

The Schwarz Alternating Method for Multiscale Coupling in Solid Mechanics

Irina Tezaur¹, Alejandro Mota¹, Coleman Alleman¹, Greg Phlipot²

¹Sandia National Laboratories, Livermore, CA, USA. ²California Institute of Technology, Pasadena, CA, USA.

COUPLED 2019

Sitges, Spain

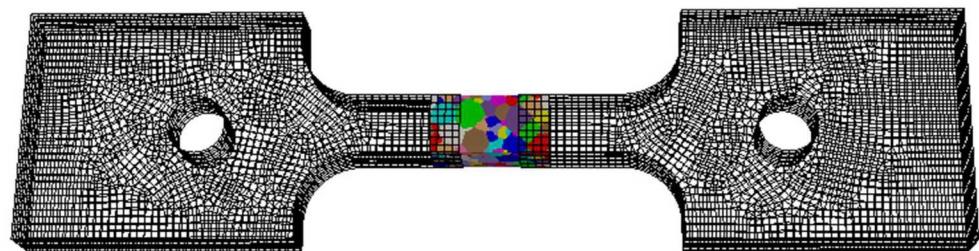
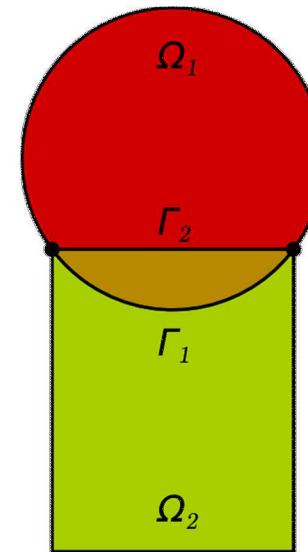
June 3-5, 2019



Sandia National Laboratories is a multi-mission laboratory managed and operated by National Technology and Engineering Solutions of Sandia, LLC, a wholly owned subsidiary of Honeywell International, Inc., for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA0003525.

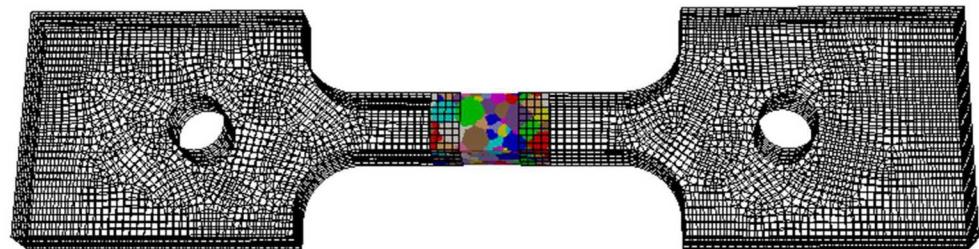
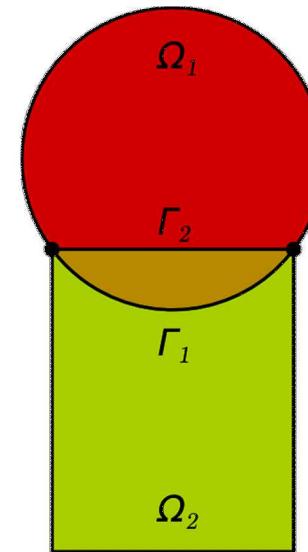
Outline

1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - Implementation
 - Numerical Examples
3. Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics
 - Formulation
 - Implementation
 - Numerical Examples
4. Summary
5. Future Work



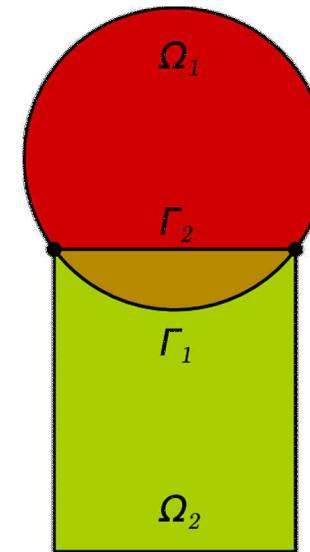
Outline

1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - Implementation
 - Numerical Examples ***COUPLED 2017 talk***
3. Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics
 - Formulation
 - Implementation
 - Numerical Examples
4. Summary
5. Future Work

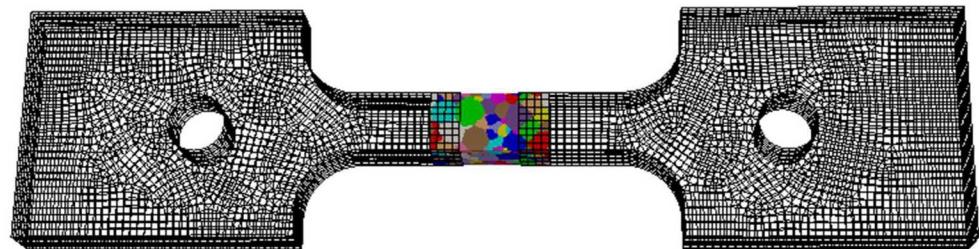


Outline

1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - Implementation
 - Numerical Examples
3. Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics
 - Formulation
 - Implementation
 - Numerical Examples
4. Summary
5. Future Work

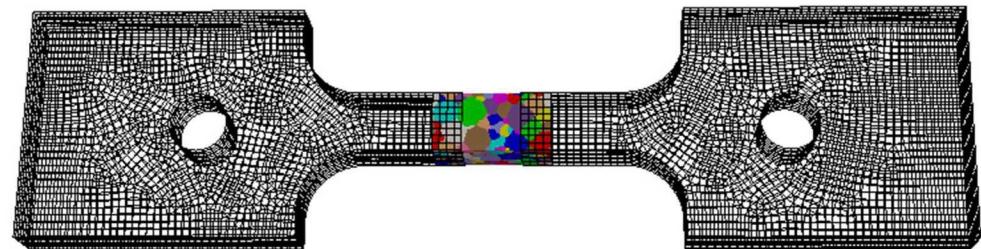
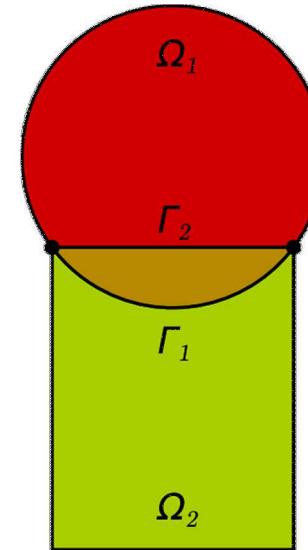


NEW!



Outline

1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - Implementation
 - Numerical Examples
3. Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics
 - Formulation
 - Implementation
 - Numerical Examples
4. Summary
5. Future Work



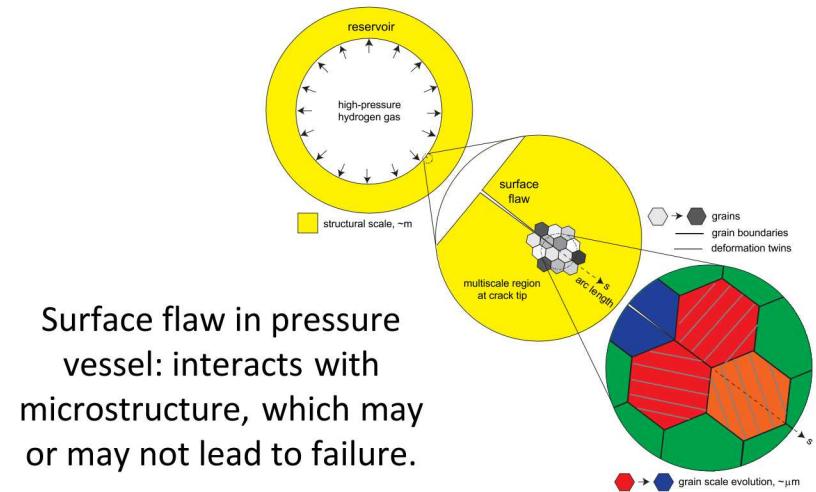
Motivation for Concurrent Multiscale Coupling

- **Large scale** structural *failure* frequently originates from **small scale** phenomena such as defects, microcracks, inhomogeneities and more, which grow quickly in unstable manner.
- Failure occurs due to **tightly coupled interaction** between small scale (stress concentrations, material instabilities, cracks, etc.) and large scale (vibration, impact, high loads and other perturbations).

Concurrent multiscale methods are **essential** for understanding and prediction of behavior of engineering systems when a **small scale failure** determines the performance of the entire system.



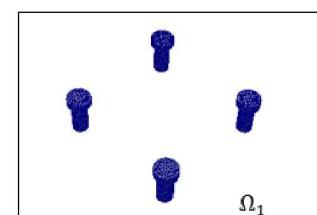
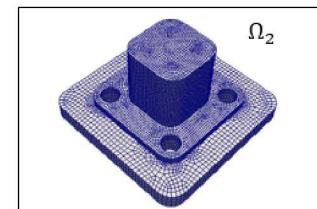
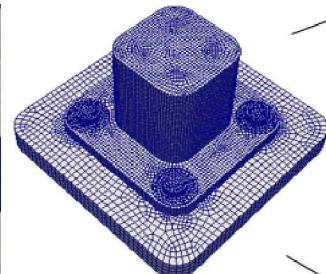
Roof failure of Boeing 737 aircraft due to fatigue cracks. From imechanica.org



Surface flaw in pressure vessel: interacts with microstructure, which may or may not lead to failure.

Requirements for Multiscale Coupling Method

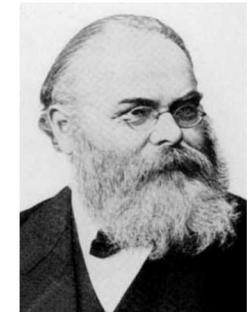
- Coupling is ***concurrent*** (two-way).
- ***Ease of implementation*** into existing massively-parallel HPC codes.
- ***Scalable, fast, robust*** (we target ***real*** engineering problems, e.g., analyses involving failure of bolted components!).
- ***"Plug-and-play" framework***: simplifies task of meshing complex geometries!
 - Ability to couple regions with ***different non-conformal meshes, different element types*** and ***different levels of refinement***.
 - Ability to use ***different solvers/time-integrators*** in different regions.
- Coupling does not introduce ***nonphysical artifacts***.
- ***Theoretical*** convergence properties/guarantees.



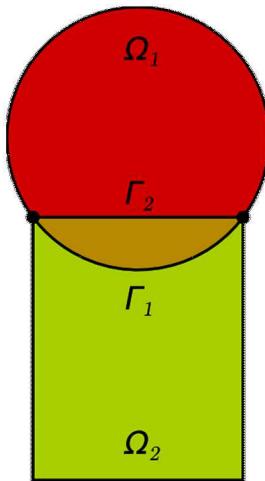
Schwarz Alternating Method for Domain Decomposition

- Proposed in 1870 by H. Schwarz for solving Laplace PDE on irregular domains.

Crux of Method: if the solution is known in regularly shaped domains, use those as pieces to iteratively build a solution for the more complex domain.



H. Schwarz (1843 – 1921)



Basic Schwarz Algorithm

Initialize:

- Solve PDE by any method on Ω_1 w/ initial guess for Dirichlet BCs on Γ_1 .

Iterate until convergence:

- Solve PDE by any method (can be different than for Ω_1) on Ω_2 w/ Dirichlet BCs on Γ_2 that are the values just obtained for Ω_1 .
- Solve PDE by any method (can be different than for Ω_2) on Ω_1 w/ Dirichlet BCs on Γ_1 that are the values just obtained for Ω_2 .

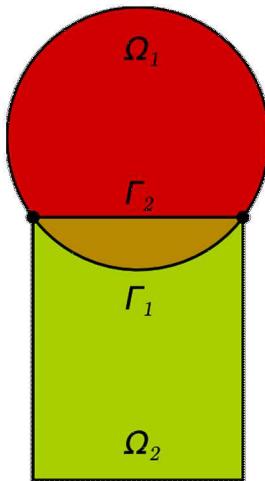
Schwarz Alternating Method for Domain Decomposition

- Proposed in 1870 by H. Schwarz for solving Laplace PDE on irregular domains.

Crux of Method: if the solution is known in regularly shaped domains, use those as pieces to iteratively build a solution for the more complex domain.



H. Schwarz (1843 – 1921)



Basic Schwarz Algorithm

Initialize:

- Solve PDE by any method on Ω_1 w/ initial guess for Dirichlet BCs on Γ_1 .

Iterate until convergence:

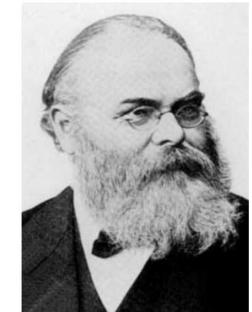
Requirement for convergence: $\Omega_1 \cap \Omega_2 \neq \emptyset$

- Solve PDE by any method (can be different than for Ω_1) on Ω_2 w/ Dirichlet BCs on Γ_2 that are the values just obtained for Ω_1 .
- Solve PDE by any method (can be different than for Ω_2) on Ω_1 w/ Dirichlet BCs on Γ_1 that are the values just obtained for Ω_2 .

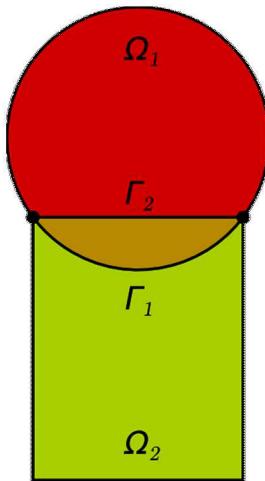
Schwarz Alternating Method for Domain Decomposition

- Proposed in 1870 by H. Schwarz for solving Laplace PDE on irregular domains.

Crux of Method: if the solution is known in regularly shaped domains, use those as pieces to iteratively build a solution for the more complex domain.



H. Schwarz (1843 – 1921)



Basic Schwarz Algorithm

Initialize:

- Solve PDE by any method on Ω_1 w/ initial guess for Dirichlet BCs on Γ_1 .

Iterate until convergence:

Requirement for convergence: $\Omega_1 \cap \Omega_2 \neq \emptyset$

- Solve PDE by any method (can be different than for Ω_1) on Ω_2 w/ Dirichlet BCs on Γ_2 that are the values just obtained for Ω_1 .
- Solve PDE by any method (can be different than for Ω_2) on Ω_1 w/ Dirichlet BCs on Γ_1 that are the values just obtained for Ω_2 .

- Schwarz alternating method most commonly used as a ***preconditioner*** for Krylov iterative methods to solve linear algebraic equations.

Schwarz Alternating Method for Domain Decomposition

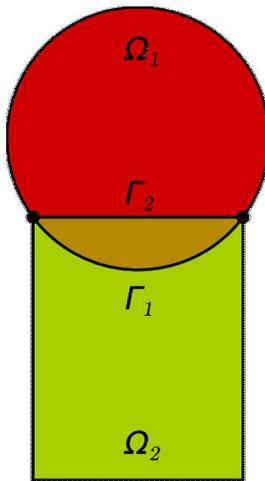


- Proposed in 1870 by H. Schwarz for solving Laplace PDE on irregular domains.

Crux of Method: if the solution is known in regularly shaped domains, use those as pieces to iteratively build a solution for the more complex domain.



H. Schwarz (1843 – 1921)



Basic Schwarz Algorithm

Initialize:

- Solve PDE by any method on Ω_1 w/ initial guess for Dirichlet BCs on Γ_1 .

Iterate until convergence:

Requirement for convergence: $\Omega_1 \cap \Omega_2 \neq \emptyset$

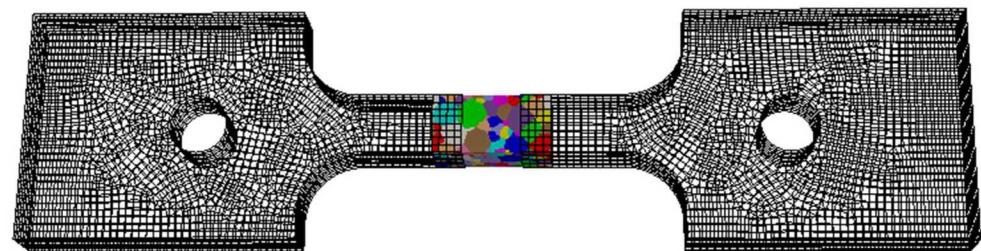
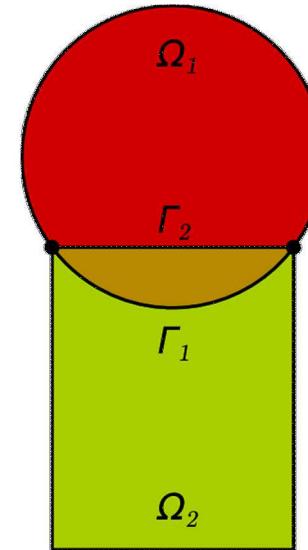
- Solve PDE by any method (can be different than for Ω_1) on Ω_2 w/ Dirichlet BCs on Γ_2 that are the values just obtained for Ω_1 .
- Solve PDE by any method (can be different than for Ω_2) on Ω_1 w/ Dirichlet BCs on Γ_1 that are the values just obtained for Ω_2 .

- Schwarz alternating method most commonly used as a **preconditioner** for Krylov iterative methods to solve linear algebraic equations.

Novel idea: using the Schwarz alternating as a **discretization method** for solving multiscale partial differential equations (PDEs).

Outline

1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - Implementation
 - Numerical Examples
3. Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics
 - Formulation
 - Implementation
 - Numerical Examples
4. Summary
5. Future Work



Schwarz Alternating Method for Multiscale Coupling in Quasistatics

1: $\varphi^{(0)} \leftarrow \text{id}_X$ in Ω_2

2: $n \leftarrow 1$

3: **repeat**

4: $\varphi^{(n)} \leftarrow \chi$ on $\partial_\varphi \Omega_i$

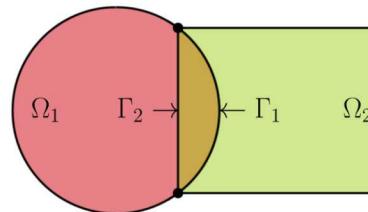
5: $\varphi^{(n)} \leftarrow P_{\Omega_j \rightarrow \Gamma_i}[\varphi^{(n-1)}]$ on Γ_i

6: $\varphi^{(n)} \leftarrow \arg \min_{\varphi \in \mathcal{S}_i} \Phi_i[\varphi]$ in Ω_i

7: $n \leftarrow n + 1$

8: **until** converged

▷ initialize to zero displacement or a better guess in Ω_2



▷ Schwarz loop

▷ Dirichlet BC for Ω_i

▷ Schwarz BC for Ω_i

▷ solve in Ω_i

Advantages:

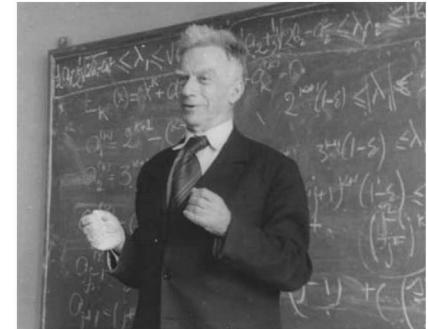
- Conceptually very *simple*.
- Allows the coupling of regions with *different non-conforming meshes, different element types*, and *different levels of refinement*.
- Information is exchanged among two or more regions, making coupling *concurrent*.
- *Different solvers* can be used for the different regions.
- *Different material models* can be coupled if they are compatible in the overlap region.
- Simplifies the task of *meshing complex geometries* for the different scales.

Theoretical Foundation

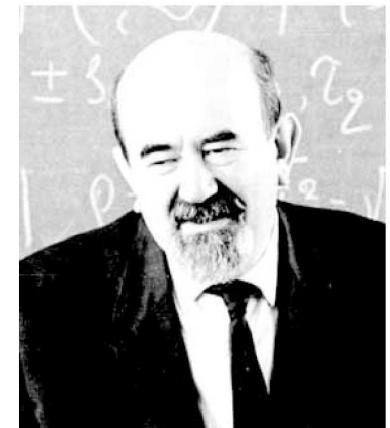
Using the Schwarz alternating as a ***discretization method*** for PDEs is natural idea with a sound ***theoretical foundation***.

- **S. L. Sobolev (1936)**: posed Schwarz method for ***linear elasticity*** in variational form and ***proved method's convergence*** by proposing a convergent sequence of energy functionals.
- **S. G. Mikhlin (1951)**: ***proved convergence*** of Schwarz method for general linear elliptic PDEs.
- **A. Mota, I. Tezaur, C. Alleman (2017)***: derived a ***proof of convergence*** of the alternating Schwarz method for the ***finite deformation quasi-static nonlinear PDEs*** (with energy functional $\Phi[\varphi]$ defined below), and determined a ***geometric convergence rate*** for the finite deformation quasi-static problem.

$$\Phi[\varphi] = \int_B W(F, Z, T) dV - \int_B \mathbf{B} \cdot \varphi dV - \int_{\partial_T B} \bar{\mathbf{T}} \cdot \varphi dS$$
$$\nabla \cdot \mathbf{P} + \mathbf{B} = \mathbf{0}$$



S. L. Sobolev (1908 – 1989)



S. G. Mikhlin (1908 – 1990)



A. Mota, I. Tezaur, C. Alleman

Four Variants* of Schwarz

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\mathbf{y}^{(1)} \leftarrow \mathbf{x}_B^{(1)}$                                          ▷ Schwarz loop
5:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$ 
6:   repeat
7:      $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$  ▷ for convergence check
8:      $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ project from  $\Omega_2$  to  $\Gamma_1$ 
9:     until  $\|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\| \leq \epsilon_{\text{machine}}$                                 ▷ Newton loop for  $\Omega_1$ 
10:     $\mathbf{y}^{(2)} \leftarrow \mathbf{x}_B^{(2)}$                                          ▷ linear system
11:     $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{P}_{21}\mathbf{x}_B^{(1)} + \mathbf{Q}_{21}\mathbf{x}_b^{(1)} + \mathbf{G}_{21}\mathbf{x}_\beta^{(1)}$            ▷ tight tolerance
12:    repeat
13:       $\Delta\mathbf{x}_B^{(2)} \leftarrow -\mathbf{K}_{AB}^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)}) \setminus \mathbf{R}_A^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)})$  ▷ for convergence check
14:       $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$                                 ▷ project from  $\Omega_1$  to  $\Gamma_2$ 
15:      until  $\|\Delta\mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\| \leq \epsilon_{\text{machine}}$                                 ▷ Newton loop for  $\Omega_2$ 
16: until  $\left[ (\|\mathbf{y}^{(1)} - \mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\|)^2 + (\|\mathbf{y}^{(2)} - \mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\|)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$  ▷ linear system

```

Full Schwarz

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$                                 ▷ Newton-Schwarz loop
5:    $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$  ▷ project from  $\Omega_2$  to  $\Gamma_1$ 
6:    $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ linear system
7:    $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{P}_{21}\mathbf{x}_B^{(1)} + \mathbf{Q}_{21}\mathbf{x}_b^{(1)} + \mathbf{G}_{21}\mathbf{x}_\beta^{(1)}$            ▷ tight tolerance
8:    $\Delta\mathbf{x}_B^{(2)} \leftarrow -\mathbf{K}_{AB}^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)}) \setminus \mathbf{R}_A^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)})$  ▷ for convergence check
9:    $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$                                 ▷ linear system
10:  until  $\left[ (\|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\|)^2 + (\|\Delta\mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\|)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$  ▷ tight tolerance

```

Modified Schwarz

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\mathbf{y}^{(1)} \leftarrow \mathbf{x}_B^{(1)}$                                          ▷ Schwarz loop
5:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$ 
6:   repeat
7:      $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$  ▷ for convergence check
8:      $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ project from  $\Omega_2$  to  $\Gamma_1$ 
9:     until  $\|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\| \leq \epsilon$                                 ▷ Newton loop for  $\Omega_1$ 
10:     $\mathbf{y}^{(2)} \leftarrow \mathbf{x}_B^{(2)}$                                          ▷ loose tolerance, e.g.  $\epsilon \in [10^{-4}, 10^{-1}]$ 
11:     $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{P}_{21}\mathbf{x}_B^{(1)} + \mathbf{Q}_{21}\mathbf{x}_b^{(1)} + \mathbf{G}_{21}\mathbf{x}_\beta^{(1)}$            ▷ for convergence check
12:    repeat
13:       $\Delta\mathbf{x}_B^{(2)} \leftarrow -\mathbf{K}_{AB}^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)}) \setminus \mathbf{R}_A^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)})$  ▷ project from  $\Omega_1$  to  $\Gamma_2$ 
14:       $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$                                 ▷ Newton loop for  $\Omega_2$ 
15:      until  $\|\Delta\mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\| \leq \epsilon$                                 ▷ solve linear system
16: until  $\left[ (\|\mathbf{y}^{(1)} - \mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\|)^2 + (\|\mathbf{y}^{(2)} - \mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\|)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$  ▷ loose tolerance, e.g.  $\epsilon \in [10^{-4}, 10^{-1}]$ 

```

Inexact Schwarz

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\left\{ \begin{array}{l} \Delta\mathbf{x}_B^{(1)} \\ \Delta\mathbf{x}_B^{(2)} \end{array} \right\} \leftarrow \begin{pmatrix} \mathbf{K}_{AB}^{(1)} + \mathbf{K}_{A\beta}^{(1)}\mathbf{H}_{11} & \mathbf{K}_{A\beta}^{(1)}\mathbf{H}_{12} \\ \mathbf{K}_{A\beta}^{(2)}\mathbf{H}_{21} & \mathbf{K}_{AB}^{(2)} + \mathbf{K}_{A\beta}^{(2)}\mathbf{H}_{22} \end{pmatrix} \setminus \left\{ \begin{array}{l} -\mathbf{R}_A^{(1)} \\ -\mathbf{R}_A^{(2)} \end{array} \right\}$  ▷ Newton-Schwarz loop
5:    $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$ 
6:    $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$ 
7: until  $\left[ (\|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\|)^2 + (\|\Delta\mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\|)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$  ▷ tight tolerance

```

Monolithic Schwarz

Four Variants* of Schwarz

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\mathbf{y}^{(1)} \leftarrow \mathbf{x}_B^{(1)}$                                          ▷ Schwarz loop
5:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$ 
6:   repeat
7:      $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$  ▷ for convergence check
8:      $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ project from  $\Omega_2$  to  $\Gamma_1$ 
9:     repeat  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{x}_\beta^{(1)}$                                          ▷ Newton loop for  $\Omega_1$ 
10:     $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$  ▷ linear system
11:     $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ tight tolerance
12:    repeat  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{x}_\beta^{(1)}$                                          ▷ for convergence check
13:       $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$  ▷ project from  $\Omega_1$  to  $\Gamma_2$ 
14:       $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ Newton loop for  $\Omega_2$ 
15:      until  $\|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\| \leq \epsilon_{\text{machine}}$                                 ▷ linear system
16:       $\left[ (\|\mathbf{y}^{(1)} - \mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\|)^2 + (\|\mathbf{y}^{(2)} - \mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\|)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$  ▷ tight tolerance

```

Full Schwarz

Least-intrusive variant: by-passes Schwarz iteration, no need for block solver.

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$                                          ▷ Newton-Schwarz loop
5:    $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$  ▷ project from  $\Omega_2$  to  $\Gamma_1$ 
6:    $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ linear system
7:    $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{P}_{21}\mathbf{x}_B^{(1)} + \mathbf{Q}_{21}\mathbf{x}_b^{(1)} + \mathbf{G}_{21}\mathbf{x}_\beta^{(1)}$                                 ▷ project from  $\Omega_1$  to  $\Gamma_2$ 
8:    $\Delta\mathbf{x}_B^{(2)} \leftarrow -\mathbf{K}_{AB}^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)}) \setminus \mathbf{R}_A^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)})$  ▷ linear system
9:    $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$                                 ▷ tight tolerance
10:  until  $\left[ (\|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\|)^2 + (\|\Delta\mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\|)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$  ▷ tight tolerance

```

Modified Schwarz

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\mathbf{y}^{(1)} \leftarrow \mathbf{x}_B^{(1)}$                                          ▷ Schwarz loop
5:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$ 
6:   repeat
7:      $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$  ▷ for convergence check
8:      $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ project from  $\Omega_2$  to  $\Gamma_1$ 
9:     repeat  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{x}_\beta^{(1)}$                                          ▷ Newton loop for  $\Omega_1$ 
10:     $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$  ▷ linear system
11:     $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ loose tolerance, e.g.  $\epsilon \in [10^{-4}, 10^{-1}]$ 
12:    repeat  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{x}_\beta^{(1)}$                                          ▷ for convergence check
13:       $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$  ▷ project from  $\Omega_1$  to  $\Gamma_2$ 
14:       $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ Newton loop for  $\Omega_2$ 
15:      until  $\|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\| \leq \epsilon$                                 ▷ solve linear system
16:       $\left[ (\|\mathbf{y}^{(1)} - \mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\|)^2 + (\|\mathbf{y}^{(2)} - \mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\|)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$  ▷ loose tolerance, e.g.  $\epsilon \in [10^{-4}, 10^{-1}]$ 

```

Inexact Schwarz

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\left\{ \begin{array}{l} \Delta\mathbf{x}_B^{(1)} \\ \Delta\mathbf{x}_B^{(2)} \end{array} \right\} \leftarrow \begin{pmatrix} \mathbf{K}_{AB}^{(1)} + \mathbf{K}_{A\beta}^{(1)}\mathbf{H}_{11} & \mathbf{K}_{A\beta}^{(1)}\mathbf{H}_{12} \\ \mathbf{K}_{A\beta}^{(2)}\mathbf{H}_{21} & \mathbf{K}_{AB}^{(2)} + \mathbf{K}_{A\beta}^{(2)}\mathbf{H}_{22} \end{pmatrix} \setminus \left\{ \begin{array}{l} -\mathbf{R}_A^{(1)} \\ -\mathbf{R}_A^{(2)} \end{array} \right\}$  ▷ Newton-Schwarz loop
5:    $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ linear system
6:    $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$                                 ▷ tight tolerance
7:   until  $\left[ (\|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\|)^2 + (\|\Delta\mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\|)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$ 

```

Monolithic Schwarz

Four Variants* of Schwarz

Most performant method: monotonic convergence, theoretical convergence guarantee.

