

# Modeling Power-Flow Using the PERSEUS/FLEXO and HYDRA Simulation Codes

Nathaniel Hamlin, Mark Hess, Matt Weis, Charles Seyler

## Motivation and outline

*Modeling power-flow with high fidelity is of utmost importance for improving the performance of experiments on present and future pulsed power facilities.*

- Hall MHD affects the magnetic field diffusion rates into electrode plasmas, and therefore the impact of electrode plasmas on power-flow.
- PERSEUS vs HYDRA comparison:
  - For 1-D B-field penetration into conducting slab, we verified agreement between the two codes and with analytic solution.
  - For 1-D B-field penetration into plasma layer, we verified agreement between the two codes.
  - For 2-D B-field penetration into a low-density plasma, we compared the impact of Hall physics on diffusion rates, as modeled by each code.

## When does the Hall Term become important for power flow?

- When the **Hall term** becomes large relative to the **dynamo term** or **resistive term** in the Generalized Ohm's Law:

- $0 = E + \mathbf{u} \times \mathbf{B} - \frac{\mathbf{J}}{n_e e} \times \mathbf{B} - \eta \mathbf{J}$
- Large **Hall term** relative to **dynamo term**: small ion inertial length relative to spatial scales, i.e. low-density plasma.
- Large **Hall term** relative to **resistive term**: strongly magnetized plasma (electron gyrofrequency large relative to collision frequency).

## Simplifying assumptions to facilitate the modeling of magnetic diffusion with Hall physics:

- 2-D axisymmetric system (MITL and plasma); only  $J_r$ ,  $J_z$ , and  $B_\theta$  are nonzero.
- Spatially uniform resistivity and density within each material, i.e. vacuum, plasma, and electrode.
- The displacement current is negligible. This implies
  - current has negligible divergence
  - small system size compared to light propagation distance over relevant time scales
- Limited motion of plasma and electrodes.

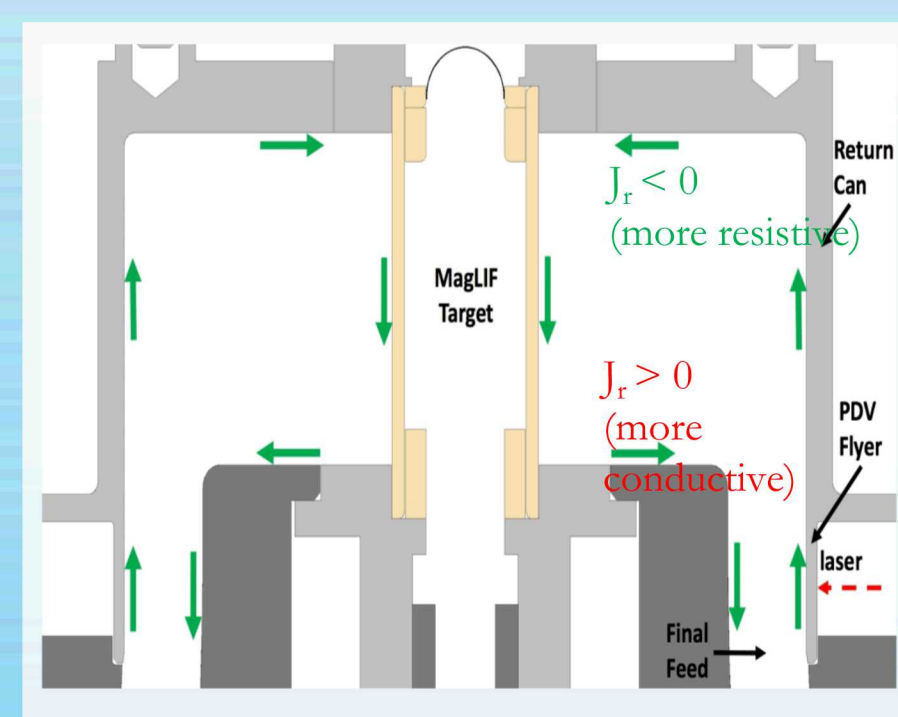
Faraday's law with Hall term:

$$-\frac{\partial \mathbf{B}}{\partial t} = \nabla \times [\eta \mathbf{J} - \mathbf{v} \times \mathbf{B} + \frac{\mathbf{J} \times \mathbf{B}}{n_e e}]$$

