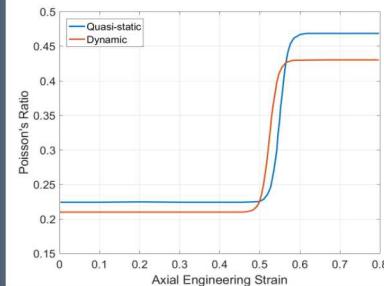
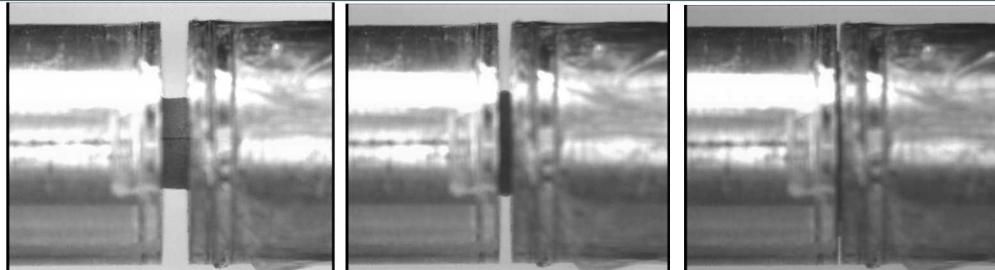


Poisson's Ratio Induced Radial Inertia During Dynamic Compression of Hyperelastic Foams



Brett Sanborn, Bo Song, Wei-Yang Lu

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Unclassified
Unlimited
Release



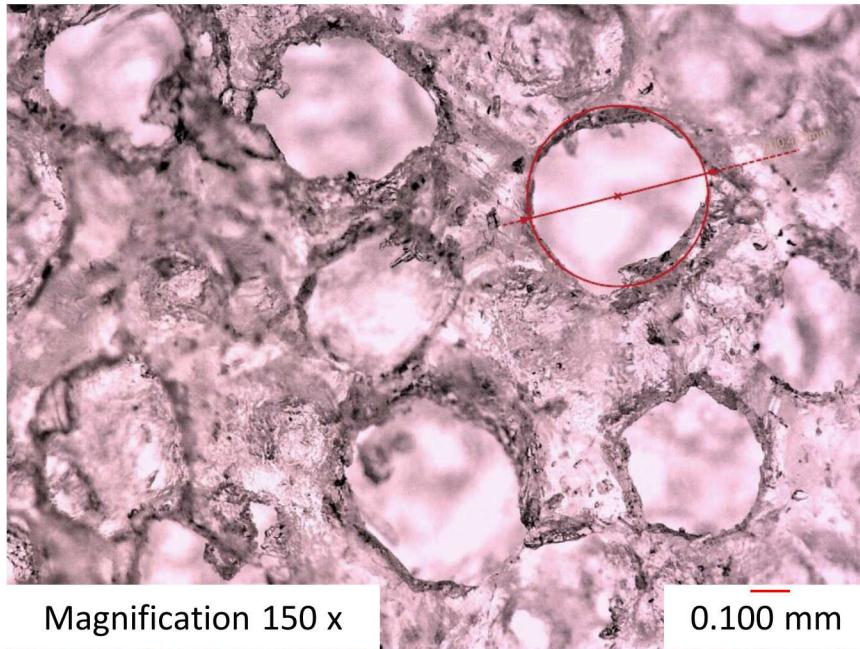
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Silicone Foam Background



- Shock Vibration Isolation and Reduction

- Light weight
- Excellent energy absorption
- Full recovery after impact loading



Open-Cell Silicone foam

Density: $608 \pm 21.85 \text{ kg/m}^3$

Average cell size: $\sim 0.5 \text{ mm}$

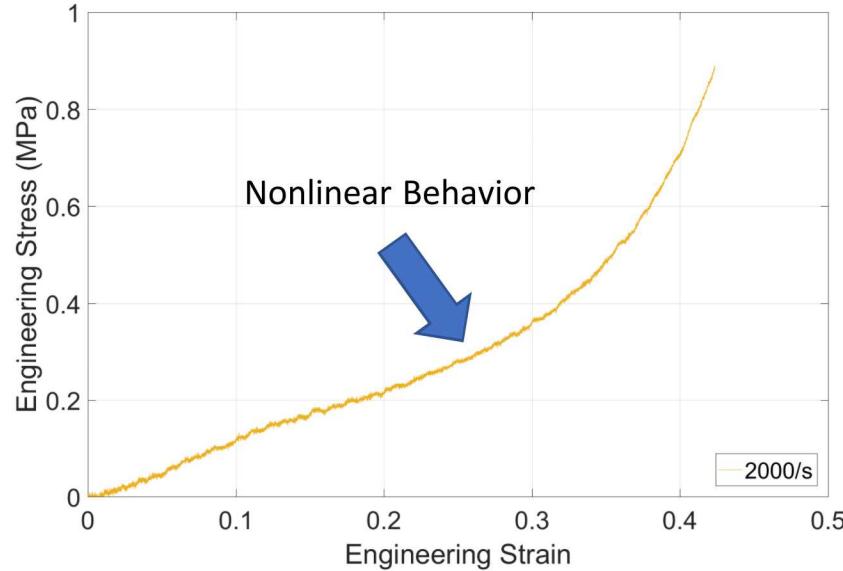
Porosity: $\sim 50\%$

At first glance, it may seem OK to neglect Poisson's Ratio or assume a small number near zero due to apparent compressibility

Silicone Foam



- Mechanical response of silicone foam may be altered by...
 - Strain rate
 - Stress-state
 - Temperature



Shock or vibration isolation pads can undergo

- High strain rate
- High densification state

Past densification, material may lose compressibility

How does Poisson's Ratio change with strain rate and densification?

