

Continuum-scale modeling of hydrogen and helium bubble growth in metals

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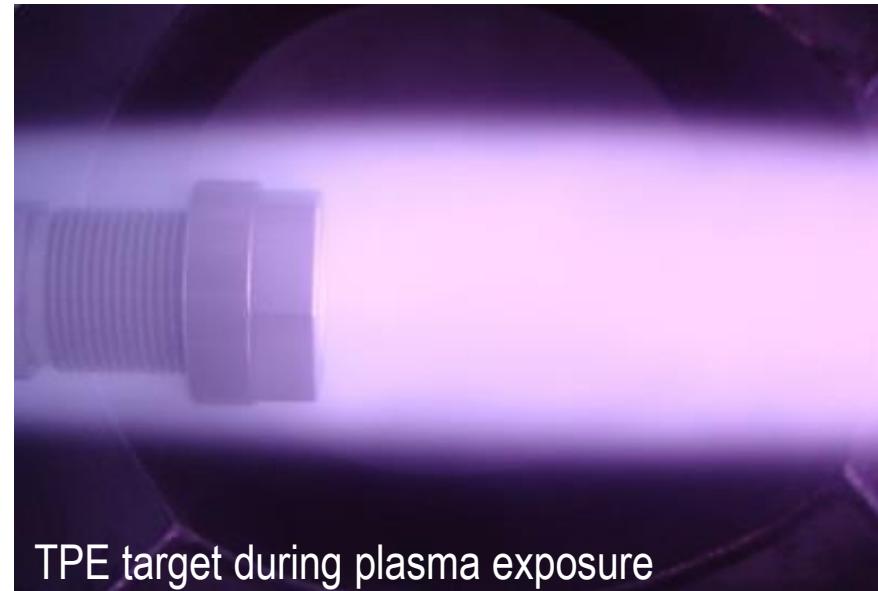
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Motivation: Analysis of bubble growth in ITER-grade W samples exposed in TPE

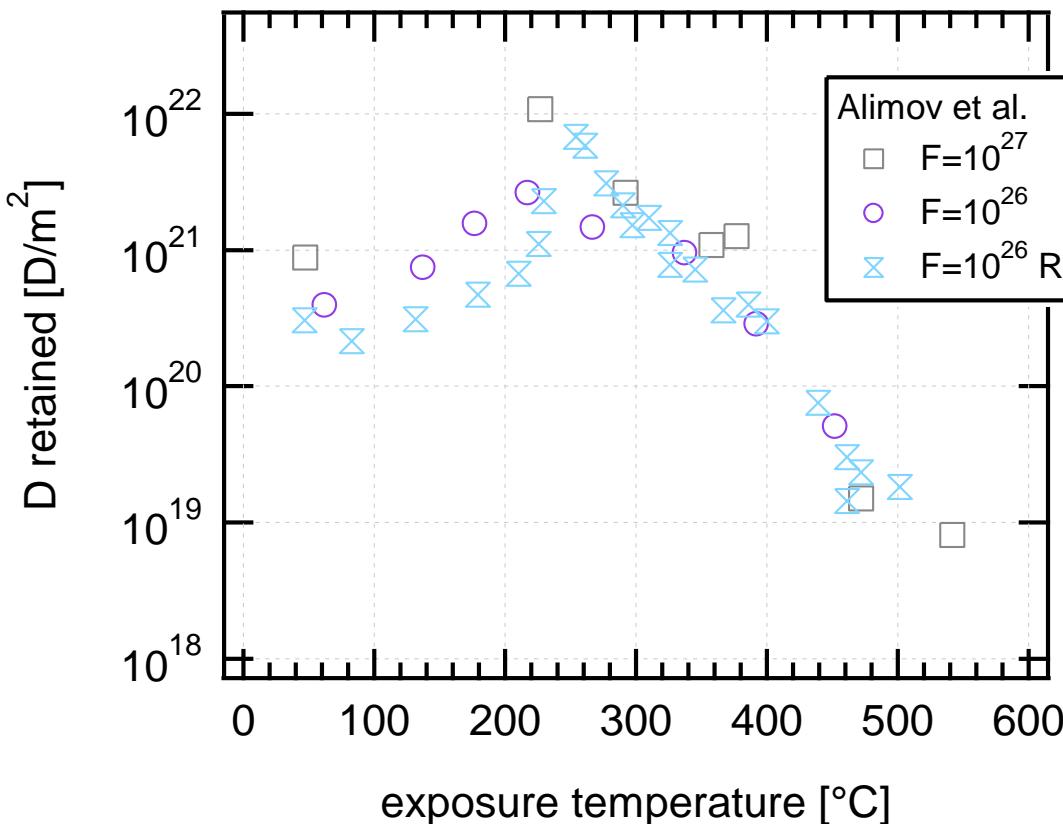
- Precipitation affects migration through material
- Bubble growth depends on microstructure
- Growth mechanisms critical to developing realistic models



| exposure type | ion energy [eV] | duration [min] | flux (Γ_i) [$\text{m}^{-2} \text{ s}^{-1}$] | fluence (Φ) [m^{-2}] |
|---------------|-----------------|----------------|--|--|
| LF | 100 | 60 | 4.9×10^{21} | 1.8×10^{25} |
| HF | 100 | 120 | 1.5×10^{22} | 1.1×10^{26} |

- TPE plasma exposures at INL
- Microscopy at Shizuoka

Retention measurements correspond closely with those obtained in other laboratories



