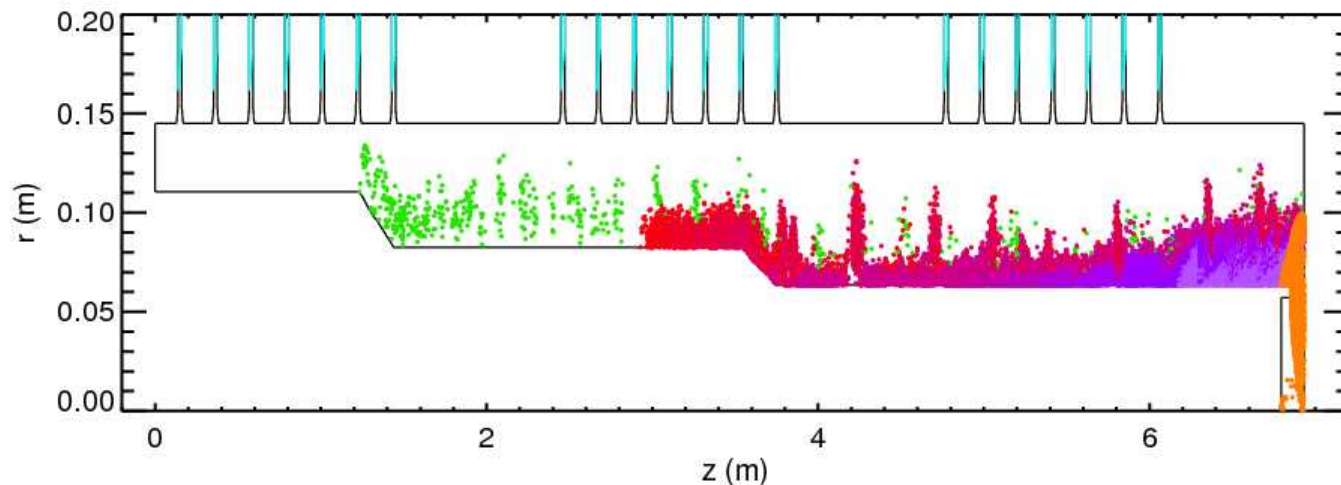


EM-PIC Simulation of Pulsed Power Systems at Sandia National Laboratories

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Tech-X Corporation

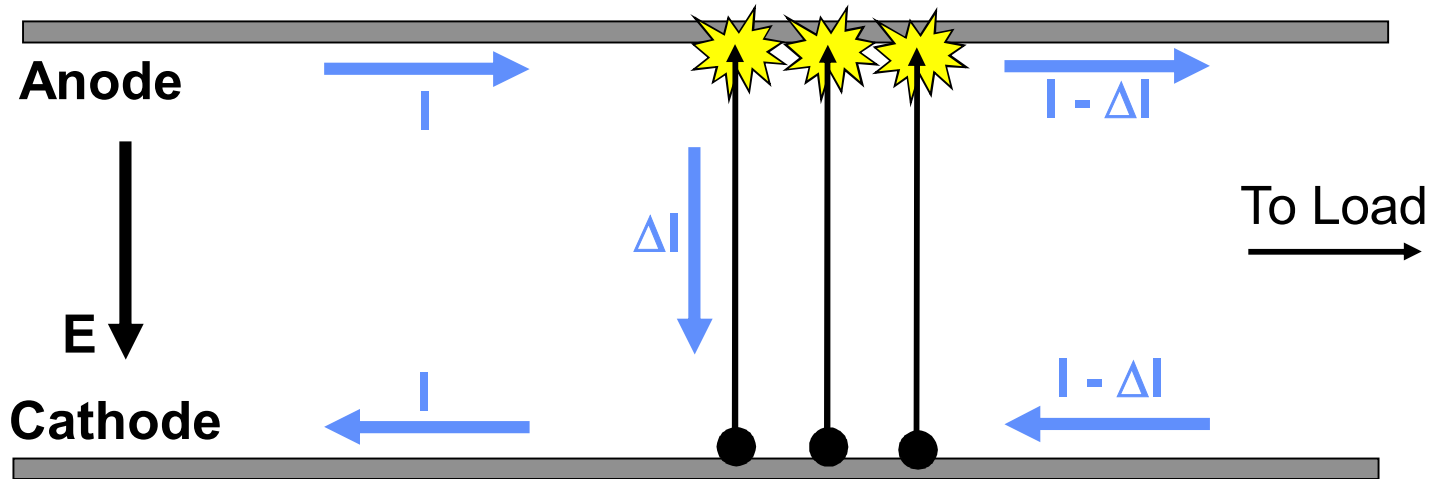
Boulder CO, June 29, 2012



Outline

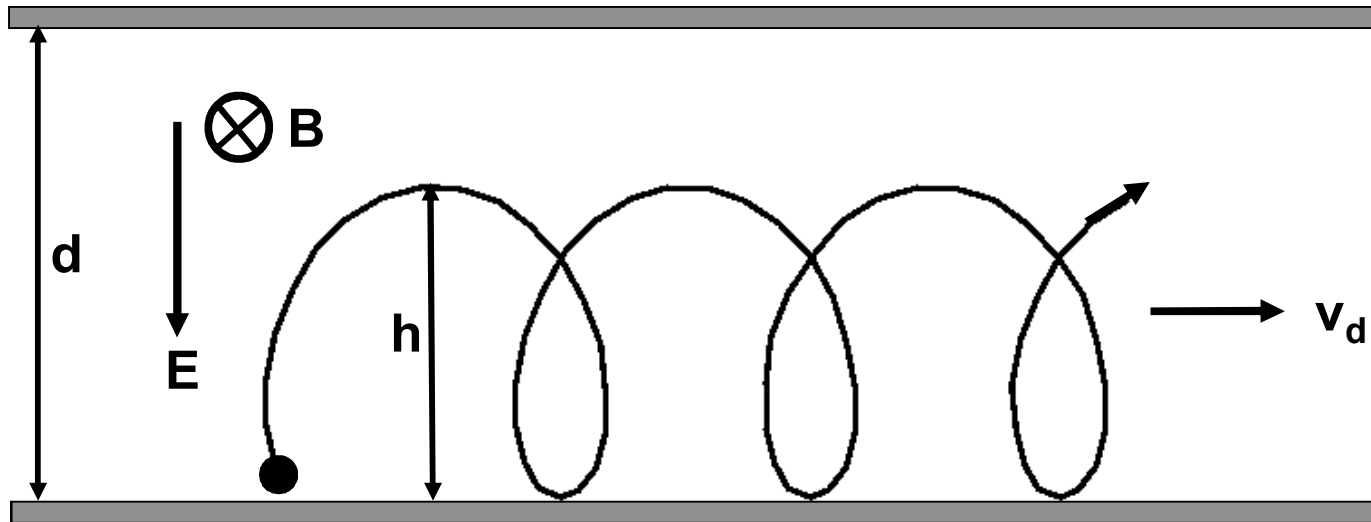
- Applications driving analysis and code development:
 - Magnetically insulated transmission lines (MITLs) (ongoing)
 - Light ion beams for fusion: applied-B ion diodes (1987 – 1997)
 - Power flow in the vacuum section of the “Z” accelerator (1997 –)
 - The 21-cavity “Ursa Minor” linear transformer driver (LTD) for radiography (2009 –)
 - Electron beam transport in gas (“gas chemistry”)
- The Quicksilver 2D/3D EM-PIC code
 - Geometry, 1D transmission lines
 - Other features - physics models
- Other PIC work (not discussed here)
 - Implicit ES hybrid fluid-PIC codes: DYNAID (1D), MUFPHI (2D)
 - Unstructured EM-PIC: tetrahedral meshes

Cathode electron emission is unavoidable for a multi-MV EM pulse in a transmission line



- ***Without special treatment***, a cathode surface emits electrons when the normal electric field E_n exceeds $\sim 50 - 200$ kV/cm
- Electrons flowing to the anode is undesirable
 - Reduces current delivered to the load
 - Heats the anode -- can create an anode plasma

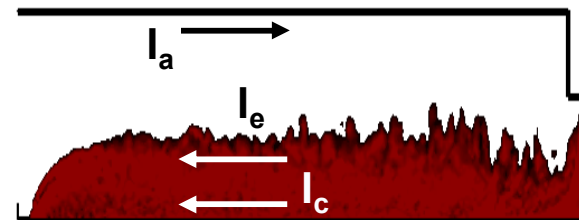
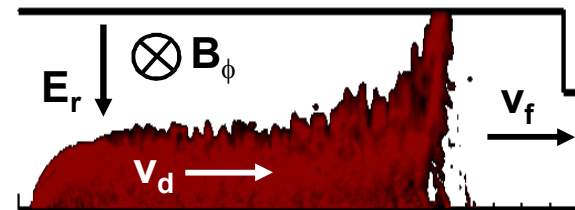
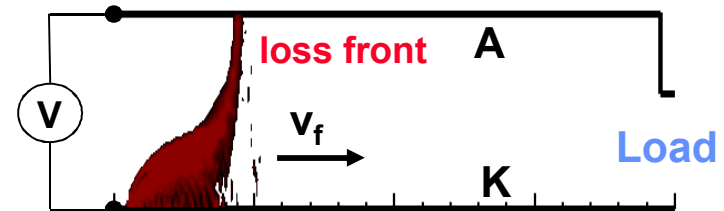
A transverse magnetic field can inhibit electrons from crossing the A-K gap



- For $cB > E$, the electrons “drift” with velocity $\mathbf{v}_d = (\mathbf{E} \times \mathbf{B})/B^2$
- For uniform \mathbf{E} and \mathbf{B} fields:
$$h = \frac{\frac{2m}{e} E}{B^2 - E^2 / c^2}$$
 - For sufficiently large B , $h < d$