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\mathbf{y}^{(1)} \leftarrow \mathbf{x}_B^{(1)}$                                          ▷ Schwarz loop
5:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$ 
6:   repeat
7:      $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$     ▷ for convergence check
8:      $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ project from  $\Omega_2$  to  $\Gamma_1$ 
9:     until  $\|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\| \leq \epsilon_{\text{machine}}$                          ▷ Newton loop for  $\Omega_1$ 
10:     $\mathbf{y}^{(2)} \leftarrow \mathbf{x}_B^{(2)}$                                          ▷ linear system
11:     $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{P}_{21}\mathbf{x}_B^{(1)} + \mathbf{Q}_{21}\mathbf{x}_b^{(1)} + \mathbf{G}_{21}\mathbf{x}_\beta^{(1)}$ 
12:    repeat
13:       $\Delta\mathbf{x}_B^{(2)} \leftarrow -\mathbf{K}_{AB}^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)}) \setminus \mathbf{R}_A^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)})$     ▷ for convergence check
14:       $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$                                 ▷ project from  $\Omega_1$  to  $\Gamma_2$ 
15:      until  $\|\Delta\mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\| \leq \epsilon_{\text{machine}}$                          ▷ Newton loop for  $\Omega_2$ 
16: until  $\left[ \left( \|\mathbf{y}^{(1)} - \mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\| \right)^2 + \left( \|\mathbf{y}^{(2)} - \mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\| \right)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$     ▷ tight tolerance
  
```

Full Schwarz

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$                                 ▷ Newton-Schwarz loop
5:    $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$     ▷ for convergence check
6:    $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ linear system
7:    $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{P}_{21}\mathbf{x}_B^{(1)} + \mathbf{Q}_{21}\mathbf{x}_b^{(1)} + \mathbf{G}_{21}\mathbf{x}_\beta^{(1)}$ 
8:    $\Delta\mathbf{x}_B^{(2)} \leftarrow -\mathbf{K}_{AB}^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)}) \setminus \mathbf{R}_A^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)})$     ▷ project from  $\Omega_1$  to  $\Gamma_2$ 
9:    $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$                                 ▷ linear system
10: until  $\left[ \left( \|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\| \right)^2 + \left( \|\Delta\mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\| \right)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$     ▷ tight tolerance
  
