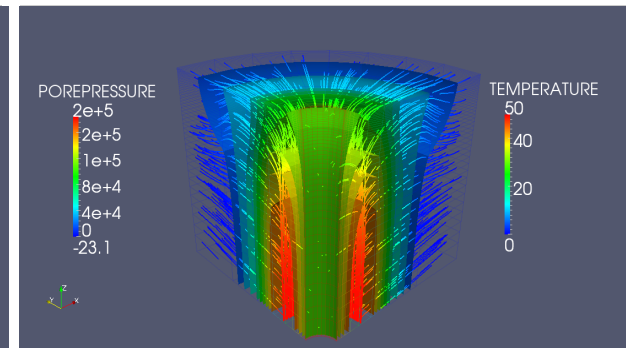
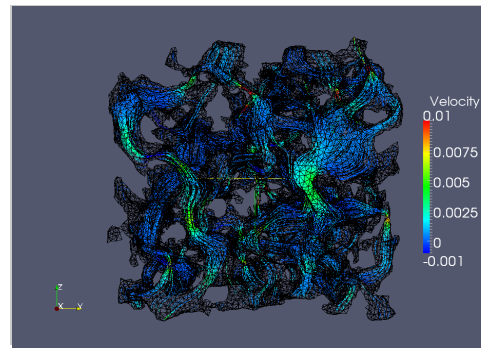
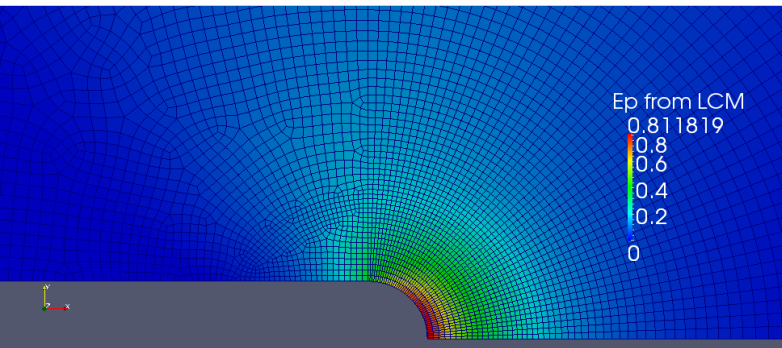


Exceptional service in the national interest



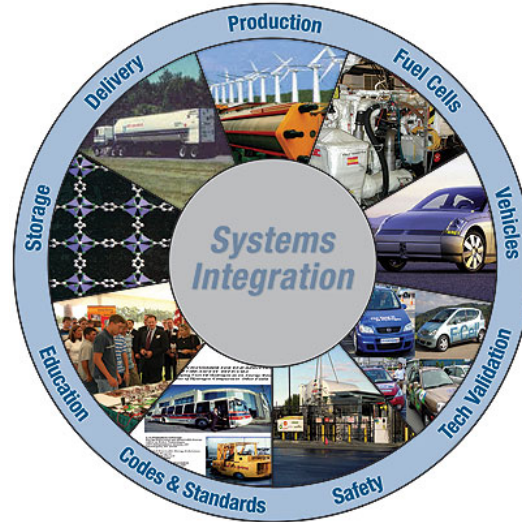
Concurrent Multi-physics Capacities in Laboratory of Computational Mechanics

Steve Sun, Jakob T. Ostien, James W. Foulk III

Motivations

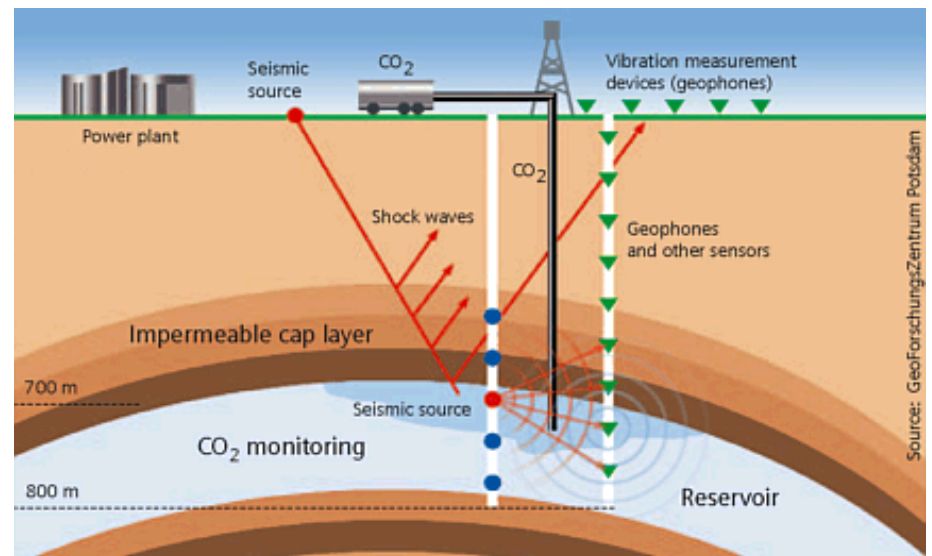
□ Hydrogen Transport

- Understand interaction between hydrogen gas and host material, in particular hydrogen enhanced fracture, embrittlement.
- Critical for hydrogen economy, especially for developing code and standards, life-cycle design..etc.



□ Poromechanics

- Understand interaction between mechanical and hydraulic coupling mechanism of solid-fluid mixture.
- Keys for CO₂ sequestration, oil recovery, geotechnical earthquake engineering, bone replacement, soft tissue modeling...etc.



Overview of Multiphysics Problems in Laboratory of Computational Mechanics

- Monolithic, implicit multiphysics capabilities currently available in LCM:
 1. Finite deformation thermo-mechanics (diffusion+ deformation)
 2. Finite deformation fully coupled hydrogen transport (advection-diffusion + deformation)
 3. Finite deformation poromechanics (diffusion + deformation)
 4. Finite deformation thermoporoplasticity (2 diffusion + deformation)
- Purpose:
 - Analyze and faithfully replicate physical processes which occur simultaneously interacting with each others.
- Challenge:
 - Additional physics requires additional equation + DOFs which may or may not have the same mathematical properties of the displacement field (e.g inf-sup condition).

