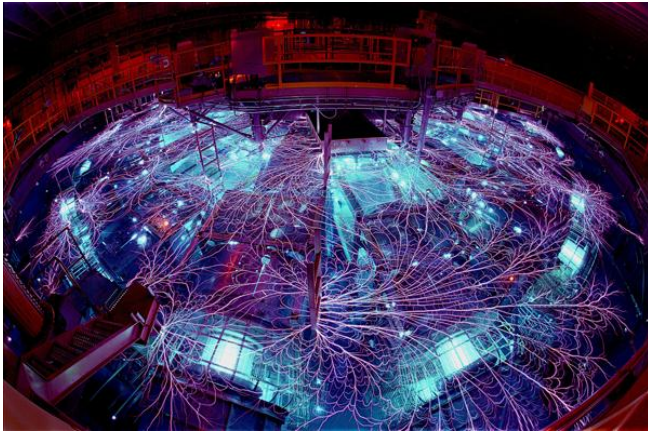


Exceptional service in the national interest

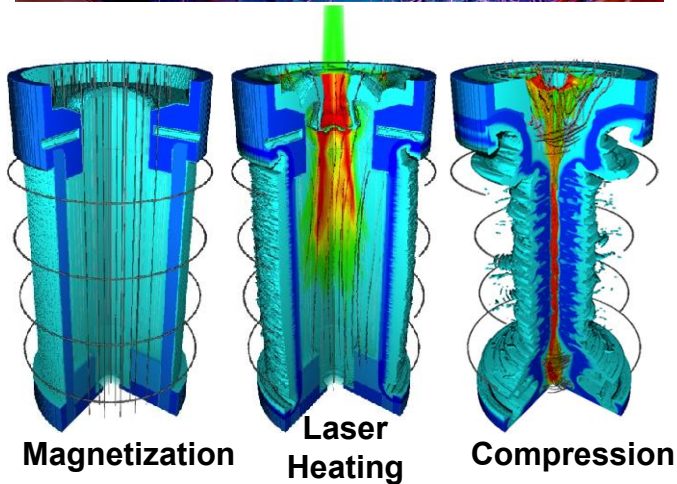


An overview of fusion, astrophysical experiments, and diagnostic development on Sandia's Z machine

Stephanie B. Hansen

Sandia National Laboratories

*Presentation to the Laboratory for Plasma Studies
August 27, 2014*



Outline

- **Magnetized Liner Inertial Fusion (MagLIF)**
Encouraging results from initial integrated experiments
- **Z Astrophysical Plasma Properties (ZAPP) Collaboration**
Studies of emission and absorption from photoionized plasmas relevant to white dwarfs, accretion disks, and solar photosphere
- **Diagnostic development on Z**
X-ray Thomson scattering and Zeeman splitting

There is a lot of interesting science going on at Z –
and many opportunities for significant contributions from Cornell's LPS.

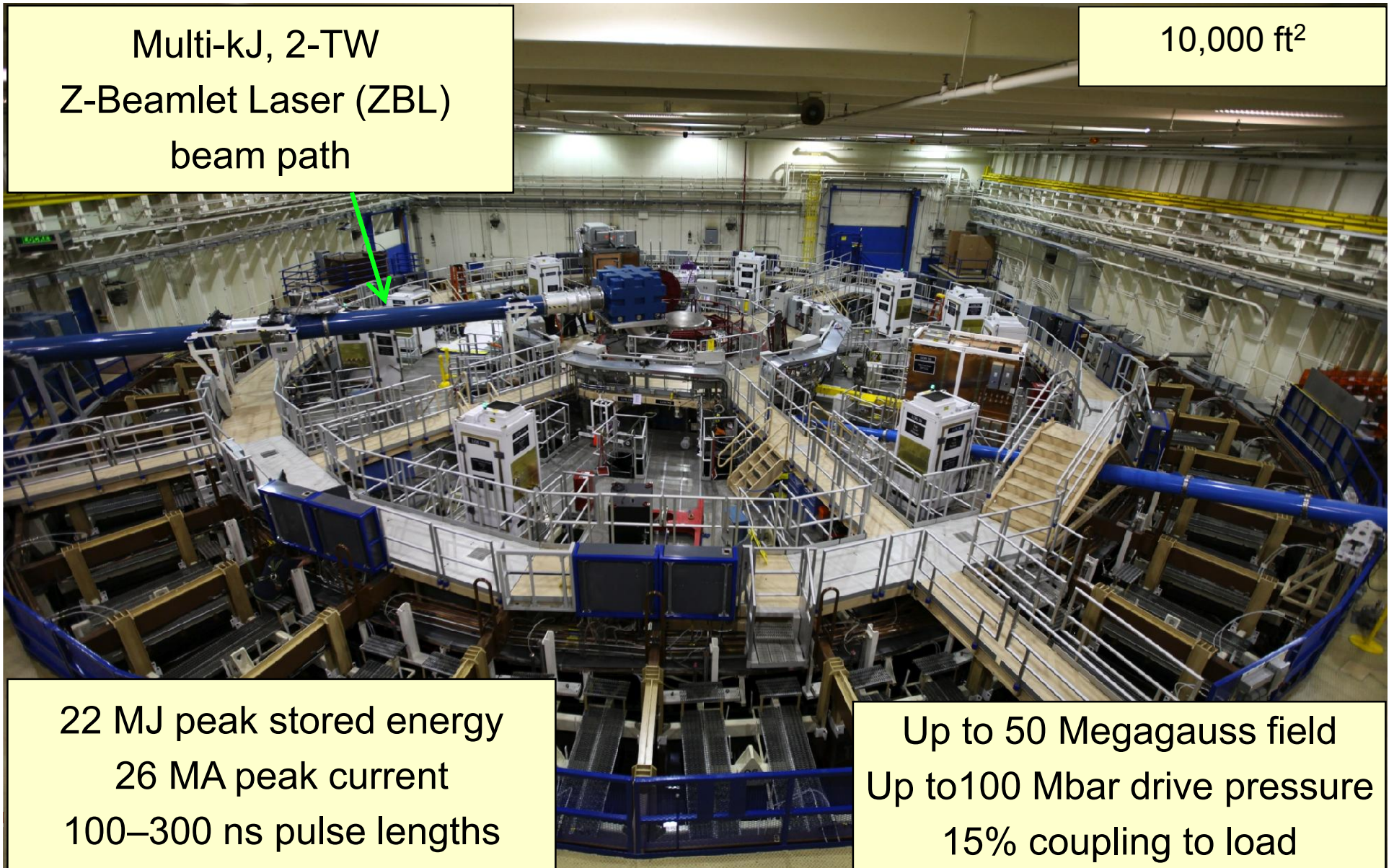
The Sandia Z pulsed power facility uses magnetic pressure to efficiently couple energy to drive relatively large targets for a wide variety of stockpile stewardship applications



Multi-kJ, 2-TW

Z-Beamlet Laser (ZBL)
beam path

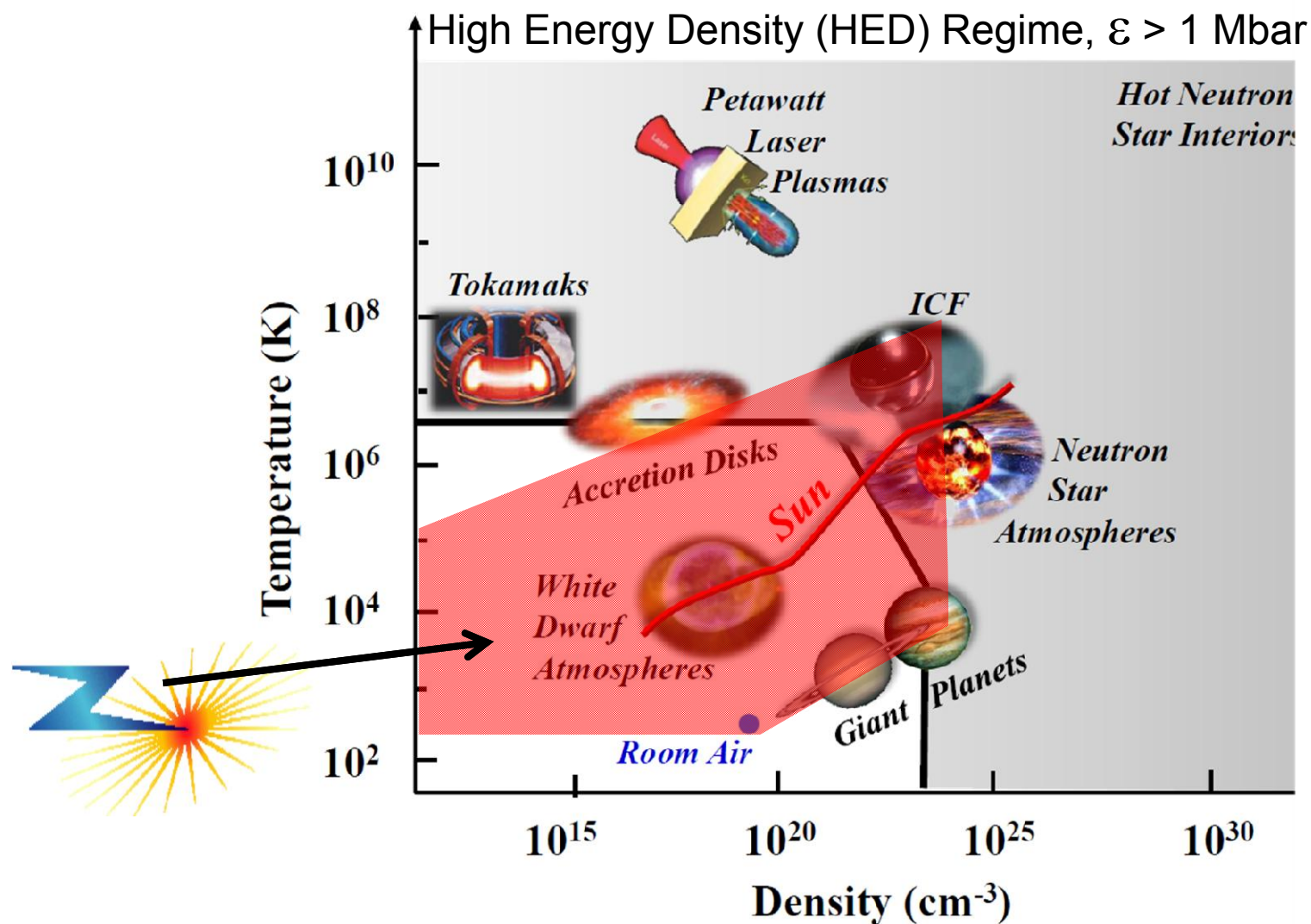
10,000 ft²



22 MJ peak stored energy
26 MA peak current
100–300 ns pulse lengths

Up to 50 Megagauss field
Up to 100 Mbar drive pressure
15% coupling to load

Experiments on Z access a broad range of the energy-density phase space



Many people are contributing to our Magnetized Liner Inertial Fusion (MagLIF) effort:



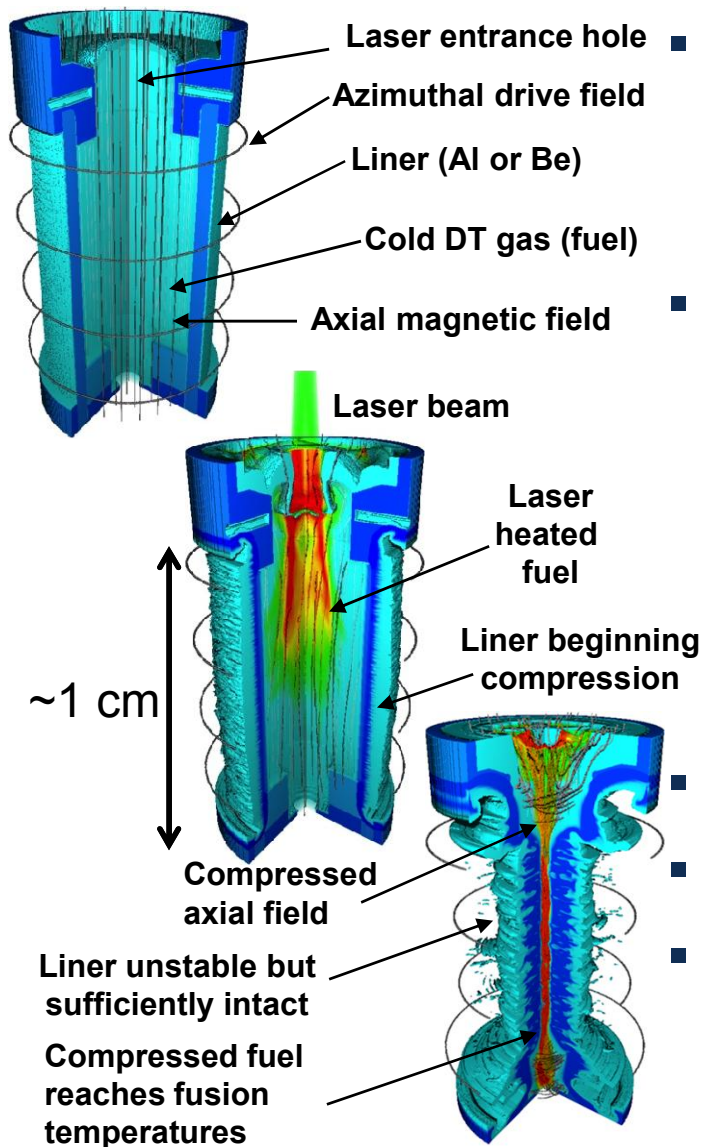
T.J. Awe, C.J. Bourdon, G.A. Chandler, P.J. Christenson, M.E. Cuneo, M. Geissel, **M.R. Gomez**, K.D. Hahn, S.B. Hansen, E.C. Harding, A.J. Harvey-Thompson, M.C. Herrmann, M.H. Hess, C.A. Jennings, B. Jones, M. Jones, R.J. Kaye, P.F. Knapp, D.C. Lamppa, M.R. Lopez, M.R. Martin, R.D. McBride, L.A. McPherson, J.S. Lash, K.J. Peterson, J.L. Porter, G.A. Rochau, D.C. Rovang, C.L. Ruiz, S.E. Rosenthal, M.E. Savage, P.F. Schmit, **A.B. Sefkow**, **D.B. Sinars**, **S.A. Slutz**, I.C. Smith, W.A. Stygar, R.A. Vesey, E.P. Yu

Sandia National Laboratories, Albuquerque, NM 87185 USA

B.E. Blue, D.G. Schroen, K. Tomlinson

General Atomics, San Diego, CA 92186 USA

We are evaluating a **Magnetized Liner Inertial Fusion (MagLIF)*** concept that may reduce fusion requirements

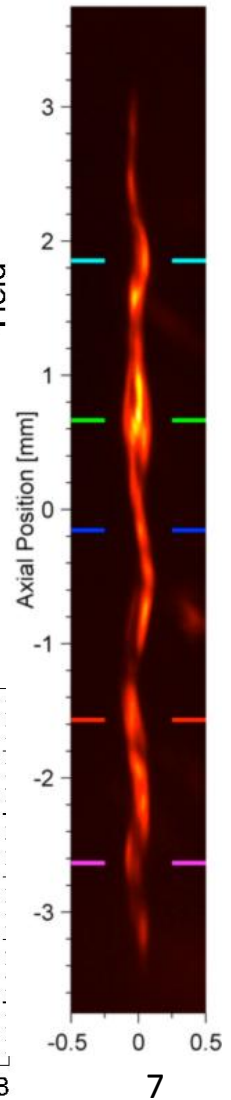
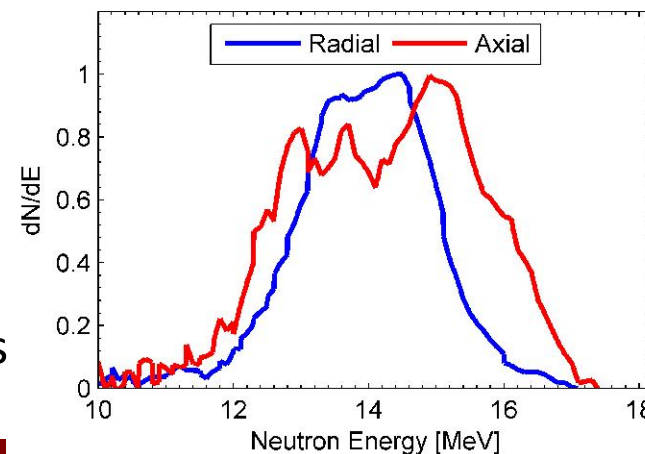
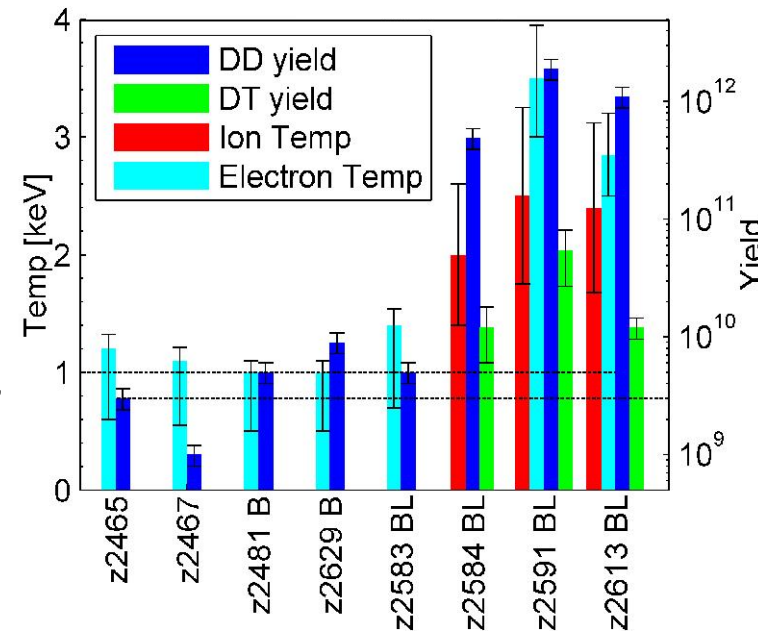


- An initial 30 T axial magnetic field is applied
 - Inhibits thermal conduction losses
 - Appears to stabilize implosion at late times
- During the ~ 100 ns implosion, the fuel is heated using the Z-Beamlet laser (about 6 kJ in designs)
 - Preheating to ~ 300 eV reduces the compression needed to obtain fusion temperatures to 23 on Z
 - Preheating reduces the implosion velocity needed to ~ 100 km/s, allowing us to use thick liners that are more robust against instabilities
- ~ 50 -250 kJ energy in fuel; 0.2-1.4% of capacitor bank
- Stagnation pressure required is ~ 5 Gbar
- DD equivalent of 100 kJ DT yield may be possible on Z in the next few years—this will require enhanced drive upgrades that are in progress, e.g., 10 T \rightarrow 30 T; 2 kJ \rightarrow >6 kJ; 19 MA \rightarrow >24 MA

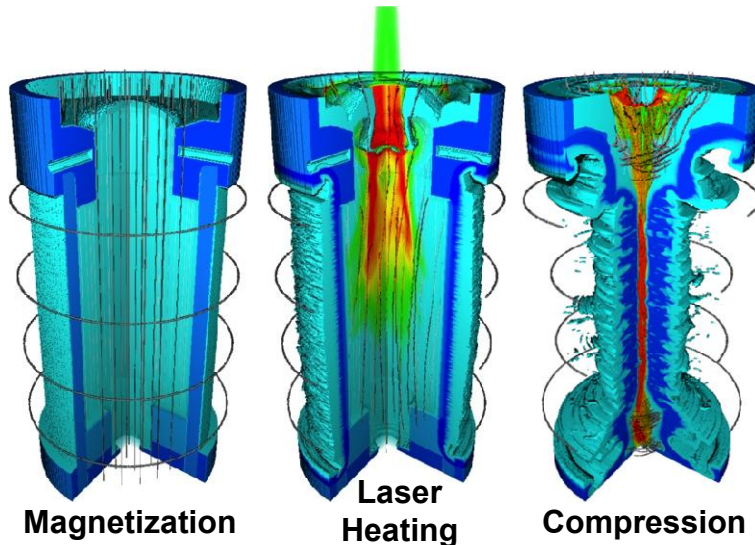
*S.A. Slutz *et al.*, Phys. Plasmas 17, 056303 (2010). S.A. Slutz and R.A. Vesey, Phys. Rev. Lett. (2012).

