

# Measuring inertial forces with ultracold neutral atoms

Grant Biedermann

Naval Postgraduate School  
April 20, 2018



**LDRD**

Laboratory Directed Research and Development

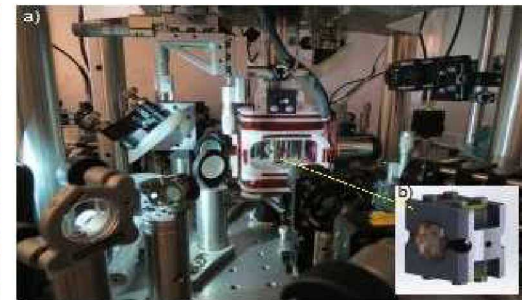
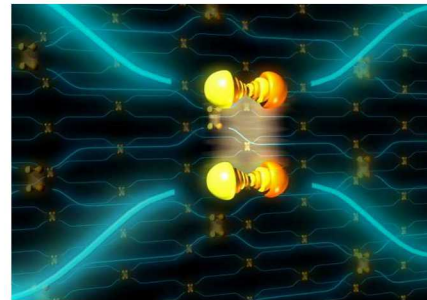


This work was supported by the Laboratory Directed Research and Development program at Sandia National Laboratories



Sandia National Laboratories is a multimission laboratory managed and operated by National Technology and Engineering Solutions of Sandia LLC, a wholly owned subsidiary of Honeywell International Inc. for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-NA0003525.

SAND2018-4252PE



SAND2018-XXXX

# Outline

- Precision sensing with atom interferometers
- Sandia atom interferometer efforts
  - Single atom interferometer
  - High data rate atom interferometer
- Next generation in sensing
  - Rydberg-dressed physics and entanglement

# Inertial Navigation

Basic idea:

$$x(t) = \iint a(t) dt$$

Uncertainty in  $a$  drives uncertainty in  $x$

However, must include gyroscope to get pointing information



U.S. Navy photo by Mass Communication Specialist 1st Class Woody Paschall—CNBC

Tactical, MEMS

Navigation, HRG

Strategic, ESG

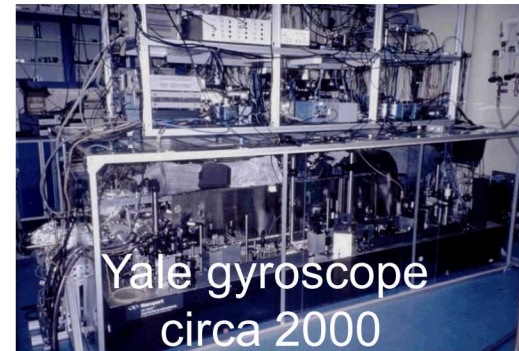
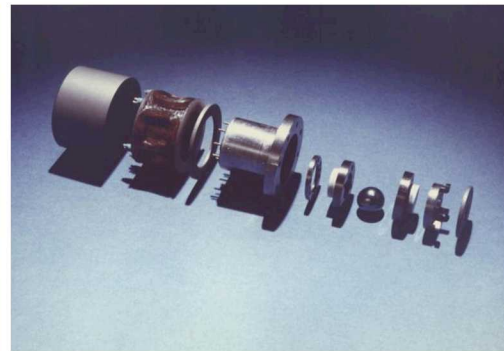
Atomic



Analog Devices  
ADIS16385



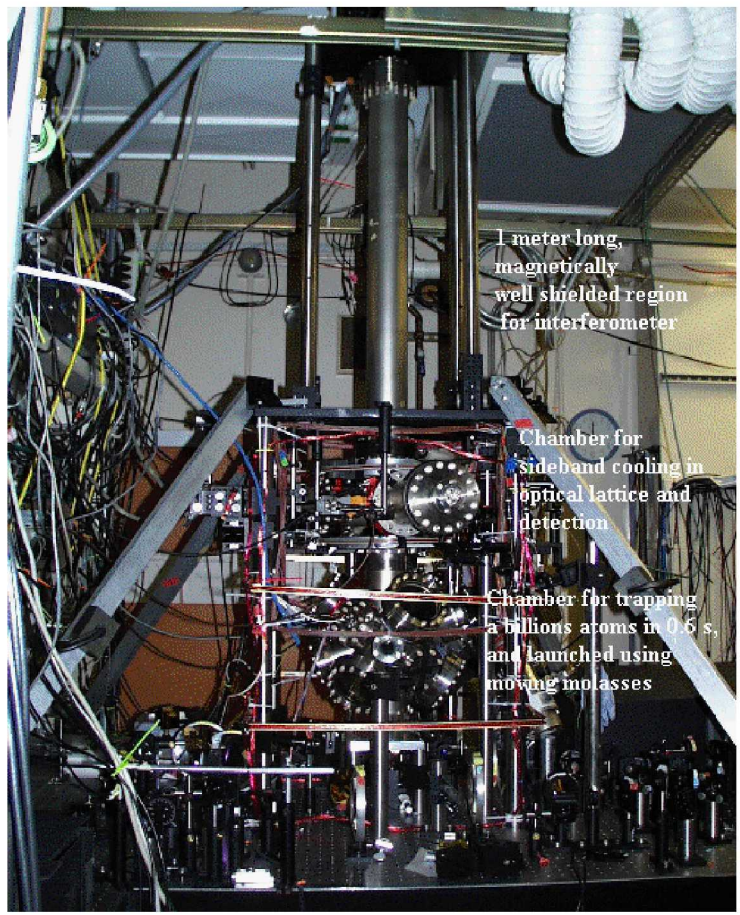
Honeywell  
HG9900



Yale gyroscope  
circa 2000

stability

# Light-pulse atom interferometers



Stanford atom interferometer gravimeter  
Steven Chu, 1990's



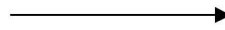
Photograph by Charles Watkins, White House photographer  
Text by GB

Superb gyroscopes and gravity gradiometers demonstrated as well

# Atom Interferometer Sensors

- AI Gravimeter

- A. Peters *et al*, Metrologia **38**, (2001)
- Capable of 20 nano-g in 1.3 s
- Recent tower work infers pico-g level per shot



10 m tower  
Mark Kasevich  
Stanford

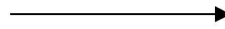


- AI Gyro

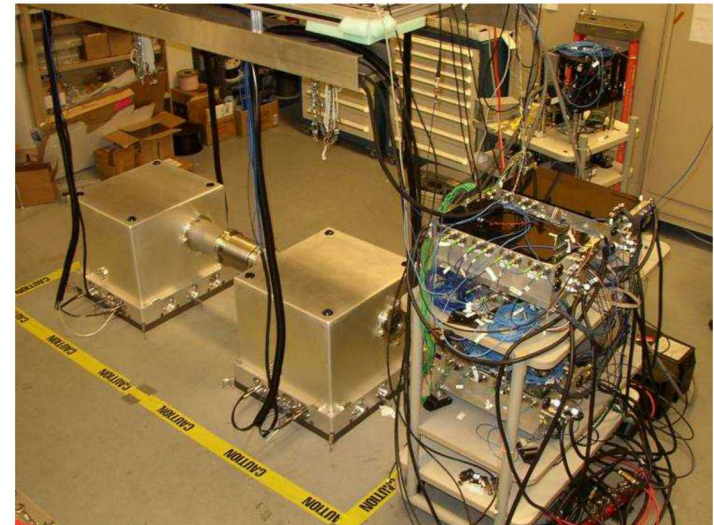
- D. Durfee *et al*, PRL **97**, (2006)
- Short-term stability  $\sim 3\mu\text{Deg}/\text{hr}^{1/2}$
- High-accuracy navigation grade
  - $< 70\mu\text{Deg}/\text{hr}$  bias stability

- AI Gradiometer

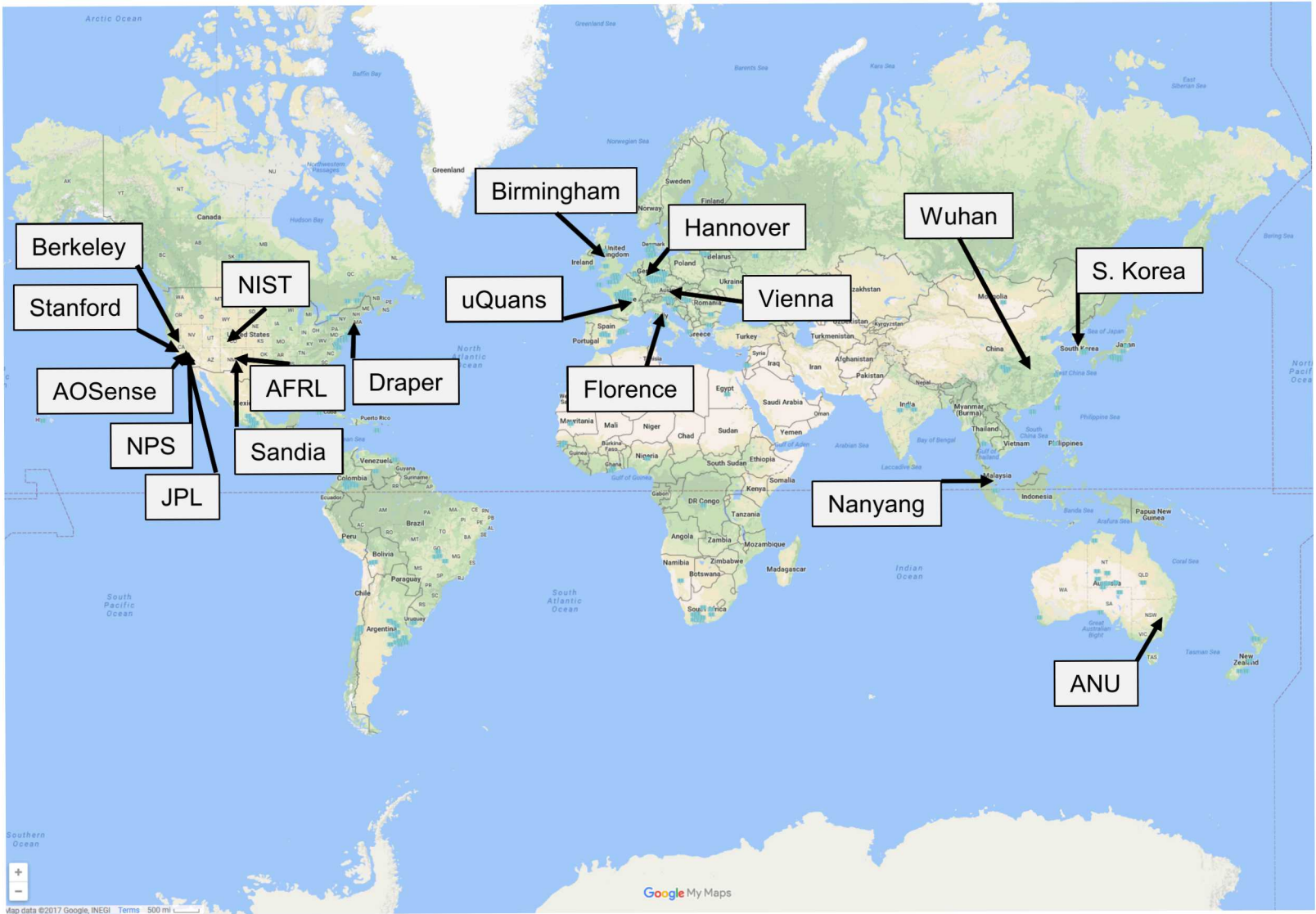
- $\Delta a = 4 \text{ nano-g Hz}^{-1/2}$ , Biedermann, PRA 2015
- Yale, 2002 (Fixler PhD thesis); Fixler, Science, 2007.
- More recent: Florence, Nature, 2014
- Precision  $\delta G/G \sim 7.7 \times 10^{-5}$
- Systematic uncertainty  $6.2 \times 10^{-5}$
- CODATA 2014 relative uncertainty:  $4.7 \times 10^{-5}$



Stanford gradiometer—G.B. and many others

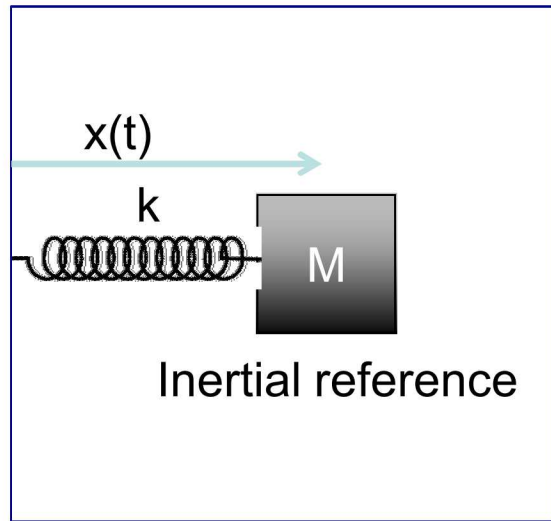


# An incomplete map of AI efforts



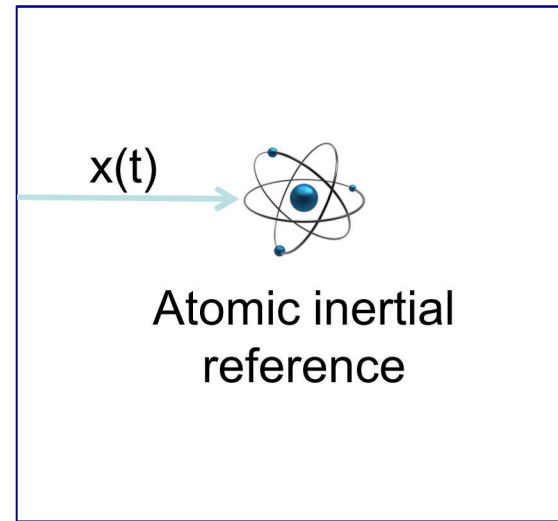
# Inertial measurements

Mass-spring accelerometer



platform

Free-fall accelerometer



platform

## Atom interferometers

- Inertial reference: laser-cooled atoms falling freely
- Positioning: laser ranging
- Truly accurate, turn-on calibration

# Atom optics, simplified

Use atom to stroboscopically measure lateral position (optical phase) at three equal-spaced points in time

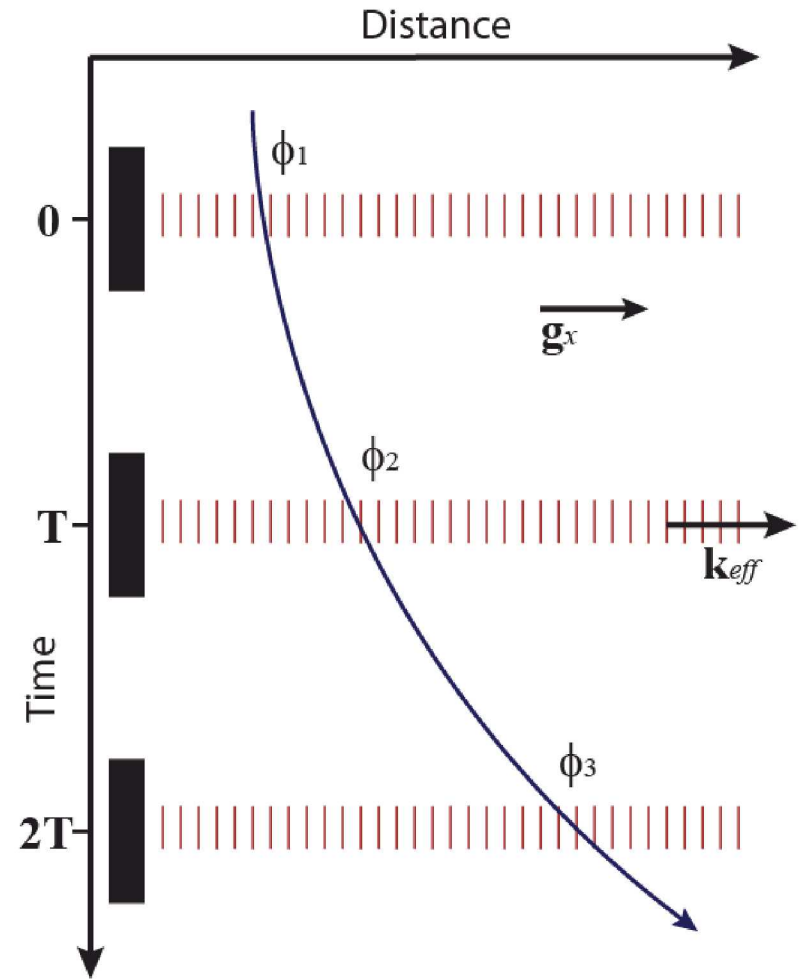
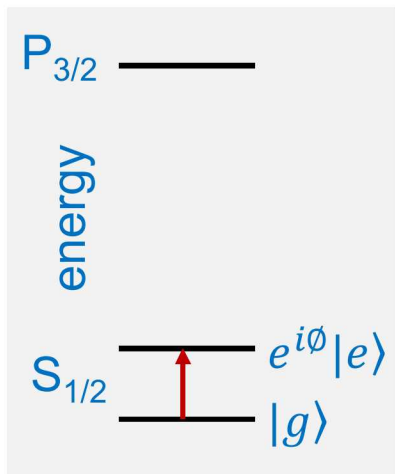
$$\varphi_i = k_{eff} \cdot x_i$$

calculate curvature via finite difference method

$$\Delta\varphi = \varphi_1 - 2\varphi_2 + \varphi_3$$

...and then

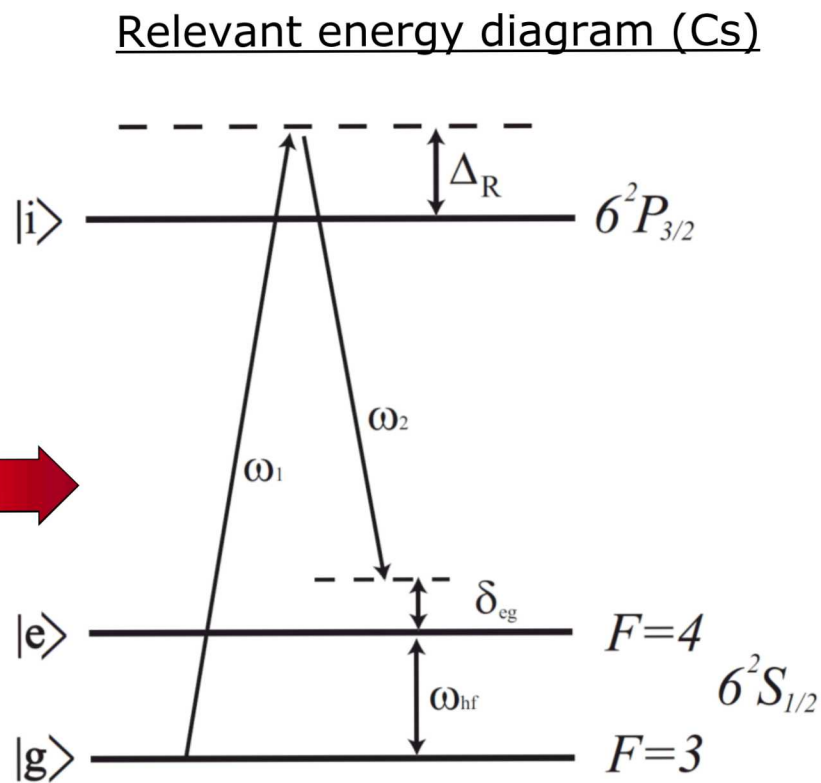
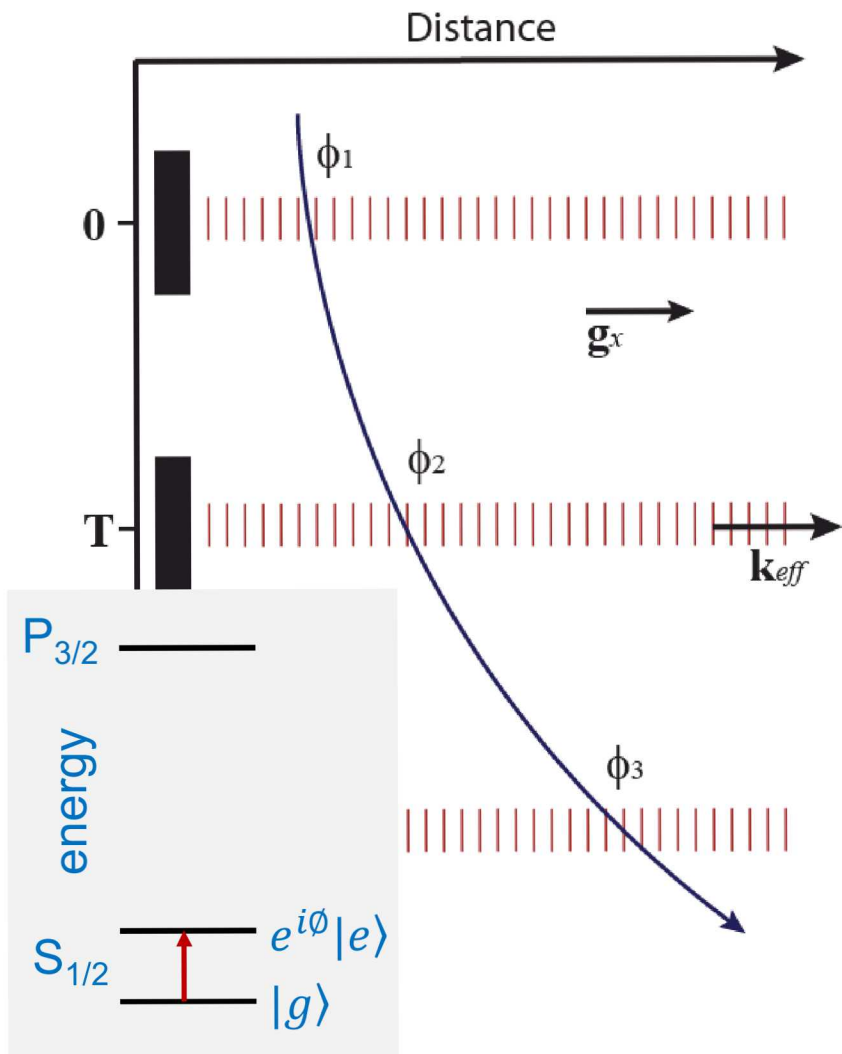
$$g_x = \frac{\Delta\varphi}{k_{eff} T^2}$$



