

IONIZATION WAVE DYNAMICS OF A PLASMA JET IN CONTACT WITH LIQUID WATER*

SAND2018-12449C

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**71st Annual Gaseous Electronics Conference
Portland, OR, USA.
8 November 2018**

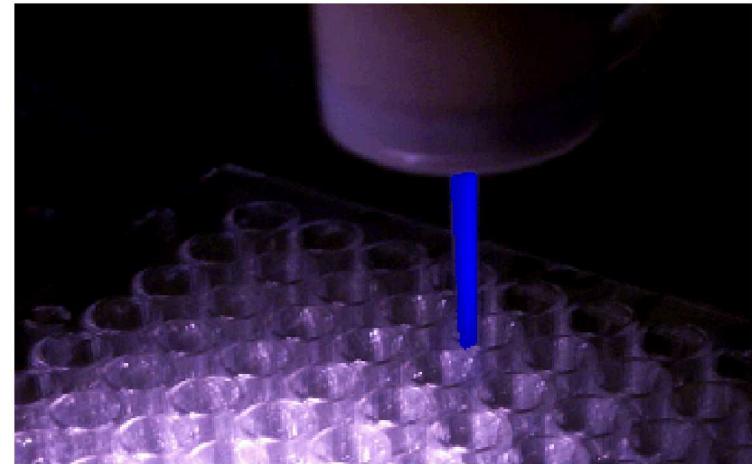
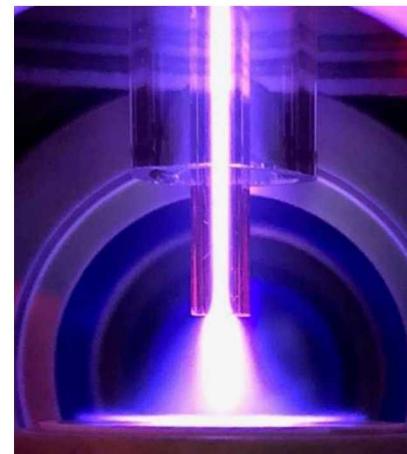
- **Work supported by the US Dept. of Energy Office of Fusion Energy Science and National Science Foundation.**
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AGENDA

- Plasma jets onto Liquids
- Description of Experiment
- Laser Collisional Induced Fluorescence (LCIF)
- Plasma Jet Contacting TiO_2 ($\epsilon_r = 80$)
- Plasma Jet Contacting Liquid Water
- Modeling and Future Work

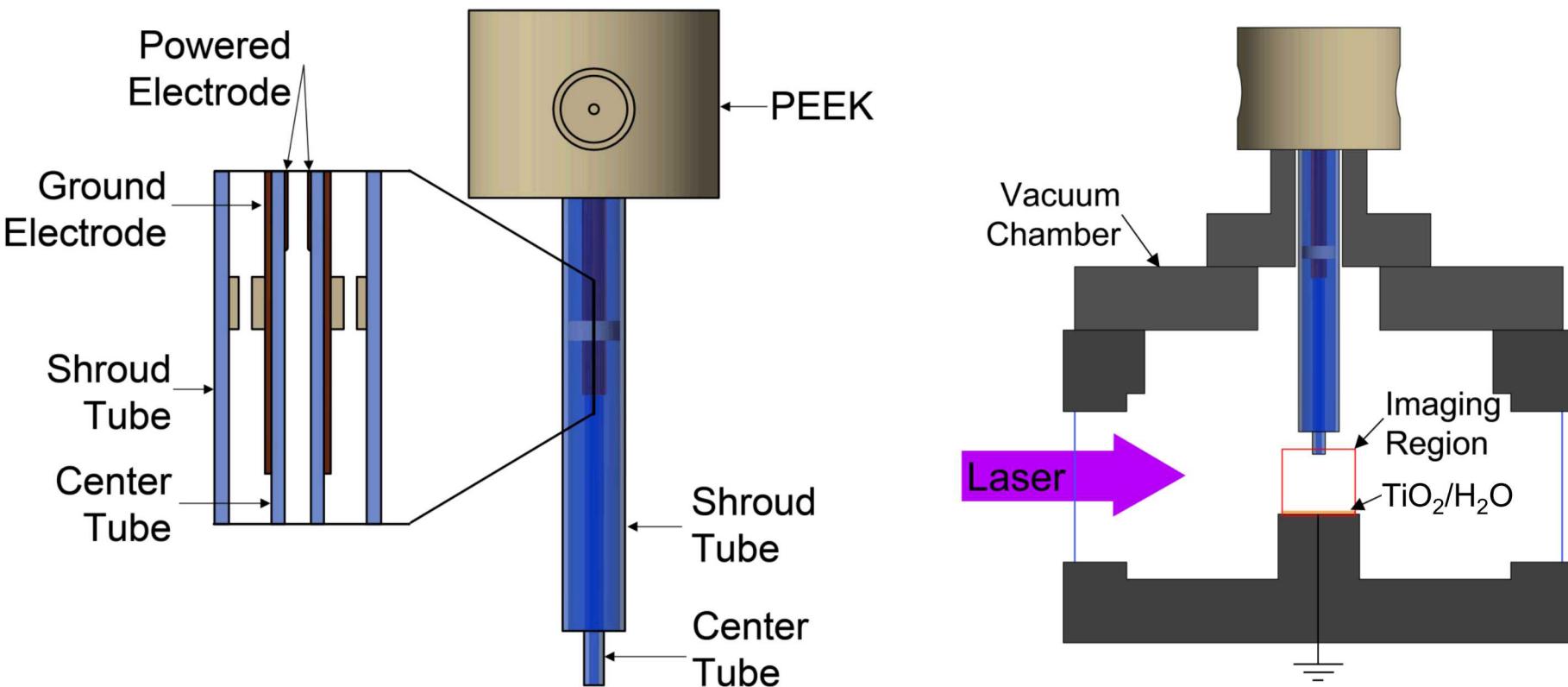
ATMOSPHERIC PRESSURE PLASMA JETS

- Atmospheric pressure plasma jets (APPJs) are a popular source of chemistry for biomedical applications.
- The plasma propagates as an ionization wave (IW) that is repetitively pulsed.
- The IW gives rise to reactive oxygen and nitrogen species (RONS) which produce the biological effect.
- **Objective:** Investigation IW dynamics in a plasma jet contacting liquid in a well controlled environment.



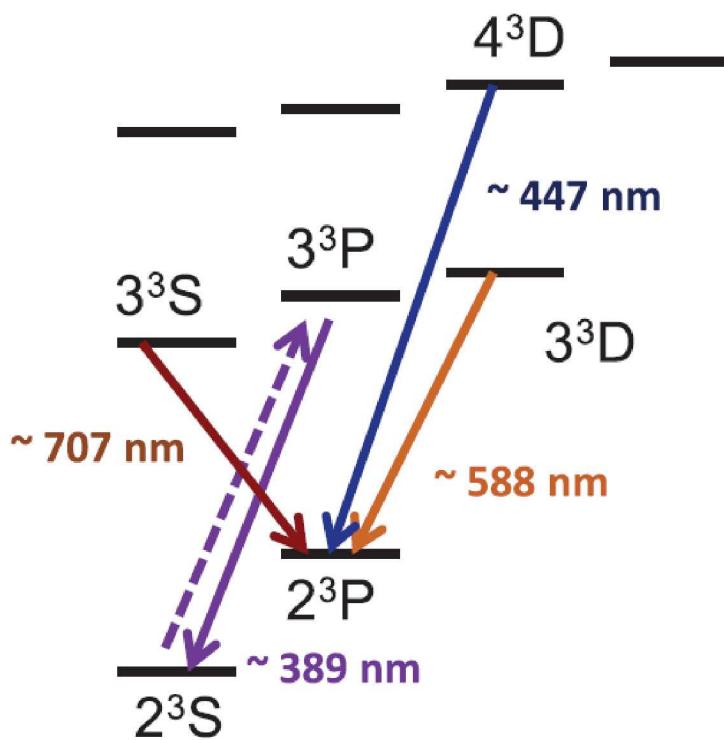
- S. Mohades, et al., Physics of Plasmas 22, 122001 (2015).