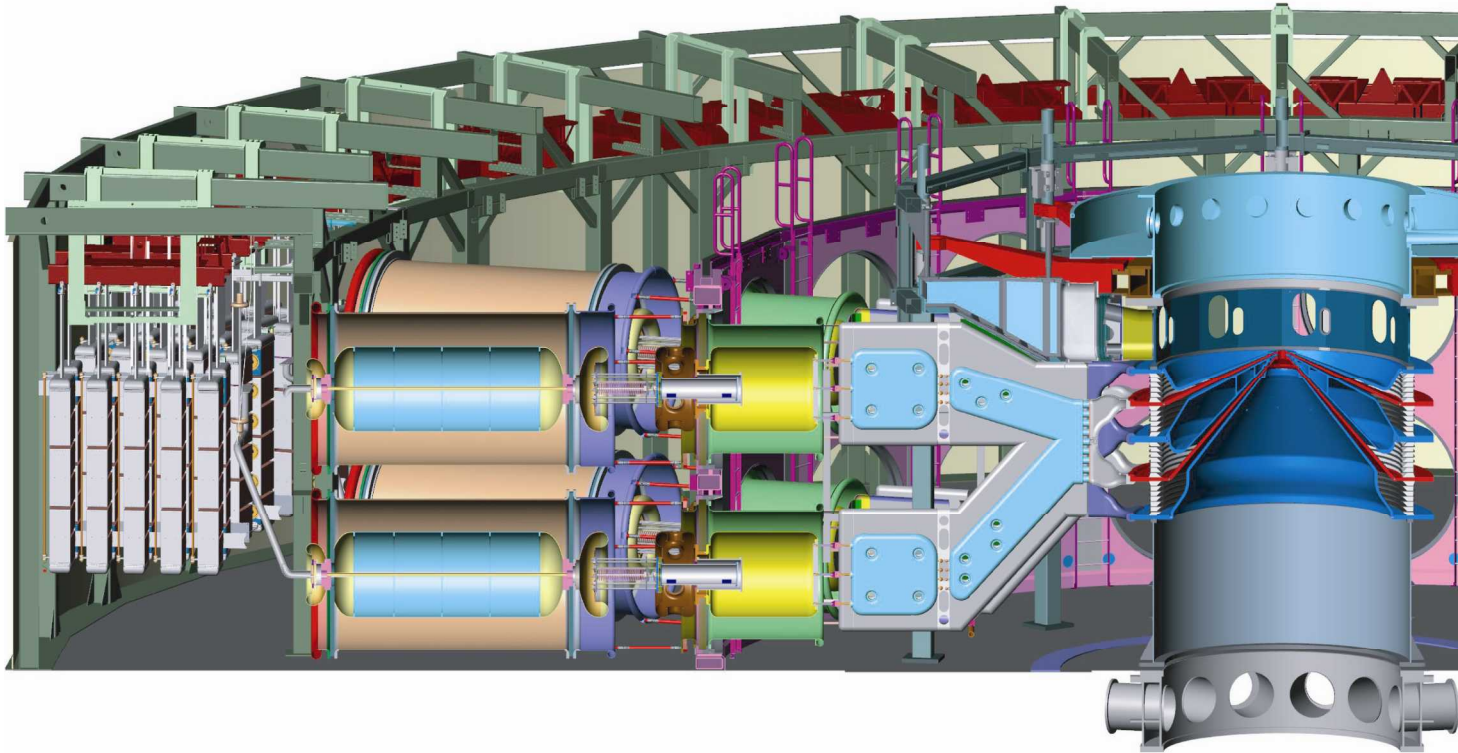
Efficient power transfer is possible with “magnetically insulated transmission lines”

- Schematic illustration of a coaxial “MITL” driving a load
- Behind the “loss front”, B-field of the current flowing in the conductors insulates electrons: they “**E** \times **B** drift” towards the load
- In equilibrium, current flow in the “electron sheath” is related to the anode and cathode conduction currents by: $I_e = I_a - I_c$
- Technical issues for real systems
 - Current loss if load impedance is too high
 - Adding lines in serial or parallel



