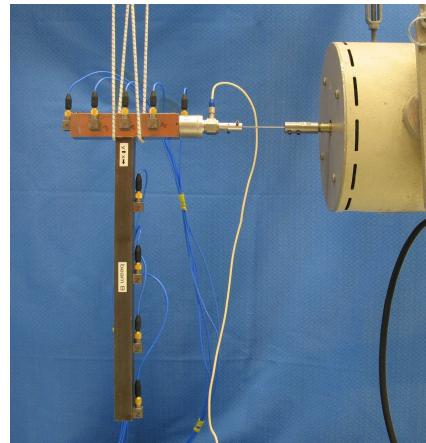


A Method for Canceling Force Transducer Mass and Inertia Effects

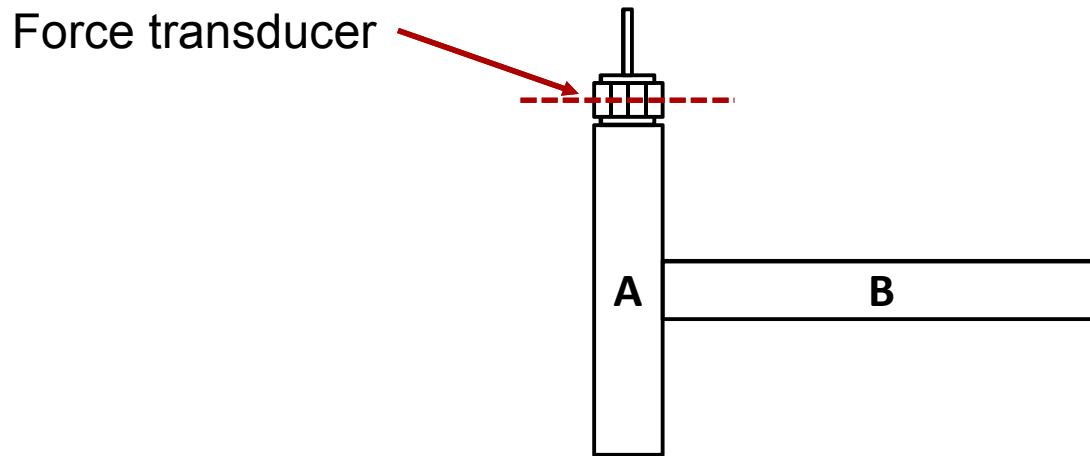
Garrett K. Lopp, Benjamin R. Pacini, Randall L. Mayes

22 August 2017

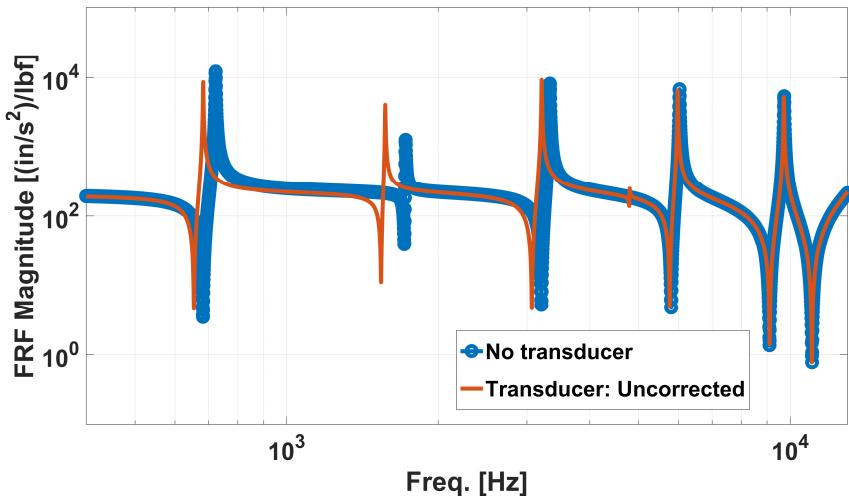


Shaker Test Setup

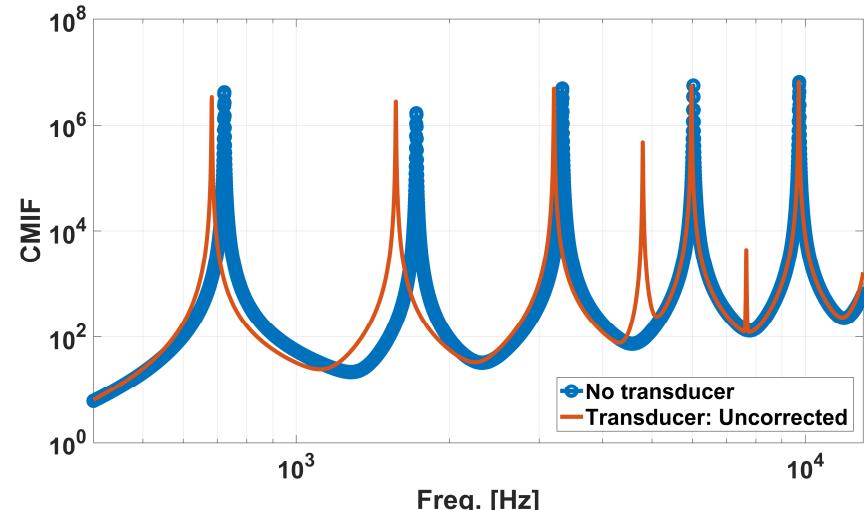
Force transducer may drastically alter system dynamics if not accounted for.



