

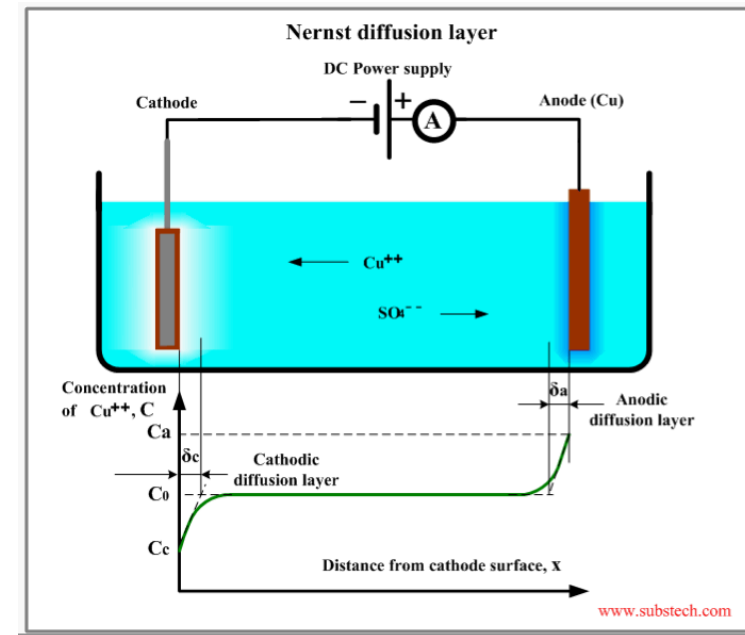
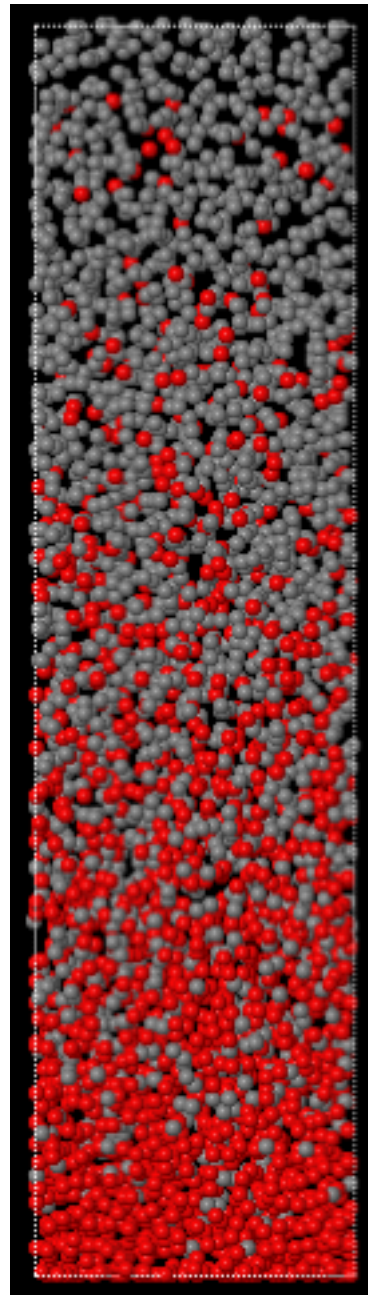
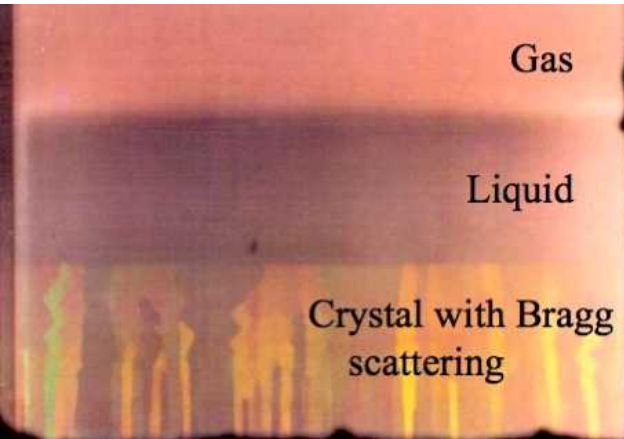
Exceptional service in the national interest



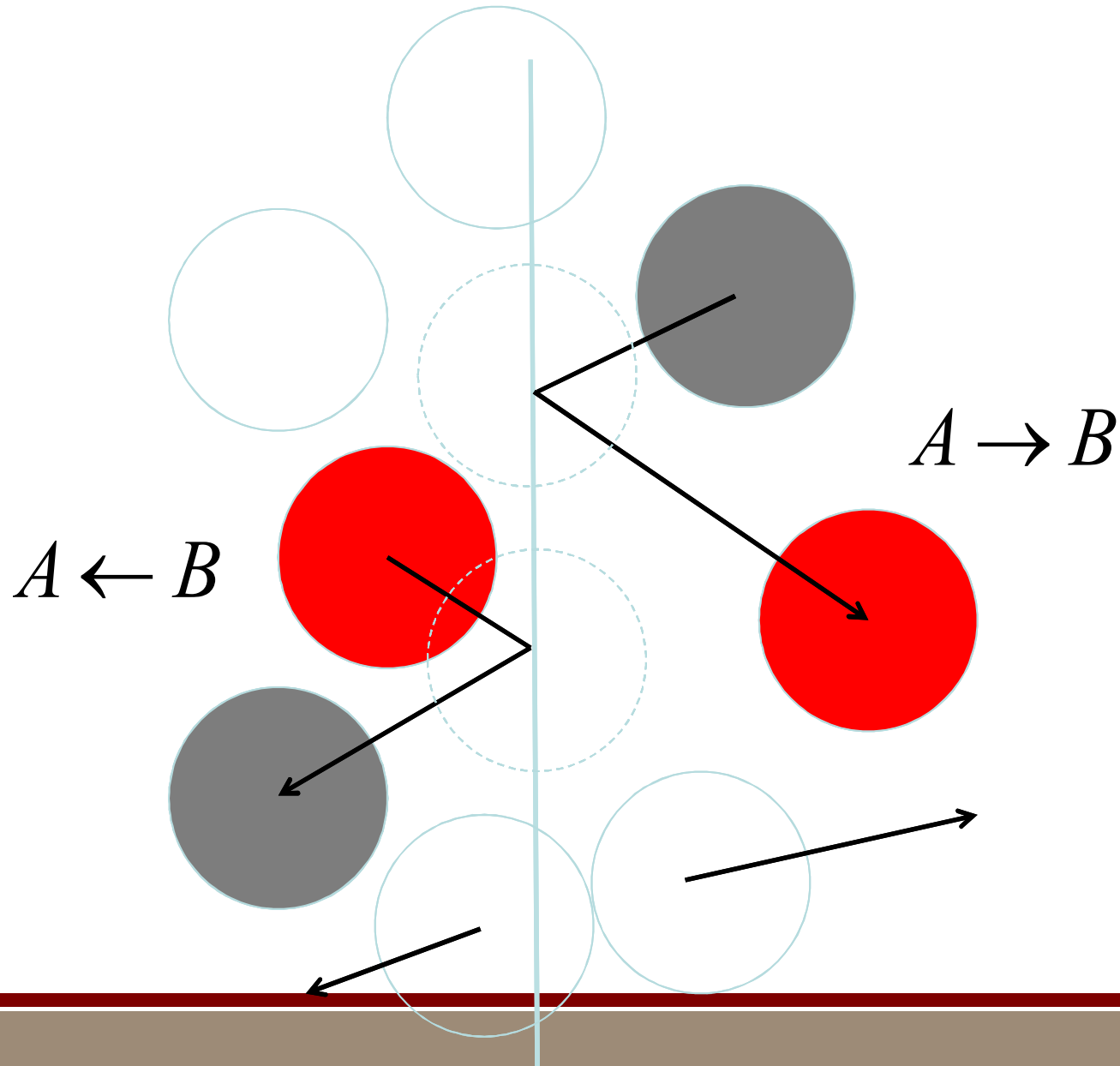
Reaction-Diffusion in Inhomogeneous Fluids

