



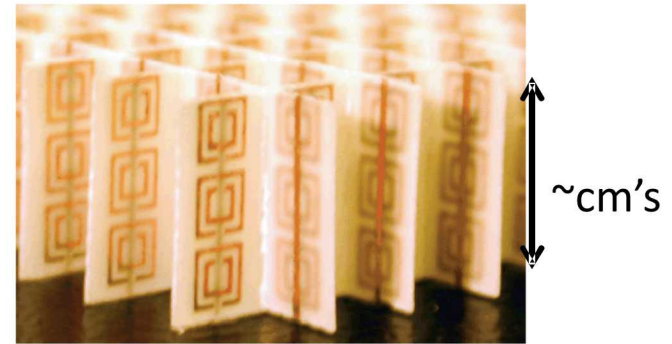
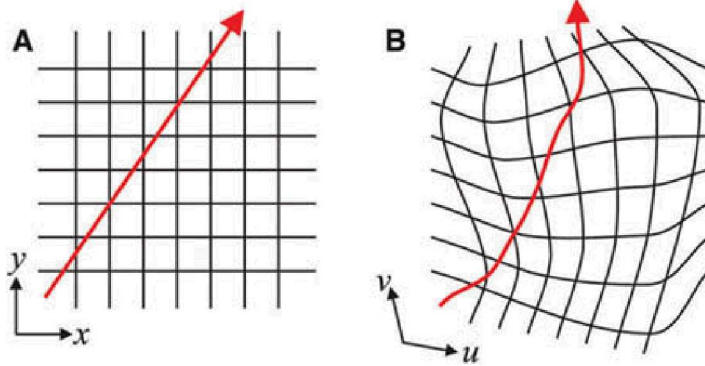
Tunable and Nonlinear Metasurfaces: Metamaterials Meet Optoelectronics

Igal Brener

**Sandia National Labs &
Center for Integrated Nanotechnologies (CINT)**

ibrener@sandia.gov, (505) 844-8097

Metamaterials



http://www.multitel.be/newsletter/mn72005/images/nl3_Metamaterials_type1.jpg

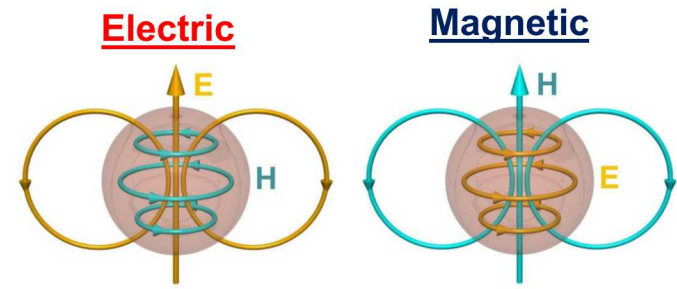
Change

The promise of metamaterials for (IR) applications was great:

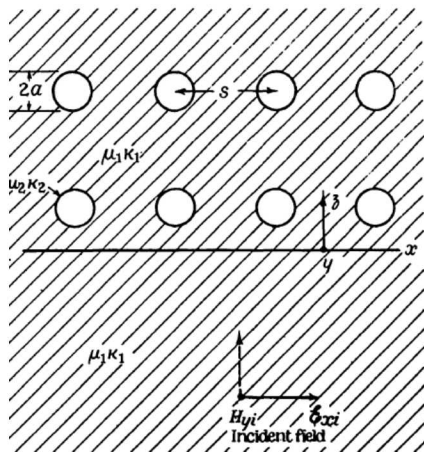
- Wavefront control (Conformal optics, novel illuminators)
- Sub- λ Field Concentration
- Absorption / Emission engineering
- “Flat” optics (thin, lightweight)

Arbitrary ϵ and μ Can be Achieved Also With Dielectrics (low loss approach)

Magnetic dipole resonance: tailor μ
Electric dipole resonance: tailor ϵ



Images: A. Miroshnichenko



THE ELECTRICAL CONSTANTS OF A MATERIAL LOAD SPHERICAL PARTICLES*

By L. LEWIN.†

(The paper was first received 4th March, and in revised form 27th September, 1946.)

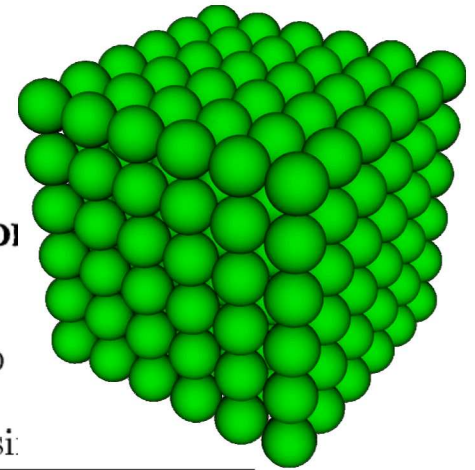
$$\mu'_{re} = \mu_{r1} \left(1 + \frac{3v_f}{\frac{F(\theta) + 2b_m}{F(\theta) - b_m} - v_f} \right)$$

$$\epsilon'_{re} = \epsilon_{r1} \left(1 + \frac{3v_f}{\frac{F(\theta) + 2b_e}{F(\theta) - b_e} - v_f} \right)$$

$$F(\theta) = \frac{2(\sin \theta)}{(\theta^2 - 1) \sin \theta + \theta \cos \theta}$$

$$\theta = k_0 a \sqrt{\epsilon'_{r2} \mu'_{r2}}$$

$$b_e = \frac{\epsilon_1}{\epsilon_2}, \quad b_m = \frac{\mu_1}{\mu_2}$$

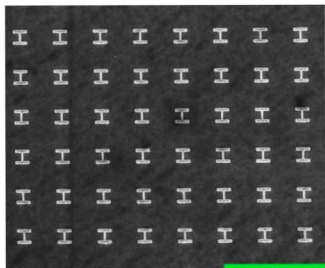


Metallic

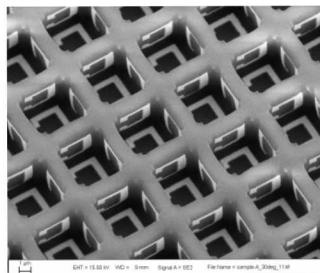
2D

3D

Planar resonators on III-V substrate



Au SRRs on resist "cavities"



>1 layer:
Incredibly difficult
Too lossy



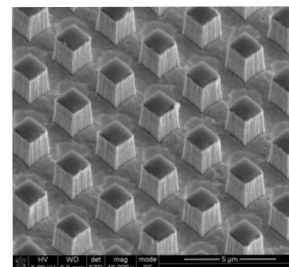
- Strong coupling to phonons, plasmons, intersubband transitions
- Voltage tuning, tunable filters
- Nonlinear optics

Dielectric

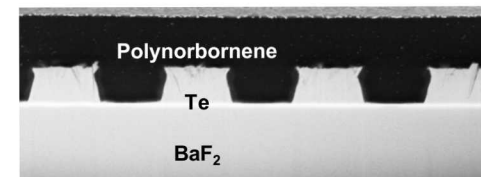
2D

3D

Tellurium cubes on BaF₂



multi-layers of Tellurium cubes?



>1 layer:
Incredibly difficult



- Optical magnetism
- Unidirectional scattering, Huygens metasurfaces
- Nonlinear optics

Metasurfaces

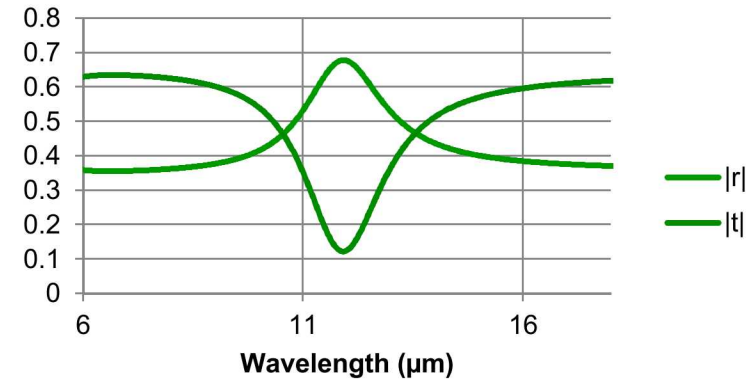
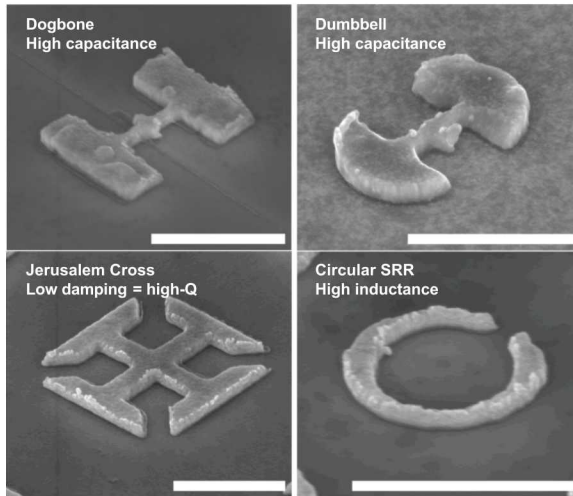
2D Metamaterials (“Metasurfaces”)

(for some applications, loss is a good thing)

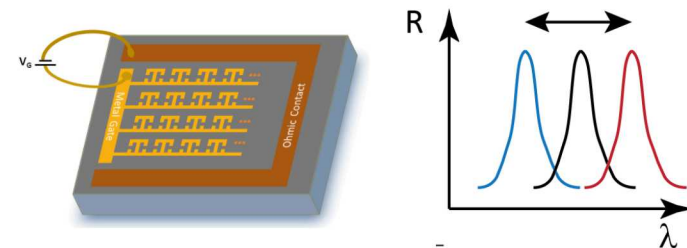
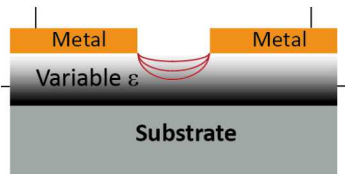
Even for different shapes, many metallic resonators behave like LCR circuits (or “antennas”)

$\sim \lambda/10$

(Sandia)

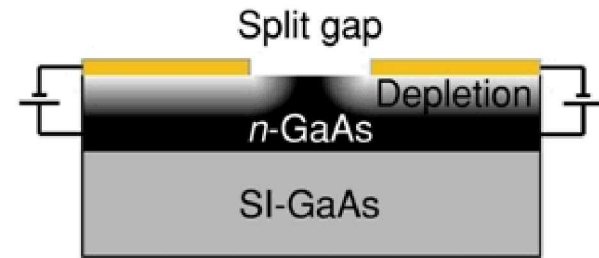
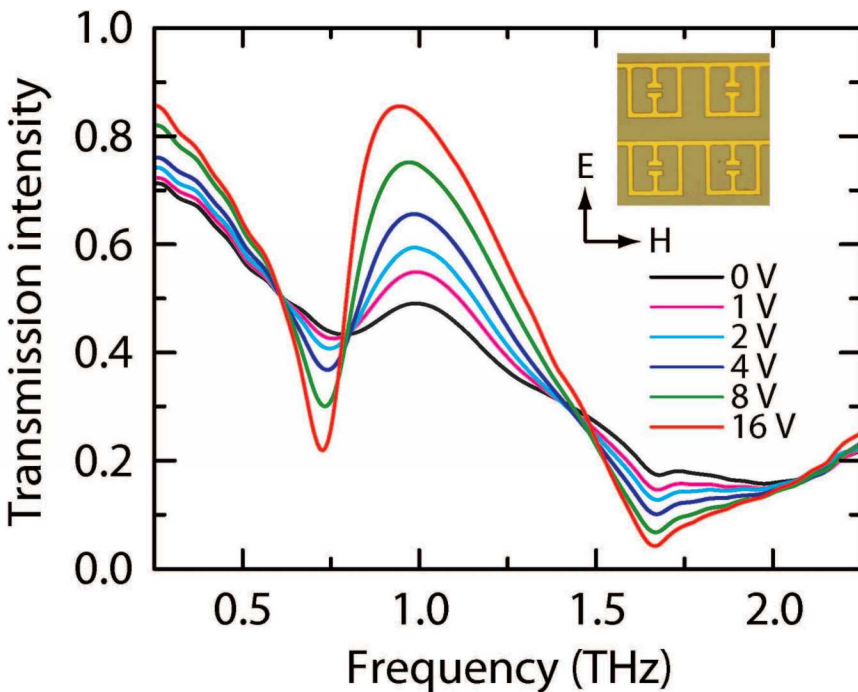


The spectral response is very sensitive to the dielectric environment





An Example at THz frequencies: Electrically Switchable Metamaterials Using Damping



This works by increasing the damping (ϵ_2).

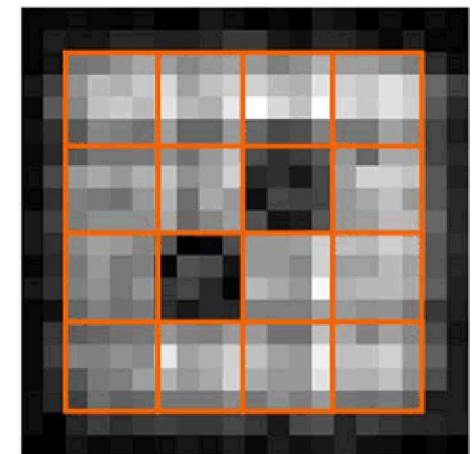
Chen, Nature 444, 597 (2006).

