

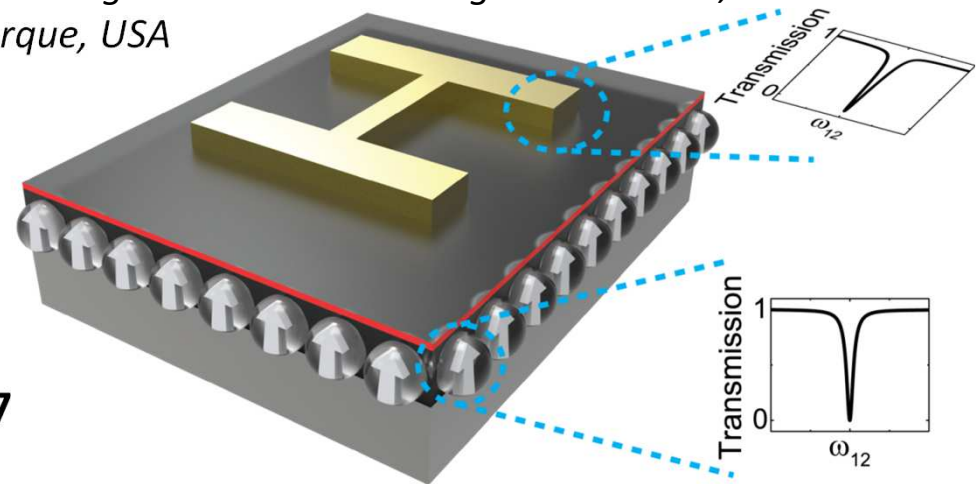


Voltage Tunable Strong Coupling with Planar Metamaterials: From Fundamentals to Optoelectronic Devices

Igal Brener

*Sandia National Laboratories & Center for Integrated Nanotechnologies NM 87185,
Albuquerque, USA*

ibrener@sandia.gov, (505) 844-8097

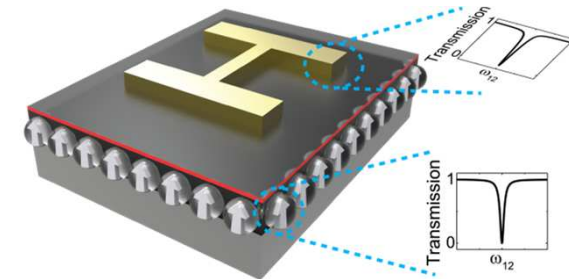
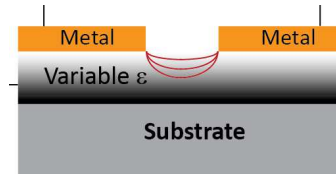


Sandia is a multiprogram laboratory operated by Sandia Corporation, a Lockheed Martin Company,
for the United States Department of Energy's National Nuclear Security Administration
under contract DE-AC04-94AL85000.



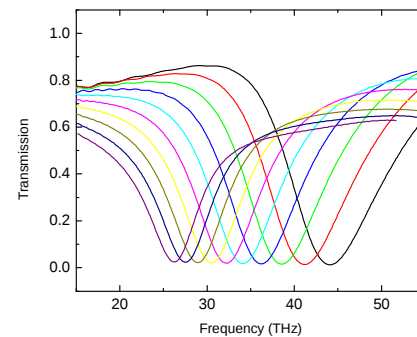
Planar Metamaterials: Strong Coupling & Tuning

MM resonators interact strongly with thin layers underneath

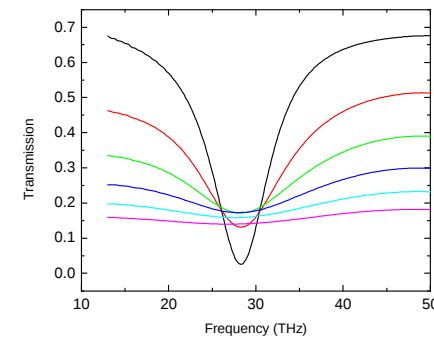


- Nonresonant

Change ϵ_1 (n)

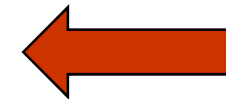
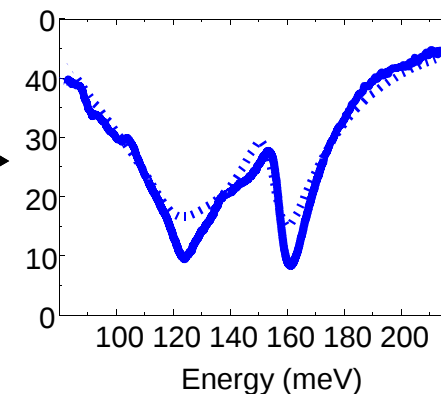
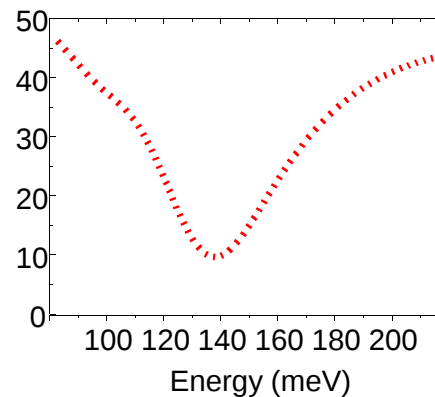


Change ϵ_2 (n)



- Resonant
(Create a dipolar resonance, example: phonon)

Nano Letters 11, 2104 (2011)





Outline

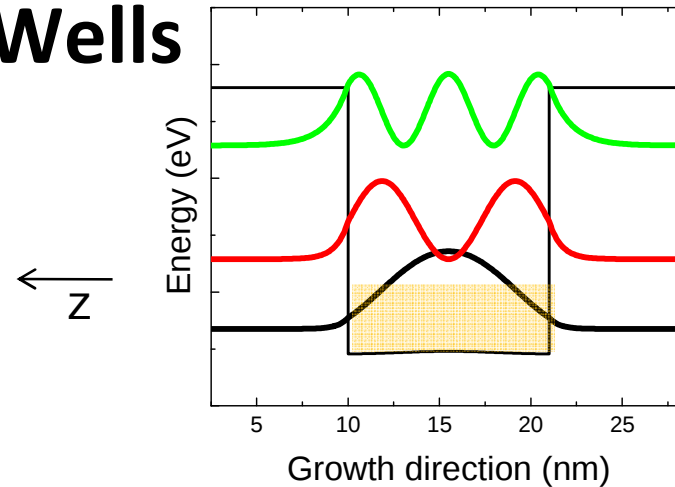
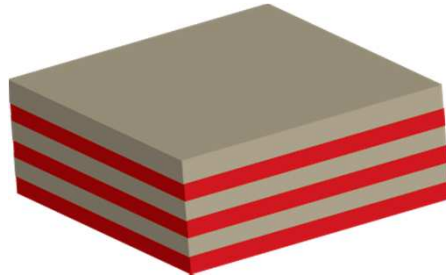
- **Strong Coupling to Intersubband Transitions**
 - Fundamentals
 - Voltage Tuning
 - Wavelength scaling
- **Strong Coupling to Epsilon-near-zero modes**
 - Fundamentals
 - Voltage Tuning
- **Combining coupling to ENZ modes and Intersubband Transitions**



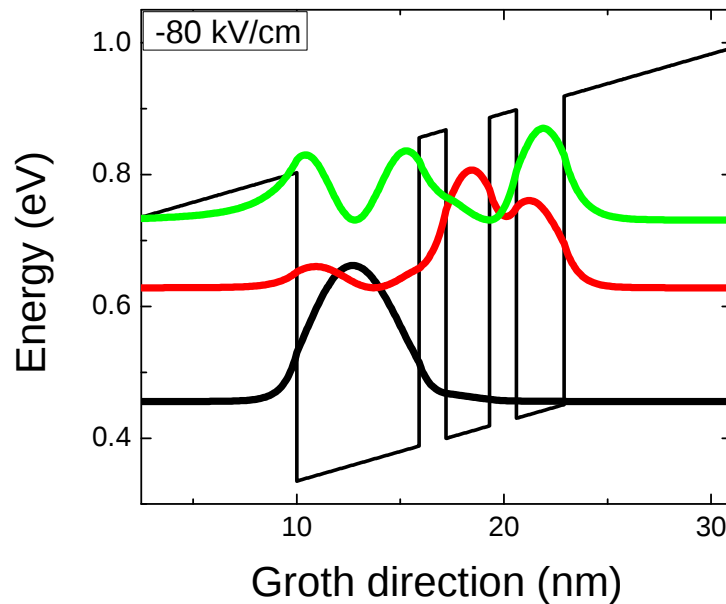
Strong Coupling to Inter-subband Transitions in Quantum Wells

- Scalable (far IR to near IR), Mature, Versatile

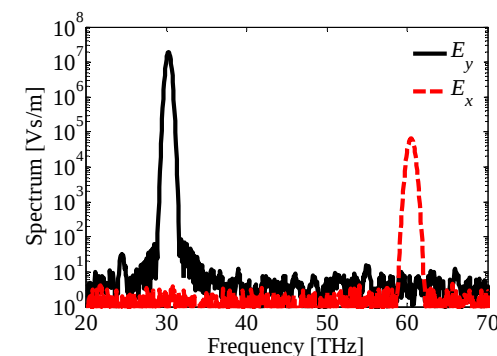
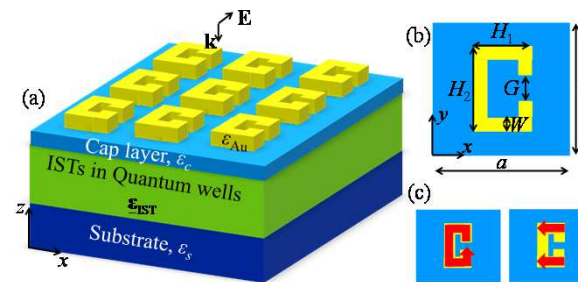
Opt. Express 20, 6584 (2012),
APL 98, 203103 (2011)



Tunable by changing levels using an electric field



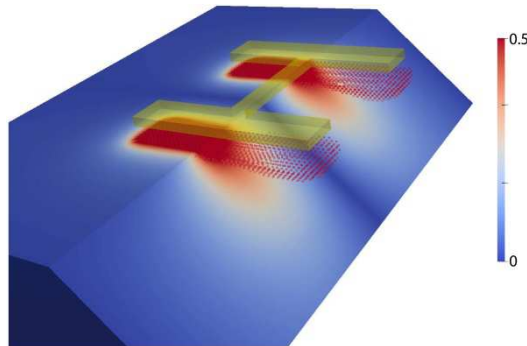
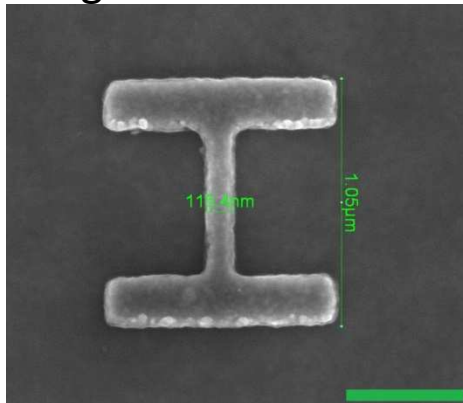
- Also good for nonlinear optics



