

# Spin Qubits Using Holes in Strained Germanium Quantum Well Heterostructures

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August 12, 2019

This work was performed, in part, at the Center for Integrated Nanotechnologies, a U.S. DOE, Office of Basic Energy Sciences, user facility. Sandia National Laboratories is a multi-mission laboratory managed and operated by National Technology and Engineering Solutions of Sandia, LLC., a wholly owned subsidiary of Honeywell International, Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA-0003525

# Outline

- Project Overview
- New Device Design
- Fabrication Challenges
- Charge Sensing
- Device Level Modeling for Holes in Ge
- Outlook

# Project Overview

Goal: Demonstrate and study hole spin qubits in strained Ge/SiGe Heterostructures

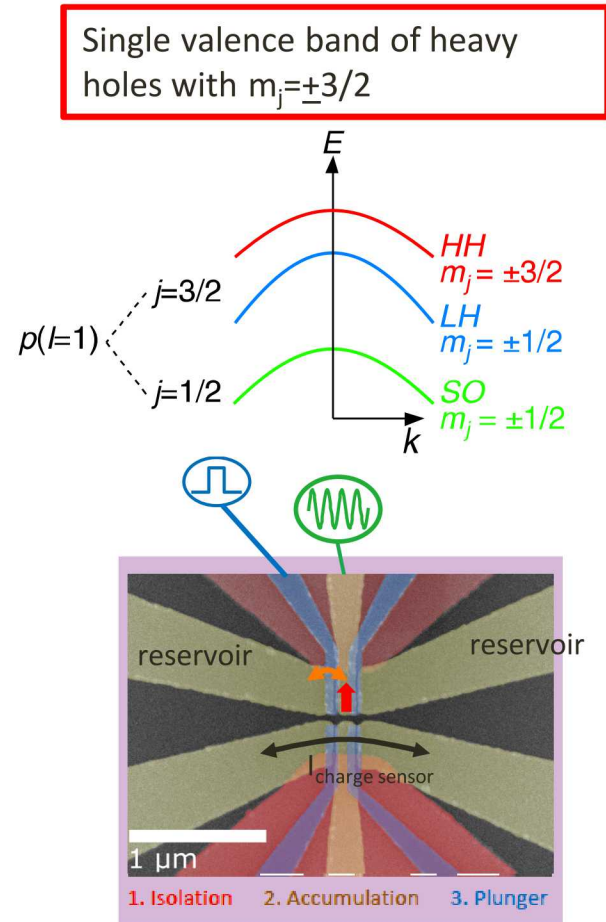
Motivation: Holes in Ge/SiGe provide a compelling alternate approach to spin based qubits. They maintain many of the advantages of silicon without valley splitting.

Team:  
Dwight Luhman (PI)  
Tzu-Ming Lu  
Will Hardy (now staff)  
Mitchell Brickson

Near Term Goals:  
Single Spin Readout  
Spin Rotations  
Better understanding of SOC in Ge dots

# Basic Idea

- Single Hole confined to lateral quantum dot in Ge heterostructure
- Spin Qubit States:  $m_j = \pm 3/2$
- Qubit readout and initialization through energy selective tunneling to reservoir
- Qubit Control through microwaves applied to gate
- Occupancy detected through nearby charge sensor
- Strong spin orbit coupling (SOC) is important. Form is  $\sim k^3$  in 2DHG.

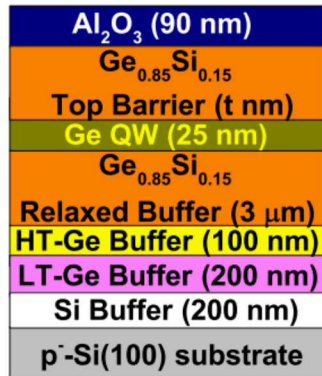


# Ge/SiGe quantum wells

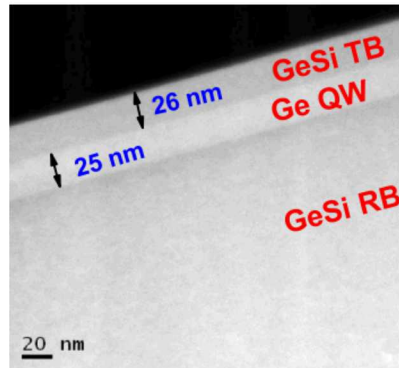
Undoped Heterostructures  
Clean epitaxial interface

Small effective mass,  $m^* \sim 0.08 m_0$

(a)

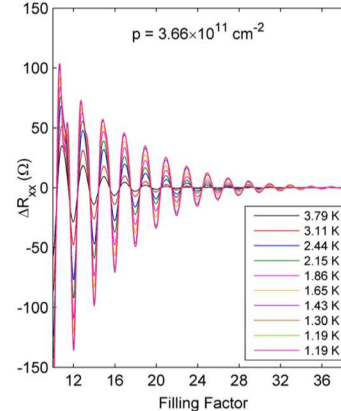


(b)

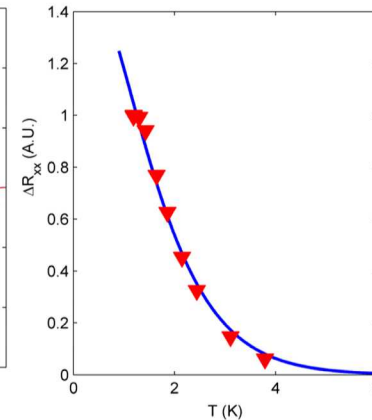


Su, *et al. Phys. Rev. Mater.* (2017)

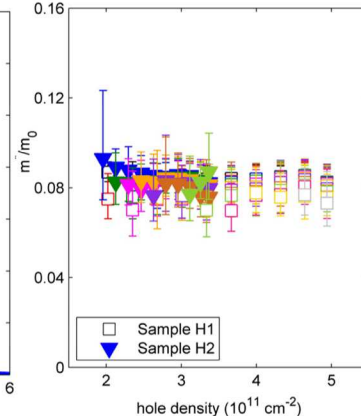
a



b



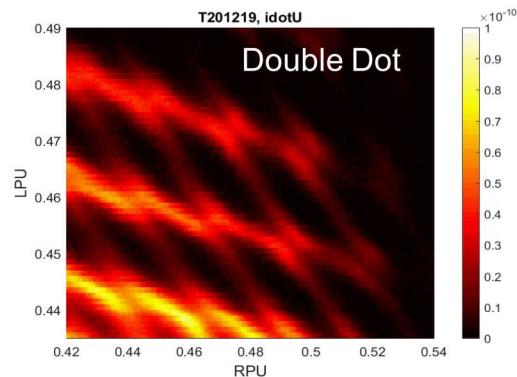
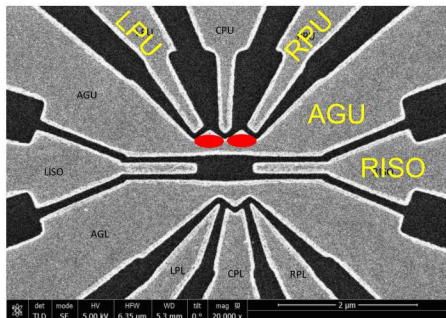
c