Hydrogen Transport

Stabilized hydrogen diffusion-deformation K-field problem for low diffusive materials

Balance of Linear Momentum

$$\nabla^{\mathbf{X}} \cdot \mathbf{P}(\mathbf{F}, \bar{\mathbf{z}}, C_T) = 0$$

Concentration Sensitive Yield Function

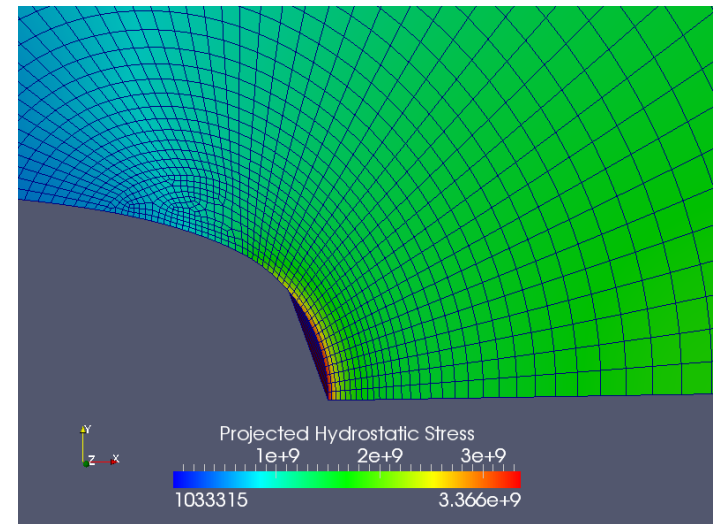
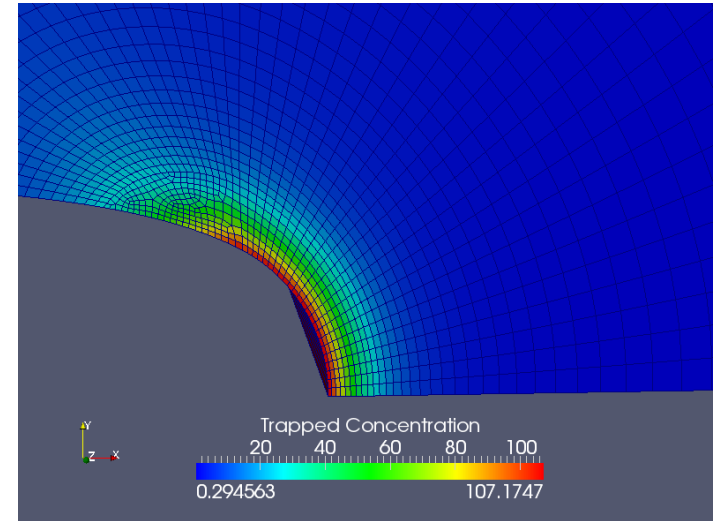
$$f(\boldsymbol{\tau}, \bar{\mathbf{z}}, C_T) = ||\text{dev}[\boldsymbol{\tau}]|| - \sqrt{\frac{2}{3}}[\sigma_Y(C_T) + K\alpha] \leq 0$$

Hydrogen Transport Theorem

$$D^* \dot{C}_L - \nabla^{\mathbf{X}} \cdot \mathbf{D}_L \nabla^{\mathbf{X}} C_L + \nabla^{\mathbf{X}} \cdot \frac{V_H}{RT} C_L \mathbf{D}_L \nabla^{\mathbf{X}} S_H + \theta_T \frac{dN_T}{d\epsilon_p} \dot{\epsilon}_p = 0$$

L2 Projection of Hydrostatic Stress

$$\int_{\mathcal{B}} N_a(\sigma_h - N_b \sigma_b) dV = 0$$



Challenges on Implementation of Hydrogen Transport Problem

1. Numerical instability may occur if (1) time step is too small and/or (2) diffusivity is too small (i.e. stainless steel) and (3) boundary layer due to the advection term is thinner than the side length of the finite element.
2. Hydrogen transport problem is highly nonlinear, thus require a consistent linearization to implicitly solve for solutions (i.e., **NUMEROUS** manual, mechanical derivations **EACH** time the problem is amended).
3. Volumetric Locking, which may occur under perfectly plastic response / isochoric deformation..etc.

Example

Say we want to find the linearized D^* (i.e., the first order term in the Taylor expansion). D^* depends on the trapped solvent, temperature, lattice concentration—but trapped solvent also depends on equivalent plastic strain and hence, displacement, yield stress...etc.

How do we derive consistent linearization with respect to all the dependent variables ?

Notice that we need to restart the entire process if we make any changes....

D^* – effective diffusion constant

$$D^* = \frac{N_T N_L}{K_T C_L^2} \left[\frac{1}{(1 + \frac{N_L}{K_T C_L})^2} \right]$$

$$C_T = \frac{N_T}{1 + \frac{N_L}{K_T C_L}}$$

$$\theta_T = \frac{1}{1 + \frac{N_L}{K_T C_L}}$$

$$C_T = \theta_T N_T$$

$$D_L = D_0 e^{-Q/RT}$$

$$N_L = N_A / V_M$$

$$K_T = e^{W_B/RT}$$

$$N_T = 10^{(A - B e^{-C \epsilon_p})} C_{L,0} = 3.47 \times 10^{-3} \text{ mol/m}^3 - \text{initial concentration}$$

$$\sigma_e = \sigma_0 + H \epsilon_p^m$$

$$\frac{\partial N_T}{\partial \epsilon_p} = 10^{(A - B e^{-C \epsilon_p})} \ln(10) B C e^{-C \epsilon_p} / N_A - \text{trapped solvent mol/m}^3$$

$D_0 = 2.0 \times 10^{-6} \text{ m}^2/\text{s}$ – pre-exponential factor

$Q = 6900 \text{ J/mol}$ – diffusion activation enthalpy

$R = 8.314 \text{ J/(molK)}$ – ideal gas constant

$T = 300 \text{ K}$ – room temperature

$D_L = 1.27 \times 10^{-8} \text{ m}^2/\text{s}$ (300K) – diffusion coefficient

$N_A = 6.0232 \times 10^{23} \text{ atoms/mol}$ – Avogadro's number

$V_M = 7.166 \times 10^{-6} \text{ m}^3/\text{mol}$ – molar volume of Fe

$N_L = 1.40 \times 10^5 \text{ solvent lattice mol/m}^3$

$W_B = 60 \times 10^3 \text{ J/mol}$ – trap binding energy

$K_T = 2.801 \times 10^{10}$ – equilibrium constant

$V_H = 2.0 \times 10^{-6} \text{ m}^3/\text{mol}$ ($\sim 300\text{K}$) – partial molar volume

$\Delta H_s = 28.6 \times 10^3 \text{ J/mol}$ – enthalpy of solution

$p_{H_2} = 1 \text{ atm}$ – hydrogen pressure

$C_{L,0} = 2.0 \times 10^{26} e^{-\Delta H_s/RT} \sqrt{p_{H_2}} \text{ atoms/m}^3$