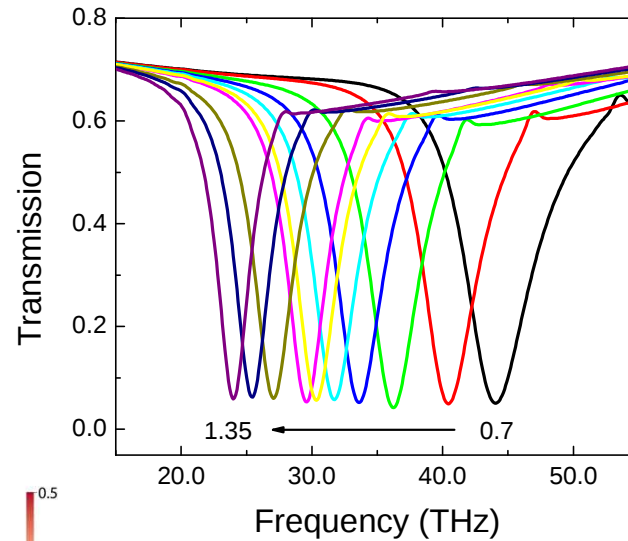


Strong Coupling Theory vs. Experiment

“Dogbone” resonators



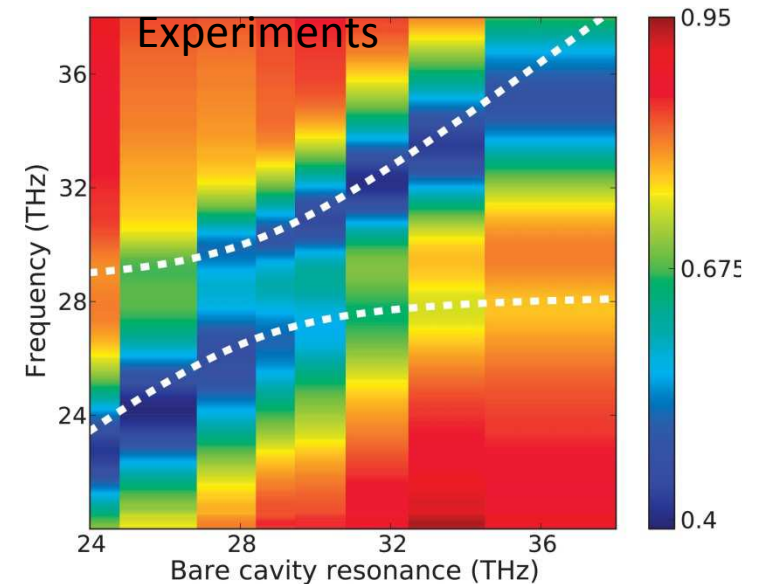
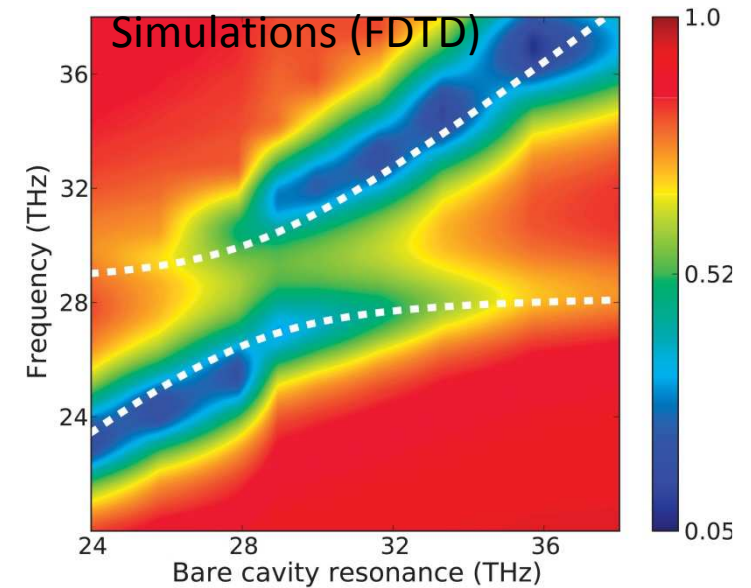
Rabi frequency depends on
simple parameters



Plasma frequency

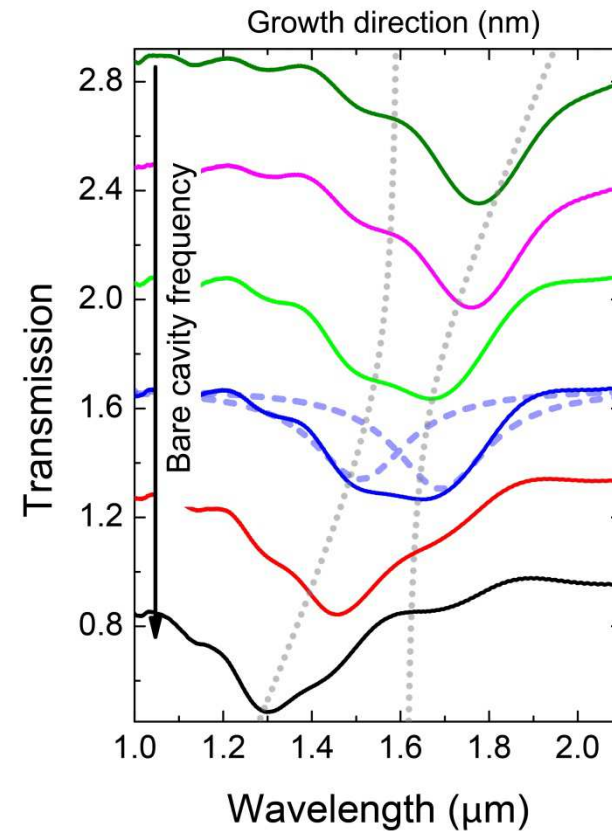
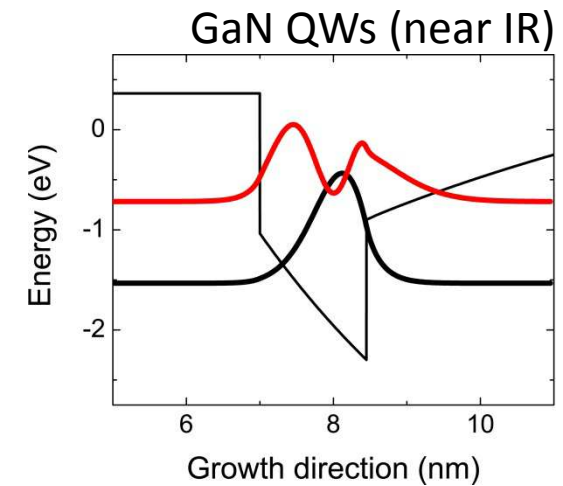
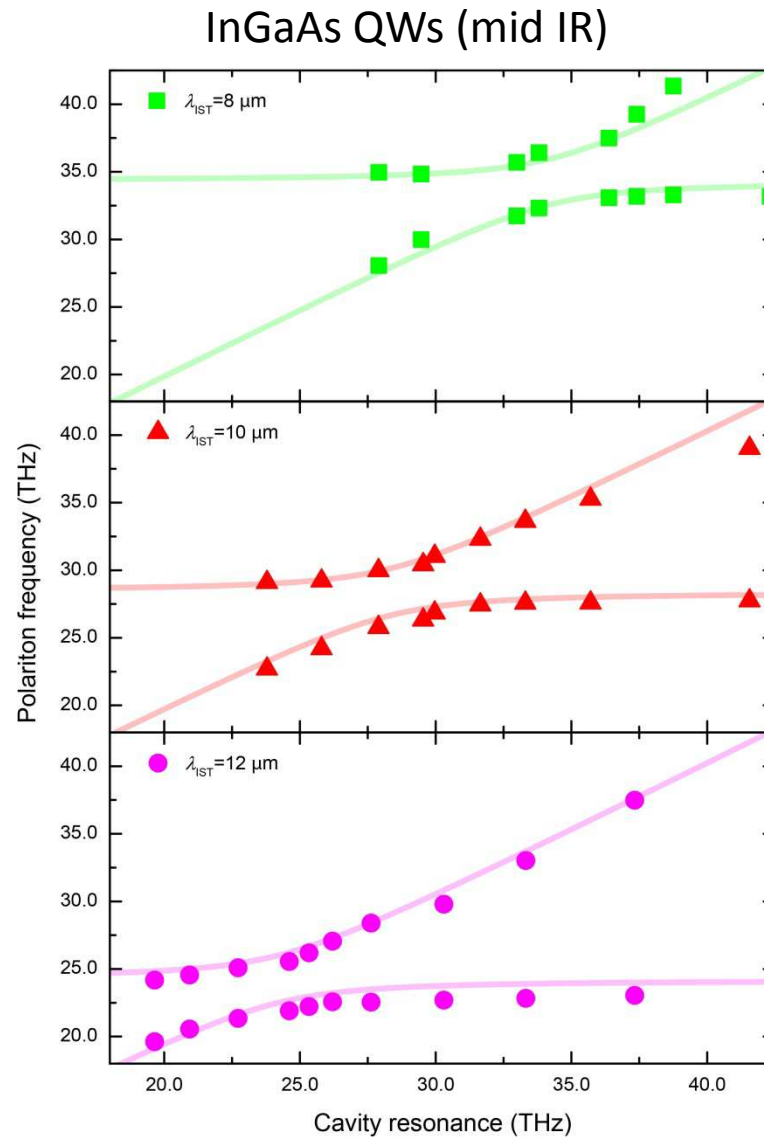
$$\Omega_R = \sqrt{f_W \omega_p} / 2$$

