

Incorporating Data from Experiments and Atomistic Simulations into Crystal Plasticity Models for BCC Metals

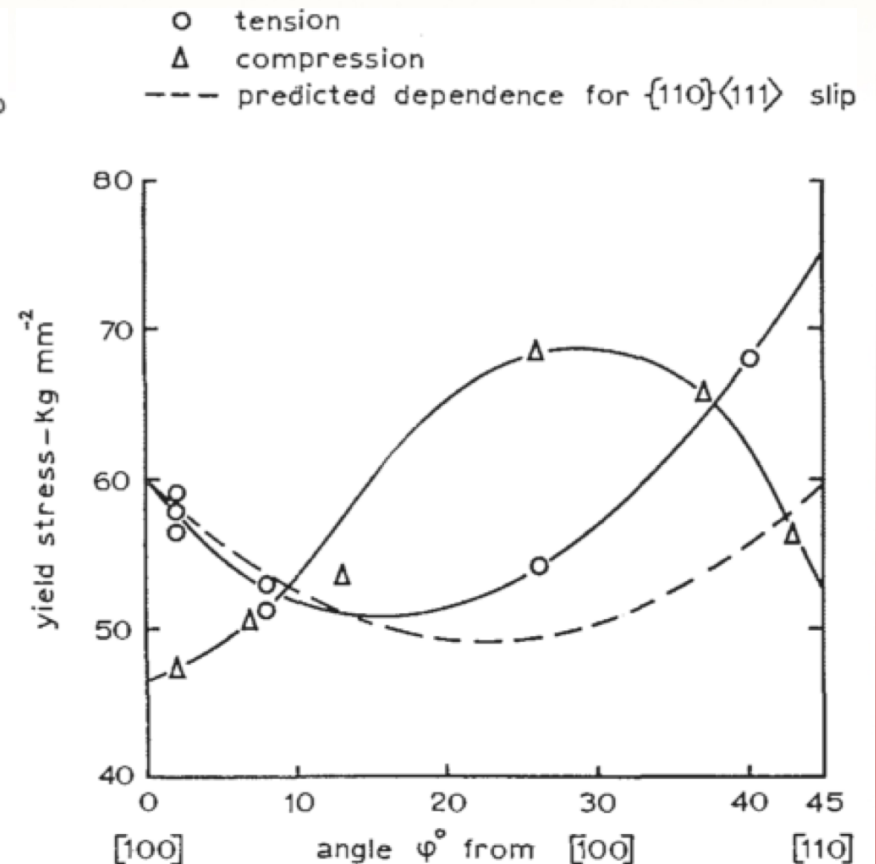
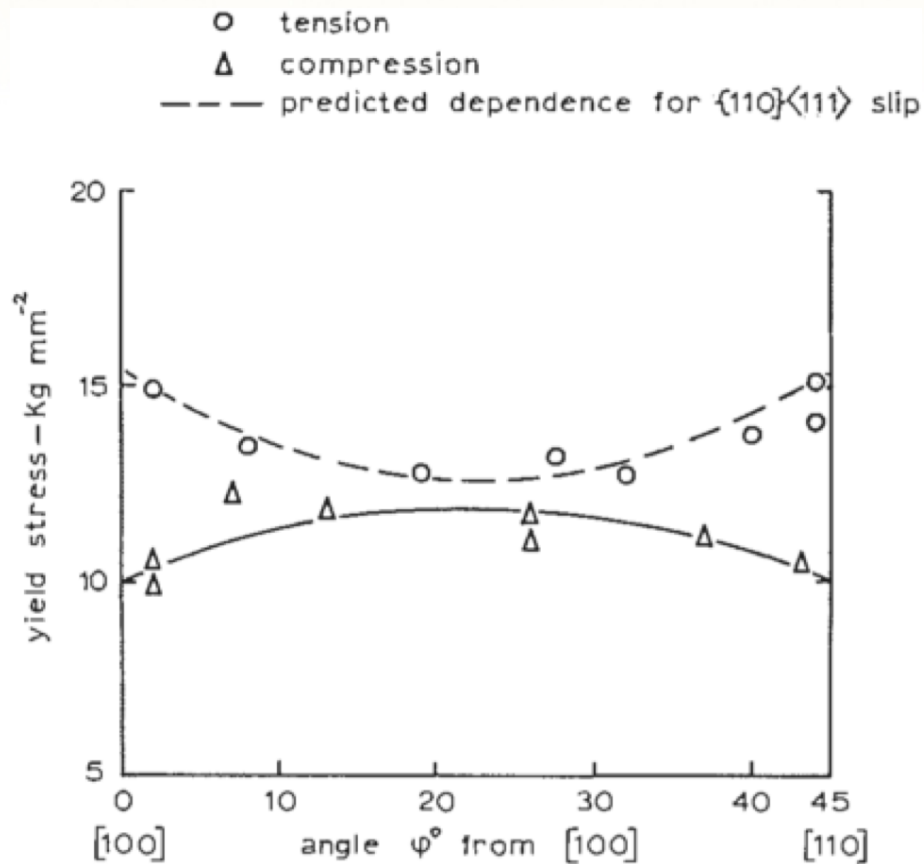
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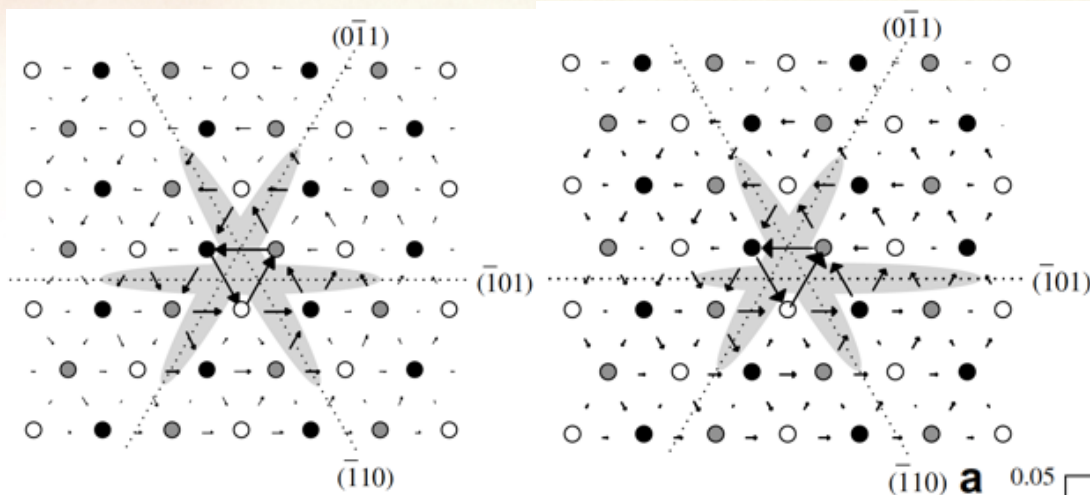
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BCC Metals Show Tension/Compression Asymmetry



Screw Dislocation Core Structure is Important



Screw dislocation cores are non-planar, leading to high lattice friction.

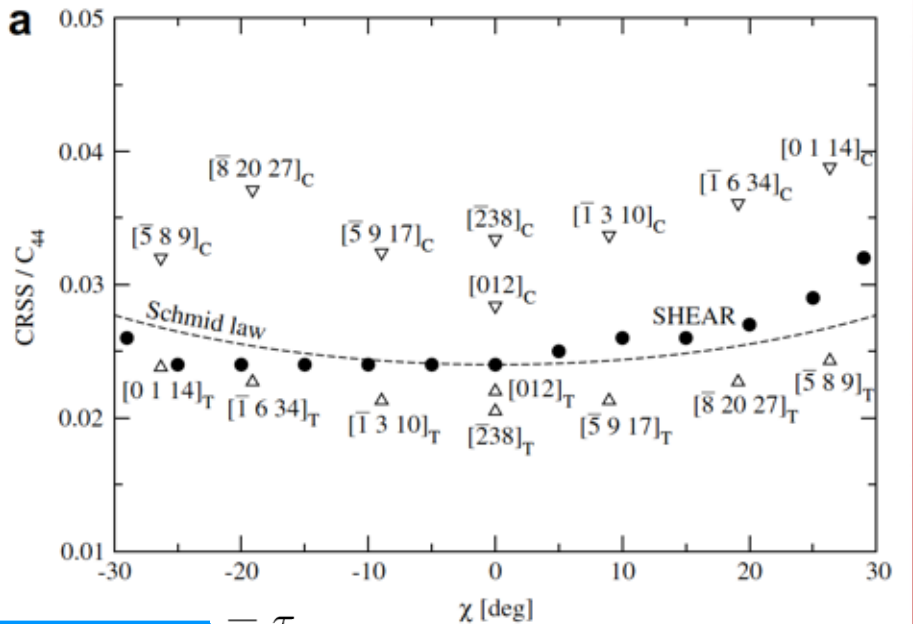
Shear stress perpendicular to the core is important.

$$\mathbf{M}^{(s)} : \sigma + \mathbf{P}_{ns}^{(s)} : \sigma = \tau_{cr}$$

$$\mathbf{m}^s \cdot \sigma \mathbf{n}^s + a_1 \mathbf{m}^s \cdot \sigma \mathbf{n}_1^s + a_2 (\mathbf{n}^s \times \mathbf{m}^s) \cdot \sigma \mathbf{n}^s + a_3 (\mathbf{n}_1^s \times \mathbf{m}^s) \cdot \sigma \mathbf{n}_1^s = \tau_{cr}$$

Gröger, Bailey, and Vitek, *Acta Mater.* **56** (2008) 5401.