Drive-point FRF

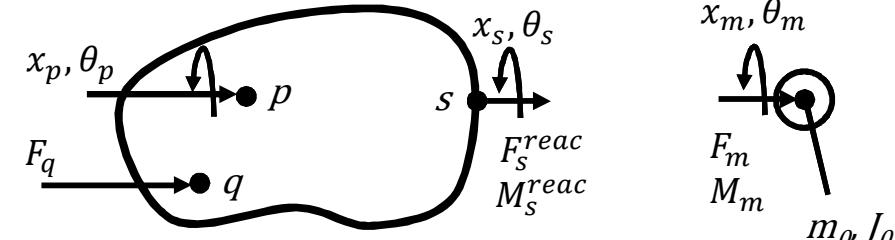
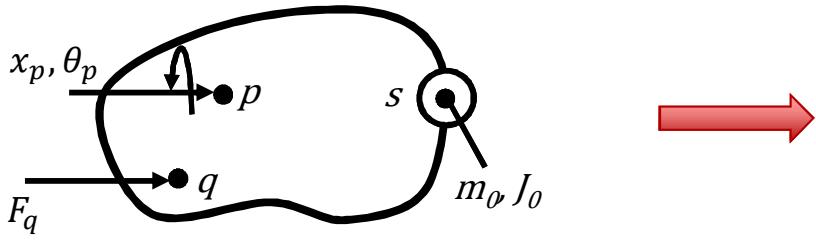


Complex Modal Indicator Function



Both translational and rotational quantities can pollute the measured FRFs

Begin with a similar approach as Ashory (1998 IMAC) but include rotation terms



Frequency response at output p :

$$\begin{Bmatrix} x_p \\ \theta_p \end{Bmatrix} = \begin{Bmatrix} A_{pq} \\ C_{pq} \end{Bmatrix} F_q + \begin{bmatrix} A_{ps} & B_{ps} \\ C_{ps} & D_{ps} \end{bmatrix} \begin{Bmatrix} F_s^{\text{reac}} \\ M_s^{\text{reac}} \end{Bmatrix}$$

Frequency response of point mass m :

$$\begin{Bmatrix} x_m \\ \theta_m \end{Bmatrix} = \begin{bmatrix} -\frac{1}{\omega^2 m_0} & 0 \\ 0 & -\frac{1}{\omega^2 J_0} \end{bmatrix} \begin{Bmatrix} F_m \\ M_m \end{Bmatrix}$$

Subject to the constraints:

$$\begin{Bmatrix} x_s \\ \theta_s \end{Bmatrix} - \begin{Bmatrix} x_m \\ \theta_m \end{Bmatrix} = 0$$

$$\begin{Bmatrix} F_s^{\text{reac}} \\ M_s^{\text{reac}} \end{Bmatrix} + \begin{Bmatrix} F_m \\ M_m \end{Bmatrix} = 0$$

Considering the drive-point ($p = q = s$)

$$\begin{Bmatrix} x_p \\ \theta_p \end{Bmatrix} = \begin{bmatrix} 1 - \omega^2 m_0 A_{pp} & -\omega^2 J_0 B_{pp} \\ -\omega^2 m_0 C_{pp} & 1 - \omega^2 J_0 D_{pp} \end{bmatrix}^{-1} \begin{Bmatrix} A_{pp} \\ C_{pp} \end{Bmatrix} F_p$$

Both translational and rotational quantities modify the true FRF

- McConnell (1999 MSSP) shows a similar result

[1] Ashory, "Correction of Mass-Loading Effects of Transducers and Suspension Effects in Modal Testing," IMAC, 1998.

[2] McConnell and Cappa, "Transducer Inertia and Stinger Stiffness Effects on FRF Measurements," MSSP, 1999.

Presentation Outline



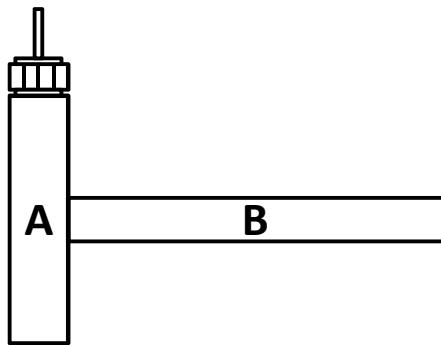
- Theoretical Development
- Numerical Results
- Experimental Results
- Conclusions

Presentation Outline

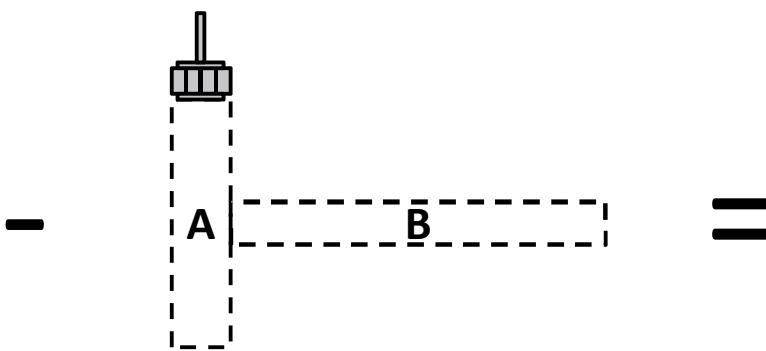


- **Theoretical Development**
 - Numerical Results
 - Experimental Results
 - Conclusions

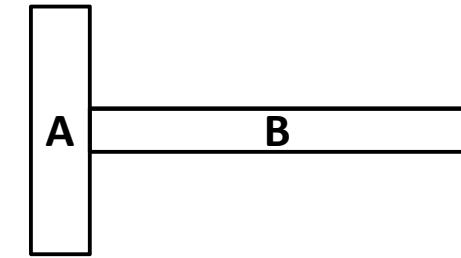
Substructuring enables removal of the force transducer from the system.



