



Optical Beat Note Readout of a Magnetic Gradient

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Sandia National Laboratories

Quantum Sensing & Magnetometry: from the
nanoscale up to geological explorations

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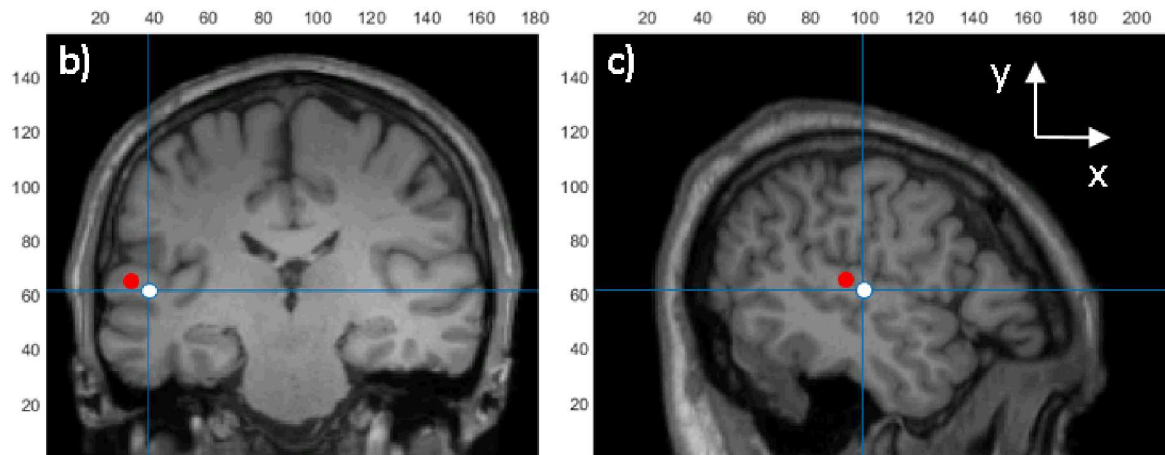
Motivation

Direct optical measurement of a magnetic field gradient

High sensitivity in the Earth's magnetic field

- High common mode rejection

Application: magnetoencephalography and magnetocardiography

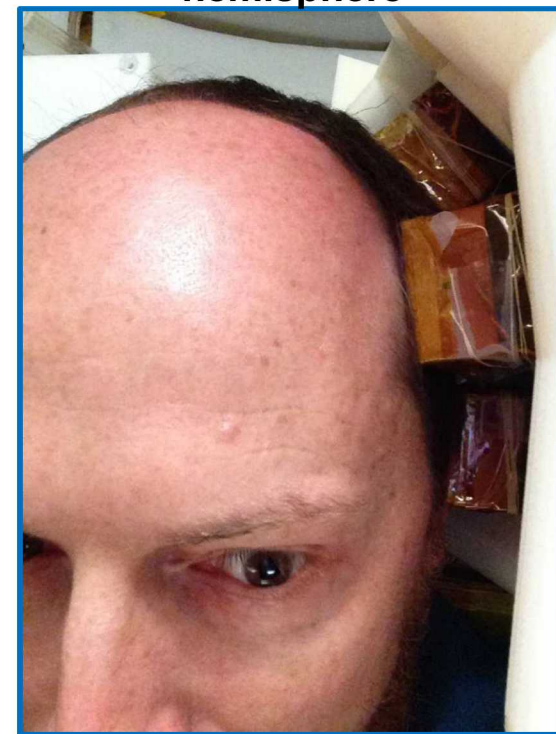
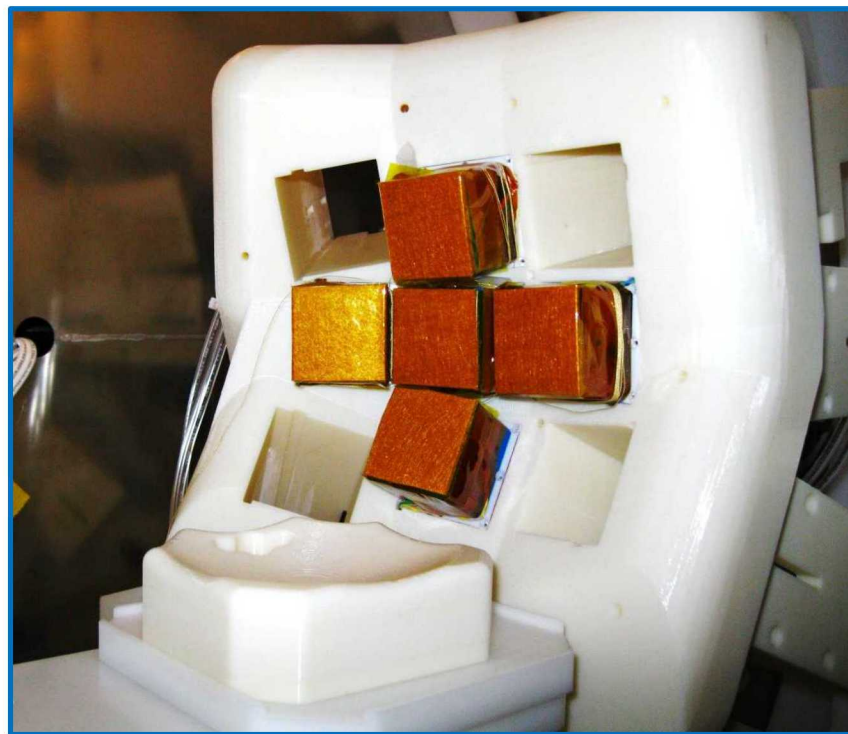


5-sensor, 20-channel array

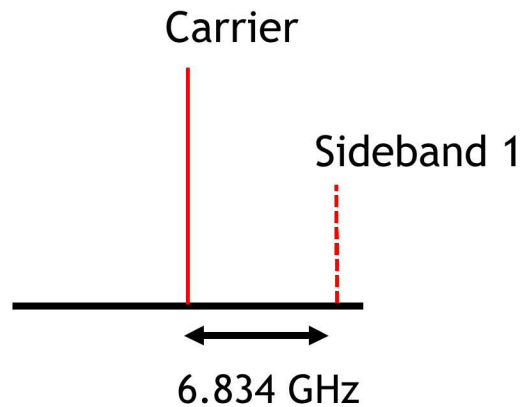
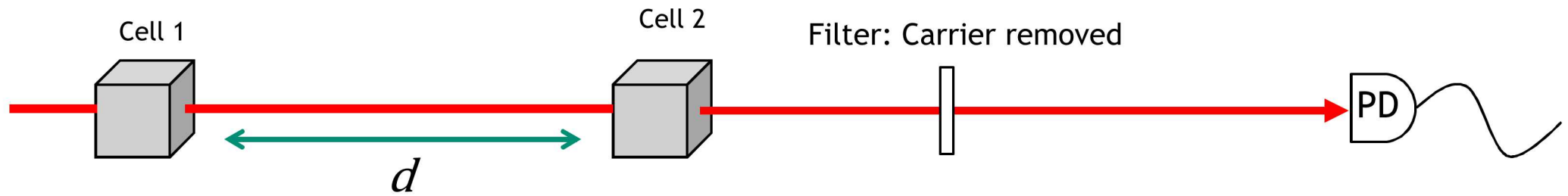
Partially covers the left hemisphere

Person-Sized Magnetic Shield

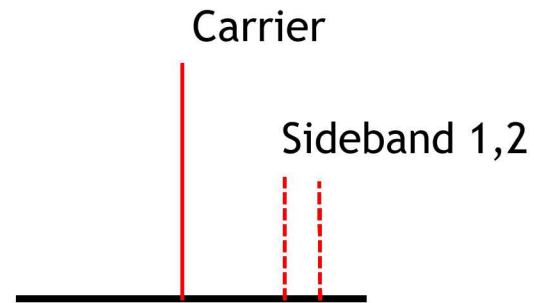
Insert
Person
Here



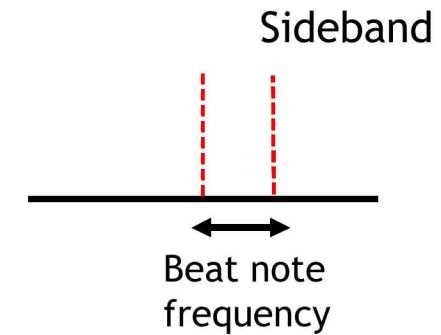
Basic Idea – Produce a beat note signal



Sideband Produced in the first cell
orthogonally polarized
to the carrier



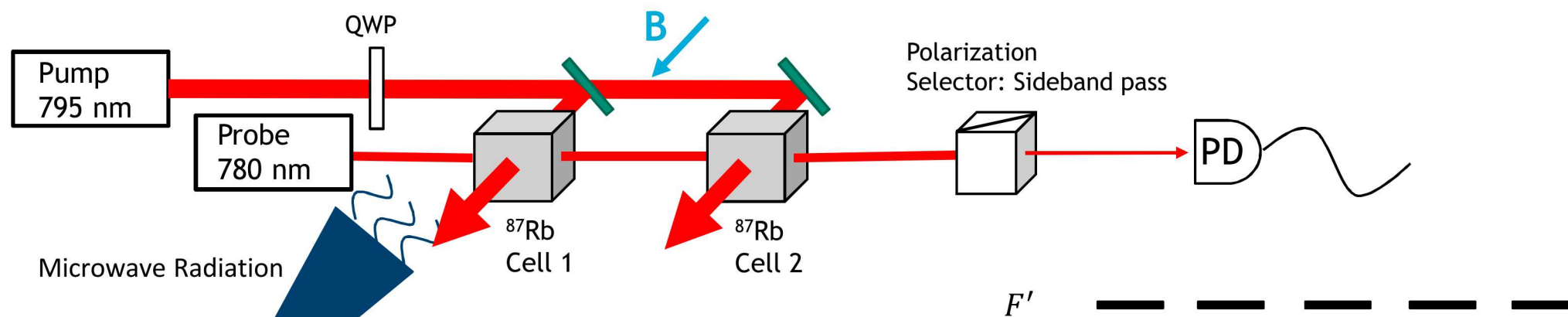
Sideband Produced in the second cell
orthogonally polarized
to the carrier



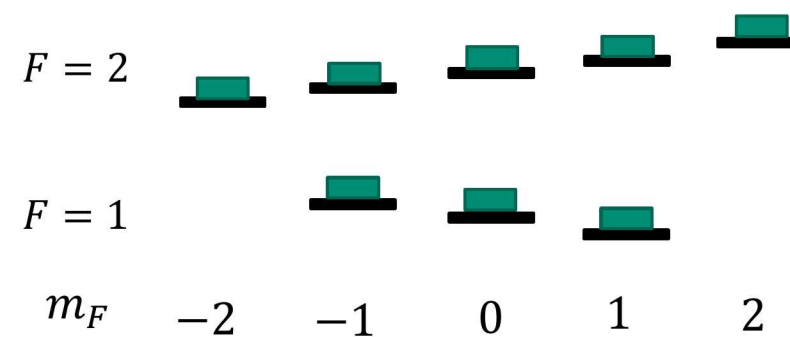
Carrier removed and
sidebands beat together

