

Infrasound and gravity waves over the Andes observed by a pressure sensor on board a stratospheric balloon

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Key Points:

- Pressure signals from mountain-generated gravity and acoustic waves are observed on a high altitude balloon
- Gravity wave observations are compared to satellite observations of AIRS instrument
- The low decrease rate of microbaroms amplitude from the coast favor global presence of this noise source in the stratosphere waveguide

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Abstract

The study of infrasound (acoustic) and gravity waves sources and propagation in the atmosphere of a planet gives us precious insight on atmosphere dynamics, climate and even internal structure. The implementation of modern pressure sensors with high rate sampling on stratospheric balloons is improving their study. We analyzed the data from the National Aeronautics and Space Administration (NASA) Ultra Long Duration Balloon (ULDB) mission (16th May to 30th June 2016). Here, we focus on the balloon's transit of the Andes Mountains. We detected gravity waves that are associated to troposphere convective activity and mountain waves. An increase of the horizontal wavelengths from 50 to 70 km with increasing distance to the mountains is favoring the presence of mountain waves. We also report on the detection of infrasounds generated by the mountains in the 0.01-0.1 Hz range with a pressure amplitude increase by a factor 2 relative background signal. Besides, we characterized the decrease of microbaroms power when the balloon was flying away from the ocean coast. These observations suggest, in a way similar to microseisms for seismometers, that microbaroms are the main background noise sources recorded in the stratosphere even far from the ocean sources. Finally, we observed a broad band signal above the Andes, between 0.45 Hz and 2 Hz, probably associated with a thunderstorm. The diversity of geophysical phenomena captured in less than a day of observation stresses the interest of high rate pressure sensors on board long duration balloon missions.

Plain language summary

The variations of pressure recorded on board a high altitude balloon are analyzed to observe various phenomena. When the balloon is passing over the Andes at 33 km altitudes, different waves emitted below it are recorded. Waves produced by storms, atmospheric flows over the mountains, oceanic waves and thunderstorms are identified in the records. Mountains waves observed by the balloon's instruments are compared to images of the same altitude range obtained by satellites. The poorly documented bass music created by winds in the mountains are observed in certain frequency range, and the sounds created by interfering oceanic waves are slowly decreasing from the coast. The variety of geophysical phenomena observed in less than one day demonstrate the power of such missions and instruments. Their interest is extending beyond Earth, because planets with dense atmosphere like Venus can be investigated with such kind of missions.

1 Introduction

Propagating waves in the atmosphere induces pressure fluctuations which can be recorded by sensitive microbarometers. Waves under 20 Hz are divided in two domains: infrasound and gravity waves. Gravity waves are buoyancy-driven and often originate from the deviation of air flow over the large planetary topographic structures (Nappo, 2002; Plougonven et al., 2008), from deep convection and from wind jets (Fritts & Alexander, 2003). Infrasound can be generated by natural or artificial events, as quakes (Hines, 1960; Lognonné et al., 2016; Martire, Brissaud, et al., 2018; Young et al., 2018), non-linear interactions between ocean waves called microbaroms (Landès et al., 2014; Bowman & Lees, 2018), volcanoes (Assink et al., 2014; Matoza & Fee, 2018), aircraft or thunderstorms (Lamb et al., 2018). The signals contain the information necessary to better understand and characterize the topography and atmospheric features of a planet and to study its internal structure. In that way, the study of infrasound is directly linked to the planetary and geophysical sciences. On Earth, balloon missions represent a great opportunity to study the atmospheric waves over inaccessible regions, such as oceans, mountains, and polar ice caps (Haase et al., 2018). Although acoustic and gravity waves captured on ground-based sensors has been well studied over several decades (Le Pichon et al., 2010; de Groot-Hedlin et al., 2014; de Groot-Hedlin et al., 2017), investigations using high altitude bal-

64 loons are more sparse. Infrasound produced by ground explosions was measured by Bowman
 65 and Albert (2018) and the issues associated with the stratospheric flight of acoustic ar-
 66 rays is described by Bowman and Lees (2015) and Lees and Bowman (2017). These stud-
 67 ies showed that the pressure sensors can be used to observe infrasound and gravity waves.
 68 One important reason for this is that, despite lower amplitude signals, passively-drifting
 69 balloons are less subject to the wind noise and the turbulence which are the main sources
 70 of noise on the ground. However, the low amount data acquired by these platforms did
 71 not allowed to quantify yet all the physical phenomena underlying the production and
 72 propagation acoustic and gravity waves.

73 Tests of balloon infrasound sensing technology on Earth is the first step to study
 74 other planets on which it is difficult or even impossible to deploy ground stations. Stevenson
 75 et al. (2015) studied the feasibility of an infrasound balloon mission on Venus. The Venus
 76 ground atmospheric conditions ($9\text{e}+6$ Pa and 700 K) do not allow any long duration de-
 77 ployment on the surface, even modern technology cannot withstand these temperatures
 78 and pressures for more than a couple of hours. In contrast, balloons with infrasound sen-
 79 sors would overpass these limitations, and may be able to probe the interior structure
 80 of a planet by acoustic pulses from seismic activity on its surface (Garcia et al., 2005)
 81 and its atmosphere structure and dynamics (Preston et al., 1986).

82 Acoustic and gravity waves captured by balloons on Earth and other planets come
 83 from a variety of sources. Here, we focus on one point of interest of ULDB (Ultra Long
 84 Duration Balloon) mission for which the data has not been analyzed in detail yet: the
 85 pass over the Andes. We start by describing the ULDB mission and the associated data
 86 which are used in our analysis (section 2). Then, we will present the various geophys-
 87 ical processes detected when the balloon flew over the Andes and quantify some of the
 88 underlying physical phenomena (section 3). Finally we conclude on the interest of such
 89 measurements to characterize wave sources and propagation effects (section 4).

90 2 Data

91 2.1 Instrumentation

92 Our study was based on the analysis of the NASA’s ULDB mission data, described
 93 in Bowman and Lees (2018). The balloon was launched from Wanaka, New Zealand on
 94 the 16 May 2016. The batteries of the sensors lasted for 19.5 days, until the 6 June 2016.
 95 The balloon’s total flight duration was 46 days, with an average altitude of 33 km. The
 96 balloon had a GPS on board, which provided the longitude, latitude, altitude and speed
 97 of the balloon every 15 minutes. The microbarometer payload was composed of one record-
 98 ing system *Omnirecs Datacube*, sampled at 200 Hz, and 3 differential microbarometers
 99 *InfraBSU*. The microbarometer configuration was the following:

- 100 • One with positive pressure polarity.
- 101 • One with negative pressure polarity.
- 102 • One mechanically disabled.

103 This setup permits us to distinguish between true pressure signals (polarity reversed be-
 104 tween the first two channels, not present on the third channel) with electromagnetic and
 105 other interference (common polarity on all three channels). It also allows us to deter-
 106 mine the non-pressure noise threshold. The corner frequency of the microphones is not
 107 precisely quantified but since the sensors are at high altitude, it is likely to be very low,
 108 around several thousand seconds, which allows us to perform our analysis in the low fre-
 109 quencies domain (Marcillo et al., 2012).

110 The ULDB data has already been used in Bowman and Lees (2018) to character-
 111 ize the energy transfer from the microbaroms signal to the thermosphere which would

112 increase its temperature by several Kelvins per day. This flight data was also used to de-
 113 scribe the acoustic signature associated with a thunderstorm (Lamb et al., 2018). We
 114 focus on a part of the data set which has not been analyzed in detail in order to illus-
 115 trate the phenomena observed as a balloon crosses a large mountain range. In this con-
 116 figuration, the winds and the balloon are flowing perpendicular to the Andes mountains
 117 ridge in a quasi two dimensional problem. It represents a perfect situation for the pro-
 118 duction of topographically-driven gravity and acoustic waves. In this manuscript, there-
 119 fore, we focus on the Andes crossing and downwind phenomena.

120 2.2 Noise characterization

121 In order to be able to distinguish between real signals and noise, it is necessary to
 122 define and characterize all the noise sources that the balloon may be subject to. The main
 123 sources of noise are shown in Figure 1. First, microbarometers are affected by wind noise,
 124 one of the main noise sources for the stations at ground level. Electronic noise is intrin-
 125 sic to the system and the electromagnetic perturbations come from thunderstorms or au-
 126 roras. The movement of the balloon and other payloads produces vibrations and the as-
 127 sociated noise. In the same way, the balloon is subject to altitude variations; thus the
 128 sensors would not sense the same level of background atmospheric pressure during the
 129 whole flight. And eventually, the last noise source is created by acoustic and gravity waves.
 130 Therefore, these variations of background pressure need to be characterized, in order to
 131 not be mistaken with an atmospheric wave of interest. The noise sources can be sum-
 132 marized by the following equation from Bowman et al. (2018):

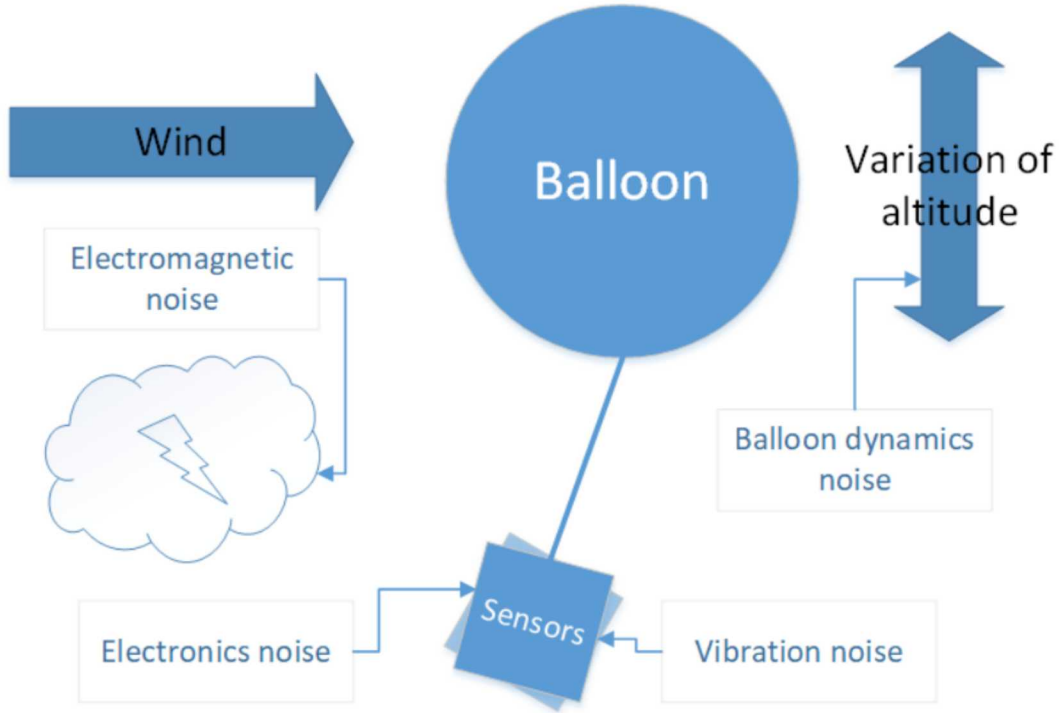
$$133 \quad \epsilon = \epsilon_v + \epsilon_e + \epsilon_s + \epsilon_w + \epsilon_p + \epsilon_a, \quad (1)$$

134 Where different constituents are:

- 135 • ϵ_v : stands for sensor motion, i.e. vibrations, which is removed from the signal by
 136 combination of two channels from active microphones (Bowman & Lees, 2018).
- 137 • ϵ_e : corresponds to the outside electromagnetic interference, the associated noise
 138 is shown on the signal coming from the mechanically disabled sensor (Bowman
 139 & Lees, 2018).
- 140 • ϵ_s : depicts the intrinsic sensor electronic noise, which is reduced by using several
 141 identical sensors and combining their signal.
- 142 • ϵ_w : stands for the wind generated noise. Because the balloon is flying in the strato-
 143 sphere, at neutral buoyancy, thus it experiences near zero differential air flow, hence
 144 the wind noise is negligible (Bowman & Lees, 2016).
- 145 • ϵ_p : depicts the non-hydrostatic pressure fluctuations, which come from the bal-
 146 loon's altitude variations. To estimate its level, we compute the pressure varia-
 147 tions sensed by the balloon function of its altitude.
- 148 • ϵ_a : describes the acoustic and gravity waves pressure signals of interest in this study.