Twinning / anti-twinning asymmetry is weak by comparison.



Generalized Yield Law Incorporating Non-Schmid Effect

The resistance to slip on a slip system α :

$$\tau_{cr}^{\alpha} = \mathbf{P}_{tot}^{\alpha} : \boldsymbol{\sigma} = \mathbf{P}_s^{\alpha} : \boldsymbol{\sigma} + \mathbf{P}_{ns}^{\alpha} : \boldsymbol{\sigma}$$

\mathbf{m}^{α} : slip direction

$$\mathbf{P}_s^{\alpha} = c_0 \frac{1}{2} (\mathbf{m}^{\alpha} \otimes \mathbf{n}^{\alpha} + \mathbf{n}^{\alpha} \otimes \mathbf{m}^{\alpha})$$

\mathbf{n}^{α} : slip plane normal

$$\mathbf{t}^{\alpha} = \mathbf{n}^{\alpha} \times \mathbf{m}^{\alpha}$$

$$\mathbf{P}_{ns}^{\alpha} = c_1 \mathbf{t}^{\alpha} \otimes \mathbf{m}^{\alpha} + c_2 \mathbf{t}^{\alpha} \otimes \mathbf{n}^{\alpha} + c_3 \mathbf{n}^{\alpha} \otimes \mathbf{n}^{\alpha} + c_4 \mathbf{t}^{\alpha} \otimes \mathbf{t}^{\alpha} + c_5 \mathbf{m}^{\alpha} \otimes \mathbf{m}^{\alpha}$$

Representation of various yield criteria

Ref.	Structure	c_0	c_1	c_2	c_3	c_4	c_5
Qin and Bassanni (1992)	FCC	$1 + \frac{\sqrt{3}}{3}B$	$\frac{\sqrt{6}}{3}B$	0	0	0	0
Steinmann et al. (1998)	FCC	α^{sm}	α^{cm}	α^{mm}	0	0	0
Qin and Bassanni (1992)	L1 ₂	$1 + \frac{\sqrt{3}}{3}B$	$\frac{\sqrt{6}}{3}B$	$-A - \frac{7}{9}Ak$	$\frac{2\sqrt{2}}{9}Ak$	$-\frac{2\sqrt{2}}{9}Ak$	0
Dao and Asaro (1993)	L1 ₂	1	$2\eta_{zs}^0$	$2\eta_{mz}^0$	$2\eta_{mm}^0$	$2\eta_{zz}^0$	$2\eta_{ss}^0$
Gröger et al. (2008)	BCC	$1 + \frac{1}{2}a_1$	$\frac{\sqrt{3}}{2}a_1$	$-a_2 + \frac{1}{2}a_3$	$\frac{\sqrt{3}}{4}a_3$	$-\frac{\sqrt{3}}{4}a_3$	0
Yalcinkaya et al. (2008)	BCC	1	η_2	η_3	η_4	η_5	η_1

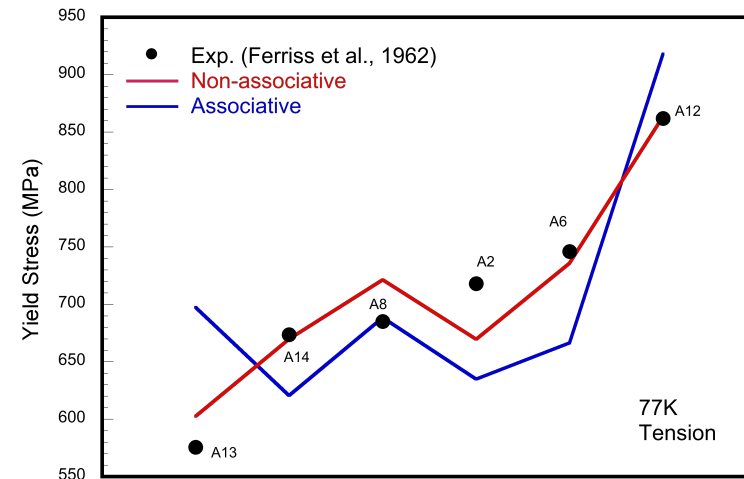
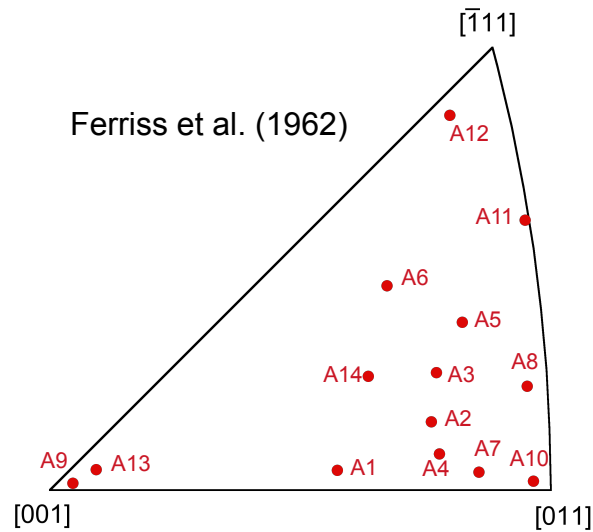


Yield Stress Prediction for Ta Single Crystals

$$\tau_{cr} + \tau_{obs}^{\alpha} = \mathbf{P}_{tot}^{\alpha} : \boldsymbol{\sigma}$$

$$\sigma_y = \frac{\tau_{cr} + \tau_{obs}^{\alpha}}{P_{tot}^T} \quad (\text{Tension})$$

$$\sigma_y = -\frac{\tau_{cr} + \tau_{obs}^{\alpha}}{P_{tot}^C} \quad (\text{Compression})$$



Ta single crystal tests

Temp.	c_1	c_2	c_3	c_4	τ_{cr}	Std. dev. (A)	Std. dev. (NA)
77 K	-0.35	0.12	-0.79	0.53	383	75.3 MPa	27.4 MPa
300 K	-1.18	0.70	0.12	0.08	57	10.9 MPa	7.7 MPa