```

Modified Schwarz

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\mathbf{y}^{(1)} \leftarrow \mathbf{x}_B^{(1)}$                                          ▷ Schwarz loop
5:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$ 
6:   repeat
7:      $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$     ▷ for convergence check
8:      $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$                                 ▷ Newton loop for  $\Omega_1$ 
9:     until  $\|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\| \leq \epsilon$                                 ▷ linear system
10:     $\mathbf{y}^{(2)} \leftarrow \mathbf{x}_B^{(2)}$                                          ▷ for convergence check
11:     $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{P}_{21}\mathbf{x}_B^{(1)} + \mathbf{Q}_{21}\mathbf{x}_b^{(1)} + \mathbf{G}_{21}\mathbf{x}_\beta^{(1)}$ 
12:    repeat
13:       $\Delta\mathbf{x}_B^{(2)} \leftarrow -\mathbf{K}_{AB}^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)}) \setminus \mathbf{R}_A^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)})$     ▷ project from  $\Omega_1$  to  $\Gamma_2$ 
14:       $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$                                 ▷ Newton loop for  $\Omega_2$ 
15:      until  $\|\Delta\mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\| \leq \epsilon$                                 ▷ solve linear system
16: until  $\left[ \left( \|\mathbf{y}^{(1)} - \mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\| \right)^2 + \left( \|\mathbf{y}^{(2)} - \mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\| \right)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$     ▷ tight tolerance
  
```

Inexact Schwarz

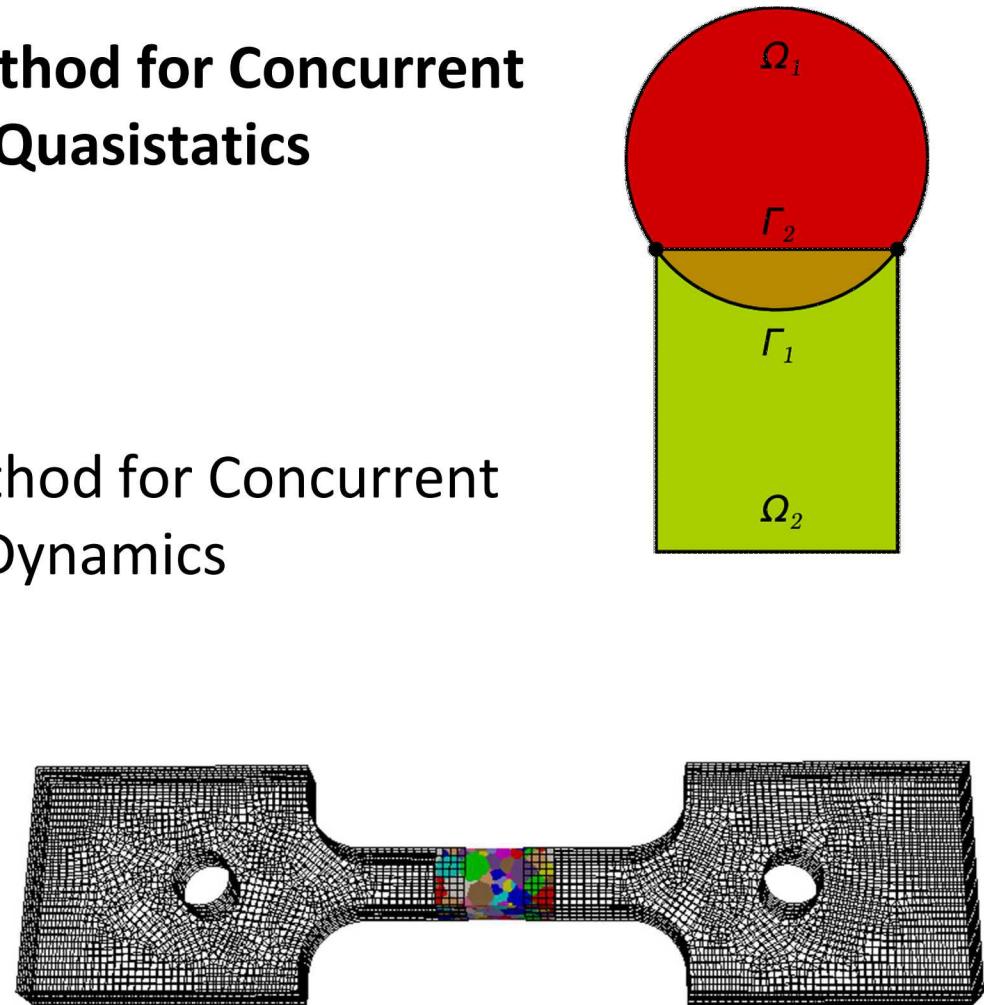
```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat
4:    $\left\{ \begin{array}{l} \Delta\mathbf{x}_B^{(1)} \\ \Delta\mathbf{x}_B^{(2)} \end{array} \right\} \leftarrow \begin{pmatrix} \mathbf{K}_{AB}^{(1)} + \mathbf{K}_{A\beta}^{(1)}\mathbf{H}_{11} & \mathbf{K}_{A\beta}^{(1)}\mathbf{H}_{12} \\ \mathbf{K}_{A\beta}^{(2)}\mathbf{H}_{21} & \mathbf{K}_{AB}^{(2)} + \mathbf{K}_{A\beta}^{(2)}\mathbf{H}_{22} \end{pmatrix} \setminus \left\{ \begin{array}{l} -\mathbf{R}_A^{(1)} \\ -\mathbf{R}_A^{(2)} \end{array} \right\}$     ▷ Newton-Schwarz loop
5:    $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$ 
6:    $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$ 
7: until  $\left[ \left( \|\Delta\mathbf{x}_B^{(1)}\|/\|\mathbf{x}_B^{(1)}\| \right)^2 + \left( \|\Delta\mathbf{x}_B^{(2)}\|/\|\mathbf{x}_B^{(2)}\| \right)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$     ▷ tight tolerance
  
```

Monolithic Schwarz

Outline

1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - **Implementation**
 - Numerical Examples
3. Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics
 - Formulation
 - Implementation
 - Numerical Examples
4. Summary
5. Future Work



Implementation within *Albany* Code

The proposed *quasistatic alternating Schwarz method* is implemented within the *LCM project* in Sandia's open-source parallel, C++, multi-physics, finite element code, *Albany*.

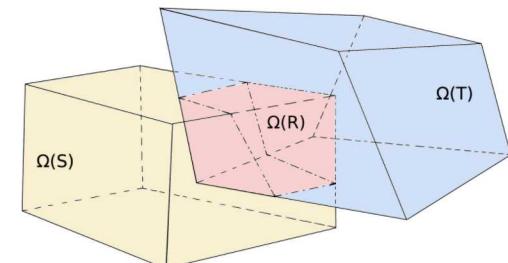


<https://github.com/gahansen/Albany>

- **Component-based** design for rapid development of capabilities.
- Contains a wide variety of **constitutive models**.
- Extensive use of libraries from the open-source *Trilinos* project.
 - Use of the *Phalanx* package to decompose complex problem into simpler problems with managed dependencies.
 - Use of the *Sacado* package for **automatic differentiation**.
 - Use of *Teko* package for block preconditioning.
- **Parallel** implementation of Schwarz alternating method uses the ***Data Transfer Kit (DTK)***.
- All software available on ***GitHub***.



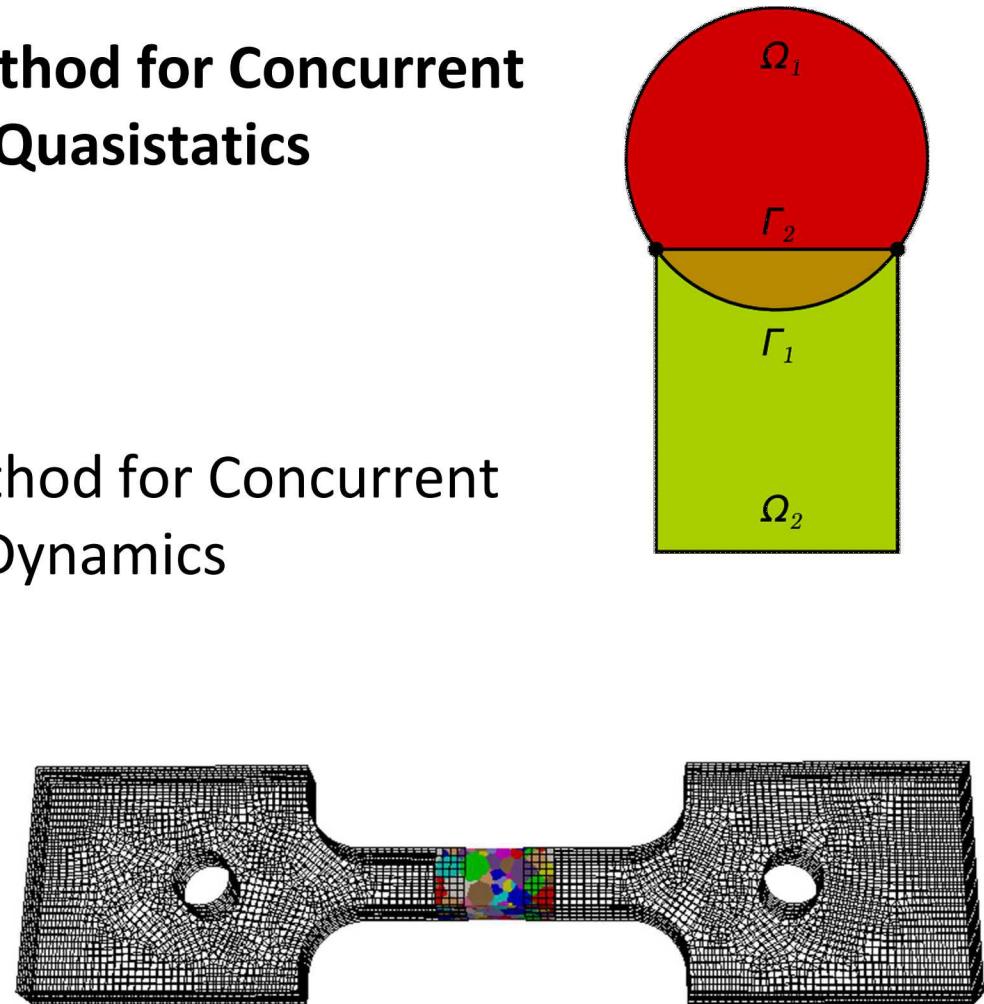
<https://github.com/trilinos/trilinos>



<https://github.com/ORNL-CEES/DataTransferKit>

Outline

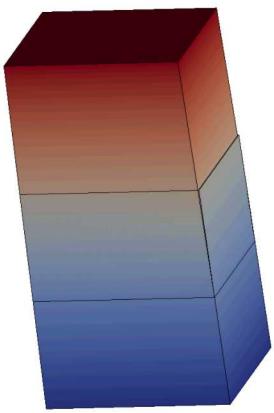
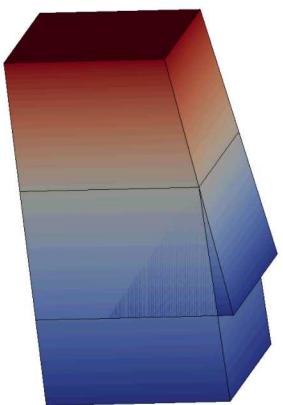
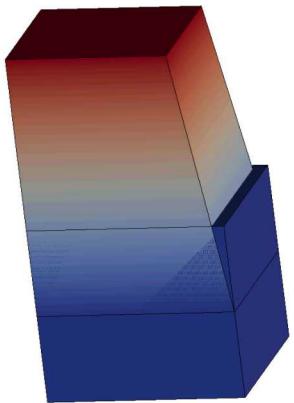
1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - Implementation
 - Numerical Examples
3. Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics
 - Formulation
 - Implementation
 - Numerical Examples
4. Summary
5. Future Work



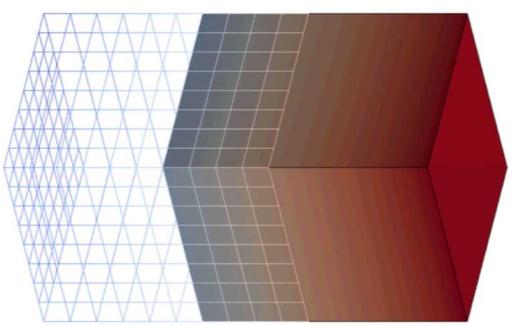
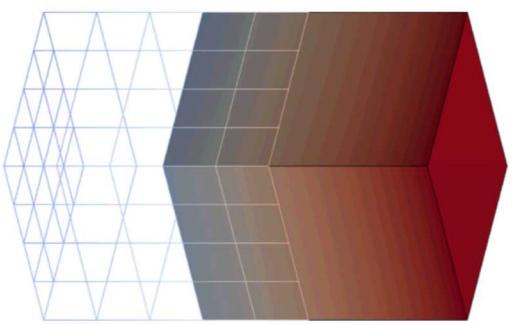
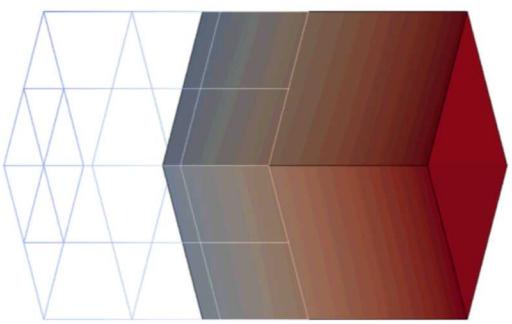
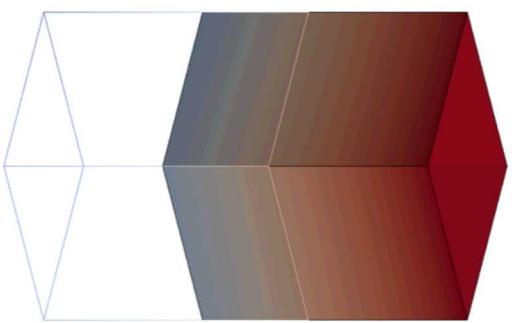
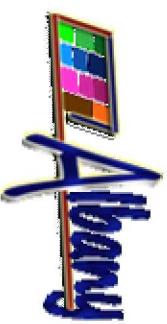
Quasistatic Example #1: Cuboid Problem



Schwarz Iteration

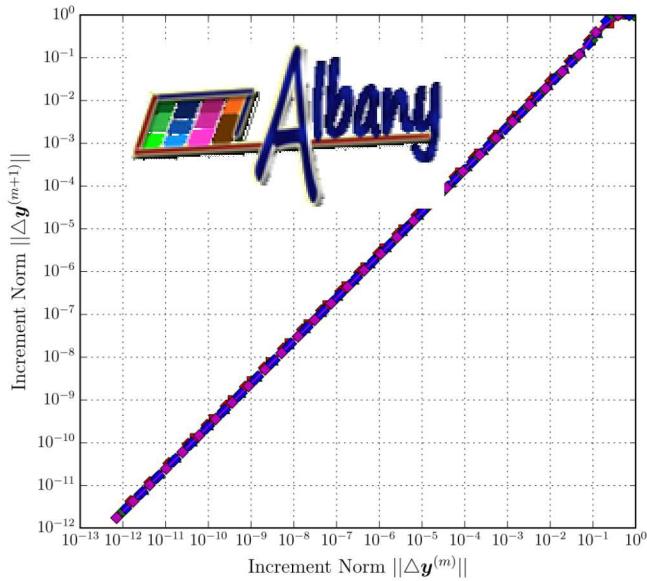


- Coupling of **two cuboids** with square base (above).
- **Neo-hookean**-type material model.

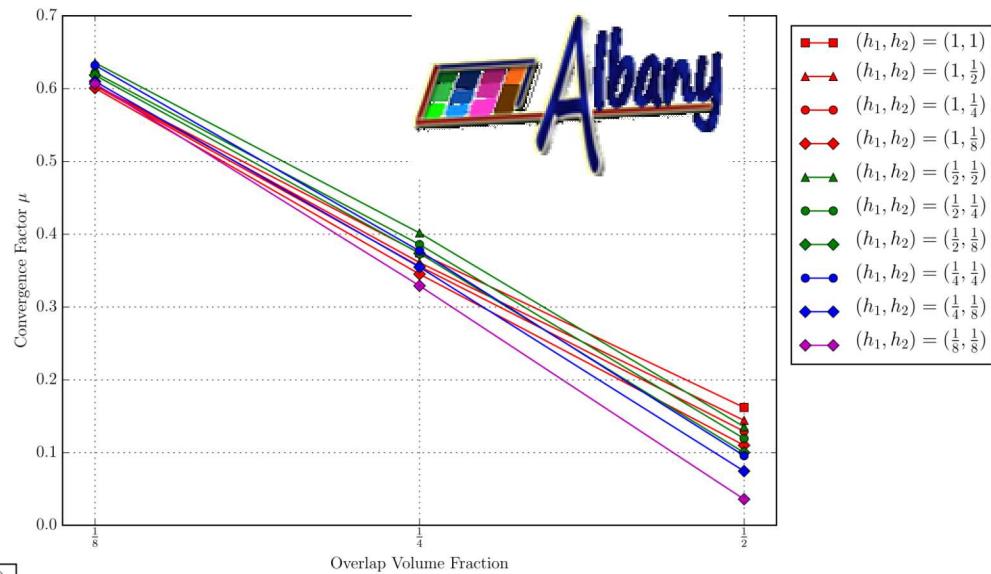


Cuboid Problem: Convergence with Overlap & Refinement

Below: Convergence of the cuboid problem for different mesh sizes and fixed overlap volume fraction. The Schwarz alternating method converges *linearly*.



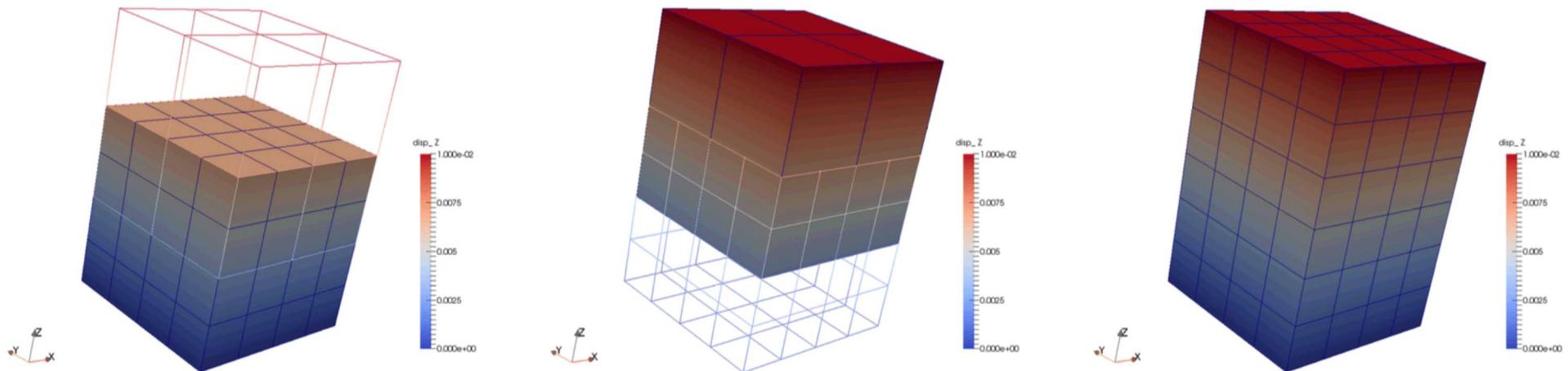
| | |
|---|---|
| ■ | $(h_1, h_2) = (1, 1)$ |
| ▲ | $(h_1, h_2) = (1, \frac{1}{2})$ |
| ● | $(h_1, h_2) = (1, \frac{1}{4})$ |
| ◆ | $(h_1, h_2) = (1, \frac{1}{8})$ |
| ▲ | $(h_1, h_2) = (\frac{1}{2}, \frac{1}{2})$ |
| ● | $(h_1, h_2) = (\frac{1}{2}, \frac{1}{4})$ |
| ◆ | $(h_1, h_2) = (\frac{1}{2}, \frac{1}{8})$ |
| ● | $(h_1, h_2) = (\frac{1}{4}, \frac{1}{4})$ |
| ◆ | $(h_1, h_2) = (\frac{1}{4}, \frac{1}{8})$ |
| ● | $(h_1, h_2) = (\frac{1}{8}, \frac{1}{8})$ |



Above: Convergence factor μ as a function of overlap volume and different mesh. There is **faster linear convergence** with increasing **overlap volume fraction**.

$$\Delta\mathbf{y}^{(m+1)} \leq \mu \Delta\mathbf{y}^{(m)}$$

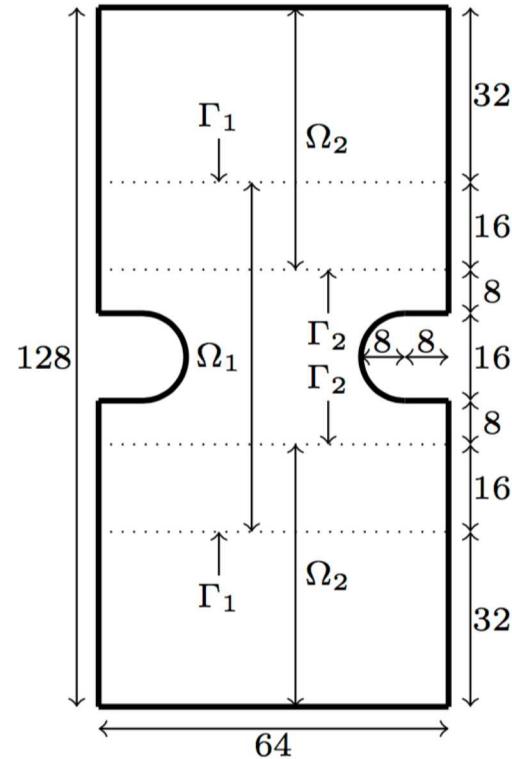
Cuboid Problem: Schwarz Error



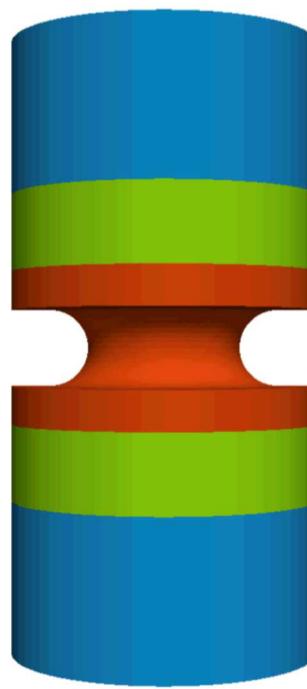
| Subdomain | u_3 relative error | σ_{33} relative error |
|------------|------------------------|------------------------------|
| Ω_1 | 1.24×10^{-14} | 2.31×10^{-13} |
| Ω_2 | 7.30×10^{-15} | 3.06×10^{-13} |



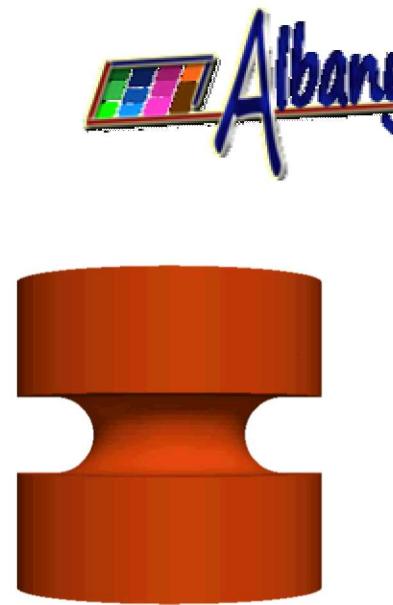
Quasistatic Example #2: Notched Cylinder



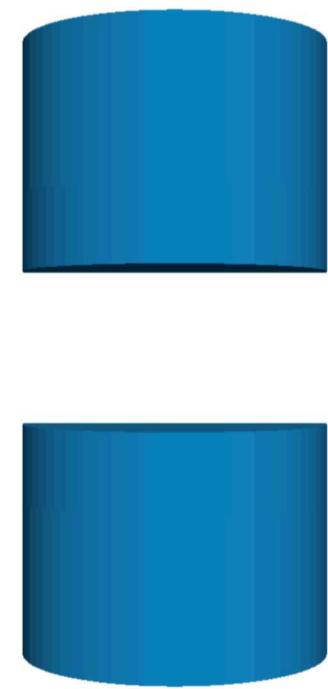
(a) Schematic



(b) Entire Domain Ω



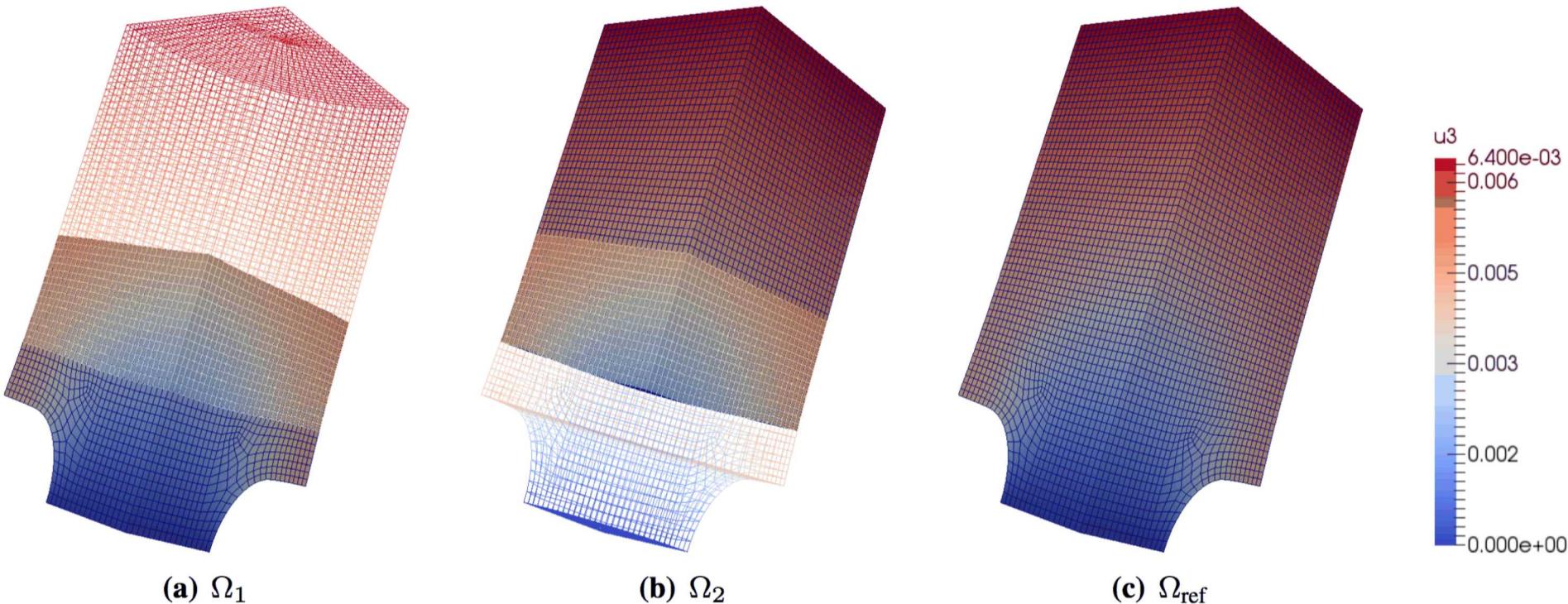
(c) Fine Region Ω_1



(d) Coarse Region Ω_2

- **Notched cylinder** that is stretched along its axial direction.
- Domain decomposed into **two subdomains**.
- **Neo-hookean**-type material model.

Notched Cylinder: Conformal HEX-HEX Coupling

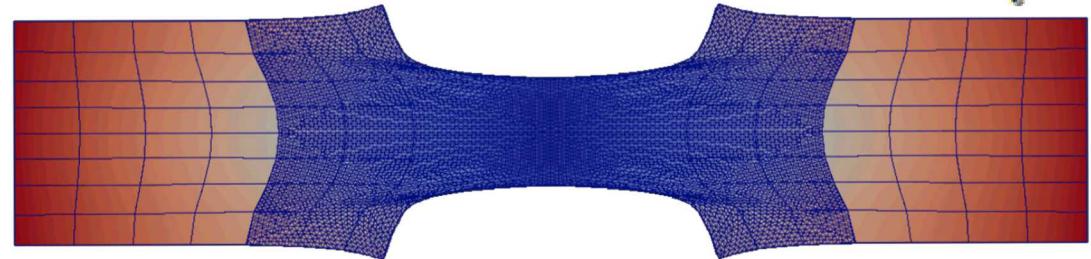
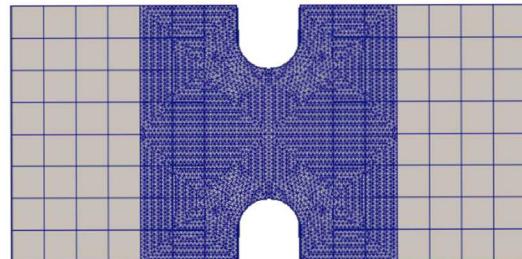
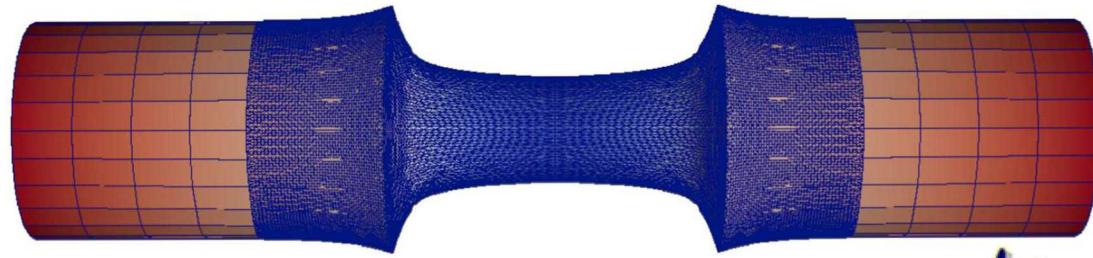
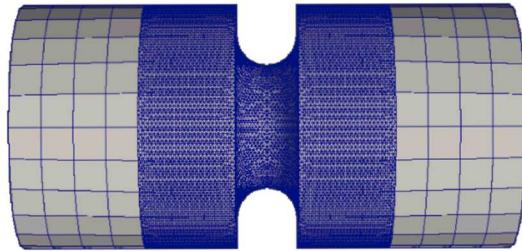


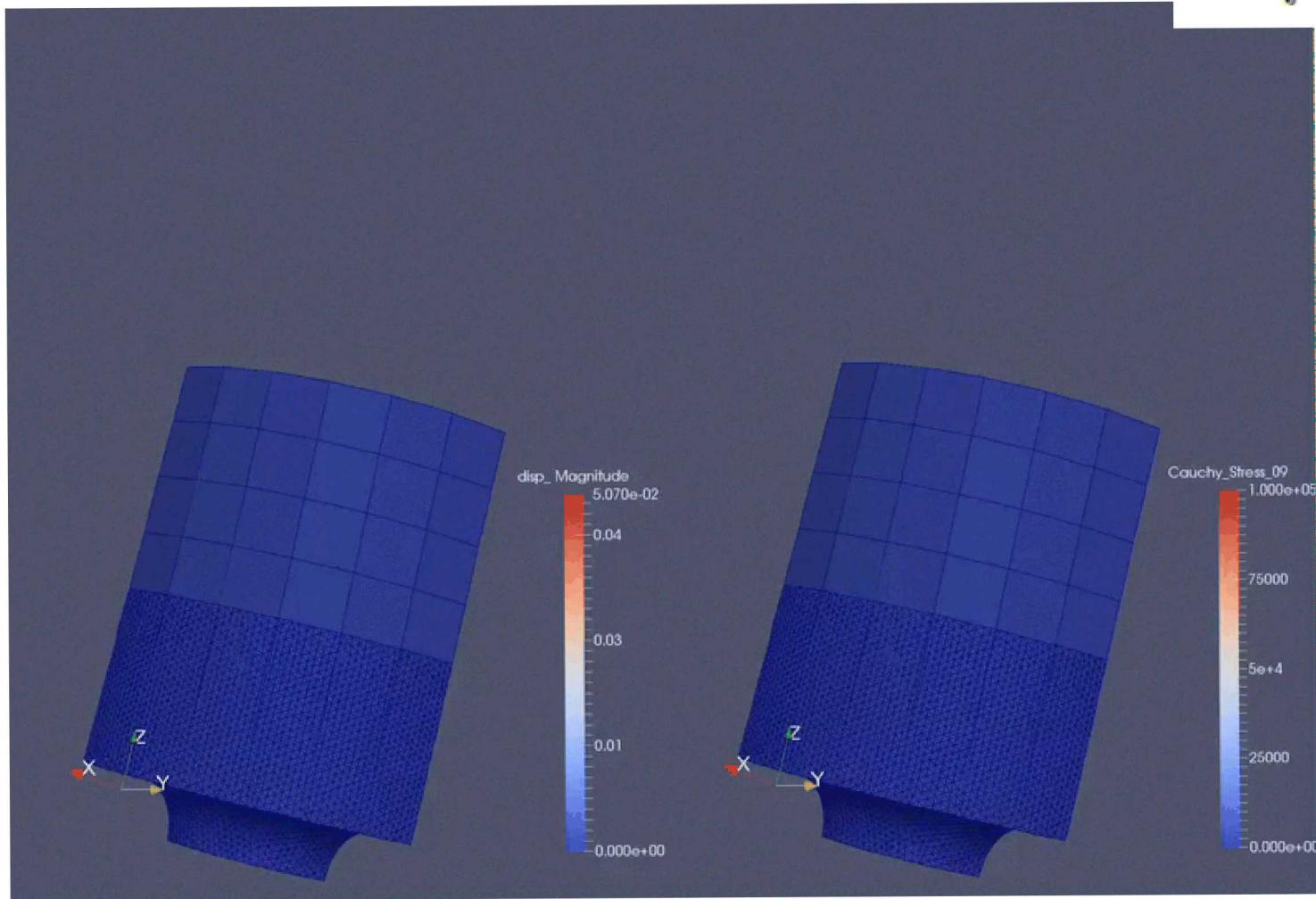
| Absolute residual tolerance | u_3 relative error | |
|-----------------------------|------------------------|------------------------|
| | Ω_1 | Ω_2 |
| 1.0×10^{-4} | 7.60×10^{-3} | 3.20×10^{-3} |
| 1.0×10^{-8} | 3.10×10^{-5} | 1.71×10^{-5} |
| 1.0×10^{-12} | 1.34×10^{-9} | 5.10×10^{-10} |
| 1.0×10^{-14} | 1.23×10^{-11} | 4.69×10^{-12} |
| 2.5×10^{-16} | 1.14×10^{-13} | 8.37×10^{-14} |



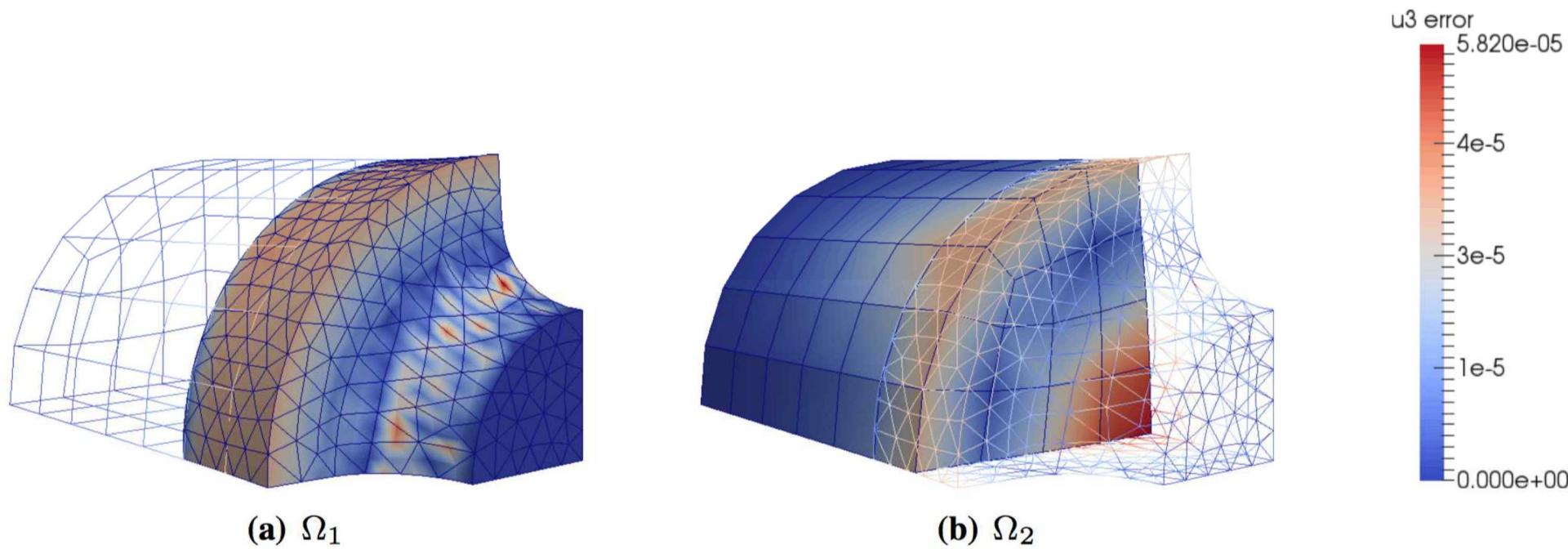
Notched Cylinder: TET-HEX Coupling

- The Schwarz alternating method is capable of coupling ***different mesh topologies***.
- The notched region, where stress concentrations are expected, is ***finely*** meshed with ***tetrahedral*** elements.
- The top and bottom regions, presumably of less interest, are meshed with ***coarser hexahedral*** elements.





Notched Cylinder: Conformal TET-HEX Coupling



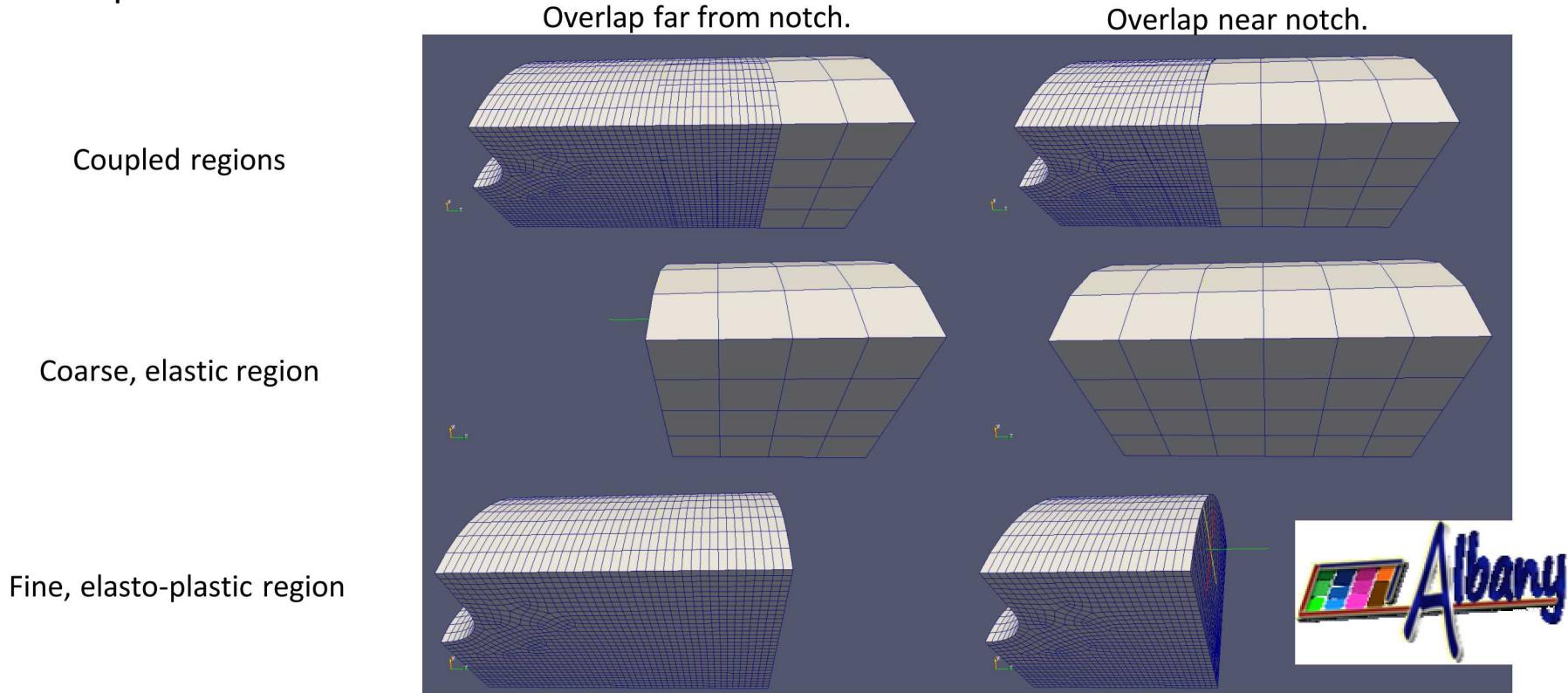
| Absolute residual tolerance | u_3 relative error | |
|-----------------------------|-----------------------|-----------------------|
| | Ω_1 | Ω_2 |
| 1.0×10^{-14} | 9.27×10^{-3} | 3.70×10^{-3} |



Notched Cylinder: Coupling Different Materials

The Schwarz method is capable of coupling regions with *different material models*.

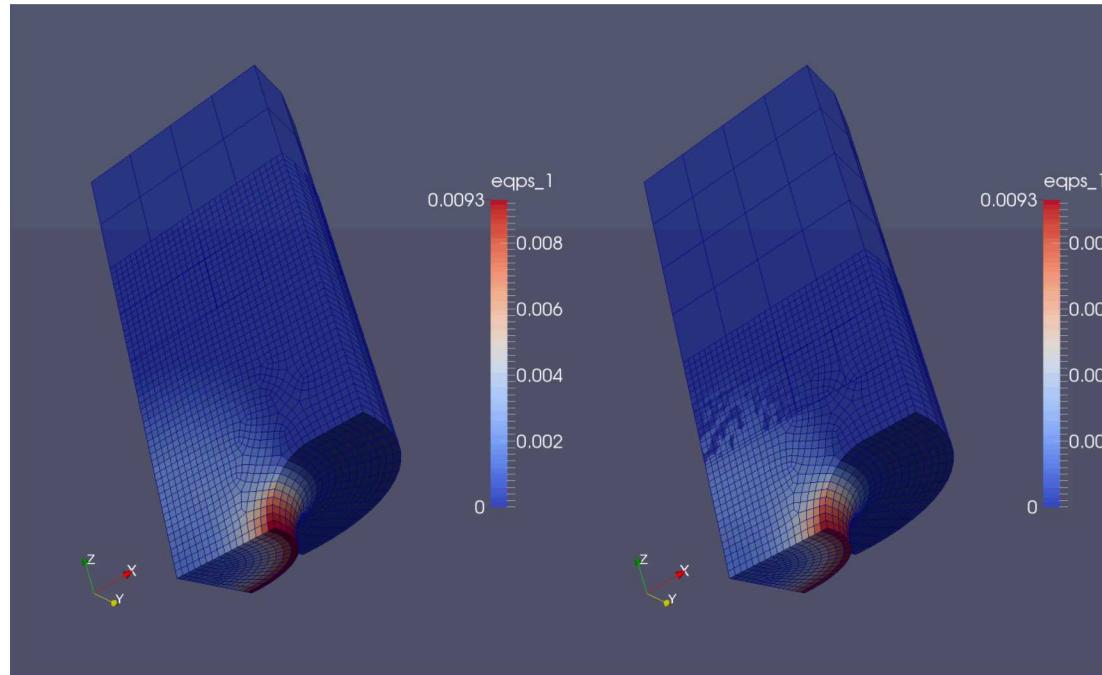
- Notched cylinder subjected to tensile load with an **elastic** and ***J*2 elasto-plastic** regions.
- **Coarse** region is **elastic** and **fine** region is **elasto-plastic**.
- The **overlap region** in the first mesh is nearer the notch, where plastic behavior is expected.



Notched Cylinder: Coupling Different Materials

Need to be careful to do domain decomposition so that material models are *consistent* in overlap region.

- When the *overlap* region is *far from the notch*, no plastic deformation exists in it: the coarse and fine regions predict the *same behavior*.
- When the *overlap* region is *near the notch*, plastic deformation spills onto it and the two models predict different behavior, affecting convergence *adversely*.

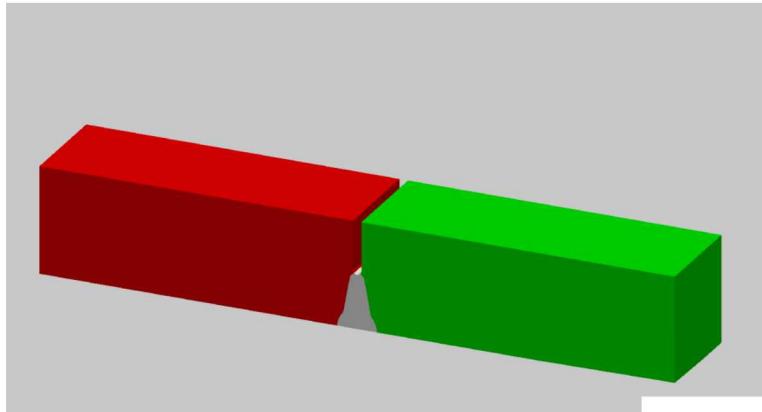


Overlap far from notch.

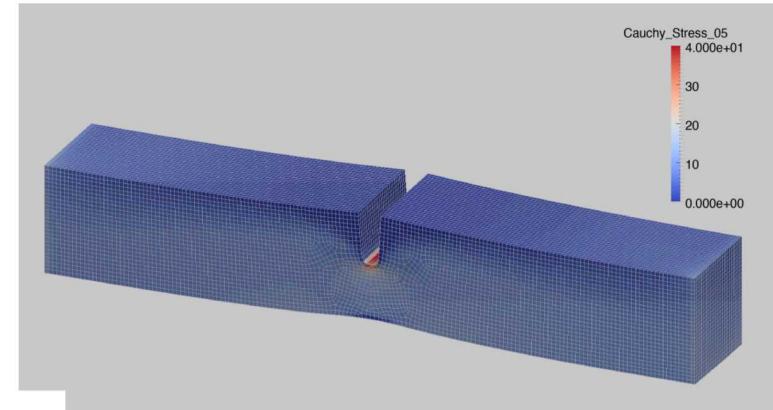
Overlap near notch.

Quasistatic Example #3: Laser Weld

Laser weld specimen

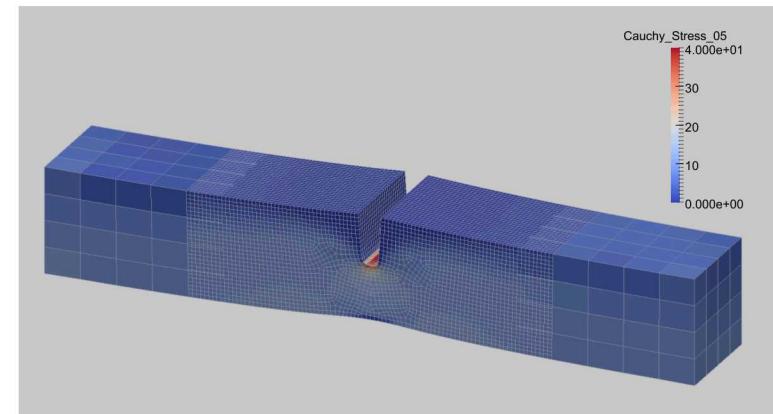


Single domain discretization

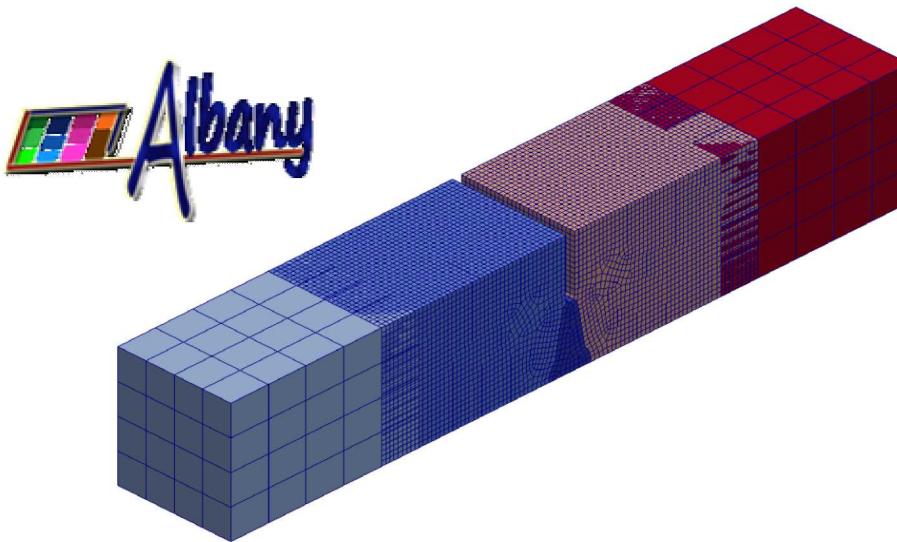


Coupled Schwarz discretization
(50% reduction in model size)

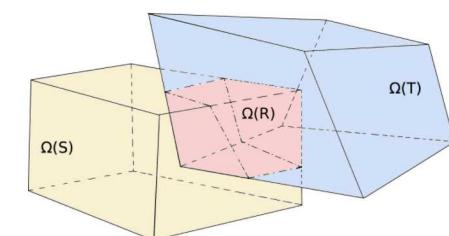
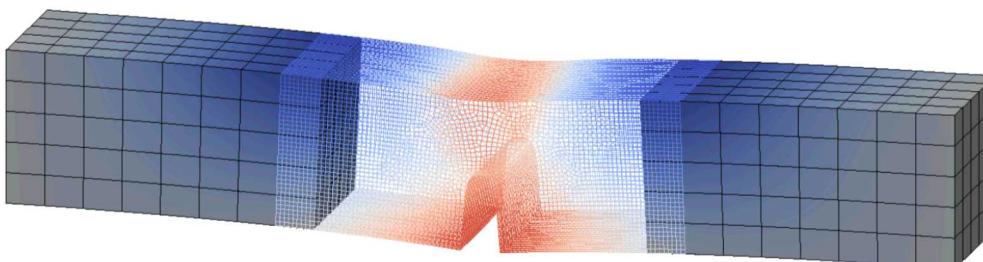
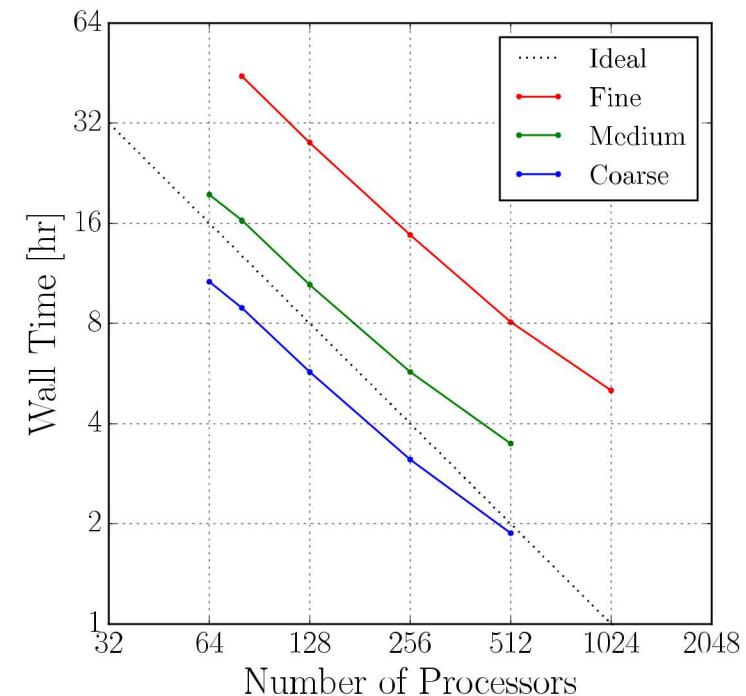
- Problem of ***practical scale*** ($\sim 200K$ *dofs*).
- ***Isotropic elasticity*** and ***J2 plasticity*** with linear isotropic hardening.
- ***Identical parameters*** for weld and base materials for proof of concept, to become independent models.



Laser Weld: Strong Scalability of Parallel Schwarz with DTK



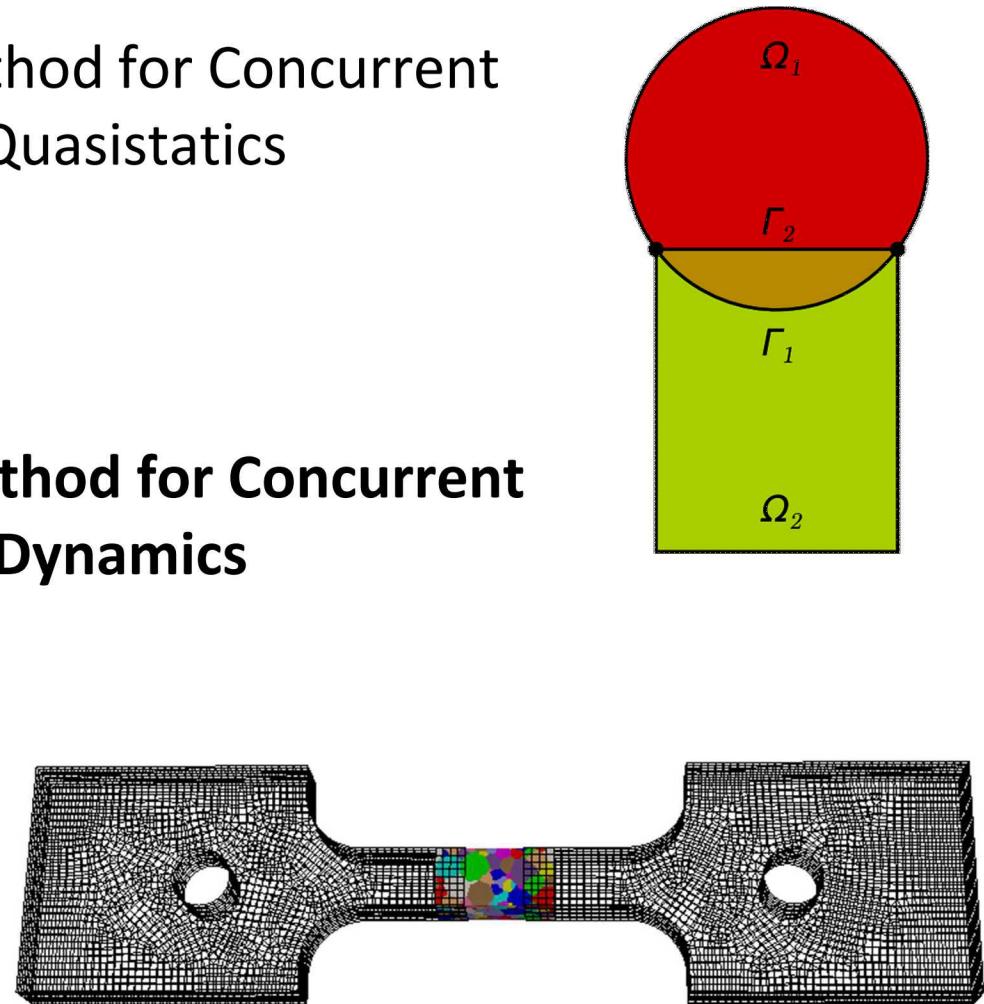
- **Near-ideal linear speedup** (64-1024 cores).



Data Transfer Kit (DTK)

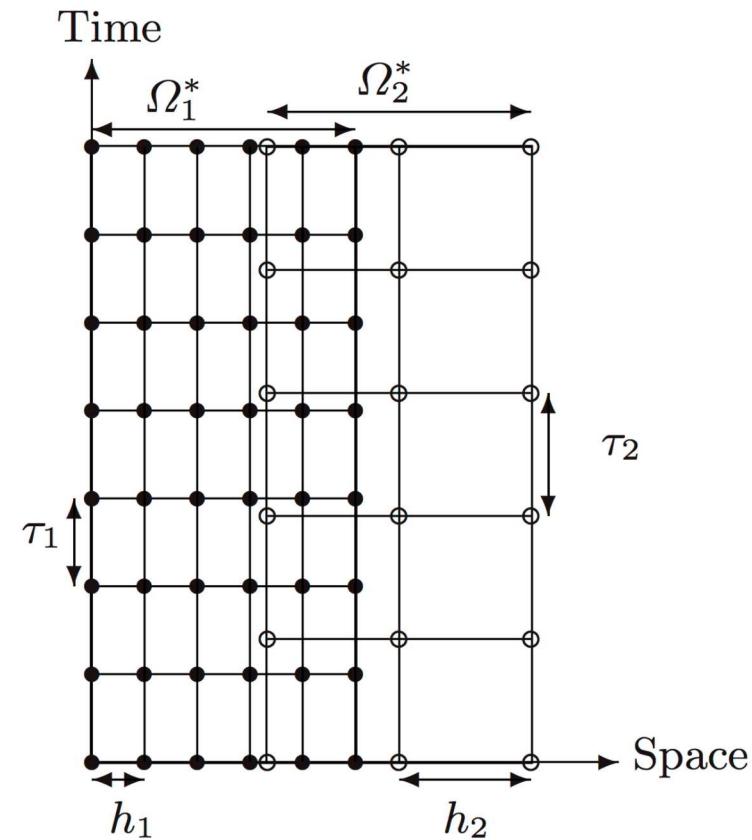
Outline

1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - Implementation
 - Numerical Examples
3. **Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics**
 - **Formulation**
 - Implementation
 - Numerical Examples
4. Summary
5. Future Work



Schwarz Alternating Method for Dynamics

- In the literature the Schwarz method is applied to dynamics by using *space-time discretizations*.



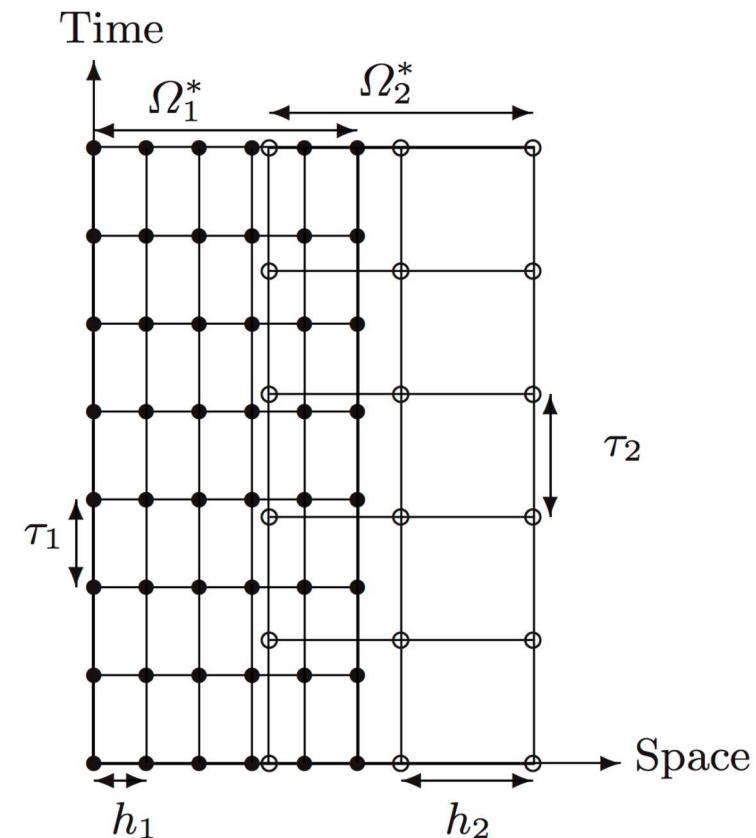
Overlapping non-matching meshes and time steps in dynamics.

Schwarz Alternating Method for Dynamics

- In the literature the Schwarz method is applied to dynamics by using ***space-time discretizations***.

Pro ☺: Can use ***non-matching*** meshes and time-steps (see right figure).

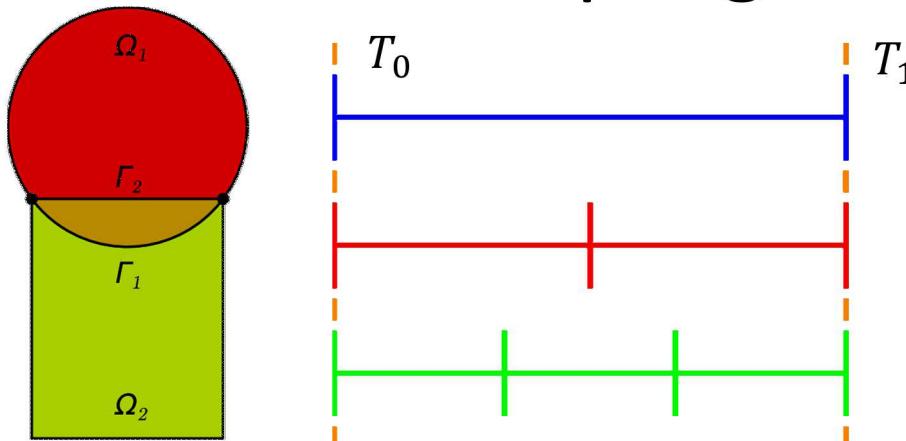
Con ☹: ***Unfeasible*** given the design of our current codes and size of simulations.



Overlapping non-matching meshes and time steps in dynamics.

Schwarz Alternating Method for Dynamic Multiscale Coupling

Controller time stepper = convenient checkpoint to facilitate implementation



Controller time stepper

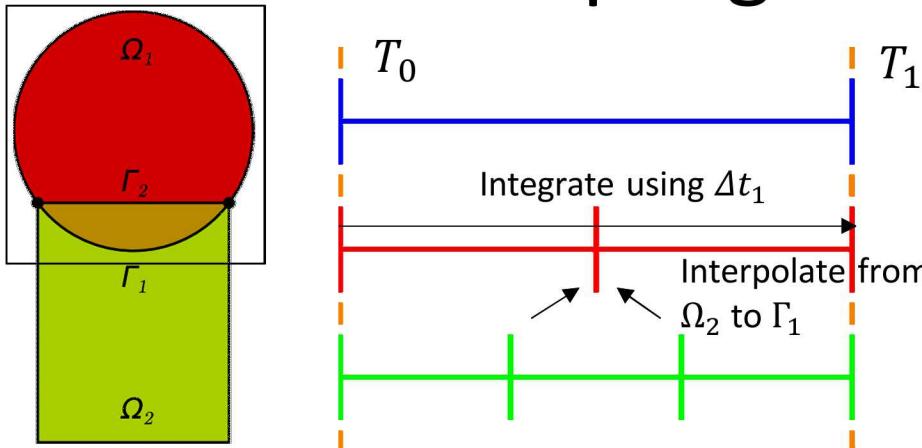
Time integrator for Ω_1

Time integrator for Ω_2

Step 0: Initialize $i = 0$ (controller time index).

Schwarz Alternating Method for Dynamic Multiscale Coupling

Controller time stepper = convenient checkpoint to facilitate implementation



Controller time stepper

Time integrator for Ω_1

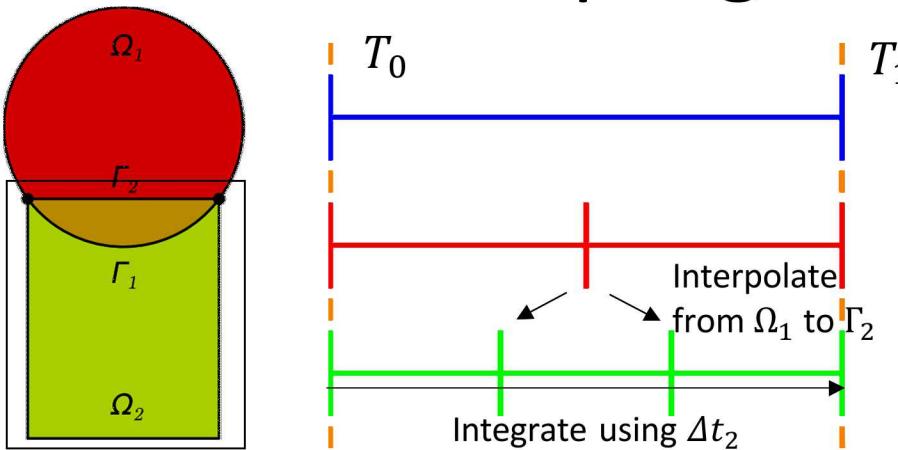
Time integrator for Ω_2

Step 0: Initialize $i = 0$ (controller time index).

Step 1: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.

Schwarz Alternating Method for Dynamic Multiscale Coupling

Controller time stepper = convenient checkpoint to facilitate implementation



Controller time stepper

Time integrator for Ω_1

Time integrator for Ω_2

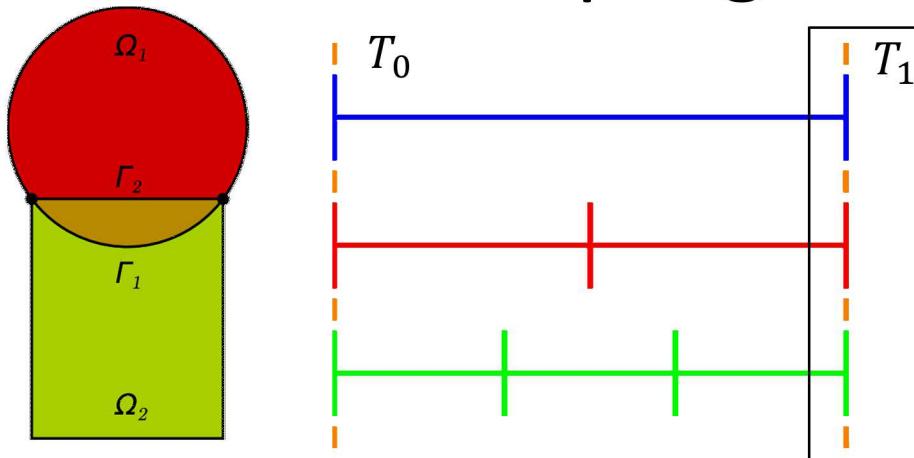
Step 0: Initialize $i = 0$ (controller time index).

Step 1: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.

Step 2: Advance Ω_2 solution from time T_i to time T_{i+1} using time-stepper in Ω_2 with time-step Δt_2 , using solution in Ω_1 interpolated to Γ_2 at times $T_i + n\Delta t_2$.

Schwarz Alternating Method for Dynamic Multiscale Coupling

Controller time stepper = convenient checkpoint to facilitate implementation



Controller time stepper

Time integrator for Ω_1

Time integrator for Ω_2

Step 0: Initialize $i = 0$ (controller time index).

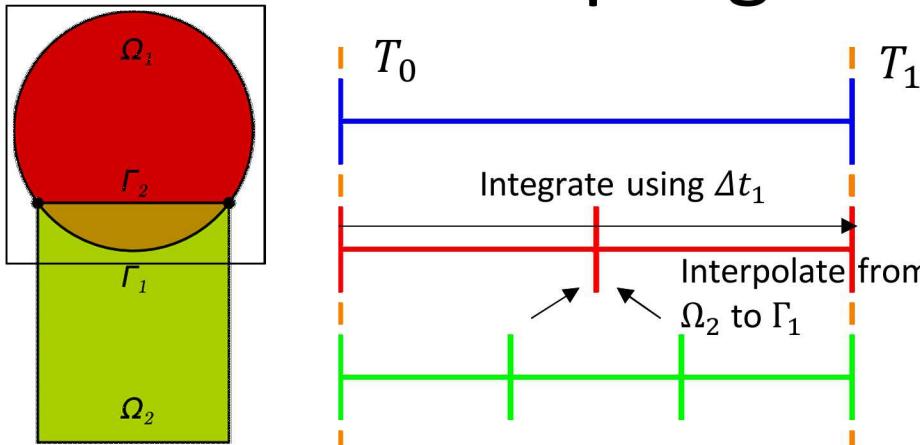
Step 1: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.

Step 2: Advance Ω_2 solution from time T_i to time T_{i+1} using time-stepper in Ω_2 with time-step Δt_2 , using solution in Ω_1 interpolated to Γ_2 at times $T_i + n\Delta t_2$.

Step 3: Check for convergence at time T_{i+1} .

Schwarz Alternating Method for Dynamic Multiscale Coupling

Controller time stepper = convenient checkpoint to facilitate implementation



Controller time stepper

Time integrator for Ω_1

Time integrator for Ω_2

Step 0: Initialize $i = 0$ (controller time index).

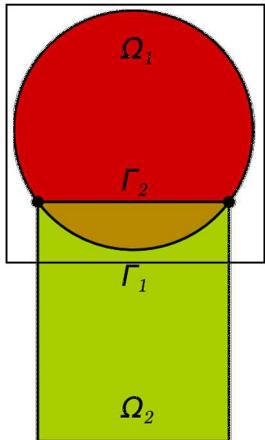
Step 1: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.

Step 2: Advance Ω_2 solution from time T_i to time T_{i+1} using time-stepper in Ω_2 with time-step Δt_2 , using solution in Ω_1 interpolated to Γ_2 at times $T_i + n\Delta t_2$.

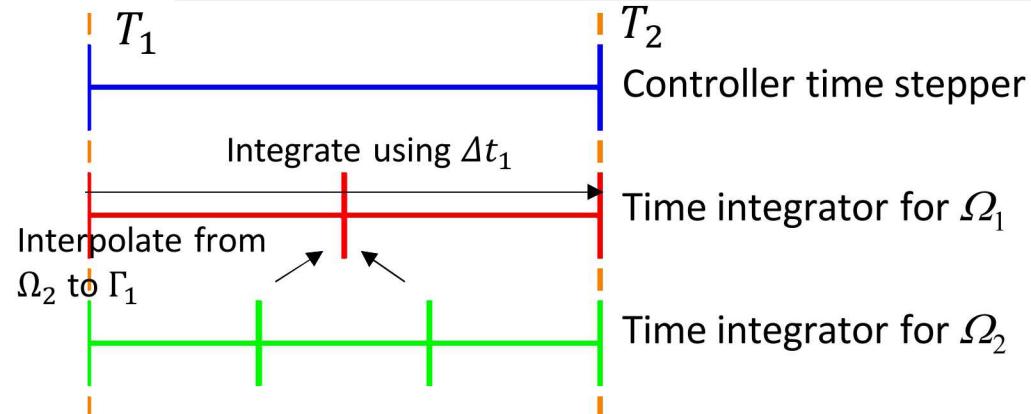
Step 3: Check for convergence at time T_{i+1} .

➤ If unconverged, return to Step 1.

Schwarz Alternating Method for Dynamic Multiscale Coupling



Controller time stepper = convenient checkpoint to facilitate implementation



Step 0: Initialize $i = 0$ (controller time index).

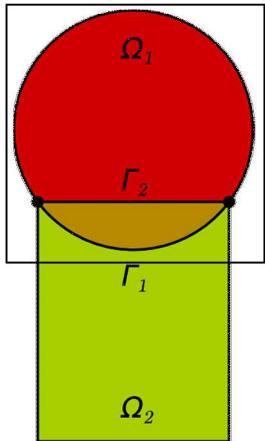
Step 1: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.

Step 2: Advance Ω_2 solution from time T_i to time T_{i+1} using time-stepper in Ω_2 with time-step Δt_2 , using solution in Ω_1 interpolated to Γ_2 at times $T_i + n\Delta t_2$.

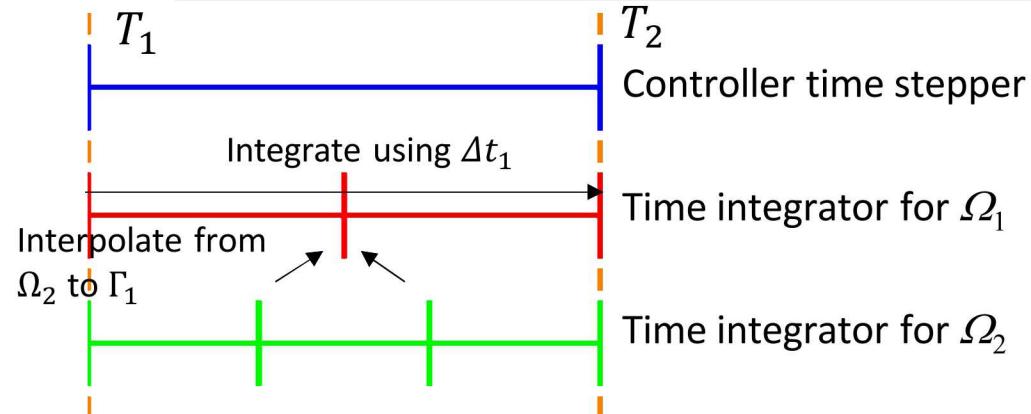
Step 3: Check for convergence at time T_{i+1} .

- If unconverged, return to Step 1.
- If converged, set $i = i + 1$ and return to Step 1.

Schwarz Alternating Method for Dynamic Multiscale Coupling



Controller time stepper = convenient checkpoint to facilitate implementation



Step 0: Initialize $i = 0$ (controller time index).

Step 1: Advance Ω_1 solution from time T_i to time T_{i+1} using time-stepper in Ω_1 with time-step Δt_1 , using solution in Ω_2 interpolated to Γ_1 at times $T_i + n\Delta t_1$.

Step 2: Advance Ω_2 solution from time T_i to time T_{i+1} using time-stepper in Ω_2 with time-step Δt_2 , using solution in Ω_1 interpolated to Γ_2 at times $T_i + n\Delta t_2$.

Step 3: Check for convergence at time T_{i+1} .

- If unconverged, return to Step 1.
- If converged, set $i = i + 1$ and return to Step 1.

Can use ***different integrators*** with ***different time steps*** within each domain!

Schwarz Alternating Method for Dynamic Multiscale Coupling: Theory

- For quasistatics, we derived a ***proof of convergence*** of the alternating Schwarz method for the ***finite deformation*** problem, and determined a ***geometric convergence rate*** [(Mota, Tezaur, Alleman, *CMAME*, 2017) and previous talk].

Theorem 1. Assume that the energy functional $\Phi[\varphi]$ satisfies properties 1–5 above. Consider the Schwarz alternating method of Section 2 defined by (9)–(13) and its equivalent form (39). Then

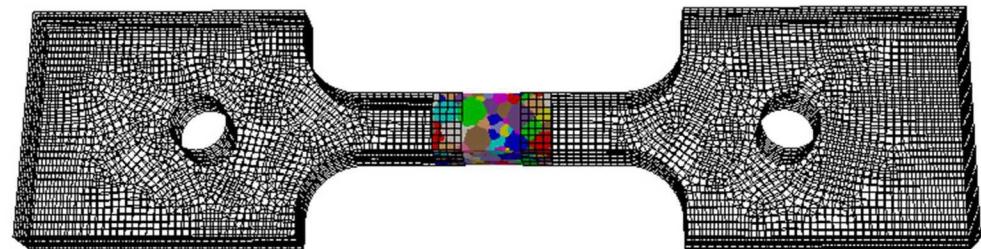
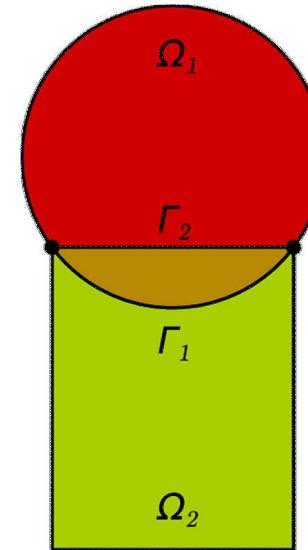
- (a) $\Phi[\tilde{\varphi}^{(0)}] \geq \Phi[\tilde{\varphi}^{(1)}] \geq \dots \geq \Phi[\tilde{\varphi}^{(n-1)}] \geq \Phi[\tilde{\varphi}^{(n)}] \geq \dots \geq \Phi[\varphi]$, where φ is the minimizer of $\Phi[\varphi]$ over \mathcal{S} .
- (b) The sequence $\{\tilde{\varphi}^{(n)}\}$ defined in (39) converges to the minimizer φ of $\Phi[\varphi]$ in \mathcal{S} .
- (c) The Schwarz minimum values $\Phi[\tilde{\varphi}^{(n)}]$ converge monotonically to the minimum value $\Phi[\varphi]$ in \mathcal{S} starting from any initial guess $\tilde{\varphi}^{(0)}$.

Extending these results to ***dynamics*** is ***work in progress***.

- Quasistatic proof ***extends naturally*** assuming conformal meshes and the same time step is used in each Schwarz subdomain.
- Some analysis of Schwarz for evolution problems was performed in (Lions, 1988) and may be possible to ***leverage***.
- Our numerical results suggest theoretical analysis is ***possible***.

Outline

1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - Implementation
 - Numerical Examples
3. Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics
 - Formulation
 - **Implementation**
 - Numerical Examples
4. Summary
5. Future Work



Implementation within *Albany* Code

The proposed ***dynamic alternating Schwarz method*** is implemented within the ***LCM project*** in Sandia's open-source parallel, C++, multi-physics, finite element code, ***Albany***.

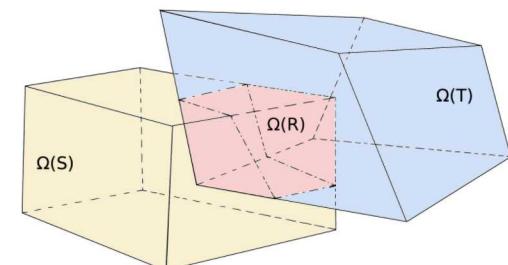


<https://github.com/gahansen/Albany>

- ***Component-based*** design for rapid development of capabilities.
- Contains a wide variety of ***constitutive models***.
- Extensive use of libraries from the open-source ***Trilinos*** project.
 - Use of the ***Phalanx*** package to decompose complex problem into simpler problems with managed dependencies.
 - Use of the ***Sacado*** package for ***automatic differentiation***.
 - Use of ***Tempus*** package for ***time-integration****.
- ***Parallel*** implementation of Schwarz alternating method uses the ***Data Transfer Kit (DTK)***.
- All software available on ***GitHub***.



<https://github.com/trilinos/trilinos>

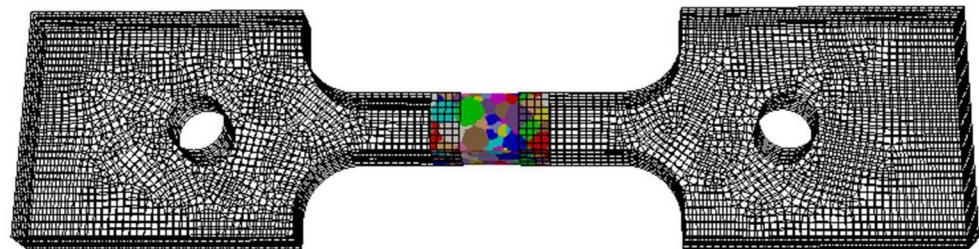
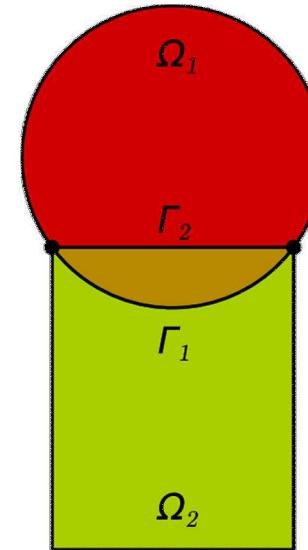


<https://github.com/ORNL-CEES/DataTransferKit>

* Current dynamic Schwarz implementation in Albany requires same Δt in different subdomains.

