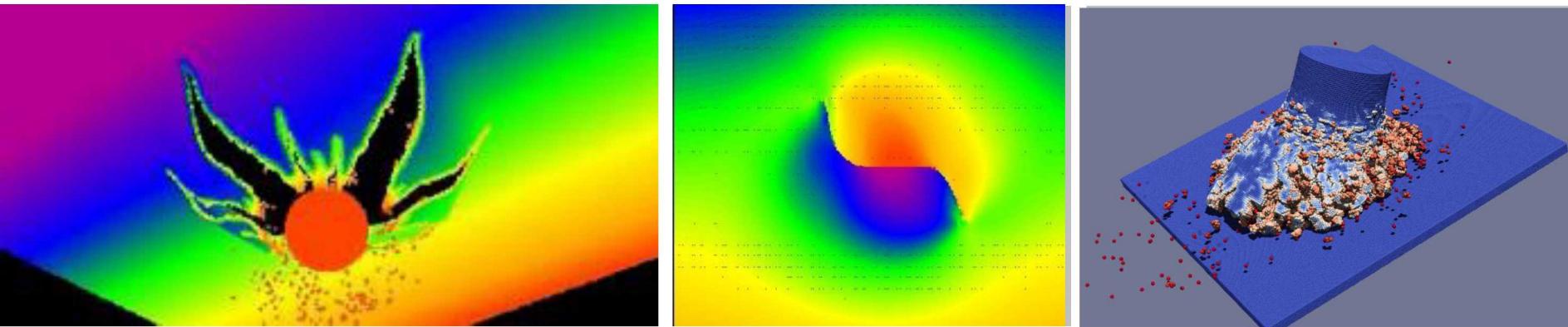


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The Mechanics of Unstable Peridynamic Materials

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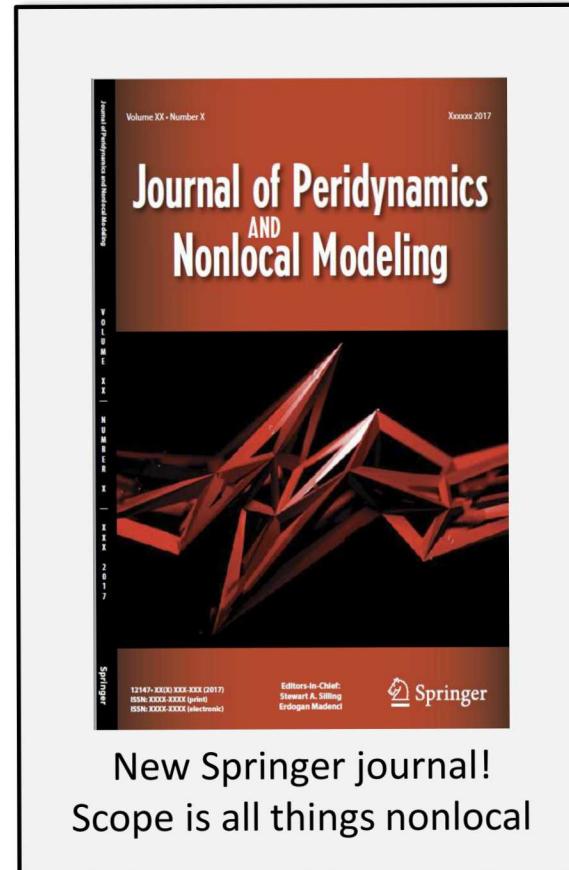
SIAM Central States Section Meeting, Ames, IA, October 19, 2019