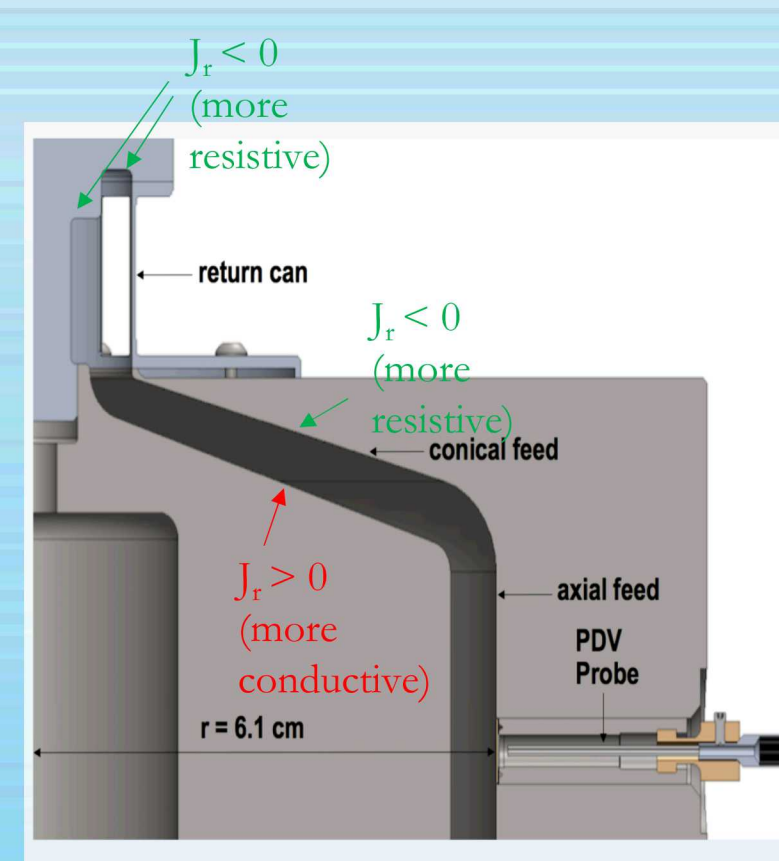
Faraday's law with Hall term and simplifying assumptions:

$$\frac{\partial (B_\theta \mathbf{e}_\theta)}{\partial t} = \eta \nabla^2 (B_\theta \mathbf{e}_\theta) - \frac{2 B_\theta J_r \mathbf{e}_\theta}{n_e r}$$

- The Hall term is only non-zero under these assumptions when  $B_\theta$  couples to  $J_r$  (Hall term is zero in 1-D simulations where  $J_r = 0$ ).
- $J_r < 0$ : material becomes more **resistive** (faster B-field diffusion) than predicted by resistive MHD.
- $J_r > 0$ : material becomes more **conductive** (slower B-field diffusion) than predicted by resistive MHD.