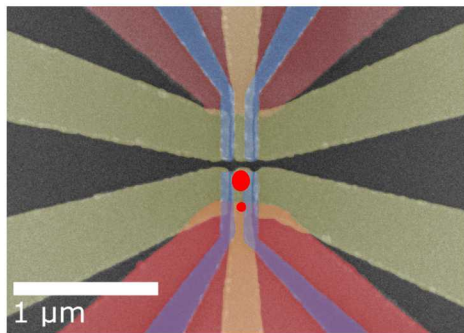
Hardy *et al., Nanotechnology* (2019)

# Summary of Previous Results

## Single Layer Design:

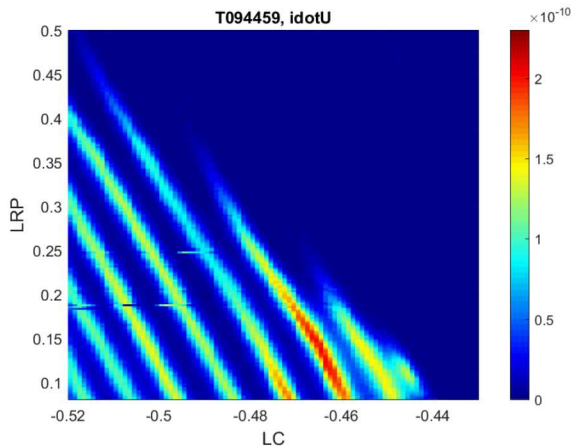


## Three-Metal Design



Good Coulomb Blockade

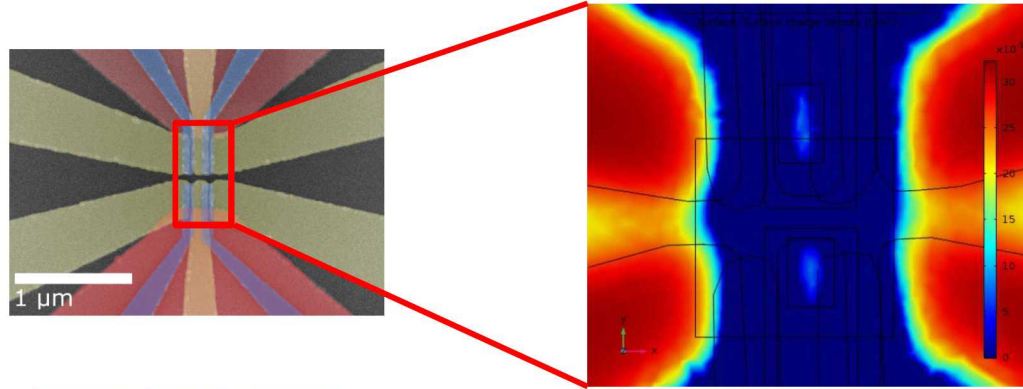
Upper and lower dots  
merge when both  
activated



Large aspect ratio of dot  
design causes multiple  
dots to be formed under  
a single gate

# Device Redesign

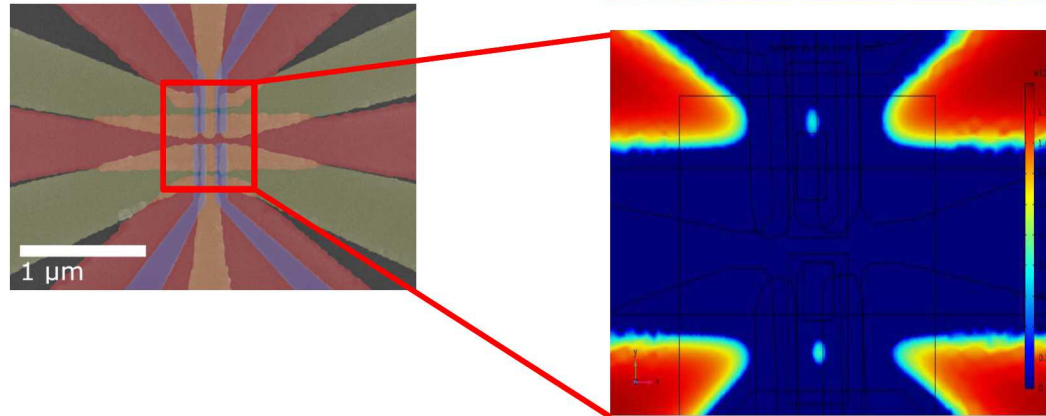
Original Design



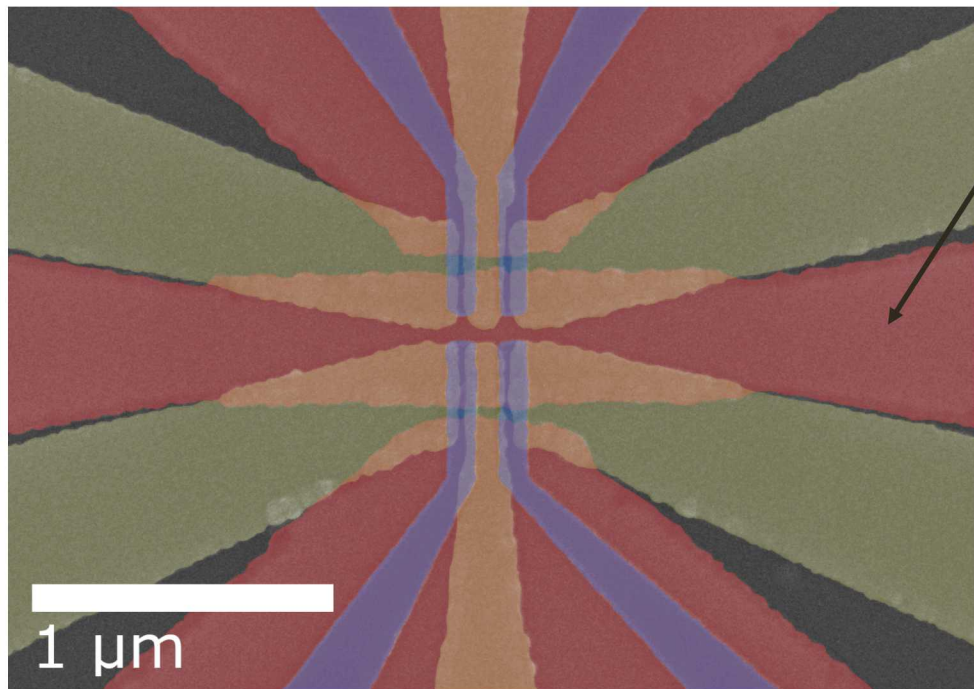
COMSOL calculations  
by M. Brickson

Central gate  
isolates 2DHGs  
between upper  
and lower dots  
and creates  
better  
confinement for  
each dot

