

Spin qubits made of quantum dots and donors in silicon

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2. Sandia National Laboratories, Albuquerque, NM, 87185, United States

APS March Meeting
March 4-8, 2019

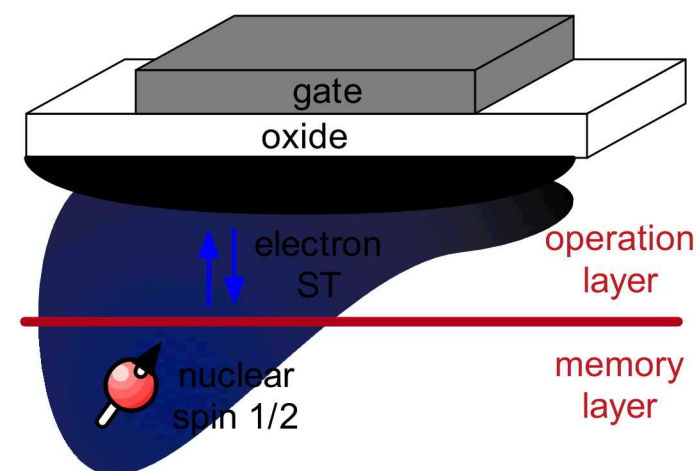
Now at QuTech, Delft University of Technology!



- Donor-Dot system
- Valley splitting, singlet-triplet splitting, coupling to donors, and 4-electron trick
- Enhanced latching readout
- Singlet-triplet qubit (hyperfine-driven rotations)

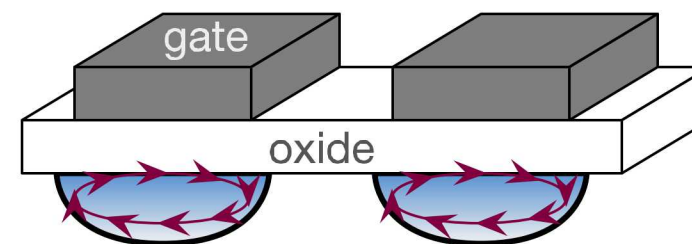
Part 1:

Donor-dot system and related topics



- Universal control using spin-orbit interaction
- Implications of the spin-orbit interaction for singlet-triplet qubits in silicon

Part 2: Spin-orbit interaction in silicon dots

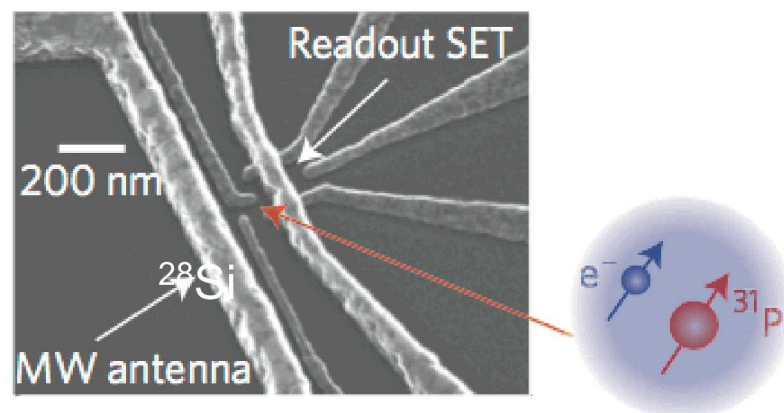


$$\begin{array}{l}
 |\uparrow\rangle \text{ — } \\
 E_Z = g_1 \mu_B B_{\text{ext}} \\
 |\downarrow\rangle \text{ — }
 \end{array}
 \quad
 \begin{array}{l}
 \otimes B_{\text{ext}} \\
 |\uparrow\rangle \text{ — } \\
 E_Z = g_2 \mu_B B_{\text{ext}} \\
 |\downarrow\rangle \text{ — }
 \end{array}$$

$$\Delta E_Z = (\Delta g) \mu_B B_{\text{ext}}$$

Can the best characteristics of qubit systems be combined?

Donor (D) qubits are high fidelity...



Electron spin 1/2:

$$T_2^* = 270 \mu\text{s} [1]$$

$$T_2, \text{CPMG} = 0.98 \text{ s} [3]$$

$$F_{\text{prep/read}} = 97\% [4, 1]$$

$$F_{\text{control}} = 99.95\%$$

Nuclear spin 1/2:

$$T_2^* = 600 \text{ ms} [3]$$

$$T_2, \text{CPMG} = 36.5 \text{ s} [3]$$

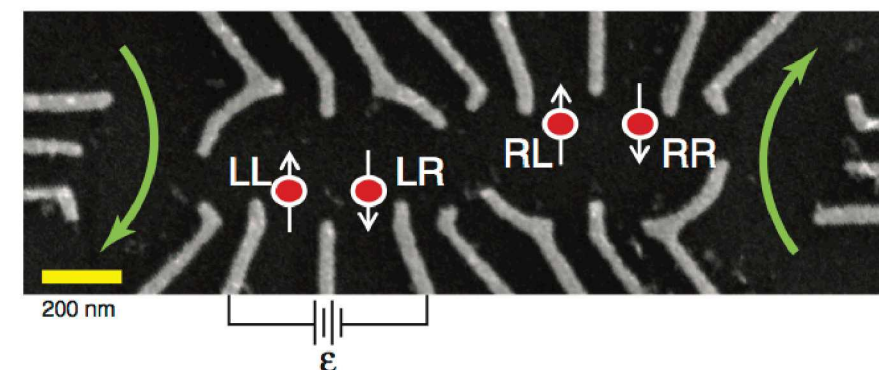
$$F_{\text{prep/read}} = 99.995\%$$

$$F_{\text{control}} = 99.99\%$$

But they are not easily coupled together!

- [1] Muhonen, J. T. et al., Nat Nano 9, 986–991 (2014).
- [2] Muhonen, J. T. et al., J. of Physics: Cond. Mat. 27, 154205 (2015).
- [3] A. Morello, unpublished.
- [4] L. A. Tracy et al, App. Phys. Lett. 108, 063101 (2016).

Quantum dot (QD) qubits are tunable...



Singlet-triplet spin qubit in GaAs:

- All-electrical control [6]
- Two-qubit gates realized [5,6,9]

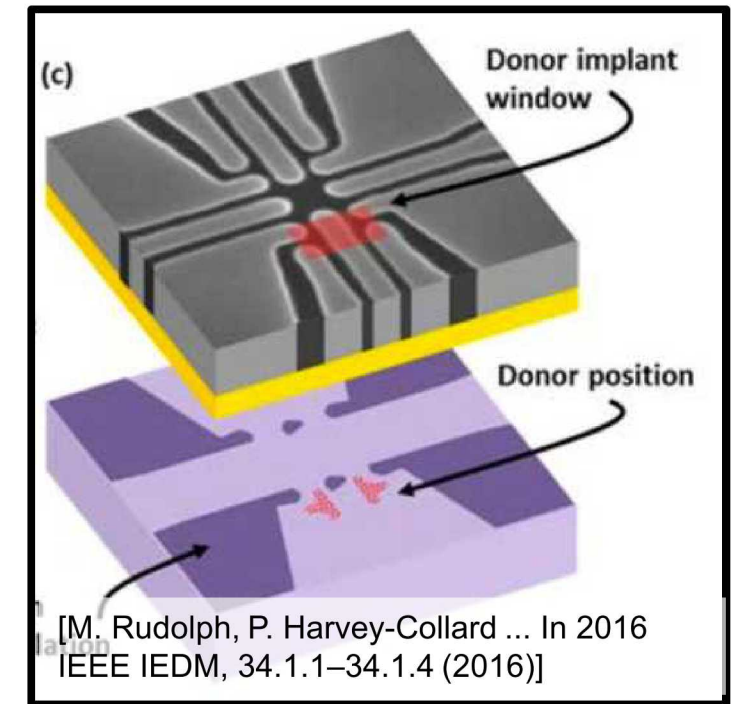
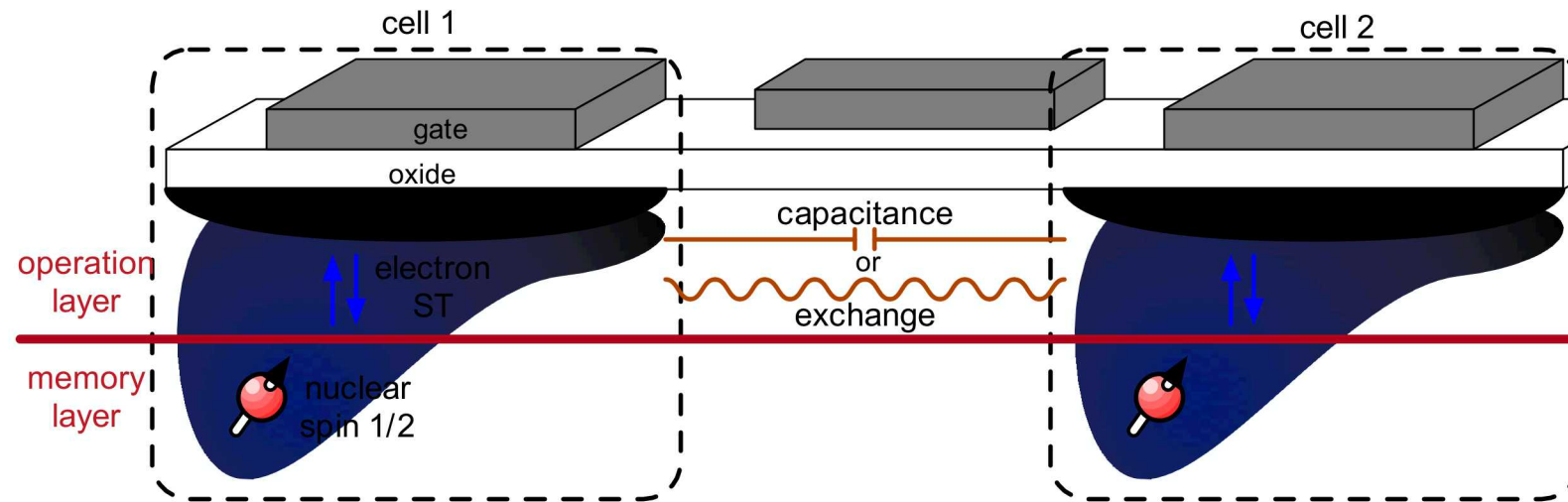
... many other GaAs and Si realizations now!

But they are not quite as high fidelity!

- [5] Shulman, M. D. et al. Science 336, 202–205 (2012).
- [6] Petta, J. R. et al. Science 309, 2180–2184 (2005).
- [7] Veldhorst, M. et al. Nat Nano 9, 981–985 (2014).
- [8] Veldhorst, M. et al. Nature 526, 410–414 (2015).
- [9] Nichol, J. M. et al. npj Quantum Information 3, 3 (2017).

The MAJIQ system

MOS (M) hyperfine (A) exchange (J) nuclear spin (I) qubit (Q)

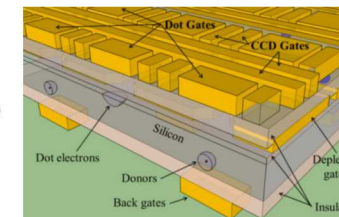


Advantages

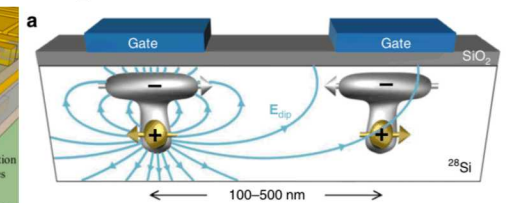
- **Fast and stable** rotation speed: Contact hyperfine interaction $A/2 = 57$ MHz gives **9 ns π rotation (potentially faster)**.
- Compact design: No need for **nuclear field bath, spin-orbit int., micromagnet**. Just **one QD**.
- ST states: good probe of coherence between the QD and the D.
- Get a **nuclear spin** for free. Nuclear spin qubit is the **"best" single qubit** in solid state.

[P. Harvey-Collard ... Nat. Commun. 8, 1029 (2017)]

Other proposals



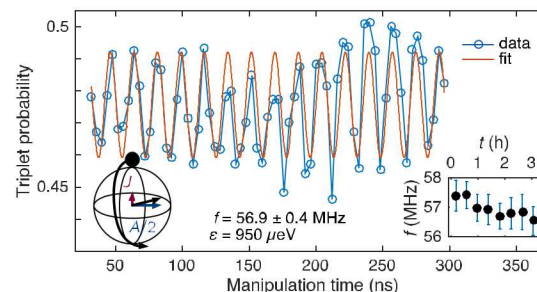
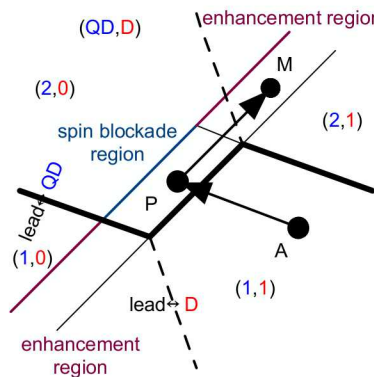
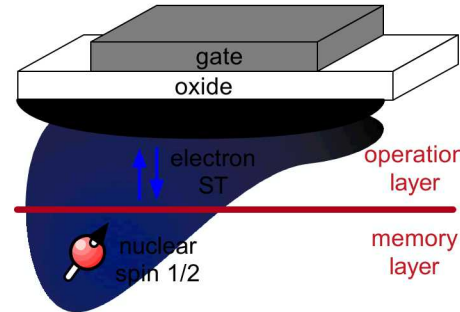
[G. Pica ... Phys. Rev. B 93, 035306 (2016)]



[G. Tosi ... Nat. Commun. 8, 450 (2017)]

Part 1: Outline

- Donor-Dot system
- Valley splitting, singlet-triplet splitting, coupling to donors, and 4-electron trick
- Enhanced latching readout
- Singlet-triplet qubit (hyperfine-driven rotations)



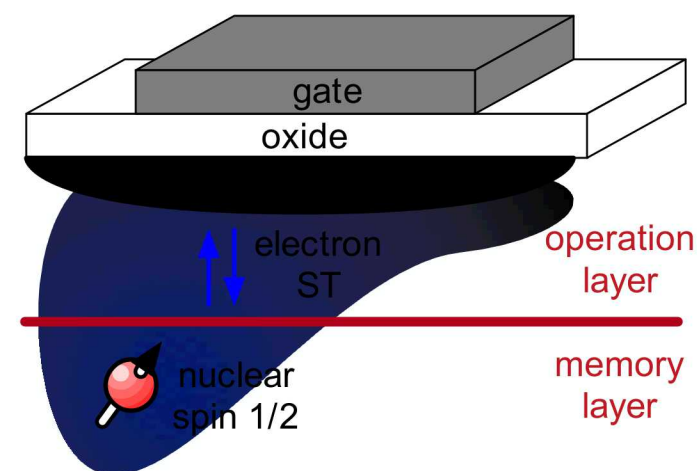
References

1. Gamble, J. K., **Harvey-Collard, P.** et al. "Valley splitting of single-electron Si MOS quantum dots" App. Phys. Lett. 109, 253101 (2016).
2. **Harvey-Collard, P.** et al., "High-fidelity single-shot readout for a spin qubit via an enhanced latching mechanism," Phys. Rev. X 8, 021046 (2018).
3. **Harvey-Collard, P.** et al. "Coherent coupling between a quantum dot and a donor in silicon" Nat. Commun. 8, 1029 (2017).
4. M. Rudolph, **P. Harvey-Collard** et al. "Coupling MOS quantum dot and phosphorous donor qubit systems," in 2016 IEEE IEDM (2016) pp. 34.1.1–34.1.4.

