

Application of peridynamics to large deformations and soft materials

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ENERGY

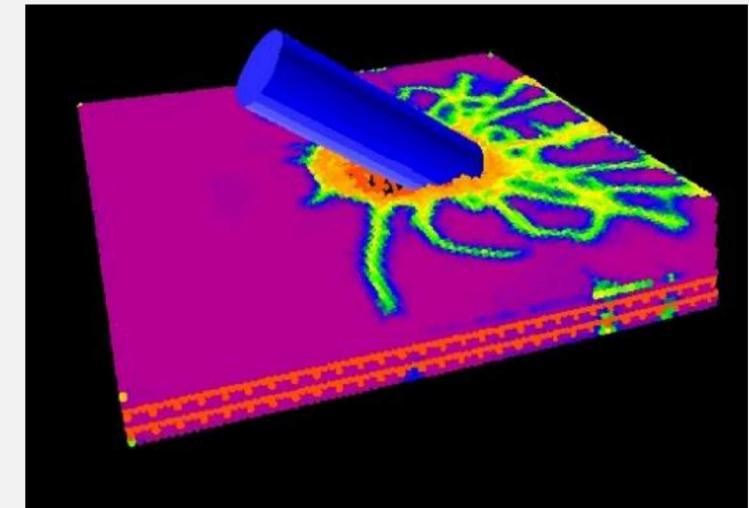


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Outline

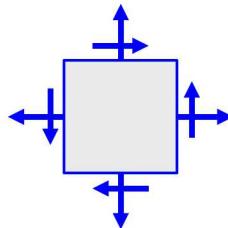
- Peridynamics background
 - Damage and fracture
- Geometric nonlinearity
 - Rubbery materials
- Thermodynamics and long-range forces
- Eulerian material models
 - Fluids and surface tension
 - Contact, friction, and wear
- Combining Lagrangian and Eulerian models
 - Gelatin
 - Bird strike simulant



Traditional application of peridynamics:
Elastic-brittle material

Traditional solid mechanics

- The traditional mathematical model for solids and structures uses partial differential equations (PDEs):



$$\frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} + b_x = 0$$
$$\frac{\partial \sigma_{yx}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + b_y = 0$$

where $\sigma_{xx} \dots$ are the stress components and b_x, b_y are the external loads.

- Is this model up to the job of predicting material failure? Key assumptions:

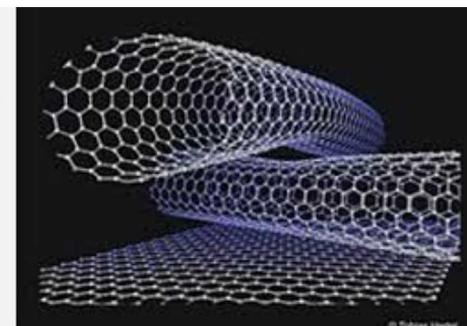
- Contact forces
- Continuity

NO!



Fracture and fragmentation

NO!



Nanoscale structures
and metamaterials



Augustin-Louis
Cauchy, 1840

These issues affect everything we “do”

- The standard PDEs are incompatible with the essential physical nature of cracks.
 - Can't apply PDEs on a discontinuity.
- Typical FE approaches implement a fracture model after numerical discretization.
 - Need supplemental kinetic relations that are understood only in idealized cases.

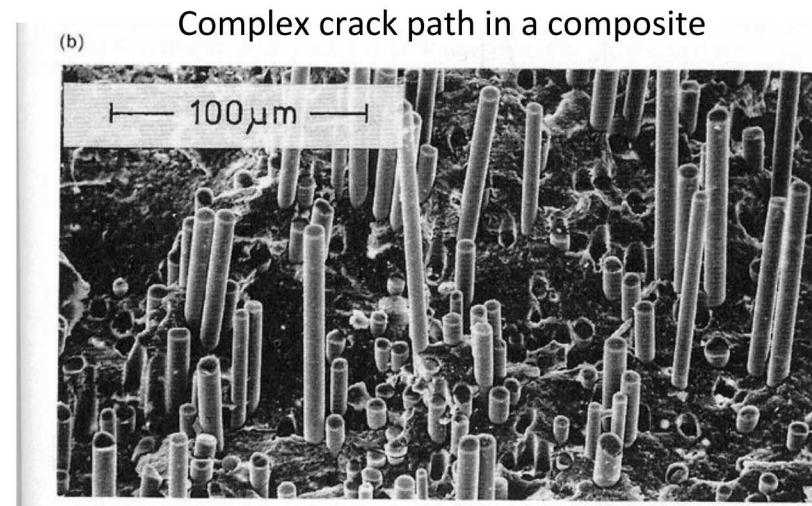
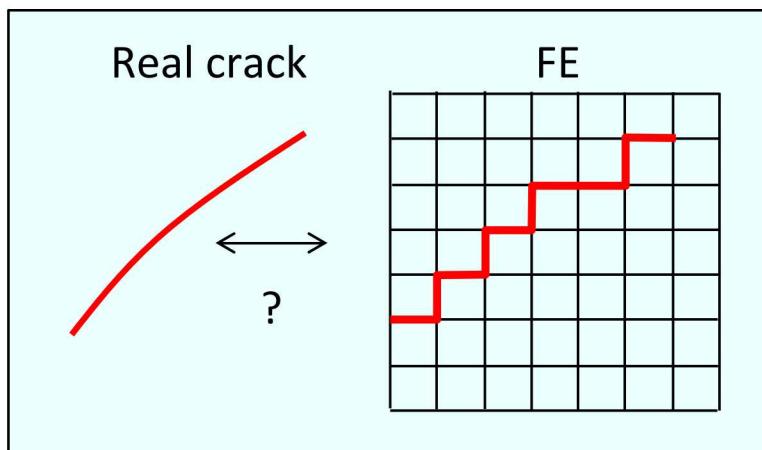
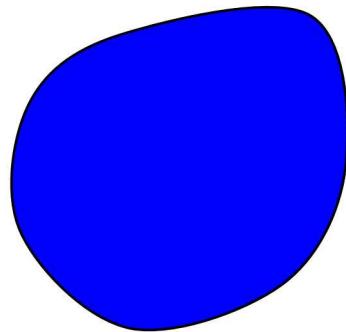


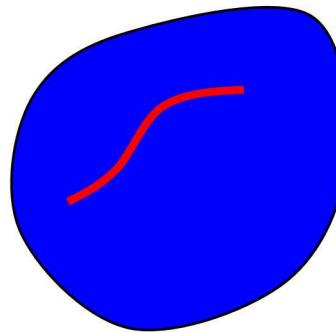
Figure 11.20 Pull-out: (a) schematic diagram; (b) fracture surface of ‘Silceram’ glass-ceramic reinforced with SiC fibres. (Courtesy H. S. Kim, P. S. Rogers and R. D. Rawlings.)

Purpose of peridynamics*

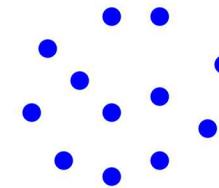
- To unify the mechanics of continuous and discontinuous media within a single, consistent set of equations.



Continuous body



Continuous body
with a defect



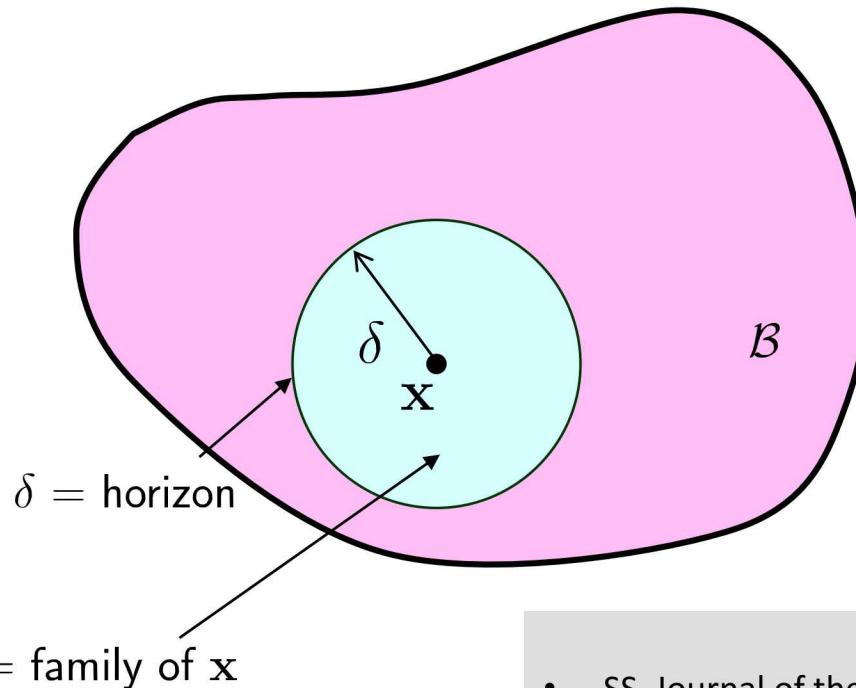
Discrete particles

- Why do this?
 - The standard theory (Cauchy, 1827) doesn't always meet the needs of technology.
 - Model complex fracture patterns.
 - Communicate across length scales.

* Peri (near) + dyn (force)

Peridynamics basics: Horizon and family

- Any point x interacts directly with other points within a distance δ called the “horizon.”
- The material within a distance δ of x is called the “family” of x , \mathcal{H}_x .

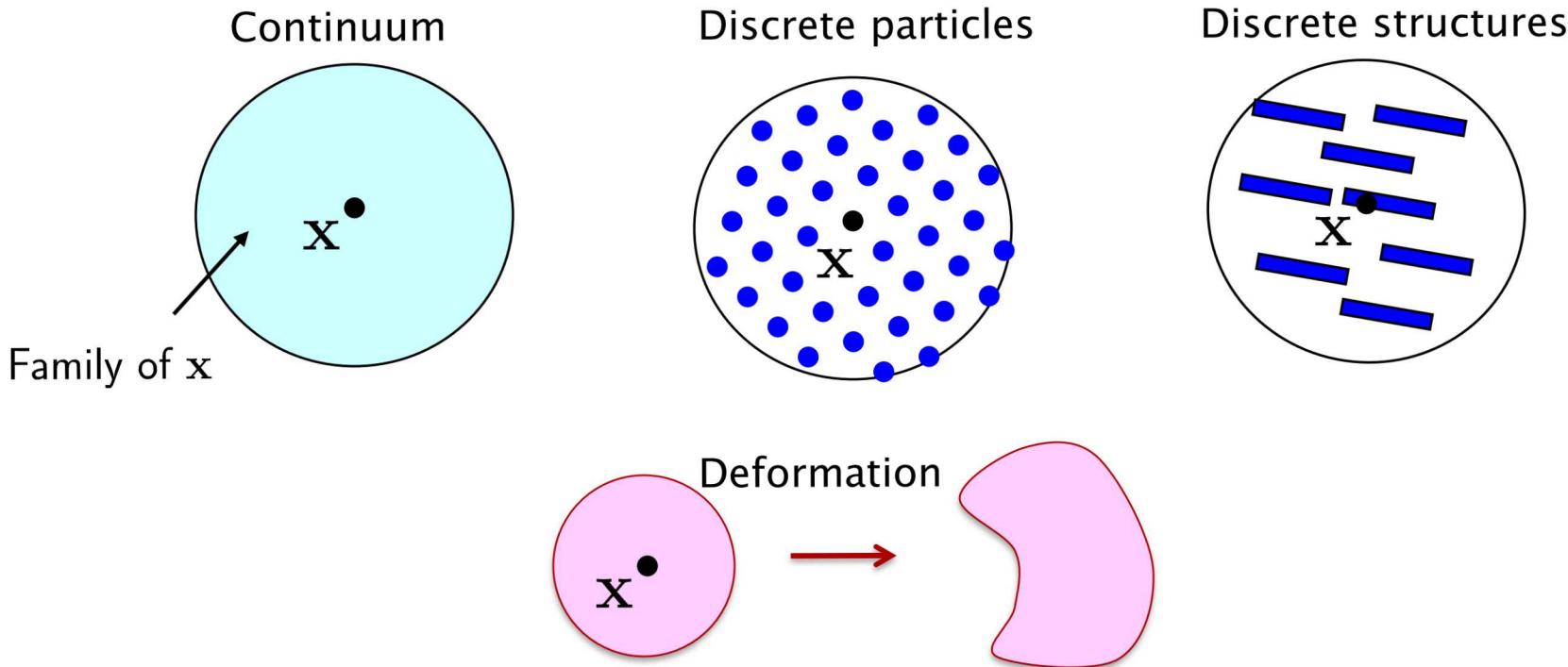


General references

- Silling, *Journal of the Mechanics and Physics of Solids* (2000)
- Silling and R. Lehoucq, *Advances in Applied Mechanics* (2010)
- Madenci & Oterkus, *Peridynamic Theory & Its Applications* (2014)

Point of departure:

Strain energy at a point



- Key assumption: the strain energy density at x is determined by the deformation of its family.