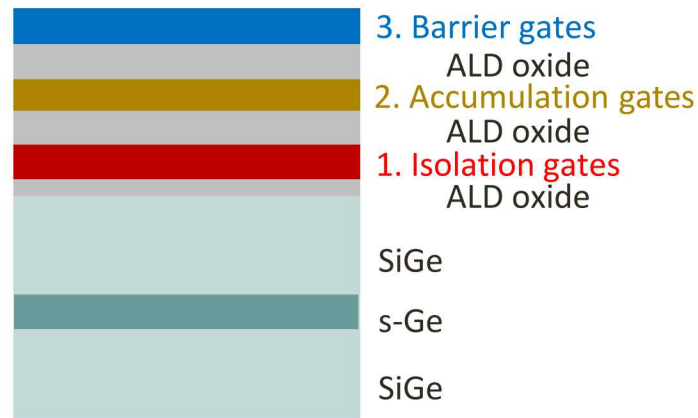
New Design



# New 3-layer device design



Includes horizontal  
isolation gate



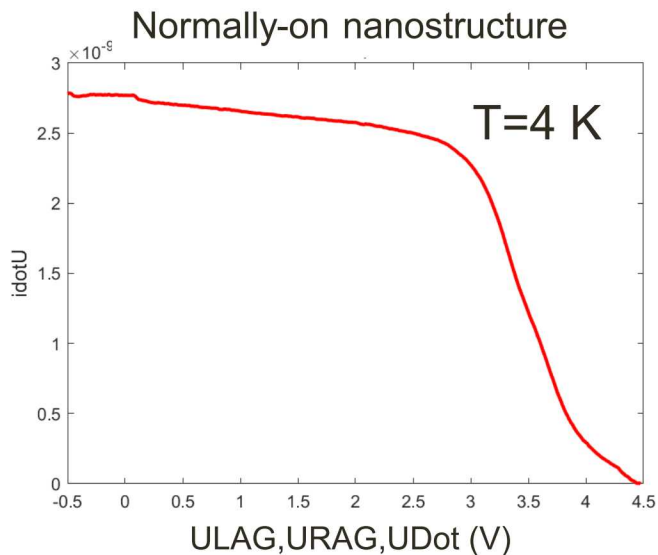
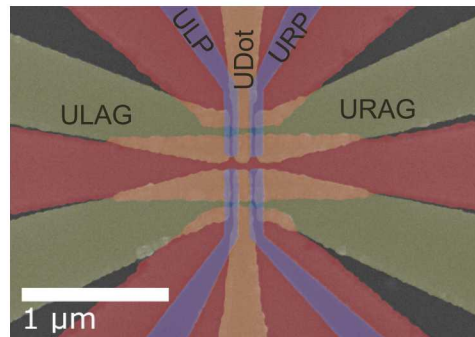
ALD Oxide = 24 nm  $\text{Al}_2\text{O}_3$  + 1 nm  $\text{HfO}_2$   
Metal = 2 nm Ti + 18 nm Pt

# Device Challenges

Lithography had good yield

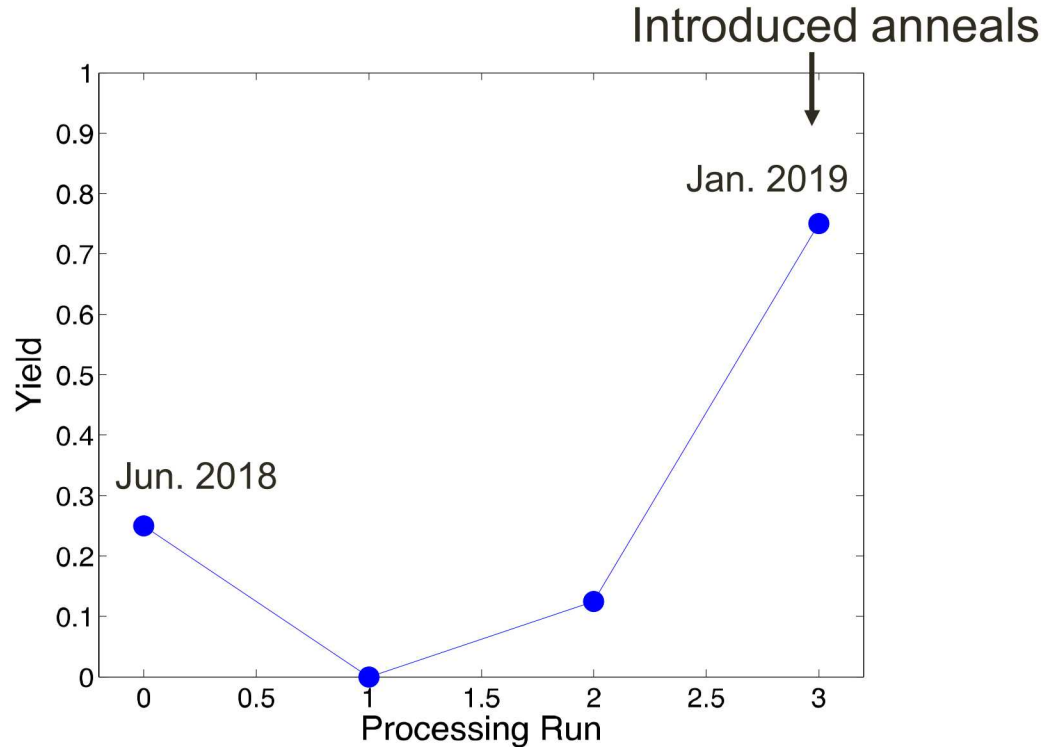
First batch of new devices had low yield

- Gate leakage
- Normally on
- Unstable
- No usable nanostructures: 0/8



Something changed between fab runs

# Improved Device Yield



- We experienced an abrupt, unexpected fabrication issue in August 2018
- The issue seems to be related to contamination of  $\text{TiCl}_4$  in the ALD machine
- A combination of cleaning the tool and forming gas anneals ( $400^\circ\text{C}$  for 30 minutes) after every oxide layer improved the issue

# New ALD Tool

## Longer term solution

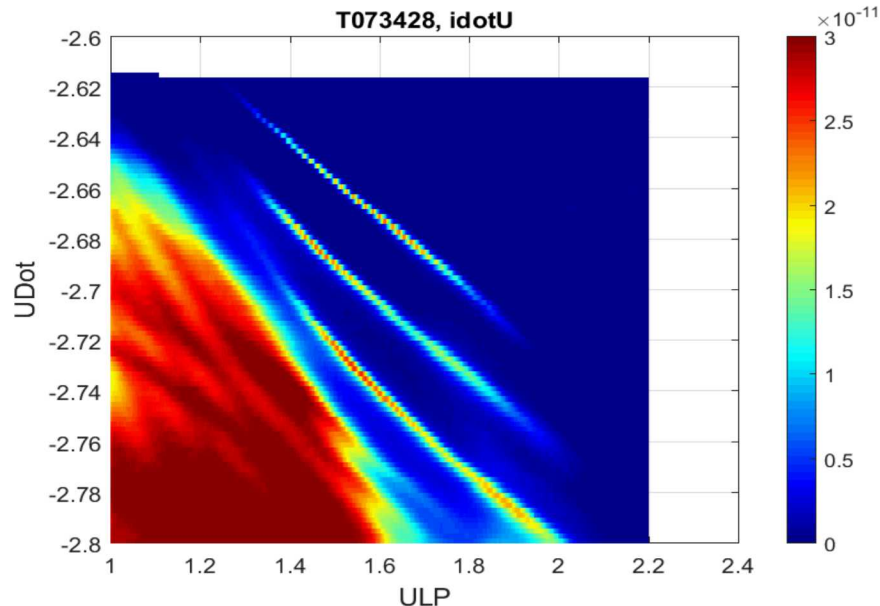
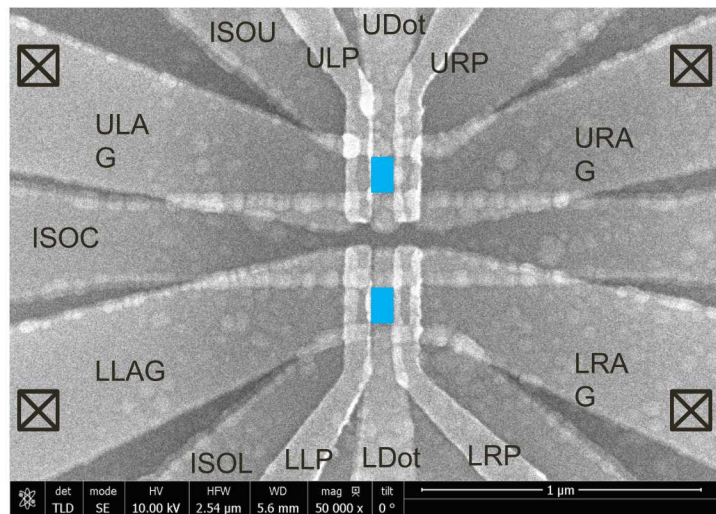