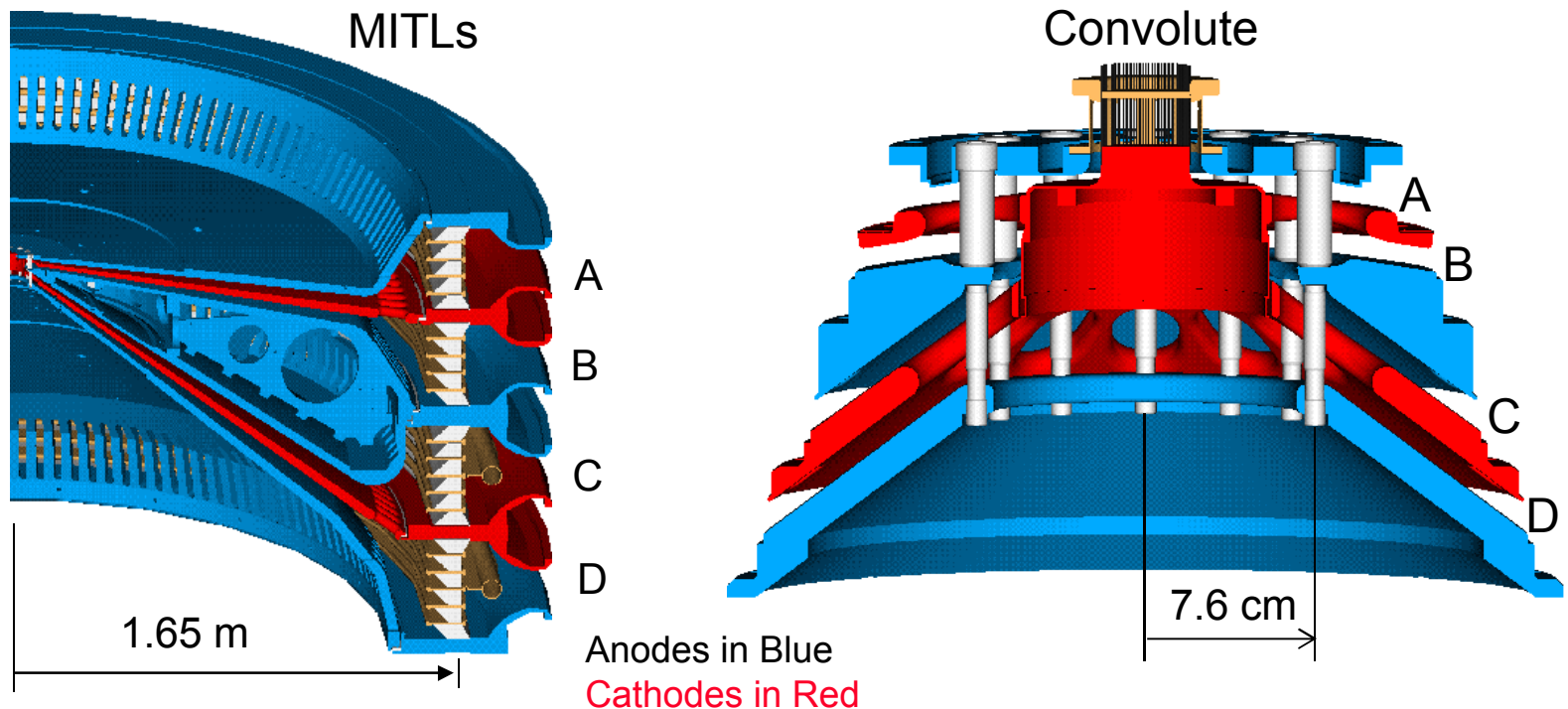
Courtesy of Steve Rosenthal

Sandia's Z accelerator is the world's most powerful high-current driver



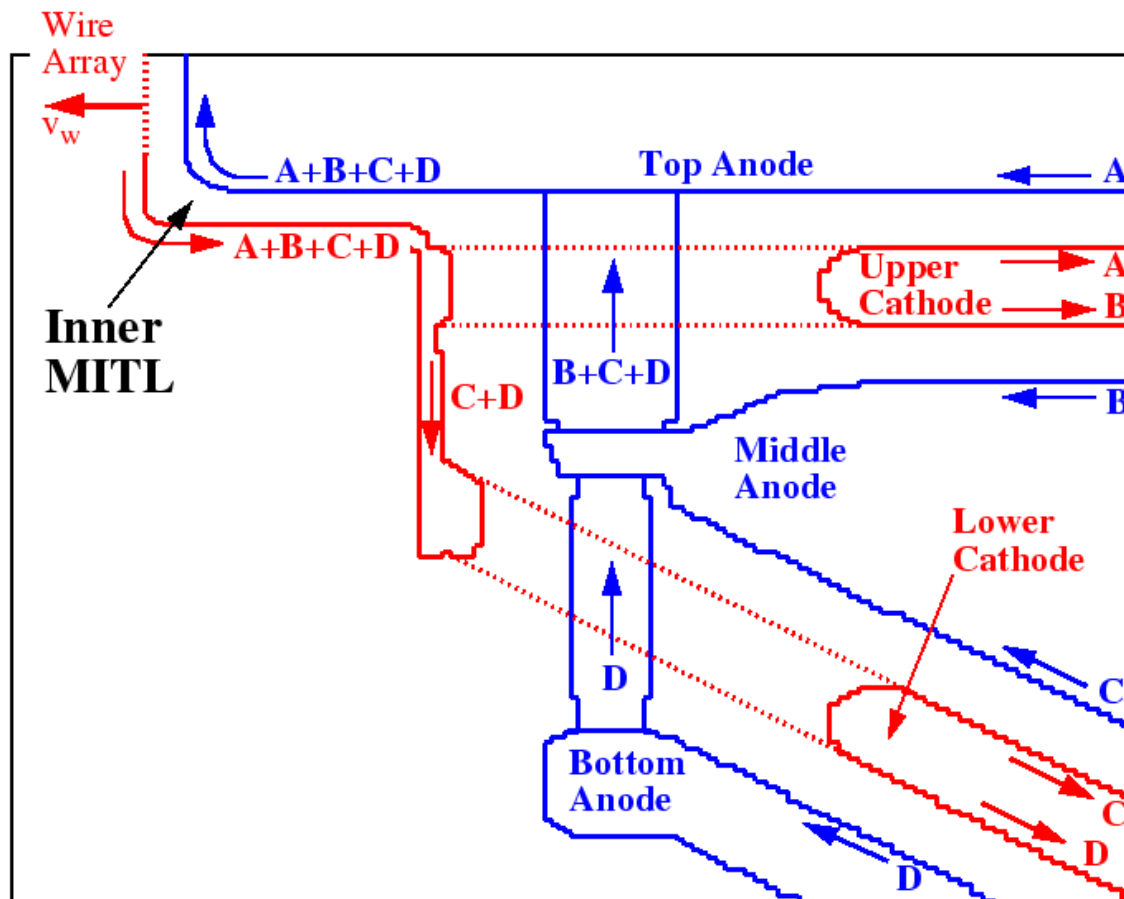
- 33 m in diameter, 6 m high, 36 pulse-forming lines
- Drives variety of loads, e.g. 26 MA, 100 ns pulse to a Z-pinch

The vacuum section conducts power from the insulator stack to the load



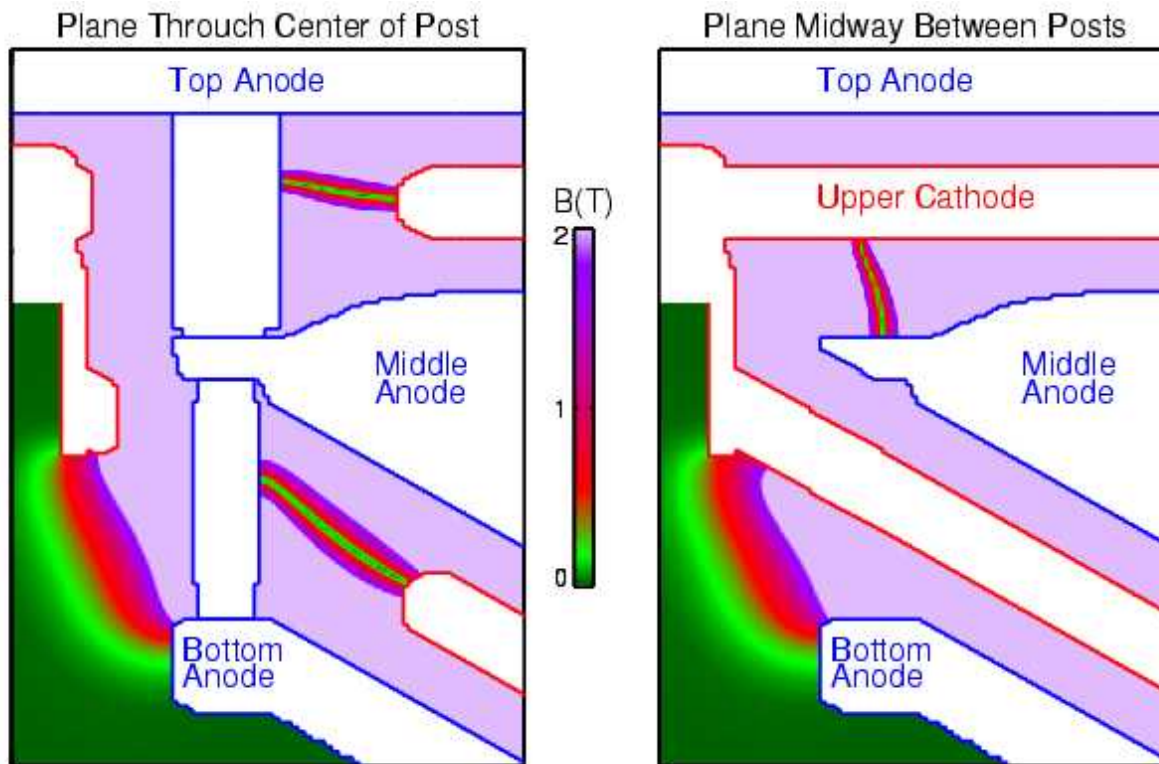
- The four MITLs are coupled in parallel at the post-hole convolute
- Electron emission in the MITLs, out to the vacuum flares at $r \sim 1.5$ m
 - Electrons $\mathbf{E} \times \mathbf{B}$ drift radially inwards into the convolute

The convolute adds the currents from the four feed lines into the inner MITL



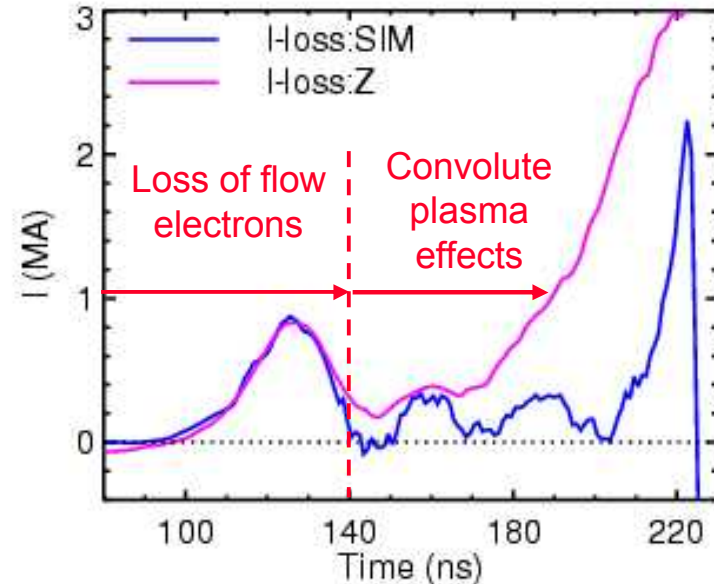
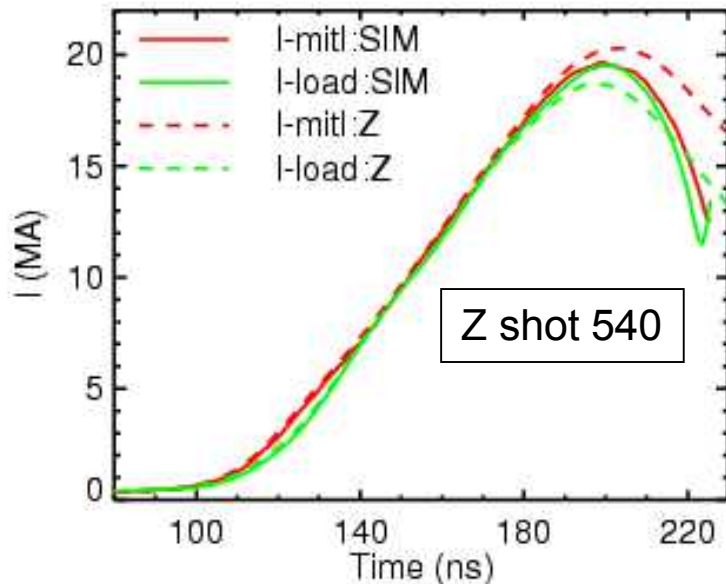
- Surface current distribution in the convolute is highly nonuniform.

Magnetic nulls in the convolute



- There will be electron losses at the nulls: “loss of magnetic insulation”
 - How much? How fast does the anode surface heat?