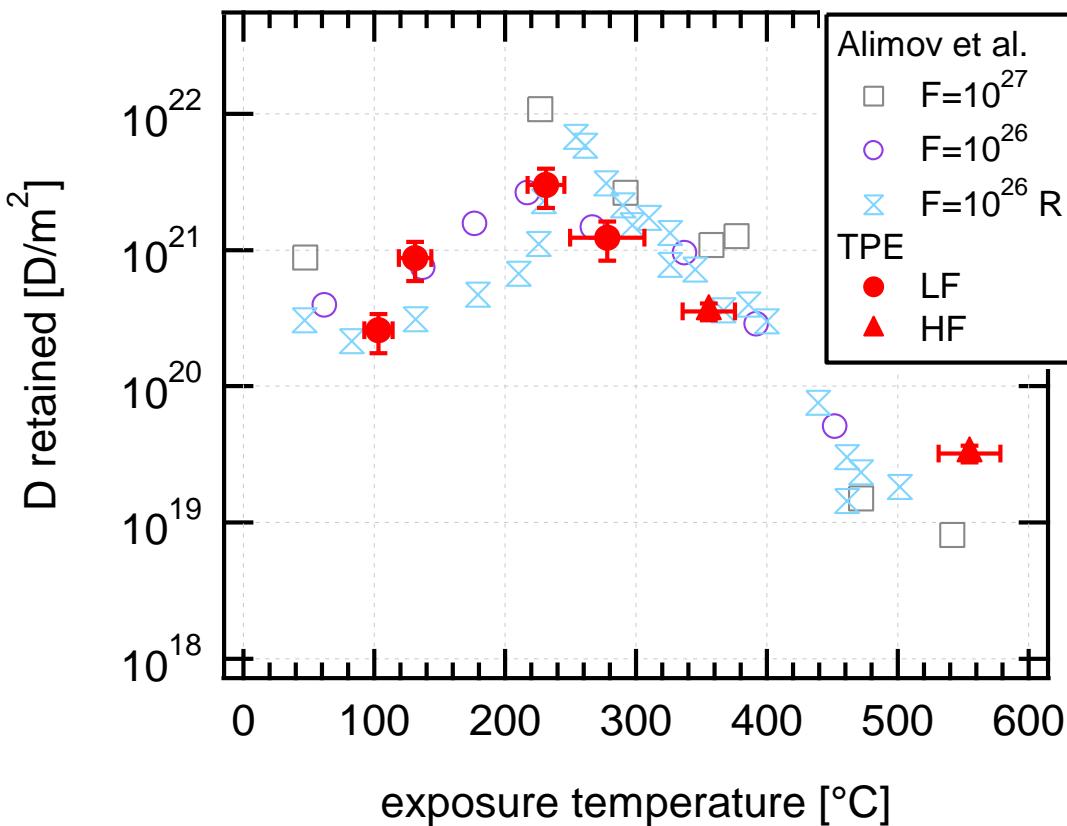
V. Kh. Alimov, et al. *J. Nucl. Mater.* **420** (2012) 519.

Previous work by Alimov et al:

- ITER-grade W
- $E = 38$ eV
- $\Phi = 10^{22} \text{ D m}^{-2} \text{ s}^{-1}$

Comparable exposure conditions

Retention measurements correspond closely with those obtained in other laboratories



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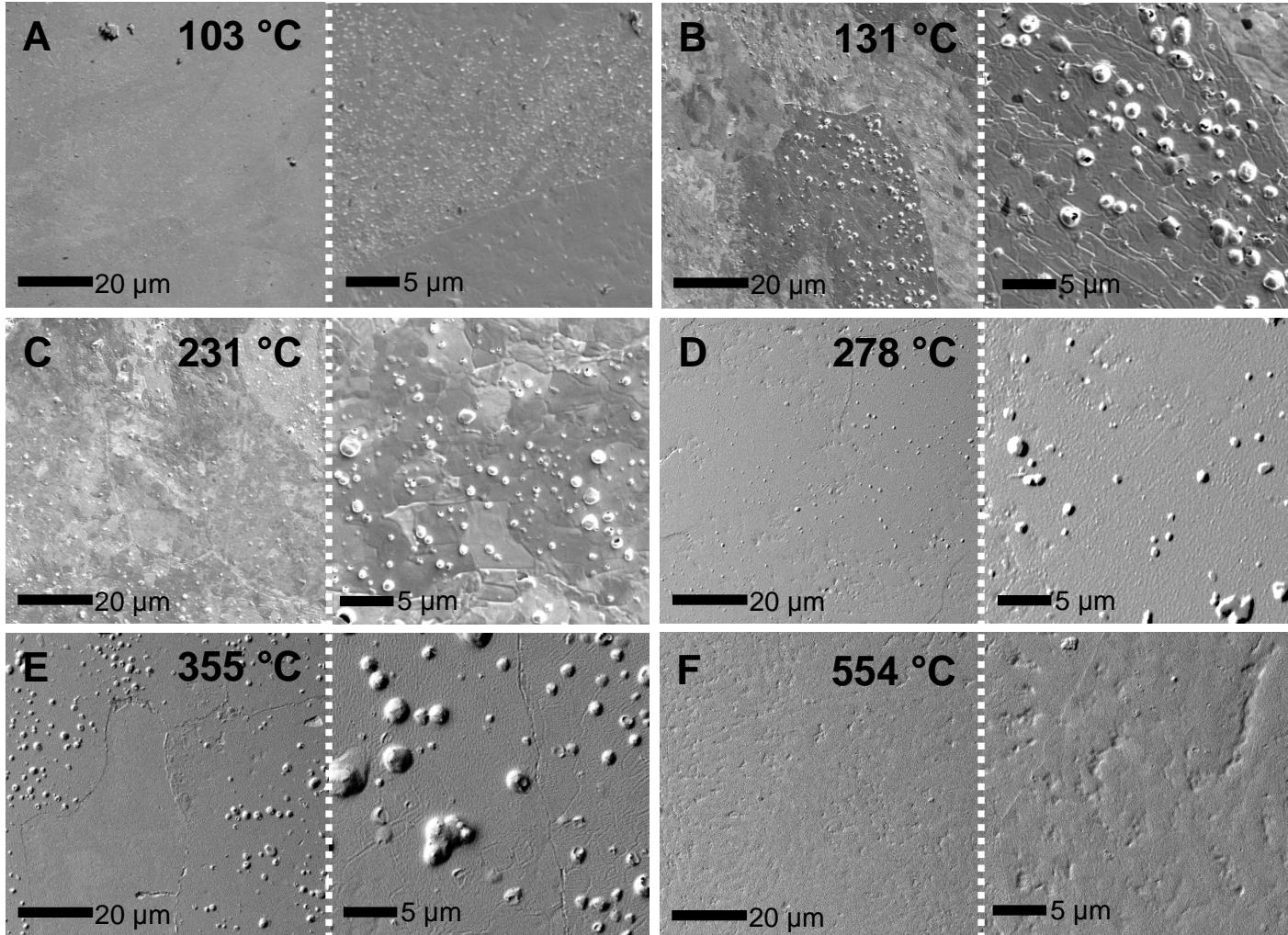
Comparable exposure conditions

TPE retention measurements:

- Correspond closely with Toyama/IPP meas.
- Confirm accepted retention temp. dependence.

