

Efficient Band-to-Trap Tunneling Model Including Heterojunction Band Offset

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Acknowledgement goes to my colleagues, Gary Hennigan, Lawrence Musson, Mihai Negoita, Jason Gates, Kara Peterson, Samuel Myers, William Wampler, and Normand Modine for their many helpful discussions; and special thanks goes to Myers for providing his numerical code.

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Traps in Semiconductor Devices

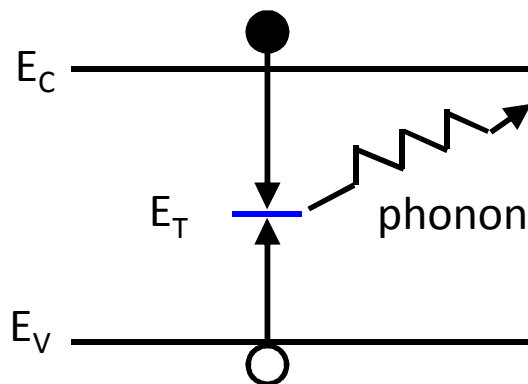
Fabrication processes

Radiation environment

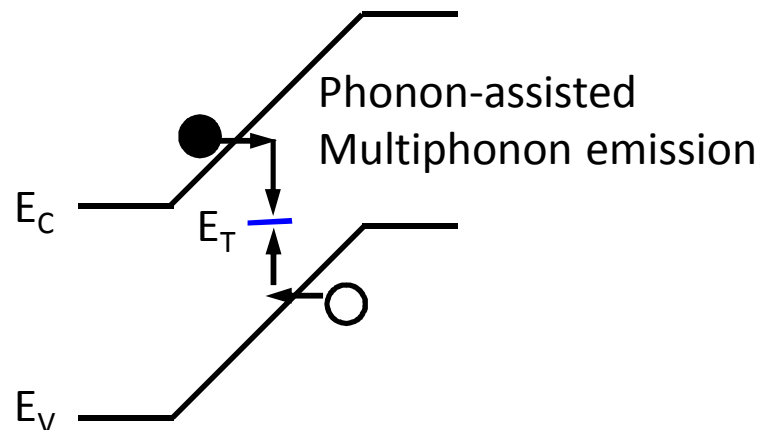
Always **traps/defects** in semiconductor devices

Other trap-related processes

Shockley-Read-Hall (SRH) Recombination



Band-to-trap tunneling (i.e., trap-assisted tunneling, TAT)



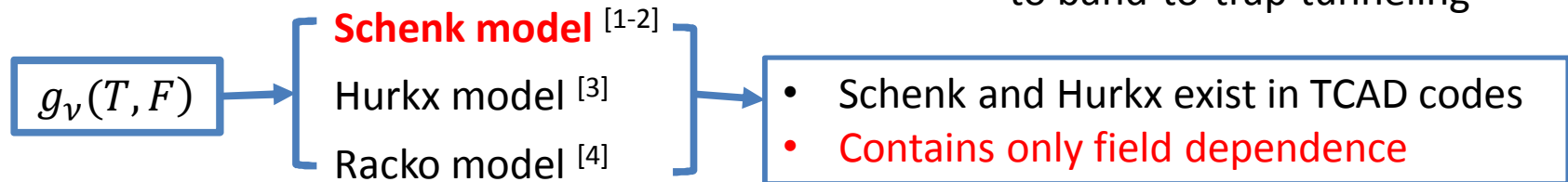
Modeling of Band-to-Trap Tunneling

Band-to-trap tunneling is widely modeled as field-dependent SRH recombination:

$$R = \frac{np - n_i^2}{\tau_p(n + n_1) + \tau_n(p + p_1)}$$

$$\tau_v = \frac{\tau_{v0}(T, F = 0)}{1 + g_v(T, F)}, v = n, p$$

Field enhancement factor due to band-to-trap tunneling



Assuming **Boltzmann statistics and deep-level traps**, the **Schenk model** is reduced to

$$g_n(T, F) = \frac{\int_0^{E_t} dE \rho_c(E) I_{E/\hbar\omega_0}(z) \exp\left(\frac{-E}{2k_B T}\right)}{\int_{E_t}^{\infty} dE \rho_c^{F=0}(E) I_{E/\hbar\omega_0}(z) \exp\left(\frac{-E}{2k_B T}\right)}$$

Density of states (DOS)

- Contains the field dependence
- Obtained assuming constant field

[1] A. Schenk, Journal of Applied Physics **71**, 3339 (1992).

[2] A. Schenk, Solid-State Electronics **35** (11), 1585 (1992).

[3] G. A. M. Hurkx, IEEE Transaction on Electron Devices **39**(2), 331 (1992).

