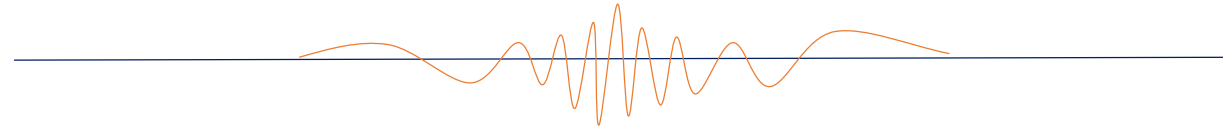


Vacancy and Mass-Impurity Phonon Scattering in Self-Irradiated Silicon



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¹University of Virginia, ²Sandia National Laboratories

Introduction

- Goal: Irradiate intrinsic Si with different isotopes to experimentally demonstrate the individual effects of mass-impurity and vacancy phonon scattering
 - Phonon: “Quantum of elastic waves” –Tamm, 1930 / Frenkel, 1932
- Si: second most abundant element in Earth’s crust, ubiquitous in electronics, thermal properties thoroughly characterized
- Stable Si isotopes:

	Isotope	Prevalence (%)
●	^{28}Si	92.22
●	^{29}Si	4.69
	^{30}Si	3.09

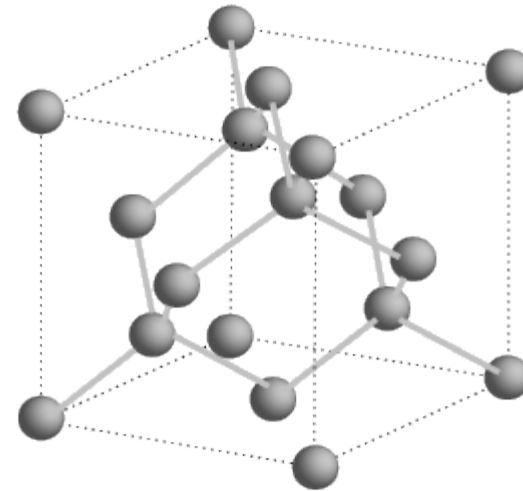
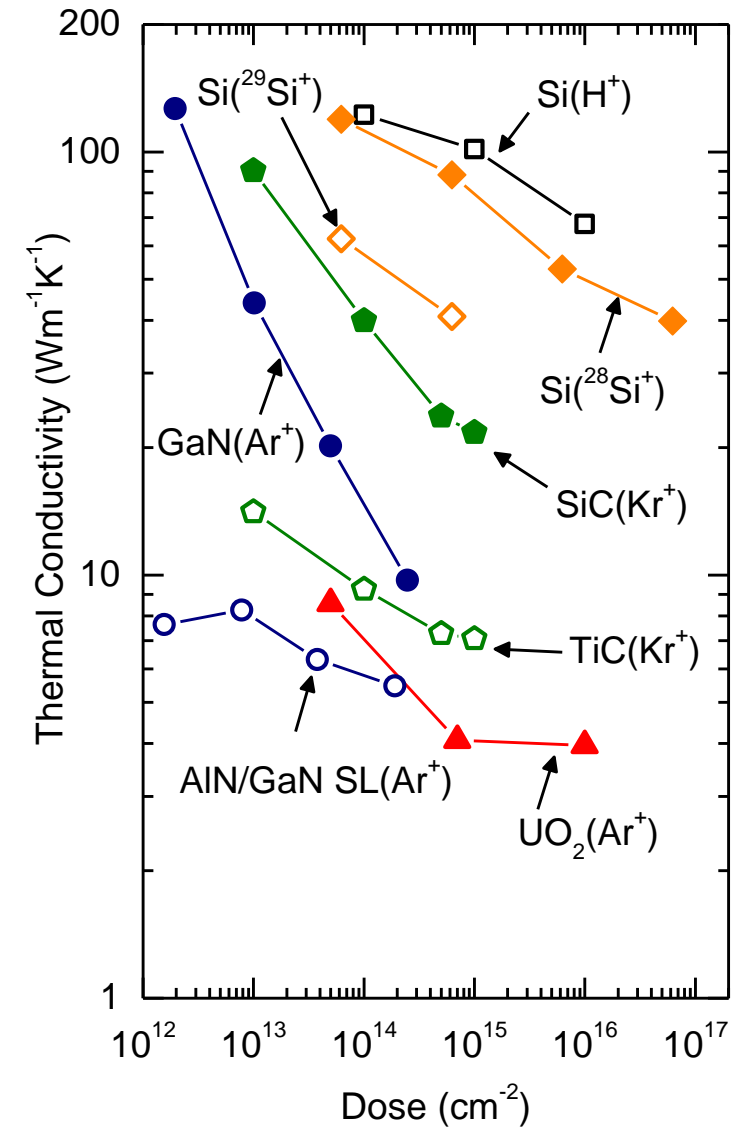
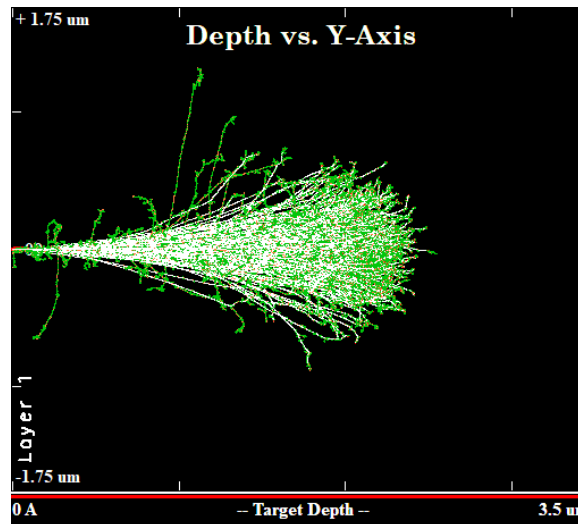
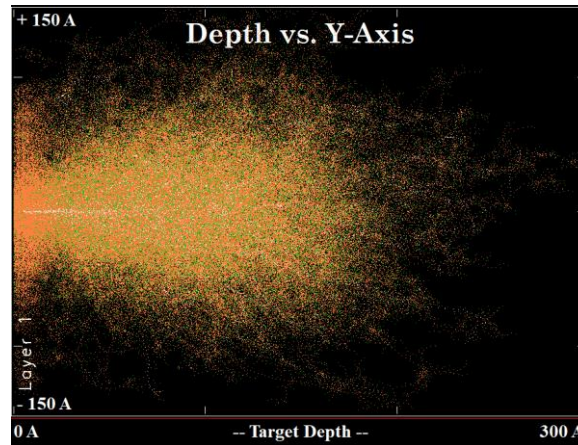
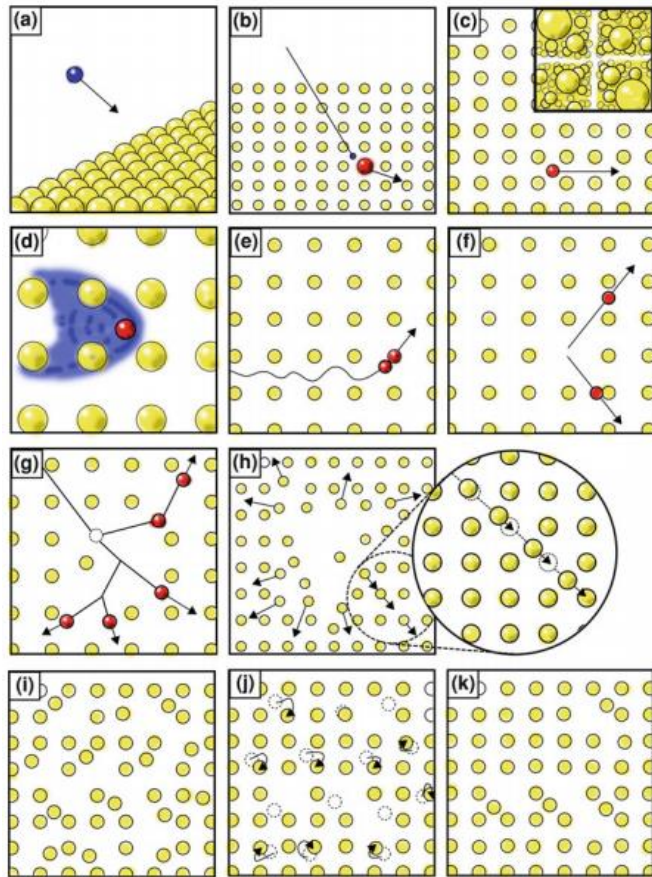