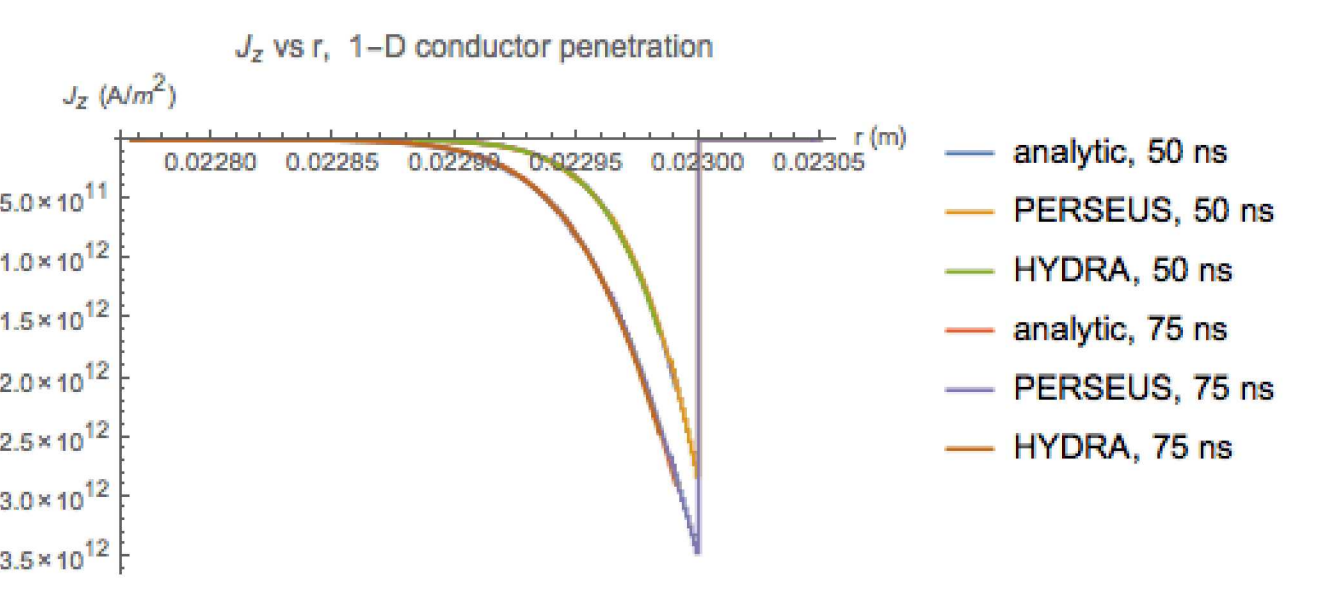
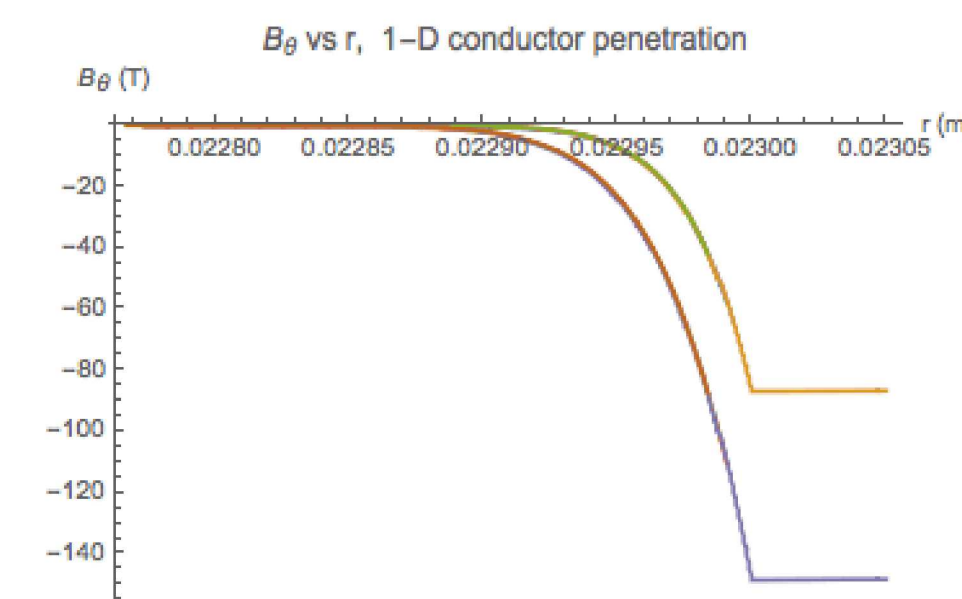
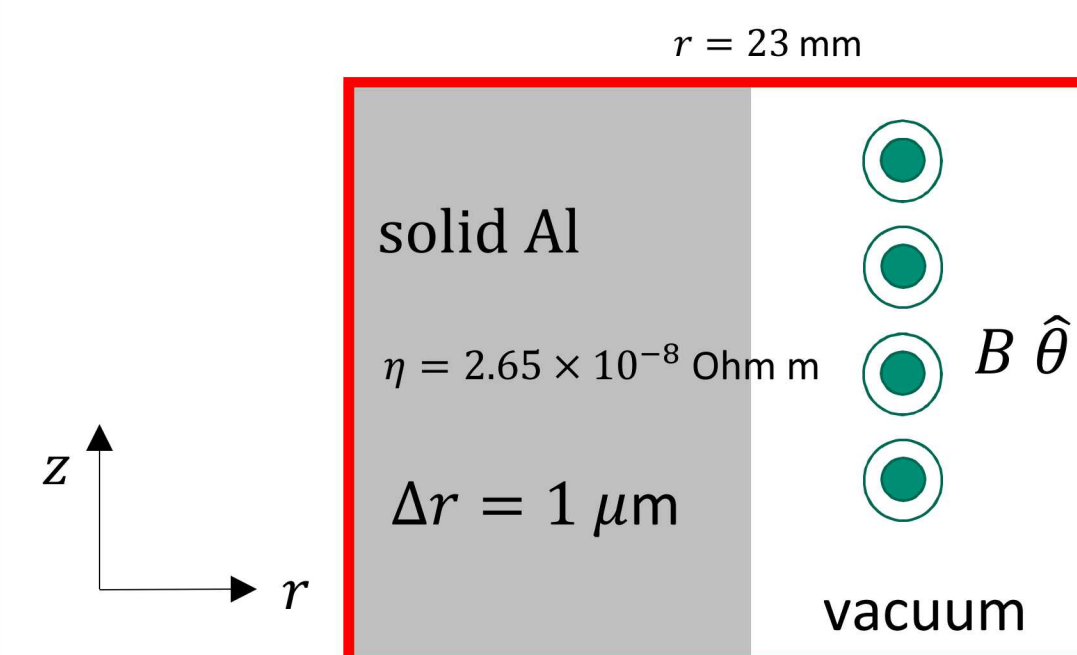
MagLIF Platform