[4] J. Racko et al., RadioEngineering **21**(1), 213 (2012).

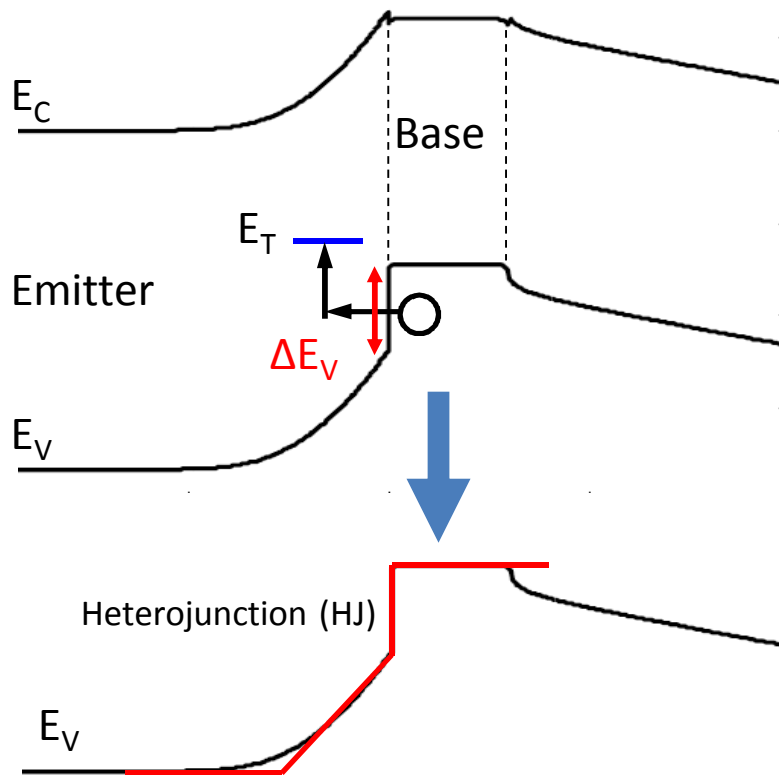
B-to-T Tunneling in Heterojunction Device

Constant field assumption for the DOS in the Schenk model

Good for **homojunction** devices

Not good for **heterojunction** devices

Consider $\text{In}_{0.49}\text{Ga}_{0.51}\text{As}/\text{GaAs}$ NP⁺N HBT:



DOS at the trap location is **greatly enhanced** [5] due to the **band offset** ΔE_V , leading to much **higher hole-to-trap tunneling**.

Sam Myers et al. [5] obtained **much higher DOS** than the constant-field DOS by **numerically solving the Schrodinger equation** given the actual E_V band profile.

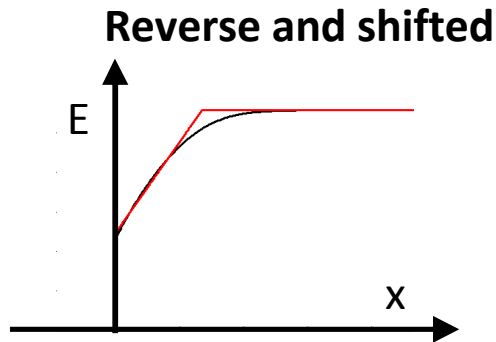
Trivial – solve 1D Schrod. eqn. numerically once

Inefficient – when solved repeatedly

Difficult – coupled to MPI-parallel PDE-based TCAD code

Goal – find **analytic** DOS model for simplified potential & easy implementation in TCAD code

Density of States Method




Goal – find analytic DOS model for simplified potential (red potential) & easy implementation into TCAD code

- Need to compute **3D DOS** (used in the Schenk model)
- Non-zero potential in the x direction only
- Three methods to compute the DOS

Method 1 – Numerical approach

$$\rho_{3D}(x, E) = \frac{m}{\pi \hbar^2} \sum_i |\psi_i(x)|^2 \theta(E - E_i) \quad \theta(\cdot) - \text{step function}$$



Wave functions **normalized to 1** **Discrete** energy spectrum

❖ Require solving the Schrodinger equation **numerically** to obtain ψ_i and E_i

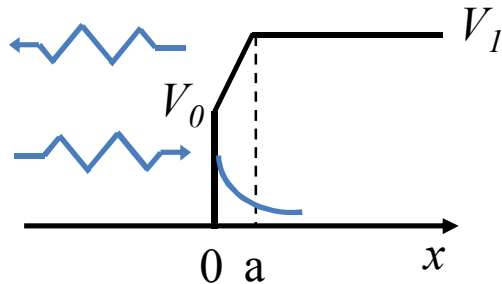
Method 2 – Green's function (GF) approach

$$\rho_{3D}(x, E) = \frac{-2}{\pi} \text{Im}[G_{3D}^R(x, E)] \quad G_{3D}^R(x, E) = \frac{m}{2\pi \hbar^2} \int_0^{+\infty} G_{1D}^R(x, E - E_{\perp}) dE_{\perp} \theta(E - E_{\perp})$$