Quicksilver accurately computes the early current loss ... but not late in time



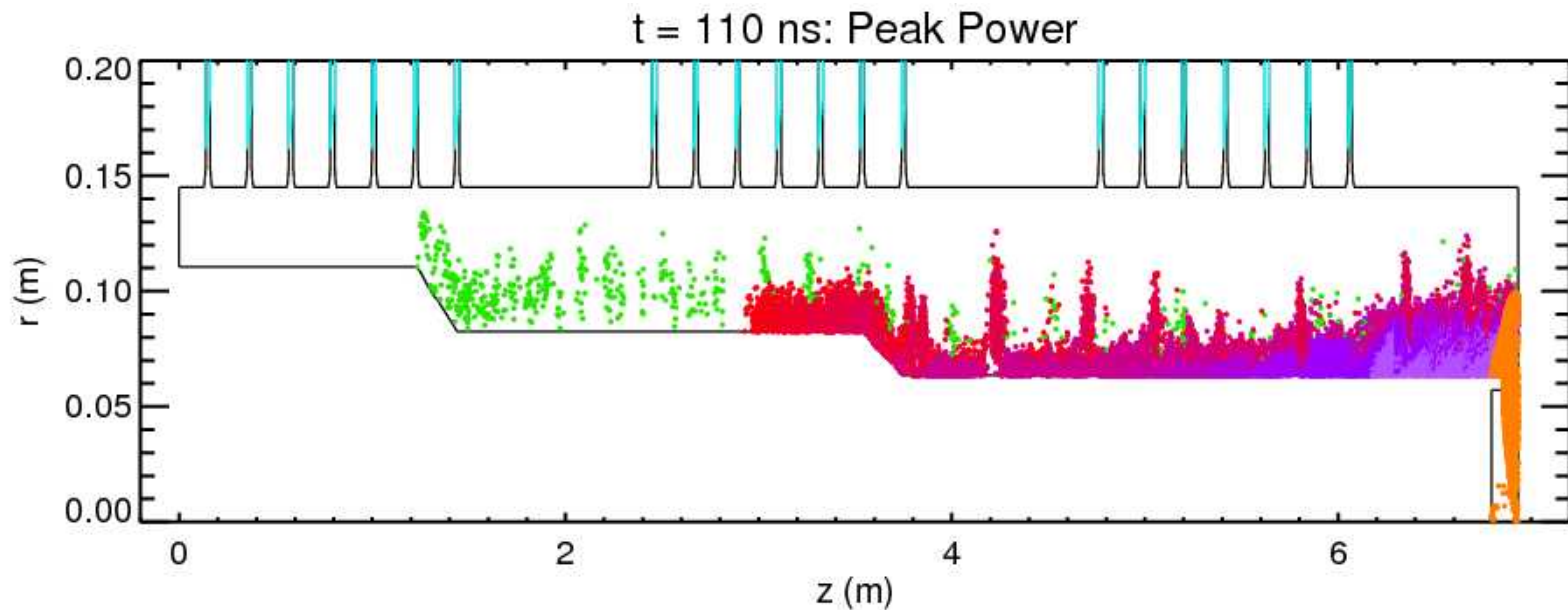
- Simulation has only vacuum flow electrons
- Early-time loss: MITL electrons lost to the anode in the convolute
- Additional late-time loss is due to electrode plasmas in the convolute
 - No detailed data: anode? cathode? both?

Ursa Minor is a testbed for adding many LTD cavities in series for voltage addition



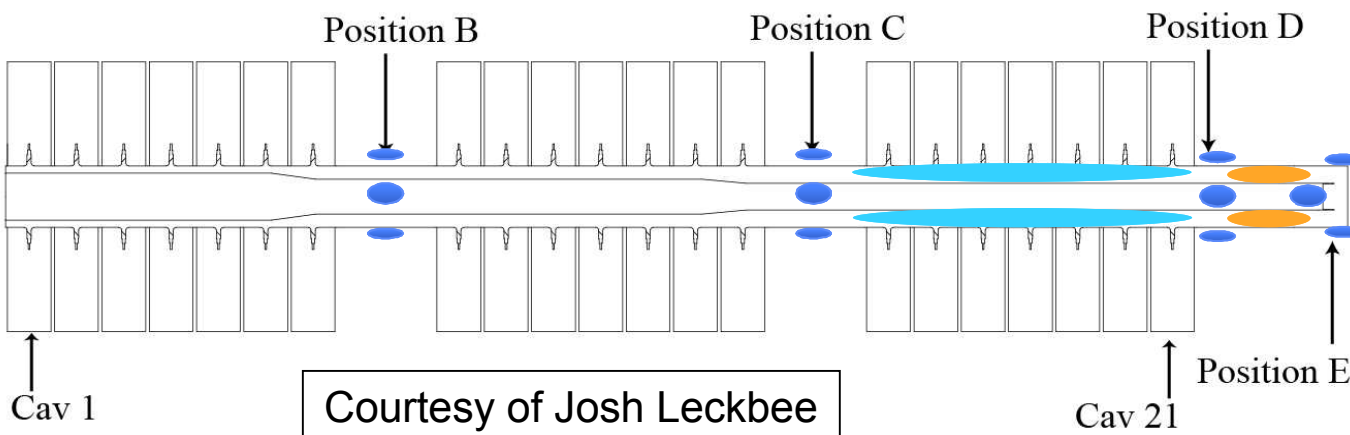
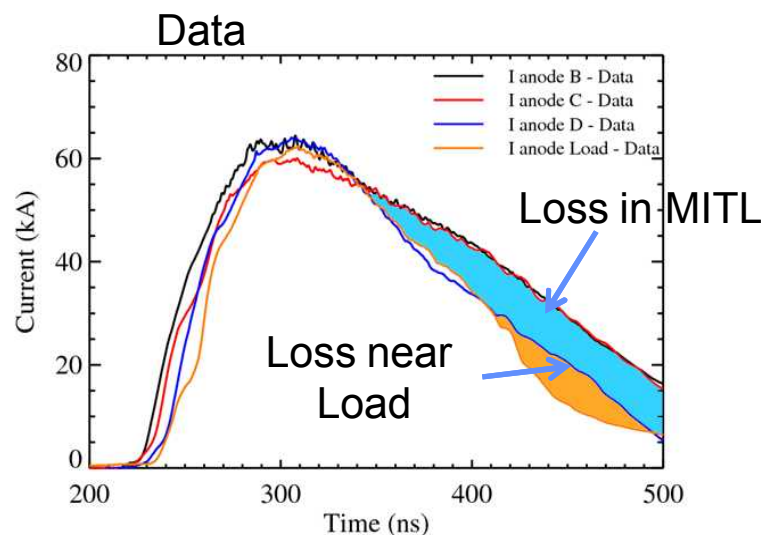
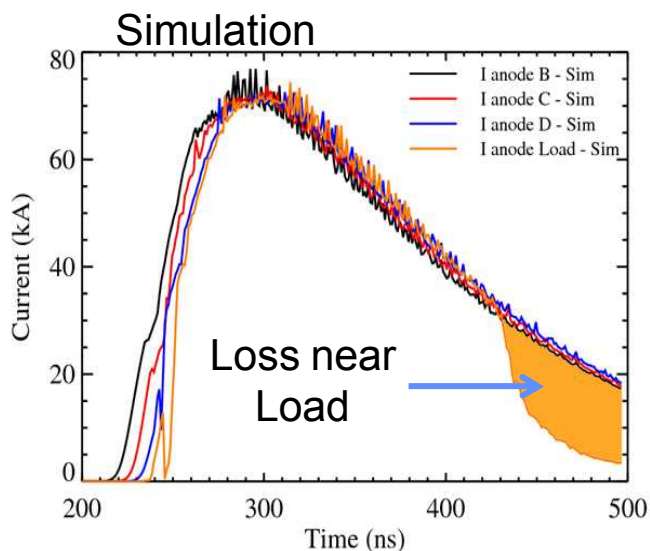
- Ultra-compact pulsed power: each cavity has 10 “bricks” in parallel, composed of two capacitors and a low-inductance switch in oil
 - No need for pulse-forming water section
- 21-cavity voltage adder, 7.5 m long, 1.5 m diameter
 - 180 kV/cavity → ~2.5 MV delivered to the load ... with unavoidable internal losses, but minimal electron flow loss

We have set up a 2-D r-z Quicksilver model of the full 21 cavity geometry



- 1 mm cell size across the 2.2 cm A-K gap of the 21 feeds
- Electrons “color-coded” by creation location
- Key questions:
 - What are electron losses to insulators (shown in cyan)?
 - What factors contribute to the loss?

Experiment has lower peak MITL current and higher loss current than simulations

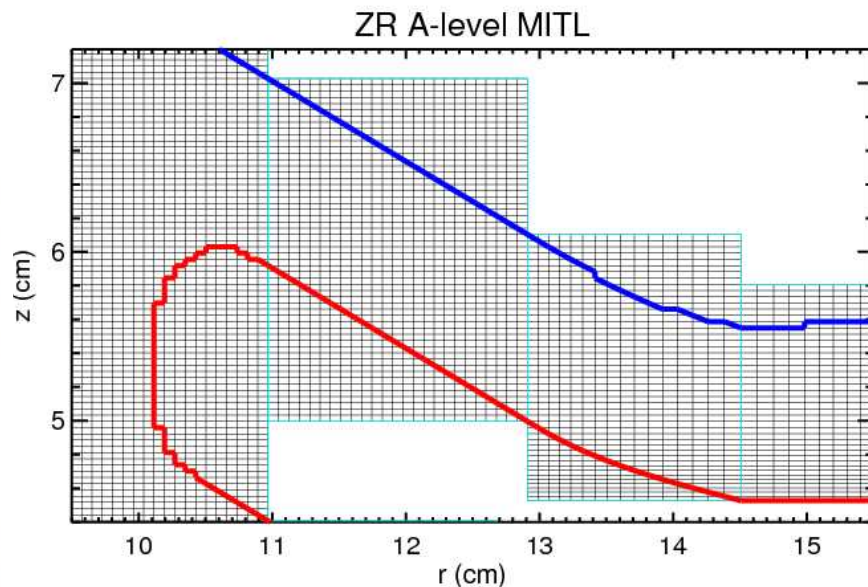
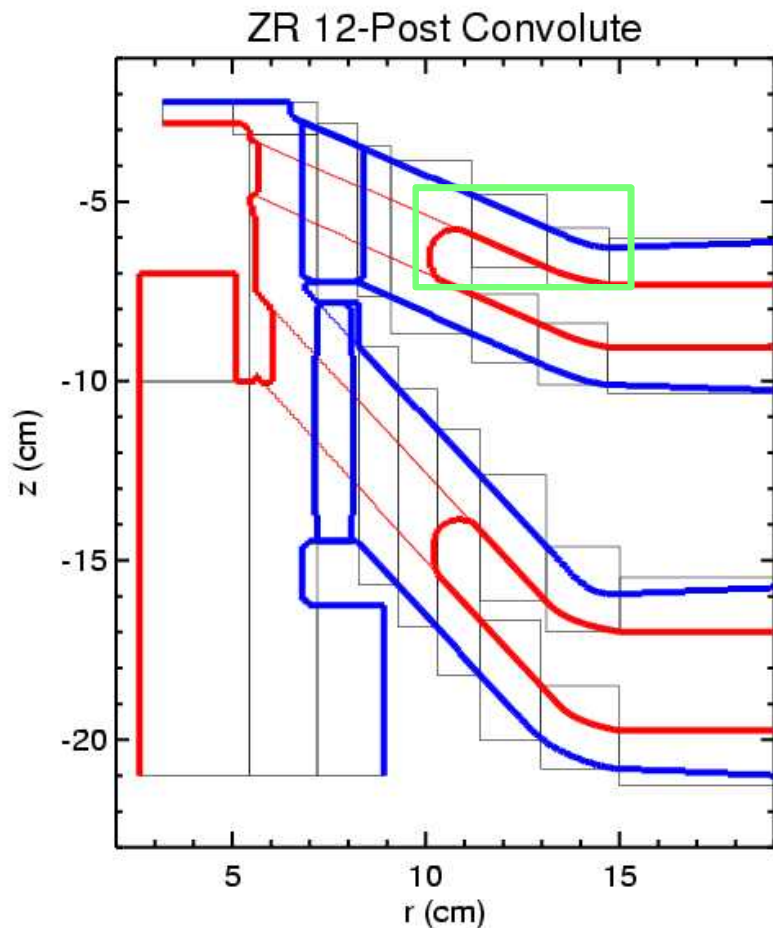