Poisson's Ratio

- Poisson's Ratio is typically regarded as a material *constant*

$$\nu = -\frac{e_r}{e_x}$$

Good for metals in Elasticity

Small strains

Materials such as foams are undergo large deformation

$$\nu = -\frac{\ln \lambda_r}{\ln \lambda_x} = -\frac{\varepsilon_r}{\varepsilon_x}$$

True strains along axial and radial directions

Secant Poisson's Ratio

Applicable to Linear Response Only



In the case of nonlinear large deformation, ***tangent Poisson's ratio*** with true strains represents actual Poisson's ratio

$$\nu = -\frac{d\varepsilon_r}{d\varepsilon_x}$$

Literature

- Most available studies on Poisson's ratio...

Conducted at quasi-static strain rates



Response might not apply to high strain rates

Limited to small deformations



Material can undergo large deformation during use

Some measure large deformation, use secant ratio to calculate

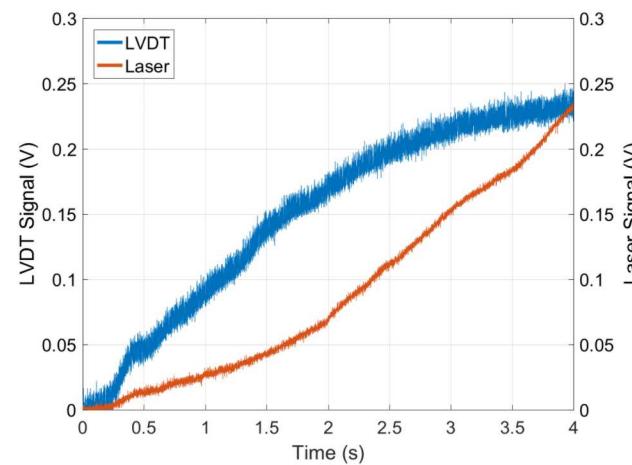
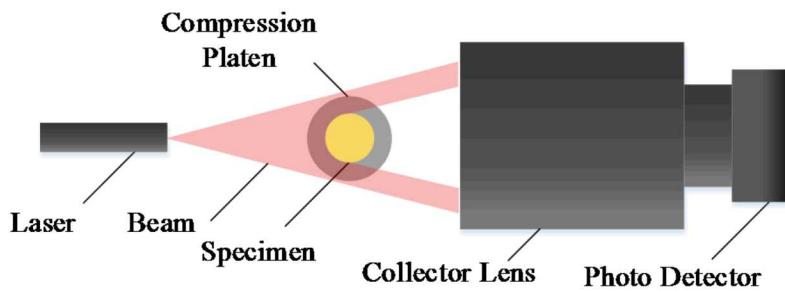
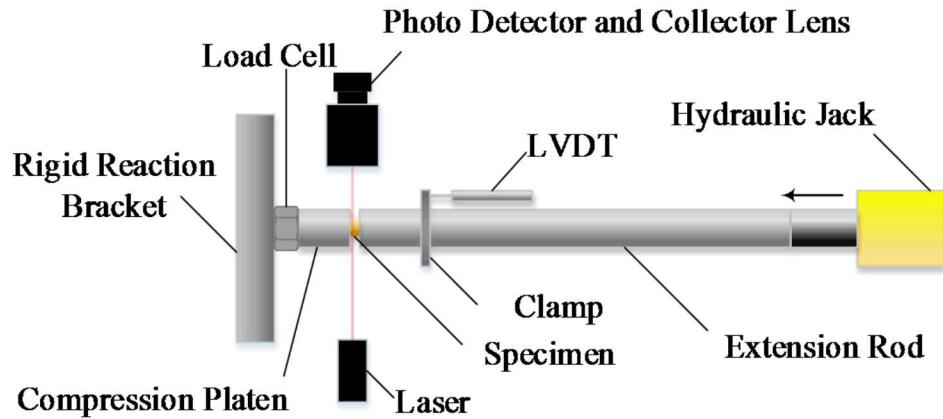


Tangent method must be used for accurate results

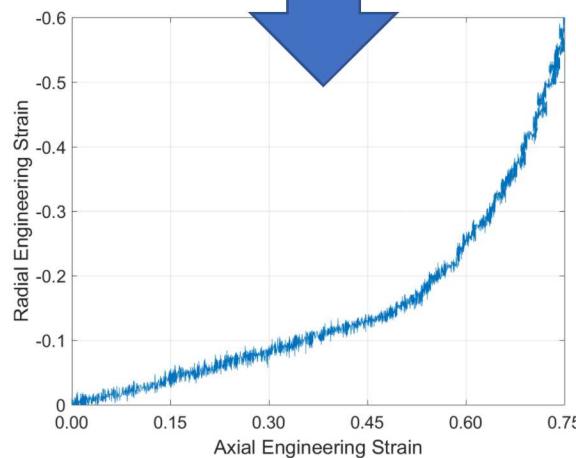
Our Approach: Measure Poisson's ratio at high rate*, large deformation, and calculate using tangent ratio for best results