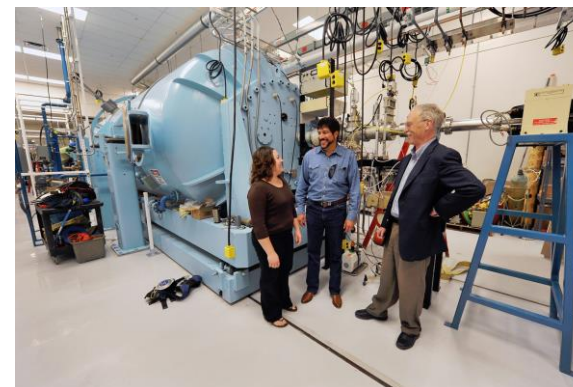
Image source: <http://hyperphysics.phy-astr.gsu.edu/hbase/Solids/sili2.html>

Ion Irradiation

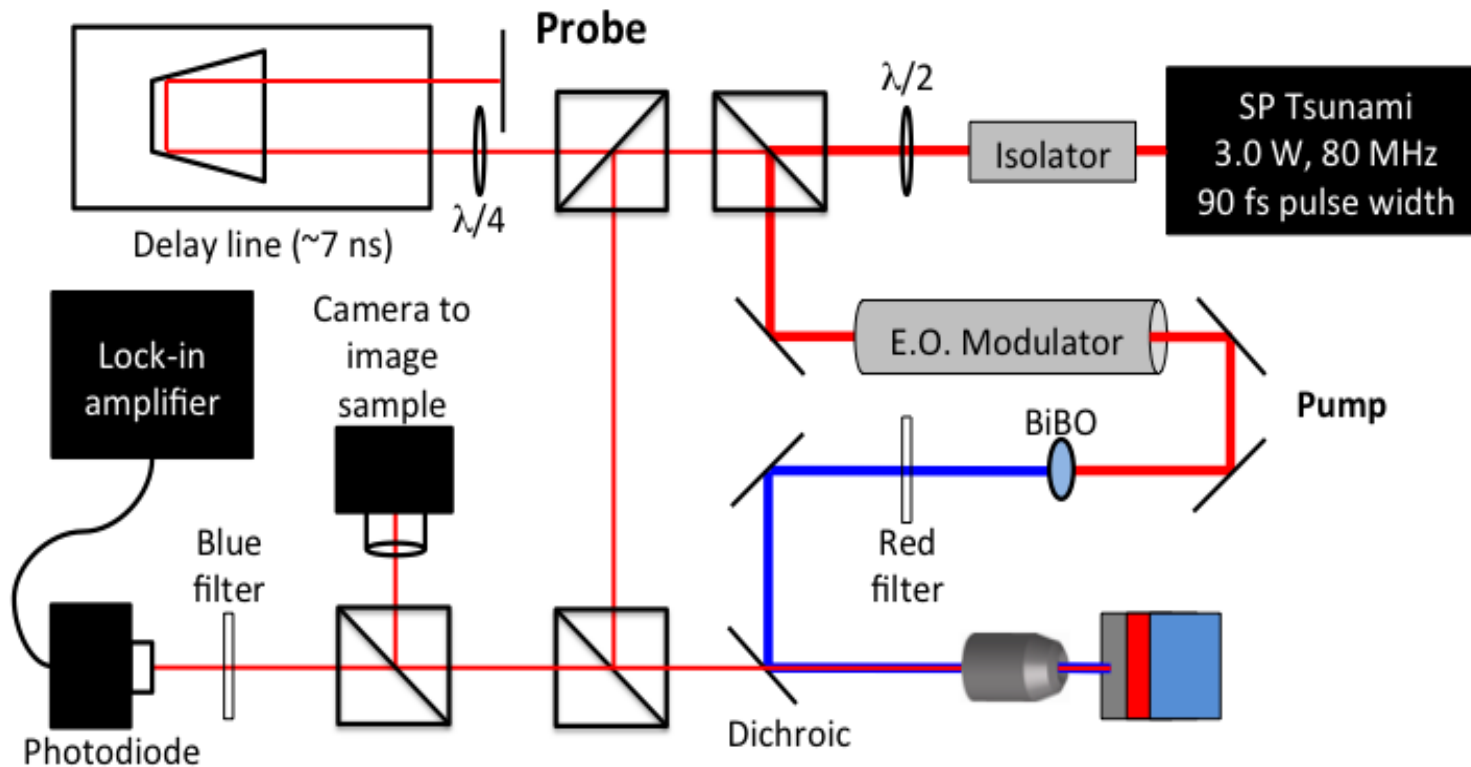


^{28}Si , ^{29}Si Irradiation

- Two sample sets:
 - ^{28}Si – Irradiated $\text{Si}_{\text{intrinsic}}$
 - Beam Energy: 3.75 [MeV]
 - Dose: $6.24\text{E}13$, $6.24\text{E}14$, $6.24\text{E}15$, $6.24\text{E}16$ [cm^{-2}]
 - Purpose: Primarily *vacancy* scattering
 - ^{29}Si – Irradiated $\text{Si}_{\text{intrinsic}}$
 - Beam energy: 3.75 [MeV]
 - Dose: $6.24\text{E}13$, $6.24\text{E}14$, $1.47\text{E}14$ [cm^{-2}]
 - Purpose: add in mass-impurity scattering
- Irradiation: 6 MV Tandem Van de Graaf-Pelletron Accelerator, Sandia National Laboratories

