

Multi-Fluid Plasma Simulation with an IMEX CG/DG Discretization



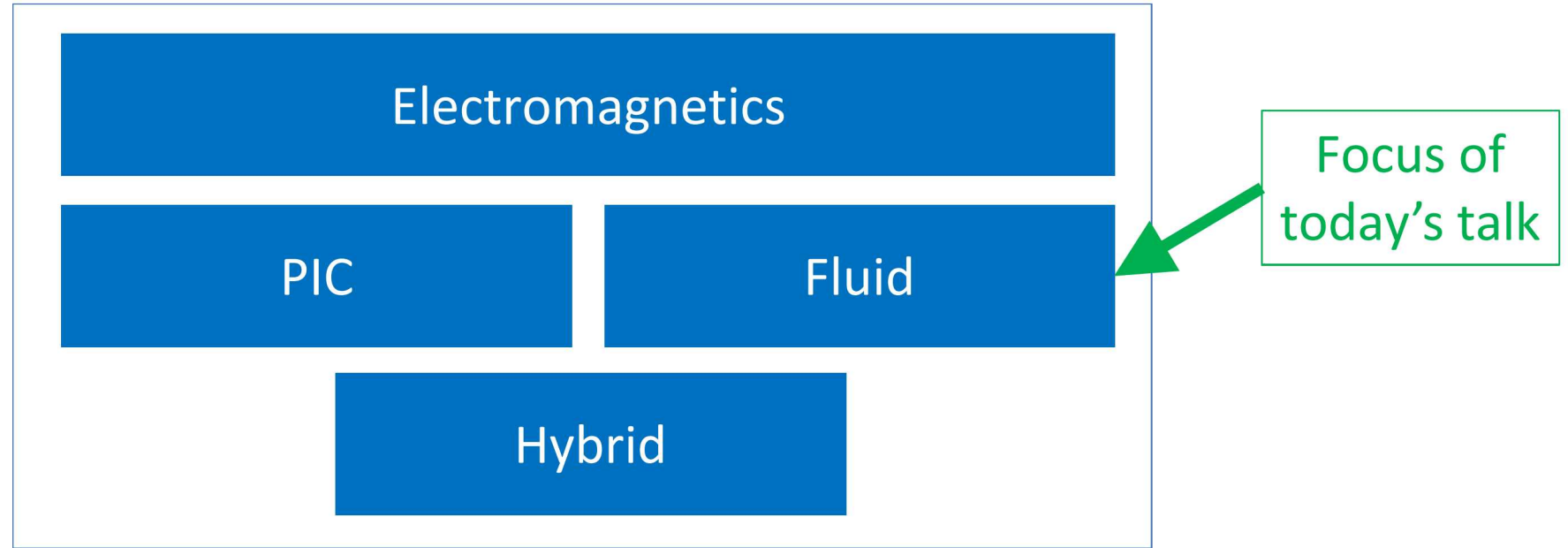
PRESENTED BY

Eric C. Cyr, Sean Miller, Edward Phillips, Kris Beckwith, Matt Bettencourt, Keith Cartwright, Richard Kramer, Roger P. Pawlowski, Nathan V. Roberts, John N. Shadid, Sidney Shields



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EMPIRE: A hierarchy of capabilities



- EMPIRE's goal: Accurately simulate plasmas across regimes on next-generation exascale computing platforms
- Expand the range of electromagnetic pulse and Z-power flow applications that we can simulate with high confidence and fidelity

Multi-Fluid Plasma Formulation



- Multi-species Euler coupled to Maxwell
- Strong collisions terms
- Maxwell involutions must be enforced

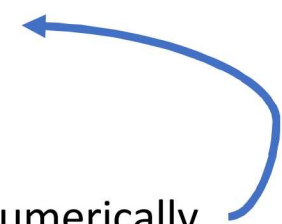
5-Moment Fluid

$$\begin{aligned}
 \frac{\partial \rho_\alpha}{\partial t} + \nabla \cdot (\rho_\alpha \mathbf{u}_\alpha) &= \sum_{\text{srcs}} m_\alpha \Gamma^{\text{src}} - \sum_{\text{sinks}} m_\alpha \Gamma^{\text{sink}} \\
 \frac{\partial (\rho_\alpha \mathbf{u}_\alpha)}{\partial t} + \nabla \cdot (\rho_\alpha \mathbf{u}_\alpha \otimes \mathbf{u}_\alpha + p_\alpha I + \Pi_\alpha) &= \frac{q_\alpha}{m_\alpha} \rho_\alpha (\mathbf{E} + \mathbf{u}_\alpha \times \mathbf{B}) \\
 &\quad + \sum_{\text{srcs}} m_\alpha \mathbf{u}_{\text{src}} \Gamma^{\text{src}} - \sum_{\text{sinks}} m_\alpha \mathbf{u}_\alpha \Gamma^{\text{sink}} + \sum_{\beta \neq \alpha} \mathbf{R}^{\alpha, \beta} \\
 \frac{\partial \mathcal{E}_\alpha}{\partial t} + \nabla \cdot ((\mathcal{E}_\alpha + p_\alpha) \mathbf{u}_\alpha + \mathbf{u}_\alpha \cdot \Pi_\alpha + \mathbf{h}_\alpha) &= \frac{q_\alpha}{m_\alpha} \rho_\alpha \mathbf{E} \cdot \mathbf{u}_\alpha + \sum_{\beta \neq \alpha} (\mathbf{u}_\alpha \mathbf{R}^{\alpha, \beta} + Q^{\alpha, \beta}) \\
 &\quad + \frac{1}{2} \sum_{\text{srcs}} m_\alpha u_{\text{src}}^2 \Gamma^{\text{src}} - \frac{1}{2} \sum_{\text{sinks}} m_\alpha u_\alpha^2 \Gamma^{\text{sink}}
 \end{aligned}$$

Maxwell Equations

$$\begin{aligned}
 \frac{\partial \mathbf{E}}{\partial t} - c^2 \nabla \times \mathbf{B} &= -\frac{1}{\epsilon_0} \sum_\alpha \frac{q_\alpha}{m_\alpha} \rho_\alpha \mathbf{u}_\alpha & \nabla \cdot \mathbf{E} &= \frac{1}{\epsilon_0} \sum_\alpha \frac{q_\alpha}{m_\alpha} \rho_\alpha \\
 \frac{\partial \mathbf{B}}{\partial t} + \nabla \times \mathbf{E} &= 0 & \nabla \cdot \mathbf{B} &= 0
 \end{aligned}$$

Important to satisfy involutions numerically



Multiple time scales (yikes!)

Plasma models are replete with multi-scale phenomena:

- Strongly dependent on species mass, density, and temperature
- Speed of light, plasma and cyclotron frequency are often stiff!
- Can be broken into **frequency**, **velocity**, and **diffusion (not used here)** scales:

Plasma frequency

$$\omega_{p\alpha} = \sqrt{\frac{q_\alpha^2 n_\alpha}{m_\alpha \epsilon_0}}$$

Cyclotron frequency

$$\omega_{c\alpha} = \frac{q_\alpha B}{m_\alpha}$$

Collision frequency

$$\nu_{\alpha\beta} \sim \frac{n_\beta}{\sqrt{m_\alpha} T_\alpha^{\frac{3}{2}}} \frac{1 + \frac{m_\alpha}{m_\beta}}{\left(1 + \frac{m_\alpha}{m_\beta} \frac{T_\beta}{T_\alpha}\right)^{\frac{3}{2}}}$$

Flow velocity

$$u_\alpha$$

Speed of sound

$$v_{s\alpha} = \sqrt{\frac{\gamma P_\alpha}{\rho_\alpha}}$$

Speed of light

$$c \gg u_\alpha, v_{s\alpha}$$

Momentum diffusivity

$$\nu_\alpha = \frac{\mu_\alpha}{\rho_\alpha}$$

Thermal diffusivity

$$\kappa_\alpha \sim \frac{k_\alpha}{\rho_\alpha}$$

Take home: These plasmas are hard to simulate!

