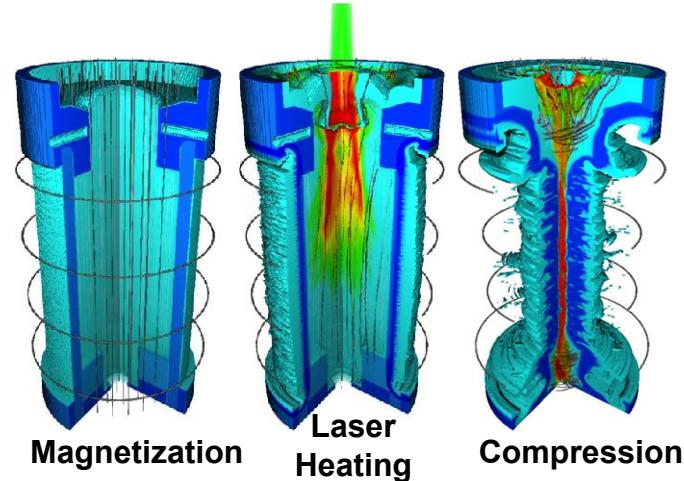
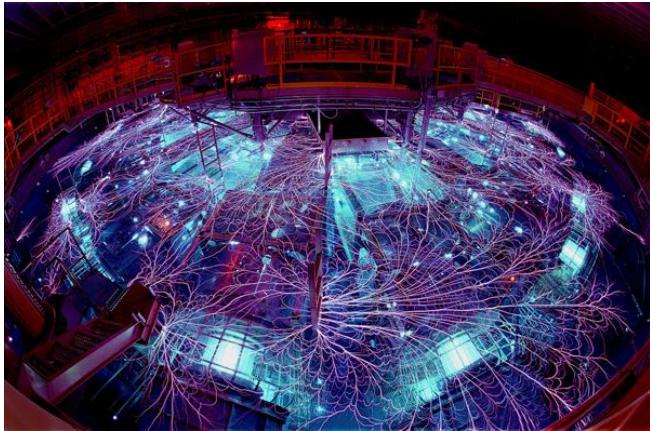


*Exceptional service in the national interest*



# An overview of fusion, astrophysical experiments, and diagnostic development on Sandia's Z machine

Stephanie B. Hansen

*Sandia National Laboratories*

*Presentation to the Laboratory for Plasma Studies  
August 27, 2014*



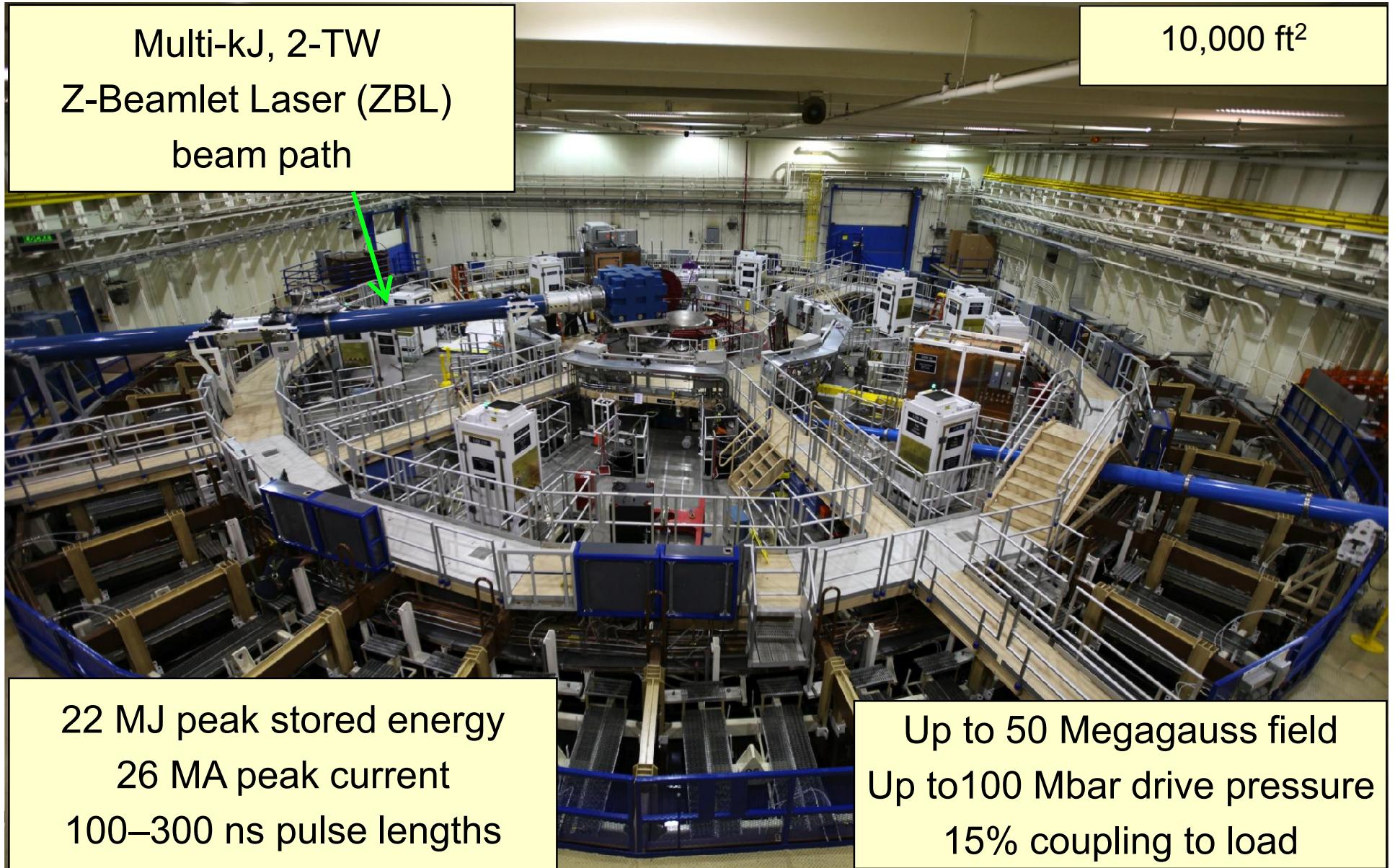
Sandia National Laboratories is a multi-program laboratory managed and operated by Sandia Corporation, a wholly owned subsidiary of Lockheed Martin Corporation, for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-AC04-94AL85000.

# Outline

- **Magnetized Liner Inertial Fusion (MagLIF)**  
Encouraging results from initial integrated experiments
- **Z Astrophysical Plasma Properties (ZAPP) Collaboration**  
Studies of emission and absorption from photoionized plasmas relevant to white dwarfs, accretion disks, and solar photosphere
- **Diagnostic development on Z**  
X-ray Thomson scattering and Zeeman splitting

There is a lot of interesting science going on at Z –  
and many opportunities for significant contributions from Cornell's LPS.

The Sandia Z pulsed power facility uses magnetic pressure to efficiently couple energy to drive relatively large targets for a wide variety of stockpile stewardship applications



Multi-kJ, 2-TW

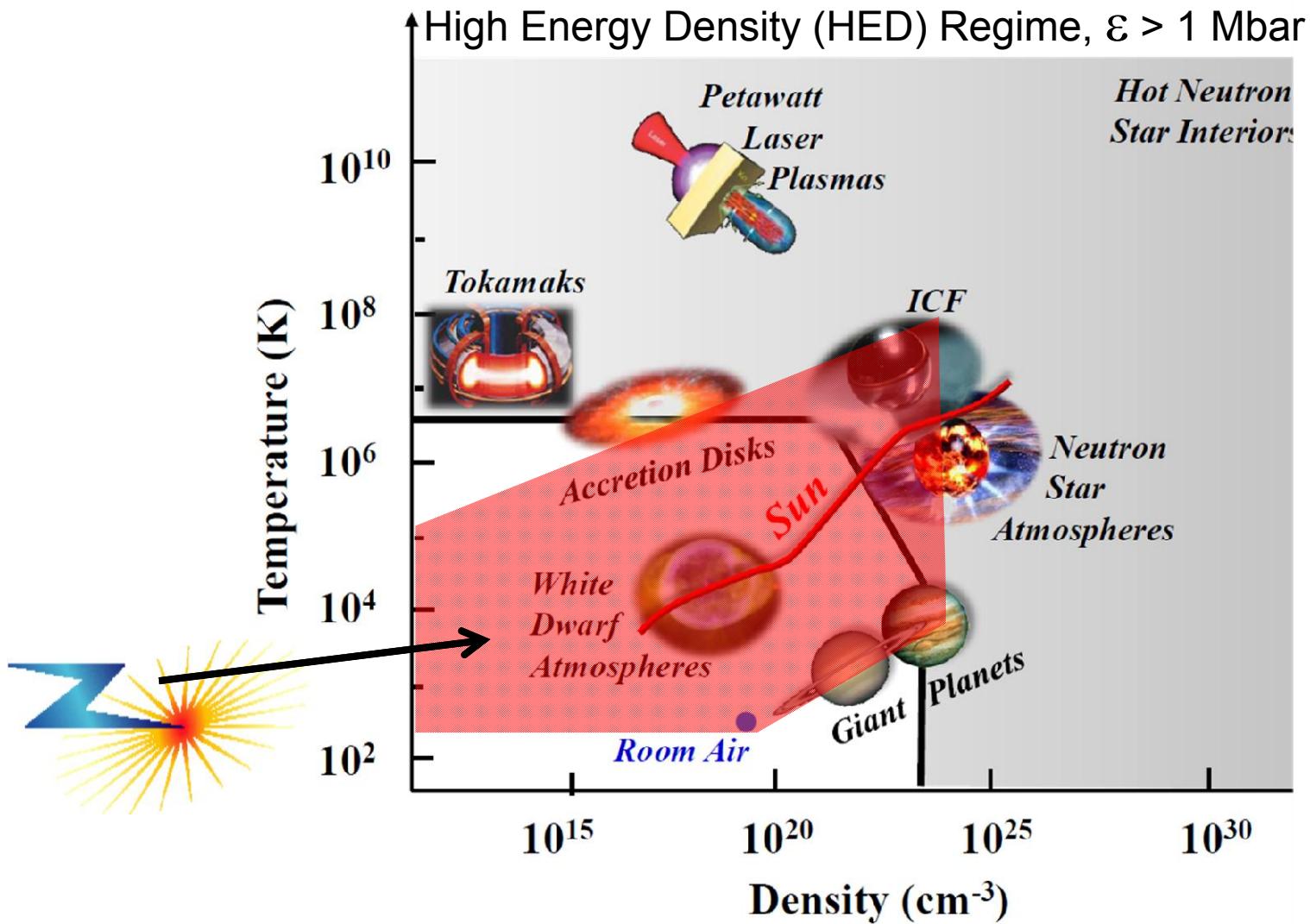
Z-Beamlet Laser (ZBL)  
beam path

10,000 ft<sup>2</sup>

22 MJ peak stored energy  
26 MA peak current  
100–300 ns pulse lengths

Up to 50 Megagauss field  
Up to 100 Mbar drive pressure  
15% coupling to load

# Experiments on Z access a broad range of the energy-density phase space



# Many people are contributing to our Magnetized Liner Inertial Fusion (MagLIF) effort:



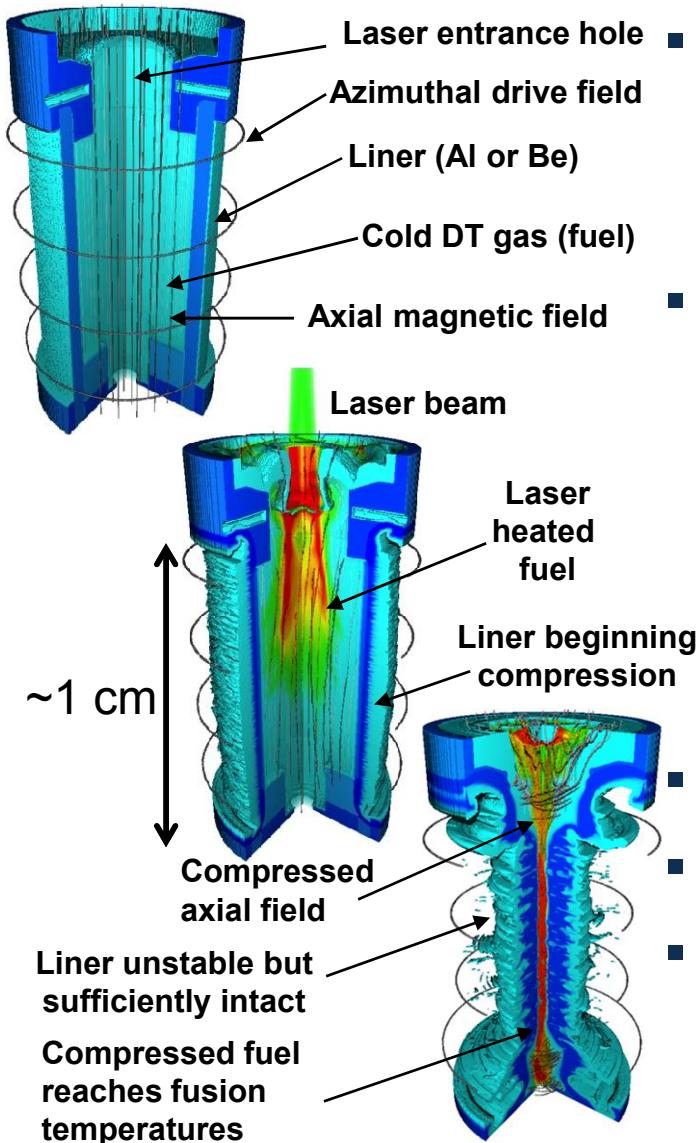
T.J. Awe, C.J. Bourdon, G.A. Chandler, P.J. Christenson, M.E. Cuneo, M. Geissel, **M.R. Gomez**, K.D. Hahn, S.B. Hansen, E.C. Harding, A.J. Harvey-Thompson, M.C. Herrmann, M.H. Hess, C.A. Jennings, B. Jones, M. Jones, R.J. Kaye, P.F. Knapp, D.C. Lamppa, M.R. Lopez, M.R. Martin, R.D. McBride, L.A. McPherson, J.S. Lash, K.J. Peterson, J.L. Porter, G.A. Rochau, D.C. Rovang, C.L. Ruiz, S.E. Rosenthal, M.E. Savage, P.F. Schmit, **A.B. Sefkow**, **D.B. Sinars**, **S.A. Slutz**, I.C. Smith, W.A. Stygar, R.A. Vesey, E.P. Yu

*Sandia National Laboratories, Albuquerque, NM 87185 USA*

B.E. Blue, D.G. Schroen, K. Tomlinson

*General Atomics, San Diego, CA 92186 USA*

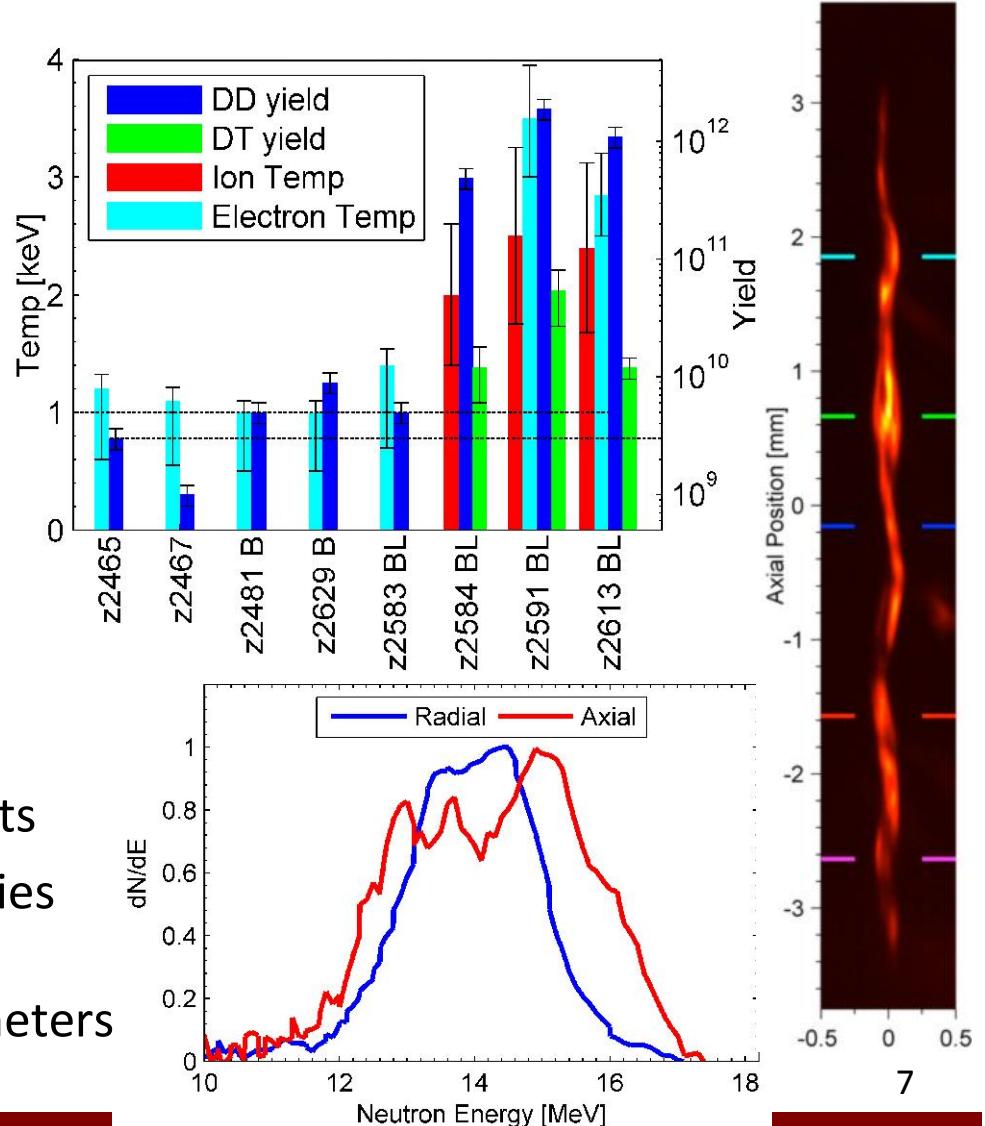
# We are evaluating a **Magnetized Liner Inertial Fusion (MagLIF)\*** concept that may reduce fusion requirements



