

The combined effects of hydrogen and aging condition on the deformation and fracture behavior of a precipitation-hardened nickel-base superalloy

Zachary D. Harris^{1*}, Jishnu J. Bhattacharyya¹, Joseph A. Ronevich², Sean R. Agnew¹, and James T. Burns¹

¹Department of Materials Science and Engineering, University of Virginia, Charlottesville, VA, USA

²Sandia National Laboratories, Livermore, CA, USA

*Corresponding author: zdh8kt@virginia.edu

Abstract

The effect of hydrogen (H) on the deformation behavior of Monel K-500 in various isothermal heat treatment conditions (non-aged, under-aged, peak-aged, and over-aged) was assessed via uniaxial mechanical testing. H-charged and non-charged specimens were strained to failure to facilitate a comparison of ductility, fracture surface morphology, strength, and work hardening behavior. For all examined heat treatment conditions, H charging leads to a significant reduction in ductility, which is accompanied by a consistent change in fracture surface morphology from ductile microvoid coalescence to brittle intergranular fracture. While H charging led to a systematic enhancement in the yield strength of all heat treatments, the three age-hardened conditions exhibited a more than 2-fold increase relative to the non-aged heat treatment. This suggests that H modifies the dislocation-precipitate interactions, which also manifest themselves through changes in work hardening metrics related to the dislocation storage and recovery rates. In particular, the H-charged peak-aged specimen exhibited a significant increase in initial hardening (dislocation storage) rate relative to the H-charged under-aged specimen. Transmission electron microscopy of these samples confirmed the onset of widespread dislocation looping in the H-charged peak-aged sample, in addition to the planar slip bands characteristic of the non-charged condition. This result suggests that hydrogen induces the particle shearing-to-looping transition at smaller particle sizes. Possible mechanistic explanations for this observed behavior are presented.

Keywords hydrogen embrittlement, Monel K-500, precipitate-dislocation interactions, work hardening

1. Introduction

Hydrogen (H)-induced degradation of structural metals remains a critical failure mode in industries spanning the aerospace, marine, energy, and transportation sectors. However, despite over a century of effort [1], understanding of the microscale processes responsible for this life-limiting effect remains incomplete. Numerous candidate mechanisms have been advanced to describe H-induced degradation, as has been extensively reviewed elsewhere [2–4], with the most commonly cited being: adsorption-induced dislocation emission (AIDE) [5], H-enhanced localized plasticity (HELP) [6], H-enhanced decohesion

35 (HEDE) [7,8], and H-enhanced vacancy stabilization (HESIV) [9]. Recently, elucidating the role of H on
36 deformation behavior in the context of the aforementioned mechanisms has received extensive attention in
37 the scientific literature [10–16], with the impetus for much of this work being Beachem’s observation of
38 small void-like features on the cleavage and intergranular facets in 4300-type steel exposed to 0.6 M NaCl
39 [17]. The presence of these ductile features was subsequently taken as evidence that ‘microscopic plasticity’
40 governed H-assisted cracking, which led to the hypothesis that H ‘unlocks’ dislocations such that they move
41 and/or multiply at lower stresses [17].

42 This postulation has subsequently been supported by extensive *in-situ* observations of dislocation
43 characteristics under a H atmosphere in an environmental transmission electron microscope (TEM) [3].
44 Specifically, evidence that H caused enhanced dislocation velocities and reduced dislocation spacing in
45 pile-ups was documented across numerous alloy systems, including Ti, Ni, Fe, 310 stainless steel, 316
46 stainless steel, Ni₃Al, 7xxx-series aluminum, and IN903 [3]. Additionally, it was suggested that H restricted
47 dislocation cross-slip based on indications that H promoted slip planarity [18]. Mechanistically, this
48 restricted cross-slip was attributed to the combined effects of (1) H-induced reductions in the stacking fault
49 energy (H was estimated to reduce the stacking fault energy by ~20% [19]), and (2) an increased repulsive
50 force between partial dislocations. Considering the latter, Delafosse [20] theoretically demonstrated that H
51 reduces the attraction between edge components in partial dislocations, leading to increased repulsion,
52 larger separation distances between the partial dislocations, and impeded cross-slip. H has also been found
53 to enhance dislocation multiplication. For example, several studies have observed a reduction in the ‘pop-
54 in’ load of H-charged specimens relative to non-charged specimens during nano-indentation experiments,
55 which were interpreted as a reduction in the energy for homogeneous dislocation nucleation [21–23]. Chen
56 *et al.* reached similar conclusions regarding H effect on dislocation nucleation [24,25]. Specifically, X-ray
57 diffraction (XRD) peak broadening analysis of cold-rolled palladium (Pd) sheets indicated a 6-fold increase
58 in dislocation density for H-charged Pd (0.75 at. % H) relative to non-charged Pd. This observation was
59 qualitatively consistent with TEM analysis showing an increased dislocation density in the rolled H-charged
60 Pd [24], though such XRD results should be interpreted cautiously because the arrangement of the
61 dislocations can induce differences in the apparent dislocation density inferred from peak broadening
62 analysis [26]. While these studies provide insights into the fundamental effects that H has on dislocation
63 processes, attempts to extend this understanding into more complicated, engineering-relevant alloys remain
64 limited. Specifically, though precipitation-hardened alloys are widely employed in structural applications,
65 the effect of H on dislocation-precipitate interactions is largely unexplored [27,28]. Such interactions are
66 acutely pertinent given the well-known sensitivity of H-assisted cracking to bulk slip behavior [29], which
67 can be significantly altered by differences in precipitate morphology.

68 One common avenue for exploring H effects on deformation behavior is through the comparison
69 of dislocation patterning in the H-charged versus non-charged condition after deformation to a given strain
70 level. For example, Wang *et al.* performed detailed TEM characterization of the deformation structure in
71 pure Ni subjected to high-pressure torsion in the H-charged and non-charged condition [13]. Specimens
72 deformed to equivalent shear strains revealed a reduced average dislocation cell diameter in the H-charged
73 specimens compared to non-charged specimens, suggesting that H accelerates dislocation multiplication
74 (in agreement with the work of Chen *et al.* on Pd [24,25]). However, other studies have shown no
75 statistically significant refinement in the dislocation patterning of H-charged and non-charged Ni
76 [11,30,31]. These contradictory results underscore the challenge of assessing deformation behavior solely
77 through transmission electron microscopy, where the field of interrogation is effectively limited to a few
78 μm^2 . Speculatively, the variability in these results could be attributed to heterogeneities in plastic strain
79 and/or variations in grain orientation between compared samples [11,32,33]. As such, assessing dislocation
80 characteristics via a coupled macroscale/microscale approach may yield more consistent insights. One
81 macroscale pathway for assessing variations in bulk deformation behavior is through differences in
82 measured work hardening parameters [34–41]. For example, Cheng *et al.* utilized variations in work
83 hardening data to understand the effect of heat treatment on deformation processes in AA6111 [42].
84 However, the use of such a framework to understand the effects of H on deformation is very limited [43],
85 especially for precipitation-hardened alloys exposed to H.

86 The objective of this study is to systematically evaluate the influence of H on the strength, strain
87 hardening, and fracture behavior of Monel K-500 as a function of aging condition. Uniaxial tensile testing
88 was conducted on four isothermal heat treatments (corresponding to the non-aged, under-aged, peak-aged,
89 and over-aged conditions) in the non-charged and H-charged conditions. The results of this testing were
90 then examined in the context of work hardening metrics, coupled with targeted transmission electron
91 microscopy, to provide mechanistic insights into H-induced modifications in dislocation-precipitate
92 interactions.

93 2. Experimental Methods

94 2.1. Materials

95 Monel K-500 is a Ni-Cu superalloy nominally composed of a face-centered cubic (fcc) Ni-Cu solid
96 solution (γ) matrix and a homogenous distribution of highly coherent (<0.1 pct misfit strain), intermetallic
97 γ' ($\text{Ni}_3(\text{Al,Ti})$) precipitates [44–47]. The γ' precipitates have an ordered L1_2 structure, and due to the low
98 misfit and interfacial energy anisotropy, γ and γ' phases form as spherical particles [44]. Two types of
99 carbides have been identified in Monel K-500: (1) a heterogeneous distribution of MC-type carbides
100 (generally TiC) are commonly found in the γ matrix and (2) isolated M_{23}C_6 -type carbides (where M: Cr,

101 Mn, Fe, or Ni) have been observed on grain boundaries [44]. Previous reports suggest that the MC-type
 102 carbide distribution does not appreciably change under typical aging conditions for Monel K-500 [48] and
 103 that their contribution to the strength of Monel K-500 is minimal [44]. Similarly, the $M_{23}C_6$ -type carbides
 104 are also not expected to significantly impact the mechanical properties of Monel K-500 [44].

105 All experiments were conducted on a single heat of Monel K-500 procured in the form of a 6-m
 106 long, 16-mm diameter bar heat-treated to the solution-annealed condition. The as-received grain structure
 107 was nominally equiaxed, with a supplier-reported average grain size of $\sim 10 \mu\text{m}$. The bulk composition was
 108 evaluated using inductively coupled plasma optical emission spectroscopy (ICP-OES), while glow
 109 discharge mass spectroscopy (GDMS) was utilized for trace elemental analysis; the measured composition
 110 results are reported in Table 1.

111 *Table 1 – Composition of the tested Monel K-500 heat*

Ni	Cu	Al	Ti	Mn	Fe	Si	Zr	P	S	B
<i>Weight percent (%)</i>						<i>Weight part per million (wppm)</i>				
62.3	31.8	2.96	0.59	0.7	0.78	850	310	69	31	1.2

112
 113 Specimens for mechanical testing were heat-treated as follows to one of four targeted aging
 114 conditions, based on a previously reported age-hardening curves (yield strength vs. isothermal aging time)
 115 [49]. Briefly, 38-mm long, 12.7-mm diameter cylindrical blanks were cut from the barstock and then
 116 solution-treated in a tube furnace for 1 hour at 1223 K, followed by an immediate water quench; this heat
 117 treatment represents the non-aged (NA) condition. Specimens targeted for the under-aged (UA), peak-aged
 118 (PA), or over-aged (OA) conditions were then age hardened at 923 K for times of 0.5, 5, or 50 hours,
 119 followed by an immediate water quench, respectively. After heat treating, each cylindrical blank was
 120 machined into two flat, rectangular tensile specimens with a nominal gage length, width, and thickness of
 121 10.2 mm, 7.1 mm, and 3.2 mm, respectively.

122 **2.2. Hydrogen Charging**

123 Atomic H was introduced into a subset of the tensile specimens *via* gaseous charging; a detailed
 124 explanation of the charging procedure is reported elsewhere [50,51]. Briefly, specimens were placed inside
 125 an autoclave mounted within a furnace, which was heated to a charging temperature of 573 K and then
 126 filled with 99.9999% pure H_2 gas until the target charging pressure of 96.5 MPa was reached. The system
 127 was held at these conditions for approximately 144 hours to ensure the development of a uniform H
 128 concentration throughout the sample thickness [52]. Upon removal from the furnace, specimens were

129 immediately stored in a cryogenic freezer maintained at 223 K and then later immersed under liquid
130 nitrogen (LN₂) to minimize H egress prior to mechanical testing.

131 The lattice H concentration (C_L) can be predicted based on Sievert's Law ($C_L = Sf^{1/2}$; S is the H
132 solubility and f is the H fugacity), where f is determined by the pressure and temperature employed during
133 gaseous charging (*i.e.*, the relationship is material agnostic) [53]. This theoretical framework has been
134 successfully employed for H charging of pure Ni [10,11,50], which offers a useful comparison to the
135 expected C_L in Monel K-500. Using a Ni-30Cu solid solution as a proxy for the matrix phase of Monel K-
136 500, the expected difference in C_L between Ni and Ni-30Cu is expected to be proportional to the change in
137 H solubility. However, reports regarding the solubility of H in a Ni-30Cu solid solution relative to pure Ni
138 are inconsistent. For example, the work of Jones and Pehlke indicates that the H solubility of pure Ni is
139 nominally identical to that of Ni-Cu alloys for <45% Cu by weight [54]; a similar observation was also
140 reported by Ko and McLellan [55]. Conversely, Hagi observed a systematic increase in the H solubility of
141 Ni-30Cu relative to pure Ni, with an approximately 2-fold increase in H solubility measured at the charging
142 temperature employed in this study (573 K) [56]. Moreover, estimates of the H solubility based on measured
143 H permeability and diffusivity values for Monel K-500 suggest a >2-fold increase in solubility relative to
144 pure Ni [57–59]. Based on these historical reports, C_L for Monel K-500 under the employed gaseous
145 charging conditions is estimated to lie between ~70 and ~160 parts per million (ppm), assuming an upper-
146 bound value of 2.25-fold of the solubility of H in pure Ni [50].

147 To determine the actual H content of the charged specimens, the H concentrations of four 6 x 5 x
148 2.75 mm pieces from each aging condition were evaluated using an inert gas fusion analyzer at a
149 commercial laboratory (Evans Analytical Group). From these four replicates, average H concentrations of
150 180, 210, 220, and 220 ppm were measured in the NA, UA, PA, and OA specimens, respectively, which
151 correspond to 1.08 (NA), 1.26 (UA), and 1.32 (PA and OA) atomic (at.) %. Recalling that Sievert's Law
152 only estimates the lattice H concentration [50,53], and therefore, does not account for H trapping at grain
153 boundaries and other microstructural features which may be present (*e.g.* TiC particles), the deviation in C_L
154 for the NA condition (which would most closely approximate a Ni-30Cu solid solution) from the theoretical
155 results is not unexpected [11,50]. Moreover, this increase in C_L over theoretical expectations is also
156 consistent with previous measurements of the diffusible hydrogen concentration ($C_{H,Diff}$), which showed
157 that the tested material heat exhibited a systematically increased $C_{H,Diff}$ relative to five other evaluated heats
158 of Monel K-500 [59–61]. Considering the increased C_L for the aged alloys, this additional H content is
159 likely attributable to H trapping at the γ/γ' interface [59,61] (or within the γ' precipitate as suggested from
160 simulations by Baskes *et al.* [62]). The variation in C_L between the aged alloys can then be explained by
161 differences in γ' volume fraction. Specifically, Dey *et al.* demonstrated that the volume fraction of γ' will

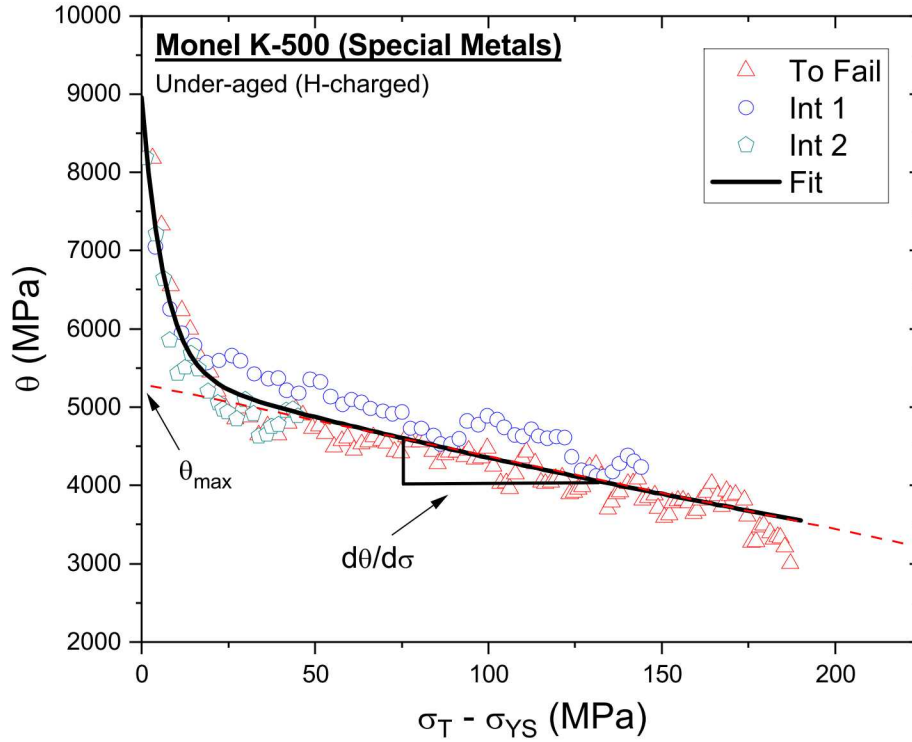
162 saturate after approximately 2 hours of aging at 923 K. As such, given that the dominant trap site in aged
163 Monel K-500 is the γ' precipitates [59,61], the saturation of the γ' volume fraction would explain the
164 identical H contents measured in the OA and PA heat treatments.

