

# Measuring inertial forces with ultracold neutral atoms

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University of Oklahoma Seminar  
November 13, 2018

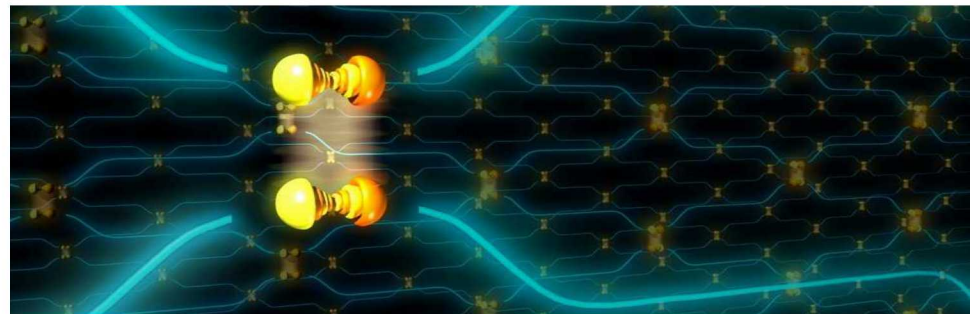
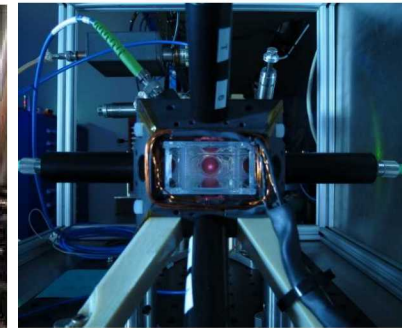
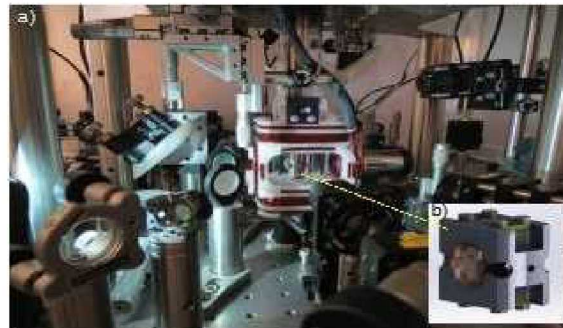


This work was supported by the Laboratory Directed Research and Development program at Sandia National Laboratories



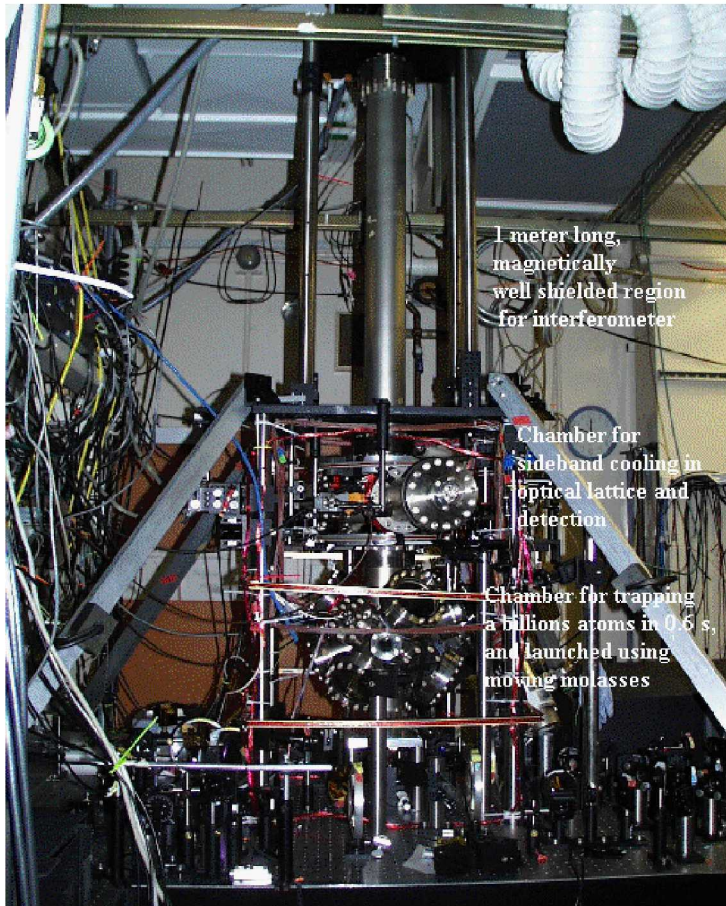
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SAND2018-13018PE

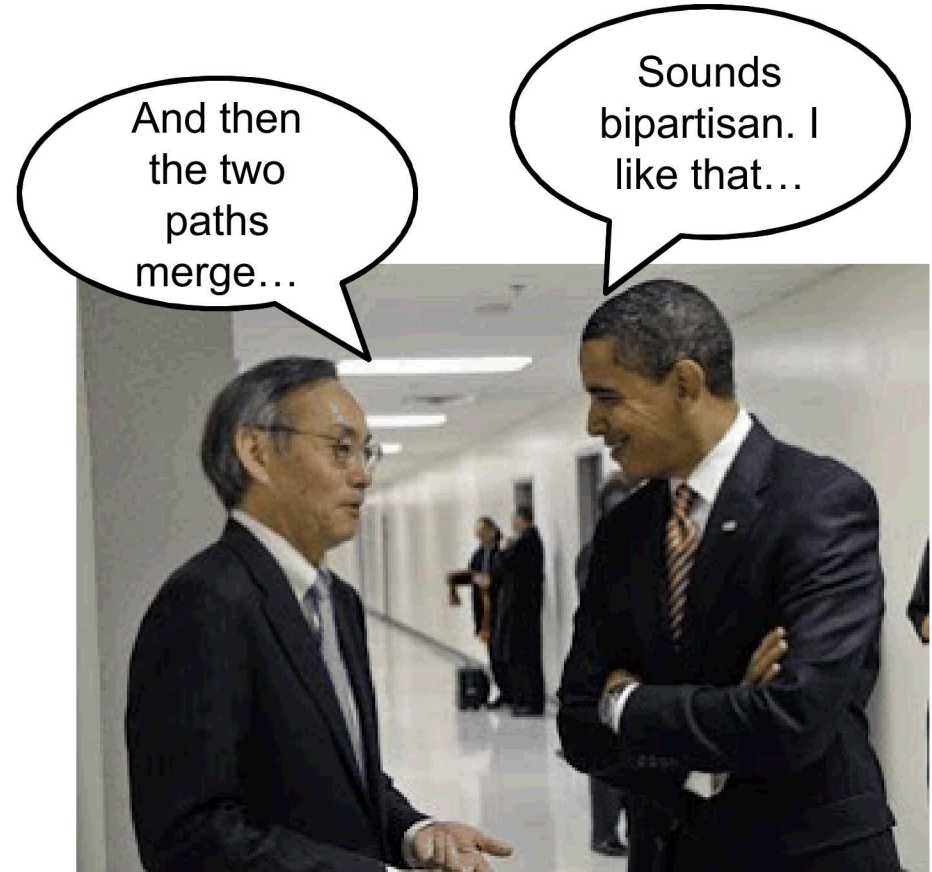


- Precision sensing with atom interferometers
- Sandia atom interferometer efforts
  - Single atom interferometer
  - Compact atom interferometer for field use
- Toward an atom interferometer with entangled spins
  - Rydberg-dressed physics and entanglement
  - Error budget and outlook

# Light-pulse atom interferometers



Stanford atom interferometer gravimeter  
Steven Chu, 1990's



Photograph by Charles Watkins, White House photographer  
Text by GB

Superb gyroscopes and gravity  
gradiometers demonstrated as well

# Applications of atom interferometer sensors

## ■ AI Gravimeter

- A. Peters *et al*, Metrologia **38**, (2001)
- Capable of 20 nano-g in 1.3 s
- Recent tower work infers pico-g level per shot
- Performance far exceeds commercial Nav. Grade

10 m tower  
Prof. Kasevich  
Stanford



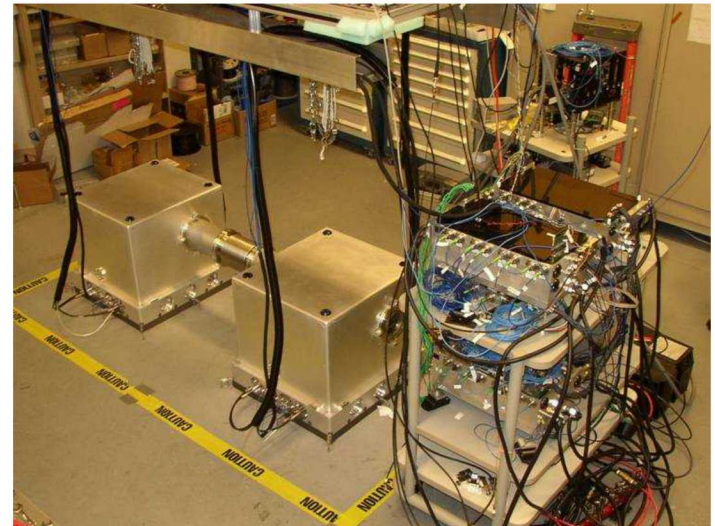
## ■ AI Gyro

- D. Durfee *et al*, PRL **97**, (2006)
- Short-term stability  $\sim 3 \mu\text{Deg/hr}^{1/2}$
- High-accuracy navigation grade
  - $< 70 \mu\text{Deg/hr}$  bias stability
- Performance exceeds commercial Nav. Grade

Stanford gradiometer—G.B. and many others

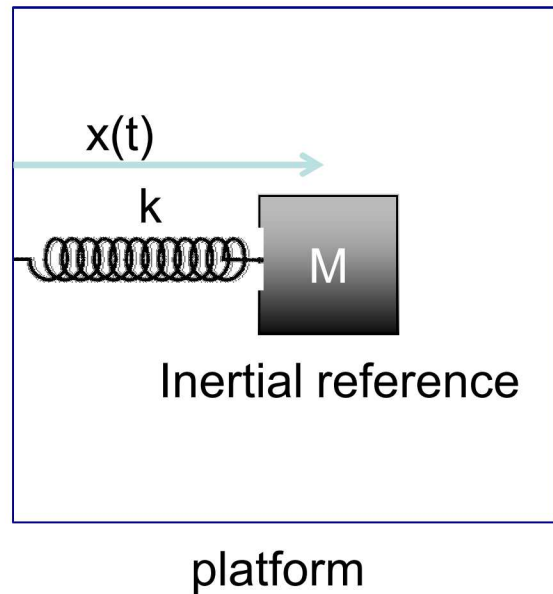
## ■ AI Gradiometer

- $\Delta a = 4 \text{ nano-g Hz}^{-1/2}$ , Biedermann, PRA 2015
- Yale, 2002 (Fixler PhD thesis); Fixler, Science, 2007.
- More recent: Florence, Nature, 2014
  - Precision  $\delta G/G \sim 7.7 \times 10^{-5}$
  - Systematic uncertainty  $6.2 \times 10^{-5}$
- CODATA 2014 relative uncertainty:  $4.7 \times 10^{-5}$

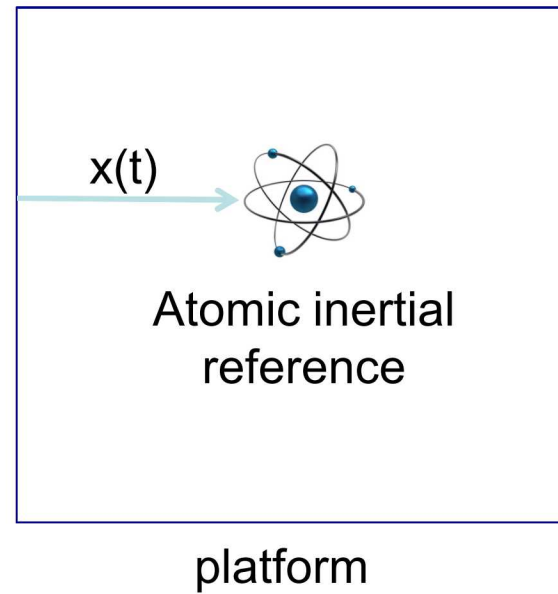


# Inertial measurements

Mass-spring accelerometer



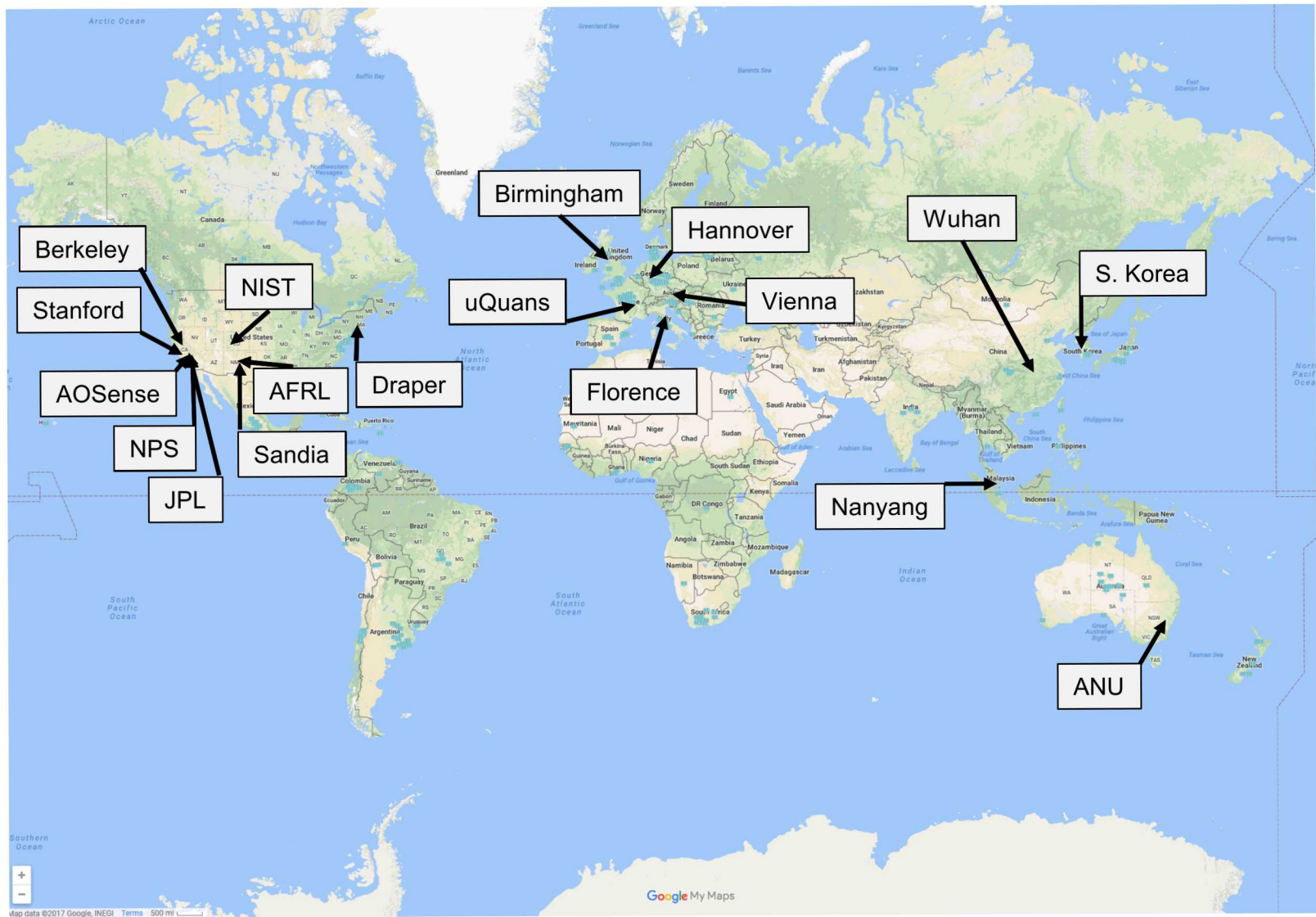
Free-fall accelerometer



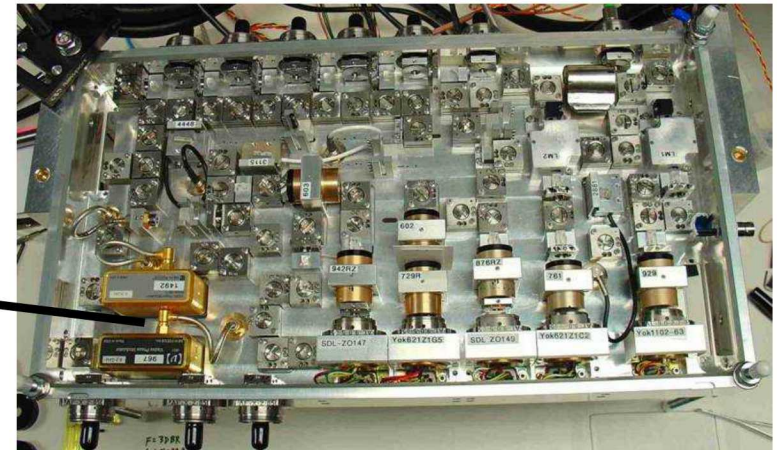
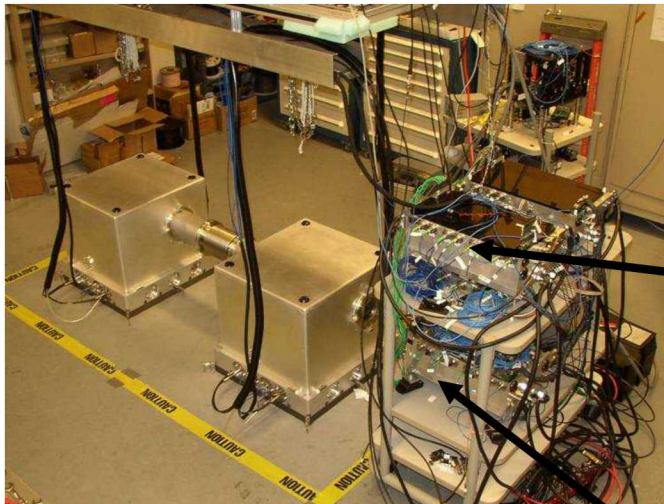
## Atom interferometers

- Inertial reference: laser-cooled atoms falling freely
- Positioning: laser ranging
- Truly accurate, turn-on calibration

# An incomplete map of AI efforts

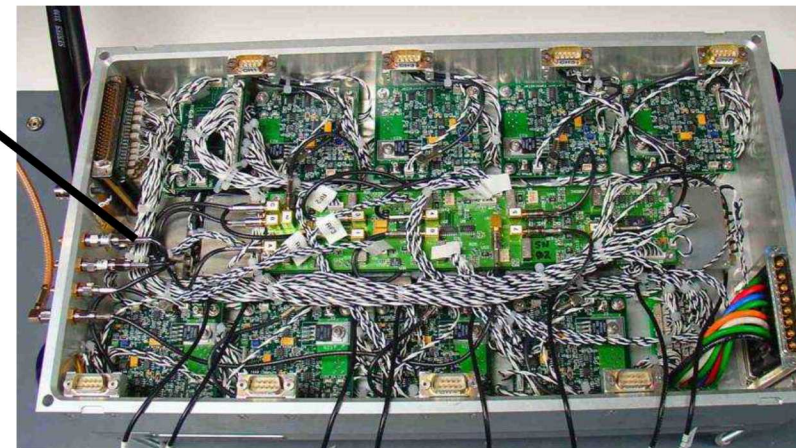
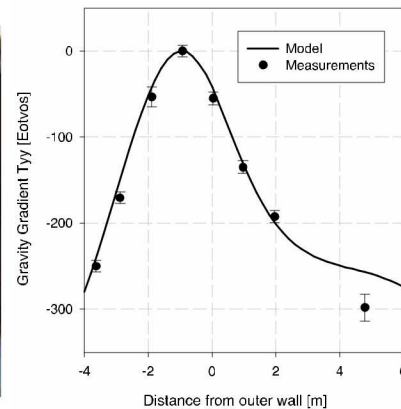


# Mobile atom interferometer



*lasers 9"x15"*

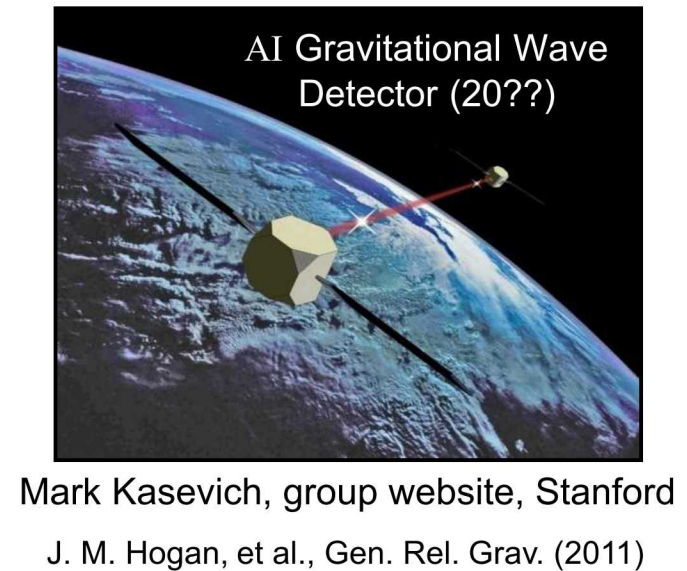
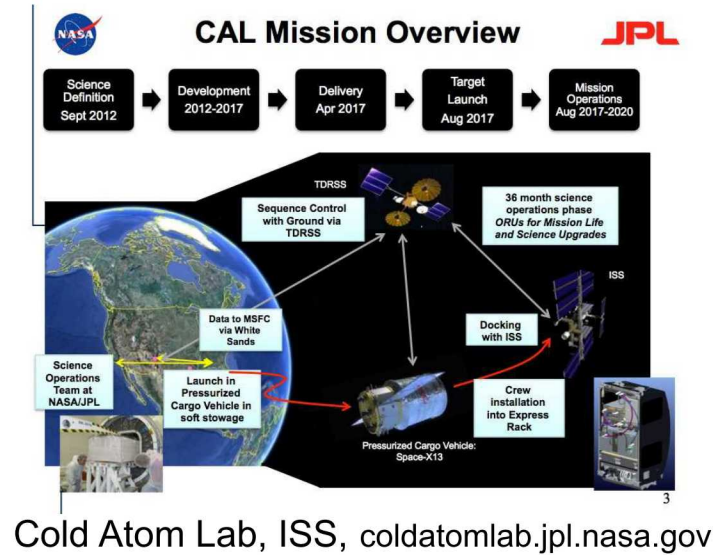
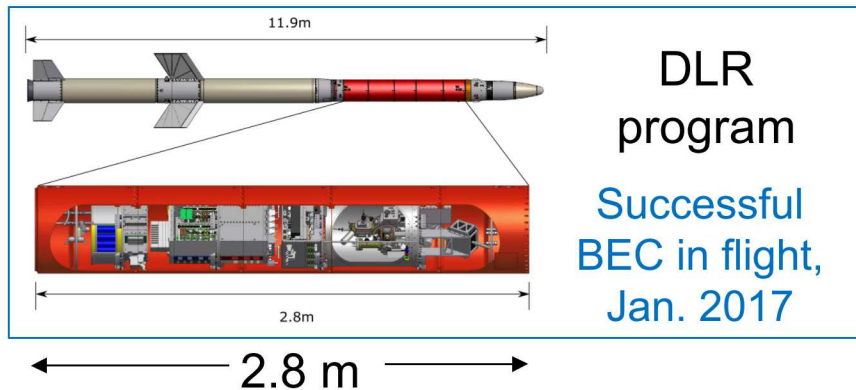
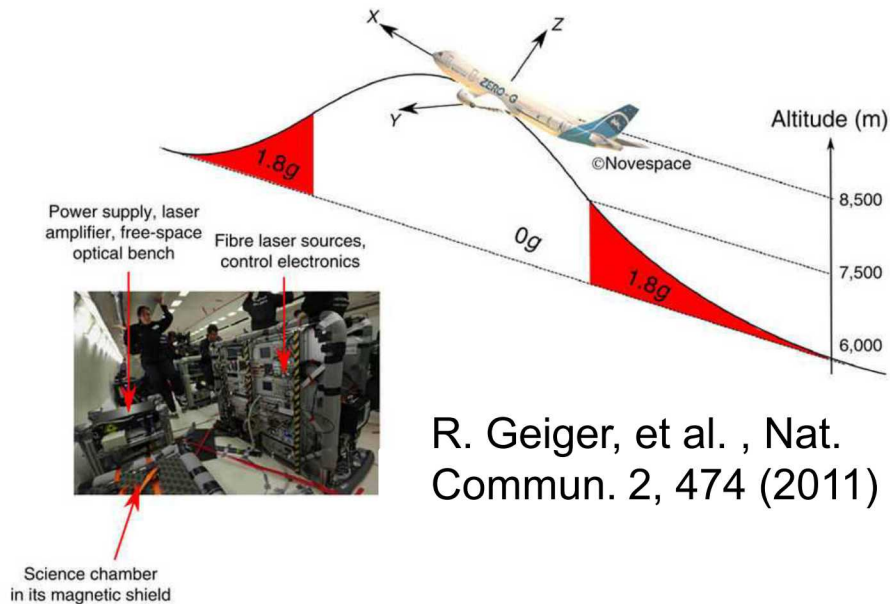
Biedermann, Ph.D. Thesis, Stanford, 2007