Geometry factor



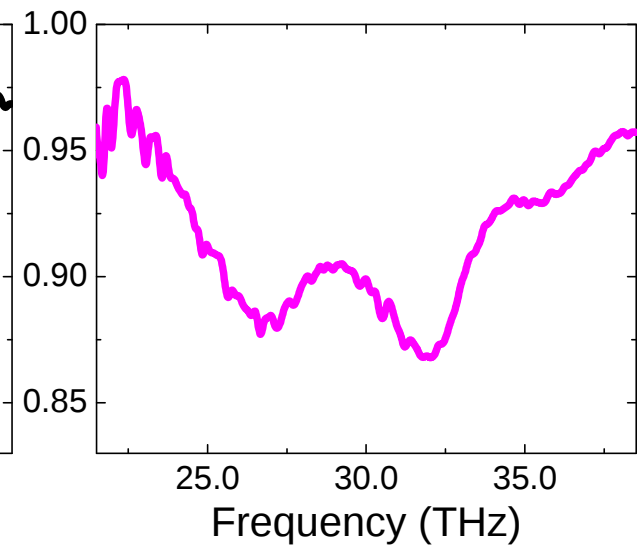
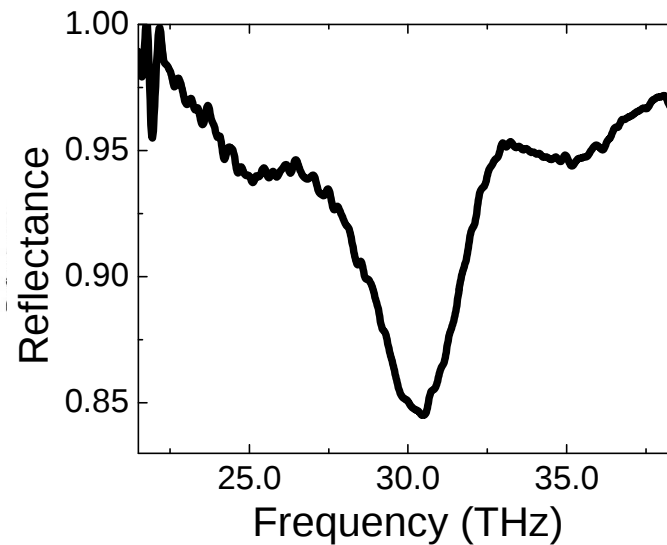
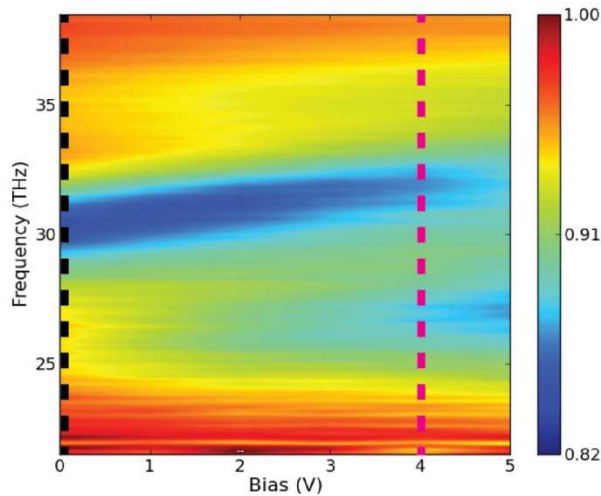
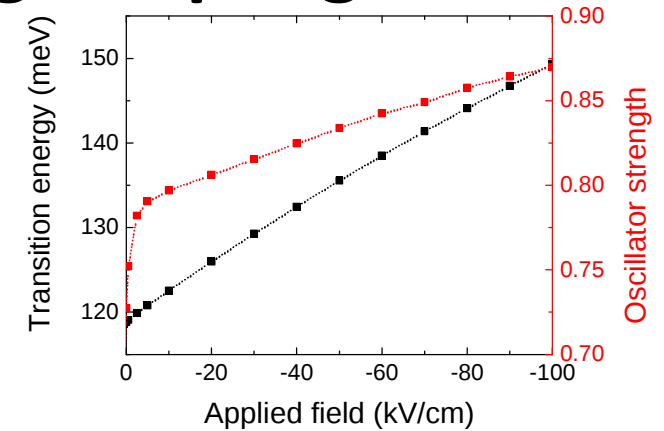
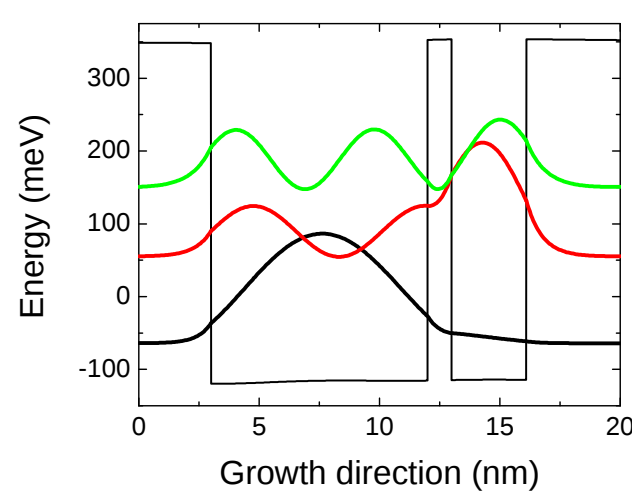
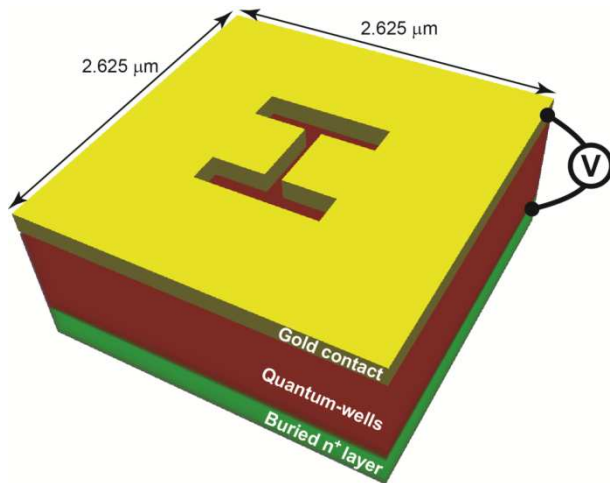


From Mid-IR to Near IR





Electrically Tunable Strong Coupling



Upper polariton shifted by 2.5 THz (=8% of center frequency or one full intersubband linewidth)

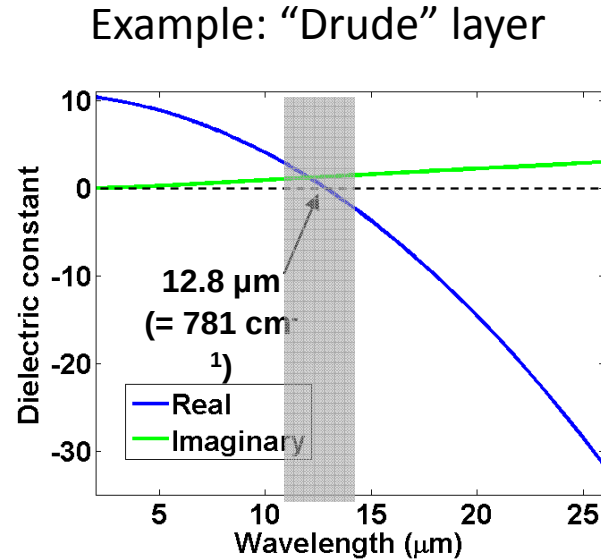
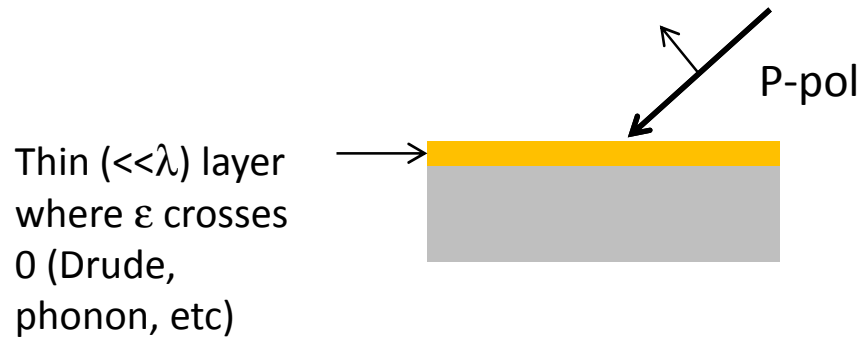


Outline

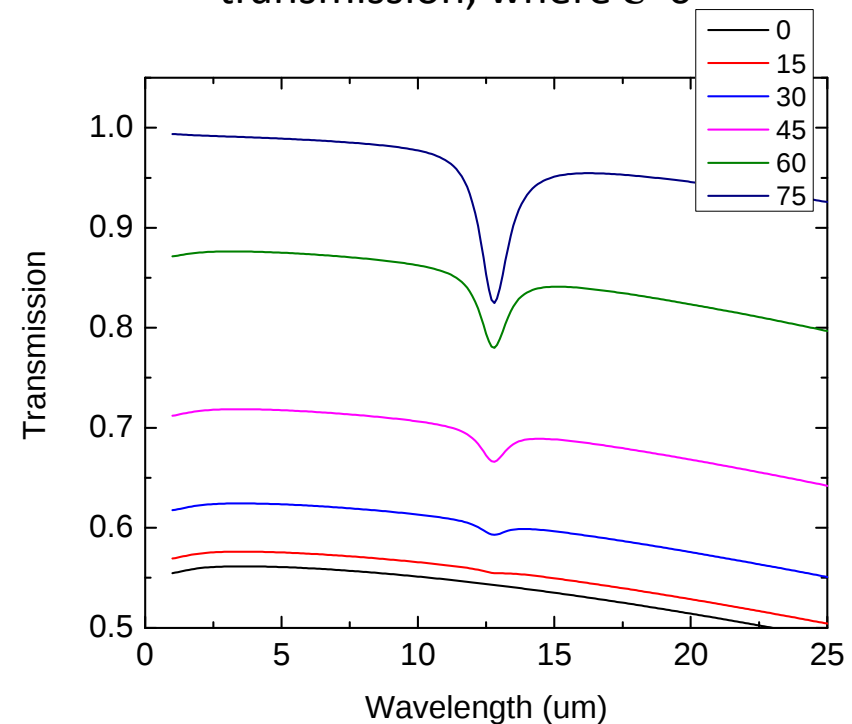
- **Strong Coupling to Intersubband Transitions**
 - Fundamentals
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 - Fundamentals
 - Voltage Tuning
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“Epsilon Near Zero” Modes



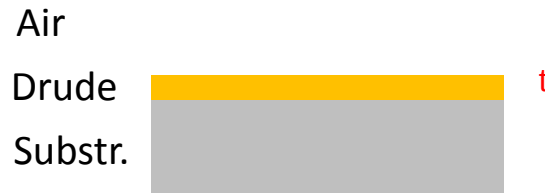
Transmission vs angle (p-pol):
A sharp dip is observed in transmission, where $\epsilon \sim 0$



