

# IONIZATION WAVE DYNAMICS OF A SAND2018-12449C PLASMA JET IN CONTACT WITH LIQUID WATER\*

**Amanda M. Lietz<sup>1</sup>, Edward V. Barnat<sup>2</sup>, Caroline Winters<sup>2</sup>,  
John E. Foster<sup>1</sup> and Mark J. Kushner<sup>1</sup>**

**<sup>1</sup>University of Michigan  
Ann Arbor, MI 48109-2122 USA  
lietz@umich.edu**

**<sup>2</sup>Sandia National Laboratories  
Albuquerque, NM 87123 USA**

**71<sup>st</sup> Annual Gaseous Electronics Conference  
Portland, OR, USA.  
8 November 2018**

- **Work supported by the US Dept. of Energy Office of Fusion Energy Science and National Science Foundation.**
- *Sandia National Laboratories is a multimission laboratory managed and operated by National Technology and Engineering Solutions of Sandia LLC, a wholly owned subsidiary of Honeywell International Inc. for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA0003525.*

# AGENDA

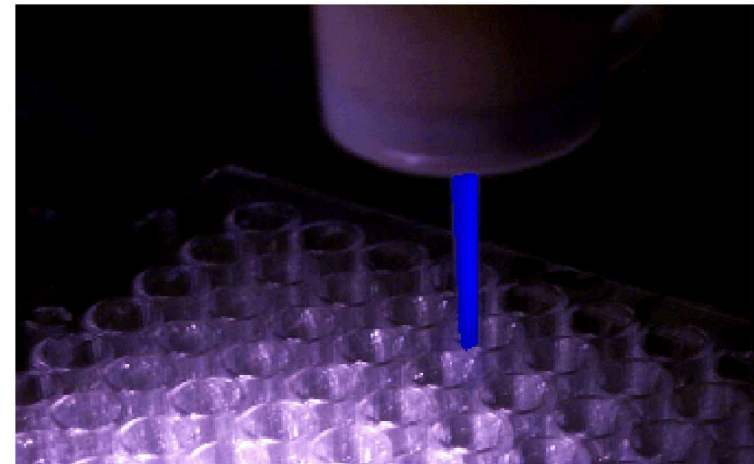
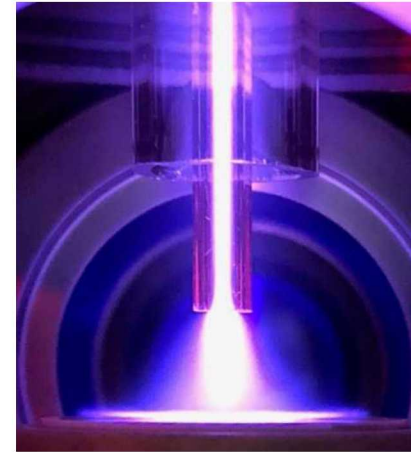
---

- Plasma jets onto Liquids
- Description of Experiment
- Laser Collisional Induced Fluorescence (LCIF)
- Plasma Jet Contacting  $\text{TiO}_2$  ( $\epsilon_r = 80$ )
- Plasma Jet Contacting Liquid Water
- Modeling and Future Work

# ATMOSPHERIC PRESSURE PLASMA JETS

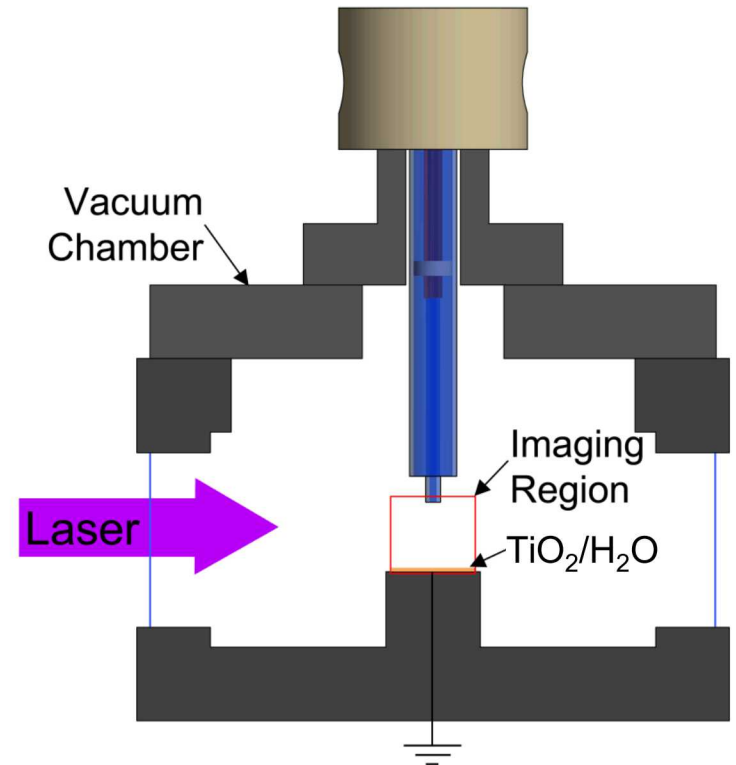
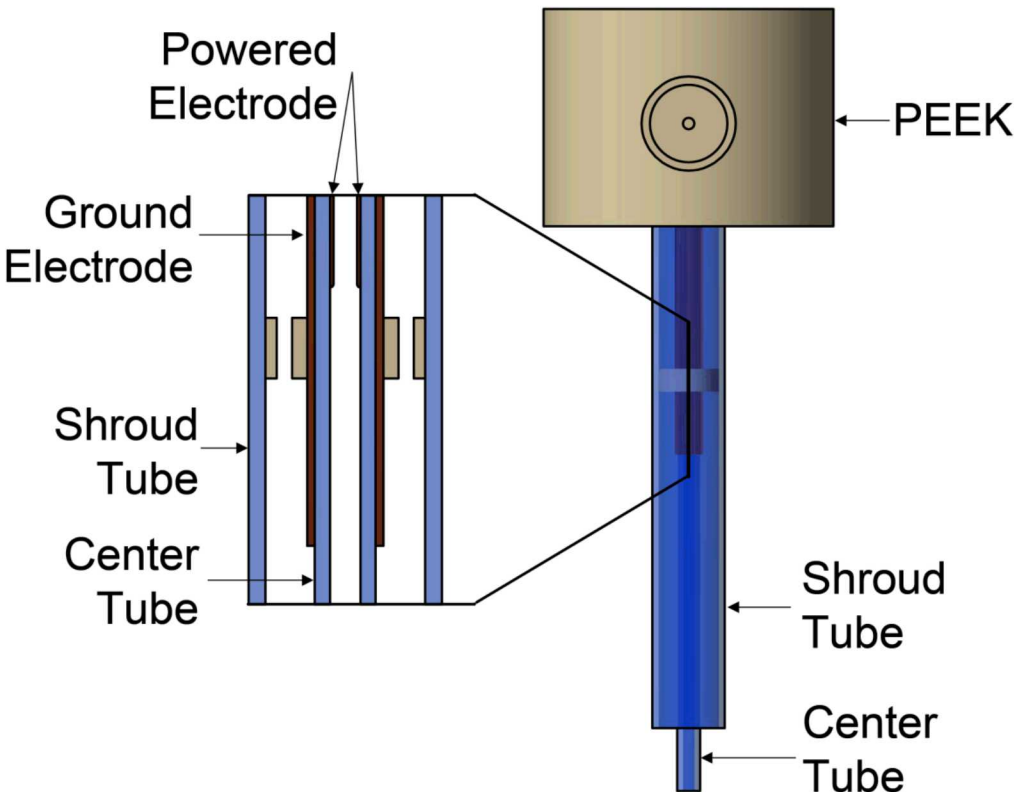
---

- Atmospheric pressure plasma jets (APPJs) are a popular source of chemistry for biomedical applications.
- The plasma propagates as an ionization wave (IW) that is repetitively pulsed.
- The IW gives rise to reactive oxygen and nitrogen species (RONS) which produce the biological effect.
- Objective: Investigation IW dynamics in a plasma jet contacting liquid in a well controlled environment.



- S. Mohades, et al., Physics of Plasmas 22, 122001 (2015).