Discretization Tools

We have (at least) two major challenges:

1. Involutions from Maxwell's equations
2. Multiple time scales

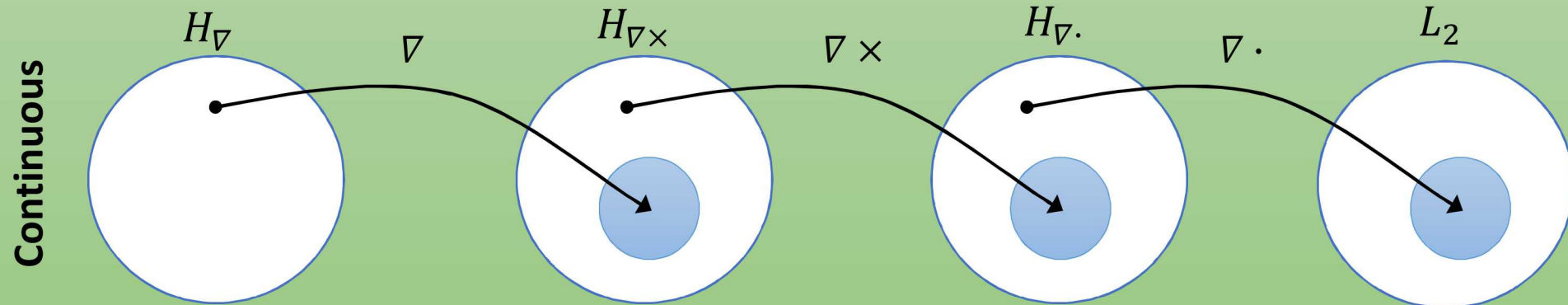
We will attack each of these in turn with two discretization tools

1. “Exact-Sequence” discretizations to structurally enforce involutions
2. Implicit-Explicit (IMEX) time integration to handle multiple time scales

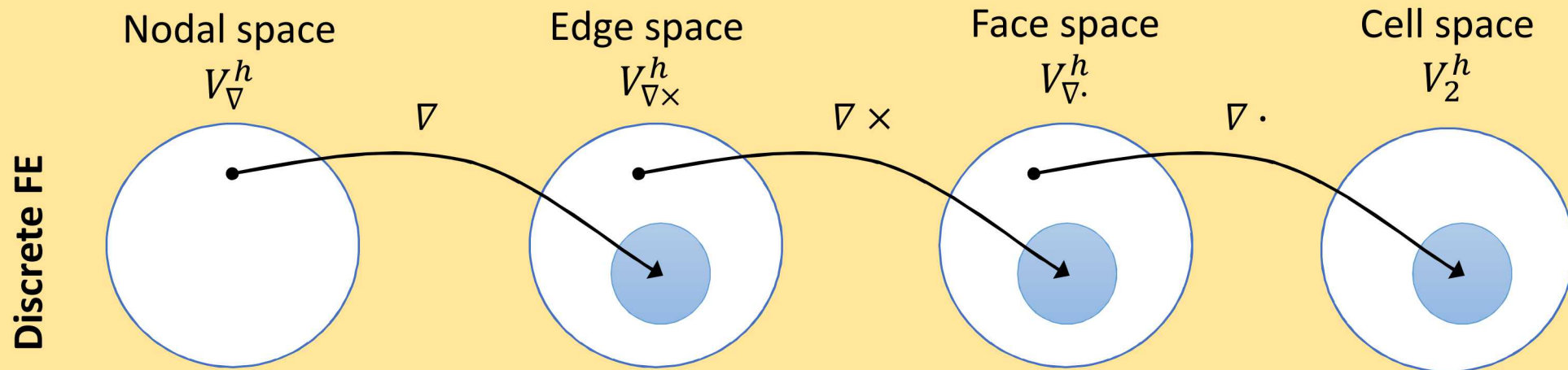
The interesting part is how these interact! (Not to mention boundary conditions! Which I won't talk about)

Exact-Sequence Discretizations

Function spaces possess an exact sequence property where the derivative maps into the next space, e.g.:



Exact sequence finite elements have been constructed¹ (note $V_*^h \subset H_*$):



¹Bochev, P., Edwards, H. C., Kirby, R. C., Peterson, K., & Ridzal, D. (2012). Solving PDEs with intrepid. *Scientific Programming*, 20(2), 151-180.

Exact-Sequence and Multi-Fluids

Discretization:
Fields continuous by
construction (CG)

$$\begin{bmatrix} \rho_\alpha \\ \rho_\alpha u_\alpha \\ \varepsilon_\alpha \end{bmatrix} \in V_\nabla^h \quad E \in V_{\nabla \times}^h \quad B \in V_{\nabla \cdot}^h$$

Discrete FE

Nodal space

V_∇^h

Edge space

$V_{\nabla \times}^h$

Face space

$V_{\nabla \cdot}^h$

Cell space

V_2^h

∇

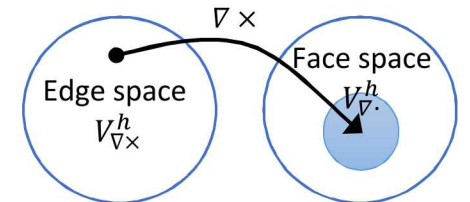
$\nabla \times$

$\nabla \cdot$

No magnetic monopoles example: Let $\mathbf{B}^h \in V_{\nabla \cdot}^h$ and $\mathbf{E}^h \in V_{\nabla \times}^h$, then the argument is straightforward and follows the continuous case:

$$\nabla \cdot \left(\frac{\partial \mathbf{B}^h}{\partial t} + \nabla \times \mathbf{E}^h \right) = 0 \Rightarrow \nabla \cdot \partial_t \mathbf{B}^h = 0$$

$$\Rightarrow \nabla \cdot \mathbf{B}^h = 0 \text{ (assuming satisfied at } t = 0 \text{)}$$



