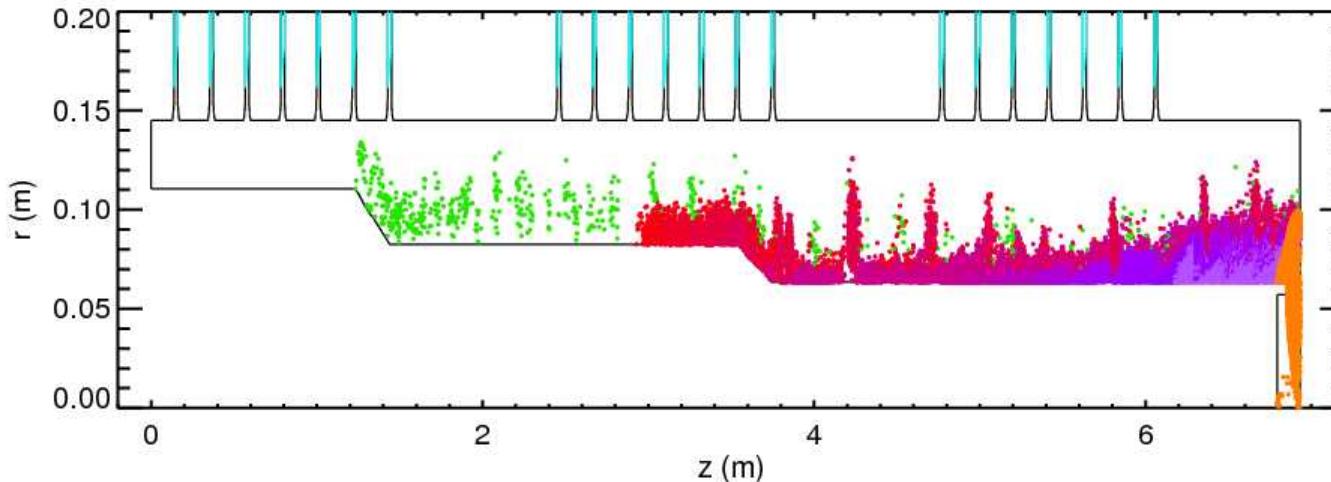


EM-PIC Simulation of Pulsed Power Systems at Sandia National Laboratories

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Tech-X Corporation
Boulder CO, June 29, 2012

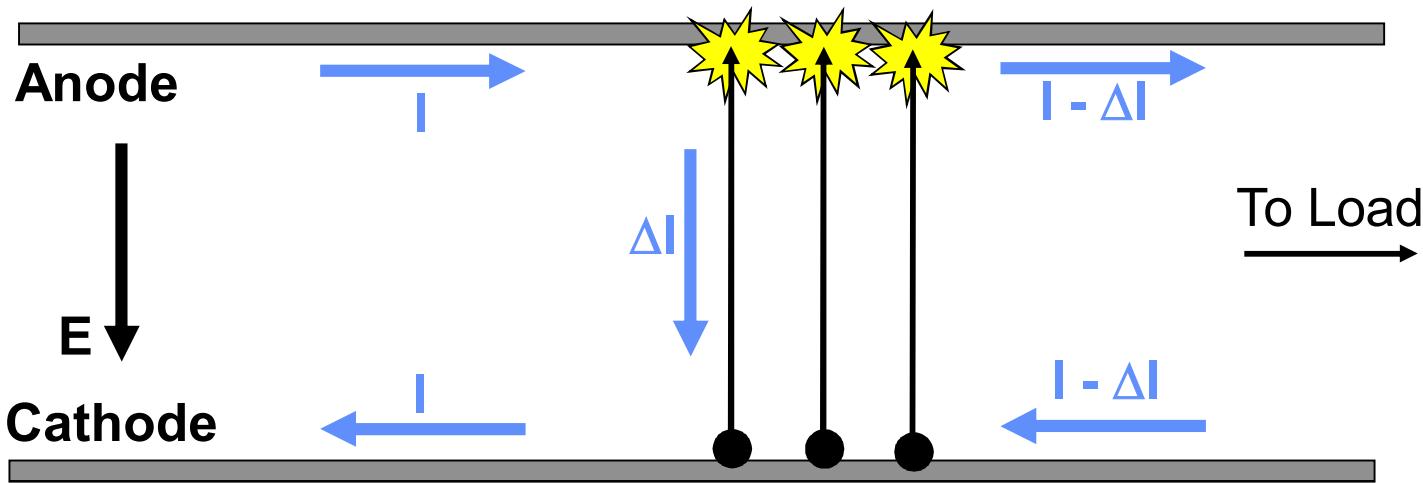
Sandia National Laboratories is a multi-program laboratory managed and operated by Sandia Corporation, a wholly owned subsidiary of Lockheed Martin Corporation, for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-AC04-94AL85000



Outline

- Applications driving analysis and code development:
 - Magnetically insulated transmission lines (MITLs) (ongoing)
 - Light ion beams for fusion: applied-B ion diodes (1987 – 1997)
 - Power flow in the vacuum section of the “Z” accelerator (1997 –)
 - The 21-cavity “Ursa Minor” linear transformer driver (LTD) for radiography (2009 –)
 - Electron beam transport in gas (“gas chemistry”)
- The Quicksilver 2D/3D EM-PIC code
 - Geometry, 1D transmission lines
 - Other features - physics models
- Other PIC work (not discussed here)
 - Implicit ES hybrid fluid-PIC codes: DYNAID (1D), MUFPHI (2D)
 - Unstructured EM-PIC: tetrahedral meshes

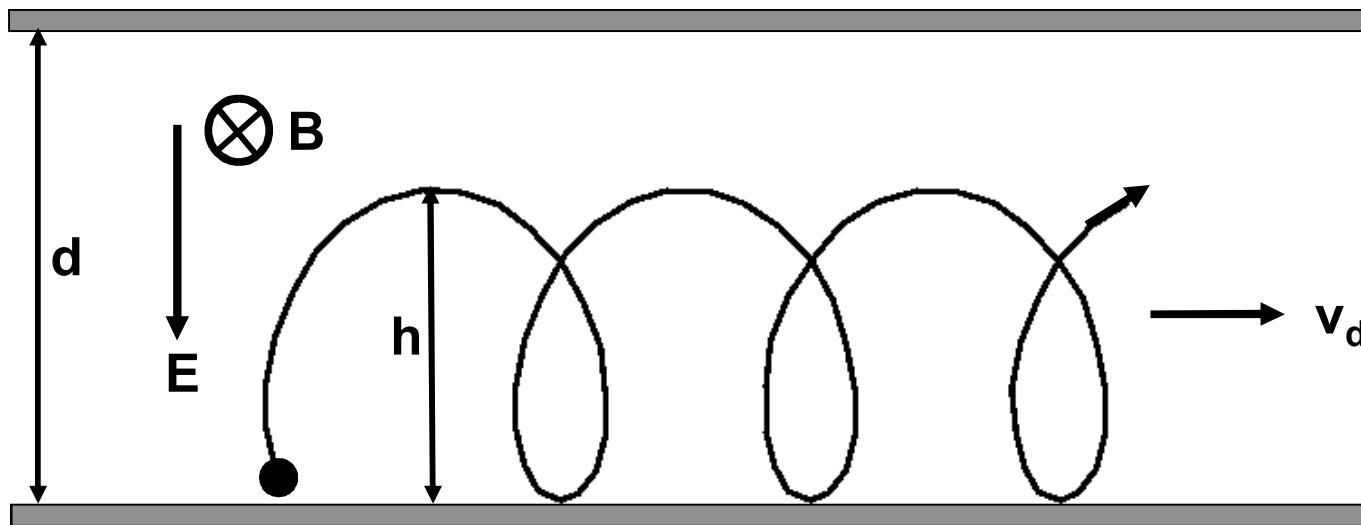
Cathode electron emission is unavoidable for a multi-MV EM pulse in a transmission line



- ***Without special treatment***, a cathode surface emits electrons when the normal electric field E_n exceeds $\sim 50 - 200$ kV/cm
- Electrons flowing to the anode is undesirable
 - Reduces current delivered to the load
 - Heats the anode -- can create an anode plasma



A transverse magnetic field can inhibit electrons from crossing the A-K gap

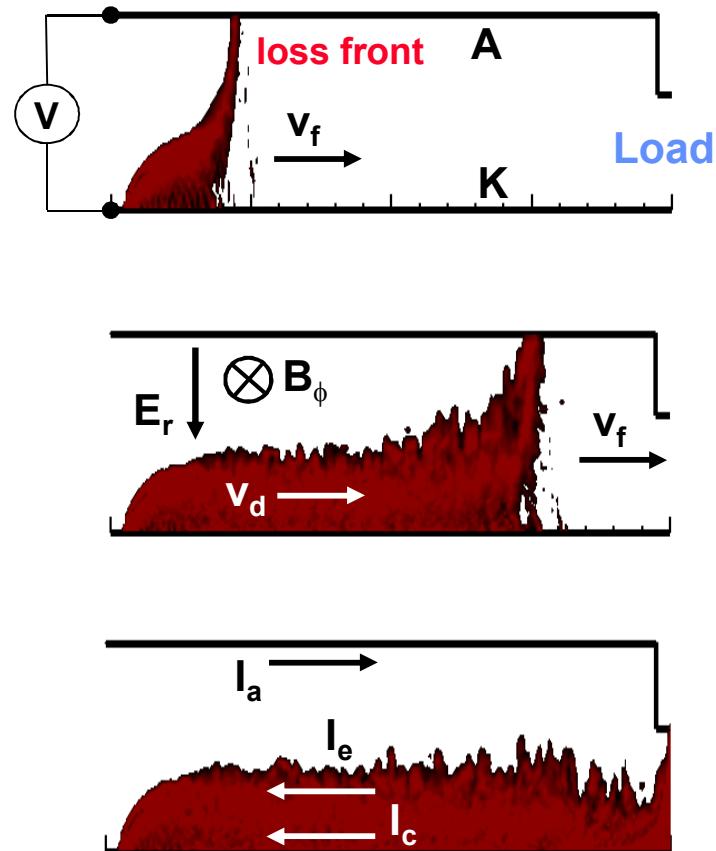


- For $cB > E$, the electrons “drift” with velocity $v_d = (E \times B) / B^2$
- For uniform \mathbf{E} and \mathbf{B} fields: $h = \frac{\frac{2m}{e} E}{B^2 - E^2 / c^2}$
 - For sufficiently large B , $h < d$