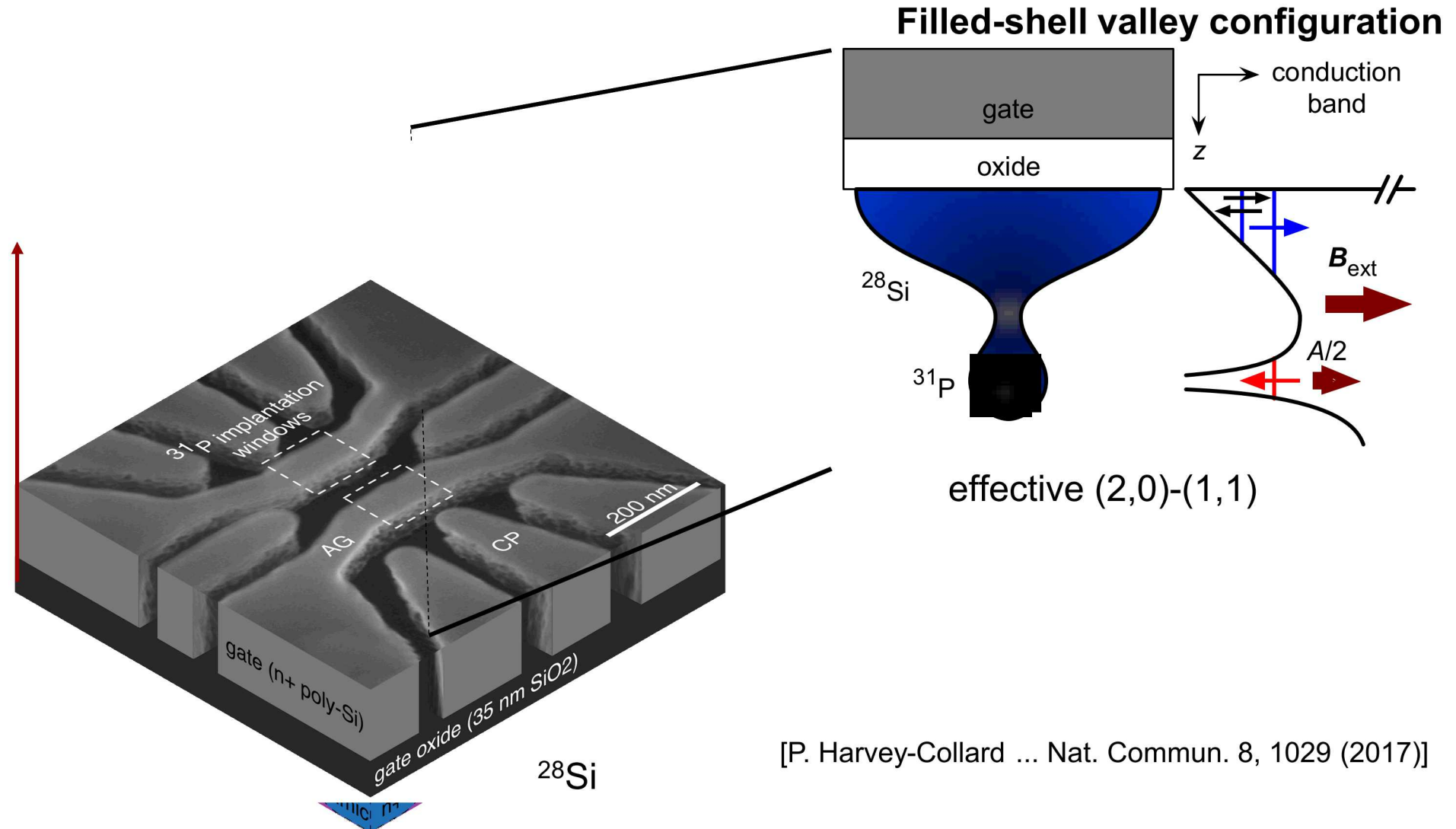
- Donor-Dot system
- Valley splitting, singlet-triplet splitting, coupling to donors, and 4-electron trick
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Part 1:

Donor-dot system and related topics

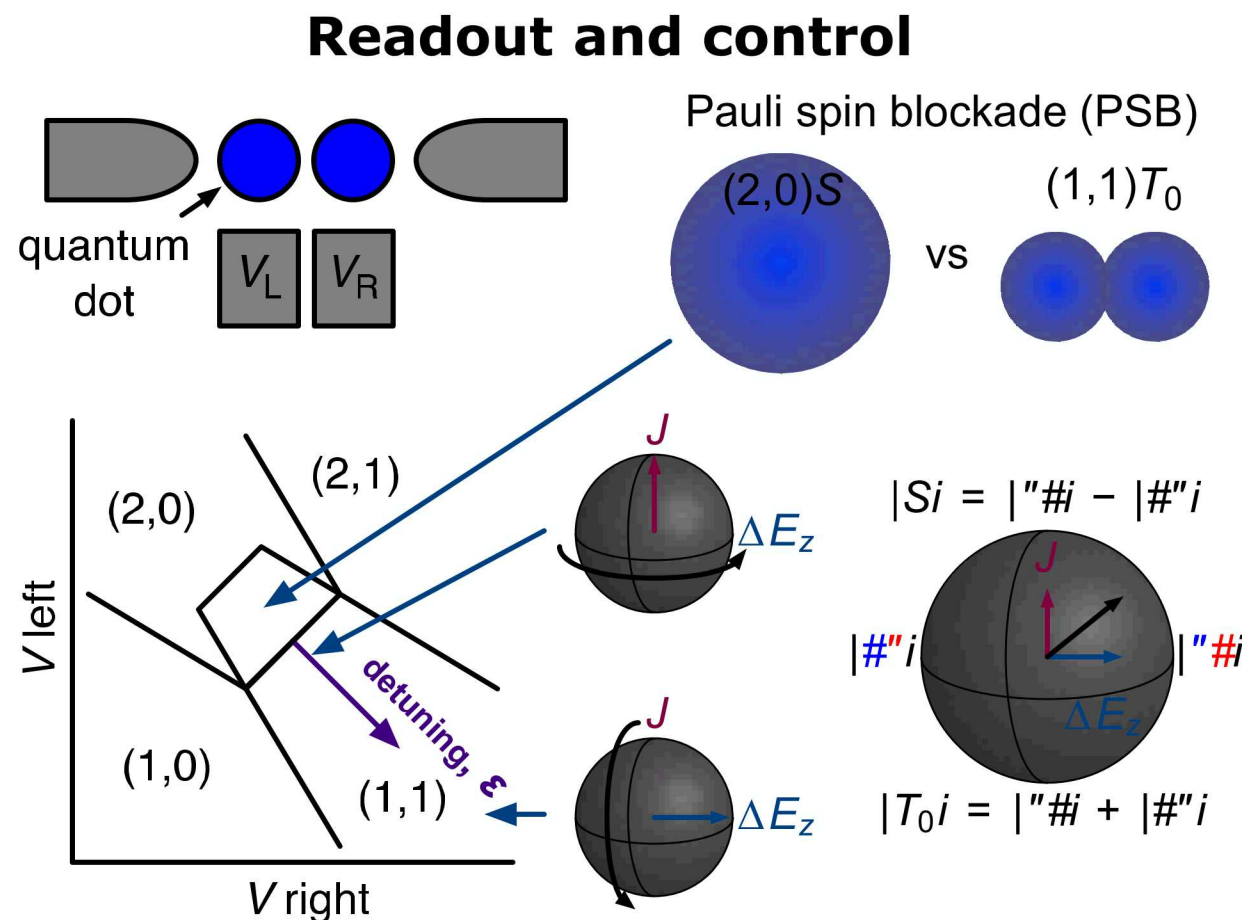


The "gated wire" device



[P. Harvey-Collard ... Nat. Commun. 8, 1029 (2017)]

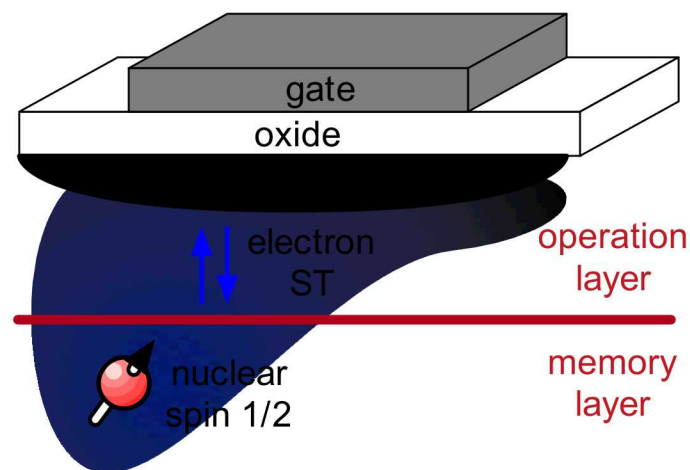
The singlet-triplet (ST) qubit



- Donor-Dot system
- Valley splitting, singlet-triplet splitting, coupling to donors, and 4-electron trick
- Enhanced latching readout
- Singlet-triplet qubit (hyperfine-driven rotations)

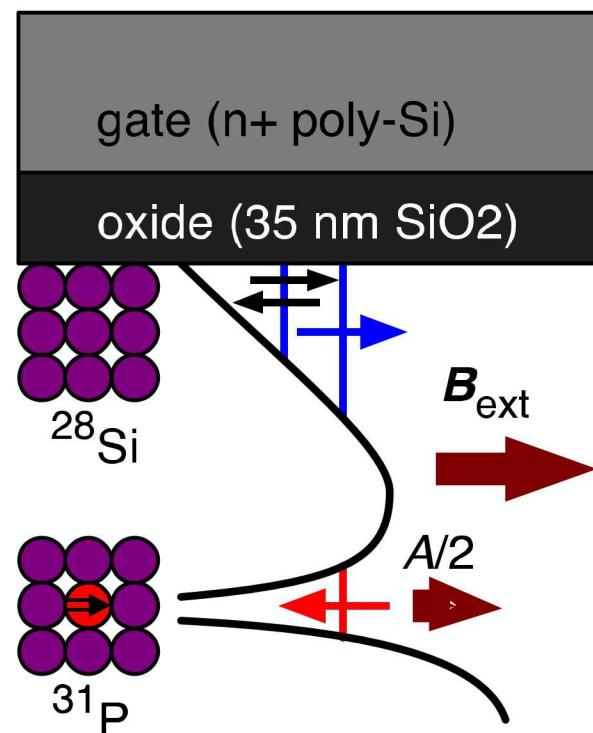
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Donor-dot system and related topics

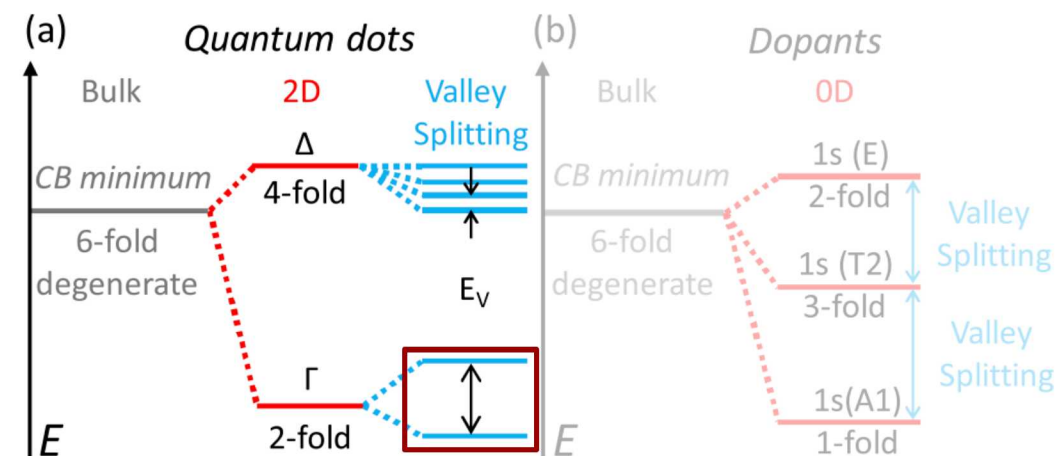


Valley splitting in Si

Filled-shell valley configuration



Valley degeneracy is lifted because of the strong confinement at the Si/SiO₂ interface [1]

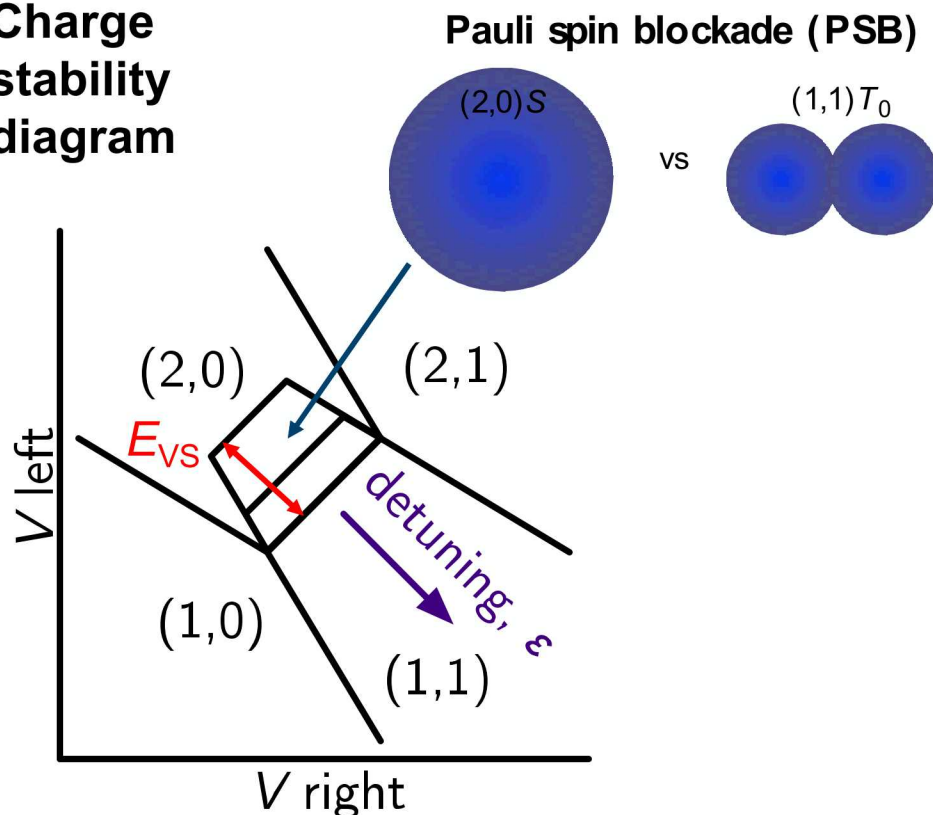


[1] Zwanenburg, F. A. ... Rev. Mod. Phys. 85, 961 (2013).

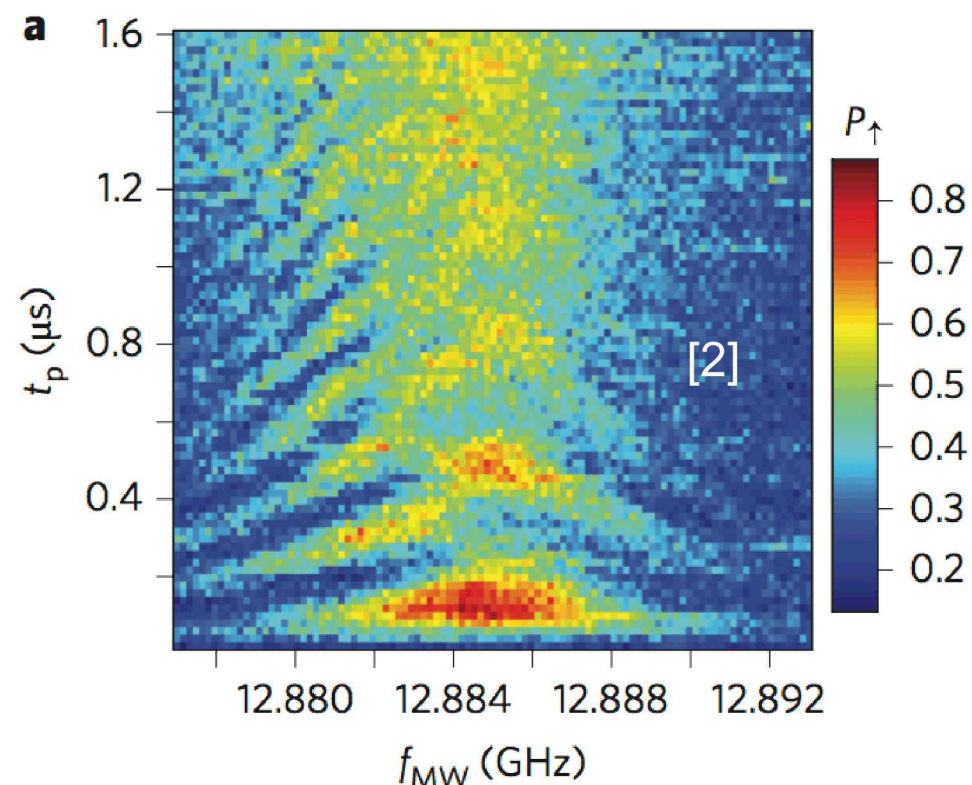
Small valley splitting is a problem

Spin blockade window can be unpractically small

Charge stability diagram



Valleys can have different g -factors [1, 2]



[1] Eng, K. ... Science Advances 1, e1500214 (2015).