149 It is important to notice that the mechanically disabled microphone is still slightly
 150 sensitive to the spatial pressure gradient caused by high frequency infrasound. Hence some
 151 of signal recorded on the mechanically disabled microphone can still be caused by pres-
 152 sure waves.

154 To compute the dynamics noise, we used the GPS altitude measurements and the
 155 reanalysis ERA5 (ECMWF, 2018). The high resolution re-analysis has a spatial reso-
 156 lution of 31 km, a temporal resolution of 1 hour and a vertical resolution of about 850
 157 m in the 30-35 km altitude range. Once the altitude is acquired, we use the atmospheric
 158 model to estimate the associated atmospheric pressure. This conversion is performed through
 159 spline interpolation. The GPS measurements are provided every 15 minutes, so we com-
 160 puted only the low-frequencies pressure variations. The sensors are differential barom-
 161 eters, thus it is impossible to subtract the computed noise from the signal. However, our



153 **Figure 1.** Schematic view of different noise sources sensed by the balloon during the flight.

162 computation allows us to estimate the power of the noise coming from the balloon dy-
 163 namics and to qualify whether the low-frequencies recorded signals are impacted by it.

164 **2.3 Methodology**

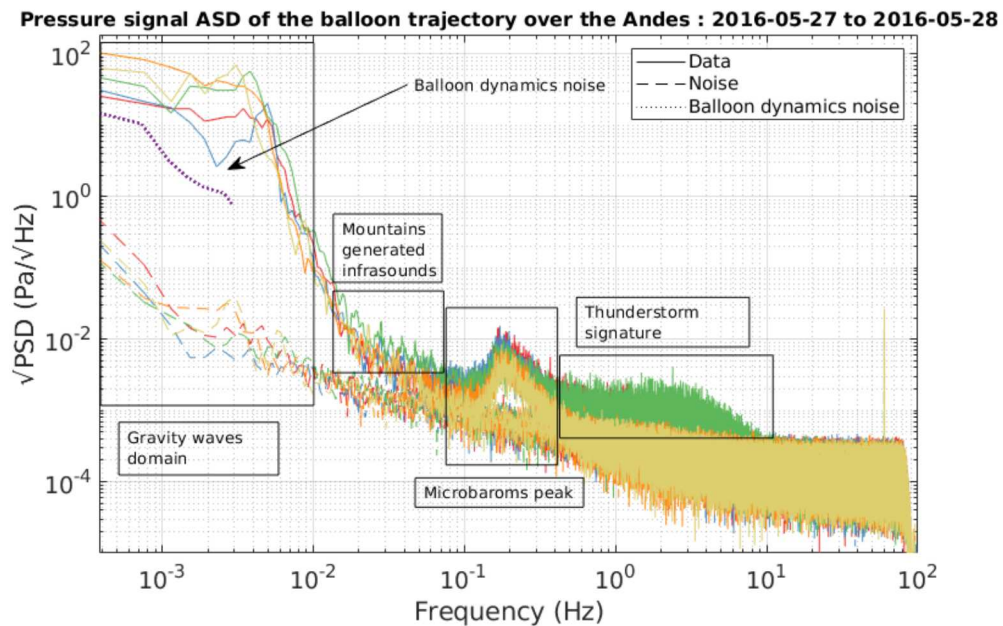
165 In Figure 2, we present the part of the balloon trajectory which was analyzed in
 166 this study. The trajectory is divided into 5 temporal windows of 1h30 each. The sub-
 167 division of the trajectory allows us to correlate the observed signal and the topographic
 168 feature over which the balloon is flying. The balloon trajectory starts at 23:00 UTC, or
 169 19:00 Local Time (LT), the 27 May 2016 and ends at 06:30 UTC (02:30 LT) the 28 May
 170 2016.

174 We conducted the data analysis by studying the Amplitude Spectral Density (ASD,
 175 defined as the square root of Power Spectral Density) of the pressure signal on each tem-
 176 poral window. The computed Pressure ASD of the balloon trajectory over the Andes is
 177 shown in Figure 3. First of all, we estimated the noise coming from balloon dynamics.
 178 This noise is a factor 2 to 10 lower than the observed long period pressure signals, thus
 179 its influence is neglected in our analysis. Secondly, we observe four different phenomena:

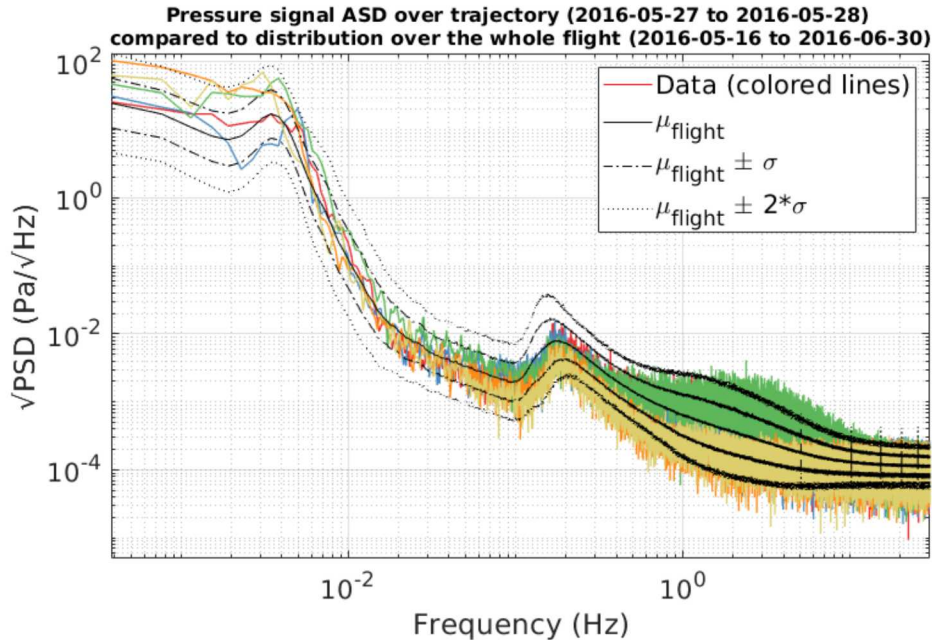
- 180 • Activity in the gravity waves domain, under 6 mHz.
- 181 • A factor 2 increase in the amplitude of pressure signal just above the mountains
- 182 between 0.02 Hz and 0.1 Hz compared to other time periods.
- 183 • Ocean microbarom signal peaking at $0.01 \text{ Pa}/\sqrt{\text{Hz}}$ between 0.12 Hz and 0.35 Hz.
- 184 • Broadband signal between 0.45 Hz and 2 Hz, above the Andes.



171 **Figure 2.** The part of the balloon flight above Andes, which was used for data analysis.
 172 The trajectory is divided into 5 temporal windows of 1h30 each, distinguishable by the different
 173 colors.



185 **Figure 3.** Pressure signal ASD of the balloon trajectory over Andes, divided into 5 temporal
 186 windows which are shown in Figure 2 with the same colors. The data (solid lines), the sensor
 187 noise (dashed lines) estimated from the mechanically disabled sensor, and the balloon dynamics
 188 noise (dotted line) are presented. Note the thin 60 Hz line corresponding to contamination by
 189 electric networks.



190 **Figure 4.** Pressure signal ASD of balloon trajectory over the Andes 3 represented with col-
 191 ored lines. ASD mean value over the whole flight (19.5 days of data mainly acquired above the
 192 ocean) μ_{flight} represented with dashed line, and its fluctuation intervals at one standard deviation
 193 (σ), represented with dash-dotted line, and two standard deviations (2σ), represented with
 194 dotted line.

195 In Figure 4, we show the previously computed ASDs, alongside the mean ASD value
 196 over the whole flight duration (19.5 days of data acquired mainly above the ocean) with
 197 its fluctuation intervals, at one and two standard deviation levels (σ and 2σ). To com-
 198 pute these statistical observations, we divided the whole data set on temporal windows
 199 of 1h30 duration each and applied the Welch spectral method in the same way as for com-
 200 putation of the Figure 3. From this statistical distribution, computed on more than 300
 201 time windows, we notice that the phenomena described above, especially peaks in the
 202 gravity waves domain, mountain infrasound and thunderstorm signature, are outside the
 203 one standard deviation fluctuation interval. Thus the observed phenomena are proba-
 204 bly not due to statistical fluctuations. In the following, we will present each of the de-
 205 tected phenomena and our analysis of associated geophysical processes.

206 3 Geophysical processes detected

207 3.1 Mountain gravity waves

208 Terrain features such as mountains and valleys can generate gravity waves. This
 209 kind of wave plays an important part in the global circulation by transporting the en-
 210 ergy away from the lower atmosphere to the higher levels. These waves have already been
 211 observed by stratospheric balloon platforms by using movements of the balloon (Hertzog
 212 et al., 2002; Vincent & Hertzog, 2014) and pressure variations (Quinn & Holzworth, 1987).
 213 The process behind the generation of these waves is the vertical displacement of the strat-
 214 ified flow when it encounters obstacles. The terrain generated waves are stationary rel-
 215 ative to the ground, because their intrinsic phase speed is equal to the background wind
 216 speed, but in opposite direction (Nappo, 2002).

217 The gravity wave domain is separated from the acoustic domain by the Brunt-Väisälä
 218 frequency N , which represents the maximum frequency allowed for gravity waves. Above
 219 this frequency, the gravity waves are evanescent. Therefore, it is necessary to estimate
 220 the value of Brunt-Väisälä frequency at the altitude of the balloon in order to be able
 221 to conclude if the observed signal has any physical meaning. We used the formula from
 222 (Lighthill, 1978):

$$223 \quad N^2 = -\frac{g}{\rho} \frac{\partial \rho}{\partial z} - \left(\frac{g}{c}\right)^2, \quad (2)$$

224 where g is the gravity, ρ is the fluid density, z is the altitude and c the sound speed. Mak-
 225 ing the assumption that the air is a perfect gas, we obtain this expression:

$$226 \quad N^2 = (\gamma - 1) \left(\frac{g}{c}\right)^2, \quad (3)$$

227 where γ is the adiabatic index. In order to compute the Brunt-Väisälä frequency at the
 228 balloon altitude, we used the ERA5 atmospheric model combined with MSIS-E-90 model
 229 (CCMC, 1990). The second model has been used because the atmospheric composition
 230 is needed by our codes to compute γ and additional wave attenuation parameters. Even-
 231 tually, we computed Brunt-Väisälä angular velocity profile along the balloon trajectory
 232 and at the altitude of 33 km, we obtained an average $N = 19$ mrad/s, corresponding
 233 to a frequency of 3.1 mHz. As shown in Figure 5, large signals with a lot of variability
 234 are observed in this frequency range and identified as gravity waves.

239 In Figure 5, we show the ASD and the temporal signal in the gravity waves spec-
 240 tral domain. First of all, we observe the presence of power peaks, which move from higher
 241 frequencies to lower frequencies with the balloon following its trajectory. The observed
 242 power peaks start at 4.969 mHz and end at 3.058 mHz. Anderson and Taback (1991)
 243 showed that the balloons can resonate at frequencies just below Brunt-Väisälä frequency,
 244 thus the observed power can be partly due to this resonance effect, which is also con-
 245 firmed by the presence of the peak in this range of frequencies on the mean value over
 246 the whole flight, as shown in Figure 4. However, in order to have the resonance, the wave
 247 energy must be present, thus we conclude that we detect gravity waves but the recorded
 248 amplitudes do not reflect the real amplitudes of the gravity waves. During the pass over
 249 the mountains, the pressure signals in the gravity wave frequency range present a pres-
 250 sure signal peak ASD of about $60 \text{ Pa}/\sqrt{Hz}$ at 3.9 mHz close to Brunt-Väisälä frequency.
 251 Directly after passing the mountains, the power peak disappears (in orange in Figure 5)
 252 and reappears in the next temporal window (in yellow in Figures 5), but with a greater
 253 peak ASD ($85 \text{ Pa}/\sqrt{Hz}$) than that recorded above mountains at a lower frequency (3
 254 mHz) (in green in Figure 5). On the temporal pass-band filtered signal, we observe the
 255 same behaviour with 4 Pa amplitudes above mountains, slightly smaller amplitudes di-
 256 rectly in the lee of the mountains and once again 4 Pa amplitude in the next time win-
 257 dow. However, with the additional hours of data, we note that the amplitudes are de-
 258 creasing with the distance from the mountains.