Efficient power transfer is possible with “magnetically insulated transmission lines”

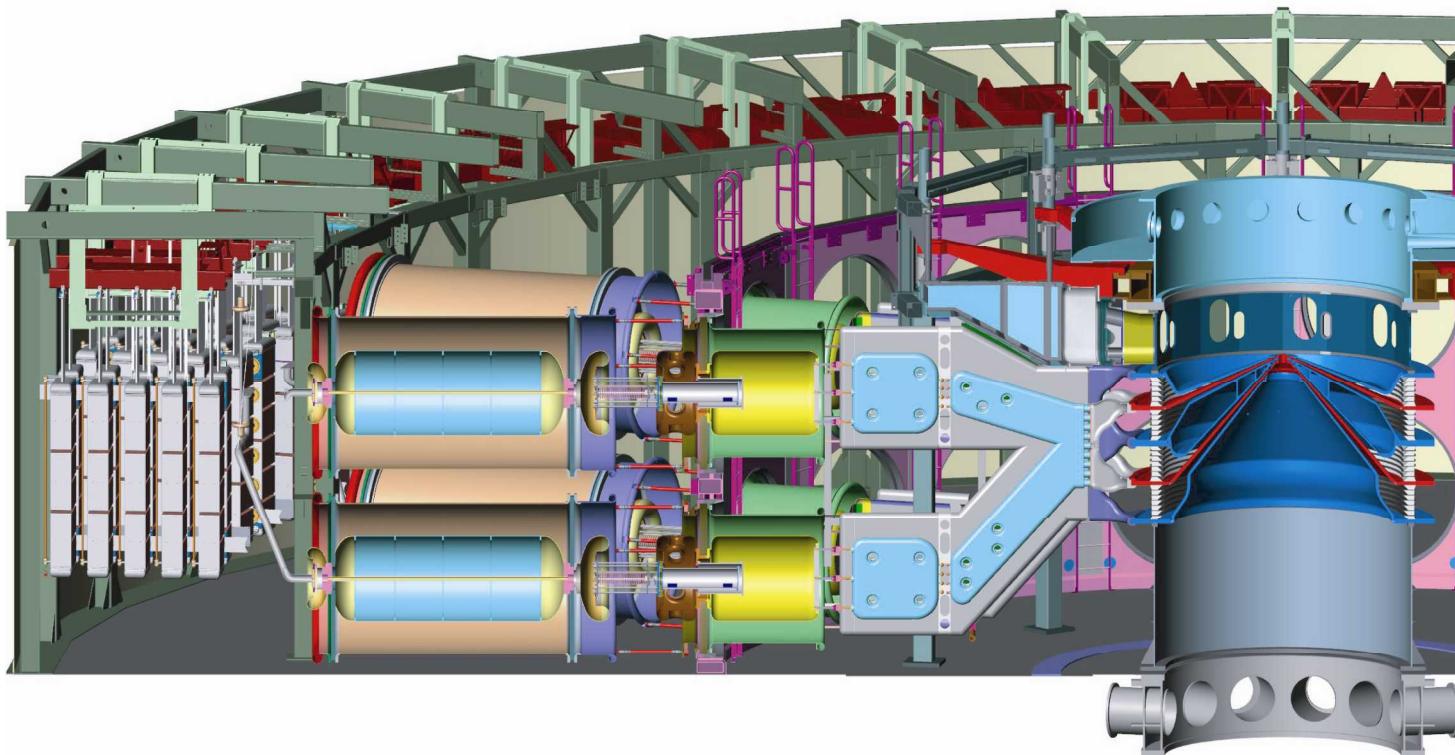
- Schematic illustration of a coaxial “MITL” driving a load
- Behind the “loss front”, B-field of the current flowing in the conductors insulates electrons: they “**ExB** drift” towards the load
- In equilibrium, current flow in the “electron sheath” is related to the anode and cathode conduction currents by: $I_e = I_a - I_c$
- Technical issues for real systems
 - Current loss if load impedance is too high
 - Adding lines in serial or parallel



Courtesy of Steve Rosenthal



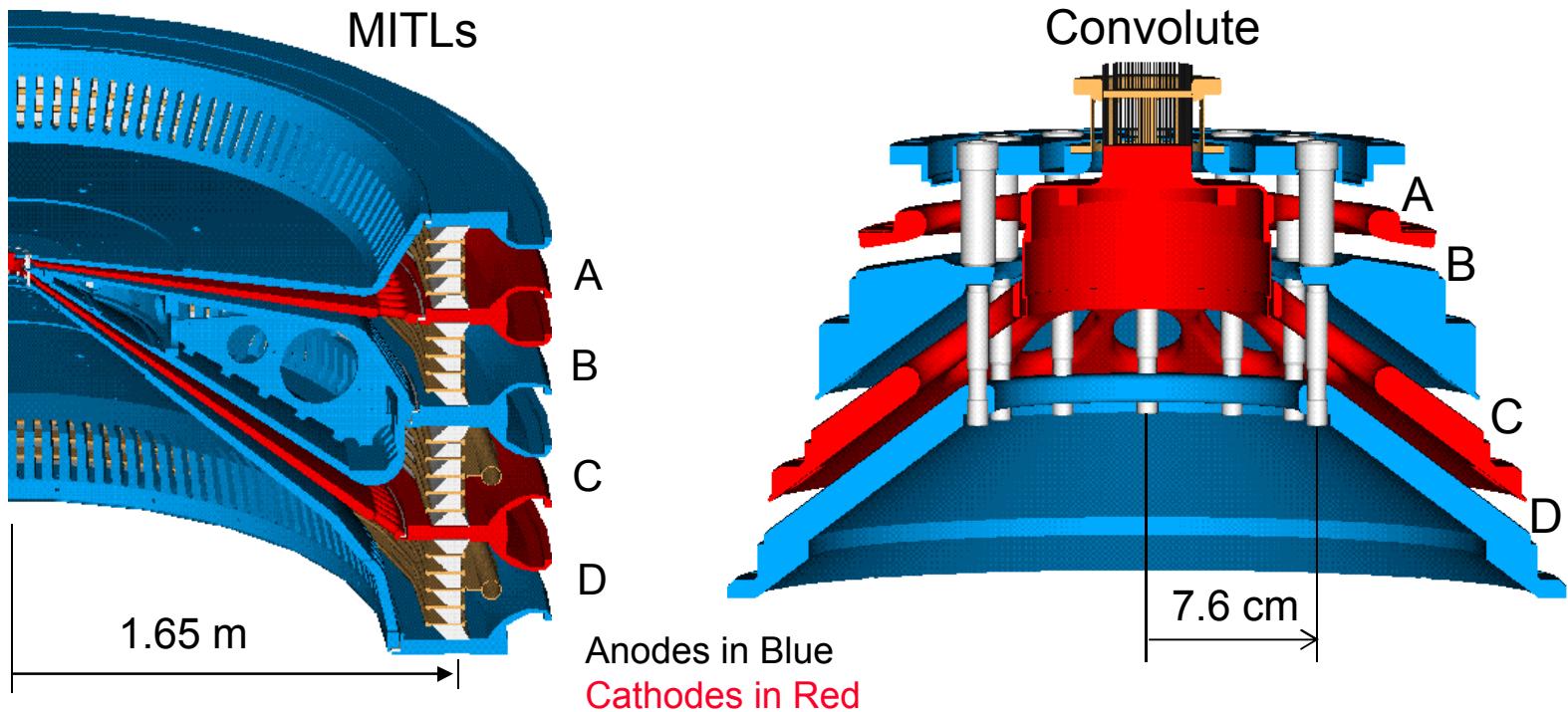
Sandia's Z accelerator is the world's most powerful high-current driver



- 33 m in diameter, 6 m high, 36 pulse-forming lines
- Drives variety of loads, e.g. 26 MA, 100 ns pulse to a Z-pinch



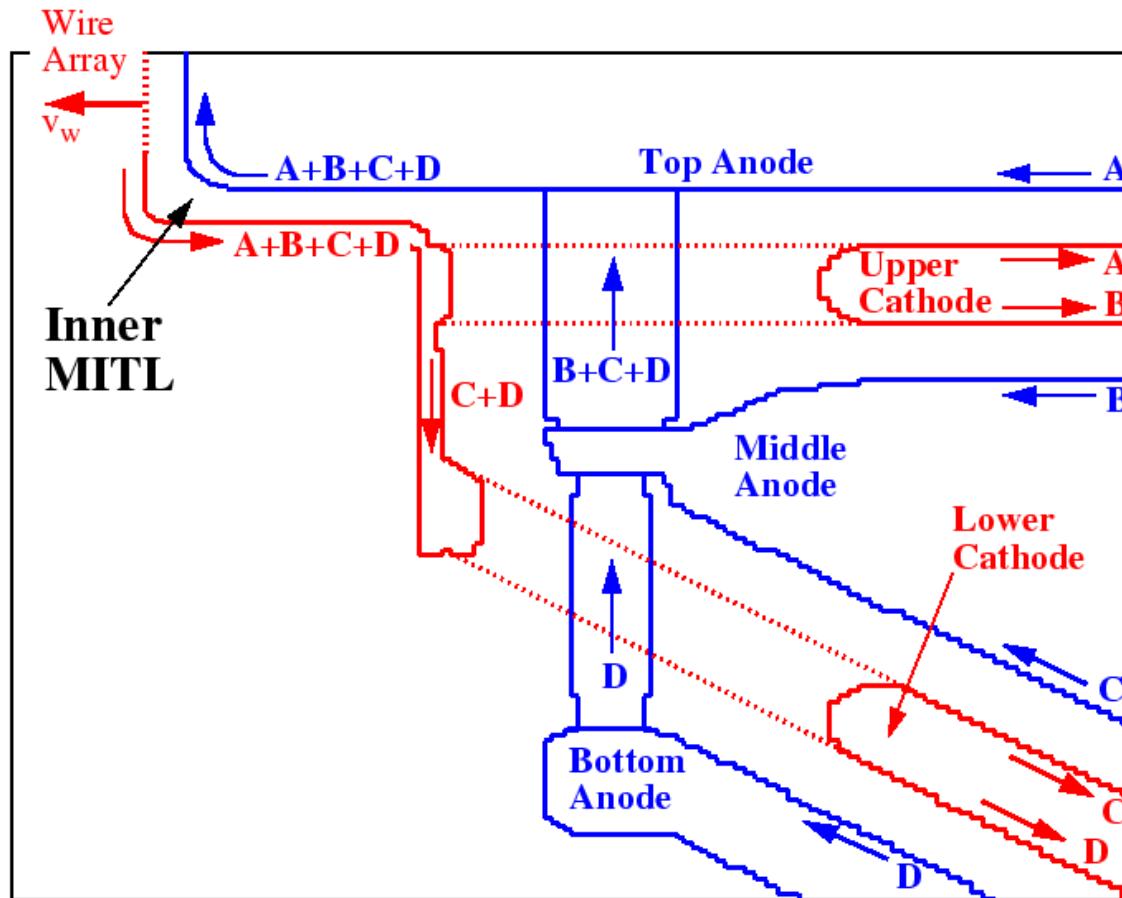
The vacuum section conducts power from the insulator stack to the load