EXPERIMENTAL SETUP



- **Annular powered electrode inside the center tube.**
- **Placing the APPJ in a vacuum chamber - consistent and controlled chemistry, ground planes, and gas flow.**
- **Coaxial tube enables a gas shroud – control environment independently of gas in main jet.**

LASER COLLISIONAL INDUCED FLUORESCENCE (LCIF)

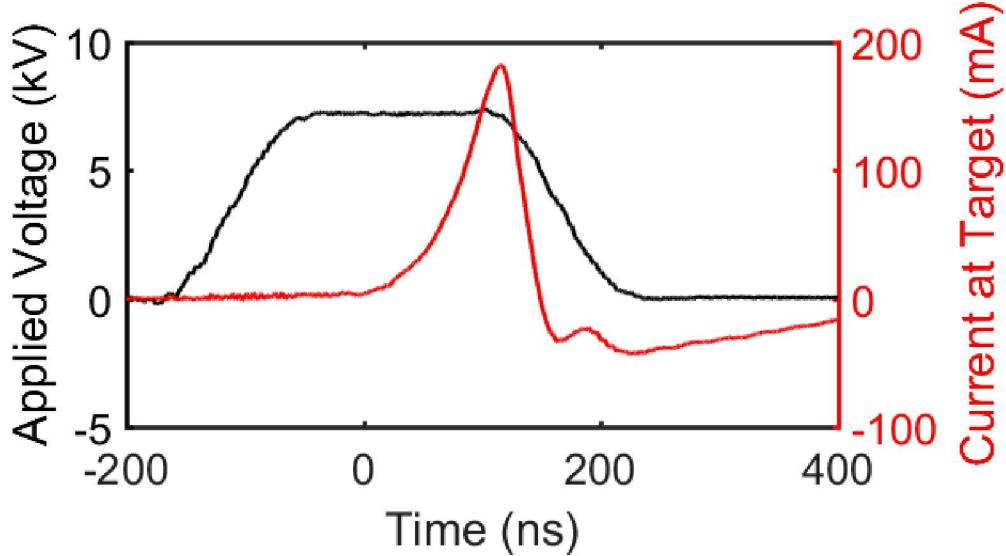
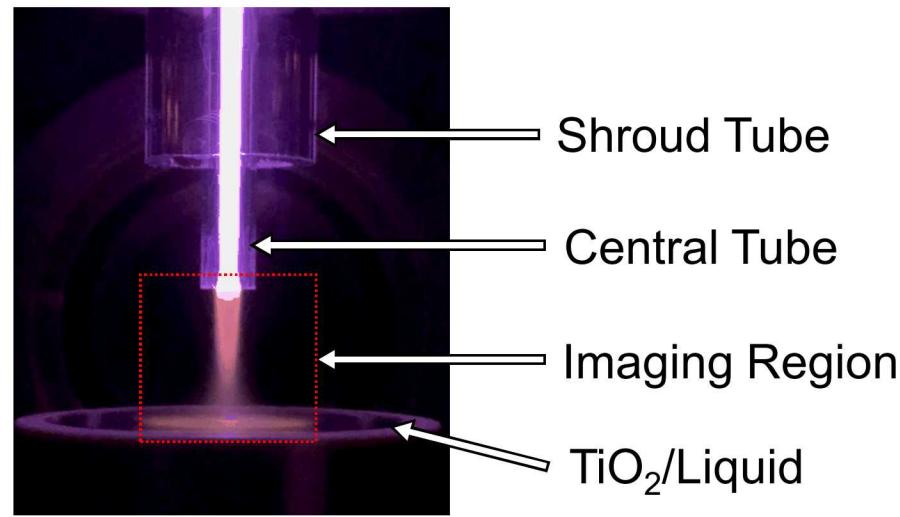


- Barnat and Fierro, J. Phys. D: Appl. Phys., 50, 14LT01 (2017).

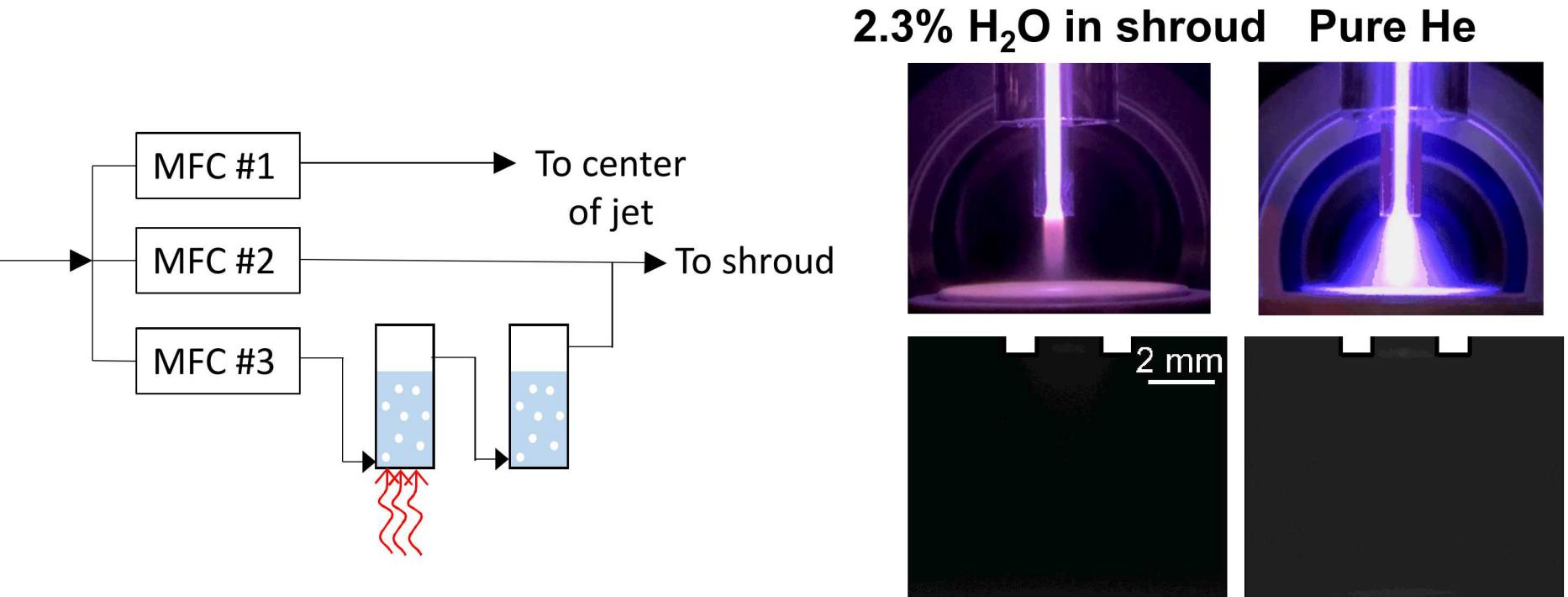
- An ultrashort pulse laser (<100 fs) was used to measure the electron density with high time resolution.
- Electrons collide with laser excited He(3^3P)
 - $e + He(3^3P) \rightarrow He(3^3D) + e$
- n_e is proportional to ratio of LIF signals (588 nm / 389 nm)
- Sufficient He(2^3S) density is critical for accurate LCIF data.
- LCIF was initially developed for pure He, and extended to mixtures for this study of APPJs.

BASE CASE

- +6 kV
- 380 ns pulse, 100 ns rise
- 200 Torr
 - Faster dynamics (for modeling)
 - Lower background LCIF signal (for experiment)
- 500 sccm He in center tube
- 0.75% H₂O in He in shroud, 500 sccm
- Gap to target: 7.5 mm
- Substrate: Liquid water or TiO₂ 2.5 mm thick (same ϵ/ϵ_0)
- Current measured at ground electrode under water/TiO₂.

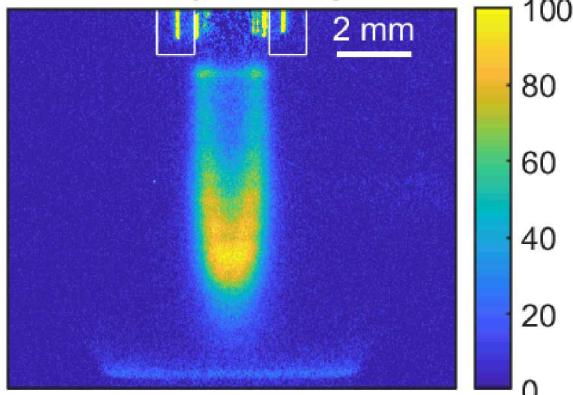


HUMID He SHROUD

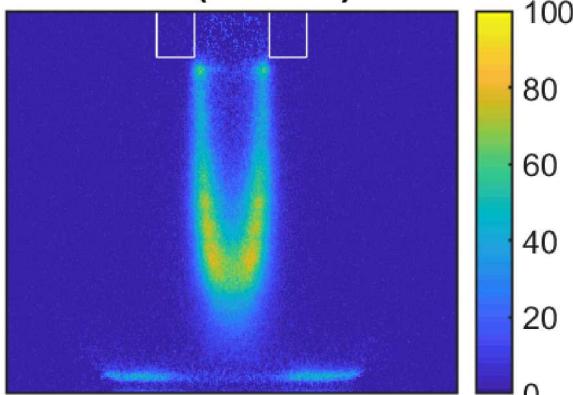


- **Humid He shroud confines the jet, much like operating with surrounding air, while being compatible with LCIF measurements.**
- **First bubbler oversaturates water vapor, second bubbler removes excess.**
- **Temperature of second bubbler determines humidity of gas.**