[2] Kawakami, E. ... Nat Nano 9, 666–670 (2014).

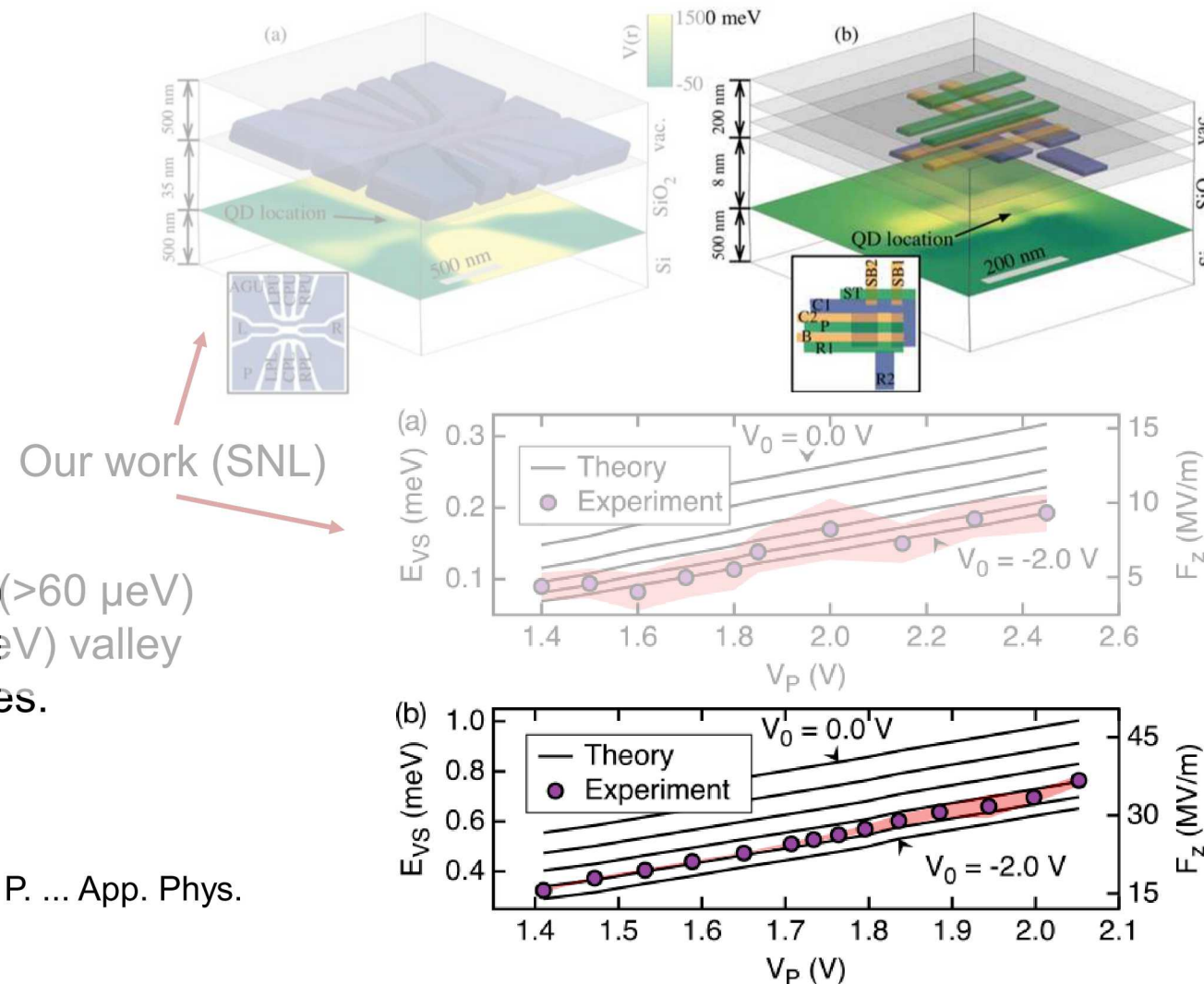
Valley splitting is tunable in MOS

Reproduced in
our devices [1].

Theoretical
understanding
[1].

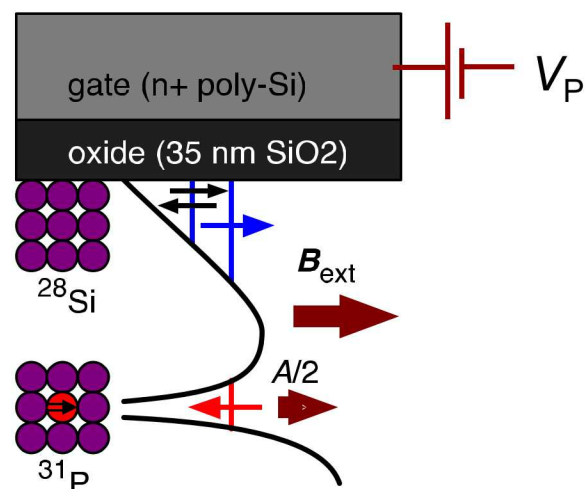
We find consistently large ($>60 \mu\text{eV}$)
and tunable (60 to $>250 \mu\text{eV}$) valley
splittings in our SNL devices.

[1] Gamble, J. K., Harvey-Collard, P. ... App. Phys.
Lett. 109, 253101 (2016).



Valley splitting in Si – donors

Filled-shell valley configuration



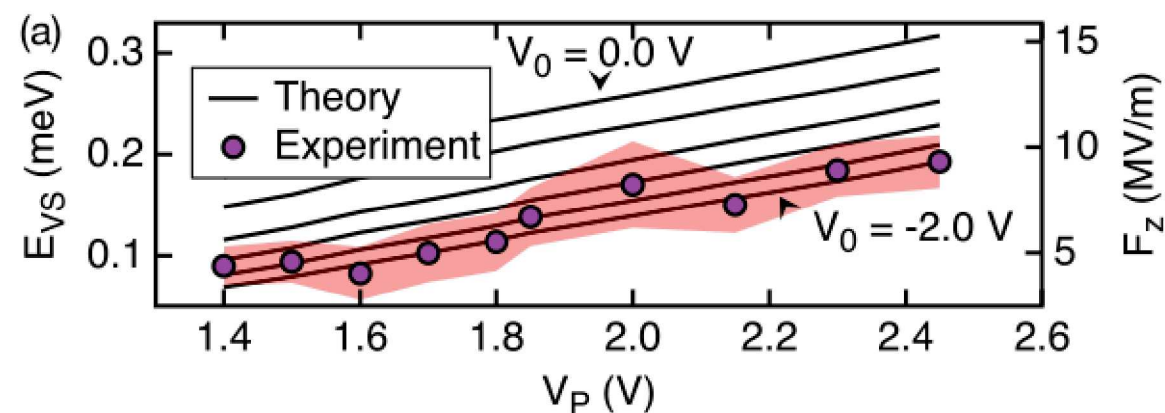
Strong tunnel coupling to donor

=

Low electric field (V_P)



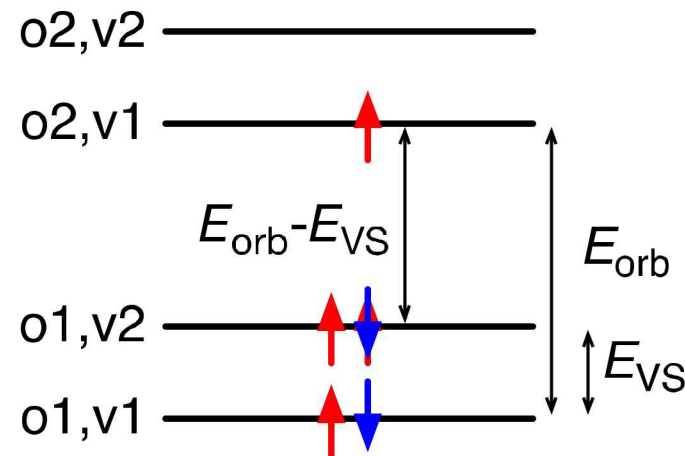
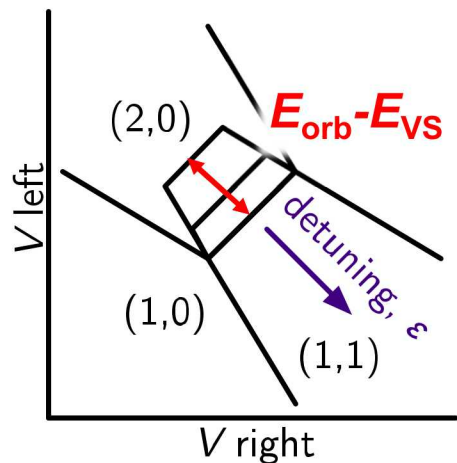
Limits the valley splitting!



Tuning the ST splitting via shell engineering

Give up? Or play some tricks!

- Spin singlets have a symmetric valley-orbit wavefunction.
- Trick: fill the lowest valley with a singlet. Then forget about it!



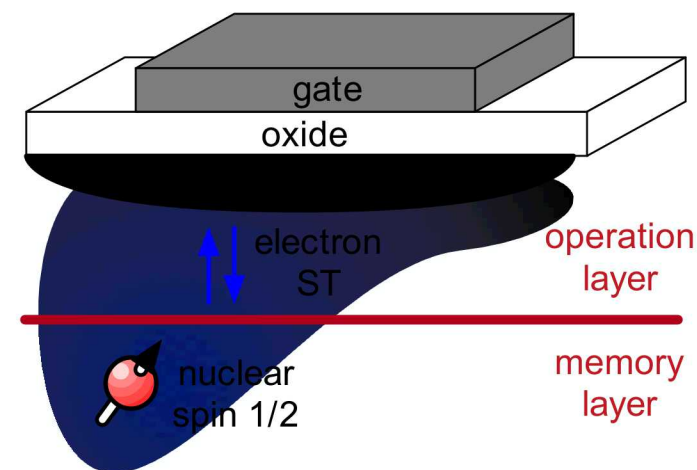
4-electron trick

- Proposed that this trick can
 - Circumvent small ST splitting
 - Change QD size to adjust tunnel coupling to donors
- Introduced in 2015 (arXiv) and [Harvey-Collard, P. ... Nat. Commun. 8, 1029 (2017).]
- Now used by other groups (HRL, U Wisc, ...)

- Donor-Dot system
- Valley splitting, singlet-triplet splitting, coupling to donors, and 4-electron trick
- Enhanced latching readout
- Singlet-triplet qubit (hyperfine-driven rotations)

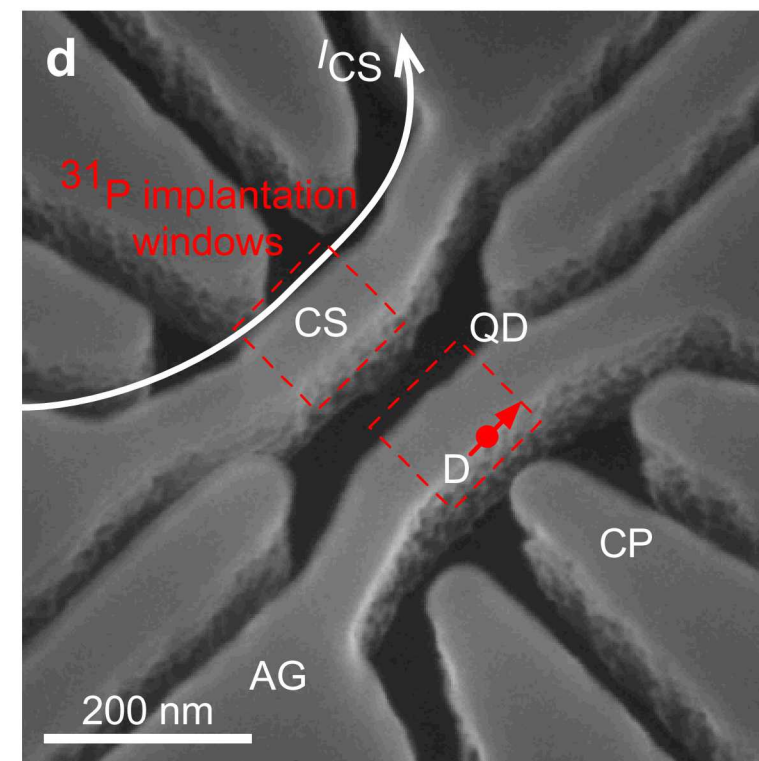
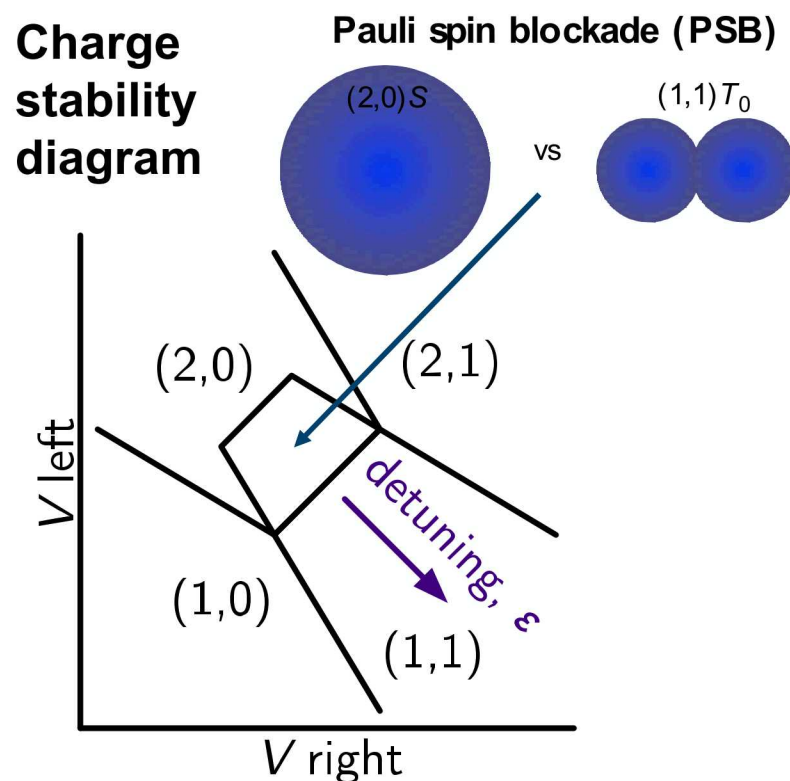
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Donor-dot system and related topics