Best-fit Parameters

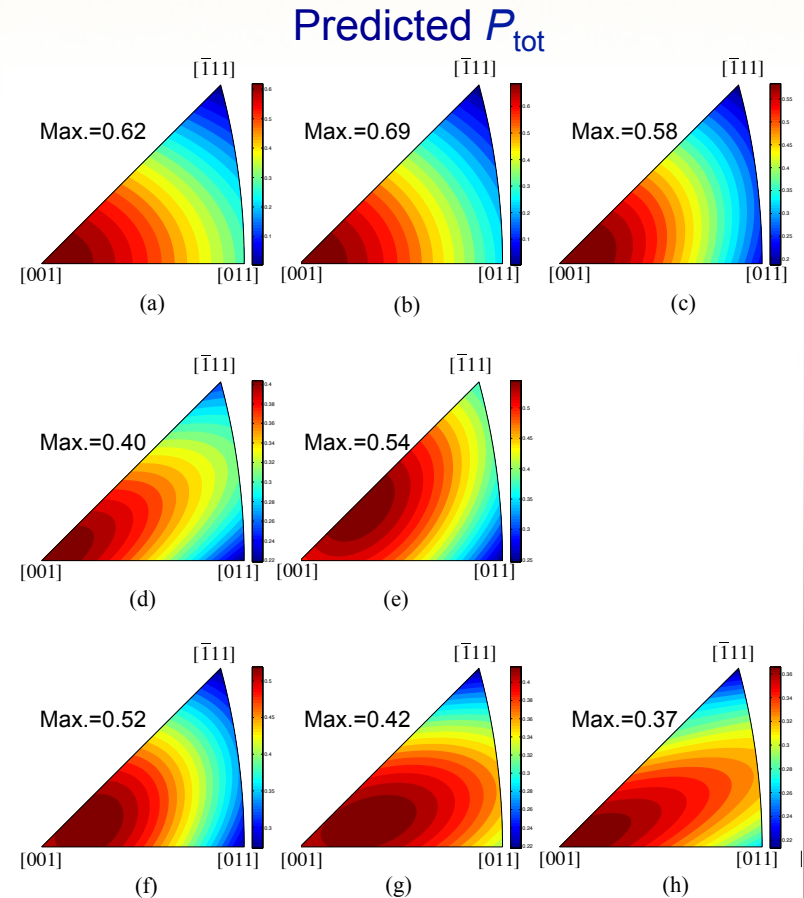


Yield Stress Prediction for Mo, W and Ta Single Crystals

Best-fit non-Schmid constants

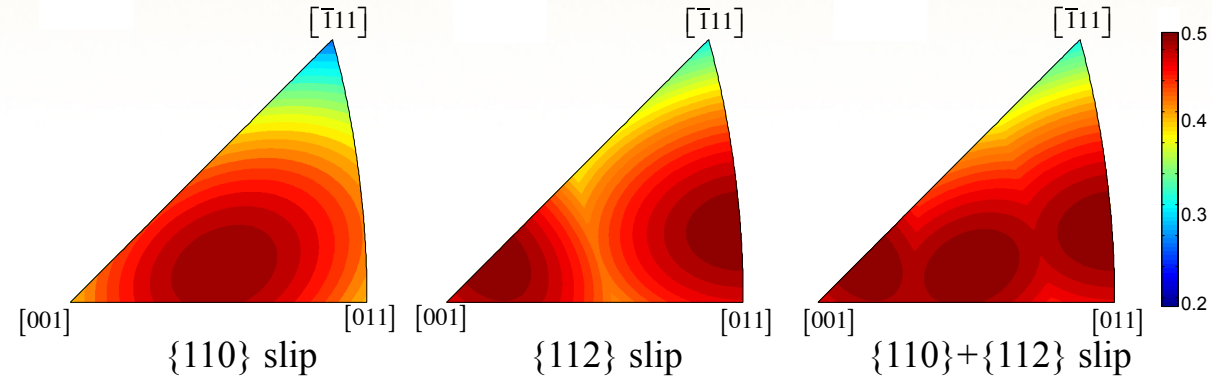
Material	c_1	c_2	c_3	c_4	τ_{cr}	Ref.	
Mo (a)	0.11	0.70	0.04	-0.04	208	Sherwood et al. (1967)	Mo
Mo (b)	0.31	0.70	0.11	-0.11	215	Guiu and Pratt (1966)	
Mo (c)	0.42	0.23	0.13	-0.02	147	Irwin et al. (1974)	
W (d)	0.15	0.05	-0.01	0.26	325	Rose et al. (1962)	W
W (e)	0.36	-0.09	0.23	-0.01	293	Argon and Maloof (1966)	
Ta (f)	0.30	0.04	0.11	0.01	173	Sherwood et al. (1967)	
Ta (g)	-0.11	0.19	-0.03	0.16	225	Hull et al. (1967)	
Ta (h)	-0.07	0.15	-0.08	0.29	211	Ferriss et al. (1962)	

$$T = 77 \text{ K}, \quad \dot{\epsilon} = 10^{-4} \text{ s}^{-1}$$



{110} vs. {112} Slip

Predicted P_{tot}



Standard deviation (MPa) between predicted and measured yield stresses

	Baseline model			Non-Schmid model			References
	{110}	{112}	{110}+{112}	{110}	{112}	{110}+{112}	
Mo	311	311	311	105	211	188	Sherwood (1967)
	388	310	297	12	109	114	Guin (1966)
	185	185	185	1	170	123	Irwin (1974)
W	322	318	314	47	68	67	Rose (1962)
	344	344	344	1	83	39	Argon (1966)
Ta	135	135	135	1	45	29	Sherwood (1967)
	101	101	101	1	1	1	Hull (1967)
	104	105	95	20	66	23	Ferriss (1962)
	107	112	103	74	69	63	Byron (1968)



BCC CP-FEM Formulation

Slip rate: $\dot{\gamma}^\alpha = \dot{\gamma}_0^\alpha \left(\frac{\tau^\alpha}{g^\alpha} \right)^{1/m}$ (Hutchinson, 1976) 24 $\langle 111 \rangle \{110\}$ slip systems

Slip resistance: $g^\alpha = \max(\tau_{\text{cr}}^\alpha - \tau_{\text{ns}}^\alpha, 0) + \tau_{\text{obs}}^\alpha$ (Weinberger, 2012)

$\xrightarrow{\text{Lattice friction}}$ $\xrightarrow{\text{Obstacle stress}}$

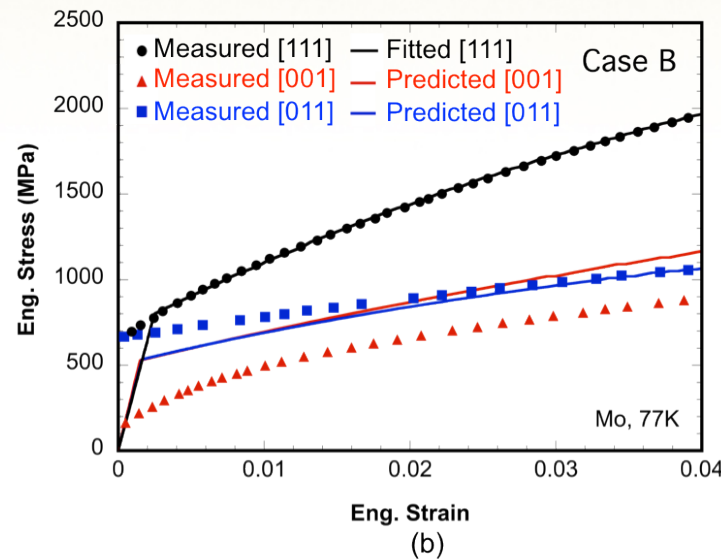
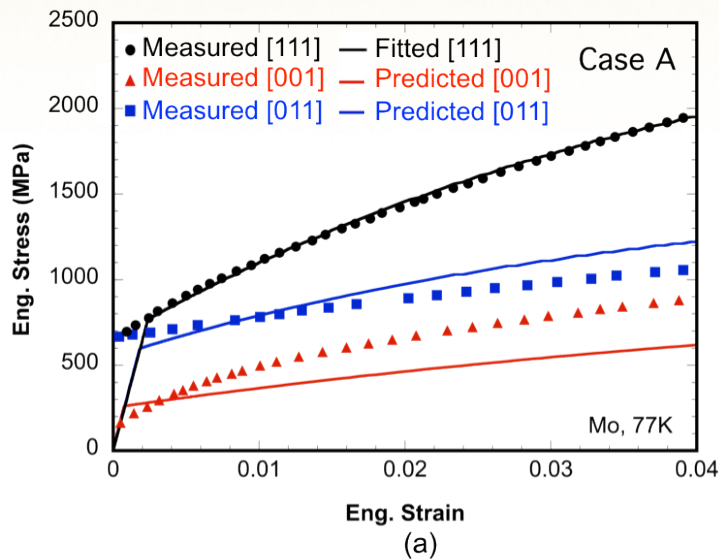
Obstacle stress: $\tau_{\text{obs}}^\alpha = \alpha \mu b \sqrt{\sum_{\beta=1}^{NS} \rho^\beta}$ (Taylor, 1934)

$$\rho^\alpha = \left(\kappa_1 \sqrt{\sum_{\beta=1}^{NS} \rho^\beta} - \kappa_2 \rho^\alpha \right) \cdot |\dot{\gamma}^\alpha| \quad (\text{Kocks, 1976})$$



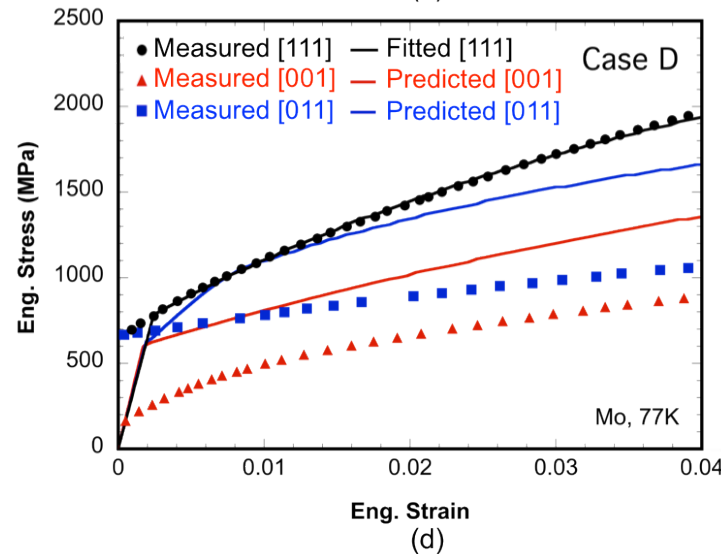
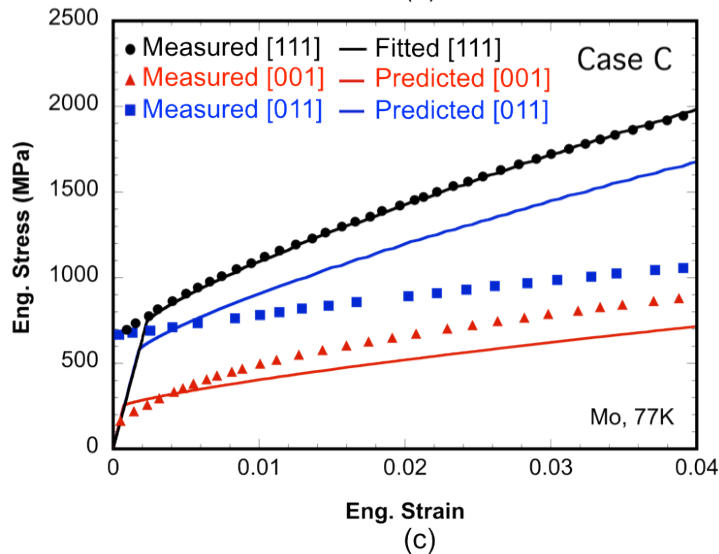
Single Crystal Stress-Strain Predictions for Mo

Orientation dependent stress-strain curve is most accurately predicted in Case A.