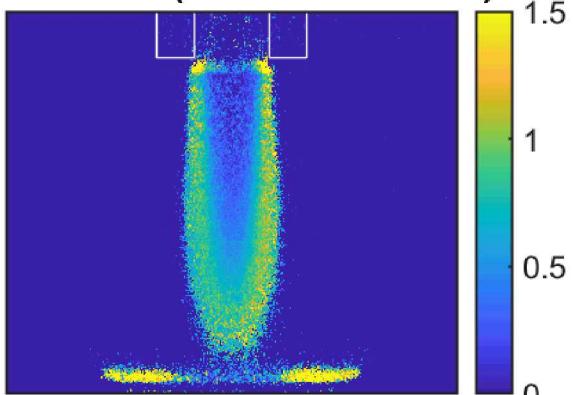
LIF (389 nm)



LCIF (588 nm)

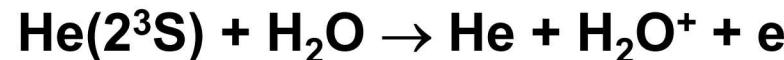


LCIF Ratio (588 nm / 389 nm)



LCIF IN HUMID He

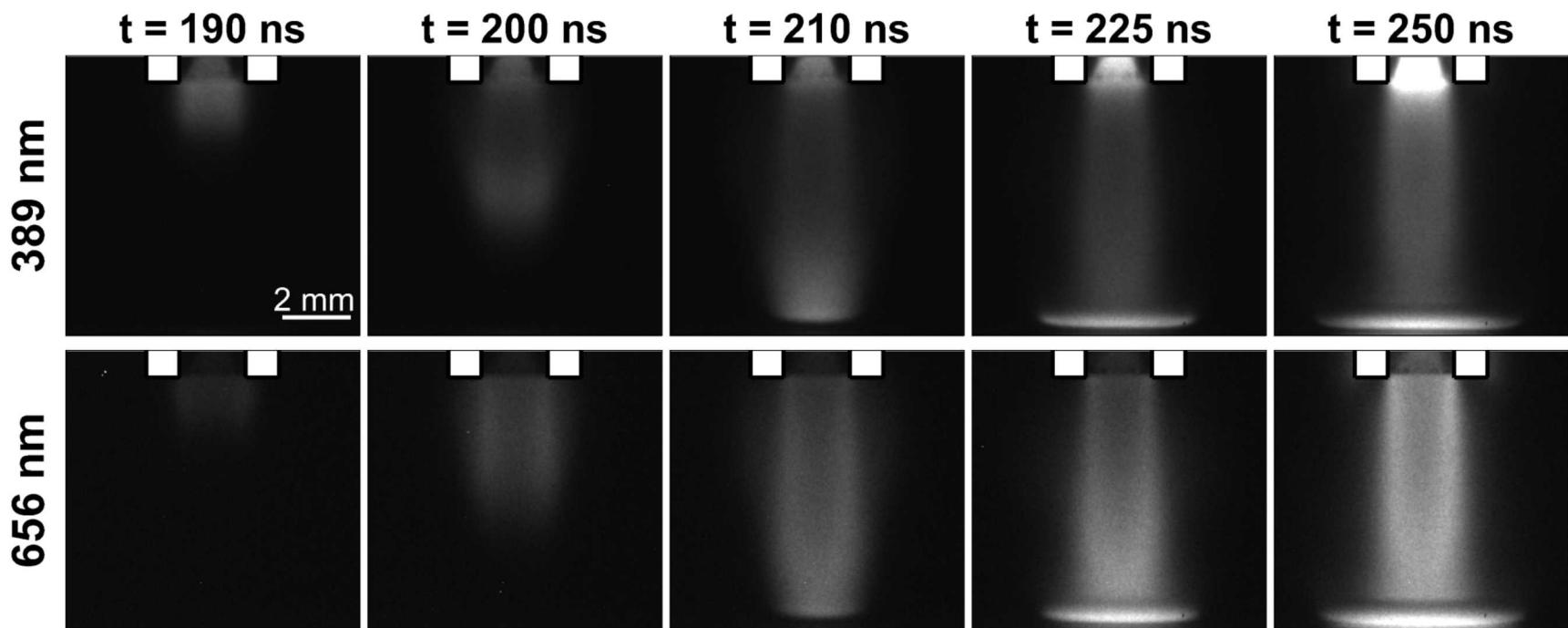
- 6 kV, 430 ns pulse
- Center: 500 sccm He
- Shroud: 500 sccm He/H₂O = 97.7/2.3
- t = 230 ns, 30 ns after IW contacts surface
- Moving away from He core, there are fewer He(2³S), LIF signal decreases.



- In regions of high H₂O concentration, there may be significant n_e which is not detectable due to low He(2³S).

- Ratio = 1 → n_e ≈ 4 × 10¹² cm⁻³

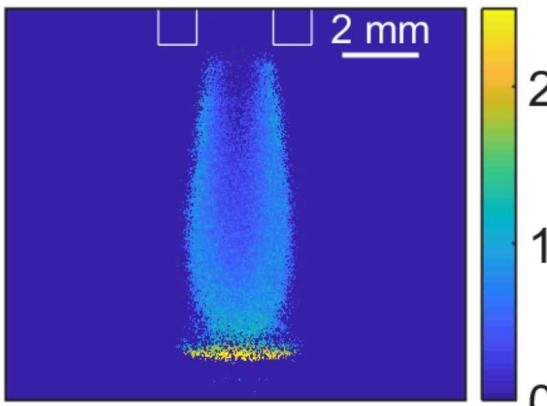
TiO₂: OPTICAL EMISSION



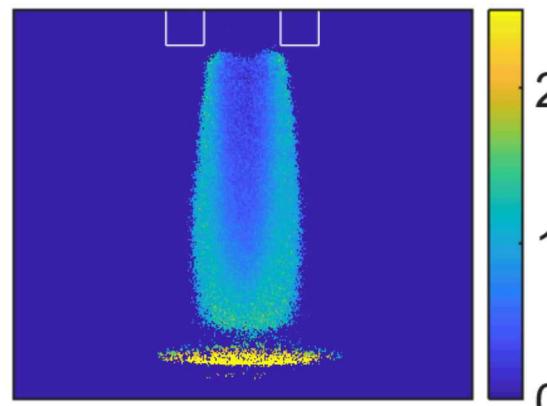
- TiO₂ substrate
- 5 ns ICCD gate, 656 nm – H_α emission, 389 nm – He(3³P) → He(2³S)
- Line of sight imaging (not Abel inverted).
- Even He(3³P) emission, indicating an annular electron density.
- Surface ionization wave (SIW) forms after IW contacts the surface.