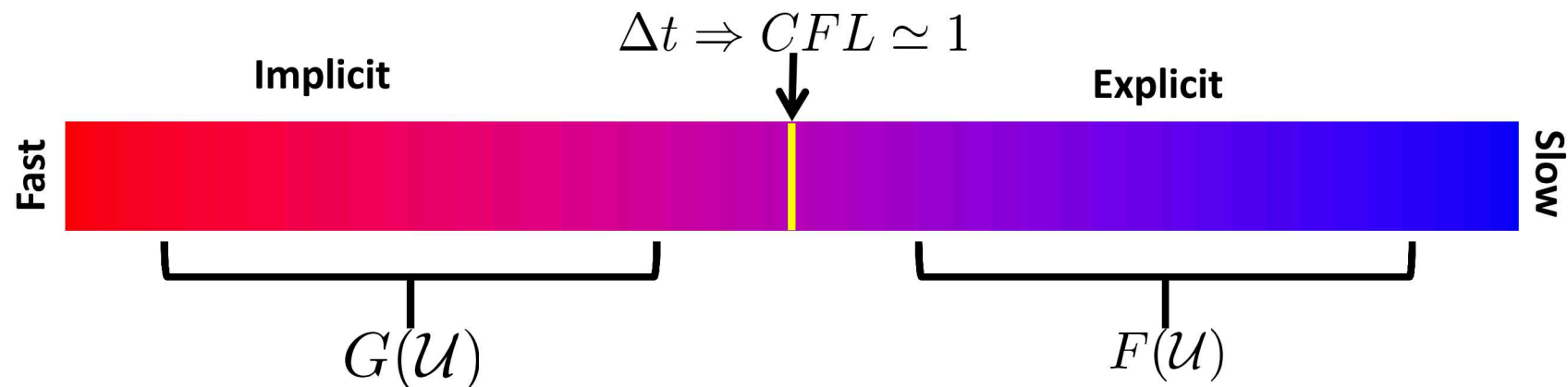
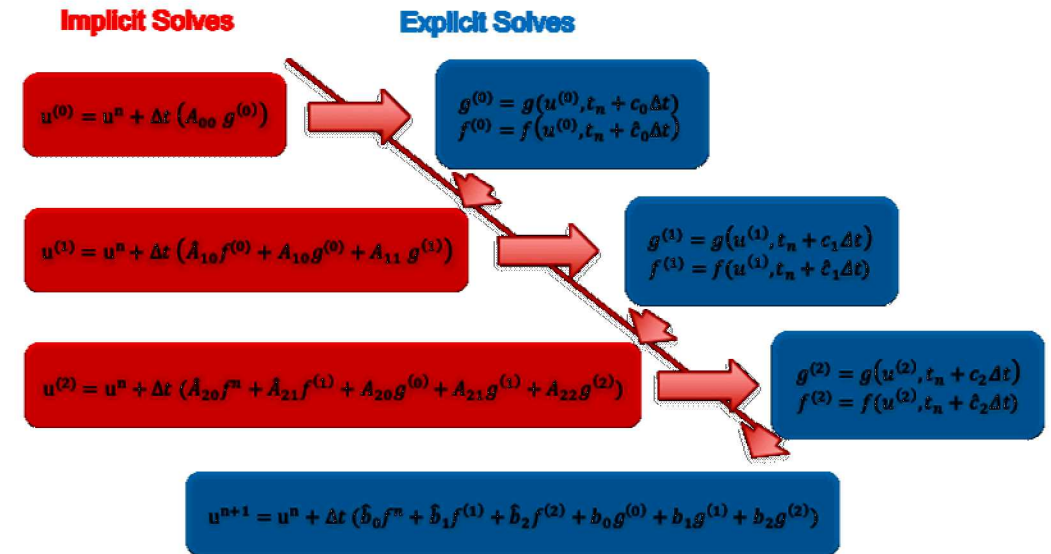
Implicit-Explicit (IMEX) Time Integration

IMEX methods split fast and slow modes

- Implicit terms solve for stiff modes (plasma oscillation, speed of light)
- Explicit terms are accurately resolved
- Combine with block/physics-based preconditioning for implicit solves
- IMEX assumes an additive decomposition:

$$\dot{\mathcal{U}} + F(\mathcal{U}) + G(\mathcal{U}) = 0$$

3 Stage IMEX-RK Algorithm



Fast/Stiff/Implicit modes in plasma model

Stiff Modes:

- Speed of light
- Plasma Oscillation
- Collisions
- Cyclotron frequency

$$\frac{\partial \rho_\alpha}{\partial t} + \nabla \cdot (\rho_\alpha \mathbf{u}_\alpha) = \sum_{\text{srcs}} m_\alpha \Gamma^{\text{src}} - \sum_{\text{sinks}} m_\alpha \Gamma^{\text{sink}}$$

$$\frac{\partial(\rho_\alpha \mathbf{u}_\alpha)}{\partial t} + \nabla \cdot (\rho_\alpha \mathbf{u}_\alpha \otimes \mathbf{u}_\alpha + p_\alpha I + \Pi_\alpha) = \frac{q_\alpha}{m_\alpha} \rho_\alpha (\mathbf{E} + \mathbf{u}_\alpha \times \mathbf{B}) + \sum_{\text{srcs}} m_\alpha \mathbf{u}_{\text{src}} \Gamma^{\text{src}} - \sum_{\text{sinks}} m_\alpha \mathbf{u}_\alpha \Gamma^{\text{sink}} + \sum_{\beta \neq \alpha} \mathbf{R}^{\alpha, \beta}$$

$$\frac{\partial \mathbf{E}}{\partial t} - c^2 \nabla \times \mathbf{B} = -\frac{1}{\epsilon_0} \sum_{\alpha} \frac{q_\alpha}{m_\alpha} \rho_\alpha \mathbf{u}_\alpha$$

$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \times \mathbf{E} = 0$$

- Speed of light arises from coupling of electromagnetic field: explicit CFL $\sim c\Delta t/\Delta x$
- Plasma oscillation arises from Ampere's law to momentum conservation: explicit CFL $\sim \Delta t$
- Collisions explicit CFL $\sim \Delta t$
- Cyclotron frequency explicit CFL $\sim |B|\Delta t$

$$\nabla \cdot (\rho_\alpha \mathbf{u}_\alpha)$$

If the plasma oscillation is implicit, then the mass flux needs to be implicit to maintain Gauss' law

Two Fluid Plasma Vortex (from Drekar)*

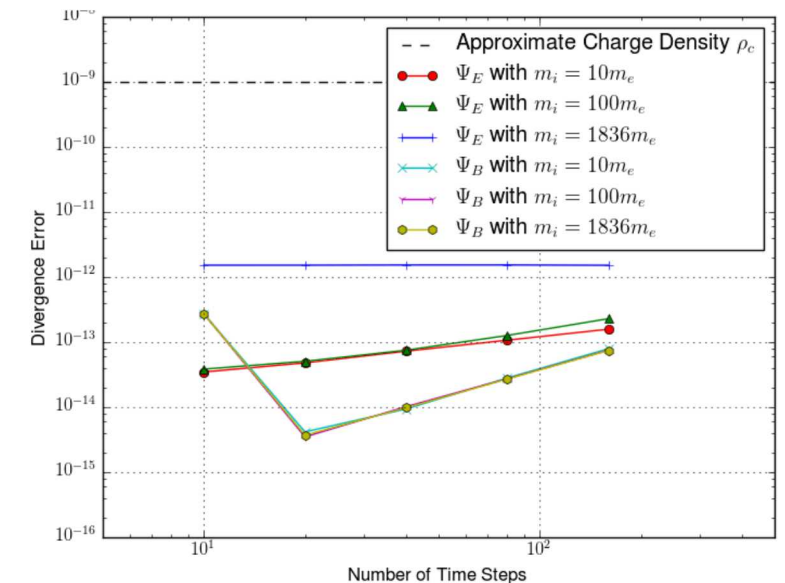
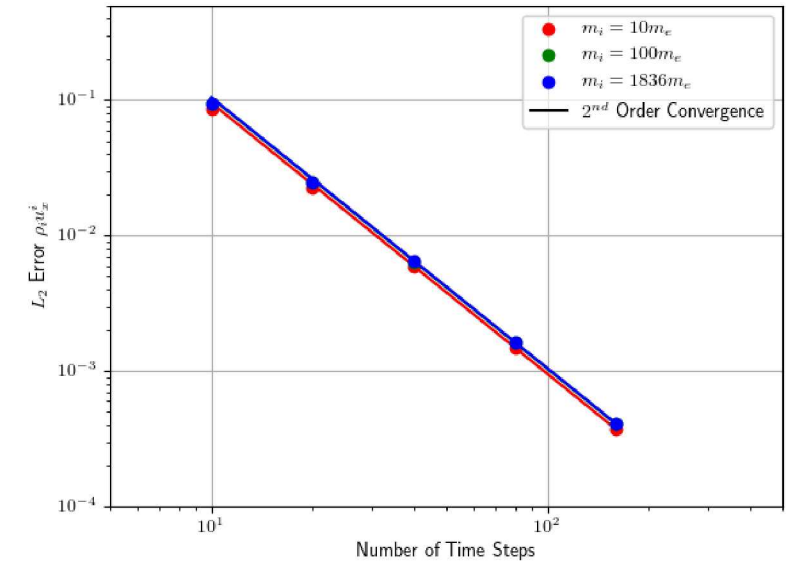
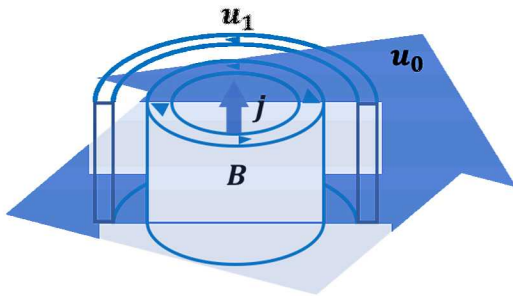
Two fluid plasma vortex in MHD limit