165 **2.3. Mechanical Testing**

166 Uniaxial tensile tests, with the loading direction oriented along the longitudinal axis of the barstock,
167 were performed on a servo-hydraulic load frame operated under actuator displacement control at a constant
168 rate of 0.0041 mm/s, yielding an initial strain rate of $\sim 4 \times 10^{-4} \text{ s}^{-1}$. Each specimen was tested until final
169 fracture, with the specimen elongation actively monitored by an attached extensometer. Prior to straining,
170 H-charged specimens, which had been continually stored under cryogenic temperatures (either in a
171 cryogenic freezer at 223 K or immersed in LN₂), were held at room temperature (RT, $\sim 298 \text{ K}$) for precisely
172 one hour, to allow H to redistribute and reach the equilibrium grain boundary H concentration for this
173 temperature [63]. In addition to the ductility obtained from the employed extensometer, the logarithmic true
174 plastic strain (ϵ_T) was determined post-test *via* the measured reduction in area (RA); the use of the true
175 plastic strain allows for direct comparison between specimens, including those which exhibited necking
176 [64]. Final area measurements were taken using calibrated optical micrographs of the fracture surface¹,
177 while initial area measurements were made using digital calipers.

178 The experimental work hardening rate (θ) was determined as follows. The true stress-true strain
179 data was binned and averaged to obtain approximately 200 data points, then these binned true stress-true
180 strain data were numerically differentiated using a 7-point ($n = 3$) polynomial method adapted from ASTM
181 Standard E647 Appendix XI [65]. The calculated θ vs. $(\sigma_T - \sigma_{YS})$ data, where σ_T is the true stress and σ_{YS}
182 is the 0.2% offset yield strength, was then fit to a sum of two exponential functions: $\theta = ae^{bx} + ce^{dx}$,
183 where $x = (\sigma_T - \sigma_{YS})$. A linear extrapolation was then performed using the fitted θ vs. $(\sigma_T - \sigma_{YS})$ values
184 over the interval of $\sim 100 \text{ MPa} < (\sigma_T - \sigma_{YS}) < 300 \text{ MPa}$ to obtain pertinent work hardening metrics (described
185 below), with the lower bound of 100 MPa selected to avoid any influence of the elastoplastic transition
186 region. In addition to the specimens tested to failure, 1-3 specimens from each heat treatment and H-
187 charging condition were deformed to specific targeted strains, where the test was interrupted for a future
188 study. The stress-strain data from these interrupted specimens were included in the analysis to (1) confirm
189 reproducibility of results and (2) improve the rigor of the fit, as shown for the UA/H condition in Figure 1.

¹ Note, due to a combination of necking and slant-type fracture, the RA values determined for the non-charged specimens should be considered an approximation. However, the flat fracture and limited necking observed in the H-charged specimens enabled an accurate measurement of the RA.



190

191 *Figure 1 –Plot of the work hardening rate (θ) versus flow stress increase ($\sigma_T - \sigma_{YS}$) illustrating the type of*
 192 *analysis performed for all of the hydrogen (H) charging and aging conditions examined. Note the curve fit*
 193 *of the experimental data as well as the simplification to linearity exhibited beyond the elastoplastic*
 194 *transition, which was used to graphically assess θ_{max} and $d\theta/d\sigma$.*

195 If the strain hardening exhibits linearity on a plot of the θ vs. $(\sigma_T - \sigma_{YS})$, the Voce approximation
 196 for Stage III hardening is appropriate, and two parameters may be used to quantitatively compare the work
 197 hardening behavior as a function of H and aging condition: the y-intercept (θ_{max}) and the slope ($d\theta/d\sigma$).
 198 These metrics are indicative of the rates of dislocation accumulation and recovery, respectively, which can
 199 be shown using the Kocks-Mecking model for the evolution of the dislocation density (ρ) with plastic strain
 200 [34]:

$$\frac{\partial \rho}{\partial \varepsilon} = \frac{\partial \rho^+}{\partial \varepsilon} - \frac{\partial \rho^-}{\partial \varepsilon} \approx k_1 \sqrt{\rho} - k_2 \rho \quad (1)$$

201 Where $\frac{\partial \rho^+}{\partial \varepsilon}$ and $\frac{\partial \rho^-}{\partial \varepsilon}$ describe the rates of dislocation accumulation and recovery, which are then dependent
 202 on the constants k_1 and k_2 , respectively. Note that numerous modifications to this basic framework have
 203 been made in the literature to account for the influence of various microstructural features (*e.g.* precipitates,
 204 solutes, grain size, free surfaces, etc. [35,37,42]) or the contribution of different populations of dislocation
 205 density (*i.e.* forest versus mobile [38,40]). A systematic description of several possible modifications is

206 presented by Keller and Hug [37], but in general, these modifications are completed through the
207 introduction of additional terms to describe either dislocation storage or recovery. Once the appropriate
208 terms are accounted for, Eqn. 1 can be considered a reasonable description for linking observed macroscale
209 metrics (θ_{\max} and $d\theta/d\sigma$) to the rates of dislocation storage and recovery. For example, as shown by Cheng
210 *et al.* for the effect of aging condition on the work hardening behavior of AA6111 [42], θ_{\max} and $d\theta/d\sigma$ are
211 directly proportional to k_1 and k_2 , respectively, though the relationship between k_1 and θ_{\max} is less
212 straightforward compared to k_2 and $d\theta/d\sigma$. Considering the current study, Voce strain hardening is
213 observed as shown by the example presented in Figure 1. However, significant scatter occurs from test-to-
214 test, thereby indicating a need to evaluate the statistical significance of the θ_{\max} and $-d\theta/d\sigma$ parameters.
215 Towards this end, the 95% confidence bands were calculated using Matlab and are reported in subsequent
216 sections [66].

217 **2.4. Characterization**

218 The fracture surface morphology for each sample was examined using a Quanta 650 FEG scanning
219 electron microscope (SEM) operated in secondary electron imaging mode with an accelerating voltage
220 between 5-10 keV. Foil specimens for transmission electron microscopy were taken from the gage section
221 of the peak-aged H-charged (PA/H) tensile specimen, with the foil plane parallel to the loading direction.
222 Specimens were prepared as followed: the gage section from the specimen was progressively thinned using
223 SiC papers, finishing at 1200 grit, to a final thickness of 100-110 μm . Care was taken to minimize polishing-
224 induced damage from lower grit sizes by removing at least 3x the maximum particle size of the preceding
225 SiC grinding paper before advancing to next grit size. For example, since the average particle size for the
226 400-grit SiC paper is $\sim 22 \mu\text{m}$, a minimum of 70 μm was removed using the 600-grit paper before advancing
227 to the 800-grit paper. After completion of this grinding procedure, 3-mm diameter discs were punched using
228 a Gatan disc puncher, further thinned to 80-90 μm using a Gatan disc grinder using 1200 grit SiC paper,
229 and then electrochemically polished in a Fischione 110 twin jet polisher using a solution of 60% H_2SO_4 and
230 40% H_2O (by volume) at ambient temperature (296 to 300 K) with an applied voltage of 5 V. Transmission
231 electron microscopy (TEM) of the obtained specimens was then conducted using an FEI Titan equipped
232 with a double-tilt specimen holder and operated at 300 keV.

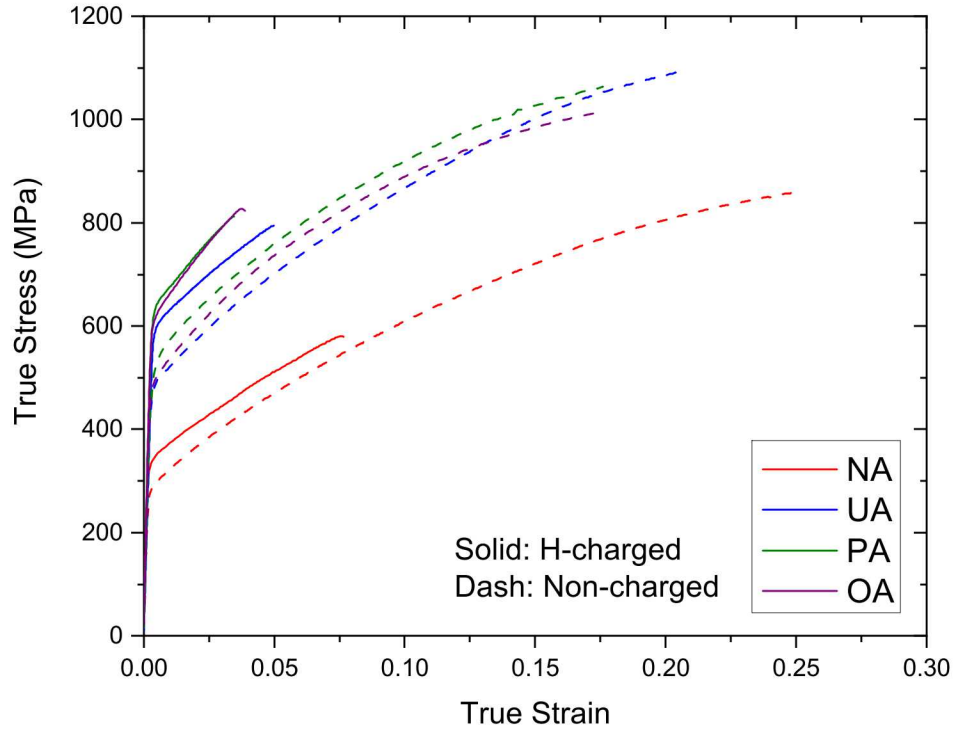
233 **3. Results**

234 **3.1. Effect of hydrogen on stress-strain behavior and fracture morphology**

235 The true stress-true strain curves for the H-charged and non-charged specimens from each aging
236 condition are shown in Figure 2. Note that the non-charged curves are truncated at the onset of necking,
237 which was assessed using Considere's criterion [67]. Mechanical properties obtained from these tests to

238 failure are listed in Table 2. Comparison of the flow curves for the H-charged versus non-charged conditions
239 demonstrates that hydrogen degrades the mechanical properties of Monel K-500, regardless of heat
240 treatment condition. First, consistent with previous literature on hydrogen effects in Monel K-500 [68–70],
241 a significant decrease in ductility was observed across all four tested aging conditions in the presence of
242 hydrogen. Second, the fracture stress (approximated by the ultimate tensile strength derived from the
243 engineering stress-strain curve for the non-charged alloys; not shown) was also reduced in the presence of
244 hydrogen by approximately 100 MPa for the NA, UA, and PA conditions, while the OA condition decreased
245 by 60 MPa. Third, the yield strength was found to increase with H-charging across all four aging conditions,
246 with increases on the order of 100 MPa observed in the UA, PA, and OA specimens, while the NA exhibited
247 an increase of 49 MPa.

248 These differences in ductility between the H-charged and non-charged conditions are consistent
249 with the observed fracture surface morphologies. As shown in Figure 3, each heat treatment displayed
250 evidence of ductile failure *via* microvoid coalescence in the absence of H, in agreement with the results of
251 prior fracture mechanics and slow-strain rate testing of Monel K-500 in laboratory air [60,61,71,72]. The
252 void morphology across the four conditions was qualitatively similar. Considering the fractography of the
253 H-charged samples, shown in Figure 4, each heat treatment exhibited widespread intergranular fracture. As
254 observed in prior studies for hydrogen-charged nickel alloys [10,11,73], slip traces were observed on the
255 intergranular facets for all heat treatment conditions (Figure 5), correlating with the fact that some plastic
256 deformation occurred prior to failure (Figure 2).