*quasi-static also included as a check

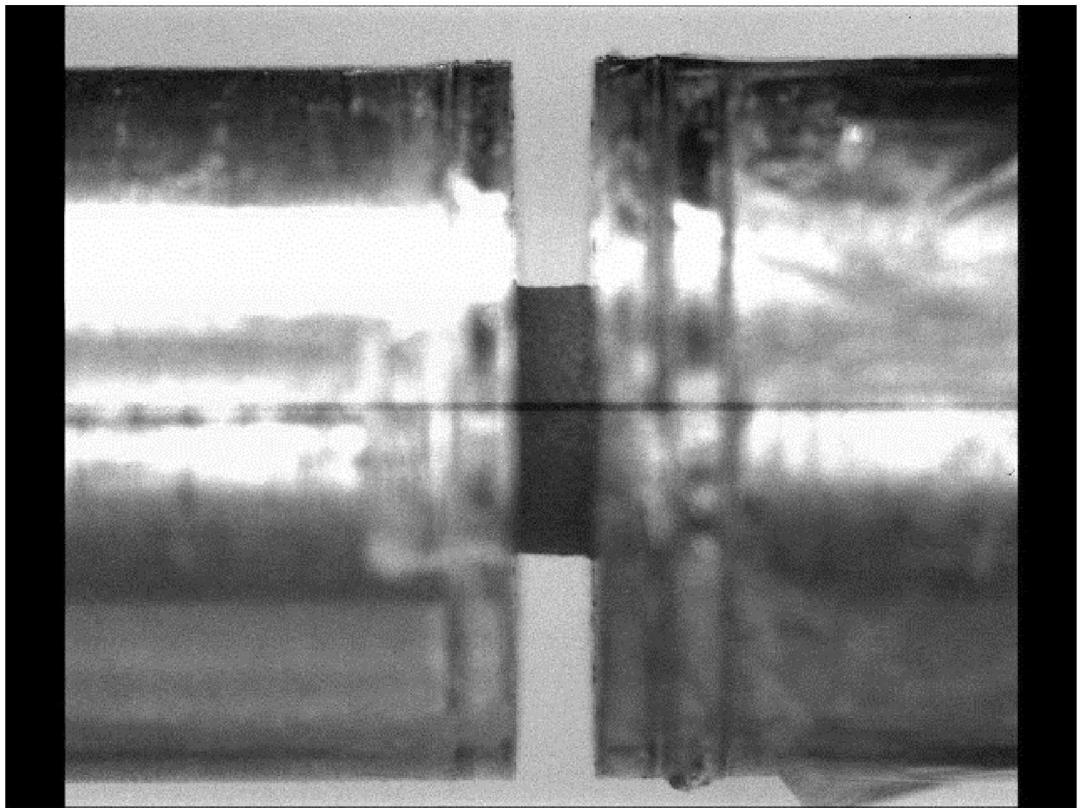
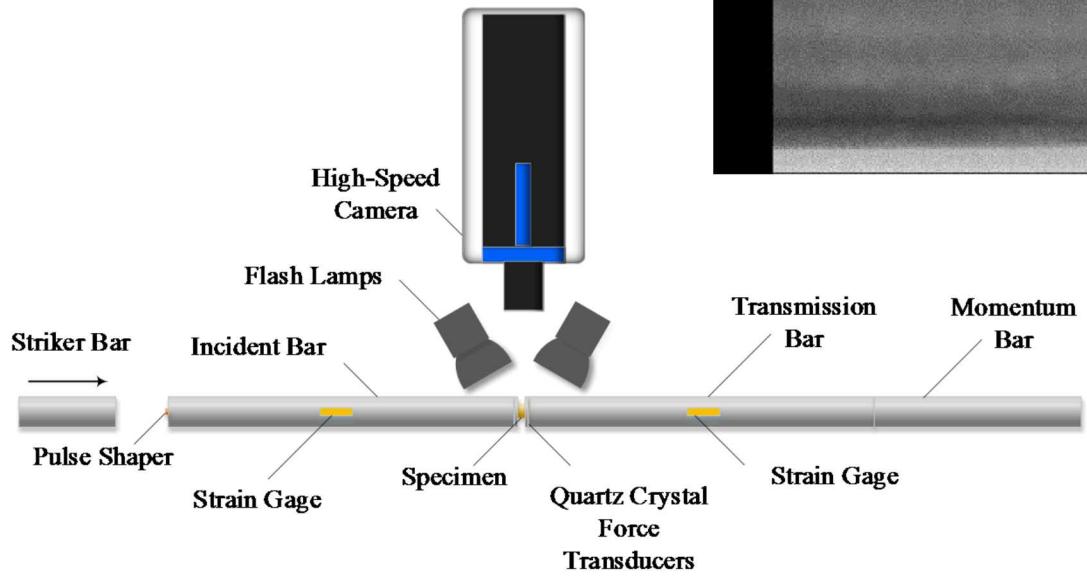
Quasi-static experimental setup and calibration



