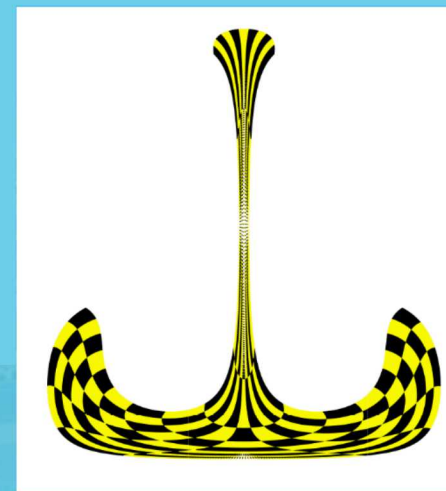
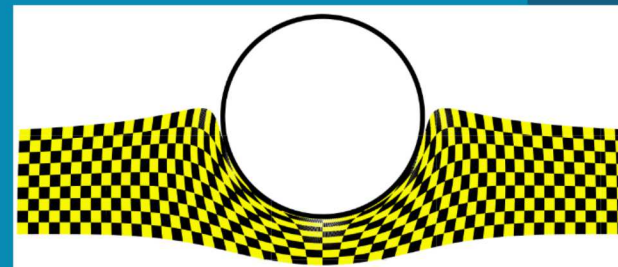


A comparison of mesh-free and mesh-based Lagrangian approximations of manufactured shear-dominated deformation fields

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Outline

1. Value proposition for meshfree
2. Modeling challenges for meshfree
3. Examples of limitations of mesh-based discretizations for extreme events
4. Verification example for exploring robustness of FE vs meshfree discretizations
5. Summary

Value proposition for meshfree

- ease of discretization
 - use of point cloud and surface representation
 - robust large deformation
 - shape functions defined on reference configuration
 - local adaptivity
 - hp adaptivity – easy to change approximation order
 - conversion from Lagrangian to semi-Lagrangian
 - “seamless” update of connectivity without tangling
- } focus of this talk

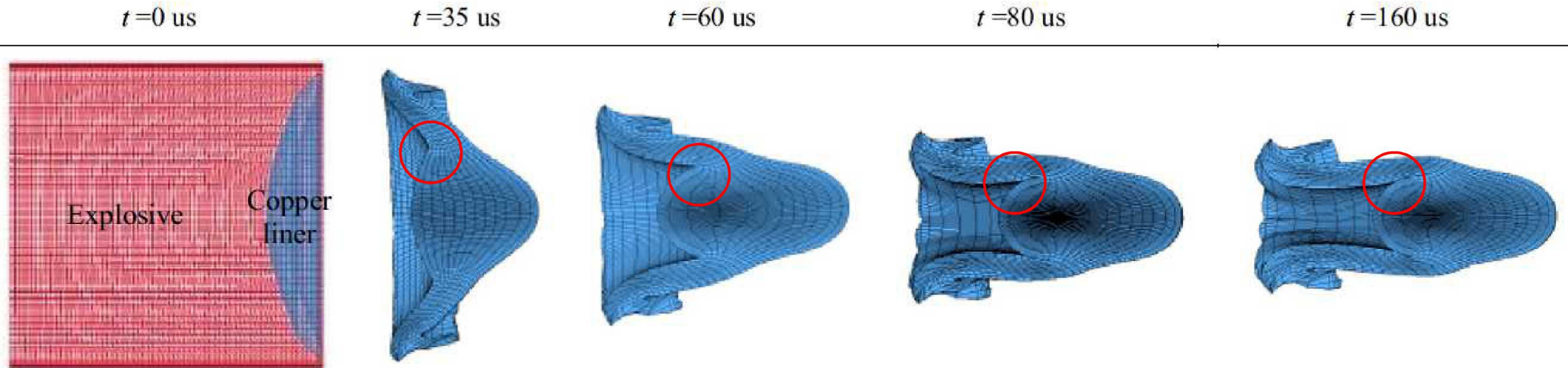
Modeling challenges for meshfree

- surface representation
 - explicit vs. implicit representation
 - adaptivity on surface
 - new surface generation for fracture
- nonconvex domains
 - weight functions around re-entrant corners
- assemblies of parts, material interfaces
- consistent quadrature, stabilization
 - stress points vs. nodal integration
 - integration consistency

Large deformation example

EFP formation process with copper liner

(Liu, et.al, 2017, International Journal of Impact Engineering, 109, 264-275)



Shear localization

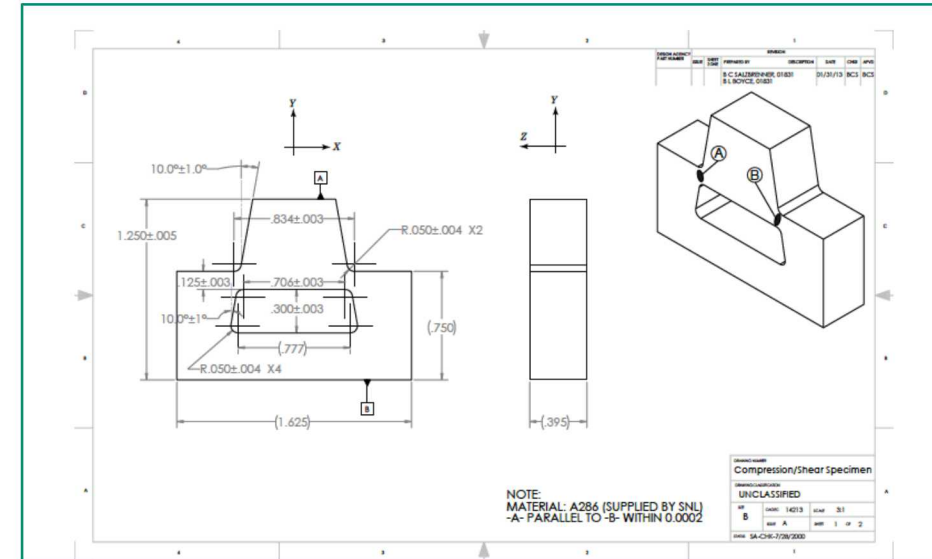
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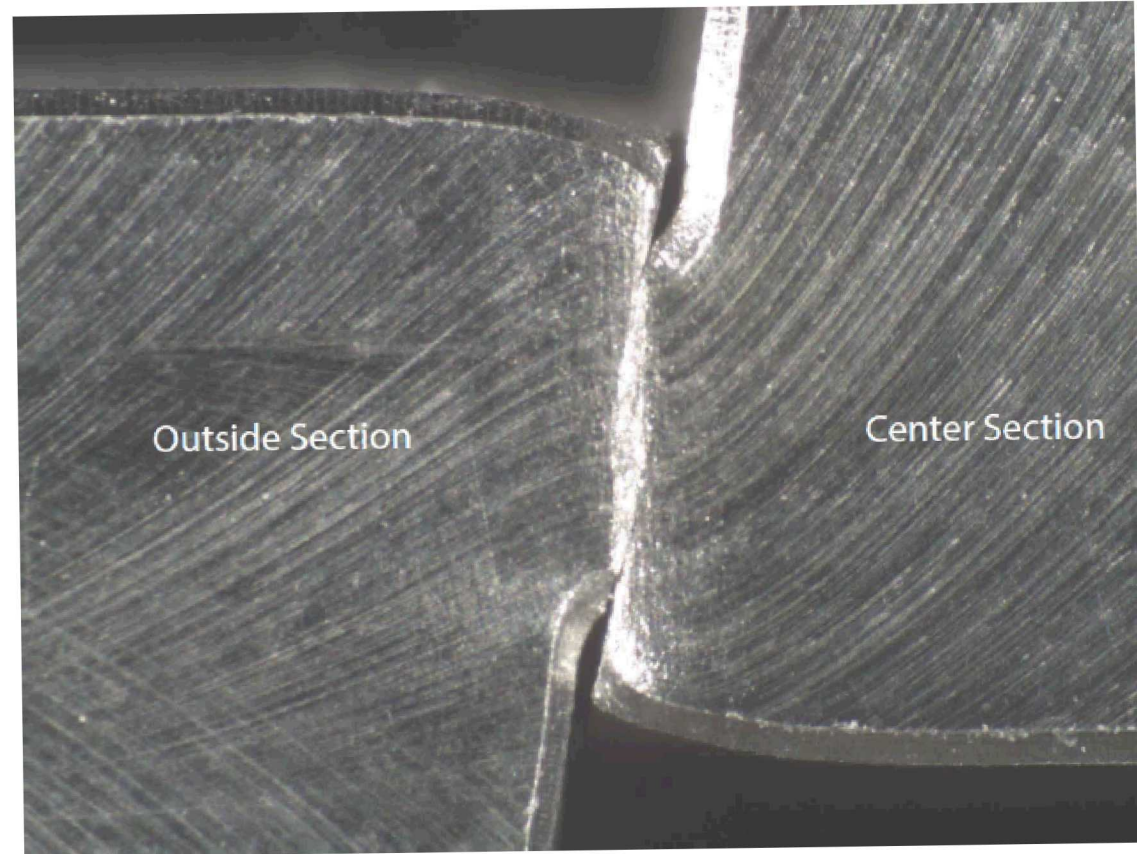
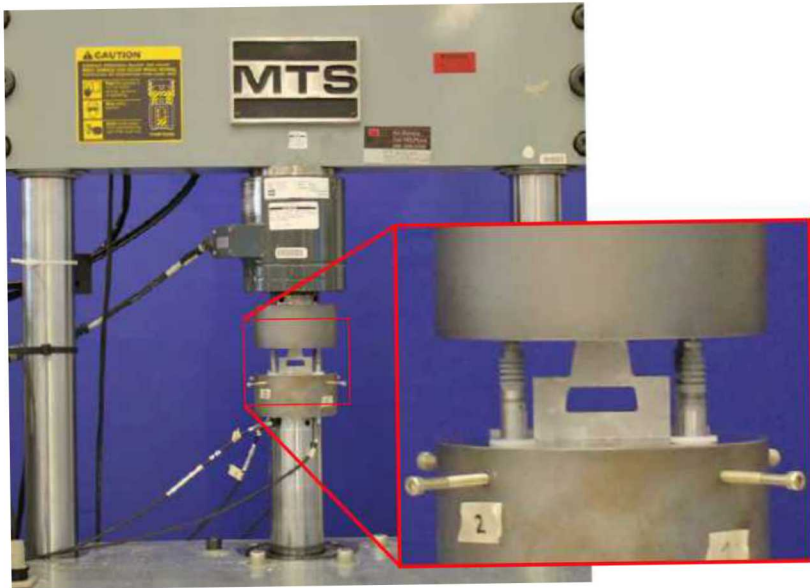
An Experimental Study of Shear-Dominated Failure in the 2013 Sandia Fracture Challenge Specimen