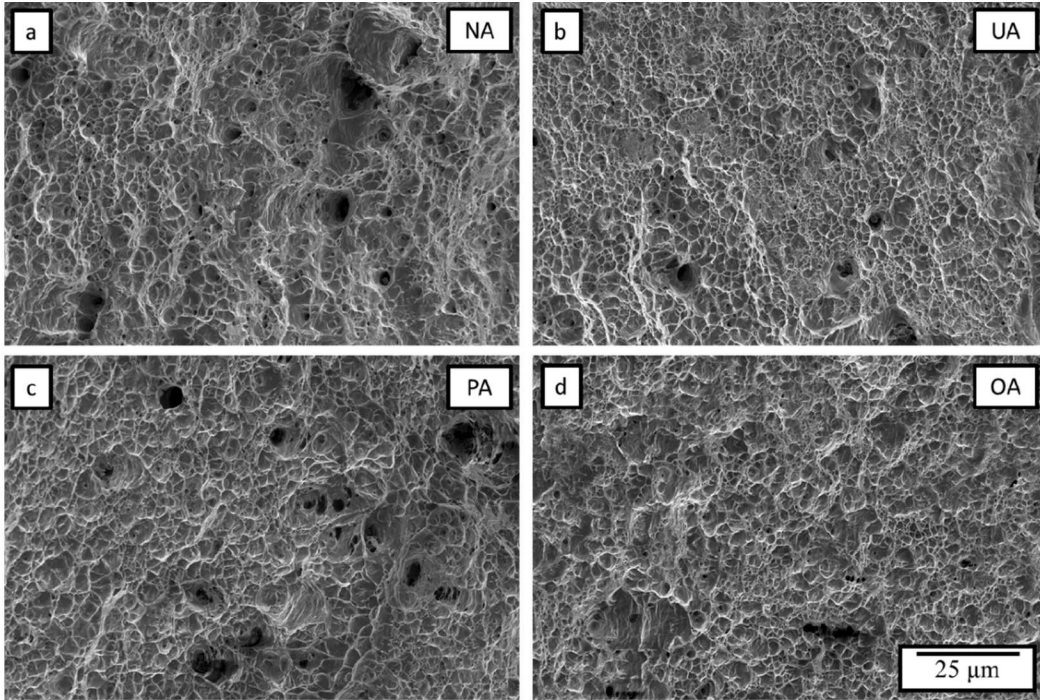


257

258 *Figure 2 – True stress-true strain curves for the hydrogen (H)-charged and non-charged specimens*
 259 *strained to failure for each heat treatment.*

260 *Table 2 – Mechanical properties as a function of heat treatment and hydrogen condition*

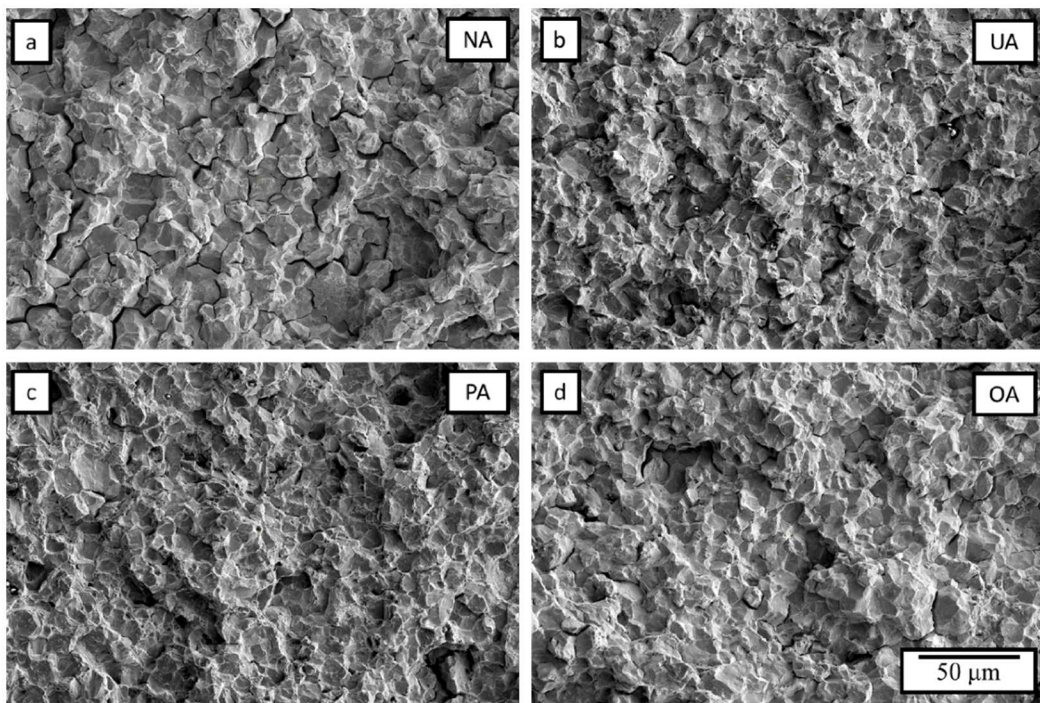
Testing Condition	Yield Strength (MPa)	Fracture Stress (MPa)	Reduction of Area (%)	True Fracture Strain
NA/H	345	538	8.7	0.091
NA/No H	294	670	50.0	0.693
UA/H	605	757	5.6	0.057
UA/No H	490	890	42.3	0.550
PA/H	641	786	5.5	0.056
PA/No H	526	894	38.1	0.479
OA/H	623	798	4.6	0.047
OA/No H	497	853	39.2	0.497



261

262
263

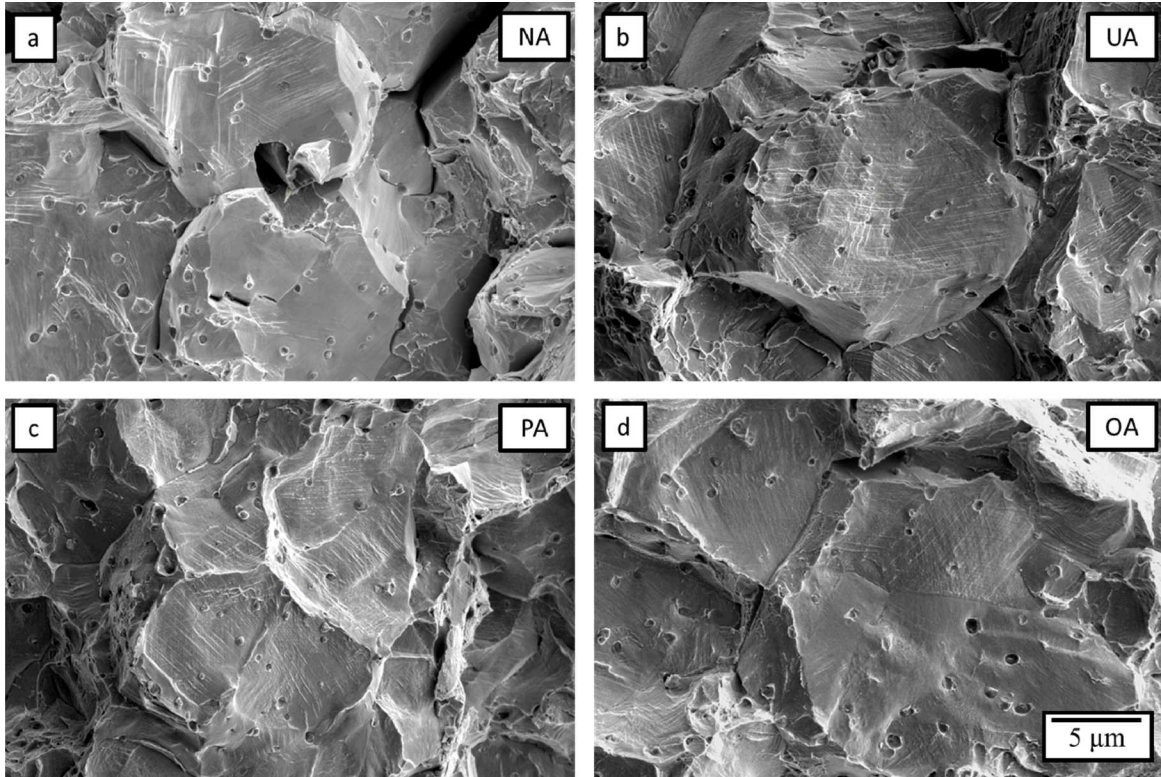
Figure 3 – Fractographs of the (a) non-aged, (b) under-aged, (c) peak-aged, and (d) over-aged heat treatments tested to failure in the non-charged condition.



264

265
266

Figure 4 – Fractographs of the (a) non-aged, (b) under-aged, (c) peak-aged, and (d) over-aged heat treatments tested to failure in the hydrogen-charged condition.



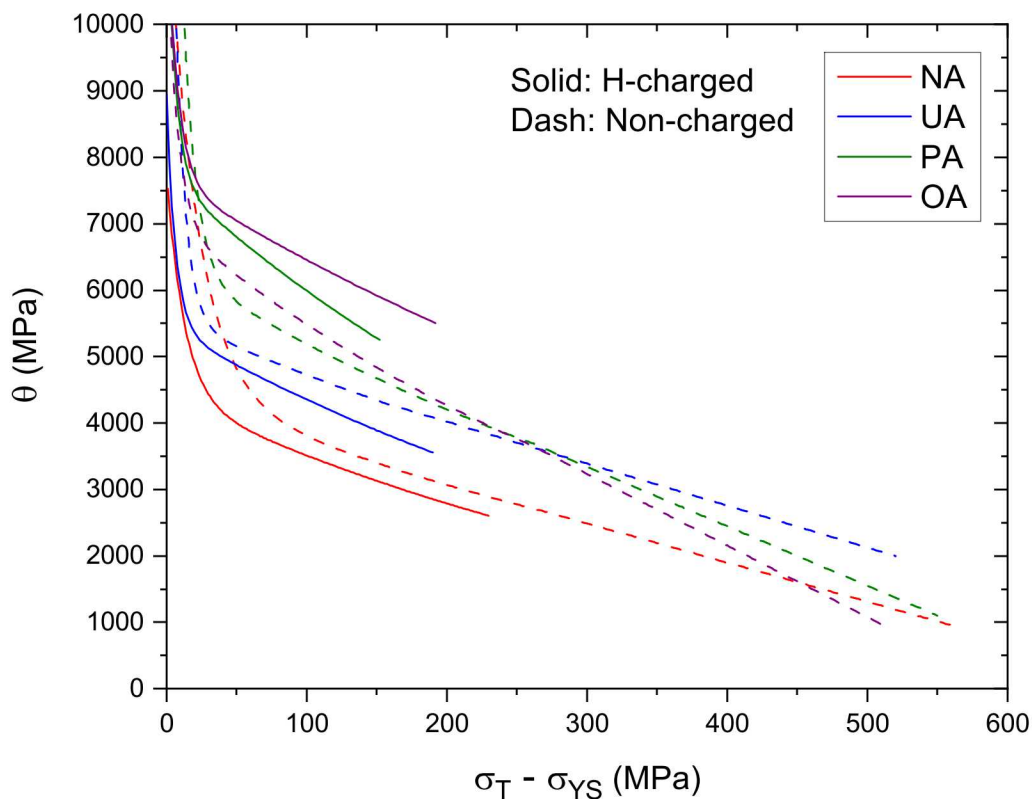
267

268 *Figure 5 – High-magnification fractographs of the slip traces observed on intergranular facets in the (a)*
 269 *non-aged, (b) under-aged, (c) peak-aged, and (d) over-aged heat treatments tested to failure in the*
 270 *hydrogen-charged condition.*

271 **3.2. Effect of hydrogen on work hardening behavior**

272 The work hardening rate (θ) versus flow stress increase ($\sigma_T - \sigma_{YS}$) relationships for the non-charged
 273 and H-charged condition in each heat treatment are presented in Figure 6. Each specimen, regardless of H
 274 and aging condition, displays an initial rapid decrease in work hardening rate at small ($\sigma_T - \sigma_{YS}$), which is
 275 related to the significant changes in slope within the elastoplastic transition of the true stress-true strain
 276 curve [42]. Beyond this transition region, each of the specimens (regardless of H content) exhibits an
 277 essentially linear decrease in the work hardening rate with increasing ($\sigma_T - \sigma_{YS}$). Clear variations in the
 278 character of this linear region arise among the tested specimens, both as a function of heat treatment and as
 279 a function of H, which can be described by θ_{max} and $-d\theta/d\sigma$. These metrics were determined from the work
 280 hardening data for each H/heat treatment combination as defined in Figure 1 and then plotted in Figure 7,
 281 along with the error bars derived from the calculated confidence bands for each testing condition. Figure
 282 7a reveals a distinct effect of H on the evolution in dislocation storage rate (indicated to the first-order by
 283 the determined θ_{max} values [42]) with heat treatment. While the NA and UA heat treatments exhibit similar
 284 θ_{max} values in the charged and non-charged condition (note that differences in $\theta_{max} < 300\text{-}400$ MPa are not
 285 generally considered significant [74]), a clear increase in $\theta_{max} \sim 1500$ MPa was observed for the H-charged

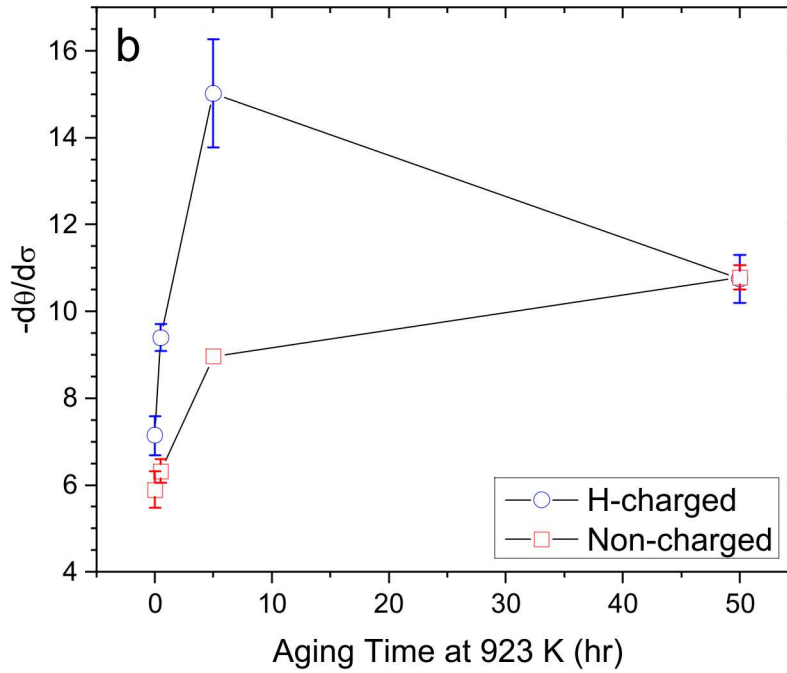
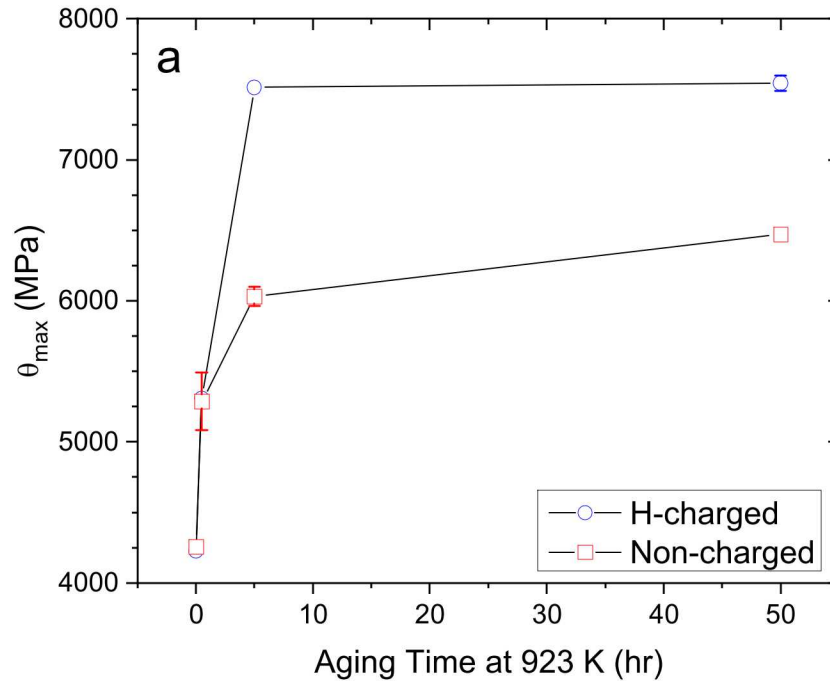
286 PA and OA conditions relative to the respective non-charged results. Figure 7b shows that H also influences
 287 the dislocation recovery rates, as demonstrated by changes in $-d\theta/d\sigma$. Considering the non-charged
 288 condition, similar values of $-d\theta/d\sigma$ (≈ 6) were observed for NA/No H and UA/No H, followed by an increase
 289 to 8.1 and 10.6 for the PA/No H and OA/No H heat treatments, respectively. Conversely, upon adding H, $-d\theta/d\sigma$
 290 was found to be increased for all heat treatments relative to the non-charged condition, with the
 291 exception of OA, which had a similar $-d\theta/d\sigma \approx 10.7$ for both conditions. In particular, $-d\theta/d\sigma$ were elevated
 292 for the H-charged UA and PA samples, compared to their respective non-charged condition results.



