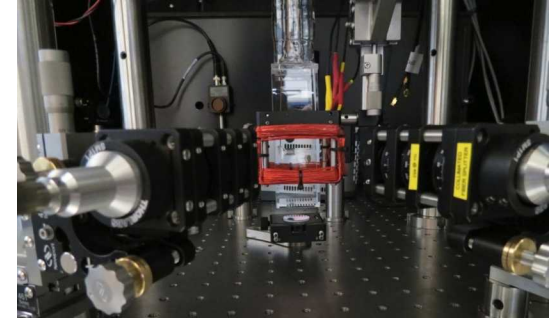
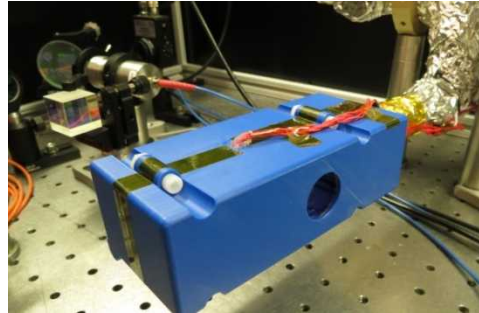


Funded by DARPA: The views expressed are those of the author(s) and do not reflect the official policy or position of the Department of Defense or the U.S. Government. Distribution "A": Approved for Public Release, Distribution Unlimited

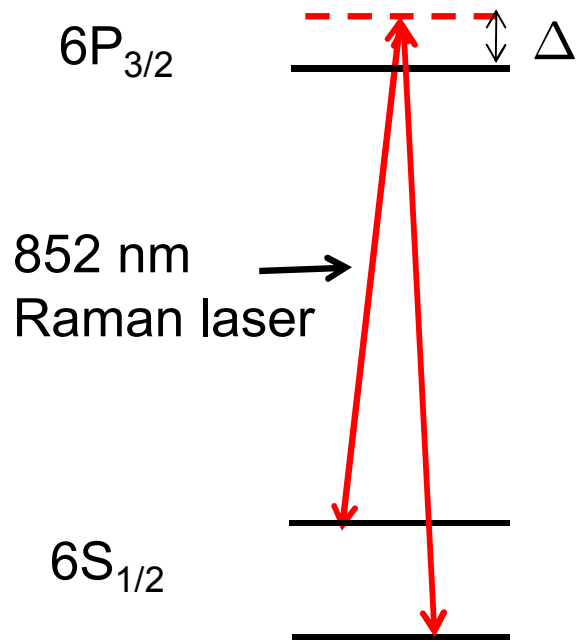


Inertial sensing with atom interferometers

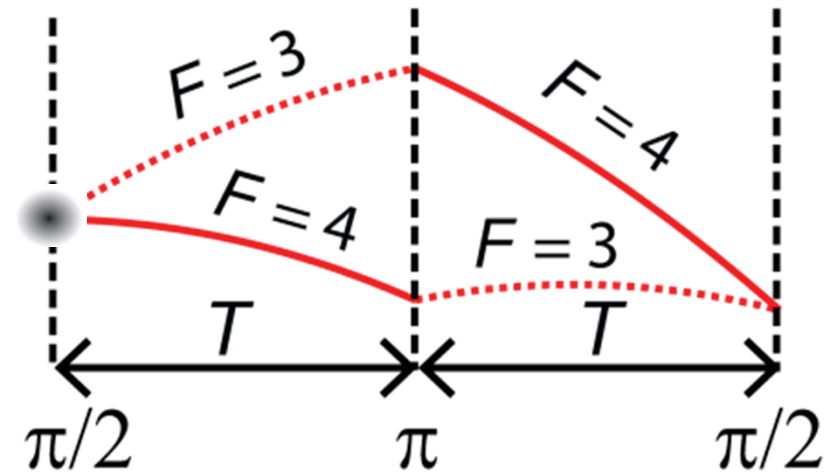
Grant Biedermann



Light-pulse atom interferometry

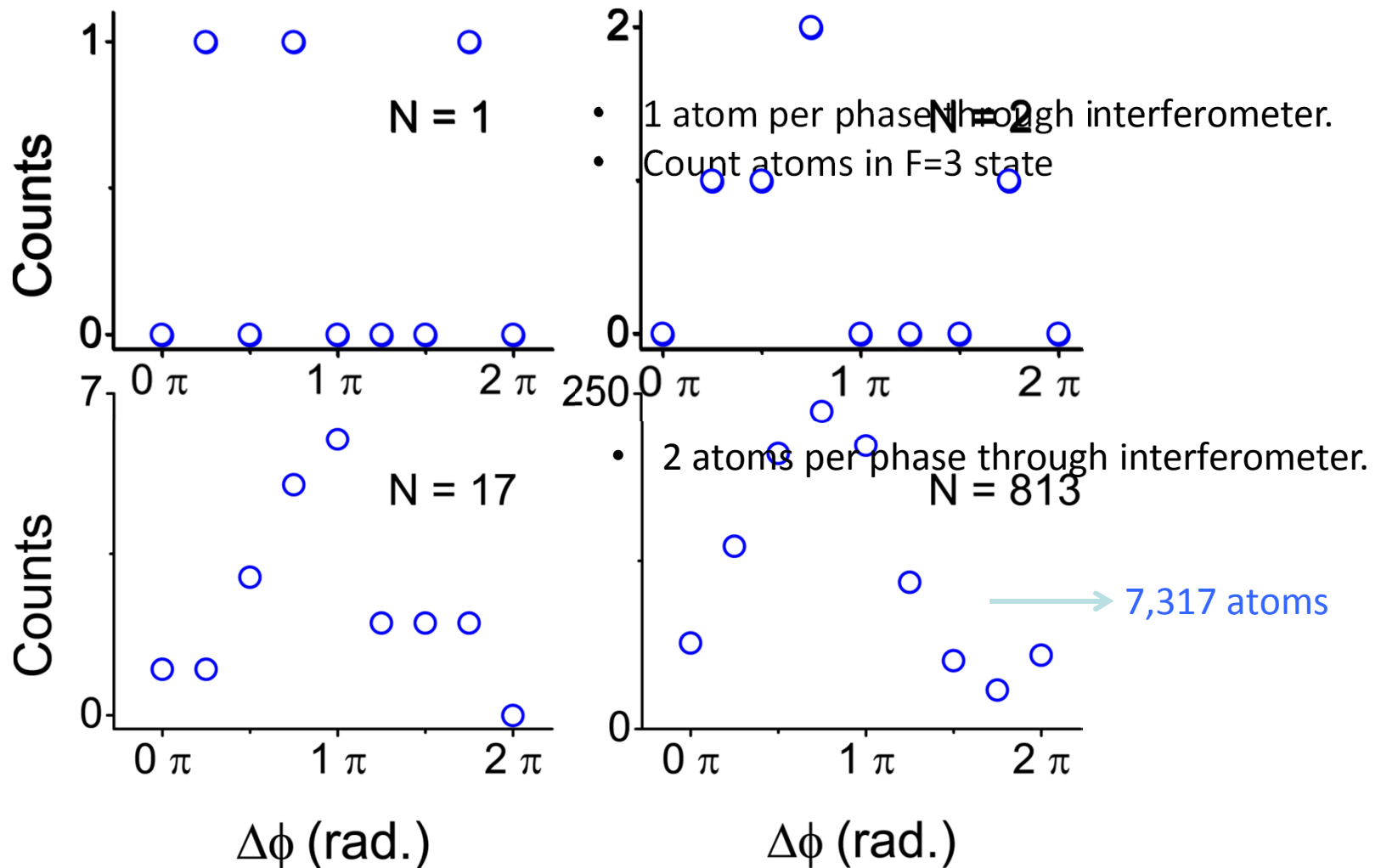


wavepacket trajectory

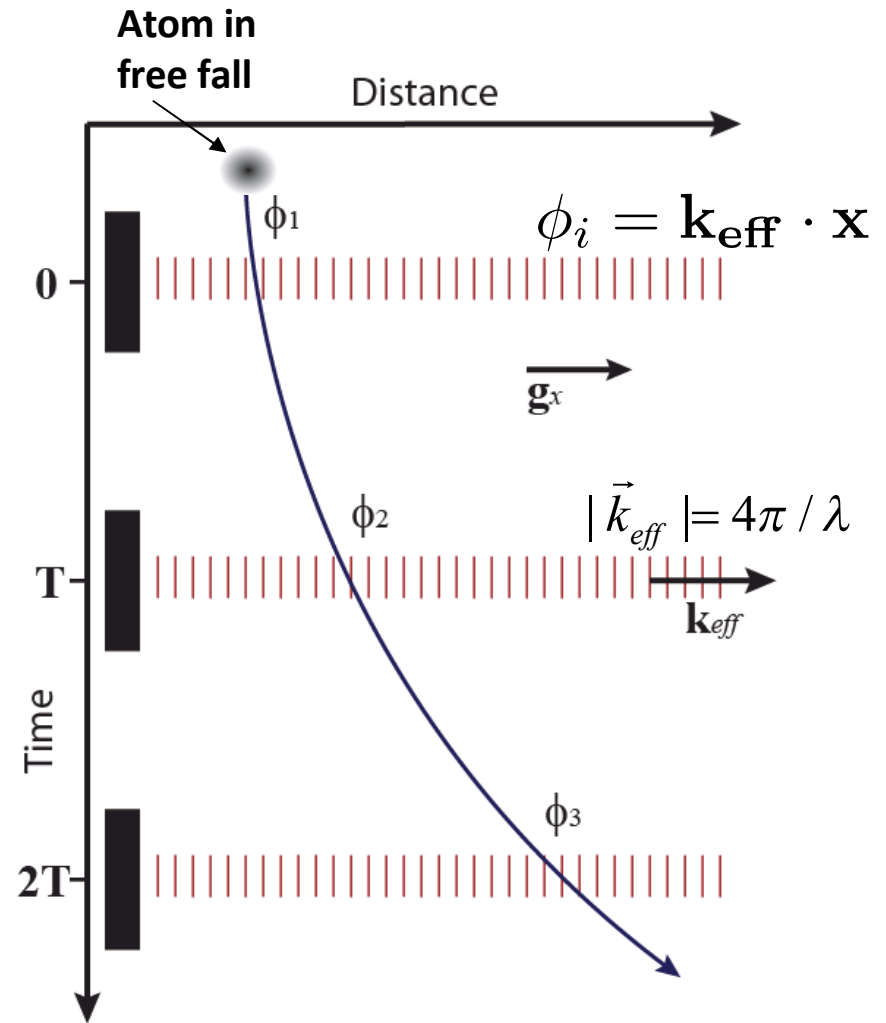
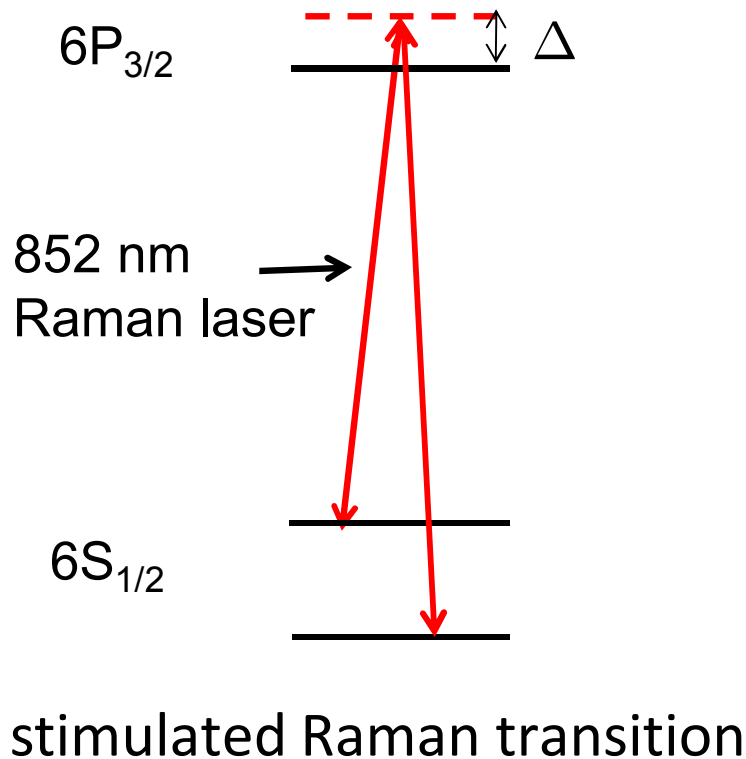


