

ANALYZING THE POTENTIAL FOR EROSION IN A SUPERCRITICAL CO₂ TURBINE NOZZLE WITH LARGE EDDY SIMULATION

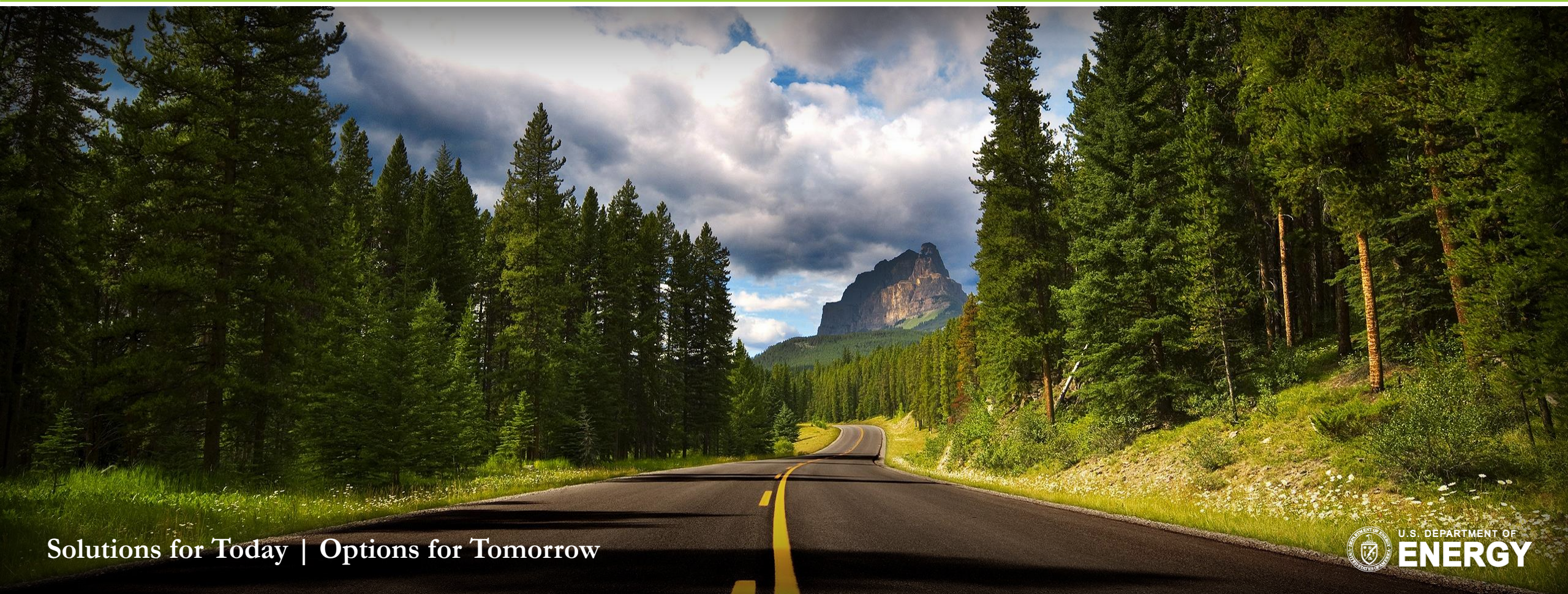
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Solutions for Today | Options for Tomorrow



Why sCO₂ power cycles?

Higher Efficiency

High working fluid temperatures

Recompression near liquid densities

High heat recuperation

Lower Capital Cost

Compact turbo machinery

Simple configurations

Lower Environmental Impact

Zero emissions

Dry cooling

Water production

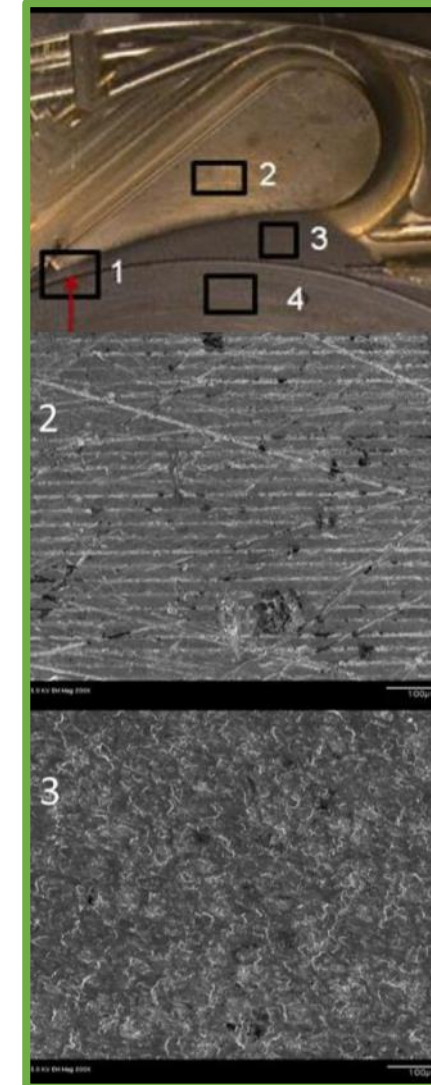
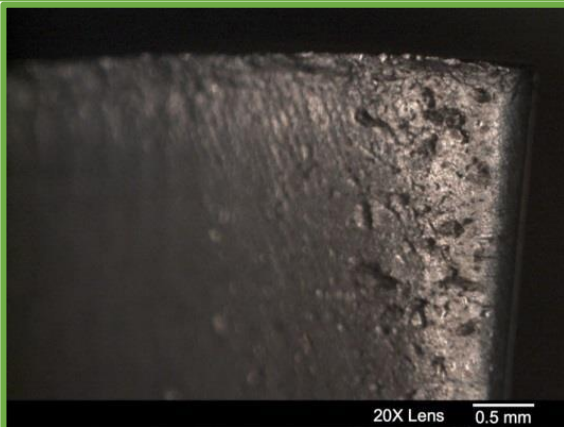
Erosion of sCO₂ Power Cycle components

Problem: Severe erosion of turbine blades and vanes has been observed in the sCO₂ cycle test loops in Sandia National Lab and Bettis Lab.



Sandia Findings

- Particle impingement on flow surfaces
- Particles in the loop:
 - Stainless steel
 - SiO₂
 - Al₂O₃



Root cause of damage?

- Solid particle impacts?
- Shear/press fluctuations?
- Phase change?