V. Kh. Alimov, et al. *J. Nucl. Mater.* **420** (2012) 519.

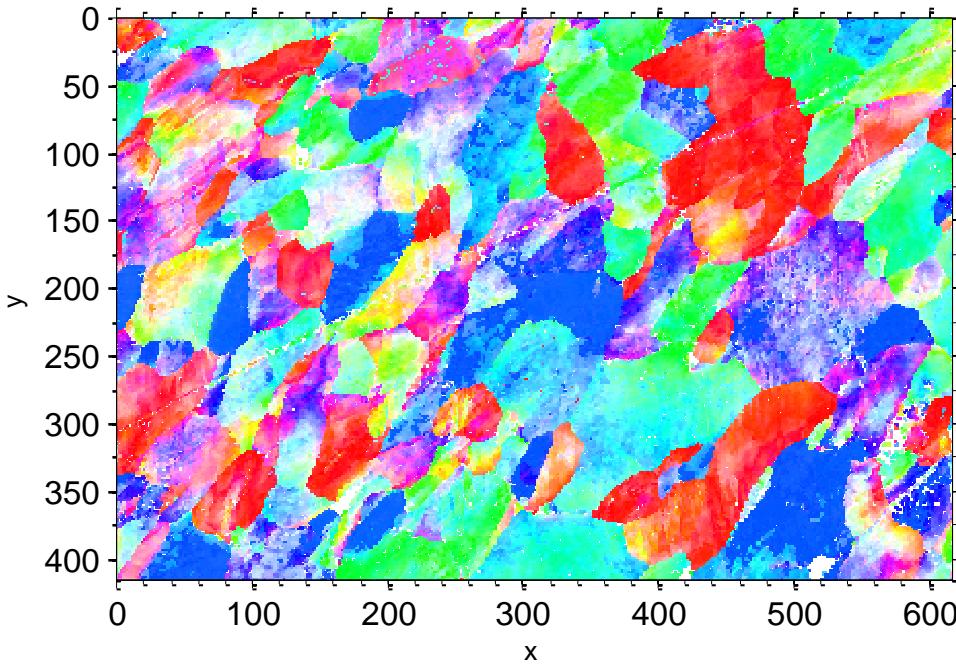
Surface morphology variation with temperature



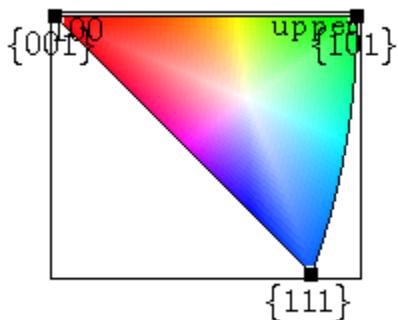
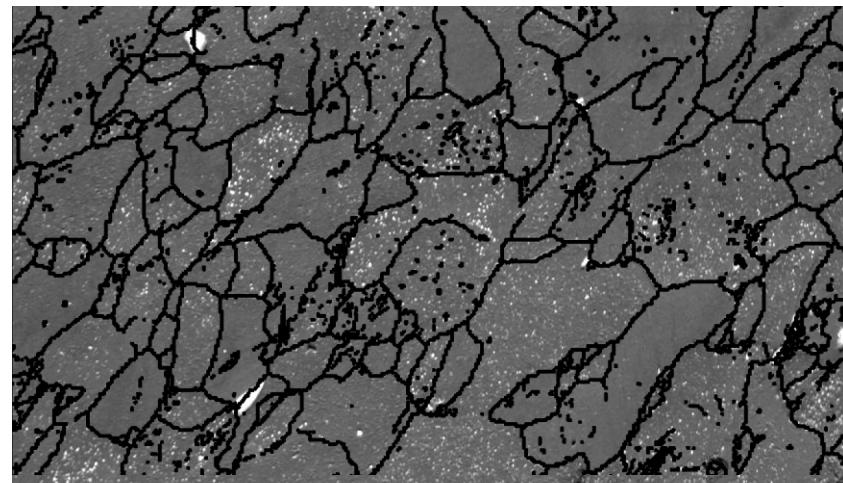
Key features:

- Non-uniform coverage
- Bubbles are small (<10 µm dia.) compared with warm-rolled W material.
- Absent at temperature extrema.

