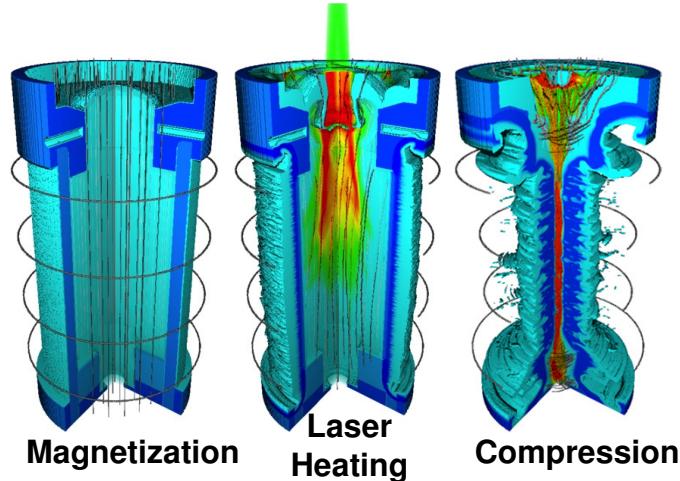
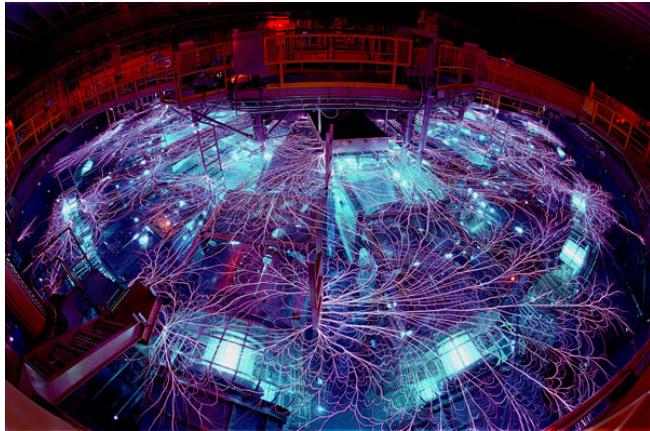


Exceptional service in the national interest



Analysis of Magnetized Liner Inertial Fusion (MagLIF) experiments on Z

Stephanie Hansen
for the MagLIF team

Sandia National Laboratories

Radiative Properties of Hot Dense Matter
Vienna, Austria
Sept 29 - Oct 3, 2014

Summary: Magnetized Liner Inertial Fusion (MagLIF) appears to be a promising ICF platform

- Magnetized Liner Inertial Fusion (MagLIF) relaxes the pressure requirements for ignition by imposing an external magnetic field that reduces heat losses and confines charged particles
- Initial integrated experiments have produced $>10^{12}$ DD neutron yields and large secondary DT yields indicating highly effective magnetic confinement
- Focused experiments are helping to increase our fundamental understanding of stability, heating, compression, and confinement
- Spectroscopic analysis of preheat and stagnation plasmas are helping to constrain the simulation tools used for target design
- Future experiments will help us understand initial conditions, platform stability, mix, and scaling

Many people are contributing to our Magnetized Liner Inertial Fusion (MagLIF) effort:

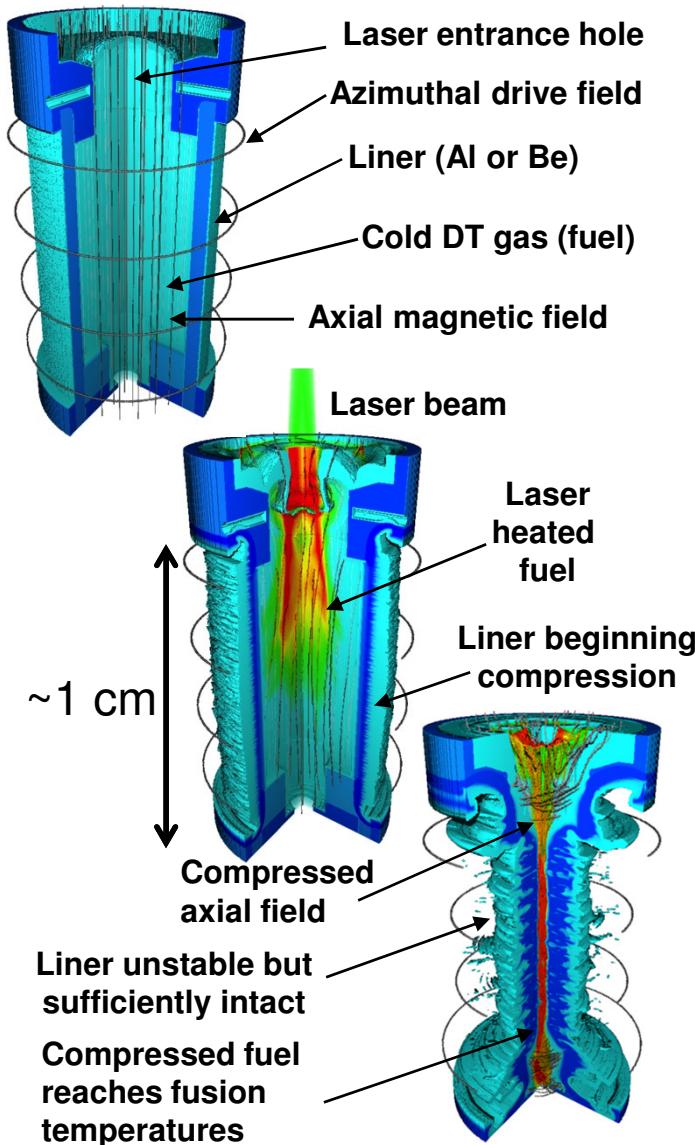
T.J. Awe, C.J. Bourdon, G.A. Chandler, P.J. Christenson, M.E. Cuneo, M. Geissel, **M.R. Gomez**, K.D. Hahn, S.B. Hansen, E.C. Harding, A.J. Harvey-Thompson, M.C. Herrmann, M.H. Hess, C.A. Jennings, B. Jones, M. Jones, R.J. Kaye, P.F. Knapp, D.C. Lamppa, M.R. Lopez, M.R. Martin, R.D. McBride, L.A. McPherson, J.S. Lash, K.J. Peterson, J.L. Porter, G.A. Rochau, D.C. Rovang, C.L. Ruiz, S.E. Rosenthal, M.E. Savage, P.F. Schmit, **A.B. Sefkow**, **D.B. Sinars**, **S.A. Slutz**, I.C. Smith, W.A. Stygar, R.A. Vesey, E.P. Yu

Sandia National Laboratories, Albuquerque, NM 87185 USA

B.E. Blue, D.G. Schroen, K. Tomlinson

General Atomics, San Diego, CA 92186 USA

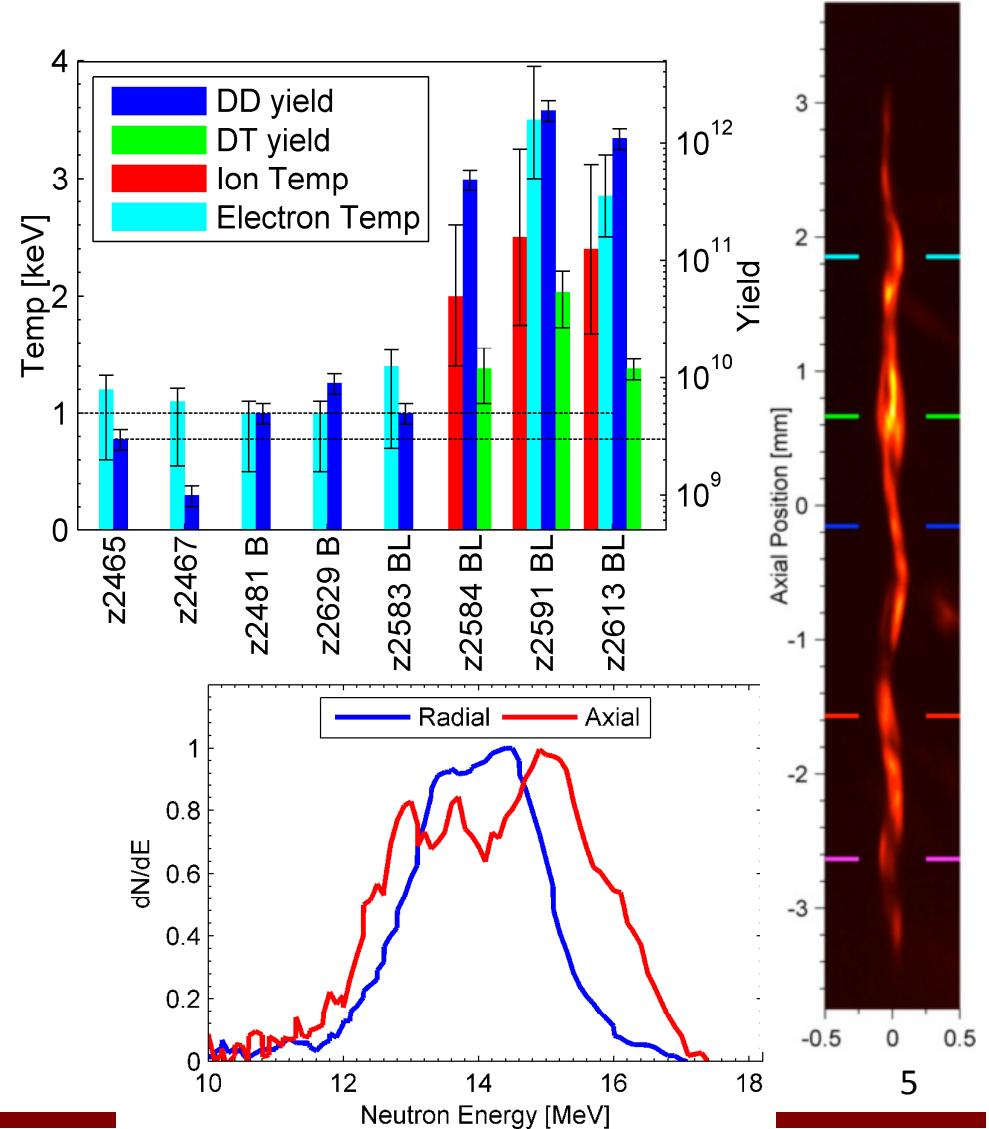
We are evaluating a **Magnetized Liner Inertial Fusion (MagLIF)** concept that may reduce fusion requirements