❖ Analytic GFs is even harder to obtain than the WFs

Proposed Density of States Method

Method 3 – Proposed scattering approach



Solve the Schrodinger equation with **open boundary condition (BC)**

- **Continuous** energy spectrum
- **Dirac-delta** normalized WF

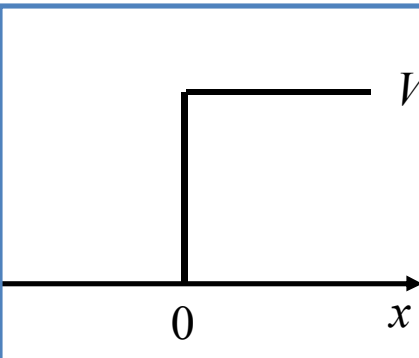
$$\int_{-\infty}^{+\infty} \psi_{E_x}^*(x) \psi_{E'_x}(x) dx = \delta(E_x - E'_x)$$

$$\rho_{3D}(x, E) = \frac{m}{\pi \hbar^2} \int_0^{+\infty} |\psi_{E_x}(x)|^2 dE_x \theta(E - E_x)$$

Proposed Density of States Method

How does the approach work ?

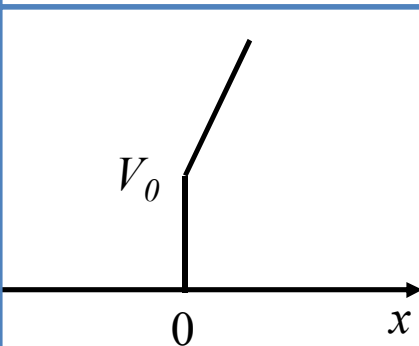
Step 1: Solve the Schrodinger equation with **open BC** to obtain **analytic** wave functions



$$\psi(x) = N \begin{cases} e^{ikx} + a e^{-ikx} & x < 0 \\ b e^{-vx} & x > 0 \end{cases}$$

$$k = \sqrt{\frac{2mE}{\hbar^2}} \quad v = \sqrt{\frac{2m(V_1 - E)}{\hbar^2}}$$

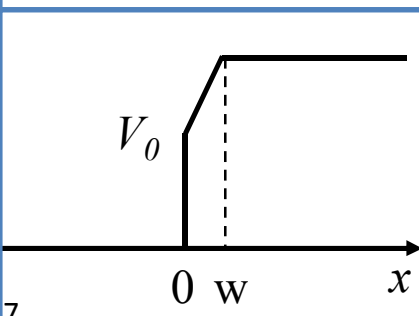
Consider $0 < E < V_1$ and $x > 0$ (relevant for b-to-t tunneling)



$$\psi(x) = N \begin{cases} e^{ikx} + a e^{-ikx} & x < 0 \\ b A_i[y(x)] & x > 0 \end{cases}$$

$$y(x) = \left(\frac{2mqF}{\hbar^2} \right)^{1/3} \left(x + \frac{V_0 - E}{qF} \right)$$

Coefficients (a, b, c, d) depend on E and are determined by continuity of ψ and $d\psi/dx$ at $x = 0$ & w



$$\psi(x) = N \begin{cases} e^{ikx} + a e^{-ikx} & x < 0 \\ b A_i[y(x)] + c B_i[y(x)] & 0 < x < w \\ d e^{-vx} & x > w \end{cases}$$

Proposed Density of States Method

Step 2: Determine the normalization factor

$$\int_{-\infty}^{+\infty} \psi_{E_x}^*(x) \psi_{E'_x}(x) dx = \delta(E_x - E'_x) \longrightarrow \text{Determine the normalization factor } \mathbf{N}$$

- Challenge task to normalize the WFs to the delta function, even for step barrier
- Discover **the same normalization factor** for all the three potentials considered
- Believe **the same factor is applicable to all other 1D potentials with open BC.**

$$\boxed{|N|^2 = \frac{m}{2\pi\hbar^2 k}} \quad \text{with} \quad k = \sqrt{\frac{2mE}{\hbar^2}}$$

Mathematically prove it for step barrier and linear potential with offset

Once **N** is known, ρ_{3D} can be computed !