The Quicksilver code has been under development since the mid-1980's

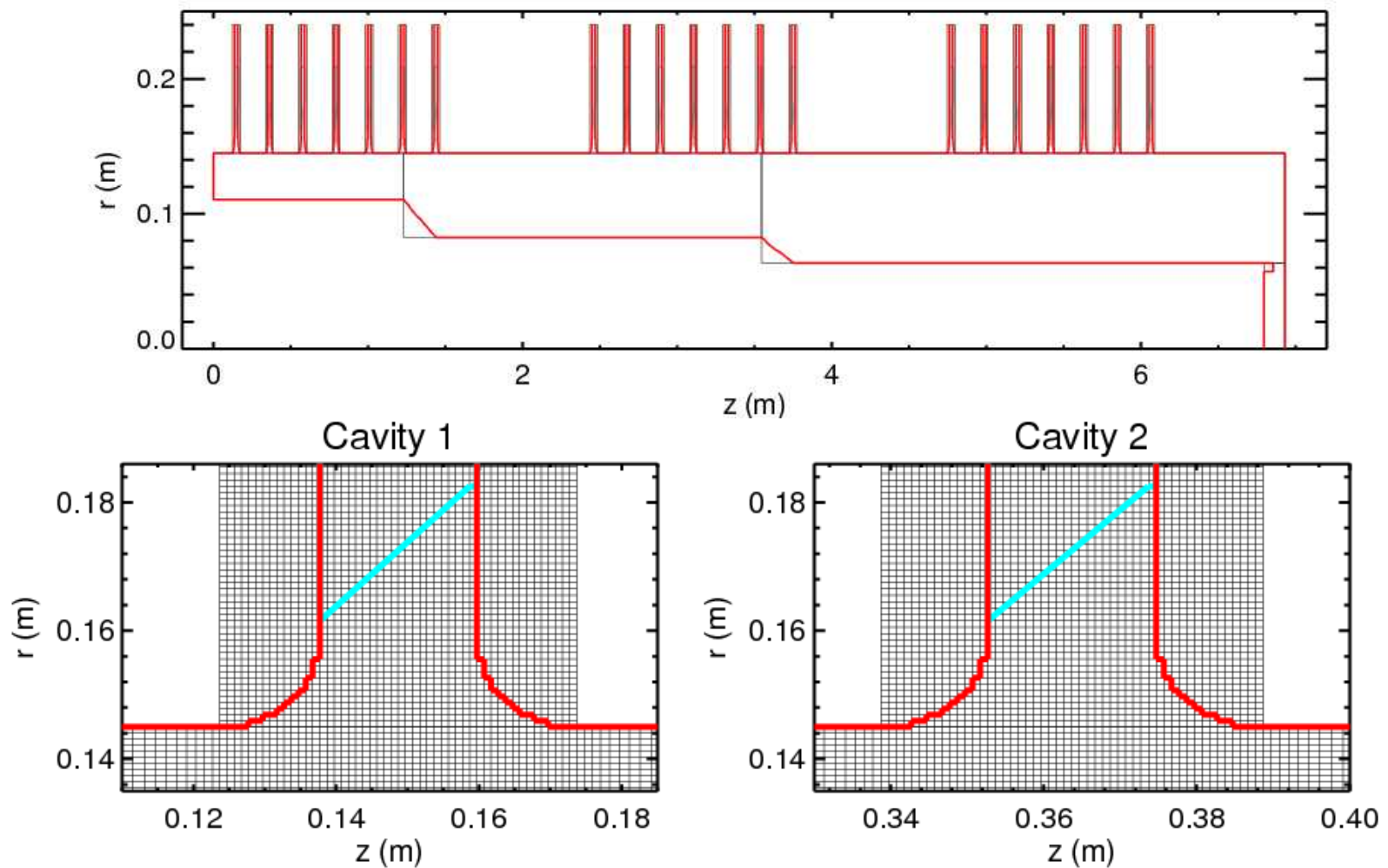
- Originally ANSI-standard F77 for portability; now has F90, C, C++
- Mercury preprocessor
 - Process input text file(s) with symbolic variables, arrays, loops, if-tests, and include files
 - Error checking and computing array sizes
- 2D/3D multi-block geometry
 - Cartesian, cylindrical or spherical coordinates
 - 1D transmission lines for external circuits
- Parallelized with MPI
 - Static decomposition for field solver
 - Dynamic load-balancing for the particle handler
- IDL for all post-processing
 - Also use IDL extensively as a preprocessing tool

QS uses multi-block geometry to efficiently grid complex systems



This geometry is built with IDL scripts and post-hole convolute geometry is imported from an acis file

Multi-block geometry is an essential capability for the Ursa Minor simulations



Augmenting PIC with 1D transmission lines is essential for modeling large systems

- TL's connect to a boundary plane of the 2D/3D domain at a “port” – two conductors (anode,cathode) with arbitrary “x-y” cross-section

- TEM mode E-field determined by:

$$\nabla_T^2 \phi = 0; \phi_{cat} = 0; \phi_{ano} = V_0; V_0 = 1 \text{ Volt}$$

$$\mathbf{E}_p = -\nabla \phi ; \text{“Poisson Field”}$$

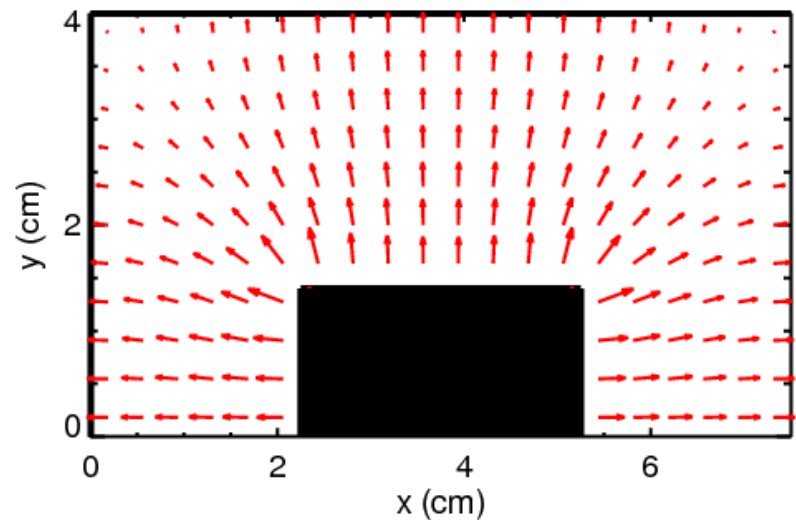
$$\mathbf{E}_{TEM}(V) = (V/V_0) \mathbf{E}_p$$

- Per-unit-length capacitance

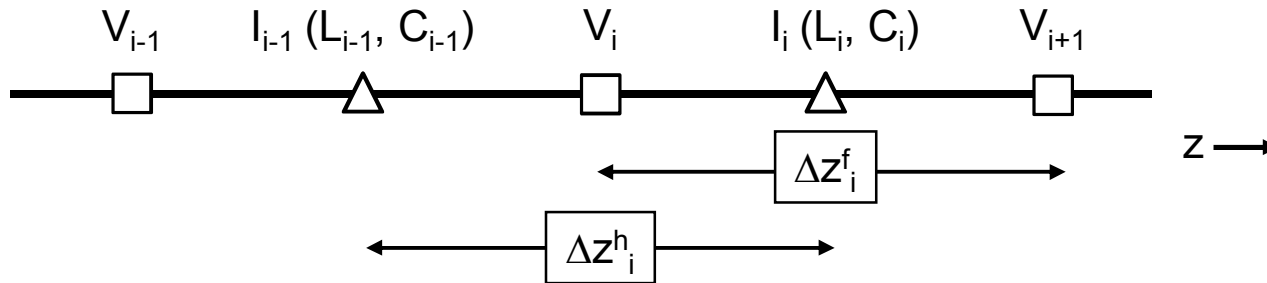
$$C = \frac{\epsilon}{V_0^2} \int E_p^2 dA$$

- Per-unit-length inductance

$$L = 1/(v_{ph}^2 C) ; v_{ph} = (\epsilon\mu)^{-1/2}$$



Voltage and current in the T-line are updated with the telegrapher's equations



- Telegrapher's equations for a lossless line:

$$\frac{\partial I}{\partial t} = -\frac{1}{L} \frac{\partial V}{\partial z} ; \quad \frac{\partial V}{\partial t} = -\frac{1}{C} \frac{\partial I}{\partial z}$$