Frank van Swol

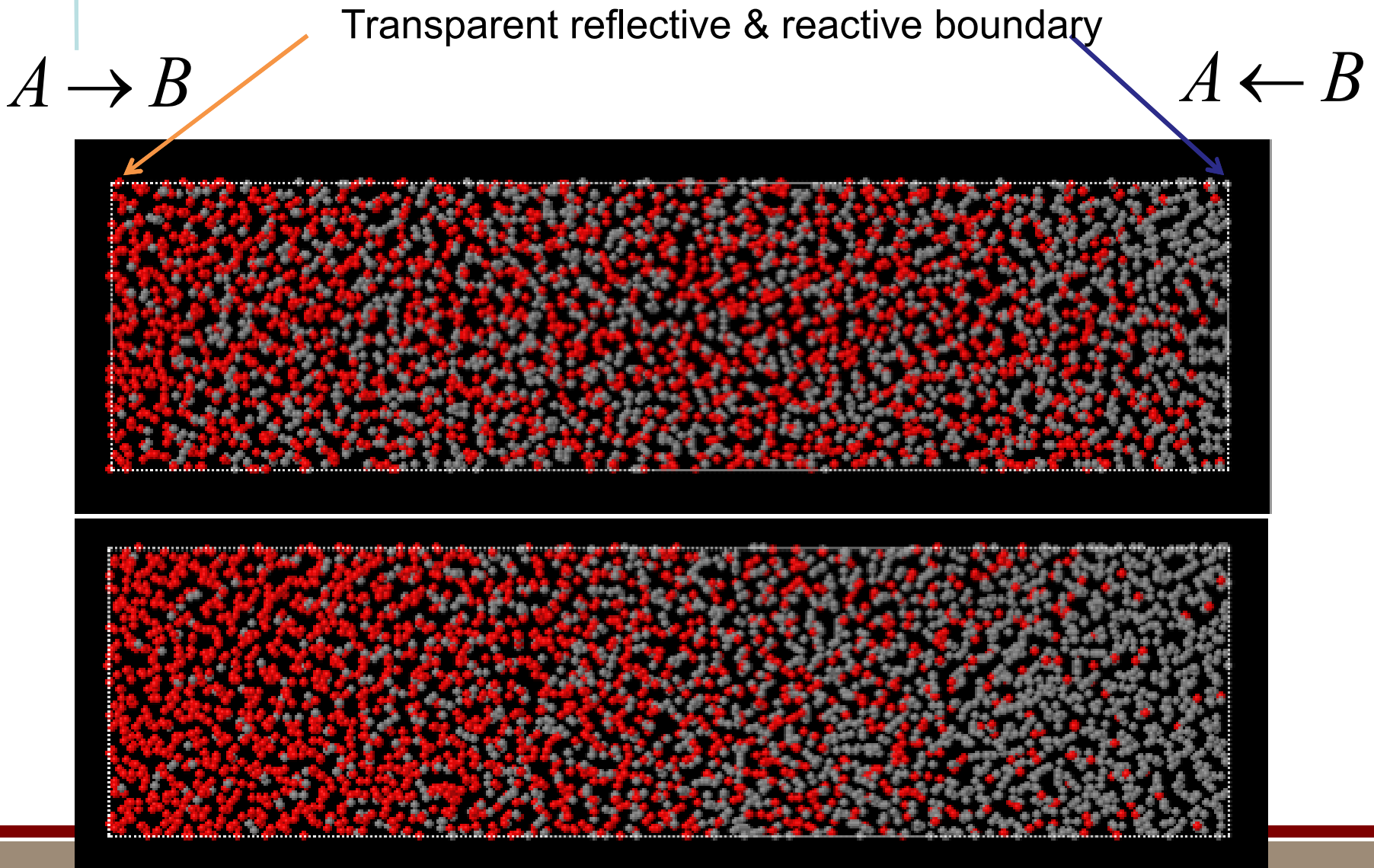
Sandia National Laboratories



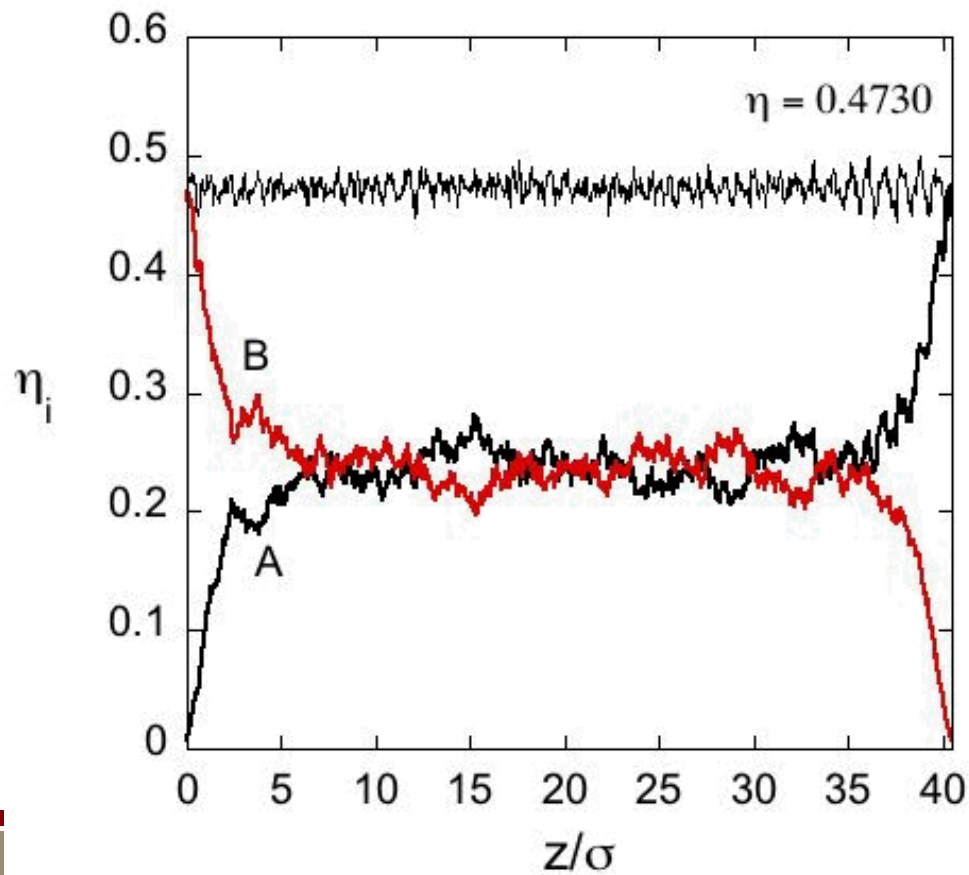
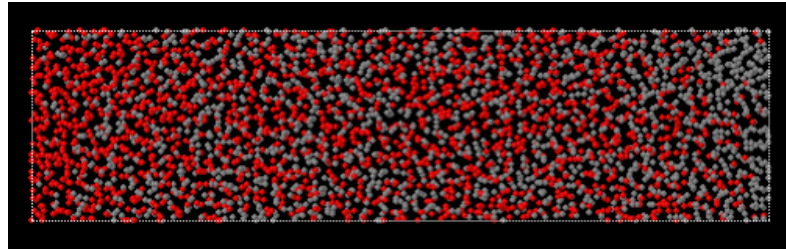
Reaction- (color) diffusion



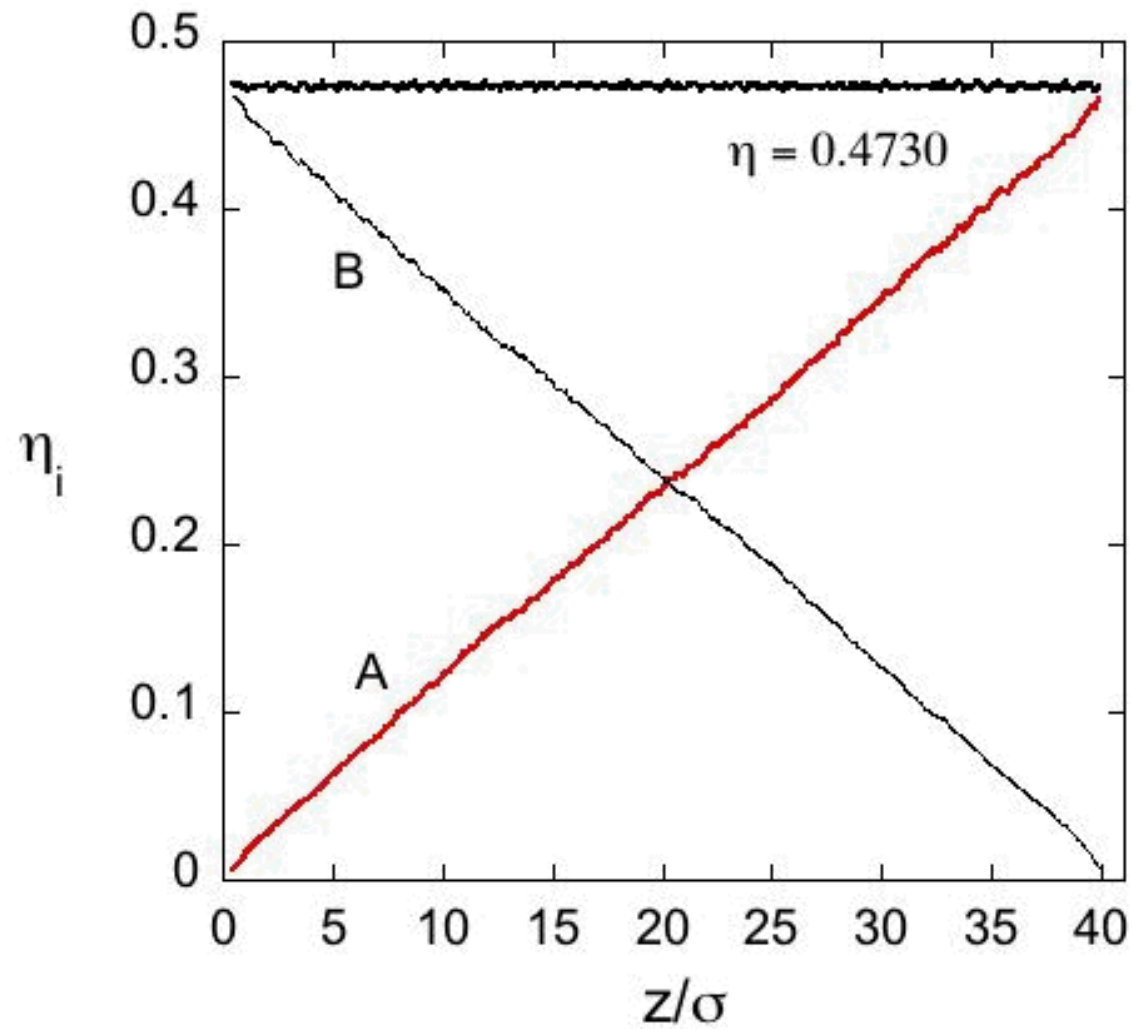
Reaction-color diffusion (PBC)



Transients



Steady-State

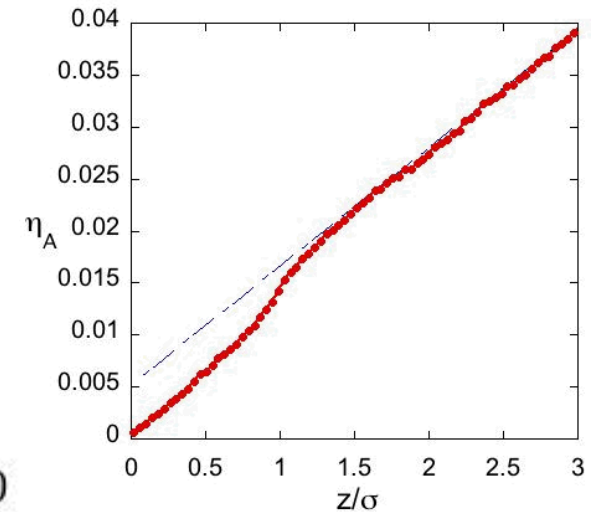
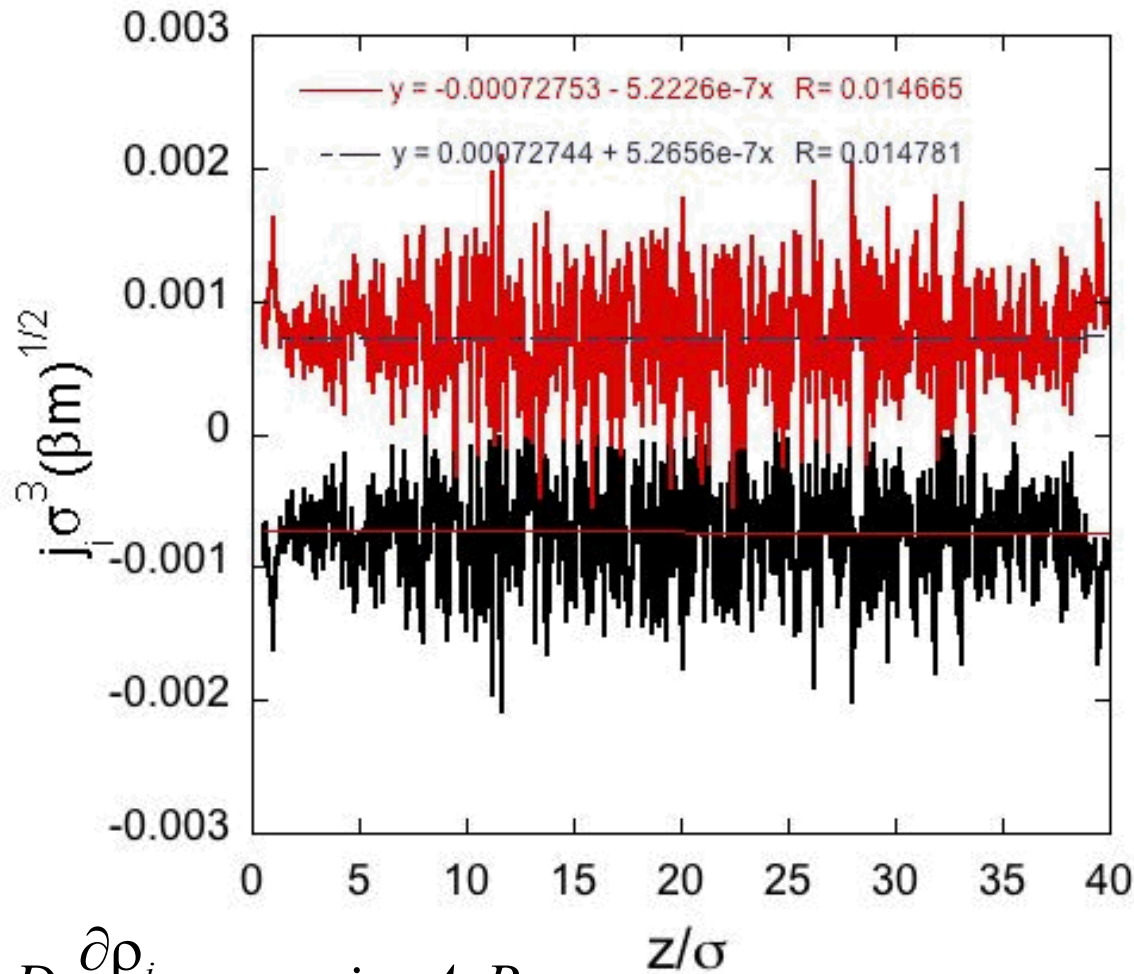