$$P_{\uparrow} = 1/2(1 - \cos\Delta\phi)$$

...think of  $S_{1/2}$  states as pristine oscillator with phase memory

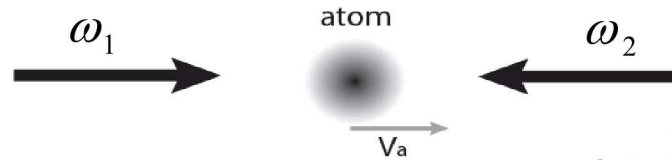
# Laser requirements: atom optics



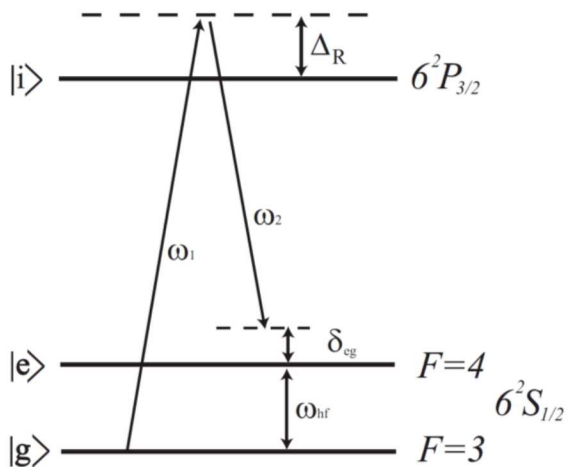
- Phase coherent laser tones separated by 6.8 GHz

# Stimulated Raman transition

## Laser configuration



### Relevant energy diagram (Cs)



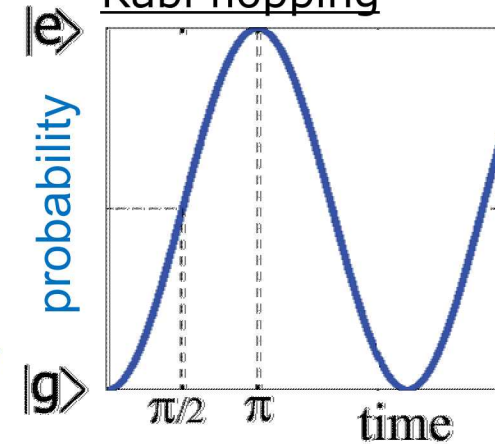
### Effective 2-level system

$$H'_{int} = -\hat{\mu}_{eff} \cdot \hat{E}_{eff} e^{i(k_{eff} \cdot \hat{x} - \omega_{hf} t)}$$

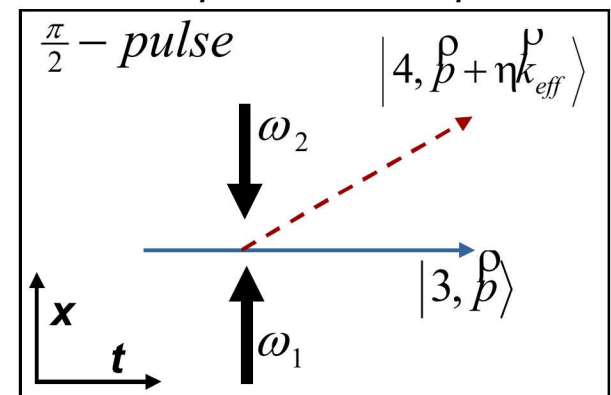
$$\vec{k}_{eff} = \vec{k}_1 - \vec{k}_2 \approx 2\vec{k}_1$$

Valid for  $\Omega_{eff} \ll \Delta$

### Rabi flopping

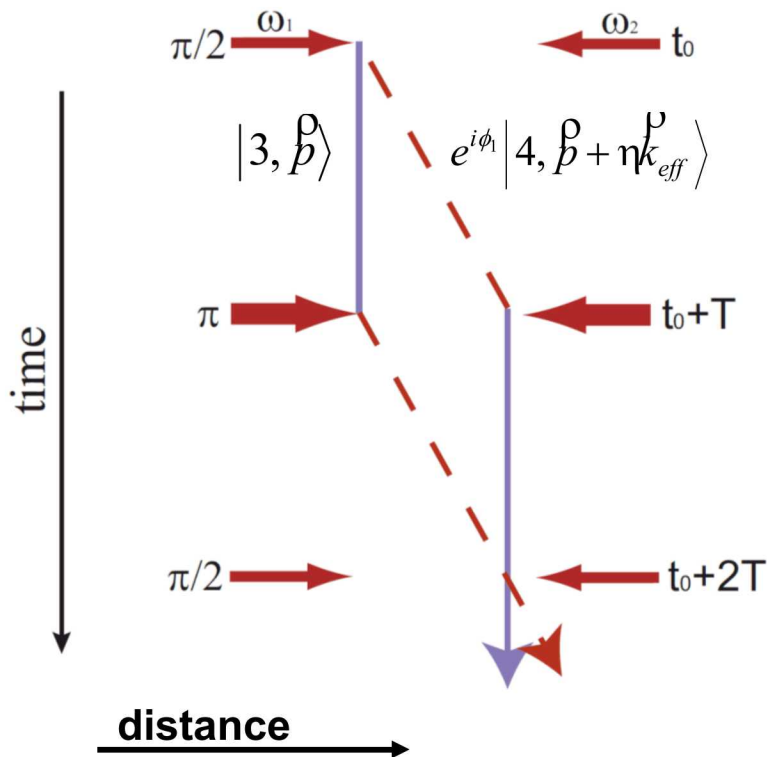


### Atom optics beamsplitter



# AI phase shift

## Interferometer Recoil diagram



## Transition rules

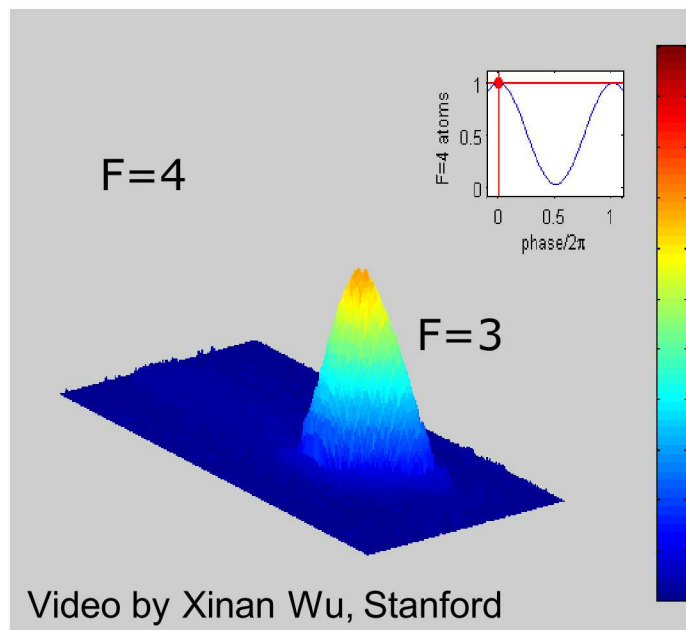
$$\left. \begin{aligned} |3, \mathbf{p}\rangle &\rightarrow e^{i\phi} |4, \mathbf{p} + \eta \mathbf{k}_{eff}^{\mathbf{p}}\rangle \\ |4, \mathbf{p} + \eta \mathbf{k}_{eff}^{\mathbf{p}}\rangle &\rightarrow e^{-i\phi} |3, \mathbf{p}\rangle \end{aligned} \right\} \Delta\phi = \phi_1 - 2\phi_2 + \phi_3$$

$$\phi = \mathbf{k}_{eff}^{\mathbf{p}} \cdot \mathbf{x} = \mathbf{k}_{eff}^{\mathbf{p}} \cdot (\mathbf{v} T^2 - 2(\mathbf{v} \times \mathbf{\Omega}) T^2)$$

Interferometer transition probability

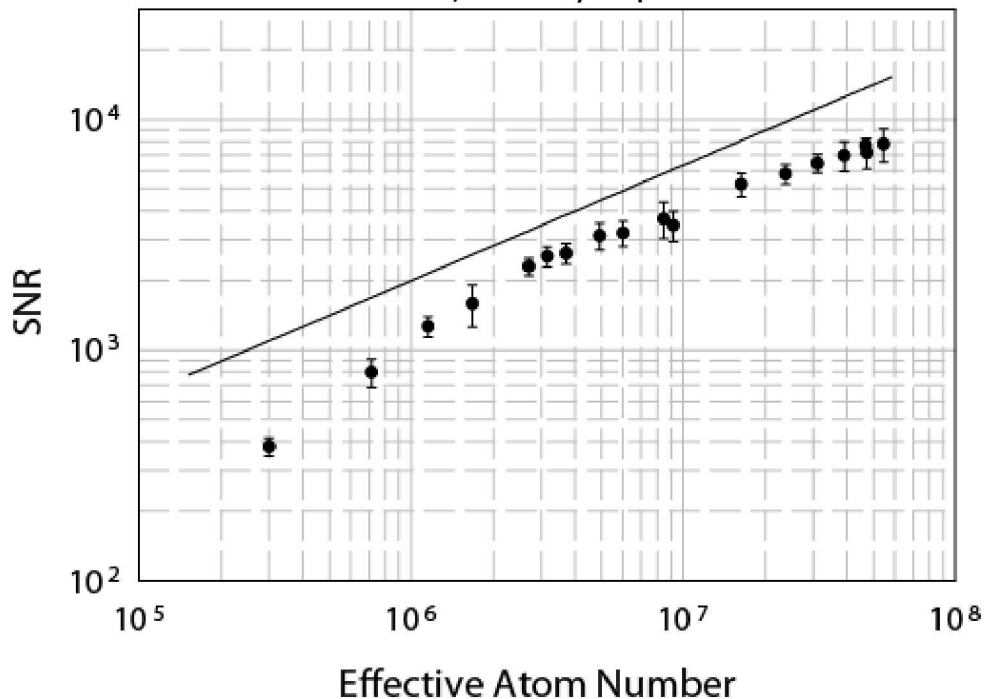
$$|\langle 4 | \Psi \rangle|^2 = \frac{1}{2}(1 - \cos \Delta\phi)$$

# State detection

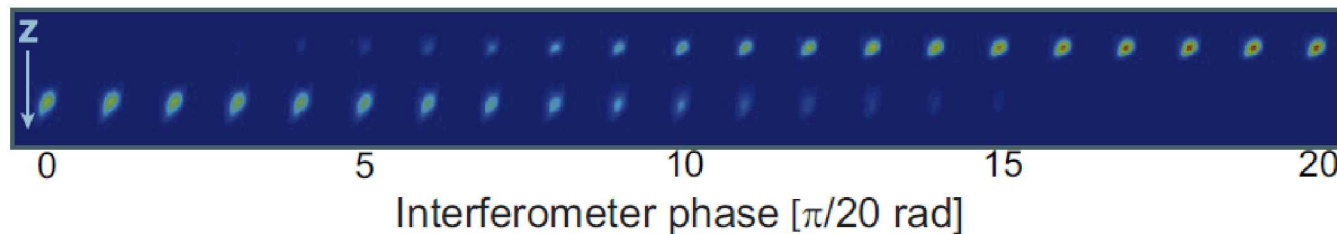


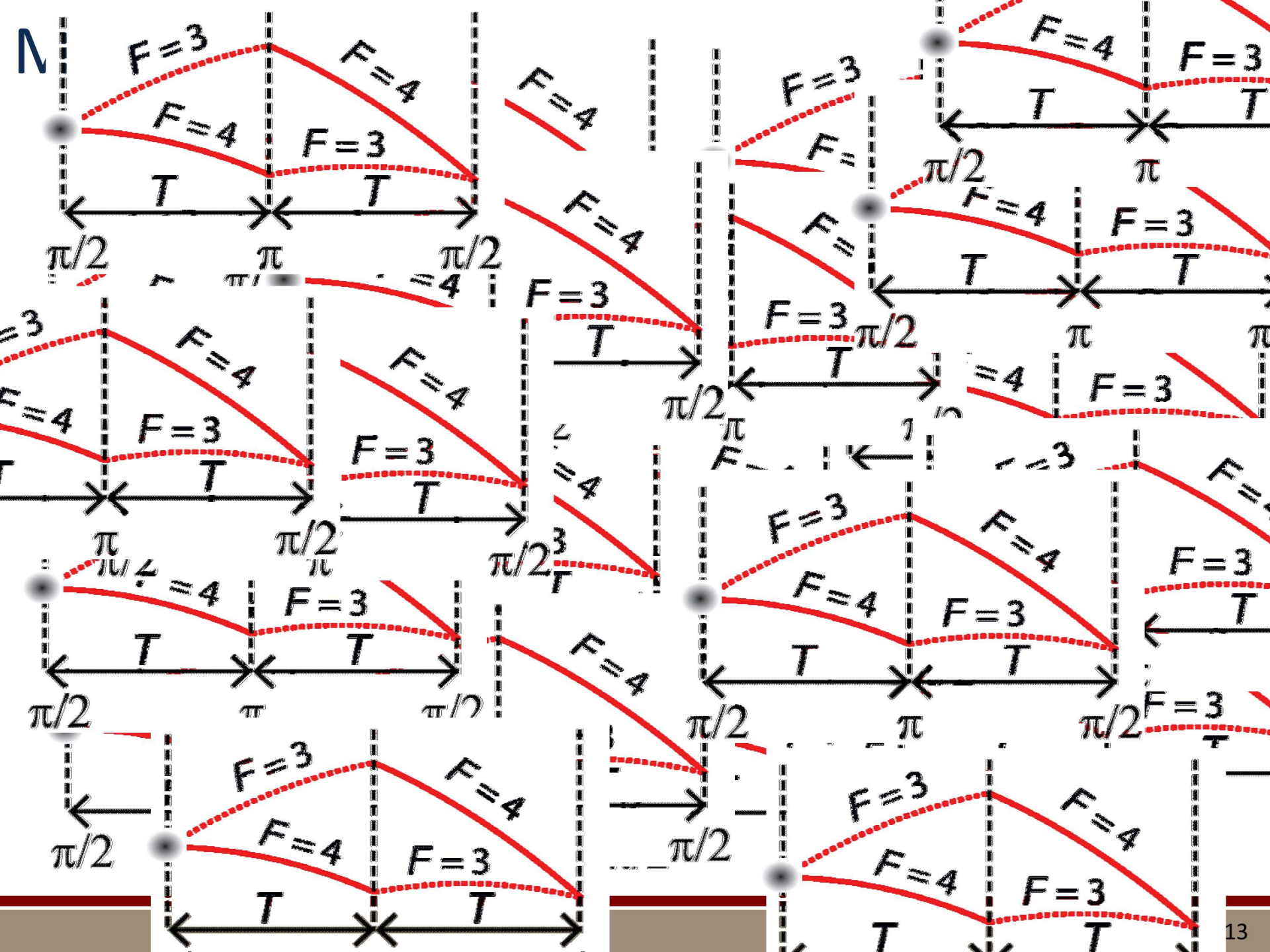
Radiation pressure spatially separates  $F=3$  and  $F=4$  atoms

Biedermann, et al., Opt. Lett. 2007

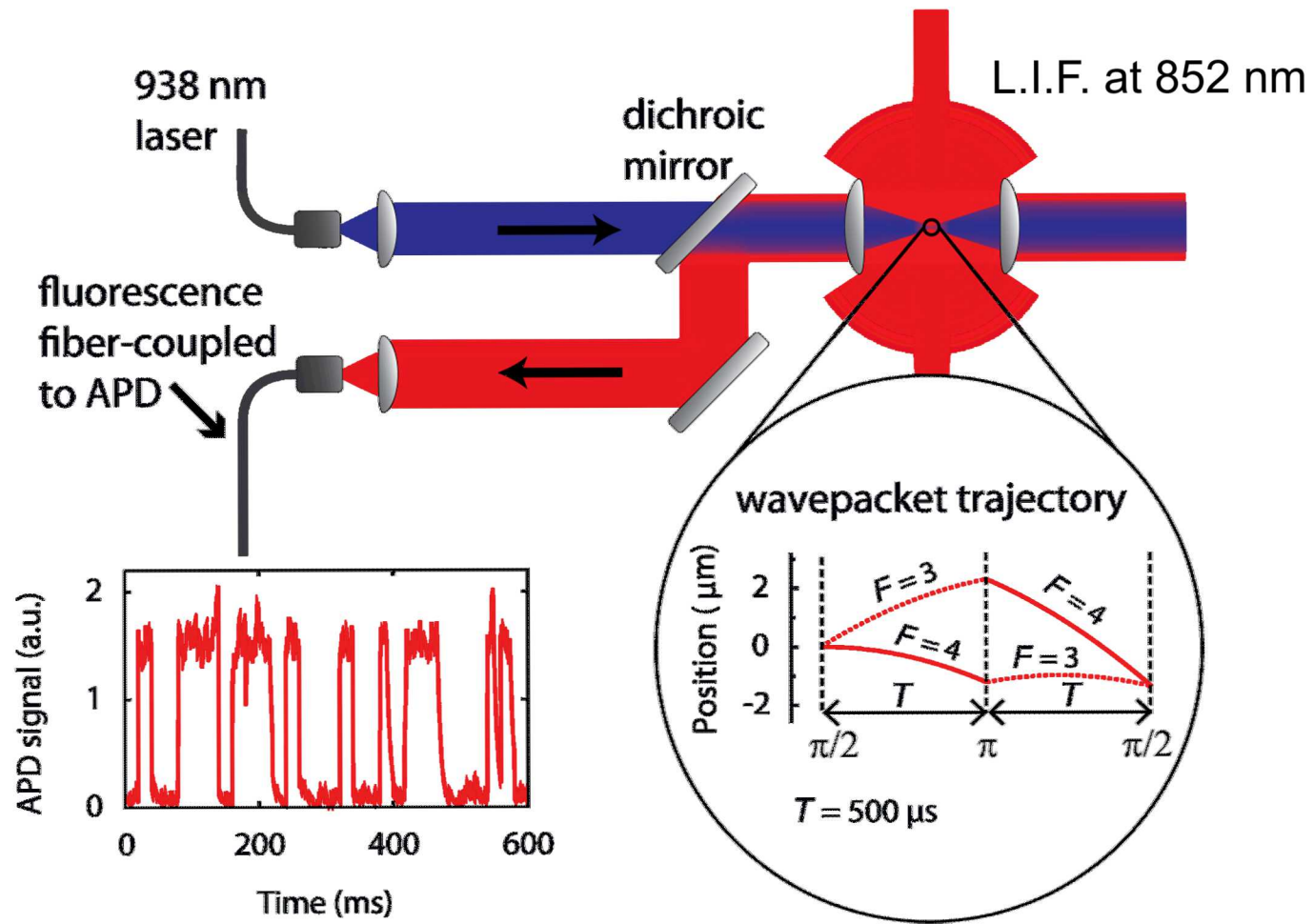


Achieves near atom shot-noise limited performance.