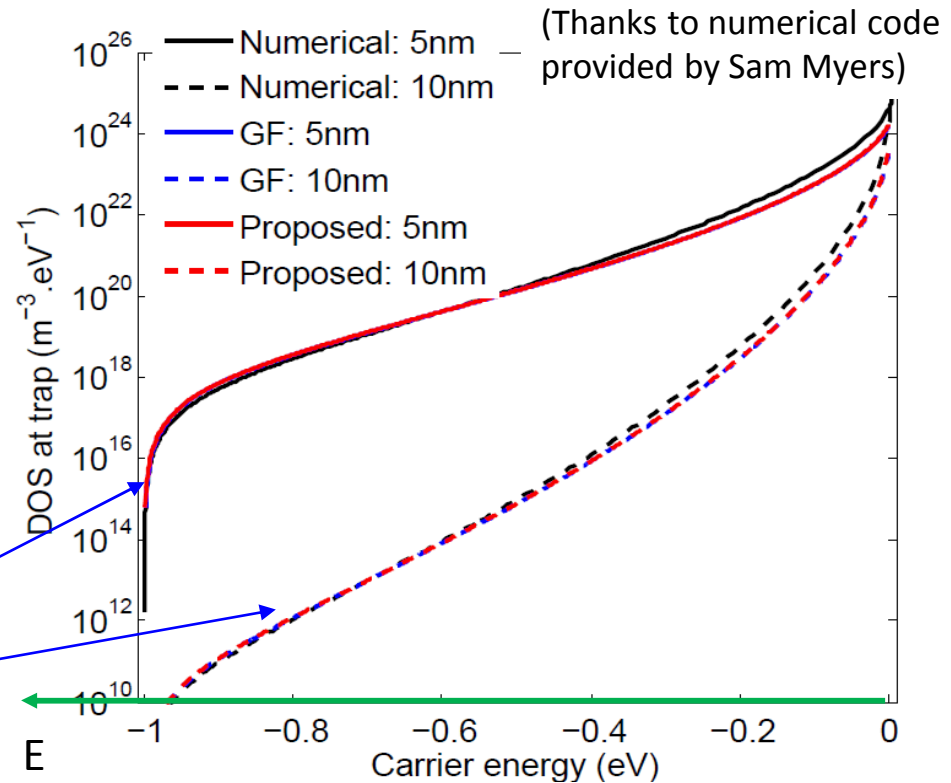
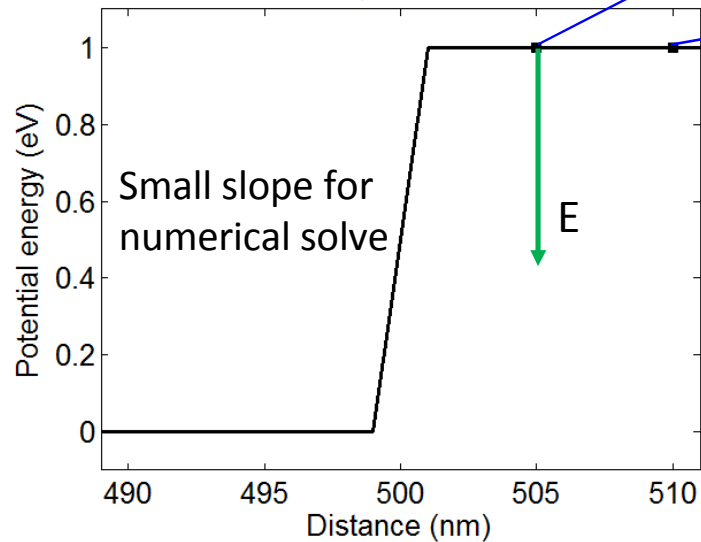
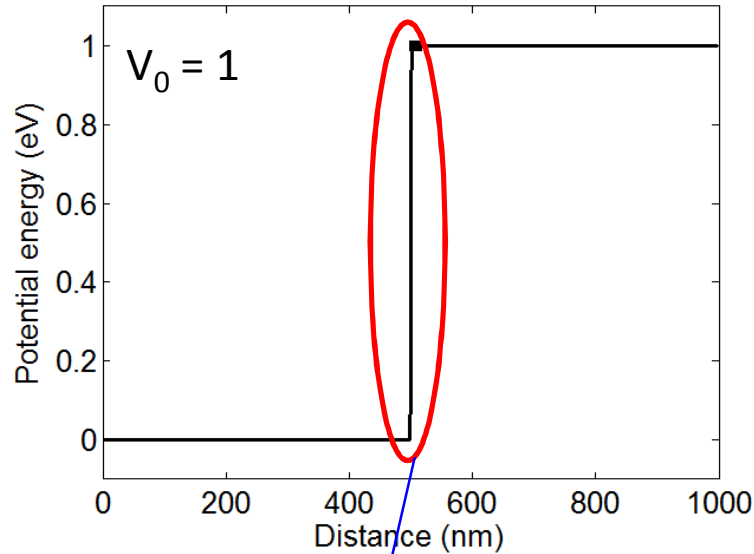
Pros:

- **Analytic** wave functions
- **Universal** normalization factor
- **Easy** implementation into TCAD

Cons:

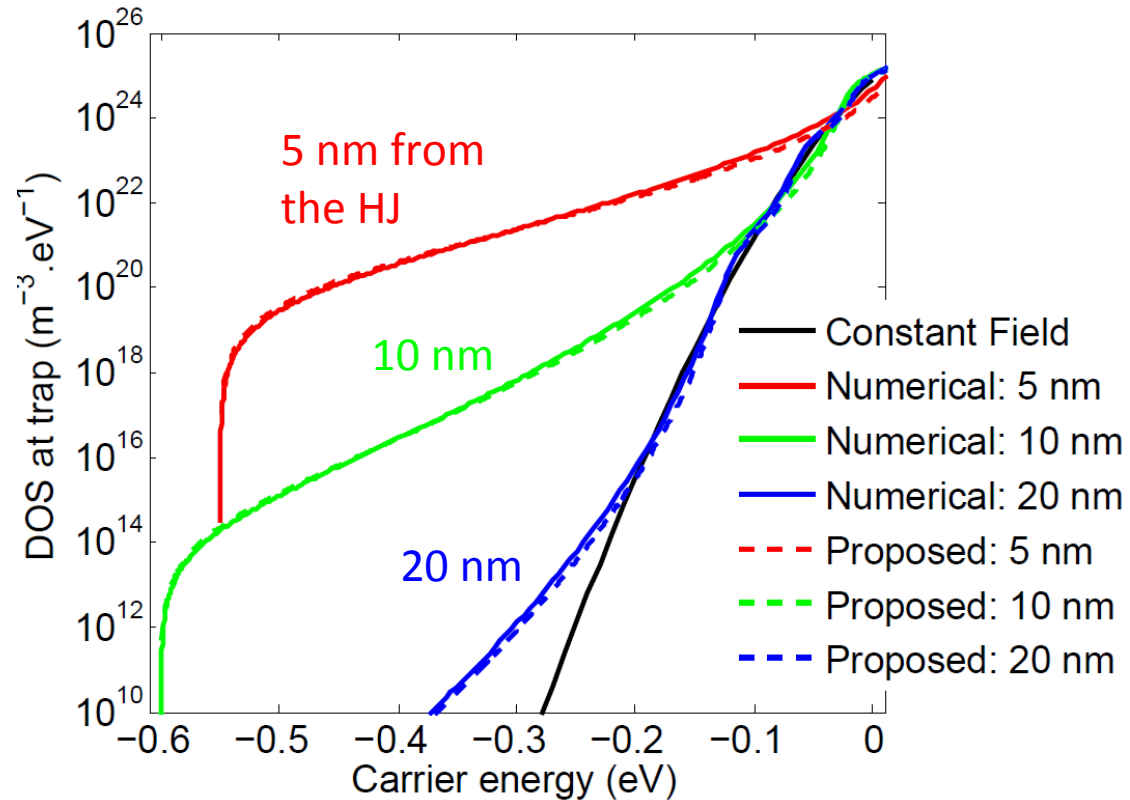
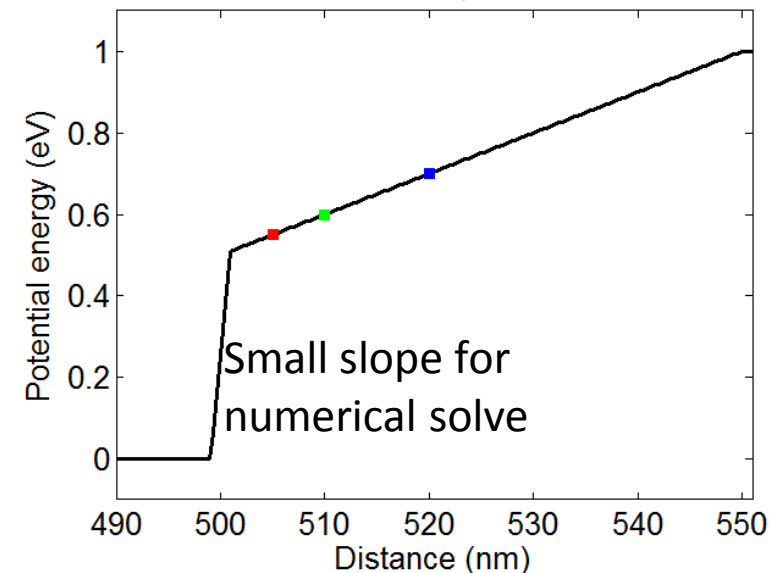
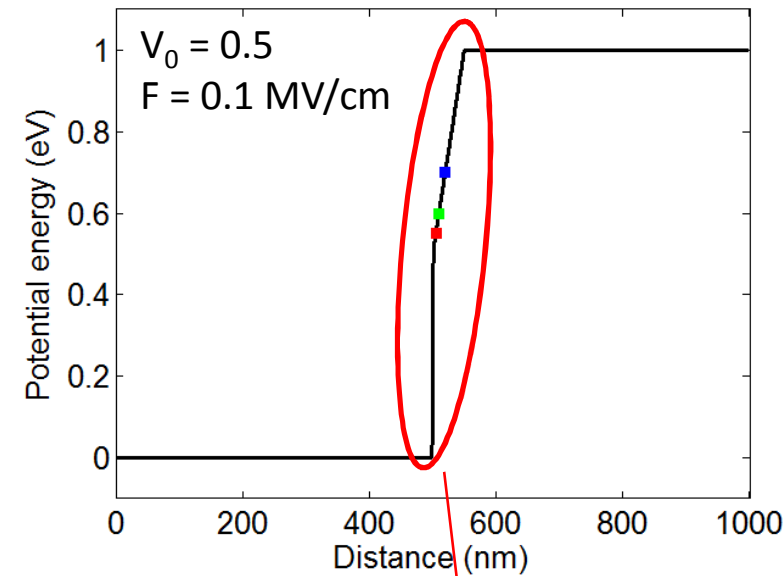
- Limited to 1D potentials that lead to piecewise analytic WFs