Equation of motion

- Write down the total potential energy in a body under the deformation \mathbf{y} :

$$\Phi = \int_{\mathcal{B}} (W - \mathbf{b} \bullet \mathbf{y}) \, dV.$$

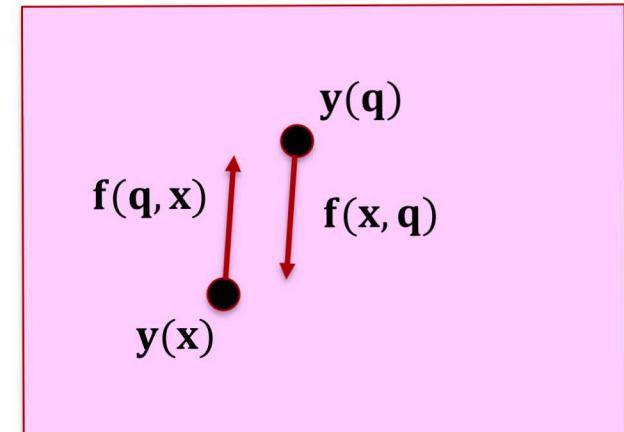
- The Euler-Lagrange equation is

$$0 = \int_{\mathcal{H}_x} \mathbf{f}(\mathbf{q}, \mathbf{x}) \, dV_{\mathbf{q}} + \mathbf{b}(\mathbf{x}),$$

which is the equilibrium equation of peridynamics.

- \mathbf{f} is the pairwise bond force density, which comes from the material model.
- Dynamics:

$$\rho \ddot{\mathbf{y}}(\mathbf{x}, t) = \int_{\mathcal{H}_x} \mathbf{f}(\mathbf{q}, \mathbf{x}, t) \, dV_{\mathbf{q}} + \mathbf{b}(\mathbf{x}, t).$$



Keeping track of the collective deformation of families: States

- A state is a mapping from bonds in a family to some other quantity. We write

$$\underline{A}[\mathbf{x}] \langle \mathbf{x}' - \mathbf{x} \rangle = \text{something.}$$

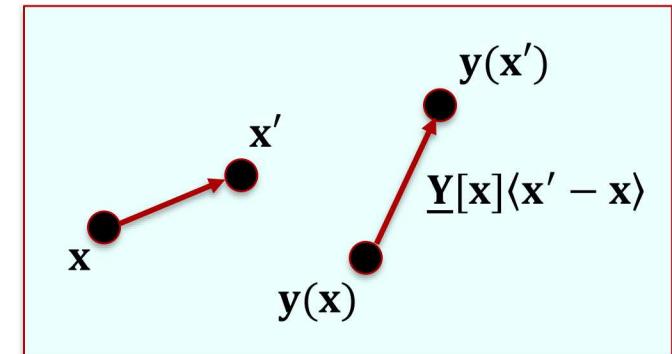
- The *deformation state* maps each bond to its deformed image:

$$\underline{Y}[\mathbf{x}] \langle \mathbf{x}' - \mathbf{x} \rangle = \mathbf{y}(\mathbf{x}') - \mathbf{y}(\mathbf{x})$$

where \mathbf{y} is the deformation.

- Dot product of two states:

$$\underline{A} \bullet \underline{B} = \int_{\mathcal{H}} \underline{A}(\xi) \underline{B}(\xi) dV_{\xi}.$$

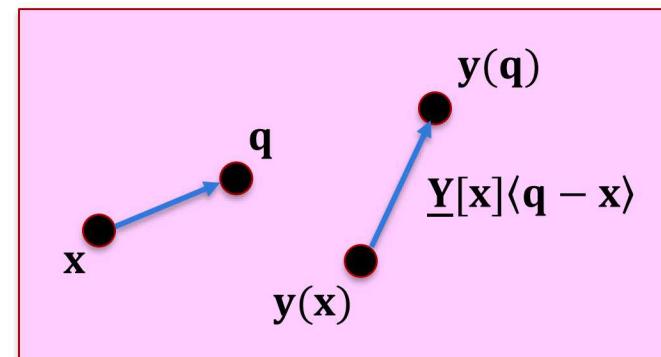
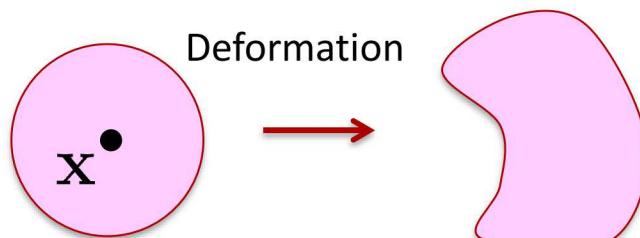


Energy density depends on the deformation state

- Strain energy at \mathbf{x} :

$$W(\underline{\mathbf{Y}}[\mathbf{x}]).$$

- Next figure out the bond forces.



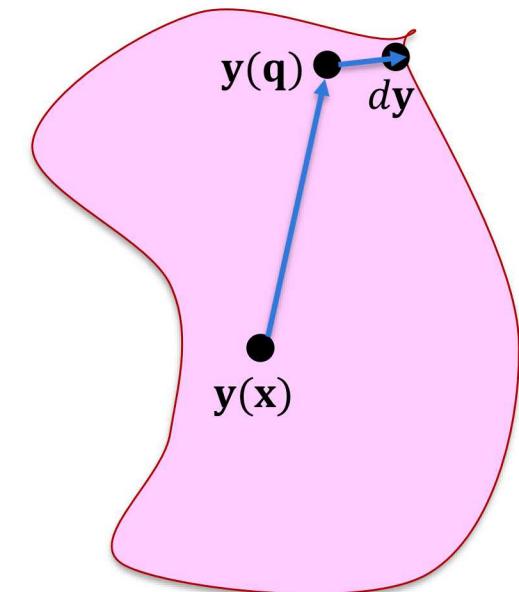
Bond forces from strain energy: Fréchet derivatives

- Perturb the deformed point $y(\mathbf{q})$ by a small additional displacement displacement $d\mathbf{y}$.
- The resulting change in strain energy is dW .
- The Fréchet derivative $W_{\underline{\mathbf{Y}}}$ is the state such that

$$dW = W_{\underline{\mathbf{Y}}} \bullet d\underline{\mathbf{Y}}.$$

- The bond forces are found from

$$\mathbf{f}(\mathbf{q}, \mathbf{x}) = W_{\underline{\mathbf{Y}}}[\mathbf{x}] \langle \mathbf{q} - \mathbf{x} \rangle - W_{\underline{\mathbf{Y}}}[\mathbf{q}] \langle \mathbf{x} - \mathbf{q} \rangle.$$



Material models

- The *force state* $\underline{\mathbf{T}}$ associates a force density vector with each bond.
- For an elastic material, this is the Fréchet derivative of strain energy density:

$$\underline{\mathbf{T}}[\mathbf{x}] \langle \mathbf{x}' - \mathbf{x} \rangle = W_{\underline{\mathbf{Y}}}(\underline{\mathbf{Y}}[\mathbf{x}]) \langle \mathbf{x}' - \mathbf{x} \rangle.$$

- More generally, a *material model* is a state-valued function of a state:

$$\underline{\mathbf{T}}[\mathbf{x}] = \hat{\underline{\mathbf{T}}}(\underline{\mathbf{Y}}[\mathbf{x}], \text{other things}).$$

- Special case: in a *bond-based* material, each bond responds independently of all the other bonds.

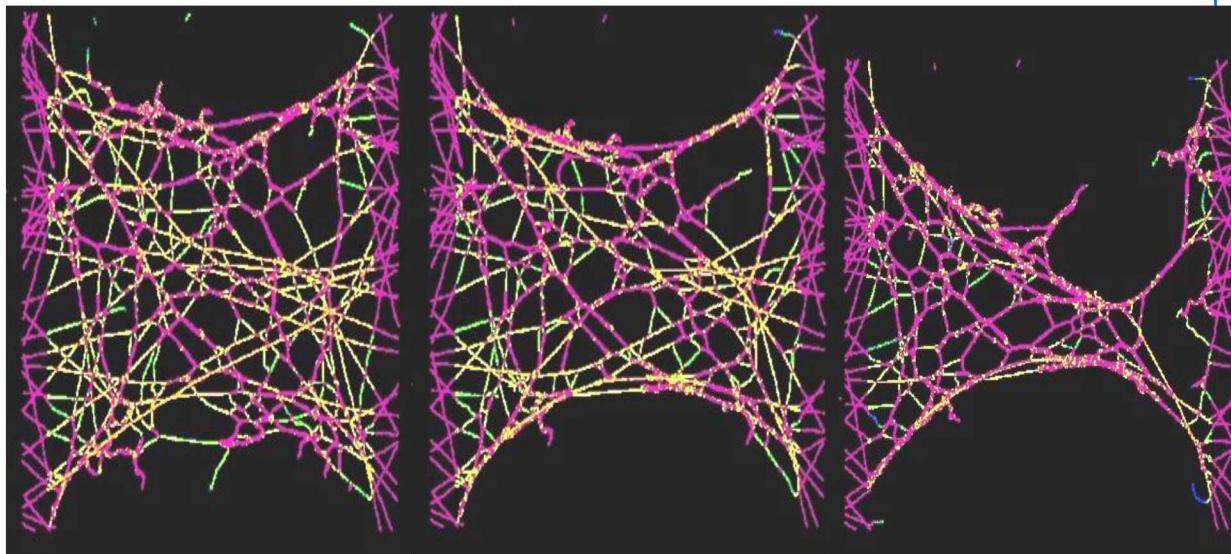
$$\hat{\underline{\mathbf{T}}}(\underline{\mathbf{Y}}[\mathbf{x}]) \langle \mathbf{x}' - \mathbf{x} \rangle = \boldsymbol{\tau}(\underline{\mathbf{Y}} \langle \mathbf{x}' - \mathbf{x} \rangle)$$

where $\boldsymbol{\tau}$ is a vector-valued function of a vector.