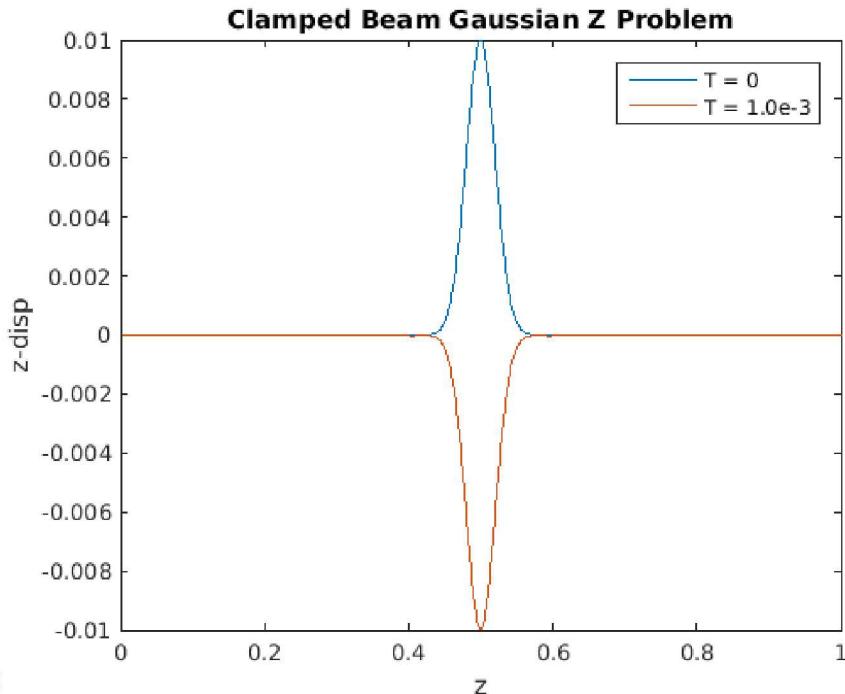
Outline

1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - Implementation
 - Numerical Examples
3. Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics
 - Formulation
 - Implementation
 - Numerical Examples
4. Summary
5. Future Work



Dynamic Example #1: Elastic Wave Propagation

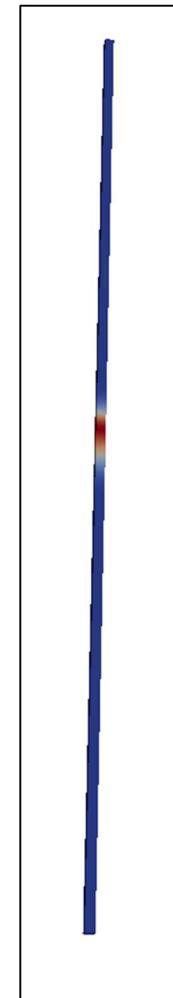
- Linear elastic **clamped beam** with Gaussian initial condition for the z-displacement (see figures to the right and below).
- Simple problem with analytical exact solution but very **stringent test** for discretization methods.
- Test Schwarz with **2 subdomains**: $\Omega_0 = (0,0.001) \times (0.001) \times (0,0.75)$, $\Omega_1 = (0,0.001) \times (0.001) \times (0.25,1)$.



Left: Initial condition (blue) and final solution (red). Wave profile is negative of initial profile at time $T = 1.0e-3$.

Time-discretizations:
Newmark-Beta (implicit, explicit) with same Δt .

Meshes: hexes, tets



Elastic Wave Propagation

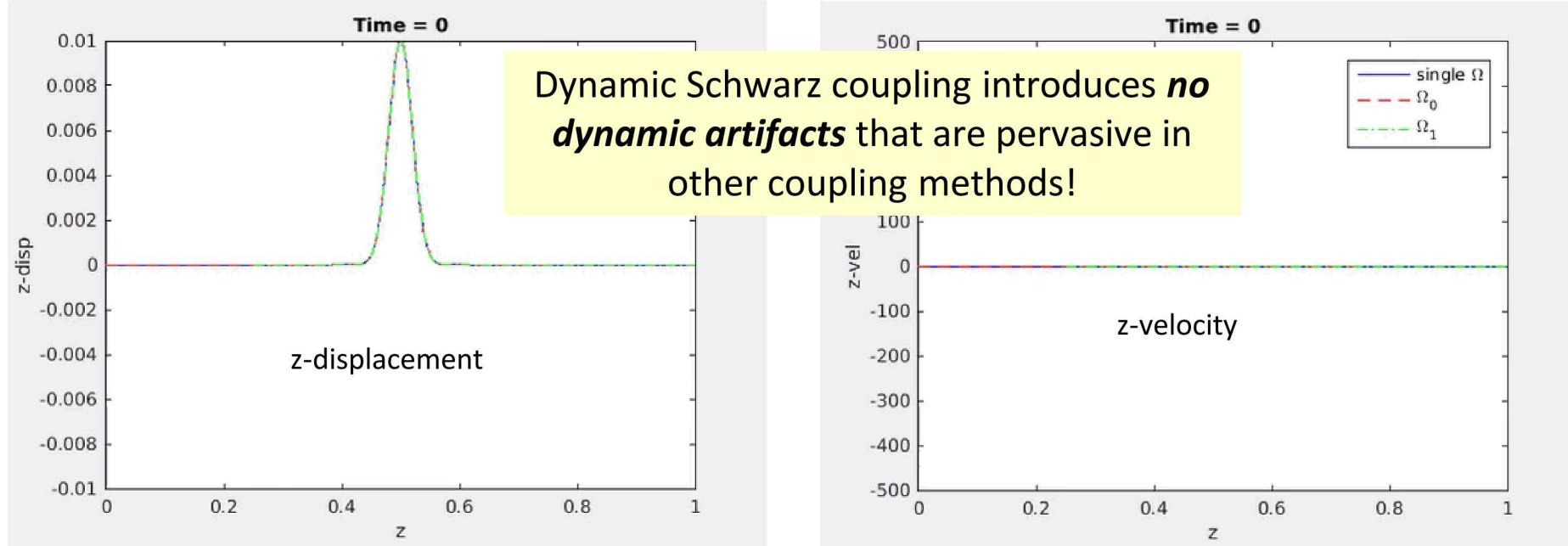
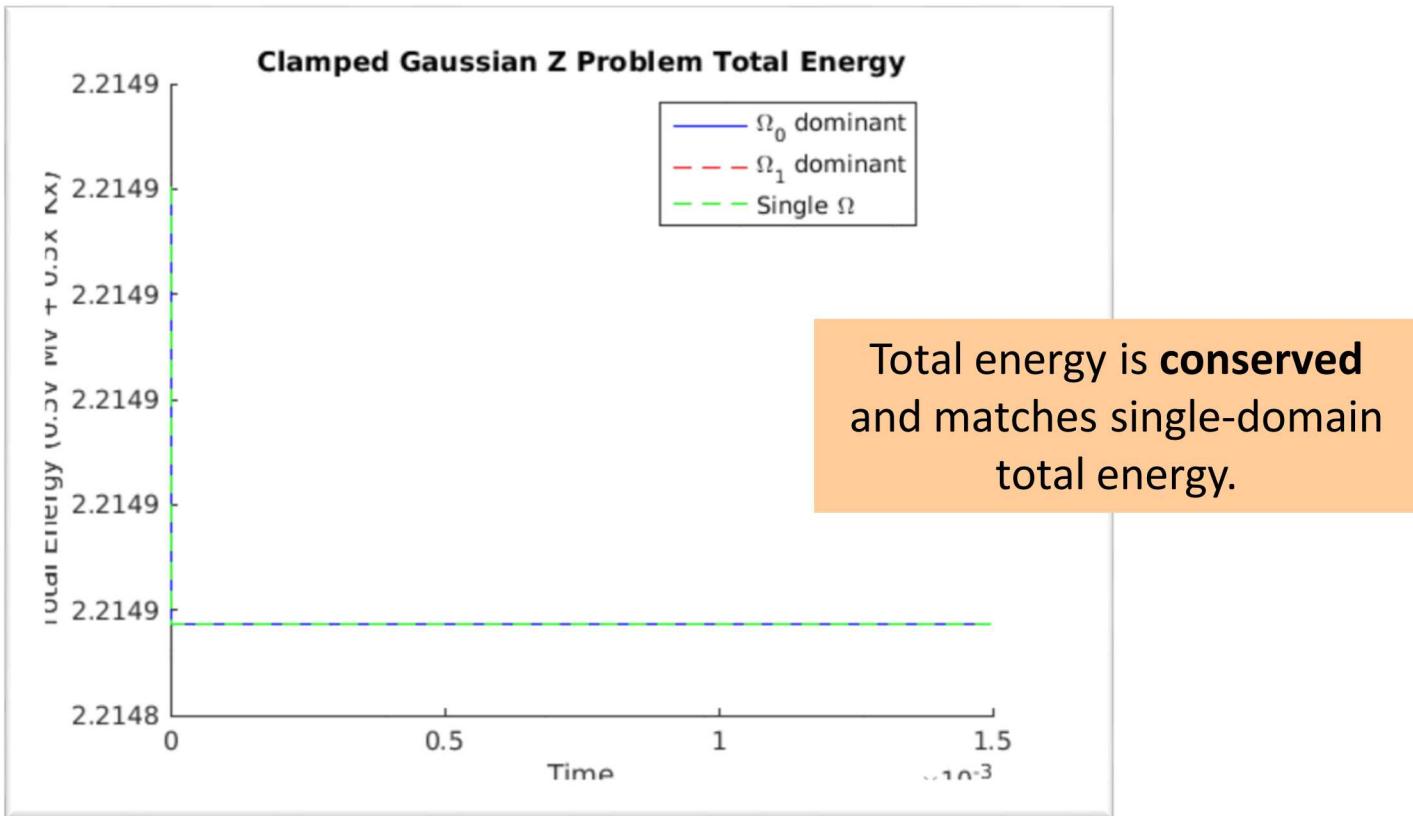


Table 1: Averaged (over times + domains) relative errors in **z-displacement** (blue) and **z-velocity** (green) for several different Schwarz couplings, 50%

| | Implicit-Implicit | | Explicit(CM)-Implicit | | Explicit(LM)-Implicit | |
|-----------------------------|-------------------|---------|-----------------------|---------|-----------------------|---------|
| Conformal hex-hex | 2.79e-3 | 7.32e-3 | 3.53e-3 | 8.70e-3 | 4.72e-3 | 1.19e-2 |
| Nonconformal hex-hex | 2.90e-3 | 7.10e-3 | 2.82e-3 | 7.29e-3 | 2.84e-3 | 7.33e-3 |
| Tet-hex | 2.79e-3 | 7.58e-3 | 3.52e-3 | 8.92e-3 | 4.72e-3 | 1.19e-2 |

Elastic Wave Propagation

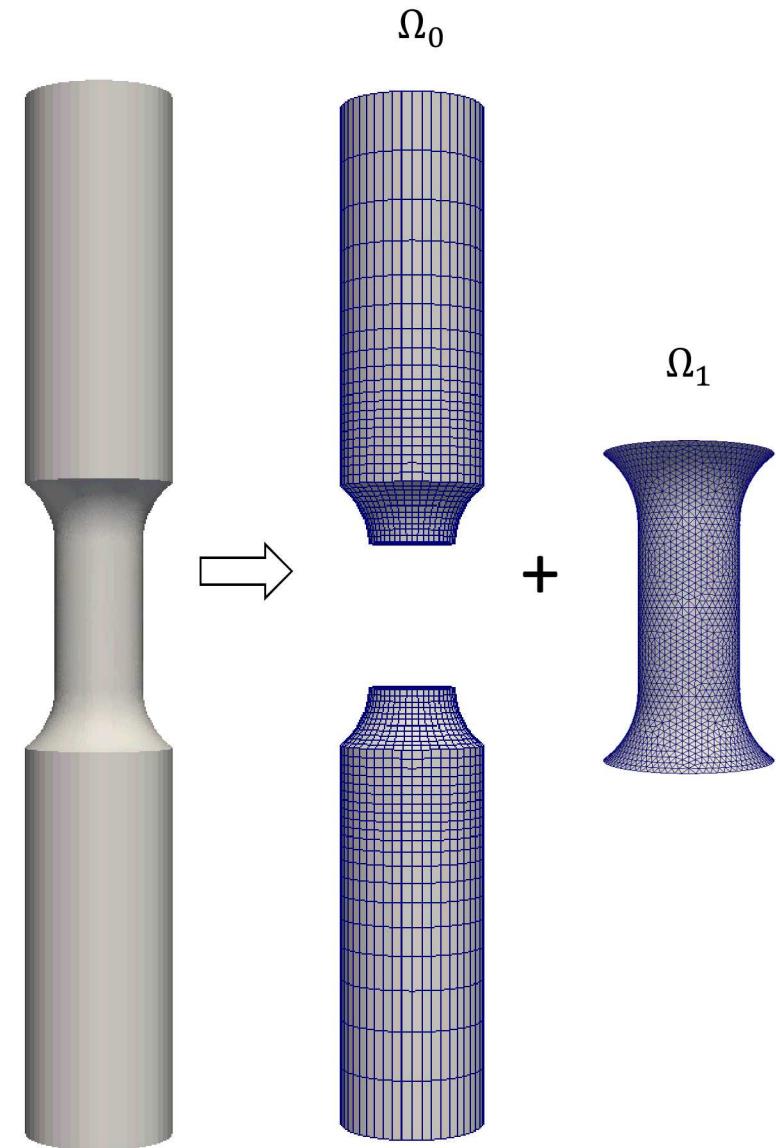
Energy Conservation



- For clamped beam problem, total energy ($TE = 0.5x^T Kx + 0.5\dot{x}^T M\dot{x}$) should be conserved.
- Total energy is calculated in 2 ways: with most of contribution from Ω_0 and from Ω_1 .

Example #2: Tension Specimen

- Uniaxial aluminum cylindrical tensile specimen with ***inelastic J₂ material model.***
- Domain decomposition into ***two subdomains*** (right): Ω_0 = ends, Ω_1 = gauge.
- ***Nonconformal hex + composite tet 10*** coupling via Schwarz.
- ***Implicit*** Newmark time-integration with ***adaptive time-stepping*** algorithm employed in both subdomains.
- Slight ***imperfection*** introduced at center of gauge to force ***necking*** upon pulling in vertical direction.

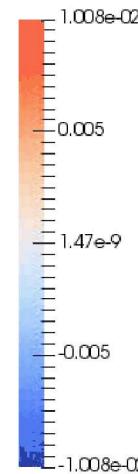


Tension Specimen

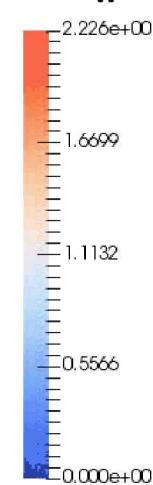
Time: 0.000000



y-displacement



Nodal eqps*



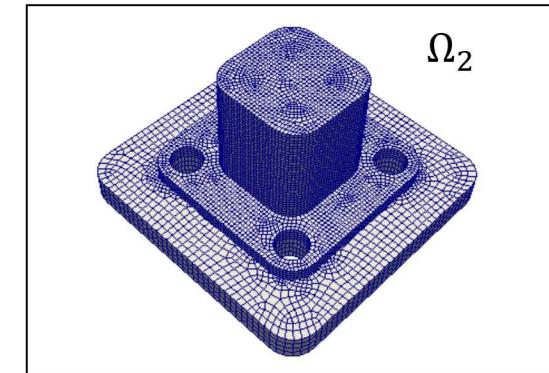
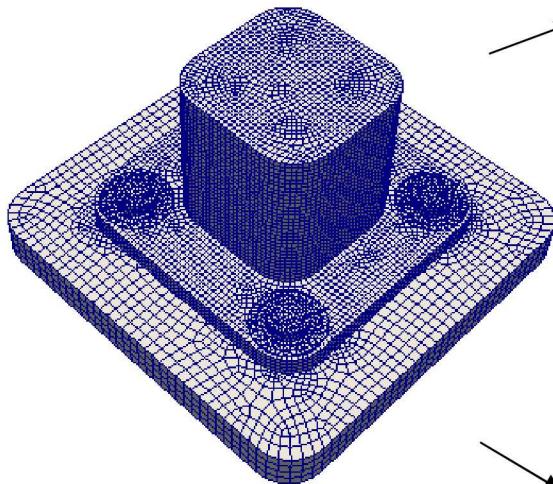
Average of ~7 Schwarz iterations/time step required for ***convergence*** to Schwarz tolerance of 1e-6.

*Nodal eqps = equivalent plastic strain computed via weighted volume average.

Example #3: Bolted Joint Problem

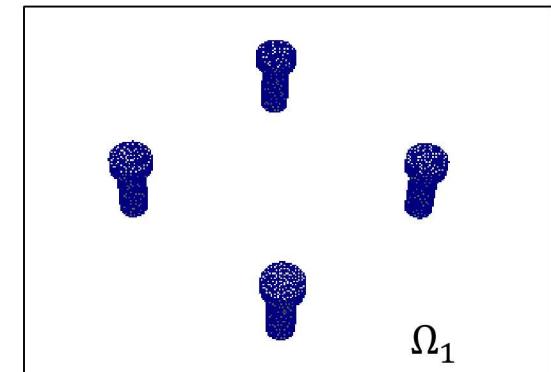
Problem of *practical scale*.

- Schwarz solution compared to single-domain solution on composite tet 10 mesh.



- BC: $x\text{-disp} = 0.02$ at $T = 1.0\text{e-}3$ on top of parts.
- Run until $T = 5.0\text{e-}4$ w/ $dt = 1\text{e-}5$ + implicit Newmark with analytic mass matrix for composite tet 10s.

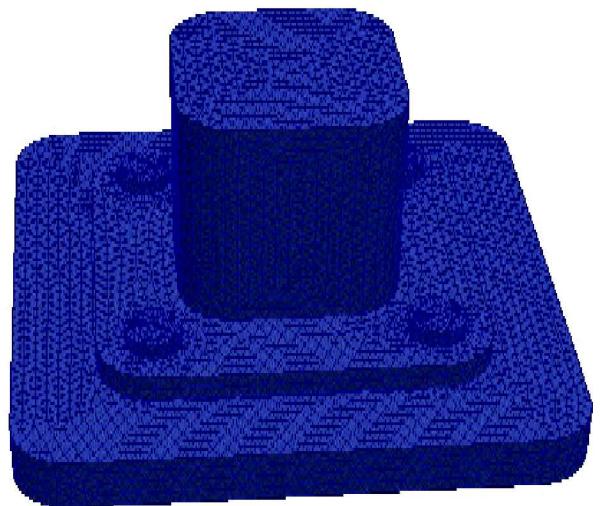
- Ω_1 = bolts (composite tet 10), Ω_2 = parts (hex).
- ***Inelastic J₂ material model*** in both subdomains.
 - Ω_1 : steel
 - Ω_2 : steel component, aluminum (bottom) plate



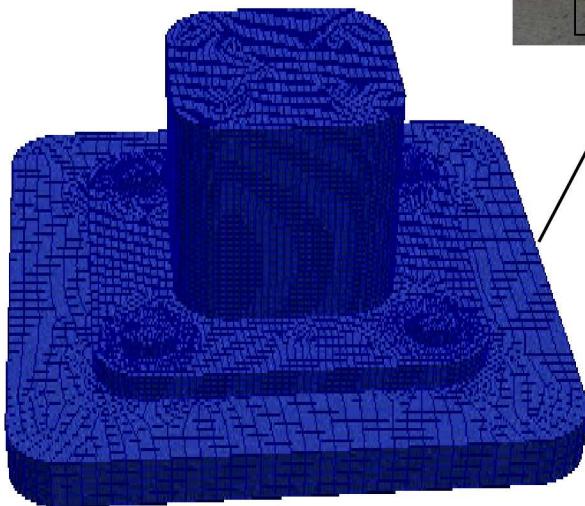
Bolted Joint Problem

Time: 0.000000

x-displacement



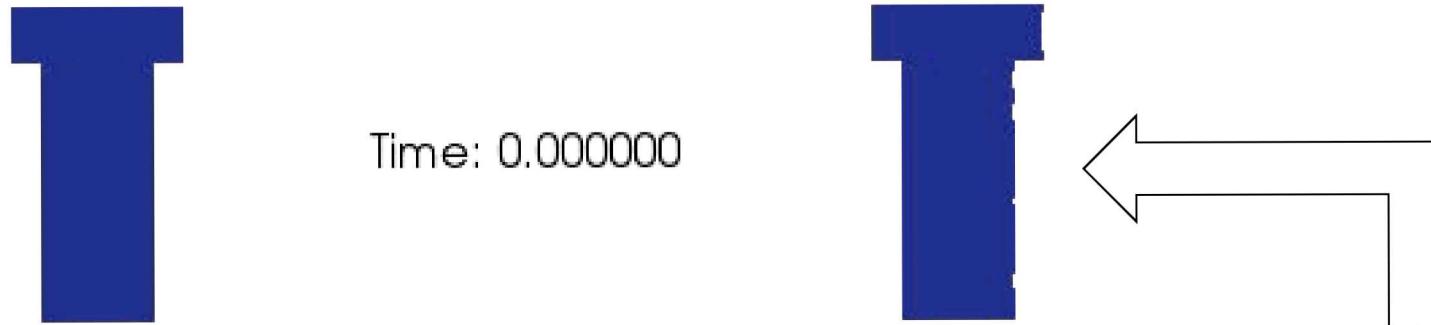
Single Ω



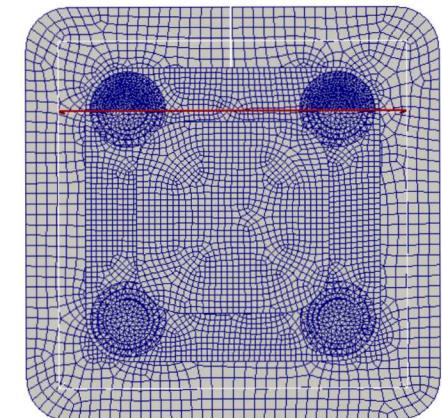
Schwarz

Bolted Joint Problem

Nodal Equivalent Plastic Strain (eqps)



Cross-section of bolts obtained via clip (right)

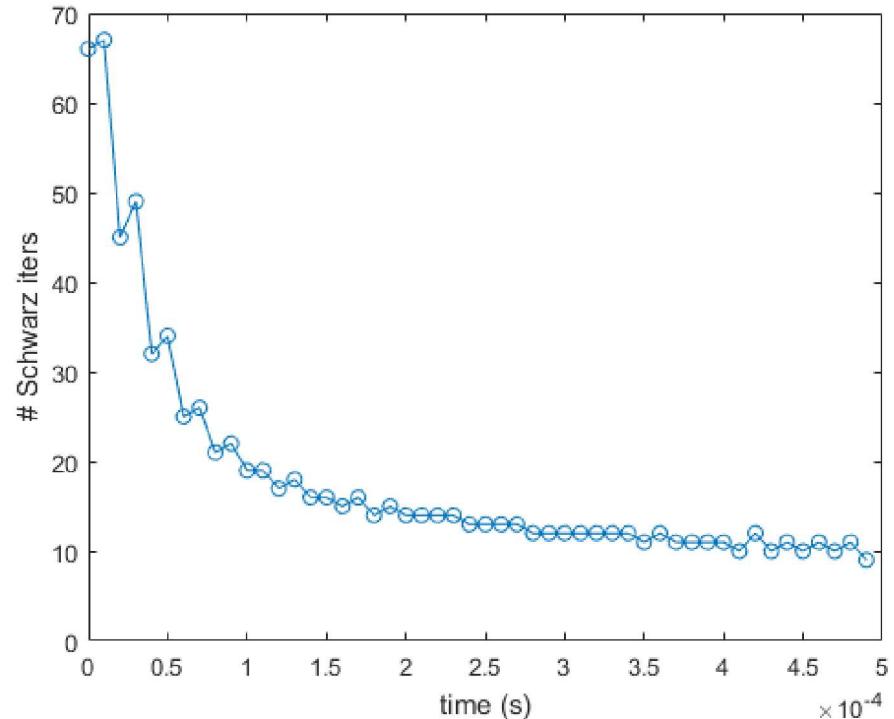


Bolted Joint Problem

Some Performance Results

Schwarz / solver settings

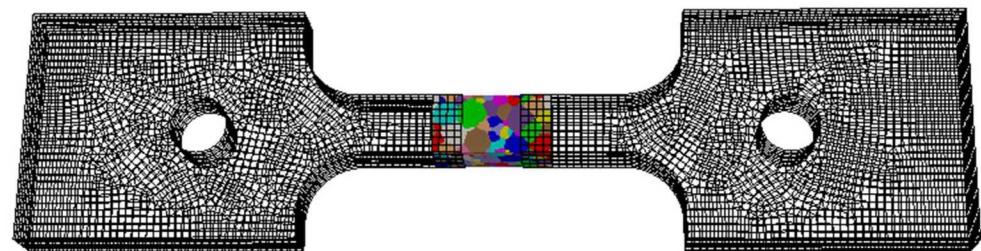
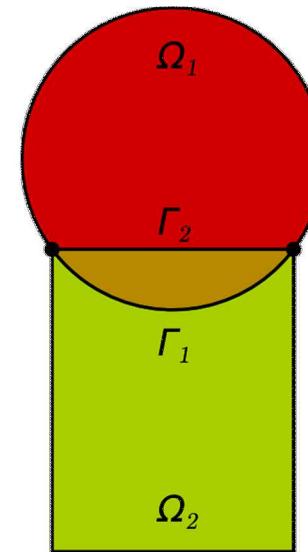
- Relatively loose Schwarz tolerances were used:
 - Relative Tolerance: $1.0e-3$.
 - Absolute Tolerance: $1.0e-4$.
- Newton tolerance on NormF: $1e-8$
- Linear solver tolerance: $1e-5$
- MueLu preconditioner



- ***Top right plot:*** # Schwarz iterations for each time step.
 - After start-up, # Schwarz iterations / time step is $\sim 9-10$. This is not bad given how small is the size of the overlap region for this problem.

Outline

1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - Implementation
 - Numerical Examples
3. Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics
 - Formulation
 - Implementation
 - Numerical Examples
4. Summary
5. Future Work



Summary

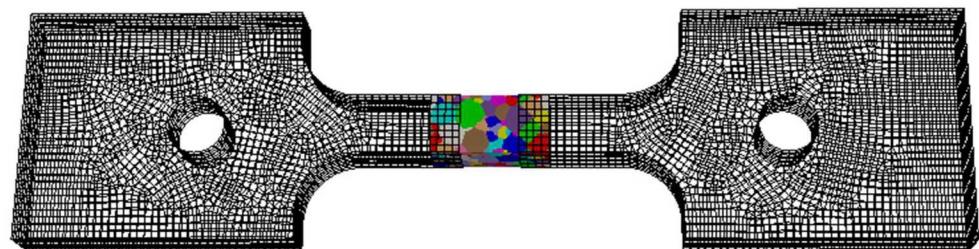
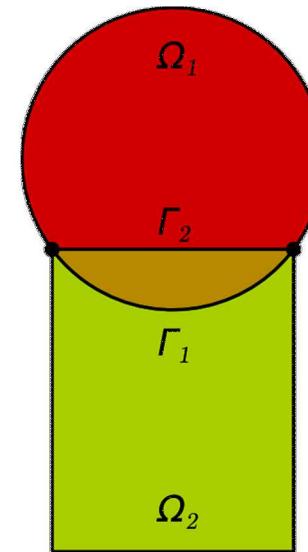
The alternating Schwarz coupling method has been developed/implemented for concurrent multiscale quasistatic & dynamic modeling in Sandia's Albany/LCM code.

- ☺ Coupling is ***concurrent*** (two-way).
- ☺ ***Ease of implementation*** into existing massively-parallel HPC codes.
- ☺ ***Scalable, fast, robust*** (we target ***real*** engineering problems, e.g., analyses involving failure of bolted components!).
- ☺ ***“Plug-and-play” framework***: simplifies task of meshing complex geometries!
- ☺ Ability to couple regions with ***different non-conformal meshes, different element types*** and ***different levels of refinement***.
- ☺ Ability to use ***different solvers/time-integrators*** in different regions.
- ☺ Coupling does not introduce ***nonphysical artifacts***.
- ☺ ***Theoretical*** convergence properties/guarantees (☺ for quasistatics).



Outline

1. Motivation
2. Schwarz Alternating Method for Concurrent Multiscale Coupling for Quasistatics
 - Formulation
 - Implementation
 - Numerical Examples
3. Schwarz Alternating Method for Concurrent Multiscale Coupling for Dynamics
 - Formulation
 - Implementation
 - Numerical Examples
4. Summary
5. Future Work



Ongoing/Future Work

- Development of ***theory*** for dynamic alternating Schwarz formulation.
- ***Journal article*** on our dynamic Schwarz formulation is in preparation.
- Extension of Albany/LCM dynamic Schwarz implementation to allow for ***different time steps*** in different subdomains.
- Application of dynamic Schwarz for problems and test cases of interest to ***production***.
- Implementation of alternating Schwarz method for concurrent multiscale coupling in Sandia ***production codes*** (Sierra Solid Mechanics), comparison to other methods (e.g., GFEM).
- Development of a ***multi-physics coupling framework*** based on variational formulations and the Schwarz alternating method.



References

- [1] M.A. Heroux *et al.* "An overview of the Trilinos project." *ACM Trans. Math. Softw.* 31(3) (2005).
- [2] A. Salinger, *et al.* "Albany: Using Agile Components to Develop a Flexible, Generic Multiphysics Analysis Code", *Int. J. Multiscale Comput. Engng.* 14(4) (2016) 415-438.
- [3] H. Schwarz. "Über einen Grenzübergang durch alternierendes Verfahren". In: Vierteljahrsschrift der Naturforschenden Gesellschaft in Zurich 15 (1870), pp. 272–286.
- [4] S.L. Sobolev. "Schwarz's Algorithm in Elasticity Theory". In: Selected Works of S.L Sobolev. Volume I: equations of mathematical physics, computational mathematics and cubature formulats. Ed. By G.V. Demidenko and V.L. Vaskevich. New York: Springer, 2006.
- [5] S. Mikhlin. "On the Schwarz algorithm". In: Proceedings of the USSR Academy of Sciences (in Russian) 77 (1951), pp. 569–571.
- [6] D.J. Evans *et al.* "The convergence rate of the Schwarz alternating procedure (II): For two-dimensional problems". In: *International Journal for Computer Mathematics* 20.3–4 (1986), pp. 325–339.
- [7] P.L. Lions. "On the Schwarz alternating method I." In: 1988, First International Symposium on Domain Decomposition methods for Partial Differential Equations, SIAM, Philadelphia.
- [8] A. Mota, I. Tezaur, C. Alleman. "The Schwarz Alternating Method in Solid Mechanics", *Comput. Meth. Appl. Mech. Engng.* 319 (2017), 19-51. [<http://www.sandia.gov/~ikalash/journal.html>].
- [9] A. Mota, I. Tezaur, G. Phlipot. "The Schwarz alternating method for dynamic solid mechanics", in prep.

Appendix. Previous Work

Comput Mech (2014) 54:803–820
DOI 10.1007/s00466-014-1034-0

ORIGINAL PAPER

A multiscale overlapped coupling formulation for large-deformation strain localization

WaiChing Sun · Alejandro Mota

Received: 18 September 2013 / Accepted: 7 April 2014 / Published online: 3 May 2014
© Springer-Verlag Berlin Heidelberg 2014

Abstract We generalize the multiscale overlapped domain framework to couple multiple rate-independent standard dissipative material models in the finite deformation regime across different length scales. We show that a fully coupled multiscale incremental boundary-value problem can be recast as the stationary point that optimizes the partitioned incremental work of a three-field energy functional. We also establish inf-sup tests to examine the numerical stability issues that arise from enforcing weak compatibility in the three-field formulation. We also devise a new block solver for the domain coupling problem and demonstrate the performance of the formulation with one-dimensional numerical examples. These simulations indicate that it is sufficient to introduce a localization limiter in a confined region of interest to regularize the partial differential equation if loss of ellipticity occurs.

strain localization may lead to the eventual failure of materials, this phenomenon is of significant importance to modern engineering applications.

The objective of this work is to introduce concurrent coupling between sub-scale and macro-scale simulations for inelastic materials that are prone to strain localization. Since it is not feasible to conduct sub-scale simulations on macroscopic problems, we use the domain coupling method such that computational resources can be efficiently allocated to regions of interest [14, 23, 24, 30]. To the best of our knowledge, this is the first work focusing on utilizing the domain coupling method to model strain localization in inelastic materials undergoing large deformation.

Nevertheless, modeling strain localization with the conventional finite element method may lead to spurious mesh-dependent results due to the loss of ellipticity at the onset of strain localization [31]. To circumvent the loss of mate-

Appendix. Previous Work

Comput Mech (2014) 54:803–820
DOI 10.1007/s00466-014-1034-0

ORIGINAL PAPER

A multiscale overlapped coupling formulation for large-deformation strain localization

WaiChing Sun · Alejandro Mota

Received: 18 September 2013 / Accepted: 7 April 2014 / Published online: 3 May 2014
© Springer-Verlag Berlin Heidelberg 2014

Abstract We generalize the multiscale overlapped domain framework to couple multiple rate-independent standard dissipative material models in the finite deformation regime across different length scales. We show that a fully coupled multiscale incremental boundary-value problem can be recast as the stationary point that optimizes the partitioned incremental work of a three-field energy functional. We also establish inf-sup tests to examine the numerical stability issues that arise from enforcing weak compatibility in the three-field formulation. We also devise a new block solver for the domain coupling problem and demonstrate the performance of the formulation with one-dimensional numerical examples. These simulations indicate that it is sufficient to introduce a localization limiter in a confined region of interest to regularize the partial differential equation if loss of ellipticity occurs.

Three-field multiscale coupling formulation with compatibility enforced weakly using ***Lagrange multipliers***.

strain localization may lead to the eventual failure of materials, this phenomenon is of significant importance to modern engineering applications.

The objective of this work is to introduce concurrent coupling between sub-scale and macro-scale simulations for inelastic materials that are prone to strain localization. Since it is not feasible to conduct sub-scale simulations on macroscopic problems, we use the domain coupling method such that computational resources can be efficiently allocated to regions of interest [14, 23, 24, 30]. To the best of our knowledge, this is the first work focusing on utilizing the domain coupling method to model strain localization in inelastic materials undergoing large deformation.

Nevertheless, modeling strain localization with the conventional finite element method may lead to spurious mesh-dependent results due to the loss of ellipticity at the onset of strain localization [31]. To circumvent the loss of mate-

Appendix. Previous Work

Comput Mech (2014) 54:803–820
DOI 10.1007/s00466-014-1034-0

ORIGINAL PAPER

A multiscale overlapped coupling formulation for large-deformation strain localization

WaiChing Sun · Alejandro Mota

Received: 18 September 2013 / Accepted: 7 April 2014 / Published online: 3 May 2014
© Springer-Verlag Berlin Heidelberg 2014

Abstract We generalize the multiscale overlapped domain framework to couple multiple rate-independent standard dissipative material models in the finite deformation regime across different length scales. We show that a fully coupled multiscale incremental boundary-value problem can be recast as the stationary point that optimizes the partitioned incremental work of a three-field energy functional. We also establish inf-sup tests to examine the numerical stability issues that arise from enforcing weak compatibility in the three-field formulation. We also devise a new block solver for the domain coupling problem and demonstrate the performance of the formulation with one-dimensional numerical examples. These simulations indicate that it is sufficient to introduce a localization limiter in a confined region of interest to regularize the partial differential equation if loss of ellipticity occurs.

Method works well, but is **difficult to implement** into existing codes.

strain localization may lead to the eventual failure of materials, this phenomenon is of significant importance to modern engineering applications.

The objective of this work is to introduce concurrent coupling between sub-scale and macro-scale simulations for inelastic materials that are prone to strain localization. Since it is not feasible to conduct sub-scale simulations on macroscopic problems, we use the domain coupling method such that computational resources can be efficiently allocated to regions of interest [14, 23, 24, 30]. To the best of our knowledge, this is the first work focusing on utilizing the domain coupling method to model strain localization in inelastic materials undergoing large deformation.

Nevertheless, modeling strain localization with the conventional finite element method may lead to spurious mesh-dependent results due to the loss of ellipticity at the onset of strain localization [31]. To circumvent the loss of mate-

Appendix. Full Schwarz Method

Classical algorithm originally proposed by Schwarz with **outer Schwarz loop** and **inner Newton loop**, each converged to a **tight tolerance** ($\epsilon_{\text{machine}}$).

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$                                 ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$                                 ▷ initialize for  $\Omega_2$ 
3: repeat                                                                                                         ▷ Schwarz loop
4:    $\mathbf{y}^{(1)} \leftarrow \mathbf{x}_B^{(1)}$                                          ▷ for convergence check
5:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$  ▷ project from  $\Omega_2$  to  $\Gamma_1$ 
6:   repeat                                                                                                         ▷ Newton loop for  $\Omega_1$ 
7:      $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$  ▷ linear system
8:      $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$ 
9:   until  $\|\Delta\mathbf{x}_B^{(1)}\| / \|\mathbf{x}_B^{(1)}\| \leq \epsilon_{\text{machine}}$                                          ▷ tight tolerance
10:   $\mathbf{y}^{(2)} \leftarrow \mathbf{x}_B^{(2)}$                                          ▷ for convergence check
11:   $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{P}_{21}\mathbf{x}_B^{(1)} + \mathbf{Q}_{21}\mathbf{x}_b^{(1)} + \mathbf{G}_{21}\mathbf{x}_\beta^{(1)}$  ▷ project from  $\Omega_1$  to  $\Gamma_2$ 
12:  repeat                                                                                                         ▷ Newton loop for  $\Omega_2$ 
13:     $\Delta\mathbf{x}_B^{(2)} \leftarrow -\mathbf{K}_{AB}^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)}) \setminus \mathbf{R}_A^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)})$  ▷ linear system
14:     $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$ 
15:  until  $\|\Delta\mathbf{x}_B^{(2)}\| / \|\mathbf{x}_B^{(2)}\| \leq \epsilon_{\text{machine}}$                                          ▷ tight tolerance
16: until  $\left[ \left( \|\mathbf{y}^{(1)} - \mathbf{x}_B^{(1)}\| / \|\mathbf{x}_B^{(1)}\| \right)^2 + \left( \|\mathbf{y}^{(2)} - \mathbf{x}_B^{(2)}\| / \|\mathbf{x}_B^{(2)}\| \right)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$  ▷ tight tolerance
  