*electronics 9"x15"*

Wu, Ph.D. Thesis, Stanford, 2009

# Micro-g atom interferometers



# State-of-the-art Atom Interferometer instrumentation



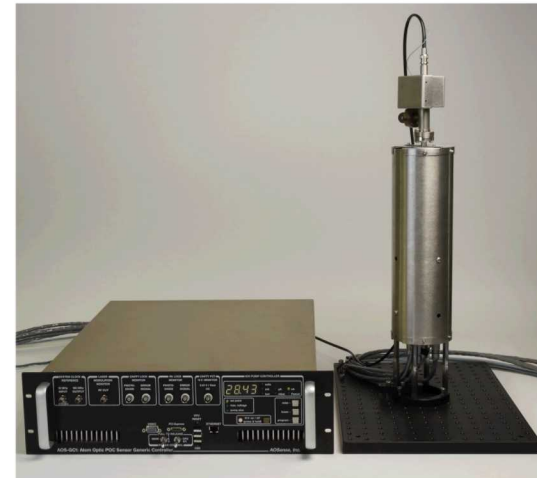
100 kg, 300 W, 50 ng/√Hz



U.K. University of Birmingham  
gravity imager



NASA/JPL ISS CAL



AOSense gravimeter

# Inertial Navigation

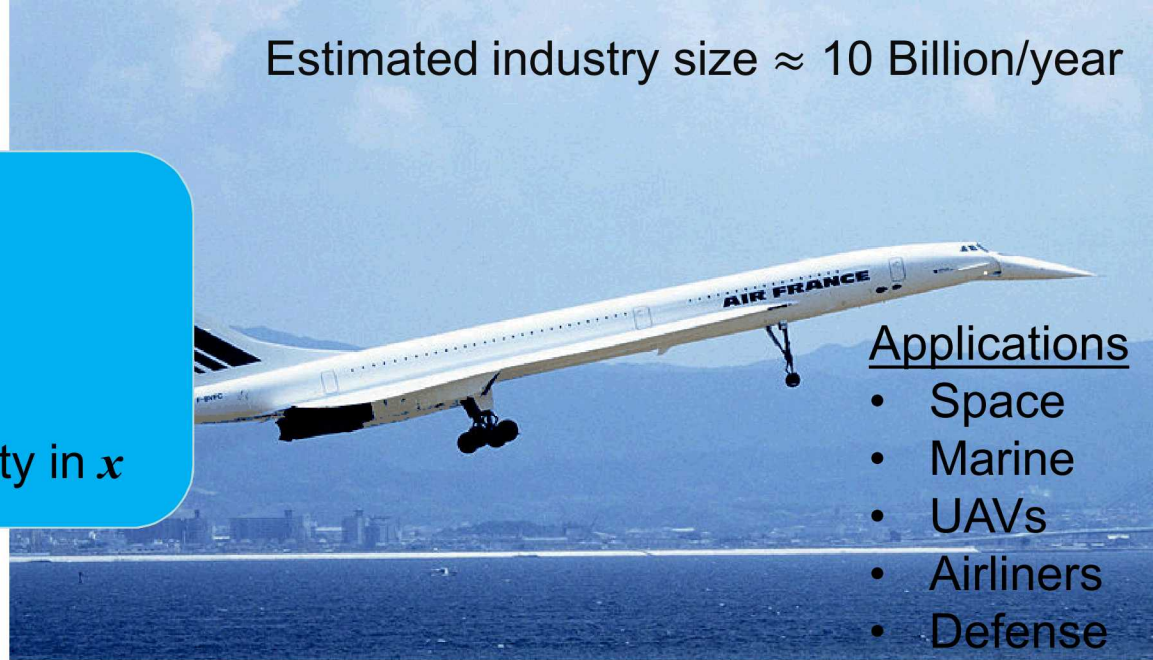
Estimated industry size  $\approx 10$  Billion/year

Basic idea:

$$x(t) = \iint a(t) dt$$

Uncertainty in  $a$  drives uncertainty in  $x$

Gyroscope gives  
heading information



## Applications

- Space
- Marine
- UAVs
- Airliners
- Defense

Concorde: Wikipedia commons

Tactical, MEMS

Navigation, HRG

Strategic, ESG

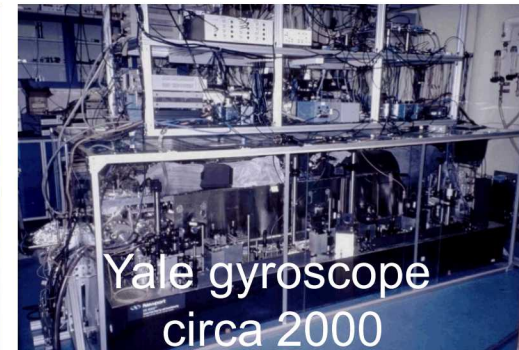
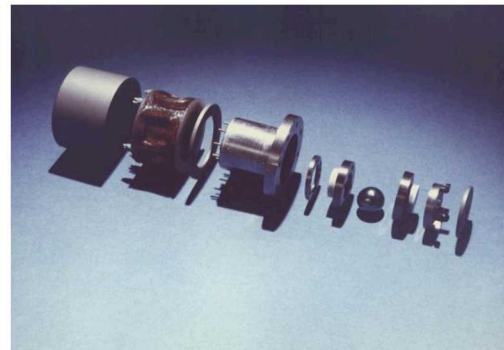
Atomic



Analog Devices  
ADIS16385



Honeywell  
HG9900



Yale gyroscope  
circa 2000

stability

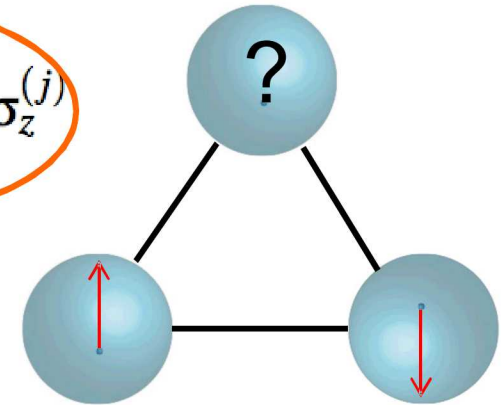
# Applications

Quantum simulation

Single spin coherence

$$H_P = \sum_{i=1}^N \tilde{h}_i \sigma_z^{(i)} + \sum_{i,j=1}^N \tilde{J}_{ij} \sigma_z^{(i)} \otimes \sigma_z^{(j)}$$

Entangling interactions



Frustrated magnetism

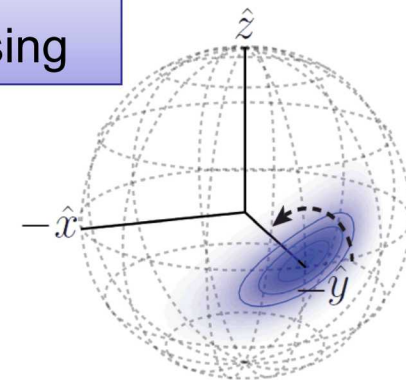
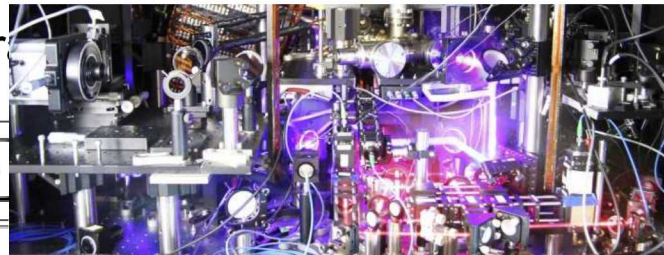
Large-scale/rapid entanglement for sensing



Sr lattice clock, Ye lab JILA

Quantum Fourier transform

$$QFT : |x\rangle \mapsto \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{2\pi i x k / N} |k\rangle$$



# AI phase shift: a simple perspective

Use atom to stroboscopically measure lateral position (optical phase) at three equal-spaced points in time

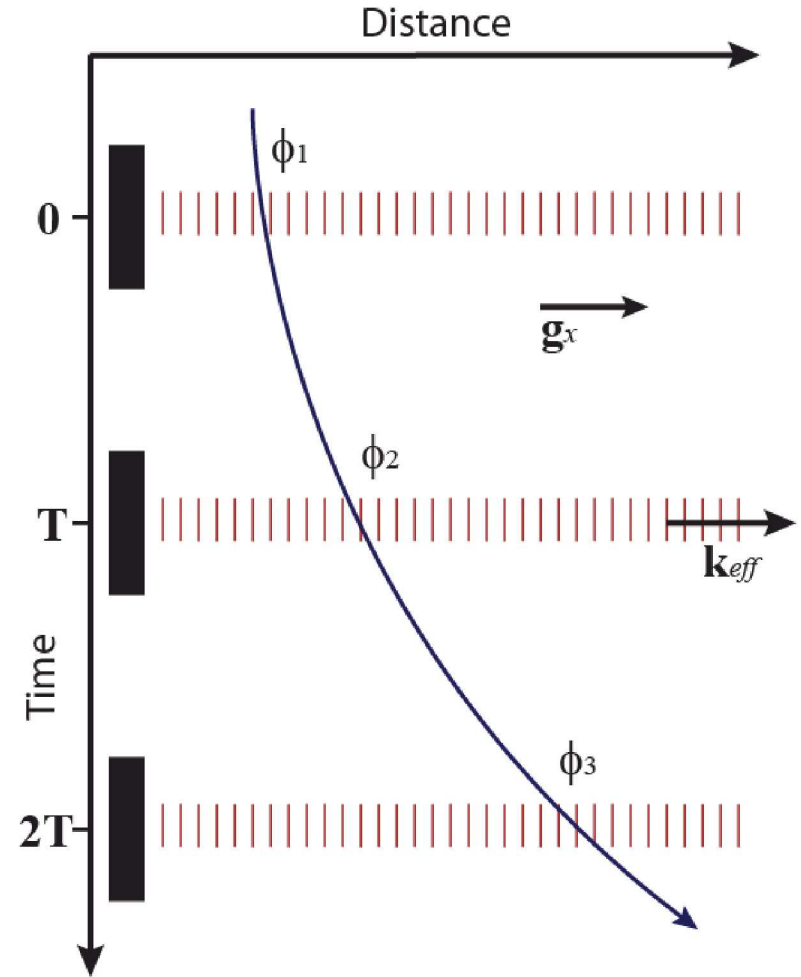
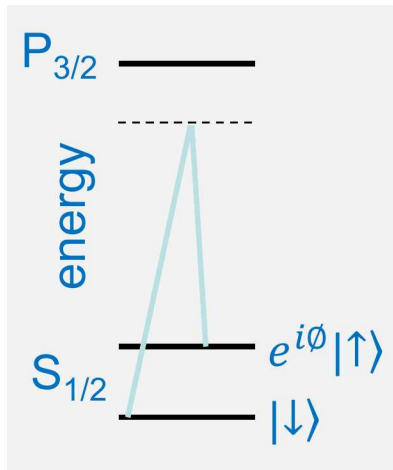
$$\varphi_i = \vec{k}_{eff} \cdot \vec{x}_i$$

calculate curvature via finite difference method

$$\Delta\varphi = \varphi_1 - 2\varphi_2 + \varphi_3$$

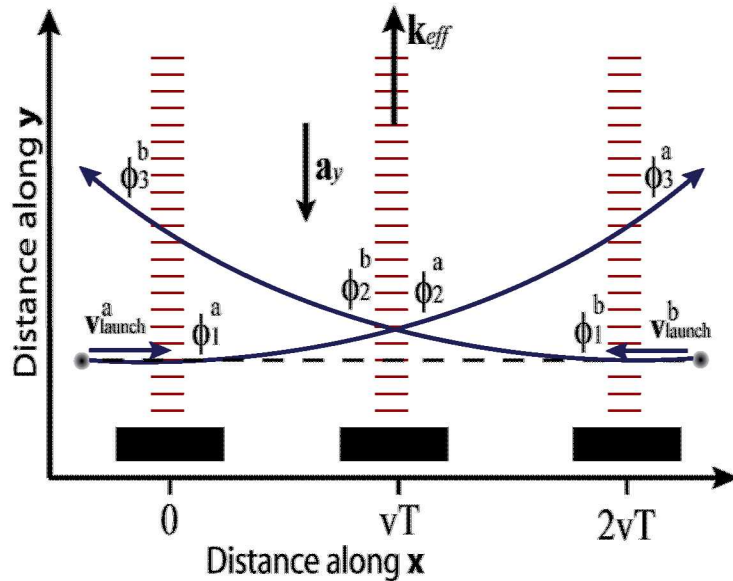
...and then

$$g_x = \frac{\Delta\varphi}{k_{eff} T^2}$$

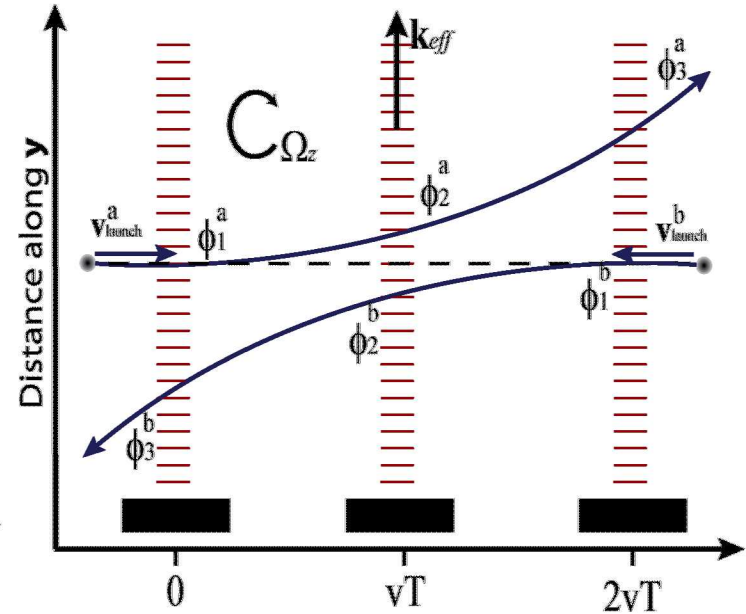


$$P_{|1\rangle} = \frac{1}{2}(1 - \cos(\Delta\phi + \phi_0)) \quad \dots S_{1/2} \text{ states behave as a pristine oscillator with phase memory}$$

# Simultaneously measure rotation and acceleration



Platform acceleration

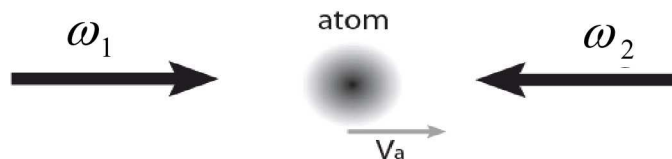


Platform rotation

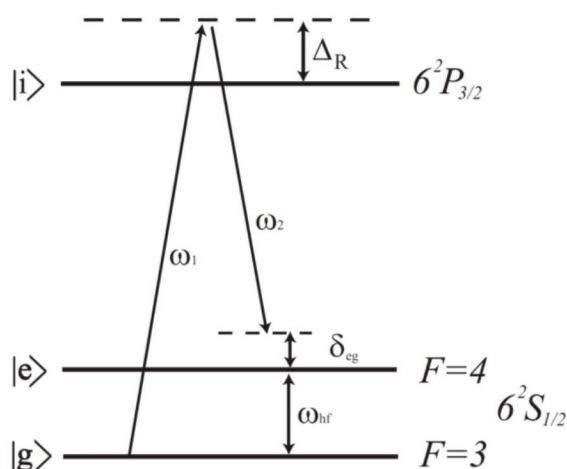
$$\Delta\phi = \vec{k}_{eff} \cdot (\vec{a}T^2 - 2(\vec{v} \times \vec{\Omega})T^2)$$

# Stimulated Raman transition

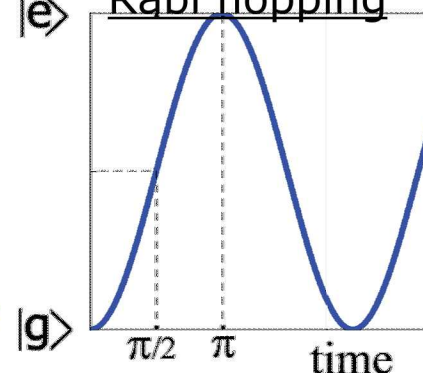
## Laser configuration



## Relevant energy diagram (Cs)



## Rabi flopping



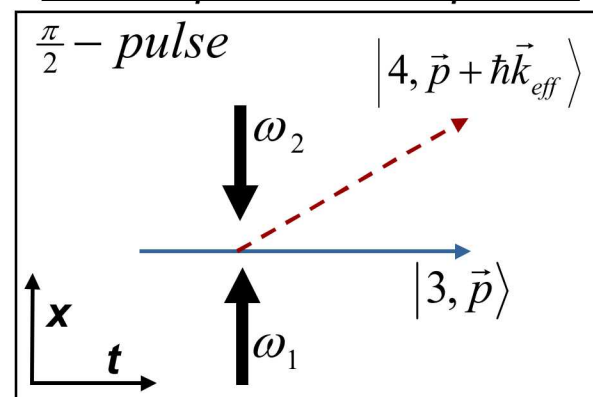
## Effective 2-level system

$$H'_{\text{int}} = -\vec{\mu}_{\text{eff}} \cdot \vec{E}_{\text{eff}} e^{i(\vec{k}_{\text{eff}} \cdot \vec{x} - \omega_{\text{hf}} t)}$$

$$\vec{k}_{\text{eff}} = \vec{k}_1 - \vec{k}_2 \approx 2\vec{k}_1$$

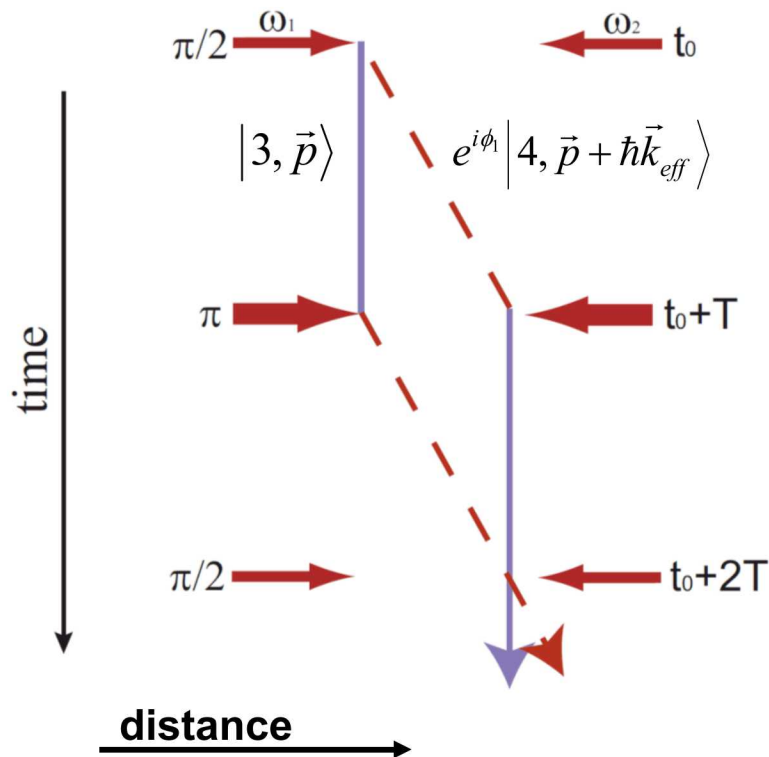
Valid for  $\Omega_{\text{eff}} \ll \Delta$

## Atom optics beamsplitter



# AI phase shift

## Interferometer Recoil diagram



## Transition rules

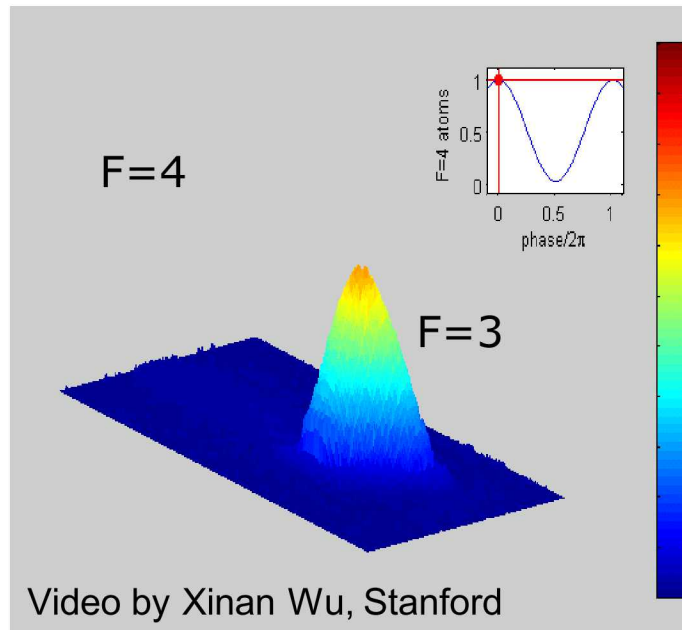
$$\left. \begin{aligned} |3, \vec{p}\rangle &\rightarrow e^{i\phi} |4, \vec{p} + \hbar \vec{k}_{eff}\rangle \\ |4, \vec{p} + \hbar \vec{k}_{eff}\rangle &\rightarrow e^{-i\phi} |3, \vec{p}\rangle \end{aligned} \right\} \begin{aligned} \Delta\phi &= \phi_1 - 2\phi_2 + \phi_3 \\ &= \vec{k}_{eff} \cdot (\vec{a}T^2 - 2(\vec{v} \times \vec{\Omega})T^2) \end{aligned}$$

$$\phi = \vec{k}_{eff} \cdot \vec{x}$$

Interferometer  
transition probability

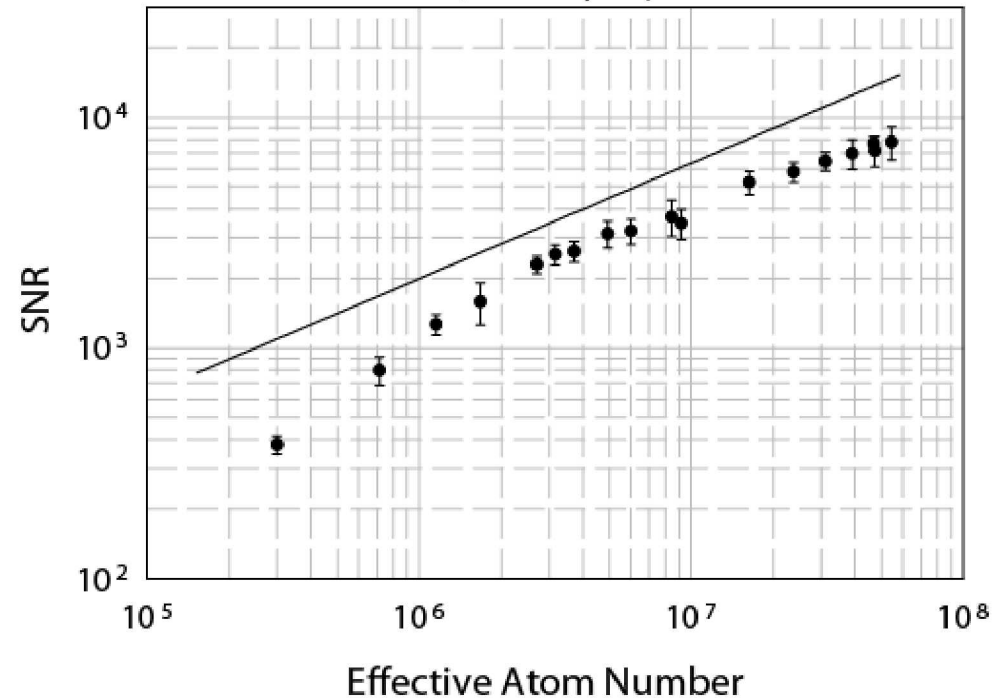
$$|\langle 4 | \Psi \rangle|^2 = \frac{1}{2}(1 - \cos \Delta\phi)$$

# State detection

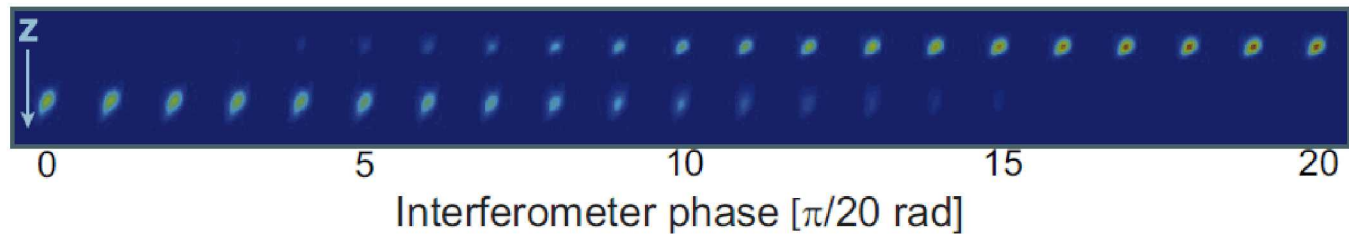


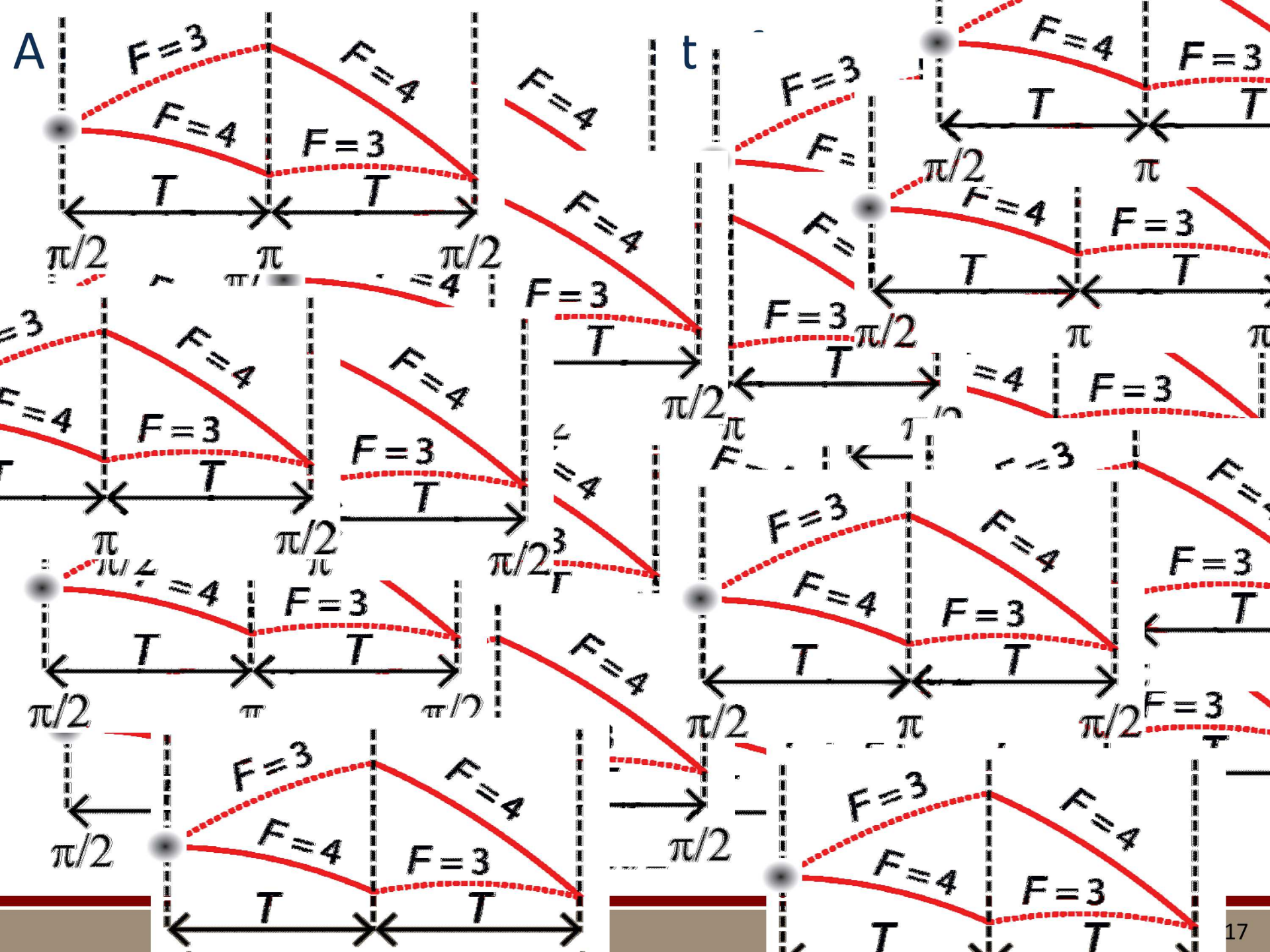
Radiation pressure spatially separates  $F=3$  and  $F=4$  atoms

Biedermann, et al., Opt. Lett. 2007

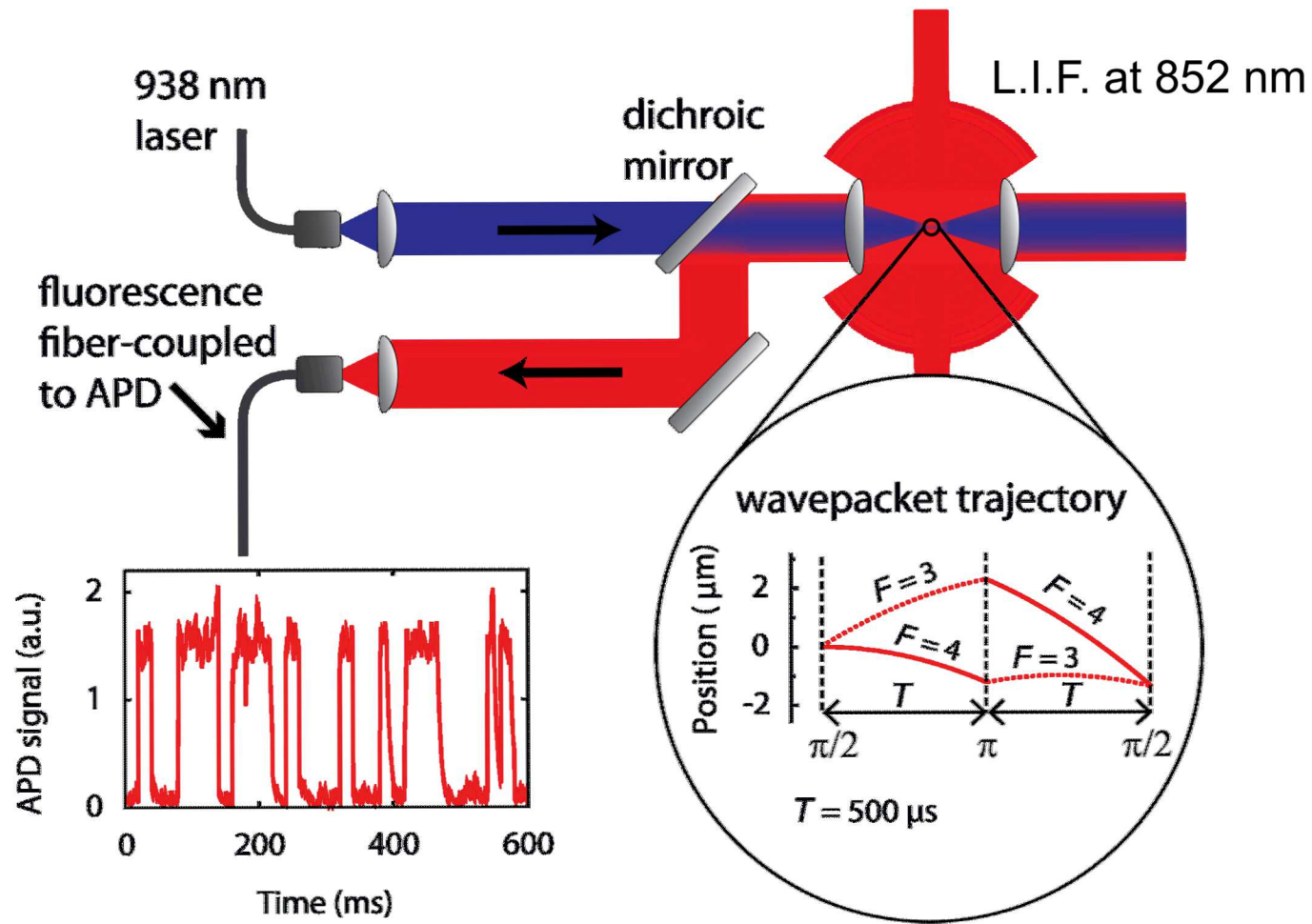


Achieves near atom shot-noise limited performance.