Adolph Fick 1855



$$j_i = -D_i \frac{\partial \rho_i}{\partial x}$$

Color fluxes



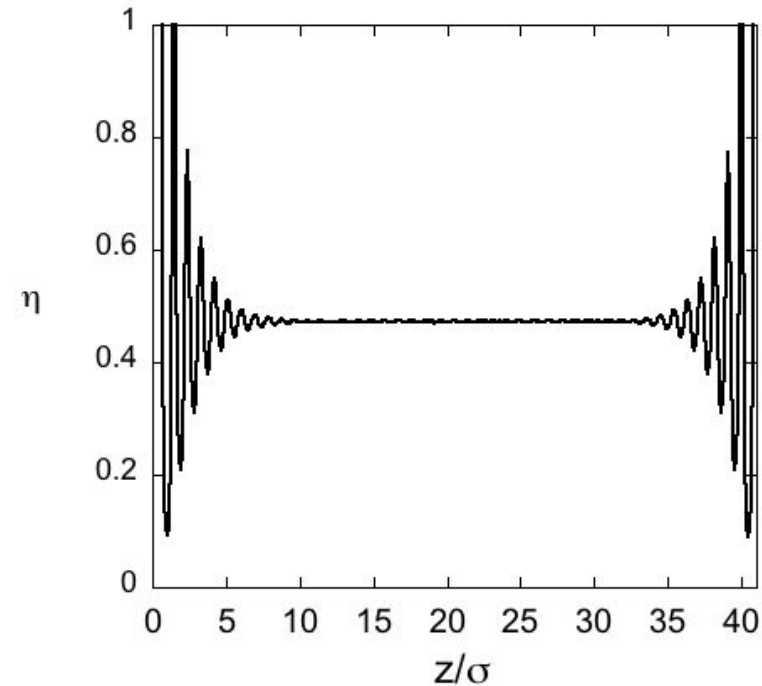
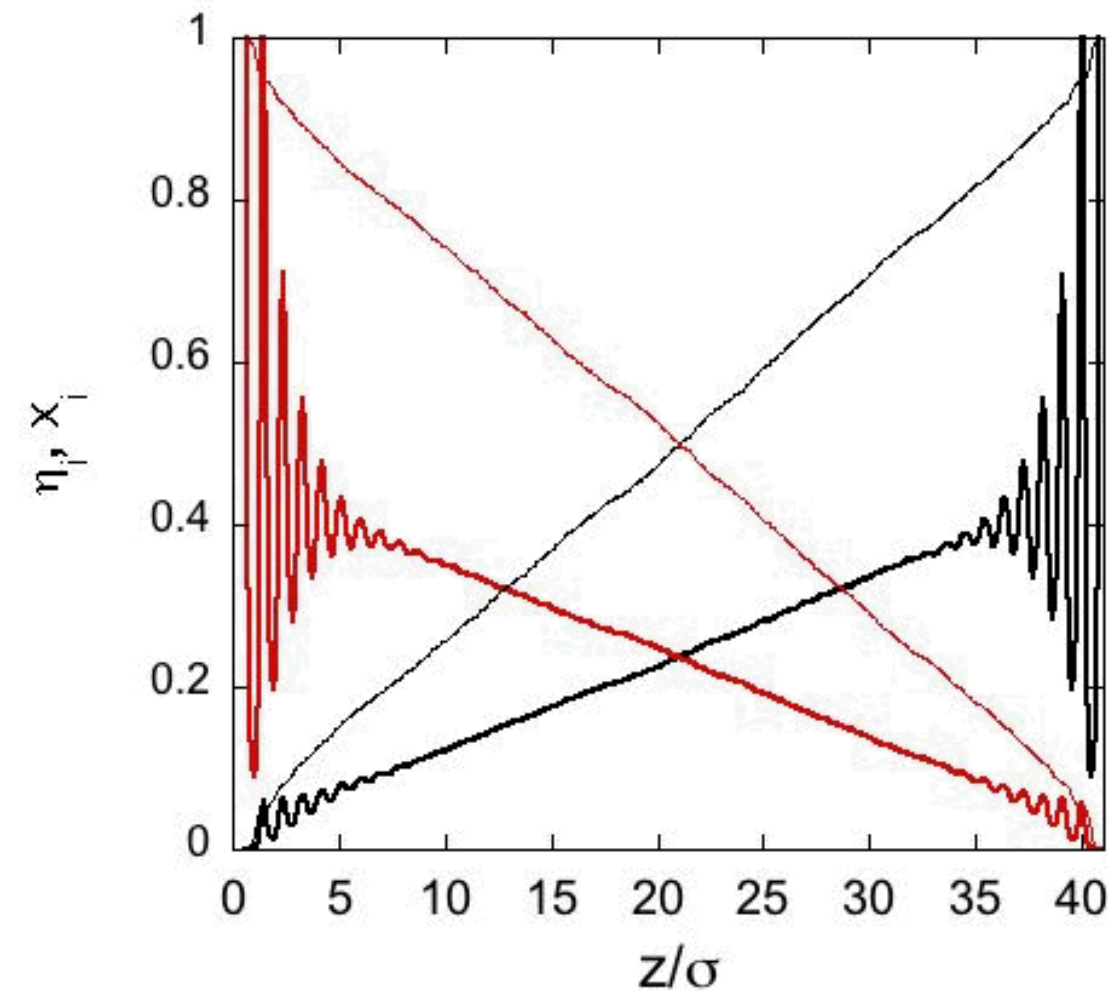
$$j_i = -D_i \frac{\partial \rho_i}{\partial x} \quad i = A, B$$

$$D_i = D_s = 0.0339 \sqrt{kT\sigma^2 / m}$$

$$R_{A \rightarrow B} = \frac{\rho_A}{\rho_A + \rho_B} p_{kin} = 0.000729$$

Inhomogeneous Hard Sphere Fluid

reactive boundary = hard wall



$$j_i = -D \frac{\partial \rho_i}{\partial x} \quad ?$$

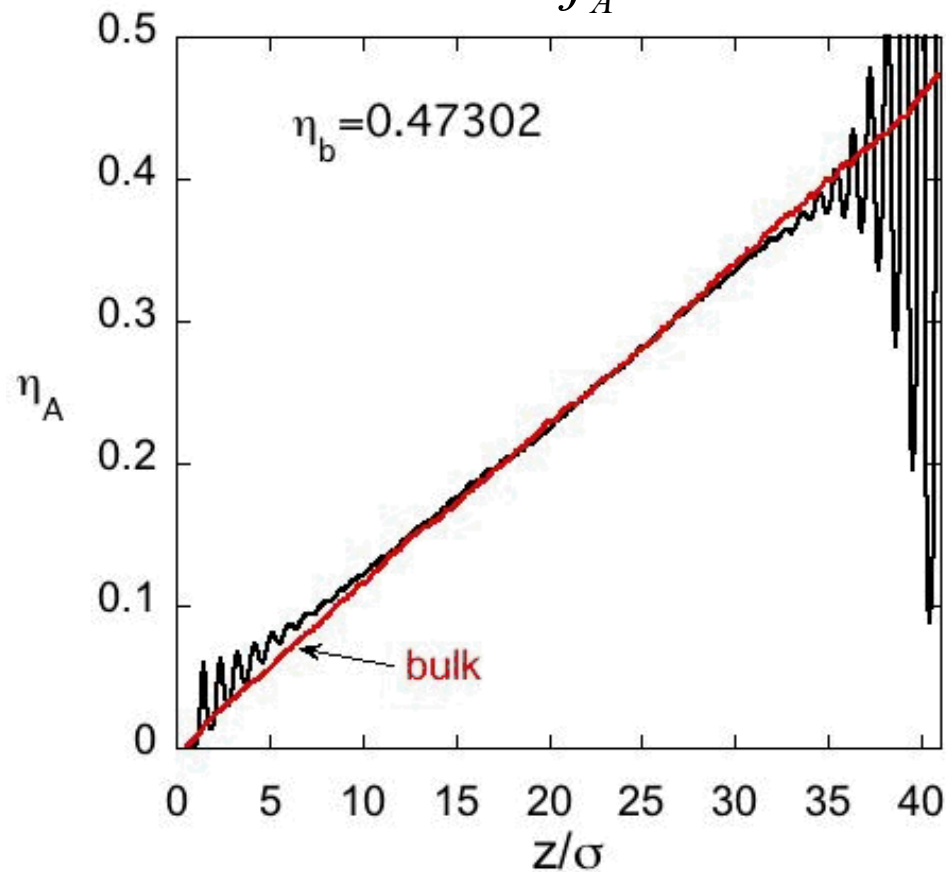
Comparing fluxes

$$j_A = 0.669 \cdot 10^{-3}$$

inhomogeneous

$$j_A = 0.937 \cdot 10^{-3}$$

bulk



What to do?

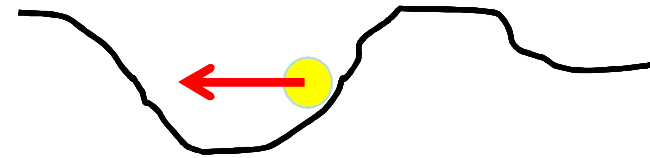
$$j_i = -D \frac{\partial \rho_i}{\partial x}$$

- Consider $D_i = D_i[\rho_i(z)]$?
- Consider a different equation?