However: this approach doesn't scale well to shorter wavelengths

0.36 THz, 4x4 SLM

(With Rice Univ.)

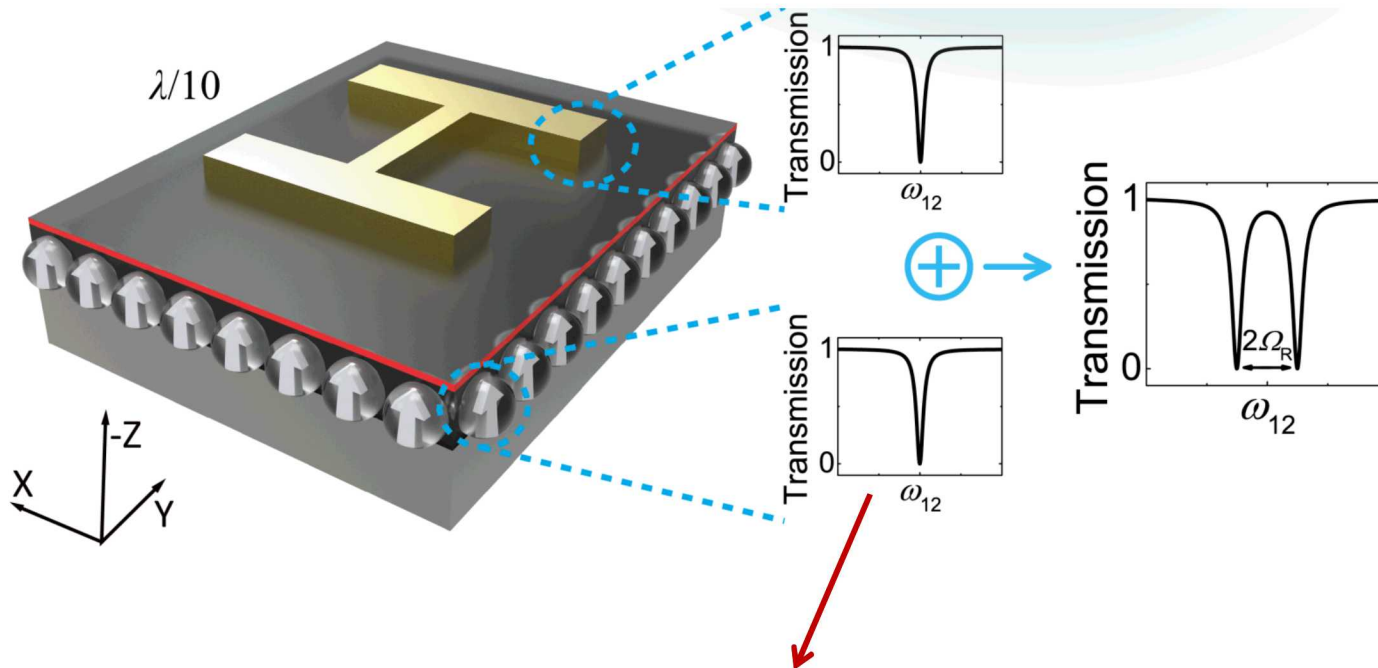
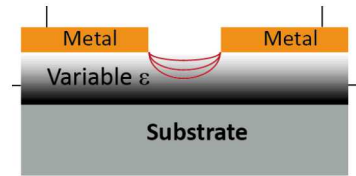
Appl. Phys. Lett. 94, 213511 (2009)



16 mm

Other Tuning Mechanisms: Coupling

MM resonators create strong optical fields that lead to strong coupling



Optical Phonons: Nano Letters 11, 2104 (2011)

Epsilon Near Zero modes: Nano Letters 13, 5391 (2013)

Intersubband Transitions: Nature Communications 4, (2013)

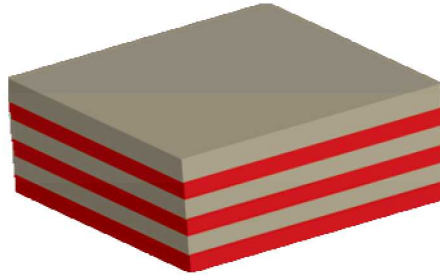
Strong Coupling to Inter-subband Transitions in Quantum Wells



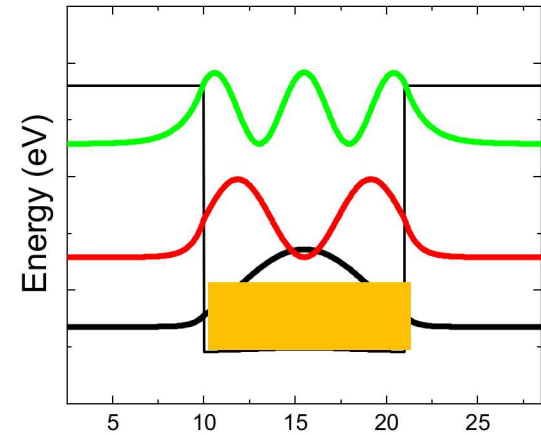
Alex Benz

- Scalable (far IR to near IR), Mature, Versatile

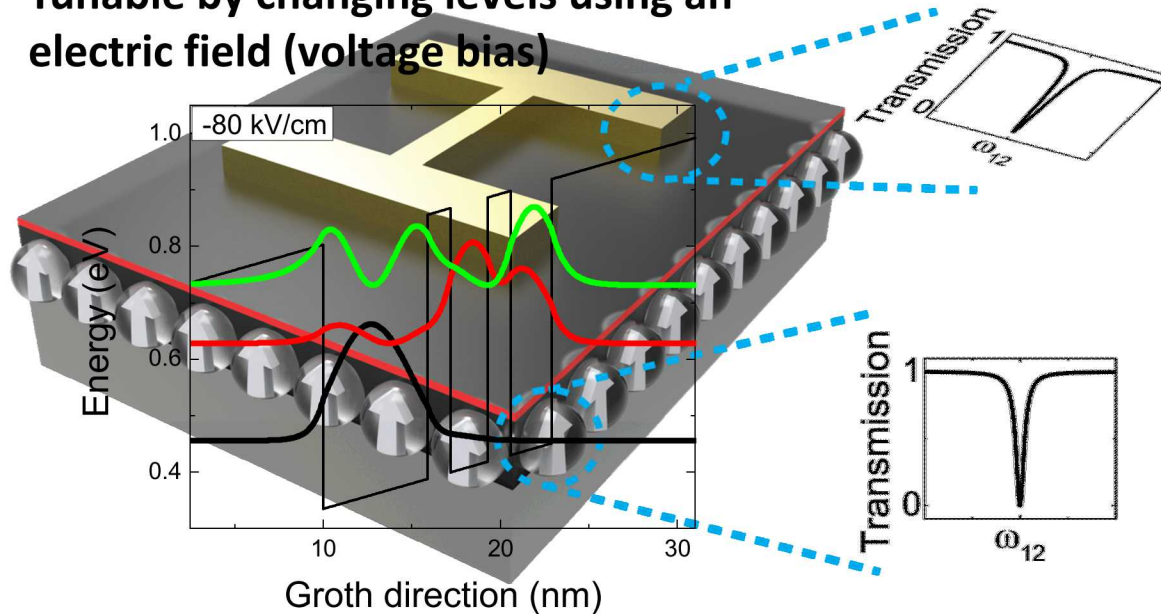
Opt. Express 20, 6584 (2012),
APL 98, 203103 (2011)



$$\epsilon_{xx} = \epsilon_{yy} \neq \epsilon_{zz}$$

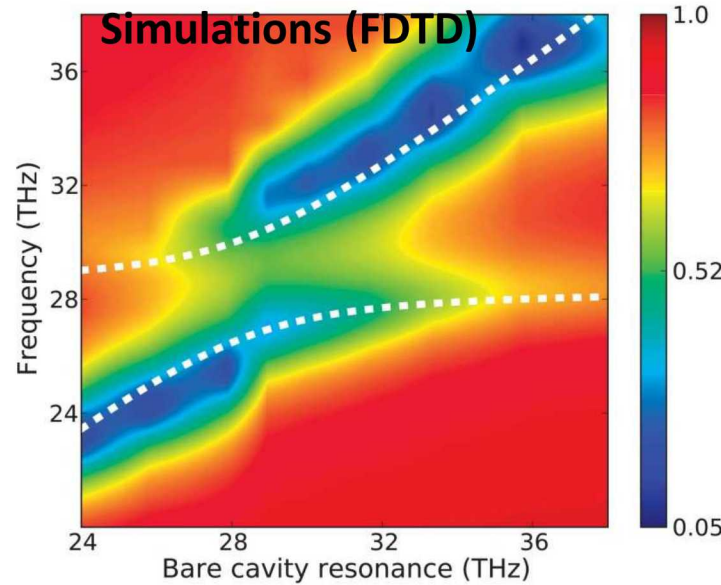
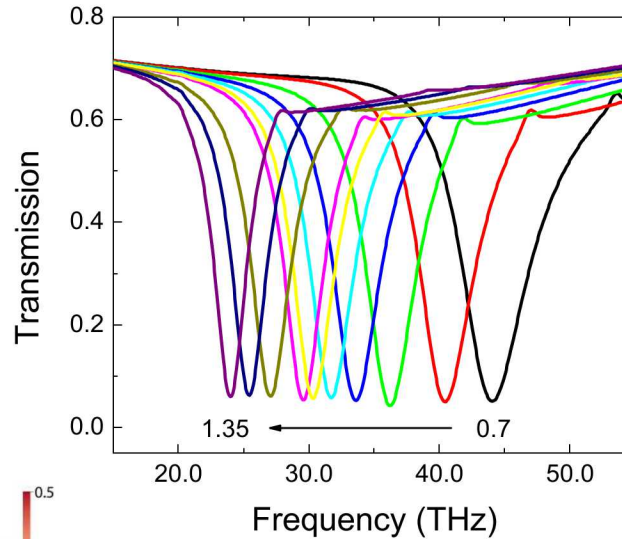
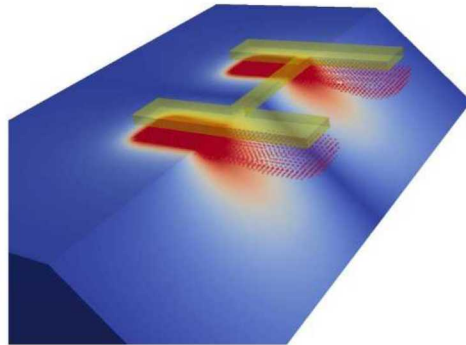
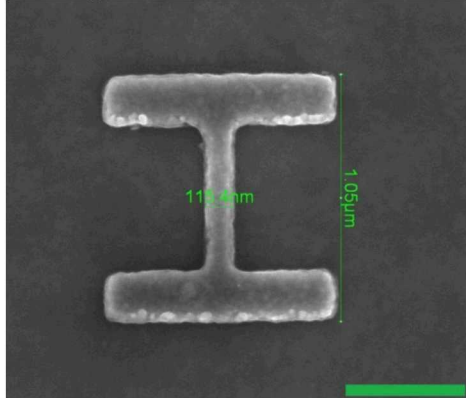


Tunable by changing levels using an electric field (voltage bias)



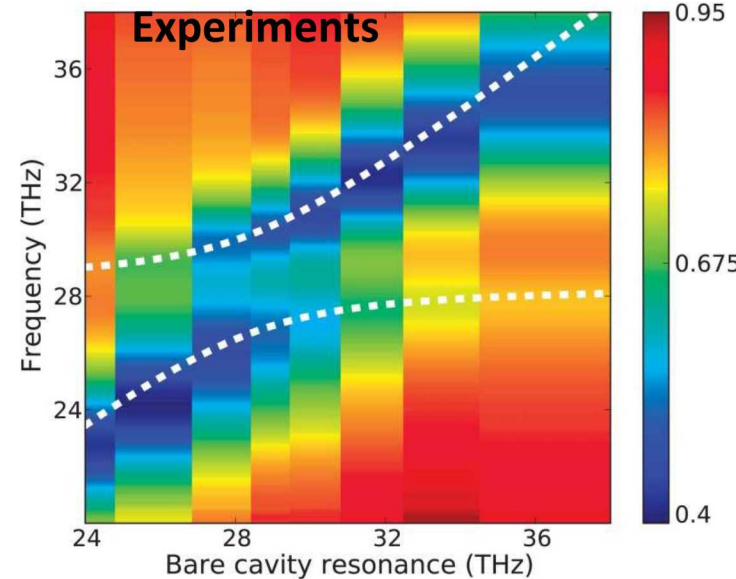
Strong Coupling Theory vs. Experiment

“Dogbone” resonators



Rabi frequency depends on simple parameters

$$\Omega_R = \sqrt{\underbrace{f_W}_{\text{Geometry factor}} \omega_p}_{\text{Plasma frequency}} / 2$$

