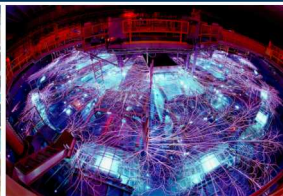


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Black Box Multigrid for MHD

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1. Black Box Multigrid (BoxMG)^{1,2} structured algebraic algorithm to build hierarchy
2. Test extension to coupled PDE systems³
3. Problems motivated by magnetohydrodynamics (MHD)
4. Use Drekar (Sandia MHD/CFD code) to assemble linear systems, solve with Matlab BoxMG code

¹J. E. Dendy. "Black Box Multigrid". In: *Journal of Computational Physics* (1982).

²J. E. Dendy Jr. and J.D. Moulton. "Black box multigrid with coarsening by a factor of three". In: *Numerical Linear Algebra with Applications* (2010). DOI: 10.1002/nla.705.

³J.E. Dendy Jr. "Black Box Multigrid for Systems". In: *Applied Mathematics and Computation* 19 (1986).

Black Box Multigrid

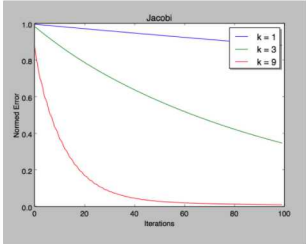
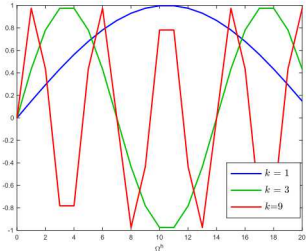
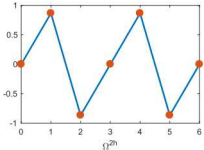
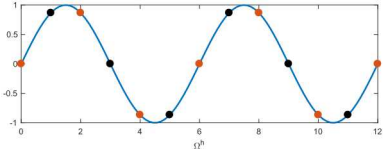
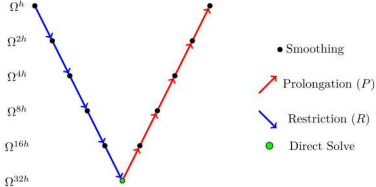
Scalar BoxMG

Vector BoxMG

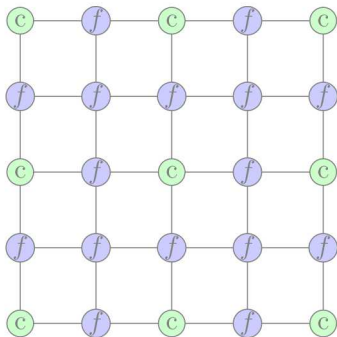
Magnetohydrodynamics (MHD)

Numerical Results

Multigrid Overview



$$A = \begin{bmatrix} A_{ff} & A_{fc} \\ A_{cf} & A_{cc} \end{bmatrix}$$



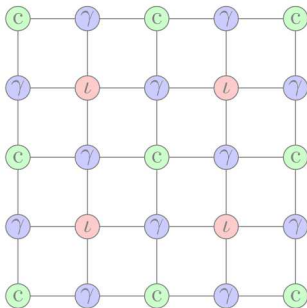
‘Perfect’ interpolation

$$P_{fc} = -A_{ff}^{-1}A_{fc}, \quad P_{cc} = I, \quad P = \begin{bmatrix} P_{fc} \\ P_{cc} \end{bmatrix}, \quad R = P^T$$

Galerkin projection produces Schur complement:

$$RAP = A_{cc} - A_{cf}A_{ff}^{-1}A_{fc}$$

Infeasible because A_{ff}^{-1} is dense!



$$A = \left[\begin{array}{cc|c} A_{ll} & A_{l\gamma} & A_{lc} \\ A_{\gamma l} & A_{\gamma\gamma} & A_{\gamma c} \\ \hline A_{cl} & A_{c\gamma} & A_{cc} \end{array} \right] \xrightarrow{\text{collapse}} \left[\begin{array}{cc|c} A_{ll} & A_{l\gamma} & A_{lc} \\ 0 & \hat{A}_{\gamma\gamma} & \hat{A}_{\gamma c} \\ \hline A_{l'c}^T & A_{\gamma c}^T & A_{cc} \end{array} \right].$$

if $\hat{A}_{\gamma\gamma}$ diagonal

$$A_{ff} = \begin{bmatrix} A_{ll} & A_{l\gamma} \\ 0 & \hat{A}_{\gamma\gamma} \end{bmatrix}$$

upper triangular \rightarrow easily inverted

Row for interior γ node with 9-point Laplacian:

$$[\dots -1 \quad -1 \quad -1 \dots -1 \quad 8 \quad -1 \dots -1 \quad -1 \quad -1 \dots]$$