EBSD measurements reveal dependence on grain orientation

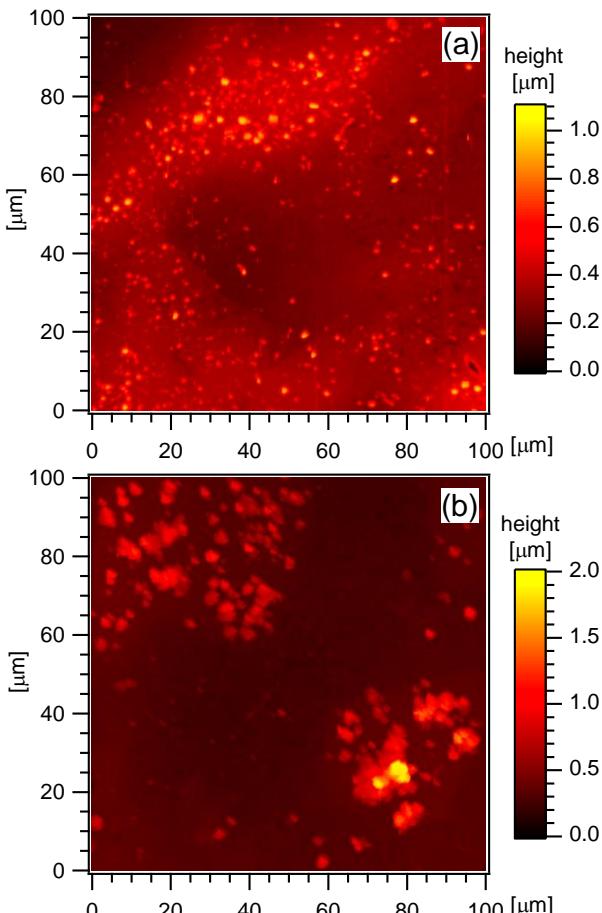


SEM image of the same area

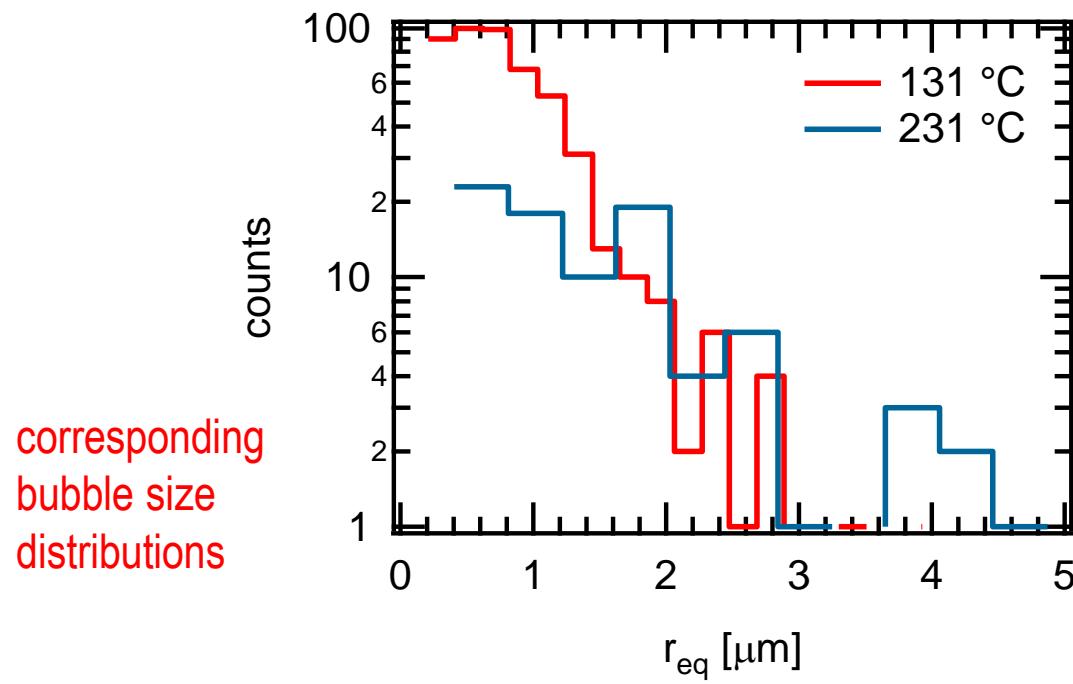


- Grain orientation indicated by inverse pole plot.
- Bubbles visible on grains with $\langle 111 \rangle$ and $\langle 110 \rangle$ directions aligned normal to surface
- Considerable distortion within individual grains
- Un-annealed sample showed increased distortion

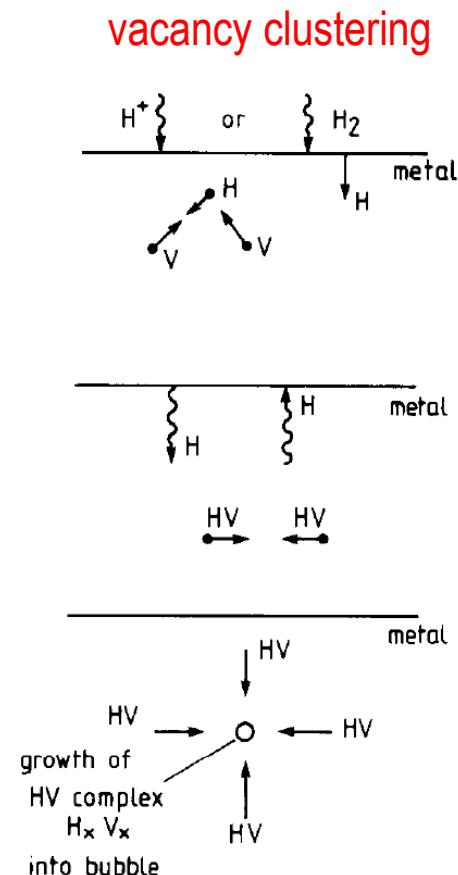
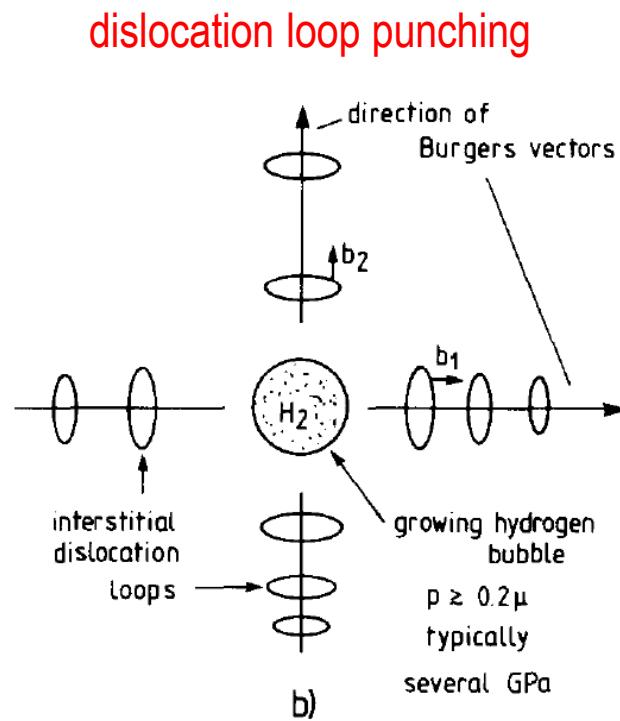
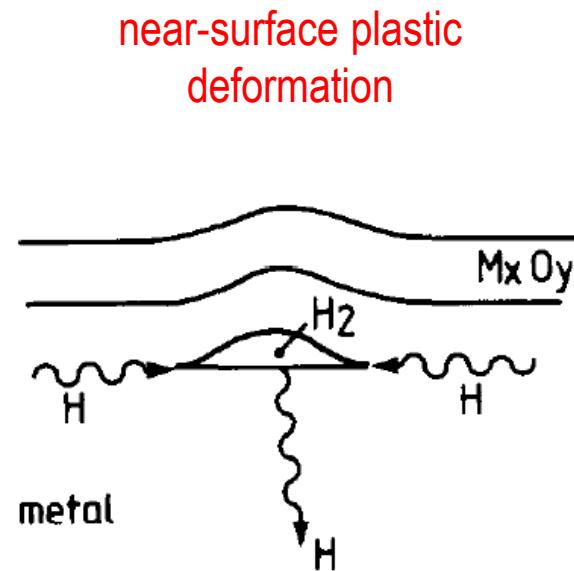
Atomic force microscopy reveals details of surface structure



- Atomic force microscopy provides information on the shape of the deformed surface.
- Individual bubbles identified and analyzed automatically.