- IMEX time discretization
- Compatible spatial discretization

| | Electrons | Ions |
|---------------------------------|-----------|---|
| $\omega_p \Delta t$ | 23 - 270 | 0.55 - 6.3 |
| $\omega_c \Delta t$ | 1 - 12 | $5.5 \cdot 10^{-4}$ - $6.3 \cdot 10^{-3}$ |
| $v_s \frac{\Delta t}{\Delta x}$ | 0.25 | $5.7 \cdot 10^{-3}$ |
| $c \frac{\Delta t}{\Delta x}$ | | 6.3 |

Achieves 2nd order conv, satisfies involutions:

- $\nabla \cdot E = \rho$ (weakly enforced)
- $\nabla \cdot B = 0$ (strongly enforced)

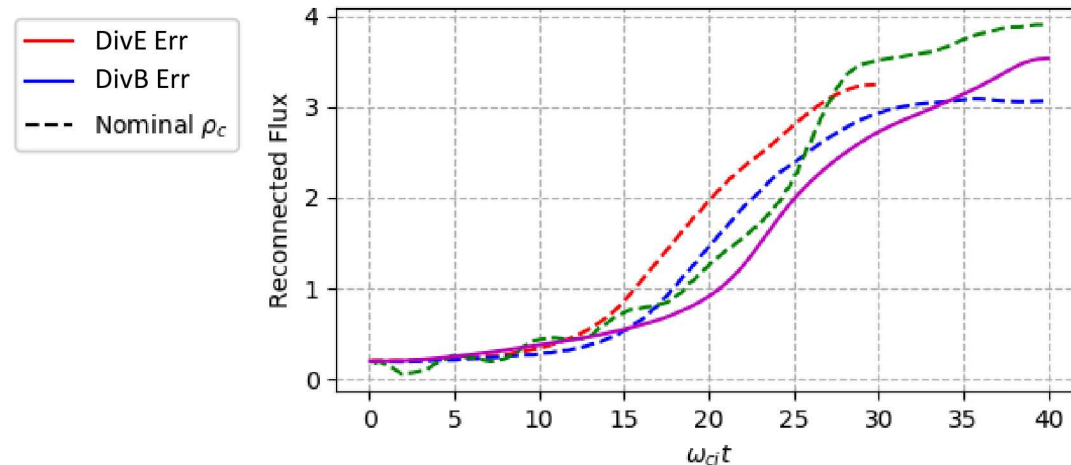
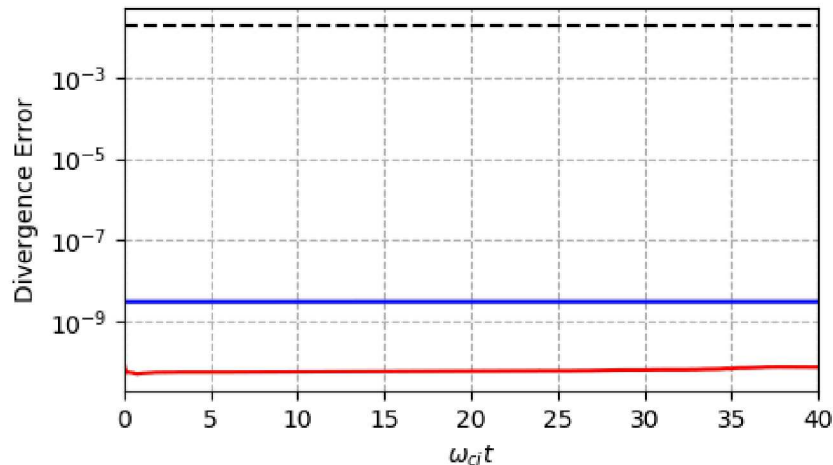
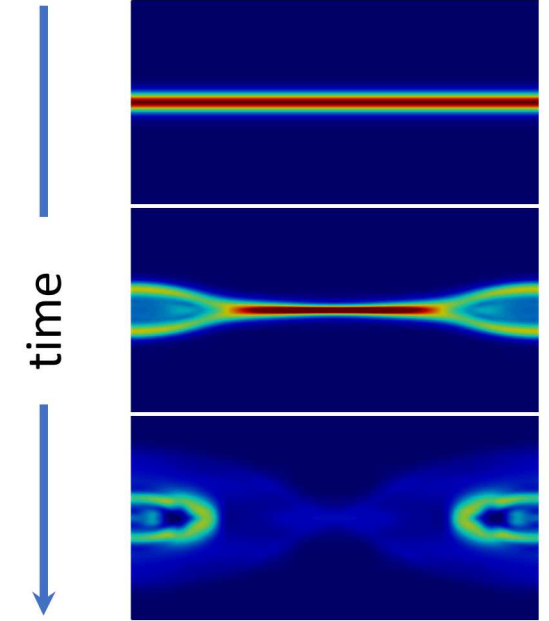


*See Miller, Cyr, Shadid, Kramer, Phillips, Conde, Pawlowski., IMEX and exact sequence discretization of the multi-fluid plasma model, in press JCP, 2019

GEM Challenge Problem (from Drekar)*

Using described spatial and temporal discretizations

- Testing magnetic reconnection using multi-fluid
- As run, unstabilized (recent improvements extend this)
- Qualitative agreement with existing reconnection results
- Preserves no magnetic monopoles and charge density involutions



--- PIC Boltzmann [Pritchett 2001]
--- DG Boltzmann [Reddell 2016]
--- DG Two-Fluid [Srinivasan 2011]
— CG Two-Fluid Exact Sequence