TiO₂: ELECTRON DENSITY

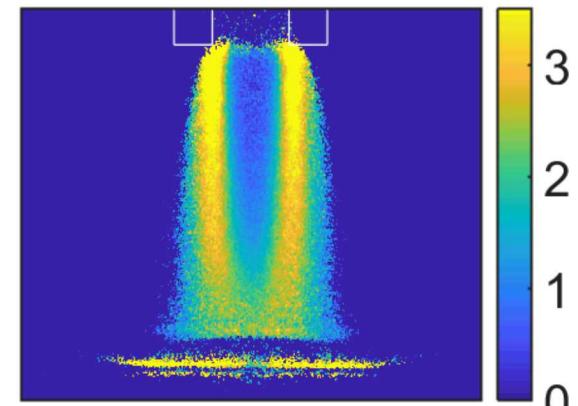
$t = 215$ ns



$t = 235$ ns

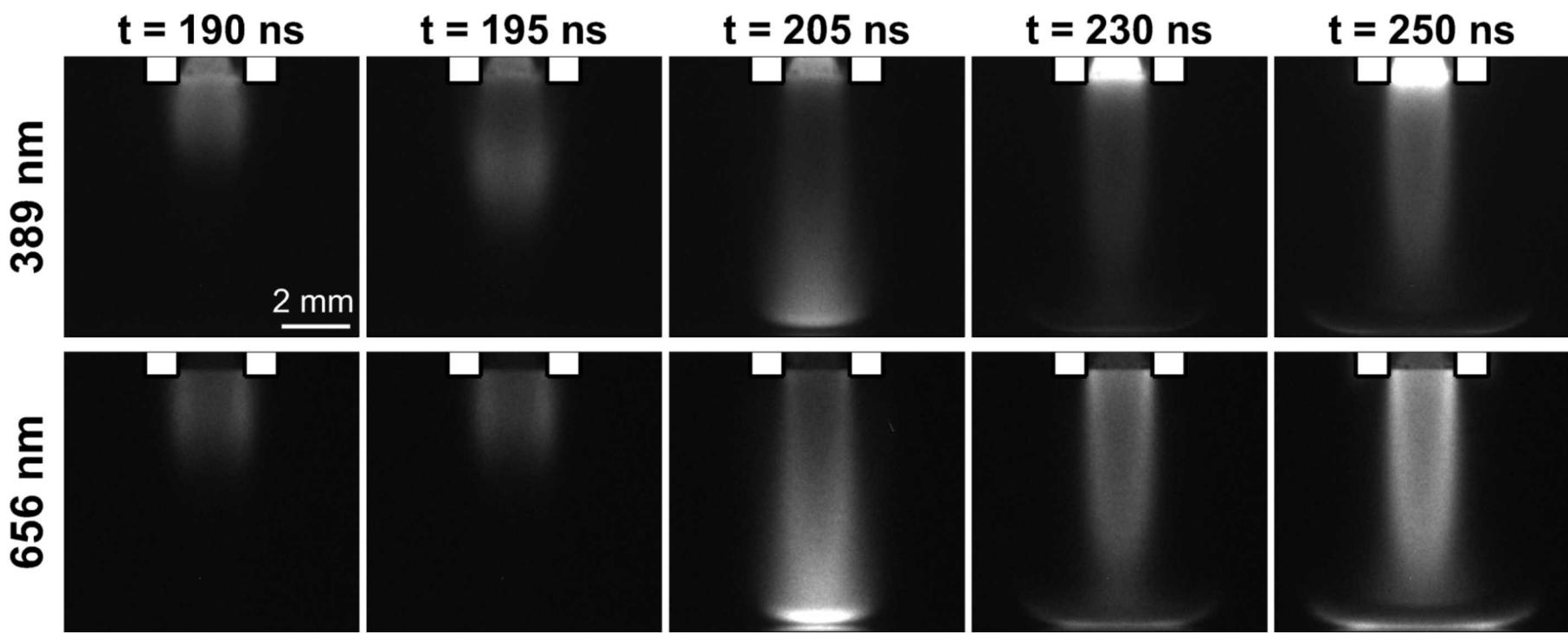


$t = 285$ ns



- Electron density profile is annular.
- Photons from the discharge (He^* or He_2^*) ionize H_2O .
- Photoionization selectively occurs in the mixing region of the pure He with the humid He.
- After IW contacts the surface and a restrike occurs, the electron density nearly doubles.
- Ratio = 1 $\rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$

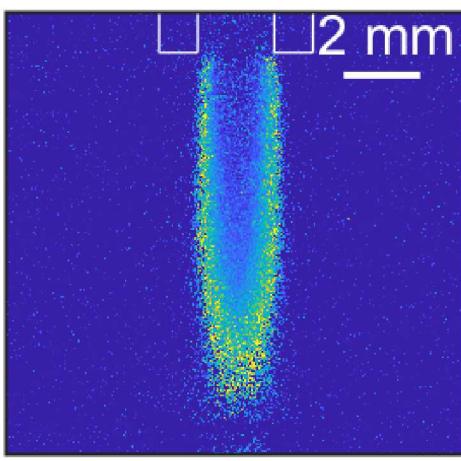
WATER: OPTICAL EMISSION



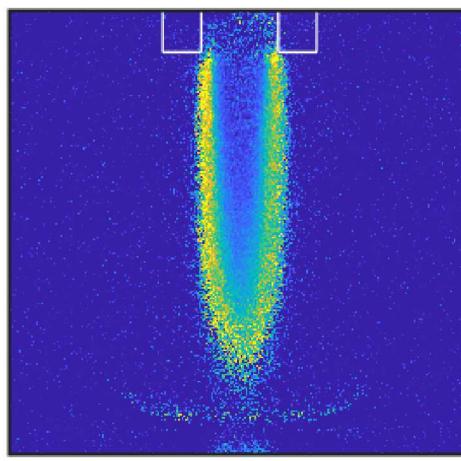
- Water substrate. 5 ns ICCD gate
- Evaporation results in much higher concentrations of H_2O enabling more photoionization and Penning ionization.
- IW contacts surface 5 ns earlier than for TiO_2 substrate.
- Plasma dims at 230 ns, after contacting the surface.

WATER: ELECTRON DENSITY

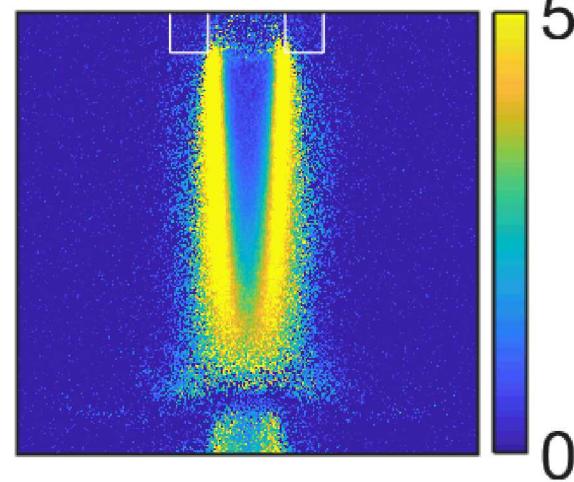
$t = 215 \text{ ns}$



$t = 235 \text{ ns}$



$t = 280 \text{ ns}$



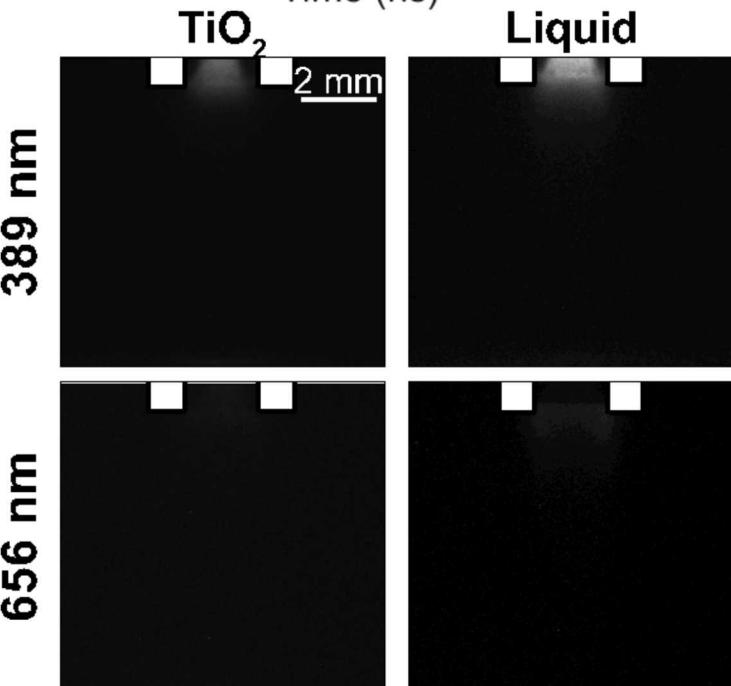
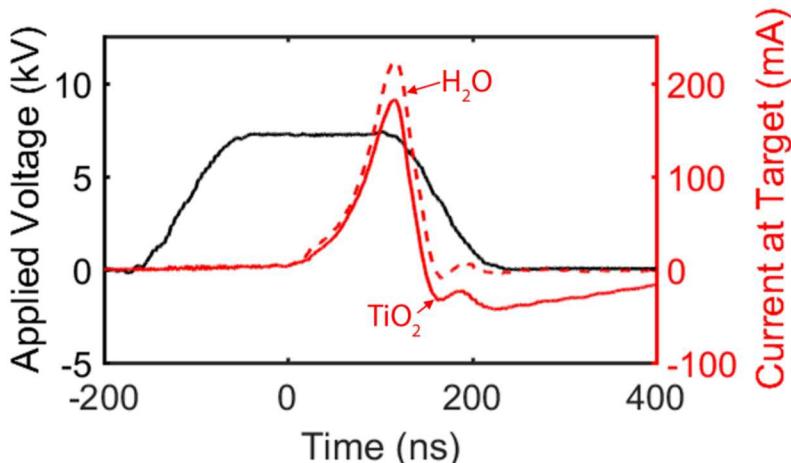
- Water substrate.
- Electron density (n_e) is again annular due to photoionization and Penning ionization of surrounding water vapor.
- SIW is not visible by LCIF because high H_2O density near the surface depletes $\text{He}(2^3\text{S})$.



- Electron density is slightly higher than for TiO_2 because of higher water vapor density.