We obtained promising initial results with MagLIF and are seeking to increase our understanding

- We achieved DD yields up to 2×10^{12} (~ 0.3 kJ DT equivalent) in our first integrated tests of Magnetized Liner Inertial Fusion (MagLIF)
- A variety of data were collected that appear to show a $< 150 \mu\text{m}$ diameter, ~ 3 keV, highly magnetized plasma was produced—remarkable for a 70-100 km/s implosion!
- We are continuing to build on these results with a balanced combination of focused and integrated experiments
- In parallel we are improving capabilities to understand how this performance will scale with increasing drive parameters



Our path to studying the underlying science is a mixture of focused and integrated experiments to address key physics



■ Key target design elements

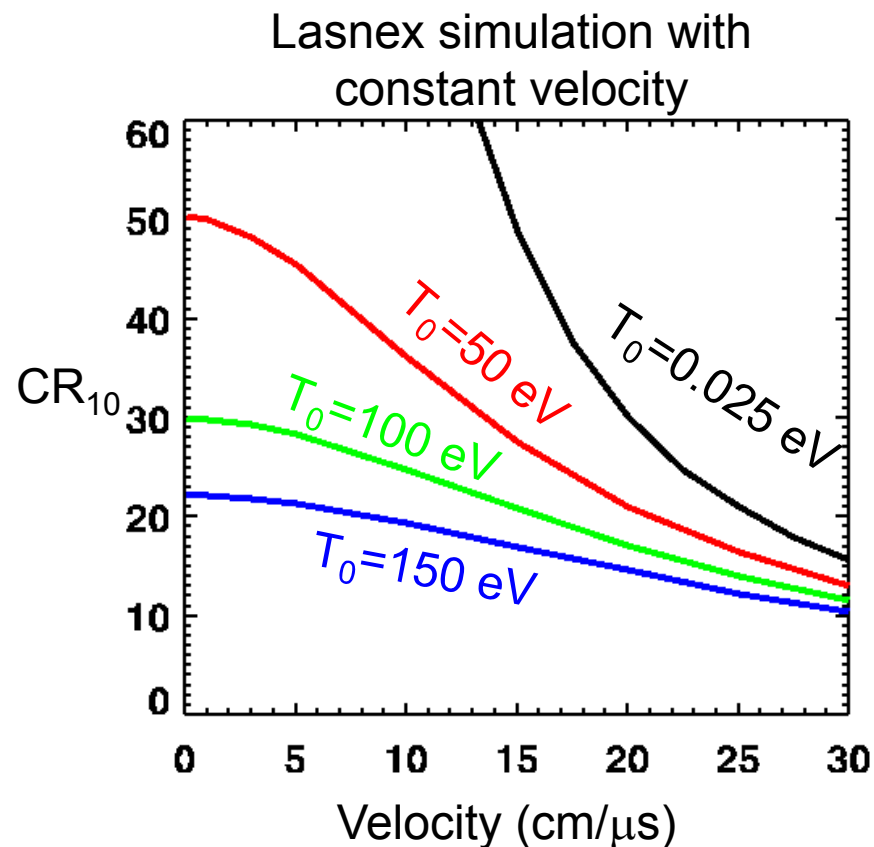
- Liner compression
- Magnetization
- Laser heating

■ Key physics uncertainties

- Liner instabilities
 - Electro-thermal
 - Magneto-Rayleigh-Taylor
 - Deceleration RT
 - Impact of 3D fuel assembly
- Liner/fuel interactions & mix
- Laser-window and laser-fuel scattering, absorption, uniformity
- Suppression of electron heat transport in dense plasma by magnetic fields
- Magnetic flux compression
- Magnetized burn

Experiments to address the key physics are planned for the Z pulsed power facility and the Z-Beamlet and Omega (and -EP) lasers.

Typical ICF implosions need high velocities to reach fusion temperatures—starting the implosion with heated fuel potentially reduces requirements



CR_{10} = Convergence Ratio (R_0/R_f) needed to obtain 10 keV (ignition) with no radiation losses or conductivity

Heating fuel to ignition temperatures is typically done with a high-velocity shock (or series of shocks)

High velocities make it easier to reach fusion temperatures and also reduce the time available for losses (e.g., electron heat conduction or radiation)

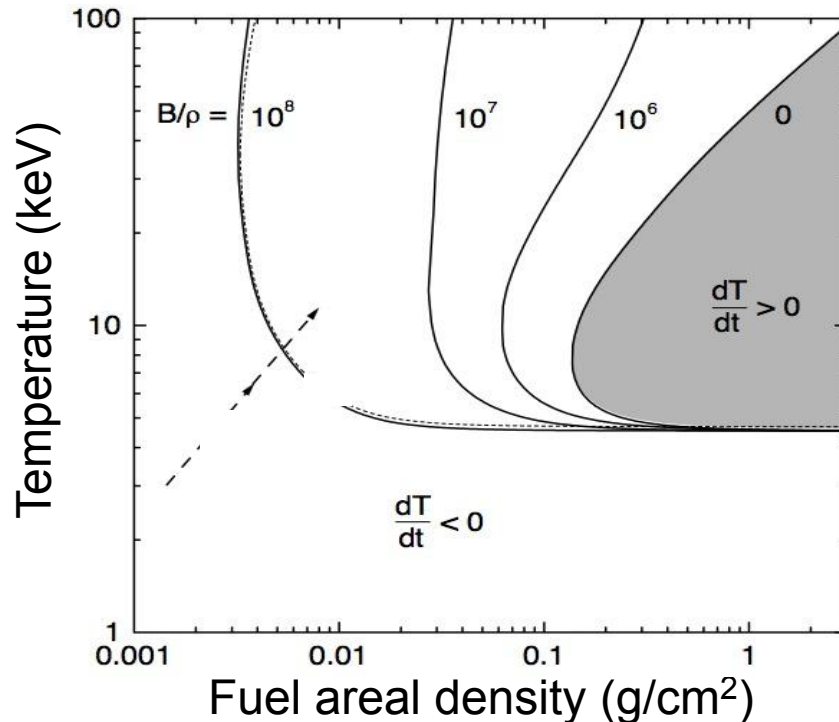
Heating the fuel prior to the implosion *in the absence of losses* can allow low-velocity, low-convergence implosions to reach ignition temperatures

Is there a way to reduce losses?

A large, embedded magnetic field can significantly reduce electron conduction losses from heated fuel



*Basko et al. *Nuc. Fusion* 40, 59 (2000)



The ρr needed for ignition can be significantly reduced by the presence of a strong magnetic field largely through inhibiting electron conduction

Lower ρr reduces the required final fuel density (e.g., $\sim 1 \text{ g/cc} \ll 100 \text{ g/cc}$), which also reduces bremsstrahlung radiation losses

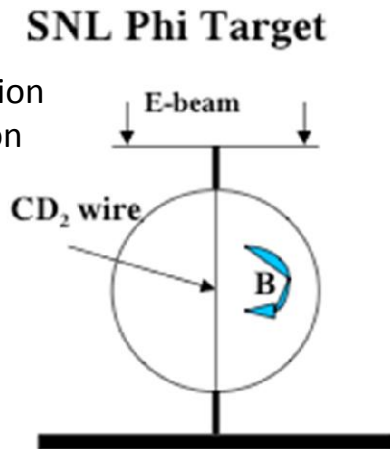
This means the stagnation plasma pressure at ignition temperatures is significantly reduced (e.g., $\sim 5 \text{ Gbar} \ll \sim 500 \text{ Gbar}$ for hot spot ignition)

Large values of B/ρ are needed and therefore large values of B are needed, $B \sim 10,000 \text{ Tesla}$ (Earth's B -field is $\sim 0.00003 \text{ Tesla}$)

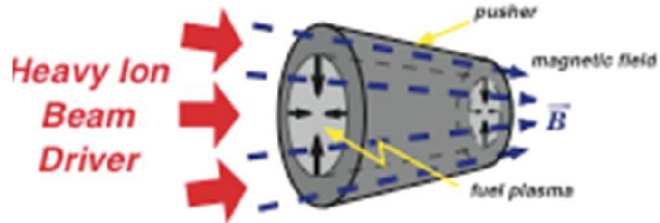
This field significantly exceeds pulsed coil technology ($B_0 \sim 10\text{-}30 \text{ T}$), therefore flux compression is needed

Many groups want to use magnetic fields to relax inertial fusion stagnation requirements

1982 Demonstration of enhanced fusion yield with magnetization (~1e6 DD yield)



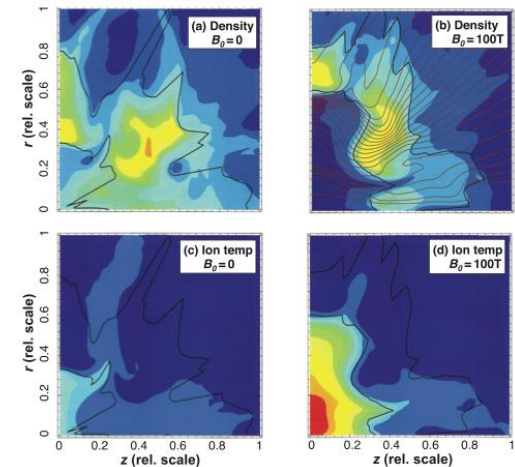
Max Planck/ITEP



Basko, Kemp, Meyer-ter-Vehn, *Nucl. Fusion* **40**, 59 (2000)
Kemp, Basko, Meyer-ter-Vehn, *Nucl. Fusion* **43**, 16 (2003)

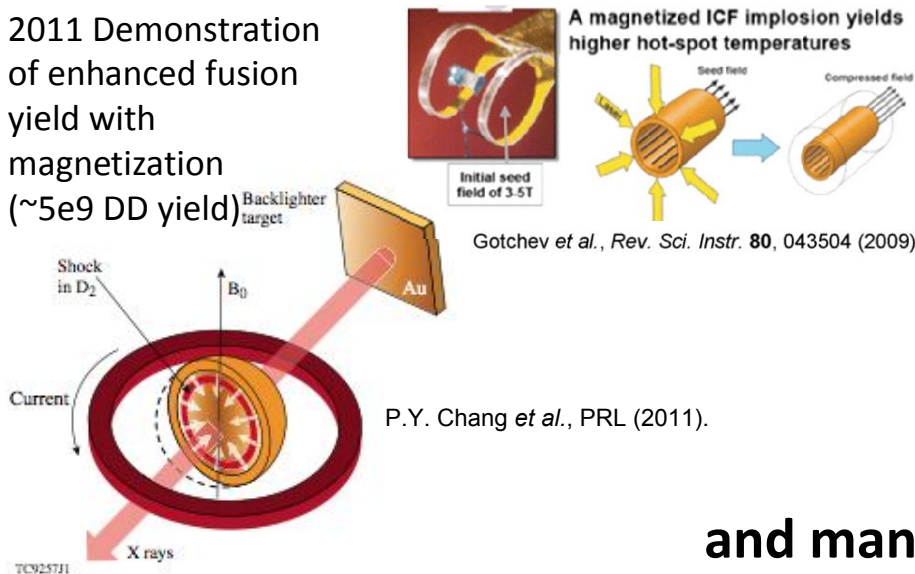
LLNL

(Perkins et al., *Phys Plasmas* 2013)



University of Rochester/LLE

2011 Demonstration of enhanced fusion yield with magnetization (~5e9 DD yield)



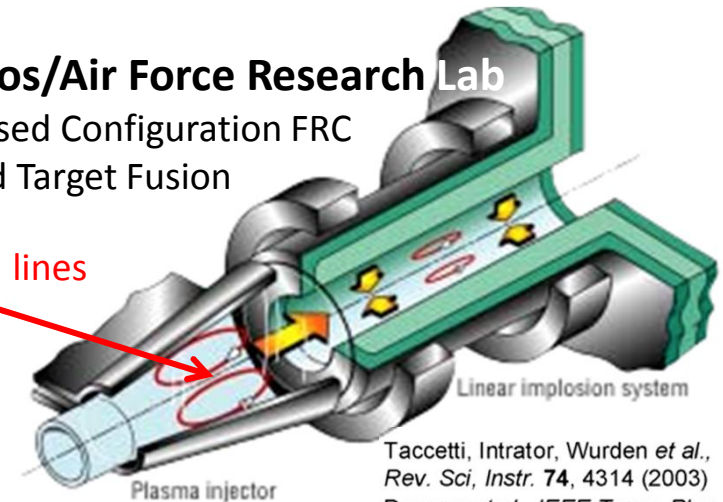
Los Alamos/Air Force Research Lab

Field Reversed Configuration FRC

Magnetized Target Fusion

Shiva Star

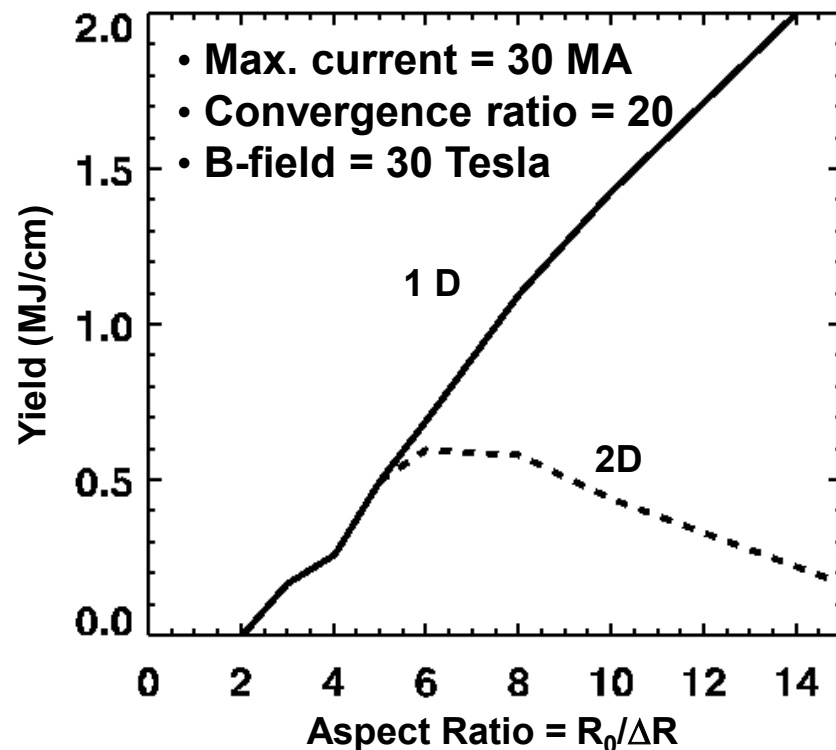
closed field lines
FRC



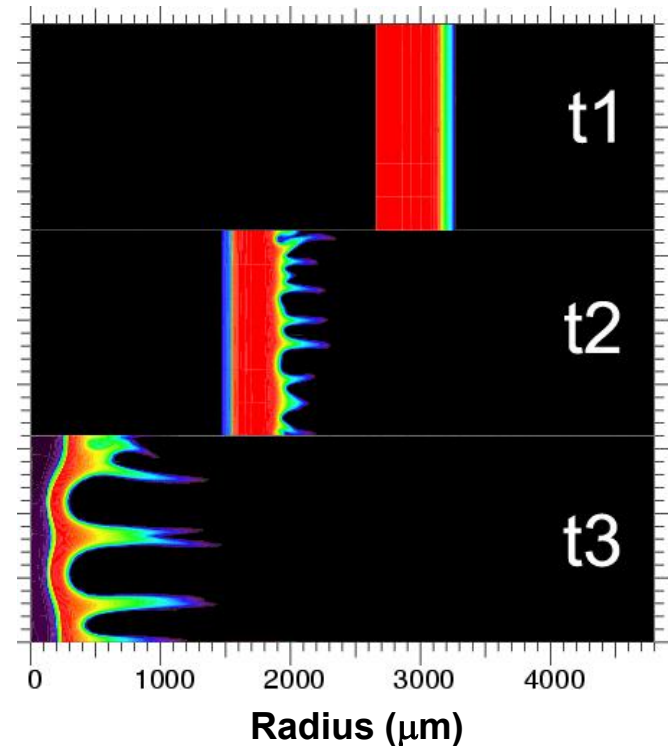
Taccetti, Intrator, Wurden et al., *Rev. Sci. Instr.* **74**, 4314 (2003)
Degnan et al., *IEEE Trans. Plas. Sci.* **36**, 80 (2008)

and many others...