Edmundo Corona, Lisa A. Deibler, Benjamin Reedlunn, Mathew D. Ingraham and Shelley Williams

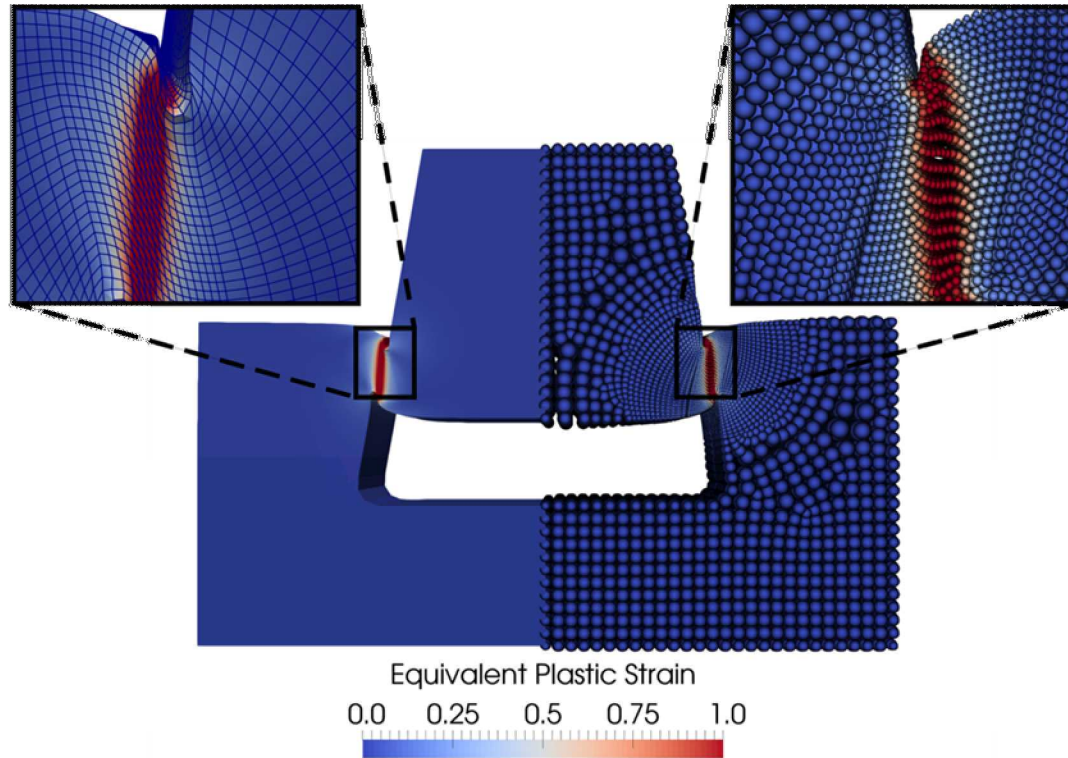
- A286 steel
- 7075 aluminum



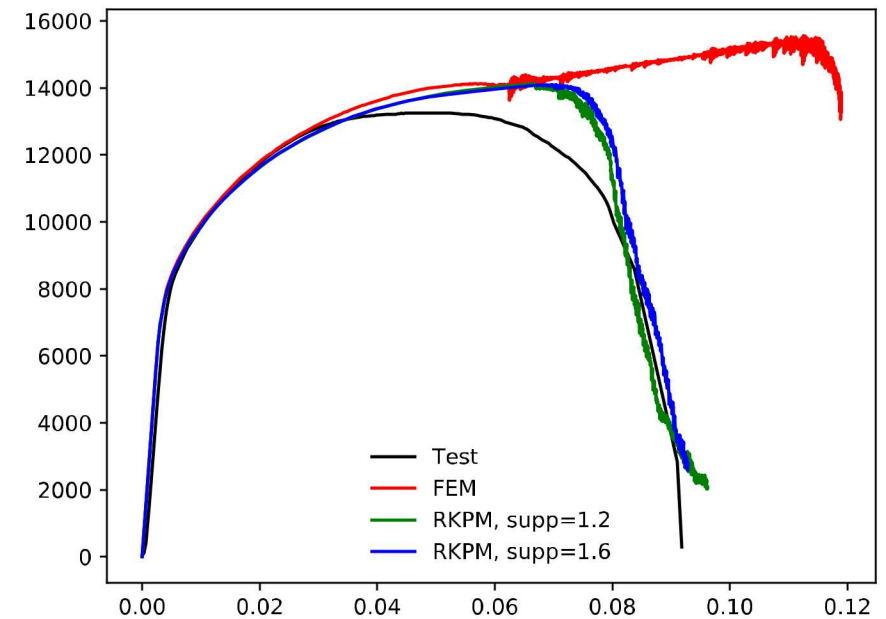
Shear localization



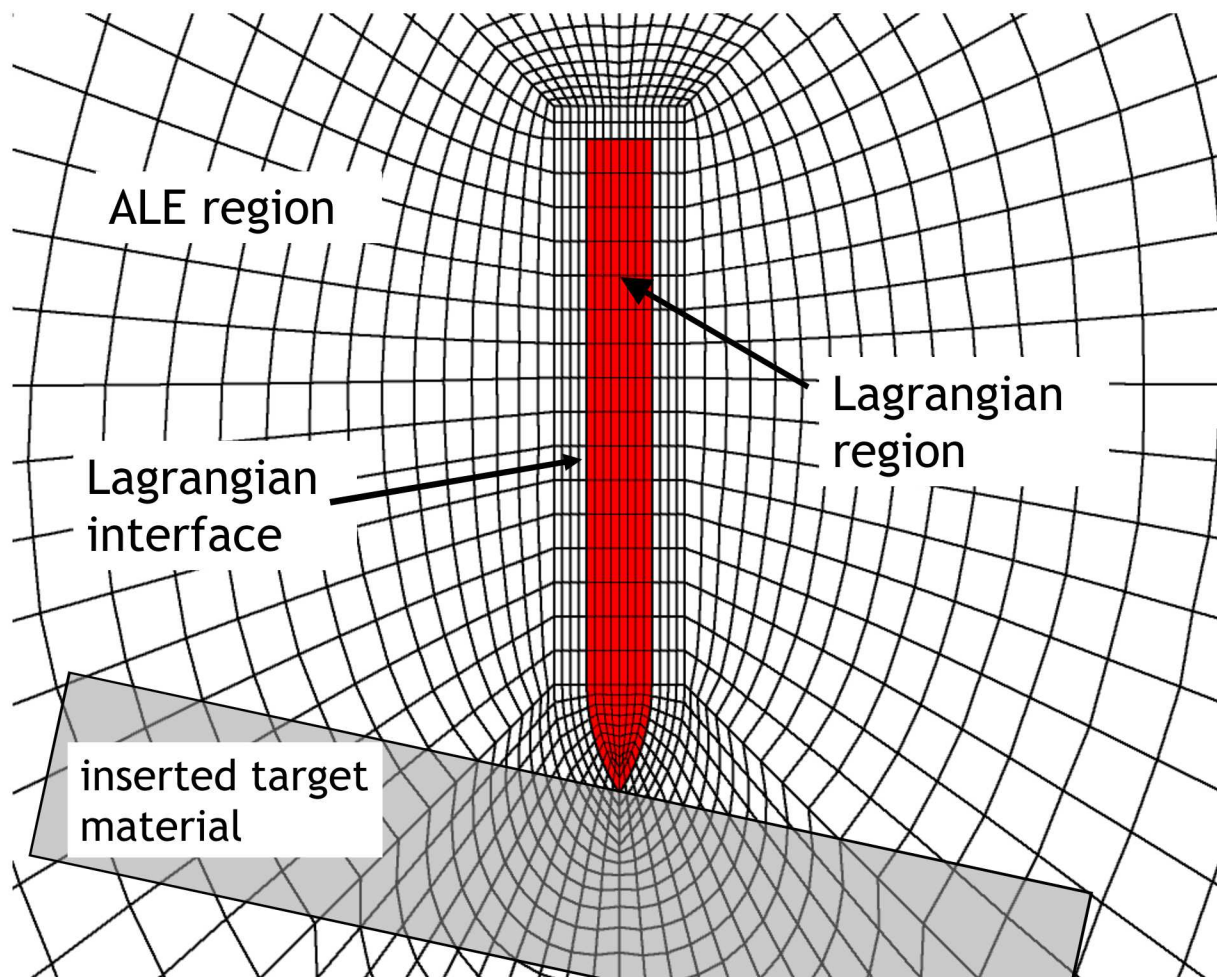
Sierra/RKPM (Jake Koester, Frank Beckwith)