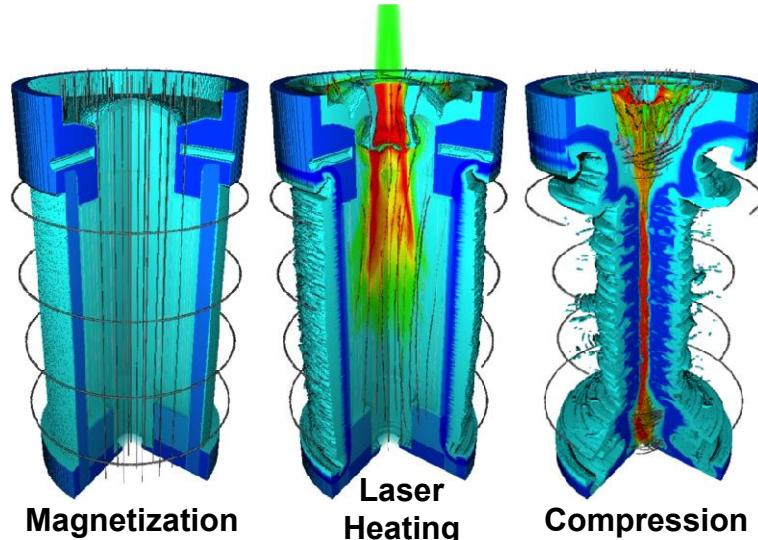
- An initial 30 T axial magnetic field is applied
  - Inhibits thermal conduction losses
  - Appears to stabilize implosion at late times
- During the  $\sim 100$  ns implosion, the fuel is heated using the Z-Beamlet laser (about 6 kJ in designs)
  - Preheating to  $\sim 300$  eV reduces the compression needed to obtain fusion temperatures to 23 on Z
  - Preheating reduces the implosion velocity needed to  $\sim 100$  km/s, allowing us to use thick liners that are more robust against instabilities
- $\sim 50$ - $250$  kJ energy in fuel; 0.2-1.4% of capacitor bank
- Stagnation pressure required is  $\sim 5$  Gbar
- DD equivalent of 100 kJ DT yield may be possible on Z in the next few years—this will require enhanced drive upgrades that are in progress, e.g.,  $10$  T  $\rightarrow$   $30$  T;  $2$  kJ  $\rightarrow$   $>6$  kJ;  $19$  MA  $\rightarrow$   $>24$  MA

# We obtained promising initial results with MagLIF and are seeking to increase our understanding

- We achieved DD yields up to  $2 \times 10^{12}$  ( $\sim 0.3$  kJ DT equivalent) in our first integrated tests of Magnetized Liner Inertial Fusion (MagLIF)
- A variety of data were collected that appear to show a  $< 150$   $\mu\text{m}$  diameter,  $\sim 3$  keV, highly magnetized plasma was produced—remarkable for a 70-100 km/s implosion!
- We are continuing to build on these results with a balanced combination of focused and integrated experiments
- In parallel we are improving capabilities to understand how this performance will scale with increasing drive parameters



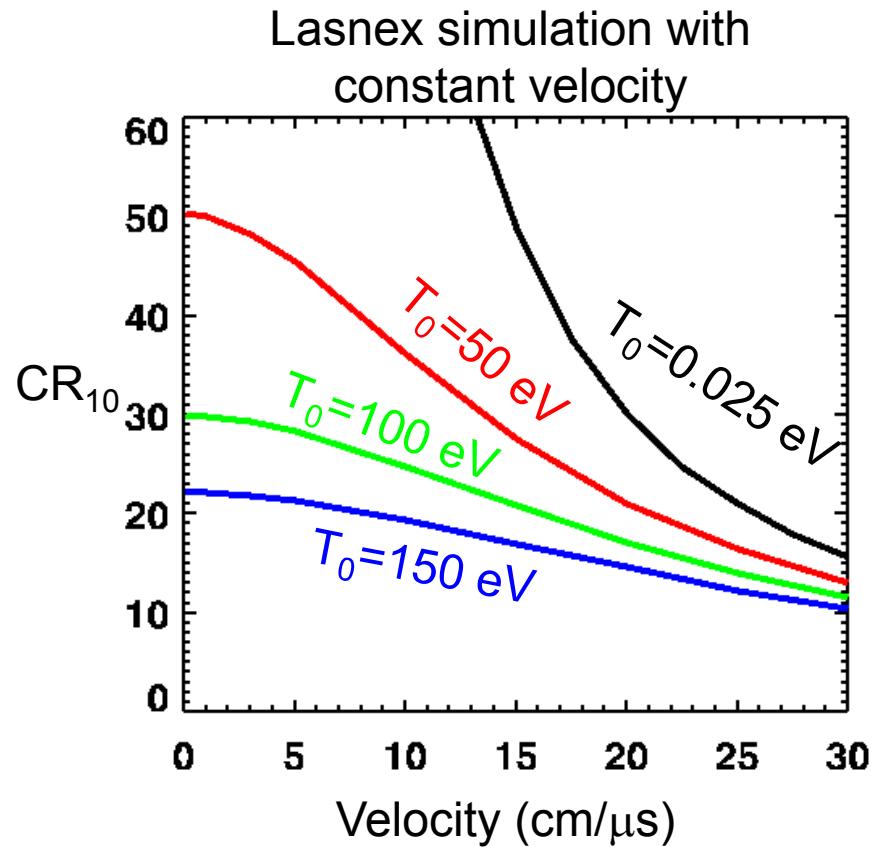
# Our path to studying the underlying science is a mixture of focused and integrated experiments to address key physics



- Key physics uncertainties
  - Liner instabilities
    - Electro-thermal
    - Magneto-Rayleigh-Taylor
    - Deceleration RT
    - Impact of 3D fuel assembly
  - Liner/fuel interactions & mix
  - Laser-window and laser-fuel scattering, absorption, uniformity
  - Suppression of electron heat transport in dense plasma by magnetic fields
  - Magnetic flux compression
  - Magnetized burn
- Key target design elements
  - Liner compression
  - Magnetization
  - Laser heating

Experiments to address the key physics are planned for the Z pulsed power facility and the Z-Beamlet and Omega (and -EP) lasers.

# Typical ICF implosions need high velocities to reach fusion temperatures—starting the implosion with heated fuel potentially reduces requirements



Heating fuel to ignition temperatures is typically done with a high-velocity shock (or series of shocks)

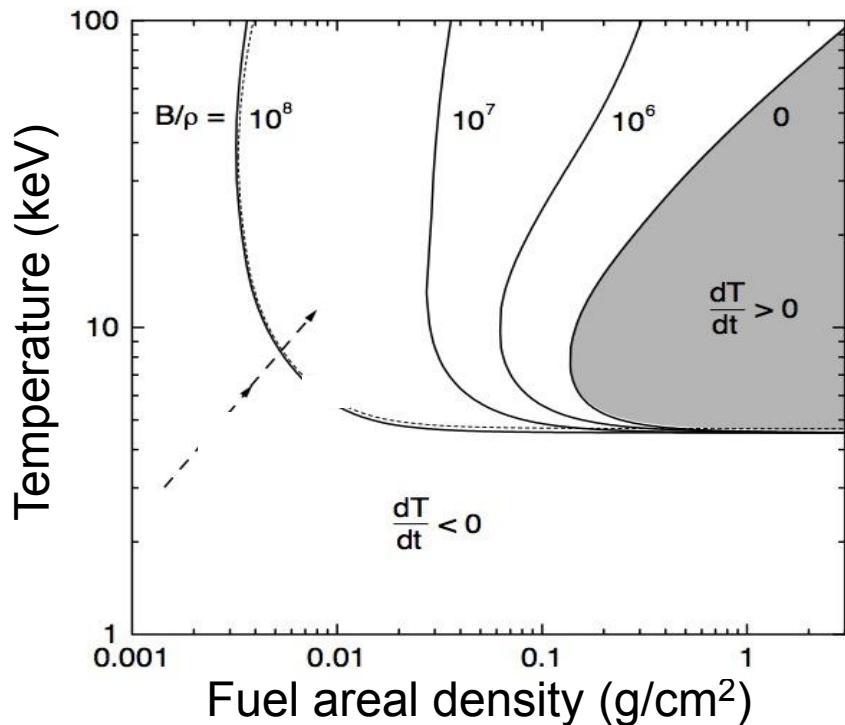
High velocities make it easier to reach fusion temperatures and also reduce the time available for losses (e.g., electron heat conduction or radiation)

Heating the fuel prior to the implosion *in the absence of losses* can allow low-velocity, low-convergence implosions to reach ignition temperatures

Is there a way to reduce losses?

# A large, embedded magnetic field can significantly reduce electron conduction losses from heated fuel

\*Basko et al. *Nuc. Fusion* 40, 59 (2000)



The  $\rho r$  needed for ignition can be significantly reduced by the presence of a strong magnetic field largely through inhibiting electron conduction

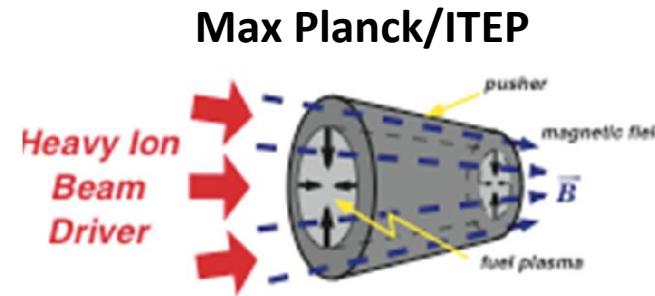
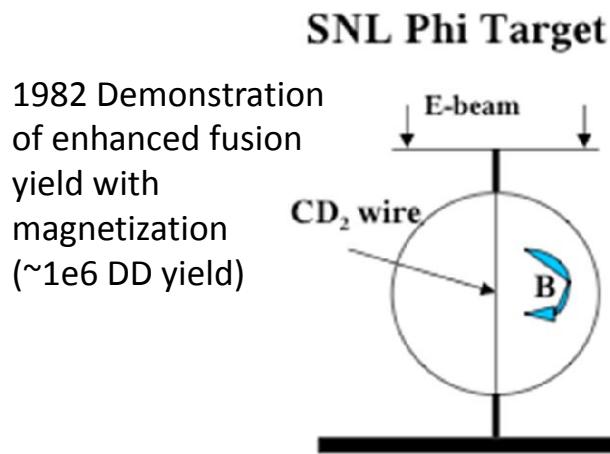
Lower  $\rho r$  reduces the required final fuel density (e.g.,  $\sim 1 \text{ g/cc} \ll 100 \text{ g/cc}$ ), which also reduces bremsstrahlung radiation losses

This means the stagnation plasma pressure at ignition temperatures is significantly reduced (e.g.,  $\sim 5 \text{ Gbar} \ll \sim 500 \text{ Gbar}$  for hot spot ignition)

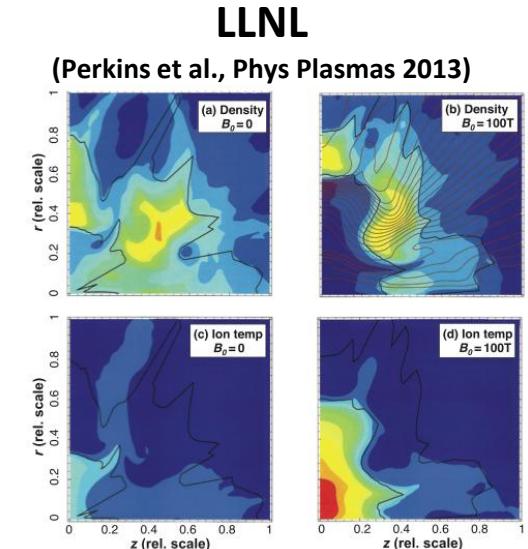
Large values of  $B/\rho$  are needed and therefore large values of  $B$  are needed,  $B \sim 10,000 \text{ Tesla}$   
(Earth's B-field is  $\sim 0.00003 \text{ Tesla}$ )

This field significantly exceeds pulsed coil technology ( $B_0 \sim 10-30 \text{ T}$ ), therefore flux compression is needed

# Many groups want to use magnetic fields to relax inertial fusion stagnation requirements

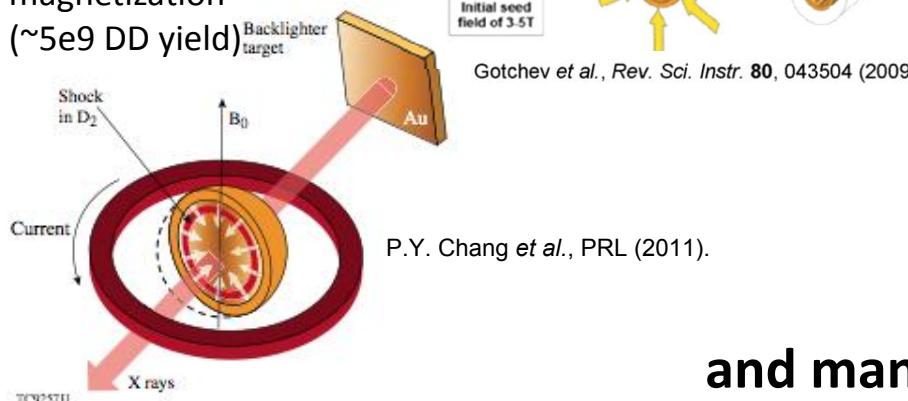


Basko, Kemp, Meyer-ter-Vehn, *Nucl. Fusion* **40**, 59 (2000)  
Kemp, Basko, Meyer-ter-Vehn, *Nucl. Fusion* **43**, 16 (2003)



## University of Rochester/LLE

2011 Demonstration of enhanced fusion yield with magnetization (~5e9 DD yield)



## Los Alamos/Air Force Research Lab

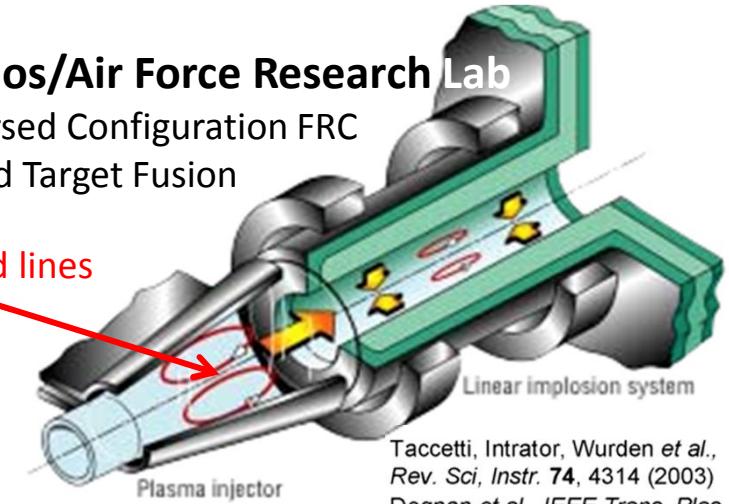
Field Reversed Configuration FRC

Magnetized Target Fusion

Shiva Star

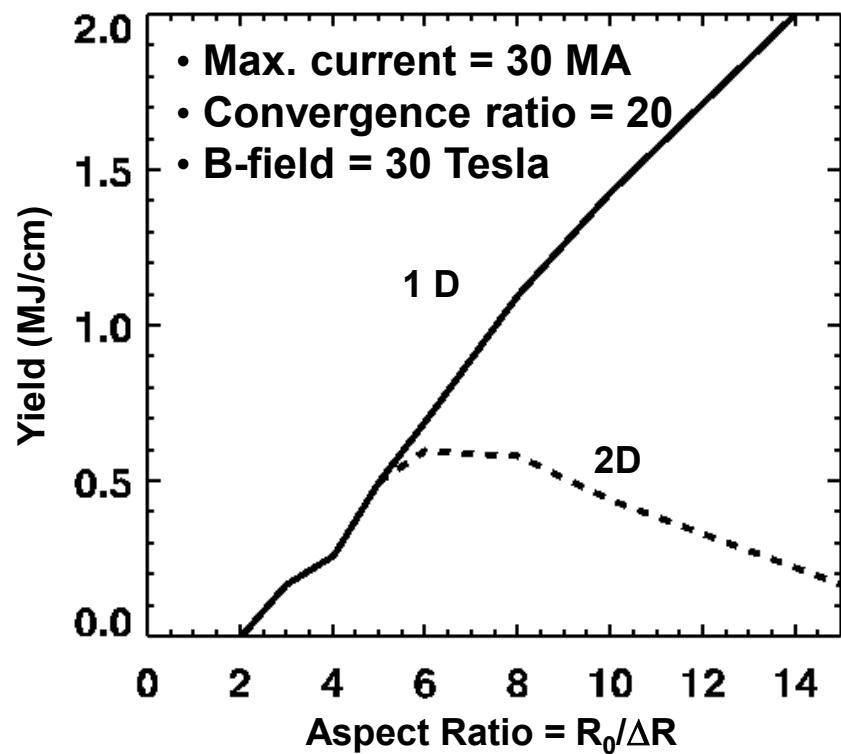
closed field lines

FRC

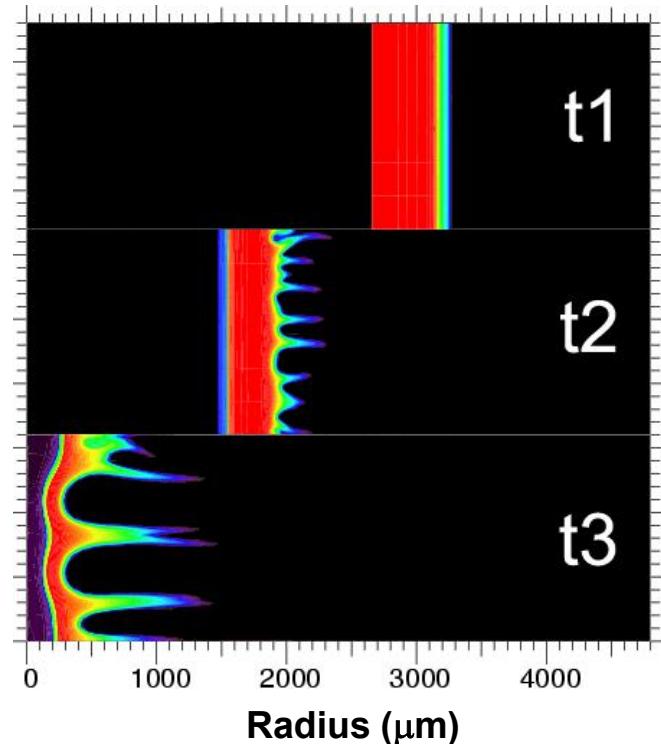


and many others...