- The four MITLs are coupled in parallel at the post-hole convolute
- Electron emission in the MITLs, out to the vacuum flares at $r \sim 1.5$ m
 - Electrons $E \times B$ drift radially inwards into the convolute

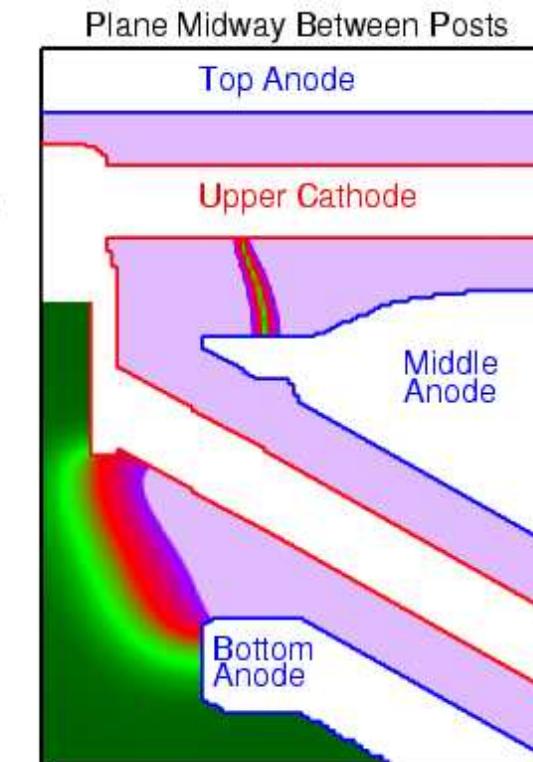
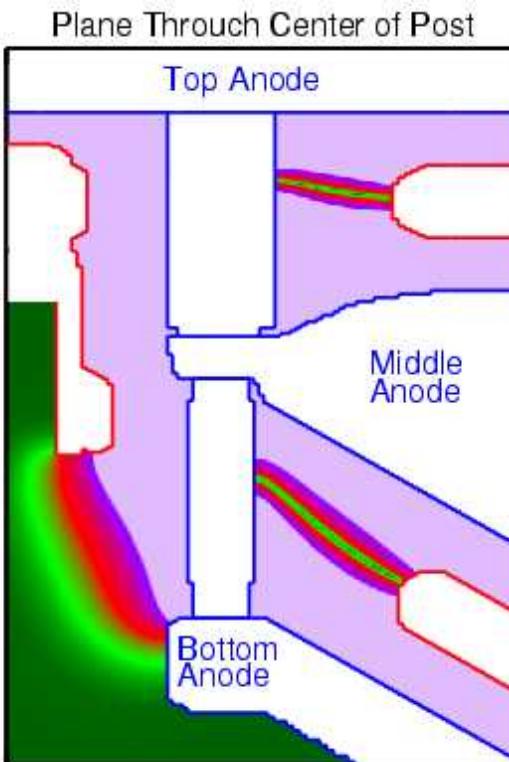


The convolute adds the currents from the four feed lines into the inner MITL



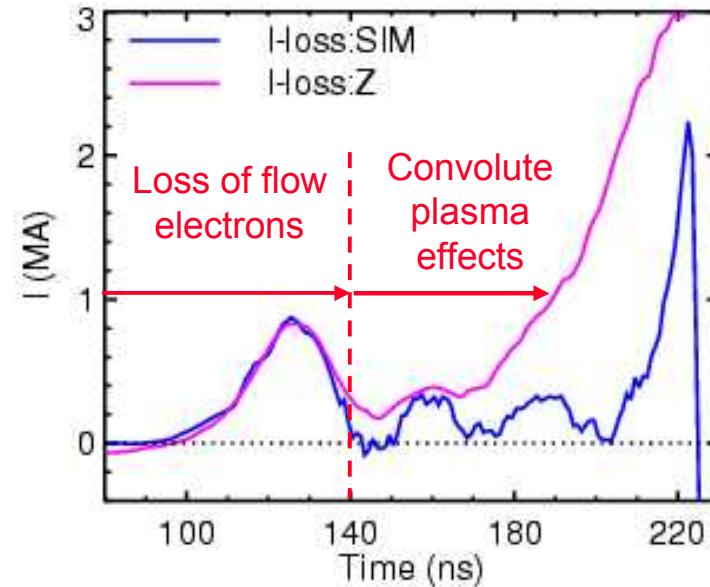
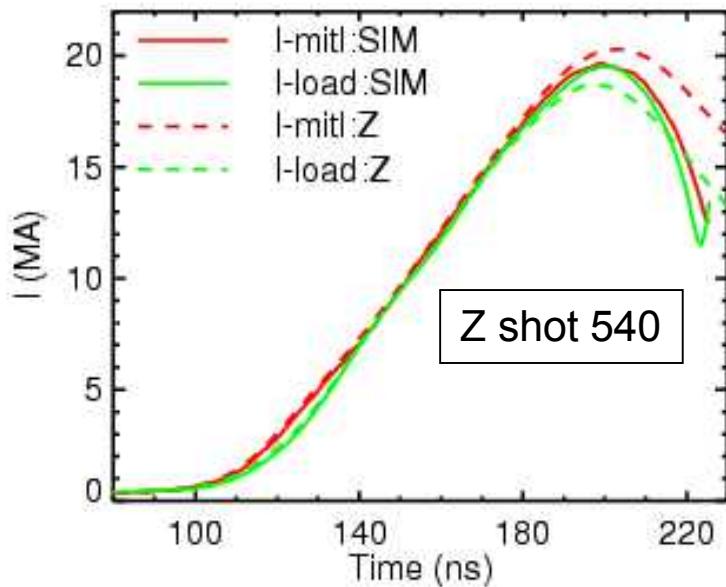
- Surface current distribution in the convolute is highly nonuniform.

Magnetic nulls in the convolute



- There will be electron losses at the nulls: “loss of magnetic insulation”
 - How much? How fast does the anode surface heat?

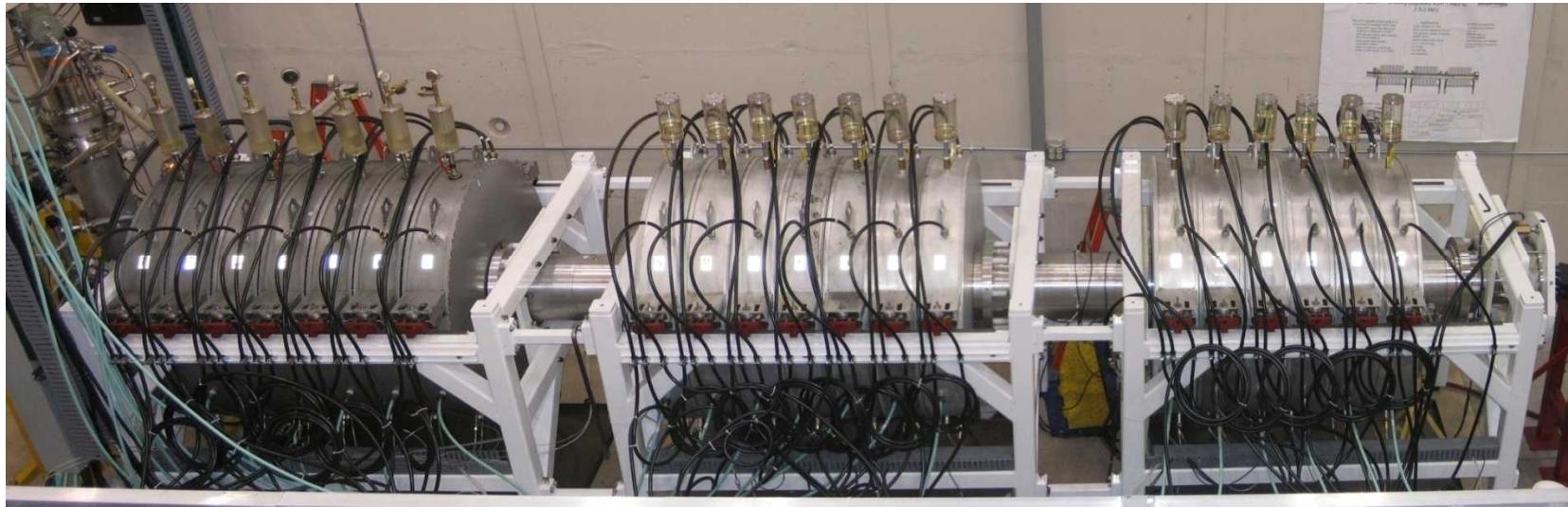
Quicksilver accurately computes the early current loss ... but not late in time



- Simulation has only vacuum flow electrons
- Early-time loss: MITL electrons lost to the anode in the convolute
- Additional late-time loss is due to electrode plasmas in the convolute
 - No detailed data: anode? cathode? both?