259 The enhanced power just above the mountains (in green in figures 3, 4 and 5) could
 260 be related to storm activity during the pass of the balloon (see section 3.4). However,
 261 the power of pressure signal is increasing again in the lee of the mountain ridge (in yel-
 262 low in figures 3, 4 and 5). This increase is associated to the presence of mountain grav-
 263 ity waves. According to various authors (Nappo, 2002; Plougonven et al., 2008; Ehard
 264 et al., 2017), the mountains gravity waves are stationary waves relative to the mountain
 265 ridge. When observed in the stratosphere, their wavelengths increase downwind and their
 266 amplitudes decrease with the distance from the ridge and with altitude. We compute the
 267 associated horizontal wavelengths by using the expression for stationary gravity waves
 268 from (Nappo, 2002):

$$269 \quad \lambda = \frac{2\pi u_0}{\Omega}, \quad (4)$$

270 where u_0 is the horizontal component of the wind and Ω is the frequency recorded by
 271 the balloon sensors. Therefore, the observed horizontal wavelengths move from 50.066

272 km to 73.653 km with increasing distance from the ridge. This trend, from smaller wave-
 273 lengths to larger wavelengths, when following the downwind balloon trajectory is in agree-
 274 ment with the expected behaviour of the mountain gravity waves. In addition, the shift
 275 of maximum energy relative to the ridge in the downwind direction (Figure 5) can be
 276 explained by the upward and downwind propagation of mountain gravity waves (Iwasaki
 277 et al., 1989; Ehard et al., 2017). This energy increase cannot be justified by the slight
 278 decrease of balloon altitude after the Andes, because only 4% amplitude increase of pres-
 279 sure variations are expected. The hypothesis of detection of mountain gravity waves is
 280 strengthened by the seasonal observations of high activity in gravity waves domain above
 281 Andes from March to October (Hoffmann et al., 2013; Liu et al., 2019).

282 Besides, we also detected these gravity waves analyzing the Atmospheric Infrared
 283 Sounder (AIRS) data (Aumann et al., 2003), shown in Figure 6. AIRS is an instrument
 284 on board of NASA's Aqua satellite measuring without interruption the mid-infrared nadir
 285 and sub-limb radiance spectra since 2002. To perform our analysis, we used AIRS level
 286 1B radiance data (NASA, 2009). The data has up to 13 km horizontal resolution. Fol-
 287 lowing the studies of Hoffmann and Alexander (2009), Hoffmann et al. (2013), Sato, Tsuchiya,
 288 Alexander, and Hoffmann (2016) and Hoffmann, Grimsdell, and Alexander (2016), we
 289 selected the data from the 4.3 μm band to sound the stratosphere. Example of spectra
 290 in this band are shown in Figure 6. As this band is excited by solar radiation during the
 291 day time (non-LTE effect), we observe that the signal is weaker during the night. To ob-
 292 tain the best signal to noise ratio, we only selected the 4.304 μm and 4.28 μm wavelengths
 293 which are associated with radiance peaks. Then, the obtained image was pass-band fil-
 294 tered, the results are also shown in Figure 6. During the daytime, we observe quite dis-
 295 tinguishable gravity waves on the AIRS data, which are much less visible during the night-
 296 time. We also represented the filtered balloon data on the same figure, along the bal-
 297 loon trajectory. The peaks does not match perfectly with the observed gravity waves due
 298 to several hours time difference between the two observations. However, the geograph-
 299 ical position and the apparent wavelength of the gravity waves from the balloon mea-
 300 surements agree with the AIRS observations.

310 **3.2 Infrasonds generated over mountains**

311 Mountain-generated infrasound is still a poorly-understood phenomenon; very few
 312 theoretical and experimental works treat this subject (Larson et al., 1971; Chimonas, 1977;
 313 Bedard, 1978; Chunchuzov, 1994; Hupe, 2018). According to these works, the mountains
 314 infrasound are probably generated by the the wind blowing into the mountain valleys
 315 or to turbulence at mountain ridges. The spectral range associated to this phenomenon
 316 is between 0.025 Hz and 0.8 Hz. Figure 7 shows that there is a factor 2 increase of pres-
 317 sure signals in this band just above the Andes, evident in both the ASD and on the tem-
 318 poral representation. Our results are in agreement with (Hupe, 2018) that is reporting
 319 a significant activity in the same frequency range in this particular region in the May-
 320 June time period. This study also suggest that these acoustic waves can be emitted all
 321 year long, their detection by ground based infrasound sensors being limited by varying
 322 atmospheric conditions. In addition, the amplitudes of about 2 mPa Root Mean Square
 323 (RMS) observed at balloon altitudes in this narrow frequency range (0.025-0.08 Hz) are
 324 similar to the ones reported from ground sensors by (Hupe, 2018) (20 mPa RMS), once
 325 the pressure amplitude decrease between ground and balloon altitude is taken into ac-
 326 count (about a factor 10). In addition, due to the low attenuation of these long period
 327 signals in the stratospheric waveguide, peaks of amplitudes observed downstream my be
 328 due to the propagation of the acoustic waves generated by the mountains in this direc-
 329 tion. This interpretation is reinforced by the apparent periodicity of these signals with
 330 about 200 km distance (Figure 7.c), which is approximately the horizontal distance cov-
 331 ered by infrasonds reflected in the stratosphere waveguide. Therefore, we hypothesize
 332 that we observed mountain-generated infrasound during this time period.

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3.3 Ocean microbarom

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The ocean microbarom originates from non linear interactions between ocean waves (Landès et al., 2014). These interactions generate stationary waves, which in turn create acoustic waves in the atmosphere. The ocean microbarom are the atmospheric equivalent of microseisms (Longuet-Higgins, 1950). The associated spectral range is between 0.12 Hz and 0.35 Hz. The signal is detectable all over the world, and consequently is part of acoustic background on Earth. It can be detected on the ground when tropospheric winds are low and stratosphere wind ducts refract them back towards the surface (Campus & Christie, 2009; Landès et al., 2014; Bowman & Lees, 2018).

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Figure 8 shows balloon data in the spectral range of microbaroms. The power in this range decreases with increasing distance from the ocean. The decrease is progressive, but the microbaroms peak does not disappear entirely. The same progressive decrease is observed on the temporal signal. A factor two difference of power before and after the mountains is observed. However, this decrease is much less severe than what is predicted by full-wave numerical simulations for direct wave paths: typically a loss factor of $\simeq 10$ over $\simeq 70$ km distance from the oceanic source (Martire, Garcia, et al., 2018). This observation demonstrates that the ocean microbarom pressure signals felt by balloon sensors are not only coming from just below the balloon, or along direct paths, but a large part of this signal is due to waves multiply reflected in the stratosphere waveguide. Numerical simulations using GeoAcGlobal software (Blom, 2014), which is using atmospheric attenuation parameters from (Sutherland & Bass, 2004), in the downstream (i.e. eastern) direction demonstrate that the reflected paths propagate over an horizontal distance of about 180 km and present attenuations by 6 to 2 dB in between each surface reflections. Assuming a continuous and extended oceanic source in the pacific ocean, these simulations predict a decrease of microbarom amplitude of about a factor 1.25 to 2 over 180 km horizontal distance inland. The microbarom amplitude decrease of about 3 dB over 200 km (Figure 8.c) is consistent with these predictions. As a consequence, our observations demonstrate that microbaroms are propagating globally in the stratospheric waveguide far from the oceanic sources, in a way similar to microseismic noise propagating globally far from the oceans.

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3.4 Thunderstorm acoustic signature

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Thunderstorms can also generate infrasound. According to Lamb et al. (2018), the acoustic signal of thunderstorms is only detected in the vicinity of the event. Besides, the thunderstorm has broad band spectral signature between 0.5 Hz and a few Hz. Since the balloon flies above a desert zone, very few meteorological stations are available. Thus, we did not obtain any precise information on the passage of the storm in this zone. However, using the data provided by the World Wide Lightning Location Network (WWLLN), three thunderbolts were detected in this region before and after the passage of the balloon:

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- Date: 2016-05-27 12:23 UTC, Lat: -36.5825° , Long: -63.915°
- Date: 2016-05-28 21:27 UTC, Lat: -36.5254° , Long: -73.3508°
- Date: 2016-05-28 23:53 UTC, Lat: -36.5254° , Long: -79.8525°

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The balloon trajectory we analyzed is between 2016-05-27 23:00 UTC and 2016-05-28 06:30 UTC, with a mean latitude of -36.5° and longitudes between -77° and -67° . The thunderstorm data support the presence of a storm during these days in the Andes. This is also supported by the presence of clouds and convection in the 8.1 m AIRS data acquired before the balloon passage on 2016-05-27 14:30 LT in this region (not shown). The exact space/time correlation between convective activity and the balloon data cannot be established but all the data sets favor strong convective activity during the passage of the balloon above the Andes.

396 On the ASD, shown in Figure 9, we observe a broadband signal, between 0.45 Hz
 397 and 2 Hz, which corresponds to the thunderstorm signal, as described by Lamb et al. (2018).
 398 As shown in Figure 9, the signal is suddenly interrupted in time. It could be explained
 399 by the fact that the acoustic signal originating from the storm can be detected only above
 400 the storm clouds, directly near to the event. That is why, the balloon lost the signal im-
 401 mediately it flew out of the direct vicinity of the thunderstorm clouds. Besides, if we ob-
 402 serve the signal of the mechanically disabled sensor, bottom panel of Figure 9, we no-
 403 tice that the acoustic signal comes along with electromagnetic perturbations. These con-
 404 firms the hypothesis of the storm origin of the signal.

405 4 Conclusion

406 Using the ULDB observations, we were able to characterize several infrasound and
 407 gravity wave sources as the balloon-borne microbarometers crossed over a large moun-
 408 tain range. We have detected gravity waves above and after the mountain ridge. The waves
 409 observed exactly above the Andes are probably due to storms in the troposphere because
 410 convective clouds and lightning activity were reported, and thunderstorms acoustic and
 411 electromagnetic perturbations are detected by the pressure sensors. However, we expect
 412 the increase of pressure signal energy below Brunt-Väisälä frequency after the Andes to
 413 be due to the mountain gravity waves because it follows frequency shifts (from 4 mHz
 414 to 3 mHz) and a decrease of amplitude as a function of distance to the mountain ridge
 415 which are expected for such waves. In addition, we have analyzed the AIRS data, which
 416 confirmed the presence of gravity waves in this zone during the passage of the balloon.
 417 We have also observed a signal, a factor 2 larger than average background signal in the
 418 0.025-0.08 Hz frequency range, which can be associated with the infrasounds generated
 419 by the interactions between the wind and the mountains. Finally, we analyzed the de-
 420 crease of ocean microbarom power with increasing distance from the coast. Because this
 421 decrease is much smaller than the one predicted for direct acoustic paths, it strongly sug-
 422 gests that microbaroms are the main part of stratospheric acoustic noise, and can be ex-
 423 plained by signals trapped in the stratospheric waveguide and propagating over large dis-
 424 tances.

425 The broad range of phenomena observed through a single stratospheric flight is strongly
 426 supporting the development of long duration high altitude balloon flights and associated
 427 payloads. This work is part of a demonstration of the interest of high altitude and high
 428 rate pressure sensors for sensing planetary atmospheres and internal structure. The phe-
 429 nomena, signals and wave sources characterized here are also expected by to be impor-
 430 tant sources of signal or noise in the atmosphere of Venus (Stevenson et al., 2015). Their
 431 quantification in the Earth environment allows to build a science case for Venus and other
 432 planets with dense atmospheres.

433 Acknowledgments

434 The authors wish to thank the NASA ULDB team, AIRS instrument team of Aqua NASA
 435 mission and World Wide Lightning Location Network (<http://wwlln.net>) for provid-
 436 ing the data used in this paper. We also acknowledge three anonymous reviewers for their
 437 contributions to the improvement of this scientific paper. This study was funded by "Di-
 438 rection Générale de l'Armement" (French DoD), "Région Occitanie" and CNES research
 439 projects. ULDB data used in our study was retrieved from online repository: [https://](https://doi.org/10.5061/dryad.40877s4)
 440 doi.org/10.5061/dryad.40877s4. AIRS data were obtained from the AIRS data repos-
 441 itory (<https://airs.jpl.nasa.gov/data/get.data>). This paper describes objective techni-
 442 cal results and analysis. Any subjective views or opinions that might be expressed in the
 443 paper do not necessarily represent the views of the U.S. Department of Energy or the
 444 United States Government.