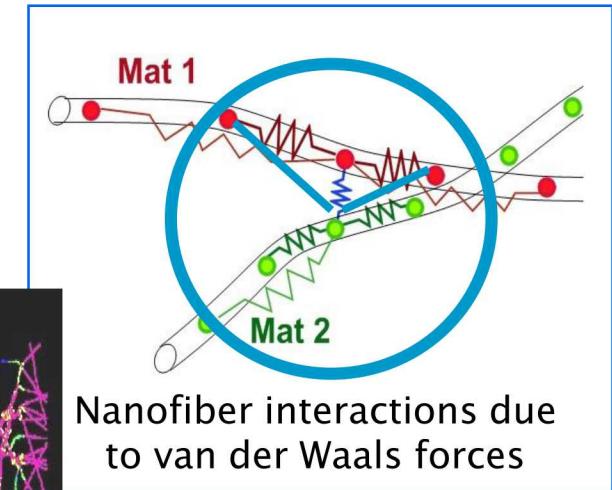
Nanofiber network:

Long-range forces dominate response

- Multiple physical interactions with different length scales can all be included in a peridynamic material model.
- This makes it a natural way to treat van der Waals and surface forces.



Nanofiber membrane (F. Bobaru, Univ. of Nebraska)



Connection with other theories

Connection to the local theory

- Start with the peridynamic equilibrium equation:

$$\int_{\mathcal{H}} (\underline{\mathbf{T}} \langle \mathbf{x}' - \mathbf{x} \rangle - \underline{\mathbf{T}}' \langle \mathbf{x} - \mathbf{x}' \rangle) dV_{\mathbf{x}'} + \mathbf{b}(\mathbf{x}) = \mathbf{0}. \quad (*)$$

- Assume the deformation is smooth. Take $\delta \rightarrow 0$.

$$(*) \rightarrow \quad \nabla \cdot \boldsymbol{\sigma} + \mathbf{b} = \mathbf{0}.$$

Connection to Kunin's nonlocal theory

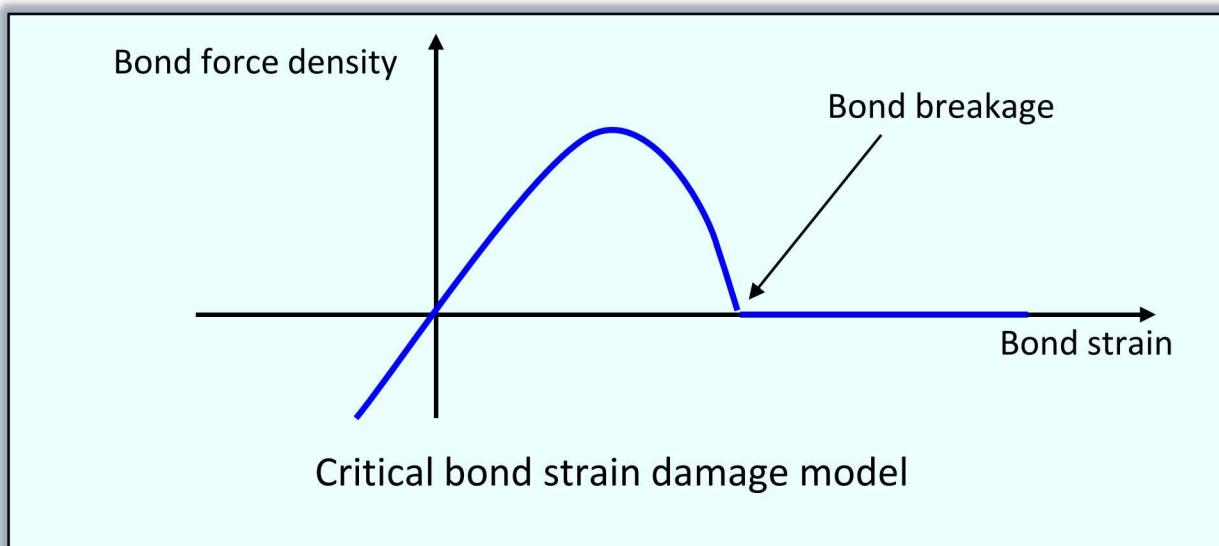
- Set $\delta = \infty$. Assume small displacements. Don't allow damage or other nonlinearity. Linearize.

$$(*) \rightarrow \quad \int_{\mathcal{B}} \mathbf{C}(\mathbf{x}', \mathbf{x}) (\mathbf{u}(\mathbf{x}') - \mathbf{u}(\mathbf{x})) dV_{\mathbf{x}'} + \mathbf{b}(\mathbf{x}) = \mathbf{0}$$

where \mathbf{C} is a tensor. This equation appears in Kunin's theory (1983).

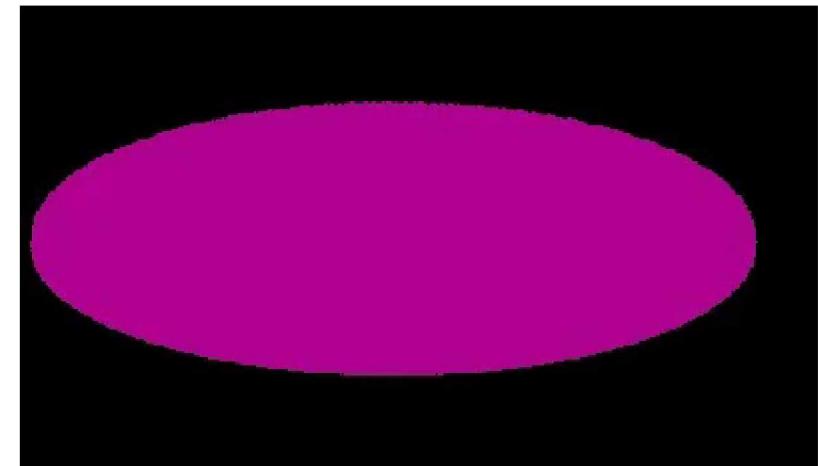
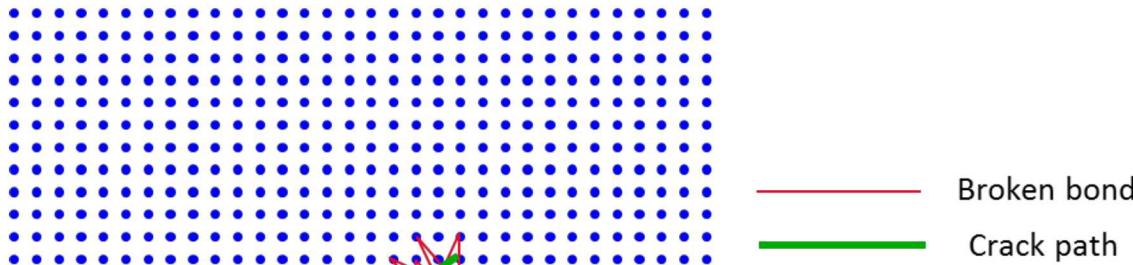
Damage due to bond breakage

- Recall: each bond carries a force.
- Damage is implemented at the bond level.
 - Bonds break irreversibly according to some criterion.
 - Broken bonds carry no force.
- Examples of criteria:
 - Critical bond strain (brittle).
 - Hashin failure criterion (composites).
 - Gurson (ductile metals).



Autonomous crack growth

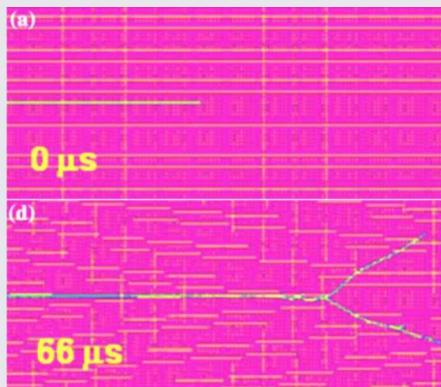
- Bond breakage leads to fracture.
- Cracks do what they want (grow, arrest, branch, curve, oscillate, ...)
- **Brittle** material model: Bond breaks when its strain reaches some critical value.



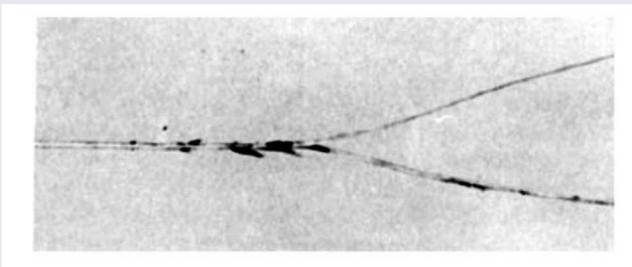
Brittle dynamic fracture

- The method reproduces many features of fracture in glass.

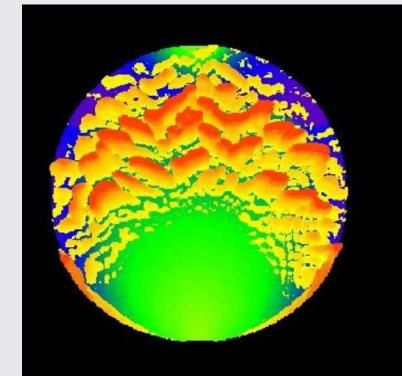
Crack branching



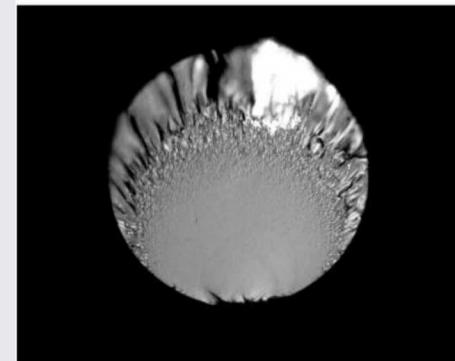
Agwai, Guven, & Madenci, Predicting crack propagation with peridynamics: a comparative study, *Int. J. Fract. Int J Fract* (2011) 171:65–78



Mirror-mist-hackle transition



3D peridynamic model



Optical fiber
(Castilone, Glaesemann & Hanson, Proc. SPIE (2002))

Particle discretization: Emu

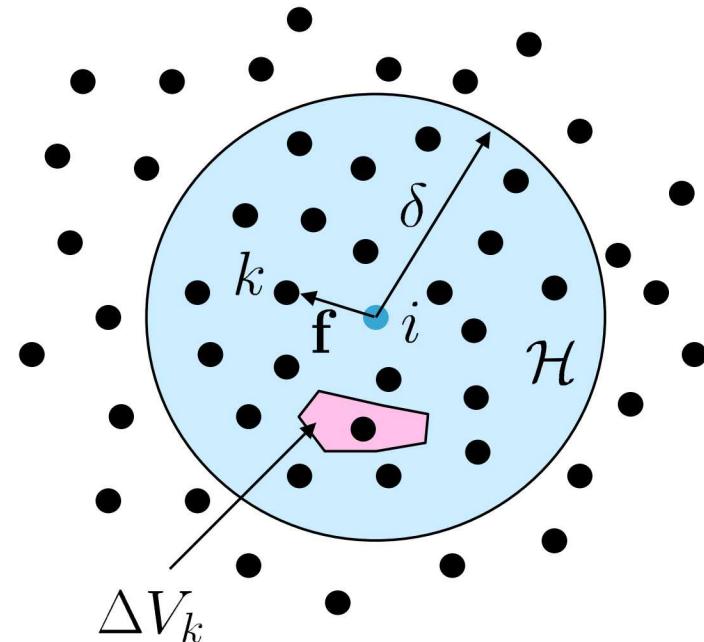
- Integral is replaced by a finite sum: resulting method is meshless and Lagrangian.

$$\rho \ddot{\mathbf{y}}(\mathbf{x}, t) = \int_{\mathcal{H}} \mathbf{f}(\mathbf{x}', \mathbf{x}, t) dV_{\mathbf{x}'} + \mathbf{b}(\mathbf{x}, t)$$

≈

$$\rho \ddot{\mathbf{y}}_i^n = \sum_{k \in \mathcal{H}} \mathbf{f}(\mathbf{x}_k, \mathbf{x}_i, t) \Delta V_k + \mathbf{b}_i^n$$

- This discretization has a special affinity with the underlying mechanics.
- Convergence properties have been studied
 - Tian & Du, SIAM J. Numerical Analysis (2014).
- Discontinuous Galerkin is also viable (LS-DYNA).
 - Chen & Gunzburger, CMAME (2011).
 - Aksoy & and Şenocak, IJNME (2011).
 - Azdoud, Han, & Lubineau, Comp. Mech. (2014).
 - Ren, Wu, & Askari, Int. J. Impact Eng. (2016).