Experimental system with transducer, CT



Analytical model of transducer, T



Experimental system, C

Equations of motion:

$$\begin{bmatrix} M_{CT} & 0 \\ 0 & -M_T \end{bmatrix} \begin{Bmatrix} \ddot{u}_{CT} \\ \ddot{u}_T \end{Bmatrix} + \begin{bmatrix} C_{CT} & 0 \\ 0 & 0 \end{bmatrix} \begin{Bmatrix} \dot{u}_{CT} \\ \dot{u}_T \end{Bmatrix} + \begin{bmatrix} K_{CT} & 0 \\ 0 & 0 \end{bmatrix} \begin{Bmatrix} u_{CT} \\ u_T \end{Bmatrix} = \begin{Bmatrix} f_{CT} \\ 0 \end{Bmatrix} + \begin{Bmatrix} g_{CT} \\ g_T \end{Bmatrix}$$

Subject to:

Compatibility at interface: $u_{CT,T} - u_T = 0 \longrightarrow B \begin{Bmatrix} u_{CT} \\ u_T \end{Bmatrix} = 0$

Interface force equilibrium: $g_{CT} + g_T = 0 \longrightarrow L^T \begin{Bmatrix} g_{CT} \\ g_T \end{Bmatrix} = 0$

System analytical model is unknown; instead, use modal equations derived from tests.

Transformation to the modal coordinate system retaining R_{CT} measured system modes and R_T transducer rigid body modes :

$$\begin{Bmatrix} u_{CT} \\ u_T \end{Bmatrix} = \begin{bmatrix} \Phi_{CT} & 0 \\ 0 & \Phi_T \end{bmatrix} \begin{Bmatrix} q_{CT} \\ q_T \end{Bmatrix}$$

Modal equations of motion:

$$\begin{bmatrix} I_{CT} & 0 \\ 0 & -I_T \end{bmatrix} \begin{Bmatrix} \ddot{q}_{CT} \\ \ddot{q}_T \end{Bmatrix} + \begin{bmatrix} [2\zeta_r \omega_r]_{CT} & 0 \\ 0 & 0 \end{bmatrix} \begin{Bmatrix} \dot{q}_{CT} \\ \dot{q}_T \end{Bmatrix} + \begin{bmatrix} [\omega_r^2]_{CT} & 0 \\ 0 & 0 \end{bmatrix} \begin{Bmatrix} q_{CT} \\ q_T \end{Bmatrix} \\ = \begin{Bmatrix} \Phi_{CT}^T f_{CT} \\ 0 \end{Bmatrix} + \begin{bmatrix} \Phi_{CT}^T & 0 \\ 0 & \Phi_T^T \end{bmatrix} \begin{Bmatrix} g_{CT} \\ g_T \end{Bmatrix}$$

Subject to:

Compatibility at interface:

$$B \begin{bmatrix} \Phi_{CT} & 0 \\ 0 & \Phi_T \end{bmatrix} \begin{Bmatrix} q_{CT} \\ q_T \end{Bmatrix} = B_p \begin{Bmatrix} q_{CT} \\ q_T \end{Bmatrix} = 0$$

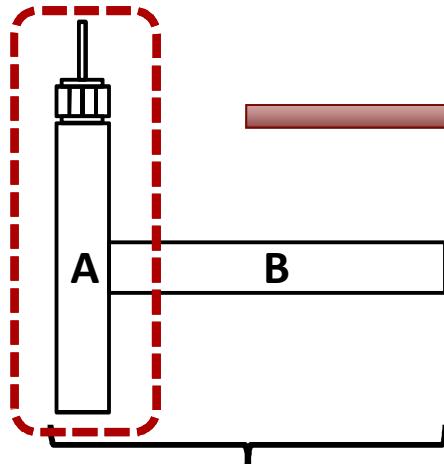
Requires rotational mode shapes at connection point which are unavailable

Interface force equilibrium:

$$L_p^T \begin{bmatrix} \Phi_{CT}^T & 0 \\ 0 & \Phi_T^T \end{bmatrix} \begin{Bmatrix} g_{CT} \\ g_T \end{Bmatrix} = L_p^T \begin{Bmatrix} g_{p,CT} \\ g_{p,T} \end{Bmatrix} = 0$$

Need to find a different set of B_p and B_q using measurable quantities!

Can estimate connection point motion using an analytical model of the local system.



Measured system CT



Analytical model of AT

Constrain translation measurements on system **CT** and local system **AT**:

$$\mathbf{u}_{AT,m} = \mathbf{u}_{CT,m}$$

Transform into the modal coordinate system:

$$\Phi_{AT,m} \mathbf{q}_{AT} = \Phi_{CT,m} \mathbf{q}_{CT}$$

Take pseudoinverse of analytical mode shape matrix

$$\mathbf{q}_{AT} = \Phi_{AT,m}^\dagger \Phi_{CT,m} \mathbf{q}_{CT}$$

Expand response to estimate full motion at transducer connection point:

$$\mathbf{u}_{AT,T} = \Phi_{AT,T} \mathbf{q}_{AT}$$

$$\mathbf{u}_{AT,T} = \underbrace{\left[\Phi_{AT,T} \Phi_{AT,m}^\dagger \Phi_{CT,m} \right]}_{\hat{\Phi}_{CT,T}} \mathbf{q}_{CT} = \hat{\mathbf{u}}_{CT,T}$$

Estimated mode shape matrix now contains rotational information!

Previously unknown constraint matrices are found using measurable quantities.