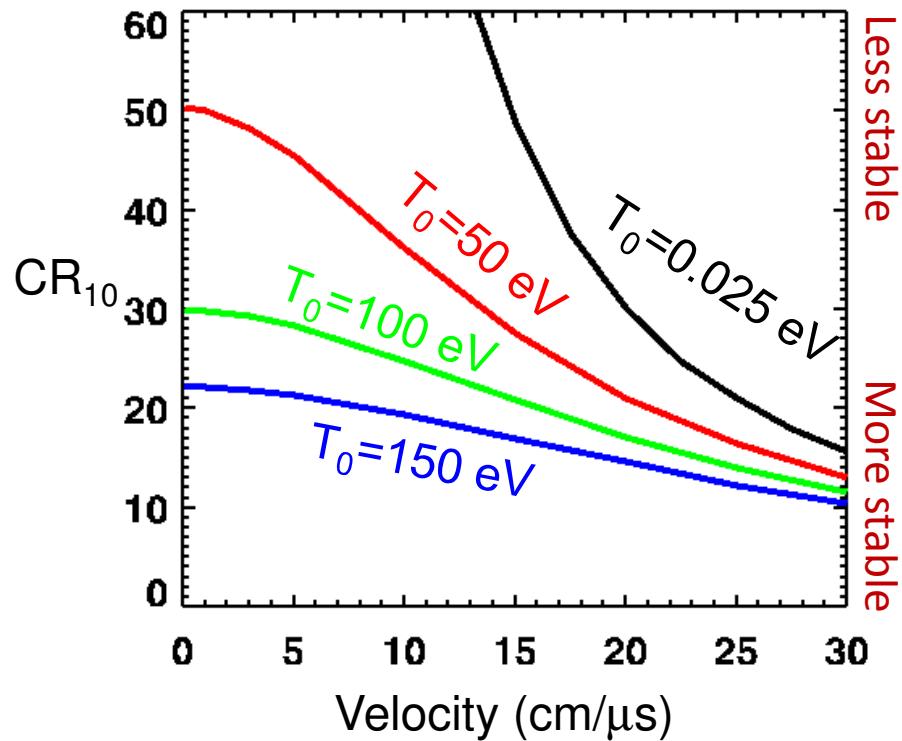
- An applied initial 10-30 T axial magnetic inhibits thermal conduction losses and appears to stabilize implosion at late times
- 30 ns before the implosion begins, the fuel is preheated using 2-6 kJ from the Z-Beamlet laser. Preheating reduces required compression and implosion velocities, increasing stability
- Z's 20 MA current efficiently drives a ~ 100 km/s implosion, delivering $\sim 1\%$ (100 kJ) of the 10 MJ stored in its capacitor banks *to the fuel*
- Each of these components must be present for the concept to work, and they work together to reduce the stagnation pressure requirement to ~ 5 Gbar from the ~ 500 Gbar of traditional ICF
- DD equivalent of 100 kJ DT yield may be possible on Z in the next few years, requiring upgrades that are in progress:
 $10\text{ T} \rightarrow 30\text{ T}$; $2\text{ kJ} \rightarrow >6\text{ kJ}$; $19\text{ MA} \rightarrow >24\text{ MA}$

We obtained promising initial results with MagLIF and are seeking to increase our understanding

- We achieved DD yields up to 2×10^{12} (~ 0.3 kJ DT equivalent) in our first integrated tests of MagLIF
- Extensive neutron and x-ray data indicate a $< 150 \mu\text{m}$ diameter, ~ 3 keV, $\sim 0.4 \text{ g/cm}^3$, highly magnetized plasma
- We are continuing to build on these results with a balanced combination of focused and integrated experiments
- In parallel, we are improving capabilities to understand how this performance will scale with increasing drive parameters



Typical ICF implosions need high velocities to reach fusion temperatures; laser preheating can reduce velocity requirements



CR_{10} = Convergence Ratio (R_0/R_f) needed to obtain 10 keV (ignition) with no radiation losses or conductivity

Heating fuel to ignition temperatures is typically done with a high-velocity shock (or series of shocks)

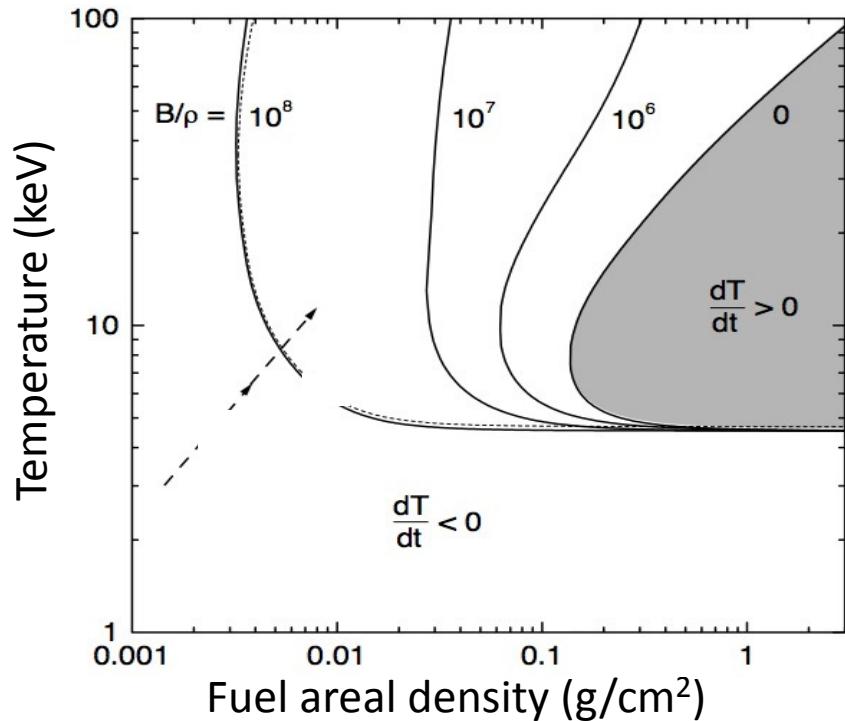
High velocities make it easier to reach fusion temperatures and also reduce the time available for losses (e.g., electron heat conduction or radiation) -- **but can lead to uncontrolled instability growth**

Heating the fuel prior to the implosion *in the absence of losses* can allow stable, low-velocity, low-convergence implosions to reach ignition temperatures

Is there a way to reduce losses?

A large, embedded magnetic field can significantly reduce electron conduction losses from heated fuel

*Basko et al. *Nuc. Fusion* 40, 59 (2000)



An external magnetic field can significantly reduce the ρr needed for ignition by inhibiting electron conduction losses

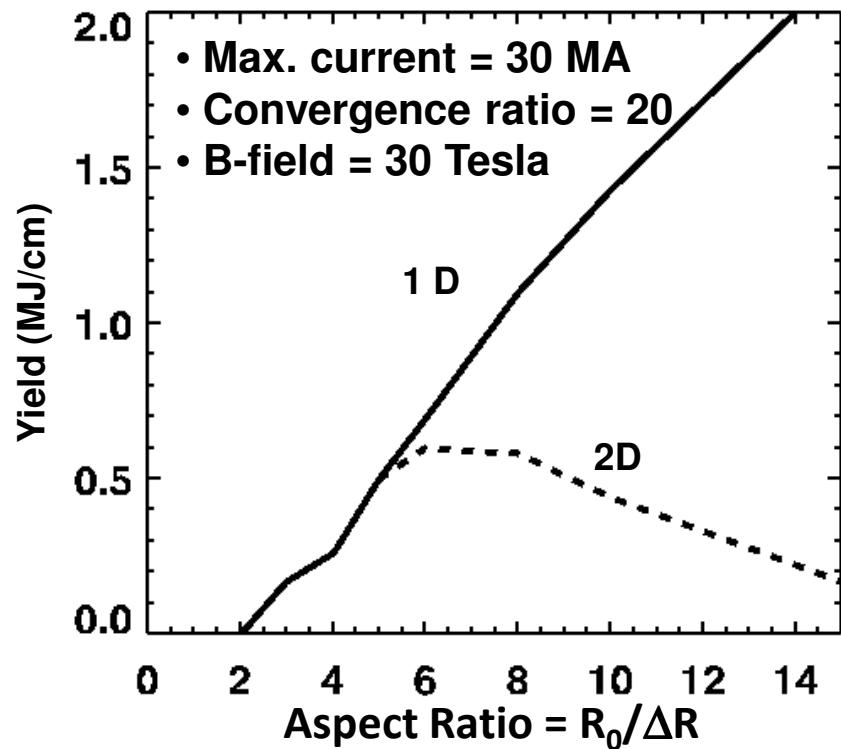
Lower ρr reduces the required final fuel density (e.g., $\sim 1 \text{ g/cc} \ll 100 \text{ g/cc}$), which also reduces bremsstrahlung radiation losses

This means the stagnation plasma pressure at ignition temperatures is significantly reduced (e.g., $\sim 5 \text{ Gbar} \ll \sim 500 \text{ Gbar}$)

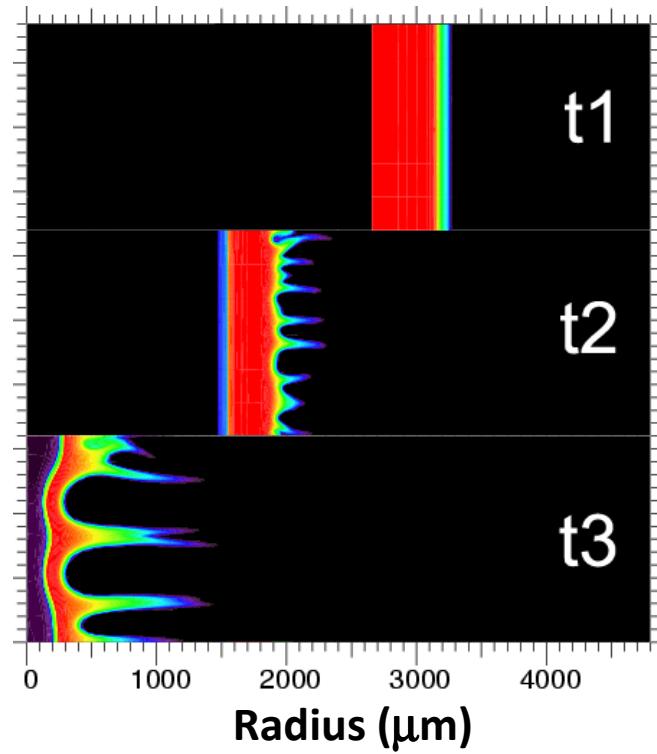
Large values of B/ρ are needed and therefore large values of B are needed:
 $B \sim 10,000 \text{ Tesla} (> 10^8 \times \text{Earth's B-field})$

This field significantly exceeds pulsed coil technology ($B_0 \sim 10\text{-}30 \text{ T}$), therefore flux compression is needed... and we have an imploding liner to do it.