Power Flow Platform

## 1-D magnetic diffusion:

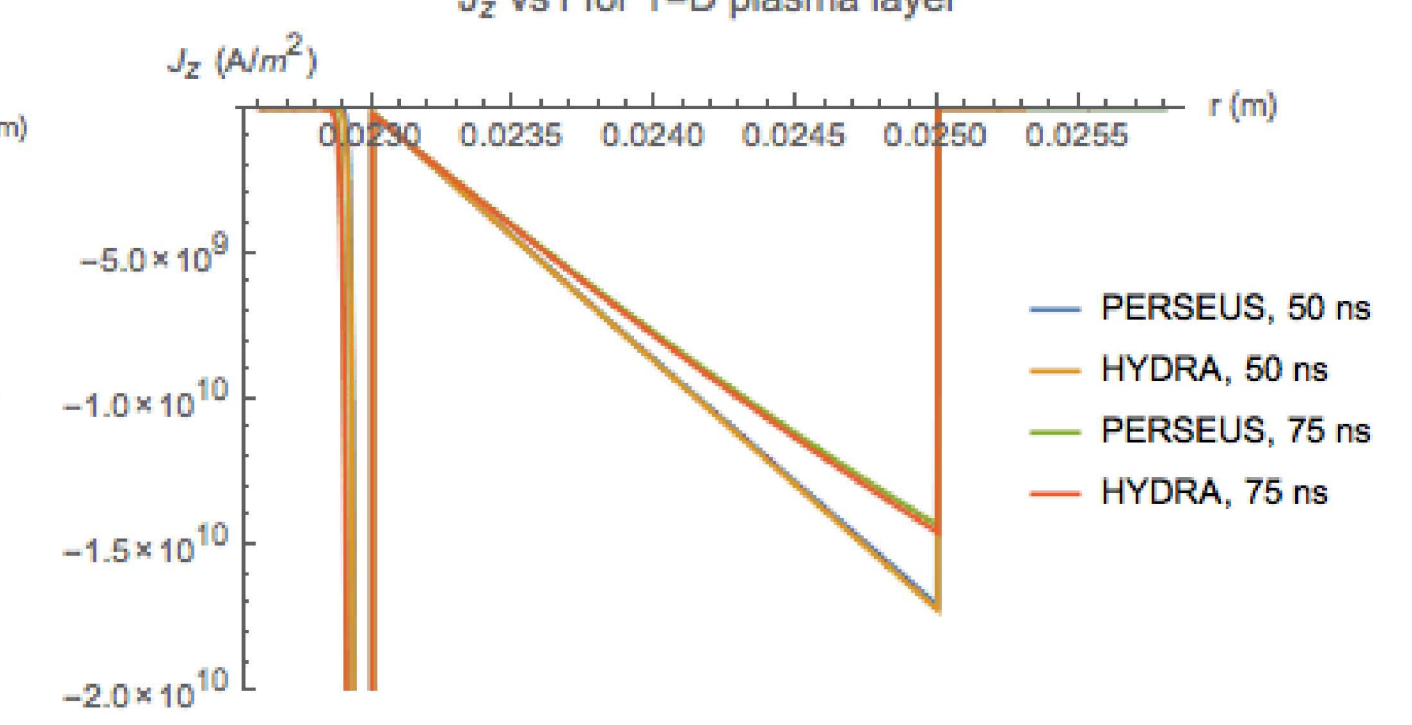