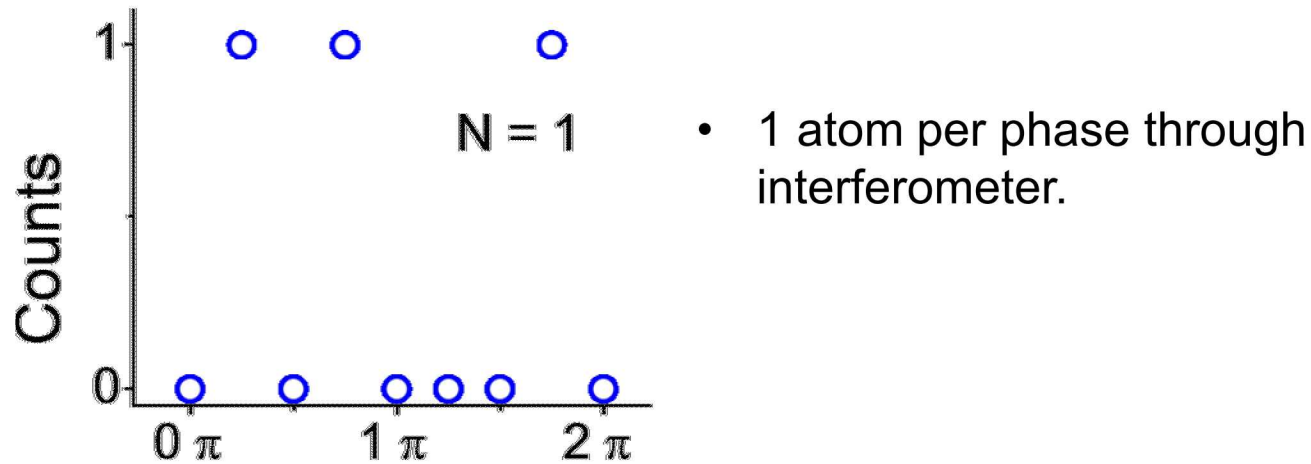




# Free-space matterwave interference with a single atom



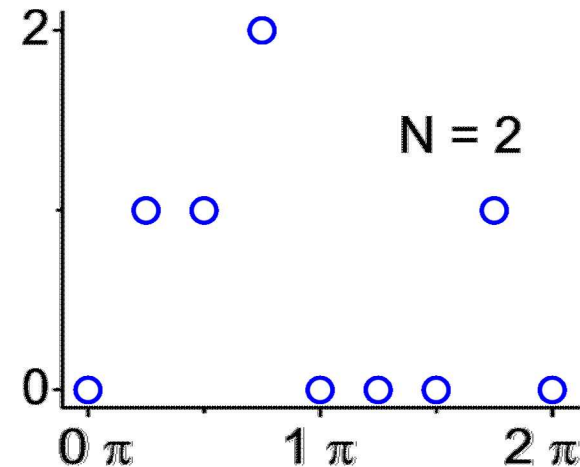
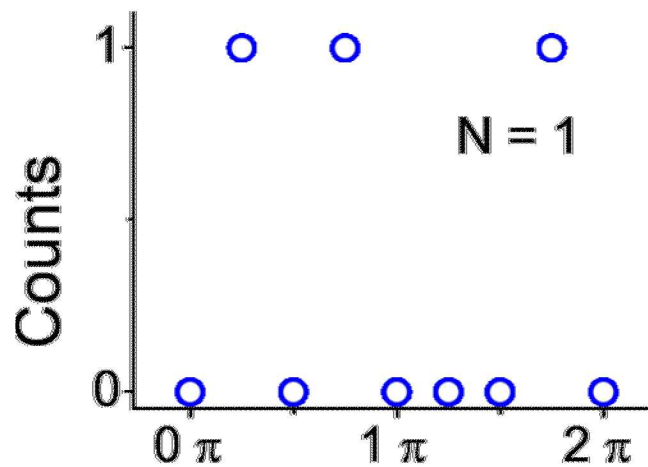
# Building fringe one atom at a time



$$H'_{\text{int}} = -\vec{\mu}_{\text{eff}} \cdot \vec{E}_{\text{eff}} e^{i(\vec{k}_{\text{eff}} \cdot \vec{x} - \omega_{\text{hf}} t)}$$

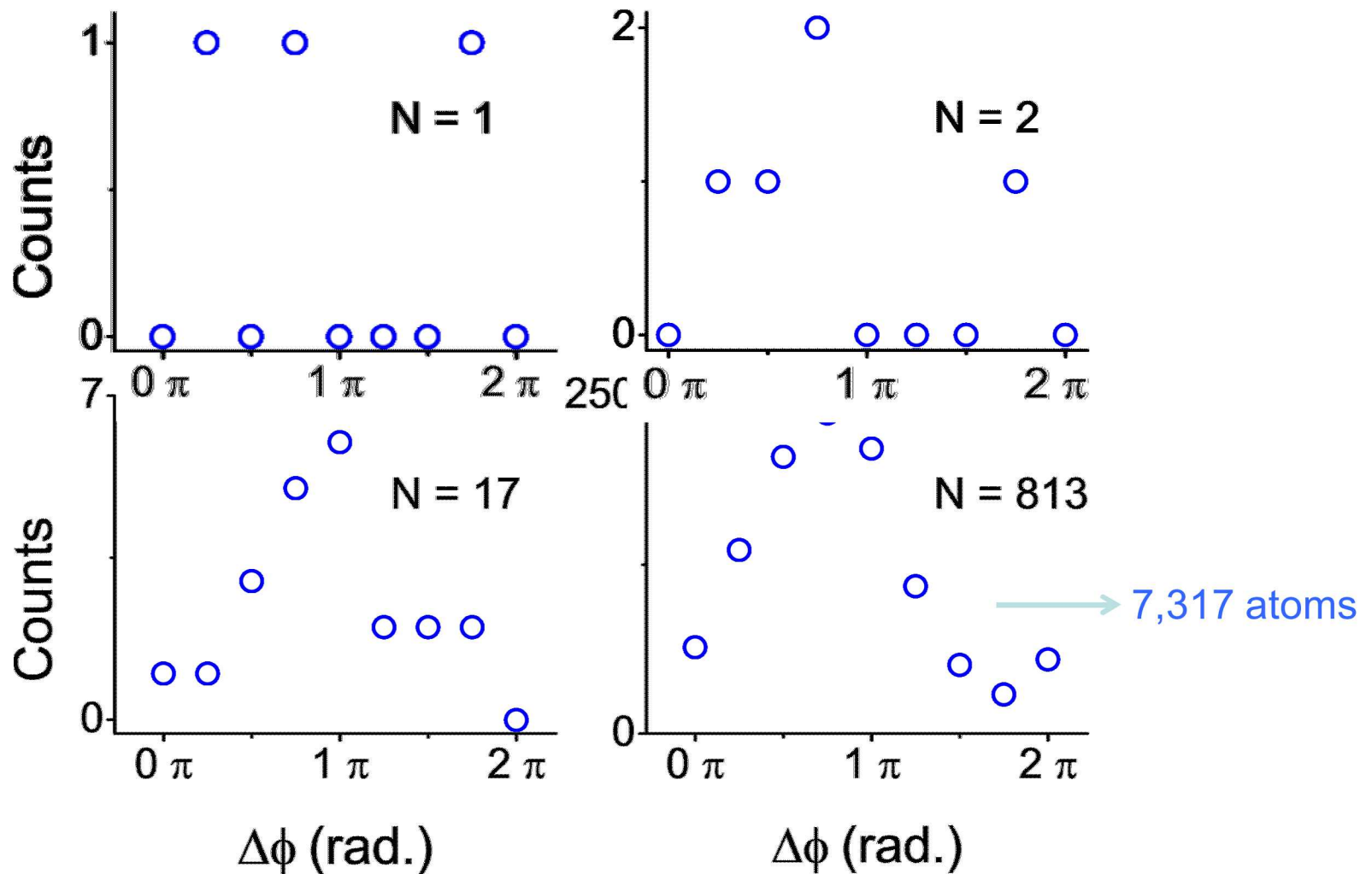
Scan laser phase

# Building fringe one atom at a time

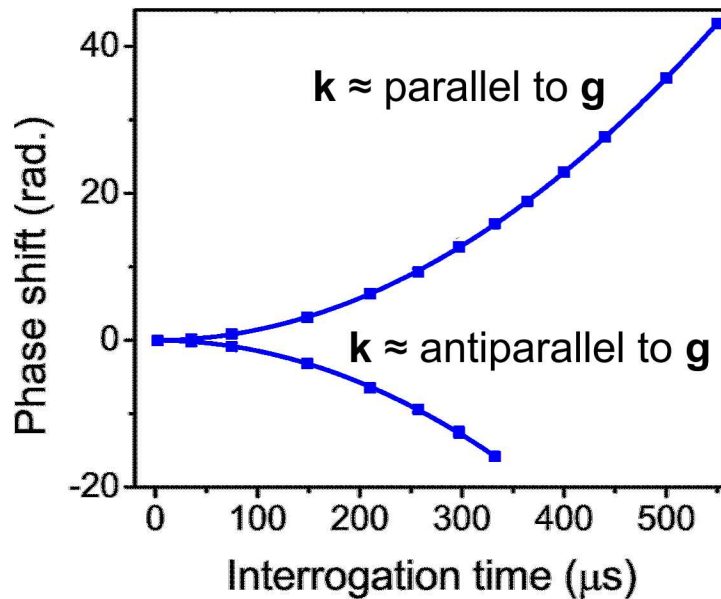
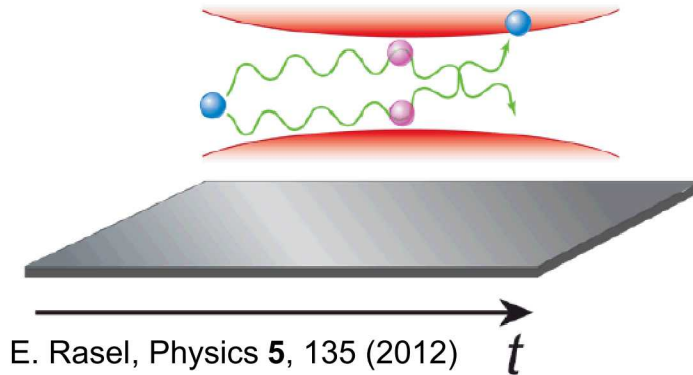


- 2 atoms per phase through interferometer

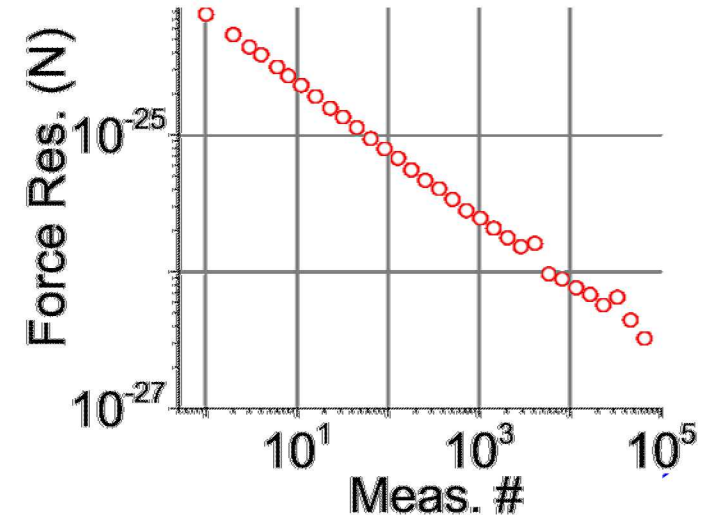
# Building fringe one atom at a time



# Force resolution of a single atom interferometer

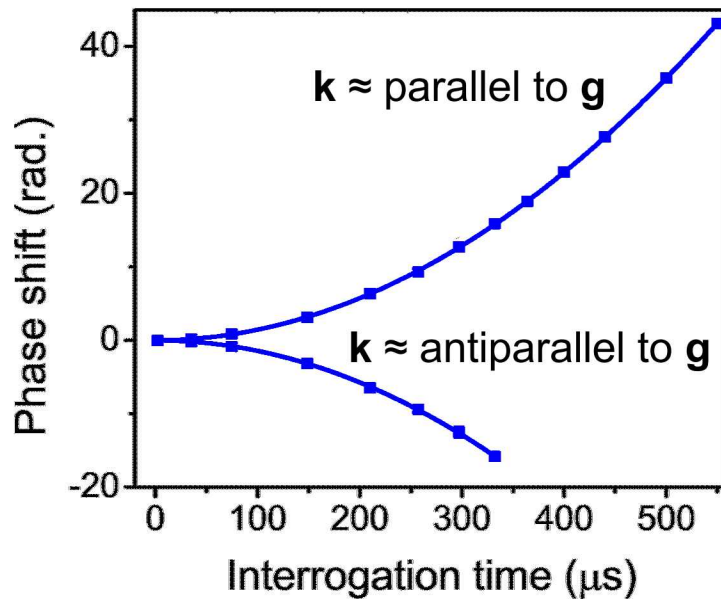
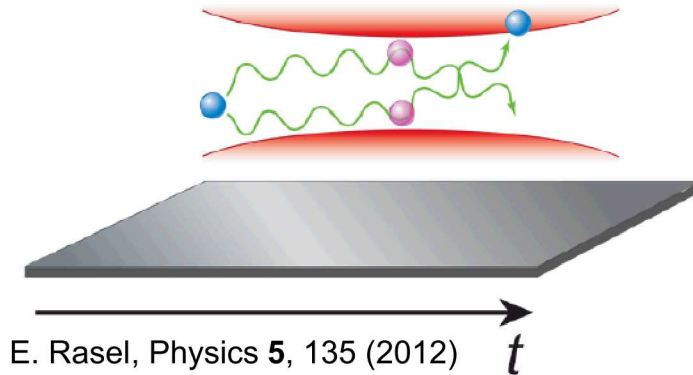


Parazzoli, et al., Phys. Rev. Lett. **109**, 230401 (2012)

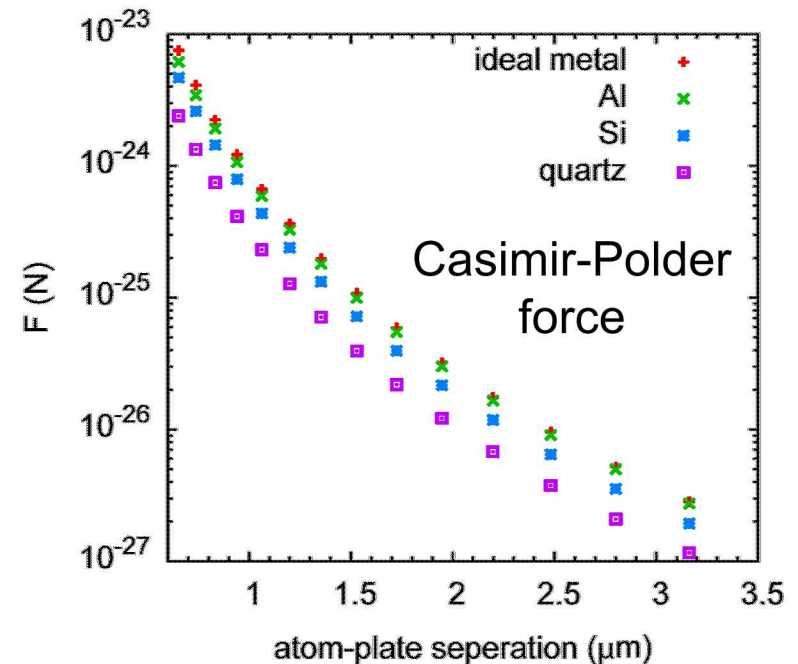
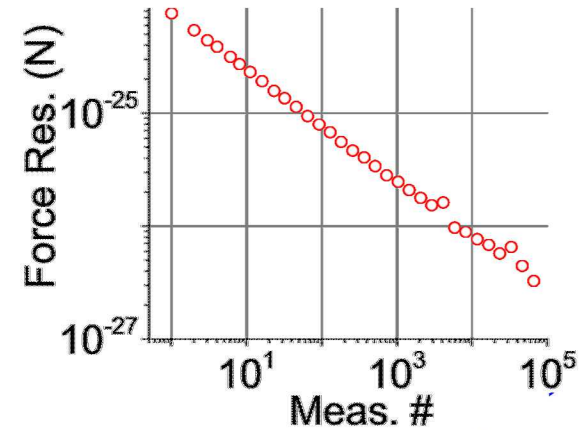


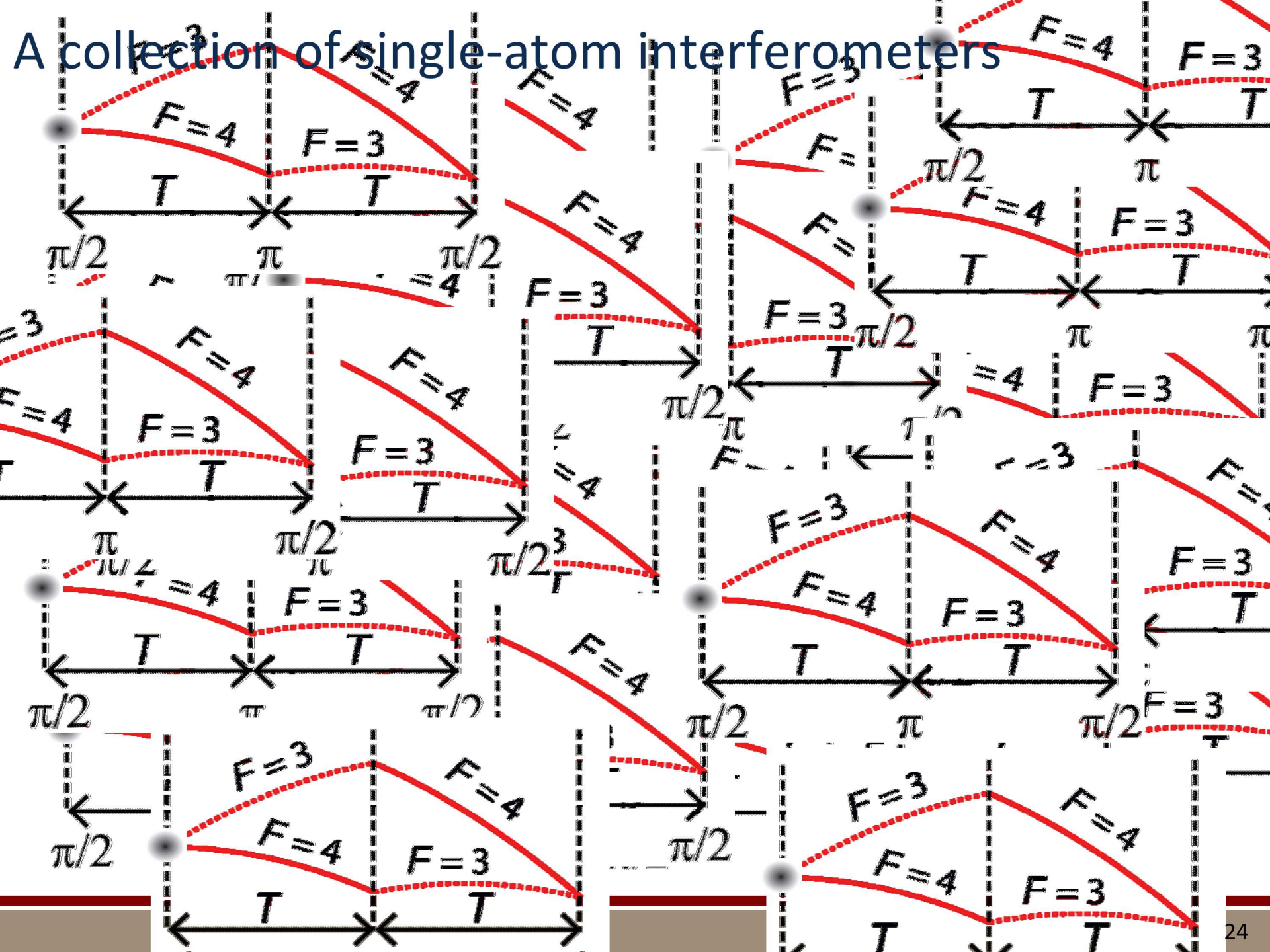
$$\Delta\phi = \vec{k}_{eff} \cdot \vec{g} T^2$$

# Force resolution of a single atom interferometer



Parazzoli, et al., Phys. Rev. Lett. **109**, 230401 (2012)

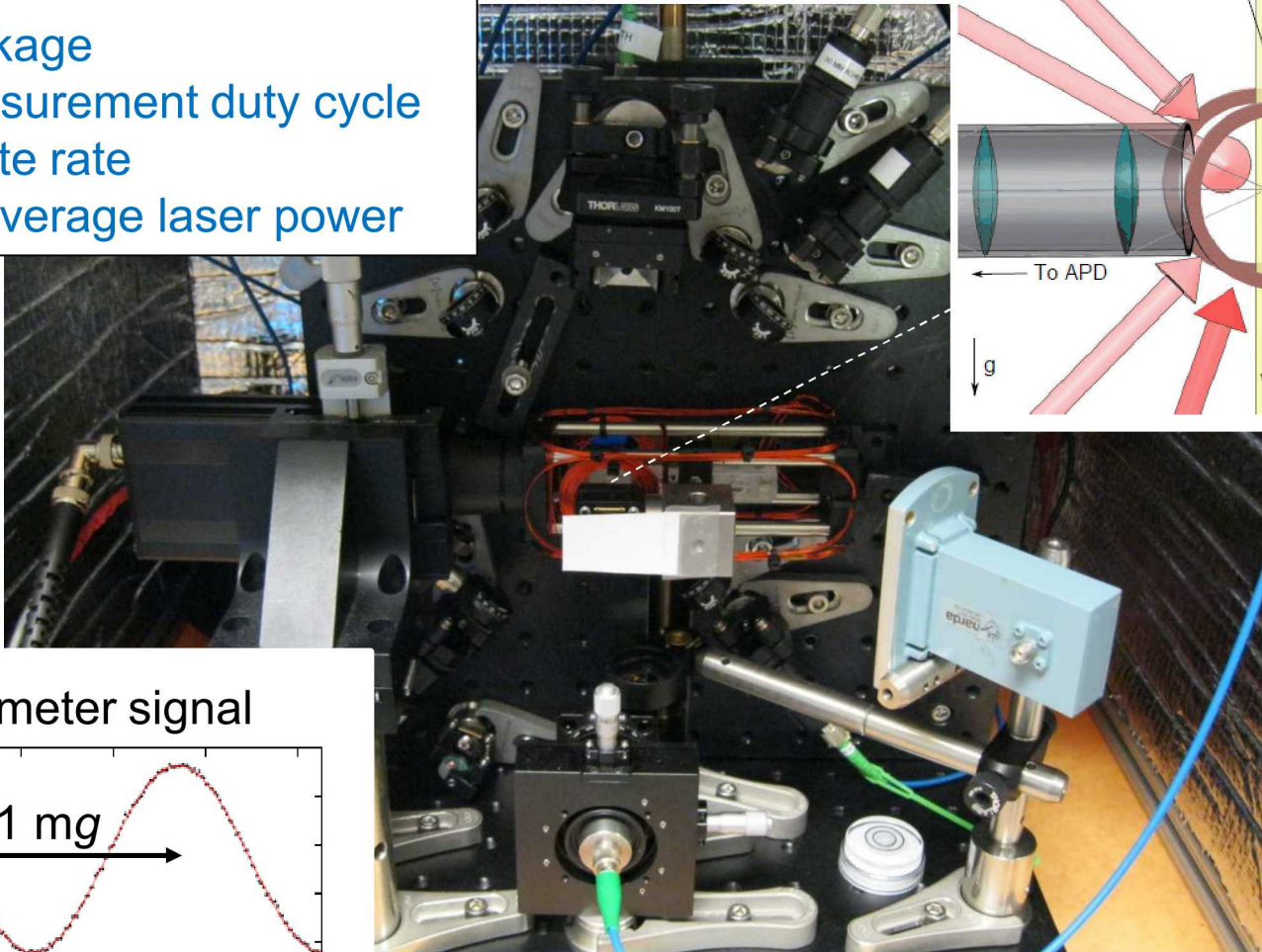
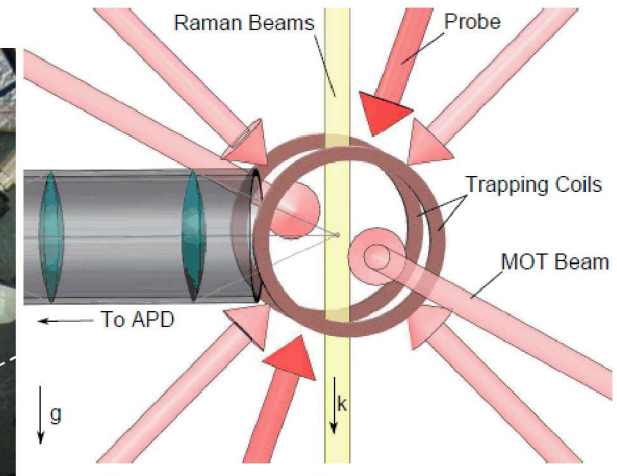