1)
$$j_i = -D \rho_i \frac{\partial \mu_i}{\partial x}$$

2) Smoluchowski equation

v. Smoluchowski (1906) Kramers(1940)-Klein(1922)

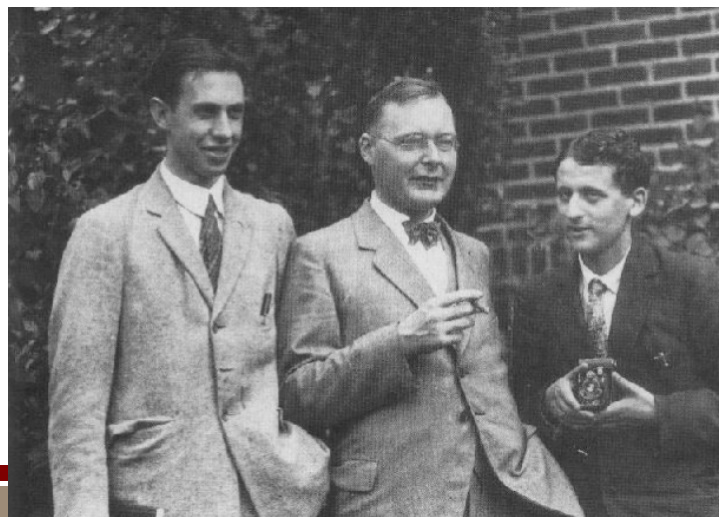


$$j_i = -D \frac{\partial \rho_i}{\partial z} - D \frac{K(z) \rho_i}{kT}; \quad i = A, B$$

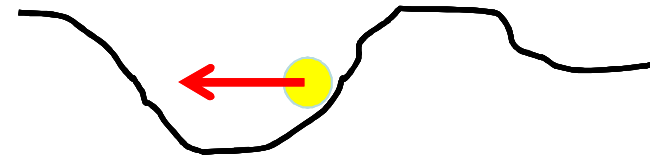
Chemical kinetics and diffusion approach: the history of the Klein-Kramers equation, S. Zambelli, *Archive for History of Exact Sciences*, 64, 395 (2010)

Brownian Motion: Fluctuations, Dynamics, and Applications

Robert M. Mazo



Extracting the force $K(z)$
from the SS profiles

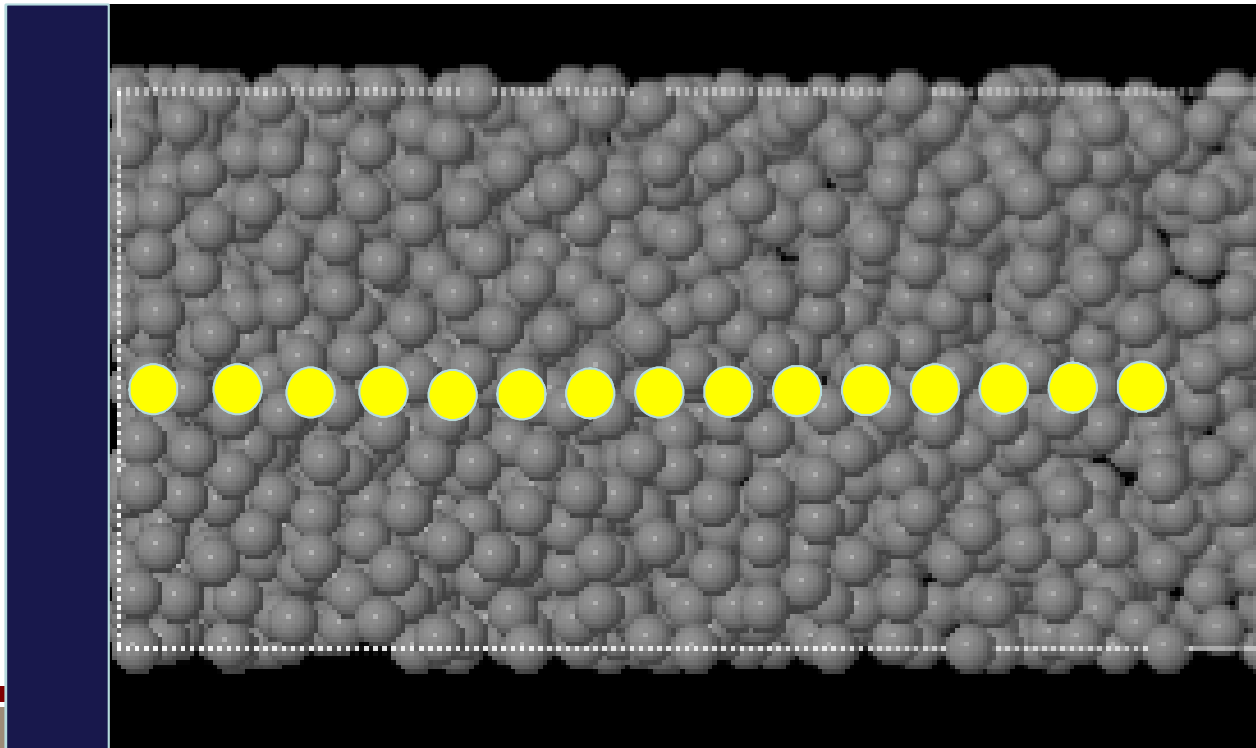
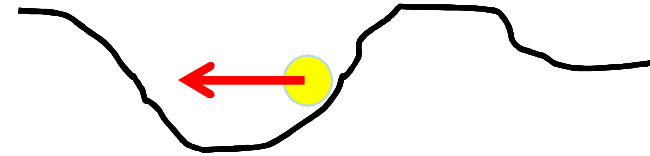


$$j_i = -D \frac{\partial \rho_i}{\partial z} - D \frac{K(z) \rho_i}{kT}; \quad i = A, B$$

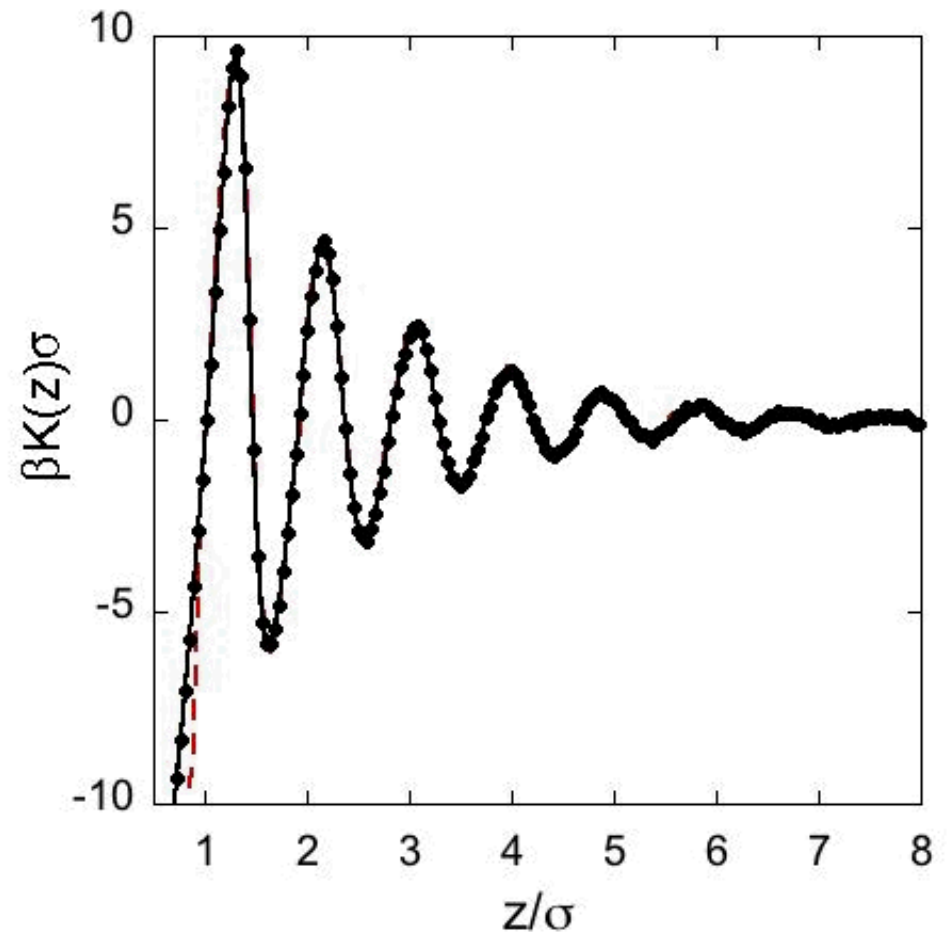


$$\frac{K(z)}{kT} = \frac{j_i}{D \rho_i} + \frac{\partial \ln \rho_i}{\partial z}; \quad i = A, B$$

Alternatively:
measure $K(z)$ directly



$K(z)$ = Solvation Force



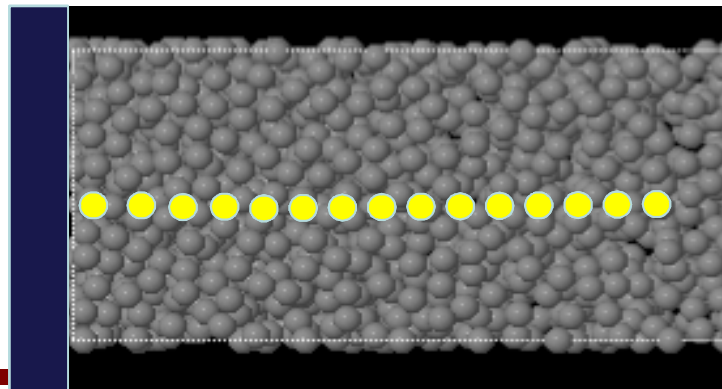
9.5 Potential Distribution Theorem (hard walls)

Pure fluid, at equilibrium:

$$\begin{aligned}\beta\mu &= \ln\rho(z) - \ln \langle e^{-\beta U_t(z)} \rangle ; 0 < z < L_z \\ &\equiv \ln\rho(z) + \beta\tilde{\mu}^{ex}(z)\end{aligned}$$

From which it follows that

$$\frac{K(z)}{kT} = \frac{\partial \ln\rho}{\partial z} = \frac{\partial \beta\tilde{\mu}^{ex}(z)}{\partial z}$$