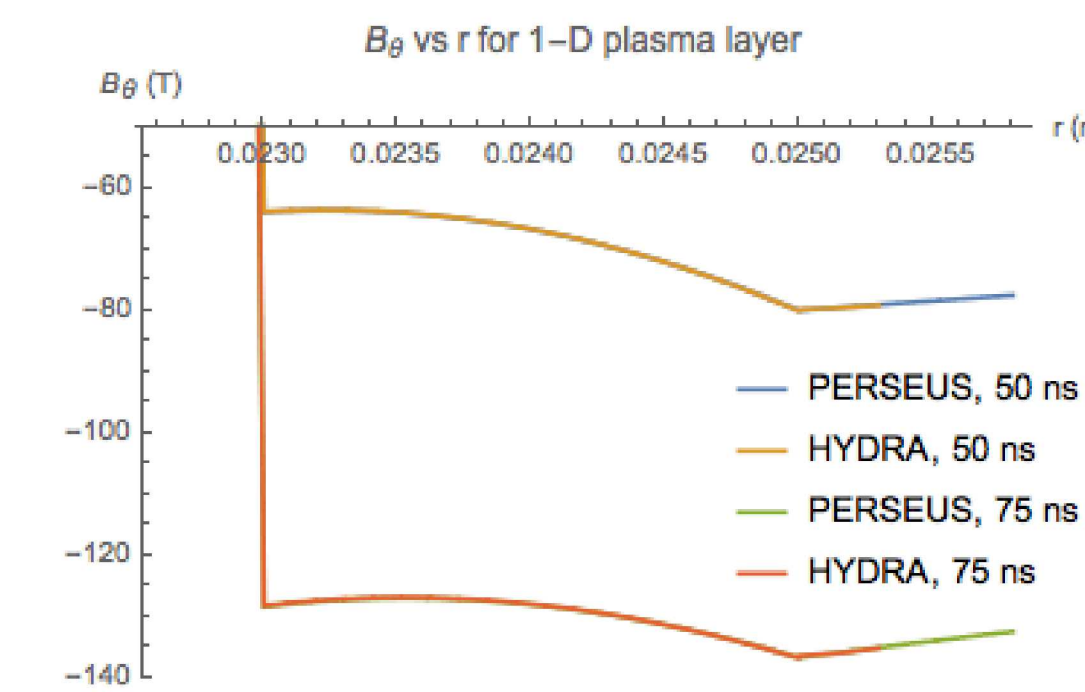
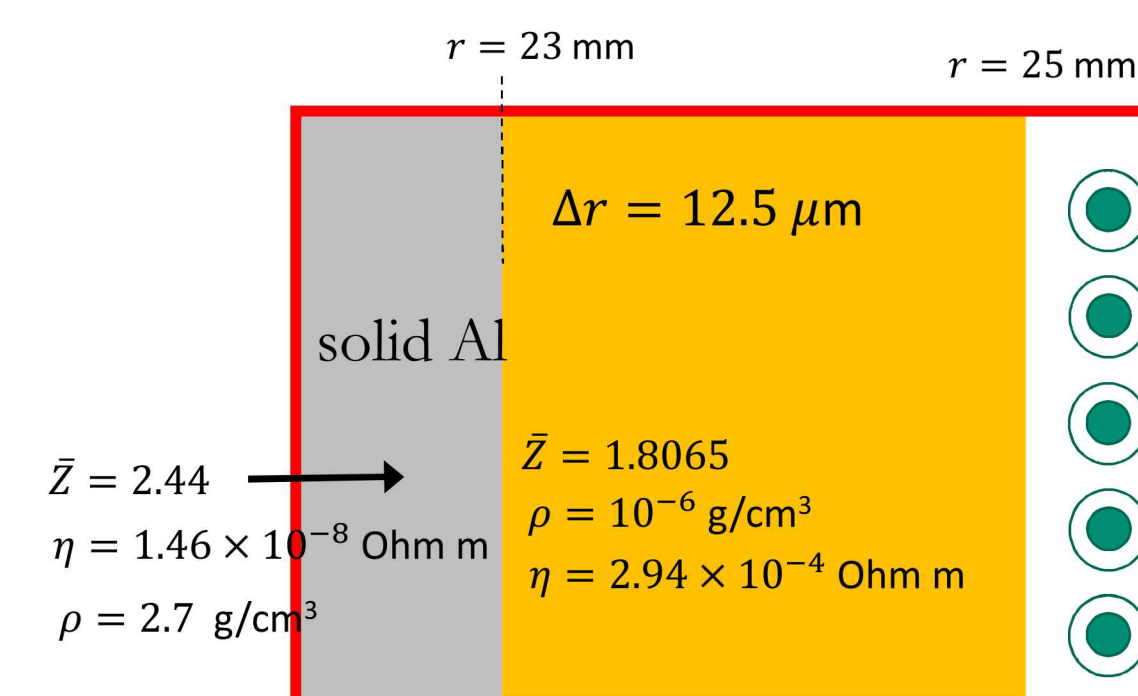
System is not influenced by Hall MHD.