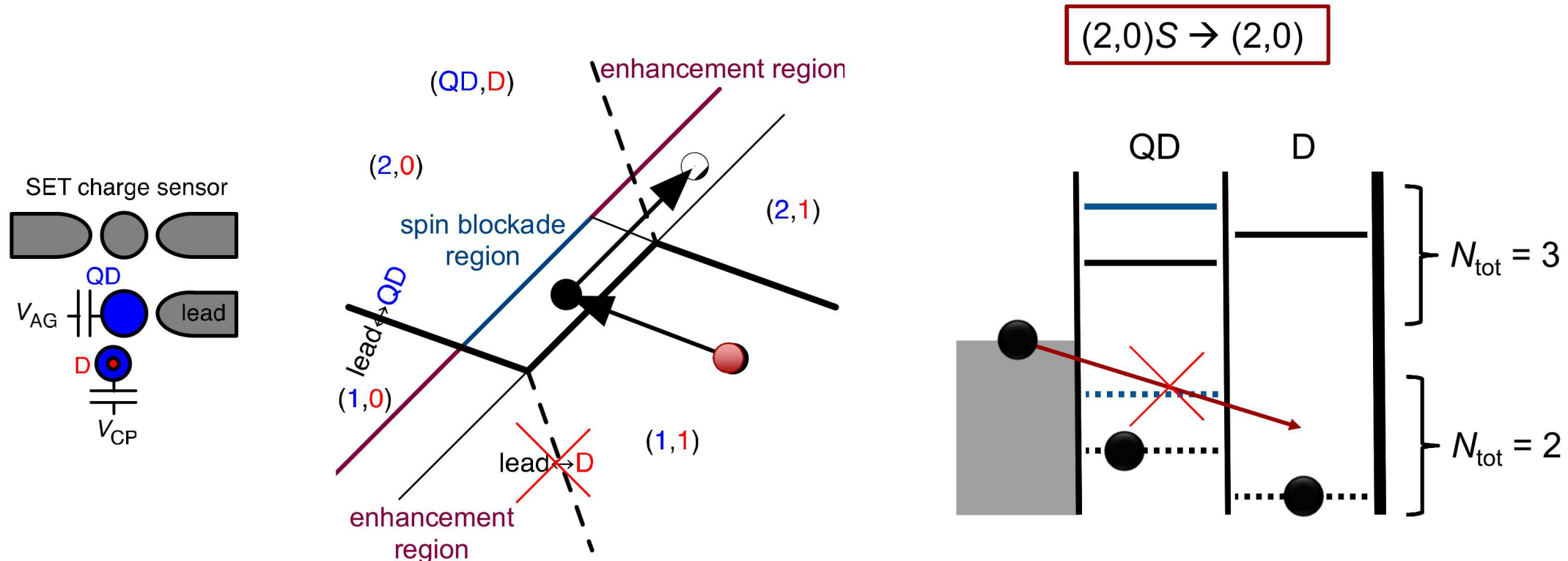


Issues with spin readouts

Example: Charge sensor not necessarily aligned with sensed dipole

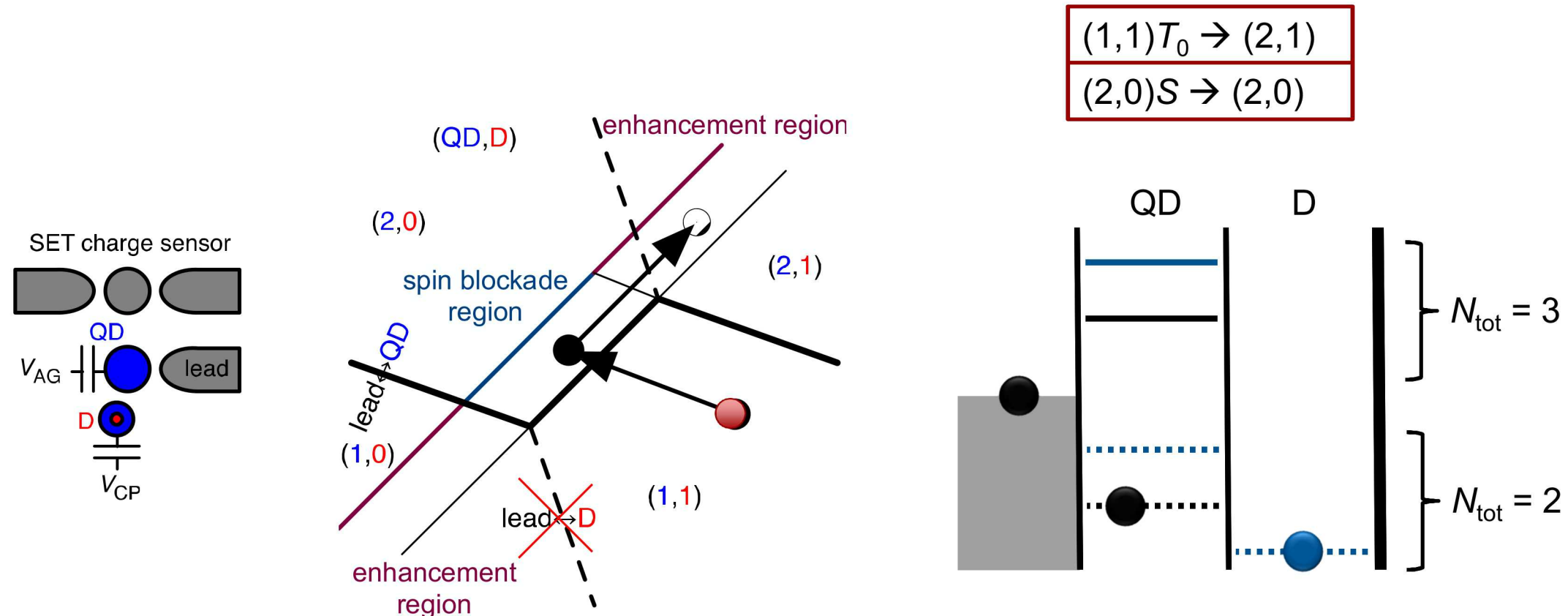


Readout concept (illustrated)

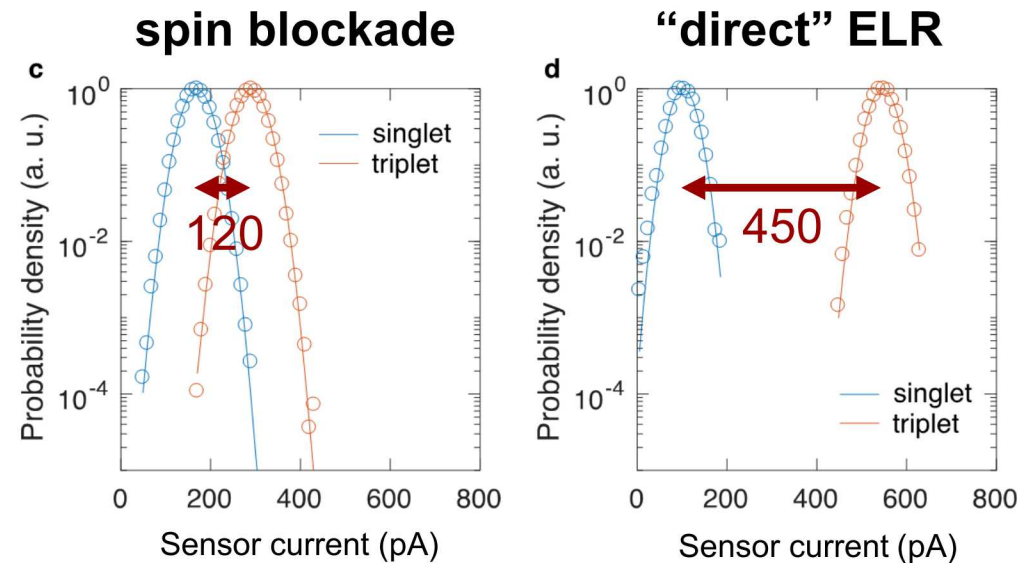




Readout concept (illustrated)

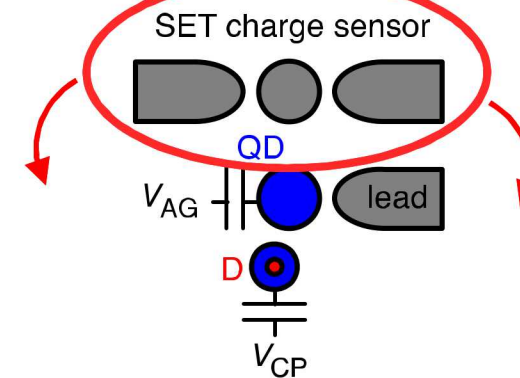


Enhanced Latching Readout enhances the signal-to-noise ratio



Noise is the same.

Relaxes layout constraints



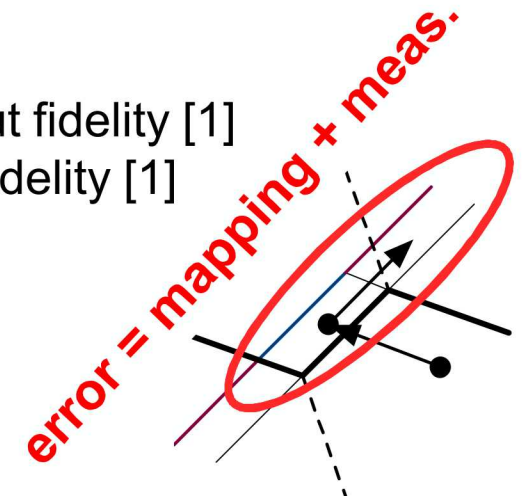
[1] P. Harvey-Collard ... Phys. Rev. X 8, 021046 (2018).

Enhances lifetime [1]

99.91% avg. charge readout fidelity [1]

99.86% avg. spin readout fidelity [1]

Still record to date?



[Inspired from:] S. A. Studenikin ... App. Phys. Lett. 101, 233101 (2012).
 [Other work:] Mason, J. D. ... Phys. Rev. B 92, 125434 (2015).
 Petersson, K. D. ... Phys. Rev. Lett. 105, 246804 (2010).



Enhanced Latching Readout (ELR) Comparison

Comparison of electron readouts

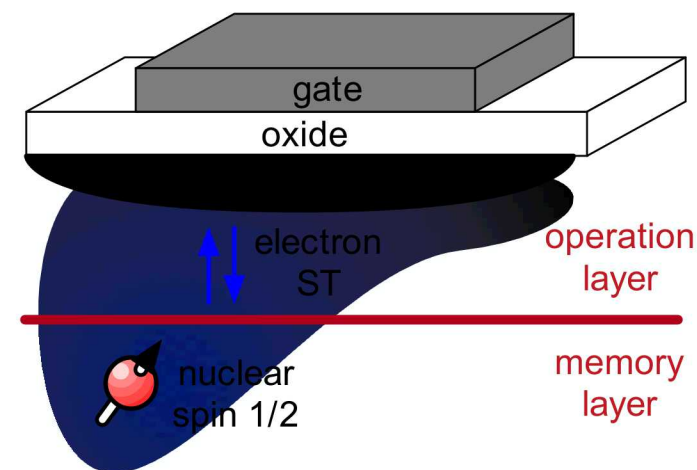
	Spin 1/2	Singlet-triplet	This work (ELR)
Mechanism	Tunneling (stochastic)	Spin blockade (deterministic)	Spin blockade (deterministic)
Speed	Slow	Fast	Fast
Signal	1 e	~0.2 e	1 e
Lifetime	Good, but long waits	~Ok (Si), Bad (GaAs)	Good

If you have ***any*** singlet-triplet system, this can work!

- Donor-Dot system
- Valley splitting, singlet-triplet splitting, coupling to donors, and 4-electron trick
- Enhanced latching readout
- Singlet-triplet qubit (hyperfine-driven rotations)

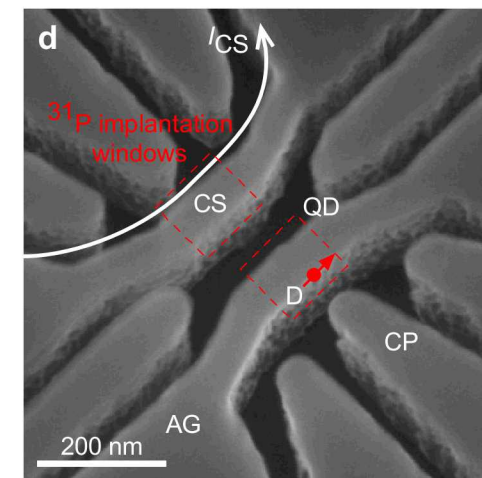
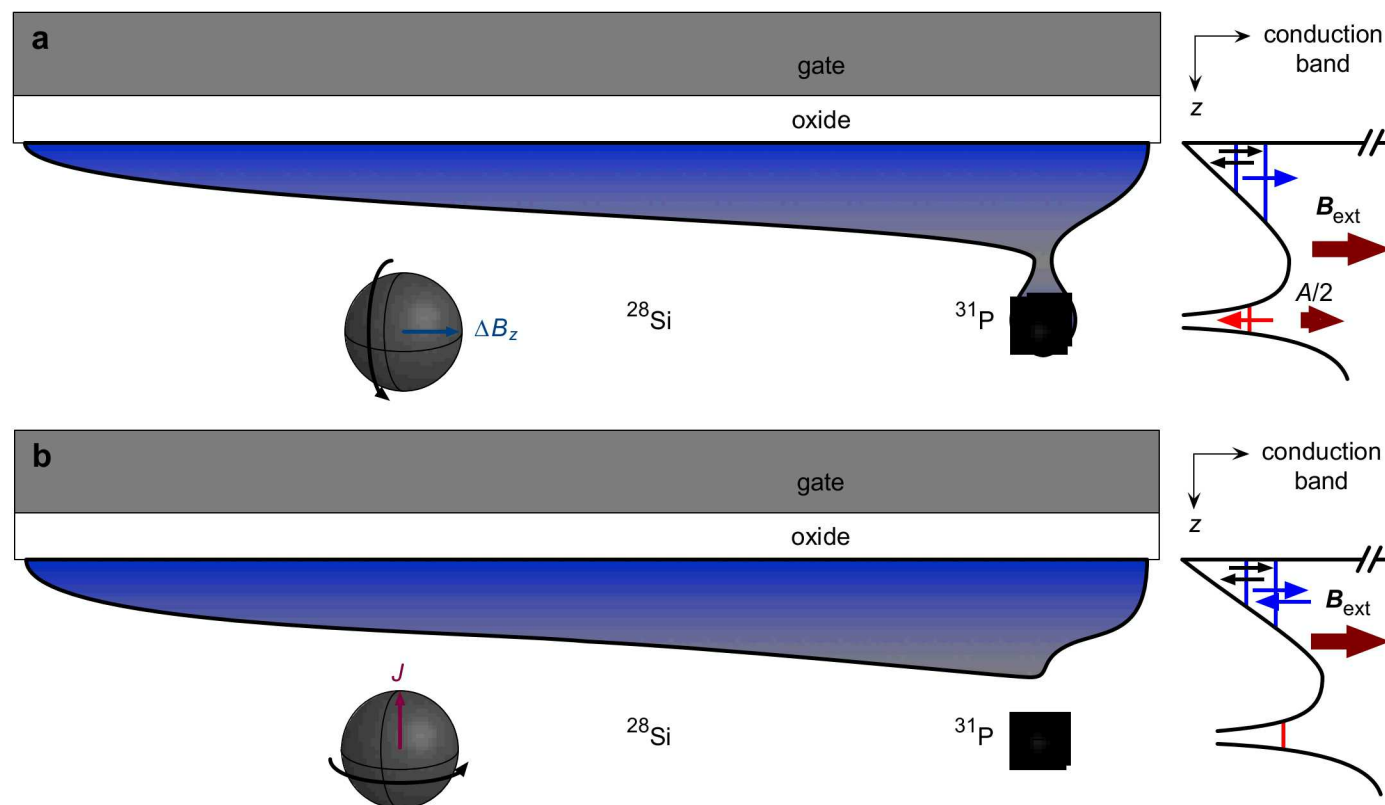
Part 1:

Donor-dot system and related topics



The MAJIQ system

MOS (M) hyperfine (A) exchange (J) nuclear spin (I) qubit (Q)



$$\langle \downarrow \downarrow | A \mathbf{S} \cdot \mathbf{I} | \downarrow \downarrow \rangle = \frac{A}{2} S$$

For electron on donor:

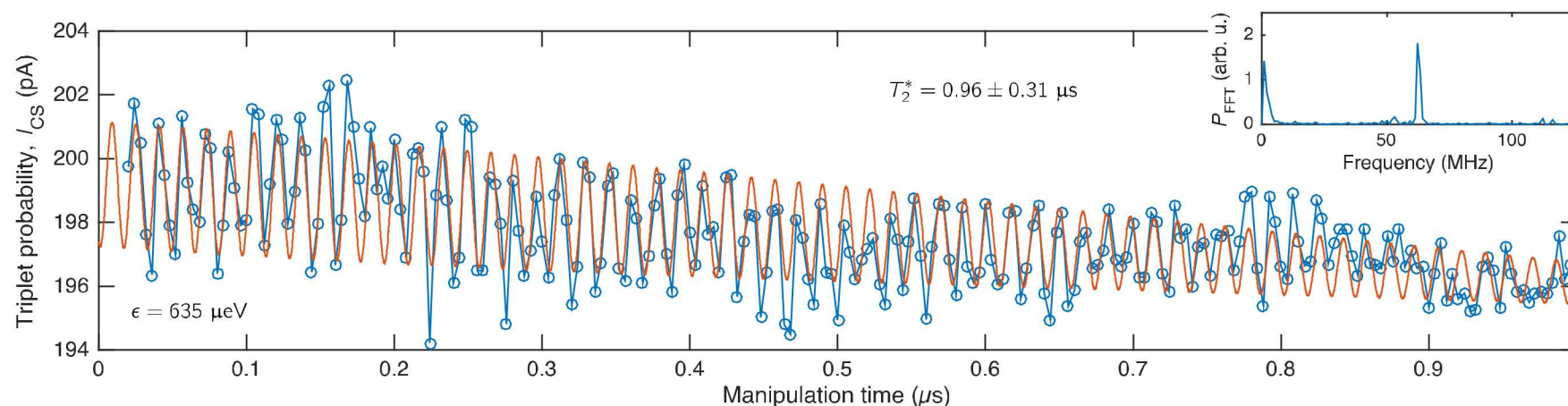
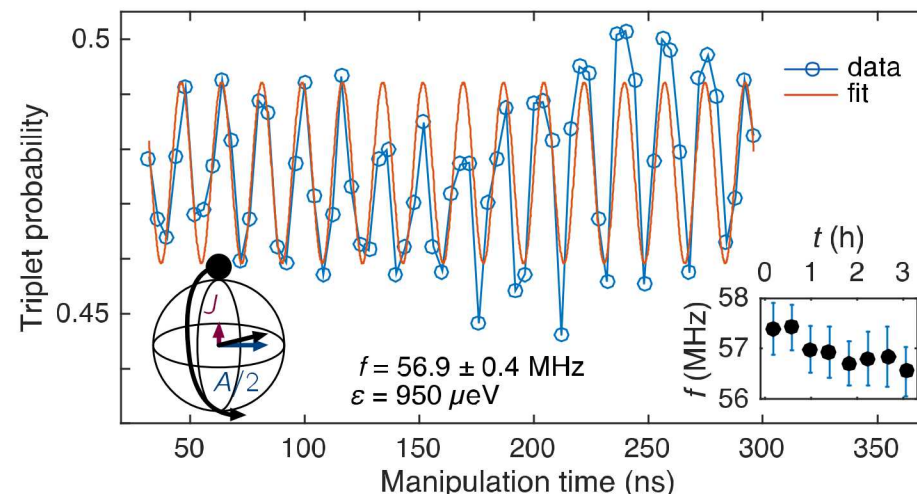
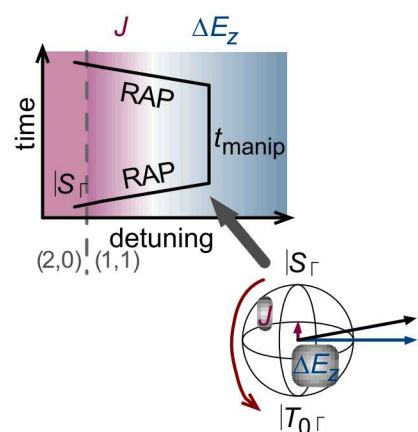
$$\hat{H} = (g\mu_B/\hbar) \mathbf{S} \cdot \mathbf{B}_{\text{ext}} + A \mathbf{S} \cdot \mathbf{I}$$