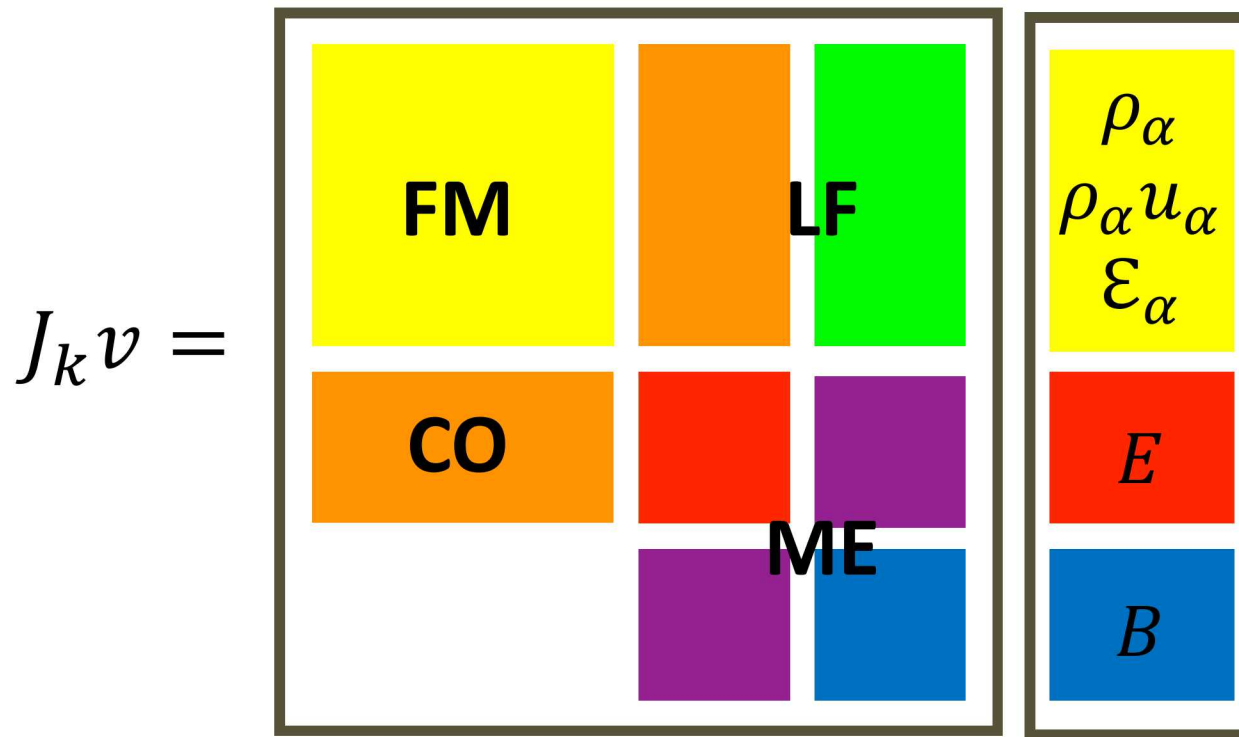
*See Miller, Cyr, Shadid, Kramer, Phillips, Conde, Pawlowski., IMEX and exact sequence discretization of the multi-fluid plasma model, in press JCP, 2019

Nonlinear Algorithm

For IMEX we have to solve a nonlinear problem:

- We are using Newton-Krylov
- To get scalability we must precondition*

$$J_k \Delta x_k = -f(x_k)$$
$$x_k = x_k + \Delta x_k$$



Nonlinear terms

- Fluid matrix (mass like)
- Lorentz Force

Linear terms

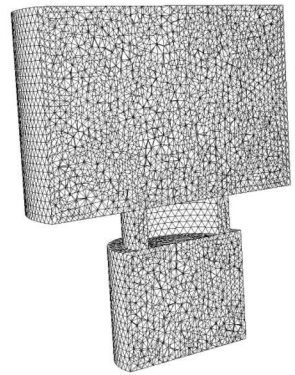
- Maxwell equations
- Current Operator

*E. G. Phillips, J. N. Shadid, E. C. Cyr, S. T. Miller, Enabling Scalable Multi-Fluid Plasma Simulations through Block Preconditioning, Accepted to Lecture Notes in Computational Science and Engineering, 2019.

Nonlinear Algorithm: Maxwell Solver

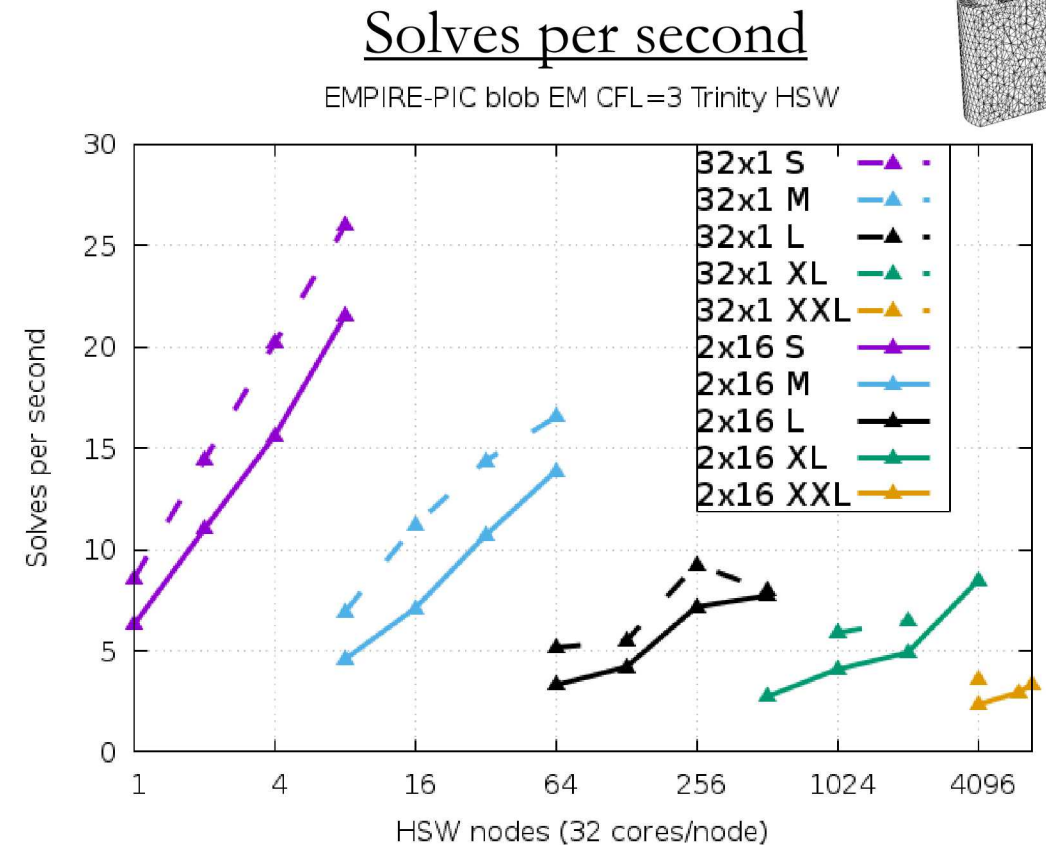
We (Sandia royal*) have made good progress in solving the Maxwell system

- Preconditioner exploits exact-sequence discretization structure



| Size | #Elts | #Nodes | #Edges | #Faces |
|------------|--------|--------|--------|--------|
| S | 337k | 60k | 406k | 683k |
| M | 2.68M | 462k | 3.18M | 5.40M |
| L | 20.7M | 3.51M | 24.4M | 41.6M |
| XL | 166M | 27.9M | 195M | 333M |
| XXL | 1.332B | 223M | 1.56B | 2.67B |

- Trinity HSW: scaling study to **entire HSW** partition (9375 nodes, 300,000 cores)
- Trinity KNL: scaling study to **99.2% KNL** partition (9900 nodes, 633,600 cores)



* **Credit (and apologies) to:** Jonathan Hu, Christian Glusa, Paul Lin, Edward Phillips, Matt Bettencourt, James Elliott, Chris Siefert, Siva Rajamanickam