very close agreement between PERSEUS, HYDRA, and analytic solution

## 1-D magnetic diffusion into plasma

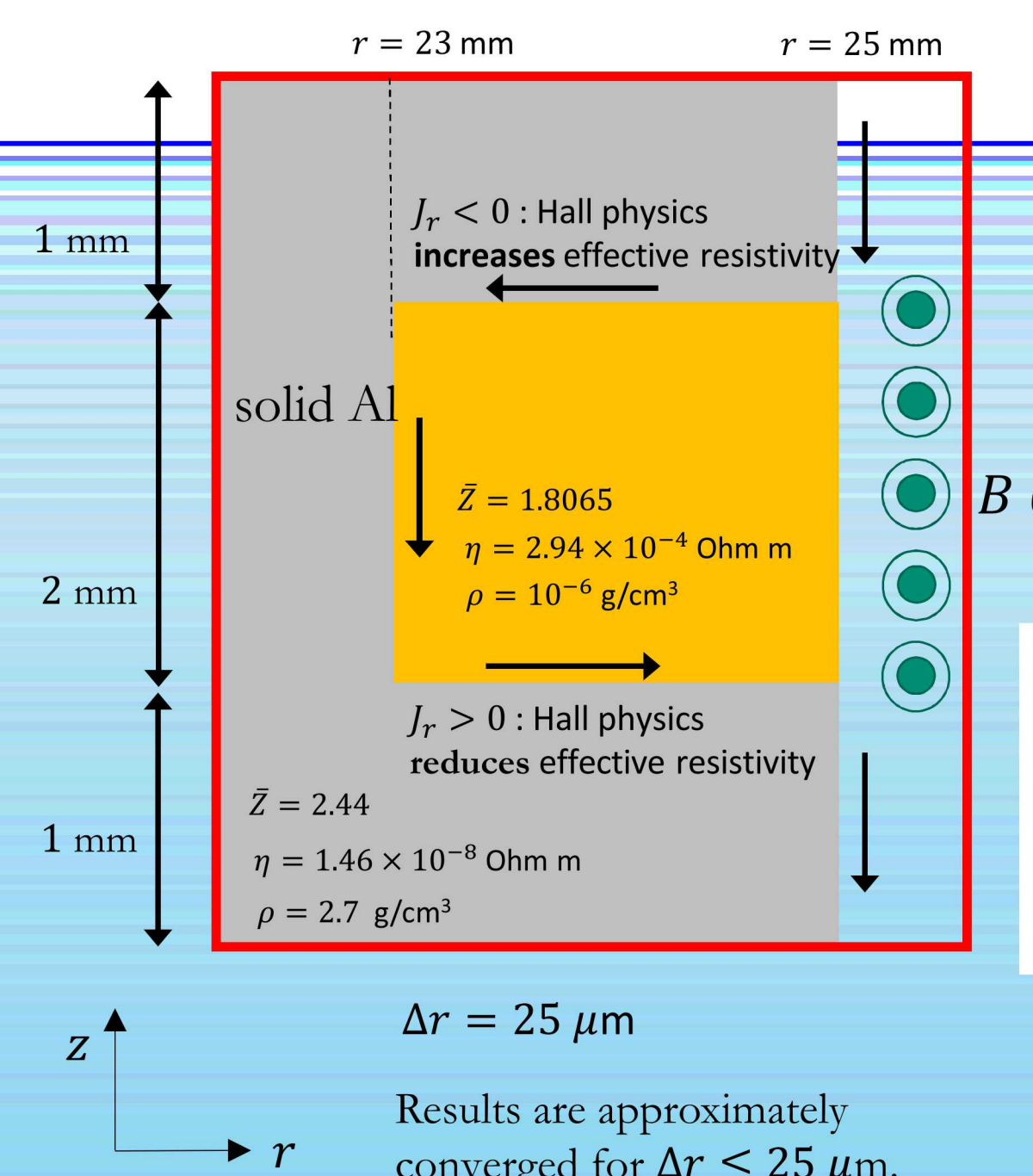
System is not influenced by Hall MHD.



very close agreement between PERSEUS and HYDRA

## 2-D magnetic diffusion into plasma

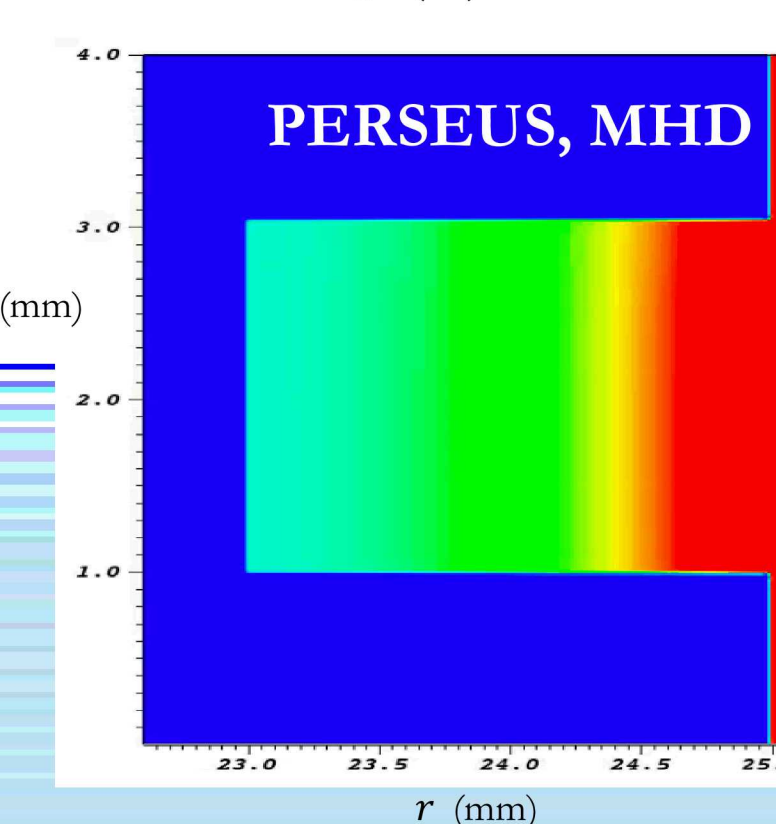
qualitative Hall MHD agreement between PERSEUS and HYDRA



