

Bevel Edge Termination for Vertical GaN Power Diodes

Andrew T. Binder*, Jeramy R. Dickerson, Mary H. Crawford, Greg W. Pickrell, Andrew A. Allerman, Paul Sharps, Robert J. Kaplar

Sandia National Laboratories, Albuquerque, NM 87185, USA

*abinder@sandia.gov

Abstract—Edge termination for vertical power devices presents a significant challenge, as improper termination can result in devices with a breakdown voltage significantly less than the ideal infinite-planar case. Edge termination for vertical GaN devices is particularly challenging due to limitations in ion implantation for GaN, and as such this work investigates a bevel edge termination technique that does not require implantation and has proven to be effective for Si and SiC power devices. However, due to key differences between GaN versus Si and SiC p-n junctions (specifically, a grown versus an implanted junction), this technology needs to be reevaluated for GaN. Simulation results suggest that by leveraging the effective bevel angle relationship, a 10-15° physical bevel angle can yield devices with 85-90% of the ideal breakdown voltage. Results are presented for a negative bevel edge termination on an ideally 2 kV vertical GaN p-n diode.

Keywords—vertical GaN diode; edge termination; junction termination extension; power devices

I. INTRODUCTION

Optimization of breakdown in vertical power devices is often limited by field crowding at the periphery of the device. For an ideal planar device, the breakdown is limited by the peak electric field in the bulk, yielding a much higher breakdown voltage than for a device limited by a large surface electric field. To address this issue, special junction terminations around the edges of the power device can be formed so that the depletion near the surface is increased, thereby reducing field crowding. The most common edge termination structures in Si and SiC (junction termination extensions {JTEs} and guard rings) are created by selective-area doping using ion implantation. However, Mg implantation in GaN aimed at achieving selective-area p-type doping suffers from low activation efficiency and a limited effective range of implantation energy and dose [1], making this method challenging to use in GaN power devices. An alternate junction termination method that does not require selective-area doping is the negative bevel edge termination which has been demonstrated both in Si and SiC power devices [2], [3]. However, epitaxially-grown p-n junctions in GaN are much more abrupt than implanted

junctions in Si and SiC power devices [4], [5], resulting in important differences in how these regions deplete, impacting the effectiveness of a bevel termination.

In this work, a vertical GaN p-n diode is used to evaluate the performance of a negative bevel edge termination. It should be noted that this edge termination technique is not in any way limited to this particular device, but can be applied to a wide range of devices.

II. THEORETICAL BACKGROUND

A. Single-Zone/Multi-Zone JTE

The design of a single-zone or multi-zone JTE hinges on Gauss's law and the charge enclosed in the Gaussian surface (the JTE) [6], [7]. According to Gauss's law, the electric flux through any closed surface is equal to the total charge inside divided by the permittivity as shown in Eq. 1:

$$\phi = E \cdot A = \frac{\sigma \cdot A}{\epsilon_r \cdot \epsilon_0} \quad (1)$$

where ϕ is the electric flux, σ is the charge per unit area, ϵ is the permittivity, A is the area of the Gaussian surface, and E is the electric field.

In this manner, the charge inside the JTE acts as a sink for the electric field in a power device at breakdown, causing the field lines to terminate into the JTE and not crowd at the periphery of the device. Premature breakdown occurs in unterminated devices when field crowding at the periphery of the device causes the surface E-field to exceed the bulk E-field by a significant margin. In properly terminated devices, at breakdown the E-field at the surface will be equal to the E-field in the bulk thereby maximizing the breakdown voltage of the device. An ideal termination is obtained when the charge in the JTE is designed to sink the critical electric field (E_{crit}).

