



Reformulating Friedman's Implicit Integrator as a Single Step Method for Maxwell's Equations as a First Order System

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Introduction

We have reformulated Friedman's integrator for wave equations to serve as a single step method for the first order system.

- Allows simpler inclusion of complex closure models
- Naturally provides an accurate \mathbf{B} for particle push
- Does not require differentiating \mathbf{J} .
- Requires two fewer edge fields ($\mathbf{J}^{n+1/2}, \mathbf{E}^{n-1}$)

Friedman's Original Integrator (for Electric Fields)

$$\frac{\mathbf{E}^{n+1} - 2\mathbf{E}^n + \mathbf{E}^{n-1}}{\Delta t^2} + c_0^2 \operatorname{curl} \operatorname{curl} \frac{\mathbf{E}^{n+1} + \frac{\theta}{2}\mathbf{E}^n + (1 - \frac{\theta}{2})\bar{\mathbf{E}}^{n-1}}{2} + \frac{1}{\epsilon_0} \frac{\mathbf{J}^{n+1/2} - \mathbf{J}^{n-1/2}}{\Delta t} = 0$$

$$\bar{\mathbf{E}}^n = (1 - \frac{\theta}{2})\mathbf{E}^n + \frac{\theta}{2}\bar{\mathbf{E}}^{n-1}$$

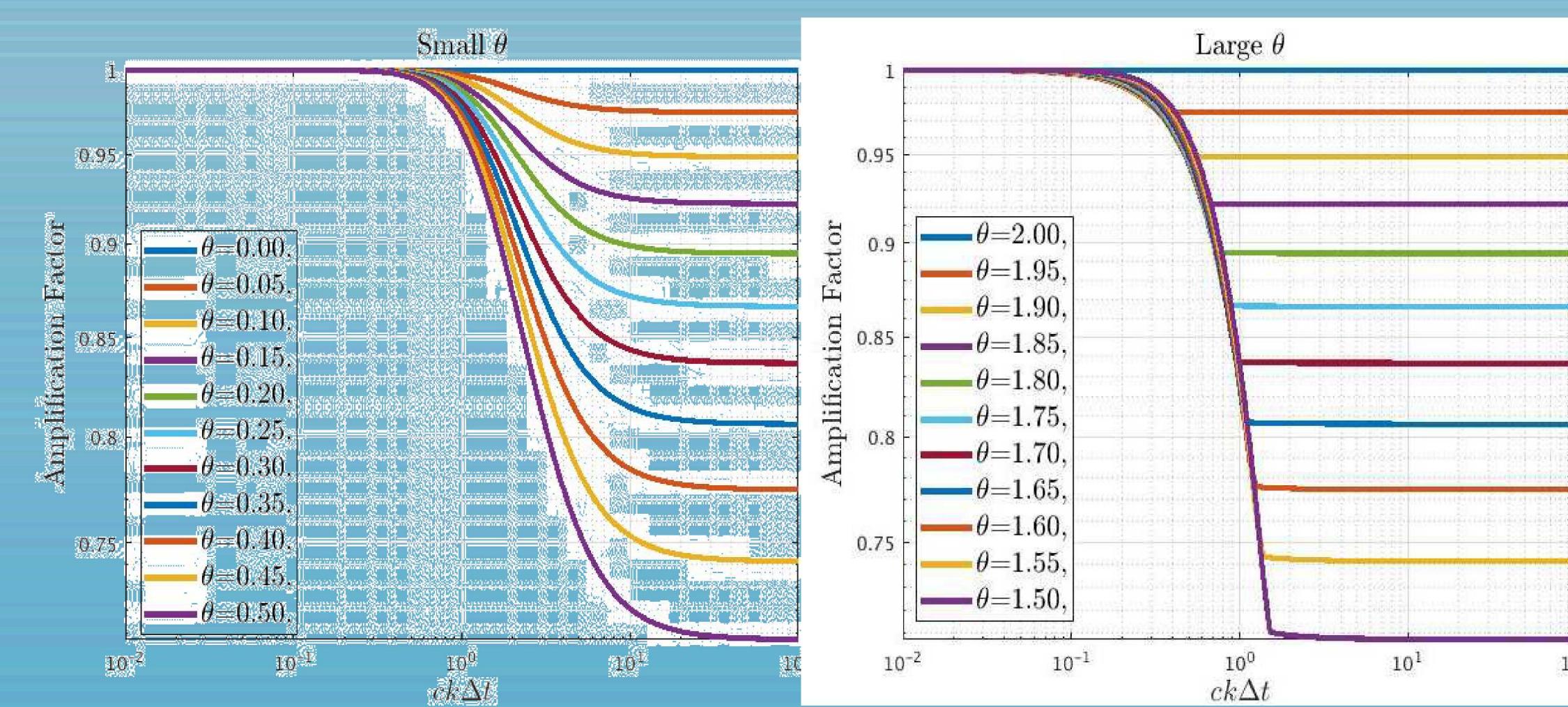
2nd order accurate for $\theta \in (0, 2)$ and quiescent initial conditions.