stimulated Raman transition

Building a fringe, one atom at a time

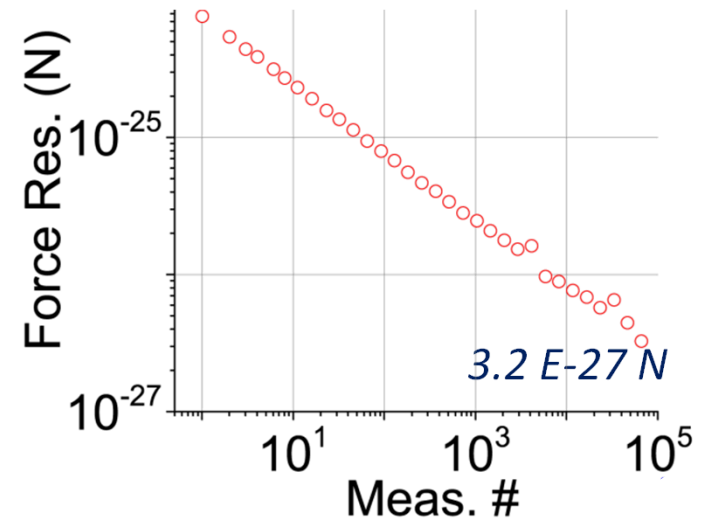
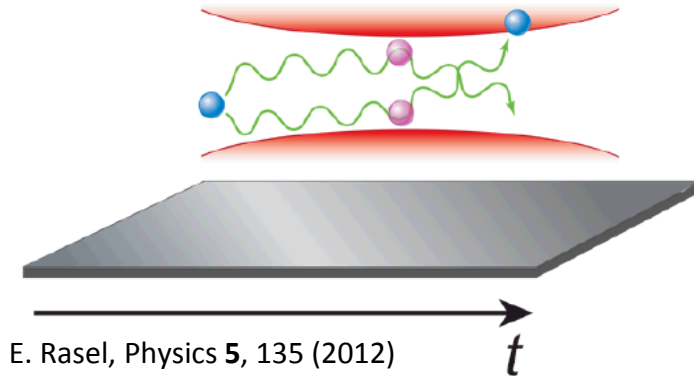


Light-pulse atom interferometry

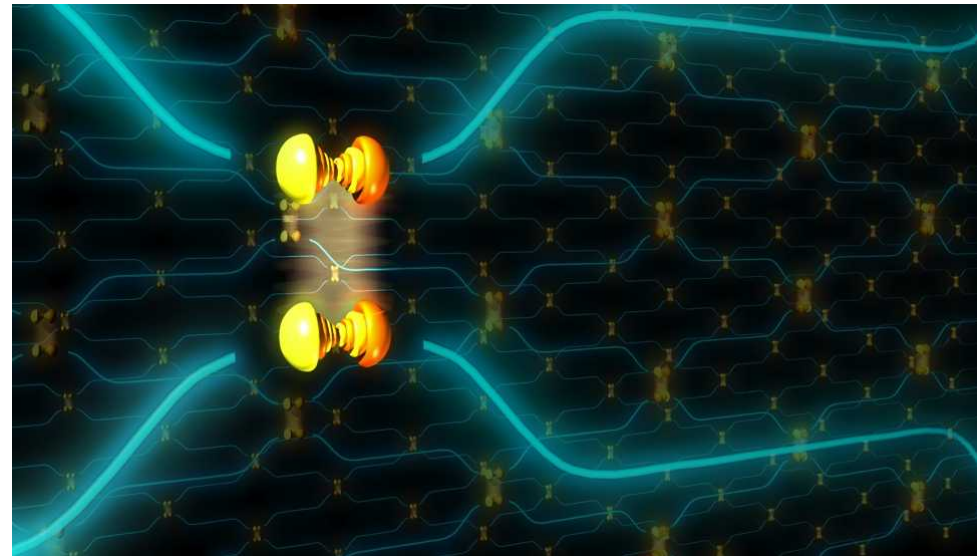
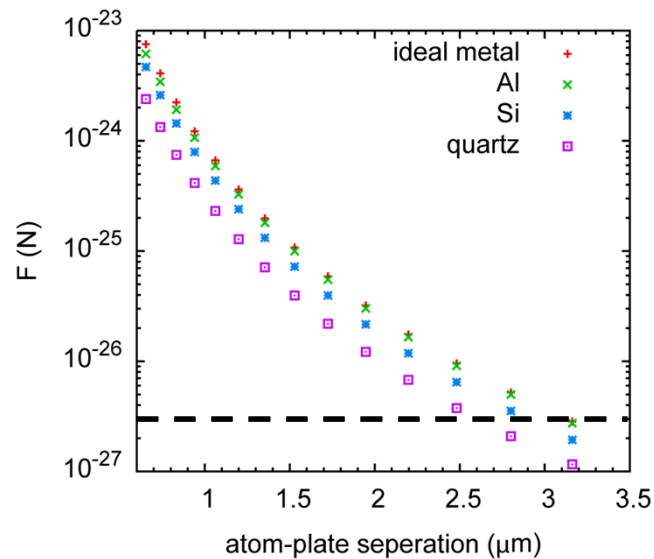


For an atom starting in $F=4$, $P_3 = \frac{1}{2} (1 - \cos\Delta\phi)$, $\Delta\phi = \phi_1 - 2\phi_2 + \phi_3$

Force resolution of a single atom interferometer



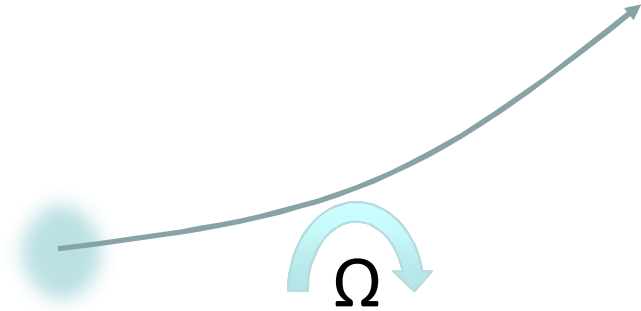
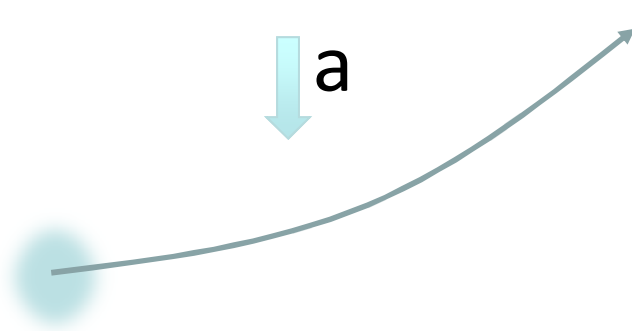
Casimir-Polder force



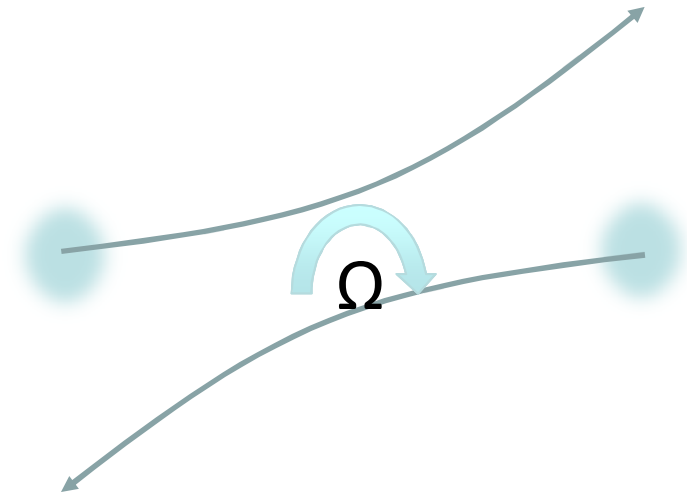
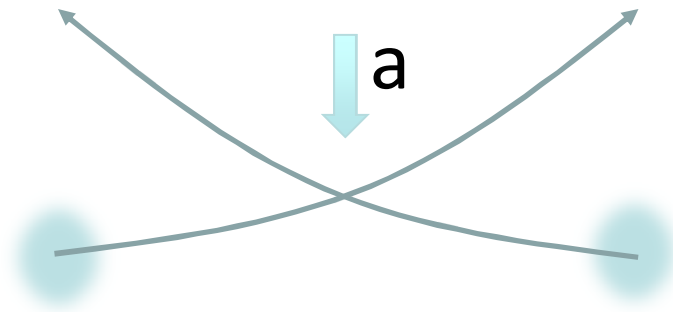
* For cesium $3.2 \text{ E-}27 \text{ N}$ is 1.5 mg

Demonstrated entanglement, arXiv:1501.03862

Rotation Measurement



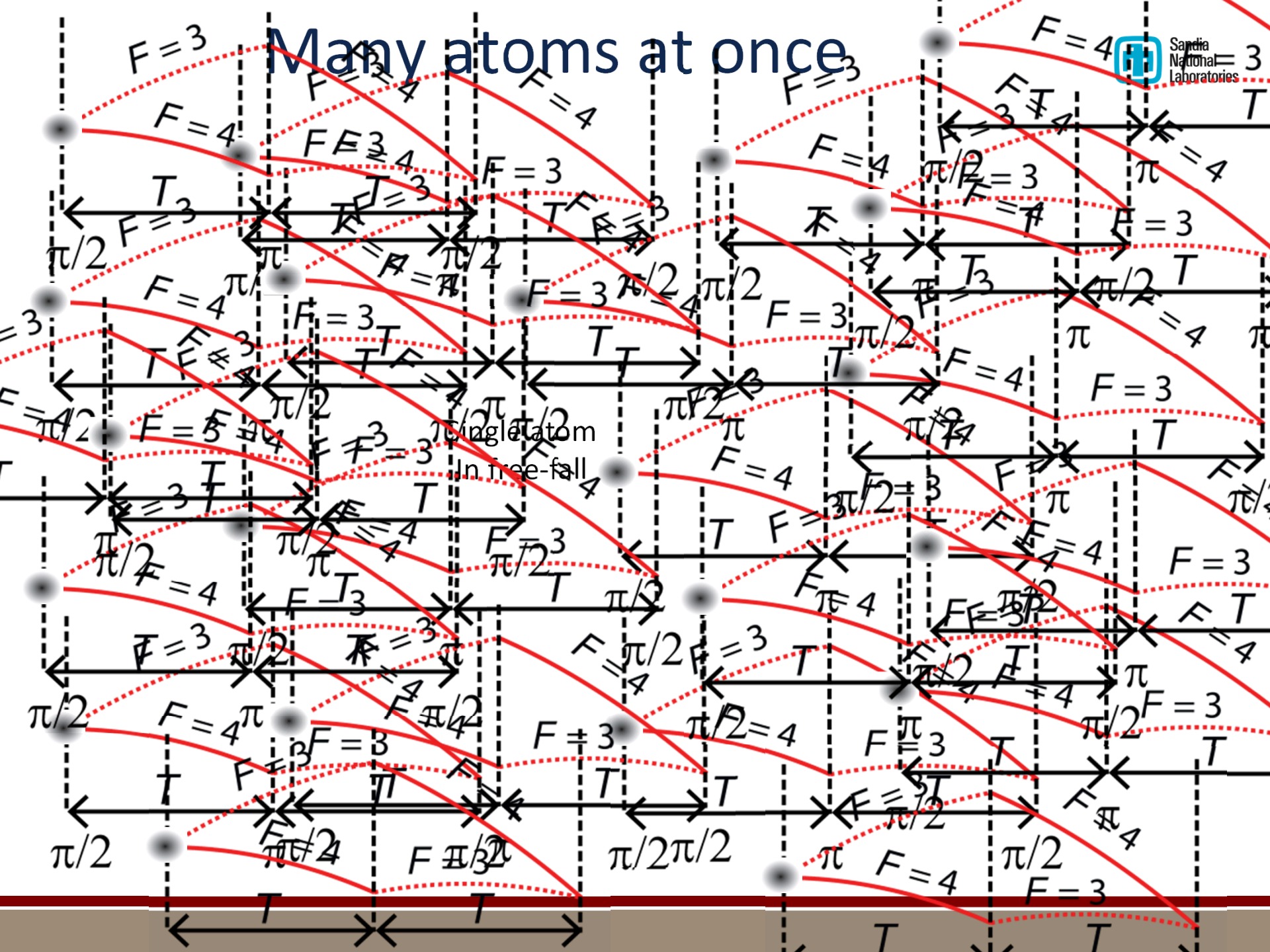
$$\Delta\phi = \vec{k}_{eff} \cdot (\vec{a}T^2 - 2(\vec{v} \times \vec{\Omega})T^2)$$



