

Chemical Dynamics of Excited Electronic States

SAND2019-0157PE

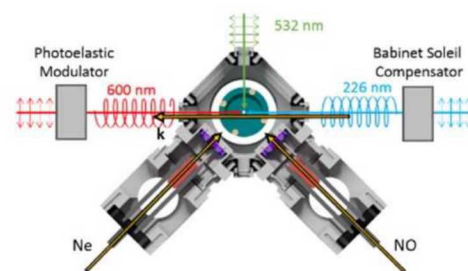
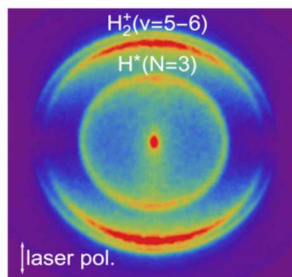
David W. Chandler

Collisional energy transfer of aligned and oriented $\text{NO}(A, v=0, N=2)$

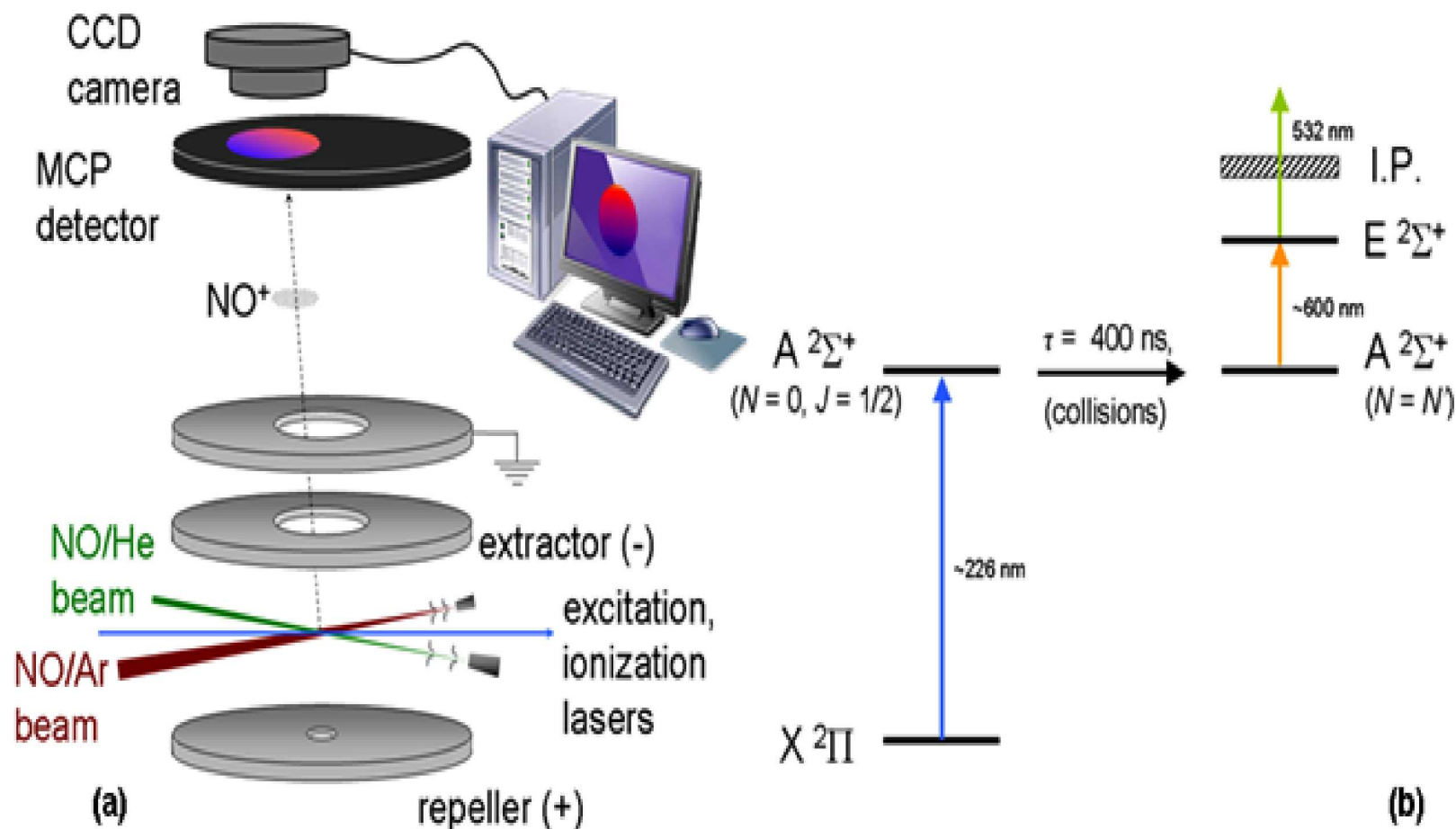
Matt Costen Tom Sharples, Joe Leng, Ken McKendrick (Heriot Watt University)
Javier Aoiz (Cuidad, Madrid)

Polarizability of $\text{H}_2(\text{EF}, v=0, N=2)$

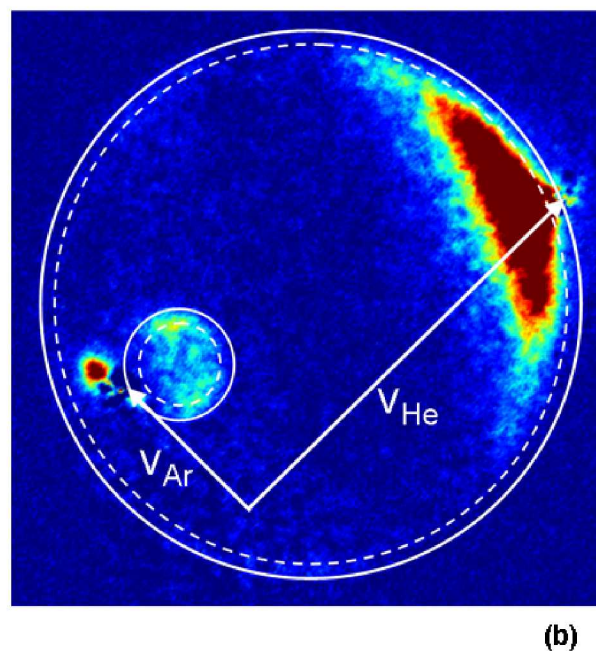
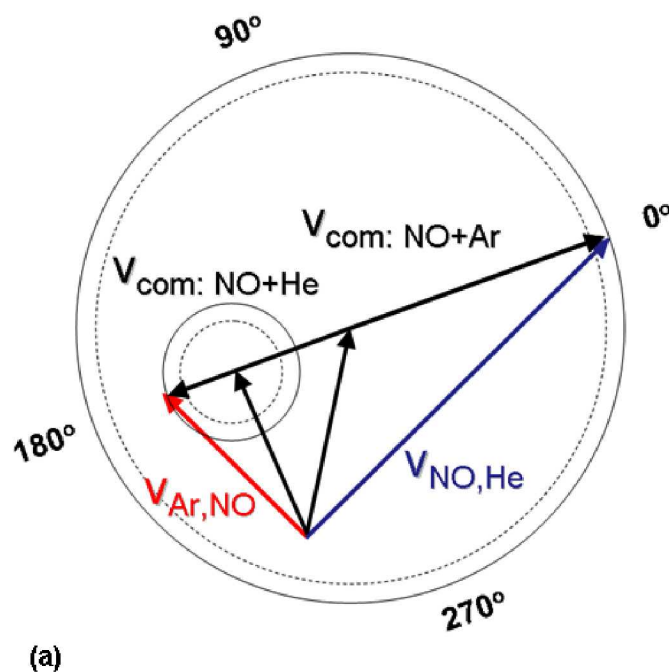
Justin Jankunis, Gary Lopez, Martin Fournier, Peter Rakitzis (University of Crete)



In the Past we Measured DCS and Alignment of NO(A) State Collisions.