Singlet-Triplet Hamiltonian:

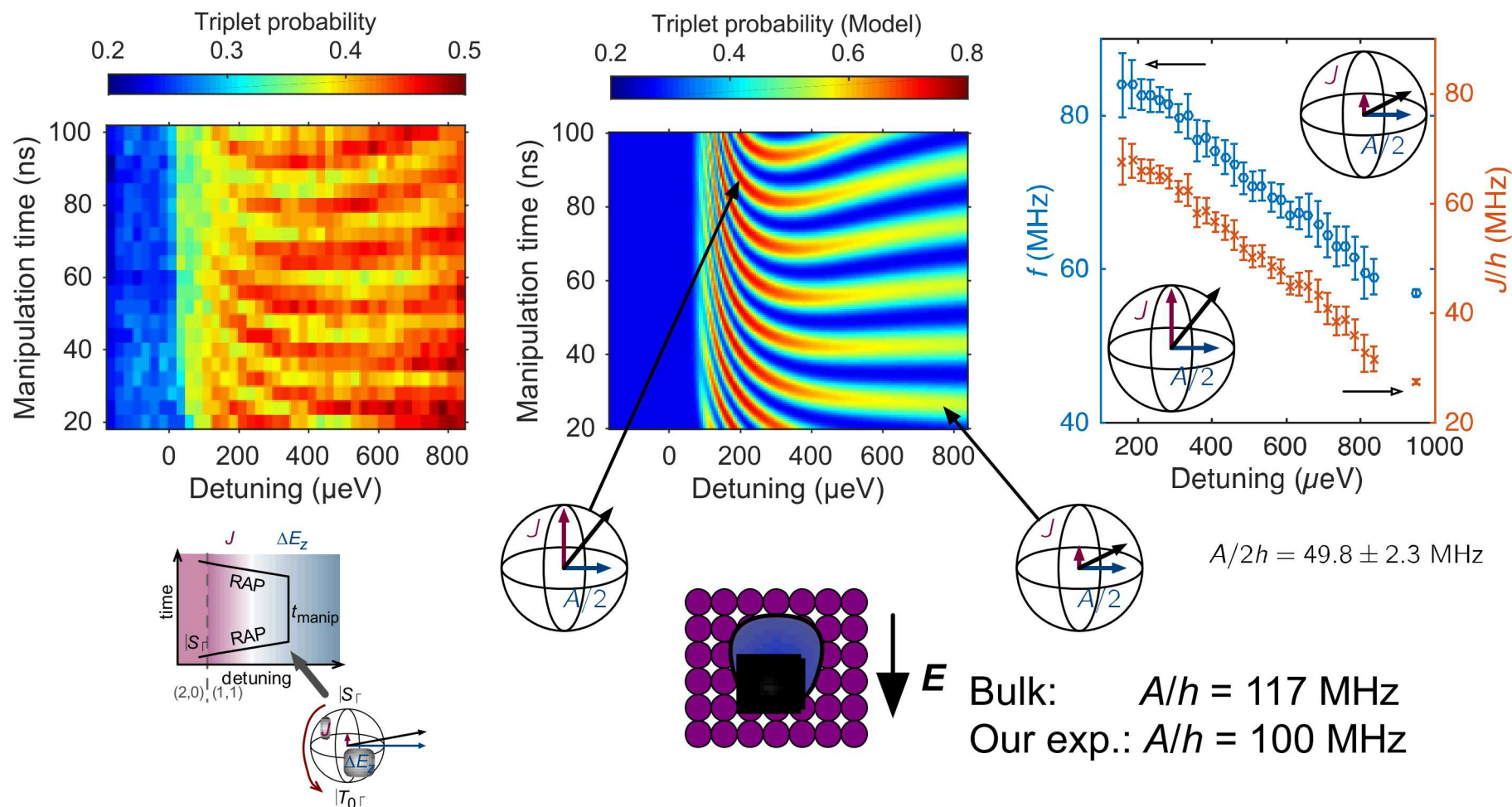
$$\hat{H}_{\text{ST}} = J(\epsilon) \hat{\sigma}_z + \Delta B_z(\epsilon) \hat{\sigma}_x$$

State manipulation

Singlet-triplet rotations driven by a single ^{31}P nucleus

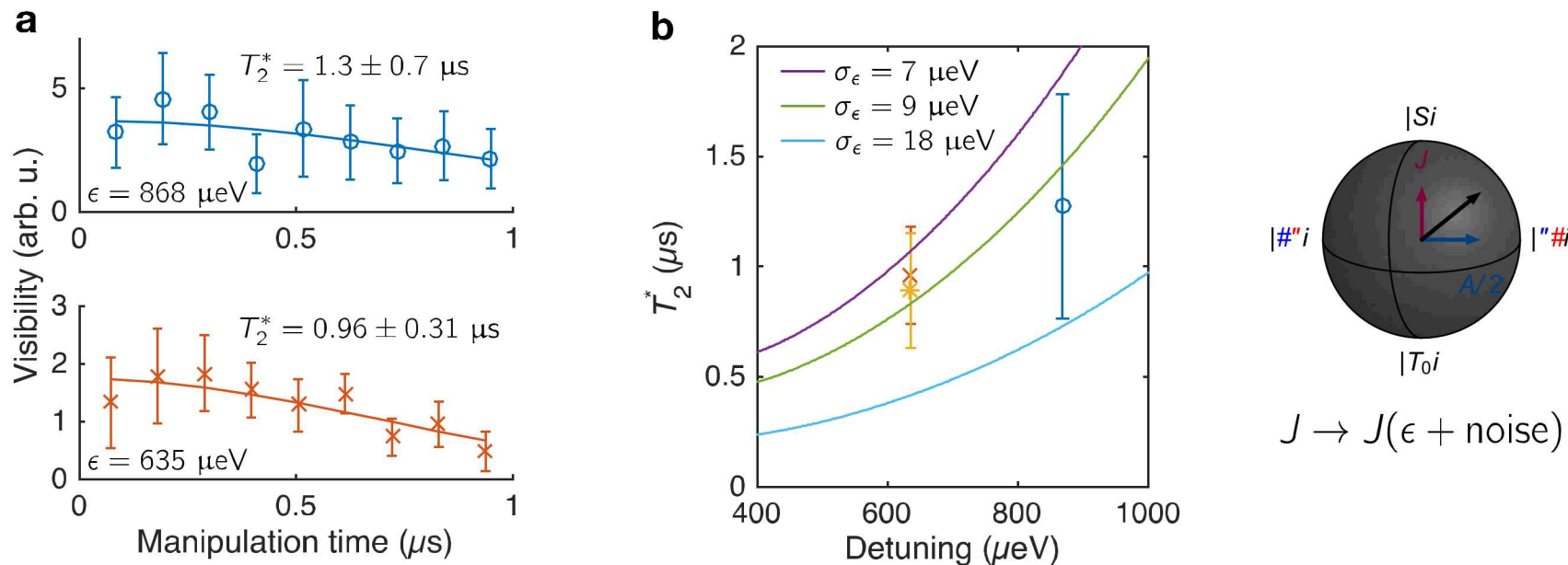


Detuning dependence of exchange interaction





Coherence time and charge noise



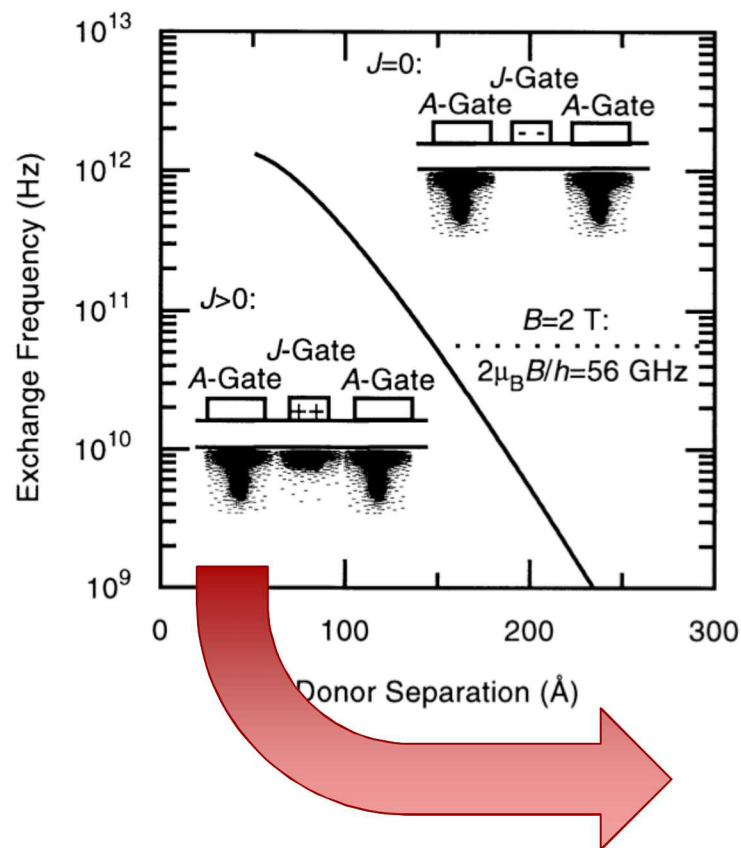
Quantitative measure
of charge noise in
Si-MOS

Reference	Petersson <i>et al.</i>	Shi <i>et al.</i>	Eng <i>et al.</i>	This work	Jock <i>et al.</i>
Material	GaAs/AlGaAs	Si/SiGe	Si/SiGe	Si (MOS)	Si (MOS)
$2\sigma_\epsilon$ (μeV)	7.4	10	9.2	18	4
$k_B T_e$ (μeV)	6.9	12	6.9 – 8.6	18	13

[1] P. Harvey-Collard ... Nat. Commun. 8, 1029 (2017)
[See also:] Rudolph, M. ... 2016 IEEE IEDM, 34.1.1–34.1.4 (2016).
R. M. Jock ... Nat. Commun. 9, 1768 (2018).

Perspectives

Modern implementation of Kane's original vision

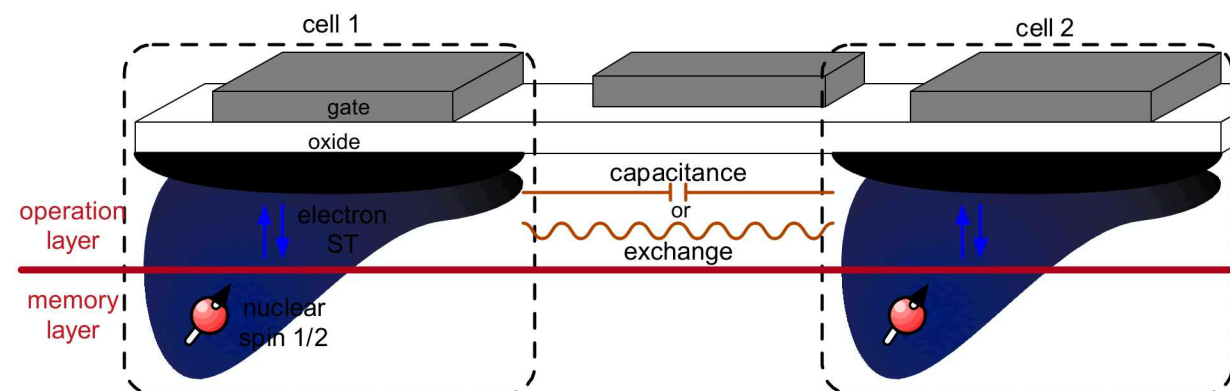


[B. Kane ...
Nature 393,
133–137
(1998)]

Over 19 years of QD-D schemes!

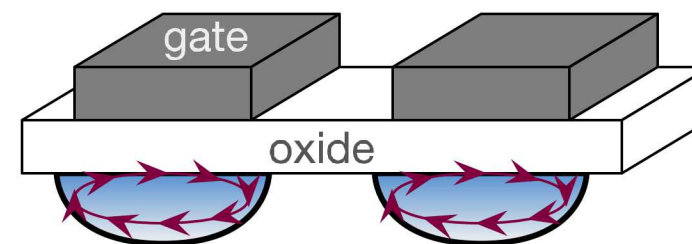
Future work:

- High-fidelity qubit operation (AC, etc.)
- Strong hyperfine donors species
- Two-MAJQ coupling
- Nuclear spin addressing



- Universal control using spin-orbit interaction
- Implications of the spin-orbit interaction for singlet-triplet qubits in silicon

Part 2: Spin-orbit interaction in silicon dots

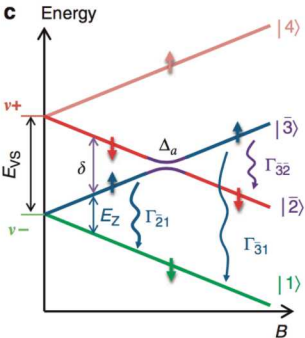


$$\begin{array}{l}
 |\uparrow\rangle \text{ —} \\
 E_Z = g_1 \mu_B B_{\text{ext}} \\
 |\downarrow\rangle \text{ —}
 \end{array}
 \quad
 \begin{array}{l}
 \otimes B_{\text{ext}} \\
 |\uparrow\rangle \text{ —} \\
 E_Z = g_2 \mu_B B_{\text{ext}} \\
 |\downarrow\rangle \text{ —}
 \end{array}$$

$$\Delta E_Z = (\Delta g) \mu_B B_{\text{ext}}$$

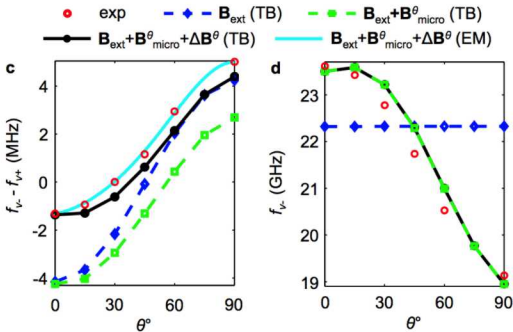
Spin-orbit interaction in silicon – Other work

Spin relaxation hot-spot in MOS



C. H. Yang ... Nat. Commun. 4, 2069 (2013).
X. Hao ... Nat. Commun. 5, 3860 (2014).

Spin resonance frequency anisotropy in Si/SiGe

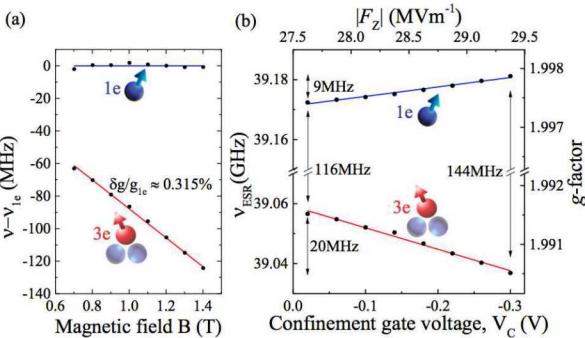


R. Ferdous ... npj Quantum Information 4, 26 (2018).

Broken symmetry at Si/material interface:

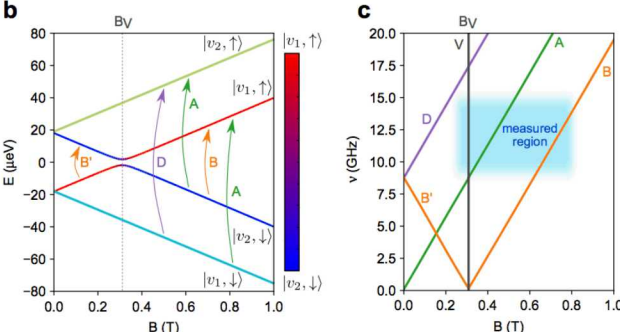
U. Rössler ... Solid State Comm. 121, 313 (2002).
L. E. Golub ... Phys. Rev. B 69, 115333 (2004).
M. Prada ... New J. of Phys. 13, 013009 (2011).