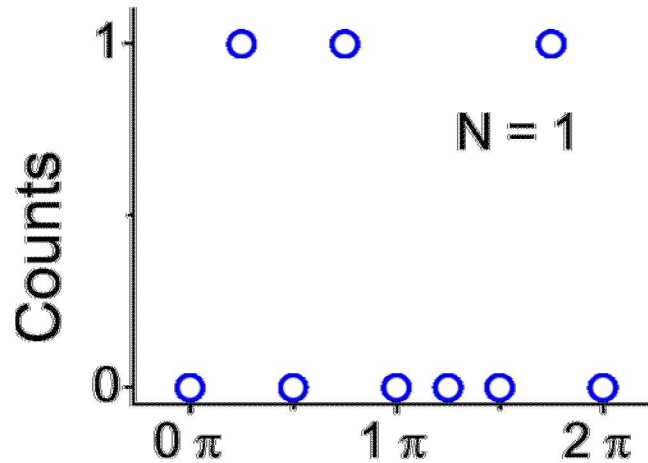




# Free-space matterwave interference with a single atom



# Building fringe one atom at a time

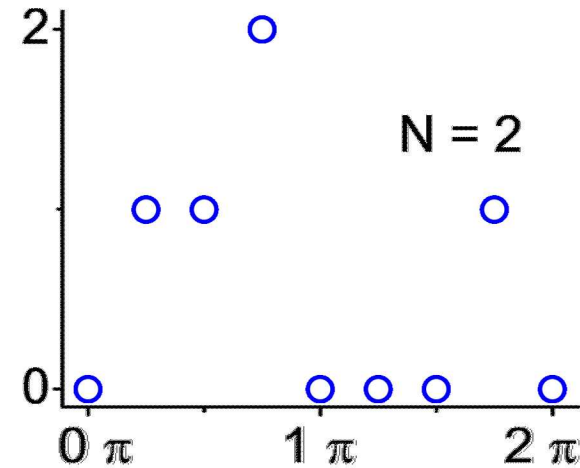
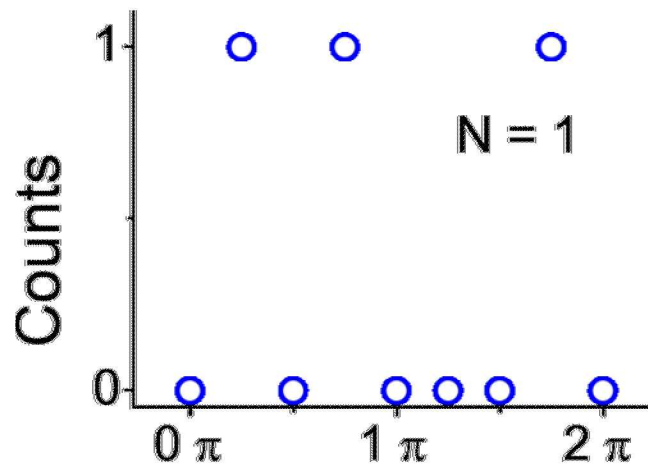


- 1 atom per phase through interferometer.

$$H'_{\text{int}} = -\hat{\mu}_{\text{eff}} \cdot \hat{E}_{\text{eff}} e^{i(k_{\text{eff}} \cdot \hat{x} - \omega_{\text{hf}} t)}$$

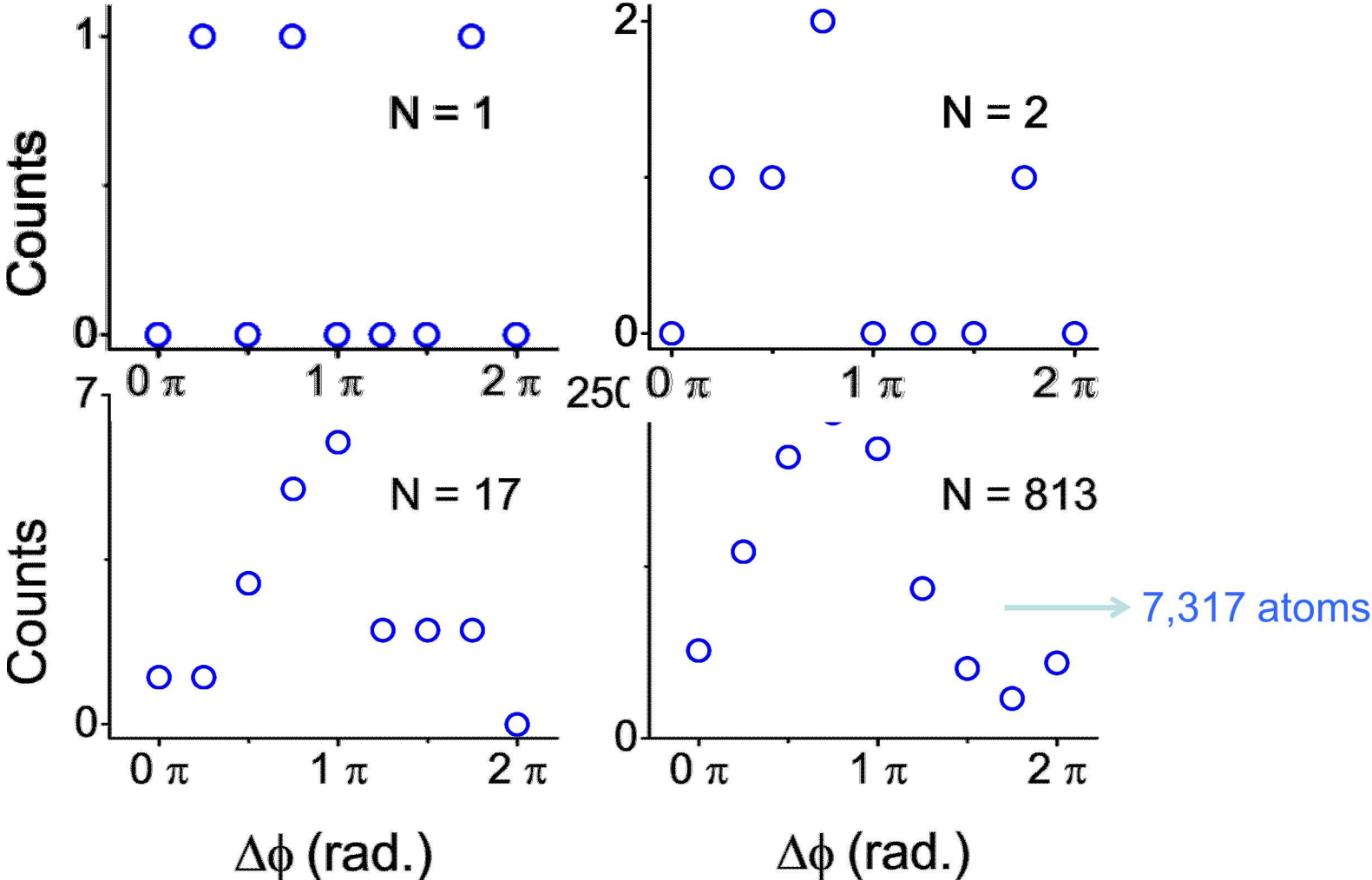
Scan laser phase

# Building fringe one atom at a time

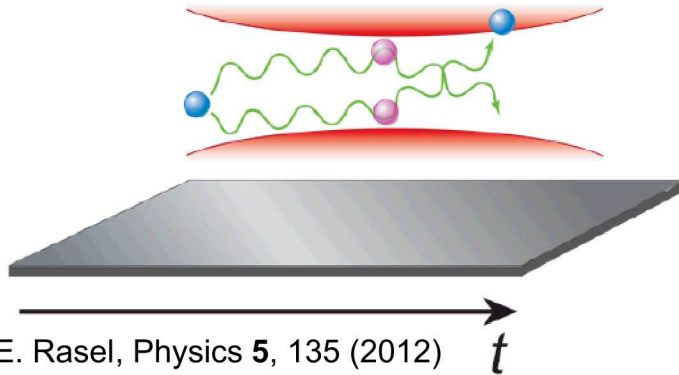


- 2 atoms per phase through interferometer

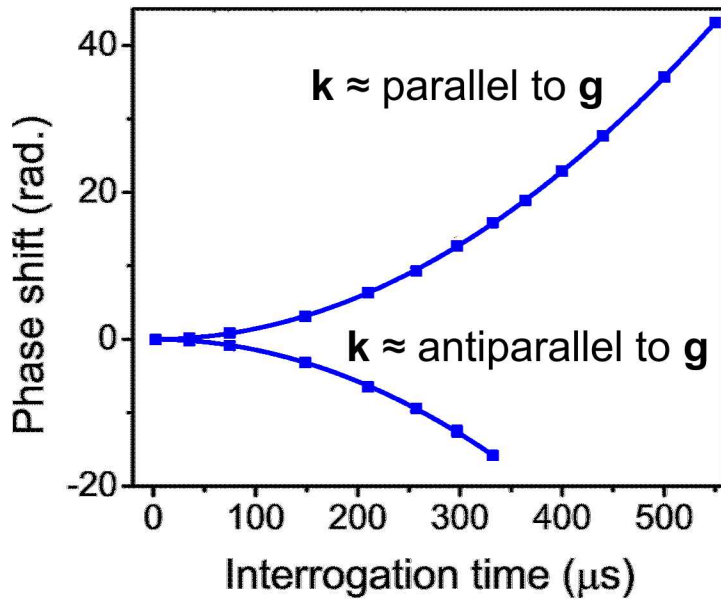
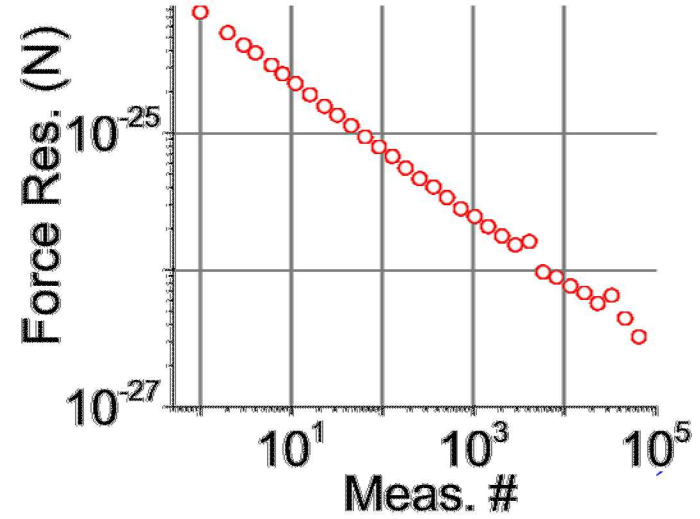
# Building fringe one atom at a time



# Force resolution of a single atom interferometer

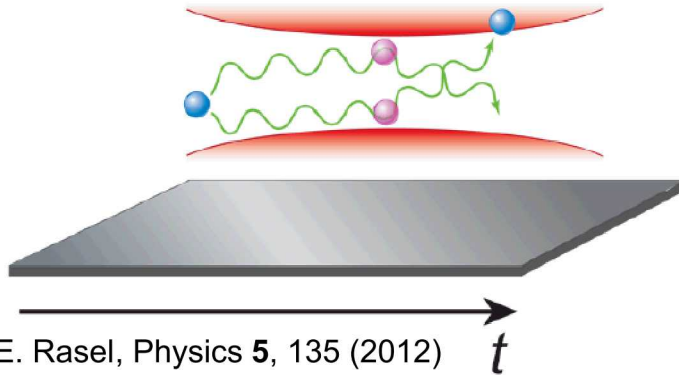


Parazzoli, et al., Phys. Rev. Lett. **109**, 230401 (2012)



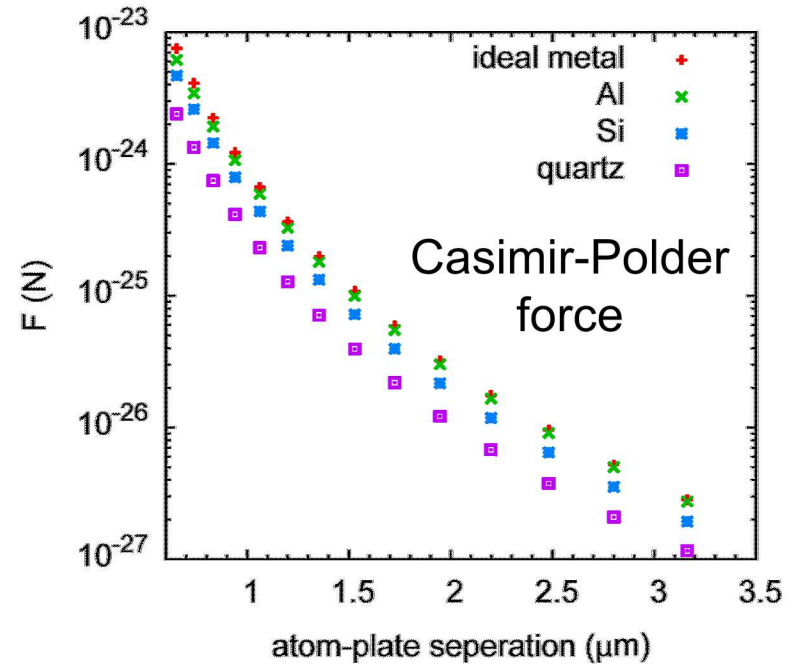
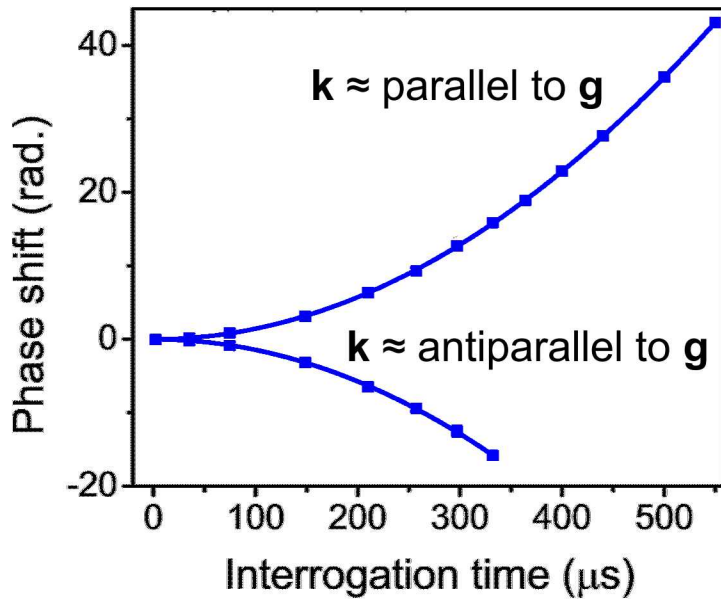
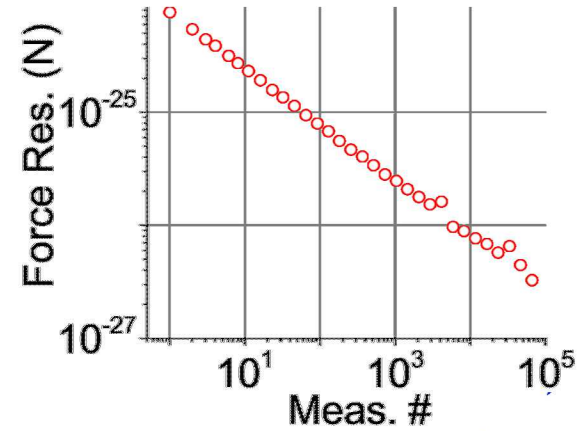
$$\Delta\Phi = k \cdot a T^2$$

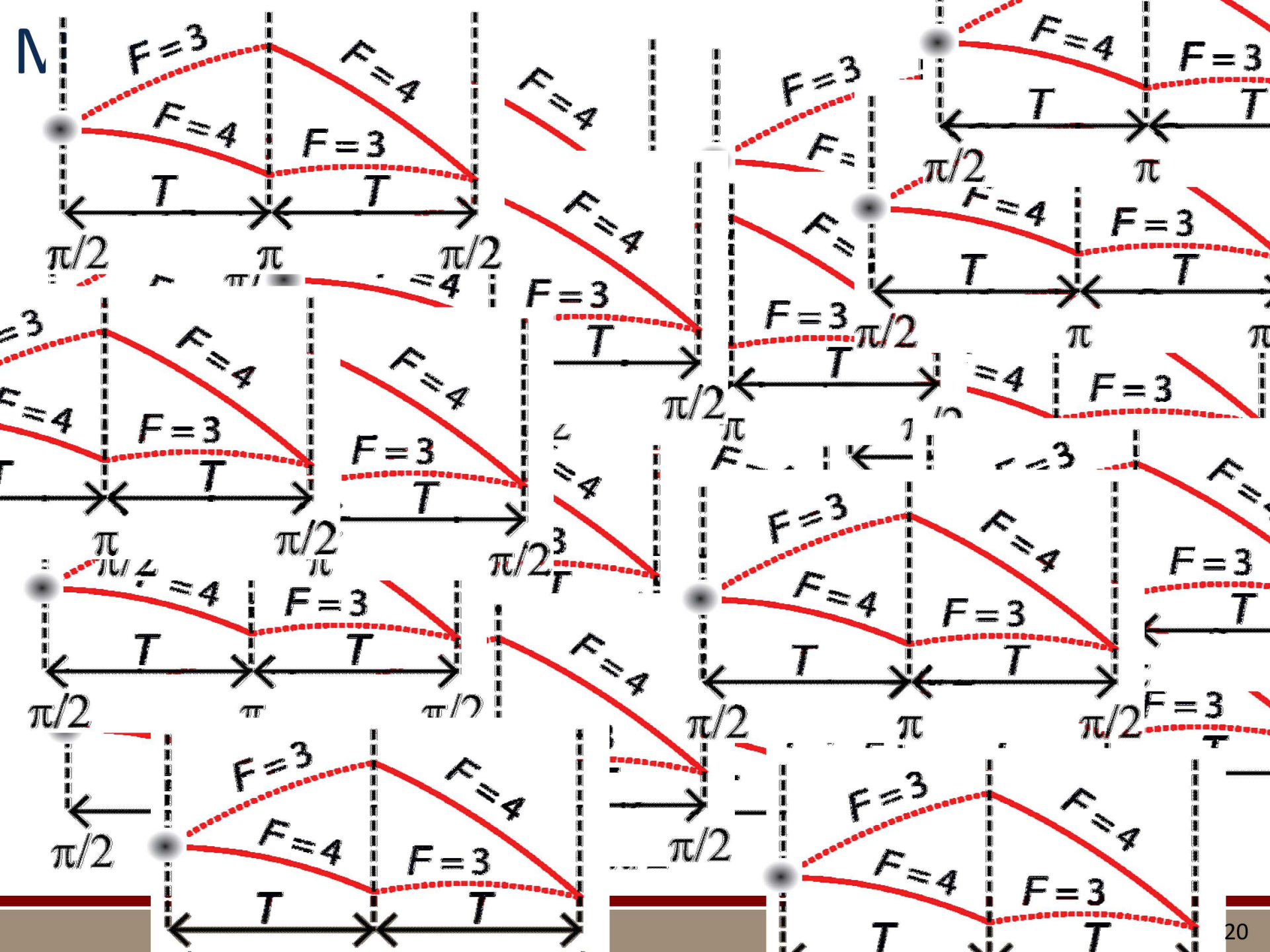
# Force resolution of a single atom interferometer