What bubble growth mechanisms are active in W during plasma exposure?



Figures from: J. B. Condon & T. Schober, *J. Nucl. Mater.* **207** (1993) 1.

Far from the free surface, dislocation loop punching is favored

Three bulk precipitate growth mechanisms considered:

- Dislocation loop punching

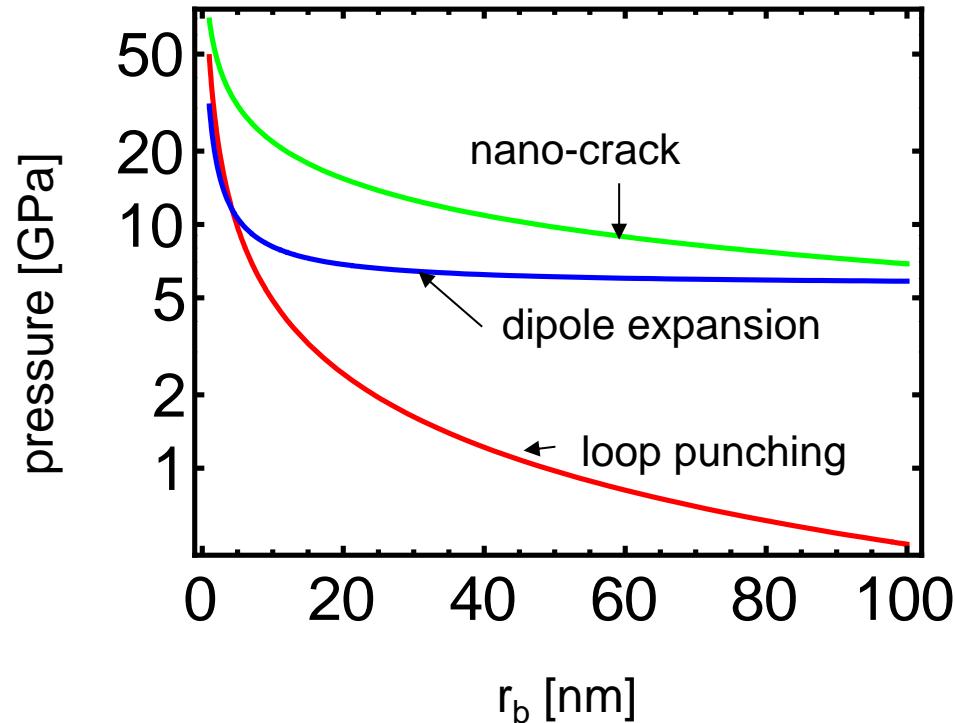
$$p_{LP} \geq \frac{2\gamma}{r} + \frac{\mu b}{r} \sim \frac{1}{r}$$

- Griffith nano-crack extension

$$p_{NC} \geq \sqrt{\frac{\pi\mu\gamma}{(1-\vartheta)r}} \sim \frac{1}{\sqrt{r}}$$

- Dislocation dipole expansion

$$p_{DE} \geq \frac{2\gamma}{s} + \frac{\mu d}{2r} \sim \frac{1}{r} + c$$



Based on methods developed in:
D. F. Cowgill, "Physics of He Platelets in
Metal Tritides," in *Effects of Hydrogen on
Materials* (2009).

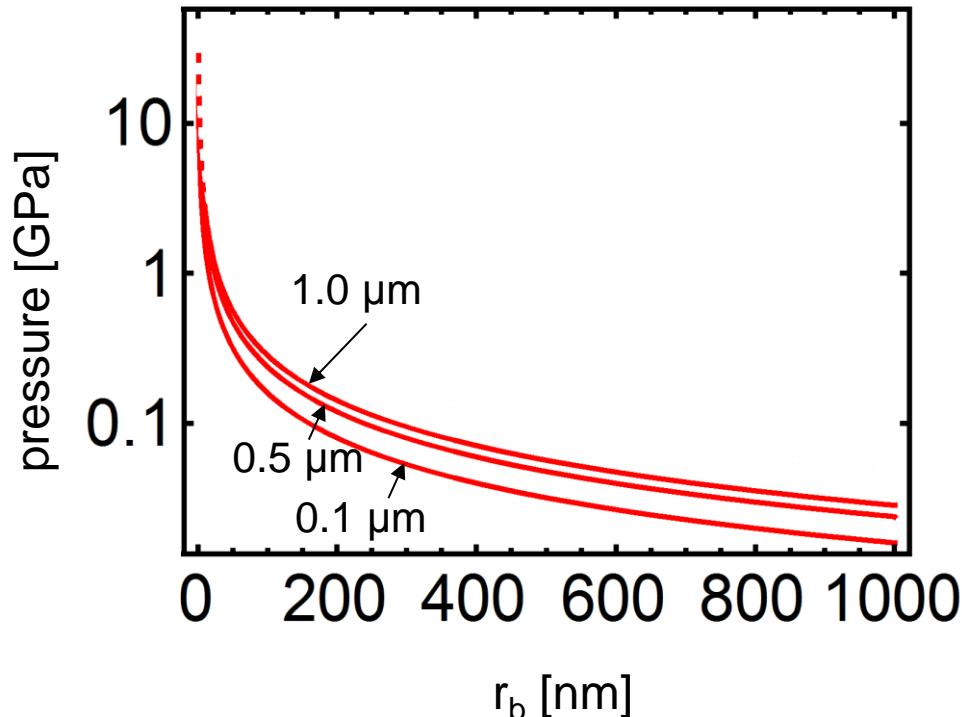
Near the free surface, bubbles may grow by crack extension

Crack extension competitive with loop punching near surface:

$$p_B \geq \frac{1}{r} \left(\frac{4\gamma(Eh)^{1/3}}{5C_1C_2} \right)^{3/4} \sim \frac{1}{r}$$