Examining the IMEX Scheme



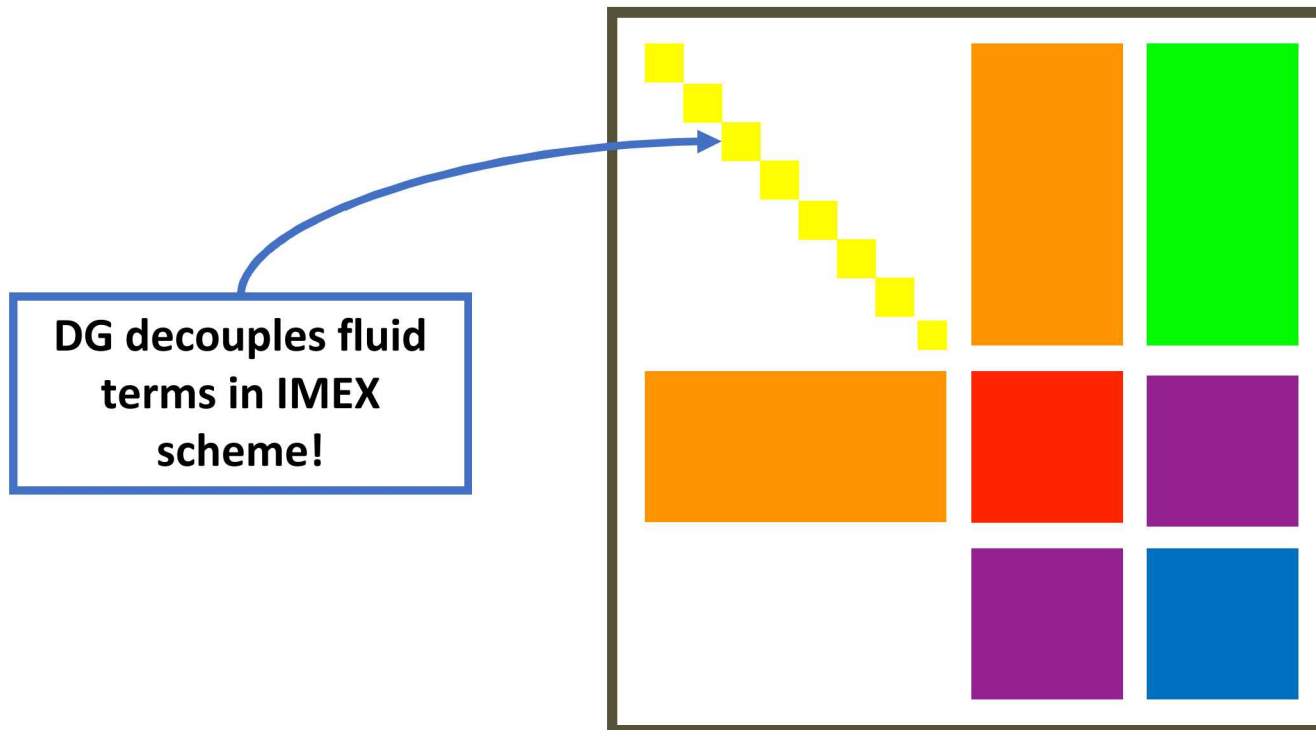
- Fluid matrix is mass matrix (CG fluids gives global coupling)
- Maxwell solver is effective (and should remain unperturbed)
 - Handles speed of light coupling
- Important to get plasma frequency and cyclotron frequency coupling
 - Handled by preconditioning
 - These are local (ODE-like) coupling terms
- Many linear operators that can be computed once and reused

We will try to construct a scheme:

- Take advantage of only local coupling in fluid operators
- Maxwell solver is effective (and should remain unperturbed)
- Handle plasma/cyclotron frequency coupling efficiently
- Reduce the number of recomputations required per nonlinear step

Introduce DG Fluids/CG Maxwell

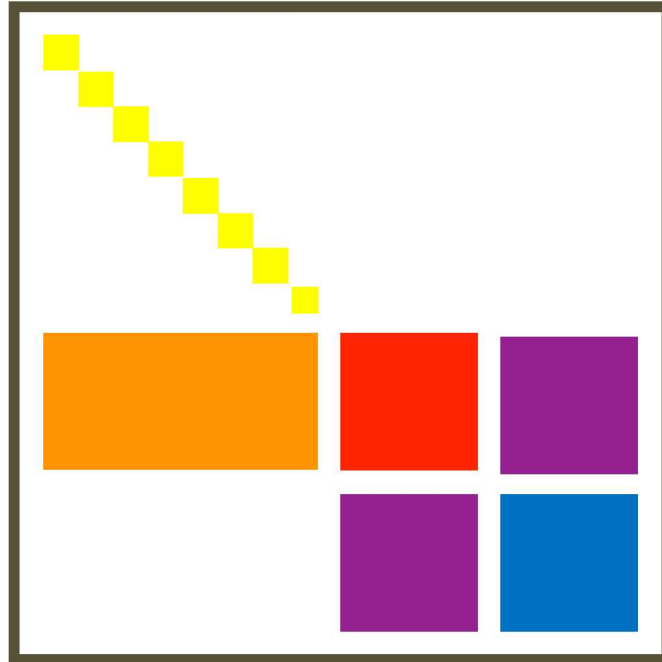
- DG Fluids will make the fluid contribution block diagonal on each element
 - Local nature of DG discretization
 - IMEX splitting choice
- Support for involutions still preserved
 - No Magnetic monopoles is the same
 - Weak enforcement of Gauss' law works (math is more complex)



Quasi-Newton Method

Typically I would do Newton-Krylov, but...

Block lower Gauss-Seidel



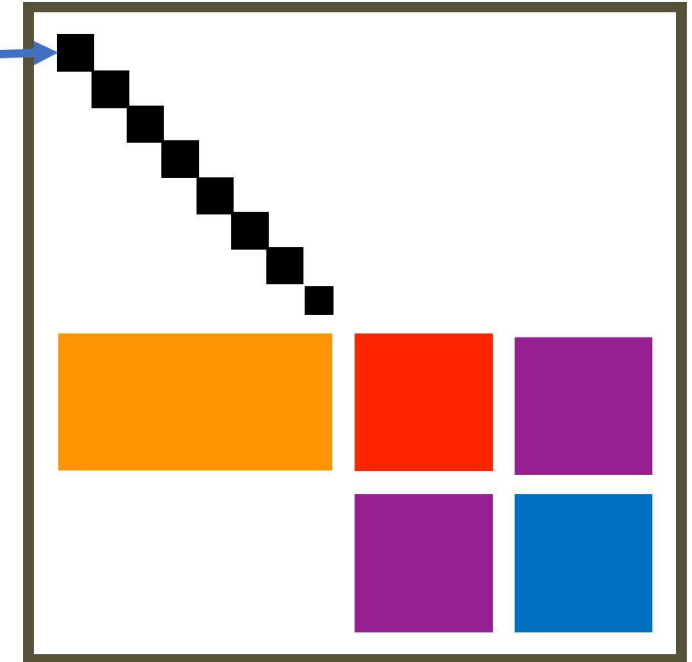
- ✓ Triangular solve
- ✓ Leverages Maxwell solver
- ✓ Block diagonal fluid solve
- ✓ Implicit cyclotron frequency
- ✗ Implicit plasma frequency

Couple in plasma frequency using Schur complement

Both schemes:

- Simplified linear construction
- Only inner Maxwell Krylov solve
- Will require more iterations than Newton
- Maybe cheaper than Newton

Block GS with Schur Complement



- ✓ Triangular solve
- ✓ Leverages Maxwell solver
- ✓ Block diagonal fluid solve
- ✓ Implicit cyclotron frequency
- ✓ Implicit plasma frequency

Plasma Frequency Schur Complement

To step over plasma frequency we must work it into the “black” part of the approximate Jacobian.