Reducing the implosion velocity requirements through fuel heating and magnetization allows us to use thicker liners for more stable implosions



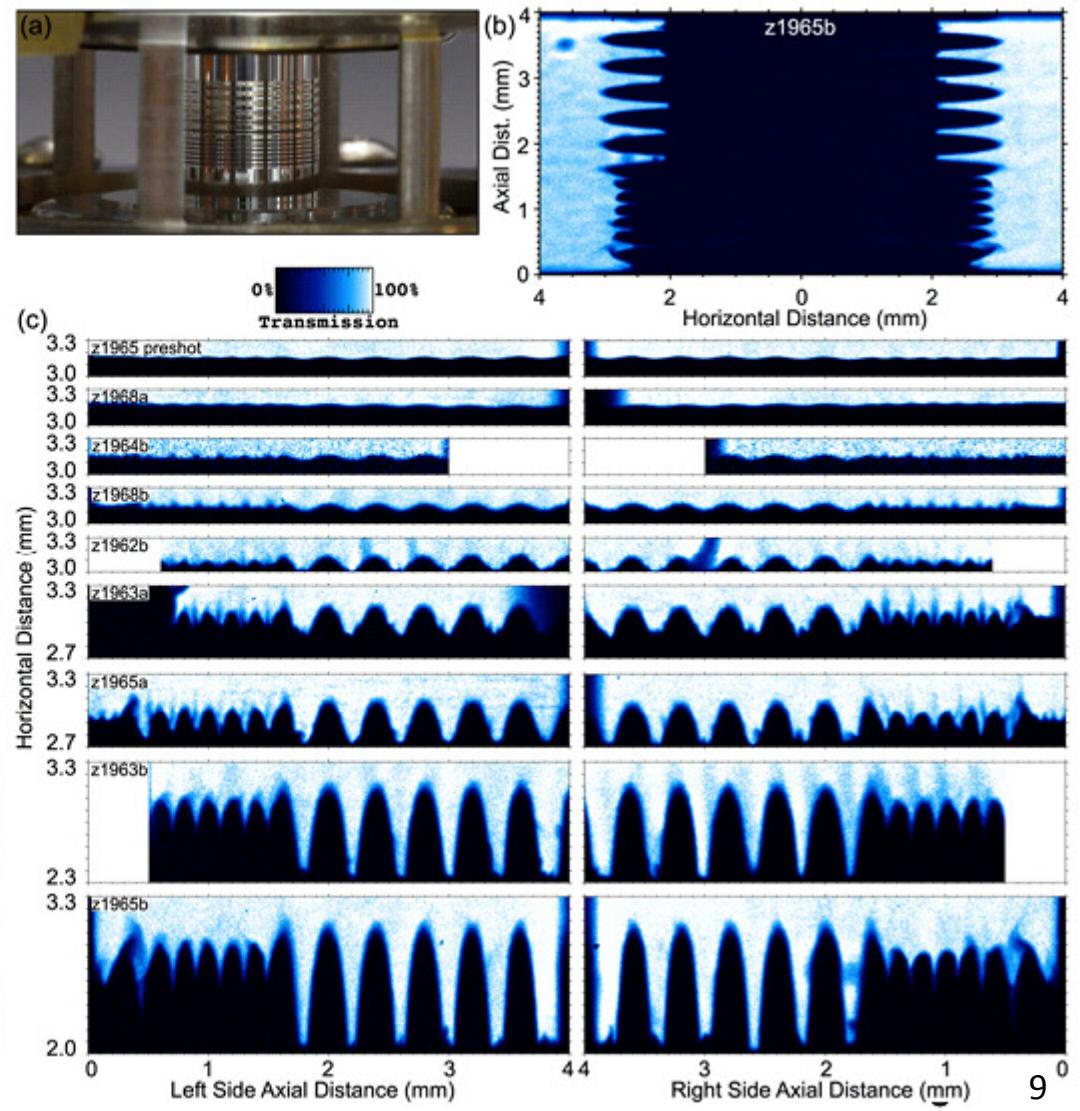
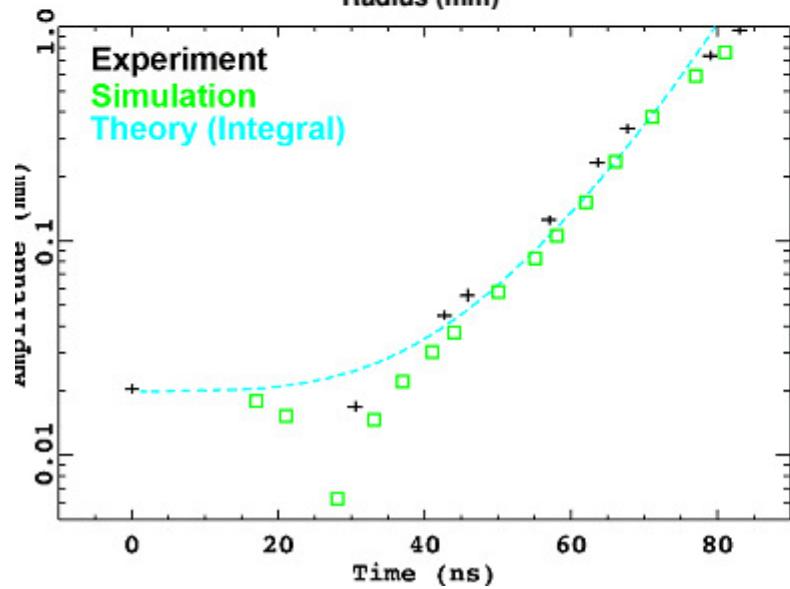
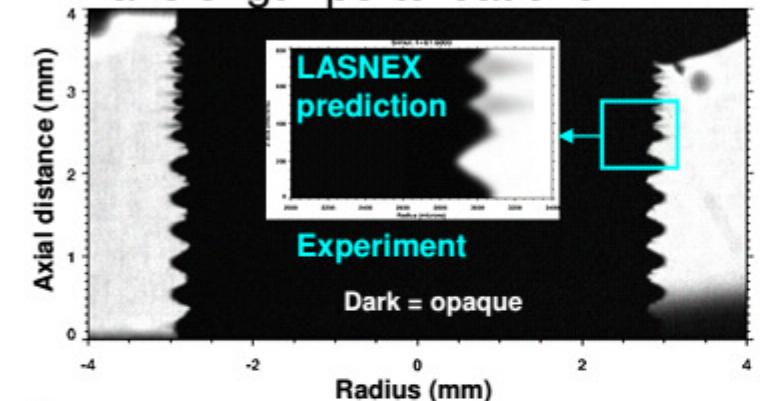
In 2D simulations, the Magneto-Rayleigh-Taylor instability degrades the yield as the aspect ratio is increased, due to decreased liner ρR .



2D simulations of AR=6 Be liner show reasonably uniform fuel compression and sufficient liner ρR at stagnation to inertially confine the fuel

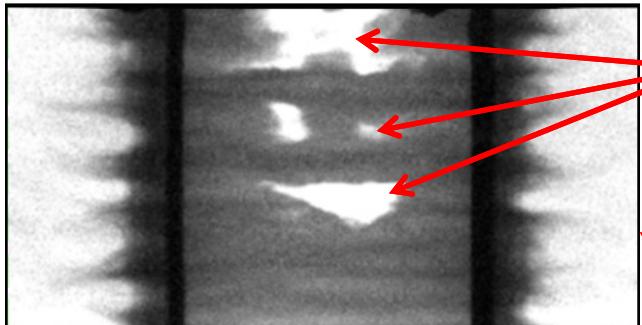
The Magneto-Rayleigh Taylor instability has been extensively studied on Z

Radiographs captured growth of intentionally-seeded 200, 400- μ m wavelength perturbations

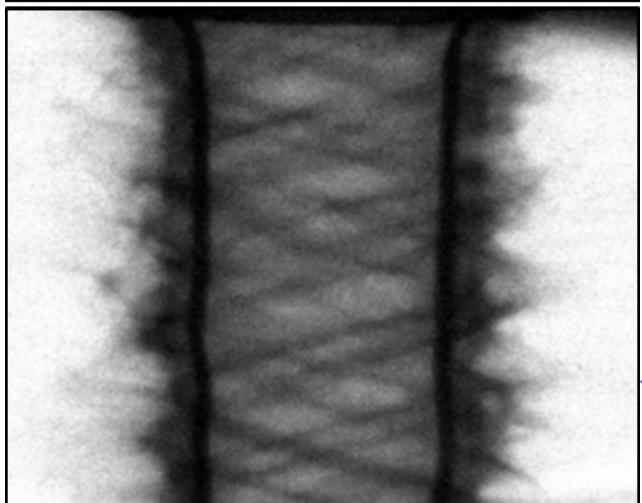


Adding an axial magnetic field appears to enhance liner stability, changing its structure from cylindrical to helical