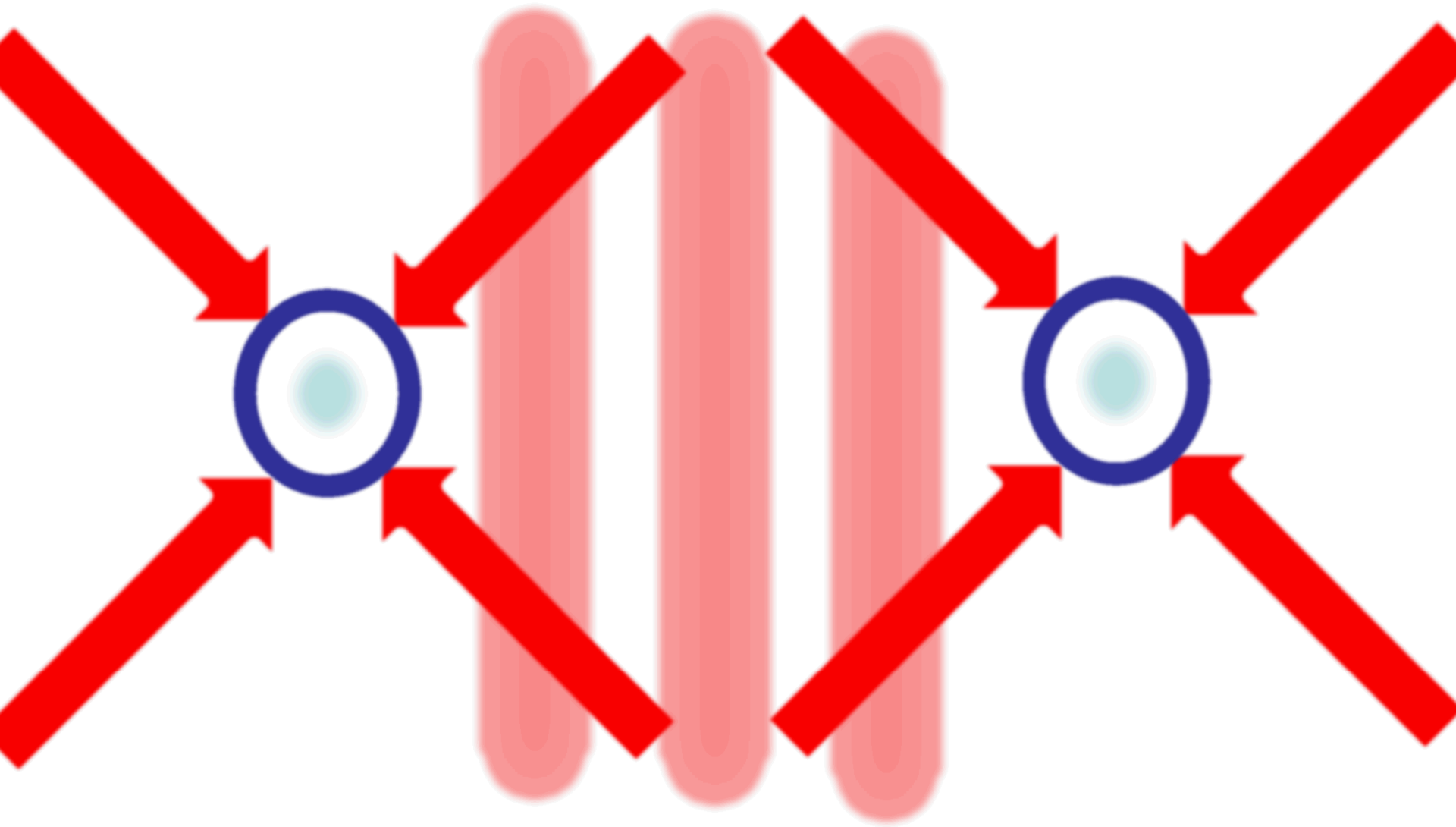
Many atoms at once



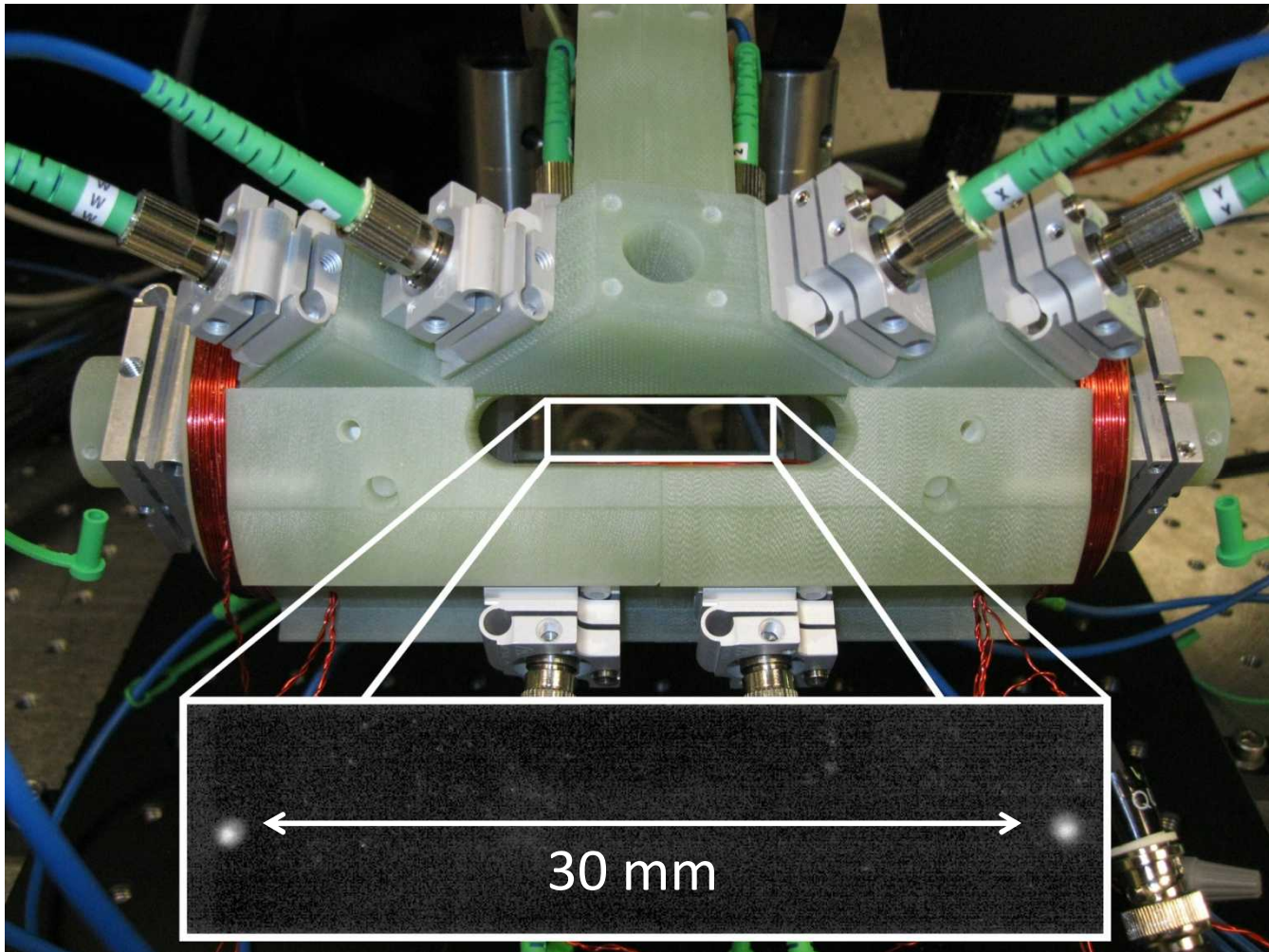
Sandia National Laboratories



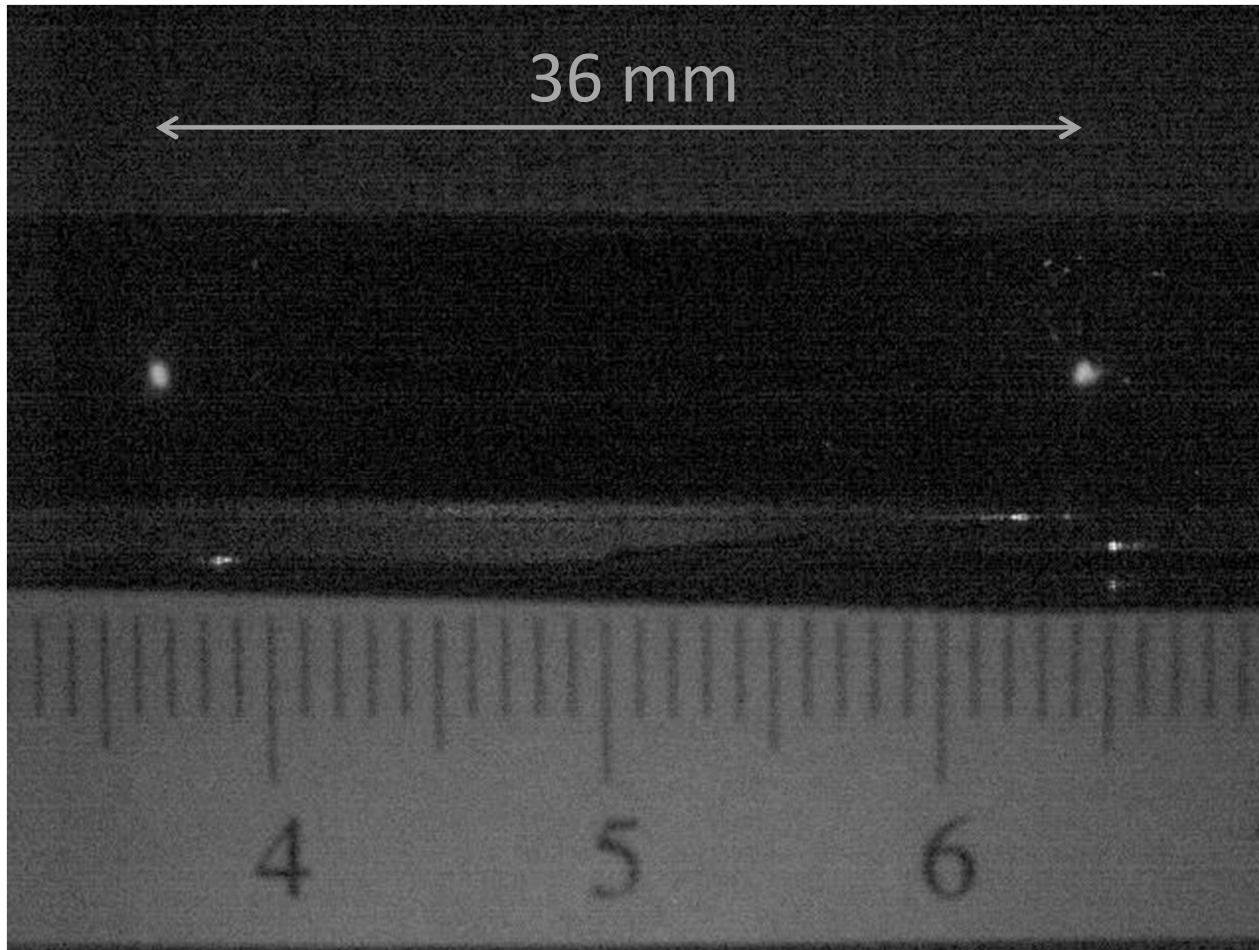
Ensemble Exchange



Dueling interferometers



LAUNCH AND RECAPTURE



LAUNCH AND RECAPTURE

Steady state atom number:

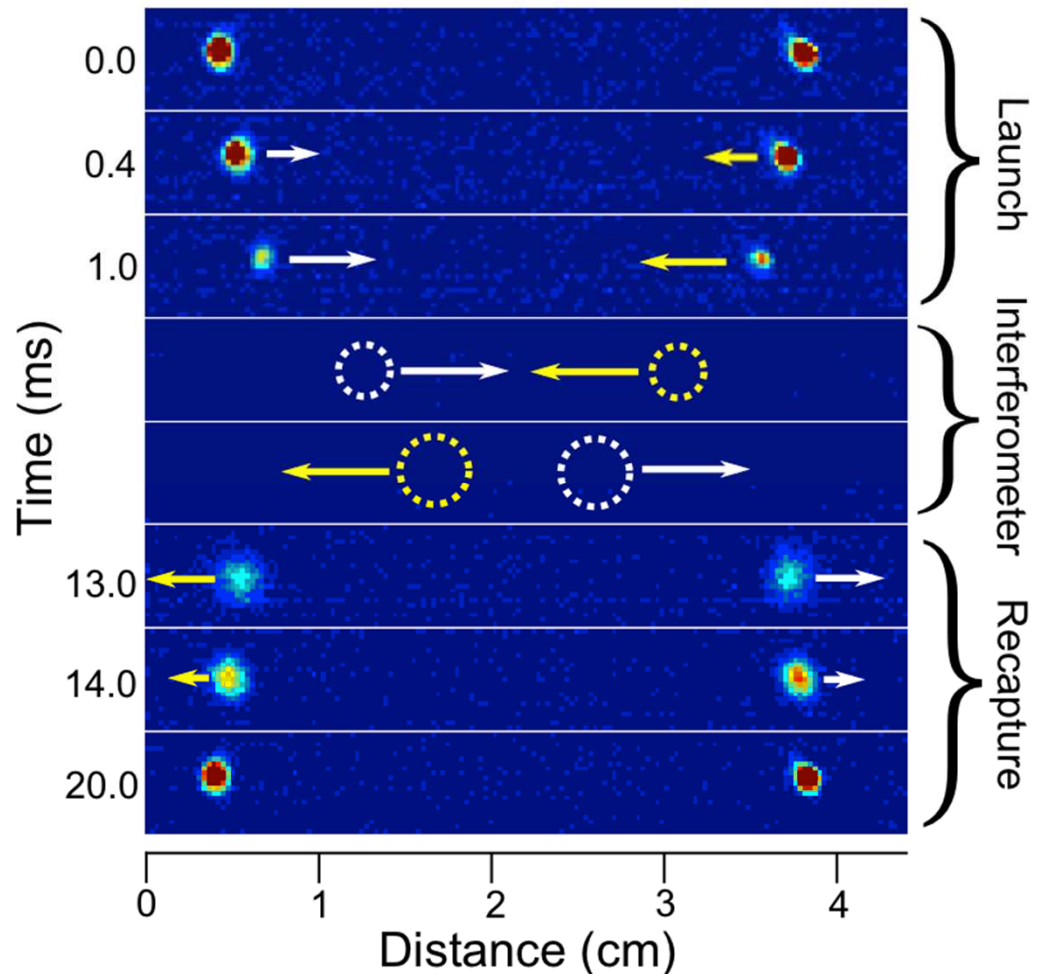
$$N_s = \frac{\alpha \eta T_c}{\beta T_c + (1 - r_0)}.$$

Base recapture efficiency $r_0 = 96\%$

Benefits

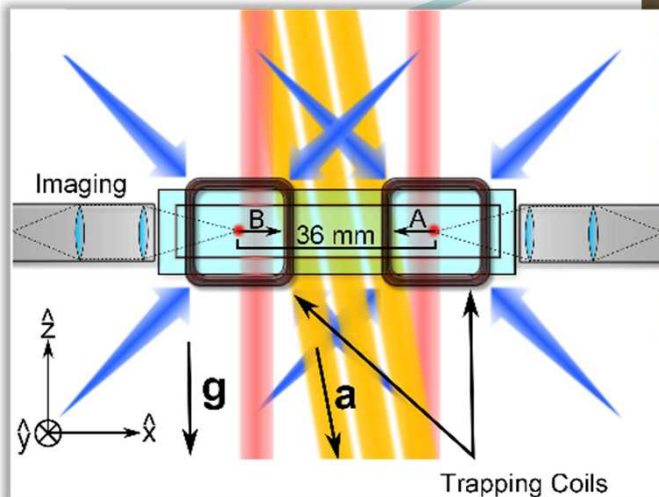
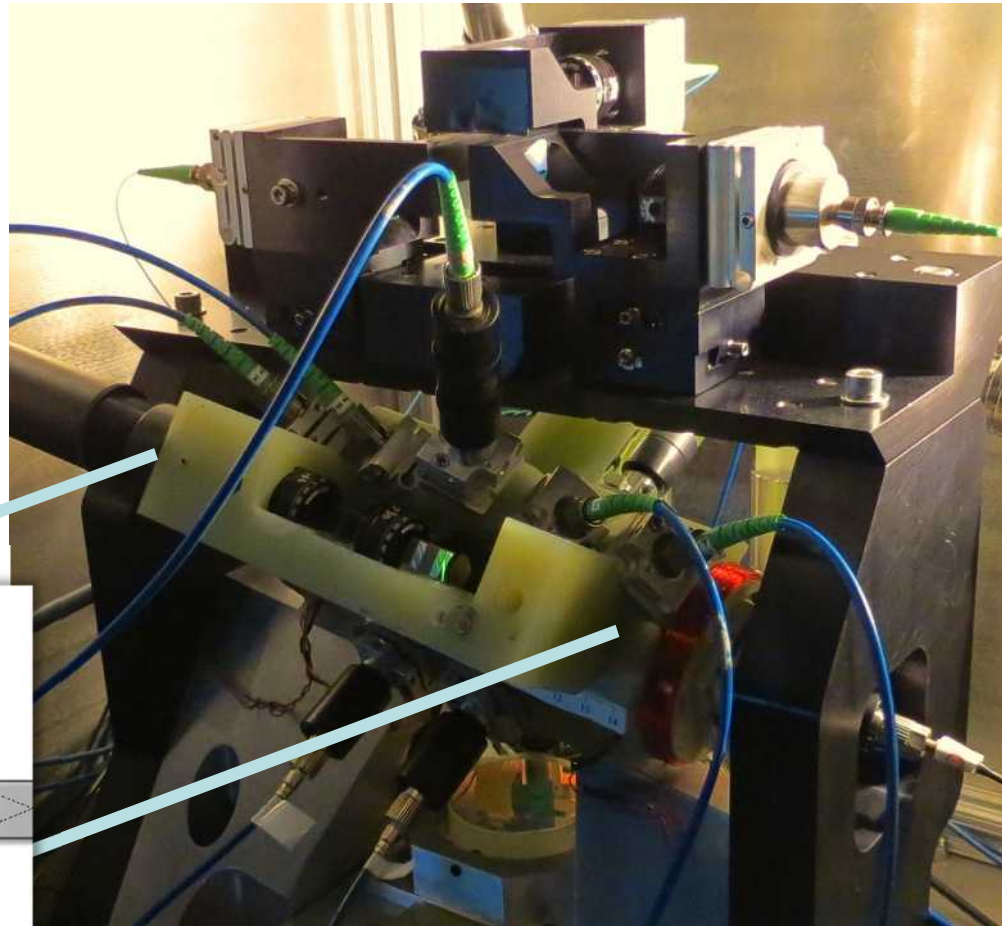
- Increases signal by 10x
- Data rates > 50 Hz
- Minimizes cycle dead time
- Reduced complexity