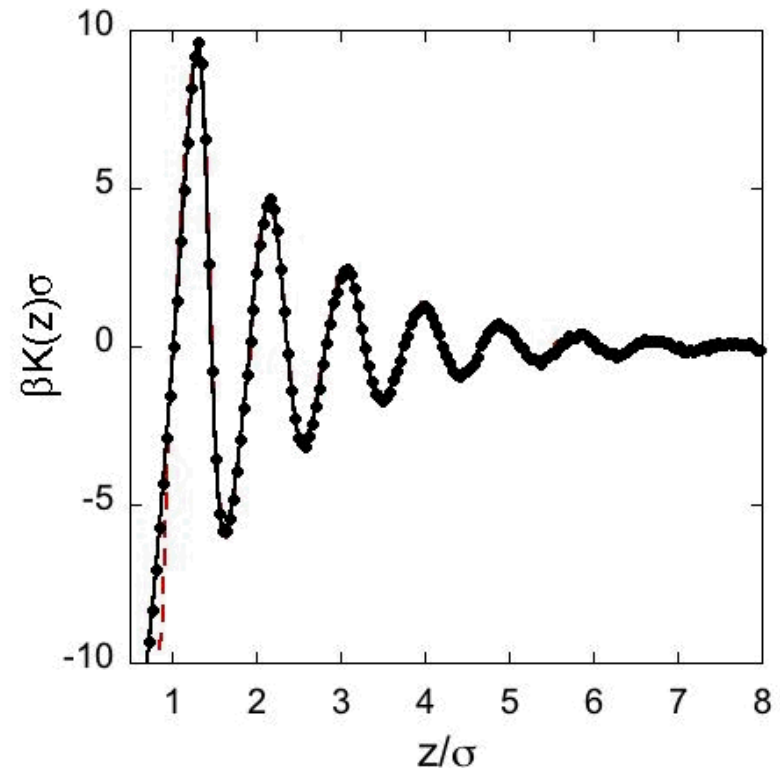


K(z) from three routes

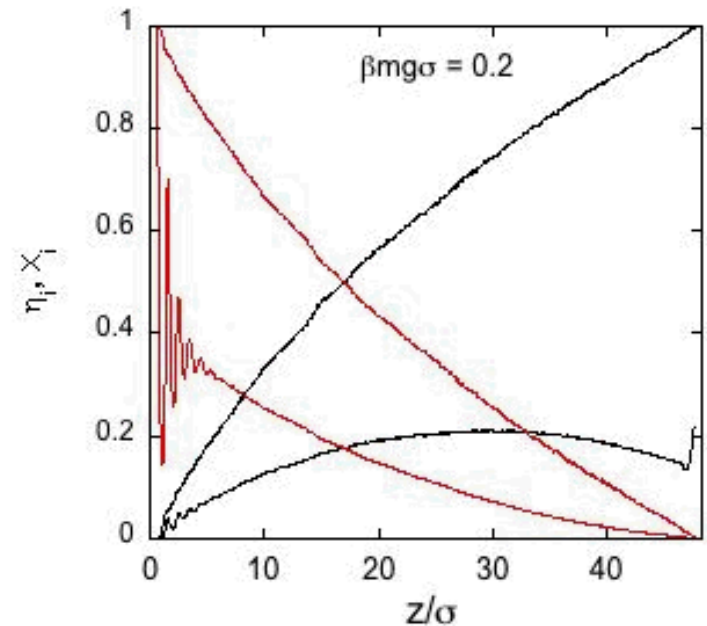
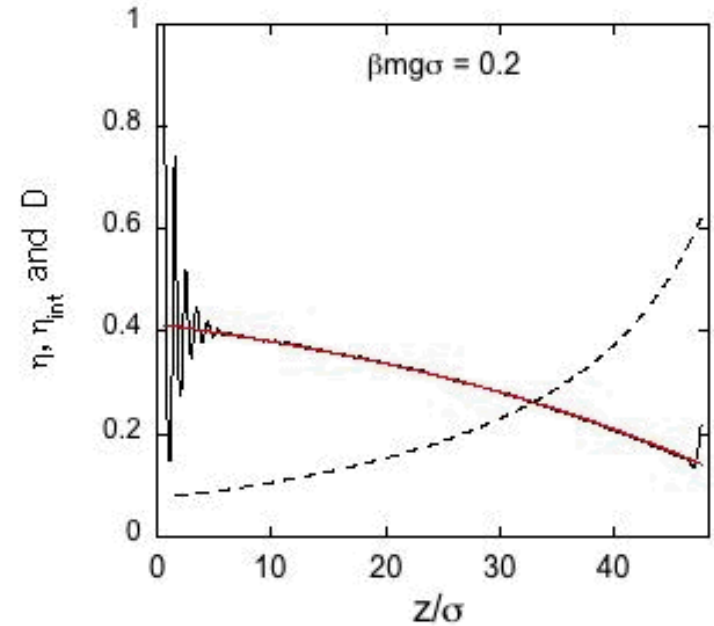
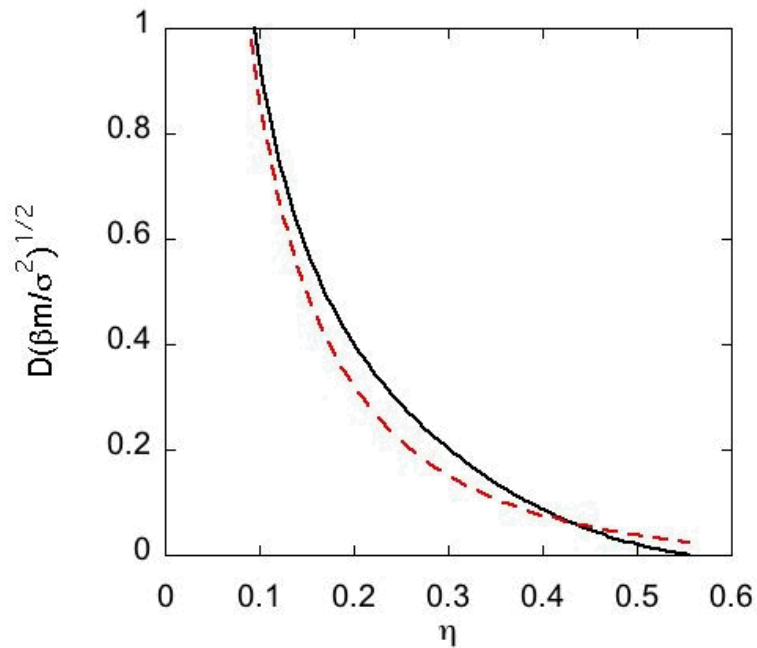
1. Force measurement
2. PDT
3. Fluxes and Profiles

Note, route 3 assumed:

$$D = \text{constant} = D(\rho_b)$$



Inhomogeneous Hard Sphere Fluid in gravity



9.6 Color Diffusion in a Slowly-Varying External Field: an alternate derivation

Color mixture, at "pure" equilibrium, and at "color" steady-state

$$\beta\mu_i(z) = \ln\rho_i(z) - \ln \langle e^{-\beta U_{t,i}(z)} \rangle + V_{ext}(z) \quad ; i = A, B \quad (53)$$

$$= \ln\rho_i(z) + \beta\tilde{\mu}^{ex}(z) + V_{ext}(z) \quad (54)$$

$$\beta\mu_i(z) = \ln x_i(z) + \beta\mu; \quad x_i \equiv \rho_i(z)/\rho(z) \quad (55)$$

$$j_i = -D_i(z)\rho_i \frac{\partial\beta\mu_i}{\partial z} \quad (56)$$

$$= -D_i(z)\frac{\partial\rho_i}{\partial z} - D_i\rho_i \frac{\partial\ln\rho}{\partial z} \quad (57)$$

given that the total local density, ρ , is also a function of z . Note that if ρ is a constant, this equation reduces to the simplest form of Fick's first law

$$-\frac{j_i}{D(z)\rho_i(z)} = \frac{\partial\ln\rho_i(z)}{\partial z} - \frac{\partial\ln\rho(z)}{\partial z} \quad ; i = A, B \quad (58)$$



Finding $D(z)$

9.7 The intrinsic chemical potential: pure fluid in gravity

$$\begin{aligned}\beta\mu &= \ln\rho(z) - \ln \langle e^{-\beta U_t(z)} \rangle + V_{ext}(z) \\ &= \underbrace{\ln\rho_i(z) + \beta\tilde{\mu}^{ex}(z)}_{\mu_{int}(z)} + V_{ext}(z)\end{aligned}$$

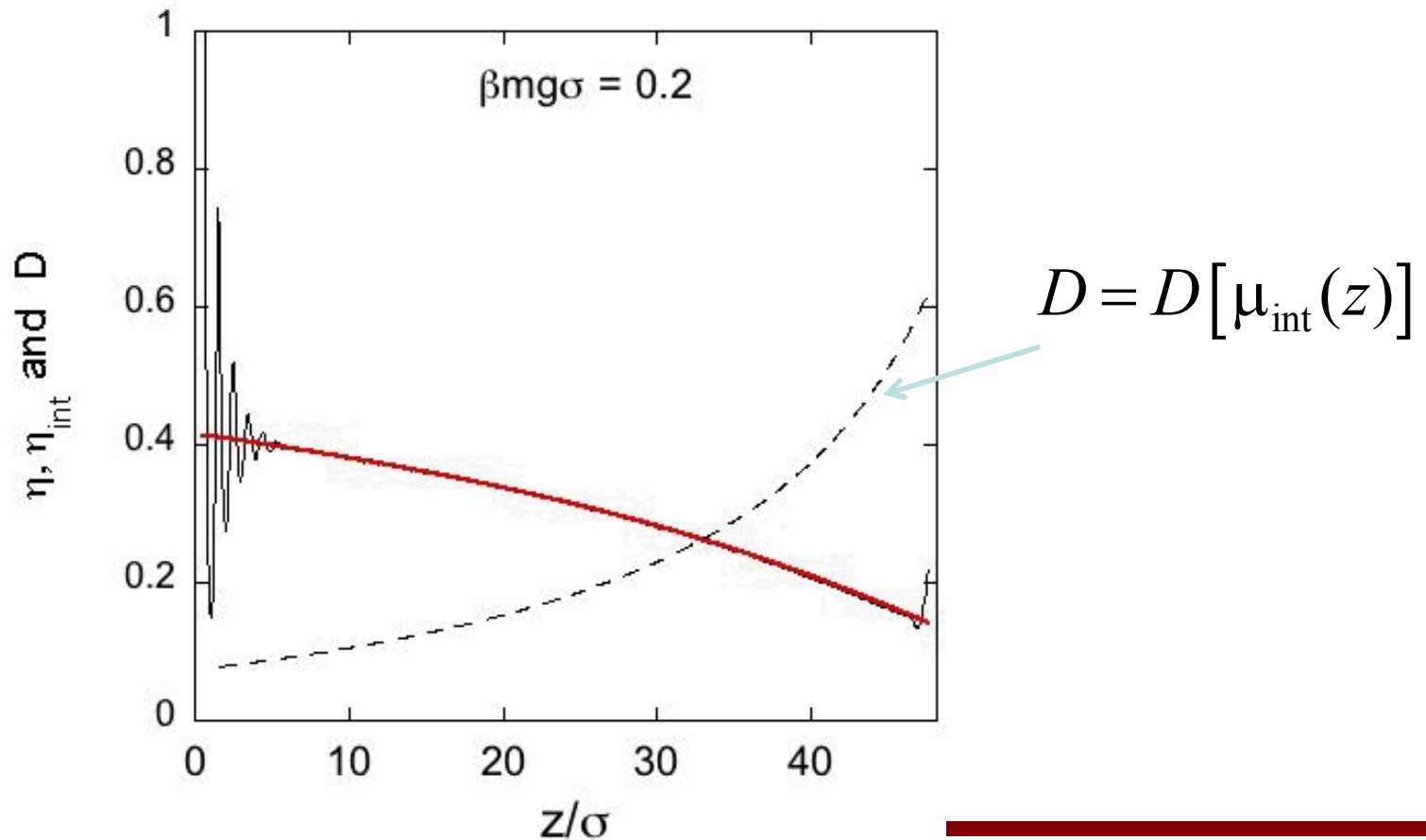
$$\begin{aligned}\mu_{int}(z) &\equiv \mu - V_{ext}(z) \\ &= \mu - mg(z - z_0); \quad \text{gravity}\end{aligned}$$