Spin resonance frequency in MOS



M. Veldhorst ... Phys. Rev. B 92, 201401 (2015).
R. Ferdous ... Phys. Rev. B 97, 241401 (2018).

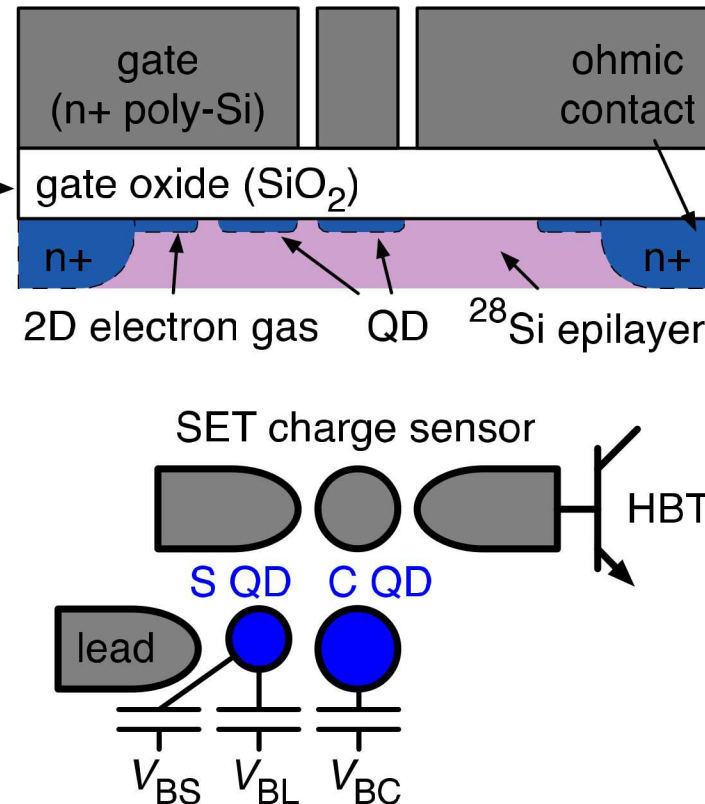
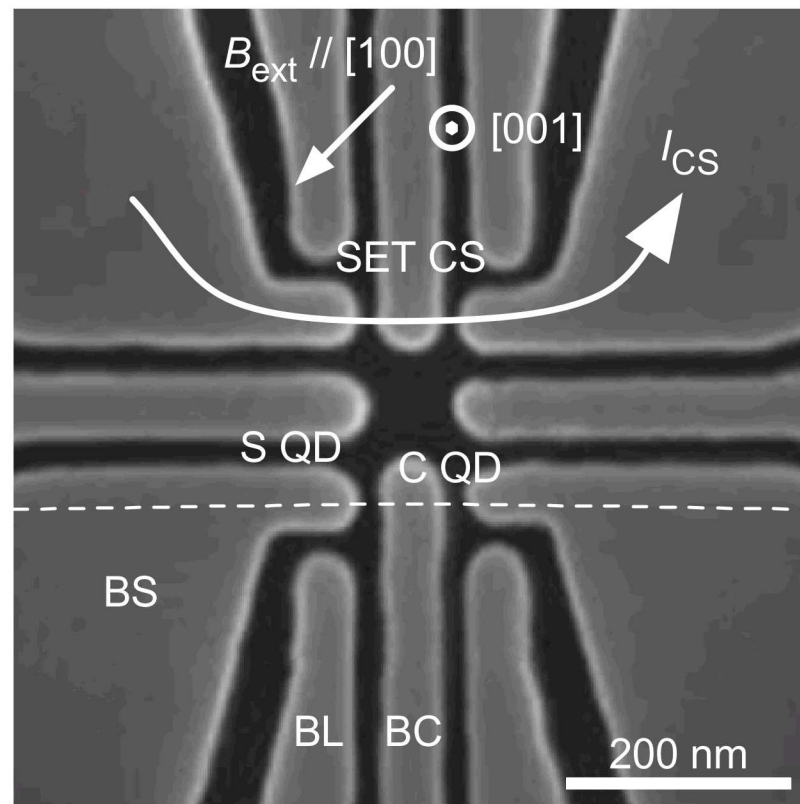
Spin-orbit-driven spin resonance in a MOS finFET



A. Corna ... npj Quantum Information 4, 6 (2018).

Simple CMOS-compatible qubit

**Pure CMOS design [1, 2, 3]
Isotopically-enriched silicon 28**



Split Wire Accumulation Gate (SWAG)

First* spin-orbit-driven e⁻ qubit in silicon.**

*: See also [Corna et al. (2018)]

** : Excluding artificial spin-orbit.

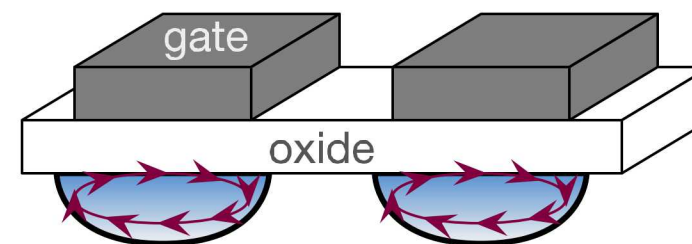
[1] P. Harvey-Collard et al. All-electrical universal control of a double quantum dot qubit in silicon MOS. *In 2017 IEEE IEDM* (2017).

[2] R. M. Jock, N. T. Jacobson, P. Harvey-Collard et al. "A silicon metal-oxide-semiconductor electron spin-orbit qubit," *Nat. Commun.* 9, 1768 (2018).

[3] P. Harvey-Collard et al., "Implications of the spin-orbit interaction for singlet-triplet qubits in silicon," *arXiv:1808.07378* (2018).

- Universal control using spin-orbit interaction
- Implications of the spin-orbit interaction for singlet-triplet qubits in silicon

Part 2: Spin-orbit interaction in silicon dots



$$\begin{array}{l} |\uparrow\rangle \text{ --- } \\ E_Z = g_1 \mu_B B_{\text{ext}} \\ |\downarrow\rangle \text{ --- } \end{array} \quad \begin{array}{l} \otimes B_{\text{ext}} \\ |\uparrow\rangle \text{ --- } \\ E_Z = g_2 \mu_B B_{\text{ext}} \\ |\downarrow\rangle \text{ --- } \end{array}$$

$$\Delta E_Z = (\Delta g) \mu_B B_{\text{ext}}$$

Initializing the qubit: Slow/Rapid Adiabatic Passage (SAP/RAP)

SAP: Slow Adiabatic Passage

Adiabatically with respect to both spin and charge.

$$|S\rangle \leftrightarrow |\uparrow\downarrow\rangle$$

$$|T_0\rangle \leftrightarrow |\downarrow\uparrow\rangle$$

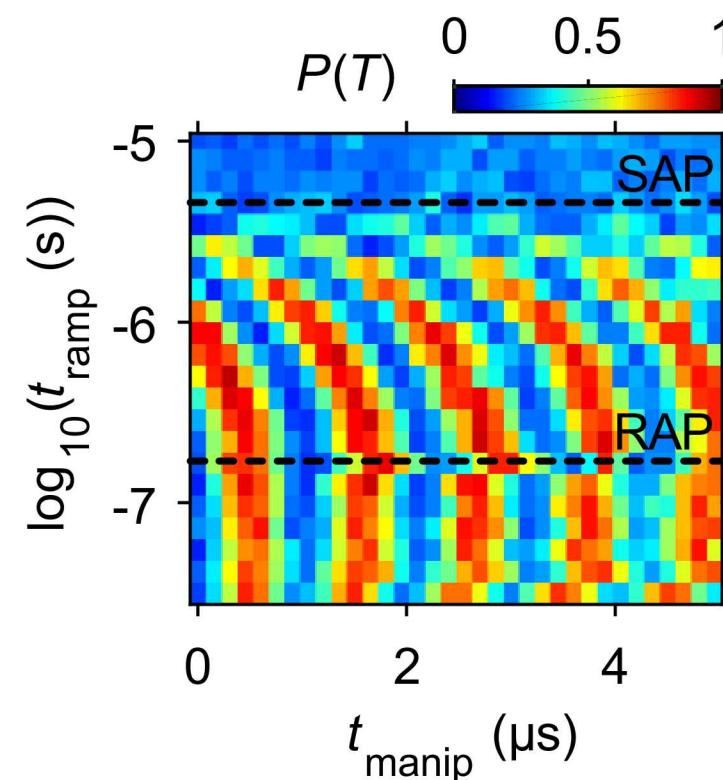
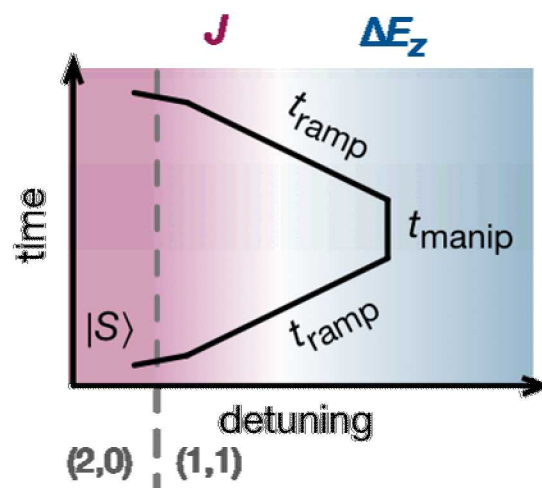
RAP: Rapid Adiabatic Passage

Adiabatically with respect to charge.

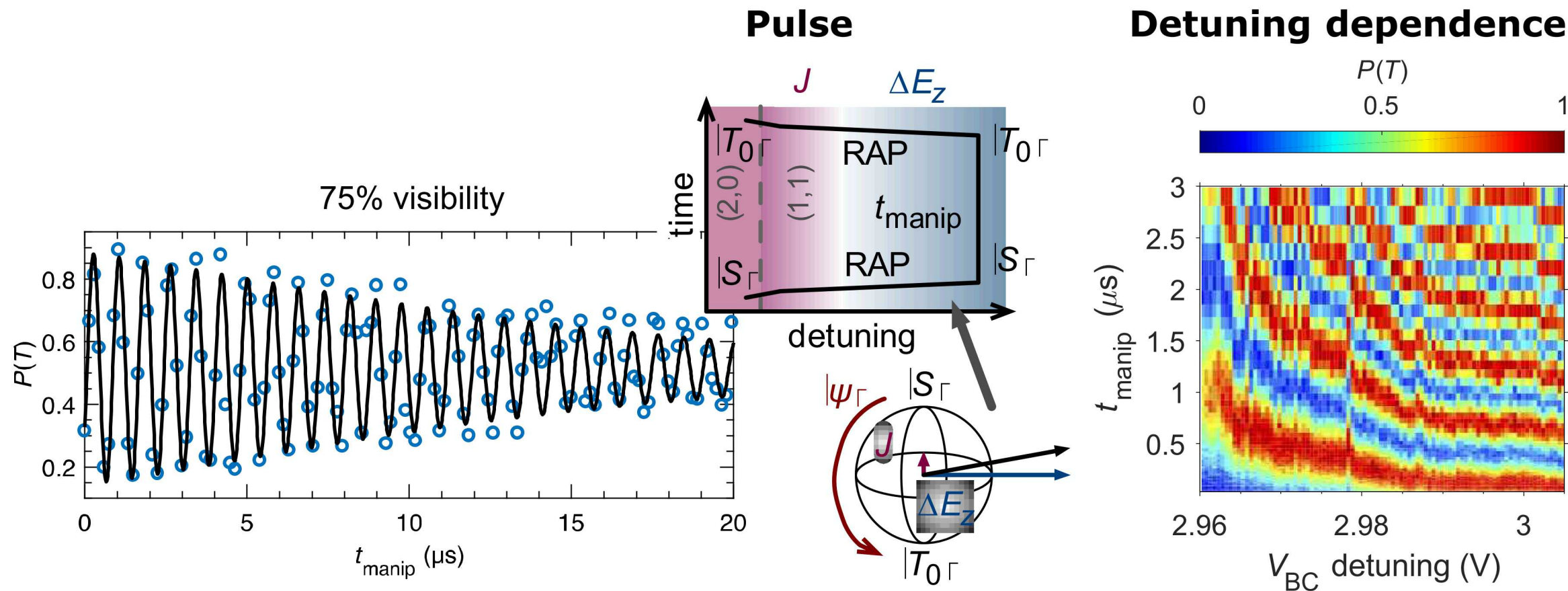
Diabatically with respect to spin.

$$|S\rangle \leftrightarrow |S\rangle$$

$$|T_0\rangle \leftrightarrow |T_0\rangle$$

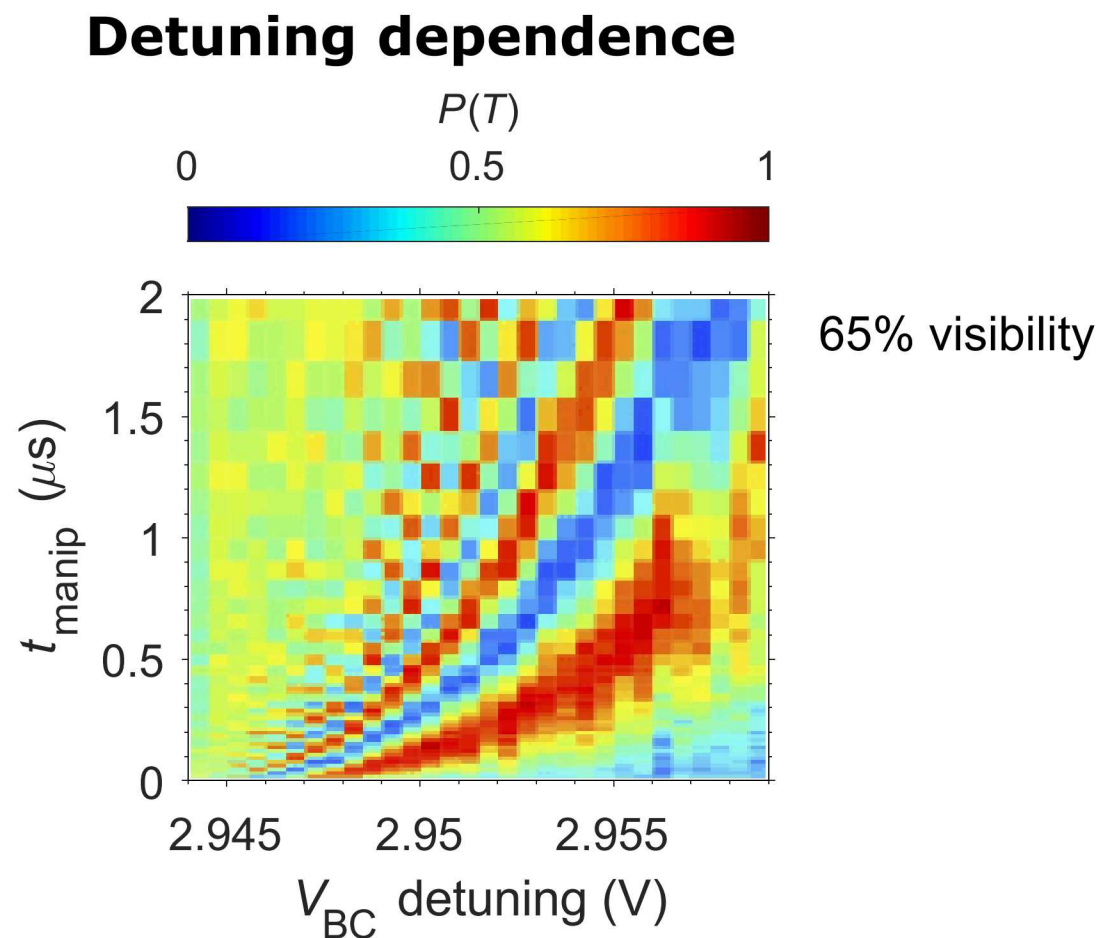
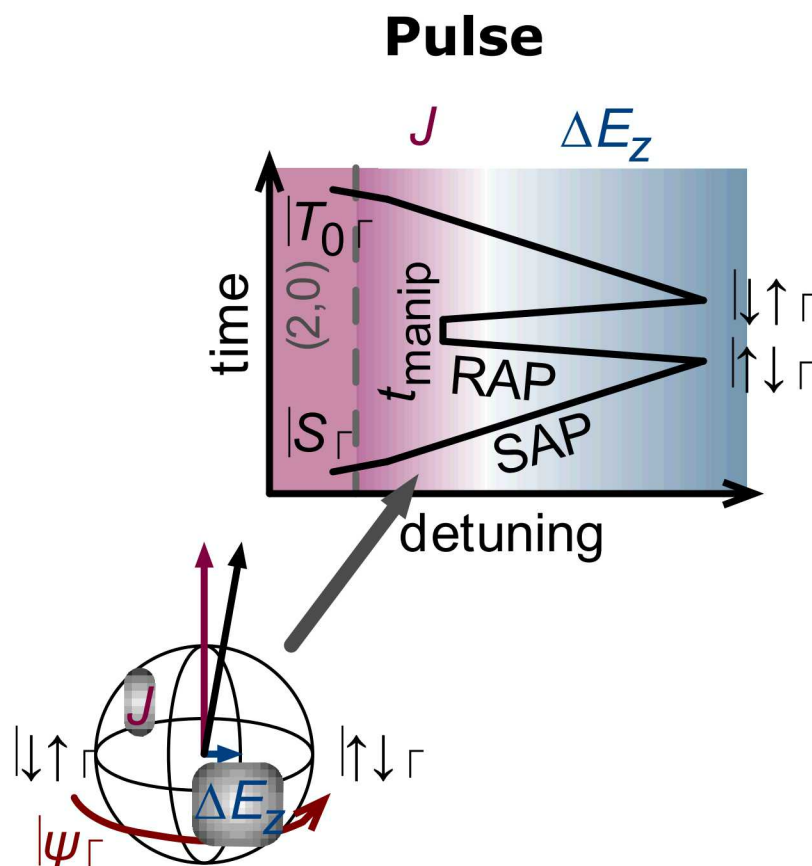