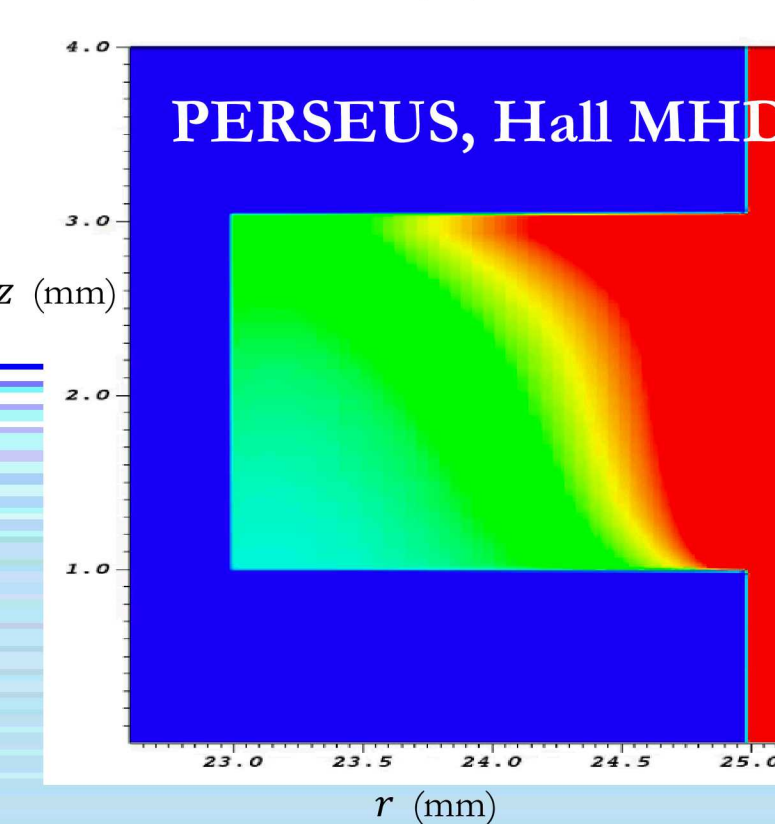
$\Delta r = 25 \mu\text{m}$

Results are approximately converged for  $\Delta r \leq 25 \mu\text{m}$ .