$$e_x = 1 - \frac{l_1}{l_0} \quad e_r = 1 - \frac{d_1}{d_0}$$

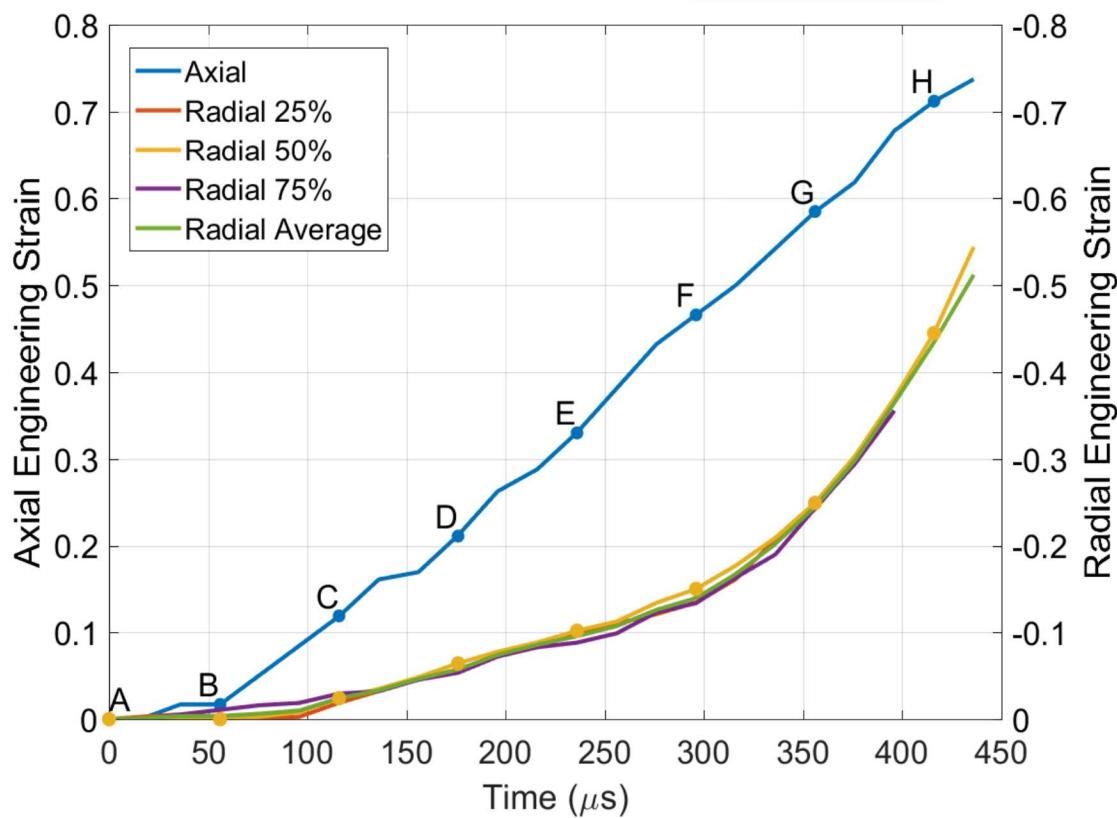
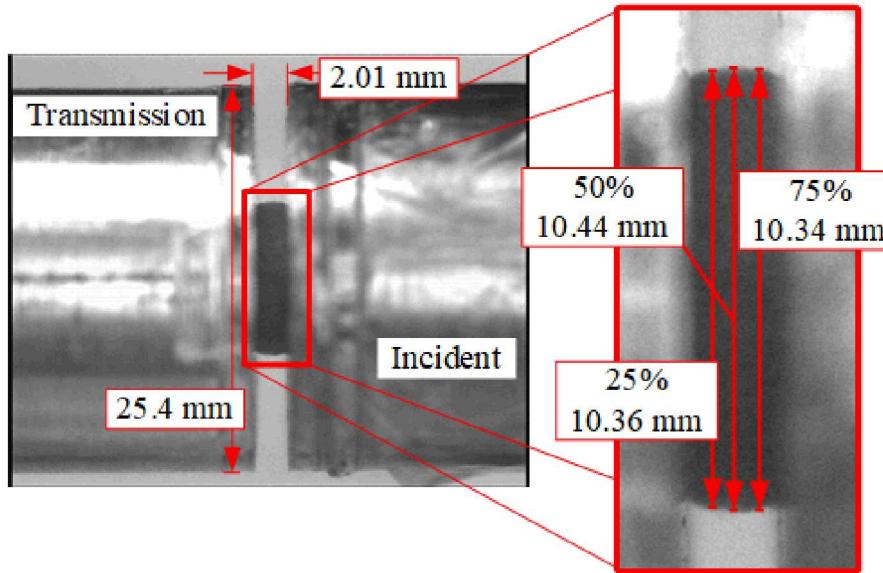
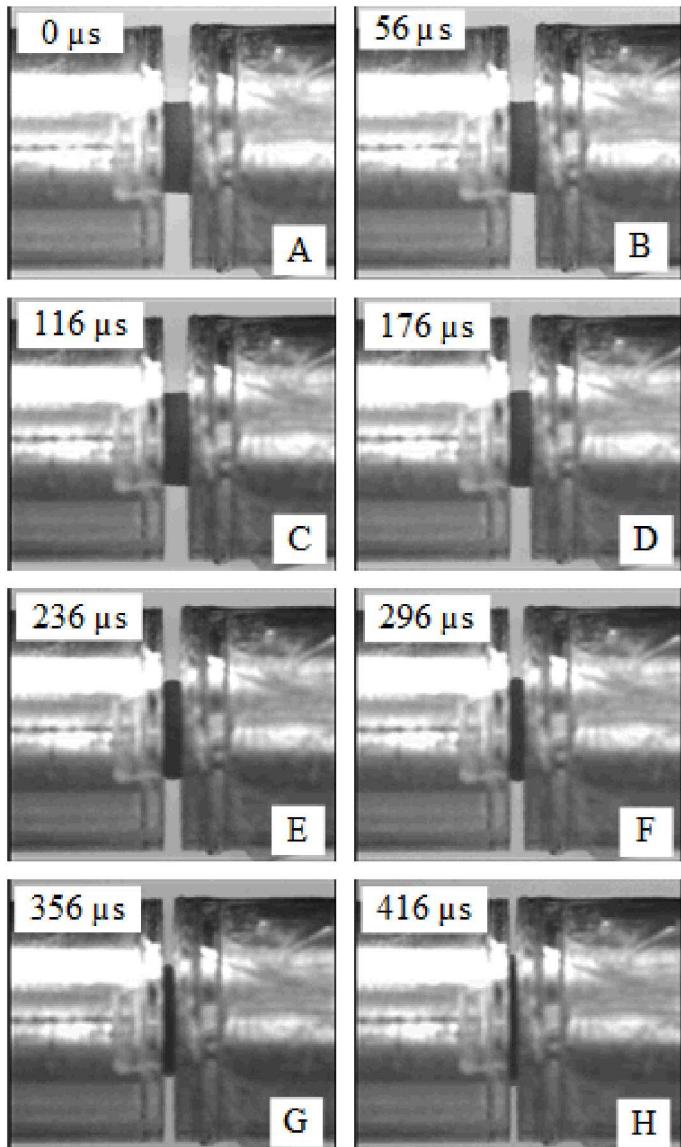