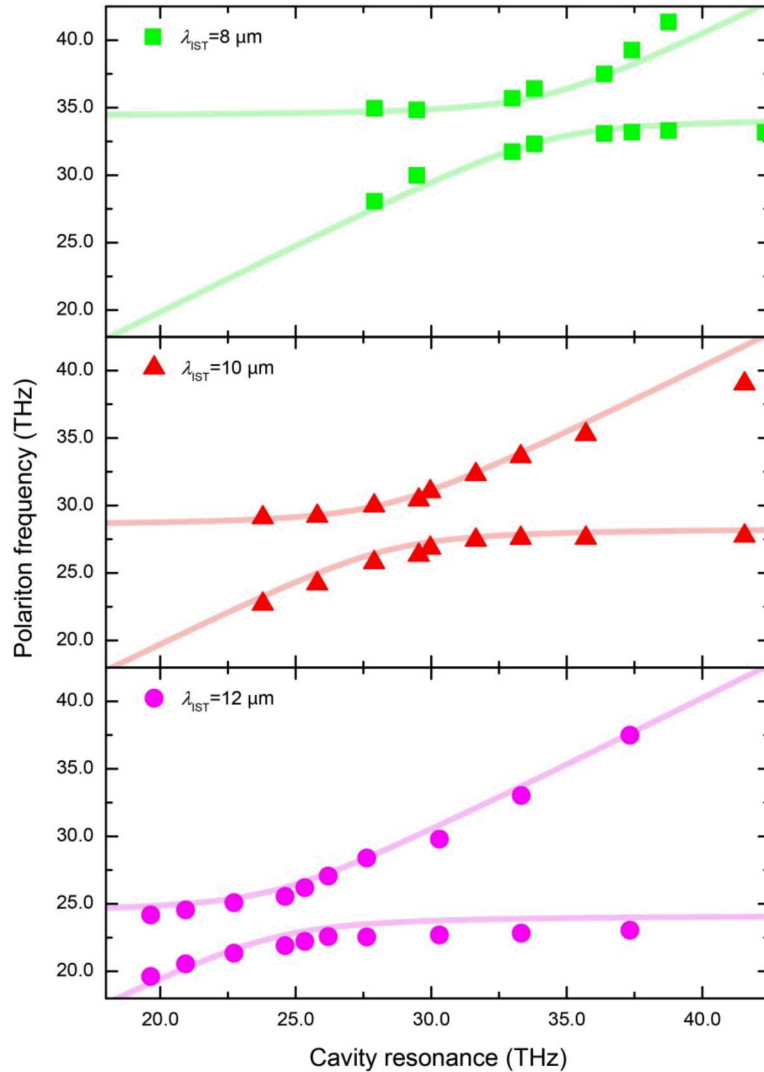




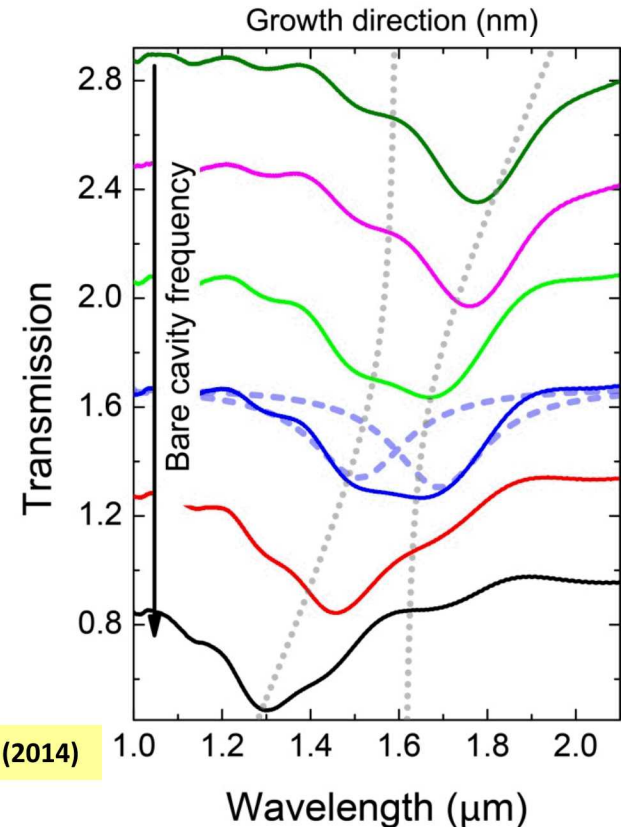
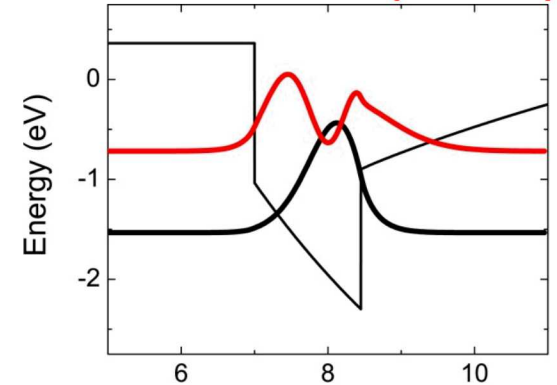
Alex Benz

From Mid-IR to Near IR

InGaAs QWs (mid IR)



GaN QWs (near IR)

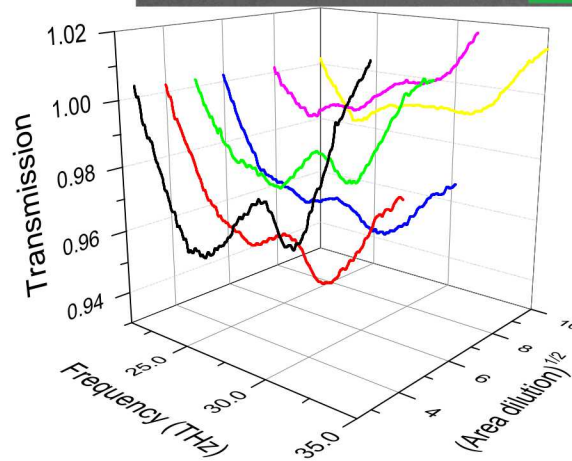
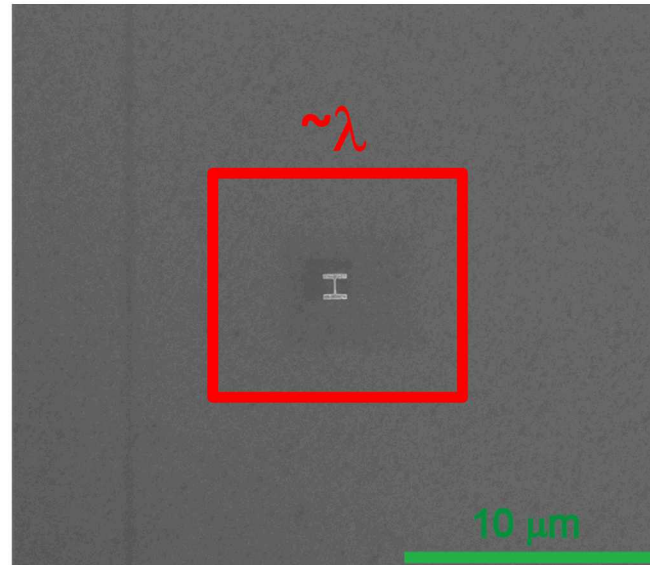
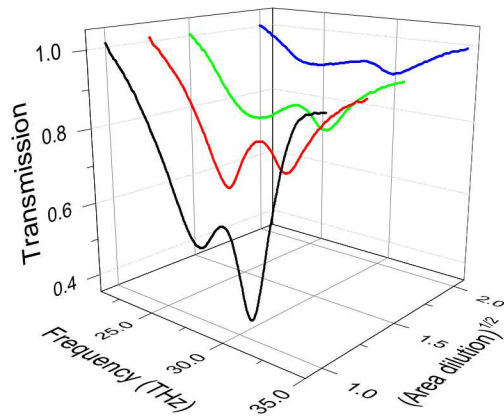
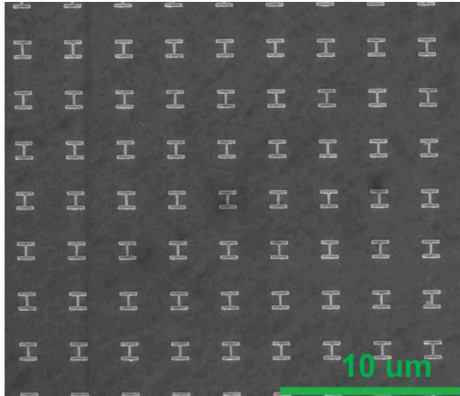


ACS Photonics (2014)

10₁₀

Nature Communications 4, (2013)

This is a “Single Resonator” Behavior

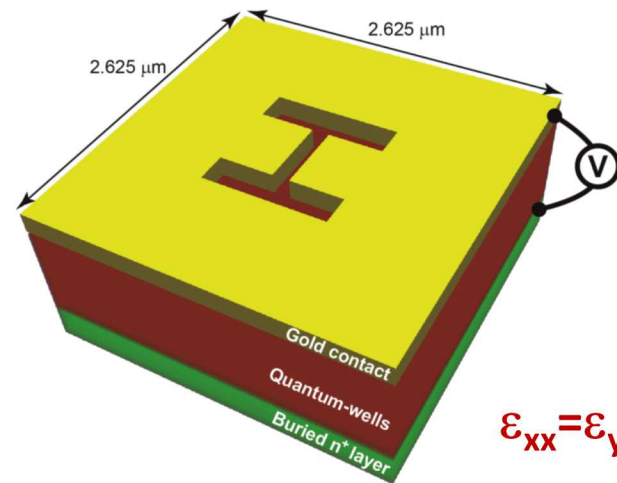


Same Rabi splitting remains down to a single resonator!

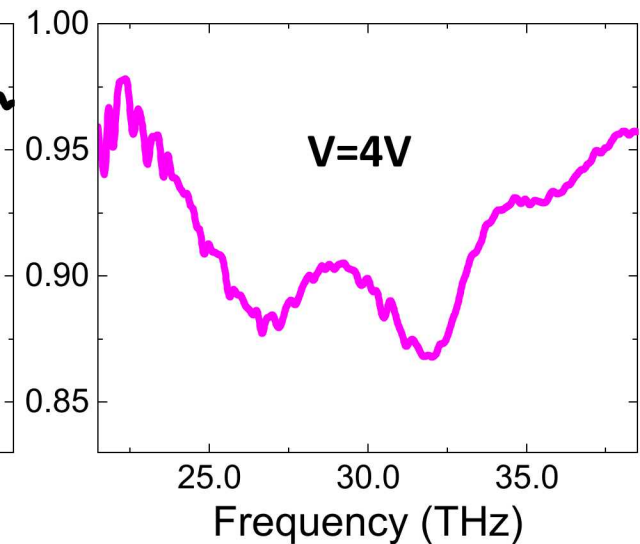
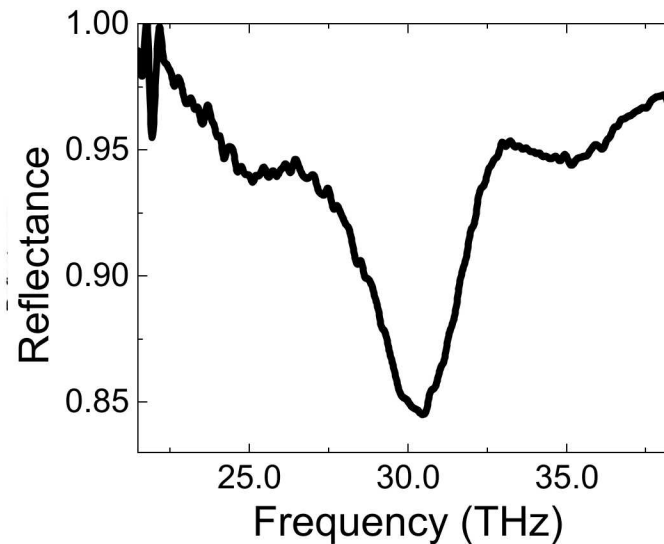
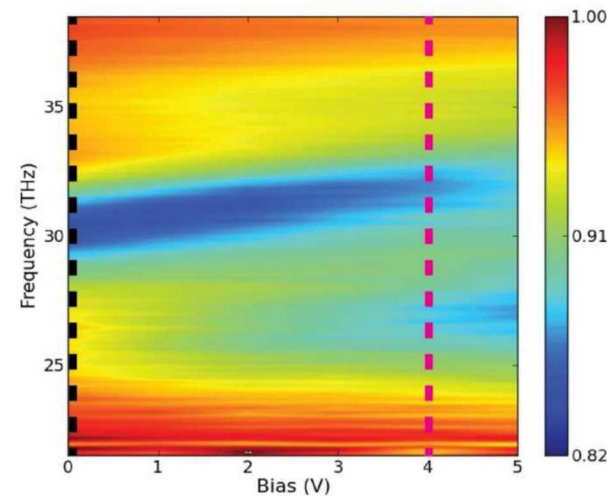
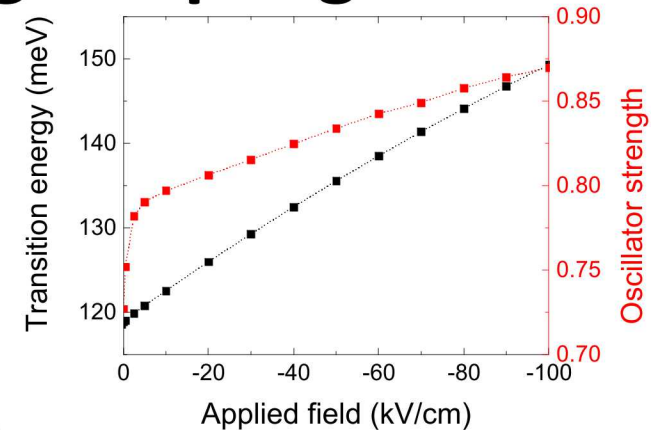
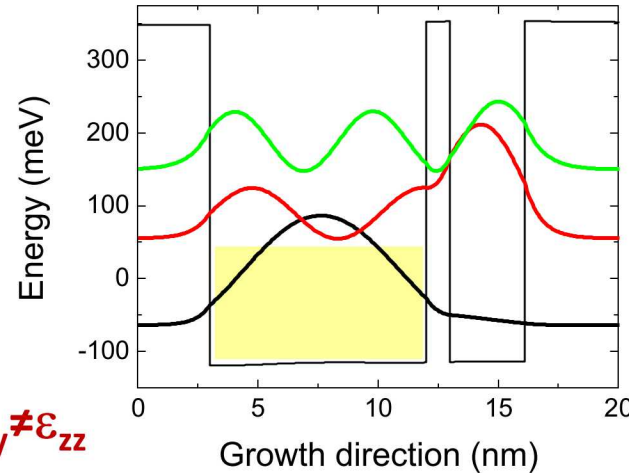
Nature communications **4**, (2013)