# Reducing the implosion velocity requirements through fuel heating and magnetization allows us to use thicker, more massive liners to compress the fuel that are more stable



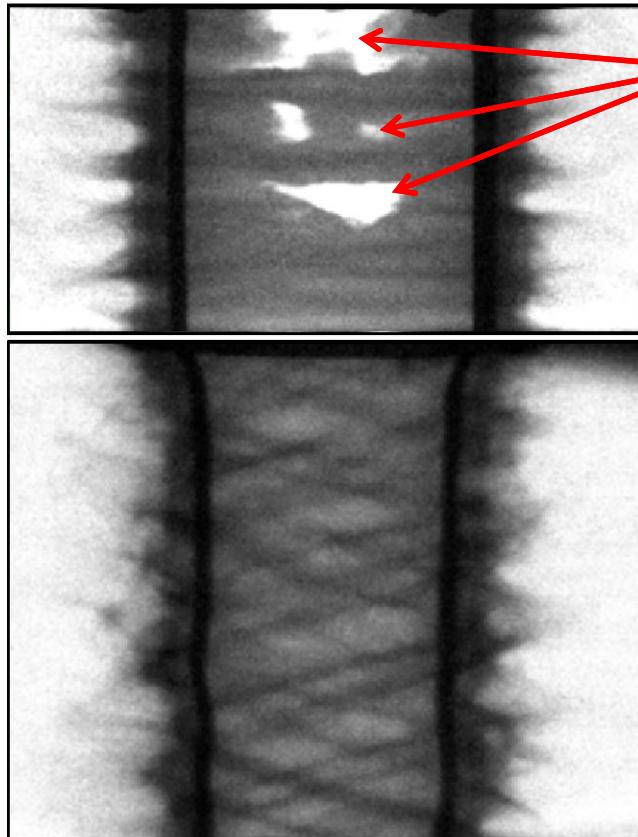
- The Magneto-Rayleigh-Taylor instability degrades the yield as the aspect ratio is increased (due to decreased liner  $\rho R$ )



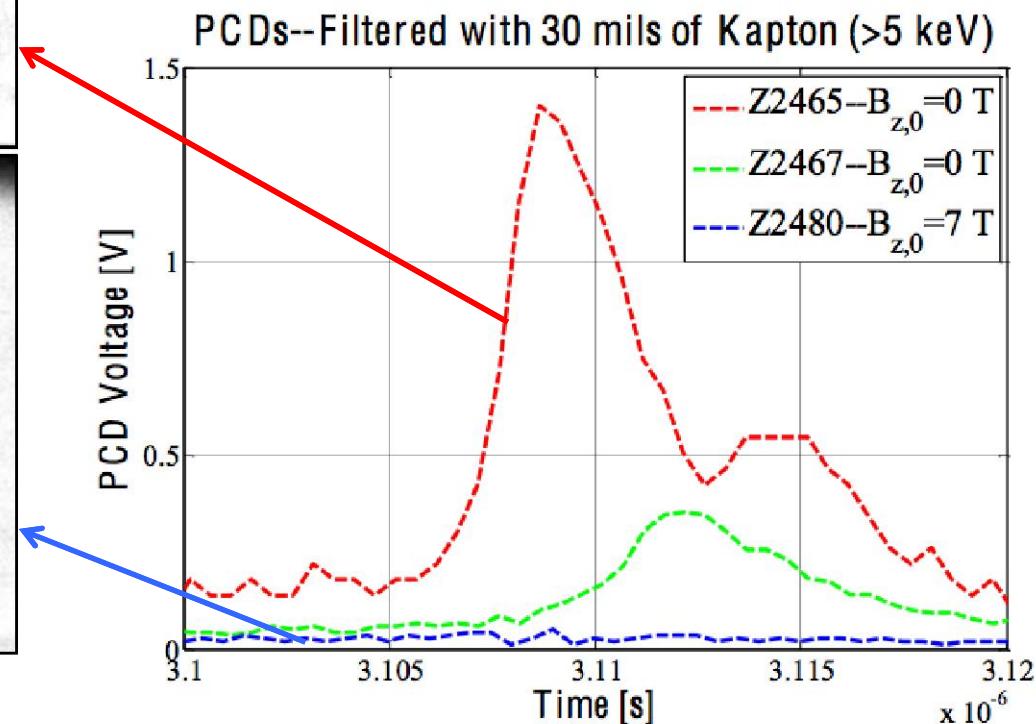
- Simulations of AR=6 Be liner show reasonably uniform fuel compression and sufficient liner  $\rho R$  at stagnation to inertially confine the fuel—important because fuel density is low!

Adding an axial magnetic field reduces hard x rays and hot spots, and changes the liner instability structure from cylindrical to helical—evidence it is doing something!

### Without Magnetic Field



Time-integrated self-emission from liner implosion at 6151 eV; missing in shots with axial field

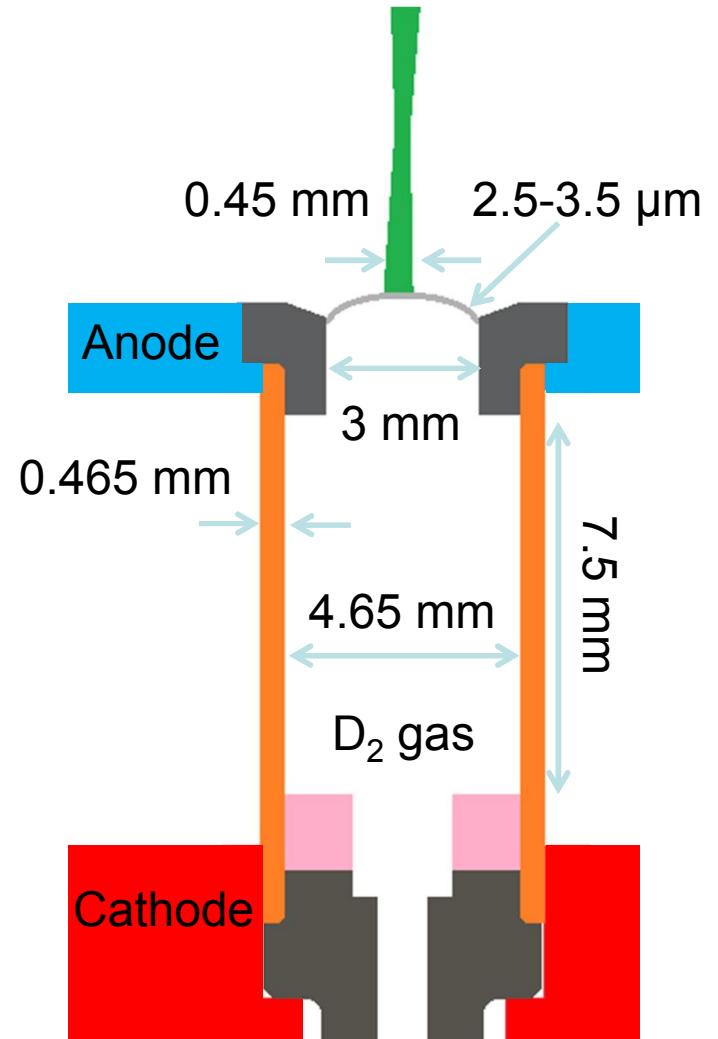


### With Magnetic Field

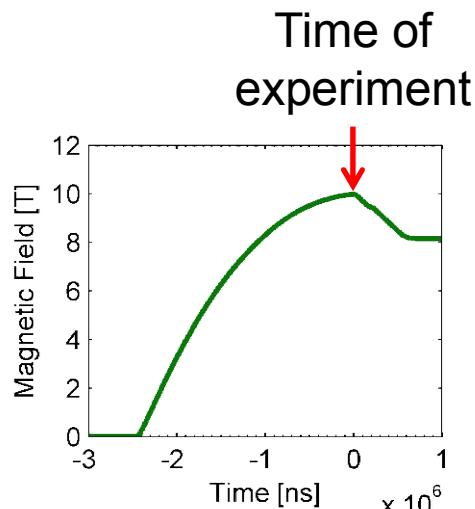
If magnetic flux roughly conserved the additional magnetic pressure from the axial field will suppress micro-pinching—this is indirect evidence for flux compression

# The target design for our initial experiments incorporates the knowledge gained from focused experiments and extensive simulations

- **Beryllium liner with aspect ratio 6**
  - Thick liner is more robust to instabilities
  - Still allows diagnostic access  $> 5$  keV
- **Top and bottom implosion cushions**
  - Mitigates wall instability
- **Standoff between LEH and imploding region**
  - Avoid window material mixing with fuel
- **Exit hole at bottom of target**
  - Avoid interaction with bottom of target

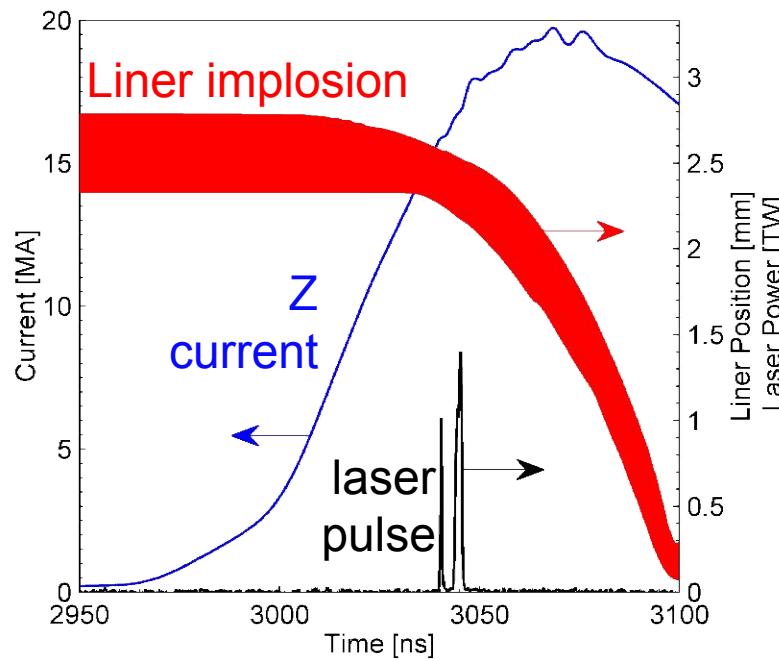


# Initial experiments were conducted at $I = 19 \text{ MA}$ , $B = 10 \text{ T}$ , and $\text{Laser} = 2.5 \text{ kJ}$



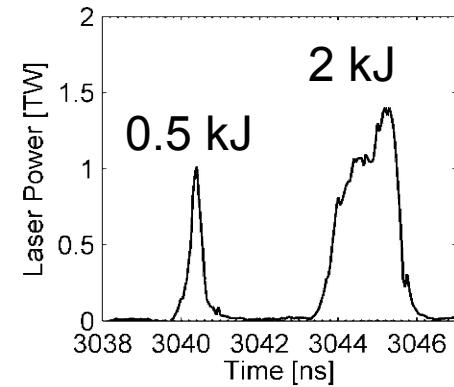
Magnetic field risetime is approximately 2 ms

$B$  is constant over the timescale of the experiment

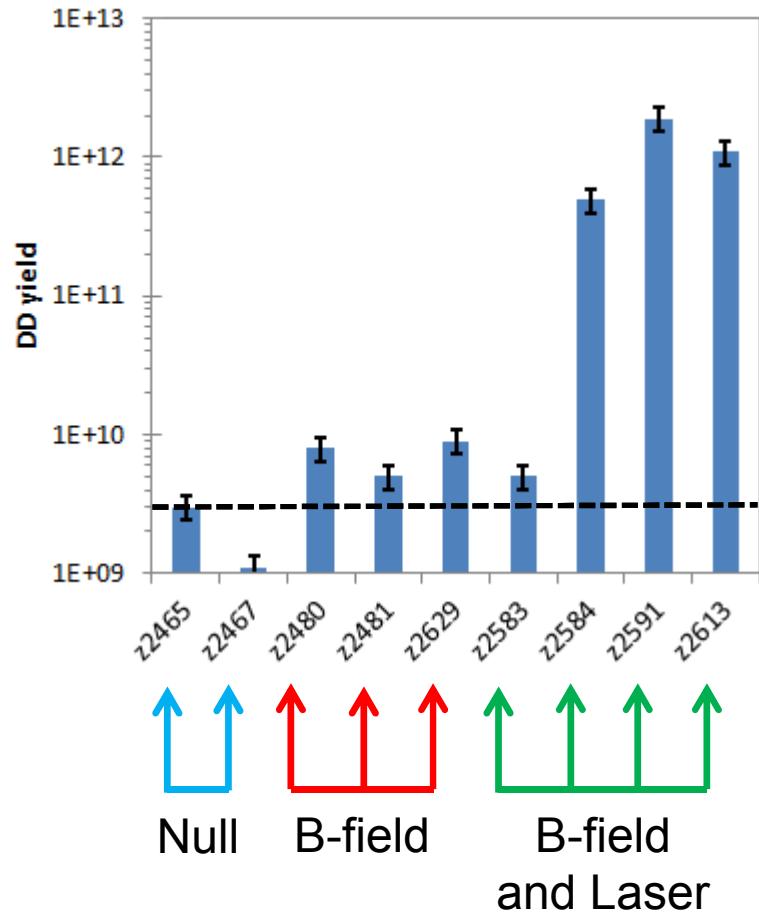


Peak current is 19 MA  
Magnetic field is 10 T  
Total laser energy is 2.5 kJ

Laser energy is split into 2 pulses:  
1<sup>st</sup> pulse intended to destroy LEH  
2<sup>nd</sup> pulse intended to heat fuel

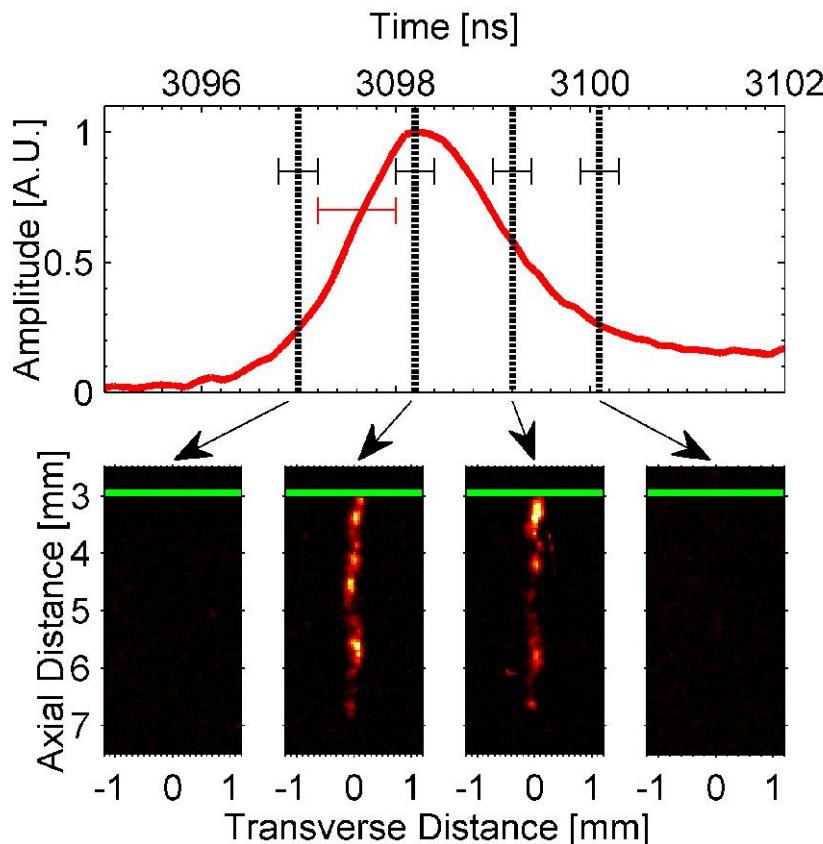


# Z shots producing DD yields in excess of $10^{12}$ were only observed in experiments with laser and B-field



- High yields were only observed on experiments incorporating both applied magnetic field and laser heating
- A series of experiments without laser and/or B-field produced yields at the background level of the measurement
- Result of z2583 is not well understood nor reproduced at this time