The disadvantage of a single-zone JTE (SZ-JTE) is that the JTE efficiency (the ratio of the breakdown voltage in a terminated device to the breakdown voltage of an ideal one-dimensional device) is very sensitive to changes in total charge within the Gaussian surface, and therefore sensitive to the thickness and doping density of the JTE [8], [9]. One method to reduce this sensitivity to process variation is to design a multi-zone JTE (MZ-JTE) which employs several JTEs each with a

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79 unique total JTE charge [9]. However, each additional zone
 80 requires additional fabrication steps, further complicating the
 81 process. In addition, both the SZ-JTE and MZ-JTE suffer from
 82 adverse effects of surface charge that strongly depend on
 83 passivation, which needs careful consideration when designing
 84 for the optimal total JTE charge. Surface charge resulting from
 85 passivation can effectively reduce voltage-blocking capability
 86 by 25-50% depending on the net interface state density [10].

87 B. Bevel Edge Termination

88 The analog approach to a step change in charge across
 89 multiple regions (MZ-JTE) is to create a bevel with a continuum
 90 of charge. With a bevel, the charge in the Gaussian surface is
 91 not a discrete value that can be hit or missed due to process
 92 variation, but rather is a continuous range which is less sensitive
 93 to changes in total JTE charge. In direct contrast to the SZ/MZ-
 94 JTE, the effectiveness of the bevel edge termination is also
 95 insensitive to surface charge, due to the continuum of charge in
 96 the bevel.

97 The design for a bevel edge termination is primarily based
 98 on two factors: the ratio of the depletion depth into the JTE to
 99 the depletion into the drift region, and the angle of the bevel. A
 100 negative bevel with a small bevel angle can be used to reduce
 101 the peak surface electric field by increasing the depletion width
 102 on the surface relative to the depletion width in the bulk as
 103 shown in Fig. 1. The increased surface depletion in turn reduces
 104 the E-field at the surface, given that a constant voltage across a
 105 greater depletion width yields a smaller E-field ($E = dv/dx$).
 106 When the bevel angle is sufficiently small the E-field at the
 107 surface will be equal to or less than the E-field in the bulk,
 108 yielding near-ideal breakdown performance in the terminated
 109 device.

110 C. Effective Bevel Angle

111 For a beveled device, the edge termination becomes more
 112 effective for the same bevel angle when the depletion depth into
 113 the bevel (x_p) increases (see Fig. 1) [11], [12]. This relationship
 114 can be expressed in terms of an effective bevel angle.

115 One fundamental difference between junctions in GaN
 116 versus junctions in SiC or Si devices is that generally GaN
 117 junctions are grown and are therefore much more abrupt than
 118 SiC junctions, which are implanted or diffused. This
 119 fundamental difference has a significant impact on the effective
 120 bevel equation. As Adler *et al.* [11] demonstrate in their early
 121 paper, for a diffused junction the effective bevel angle is a
 122 function of the ratio of the depletions widths squared (Eq. 2),
 123 but for an abrupt junction it is simply the ratio of the depletion
 124 widths (Eq. 3):

$$126 \text{ Diffused Junction: } \theta_{eff} = 0.04 \cdot \theta_{act} \cdot \left(\frac{x_{DB}^-}{x_{DB}^+} \right)^2 \quad (2)$$

$$128 \text{ Abrupt Junction: } \theta_{eff} = 0.04 \cdot \theta_{act} \cdot \left(\frac{x_{DB}^-}{x_{DB}^+} \right) \quad (3)$$

130 where θ_{act} is the physical bevel angle, x_{DB}^- is the depletion width
 131 into the lightly-doped side at breakdown (i.e. x_n), and x_{DB}^+ is the

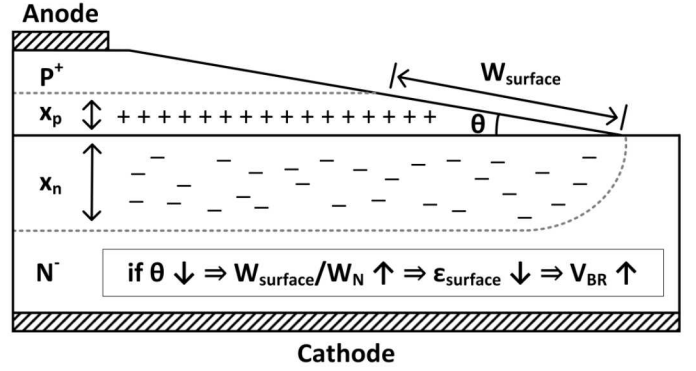


Fig. 1. Representation of fundamental theory for a negative bevel edge termination.