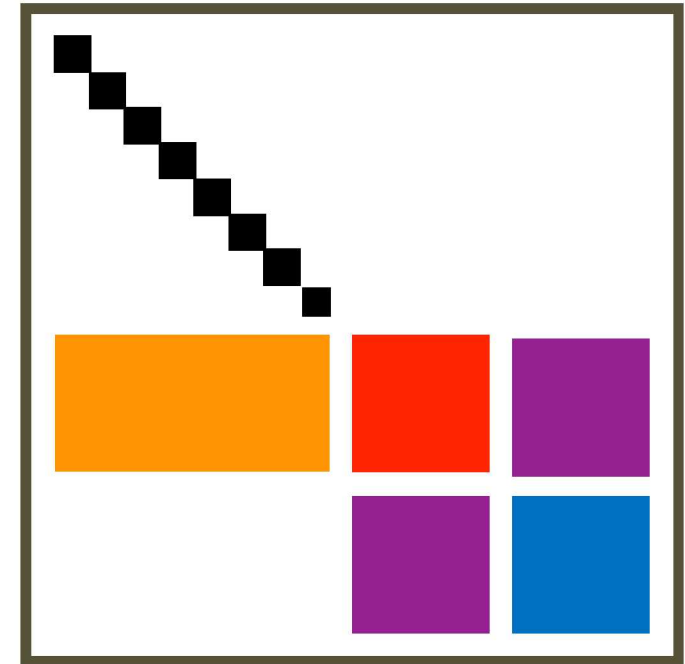
- Mode derived from coupling Ampere’s law and momentum equation

$$\left. \begin{aligned} \frac{\partial(\rho_\alpha \mathbf{u}_\alpha)}{\partial t} &= \frac{q_\alpha}{m_\alpha} \rho_\alpha \mathbf{E} \\ \frac{\partial \mathbf{E}}{\partial t} &= -\frac{1}{\epsilon_0} \frac{q_\alpha}{m_\alpha} (\rho_\alpha \mathbf{u}_\alpha) \end{aligned} \right\} \frac{\partial^2(\rho_\alpha \mathbf{u}_\alpha)}{\partial t^2} = -\frac{1}{\epsilon_0} \frac{q_\alpha^2}{m_\alpha^2} \rho_\alpha (\rho_\alpha \mathbf{u}_\alpha)$$

- We apply the local Schur complement to the fluid contribution as a correction

$$\frac{(\rho_\alpha \mathbf{u}_\alpha)}{\Delta t} + \underbrace{\Delta t \frac{1}{\epsilon_0} \frac{q_\alpha^2}{m_\alpha^2} \rho_\alpha (\rho_\alpha \mathbf{u}_\alpha)}_{\text{Correction}}$$

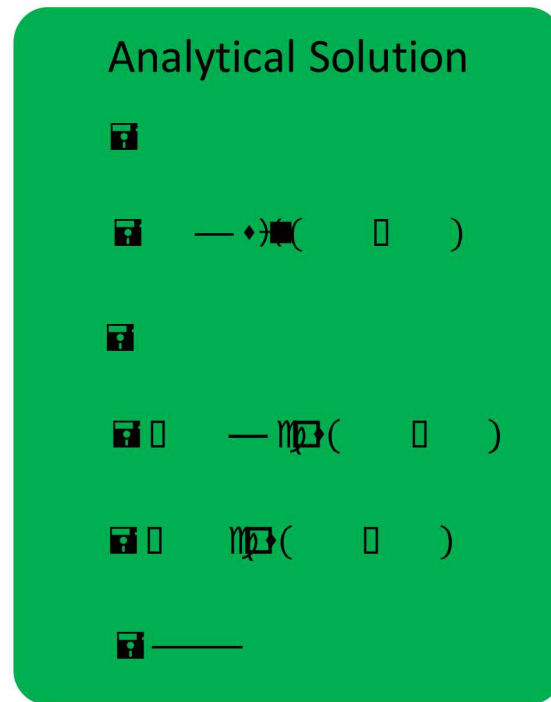
Block GS with Schur Complement



O-Wave Convergence Results (EMPIRE-Fluid)

A linear wave verification test*

- Refining in space and time
- Running IMEX SSPRK2



* S. Miller, J. Niederhaus, R.M.J. Kramer, and G. Radtke, Robust Verification of the Multi-Fluid Plasma Model in Drekar.. United States: N. p., 2017. Web.

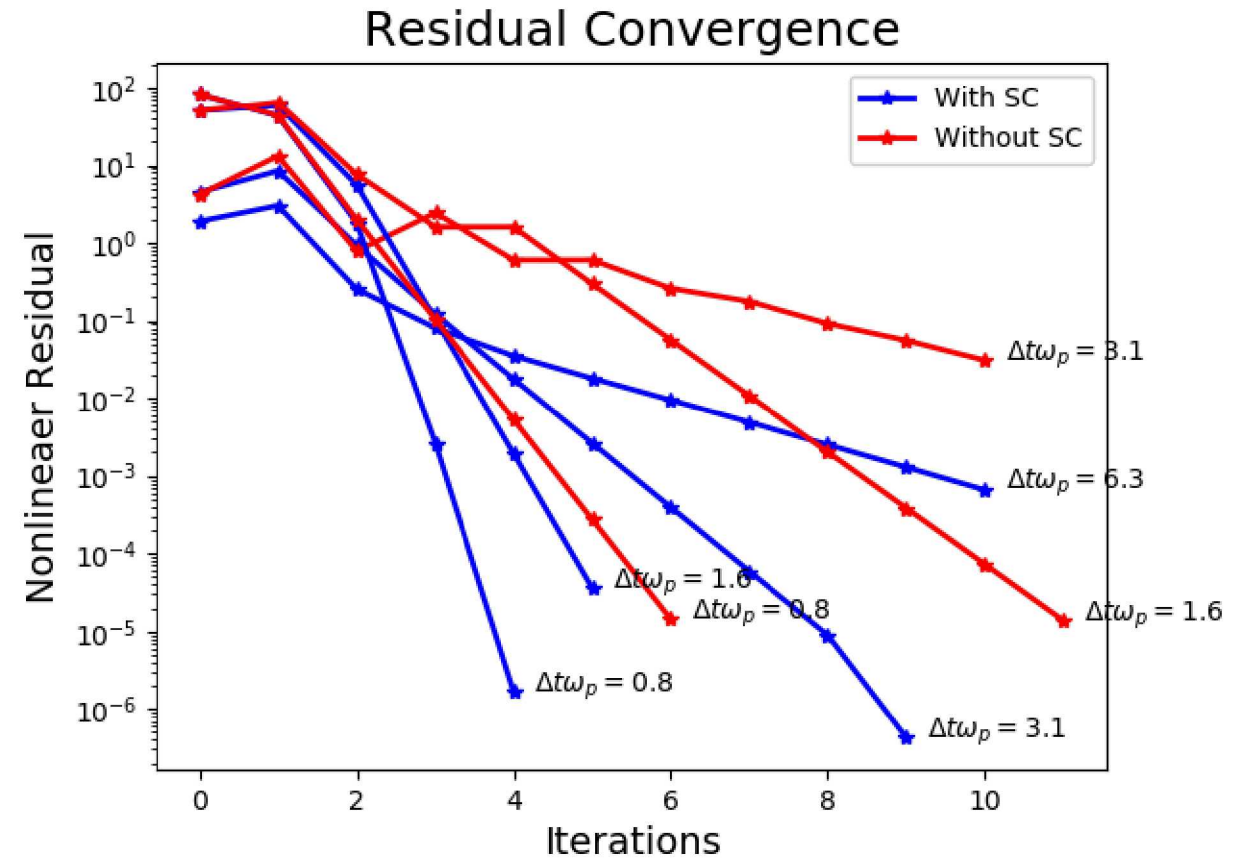
O-Wave Nonlinear Solver (EMPIRE-Fluid)