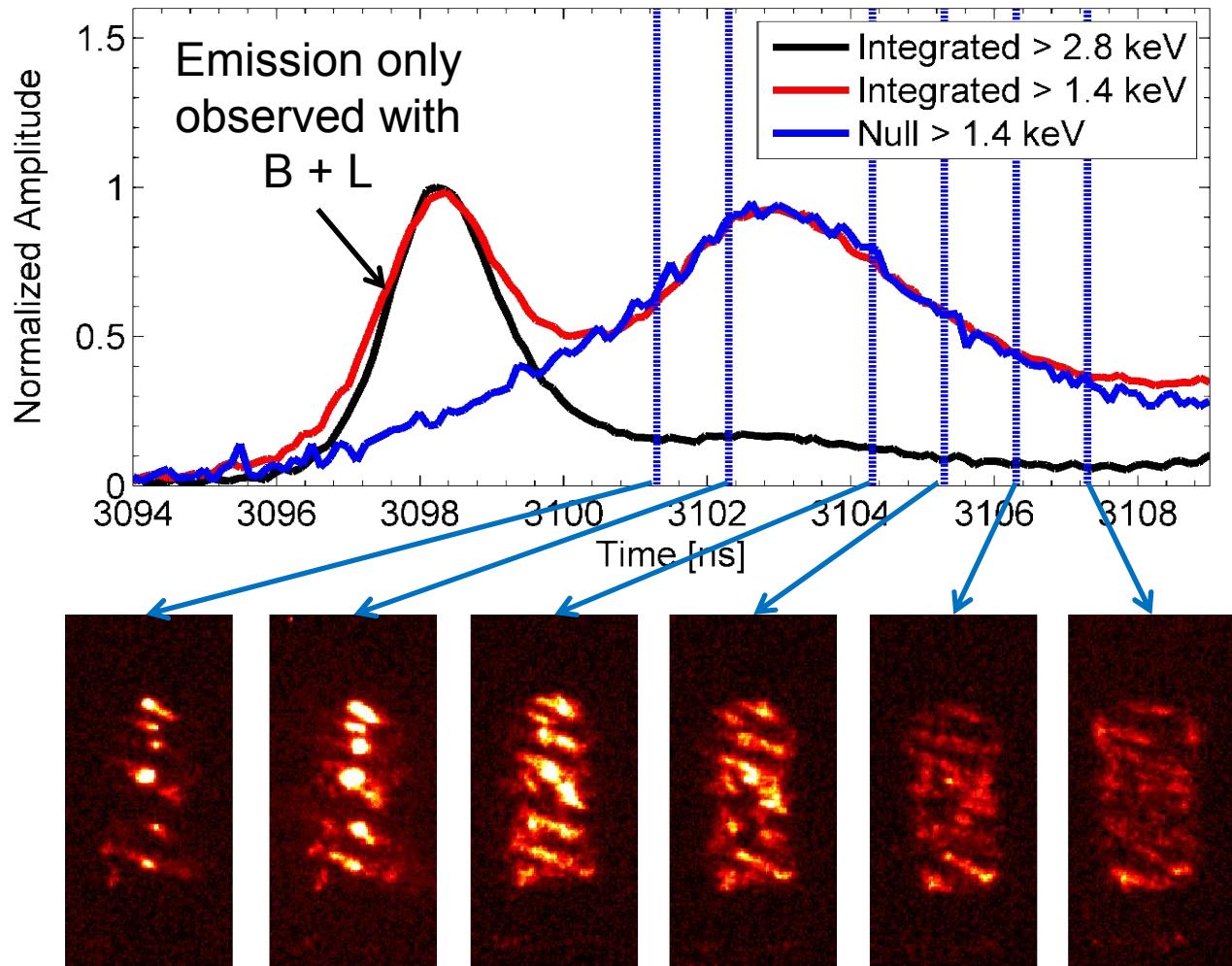
# Time-resolved x-ray pinhole imaging ( $h\nu > 2.8$ keV) shows a narrow emission column during peak in X-ray signal



- Emission column is observed only during the peak in the x-ray signal
- Emission column is only observed on experiments with high neutron yield
- Stagnation column width is at the resolution limit of this instrument ( $\sim 150$  microns)

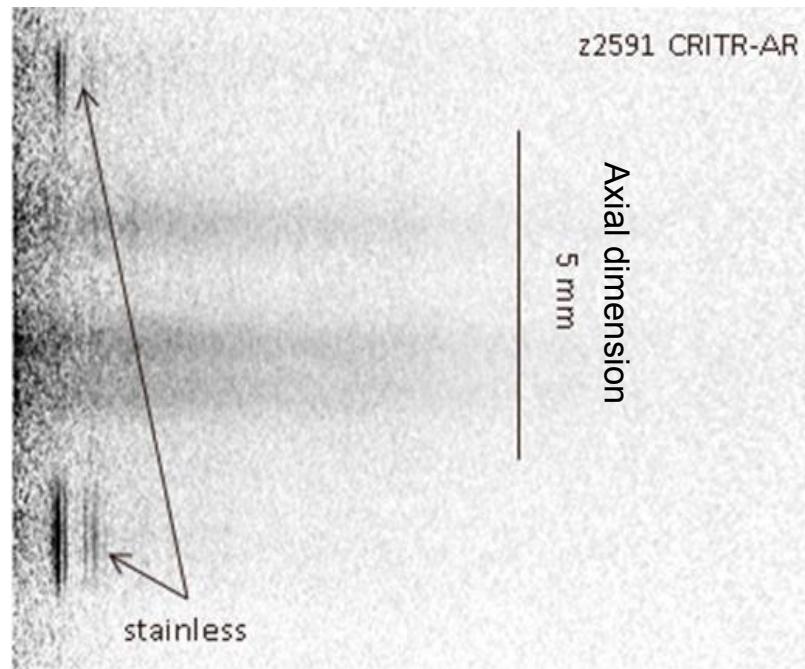
# High energy x-ray signal and narrow emission region are absent in null experiments

- Liner emission is observed in all experiments
- Liner emission is at a lower photon energy ( $< 2.8$  keV)
- Liner emission is getting larger at late times

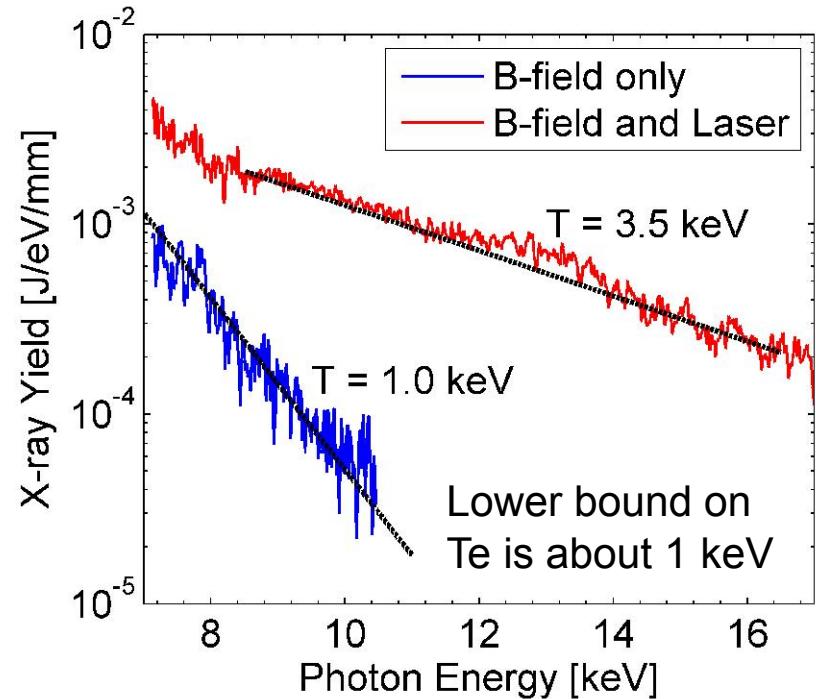


# High-energy spectra show axial variations in temperature and composition, with $\sim 3.5$ keV electron temperature in the pinch region—remarkable for a 70-100 km/s implosion!

Emission lines from stainless steel (Fe, Cr, Ni) appear at the anode and cathode, but minimal high-Z contamination is observed in hot central regions



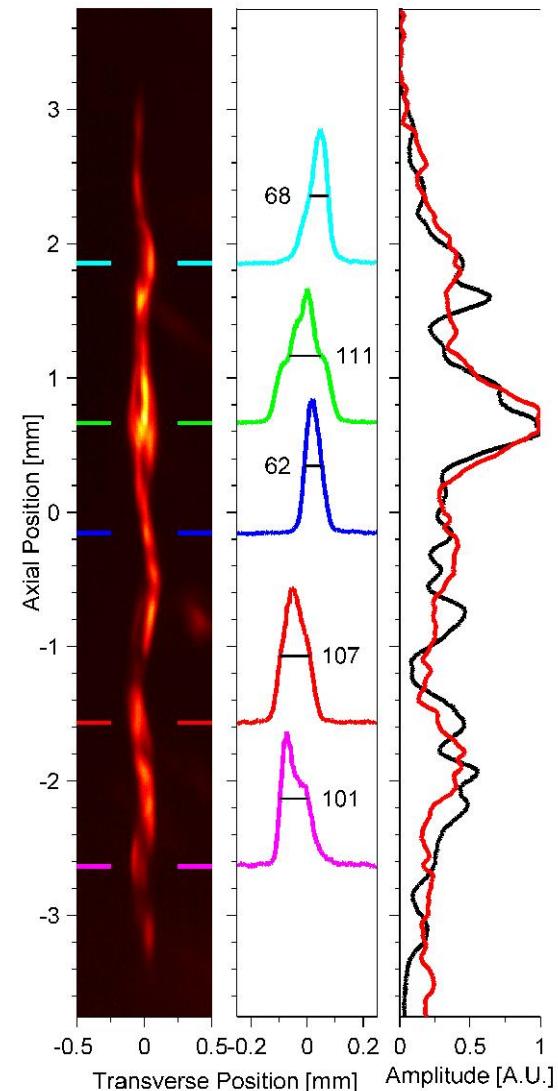
The slope of the high-energy continuum emission implies  $Te \sim 1.5$  keV at the anode and cathode, and  $T \sim 3.5$  keV in the central regions



The measured electron temperature is close to the ion temperature obtained from neutron time-of-flight data; x-ray emission yields are consistent with fuel  $\rho \sim 0.4$  g/cm<sup>3</sup>

# High-resolution monochromatic imaging of the x-ray emission shows a narrow, hot plasma column with weakly helical structure

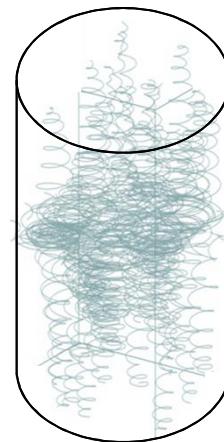
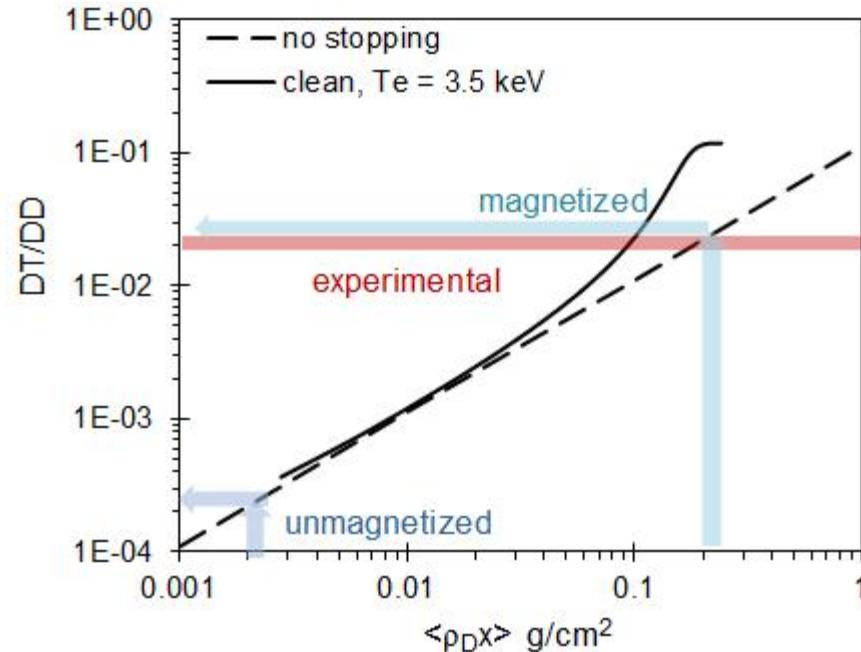
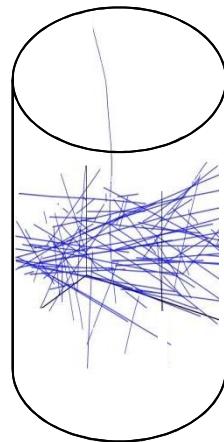
- Lineouts of stagnation column vary from 60 to 120  $\mu\text{m}$  FWHM (resolution about 60  $\mu\text{m}$ )
- Emission is observed from about 6 mm of the 7.5 mm axial extent
- Note that the emission doesn't necessarily define the fuel-liner boundary, but only the hot fuel region
- The stagnation column is weakly helical with a wavelength of about 1.3 mm and a 0.05 mm horizontal offset
- Axial lineouts of image (black) agree with 9.3 keV 1D spectrometer lineouts (red), suggesting features are due to emission and not liner opacity (Be opacity >9 keV small).
- With  $\rho \sim 0.4 \text{ g/cm}^3$ ,  $\rho_r \sim 2 \text{ mg/cm}^2$



# In addition to the significant $\sim 2 \times 10^{12}$ DD neutron yields, we measure a remarkable $\sim 5 \times 10^{10}$ DT neutrons



“Secondary” 14 MeV neutrons can be produced by 1 MeV tritons interacting with D fuel

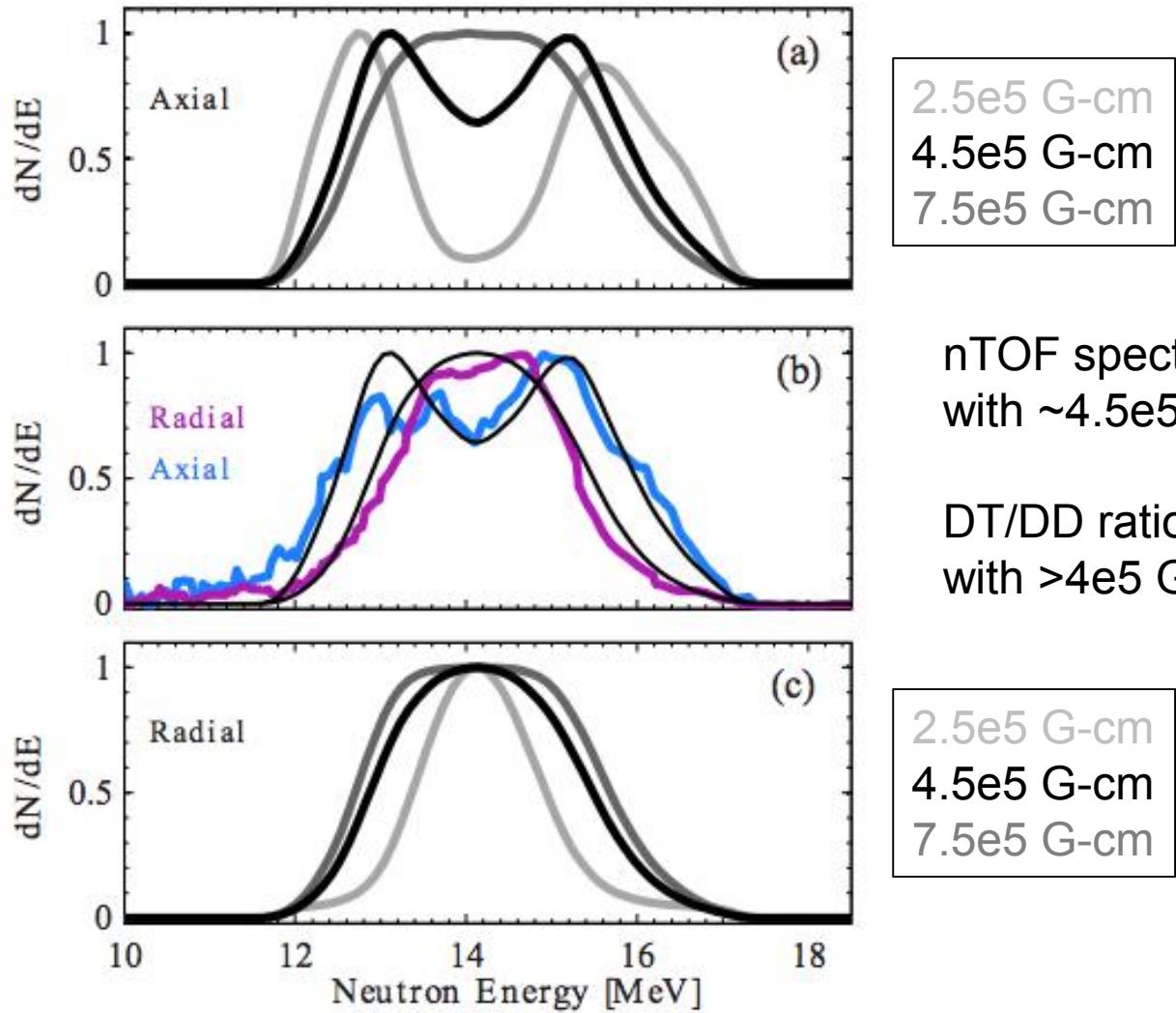


Unmagnetized plasmas must reach pressures of  $\sim 500$  Gbar and  $\rho R > 0.2$  g/cm<sup>2</sup> to achieve the  $\alpha$ -particle confinement required for ignition

In magnetized plasmas, thermal confinement and  $\alpha$ -deposition are both enhanced by  $B$ , reducing pressure and  $\rho R$  requirements by factors of  $\sim 100$ .

A field that confines tritons also confines electrons -- and will confine alphas!

# Neutron time-of-flight data are consistent with high magnetization



2.5e5 G-cm  
4.5e5 G-cm  
7.5e5 G-cm

nTOF spectra consistent with  $\sim 4.5 \times 10^5$  G-cm

DT/DD ratio consistent with  $> 4 \times 10^5$  G-cm

2.5e5 G-cm  
4.5e5 G-cm  
7.5e5 G-cm

# MagLIF Summary

- Results from initial MagLIF experiments have been encouraging, with significant DD and DT yields and strong evidence for good stability, confinement, and scaling
- A helical stagnation column with  $T \sim 3$  keV,  $r \sim 0.4$  g/cc,  $r \sim 50$  um, and  $B_z \sim 10$  kT is consistent with an extensive collection of neutron and x-ray data
- Both integrated and focused experiments are ongoing
- Better understanding of how high magnetic fields affect thermal transport and stopping power will increase confidence in our predictions for yield scaling with increasing current, external field, and laser power

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Studies of emission and absorption from photoionized plasmas relevant to white dwarfs, accretion disks, and solar photosphere
- **Diagnostic development on Z**  
X-ray Thomson scattering and Zeeman splitting

There is a lot of interesting science going on at Z –  
and many opportunities for significant contributions from Cornell's LPS.