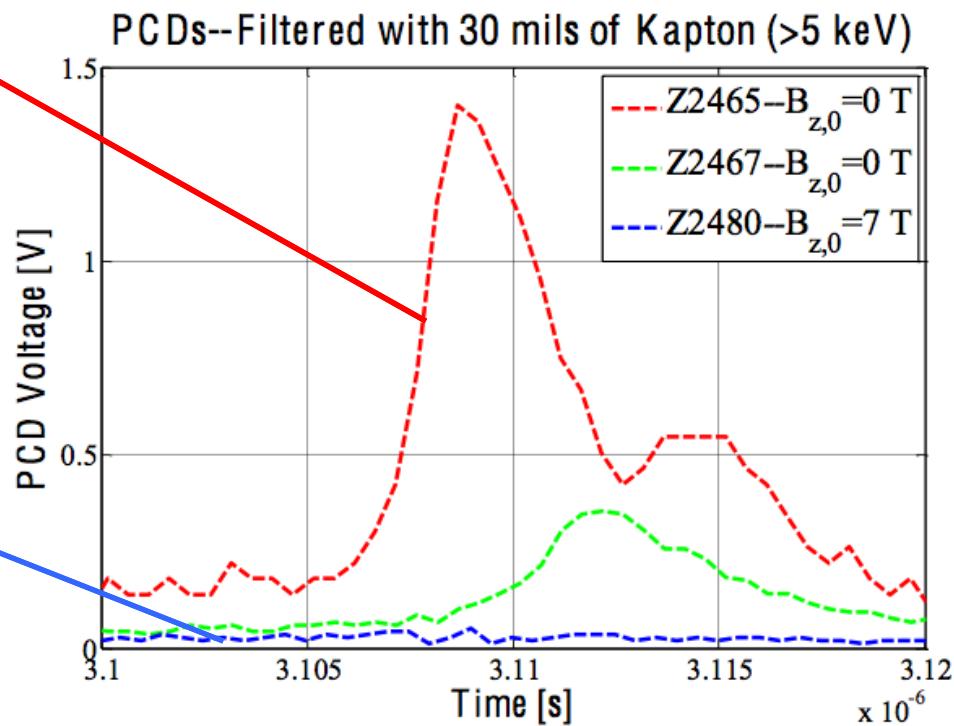
Without Magnetic Field



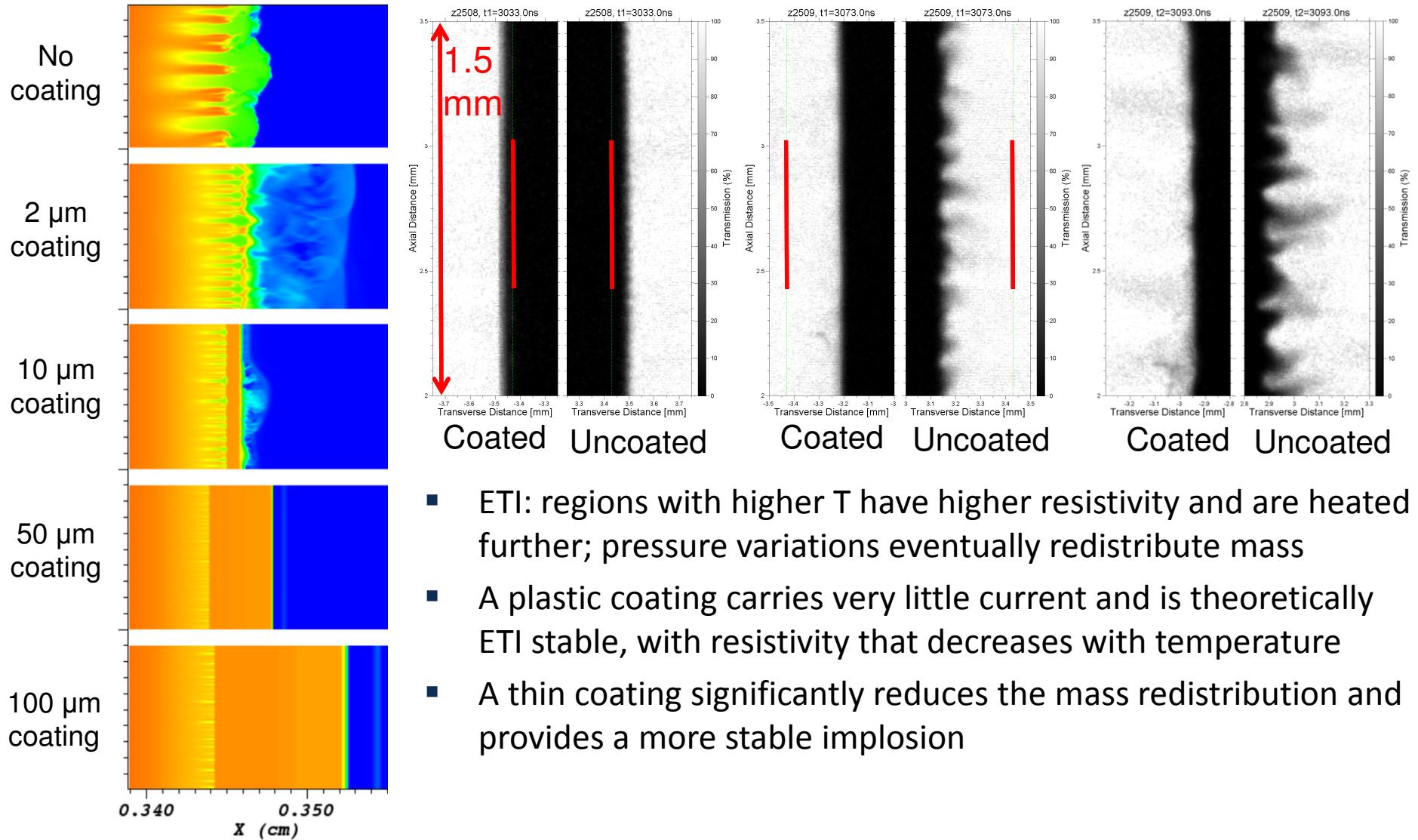
Time-integrated 6 keV self-emission from liner implosion is absent from shots with axial field, indicating suppression of micro-pinching by flux-compressed B_z



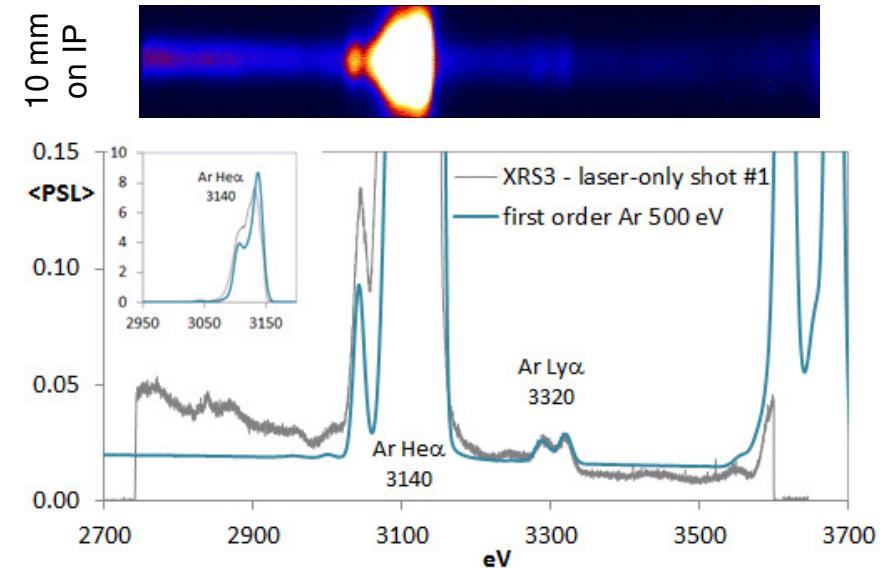
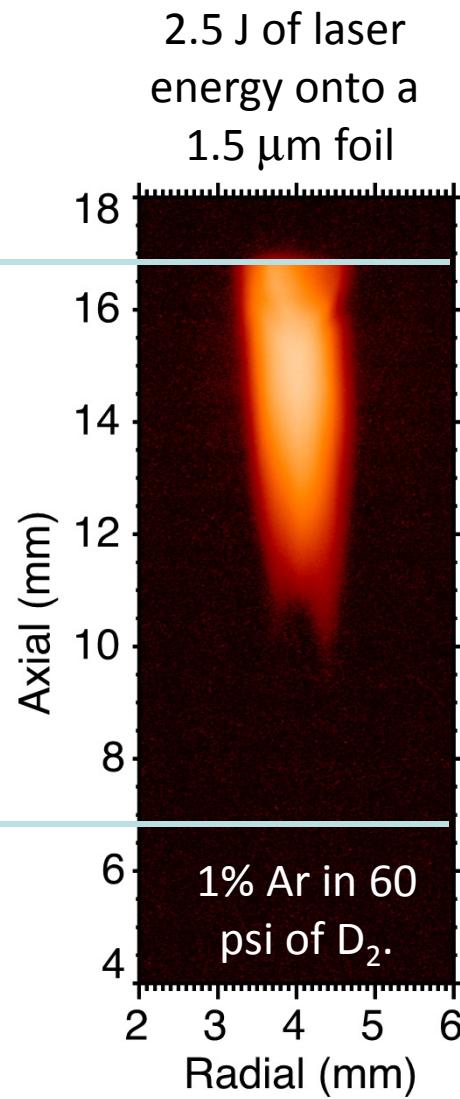
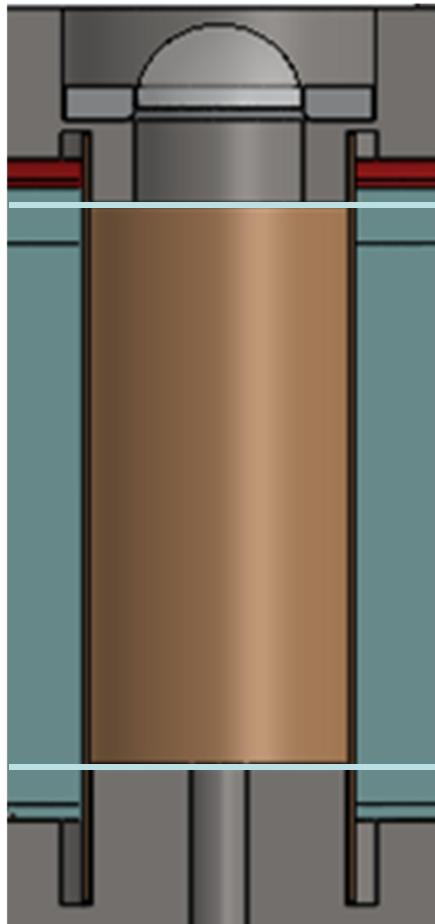
With Magnetic Field