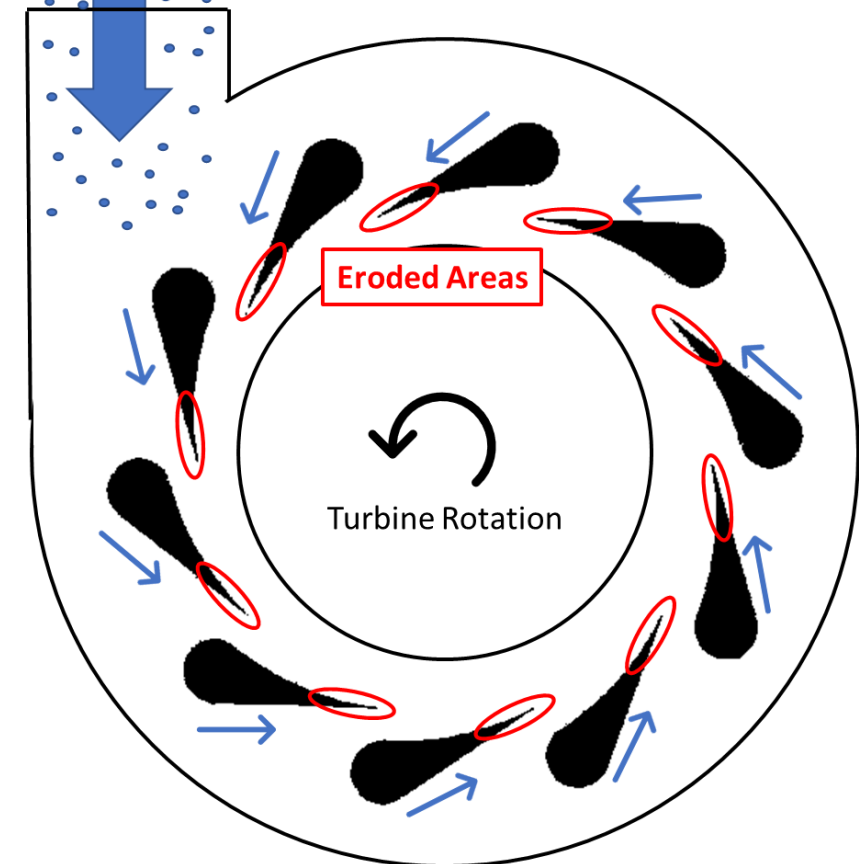
Image: Sandia turbine assembly. Inset shows damage to nozzle after operation with $s\text{CO}_2$ [1].



This Talk: Potential of entrained oxide particles to erode radial turbine nozzle.

- LES + Particle Tracking
- SNL nozzle geometry
- Semi-empirical erosion model
- Scaling to operating conditions

$s\text{CO}_2$ + Solids



Methodology: LES + Lagrangian particle tracking

Fluid phase (sCO₂)

- Spatially filtered, incompressible N-S for Large Eddy Simulation

$$\frac{\partial \bar{u}_j}{\partial x_j} = 0$$
$$\frac{\partial \bar{u}_j}{\partial t} + \frac{\partial \bar{u}_i \bar{u}_j}{\partial x_j} = -\frac{1}{\rho} \frac{\partial \bar{P}}{\partial x_i} + \nu \frac{\partial^2 \bar{u}_i}{\partial x_j^2} - \boxed{\frac{\partial q_{ij}}{\partial x_j}}$$

- Dynamic Smagorinsky closure for sub-grid scale stresses

Key Assumptions

- Incompressible, low Mach number (for now)
- Constant properties (density, viscosity)
- One-way coupling of particles to flow (dilute assumption)

Solid phase (porous oxide particles)

- Large Eddy Simulation (LES) for turbulence
- Particle tracking/collision prediction

$$\frac{d}{dt}(\mathbf{x}_p) = \mathbf{u}_p \quad (6)$$

$$\frac{d}{dt}(\mathbf{u}_p) = -\frac{1}{\rho_p} \nabla \bar{P} + \frac{f(Re_p)}{\tau_p} (\bar{\mathbf{u}} - \mathbf{u}_p) \quad (7)$$

Numerical Solution

- Second order, co-located finite volume discretization for arbitrary, unstructured grids¹
- Fractional step, pressure projection time advancement
- Fully parallelized, capable of variable density (future work)
- Validation for particle laden flows in complex geometries^{2,3}

Operating conditions

Assumed Operating Conditions (125kWe SNL test facility)

- Nozzle chord length $L \approx 2.5\text{cm}$, height $h \approx 0.4\text{cm}$
- Flow rates in the neighborhood of $\dot{V} = 600\text{GPM}$
 - Velocity at nozzle tip, $u_{bulk} = 45\text{m/s}$
- Turbine inlet conditions : $T = 550^\circ\text{C}$, $P = 29.11\text{MPa}$
 - $\rho = 177 \frac{\text{kg}}{\text{m}^3}$, $\mu = 39 \times 10^{-6}\text{Pa} \cdot \text{s}$
- Reynolds number: $Re_c = \frac{\rho \cdot u_{bulk} \cdot L}{\mu} \approx 5 \times 10^6$

Simulation conditions

- Order of magnitude reduction in Re_c to accommodate LES, without wall model.
- Solid particles: Spalled porous oxide scales
 - $\rho_p = 2500 \text{ kg/m}^3$, $d_p = 30\mu\text{m}$
- Stokes Numbers (measure of particle inertia)

$$St_p = \frac{\tau_p}{\tau_k} = 89, \quad St_b = \frac{u_{bulk} \cdot \tau_p}{R_c} = 0.08.$$

Table: Simulation parameters

	sCO ₂		Particles
ρ_f	$177 \text{ kg} \cdot \text{m}^{-3}$	ρ_p	$2500 \text{ kg} \cdot \text{m}^{-3}$
μ	$39 \times 10^{-6} \text{ Pa} \cdot \text{s}$	d_p	$30\mu\text{m}$
θ	20°	St_+	89
U_{in}	$8.37 \text{ m} \cdot \text{s}^{-1}$	St_b	0.08
u_{bulk}	$4.67 \text{ m} \cdot \text{s}^{-1}$		
Re	529,865	$\langle\langle \varepsilon \rangle\rangle$	10^{-3}
dt_f	$2.5 \times 10^{-8} \text{ s}$	dt_p	$2.5 \times 10^{-8} \text{ s}$