Reducing the implosion velocity requirements through fuel heating and magnetization allows us to use thicker, more massive liners to compress the fuel that are more stable



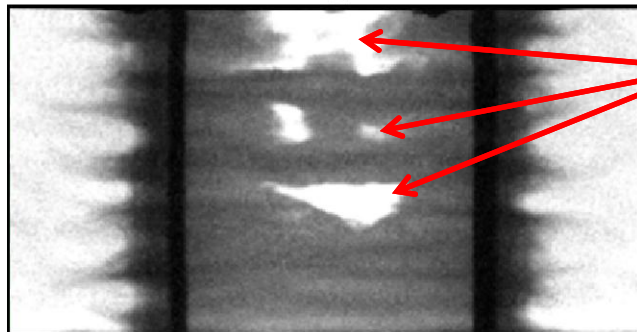
- The Magneto-Rayleigh-Taylor instability degrades the yield as the aspect ratio is increased (due to decreased liner ρR)



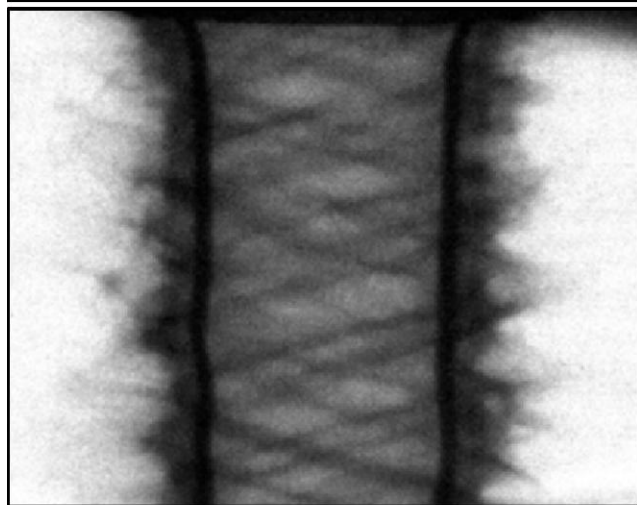
- Simulations of AR=6 Be liner show reasonably uniform fuel compression and sufficient liner ρR at stagnation to inertially confine the fuel—important because fuel density is low!

Adding an axial magnetic field reduces hard x rays and hot spots, and changes the liner instability structure from cylindrical to helical—evidence it is doing something!

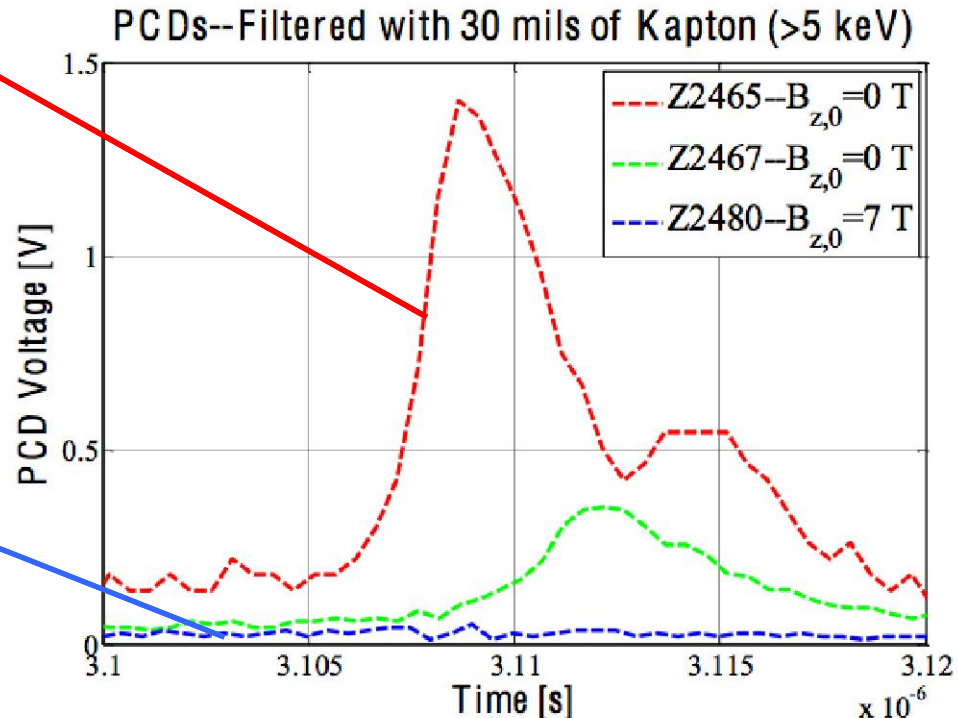
Without Magnetic Field



Time-integrated self-emission from liner implosion at 6151 eV; missing in shots with axial field



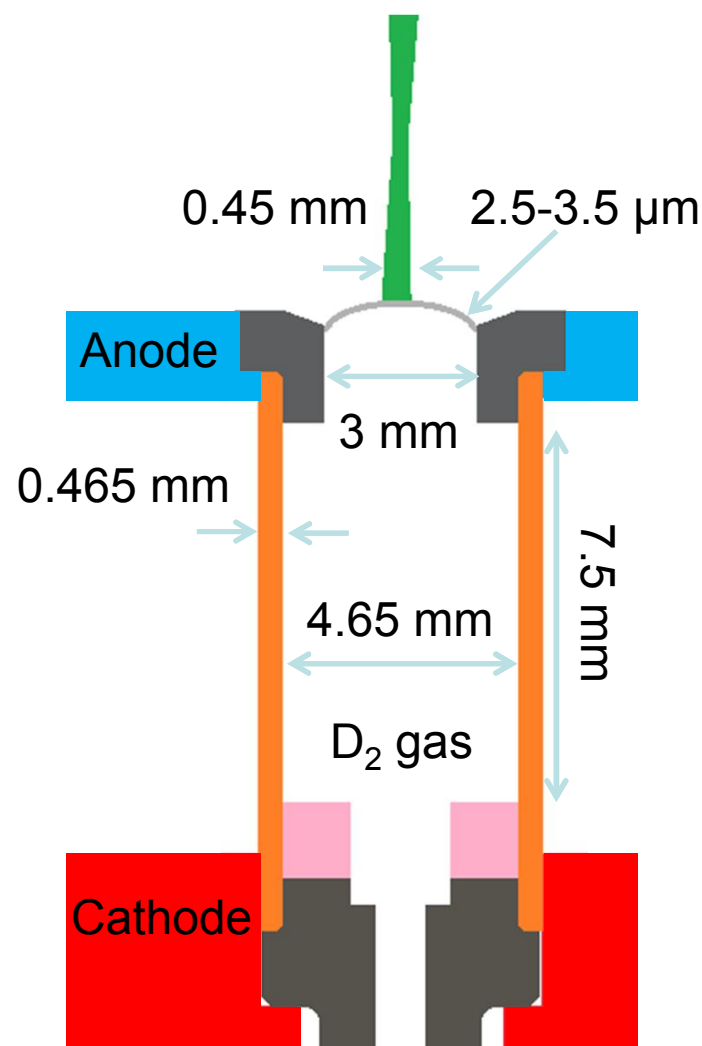
With Magnetic Field



If magnetic flux roughly conserved the additional magnetic pressure from the axial field will suppress micro-pinching—this is indirect evidence for flux compression 13

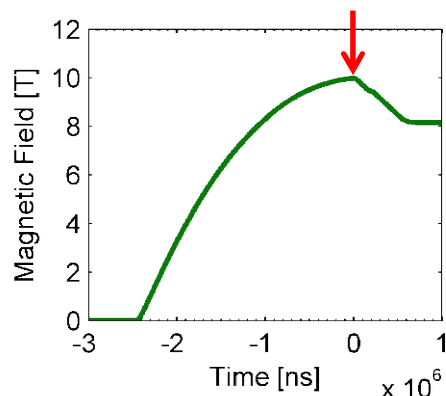
The target design for our initial experiments incorporates the knowledge gained from focused experiments and extensive simulations

- Beryllium liner with aspect ratio 6
 - Thick liner is more robust to instabilities
 - Still allows diagnostic access > 5 keV
- Top and bottom implosion cushions
 - Mitigates wall instability
- Standoff between LEH and imploding region
 - Avoid window material mixing with fuel
- Exit hole at bottom of target
 - Avoid interaction with bottom of target



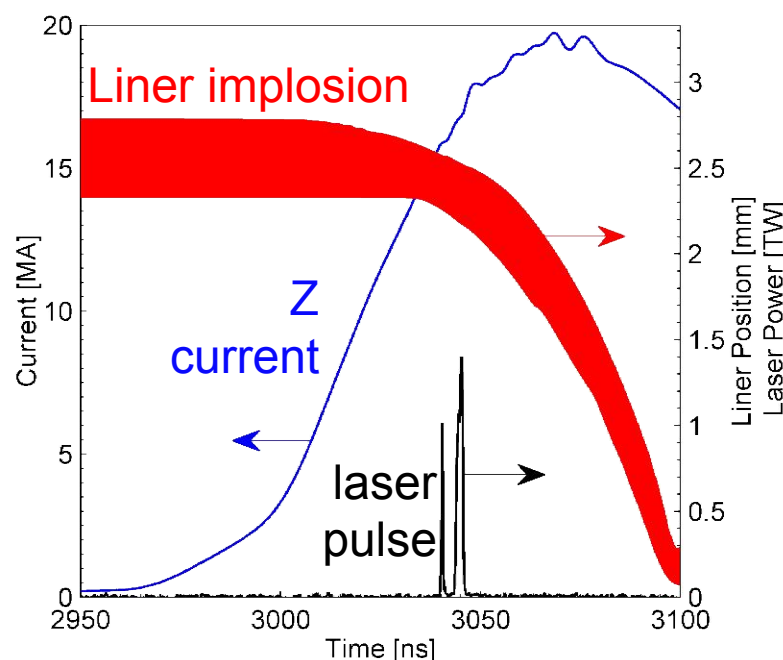
Initial experiments were conducted at $I = 19 \text{ MA}$, $B = 10 \text{ T}$, and Laser = 2.5 kJ

Time of
experiment



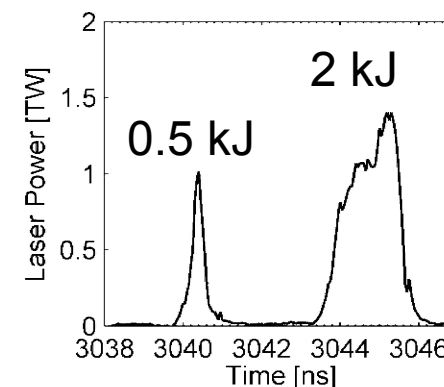
Magnetic field risetime
is approximately 2 ms

B is constant over the
timescale of the
experiment

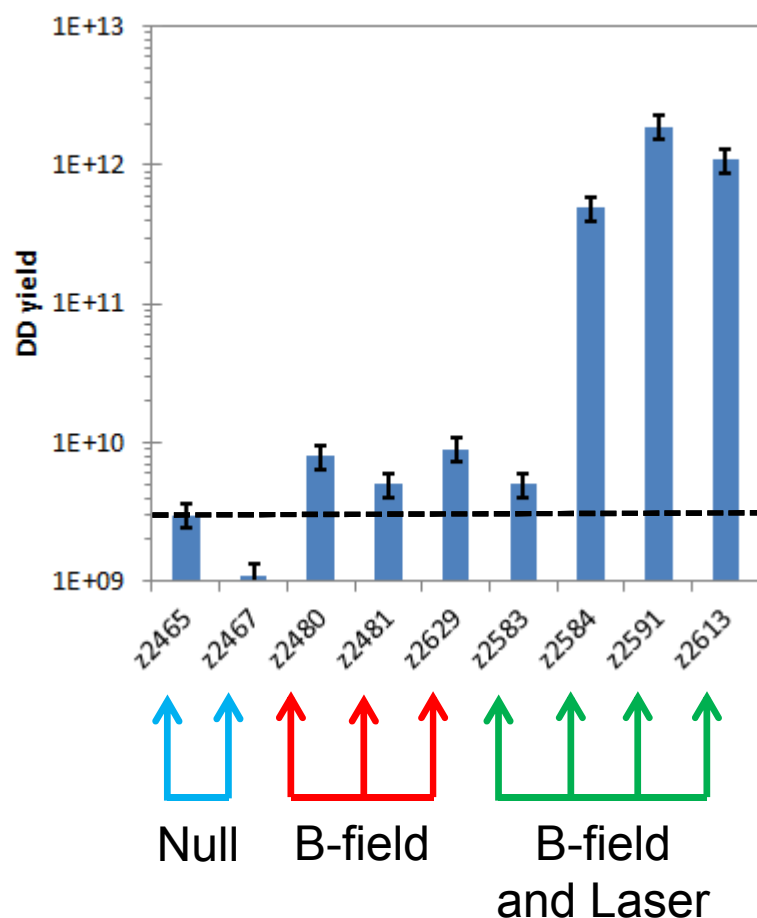


Peak current is 19 MA
Magnetic field is 10 T
Total laser energy is 2.5 kJ

Laser energy is split
into 2 pulses:
1st pulse intended to
destroy LEH
2nd pulse intended to
heat fuel

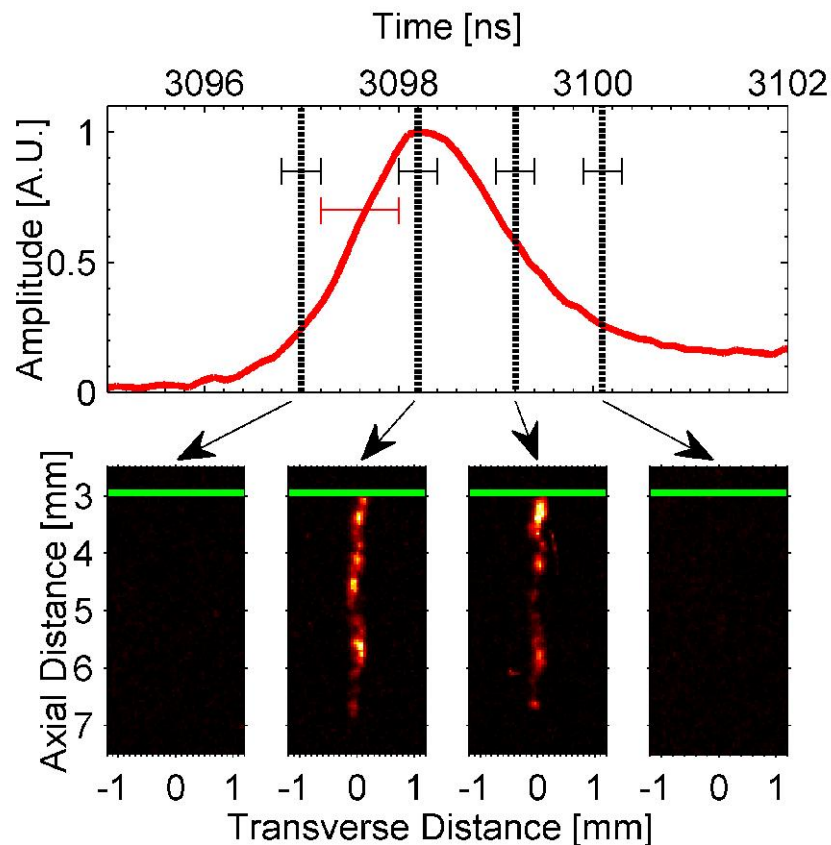


Z shots producing DD yields in excess of 10^{12} were only observed in experiments with laser and B-field



- High yields were only observed on experiments incorporating both applied magnetic field and laser heating
- A series of experiments without laser and/or B-field produced yields at the background level of the measurement
- Result of z2583 is not well understood nor reproduced at this time