(a) Case A
Non-associative model
{110} slip
Taylor hardening

(b) Case B
Associative model
{110} slip
Taylor hardening



(c) Case C
Non-associative model
{110} slip
Power-law hardening

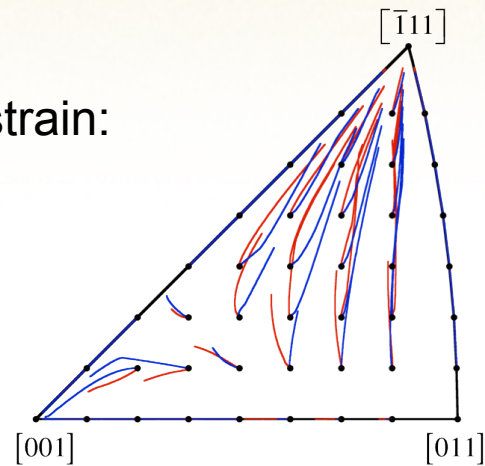
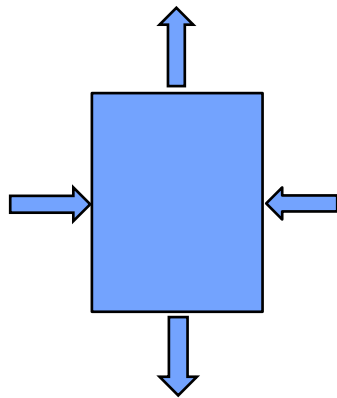
(d) Case D
Non-associative model
{112} slip
Taylor hardening



Texture Evolution in Mo

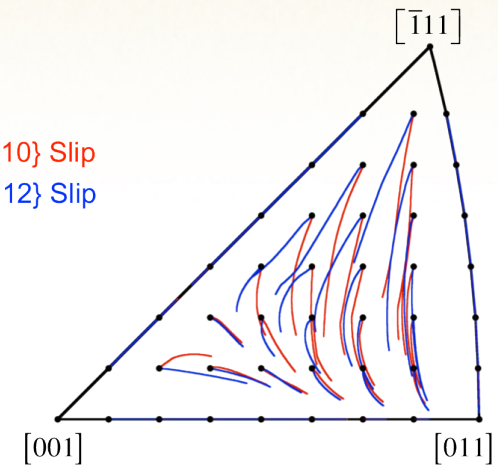
Isochoric deformation to 20% strain:

$$\dot{\epsilon} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1/2 & 0 \\ 0 & 0 & -1/2 \end{pmatrix} dt$$