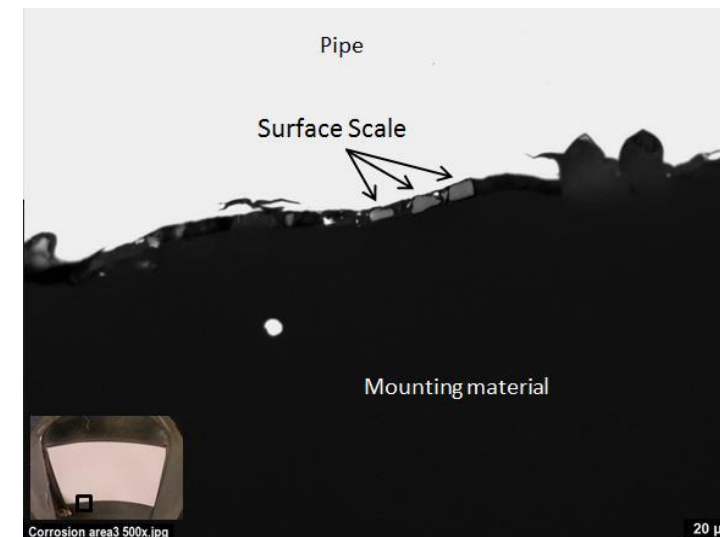
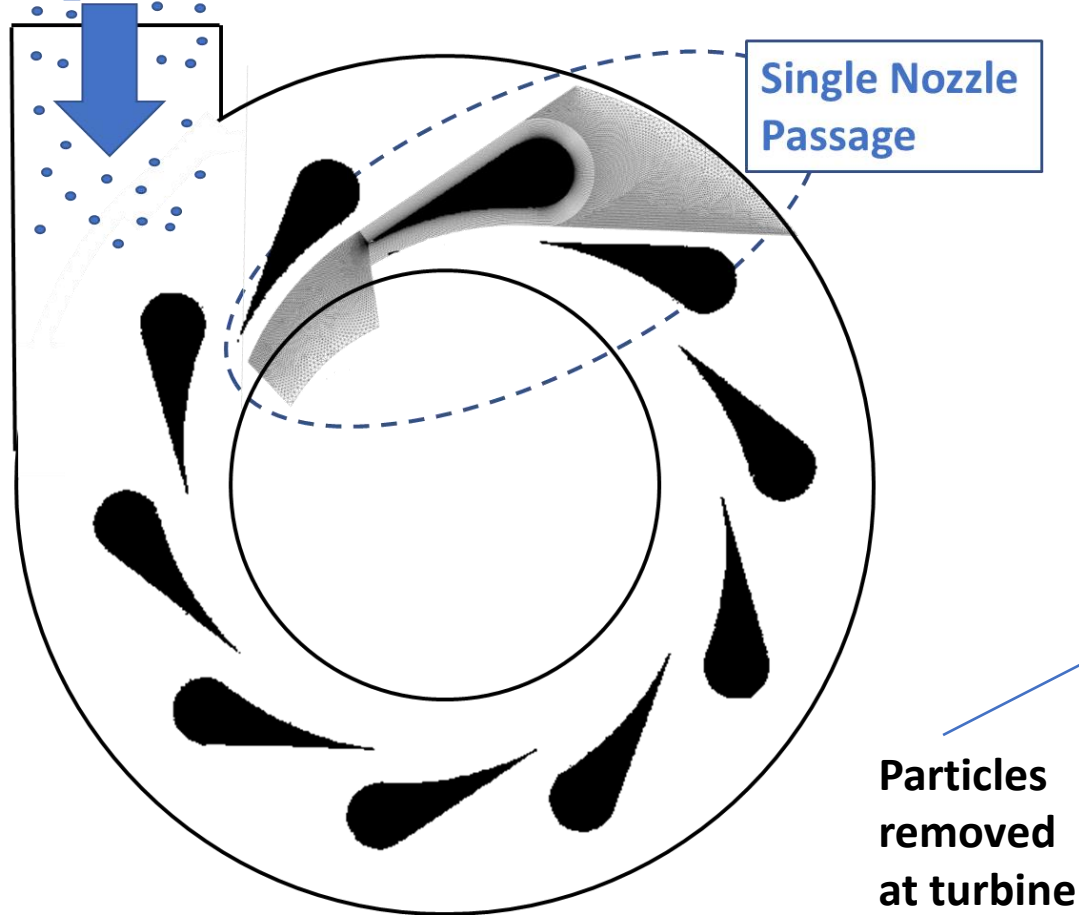


FIG: Evidence of intergranular corrosion and surface scale on sCO₂ flow loop piping [1]

Simulation Setup: Domain, Mesh, and B.C's

sCO₂ + Solids



A-Priori grid spacing:
 $\Delta_n^+ \approx 1, \Delta_s^+ \approx 20, \Delta_z^+ \approx 50,$