$$f_B \propto \frac{\Delta B}{d}$$

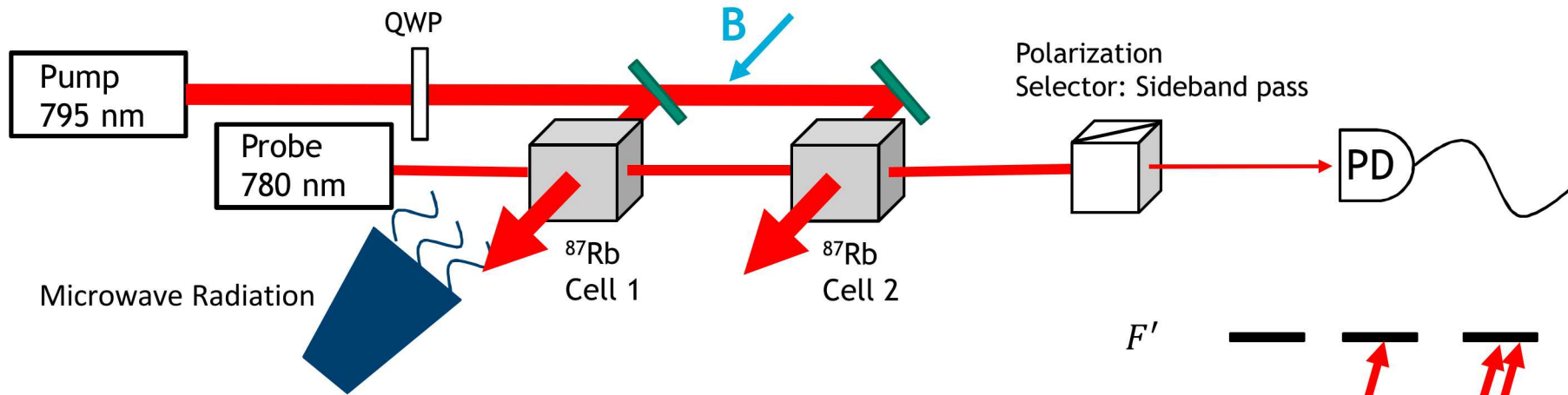
More detailed description



Gradiometer Process

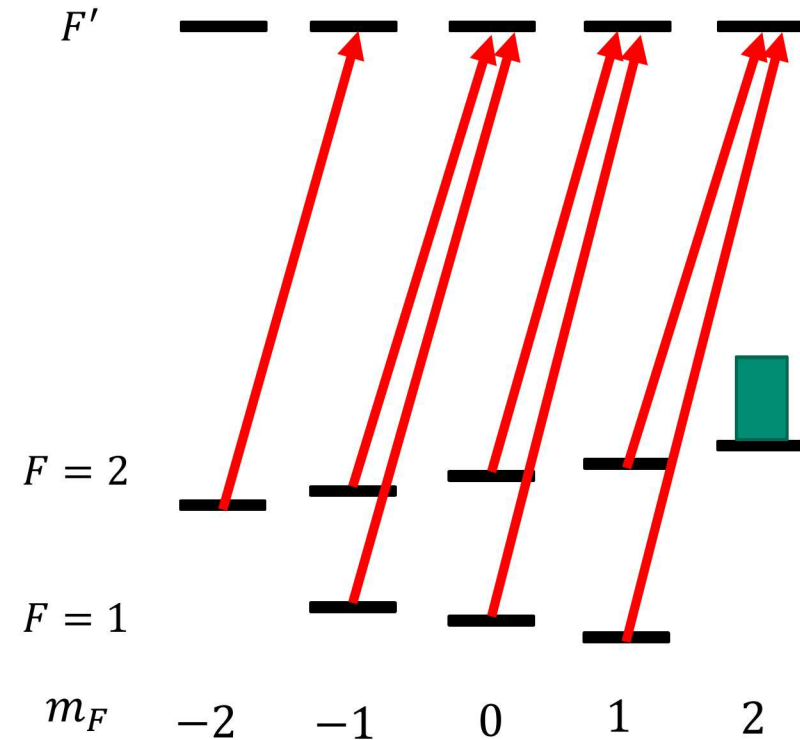


More detailed description

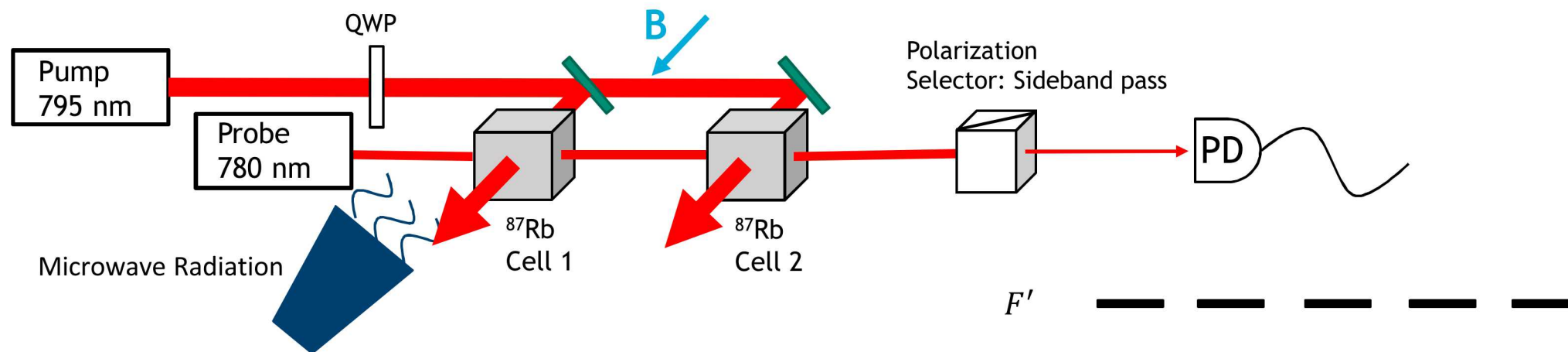


Gradiometer Process

- 1) Optically pump atoms to the end-state

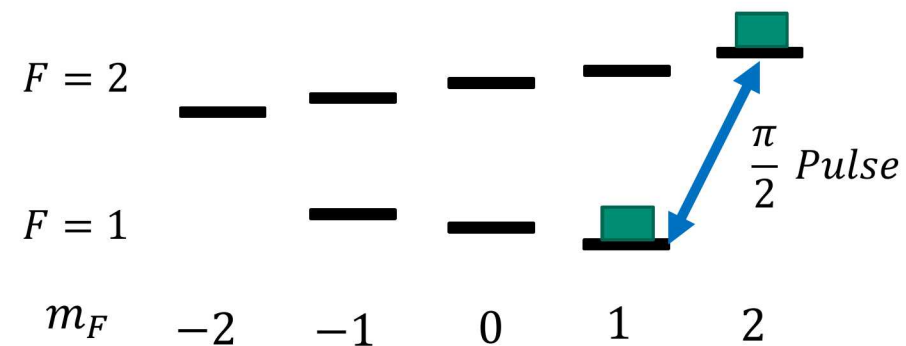


More detailed description

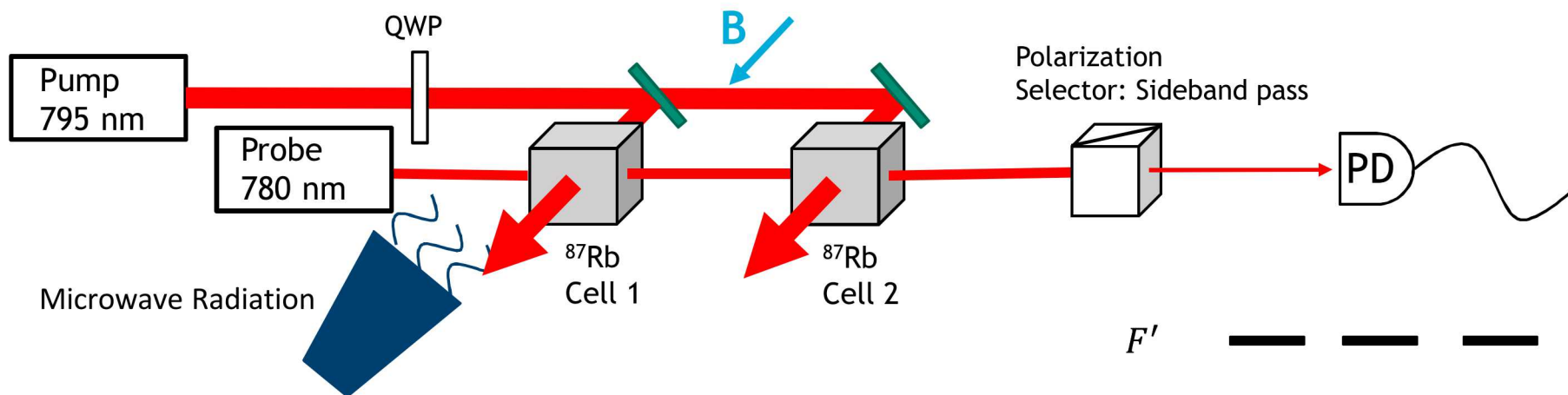


Gradiometer Process

- 1) Pump atoms to the end-state
- 2) Apply a $\frac{\pi}{2}$ pulse of magnetic energy to put the atoms in a coherent superposition

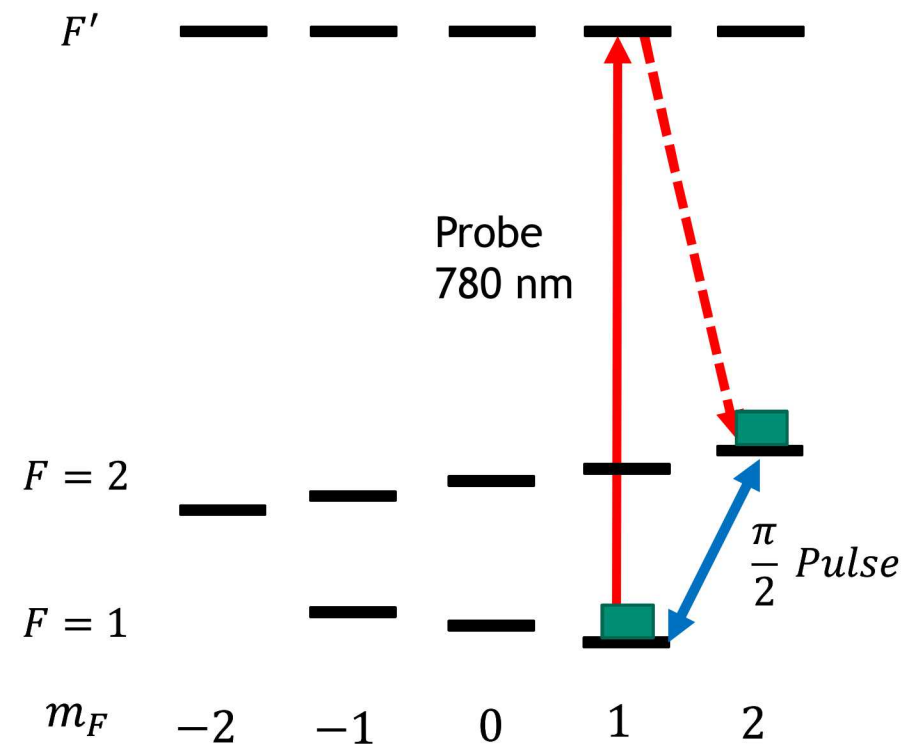


More detailed description

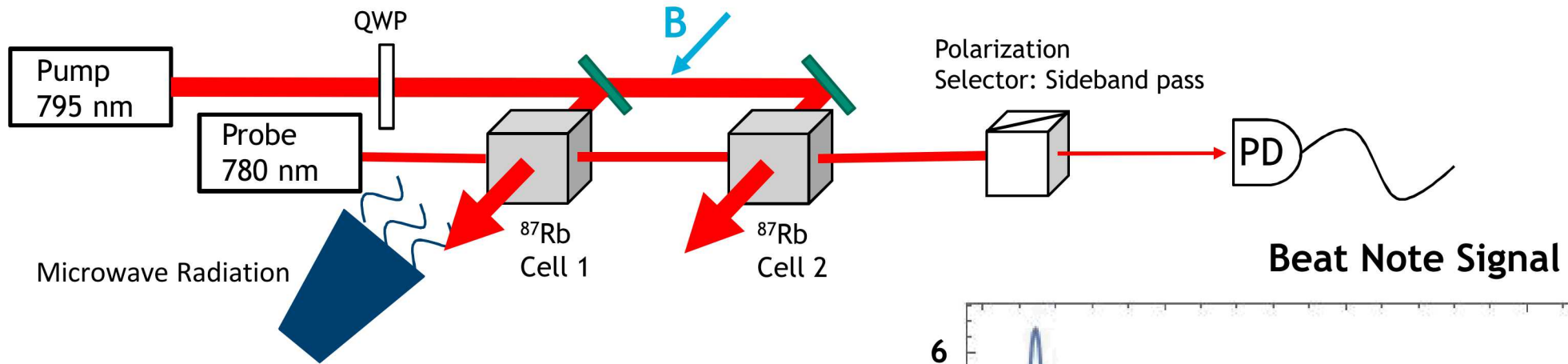


Gradiometer Process

- 1) Pump atoms to the end-state
- 2) Apply a $\frac{\pi}{2}$ pulse of magnetic energy to put the atoms in a coherent superposition
- 3) Send in Probe light to generate a sideband in each cell

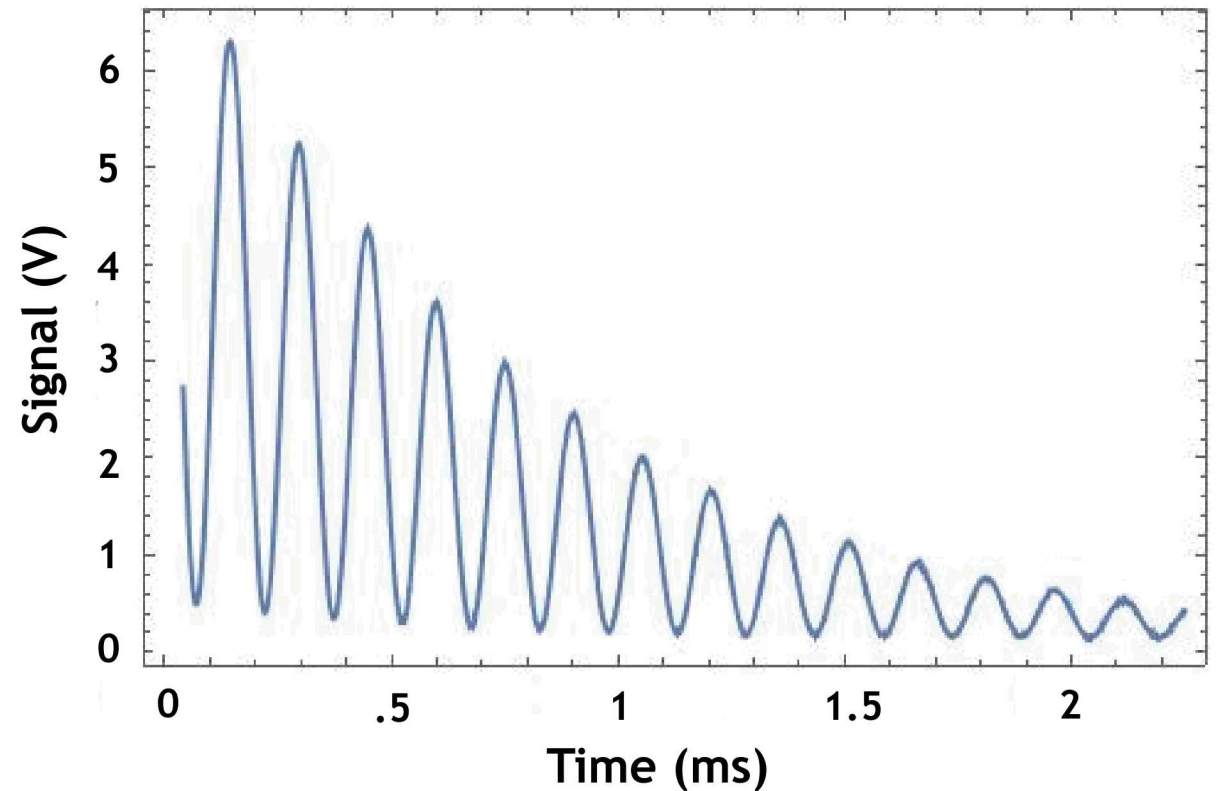


More detailed description

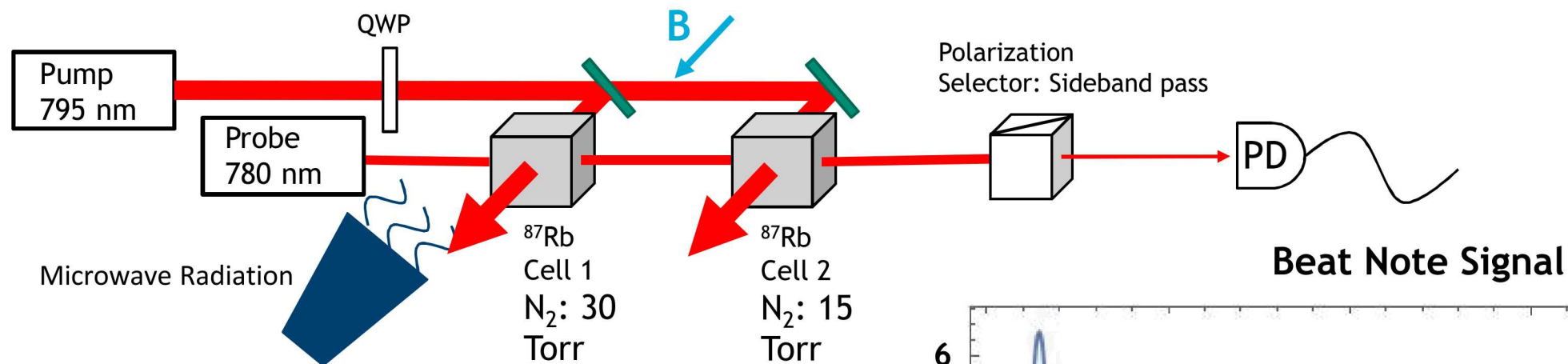


Gradiometer Process

- 1) Pump atoms to the end-state
- 2) Apply a $\frac{\pi}{2}$ pulse of magnetic energy to put the atoms in a coherent superposition
- 3) Send in Probe light to generate a sideband in each cell
- 4) Beat the sidebands together to produce a beat note

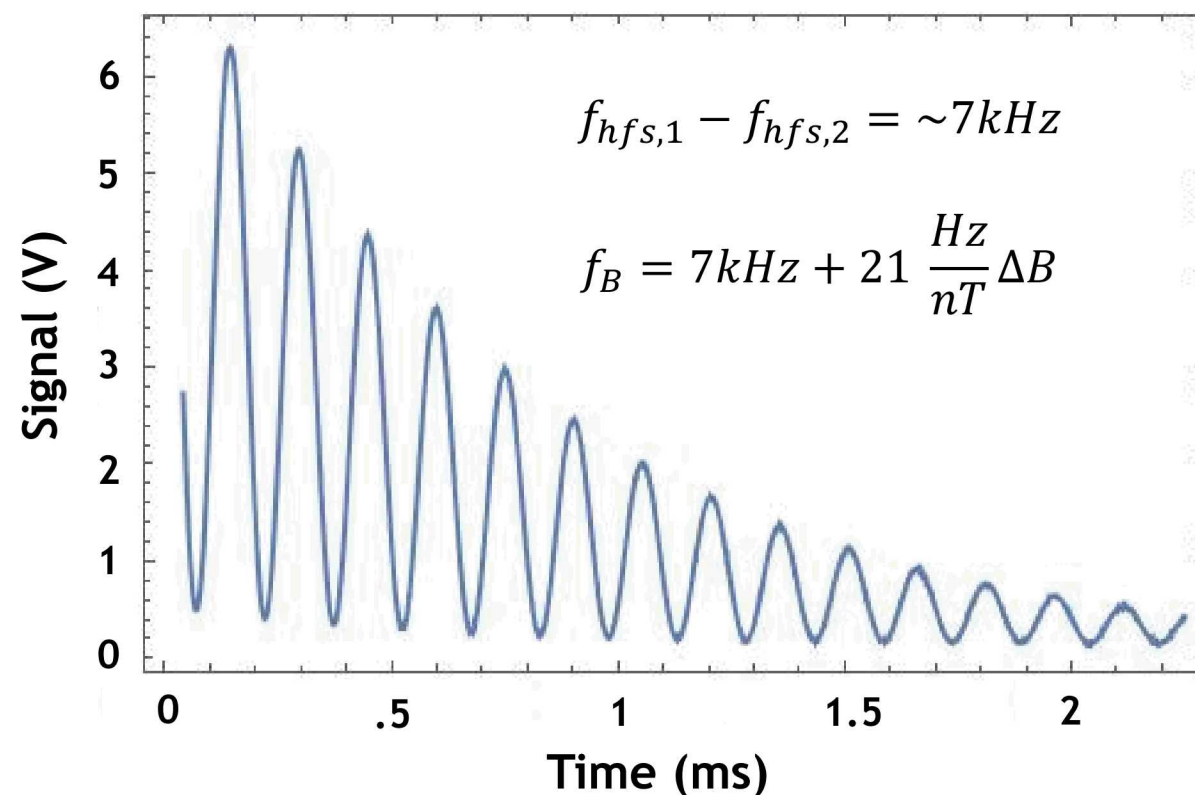


More detailed description



Gradiometer Process

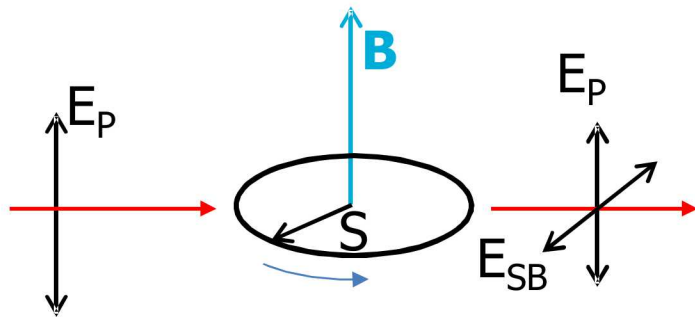
- 1) Pump atoms to the end-state
- 2) Apply a $\frac{\pi}{2}$ pulse of magnetic energy to put the atoms in a coherent superposition
- 3) Send in Probe light to generate a sideband in each cell
- 4) Beat the sidebands together to produce a beat note
- 5) Measure the frequency of the beat note to determine the magnetic gradient between the cells