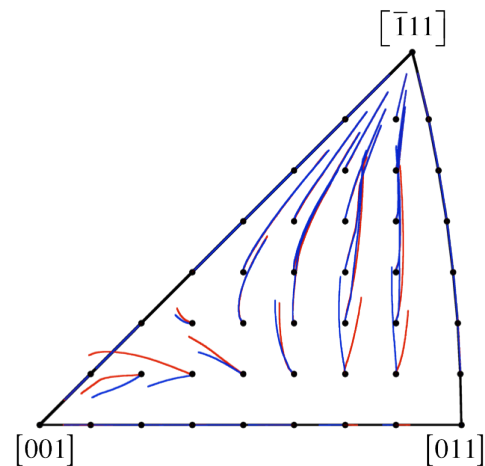


Compression

— {110} Slip
— {112} Slip

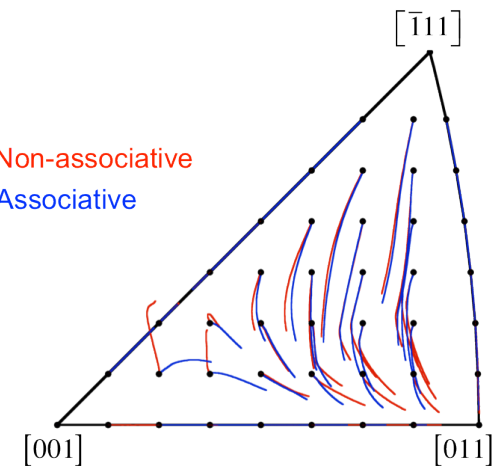


Tension



Compression

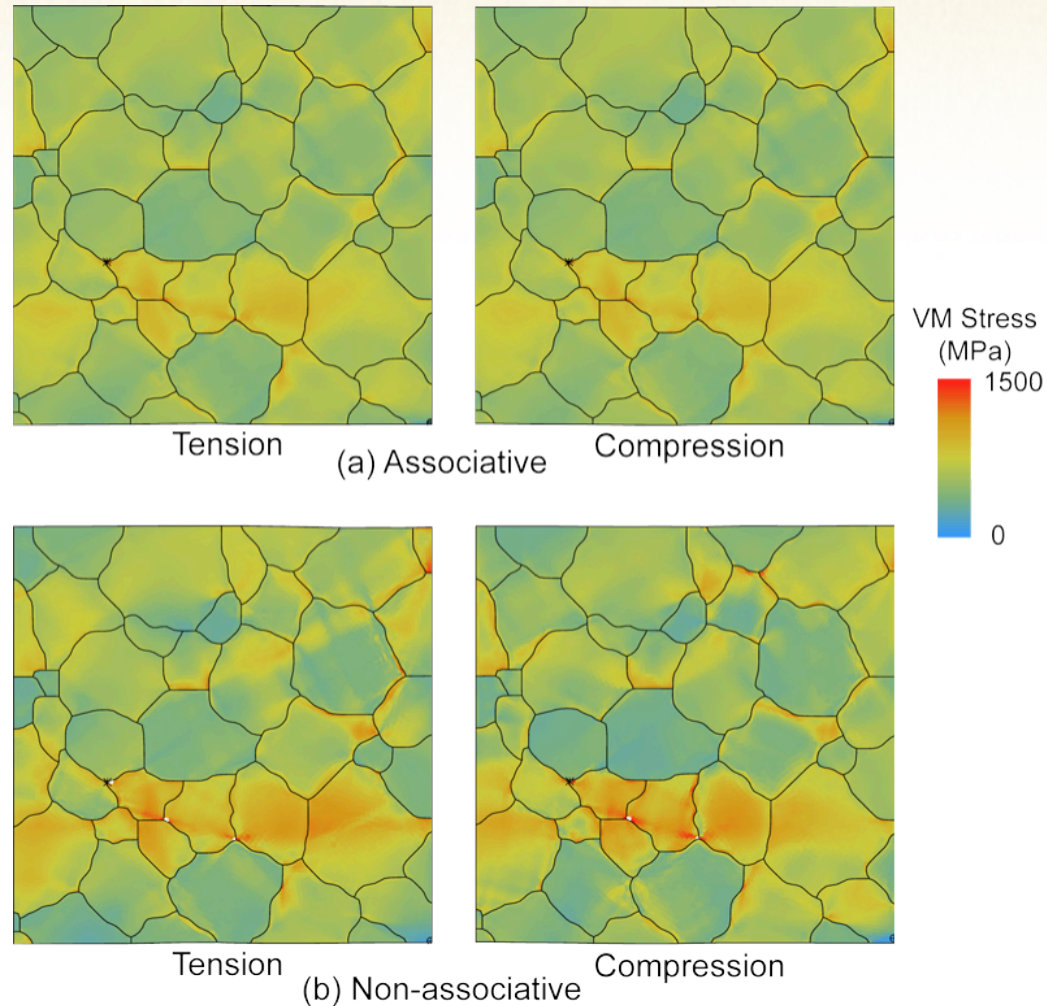
— Non-associative
— Associative



Tension



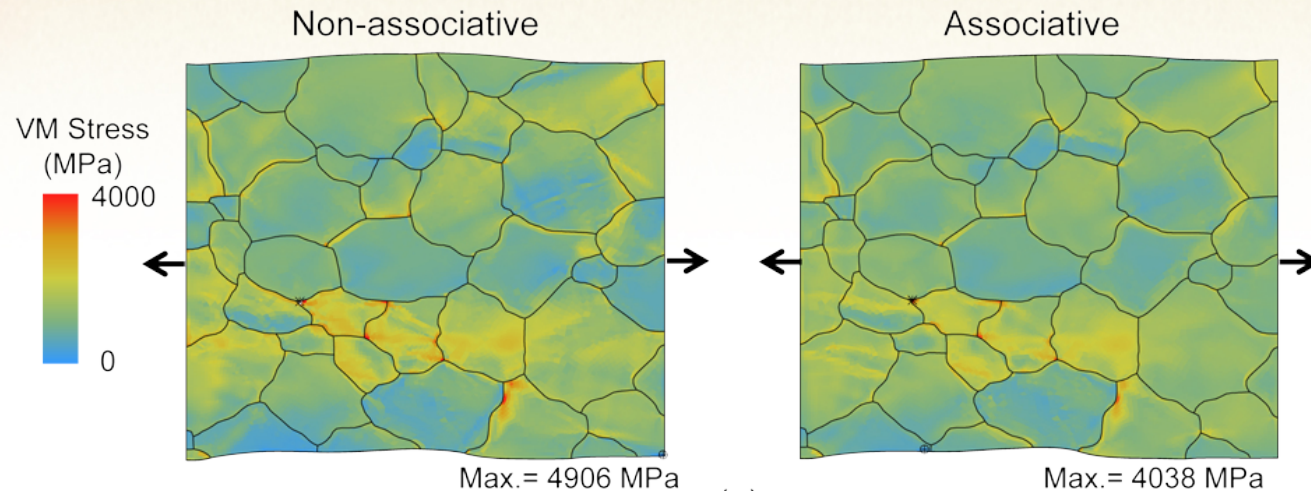
Stress Maps in Mo



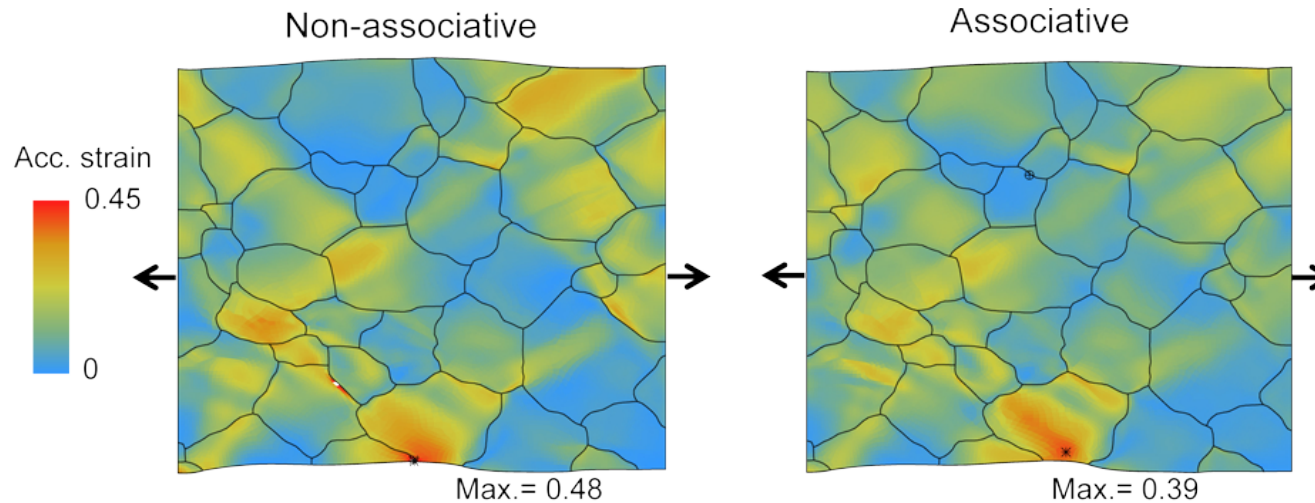
Similar Stress Distributions
Non-Associative Model has higher and more localized stresses,
greater tendency for failure