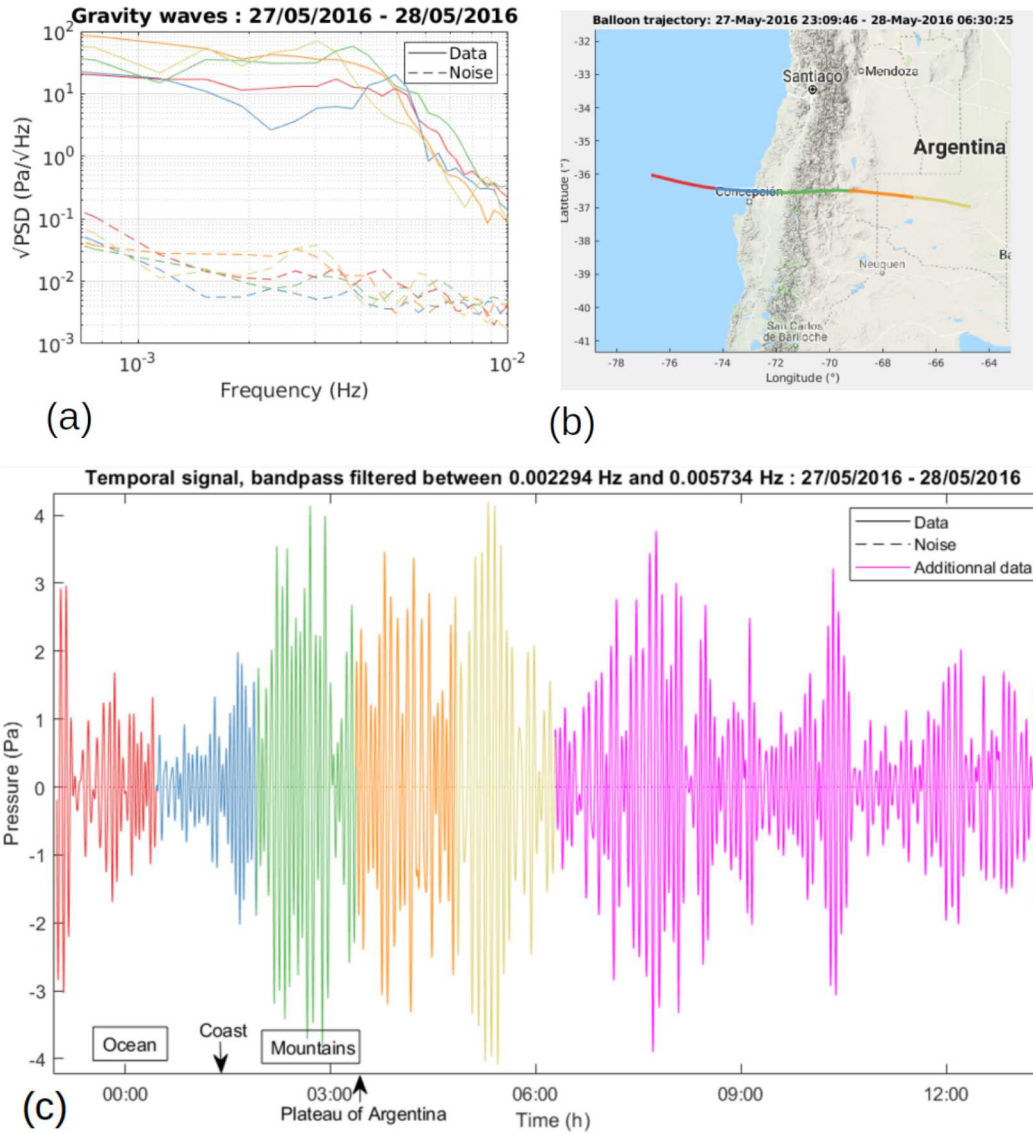
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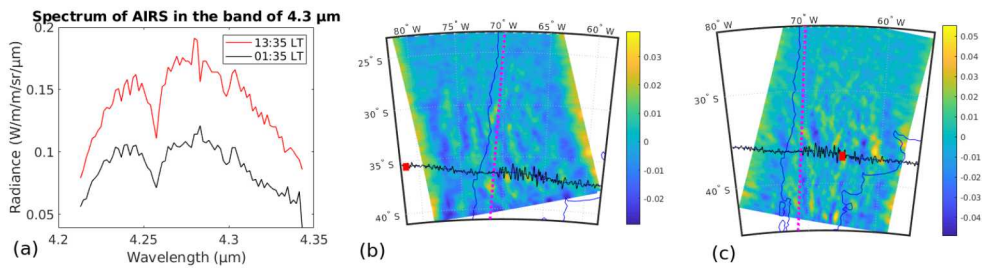
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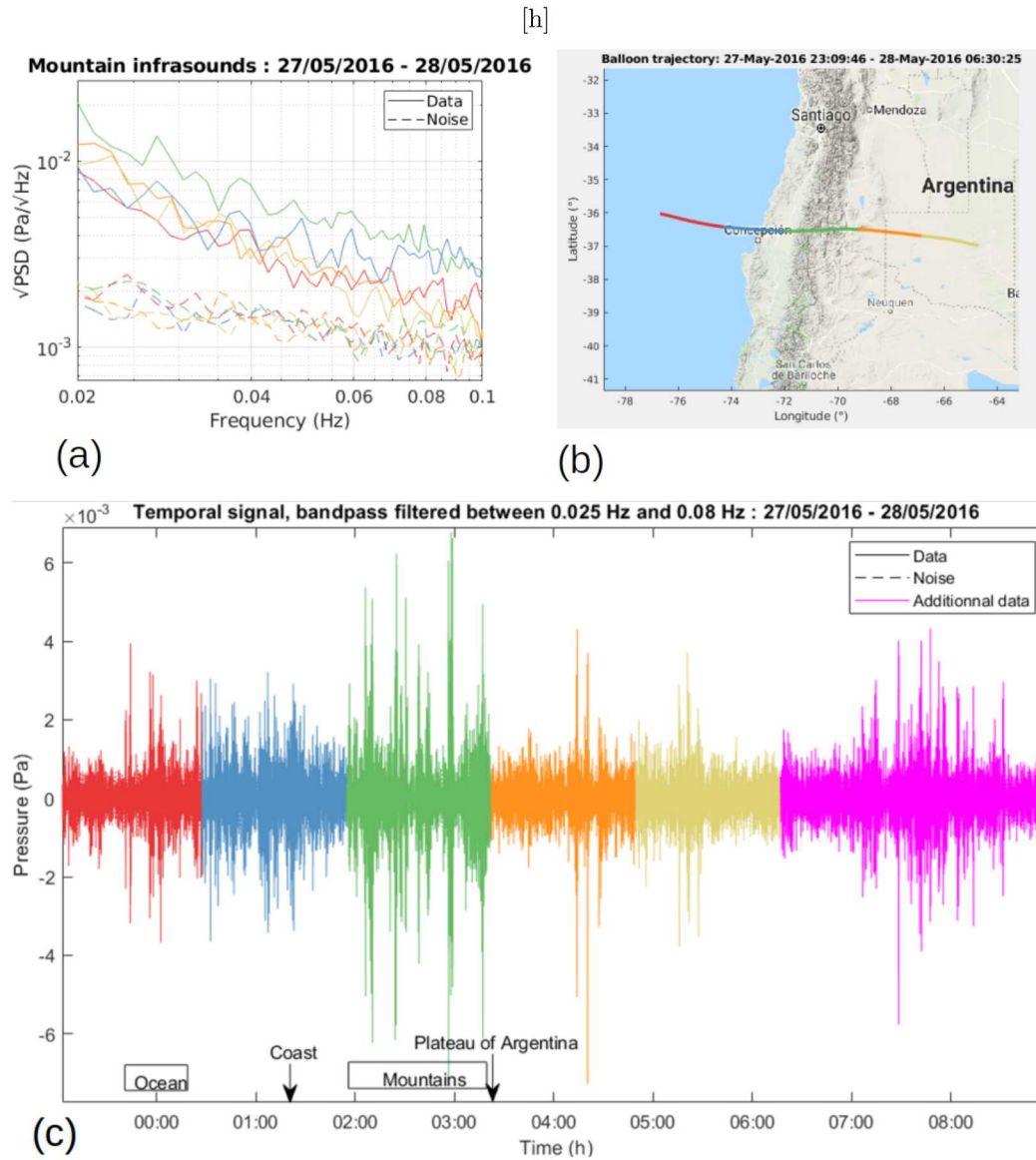
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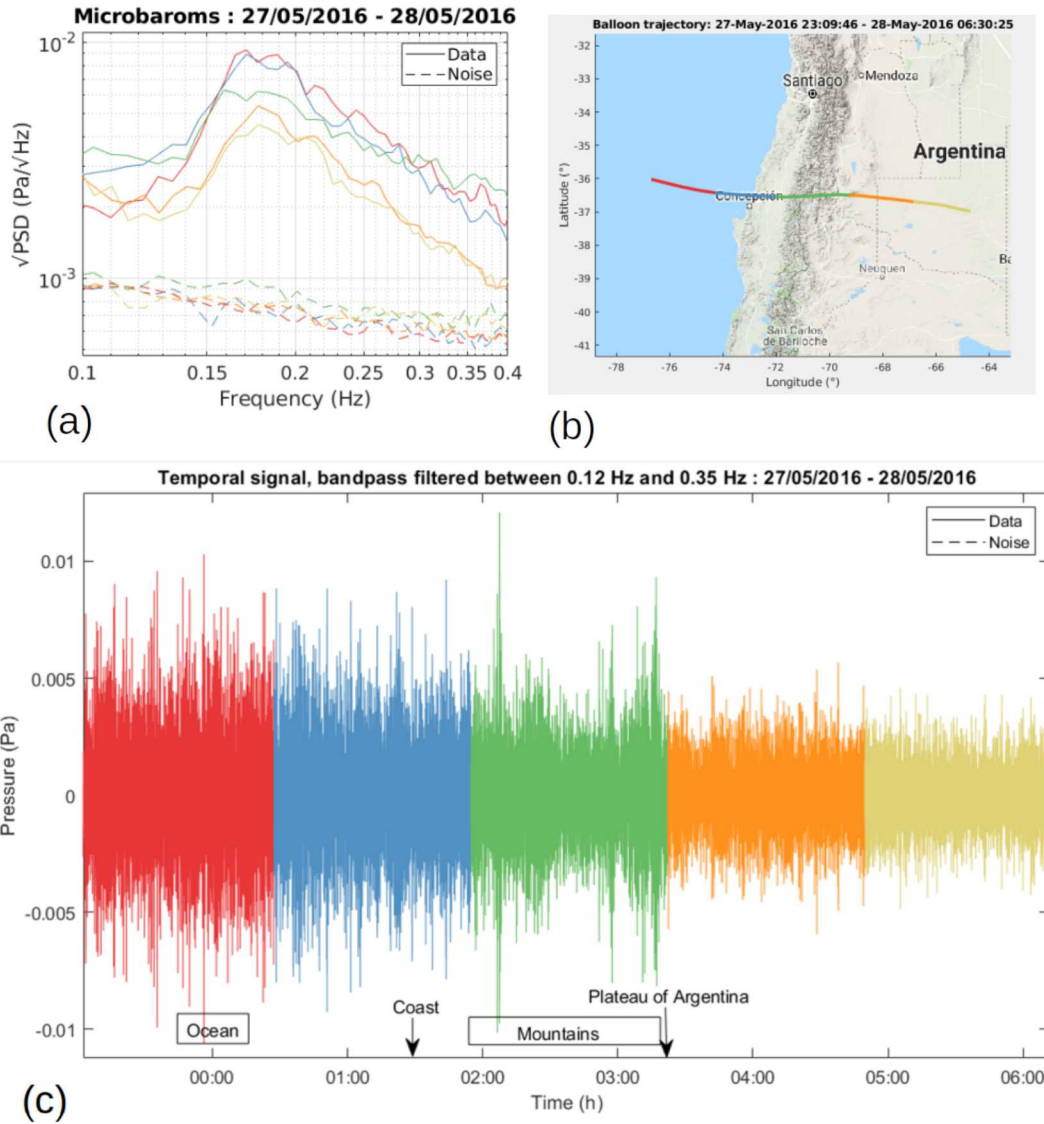
235 **Figure 5.** Gravity waves analysis. Top left (a): the computed ASD over the balloon trajec-
 236 tory, between 0.7 mHz and 10 mHz. Top right (b): the trajectory of the balloon. Bottom (c): the
 237 temporal acoustic signal of the studied balloon trajectory with a few hours of additional data,
 238 pass-band filtered between 2.294 mHz and 5.734 mHz.