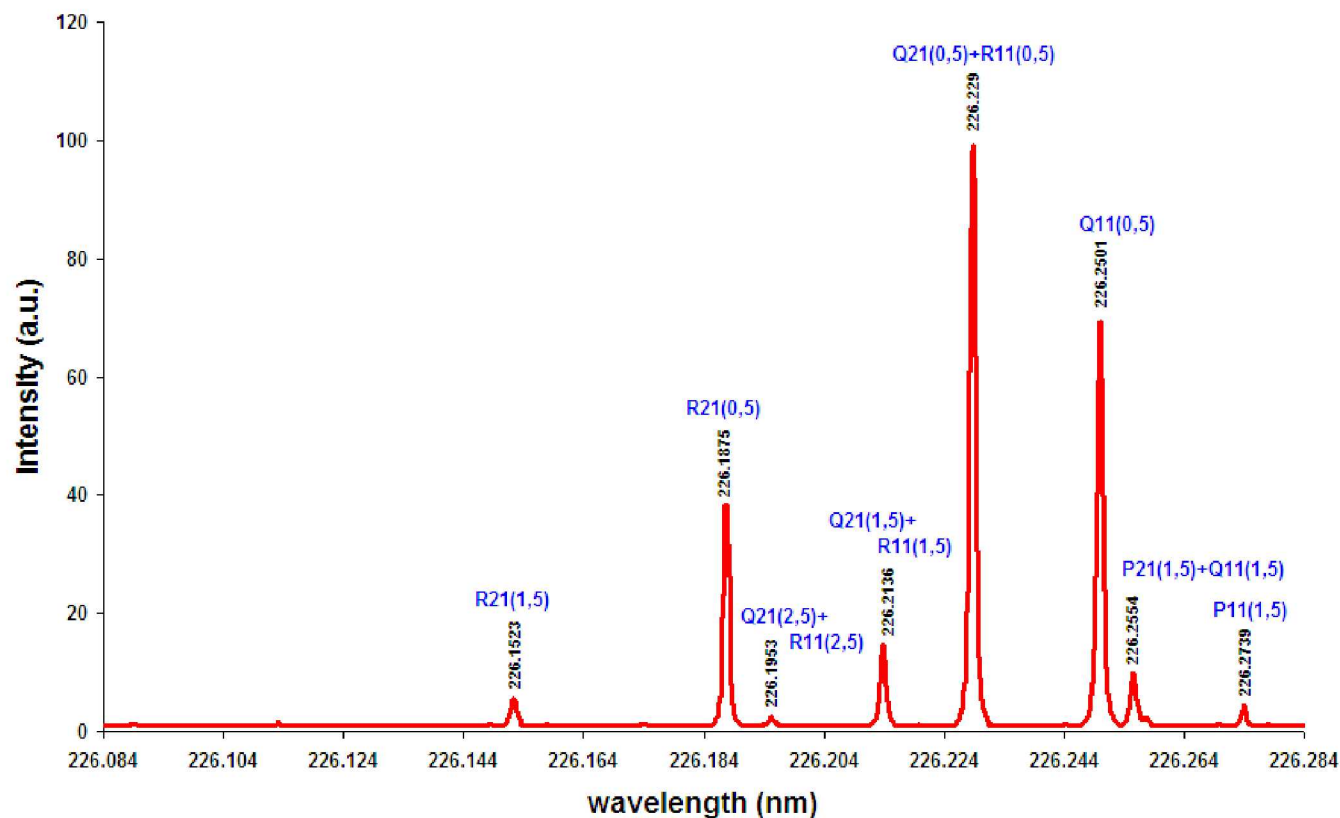


Crossing two beam: NO/He with NO/Ar. NO is Scattered from Each Beam Simultaneously.

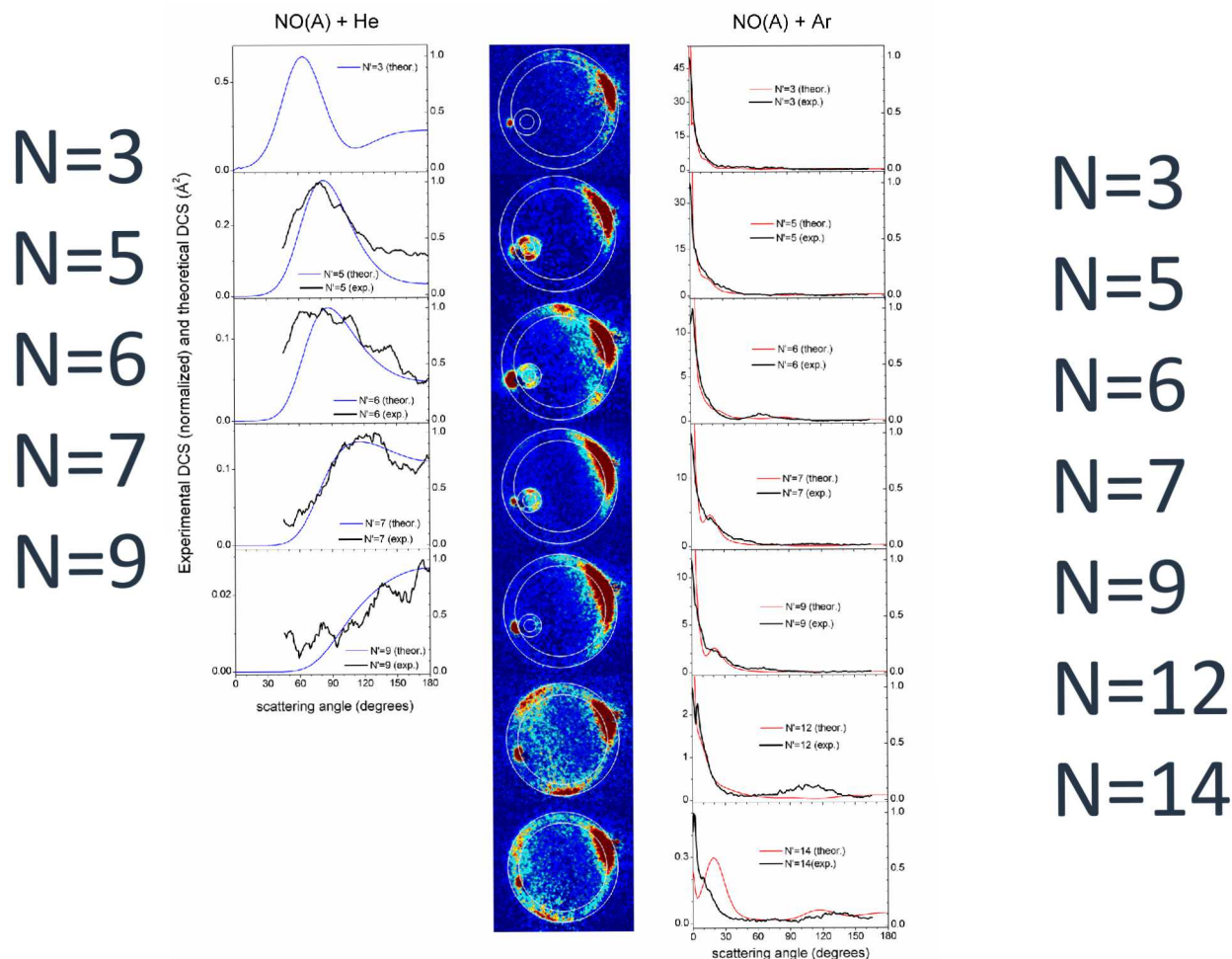


We Excite Different Rotational Levels of the A state Using Different Transitions.

NO($v=0$) A $\leftarrow$ X LIF spectrum

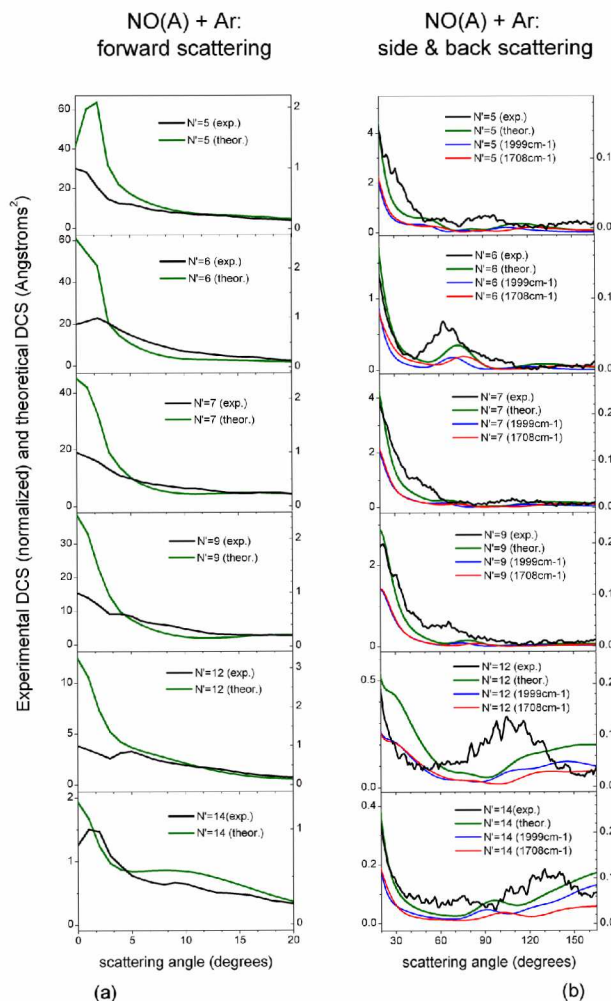


Easier to Compare the Calculations of the NO He and NO Ar scattering when performed in this manner.



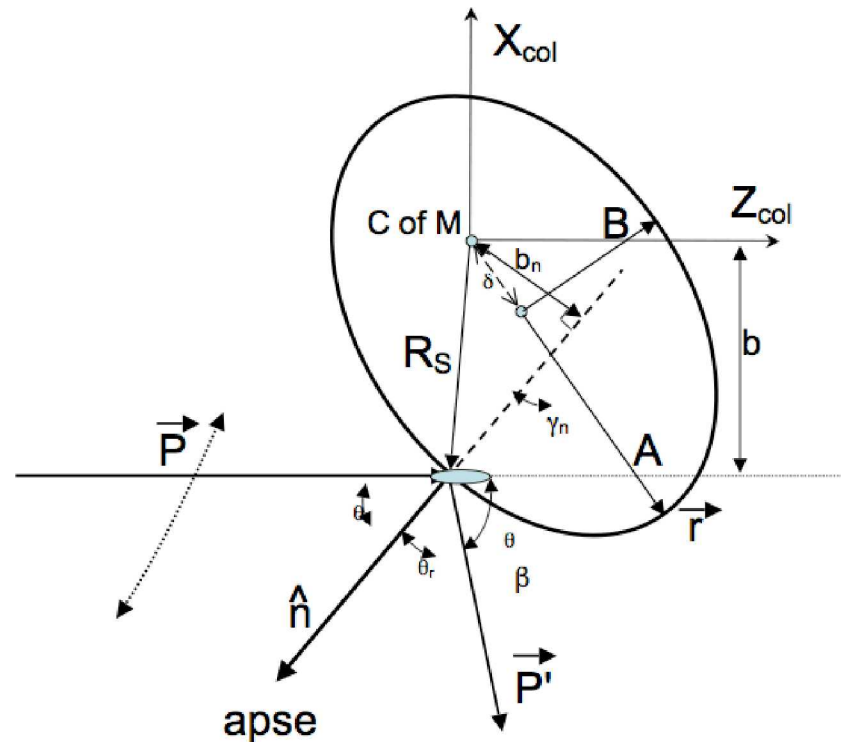
QM Scattering Calculations do not Reproduce Exactly the Results but Capture the Trends.