293

294 *Figure 6 – Work hardening rate versus flow stress increase for the hydrogen (H)-charged and non-*
 295 *charged specimens strained to failure for each heat treatment.*

296



297

298 *Figure 7 – Plot of (a) θ_{max} and (b) $-d\theta/d\sigma$ as a function of aging time and hydrogen (H). Parameters were*
 299 *obtained using the graphical method shown in Figure 1, with error bars defined based upon a linear fit to*
 300 *the bounds of the confidence bands for each heat treatment/hydrogen content combination. Note that the*
 301 *error bars were removed for plot clarity if they fell within the width of the respective data symbol.*

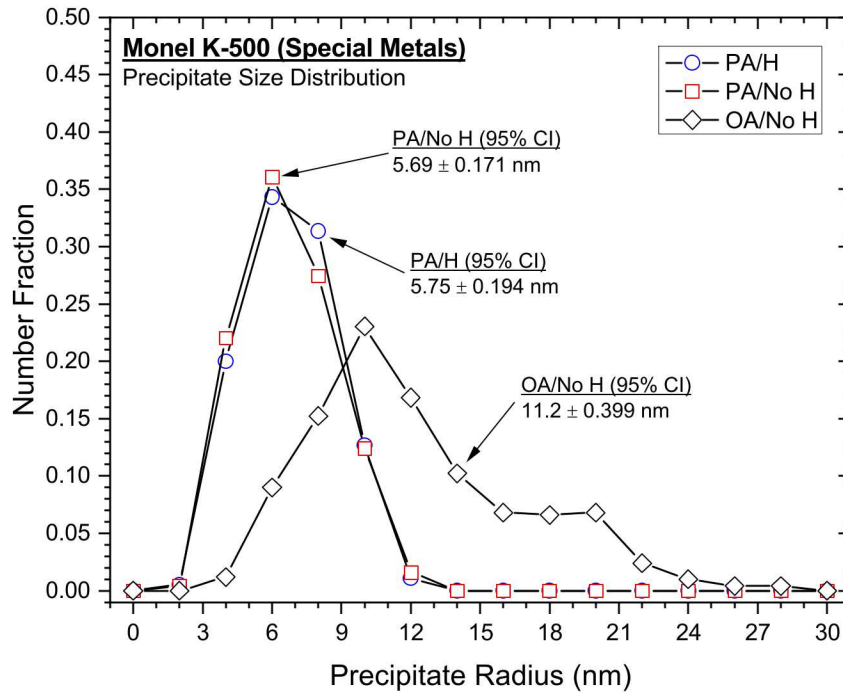
302

3.3. Verification of similar precipitate morphologies between the non-charged and H-charged conditions

A critical assumption of the current work is that the γ' morphology (*e.g.* size) and distribution is not appreciably modified during the thermal charging procedure (at 573 K) utilized to introduce hydrogen into the tensile specimens. As discussed by Cheng et al. [42], variations in precipitate size can induce changes in the work hardening behavior, which would then complicate direct assessment of the role of H for a given heat treatment condition. The extrapolation of γ' coarsening kinetics data previously obtained for Monel K-500 [44] suggests that negligible growth (<0.01 nm) should be expected after 150 hours at 573 K. However, while limited growth is expected in the age-hardened alloys based on these calculations, the extended time at elevated temperature could be sufficient to enable the nucleation of γ' in the NA heat treatment. To assess this possibility, a 12.7-mm diameter, 2-mm thick disc specimen in the NA condition was aged at 573K for 335 hours, then immediately water quenched. Ten Rockwell B (HRB) hardness measurements were taken from polished samples in the (1) baseline NA condition, and (2) the 573K annealed NA condition, which yielded an average hardness of 83.3 ± 2.7 and 82.8 ± 2.1 HRB, respectively. These similarities in measured hardness suggest that γ' nucleation would not occur due to the extended time at 573 K, which is significant given that the NA condition will have the highest driving force for change in the precipitate microstructure.

In addition to the effect of prolonged exposure to elevated temperatures on the precipitate morphology, it has previously been shown that H can stabilize the vacancy concentration above that predicted by thermodynamics [10,15], which may cause an acceleration in coarsening kinetics or assist in the nucleation of precipitates in the NA condition. Regarding the latter, an increase of ~ 50 MPa in yield strength (see Table 2) is observed for the NA sample, after H-charging, which could be attributed to the nucleation of γ' . However, a similar increase in yield strength was observed for 304 stainless steel charged with ~ 1 at. % H [75], suggesting that the observed increase may be due to solute strengthening effects and not the nucleation of γ' precipitates. To assess the possibility of H-enhanced particle coarsening of the aged heat treatments during the H charging process at 573 K, an assessment of the precipitate size distribution was completed using transmission electron microscopy on foil specimens prepared from the PA/H tensile specimen. To ensure a representative evaluation of the γ' size, a total of 375 precipitates were examined from multiple locations on two separate foil specimens prepared from the PA/H tensile specimen. This measured γ' size distribution was then compared to a distribution measured for the PA/No H condition in prior work [49]; the distribution for OA/No H was also included for comparison. As shown by Figure 8, the distributions for the non-charged and hydrogen-charged PA specimens are nominally identical, with the difference in the average particle radius between the two conditions within the calculated 95% confidence

336 interval. Additionally, clear differences are observed between the two PA conditions and the OA/No H
 337 specimen, with the latter exhibiting a ~2-fold increase in the average particle size and a significantly wider
 338 distribution of observed particle sizes, as indicated by the ~2-fold increase in the 95% confidence interval.



339
 340 *Figure 8 – Comparison of the γ' size distribution for the non-charged over-aged condition (OA/No H)*
 341 *and both the non-charged (PA/No H) and hydrogen-charged (PA/H) peak-aged conditions, along with the*
 342 *average and 95% confidence interval of the precipitate size for each condition.*

343 4. Discussion

344 The presented flow curves (Figure 2) demonstrate that the ductility and fracture stress (Table 2) of
 345 Monel K-500 are decreased in the presence of H across all four tested aging conditions. Fractography
 346 corroborated this assessment based on a systematic transition in fracture morphology from ductile
 347 microvoid coalescence to intergranular failure for all H-charged specimens (Figures 3-5). Analysis of work
 348 hardening data (Figure 6) revealed a clear effect of H on the deformation behavior of Monel K-500, as
 349 quantified by θ_{max} and $-d\theta/d\sigma$ (Figure 7). In particular, systematic increases in θ_{max} and $-d\theta/d\sigma$ were
 350 observed for the PA/H relative to PA/No H. Lastly, examination of the precipitate size distribution for PA/H
 351 and PA/No H conditions indicates that the H charging procedure did not induce statistically significant
 352 differences in the precipitate morphology (Figure 8). Based on these observations, the following discussion
 353 will: (1) assess the role of H on the mechanical behavior of Monel K-500 as a function of aging condition
 354 and (2) comment on the mechanistic implications of the presented results as well as identify research
 355 questions for future study.

356 **4.1. Effect of hydrogen on the mechanical behavior of Monel K-500 as a function of aging**
357 **condition**

358 *4.1.1 Effect of hydrogen on yield strength*

359 Comparison of the *change* in yield strength due to H-charging indicates a systematic difference
360 between the NA condition and the three age-hardened heat treatments (Figure 2, Table 2). Specifically, a
361 difference in yield strength of 51 MPa was observed for NA, while increases of greater than 115 MPa were
362 observed for UA, PA, and OA. The precipitate size distributions and analysis presented in Section 3.3
363 demonstrates that it is unlikely that this increase is caused by modification of the precipitates during the
364 elevated temperature gaseous charging procedure, suggesting that H is responsible for the observed increase
365 in yield strength. It is useful then to consider how H may impact the operative strengthening mechanisms
366 for each aging condition. As shown in Eqn. 2, the yield strength can be represented as the sum of the
367 contributions from each potential strengthening mechanism: lattice friction (σ_0), cold work (σ_D), grain size
368 (σ_{GB}), precipitation hardening (σ_{PH}), and solid solution strengthening (σ_{SS}).

$$\sigma_{YS} = \sigma_0 + \sigma_D + \sigma_{GB} + \sigma_{PH} + \sigma_{SS} \quad (2)$$

369 For the NA condition, the effect of H on precipitation hardening and cold work can be reasonably
370 considered negligible given that (1) precipitates are not present in the NA heat treatment, and (2) the
371 specimens were received in the annealed condition from the supplier, re-solutionized at 1223 K and
372 quenched, and then effectively annealed again during the H charging process (which was completed at 573
373 K). Similarly, it is unlikely that an effect of H on the lattice resistance is responsible since σ_0 is a small
374 contributor to the strength of fcc alloys and any subtle effect of H on the lattice resistance can be subsumed
375 into the solid solution strengthening term (σ_{SS}).

376 The effectiveness of grain size strengthening could be modified by H-charging in two ways. First,
377 it is possible that the strength could be reduced by grain growth that could have occurred during the H
378 charging. However, prior work on the same heat of Monel K-500 demonstrated negligible grain growth
379 after 50 hours at 923 K [49], where the driving force for grain growth would be much higher than the H
380 charging temperature employed herein (573 K) [76,77]. Second, H could affect the difficulty of grain
381 boundary transmission by dislocations (*i.e.* increase the locking constant of the Hall-Petch relationship
382 [78]), thereby increasing the contribution of a given grain size to the yield strength. Such an effect has been
383 documented for carbon in steel, though the efficacy appears to depend on the solute, as nitrogen was found
384 to have minimal influence on grain boundary strengthening [79,80]. Considering H, reports are inconsistent
385 and predominantly conducted in Fe-based alloys. For example, Mine *et al.* found no change for metastable
386 austenitic stainless steels [81], while Bernstein reported no influence in Fe-0.15Ti, but did note an increase

387 for decarburized Fe [82]. Similarly, Caskey observed a systematic increase in the locking constant for H-
 388 charged 304L [83]. In the case of pure Ni, Wilcox and Smith found negligible difference in the grain size-
 389 dependent yield strength with H exposure [84], while the data of Lawrence *et al.* suggests a slight increase
 390 in the locking constant for H-charged Ni (albeit on a sample size of two different grain sizes) [10]. Based
 391 on these reports, H may modify the efficacy of grain size strengthening, but not at a sufficient level to fully
 392 account for the observed increase in yield strength for the NA condition. Such a conclusion is consistent
 393 with the findings of a recent literature review on the Hall-Petch effect [78], which attributed most of the
 394 strength contribution from interstitials in Nb to solid solution hardening as opposed to modifications in the
 395 efficacy of grain size strengthening. As such, H-induced solid solution strengthening is likely the dominant
 396 contributor to the 51 MPa increase observed in the NA condition upon H charging.

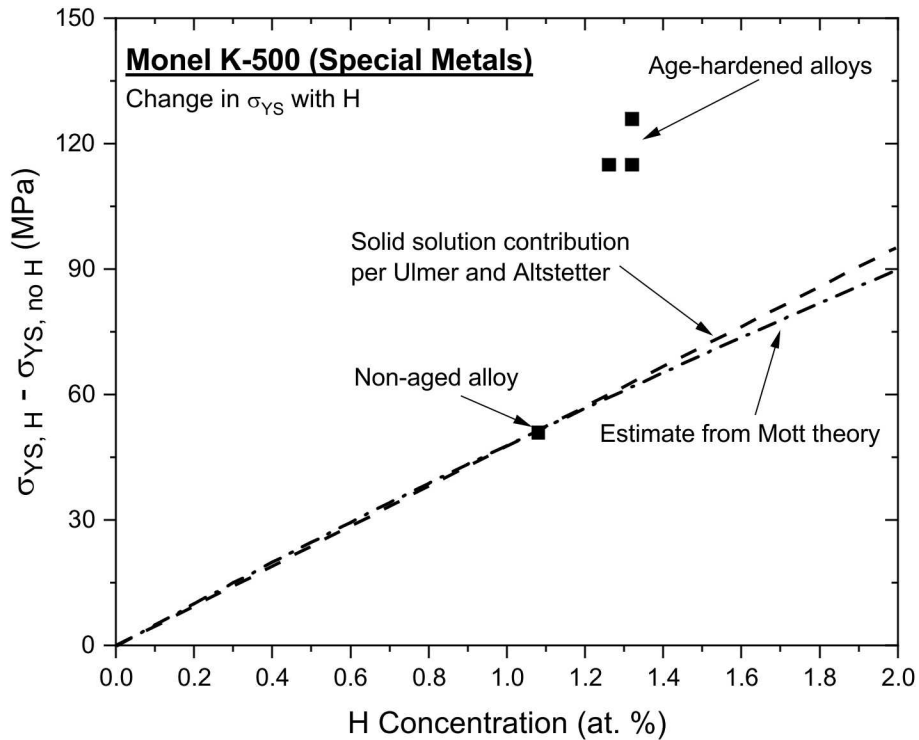
397 Interestingly, the age-hardened specimens all exhibited a >2-fold increase in yield strength
 398 compared to the NA condition (Table 2). As the arguments for the lack of H effects on the lattice friction,
 399 cold work, and grain size contributions to yield strength set forth for the NA condition are also pertinent to
 400 the aged conditions, such factors would not explain this additional increase. This suggests that the increase
 401 in strength for these age-hardened alloys is due to either (1) additional H-induced increases in the solid
 402 solution strengthening relative to the NA condition or (2) an effect of H on precipitate strengthening
 403 mechanisms. Regarding the former, though the aged samples contained roughly 20% more H than the NA
 404 sample, empirically-based estimates and theoretical calculations (shown in Figure 9) suggest that the
 405 increased solid solution strengthening associated with this added H content does not explain the observed
 406 increase in yield strength for these alloys as compared to the NA condition. For example, Ulmer and
 407 Altstetter conducted a systematic evaluation of the increase in yield strength of single phase, H-charged
 408 304 stainless steel, which revealed a linear dependence on the yield strength with increasing H
 409 concentration where the yield strength increased by 40-50 MPa/at. % of H up to 10 at. % [75]. Assuming a
 410 similar linear approximation based on the increase in yield strength for the NA condition (shown by the
 411 dashed line in Figure 9), an increase on the order of 60-65 MPa would be expected between the H-charged
 412 and non-charged conditions for the UA, PA, and OA heat treatments. Considering theoretical predictions,
 413 while more rigorous models for solute strengthening (both analytical [85–87] and computational [16,88–
 414 90]) exist for solid solution strengthening, the Mott model [86] can be employed as a first-order assessment,
 415 as demonstrated by Nakada *et al.* for the effect of nitrogen and carbon on the yield strength of Fe [91]:

$$\Delta\sigma_{YS} = \frac{1}{2} M\mu\epsilon^{\frac{4}{3}} f \left[\frac{f^{\frac{2}{3}}(\log f)^4}{0.5} \right]^{\frac{1}{3}} \quad (3)$$

416 Where M is the Taylor factor for an fcc material, μ is the shear modulus (taken to be 72 GPa based on a Ni-
 417 30Cu matrix [92]), f is the atom fraction of solute, and ϵ is the misfit strain defined by:

$$\epsilon = \frac{1}{a} \frac{da}{df} \quad (4)$$

418 Where a is the lattice parameter and $\frac{da}{df}$ represents the change in lattice parameter with solute concentration.
 419 Using the data of Baranowski *et al.* [93] for pure Ni, and assuming weak coupling between the interstitial
 420 H and substitutional Cu solute (*i.e.* Baranowski's relation for Ni-H holds for Ni-Cu-H), $\frac{da}{df}$ can be calculated
 421 as $\sim 0.3 \text{ \AA/unit concentration}$ and the lattice parameter can be approximated as that of uncharged Ni-30Cu
 422 based on the work of Sakamoto ($a = 3.55 \text{ \AA}$) [58]. As shown by the dotted-dashed line in Figure 9, the
 423 prediction of Mott's simple model is in reasonable agreement with the empirical estimate based on the
 424 results of Ulmer and Altstetter [75]. Critically, this analysis also suggests that the increased magnitude of
 425 strengthening observed for the H-charged aged conditions is not attributable to solid solution effects,
 426 thereby suggesting an H-induced modification of dislocation-precipitate interactions.