Area of a resonator * carrier density ~ 1000 -3000 electrons

Electrically Tunable Strong Coupling

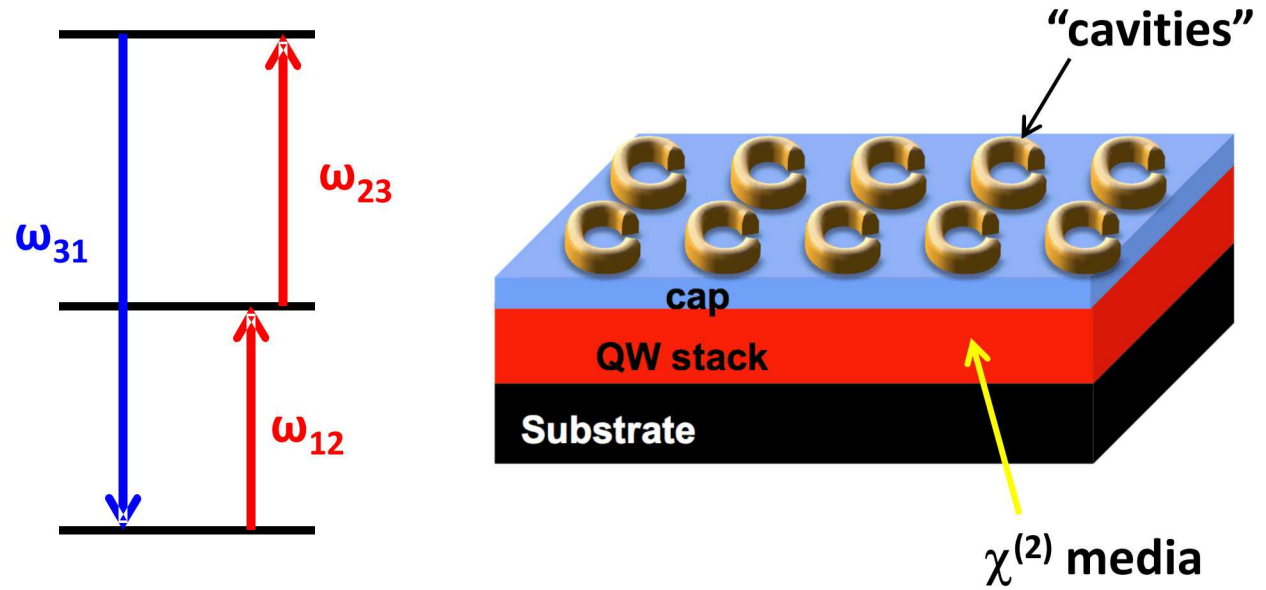


$$\epsilon_{xx} = \epsilon_{yy} \neq \epsilon_{zz}$$



Upper polariton shifted by 2.5 THz (~8% of center frequency)

Resonant $\chi^{(2)}$ in a Three-Level System



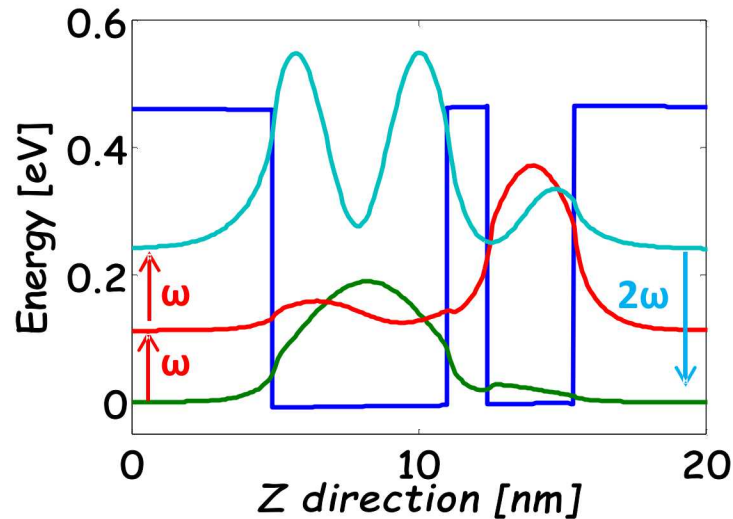
$$\chi^{(2)}(\omega) \propto \frac{N \cdot z_{12} z_{23} z_{13}}{(\omega - \omega_{12} - i\Gamma_{12})(2\omega - \omega_{13} - i\Gamma_{13})}$$

$$z_{ij} \propto \langle \Psi_i | \vec{R} | \Psi_j \rangle$$

Design of the $\chi^{(2)}$ and Resonators

Design: $10\mu\text{m} \rightarrow 5\mu\text{m}$ SHG

- InGaAs/AlInAs asymmetric system

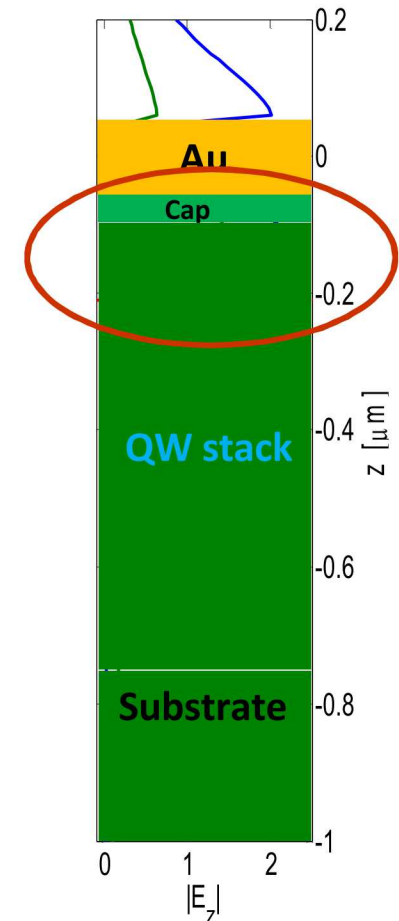
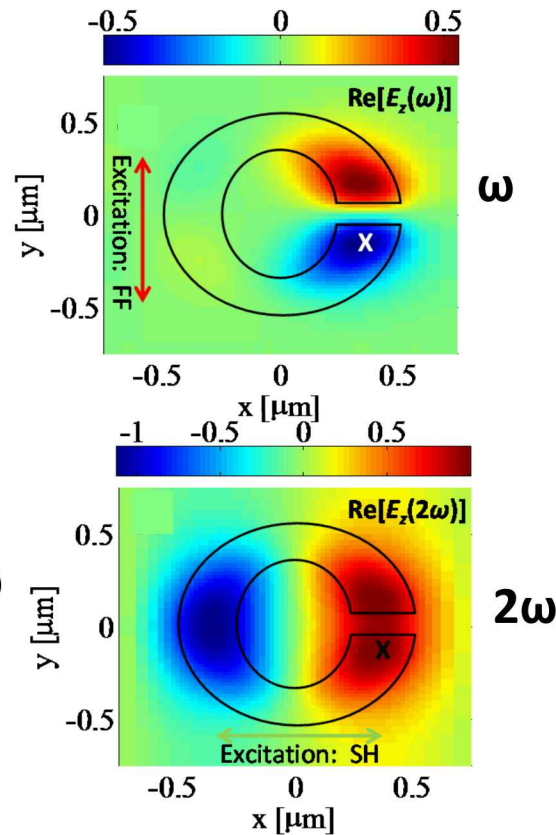


F. Capasso, C. Sirtori, and A.Y. Cho, IEEE J. Quantum Electron. 30, 1313 (1994)

$\sim 250 \text{ nm/V}$ - QWs

10s of pm/V - LiNbO_3

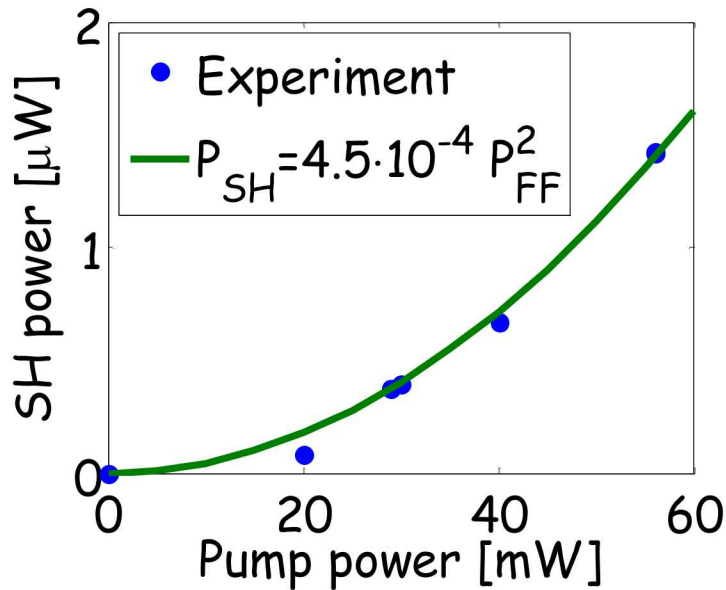
Resonators are designed to have resonances at 30 & 60 THz (5 & 10 μm)



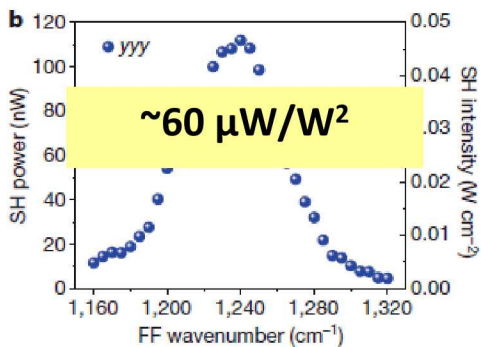
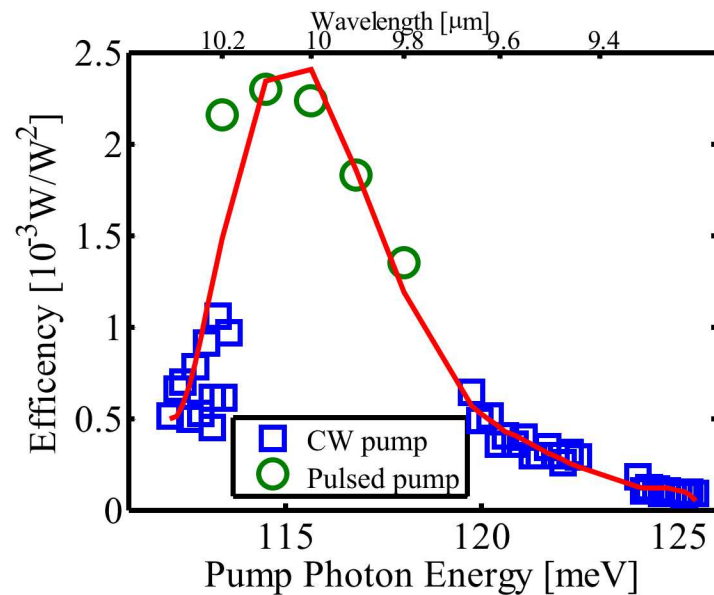
APL 104, 131104 (2014)

