}$

$\Rightarrow ff_{\text{air}} \sim 0.2$ & $L = 500\text{nm}$

$\sigma_{\text{PnC}} \approx (1 - ff_{\text{air}}) \sigma_{\text{SC-Bulk}} \approx 0.8 \sigma_{\text{SC-Bulk}}$

$\kappa_{\text{PnC}} \sim 0.2 \kappa_{\text{SC-Bulk}}$

$S_{\text{PnC}} \sim S_{\text{SC-Bulk}}$

A factor of 4 enhancement in $ZT@1000\text{K}$!

Increasing σ While Maintaining Low κ Semiconductor/Metal PnC

- Effective σ :
- Increased doping
→ decreases the interface potential width
→ Electron tunneling
- “Ohmic Contact”!

- At degenerate doping:

$$\rho_{PnC+Metal} \approx ff_{Metal} \rho_{Metal} + (1 - ff_{Metal}) \rho_{SC-Bulk} \approx (1 - ff_W) \rho_{SC-Bulk}$$

- Effective κ ($1/\mathcal{R}$):

- What happens to the thermal resistance now that we have plugged the holes with “metal”?

- Mechanical impedance mismatch is still very high
- Incoherent scattering almost unaffected
- Coherent scattering slightly enhanced

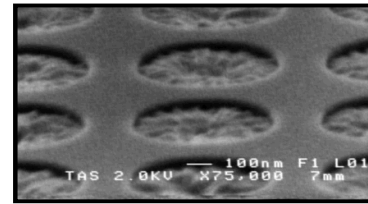
$$\begin{aligned} \mathcal{R}_{PnC+W} &\approx ff_{Metal} \mathcal{R}_{Metal} + (1 - ff_{Metal}) \mathcal{R}_{PnC} \\ &\approx (1 - ff_{Metal}) (5 \mathcal{R}_{SC-Bulk}) \approx \mathcal{R}_{SC-Bulk} \end{aligned}$$

- Effective ZT:

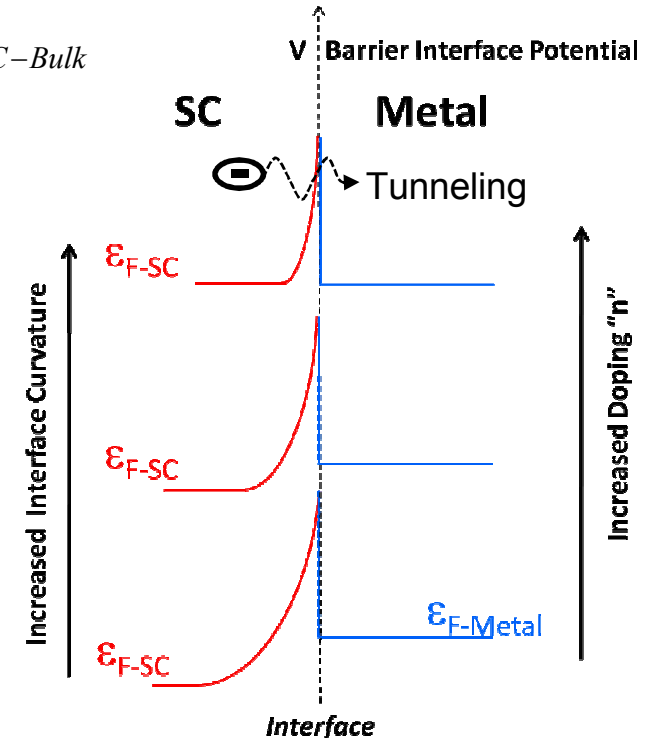
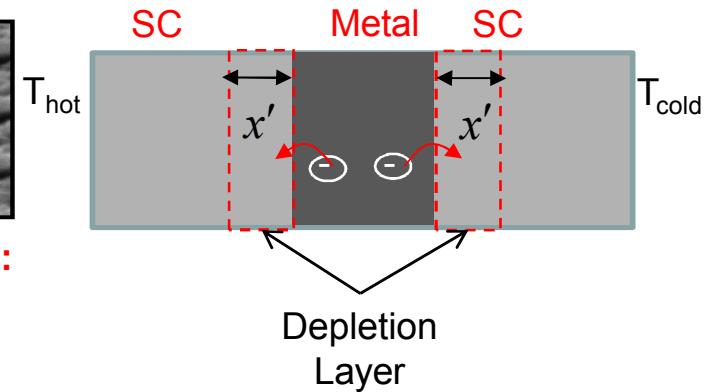
- Assuming: $a = 1 \mu m$ & $r = 250 nm \Rightarrow ff_W \sim 0.2$ & $L = 500 nm$

$$\Rightarrow \kappa_{PnC+Metal} = \frac{1}{\mathcal{R}_{PnC+Metal}} \approx \frac{1}{\mathcal{R}_{SC-Bulk}} \approx \kappa_{SC-Bulk}$$

$$\Rightarrow \rho_{PnC+Metal} \approx 0.2 \rho_{SC-Bulk}$$



Fabricated Example:
Si/W PnC



A factor of **5** enhancement in **ZT@1000K!**

Increasing S (quadratic) Metallic Nano-Particle implantation (Up to a factor of 8 total enhancement!)