427
 428 *Figure 9 - Observed increase in yield strength versus hydrogen concentration. Note that the dashed line*
 429 *is adapted based on the work of Ulmer and Altstetter [75] and the dotted line is the prediction based on*
 430 *Equation 4.*

431

432 4.1.2. Effect of hydrogen on the initial hardening rate (θ_{max})

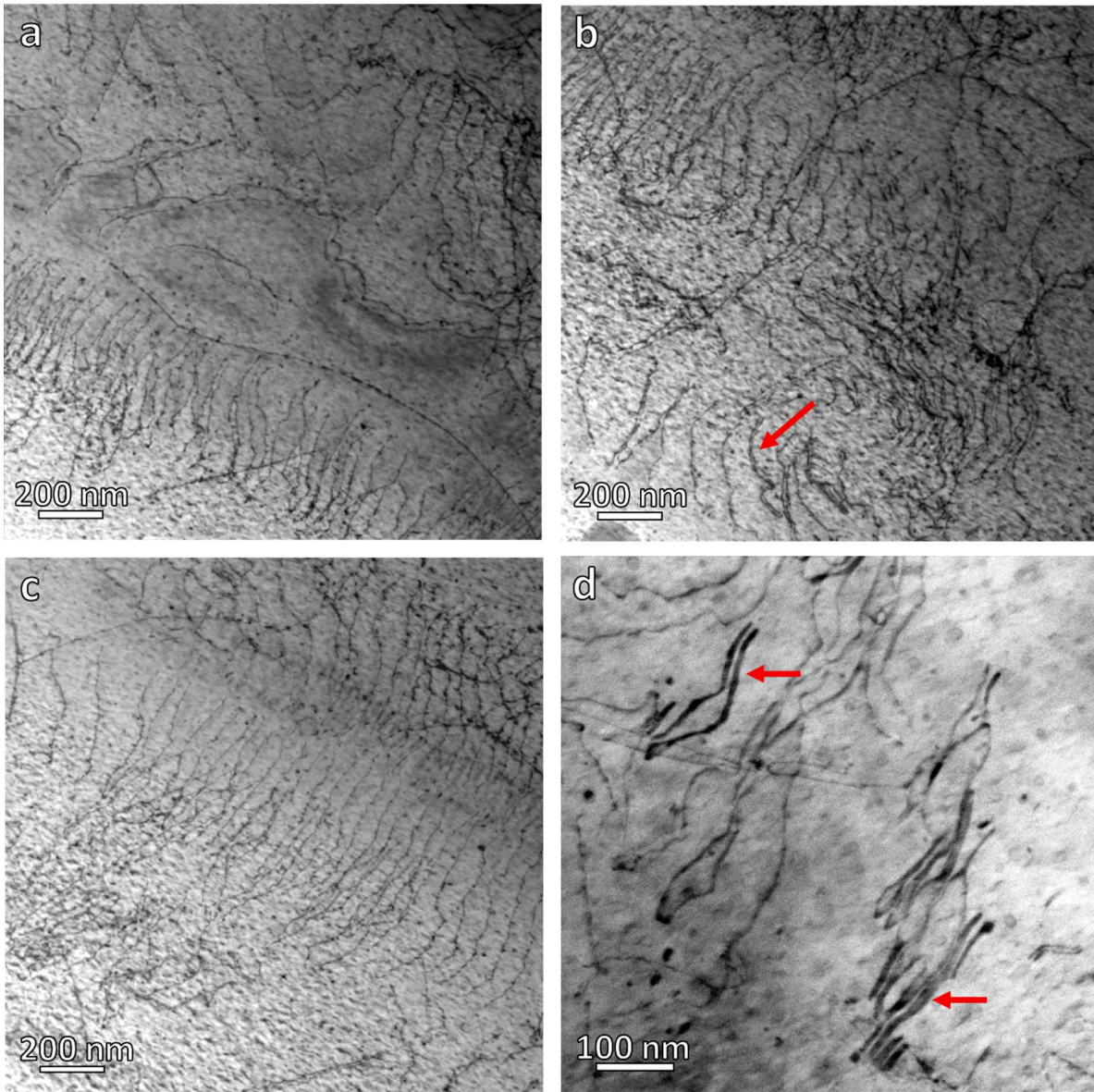
433 Previous studies regarding the sensitivity of θ_{max} to the presence of solute atoms suggest either a
434 limited influence of solute [34,42] or the sensitivity being dependent on alloy/solute combination [94].
435 Based on the results in Figure 7a, θ_{max} appears to be independent of H-charging in the NA and UA alloys,
436 while large increases in θ_{max} were observed for the PA and OA conditions. Historically, such increases in
437 the work hardening rate for different aging conditions in alloys hardened by secondary phases are generally
438 attributed to two factors: (1) image stresses from concentric dislocation loops that form around precipitates
439 (*i.e.* looping processes) [95,96] and/or (2) constraint effects when particles no longer co-deform with the
440 matrix (*i.e.* load partitioning) [42,95,97]. Particles which no-longer co-deform then contribute to the work
441 hardening rate *via* elastic constraint (due to the mismatch in accommodated plastic strain between the matrix
442 and precipitate), with the magnitude of this effect being proportional to βGf [95], where β is a scaling
443 factor on the order of 0.5 to account for precipitate shape *via* Eshelbian micromechanics, G is the shear
444 modulus of the particle, and f is the particle volume fraction. However, it is important to recognize that
445 these looping and load partitioning contributions are not mutually exclusive, as the onset of particle
446 bypassing inherently induces non-uniform load accommodation between the matrix and precipitate phases.
447 This notion of a coupled effect on the work hardening rate is supported by neutron diffraction experiments
448 on Ni-base superalloys during deformation [98–100]. For example, modeling by Francis *et al.* of the phase-
449 specific (γ and γ') deformation based on neutron diffraction measurements on RR1000 with varying sizes
450 of γ' revealed that the hardening rate of the γ matrix increased as the γ' size increased [98]. As such, though
451 Francis *et al.* demonstrate that load partitioning is the dominant contribution to work hardening in RR1000
452 [98], this need for an increased hardening rate in the matrix phase in order to capture the deformation
453 behavior of the aggregate material ($\gamma + \gamma'$) implies that both load partitioning and looping-associated back-
454 stresses (to an unknown degree) are operative in γ' -strengthened alloys.

455 It is useful then to consider both possible contributions to work hardening in order to understand
456 the observed variation in θ_{max} for the evaluated heat treatments and H contents. Given a shear modulus of
457 77 GPa for Ni₃Al [101] and a maximum γ' volume fraction of 6.5% [44], the largest work hardening
458 contribution from load partitioning for Monel K-500 can be estimated as ~2500 MPa, assuming the entire
459 volume fraction of γ' is no longer co-deforming [95,97]. Regarding the non-charged alloys, prior work on
460 Monel K-500 demonstrates that the transition from particle cutting to bypassing occurs at an aging time
461 between the current PA and OA heat treatments [46,49]. As such, looping-associated back-stresses are not
462 expected to contribute to θ_{max} until the OA condition, since shearable particles have minimal effect on θ_{max}
463 [35,42]. Given that several authors have reported load partitioning for even ‘fine’ γ' -containing Ni-base
464 superalloys [98–100], which exhibit particle shearing, it is reasonable to suggest that the observed ~1000

465 MPa increase in θ_{\max} from the NA/No H to UA/No H condition is likely due to load partitioning effects.
466 Similarly, the increase in θ_{\max} of ~ 700 MPa between the UA/No H and PA/No H is also assumed to be due
467 to load partitioning, with this additional increase arising from a larger volume fraction of particles that do
468 not co-deform. Lastly, considering the increase in θ_{\max} of ~ 500 MPa between the PA/No H and OA/No H,
469 the prior observation of dislocation loops in the OA/No H condition [49] suggests that both possible
470 contributions to work hardening may be active. Based on the work of Francis *et al.* [98] it is speculated that
471 this 500 MPa increase is principally due to load partitioning effects, however it is likely there is a tangible
472 contribution of looping-induced back-stresses. Regarding the H-charged alloys, the increase of ~ 1000 MPa
473 in θ_{\max} between the NA/H and UA/H conditions can likely be attributed to load partitioning effects, given
474 the similarities with the non-charged condition. However, θ_{\max} is then found to strongly increase by ~ 2250
475 MPa between the UA/H and PA/H conditions. This increase of ~ 3250 MPa over the NA/H condition is
476 larger than the maximum contribution predicted for load partitioning effects (~ 2500 MPa), which implies
477 that H-charging has induced an early transition to bypassing mechanisms (in order to generate the looping-
478 induced back-stresses). Interestingly, θ_{\max} was found to be similar for the PA/H and OA/H conditions;
479 speculatively, this result may indicate that a saturation of the looping contribution to work hardening has
480 occurred.

481 Based on the discussion above, the most tractable mechanistic explanation for the increased work
482 hardening rate (as quantified by θ_{\max}) between the PA/H and PA/No H conditions is that H induces the
483 transition from particle shearing to bypassing at smaller precipitate sizes. This hypothesis is supported by
484 (TEM) of the bulk slip morphology in a PA/No H specimen compressively deformed to 2% plastic strain
485 in prior work [49] and the PA/H specimen strained to failure, shown in Figures 10 and 11, respectively.
486 Specifically, the PA/No H specimen exhibits large planar arrays with evidence of dislocation coupling
487 (highlighted by the red arrows in Figures 10b and 10d), with limited evidence of dislocation bypassing.
488 Conversely, while the PA/H specimen contained planar bands (red arrows in Figure 11a) oriented along the
489 trace of a $\{111\}$ plane, particle bypassing was also observed (*e.g.* dislocation loops indicated by the red
490 arrows in Figure 11c-d). Alternate explanations for the observed dislocation loops are not compelling, but
491 include changes in particle size distribution between PA/No H and PA/H or the injection of dislocation
492 loops during the TEM analysis. Specifically, the particle size distributions plotted in Figure 8 demonstrate
493 that the average precipitate size between PA/No H and PA/H is sufficiently similar to suggest that
494 modifications in precipitate size are not responsible. Similarly, the possibility that the dislocation loops
495 were injected by the TEM is unlikely as dislocation loops in Ni alloys generally require long exposure times
496 and/or accelerating voltages on the order of 1 MeV [102]. As such, it is concluded that H acts to decrease
497 the precipitate size required for the onset of particle bypassing mechanisms. However, *why* H induces a

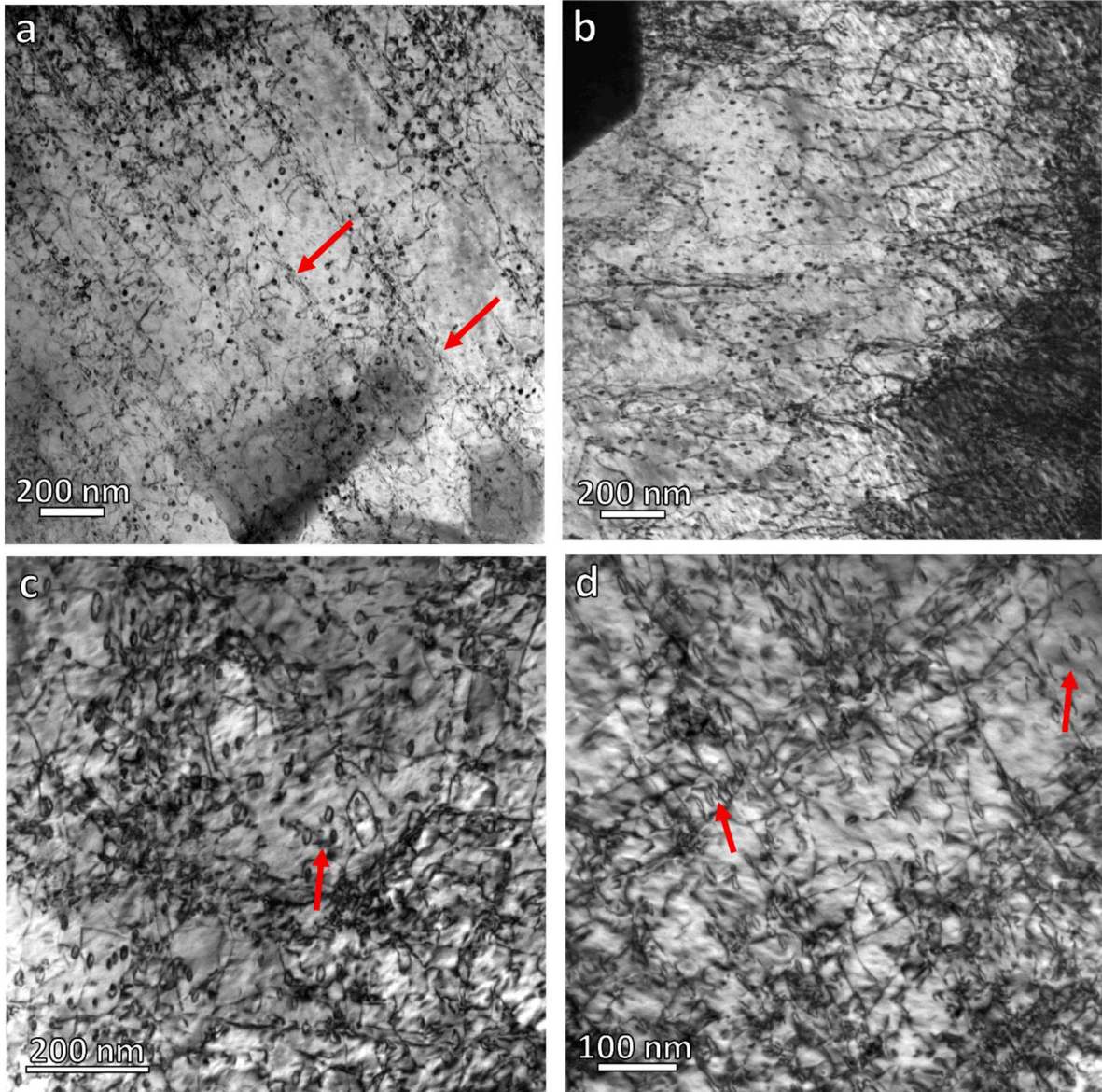
498 transition from dislocation shearing to bypassing of the precipitates at an earlier stage of age hardening will
499 be explored in the following discussion on dislocation recovery.