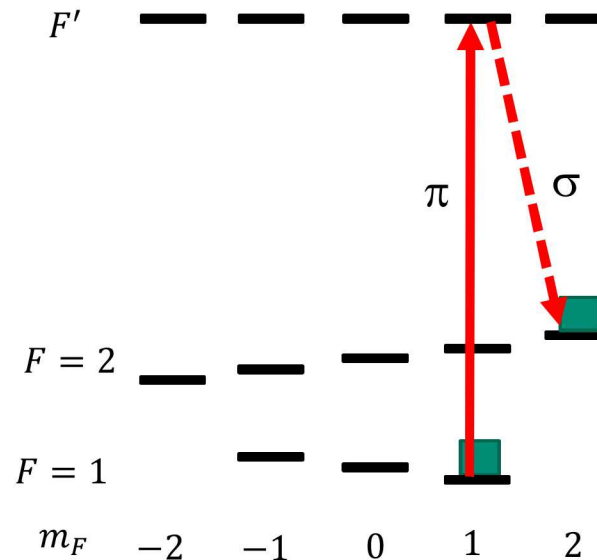
Sideband Generation

Modulation of the index of refraction

- Polarization rotation at the hyperfine frequency

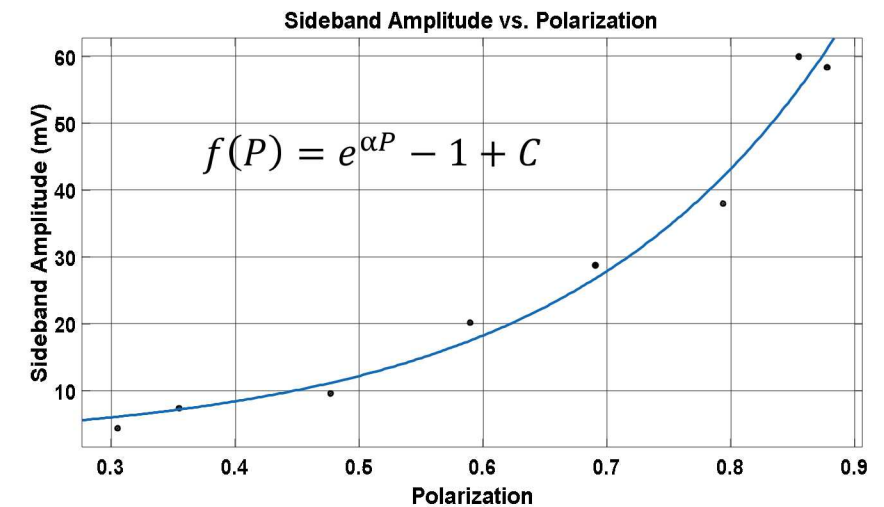


Stimulated Raman Transitions



$$\frac{d\hat{\rho}}{dt} = -\frac{i}{\hbar}[\hat{H}, \hat{\rho}] + R,$$

$$\frac{d\Omega_i}{dz} = -\alpha_i \times \text{Im}[\rho_{i3}],$$



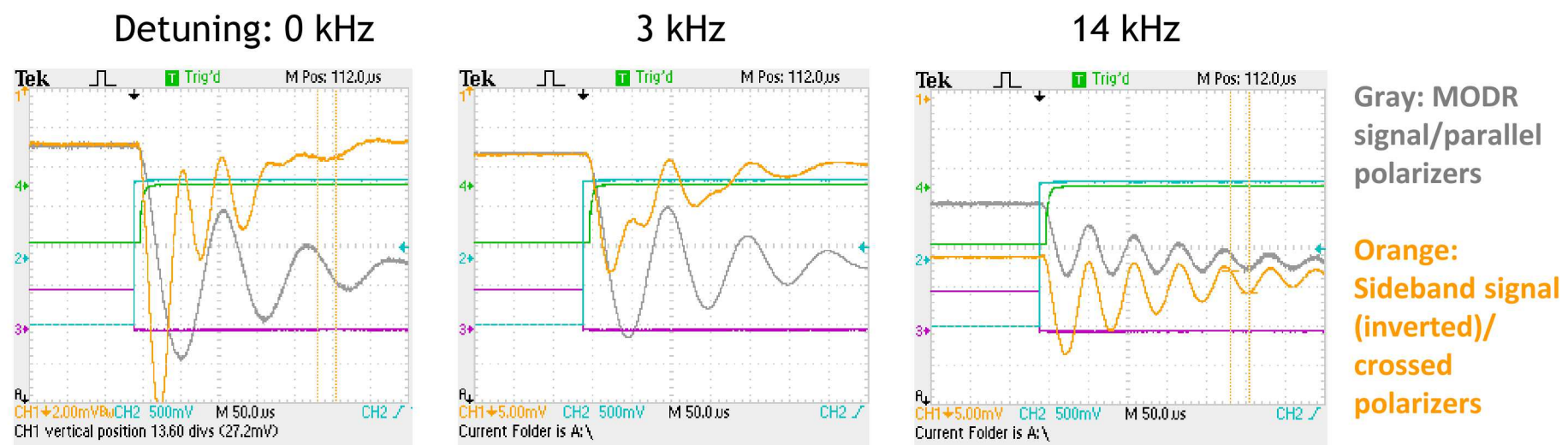
H. Tang and W. Happer, "Parametric Frequency Conversion of Resonance Radiation in Optically Pumped Rubidium-87 Vapor," *Phys. Rev. Lett.* **24**, 551 (1970).

Henry Tang, "Parametric Frequency Conversion of Resonance Radiation in Optically Pumped Rb⁸⁷ Vapor," *Phys. Rev. A* **7**, 1010 (1973).

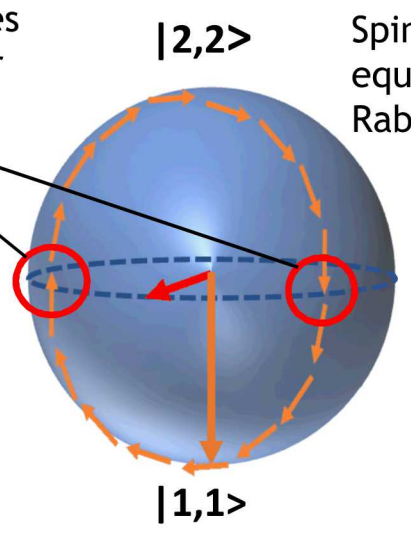
Vishal Shah, *System and Method for Measuring a Magnetic Gradient Field*. Patent. US10088535 (2018).

Sideband generation from coherence

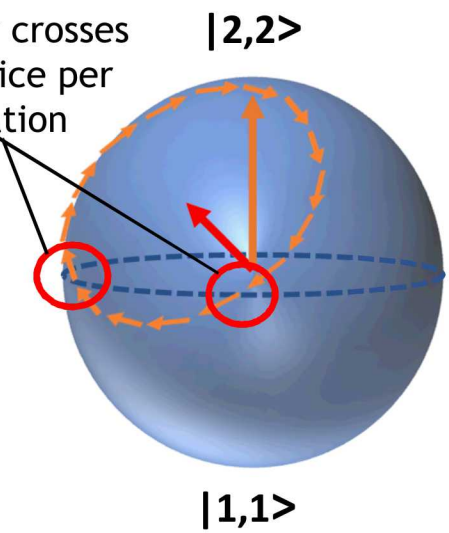
- Microwave optical double resonance (MODR) signal oscillates at the Rabi frequency, Ω .
- Sideband signal oscillates at 2Ω on resonance.
- Sidebands maximized when coherence is maximized.



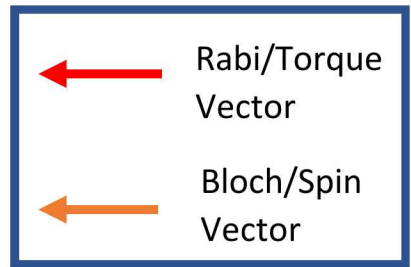
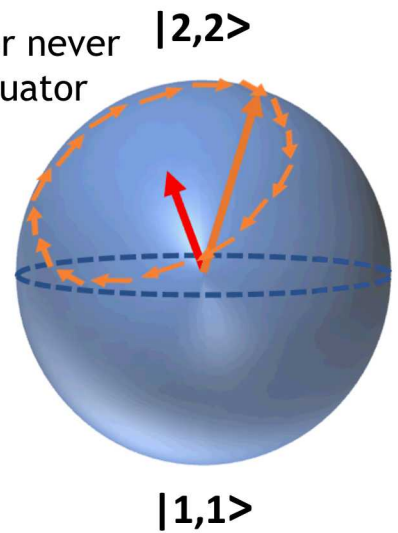
Spin vector crosses equator twice per Rabi oscillation



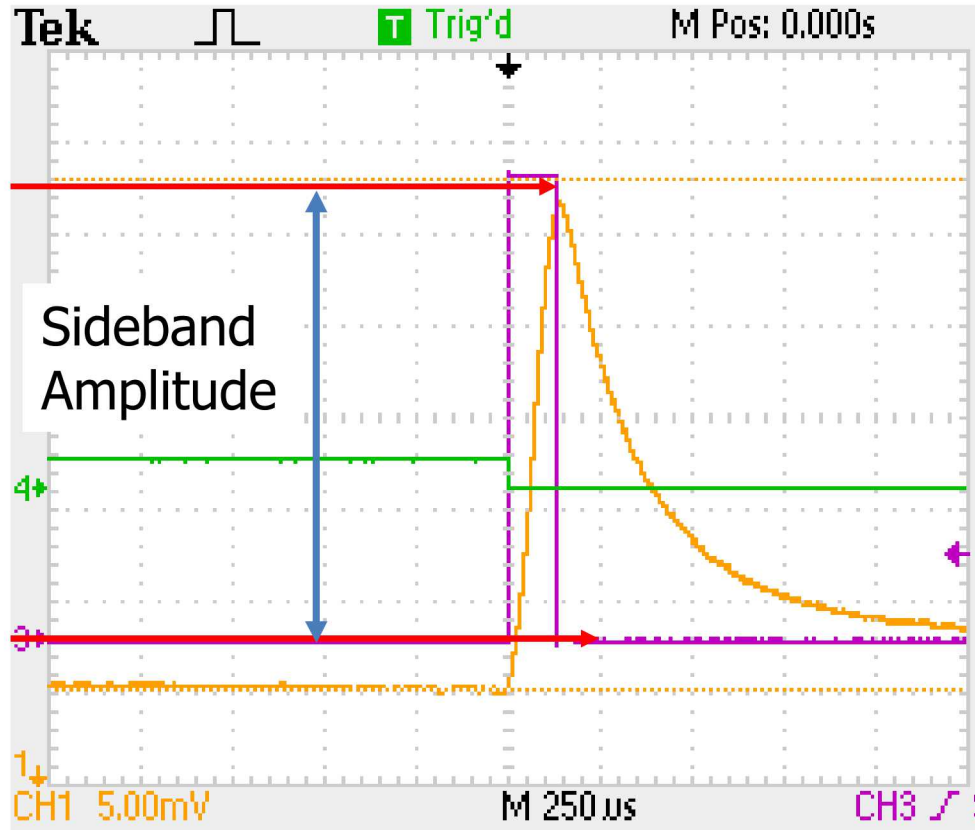
Spin vector crosses equator twice per Rabi oscillation



Spin vector never crosses equator

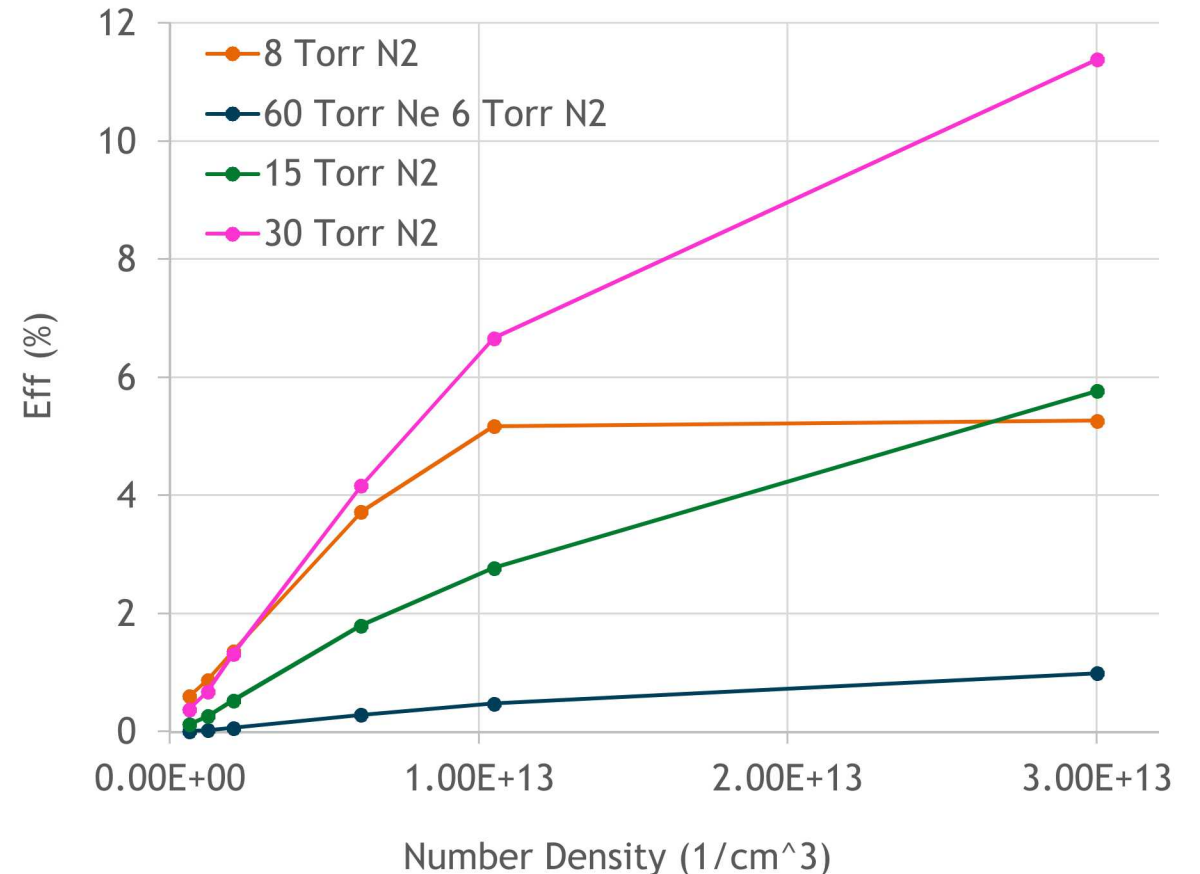


Sideband Optimization: Buffer Gas Pressure



$$Eff = \frac{\text{Sideband Amplitude } (\mu W)}{\text{Input Probe Power } (\mu W)}$$