```

Appendix. Inexact Schwarz Method

Classical algorithm originally proposed by Schwarz with **outer Schwarz loop** and **inner Newton loop**, with Newton step converged to a **loose tolerance**.

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat                                                 ▷ Schwarz loop
4:    $\mathbf{y}^{(1)} \leftarrow \mathbf{x}_B^{(1)}$                                ▷ for convergence check
5:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$            ▷ project from  $\Omega_2$  to  $\Gamma_1$ 
6:   repeat                                                 ▷ Newton loop for  $\Omega_1$ 
7:      $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$            ▷ linear system
8:      $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$ 
9:   until  $\|\Delta\mathbf{x}_B^{(1)}\| / \|\mathbf{x}_B^{(1)}\| \leq \epsilon$            ▷ loose tolerance, e.g.  $\epsilon \in [10^{-4}, 10^{-1}]$ 
10:   $\mathbf{y}^{(2)} \leftarrow \mathbf{x}_B^{(2)}$                                ▷ for convergence check
11:   $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{P}_{21}\mathbf{x}_B^{(1)} + \mathbf{Q}_{21}\mathbf{x}_b^{(1)} + \mathbf{G}_{21}\mathbf{x}_\beta^{(1)}$            ▷ project from  $\Omega_1$  to  $\Gamma_2$ 
12:  repeat                                                 ▷ Newton loop for  $\Omega_2$ 
13:     $\Delta\mathbf{x}_B^{(2)} \leftarrow -\mathbf{K}_{AB}^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)}) \setminus \mathbf{R}_A^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)})$            ▷ solve linear system
14:     $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$ 
15:  until  $\|\Delta\mathbf{x}_B^{(2)}\| / \|\mathbf{x}_B^{(2)}\| \leq \epsilon$            ▷ loose tolerance, e.g.  $\epsilon \in [10^{-4}, 10^{-1}]$ 
16: until  $\left[ \left( \|\mathbf{y}^{(1)} - \mathbf{x}_B^{(1)}\| / \|\mathbf{x}_B^{(1)}\| \right)^2 + \left( \|\mathbf{y}^{(2)} - \mathbf{x}_B^{(2)}\| / \|\mathbf{x}_B^{(2)}\| \right)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$            ▷ tight tolerance
  
```

Appendix. Monolithic Schwarz Method

Combines Schwarz and Newton loop into ***since Newton-Schwarz loop***, with ***elimination of Schwarz boundary DOFs***, and tight convergence tolerance.

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,           ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,           ▷ initialize for  $\Omega_2$ 
3: repeat                                                               ▷ Newton-Schwarz loop
4:    $\begin{Bmatrix} \Delta \mathbf{x}_B^{(1)} \\ \Delta \mathbf{x}_B^{(2)} \end{Bmatrix} \leftarrow \begin{pmatrix} \mathbf{K}_{AB}^{(1)} + \mathbf{K}_{A\beta}^{(1)} \mathbf{H}_{11} & \mathbf{K}_{A\beta}^{(1)} \mathbf{H}_{12} \\ \mathbf{K}_{A\beta}^{(2)} \mathbf{H}_{21} & \mathbf{K}_{AB}^{(2)} + \mathbf{K}_{A\beta}^{(2)} \mathbf{H}_{22} \end{pmatrix} \backslash \begin{Bmatrix} -\mathbf{R}_A^{(1)} \\ -\mathbf{R}_A^{(2)} \end{Bmatrix}$            ▷ linear system
5:    $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta \mathbf{x}_B^{(1)}$ 
6:    $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta \mathbf{x}_B^{(2)}$ 
7: until  $\left[ \left( \|\Delta \mathbf{x}_B^{(1)}\| / \|\mathbf{x}_B^{(1)}\| \right)^2 + \left( \|\Delta \mathbf{x}_B^{(2)}\| / \|\mathbf{x}_B^{(2)}\| \right)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$            ▷ tight tolerance

```

Advantages:

- By-passes Schwarz loop.