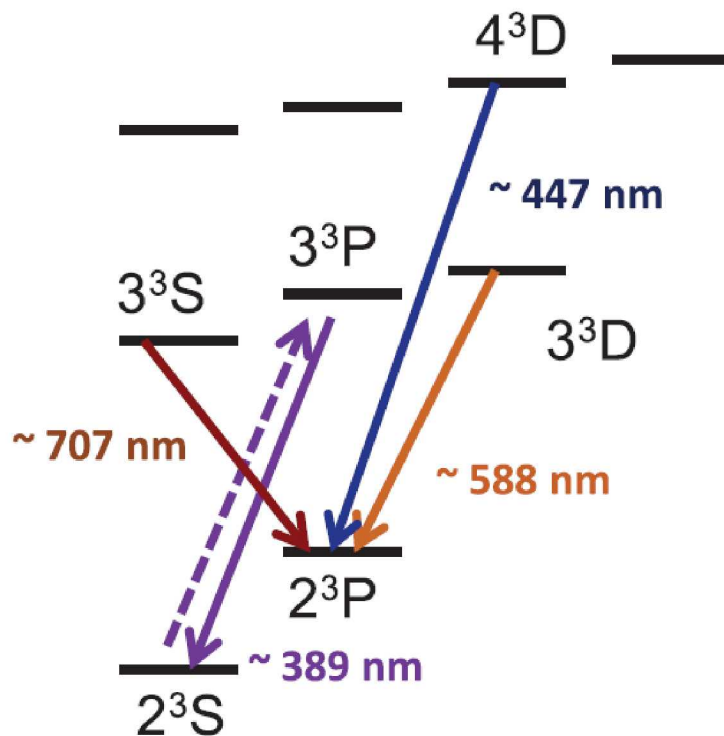
# EXPERIMENTAL SETUP



- **Annular powered electrode inside the center tube.**
- **Placing the APPJ in a vacuum chamber - consistent and controlled chemistry, ground planes, and gas flow.**
- **Coaxial tube enables a gas shroud – control environment independently of gas in main jet.**

# LASER COLLISIONAL INDUCED FLUORESCENCE (LCIF)

---

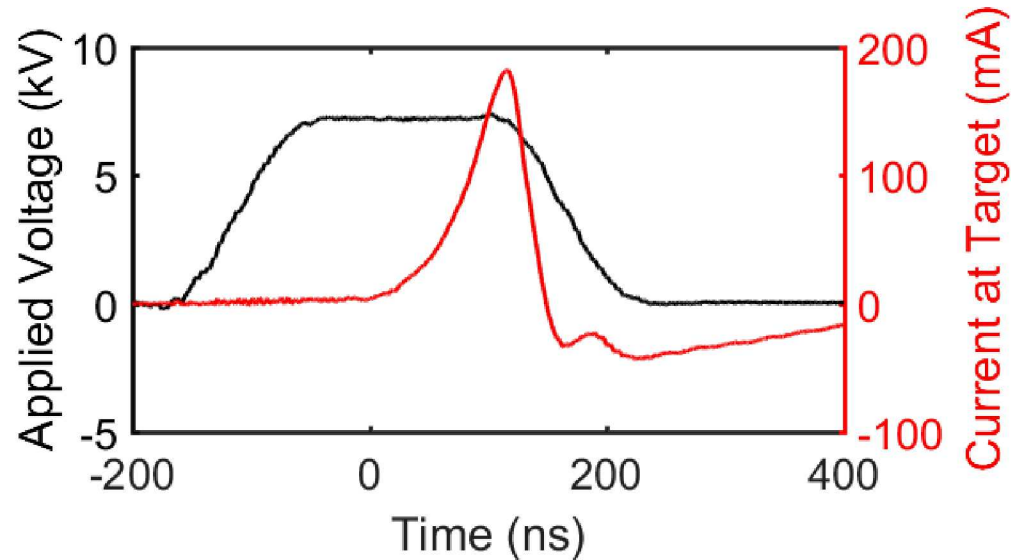
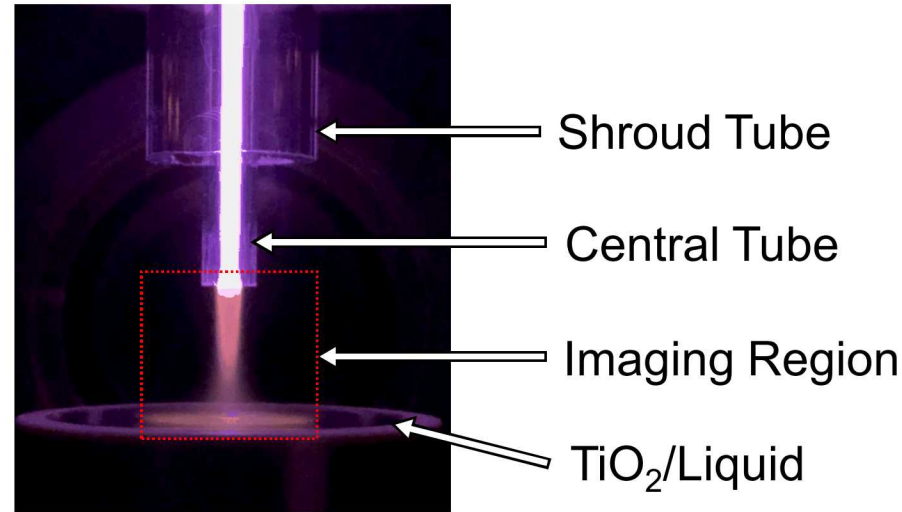


- Barnat and Fierro, J. Phys. D: Appl. Phys., 50, 14LT01 (2017).

- An ultrashort pulse laser ( $<100 \text{ fs}$ ) was used to measure the electron density with high time resolution.
- Electrons collide with laser excited  $\text{He}(3^3P)$ 
  - $e + \text{He}(3^3P) \rightarrow \text{He}(3^3D) + e$
- $n_e$  is proportional to ratio of LIF signals ( $588 \text{ nm} / 389 \text{ nm}$ )
- Sufficient  $\text{He}(2^3S)$  density is critical for accurate LCIF data.
- LCIF was initially developed for pure He, and extended to mixtures for this study of APPJs.

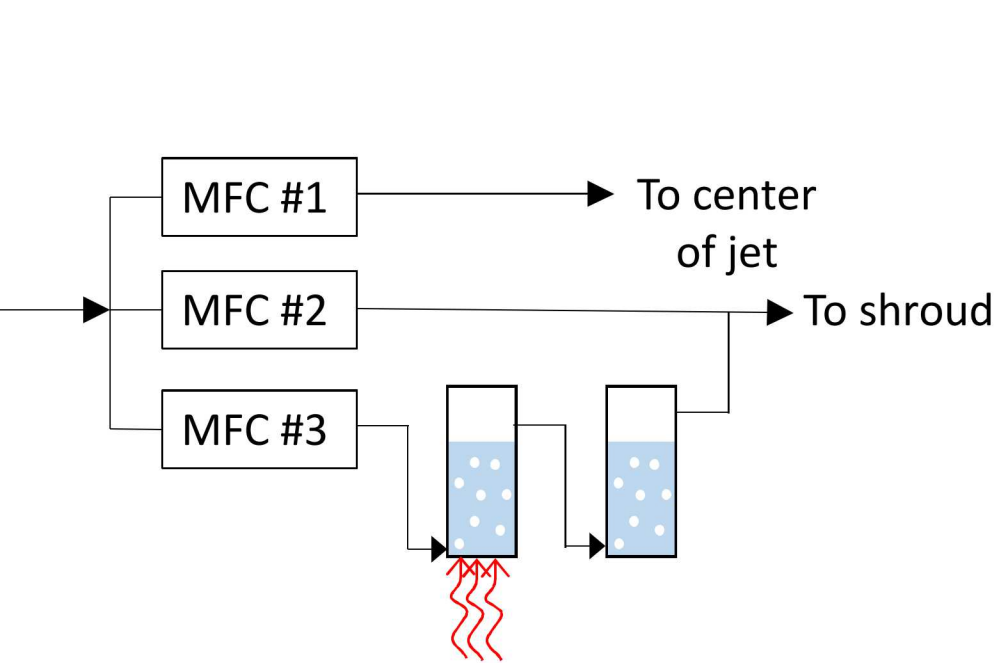
# BASE CASE

- +6 kV
- 380 ns pulse, 100 ns rise
- 200 Torr
  - Faster dynamics (for modeling)
  - Lower background LCIF signal (for experiment)
- 500 sccm He in center tube
- 0.75% H<sub>2</sub>O in He in shroud, 500 sccm
- Gap to target: 7.5 mm
- Substrate: Liquid water or TiO<sub>2</sub> 2.5 mm thick (same  $\epsilon/\epsilon_0$ )
- Current measured at ground electrode under water/TiO<sub>2</sub>.

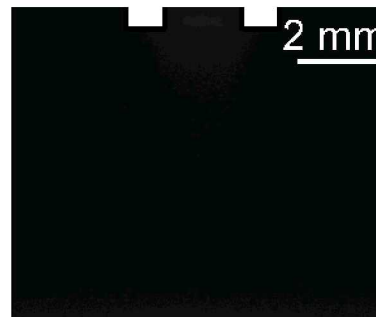
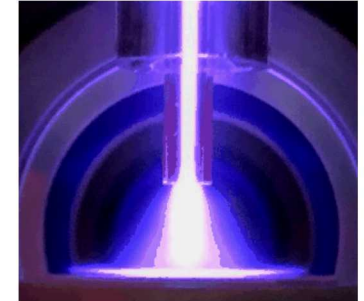
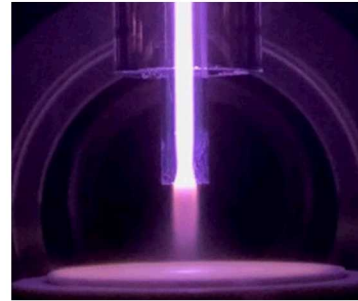




# HUMID He SHROUD



2.3% H<sub>2</sub>O in shroud    Pure He



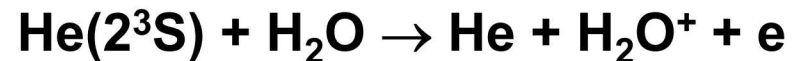
- Humid He shroud confines the jet, much like operating with surrounding air, while being compatible with LCIF measurements.
- First bubbler oversaturates water vapor, second bubbler removes excess.
- Temperature of second bubbler determines humidity of gas.