DC control – Zeeman rotations



[P. Harvey-Collard ... In 2017 IEEE IEDM (2017)]

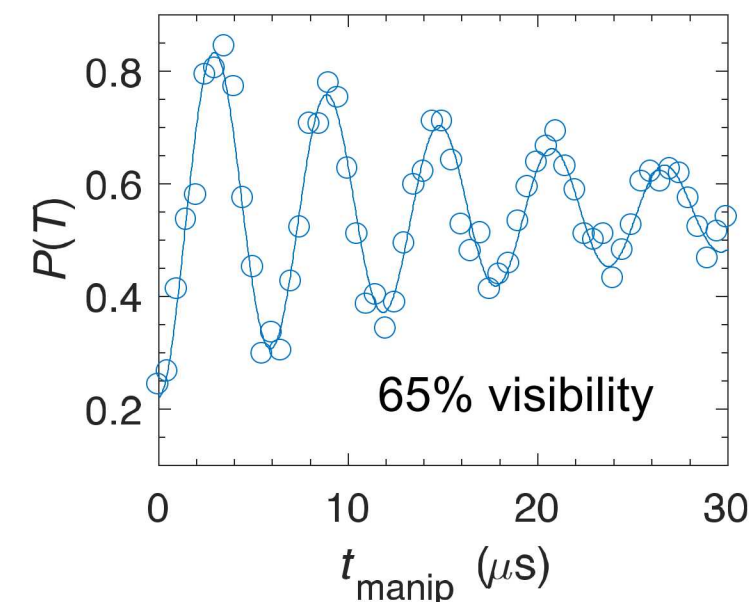
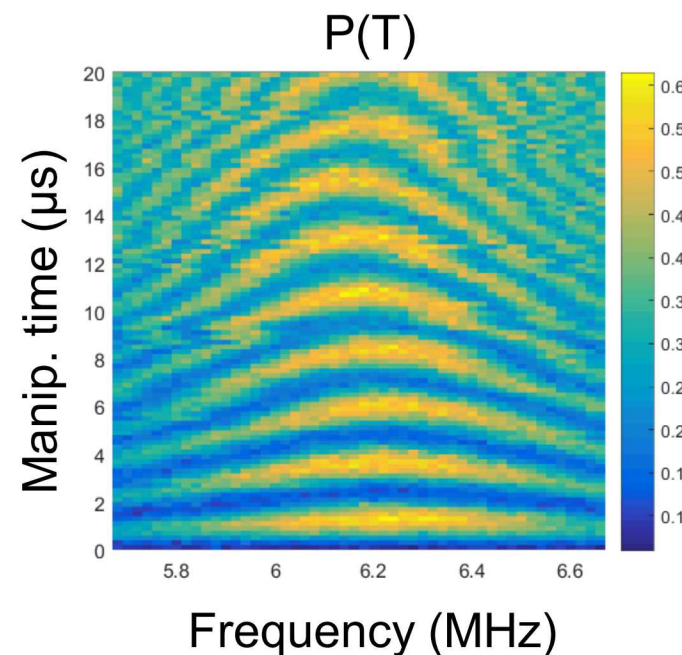
DC control — Exchange rotations



AC-controlled rotations [1, 2, 3]

Access any 1 qubit state

Figure 1: Schematic of the experimental setup. The top panel shows a time evolution plot with 'time' on the vertical axis and 'detuning' on the horizontal axis. A purple shaded region labeled 'SAP' (Spin Anomalous Phase) is bounded by two lines that converge at a point labeled 't_manip'. The region is labeled with 'J' (red) and 'ΔE_Z' (blue). The initial state is labeled '|S_Γ' and the final state is labeled '↑↓_Γ'. The bottom panel shows a Bloch sphere with a red arrow labeled 'J_ac' and a blue arrow labeled 'ΔE_Z'.

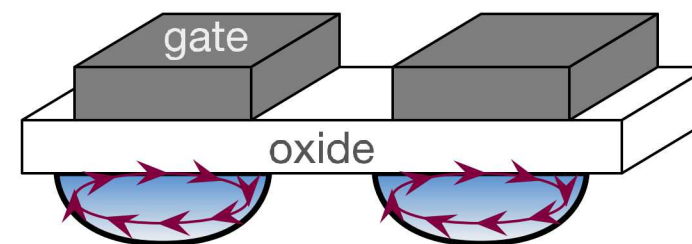


Data: C. Bureau-Oxton

- [1] Klauser, D. ... Phys. Rev. B 73, 205302 (2006).
 [2] Shulman, M. D. ... Nat Commun 5, 233101 (2014).
 [3] P. Harvey-Collard ... In 2017 IEEE IEDM (2017).

- Universal control using spin-orbit interaction
- Implications of the spin-orbit interaction for singlet-triplet qubits in silicon

Part 2: Spin-orbit interaction in silicon dots



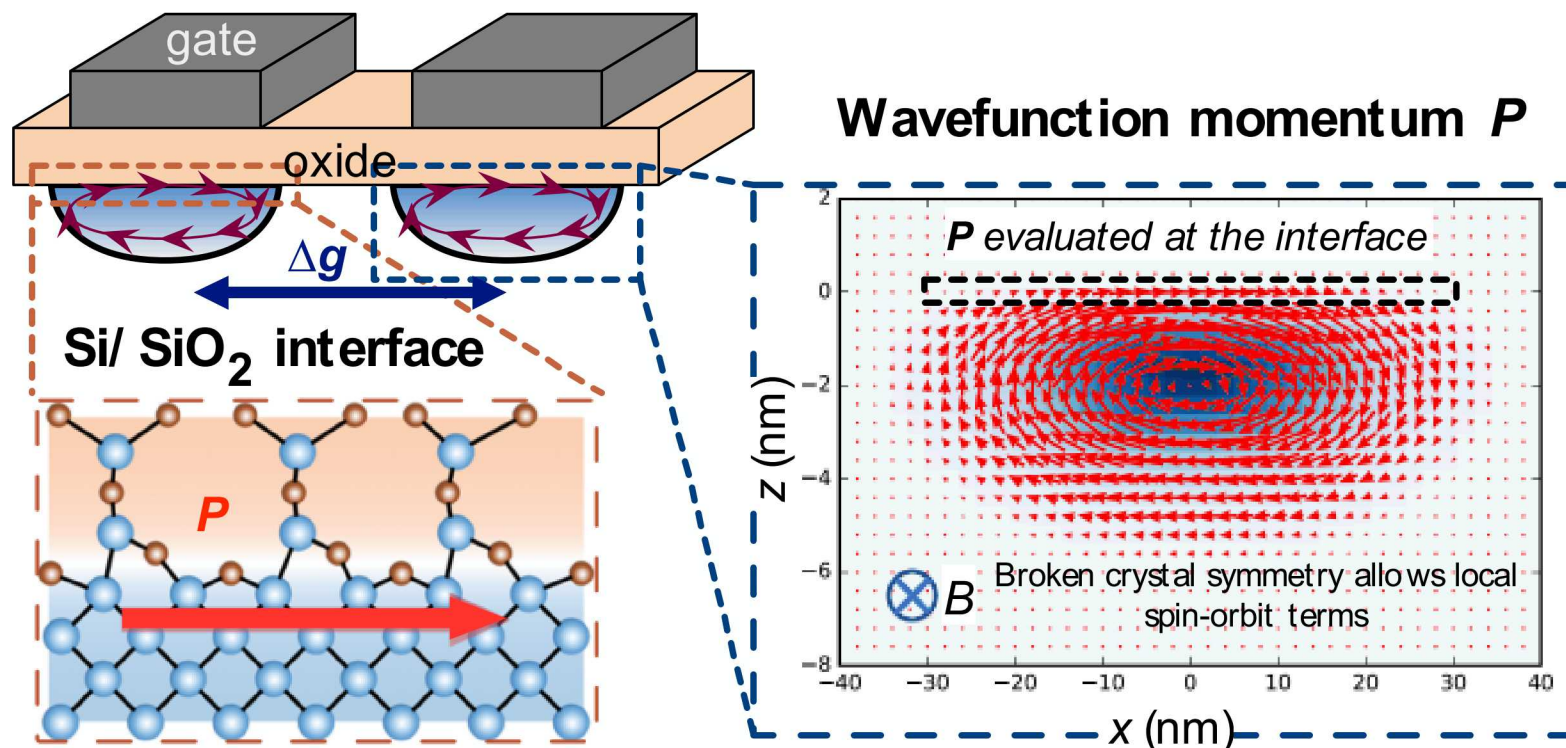
$$\begin{array}{l}
 |\uparrow\rangle \text{ —} \\
 E_Z = g_1 \mu_B B_{\text{ext}} \\
 |\downarrow\rangle \text{ —}
 \end{array}
 \quad
 \begin{array}{l}
 \otimes B_{\text{ext}} \\
 |\uparrow\rangle \text{ —} \\
 E_Z = g_2 \mu_B B_{\text{ext}} \\
 |\downarrow\rangle \text{ —}
 \end{array}$$

$$\Delta E_Z = (\Delta g) \mu_B B_{\text{ext}}$$

Spin-orbit interaction in silicon MOS

– Interface mechanism

Interface mechanism due to broken crystal symmetry*



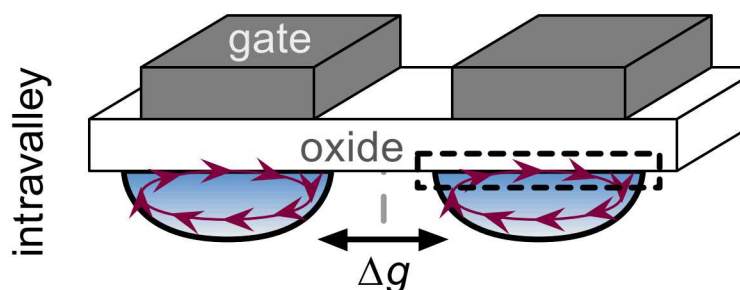
* Different from E-field-induced broken symmetry in 2DEGs
[L. E. Golub ... Phys. Rev. B 69, 115333 (2004)]

-> First-order effect vs P

Spin-orbit interaction in silicon MOS

– Intravalley effect

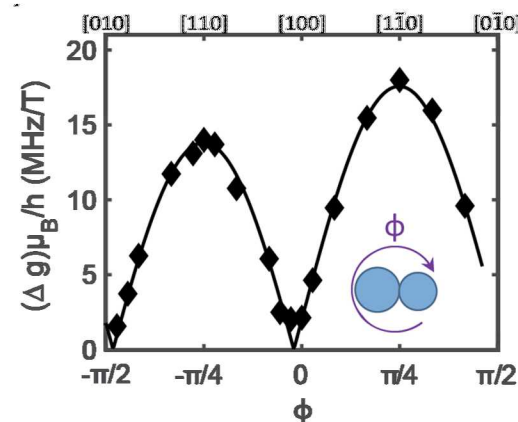
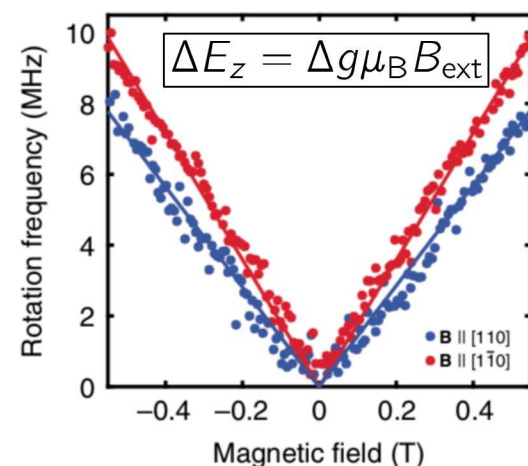
g is renormalized
in each dot



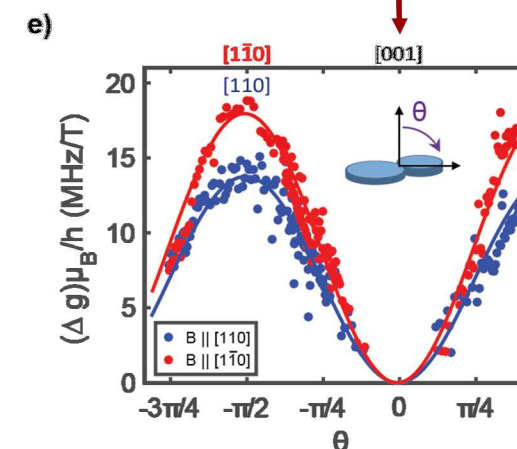
g -tensor

$$g = g_0 + g_{\text{SO}} = \begin{pmatrix} g_{\perp} - \frac{2}{\mu_B}\alpha & \frac{2}{\mu_B}\beta & 0 \\ \frac{2}{\mu_B}\beta & g_{\perp} - \frac{2}{\mu_B}\alpha & 0 \\ 0 & 0 & g_{\parallel} \end{pmatrix}$$