load-deflection response



ALEGRA / SHISM (Soft-Hard Interface Surface Mechanics)



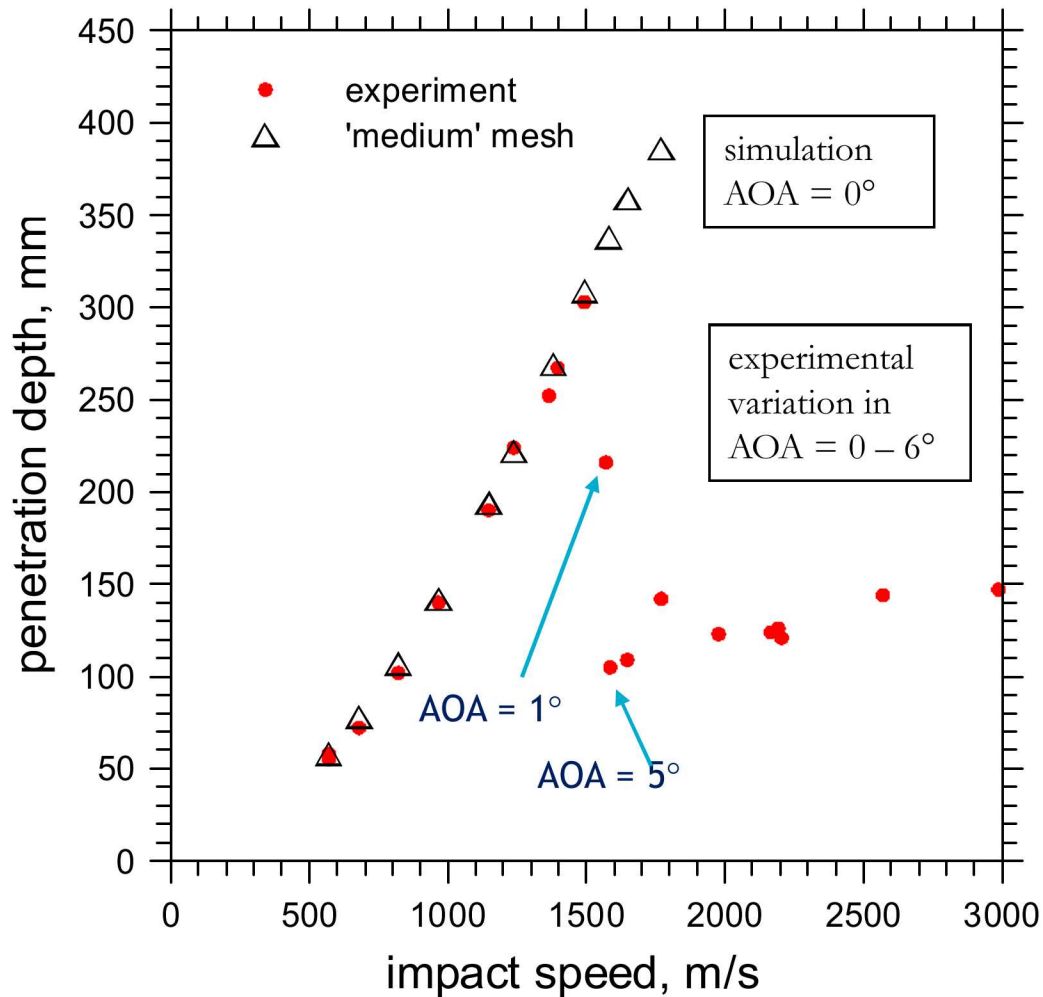
- External ALE mesh 'tracks' the c.m. of the penetrator (translation and rotation).
- Target material 'flows' around the penetrator.
- External mesh interacts with penetrator via contact algorithm.

Assumptions

1. The penetrator is much harder than the target.
2. The target material may flow extensively around the penetrator.
3. There is no penetrator erosion or breakup, but the penetrator may exhibit significant deformation.

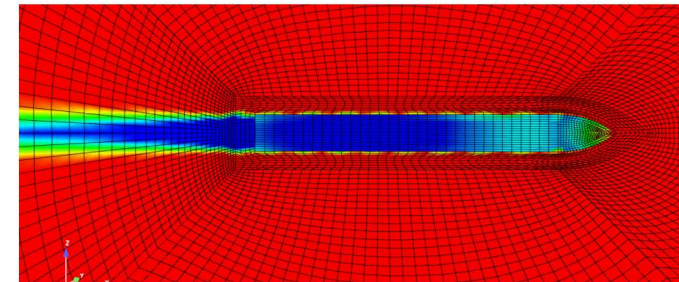
Effect of AOA at high striking velocities

Bishop and Voth, "Semi-Infinite Target Penetration by Ogive-Nose Penetrators: ALEGRA/SHISM Code Predictions for Ideal and Nonideal Impacts," JPVT, 2009. 131

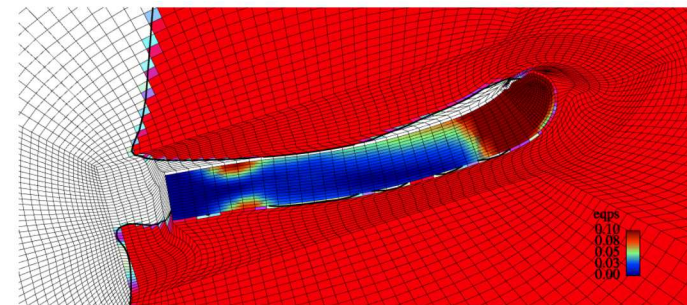


Modest changes in AOA can have catastrophic effects on penetrator deformation.

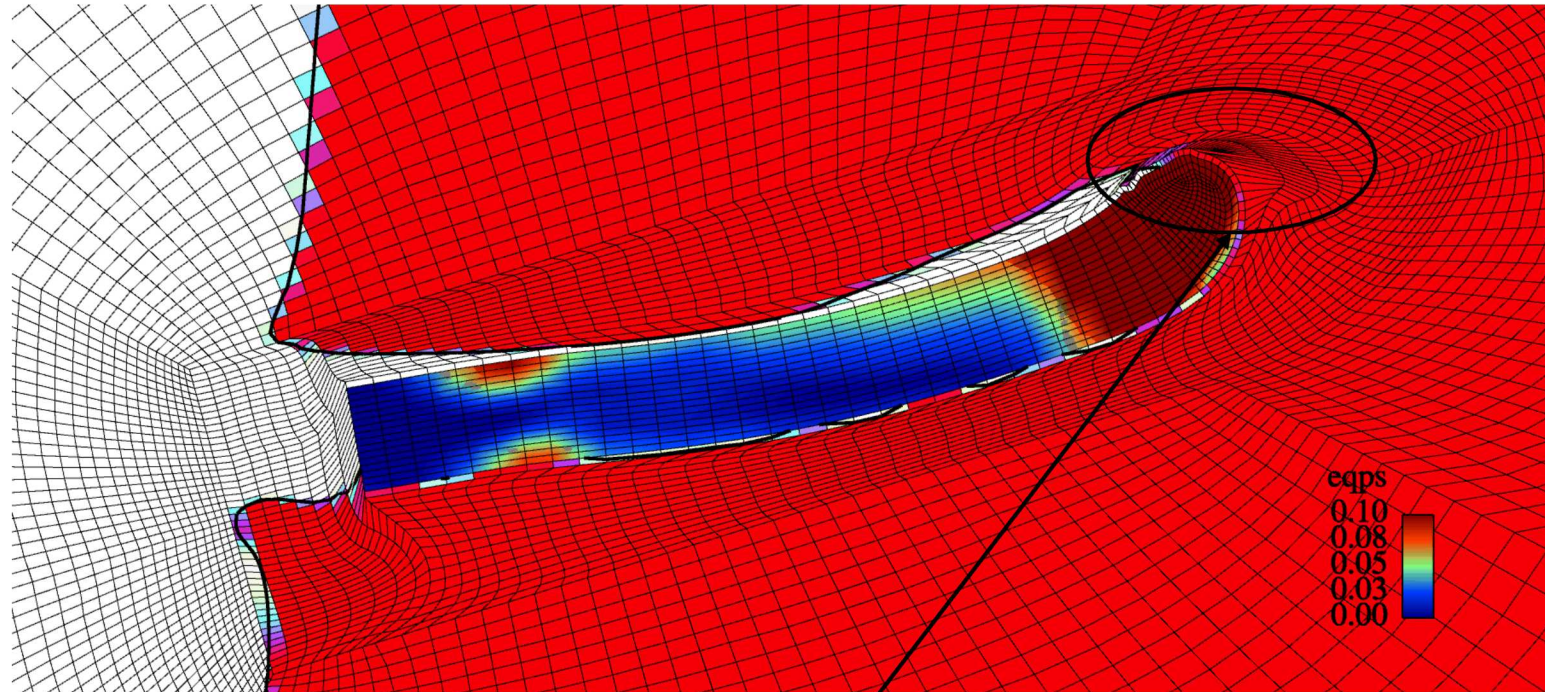
$V_s = 1580 \text{ m/s}$, $\text{AOA} = 0^\circ$