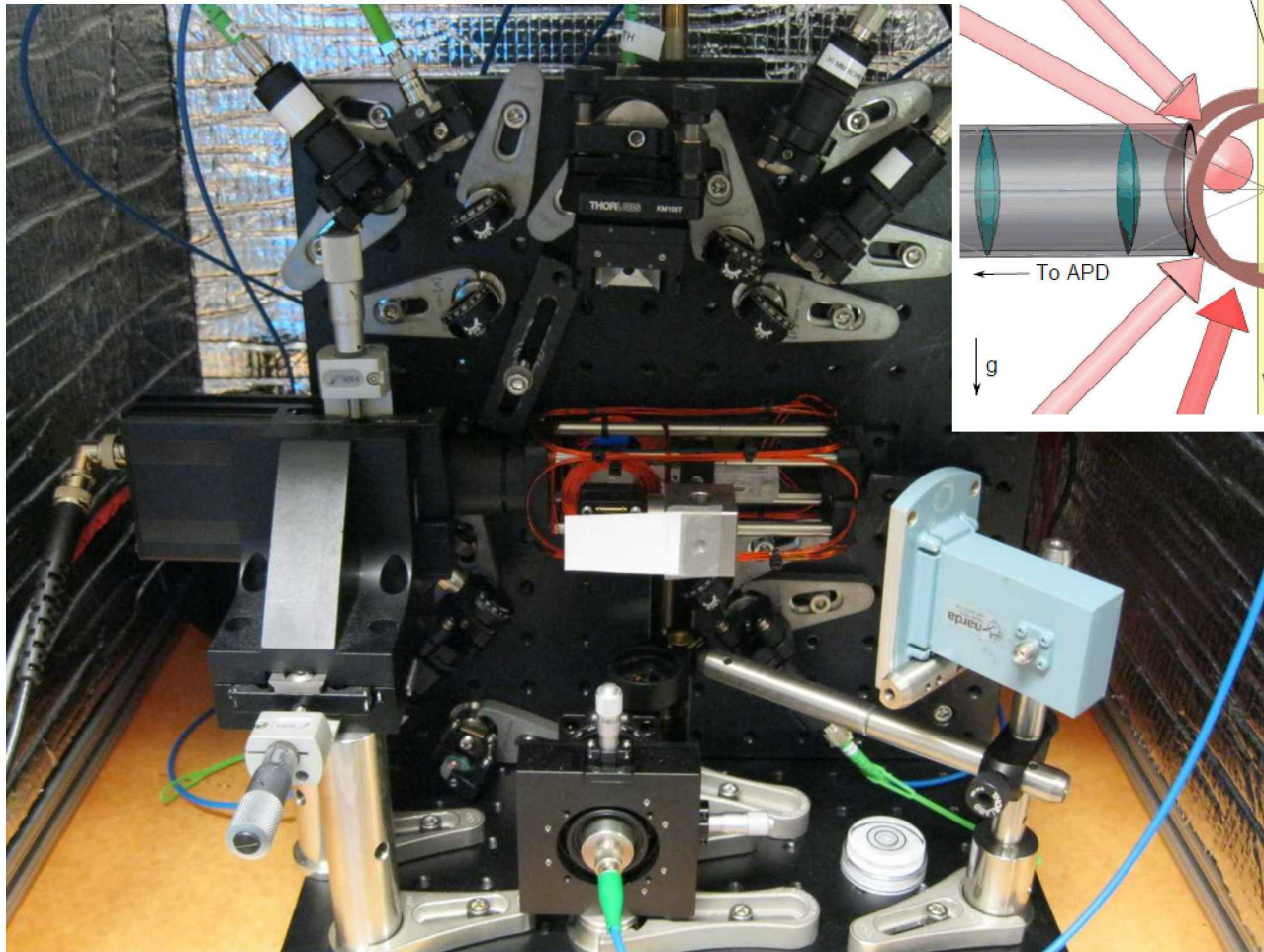
E. Rasel, Physics **5**, 135 (2012)

Parazzoli, et al., Phys. Rev. Lett. **109**, 230401 (2012)





# High-bandwidth interferometer

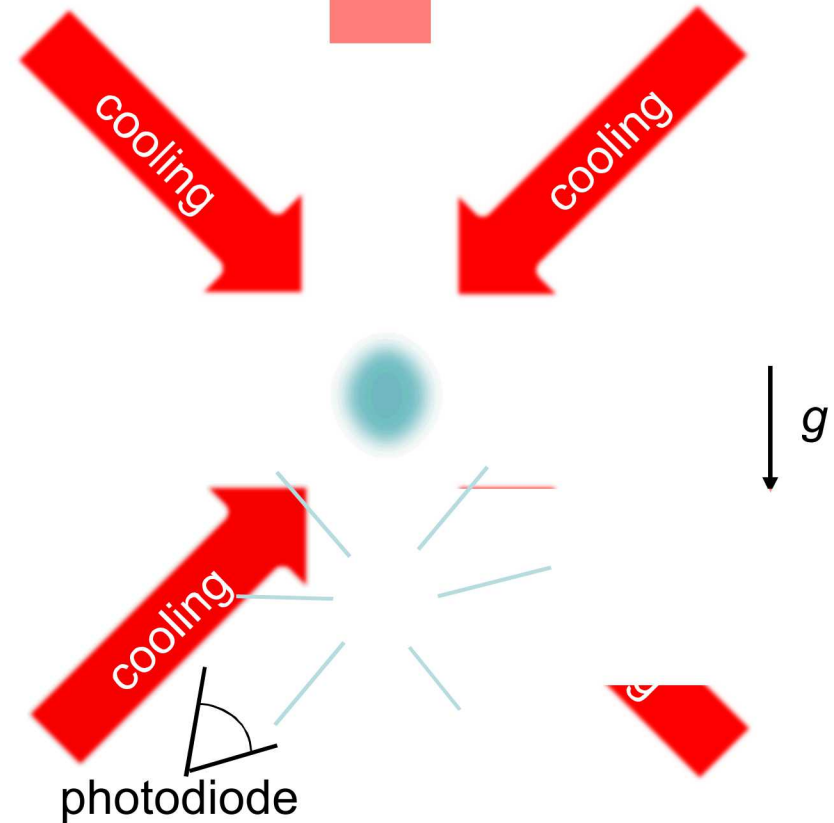


$600 \text{ ng}/\sqrt{\text{Hz}}$  at 50 Hz

# One measurement cycle

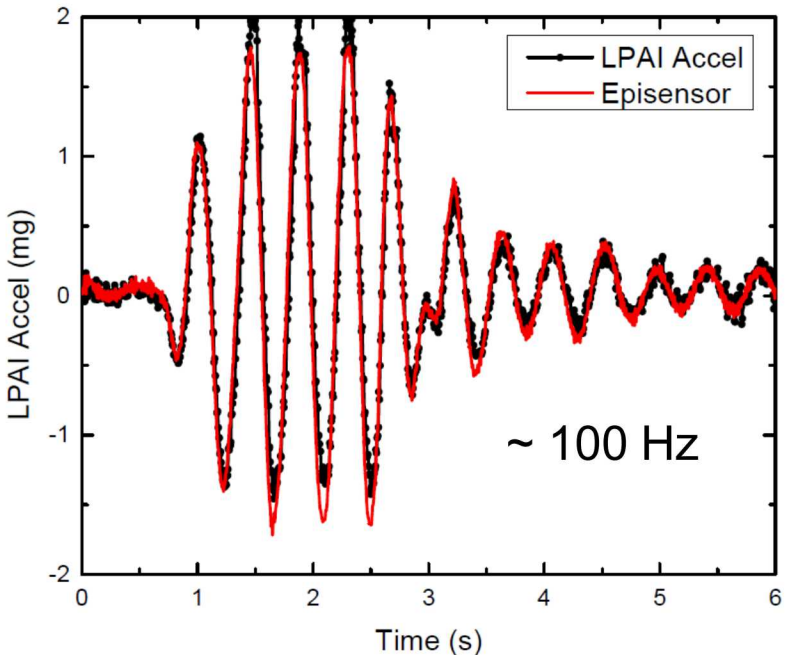
*Example,  $50 \text{ Hz}^{-1}$  cycle:*

- Laser cool  $10^6$  atoms (5 ms)
  - $T \approx 5 \mu\text{K}$
- Release atoms
- Raman pulse sequence (14 ms,  $T = 7 \text{ ms}$ )
- Detect
- Recapture (1.7 ms)

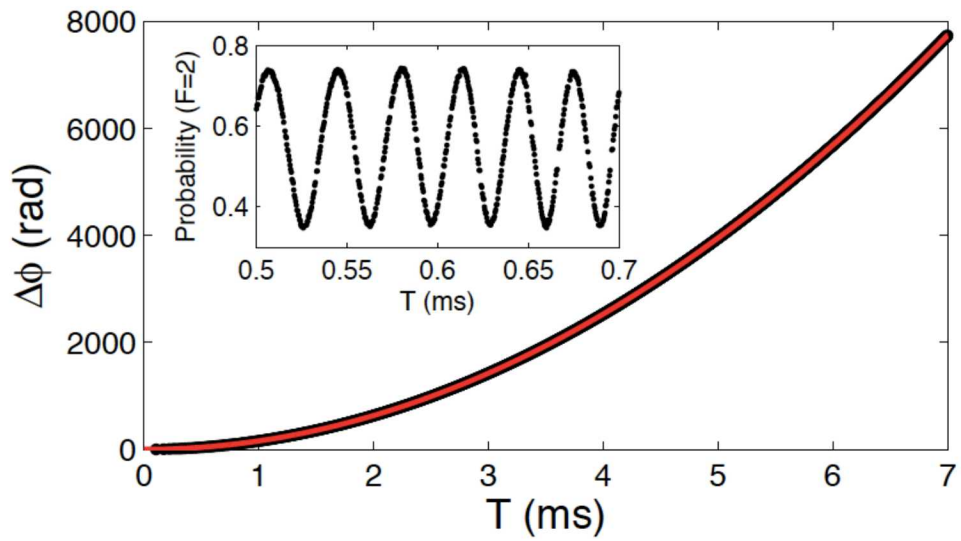


# High bandwidth interferometer results

Operating under dynamics



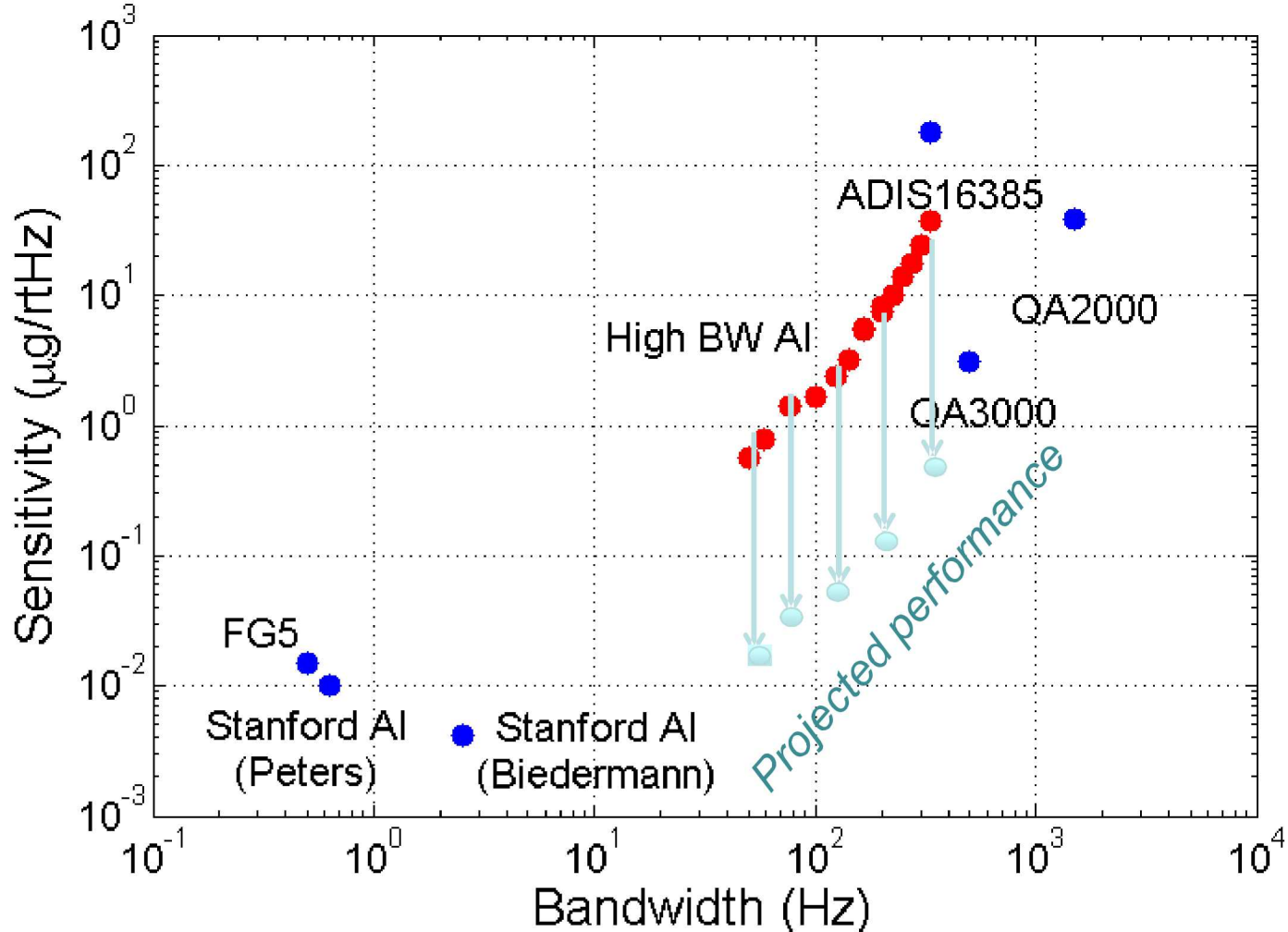
Measurement of g



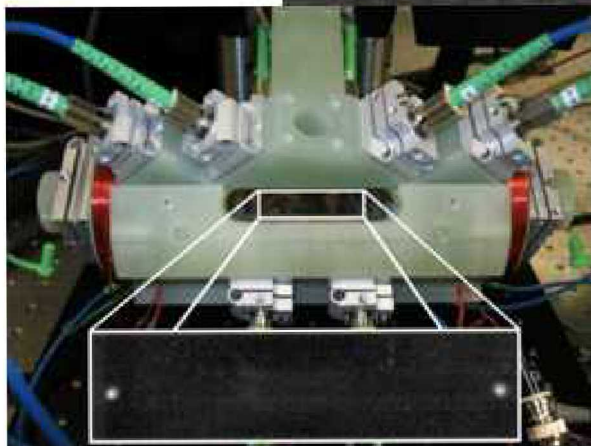
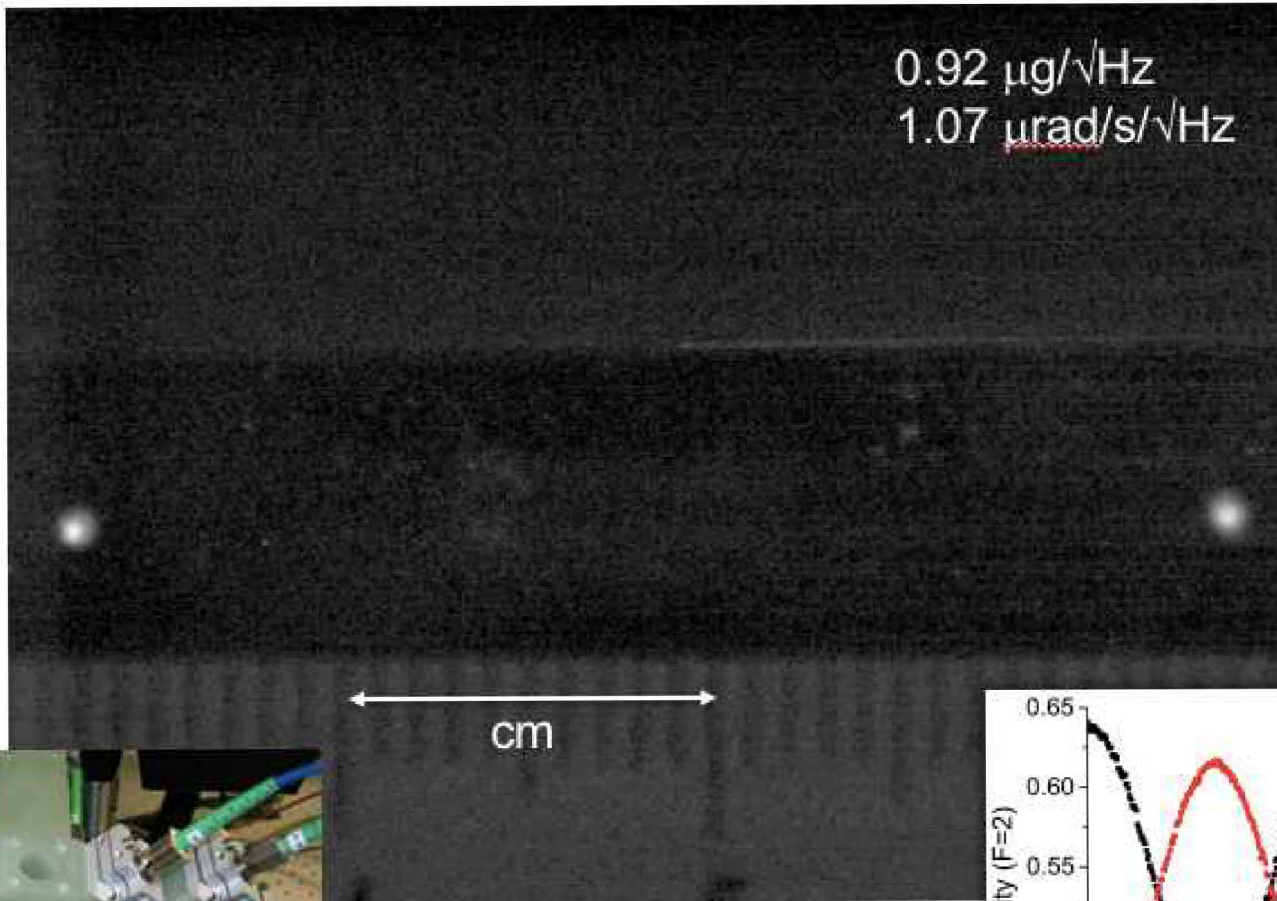
$$\Delta\phi = k_{eff} \cdot aT^2$$

$$P_e(\Delta\phi) = \frac{1}{2} (1 - \cos(\Delta\phi))$$

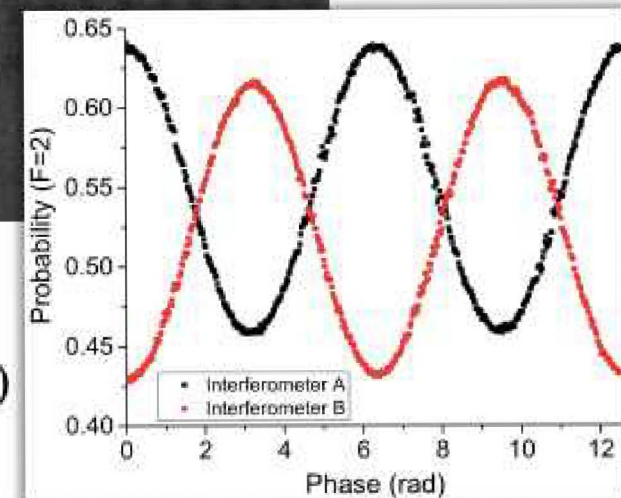
# Interferometer comparison



# Dueling interferometers



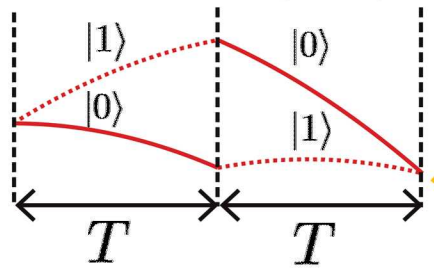
$$\Delta\phi = \vec{k}_{eff} \cdot (\vec{a}T^2 - 2(\vec{v} \times \vec{\Omega})T^2)$$



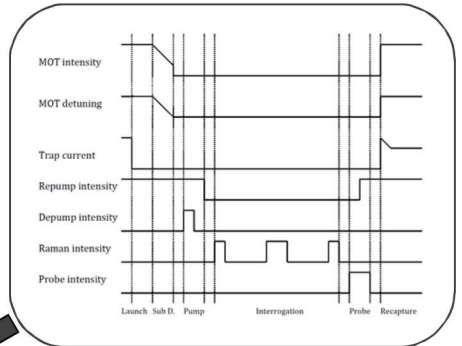
# What does it take?