Further enhancement of stability appears possible through mitigation of Electro-thermal instability growth



Focused studies of the laser preheat performance are ongoing at both Z/Beamlet and Omega



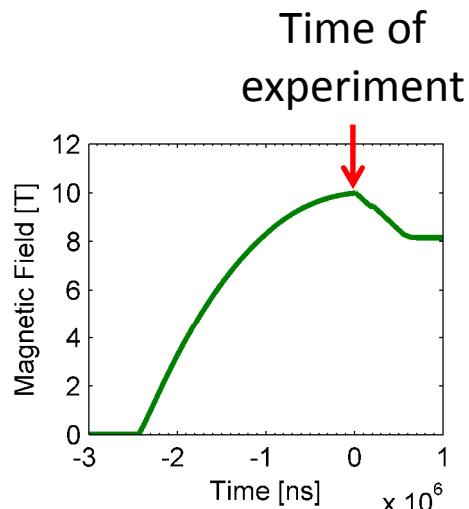
-3.1 keV crystal imager and spectral data* indicate that about half of the 10-mm axial dimension is heated to $\text{Te} \sim 500$ eV

-A laser-only shot with no gas fill and 4.5 kJ laser energy showed $\text{Te} \sim 600$ eV emission from a Cl dopant at the bottom of the target

-X-ray emission from targets at Omega indicated $T > 200$ eV only with $B_z > 0$

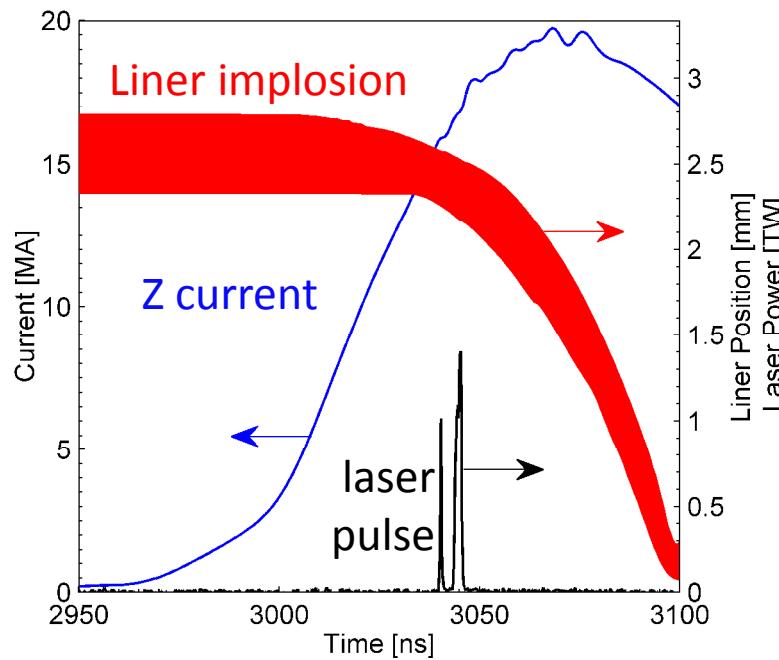
The initial integrated experiments used

$I = 19 \text{ MA}$, $B = 10 \text{ T}$, and $\text{Laser} = 2.5 \text{ kJ}$



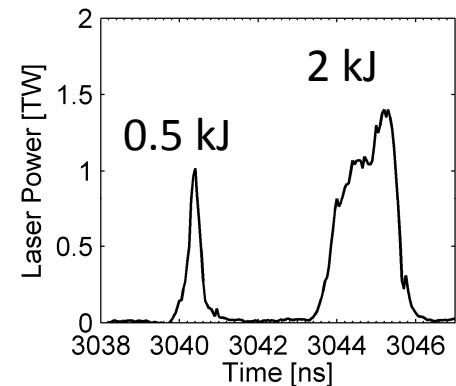
Magnetic field risetime is approximately 2 ms

B is constant over the timescale of the experiment

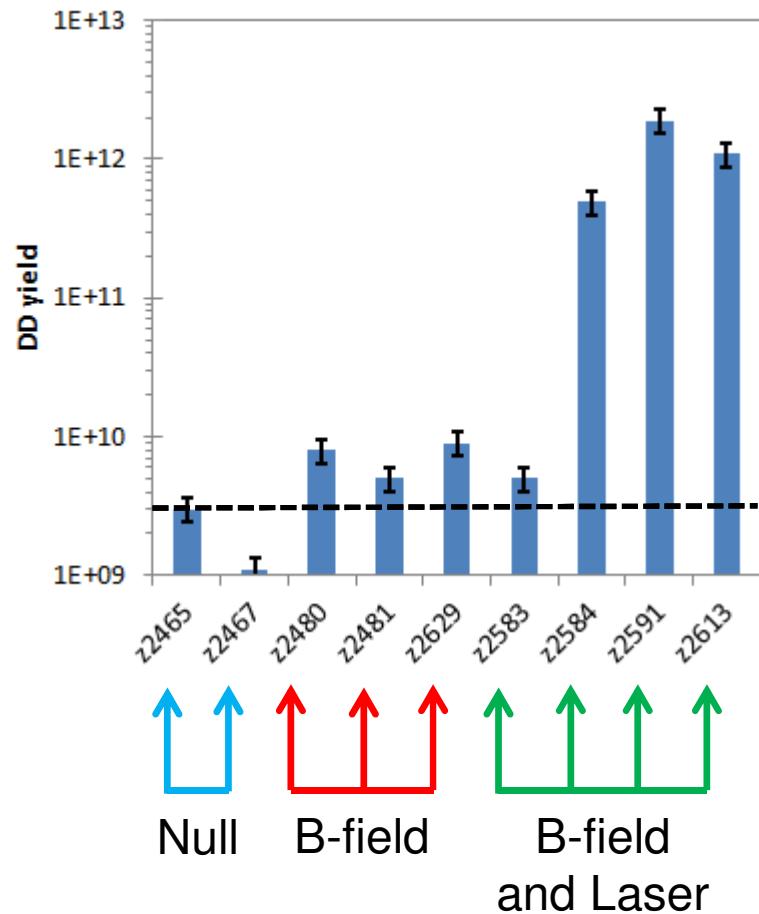


Peak current is 19 MA
Magnetic field is 10 T
Total laser energy is 2.5 kJ

Laser energy is split into 2 pulses:
1st pulse intended to destroy LEH
2nd pulse intended to heat fuel

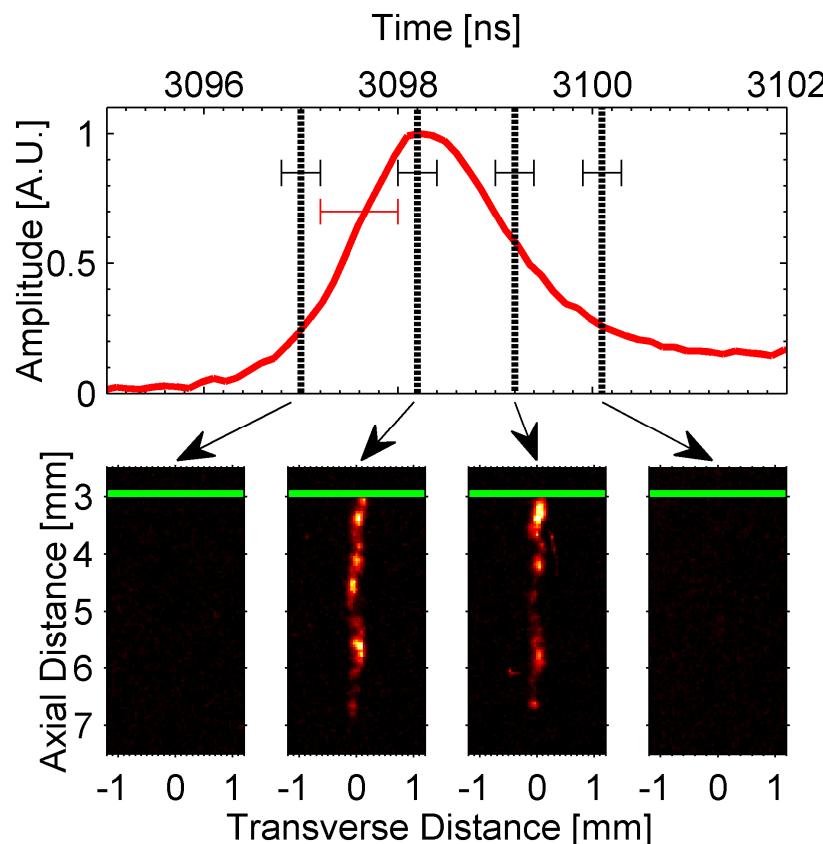


Z shots producing DD yields in excess of 10^{12} were only observed in experiments with laser and B-field