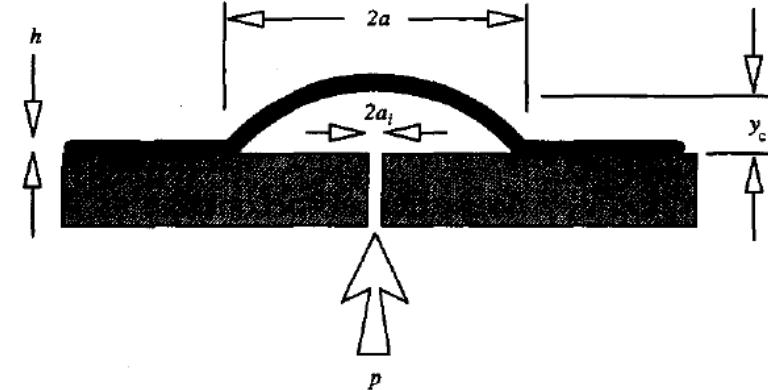
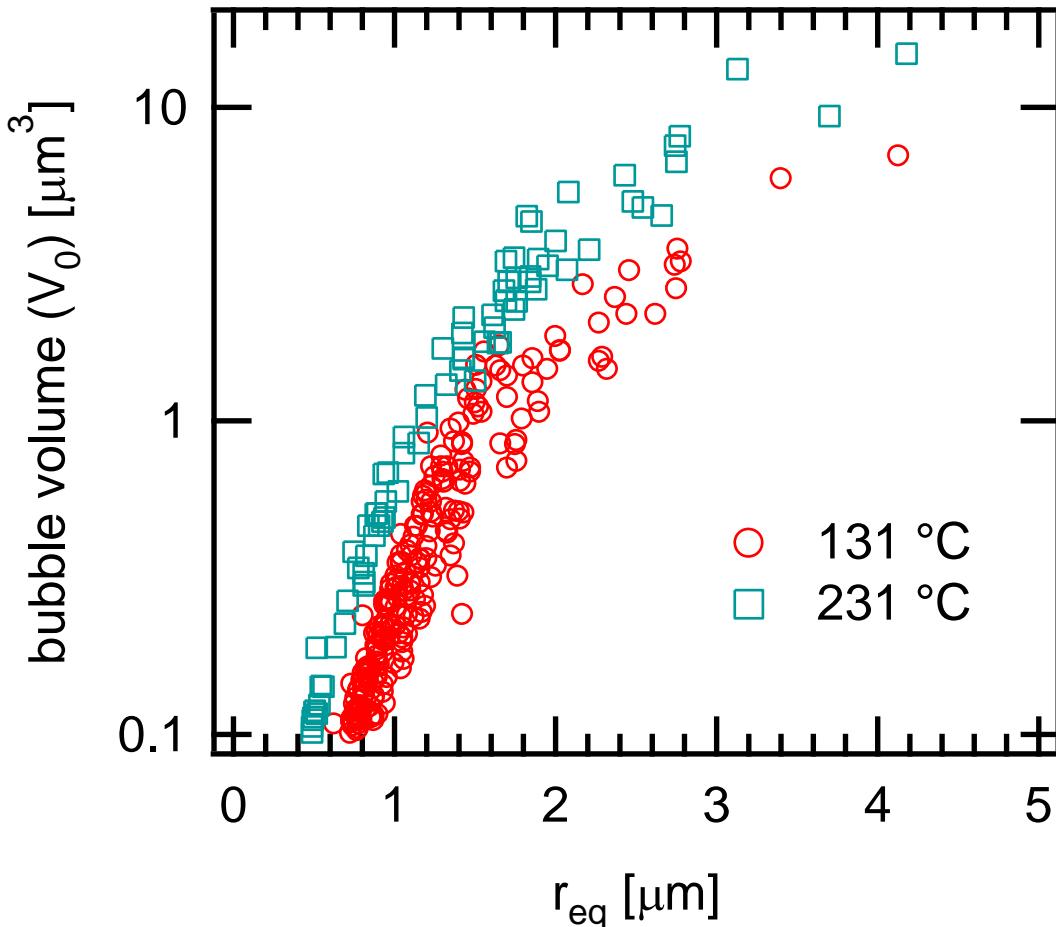
Limitations:

- Correction for thick blisters
- Effect of plasticity (blunting of crack tip)
- Hydrogen effects



Stress calculations based on calculations by K. Wan & Y. Mai, *Acta metall. mater.* **43** (1995) 4109.