Density of States Comparison



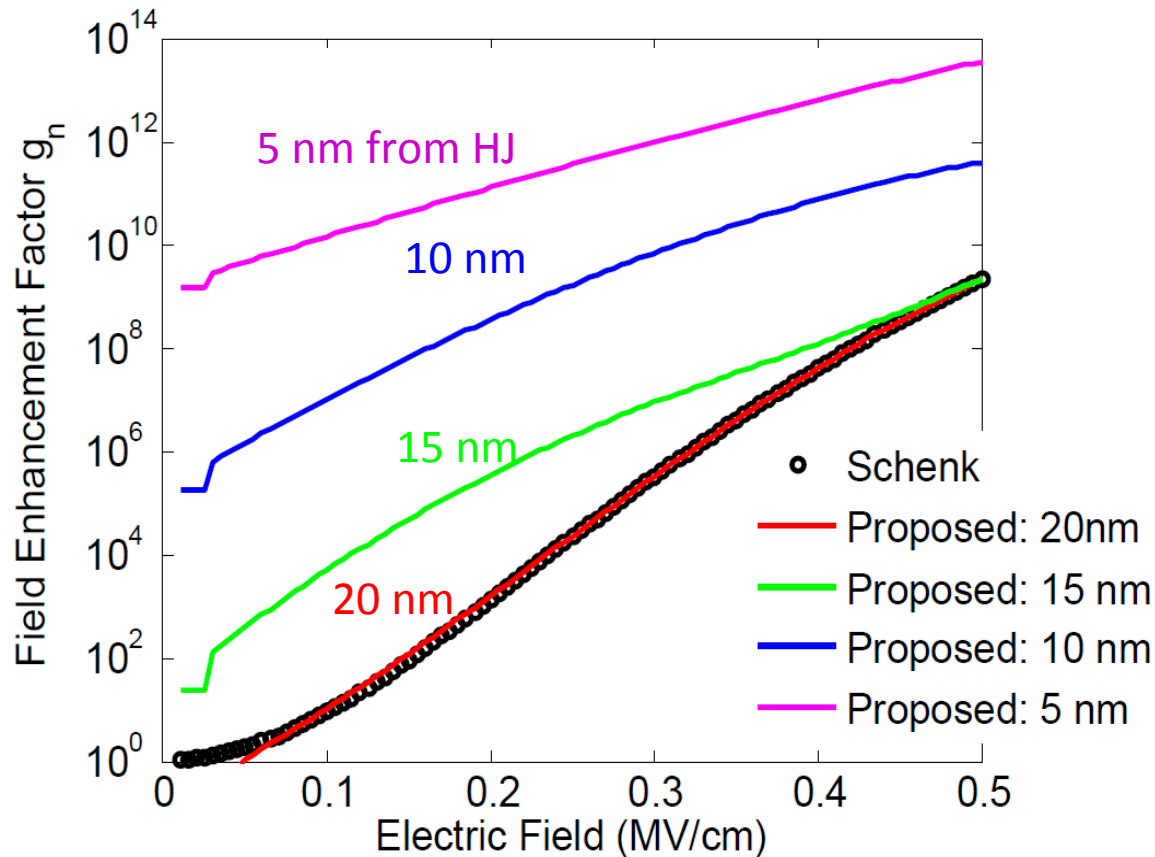
- The Green's function and proposed methods produce **the same DOS**
- Small difference between the numerical and proposed methods is due to the **non-ideal step** used in the numerical code

