

IONIZATION WAVE DYNAMICS OF A PLASMA JET IN CONTACT WITH A DIELECTRIC SURFACE

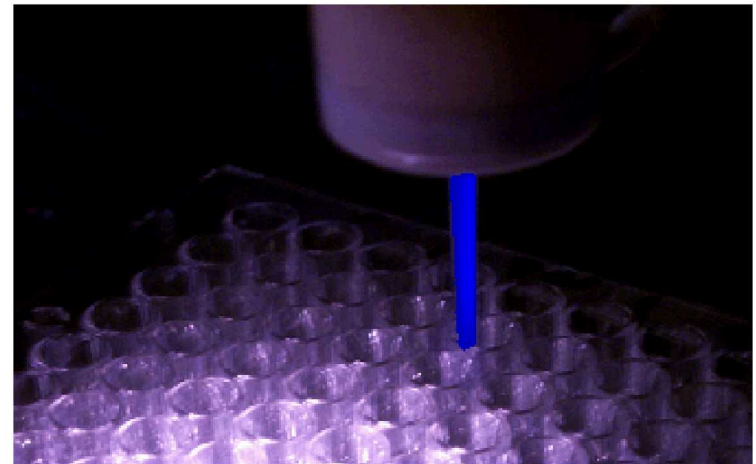
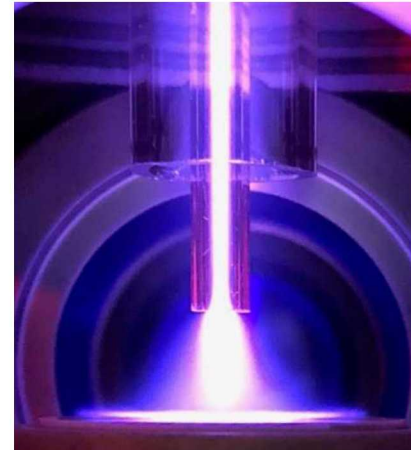
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DOE Center on Control of Plasma Kinetics
Annual Meeting May 2018

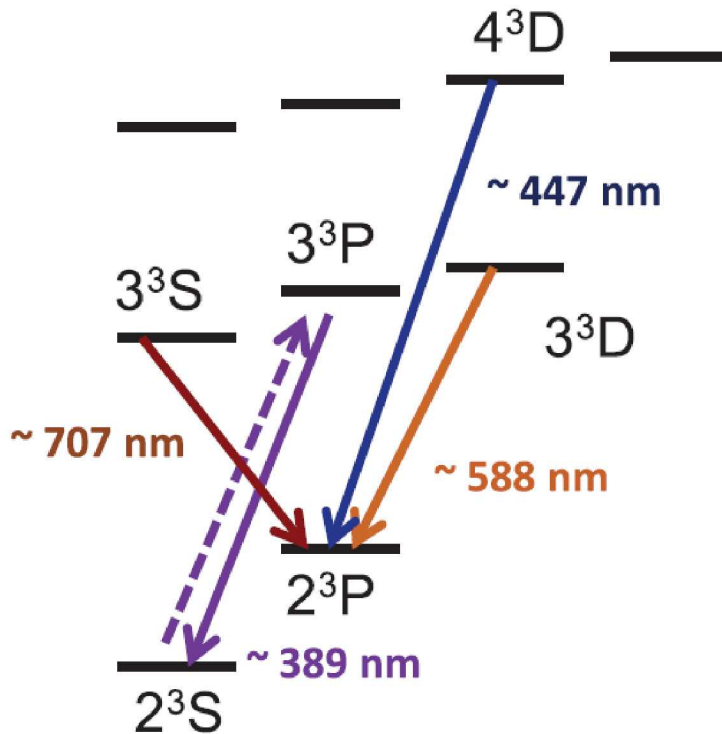
ATMOSPHERIC PRESSURE PLASMA JETS

- Atmospheric pressure plasma jets (APPJs) are a popular source of chemistry for biomedical applications.
- The plasma propagates as an ionization wave (IW) that is repetitively pulsed.
- The IW gives rise to reactive oxygen and nitrogen species (RONS) which produce the biological effect.



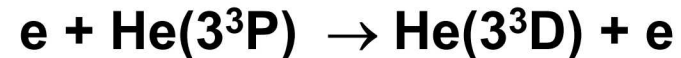
S. Mohades, et al., *Physics of Plasmas* 22, 122001 (2015).

LASER COLLISIONAL INDUCED FLUORESCENCE (LCIF)



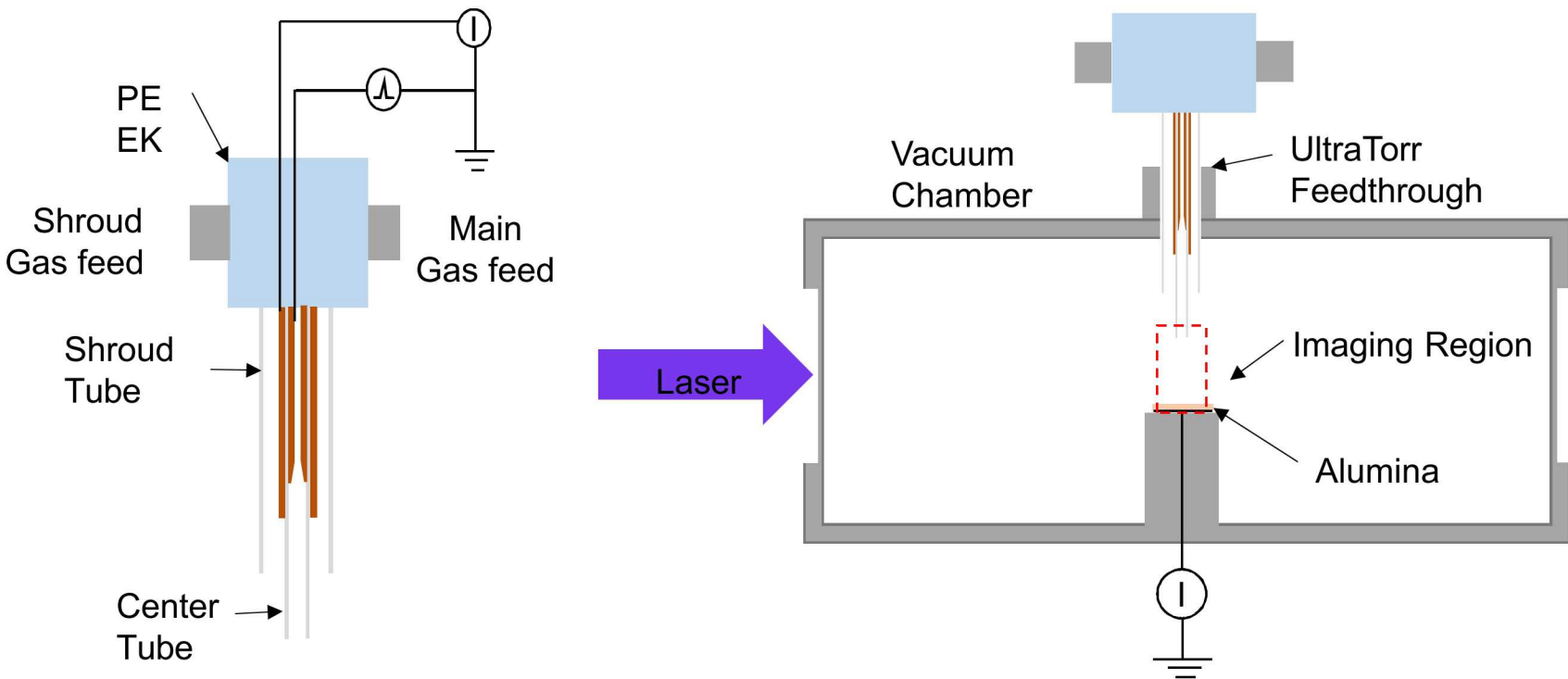
Barnat and Fierro, J. Phys. D: Appl. Phys., 50, 14LT01 (2017).

- An ultrafast laser (<100 fs) is used to measure the electron density with a high time resolution.
- Electrons collide with laser excited He(3^3P)



- $n_e \sim 588 \text{ nm} / 389 \text{ nm}$
- Sufficient He(2^3S) density is critical for accurate LCIF data.
- LCIF was developed for a pure He environment, but extending this diagnostic to mixtures is important for APPJs.

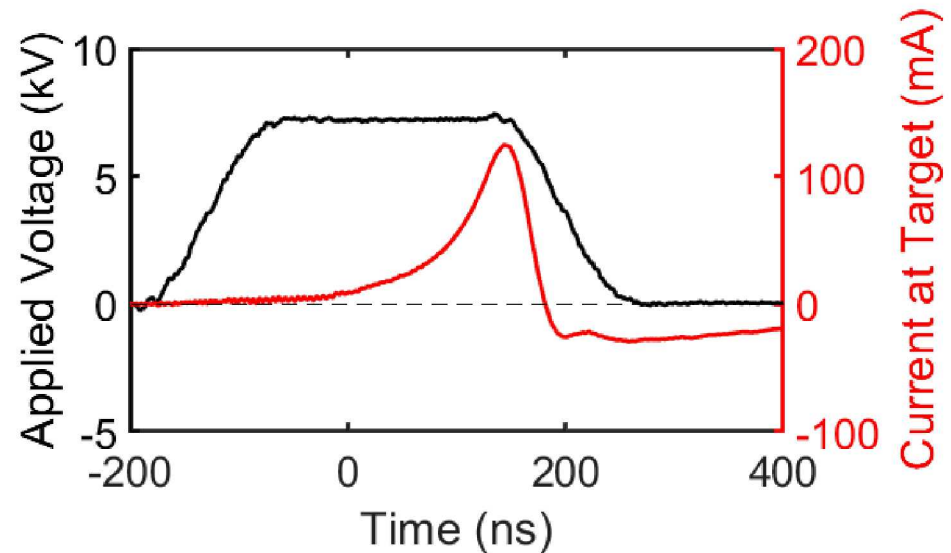
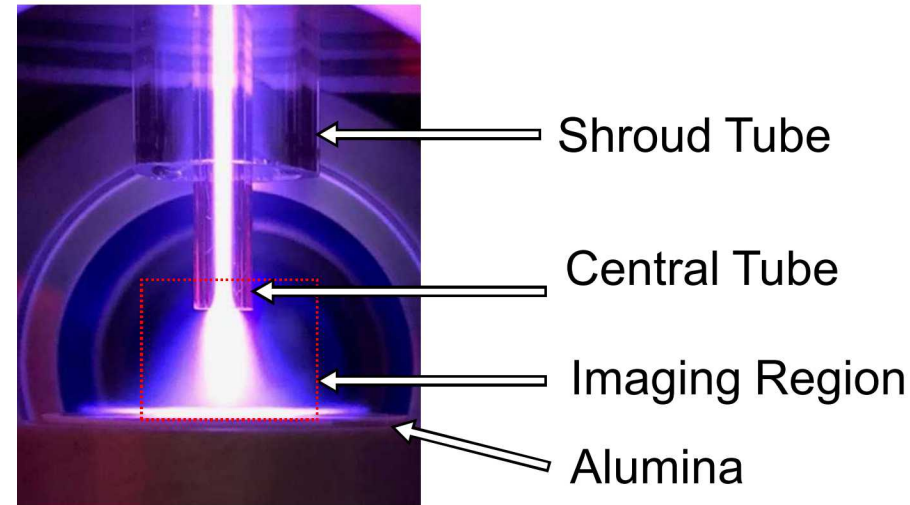
EXPERIMENTAL SETUP