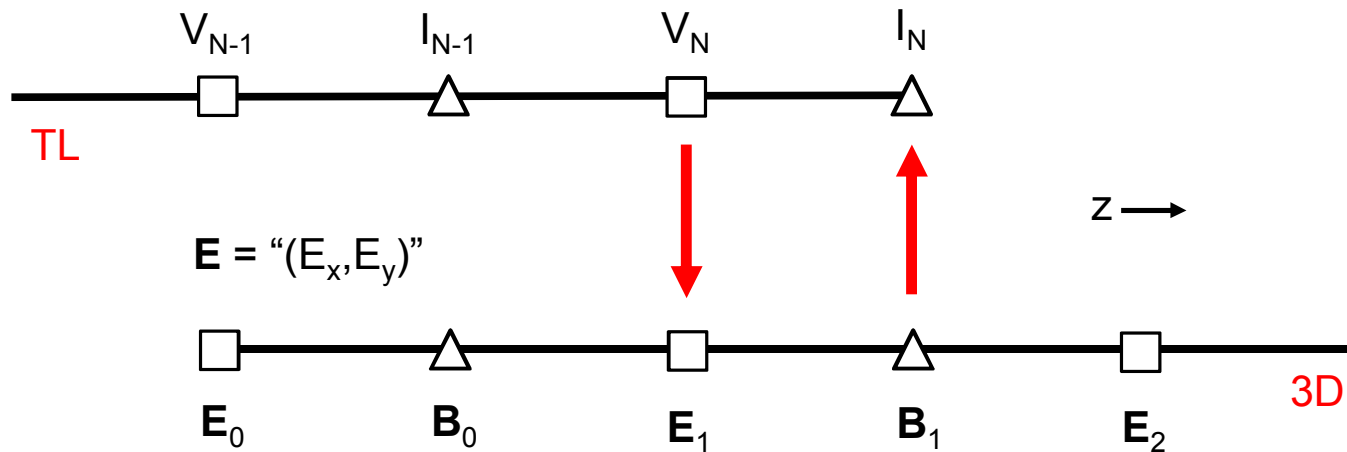
- Leapfrog finite difference, $\{V^{n-1}, I^{n-1/2}\} \rightarrow \{V^n, I^{n+1/2}\}$:

$$V_i^n = V_i^{n-1} - \gamma_i (I_i^{n-1/2} - I_{i-1}^{n-1/2}) ; \quad \gamma_i = 2\Delta t / (\Delta z_i^h [C_{i-1} + C_i])$$

$$I_i^{n+1/2} = I_i^{n-1/2} - \lambda_i (V_{i+1}^n - V_i^n) ; \quad \lambda_i = \Delta t / (\Delta z_i^f L_i)$$

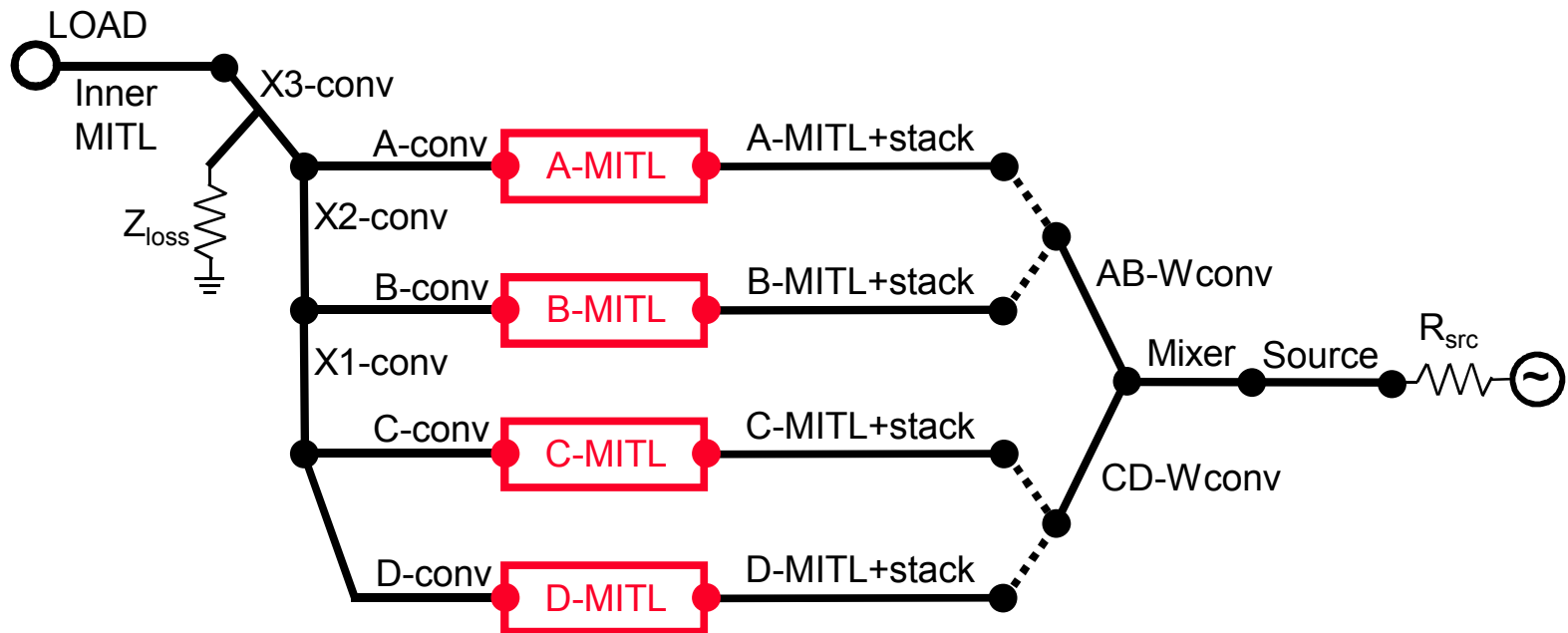
- Blazingly fast compared to 2D/3D field solver update
- Easy to add BC's at the end of a TL: generators, Z-pinch loads, etc.

The transmission line couples naturally to the TEM mode in the 2D/3D system



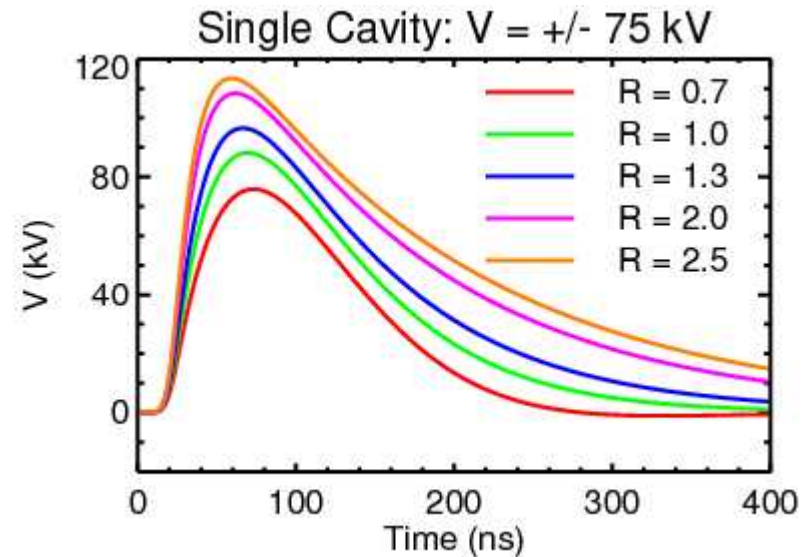
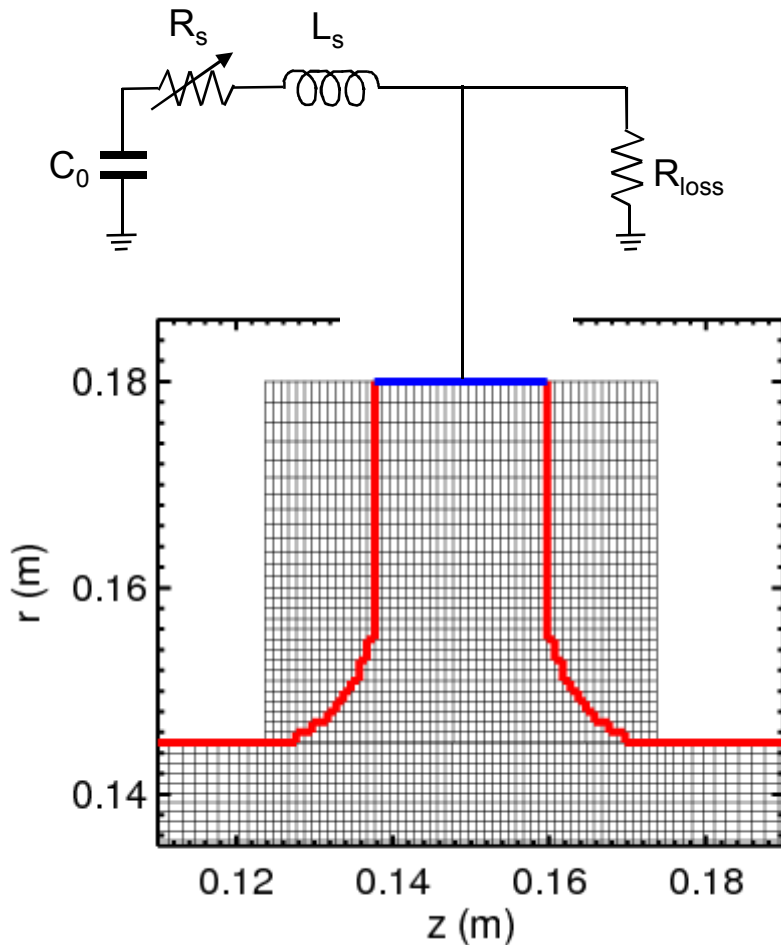
- TL voltage V_N^n determines \mathbf{E}_T^n on boundary plane of the 3D system
 - $\mathbf{E}_1^n = \mathbf{E}_{\text{TEM}}(V_N^n) + \mathbf{E}_{\text{non-TEM}}^n$
 - Hard part is specifying $\mathbf{E}_{\text{non-TEM}}^n$
- 3D magnetic field $\mathbf{B}_1^{n+1/2}$ determines boundary current $I_N^{n+1/2}$
 - $I_N = \frac{1}{\mu V_0} \int \mathbf{e}_z \times \mathbf{E}_p \cdot \mathbf{B}_1 dA$ (filters out non-TEM part of \mathbf{B}_1)
 - This defines V_N^{n+1} on next timestep