Simulation of impact damage

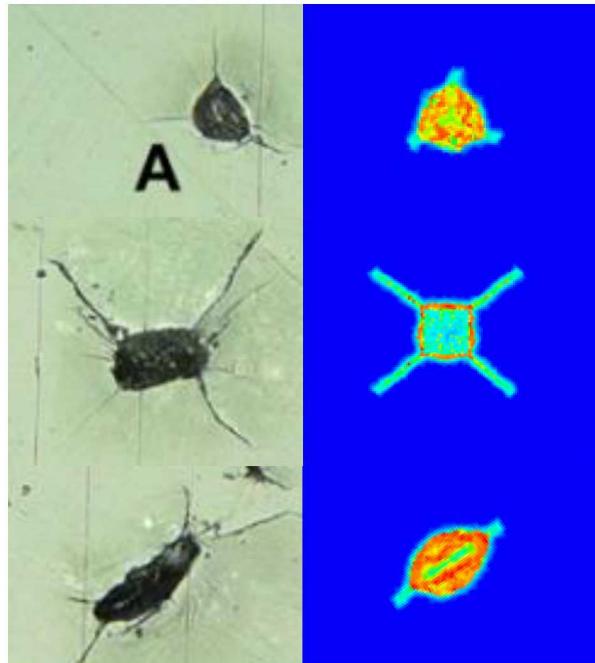
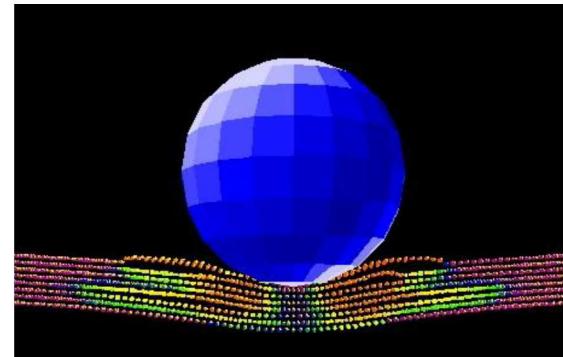
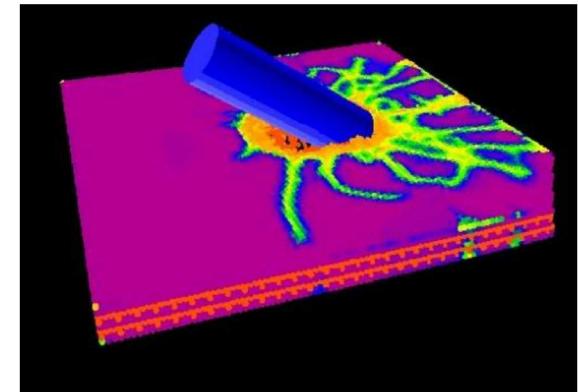


Image taken from
Maekawa, 1991.

Particle impact on damage in glass (Guven)



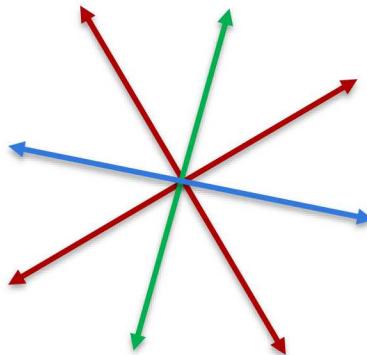
Hail against composite



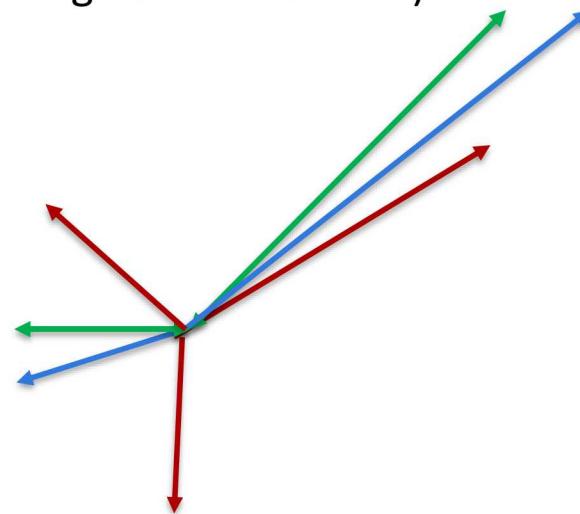
Impact on reinforced concrete

Soft materials: Geometrical nonlinearity comes “for free”

- Bond forces rotate with the bonds as the body deforms.
- Material models must be objective (invariant with respect to rotations).
- Material models must be nonpolar (balanced angular momentum).



Undeformed bonds

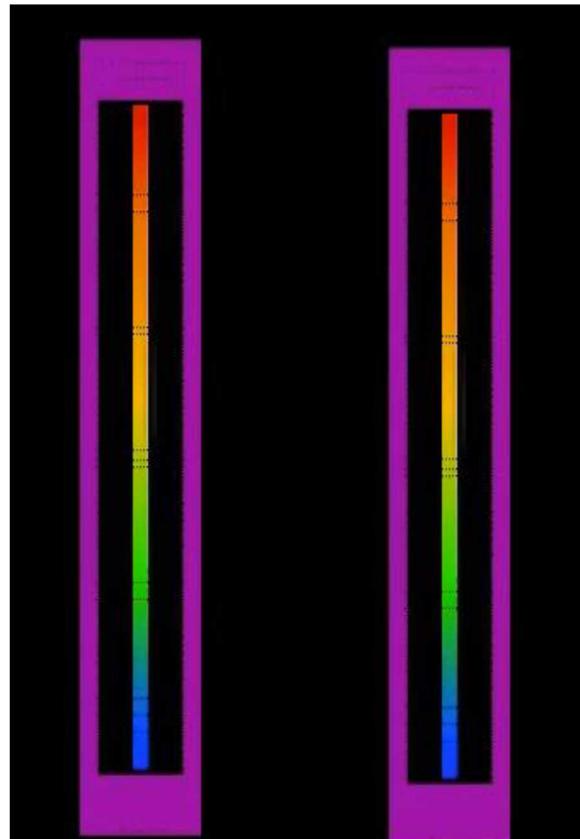


Deformed bonds

Example: Buckling and folding

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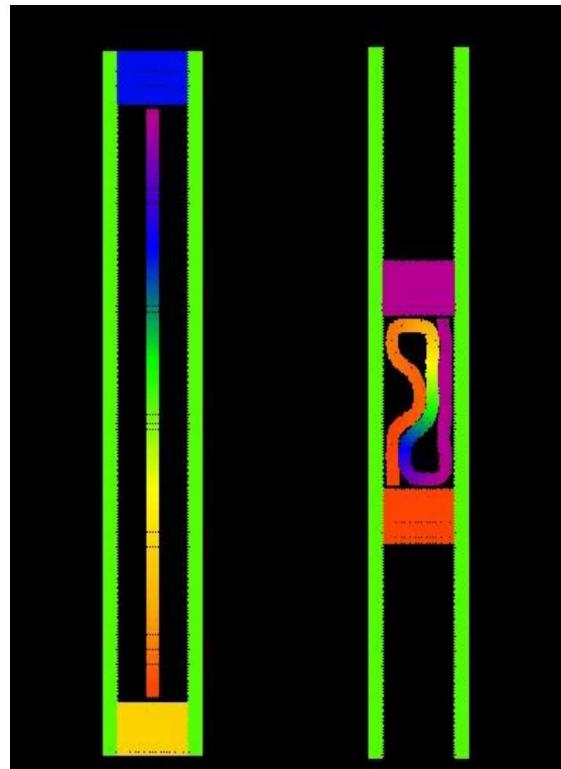
Axial compression of a beam



Linear peridynamic
solid (LPS)

Microplastic

Example: Buckling and folding

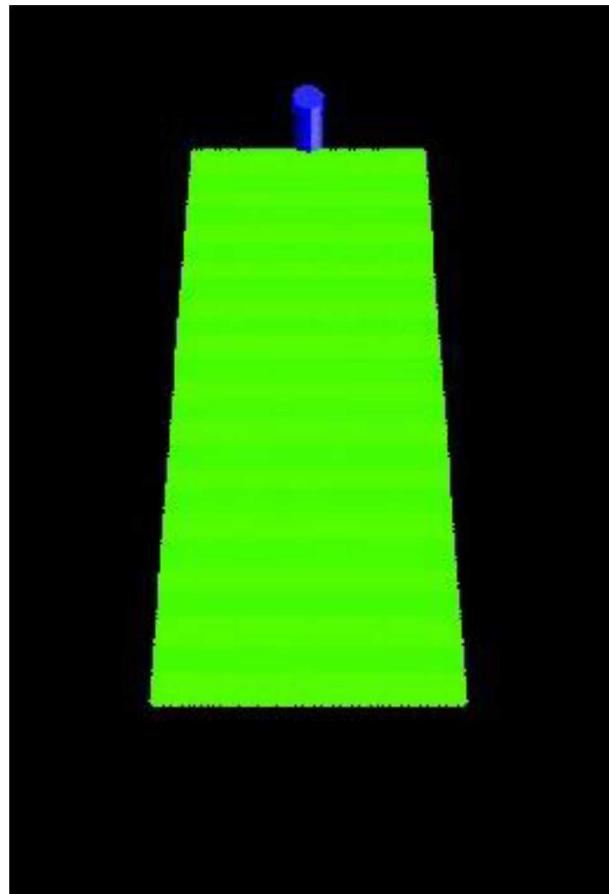


Compression of an elastic strip
State-based material model:
Linear peridynamic solid (LPS)