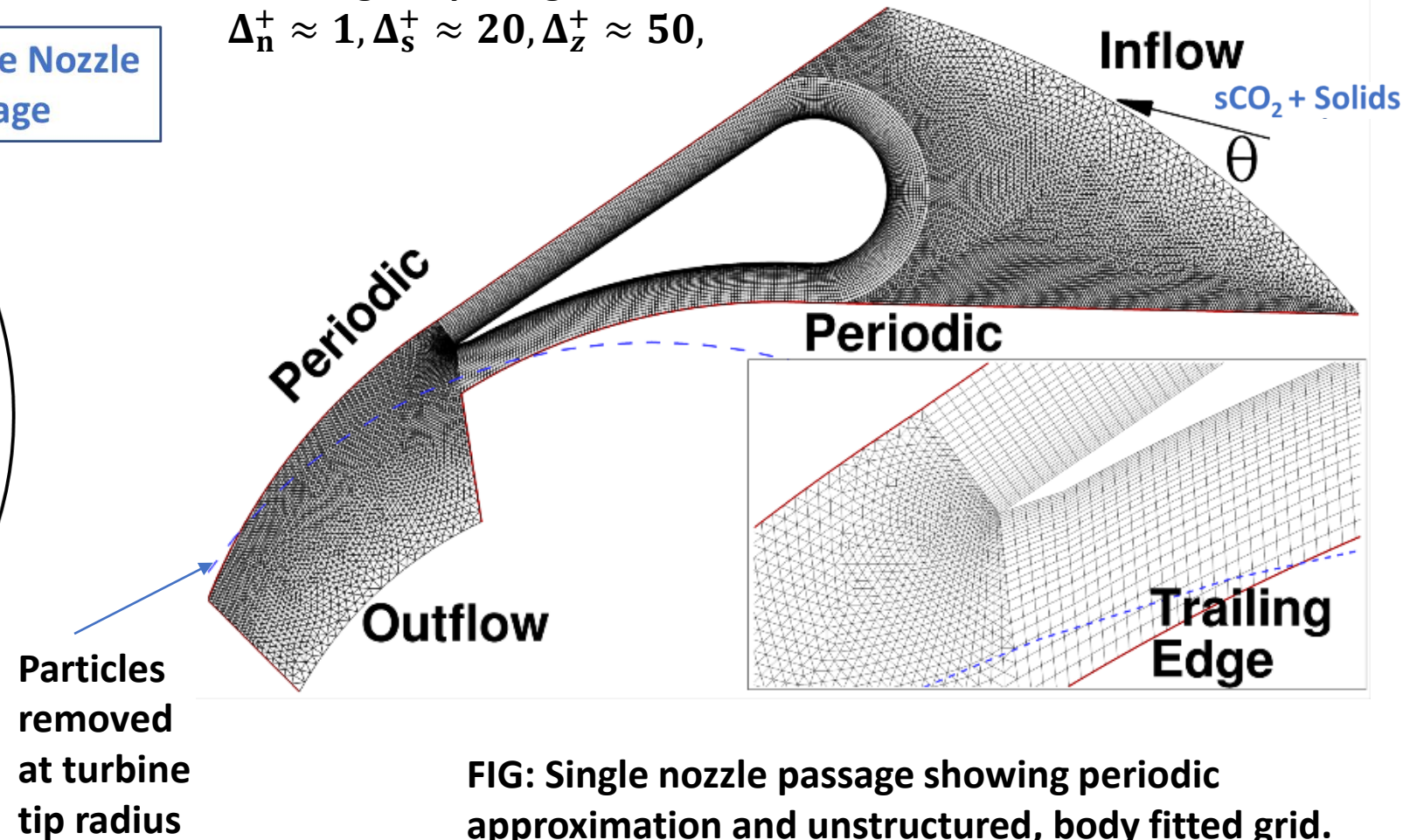


FIG: Single nozzle passage showing periodic approximation and unstructured, body fitted grid.

Turbulent inflow conditions

- Synthetic eddy method (Jarrin et al, 2006) used to specify homogeneous velocity fluctuations on top of mean flow:

$$\overline{u}_i = \langle u_i \rangle + a_{ij} u'_i$$

$$u'_j(x, t) = \sum_{i=1}^N \epsilon_{ij} g(x - x_i)$$

$$a_{ij} = \begin{pmatrix} \sqrt{R_{11}} & 0 & 0 \\ R_{21}/a_{11} & \sqrt{R_{22} - a_{21}^2} & 0 \\ R_{31}/a_{11} & \left(R_{32} - \frac{a_{21}a_{31}}{a_{22}} \right) & \sqrt{R_{33} - a_{31}^2 - a_{32}^2} \end{pmatrix}$$