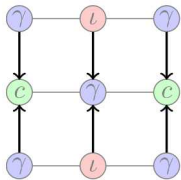


Figure: γ point on x coarse line

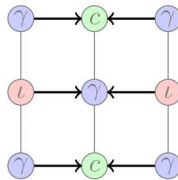


Figure: γ point on y coarse line

$$\begin{bmatrix} -1 & -1 & -1 \\ -1 & 8 & -1 \\ -1 & -1 & -1 \end{bmatrix} \xrightarrow{\text{collapse } y} \begin{bmatrix} 0 & 0 & 0 \\ -3 & 6 & -3 \\ 0 & 0 & 0 \end{bmatrix}$$

$$\begin{bmatrix} -1 & -1 & -1 \\ -1 & 8 & -1 \\ -1 & -1 & -1 \end{bmatrix} \xrightarrow{\text{collapse } x} \begin{bmatrix} 0 & -3 & 0 \\ 0 & 6 & 0 \\ 0 & -3 & 0 \end{bmatrix}$$

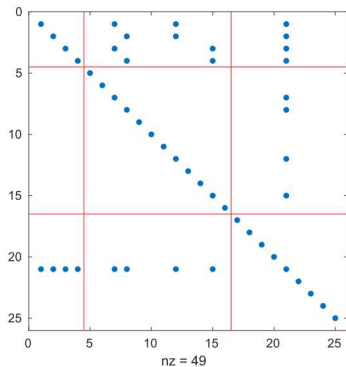
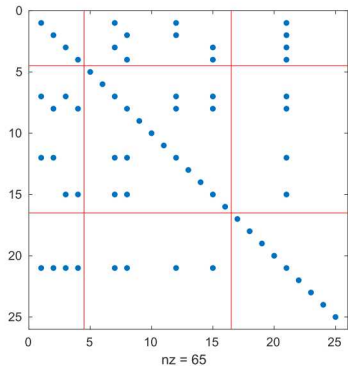
Accomplish two things:

- ▶ Decoupling: $\widehat{A}_{\gamma\gamma}$ now diagonal, $A_{\gamma l} = 0$
- ▶ Preserve action of matrix on constant vector: for Laplacian preserves nullspace

Take ∇^2 with 9-point stencil on a 5x5 mesh in 2d

$$\left[\begin{array}{cc|c} A_{\ell\ell} & A_{\ell\gamma} & A_{\ell c} \\ A_{\ell\gamma}^T & A_{\gamma\gamma} & A_{\gamma c} \\ \hline A_{\ell c}^T & A_{\gamma c}^T & A_{cc} \end{array} \right]$$

$$\left[\begin{array}{cc|c} A_{\ell\ell} & A_{\ell\gamma} & A_{\ell c} \\ 0 & \hat{A}_{\gamma\gamma} & \hat{A}_{\gamma c} \\ \hline A_{\ell c}^T & A_{\gamma c}^T & A_{cc} \end{array} \right]$$



$$A = \begin{bmatrix} A_{11} & A_{12} \\ A_{21} & A_{22} \end{bmatrix}$$

Block prolongators,

- ▶ $P_1 = \text{BoxMG}(A_{11})$, $R_1 = P_1^T$
- ▶ $P_2 = \text{BoxMG}(A_{22})$, $R_2 = P_2^T$

monolithic operators

$$P = \begin{bmatrix} P_1 & 0 \\ 0 & P_2 \end{bmatrix}, \quad R = \begin{bmatrix} R_1 & 0 \\ 0 & R_2 \end{bmatrix}$$

So

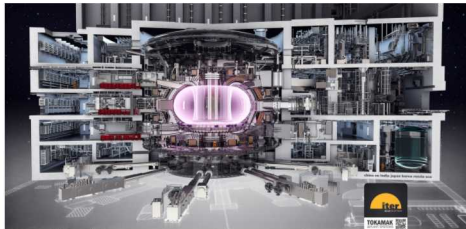
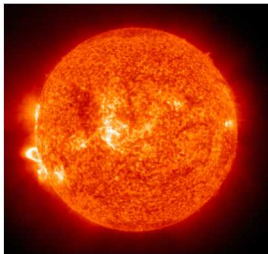
$$RAP = \begin{bmatrix} R_1 A_{11} P_1 & R_1 A_{12} P_2 \\ R_2 A_{21} P_1 & R_2 A_{22} P_2 \end{bmatrix}$$

Diagonal operators are being projected well

Black Box Multigrid
Scalar BoxMG
Vector BoxMG

Magnetohydrodynamics (MHD)

Numerical Results



- ▶ Continuum model of plasma states
- ▶ Mathematical Model: Continuum representation of collisional plasma systems.
- ▶ Structure of simple resistive MHD is Navier-Stokes + Lorentz force coupled with low frequency Maxwell equations.