$C_{L,0} = 3.47 \times 10^{-3} \text{ mol/m}^3$ – initial concentration

$N_T = 10^{(23.3 - 2.33 e^{-4.0 \epsilon_p})} / N_A$ – trapped solvent mol/m³

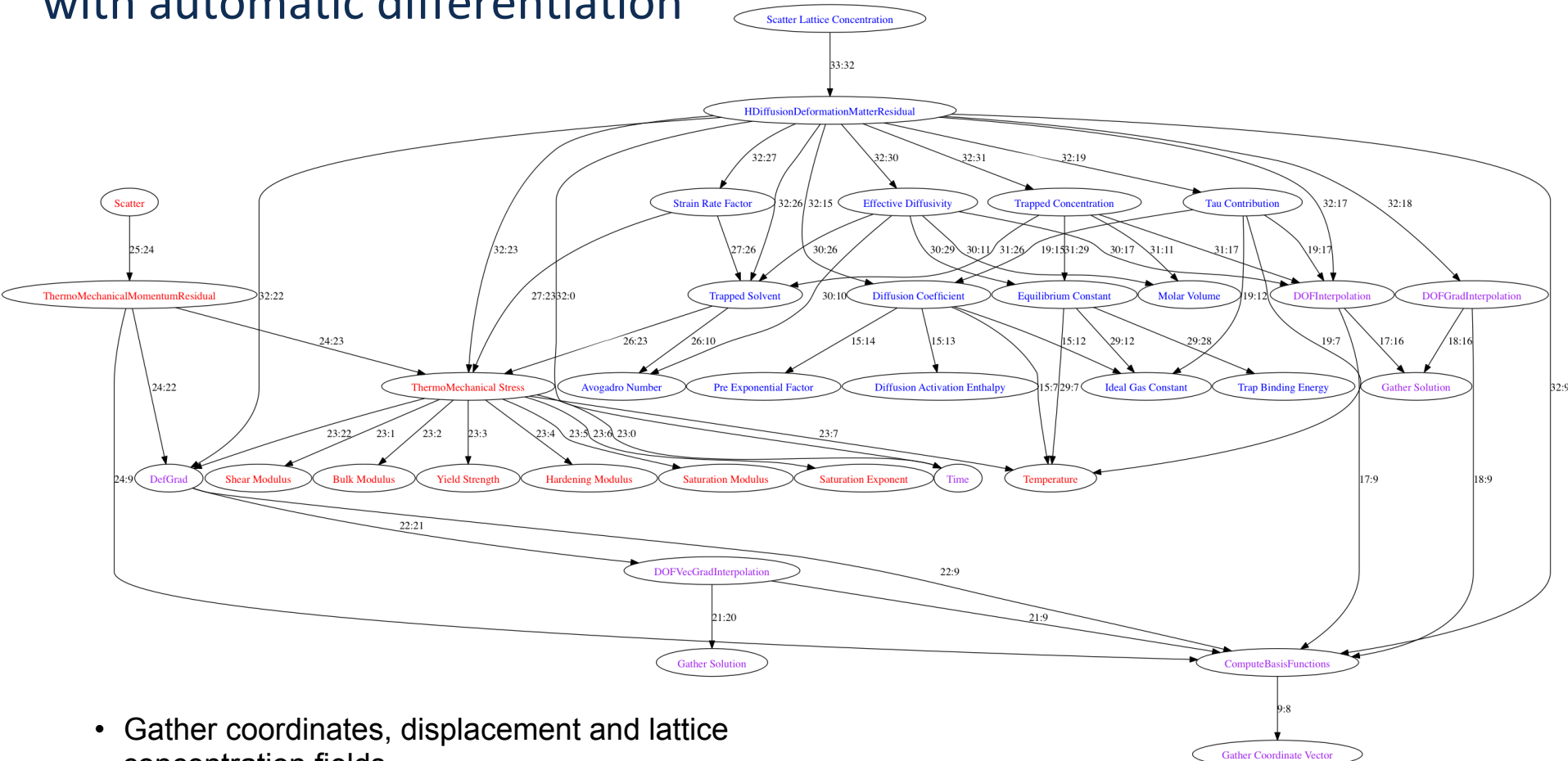
$N_{T,0} = 1.549 \times 10^{-3}$ – trapped solvent mol/m³

$\theta_{T,0} = 0.9986$

$C_{T,0} = 1.547 \times 10^{-3} \text{ mol/m}^3$ – initial concentration

$C_{total,0} = C_{L,0} + C_{T,0} = 5.017 \times 10^{-3} \text{ mol/m}^3$

Implementation of Hydrogen Diffusion-Mechanics Problem with automatic differentiation



- Gather coordinates, displacement and lattice concentration fields
- Interpolate fields and gradients to integration points
- Chain together Evaluators to compute Momentum and Conservation of Hydrogen Residuals
- Scatter back to the global system of equations

Blue = Hydrogen Transport
 Red = Solid Mechanics (J2 Plasticity)
 Purple = coupled terms

Combined F-bar formulation

Isochoric-volumetric split

$$\mathbf{F} = \mathbf{F}_{\text{vol}} \cdot \mathbf{F}_{\text{iso}}$$

Replacing volumetric split with assumed term

$$\bar{\mathbf{F}} = \bar{J}^{1/3} \mathbf{F}_{\text{iso}} = \bar{J}^{1/3} J^{-1/3} \mathbf{F}$$

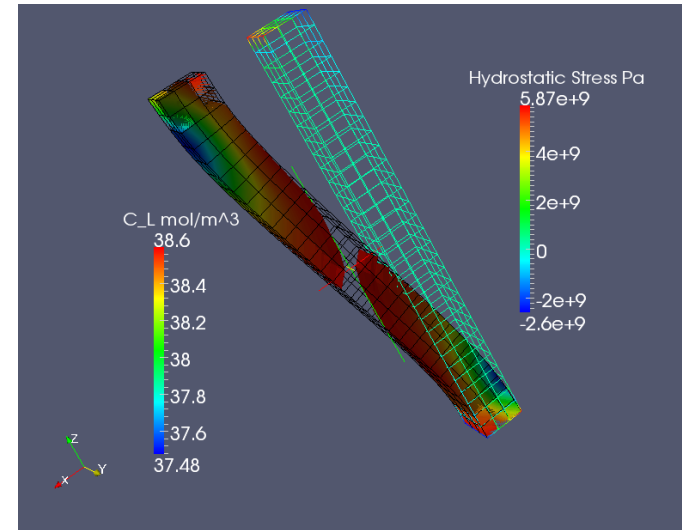
Classical Combined F-bar approach

$$\tilde{\mathbf{F}} = \alpha \mathbf{F} - (1 - \alpha) \bar{\mathbf{F}}.$$

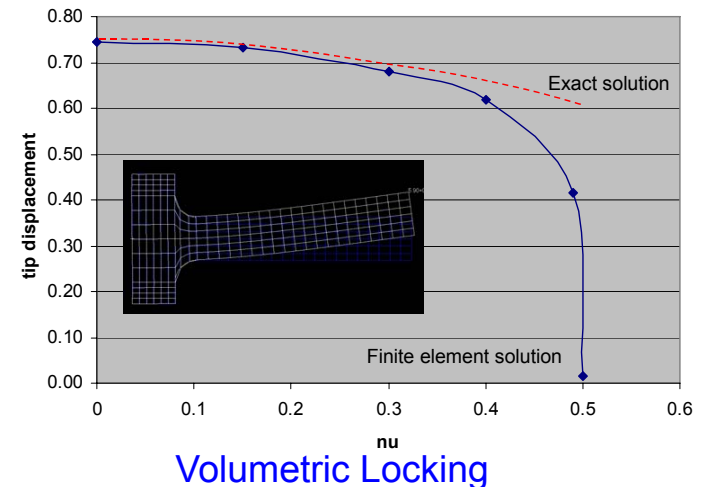
Current Approach via Lie algebra

$$J = J^e J^p J^C \quad ; \quad J^C = 1 + \lambda(C - C_o)$$

$$\tilde{J}(X) = \exp \left(\frac{1 - \alpha}{V_{\mathcal{B}^e}} \int_{\mathcal{B}^e} \log J(X) \, dV + \alpha \log J(X) \right)$$