Animation Slide

University of Michigan  
Institute for Plasma Science & Engr.

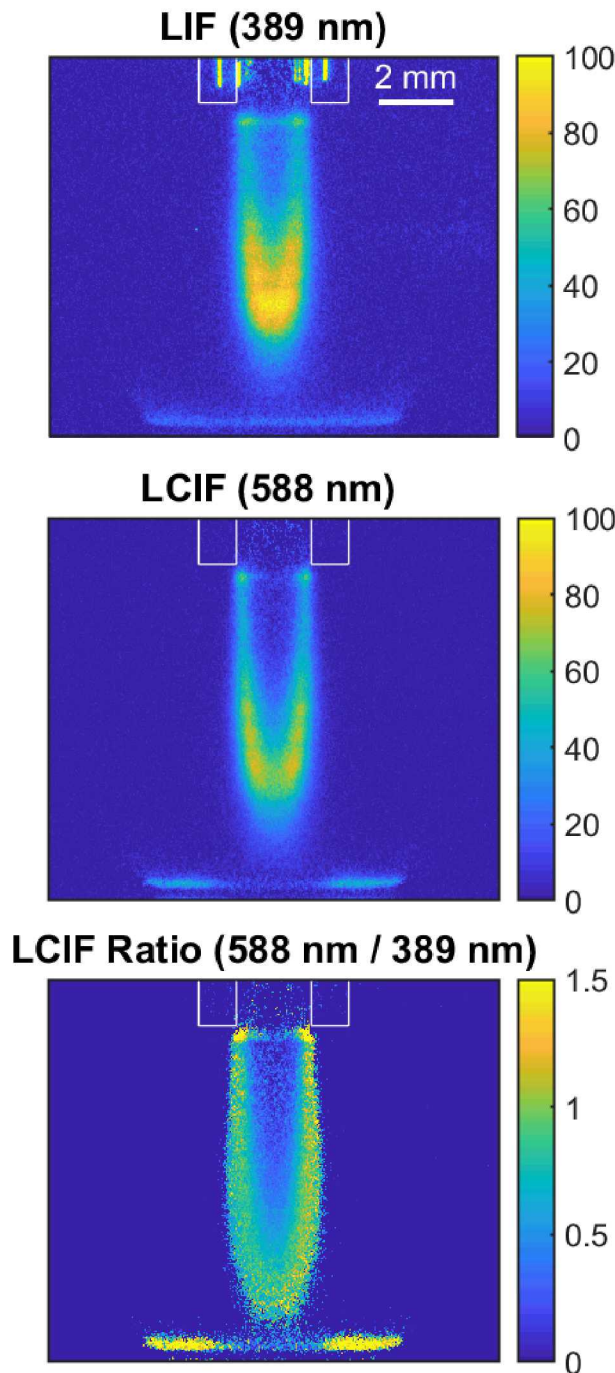
# LCIF IN HUMID He

- 6 kV, 430 ns pulse
- Center: 500 sccm He
- Shroud: 500 sccm He/H<sub>2</sub>O = 97.7/2.3
- $t = 230$  ns, 30 ns after IW contacts surface
- Moving away from He core, there are fewer He(2<sup>3</sup>S), LIF signal decreases.



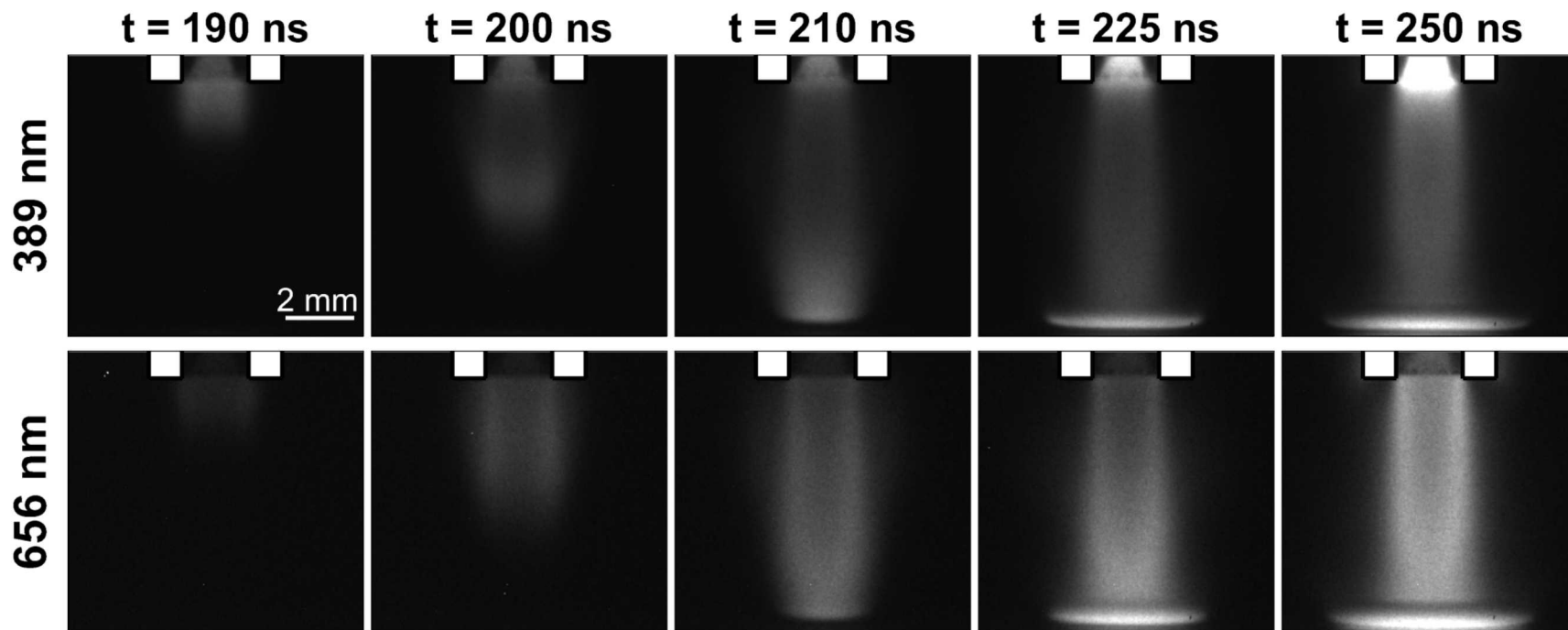
- In regions of high H<sub>2</sub>O concentration, there may be significant  $n_e$  which is not detectable due to low He(2<sup>3</sup>S).