SHG: Power and Frequency Dependence

Pump: CW CO₂ laser



max efficiency: $\sim 2.3 \text{ mW/W}^2$

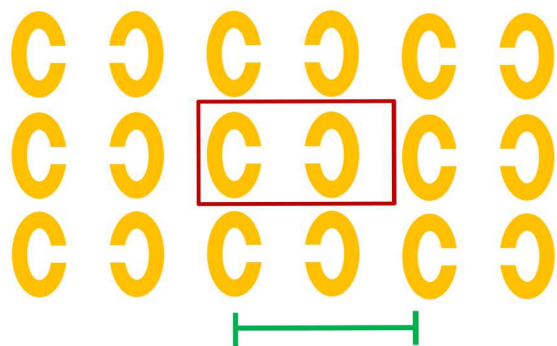


- Saturation is still an issue
- Max. conversion efficiencies of few % are possible
- Added functionality is the advantage

SHG Beam Manipulation



flipping induces π phase shift



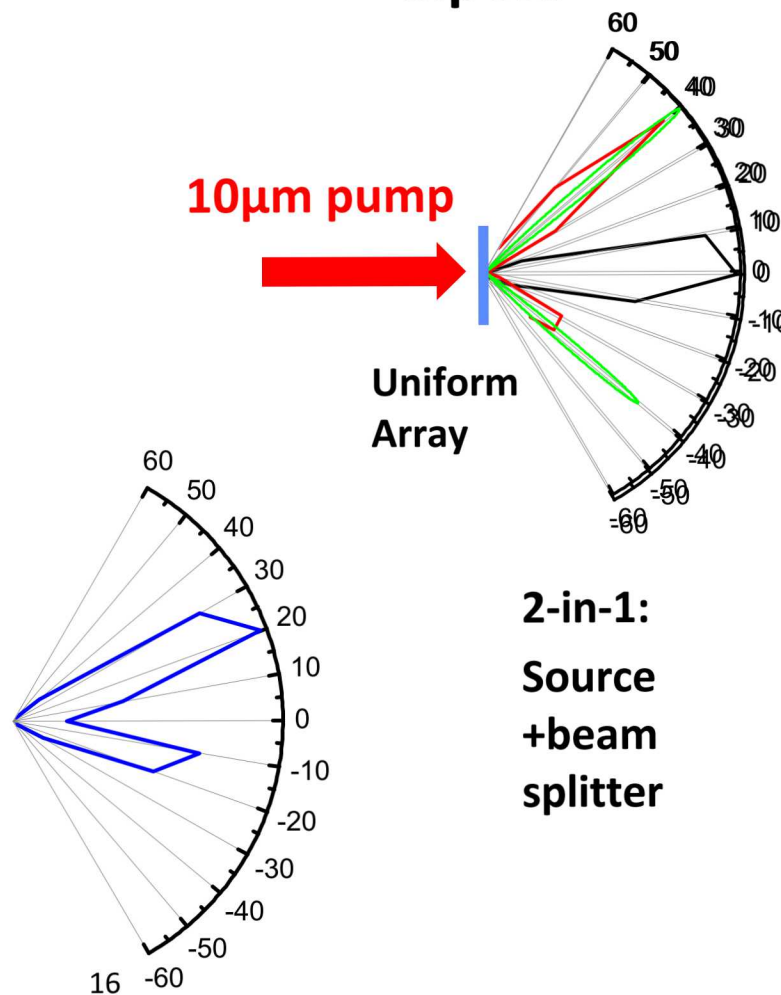
Period determines
angular separation

10 μ m pump



Uniform
Array

5 μ m
beam



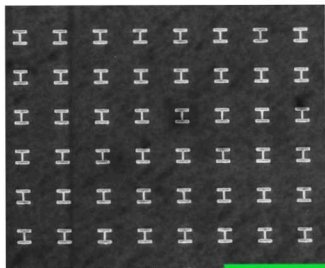
2-in-1:
Source
+beam
splitter

Metallic

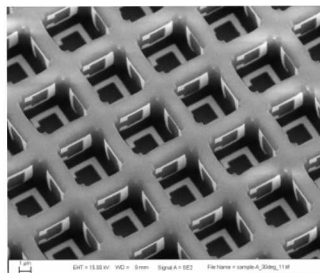
2D

3D

Planar resonators on III-V substrate



Au SRRs on resist "cavities"



>1 layer:
Incredibly difficult
Too lossy



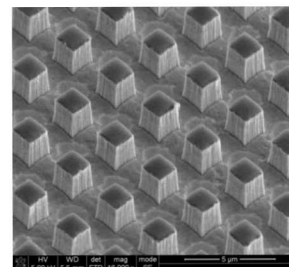
- Strong coupling to phonons, plasmons, intersubband transitions
- Voltage tuning, tunable filters
- Nonlinear optics

Dielectric

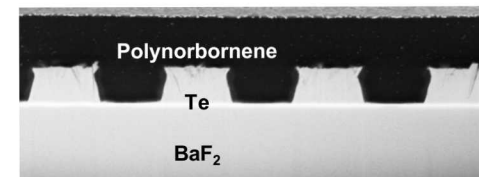
2D

3D

Tellurium cubes on BaF₂



multi-layers of Tellurium cubes?



>1 layer:
Incredibly difficult



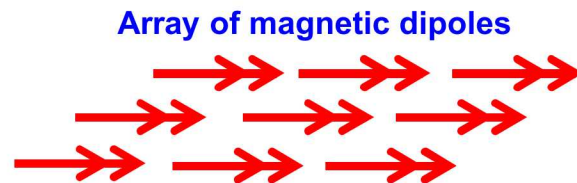
- Optical magnetism
- Unidirectional scattering, Huygens metasurfaces
- Nonlinear optics

Metasurfaces



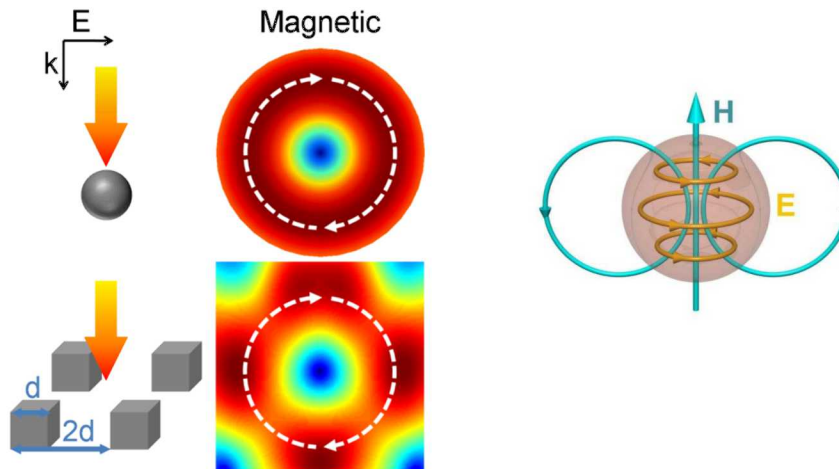
Creating a Magnetic Mirror from a Collection of Magnetic Dipoles

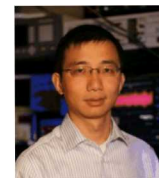
A perfect magnetic conductor does not exist in nature



Because the magnetic dipole responds in phase with the electric field, this represents an **artificial magnetic conductor**

We can create magnetic dipoles with Dielectric Resonators



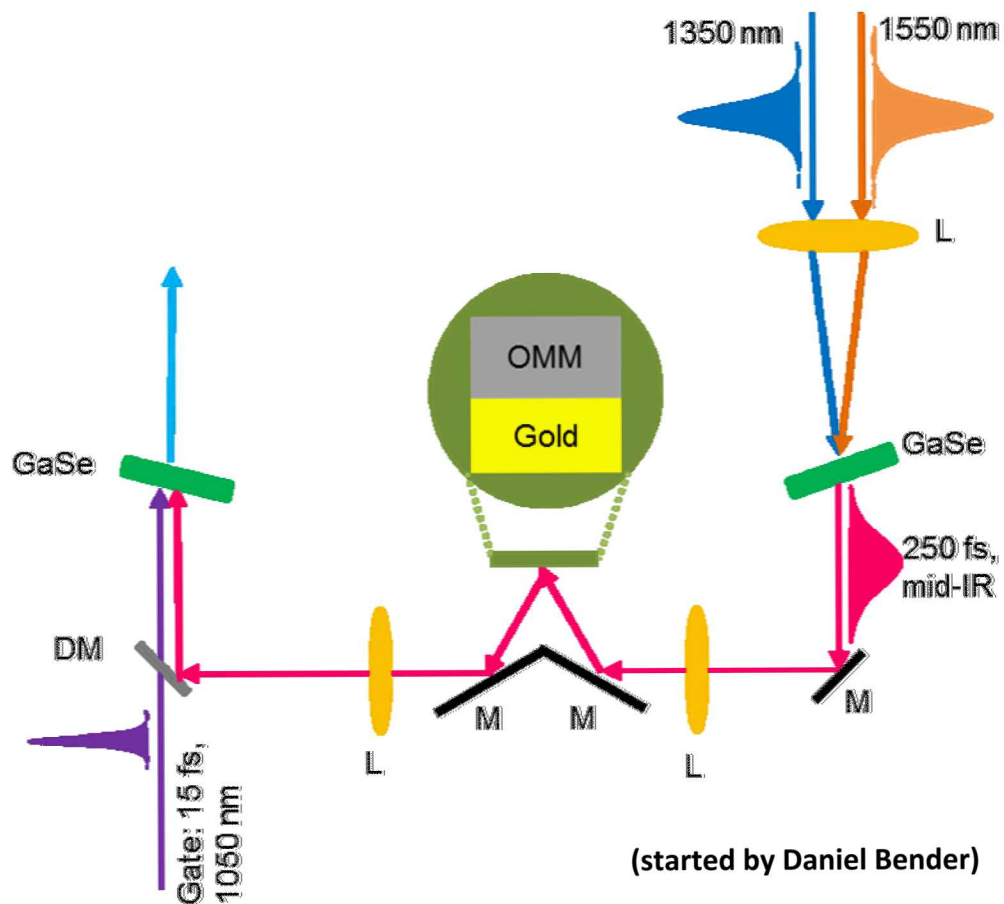
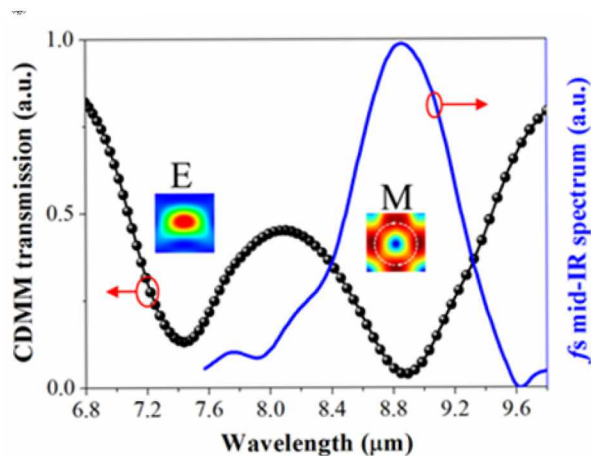
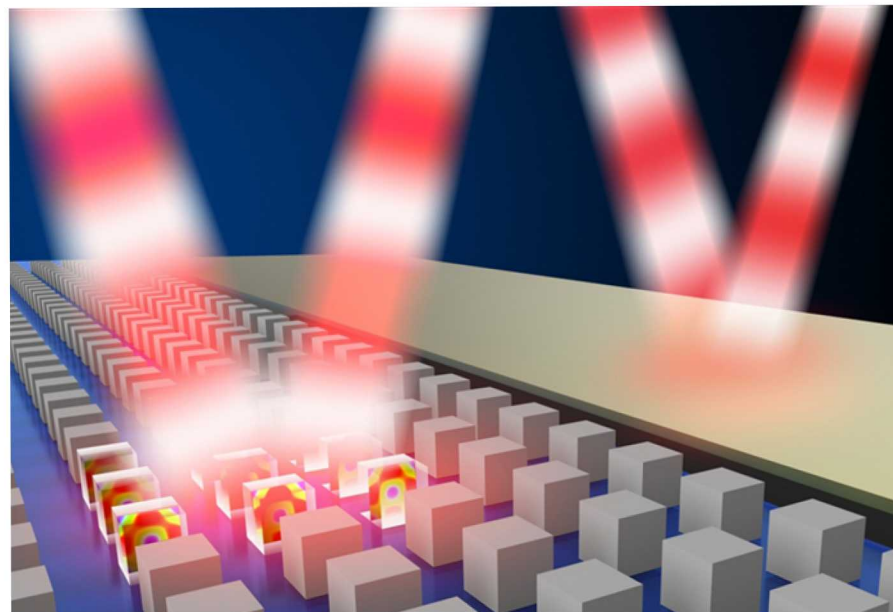