Concentration prediction affected by locking





Stabilization at small time limit


Hydrogen Transport Theorem

$$D^* \dot{C}_L - \nabla^{\mathbf{x}} \cdot \mathbf{D}_L \nabla^{\mathbf{x}} C_L + \nabla^{\mathbf{x}} \cdot \frac{V_H}{RT} C_L \mathbf{D}_L \nabla^{\mathbf{x}} S_H + \theta_T \frac{dN_T}{d\epsilon_p} \dot{\epsilon}_p = 0$$


 transient
term


 Hydrogen diffusion
term


 Advection coupling
term


 Plastic strain nonlinear
coupling
term

- Spurious oscillations may occur when
 - D^* is large, which means local rate of change dominates
 - The mesh size h is large (relative to the advection and diffusion length scale)
 - The time step is small (relative to the advection and diffusion time scale).
 - Notice that Peclet number measures whether advection of diffusion is more important, but did not tell much about the transient term!
- Examples of stabilization scheme
 - Petrov-Galerkin/streamline upwind method (Hughes, 1978, Johnson, 1984)
 - Space-time finite incremental calculus method (Onate and Manzan 2000)
 - SUPG with adaptive stabilization parameters (Tezduyar 2003)
 - Spurious oscillations at layers diminishing method (Volker and Schmeyer, 2008).
 - Artificial diffusivity (Onate and Manzan, 2000).

Stabilization for hydrogen convection-diffusion problem

- For transient problem, stability criteria must be satisfied for the pair of time step **and** element size h used in simulations. γ is the parameter for backward Euler time integrator (Harari, 2004).

$$\frac{h^2 D^*}{6\gamma D_L \Delta t} < 1$$

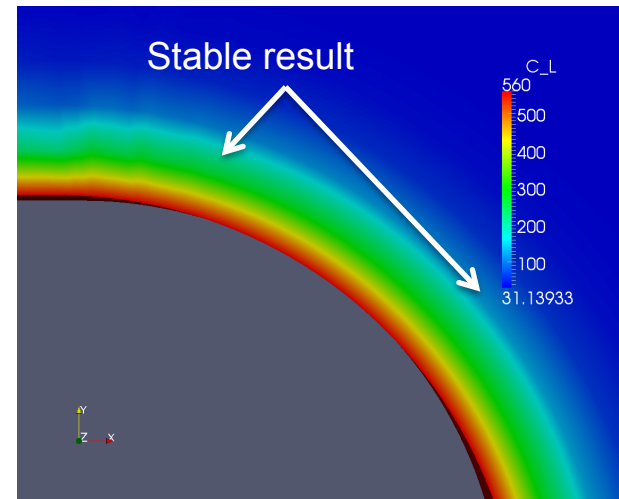
- Meanwhile, stability of the steady solution can be predicted by the Peclet number. If Pe is less than 1 and D^*/D_L is reasonably close to h^2/dt

$$Pe = \frac{V_H}{RT} |\nabla^{\mathbf{X}} S_H| h < 1$$

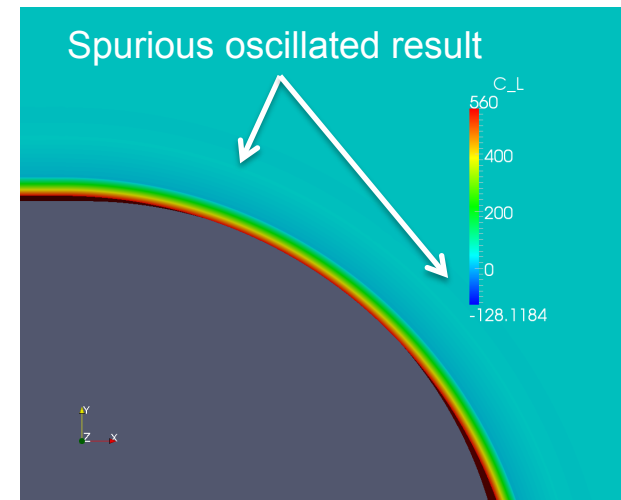
- add stabilization term to penalize the deficiency

$$R_{stab} = \int_{\Omega} \left[\tau (\eta^h - \frac{1}{V_{\Omega}} \int_{\Omega} \eta^h d\Omega) (\dot{C}_L^h - \frac{1}{V_{\Omega}} \int_{\Omega} \dot{C}_L^h d\Omega) \right] d\Omega$$

$$R_{stab} = \int_{\Omega} \tau \nabla^{\mathbf{X}} \eta^h \nabla^{\mathbf{X}} \dot{C}_L^h d\Omega$$