$V_s = 1580 \text{ m/s}$, $\text{AOA} = 2^\circ$



Normal Impact, $V_s = 1580$ m/s, AOA = 2°



severe mesh distortion at $t = 50 \mu\text{s}$
prevents further simulation

Comparison of mesh-free and mesh-based Lagrangian approximations of a manufactured shear-dominated deformation field

- Would like to develop an exact solution to compare FE to meshfree discretizations.
 - Use to compare robustness, sensitivity to initial distortion.
 - Use to explore surface representations.
 - Use to explore adaptivity.
-
- In this talk, explore best-approximation error, FE vs. MLS

Interpolation bounds for isoparametric finite element spaces

for linear elements
(shape regular)

$$||u - \underbrace{I^h}_{\text{interpolation operator}} u||_0 \leq c h^2 |u|_2$$

L^2 norm

$$||u - I^h u||_1 \leq c h |u|_2$$

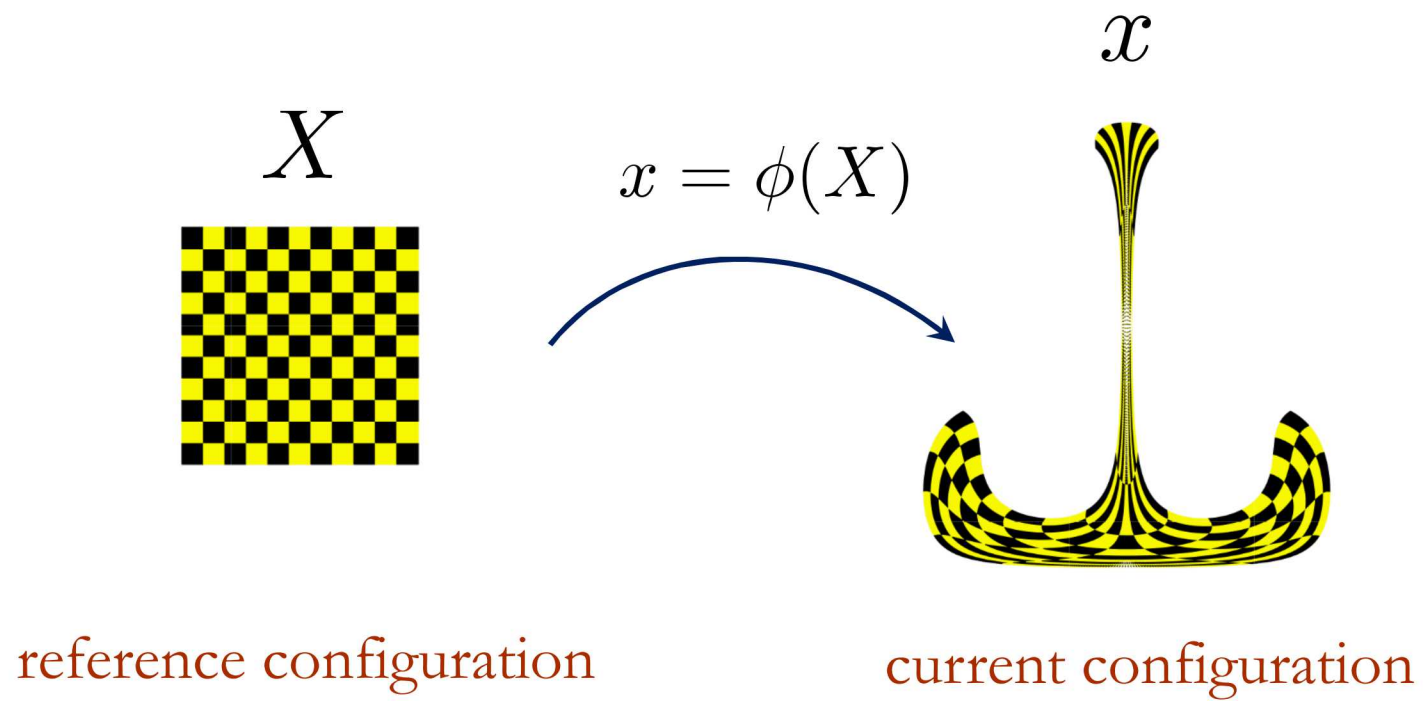
H^1 norm

similar bounds for MLS space

Best approximation

best approximation in H^1 $\hat{u} = \arg \min_{u \in V^h} \|u - u^{\text{exact}}\|_{H^1}$

so $\|\hat{u} - u^{\text{exact}}\|_{H^1} \leq \|I^h u^{\text{exact}} - u^{\text{exact}}\|_{H^1}$



Define deformation map using stream function

$$\Psi = \frac{1}{4\pi} \sin[4\pi(x_1 + 1/2)] \cos[4\pi(x_2 + 1/2)]$$

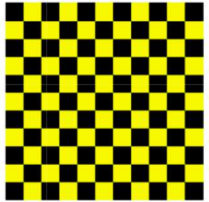
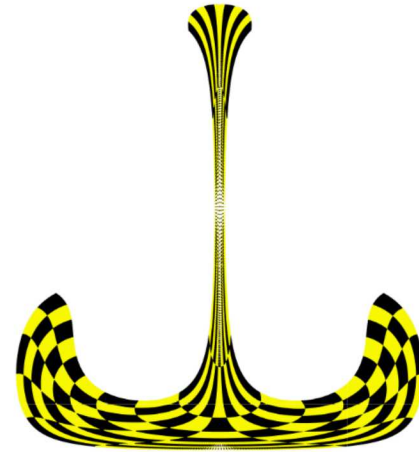
$$v_1 = -\frac{\partial \Psi}{\partial x_2}, \quad v_2 = \frac{\partial \Psi}{\partial x_1}$$