High-resolution simulations of the Z MITLs use four 2D PIC regions coupled with TL's



- 2-D PIC region for each MITL, spherical coords., $0.1 < r < 0.6$ m, shown in red
- 1-D transmission lines in black
 - Outer lines extend out to $r \sim 3$ m
 - Inner convolute line parameters obtained from 3-D simulations
 - Optional, time-varying convolute Z_{loss} element
 - Either a Z-pinch or ICE load

Ursa Minor simulations use an external RLC circuit with core loss for each cavity



- $C_0 = 100$ nF, $L_s = 25$ nH, $R_{loss} = 4$ Ω
- R_s : $10^5 \rightarrow 0.2$ Ω with 2 ns exponential decay
- Single cavity into resistive load R agrees to ~5% with more detailed circuit model



Other features in Quicksilver

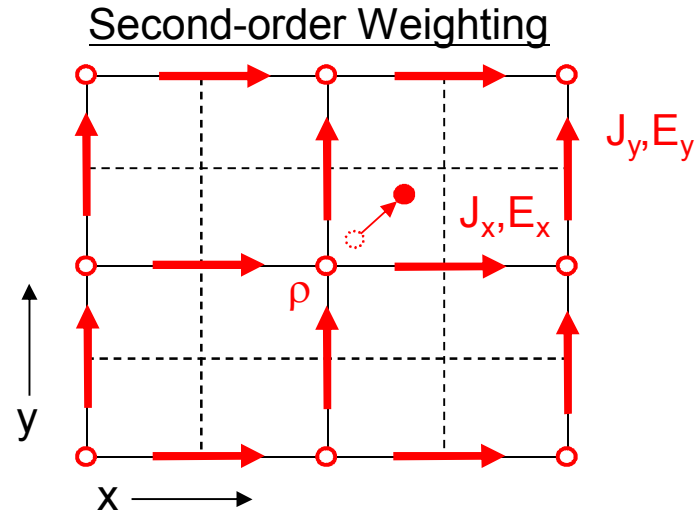
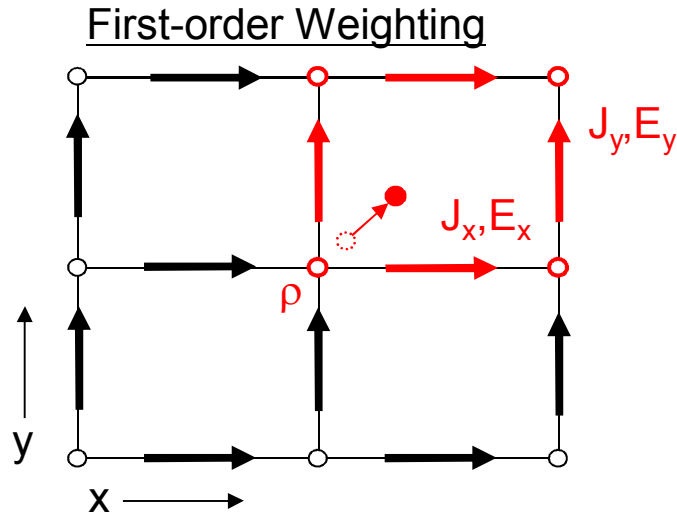
- Friedman's explicit EM field solver with adjustable damping, $0 \leq \theta \leq 1$
 - Also use this scheme to add damping to the explicit particle pusher
- Exact, second-order \mathbf{J} and ρ accumulation, handling boundaries
 - At conductors, morph 2nd \rightarrow 1st order within half a cell of boundary
- Energy-conserving particle pusher: interpolate \mathbf{E} directly from the Yee mesh using exactly the same weighting that \mathbf{J} is laid down on the mesh
 - Enables stability at $\Delta x/\lambda_D \gg 1$, $\omega_p \Delta t \sim 1$ (little numerical heating)
 - Accuracy? *Caveat emptor!*
- Floating-point precision issues:
 - Particle push in local coord. system relative to lower corner of cell
 - Build executables with 4- or 8-byte reals
 - ~ 1.6 slowdown for r8 – memory access



Other features in Quicksilver, contd.

- Sort particles by cell and species:
 - Required for Coulomb collisions and particle merger
 - Improved performance by limiting cache-thrashing
- Particle-particle Coulomb collisions for cold, dense plasmas
- Electron-surface interactions
 - Heating
 - Reflection and secondary emission
- Kinetic gas-chemistry model
 - Gas breakdown plasma modeled with electron-ion pairs
 - Robust particle merger to deal with exponentially-growing particle count

Energy-conserving PIC: Specific scheme for E-field interpolation to the particles



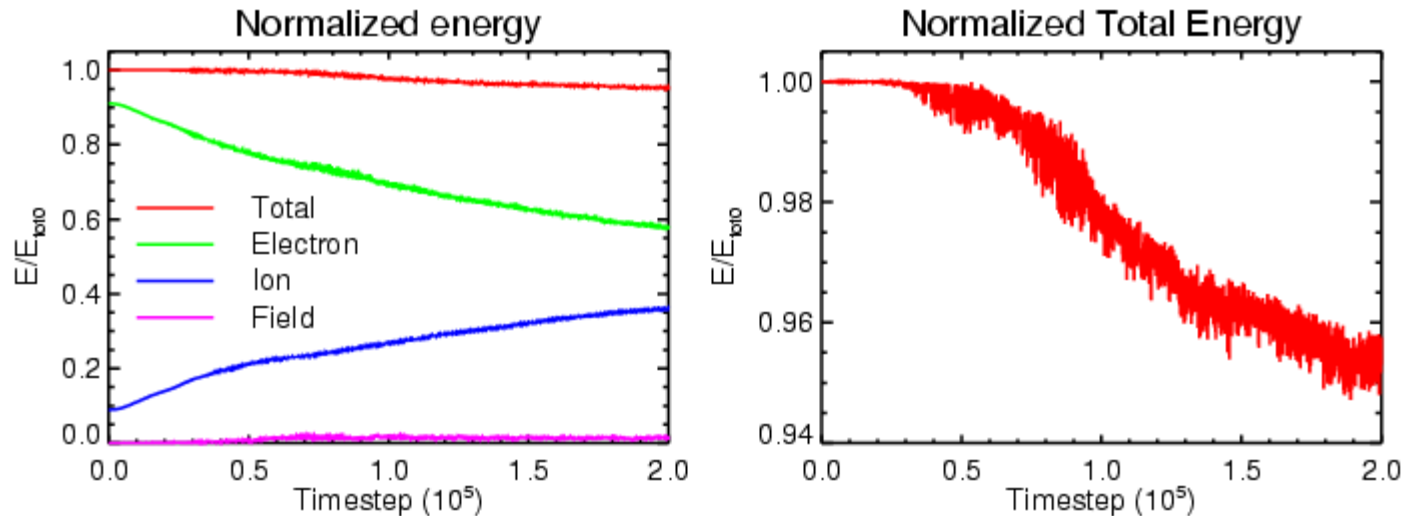
- For EM codes: interpolate \mathbf{E} directly from the Yee mesh to the particles, using exactly the same weighting scheme as \mathbf{J} from the particles to the grid
- First-order weighting: Interpolation of E_x is nearest-grid-point in x , linear in y
 - Although it appears crude, we have found it very useful
- Second-order weighting: Interpolation of E_x is linear in x , quadratic in y
 - Handle boundaries: T. Pointon, Comput. Phys. Commun. **179** (2008) 535.