CCD images of ensemble exchange



EXPERIMENT PLATFORM

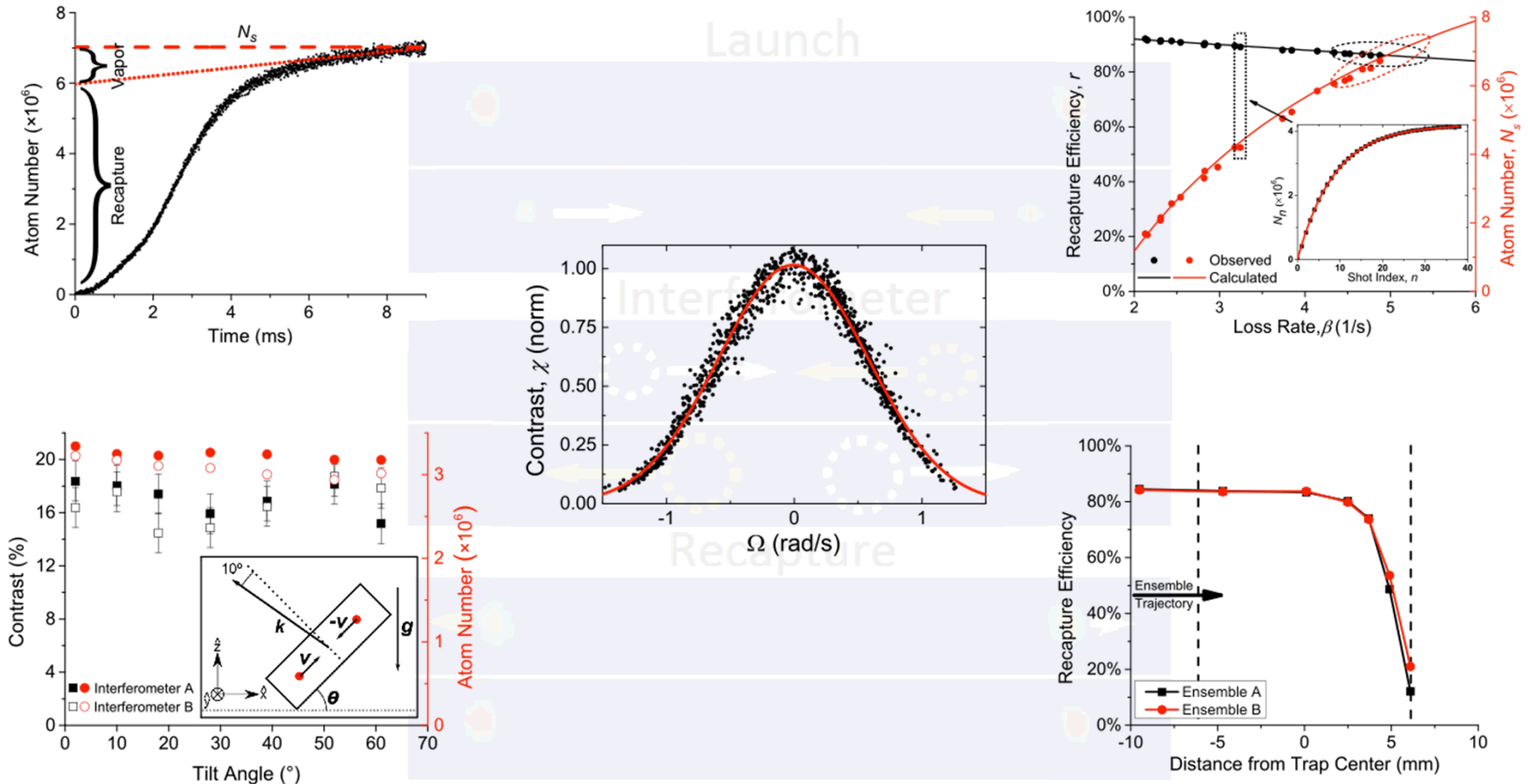
Picture of interferometer sensor



ENSEMBLE EXCHANGE TECHNIQUE

Sandia National Laboratories

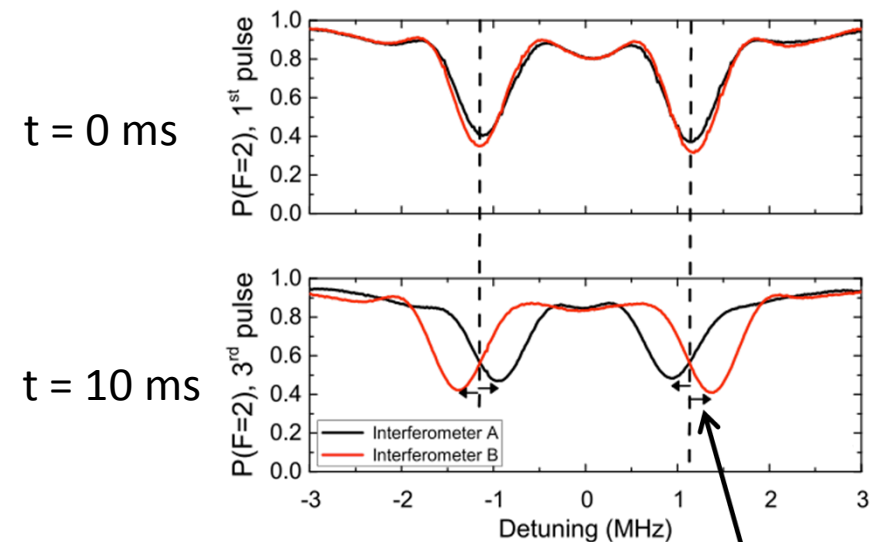
- Dynamic aspects of Ensemble Exchange characterized in previous work
- Robust to rotations, tilts and displacements



TECHNICAL CHALLENGE

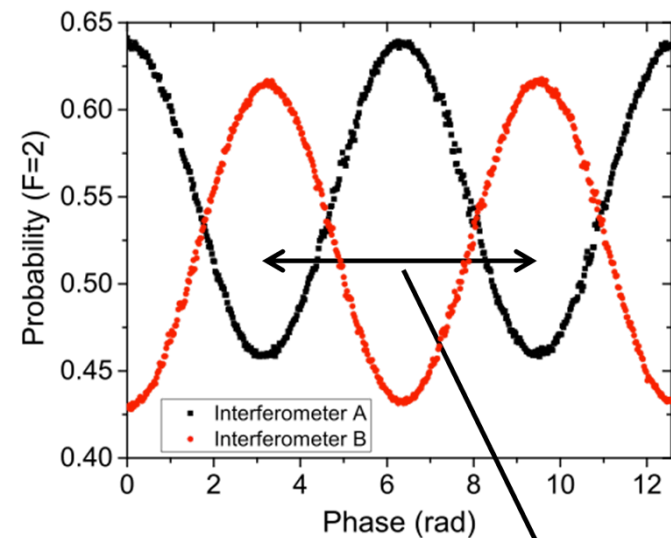
Controlling frequency and phase in a dynamic environment

Raman laser frequency scan



Significant detuning develops over 10 ms due to gravitationally-driven Doppler shift

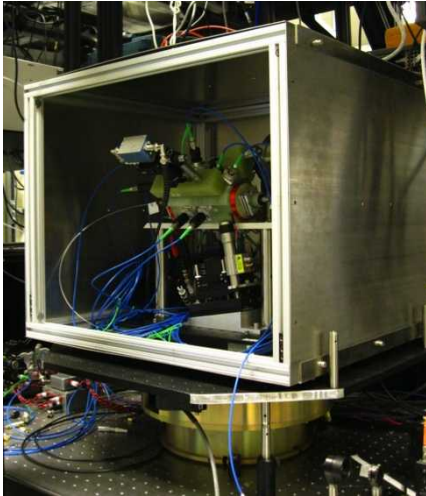
Raman laser phase scan



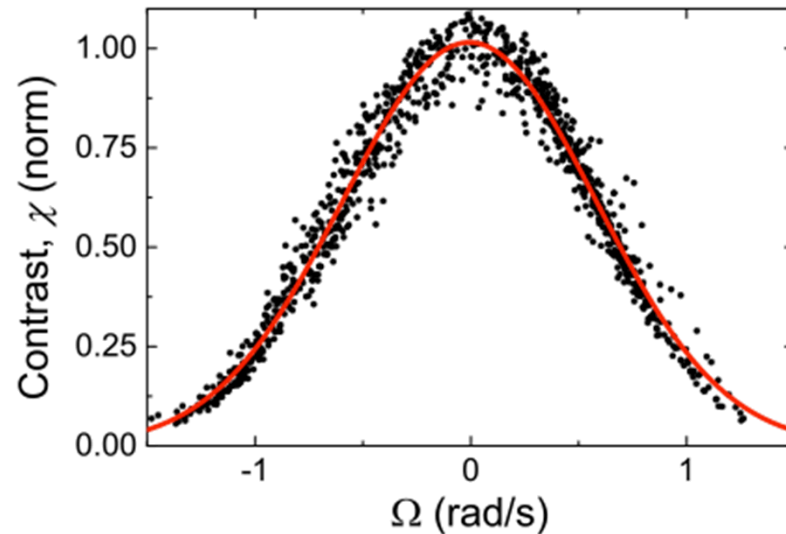
\approx mg OR \approx mrad/s (400 deg/hr)

Small changes in acceleration or rotation during one interferometer cycle cause phase slip

Rotation rate limit



Sensor on rotation stage



$$\Delta\phi = \vec{k}_{eff} \cdot (\vec{a}T^2 - 2(\vec{v} \times \vec{\Omega})T^2)$$

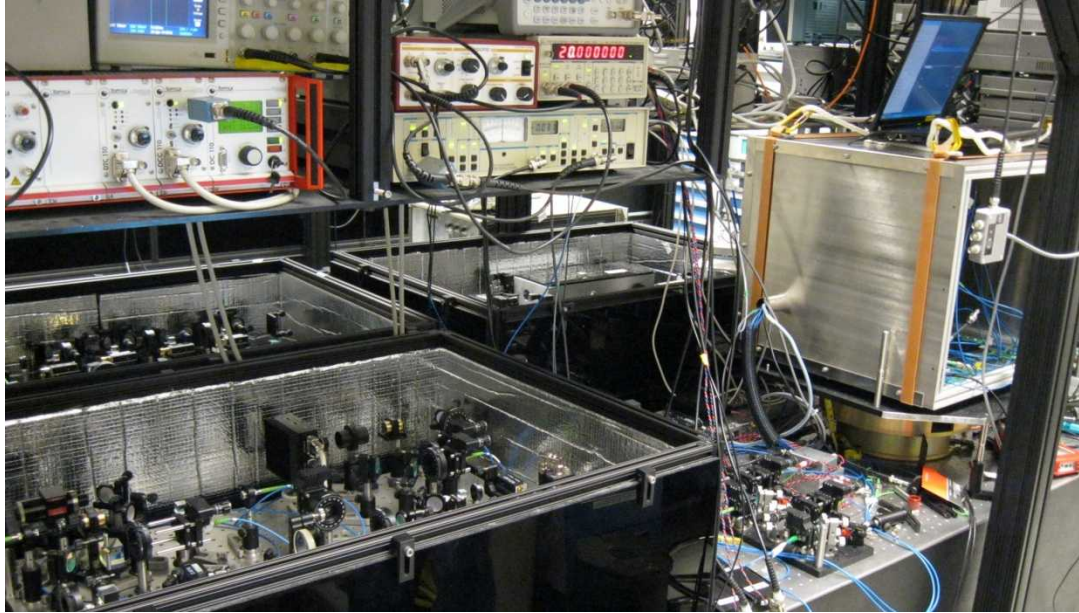
Contrast limited by temperature

$$\chi(\Omega_z) = \exp\left(\frac{-\Omega_z^2}{2\sigma_\Omega^2}\right)$$

$$\sigma_\Omega = \frac{1}{2} \sqrt{\frac{m}{kT}} \frac{1}{k_e T^2 \sin \theta}$$

Solutions: 3D readout or ultra-short T

Exploring ultra-short T



Sandia atom interferometer

- Laser cooled ensemble
- >1,000,000 cc

- Want a small, low-power, low-cost atomic accelerometer
- Warm vapor approach has historically made excellent gyros, clocks and magnetometers.
- What about accelerometers?

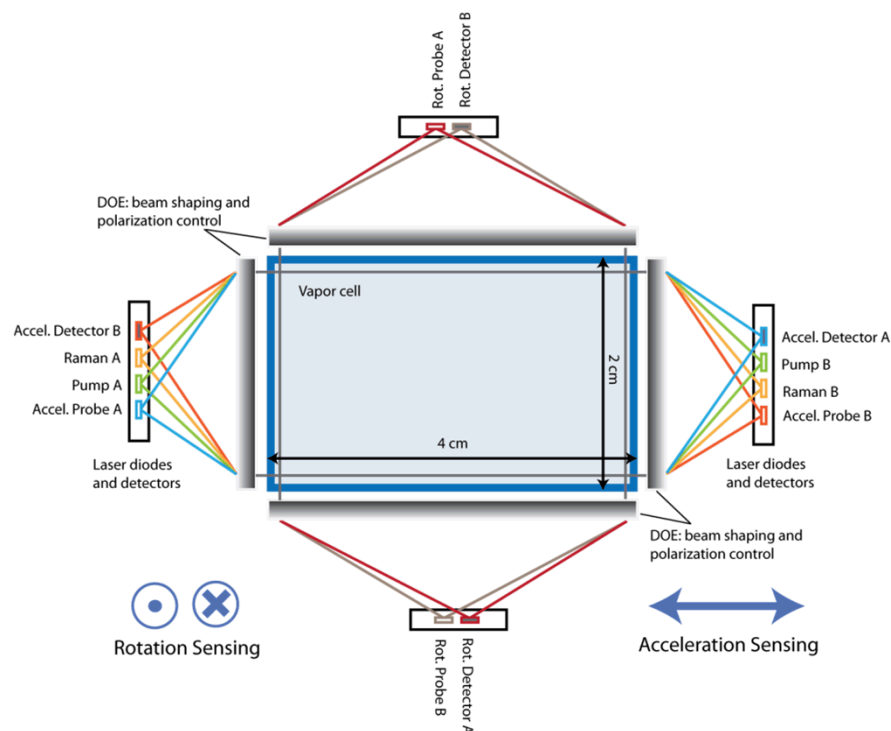


Symmetricom SA.45s CSAC

- Warm vapor ensemble
- 17 cc

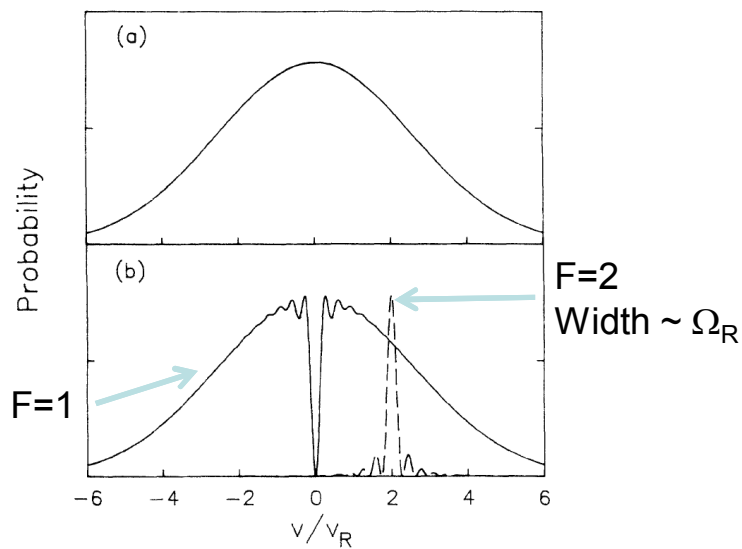
Warm vapor concept

- Using the Doppler sensitivity of Raman transitions, and ultra-short duration atom interferometry, LPAI is possible in a warm vapor
- The challenge: Target atom shot noise limit on 10^8 atoms



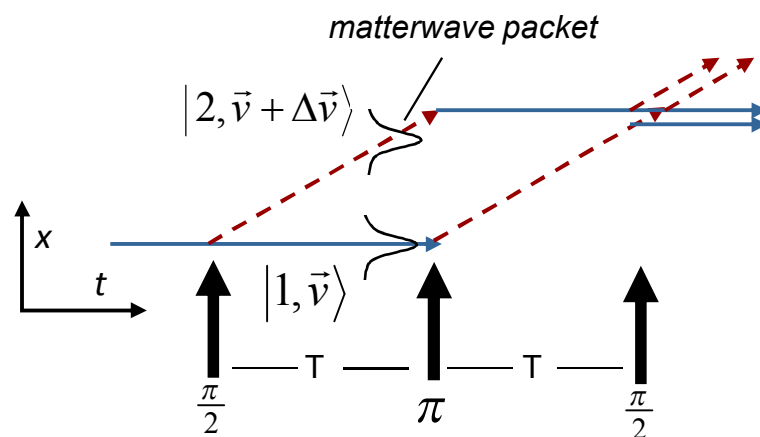
Motivation: potentially highly compact and simplified. Conceptual diagram (not to scale) of a 2-axis atomic sensor.