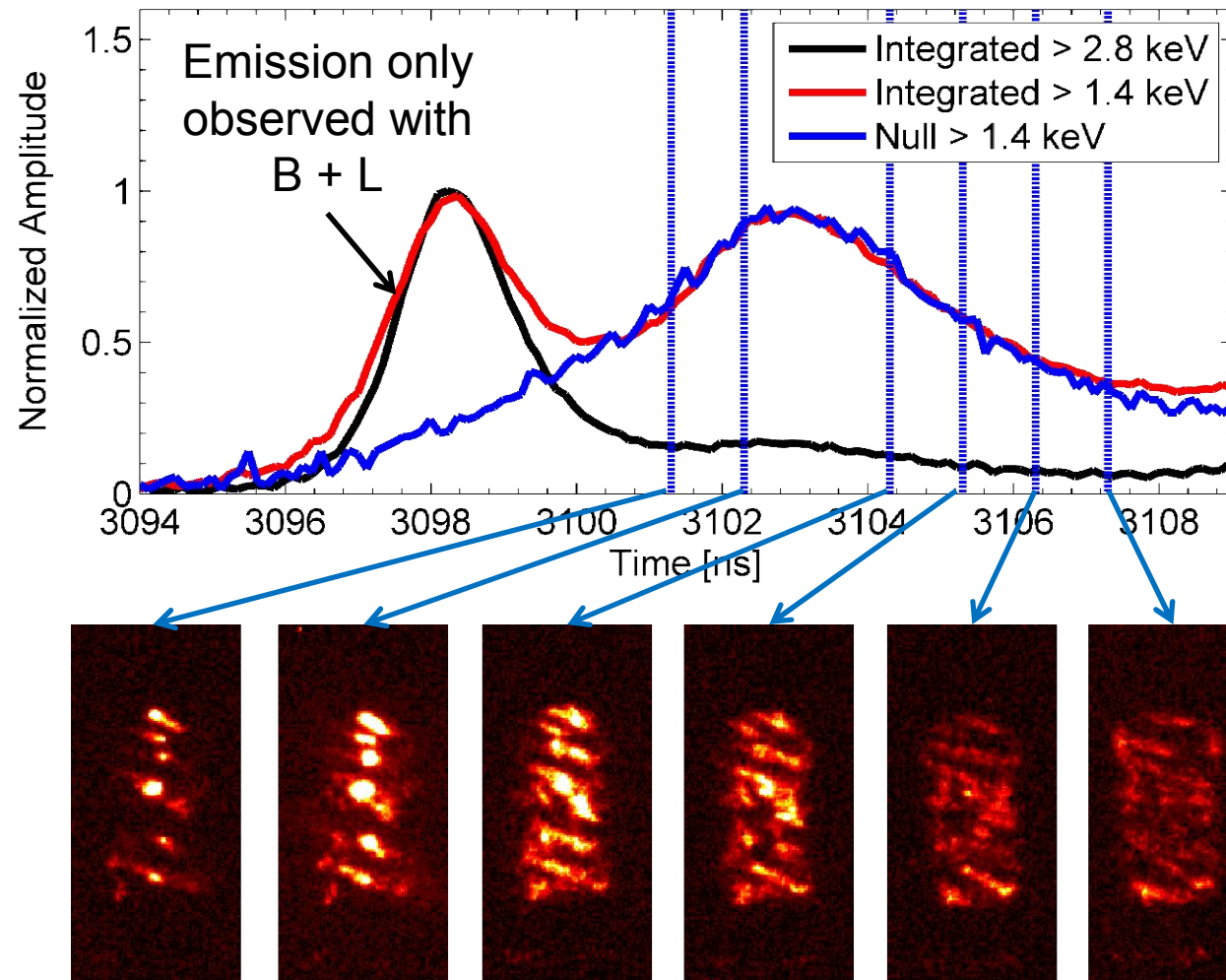
Time-resolved x-ray pinhole imaging ($h\nu > 2.8$ keV) shows a narrow emission column during peak in X-ray signal



- Emission column is observed only during the peak in the x-ray signal
- Emission column is only observed on experiments with high neutron yield
- Stagnation column width is at the resolution limit of this instrument (~ 150 microns)

High energy x-ray signal and narrow emission region are absent in null experiments

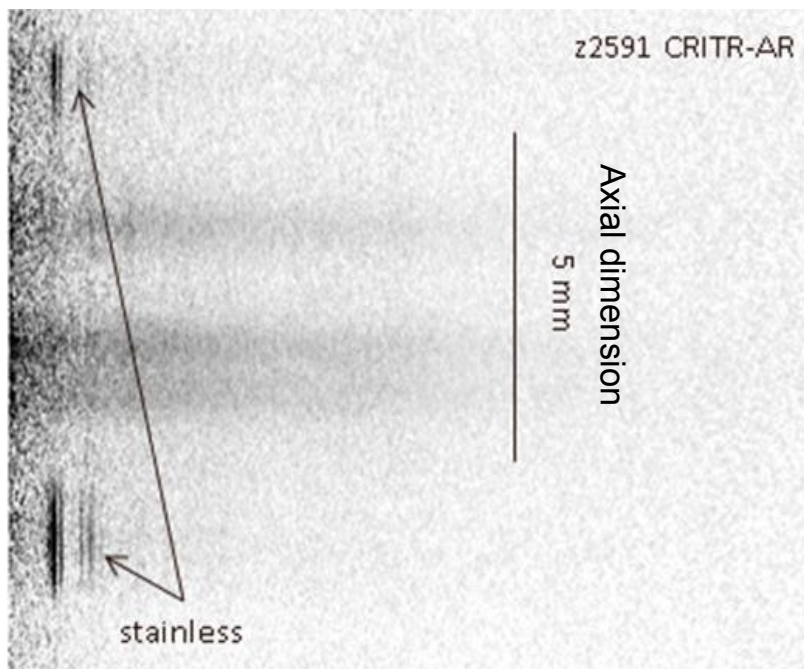
- Liner emission is observed in all experiments
- Liner emission is at a lower photon energy (< 2.8 keV)
- Liner emission is getting larger at late times



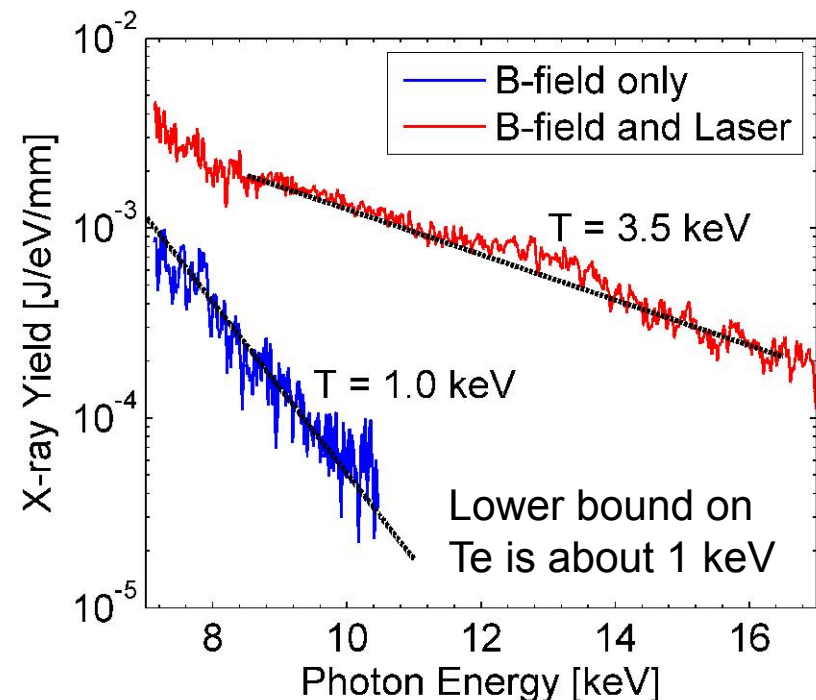
High-energy spectra show axial variations in temperature and composition, with ~ 3.5 keV electron temperature in the pinch region—remarkable for a 70-100 km/s implosion!



Emission lines from stainless steel (Fe, Cr, Ni) appear at the anode and cathode, but minimal high-Z contamination is observed in hot central regions



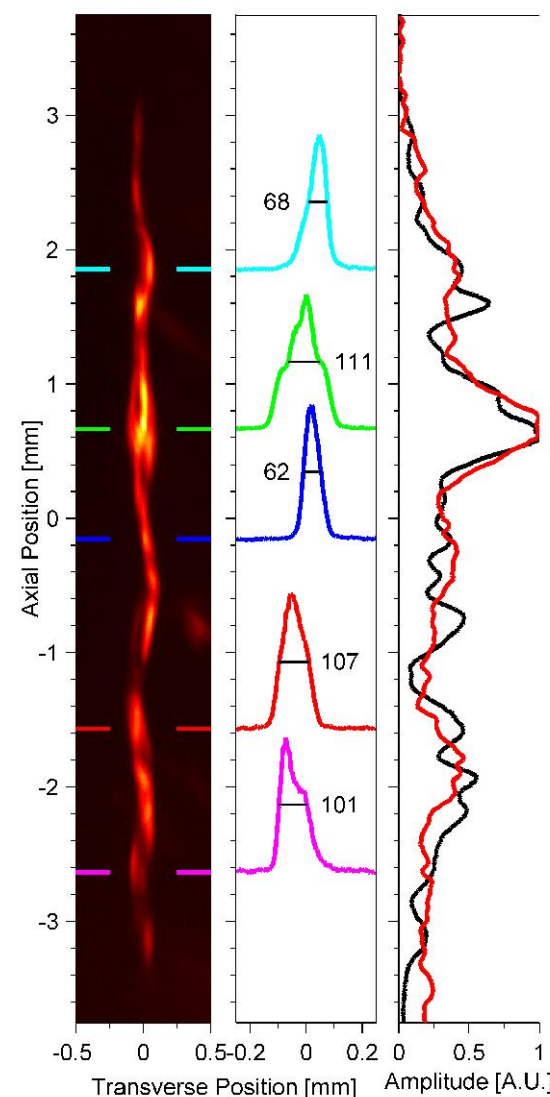
The slope of the high-energy continuum emission implies $T_e \sim 1.5$ keV at the anode and cathode, and $T \sim 3.5$ keV in the central regions



The measured electron temperature is close to the ion temperature obtained from neutron time-of-flight data; x-ray emission yields are consistent with fuel $\rho \sim 0.4$ g/cm³

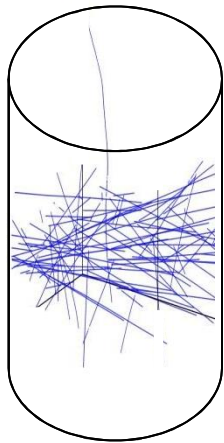
High-resolution monochromatic imaging of the x-ray emission shows a narrow, hot plasma column with weakly helical structure

- Lineouts of stagnation column vary from 60 to 120 μm FWHM (resolution about 60 μm)
- Emission is observed from about 6 mm of the 7.5 mm axial extent
- Note that the emission doesn't necessarily define the fuel-liner boundary, but only the hot fuel region
- The stagnation column is weakly helical with a wavelength of about 1.3 mm and a 0.05 mm horizontal offset
- Axial lineouts of image (black) agree with 9.3 keV 1D spectrometer lineouts (red), suggesting features are due to emission and not liner opacity (Be opacity >9 keV small).
- With $\rho \sim 0.4 \text{ g/cm}^3$, $\rho_r \sim 2 \text{ mg/cm}^2$

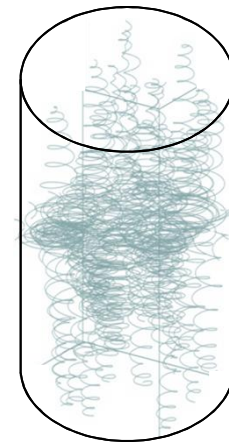
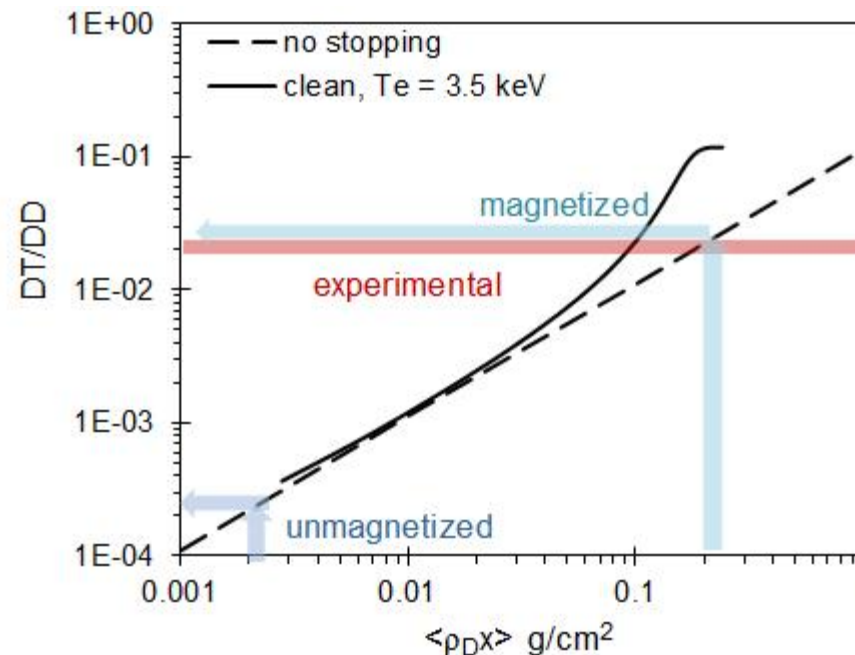


In addition to the significant $\sim 2 \times 10^{12}$ DD neutron yields, we measure a remarkable $\sim 5 \times 10^{10}$ DT neutrons

“Secondary” 14 MeV neutrons can be produced by 1 MeV tritons interacting with D fuel



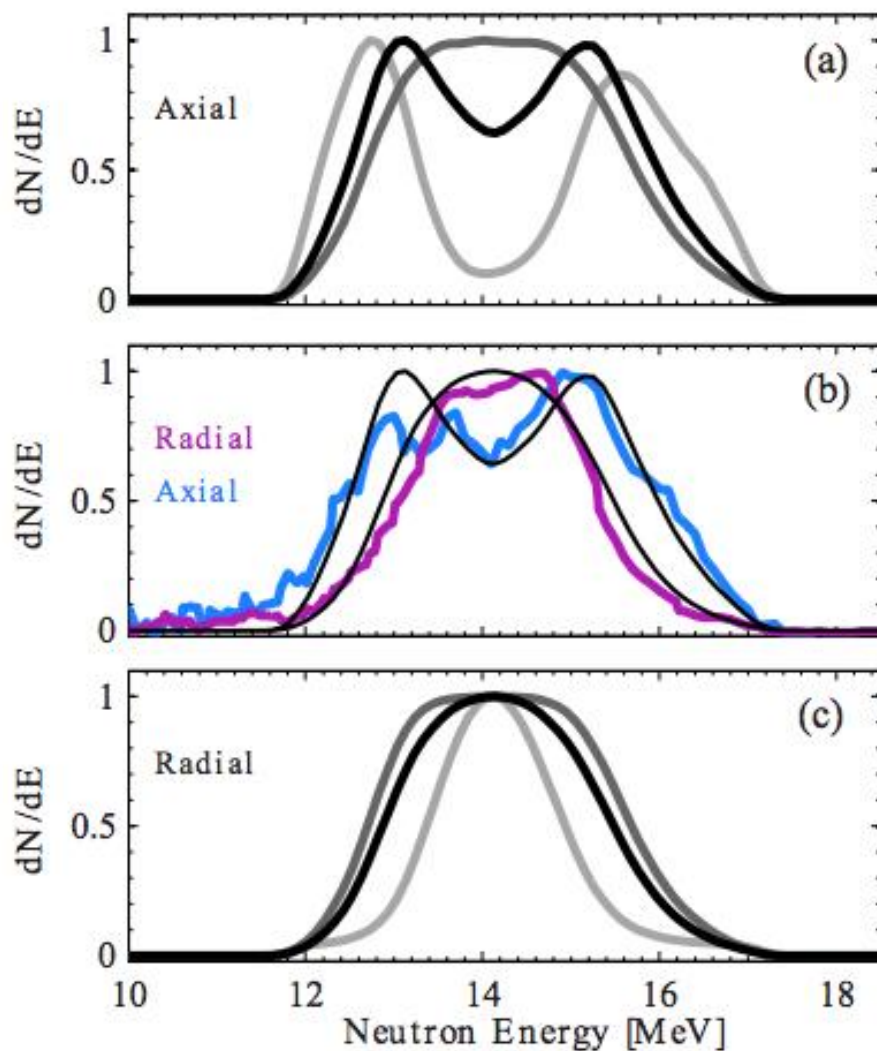
Unmagnetized plasmas must reach pressures of ~ 500 Gbar and $\rho R > 0.2$ g/cm² to achieve the α -particle confinement required for ignition



In magnetized plasmas, thermal confinement and α -deposition are both enhanced by B, reducing pressure and ρR requirements by factors of ~ 100 .

A field that confines tritons also confines electrons -- and will confine alphas!

Neutron time-of-flight data are consistent with high magnetization



$2.5e5$ G-cm
 $4.5e5$ G-cm
 $7.5e5$ G-cm

nTOF spectra consistent with $\sim 4.5e5$ G-cm

DT/DD ratio consistent with $>4e5$ G-cm

$2.5e5$ G-cm
 $4.5e5$ G-cm
 $7.5e5$ G-cm

MagLIF Summary

- Results from initial MagLIF experiments have been encouraging, with significant DD and DT yields and strong evidence for good stability, confinement, and scaling
- A helical stagnation column with $T \sim 3$ keV, $\rho \sim 0.4$ g/cc, $r \sim 50$ μ m, and $B_z \sim 10$ kT is consistent with an extensive collection of neutron and x-ray data
- Both integrated and focused experiments are ongoing
- Better understanding of how high magnetic fields affect thermal transport and stopping power will increase confidence in our predictions for yield scaling with increasing current, external field, and laser power

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Encouraging results from initial integrated experiments
- **Z Astrophysical Plasma Properties (ZAPP) Collaboration**
Studies of emission and absorption from photoionized plasmas relevant to white dwarfs, accretion disks, and solar photosphere
- **Diagnostic development on Z**
X-ray Thomson scattering and Zeeman splitting

There is a lot of interesting science going on at Z –
and many opportunities for significant contributions from Cornell's LPS.