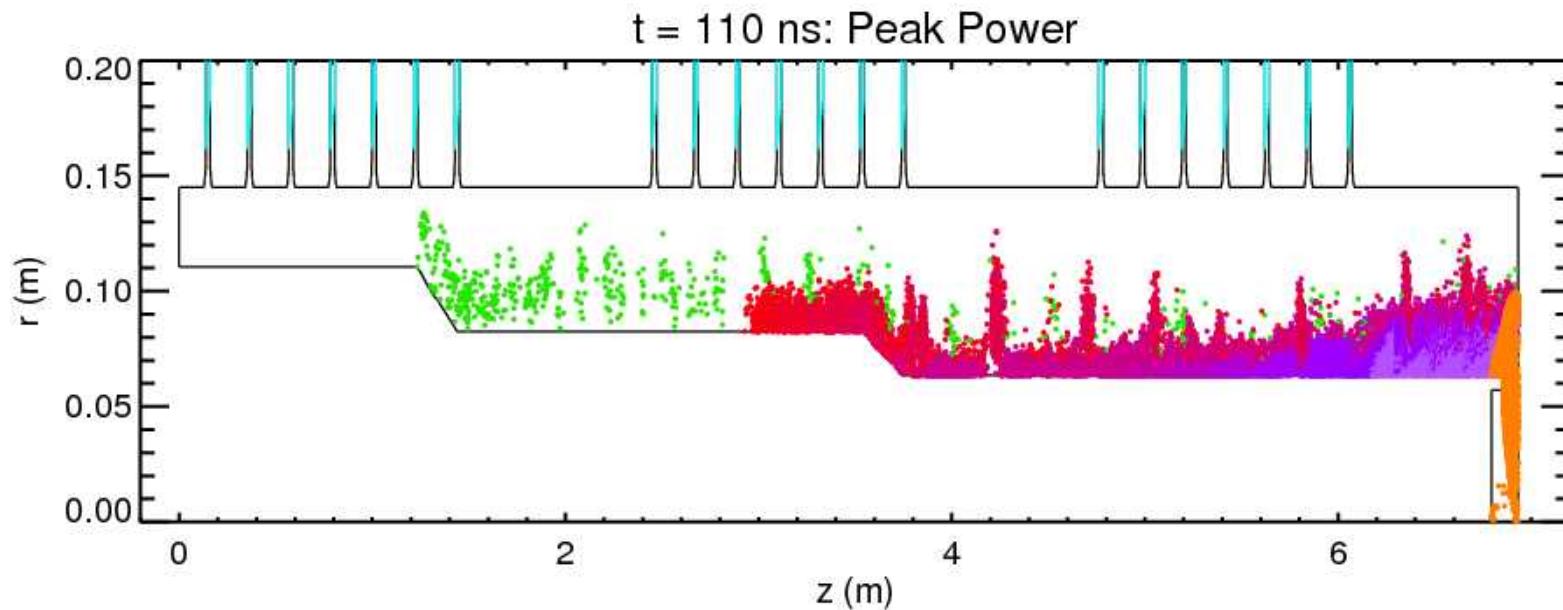


Ursa Minor is a testbed for adding many LTD cavities in series for voltage addition



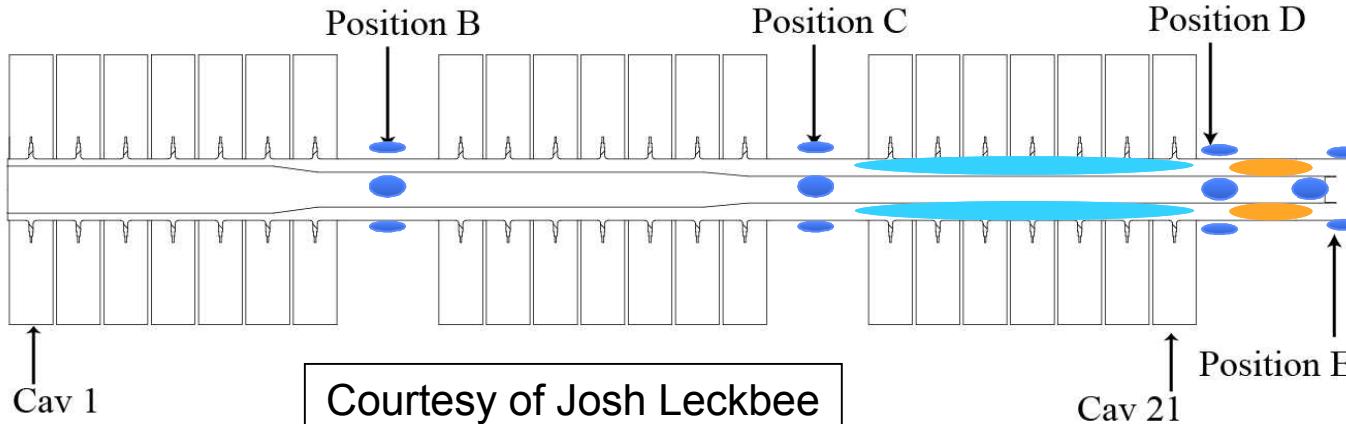
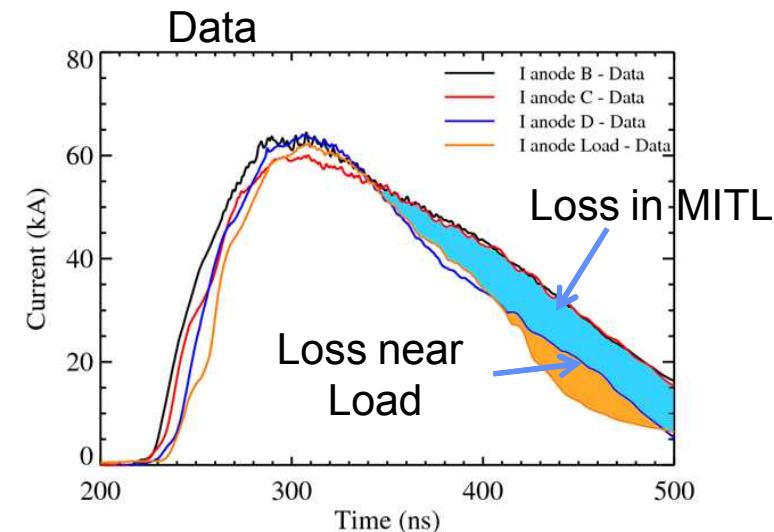
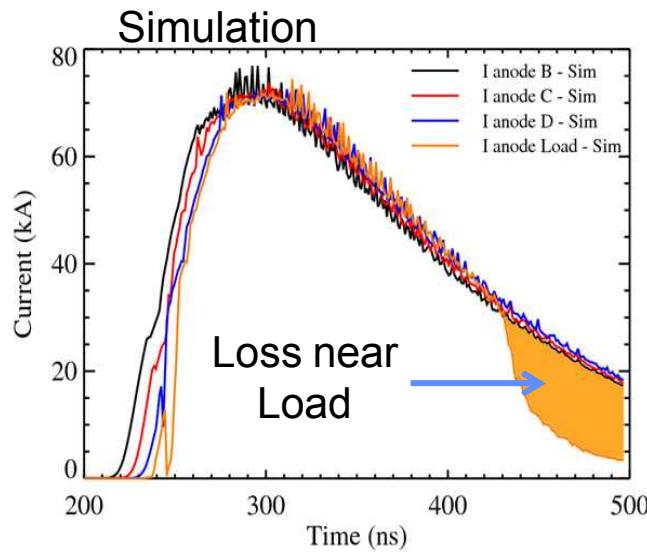
- Ultra-compact pulsed power: each cavity has 10 “bricks” in parallel, composed of two capacitors and a low-inductance switch in oil
 - No need for pulse-forming water section
- 21-cavity voltage adder, 7.5 m long, 1.5 m diameter
 - 180 kV/cavity → ~2.5 MV delivered to the load ... with unavoidable internal losses, but minimal electron flow loss

We have set up a 2-D r-z Quicksilver model of the full 21 cavity geometry



- 1 mm cell size across the 2.2 cm A-K gap of the 21 feeds
- Electrons “color-coded” by creation location
- Key questions:
 - What are electron losses to insulators (shown in cyan)?
 - What factors contribute to the loss?