# These efforts represents a large collaboration between the NNSA labs and the academic community



Jim Bailey, Taisuke Nagayama,  
Guillaume Loisel, Stephanie Hansen,  
Dave Bliss, Tom Nash, Tom Ao, Eric  
Harding, Greg Rochau, Matt Gomez,  
Michael Desjarlais  
**Sandia National Laboratories**



Roberto Mancini, Iain Hall, Tom  
Lockard, Dan Mayes  
**University of Nevada – Reno**



Don Winget, Mike Montgomery, Ross  
Falcon, Thomas Gomez, Alan Wootton,  
Jennifer Ellis, Sean Moorhead, Roger  
Bengtson  
**University of Texas – Austin**



Anhil Pradhan, C. Orban, Mark  
Pinsonneault, and S.N. Nahar  
**Ohio State University**



Mark Koepke, Ted Lane, Matt Flaugh  
**West Virginia University**



Duane Leidahl, Carlos Iglesias, Brian Wilson  
**Lawrence Livermore National Laboratory**



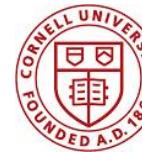
Manolo Sherrill, Heidi Tierney ,Chris Fontes,  
James Colgan, Dave Kilcrease  
**Los Alamos National Laboratory**



C. Blancard, Ph. Cosse, G. Faussurier, F.  
Gilleron, J.C. Pain  
**French Alternative Energies and Atomic  
Energy Commission (CEA)**



Joe MacFarlane, Igor Golovkin  
**Prism Computational Sciences**

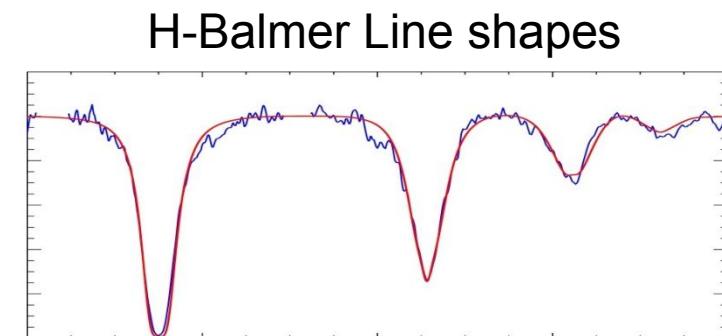
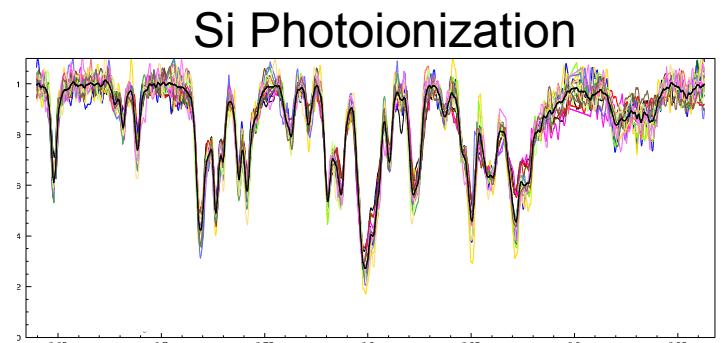
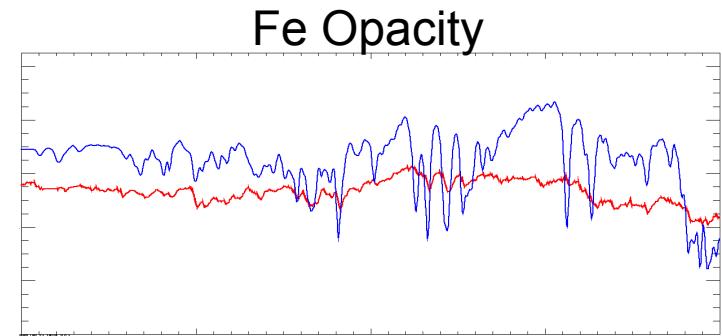


Laura Johnson  
**Cornell University**

We are interested in developing  
new collaborations

# ZAPP experiments measure the fundamental properties of atoms in plasmas to solve important astrophysical puzzles:

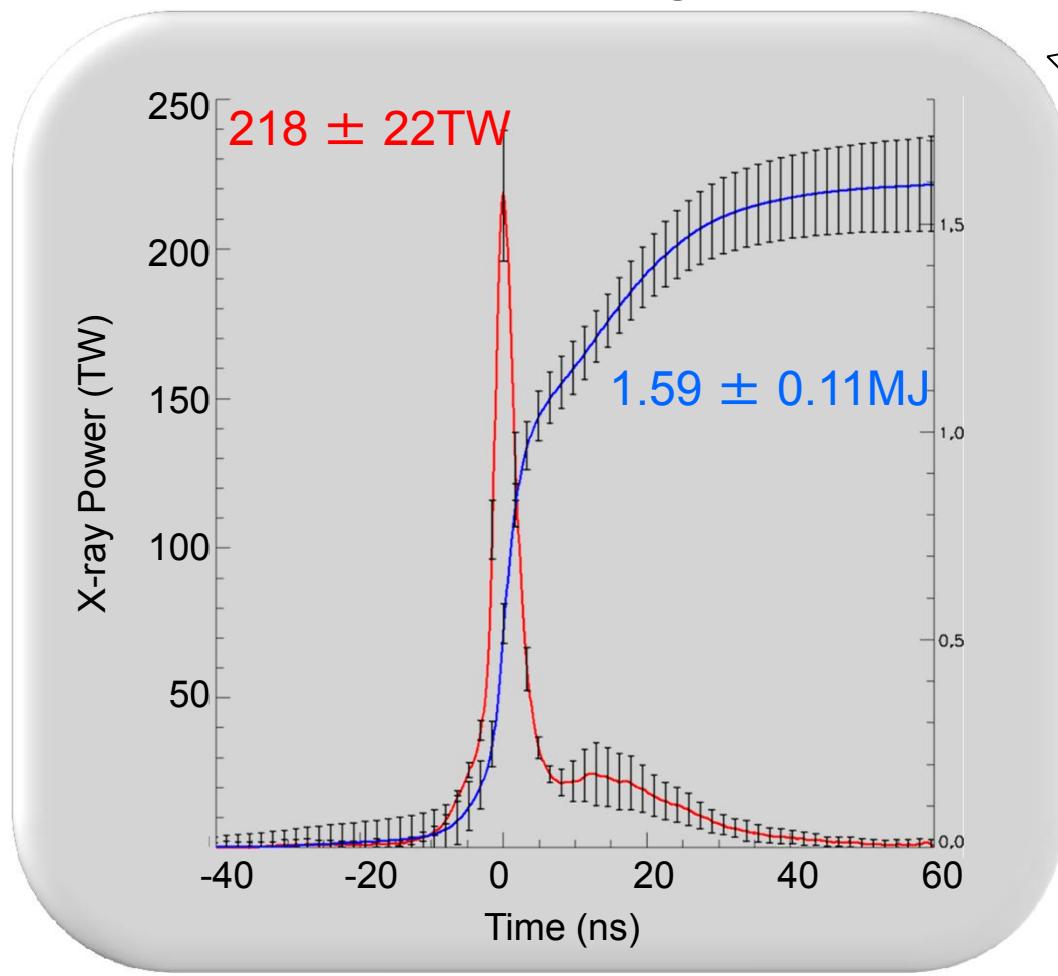
- Why can't we predict the location of the convection zone boundary in the Sun?
  - Opacity of Fe at  $T \sim 200$  eV
- How does ionization and line formation occur in accreting objects and warm absorbers?
  - Ionization distribution and spectral properties of photoionized Ne and Si
- Why doesn't spectral fitting provide the correct properties for White Dwarfs?
  - Stark-broadened H-Balmer line profiles



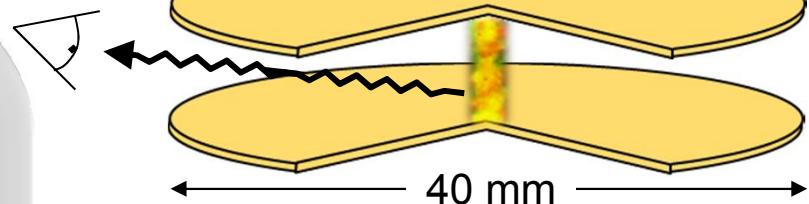
# The ZPDH x-ray emission is reproducible to $\pm 10\%$ in peak power and $\pm 7\%$ in energy



Radial X-ray Power and Energy  
(20 shot average)



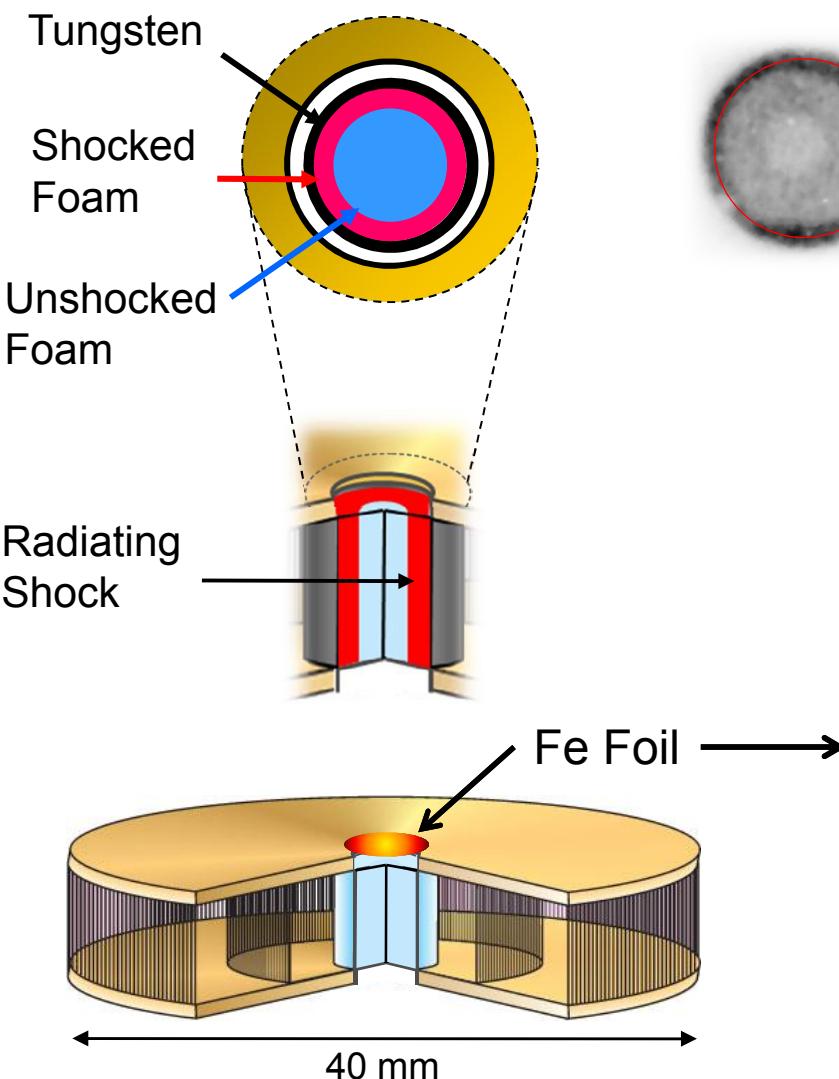
Z-pinch Dynamic Hohlraum



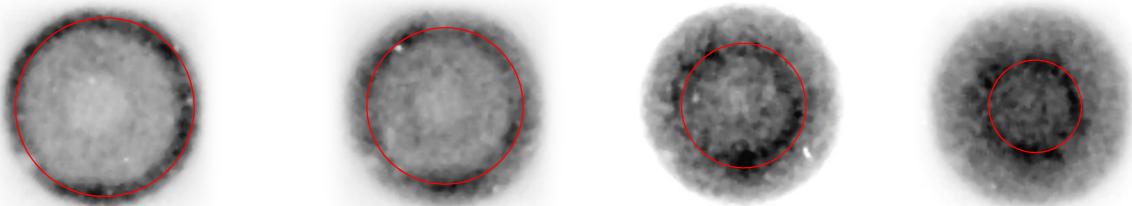
	ZR >2011	Z <2007
Marx Energy	20.3 MJ	11.4 MJ
Ipeak	25.8 MA (1.5%)	21.7 MA* (2.1%)
Mass	8.5 mg	3.8 mg
Peak Power	220 TW (10%)	120 TW (14%)
Radiated Energy	1.6 MJ (7%)	0.82 MJ (17%)

\*Wagoner, PRSTAB 11 (2008)

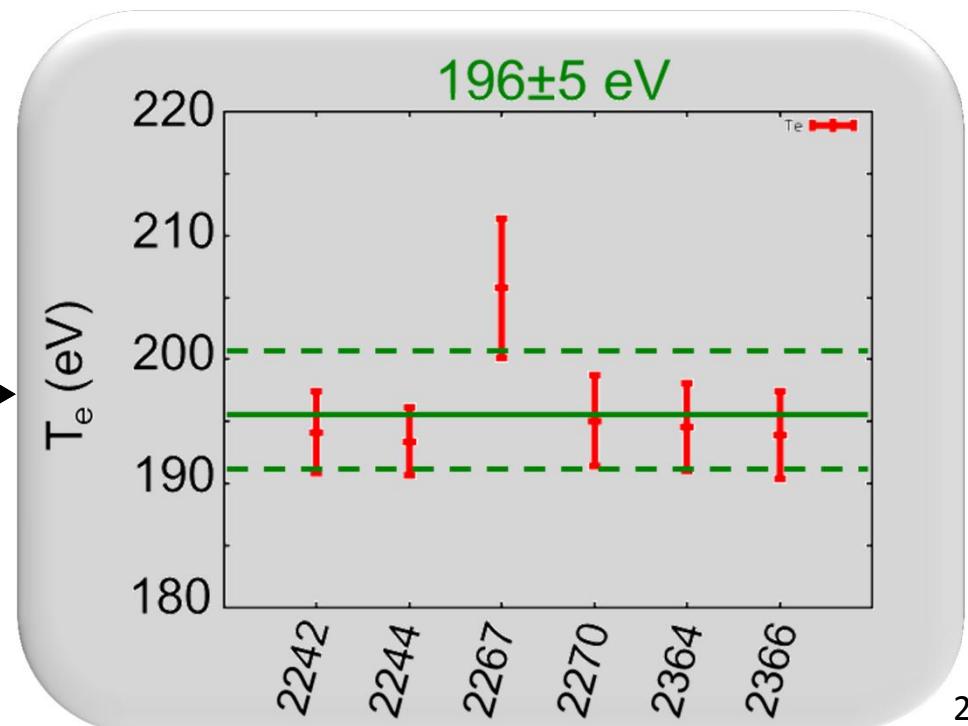
The ZPDH can also radiatively heat samples placed above the z-pinch to  $T_e \sim 200$  eV.