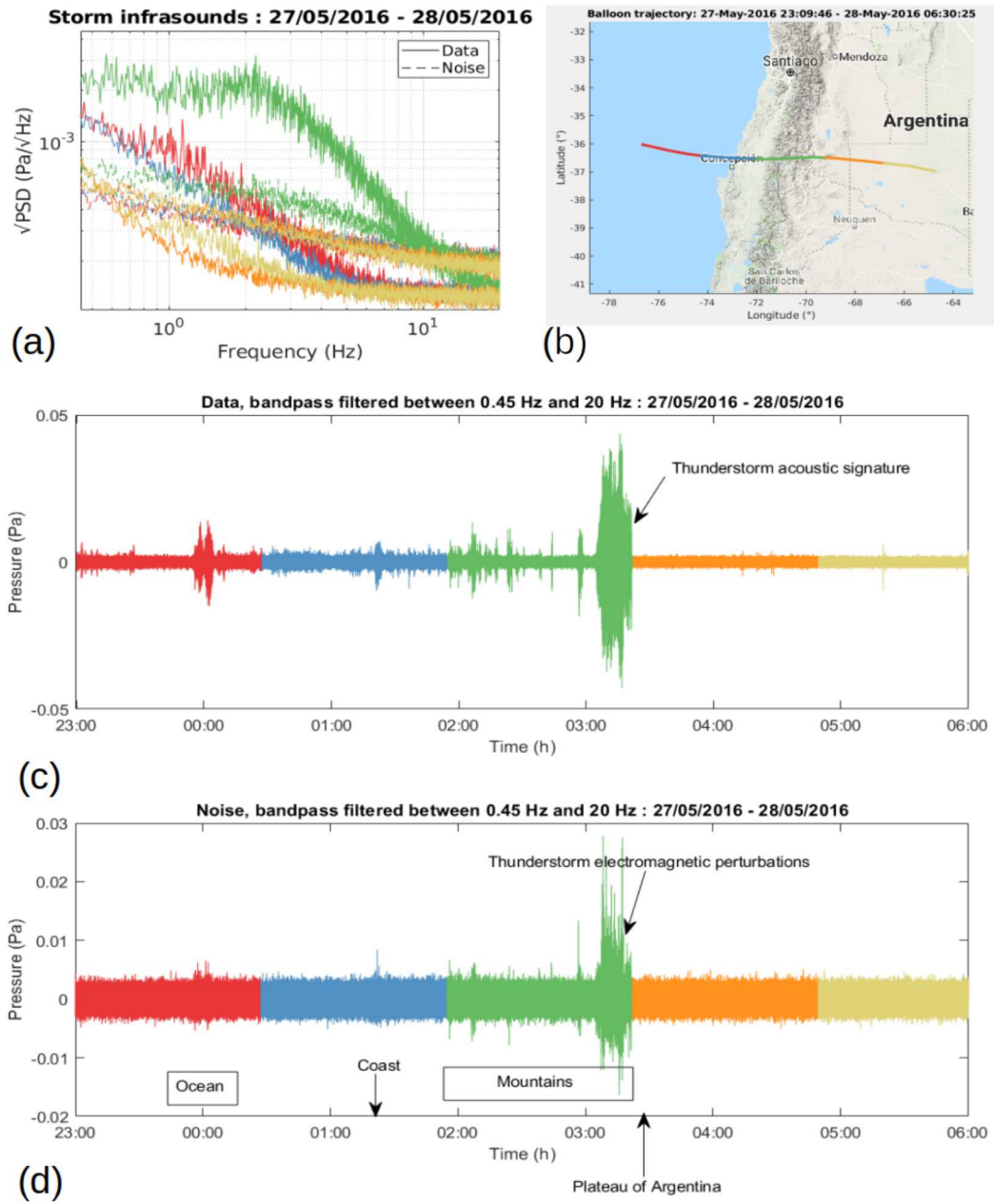
301 **Figure 6.** AIRS radiance data (in $\text{W}\cdot\text{m}^{-2}\cdot\text{sr}^{-1}\cdot\mu\text{m}^{-1}$). (a) Spectrum of AIRS data in the 4.3
 302 μm band taken above the Andes at 13:35 LT (day time) and 01:35 LT (night time). (b) and (c)
 303 short wavelength perturbations in AIRS data at 4.3 μm for images acquired respectively at 13:35
 304 LT (day time) and 01:35 LT (night time). Color bar indicates the amplitude of 4.3 μm radiance
 305 variations (in $\text{W}\cdot\text{m}^{-2}\cdot\text{sr}^{-1}\cdot\mu\text{m}^{-1}$). Black line represents the pressure perturbations (band pass
 306 filtered in the 0.4-2 mHz range) along the balloon's trajectory between 2016-05-27 19:06 GMT
 307 (13:06 LT) and 2016-05-28 10:15 GMT (04:15 LT). Red square indicates the balloon position at
 308 the time of AIRS data acquisition. The magenta dotted line represents the mean longitude of the
 309 Andes.



333 **Figure 7.** Infrasound generated by the mountains. Top left (a): the computed ASD over the
 334 balloon trajectory, between 20 mHz and 100 mHz. Top right (b): the trajectory of the balloon.
 335 Bottom (c): the temporal acoustic signal of the studied balloon trajectory, pass-band filtered
 336 between 0.025 Hz and 0.08 Hz.



367 **Figure 8.** Microbaroms signal. Top left (a): the computed ASD over the balloon trajectory,
 368 between 0.1 Hz and 0.4 Hz. Top right (b): the trajectory of the balloon. Bottom (c): the tempo-
 369 ral acoustic signal of the studied balloon trajectory, pass-band filtered between 0.12 Hz and 0.35
 370 Hz.



391 **Figure 9.** Thunderstorm signal. Top left (a): the computed ASD over the balloon trajec-
 392 tory, between 0.4 Hz and 20 Hz. Top right (b): the balloon trajectory. Middle panel (c): the
 393 temporal signal of the studied balloon trajectory, pass band filtered between 0.45 Hz and 20 Hz.
 394 Bottom panel (d): temporal signal of mechanically disabled sensor used to estimate sensor and
 395 electromagnetic noises.

Figure 1.

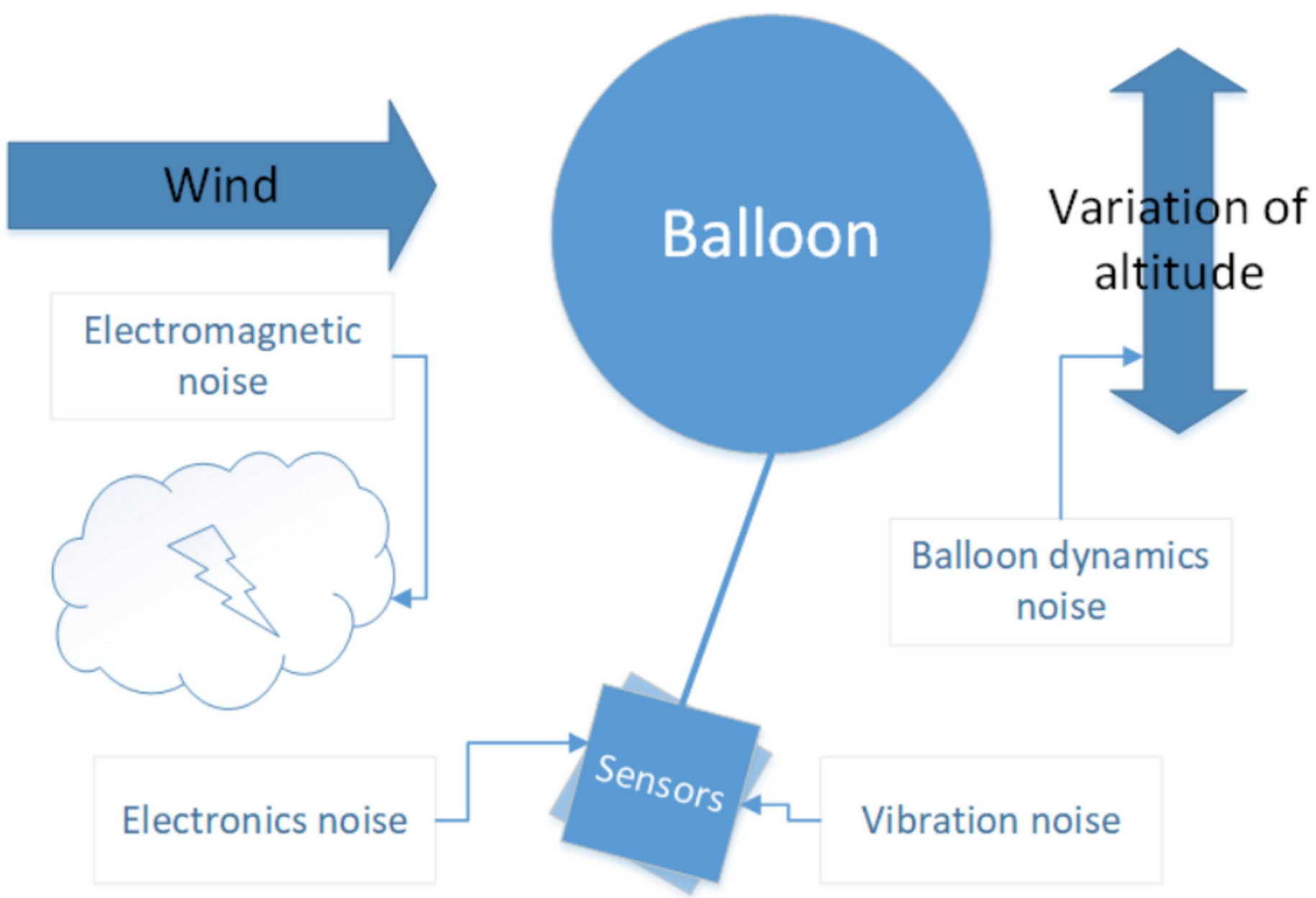


Figure 2.