This determines the "state" of the local fluid. Now approximate

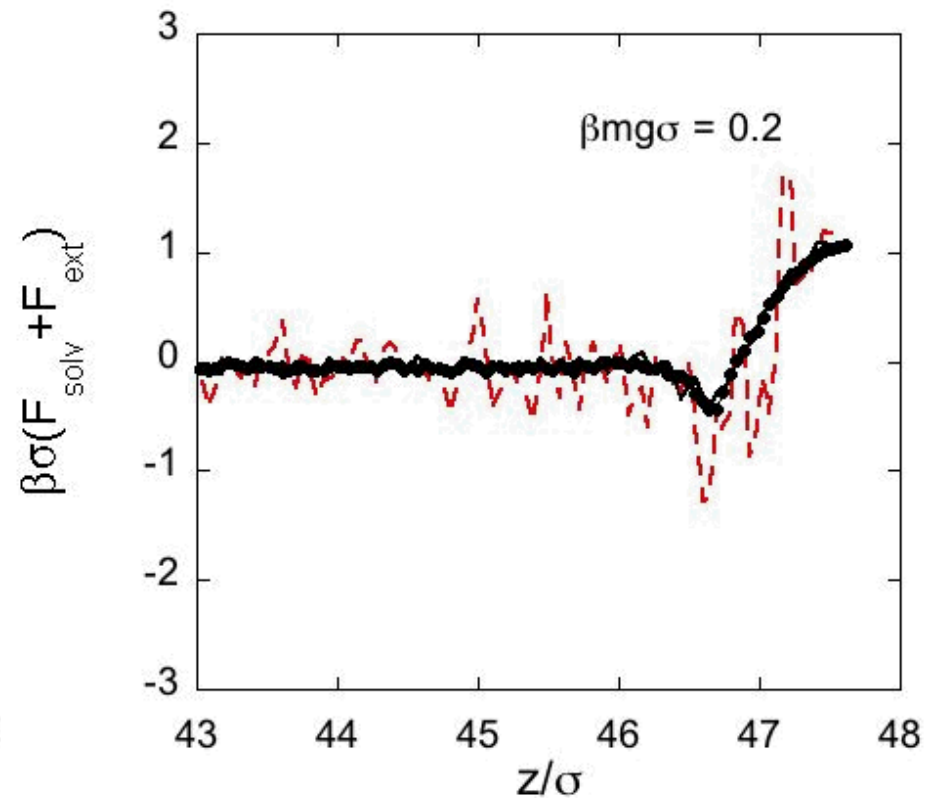
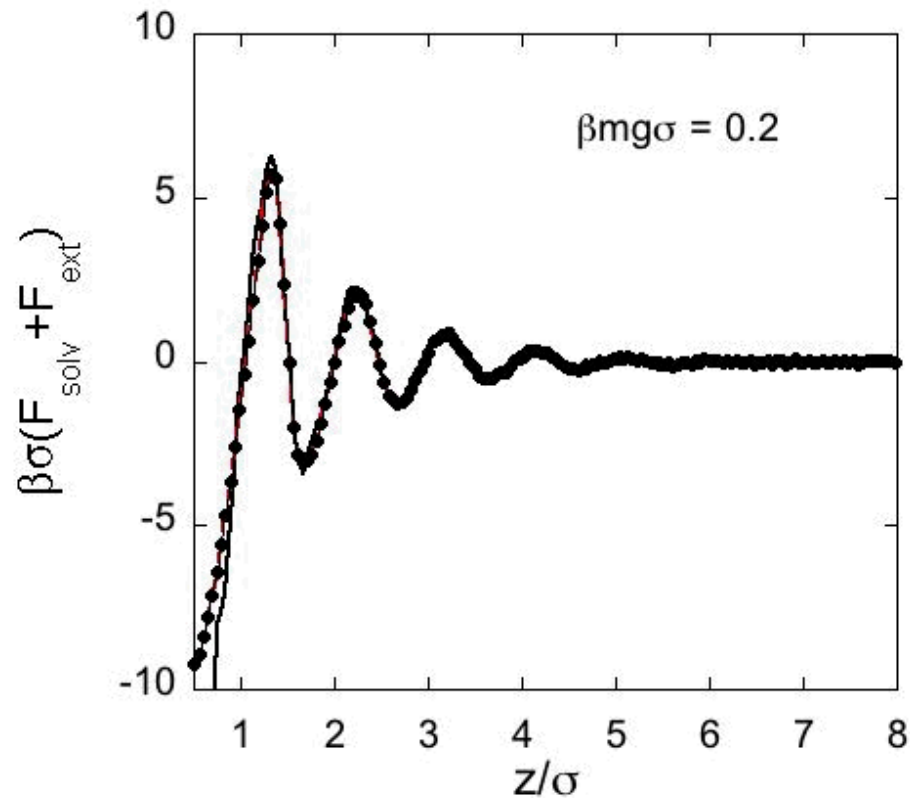
$$D = D[\mu_{int}(z)]$$

The intrinsic density profile

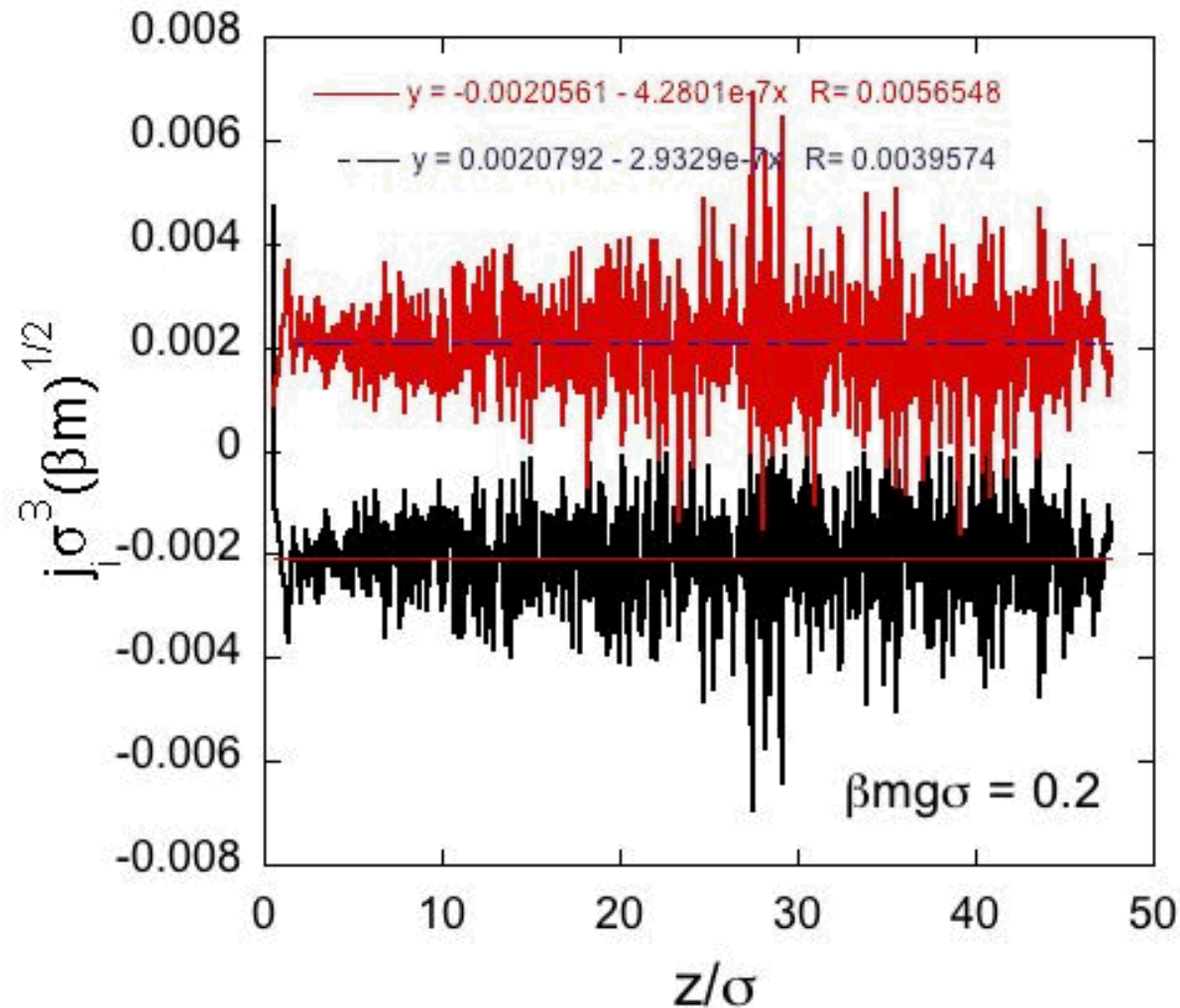
$$\mu_{\text{int}}(z) = \mu(z_0) - mg(z - z_0)$$



Checking $K(z)$ in gravity



Checking the fluxes, j_A & j_B



Concusion

For inhomogeneous fluids use

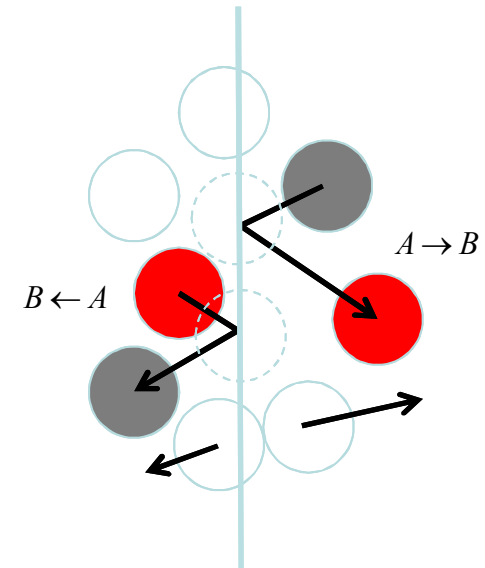
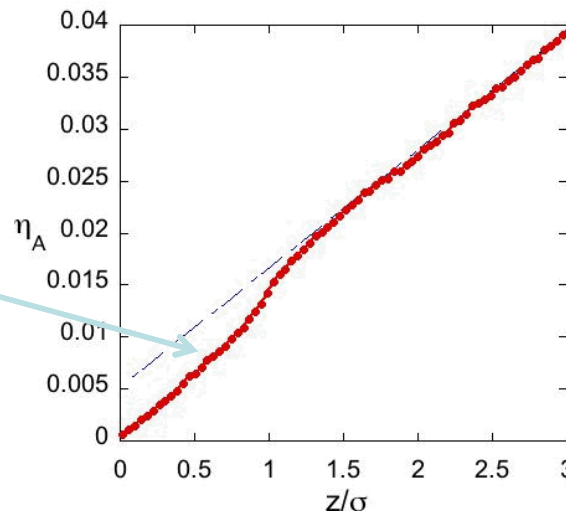
$$D(z) = D[\mu_{\text{int}}(z)]$$

and combine with the solvation force.