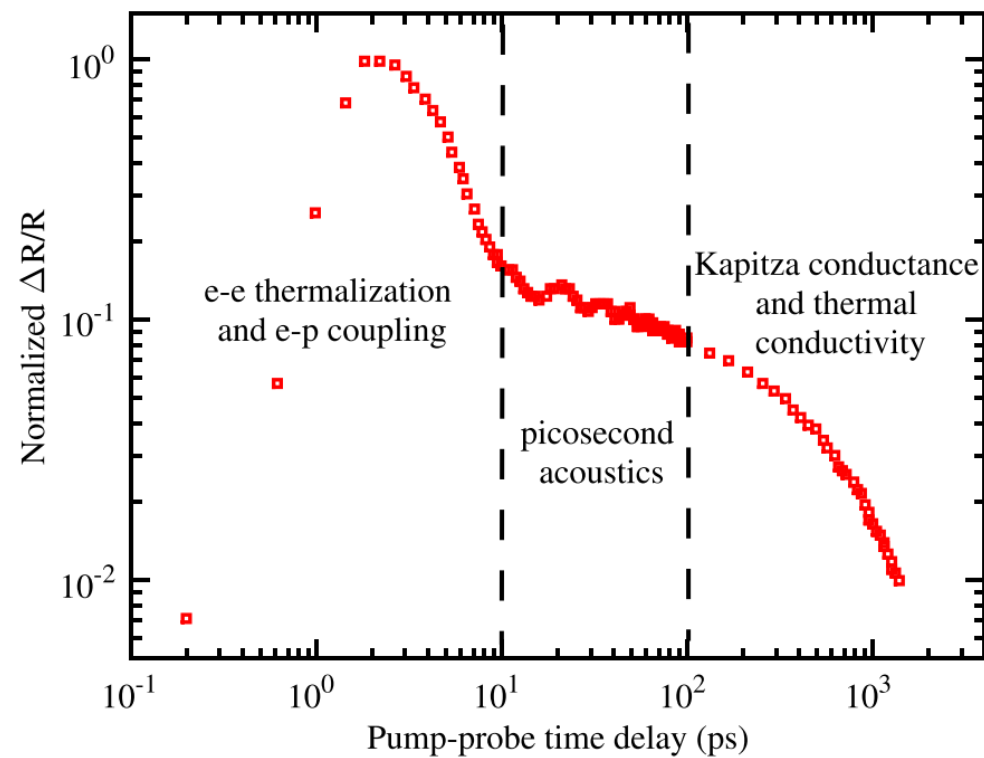
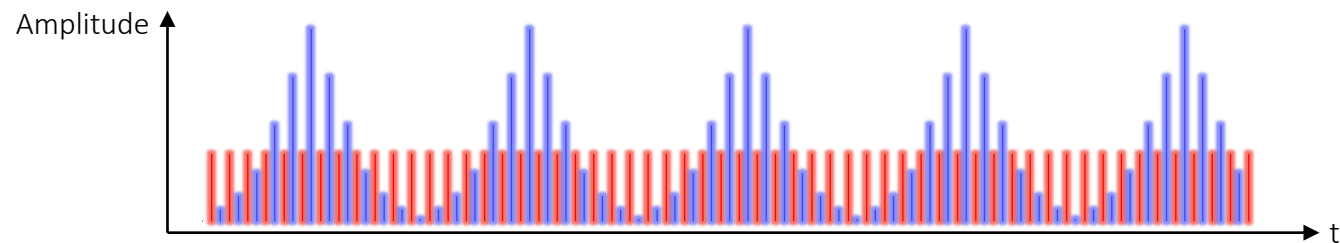
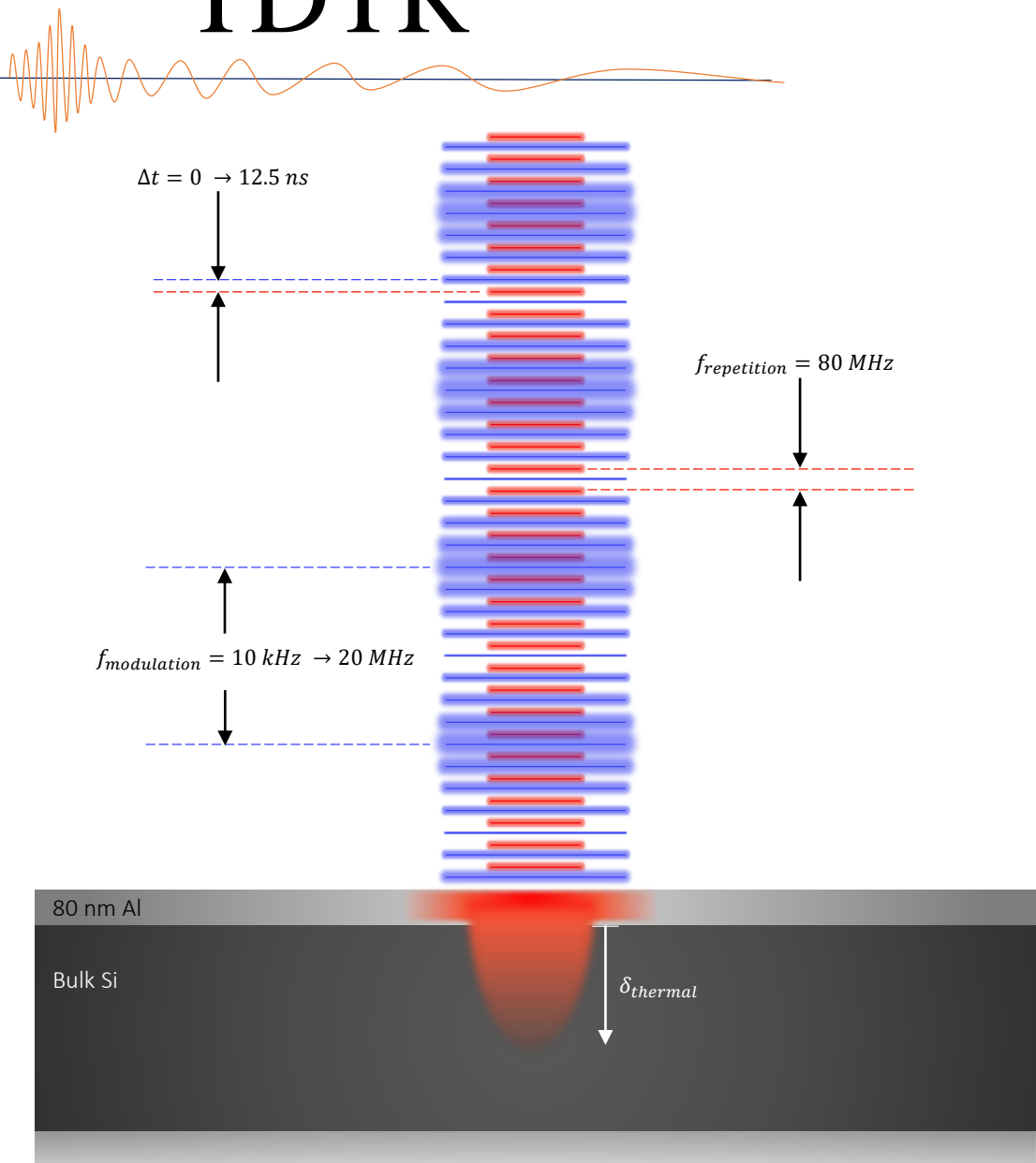