Effect of number density and buffer gas pressure



- Percentage of probe power being converted to sideband by the vapor cell.
- Input power = 12 μW

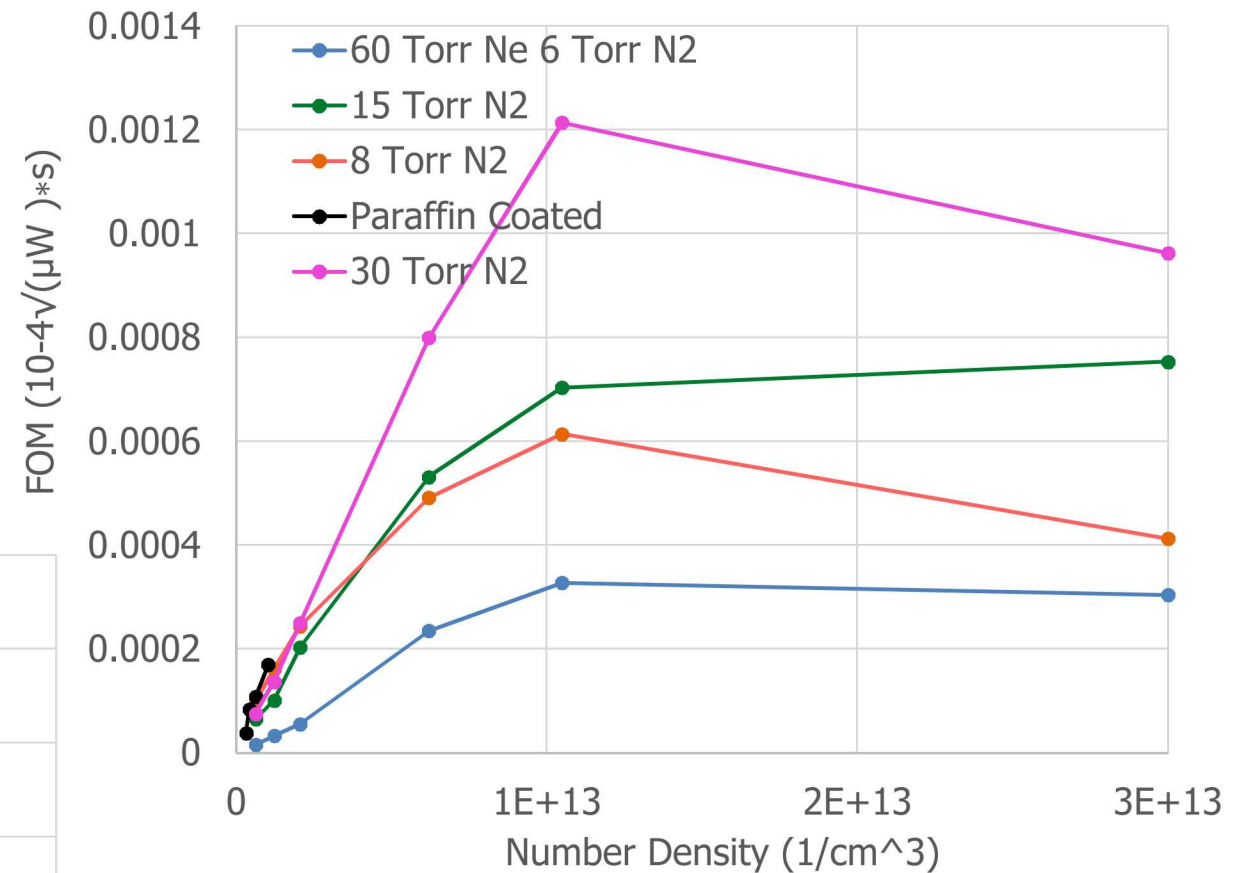
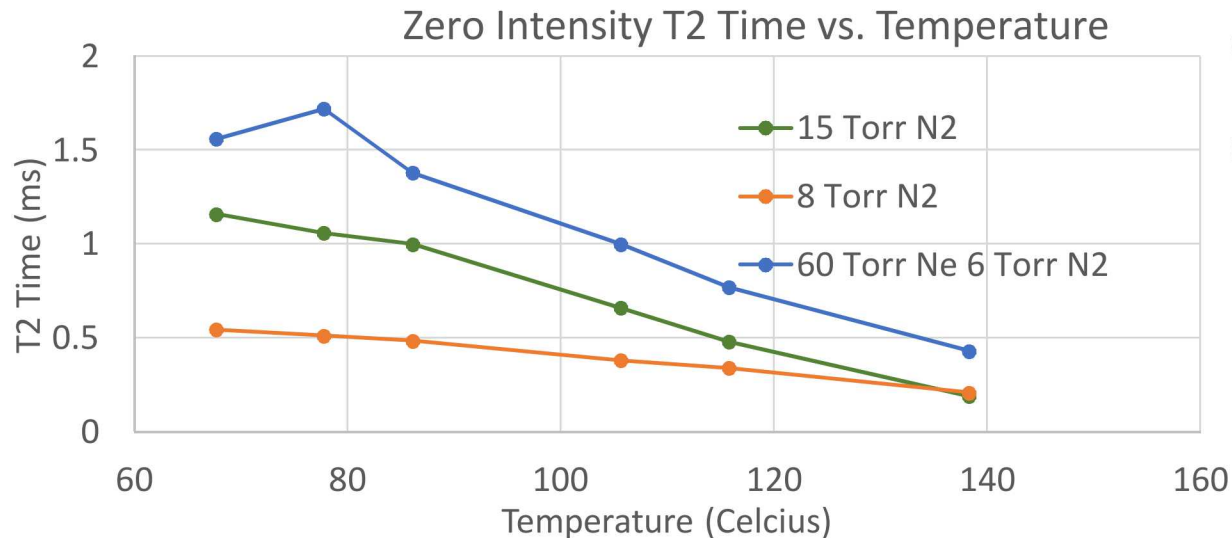
Figure of Merit

Figure of merit

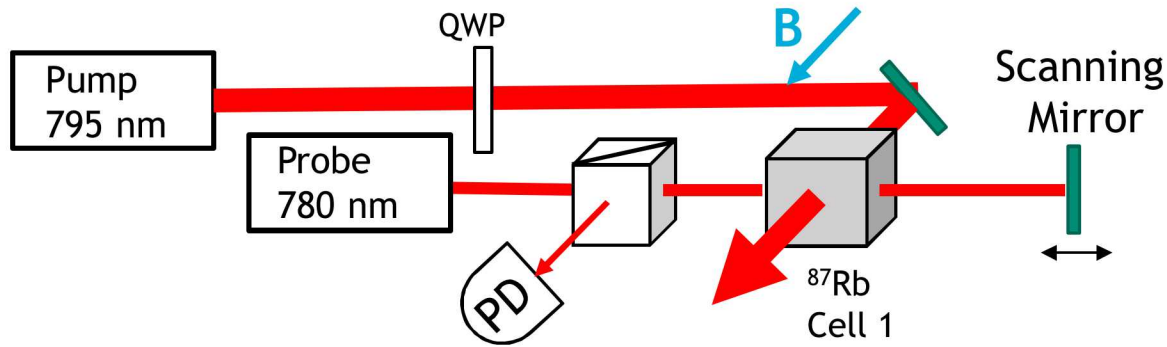
- Reflects the performance of the gradiometer
- $FOM = \sqrt{S}/w_2$
- S = sideband signal amplitude
- The linewidth $w_2 = 1/(\pi T_2)$.

Probe on continuously throughout the measurement.

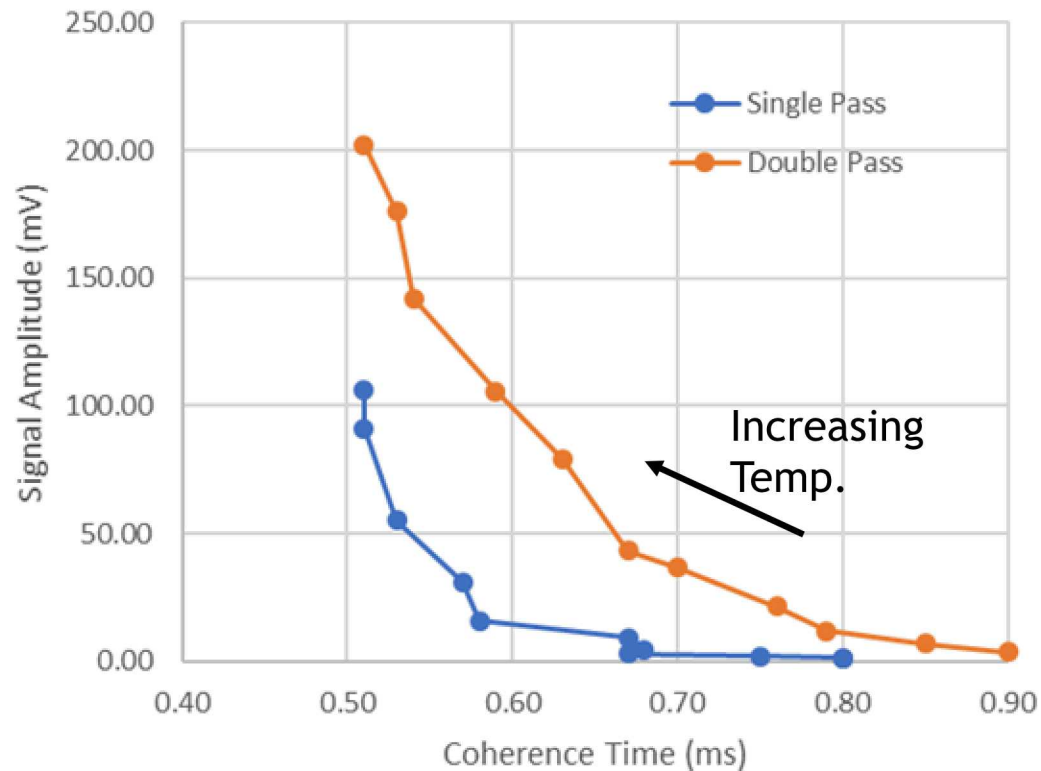
- The probe power is 12 uW.



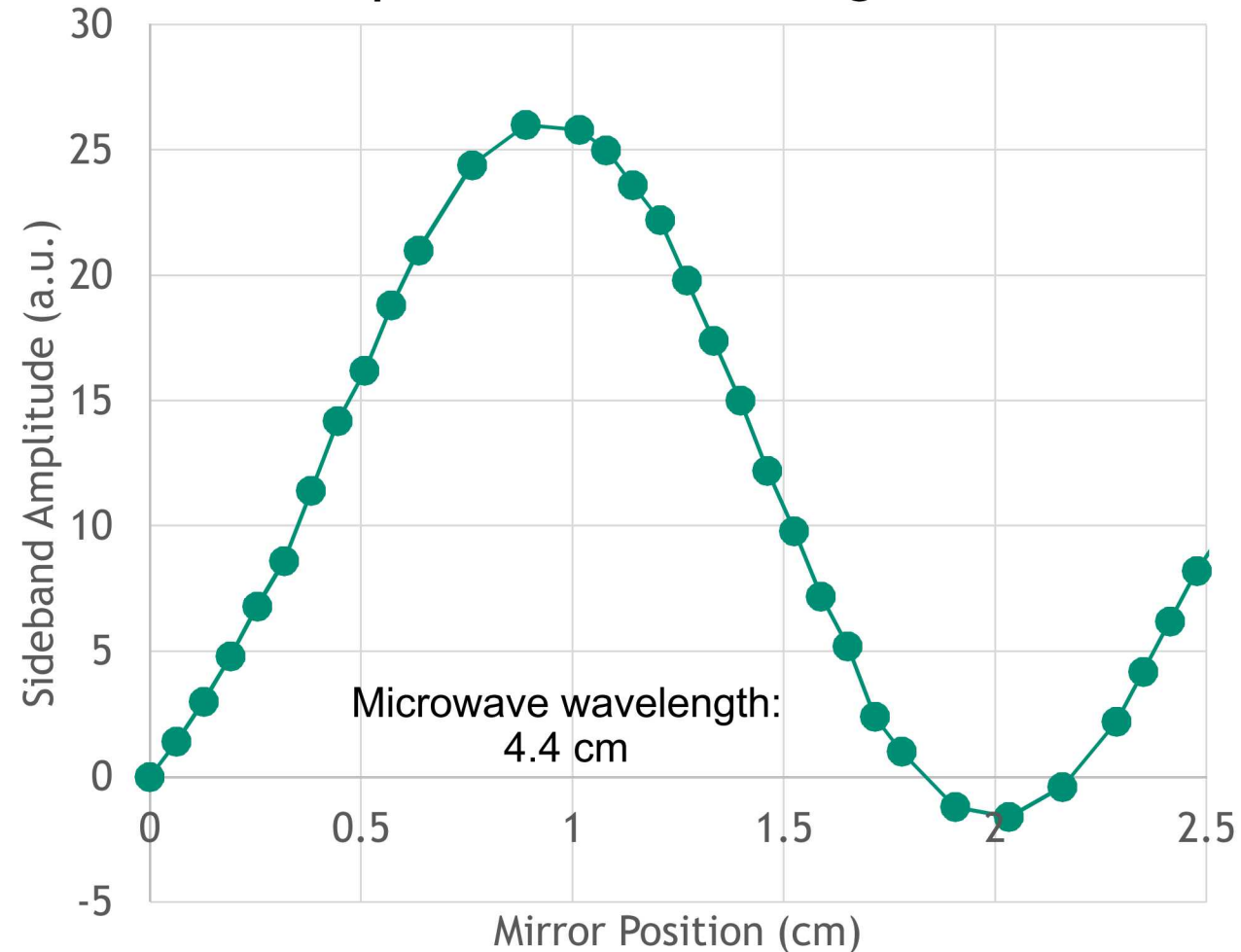
Sideband enhancement: Retroreflected probe laser



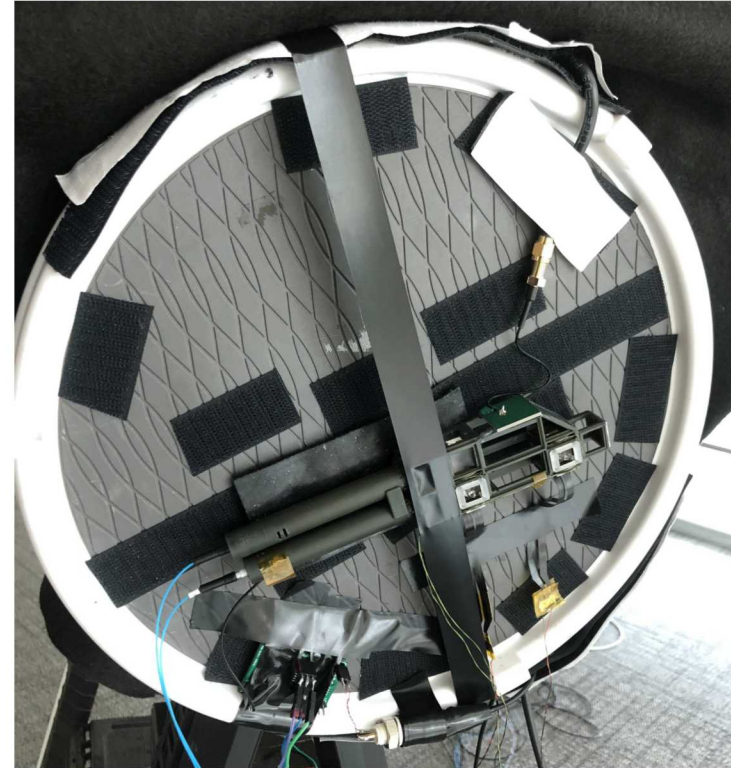
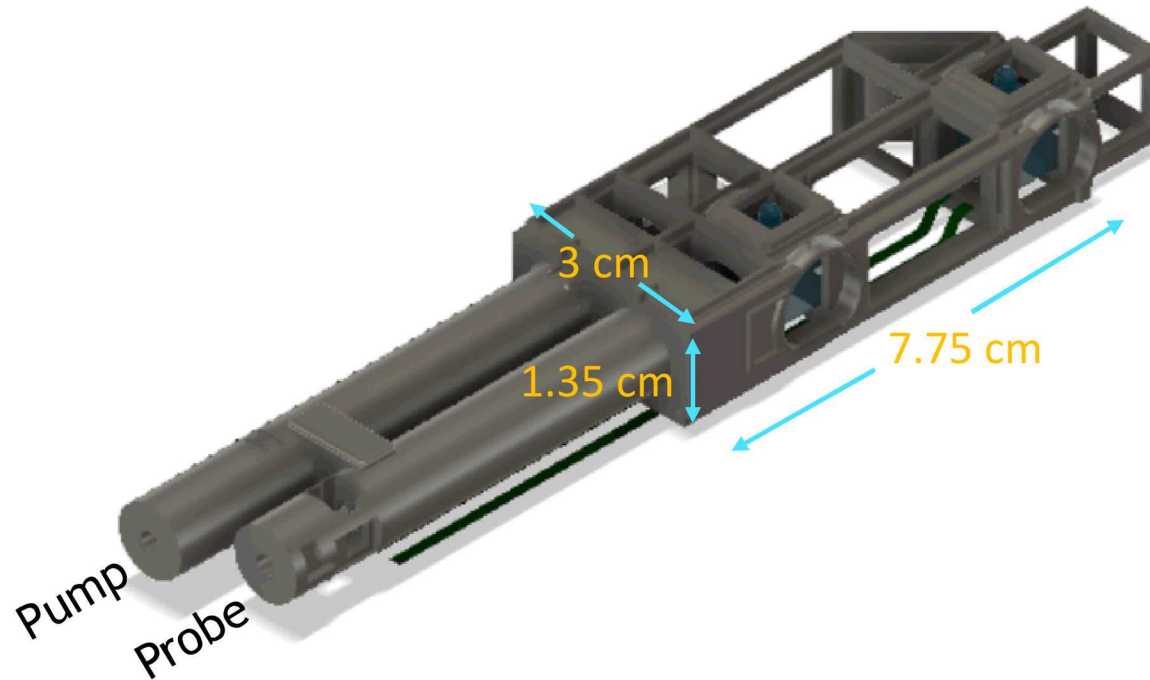
Single vs Double Pass