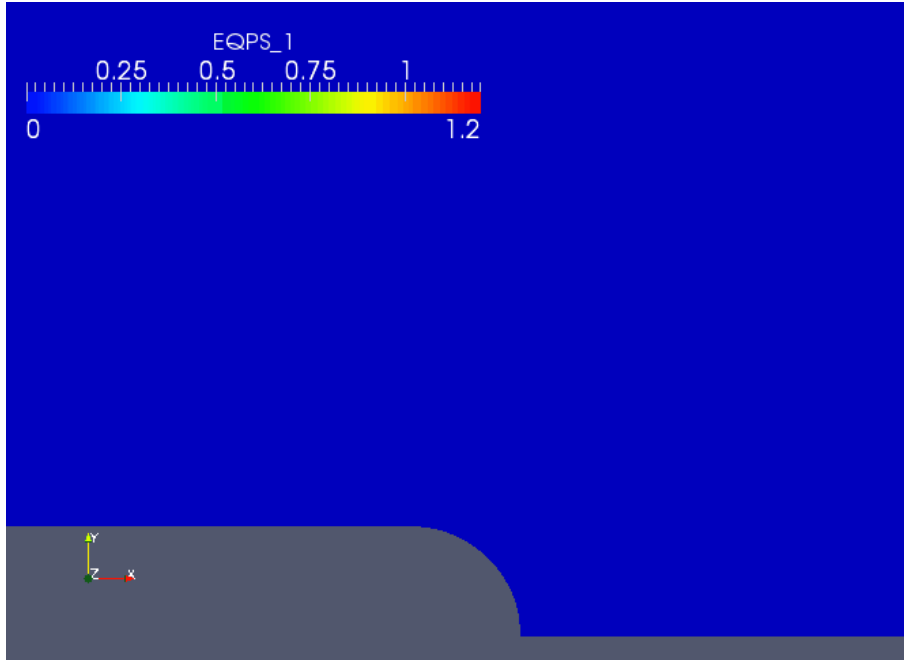


Stabilized Implicit Galerkin Formulation

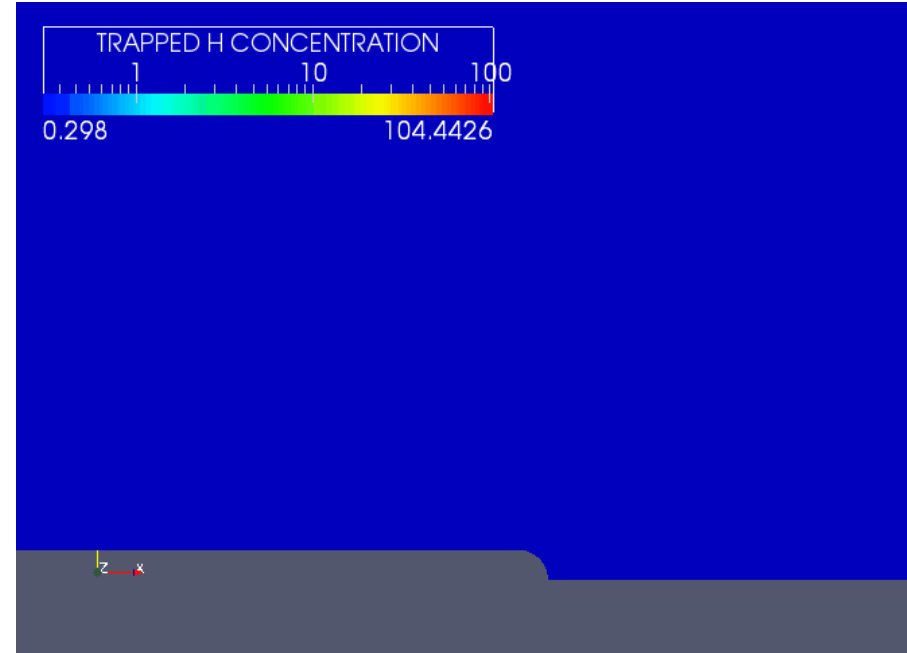


Standard Implicit Galerkin Formulation

K-field blunt tip problem



Equivalent Plastic Strain

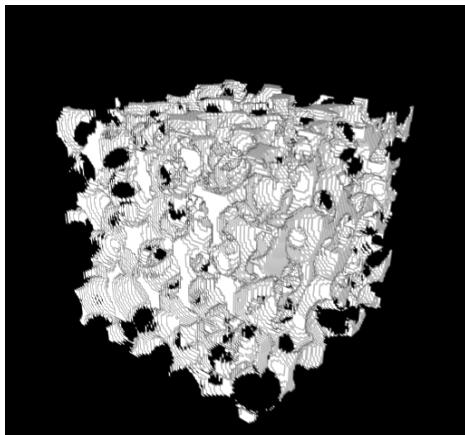


Trapped Hydrogen Concentration

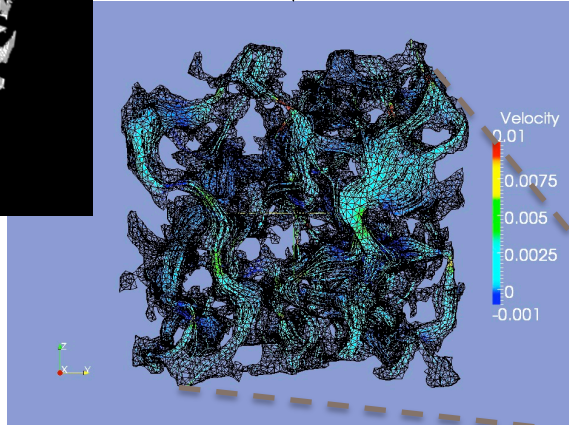
Conclusion

- We have implemented a fully implicit hydrogen transport model in the Laboratory of Computational Mechanics package with the following desirable features
 - Open-source, open for collaborations.
 - Capacity to conduct fully coupled, fully implicit simulations with stabilization scheme.
 - Automatic differentiation, which makes it fast and easy to make amendment to existing model and eliminate chance of making error in derivation
 - Stabilization scheme available to handle material with extremely low diffusivity, thin boundary layer...etc.
 - A L2 projection scheme to obtain a C^0 -continuous stress gradient term that enables the advection term to be correctly modeled without introducing errors during the extrapolation process.

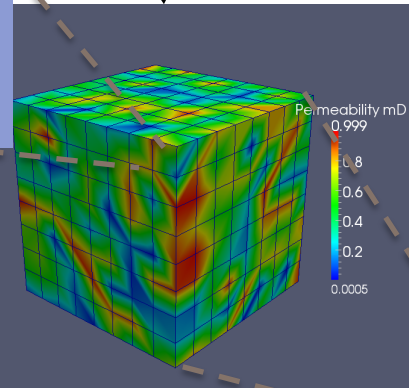
Poromechanics



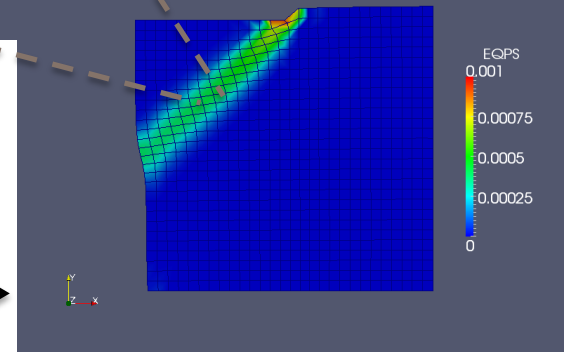
Level set scheme to
obtain signed distance function,
iso-surface and mesh



Up-scaling hydraulic
and mechanical
parameters by solving
inverse problem



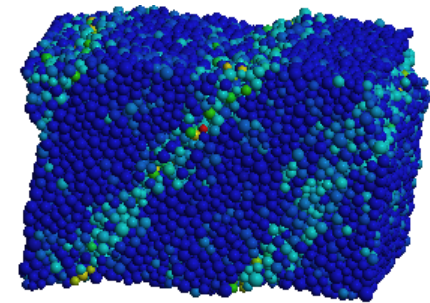
Finite Element models
incorporating Spatial
Variability and length
scale



Meso-Scale homogenization

Arlequin domain coupling

Field-scale Boundary Value Problem



Grain Scale Discrete
Element Simulation

Pore-scale Calculation

1nm

1mm

1cm

1m

U-P Formulation of Poromechanics Problem

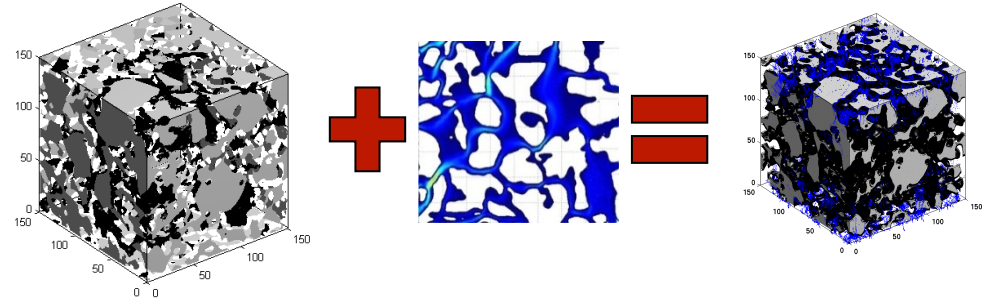
Balance of Linear Momentum

$$\nabla^{\mathbf{X}} \cdot \mathbf{P} + J(\rho^s + \rho^f)\mathbf{G} = \mathbf{0}$$

Effective Stress Concept

$$\mathbf{P}(\mathbf{F}, z, p^f) = \mathbf{P}'(\mathbf{F}, z) - J\left(1 - \frac{K}{K_s}\right)p^f \mathbf{F}^{-T}$$