These efforts represents a large collaboration between the NNSA labs and the academic community



Jim Bailey, Taisuke Nagayama,
Guillaume Loisel, Stephanie Hansen,
Dave Bliss, Tom Nash, Tom Ao, Eric
Harding, Greg Rochau, Matt Gomez,
Michael Desjarlais
Sandia National Laboratories



Roberto Mancini, Iain Hall, Tom
Lockard, Dan Mayes
University of Nevada – Reno



Don Winget, Mike Montgomery, Ross
Falcon, Thomas Gomez, Alan Wootton,
Jennifer Ellis, Sean Moorhead, Roger
Bengtson
University of Texas – Austin



Anhil Pradhan, C. Orban, Mark
Pinsonneault, and S.N. Nahar
Ohio State University



Mark Koepke, Ted Lane, Matt Flaugh
West Virginia University



Duane Leidahl, Carlos Iglesias, Brian Wilson
Lawrence Livermore National Laboratory



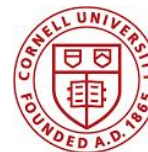
Manolo Sherrill, Heidi Tierney, Chris Fontes,
James Colgan, Dave Kilcrease
Los Alamos National Laboratory



C. Blancard, Ph. Cosse, G. Faussurier, F.
Gilleron, J.C. Pain
**French Alternative Energies and Atomic
Energy Commission (CEA)**



Joe MacFarlane, Igor Golovkin
Prism Computational Sciences



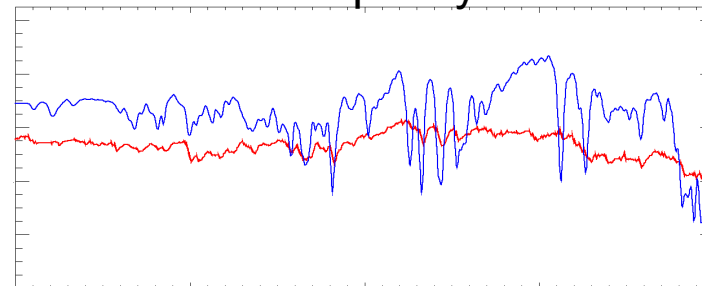
Laura Johnson
Cornell University

We are interested in developing
new collaborations

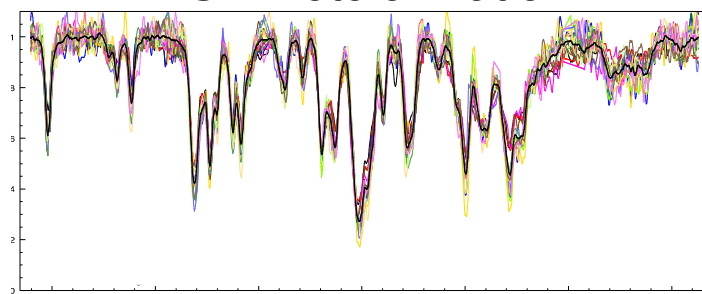
ZAPP experiments measure the fundamental properties of atoms in plasmas to solve important astrophysical puzzles:

- Why can't we predict the location of the convection zone boundary in the Sun?
 - Opacity of Fe at $T \sim 200$ eV
- How does ionization and line formation occur in accreting objects and warm absorbers?
 - Ionization distribution and spectral properties of photoionized Ne and Si
- Why doesn't spectral fitting provide the correct properties for White Dwarfs?
 - Stark-broadened H-Balmer line profiles

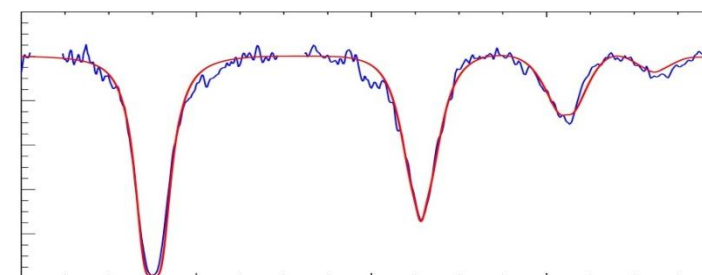
Fe Opacity



Si Photoionization



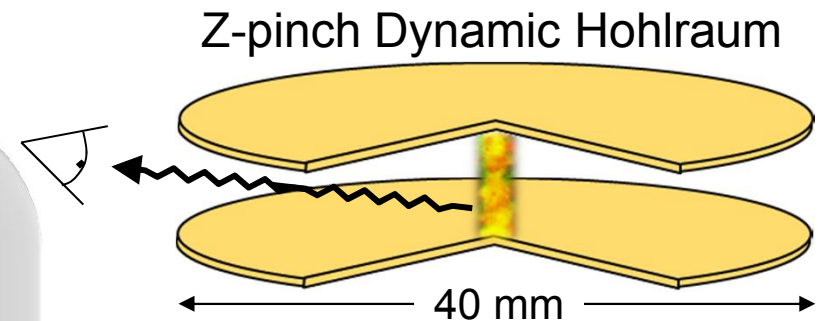
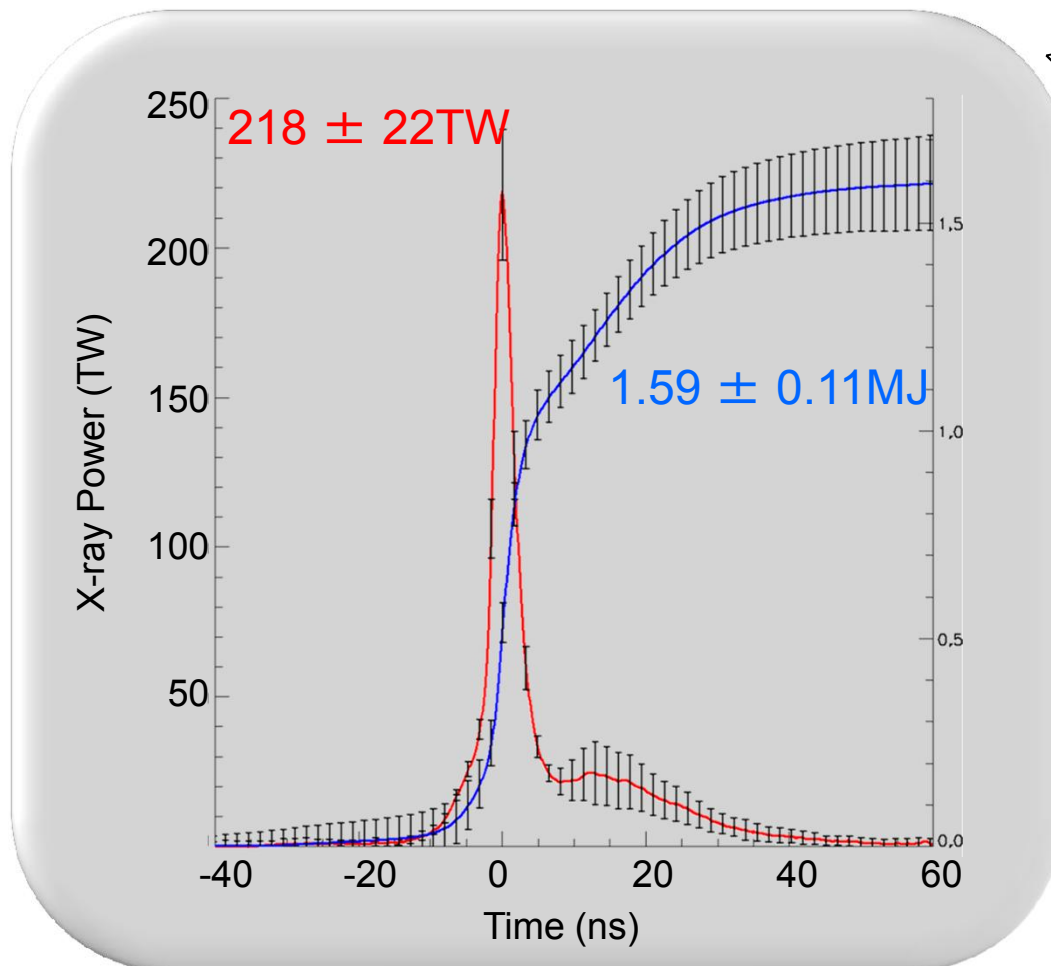
H-Balmer Line shapes



The ZPDH x-ray emission is reproducible to $\pm 10\%$ in peak power and $\pm 7\%$ in energy



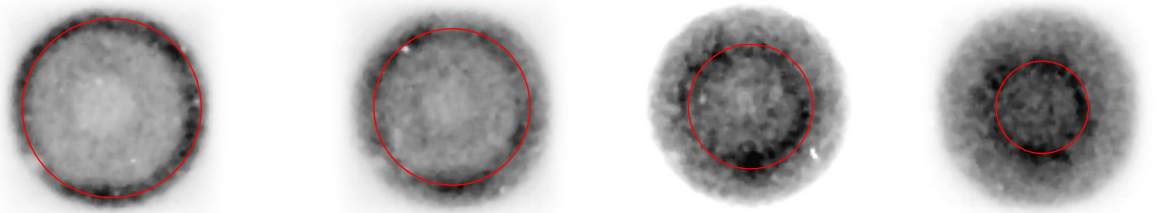
Radial X-ray Power and Energy (20 shot average)



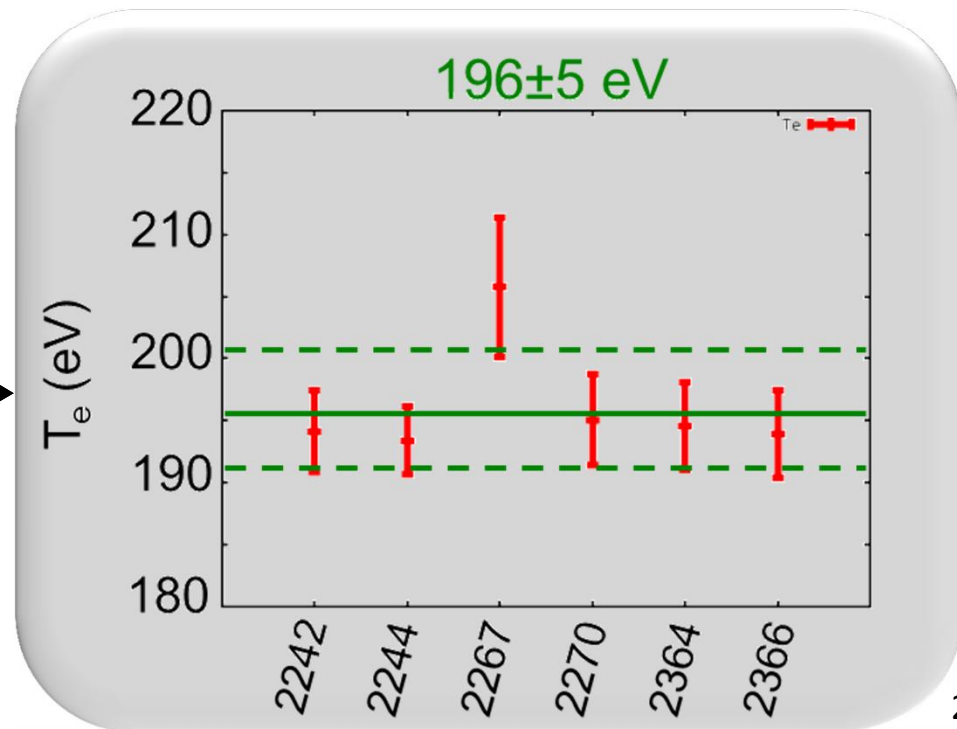
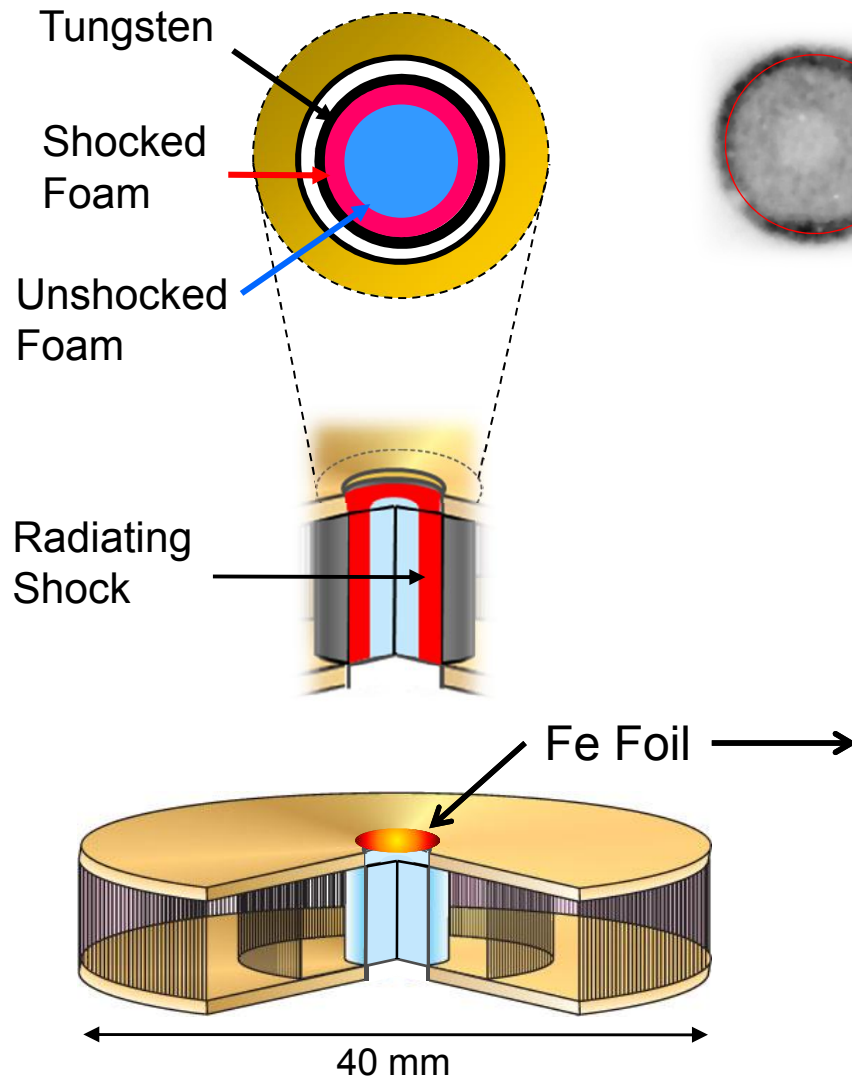
	ZR >2011	Z <2007
Marx Energy	20.3 MJ	11.4 MJ
I _{peak}	25.8 MA (1.5%)	21.7 MA* (2.1%)
Mass	8.5 mg	3.8 mg
Peak Power	220 TW (10%)	120 TW (14%)
Radiated Energy	1.6 MJ (7%)	0.82 MJ (17%)

The ZPDH can also radiatively heat samples placed above the z-pinch to $T_e \sim 200$ eV.

Framing Pinhole Camera Images

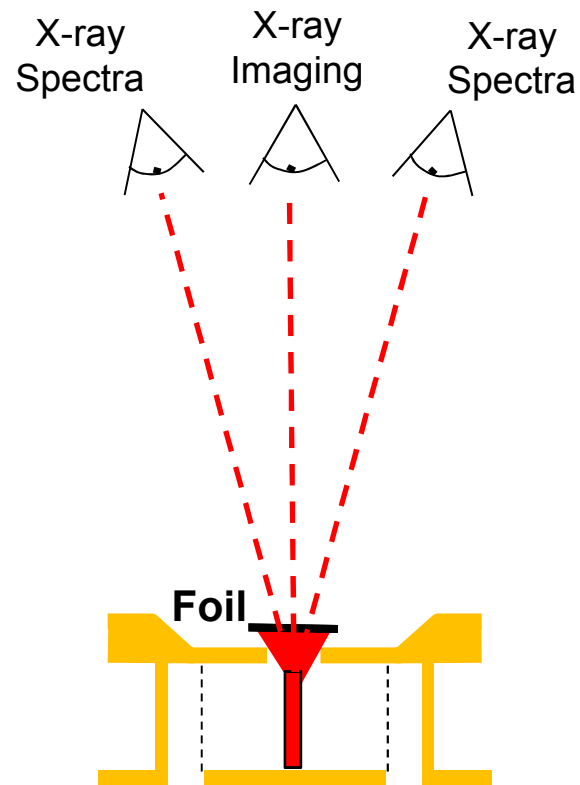


Axial Fe Foil Temperature

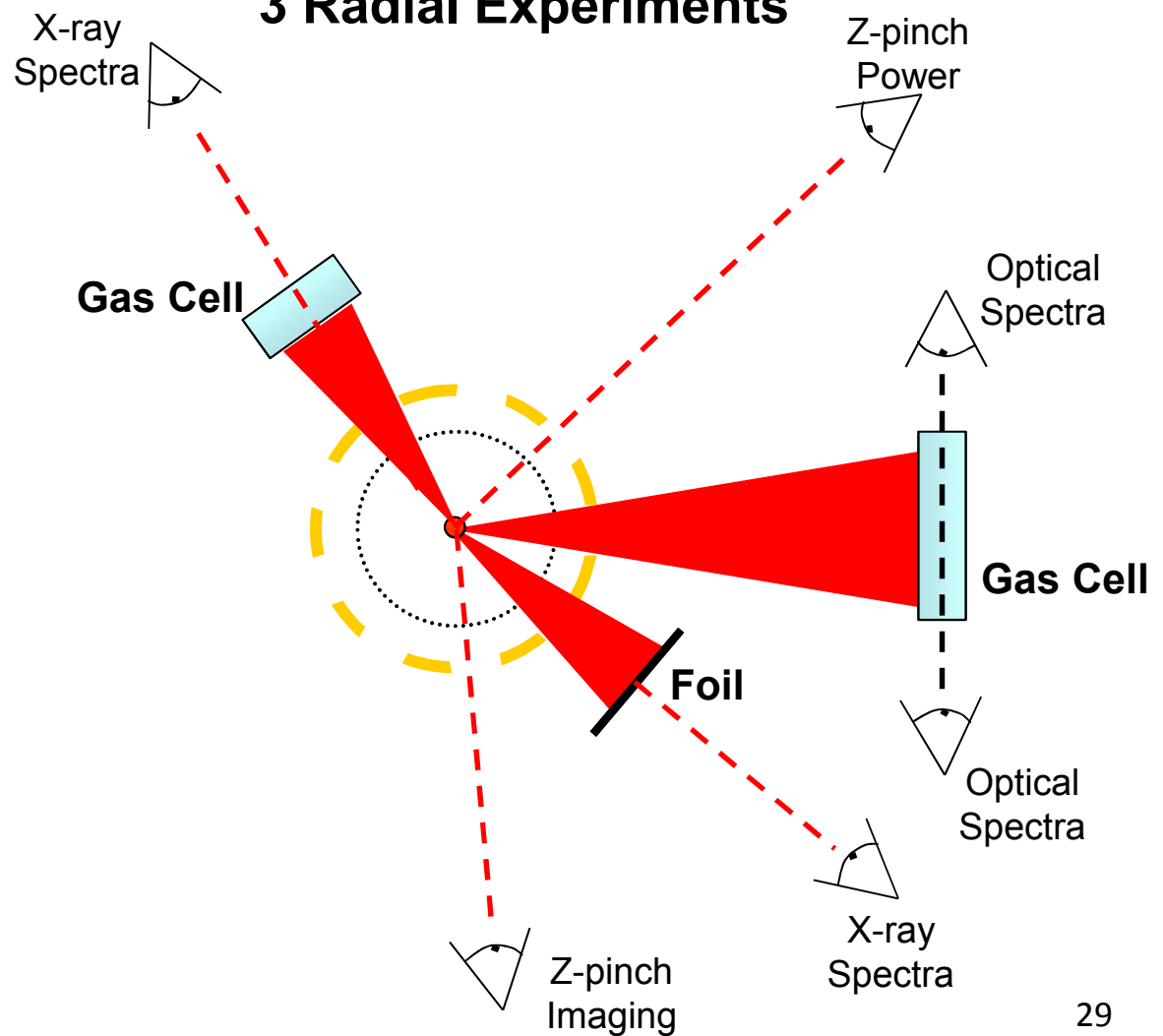


The ZPDH simultaneously drives four independent experiments on a single ZAPP shot