Constrain modal motion of transducer \mathbf{T} and measured system \mathbf{CT}

$$\hat{\Phi}_{CT,T} \mathbf{q}_{CT} - \Phi_T \mathbf{q}_T = \mathbf{0} \quad \Rightarrow$$

$$\boxed{\underbrace{[\hat{\Phi}_{CT,T} - \Phi_T]}_{\mathbf{B}_p} \begin{Bmatrix} \mathbf{q}_{CT} \\ \mathbf{q}_T \end{Bmatrix} = \mathbf{0}}$$

\mathbf{B}_p enforces as many constraints as there are transducer rigid body modes

Have R_T constraints and $R_{CT} + R_T$ modal equations of motion – can transform to a set unconstrained equations by simply choosing \mathbf{q}_{CT} as the unconstrained modes:

$$\boxed{\begin{Bmatrix} \mathbf{q}_{CT} \\ \mathbf{q}_T \end{Bmatrix} = \underbrace{\begin{bmatrix} \mathbf{I}_{CT} \\ \Phi_T^{-1} \hat{\Phi}_{CT,T} \end{bmatrix}}_{\mathbf{L}_p} \mathbf{q}_{CT}}$$

Force transducer can now be immediately removed from the experimental system.

Apply the transformation to the modal EOM

$$\begin{Bmatrix} \mathbf{q}_{CT} \\ \mathbf{q}_T \end{Bmatrix} = \mathbf{L}_p \mathbf{q}_{CT}$$

Unconstrained equations of motion

$$\tilde{\mathbf{M}}\ddot{\mathbf{q}}_{CT} + \tilde{\mathbf{C}}\dot{\mathbf{q}}_{CT} + \tilde{\mathbf{K}}\mathbf{q}_{CT} = \tilde{\mathbf{f}} + \tilde{\mathbf{g}}$$

$$\tilde{\mathbf{M}} = \mathbf{L}_p^T \begin{bmatrix} \mathbf{I}_C & \mathbf{0} \\ \mathbf{0} & -\mathbf{I}_T \end{bmatrix} \mathbf{L}_p = \mathbf{I}_{CT} - \hat{\Phi}_{CT,T}^T [\Phi_T \Phi_T^T]^{-1} \hat{\Phi}_{CT,T}$$

$$\tilde{\mathbf{C}} = -\mathbf{L}_p^T \begin{bmatrix} [\sqrt{2\zeta_n \omega_n}]_{CT} & \mathbf{0} \end{bmatrix} = -\mathbf{L}_p^T \begin{bmatrix} \mathbf{0} & \mathbf{0} \end{bmatrix}$$

$$\tilde{\mathbf{K}} = \mathbf{I}_{CT} - \Phi_{CT,m}^T \underbrace{\left[(\Phi_{AT,T} \Phi_{AT,m}^\dagger)^T (\Phi_T \Phi_T^T)^{-1} \Phi_{AT,T} \Phi_{AT,m}^\dagger \right]}_{M_{T,m}} \Phi_{CT,m}$$

This process creates equivalent transducer mass with zero inertia at measurement locations allowing for immediate removal!

$$\mathbf{L}_p^T \begin{bmatrix} \mathbf{0} & \Phi_T^T \end{bmatrix} \begin{bmatrix} \mathbf{g}_T \end{bmatrix}$$

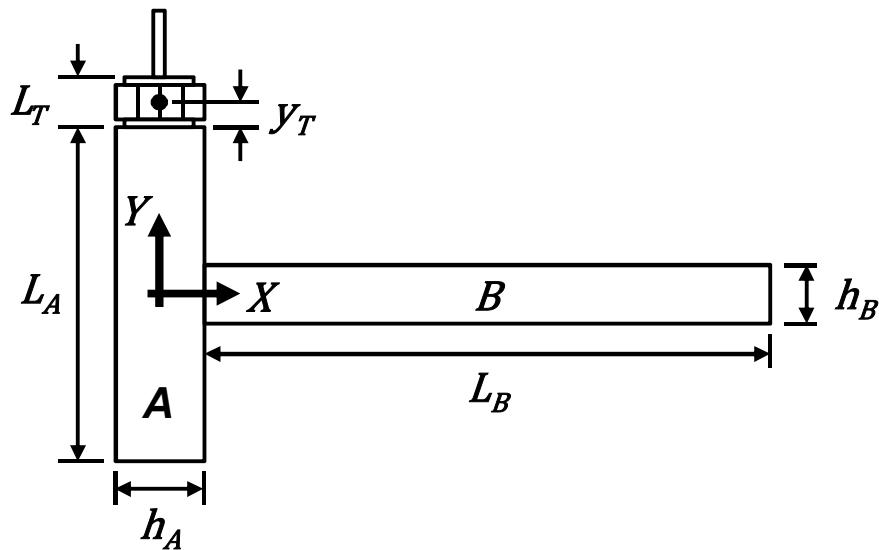
Presentation Outline



- Theoretical Development
- **Numerical Results**
- Experimental Results
- Conclusions

Overview of numerical model

Beam system modeled using FEM

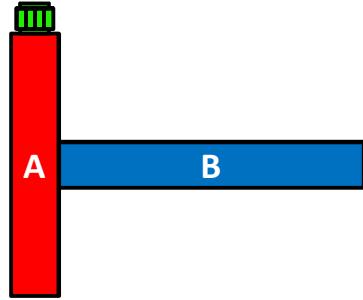


	Beam A	Beam B
Material	Steel	Steel
Length [mm]	114	305
Width [mm]	25.4	25.4
Thickness [mm]	25.4	19.1