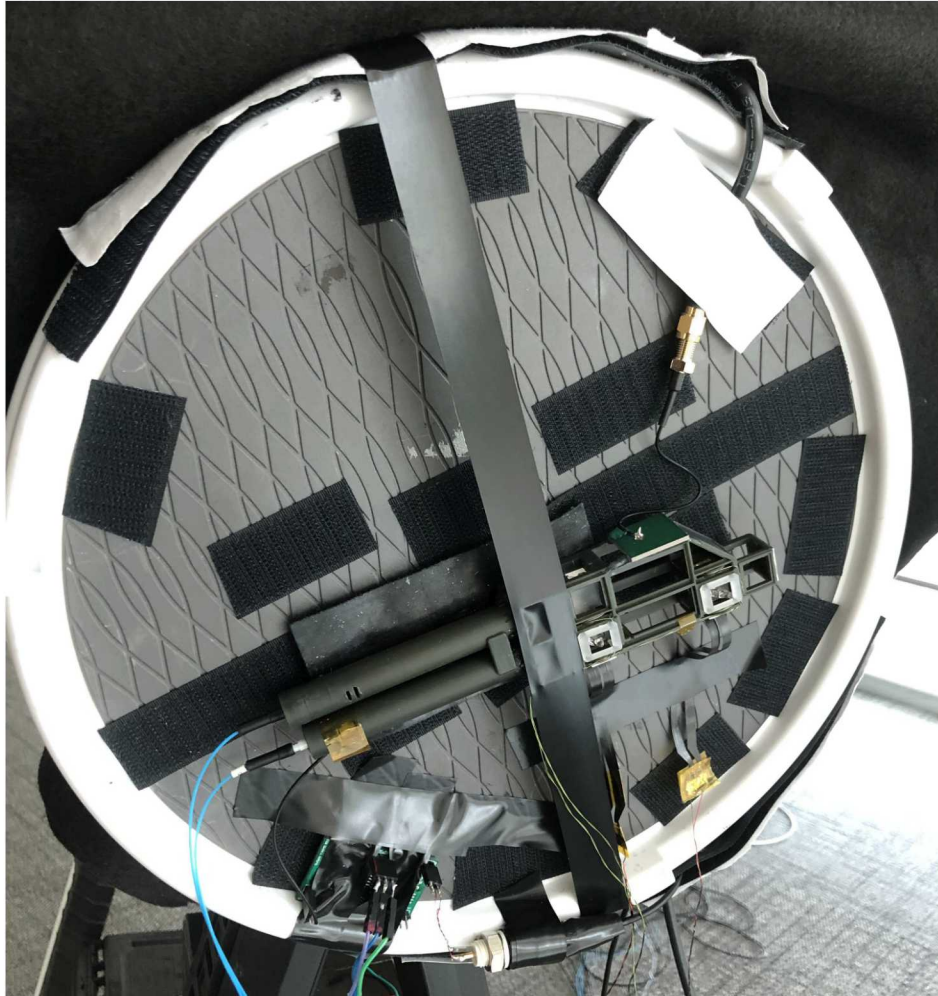
Double pass while moving the mirror



Prototype Design

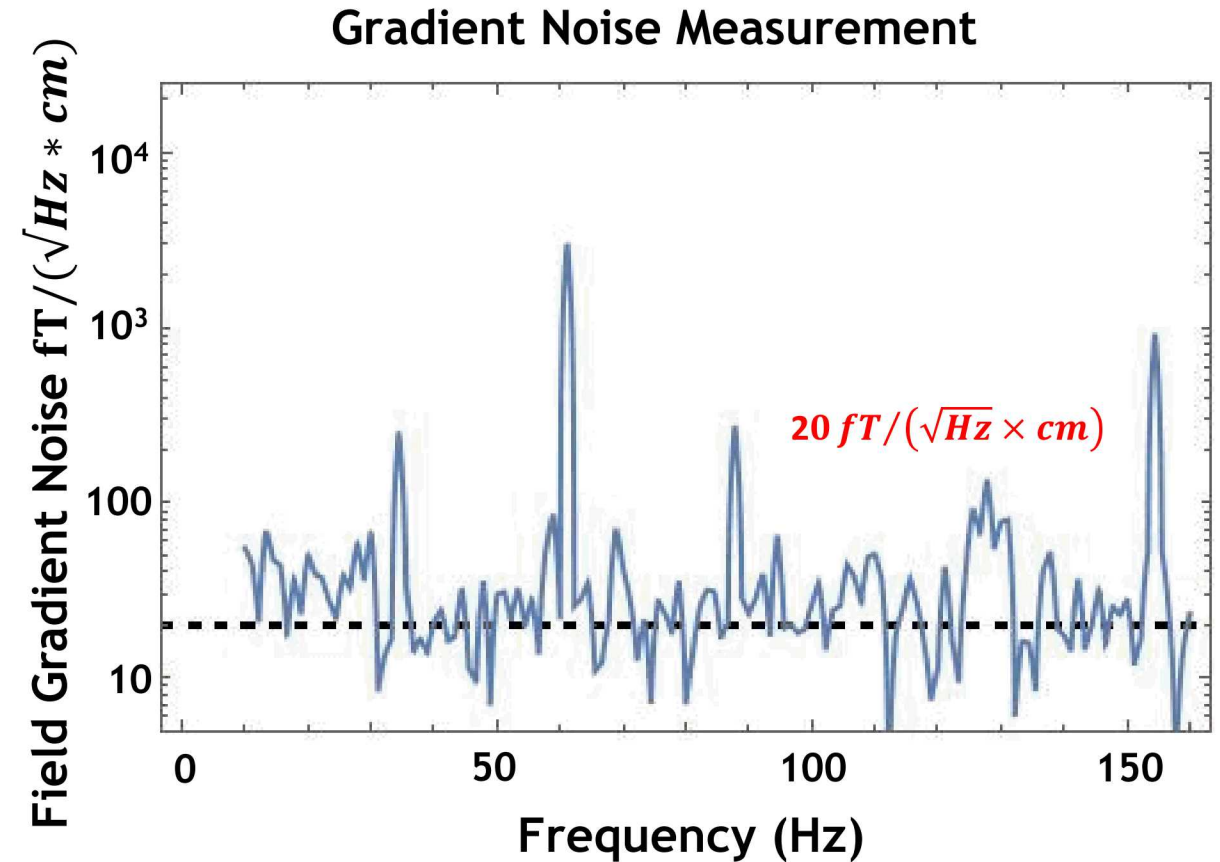


- All components are dropped-in design. Requires minimum alignment.
- Improved gradient cancellation coil design.



***^{87}Rb cells with 30 torr N_2 and 15 torr N_2
Separation: 4.4 cm***

Data and photo courtesy of QuSpin



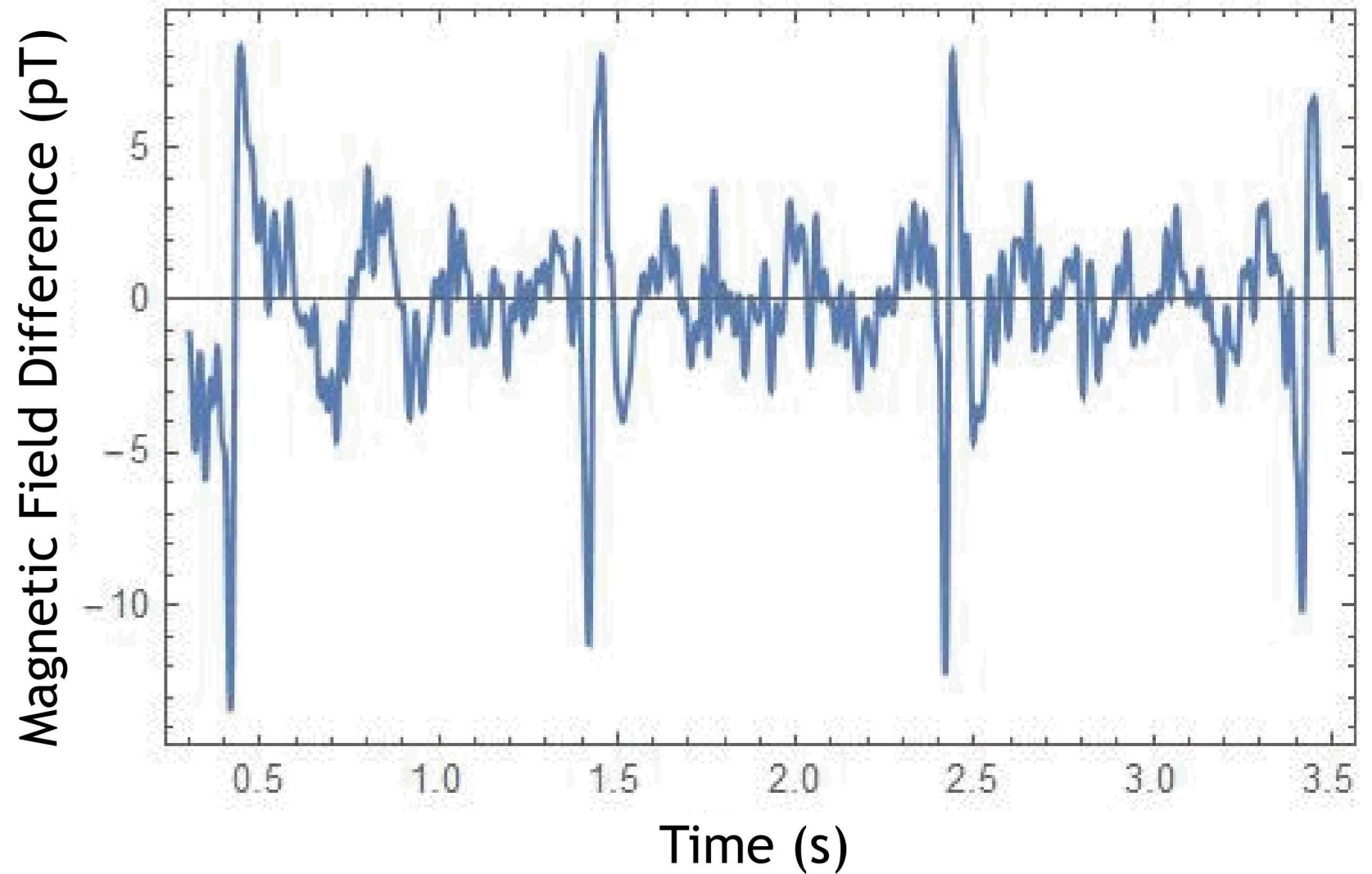
Data rate: 320 Hz

Photon Shot Noise: $\sim 6 \text{ fT}/\text{cm}/\text{rt-Hz}$

Heartbeat Measurement

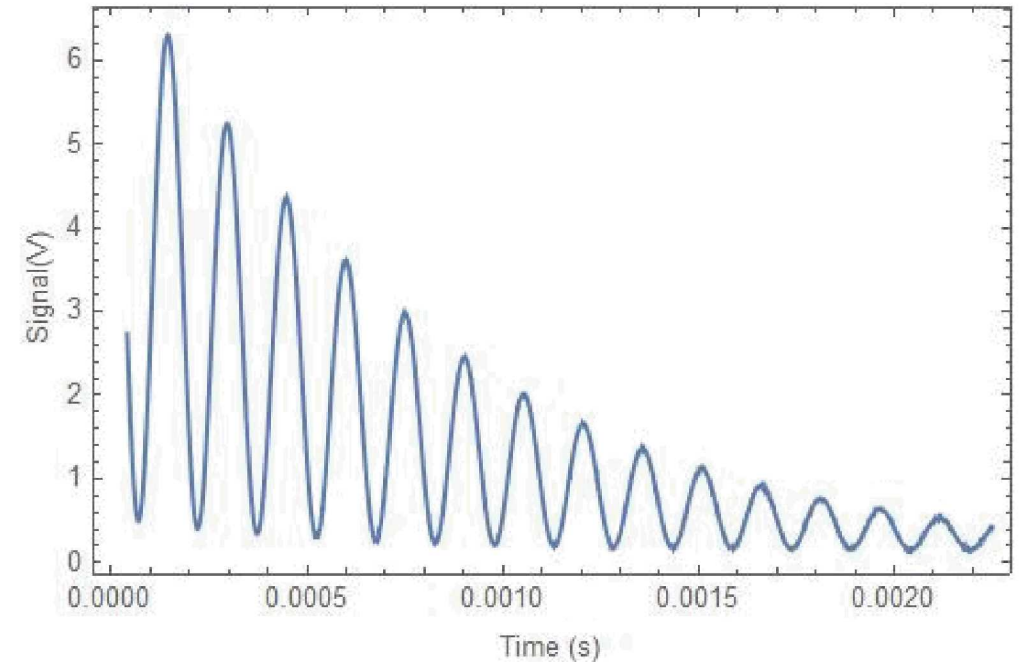
Heartbeat
measurement of an
adult male

Bandpass: 3-45 Hz



Conclusion

- Demonstrated beat note detection of a gradient
 - Use hyperfine structure
- Optimized sensitivity versus buffer gas pressure and temperature
- Earth's field sensitivity as low as $\frac{20 \text{ fT}}{\sqrt{\text{Hz} \cdot \text{cm}}}$.
- Next steps:
 - Make more compact
 - Dead-zone free



Acknowledgements

Experiment

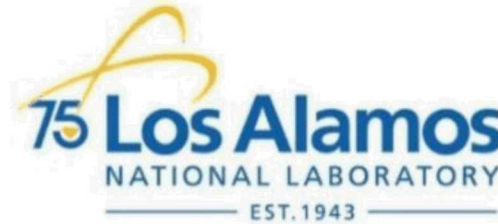


-Kaleb Campbell
-Peter D.D Schwindt



-Ying-Ju Wang
-Vishal Shah-PI

Theory



-Igor Savukov

Funding

**DARPA -Atomic Magnetometer for
Biological Imaging In Earth's Native
Terrain (AMBIENT)**

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Optical Beat Note Readout of a Magnetic Gradient

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²*QuSpin Inc., Louisville, CO, USA*