Sheng Liu

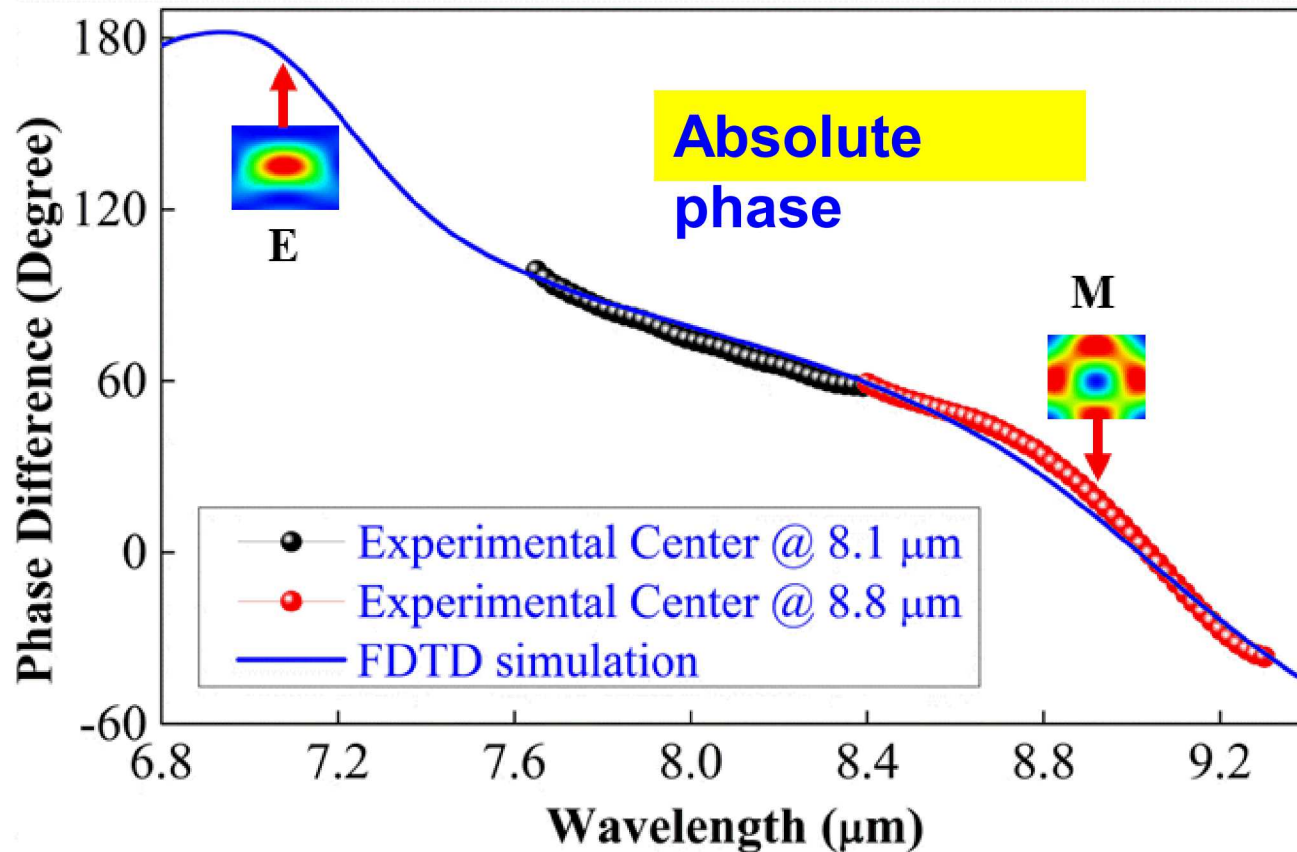
Proving Optical Magnetism: Measure Absolute Phase of Reflected Wave

a

Phase-locked Time Domain Spectroscopy



Fourier Transform of Experimental Results (Phase)



- Phase difference between magnetic (~170 degrees) and electric (10 degree) resonance is ~160 degrees (close to π)

Shorter Wavelengths: Limitations of Si for Dielectric Metamaterials

- Challenging to realize real **3D** metamaterials (multi-layers)—different etchant for Si and SiO₂

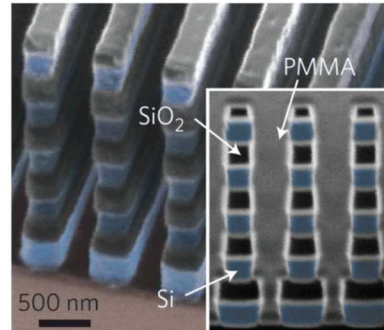
nature
photonics

LETTERS

PUBLISHED ONLINE: 25 AUGUST 2013 | DOI: 10.1038/NPHOTON.2013.214

Realization of an all-dielectric zero-index optical metamaterial

Parikshit Moitra¹, Yuanmu Yang¹, Zachary Anderson², Ivan I. Kravchenko³, Dayl P. Briggs³ and Jason Valentine^{4*}

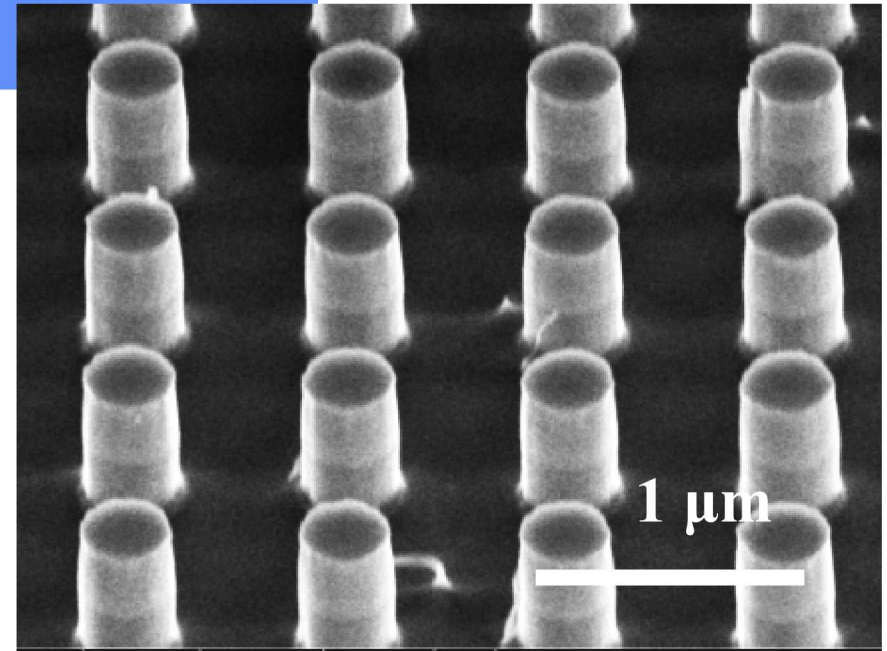
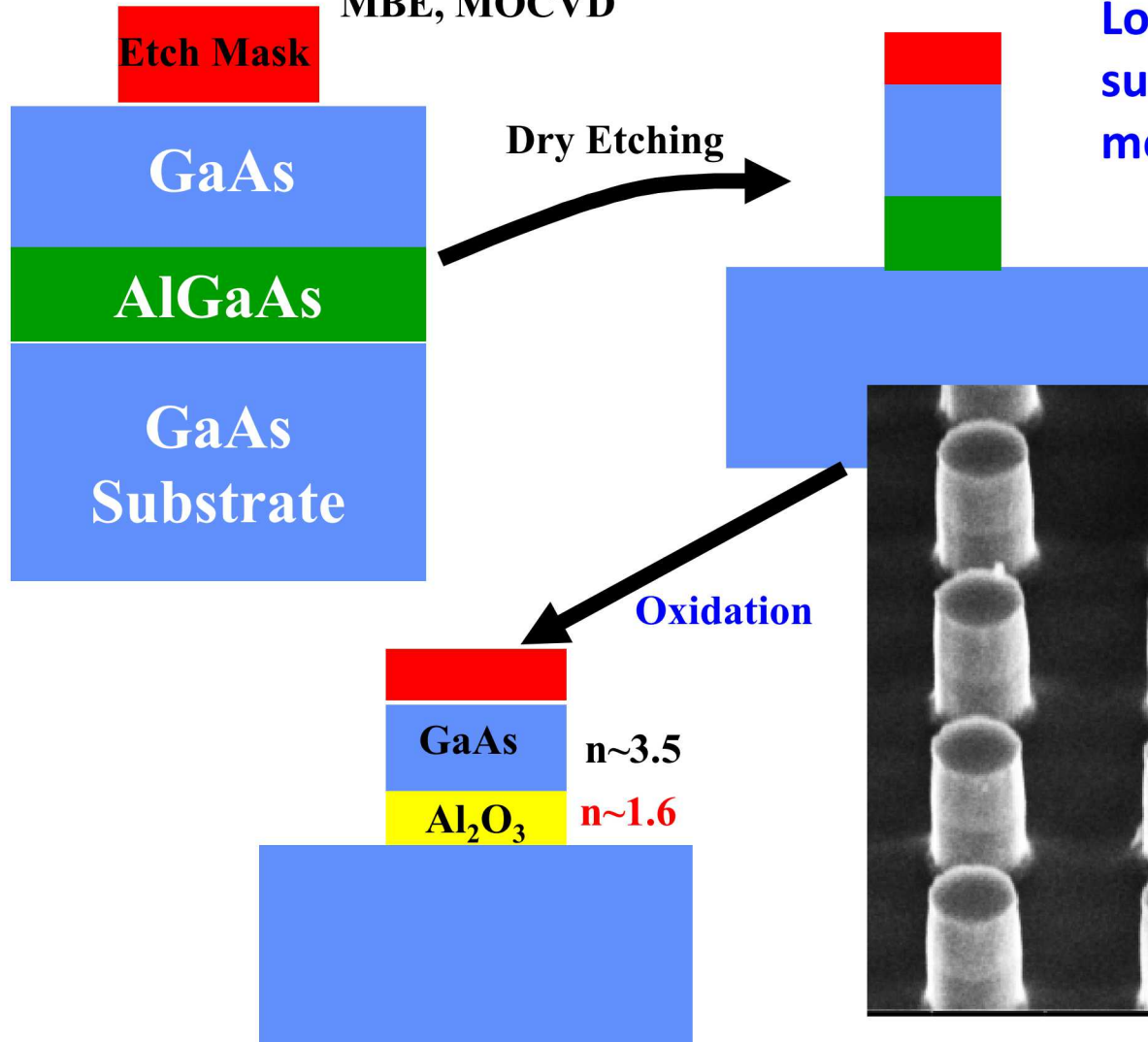


- Free carrier absorption of Silicon (long electron lifetime)
- Silicon is centrosymmetric → NO second order nonlinearity
- Inefficient light emission (hard to incorporate active media)

**Direct bandgap semiconductors: III-V
(AlGaAs, InGaAs, GaAs)**

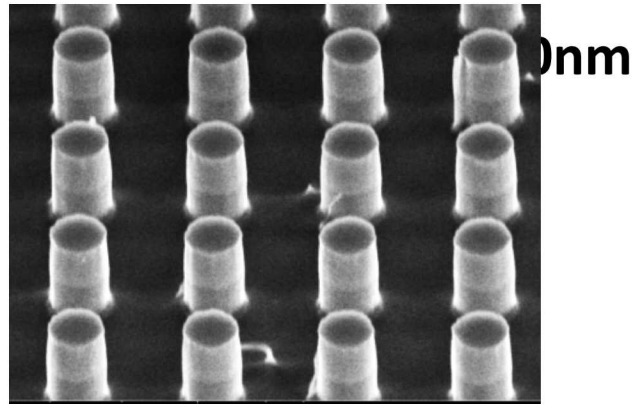
Fabrication of Al(In)GaAs Based Dielectric MM

Epitaxially grown:
MBE, MOCVD

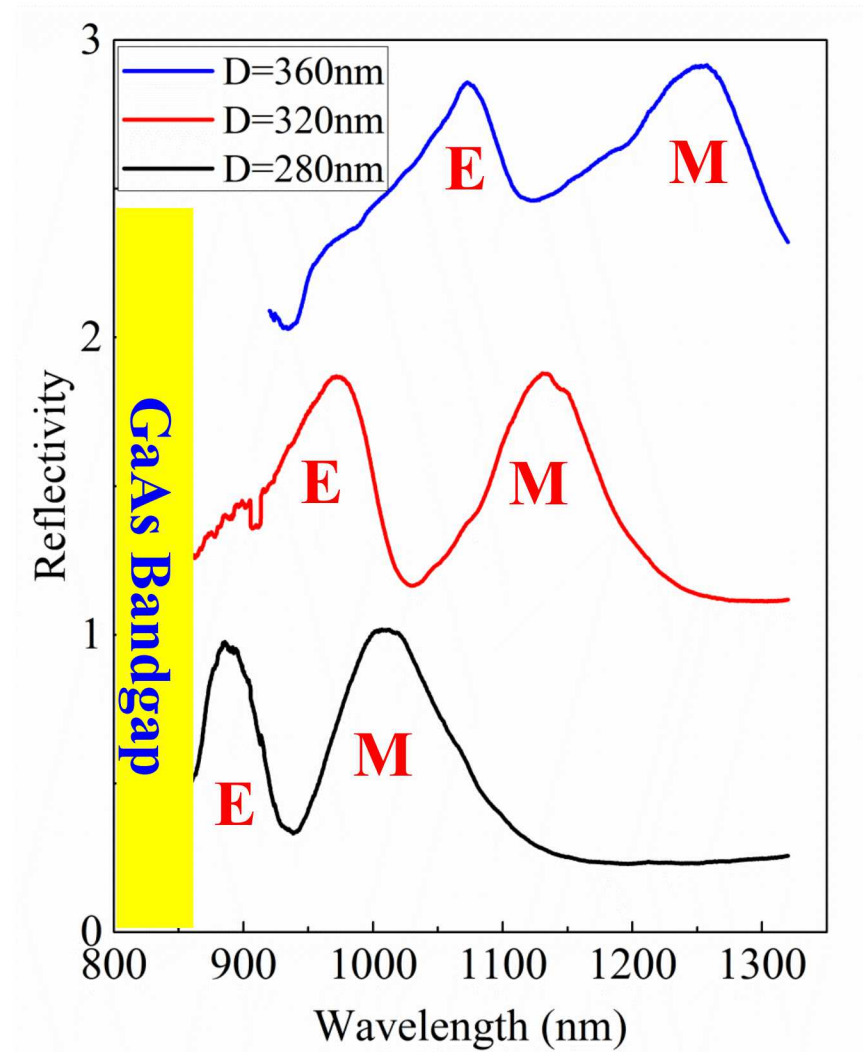


Sandia VCSEL: IEEE JSTQE, 3, 916 (1997)