Experiment has lower peak MITL current and higher loss current than simulations



Courtesy of Josh Leckbee

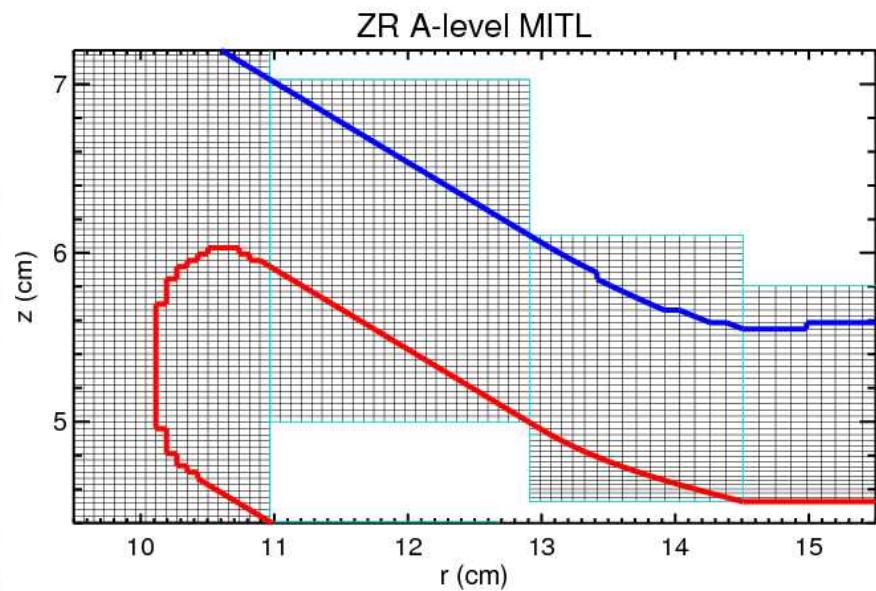
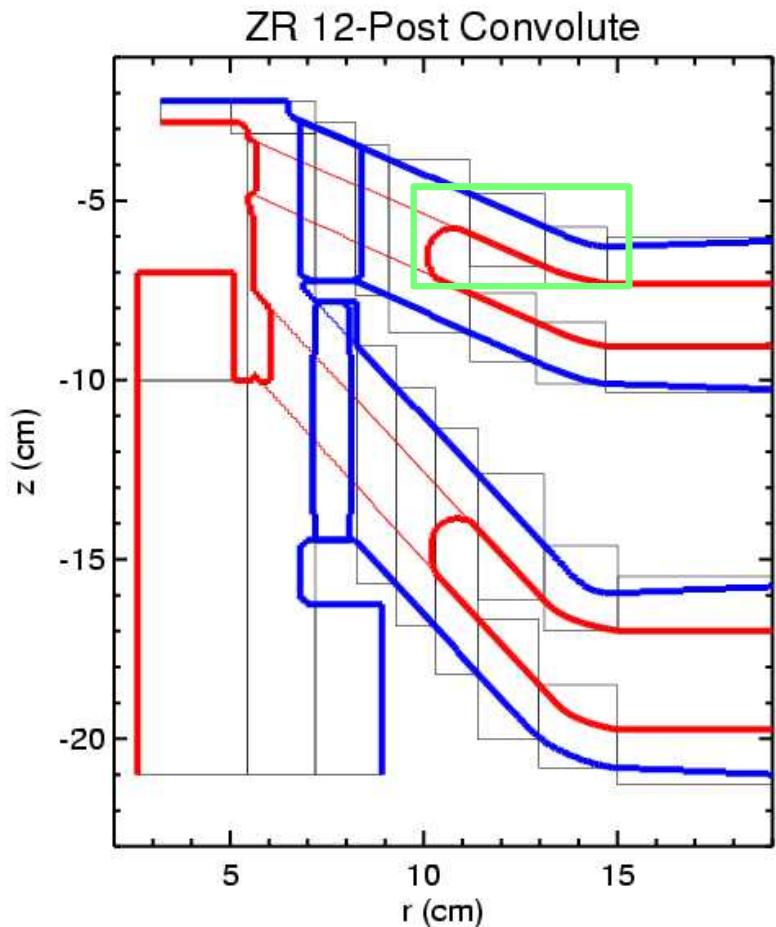


The Quicksilver code has been under development since the mid-1980's

- Originally ANSI-standard F77 for portability; now has F90, C, C++
- Mercury preprocessor
 - Process input text file(s) with symbolic variables, arrays, loops, if-tests, and include files
 - Error checking and computing array sizes
- 2D/3D multi-block geometry
 - Cartesian, cylindrical or spherical coordinates
 - 1D transmission lines for external circuits
- Parallelized with MPI
 - Static decomposition for field solver
 - Dynamic load-balancing for the particle handler
- IDL for all post-processing
 - Also use IDL extensively as a preprocessing tool

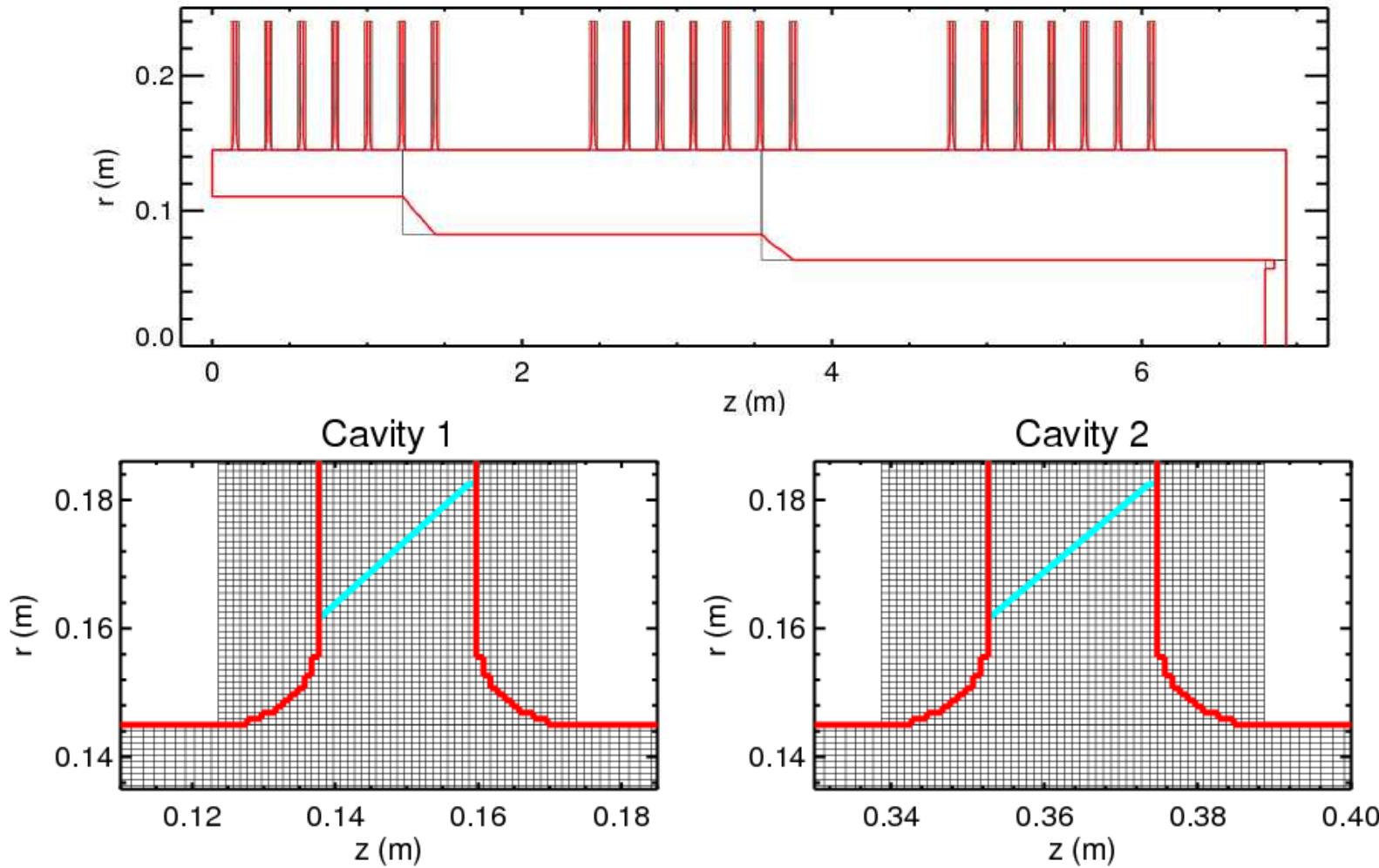


QS uses multi-block geometry to efficiently grid complex systems



This geometry is built with IDL scripts and post-hole convolute geometry is imported from an acis file

Multi-block geometry is an essential capability for the Ursa Minor simulations

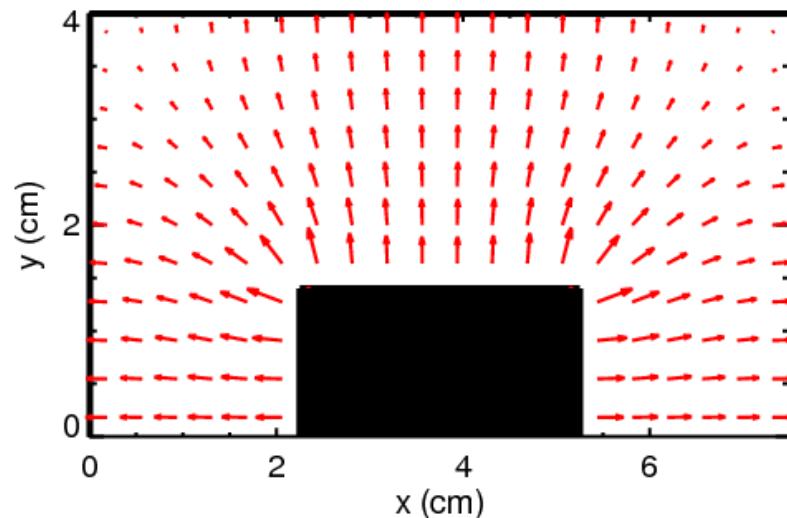


Augmenting PIC with 1D transmission lines is essential for modeling large systems

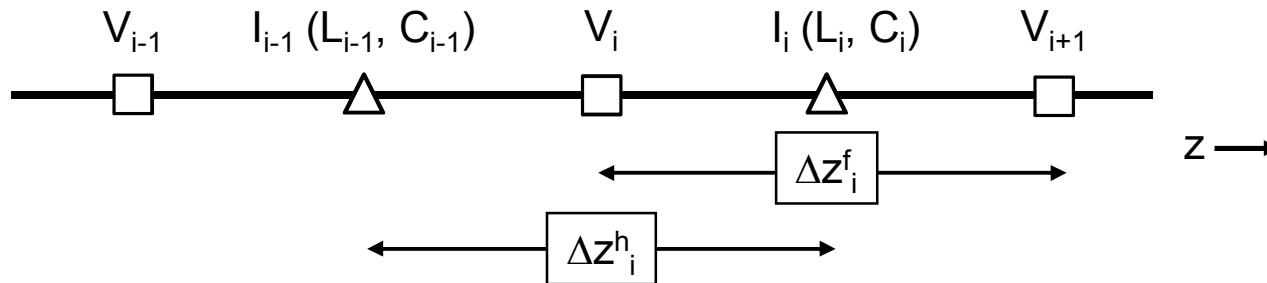
- TL's connect to a boundary plane of the 2D/3D domain at a “port” – two conductors (anode,cathode) with arbitrary “x-y” cross-section
- TEM mode E-field determined by:
$$\nabla_T^2 \phi = 0; \phi_{cat} = 0; \phi_{ano} = V_0; V_0 = 1 \text{ Volt}$$

$$\mathbf{E}_p = -\nabla \phi; \text{ “Poisson Field”}$$

$$\mathbf{E}_{TEM}(V) = (V/V_0) \mathbf{E}_p$$
- Per-unit-length capacitance
$$C = \frac{\epsilon}{V_0^2} \int E_p^2 dA$$
- Per-unit-length inductance
$$L = 1/(v_{ph}^2 C) ; v_{ph} = (\epsilon \mu)^{-1/2}$$