Bubble volumes measured with AFM correlate well with blister model

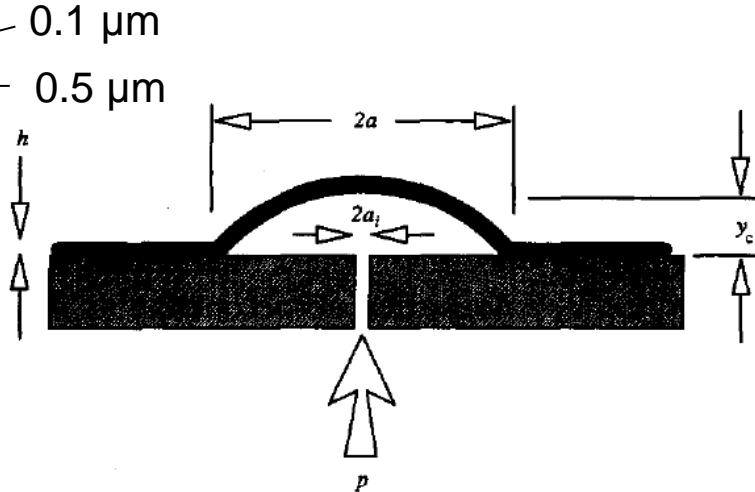
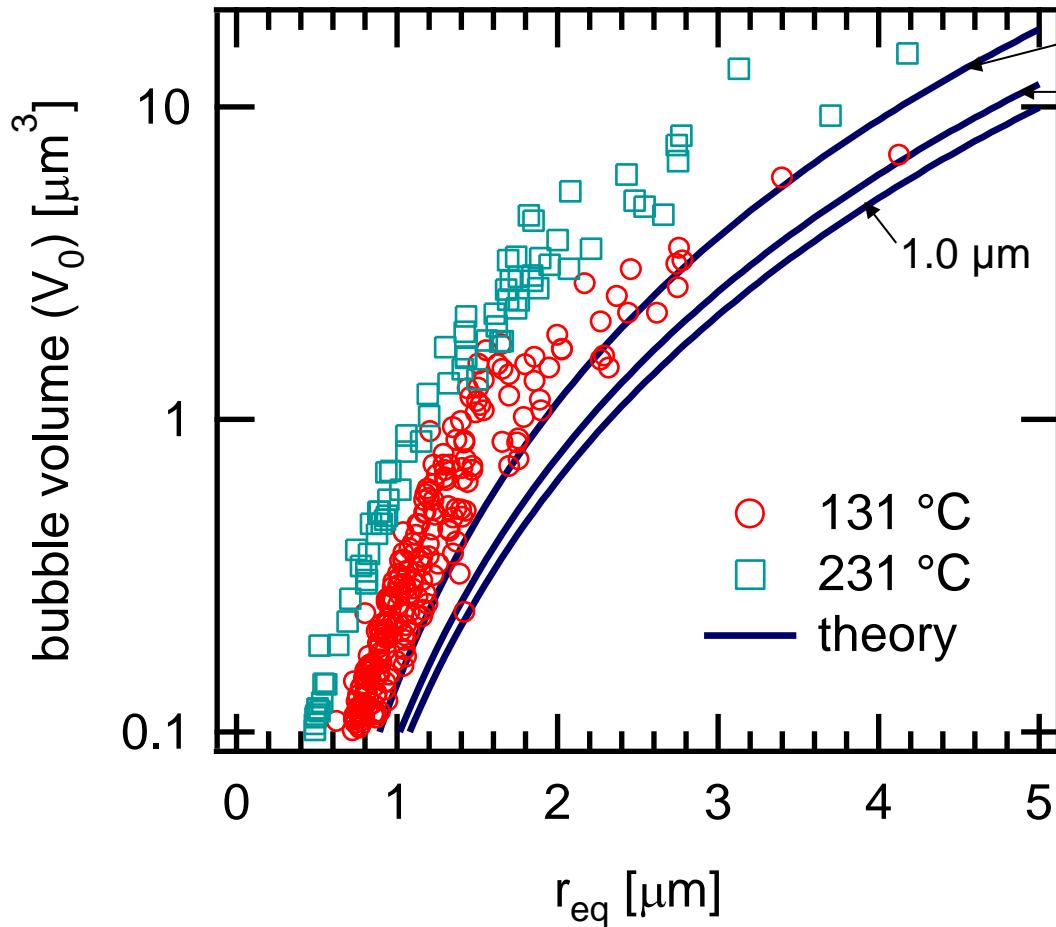


Volume modeled using blister test for thin film adhesion:

$$V = \int y(r) 2\pi r dr = C_1 \pi a^2 y_c$$

K. Wan & Y. Mai, *Acta metall. mater.* **43** (1995) 4109.

Bubble volumes measured with AFM correlate well with deflection model



Volume modeled using blister test for thin film adhesion:

$$V = \int y(r) 2\pi r dr = C_1 \pi a^2 y_c$$

K. Wan & Y. Mai, *Acta metall. mater.* **43** (1995) 4109.

Diffusion and trapping modeled with a continuum-scale approach

Diffusion: 1-D, uniform temperature:

$$\begin{aligned}\partial u(x, t) / \partial t \\ = D(t) \partial^2 u(x, t) / \partial x^2 - q_T(x, t) - q_B(x, t)\end{aligned}$$

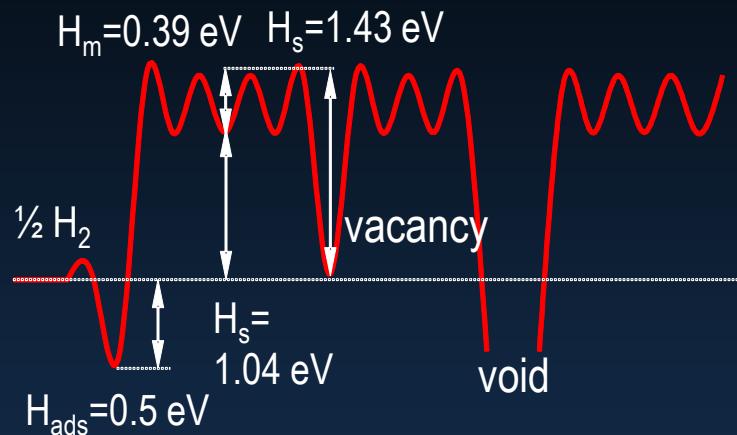
Point defects:

- 1.4 eV saturable traps, no nucleation.
- Used approach of Ogorodnikova [J. Nucl. Mater. (2009)] to address trapping and release.

Bubbles:

Modeled using a approach of Mills [J. Appl. Phys. (1959)].

$$\begin{aligned}q_B(x, t) &= \partial u_B(x, t) / \partial t \\ &= 4\pi D(t) r_B(x, t) N_B(x) [u(x, t) - u_{eq}(x, t)]\end{aligned}$$



Enthalpies for H migrating through W.

Dissolution of H in W is highly endothermic.