1 Axial Experiment

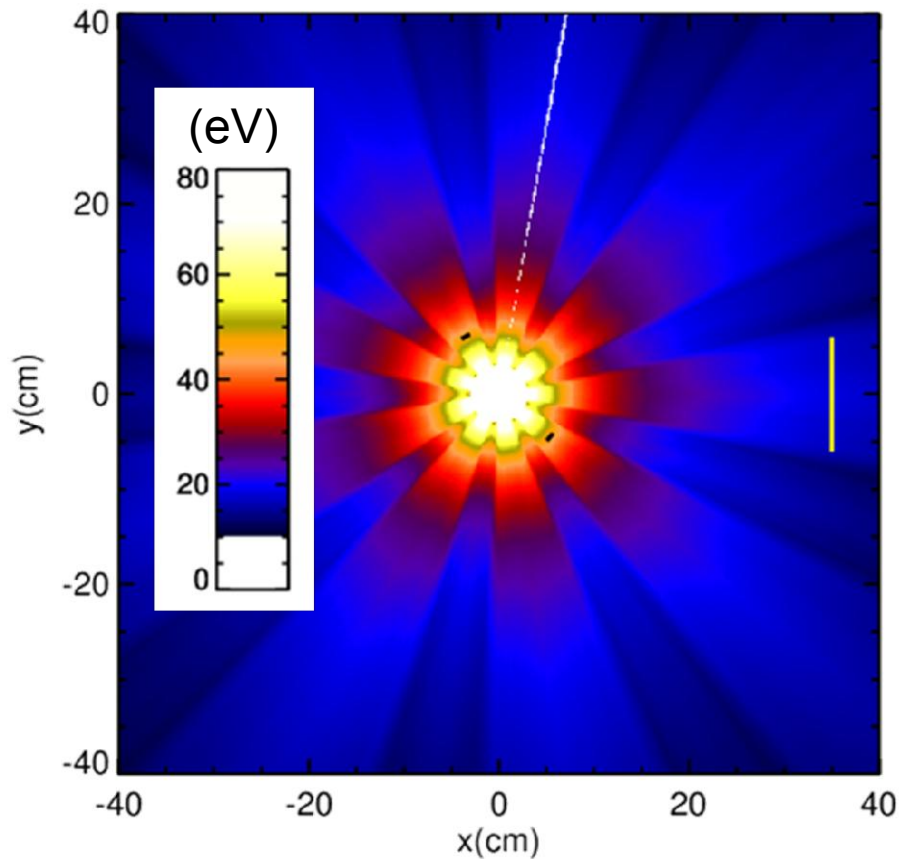


3 Radial Experiments

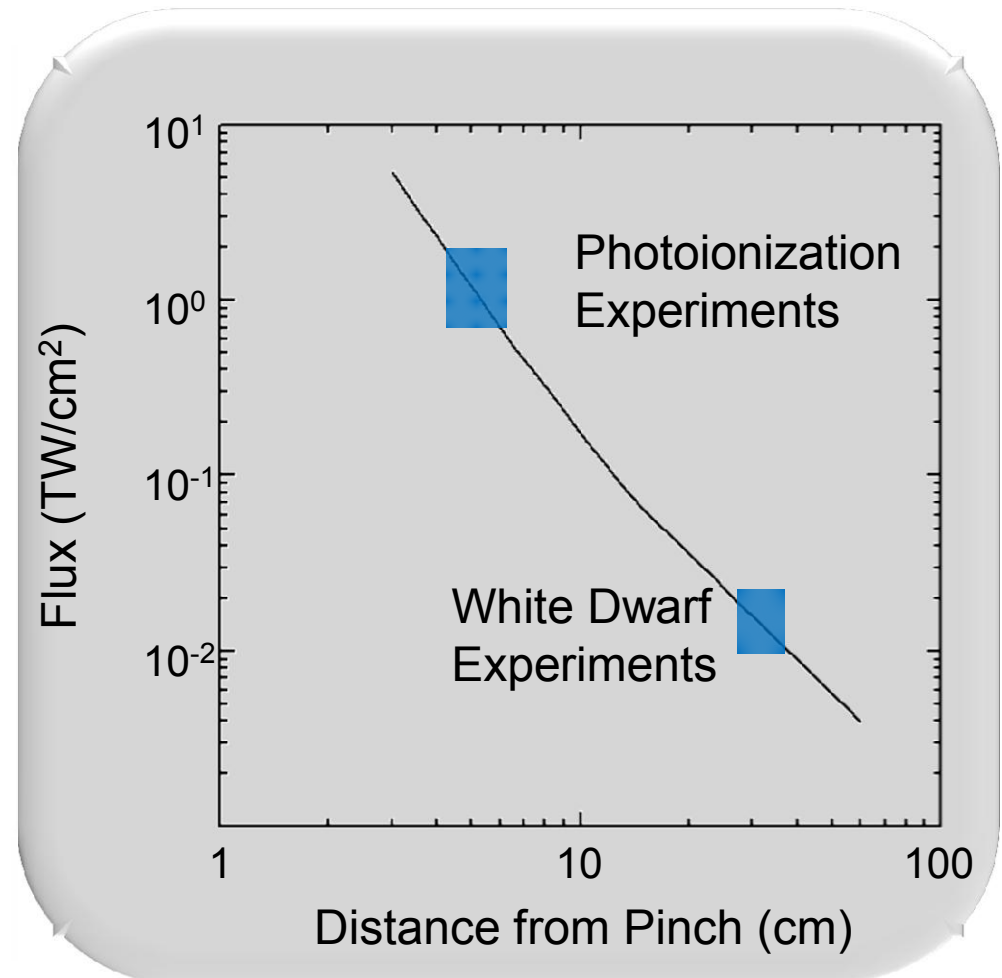


Placing samples at multiple distances from the z pinch provides a broad range of drive flux.

**r- θ Peak Brightness Temperature
Contours Around Z Pinch**



Peak Drive Flux on a Sample



ZAPP campaigns simultaneously study multiple issues spanning 200x in temperature and 10^6 x in density



Solar Opacity



Question:

Why can't we predict the location of the convection zone boundary in the Sun?

Achieved Conditions:

$T_e \sim 200$ eV, $n_e \sim 10^{23}$ cm $^{-3}$



Photoionized Plasmas



Question:

How does ionization and line formation occur in accreting objects?

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$T_e \sim 20$ eV, $n_e \sim 10^{18}$ cm $^{-3}$



White Dwarf Line-Shapes



Question:

Why doesn't spectral fitting provide the correct properties for White Dwarfs?

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$T_e \sim 1$ eV, $n_e \sim 10^{17}$ cm $^{-3}$



ZAPP campaigns simultaneously study multiple issues spanning 200x in temperature and 10^6 x in density



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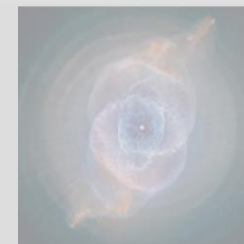
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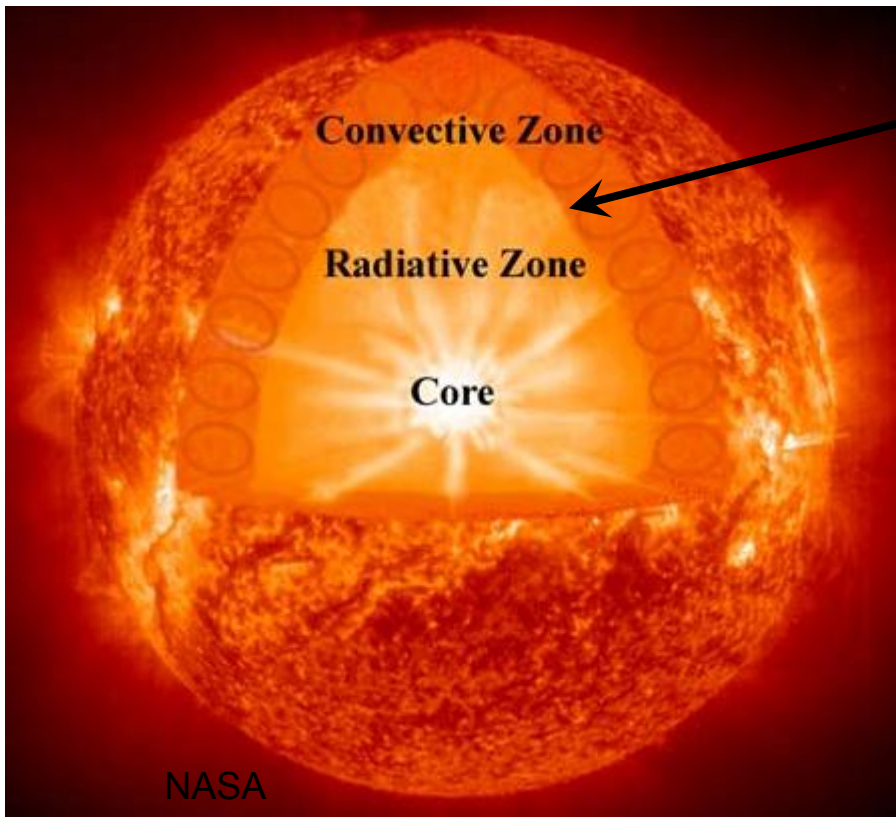
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Models for solar interior structure disagree with helioseismology observations.



Convection-Zone (CZ) Boundary
Models are off by 10-30 σ

Models depend on:

- Composition (revised in 2005*)
- EOS as a function of radius
- The solar matter *opacity*
- Nuclear cross sections

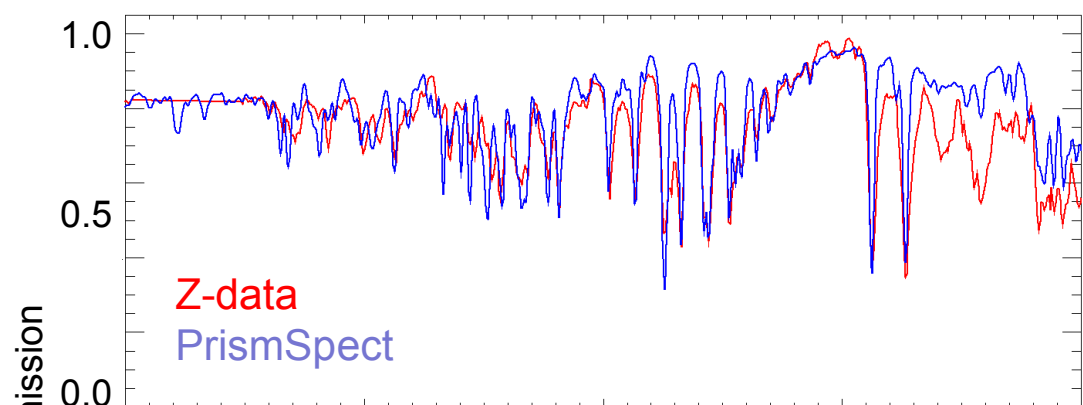
Question: Is opacity uncertainty the cause of the disagreement?

Objective: Measure Fe opacity at CZ base conditions.

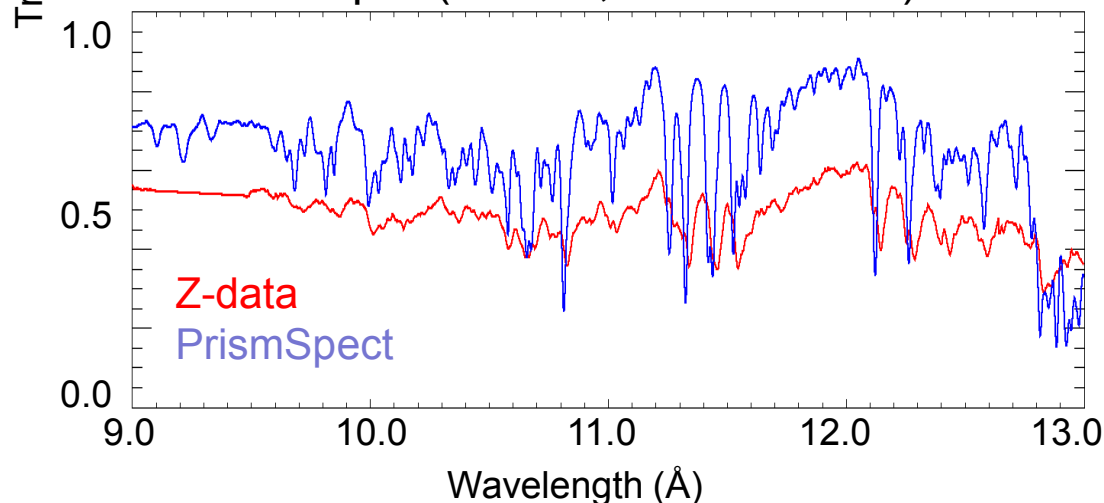
*M. Asplund *et al*, Annu. Rev. Astro. Astrophys. **43**, 481 (2005).

Modern computations of Fe opacity show large disagreements with data at CZ base conditions

Thin Tamper* (156 eV, $6.9 \times 10^{21} \text{ cm}^{-3}$)



Thick Tamper (182 eV, $31 \times 10^{21} \text{ cm}^{-3}$)



Present Status

- Agreement between data and computation becomes worse at increasing temp. and dens.
- Disagreements at CZ base conditions can partially explain the CZ boundary problem.
- The differences are probably not unique to Fe... more scrutiny of the data is prudent.