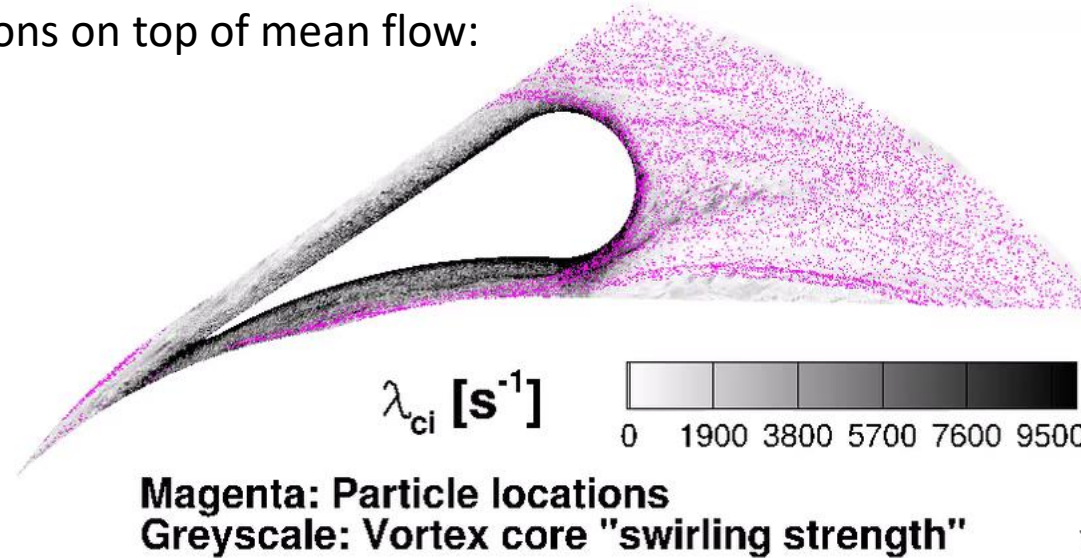
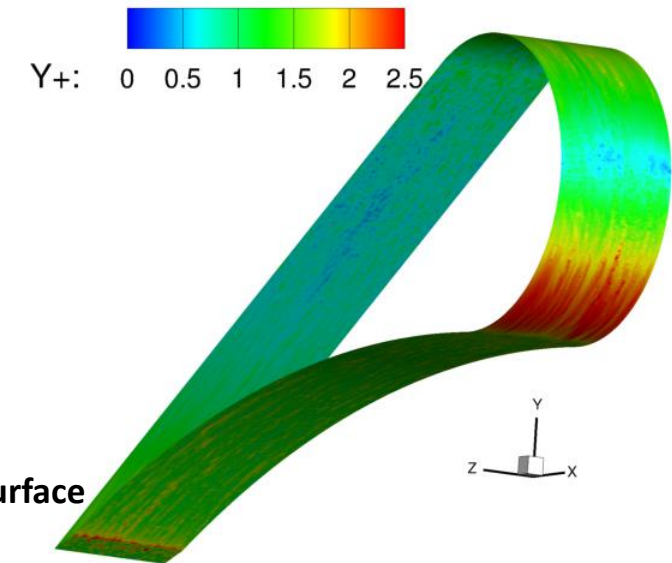
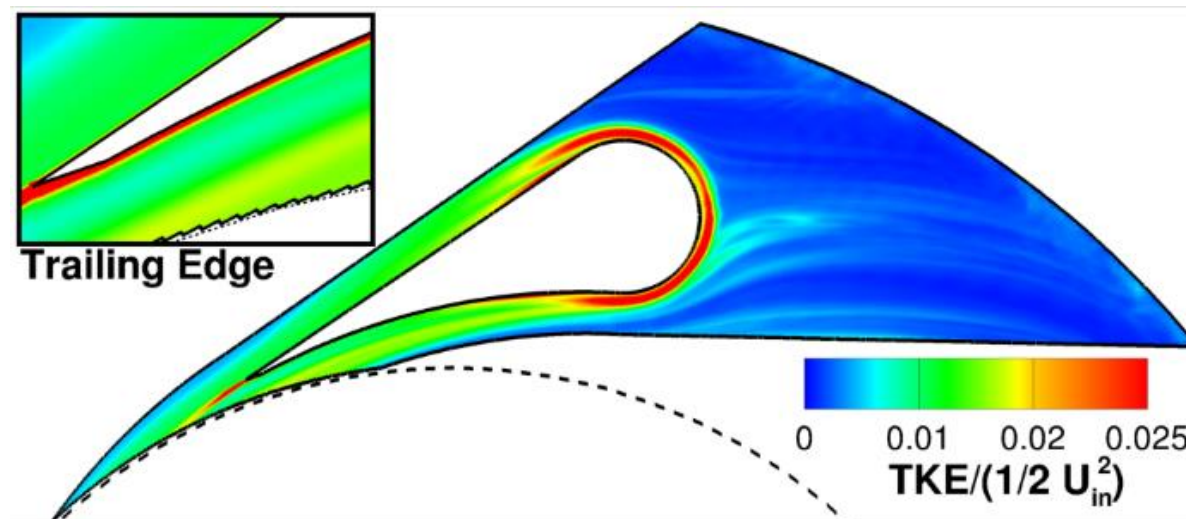
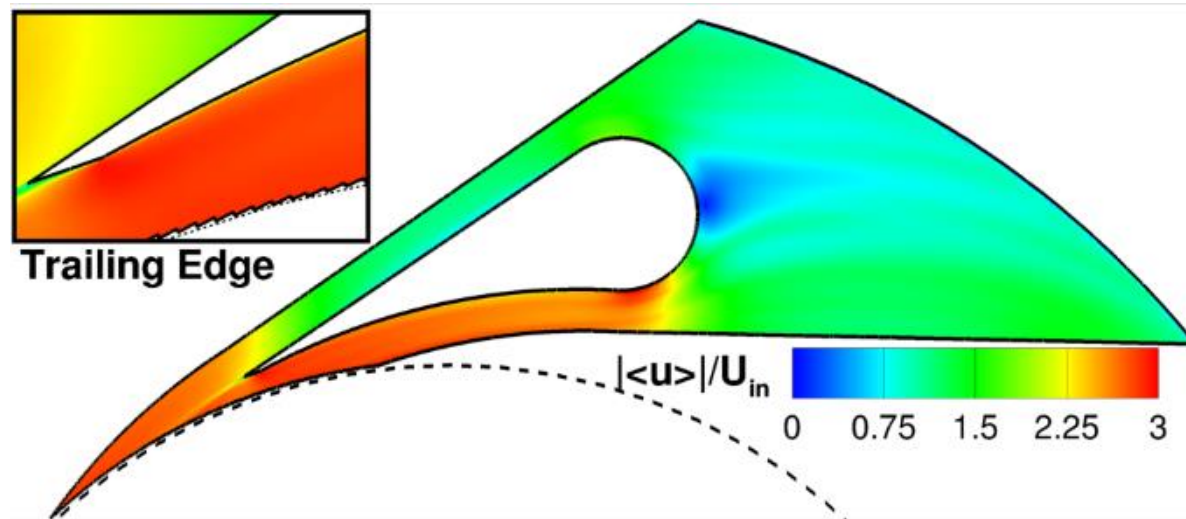


FIG: Simulated particle trajectories (magenta) and vortex swirling strength (greyscale) on center plane



Results: Time averaged flow



Acceleration through the passage:

- Trailing edge velocity $\sim 3\times$ greater than leading edge
- Increasing TKE along pressure-side of passage