- High yields were only observed on experiments incorporating both applied magnetic field and laser heating
- A series of experiments without laser and/or B-field produced yields at the background level of the measurement
- A high-pressure shot and a shot with high-Z-doped fuel did not produce measurable yield
- Shots in June 2014 with 4 kJ laser energy did not produce measurable yield
- Shots in September 2014 with thinner LEH windows produced few- 10^{11} yields

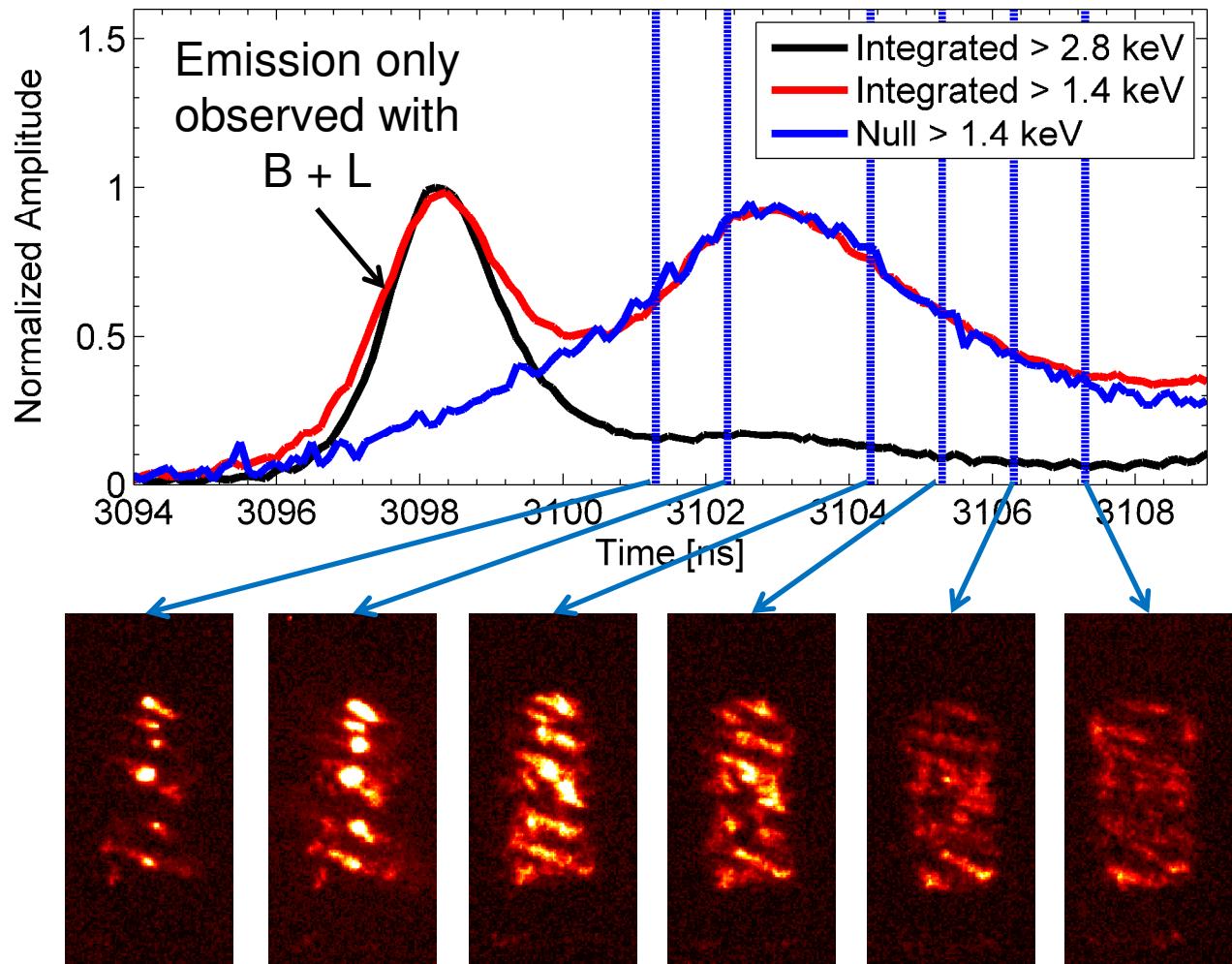
Time-resolved x-ray pinhole imaging ($h\nu > 2.8$ keV) shows a narrow emission column during peak in X-ray signal



- Emission column is observed only during the peak in the x-ray signal, and only on experiments with high neutron yield
- Stagnation column width is at the resolution limit of this instrument (~150 microns)
- Absolute x-ray powers suggest stagnation density of 0.4 ± 0.2 g/cm³ and Be mix < 10%

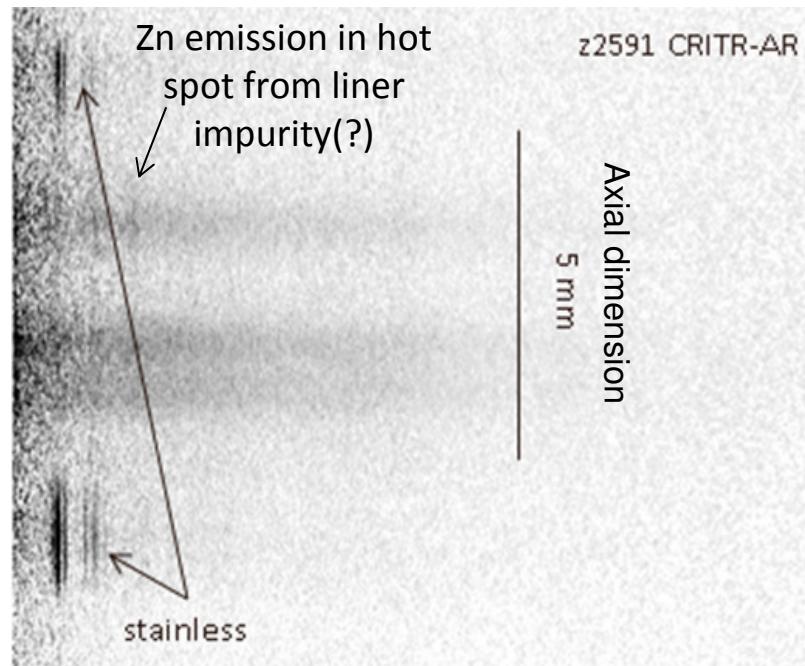
High energy x-ray signal and narrow emission region are absent in null experiments

- Liner emission is observed in all experiments
- Liner emission is at a lower photon energy (< 2.8 keV)
- Liner emission region expands after stagnation & exhibits helical structure

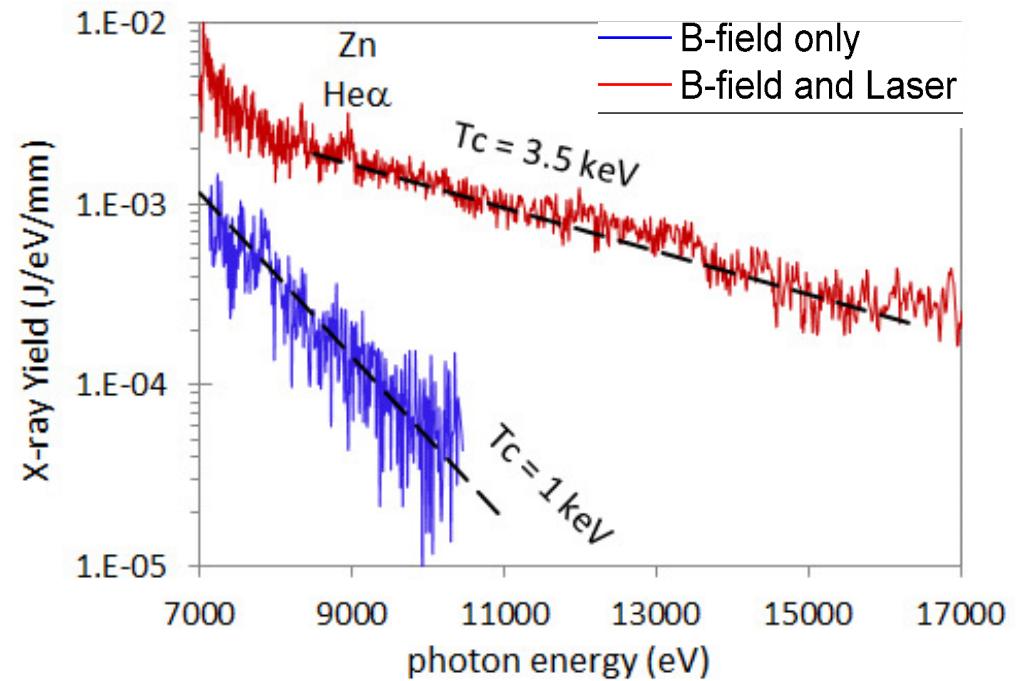


High-energy spectra show axial variations in temperature and composition, with ~ 3.5 keV electron temperature in the pinch region—remarkable for a 70-100 km/s implosion!