- $\text{Ratio} = 1 \rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$



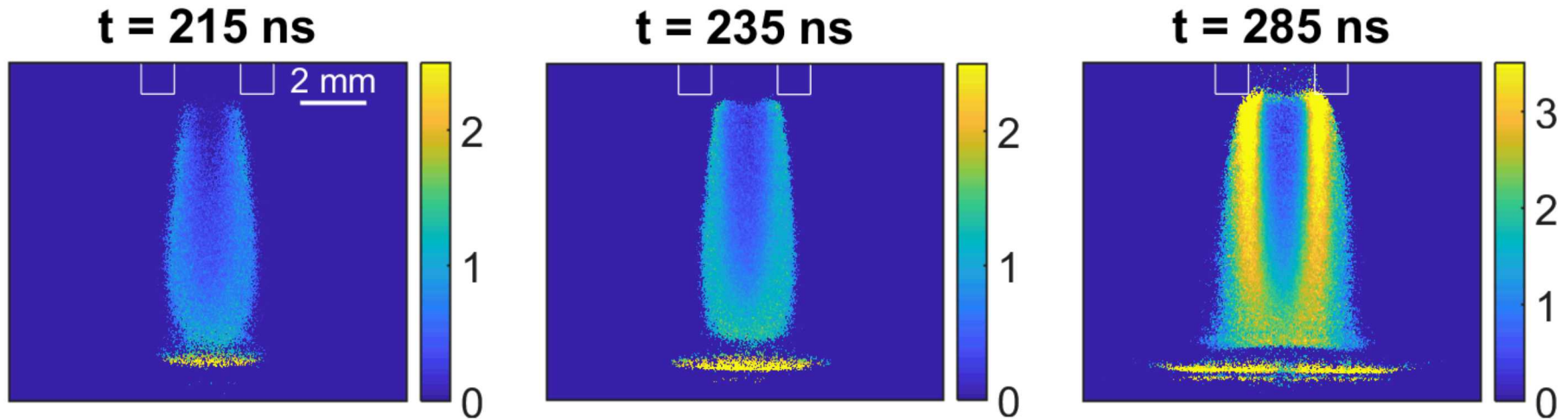


# TiO<sub>2</sub>: OPTICAL EMISSION



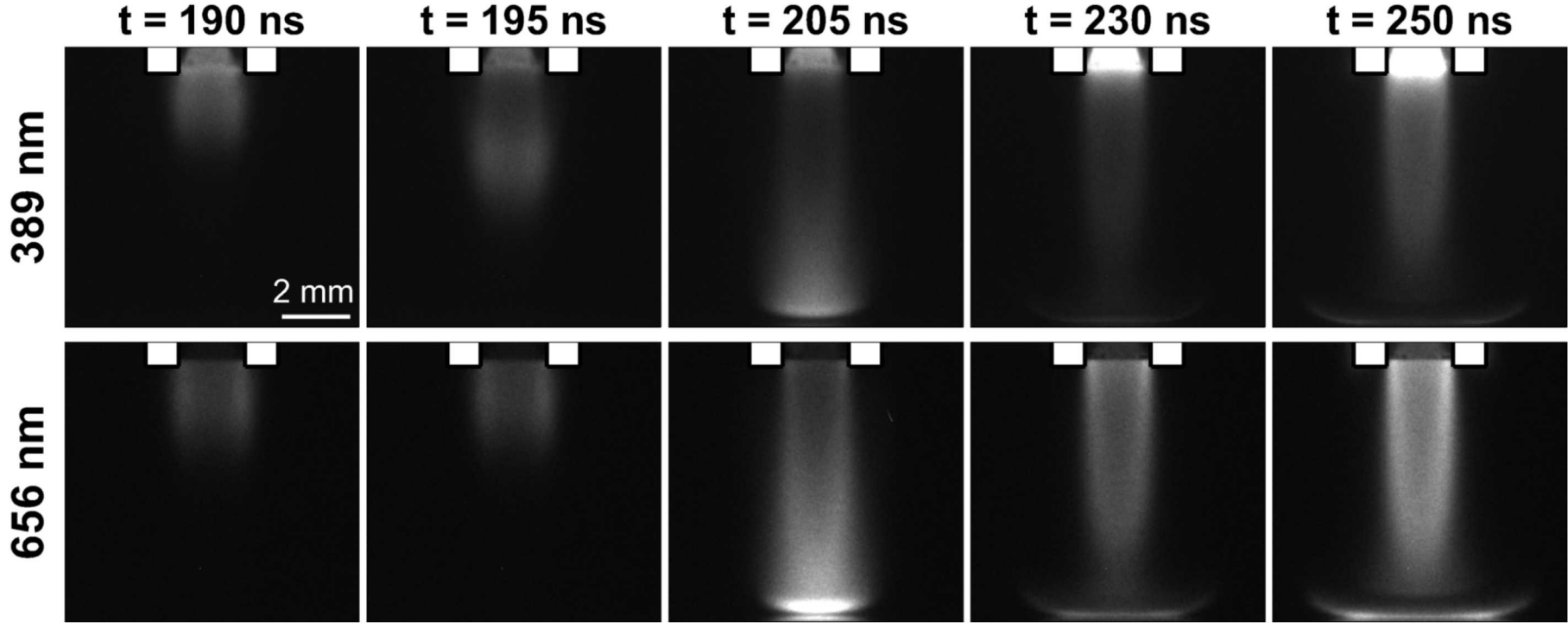
- TiO<sub>2</sub> substrate
- 5 ns ICCD gate, 656 nm – H <sub>$\alpha$</sub>  emission, 389 nm – He(3<sup>3</sup>P) → He(2<sup>3</sup>S)
- Line of sight imaging (not Abel inverted).
- Even He(3<sup>3</sup>P) emission, indicating an annular electron density.
- Surface ionization wave (SIW) forms after IW contacts the surface.

# TiO<sub>2</sub>: ELECTRON DENSITY



- Electron density profile is annular.
  - Photons from the discharge ( $\text{He}^*$  or  $\text{He}_2^*$ ) ionize  $\text{H}_2\text{O}$ .
  - Photoionization selectively occurs in the mixing region of the pure He with the humid He.
  - After IW contacts the surface and a restrike occurs, the electron density nearly doubles.
- Ratio = 1  $\rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$

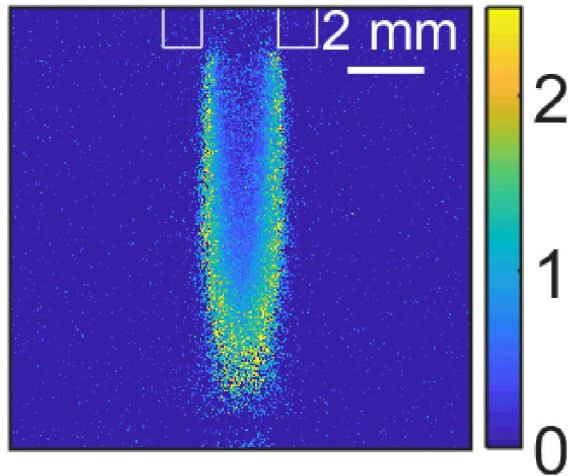
# WATER: OPTICAL EMISSION



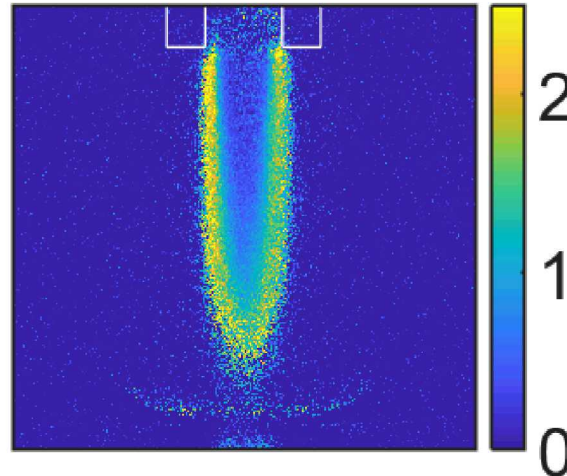
- Water substrate. 5 ns ICCD gate
- Evaporation results in much higher concentrations of  $\text{H}_2\text{O}$  enabling more photoionization and Penning ionization.
- IW contacts surface 5 ns earlier than for  $\text{TiO}_2$  substrate.
- Plasma dims at 230 ns, after contacting the surface.

# WATER: ELECTRON DENSITY

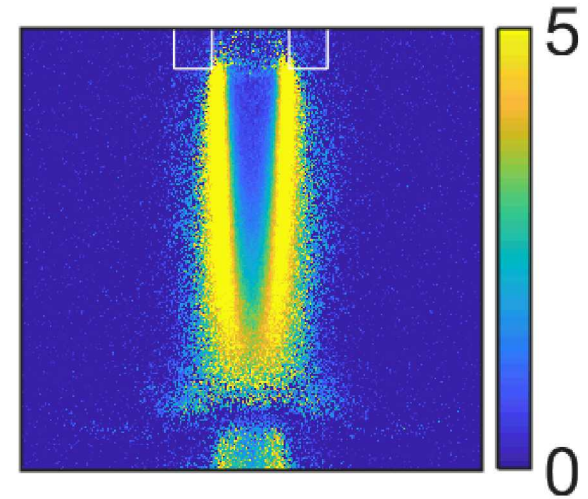
**t = 215 ns**



**t = 235 ns**



**t = 280 ns**



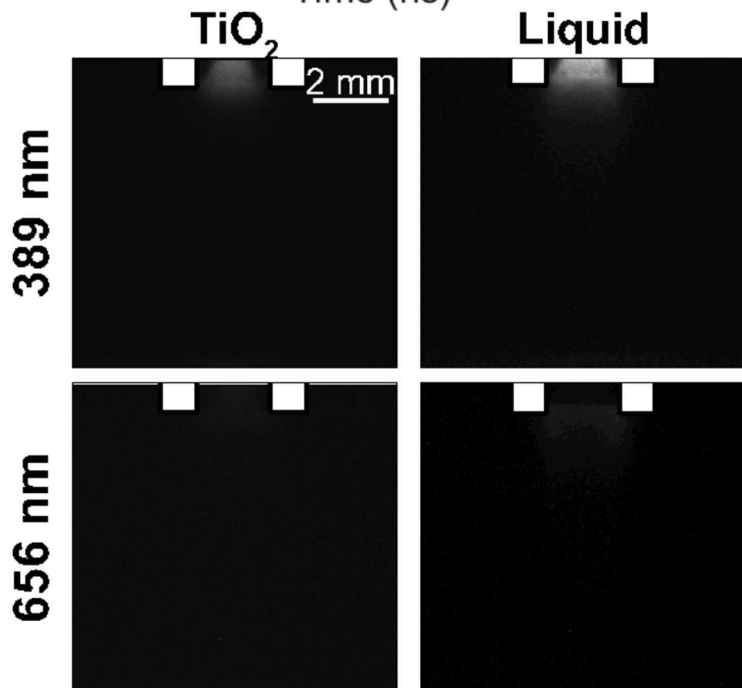
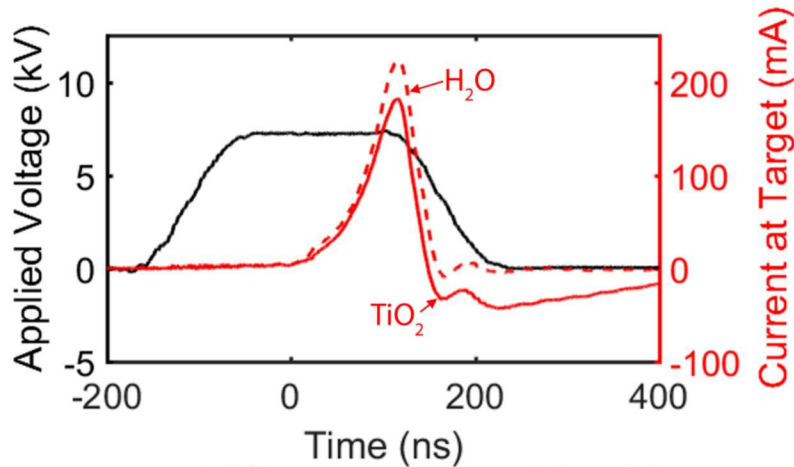
- Water substrate.
- Electron density ( $n_e$ ) is again annular due to photoionization and Penning ionization of surrounding water vapor.
- SIW is not visible by LCIF because high  $\text{H}_2\text{O}$  density near the surface depletes  $\text{He}(2^3\text{S})$ .