Voltage and current in the T-line are updated with the telegrapher's equations



- Telegrapher's equations for a lossless line:

$$\frac{\partial I}{\partial t} = -\frac{1}{L} \frac{\partial V}{\partial z} ; \quad \frac{\partial V}{\partial t} = -\frac{1}{C} \frac{\partial I}{\partial z}$$

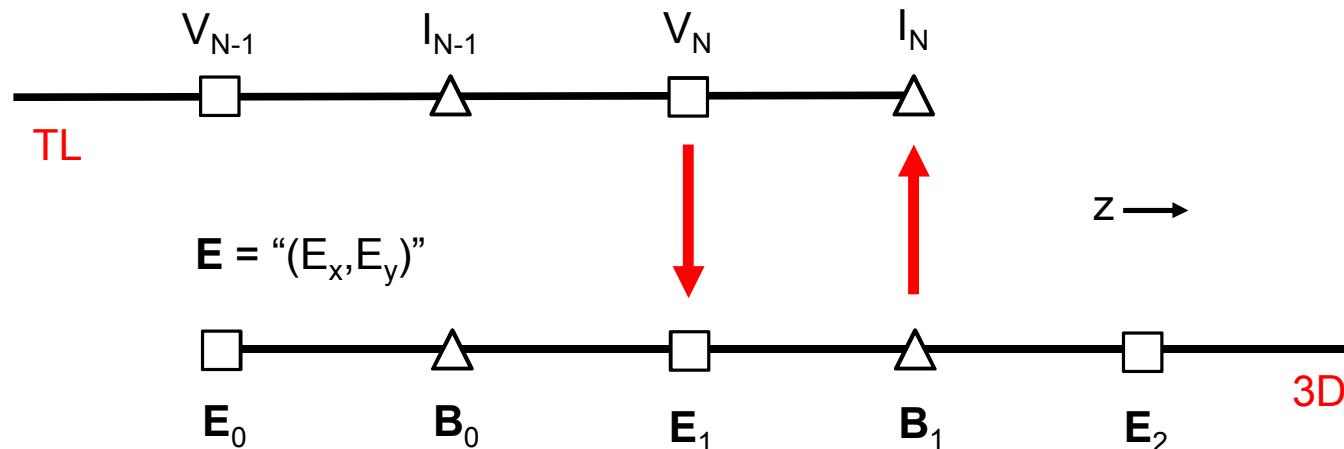
- Leapfrog finite difference, $\{V^{n-1}, I^{n-1/2}\} \rightarrow \{V^n, I^{n+1/2}\}$:

$$V_i^n = V_i^{n-1} - \gamma_i (I_i^{n-1/2} - I_{i-1}^{n-1/2}) ; \quad \gamma_i = 2\Delta t / (\Delta z_i^h [C_{i-1} + C_i])$$

$$I_i^{n+1/2} = I_i^{n-1/2} - \lambda_i (V_{i+1}^n - V_i^n) ; \quad \lambda_i = \Delta t / (\Delta z_i^f L_i)$$

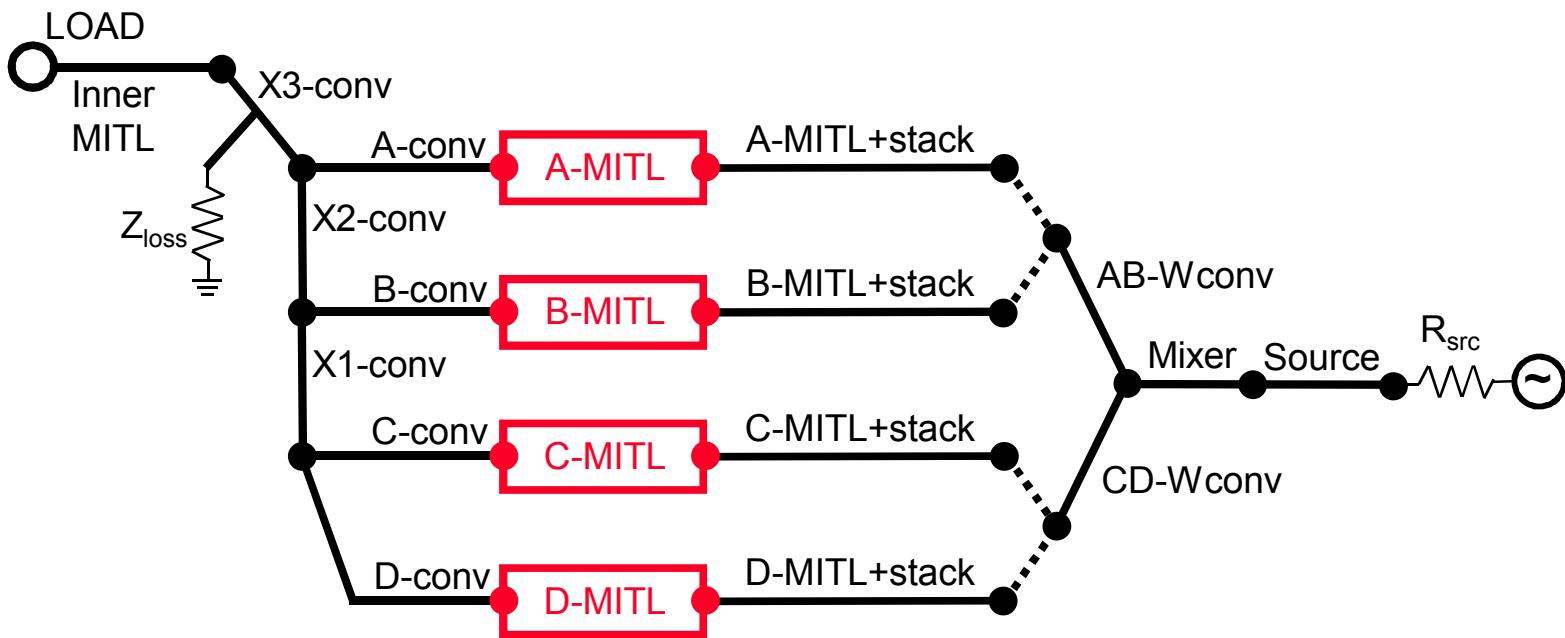
- Blazingly fast compared to 2D/3D field solver update
- Easy to add BC's at the end of a TL: generators, Z-pinch loads, etc.

The transmission line couples naturally to the TEM mode in the 2D/3D system



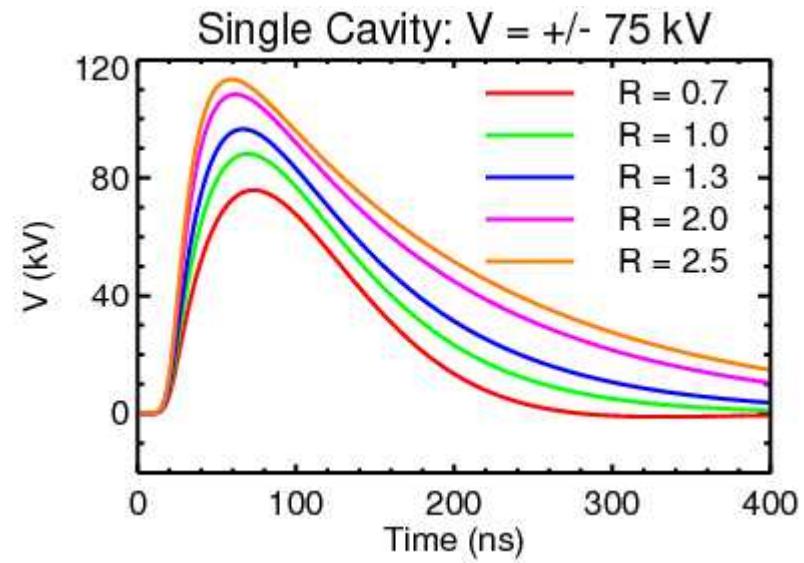
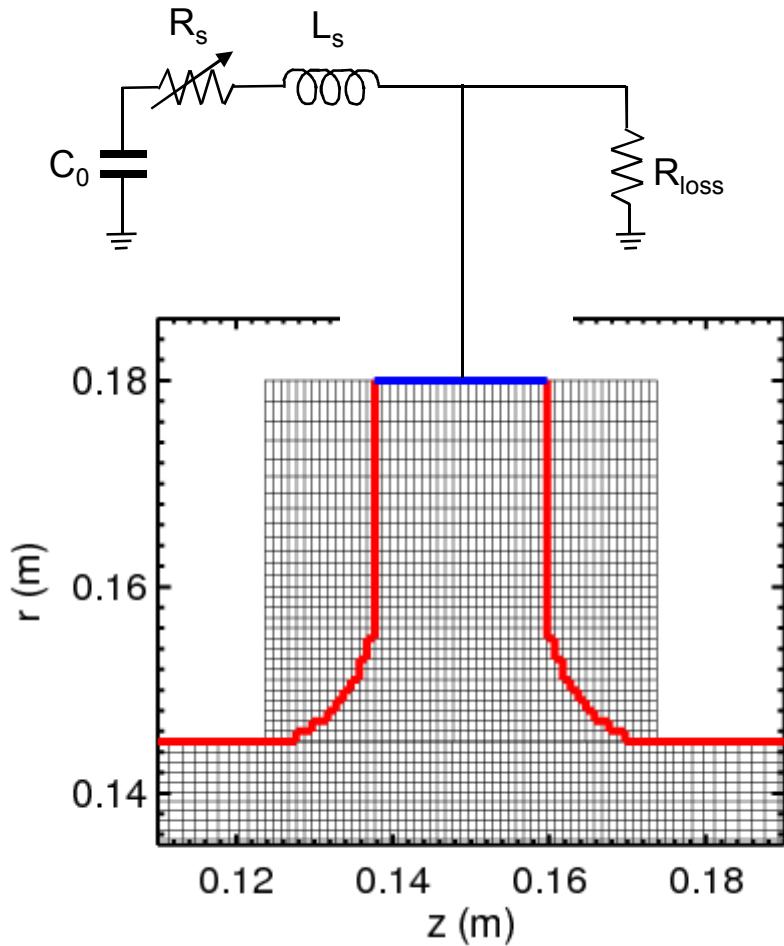
- TL voltage V_N^n determines E_T^n on boundary plane of the 3D system
 - $E_1^n = E_{TEM}(V_N^n) + E_{non-TEM}^n$
 - Hard part is specifying $E_{non-TEM}^n$
- 3D magnetic field $B_1^{n+1/2}$ determines boundary current $I_N^{n+1/2}$
 - $I_N = \frac{1}{\mu V_0} \int e_z \times E_p \cdot B_1 dA$ (filters out non-TEM part of B_1)
 - This defines V_N^{n+1} on next timestep