$$\frac{\mathbf{E}^{n+1} + \frac{\theta}{2}\mathbf{E}^n + (1 - \frac{\theta}{2})\bar{\mathbf{E}}^{n-1}}{2} = \frac{\mathbf{E}^{n+1} + \frac{\theta}{2}\mathbf{E}^n + (1 - \frac{\theta}{2})^2 \sum_{j=0}^{\infty} (\frac{\theta}{2})^j \mathbf{E}^{n-j-1}}{2}$$

$$\approx \frac{1 + \frac{\theta}{2} + (1 - \frac{\theta}{2})^2 \sum_{j=0}^{\infty} (\frac{\theta}{2})^j}{2} \mathbf{E}(n\Delta t)$$

$$+ \frac{1 - (1 - \frac{\theta}{2})^2 \sum_{j=0}^{\infty} (j+1)(\frac{\theta}{2})^j}{2} \Delta t \mathbf{E}'(n\Delta t) + \mathcal{O}(\Delta t^2)$$

For $\theta \in (0, 2)$, the method is unconditionally stable with third order damping. Amplification factors can be determined by calculating the spectral radius of the amplification matrix.



Staggered Formulation for Electric & Magnetic Fields (I)

From Friedman's original paper: Apply Yee staggering to Ampere's Law and correct Faraday's Law to match the 2nd order stencil for \mathbf{E} .

$$\begin{cases} \frac{\mathbf{E}^{n+1} - \mathbf{E}^n}{\Delta t} - c_0^2 \operatorname{curl} \mathbf{B}^{n+1/2} + \frac{1}{\epsilon_0} \mathbf{J}^{n+1/2} = 0 \\ \frac{\mathbf{B}^{n+1/2} - \mathbf{B}^{n-1/2}}{\Delta t} + \operatorname{curl} \frac{\mathbf{E}^{n+1} + \frac{\theta}{2}\mathbf{E}^n + (1 - \frac{\theta}{2})\bar{\mathbf{E}}^{n-1}}{2} = 0 \\ \bar{\mathbf{E}}^n = (1 - \frac{\theta}{2})\mathbf{E}^n + \frac{\theta}{2}\bar{\mathbf{E}}^{n-1} \end{cases}$$

Eliminate \mathbf{E} from Ampere's Law

$$\begin{aligned} & \mathbf{B}^{n+1/2} - 2\mathbf{B}^{n-1/2} + \mathbf{B}^{n-3/2} \\ & + \frac{\Delta t^2}{2} \mathbf{B}^{n+1/2} + \frac{\theta}{2} \mathbf{B}^{n-1/2} + (1 - \frac{\theta}{2})^2 \sum_{j=0}^{\infty} (\frac{\theta}{2})^j \mathbf{B}^{n-3/2-j} \\ & + c_0^2 \operatorname{curl} \operatorname{curl} \frac{\mathbf{B}^{n+1/2} + \frac{\theta}{2} \mathbf{B}^{n-1/2} + (1 - \frac{\theta}{2})^2 \sum_{j=0}^{\infty} (\frac{\theta}{2})^j \mathbf{B}^{n-3/2-j}}{2} = 0 \end{aligned}$$

2nd order accurate at half steps!