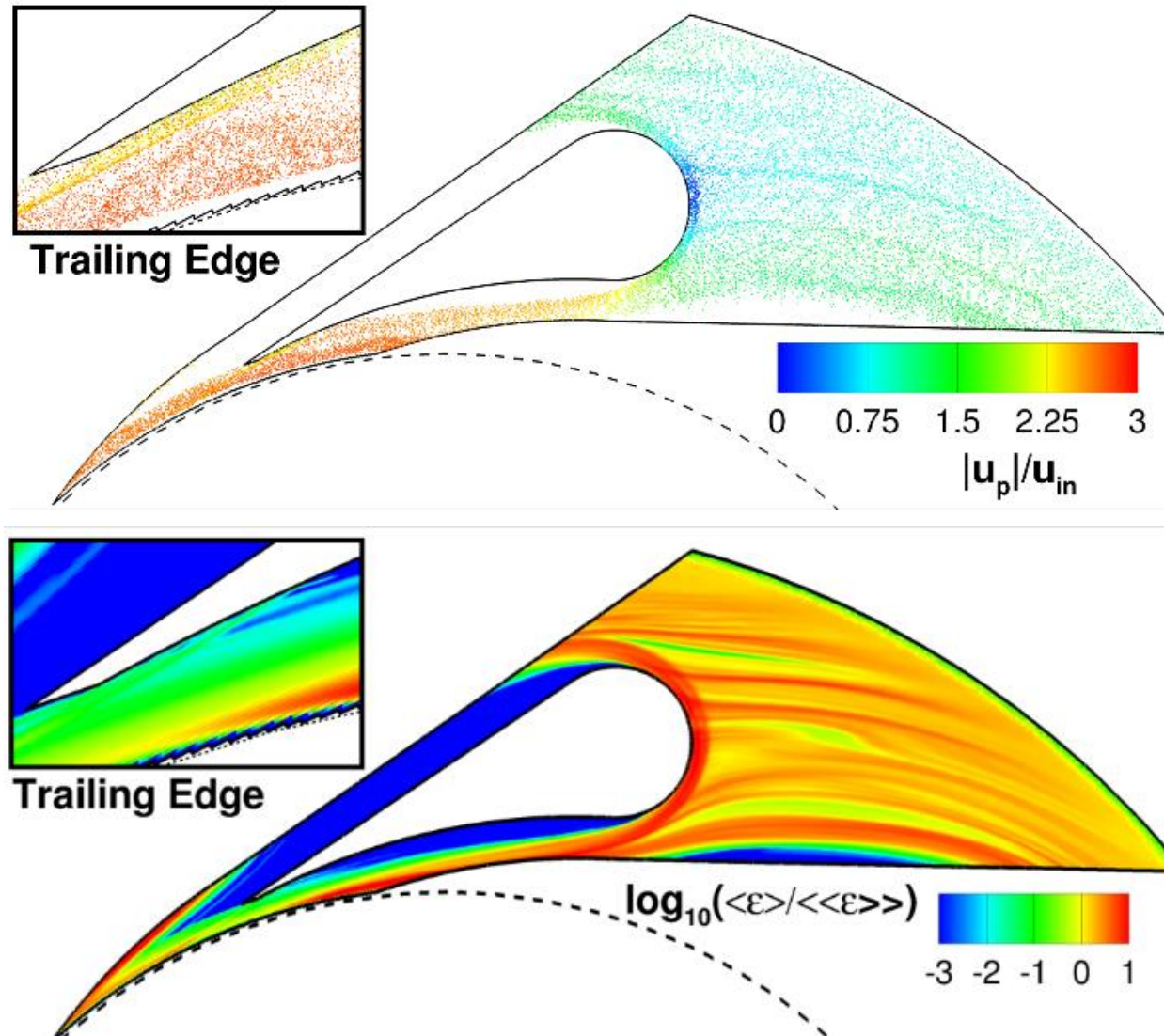
Particle dispersion

Increased slip velocities

High velocity impacts

FIG: Time averaged flow velocity magnitude (top), and Turbulent Kinetic Energy (bottom).

Results: Solid particle trajectories



High Stokes number particles *NOT* evenly distributed

- Stagnation at leading edge, separation around mid span
- Acceleration through the passage
- Particles 10x more concentrated in passage than upstream
- Small amount of high vel. particles approaching trailing edge

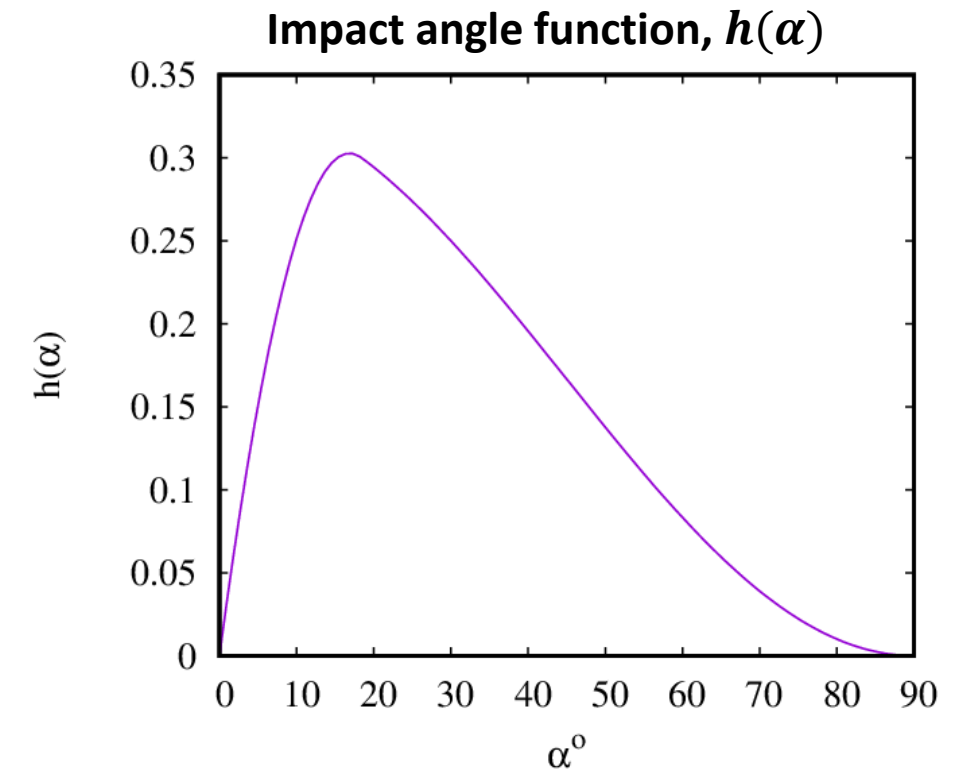
FIG: Snapshot of particle distribution, colored by velocity (top), and time averaged particle concentration field (bottom).

Erosion model

- Oxide scale collisions with nozzle
 - Specular reflections
 - Constant coefficient of restitution, $e = 0.8$
- Erosion rate calculated from collisions using a Finnie model [Finnie,1960]
- Erosion calculated over several flow through times, normalized by local maximum.