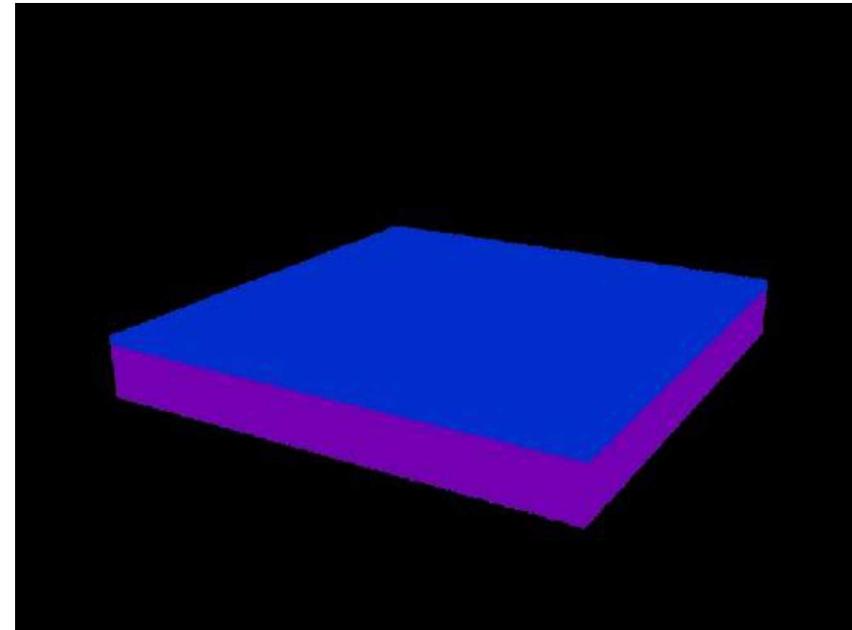
Method reproduces some subtleties in the fracture and debonding of membranes

Unstable crack path in a polyethylene membrane
(Silling & Bobaru, Int. J. Nonlin. Mech. 2005)

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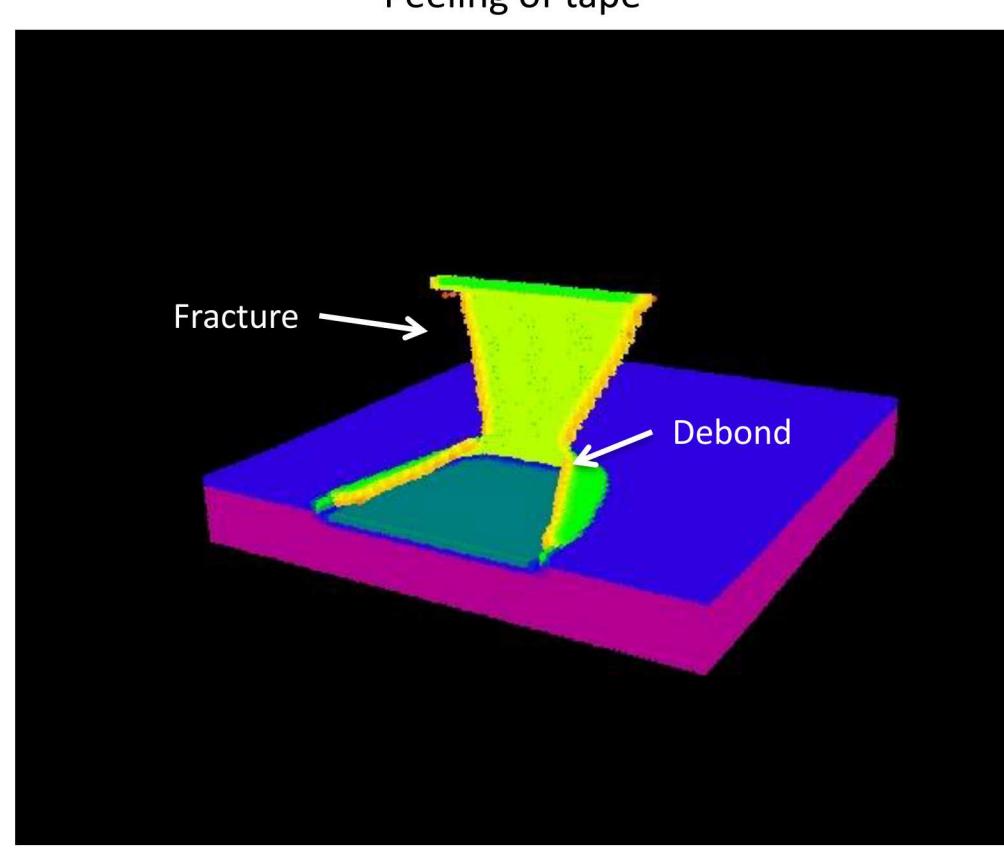
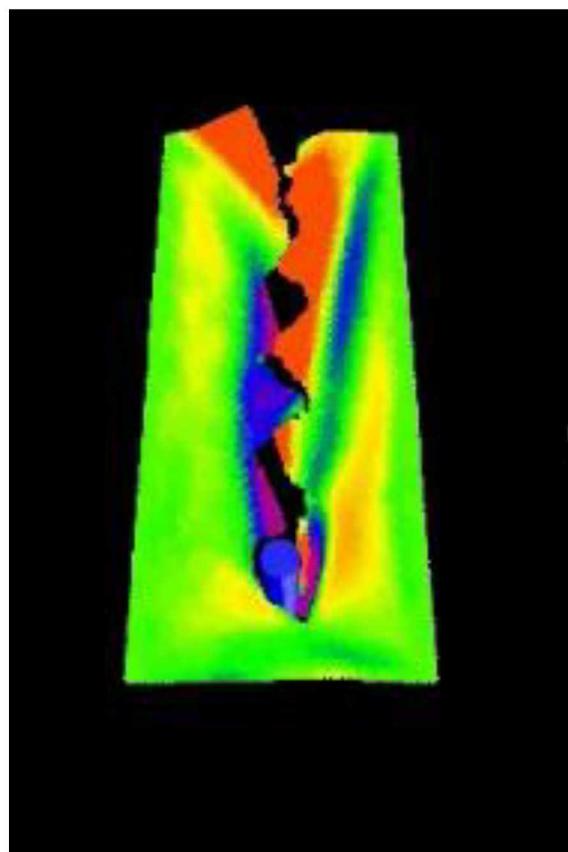


Peeling of tape



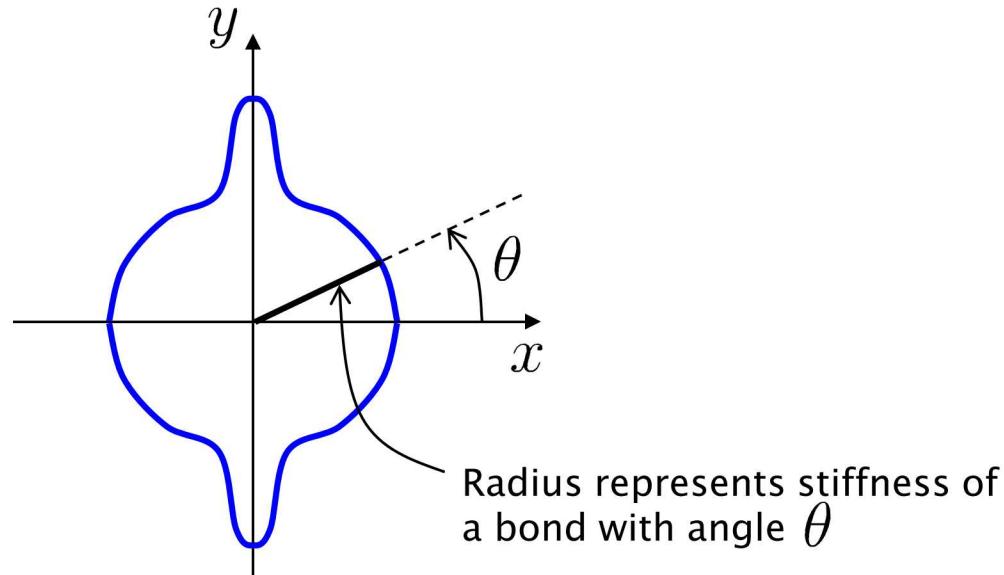
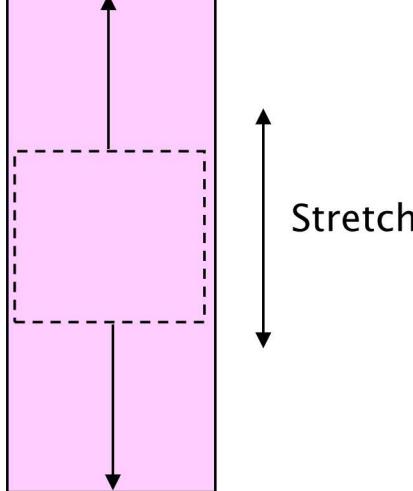
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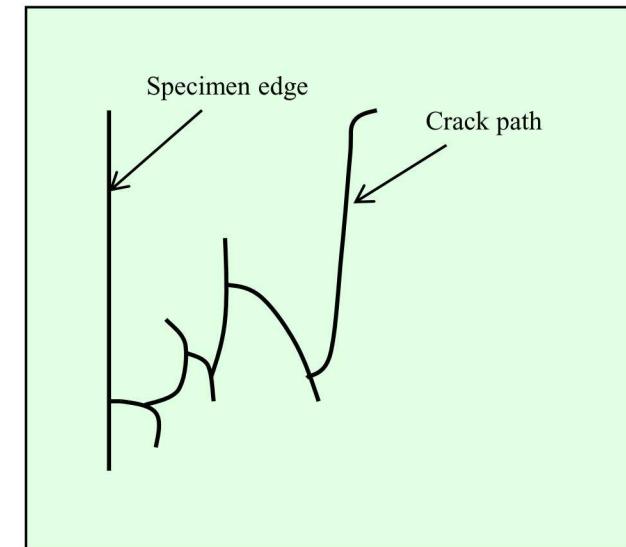
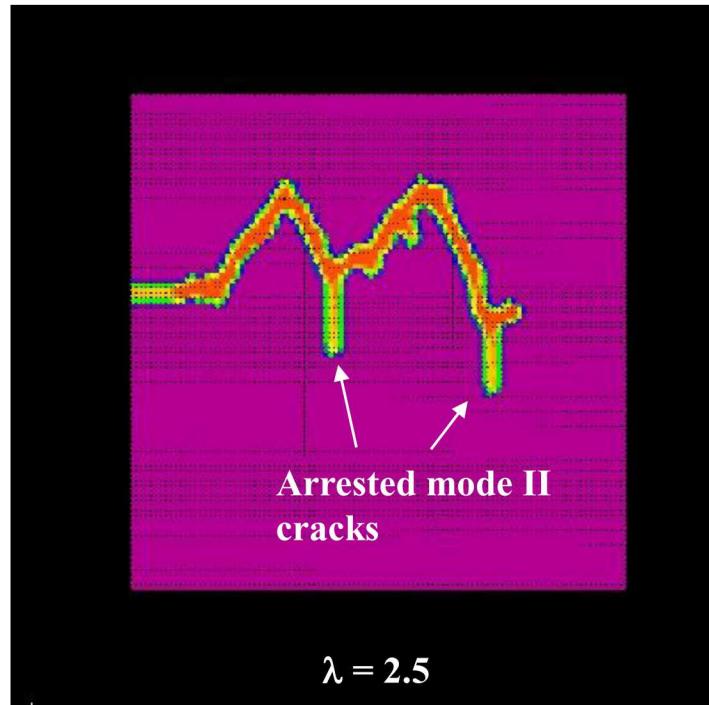
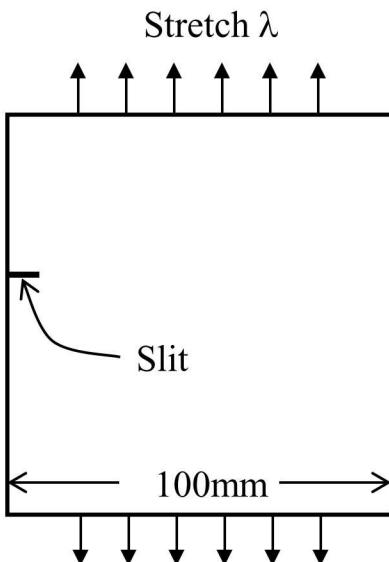
Nonlinearity in bond response can induce anisotropy: Rubber fracture

$$|\mathbf{f}| = c \left[(1 - a) + a |\sin \theta|^n \right] s$$



Crack turning in rubber

- Due to anisotropy, crack growth in mode II (parallel to loading) is competitive with mode I (straight ahead).
 - But the fields seen by a crack tip change as it grows.
 - Result: Cracks change direction and branches appear (then stop).