Framing Pinhole Camera Images

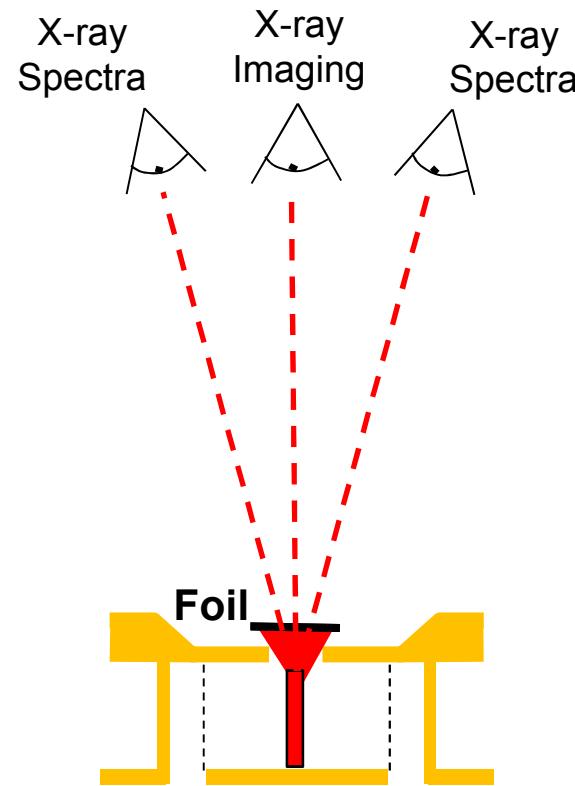


Axial Fe Foil Temperature

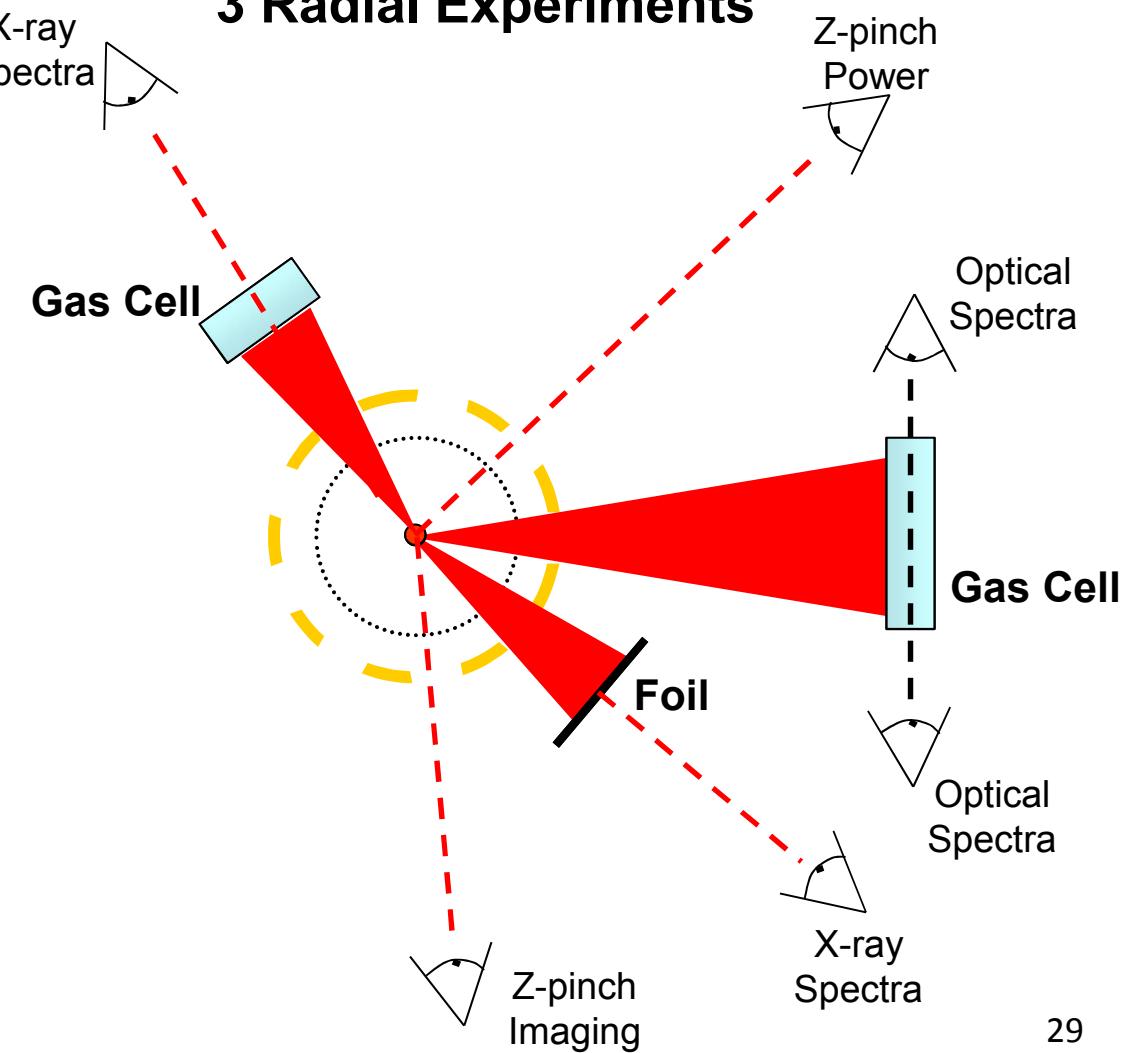


# The ZPDH simultaneously drives four independent experiments on a single ZAPP shot

## 1 Axial Experiment



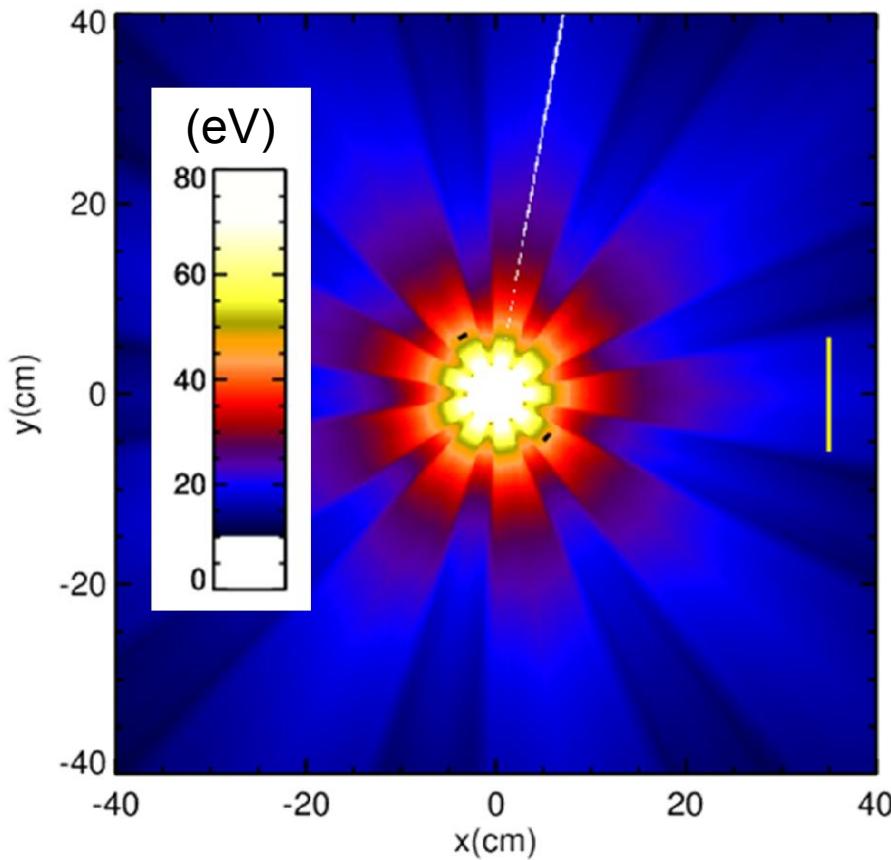
## 3 Radial Experiments



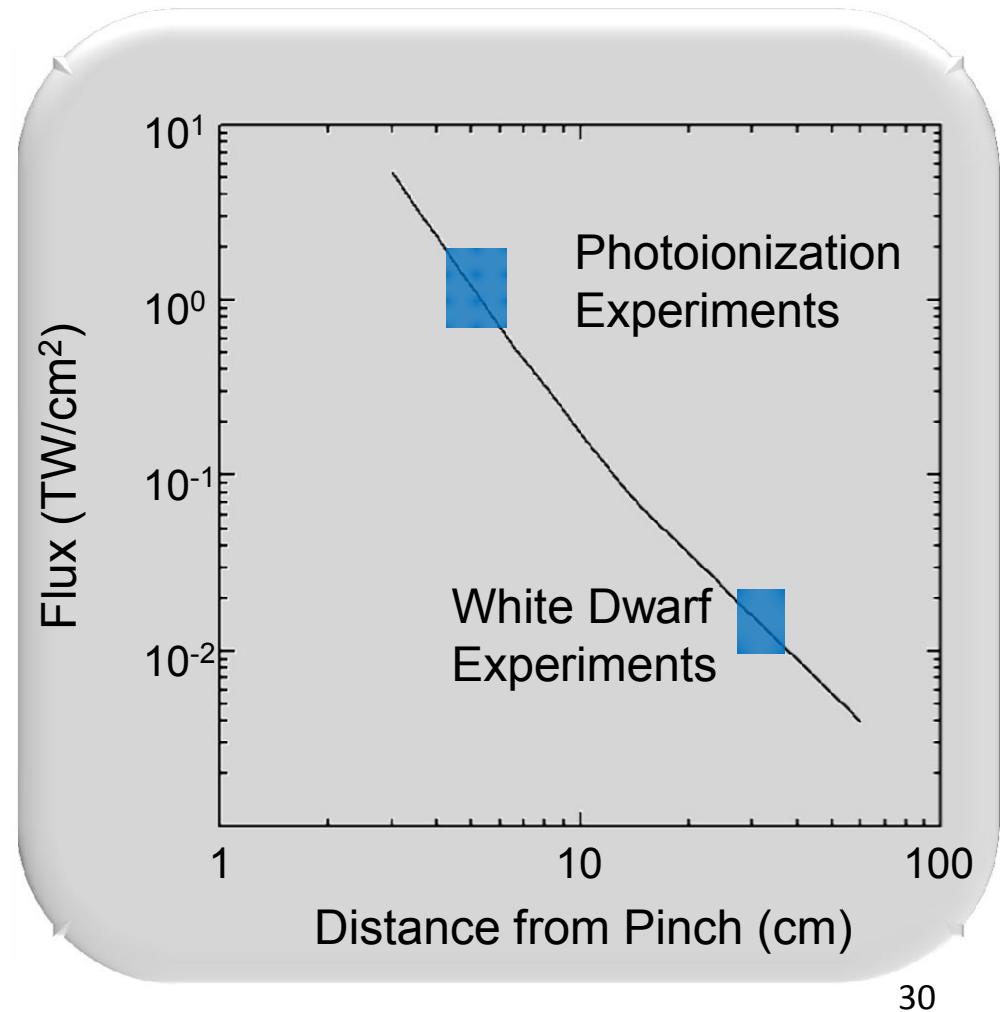
# Placing samples at multiple distances from the z pinch provides a broad range of drive flux.



**r-θ Peak Brightness Temperature Contours Around Z Pinch**



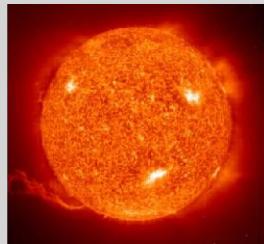
**Peak Drive Flux on a Sample**



# ZAPP campaigns simultaneously study multiple issues spanning 200x in temperature and $10^6$ x in density



## Solar Opacity



### Question:

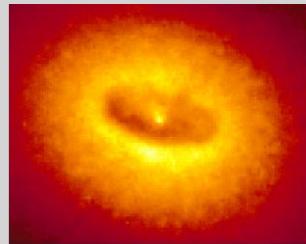
Why can't we predict the location of the convection zone boundary in the Sun?

### Achieved Conditions:

$T_e \sim 200$  eV,  $n_e \sim 10^{23}$  cm $^{-3}$



## Photoionized Plasmas



### Question:

How does ionization and line formation occur in accreting objects?

### Achieved Conditions:

$T_e \sim 20$  eV,  $n_e \sim 10^{18}$  cm $^{-3}$



## White Dwarf Line-Shapes



### Question:

Why doesn't spectral fitting provide the correct properties for White Dwarfs?

### Achieved Conditions:

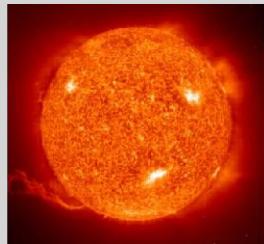
$T_e \sim 1$  eV,  $n_e \sim 10^{17}$  cm $^{-3}$



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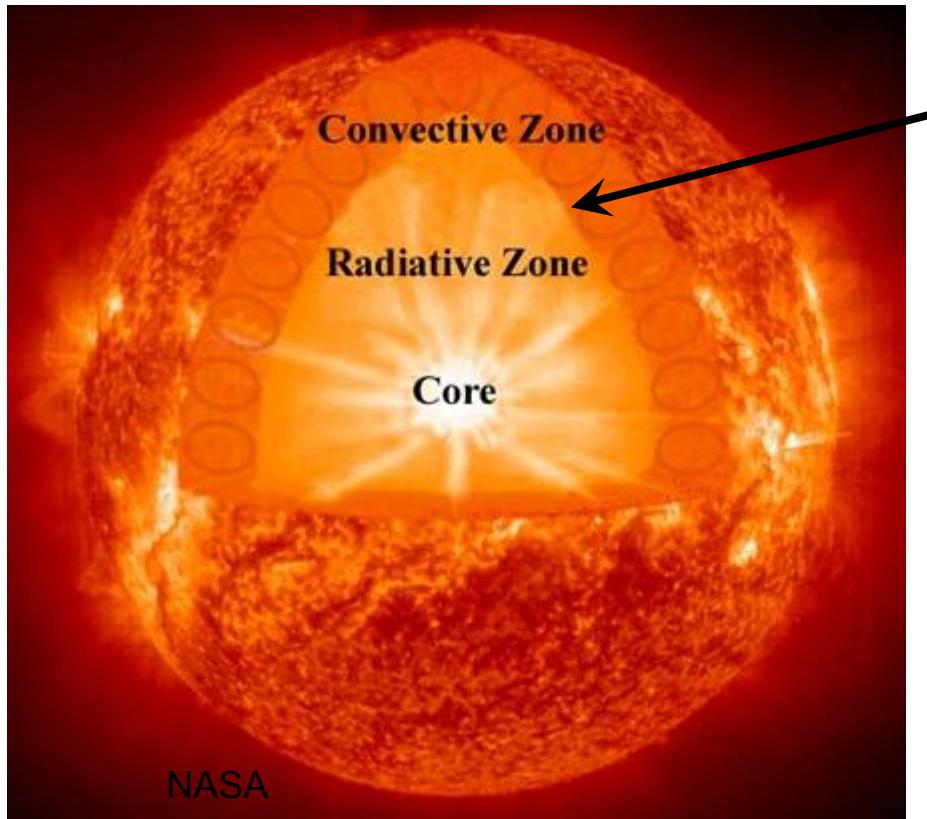
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# Models for solar interior structure disagree with helioseismology observations.



**Convection-Zone (CZ) Boundary**  
Models are off by  $10\text{-}30 \sigma$

**Models depend on:**

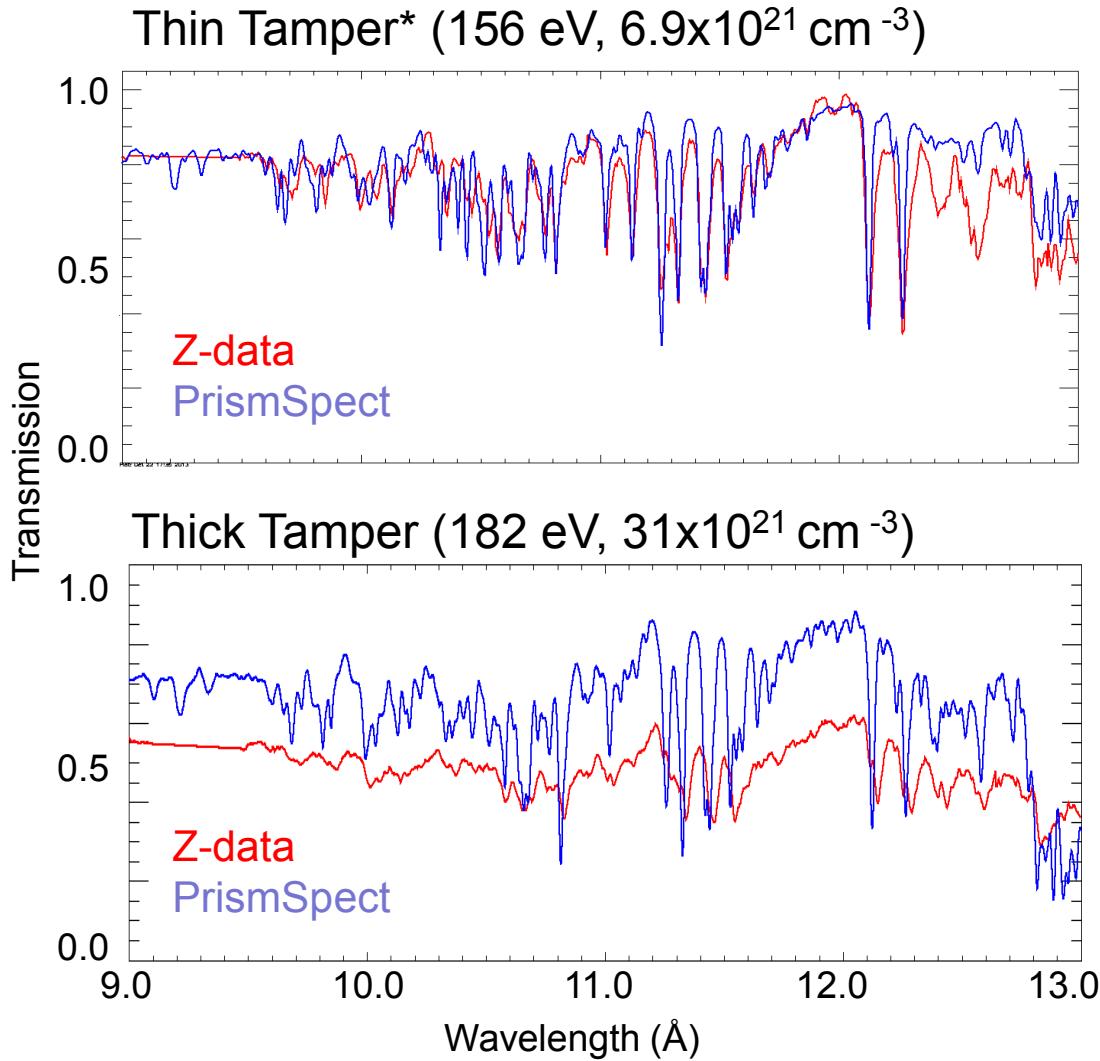
- Composition (revised in 2005\*)
- EOS as a function of radius
- The solar matter *opacity*
- Nuclear cross sections

**Question:** Is opacity uncertainty the cause of the disagreement?

**Objective:** Measure Fe opacity at CZ base conditions.

\*M. Asplund *et al*, Annu. Rev. Astro. Astrophys. **43**, 481 (2005).

# Modern computations of Fe opacity show large disagreements with data at CZ base conditions



## Present Status

- Agreement between data and computation becomes worse at increasing temp. and dens.
- Disagreements at CZ base conditions can partially explain the CZ boundary problem.
- The differences are probably not unique to Fe... more scrutiny of the data is prudent.