This also works in the presence of slowly-varying fields (e.g., gravity)

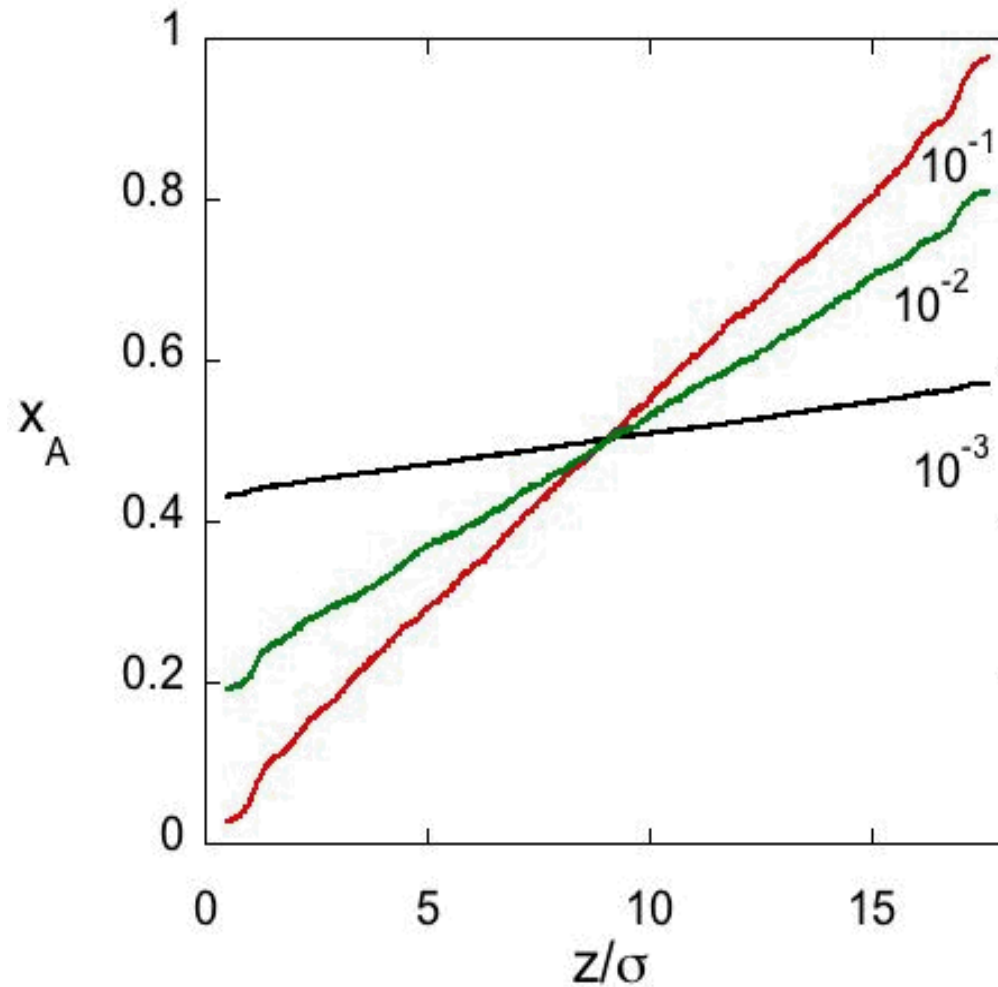
Question

What is this?

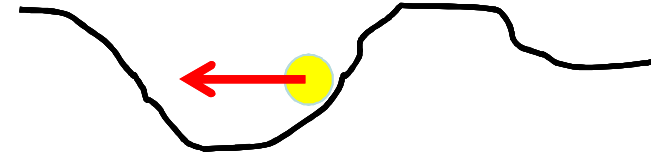


This work is supported by the DOE office of Basic Energy Sciences, Division of Material Sciences and Engineering. Sandia is a multiprogram laboratory operated by Sandia Corporation, a Lockheed- Martin Company, for the U.S. DOE under Contract No. DE-AC04-94AL85000.

The effect of Reaction Probabilities

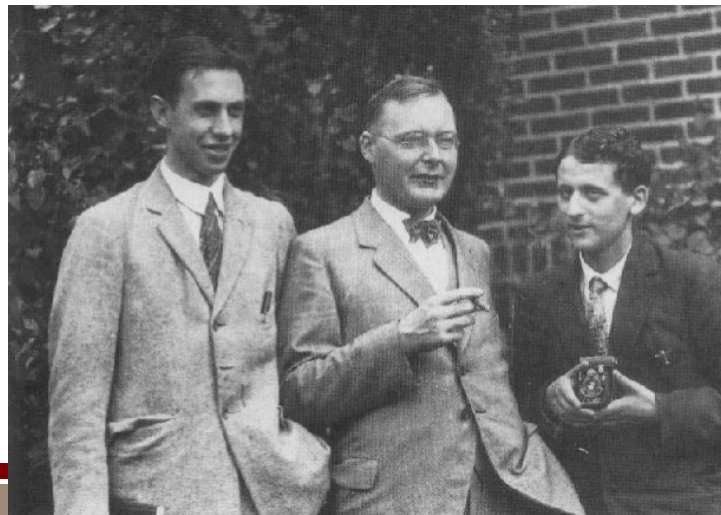


v. Smoluchowski (1906) Kramers(1940)-Klein(1922)

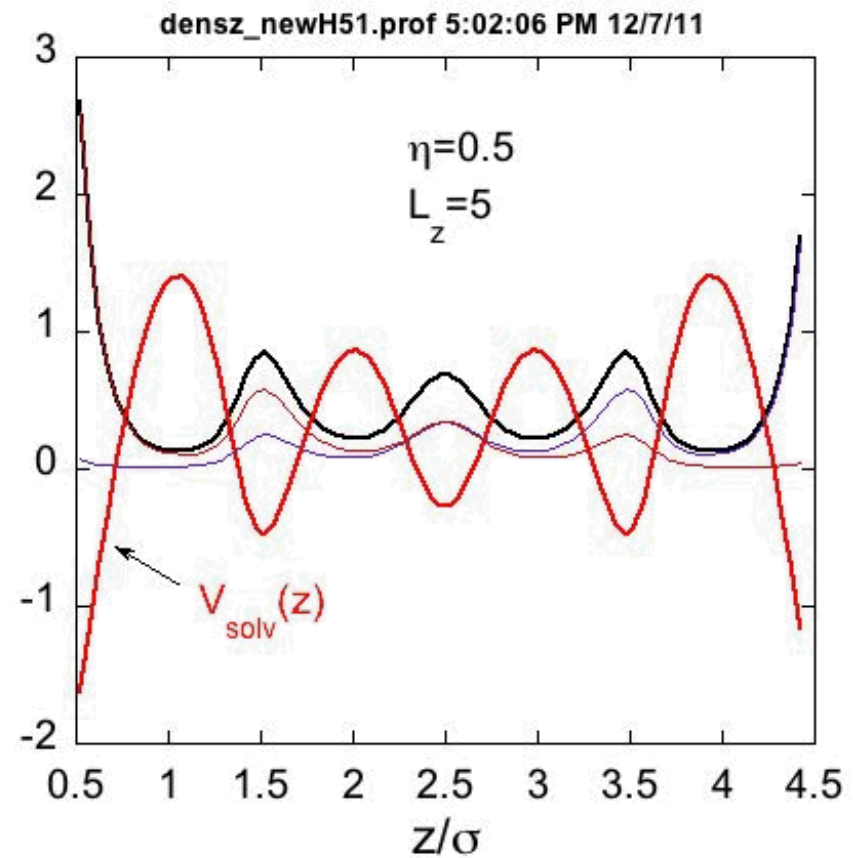
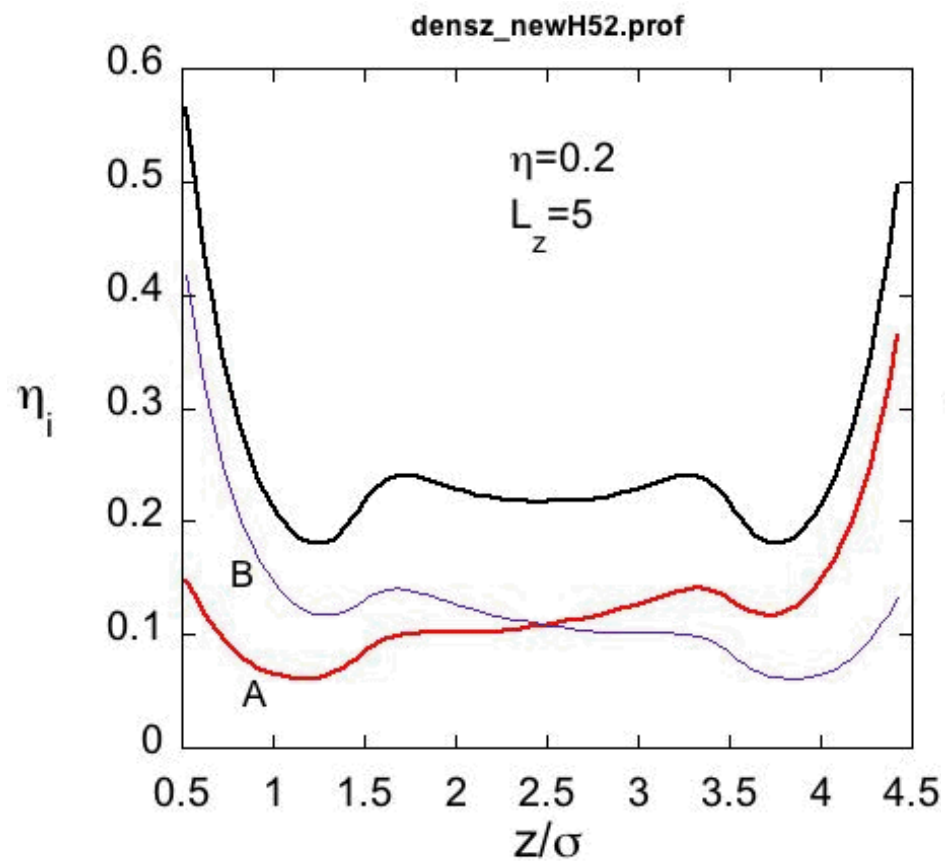


$$\dot{j}_i = -D \frac{\partial \rho_i}{\partial x} - D \frac{K(z) \rho_i}{L T}; \quad i = A, B$$