Experiment
(After Hamed et al. , 1996)

Nonlocal thermodynamics: Internal energy and stress power*

- First law statement:

$$\dot{\varepsilon} = \underline{\mathbf{T}} \bullet \dot{\underline{\mathbf{Y}}} + h + r$$

where ε =internal energy density, r =energy source rate, h =energy transport rate.

- Compare this with the statement in the standard theory:

$$\dot{\varepsilon} = \boldsymbol{\sigma} \cdot \dot{\mathbf{F}} + h + r.$$

- The stress power term sums up the work done on individual bonds:

$$\underline{\mathbf{T}} \bullet \dot{\underline{\mathbf{Y}}} = \int_{\mathcal{H}} \underline{\mathbf{T}}(\xi) \cdot \dot{\underline{\mathbf{Y}}}(\xi) \, dV_{\xi}.$$

- If the material is elastic, all of this work goes into the strain energy density:

$$W_{\underline{\mathbf{Y}}} \bullet \dot{\underline{\mathbf{Y}}} = \dot{W}.$$

* Joint work with Rich Lehoucq. See SS & Lehoucq, Adv. Appl. Mech. (2010)

Heat transport in bonds

- A bond-based nonlocal heat transport law:

$$h(\mathbf{x}, t) = \int_{\mathcal{H}} K(\mathbf{q}, \mathbf{x})(\theta(\mathbf{q}, t) - \theta(\mathbf{x}, t)) dV_{\mathbf{q}}$$

where θ is temperature, K is the bond conductivity.

-

$$K(\mathbf{x}, \mathbf{q}) = K(\mathbf{q}, \mathbf{x})$$

ensures conservation of energy.

- Second law

$$\theta \dot{\eta} \geq h + r$$

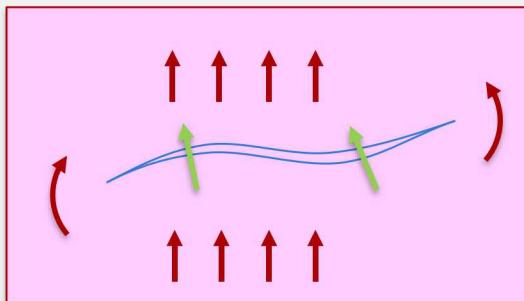
where η is the entropy.

- Can show that the 2nd law implies $K \geq 0$ always.

Cracks and other singularities in heat transport

Heat conduction near a crack:

- Some heat goes around the corners where PDE heat equation blows up.
- Some heat jumps across the crack.



- F. Bobaru, & M. Duangpanya. "A peridynamic formulation for transient heat conduction in bodies with evolving discontinuities." *Journal of Computational Physics* (2012).
- S. Oterkus, E. Madenci, and A. Agwai. "Peridynamic thermal diffusion." *Journal of Computational Physics* (2014).

“Rewetting problem”

Water

Hot nuclear fuel rod

Fluids:

Effectively Eulerian material models

- A Lagrangian material model involves both the undeformed and deformed bond vectors. Example:

$$\underline{\mathbf{T}}(\xi) = (|\underline{\mathbf{Y}}(\xi)| - |\xi|) \frac{\underline{\mathbf{Y}}(\xi)}{|\underline{\mathbf{Y}}(\xi)|}.$$

This term makes the model Lagrangian

- An Eulerian material model has bond forces that depend only on the deformed bond vectors. Example:

$$\underline{\mathbf{T}}(\xi) = |\underline{\mathbf{Y}}(\xi)|^{-n} \frac{\underline{\mathbf{Y}}(\xi)}{|\underline{\mathbf{Y}}(\xi)|},$$

$$n > 0.$$

Using an equation of state to find the bond forces

- Define a nonlocal density by

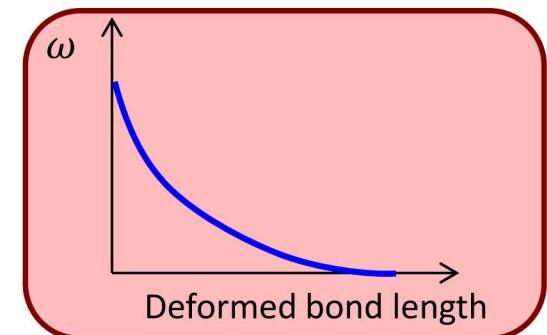
$$\rho = \rho_0 \int_{\mathcal{B}} \omega(|\underline{\mathbf{Y}}(\xi)|) dV_\xi$$

where ρ_0 is the reference density and ω is a weighting function such that $\int \omega = 1$. Integration is in the reference configuration.

- Compute the pressure from

$$p = -\frac{1}{\rho^2} \frac{\partial \psi}{\partial \rho}.$$

where ψ is the free energy density.



- The force state is found from the Frechet derivative of ψ to be

$$\underline{\mathbf{T}}(\xi) = \frac{\partial \psi}{\partial \underline{\mathbf{Y}}} = \frac{\partial \psi}{\partial \rho} \frac{\partial \rho}{\partial \underline{\mathbf{Y}}} = \frac{p \omega'(\xi)}{\rho^2} \frac{\underline{\mathbf{Y}}(\xi)}{|\underline{\mathbf{Y}}(\xi)|}.$$

Surface tension is implemented through nonlocal forces

- Surface tension arises from nonlocal forces between molecules.
- Peridynamic Eulerian model:

$$\underline{\mathbf{T}} = \underline{\mathbf{T}}^{\text{eos}} + \underline{\mathbf{T}}^{\text{surf}}$$

where

$$\underline{\mathbf{T}}^{\text{surf}} = \gamma \frac{\underline{\mathbf{Y}}}{|\underline{\mathbf{Y}}|}$$

and γ is a constant.

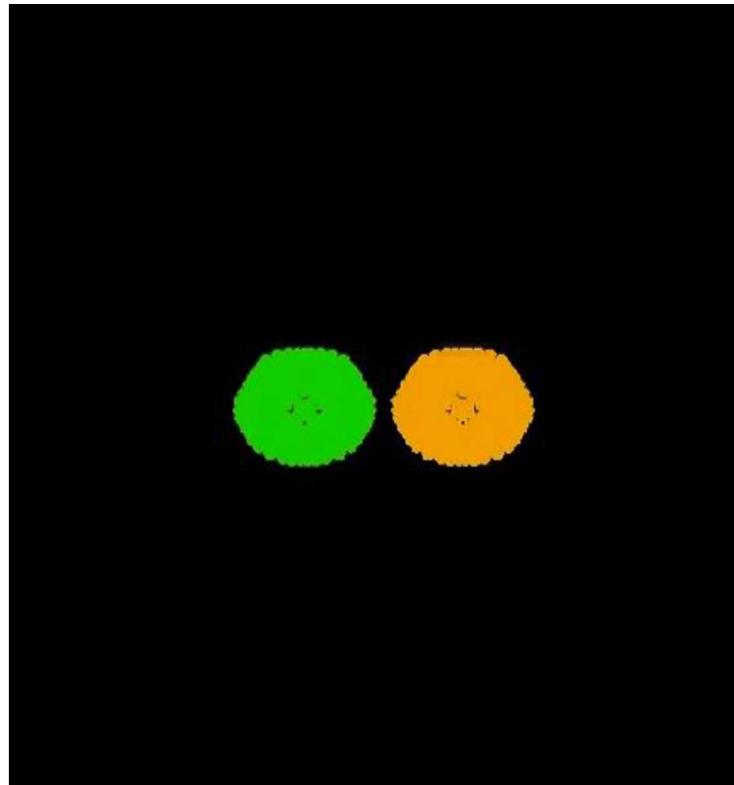
- Each pair of material particles within each other's family attracts.

Surface tension examples

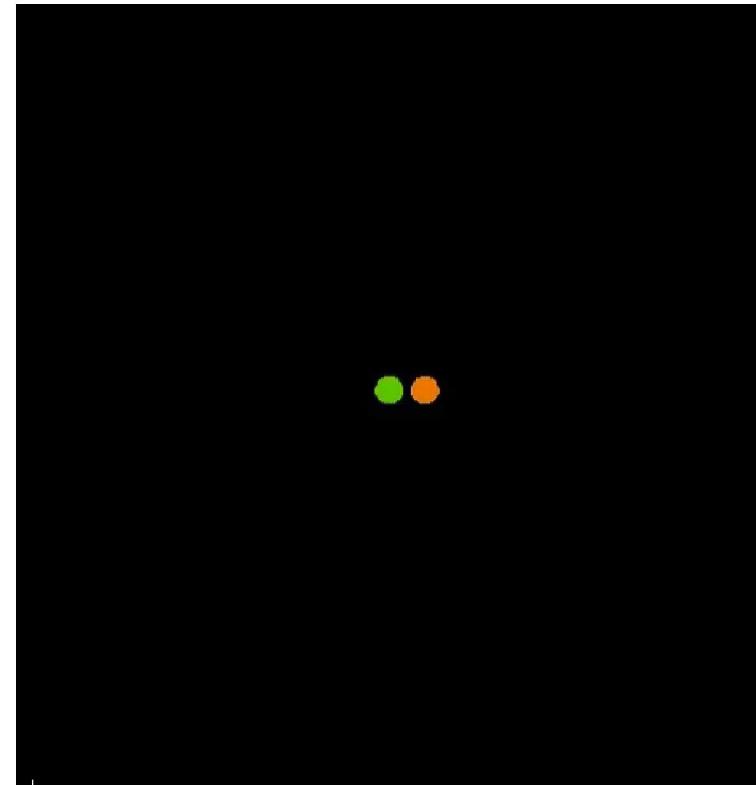
- Two droplets collide.
- Mie-Gruneisen EOS is used.

VIDEOS

THE BLOB ($v=10$)



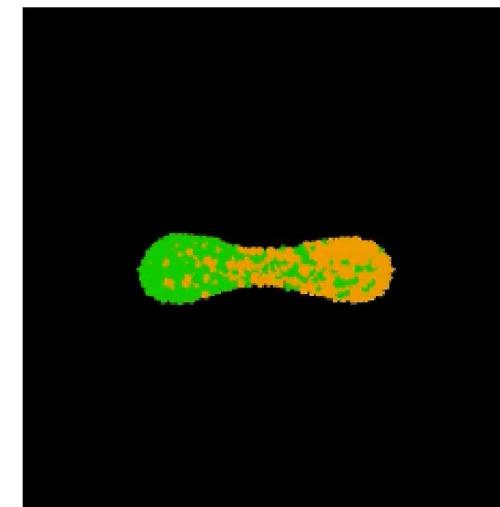
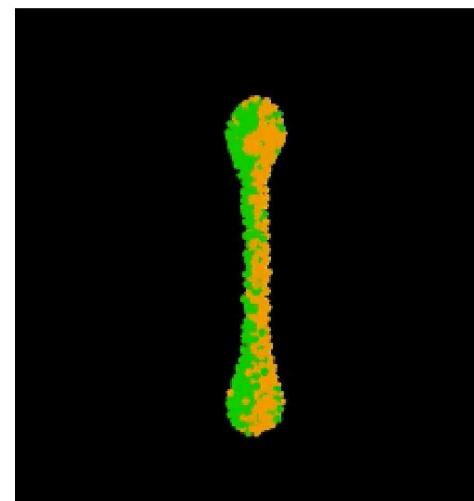
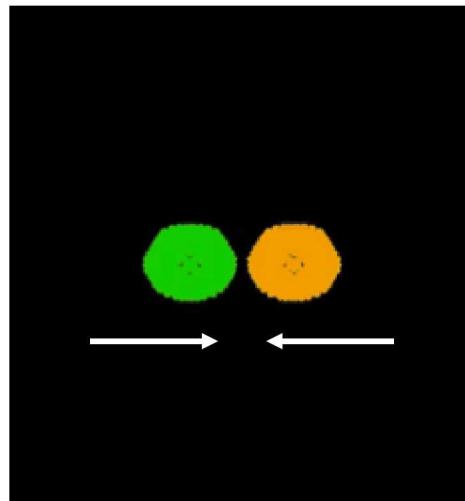
REVENGE OF THE BLOB ($v=100$)



Surface tension examples

- Two droplets collide.
- Mie-Gruneisen EOS is used.