# ZAPP campaigns simultaneously study multiple issues spanning 200x in temperature and $10^6$ x in density



## Solar Opacity



### Question:

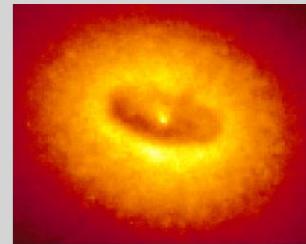
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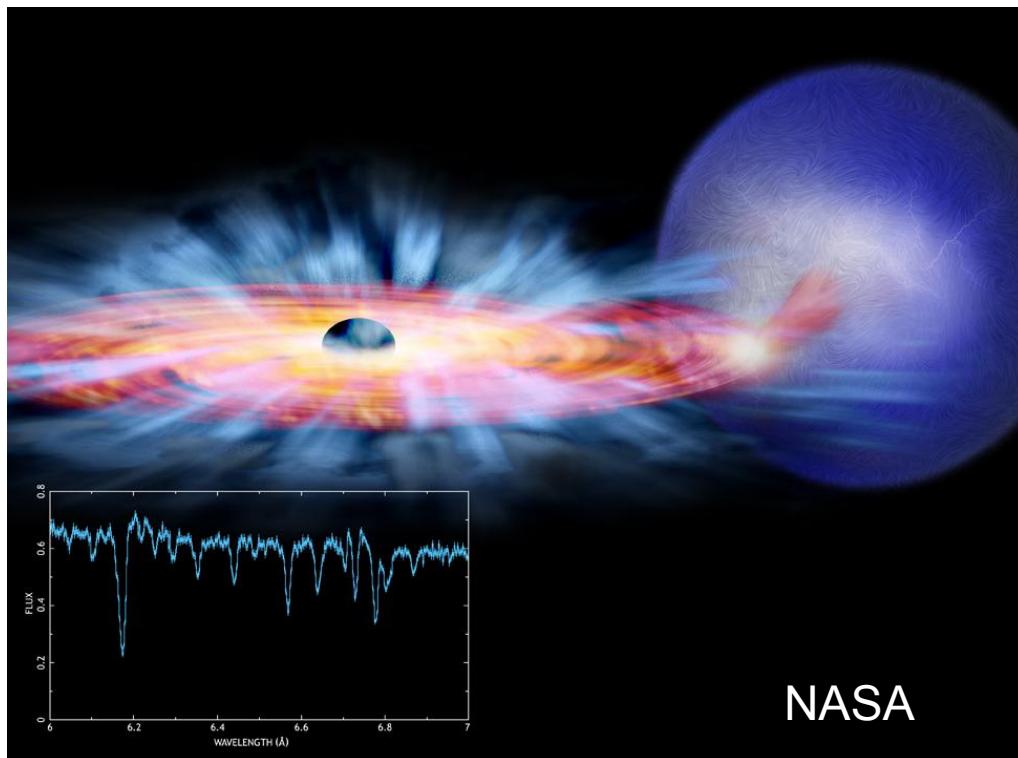


# We learn about black holes from the matter falling into them – these are photoionized plasmas



Conceptual Picture of a Black-Hole Accretion Disk

$$\xi \sim 10 - 10,000 \text{ erg.cm.s}^{-1}$$



Photoionization parameter

$$\xi \equiv \frac{4\pi F}{n_e} \text{ [erg.cm.s}^{-1}\text{]}$$

Laboratory Plasmas

$$n_e \sim 10^{19} \text{ cm}^{-3}$$

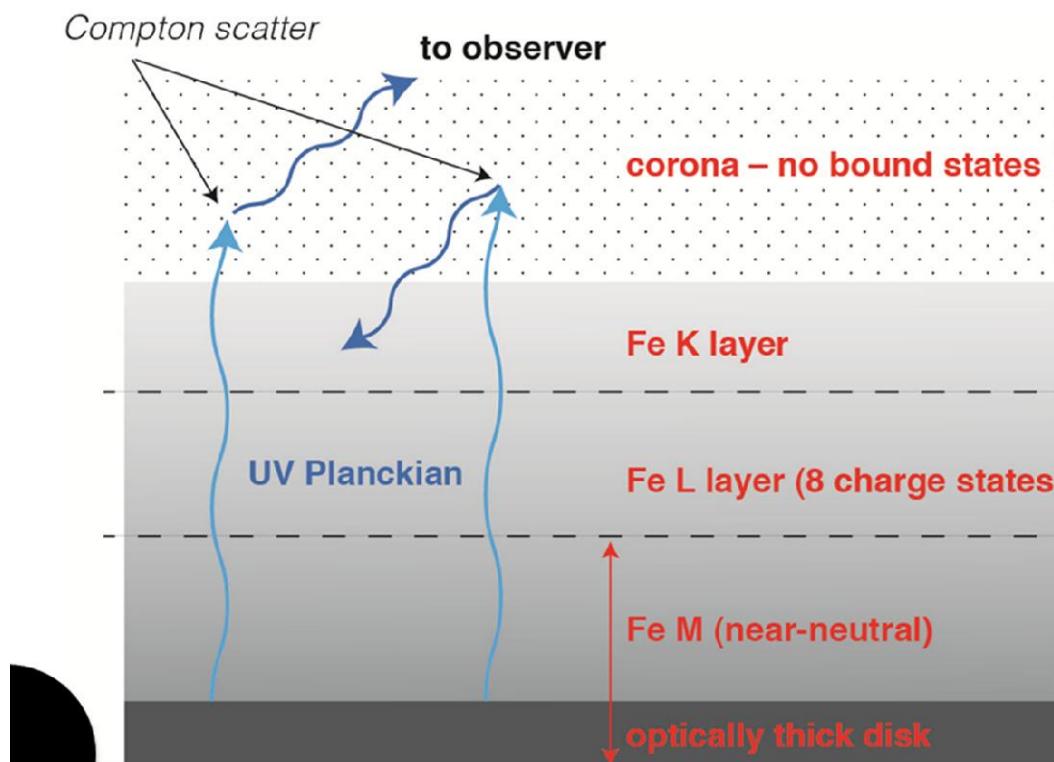
$$F > 1 \text{ TW/cm}^2 \text{ for } \xi > 10$$

- Can we model the ionization?
- Can we model the line emission?

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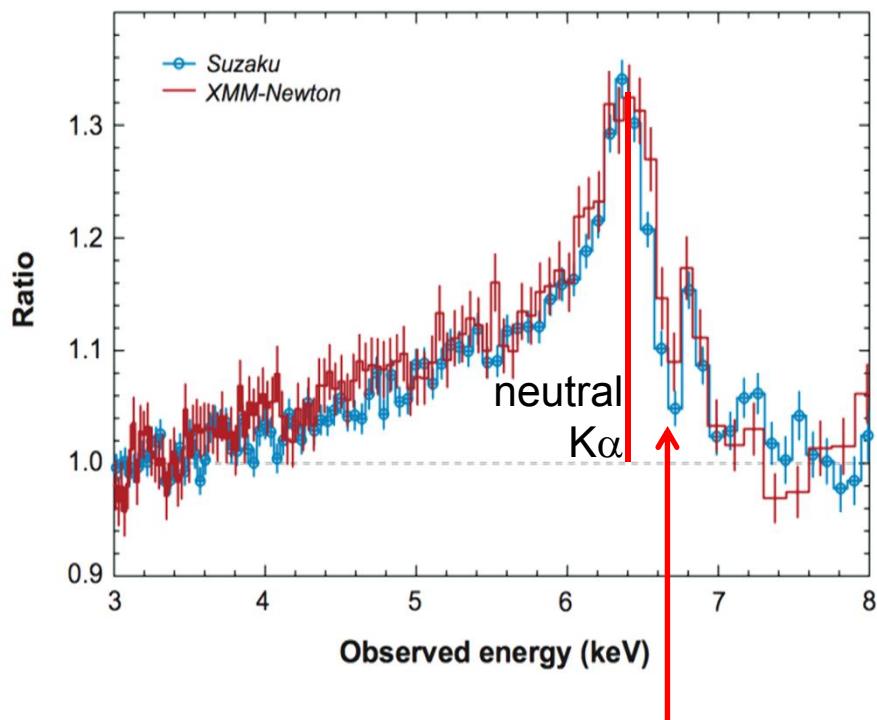
$$n_e \sim 10^{19} \text{ cm}^{-3}$$

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- Can we model the ionization?
- Can we model the line emission?

# A Specific Problem: Emission from L-shell ions is not seen in some prominent black-hole accretion disks.

## Measured Fe Emission from MCG 6-30-15

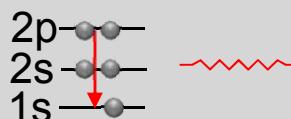


No observed emission from Fe ionized to the L-shell

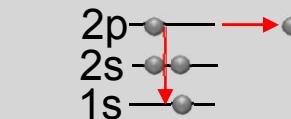
## Resonant Auger Destruction (RAD) was accepted as the reason\*

- 2 competing processes for the de-excitation of L-shell ions:

*Radiative Decay*

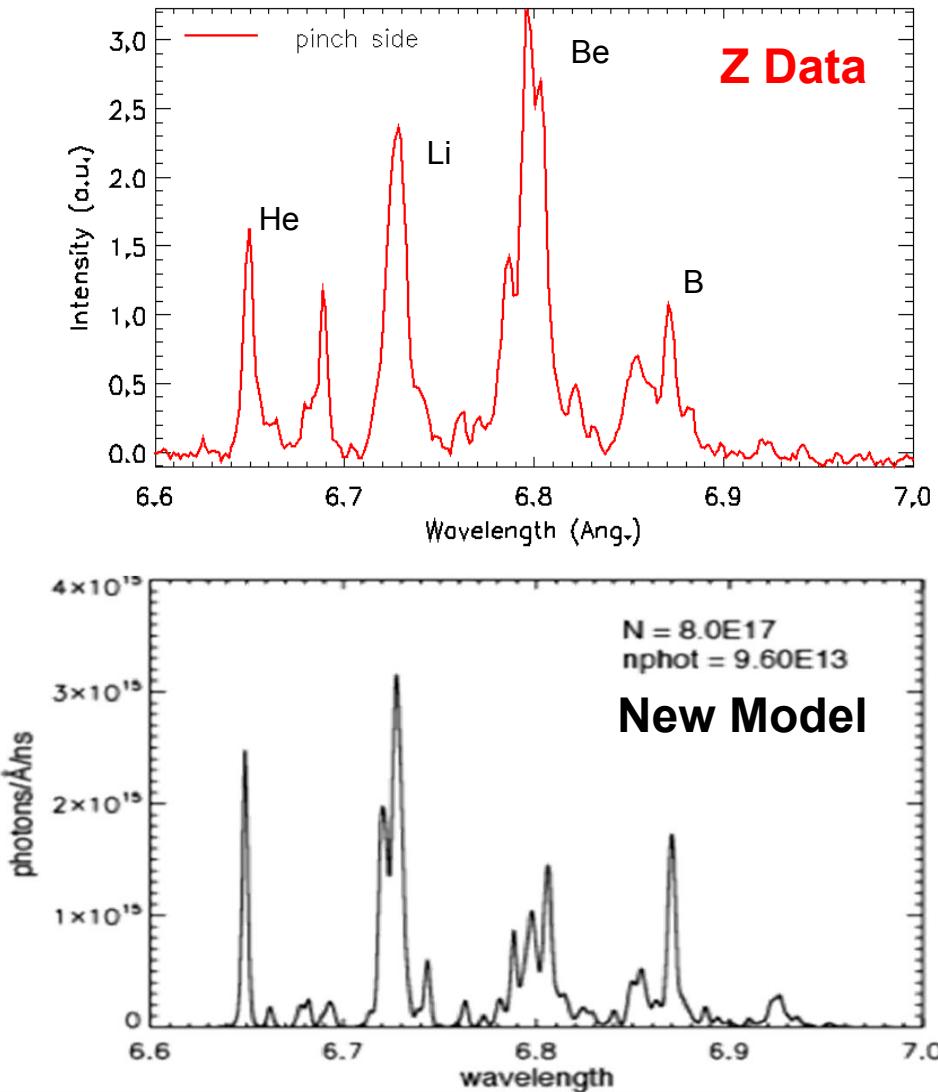


*Auger Decay*



- **Thin Plasma:** high probability of observing the photon
- **Thick Plasma:** high probability of the photon being resonantly absorbed  
→ Higher probability of Auger Decay for the ensemble

# Recent emission measurements demonstrate that L-shell emission is not 100% quenched by RAD.



## Present Status

- Z Data demonstrates that L-shell emission does escape at column depths  $>1E17$  at/cm<sup>2</sup>.
- Present data can discriminate between models of the ionization distribution AND relative line strengths.
- Absolute intensity is needed to determine efficiency of RAD process.

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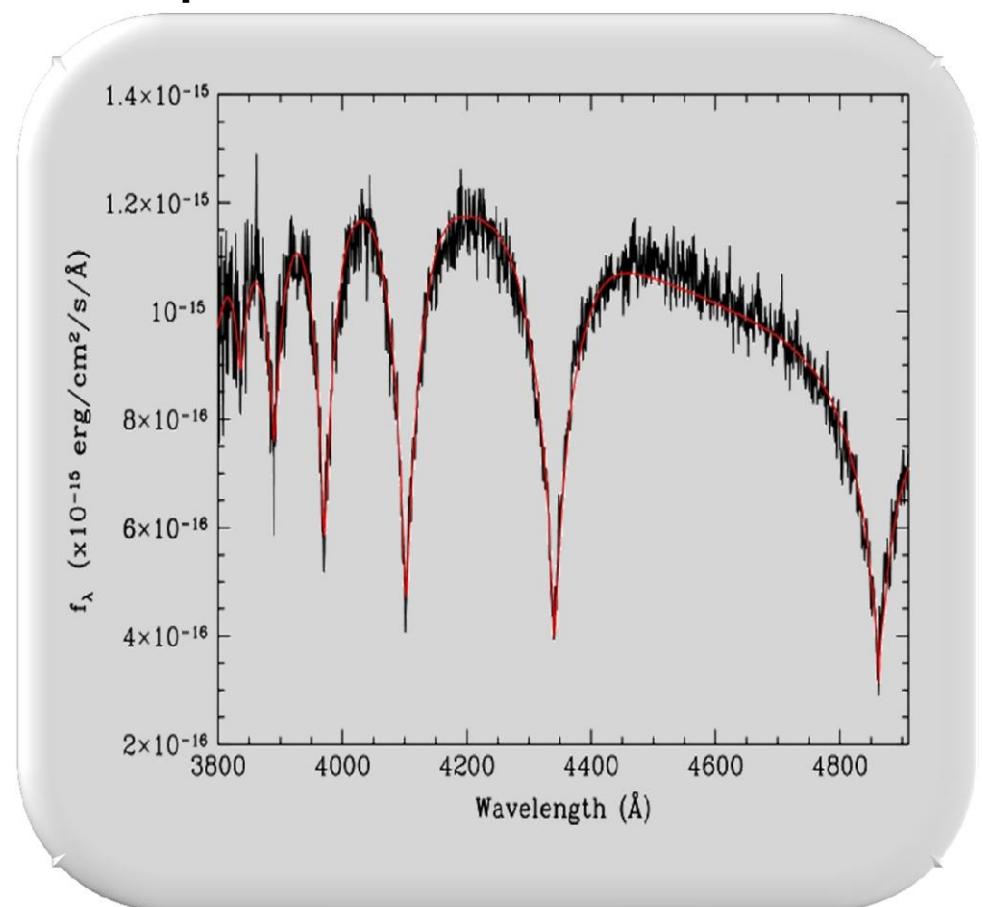
$T_e \sim 1 \text{ eV}$ ,  $n_e \sim 10^{17} \text{ cm}^{-3}$



# The properties of White Dwarfs are determined by spectral fitting, but disagrees with other methods

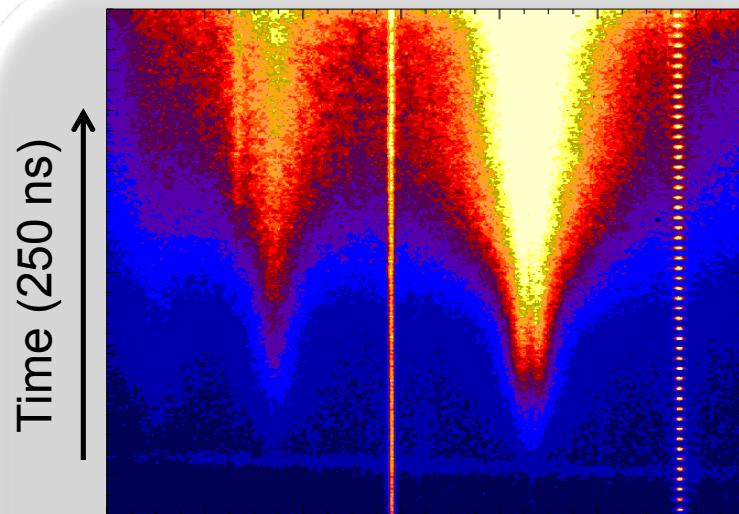
- White Dwarfs are fundamentally important
  - Evolutionary endpoint for ~98% of stars
  - Simple in structure and evolution
  - Cosmic laboratories (cosmochronology)
- WD surface temperature and total mass are usually determined by fitting the observed spectra
- The spectroscopic method and gravitational redshift disagree by >10% in the stellar mass

Spectral fit of WD J1916+3938\*

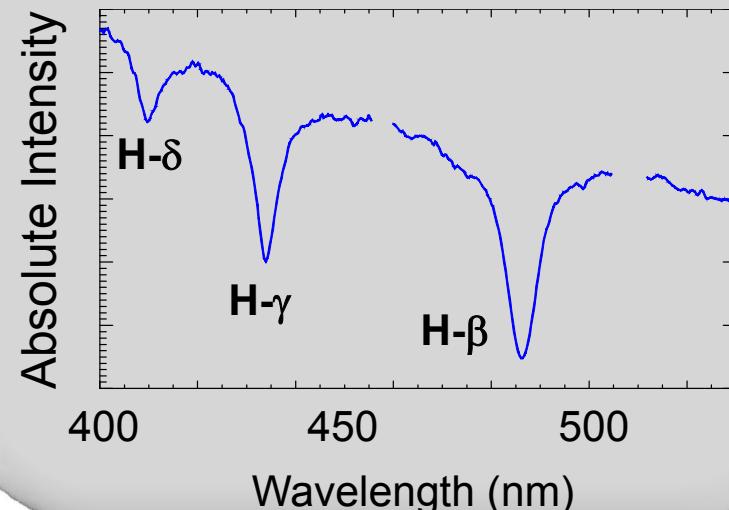
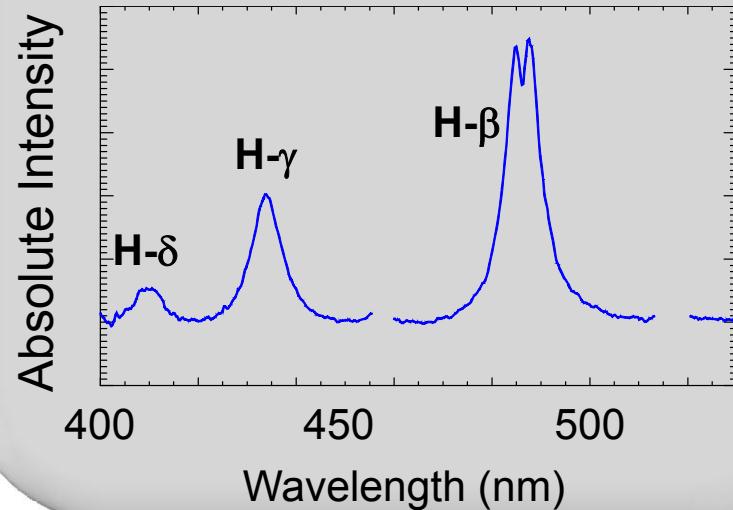
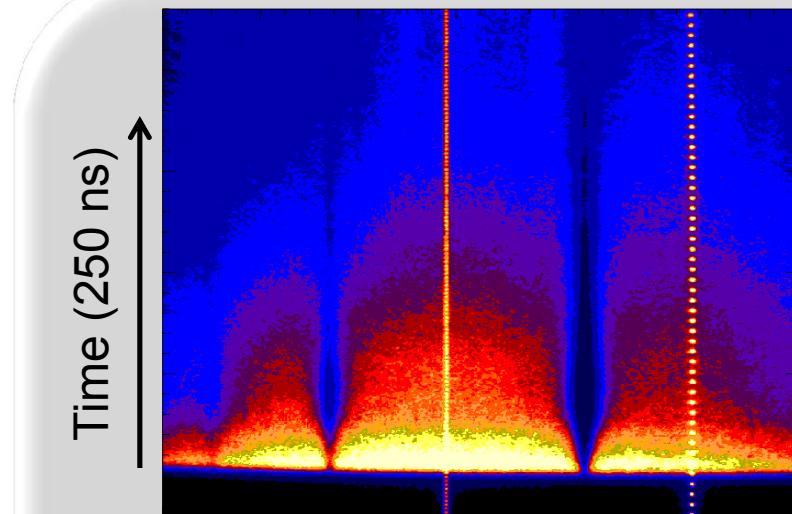