(“Berreman” dip)



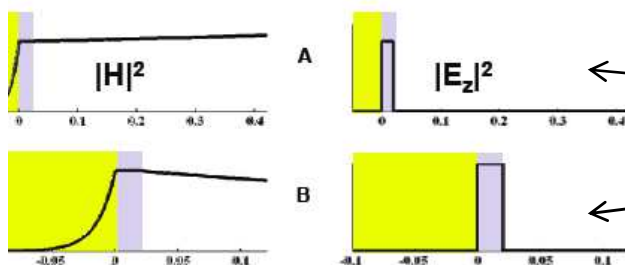
“Epsilon Near Zero” vs Berreman Modes



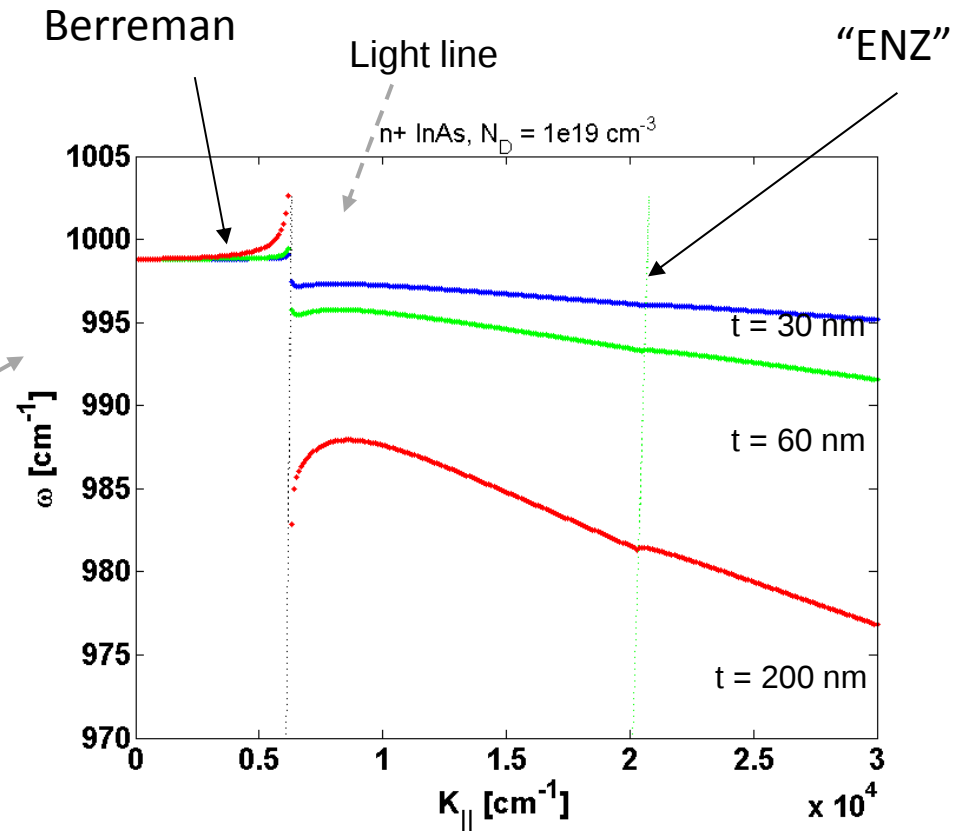
Dispersion relation of the TM mode in a 3-layer

$$\left(1 + \frac{\epsilon_1 k_{z,1}}{\epsilon_3 k_{z,3}}\right) = i \tan(k_{z,2}d) \left(\frac{\epsilon_2 k_{z,3}}{\epsilon_3 k_{z,2}} + \frac{\epsilon_1 k_{z,2}}{\epsilon_2 k_{z,1}}\right)$$

Numerical solution in complex domain
(assuming complex freq, real k)



- Berreman mode (leaky mode),
- ENZ mode (bound mode)
- ENZ mode is a special type of a plasmon polariton

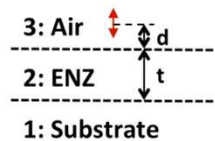




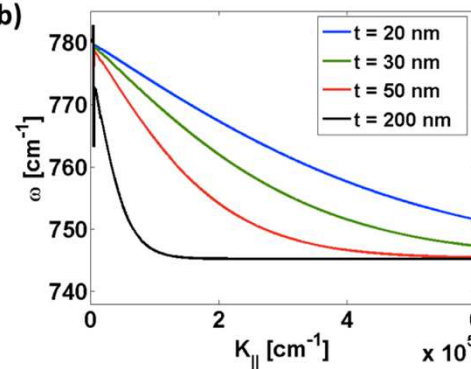
MM Resonators Provide k-vectors to Excite ENZ Modes

ENZ dispersion

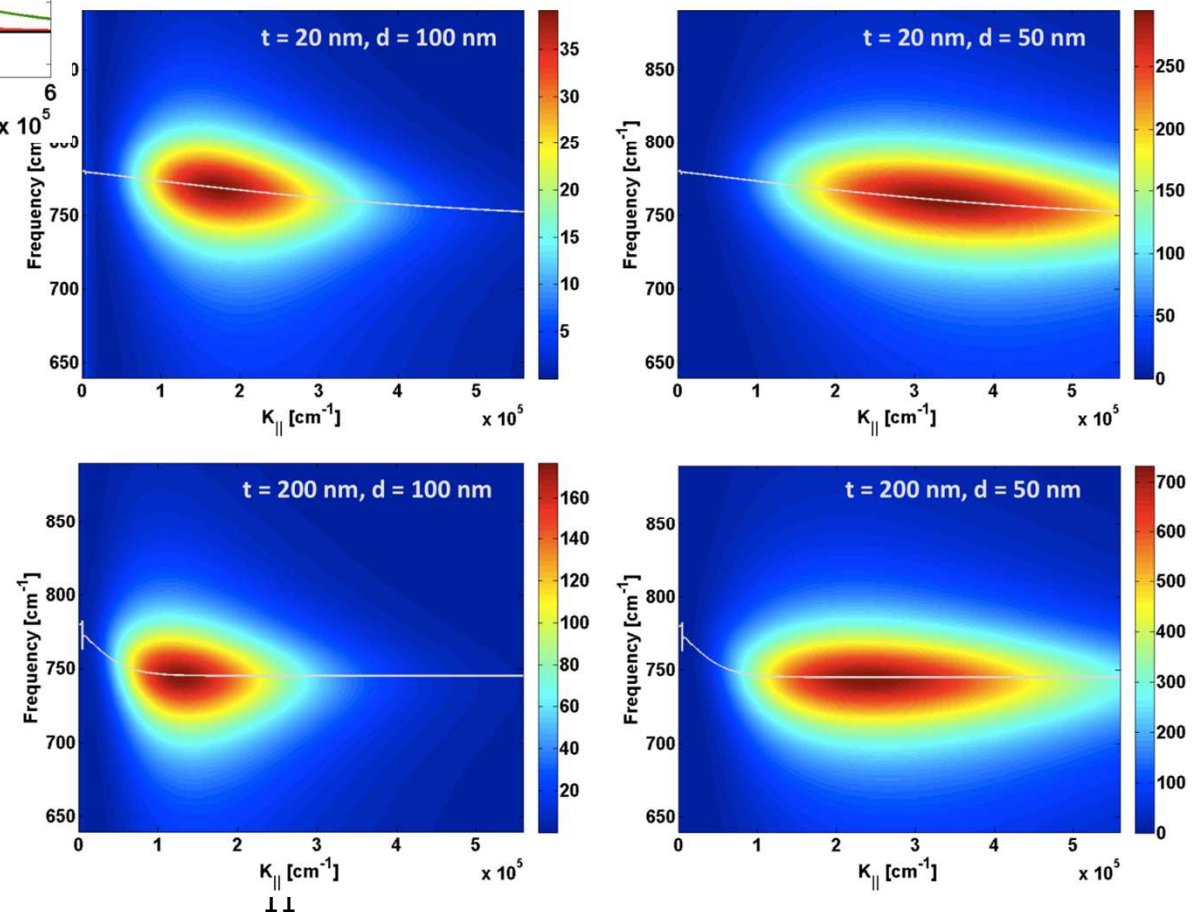
(a)



(b)