THE BLOB ($\nu=10$)



Contact and friction forces as Eulerian material response

- Short-range contact and friction forces can be included within an Eulerian material model:

$$\underline{\mathbf{T}}^E = \underline{\mathbf{T}}^{\text{eos}} + \underline{\mathbf{T}}^{\text{contact}} + \underline{\mathbf{T}}^{\text{friction}}$$

where

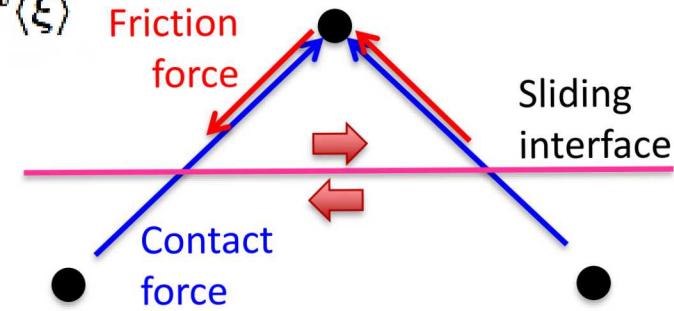
$$\underline{\mathbf{T}}^{\text{contact}}(\xi) = \begin{cases} -c(r_c - |\underline{\mathbf{Y}}(\xi)|)\mathbf{M} & \text{if } |\underline{\mathbf{Y}}(\xi)| < r_c, \\ 0 & \text{otherwise.} \end{cases}, \quad \mathbf{M} = \frac{\underline{\mathbf{Y}}(\xi)}{|\underline{\mathbf{Y}}(\xi)|}$$

where r_c is a cut-off distance for contact forces.

- The friction force state is

$$\underline{\mathbf{T}}^{\text{friction}}(\xi) = -F \operatorname{sgn} \left(\frac{\partial}{\partial t} |\underline{\mathbf{Y}}(\xi)| \right) \underline{\mathbf{T}}^{\text{contact}}(\xi)$$

where F is the friction coefficient.



Combining Lagrangian and Eulerian response in a single material model

- We'd like to model both fluid-like and solid-like response in the same material model.
- Combine the two as a linear combination of force states:

$$\underline{\mathbf{T}} = \beta(p) \underline{\mathbf{T}}^E + (1 - \beta(p)) \underline{\mathbf{T}}^L$$

where $\underline{\mathbf{T}}^E$ and $\underline{\mathbf{T}}^L$ are the Eulerian (fluid-like) and Lagrangian (solid-like) contributions respectively.

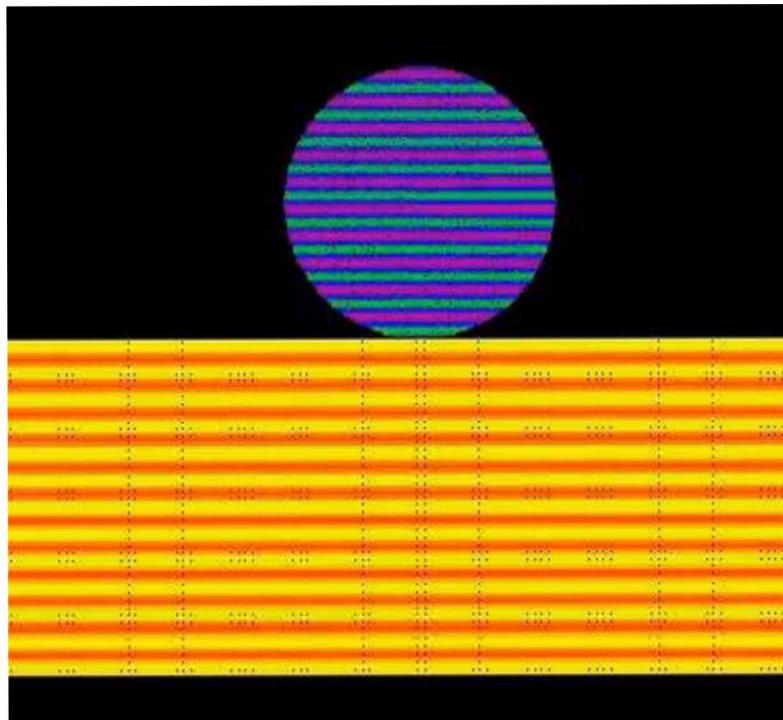
- $\beta(p)$ is a pressure-dependent interpolation parameter, $0 \leq \beta \leq 1$.
- Example: EOS & bond-based:

$$\underline{\mathbf{T}}(\xi) = \left(\frac{\beta(p)p\omega'(\xi)}{\rho} + (1 - \beta(p))C(\xi)(|\underline{\mathbf{Y}}(\xi)| - |\xi|) \right) \frac{\underline{\mathbf{Y}}(\xi)}{|\underline{\mathbf{Y}}(\xi)|}$$

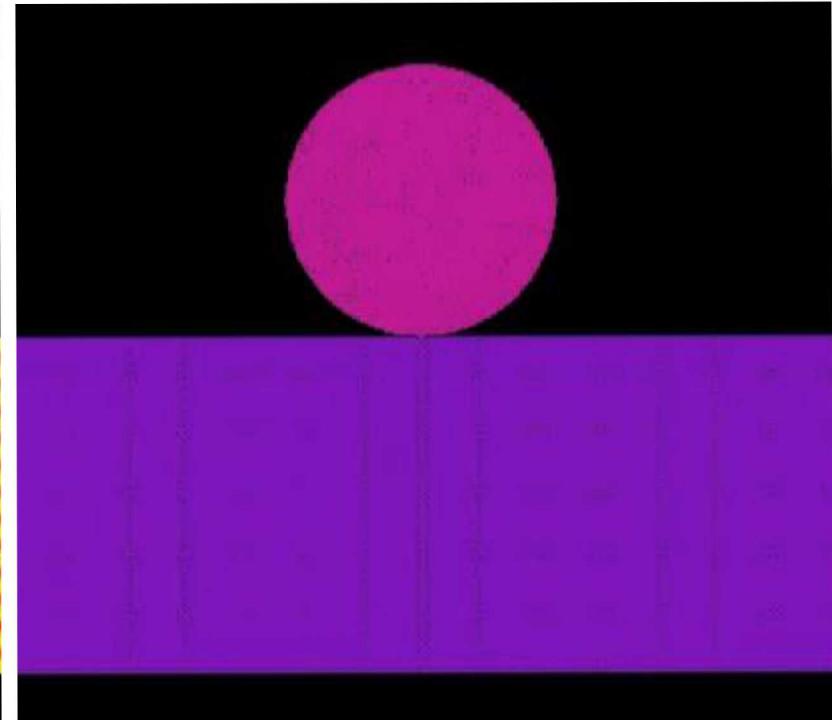
Example of Eulerian + Lagrangian material models: Wear