Experiment
under
Estimates the
forward
scattering due
background
subtraction
procedure



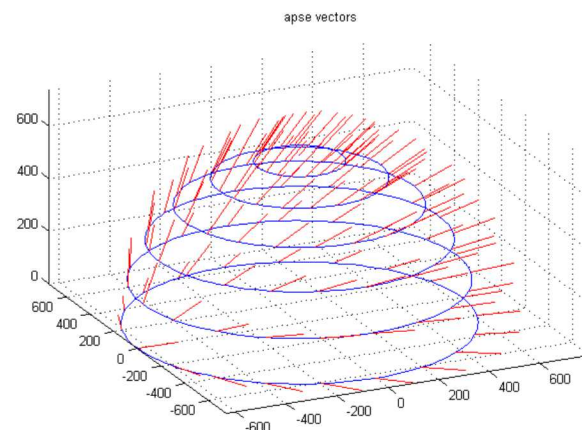
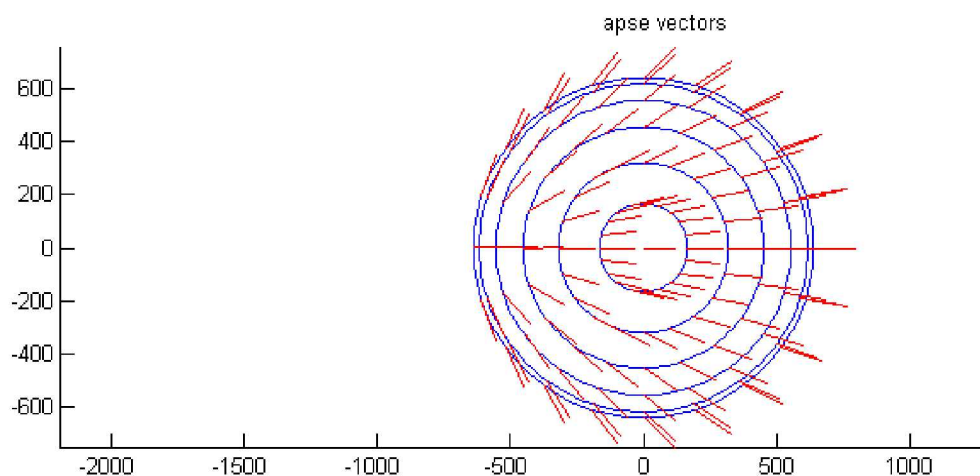
Theory to
predict the
position of the
Rainbow
Scattering at too
large and angle.

Collisions at High Collision Energy Can be Roughly Thought of as Hard and Sudden.



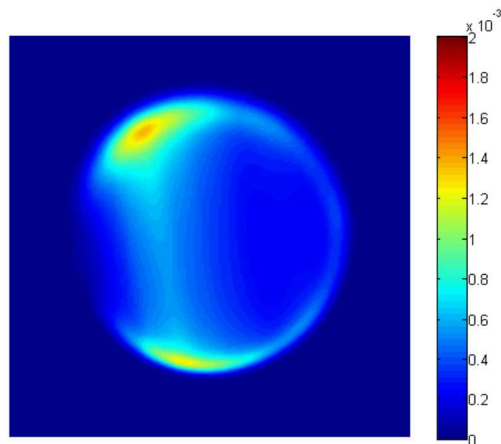
The classical dynamics associated with hard ellipse scattering can be analytically described. The Apse vector is the vector along which angular momentum is transferred

The Apse Model States That the Final J Vector Cannot Lie Along the Apse Vector

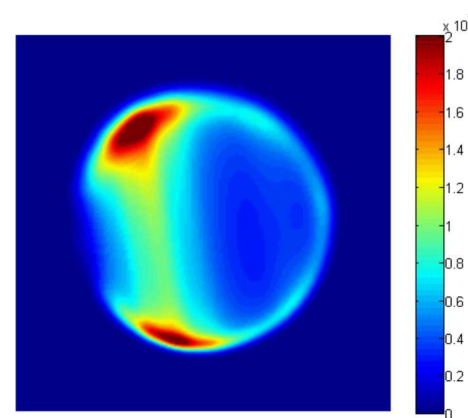


This Simple Rule Implies an Alignment to the Collisions

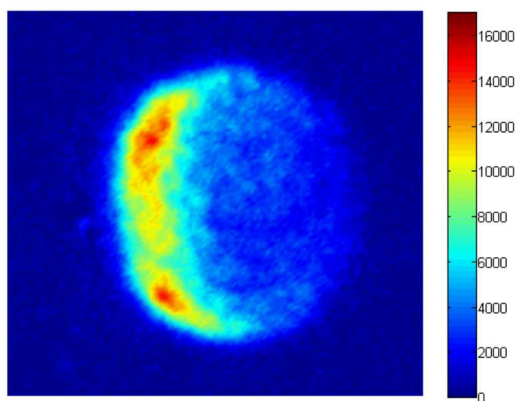
Detection with Linear Polarized Light Show the Alignment Inherent in the Images. NO(N=9)



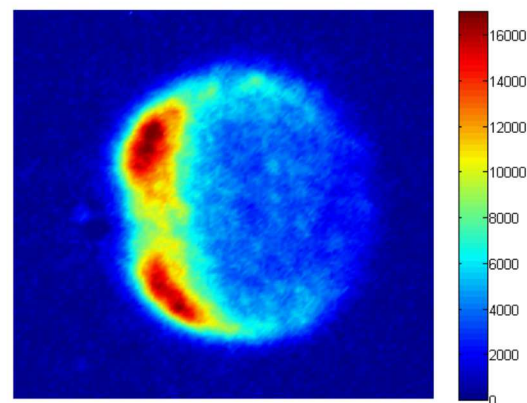
Exact model: V



Exact Model: H

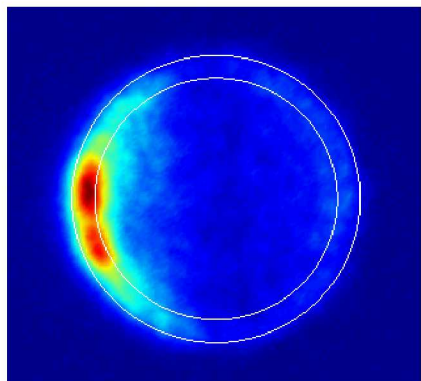


Experiment: V

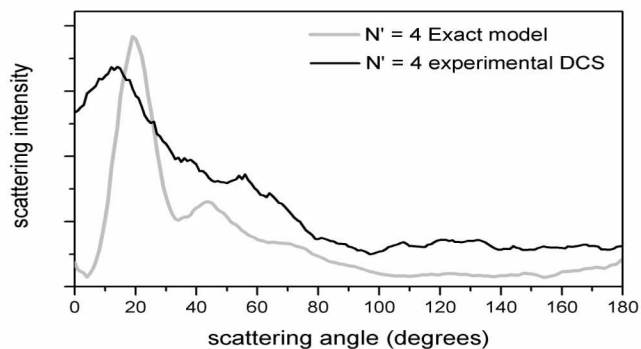


Experiment: H

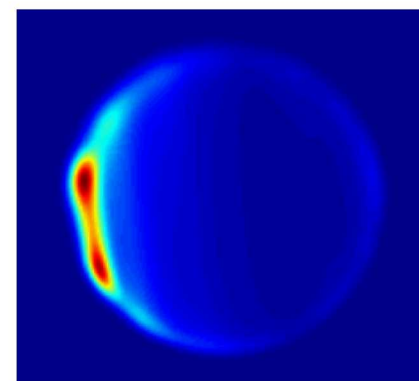
Calculated and Experimental Images are in Good Agreement.