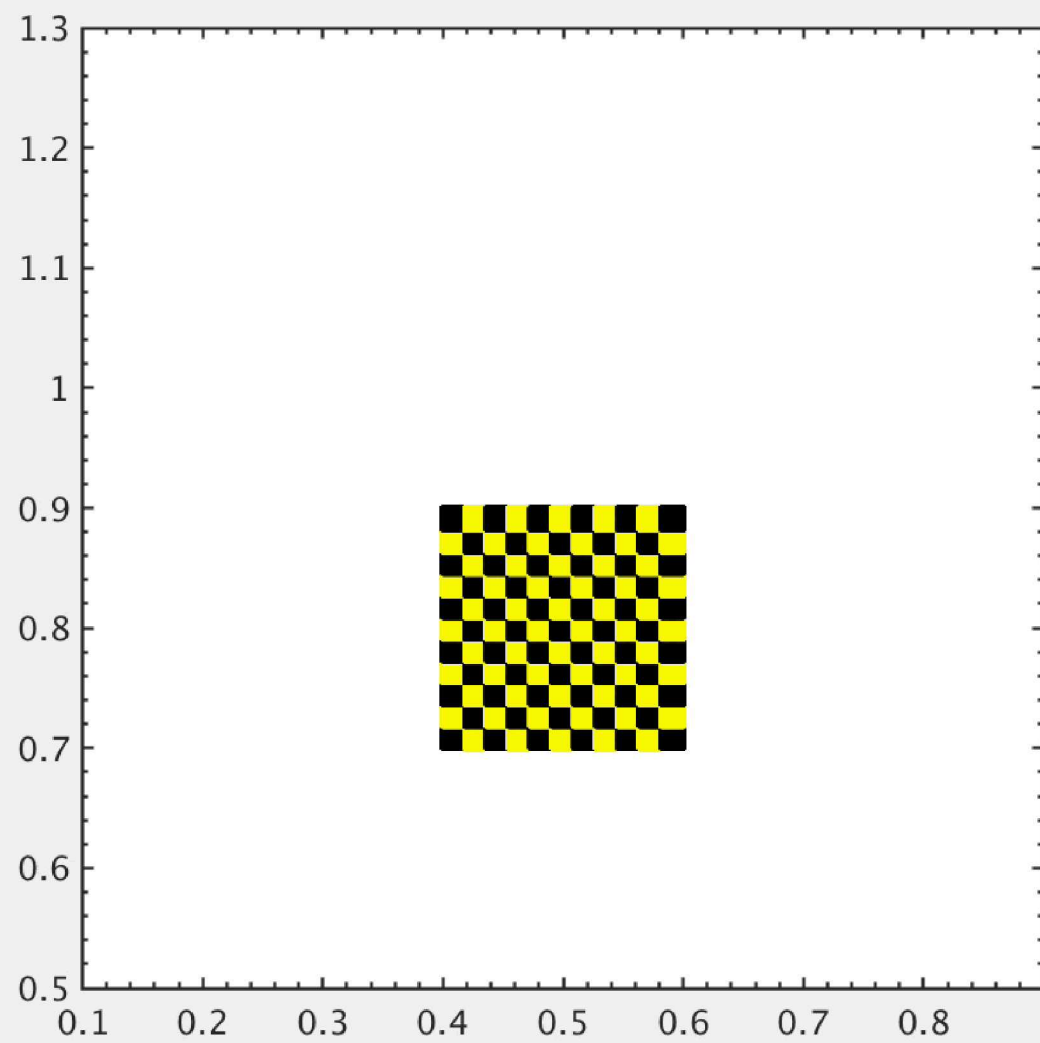
divergence free since

$$\nabla \cdot v = \frac{\partial v_1}{\partial x_1} + \frac{\partial v_2}{\partial x_2} = -\frac{\partial^2 \Psi}{\partial x_1 \partial x_2} + \frac{\partial^2 \Psi}{\partial x_2 \partial x_1} = 0$$

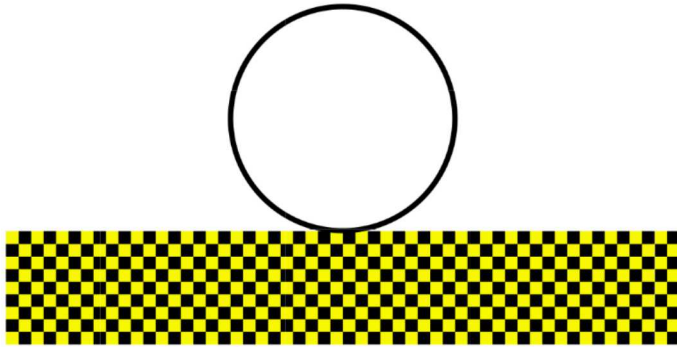
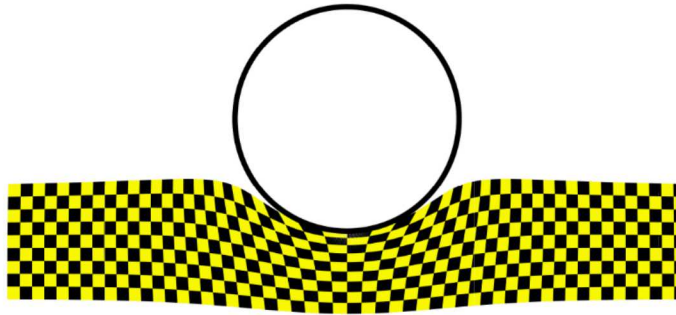
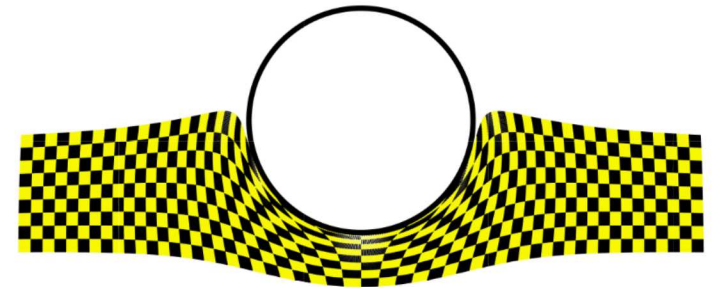
isochoric $J = \det(F) = 1$ for all time F is deformation gradient

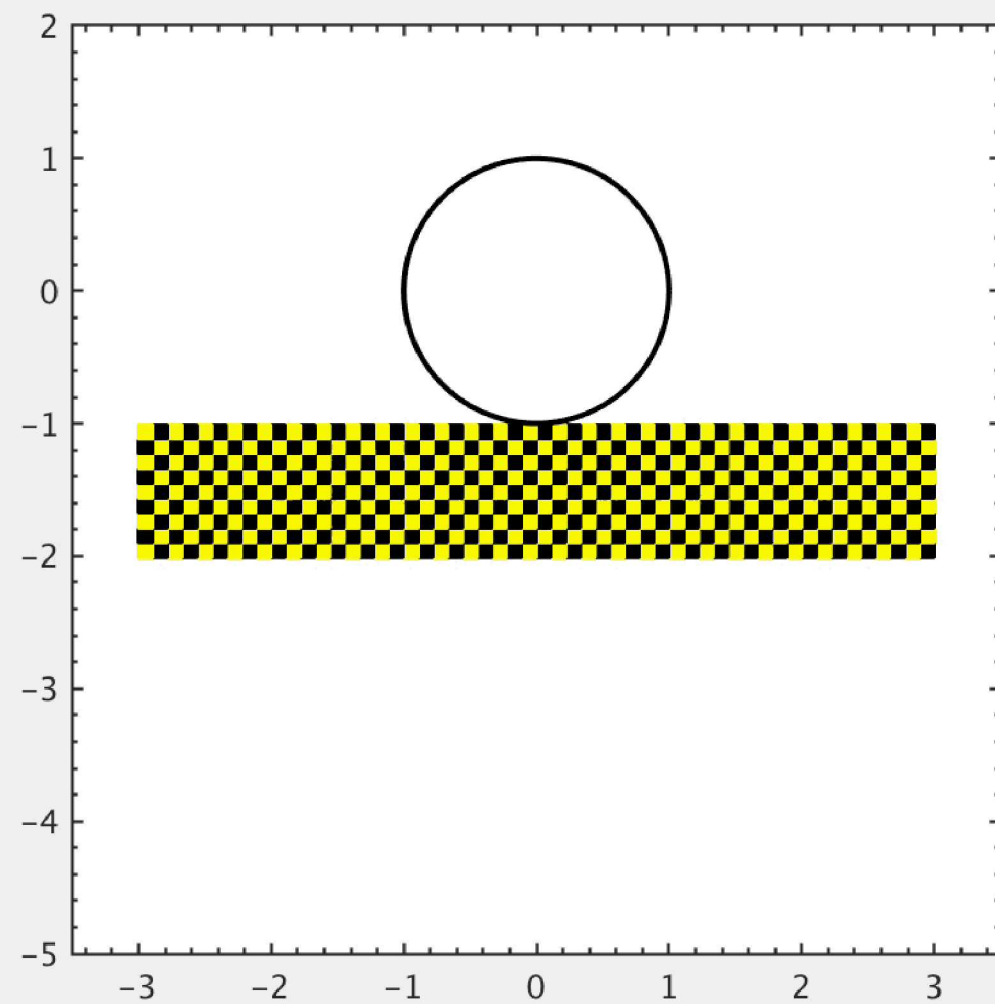
How to convert this Eulerian description to Lagrangian?

$t = 0$  $t = 0.125$  $t = 0.25$ 



$$\Psi = r \sin \theta \left(1 - \frac{1}{r^2} \right) \quad \begin{aligned} r^2 &= x_1^2 + x_2^2 \\ \tan \theta &= x_2/x_1 \end{aligned}$$

 $t = 0$  $t = 0.4$  $t = 0.8$ 



How to convert Eulerian description to Lagrangian?