Balloon trajectory: 27-May-2016 23:09:46 - 28-May-2016 06:30:25

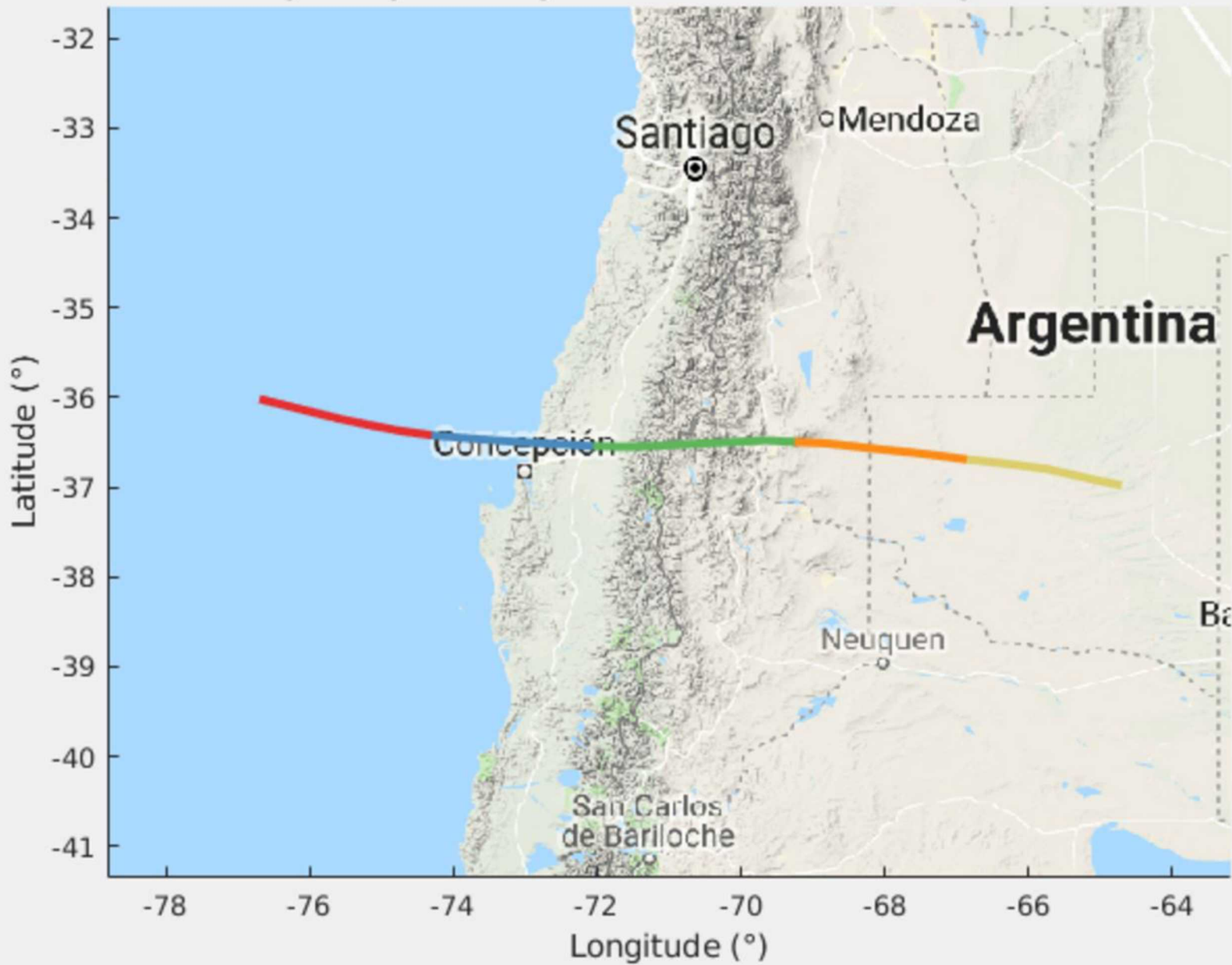


Figure 3.

Pressure signal ASD of the balloon trajectory over the Andes : 2016-05-27 to 2016-05-28

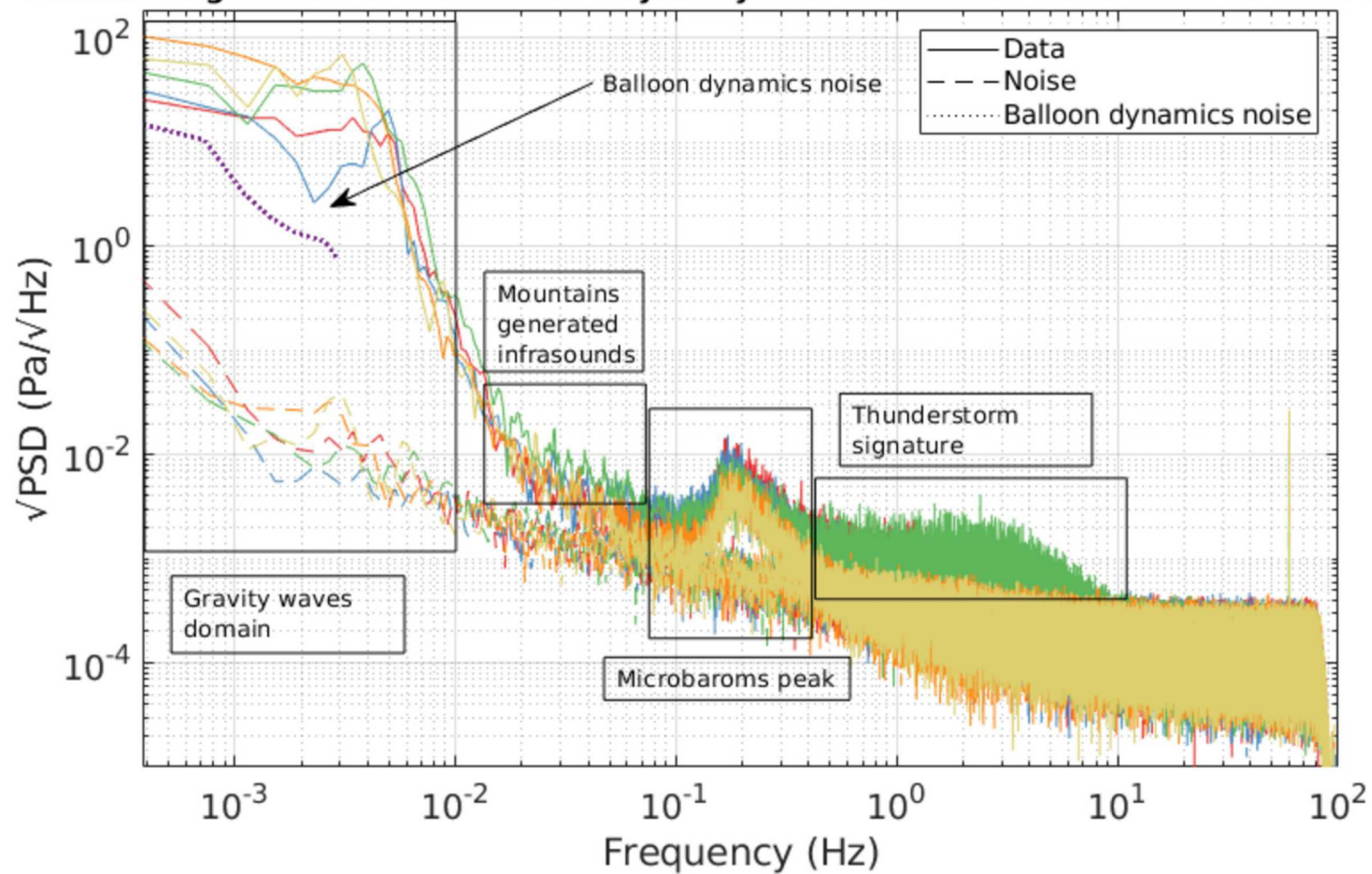


Figure 4.

Pressure signal ASD over trajectory (2016-05-27 to 2016-05-28)
compared to distribution over the whole flight (2016-05-16 to 2016-06-30)

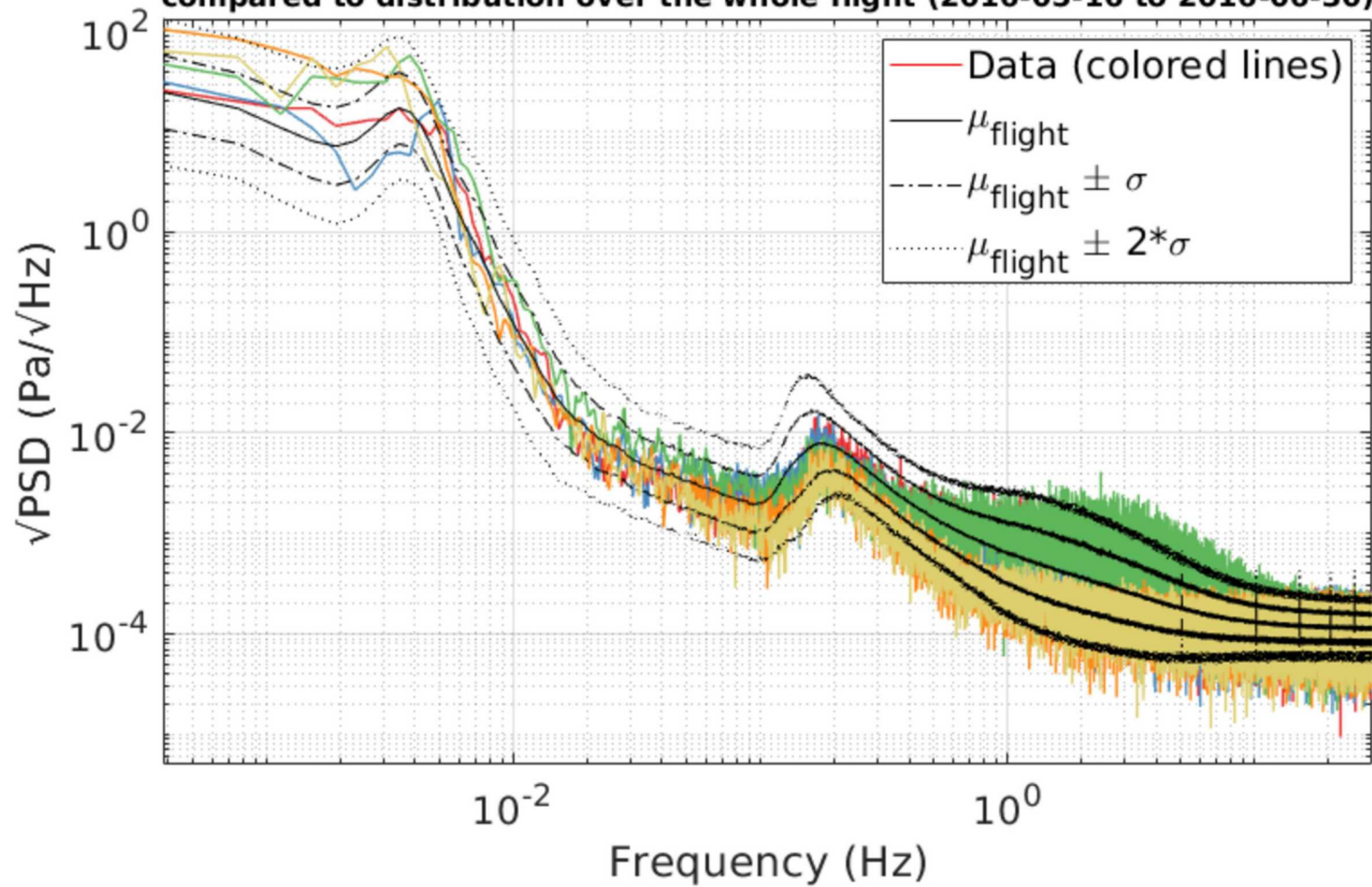
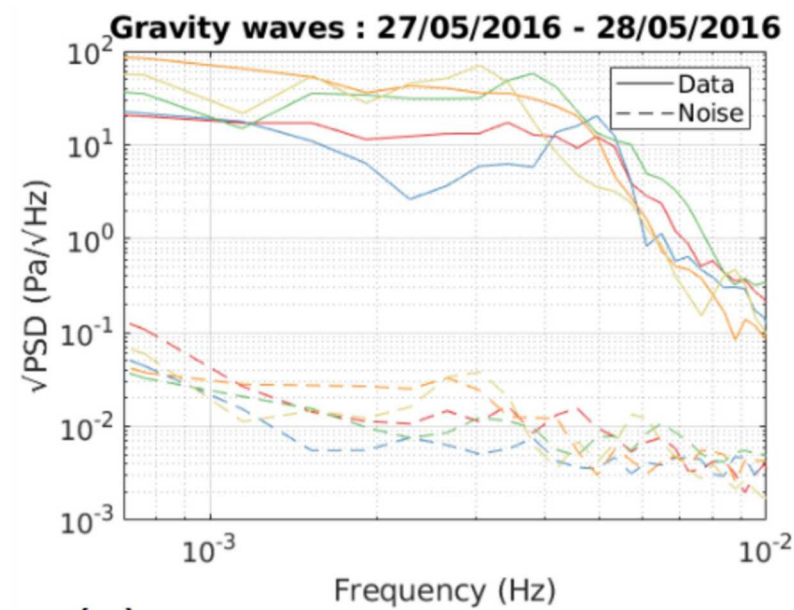
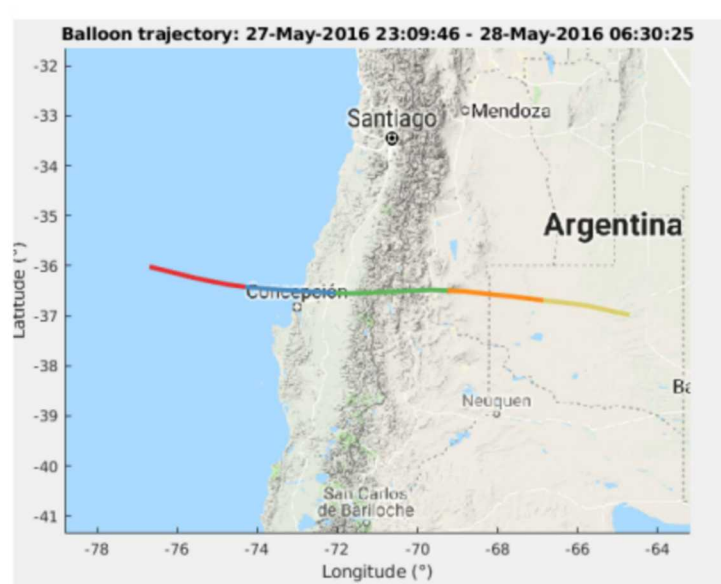