Solid skeleton Solid skeleton POROUS MEDIA



Balance of Mass

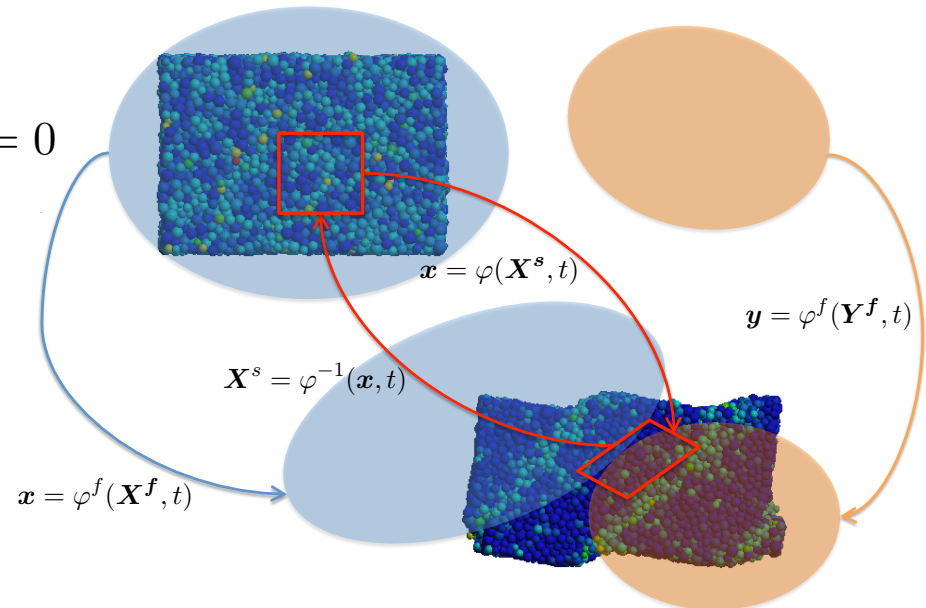
$$\frac{DB}{Dt} \left(\log J + \frac{p^f}{K_s} \right) + \frac{B}{J} \frac{DJ}{Dt} + \frac{1}{M} \frac{Dp^f}{Dt} + \nabla^{\mathbf{X}} \cdot \mathbf{Q} = 0$$

Darcy's Law

$$\frac{1}{\rho_f} \mathbf{W} = \mathbf{K} \cdot (-\nabla^{\mathbf{X}} p^f + \rho_f \mathbf{F}^T \cdot \mathbf{G})$$

$$\mathbf{K} = J \mathbf{F}^{-1} \cdot \mathbf{k} \cdot \mathbf{F}^{-T}$$

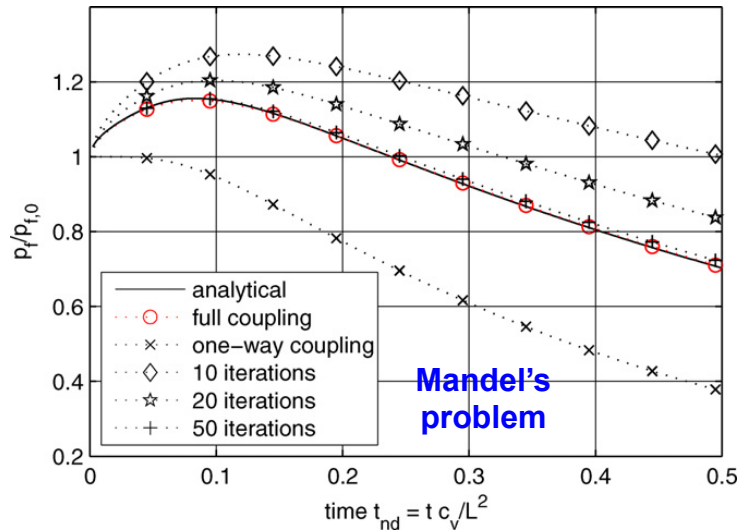
$$\mathbf{Q} = (1/\rho_f) \mathbf{W}$$



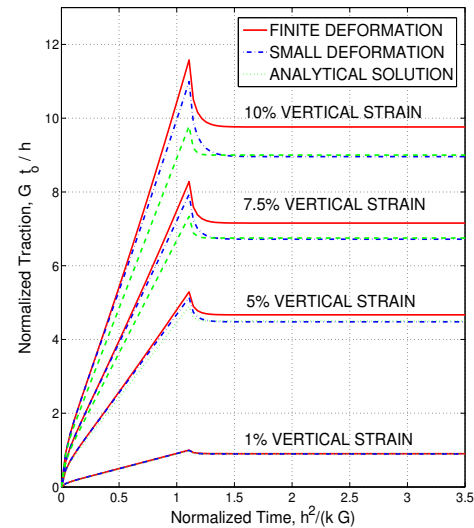
Challenges on Implementation of Poromechanics Problem

1. Preserve Mandel-Cryer effect (pore pressure changes non-monotonically in time upon subjected to mechanical loading).
2. Only a selection of finite element space satisfying inf-sup condition and equal-order finite element discretization does not.
3. Volumetric Locking may occur under perfectly plastic response / isochoric deformation..etc.
4. Often requires addition field equations to capture essential physical processes (heat transport, chemo-mechanics...etc)

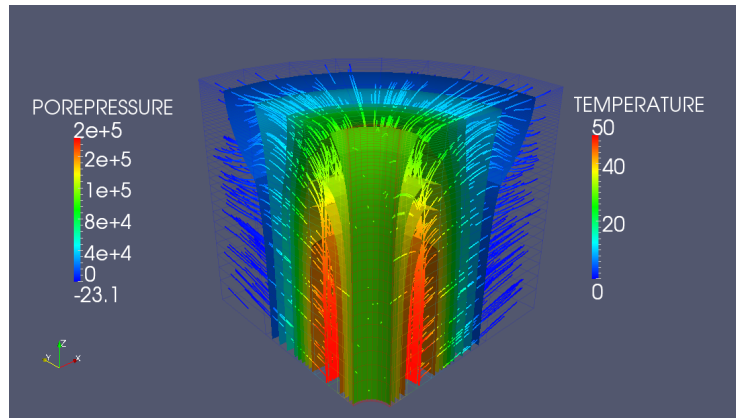
Key Features available in Poromechanics Models



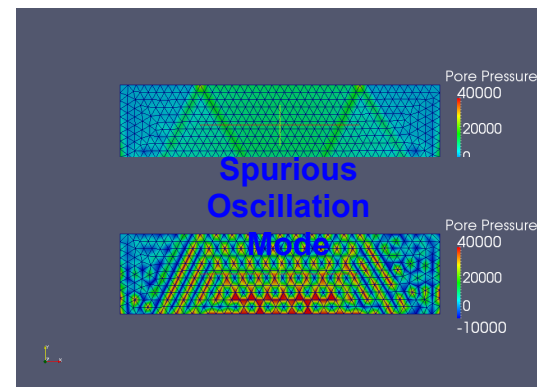
Monolithic Coupled Hydromechanics



Total Lagrangian Finite Strain Formulation



Thermomechanical effect of mixture



Stabilized Equal-order Discretization

Stabilization for Equal Order Mixed Finite Element

Inf-sup Condition (**not satisfied**)

$$\sup_{\mathbf{u}_h \in V_{\mathbf{u}}^h, \mathbf{u}_h \neq 0} \frac{\int p_h \nabla^{\mathbf{x}} \cdot \mathbf{u}_h d\Omega}{\|\mathbf{u}_h\|_1} \geq \gamma \|p^h\|_0, \forall p^h \in V_p^h$$

Weaker Inf-sup Condition (**still satisfied**)

$$\sup_{\mathbf{u}_h \in V_{\mathbf{u}}^h, \mathbf{u}_h \neq 0} \frac{\int p_h \nabla^{\mathbf{x}} \cdot \mathbf{u}_h d\Omega}{\|\mathbf{u}_h\|_1} \geq \gamma_1 \|p^h\|_0 - \gamma_2 h \|\nabla^{\mathbf{x}} p^h\|_0, \forall p^h \in V_p^h$$

Pressure Projection Stabilization

- add stabilization term to penalize the deficiency in the displacement-pore pressure approximation pair.
- We use an adaptive scheme to turn on/off stabilization such that no excess diffusion is introduced.