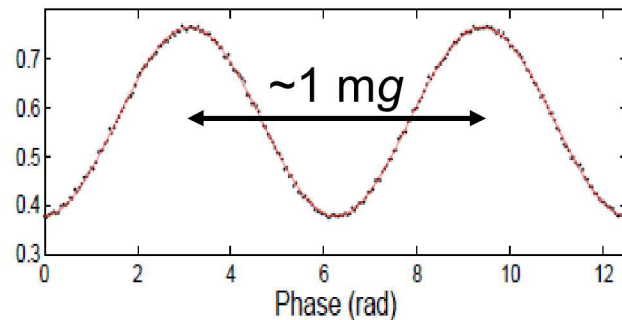
# High-bandwidth interferometer

## Key features

- small package
- large measurement duty cycle
- high update rate
- reduced average laser power



Interferometer signal

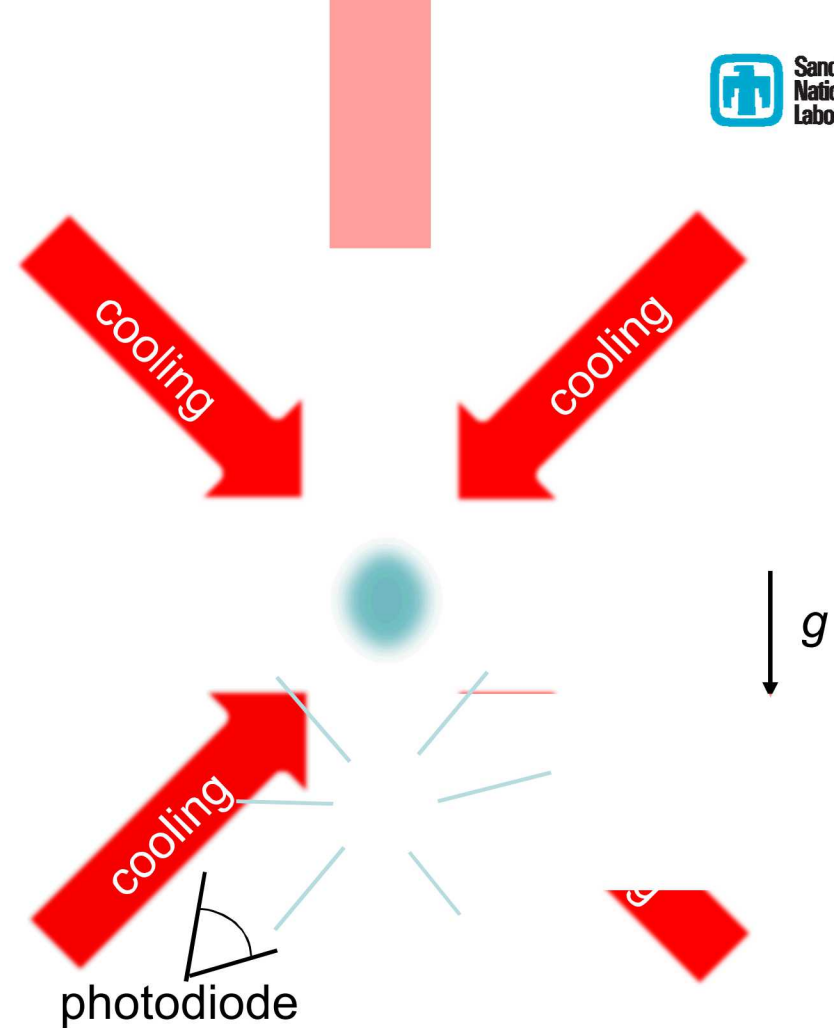


$$600 \text{ ng}/\sqrt{\text{Hz}} \text{ at } 50 \text{ Hz}$$

# One measurement cycle

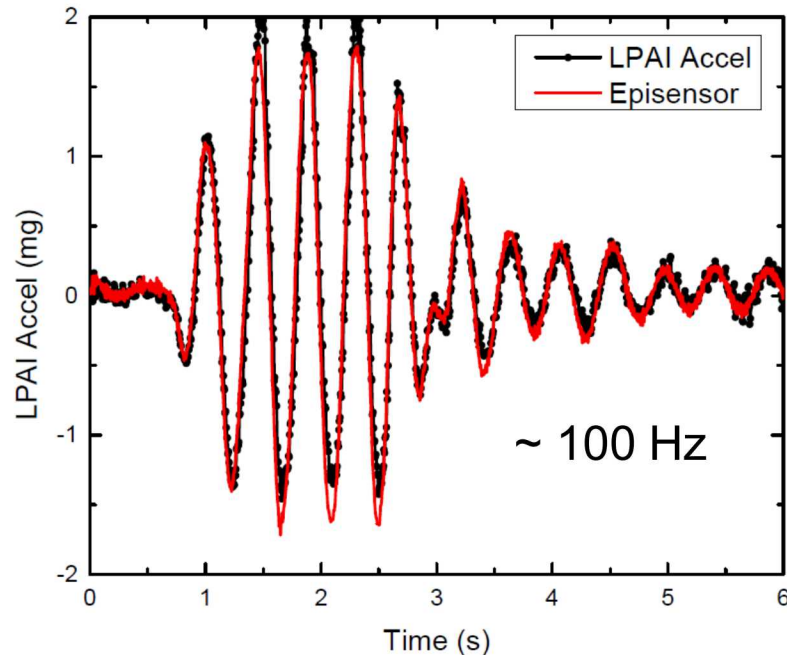
*Example,  $(50 \text{ Hz})^{-1}$  cycle:*

- Laser cool  $10^6$  atoms (4.3 ms)
  - $T \approx 5 \mu\text{K}$
- Release atoms
- Raman pulse sequence (14 ms,  $T = 7 \text{ ms}$ )
- Detect
- Recapture (1.7 ms)

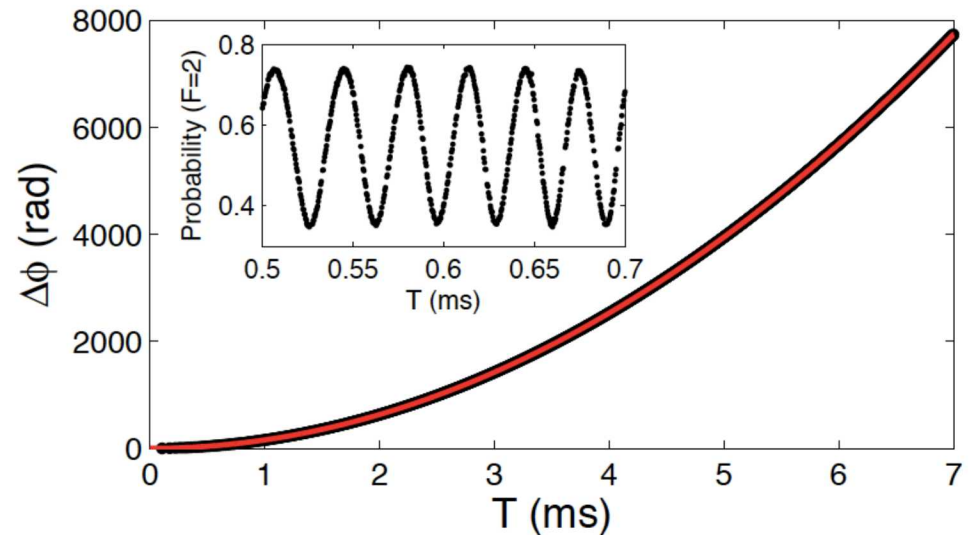


# High bandwidth interferometer results

Operating under dynamics



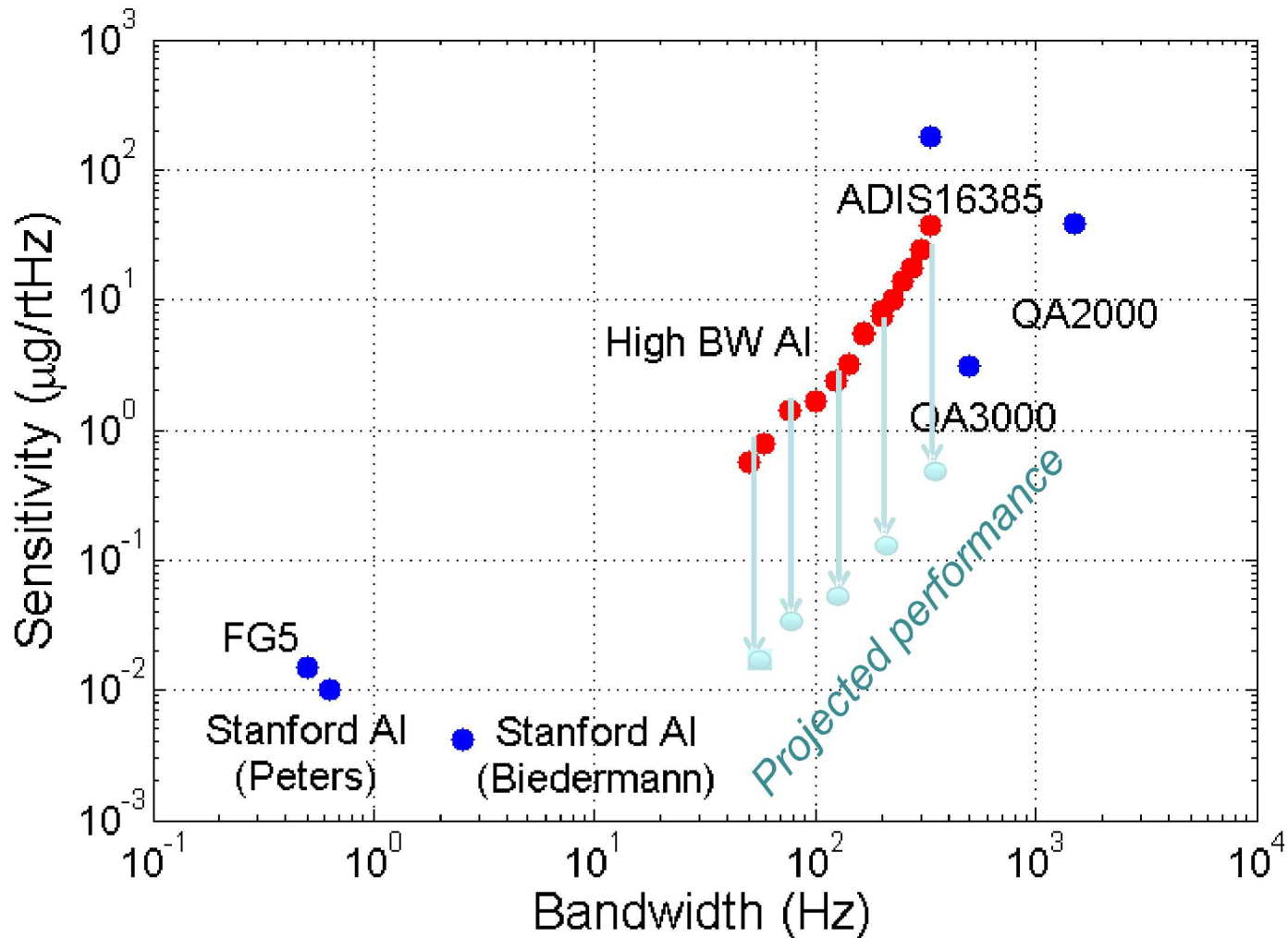
Measurement of  $g$



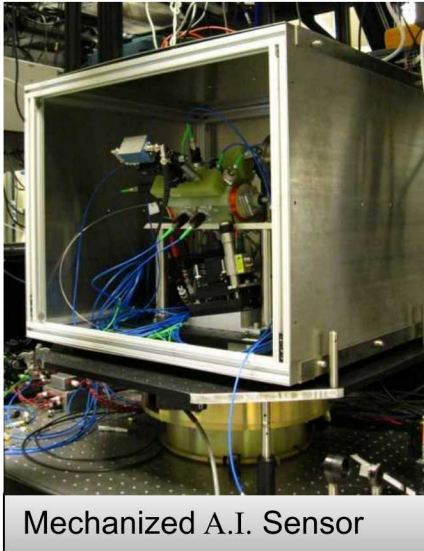
$$\Delta\phi = k_{eff} \cdot aT^2$$

$$P_e(\Delta\phi) = \frac{1}{2} (1 - \cos(\Delta\phi))$$

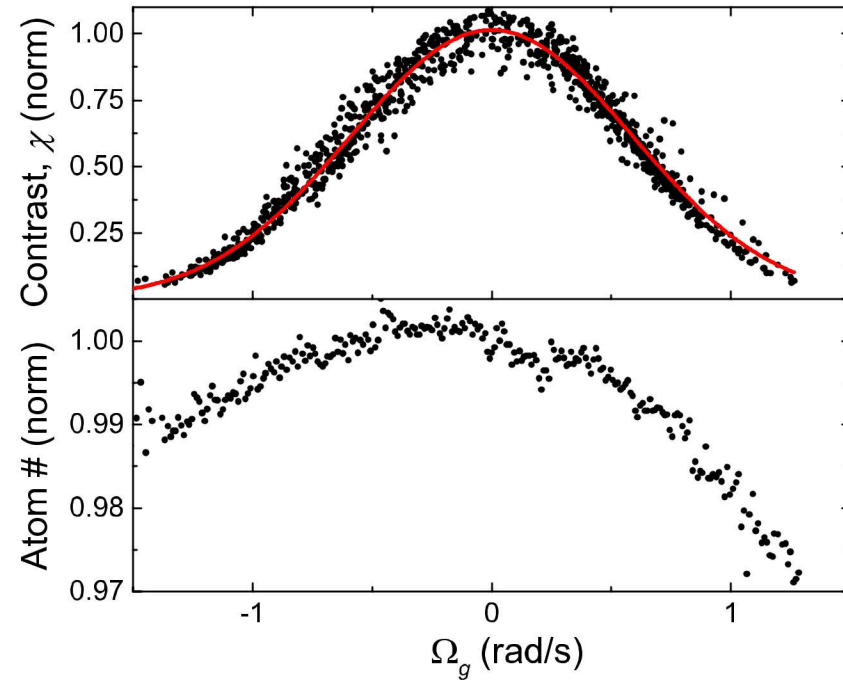
# Interferometer comparison



# Rotation rate limit



Mechanized A.I. Sensor



$$\Delta\phi = \vec{k}_{eff} \cdot (\vec{a}T^2 - 2(\vec{v} \times \vec{\Omega})T^2)$$

Limited by temperature  
&  
Wavepacket overlap

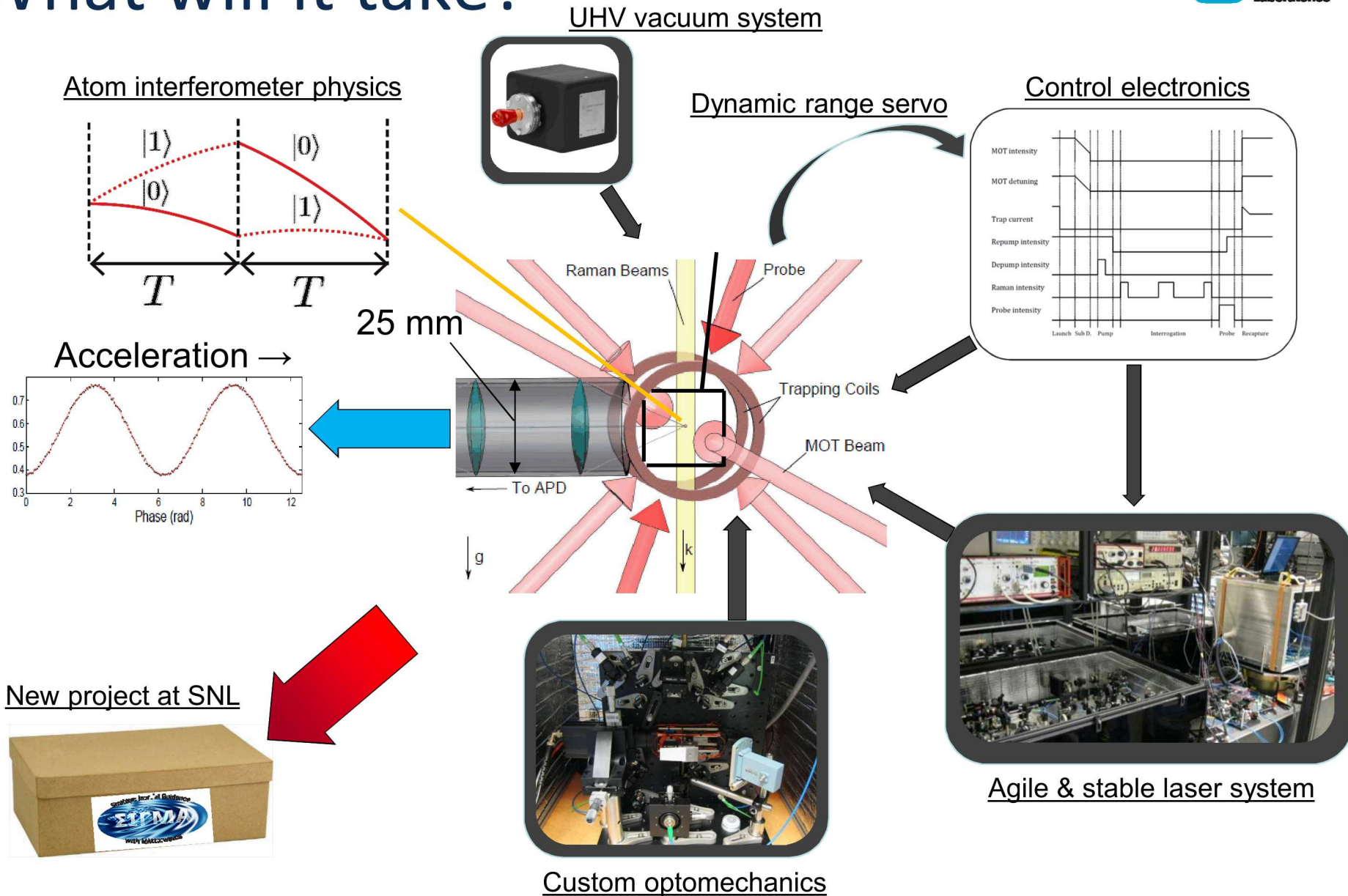
Contrast  
envelope

$$\chi(\Omega_z) = \exp\left(\frac{-\Omega_z^2}{2\sigma_\Omega^2}\right)$$

Envelope  
width

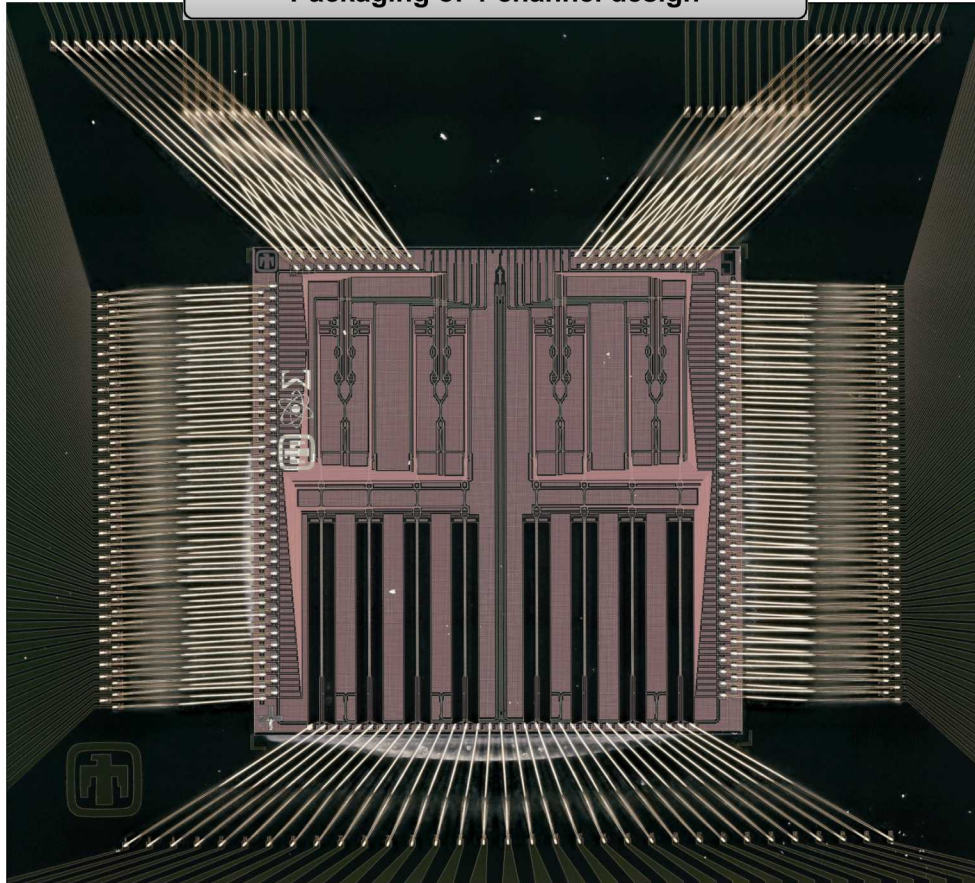
$$\sigma_\Omega = \frac{1}{2} \sqrt{\frac{m}{kT}} \frac{1}{k_e T^2 \sin \theta}$$

# What will it take?

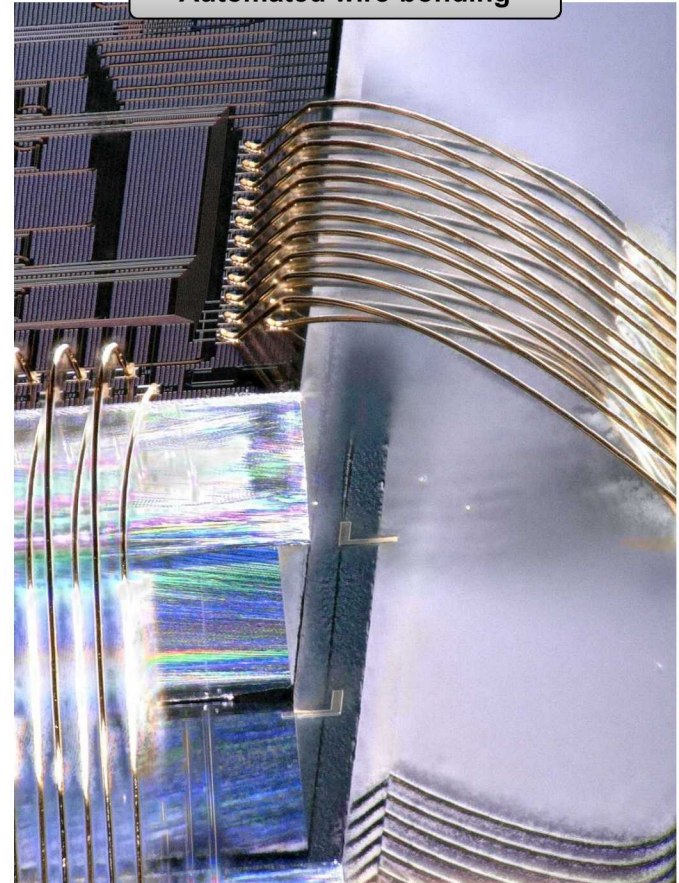


# A multi-channel laser system front end

Packaging of 4-channel design



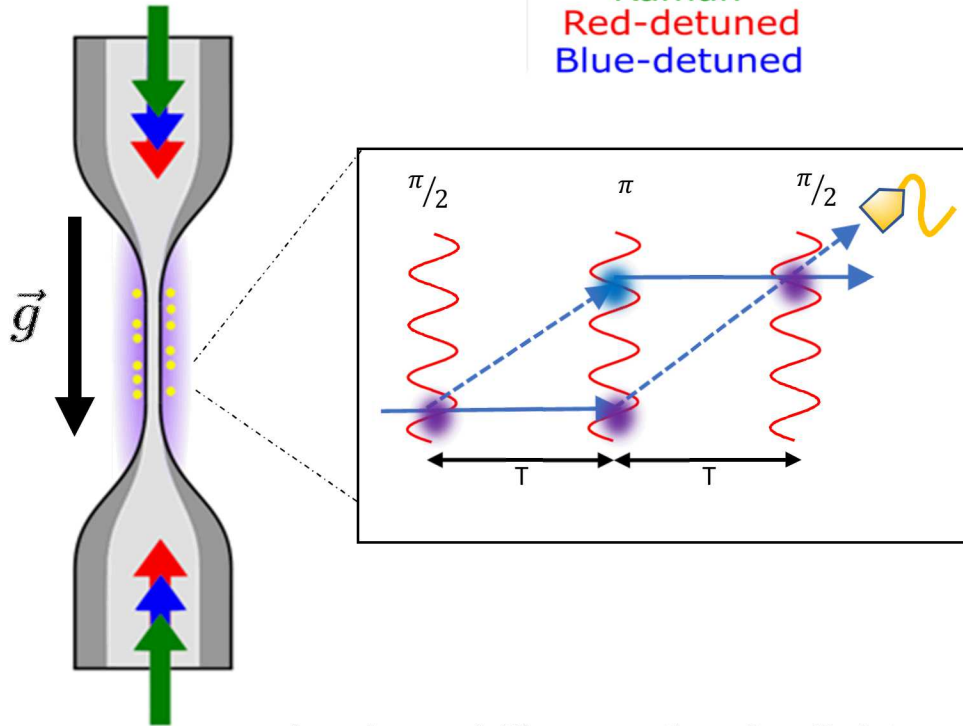
Automated wire-bonding



# A nanofiber guide for matterwaves

Atom Interferometer guided via tapered optical fiber (TOF)