500

501 *Figure 10 – Representative bright-field micrographs of the bulk slip morphology in the non-charged*
502 *peak-aged specimen strained to 2% in compression. The red arrows in (b) and (d) indicate evidence of*
503 *coupled dislocation pairs.*

504



505

506 *Figure 11 – Representative bright-field micrographs of the bulk slip morphology in the hydrogen-charged*
 507 *peak-aged specimens strained to failure. The red arrows in (a) indicate planar slip bands oriented along*
 508 *the trace of the {111} plane, while the red arrows in (c) and (d) denote examples of dislocation looping.*

509 *4.1.3. Effect of hydrogen on dislocation recovery ($-d\theta/d\sigma$)*

510 Figure 7b demonstrates that $-d\theta/d\sigma$ is increased in the presence of H for all heat treatments except
 511 the OA condition, where the non-charged and H-charged alloys exhibited a similar $-d\theta/d\sigma$. While the
 512 physical basis for this ‘dynamic recovery’ parameter is widely debated in the literature [34–36,41,103], it
 513 is consistently related to the annihilation of dislocation line length, with larger values of $-d\theta/d\sigma$ implying
 514 increased rates of annihilation [34–36,42]. Moreover, unlike the θ_{\max} parameter, $-d\theta/d\sigma$ can be
 515 unambiguously linked to the rate of dislocation annihilation (k_2) in the Kocks-Mecking model for work

516 hardening (Eqn. 1), as demonstrated by Cheng *et al.* in AA6111 [42]. It therefore follows that H must be
517 increasing the rate of dislocation annihilation. Mechanistically, the two most commonly invoked pathways
518 by which this increase in $-d\theta/d\sigma$ would be achieved are: (1) increased propensity for cross-slip or (2) easier
519 activation of dislocation climb [34–36]. That H may act to increase the propensity for dislocation cross-slip
520 is inconsistent with prior literature, which widely indicates that H reduces the probability of cross-slip either
521 by stabilizing the edge component of dislocations and/or reducing the stacking fault energy [3,18–20]. That
522 being said, the observation of loop ‘stacks’ in Figure 11d (one example highlighted by the left red arrow)
523 in the PA/H alloy are consistent with an increased propensity for cross-slip, since these dislocation
524 arrangements are commonly considered to involve multiple cross-slip steps [95,104,105]. These opposing
525 observations are consistent with the findings of recent simulations examining the effect of H on the
526 propensity for cross-slip in Ni [106], which illustrates the complexity of the problem and highlights that
527 there may not be a simple answer regarding the role of cross-slip.

528 Considering vacancies, recent experimental and atomistic simulations agree that H lowers the
529 vacancy formation energy, thus stabilizing an increased vacancy concentration above that expected based
530 on thermal equilibrium in the absence of H [9,10,16,107–110]. For example, the simulations of Tanguy *et al.*
531 *al.* indicate that the total vacancy concentration in Ni at 300 K would be increased by six orders of
532 magnitude for the concentration of 200 wppm H observed in the current study [107]. A similar increase in
533 expected vacancy concentration was also found by Metsue *et al.* for calculations on Ni exposed to ~ 100
534 MPa H₂ [108]. This elevated vacancy concentration may then act to increase the propensity for dislocation
535 climb, thereby offering a potential explanation for both the observation of particle bypassing at smaller
536 precipitate sizes and the increased rates of recovery ($-d\theta/d\sigma$) after H charging. Speculatively, it is
537 hypothesized that the dislocation-precipitate interaction changes from the typical weak/strong-pair coupling
538 to a mechanism that is local climb-mediated [111]. Since dislocations overcome precipitates by thermal
539 activation [103], it is possible that while the dislocation is waiting for a successful activation event, it can
540 interact with H-stabilized vacancies, vacancy clusters or vacancy-H complexes. This interaction is then
541 postulated to enable the dislocation to overcome the precipitate by locally climbing over it rather than
542 shearing or bowing around it. Critically, it has been shown that H-vacancy clusters in the glide plane will
543 interact with dislocations leading to the formation of jog-pairs, thus facilitating the climb process as well
544 as dipole debris formation [16]. Both of these dislocation-vacancy interactions would then explain the
545 observed dislocation storage rate (θ_{\max}) for the PA/H and OA/H conditions since they would lead to the
546 formation of a significant amount of debris, such as that shown in Figure 11d. This hypothesis is also
547 consistent with the observed increased recovery rate (k_2), especially in the case of the PA/H sample. As
548 previous researchers have noted, increased recovery under conditions leading to particle bypass was

549 attributed to such debris annihilation reactions [42]. Finally, it is important to consider why this increased
550 recovery rate is not observed in the H-charged OA sample. It has been shown, albeit for steady-state creep,
551 that the climb-mediated bypass of coherent precipitates is greatly facilitated for small particles [112,113].
552 Thus, as the size of the precipitates increase, this climb mediated mechanism is replaced by the usual
553 weak/strong pair coupling mechanism or Orowan looping, depending on the precipitate size.

554 **4.2. Other possible contributions of hydrogen to dislocation-precipitate interactions**

555 The above results indicate that H modifies dislocation-precipitate interactions, resulting in the onset
556 of particle bypassing at smaller precipitate sizes (Figures 10-11). Based on the work hardening results
557 presented in Figure 7, it is postulated that this transition in dislocation-precipitate interactions may be
558 attributed to local climb enabled by H-induced increases in the equilibrium vacancy concentration.
559 However, it is important to consider other possible intrinsic H effects that could result in the observed
560 behavior. Specifically, the transition from particle shearing to bypassing could also be attributed to either
561 (a) an H-induced increase in the resistance to particle shearing (*i.e.* a modification of the precipitate itself)
562 and/or (b) a modification of the dislocation properties which would enable easier cross-slip or climb.
563 Regarding (a), the resistance to particle shearing could be promoted by modifications in: (1) particle size
564 (already refuted in Figure 8), (2) chemical (γ/γ' interface) strengthening, (3) stacking fault strengthening,
565 (4) coherency strengthening *via* distinct changes in the lattice constants of the matrix and precipitate phases,
566 (5) modulus strengthening, and (6) order strengthening [103,114–116]. It is unlikely that H segregation to
567 the γ/γ' interface would act to increase interfacial energy [117,118]. Stacking fault strengthening occurs
568 when the stacking fault energy (SFE) of the precipitate is much lower than the matrix [103]; however, given
569 that the SFE in Ni₃Al is likely larger than the Ni-Cu matrix [119], such an effect can likely be neglected
570 [46,120]. Differences from coherency strengthening can also be neglected since the stresses due to misfit
571 strains are small (the misfit strain between γ/γ' is on the order of <0.1% for Monel K-500 [44,46,49]) and
572 the presence of H would only reduce them further (since H has been shown to “shield” stresses [6]).
573 Considering modulus strengthening, the shear moduli of the γ and γ' phases (assuming Ni-30Cu for γ) are
574 quite similar, suggesting that such an effect would be negligible. Experimentalists recently suggested that
575 hydrogen can reduce the elastic modulus in Ni [121], but these assertions are counter to atomistic modeling
576 [122] and experiments by Hachet *et al.*, which have shown that the previously measured decreases in elastic
577 properties with H in Ni are likely due to enhanced vacancy concentrations and vacancy clustering induced
578 during H charging, thereby implying that H does not strongly affect elastic properties [123]. Finally,
579 consider order hardening, which is the dominant contribution to strength in γ' -strengthened alloys [46,103].
580 As documented by Ardell [114], this strengthening mechanism is sensitive to variations in the antiphase
581 boundary energy (γ_{APB}), with larger γ_{APB} values resulting in an increased resistance to shearing. Although

582 the effect of H on the γ_{APB} for Ni₃Al has not been explicitly studied, the defactant concept [124] suggests
583 that H will reduce the energy of such interfaces, which would then *increase* the propensity for particle
584 shearing [103,114,120]. In short, though the topic merits further detailed study, all of the possible means
585 by which H would increase the difficulty of shearing seem unlikely, suggesting that H-induced
586 modifications in precipitate properties are not responsible for the observed transition in behavior.

587 Considering (b), it is possible that H-induced modifications in dislocation properties could explain
588 the onset of dislocation bypassing in the PA/H alloy. Per the “defactant” concept developed by Kirchheim
589 [124] and recently considered by Sills and Cai [125], H could reduce the line energy of dislocations, thereby
590 enabling easier kink- or jog-pair nucleation required for particle bypassing by cross-slip or climb,
591 respectively. Coupled with the expected H-induced increases in the vacancy concentration [9,10,16,107–
592 110], especially if localized at the precipitate interface, this reduction in the dislocation line energy would
593 act to further ease the propensity for dislocation climb. As discussed in the preceding section, such an
594 increased ease of bypass would logically explain both the early transition from shearing to bypass, increased
595 initial hardening rates, and increased rates of recovery observed for H-charged samples in the UA and PA
596 conditions. However, detailed vetting of such synergistic influences will require further targeted modeling
597 and experimental efforts.

598 5. Conclusions

599 The influences of H on dislocation-precipitate interactions in a γ' -strengthened Ni-Cu superalloy were
600 systematically evaluated through mechanical testing on four isothermal heat treatments. Analysis of these
601 results, coupled with targeted transmission electron microscopy of the PA heat treatments, revealed several
602 important insights into H effects on deformation processes in precipitation-hardened alloys:

- 603 1. For all heat treatments, specimens pre-charged with H exhibit increased yield strength and strongly
604 reduced ductility compared to non-charged specimens.
- 605 2. All heat treatments exhibited a clear transition in fracture morphology from ductile, microvoid
606 coalescence to brittle, intergranular failure when pre-charged with H.
- 607 3. Analysis of the strain hardening behavior revealed a strong increase in θ_{max} for the H-charged
608 specimens relative to the non-charged specimens for the PA and OA conditions. An increased -
609 $d\theta/d\sigma$ was observed for the H-charged condition in the NA, UA, and PA conditions, while - $d\theta/d\sigma$
610 was unaffected by H-charging in the OA condition.
- 611 4. An analysis of the change in yield strength with H-charging as a function of heat treatment
612 suggested that H modifies dislocation-precipitate interactions. This postulation was confirmed by
613 TEM analysis of the deformation structure in the non-charged and H-charged PA specimens,

614 which showed that H induces a transition from particle shearing to bypassing at a shorter aging
615 time.

616 5. While additional experiments/modeling are needed to further explore the effect of H on
617 dislocation-precipitate interactions, it is hypothesized that the current results can be explained by
618 a local climb-mediated bypass mechanism facilitated by H-induced increases in the equilibrium
619 vacancy concentration.

620 6. Acknowledgements

621 The assistance of Dr. Chris San Marchi and Mr. Jeff Campbell with H charging, Mr. Michael Ritzo with
622 TEM sample preparation and mechanical testing, as well as helpful discussions with Prof. Richard
623 Gangloff, Prof. John Scully, Dr. Brian Somerday, Prof. Petros Sofronis, and Prof. David McDowell are
624 all gratefully acknowledged. ZDH acknowledges the support of the ALCOA Graduate Fellowship and the
625 Volkswagen Group of North America Graduate Fellowship administered by the School of Engineering
626 and Applied Science at the University of Virginia. SRA and JJB were supported by the Army Research
627 Laboratory under Cooperative Agreement Number W911NF-15-2-0025. Sandia National Laboratories is
628 a multimission laboratory managed and operated by National Technology and Engineering Solutions of
629 Sandia, LLC., a wholly owned subsidiary of Honeywell International, Inc., for the U.S. Department of
630 Energy's National Nuclear Security Administration under contract DE-NA-0003525. **Disclaimer:** The
631 views and conclusions contained in this document are those of the authors and should not be interpreted
632 as representing the official policies, either expressed or implied, of the Army Research Laboratory or the
633 U.S. Government. The U.S. government is authorized to reproduce and distribute reprints for government
634 purposes notwithstanding any copyright notation hereon.

635 7. Data Availability Statement

636 The raw and processed data generated during this study will be made available upon reasonable
637 request.

638 8. References

- 639 [1] W.H. Johnson, On some remarkable changes produced in iron and steel by the action of hydrogen
640 and acids, Proc. R. Soc. London. 23 (1874) 168–179. doi:10.1098/rspl.1874.0024.
- 641 [2] S.P. Lynch, Mechanisms and Kinetics of Environmentally Assisted Cracking: Current Status,
642 Issues, and Suggestions for Further Work, Metall. Mater. Trans. A. 44A (2013) 1209–1229.
643 doi:10.1007/s11661-012-1359-2.
- 644 [3] I.M. Robertson, P. Sofronis, A. Nagao, M.L. Martin, S. Wang, D.W. Gross, K.E. Nygren,
645 Hydrogen Embrittlement Understood, Metall. Mater. Trans. B. (2015). doi:10.1007/s11663-015-
646 0325-y.