Anisotropy of Δg



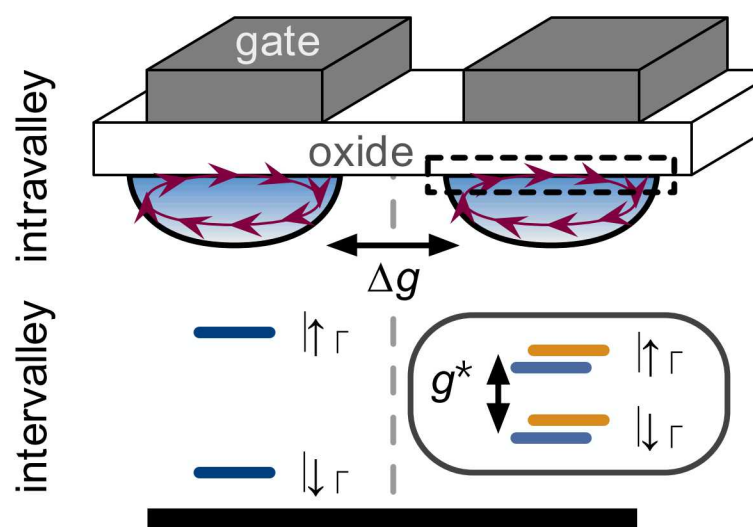
Δg vanishes for
out-of-plane B-field



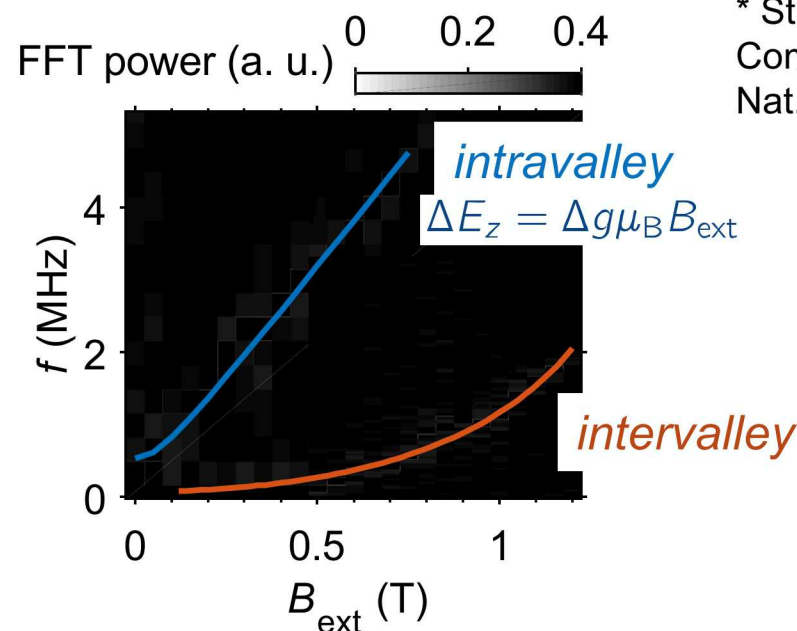
Spin-orbit interaction in silicon MOS

– Intervalley effect

Interaction with excited valley state



Non-linear effect

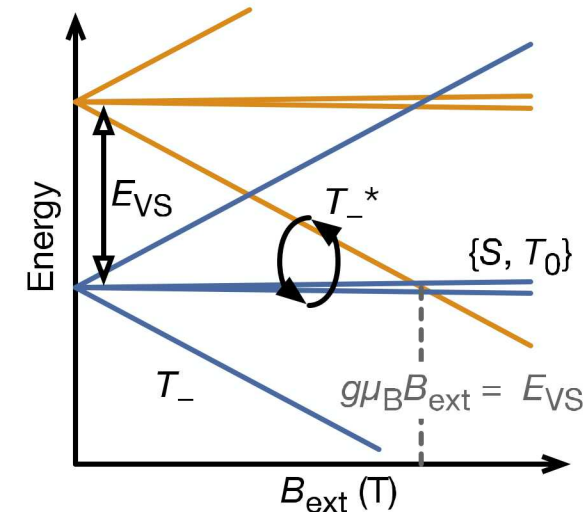


$$\Delta E_z = \frac{\bar{\beta}_C^2 B_{\text{ext}}^2}{2(E_{\text{VS}} - g\mu_B B_{\text{ext}})}$$

$$\bar{\beta}_C = 0.7 \mu\text{eV T}^{-1}$$

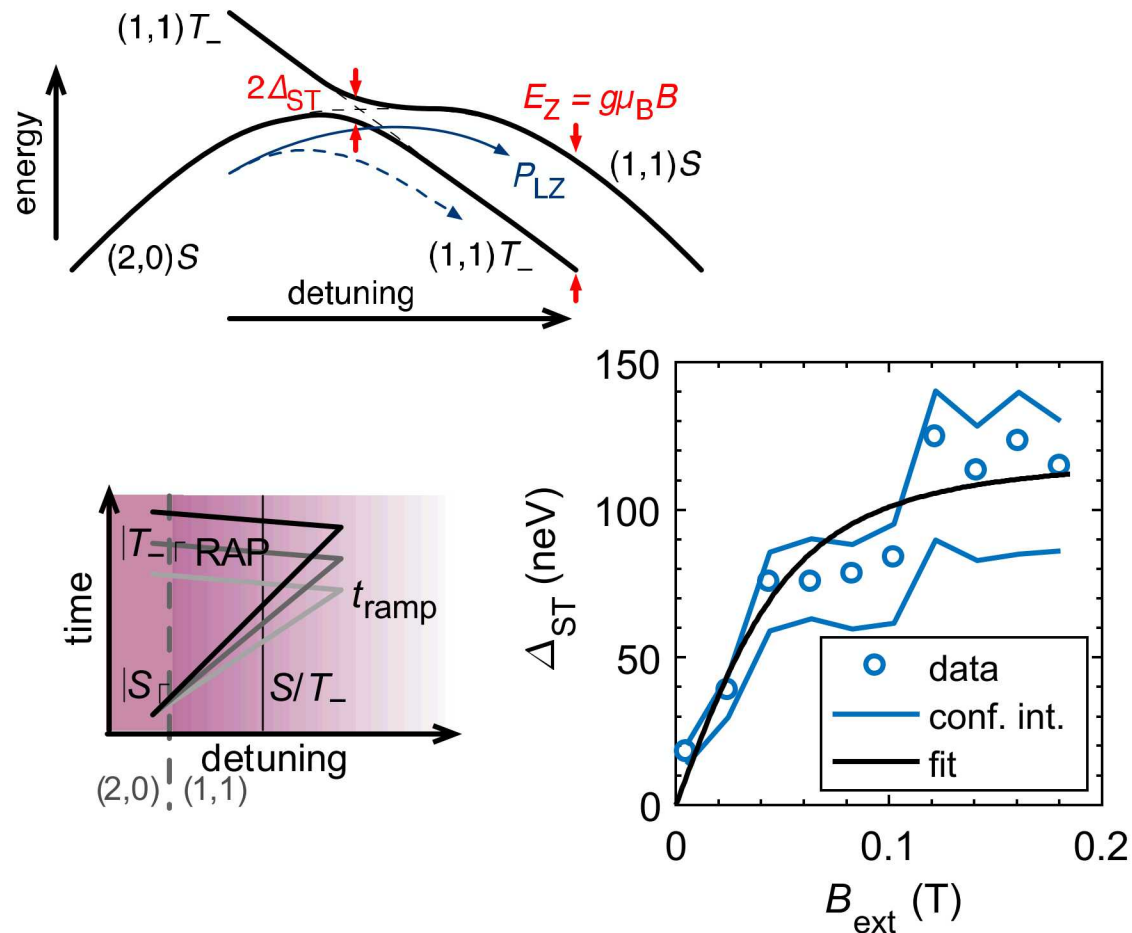
Second-order intervalley mechanism*

* Strength consistent with [C. H. Yang ... Nat. Commun. 4, 2069 (2013)] and [X. Hao ... Nat. Commun. 5, 3860 (2014)]



Spin-orbit interaction in silicon MOS

– Tunneling + spin-flip mechanism



Spin-orbit length [1]:

Result: $\Omega_{SO} \sin \Phi = 113 \pm 22$ neV $B_{ext} // [100]$

$\lambda_{SO} = t_C d / \Omega_{SO} \sqrt{2} \sim 1 \mu\text{m}$

Bulk Si: 20 μm !

Bulk GaAs: 1.1 μm !

Affects dynamic nuclear polarization (DNP)

Hyperfine interaction [1]:

Find bound: $\sigma_{HF} < 5$ neV

Expect: $\sigma_{HF} \approx 0.6$ neV

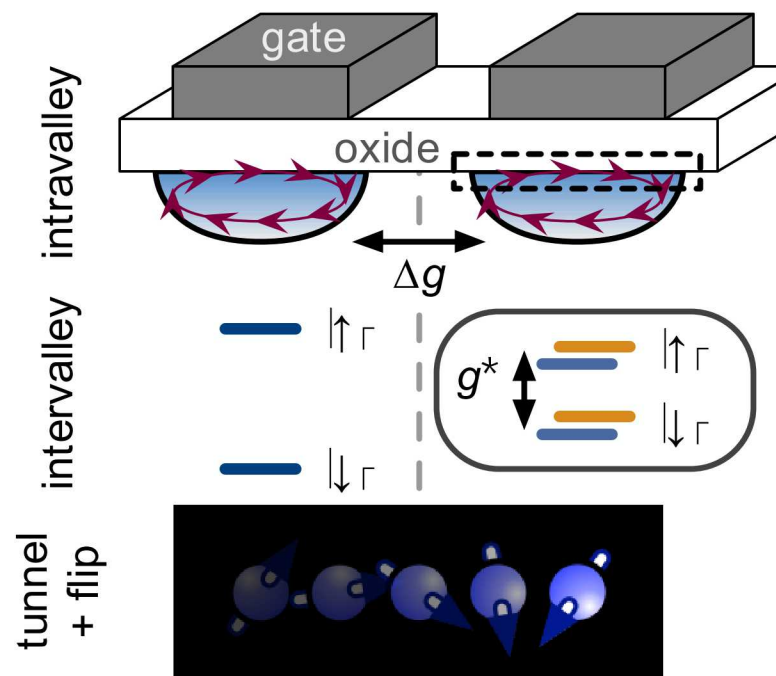
BUSTED

[Method:] J. M. Nichol ... Nat Commun 6, 7682 (2015).

[1] P. Harvey-Collard ... arXiv:1808.07378 (2018).

How should I mount my device?

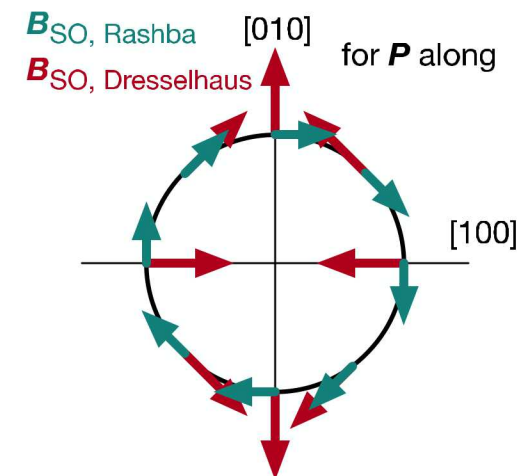
3 spin-orbit effects



$S-T_0$ effects caused by $\Delta B_{\parallel}$

- Static (1,1)
- intravalley, intervalley
- momentum P :
 B_{ext} via cyclotron orbit
- component of $B_{\text{SO}} \parallel B_{\text{ext}}$

Δg vanishes for
out-of-plane B-field



$S-T_{\perp}$ effect caused by $\Delta B_{\perp}$

- Dynamic (2,0) $\leftrightarrow$ (1,1)
- tunneling + spin-flip
- momentum P :
double quantum dot axis
- component of $B_{\text{SO}} \perp B_{\text{ext}}$

Δ_{ST} persists for
out-of-plane B-field



Conclusion & Summary

Acknowledgements & Team

Experiments

- Martin Rudolph
- Ryan M. Jock
- Chloé Bureau-Oxton
- Matthew J. Curry
- Andrew M. Mounce
- Michael P. Lilly
- Malcolm S. Carroll

Theory

- N. Tobias Jacobson
- Andrew D. Baczewski
- Vanita Srinivassa
- Wayne M. Witzel
- John King Gamble

Fabrication

- Daniel R. Ward
- John M. Anderson
- Ronald P. Manginell
- Joel R. Wendt
- Jason Dominguez
- Tammy Pluym

Université de Sherbrooke

- Michel Pioro-Ladrière

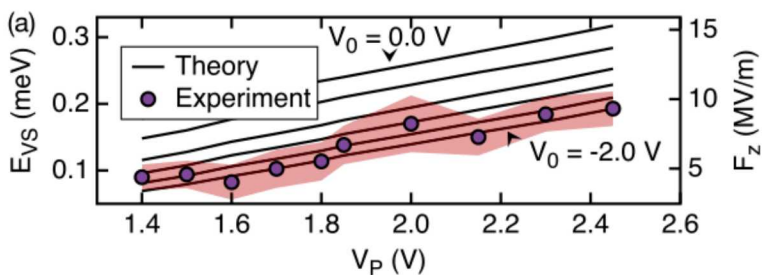


And the rest
of the QIST
team!



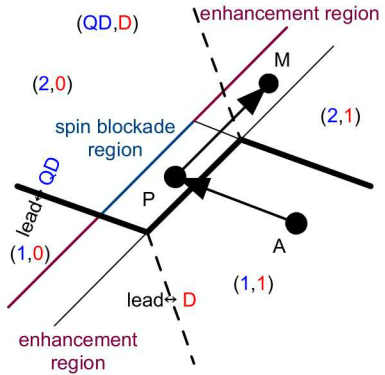
This work was performed, in part, at the Center for Integrated Nanotechnologies, an Office of Science User Facility operated for the U.S. Department of Energy (DOE) Office of Science. Sandia National Laboratories is a multimission laboratory managed and operated by National Technology and Engineering Solutions of Sandia, LLC, a wholly owned subsidiary of Honeywell International, Inc., for the DOE's National Nuclear Security Administration under contract DE-NA0003525.

Valley splitting



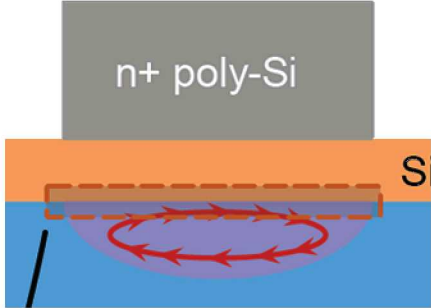
J. K. Gamble, P. Harvey-Collard et al., App. Phys. Lett. 109, 253101 (2016).

Enhanced latching readout



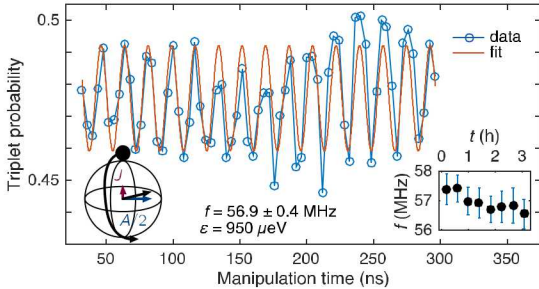
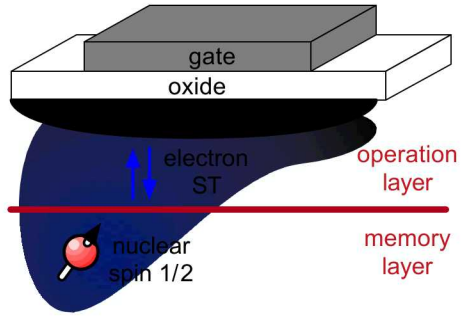
P. Harvey-Collard et al., Phys. Rev. X 8, 021046 (2018).

Spin-orbit



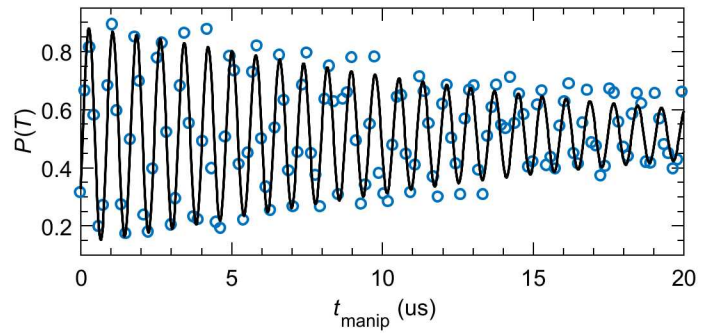
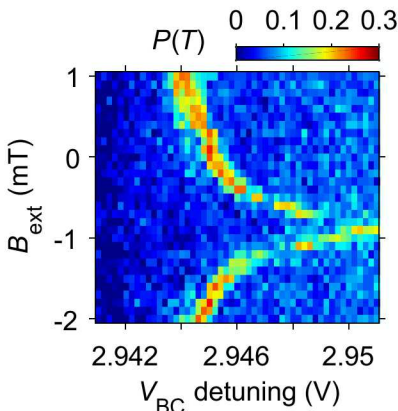
R. M. Jock, N. T. Jacobson,
P. Harvey-Collard et al., Nat. Commun. 9, 1768 (2018).

MAJIQ (donor-dot)



- 1. P. Harvey-Collard et al., Nat. Commun. 8, 1029 (2017).
- 2. M. Rudolph, P. Harvey-Collard et al. In 2016 IEEE IEDM (2016) pp. 34.1.1–34.1.4.

More spin-orbit



- 1. P. Harvey-Collard et al. In 2017 IEEE IEDM (2017).
- 2. P. Harvey-Collard et al., arXiv:1808.07378 (2018).