- Electron density is slightly higher than for  $\text{TiO}_2$  because of higher water vapor density.

- Ratio = 1  $\rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$

# WATER vs. $\text{TiO}_2$

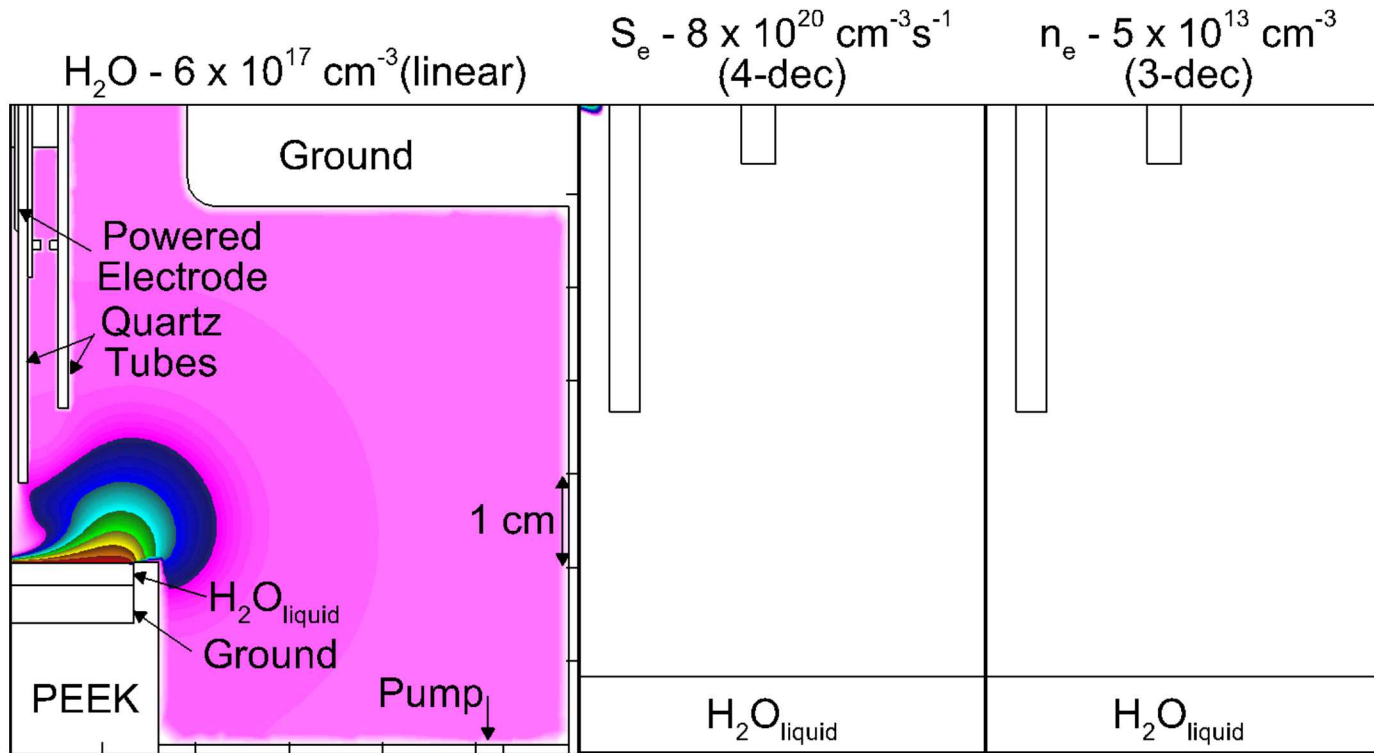


Animation Slide

- With water base, magnitude of forward current is higher, reverse current is lower.
- Forward current enhanced by humidity.
- Electrons/ions solvate into the water while being physisorbed on  $\text{TiO}_2$ .
- Water less likely to support electron emission to support reverse of current.
- Evaporation produces higher water concentrations near surface resulting in lower mobility, thinner SIW.
- Plasma dims after IW contacts liquid surface.



# MODELING LIQUID INTERACTIONS



- **Cylindrically symmetric.**
- **Modeling work ongoing.**
- **Same parameters as experiment.**

- ***nonPDPSIM* – 2-dimensional plasma hydrodynamics model**
- **Photoionization and photoemission from surfaces are critical in positive IW propagation.**
- **Surface ionization wave develops along the liquid surface.**
- **Electron density profile is annular.**

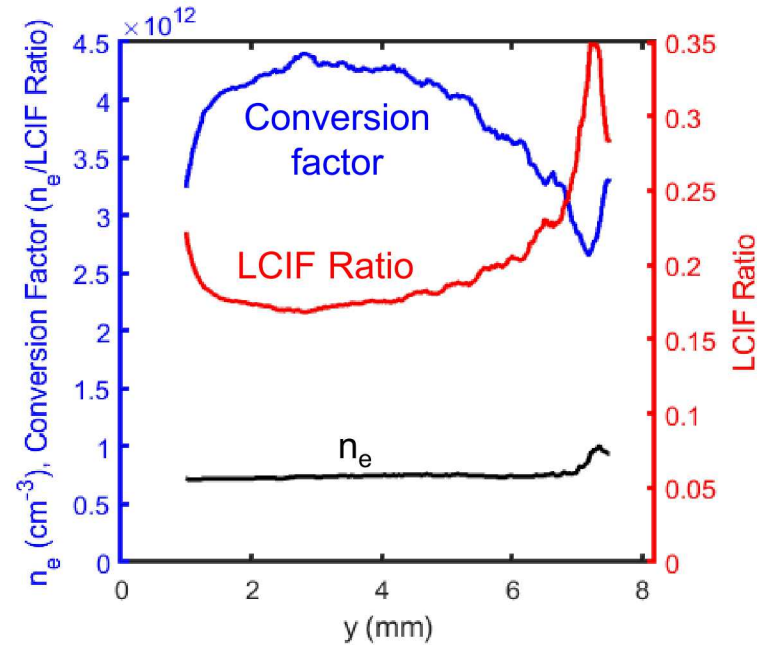
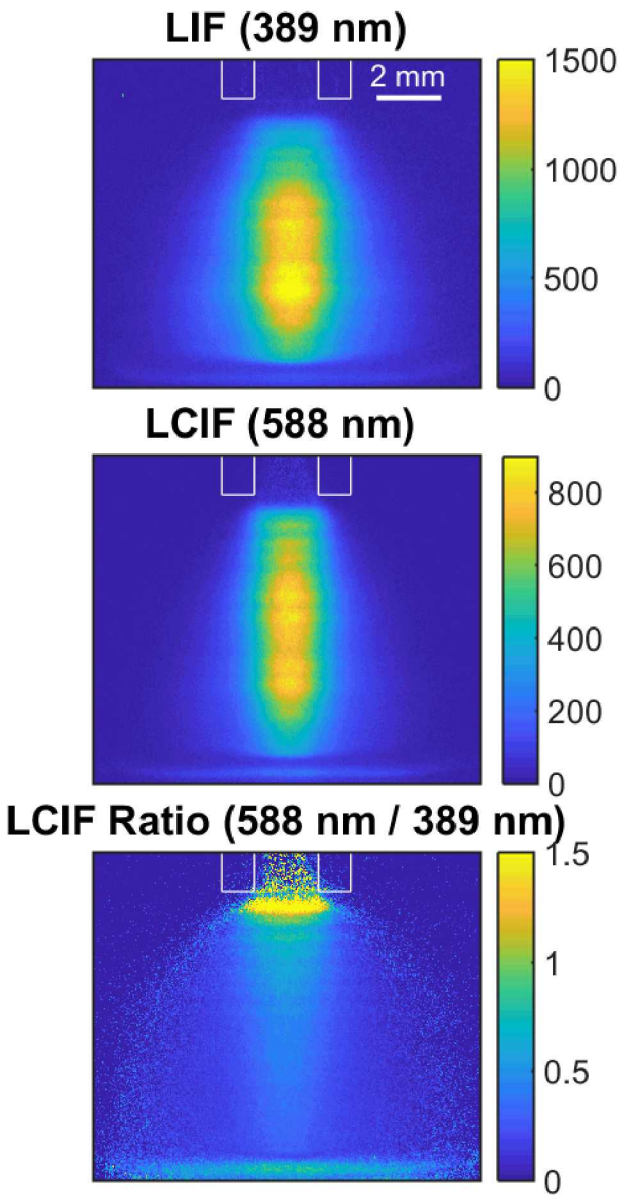
# CONCLUDING REMARKS

---

- An ionization wave exiting a He plasma jet into a humid environment is annular due to Penning ionization and photoionization at the boundary
- Plasma in contact with water vs  $\text{TiO}_2$  (same  $\epsilon/\epsilon_0$ ) are distinguishable from electrical I-V traces.
  - Mild humidity above the water produces larger forward current.
  - Solvation of electrons/ions in water and larger humidity above water result in lower reverse current.
- SIWs propagating over a liquid surface are thinner due to the lower mobility of electrons in a saturated water vapor environment.