Stress and Strain Maps in Mo



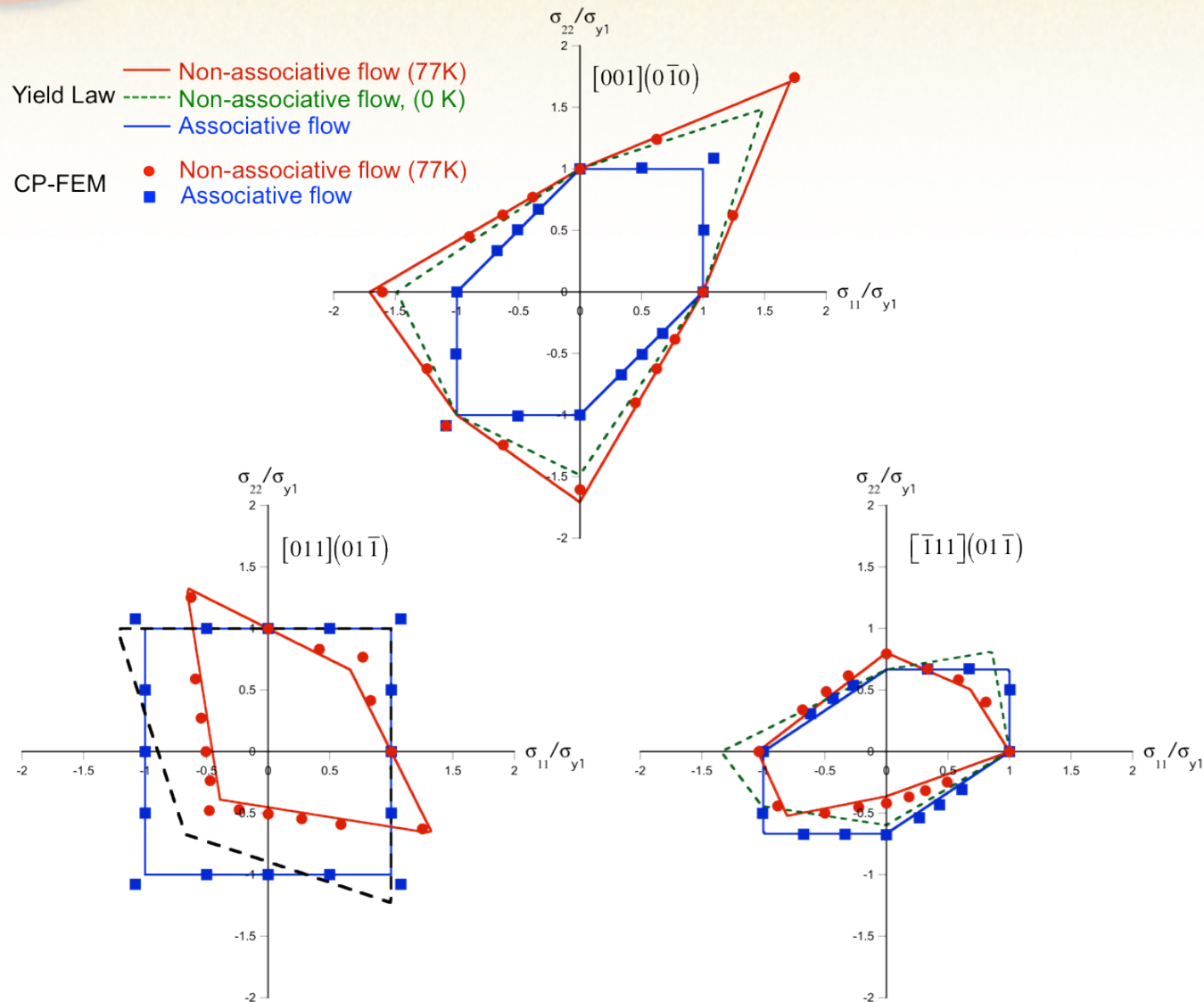
Von Mises Stress Distribution



Accumulated Strain Distribution

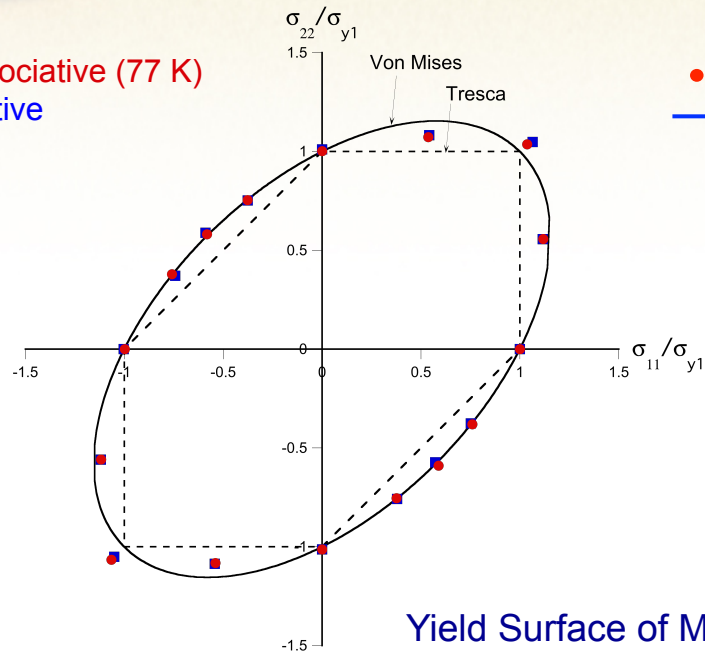


Yield Surface Predictions: Mo Single Crystals

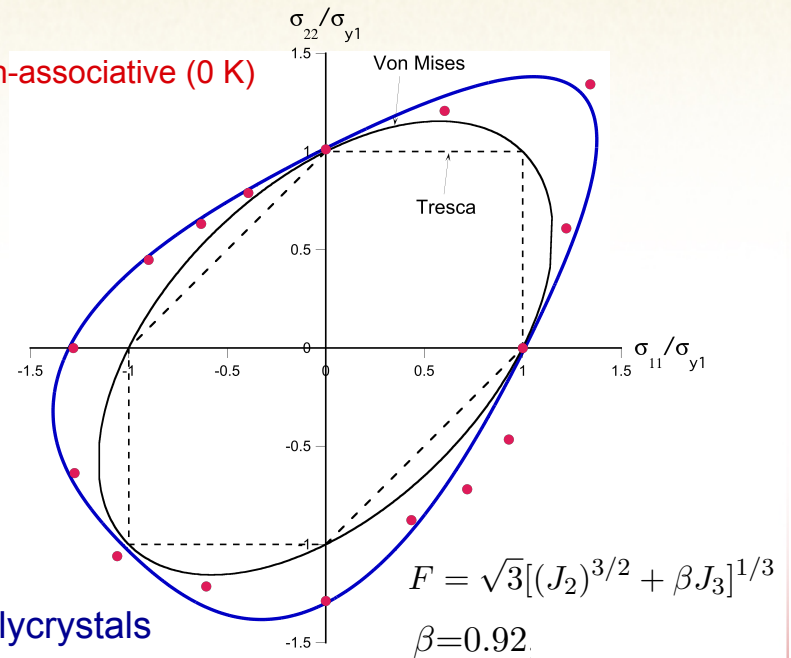


Yield Surface Predictions: Mo Polycrystals

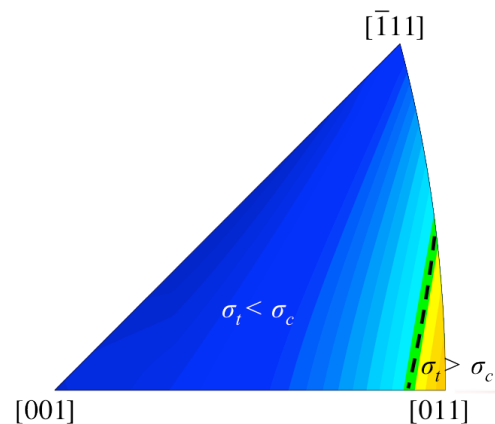
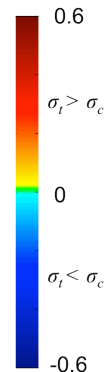
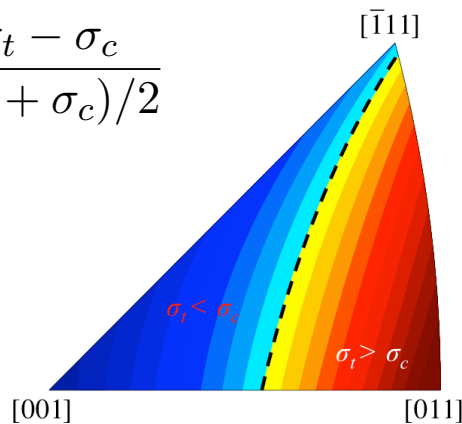
- Non-associative (77 K)
- Associative



- Non-associative (0 K)
- Fit

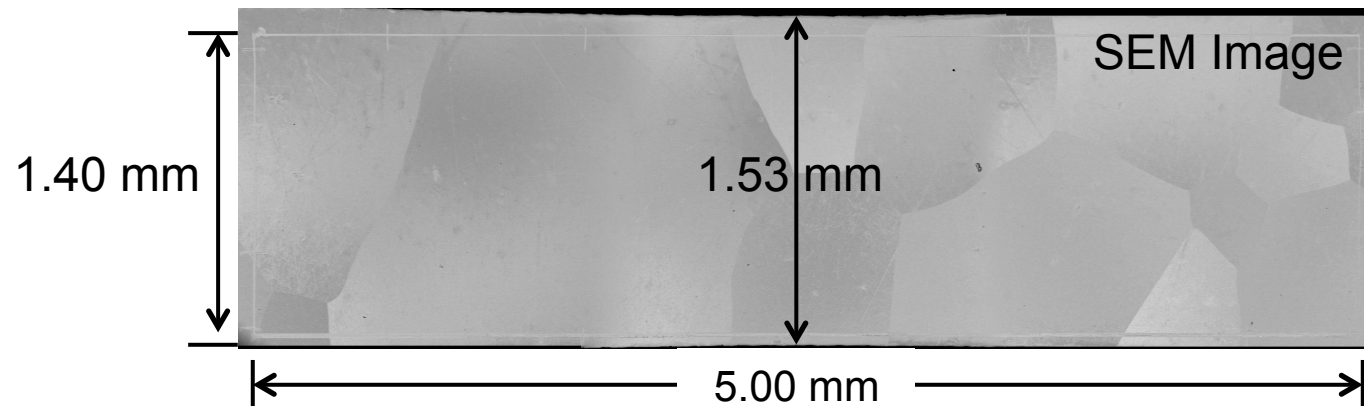
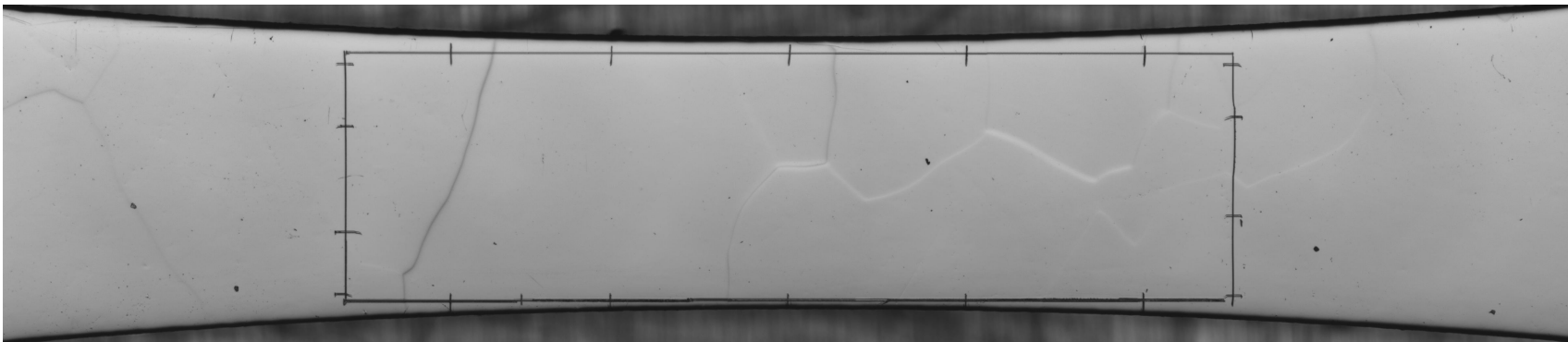
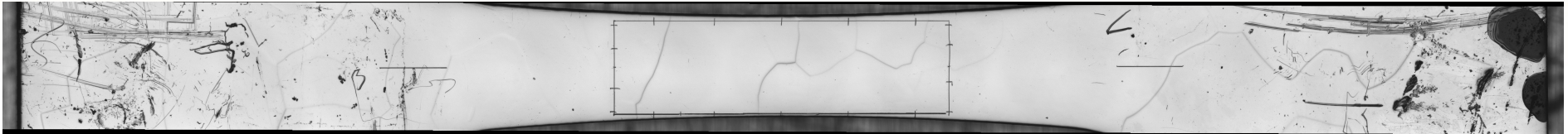