$$\dot{F} = \frac{\partial v}{\partial X} = \frac{\partial v}{\partial x} \frac{\partial x}{\partial X} = lF \quad l \text{ is velocity gradient}$$

$$l = \begin{bmatrix} l_{11} & l_{12} \\ l_{21} & l_{22} \end{bmatrix} = \begin{bmatrix} \partial v_1 / \partial x_1 & \partial v_1 / \partial x_2 \\ \partial v_2 / \partial x_1 & \partial v_2 / \partial x_2 \end{bmatrix}$$

$$F = \begin{bmatrix} F_{11} & F_{12} \\ F_{21} & F_{22} \end{bmatrix} = \begin{bmatrix} \partial x_1 / \partial X_1 & \partial x_1 / \partial X_2 \\ \partial x_2 / \partial X_1 & \partial x_2 / \partial X_2 \end{bmatrix}$$

$$\begin{pmatrix} \dot{F}_{11} \\ \dot{F}_{12} \\ \dot{F}_{21} \\ \dot{F}_{22} \end{pmatrix} = \begin{pmatrix} l_{11}F_{11} + l_{12}F_{21} \\ l_{11}F_{12} + l_{12}F_{22} \\ l_{21}F_{11} + l_{22}F_{21} \\ l_{21}F_{12} + l_{22}F_{22} \end{pmatrix}$$

$$\begin{pmatrix} \dot{x}_1 \\ \dot{x}_2 \\ \dot{F}_{11} \\ \dot{F}_{12} \\ \dot{F}_{21} \\ \dot{F}_{22} \end{pmatrix} = \begin{pmatrix} v_1 \\ v_2 \\ l_{11}F_{11} + l_{12}F_{21} \\ l_{11}F_{12} + l_{12}F_{22} \\ l_{21}F_{11} + l_{22}F_{21} \\ l_{21}F_{12} + l_{22}F_{22} \end{pmatrix}$$

system of ODEs with initial conditions

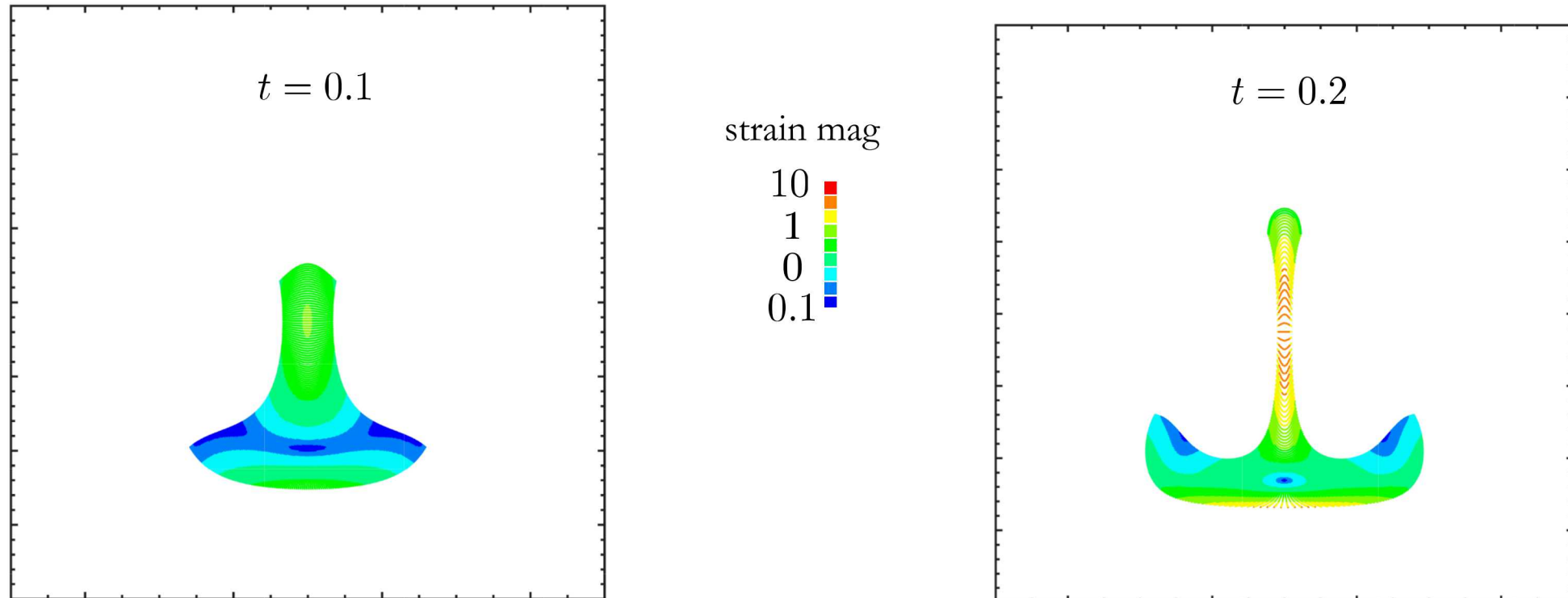
$$x_1 = X_1, \quad x_2 = X_2, \quad F(0) = I$$

This system can be integrated in time to get
Lagrangian description of motion.

Strain magnitude

Lagrange strain $E = \frac{1}{2}(F^T F - I)$

Frobenius norm $\|E\| = \sqrt{E : E}$



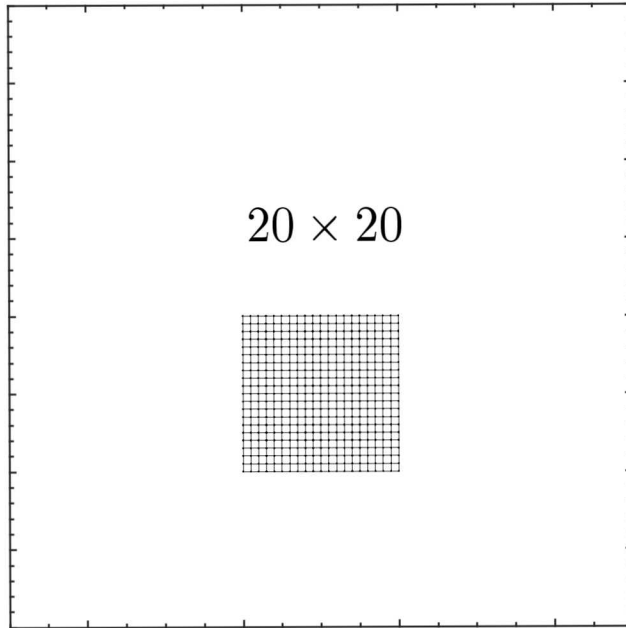


Compare best approximation in H^1 : FE vs. MLS

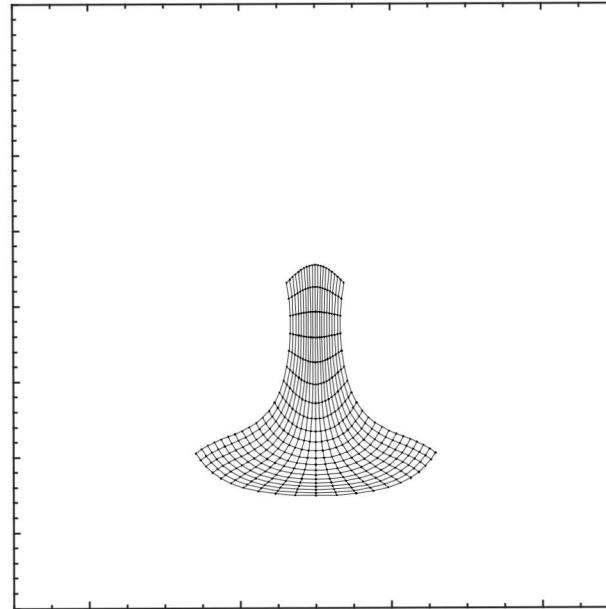
- FE bilinear quad
- MLS: linear reproduction $\{1 \times y\}$, support size = $2 \, dx$
- look at effect of initial mesh distortion

Lagrangian motion (mesh interpolation)