³*Los Alamos National Laboratories, Los Alamos, NM, USA*

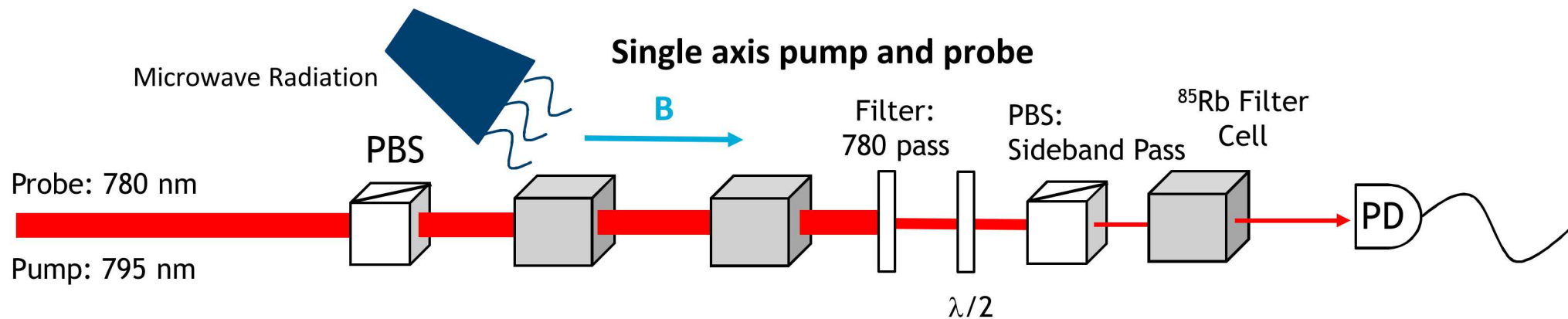
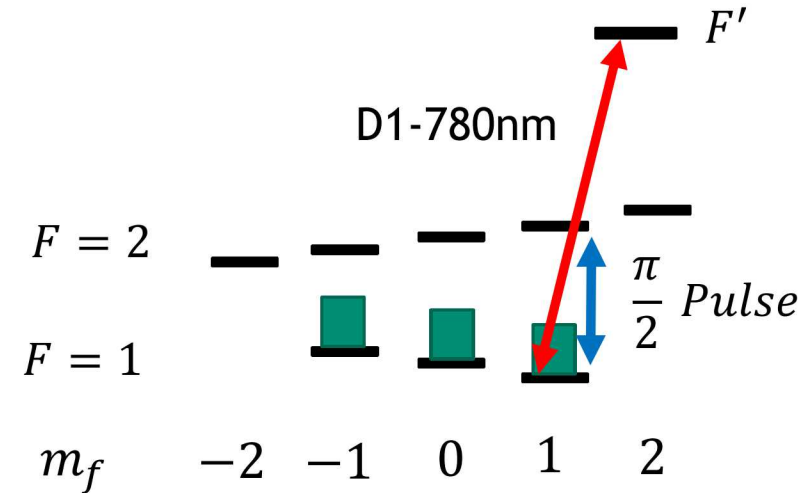
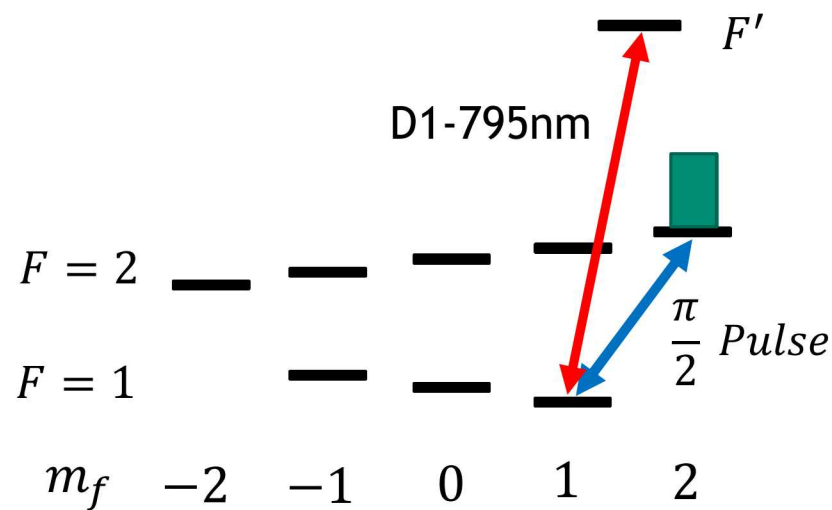
Typical gradient measurements using optically pumped magnetometers consist of making two independent measurements and then subtracting the results electronically. We present a technique where the gradient of the magnetic field is derived directly from the optical signal [1]. Using two ^{87}Rb vapor cells, the atoms are pumped into the $|F = 2, m_F = 2\rangle$ stretched state. Then, a resonant microwave pulse is applied to make a superposition between the $|2,2\rangle$ and $|1,1\rangle$ levels, and a resonant 780-nm probe beam is passed through the two vapor cells. With the atomic superposition precessing at the hyperfine splitting frequency, the probe laser will be modulated, parametrically generating an optical sideband [2]. If there is a magnetic field gradient between the two vapor cells, the sidebands will have a frequency difference and generate a beat note. Thus, the beat note frequency will be proportional the magnetic gradient. We will present an experimental implementation of this technique and describe efforts to improve the sensitivity and to eliminate dead zones of the gradiometer.

References

1. Vishal Shah, System and Method for Measuring a Magnetic Gradient Field. Patent. US10088535 (2018).
2. Tang, H. Parametric Frequency Conversion of Resonance Radiation in Optically Pumped Rb^{87} Vapor. Phys. Rev. A, **7**, 2010–2032 (1973).

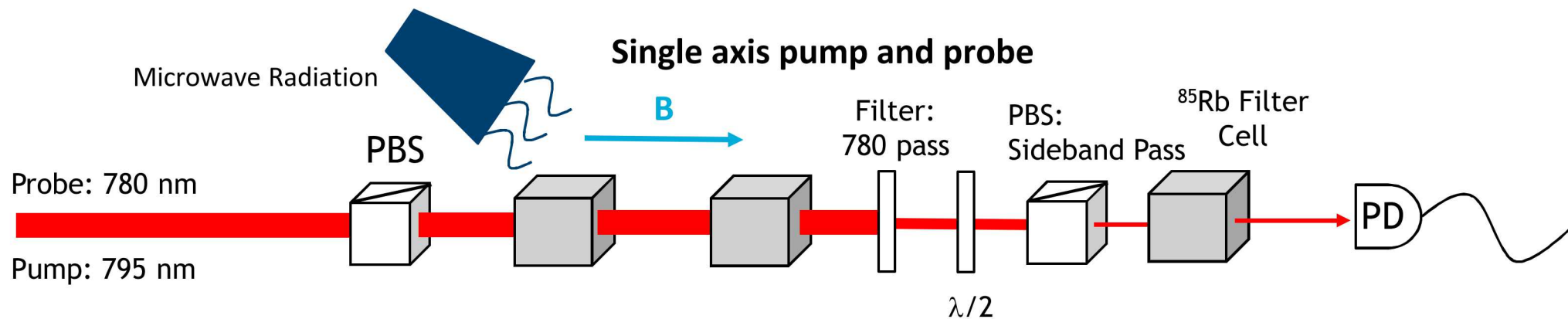
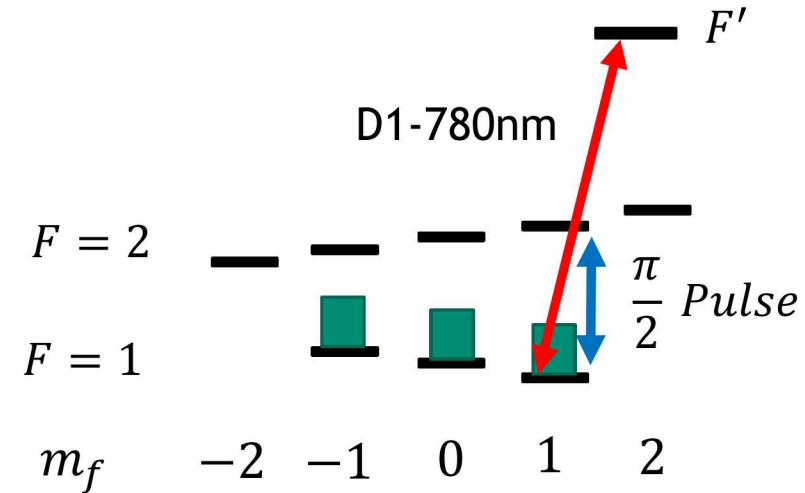
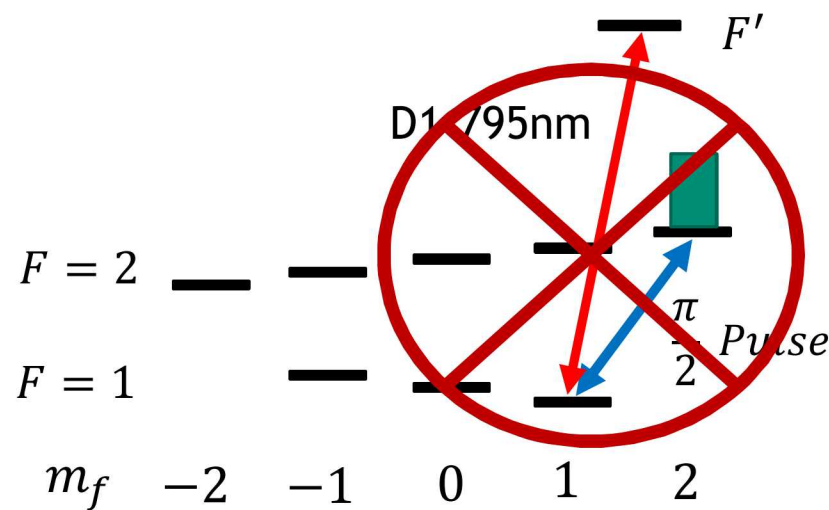
Can we make the sensor more compact?

- Pump and Carrier along the same axis?
- Selection rules prohibit sideband generation on the $|F=1, m_f = 1\rangle$ to $|F=2, m_f = 2\rangle$ transition.
- Selection rules allow only $\Delta m_F = 0$ transitions.



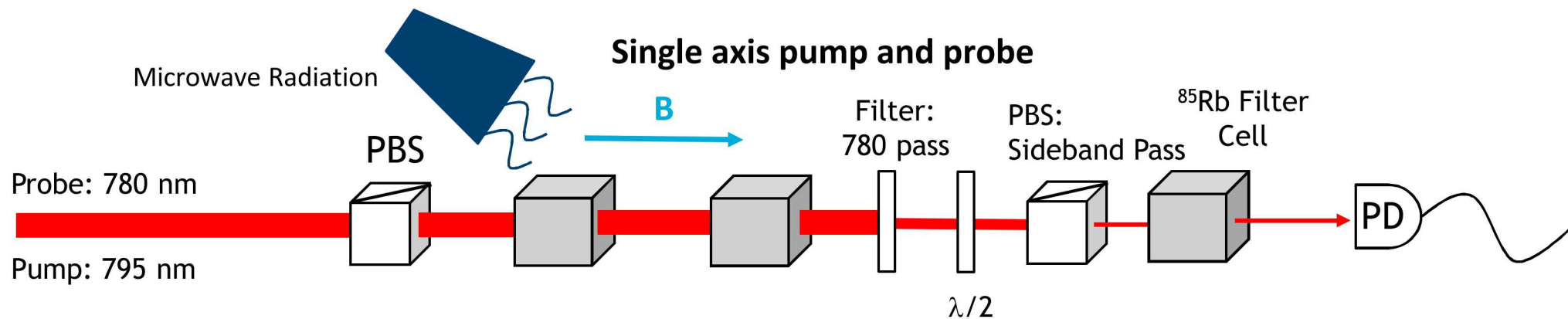
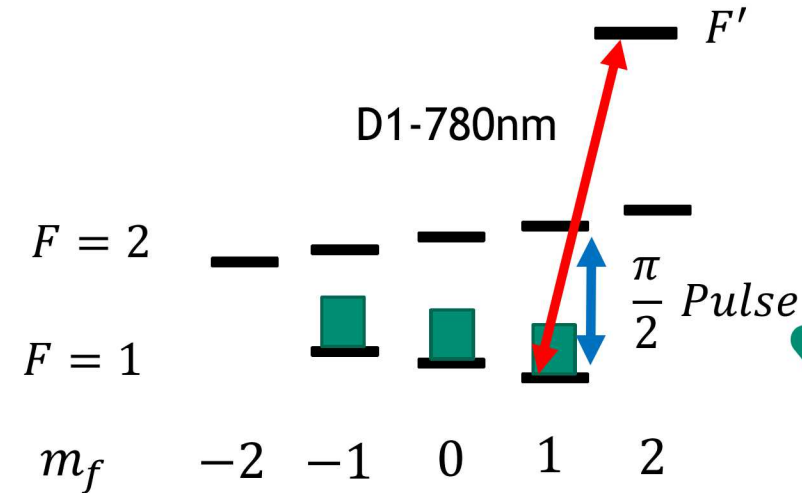
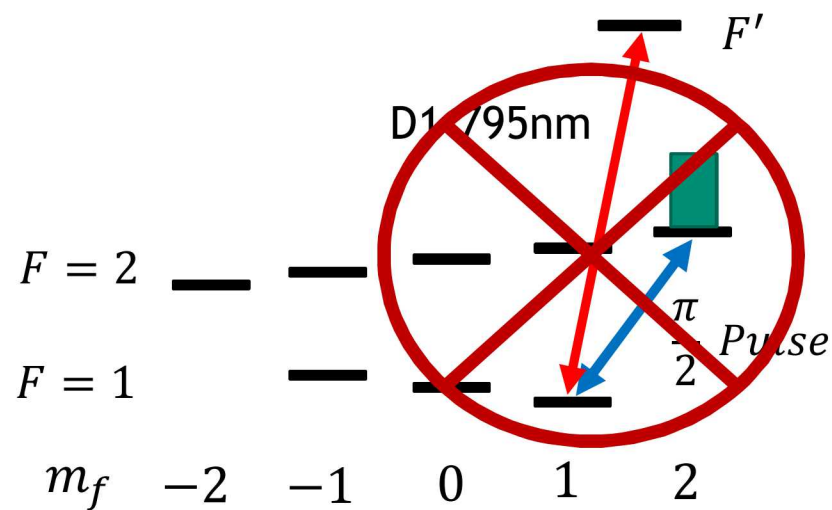
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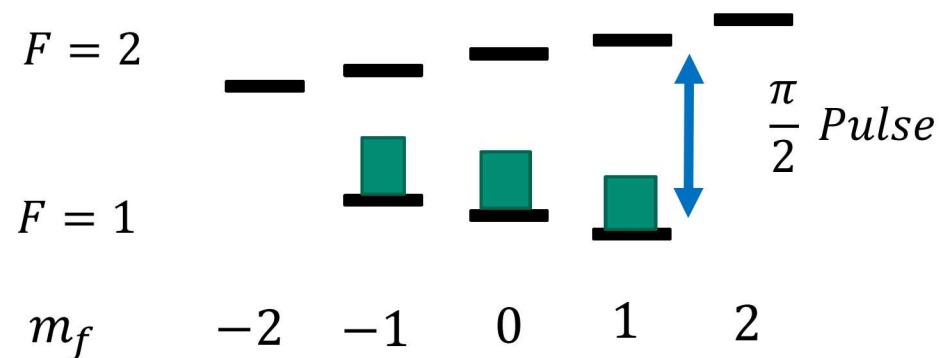
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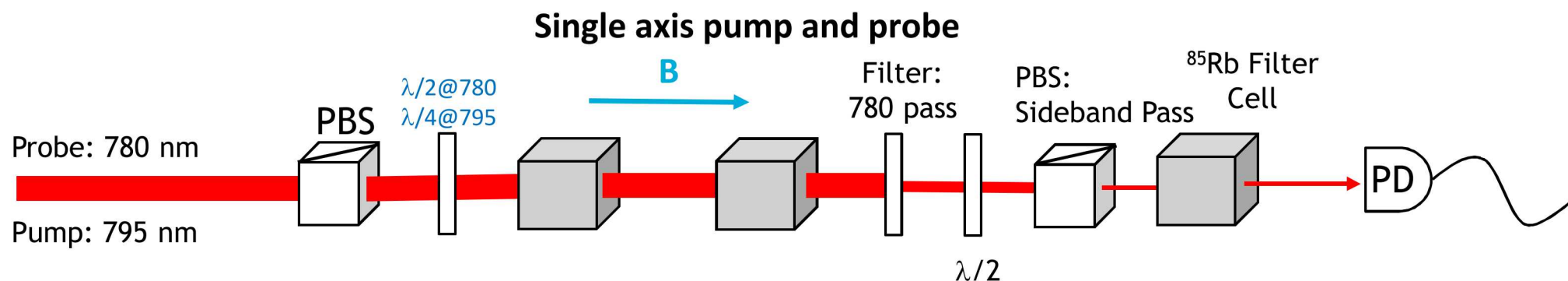