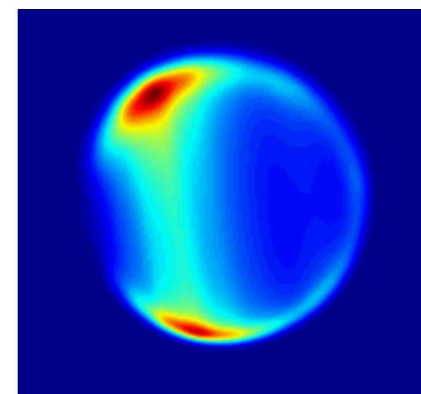
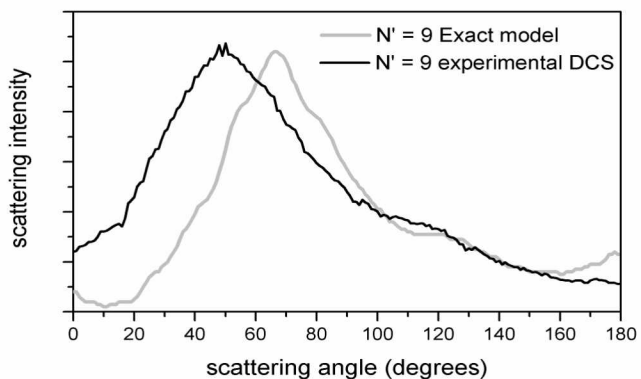
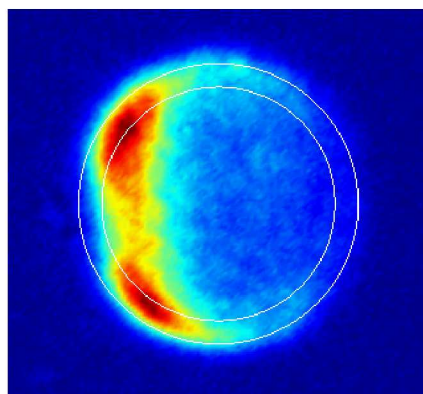
Experimental images



Annular profile comparison

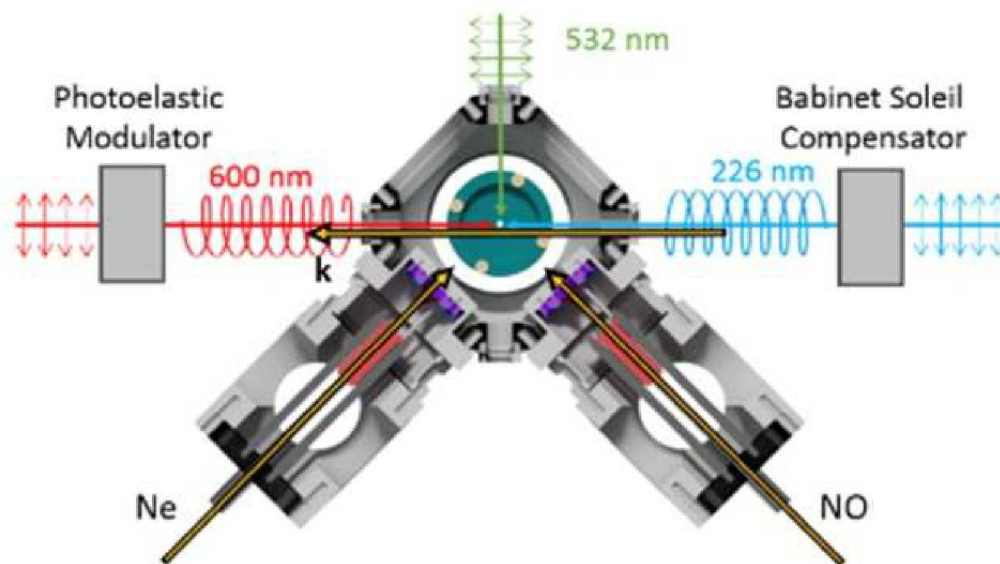


Exact model

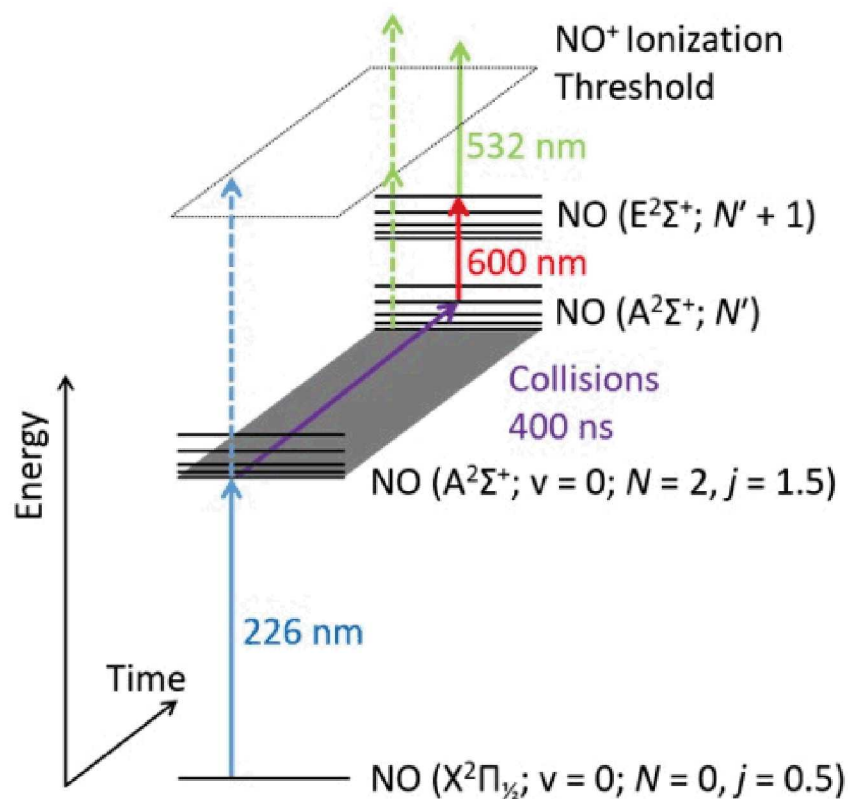


Collisional Energy Transfer of Oriented NO (A, $v=0$, $N=2$)

- 1) Prepare NO(A) with circularly polarized light at ~ 226 nm
- 2) Orient K vector with relative velocity vector of collision
- 3) Detect with circularly polarized light



Scheme for production of NO(A) and detection of NO (A).