Time-Domain Thermoreflectance (TDTR)



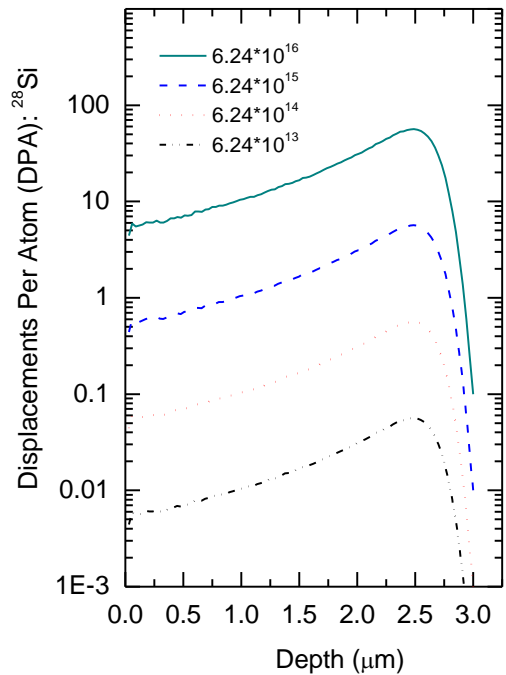
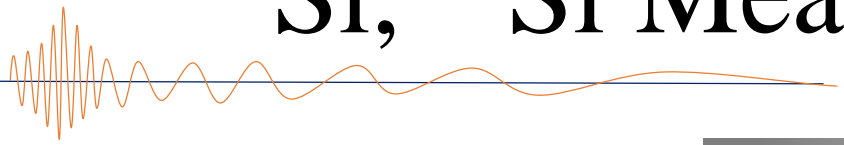
- Non-destructive, pump-probe technique
- Sub-picosecond temporal resolution
- Spot sizes as low as $\sim 1 \mu\text{m}$
- Modulation frequencies: 10 kHz – 20 MHz

TDTR

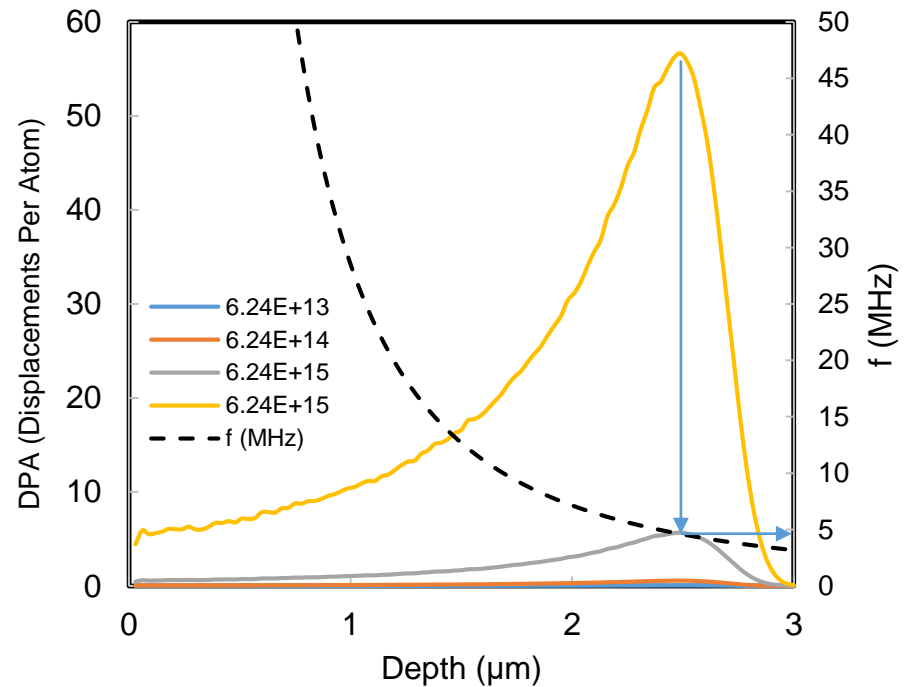


Giri et al., *Journal of Applied Physics*. 117, 105105 (2015)