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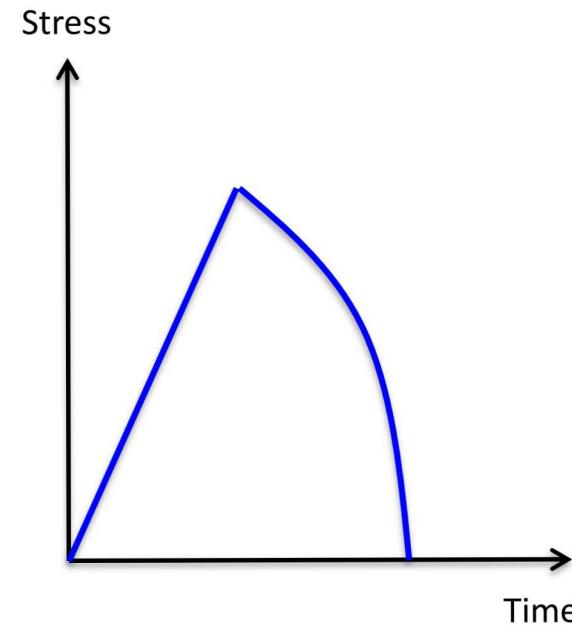
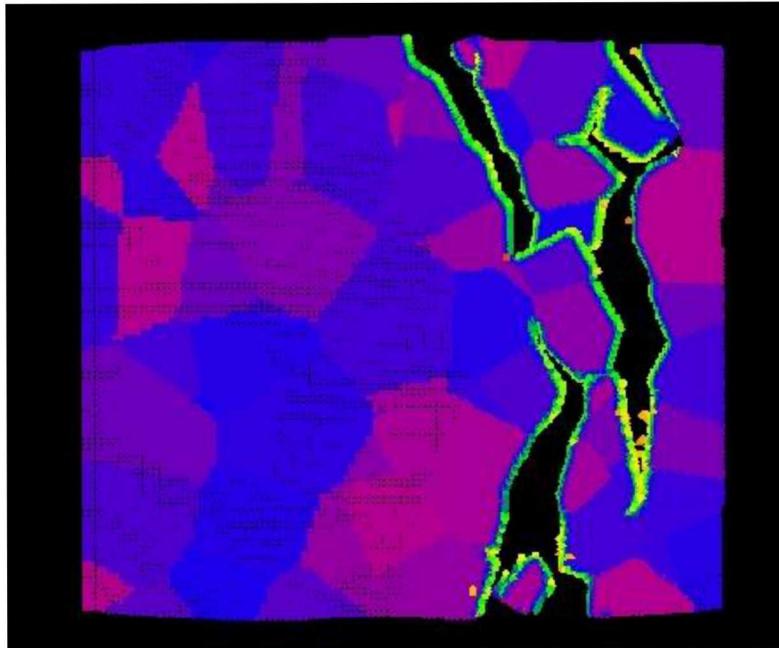
Outline

- Kinetics of material failure
- Peridynamic theory
- Non-propagating waveforms
- Growth of a small pulse
- Time to failure
- Rate effect



Failure kinetics: time it takes to fail

- If we pull slowly on a test specimen, it breaks when it reaches some strain.
- Once it reaches the failure strain, how much **time** does it take to fail?
- Failure involves small-scale processes that occur at a finite rate.
- How to model this in a homogenized material?
- **We will study this problem using an unstable peridynamic material model.**



Peridynamic mechanics

- General peridynamic equation of motion:

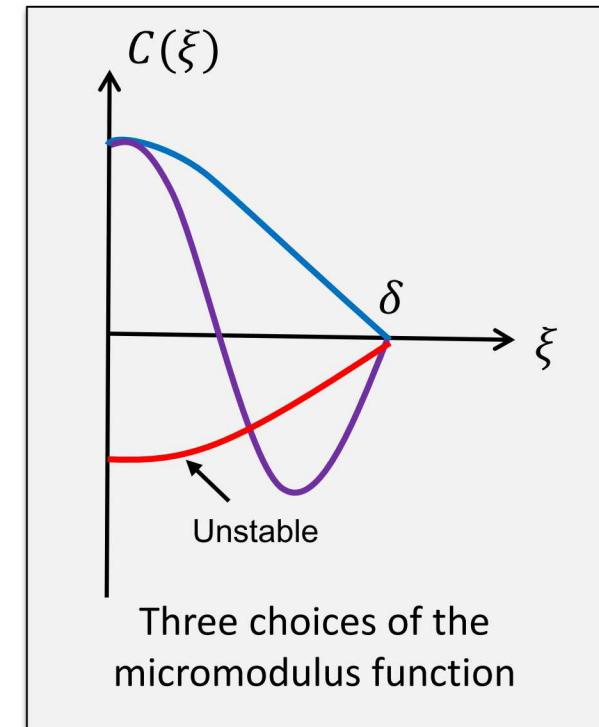
$$\rho(x)\ddot{u}(x, t) = \int_{-\delta}^{\delta} f(\eta, \xi) d\xi + b(x, t), \quad \xi = q - x, \quad \eta = u(q, t) - u(x, t)$$

where ρ =density, u =displacement, b =external force density, δ =horizon, ξ =bond, η =bond elongation.

- The horizon is a cutoff distance for nonlocal force interactions.
- Linear microelastic material:

$$f(\eta, \xi) = C(\xi)\eta, \quad C(\xi) = \text{micromodulus.}$$

- The micromodulus contains the material properties.
- Assume an infinite homogeneous bar with zero external force.



Dispersion and wave speeds

- Assume a wave of the form

$$u(x, t) = A e^{i(kx - \omega t)}$$

where A =amplitude, k =wavenumber, ω =frequency.