UHV vacuum system

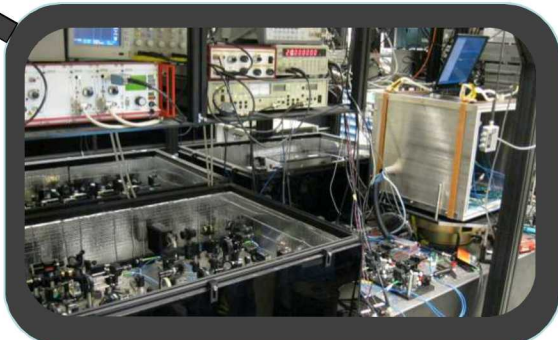
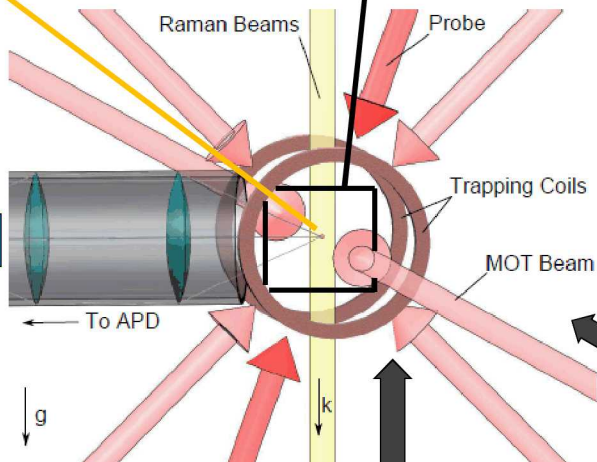
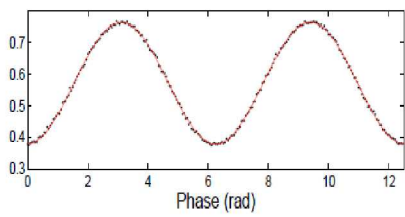
Matterwave trajectory



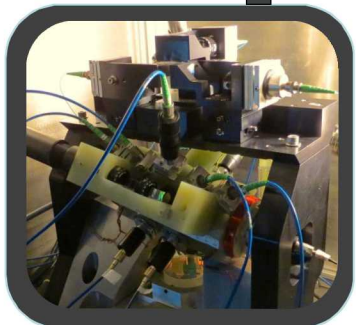
Control electronics



Acceleration →



Agile & stable laser system

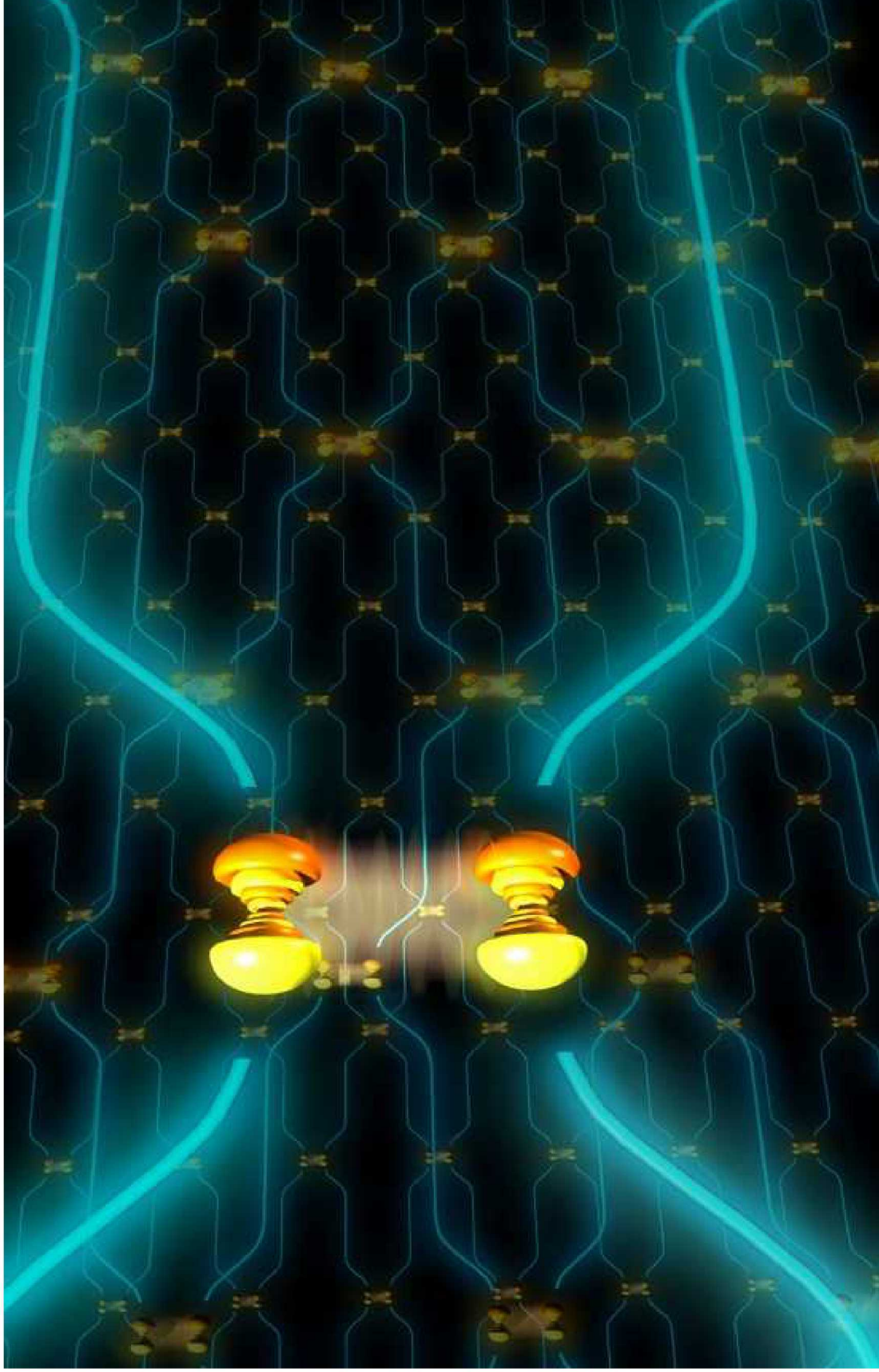


Custom optomechanics

New effort



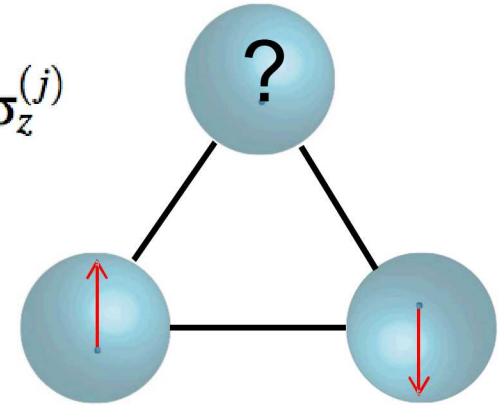
# Entangled states of neutral atoms



# Applications

Quantum simulation

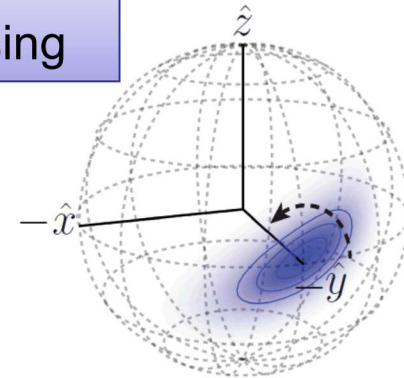
$$H_P = \sum_{i=1}^N \tilde{h}_i \sigma_z^{(i)} + \sum_{i,j=1}^N \tilde{J}_{ij} \sigma_z^{(i)} \otimes \sigma_z^{(j)}$$



Frustrated magnetism



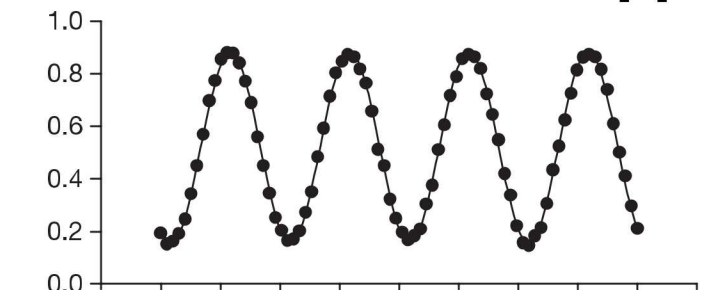
Large-scale/rapid entanglement for sensing



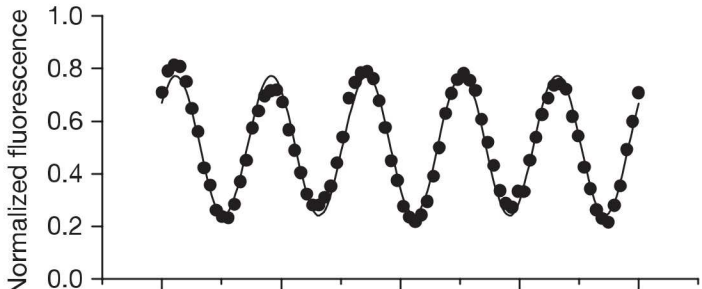
# Interferometry with GHZ states

GHZ states with ions [1]

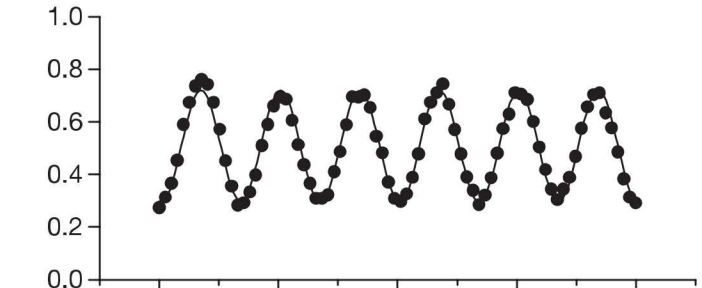
# Ions = 4



5



6

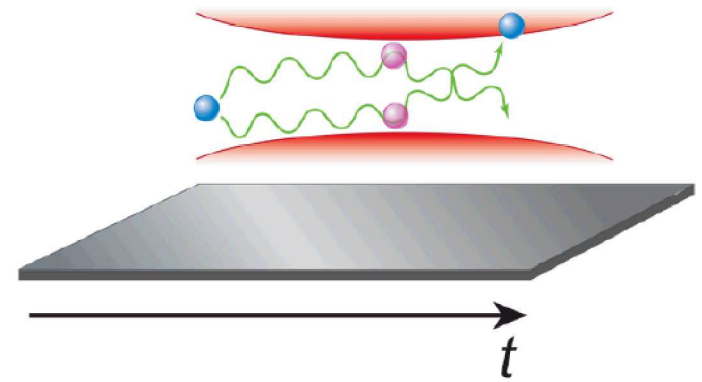


Decoding phase,  $\phi$



- Sensitivity to a Ramsey interferometer phase shift scales with N

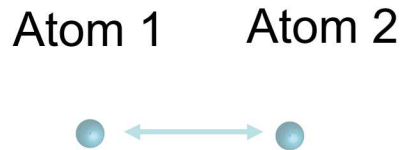
E. Rasel, Physics 5, 135 (2012)



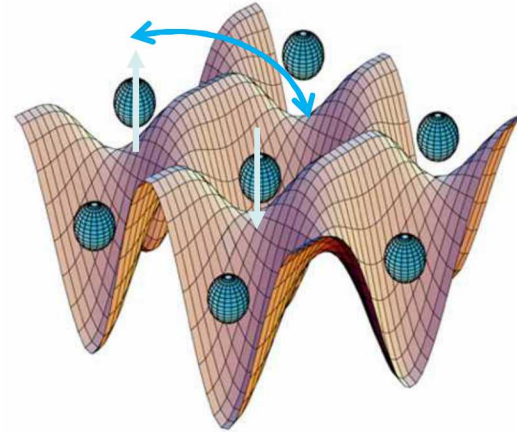
- Apply to single atom interferometers

[1] Leifried, et al., "Creation of a six-atom 'Schrödinger cat' state", *Nature* 438, 639 (2005)

# Interaction between *neutral* atoms



$$J_{ex} \propto J^2 / U$$



- Interaction between ground state atoms is small  $\sim 100$  Hz
- Thermal energy scales too large (e.g., QSIM)
- Long gate times (e.g., QIP)

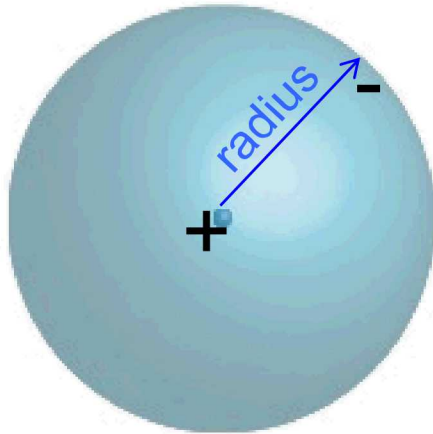
One solution: use Rydberg states

S. Trotzky *et al.*, Science **319**, 295-299 (2008)

I. Bloch, J. Dalibard, and S. Nascimbène, Nat. Phys. **8**, 267-276 (2012)

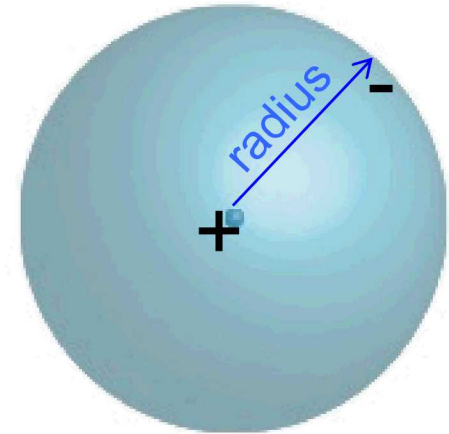
# Interaction between *neutral* atoms

Valence electron  
in Rydberg state



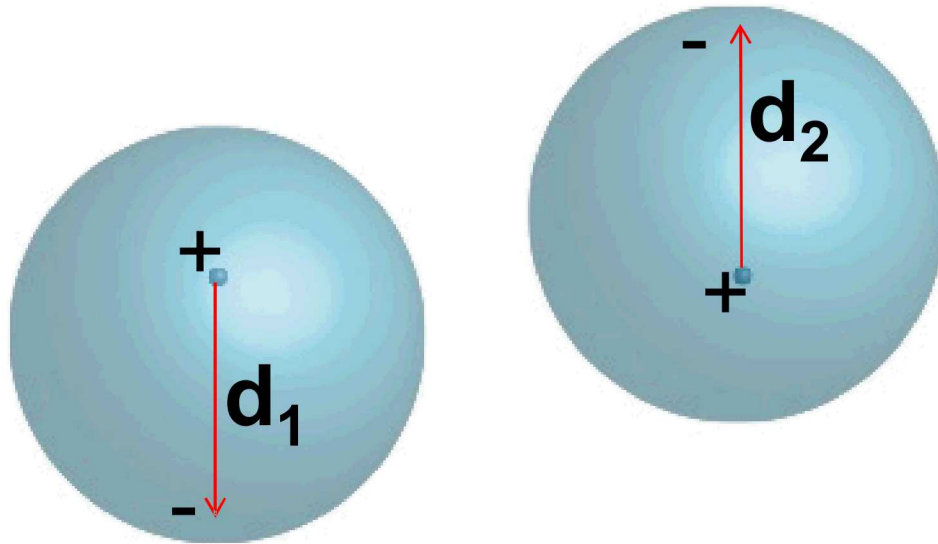
orbital radius  $\propto n^2$

Valence electron  
in Rydberg state



- Excite valence electron to Rydberg state—nearly ionized
- Atom becomes highly polarizable—strong interactions

# Interaction between *neutral* atoms



van der Waals interaction

## Parameter scaling

van der Waals

$$U \propto n^{11}$$

Lifetime

$$\tau \propto n^3$$

DC polarizability

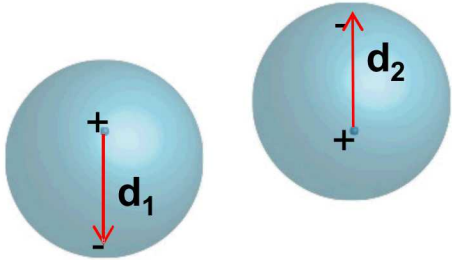
$$\alpha(0) \propto n^7$$

- Even the presence of another atom can cause a massive response  $\gg 10$  MHz
- Induced Electric Dipole-Dipole Interaction  $\propto 1/r^6$

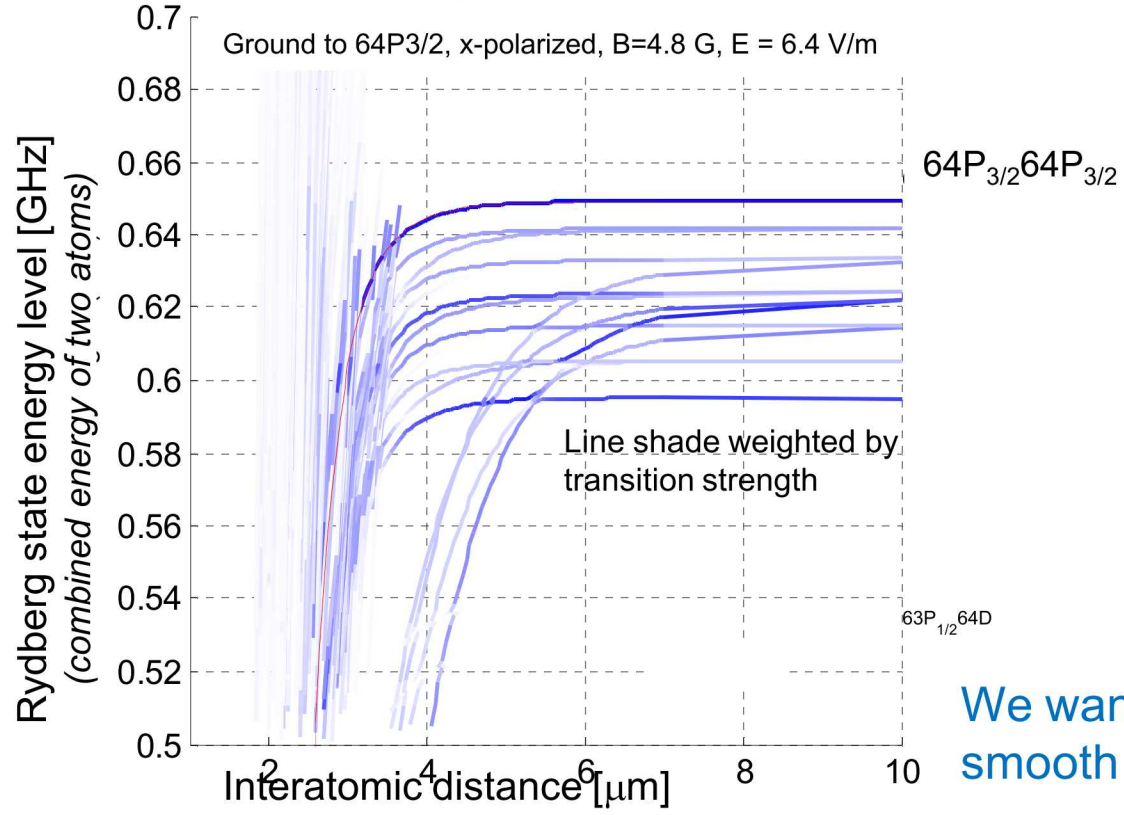
Entanglement demonstrations

- Madison: Phys. Rev. Lett. 104, 010503 (2010)
- Paris: Phys. Rev. Lett. 104, 010502 (2010)