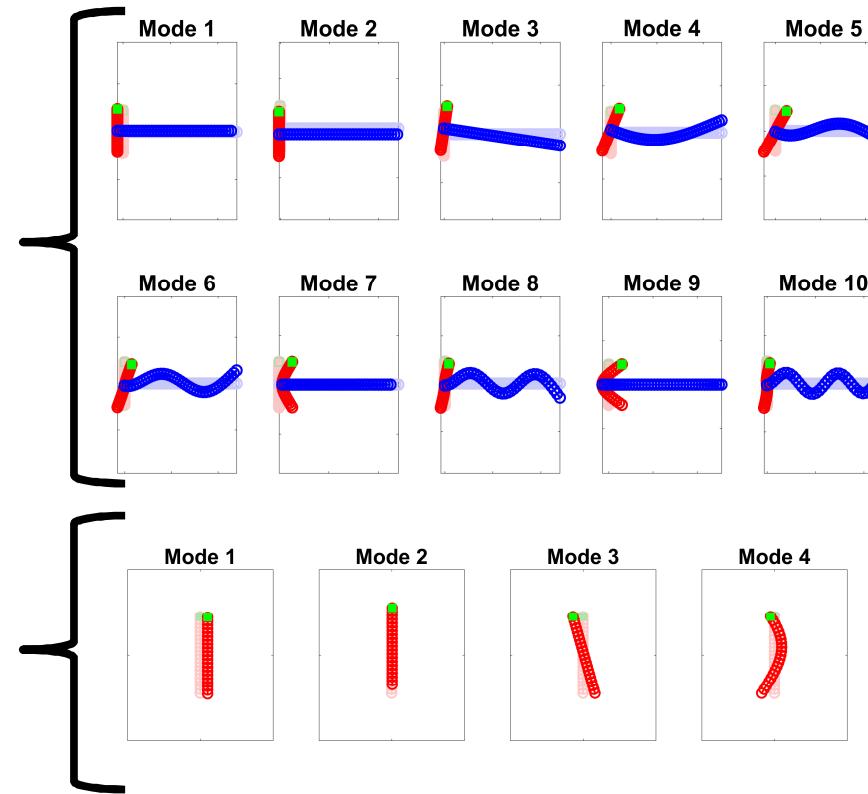
Transducer modeled as rigid body

$$\mathbf{M}_T = \begin{bmatrix} m_T & 0 & -m_T y_T \\ 0 & m_T/2 & 0 \\ -m_T y_T & 0 & I_{zz} \end{bmatrix}$$

Correction approach requires a sufficient number of analytical modes



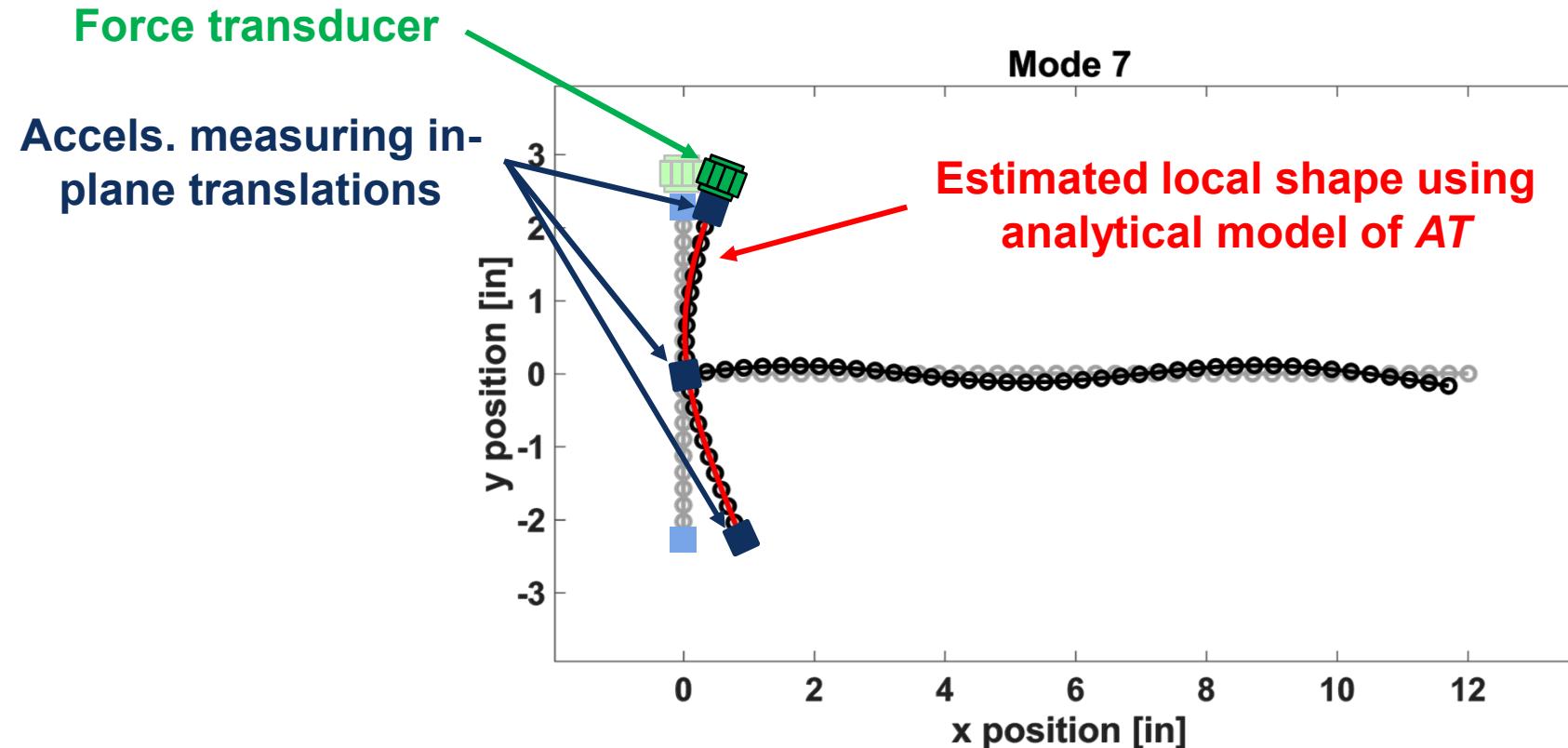
Simulated experimental system with transducer, CT