- Equation of motion:

$$\rho \ddot{u}(x, t) = \int_{-\delta}^{\delta} C(\xi) [u(x + \xi, t) - u(x, t)] d\xi$$

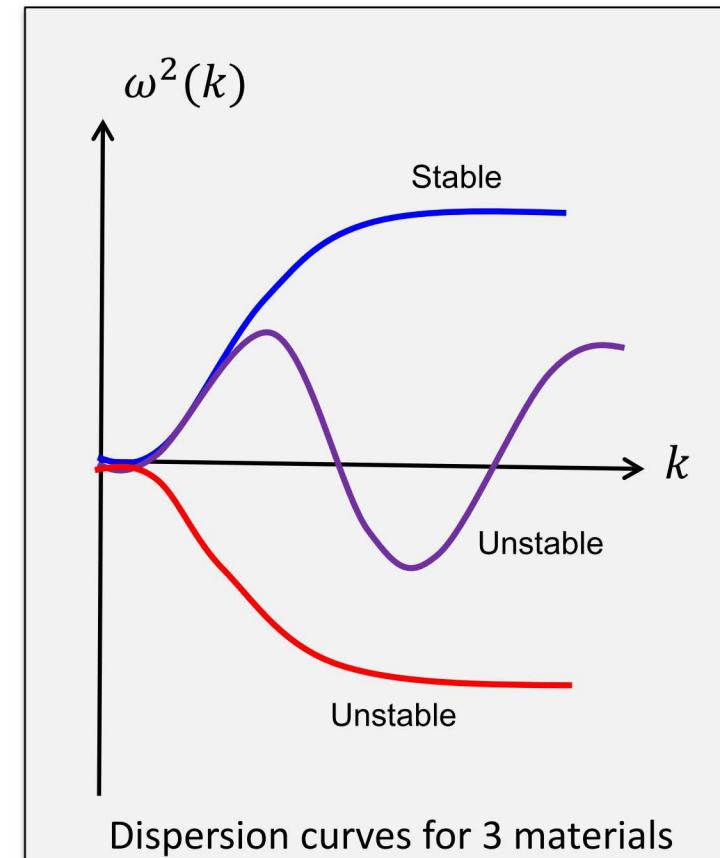
- leads to

$$-\rho \omega^2 = \int_{-\delta}^{\delta} C(\xi) [e^{ik\xi} - 1] d\xi$$

- Therefore the dispersion relation is

$$\omega(k) = \sqrt{\frac{P - \bar{C}(k)}{\rho}}, \quad P = \bar{C}(0)$$

where $\bar{C}(k)$ is the Fourier transform of $C(\xi)$.



Dispersion curves for 3 materials

Unstable waveforms grow exponentially

- Initial data in the infinite bar:

$$u(x, 0) = A \cos kx, \quad \dot{u}(x, 0) = 0 \quad \forall x.$$

- If $\omega^2(k) > 0$:

$$u(x, t) = \frac{A}{2} [\cos(kx - \omega(k)t) + \cos(kx + \omega(k)t)].$$

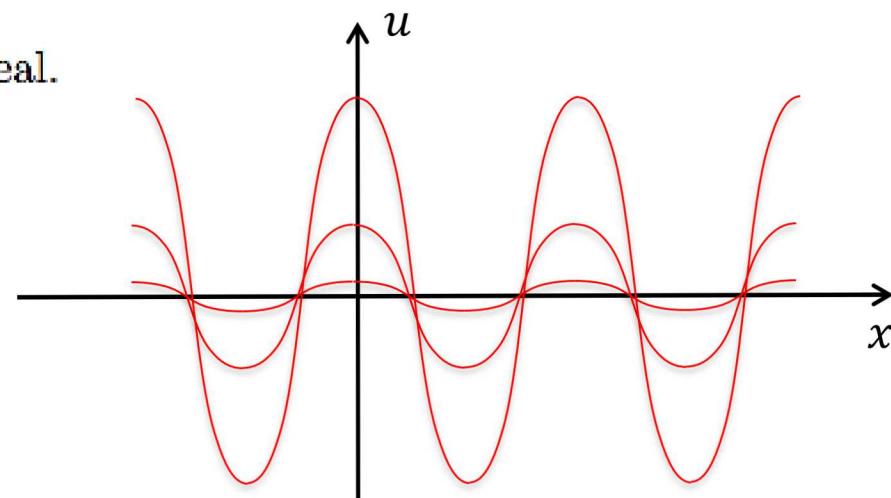
- If $\omega^2(k) < 0$:

$$u(x, t) = A \cos(kx) \cosh(\lambda(k)t)$$

where

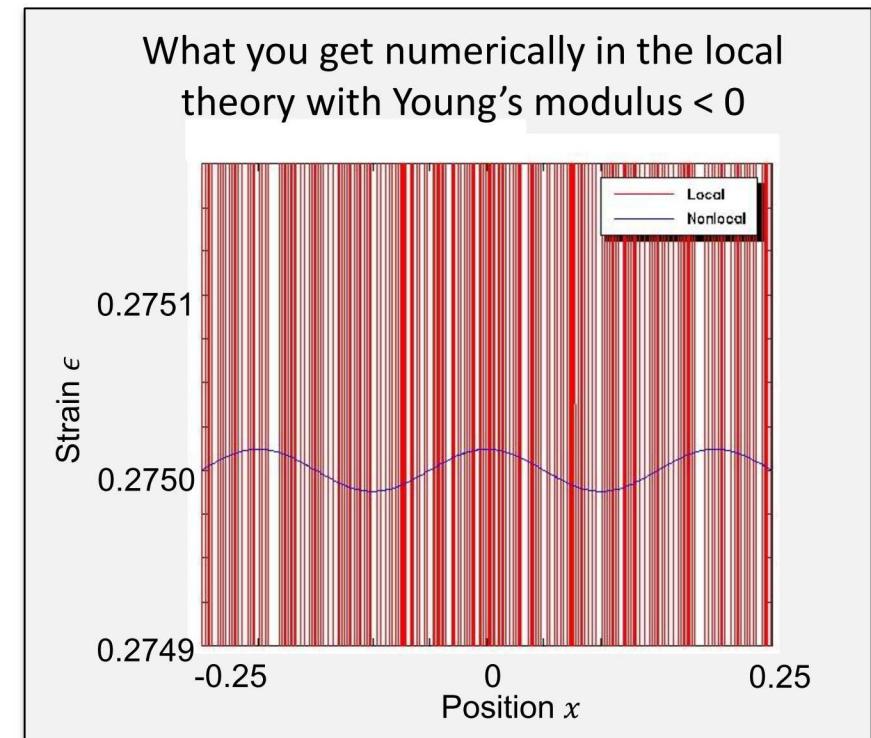
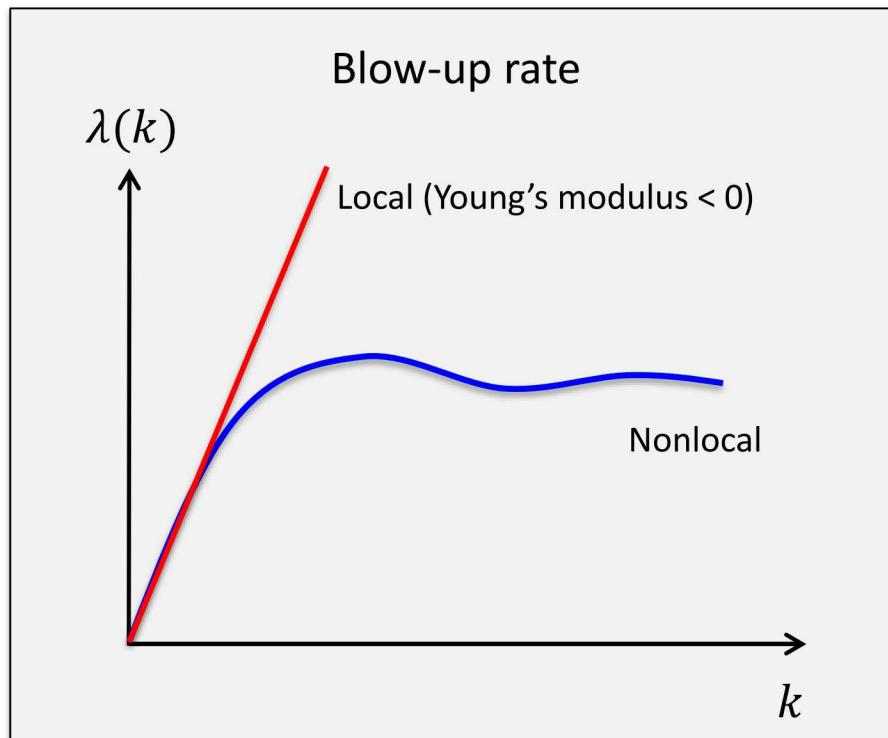
$$\lambda(k) = \sqrt{-\omega^2} \quad \text{real.}$$

- λ is called the *blow-up rate*.



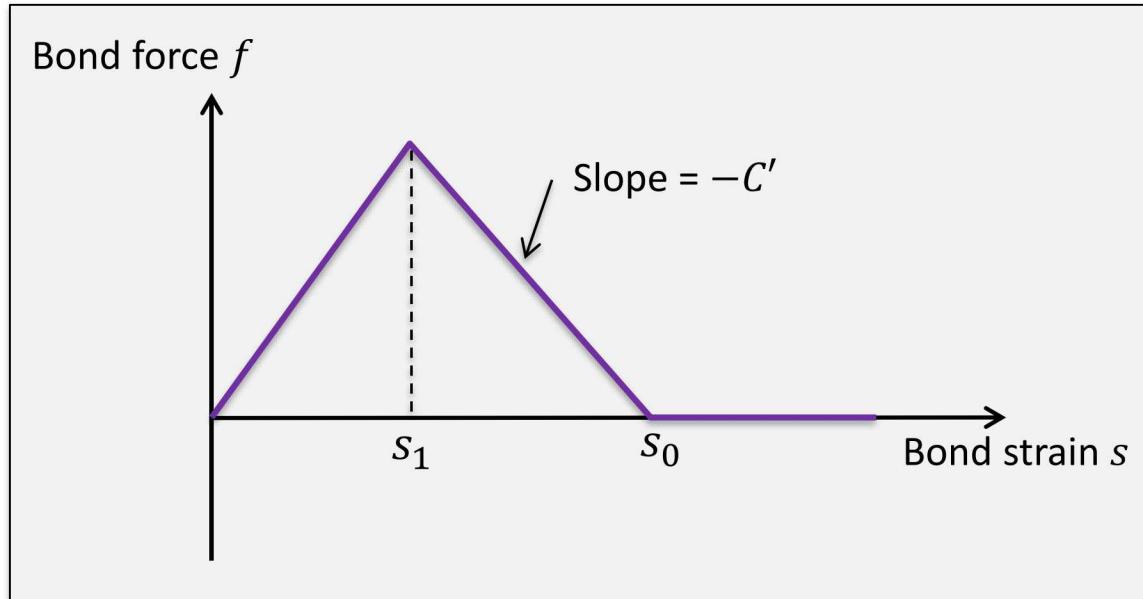
Unstable mechanics is fundamentally different in local & peridynamic theories