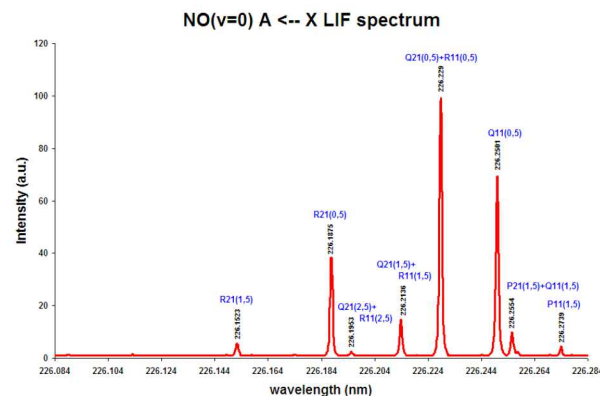


NO(A) has lifetime of 200ns

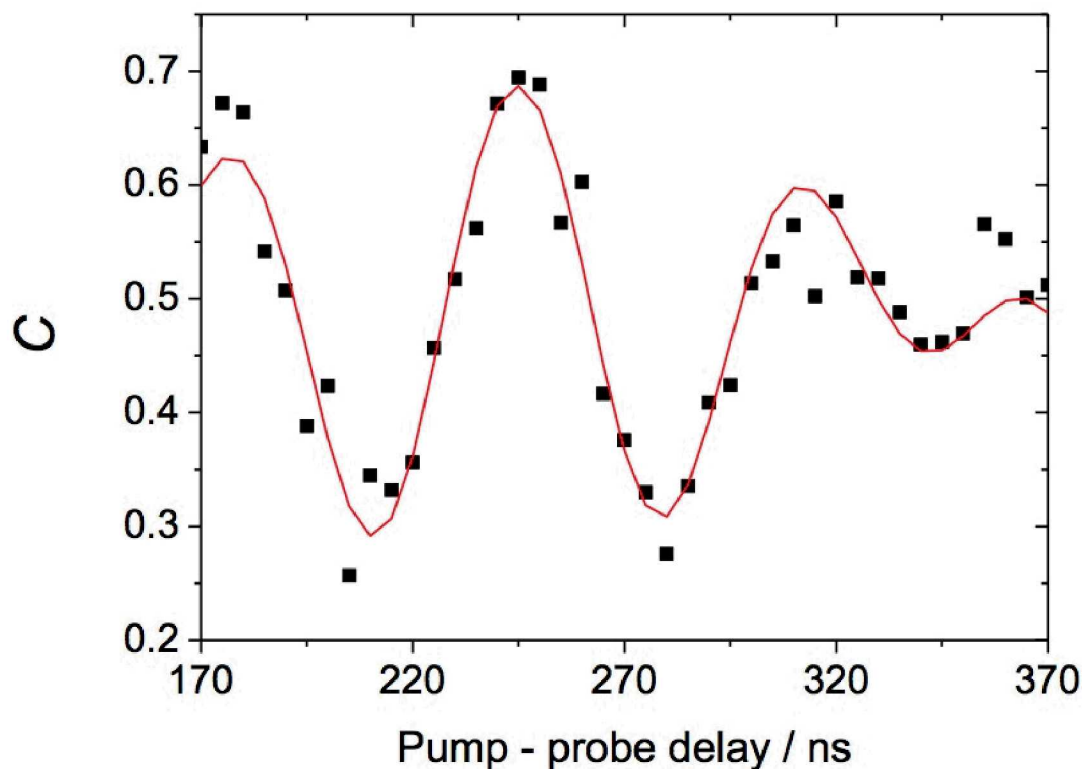
Detect NO (A) after 400 ns

Excite from $N=0, j=0.5$ to $N=2, j=1.5$

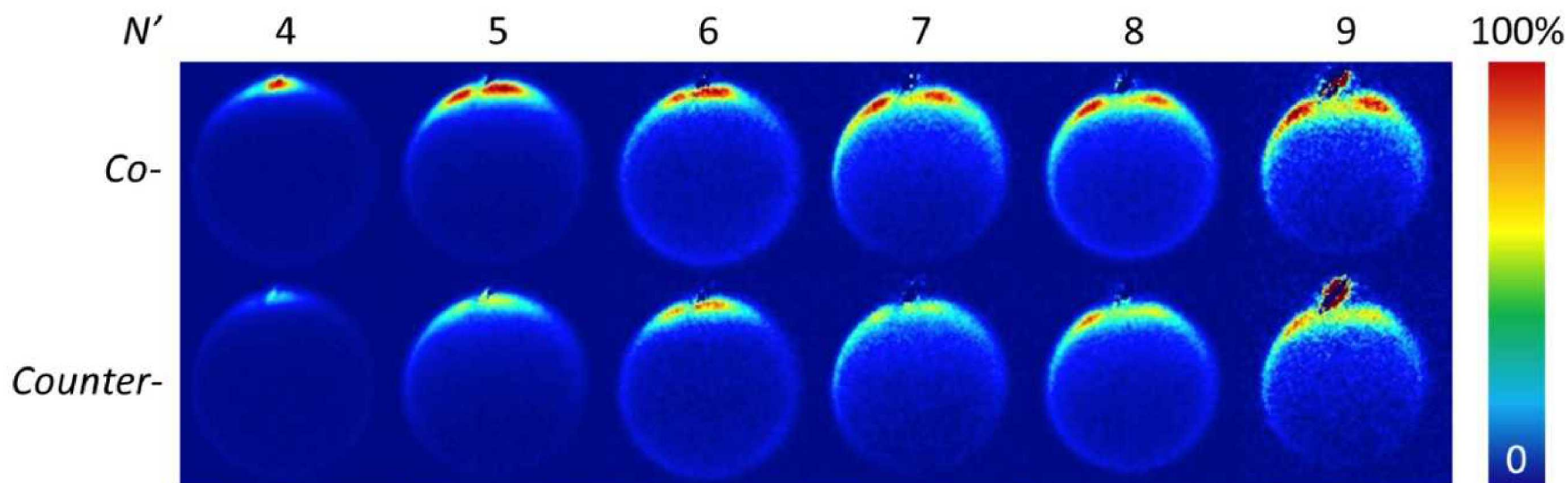
Detect different N states and compare
Co-rotating
versus
Counter rotating



Hyperfine Depolarization Makes the Orientation Oscillate but an Average Value of $A_{10} = 0.6$ is Observed

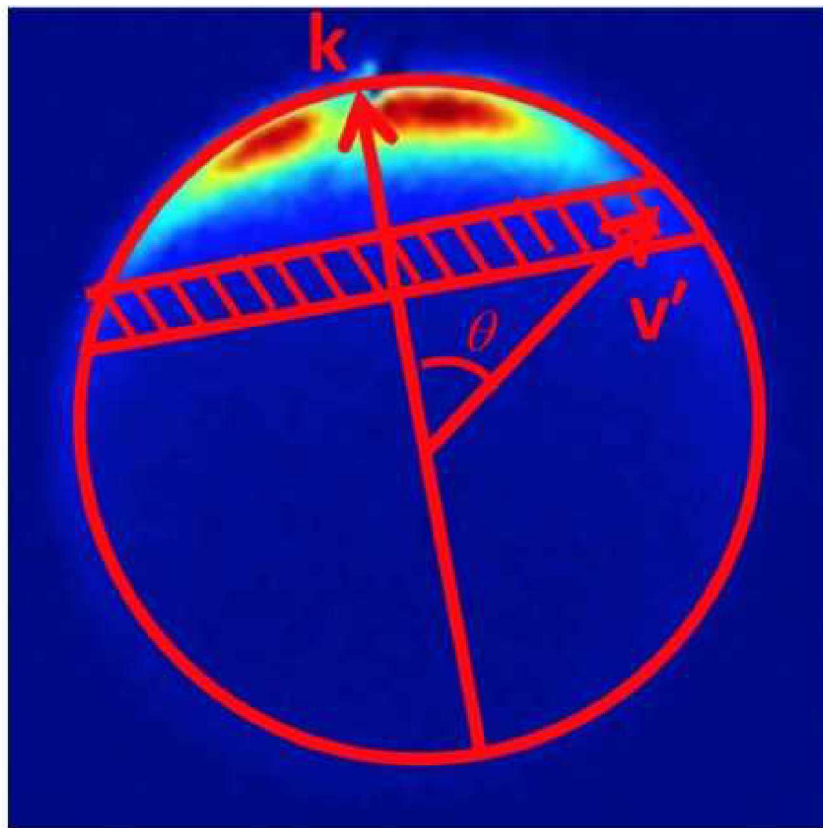


Images of Co- and Counter- Rotating NO After Collision With NO($N=2$)



Scattering images for each final N' in the two geometries. Top row: co-rotating geometry, bottom row: counter-rotating geometry.

The intensity of the images as a function of azimuthal angle are recorded and compared



$$C = \frac{I_{co} - I_{counter}}{I_{co} + I_{counter}}$$

$$C = \frac{3h^{\{1\}}(j)g^{\{1\}}(j,t)A_0^{\{1\}}}{2 - h^{\{2\}}(j)g^{\{2\}}(j,t)A_0^{\{2\}}}$$

C is measured and compared to apse model and Quantum Calculations.

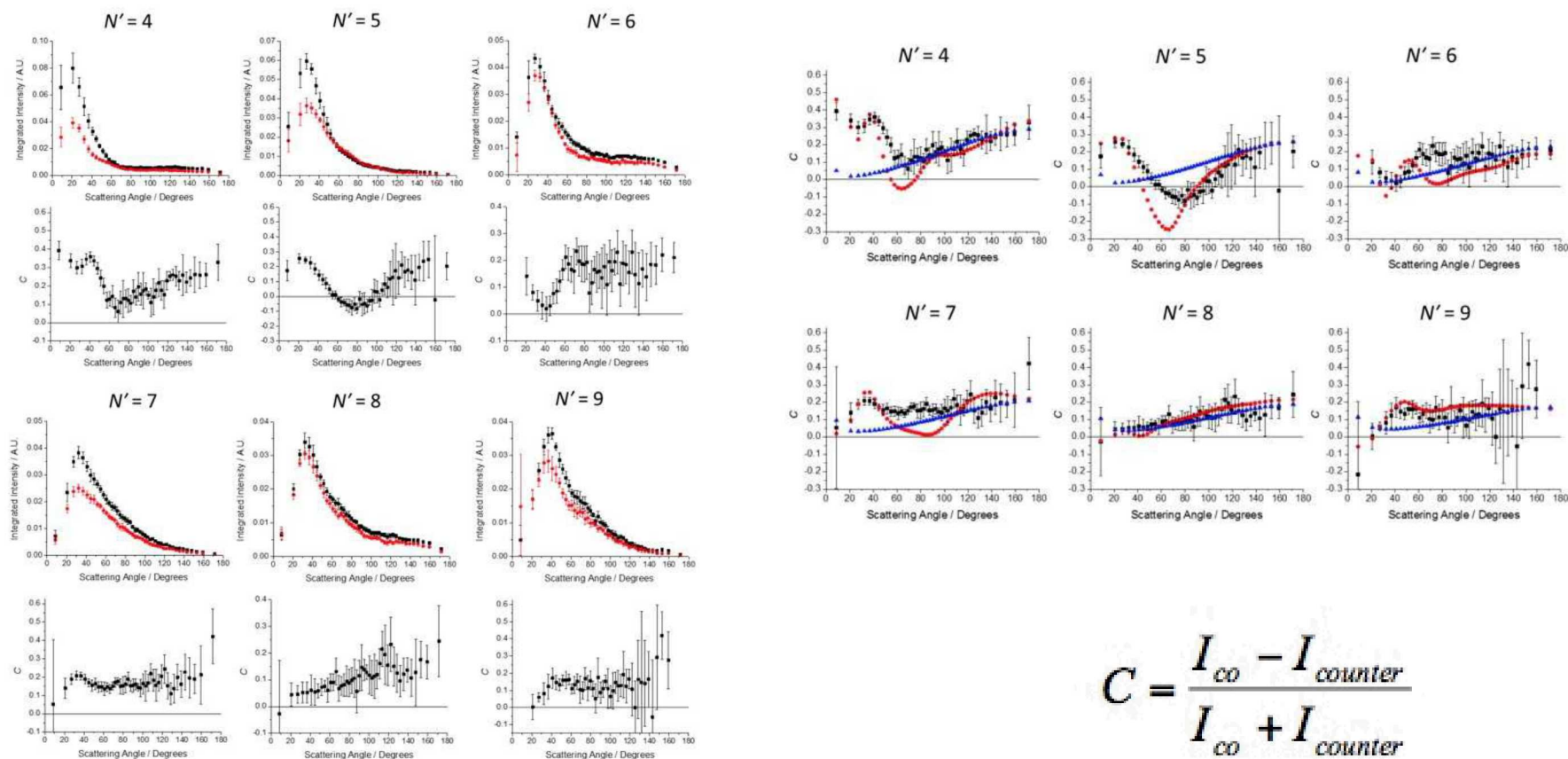
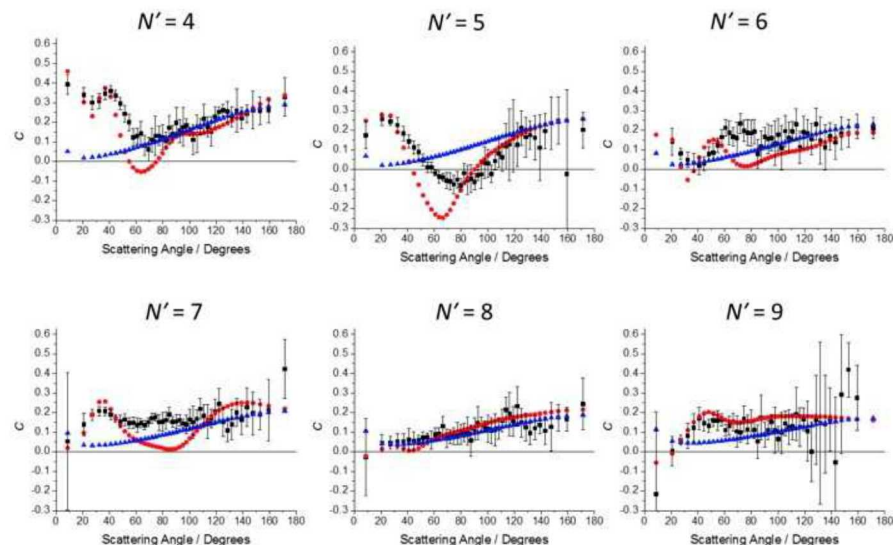
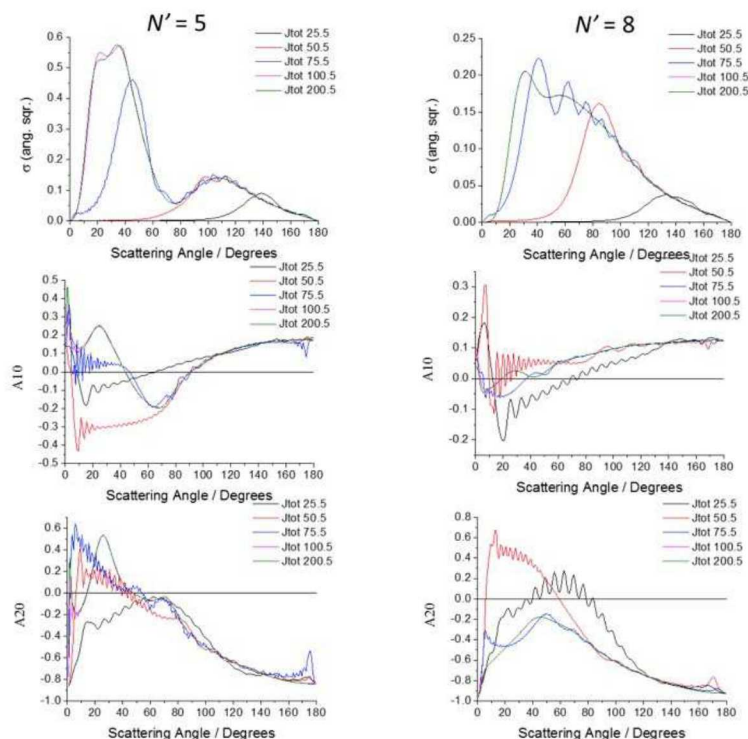