$$SD = \frac{\sigma_t - \sigma_c}{(\sigma_t + \sigma_c)/2}$$

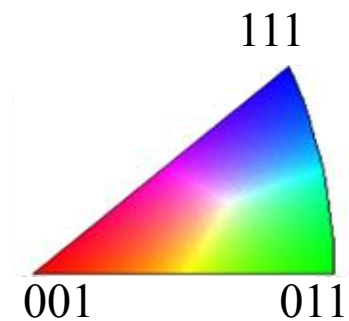
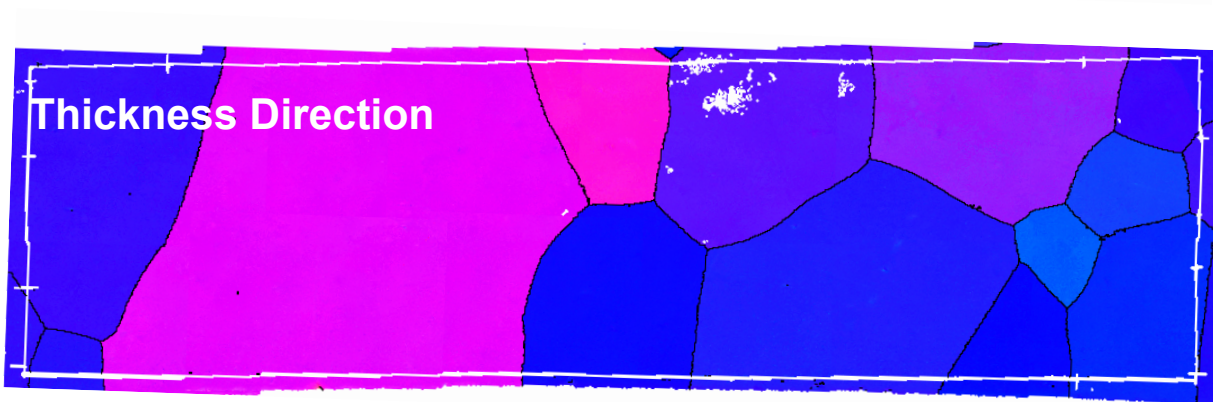
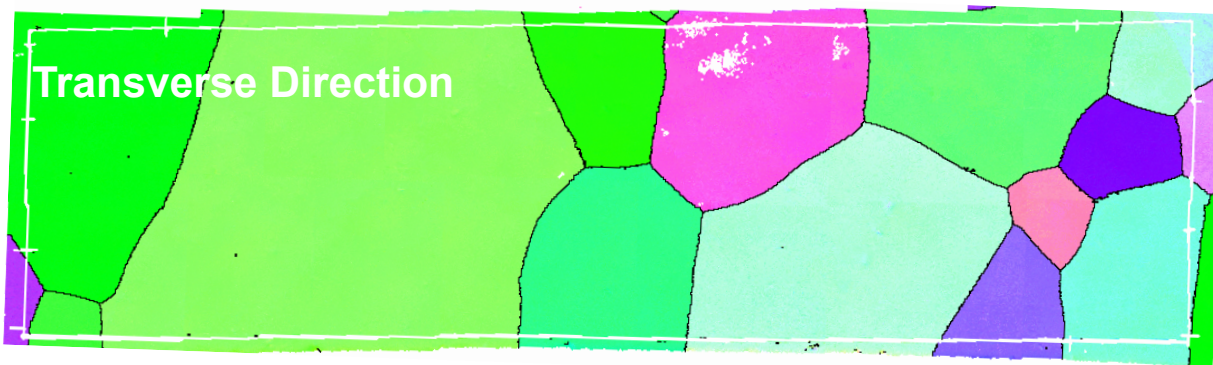
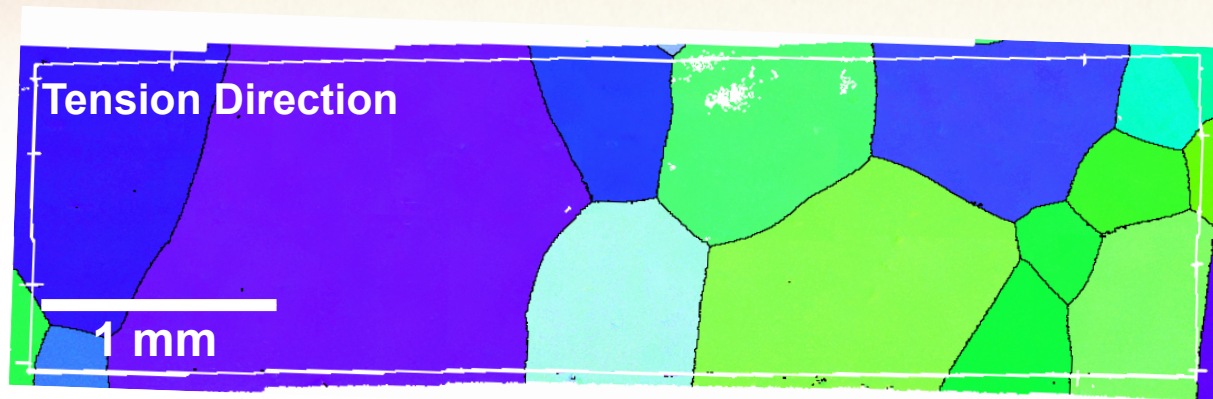


Digital Image Correlation of Ta Oligocrystal

23 mm

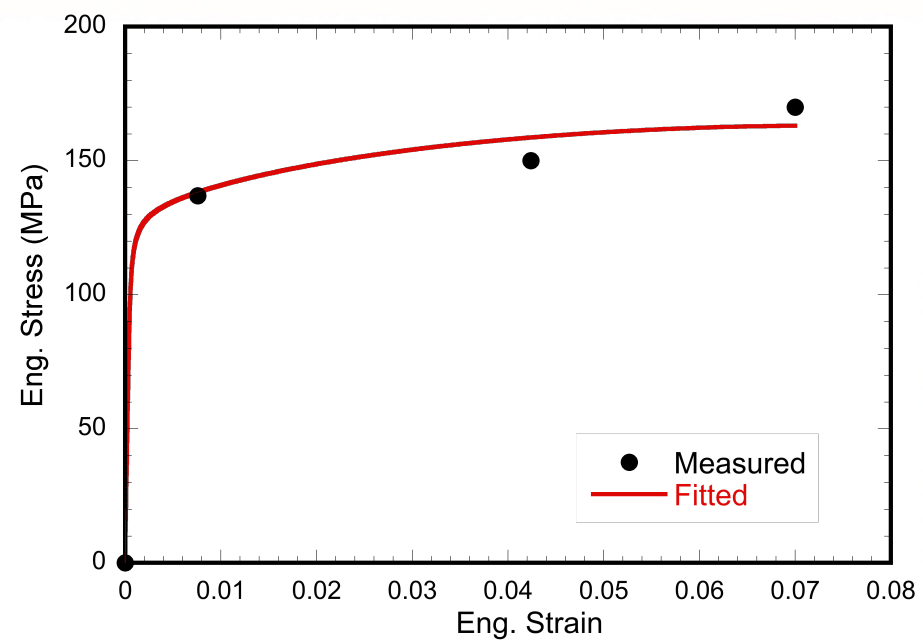
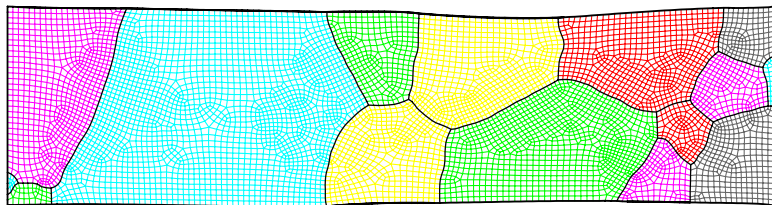
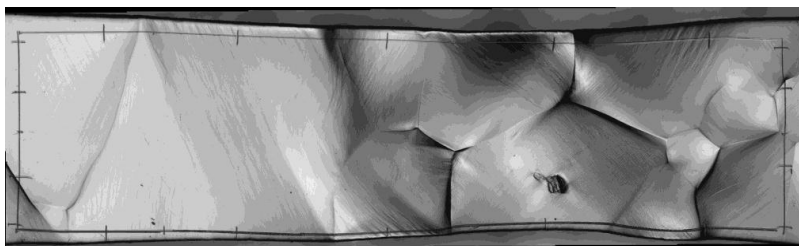
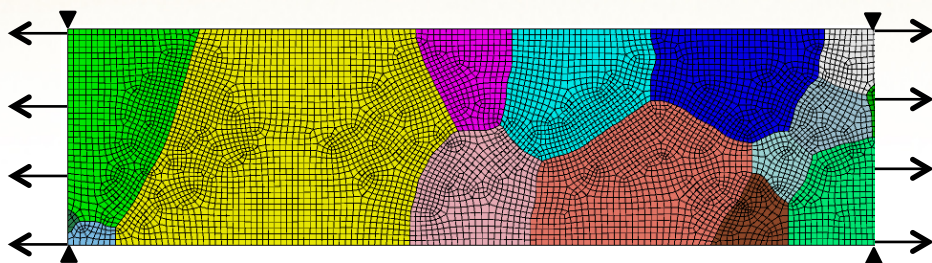


Electron Backscatter Diffraction Maps



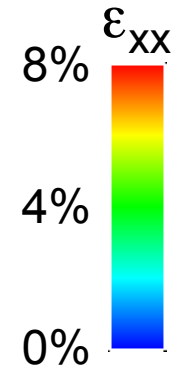
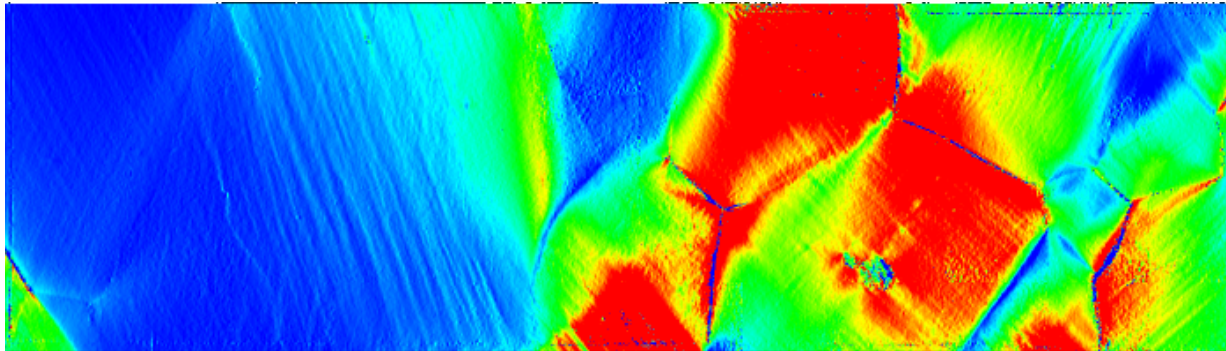
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Finite Element Mesh