Disadvantages:

- Off-diagonal coupling terms \rightarrow block linear solver is needed.

Appendix. Modified Schwarz Method

Combines Schwarz and Newton loop into *since Newton-Schwarz loop*, with *Schwarz boundaries* at *Dirichlet boundaries* and tight convergence tolerance.

```

1:  $\mathbf{x}_B^{(1)} \leftarrow \mathbf{X}_B^{(1)}$  in  $\Omega_1$ ,  $\mathbf{x}_b^{(1)} \leftarrow \chi(\mathbf{X}_b^{(1)})$  on  $\partial_\varphi \Omega_1$ ,  $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{X}_\beta^{(1)}$  on  $\Gamma_1$            ▷ initialize for  $\Omega_1$ 
2:  $\mathbf{x}_B^{(2)} \leftarrow \mathbf{X}_B^{(2)}$  in  $\Omega_2$ ,  $\mathbf{x}_b^{(2)} \leftarrow \chi(\mathbf{X}_b^{(2)})$  on  $\partial_\varphi \Omega_2$ ,  $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{X}_\beta^{(2)}$  on  $\Gamma_2$            ▷ initialize for  $\Omega_2$ 
3: repeat                                                 ▷ Newton-Schwarz loop
4:    $\mathbf{x}_\beta^{(1)} \leftarrow \mathbf{P}_{12}\mathbf{x}_B^{(2)} + \mathbf{Q}_{12}\mathbf{x}_b^{(2)} + \mathbf{G}_{12}\mathbf{x}_\beta^{(2)}$            ▷ project from  $\Omega_2$  to  $\Gamma_1$ 
5:    $\Delta\mathbf{x}_B^{(1)} \leftarrow -\mathbf{K}_{AB}^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)}) \setminus \mathbf{R}_A^{(1)}(\mathbf{x}_B^{(1)}; \mathbf{x}_b^{(1)}; \mathbf{x}_\beta^{(1)})$            ▷ linear system
6:    $\mathbf{x}_B^{(1)} \leftarrow \mathbf{x}_B^{(1)} + \Delta\mathbf{x}_B^{(1)}$ 
7:    $\mathbf{x}_\beta^{(2)} \leftarrow \mathbf{P}_{21}\mathbf{x}_B^{(1)} + \mathbf{Q}_{21}\mathbf{x}_b^{(1)} + \mathbf{G}_{21}\mathbf{x}_\beta^{(1)}$            ▷ project from  $\Omega_1$  to  $\Gamma_2$ 
8:    $\Delta\mathbf{x}_B^{(2)} \leftarrow -\mathbf{K}_{AB}^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)}) \setminus \mathbf{R}_A^{(2)}(\mathbf{x}_B^{(2)}; \mathbf{x}_b^{(2)}; \mathbf{x}_\beta^{(2)})$            ▷ linear system
9:    $\mathbf{x}_B^{(2)} \leftarrow \mathbf{x}_B^{(2)} + \Delta\mathbf{x}_B^{(2)}$ 
10:  until  $\left[ \left( \|\Delta\mathbf{x}_B^{(1)}\| / \|\mathbf{x}_B^{(1)}\| \right)^2 + \left( \|\Delta\mathbf{x}_B^{(2)}\| / \|\mathbf{x}_B^{(2)}\| \right)^2 \right]^{1/2} \leq \epsilon_{\text{machine}}$            ▷ tight tolerance

```

Advantages:

- By-passes Schwarz loop.
- No diagonal coupling (conventional linear solver can be used in each subdomain).

Least-intrusive variant: by-passes Schwarz iteration, no need for block solver.

Appendix. Convergence Proof



A. Mota, J. Tezaur, C. Allerman

Schwarz Alternating Method in Solid Mechanics

2 Formulation of the Schwarz Alternating Method

We start by defining the standard finite element variational formulation to establish notation before presenting the formulation of the coupling method.

2.1 Variational Formulation on a Single Domain

Consider a domain $\Omega \subset \mathbb{R}^2$ representing a region described by the mapping $\varphi: \varphi(\mathbf{X}) \subset \Omega \rightarrow \mathbb{R}^2$, $\mathbf{X} \in \mathbb{R}^2$. $\varphi(\mathbf{X})$ is a displacement in \mathbb{R}^2 , $\varphi(\mathbf{X})$ is a function in \mathbb{R}^2 , and $\varphi(\mathbf{X}, \Omega) \subset \Omega \subset \mathbb{R}^2$. The prescribed boundary displacements or Dirichlet boundary conditions are $\varphi|_{\Gamma_D} = \mathbf{0}$. The traction boundary conditions are $\mathbf{F}|_{\Gamma_T} = \mathbf{0}$. Let $\mathbf{F} = \mathbf{G}$ be the displacement gradient. Let also $\mathbf{P} = \mathbf{F}^T \mathbf{F}$ be the body force, with the mass matrix and the corresponding stiffness matrix. Likewise, introduce the energy functional

$$S(\varphi) = \int_{\varphi(\mathbf{X})} A(\mathbf{F}, \mathbf{Z}) \, d\mathbf{Z} - \int_{\Gamma_T} \mathbf{F} \cdot \mathbf{v} - \mathbf{v}^T \cdot \mathbf{V} - \int_{\Gamma_D} \mathbf{v} \cdot \mathbf{p}, \quad (1)$$

in which $A(\mathbf{F}, \mathbf{Z})$ is the Helmholtz energy density in \mathbb{R}^2 as a function of material variables. The weak form of the problem is obtained by minimizing the energy functional $S(\varphi)$ over the Sobolev space $W(\Omega)$ that is composed of functions φ such that the corresponding displacement gradient \mathbf{F} and traction $\mathbf{F}^T \mathbf{F}$ are square-integrable functions.

$$\mathcal{S} = \{\varphi \in W(\Omega) : \varphi|_{\Gamma_D} = \mathbf{0}, \varphi \in \mathcal{C}_0(\Omega)\}, \quad (2)$$

and

$$\mathcal{V} = \{\mathbf{v} \in W(\Omega) : \mathbf{v}|_{\Gamma_D} = \mathbf{0}\}, \quad (3)$$

where $\mathcal{C}_0(\cdot)$ is a test function. The potential energy is minimized if and only if $\mathbf{F}^T \mathbf{F} \cdot \mathbf{v} = \mathbf{0}$ for all $\mathbf{v} \in \mathcal{V}$ and $\mathbf{v} \in \mathcal{S}$. It is straightforward to show that the minimum of $S(\varphi)$ is the mapping $\varphi \in \mathcal{S}$ that satisfies

$$\int_{\varphi(\mathbf{X})} \{\mathbf{v}, \mathbf{Z}\} \, d\mathbf{Z} = \int_{\Gamma_T} \mathbf{F} \cdot \mathbf{v} - \mathbf{v}^T \cdot \mathbf{V} - \int_{\Gamma_D} \mathbf{v} \cdot \mathbf{p} = 0, \quad (4)$$

where $\mathbf{P} = \mathbf{F}^T \mathbf{F}$ denotes the first Piola-Cauchy tensor. This is the discrete equation corresponding to the variational statement (1).

$$\begin{aligned} \mathbf{D}(\mathbf{F}, \mathbf{Z}, \mathbf{p}, \mathbf{v}) &= \mathbf{0}, \\ \mathbf{D}(\mathbf{F}, \mathbf{Z}, \mathbf{p}, \mathbf{v}) &= \mathbf{0}, \\ \mathbf{F} \cdot \mathbf{v} &= \mathbf{0}. \end{aligned} \quad (5)$$

2.2 Coupling Two or More Subdomains via the Schwarz Alternating Method

In this section, we describe the Schwarz alternating method for coupling two overlapping subdomains. Consider subdomains Ω_1 and Ω_2 that overlap in the region $\Omega \subset \Omega_1 \cup \Omega_2$ as shown in Figure 1.

In order to prove the convergence of the Schwarz alternating method, possibly we ask for some convergence for the finite difference or finite element problem. We form a set of indices that alternate between the subdomains as

$$n \in \{1, 2, 1, 2, \dots, 2\} = \{1, 2, 1, 2, \dots, 2, 1, 2, \dots, 2\} \subset \{1, 2\}. \quad (6)$$

5

A. Mota, J. Tezaur, C. Allerman

Schwarz Alternating Method in Solid Mechanics

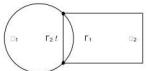


Figure 1: Two subdomains, Ω_1 and Ω_2 , and the corresponding boundaries Γ_1 and Γ_2 , used by the Schwarz alternating method.

that is $i = 1$ and $j = 2$ if n is odd, and $i = 2$ and $j = 1$ if n is even. Introduce the following definitions for each subdomain i :

- 1. Domain: $\Omega_i \subset \mathbb{R}^2$, $i \in \{1, 2\}$,
- 2. Mesh boundary: $\Gamma_i \subset \mathbb{R}^1$, $i \in \{1, 2\}$,
- 3. Schwarz boundary: $\Gamma_{ij} \subset \mathbb{R}^1$, $i, j \in \{1, 2\}$.

Note that with these definitions we guarantee that $\mathbb{1} \otimes \mathbb{1} = \mathbb{1} \otimes \mathbb{1}$. Now define the space

$$\mathcal{S} = \{x \in \mathbb{W}(\Omega_1 \cup \Omega_2) : x = \mathbf{0} \otimes \mathbb{1}_{\Omega_i} \otimes \mathbf{P}_{ij} \otimes \mathbf{v}_j \text{ on } \Gamma_{ij}\}, \quad (7)$$

and

$$V = \{v \in \mathbb{W}(\Omega_1 \cup \Omega_2) : v = 0 \otimes \mathbb{1}_{\Omega_i} \text{ on } \Gamma_{ij}\}, \quad (8)$$

where the symbol P_{ij} denotes the projection from the subdomain i onto the Schwarz boundary Γ_{ij} . This projection operator plays a central role in the Schwarz alternating method. Its form and implementation are discussed in subsequent sections. For the moment, it is sufficient to assume that the operator is able to project a function \mathbf{v}_j from the Schwarz boundary of the i th subdomain to the Schwarz boundary of the j th subdomain.

The Schwarz alternating method solves a sequence of problems on $\Omega_1 \cup \Omega_2$. The solution $\mathbf{u}^{(n)}$ for the n th problem is given by

$$\mathbf{u}^{(n)} = \arg \min_{\mathbf{u} \in \mathcal{S}} \mathcal{S}(\mathbf{u}), \quad \text{for } n \geq 0, \quad (9)$$

where \mathbf{u}_0 is the identity map that maps X onto itself, i.e., zero at point x and

$$\Phi^{(0)} = \{A(\mathbf{F}, \mathbf{Z}) - \mathbf{F} \mathbf{B}^T - \mathbf{B} \mathbf{D}^T - \mathbf{D} \mathbf{C}^T\} \, d\mathbf{Z} - \mathbf{T} \, d\mathbf{S}. \quad (10)$$

A better guess, if available, may be used to initialize (9) on Ω_1 rather than the identity map \mathbf{u}_0 . The minimization of the functional (10) leads to a variational formulation of the form (4)–(5) for each subdomain i :

$$\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}) = \mathbf{P} \cdot \mathbf{G}(\mathbf{F}^{(n)}) \mathbf{D}^{(n)} \mathbf{V}^{(n)} - \mathbf{R} \mathbf{B}^{(n)} \mathbf{D}^{(n)} \mathbf{Z}^{(n)} - \mathbf{T}^{(n)} \mathbf{D}^{(n)} \mathbf{q} \otimes \mathbf{q}, \quad (11)$$

and

$$\{\mathbf{v}^{(n)}, \mathbf{Z}^{(n)}\} = \{\mathbf{P}^{(n)} \mathbf{v}^{(n)}, \mathbf{Z}^{(n)}\} \text{ on } \Gamma_{ij}, \quad (12)$$

where $\mathbf{v}^{(n)}$ is the identity map that maps X onto itself, i.e., zero at point x and

$$\Phi^{(n)} = \{A(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}) - \mathbf{F}^{(n)} \mathbf{B}^{(n)T} - \mathbf{B}^{(n)} \mathbf{D}^{(n)T} - \mathbf{D}^{(n)} \mathbf{C}^{(n)T}\} \, d\mathbf{Z} - \mathbf{T}^{(n)} \, d\mathbf{S}. \quad (13)$$

Again using (12) and also (5) in (13) leads to

$$\{\mathbf{P}^{(n)} \mathbf{v}^{(n)}, \mathbf{Z}^{(n)}\} = \{\mathbf{P}^{(n)} \mathbf{v}^{(n)}, \mathbf{Z}^{(n)}\} - \{\mathbf{P}^{(n)} \mathbf{v}^{(n)}, \mathbf{P}^{(n-1)} \mathbf{v}^{(n-1)}\} \otimes \mathbf{C}_n, \quad (14)$$

and substituting (14) into (13) we finally obtain that

$$\{\mathbf{v}^{(n)}, \mathbf{Z}^{(n)}\} \otimes \mathbf{C}_n \otimes \{\mathbf{P}^{(n)} \mathbf{v}^{(n)}, \mathbf{Z}^{(n)}\} \otimes \mathbf{C}_n. \quad (15)$$

Now, for $i = 1, 2, \dots, 2$ we have

$$\mathbf{v}_i := \mathbf{v}^{(n)} \otimes \mathbf{C}_i. \quad (16)$$

Remark 4: For property 3, the uniform continuity of $\Psi(\varphi)$, there exists a modulus of continuity $\omega > 0$ such that

$$|\mathbf{v}(\mathbf{X}) - \mathbf{v}(\mathbf{Y})| \leq \omega \|\mathbf{X} - \mathbf{Y}\|, \quad \forall \mathbf{X}, \mathbf{Y} \in \Omega. \quad (17)$$

where $\mathbf{v}(\mathbf{X}) = \mathbf{v}(\mathbf{X}, \Omega)$. From (16), we have $|\mathbf{v}_i(\mathbf{X}) - \mathbf{v}_i(\mathbf{Y})| \leq \omega \|\mathbf{X} - \mathbf{Y}\|$.

Remark 5: It was shown in [3] that in the case $\mathbf{v}_1(\mathbf{X}) = \mathbf{v}_2(\mathbf{X}) = \mathbf{0}$, $\forall \mathbf{X} \in \Omega$, there exist $\zeta_1 \in \Omega_1$ and $\zeta_2 \in \Omega_2$ such that

$$\mathbf{v}(\mathbf{X}) = \mathbf{v}_1(\mathbf{X}) + \mathbf{v}_2(\mathbf{X}), \quad (18)$$

and

$$\mathbf{v}(\mathbf{X}) = \mathbf{v}_1(\mathbf{X}) \otimes \mathbf{C}_1 + \mathbf{v}_2(\mathbf{X}) \otimes \mathbf{C}_2. \quad (19)$$

Remark 6: Note that for $\omega > 0$ we have

$$|\mathbf{v}(\mathbf{X}) - \mathbf{v}(\mathbf{Y})| = \omega \|\mathbf{X} - \mathbf{Y}\|, \quad (20)$$

for $\mathbf{X}, \mathbf{Y} \in \Omega$.

Remark 7: Let $\Psi(\varphi) = 0$. Then φ is the solution of the equation $\mathbf{D}(\mathbf{F}, \mathbf{Z}, \mathbf{p}, \mathbf{v}) = 0$.

Remark 8: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 9: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 10: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 11: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 12: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 13: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 14: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 15: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 16: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 17: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 18: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 19: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 20: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 21: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 22: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 23: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 24: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 25: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 26: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 27: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 28: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 29: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 30: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 31: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 32: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 33: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 34: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 35: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 36: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 37: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 38: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 39: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 40: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 41: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 42: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 43: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 44: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 45: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 46: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 47: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 48: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 49: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 50: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 51: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{(n)}, \mathbf{v}^{(n)}) = 0$.

Remark 52: Let $\mathbf{v}^{(n)} = \mathbf{0}$. Then $\mathbf{v}^{(n)} = \mathbf{0}$ is the solution of the equation $\mathbf{D}(\mathbf{F}^{(n)}, \mathbf{Z}^{(n)}, \mathbf{p}^{($

Appendix. Foulk's Singular Bar

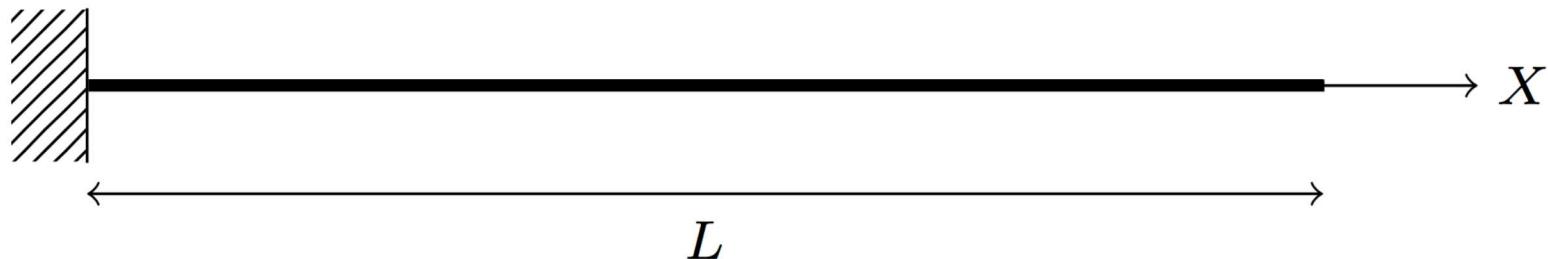
- **1D proof of concept** problem:
 - **1D bar** with area proportional to square root of length.
 - Strong **singularity** on left end of bar.
 - Simple **hypereelastic** material model with no damage.
 - **MATLAB** implementation.



$$u(0) = 0$$

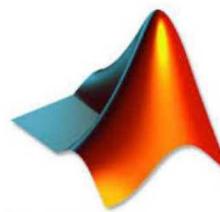
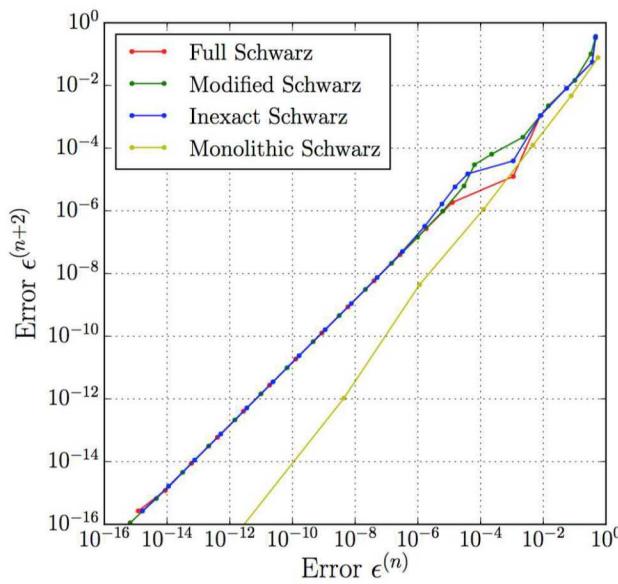
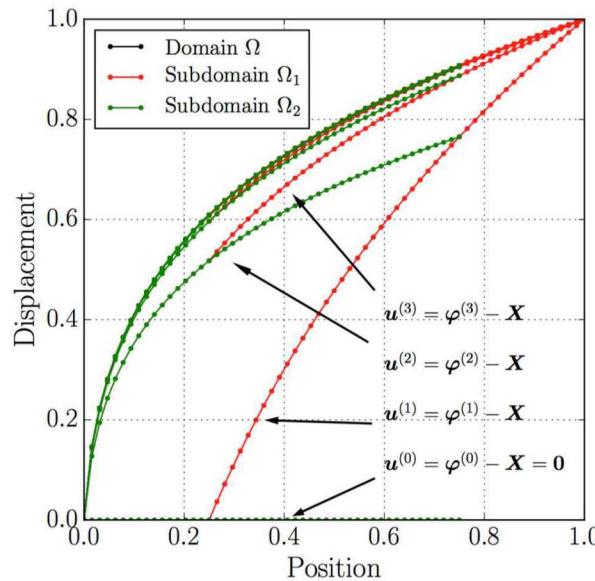
$$A(X) = A_0 \sqrt{X/L}$$

$$u(L) = \Delta$$

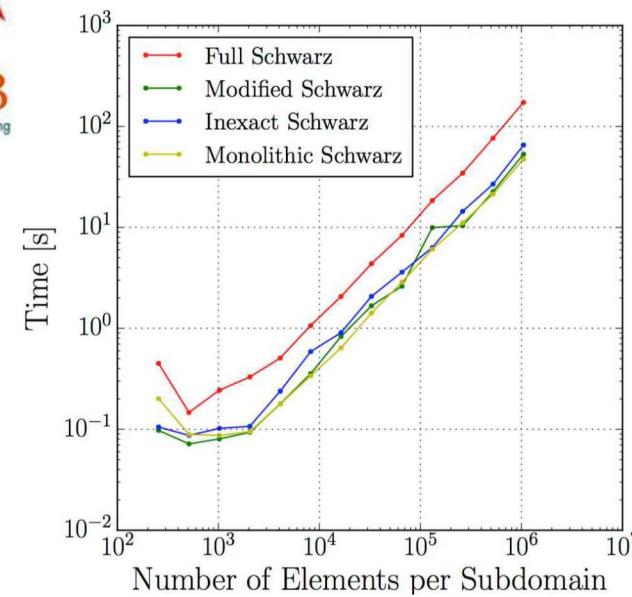
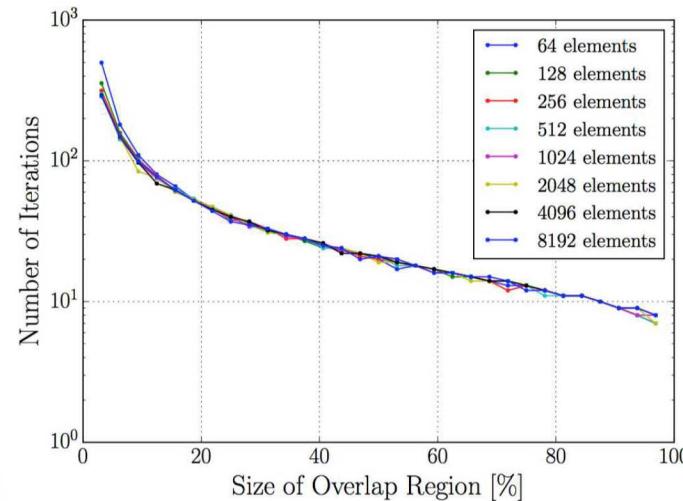


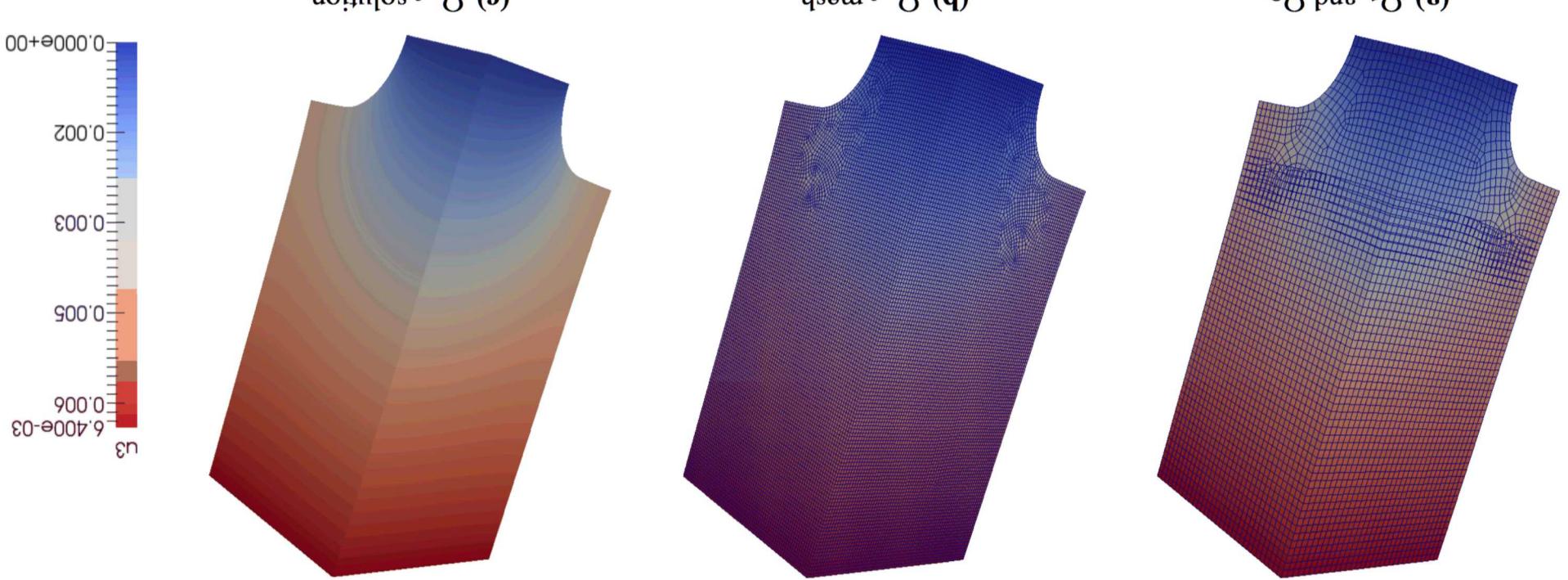
- **Problem goals:**
 - Explore **viability** of **4 variants** of the Schwarz alternating method.
 - Test **convergence** and compare with literature (Evans, 1986).
 - Expect **faster convergence** in **fewer iterations** with **increased overlap**.

Appendix. Singular Bar and Schwarz Variants



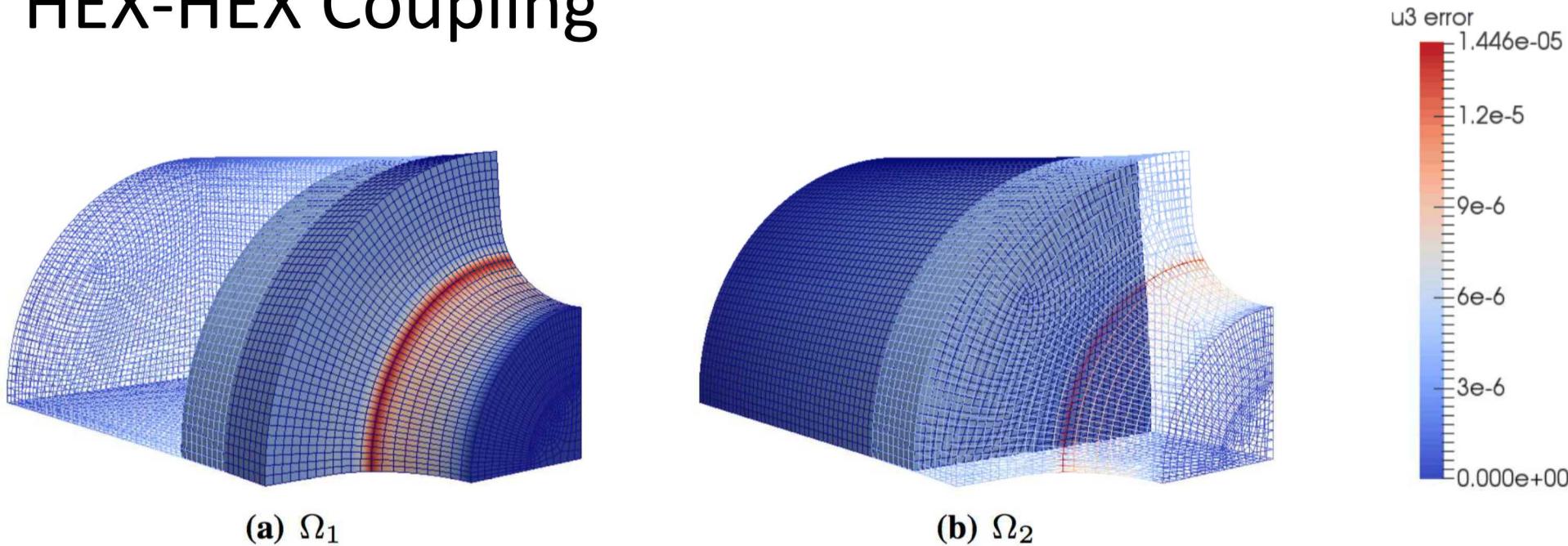
MATLAB
The Language of Technical Computing





Appendix. Notched Cylinder: Nonconforming
HEX-HEX Coupling

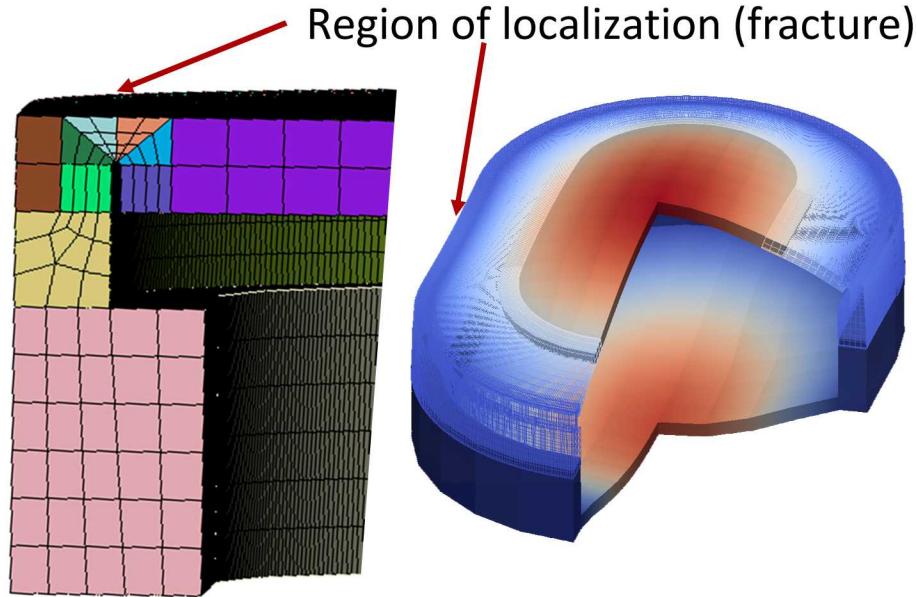
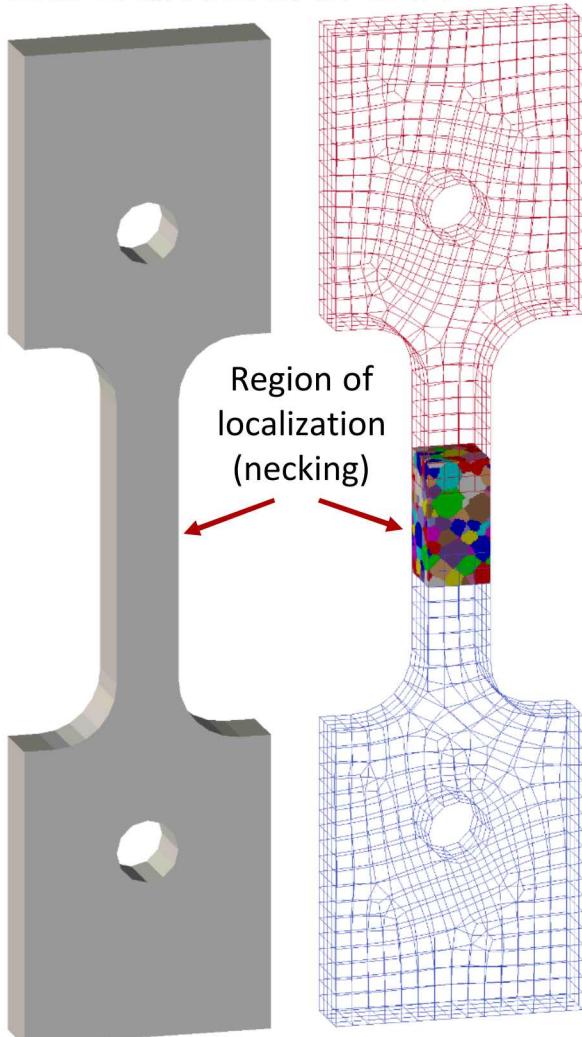
Appendix. Notched Cylinder: Nonconformal HEX-HEX Coupling



| Absolute residual tolerance | u_3 relative error | |
|-----------------------------|-----------------------|-----------------------|
| | Ω_1 | Ω_2 |
| 1.0×10^{-8} | 1.31×10^{-3} | 4.45×10^{-4} |
| 1.0×10^{-12} | 1.30×10^{-3} | 4.43×10^{-4} |
| 1.0×10^{-14} | 1.30×10^{-3} | 4.43×10^{-4} |
| 2.5×10^{-16} | 1.30×10^{-3} | 4.43×10^{-4} |