New dedicated ALD tool for Si-based quantum devices

- Veeco Savannah G2
- $\text{Al}_2\text{O}_3$  and  $\text{HfO}_2$  available
- Controlled access to maintain cleanliness
- Status
  - Tool installed and qualified
    - 0.93 Å/cycle
    - 1.6% variation across 6" wafer
  - Awaiting final administrative approval (expected soon)

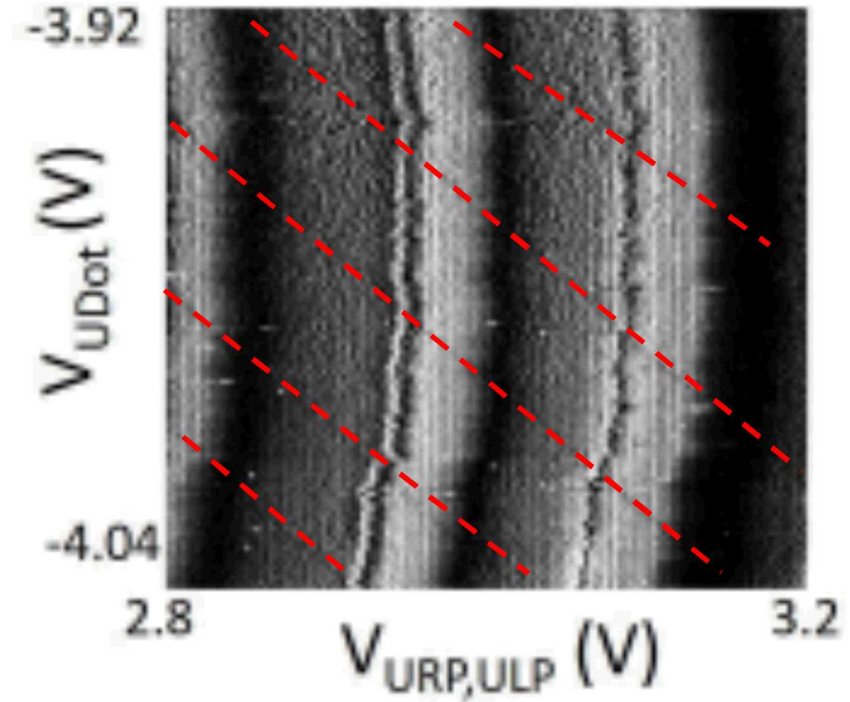
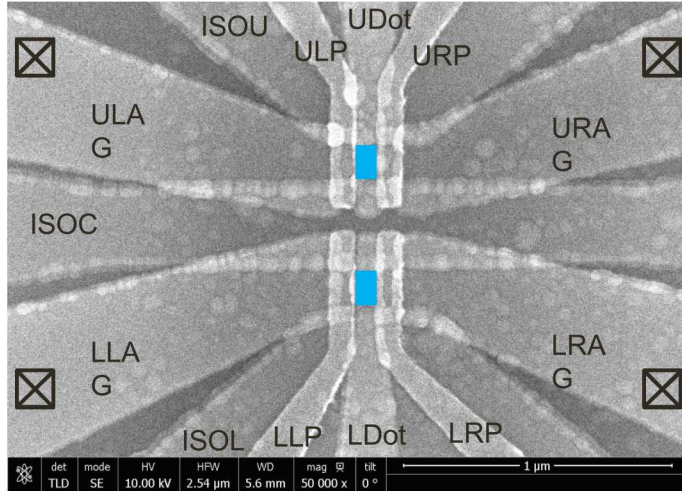
# Successful Three Layer Devices