High Rate Experiments



250 kFPS
2000 s^{-1}

High Rate

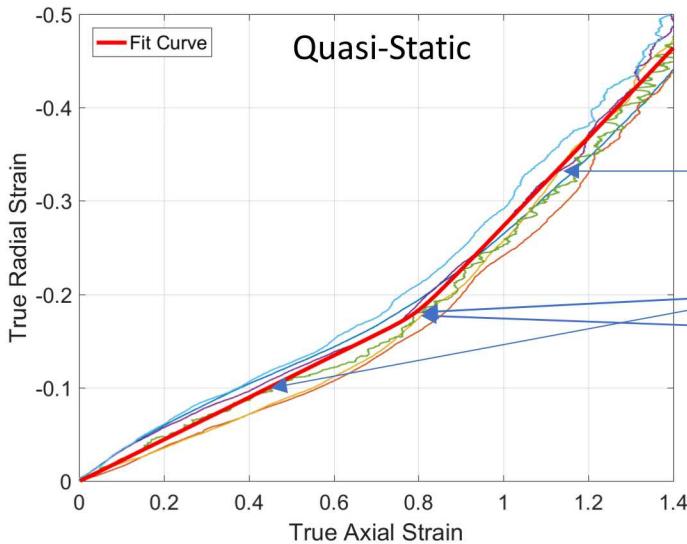


Poisson's Ratio Calculation

$$\nu = -\frac{d\varepsilon_r}{d\varepsilon_x}$$

- Tangent Poisson's Ratio for large, nonlinear deformation

$\varepsilon = -\ln(1-e)$ True strain: Axial compression (+) and Radial Tension (-)