Appendix. Multiscale Modeling of Localization

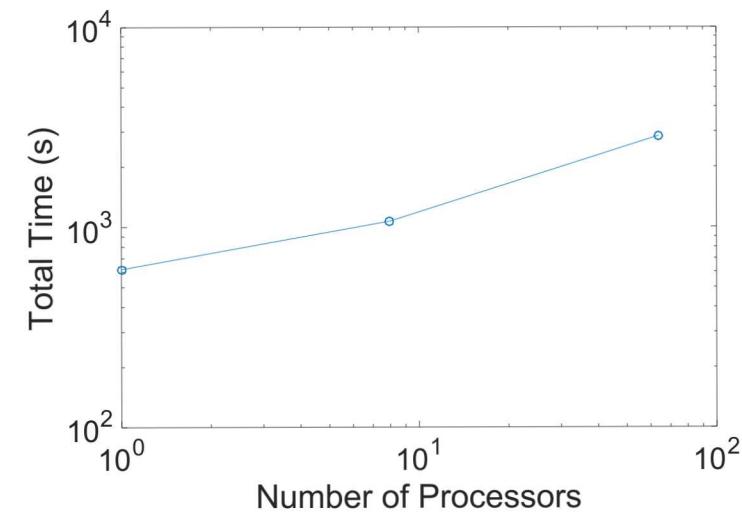
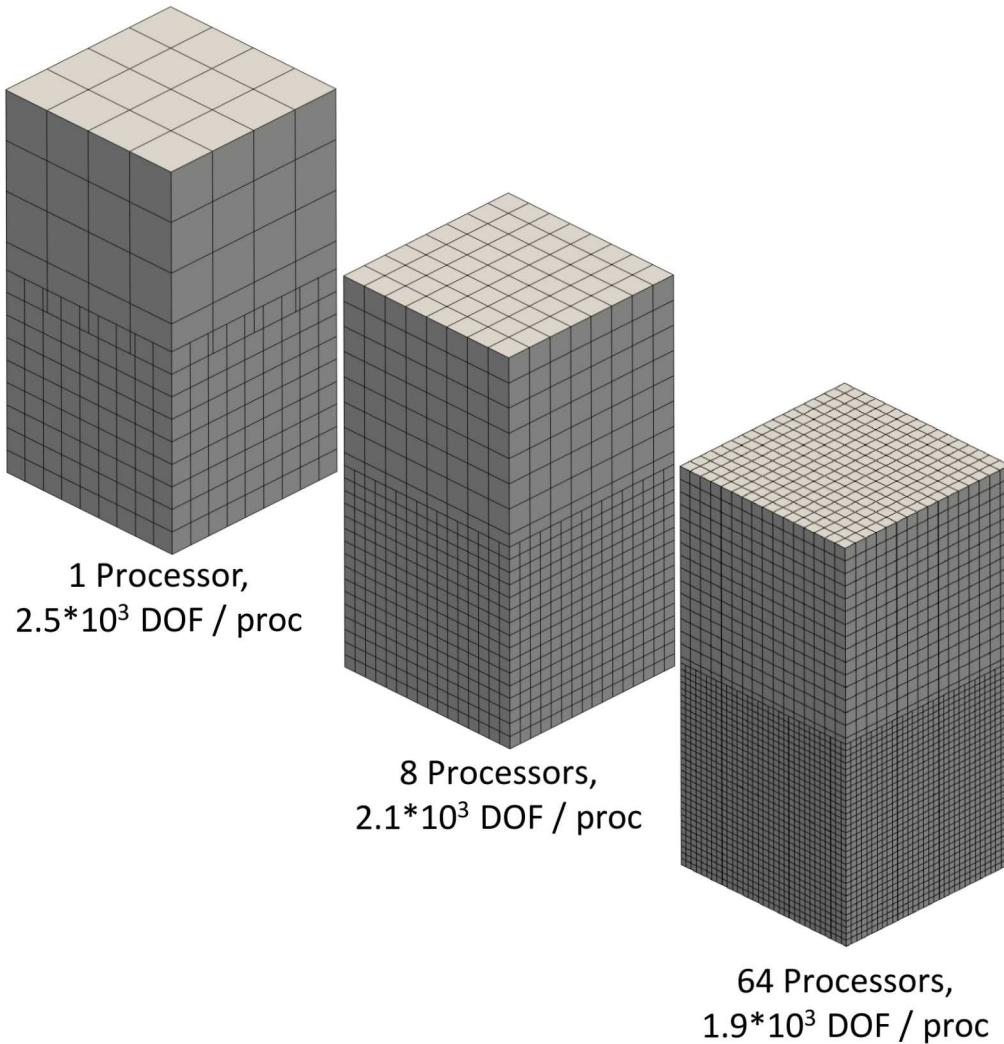


Strain localization can cause ***localized necking*** (left) and ultimately ***fracture*** (above).

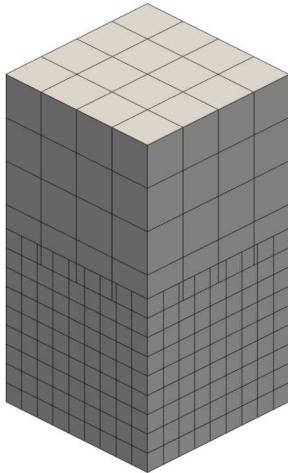
Goals:

- Connect ***physical length scales*** to ***engineering scale models***.
- Investigate importance of ***microstructural detail***.
- Develop bridging technologies for ***spatial multiscale/multiphysics***.

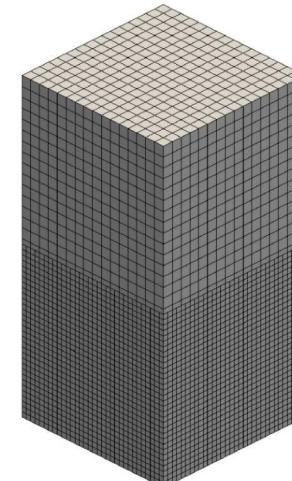
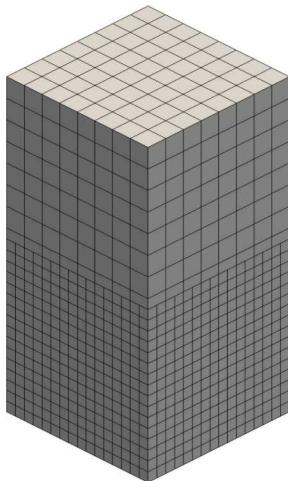
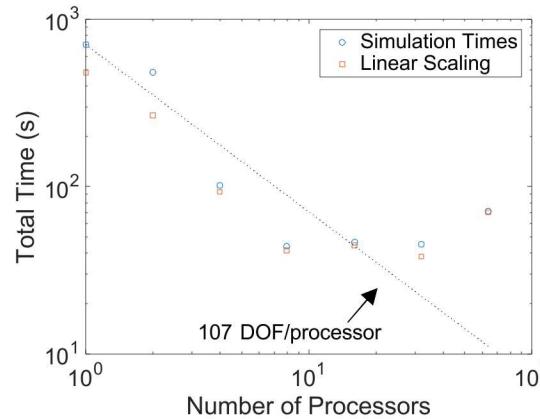
Appendix. Parallelization via DTK: Weak Scaling on Cubes Problem



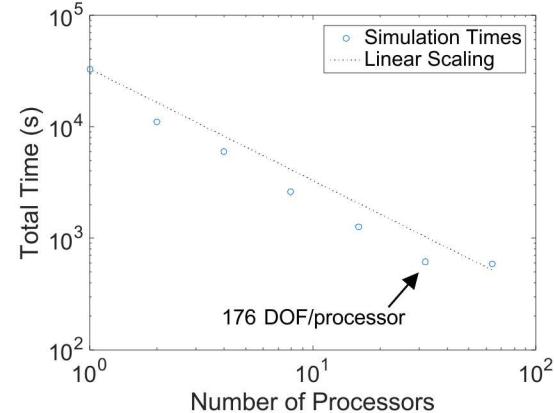
Appendix. Parallelization via DTK: Strong Scaling on Cubes Problem



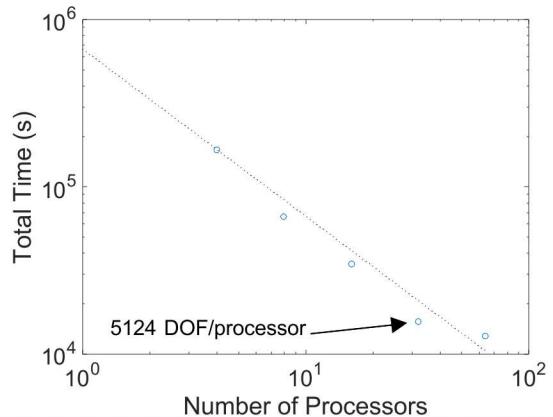
Small problem (2.5×10^3 DOFs)



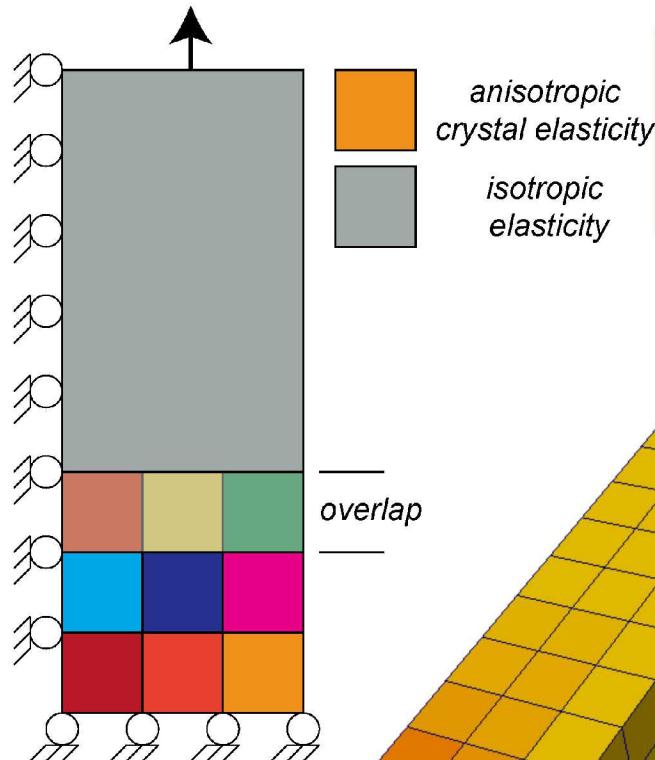
Medium problem (1.7×10^4 DOFs)



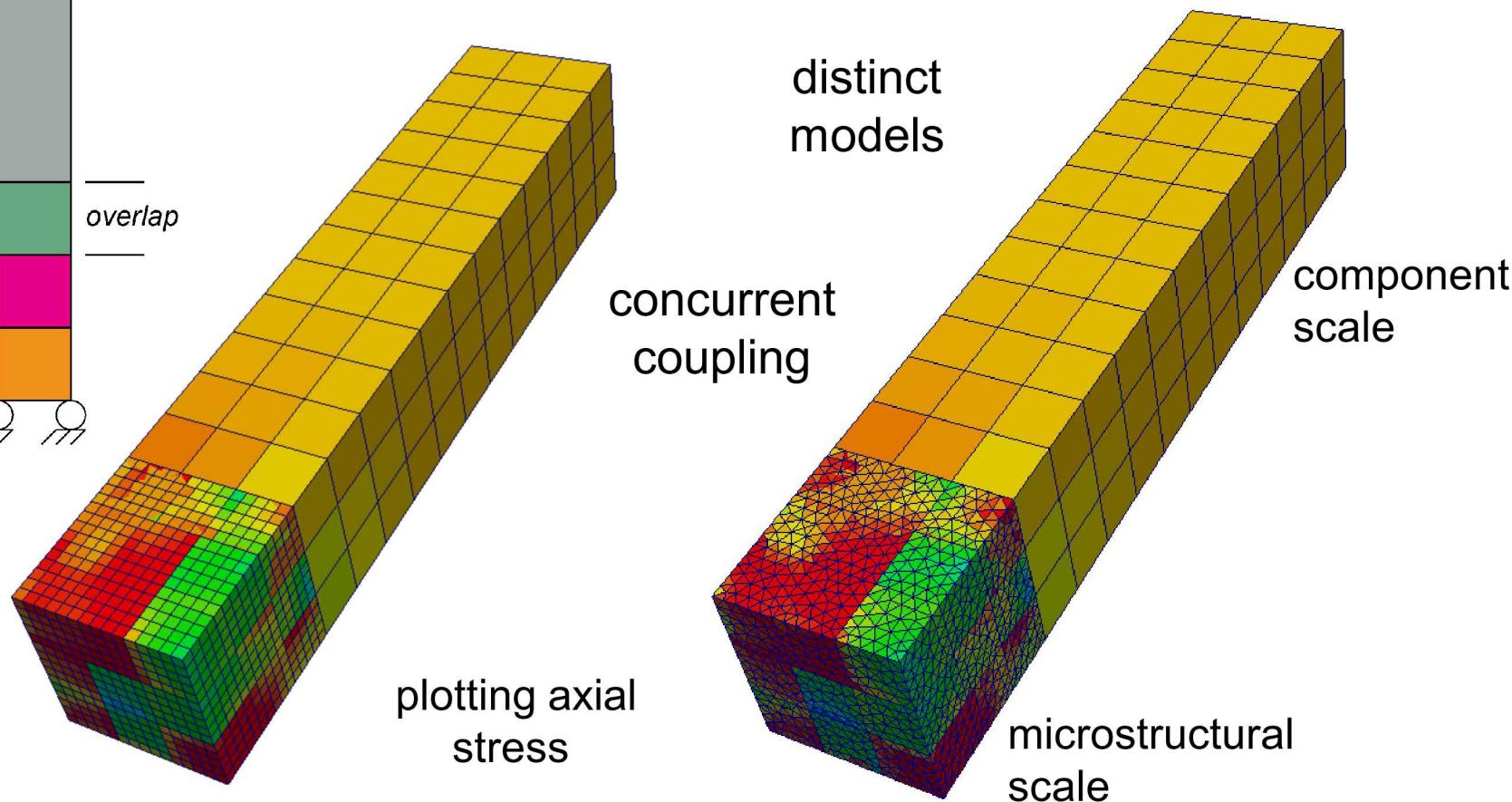
Large problem (1.6×10^5 DOFs)



Appendix. Rubiks Cube Problem

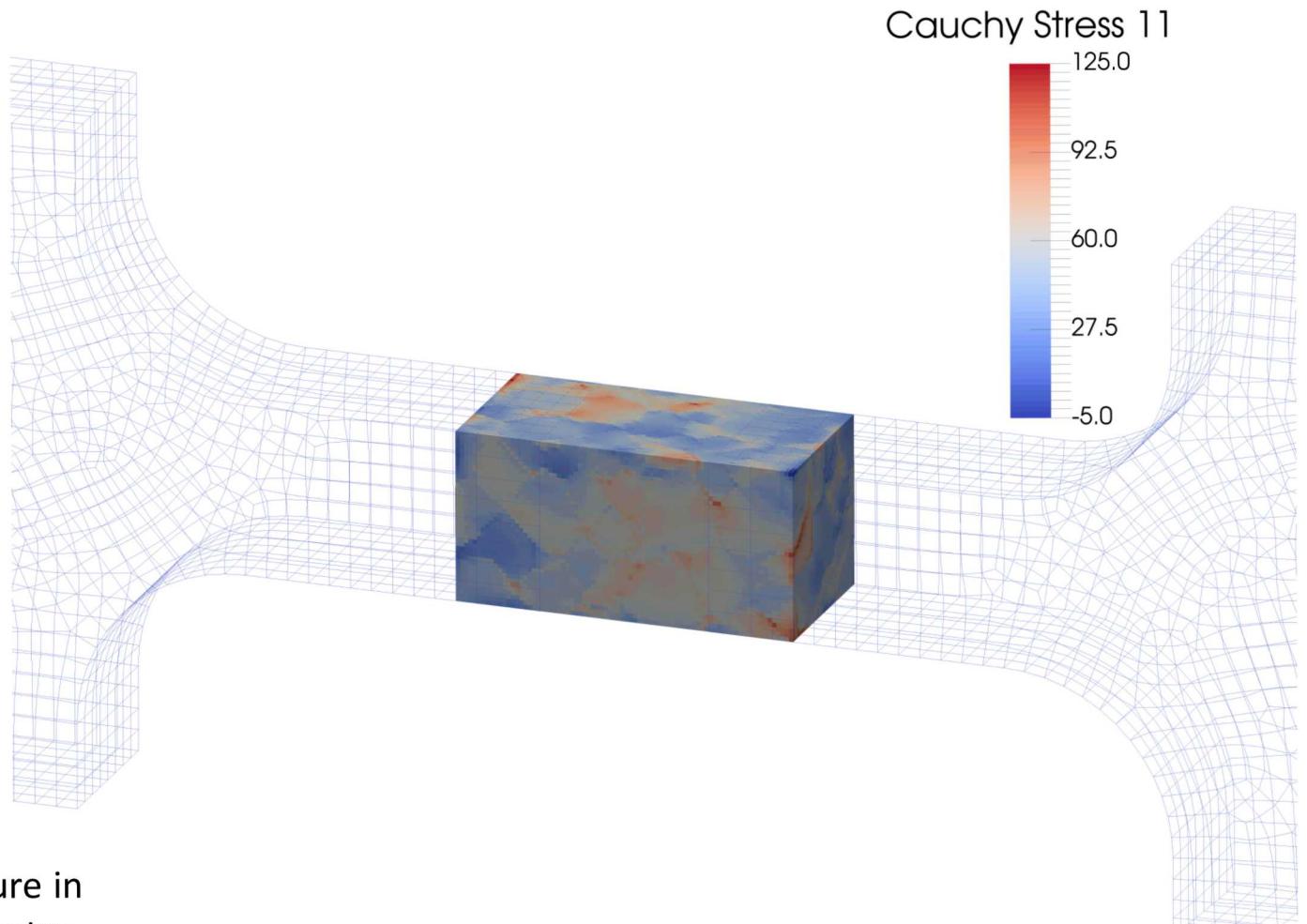
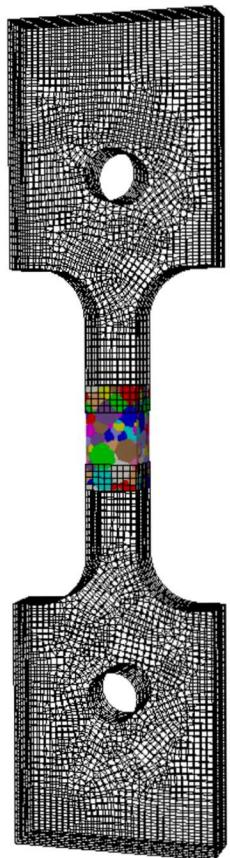


Two distinct bodies, the component scale and the microstructural scale, are coupled iteratively with alternating Schwarz



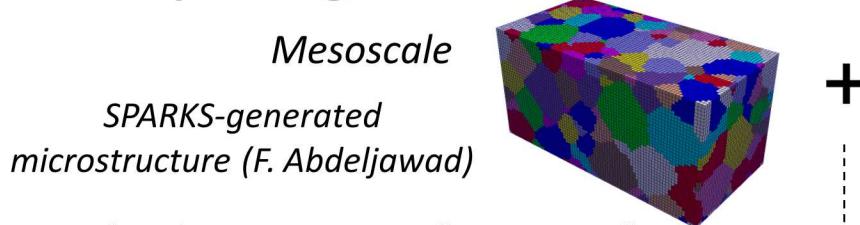
Work by J. Foulk, D. Littlewood, C. Battaile, H. Lim

Appendix. Tensile Bar



Embed microstructure in
ASTM tensile geometry

Appendix. Tensile Bar: Meso-Macroscopic Coupling



cubic elastic constant : $C_{11} = 204.6$ GPa

cubic elastic constant : $C_{12} = 137.7$ GPa

cubic elastic constant : $C_{44} = 126.2$ GPa

reference shear rate : $\dot{\gamma}_0 = 1.0$ 1/s

rate sensitivity factor : $m = 20$

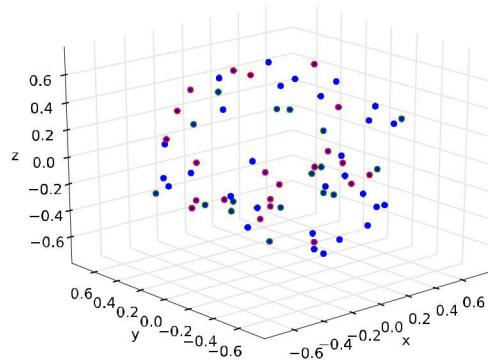
hardening rate parameter : $g_0 = 2.0 \times 10^4$ 1/s

initial hardness : $g_0 = 90$ MPa

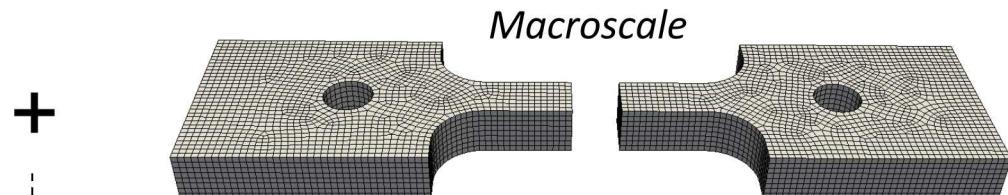
saturation hardness : $g_s = 202$ MPa

saturation exponent : $\omega = 0.01$

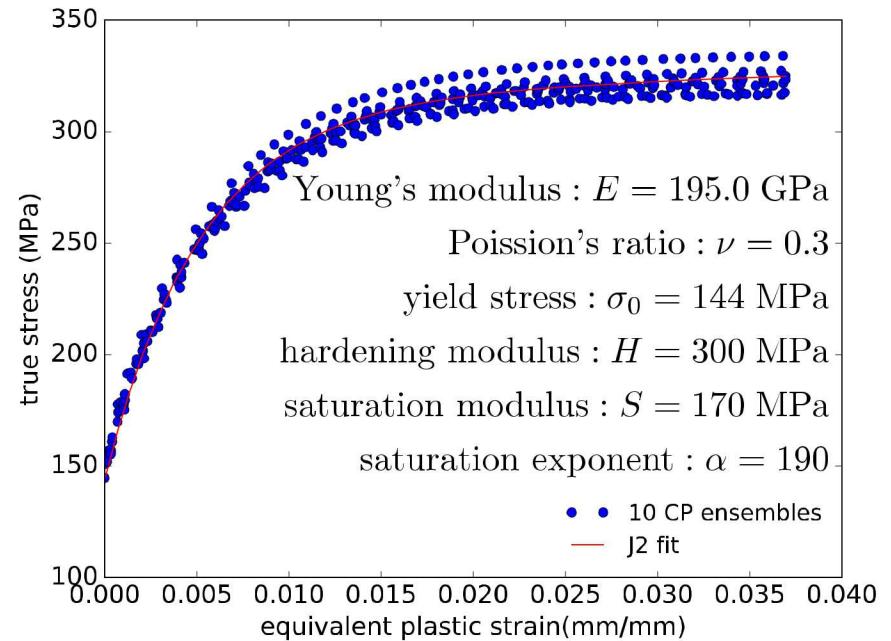
Fix microstructure, investigate ensembles



151 axial vectors
from 3 of the 10
ensembles of
random rotations
(blue, green, red)



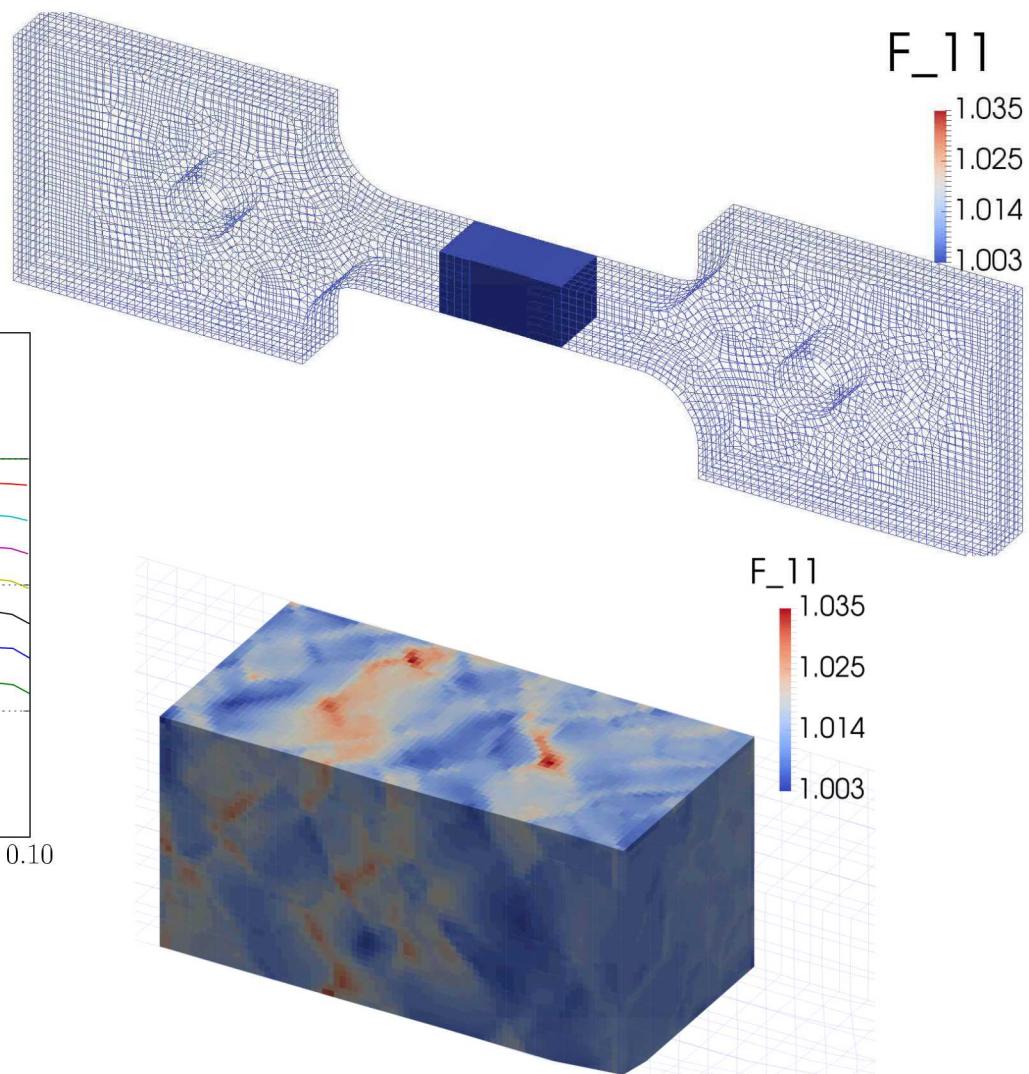
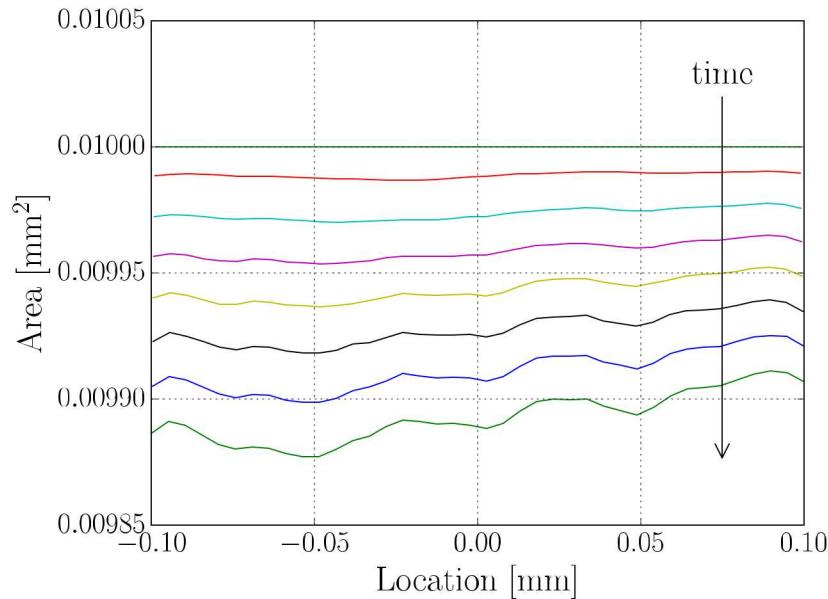
- Load microstructural ensembles in uniaxial stress
- Fit flow curves with a macroscale J_2 plasticity model



$$\sigma_y = \sigma_0 + H\epsilon_p + S(1 - e^{-\alpha\epsilon_p})$$

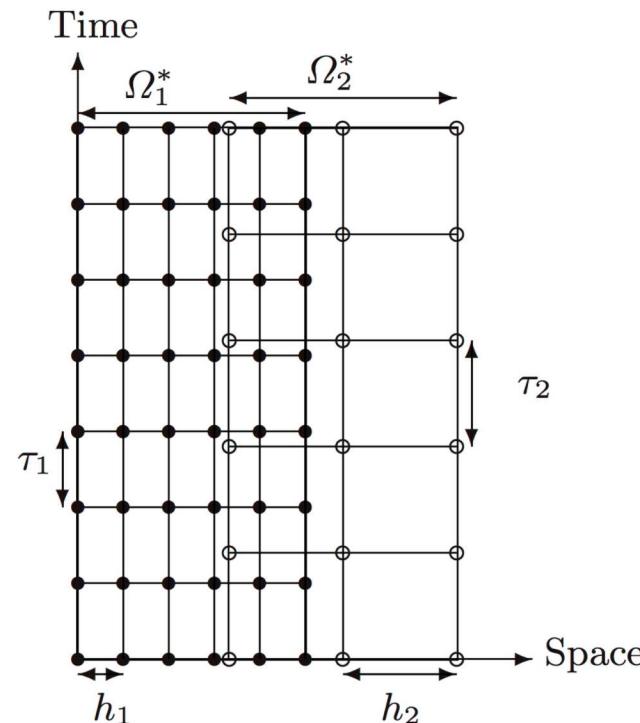
Appendix. Tensile Bar: Results

Reduction in cross-sectional area over time



Appendix. Schwarz Alternating Method for Dynamics

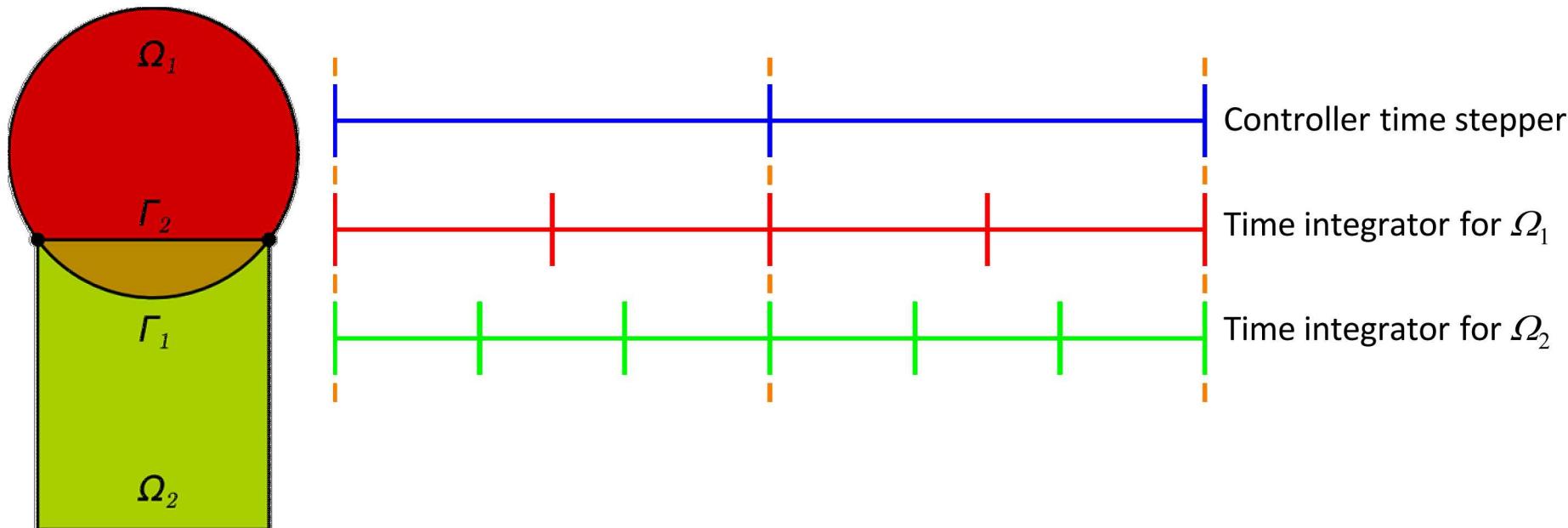
- In the literature the Schwarz method is applied to dynamics by using ***space-time discretizations***.
- This was deemed ***unfeasible*** given the design of our current codes and size of simulations.