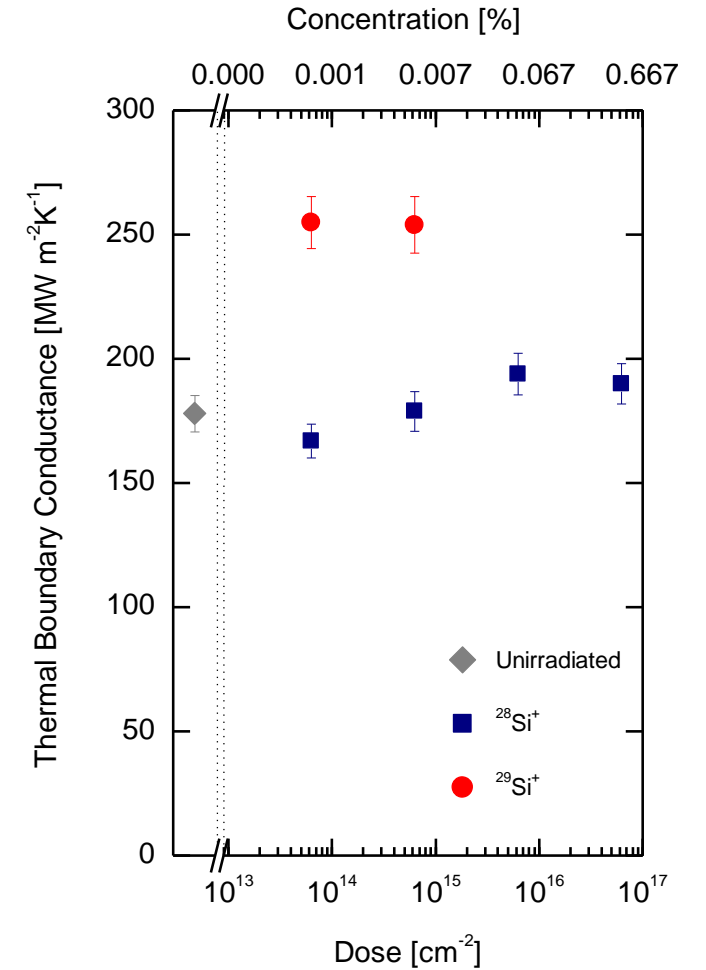
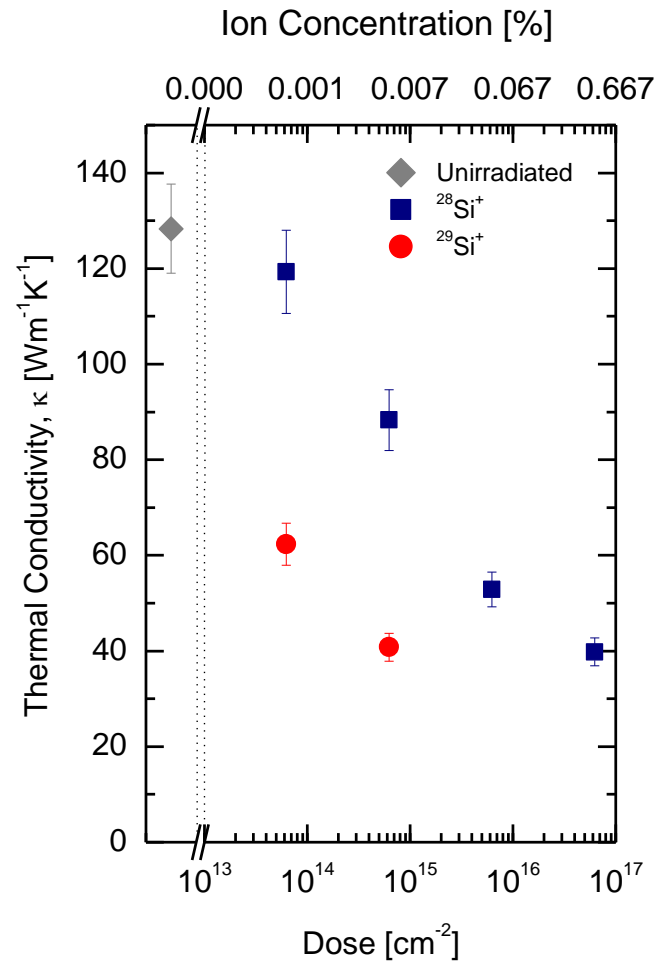
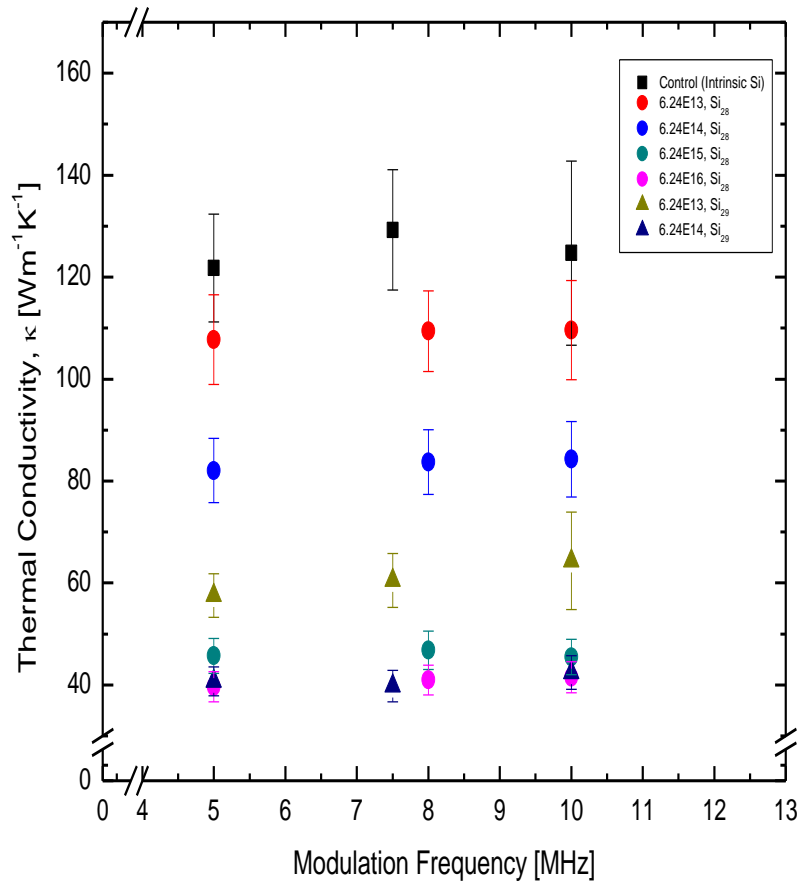
^{28}Si , ^{29}Si Measurement



$$\delta_{thermal} = \sqrt{\frac{\kappa}{\pi C f}}$$



Results



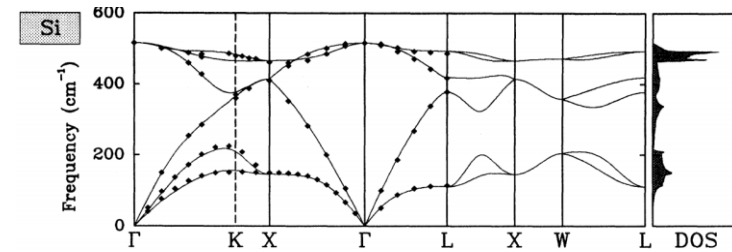
Thermal Conductivity

- From kinetic theory:

$$\kappa = \frac{1}{3} C_v v l = \frac{1}{3} C_v v^2 \tau$$

$$\kappa = \frac{1}{6\pi^2} \sum_j \int \hbar \omega_j(k) k^2 \frac{\partial f}{\partial T} v_j^2(k) \tau_j(k) dk$$

- All parameters can be assumed from literature dispersion of Si
- Scattering time, τ is unique from sample to sample



P. Giannozzi, S. De Gironcoli, P. Pavone, and S. Baroni, Phys. Rev. B **43**, 7231 (1991).

First, determine baseline scattering:

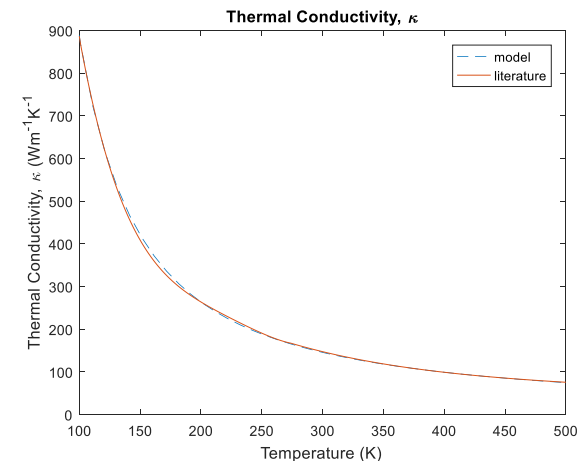
$$\frac{1}{\tau} = \frac{1}{\tau_{ph}} + \frac{1}{\tau_{Imp}}$$

- Phonon-phonon (Umklapp):

$$\tau_{ph}^{-1} = BT\omega(k)e^{-\frac{C}{T}}$$

- Baseline impurity:

$$\tau_{Imp}^{-1} = A\omega(k)^4$$



Disorder Scattering – Klemens Model

- Phenomenological model developed in 1955, extensively used for thermal conductivity modeling

$$\tau_{Def}^{-1} = \frac{\omega(k)^4 \delta^3 \Gamma_i}{4\pi v_j(k)} \rightarrow \text{Def can be vacancies or mass – impurities}$$