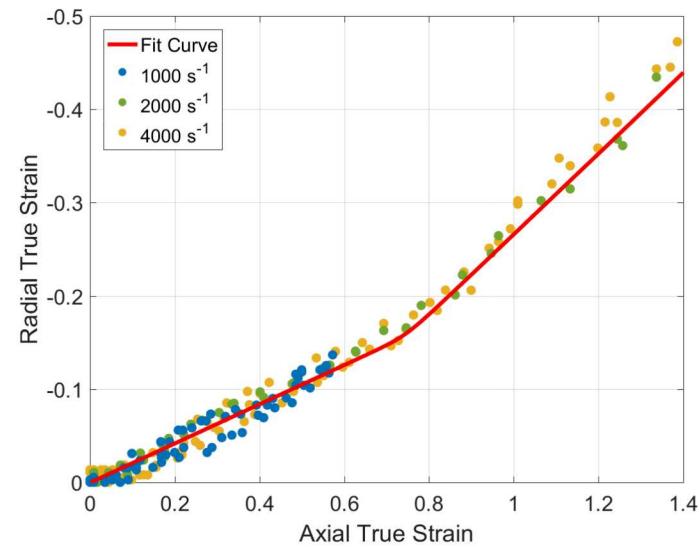


Bilinear

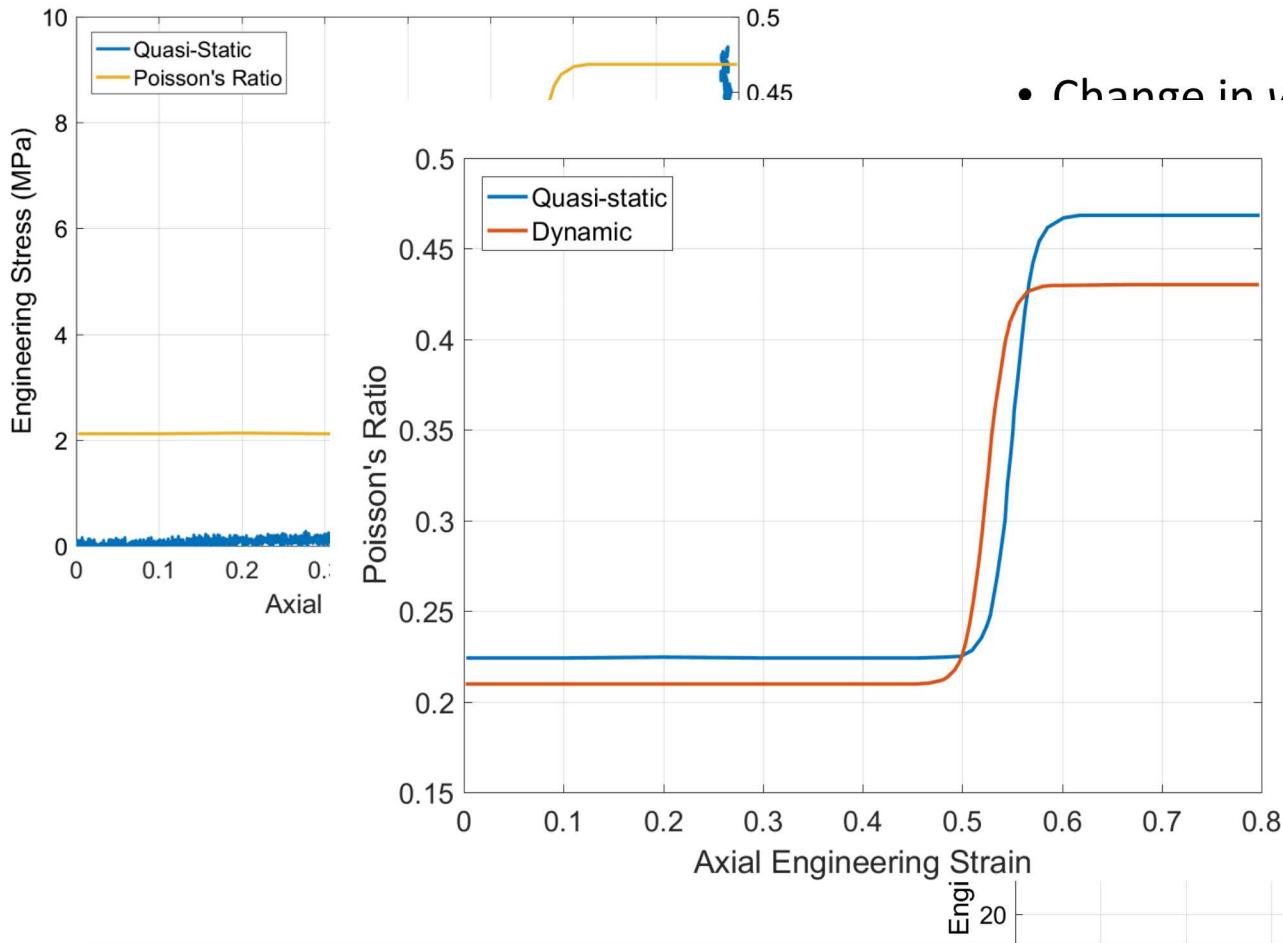
Boltzmann Sigmoidal Function

$$\nu(e_x) = \frac{\nu_1 - \nu_2}{1 + \exp\left(\frac{e_x - e_{x0}}{\delta}\right)} + \nu_2$$

Parameter	Quasi-static	Dynamic
ν_1	0.22	0.21
ν_2	0.47	0.43
e_{x0}	0.550	0.525
δ	0.01	0.01



QS and High Rate Poisson's Ratio



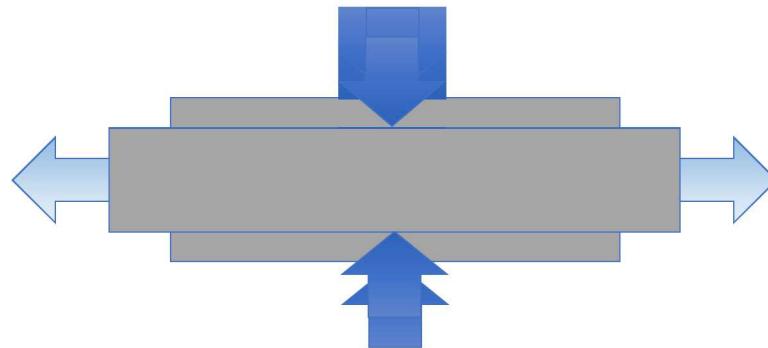
- Change in ν coincided with
or both rate

Change in Poisson's ratio from compressible to nearly incompressible needs to be included in models for most accurate results

Radial Inertia during Dynamic Compressive Loading

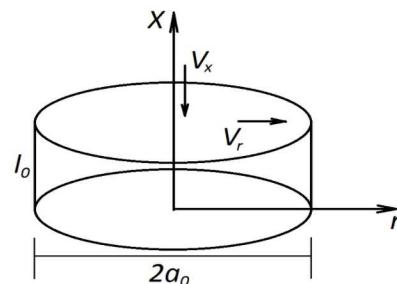


- Change in Poisson's Ratio affects the measured response of the material due to **Radial Inertia**



Dharan and Hauser (1970): In a compression test with constant axial velocity:

- At a specific location, particle velocity along lateral (radial) direction increases with time (or increasing strain/decreasing specimen thickness)
- At a specific time, particle at outer diameter moves faster along lateral (radial) direction than the particle at inner diameter



$$V_r(t) = \frac{r}{2l} V_x(t)$$

Result: Radial Confinement

Consequence of Radial Inertia

- Abnormal axial stress history  May overshadow material response (soft materials)

Abnormal spike in stress history due to radial inertia

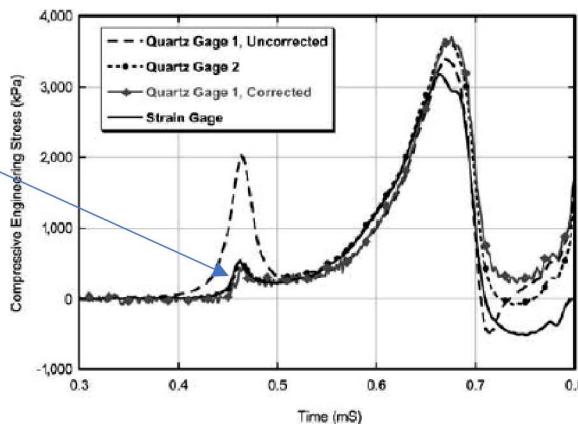


Fig. 1 The abnormal initial spike-like features in the stress histories of a gelatin specimen

The bump is stress measured in axial direction due to radial-inertia-induced confinement in the sample

Could be incorrectly attributed to material response

Currently Existing Analysis

- Compressible Solid
- Small deformation

- Incompressible Solid
- Large deformation

➤ Kolsky (1949)

$$\sigma_z = \frac{\nu^2 a_0^2 \rho_0}{2} \cdot \ddot{e}$$



$$\bar{\sigma}_z = \frac{1}{\pi a^2} \int_0^a \int_0^{2\pi} \sigma_z(r) r dr d\theta$$

❖ Dharan and Hauser (1970)

$$\bar{\sigma}_z = \frac{\rho_0 a_0^2}{4l_0 (1-e_x)^2} \left[\frac{3V_x^2}{2l_0 (1-e_x)} + \frac{dV_x}{dt} \right]$$

➤ Forrestal et al. (2007)

$$\sigma_z = \frac{\nu^2 (3-2\nu)}{4(1-\nu)} \left[a_0^2 - \frac{2r^2}{(3-2\nu)} \right] \rho_0 \cdot \ddot{e}$$



➤ Warren and Forrestal (2010)

$$\sigma_z = \frac{\rho}{4(1-e_x)^2} \left[\frac{3\dot{e}^2}{2(1-e_x)} + \ddot{e}_x \right] (a_0^2 - r^2)$$

Factors that affect radial inertia stress

Specimen Geometry:

Specimen Diameter
Specimen Thickness

Loading Conditions:

Impact Velocity
Strain Rate
Strain

Specimen Properties:

Density
Poisson's Ratio

Comprehensive Radial Inertia Analysis

- Re-deriving the previous analysis using **mass** and **momentum** conservation

Mass:

$$\frac{1}{\rho(t)} \cdot \frac{d\rho(t)}{dt} = -\left(\frac{V_x(t)}{l(t)} + 2 \frac{V_r(r,t)}{r(t)} \right)$$

Momentum:

$$\frac{\partial \sigma_r(r,t)}{\partial r} = -\rho(t) \left(\frac{\partial V_r(r,t)}{\partial t} + V_r(r,t) \frac{\partial V_r(r,t)}{\partial r} \right)$$

$$\frac{\partial \sigma_r(r,t)}{\partial r} = -\rho(t) \cdot \frac{r}{1-e_x(t)} \cdot \left[v(t) \ddot{e}_x(t) + v(t) \cdot (v(t)+1) \cdot \frac{\dot{e}_x^2(t)}{1-e_x(t)} + \frac{dv(e_x)}{de_x} \dot{e}_x^2(t) \right]$$



$$\sigma_r(r,t) = \rho(t) \cdot \frac{a^2(t) - r^2(t)}{2(1-e_x(t))} \cdot \left[v(t) \ddot{e}_x(t) + v(t) \cdot (v(t)+1) \frac{\dot{e}_x^2(t)}{1-e_x(t)} + \frac{dv(e_x)}{de_x} \dot{e}_x^2(t) \right]$$



★ $\sigma_r(r_0,t) = \rho_0 \cdot \frac{a_0^2 - r_0^2}{4(1-e_x(t))^2} \cdot \left[\ddot{e}_x(t) + \frac{3}{2} \cdot \frac{\dot{e}_x^2(t)}{1-e_x(t)} \right]$ Recover equation from Warren and Forrestal (2010)

0.5

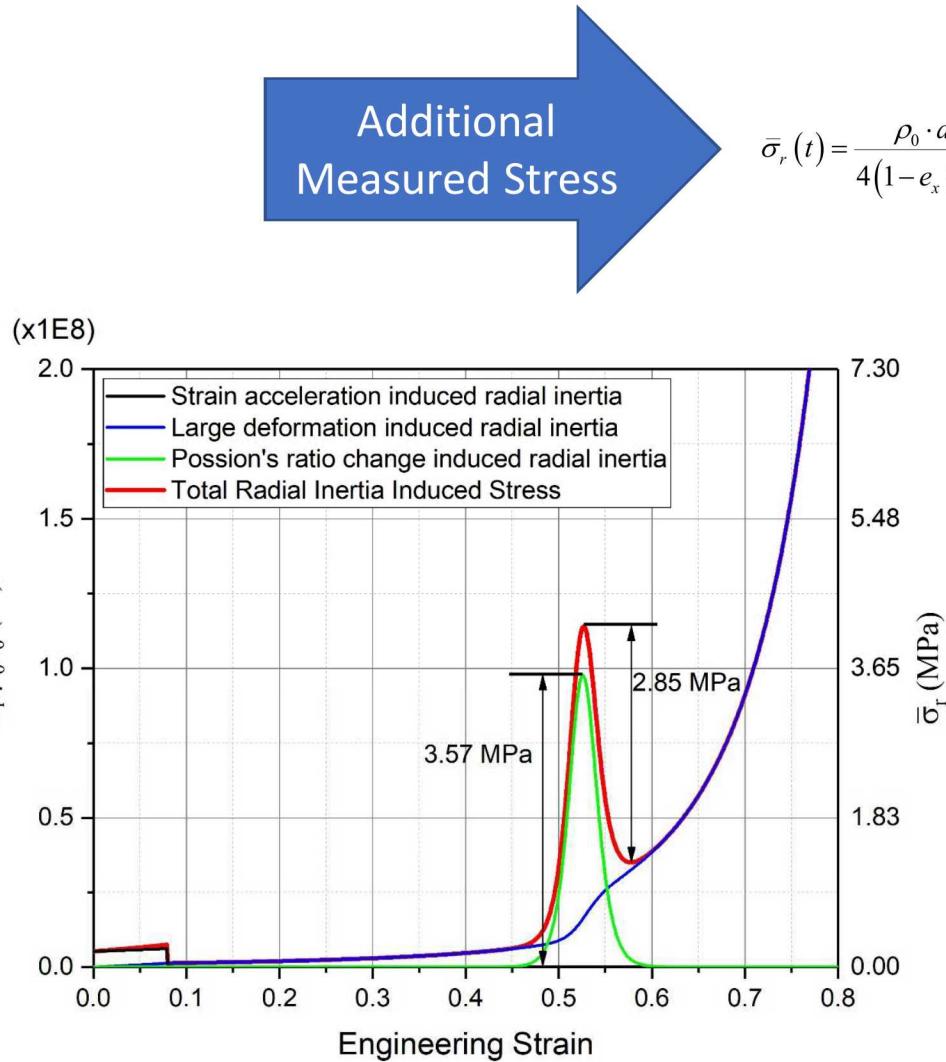
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Radial Inertia in Silicone Foam