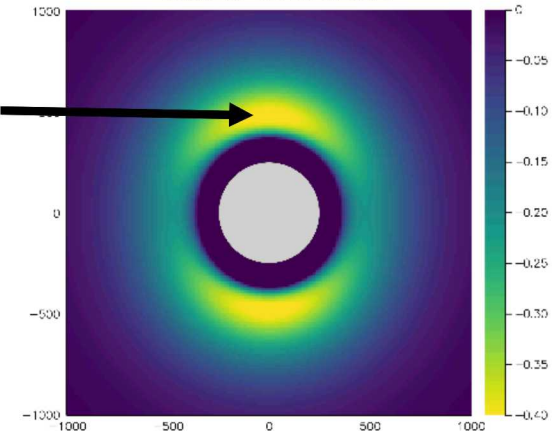
Raman  
Red-detuned  
Blue-detuned



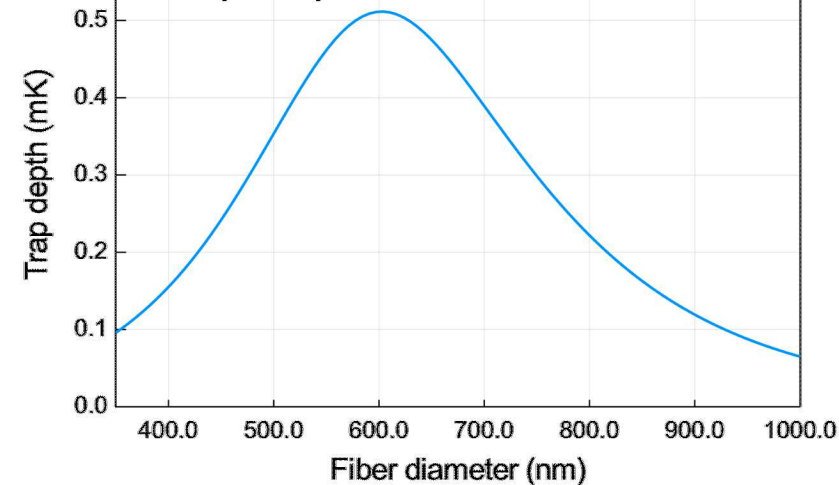
- Analog of fiber optics for light
- Guide atoms near surfaces
  - Casimir-Polder effect
  - Constrain 5<sup>th</sup> force theories

Calculated intensity distribution

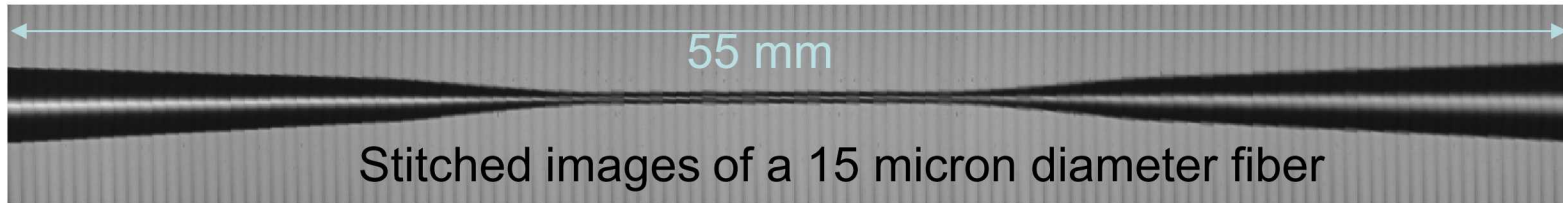
Atoms guided here



Trap depth vs. fiber diameter

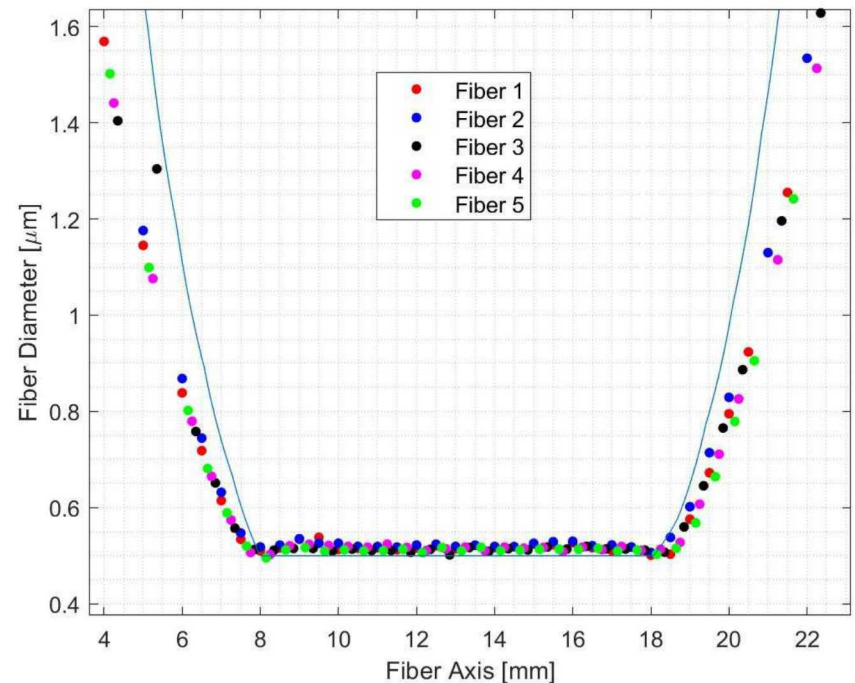
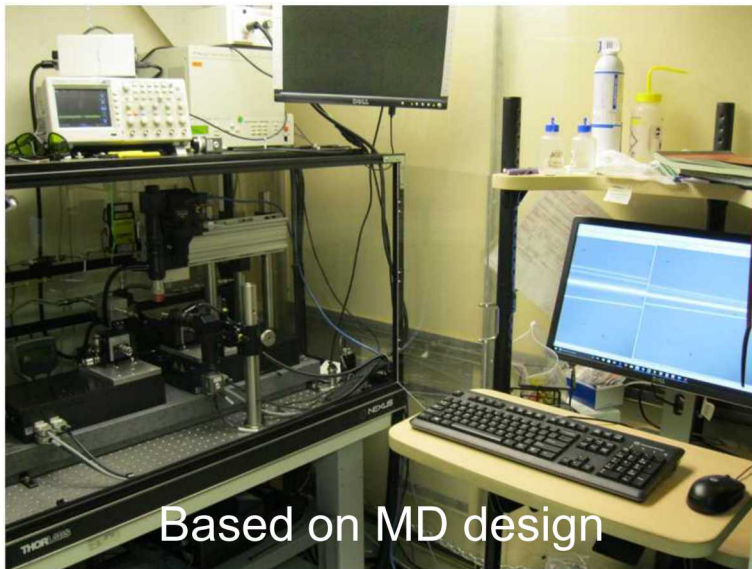


# A nanofiber guide for matterwaves

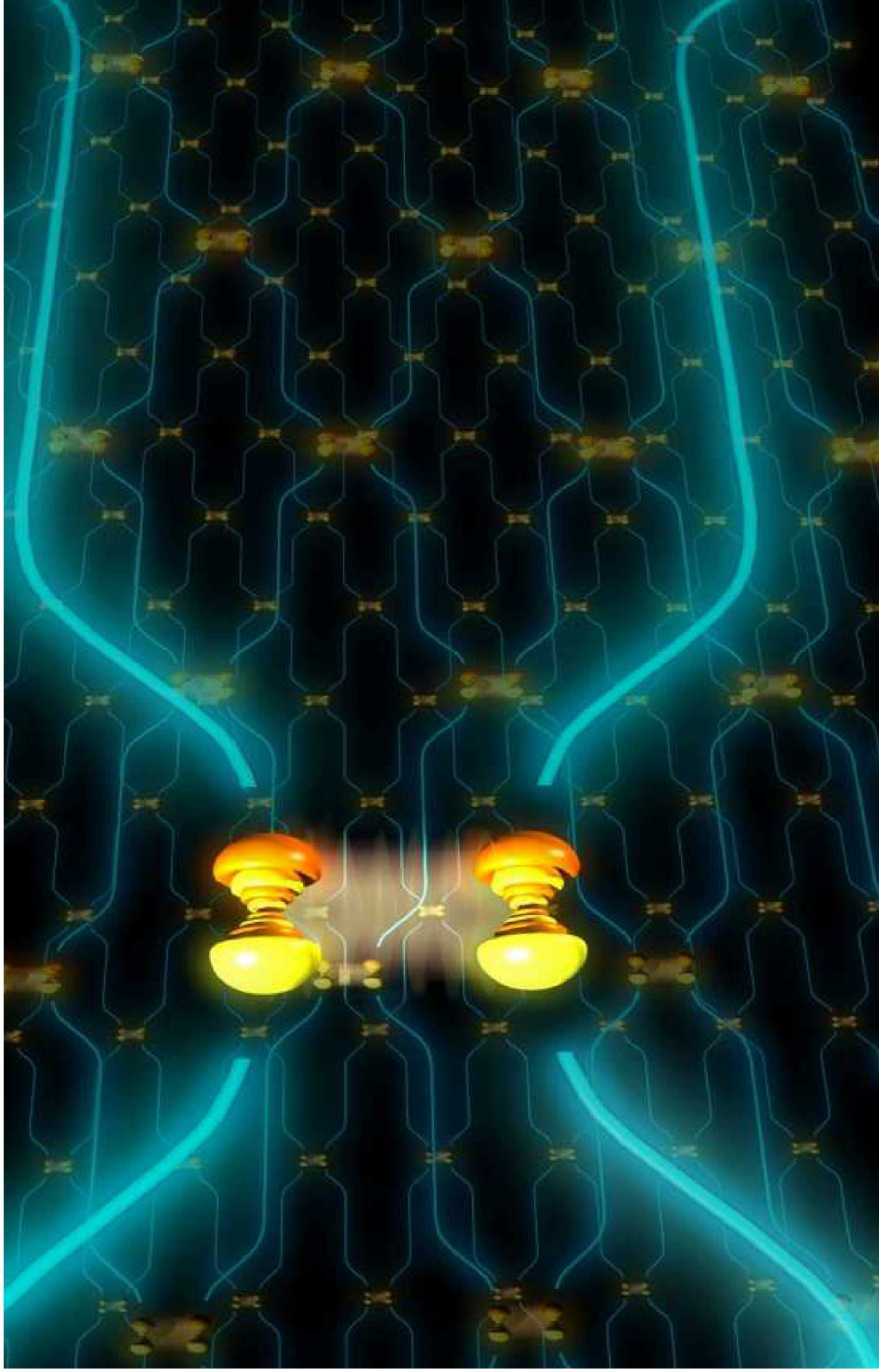


Stitched images of a TOF

## Sandia nanofiber rig—NEW

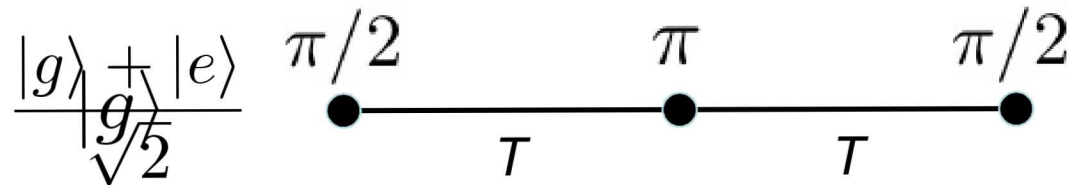


# Entangled states of neutral atoms



# AI operated with entangled atoms

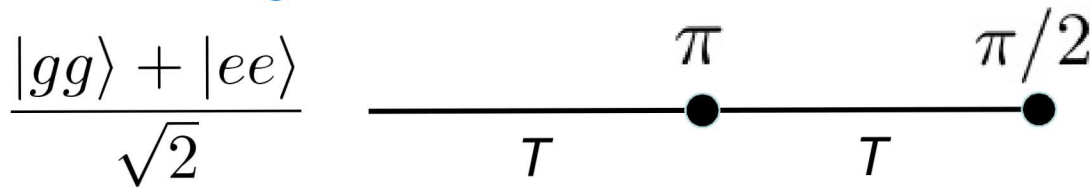
## Single atom



$$\langle \Pi \rangle = \cos \Theta_1$$

$$\Theta_1 = \mathbf{K} \cdot \mathbf{a} T^2$$

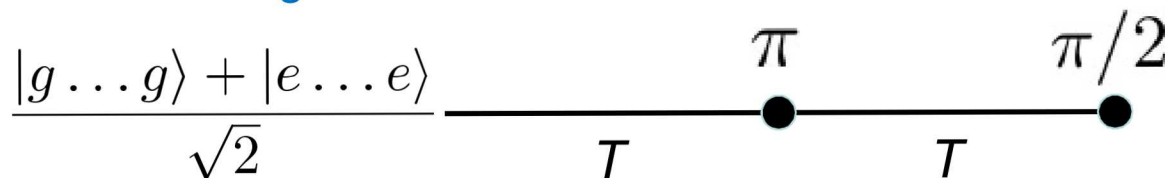
## Two entangled atoms



$$\langle \Pi \rangle = \cos \Theta_2$$

$$\Theta_2 = \mathbf{K} \cdot (\mathbf{a}_1 + \mathbf{a}_2) T^2$$

## N entangled atoms



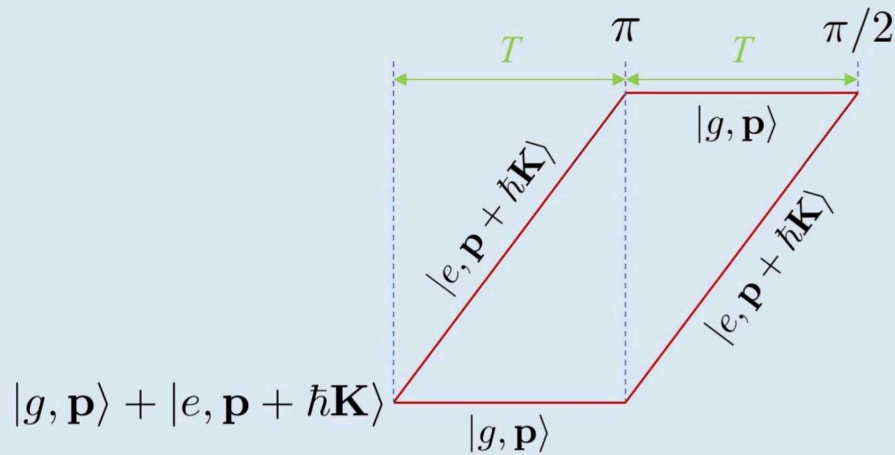
$$\langle \Pi \rangle = \cos \Theta_N$$

$$\Theta_N = \mathbf{K} \cdot \sum_{\alpha=1}^N \mathbf{a}_{\alpha} T^2$$

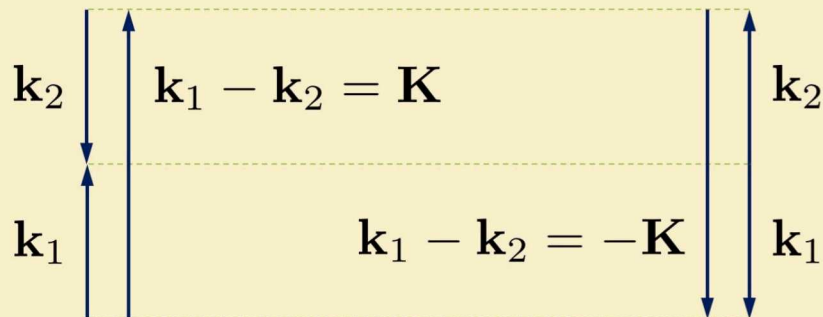
$$\langle \Pi \rangle = (-1)^{N_{\uparrow}}$$

# New pulse sequences for AI operation

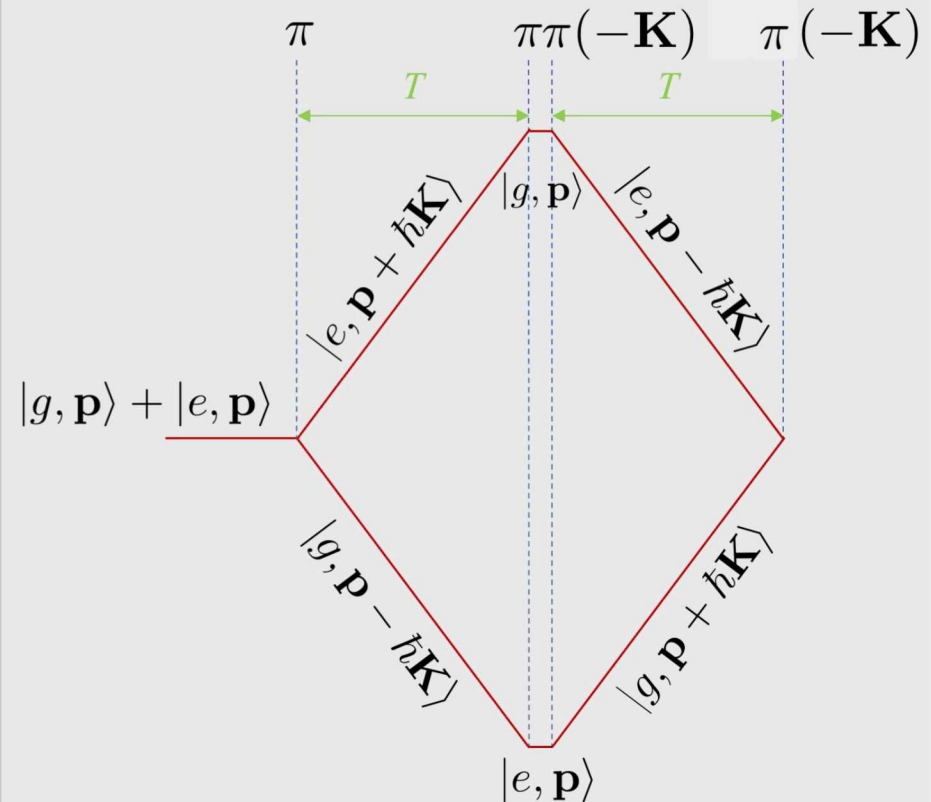
AI with Doppler-sensitive initial entangled state



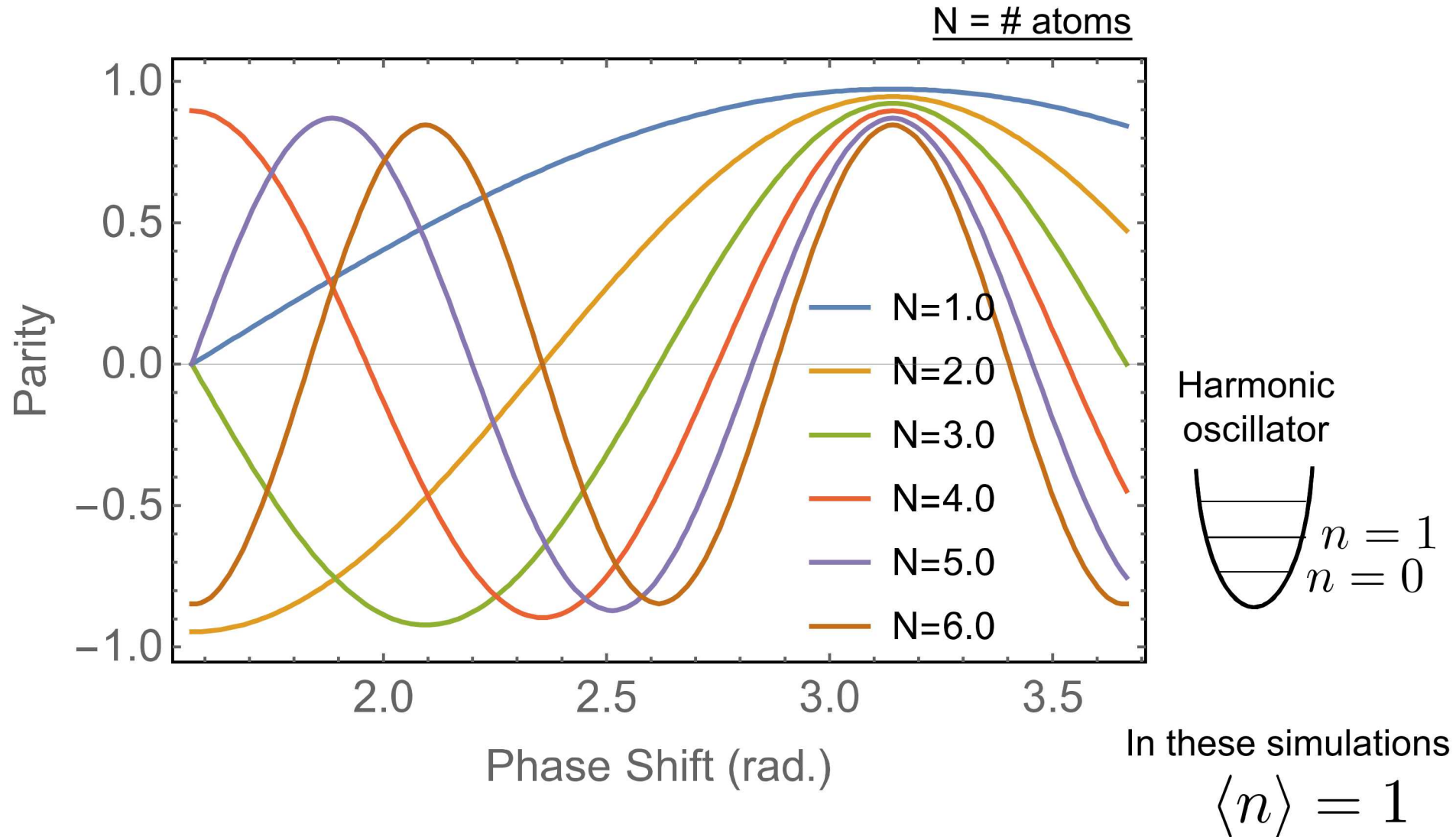
Raman beam configurations for  $\mathbf{K}$  and  $-\mathbf{K}$  pulses



AI with Doppler-free initial entangled state and two additional pulses:  $\pi(\mathbf{K})$  and  $\pi(-\mathbf{K})$



# Effect of temperature

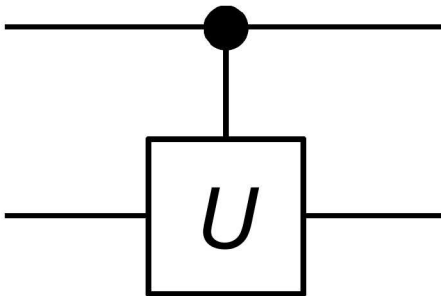


# Applications

Quantum simulation

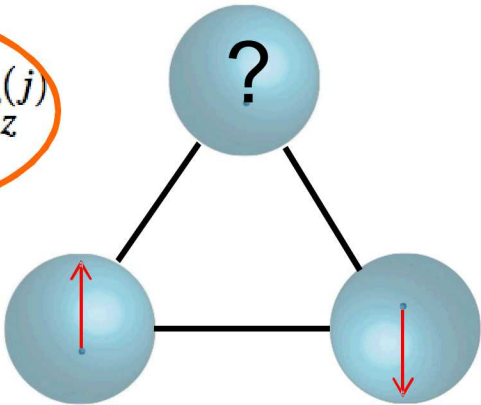
$$H_P = \sum_{i=1}^N \tilde{h}_i \sigma_z^{(i)} + \sum_{i,j=1}^N \tilde{J}_{ij} \sigma_z^{(i)} \otimes \sigma_z^{(j)}$$

Pairwise entangling gates between two neutral atoms

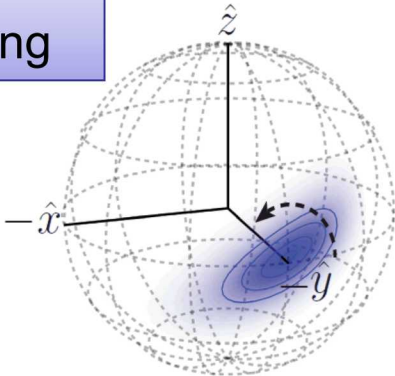


Goal of our research  
Achieve exquisite control of interactions

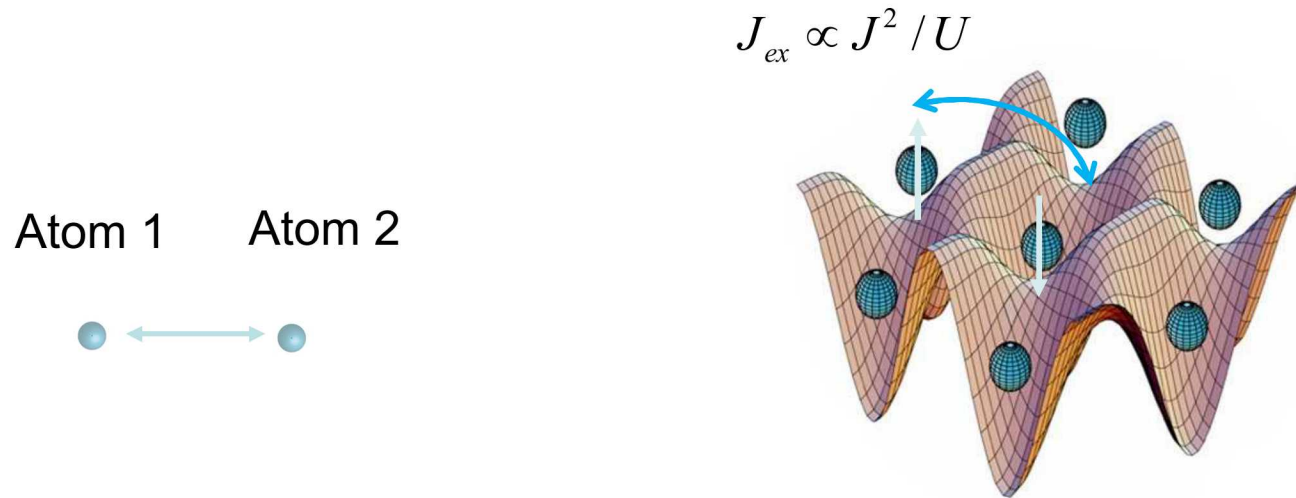
Large-scale/rapid entanglement for sensing



Frustrated magnetism



# Interaction between *neutral* atoms



- Interaction between ground state atoms is small  $\sim 100$  Hz
- Thermal energy scales too large (e.g., QSIM)
- Long gate times (e.g., QIP)