Basic Idea



Result of Raman transition used for LPAI

K. Moler, et al., *Phys. Rev. A*, 45, 342 (1992)

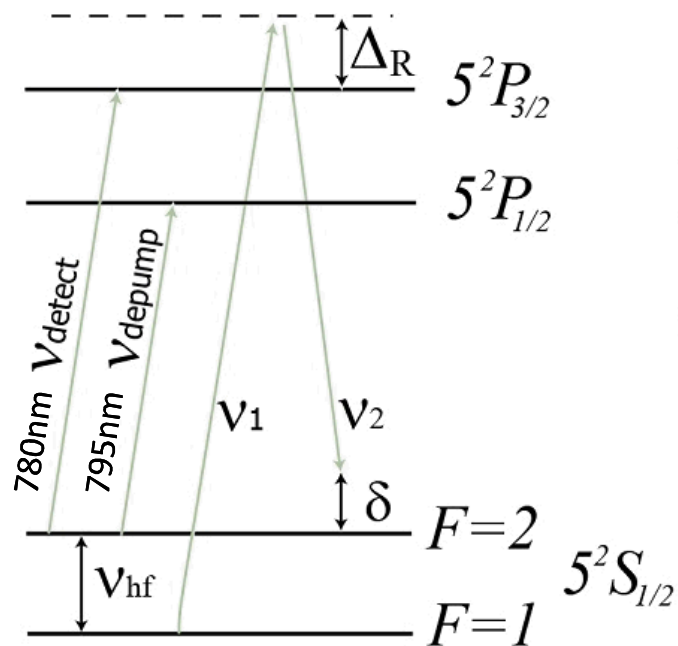


Analyze in momentum space

- each atom interferes with itself
- Non-ideal paths do not contribute to fringe

$$\Delta\phi = \mathbf{k}_{\text{eff}} \cdot (\mathbf{g} - 2\mathbf{v} \times \boldsymbol{\Omega}) T^2$$

Basic Idea

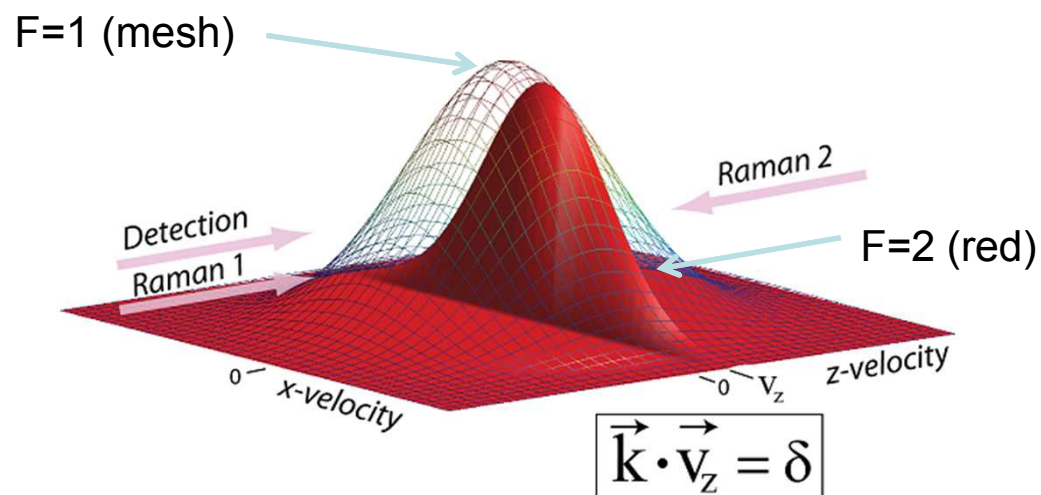
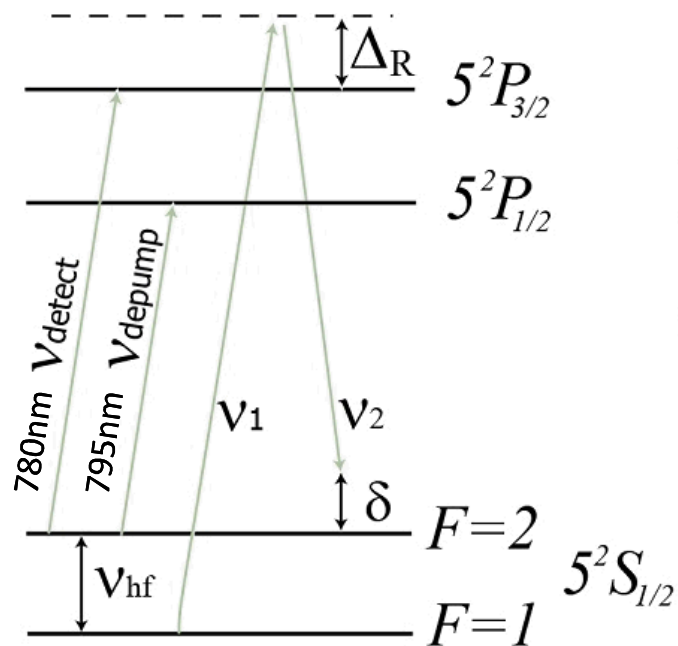


Sequence of events

- Prepare atoms in $F=1$ ground state
- Drive atom interferometer pulse sequence
- Detect in previously vacant state

Simplified laser frequency requirements:
potential for miniaturization

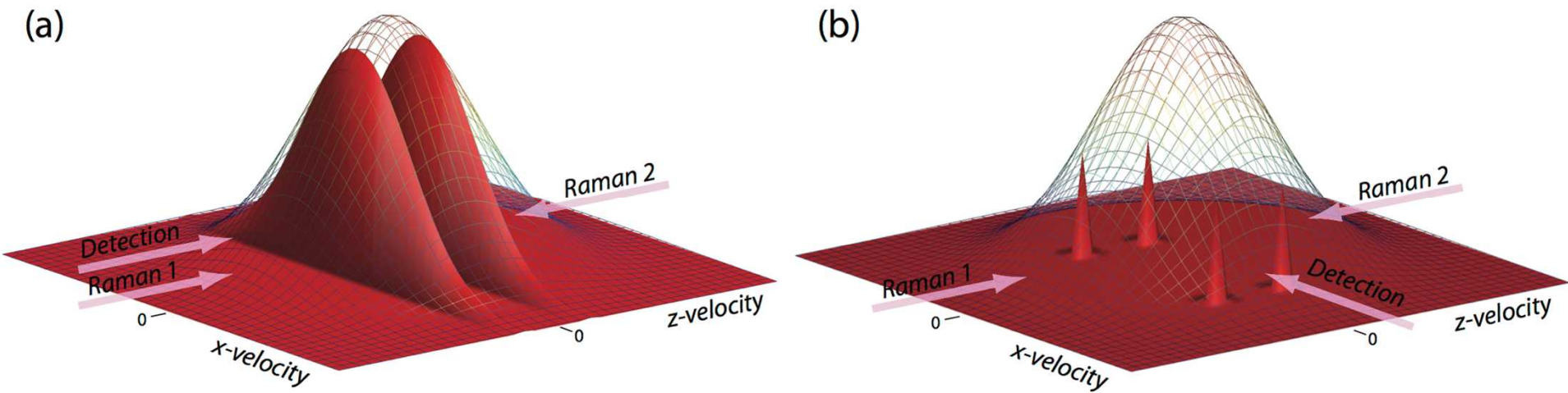
Basic Idea



MB distribution with laser fields

- use detuning δ to drive and probe nonzero velocity class
- Width set by Rabi frequency

Basic Idea



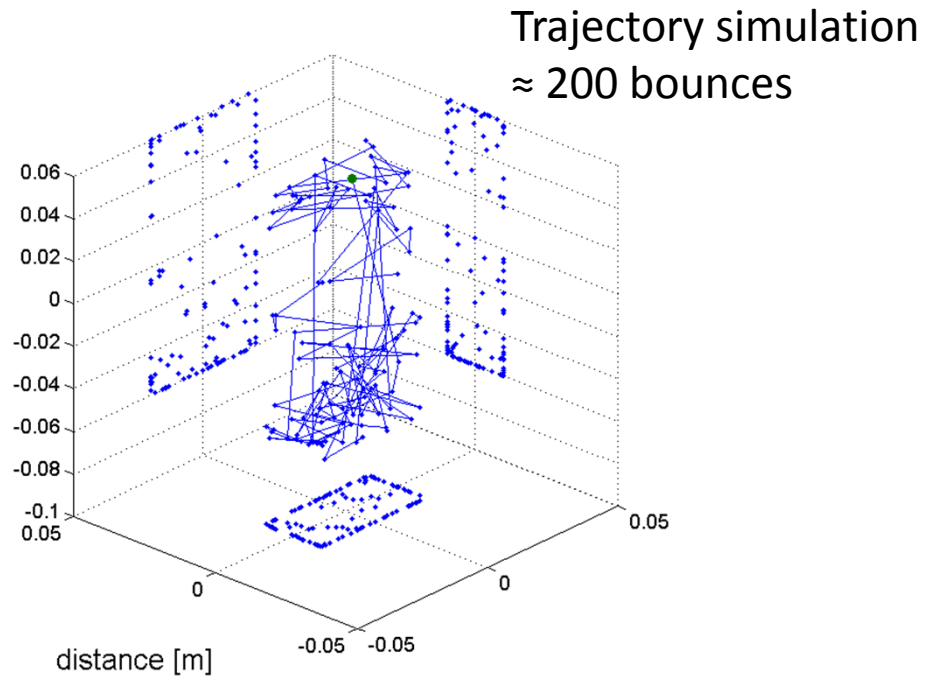
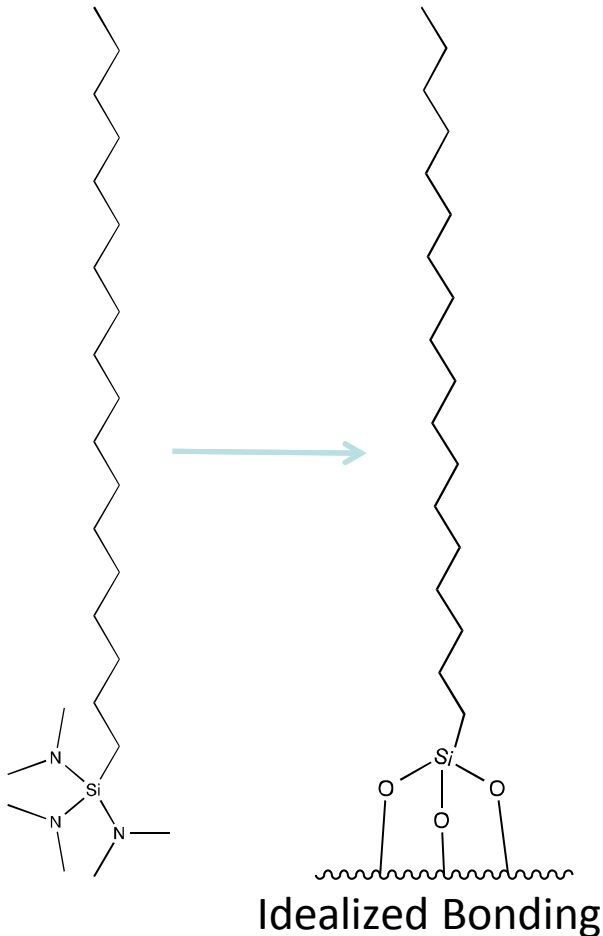
State preparation