Figure 6: Integrals and values of C for each quantum state. For each quantum state the integrals for the co- and counter rotating geometries are shown in the upper plot, in black and red, respectively. The values of C are plotted in the lower plot.

$$C = \frac{I_{co} - I_{counter}}{I_{co} + I_{counter}}$$

C values are seen to be dominated by the A10 moments.



$$C = \frac{3h^{[1]}(j)g^{[1]}(j,t)A_0^{[1]}}{2 - h^{[2]}(j)g^{[2]}(j,t)A_0^{[2]}}$$

$$C = \frac{I_{co} - I_{counter}}{I_{co} + I_{counter}}$$

Conclusions

Oriented collisions allow one to test the details of the potential energy surfaces
And inform the scattering calculations.

Co-rotating products are more probable than counter-rotating for low values of N .

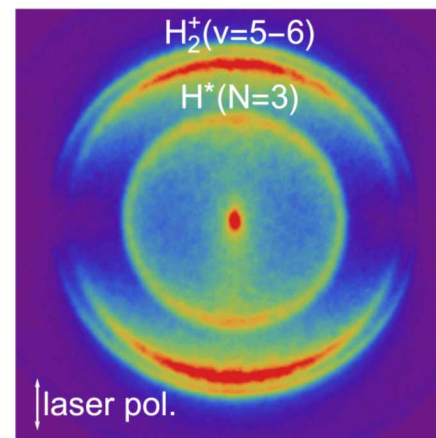
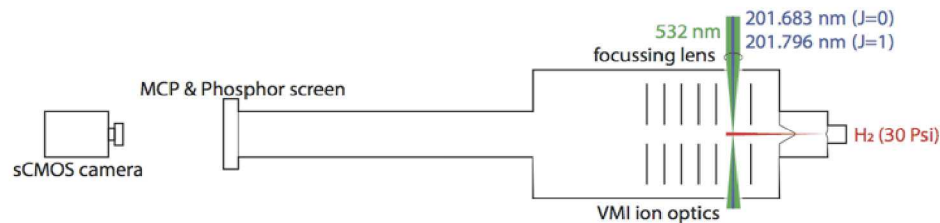
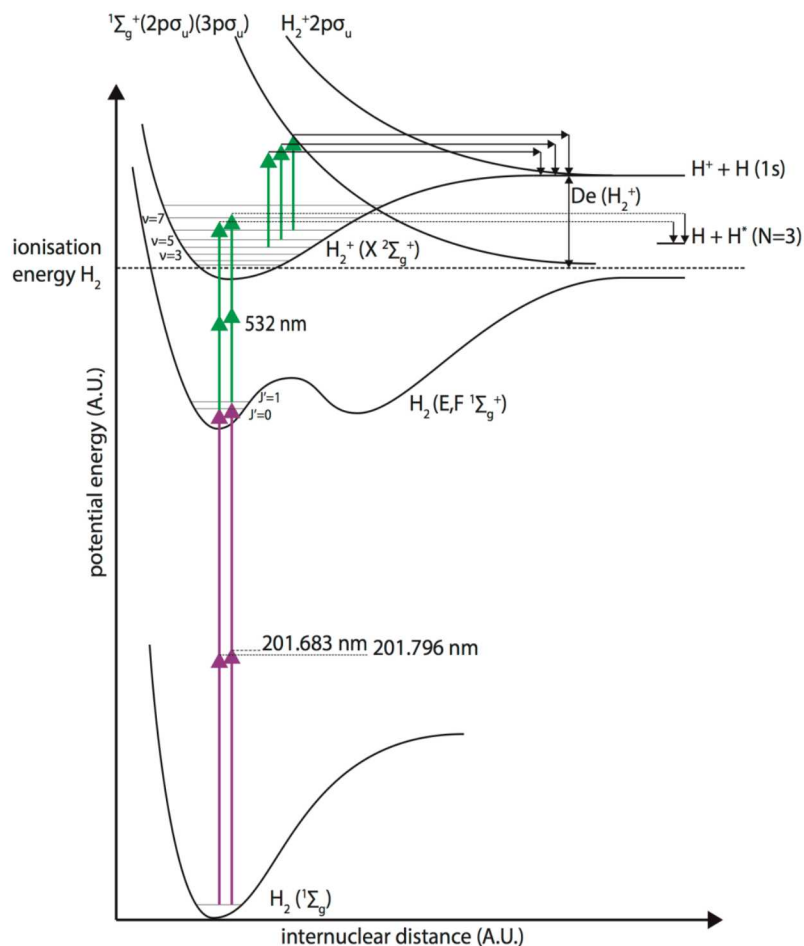
This is a four vector correlation.

H₂ Alignment utilizing double resonance and studied with Velocity-mapped imaging of photo-dissociation products

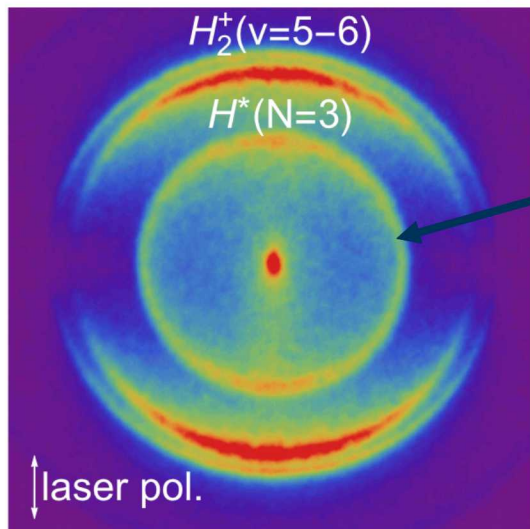
- The distribution of products VMI imaging provides a measurement of the degree of molecular alignment prior to photodissociation.
- Even as the simplest molecule, H₂ in strong laser fields remains a challenge for high level quantum theory.
- Single quantum states may be excited with narrowband tunable lasers, providing a high level of control, minimizing the number of coupled rotational states, and thus simplifying the analysis of data.