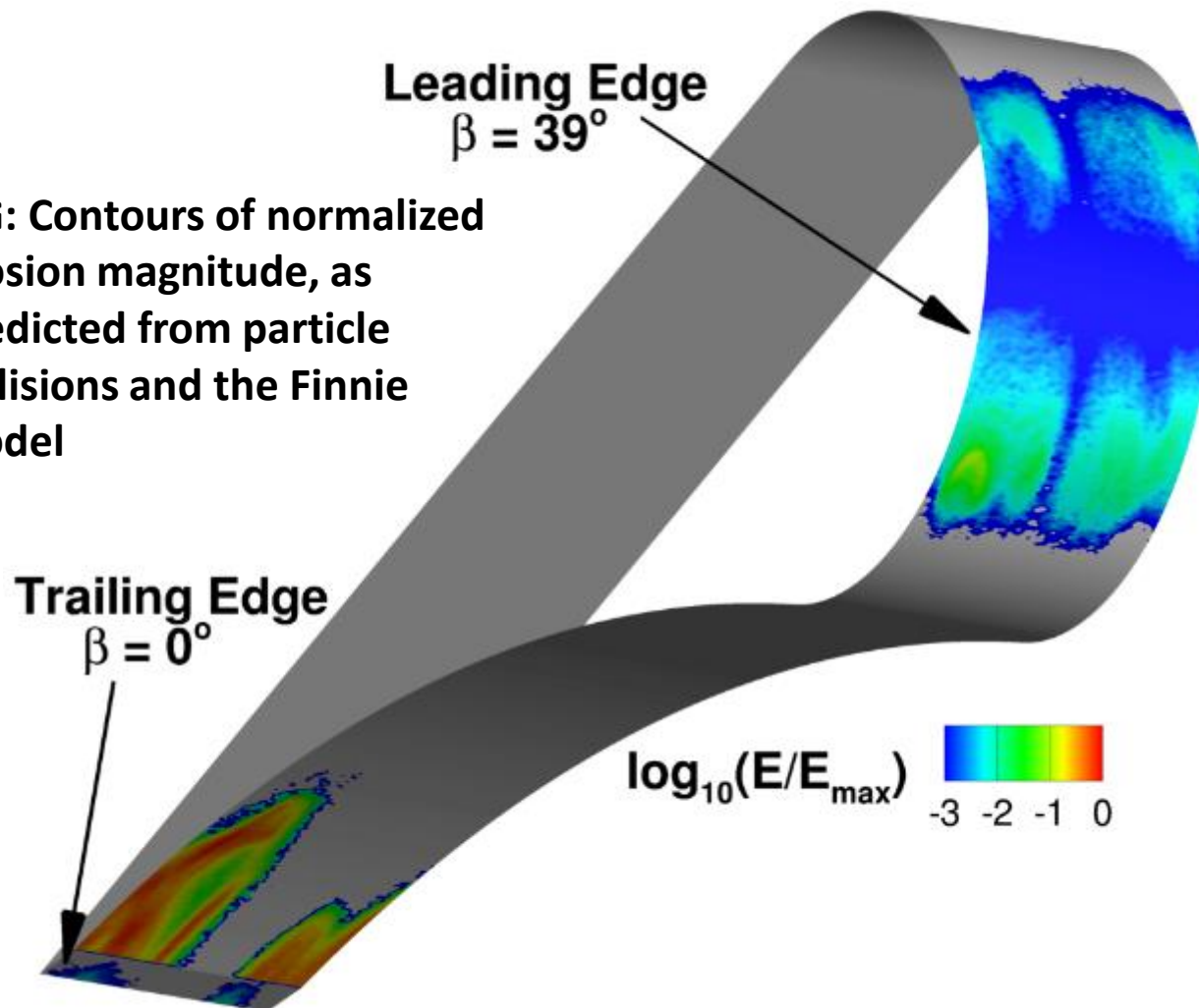
$$E \propto m_p |\mathbf{u}_p|^2 g(\alpha).$$

$$h(\alpha) = \begin{cases} \frac{1}{3} \cos^2(\alpha) & \text{for } \tan(\alpha) \leq \frac{1}{3} \\ \sin(2\alpha) - 3 \sin^2(\alpha) & \text{for } \tan(\alpha) > \frac{1}{3} \end{cases}$$



Results: Erosion Magnitudes

FIG: Contours of normalized erosion magnitude, as predicted from particle collisions and the Finnie model

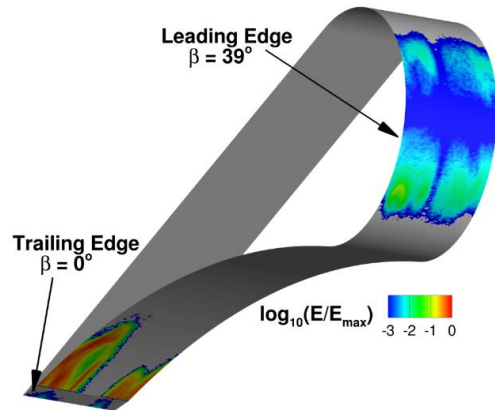


- No particle collisions over much of surface
- Trailing edge erosion roughly 2 orders of magnitude larger than leading edge
- Consistent with location of damage observed at SNL