- Blow-up rate is bounded in peridynamics for “most” materials for all k .
- Unbounded in the local theory for large k .
 - Hadamard instability (loss of ellipticity, “imaginary wave velocity”).
 - Initial value problems generally don’t have sensible solutions.



Failure kinetics: a peridynamic material that changes from stable to unstable

- As the microelastic material is stretched, the bond force passes through a peak into an unstable region.
- Lipton showed problems are well posed. They may spontaneously produce Griffith cracks.

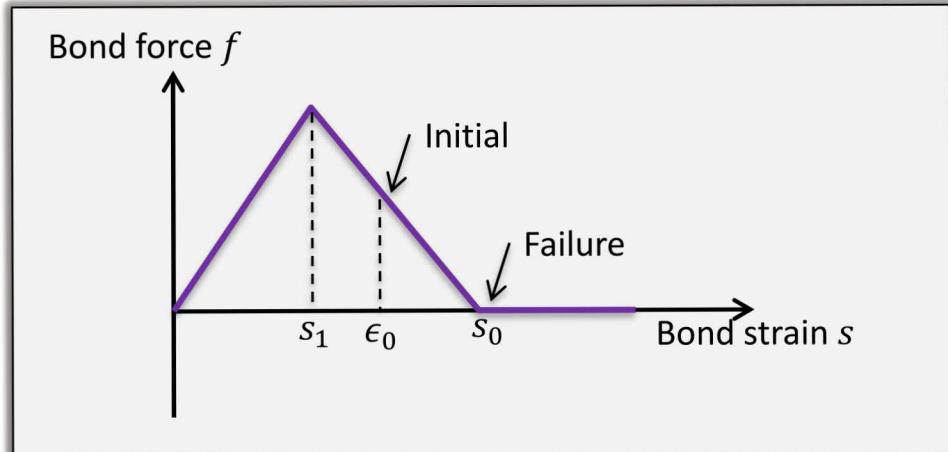


- K. Dayal & K. Bhattacharya *JMPS* (2006)
- R. Lipton, *Journal of Elasticity* (2014, 2015)
- R. Lipton, R. B. Lehoucq, & P. K. Jha, *J. Peridynamics & Nonlocal Modeling* (2019)
- R. Lipton, E. Said, & P. K. Jha. in *Handbook of Nonlocal Continuum Mechanics for Materials and Structures* (2018).
- P. K. Jha, R. Lipton, *IJNME* (2018)

Time to failure with initial data inside the unstable region

- Material model for bond force as shown.
- Initial data in the infinite bar:

$$u(x, 0) = \epsilon_0 x + \frac{\tilde{\epsilon}_0}{k} \cos kx, \quad \dot{u}(x, 0) = 0 \quad \forall x$$



where

$$\epsilon_0 = \frac{s_0 + s_1}{2}$$

and $\tilde{\epsilon}_0$ is a small strain ($\approx 10^{-6} s_0$).

- Use the blow-up rate $\lambda(k)$ to find how much time it takes for any bond to fail (bond strain $\geq s_0$).
- Result:

$$t_{fail} = \frac{1}{\lambda(k)} \cosh^{-1} \left(\frac{s_0 - \epsilon_0}{\tilde{\epsilon}_0} \right).$$

Initial data with a spectrum of wavenumbers

- Recall for a single wavenumber k in the unstable region:

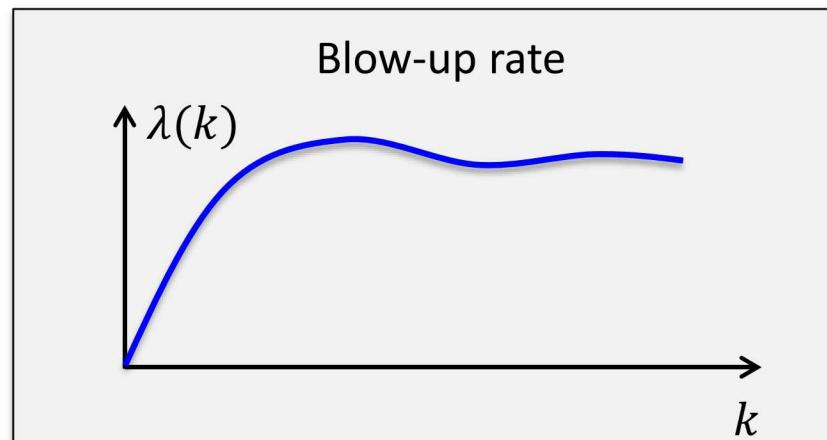
$$u(x, t) = A \cos(kx) \cosh(\lambda(k)t).$$

- For general initial data (even function of x for simplicity):

$$u(x, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \bar{u}_0(k) \cos(kx) \cosh(\lambda(k)t) dk$$

where \bar{u}_0 is the Fourier transform of the initial displacement.

- Usually large k wins the race among Fourier components over long time periods.



Time to failure with a Gaussian initial perturbation in strain

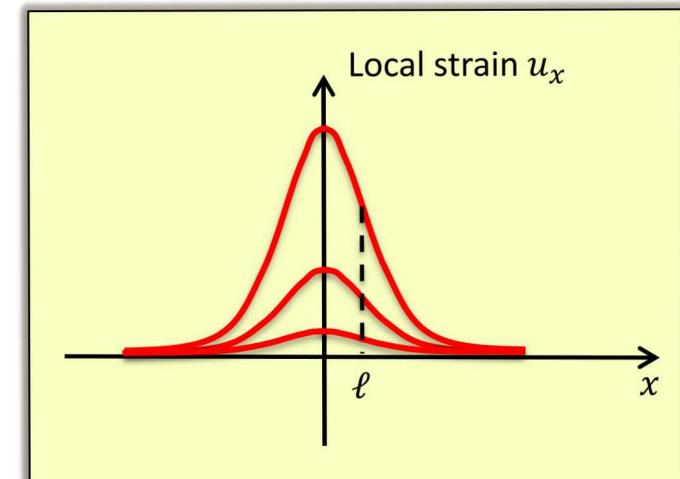
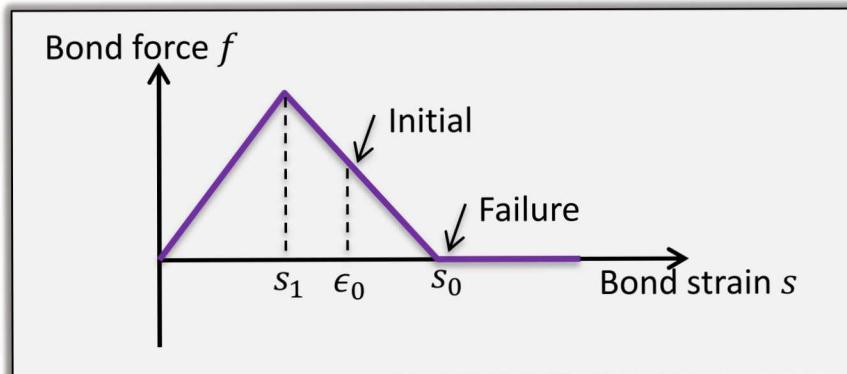
- Initial data in the infinite bar:

$$u(x, 0) = \epsilon_0 x + \frac{\sqrt{\pi} \ell \tilde{\epsilon}_0}{2} \operatorname{erf}\left(\frac{x}{\ell}\right), \quad \dot{u}(x, 0) = 0 \quad \forall x$$

where $\tilde{\epsilon}_0$ is a small strain, erf is the error function and $\ell \ll \delta$ is the half-width of the Gaussian.

- This differs from the previous example in that there are many Fourier components.
- To solve, use the Fourier transform equation of motion and initial data.
- Result:

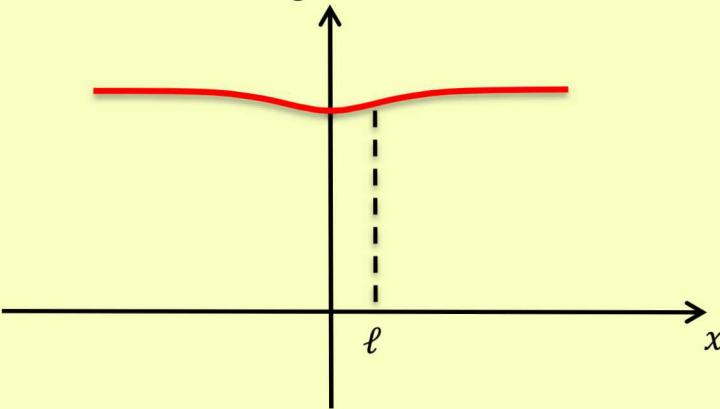
$$t_{fail} \approx \frac{\delta}{C} \sqrt{\frac{s_0 - s_1}{6s_1}} \cosh^{-1} \left(\frac{s_0 - \epsilon_0}{\tilde{\epsilon}_0} \right).$$