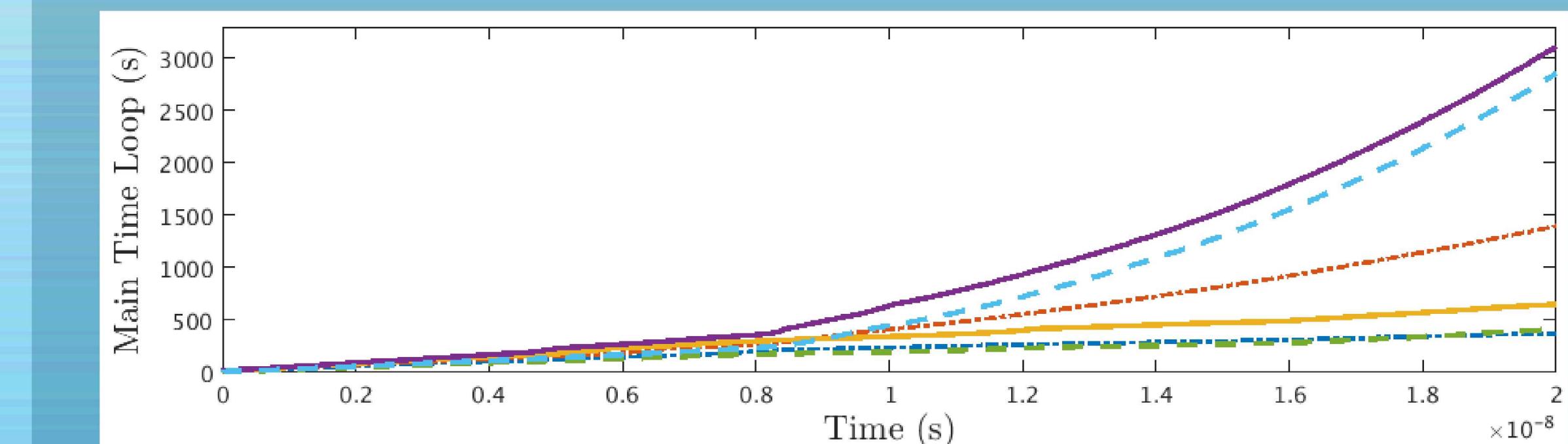
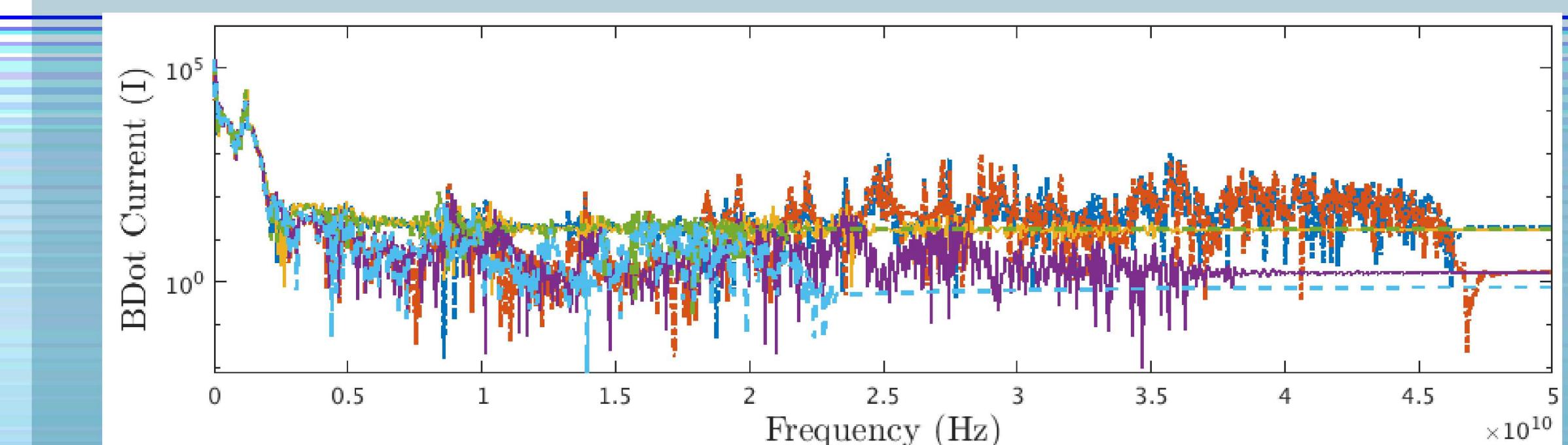
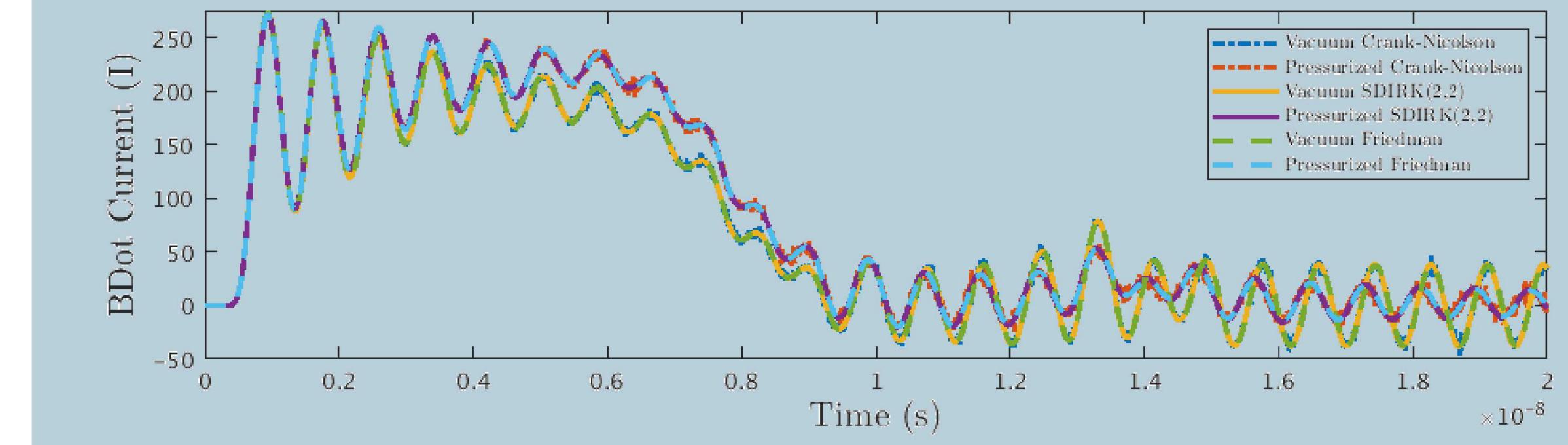
Code Verification for the single step formulation