# Simultaneous streaked absorption and emission provide a unique capability to measure lineshapes

Emission



Absorption



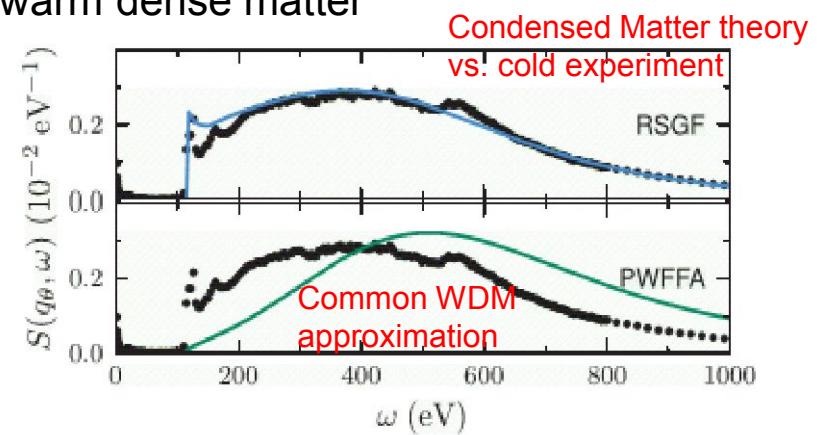
# Outline

- **Magnetized Liner Inertial Fusion (MagLIF)**  
Encouraging results from initial integrated experiments
- **Z Astrophysical Plasma Properties (ZAPP) Collaboration**  
Studies of emission and absorption from photoionized plasmas relevant to white dwarfs, accretion disks, and solar photosphere
- **Diagnostic development on Z**  
X-ray Thomson scattering and Zeeman splitting

There is a lot of interesting science going on at Z –  
and many opportunities for significant contributions from Cornell's LPS.

# X-ray Thomson Scattering has recently emerged as a potential diagnostic for warm dense matter (WDM)

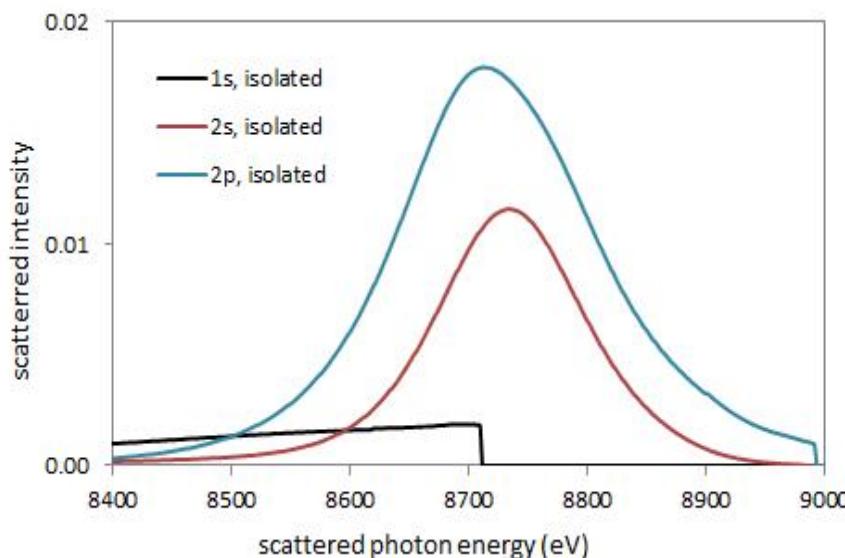
- Warm Dense Matter is difficult to both model and diagnose:
  - Since the kinetic energy of the plasma particles is of the same order as their potential energy, one cannot make the simplifying assumptions appropriate for Condensed Matter or Ideal Plasmas
  - Since Warm Dense Matter emits weakly and absorbs strongly, optical diagnostics only probe its surface; an external x-ray source is required for bulk characterization
- Seven recent PRLs have been devoted to the use of XRTS as a diagnostic for electron density, temperature, and ionization
  - However, the models used to predict scattering signals make significant assumptions about the electronic and ionic properties of warm dense matter
- Experimental studies with uniform, well-characterized sample conditions and high signal-to-noise scattering measurements are needed to augment existing cold scattering data



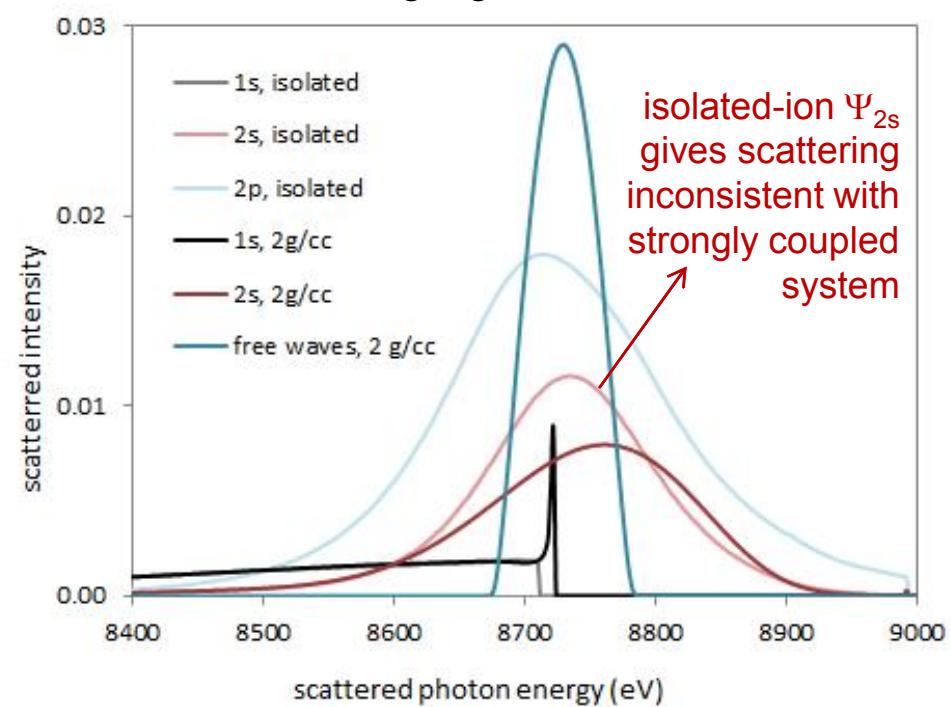
B.A. Mattern & G.T. Seidler,  
*Phys Plas* **20**, 022706 (2013)

# Key question: what happens to the scattering signal as bound electrons become pressure ionized?

135° scattering from cold, isolated carbon:  
Valence 2p state is bound



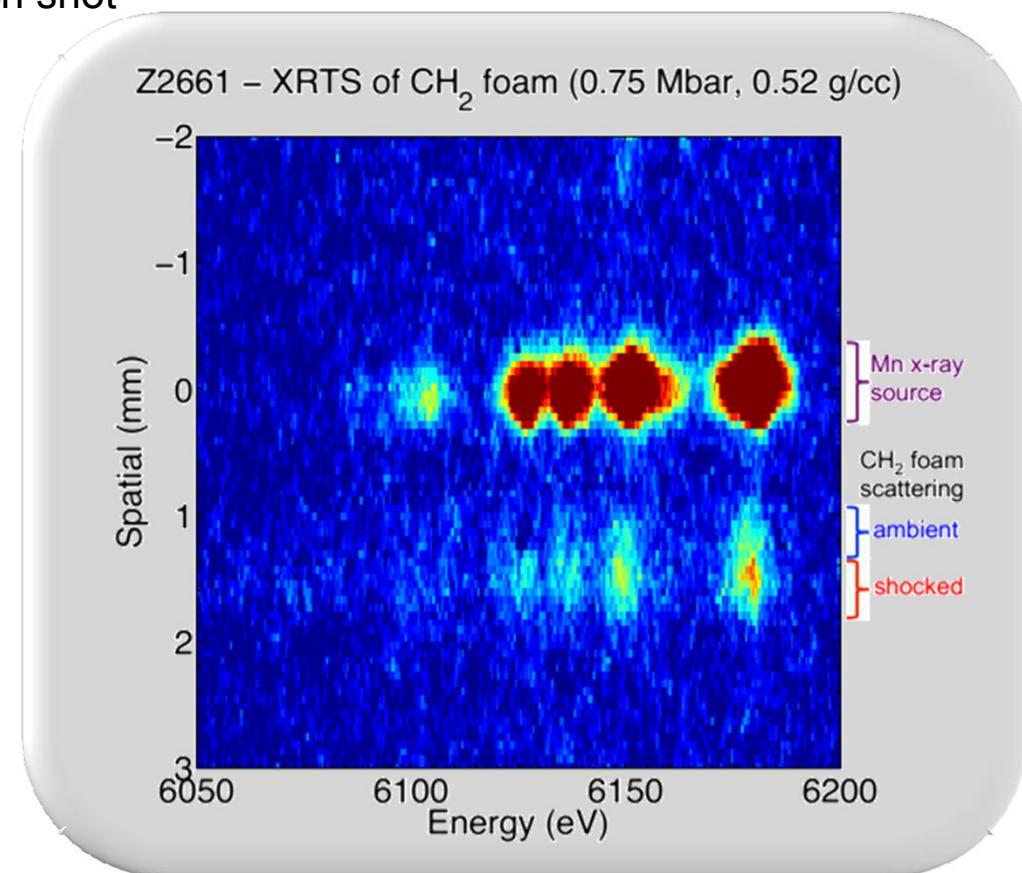
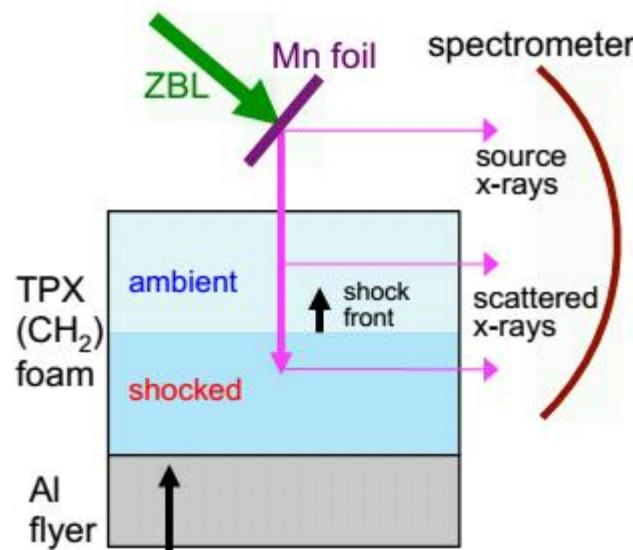
Scattering from cold carbon at 2 g/cc:  
2p is pressure ionized (but not to a plane wave!)  
1s and 2s scattering signals are modified



Laura Johnson is investigating the effect of continuum-wave distortion on free-free scattering signals.

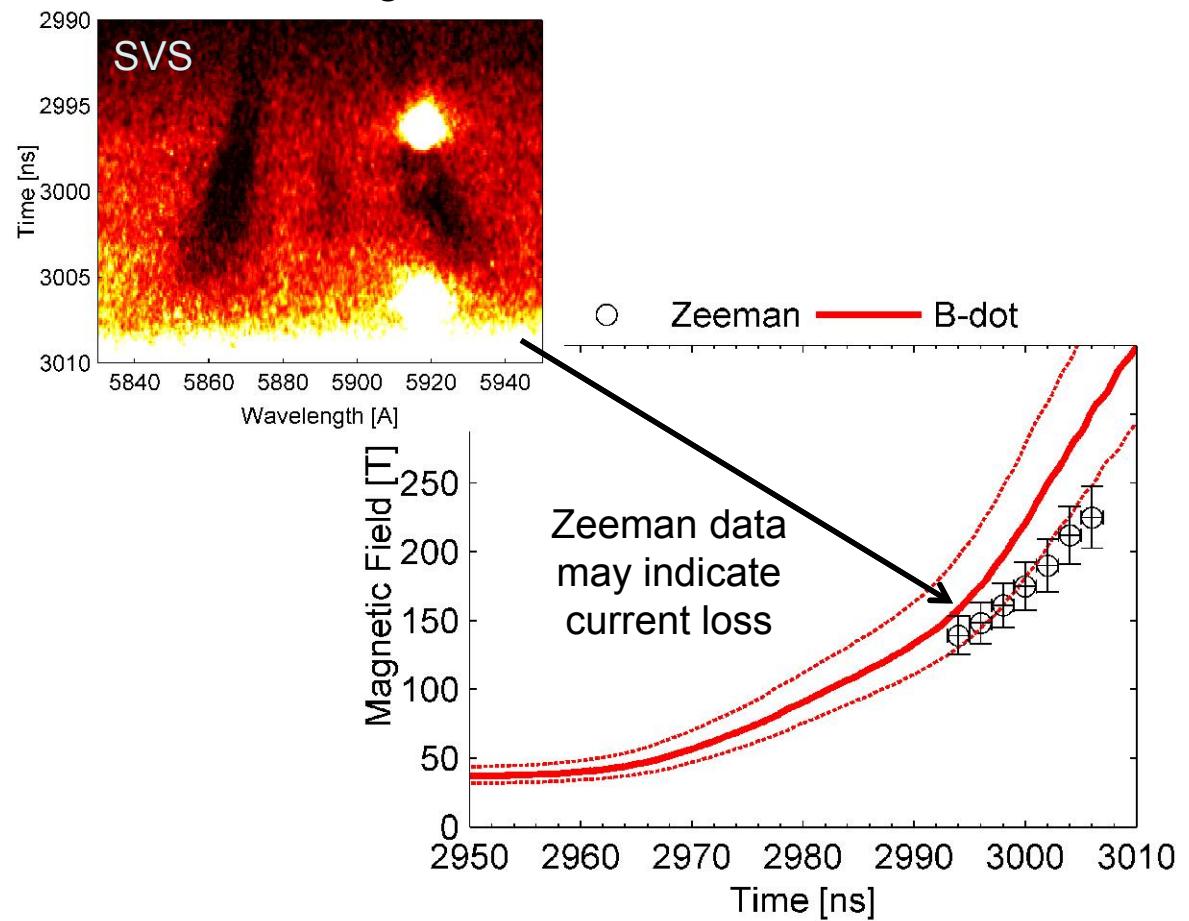
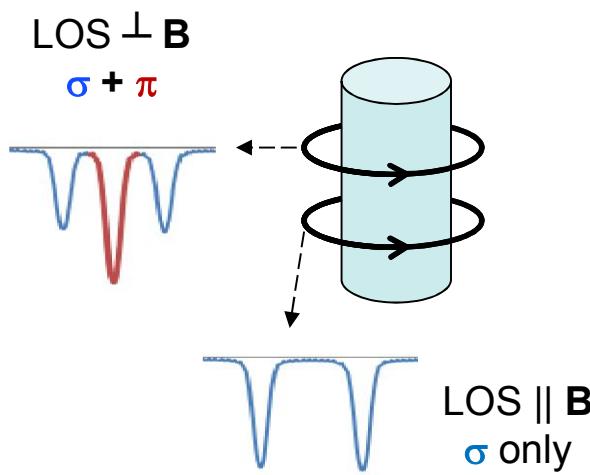
# We recently obtained scattering data from both shocked and unshocked foam on Z

- Z experiments provide:
  - Uniform, long-lived, well-defined shock state
  - *In-situ* comparison with ambient state
  - Source spectrum measured for each shot
  - Potential for high S/N



# Zeeman splitting is being used to characterize Z's current drive and flux compression in Magnetized Liner Inertial Fusion (MagLIF) experiments

- Sodium deposits vaporized and backlit by current-carrying surfaces signal both the magnitude and direction of the local magnetic field:



# Summary: We are using Z to explore new frontiers in laboratory astrophysics and diagnostics



- The Z Astrophysical Plasma Properties (ZAPP) collaboration uses the intense x-ray flux from wire arrays to create mm-scale regions of material with properties similar to those found in accretion objects, white dwarfs, and the solar photosphere, and measures their properties with extensive x-ray and optical spectrometers
- X-ray spectrometers are being fielded on dynamic materials experiments to help advance understanding of x-ray scattering in warm dense matter, with the potential to provide critical tests of scattering theory and temperature diagnostics
- Streaked optical spectrometers are being used to measure Zeeman splitting of optical absorption features to diagnose local magnetic fields important for understanding flux compression and current loss

We welcome ideas for additional collaborations!