Quantum dot size is  $\sim 130 \times 80$  nm



Regular Coulomb blockade observed in multiple devices

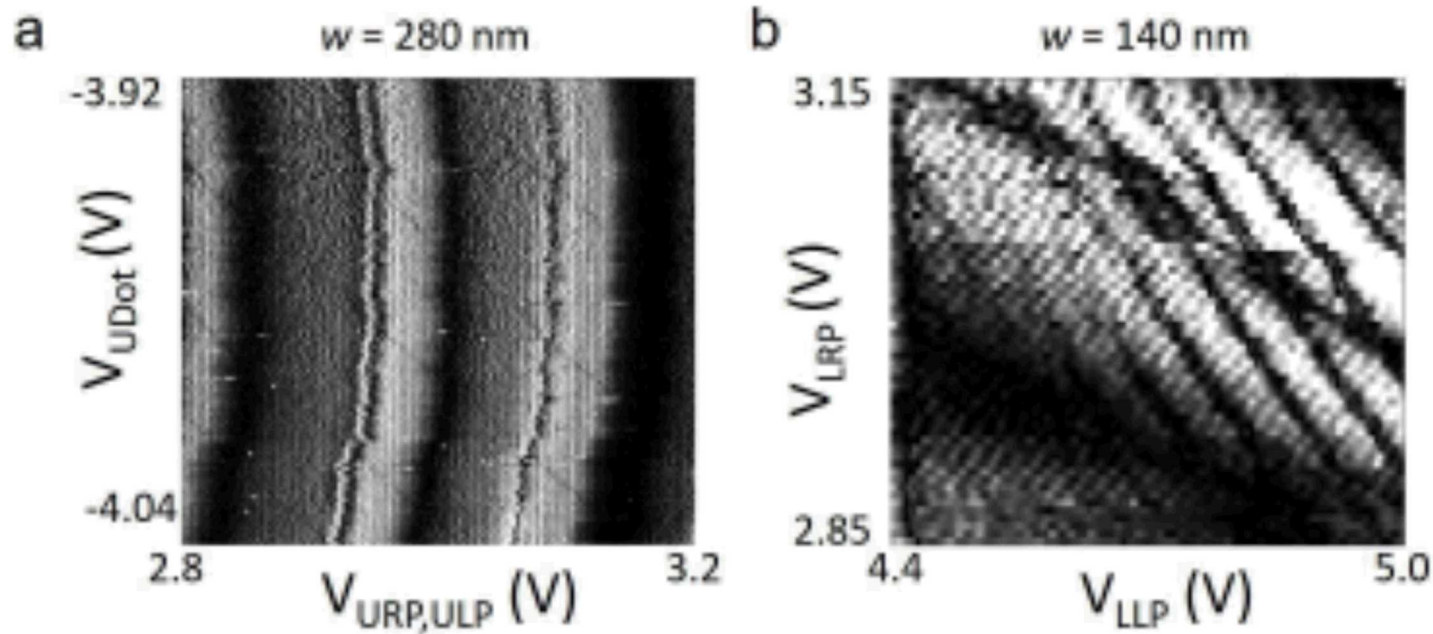
# Charge Sensing



- New design allows charge sensing
- Charge sensed signal is very weak

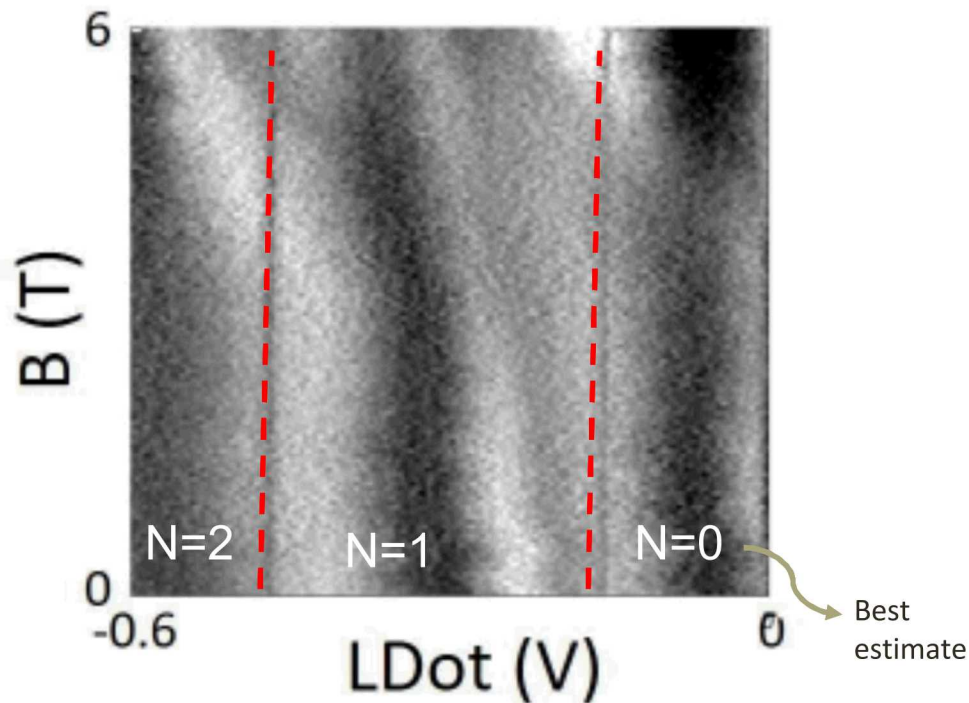
# Improved Charge Sensing

- Reduce center gate from  $w=280$  nm to 140 nm in width
- 150% increase in charge sensed signal



# Magnetospectroscopy

- Tune to the low electron regime
- Observe charge transition lines versus magnetic field
- Slope of the lines gives an estimate of the g-factor (field in-plane)
- $g_{||} \sim 0.3$  (rough estimate), consistent with Delft result\* for a quantum dot



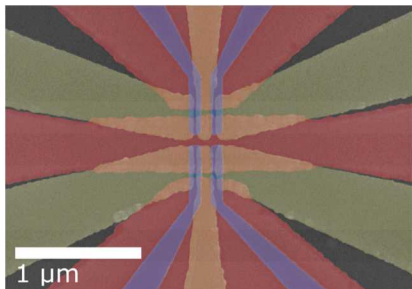
# Modeling SOC in Ge QDs

Overall Goal: Develop accurate device-level models of SOC coupling in quantum dots to guide & interpret qubit experiments

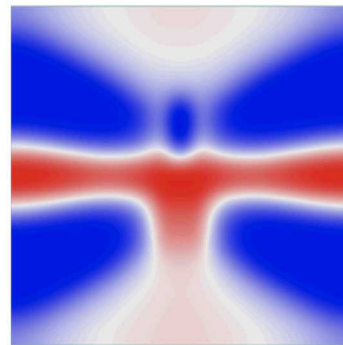
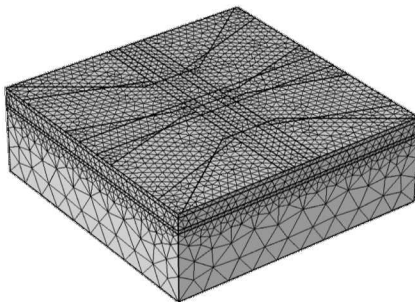
Near Term Goal: Understand the magnetic field angular dependence of the Rabi frequency in Ge quantum dots

# Modeling Sequence

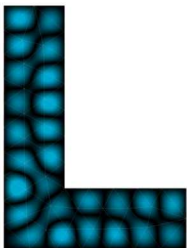
Real Device



Electrostatic Potential in COMSOL



Laconic (Andrew Baczewski)



- Mesh based solver for multi-scale quantum models
- Locally enriched to capture difficult physics



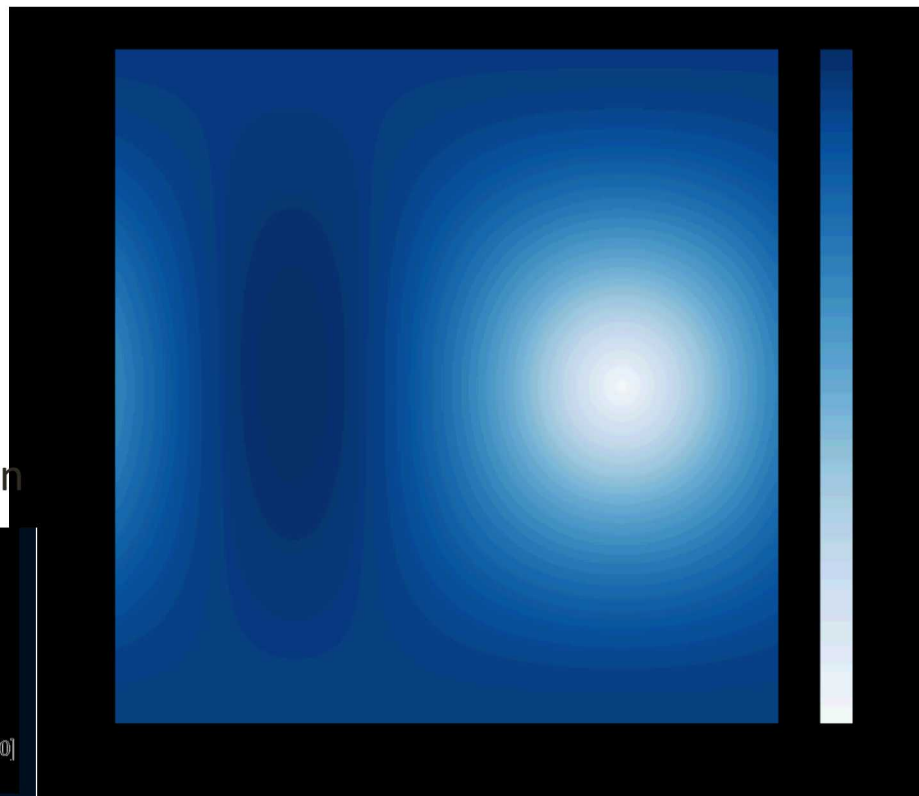
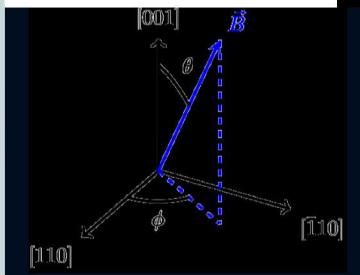
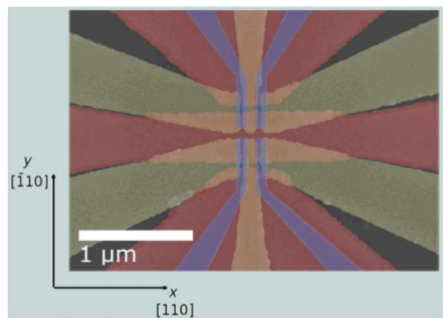
# Single Band Model