# Appendix

# CONVERT LCIF TO $n_e$

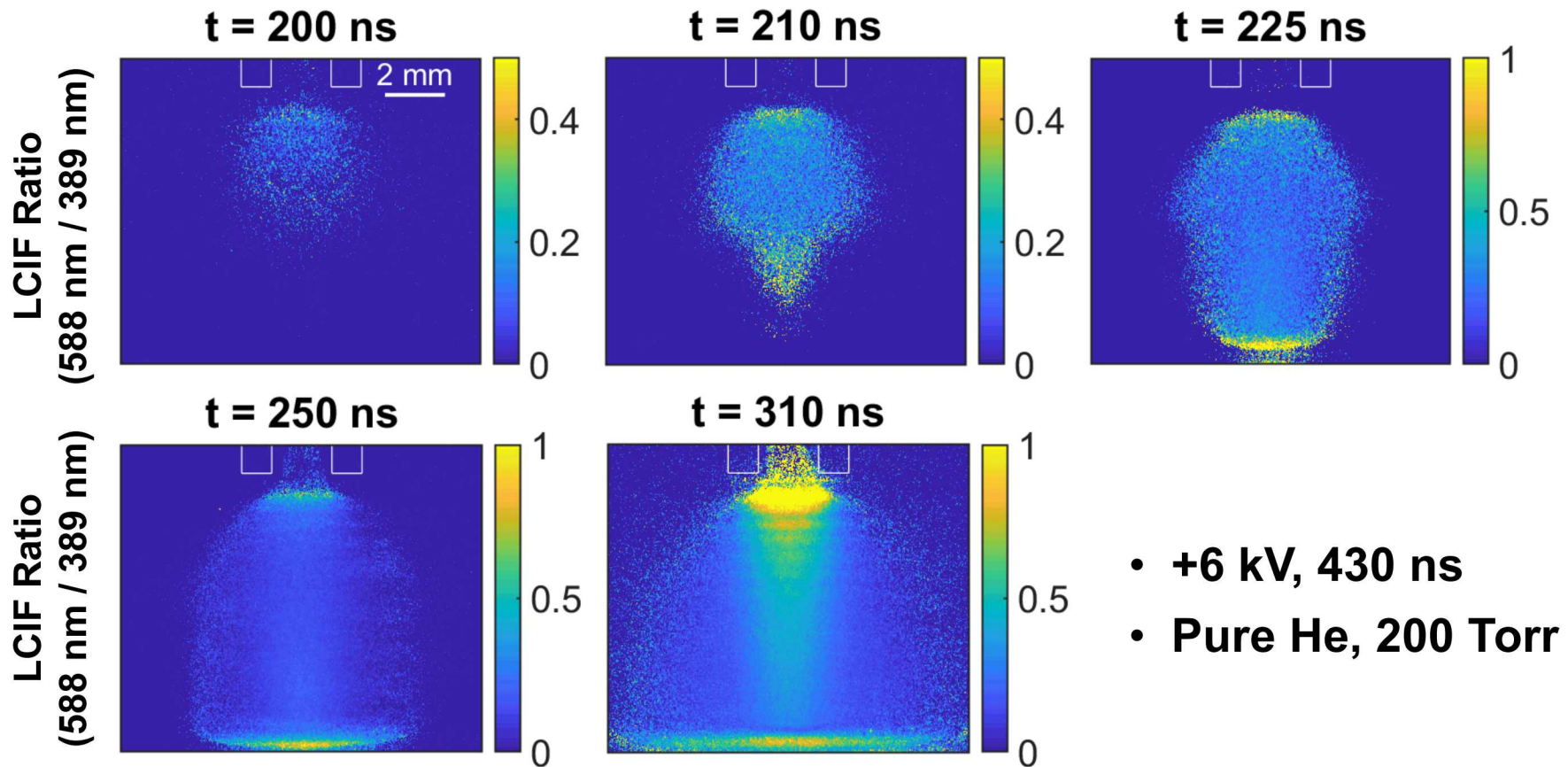


- Calculate  $n_e$  using Ohm's law and  $E/N$ .

$$I = \frac{e^2}{m_e} \frac{N}{v_m} \frac{E}{N} n_e A$$

- An LCIF ratio of 1, is approximately  $4 \times 10^{12} \text{ cm}^{-3}$  electrons.
- Previously, conversion factor estimated at  $1.5 \times 10^{13} \text{ cm}^{-3}$  at 600 Torr.

# BASE CASE LCIF



- Before IW reaches the surface, He( $2^3S$ ) densities are low.
- $n_e$  in the SIW is nearly double that of the bulk.
- Elevated  $n_e$  in IW front may be due to Stark mixing.

$$\text{Ratio} = 1 \rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$$



# VARY PRESSURE – EMISSION

---

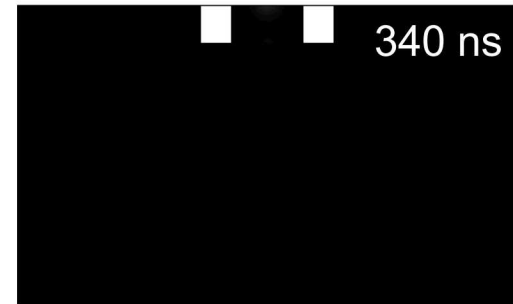
150 Torr



200 Torr (Base)



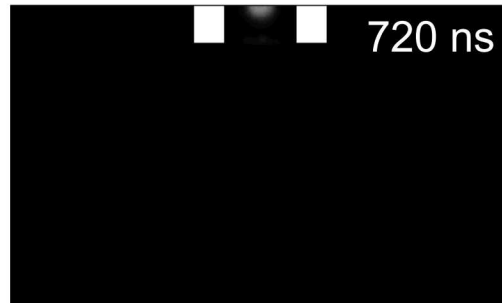
300 Torr



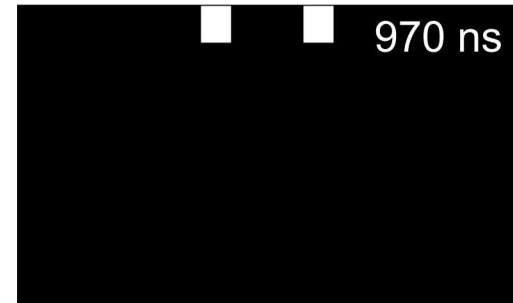
400 Torr



500 Torr



600 Torr



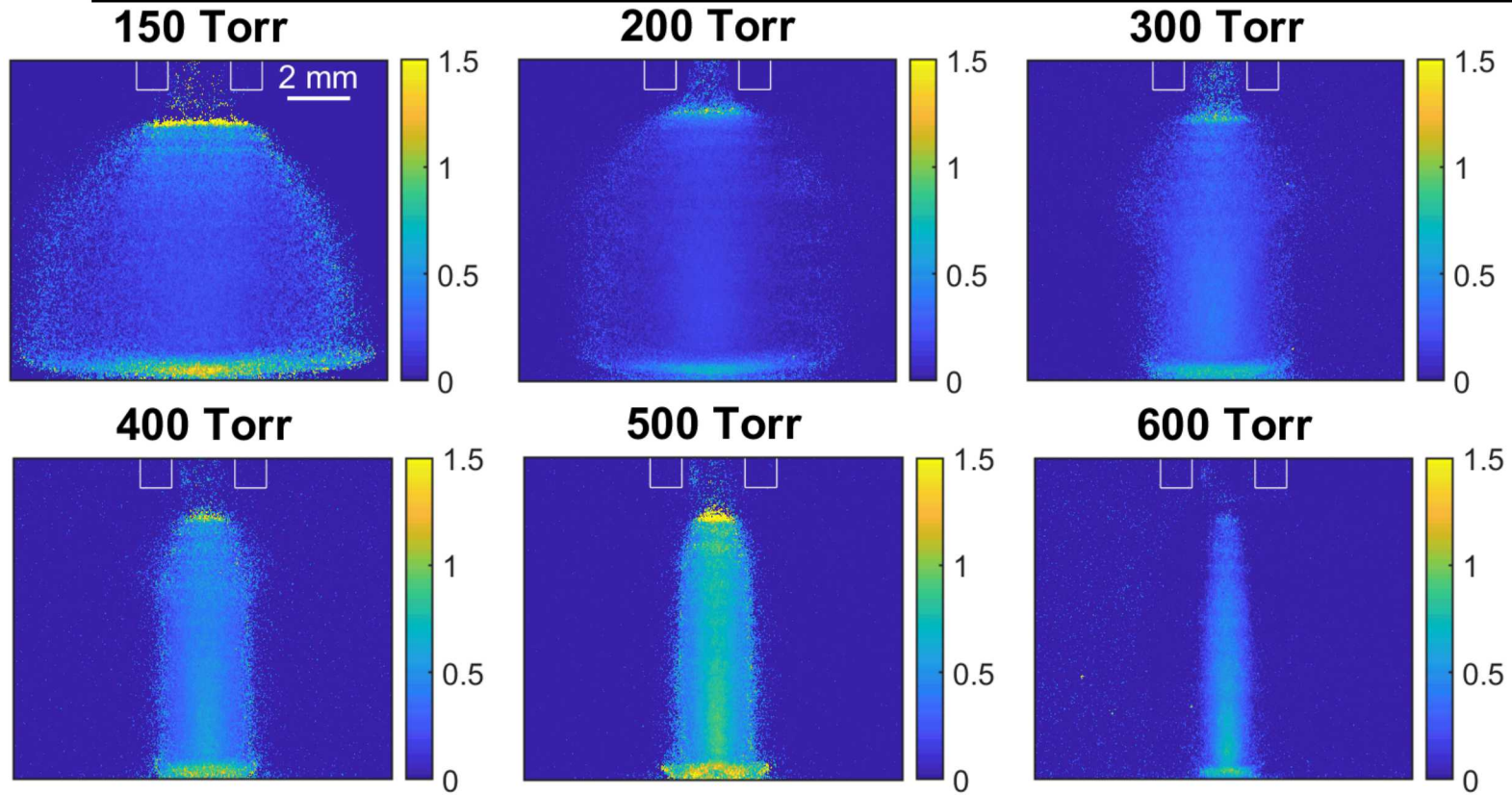
- 6 kV, 500 sccm He, 390 nm plasma emission.
- 350 ns animated, time of first frame indicated.
- Varied pulse duration so voltage is on for 80 ns after contact.
- IW propagates slower and SIW becomes thinner for higher pressures. (lower electron mobility)