Density of States Comparison



- The constant-field DOS is **valid only for more than 20 nm** away from the heterojunction (HJ)
- The DOS computed using the proposed method **agree well** with numerical results
- The **band offset has a strong effect** on the DOS for distances within 20 nm from the HJ

Field Enhancement Factor g_n



- For locations ≥ 20 nm from the HJ, proposed & Schenk models produces **similar** g_n
- g_n by the proposed model **increases with decreasing distance** from the HJ
- Field dep. of g_n by the proposed model **reduces with decreasing distance** from the HJ

Charon Device Simulator

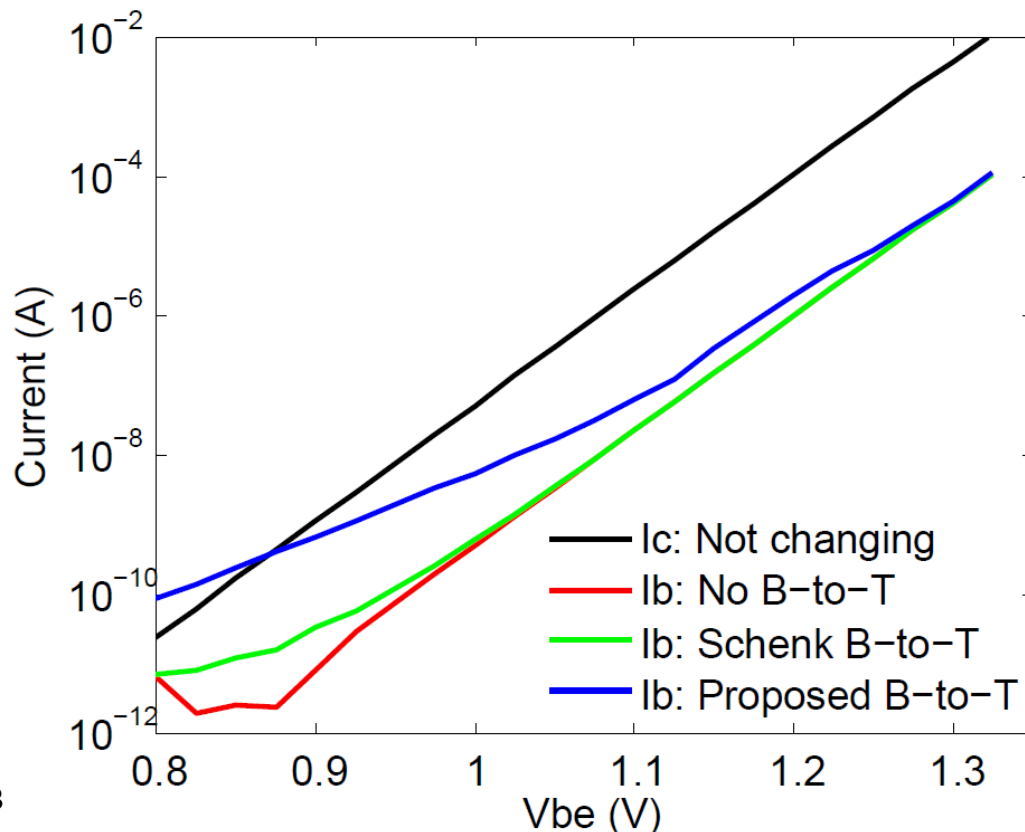
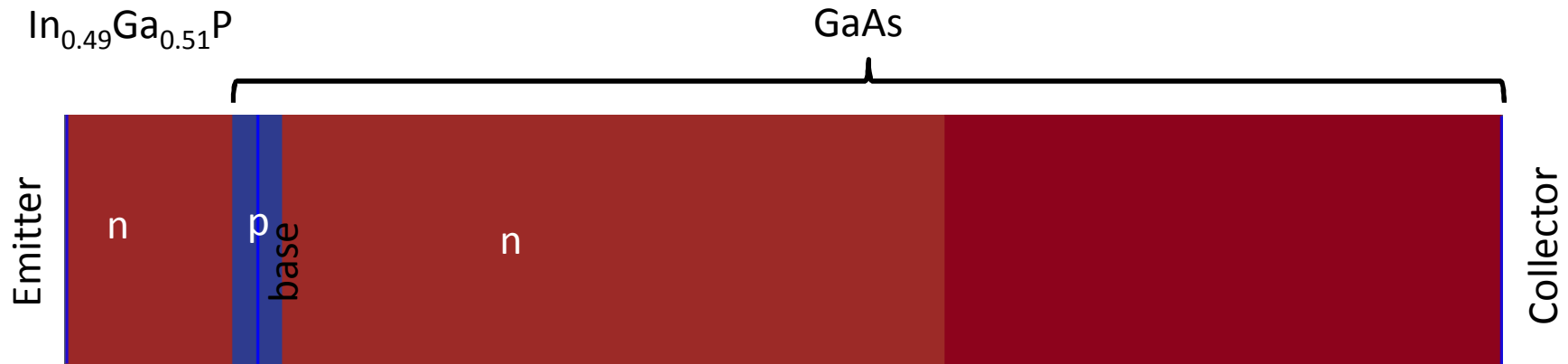
The proposed band-to-trap tunneling model is implemented in **Charon**.