To eliminate background signal, must depump all velocity classes

In-house coating development

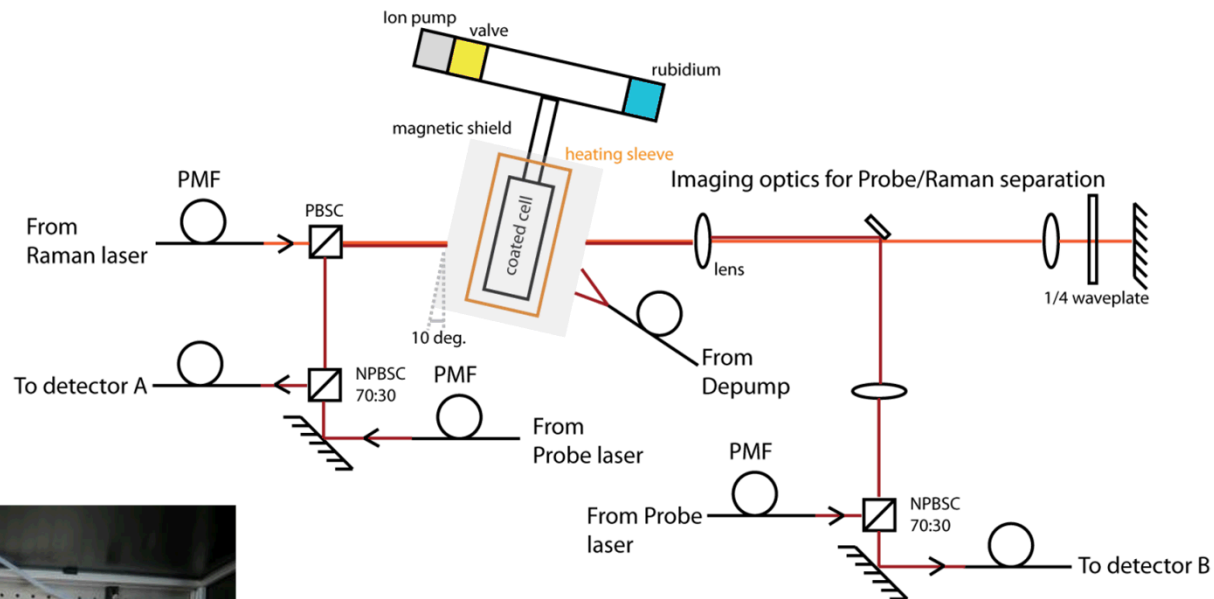
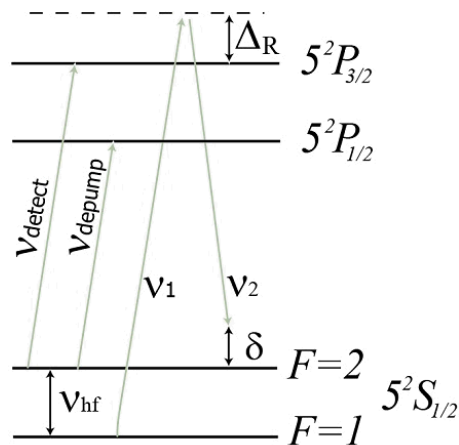
Tris(N,N-dimethylamino)octadecylsilane

measured: $\tau = 23$ ms hyperfine state lifetime

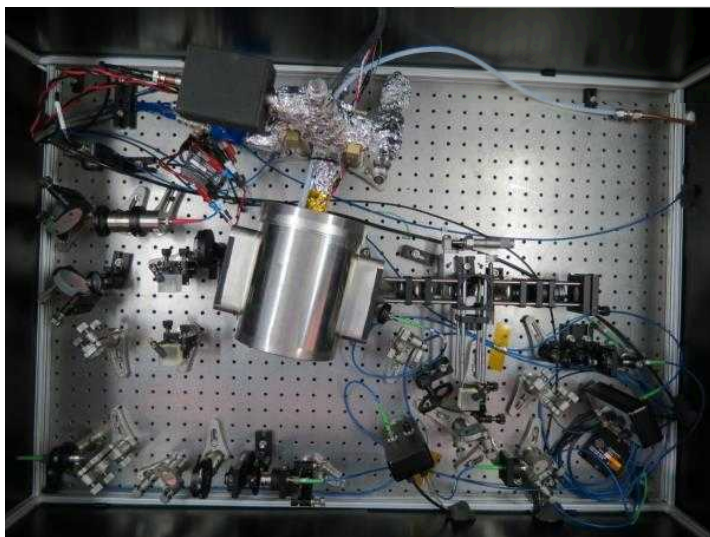


Must be inertially pure—no collisions!
Coating outgassing $< 10^{-4}$ Torr
No buffer gas

The experiment



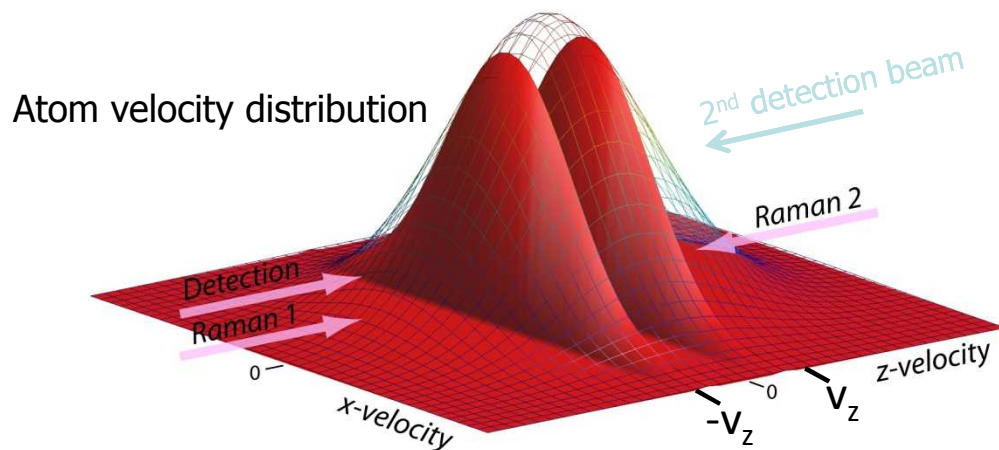
Experiment schematic



Experiment Photograph

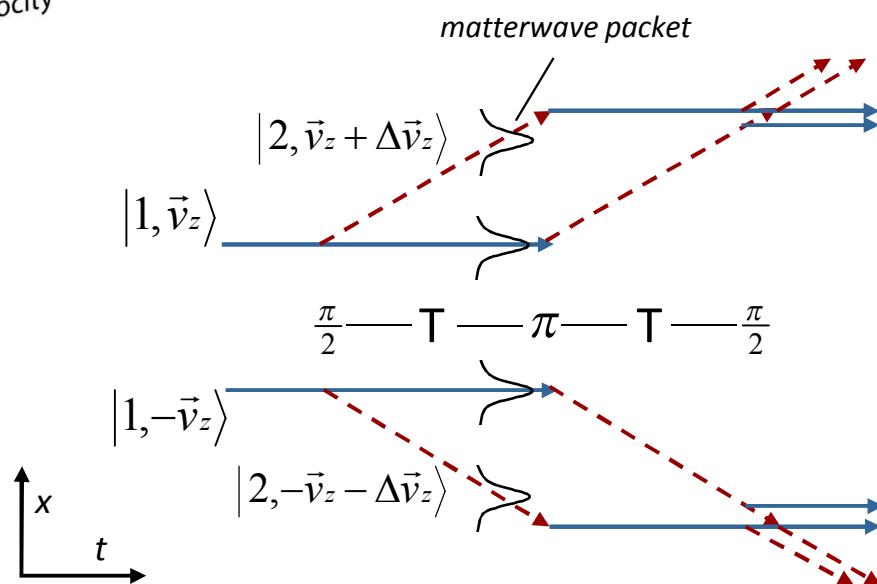
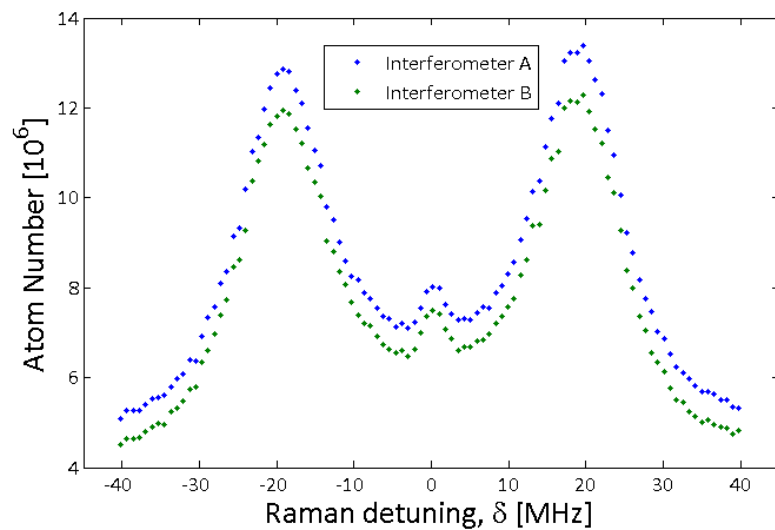
New possibilities