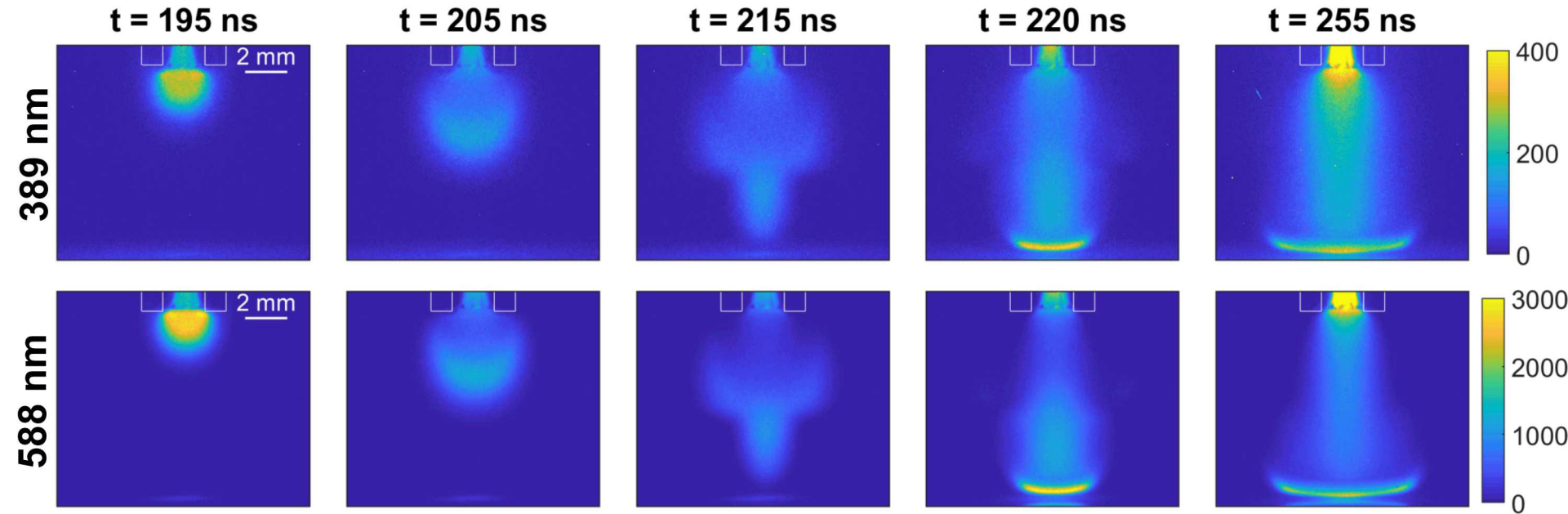
- **Annular powered electrode within the center tube**
- **Placing the APPJ in a vacuum chamber - consistent and controlled chemistry, ground planes, and gas flow.**
- **Coaxial tube allows a gas shroud – control environment independently of gas in main jet.**

BASE CASE

- +6 kV
- 430 ns pulse, 100 ns rise
- 200 Torr
 - Faster dynamics (for modeling)
 - Lower background LCIF signal (for experiment)
- 500 sccm He in center tube
- Base pressure 20 mTorr
- Gap to target = 7.5 mm
- 65 μm thick alumina disk
- Current measured at ground electrode under the alumina.

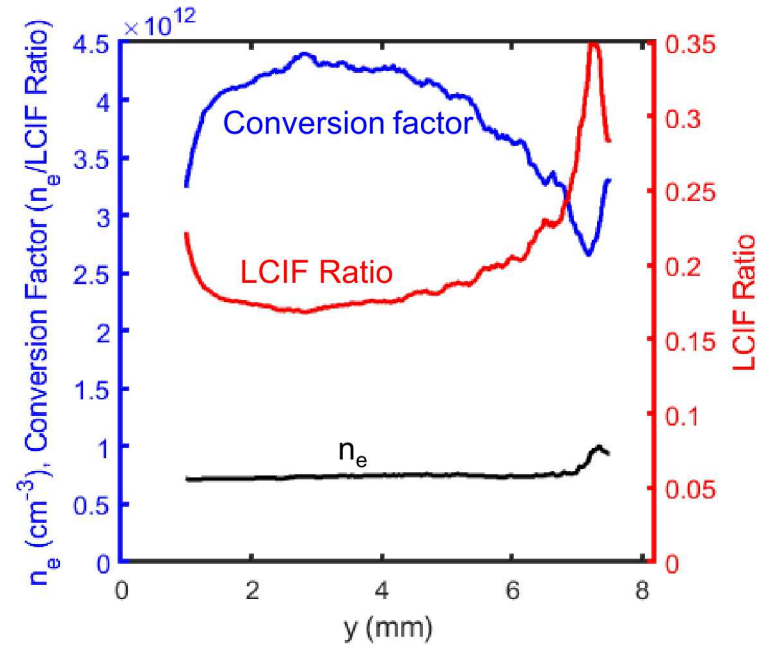
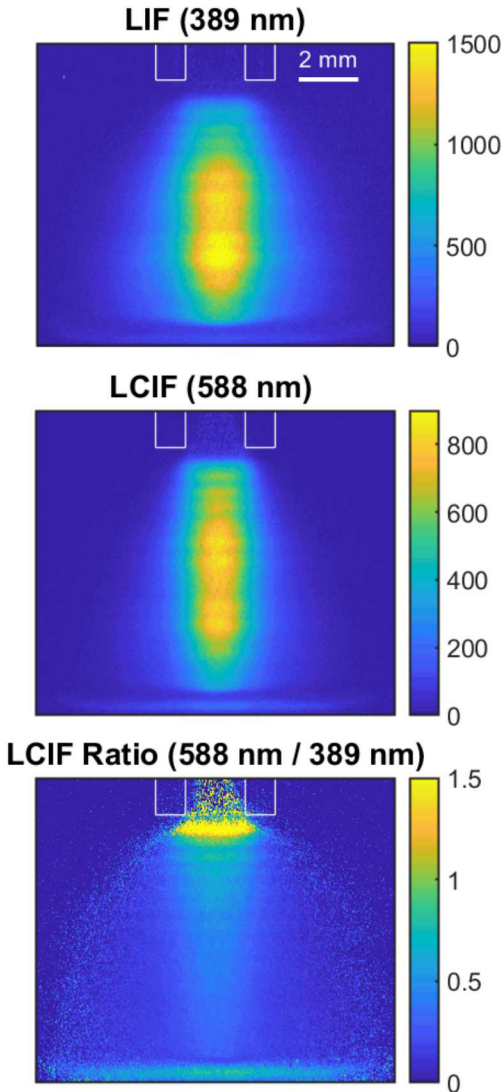


BASE CASE – PLASMA EMISSION



- 5 ns ICCD gate
- Ionization wave expands as it exits the tube.
- Approaching the alumina, it becomes more directed.
- Upon contacting the alumina, forms a surface IW (SIW) and spreads, charging the surface.
- As voltage falls, there a restrike occurs.

CONVERT LCIF TO n_e



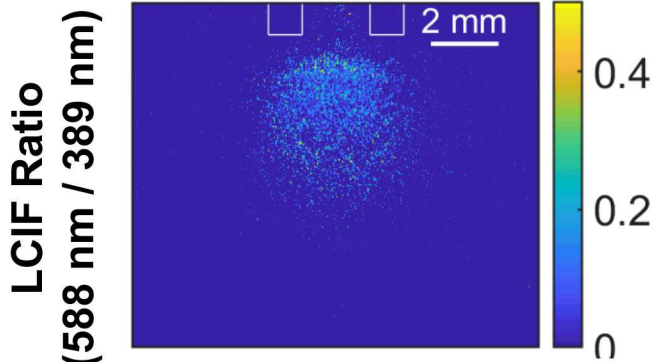
- Calculate n_e using Ohm's law and measured E/N.

$$I = \frac{e^2 N E}{m_e v_m N} n_e A$$

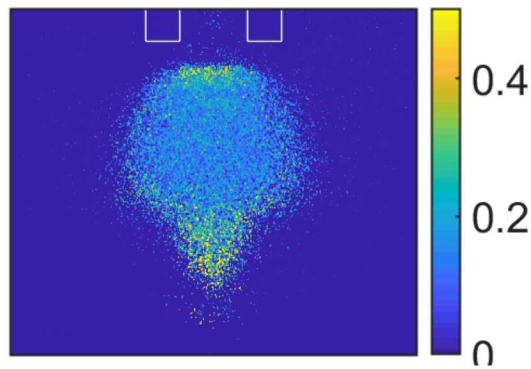
- An LCIF ratio of 1, is approximately $4 \times 10^{12} \text{ cm}^{-3}$ electrons.
- Previously, conversion factor estimated at $1.5 \times 10^{13} \text{ cm}^{-3}$.