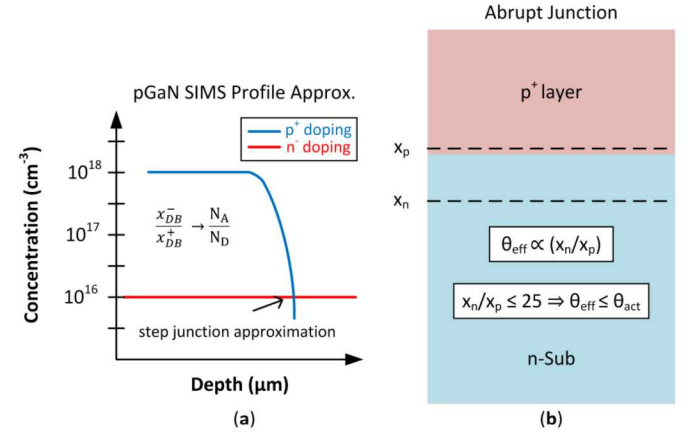


Fig. 2. Simplified representation of (a) the doping profile for an epitaxially grown p-n junction in GaN, and (b) an example of an abrupt junction and the associated bevel angle relationships.

depletion width into the heavily-doped side at breakdown (i.e. x_p).

Accordingly, an abrupt junction profile can be more easily leveraged to produce an effective bevel angle that is smaller than the physical bevel angle, making GaN an inherently better candidate for a bevel edge termination compared to Si or SiC. For an epitaxially-grown GaN p-n junction the doping profile approximates a step junction as depicted in Fig. 2(a). In the case of a step junction the ratio of the depletion widths is equal to the inverse of the ratio of the doping density in the two regions (Eq. 4).

$$\frac{x_n}{x_p} = \frac{N_A}{N_D} \quad (4)$$

The effective bevel angle equation acts as a normalization factor to evaluate the impact of bevel angle considering variation in doping on either side of the junction. In practice, however, it is also possible to use this equation to leverage the physical bevel angle by manipulating the doping such that the effective bevel angle is smaller than the physical bevel angle. This results in high breakdown performance but with a larger bevel angle, thus increasing the manufacturability of such a structure. The effective bevel angle will be smaller than the

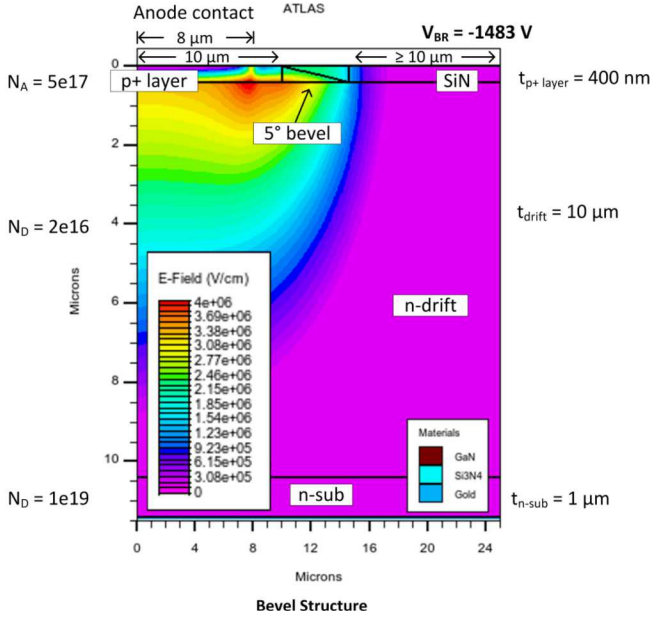


Fig. 3. Baseline simulation structure used to evaluate the negative bevel edge termination.