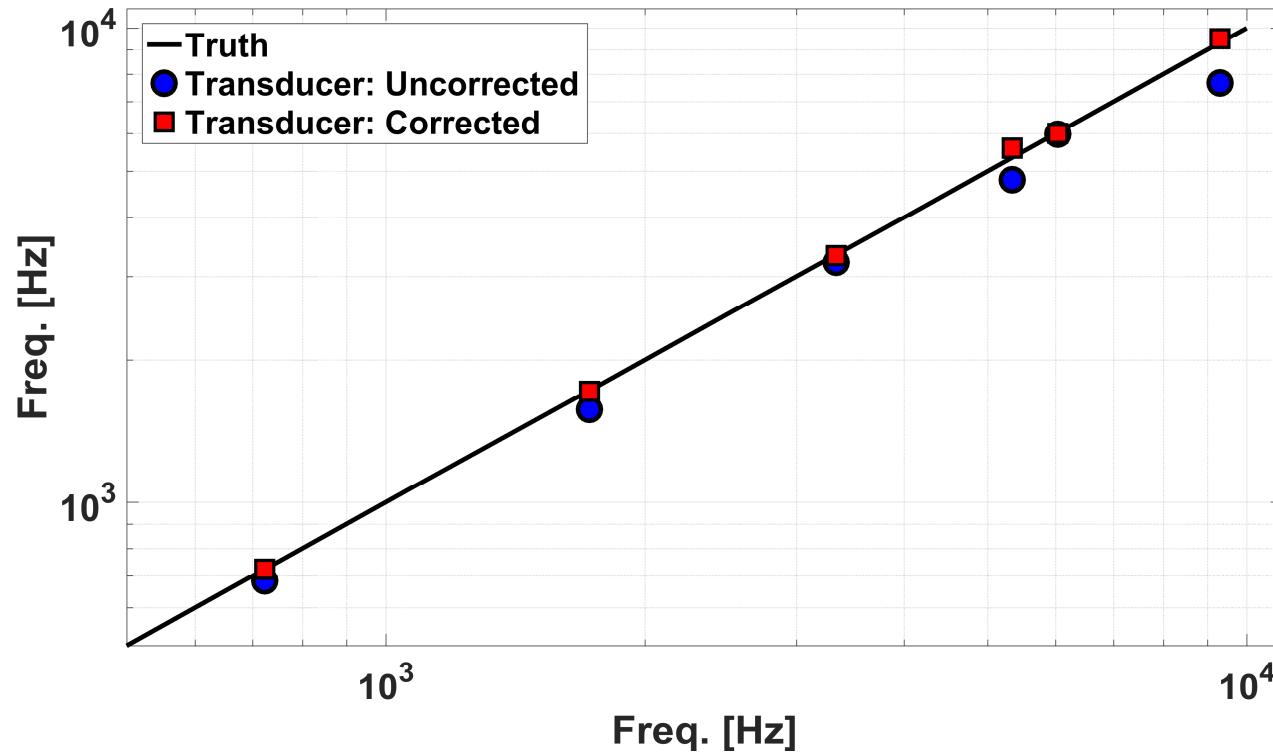
Analytical model of local system with transducer, AT

First 4 analytical modes (3 rigid body, 1 bending) sufficiently capture the full system modal motion

The analytical model of beam AT estimates full motion at transducer connection point.

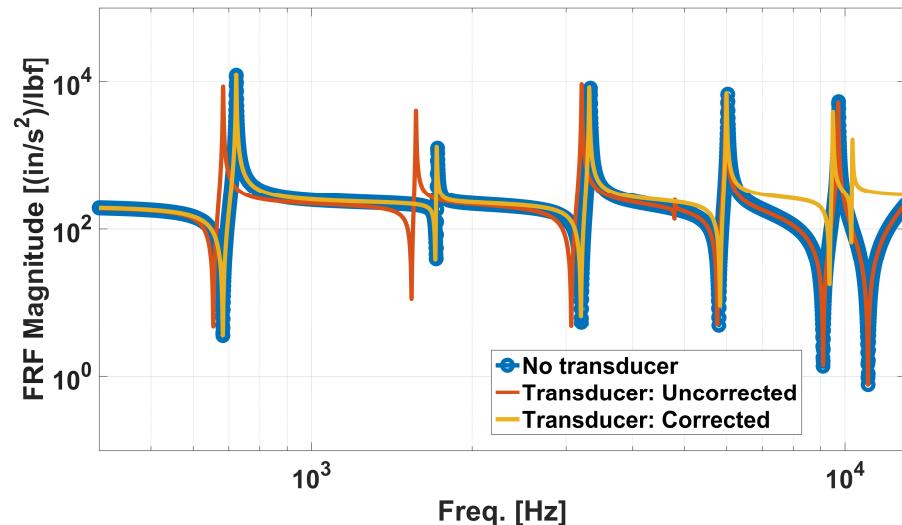


The corrected frequencies show much better agreement with the truth case.