Energy-conserving PIC removes the restriction of having to run at $\Delta x/\lambda_D < \sim 1$

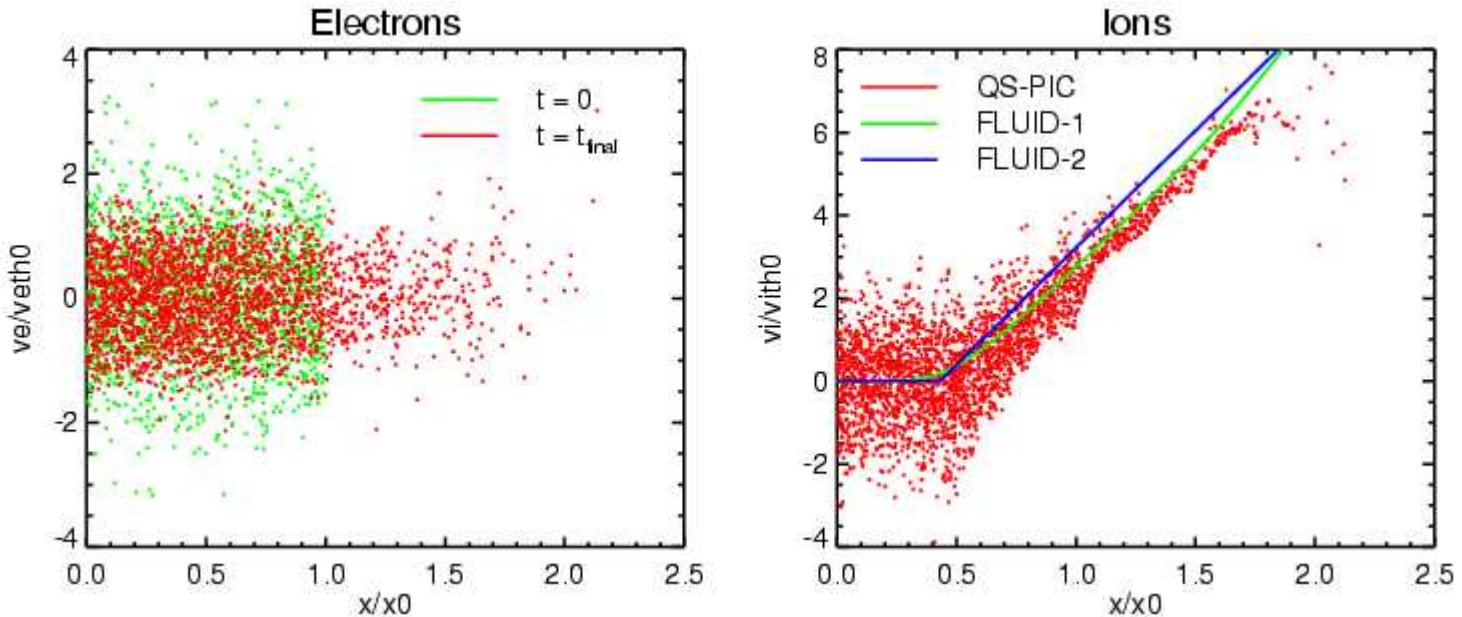
- Explicit, energy-conserving PIC runs are stable for $\Delta x/\lambda_D \gg 1$
 - Minimal numerical heating
- Stability does not imply accuracy!
 - 1st-order algorithm cannot handle extremely large E-field gradients
 - Neither algorithm can accurately model a plasma at temperature T in presence of a large external E-field satisfying $qE\Delta x \gg T$
- For well-chosen parameters, neither of these factors is an issue, and we can run at $\Delta x/\lambda_D \gg 1$ with reasonable accuracy, limited only by $\omega_p \Delta t < \sim 1$
 - For $\Delta t = 1$ ps, $\omega_p \Delta t = 1 \rightarrow n < 3.1 \times 10^{14} \text{ cm}^{-3}$ ($\sim 1\%$ ionization at 1 Torr)
- However ... ‘On two occasions I have been asked “Pray Mr. Babbage, if you put into the machine wrong figures, will the right answer come out?” I am not able rightly to apprehend the kind of confusion of ideas that could provoke such a question’, Charles Babbage, 1864

Canonical test for high-density plasma: 1-D expansion of plasma slab into vacuum



- Use $\Delta x = 0.5$ mm, $\Delta t = 0.5$ ps; characteristic of Z-convolute runs
- Plasma density near explicit limit: $n = 10^{15}$ cm $^{-3}$; $T_e = 1$ eV, $T_i = 0.1$ eV
 - $\omega_{pe}\Delta t = 0.89$, $\Delta x/\lambda_{De} = 2130$
- ~5% numerical cooling after 200,000 timesteps
(~4% with no damping in the field solver)

Particle phasespace plots show the plasma expansion is modeled quite well



- Transfer thermal electron energy to directed ion energy at the front
- Results agree well with 1-D multi-fluid (electron/ion) simulations
 - 1st multi-fluid simulation uses same Δx as the QS run
 - 2nd m-fluid simulation uses Δx 2000 times smaller ($\Delta x/\lambda_{De} \sim 1$)



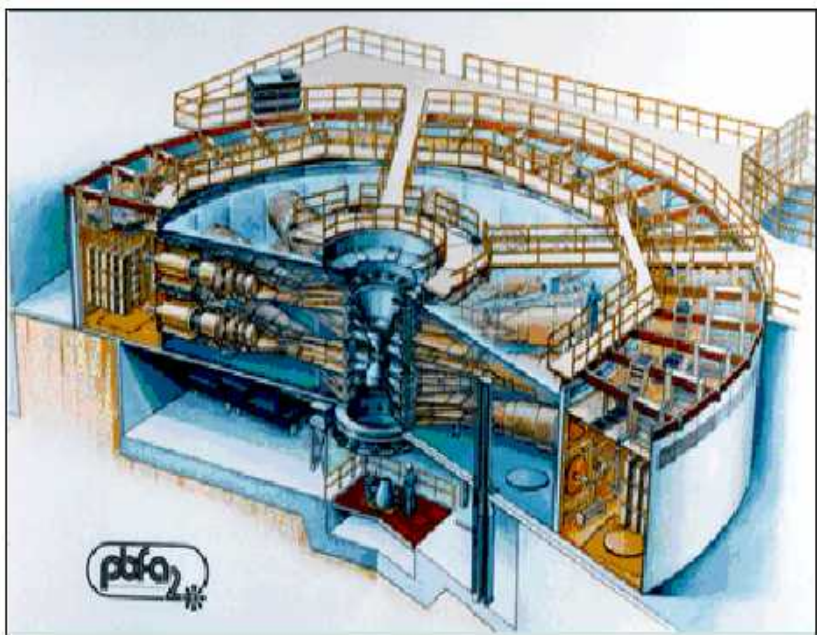
Summary

- Quicksilver is a mature 2D/3D EM-PIC code with a set of features driven by it's application space
 - Pulsed power systems
 - EM response to gas-filled cavities with electron emission from walls
- Quicksilver features
 - Excellent capabilities for handling complex geometry
 - Energy-conserving PIC
 - Stability at $\Delta x/\lambda_D \gg 1$, $\omega_p \Delta t \sim 1$ (minimal numerical heating)
- Post-processing with IDL

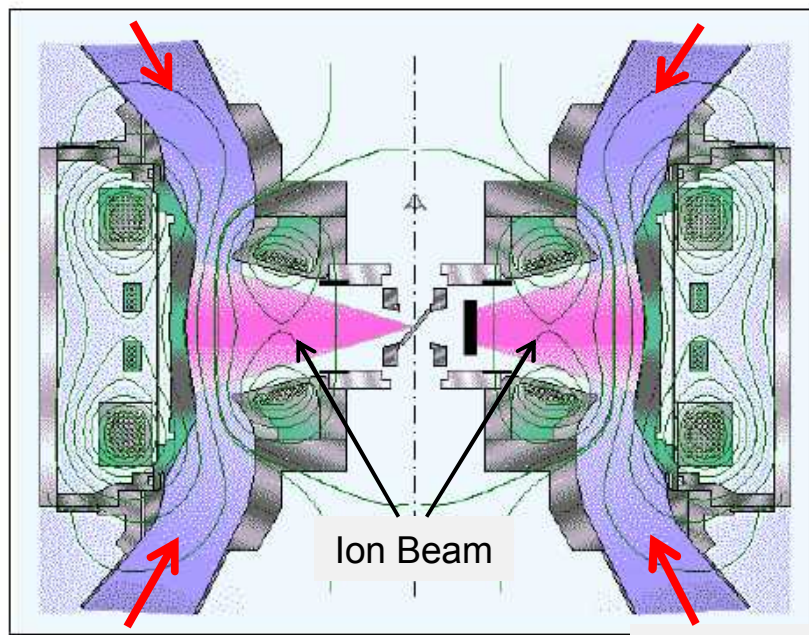


Backup Slides

The applied-B ion diode on PBFA II produced 10 MV, 1 MA, 40 ns Li^+ ion beams



← 33 m →



← 15 cm →

EM Power +
Electrons

- 2D PIC simulations: coupling 4 MITLs in series to achieve 10 MV
- 3D PIC simulations: EM instabilities → beam divergence
- 1D DYNAID simulations: anode plasma expansion → gap closure



Quicksilver and IDL use PFF (portable file format) to handle floating-point data

- PFF was developed by David Seidel in late 1980's for QS
 - Save disk space by encoding floating-point arrays into 16-bit integer arrays using an offset and multiplier
 - Full precision is now available if necessary
 - Pre-defined “dataset types” for data relevant to PIC
- User interface at the dataset level is machine-independent
 - UF1: uniformly spaced 1-D data (time histories)
 - NF3, NV3: Scalars and 3-D vectors on 3-D multi-block domain (field snapshots)
 - VTX: List of vertices on an m-D space with n attributes per vertex (particle snapshots, surface field snapshots)
 - NGD: m-D vectors on a single n-D block
- Public domain: <http://sourceforge.net/projects/hermes-util>



Quicksilver uses IDL for post-processing ... and for more and more pre-processing

- “PFIDL” layer of IDL routines (Paul Mix)
 - PFF dataset read/write
 - Waveform (1D time history) processing: plot, arithmetic, FFT, *etc.*
 - “Structure” processing: 2D/3D field snapshots, particle snapshots
- My IDL routines for post-processing
 - Time history modification: filtering, *etc.*
 - Field, lineout, and particle snapshot animation
- My IDL routines for QS pre-processing
 - Generation of simulation geometry
 - Processing rad-transport simulation output for EM simulations: electron reflection, secondary emission and surface heating
 - Processing gas-chemistry cross-section data

Quicksilver/Mercury/IDL flowchart