- Scattering cross section of impurity, i :

$$\Gamma_i = x_i \left[\left(\frac{\Delta M_i}{M} \right)^2 + 2 \left(\left(\frac{\Delta G_i}{G} \right) - 2 * 3.2\gamma \left(\frac{\Delta \delta_i}{\delta} \right) \right)^2 \right]$$

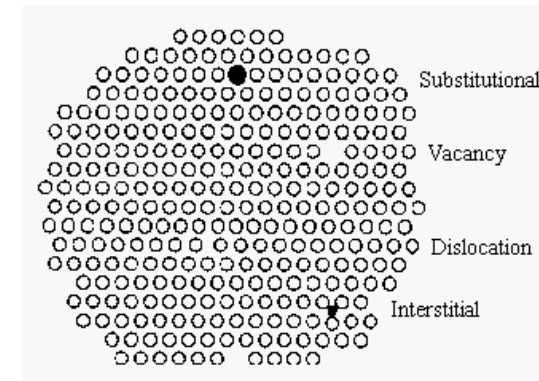
- x_i : defect concentration
- M : mass of lattice atom
- G : stiffness constant of nearest neighbor atoms
- δ : atomic radii
- We expect two kinds of impurities → vacancies, interstitials
- Since we self-irradiated: $\Delta G_i \approx \Delta \delta_i \approx 0$
- For $M_{Vac} = M_{Host}$: $\Gamma_i = x_i [(-3)^2 + 0]$
- But what is x_i ?? Cannot Predict with TRIM

The Scattering of Low-Frequency Lattice Waves by Static Imperfections

By P. G. KLEMENS

Division of Physics, Commonwealth Scientific and Industrial Research Organization,
Sydney

Communicated by G. H. Briggs; MS. received 7th December 1954 and in amended form
12th July 1955

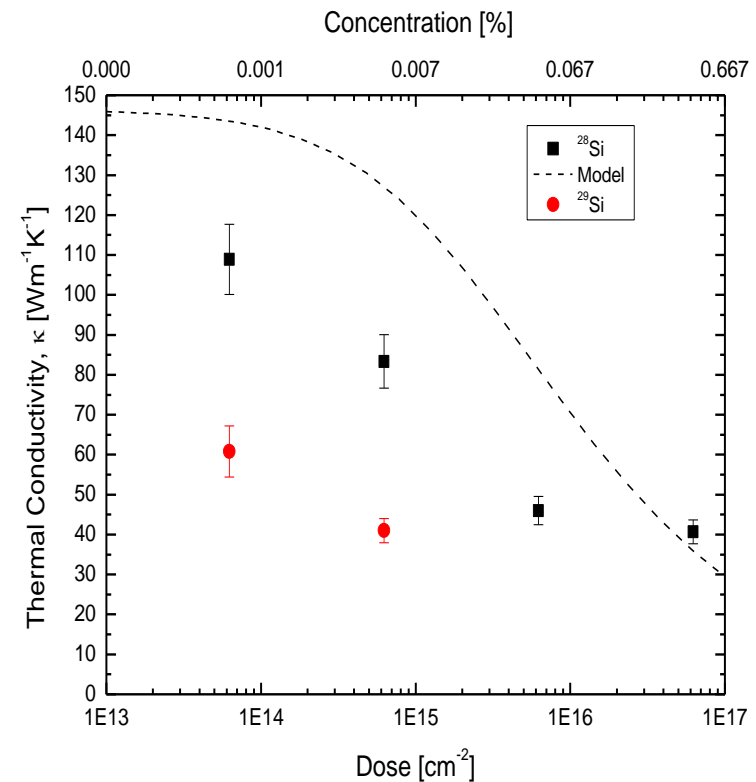
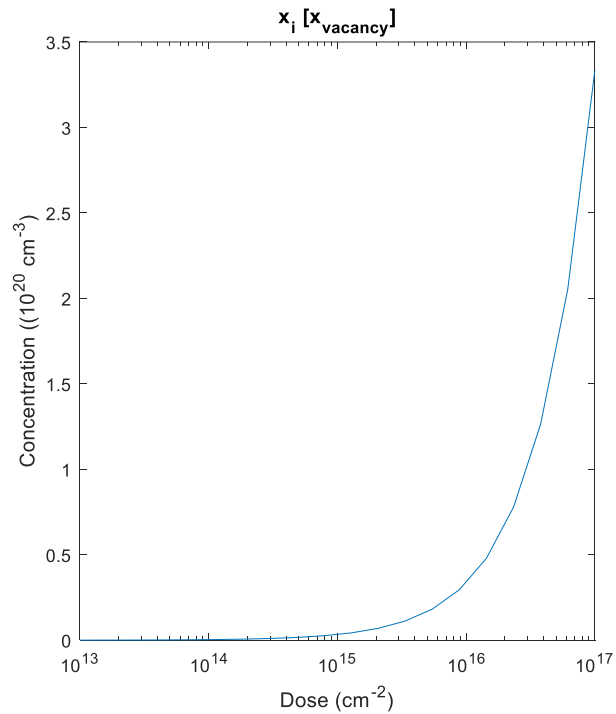


Defect Concentration, x_i

- Ion concentration is known, vacancy concentration is not.
 - If a 1:1 vacancy to ion concentration is assumed:

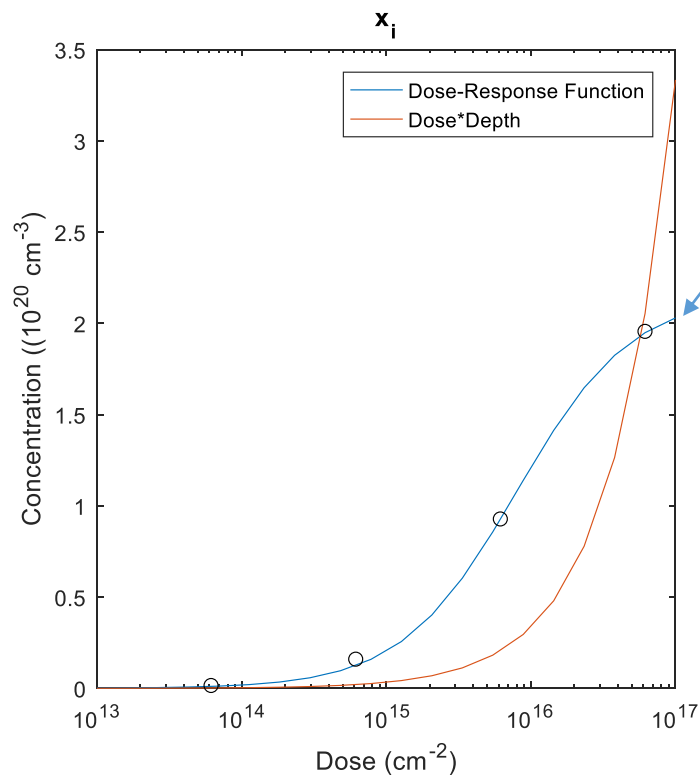
$$\tau_{Def}^{-1} = \frac{\omega(k)^4 \delta^3 \Gamma_i}{4\pi v_j(k)}$$

$$\Gamma_i = x_i \left[\left(\frac{\Delta M_i}{M} \right)^2 + 2 \left(\left(\frac{\Delta G_i}{G} \right) - 2 * 3.2\gamma \left(\frac{\Delta \delta_i}{\delta} \right) \right)^2 \right]$$