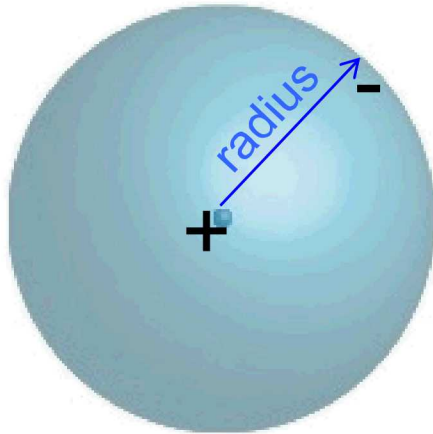
One solution: use Rydberg states

S. Trotzky *et al.*, Science **319**, 295-299 (2008)

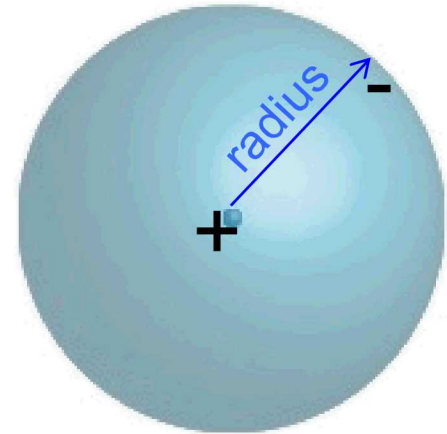
I. Bloch, J. Dalibard, and S. Nascimbène, Nat. Phys. **8**, 267-276 (2012)

# Interaction between *neutral* atoms

Valence electron  
in Rydberg state



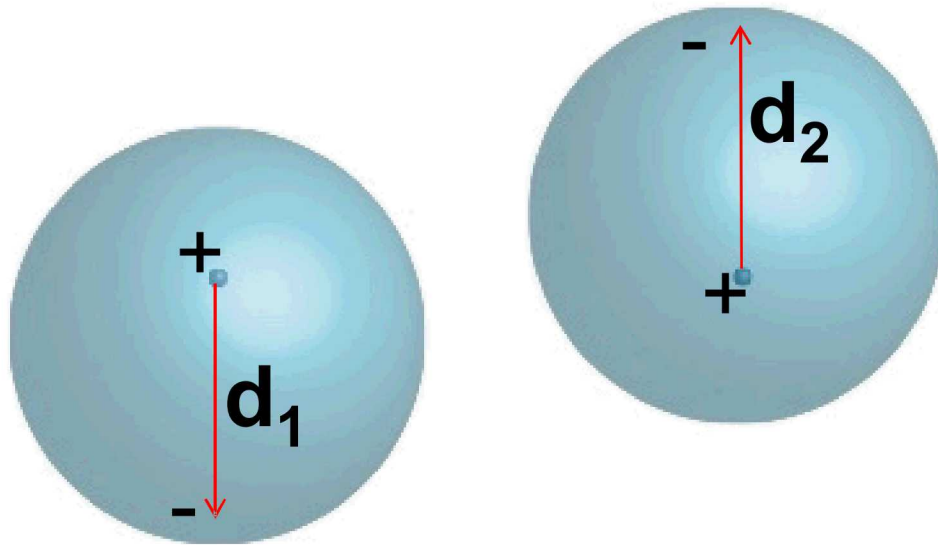
Valence electron  
in Rydberg state



orbital radius  $\propto n^2$

- Excite valence electron to Rydberg state—nearly ionized
- Atom becomes highly polarizable—strong interactions

# Interaction between *neutral* atoms



van der Waals interaction

## Parameter scaling

van der Waals

$$U \propto n^{11}$$

Lifetime

$$\tau \propto n^3$$

DC polarizability

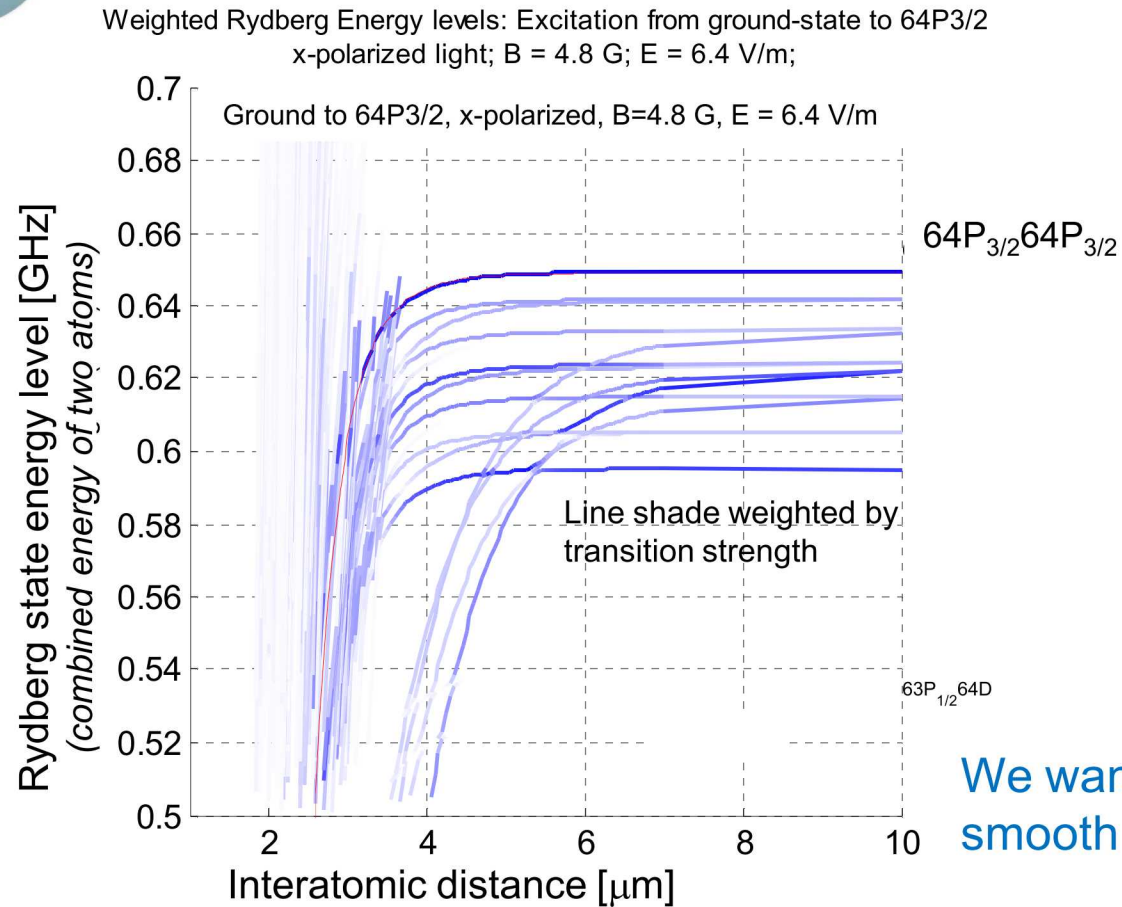
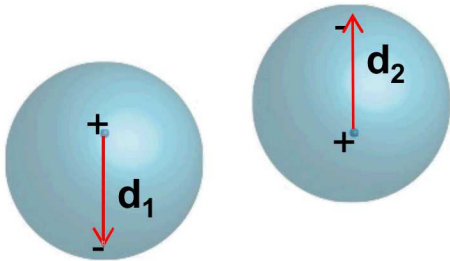
$$\alpha(0) \propto n^7$$

- Even the presence of another atom can cause a massive response  $\gg 10$  MHz
- Induced Electric Dipole-Dipole Interaction  $\propto 1/r^6$

Entanglement demonstrations

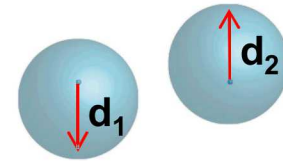
- Madison: Phys. Rev. Lett. 104, 010503 (2010)
- Paris: Phys. Rev. Lett. 104, 010502 (2010)

# Rydberg blockade—the nitty gritty

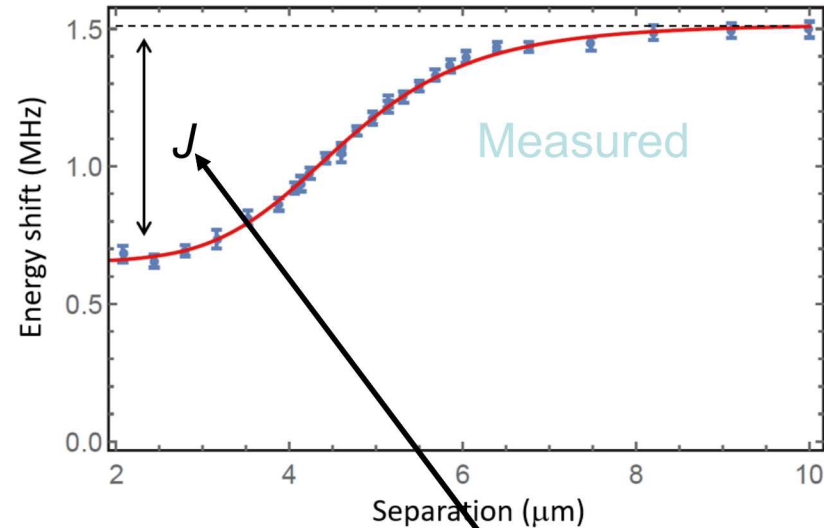
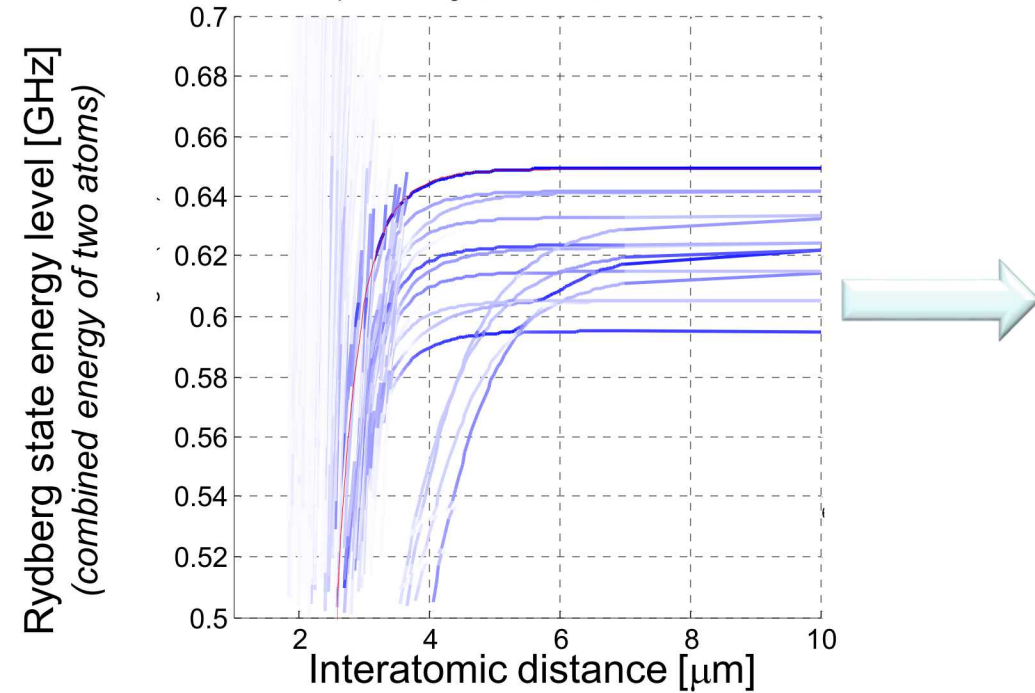


We want something  
smooth and tunable

# Direct Rydberg → Rydberg-Dressed

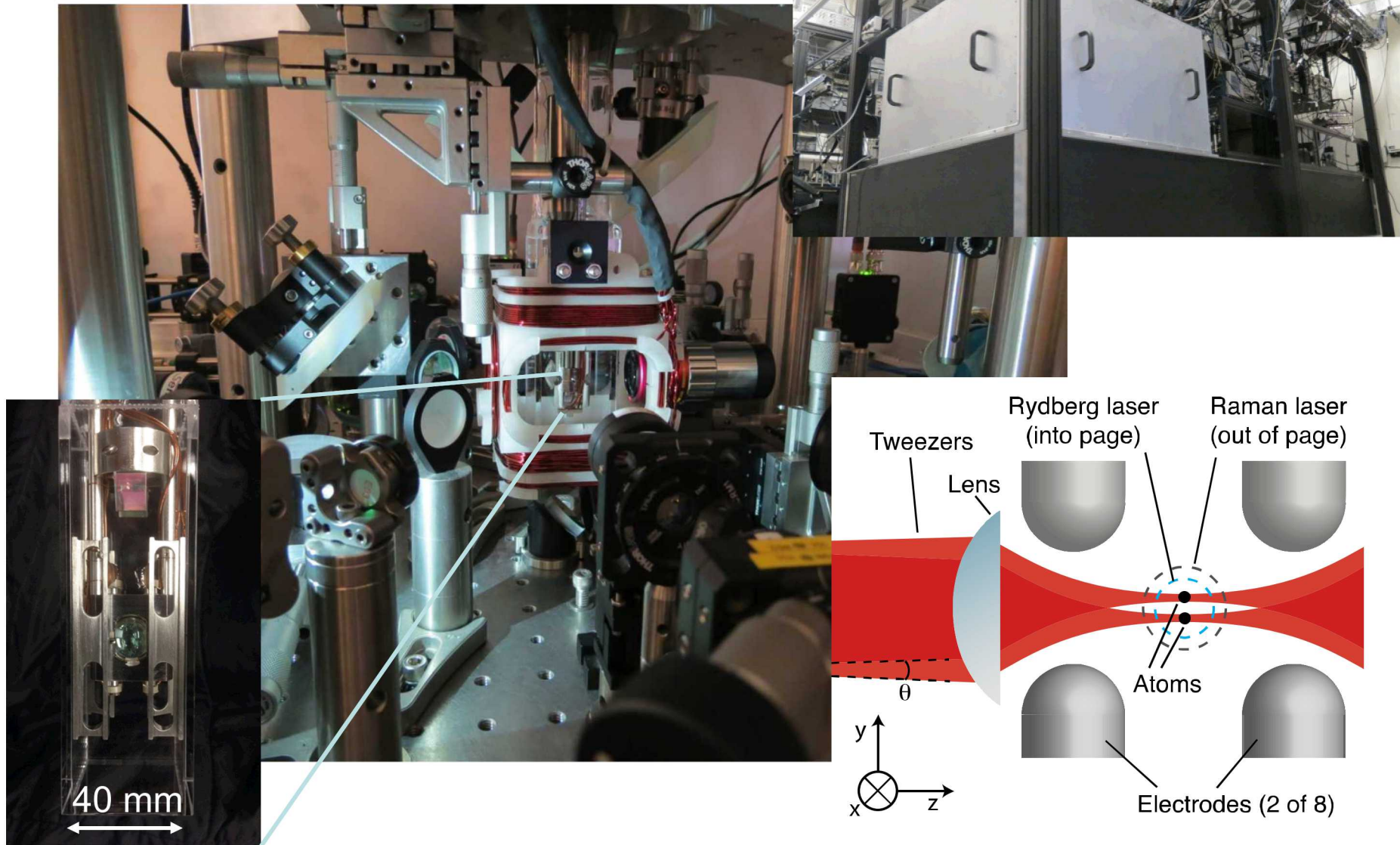


Ground to 64P3/2, x-polarized, B=4.8 G, E = 6.4 V/m

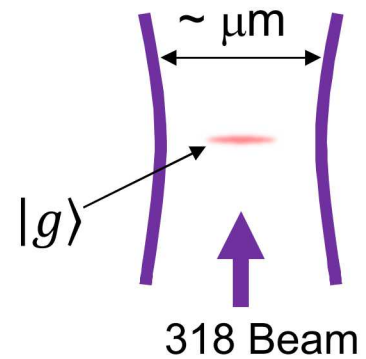
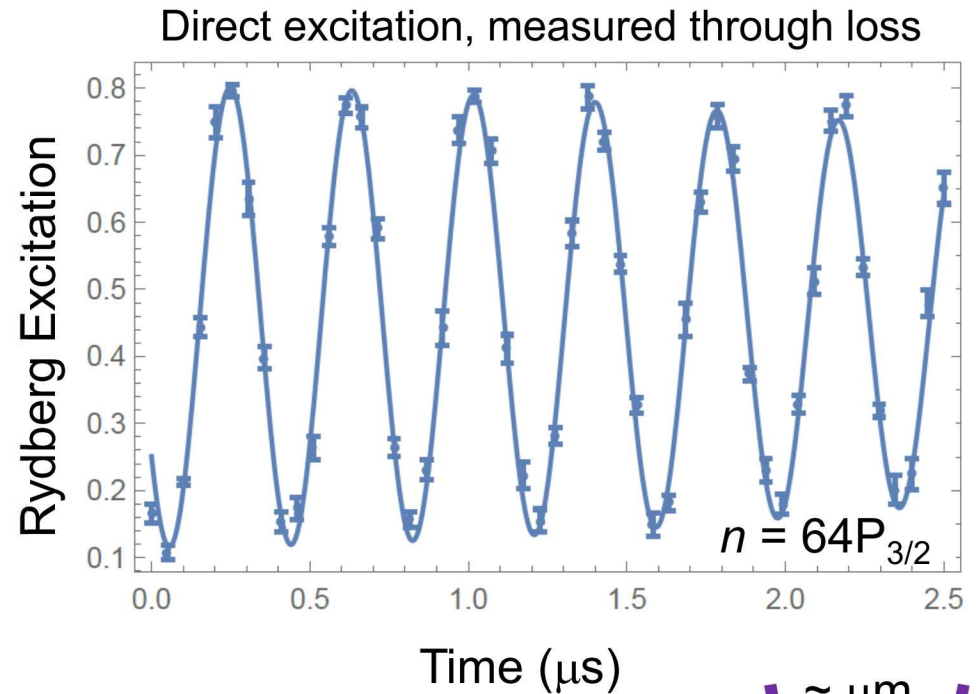
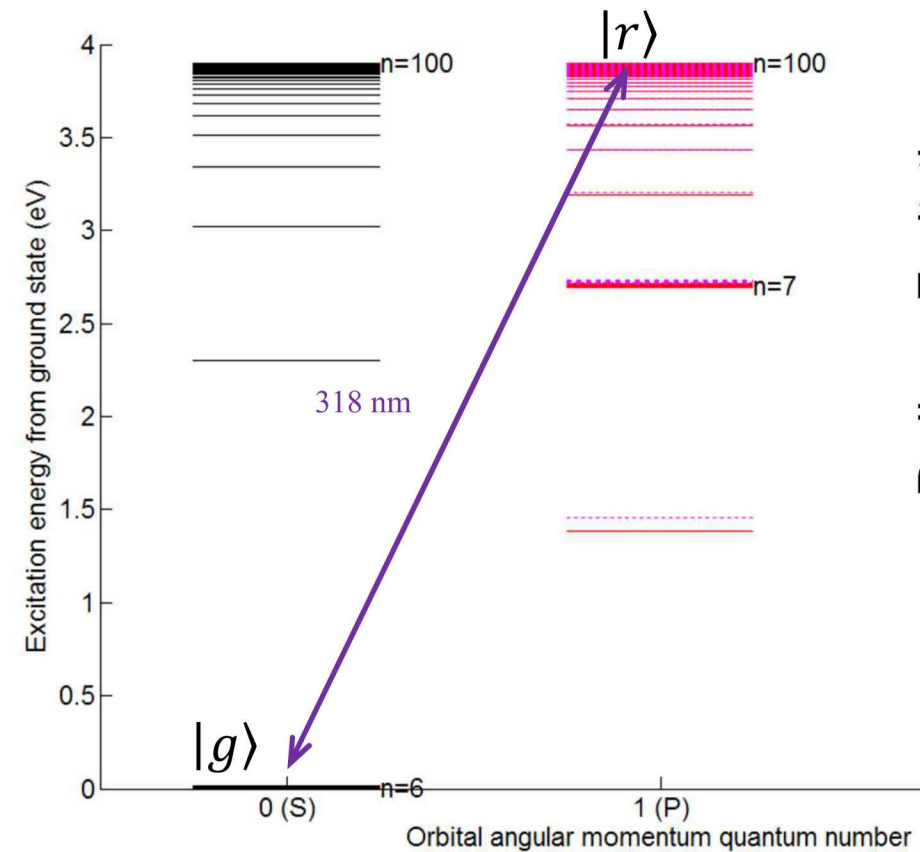


$$H_P = \sum_{i=1}^N \tilde{h}_i \sigma_z^{(i)} + \sum_{i,j=1}^N \tilde{J}_{ij} \sigma_z^{(i)} \otimes \sigma_z^{(j)}$$

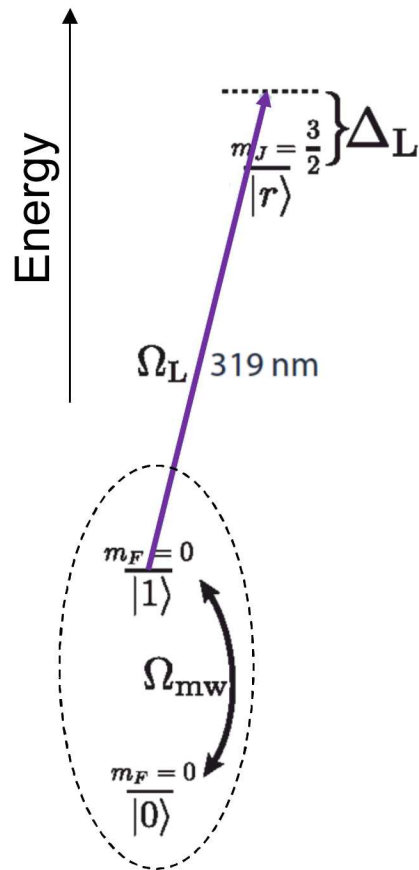
# Apparatus



# Rydberg Rabi oscillations with 318 nm laser



# Creating Rydberg-Dressed states

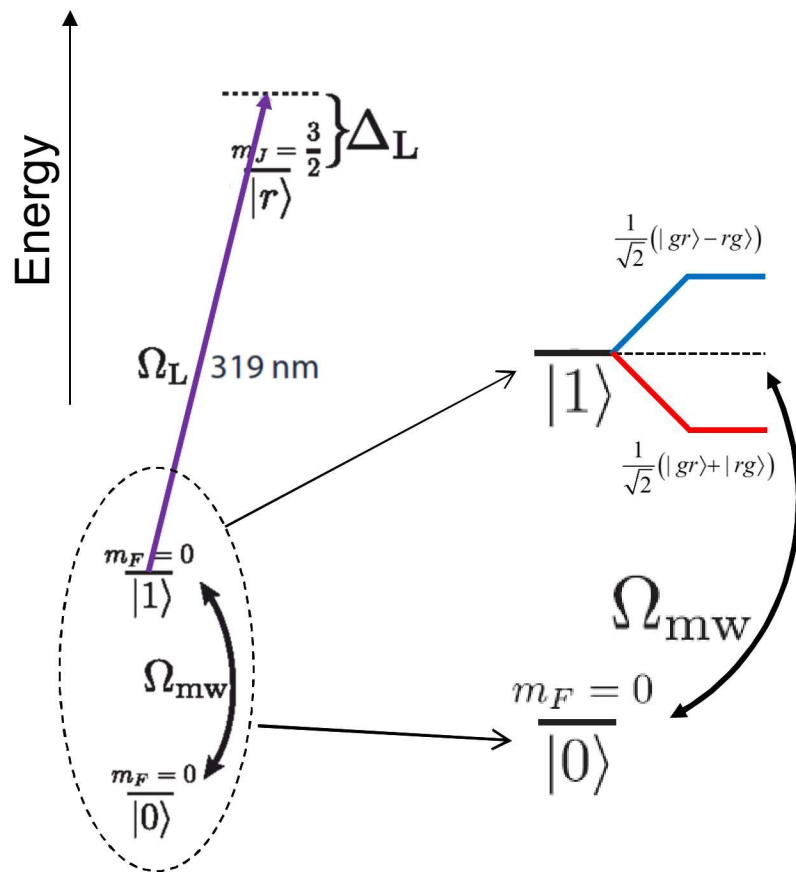


Light-shift Hamiltonian

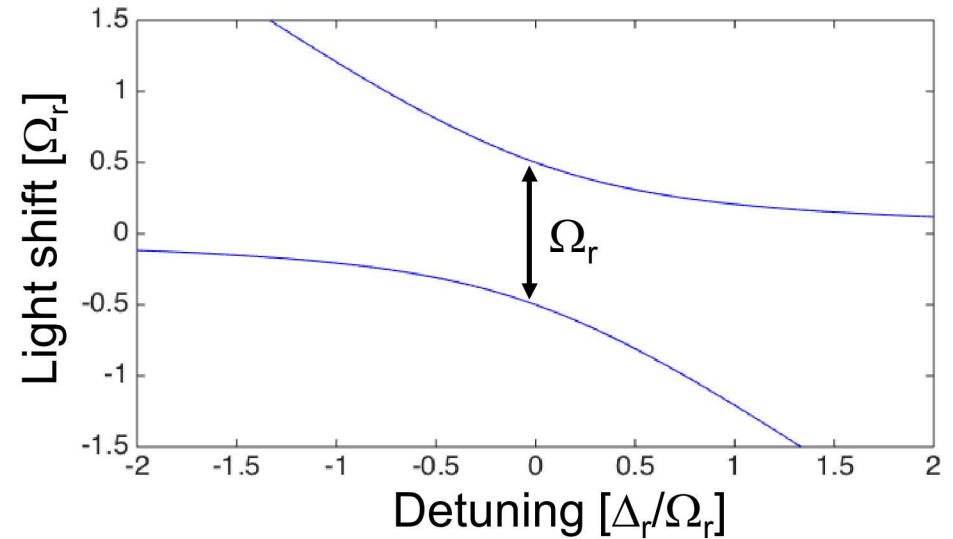
$$H = \frac{\hbar}{2} \begin{pmatrix} -2\Delta_L & \Omega_L \\ \Omega_L & 0 \end{pmatrix}$$

$$\Delta = \frac{\hbar \Omega_L^2}{4 \Delta_L}$$

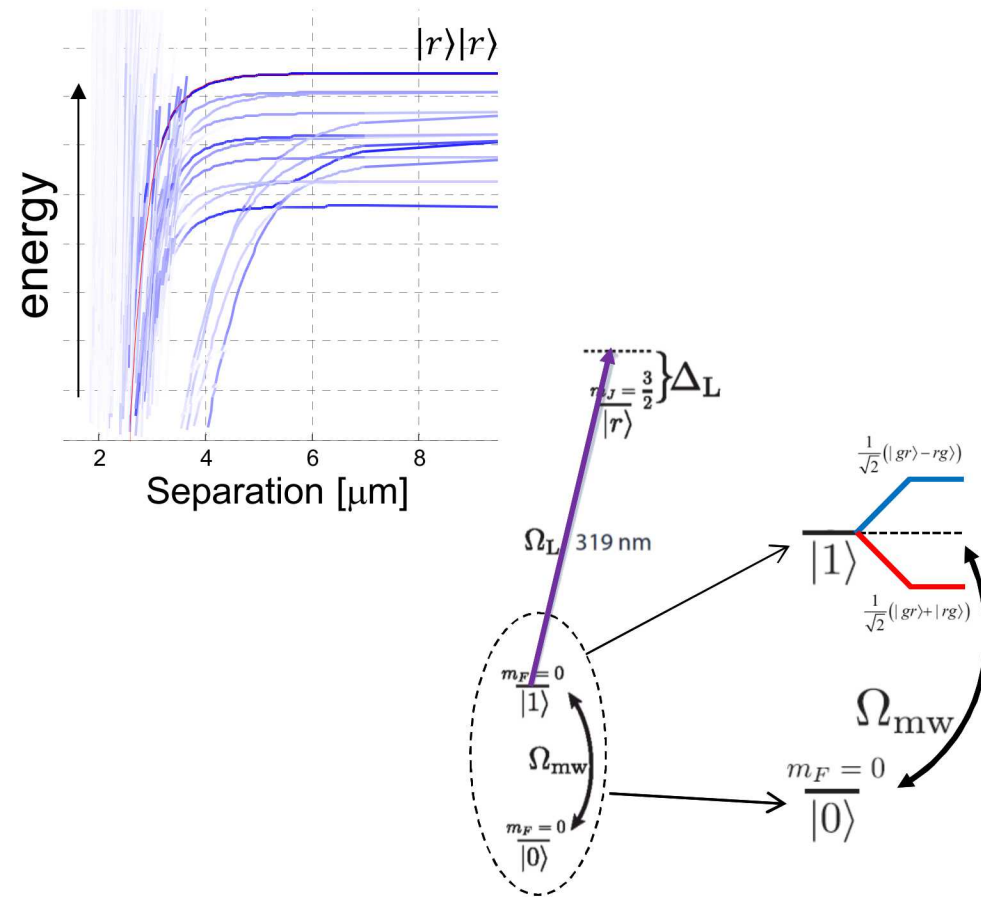
# Creating Rydberg-Dressed states



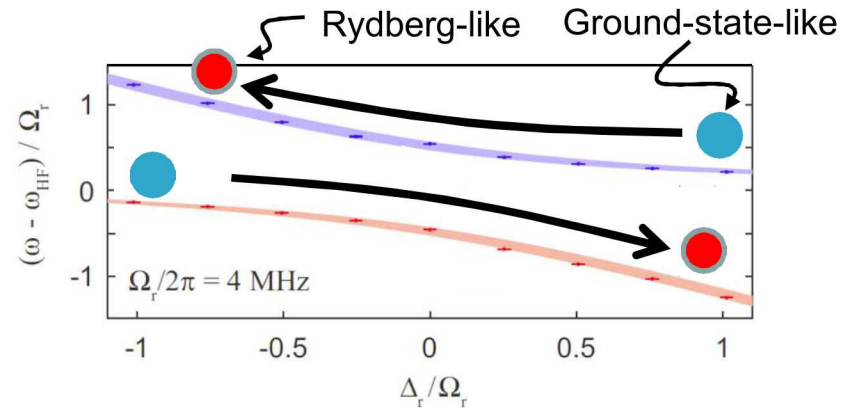
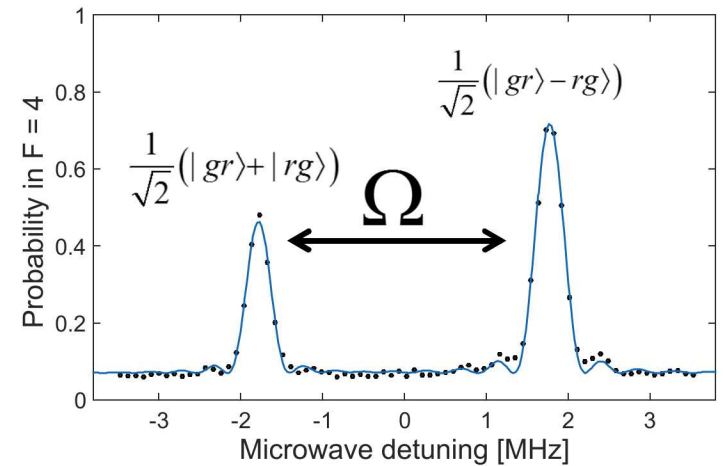
Ground state response