Simple angular dependence

Capabilities added to Laconic

- Magnetic Fields (vector potential)
- Zeeman Splitting
- Momentum Operator
- Spin Orbit Coupling

$V_{AC}$  along the  $y$   $[\bar{1}10]$  direction



# Multi Band Model

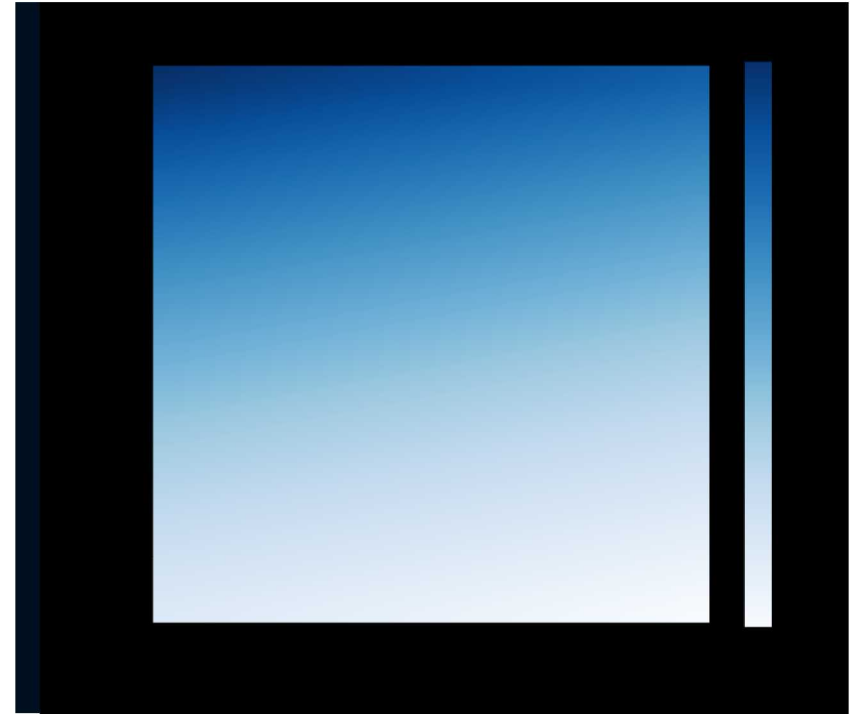
- Luttinger Hamiltonian to describe light and heavy hole mixing
- Expand in effective mass envelope function basis
- Ground and excited states are a combination of heavy hole and light hole terms
- Complex behavior from the Luttinger Hamiltonian
  - Natural spin quantization along confinement direction competes with Zeeman splitting
  - Near degeneracy in heavy hole term near  $\theta=\pi/2$  causes a spike in frequency

Predicted Rabi frequency for our design



# Dependence on Dot Shape

- Increasing the size of the quantum dot in the y direction leads to a larger Rabi frequency due to a larger dipole
- Increasing the size of the quantum dot in the x direction leads to a modestly smaller Rabi frequency due to a decrease in extent of k-space



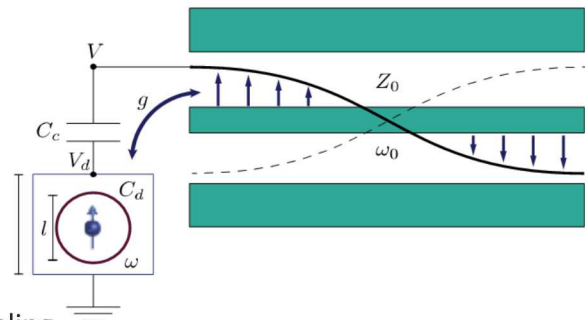
# Spin photon interaction

Understand the spin-photon interaction in Ge holes in low photon limit

- Perpendicular magnetic field
- Cubic SOC as perturbation
- Quantized photon Interaction

Charge Photon Coupling

$$g_0 \equiv \frac{elv}{2s} \omega_0 \sqrt{\frac{Z_0}{\pi \hbar}}$$



$g=20$   
 $f=4.5$  GHz  
 $B=5$  mT

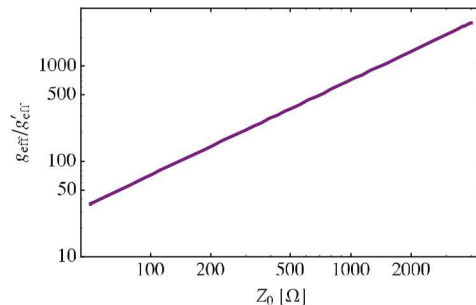
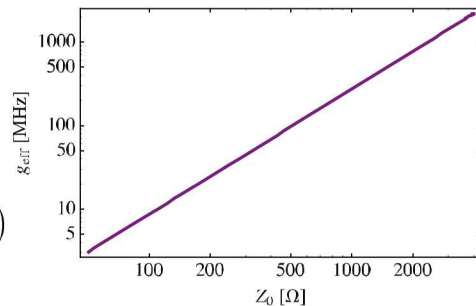
$$H_{\text{SO}} = i\hbar\omega_{\text{SO}}l^3 (k_-^3\sigma_+ - k_+^3\sigma_-) \quad H'_{\text{SO}} = -i\hbar\omega'_{\text{SO}}l^3 (k_+k_-k_+\sigma_+ - k_-k_+k_-\sigma_-)$$

$$g_{\text{eff}} = \left| \langle \widetilde{0, 0, \uparrow}; 3 | H_{\text{int, eff}} | \widetilde{0, 0, \downarrow}; 0 \rangle \right|$$

$$\approx \frac{3\sqrt{6}g_0^3\eta}{4\Omega^2} \left| \tilde{\Lambda}_+ + \tilde{\Lambda}_- \right|$$

$$g'_{\text{eff}} = \left| \langle \widetilde{0, 0, \uparrow}; 1 | H_{\text{int}} | \widetilde{0, 0, \downarrow}; 0 \rangle \right|$$