- Solid response and damage: Lagrangian
- Contact and friction: Eulerian

VIDEOS

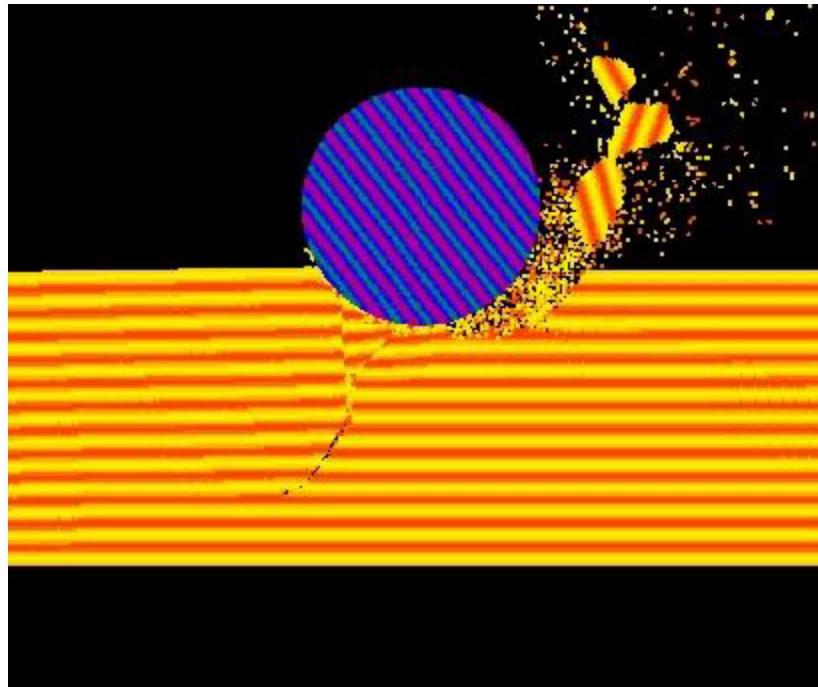


Material deformation

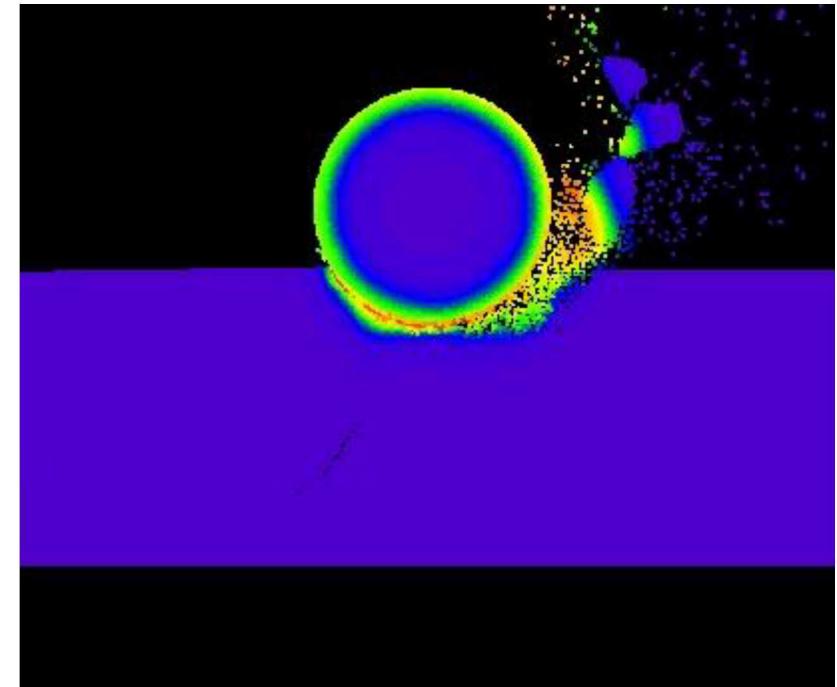


Damage

Friction forces appear in 1st law expression leading to heating



Material deformation

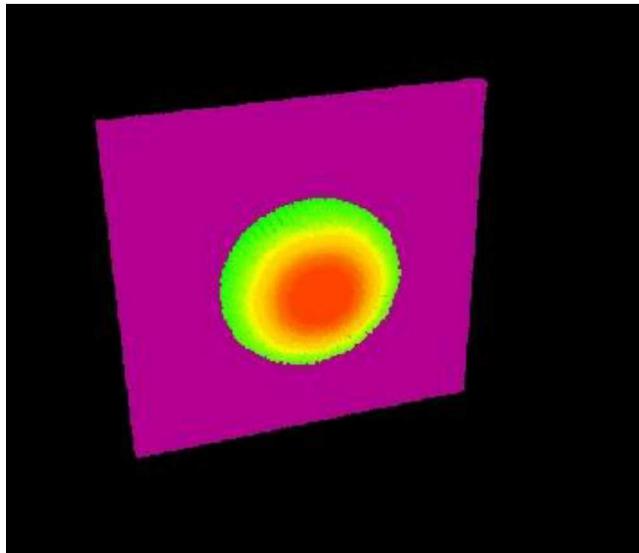


Temperature

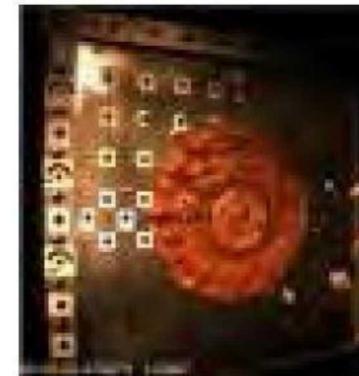
Bird strike*

Peridynamics compared with SPH

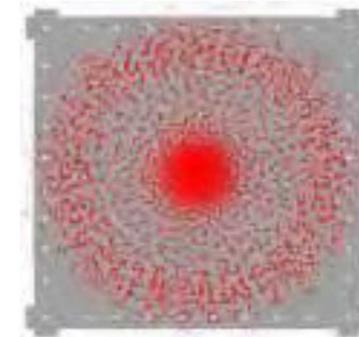
- Bird simulant (gelatin) vs. heavy plate
- The peridynamic model helps reduce the “spray” that is sometimes seen with SPH.



Peridynamics



Test - LG 997



SPH

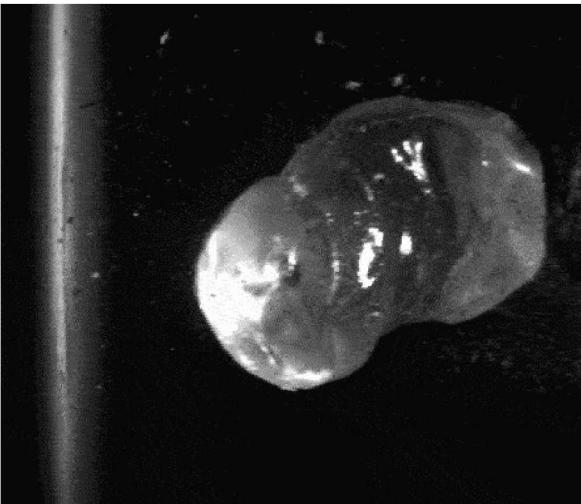
Olivares, NIS Document 09-039 (2010)

*Joint work with Boeing Research & Technology

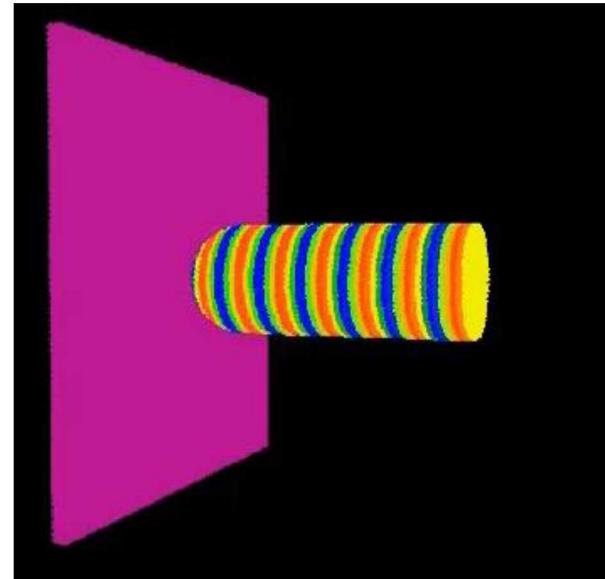
Bird strike simulant (gelatin)



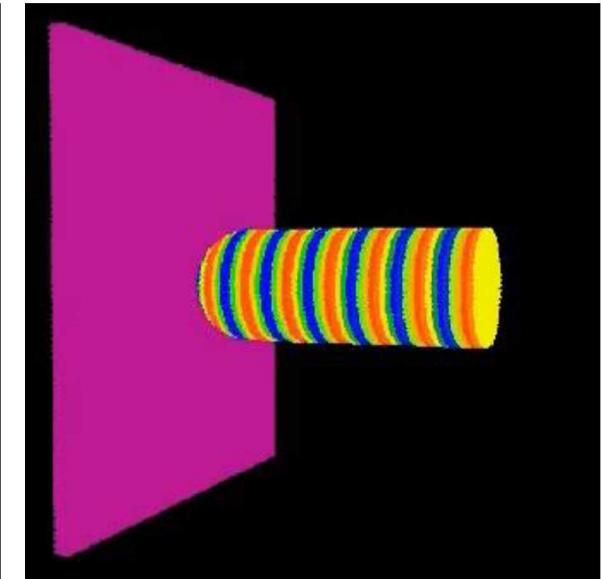
- Who cares?
 - The splash pattern helps determine loading on the structure, especially when the structure is itself highly deformable.



Typical test
(credit: Arthur Core)

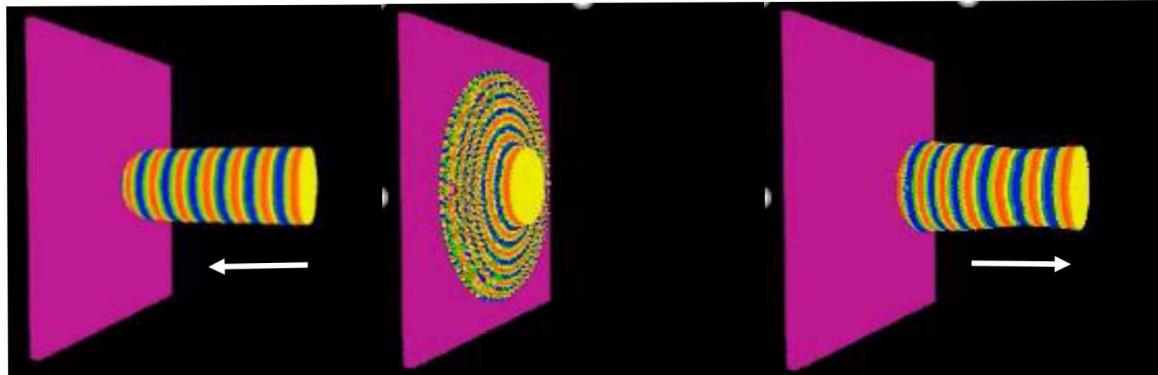


Meshless PD

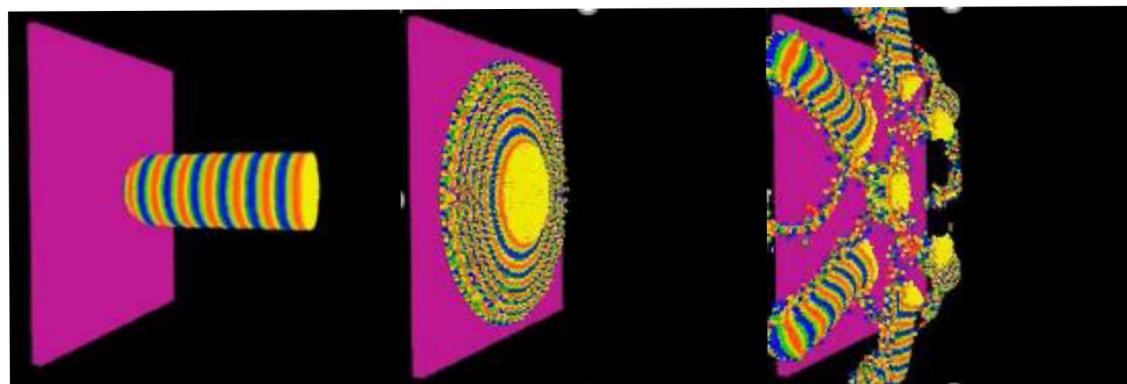


Meshless PD
with bond damage

Bird strike simulant (gelatin)



No fracture

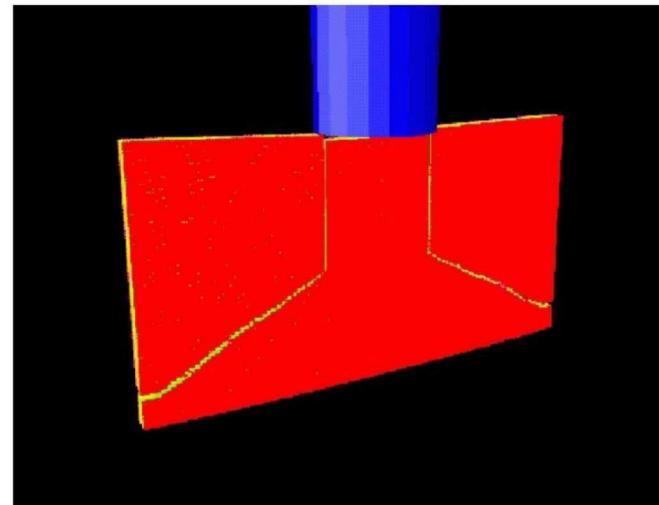


With fracture

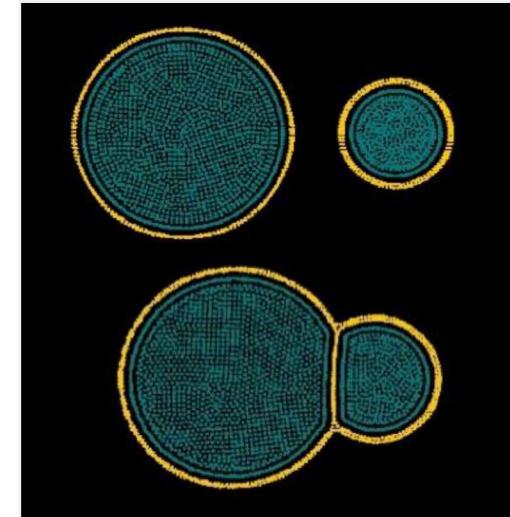
Summary

- Some techniques within peridynamics for studying large deformation and fracture in soft materials have been developed:
 - Eulerian material models.
 - Surface forces, contact, and friction as part of an Eulerian material model.
 - Combining Lagrangian & Eulerian material response.

The same theory encompasses a wide spectrum of phenomena, depending only on the choice of an appropriate material model.



Kalthoff-Winkler test



Soap bubbles