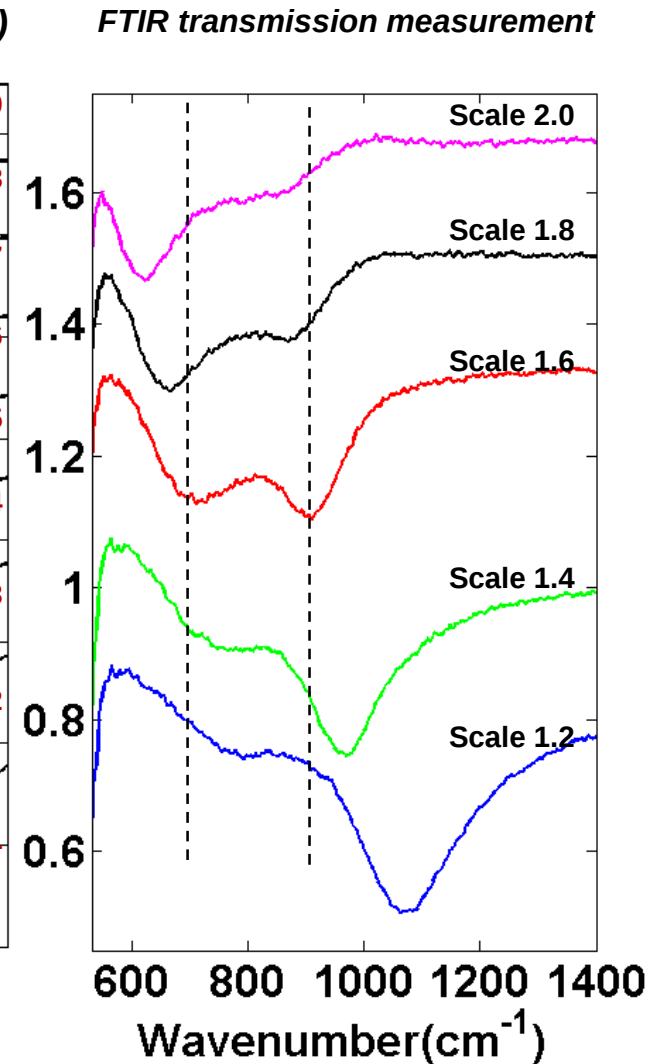
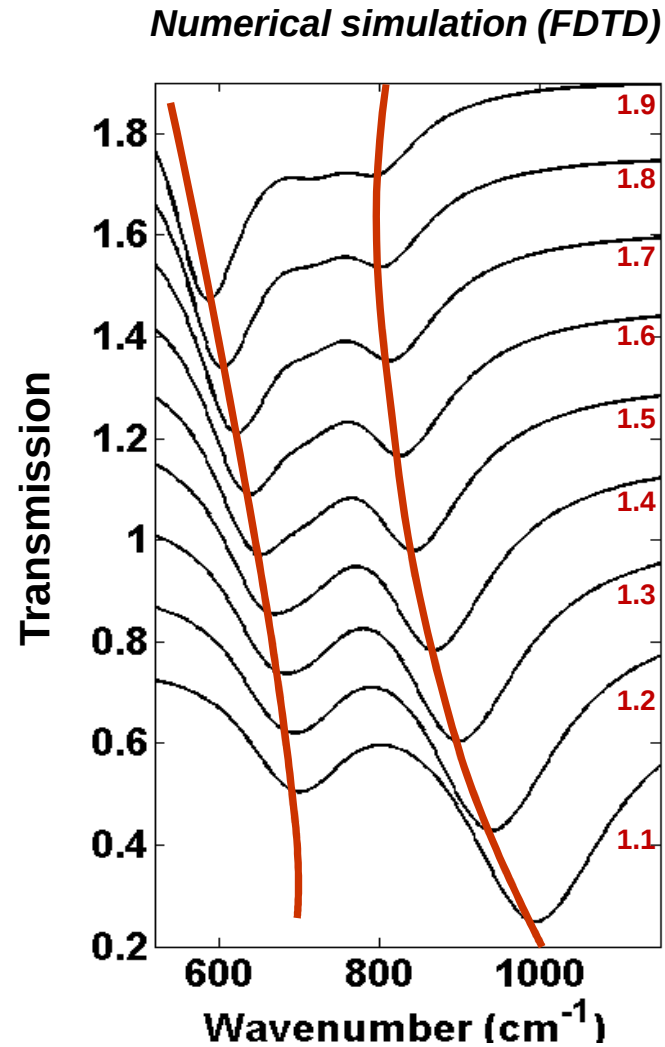
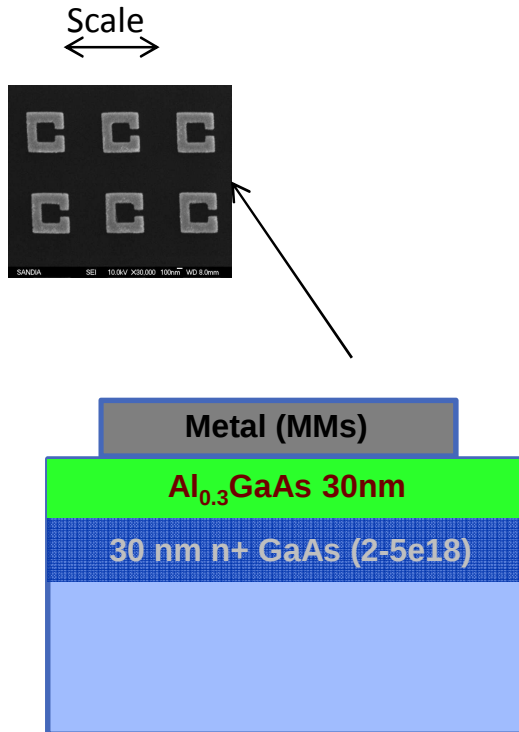


Superimpose dipole emission spectra on
the ENZ dispersion



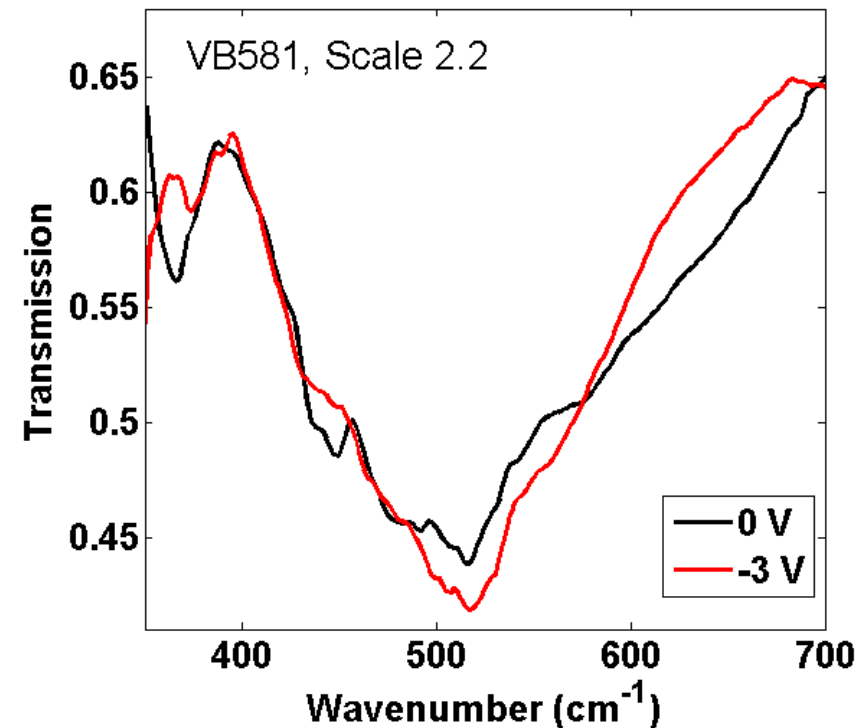
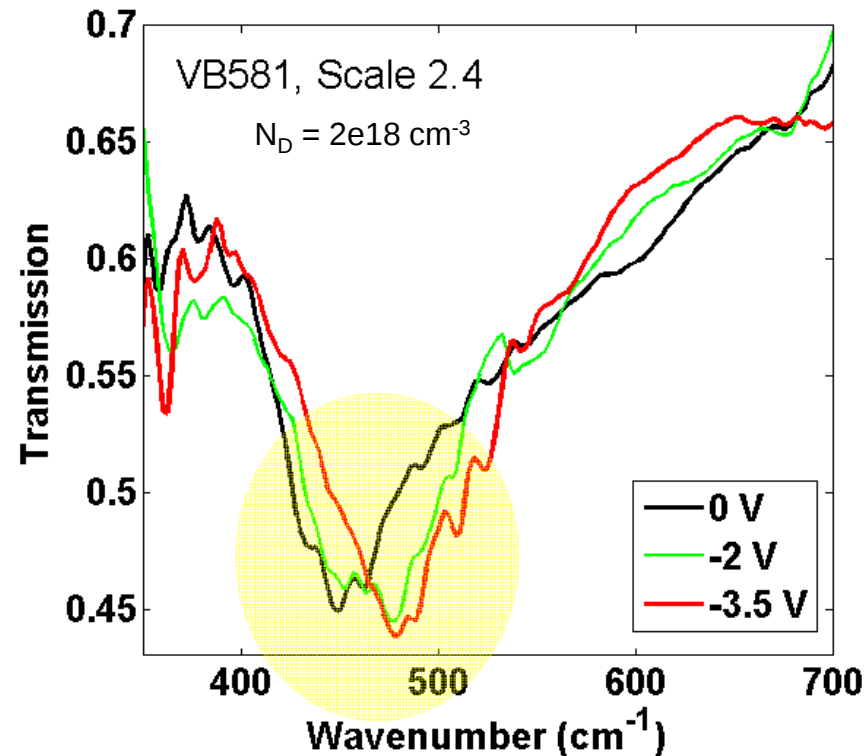


Strong Coupling to ENZ Mode: Theory vs. Experiment





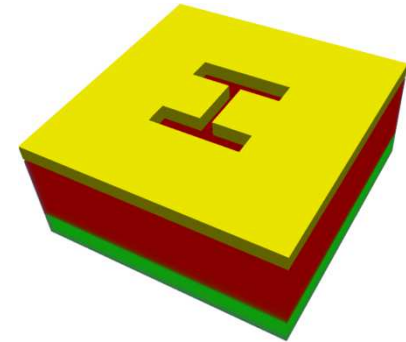
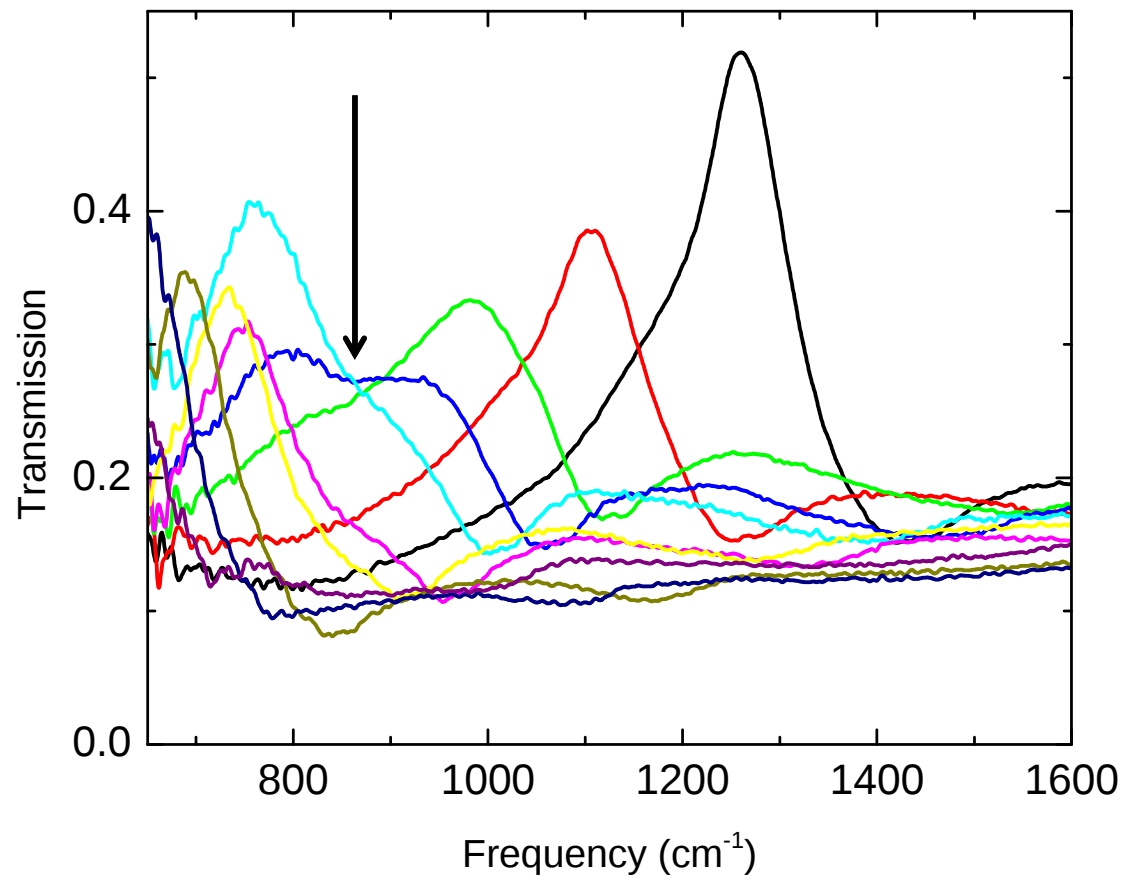
Electrically Tuning the Coupling to the ENZ Mode