TABLE 1
SIMULATION PARAMETERS AND VALUES

Symbol	Quantity	Value
E_g	GaN Bandgap	3.44 eV
E_{crit}	Critical Electric Field	4 MV/cm
ϵ_s	Relative Permittivity	9.7
T	Temperature	300 K
μ_p	Hole Mobility	11 cm ² /Vs
μ_n	Electron Mobility	1200 cm ² /Vs

III. DEVICE DESIGN AND SIMULATION MODEL

To evaluate the impact of a bevel edge termination on vertical GaN power diode, we evaluate breakdown performance focusing specifically on the influence of bevel angle, p-GaN doping concentration, and p-GaN thickness. These results are compared to the ideal one-dimensional device breakdown voltage in the terminated device to the ideal breakdown voltage ($V_{B(\text{beveled})}/V_{B(\text{planar})}$), also referred to as the JTE efficiency. The device under study is nominally a 2 kV p-n junction power diode with a drift thickness and doping of 10 μm and $2 \times 10^{16} \text{ cm}^{-3}$ respectively. The baseline simulation structure used for evaluation is shown in Fig. 3 with all listed parameters held constant with the exception of bevel angle, p-layer doping, and p-layer thickness.

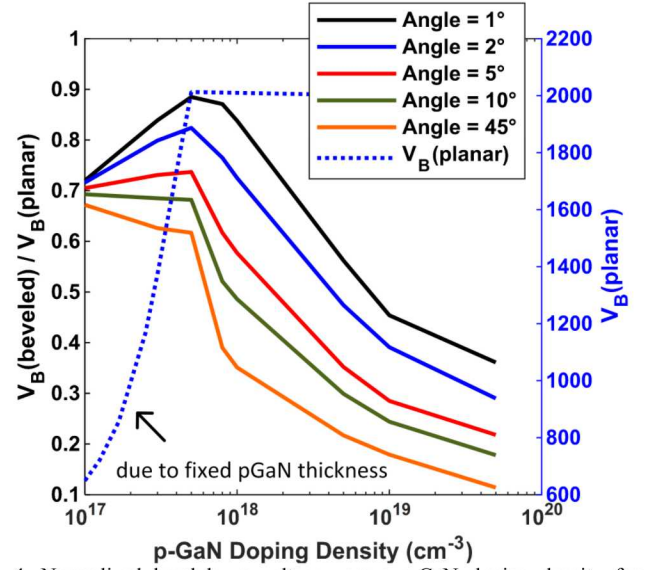


Fig. 4. Normalized breakdown voltage versus p-GaN doping density for varying bevel angles at $t_{\text{pGaN}} = 400 \text{ nm}$ (left y-axis), and ideal breakdown voltage versus p-GaN doping (right y-axis).

Device simulations were performed with Silvaco TCAD [13] to determine the reverse-bias avalanche breakdown voltage. Breakdown for the simulation is defined as the point at which the reverse-bias current density reaches 10 mA/mm², indicating the onset of avalanche. Impact ionization coefficients were specified such that at a drift doping of $2 \times 10^{16} \text{ cm}^{-3}$ the critical electric field (E_{crit}) is 4 MV/cm [14] which determines the point at which breakdown occurs. Further simulation parameters are given in Table 1 and were developed previously in [15] and [16] for similar vertical-GaN-based simulation work.

IV. RESULTS AND DISCUSSION

For a fixed drift layer design (fixed drift layer thickness and doping) the design of the bevel is dependent on three parameters: bevel p-layer thickness, p-layer doping, and bevel angle. The design considerations for bevel thickness and p-layer doping should be linked, with the desired design outcome being to maximize the depletion depth of the bevel, which serves to leverage the effective bevel angle relationship. By increasing the depletion depth into the bevel, the physical bevel angle can be increased while still yielding the same breakdown performance. This is very desirable considering that shallow ($< 10^\circ$) bevel angles can not only consume a considerable amount of wafer area but are quite difficult to manufacture.