Stretching of a bar with a small defect

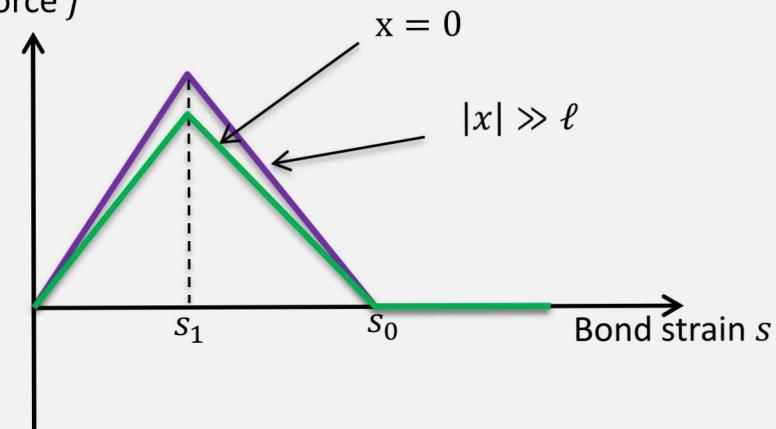
- Bar is slightly more compliant in a small region near $x = 0$.
- Bar is initially stretching at a constant rate r_0 .
- Effect of the defect is amplified as the strain passes the peak in the material model.

Young's modulus



Bar with a small heterogeneity

Bond force f

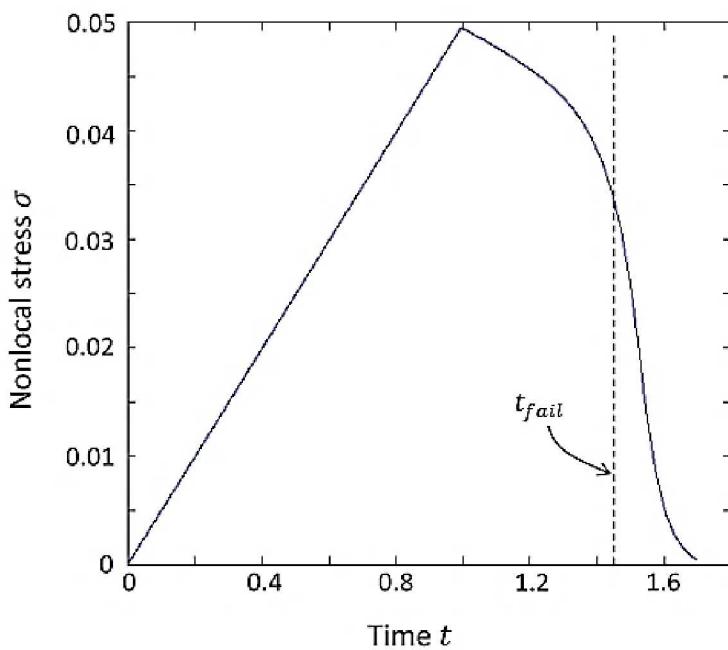


Position-dependent material model

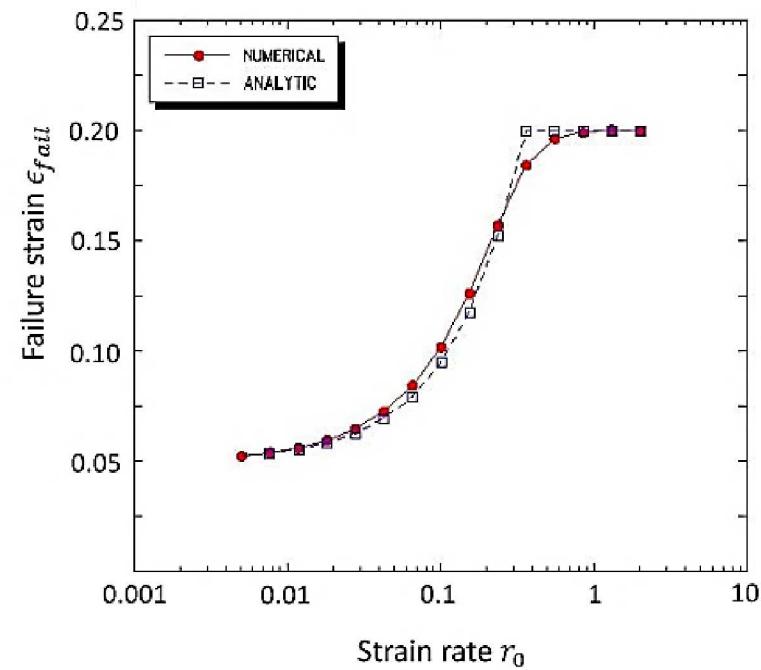
Defect nucleates failure

- Finiteness of the failure time after the peak results in a strain rate effect on the remote failure strain.
- Approximate analysis and Emu code give nearly the same results.