***Fundamentally different than tuning just by changing a local permittivity!
Removal of carriers -> removal of ENZ mode***



Inverted Dogbone + ENZ



Tunable Passband Filter

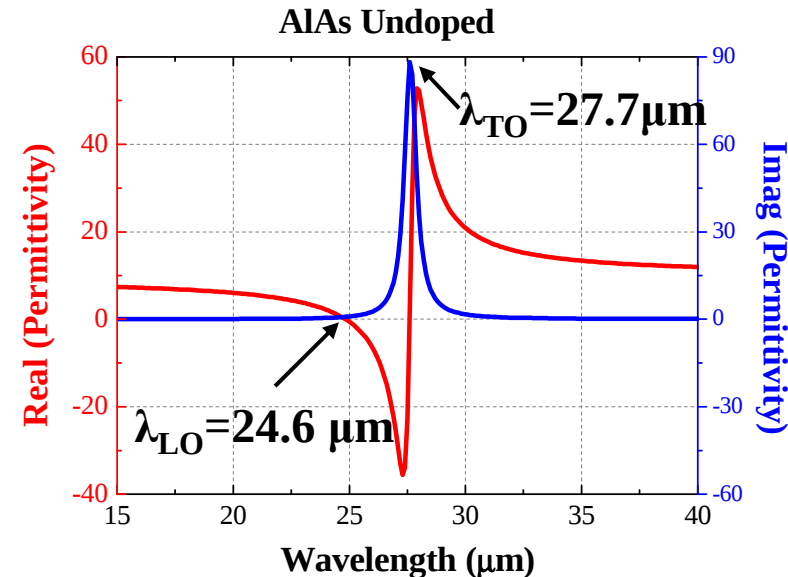
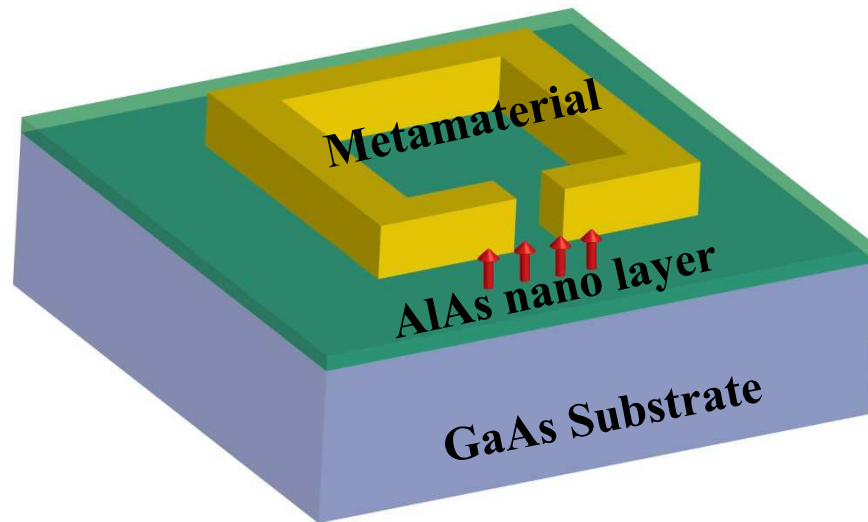


Outline

- **Strong Coupling to Intersubband Transitions**
 - Fundamentals
 - Voltage Tuning
 - Wavelength scaling
- **Strong Coupling to Epsilon-near-zero modes**
 - Fundamentals
 - Voltage Tuning
- **Combining coupling to ENZ modes and Intersubband Transitions**



Coupling of Metamaterial Resonators to Dipole Resonances + ENZ modes



Lorentz model

$$\epsilon = \epsilon_{\infty} \left(1 + \frac{\omega_L^2 - \omega_T^2}{\omega_T^2 - \omega^2 + i\omega\Gamma} - \frac{\omega_p^2}{\omega(\omega - i\gamma)} \right) = (\eta - i\kappa)^2$$

Γ : Phonon damping constant

ω_p : Plasma frequency

γ : Free carrier damping constant

One thin layer contains:

- $\epsilon \sim 0$ crossing $\rightarrow$ ENZ mode
- TO Phonon: “dipole” transition

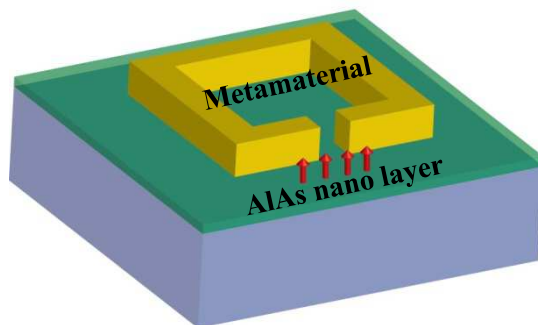
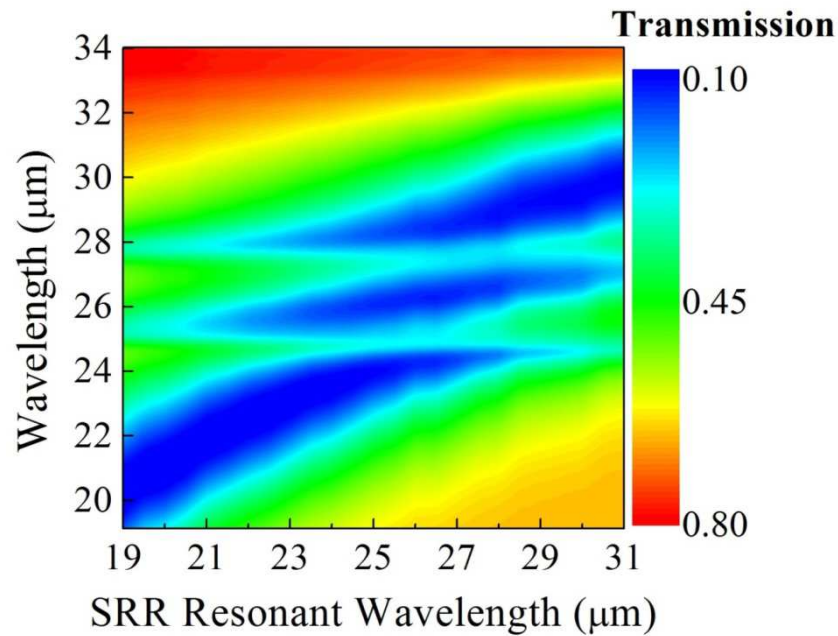
J. Appl. Phys., **48**, 212 (1976).

Solid State Communications, **62**, 645-647, (1987).



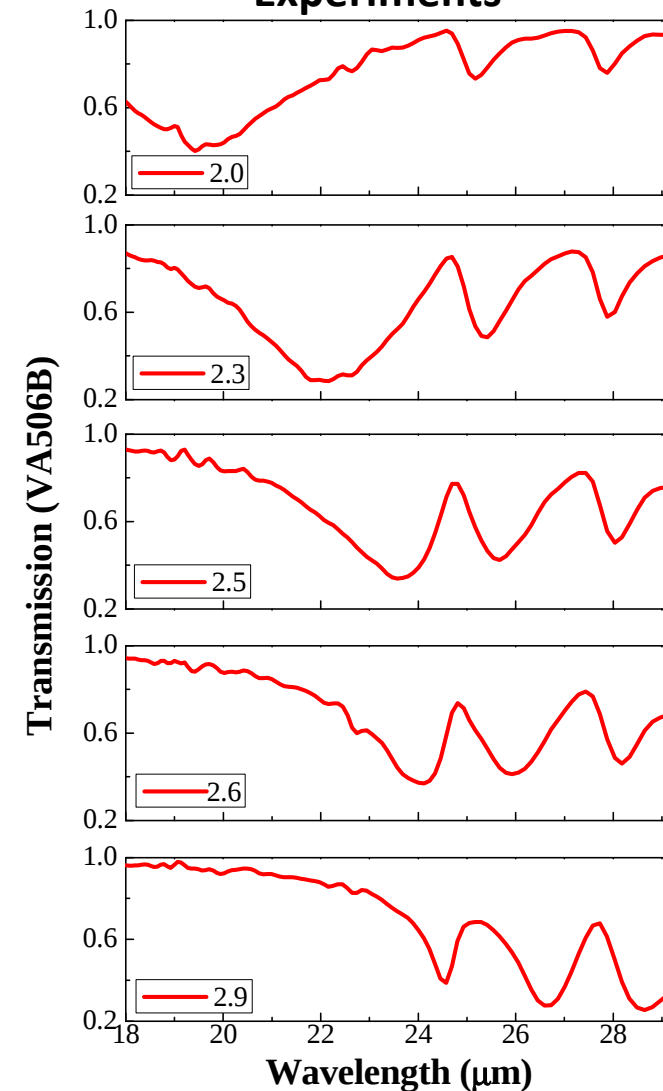
Strong Coupling: 2 Polariton Branches

Simulations



Increasing SRR size

Experiments





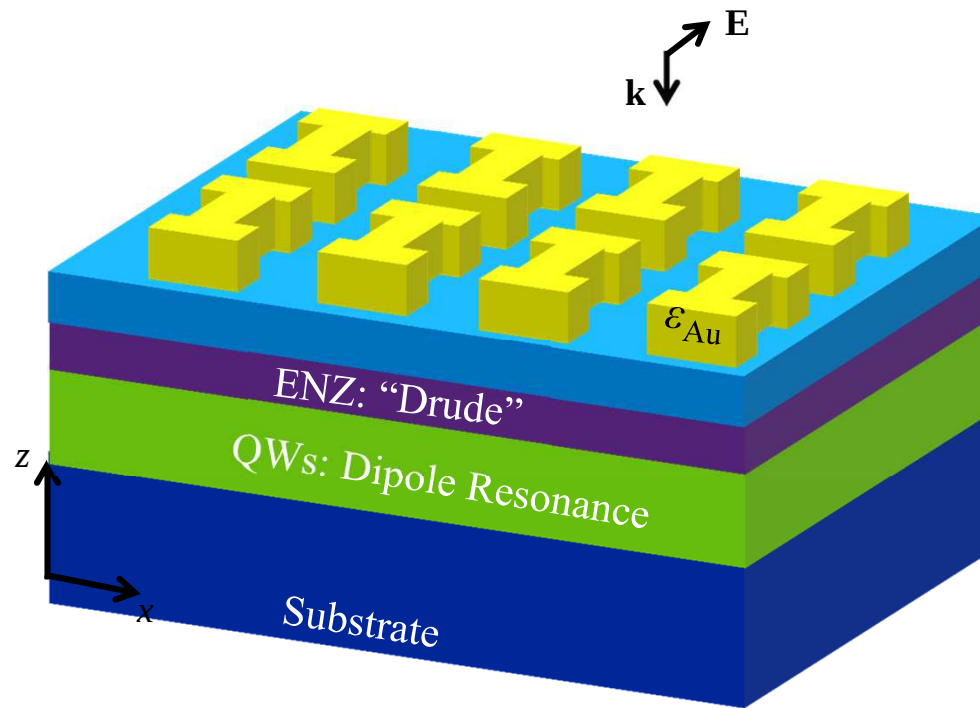
Let's Separate the ENZ Layer and the Dipole Resonance Layer