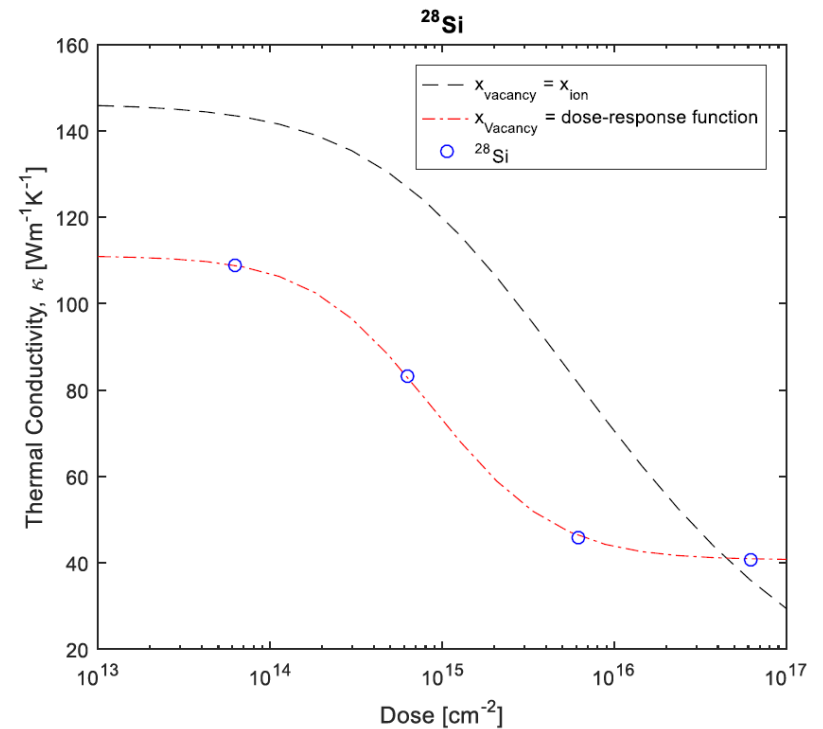
Defect Concentration, x_i

- Instead, calculate x_i for each measured thermal conductivity – then, fit to the calculated values:
 - “Dose Response Function”