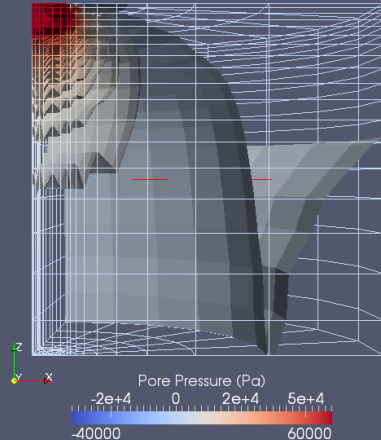
$$R_{stab} = \int_{\Omega} \left[\tau \left(\eta^h - \frac{1}{V_{\Omega}} \int_{\Omega} \eta^h d\Omega \right) \left(p^h - \frac{1}{V_{\Omega}} \int_{\Omega} p^h d\Omega \right) \right] d\Omega$$

↑
Stabilization
parameter

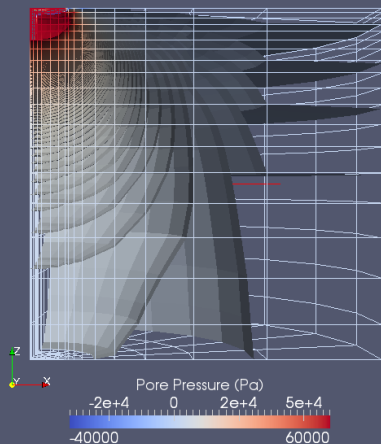
↑
Interpolated pore
pressure field

↑
Projected pore
pressure field

No stabilization



With stabilization



Stabilized F-bar Poromechanics with Monolithic Coupling at finite strain

Balance Law

$$\nabla^{\mathbf{X}} \cdot \mathbf{P} + J(\rho^s + \rho^f) \mathbf{G} = 0$$

$$\frac{DB}{Dt} \left(\log J + \frac{p^f}{K_s} \right) + \frac{B}{J} \frac{DJ}{Dt} + \frac{1}{M} \frac{Dp^f}{Dt} + \nabla^{\mathbf{X}} \cdot \mathbf{Q} = 0$$

Regularization Energy

$$W^{pen}(p_\tau^f) = \frac{1}{2} \sum_{K \in \Omega} \int_K (p_\tau^f - \Pi p_\tau^f) \frac{\gamma}{M} (p_\tau^f - \Pi p_\tau^f) dV$$

$$W^{pen}(p_\tau^f) = \frac{1}{2} \sum_{K \in \Omega} h_K \int_K \nabla^{\mathbf{X}} p_\tau^f \cdot \beta \mathbf{K} \nabla^{\mathbf{X}} p_\tau^f dV$$

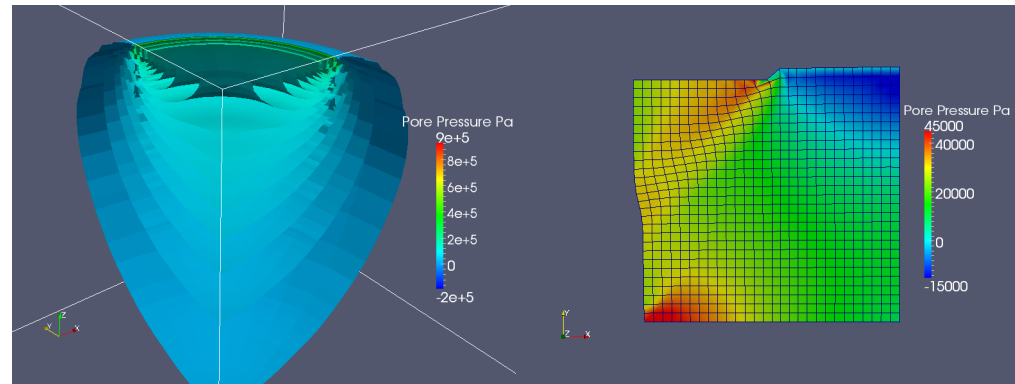
Harari's analysis

$$\gamma = \left(2 - 6 \frac{\vartheta c \Delta t}{h^2} \right) \left(\frac{1}{2} + \frac{1}{2} \tanh \left(2 - 12 \frac{\vartheta c \Delta t}{h^2} \right) \right)$$

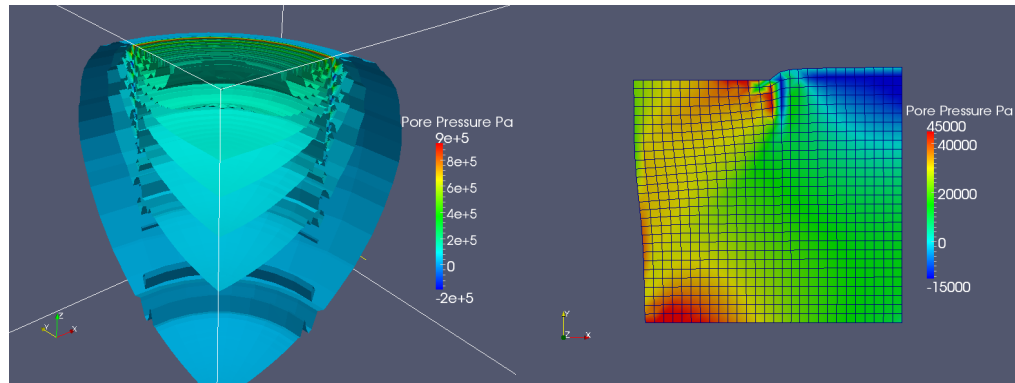
Weighted F-bar

$$\tilde{\mathbf{F}} = \tilde{J}^{1/3} J^{-1/3} \mathbf{F}$$

$$\tilde{J}(X) = \exp \left(\frac{1 - \alpha}{V_{\mathcal{B}^e}} \int_{\mathcal{B}^e} \log J(X) dV + \alpha \log J(X) \right)$$