High-resolution simulations of the Z MITLs use four 2D PIC regions coupled with TL's

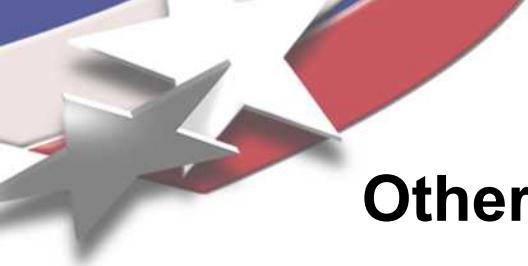


- 2-D PIC region for each MITL, spherical coords., $0.1 < r < 0.6$ m, shown in red
- 1-D transmission lines in black
 - Outer lines extend out to $r \sim 3$ m
 - Inner convolute line parameters obtained from 3-D simulations
 - Optional, time-varying convolute Z_{loss} element
 - Either a Z-pinch or ICE load

Ursa Minor simulations use an external RLC circuit with core loss for each cavity



- $C_0 = 100$ nF, $L_s = 25$ nH, $R_{loss} = 4$ Ω
- R_s : $10^5 \rightarrow 0.2$ Ω with 2 ns exponential decay
- Single cavity into resistive load R agrees to $\sim 5\%$ with more detailed circuit model



Other features in Quicksilver

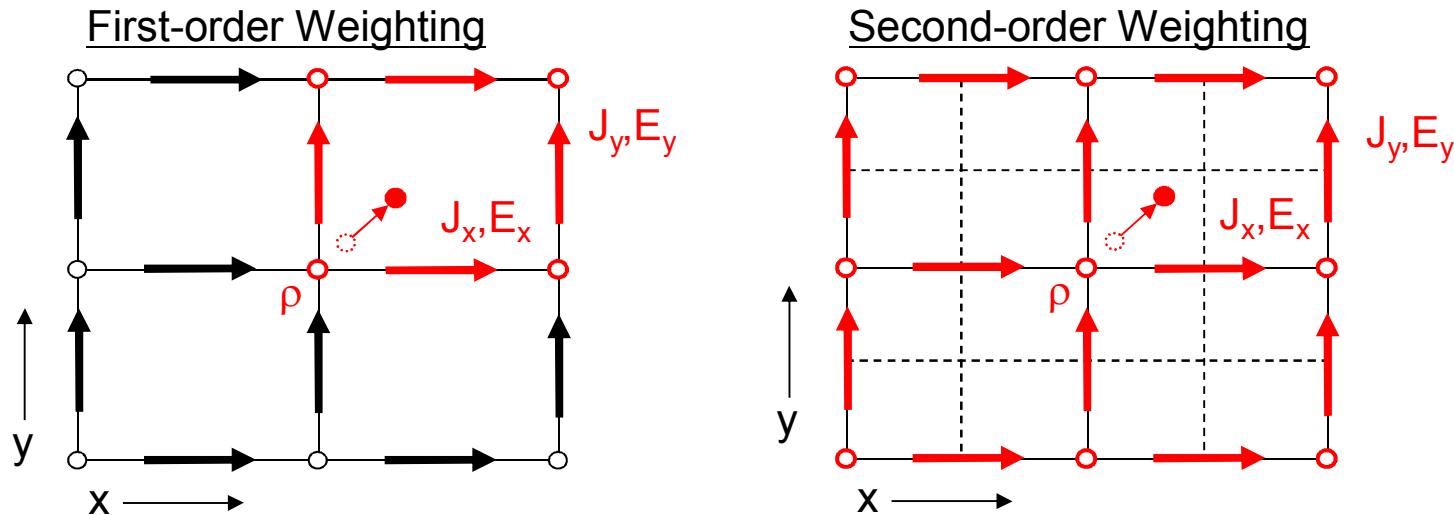
- Friedman's explicit EM field solver with adjustable damping, $0 \leq \theta \leq 1$
 - Also use this scheme to add damping to the explicit particle pusher
- Exact, second-order \mathbf{J} and ρ accumulation, handling boundaries
 - At conductors, morph 2nd \rightarrow 1st order within half a cell of boundary
- Energy-conserving particle pusher: interpolate \mathbf{E} directly from the Yee mesh using exactly the same weighting that \mathbf{J} is laid down on the mesh
 - Enables stability at $\Delta x / \lambda_D \gg 1$, $\omega_p \Delta t \sim 1$ (little numerical heating)
 - Accuracy? *Caveat emptor!*
- Floating-point precision issues:
 - Particle push in local coord. system relative to lower corner of cell
 - Build executables with 4- or 8-byte reals
 - ~ 1.6 slowdown for r8 – memory access



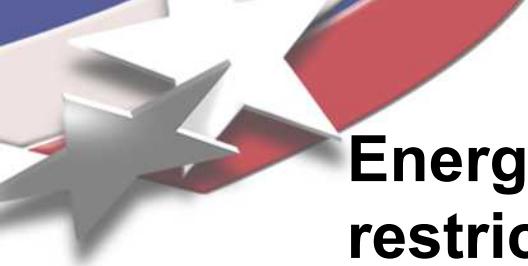
Other features in Quicksilver, contd.

- Sort particles by cell and species:
 - Required for Coulomb collisions and particle merger
 - Improved performance by limiting cache-thrashing
- Particle-particle Coulomb collisions for cold, dense plasmas
- Electron-surface interactions
 - Heating
 - Reflection and secondary emission
- Kinetic gas-chemistry model
 - Gas breakdown plasma modeled with electron-ion pairs
 - Robust particle merger to deal with exponentially-growing particle count

Energy-conserving PIC: Specific scheme for E-field interpolation to the particles



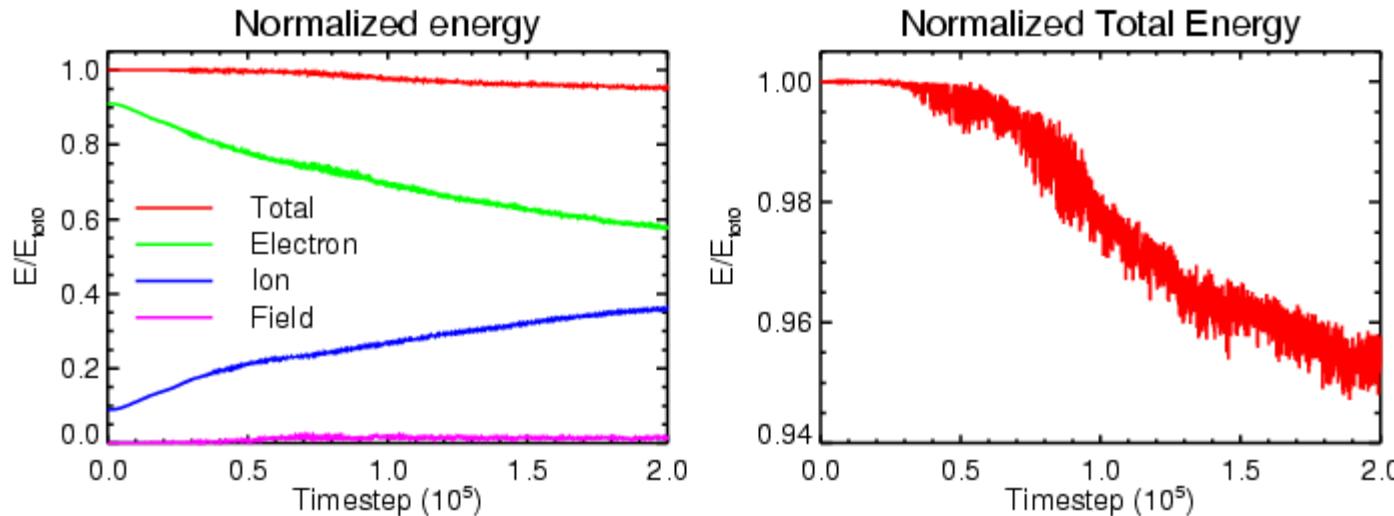
- For EM codes: interpolate **E** directly from the Yee mesh to the particles, using exactly the same weighting scheme as **J** from the particles to the grid
- First-order weighting: Interpolation of E_x is nearest-grid-point in x, linear in y
 - Although it appears crude, we have found it very useful
- Second-order weighting: Interpolation of E_x is linear in x, quadratic in y
 - Handle boundaries: T. Pointon, Comput. Phys. Commun. **179** (2008) 535.