Stationary incompressible MHD system

$$-\eta \Delta \mathbf{u} + \rho \mathbf{u} \cdot \nabla \mathbf{u} + \nabla p - \frac{1}{\mu_0} (\nabla \times \mathbf{B}) \times \mathbf{B} = \mathbf{f},$$

$$\nabla \cdot \mathbf{u} = 0,$$

$$\nabla \times \left(\frac{\eta}{\mu_0} \nabla \times \mathbf{B} \right) - \nabla \times (\mathbf{u} \times \mathbf{B}) = \mathbf{0},$$

$$\nabla \cdot \mathbf{B} = 0.$$

- ▶ Unknowns: magnetic field \mathbf{B} , velocity \mathbf{u} , pressure p
- ▶ Physical parameters: resistivity η , density ρ , magnetic permeability of free space μ_0 .

Take induction equation as PDE system

$$\nabla \times \left(\frac{\eta}{\mu_0} \nabla \times \mathbf{B} \right) - \nabla \times (\mathbf{u} \times \mathbf{B}) = \mathbf{0}$$

- ▶ $\nabla \times \nabla \times$ has large nullspace (gradients of all scalars)
- ▶ Introduce **penalty term** in weak form

$$\int_{\Omega} \frac{\eta}{\mu_0} (\nabla \cdot \mathbf{B})(\nabla \cdot \mathbf{C}) - \frac{\eta}{\mu_0} (\nabla \times \mathbf{B}) \cdot (\nabla \times \mathbf{C}) + \mathbf{C} \cdot [\nabla \times (\mathbf{u} \times \mathbf{B})] \, dx = 0$$

- ▶ Well posedness proved in [Gunzburger et. al]⁴
- ▶ Augmented strong form,

$$\nabla \times \left(\frac{\eta}{\mu_0} \nabla \times \mathbf{B} \right) - \frac{\eta}{\mu_0} \nabla (\nabla \cdot \mathbf{B}) - \nabla \times (\mathbf{u} \times \mathbf{B}) = \mathbf{0}$$

⁴Max D. Gunzburger, Ammon J. Meir, and Janet S. Peterson. "On the Existence, Uniqueness, and Finite Element Approximation of Solutions of the Equations of Stationary, Incompressible Magnetohydrodynamics". In: *Mathematics of Computation* 56 (1991).

$$\nabla \times \left(\frac{\eta}{\mu_0} \nabla \times \mathbf{B} \right) - \frac{\eta}{\mu_0} \nabla (\nabla \cdot \mathbf{B}) - \nabla \times (\mathbf{u} \times \mathbf{B}) = \mathbf{0}$$

- ▶ What have we gained?
- ▶ Vector identity $\nabla \times (\nabla \times \mathbf{B}) - \nabla (\nabla \cdot \mathbf{B}) = -\nabla^2 \mathbf{B}$
- ▶ We have “completed” vector Laplacian
- ▶ Nullspace reduction to constants: good for multigrid

Black Box Multigrid
Scalar BoxMG
Vector BoxMG

Magnetohydrodynamics (MHD)

Numerical Results

Let $\Omega = (0, 1)^2$, $\Gamma = \partial\Omega$,
 $f(x, y) = 4\pi^2 \sin(2\pi x) \sin(2\pi y)$.

$$\begin{cases} \nabla^2 u = f, & x \in \Omega \\ u = 0, & x \in \Gamma \end{cases}$$

True solution

$$u(x, y) = -\sin(2\pi x) \sin(2\pi y).$$

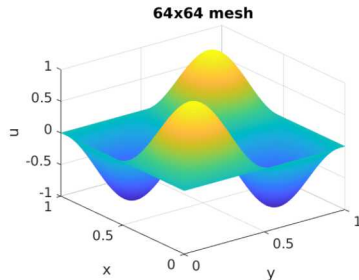


Figure: Solution plot

BoxMG for Scalar Poisson (tol = 1e-7)				
# Elements	Total DOFs	V Cycles	Jacobi Sweeps	Levels
32^2	1089	7	(3,3)	5
64^2	4225	8	(3,3)	6
128^2	16641	8	(3,3)	7
256^2	66049	8	(3,3)	8
512^2	264169	8	(3,3)	9