If W effect on κ is drastic:

- We can increase S and mildly affect σ and κ by ion implantation of metallic nanoparticles:
 - Al ions implanted in SiC-PnC at IBL-SNL
 - Anneal $\rightarrow$ congregate islands $\sim 5\text{-}20\text{nm}$
 - Conceptual proof data is for PbTe: same arguments hold for a highly-doped SC, although quantum mechanical treatment would be best (SNL-California)

SiC-PnC + Metallic-nanoparticles

$R=1.5\text{nm}$; $x=5\%$

Solid circles:

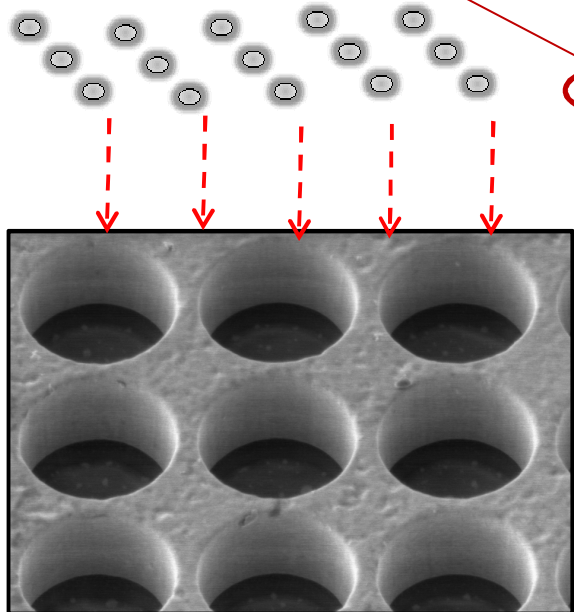
Electron & phonon scattering on the inclusions;

Open circles:

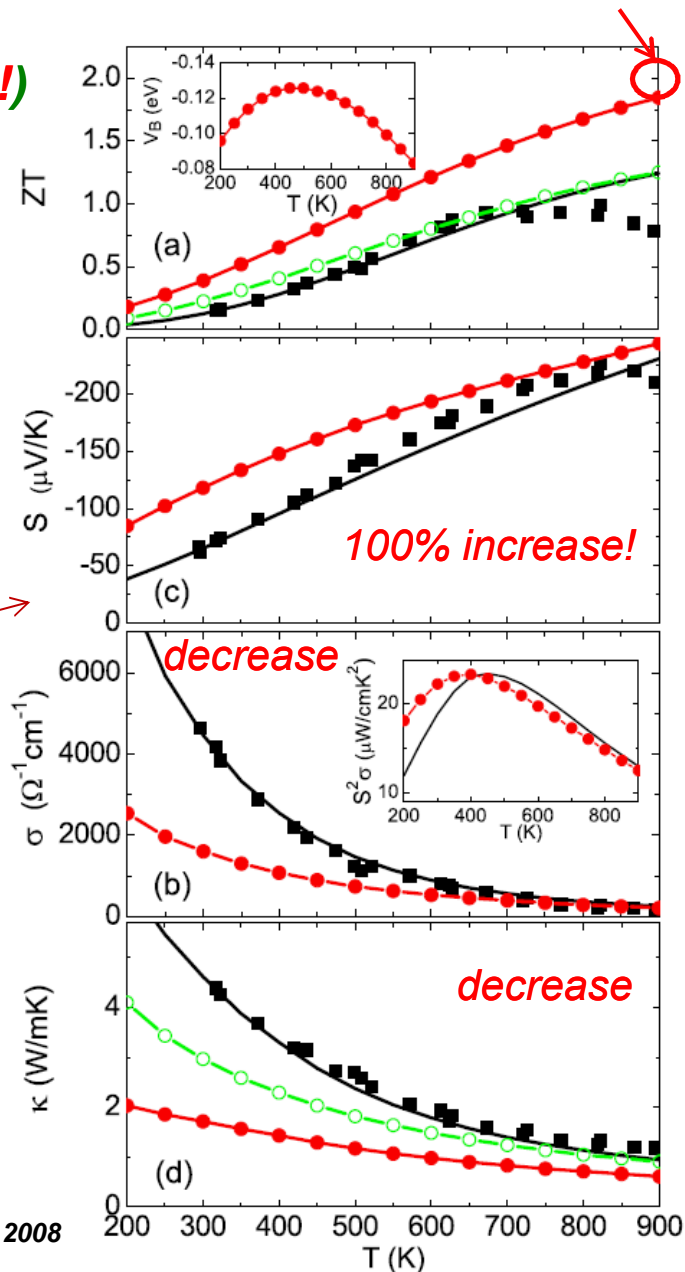
Include electron scattering on the inclusions but with bulk values of κ ;