- Ratio = 1 $\rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$

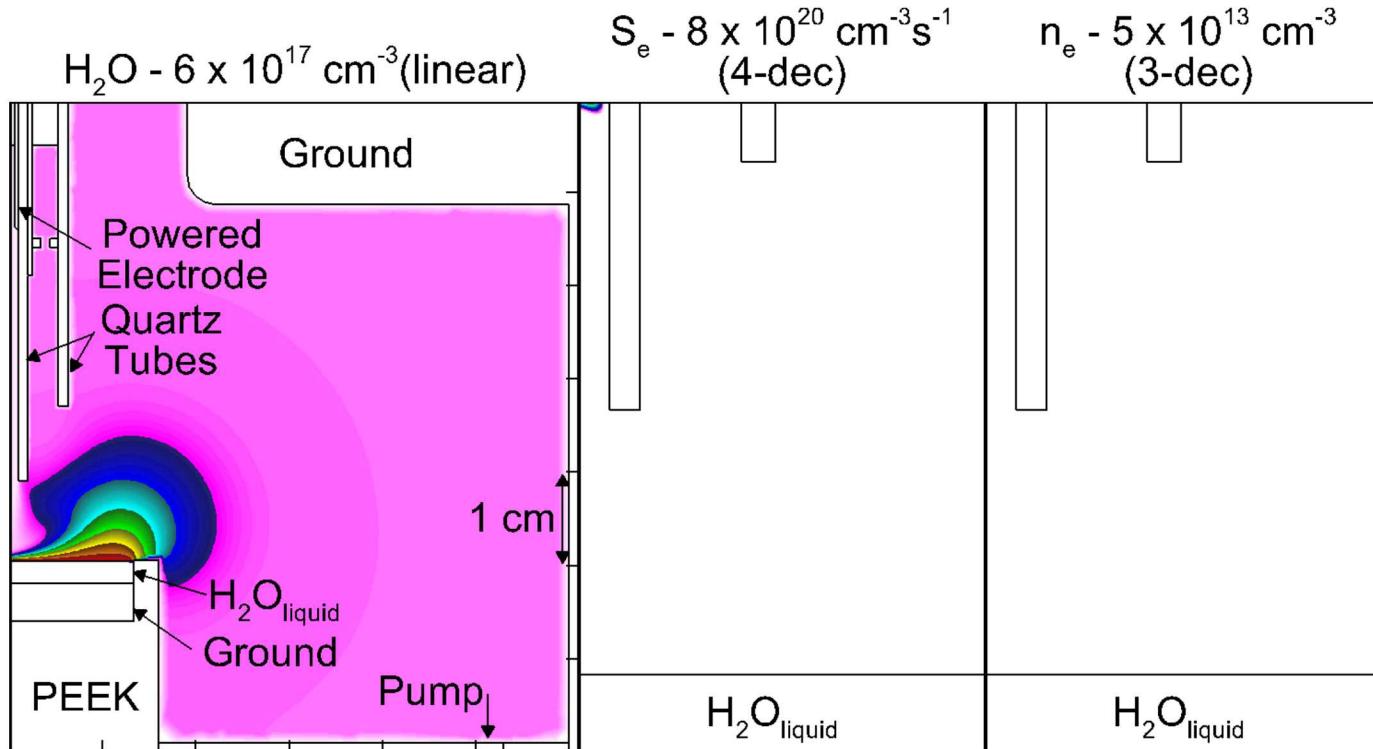
WATER vs. TiO_2



- With water base, magnitude of forward current is higher, reverse current is lower.
- Forward current enhanced by humidity.
- Electrons/ions solvate into the water while being physisorbed on TiO_2 .
- Water less likely to support electron emission to support reverse of current.
- Evaporation produces higher water concentrations near surface resulting in lower mobility, thinner SIW.
- Plasma dims after IW contacts liquid surface.

Animation Slide

MODELING LIQUID INTERACTIONS



- Cylindrically symmetric.
- Modeling work ongoing.
- Same parameters as experiment.

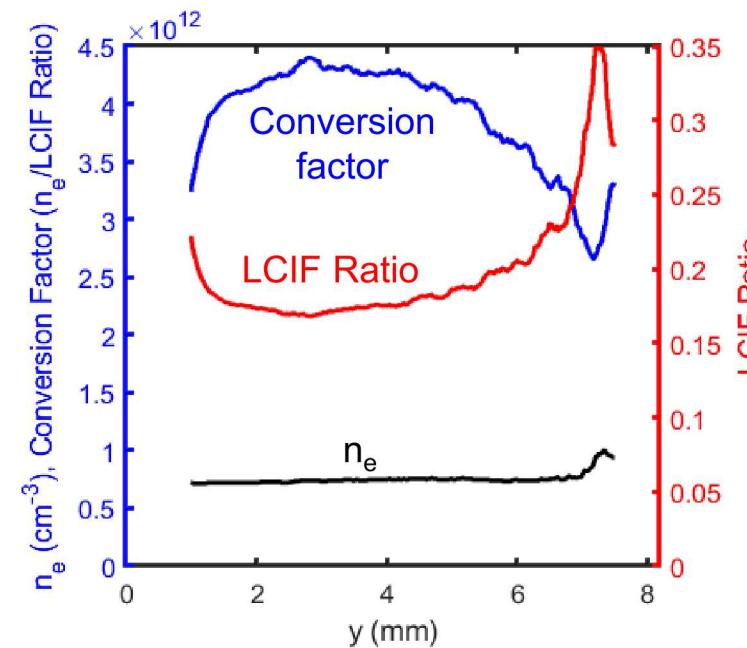
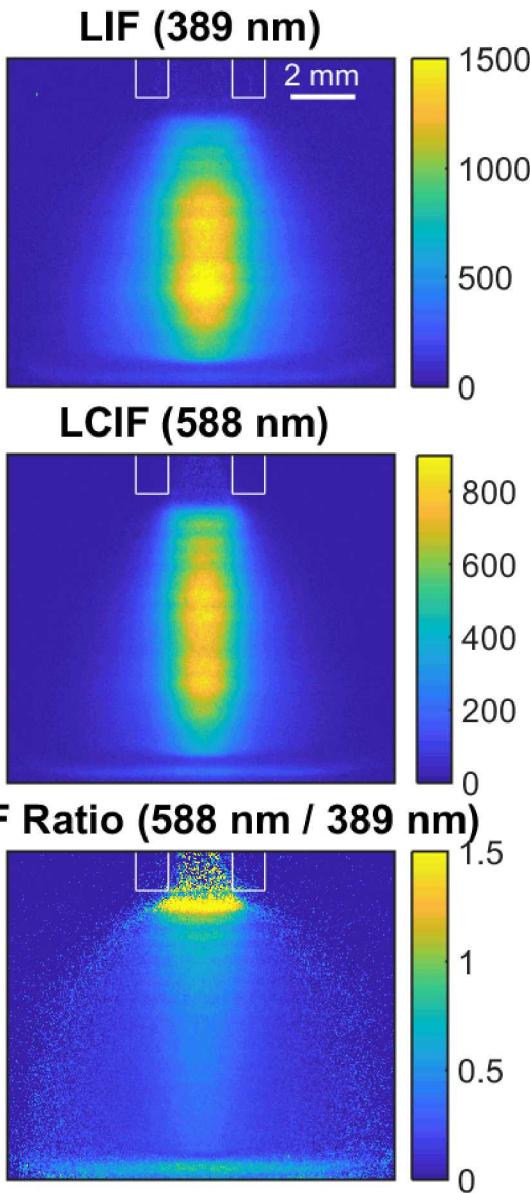
- *nonPDPSIM* – 2-dimensional plasma hydrodynamics model
- Photoionization and photoemission from surfaces are critical in positive IW propagation.
- Surface ionization wave develops along the liquid surface.
- Electron density profile is annular.

CONCLUDING REMARKS

- An ionization wave exiting a He plasma jet into a humid environment is annular due to Penning ionization and photoionization at the boundary
- Plasma in contact with water vs TiO_2 (same ϵ/ϵ_0) are distinguishable from electrical I-V traces.
 - Mild humidity above the water produces larger forward current.
 - Solvation of electrons/ions in water and larger humidity above water result in lower reverse current.
- SIWs propagating over a liquid surface are thinner due to the lower mobility of electrons in a saturated water vapor environment.