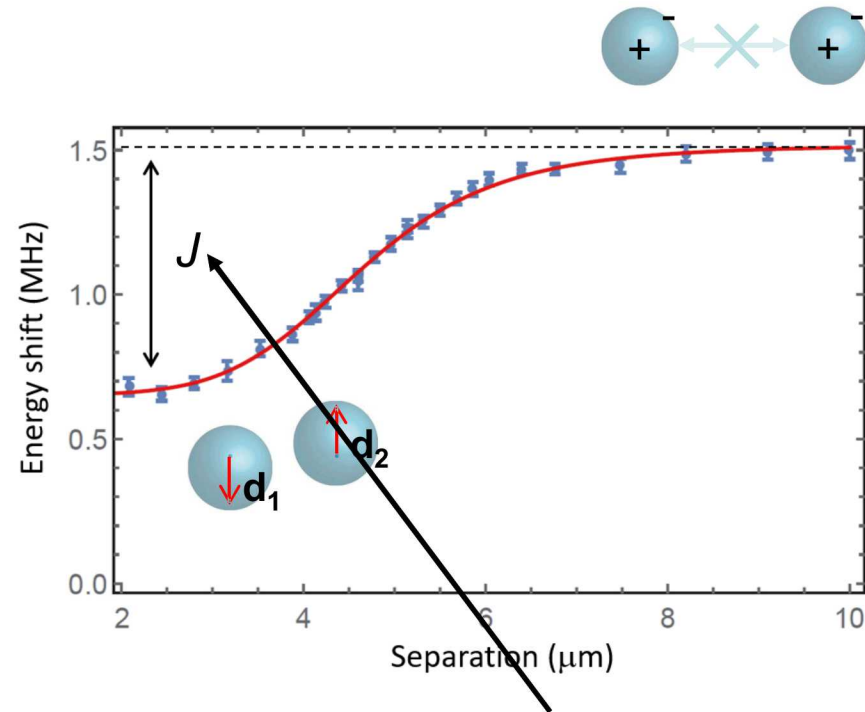
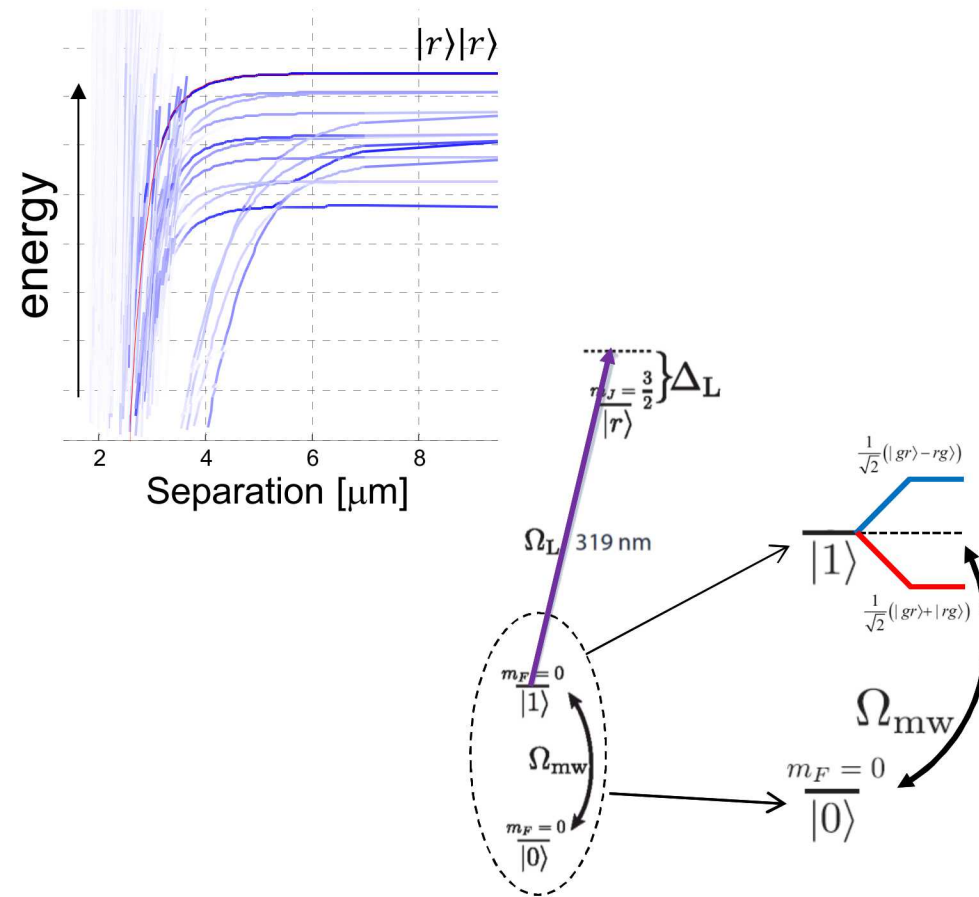
# 1-atom Rydberg-dressed states



Dressed F=4 state Autler-Townes splitting



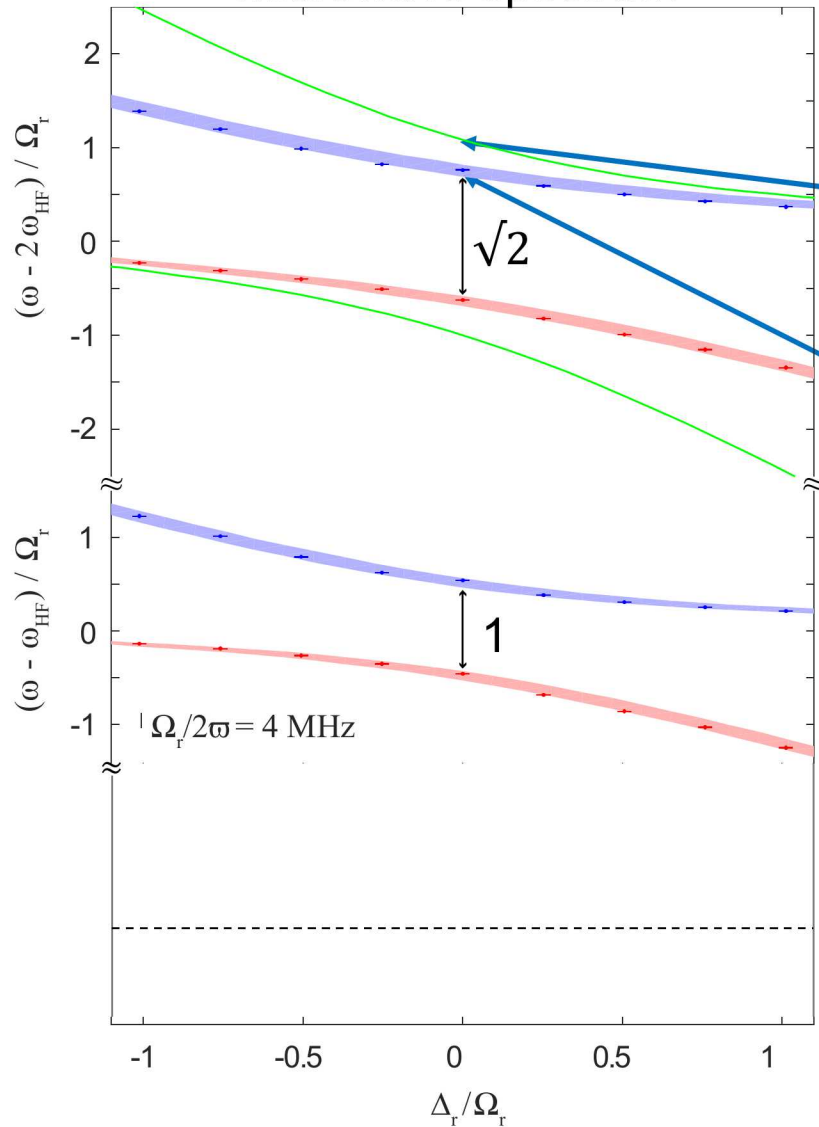
# Rydberg-Dressed interaction



$$H_P = \sum_{i=1}^N \tilde{h}_i \sigma_z^{(i)} + \sum_{i,j=1}^N \tilde{J}_{ij} \sigma_z^{(i)} \otimes \sigma_z^{(j)}$$

# 2-atom Rydberg-dressed states

Microwave spectrum



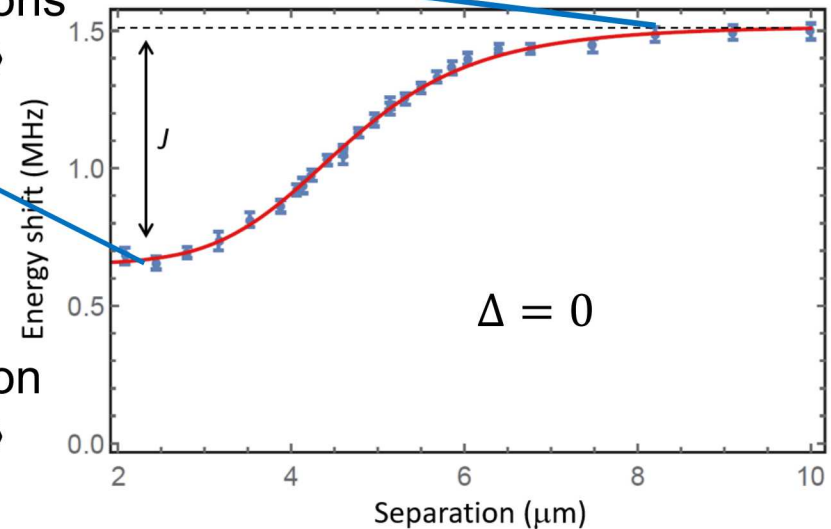
2 excitations to  $|F = 4\rangle$

1 excitation to  $|F = 4\rangle$

All in  $|F = 3\rangle$

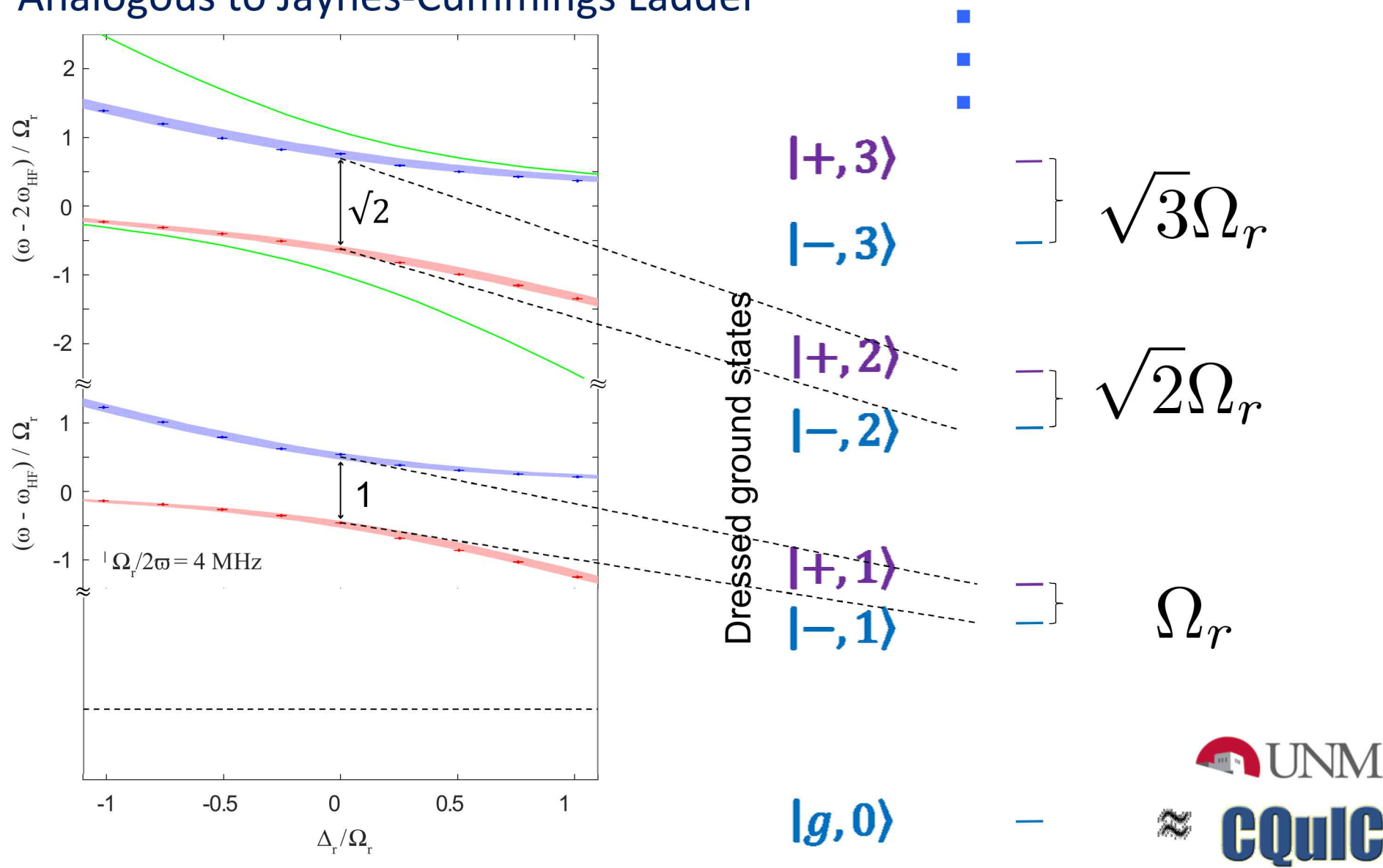
Spatial dependence

Light shift of Rydberg-dressed state



# Extending to $N$ atoms—symmetric

Analogous to Jaynes-Cummings Ladder



# Sample Control Task

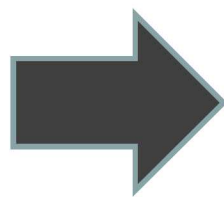
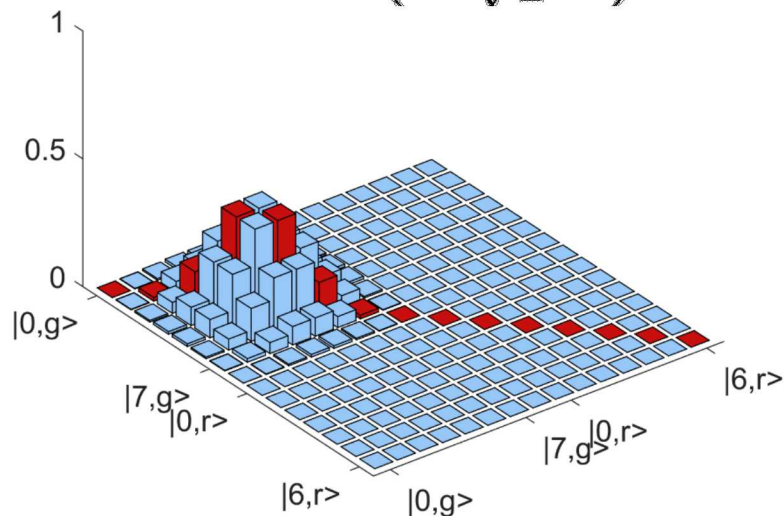
- Create a 7-atom cat state:

$$2\pi * \Omega_r = 5 \text{ MHz}$$

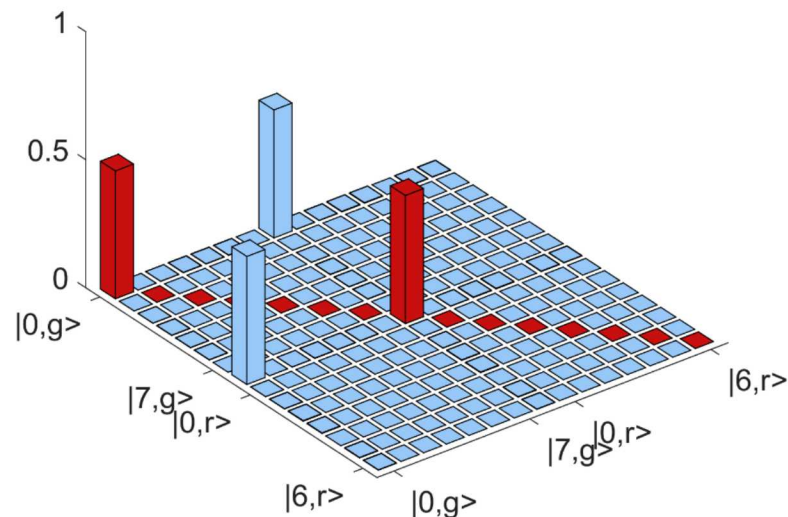
$$2\pi * \delta_r = -2.5 \text{ MHz}$$

$$2\pi * \kappa \approx 600 \text{ kHz}$$

$$\rho_{in} = \left( \frac{|0\rangle + |1\rangle}{\sqrt{2}} \right)^{\otimes 7}$$



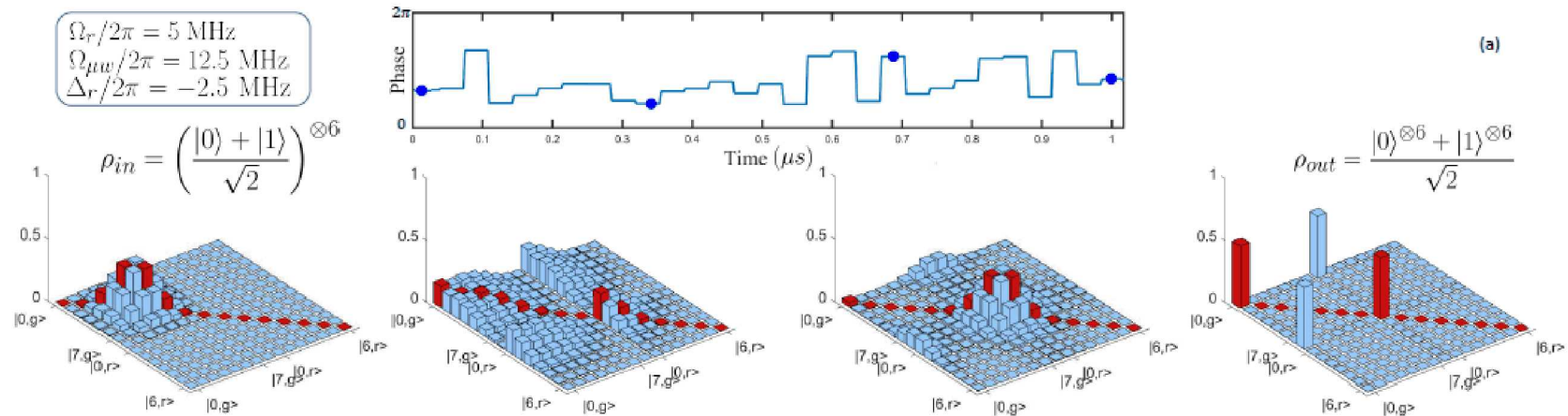
$$\rho_{out} = \frac{|0\rangle^{\otimes 7} + |1\rangle^{\otimes 7}}{\sqrt{2}}$$



- Find the sequence of microwave phases that does this with maximum fidelity

# Arbitrary control

T. Keating *et al.*, Phys. Rev. Lett **117**, 213601 (2016)



$$\hat{H}_c[\phi_{\mu w}(t)] = \frac{\hbar\Omega_{\mu w}}{2} \sum_{i=1}^N \left( e^{-i\phi_{\mu w}(t)} |1\rangle_i \langle 0| + e^{+i\phi_{\mu w}(t)} |0\rangle_i \langle 1| \right) = \hbar\Omega_{\mu w} (\cos[\phi_{\mu w}(t)] \hat{J}_x + \sin[\phi_{\mu w}(t)] \hat{J}_y)$$

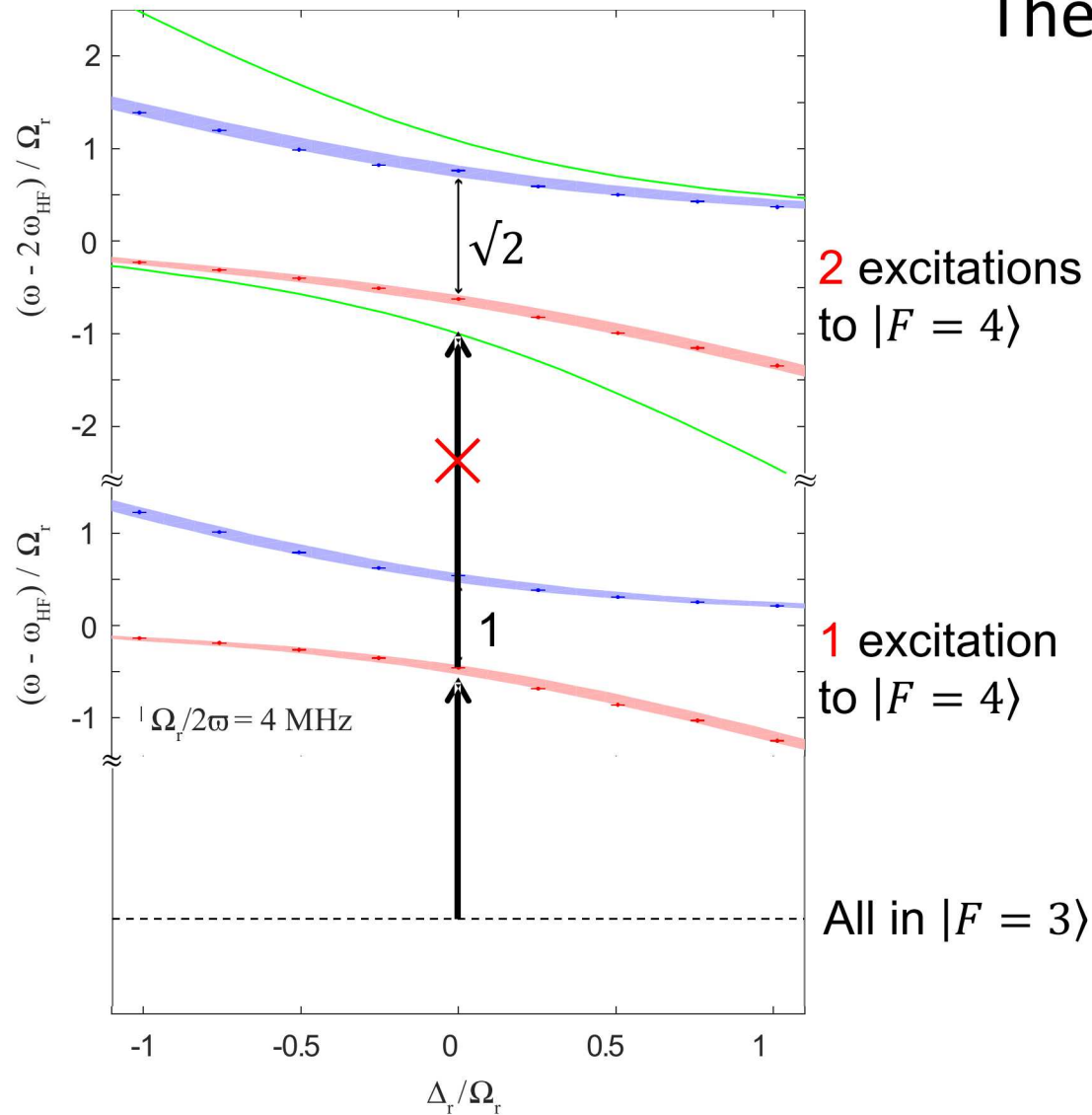
control

$$\phi_{\mu w}(t)$$

The nonlinearity of the JCM, together with externally applied fields makes the system fully controllable on the symmetric Dicke space; we can generate an arbitrary superposition state.