3 Interacting Systems

MM
Resonators

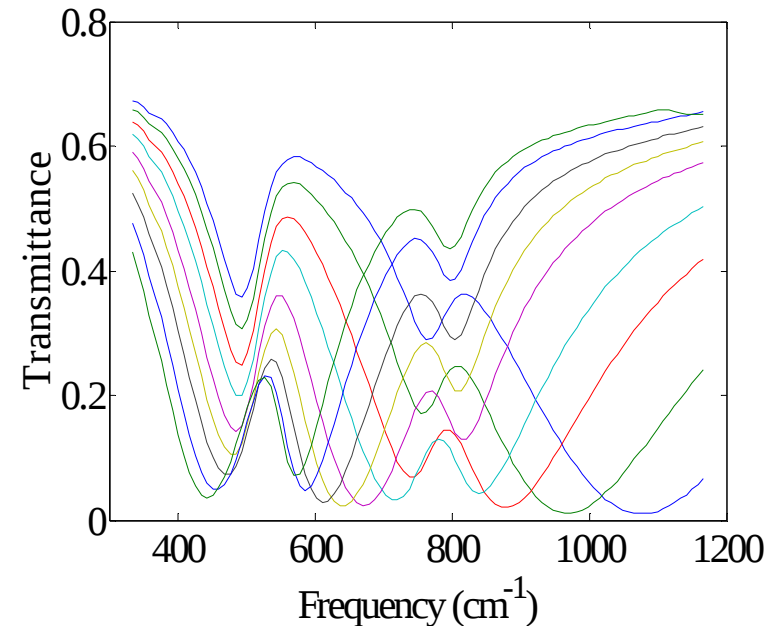
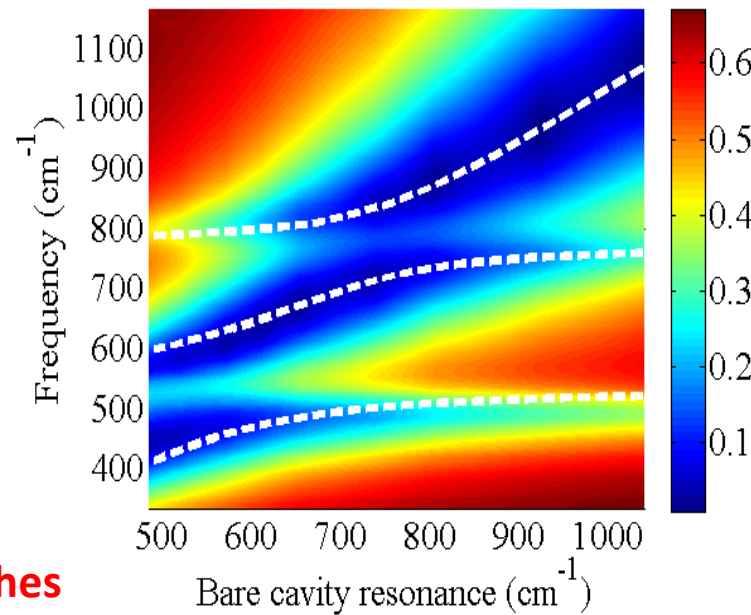
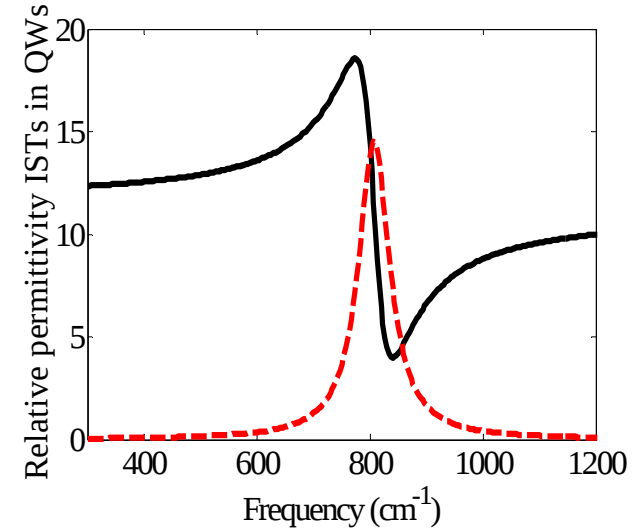
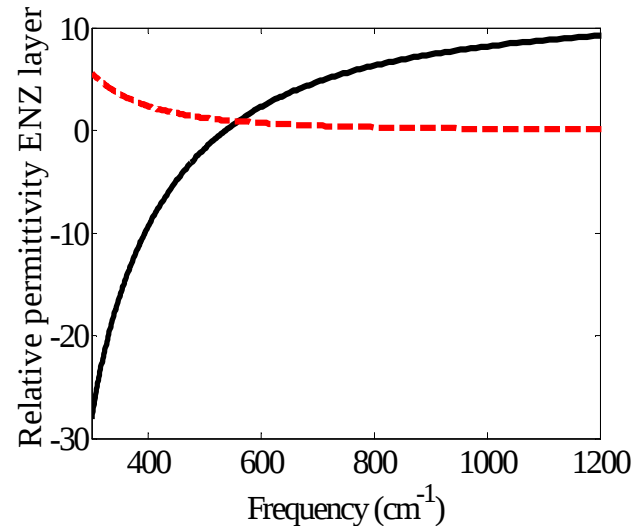
ENZ Mode
(Drude layer)

Dipole Transition
(ISTs in QWs)





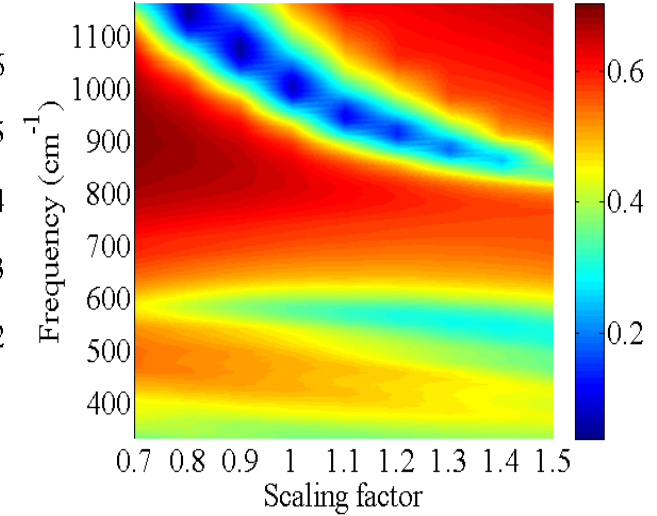
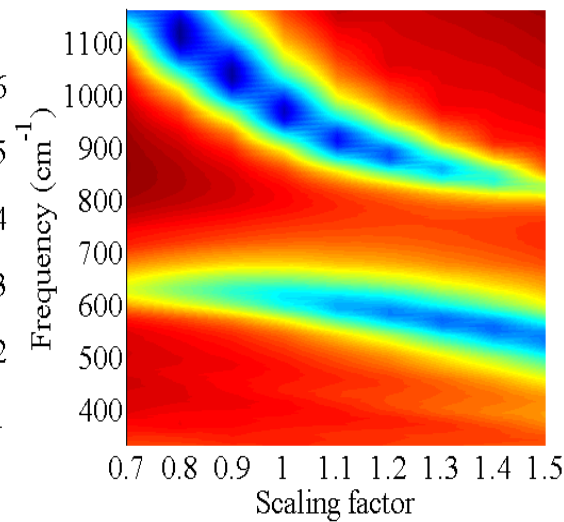
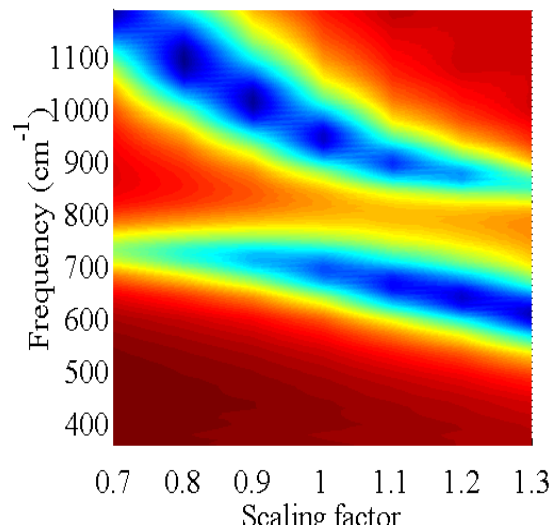
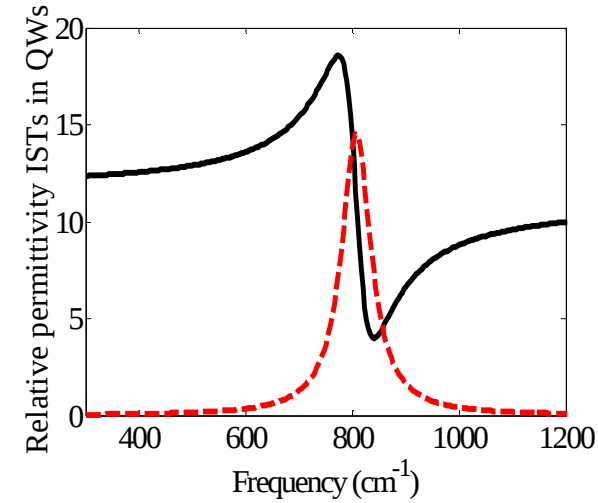
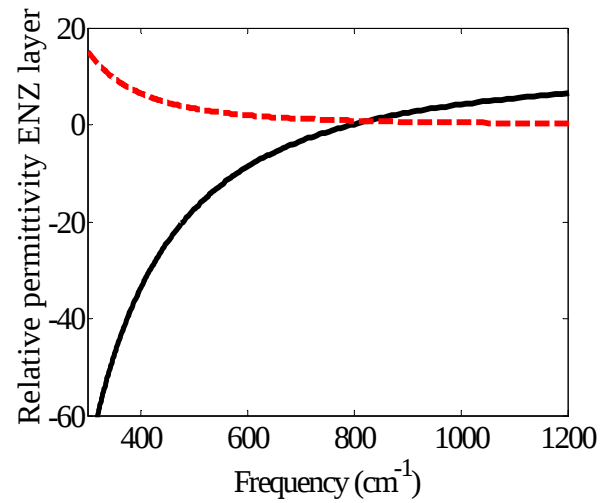
Case 1: Different Frequencies for Dipole Transition and ENZ