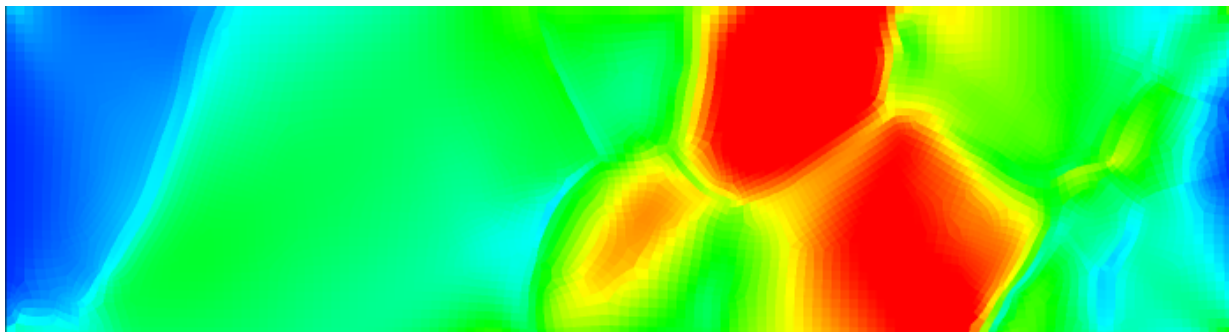


Validating CP-FEM Simulations w/ DIC Experiments

Experiment



CP-FEM Simulation

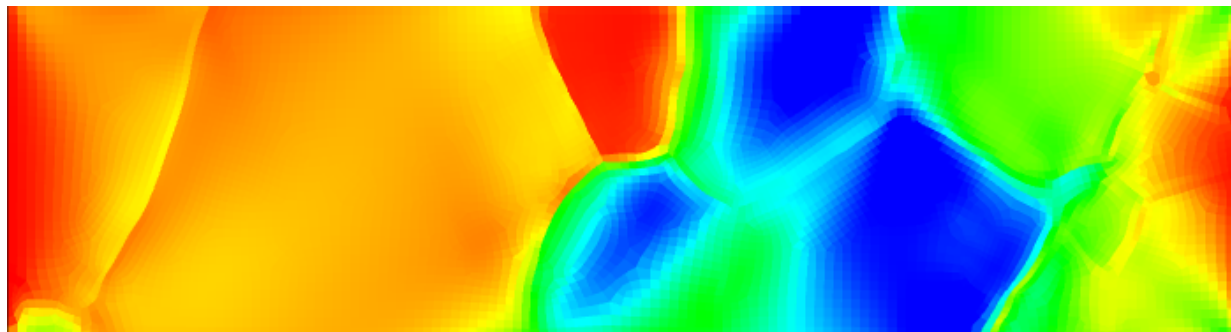
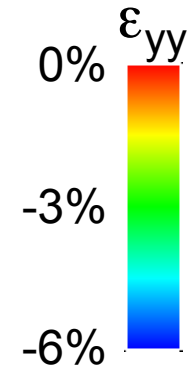
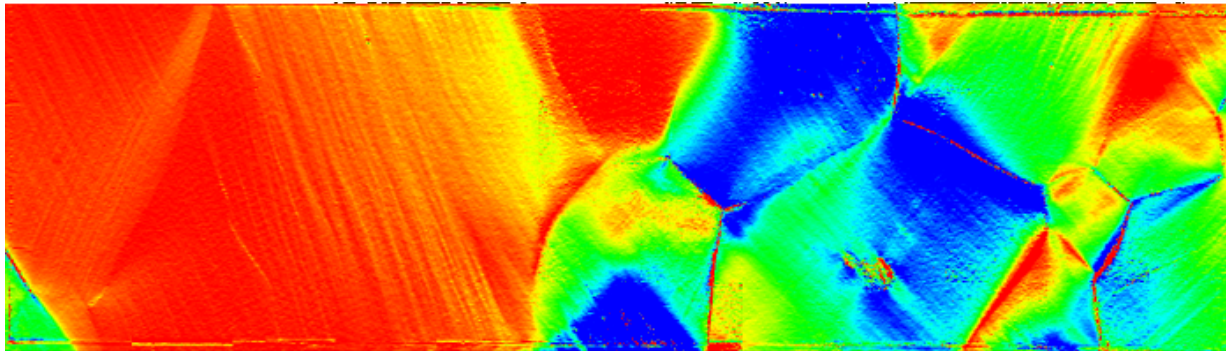


4.2% Applied Strain



Validating CP-FEM Simulations w/ DIC Experiments

Experiment



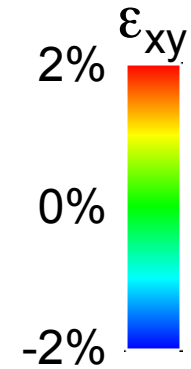
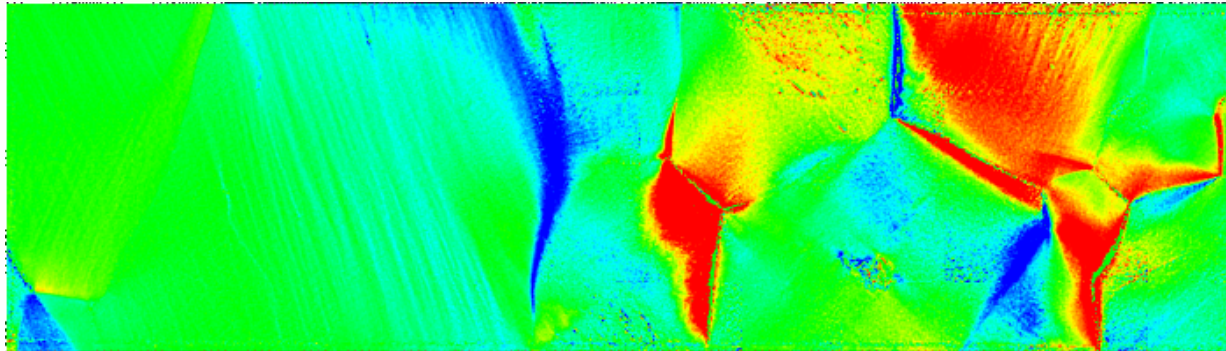
CP-FEM Simulation

4.2% Applied Strain

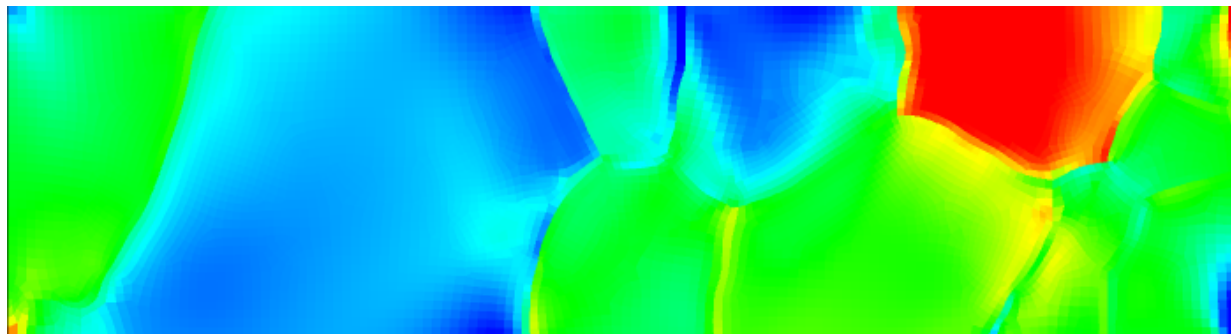


Validating CP-FEM Simulations w/ DIC Experiments

Experiment



CP-FEM Simulation



4.2% Applied Strain



Summary

- Dislocation plasticity in BCC metals (and most materials, for that matter) is significantly more complex than in FCC.
- We have developed a generalized yield criterion for BCC transition metals, and implemented it into a BCC CP-FEM model.
- Yield criteria are calibrated to single crystal experiments on Mo, Ta, and W.
- Non-Schmid effects are clearly reflected in the stress-strain response, texture evolution, damage localization, and yield surfaces of single and polycrystals.
- Early CP-FEM predictions show good qualitative agreement with digital image correlation experiments on Ta oligocrystals.