Adding Schur complement improves nonlinear convergence

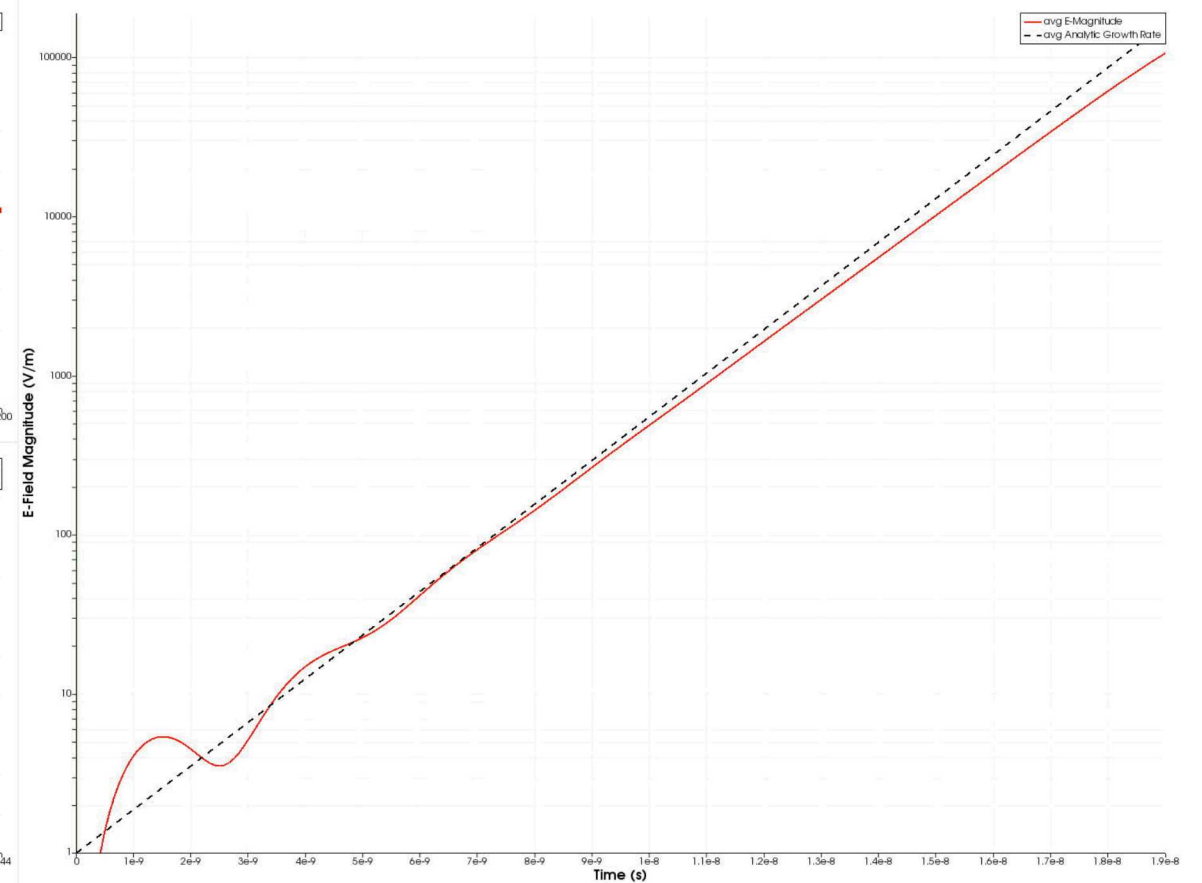
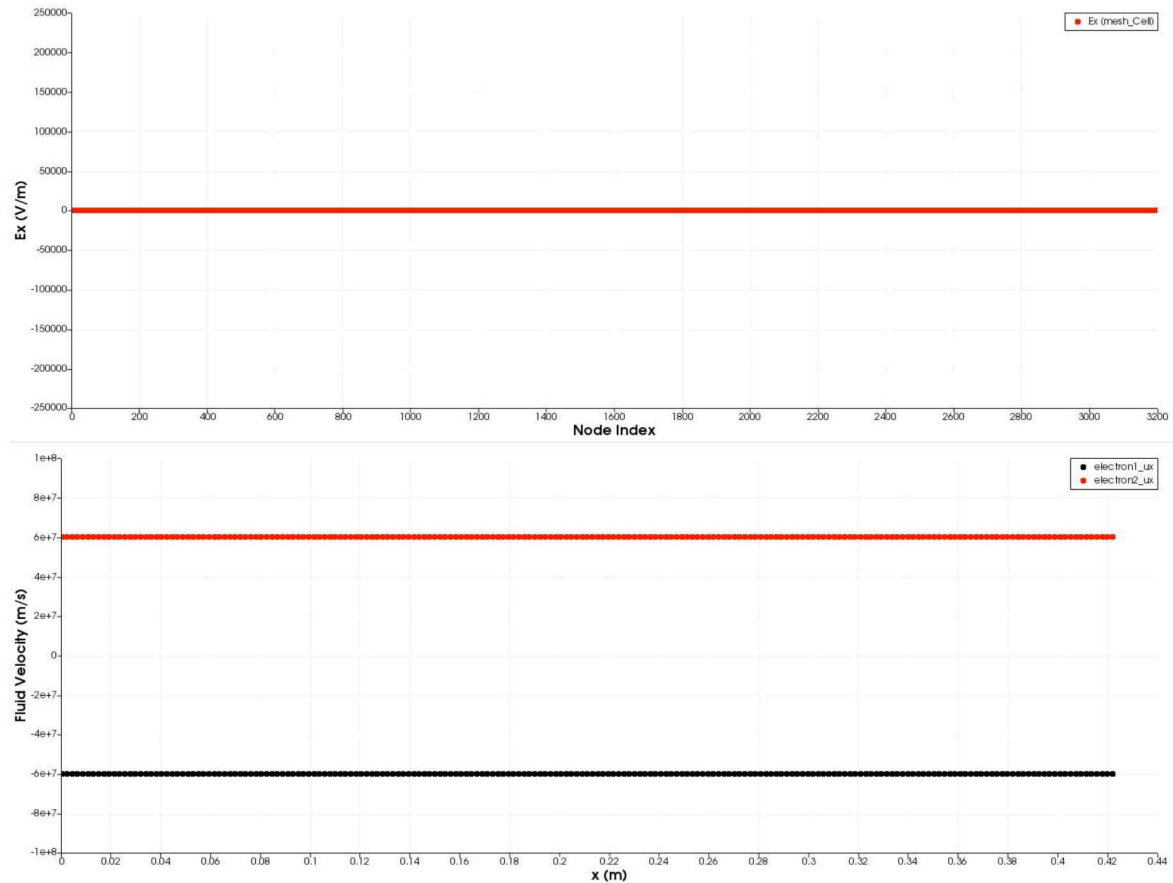
- Still has strong growth in iteration count with increasing time steps
- Cost/benefit tradeoff study against Newton-Krylov with similar preconditioner required

Each nonlinear iteration requires:

1. Reconstruction of fluid Jacobian inverse
2. Solve of Maxwell system



Two Stream Instability*



* We run this only through the linear growth regime

Final Thoughts

In Review:

- Mixture of IMEX temporal and Exact-Sequence discretizations has been shown to enforce involutions
- Extending previous work to DG fluid/CG Maxwell discretization
- Developed quasi-Newton solver for DG/CG discretization
 - Takes advantage of IMEX, DG structure, and linearity
 - Schur complement correction required for improved performance
- Early results for DG/CG discretizations are encouraging

Open questions:

1. How does quasi-Newton compare with Newton-Krylov?
2. How does quasi-Newton converge in other scenarios? (e.g. scalability)
3. How does stabilization interact with Exact-Sequence/IMEX discretization?