Emission from high-Z components of stainless steel is consistent with late-time emission yields

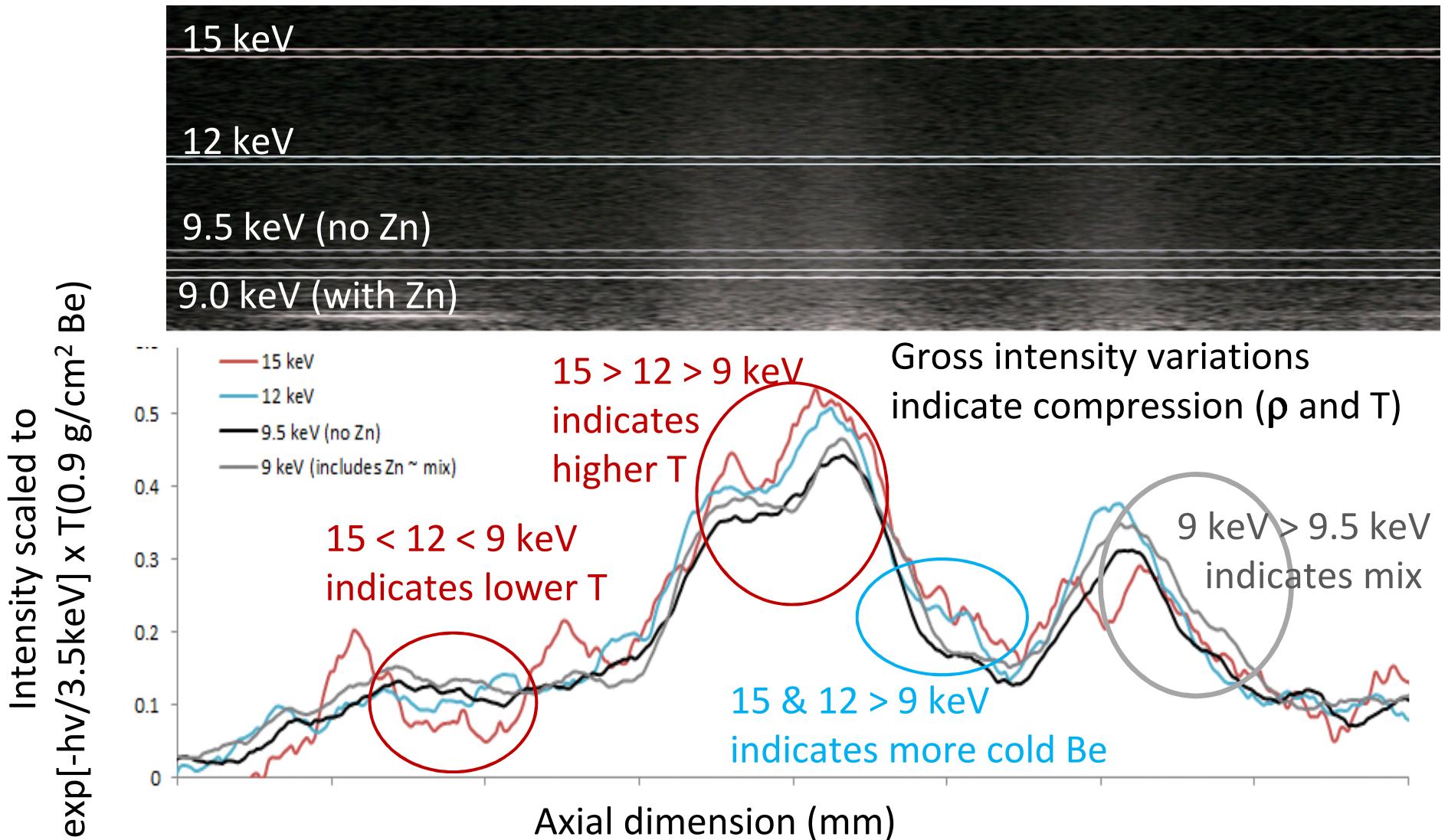


High-energy spectra from hot spots indicate $T_e \sim 3.0$ -3.5 keV from high-yield shots and constrain Be mix to < 10%



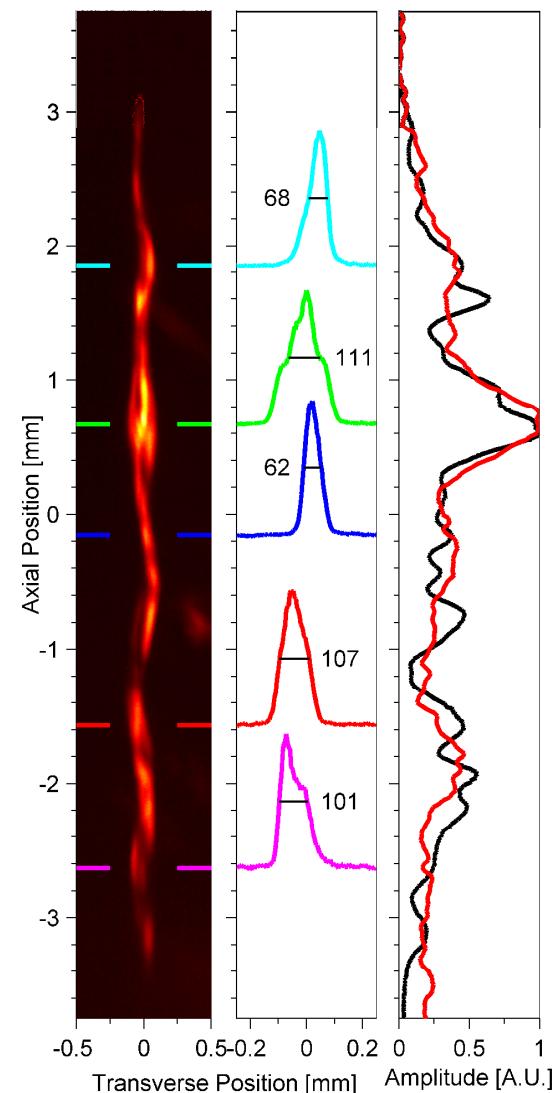
The measured T_e are close to the T_{ion} obtained from neutron time-of-flight data and tend to scale with the measured neutron yields.

Lineouts of the high-energy spectrum at different energies can be used to infer axial variations in liner opacity, compression, temperature, and mix



High-resolution monochromatic imaging of the x-ray emission shows a narrow, hot plasma column with weakly helical structure

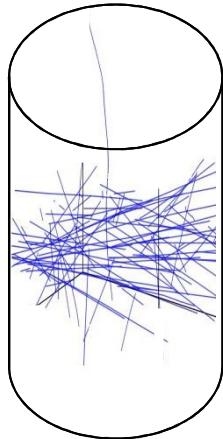
- Lineouts of stagnation column vary from 60 to 120 μm FWHM (resolution about 60 μm)
- Emission is observed from about 6 mm of the 7.5 mm axial extent
- Note that the emission doesn't necessarily define the fuel-liner boundary, but only the hot fuel region
- The stagnation column is weakly helical with a wavelength of about 1.3 mm and a 0.05 mm horizontal offset
- Axial lineouts of image (black) agree with 9.3 keV 1D spectrometer lineouts (red), suggesting features are due to emission and not liner opacity (Be opacity is small >9 keV)
- **With $\rho \sim 0.4 \text{ g/cm}^3 \rightarrow \rho r \sim 2 \text{ mg/cm}^2$**



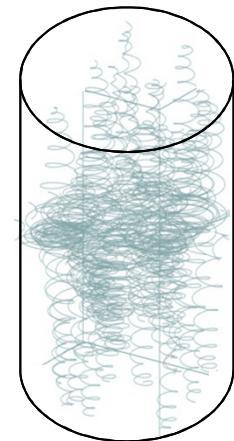
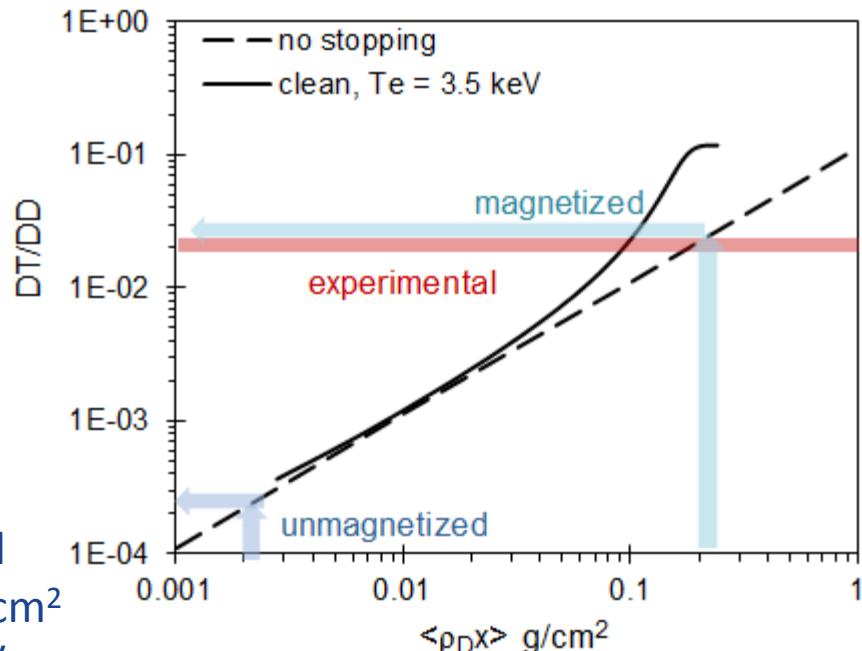
In addition to the significant $\sim 2 \times 10^{12}$ DD neutron yields, we measure a remarkable $\sim 5 \times 10^{10}$ DT neutrons



“Secondary” 14 MeV neutrons are produced by 1 MeV tritons interacting with D fuel:



In an unmagnetized plasma, $\rho R > 200 \text{ mg/cm}^2$ is required for triton/α confinement



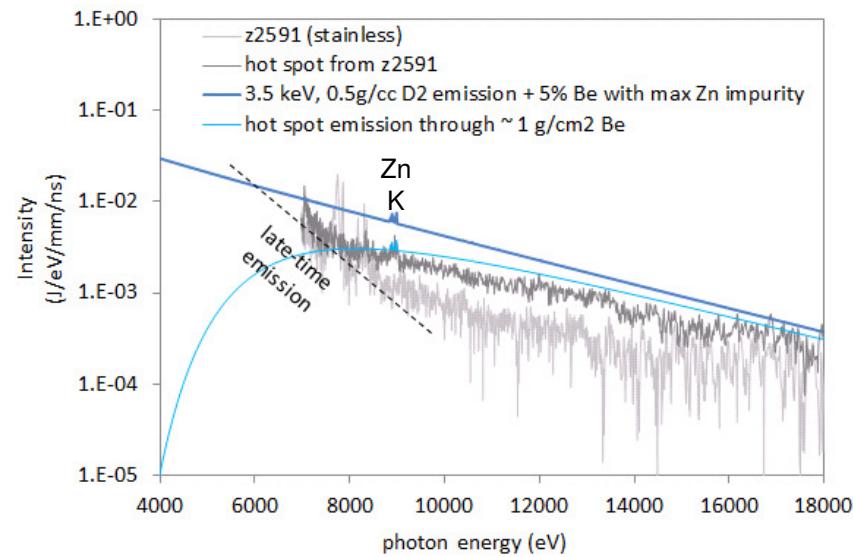
In a magnetized plasma, $\rho R \sim 2 \text{ mg/cm}^2$ is sufficient to confine 1 MeV tritons

This demonstration that B_z confines 1 MeV tritons is critically important for MagLIF – which requires magnetic enhancement of thermal confinement and a-deposition – because a field that confines 1 MeV tritons also confines electrons and 3 MeV alphas.