BASE CASE LCIF

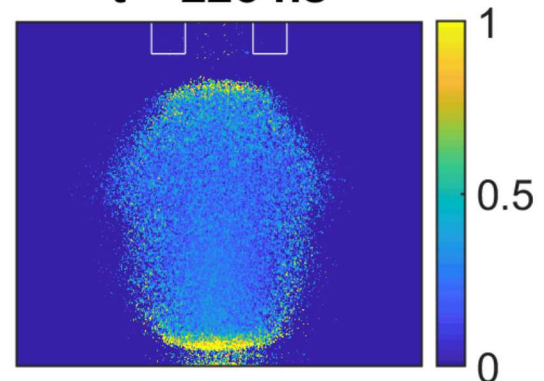
t = 200 ns



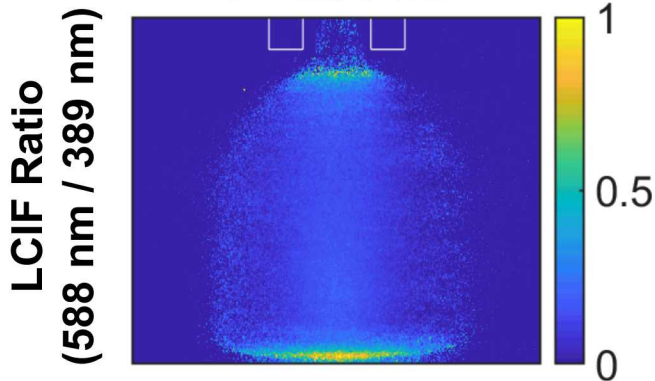
t = 210 ns



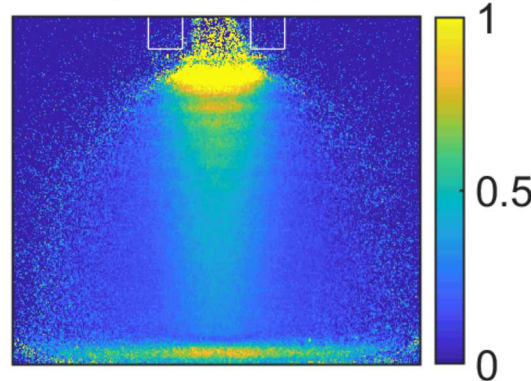
t = 225 ns



t = 250 ns



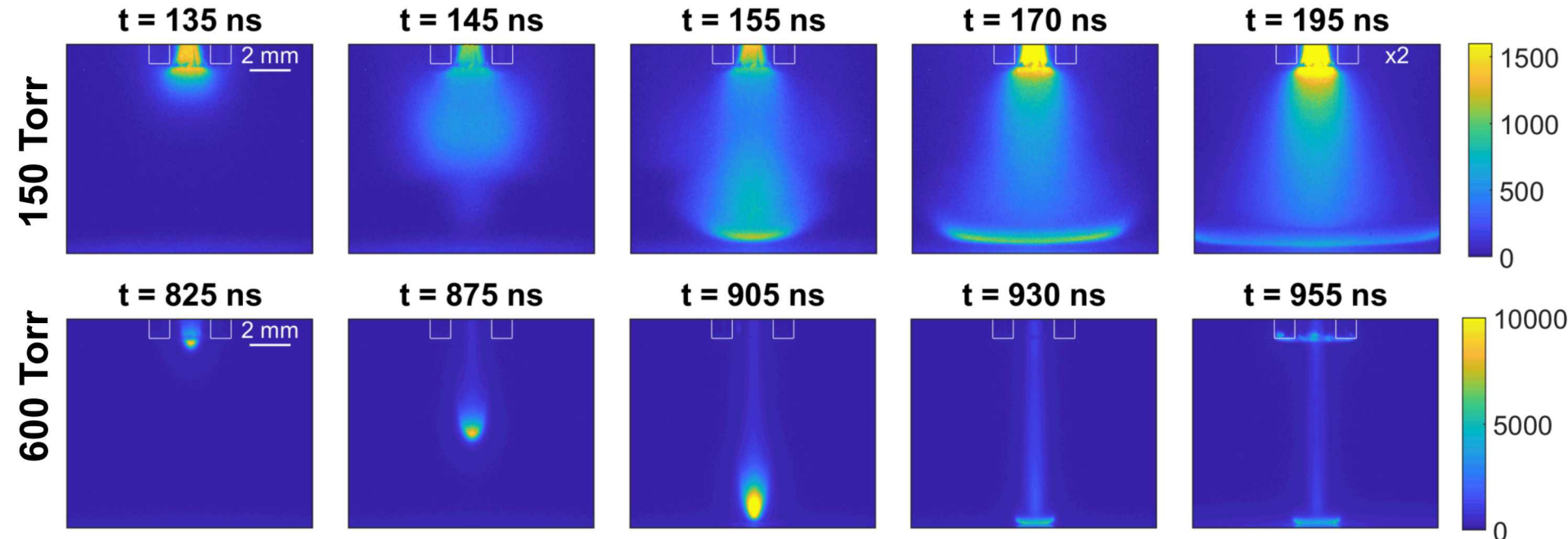
t = 310 ns



- +6 kV, 430 ns
- Pure He, 200 Torr

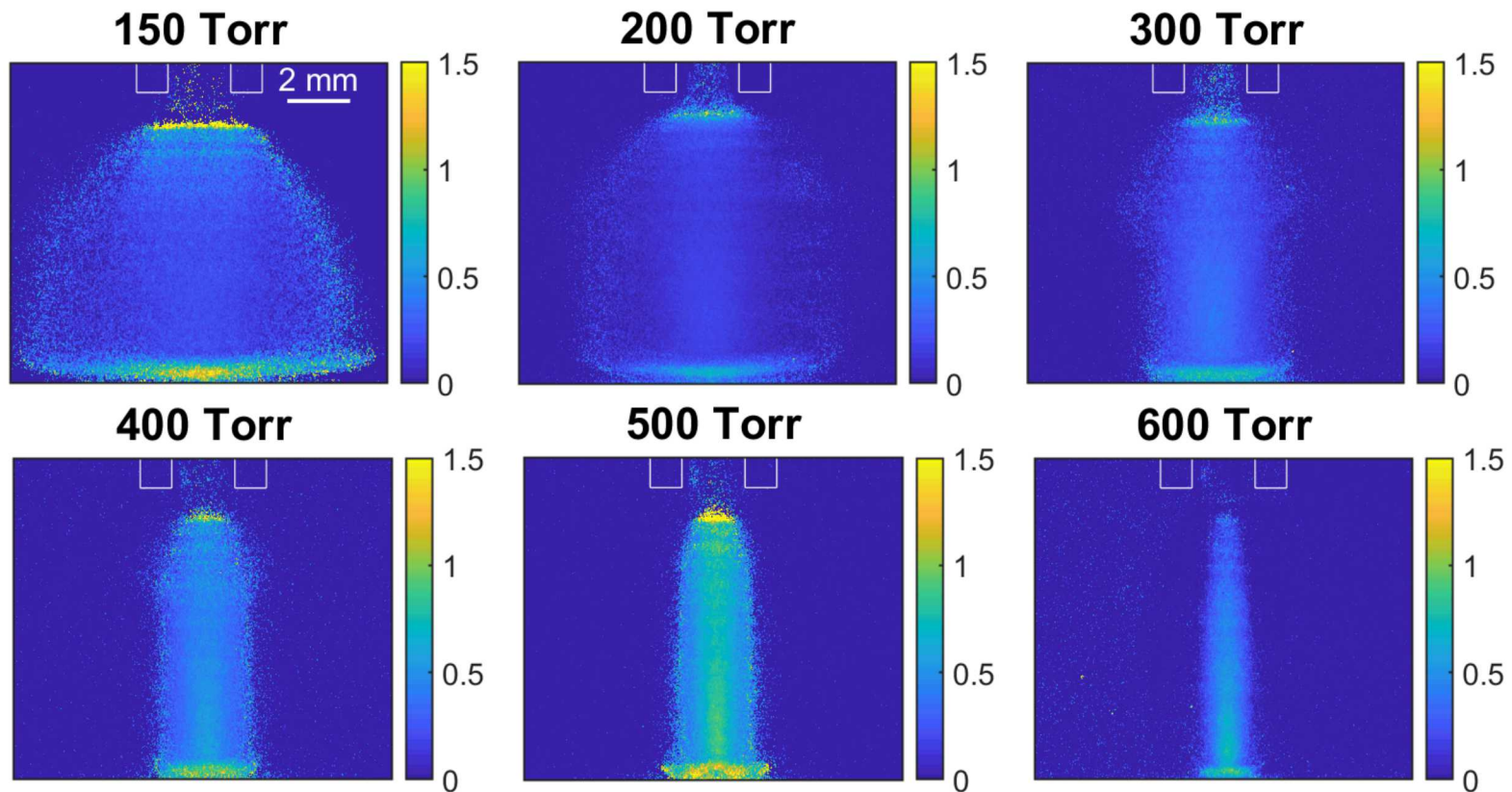
- Before IW reaches the surface, He(2^3S) densities are low.
- n_e in the SIW is nearly double that of the bulk.
- Elevated n_e in IW front may be due to Stark mixing.

VARY PRESSURE – EMISSION



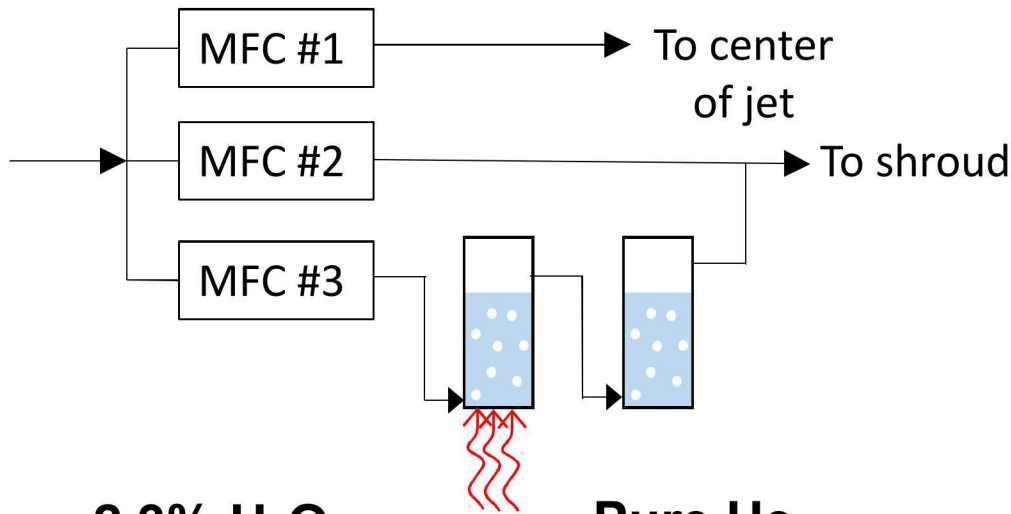
- 6 kV, 500 sccm He, 390 nm plasma emission
- Varied pulse duration so that voltage on for 80 ns after contact.
- Increasing pressure reduces electron mobility, preventing the IW from spreading as it exits the jet.
- IW propagates much slower for higher pressures.
- SIW becomes thinner at higher pressure.