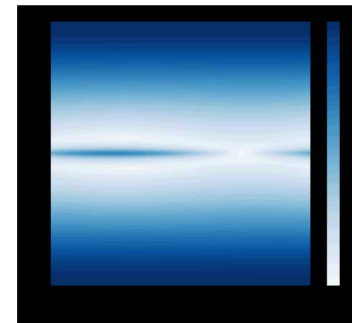
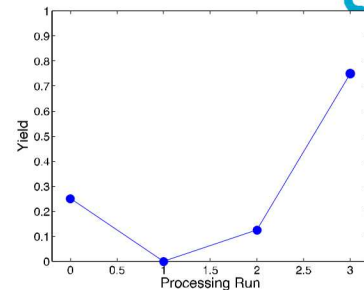
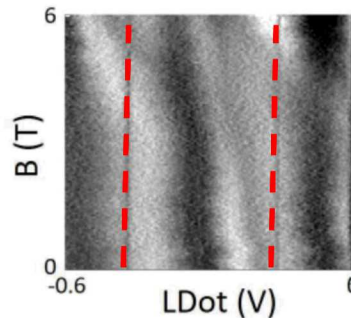
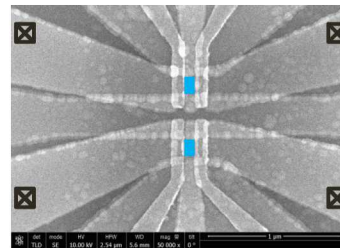
$$\approx 4g_0\eta'\tilde{\Lambda}'$$



Complementary internally funded work

# Progress and Outlook

- Improved Device Design
- Overcame major device fabrication hurdles
- Successful Quantum Dots
  - Charge Sensing
  - Single Hole Occupancy
  - Magnetospectroscopy
- Modeling
  - Multiband model for realistic quantum dots



In Progress

Future

Future

✓  
Demonstrate  
Quantum Dot

✓  
Single Hole  
Occupation

In Progress  
Spin Readout  
& Initialization

Future  
Qubit Control  
(EDSR)

Future  
Qubit  
Characterization

# Acknowledgements

## Modeling:

- Leon Maurer
- Toby Jacobson
- Andrew Baczewski

## Center for Integrated Nanotechnologies:

- Tom Harris
- Mike Lilly

## National Taiwan University:

- Jiun-Yun Li
- C.-Y. Liu
- C.-T. Chou



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# Backup Slides

# Venitucci and Niquet model

- Particle in a box with a static E-field along y-direction
- Calculate Rabi frequency as a function of B-field direction

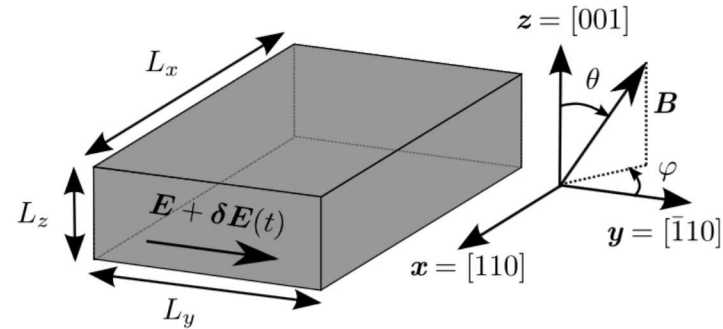
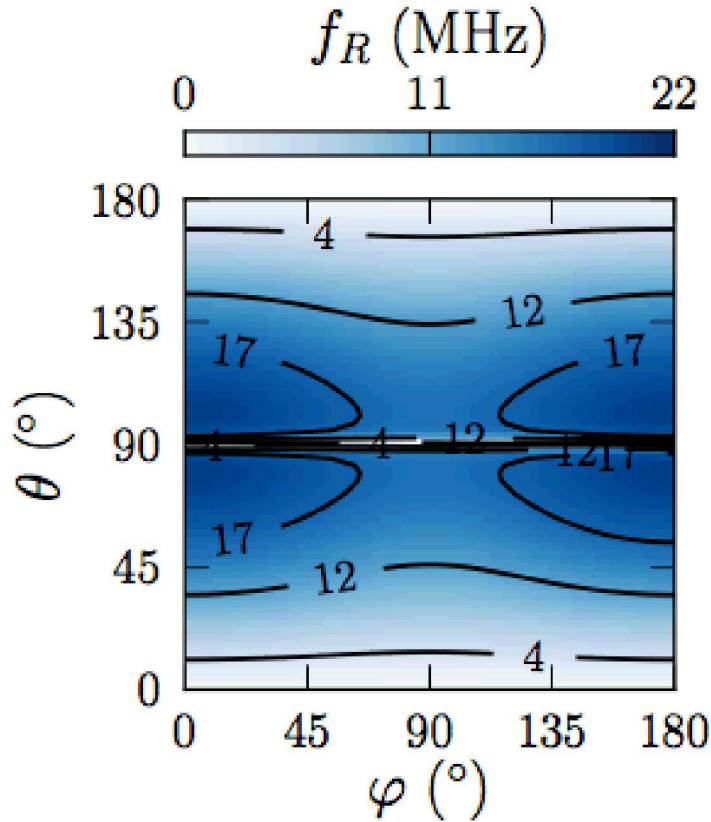


FIG. 1. The model system. A rectangular box with sides  $L_x$ ,  $L_y$  and  $L_z$  is subjected to a static magnetic field  $\mathbf{B}$ , a static electric field  $\mathbf{E} = E_0 \mathbf{y}$  and a radio-frequency electric field modulation  $\delta \mathbf{E}(t) = E_{ac} \sin(2\pi f_L t + \phi) \mathbf{y}$ . The orientation of  $\mathbf{B}$  is characterized by the polar angle  $\theta$  and the azimuthal angle  $\varphi$ .

Venitucci and Niquet, arXiv:1901.09563

# Venitucci and Niquet results



Approximate eigenstates of the Luttinger Hamiltonian:

$$|g\rangle \approx \alpha_g | +3/2 \rangle + \beta_g | -1/2 \rangle,$$

$$|e\rangle \approx \alpha_e | -3/2 \rangle + \beta_e | +1/2 \rangle.$$

The dipole drive is diagonal in spin. Thus if we have a magnetic field along  $z$ , which also doesn't further mix spins, we have no Rabi oscillations.

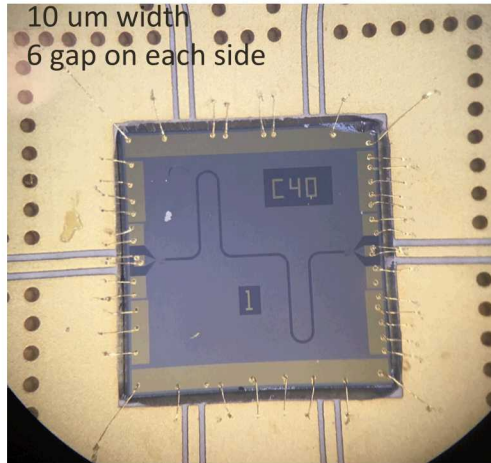
The Rabi frequency increases mostly as the magnetic field drifts away from  $z$ , until it comes to the  $xy$ -plane. With  $B_z = 0$ , eigenstates are like