# Rydberg blockade—the nitty gritty

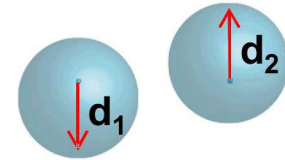


Weighted Rydberg Energy levels: Excitation from ground-state to  $64P_{3/2}$   
x-polarized light;  $B = 4.8$  G;  $E = 6.4$  V/m;



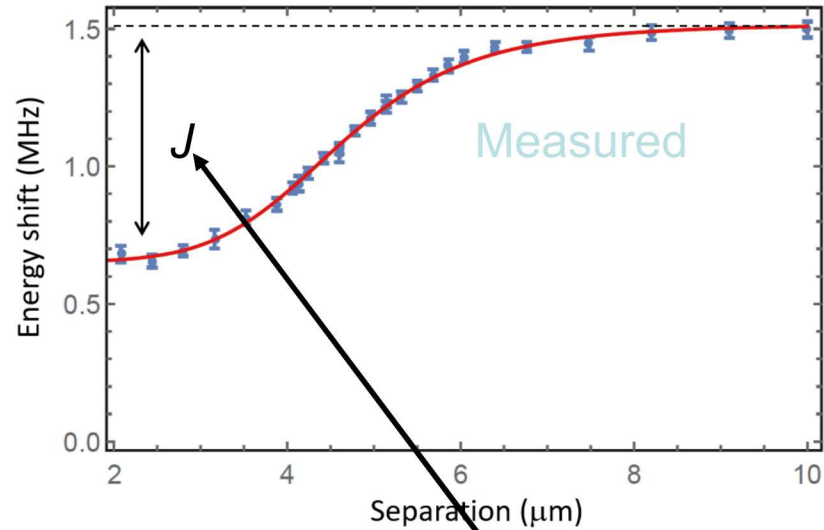
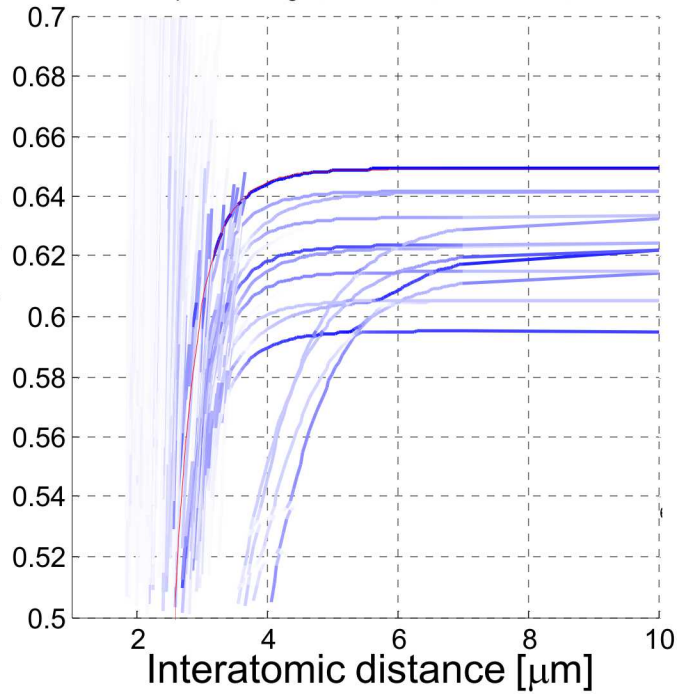
We want something smooth and tunable

# Direct Rydberg $\rightarrow$ Rydberg-Dressed



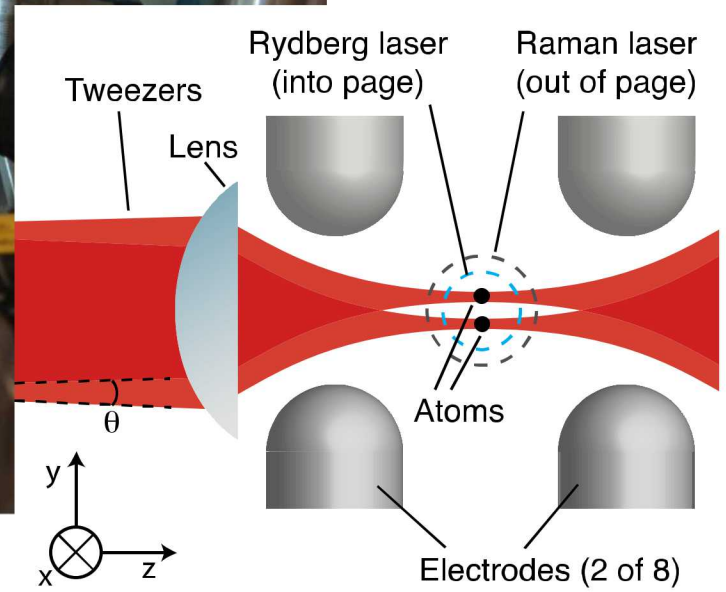
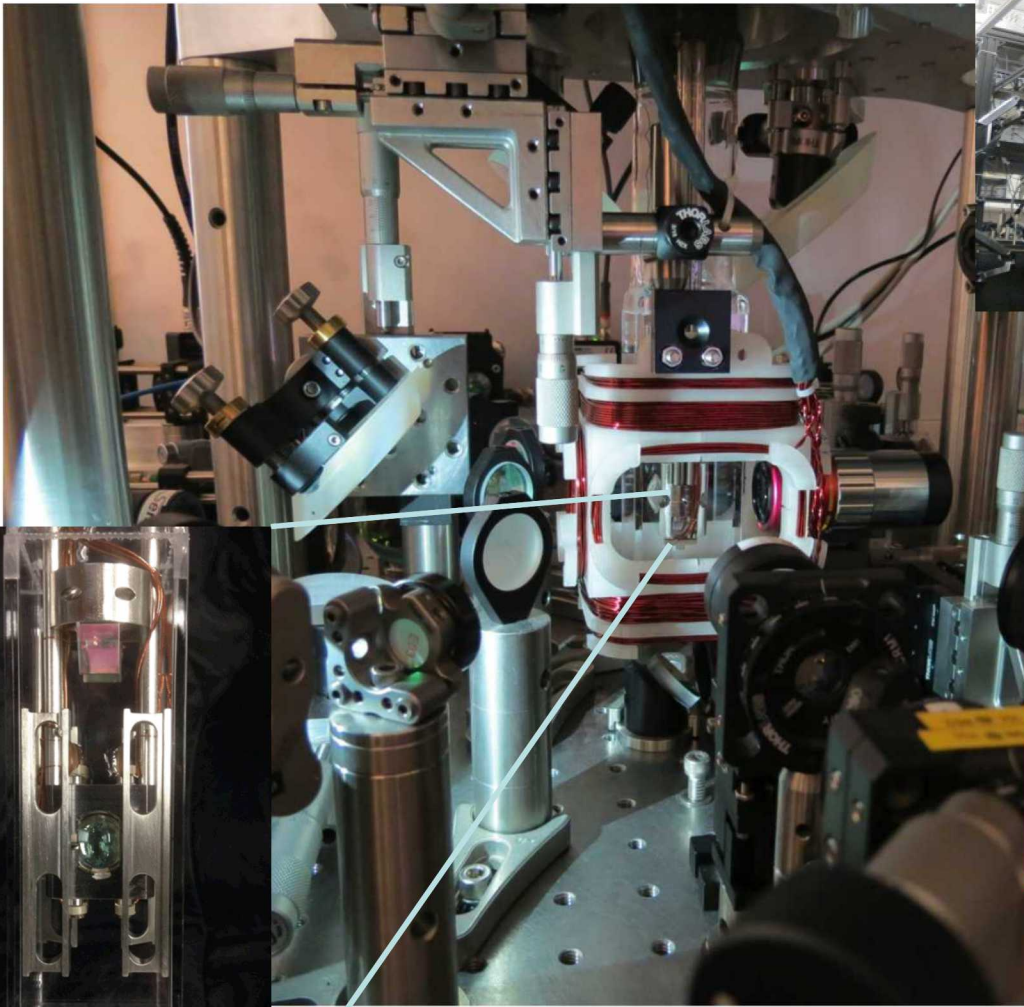
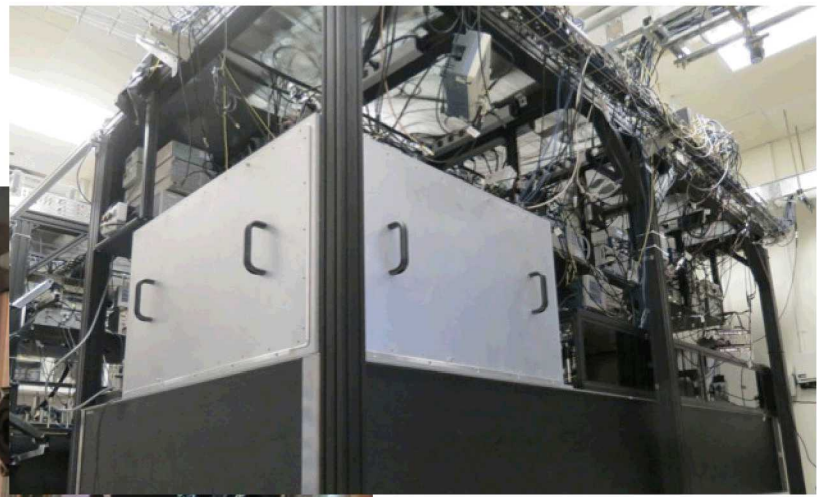
Ground to 64P3/2, x-polarized, B=4.8 G, E = 6.4 V/m

Rydberg state energy level [GHz]  
(combined energy of two atoms)

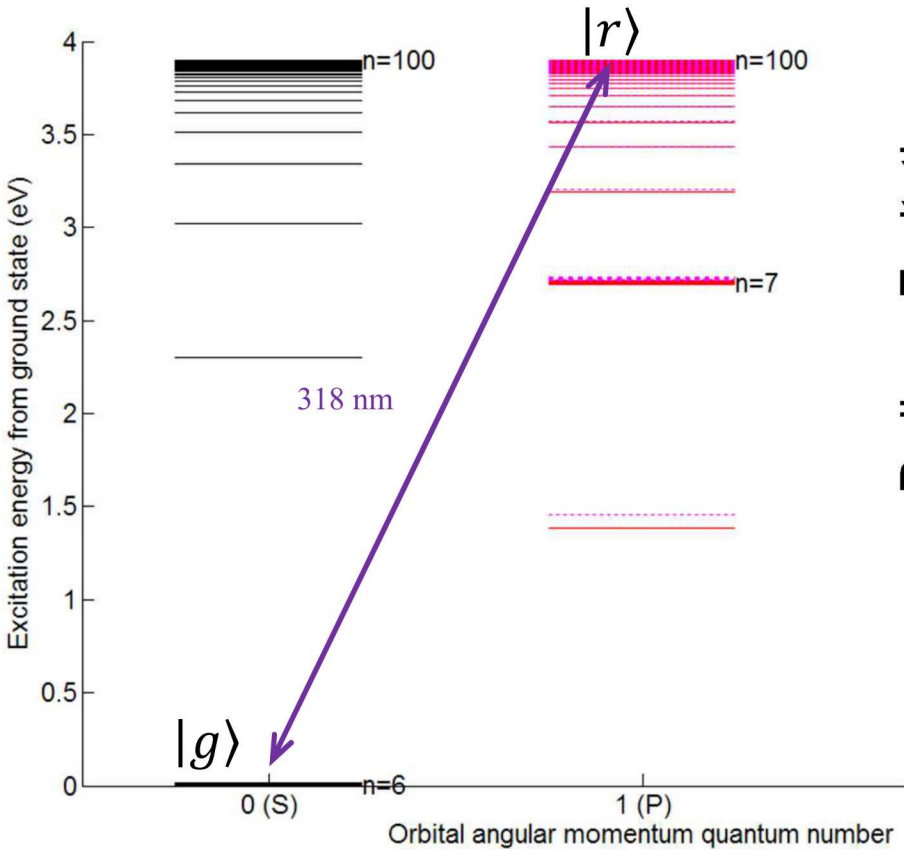


$$H_P = \sum_{i=1}^N \tilde{h}_i \sigma_z^{(i)} + \sum_{i,j=1}^N \tilde{J}_{ij} \sigma_z^{(i)} \otimes \sigma_z^{(j)}$$

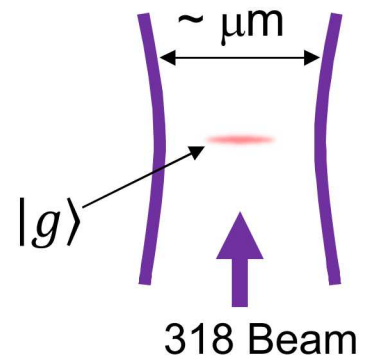
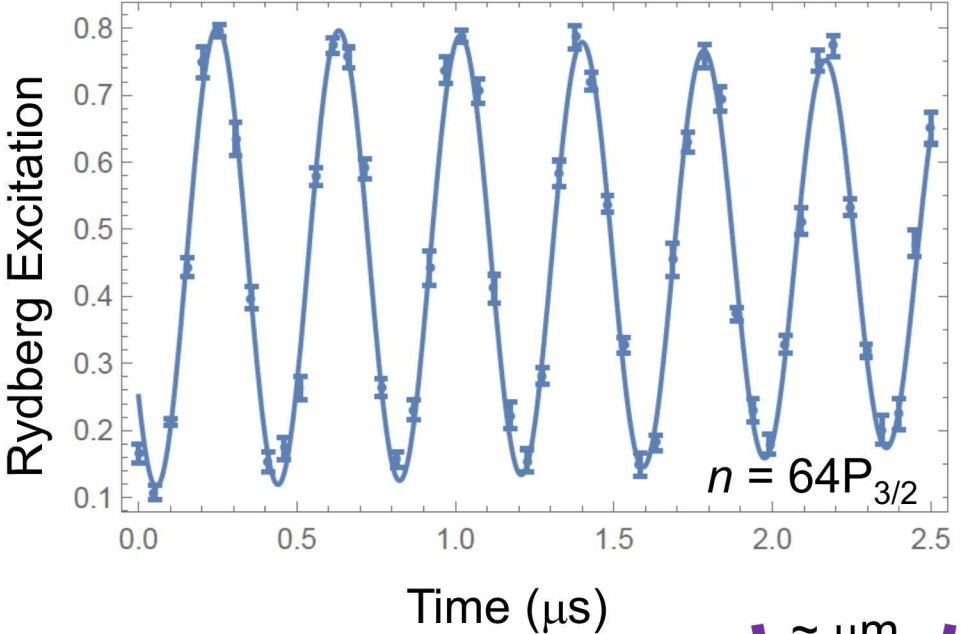
# Apparatus



# Rydberg Rabi flopping with 318 nm laser

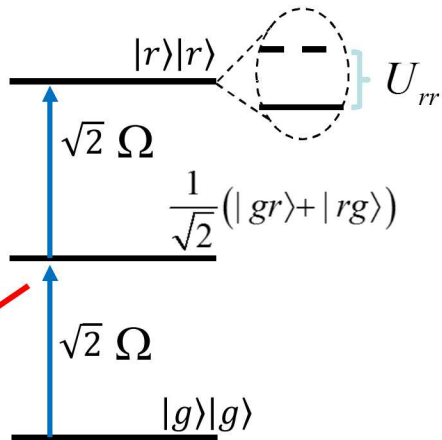
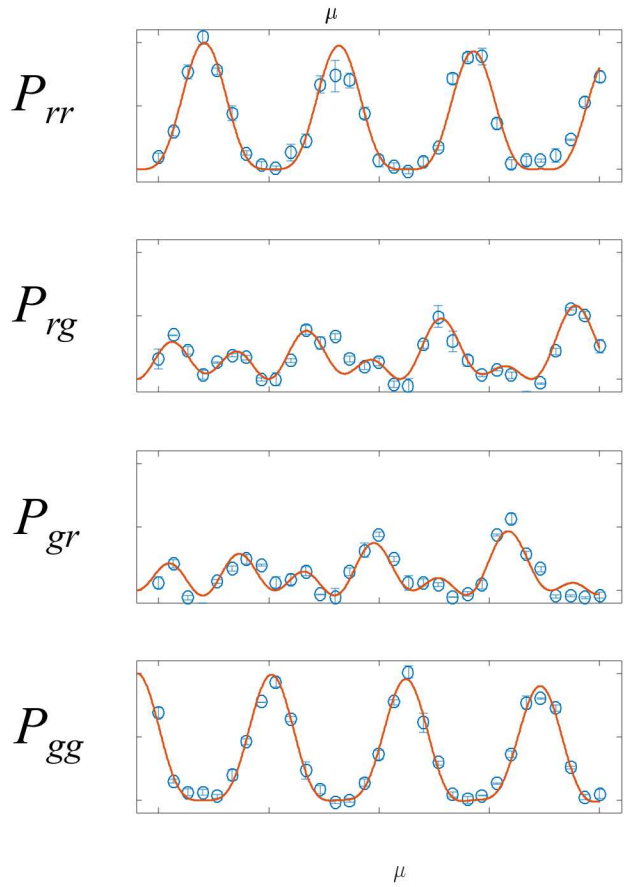


Direct excitation, measured through loss

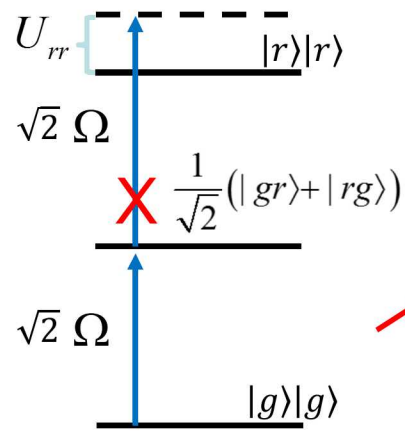
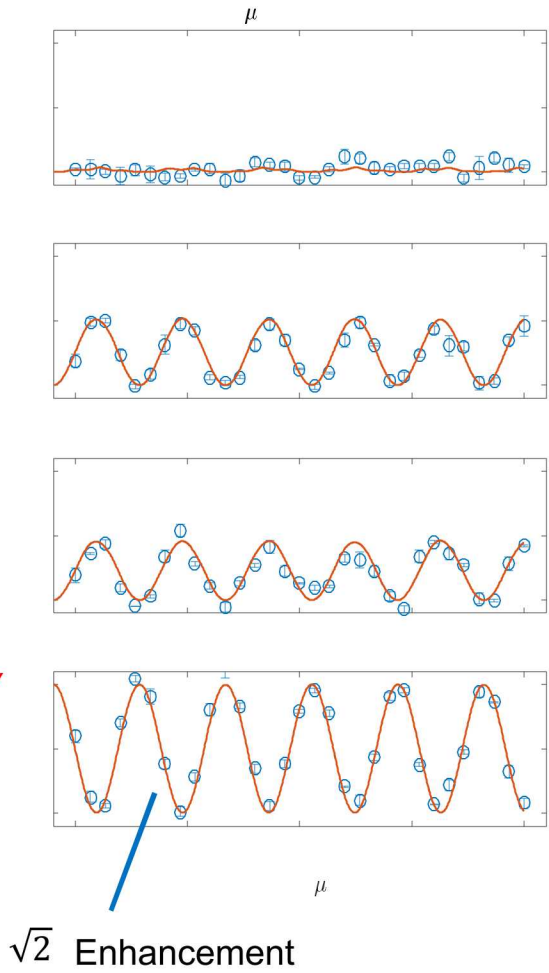