Summary: Magnetized Liner Inertial Fusion (MagLIF) appears to be a promising ICF platform

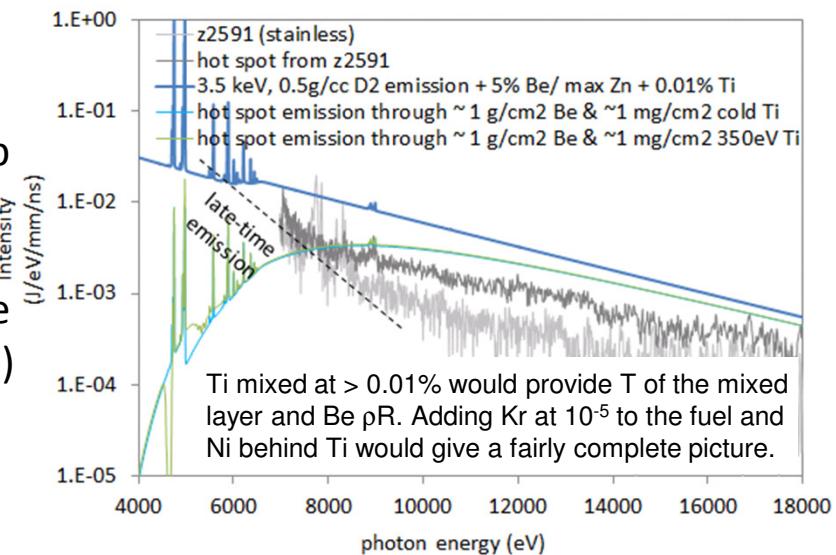
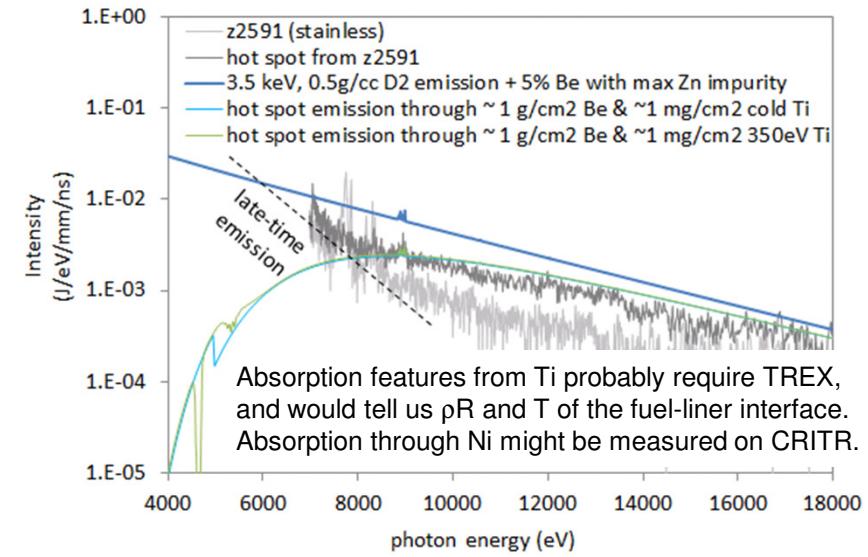
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- Spectroscopic analysis of preheat and stagnation plasmas are helping to constrain the simulation tools used for target design
- Future experiments will help us understand initial conditions, platform stability, mix, and scaling

While fuel dopants lead to radiative losses during preheat that can add stagnation, interior coatings that mix with fuel at stagnation would be less detrimental.

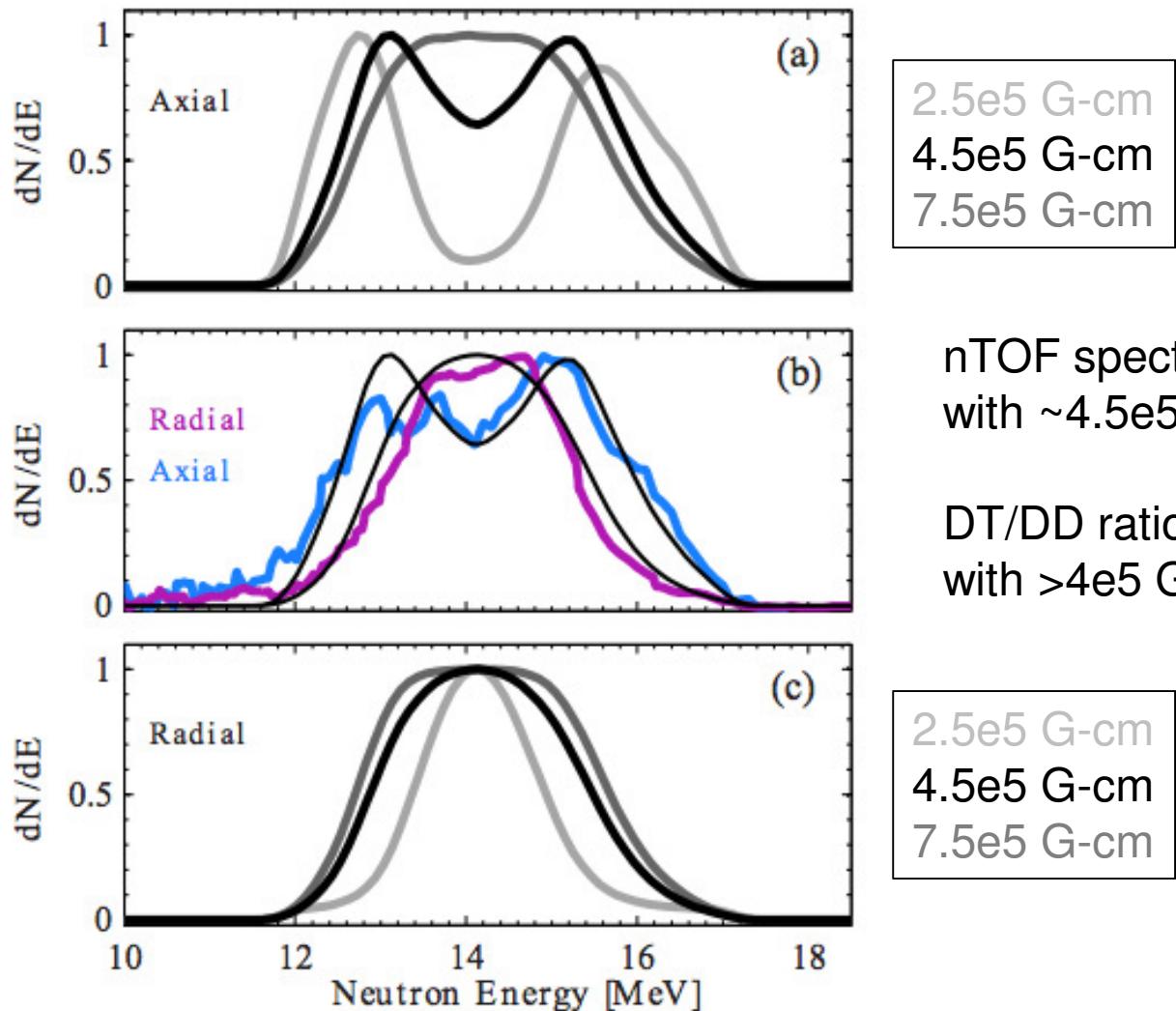


Uncoated Be (< 0.005% Zn)
Measured Zn K-shell suggests a hot core surrounded by a slightly cooler layer mixed layer with ~10% Be (~doubles rad loss)

300 nm Ti would absorb emission from the hot core and may mix with cooler layer and/or core (0.03% doubles rad loss)



Neutron time-of-flight data are consistent with high magnetization

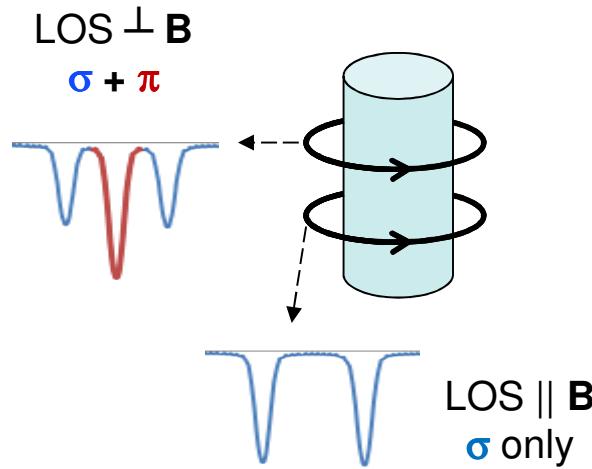


nTOF spectra consistent with $\sim 4.5e5$ G-cm

DT/DD ratio consistent with $>4e5$ G-cm

Zeeman splitting is being used to characterize Z's current drive and flux compression in Magnetized Liner Inertial Fusion (MagLIF) experiments

- Sodium deposits vaporized and backlit by current-carrying surfaces signal both the magnitude and direction of the local magnetic field:



The relative strength of σ and π components indicates field direction