VARY PRESSURE – n_e

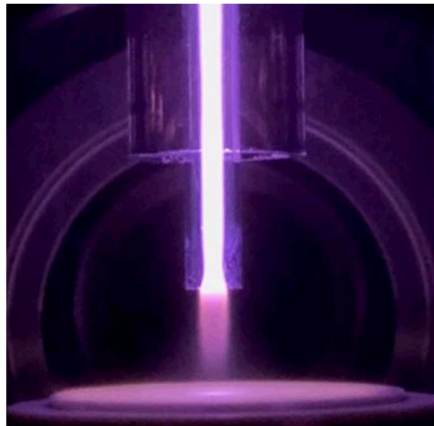


- 6 kV, 500 sccm He, 30 ns after IW contacts surface.
- Plasma is more confined at higher pressure, n_e increases.
- Current and energy deposition outside of the tube decrease with increasing pressure.
- Above 500 Torr, n_e is collisional enough that ionization rate drops.

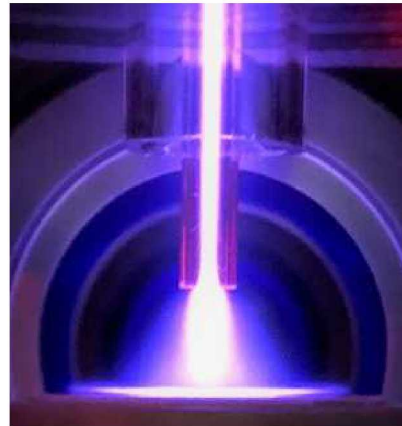
HUMID He SHROUD



2.3% H₂O
in shroud

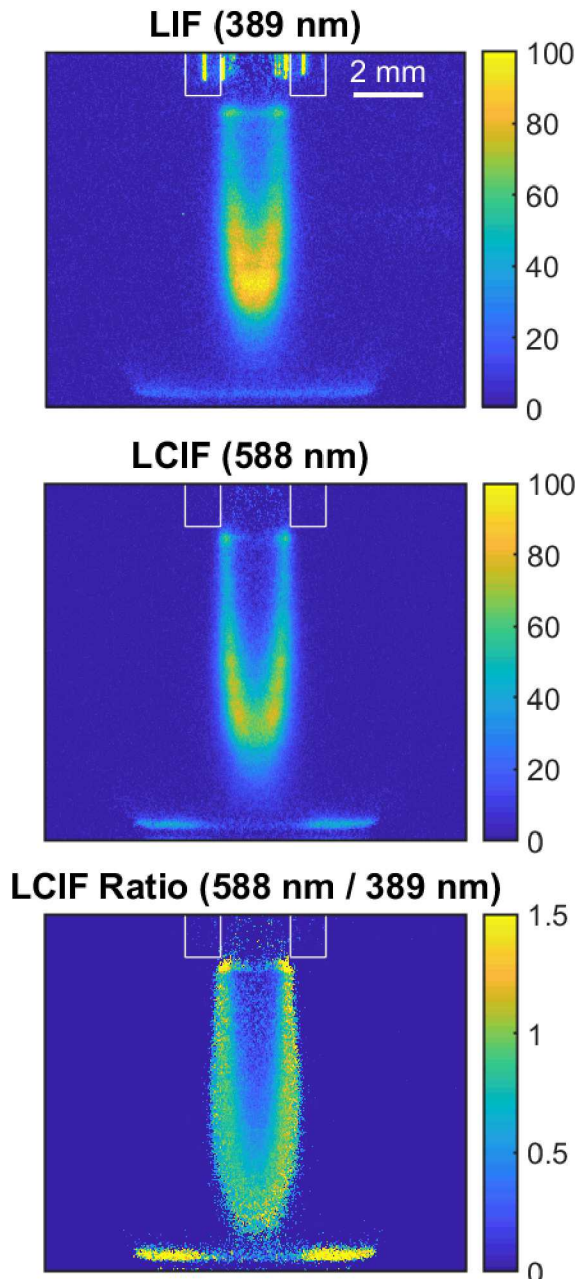


Pure He
(base case)

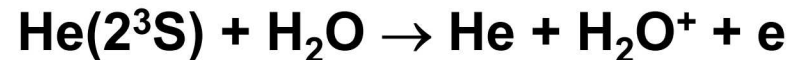


- Humid He shroud allows an investigation of molecular gasses surrounding the jet, without interfering with LCIF measurements.
- First bubbler oversaturates water vapor, second bubbler removes excess.
- Temperature of second bubbler determines humidity of gas.

HUMID He SHROUD

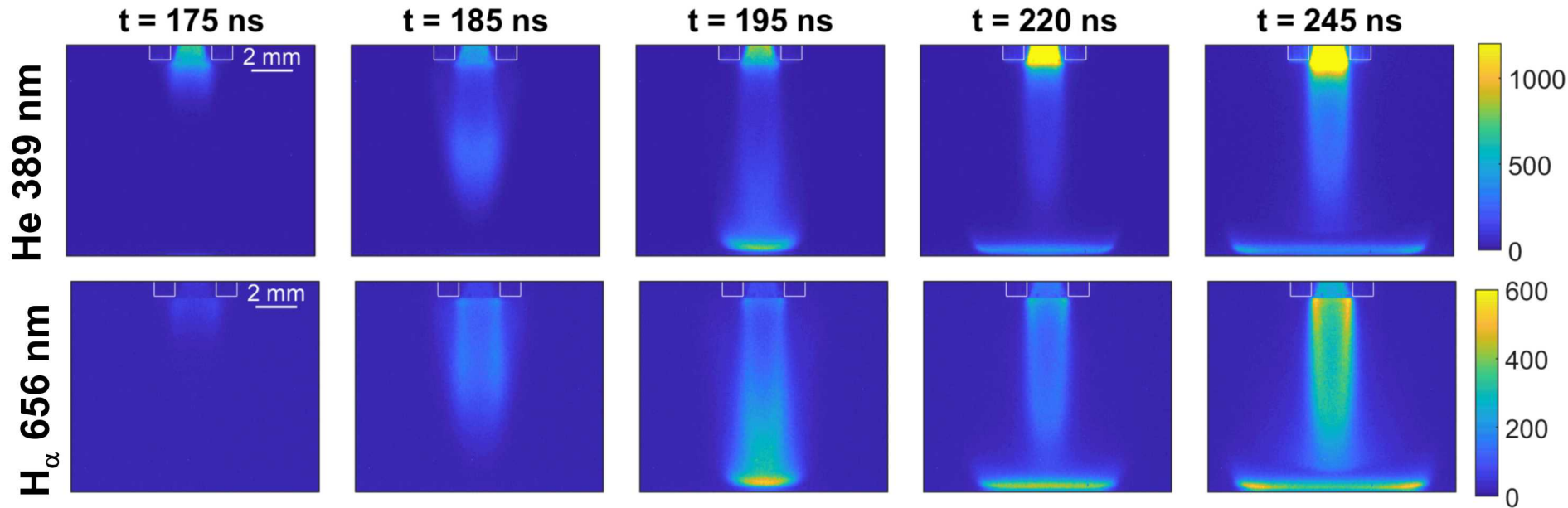


- 6 kV, 430 ns pulse
- Center: 500 sccm He
- Shroud: 500 sccm He/H₂O = 97.7/2.3
- $t = 230$ ns, 30 ns after IW contacts surface
- Moving away from He core, there are fewer He(2³S), LIF signal decreases.



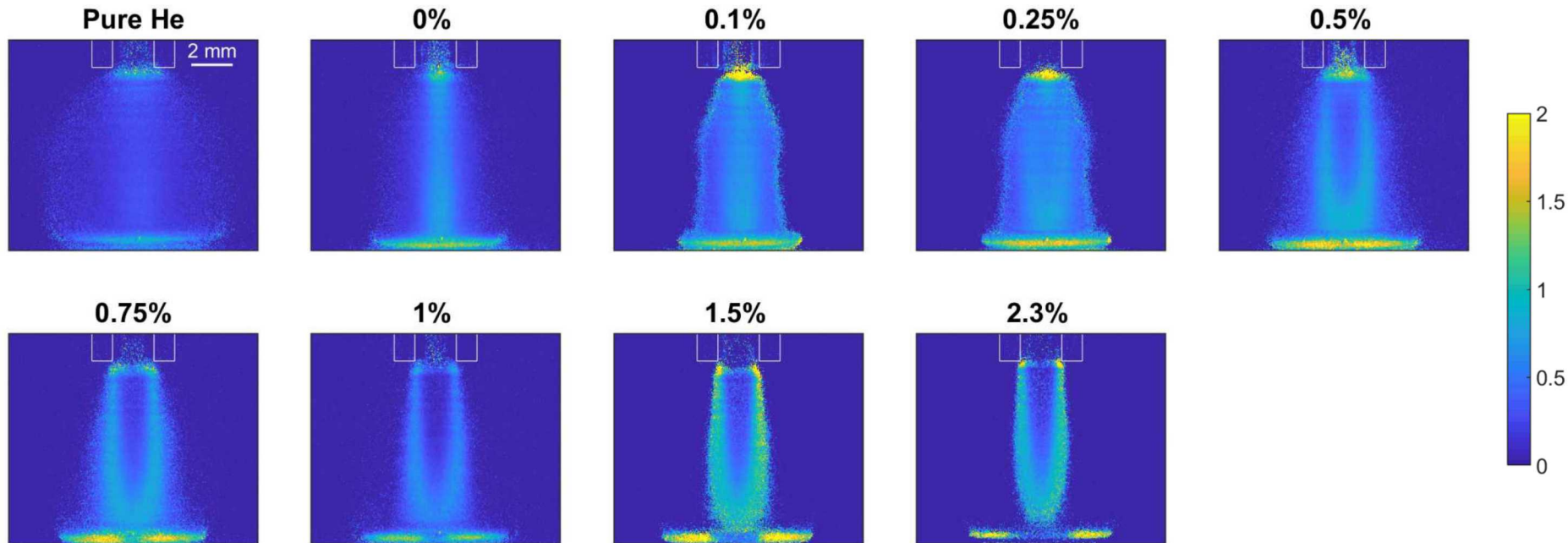
- In regions of high H₂O concentration, there may be significant n_e which is not detectable due to low He(2³S).