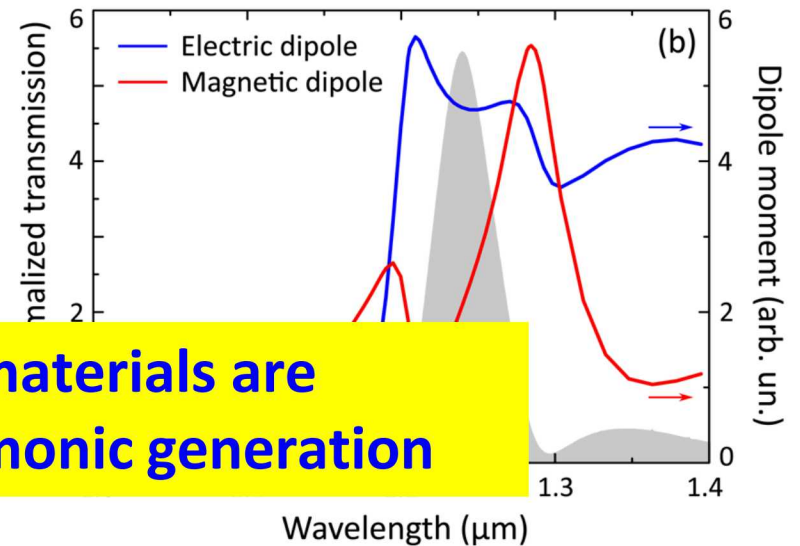
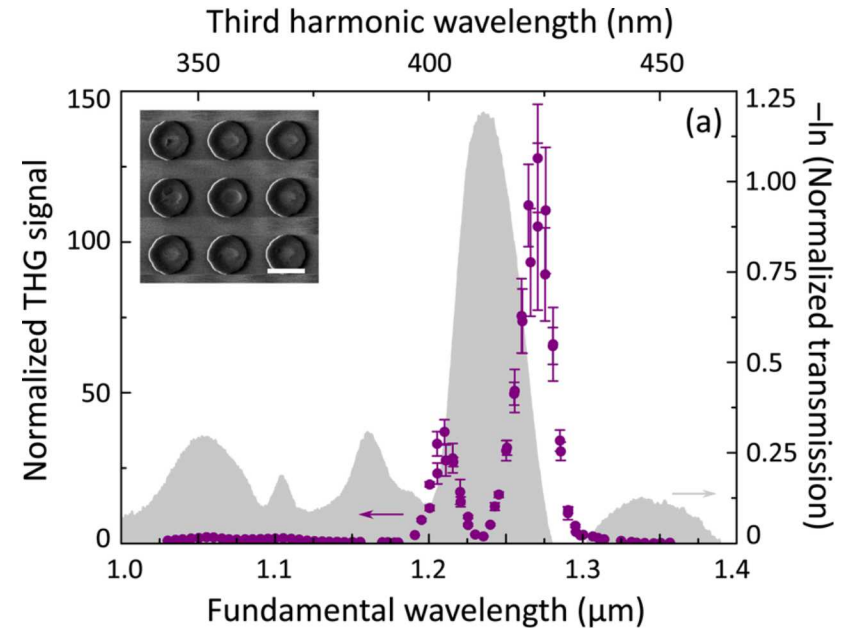
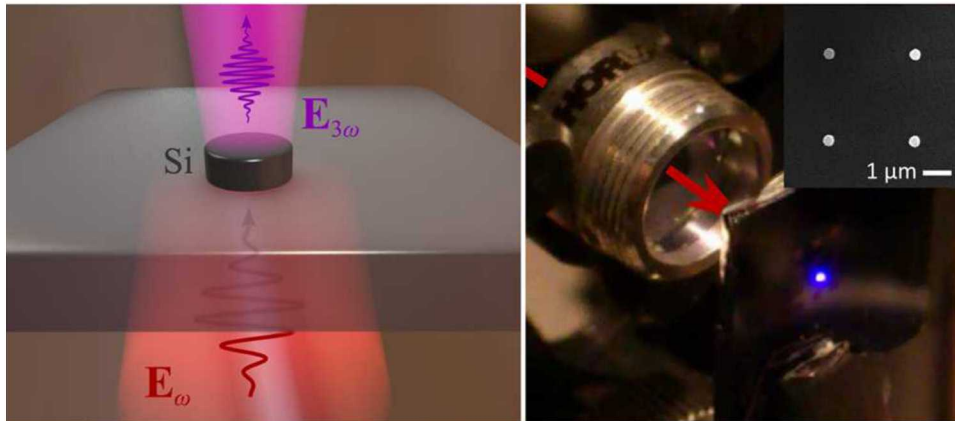
GaAs Dielectric Resonators (1 layer)



- Extremely low loss below bandgap
- Crystalline



Optical Nonlinearities in Dielectric Metasurfaces



(With ANU & Moscow State)

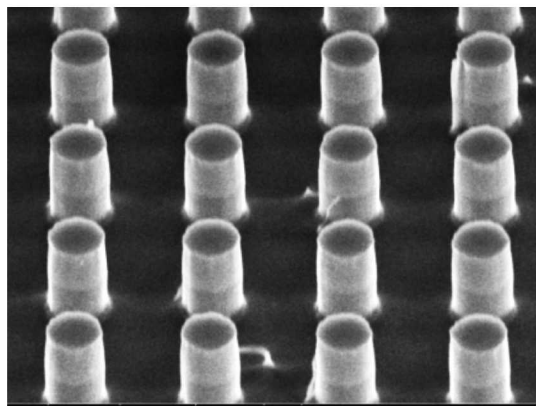
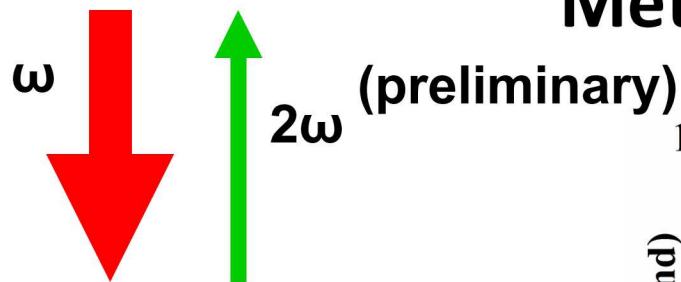
Non-centrosymmetric materials are needed for second harmonic generation

Nano Letters (2015)

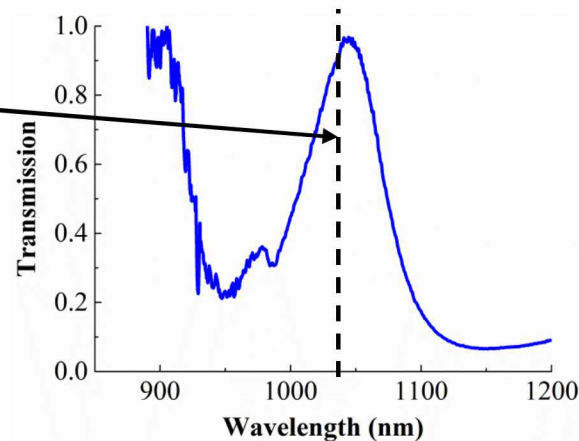
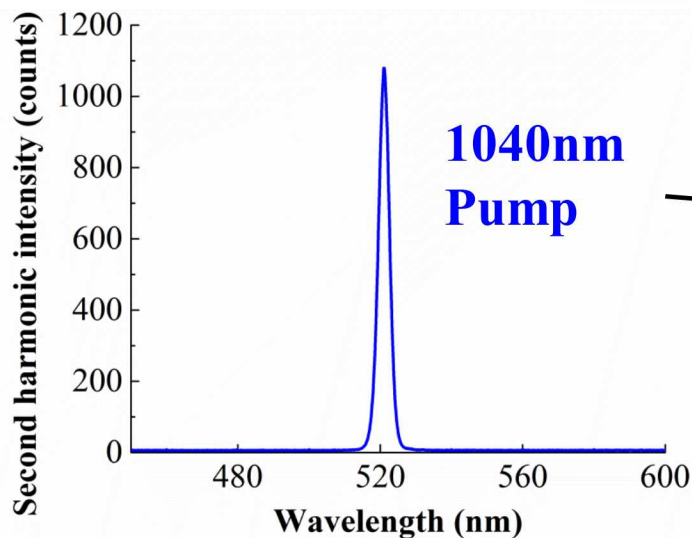
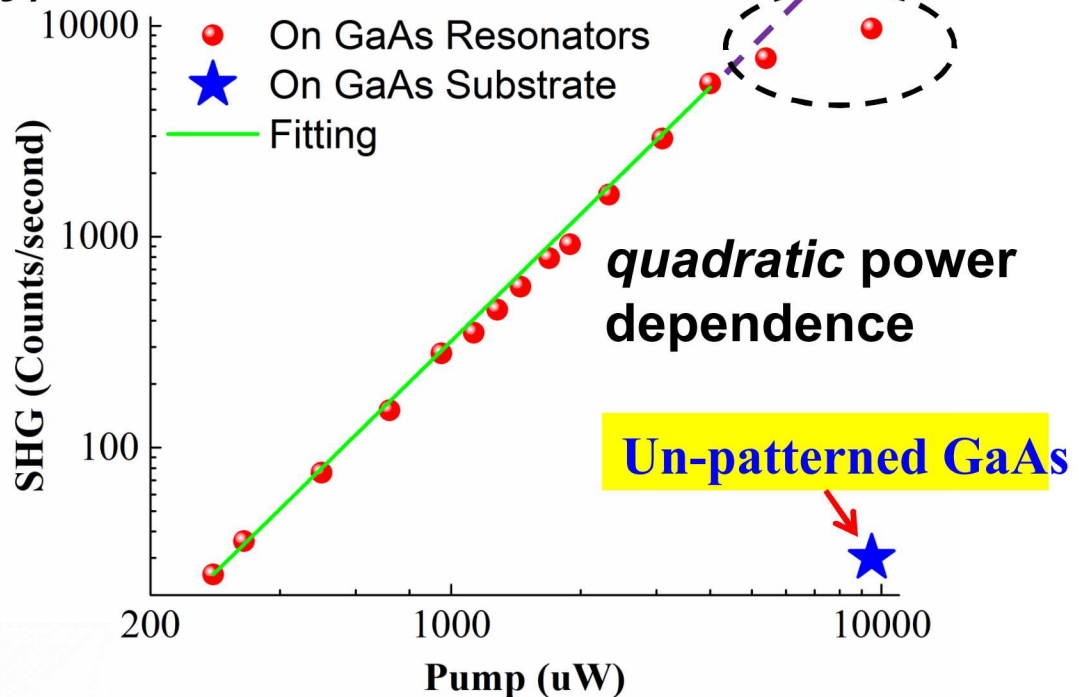