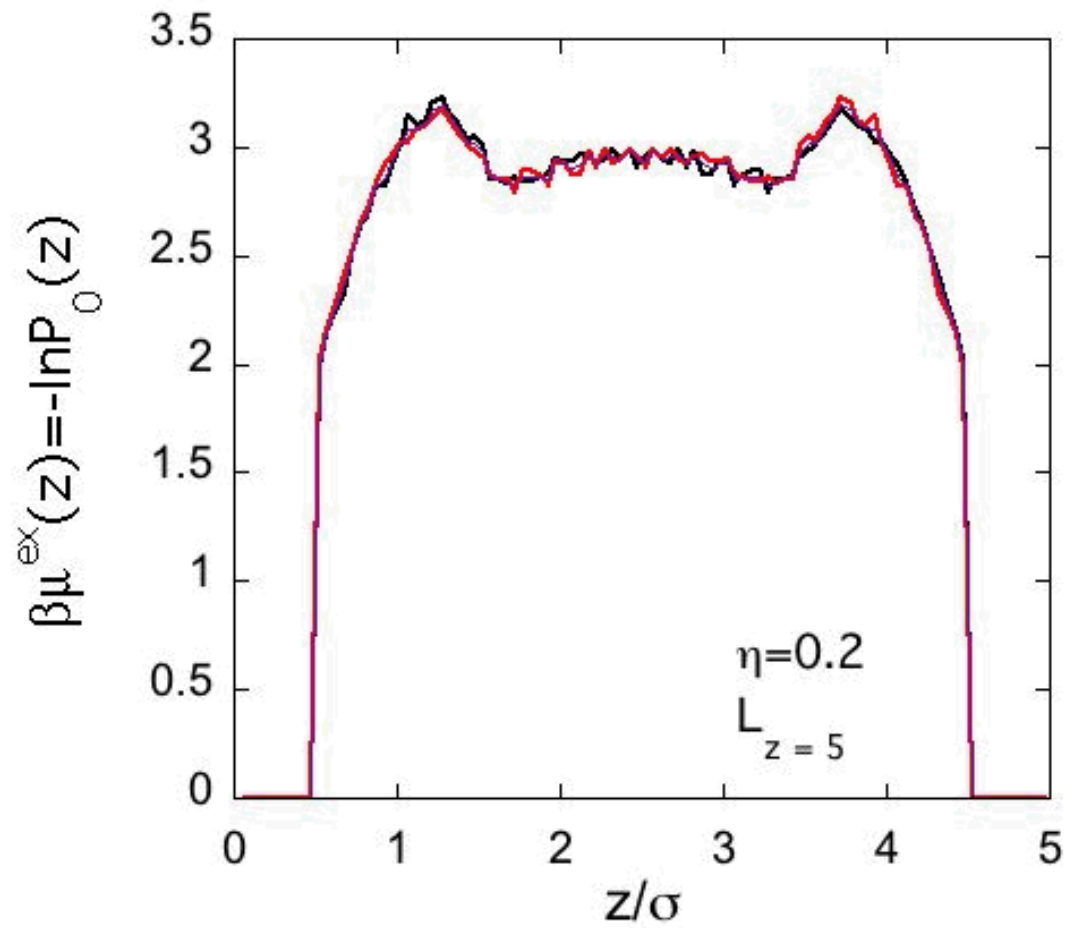
$$\frac{\partial P(x, x_0, t)}{\partial t} = D \frac{\partial^2 P(x, x_0, t)}{\partial^2 x^2} - \frac{\partial}{\partial x} \left[\frac{F(x)}{\xi} P(x, x_0, t) \right]$$



profiles



$$\beta\mu^{\text{ex}}(z)$$

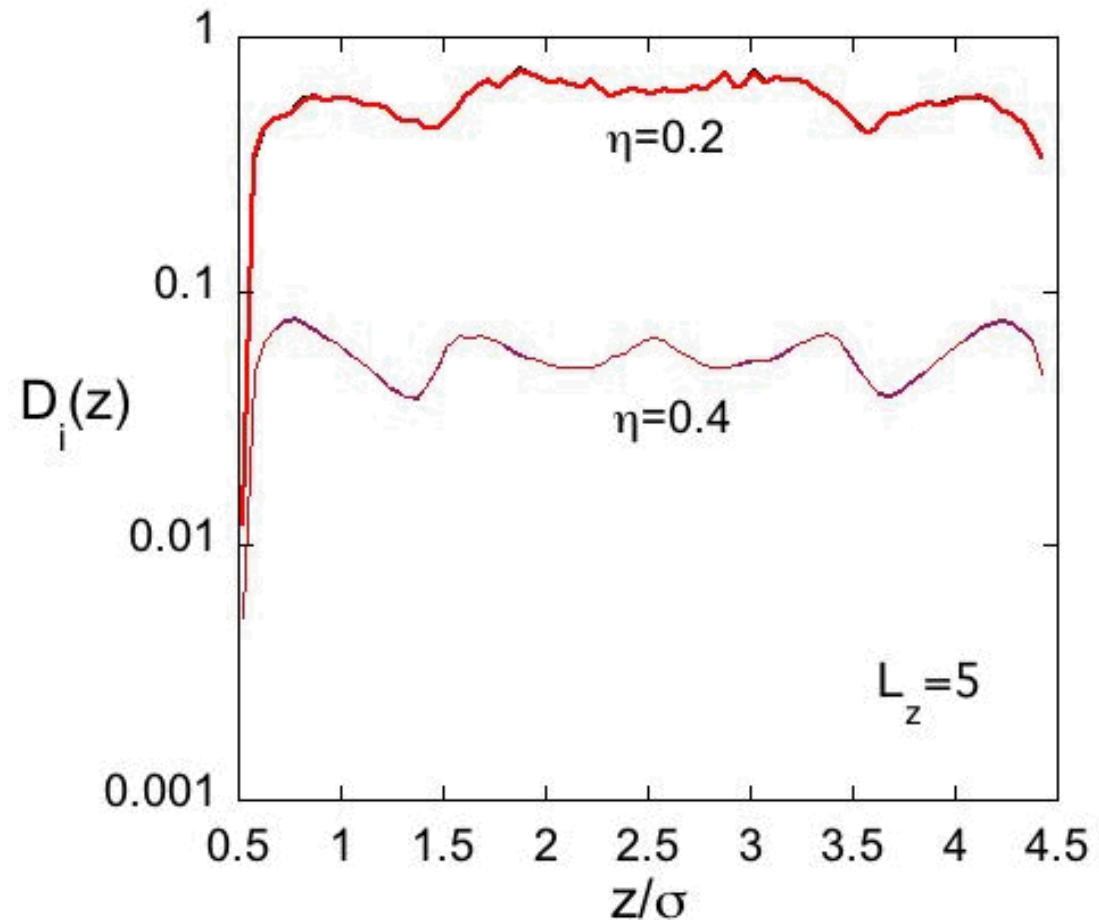


$$D(z) = -j_i \left[\rho_i \left(\frac{K(z)}{kT} + \frac{\partial \ln \rho_i}{\partial z} \right) \right]^{-1}; i = A, B$$

Ideal gas:

$$D(z) = -j_i \left[\left(\frac{\partial \rho_i}{\partial z} \right) \right]^{-1}; i = A, B$$

$D(z) = \text{constant}$

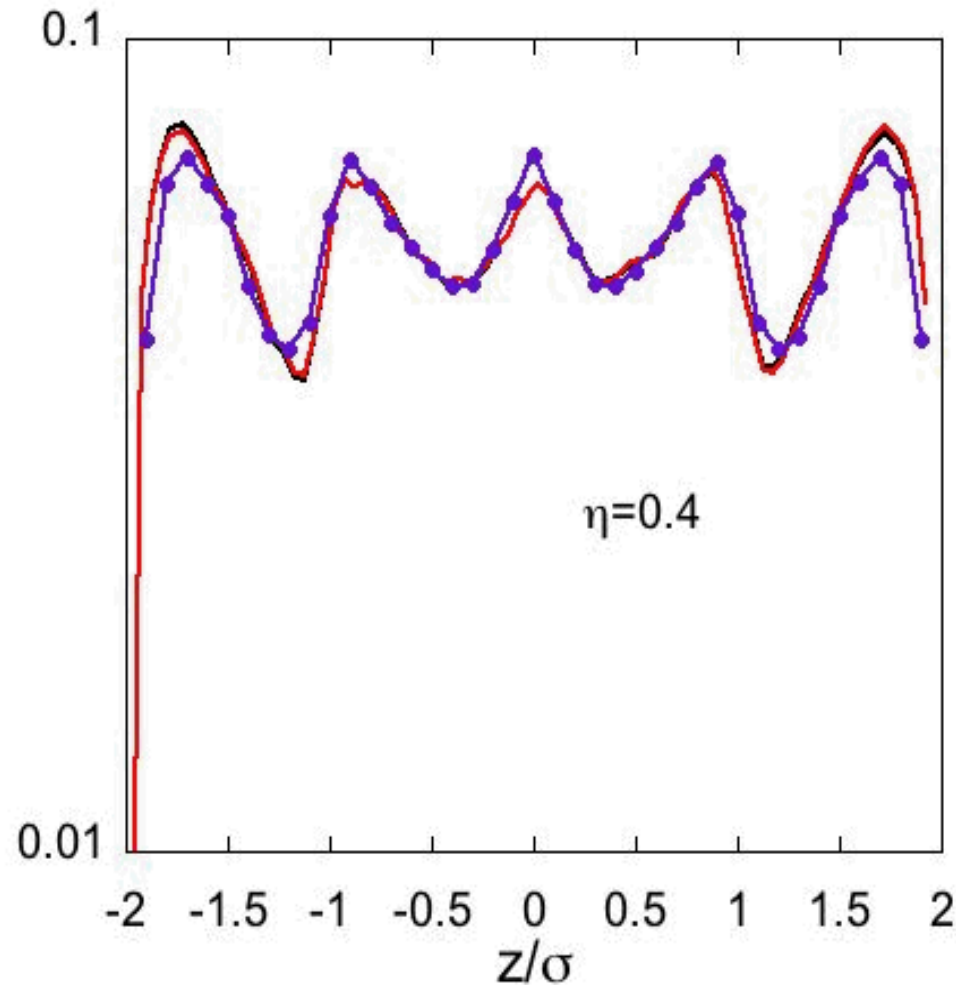


$$D(z) = -j_i \left[\rho_i \left(\frac{K(z)}{4\pi} + \frac{\partial \ln \rho_i}{\partial z} \right) \right]^{-1}; i = A, B$$

Ideal gas:

$$D(z) = -j_i \left[\left(\frac{\partial \rho_i}{\partial z} \right) \right]^{-1}$$

$D(z) = \text{constant}$



$\beta\mu^{\text{ex}}(z)$ in gravity