HUMID He SHROUD EMISSION



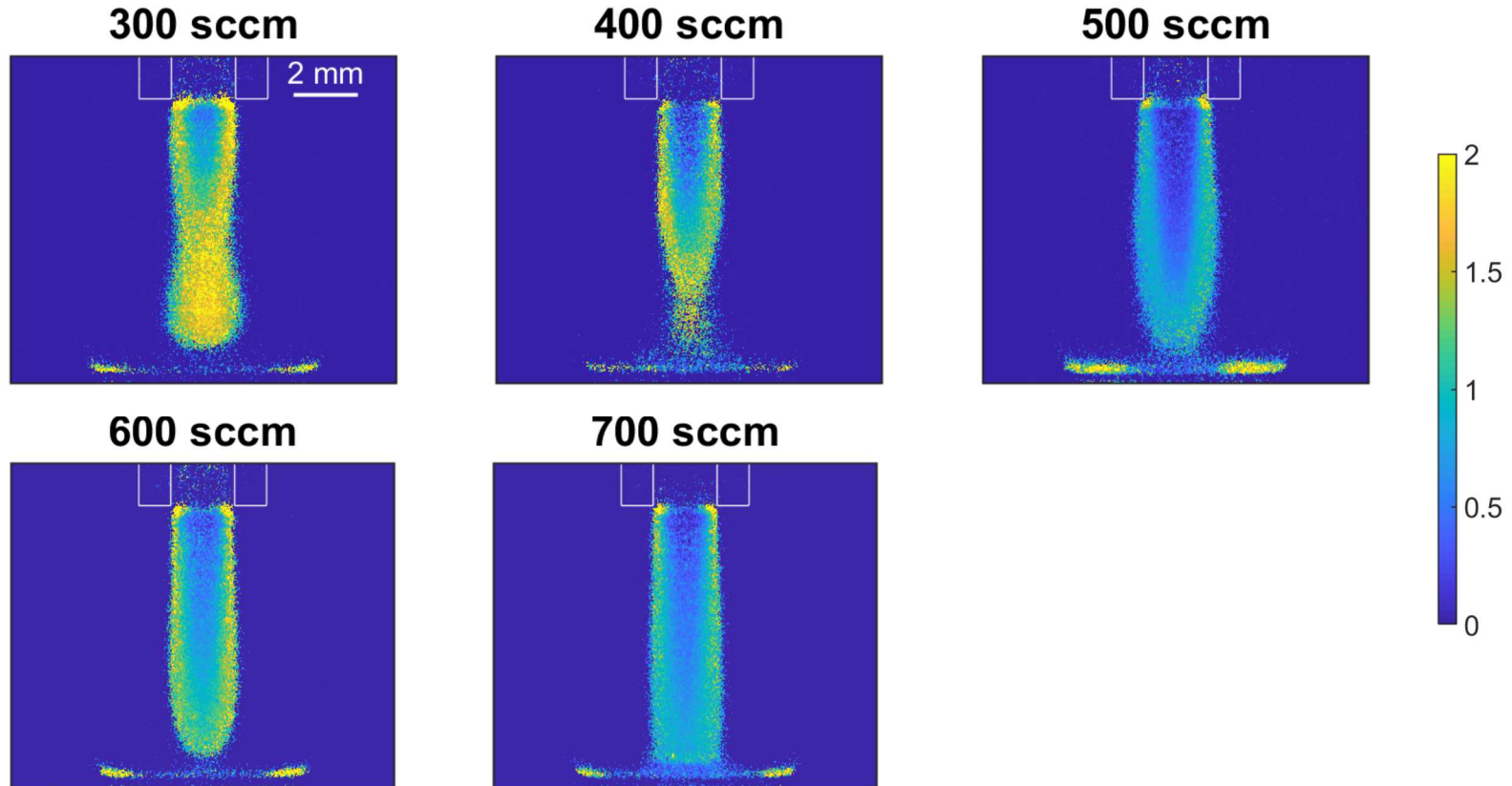
- 2.3% H₂O in shroud
- IW reaches outlet of the tube earlier than in base case – photoionization from He₂* causes non-local seed ionization.
- Photoionization and Penning ionization promote IW speed.
- H_α emission appears more annular – dominates at the interface of the center and shroud flow.

SHROUD HUMIDITY



- 30 ns after IW reaches surface.
- Transition from diffuse in pure helium case to confined by humid shroud.
- Higher electron energy loss rates with H_2O because of vibrational and rotational excitation.
- n_e increases with humidity due to Penning ionization.

VARY He FLOW RATE - n_e



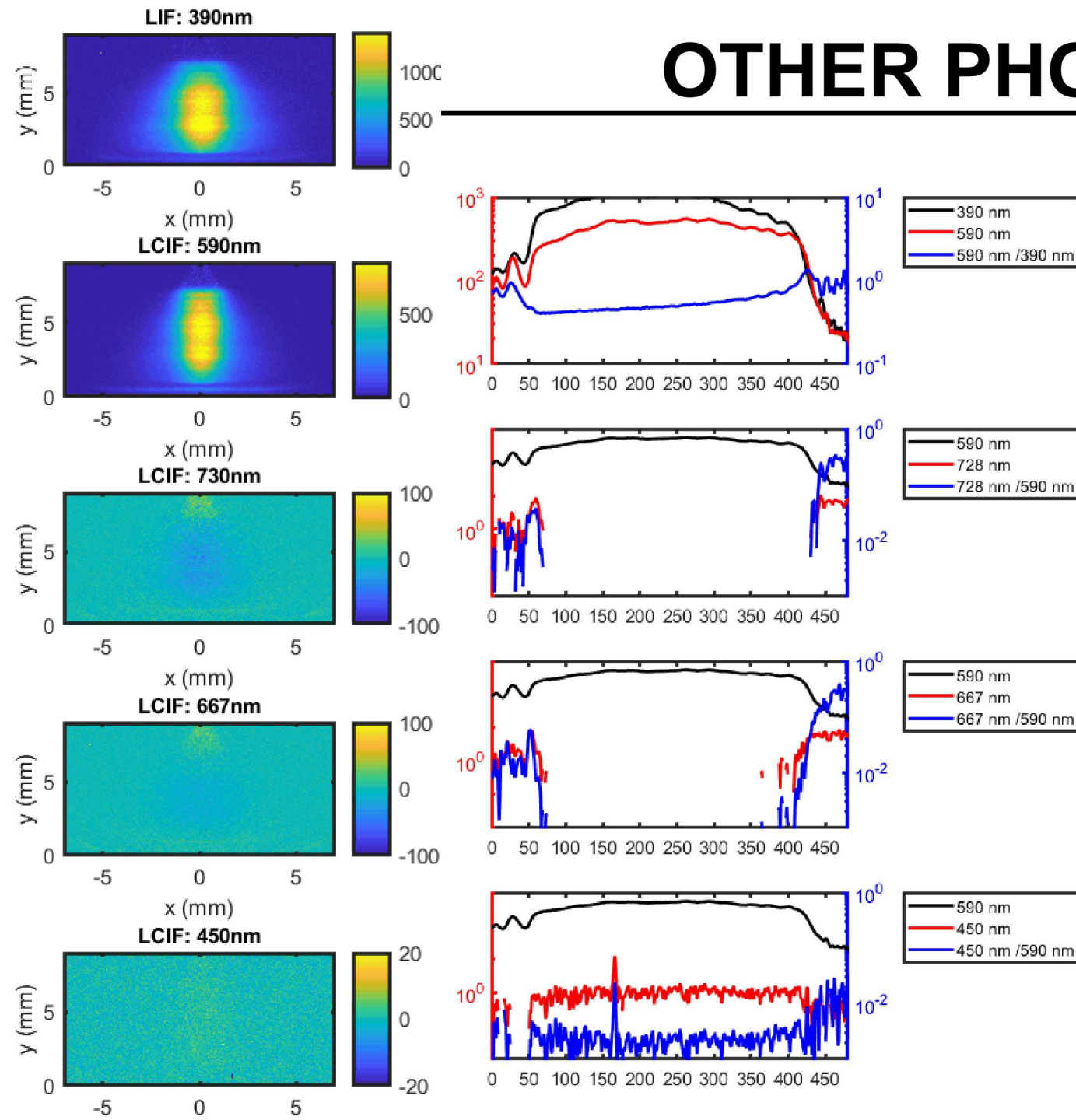
- Vary center flow rate. Shroud: 500 sccm, He/H₂O = 98.7/2.3
- Higher He flow rates more rapidly convect in-diffusing H₂O.

CONCLUDING REMARKS

- A plasma jet interacting with a dielectric surface produces a surface IW as it charges the dielectric.
- Molecular gases surrounding an APPJ confine the IW due to a lower electron mobility and higher electron energy loss rates.
- In the presence of H₂O, the IW speed increase, and the plasma and electron density become more annular.
- LCIF can be used in humid He, as long as the He(2³S) density is sufficient.