- 647 [4] W.W. Gerberich, Modelling hydrogen induced damage mechanisms in metals, in: R.P. Gangloff,
648 B.P. Somerday (Eds.), *Gaseous Hydrog. Embrittlement Mater. Energy Technol.*, Woodhead
649 Publishing, 2012: pp. 209–246.
- 650 [5] S.P. Lynch, Environmentally assisted cracking: Overview of evidence for an adsorption-induced
651 localised-slip process, *Acta Metall.* 36 (1988) 2639–2661. doi:10.1016/0001-6160(88)90113-7.
- 652 [6] H.K. Birnbaum, P. Sofronis, Hydrogen-enhanced localized plasticity—a mechanism for hydrogen-
653 related fracture, *Mater. Sci. Eng. A.* 176 (1994) 191–202. doi:10.1016/0921-5093(94)90975-X.
- 654 [7] R.A. Oriani, A Mechanistic Theory of Hydrogen Embrittlement of Steels, *Berichte Der Bunsen-
655 Gesellschaft Fur Phys. Chemie.* 76 (1972) 848–857.
- 656 [8] A.R. Troiano, The Role of Hydrogen and Other Interstitials in the Mechanical Behavior of Metals:
657 (1959 Edward De Mille Campbell Memorial Lecture), *Metallogr. Microstruct. Anal.* 5 (2016)
658 557–569. doi:10.1007/s13632-016-0319-4.
- 659 [9] M. Nagumo, Hydrogen related failure of steels – a new aspect, *Mater. Sci. Technol.* 20 (2004)
660 940–950. doi:10.1179/026708304225019687.
- 661 [10] S.K. Lawrence, Y. Yagodzinskyy, H. Hanninen, E. Korhonen, F. Tuomisto, Z.D. Harris, B.P.
662 Somerday, Effects of grain size and deformation temperature on hydrogen-enhanced vacancy
663 formation in Ni alloys, *Acta Mater.* 128 (2017) 218–226. doi:10.1016/j.actamat.2017.02.016.
- 664 [11] Z.D. Harris, S.K. Lawrence, D.L. Medlin, G. Guetard, J.T. Burns, B.P. Somerday, Elucidating the
665 contribution of mobile hydrogen-deformation interactions to hydrogen-induced intergranular
666 cracking in polycrystalline nickel, *Acta Mater.* (2018). doi:10.1016/j.actamat.2018.07.043.
- 667 [12] Y. Deng, A. Barnoush, Hydrogen embrittlement revealed via novel in situ fracture experiments
668 using notched micro-cantilever specimens, *Acta Mater.* 142 (2018) 236–247.
669 doi:10.1016/j.actamat.2017.09.057.
- 670 [13] S. Wang, A. Nagao, K. Edalati, Z. Horita, I.M. Robertson, Influence of hydrogen on dislocation
671 self-organization in Ni, *Acta Mater.* 135 (2017) 96–102. doi:10.1016/j.actamat.2017.05.073.
- 672 [14] M.L. Martin, M. Dadfarnia, A. Nagao, S. Wang, P. Sofronis, Enumeration of the hydrogen-
673 enhanced localized plasticity mechanism for hydrogen embrittlement in structural materials, *Acta
674 Mater.* 165 (2019) 734–750. doi:10.1016/j.actamat.2018.12.014.
- 675 [15] M. Nagumo, K. Takai, The predominant role of strain-induced vacancies in hydrogen
676 embrittlement of steels: Overview, *Acta Mater.* 165 (2019) 722–733.
677 doi:10.1016/j.actamat.2018.12.013.
- 678 [16] A. Tehranchi, W.A. Curtin, The role of atomistic simulations in probing hydrogen effects on
679 plasticity and embrittlement in metals, *Eng. Fract. Mech.* (2019).
680 doi:10.1016/j.engfracmech.2019.106502.
- 681 [17] C.D. Beachem, A new model for hydrogen-assisted cracking (hydrogen “embrittlement”), *Metall.
682 Trans.* 3 (1972) 441–455. doi:10.1007/BF02642048.
- 683 [18] P.J. Ferreira, I.M. Robertson, H.K. Birnbaum, Hydrogen effects on the interaction between
684 dislocations, *Acta Mater.* 46 (1998) 1749–1757. doi:10.1016/S1359-6454(97)00349-2.
- 685 [19] I.M. Robertson, The effect of hydrogen on dislocation dynamics, *Eng. Fract. Mech.* 65 (2001)

- 686 671–692.
- 687 [20] D. Delafosse, Hydrogen effects on the plasticity of face centred cubic (fcc) crystals, in: R.P.
688 Gangloff, B.P. Somerday (Eds.), *Gaseous Hydrog. Embrittlement Mater. Energy Technol.*,
689 Woodhead Publishing, 2012.
- 690 [21] A. Barnoush, H. Vehoff, Recent developments in the study of hydrogen embrittlement: Hydrogen
691 effect on dislocation nucleation, *Acta Mater.* 58 (2010) 5274–5285.
692 doi:10.1016/j.actamat.2010.05.057.
- 693 [22] E. Tal-Gutelmacher, R. Gemma, C.A. Volkert, R. Kirchheim, Hydrogen effect on dislocation
694 nucleation in a vanadium (1 0 0) single crystal as observed during nanoindentation, *Scr. Mater.* 63
695 (2010) 1032–1035. doi:10.1016/j.scriptamat.2010.07.039.
- 696 [23] K.A. Nibur, D.F. Bahr, B.P. Somerday, Hydrogen effects on dislocation activity in austenitic
697 stainless steel, *Acta Mater.* 54 (2006) 2677–2684. doi:10.1016/j.actamat.2006.02.007.
- 698 [24] Y.Z. Chen, H.P. Barth, M. Deutges, C. Borchers, F. Liu, R. Kirchheim, Increase in dislocation
699 density in cold-deformed Pd using H as a temporary alloying addition, *Scr. Mater.* 68 (2013) 743–
700 746. doi:10.1016/j.scriptamat.2013.01.005.
- 701 [25] Y.Z. Chen, X.Y. Ma, X.H. Shi, T. Suo, C. Borchers, K.H. Zhang, F. Liu, R. Kirchheim, Hardening
702 effects in plastically deformed Pd with the addition of H, *Scr. Mater.* 98 (2015) 48–51.
703 doi:10.1016/j.scriptamat.2014.11.012.
- 704 [26] L. Balogh, L. Capolungo, C.N. Tomé, On the measure of dislocation densities from diffraction line
705 profiles: A comparison with discrete dislocation methods, *Acta Mater.* 60 (2012) 1467–1477.
706 doi:10.1016/j.actamat.2011.10.037.
- 707 [27] A.M. Brass, D. Roux, J. Chêne, Role of secondary γ' precipitation and of hydrogen in the first
708 stage of the plastic deformation of the ν matrix of a Ni base superalloy single crystal, *Mater. Sci.*
709 *Eng. A.* 323 (2002) 97–102. doi:10.1016/S0921-5093(01)01349-1.
- 710 [28] Z. Guo, M. Zhao, C. Li, S. Chen, L. Rong, Mechanism of hydrogen embrittlement in a gamma-
711 prime phase strengthened Fe-Ni based austenitic alloy, *Mater. Sci. Eng. A.* 555 (2012) 77–84.
712 doi:10.1016/j.msea.2012.06.036.
- 713 [29] A.W. Thompson, I.M. Bernstein, The Role Of Metallurgical Variables In Hydrogen-Assisted
714 Environmental Fracture, *Adv. Corros. Sci. Technol.* (1980) 53–175. doi:10.1007/978-1-4615-
715 9065-1_2.
- 716 [30] I.M. Robertson, H.K. Birnbaum, Effect of hydrogen on the dislocation structure of deformed
717 nickel, *Scr. Metall.* 18 (1984) 269–274. doi:10.1016/0036-9748(84)90521-0.
- 718 [31] W.A. McInteer, A.W. Thompson, I.M. Bernstein, The effect of hydrogen on the slip character of
719 nickel, *Acta Metall.* 28 (1980) 887–894. doi:10.1007/BF02655107.
- 720 [32] J.A. Wert, X. Huang, G. Winther, W. Pantleon, H.F. Poulsen, Revealing deformation
721 microstructures, *Mater. Today.* 10 (2007) 24–32. doi:10.1016/S1369-7021(07)70206-7.
- 722 [33] J. Jiang, T. Ben Britton, A.J. Wilkinson, The orientation and strain dependence of dislocation
723 structure evolution in monotonically deformed polycrystalline copper, *Int. J. Plast.* 69 (2015) 102–
724 117. doi:10.1016/j.ijplas.2015.02.005.

- 725 [34] U.F. Kocks, H. Mecking, Physics and phenomenology of strain hardening: The FCC case, *Prog.*
726 *Mater. Sci.* 48 (2003) 171–273. doi:10.1016/S0079-6425(02)00003-8.
- 727 [35] Y. Estrin, Dislocation-Density-Related Constitutive Modeling, in: A.S. Krausz, K. Krausz (Eds.),
728 *Unified Const. Laws Plast. Deform.*, Academic Press, 1996: pp. 69–107.
- 729 [36] E. Nes, Modelling of work hardening and stress saturation in FCC metals, *Prog. Mater. Sci.* 41
730 (1997) 129–193. doi:10.1016/S0079-6425(97)00032-7.
- 731 [37] C. Keller, E. Hug, Kocks-Mecking analysis of the size effects on the mechanical behavior of
732 nickel polycrystals, *Int. J. Plast.* 98 (2017) 106–122. doi:10.1016/j.ijplas.2017.07.003.
- 733 [38] Y. Estrin, L.P. Kubin, Local strain hardening and nonuniformity of plastic deformation, *Acta*
734 *Metall.* 34 (1986) 2455–2464.
- 735 [39] J. Christopher, B.K. Choudhary, Kinetics of Uniaxial Tensile Flow and Work Hardening Behavior
736 of Type 316L(N) Austenitic Stainless Steel in the Framework of Two-Internal-Variable Approach,
737 *Metall. Mater. Trans. A Phys. Metall. Mater. Sci.* 46 (2015) 674–687. doi:10.1007/s11661-014-
738 2660-z.
- 739 [40] F. Barlat, M.V. Glazov, J.C. Brem, D.J. Lege, A simple model for dislocation behavior, strain and
740 strain rate hardening evolution in deforming aluminum alloys, *Int. J. Plast.* 18 (2002) 919–939.
741 doi:10.1016/S0749-6419(01)00015-8.
- 742 [41] D. Kuhlmann-Wilsdorf, Theory of plastic deformation: - properties of low energy dislocation
743 structures, *Mater. Sci. Eng. A.* 113 (1989) 1–41. doi:10.1016/0921-5093(89)90290-6.
- 744 [42] L.M. Cheng, W.J. Poole, J.D. Embury, D.J. Lloyd, The influence of precipitation on the work-
745 hardening behavior of the aluminum alloys AA6111 and AA7030, *Metall. Mater. Trans. A.* 34
746 (2003) 2473–2481. doi:10.1007/s11661-003-0007-2.
- 747 [43] K.M. Bertsch, Hydrogen effects on the evolution of plastic deformation in polycrystalline nickel: a
748 mechanism for intergranular fracture, University of Illinois at Urbana-Champaign, 2017.
- 749 [44] G.K. Dey, P. Mukhopadhyay, Precipitation in the Ni-Cu-base alloy Monel K-500, *Mater. Sci. Eng.*
750 84 (1986) 177–189. doi:10.1016/0025-5416(86)90236-3.
- 751 [45] L.E. Shoemaker, G.D. Smith, A Century of Monel Metal: 1906 – 2006, *JOM.* 58 (2006) 22–26.
- 752 [46] G.K. Dey, R. Tewari, P. Rao, S.L. Wadekar, P. Mukhopadhyay, Precipitation hardening in nickel-
753 copper base alloy monel K 500, *Metall. Trans. A.* 24 (1993) 2709–2719.
754 doi:10.1007/BF02659495.
- 755 [47] J.R. Davis, ed., Nickel, Cobalt, and Their Alloys: ASM Specialty Handbook, in: 2000: pp. 167–
756 188.
- 757 [48] Å. Gustafson, Coarsening of TiC in austenitic stainless steel - Experiments and simulations in
758 comparison, *Mater. Sci. Eng. A.* 287 (2000) 52–58. doi:10.1016/S0921-5093(00)00827-3.
- 759 [49] Z.D. Harris, J.T. Burns, The effect of isothermal heat treatment on hydrogen environment-assisted
760 cracking susceptibility in Monel K-500, *Mater. Sci. Eng. A.* 764 (2019) 138249.
761 doi:10.1016/j.msea.2019.138249.
- 762 [50] S. Bechtle, M. Kumar, B.P. Somerday, M.E. Launey, R.O. Ritchie, Grain-boundary engineering

- 763 markedly reduces susceptibility to intergranular hydrogen embrittlement in metallic materials,
764 *Acta Mater.* 57 (2009) 4148–4157. doi:10.1016/j.actamat.2009.05.012.
- 765 [51] B.P. Somerday, J.A. Campbell, K.L. Lee, J.A. Ronevich, C. San Marchi, Enhancing safety of
766 hydrogen containment components through materials testing under in-service conditions, *Int. J.*
767 *Hydrogen Energy.* (2016) 4–11. doi:10.1016/j.ijhydene.2016.04.189.
- 768 [52] J. Crank, *Mathematics of Diffusion*, 1975. doi:10.1016/0306-4549(77)90072-X.
- 769 [53] C.S. Marchi, B.P. Somerday, S.L. Robinson, Permeability, solubility and diffusivity of hydrogen
770 isotopes in stainless steels at high gas pressures, *Int. J. Hydrogen Energy.* 32 (2007) 100–116.
771 doi:10.1016/j.ijhydene.2006.05.008.
- 772 [54] F. Jones, R. Pehlke, Solubility of hydrogen in solid Ni–Co and Ni–Cu alloys, *Metall. Trans.* 2
773 (1971). <http://link.springer.com/article/10.1007/BF02814909>.
- 774 [55] C. Ko, R.B. McLellan, The thermodynamics of Ni-Cu-H solid solutions, *Acta Metall. Mater.* 41
775 (1993) 2473–2476. doi:10.1016/0956-7151(93)90327-O.
- 776 [56] H. Hagi, Diffusion Coefficients of Hydrogen in Ni-Cu and Ni-Co Alloys, *Trans. Japan Inst. Met.*
777 27 (1986) 233–240. doi:10.2320/matertrans1960.27.233.
- 778 [57] D.J. Mitchell, E.M. Edge, Permeation characteristics of some iron and nickel based alloys, *J. Appl.*
779 *Phys.* 57 (1985) 5226–5235. doi:10.1063/1.335261.
- 780 [58] Y. Sakamoto, O. Shimizu, K. Hirayama, K. Baba, The diffusivity of hydrogen in nickel-based
781 solid solutions, *J. Phys. Chem. Solids.* 49 (1988) 897–903. doi:10.1016/0022-3697(88)90006-6.
- 782 [59] J. Ai, H.M. Ha, R.P. Gangloff, J.R. Scully, Hydrogen diffusion and trapping in a precipitation-
783 hardened nickel–copper–aluminum alloy Monel K-500 (UNS N05500), *Acta Mater.* 61 (2013)
784 3186–3199. doi:10.1016/j.actamat.2013.02.007.
- 785 [60] R.P. Gangloff, H.M. Ha, J.T. Burns, J.R. Scully, Measurement and Modeling of Hydrogen
786 Environment-Assisted Cracking in Monel K-500, *Metall. Mater. Trans. A.* 45 (2014) 3814–3834.
787 doi:10.1007/s11661-014-2324-z.
- 788 [61] B.C.R. Troconis, Z.D. Harris, H. Ha, J.T. Burns, J.R. Scully, The effect of heat-to-heat variations
789 in metallurgy and hydrogen-metal interactions on the hydrogen embrittlement of Monel K-500,
790 *Mater. Sci. Eng. A.* 703 (2017) 533–550. doi:10.1016/j.msea.2017.07.019.
- 791 [62] M.I. Baskes, J.E. Angelo, N.R. Moody, SAND-94-8741C: Atomistic calculations of hydrogen
792 interactions with Ni3Al grain boundaries and Ni/Ni3Al interfaces, 1994.
- 793 [63] D.H. Lassila, H.K. Birnbaum, The effect of diffusive hydrogen segregation on fracture of
794 polycrystalline nickel, *Acta Metall.* 34 (1986) 1237–1243. doi:10.1016/0001-6160(86)90010-6.
- 795 [64] T.H. Courtney, *Mechanical Behavior of Materials*, 2nd ed., McGraw-Hill, 2000.
- 796 [65] ASTM Standard E647-13: Standard Test Method for Measurement of Fatigue Crack Growth
797 Rates, (2013). doi:10.1520/E0647-13E01.2.
- 798 [66] G.A.F. Seber, C.J. Wild, *Nonlinear Regression*, Wiley, 1989.
- 799 [67] R.E. Reed-Hill, R. Abbaschian, *Physical metallurgy principles*, Van Nostrand, Princeton, 1994.