Appendix

CONVERT LCIF TO n_e

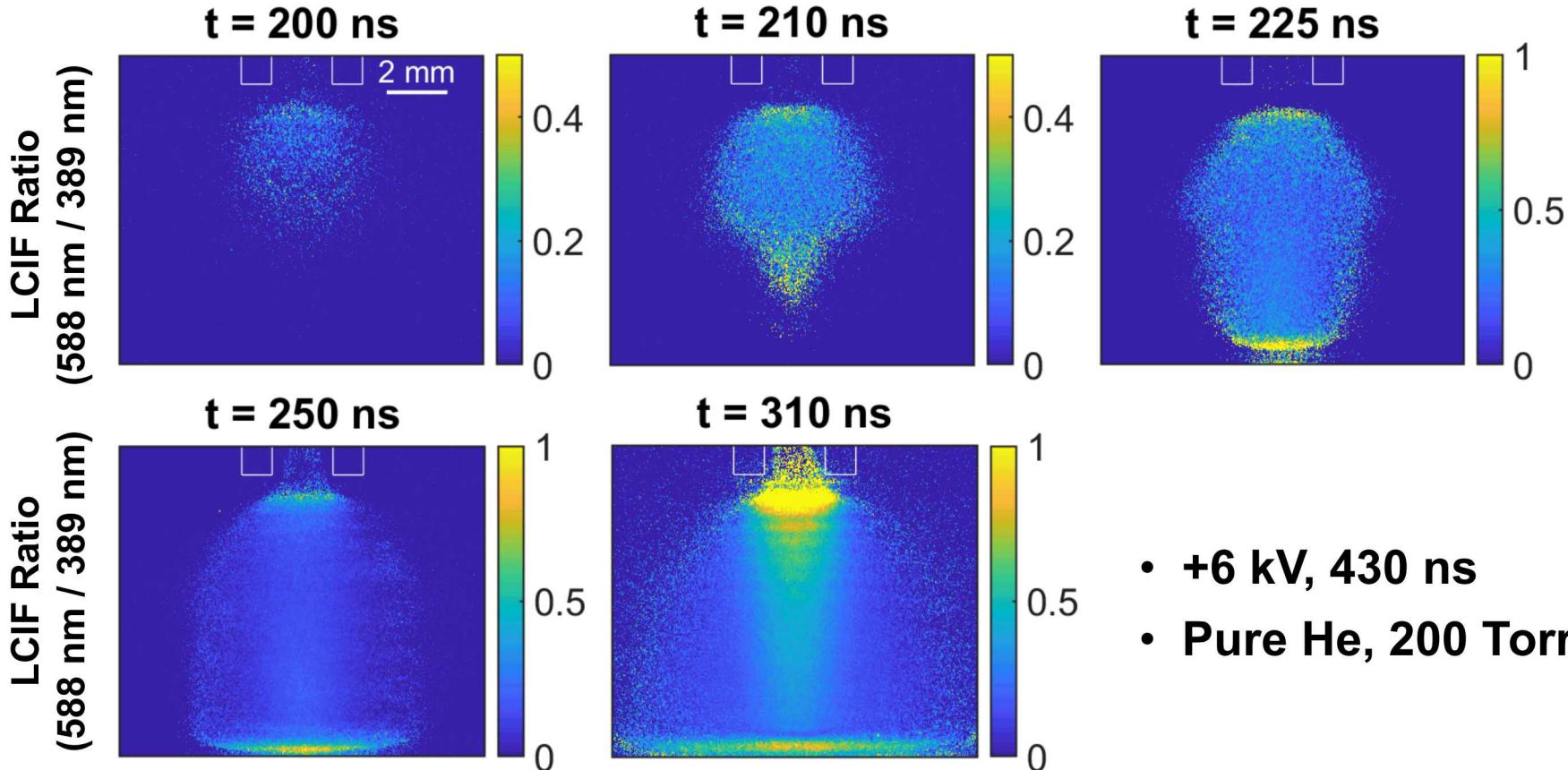


- Calculate n_e using Ohm's law and E/N.

$$I = \frac{e^2}{m_e} \frac{N}{v_m} \frac{E}{N} n_e A$$

- An LCIF ratio of 1, is approximately 4×10^{12} cm $^{-3}$ electrons.
- Previously, conversion factor estimated at 1.5×10^{13} cm $^{-3}$ at 600 Torr.

BASE CASE LCIF

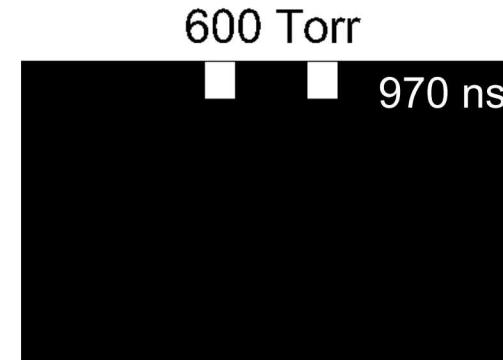
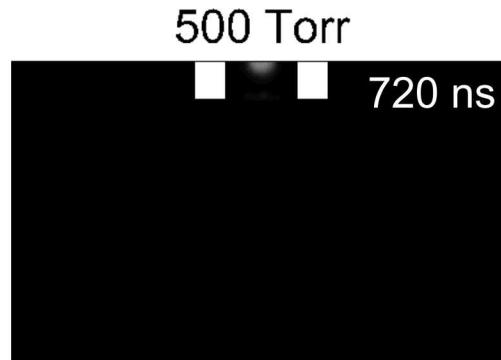
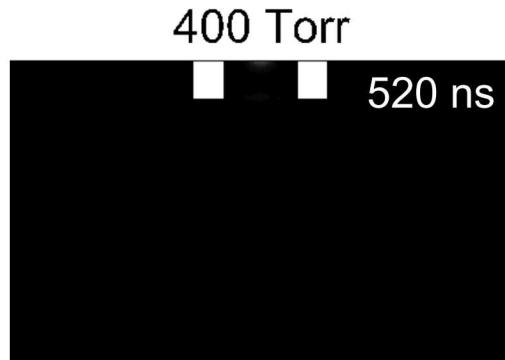
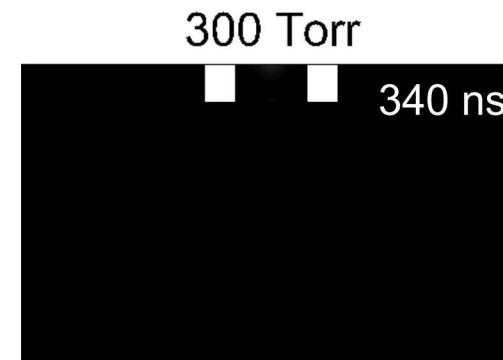
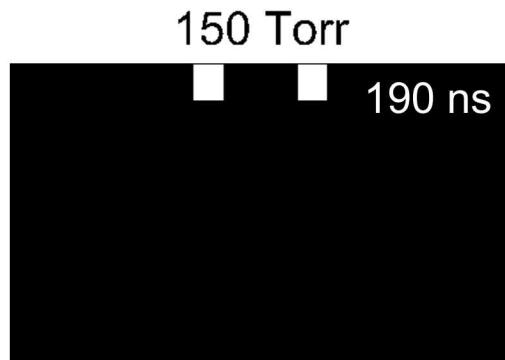


- +6 kV, 430 ns
- Pure He, 200 Torr

- Before IW reaches the surface, He(2^3S) densities are low.
- n_e in the SIW is nearly double that of the bulk.
- Elevated n_e in IW front may be due to Stark mixing.

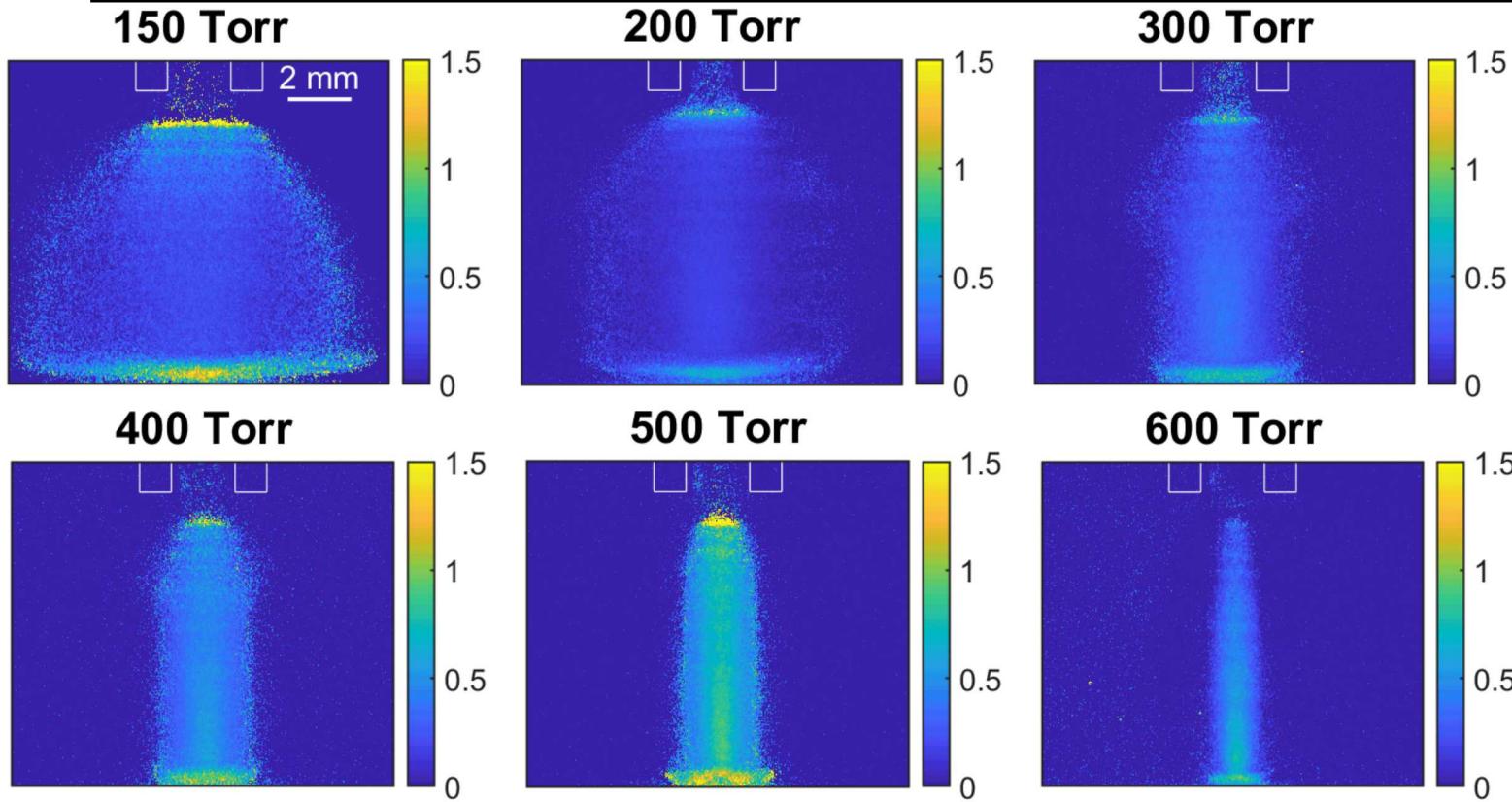
$$\text{Ratio} = 1 \rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$$