Sandia-developed TCAD code with support for radiation effects modeling

Unique Capabilities provided by Charon

- **Open source** (just approved)
- Two & three dimensional + **MPI parallel** capability
- Various governing PDEs (Poisson, drift-diffusion for e/h/**ions**, lattice heating)
- **Different discretization** schemes (e.g., finite volume, finite element)
- Advanced physics models (Fermi-Dirac, transport across heterojunction, band-to-trap tunneling, etc.)
- Numerous devices (diodes, BJTs, HBTs, MOSFETs, **GaN** devices, **memristors**, etc.)

Application of B-to-T Model to NPN HBT

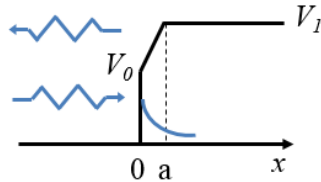


- Mid-band gap traps in the **emitter**
- $\tau_{n0} = 1 \text{ ns}$, $\tau_{p0} = 0.1 \text{ us}$
- Only hole-to-trap tunneling

The proposed b-to-t tunneling model **produces much larger base current** than the original Schenk model, due to the **hole-to-trap tunneling** enhanced by the emitter-base band offset.

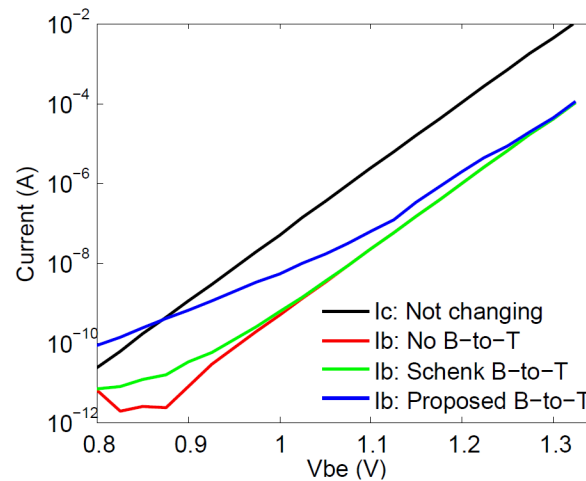
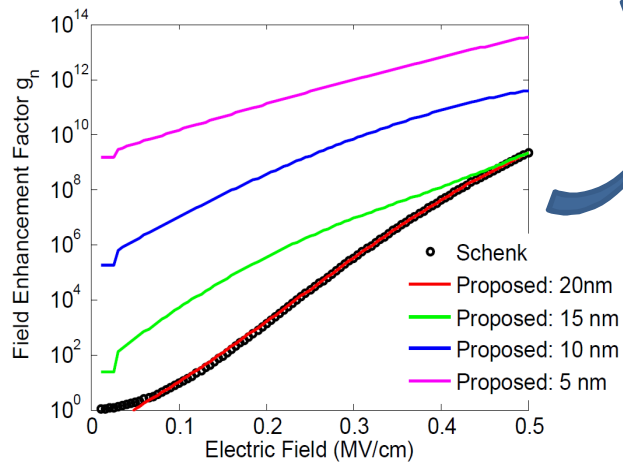
Conclusion

- ❑ Develop an **analytic DOS model** that includes the effects of **both electric field and HJ band offset** based on the open boundary scattering approach



$$\rho_{3D}(x, E) = \frac{m}{\pi \hbar^2} \int_0^{+\infty} |\psi_{E_x}(x)|^2 dE_x \theta(E - E_x)$$

- ❑ The proposed band-to-trap tunneling model can be **easily implemented into TCAD** codes thanks to its analytic form
- ❑ Demonstrate the **dramatically increased band-to-trap tunneling** strength (via the field enhancement factor) at **locations close to the HJ** using the proposed model



- ❑ Demonstrate the **dramatically increased base current** in a InGaP/GaAs NPN HBT due to the hole-to-trap tunneling **enhanced by the emitter-base junction band offset**

Backup Slides

Charon Device Simulator