Appendix

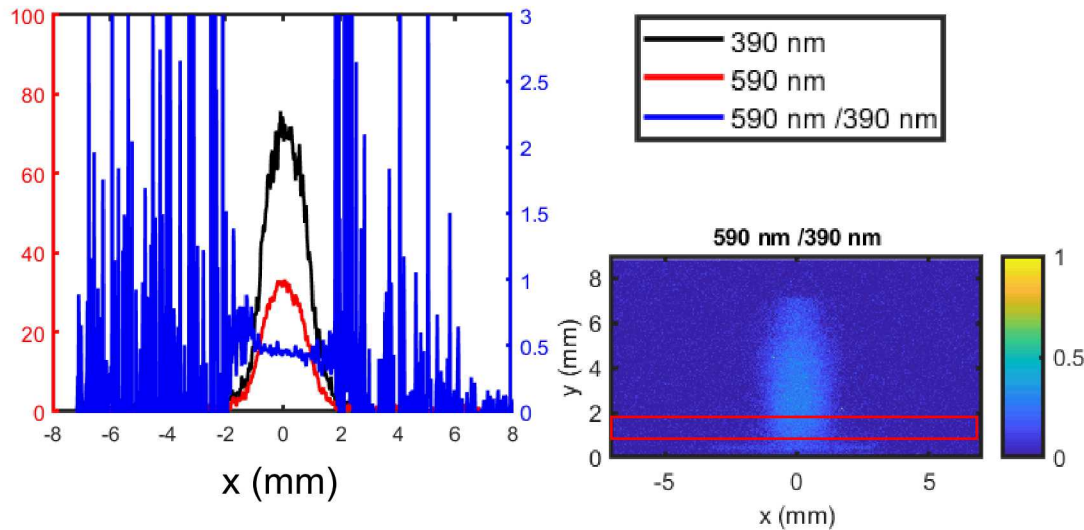
OTHER PHOTONS



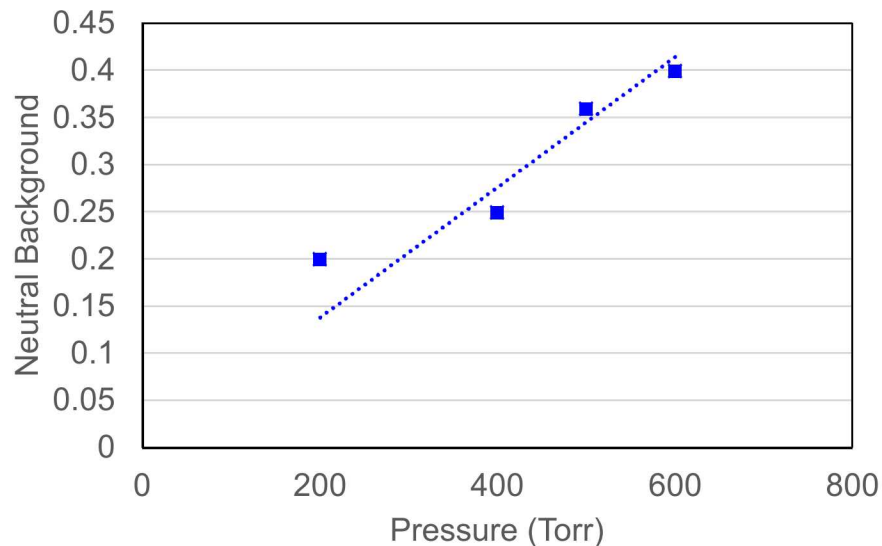
- $t = 310$ ns
- 730 nm and 667 nm are consistently negative (depletion of $\text{He}2^3\text{S}$ by laser?)
- Ratios indicate $E/N \sim 2$ Td

NEUTRAL BACKGROUND

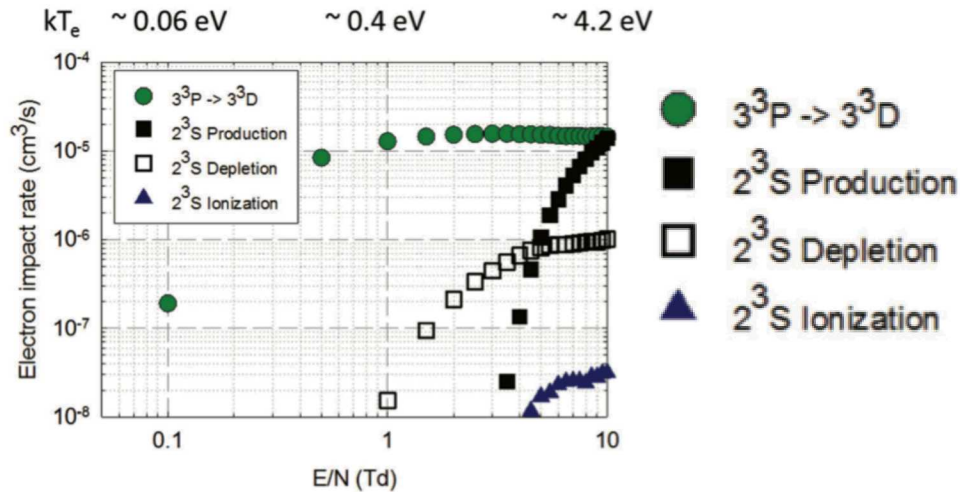
600 Torr



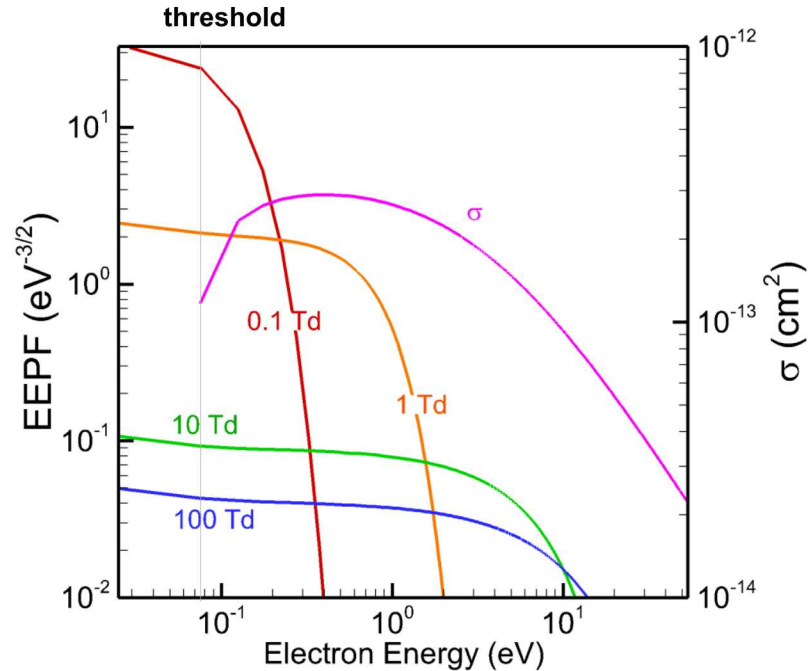
- There is some portion of the LCIF signal which is caused by collisions with the background He gas.
- Should be proportional to pressure.
- Images take $>4.5 \mu\text{s}$ after pulse ends.
- In the late afterglow electrons are thermalized and no longer contribute to LCIF.



BOLTZMANN CALCULATIONS

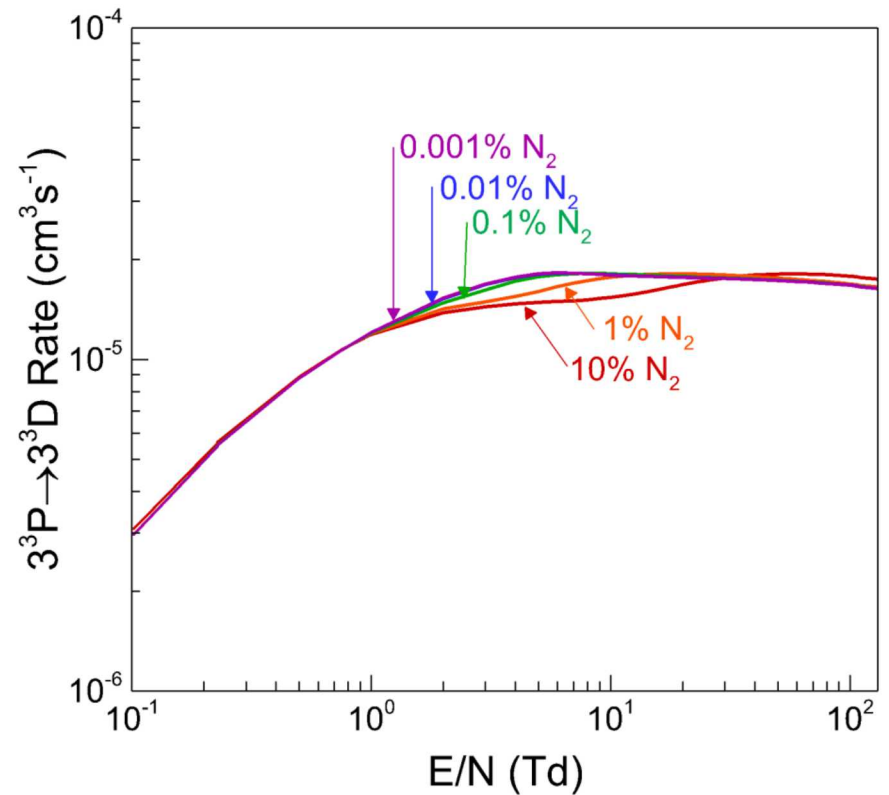
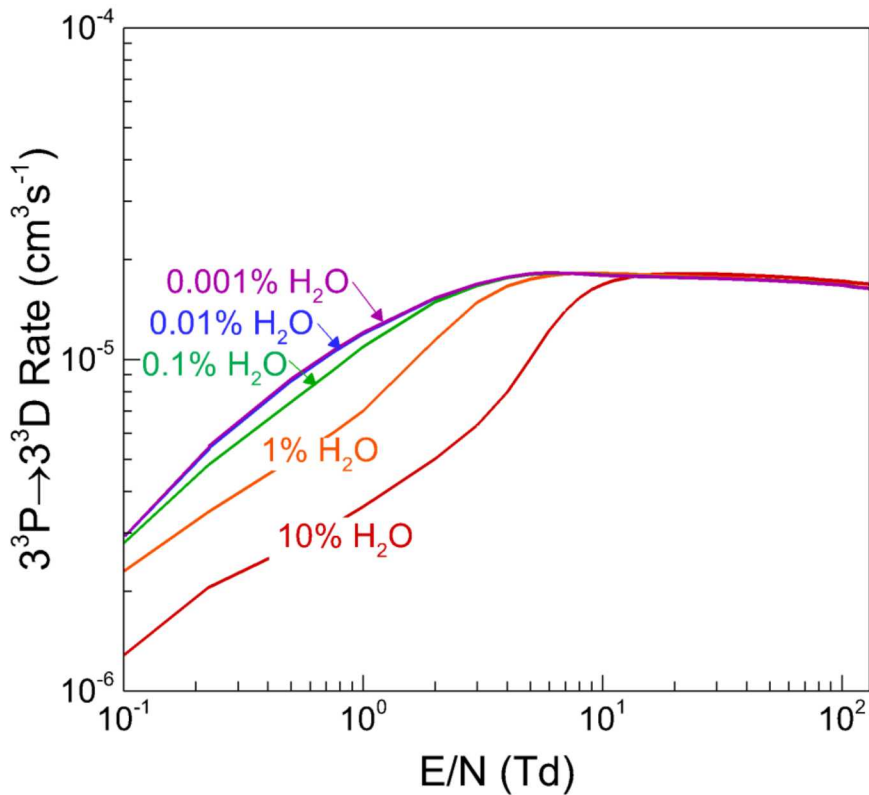


(b)