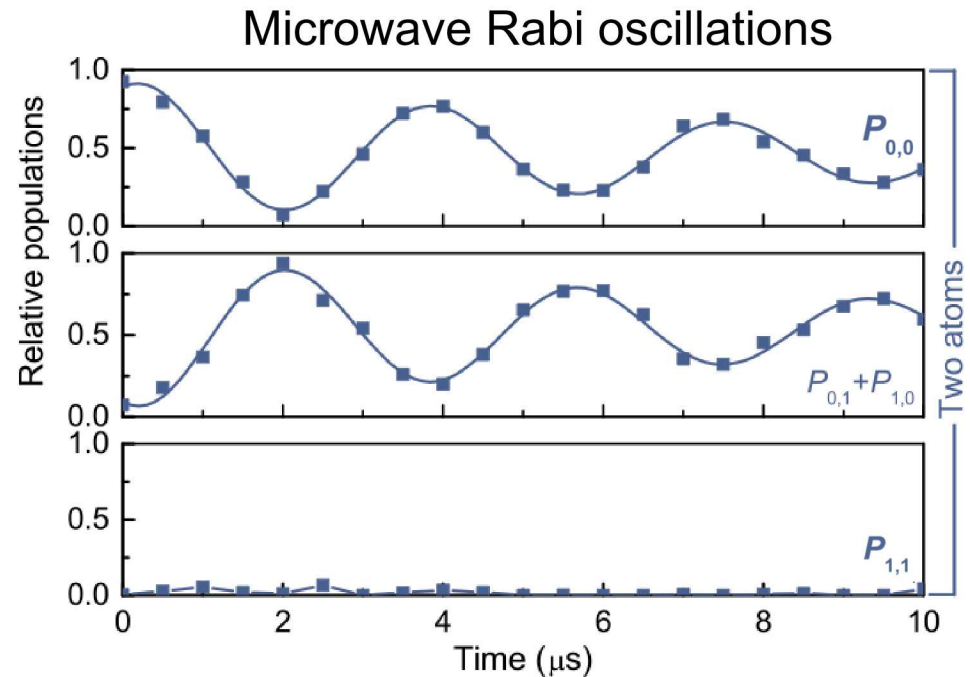
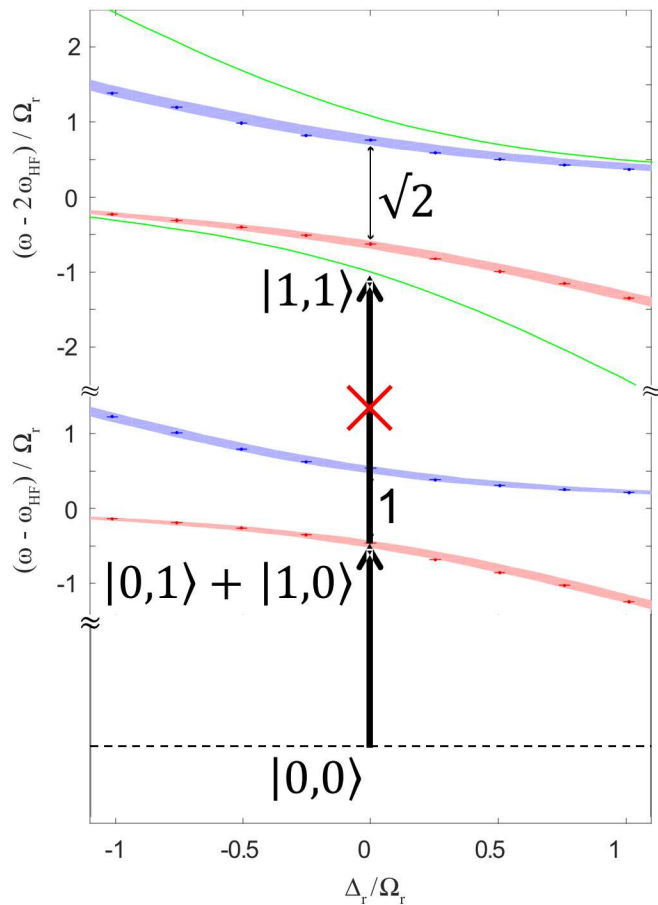
# Creating entanglement with a single step

The “spin-flip blockade”



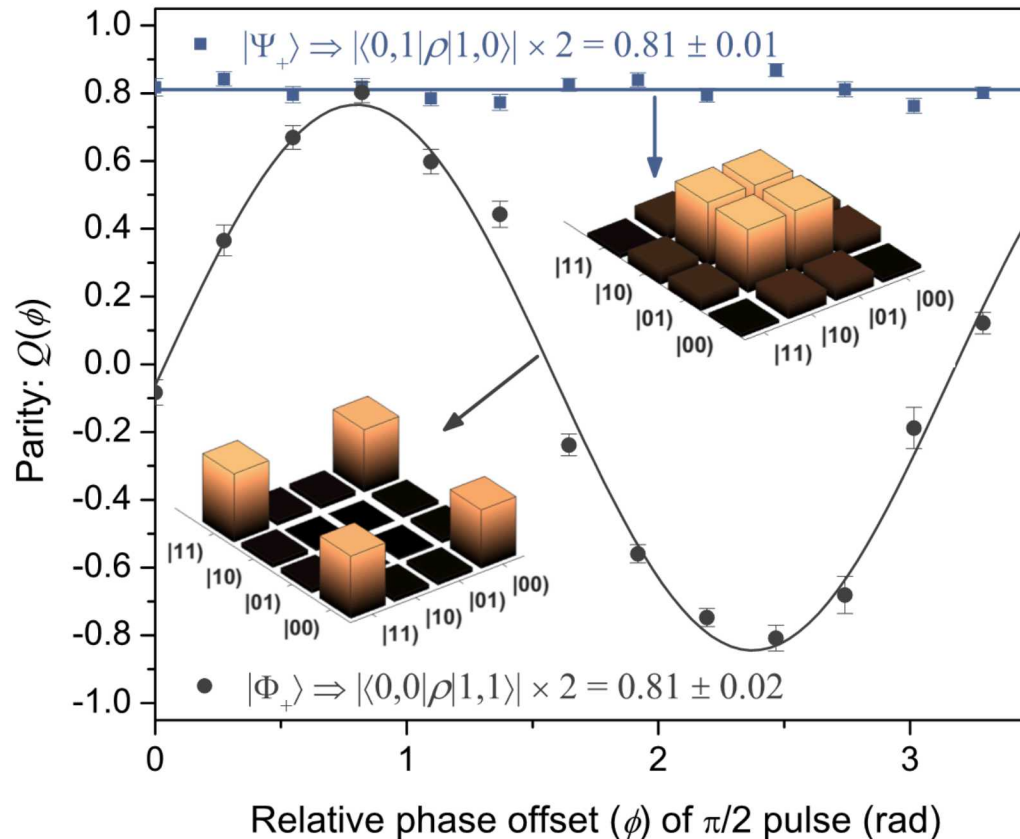
# Creating entanglement with a single step

## The “spin-flip blockade”



# Spin-flip blockade

Verify the entanglement via parity measurements



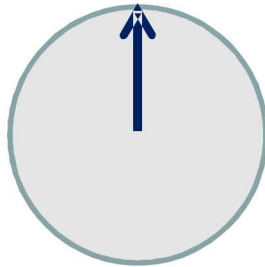
1. Prepare Bell state
2. Apply global  $\pi/2$  with given phase
3. Measure parity  $Q$
4. Obtain bound on fidelity = 0.81(2)

# Generating Entanglement

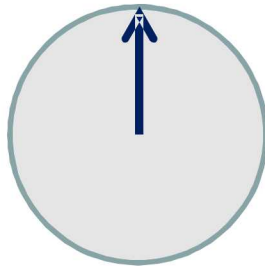
Phase:

$$\hat{H}_{DD} = J |11\rangle\langle 11|$$

Laser **ON**F

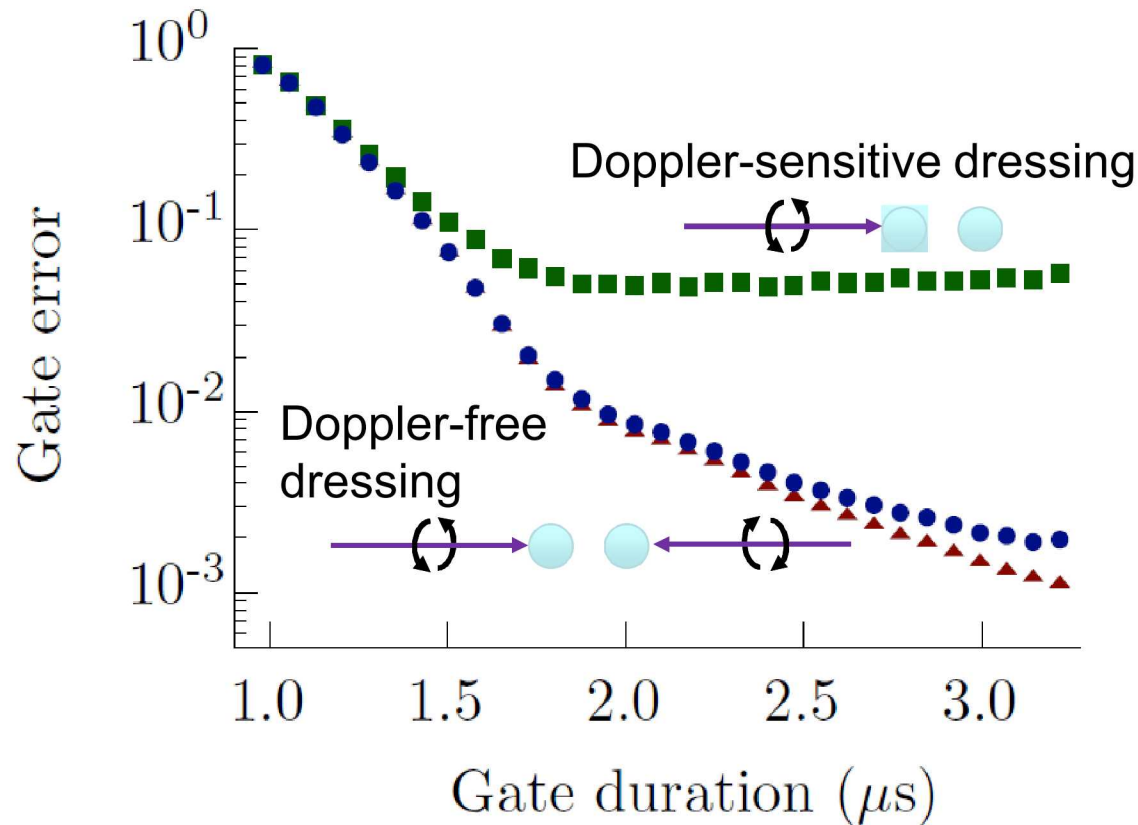


—————  $|11\rangle$



—————  $|01\rangle |10\rangle$   
and  $|00\rangle$

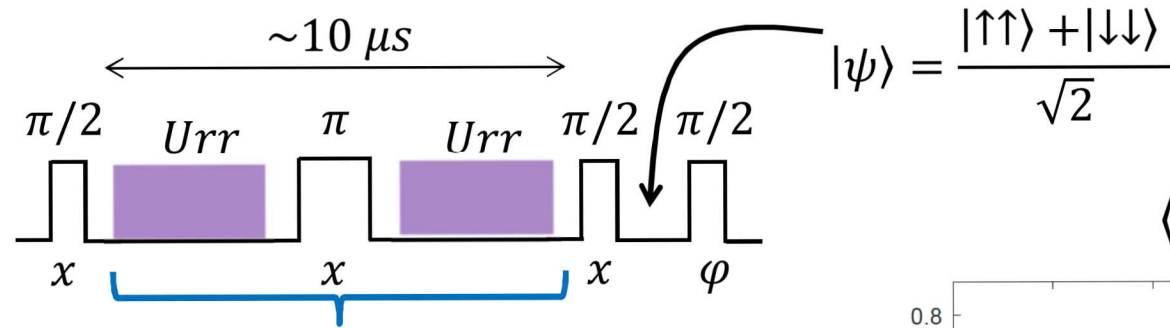
# Simulated CPHASE gate fidelities



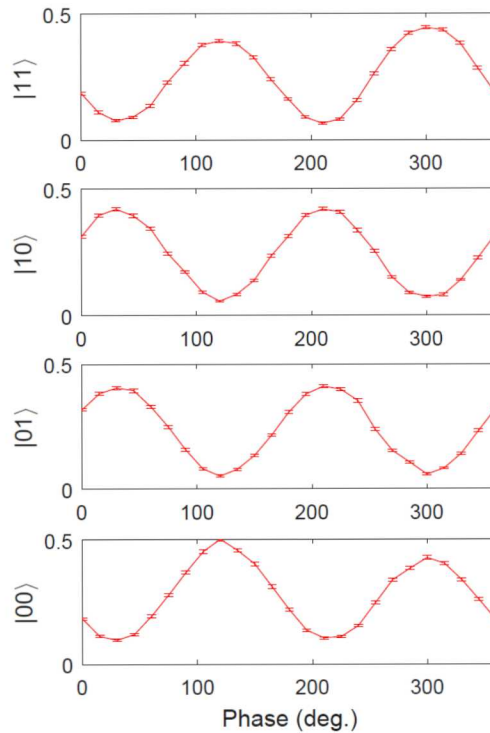
$\Omega = 0 \rightarrow 3 \text{ MHz}$   
 $\Delta/2\pi = 6 \rightarrow 0 \text{ MHz}$   
 $\Gamma = 3.7 \text{ kHz}$   
 $T = 16 \text{ } \mu\text{K}$

- Motional errors set a high floor on error for the single-beam scheme.
- The Doppler-free scheme is limited by the much smaller photon scattering rate.

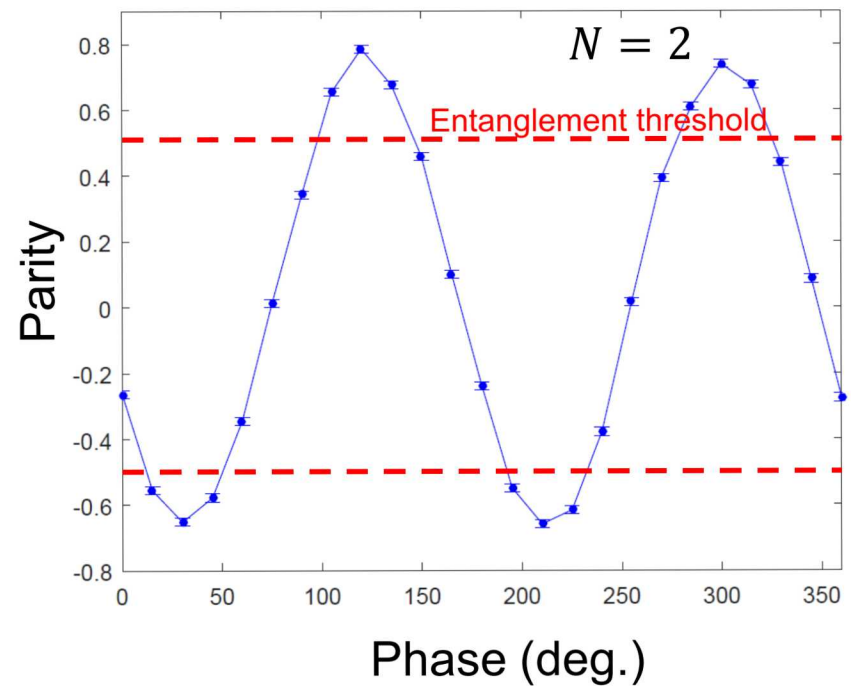
# Dressed CPHASE gate



$$\hat{H}_{\text{eff}} = J \times (\hat{S}_z)^2$$



$$\langle \hat{\Pi} \rangle \propto \sin N\phi$$

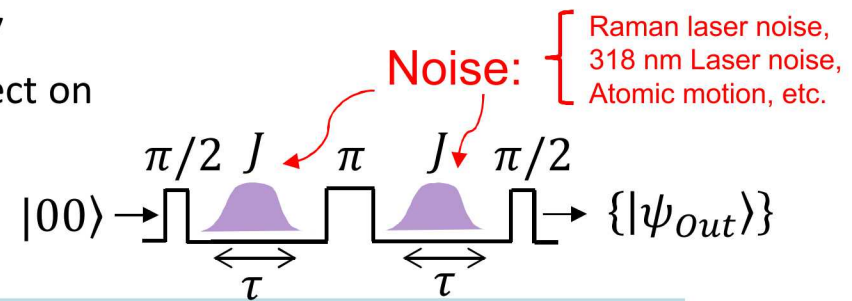


# Error budget for CPHASE gate

## Error budget: informing an experimental strategy

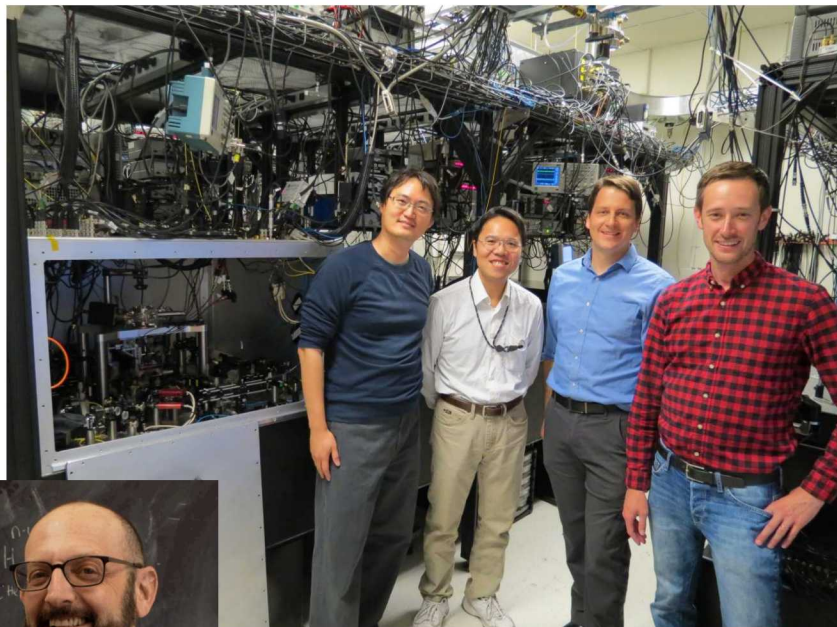
- Trace over external parameters to calculate effect on state fidelity, using measured parameters.

$$\text{Fidelity} = \langle \psi_{tar} | \rho_{out} | \psi_{tar} \rangle$$



| Effect                       | Fidelity reduction                      | Mitigation                                                           |
|------------------------------|-----------------------------------------|----------------------------------------------------------------------|
| LO noise                     | 10% $\pm$ 0.1%(stat.)<br>$\pm$ 2%(sys.) | Clean Raman laser/ $\mu$ Wave cavity                                 |
| State purity                 | <3%                                     | Clean Raman laser/ $\mu$ Wave cavity                                 |
| Atomic position spread       | 3% $\pm$ 0.5%                           | Sideband cooling to ground state                                     |
| Wave-packet overlap          | <0.1%                                   | Sideband cooling to ground state                                     |
| Atomic velocity spread       | <0.1%                                   | Sideband cooling to ground state                                     |
| 318 nm Laser frequency noise | 0.2% $\pm$ 0.1%                         | Pre-stabilized seed lasers, different detuning, dynamical decoupling |
| Spontaneous emission         | 0.4% $\pm$ 0.2%                         | Higher principal quantum #                                           |
| 318 nm laser amplitude noise | <0.1%                                   | Install "noise eater" on laser                                       |

# CQuIC-Sandia Team



J. Lee, Y-Y. Jau, M. Martin, G. B.



I. Deutsch (UNM)

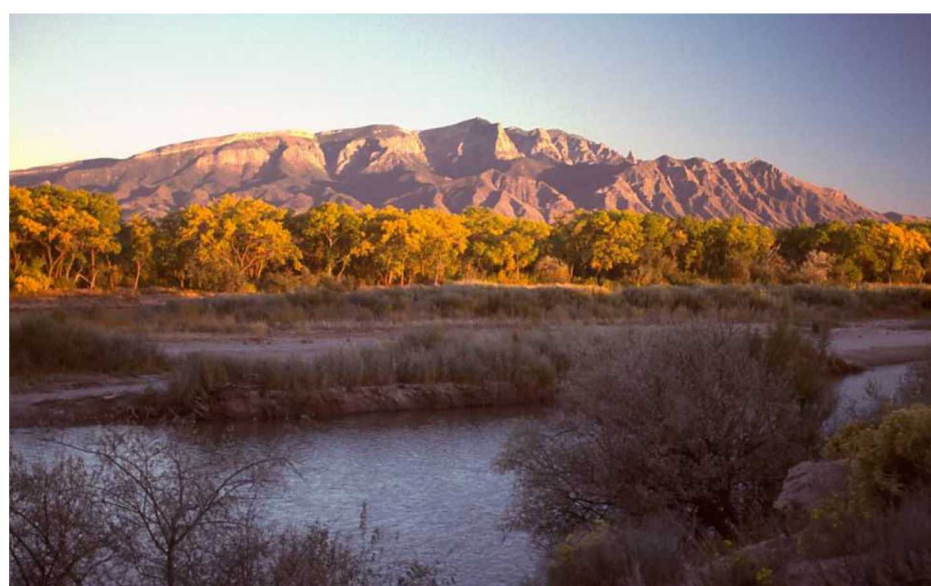
Current  
team  
members



J. Bainbridge (UNM)



A. Orozco (UNM)



Sandia mountains viewed from the Rio Grande near Albuquerque, New Mexico

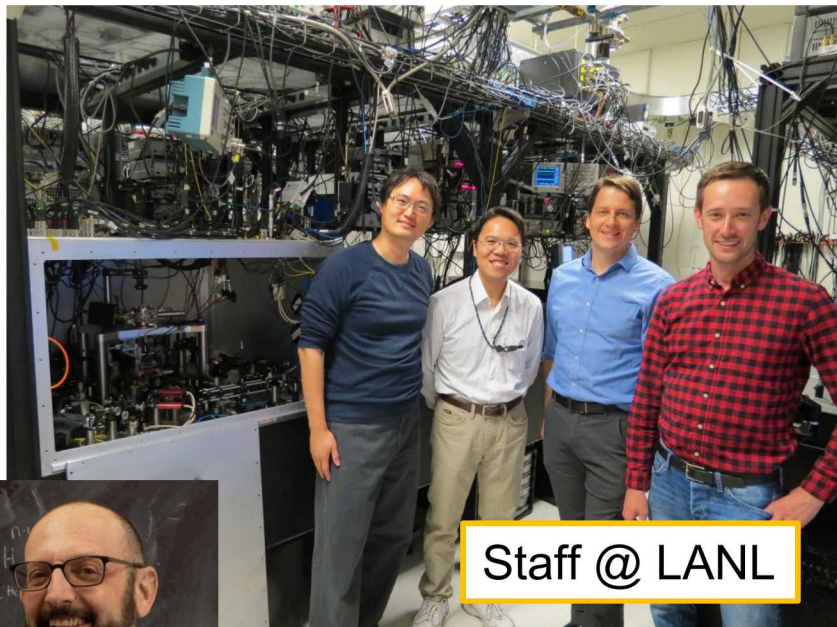


**LDRD**

Laboratory Directed Research and Development



# CQuIC-Sandia Team



Staff @ LANL

J. Lee, Y-Y. Jau, M. Martin, G. B.



I. Deutsch (UNM)

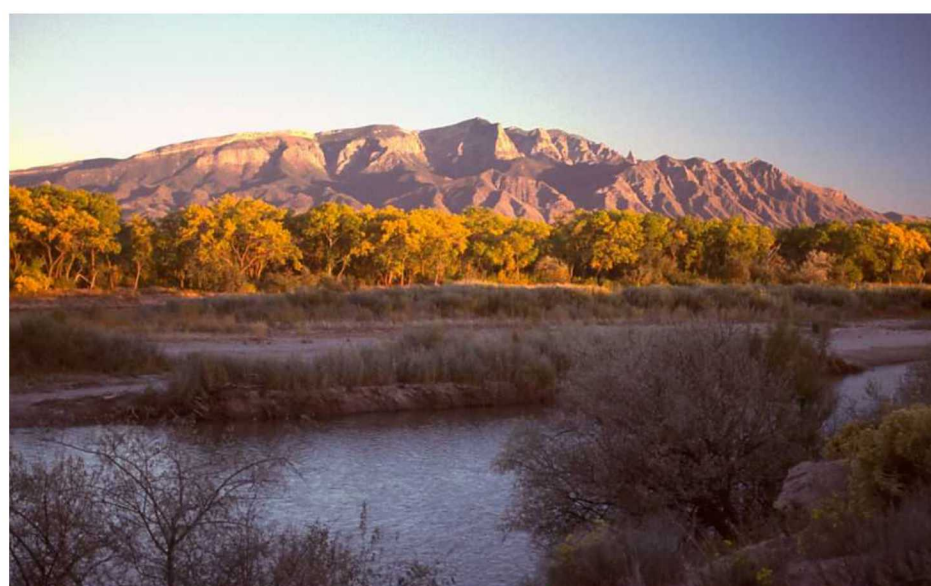
Previous  
team  
members



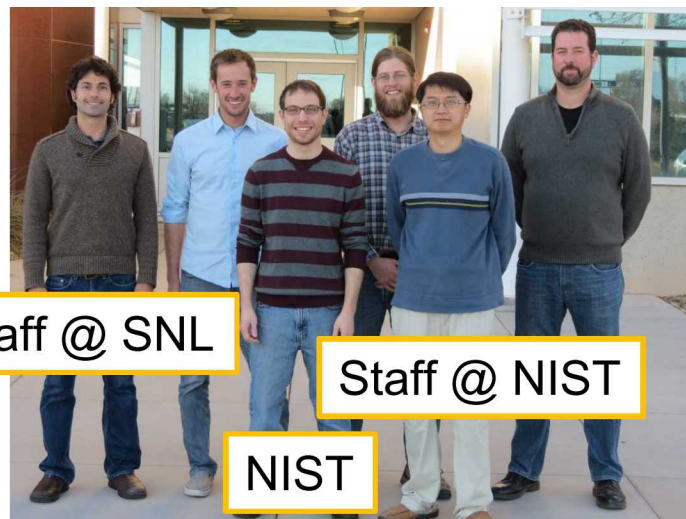
Staff @ SNL

AOSense

G. B., A. Rakholla, H. McGuinness



Sandia mountains viewed from the Rio Grande near Albuquerque, New Mexico



Staff @ SNL

Staff @ NIST

NIST

L. Parazzoli, G. B., A. Hankin, A. Ferdinand, J. Chou, G. Burns

!!!Wanted: Postdocs and students!!!