- 800 [68] J.A. Harris, R.C. Scarberry, C.D. Stephens, Effects of Hydrogen on Engineering Properties of
801 MONEL Nickel-Copper Alloy K-500, *Corrosion*. 28 (1972) 57–62.
- 802 [69] K.D. Efind, Failure of Monel Ni-Cu-Al Alloy K-500 Bolts in Seawater, *Mater. Perform.* 24 (1985)
803 37–40.
- 804 [70] L.H. Wolfe, M.W. Joosten, Failures of Nickel/Copper Bolts in Subsea Applications, *SPE Prod.*
805 *Eng.* 3 (1988) 382–386. doi:10.2118/16632-PA.
- 806 [71] Z.D. Harris, J.D. Dolph, G.L. Pioszak, B.C. Rincon Troconis, J.R. Scully, J.T. Burns, The Effect
807 of Microstructural Variation on the Hydrogen Environment-Assisted Cracking of Monel K-500,
808 *Metall. Mater. Trans. A.* 47 (2016) 3488–3510. doi:10.1007/s11661-016-3486-7.
- 809 [72] J.T. Burns, Z.D. Harris, J.D. Dolph, R.P. Gangloff, Measurement and Modeling of Hydrogen
810 Environment Assisted Cracking in a Ni-Cu-Al-Ti Superalloy, *Metall. Mater. Trans. A.* 47 (2016)
811 990–997. doi:10.1007/s11661-015-3315-4.
- 812 [73] M.L. Martin, B.P. Somerday, R.O. Ritchie, P. Sofronis, I.M. Robertson, Hydrogen-induced
813 intergranular failure in nickel revisited, *Acta Mater.* 60 (2012) 2739–2745.
814 doi:10.1016/j.actamat.2012.01.040.
- 815 [74] Z.D. Harris, Personal communication with Warren Poole and Sean Agnew, (2017).
- 816 [75] D.G. Ulmer, C.J. Altstetter, Hydrogen-induced strain localization and failure of austenitic stainless
817 steels at high hydrogen concentrations, *Acta Metall. Mater.* 39 (1991) 1237–1248. doi:Doi:
818 10.1016/0956-7151(91)90211-i.
- 819 [76] V. Randle, D. Horton, Grain growth phenomena in nickel, *Scr. Metall. Mater.* 31 (1994) 891–895.
820 doi:10.1016/0956-716X(94)90498-7.
- 821 [77] A.W. Thompson, Yielding in nickel as a function of grain or cell size, *Acta Metall.* 23 (1975)
822 1337–1342.
- 823 [78] Z.C. Cordero, B.E. Knight, C.A. Schuh, Six decades of the Hall–Petch effect – a survey of grain-
824 size strengthening studies on pure metals, *Int. Mater. Rev.* 61 (2016) 495–512.
825 doi:10.1080/09506608.2016.1191808.
- 826 [79] K. Takeda, N. Nakada, T. Tsuchiyama, S. Takaki, Effect of Interstitial Elements on Hall–Petch
827 Coefficient of Ferritic Iron, *ISIJ Int.* 48 (2008) 1122–1125. doi:10.2355/isijinternational.48.1122.
- 828 [80] S. Takaki, D. Akama, N. Nakada, T. Tsuchiyama, Effect of Grain Boundary Segregation of
829 Interstitial Elements on Hall-Petch Coefficient in Steels, *Mater. Trans.* 55 (2013) 28–34.
830 doi:10.2320/matertrans.ma201314.
- 831 [81] Y. Mine, N. Horita, Z. Horita, K. Takashima, Effect of ultrafine grain refinement on hydrogen
832 embrittlement of metastable austenitic stainless steel, *Int. J. Hydrogen Energy.* 42 (2017) 15415–
833 15425. doi:10.1016/j.ijhydene.2017.04.249.
- 834 [82] I.M. Bernstein, The effect of hydrogen on the deformation of iron, *Scr. Metall.* 8 (1974) 343–349.
835 doi:10.1016/0036-9748(74)90136-7.
- 836 [83] G.R. Caskey, Effect of hydrogen on work hardening of Type 304L austenitic stainless steel, *Scr.*
837 *Metall.* 15 (1981) 1183–1186.

- 838 [84] B.A. Wilcox, G.C. Smith, Intercrystalline fracture in hydrogen-charged nickel, *Acta Metall.* 13
839 (1965) 331–343.
- 840 [85] Y. Mishima, S. Ochiai, N. Hamao, M. Yodogawa, T. Suzuki, Solid Solution Hardening of Nickel -
841 Role of Transition Metal and B-subgroup Solute, *Trans. Japan Inst. Met.* 27 (1986) 656–664.
- 842 [86] N.F. Mott, Mechanical strength and creep in metals, in: *Imperfections Nearly Perfect Cryst.*, 1952:
843 pp. 173–196.
- 844 [87] M.Z. Butt, P. Feltham, Solid-solution hardening, *Acta Metall.* 26 (1978) 167–173.
- 845 [88] R.J. Arsenault, S. Patu, D.M. Esterling, Computer simulation of solid solution strengthening in
846 Fcc alloys: Part I. Friedel and mott limits, *Metall. Trans. A.* 20 (1989) 1411–1418.
847 doi:10.1007/bf02665498.
- 848 [89] D. Ma, M. Friák, J. Von Pezold, D. Raabe, J. Neugebauer, Computationally efficient and
849 quantitatively accurate multiscale simulation of solid-solution strengthening by ab initio
850 calculation, *Acta Mater.* 85 (2015) 53–66. doi:10.1016/j.actamat.2014.10.044.
- 851 [90] G.P.M. Leyson, L.G. Hector, W.A. Curtin, Solute strengthening from first principles and
852 application to aluminum alloys, *Acta Mater.* 60 (2012) 3873–3884.
853 doi:10.1016/j.actamat.2012.03.037.
- 854 [91] Y. Nakada, S. Keh, Solid solution strengthening in Fe-N single crystals, *Acta Metall.* 16 (1968)
855 903–914.
- 856 [92] W. Rauscher, W. Köster, RELATIONS BETWEEN THE MODULUS OF ELASTICITY OF
857 BINARY ALLOYS AND THEIR STRUCTURE, Stuttgart, Germany, 1951.
- 858 [93] B. Baranowski, S. Majchrzak, T.B. Flanagan, The volume increase of fcc metals and alloys due to
859 interstitial hydrogen over a wide range of hydrogen contents, *J. Phys. F Met. Phys.* 1 (1971) 258–
860 261. doi:10.1088/0305-4608/1/3/307.
- 861 [94] C.G. Schmidt, A.K. Miller, The effect of solutes on the strength and strain hardening behavior of
862 alloys, *Acta Metall.* 60 (1982) 615–625.
- 863 [95] J.W. Martin, *Micromechanisms in particle-hardened alloys*, Cambridge University Press,
864 Cambridge, UK, 1980.
- 865 [96] J.C. Fisher, E.W. Hart, R.H. Pry, The hardening of metal crystals by precipitate particles, *Acta*
866 *Metall.* 1 (1953) 336–339.
- 867 [97] L.M. Brown, D.R. Clarke, Work hardening due to internal stresses in composite materials, *Acta*
868 *Metall.* 23 (1975) 821–830. doi:10.1016/0001-6160(75)90198-4.
- 869 [98] E.M. Francis, B.M.B. Grant, J.Q. Da Fonseca, P.J. Phillips, M.J. Mills, M.R. Daymond, M.
870 Preuss, High-temperature deformation mechanisms in a polycrystalline nickel-base superalloy
871 studied by neutron diffraction and electron microscopy, *Acta Mater.* 74 (2014) 18–29.
872 doi:10.1016/j.actamat.2014.04.028.
- 873 [99] M. Preuss, J.Q. Da Fonseca, B. Grant, E. Knoche, R. Moat, M. Daymond, The effect of γ' particle
874 size on the deformation mechanism in an advanced polycrystalline nickel-base superalloy, in:
875 *Superalloys 2008*, 2008: pp. 405–414.

- 876 [100] H.J. Stone, T.M. Holden, R.C. Reed, On the generation of microstrains during the plastic
877 deformation of Waspaloy, *Acta Mater.* 47 (1999) 4435–4448. doi:10.1016/S1359-6454(99)00314-
878 6.
- 879 [101] D.E. Kim, S.L. Shang, Z.K. Liu, Effects of alloying elements on elastic properties of Ni₃Al by
880 first-principles calculations, *Intermetallics.* 18 (2010) 1163–1171.
881 doi:10.1016/j.intermet.2010.02.024.
- 882 [102] A. Bourret, Irradiation damage in nickel and iron in a high-voltage electron microscope and
883 threshold energy determination, *Phys. Status Solidi.* 4 (1971) 813–825.
- 884 [103] A.S. Argon, *Strengthening mechanisms in crystal plasticity*, Oxford University Press, 2008.
- 885 [104] F.J. Humphreys, P.B. Hirsch, The Deformation of Single Crystals of Copper and Copper-Zinc
886 Alloys Containing Alumina Particles. II. Microstructure and Dislocation-Particle Interactions,
887 *Proc. R. Soc. A Math. Phys. Eng. Sci.* 318 (1970) 73–92. doi:10.1098/rspa.1970.0134.
- 888 [105] Y. Xiang, D.J. Srolovitz, L.T. Cheng, W. E, Level set simulations of dislocation-particle bypass
889 mechanisms, *Acta Mater.* 52 (2004) 1745–1760. doi:10.1016/j.actamat.2003.12.016.
- 890 [106] Y. Tang, J.A. El-Awady, Atomistic simulations of the interactions of hydrogen with dislocations
891 in fcc metals, *Phys. Rev. B - Condens. Matter Mater. Phys.* 86 (2012) 1–13.
892 doi:10.1103/PhysRevB.86.174102.
- 893 [107] D. Tanguy, Y. Wang, D. Connétable, Stability of vacancy-hydrogen clusters in nickel from first-
894 principles calculations, *Acta Mater.* 78 (2014) 135–143. doi:10.1016/j.actamat.2014.06.021.
- 895 [108] A. Metsue, A. Oudriss, X. Feaugas, Hydrogen solubility and vacancy concentration in nickel
896 single crystals at thermal equilibrium: New insights from statistical mechanics and ab initio
897 calculations, *J. Alloys Compd.* 656 (2016) 555–567. doi:10.1016/j.jallcom.2015.09.252.
- 898 [109] S. Harada, S. Yokota, Y. Ishii, Y. Shizuku, M. Kanazawa, Y. Fukai, A relation between the
899 vacancy concentration and hydrogen concentration in the Ni-H, Co-H and Pd-H systems, *J. Alloys
900 Compd.* 404–406 (2005) 247–251. doi:10.1016/j.jallcom.2005.02.077.
- 901 [110] Y. Fukai, Formation of superabundant vacancies in M-H alloys and some of its consequences: A
902 review, *J. Alloys Compd.* 356–357 (2003) 263–269. doi:10.1016/S0925-8388(02)01269-0.
- 903 [111] J. Cadek, *Creep in Metallic Materials*, Elsevier, 1988.
- 904 [112] E.A. Marquis, D.C. Dunand, Model for creep threshold stress in precipitation-strengthened alloys
905 with coherent particles, *Scr. Mater.* 47 (2002) 503–508. doi:10.1016/S1359-6462(02)00165-3.
- 906 [113] D.N. Seidman, E.A. Marquis, D.C. Dunand, Precipitation strengthening at ambient and elevated
907 temperatures of heat-treatable Al(Sc) alloys, *Acta Mater.* 50 (2002) 4021–4035.
908 doi:10.1016/S1359-6454(02)00201-X.
- 909 [114] A.J. Ardell, Precipitation hardening, *Metall. Trans. A.* 16 (1985) 2131–2165.
910 doi:10.1007/BF02670416.
- 911 [115] T. Gladman, Precipitation hardening in metals, *Mater. Sci. Technol.* 15 (2004) 30–36.
912 doi:10.1179/026708399773002782.
- 913 [116] A. Kelly, R.B. Nicholson, Precipitation hardening, *Prog. Mater. Sci.* 10 (1963) 151–391.

- 914 [117] S. Jothi, T.N. Croft, S.G.R. Brown, Influence of grain boundary misorientation on hydrogen
915 embrittlement in bi-crystal nickel, *Int. J. Hydrogen Energy*. (n.d.).
916 doi:<http://dx.doi.org/10.1016/j.ijhydene.2014.07.020>.
- 917 [118] A. Alvaro, I. Thue Jensen, N. Kheradmand, O.M. Løvvik, V. Olden, Hydrogen embrittlement in
918 nickel, visited by first principles modeling, cohesive zone simulation and nanomechanical testing,
919 *Int. J. Hydrogen Energy*. 40 (2015) 16892–16900. doi:[10.1016/j.ijhydene.2015.06.069](https://doi.org/10.1016/j.ijhydene.2015.06.069).
- 920 [119] Y.F. Wen, J. Sun, J. Huang, First - Principles study of stacking fault energies in Ni₃Al
921 intermetallic alloys, *Trans. Nonferrous Met. Soc. China (English Ed.* 22 (2012) 661–664.
922 doi:[10.1016/S1003-6326\(11\)61229-6](https://doi.org/10.1016/S1003-6326(11)61229-6).
- 923 [120] D. Raynor, J.M. Silcock, Strengthening Mechanisms in γ' Precipitating Alloys, *Met. Sci.* 4 (1970)
924 121–130. doi:[10.1179/msc.1970.4.1.121](https://doi.org/10.1179/msc.1970.4.1.121).
- 925 [121] S.K. Lawrence, B.P. Somerday, M.D. Ingraham, D.F. Bahr, Probing the Effect of Hydrogen on
926 Elastic Properties and Plastic Deformation in Nickel Using Nanoindentation and Ultrasonic
927 Methods, *Jom*. 70 (2018) 1068–1073. doi:[10.1007/s11837-018-2850-z](https://doi.org/10.1007/s11837-018-2850-z).
- 928 [122] M. Wen, A. Barnoush, K. Yokogawa, Calculation of all cubic single-crystal elastic constants from
929 single atomistic simulation: Hydrogen effect and elastic constants of nickel, *Comput. Phys.*
930 *Commun.* 182 (2011) 1621–1625. doi:[10.1016/j.cpc.2011.04.009](https://doi.org/10.1016/j.cpc.2011.04.009).
- 931 [123] G. Hachet, A. Metsue, A. Oudriss, X. Feugas, Influence of hydrogen on the elastic properties of
932 nickel single crystal: A numerical and experimental investigation, *Acta Mater.* 148 (2018) 280–
933 288. doi:[10.1016/j.actamat.2018.01.056](https://doi.org/10.1016/j.actamat.2018.01.056).
- 934 [124] R. Kirchheim, Reducing grain boundary, dislocation line and vacancy formation energies by solute
935 segregation. I. Theoretical background, *Acta Mater.* 55 (2007) 5129–5138.
936 doi:[10.1016/j.actamat.2007.05.047](https://doi.org/10.1016/j.actamat.2007.05.047).
- 937 [125] R.B. Sills, W. Cai, Free energy change of a dislocation due to a Cottrell atmosphere, *Philos. Mag.*
938 98 (2018) 1491–1510. doi:[10.1080/14786435.2018.1441560](https://doi.org/10.1080/14786435.2018.1441560).
- 939