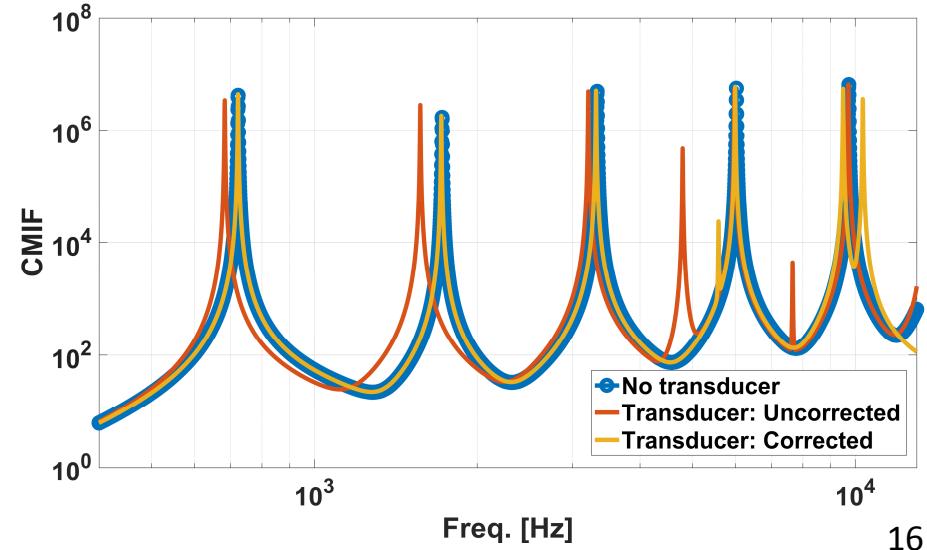
The corrected FRF and CMIF show much better agreement to the truth case.

Drive-point FRF



Only modes with frequencies < 10 kHz were “measured” -- Leads to the deviations in the FRF starting around 8 kHz

Complex Modal Indicator Function



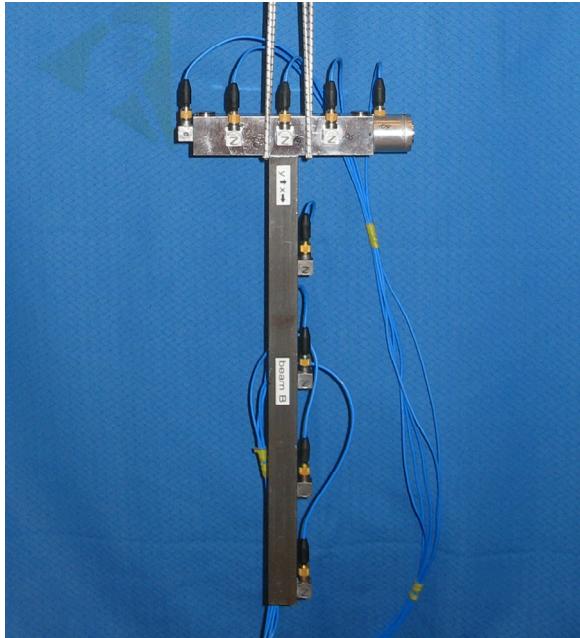
The spurious peaks in the CMIF were mostly removed – increasing measurement points showed further improvements

Presentation Outline

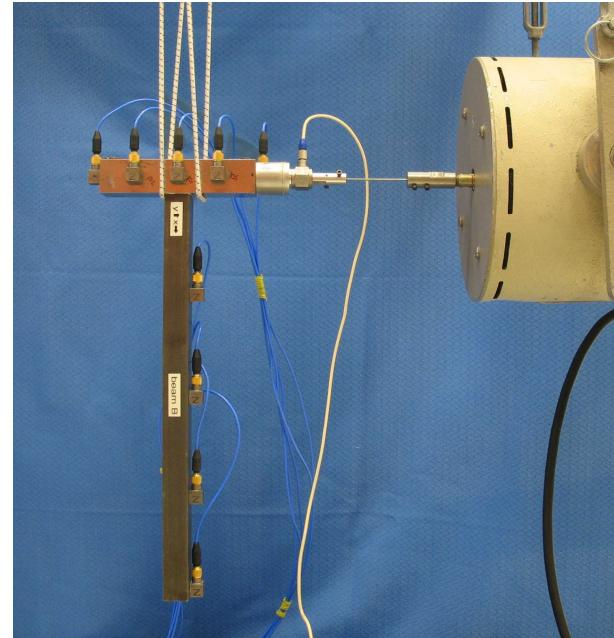


- Theoretical Development
- Numerical Results
- **Experimental Results**
- Conclusions

Hammer test with no attached transducer used for the truth baseline.