Normalized breakdown voltage versus p-layer doping is shown in Fig. 4 for a fixed p-GaN thickness of 400 nm. In addition, the breakdown voltage for the ideal device is shown on the right axis for comparison. Notice that for $N_A \leq 5 \times 10^{17} \text{ cm}^{-3}$ the ideal breakdown voltage drops sharply. This is an example of what happens when the p-region under the anode contact fully depletes and can be corrected simply by increasing the p-GaN thickness or doping so that the layer does not fully deplete. From this figure it is seen that at $N_A = 5 \times 10^{17} \text{ cm}^{-3}$ a bevel angle of 1° approaches 90% of the ideal breakdown voltage, but a 5° bevel is only 75% ideal. It is important to note

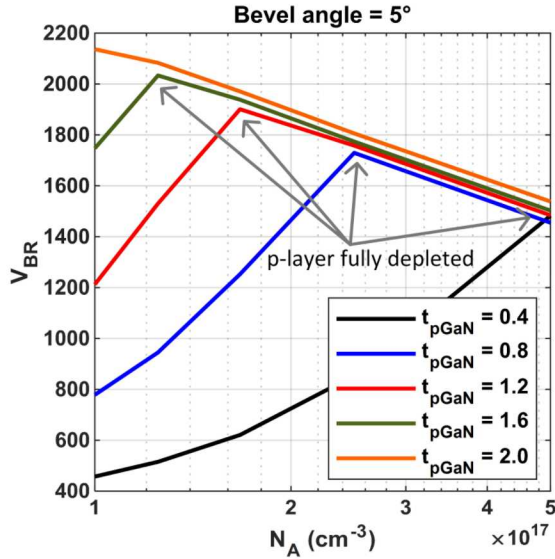


Fig. 6. Breakdown voltage versus p-GaN doping density for varying p-GaN thickness (in μm).

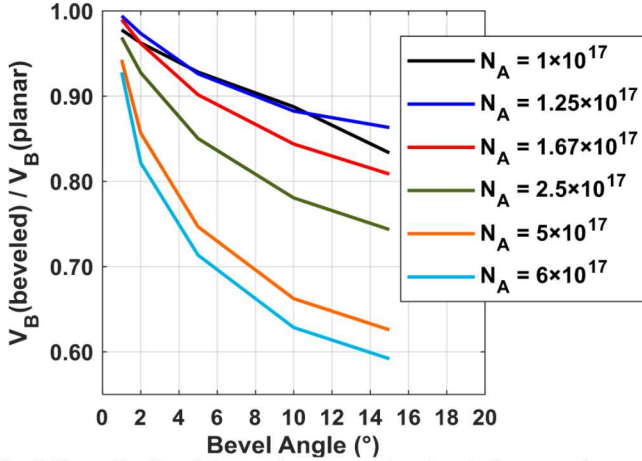


Fig. 5. Normalized breakdown voltage versus bevel angle for $t_{\text{pGaN}} = 2 \mu\text{m}$.

that for these doping levels, the effective bevel angle is equal to the physical angle, with the drift doping set to $2 \times 10^{16} \text{ cm}^{-3}$ and the p-layer doping at $5 \times 10^{17} \text{ cm}^{-3}$ (see Eq. 3). This means that the effective bevel angle relationship is completely unleveraged.

To leverage the effective bevel angle relationship, the p-layer doping should be reduced below $5 \times 10^{17} \text{ cm}^{-3}$ for the given drift doping. Since the depletion depth into the p-layer is directly proportional to the doping concentration, reducing the doping by a factor of five requires a p-layer that is five times thicker. As mentioned previously, while a single-zone JTE is very sensitive to both doping density and thickness of the JTE, a bevel edge termination does not require precise control of these variables to achieve good breakdown performance. This is highlighted in Fig. 5 which shows similar expected breakdown voltage as the p-layer thickness is varied, and an overall increase in breakdown voltage as the doping is reduced so long as the doping is not reduced to the point that the layer under the anode contact fully depletes.