Solid lines (filled squares):

Experimental & calculated Bulk values

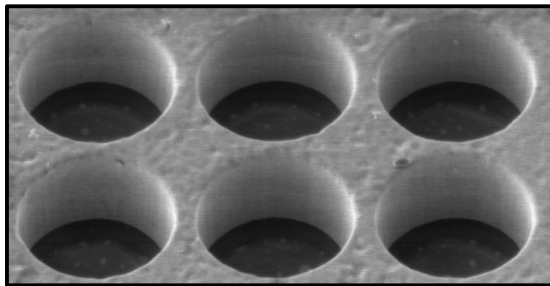


Additional Factor ≥ 2
enhancement in $ZT@1000\text{K}$



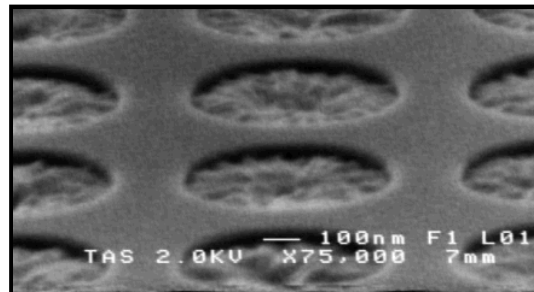
Overall Picture: 3 Possible Approaches in Order of Difficulty

#1 SC/Air PnC



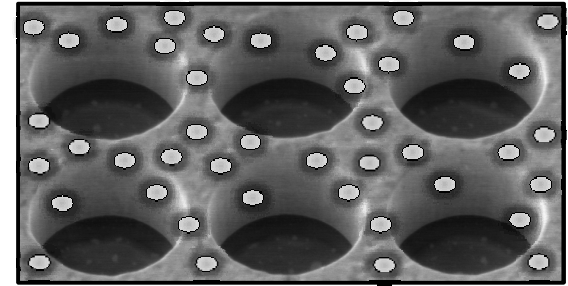
ZT up by 4 @1000K

#2 SC/Metal PnC



ZT up by 5 @1000K

#3 SC/Air PnC + MNPs



ZT up by 8 @1000K

These are **NOT** independent routes, but rather **different** post-processing techniques for the **same PnC** device with **different** optimal parameters!

Possible Material Systems (test beds)

- Si/Air & Si/W:
 - ✓ Easy to process
 - ✓ CMOS fab. compatible
 - ✓ Cheap!
 - ✓ Provides a proof of concept!
 - ✓ Si has a large power factor
 - ✓ Suitable for room T applications
- SiC/Air & SiC/W:
 - ✓ Corrosion resistant
 - ✓ Resilient in harsh environments
 - ✓ Withstands cycling at high T
 - ✓ CMOS fab. compatible
 - ✓ Know how to process well and control minimum feature size
 - ✗ Hard to dope (p -type)

Surpassing the High-Temperature Thermoelectric Figure-of-Merit by Nano-Structured Phononic Crystals

Key Team Members:

PI:

- Ihab El-Kady (PI -1725) ielkady@sandia.gov, (505) 284-4308

Theory:

- Francois Leonard (Electronic Transport - 8656)
- Dan Catalin Spataru (First Principles Calculations - 8656)
- Charles Reinke (PnC Dispersion Engineering - 1725)
- Alan McGaughey (Molecular Dynamics - Carnegie Mellon U.)
- Mahmoud Hussein (Lattice Dynamics - UC-Boulder)

Materials:

- Peter Sharma (TE Material Science - 8626)

Experimental Measurements:

- Thomas Beechem (Phonon Life Time Raman Spectroscopy - 1112)
- Blythe Clark (High Resolution TEM - 1111)
- Zayd Leseman (ZT Temperature-Dependant Characterization – U. New Mexico)
- Patrick Hopkins (Pump Probe thermal conductivity measurements – U. Virginia)

Fabrication:

- Roy Olsson (MDL – Fab - 1749)
- Khalid Hattar (Nanoparticle Implantation and Lattice Defects using FIB - 1111)
- Bongsang Kim (Electronic Characterization and Device Layout - 1749)

Device Models:

- Darren Branch (PNC – TE Device Multi-Physics Modeling - 1714)



Reducing κ : Planar PnC Approach

- Evidence of Coherent Scattering:
- Sensitive to the lattice periodic topology