Hammer Test Setup

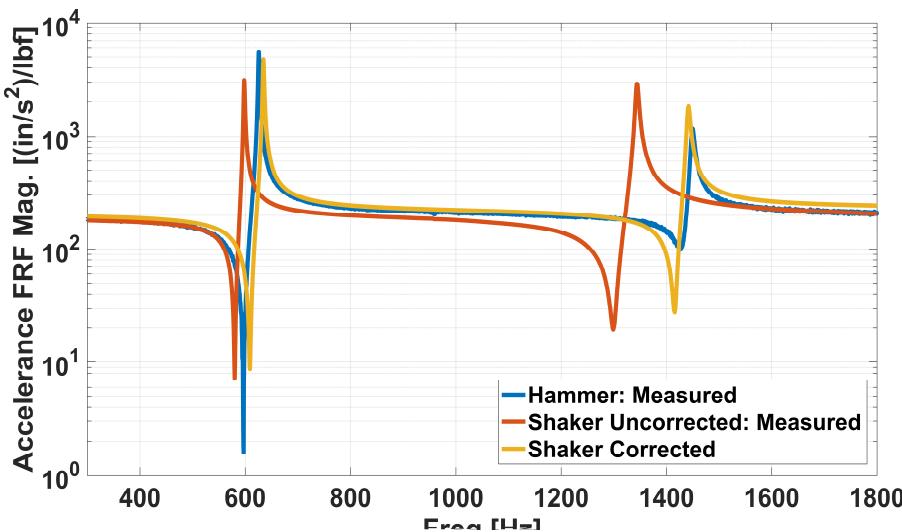


Shaker Test Setup

Excitation up to ~ 3 kHz provided information for the first 2 flexible modes

The corrected FRFs shows much better agreement to the truth case.

Drive-point FRF



Correction procedure improves modal frequency estimates:

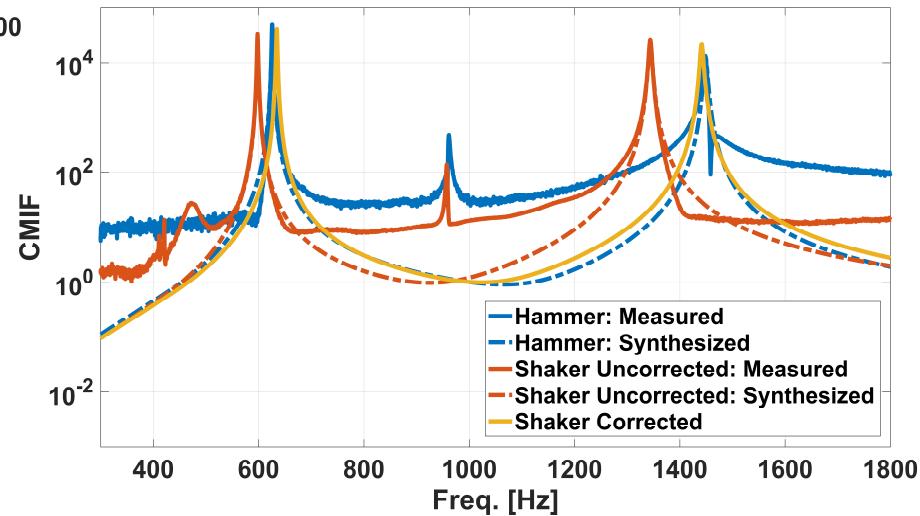
Uncorrected freq. error:

-4.44%, -7.25%

Corrected freq. error:

1.39%, -0.51%

Complex Modal Indicator Function



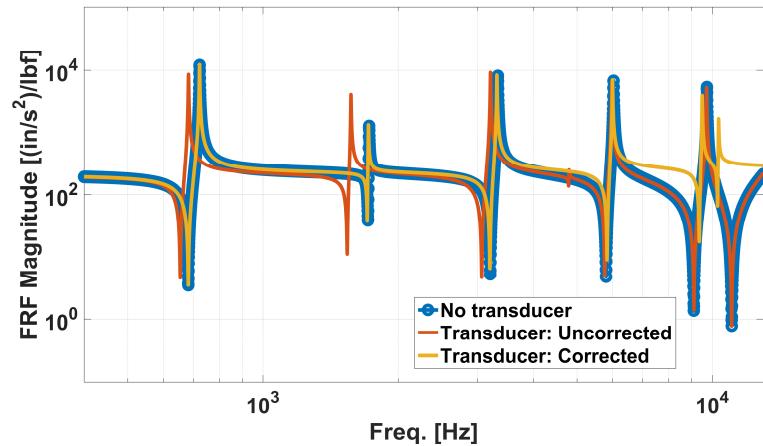
Presentation Outline



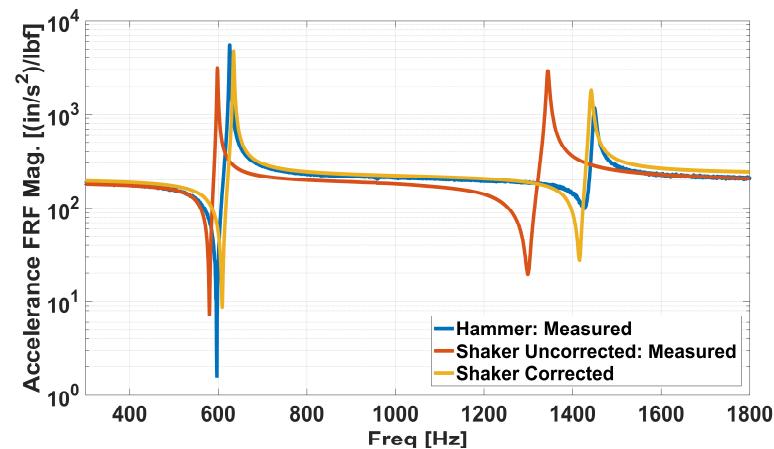
- Theoretical Development
- Numerical Results
- Experimental Results
- **Conclusions**

Successfully developed a method to remove force transducer effects during shaker testing.

- **A local analytical model of experimental system facilitated estimation of connection point motion**
 - Removed requirement for measurements of rotation to correct for mass moment of inertia
- **Numerical results show the correction method provides excellent agreement with the case of no attached transducer.**
- **Experimental tests provide validation for the method, especially for correcting the freqs.**



Numerical Results



Experimental Results