Animation Slide

---

University of Michigan  
Institute for Plasma Science & Engr.

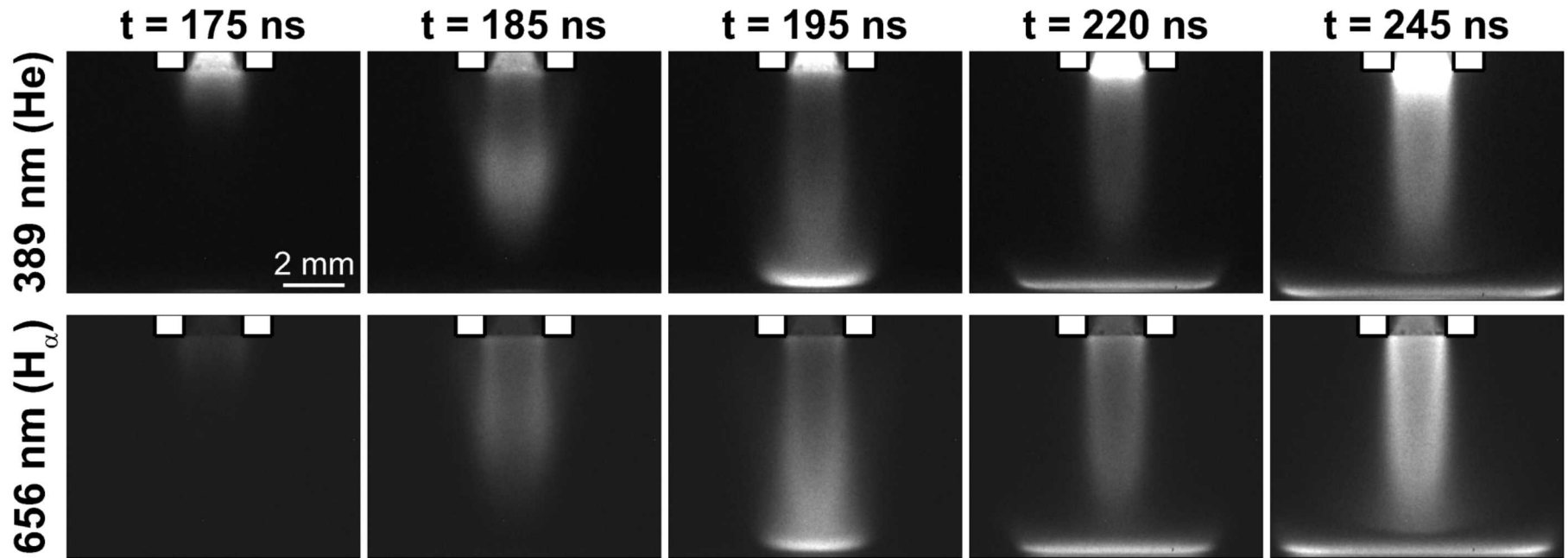
# VARY PRESSURE – $n_e$



- 6 kV, 500 sccm He, 30 ns after IW contacts surface.
- Plasma is more confined at higher pressure,  $n_e$  increases.
- Current and energy deposition decrease with increasing pressure.
- Above 500 Torr,  $n_e$  is collisional enough that ionization rate drops.

$$\text{Ratio} = 1 \rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$$

# HUMID He SHROUD EMISSION



- 2.3% H<sub>2</sub>O in shroud.
- Images have not been Abel inverted.
- IW reaches outlet of the tube earlier than in base case – photoionization from He<sub>2</sub><sup>\*</sup> causes non-local seed ionization.
- Photoionization and Penning ionization promote IW speed.
- H $\alpha$  emission appears more annular – dominates at the interface of the center and shroud flow.

Animation Slide

University of Michigan  
Institute for Plasma Science & Engr.

# HUMID He SHROUD

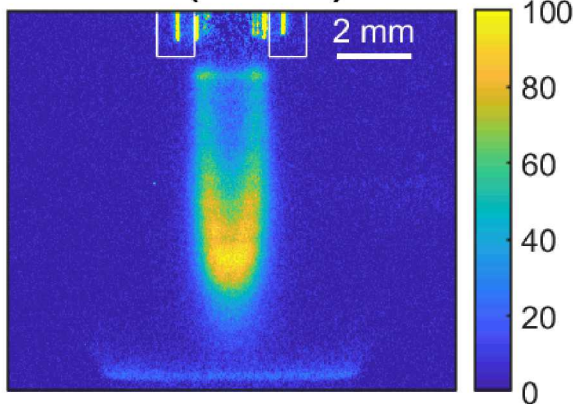
- 6 kV, 430 ns pulse
- Center: 500 sccm He
- Shroud: 500 sccm He/H<sub>2</sub>O = 97.7/2.3
- $t = 230$  ns, 30 ns after IW contacts surface
- Moving away from He core, there are fewer He(2<sup>3</sup>S), LIF signal decreases.  
$$\text{He}(2^3\text{S}) + \text{H}_2\text{O} \rightarrow \text{He} + \text{H}_2\text{O}^+ + e$$
- In regions of high H<sub>2</sub>O concentration, there may be significant  $n_e$  which is not detectable due to low He(2<sup>3</sup>S).

$$\text{Ratio} = 1 \rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$$

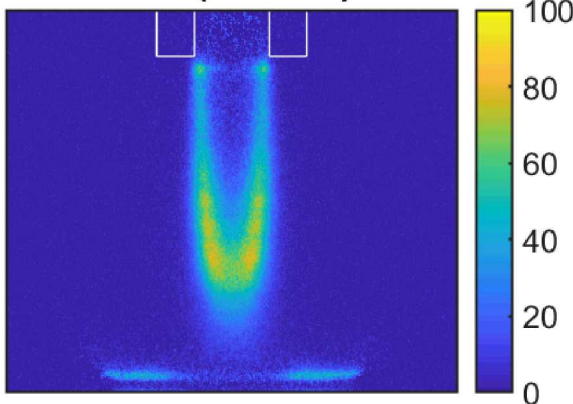
---

University of Michigan  
Institute for Plasma Science & Engr.

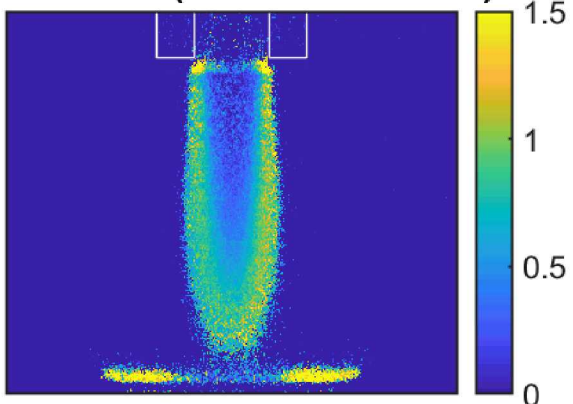
LIF (389 nm)