Common mode noise rejection

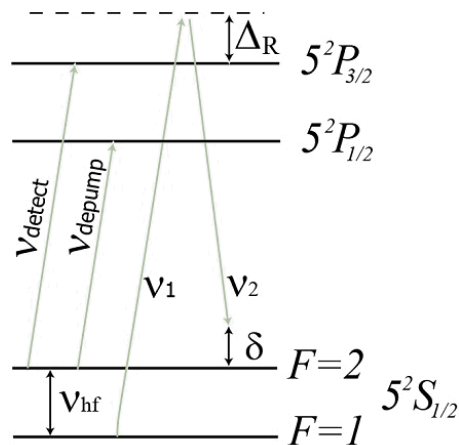


$$\Delta\Phi = \mathbf{k} \cdot \mathbf{a} T^2$$

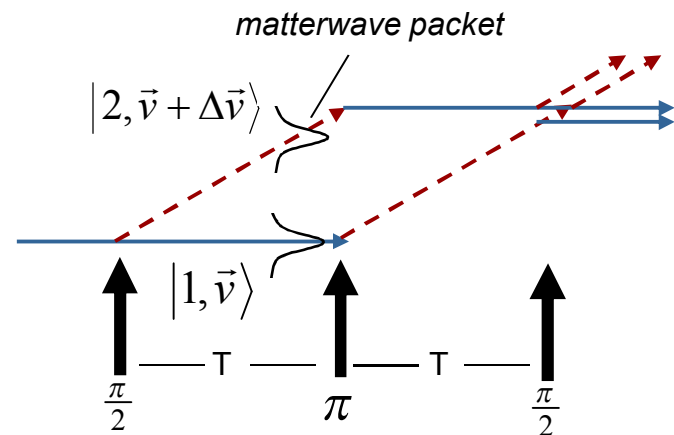
Simultaneous interferometers



Experiment timing



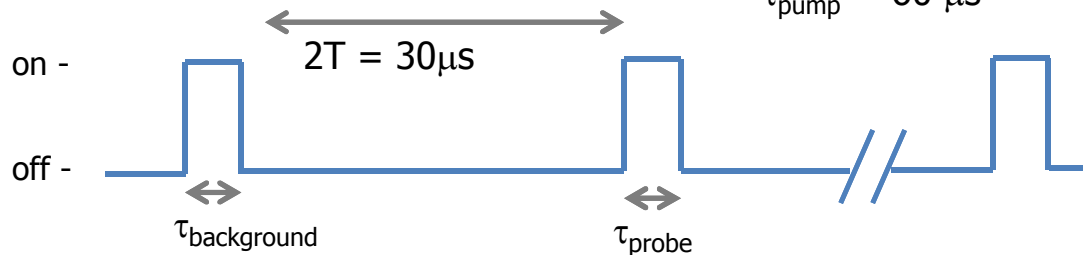
Timing diagram



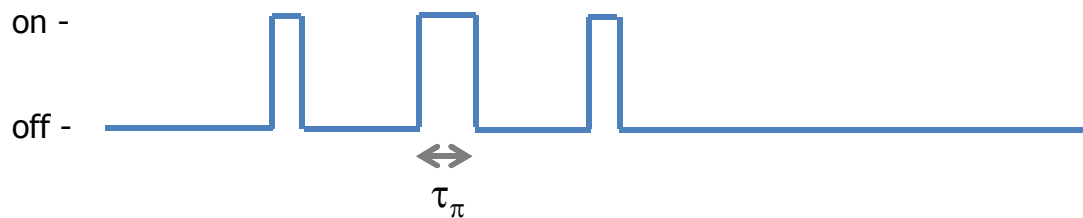
Depumping laser



Probe laser



Raman laser



Wavepacket overlap in a “warm” vapor

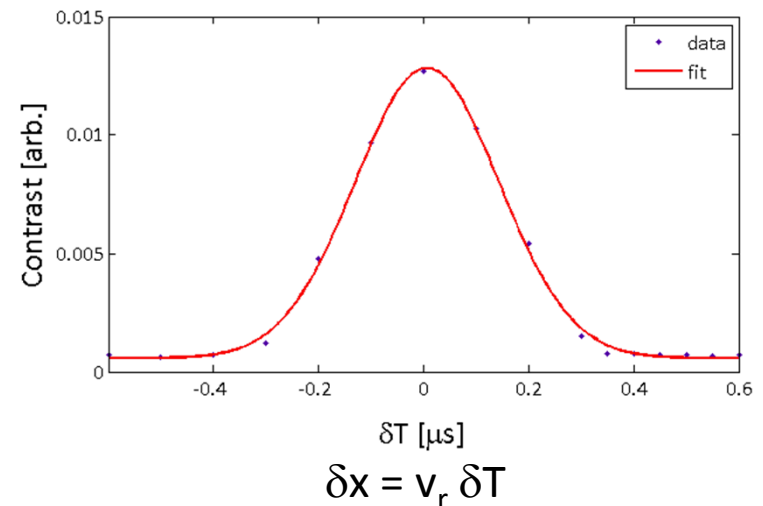
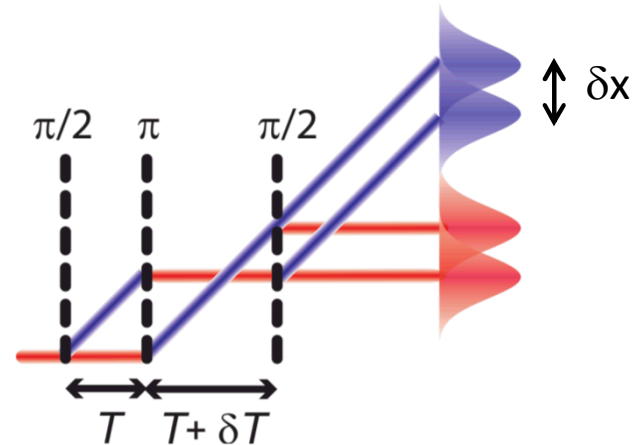
The contrast, $\chi(\delta T)$, is given by

$$\chi(\delta T) = \exp\left(-\frac{v_r^2 \delta T^2}{8x_a^2}\right)$$

where v_r is the recoil velocity (11.8 mm/s)
and x_a is the average coherence length x_a
 $= 0.81(3)$ nm.

This corresponds to a temperature of
1.44(6) mK

Consistent with a Doppler width of 0.95(4) MHz
Note: Rabi frequency is 1.61 MHz



Wavepacket overlap in a “warm” vapor

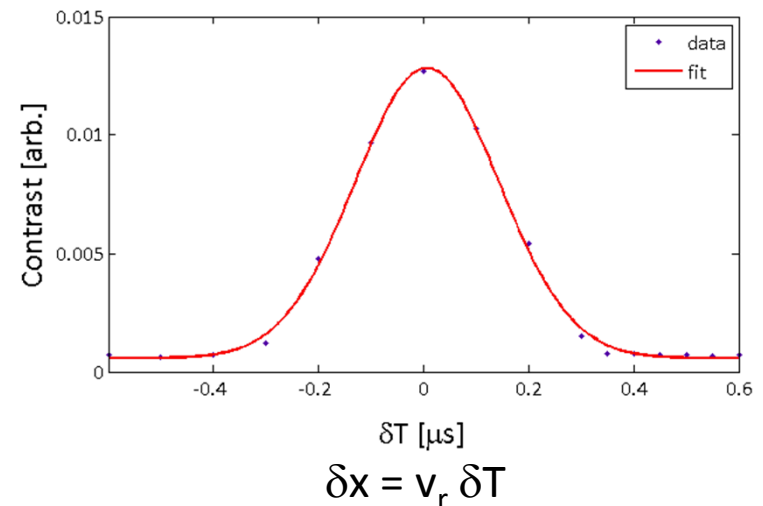
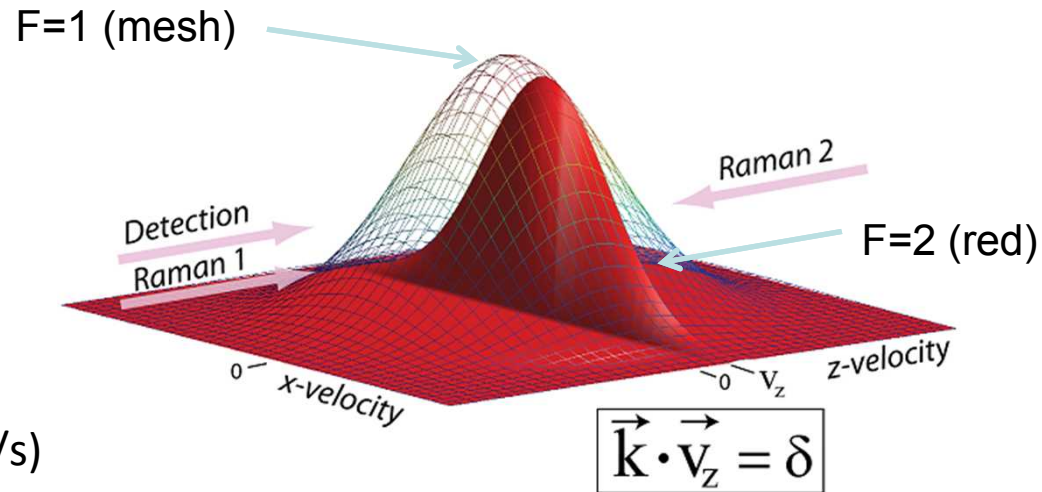
The contrast, $\chi(\delta T)$, is given by

$$\chi(\delta T) = \exp\left(-\frac{v_r^2 \delta T^2}{8x_a^2}\right)$$

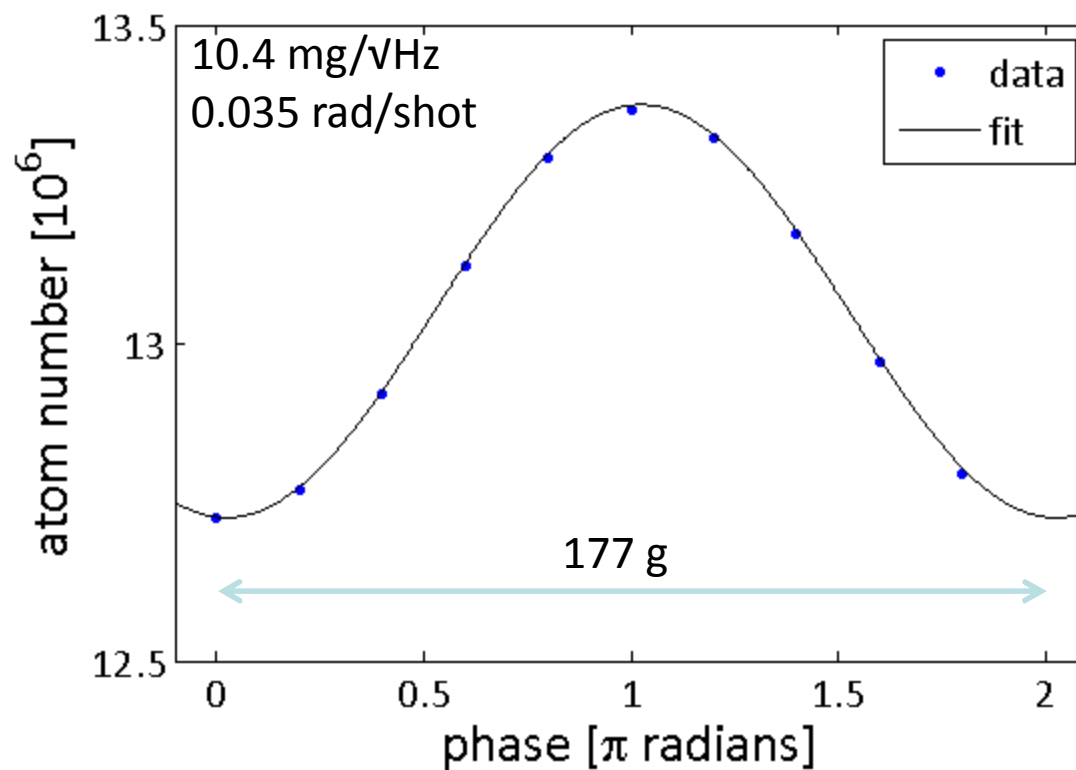
where v_r is the recoil velocity (11.8 mm/s) and x_a is the average coherence length $x_a = 0.81(3)$ nm.

This corresponds to a temperature of
1.44(6) mK

Consistent with a Doppler width of 0.95(4) MHz
Note: Rabi frequency is 1.61 MHz



Results

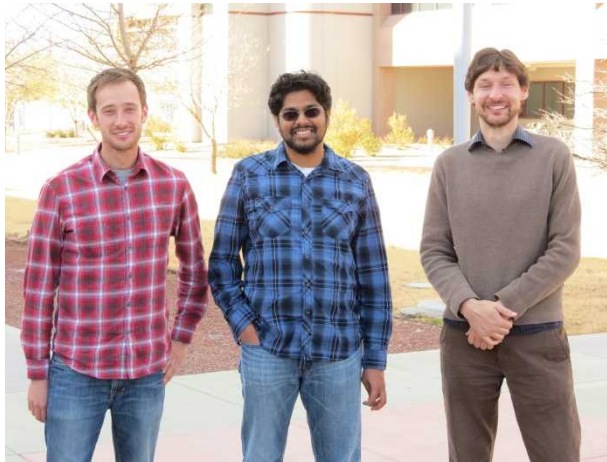


Measured:

- without CMR: 40 mg/VHz
- With CMR: 10.4 mg/VHz

Atom shot noise: 3.1 mg/VHz

Thank you



G. Biedermann, A. Rakholia, H. McGuinness

Not pictured

- Yuan-Yu Jau
- David Wheeler