H equation of state takes into account non-ideal gas effects

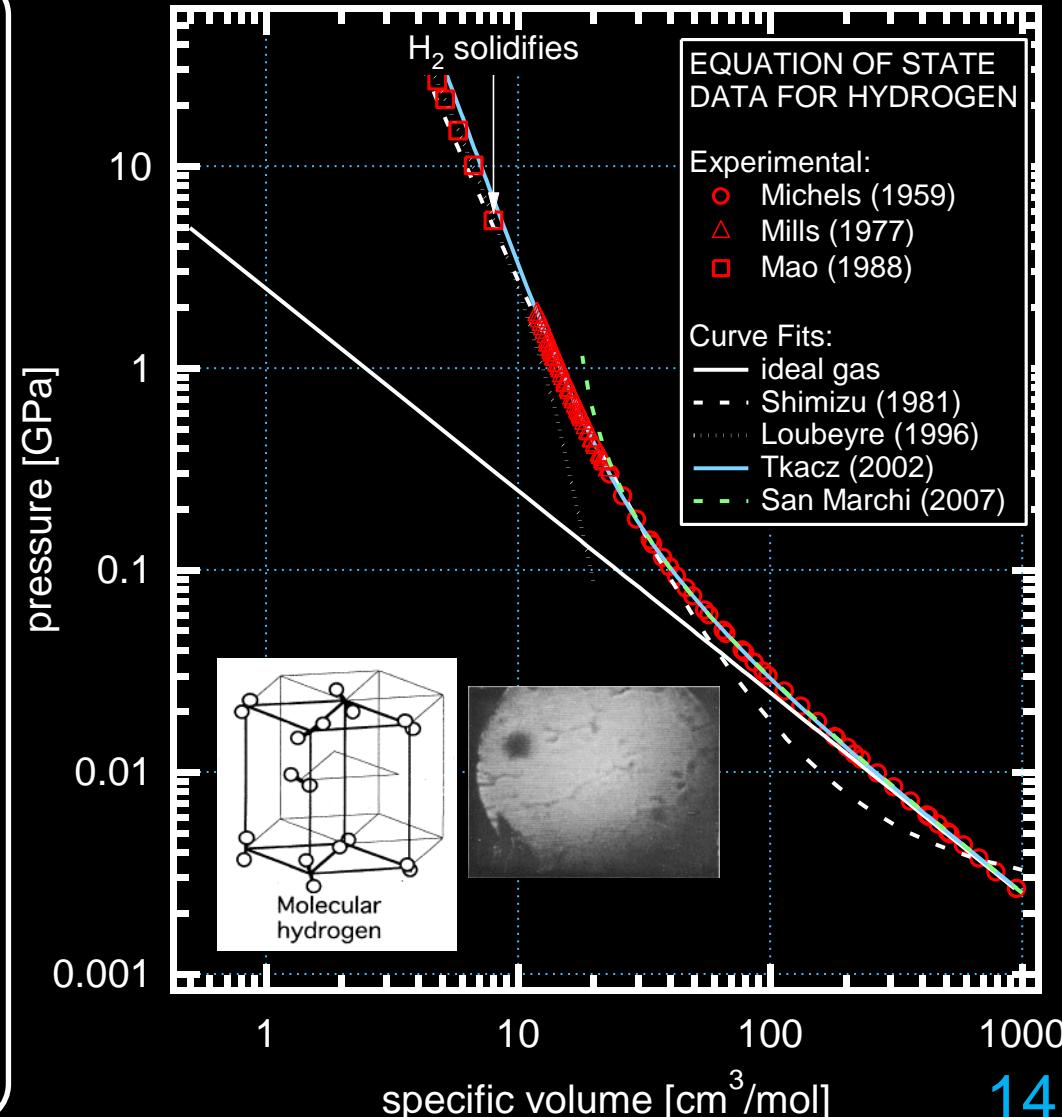
H_2 equation of state (EOS):

- $P > 1$ GPa expected within small bubbles.
- At 300 K, H_2 solidifies at $p=5.7$ GPa.
- Tkacz's [J. Alloys & Compounds (2002)] EOS to provide the best fit:

$$v = Ap^{-1/3} + Bp^{-2/3} + Cp^{-4/3} + (D + ET)p^{-1}$$

- San Marchi's simplified EOS better at low pressure:

$$v = \frac{RT}{p} + b$$



When is bubble growth favorable?

Calculation of equilibrium press.

When is precipitate in equilibrium with mobile conc.?

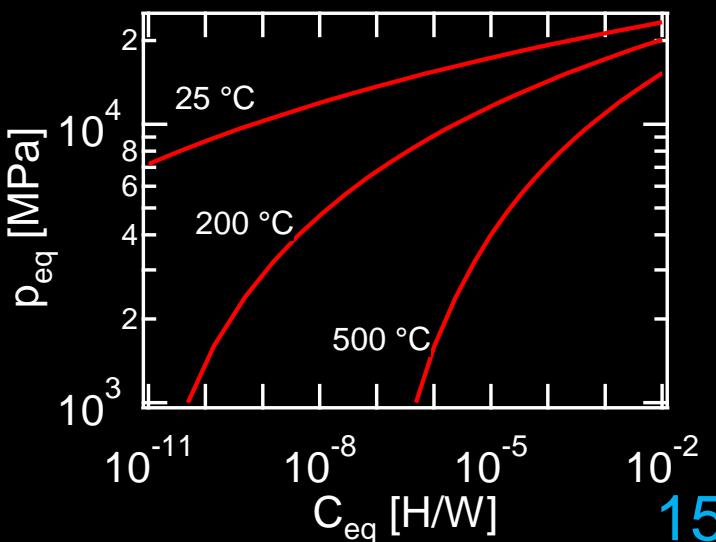
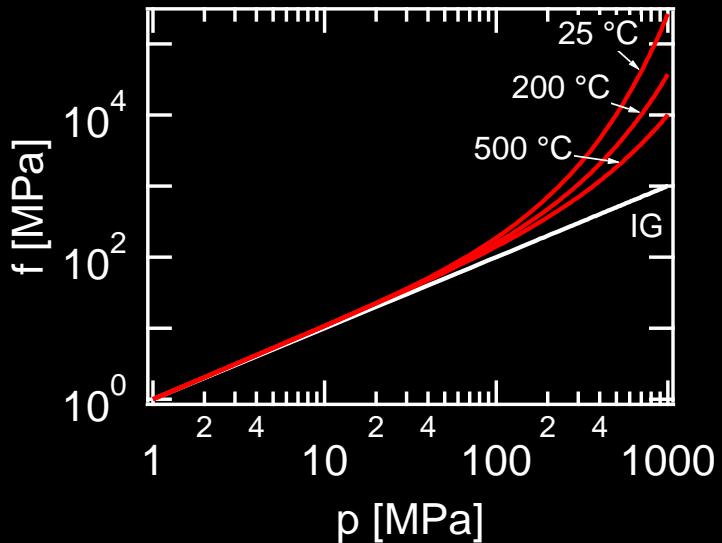
- Equate chemical potentials of gas and solution phase.
- Calculate fugacity to account for non-ideal behavior:

$$\ln(f/p) = \int_0^p \left(\frac{v(p, T)}{RT} - \frac{1}{p} \right) dp$$

- Equilibrium conc. given by:

$$u_{eq} = \sqrt{f} S_0 \exp(-H_S/RT)$$

S_0 and H_S from Frauenfelder [JVST, 1969].



Summary of surface morphology findings

- ITER-grade W sample exposed in TPE show similar retention to Toyama/IPP studies.
- Analysis of surface morphology:
 - XPS shows implanted C reduced considerably
 - SEM/EBSD illustrate non-uniform bubble growth over surface
 - Bubble grow on (110) and (111) crystal planes
 - AFM analysis provide bubble volumes
- Modeling of bubbles:
 - Thin film adhesion model adapted to model blister grown on tungsten.
 - Model reproduces bubble sizes observed with AFM

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