VARY PRESSURE – EMISSION



- 6 kV, 500 sccm He, 390 nm plasma emission.
- 350 ns animated, time of first frame indicated.
- Varied pulse duration so voltage is on for 80 ns after contact.
- IW propagates slower and SIW becomes thinner for higher pressures. (lower electron mobility)

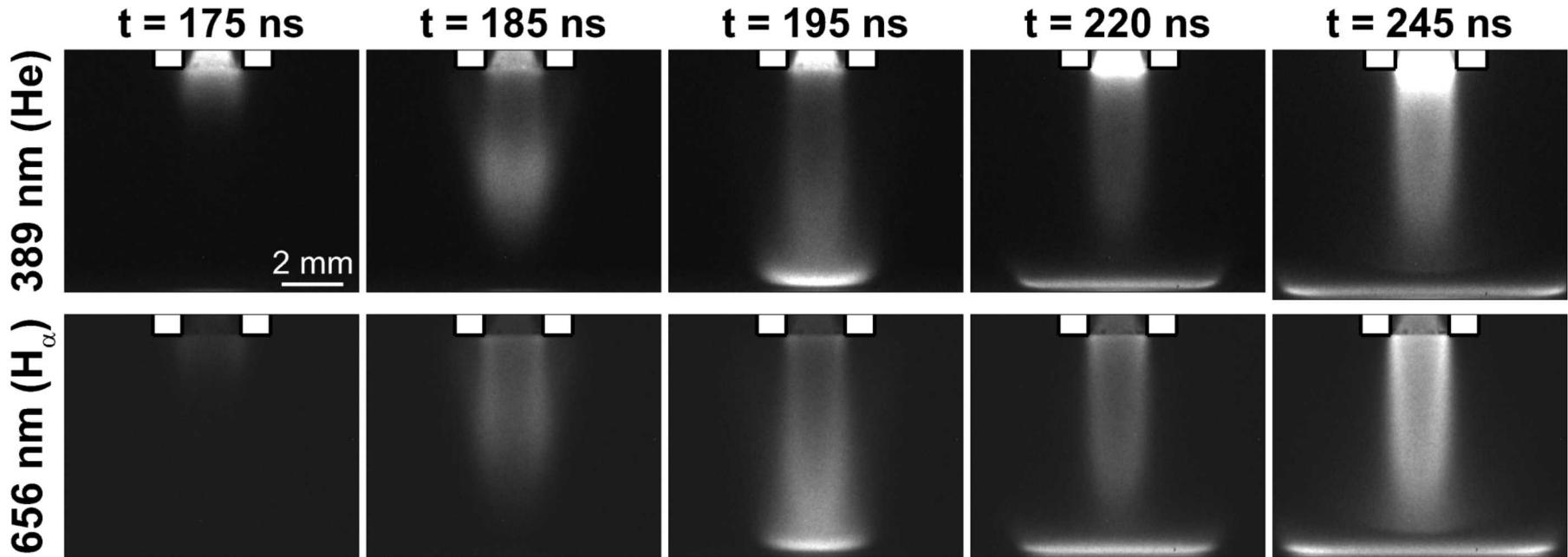
VARY PRESSURE – n_e



- 6 kV, 500 sccm He, 30 ns after IW contacts surface.
- Plasma is more confined at higher pressure, n_e increases.
- Current and energy deposition decrease with increasing pressure.
- Above 500 Torr, n_e is collisional enough that ionization rate drops.

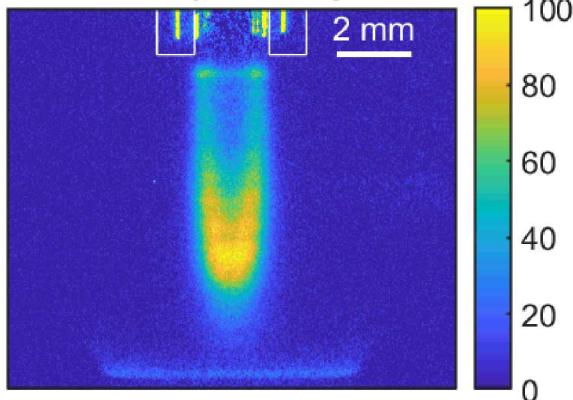
$$\text{Ratio} = 1 \rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$$

HUMID He SHROUD EMISSION

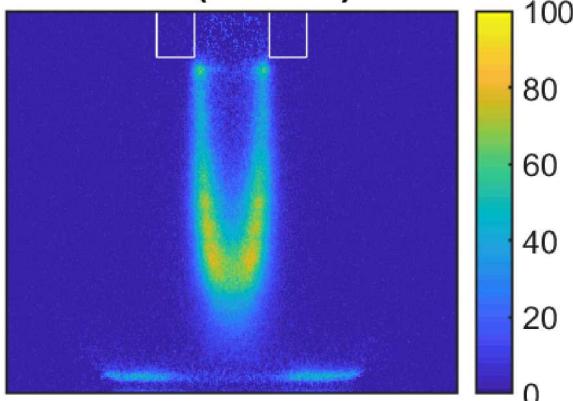


- 2.3% H_2O in shroud.
- Images have not been Abel inverted.
- IW reaches outlet of the tube earlier than in base case – photoionization from He_2^* causes non-local seed ionization.
- Photoionization and Penning ionization promote IW speed.
- H_α emission appears more annular – dominates at the interface of the center and shroud flow.

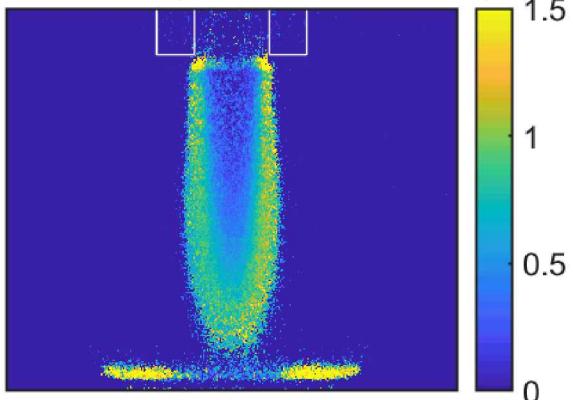
LIF (389 nm)



LCIF (588 nm)



LCIF Ratio (588 nm / 389 nm)