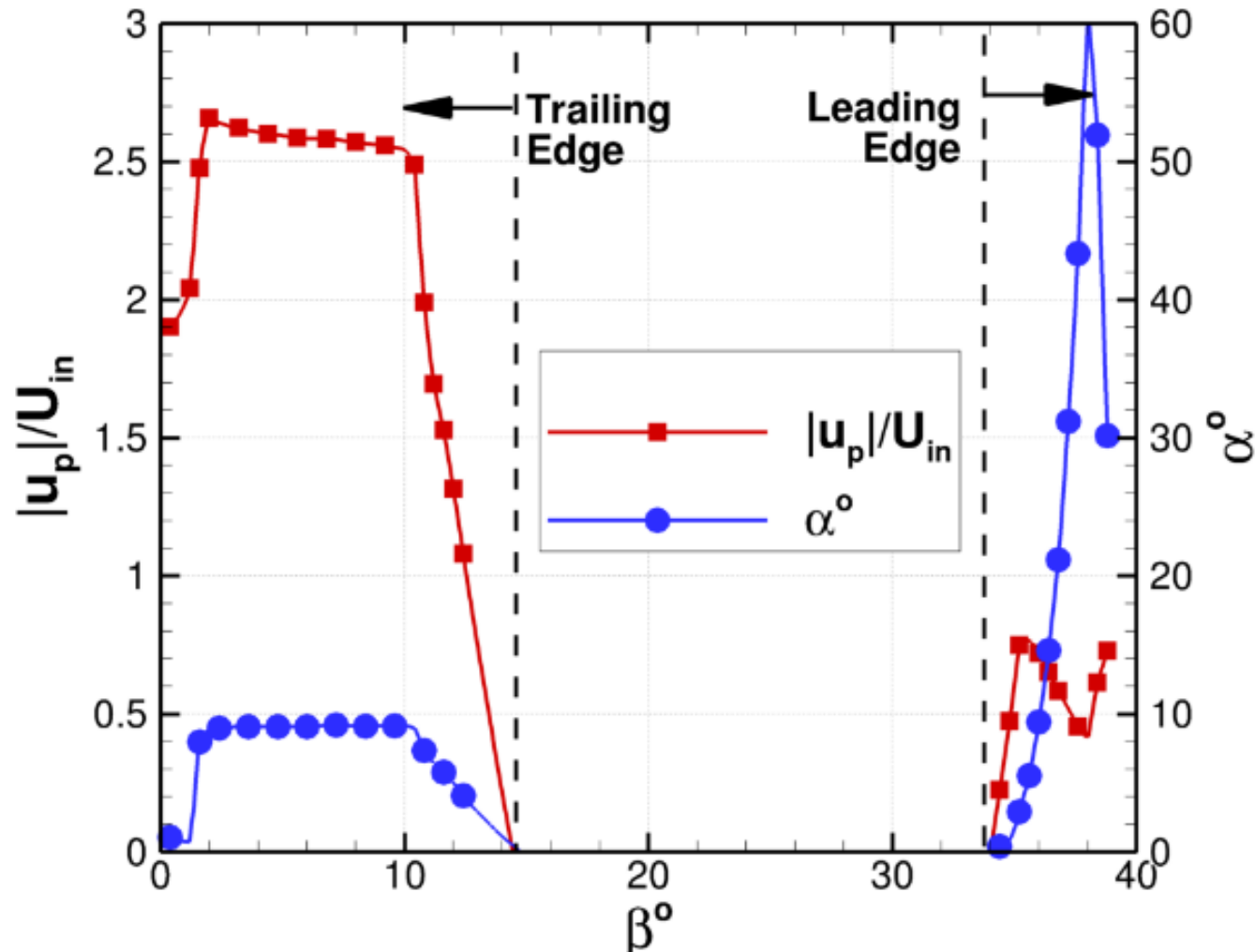
FIG: Eroded nozzle at SNL

Results: Impact statistics



Trailing Edge:

- High velocities
- Consistently low impact angles

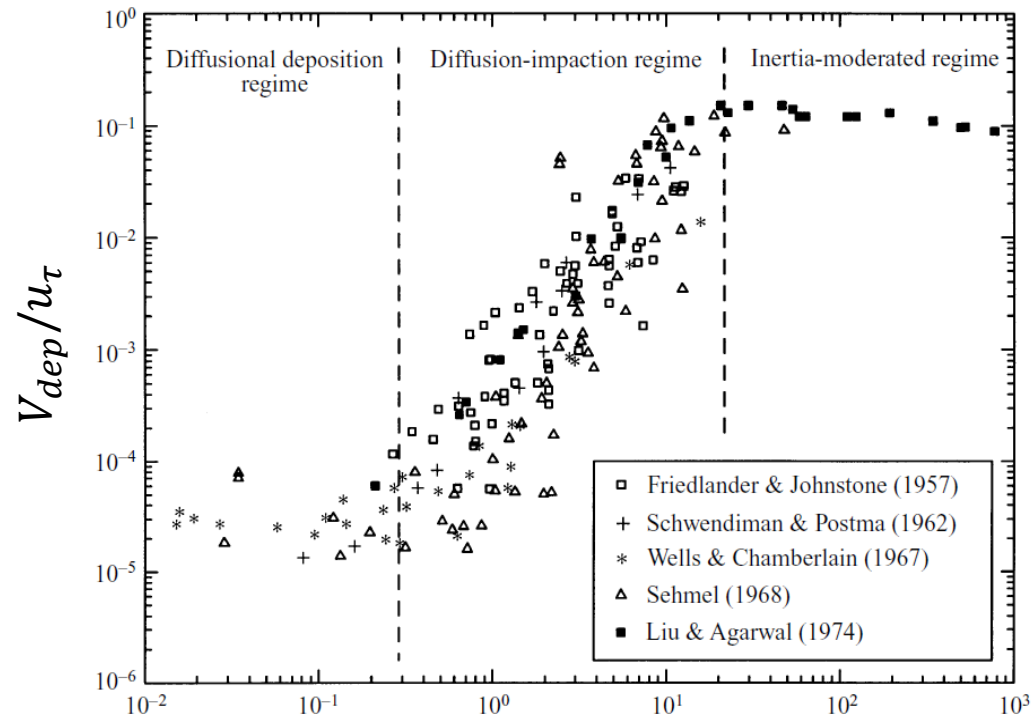


Leading Edge:

- Low impact velocities near stagnation point
- Wider distribution of oblique impact angles

Implications for operational Re.

Particle Deposition velocity depends primarily on Stokes number¹

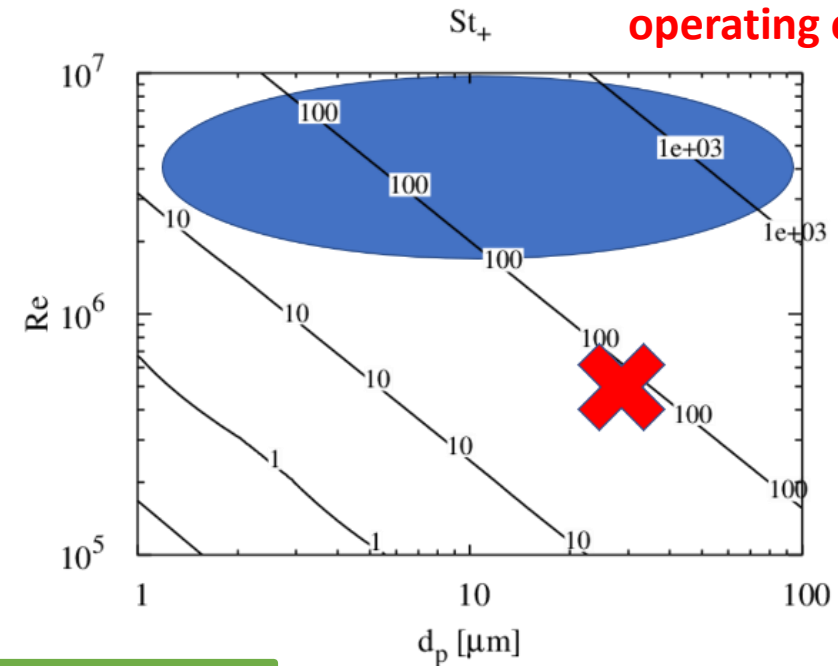


$$St_p = \frac{\tau_p}{\tau_k}$$

May need to filter very small particles from flow loop to avoid damaging impacts!

In our system, St_p scales with d_p and Re_c

$St_p = \mathcal{O}(10^3)$ for operating conditions!



Conclusions

- Potential for **solid particle erosion** in sCO₂ turbine nozzle has been investigated
 - Spalled oxide scales entrained in closed flow loop
 - High velocity impacts cause damage to certain components
- **Approach:** Large eddy simulation with Lagrangian particle tracking
 - Periodic domain consisting of 1 nozzle passage
 - Synthetic turbulence specified upstream
- Uniformly distributed particles become **highly concentrated** as they traverse the passage
- **High velocity, oblique impacts** near trailing edge lead to relatively large erosion potential in region that was damaged at SNL
- **Ongoing effort:** More simulations to understand dependence on St_p

Acknowledgments



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Extra Slides

Animation at lower Reynolds Number, $Re=18,000$