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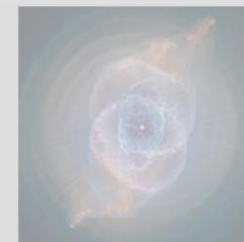
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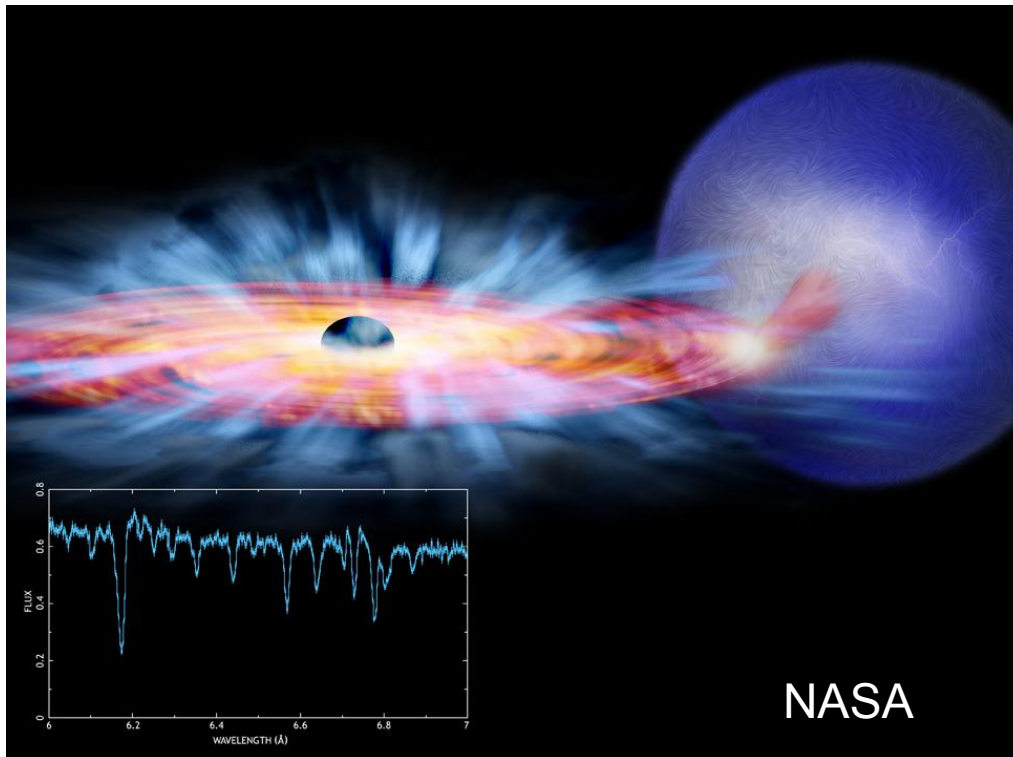
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We learn about black holes from the matter falling into them – these are photoionized plasmas

Conceptual Picture of a Black-Hole Accretion Disk
 $\xi \sim 10 - 10,000 \text{ erg.cm.s}^{-1}$



Photoionization parameter

$$\xi \equiv \frac{4\pi F}{n_e} [\text{erg.cm.s}^{-1}]$$

Laboratory Plasmas

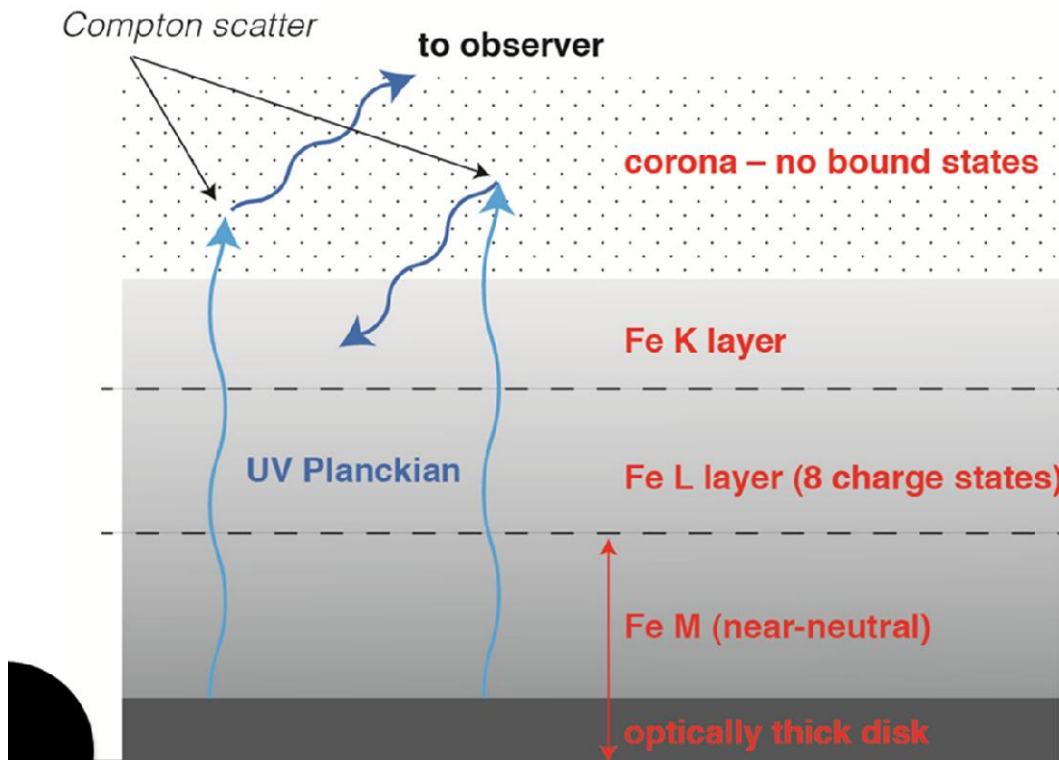
$$n_e \sim 10^{19} \text{ cm}^{-3}$$

$$F > 1 \text{ TW/cm}^2 \text{ for } \xi > 10$$

- Can we model the ionization?
- Can we model the line emission?

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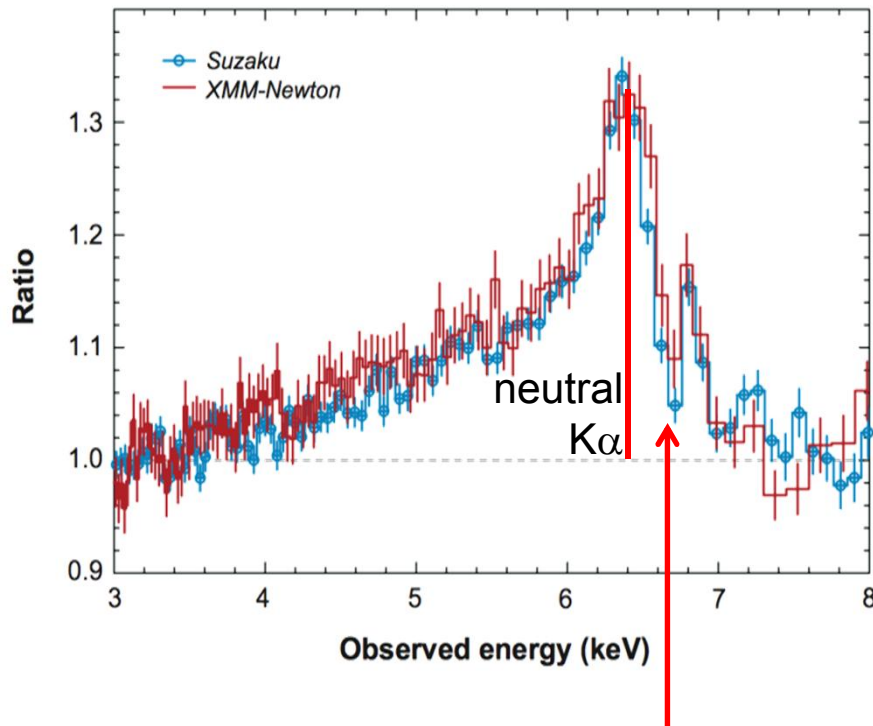
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A Specific Problem: Emission from L-shell ions is not seen in some prominent black-hole accretion disks.

Measured Fe Emission from MCG 6-30-15

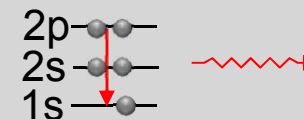


No observed emission from Fe ionized to the L-shell

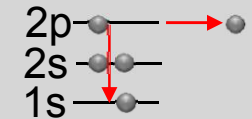
Resonant Auger Destruction (RAD) was accepted as the reason*

- 2 competing processes for the de-excitation of L-shell ions:

Radiative Decay



Auger Decay

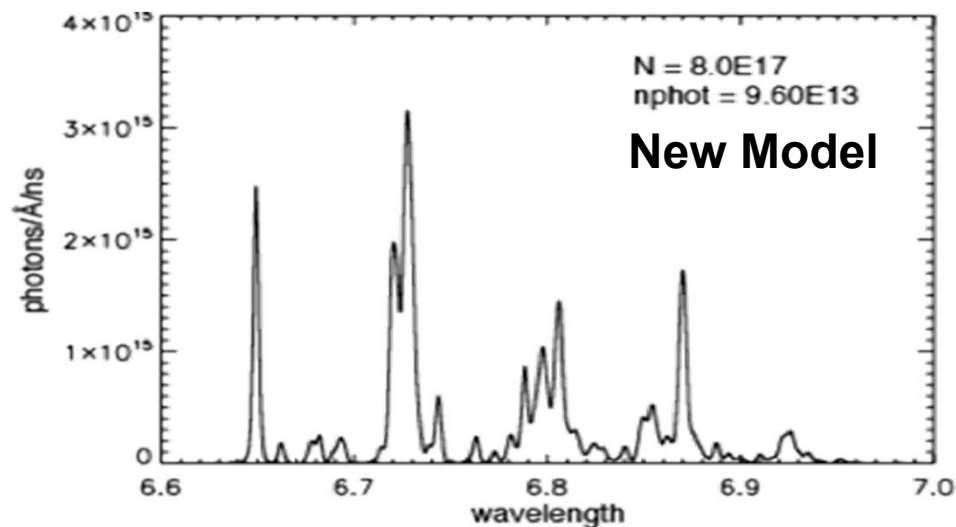
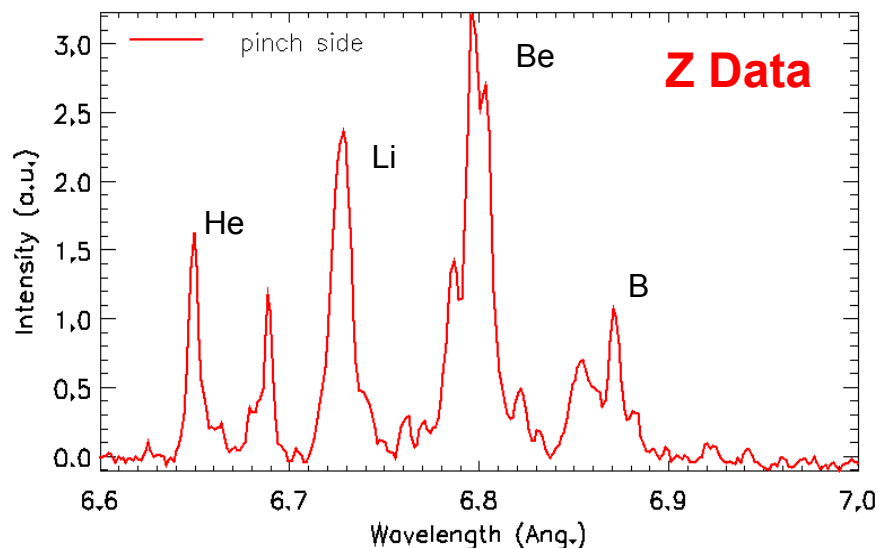


- Thin Plasma:** high probability of observing the photon
- Thick Plasma:** high probability of the photon being resonantly absorbed
→ Higher probability of Auger Decay for the ensemble

Recent emission measurements demonstrate that L-shell emission is not 100% quenched by RAD.

Present Status

- Z Data demonstrates that L-shell emission does escape at column depths $>1\text{E}17$ at/cm².
- Present data can discriminate between models of the ionization distribution AND relative line strengths.
- Absolute intensity is needed to determine efficiency of RAD process.



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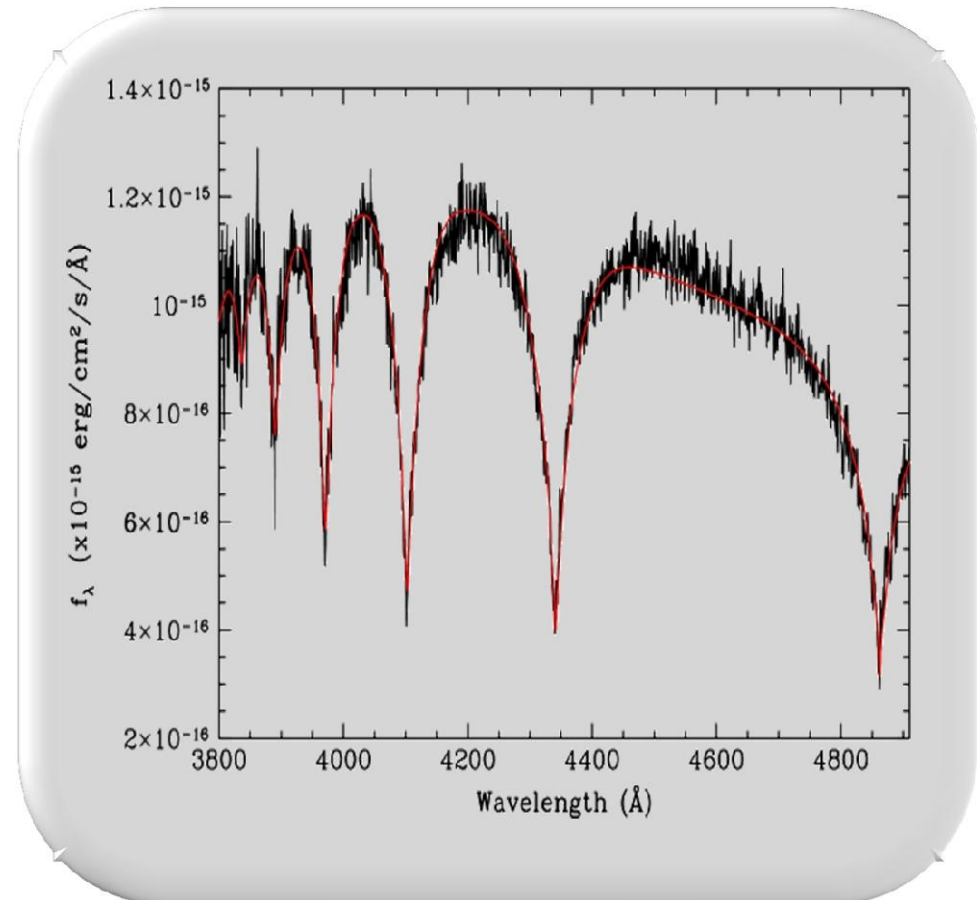
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The properties of White Dwarfs are determined by spectral fitting, but disagrees with other methods

Spectral fit of WD J1916+3938*

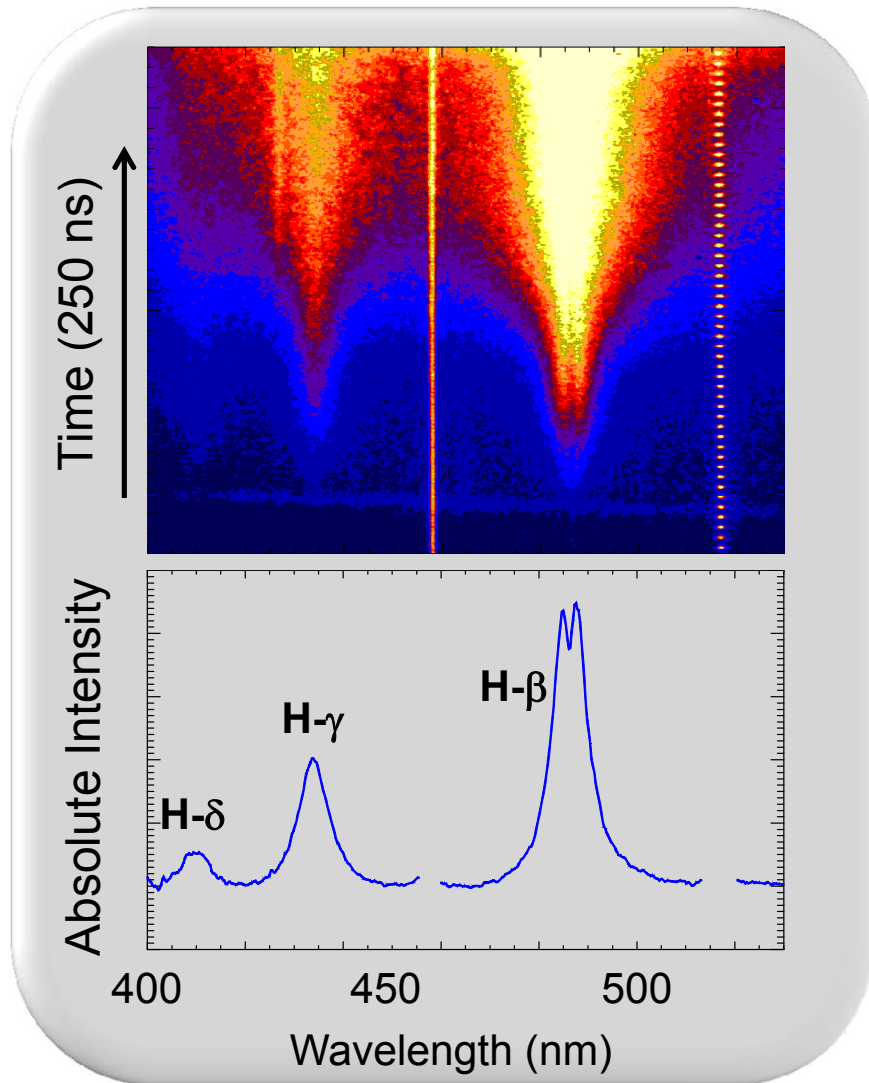
- White Dwarfs are fundamentally important
 - Evolutionary endpoint for ~98% of stars
 - Simple in structure and evolution
 - Cosmic laboratories (cosmochronology)
- WD surface temperature and total mass are usually determined by fitting the observed spectra
- The spectroscopic method and gravitational redshift disagree by >10% in the stellar mass



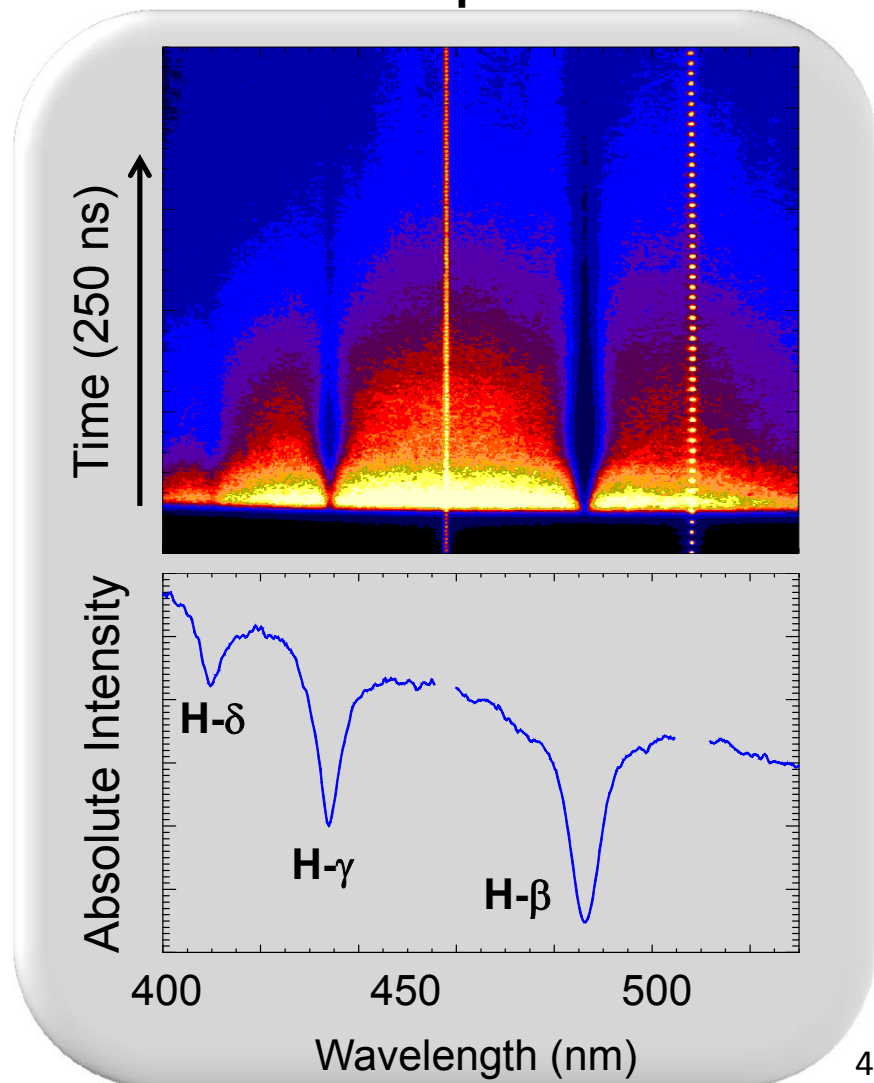
*Hermes et al. (2011)