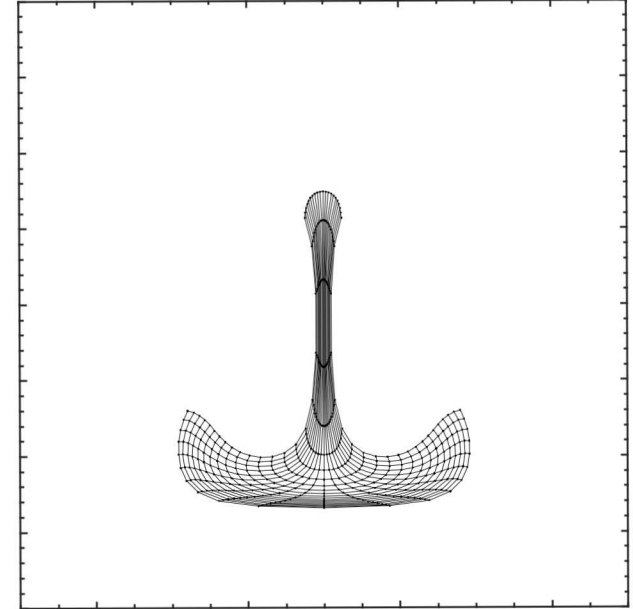
$t = 0$

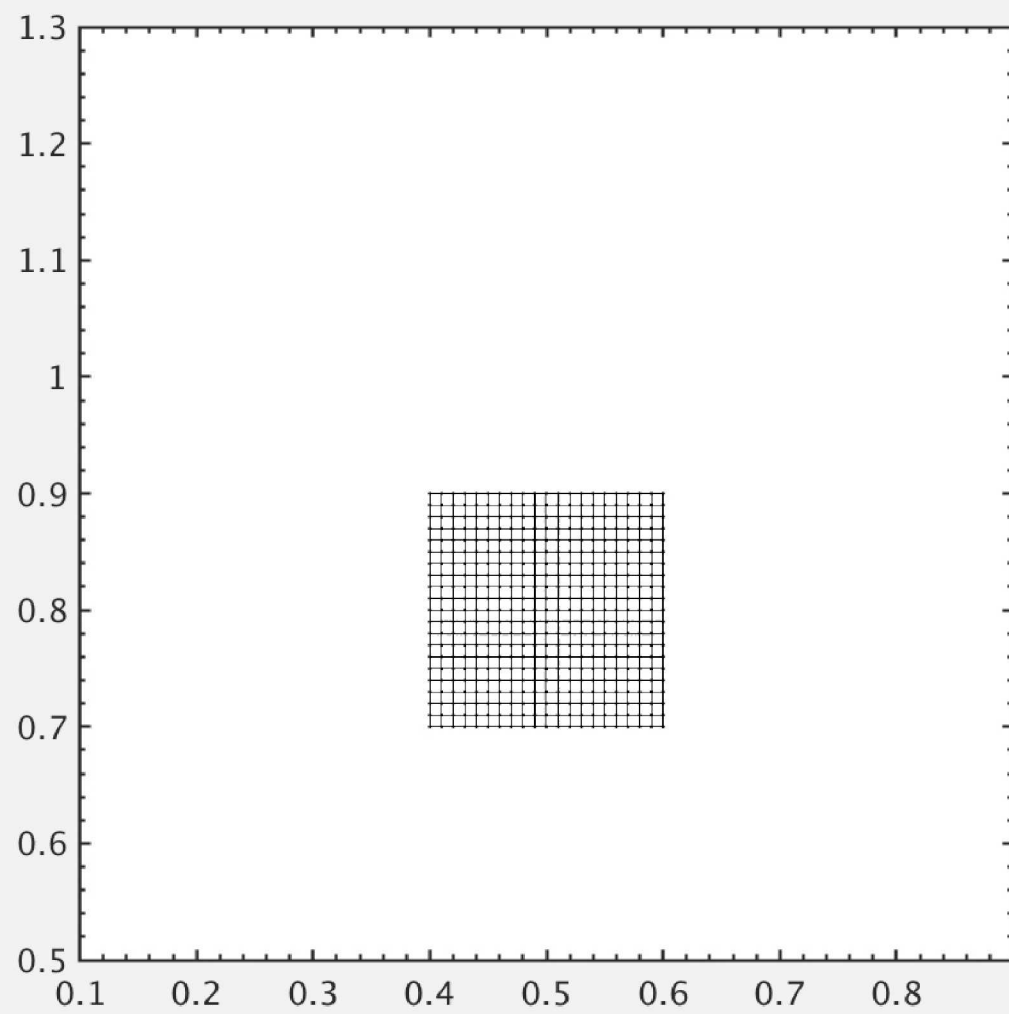


$t = 0.1$



$t = 0.2$





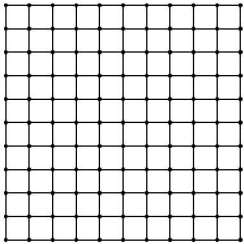
FE vs. meshfree (MLS)

increasing random distortion



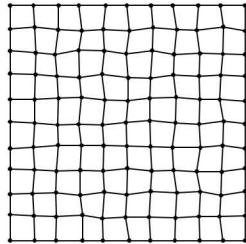
FEA

$r = 0.0$

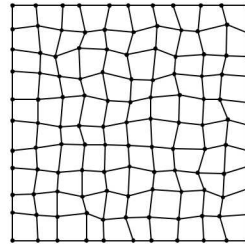


10×10

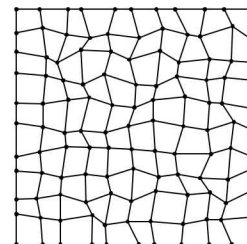
$r = 0.1$



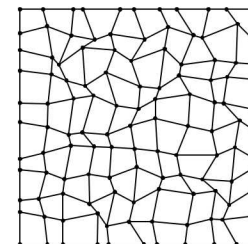
$r = 0.2$



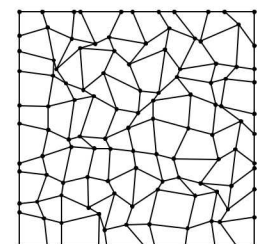
$r = 0.3$



$r = 0.4$

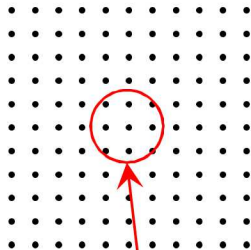


$r = 0.5$

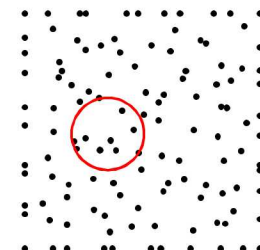
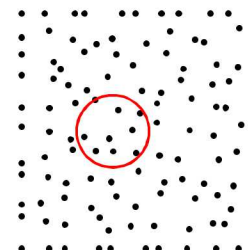
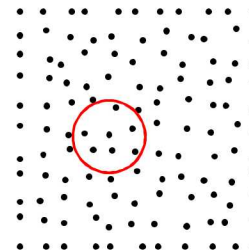
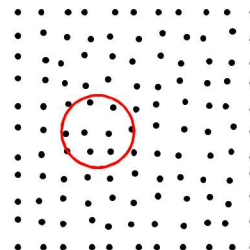
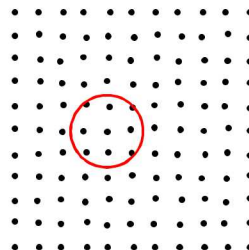