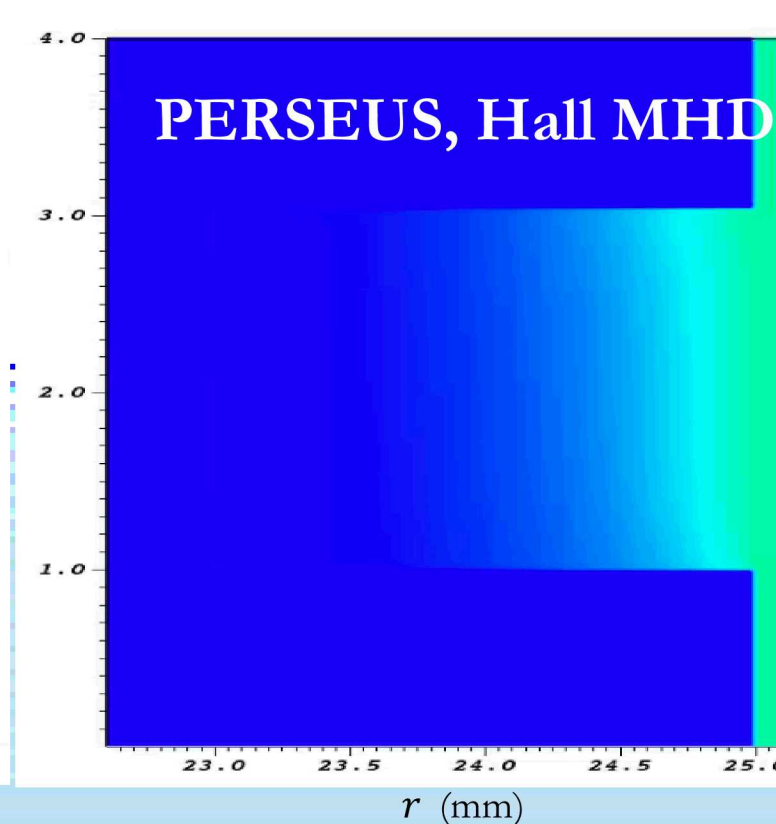
$B_\theta$  (T) at 10 ns



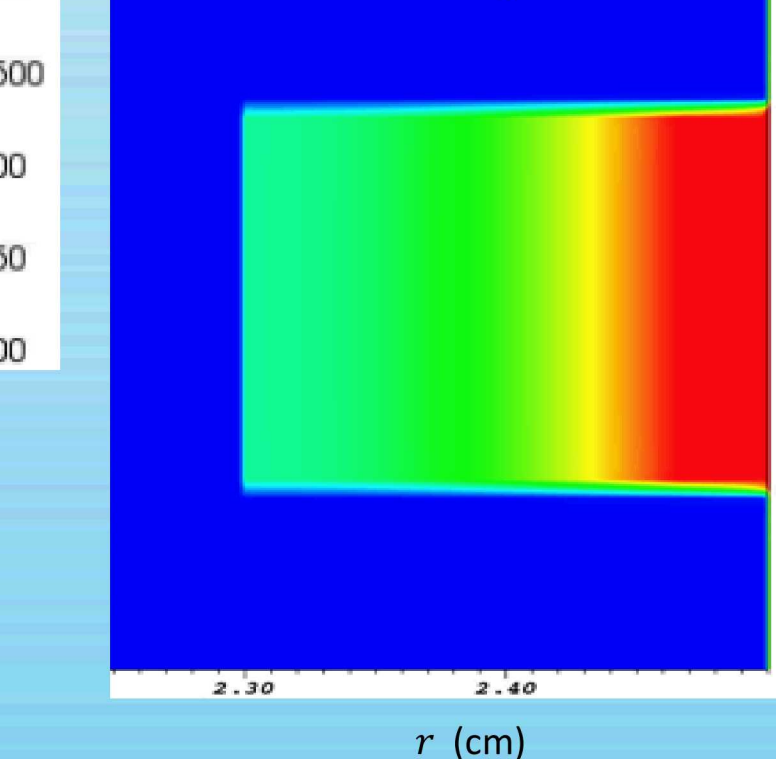
$B_\theta$  (T) at 10 ns



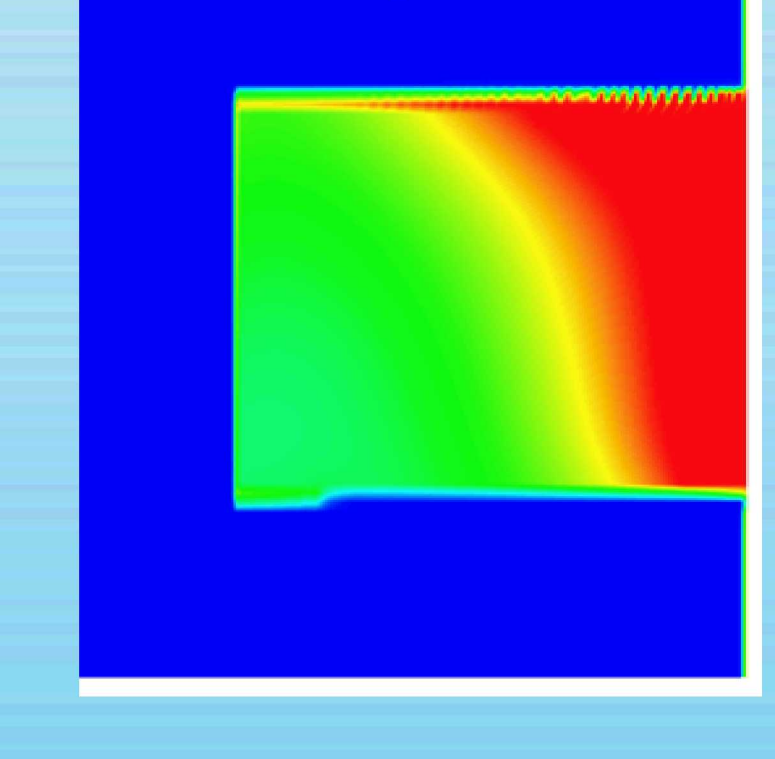
$B_\theta$  (T) at 5 ns



HYDRA, MHD



HYDRA, Hall MHD



HYDRA, Hall MHD