Simultaneous streaked absorption and emission provide a unique capability to measure lineshapes

Emission



Absorption



Outline

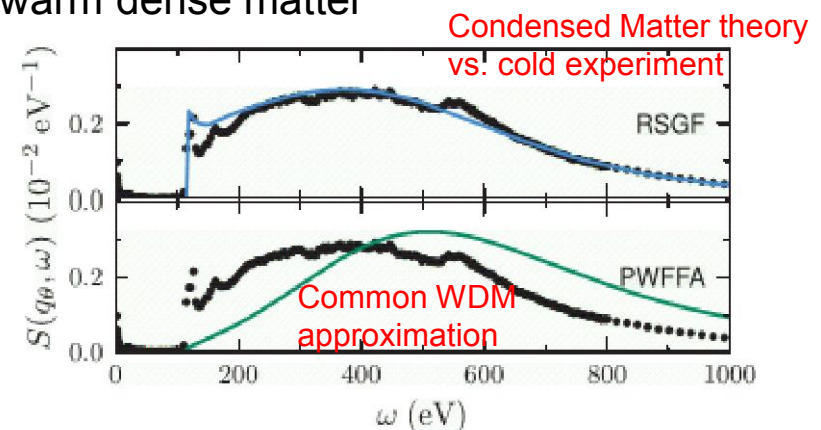
- **Magnetized Liner Inertial Fusion (MagLIF)**
Encouraging results from initial integrated experiments
- **Z Astrophysical Plasma Properties (ZAPP) Collaboration**
Studies of emission and absorption from photoionized plasmas relevant to white dwarfs, accretion disks, and solar photosphere
- **Diagnostic development on Z**
X-ray Thomson scattering and Zeeman splitting

There is a lot of interesting science going on at Z –
and many opportunities for significant contributions from Cornell's LPS.

X-ray Thomson Scattering has recently emerged as a potential diagnostic for warm dense matter (WDM)



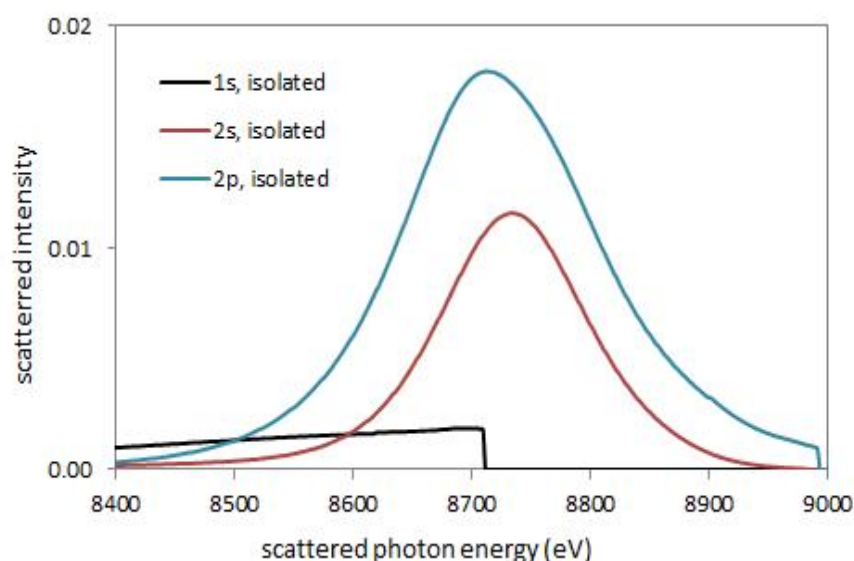
- Warm Dense Matter is difficult to both model and diagnose:
 - Since the kinetic energy of the plasma particles is of the same order as their potential energy, one cannot make the simplifying assumptions appropriate for Condensed Matter or Ideal Plasmas
 - Since Warm Dense Matter emits weakly and absorbs strongly, optical diagnostics only probe its surface; an external x-ray source is required for bulk characterization
- Seven recent PRLs have been devoted to the use of XRTS as a diagnostic for electron density, temperature, and ionization
 - However, the models used to predict scattering signals make significant assumptions about the electronic and ionic properties of warm dense matter
- Experimental studies with uniform, well-characterized sample conditions and high signal-to-noise scattering measurements are needed to augment existing cold scattering data



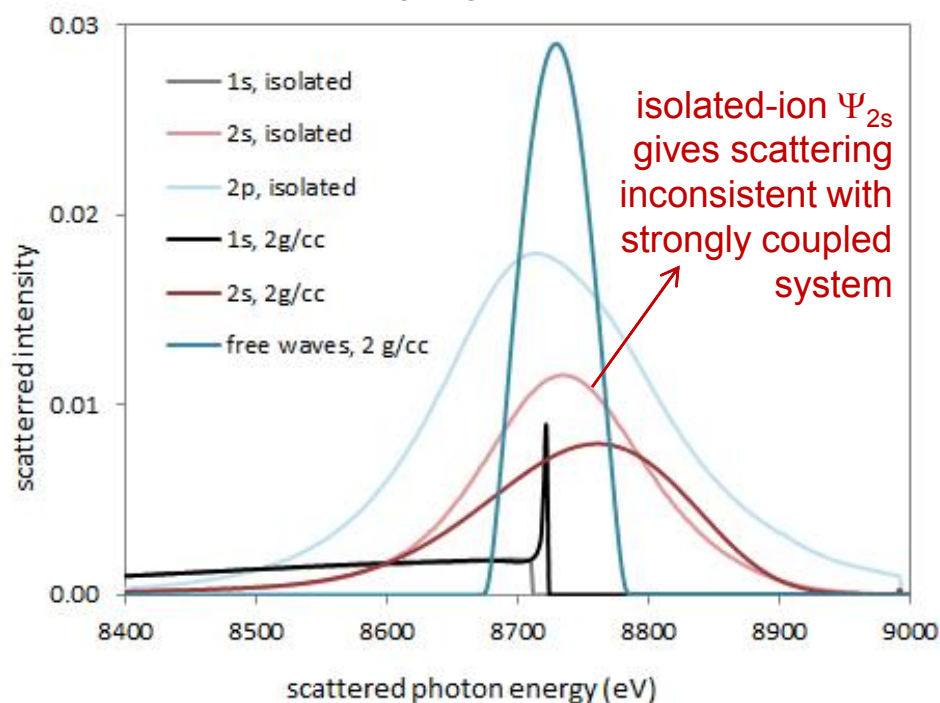
B.A. Mattern & G.T. Seidler,
Phys Plas **20**, 022706 (2013)

Key question: what happens to the scattering signal as bound electrons become pressure ionized?

135° scattering from cold, isolated carbon:
Valence 2p state is bound



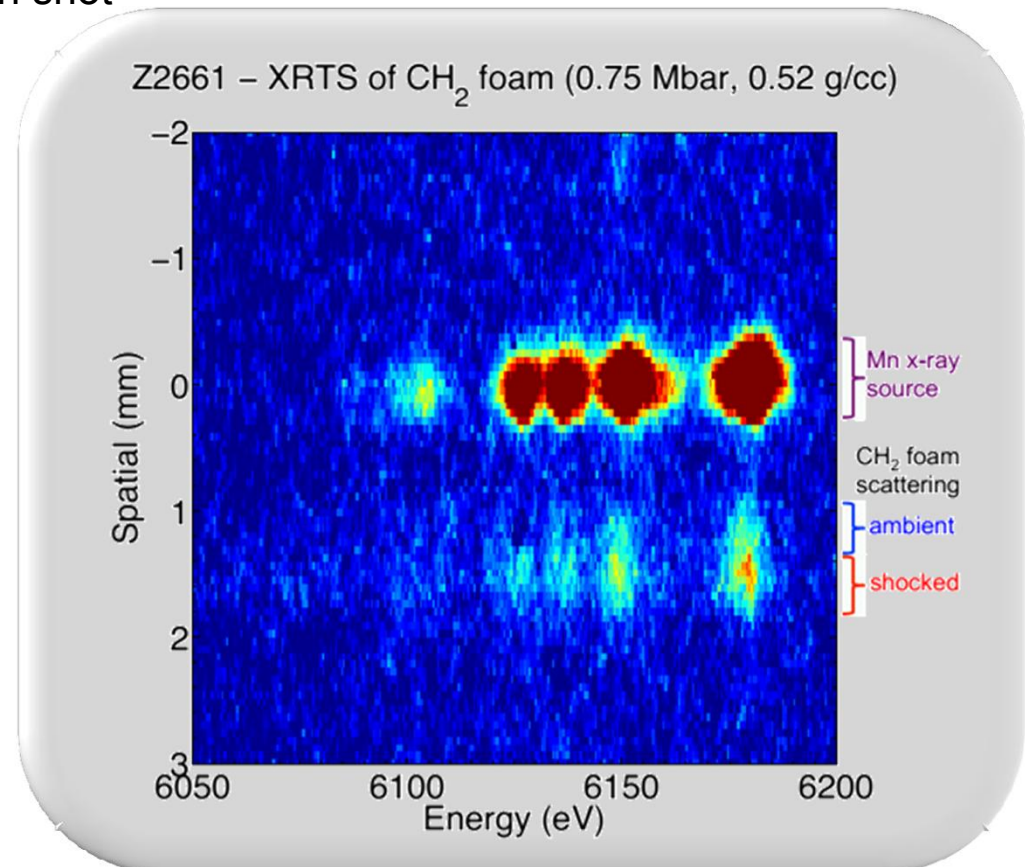
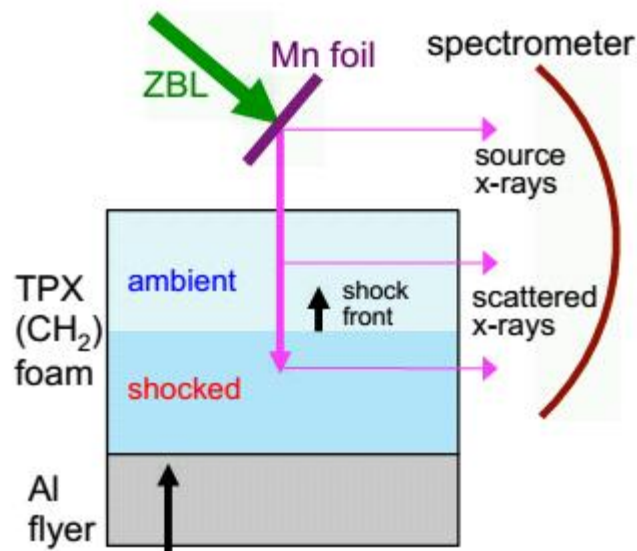
Scattering from cold carbon at 2 g/cc:
2p is pressure ionized (but not to a plane wave!)
1s and 2s scattering signals are modified



Laura Johnson is investigating the effect of continuum-wave distortion on free-free scattering signals.

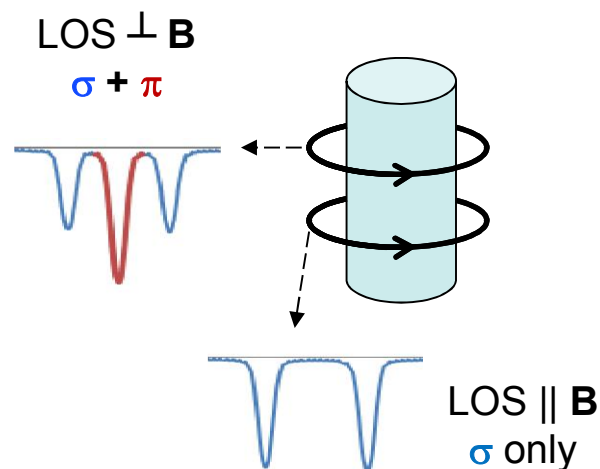
We recently obtained scattering data from both shocked and unshocked foam on Z

- Z experiments provide:
 - Uniform, long-lived, well-defined shock state
 - *In-situ* comparison with ambient state
 - Source spectrum measured for each shot
 - Potential for high S/N

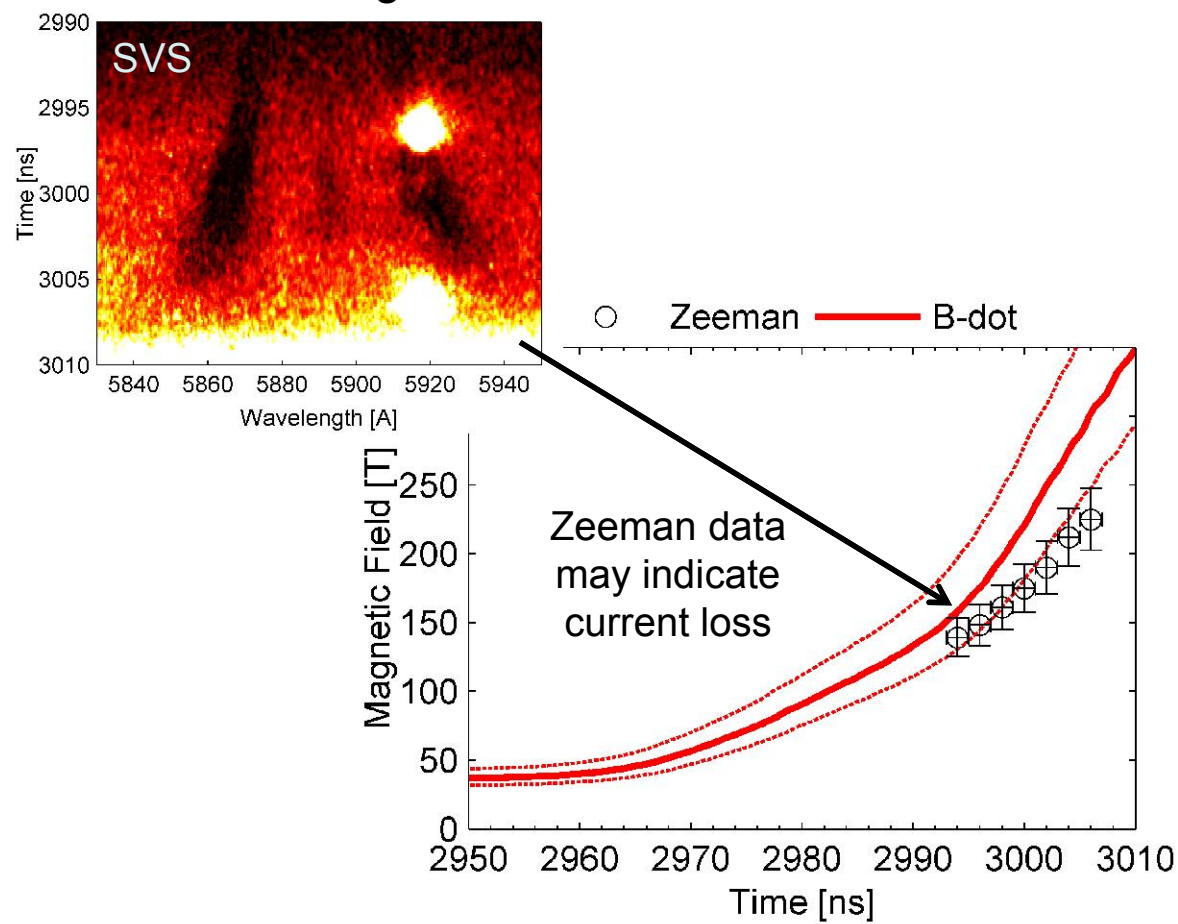


Zeeman splitting is being used to characterize Z's current drive and flux compression in Magnetized Liner Inertial Fusion (MagLIF) experiments

- Sodium deposits vaporized and backlit by current-carrying surfaces signal both the magnitude and direction of the local magnetic field:



The relative strength of σ and π components indicates field direction



Summary: We are using Z to explore new frontiers in laboratory astrophysics and diagnostics



- The Z Astrophysical Plasma Properties (ZAPP) collaboration uses the intense x-ray flux from wire arrays to create mm-scale regions of material with properties similar to those found in accretion objects, white dwarfs, and the solar photosphere, and measures their properties with extensive x-ray and optical spectrometers
- X-ray spectrometers are being fielded on dynamic materials experiments to help advance understanding of x-ray scattering in warm dense matter, with the potential to provide critical tests of scattering theory and temperature diagnostics
- Streaked optical spectrometers are being used to measure Zeeman splitting of optical absorption features to diagnose local magnetic fields important for understanding flux compression and current loss

We welcome ideas for additional collaborations!