$$\begin{cases} \mathbf{A}_0(x, y) = (0, 0, \sin(k_x x) \sin(k_y y)) \\ \mathbf{J} = J_0 \sin(\omega t) \mathbf{A}_0 \\ \mathbf{E} = E_0 (\cos(c_0 k t) - \cos(\omega t)) \mathbf{A}_0 \\ \mathbf{B} = E_0 (\frac{1}{\omega} \sin(\omega t) - \frac{1}{c k} \sin(c k t)) \operatorname{curl} \mathbf{A}_0 \end{cases}$$

Compute Relative L2 Errors at $6\pi\omega^{-1}$

$\frac{2\pi}{kh}$	$\frac{2\pi}{\omega\Delta t}$	$\frac{\ \mathbf{E} - \mathbf{E}_h\ _{L^2}}{\ \mathbf{E}\ _{L^2}}$	rate	$\frac{\ \mathbf{B} - \mathbf{B}_h\ _{L^2}}{\ \mathbf{B}\ _{L^2}}$	rate
20	10	3.53e-2	—	1.80e-2	—
40	20	8.93e-3	1.99	4.49e-3	2.00
80	40	2.24e-3	2.00	1.12e-3	2.00
160	80	5.60e-4	2.00	2.80e-4	2.00

Radiation Driven BDot Cavity at 0.0 and 0.01 Torr



Single Step Formulation for Electric & Magnetic Fields (II)

Advance Faraday's Law with Crank-Nicolson and correct Ampere's Law to match 2nd order stencil for \mathbf{E} .

$$\begin{cases} \frac{\mathbf{E}^{n+1} - \mathbf{E}^n}{\Delta t} - c_0^2 \operatorname{curl} \mathbf{B}^{n+1} - \frac{c_0^2 \Delta t}{2} \operatorname{curl} \operatorname{curl} \bar{\mathbf{E}}^n + \frac{1}{\epsilon_0} \mathbf{J}^{n+1/2} = 0 \\ \frac{\mathbf{B}^{n+1} - \mathbf{B}^n}{\Delta t} + \operatorname{curl} \frac{\mathbf{E}^{n+1} + \mathbf{E}^n}{2} = 0 \\ \bar{\mathbf{E}}^{n+1} = (1 - \frac{\theta}{2})\mathbf{E}^{n+1} + \frac{\theta}{2}\bar{\mathbf{E}}^n \end{cases}$$

Eliminate \mathbf{E} from Ampere's Law

$$\begin{aligned} & \mathbf{B}^{n+1} - 2\mathbf{B}^n + \mathbf{B}^{n-1} \\ & + \frac{\Delta t^2}{2} \mathbf{B}^{n+1} + \frac{\theta}{2} \mathbf{B}^n + (1 - \frac{\theta}{2})^2 \sum_{j=0}^{\infty} (\frac{\theta}{2})^j \mathbf{B}^{n-j-1} \\ & - \frac{1}{\epsilon_0} \operatorname{curl} \frac{\mathbf{J}^{n+1/2} + \mathbf{J}^{n-1/2}}{2} = 0 \end{aligned}$$

2nd order accurate on integer steps!

References

A. Friedman, "A second-order implicit particle mover with adjustable damping." J. Comput. Phys. 1990