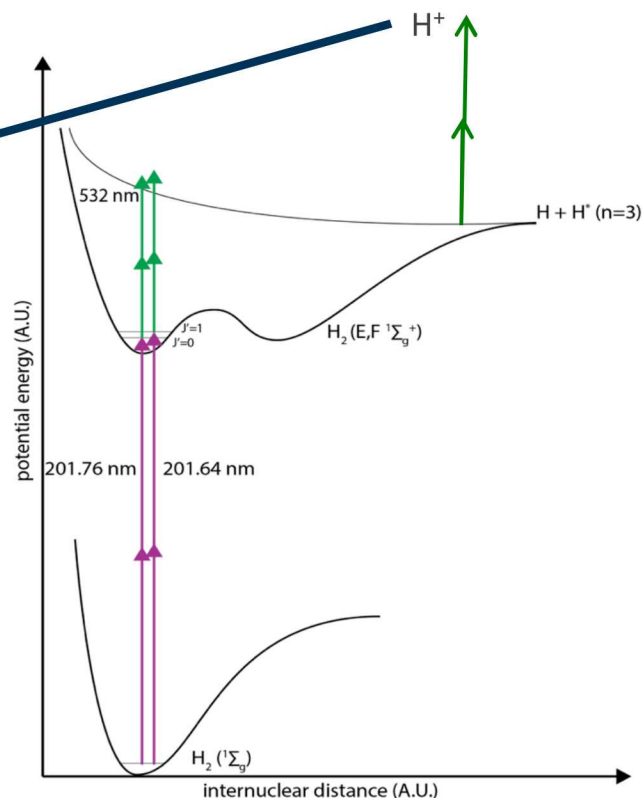
To measure the alignment of the H_2 (E,F) we use a 201.7-nm laser beam to produce the ro-vibrational state we desire, then align and detect with 532-nm light.



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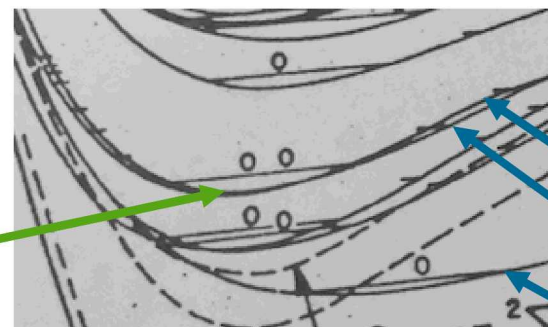
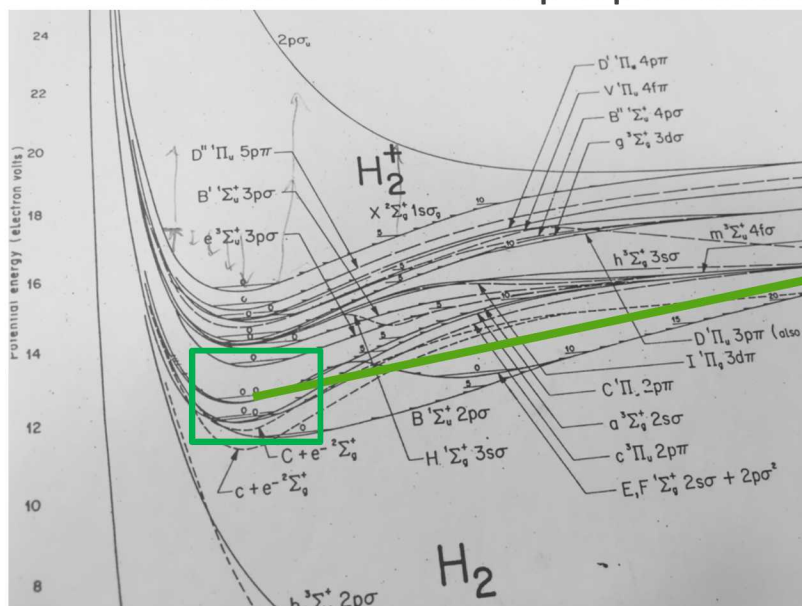


This is the image you obtain of H^+ coming from 532-nm Photo-dissociation of the H_2 EF ($v=0, j=0$) quantum state.



The E,F state lies on top of the B and C states. Those States Couple to the E,F state via the Electronic Dipole Operator.

Coupling to the B state is along the molecular axis and to the C state is perpendicular to molecular axis

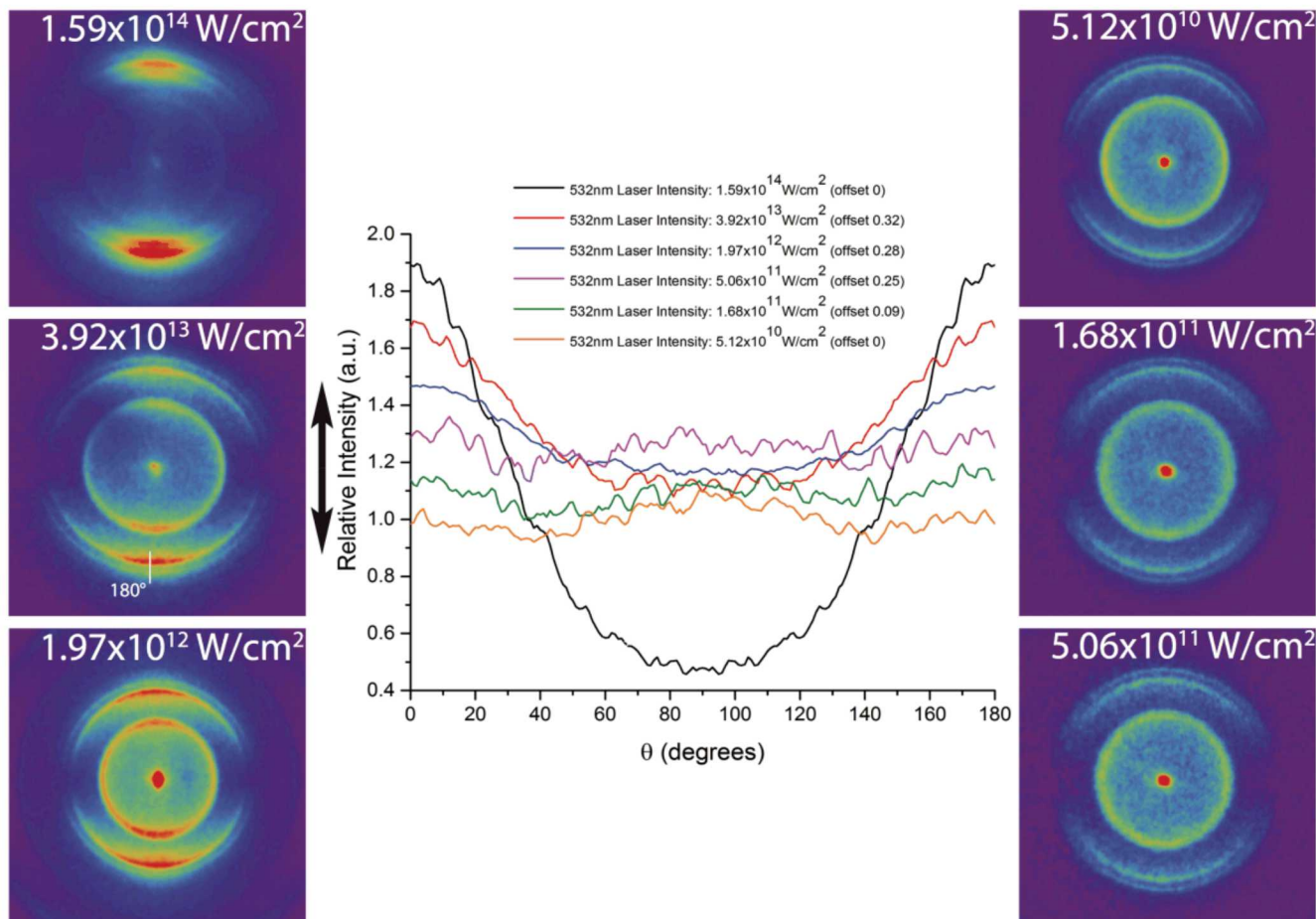


The Electric field of the laser beam mixes the EF, B and C states creating a superposition state That has 100 times higher polarizability difference than ground state H_2