Nonlinear Harmonic Generation using III-V

Metasurfaces

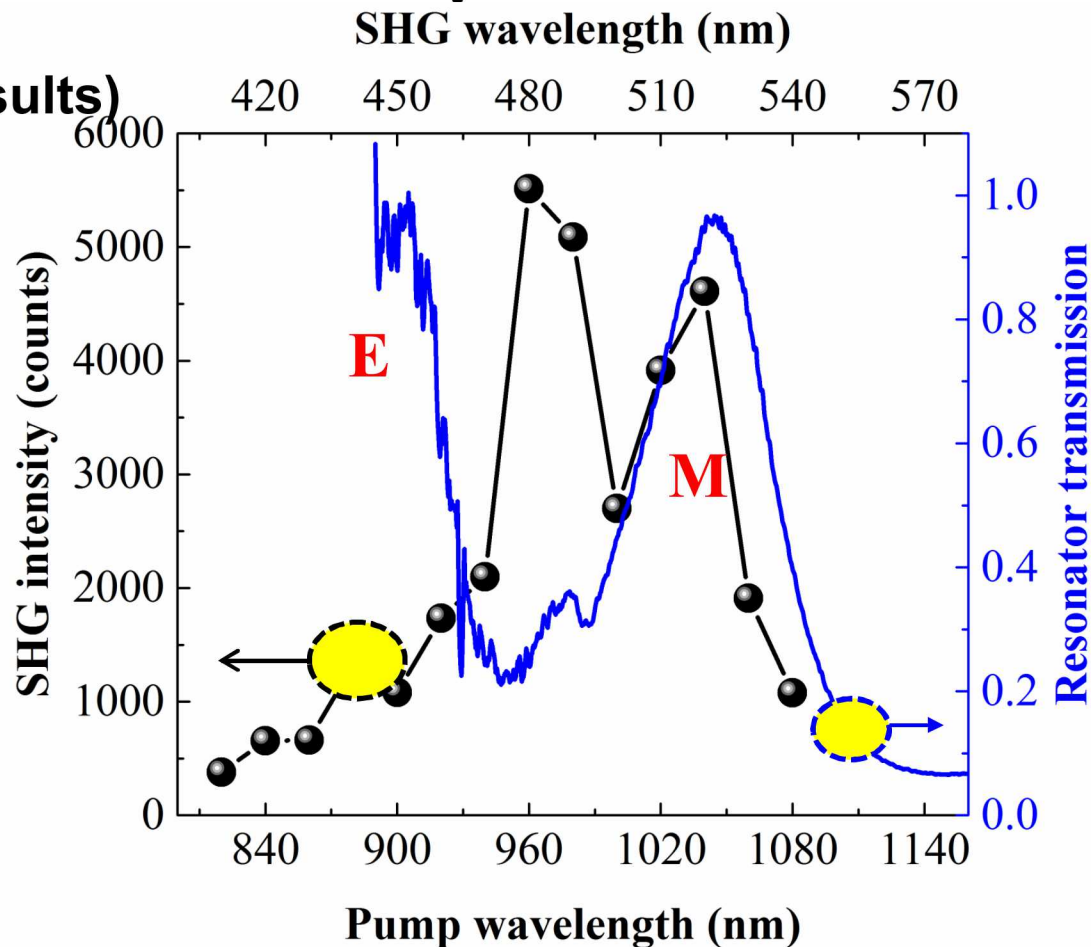


Pump: 1040nm



Second Harmonic– Pump Wavelength Dependence

(preliminary results)



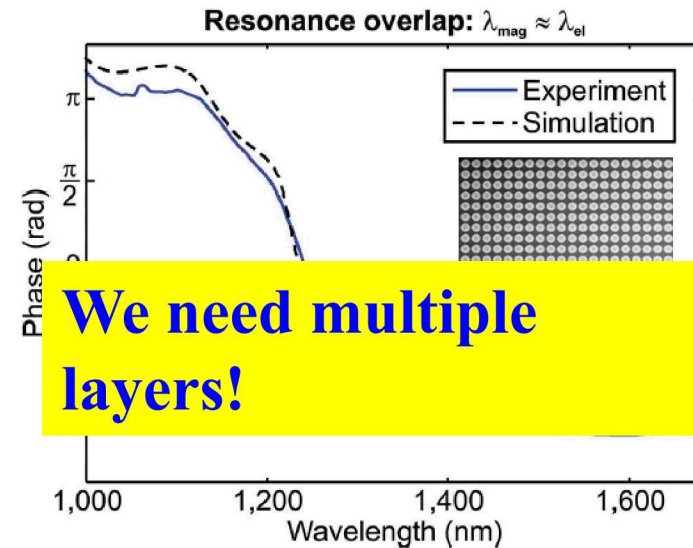
- Strong pump wavelength dependence—resonant enhancement
- @ same pump wavelength—different diameter resonators show different SHG efficiency
- More work is under way

Multilayer Dielectric Metamaterials:

Example, Femtosecond Pulse Compression

Huygens sources radiate unidirectional as superposition of electric and magnetic dipoles

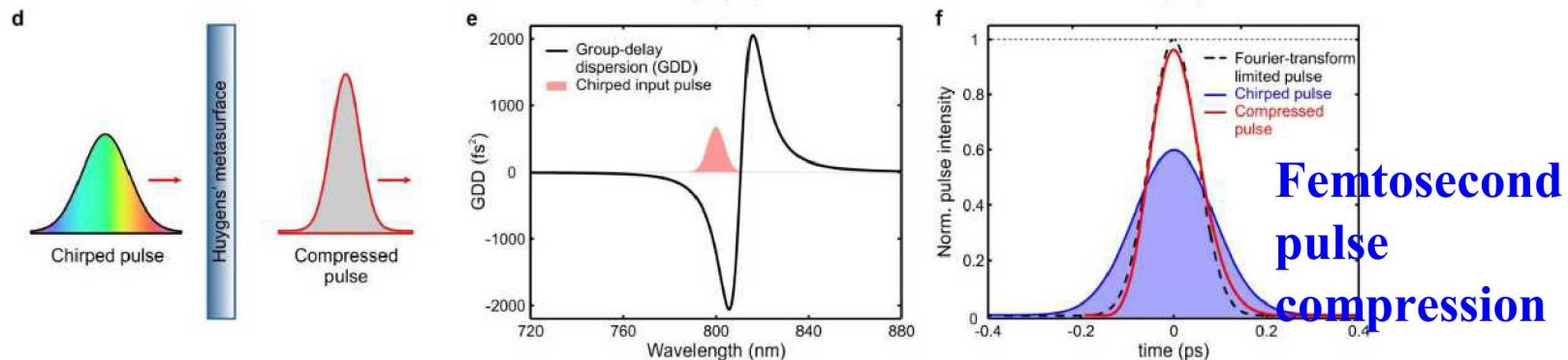
Single layer is not enough to generate sufficient phase difference within the spectrum of femtosecond pulses.



We need multiple layers!

Huygens Metasurface

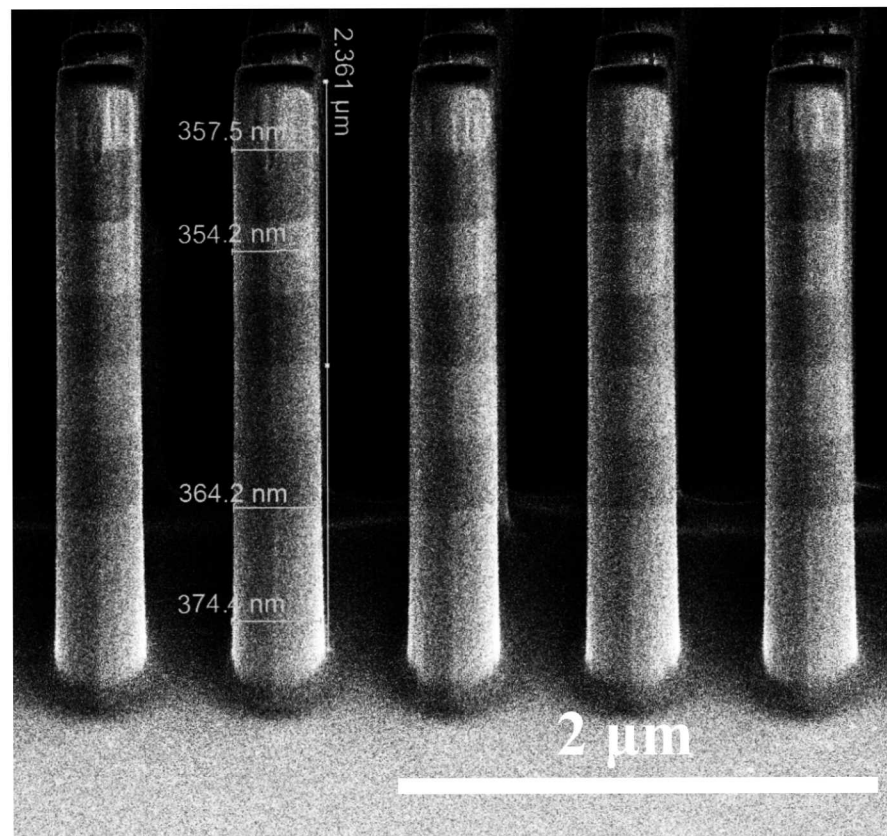
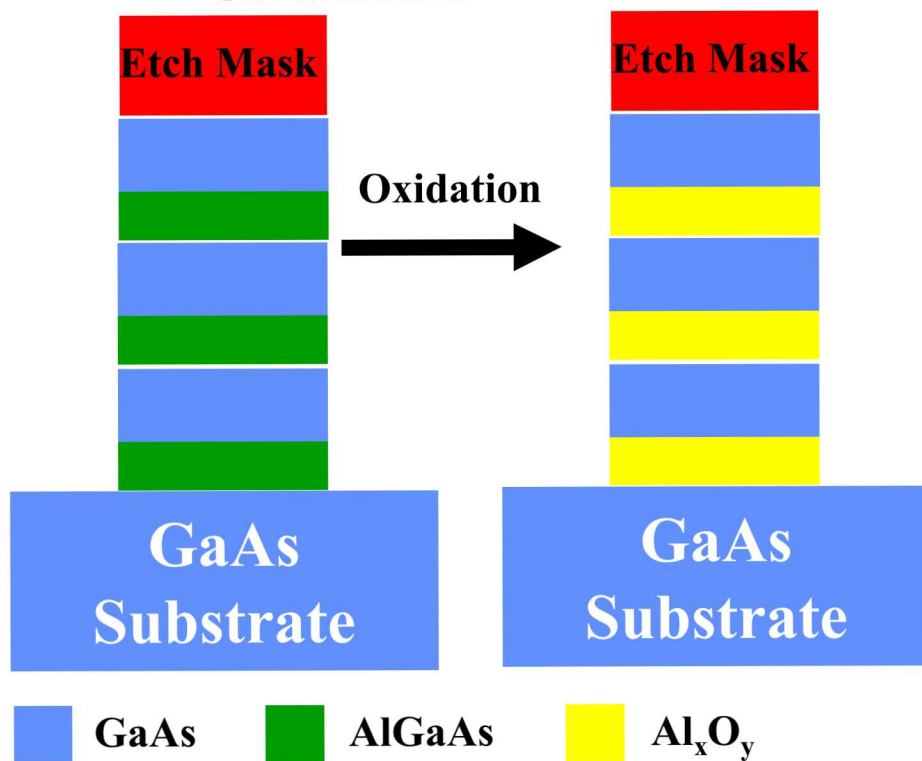
Complete 2π phase range in transmission



“High-Efficiency Dielectric Huygens’ Surfaces”, Advanced optical materials, 3, 813 (2015)

Multilayer III-V Dielectric Metamaterials

Epitaxially grown multi-layer
GaAs/AlGaAs





Predictions.... Thoughts...

- **Voltage tunable metasurfaces will be incorporated with III-V devices and will enable next generation of focal plane arrays**
- **Work in nonlinear metasurfaces will increase: opportunities in quantum information science, etc.**
- **Advances in low-loss dielectric metasurfaces could enable practical optical devices in most of the IR (phase plates, holographic elements, etc)**
- **New types of III-V devices that combine dielectric metasurfaces and III-V optoelectronics?**



Acknowledgments

Current Postdocs: Salvo Campione, Omri Wolf, Sheng Liu

Departed Postdocs: Alex Benz, Young-Chul Jun, James Ginn, Alon Gabbay, Brandon Passmore, Daniel Bender

CINT + Sandia: John Reno, Mike Sinclair, Ines Montano, John Klem, Eric Shaner, Mike Wanke, Joel Wendt, Paul Clem, Jon Ihlefeld, Bruce Burckel, Jason Domingues, Andy Allerman, Jon Wierer,

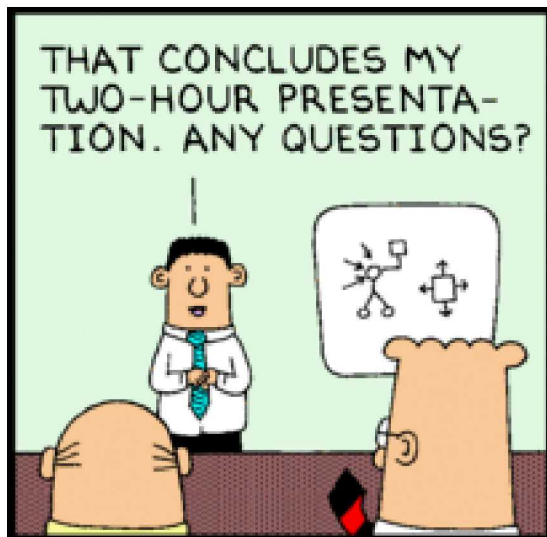
ANU: Y. Kivshar, I. Staude, M. Decker, D. Neshev,

Lomonosov Moscow State: M. Shcherbakov,...



U.S. DEPARTMENT OF
ENERGY

Office of
Science



www.dilbert.com scottadams@aol.com



8/4/03 © 2003 United Feature Syndicate, Inc.