$$|g\rangle = \frac{1}{\sqrt{2}} (| +3/2 \rangle + e^{i\alpha} | -3/2 \rangle)$$

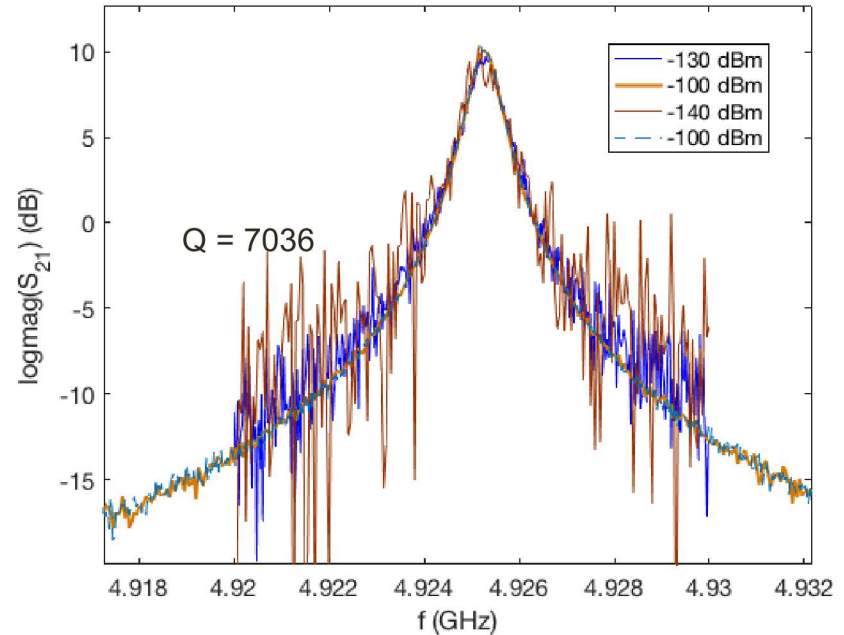
with too large of a level spacing between LH and HH to mix better

# Nb 50 Ohm CPW sample from Star Cryo

Metal patterned on 25 nm  $\text{Al}_2\text{O}_3$ , similar to what we expect to use for quantum dots



6 pF coupling capacitors



- High quality superconducting films can be obtained from Star Cryo
- High Q resonators are possible on our quantum dot oxide

# NbTiN summary

T=4 K



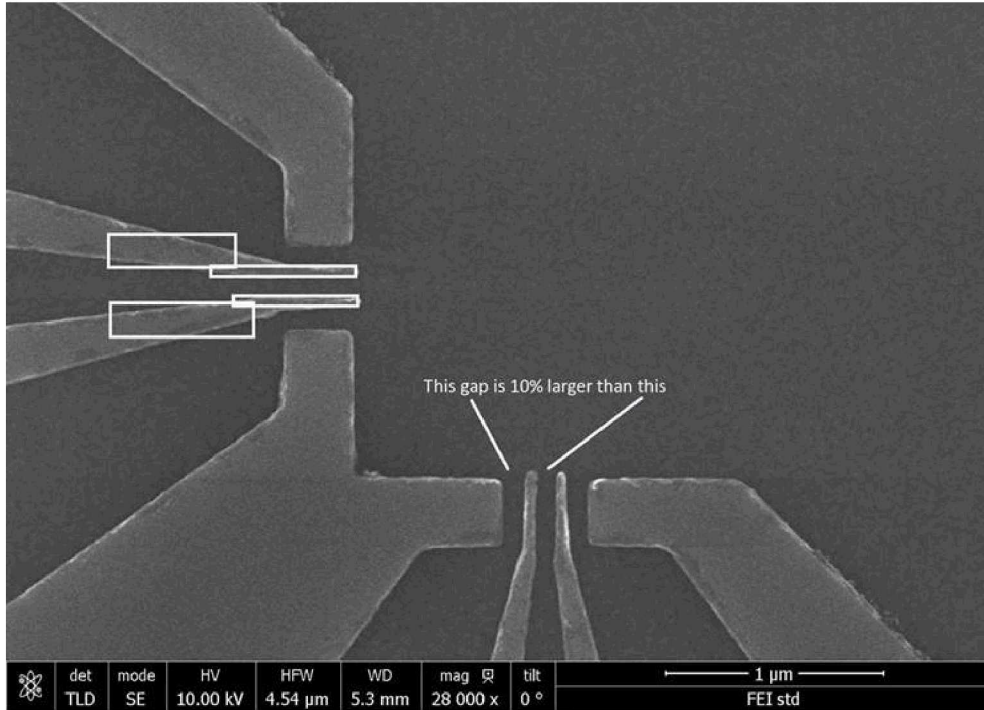
From AppCad

Sample	fp1	fwhm	Q	Z non SC	F design	fm/fd	Z m
5 um line	4.699 GHz	1.76 MHz	2670	94.3 Ohms	5.321 GHz	.883	105
3 um line	3.5428 GHz	1.26 MHz	2806	106.6	5.400	.656	163
2 um line	1.963 GHz	281 KHz	6981	117.4	5.506	.3565	322
1 um line	not measured yet			139.0	5.8		

$$\left(\frac{Z_g}{Z_m}\right)^2 = \frac{L_g}{L_g + L_k}$$
$$= 1 - \alpha$$

$$\alpha = \frac{L_k}{L_k + L_g}$$

$$\alpha = 1 - \left(\frac{f_m}{f_g}\right)^2$$



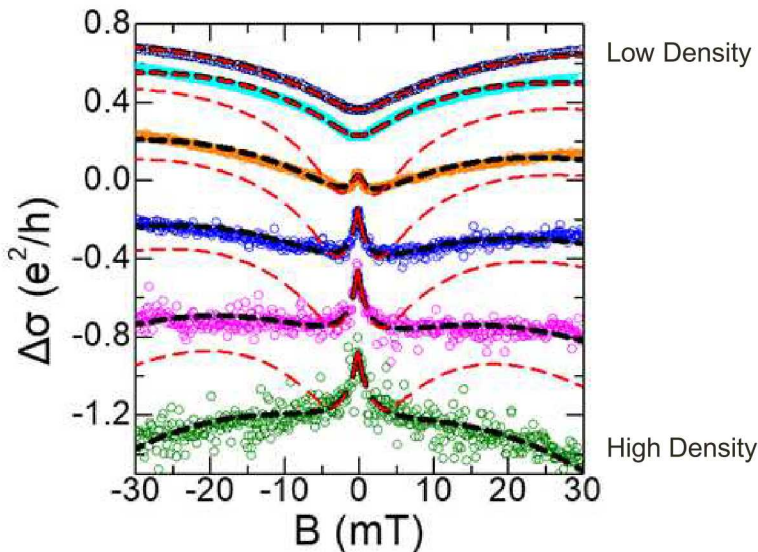
# Hole spins in Ge/SiGe provide a compelling alternative to electron spin qubits

- Absence of nearly degenerate states (i.e. valley states)
- Low Disorder (heterostructure similar to Si/SiGe)
- Enhanced Quantum Dot-Quantum Dot coupling due to a small effective mass
- Natural way to electrically control the spin (strong spin-orbit coupling) without additional components, such as micro-magnets
- Ge and Si have spin free isotopes and can be enriched.
- Potential for weaker hyperfine coupling because p-type wavefunctions vanish at the nucleus.
- Compatible with silicon processing techniques
- Can leverage designs and techniques already developed for semiconductor qubits.
- Challenge: Charge Noise

# 2DHG Properties

Strong Spin-Orbit Coupling

Weak antilocalization peak emerges with increasing density



- - - Linear Rashba Spin Orbit Coupling
- - -  $k^3$ -Rashba Spin Orbit Coupling