Energy-conserving PIC removes the restriction of having to run at $\Delta x/\lambda_D < \sim 1$

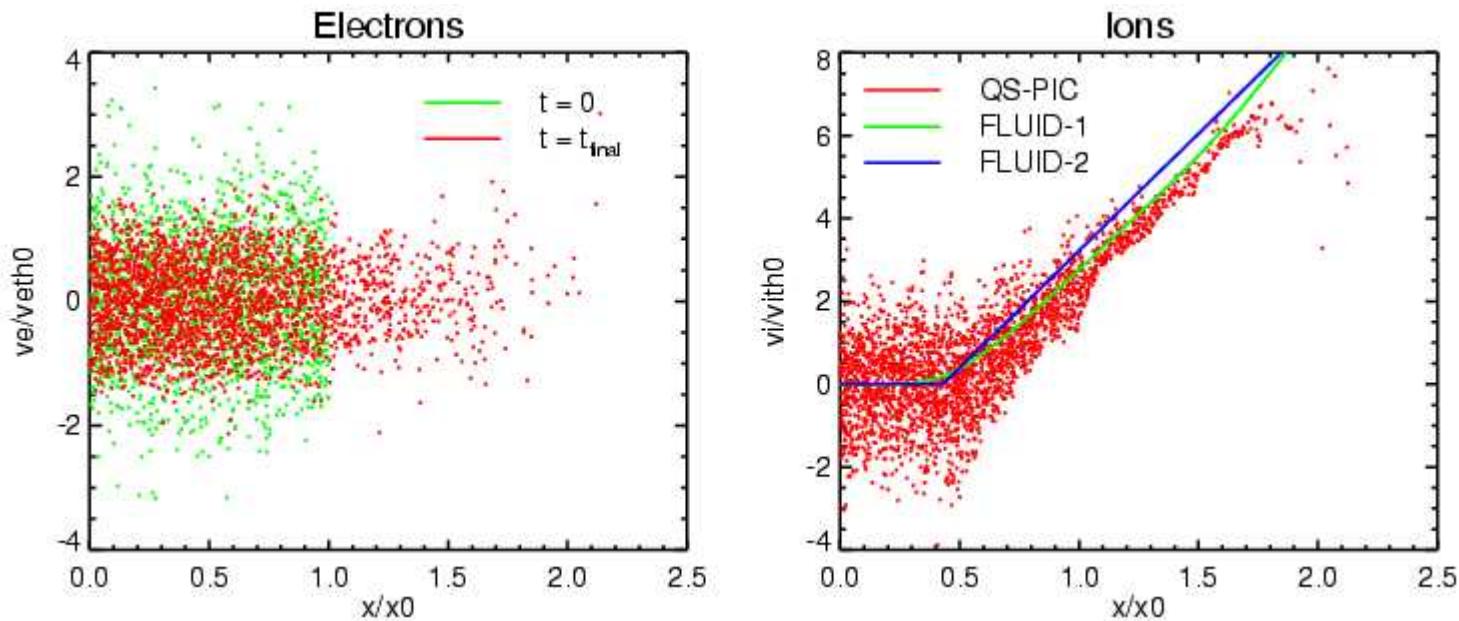
- Explicit, energy-conserving PIC runs are stable for $\Delta x/\lambda_D \gg 1$
 - Minimal numerical heating
- Stability does not imply accuracy!
 - 1st-order algorithm cannot handle extremely large E-field gradients
 - Neither algorithm can accurately model a plasma at temperature T in presence of a large external E-field satisfying $qE\Delta x \gg T$
- For well-chosen parameters, neither of these factors is an issue, and we can run at $\Delta x/\lambda_D \gg 1$ with reasonable accuracy, limited only by $\omega_p\Delta t < \sim 1$
 - For $\Delta t = 1$ ps, $\omega_p\Delta t = 1 \rightarrow n < 3.1 \times 10^{14}$ cm⁻³ ($\sim 1\%$ ionization at 1 Torr)
- However ... ‘On two occasions I have been asked “Pray Mr. Babbage, if you put into the machine wrong figures, will the right answer come out?” I am not able rightly to apprehend the kind of confusion of ideas that could provoke such a question’, Charles Babbage, 1864

Canonical test for high-density plasma: 1-D expansion of plasma slab into vacuum



- Use $\Delta x = 0.5$ mm, $\Delta t = 0.5$ ps; characteristic of Z-convolute runs
- Plasma density near explicit limit: $n = 10^{15} \text{ cm}^{-3}$; $T_e = 1 \text{ eV}$, $T_i = 0.1 \text{ eV}$
 - $\omega_{pe}\Delta t = 0.89$, $\Delta x/\lambda_{De} = 2130$
- ~5% numerical cooling after 200,000 timesteps (~4% with no damping in the field solver)

Particle phasespace plots show the plasma expansion is modeled quite well

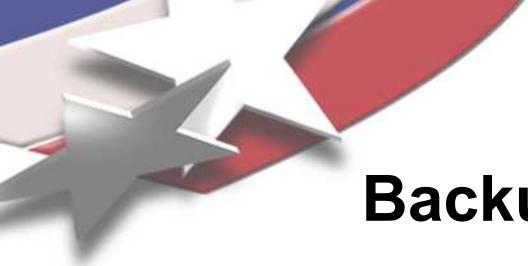


- Transfer thermal electron energy to directed ion energy at the front
- Results agree well with 1-D multi-fluid (electron/ion) simulations
 - 1st multi-fluid simulation uses same Δx as the QS run
 - 2nd m-fluid simulation uses Δx 2000 times smaller ($\Delta x/\lambda_{De} \sim 1$)



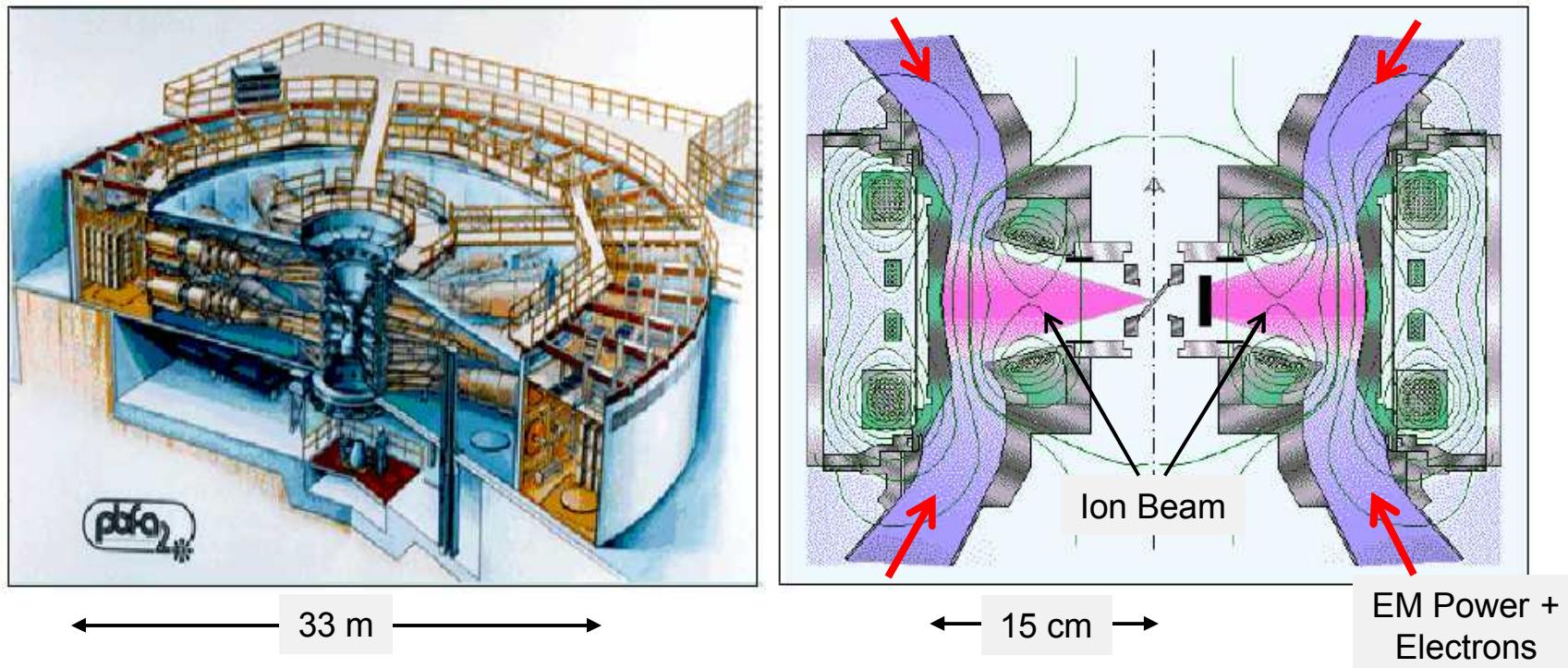
Summary

- Quicksilver is a mature 2D/3D EM-PIC code with a set of features driven by its application space
 - Pulsed power systems
 - EM response to gas-filled cavities with electron emission from walls
- Quicksilver features
 - Excellent capabilities for handling complex geometry
 - Energy-conserving PIC
 - Stability at $\Delta x/\lambda_D \gg 1$, $\omega_p \Delta t \sim 1$ (minimal numerical heating)
- Post-processing with IDL



Backup Slides

The applied-B ion diode on PBFA II produced 10 MV, 1 MA, 40 ns Li⁺ ion beams



- 2D PIC simulations: coupling 4 MITLs in series to achieve 10 MV
- 3D PIC simulations: EM instabilities → beam divergence
- 1D DYNAID simulations: anode plasma expansion → gap closure



Quicksilver and IDL use PFF (portable file format) to handle floating-point data

- PFF was developed by David Seidel in late 1980's for QS
 - Save disk space by encoding floating-point arrays into 16-bit integer arrays using an offset and multiplier
 - Full precision is now available if necessary
 - Pre-defined “dataset types” for data relevant to PIC
- User interface at the dataset level is machine-independent
 - UF1: uniformly spaced 1-D data (time histories)
 - NF3, NV3: Scalars and 3-D vectors on 3-D multi-block domain (field snapshots)
 - VTX: List of vertices on an m-D space with n attributes per vertex (particle snapshots, surface field snapshots)
 - NGD: m-D vectors on a single n-D block
- Public domain: <http://sourceforge.net/projects/hermes-util>



Quicksilver uses IDL for post-processing ... and for more and more pre-processing

- “PFIDL” layer of IDL routines (Paul Mix)
 - PFF dataset read/write
 - Waveform (1D time history) processing: plot, arithmetic, FFT, etc.
 - “Structure” processing: 2D/3D field snapshots, particle snapshots
- My IDL routines for post-processing
 - Time history modification: filtering, etc.
 - Field, lineout, and particle snapshot animation
- My IDL routines for QS pre-processing
 - Generation of simulation geometry
 - Processing rad-transport simulation output for EM simulations: electron reflection, secondary emission and surface heating
 - Processing gas-chemistry cross-section data

Quicksilver/Mercury/IDL flowchart