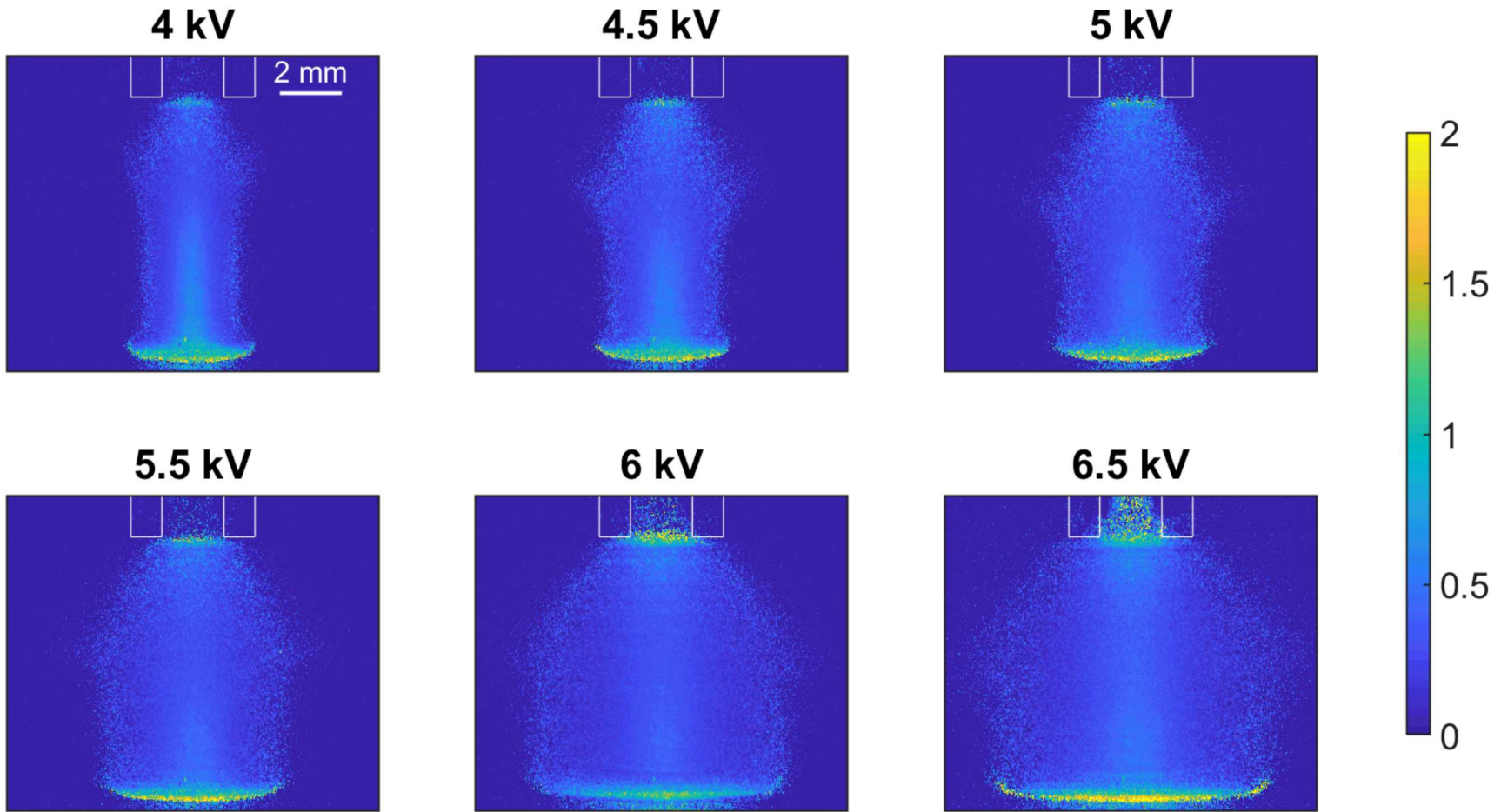
- For LCIF measurement to be a linear representation of n_e , rate of $3^3\text{P} \rightarrow 3^3\text{D}$ must be independent of T_e .
- In pure He, this occurs when $E/N > \sim 0.8 \text{ Td}$
- Threshold = 0.06 eV

BOLTZMANN CALCULATIONS



- Cross section set to 0 for $E < 0.15 \text{ eV}$
- For $\text{H}_2\text{O} < 1\%$, LCIF is valid $\sim 1 \text{ Td}$
- For all N_2 concentrations, LCIF is valid for $> 0.8 \text{ Td}$

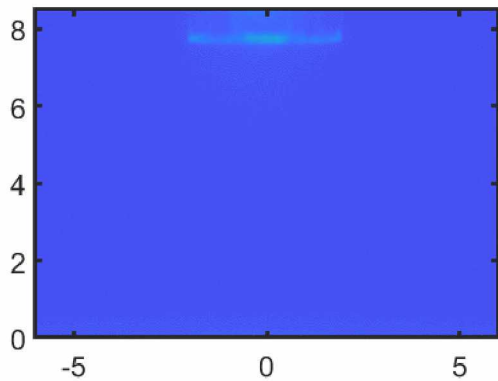
VARY VOLTAGE – n_e



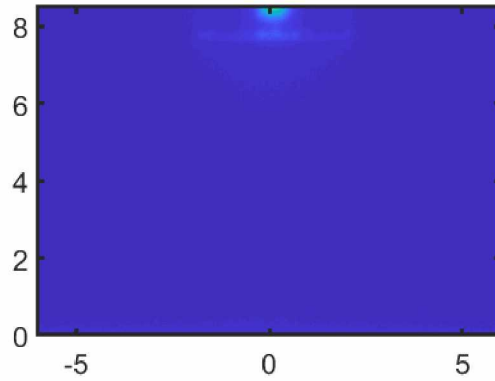
- LCIF ratio 30 ns after IW contacts surface.
- SIW spreads more rapidly for higher voltage, but n_e is only slightly higher.

VARY VOLTAGE – Emission

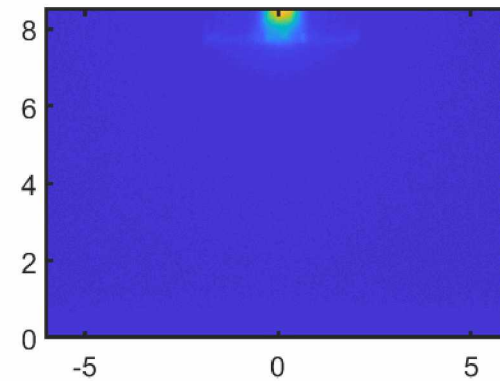
4 kV



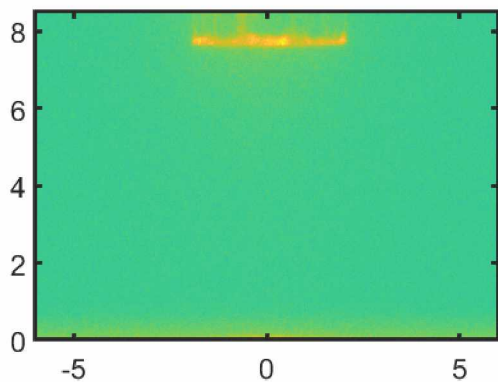
4.5 kV



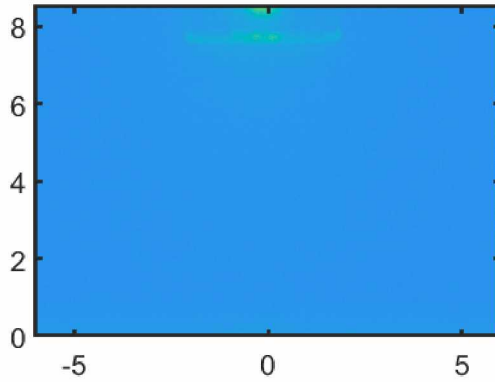
5 kV



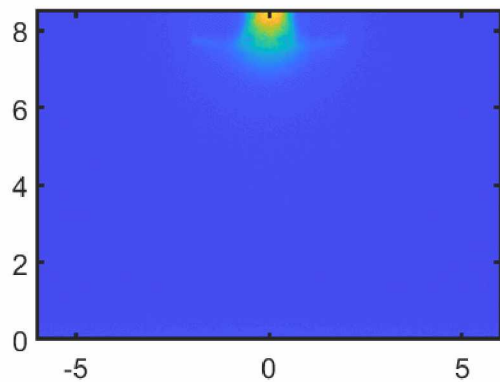
5.5 kV



6 kV (base)



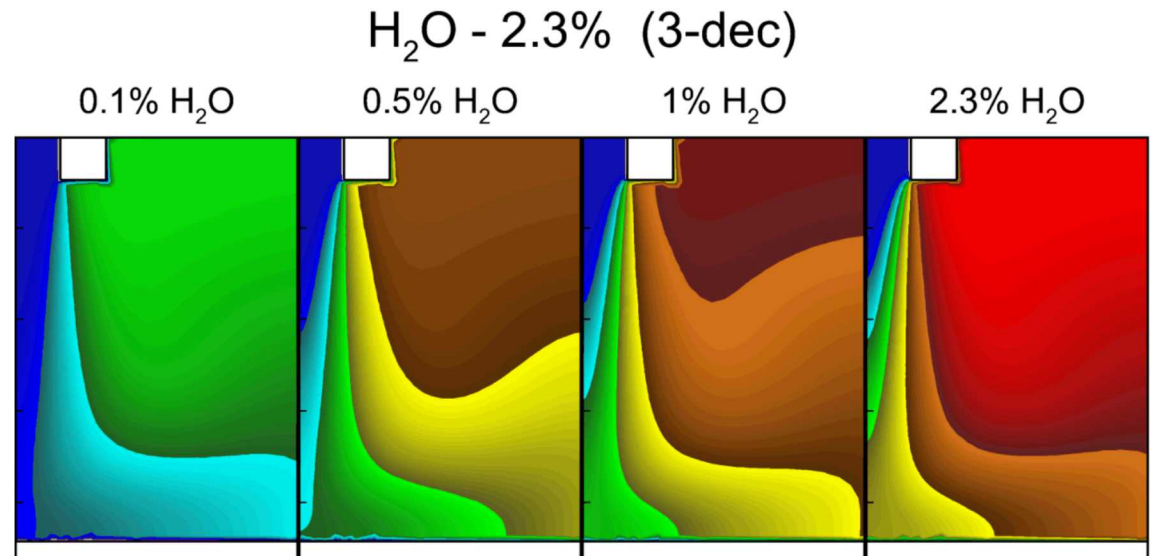
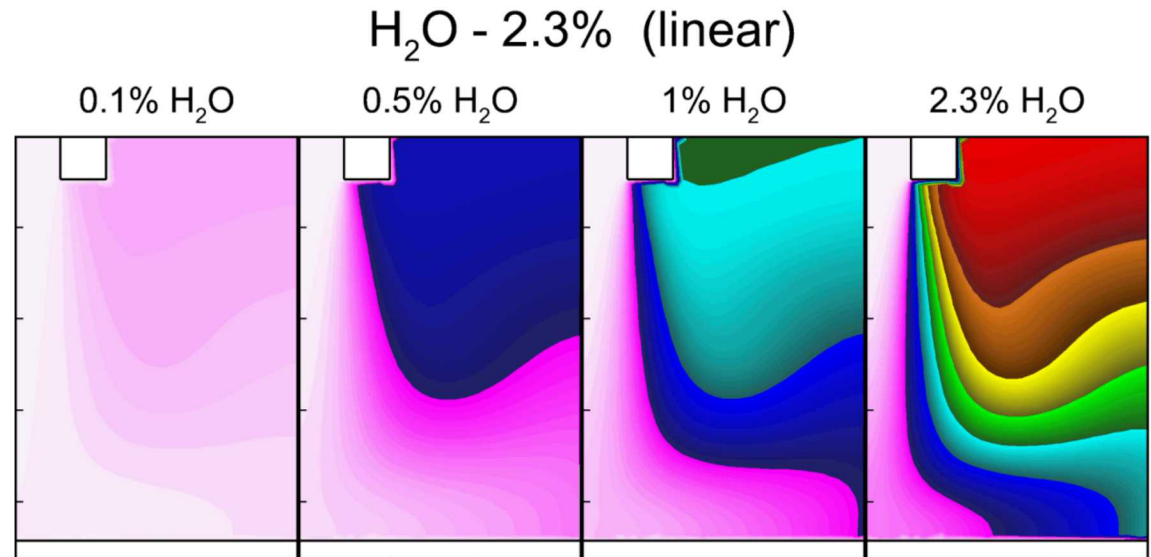
6.5 kV