# Rydberg blockade

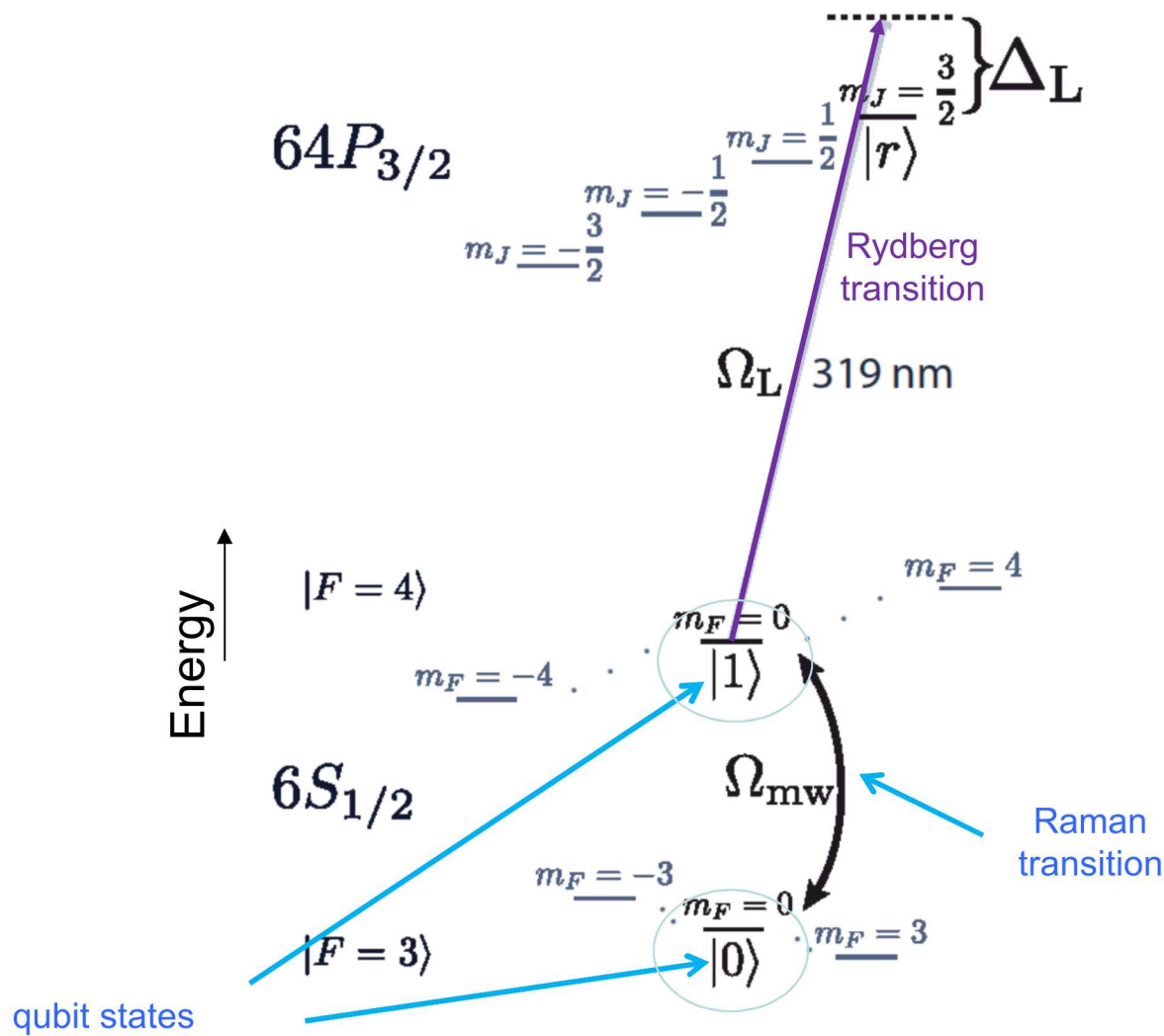
Weak ( $U_{RR} < 100$  kHz)



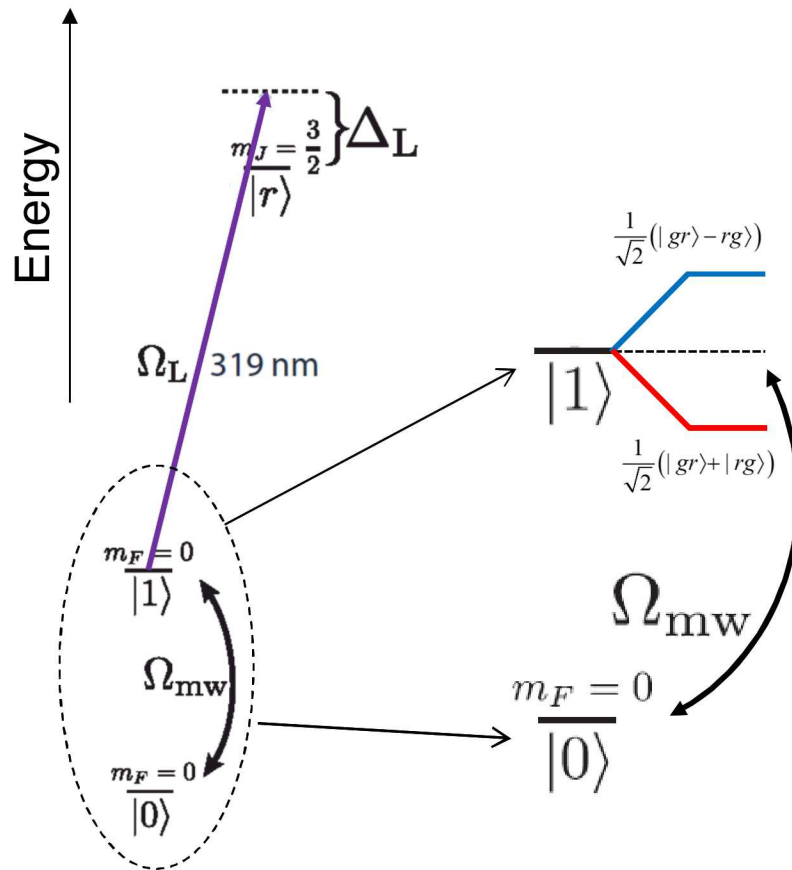
Strong ( $U_{RR} > 6$  MHz)



# Creating Rydberg-Dressed states



# Creating Rydberg-Dressed states

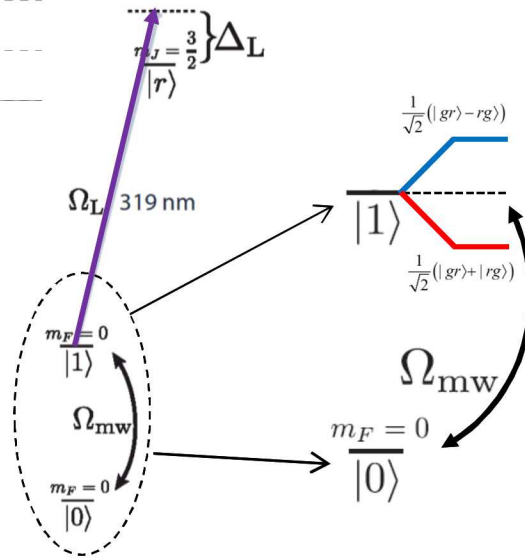
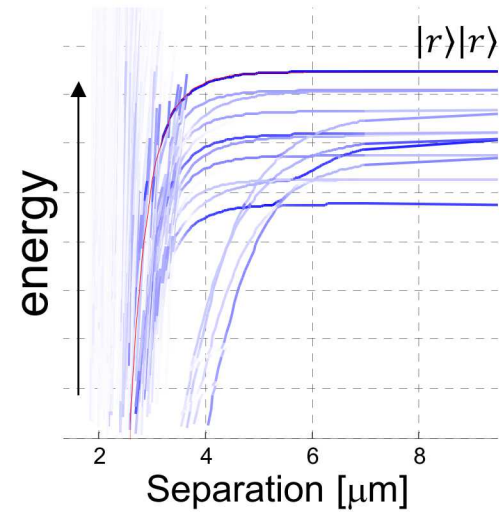


## Light-shift Hamiltonian

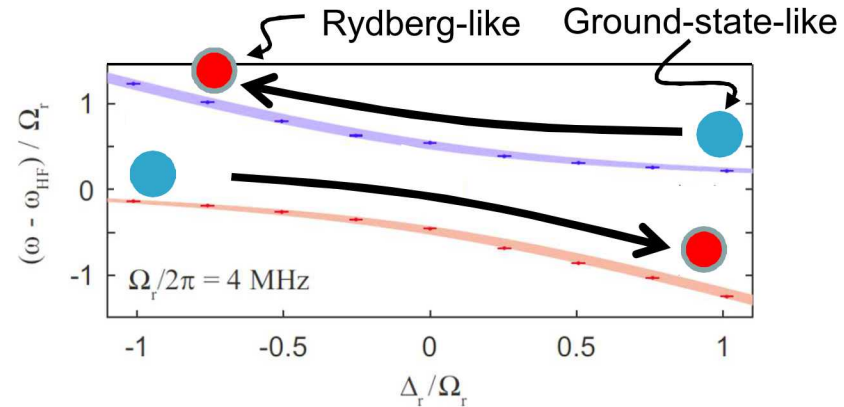
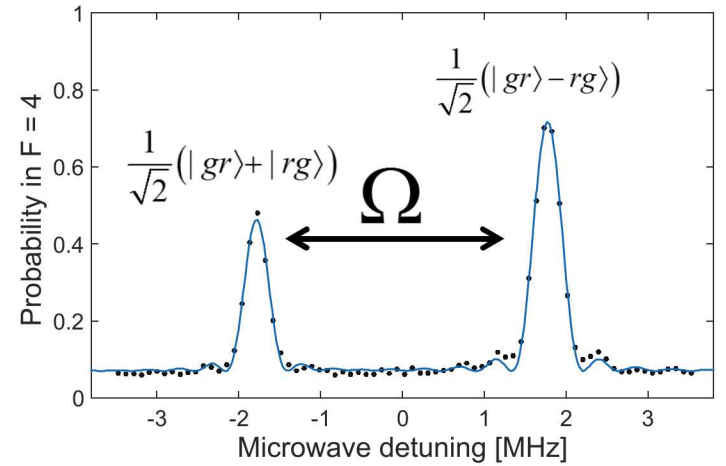
$$H = \frac{\hbar}{2} \begin{pmatrix} -2\Delta_L & \Omega_L \\ \Omega_L & 0 \end{pmatrix}$$

$$\Delta = \frac{\hbar \Omega_L^2}{4 \Delta_L}$$

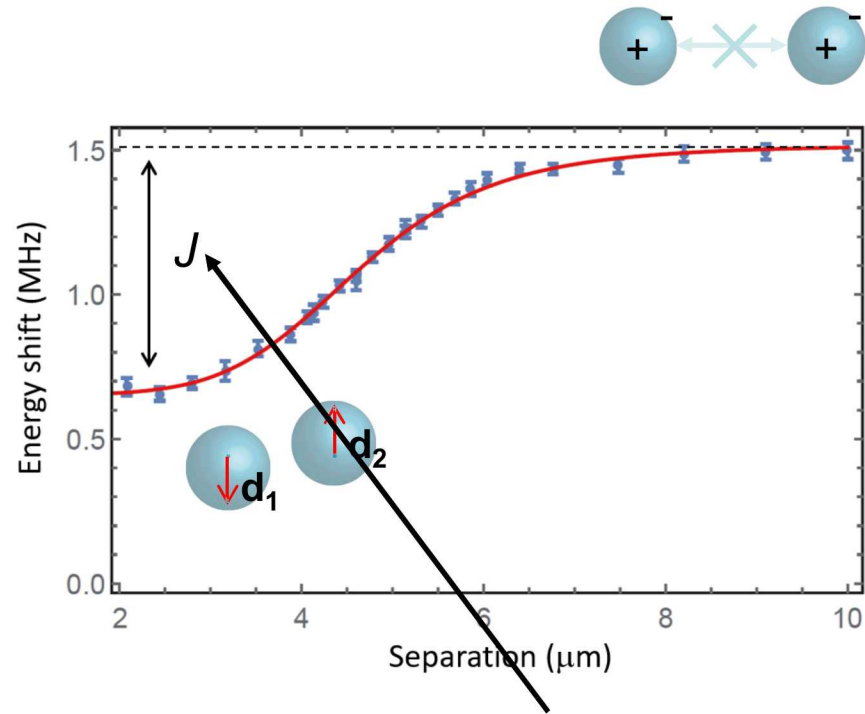
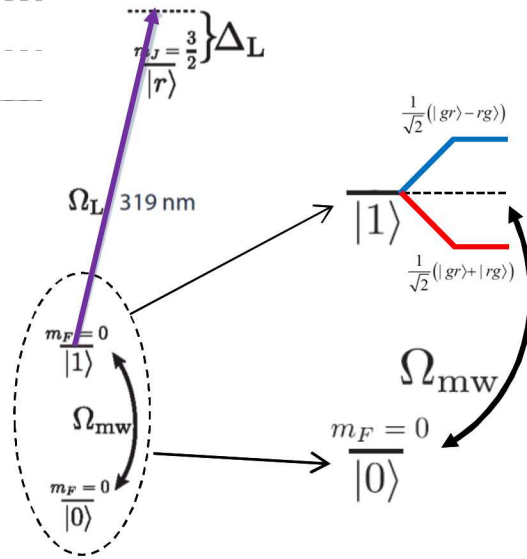
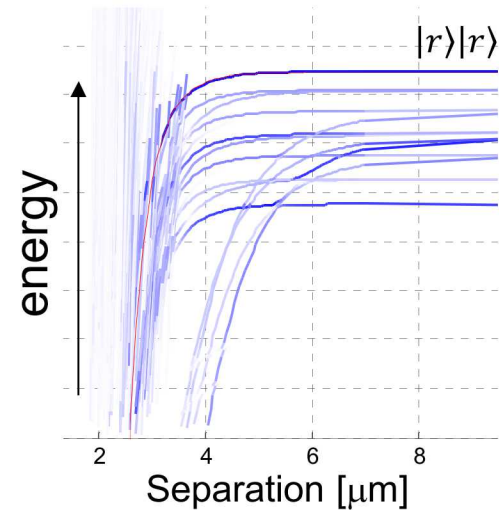
# 1-atom Rydberg-dressed states



Dressed F=4 state Autler-Townes splitting



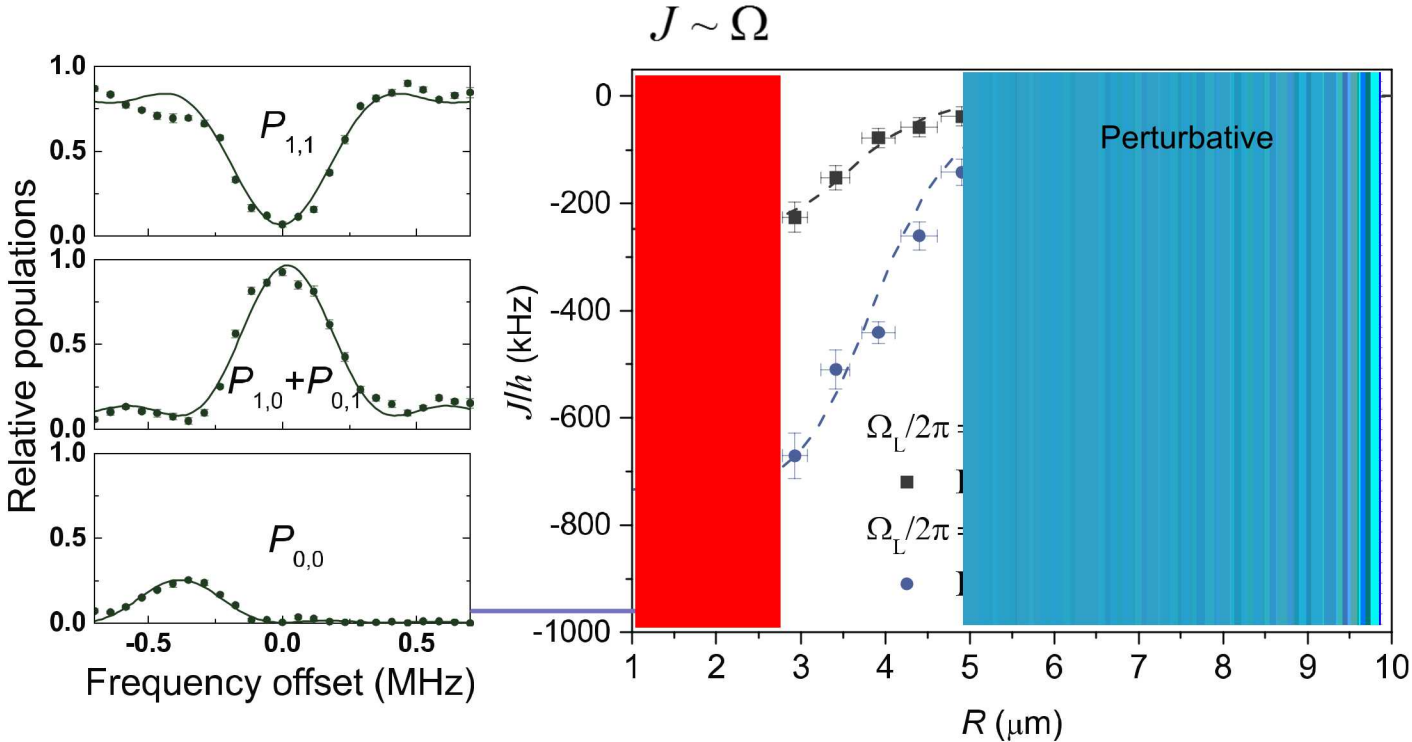
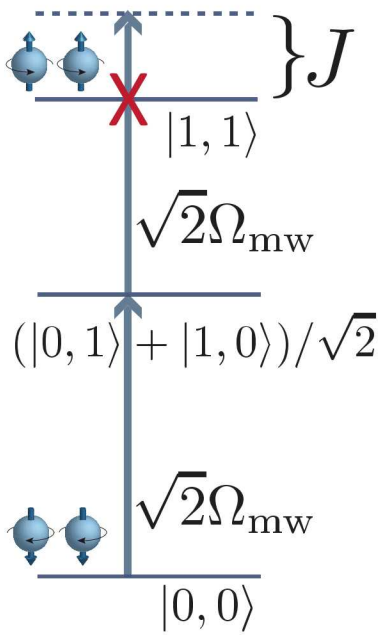
# Rydberg-Dressed interaction



$$H_P = \sum_{i=1}^N \tilde{h}_i \sigma_z^{(i)} + \sum_{i,j=1}^N \tilde{J}_{ij} \sigma_z^{(i)} \otimes \sigma_z^{(j)}$$

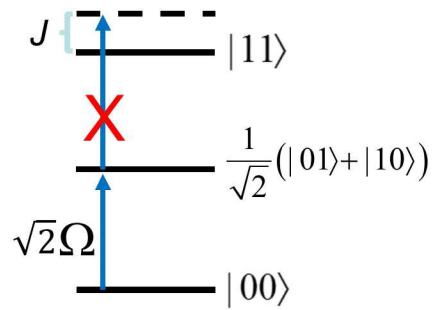
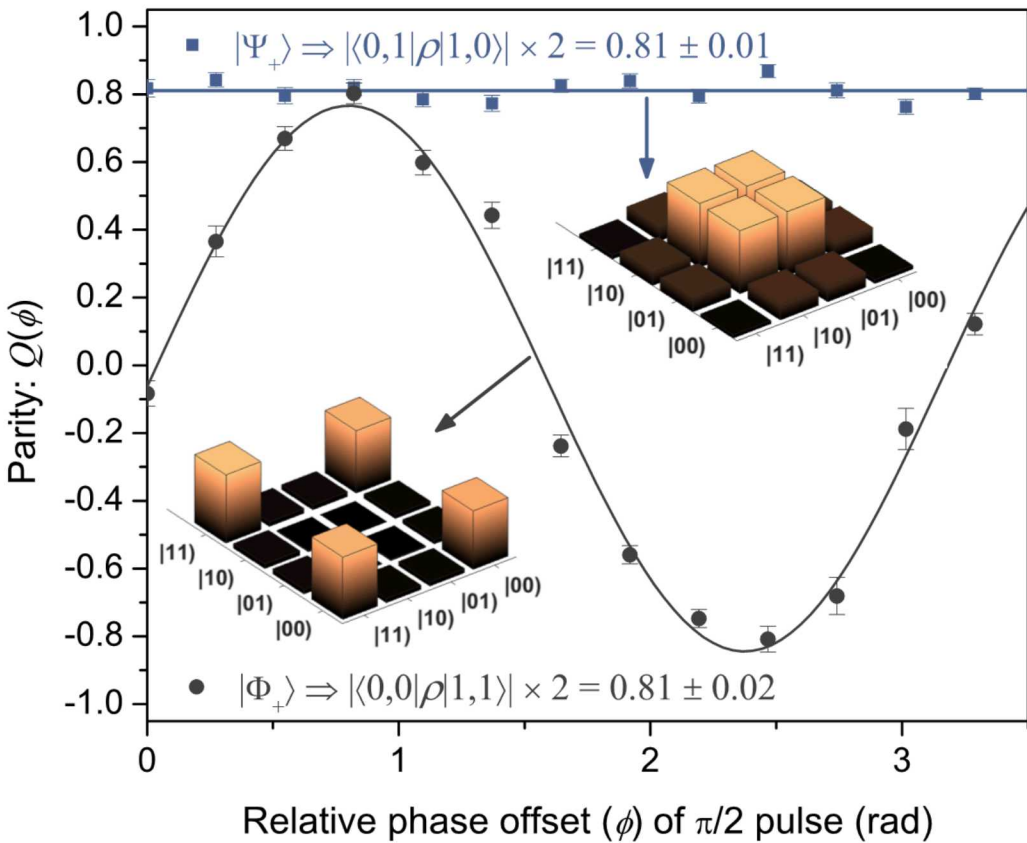
# Observing blockade on hyperfine qubit

Direct measurement of two-qubit interaction strength  $J$  as a function of two-atom separation with different dressing conditions.