Overlapping non-matching meshes and time steps in dynamics.

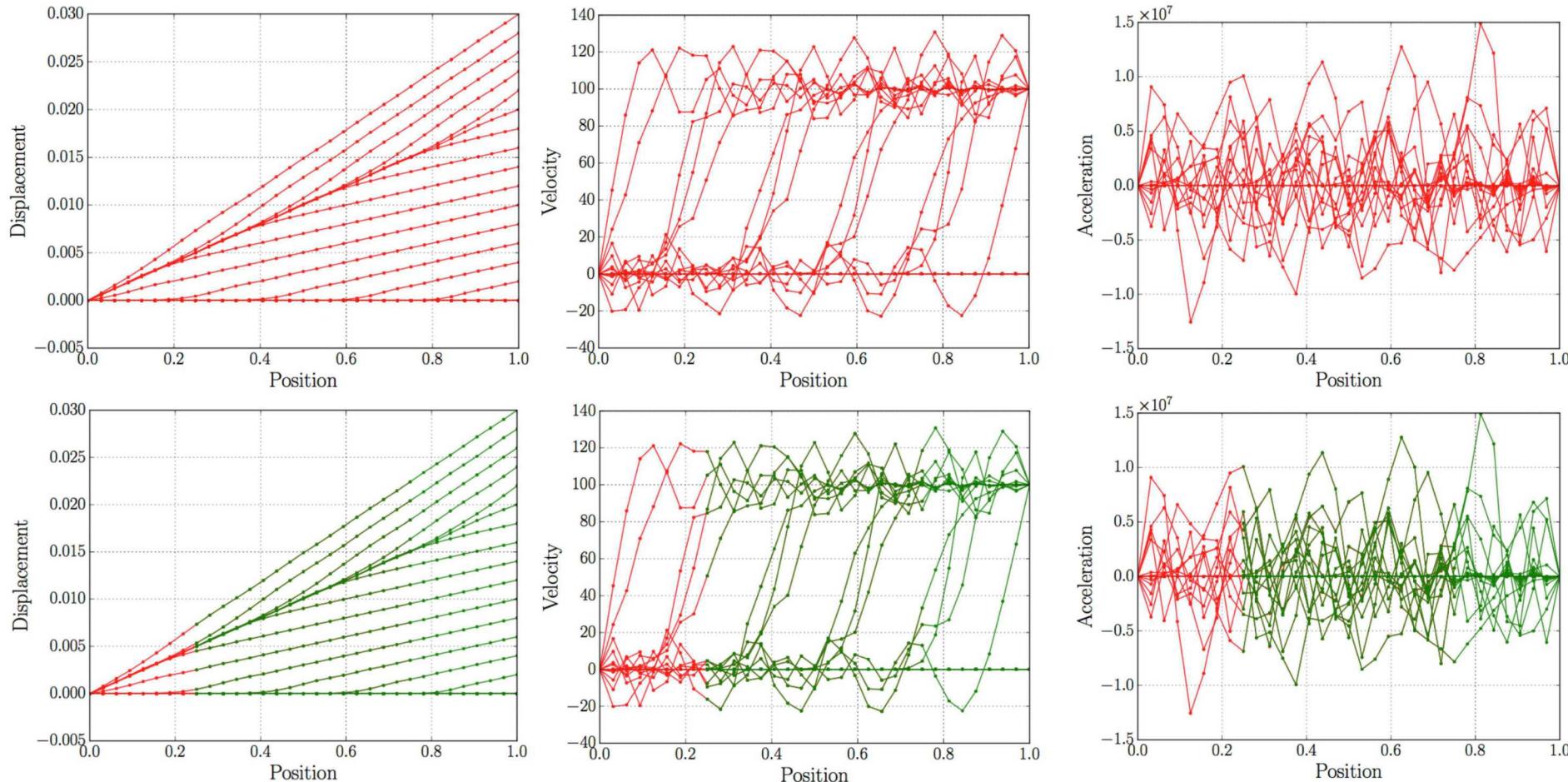
Appendix. A Schwarz-like Time Integrator

- We developed an ***extension of Schwarz coupling*** to ***dynamics*** using a governing time stepping algorithm that controls time integrators within each domain.
- Can use ***different integrators*** with ***different time steps*** within each domain.
- 1D results show ***smooth coupling without numerical artifacts*** such as spurious wave reflections at boundaries of coupled domains.



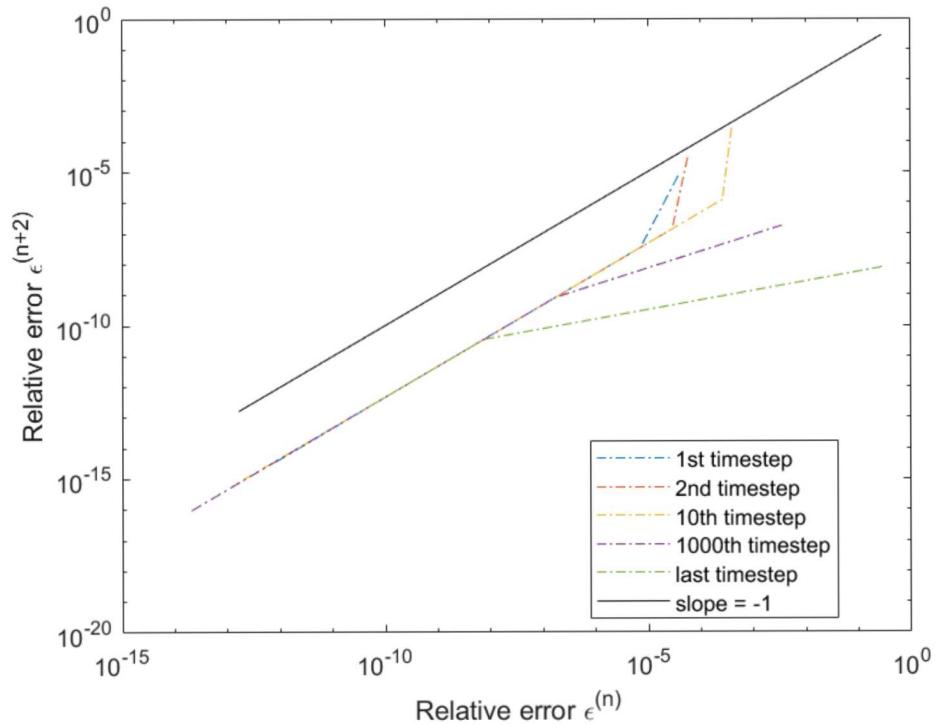
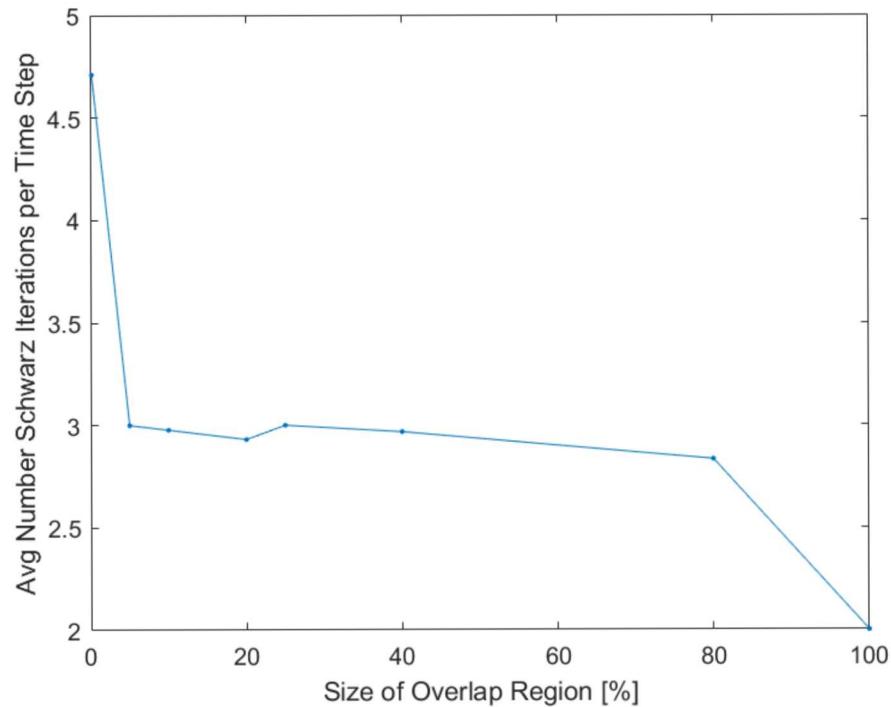
Appendix. Dynamic Singular Bar

- Inelasticity masks problems by introducing *energy dissipation*.
- Schwarz does *not* introduce *numerical artifacts*.
- Can couple domains with *different time integration schemes* (*Explicit-Implicit* below).



Appendix. Elastic Wave Propagation

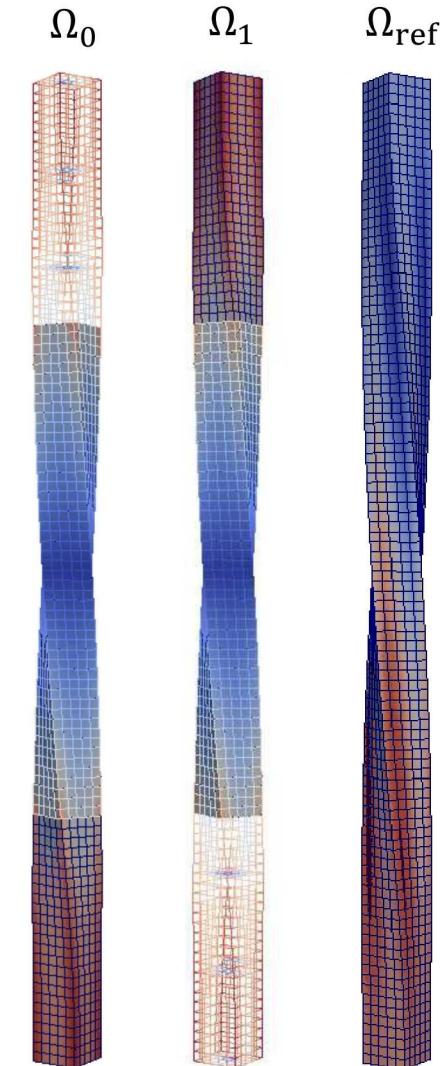
Some Performance Results



- Left figure shows **# of iterations** as a function of **overlap region size** for 2 subdomains. The method does not converge for 0% overlap. If the overlap is 100% then the single-domain solution is recovered for each of the subdomains.
- Right figure shows **linear convergence rate** of dynamic Schwarz implementation (for small overlap fraction of 0.2%).

Appendix. Torsion

- Nonlinear elastic bar (Neo-hookean material model) subjected to a high degree of ***torsion***.
- The ***domain*** is $\Omega = (-0.025, 0.025) \times (-0.025, 0.025) \times (-0.5, 0.5)$.
- We evaluate ***dynamic Schwarz*** with 2 subdomains:
 $\Omega_0 = (-0.025, 0.025) \times (-0.025, 0.025) \times (-0.5, 0.25)$, $\Omega_1 = (-0.025, 0.025) \times (-0.025, 0.025) \times (-0.25, 0.5)$.
- ***Time-discretizations***: Newmark-Beta (implicit, explicit) with same Δt .
- ***Meshes***: hexes, composite tet 10s.

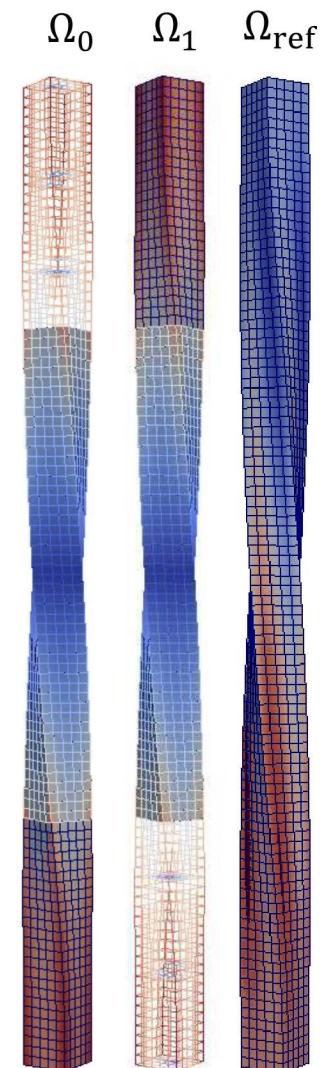
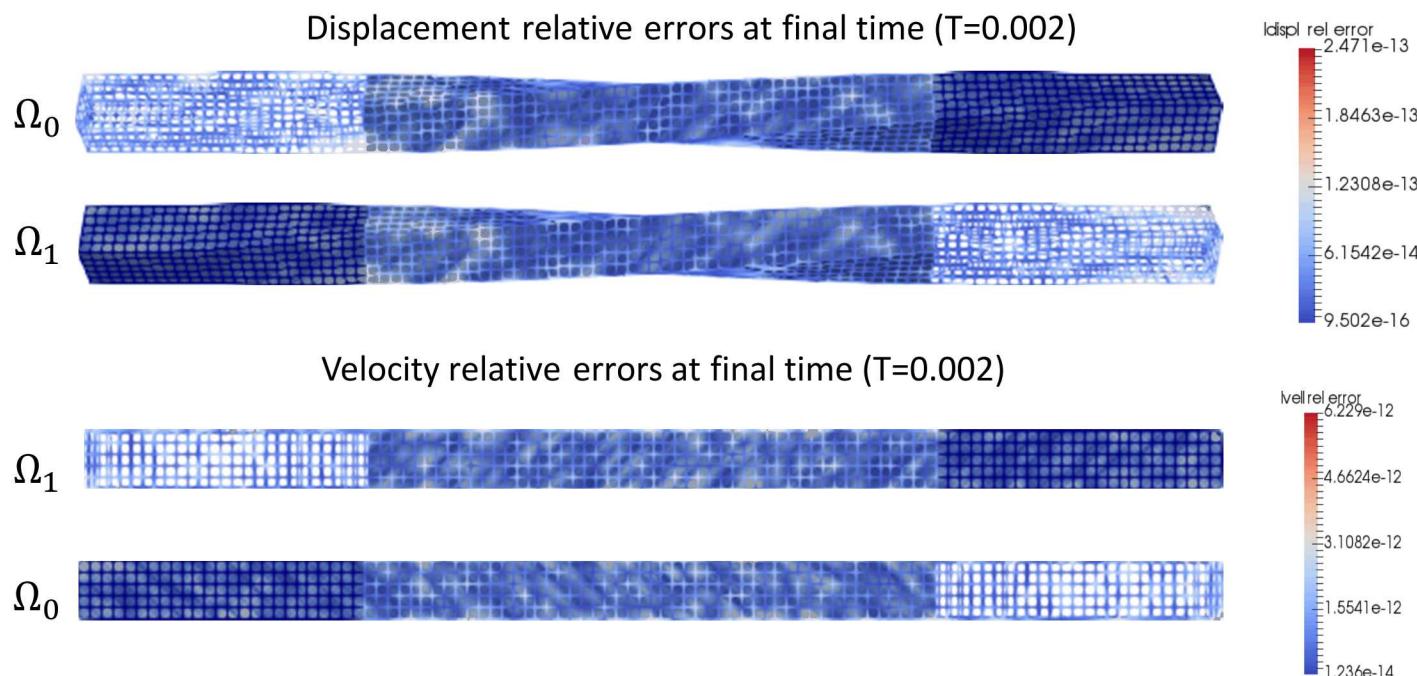


Appendix. Torsion

Schwarz and single-domain results agree to almost *machine-precision*!

Conformal Hex + Hex Coupling

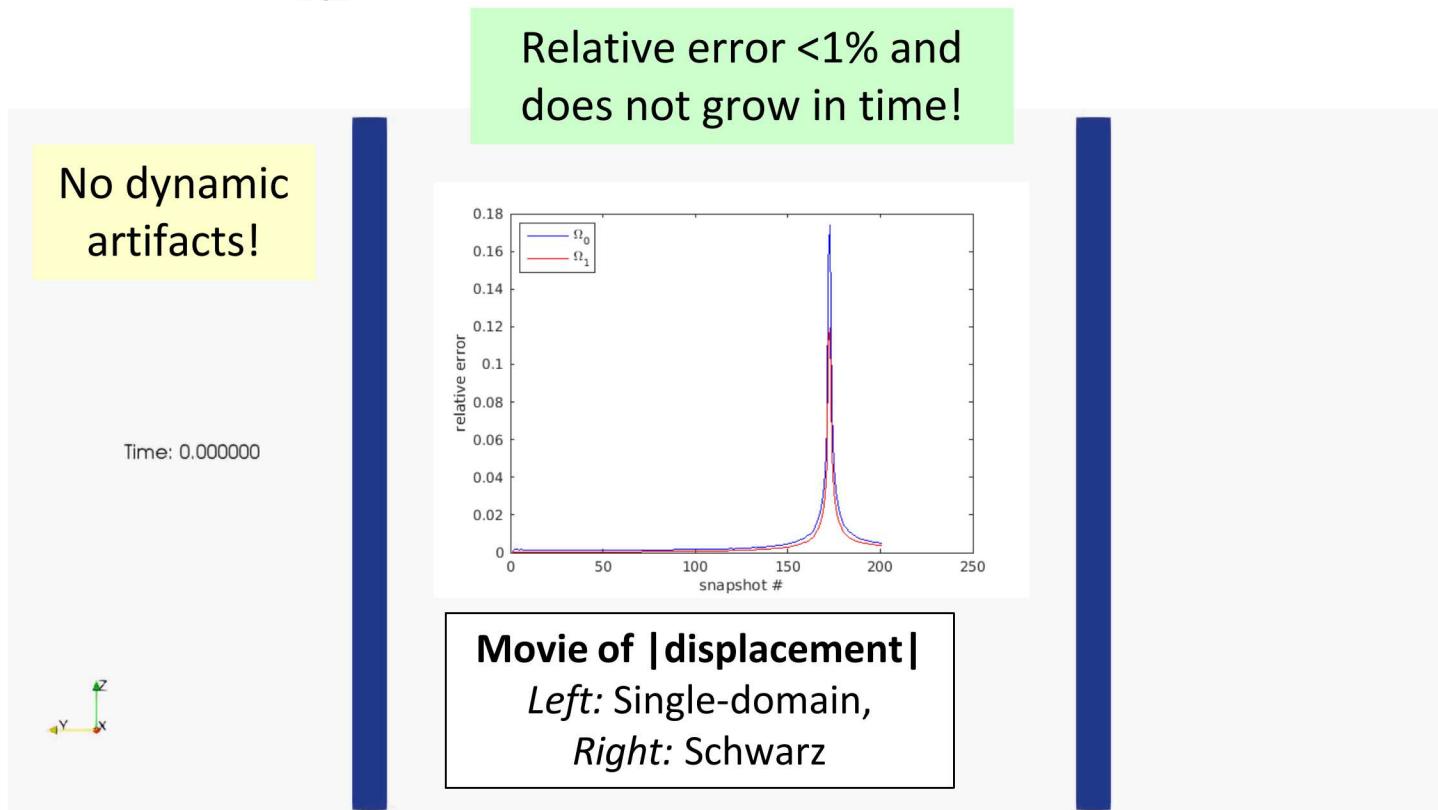
- Each subdomain discretized using **uniform hex mesh** with $\Delta x_i = 0.01$, and advanced in time using implicit Newmark-Beta scheme with $\Delta t = 1e-6$.
- Results compared to single-domain solution on mesh **conformal** with Schwarz domain meshes.



Appendix. Torsion

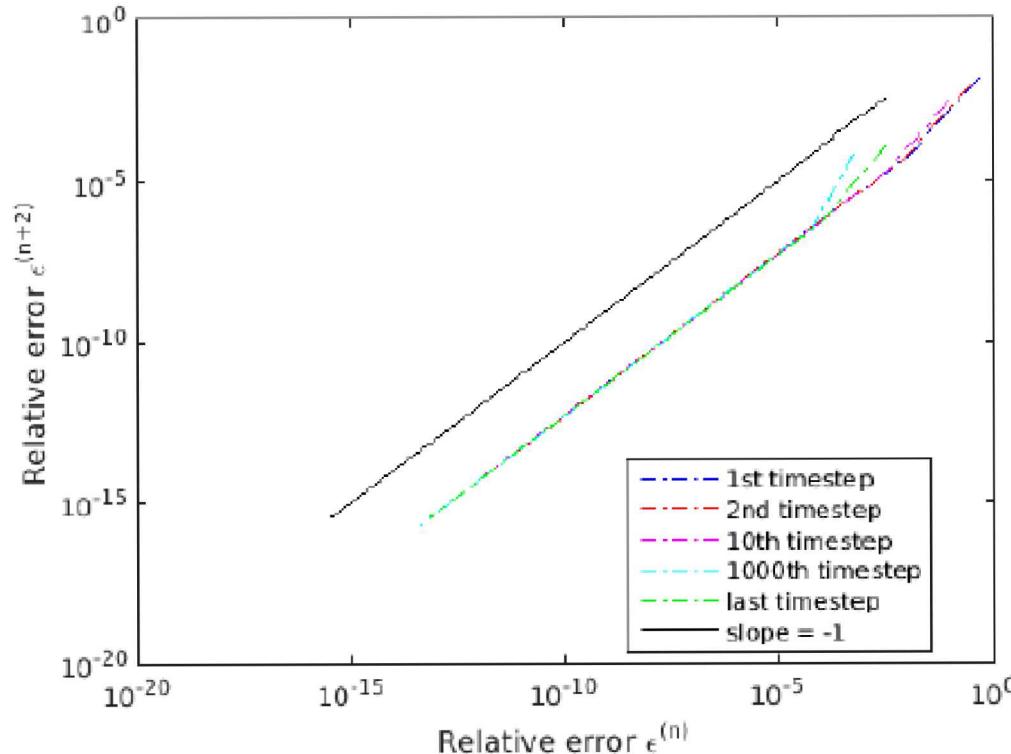
Hex + Composite Tet 10 Coupling

- Coupling of composite tet 10s + explicit Newmark with consistent mass in Ω_0 with hexes + implicit Newmark in Ω_1 .
- Reference solution is computed on fine hex mesh + implicit Newmark Ω_{ref}



Appendix. Torsion

Some Performance Results



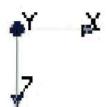
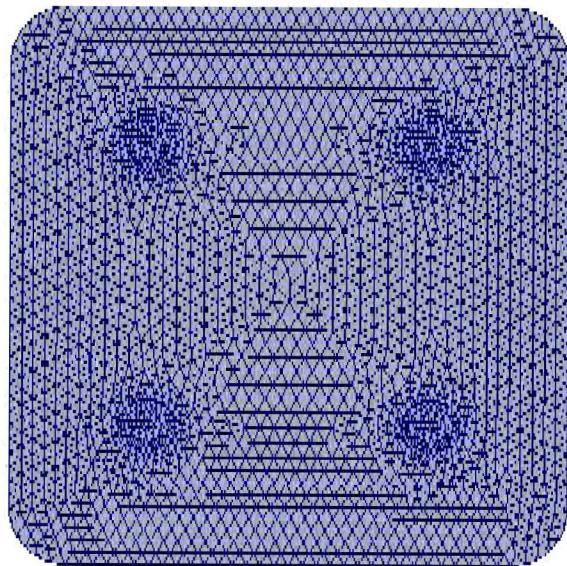
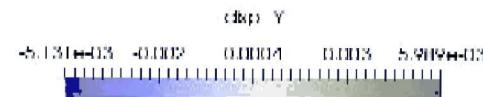
- Convergence behavior of the dynamic Schwarz algorithm for the torsion problem for small overlap volume fraction (2%) in which each subdomain is discretized using a hexahedral mesh. The plot shows that a ***linear convergence rate*** is achieved.

Appendix. Bolted Joint Problem

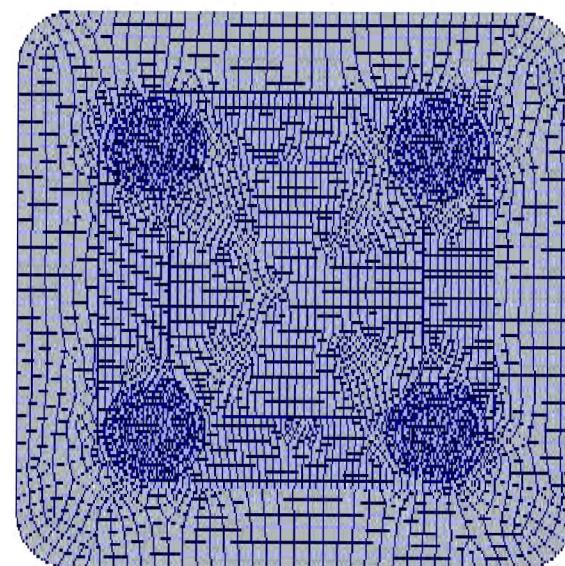
1es

y-displacement

Time: 0.000000



Single Ω



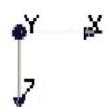
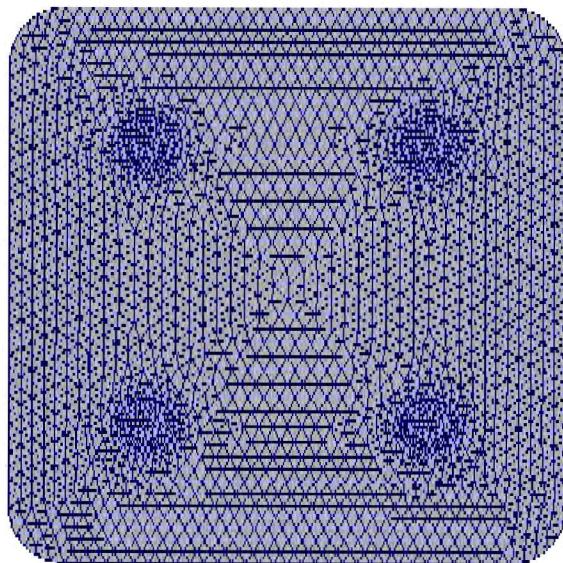
Schwarz

Appendix. Bolted Joint Problem

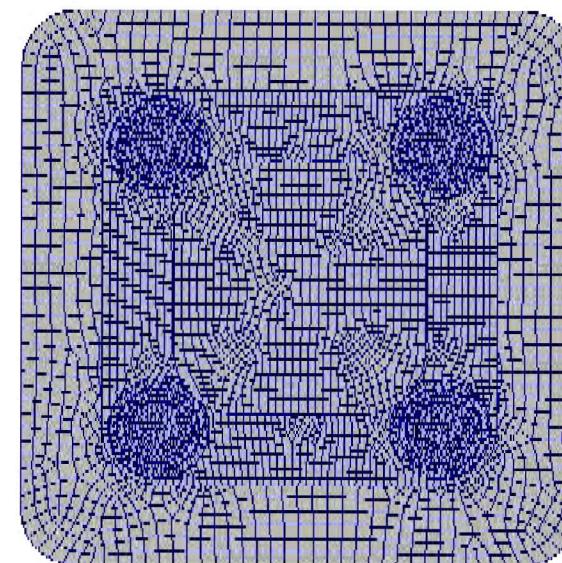
ies

z-displacement

Time: 0.000000



Single Ω



Schwarz