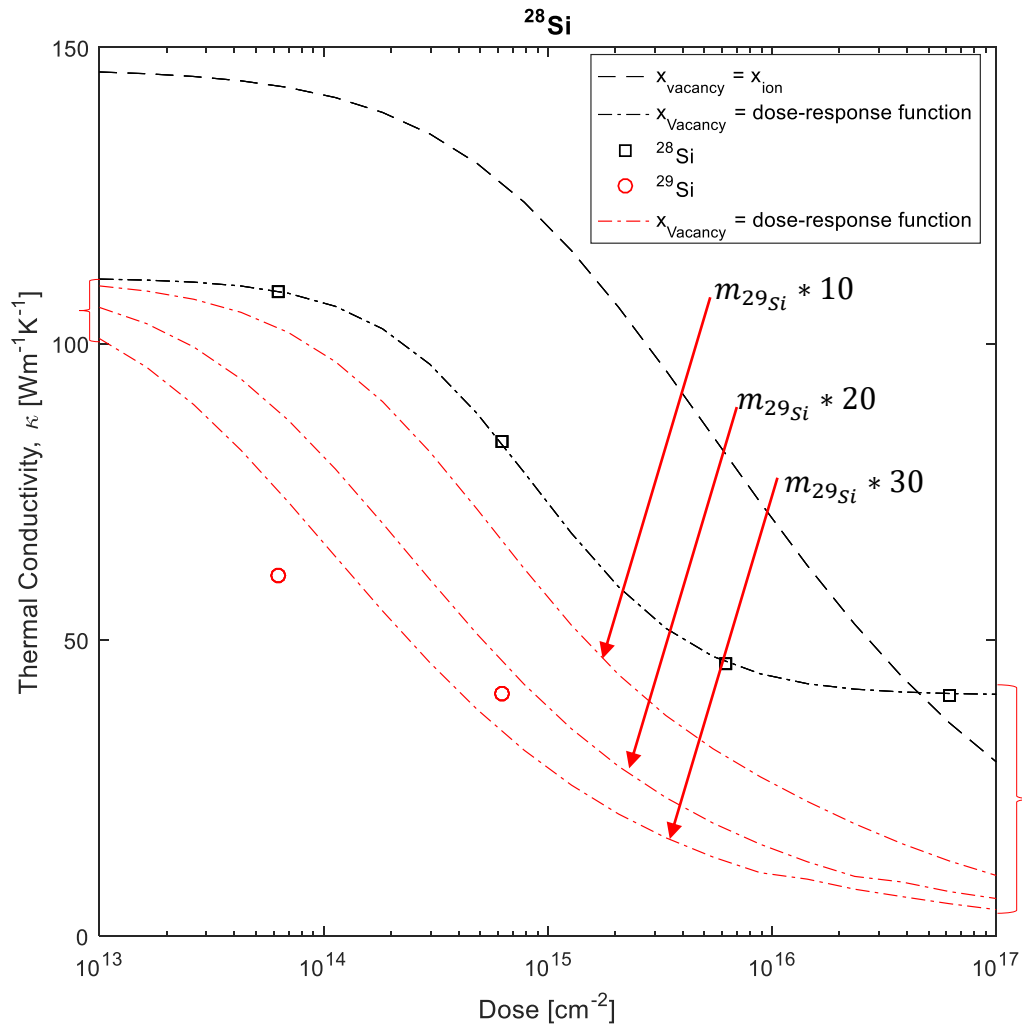


$$x_i = \frac{b-a}{1+\left(\frac{c}{x}\right)^d}$$

- a = min. asymptote
- b = max. asymptote
- c = middle dose value
- d = slope



Mass-Impurity Contribution

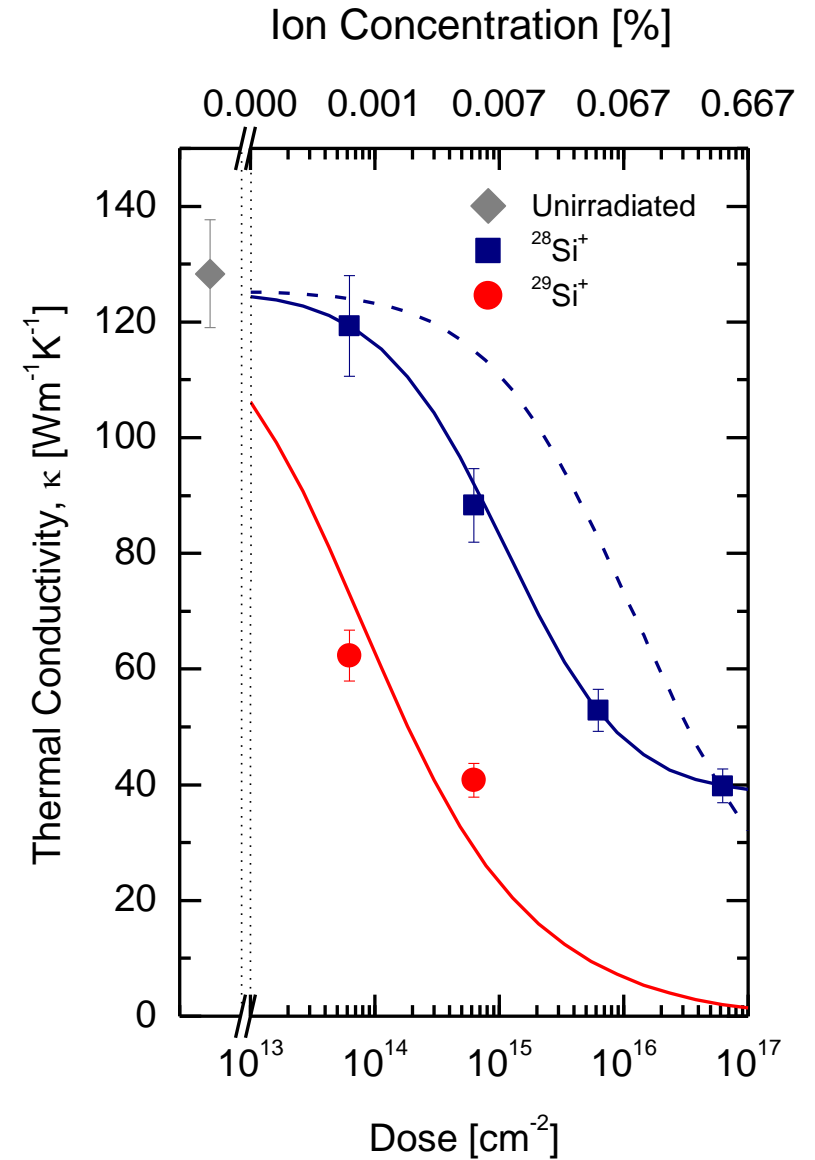


- Clearly, simply adjusting the mass-impurity term does not capture the trend in the data

- Also consider dislocations and dislocation cores (ω, ω^3)

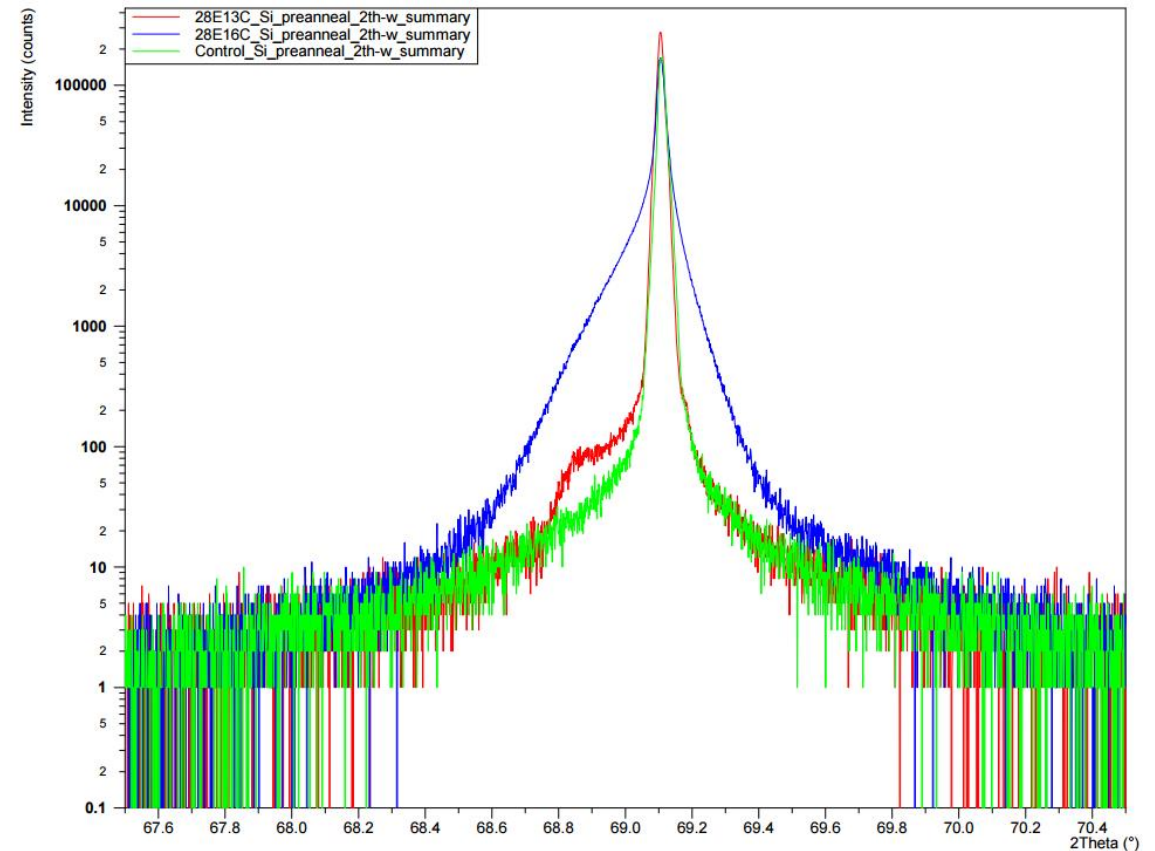
Assume:

$$\tau_{\text{Def}}^{-1} = x_{\text{ion}}(A\omega^4 + B\omega^2 + C\omega^3 + D\omega)$$



Strain, Amorphization

- Thermal conductivity model still not within uncertainty of model
 - Disparity due to presence of amorphous regions
 - XRD used to characterize amorphization
- Clear trends are evident with increasing dose
- ^{28}Si :
 - No dose – clean, single peak
 - Low dose ($6.24\text{E}13$) – secondary peak formation indicating strain
 - High dose ($6.24\text{E}16$) – significant peak broadening, indicative of amorphization of sample



Summary / Future Work

- Significant decrease in thermal conductivity when irradiating with heavier isotope
- Decrease cannot solely be attributed to mass differences as predicted by Klemens model – must be a change in vacancy concentration / amorphization
 - Annealing and XRD to yield further insight
- Thermal boundary conductance increases with dose
 - Shown previously in literature (Gorham, et. al) – attributed to spatial gradation of native oxide into Si, leading to region of average vibrational spectra
 - Further testing needed to verify that oxide mixing is responsible for increase