# Spin-flip blockade

Verify the entanglement via parity measurements



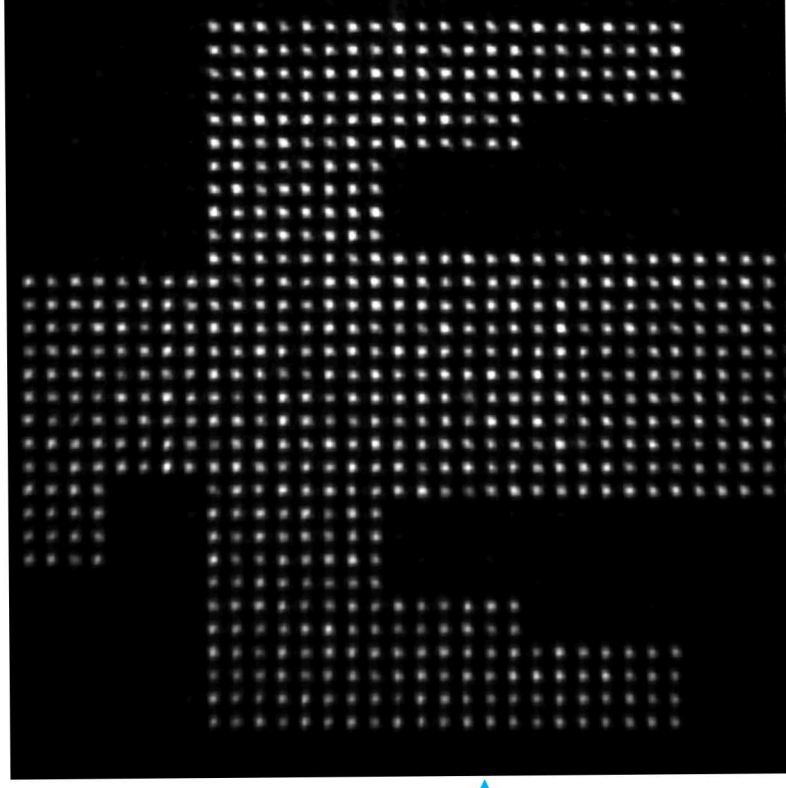
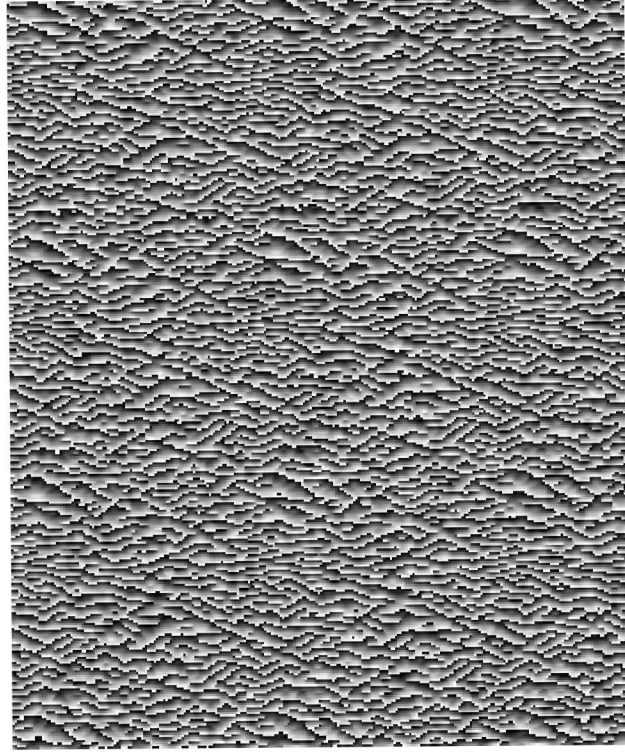
1. Prepare Bell state
2. Apply global  $\pi/2$  with given phase
3. Measure parity  $Q$
4. Obtain bound on fidelity = 0.81(2)

# Generation and control of >500 of individual traps

Intensity profile (measured)

Gerchberg-Saxton algorithm  
With GPU compute acceleration  
(~50 Hz hologram calculation)

Phase to SLM

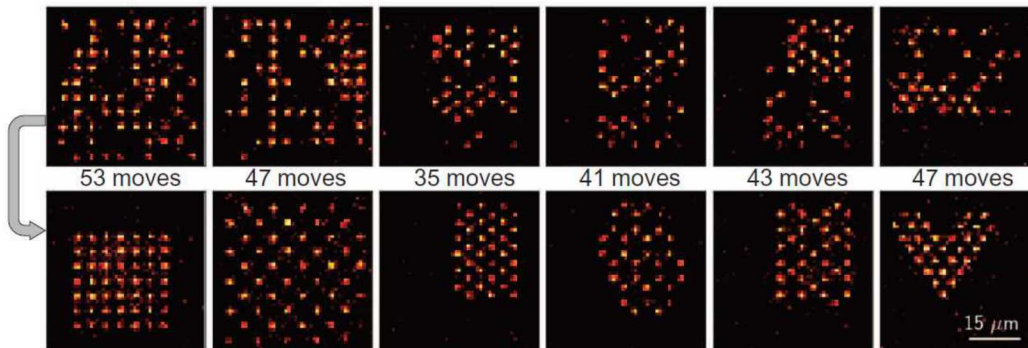


# Defect-free neutral atom arrays

## Lukin group, Harvard<sup>1</sup>

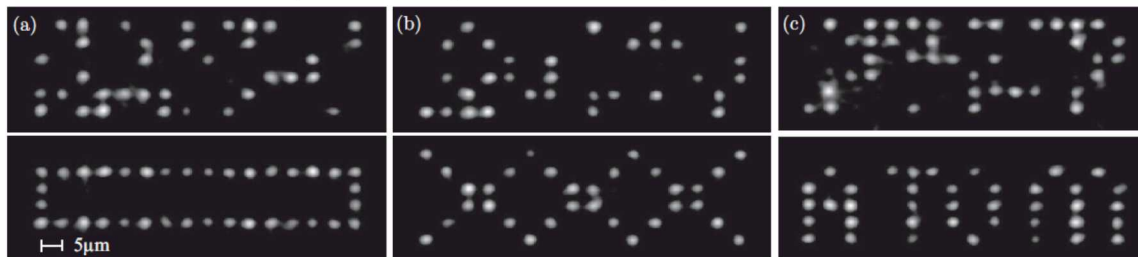


## Browaeys group, Institut d'Optique (France)<sup>2</sup>



- There has been a revolution in creating defect free, controllable arrays of neutral atoms.
- 1D, 2D & 3D arrays demonstrated.
- This capability makes scaling to a many qubit simulator a possibility.

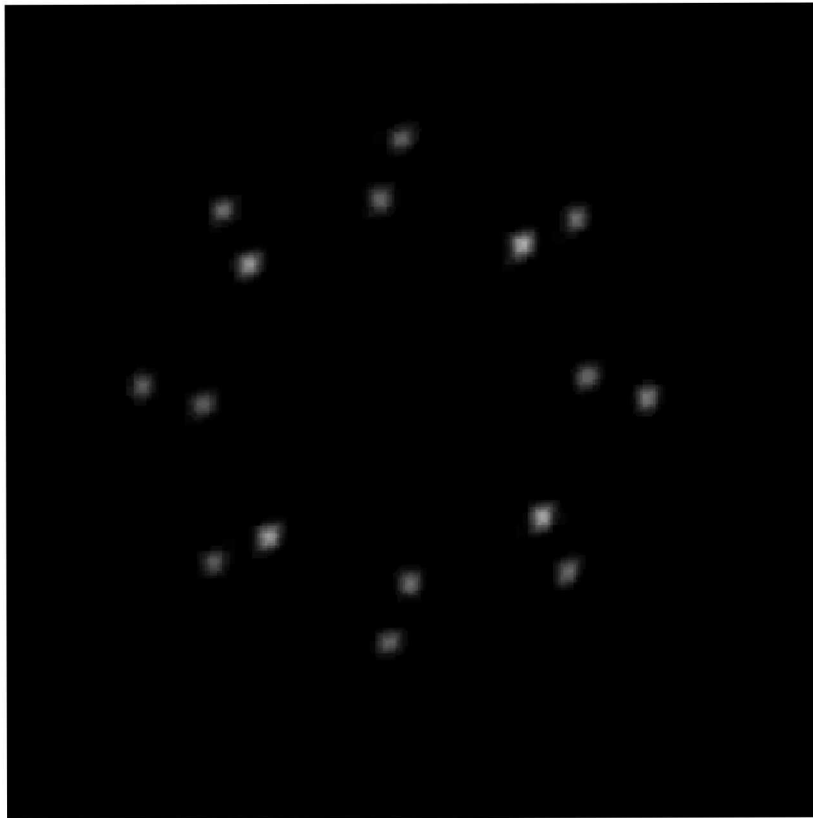
## Ahn Group, KAIST (South Korea)<sup>3</sup>



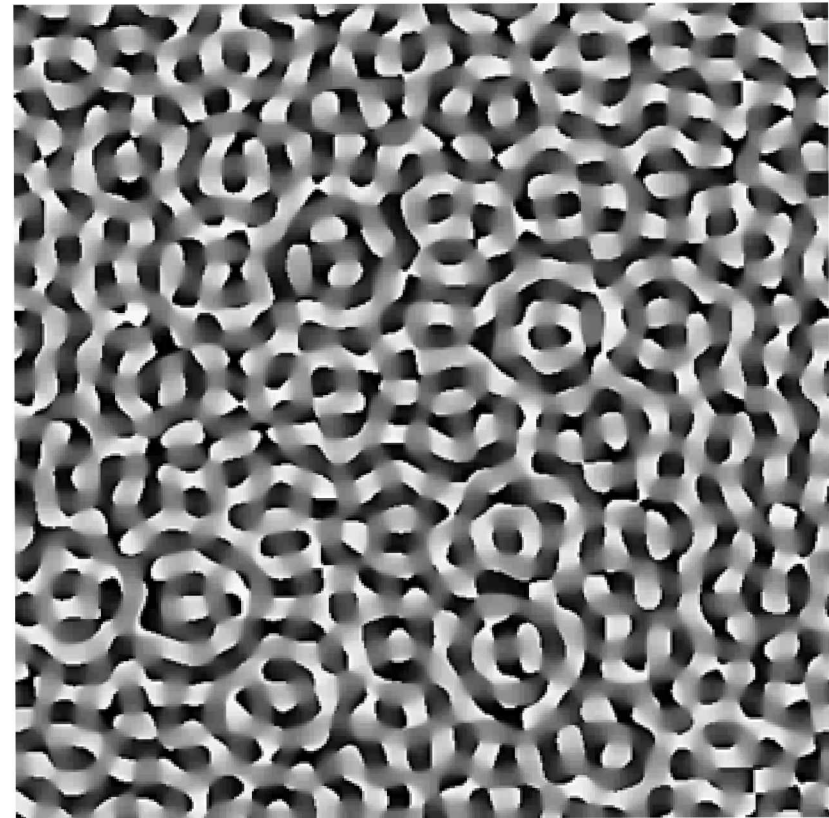
1. M. Endres et al., Science 354, 1024 (2016)
2. D. Barredo et al., Science 354, 1021 (2016)
3. W. Lee et al., Phys. Rev. A **95**, 053424 (2017)

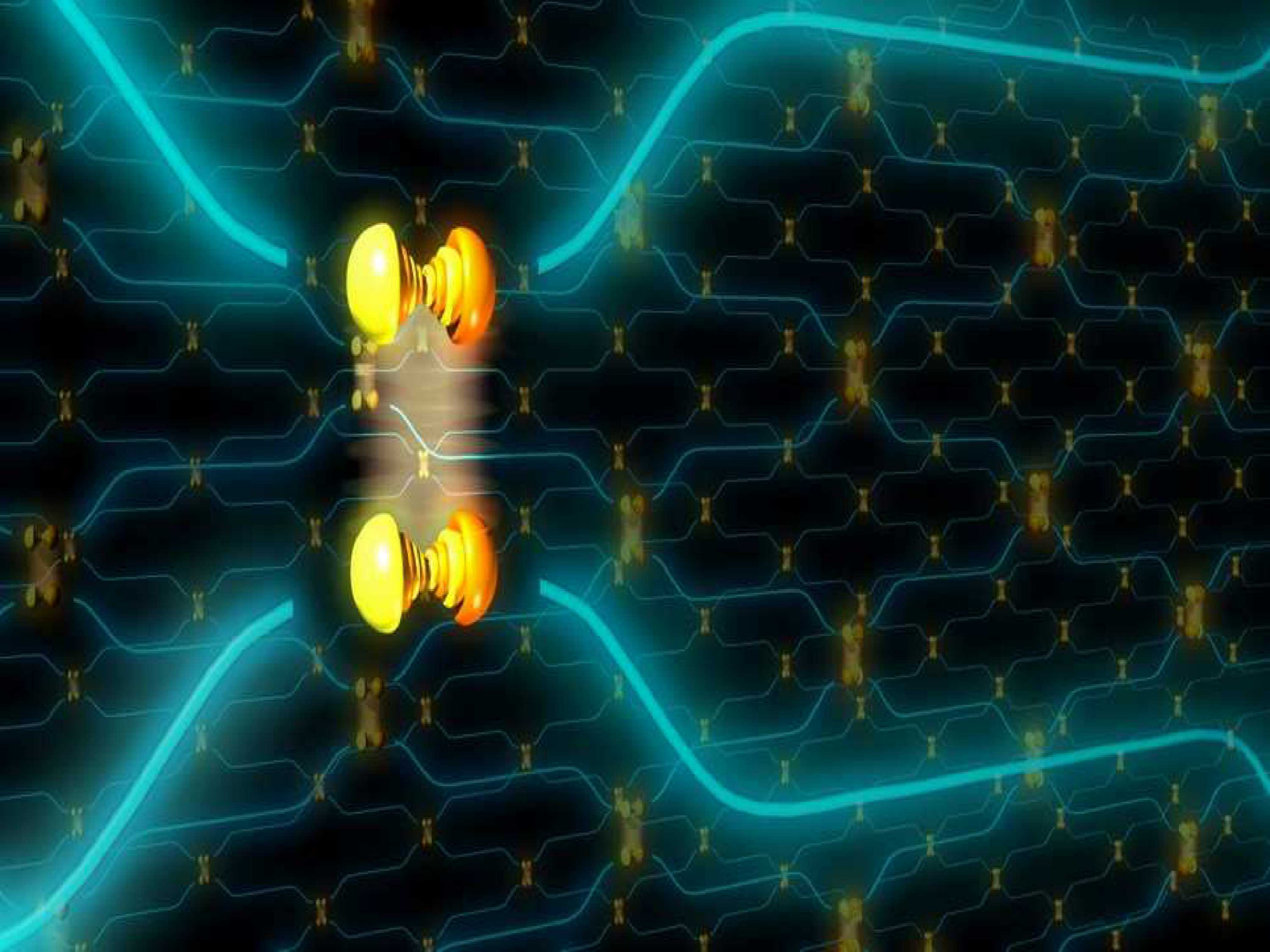
# Dynamic control of hologram-generated traps

30 fps, real time

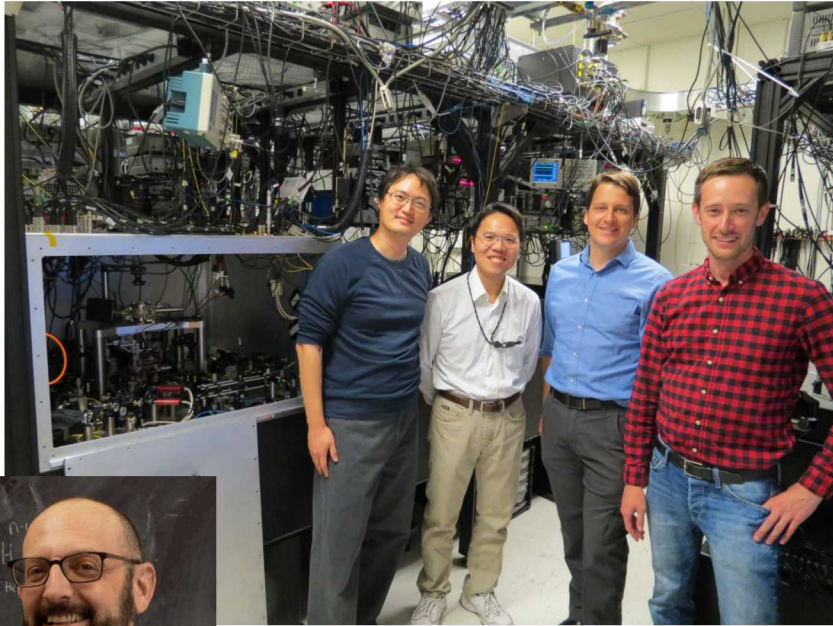


Phase hologram





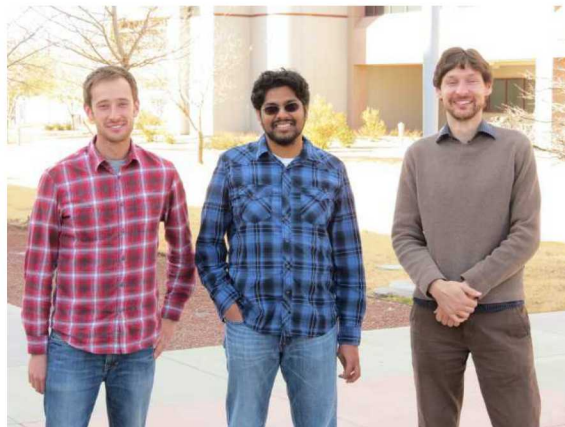
# Sandia Team



J. Lee, Y-Y. Jau, M. Martin, G. B.



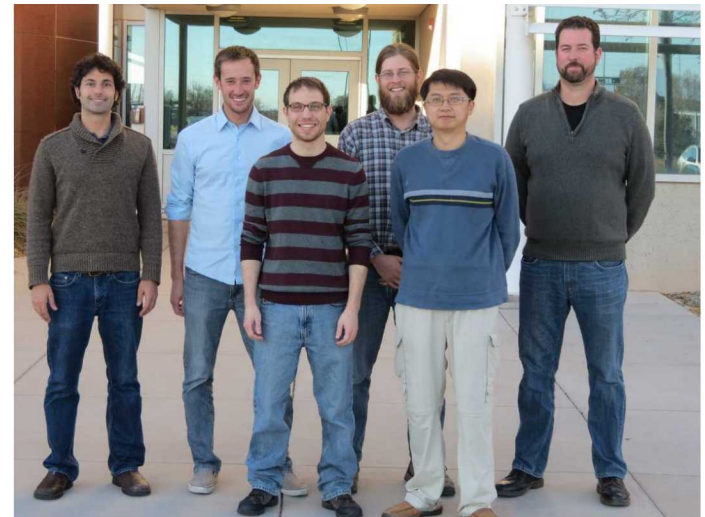
I. Deutsch (UNM)



G. B., A. Rakholia, H. McGuinness



Sandia mountains viewed from the Rio Grande near Albuquerque, New Mexico



L. Parazzoli, G. B., A. Hankin, A. Ferdinand, J. Chou, G. Burns

Postdoc position available. Contact me!