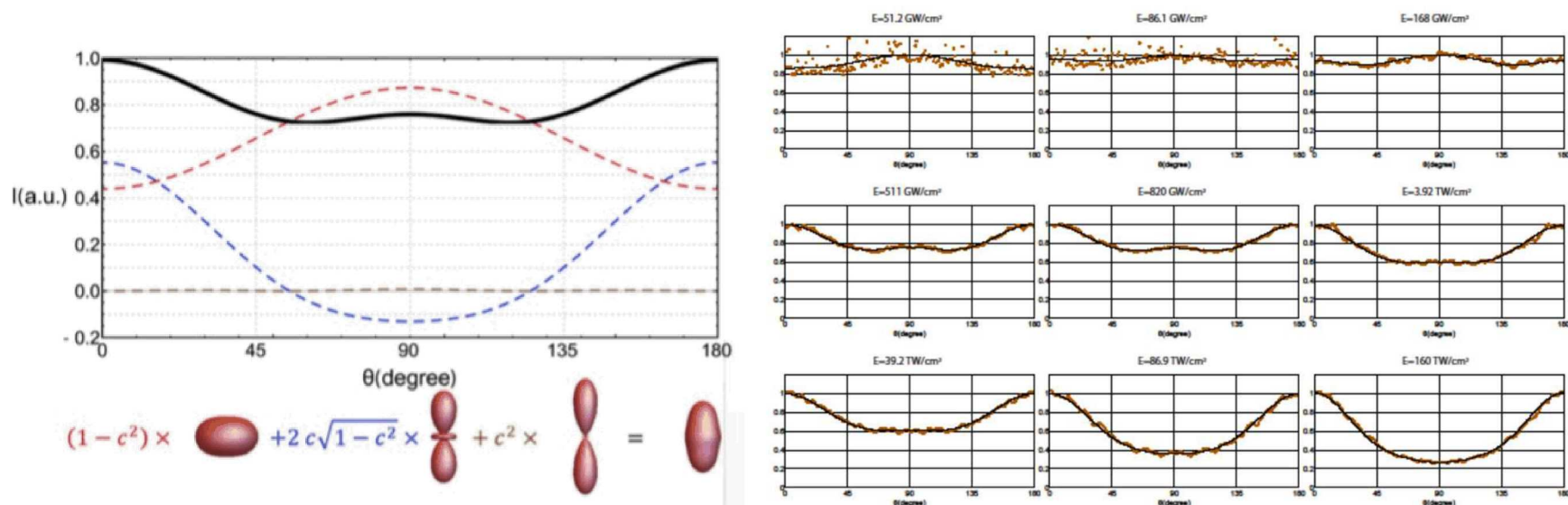


The alignment of the H_2 by intense 532-nm light is quantified by taking images at different 532-nm intensities.

H_2 (E,F; J = 0)

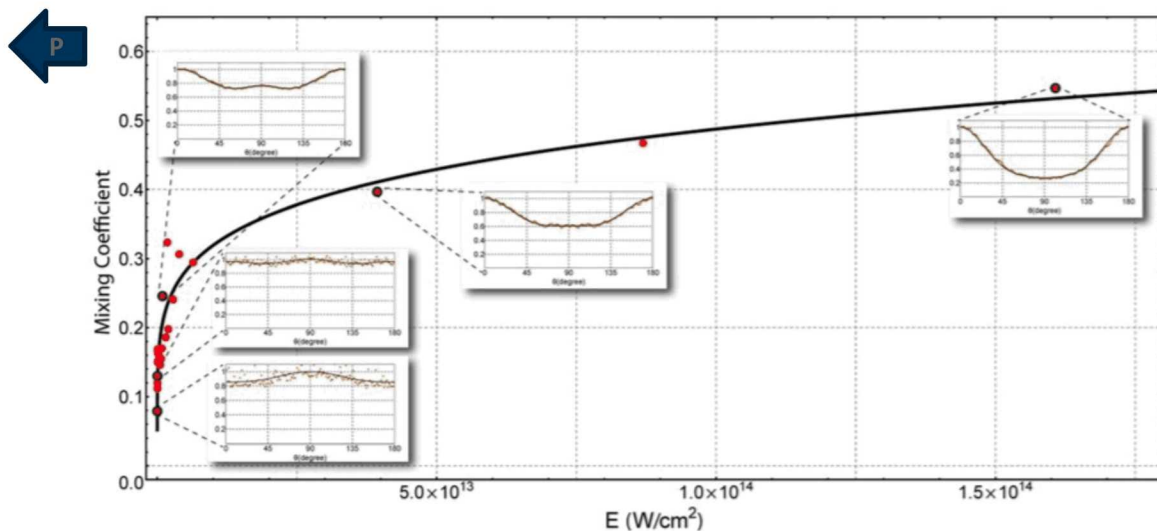


The alignment of the H₂ (E,F J=0) as a function of 532-nm laser intensity is plotted and fitted using an Adiabatic Alignment model Employing J=0, J=2 and a cross term.



The angular distribution of the inner ring can be well fit with a model that involves mixing of rotational wave functions as a function of laser intensity. At the lowest laser power at which we observe signal we can see evidence of H₂(E,F J=2) quantum state mixing

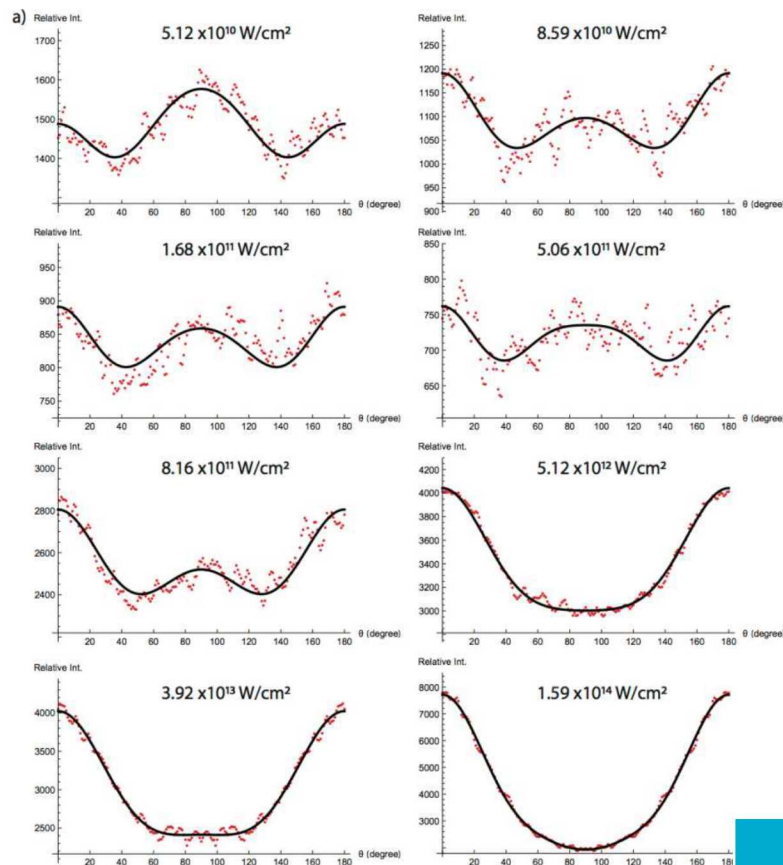
Analysis of the alignment of the H₂ (E,F J=0 and J=1) as a function of 532-nm laser intensity resulted in polarizability measurement



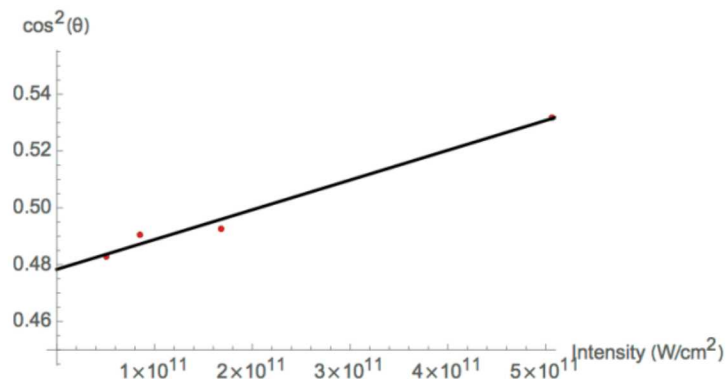
Mixing coefficient is measure of the mixing between the different rotational states of the H₂ molecule. From the slope of this line at low laser intensity we extract the polarizability

Proposed Work: To remove ambiguity for comparison to simulations, two different wavelengths will be used to prepare and probe the molecular alignment.

By modeling the data, one can extract the polarizability anisotropy of H₂(E,F) state when irradiated with 532-nm light.



$$\langle \cos^2(\theta) \rangle_{\Delta\omega} = \frac{1}{3} + \boxed{\frac{10^{13} \delta\alpha}{135 B}} I$$



J state	Linear Fit		Anisotropy (a.u.)
	Intercept	slope	
J0	0.5325	1.032E-13	312 ± 82

How Polarizable is the H₂(E,F) state?

J state	Linear Fit		Anisotropy (a.u.)
	Intercept	slope	
J0	0.5325	1.032E-13	312 ± 82

Theoretical values in static fields:

$\delta\alpha$ ground state = 1.9 a.u.*

$\delta\alpha$ E,F state = -9600 a.u.†

* E. Ishiguro, et al; Proc. Phys. Soc. A 65, 178 (1951)

† Komasa, J.; Adv. Quant. Chem. 48, 151 (2005).

- Anisotropy ($\delta\alpha$) values of 312 a.u. is almost 46, Å³, respectively.
- By comparison CS₂ has 10 Å³, I₂ has 7 Å³ and 1,4 diiodobenzene has about 18 Å³.
- The most polarizable state of any molecule to date.

Conclusions

Doing dynamics on transient excited states allows one to have better preparation of the initial conditions.

We have demonstrated this preparation for NO(A) and H₂ (EF) both having 200 ns lifetimes.

Dynamics not observed on the ground state are studied. In Hydrogen he mixing with nearby electronic states creates superposition states with strange characteristics. For NO (A) an oscillating orientation of the J vector is observed and utilized for collisional studies.