Problems with this orientation

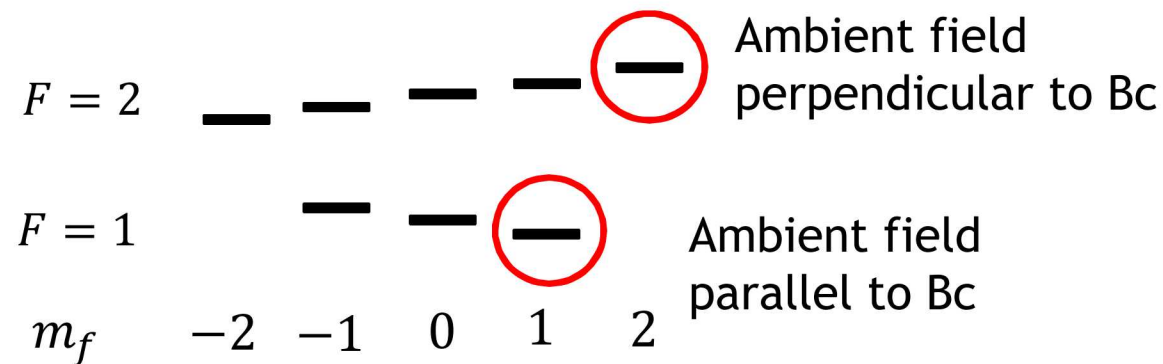
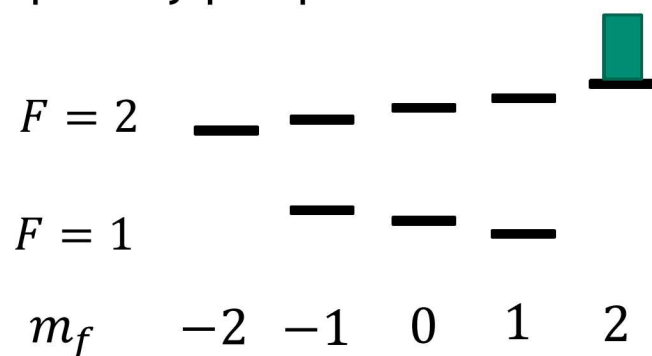
- We can only pump the atoms to a single ground state, not an individual Zeeman sublevel
- Reduced sideband amplitude



Solutions?

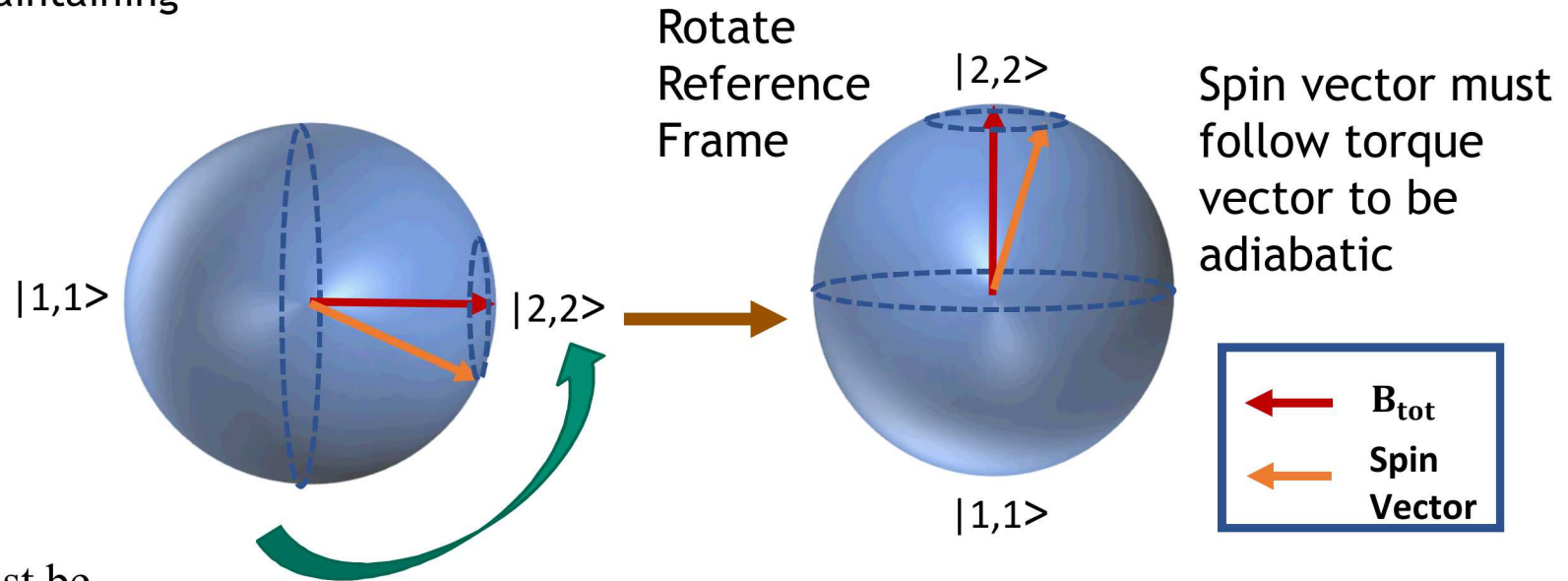
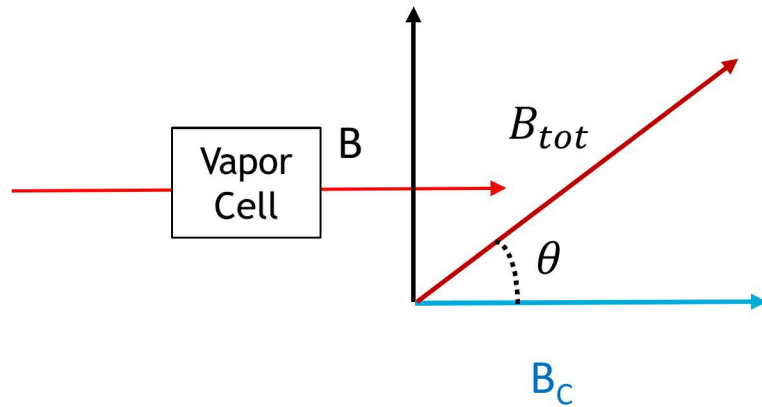


Atoms optically pumped to the end-state

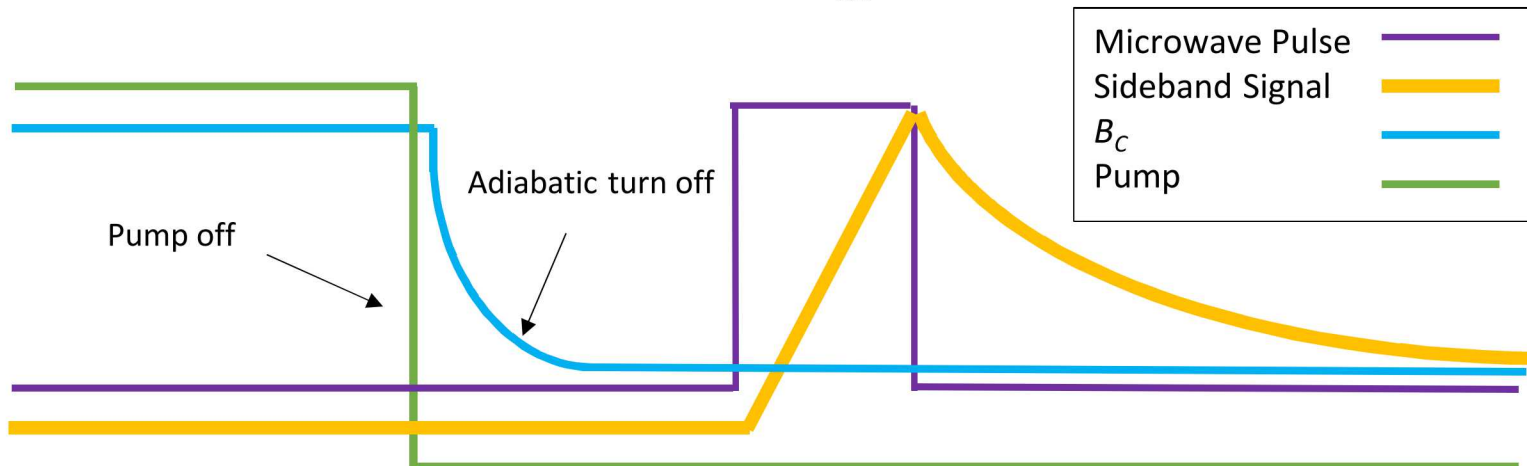


Case I: Ambient field perpendicular to the laser axis

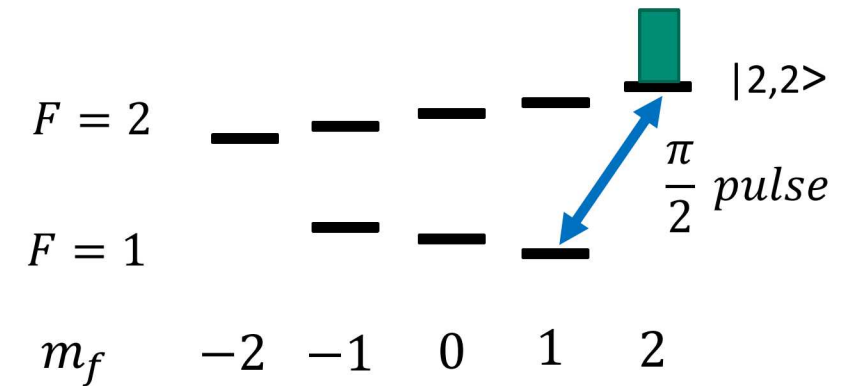
We rotate the quantization axis, while maintaining the atomic population in the $|2,2\rangle$ state.



To be adiabatic, the rotation rate of the field must be less than the Larmor precession frequency $\frac{d\theta}{dt} \ll \omega_L$.



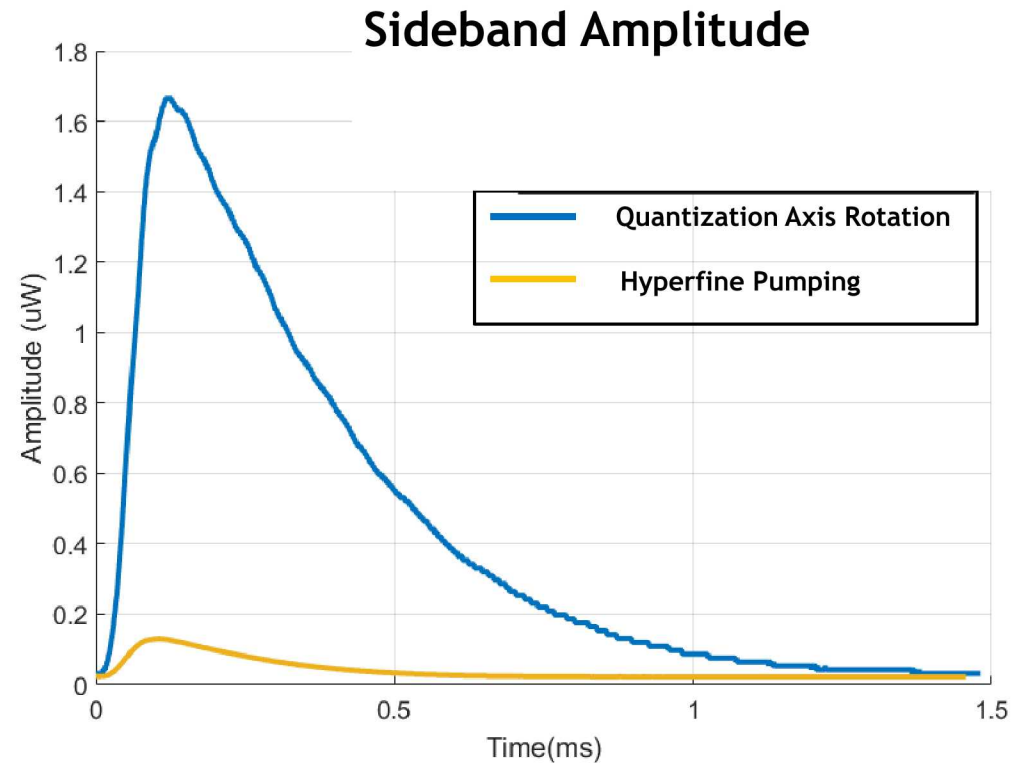
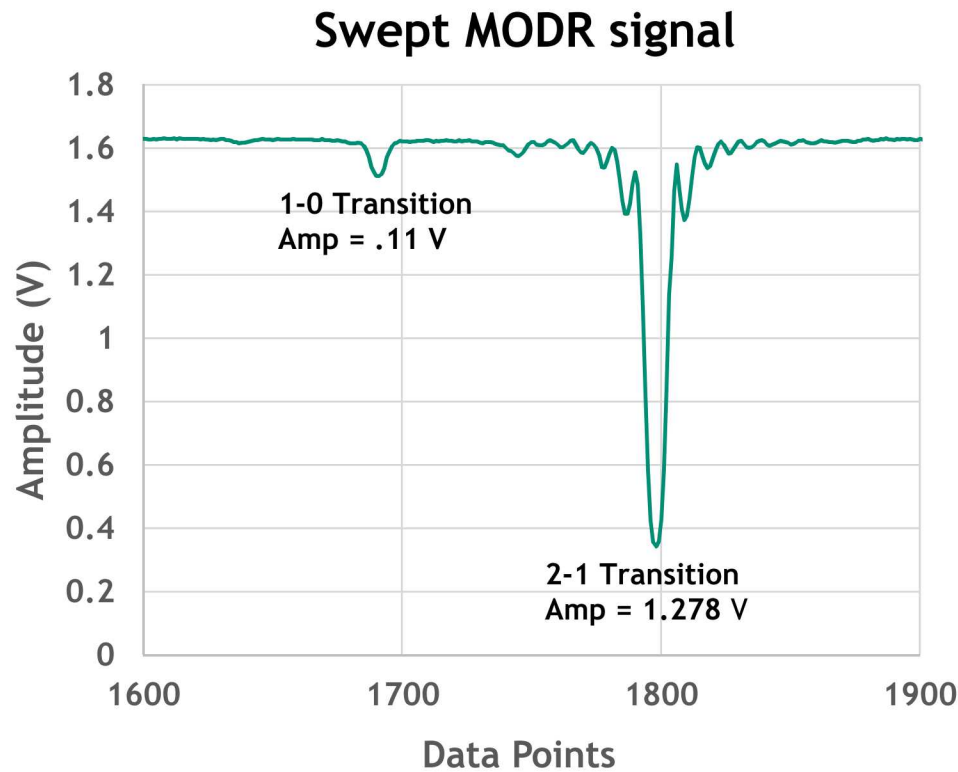
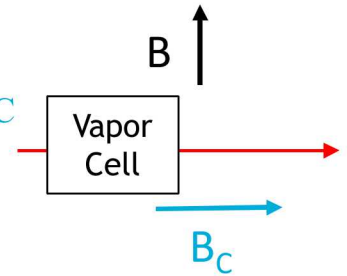
Selection rules now allow $\Delta m_F = 1$ transitions



Case I: Results

Compare adiabatic field switching to the old method of hyperfine pumping.