Let $\Omega = (0, 1)^2$, $\Gamma = \Gamma_D \cap \Gamma_N$, where $\Gamma_D = \{x = 0 \text{ or } x = 1\}$, and $\Gamma_P = \{y = 0 \text{ or } y = 1\}$. $\mathbf{B} = (B_1, B_2)^T$, $\frac{\eta}{\mu_0} = 1$,

$$\begin{cases} \nabla \times \left(\frac{\eta}{\mu_0} \nabla \times \mathbf{B} \right) - \frac{\eta}{\mu_0} \nabla (\nabla \cdot \mathbf{B}) = \mathbf{g}(\mathbf{x}), & \mathbf{x} \in \Omega \\ B_1 \equiv 0, & \mathbf{x} \in \Gamma \\ B_2 \equiv 0, & \mathbf{x} \in \Gamma_D \\ B_2(0, y) = B_2(1, y), & \mathbf{x} \in \Gamma_P \end{cases}$$

$\mathbf{g}(\mathbf{x}) = [0, -2]^T$ to produce analytic solution

$$\mathbf{B} = [0, x(x-1)]^T.$$

$\nabla \times (\mathbf{u} \times \mathbf{B}) = \mathbf{0} \implies$ no off diagonal coupling because

$\nabla \times (\nabla \times \mathbf{B}) - \nabla (\nabla \cdot \mathbf{B}) = -\nabla^2 \mathbf{B}$ so block structure

$$A = \begin{bmatrix} A_{11} & 0 \\ 0 & A_{22} \end{bmatrix}$$

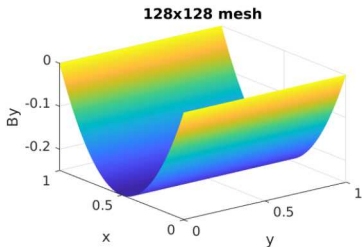


Figure: Solution plot

BoxMG for Uncoupled Problem (tol = 1e-7)				
# Elements	Total DOFs	V Cycles	Jacobi Sweeps	Levels
32^2	2178	8	(3,3)	5
64^2	8450	9	(3,3)	6
128^2	33282	9	(3,3)	7
256^2	132098	9	(3,3)	8
512^2	526338	9	(3,3)	9

$$\Omega = (0, 1)^2, \Gamma = \partial\Omega, \frac{\eta}{\mu_0} = 1.$$

Fixed velocity field $\mathbf{u} = (1, 1)^T$ in 2D produces

$$\nabla \times (\mathbf{u} \times \mathbf{B}) = \left[\frac{\partial(B_2 - B_1)}{\partial y}, -\frac{\partial(B_2 - B_1)}{\partial x} \right]^T.$$

Take

$$\mathbf{g}(\mathbf{x}) = \left[4\pi^2 \sin(2\pi y) - 2\pi \cos(2\pi x), 4\pi^2 \sin(2\pi x) - 2\pi \cos(2\pi y) \right]^T$$

chosen to manufacture analytic solution:

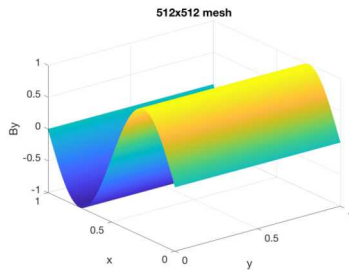
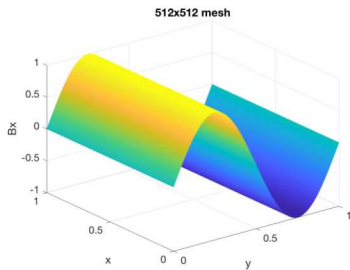
$$\mathbf{B}_{analytic} = \left[\sin(2\pi y), B_2 = \sin(2\pi x) \right]^T$$

$$\begin{cases} \nabla \times \left(\frac{\eta}{\mu_0} \nabla \times \mathbf{B} \right) - \frac{\eta}{\mu_0} \nabla (\nabla \cdot \mathbf{B}) - \nabla \times (\mathbf{u} \times \mathbf{B}) = \mathbf{g}(\mathbf{x}), & \mathbf{x} \in \Omega \\ \mathbf{B} = \mathbf{B}_{analytic}, & \mathbf{x} \in \partial\Omega \end{cases}$$

Block structure,

$$A = \begin{bmatrix} A_{11} & A_{12} \\ A_{21} & A_{22} \end{bmatrix}$$

2D Coupled Magnetic Field Convection-Diffusion Results



BoxMG for Uncoupled Problem (tol = 1e-7)

# Elements	Total DOFs	V Cycles	Jacobi Sweeps	Levels
32^2	2178	8	(3,3)	5
64^2	8450	8	(3,3)	6
128^2	33282	8	(3,3)	7
256^2	132098	8	(3,3)	8
512^2	526338	8	(3,3)	9

- ▶ Encouraging result for coupled system
- ▶ Like to perform experiment for full MHD system
- ▶ Also explore other PDE systems (elasticity)
- ▶ Current solver written in Matlab; developing system version in Trilinos/MueLu
- ▶ Possibly explore extension to Nédélec elements