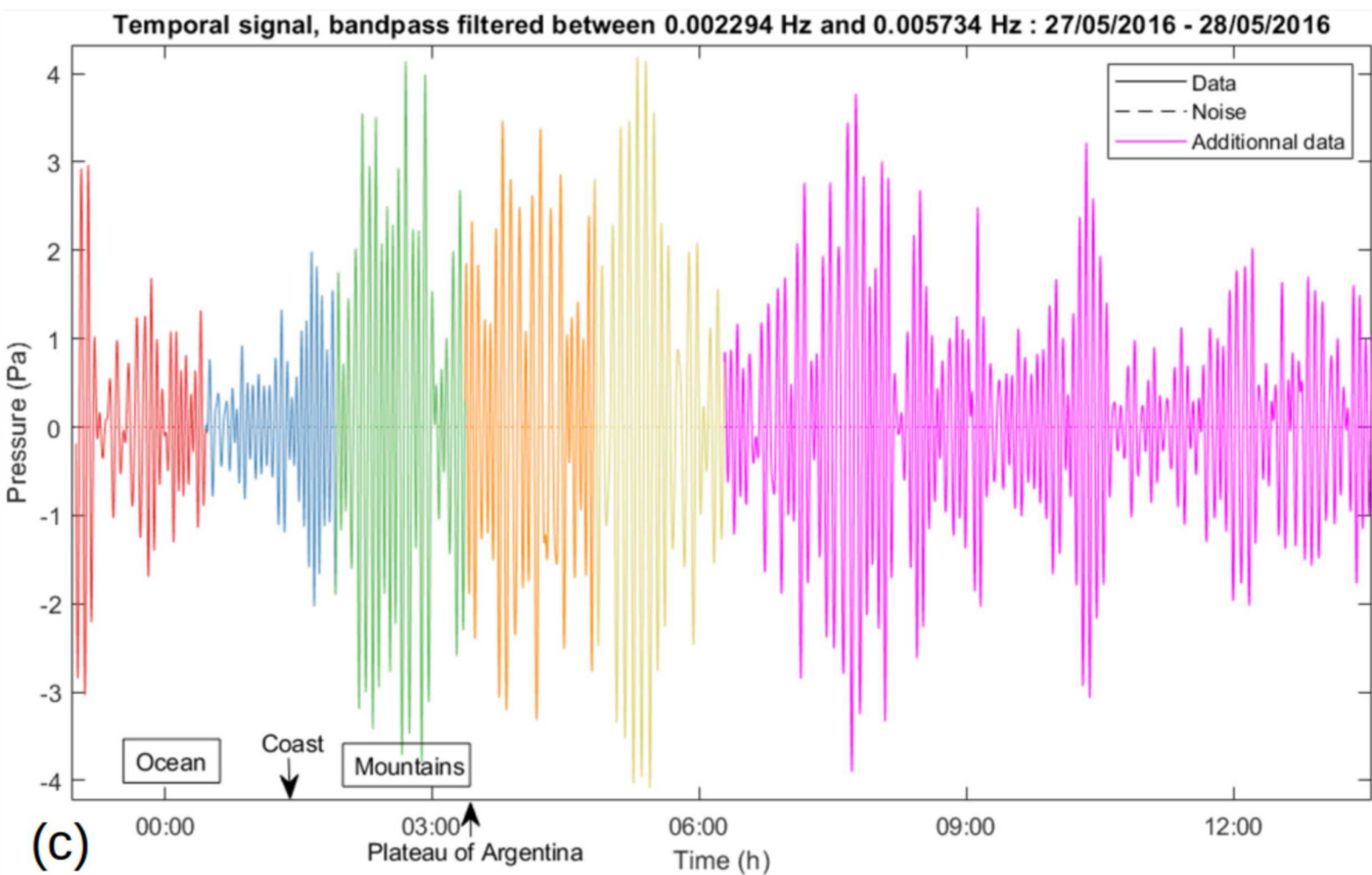
Figure 5.



(a)



(b)



(c)

Figure 6.

Spectrum of AIRS in the band of 4.3 μm

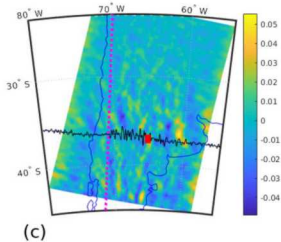
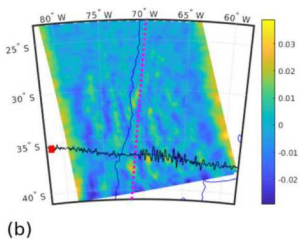
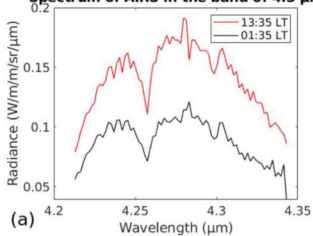
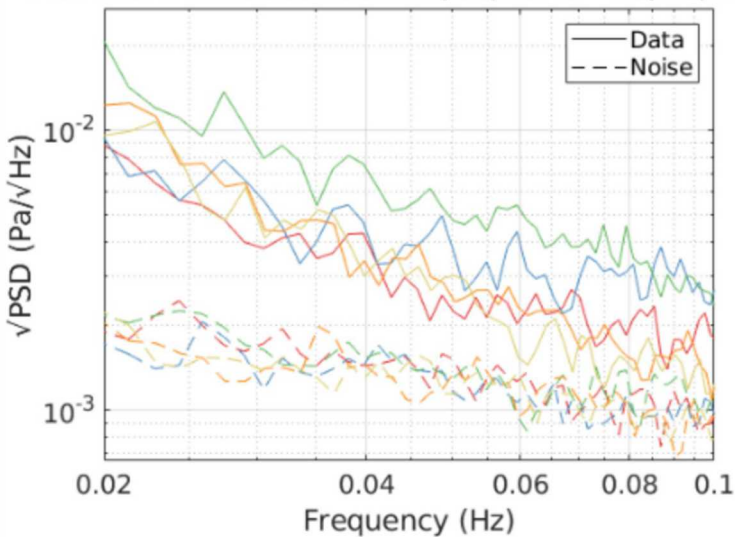


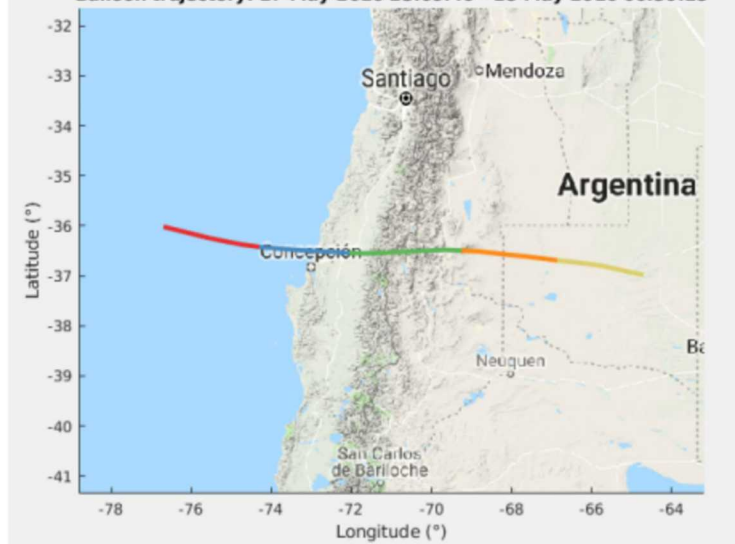
Figure 7.

Mountain infrasounds : 27/05/2016 - 28/05/2016



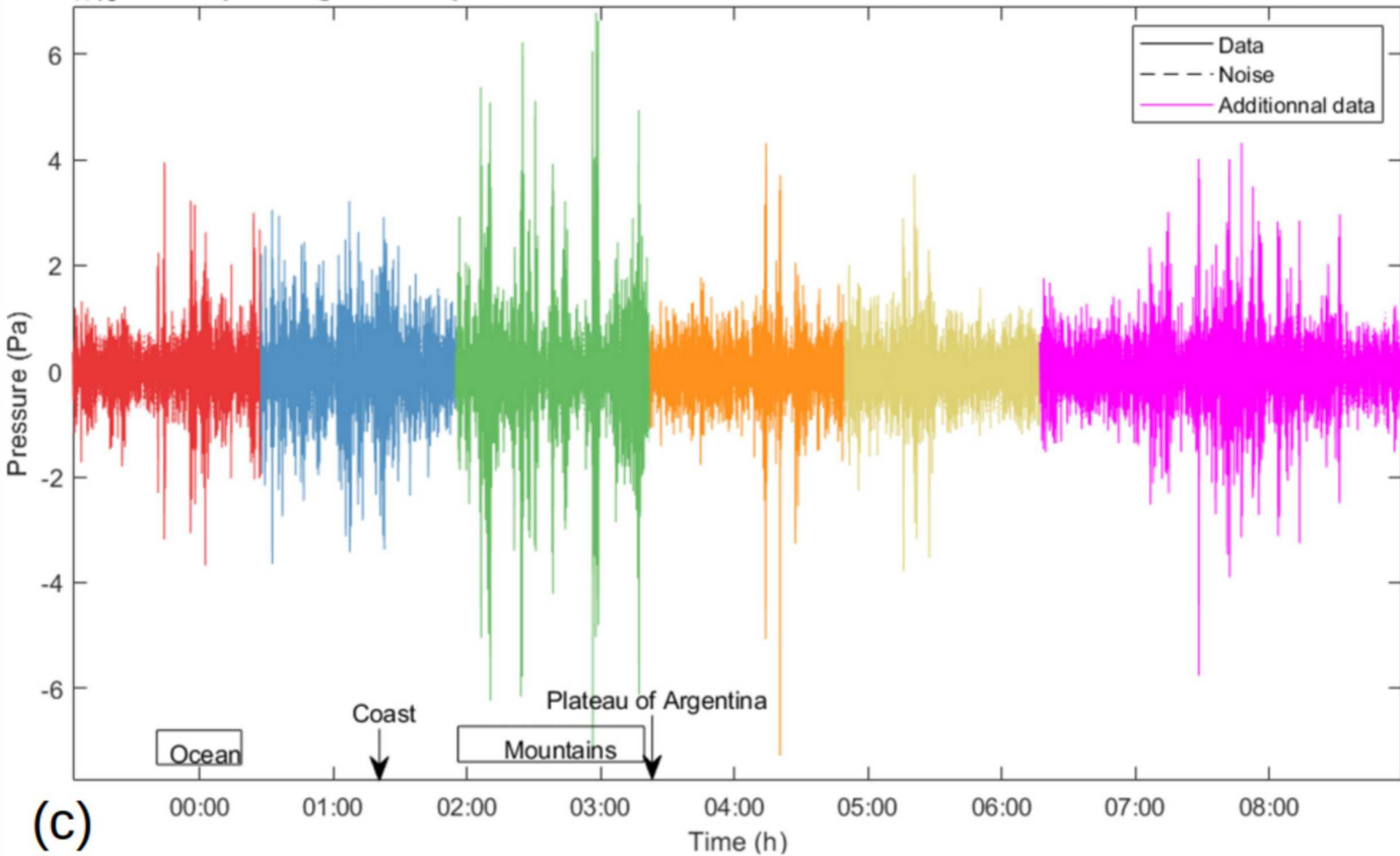
(a)

Balloon trajectory: 27-May-2016 23:09:46 - 28-May-2016 06:30:25



(b)

Temporal signal, bandpass filtered between 0.025 Hz and 0.08 Hz : 27/05/2016 - 28/05/2016



(c)

Figure 8.