Stress at $x = 0$

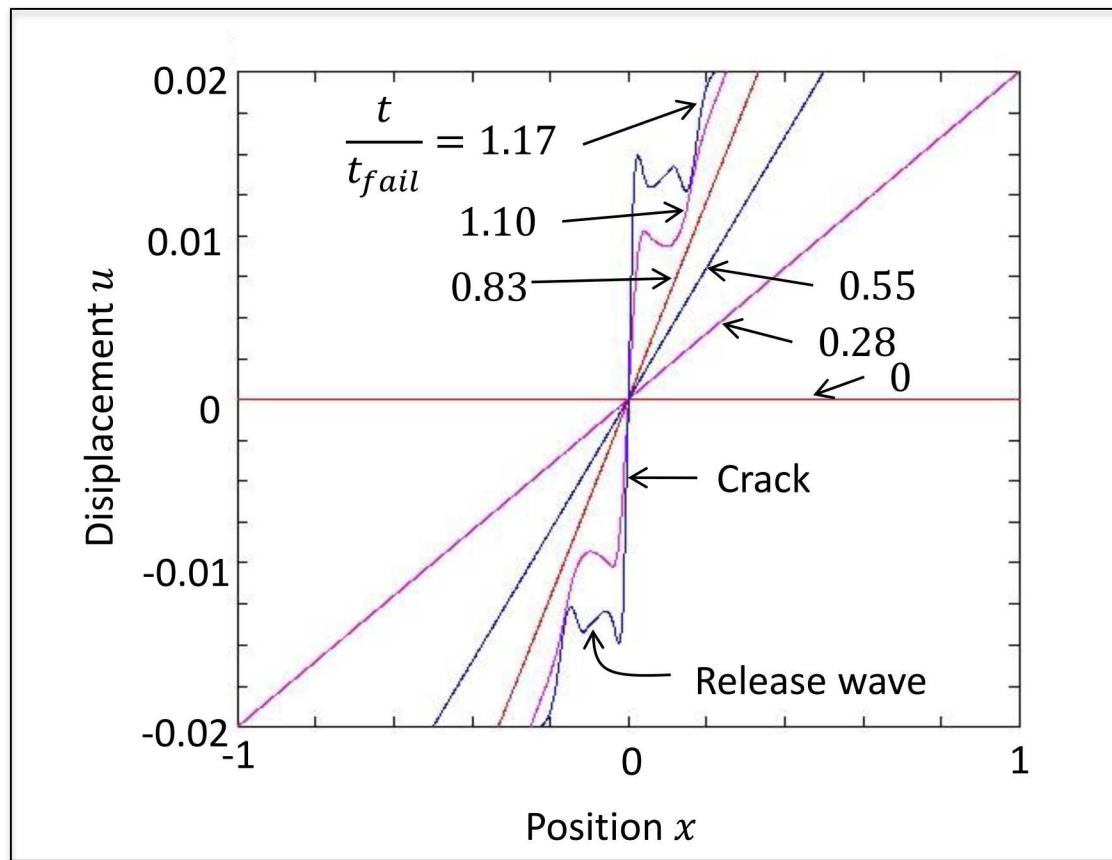


Remote strain at time of failure



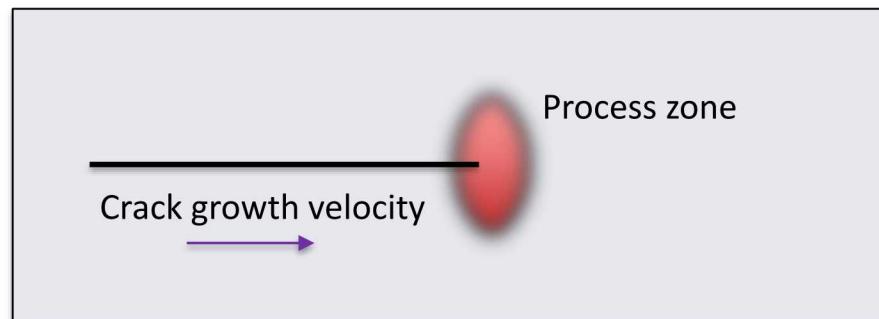
Defect nucleates failure: Numerical simulation

- A crack opens up at $x = 0$ after the bar enters the unstable region.



Discussion

- We get a kinetic effect on material failure “for free” with a nonconvex peridynamic material model.
- Conjectures:
 - This stuff happens near a growing crack tip.
 - Size of the (unstable) process zone influences the blow-up rate.
- Need more research on numerics further for unstable peridynamic materials.
- Need more general theories of convexity and material stability with peridynamics.



- This work: S. S., Kinetics of failure in an elastic peridynamic material, SAND2019-5984
- P. K. Jha & R. Lipton, IJNME (2018)
- P. K. Jha & R. Lipton, CMAME (2019)