- Pure He, 500 sccm
- Bubbler installed (leak rate slightly higher)
- 250 ns animated
- 390 nm emission
- 20180410, 20180411, 6kV on 20180403

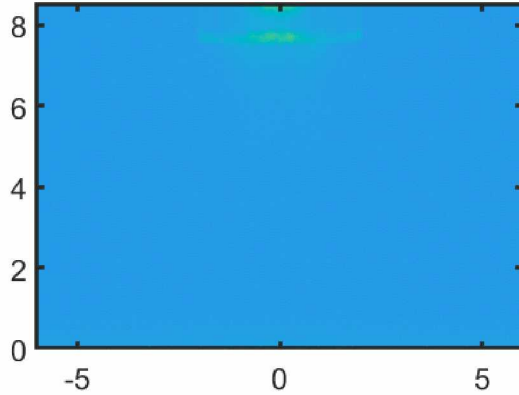
VARY SHROUD HUMIDITY

- 24 ms fluid simulation
- Region of high ne does not correspond to a particular % H₂O
- Chemistry and electron energy losses change with % H₂O, but mobility changes very little

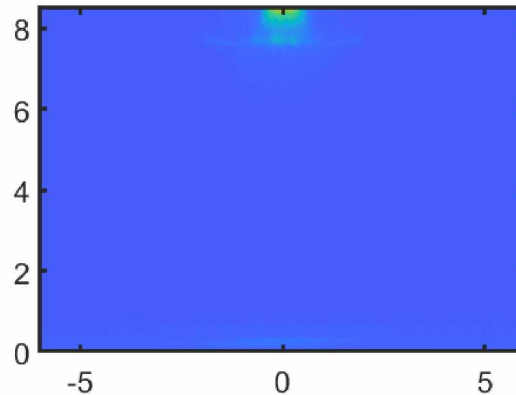


Emission - Vary Shroud Humidity

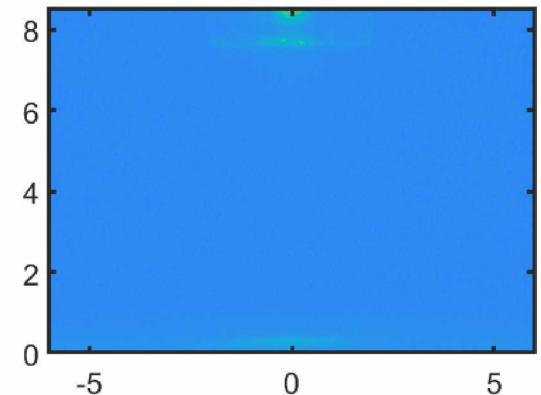
2.3% H₂O



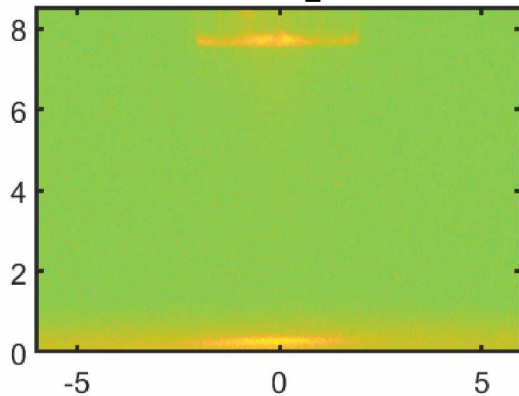
1% H₂O



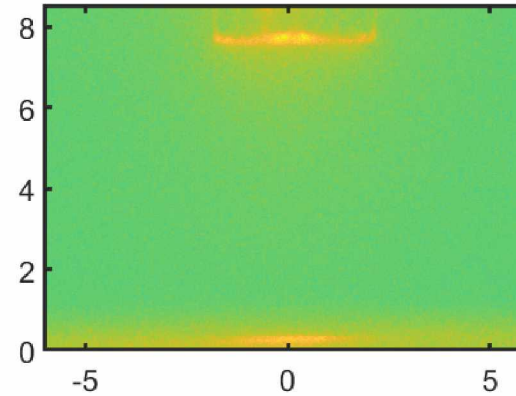
0.5% H₂O



0.1% H₂O

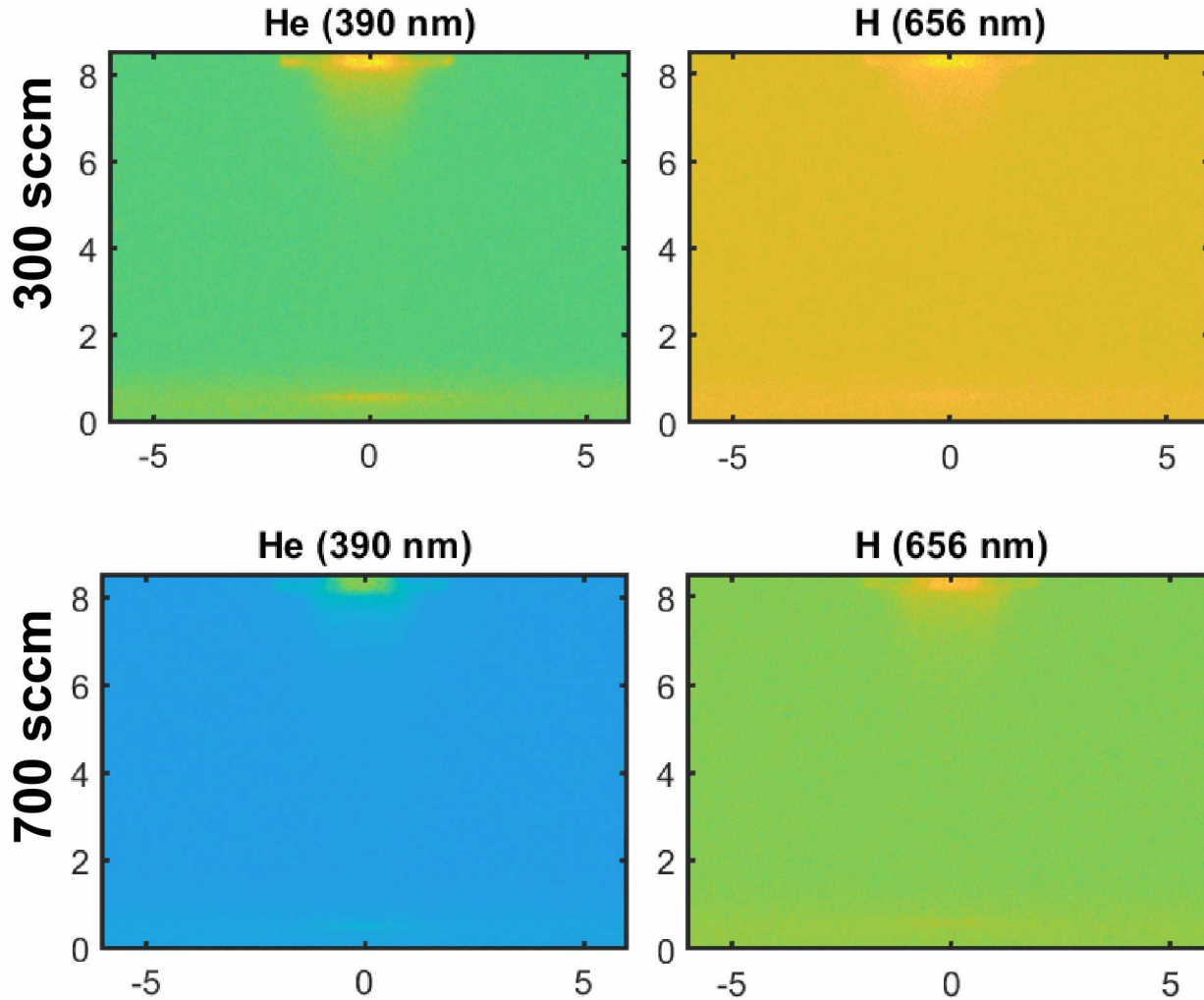


Pure He



- 390 nm emission, 250 ns animated
- IW propagates across the gap faster for higher H₂O – penning ionization, photoionization
- SIW spreads less rapidly for higher H₂O
- 20180404 (2.3%), 20180327(1%,0.1%,pure He), 20180329 (0.5%)

VARY He FLOW RATE – EMISSION



- 250 ns animated
- OH emission was not visible – UV filter has lower efficiency, camera is less sensitive
- H_α is much more accessible
- H_α emission is more annular at greater distances for 700 sccm