not computable

MLS

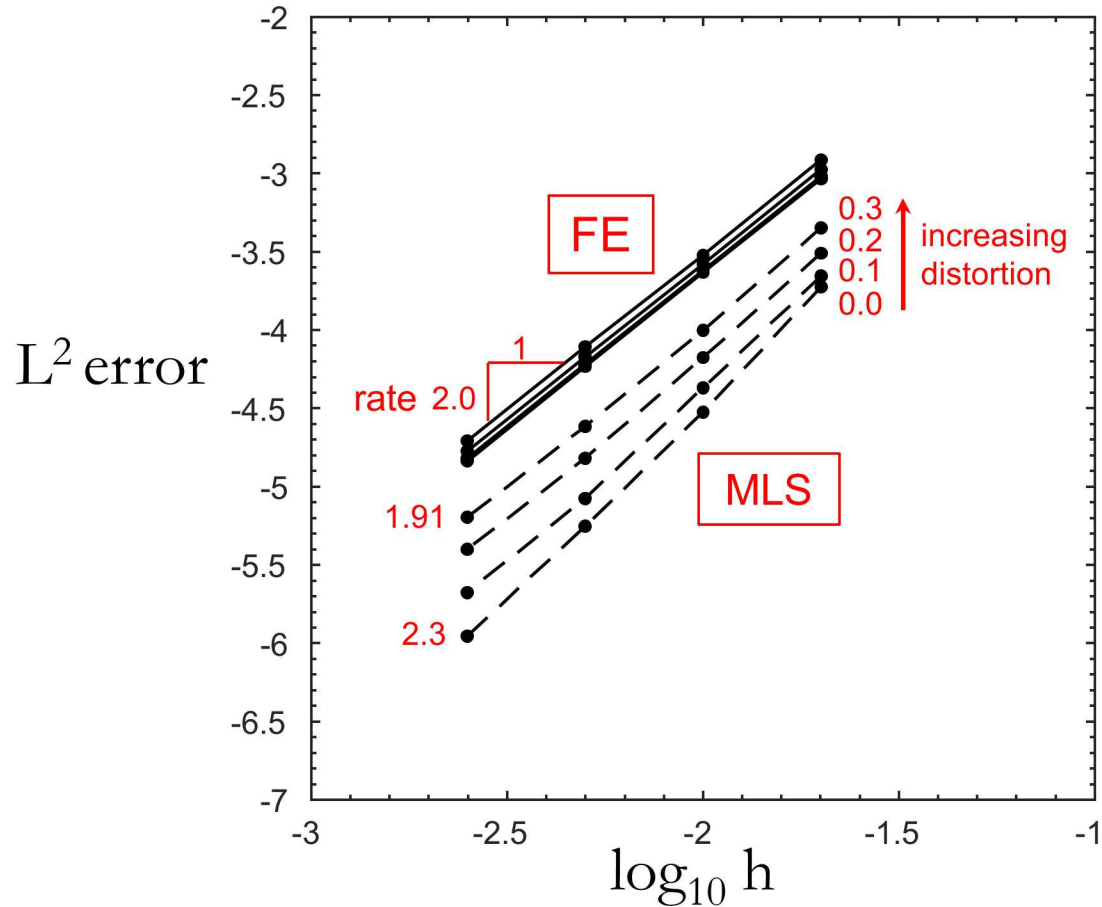


node support

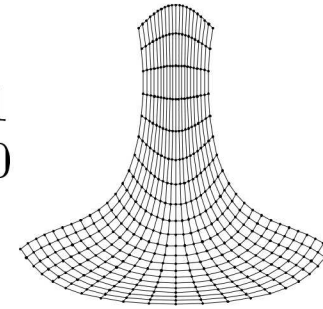


Best-approximation error comparison

L^2 convergence at a fixed time



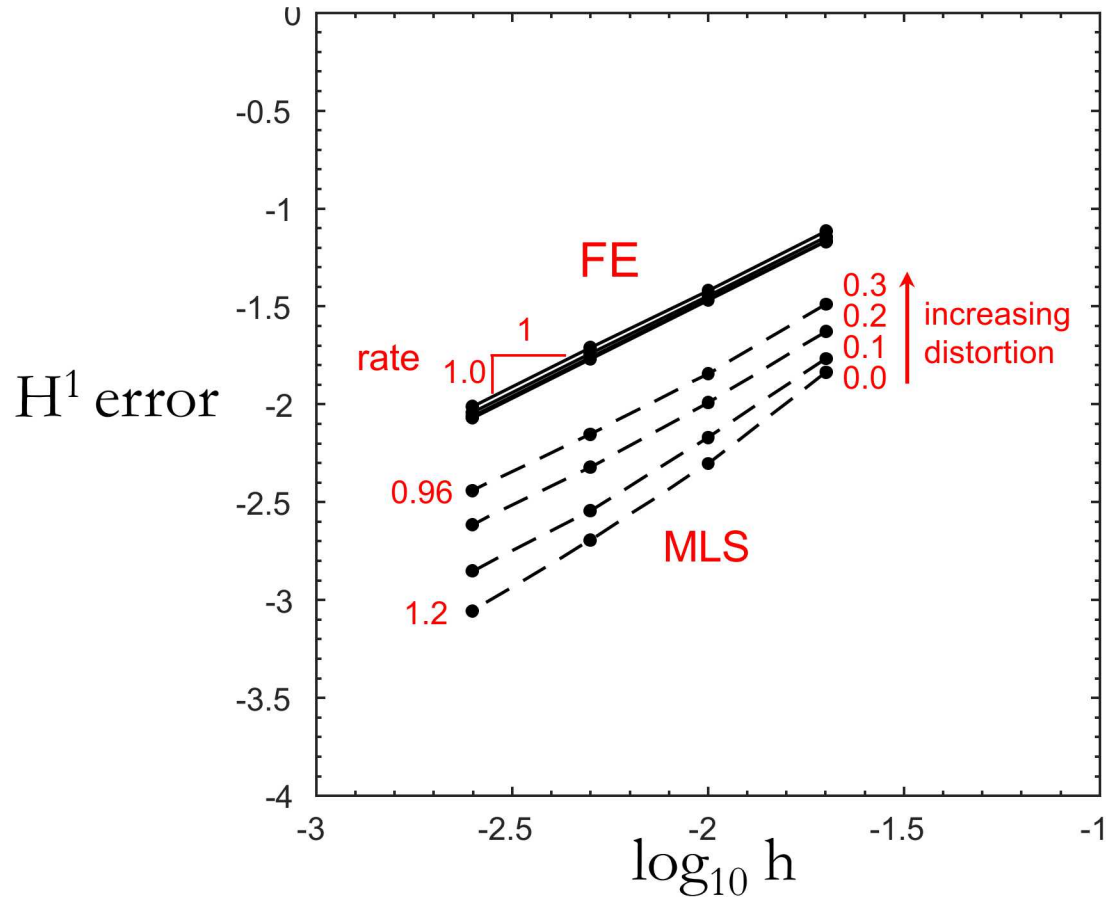
$t = 0.1$
 20×20



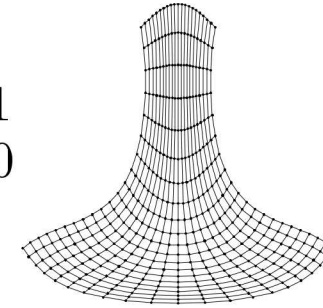
- MLS is significantly more accurate than FE.
- MLS error is sensitive to “distortion” compared to FE.

Best-approximation error comparison

H^1 convergence at a fixed time



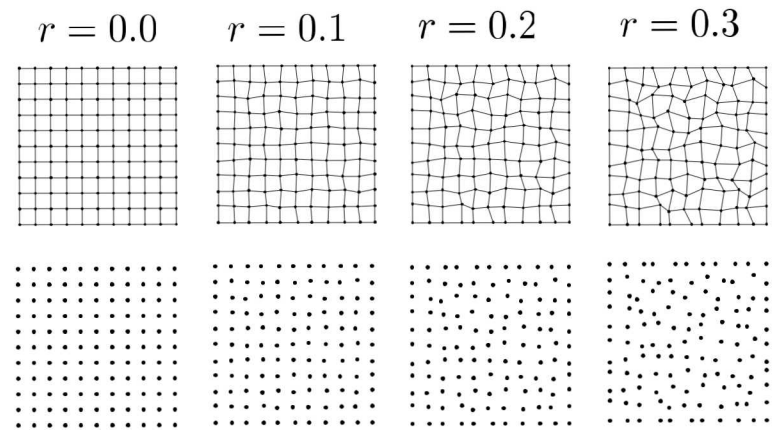
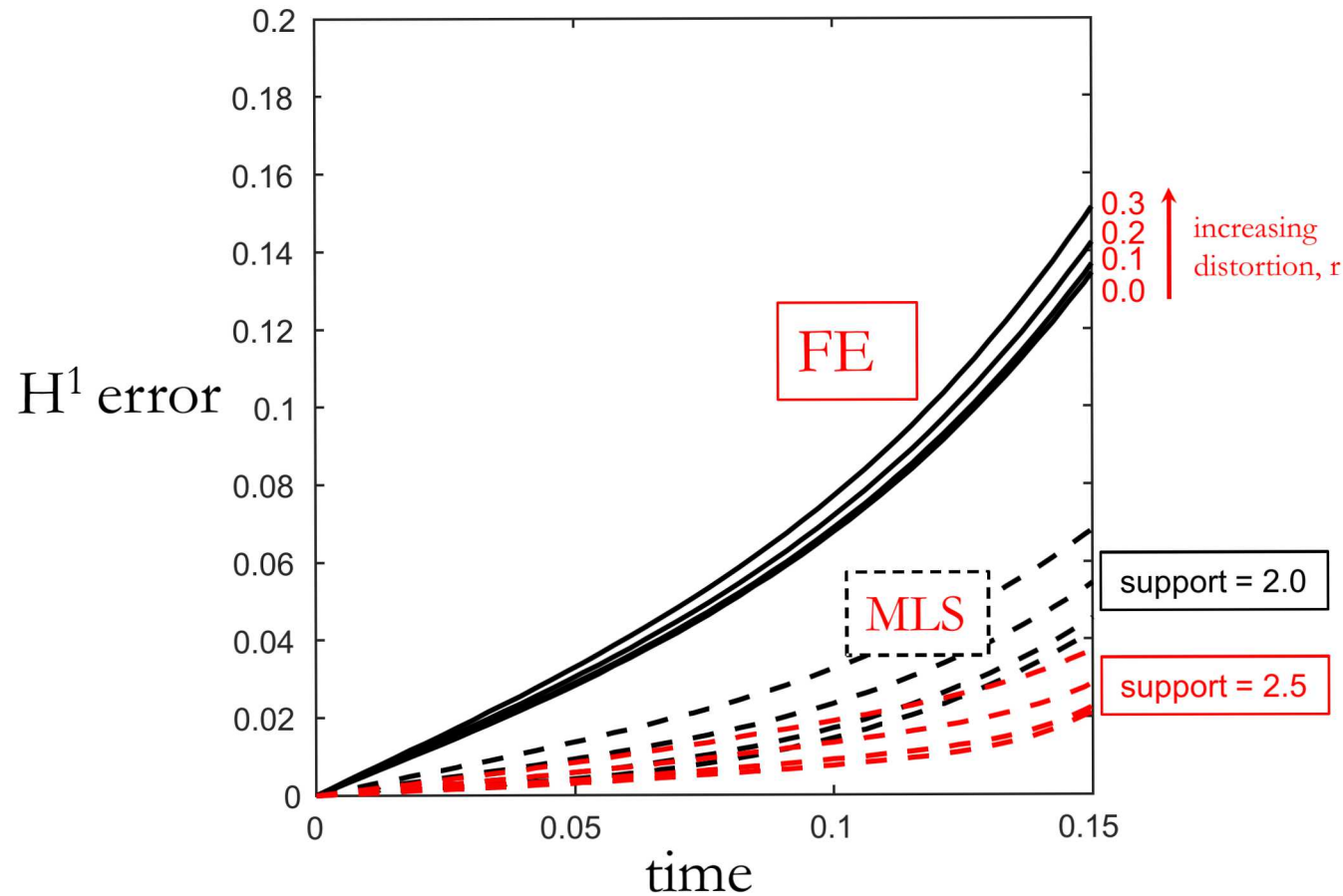
$t = 0.1$
 20×20



- MLS is significantly more accurate than FE.
- MLS error is sensitive to “distortion” compared to FE.

Robustness

H^1 error in time for a fixed mesh



- MLS is significantly more accurate than FE.
- Increasing support size of MLS increases accuracy.

Summary

1. Examples of limitations of mesh-based discretizations for extreme events.
2. Developing manufactured solutions to compare robustness of FE vs. meshfree discretizations in large deformation regime.
3. Exploring best-approximation error in L^2 and H^1 norms.
4. Will also explore surface representations and adaptivity.