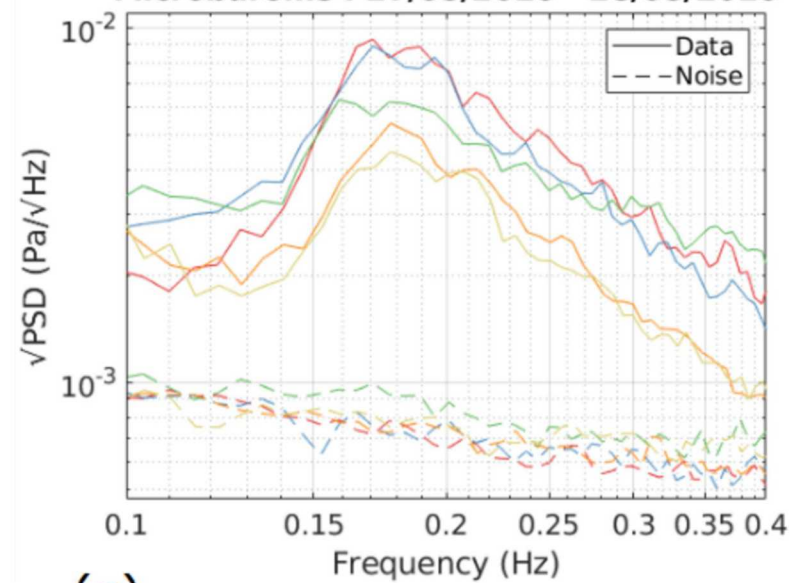
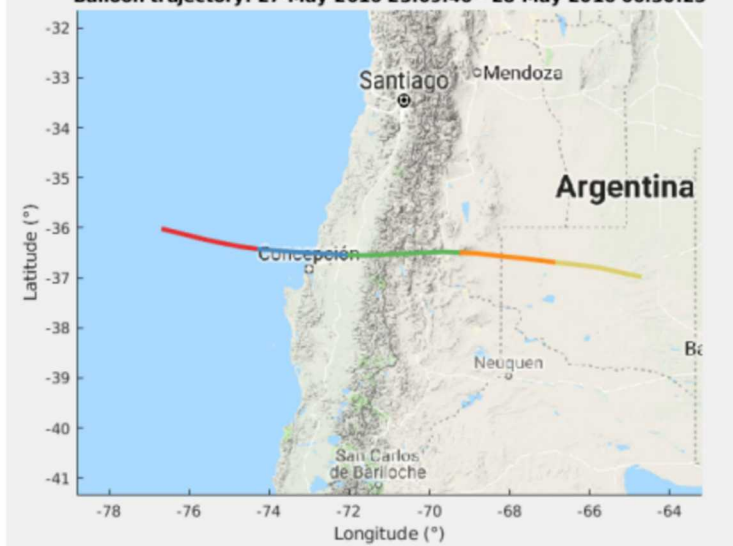
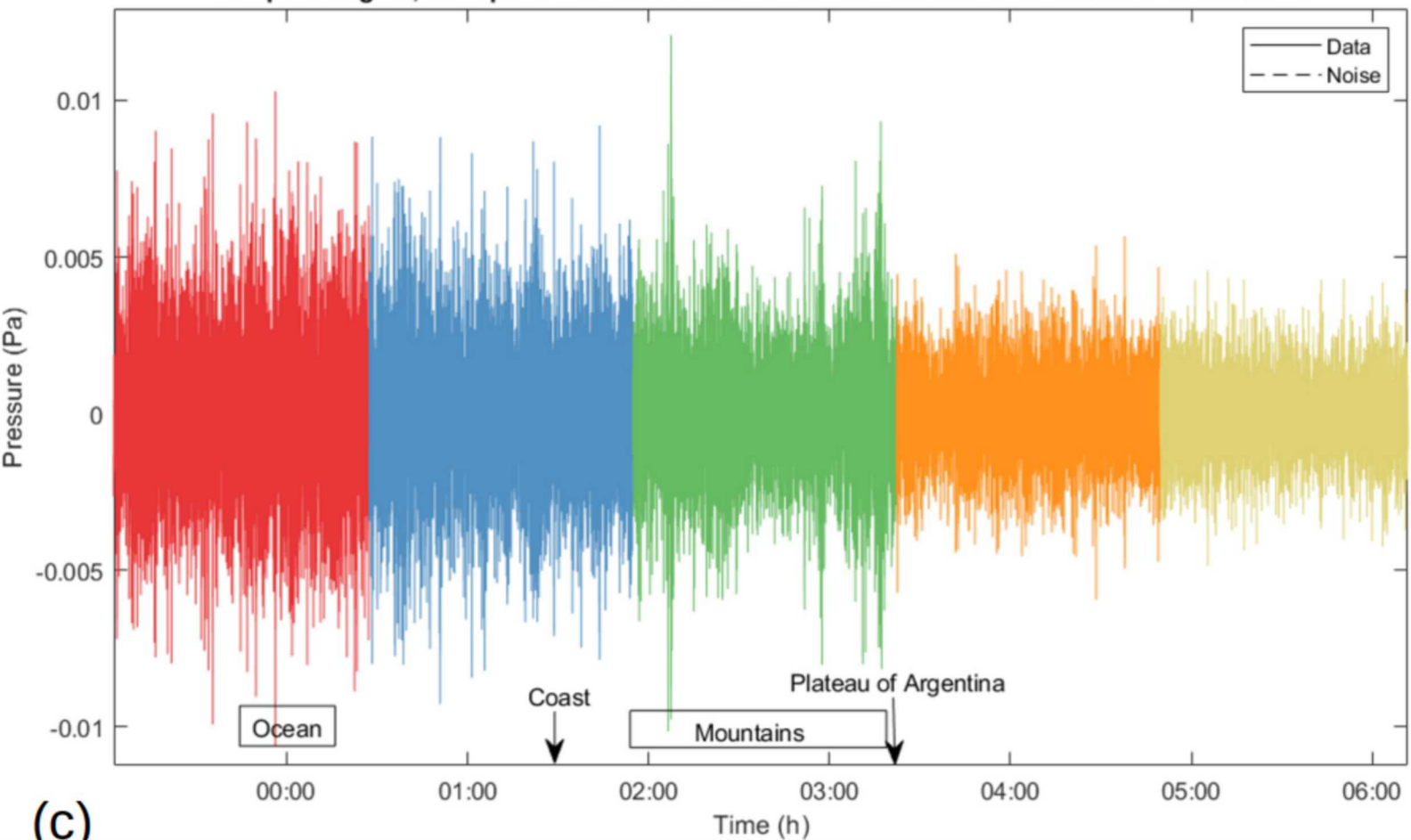
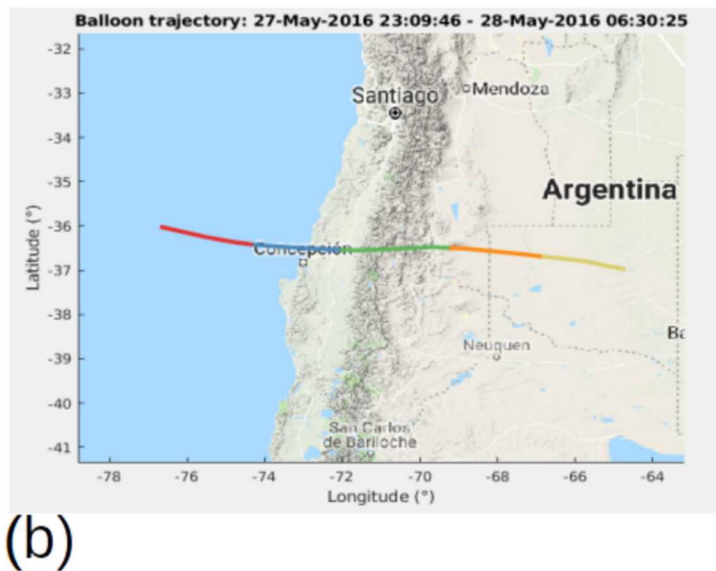
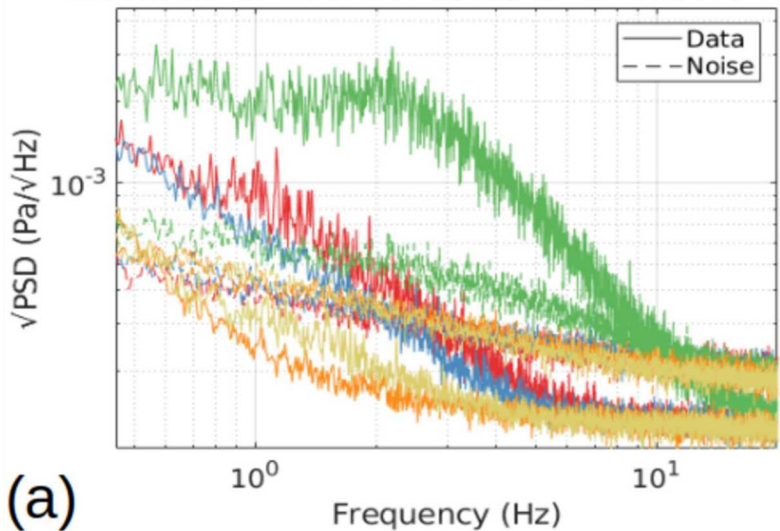
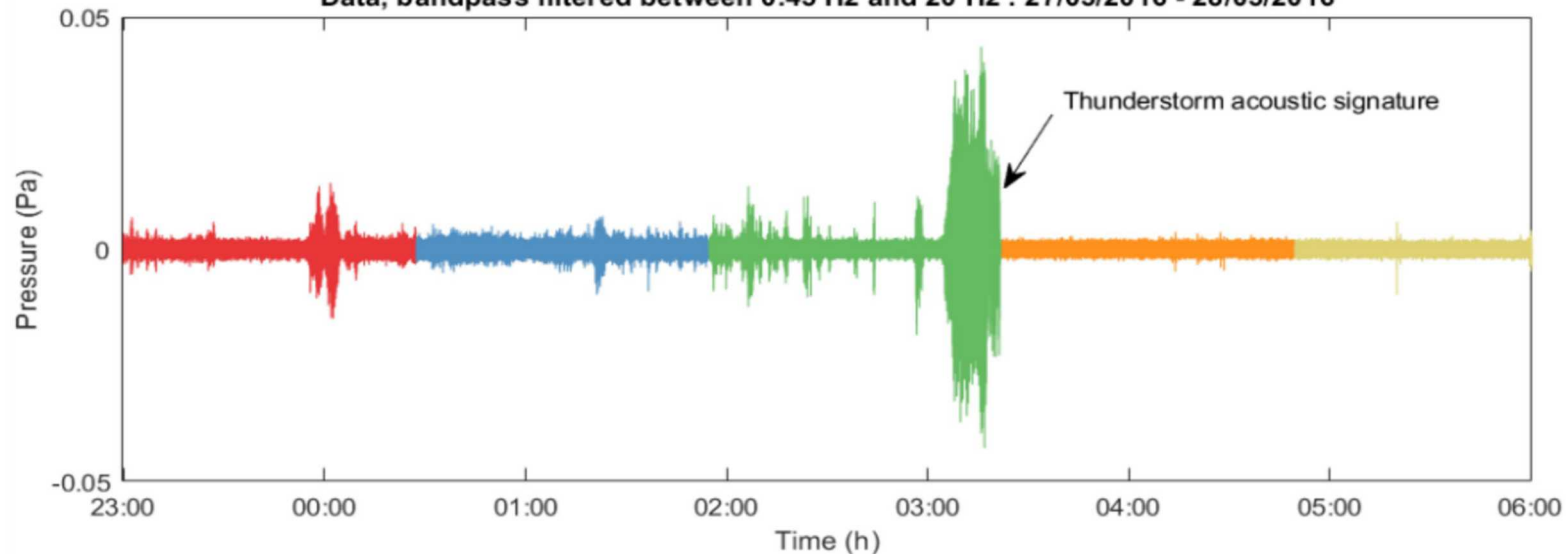
Microbaroms : 27/05/2016 - 28/05/2016**(a)****Balloon trajectory: 27-May-2016 23:09:46 - 28-May-2016 06:30:25****(b)****Temporal signal, bandpass filtered between 0.12 Hz and 0.35 Hz : 27/05/2016 - 28/05/2016****(c)**

Figure 9.

Storm infrasounds : 27/05/2016 - 28/05/2016



Data, bandpass filtered between 0.45 Hz and 20 Hz : 27/05/2016 - 28/05/2016



Noise, bandpass filtered between 0.45 Hz and 20 Hz : 27/05/2016 - 28/05/2016