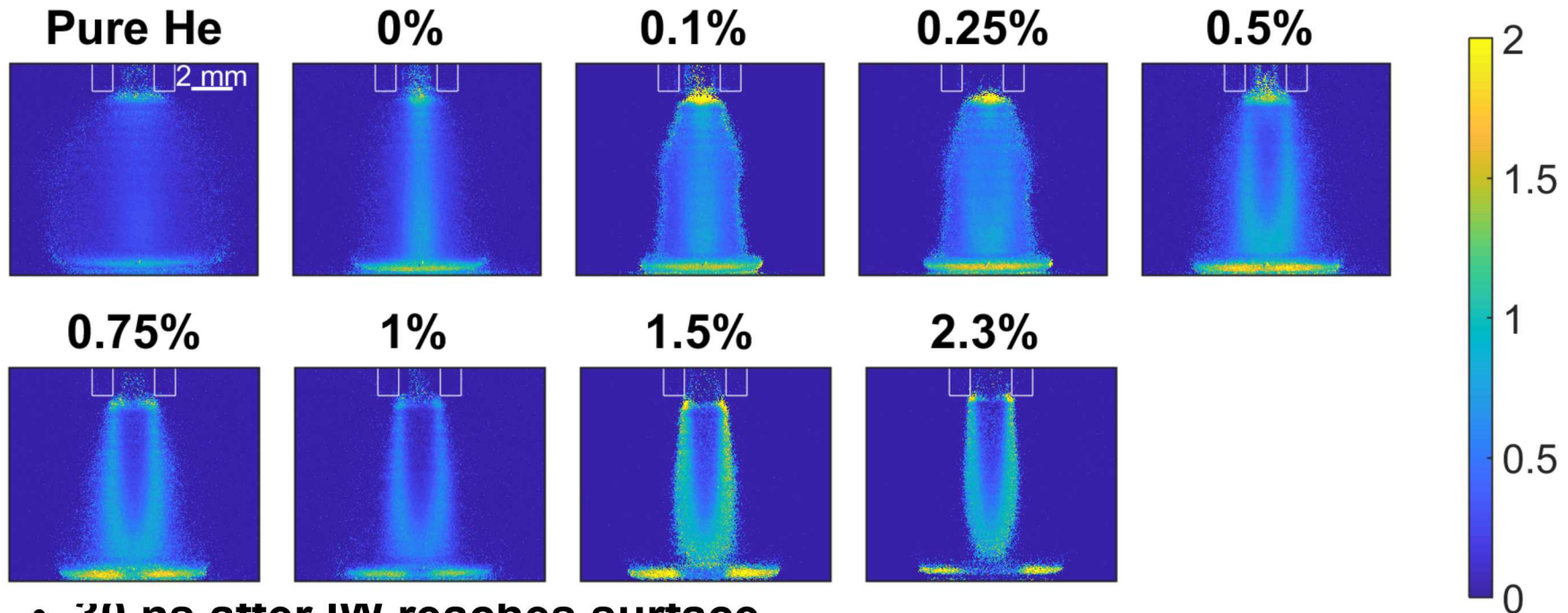
LCIF (588 nm)



LCIF Ratio (588 nm / 389 nm)



# SHROUD HUMIDITY

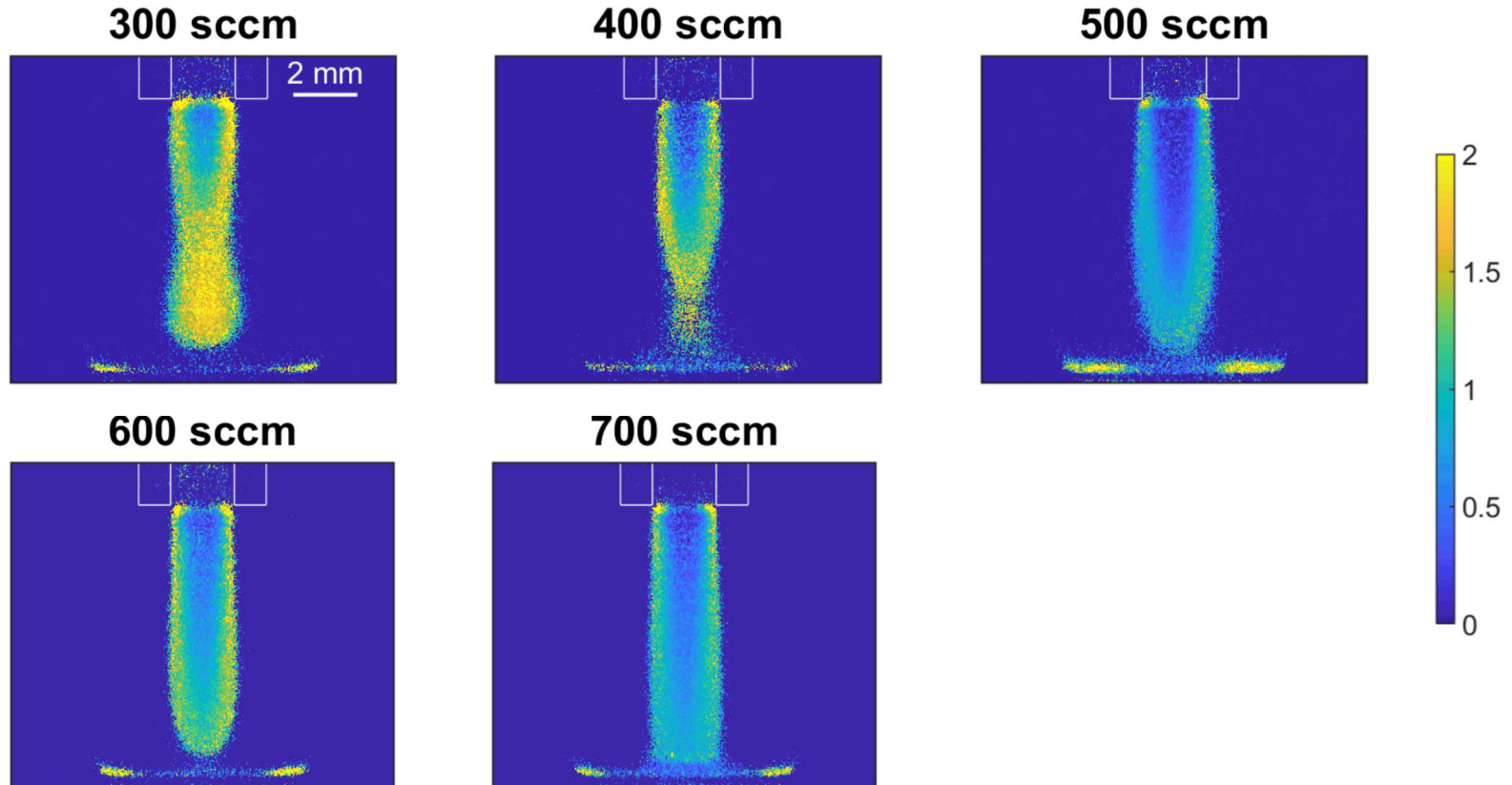


- 30 ns after IW reaches surface.
- Transition from diffuse in pure helium case to confined by humid shroud.
- Higher electron energy loss rates with  $\text{H}_2\text{O}$  because of vibrational and rotational excitation.
- $n_e$  increases with humidity due to Penning ionization.

$$\text{Ratio} = 1 \rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$$



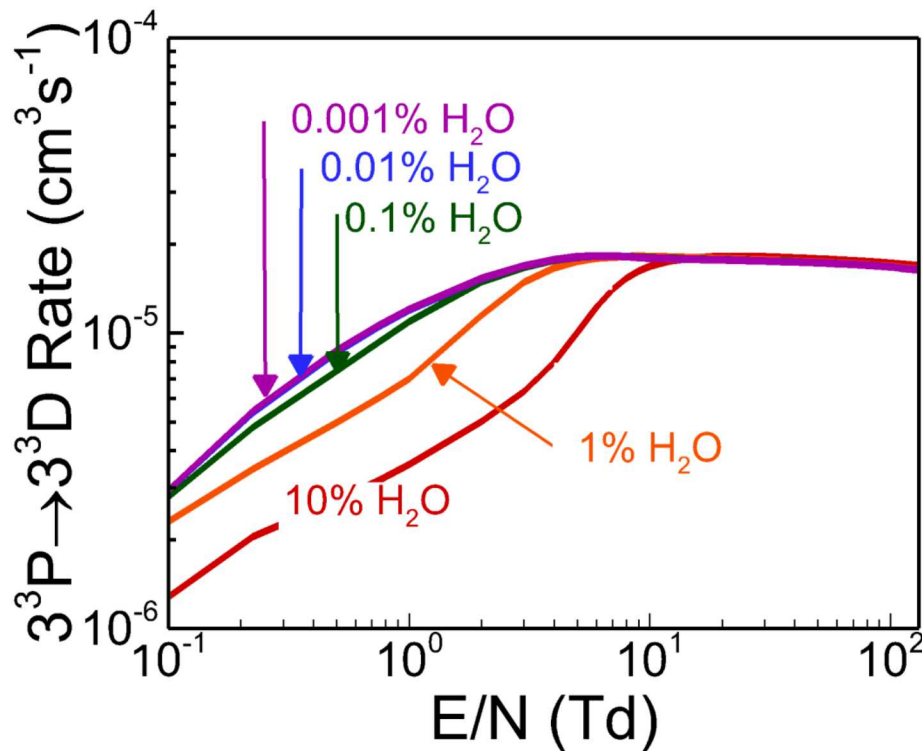
# VARY He FLOW RATE - $n_e$



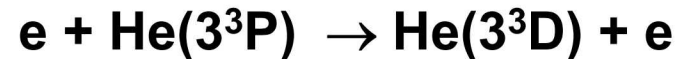
- Vary center flow rate.
- Shroud: 500 sccm,  $\text{He}/\text{H}_2\text{O} = 98.7/2.3$
- Higher He flow rates more rapidly convect in-diffusing  $\text{H}_2\text{O}$ .

$$\text{Ratio} = 1 \rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$$

# BOLTZMANN CALCULATIONS



- For LCIF measurement to be a linear representation of  $n_e$ , this rate must be independent of  $T_e$ :



- Threshold = 0.06 eV
- In pure He, this occurs when  $E/N > \sim 0.8$  Td
- For  $\text{H}_2\text{O} < 1\%$ , LCIF is valid  $\sim 1$  Td