Stabilized F-bar Mixed FEM

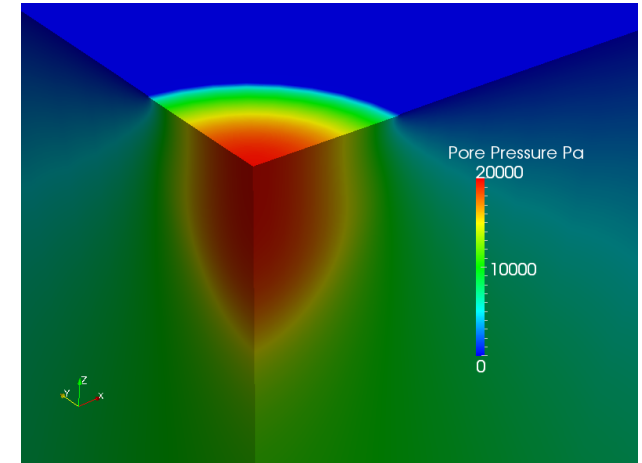
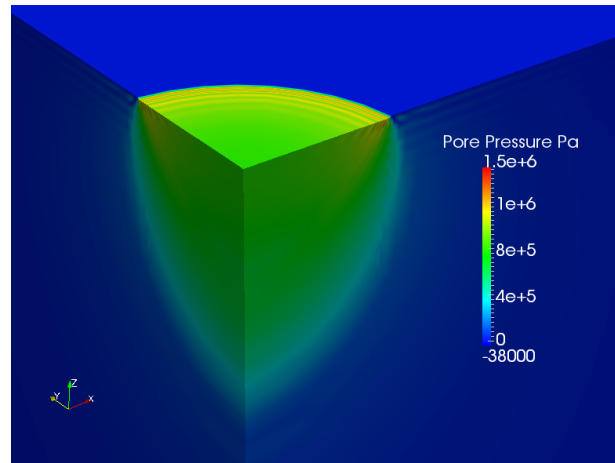


Standard Galerkin Method

Footing placed on an elastoplastic porous medium

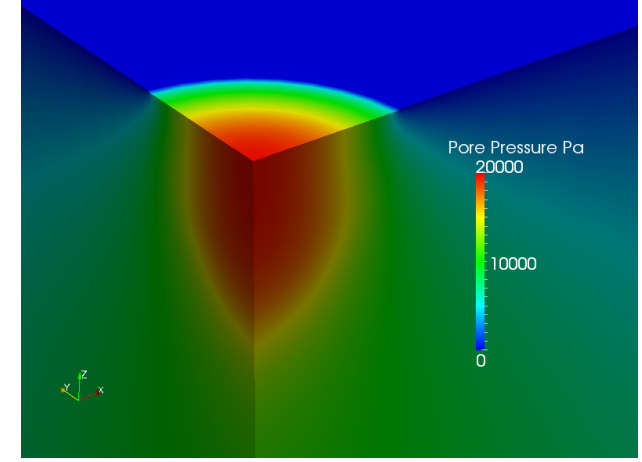
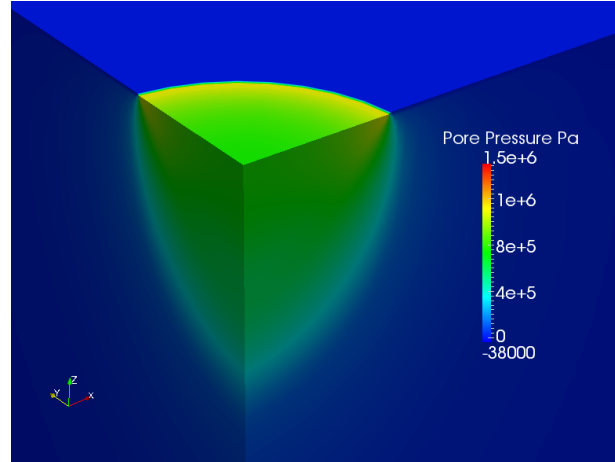
Without Stabilization

In low diffusivity case, the stabilization scheme is able to eliminate spurious oscillation.



With Stabilization

In high diffusivity case, the stabilization scheme does not introduce extra diffusivity that cause error.

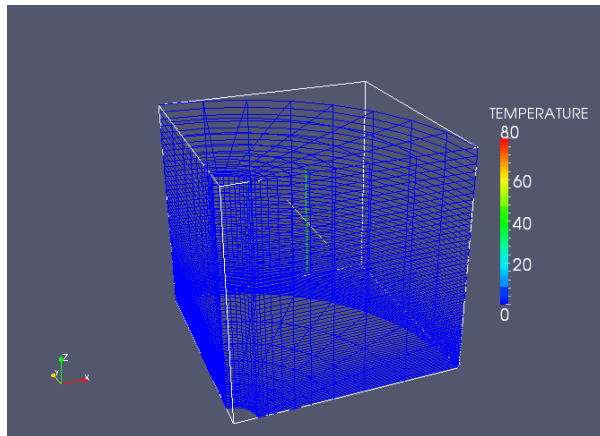


Low diffusivity case

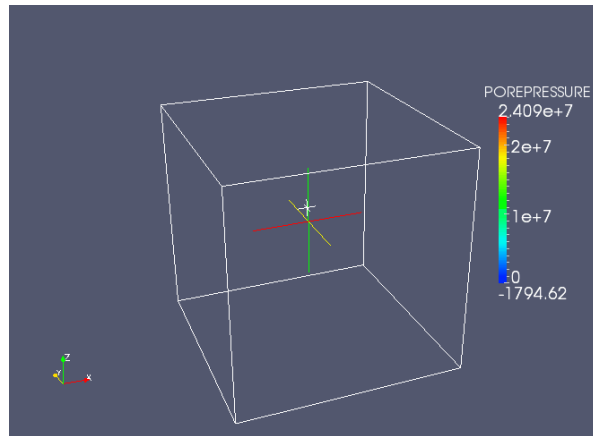
High diffusivity case

Heat pump problem

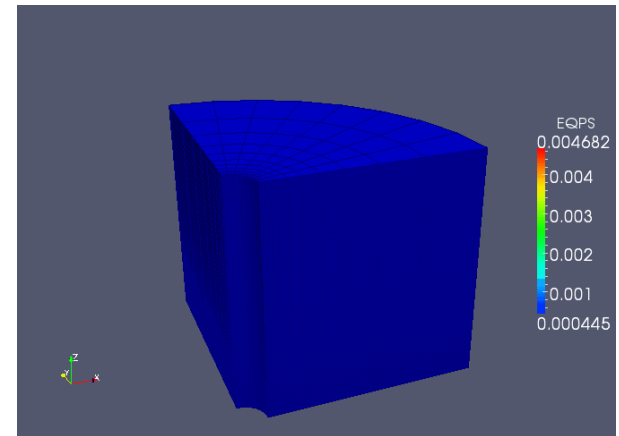
1. Hot liquid is injected into a elasto-plastic porous medium
2. Pore-fluid diffusion and heat diffusion occur at different rate
3. Porous medium expands even though no mechanical load is applied.



Temperature



Pore pressure



equivalent plastic strain

Conclusion

- A fully coupled, finite deformation, stabilization poromechanics finite element model is implemented.
- This model preserve Mandel-Cryer effect, and is able to eliminate spurious oscillation due to the lack of inf-sup condition.
- Thermo-poro-plasticity model is extended and tested.
- Unsaturated flow and fully coupled thermo-poroelasticity will be further tested via analytical solutions.

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