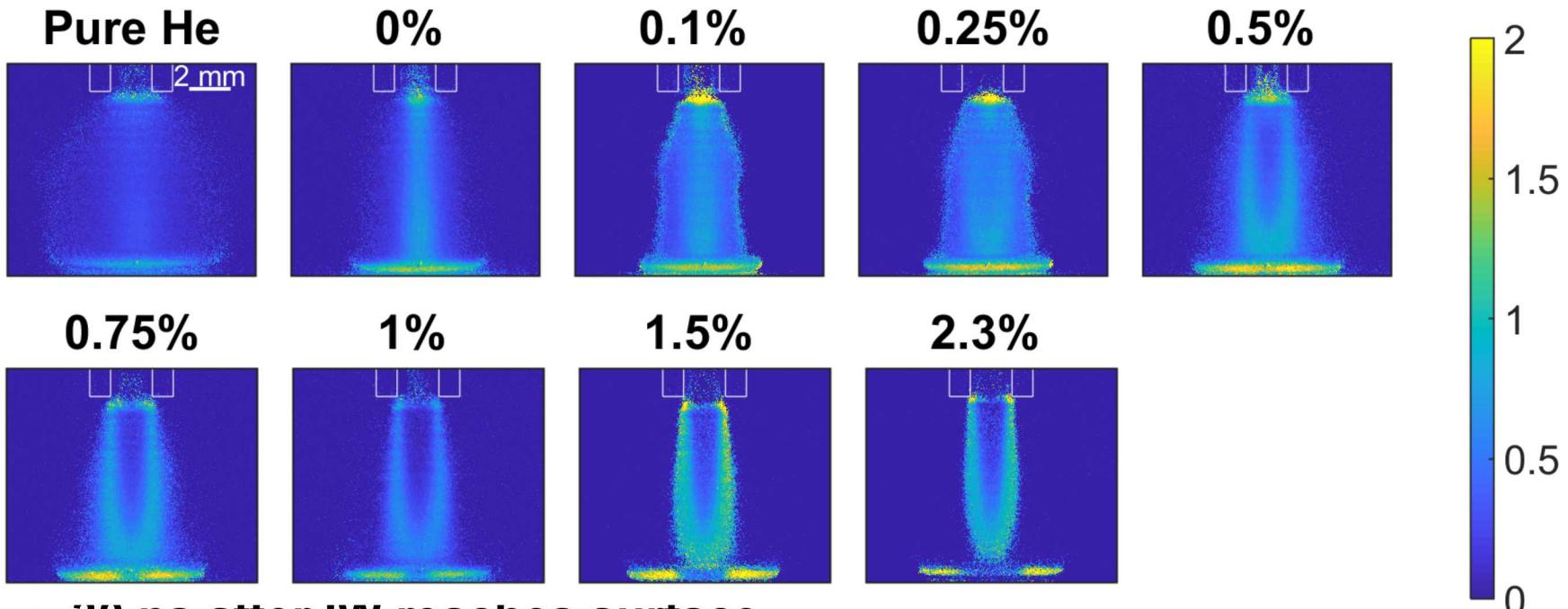


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- $t = 230$ ns, 30 ns after IW contacts surface
- Moving away from He core, there are fewer He(2³S), LIF signal decreases.
- $$\text{He}(2^3\text{S}) + \text{H}_2\text{O} \rightarrow \text{He} + \text{H}_2\text{O}^+ + \text{e}^-$$
- In regions of high H₂O concentration, there may be significant n_e which is not detectable due to low He(2³S).

$$\text{Ratio} = 1 \rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$$

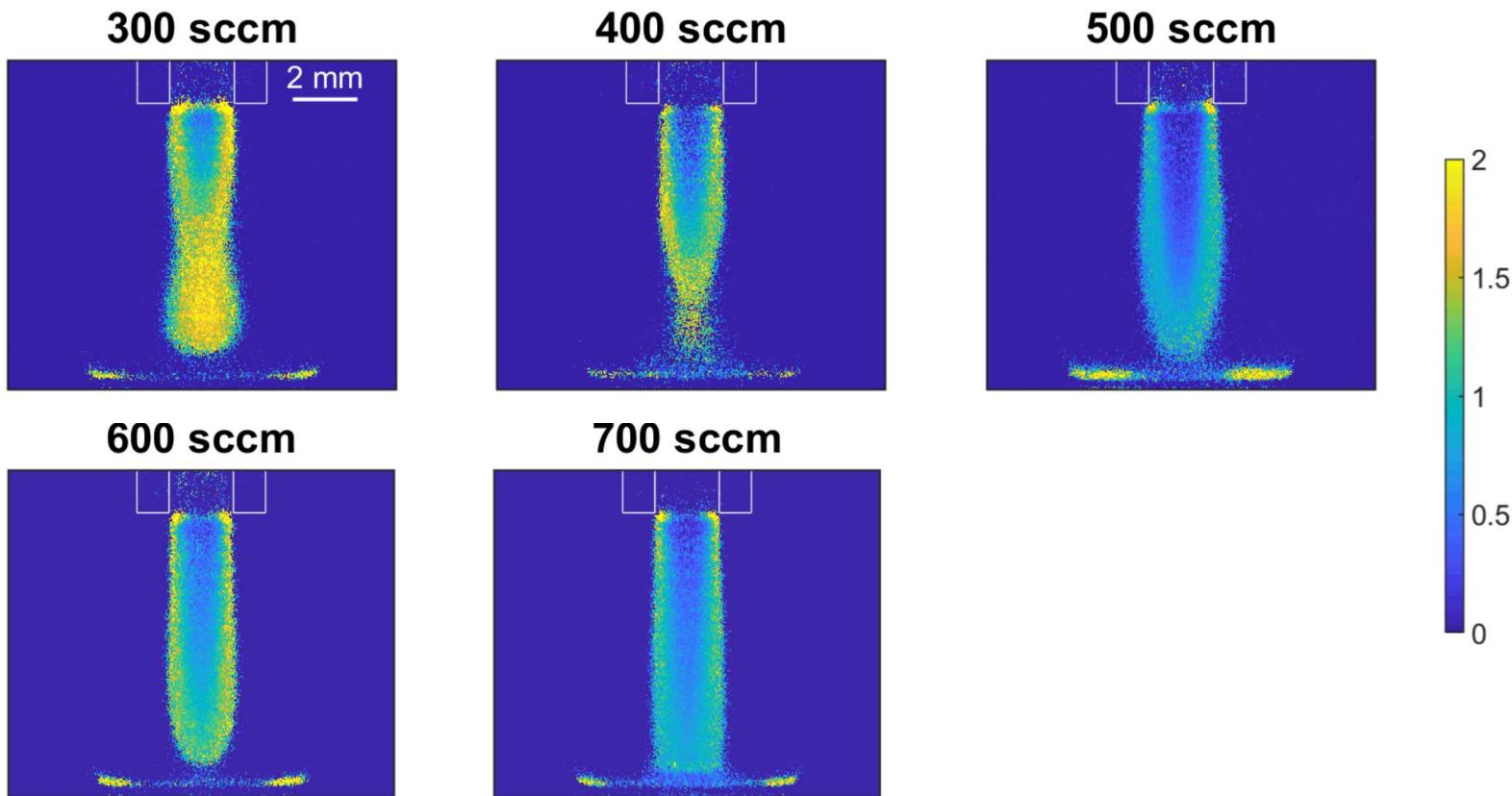
SHROUD HUMIDITY



- 30 ns after IW reaches surface.
- Transition from diffuse in pure helium case to confined by humid shroud.
- Higher electron energy loss rates with H_2O because of vibrational and rotational excitation.
- n_e increases with humidity due to Penning ionization.

$$\text{Ratio} = 1 \rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$$

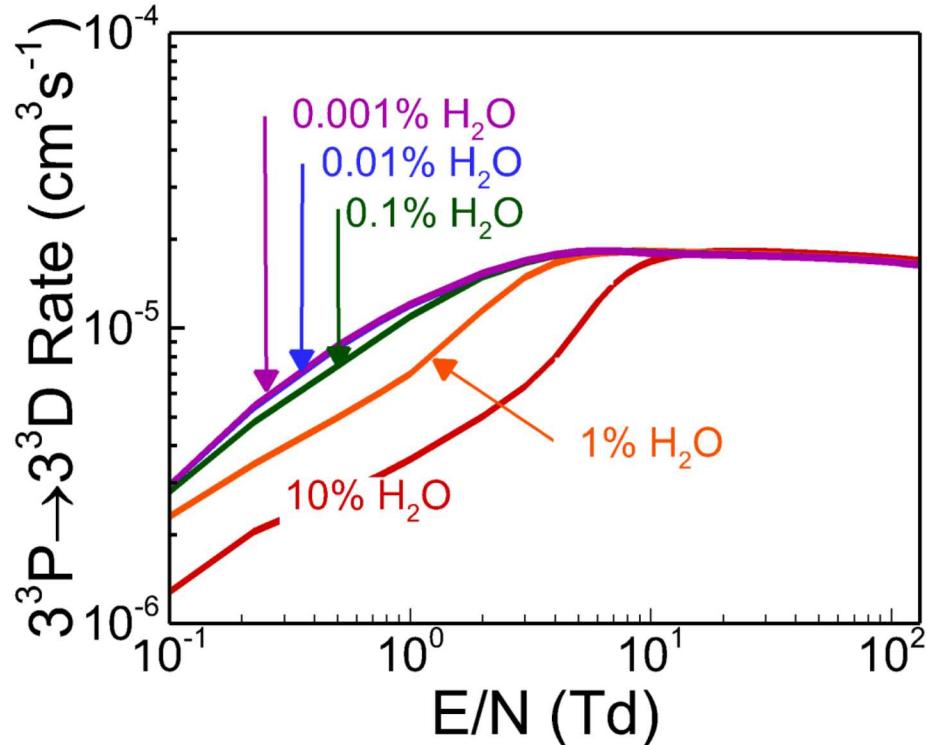
VARY He FLOW RATE - n_e



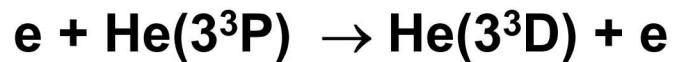
- Vary center flow rate.
- Shroud: 500 sccm, $\text{He}/\text{H}_2\text{O} = 98.7/2.3$
- Higher He flow rates more rapidly convect in-diffusing H_2O .

$$\text{Ratio} = 1 \rightarrow n_e \approx 4 \times 10^{12} \text{ cm}^{-3}$$

BOLTZMANN CALCULATIONS



- For LCIF measurement to be a linear representation of n_e , this rate must be independent of T_e :



- Threshold = 0.06 eV
- In pure He, this occurs when $E/N > \sim 0.8$ Td
- For $H_2O < 1\%$, LCIF is valid ~ 1 Td