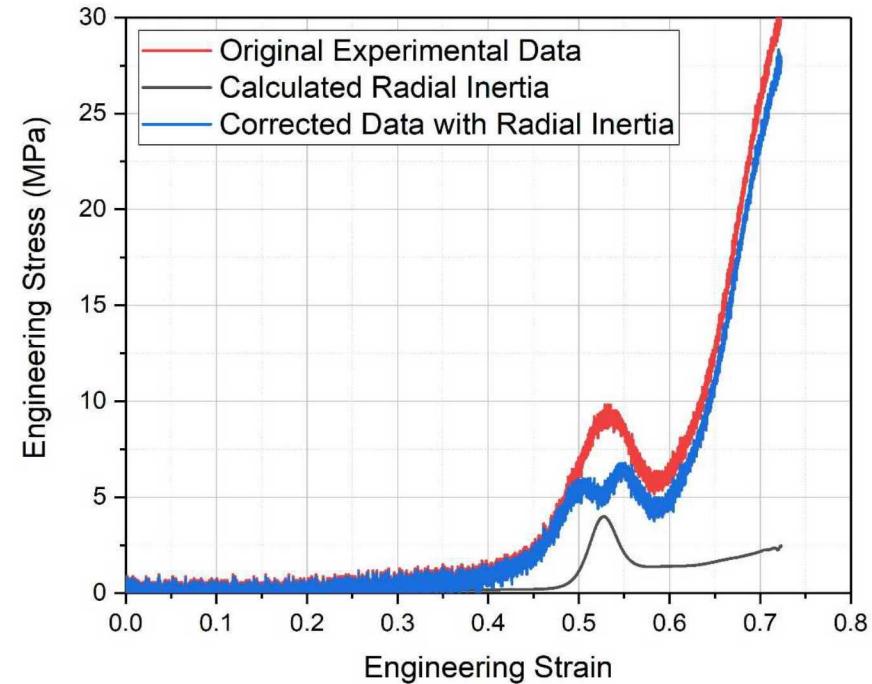
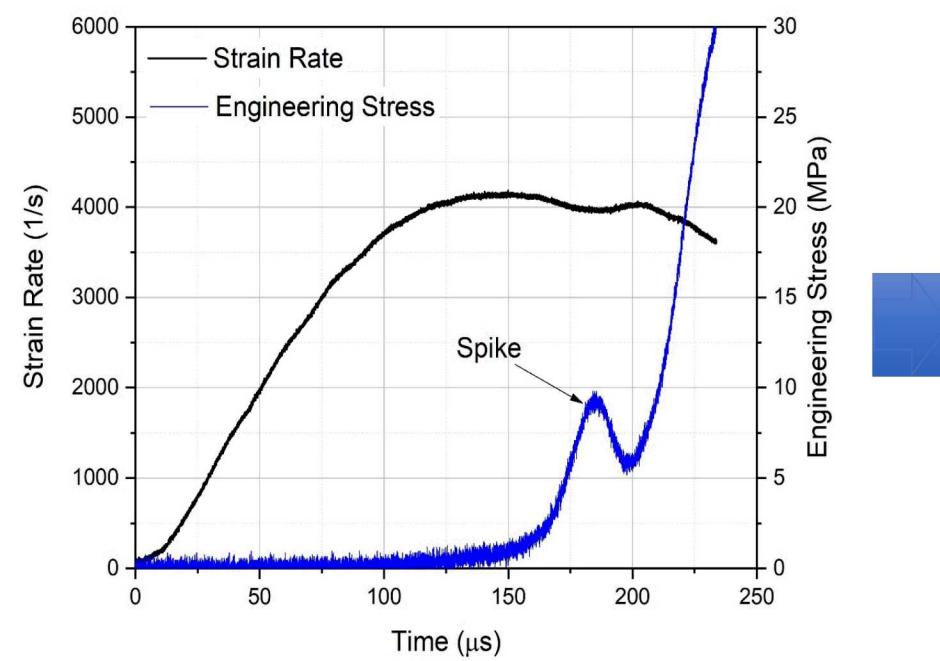
Extra Stress Including Poisson's Ratio Change Using Boltzmann Fit

$$\bar{\sigma}_r(t) = \frac{\rho_0 \cdot a_0^2}{4(1-e_x(t))^2} \cdot \left[\ddot{e}_x(t) \cdot \left(\nu_2 + \frac{\nu_1 - \nu_2}{1 + \exp\left(\frac{e_x(t) - e_{x0}}{\delta}\right)} \right) + \left(\nu_2 + \frac{\nu_1 - \nu_2}{1 + \exp\left(\frac{e_x(t) - e_{x0}}{\delta}\right)} \right) \cdot \left(1 + \nu_2 + \frac{\nu_1 - \nu_2}{1 + \exp\left(\frac{e_x(t) - e_{x0}}{\delta}\right)} \right) \cdot \frac{\dot{e}_x^2(t)}{1 - e_x(t)} + \dot{e}_x^2(t) \cdot \frac{\nu_2 - \nu_1}{\delta} \cdot \frac{\exp\left(\frac{e_x(t) - e_{x0}}{\delta}\right)}{\left[1 + \exp\left(\frac{e_x(t) - e_{x0}}{\delta}\right)\right]^2} \right]$$

Contributors to Inertia Stress (Strain Acceleration)

- Large Deformation
- Poisson's Ratio Change

Experimental Verification and Conclusion



Inertia stress induced by Poisson's ratio change is dominated during onset of densification (large strains)

Effect is worse at higher strain rates: almost no spike seen at 2000 s^{-1} , larger spike at 4000 s^{-1} ($\dot{e}_x^2(t)$)

This method can be used to correct inertia-induced stress spike seen in foam data

New radial inertia analysis is comprehensive especially accounting for Poisson's ratio change

Conclusions



- Optical methods were used to measure deformation of silicone foam
- Radial-axial strains were collected and fit using a Boltzmann sigmoid function
- Tangent method to calculate Poisson's ratio for large strain, nonlinear deformation
- Overall, foam transitioned from being compressible to nearly incompressible with densification
- Unique Poisson's ratio behavior resulted in radial inertia during tests
- Comprehensive radial inertia analysis for SHPB experiment
- Specimen strain (large strains) and Poisson's ratio change contributed to radial inertia in sample
- Numerical correction of stress-strain curve is possible but must be refined

Questions



Radial Laser Calibration