Given that a much thicker p-layer is needed to accommodate doping down to $1 \times 10^{17} \text{ cm}^{-3}$, Fig. 6 shows the impact on

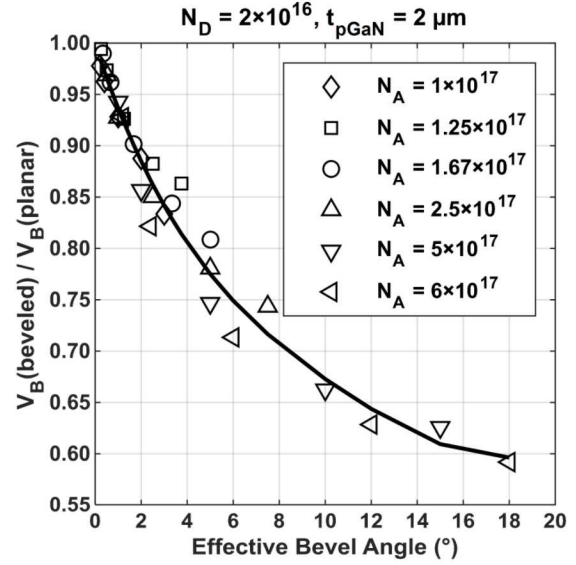


Fig. 7. Effective bevel angle plot normalizes the bevel angle to account for variations in doping.

breakdown performance of bevel angles less than 15° given $t_{\text{pGaN}} = 2 \mu\text{m}$. Note that as the bevel approaches the fully depleted scenario (at $1 \times 10^{17} \text{ cm}^{-3}$) the increase in breakdown performance begins to saturate, with results for $1\text{--}1.25 \times 10^{17} \text{ cm}^{-3}$ yielding similar breakdown performance. However, at this doping level ($1\text{--}1.25 \times 10^{17} \text{ cm}^{-3}$) it is possible to achieve nearly 90% of the ideal breakdown voltage for a bevel angle of 10° . This is a considerable improvement over the 1° bevel angle necessary to approach 90% ideality at $N_A = 5 \times 10^{17} \text{ cm}^{-3}$.

The data from Fig. 6 can be reduced to a single function by normalizing the data using the effective bevel relationship in Eq. 3. This normalization accounts for variation in doping and demonstrates a very good fit, showcasing the accuracy of the analytical equation for an abrupt junction (Eq. 3). This normalized data is presented in Fig. 7 demonstrate that for an effective bevel angle of less than 1° it is possible to reach near 100% ideality for breakdown performance. The symbols in Fig. 7 represent the discrete datapoints from the simulations represented in Fig. 6 normalized using Eq. 3.

To obtain such a high level of breakdown performance, it is necessary to significantly leverage the effective bevel angle equation by using a very low-doped and consequently very thick p-layer, or by fabricating very shallow bevel angles. Fabricating a very thick p-layer with low doping is generally undesirable due to the negative impact on forward resistance (R_{on}). However, by employing a moderate approach with respect to p-layer doping and thickness, it is possible to take advantage of the effective bevel relationship without significantly increasing R_{on} .

V. CONCLUSION

The bevel edge termination is fundamentally a very attractive approach to realizing the full breakdown performance for a vertical power device. In contrast to a single/multi-zone JTE, the bevel design is robust to changes in p-layer doping and thickness. While surface charge from passivation can kill a MZ-JTE due to the discrete total charge in each zone, the bevel's

continuum of charge makes it more tolerant of surface charge. However, from a fabrication standpoint there are significant challenges to overcome to produce such a shallow (10-15°) bevel angle. In contrast, a multi-zone JTE is relatively straightforward to process, although adding more zones comes at the cost of higher complexity.

The simulation results presented in this work demonstrate the capability of a negative bevel edge termination to achieve near 90% breakdown ideality for a 10-15° physical bevel angle. This is achieved by leveraging the effective bevel angle relationship. Increasing the depletion depth into the bevel compared to the drift depletion yields higher breakdown performance for the same given bevel angle. In this way, leveraging the effective bevel angle relationship can serve to increase the manufacturability of the bevel by relaxing the bevel angle requirement necessary for high breakdown performance. Not only does a 10° bevel angle take up less wafer real estate than a 1° bevel, but it is also much more feasible to fabricate.

Doping profiles for epitaxially-grown GaN p-n junctions are much more abrupt compared to implanted junctions characteristic of SiC power devices. As demonstrated by the analytical equations, the effective bevel angle relationship can be more easily leveraged with an abrupt junction (GaN) than with an implanted junction (SiC). This gives the bevel edge termination even more significance for GaN based power devices.

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