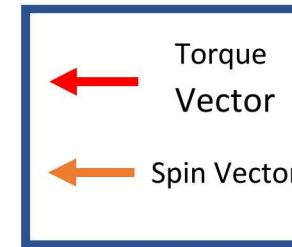
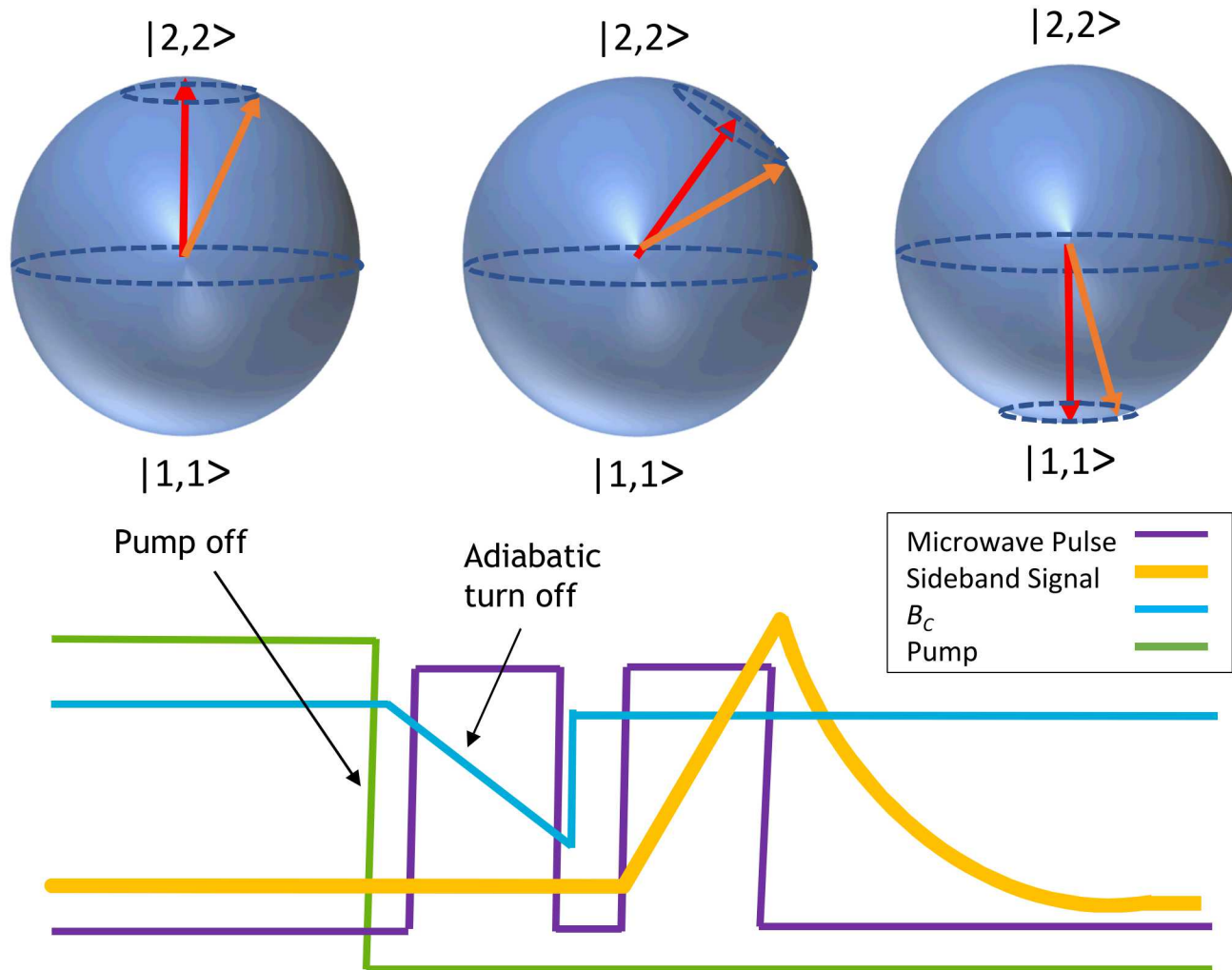
- Observe $\sim 90\%$ in the $|2,2\rangle$ state with optical pumping: expect improvement with larger B_C
- Adiabatic switch does not degrade populations
- See a 10x improvement in sideband size



Probe Power = $\sim 20 \mu\text{W}$; Detuning near $F = 1$; Using Rb-85 filter cell;
Buffer gas: 30 Torr N₂; Cell Temp = 98 °C;

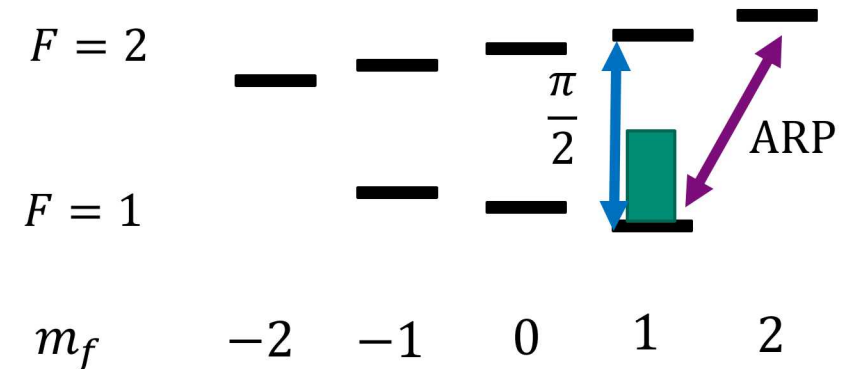
Case 2: Ambient field parallel to the laser axis

- We perform adiabatic rapid passage to transfer the population from the $|2,2\rangle$ state to the $|1,1\rangle$ state.
- We use a magnetic field ramp to simplify microwave and cover both cells.



Magnetic field must be ramped slow enough that the Spin vector follows the Torque vector

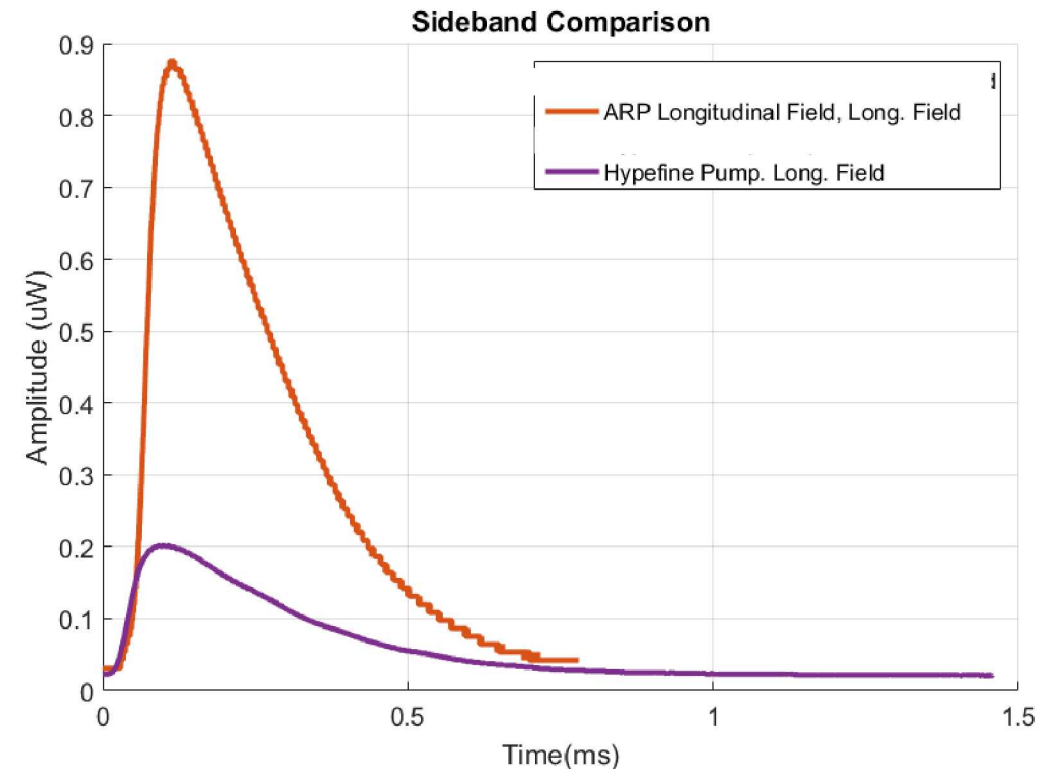
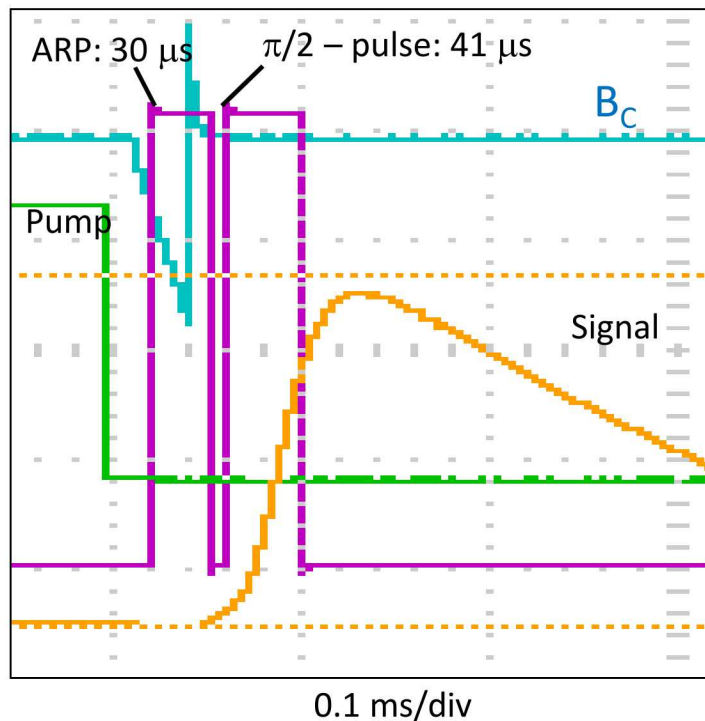
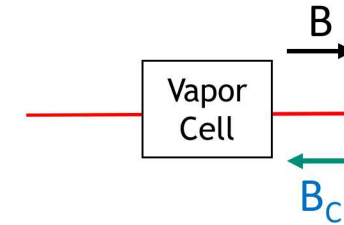
Selection rules allow $\Delta m_F = 0$ transitions



Field parallel to the laser axis

Compare adiabatic field switching to the old method of hyperfine pumping.

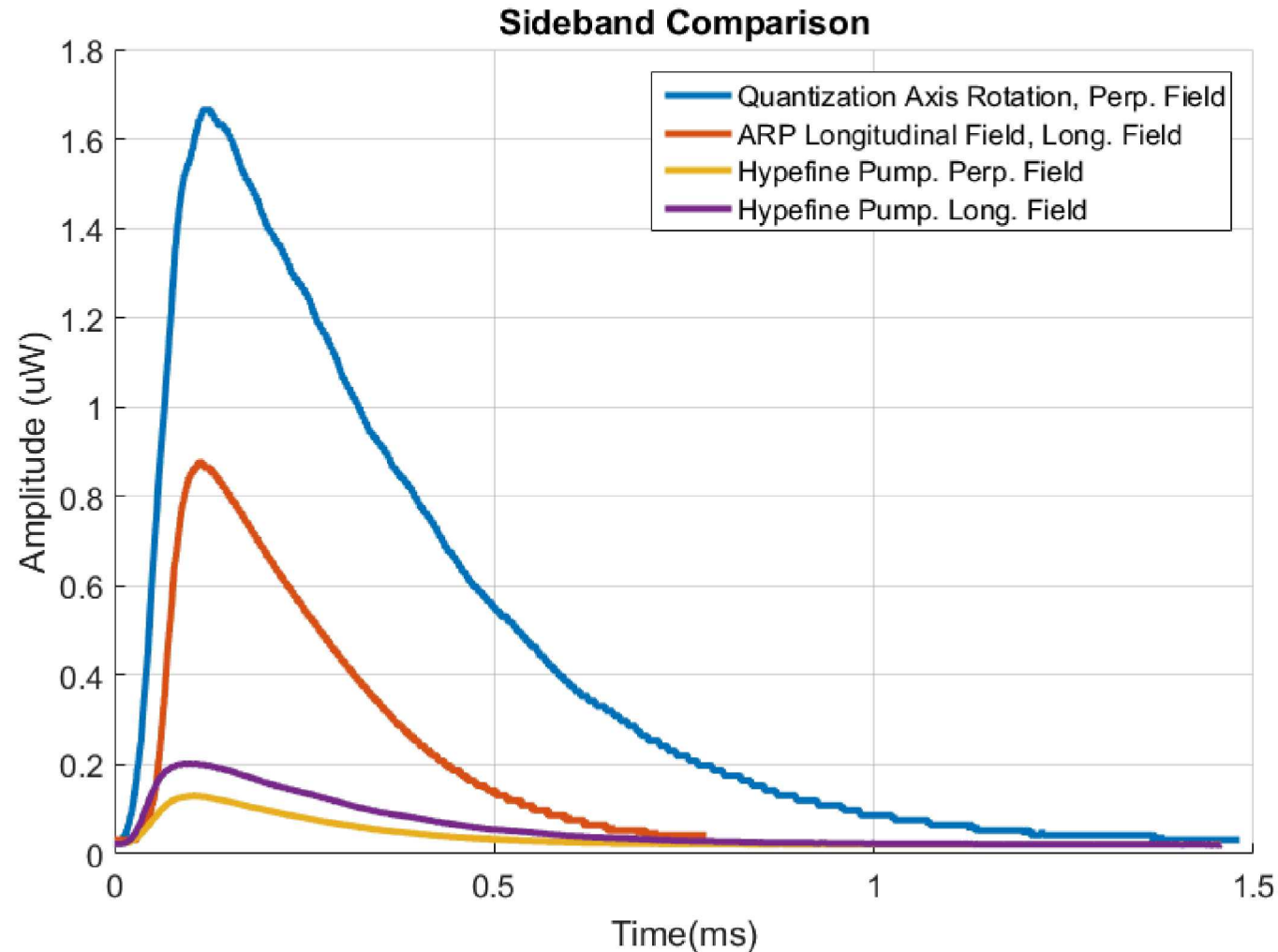
- Observe ARP from $|2,2\rangle$ to $|1,1\rangle$
 - Use a single microwave synthesizer
 - The percentage of the transfer uncertain
- See a 4x improvement in sideband size on the $|1,1\rangle$ to $|2,1\rangle$ transition.



Comparison of four experiments

Quantization axis rotation currently working better than microwave ARP

- Why is this?



Can we work towards dead-zone free operation?

Minimum field

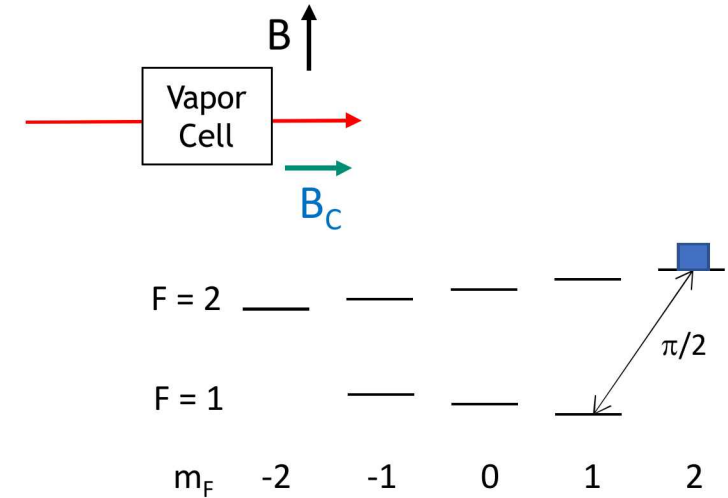
- $f_{1 \rightarrow 2} - f_{1 \rightarrow 1} \gg \frac{1}{2\pi T_{\pi/2}}$
- For $T_{\pi/2} = 0.1 \text{ ms}$, $B_{min} \gg 230 \text{ nT}$

How to switch between the two schemes

- Start up:
 1. Determine ambient field and direction using a field zeroing scheme
 2. Select scheme, microwave frequency, and direction of B_C
 3. Begin operation.
- Continuous operation:
 1. Monitor signal size
 2. If signal size drops below threshold, switch scheme.
 1. If this fails, re-zero field.

Need to understand better how to ramp B_C for ARP

Quantization axis rotation



ARP with longitudinal field