**2 Different
Polariton Branches**



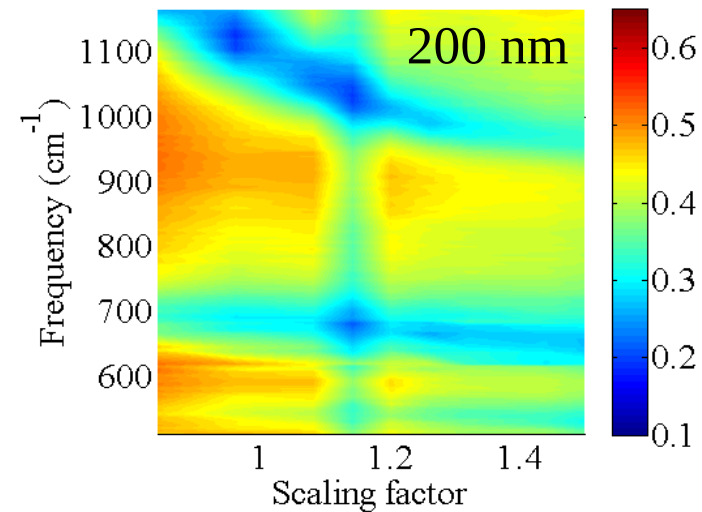
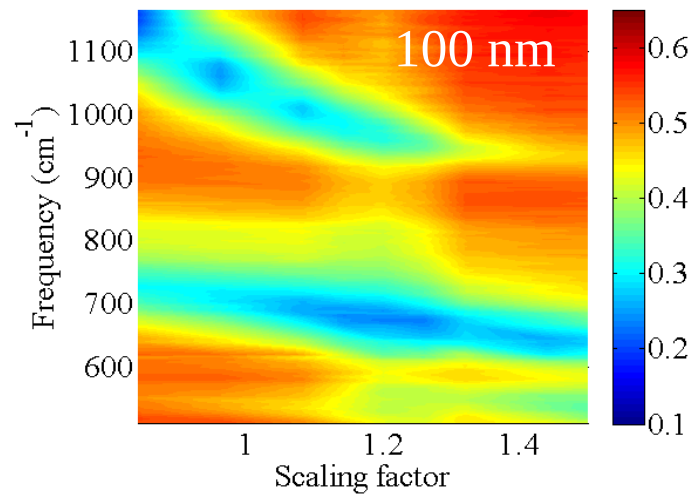
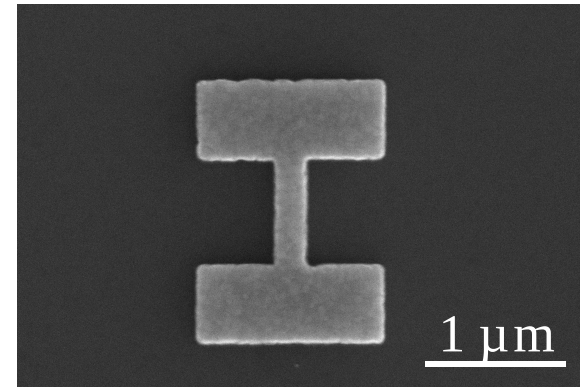
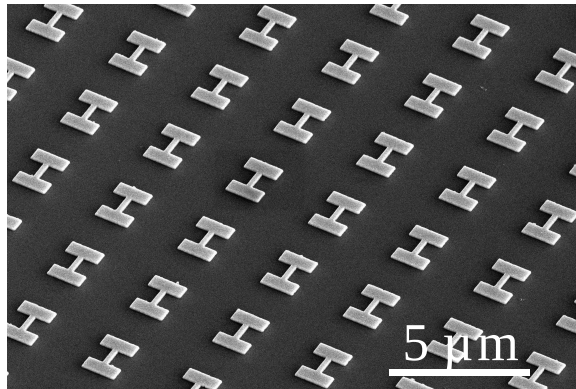
Case 2: Same Frequencies for Dipole Transition and ENZ



Only 1 Polariton, larger Rabi Splitting!

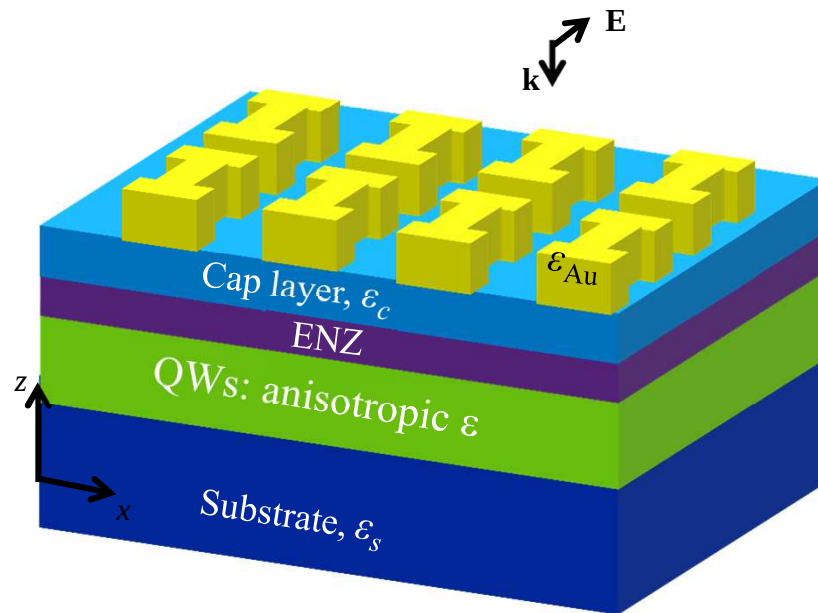


Experimental verification for 100 nm and 200 nm ENZ + QWs

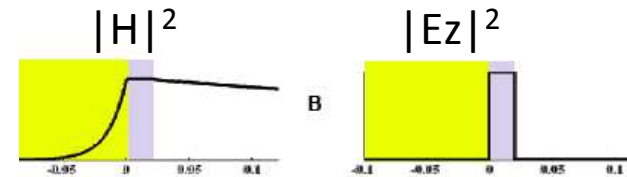




Combining ENZ Modes and Intersubband Transitions for Stronger Coupling

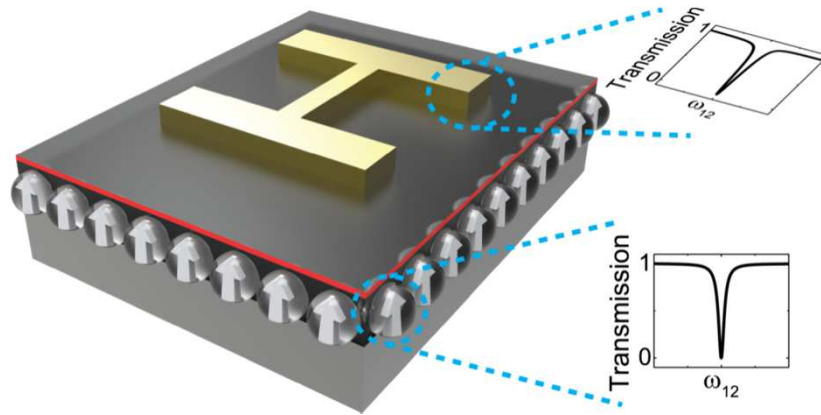


- Field decays exponentially with distance from MM resonators
- An ENZ thin layer acts as a “transducer” of E_z field





Summary & Acknowledgments



Strong coupling to
Intersubband Transitions

Postdocs: Young Chul Jun (n+),



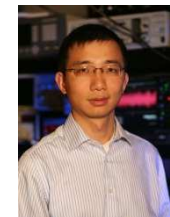
Salvo Campione,



Alex Benz (QWs),



Sheng Liu (TDS),



CINT + Sandia: John Reno, Mike Sinclair, Ines Montano, John Klem (InGaAs), Andy Allerman (GaN)

UC Irvine: Filippo Capolino

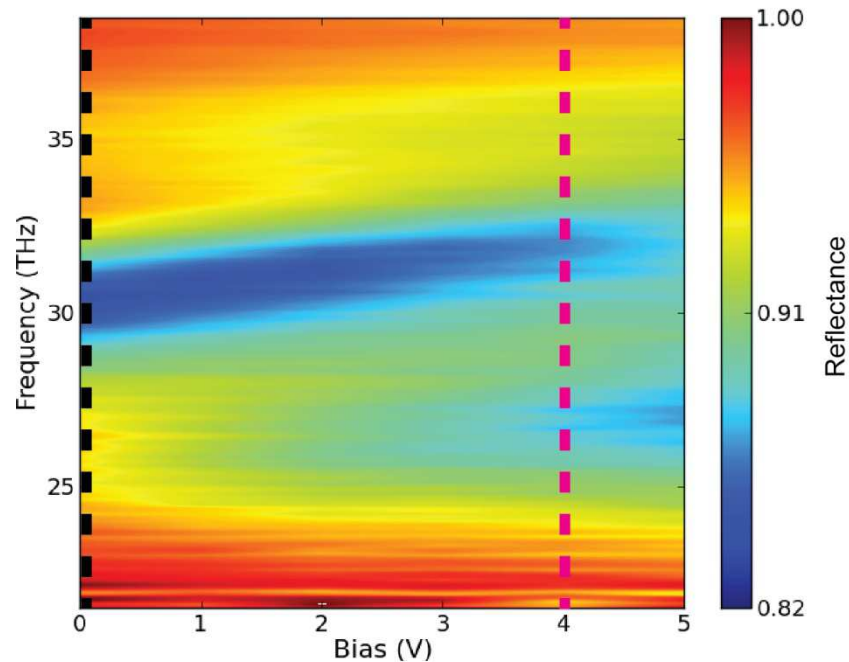
France: J.J. Greffet, F. Marquier (ENZ)



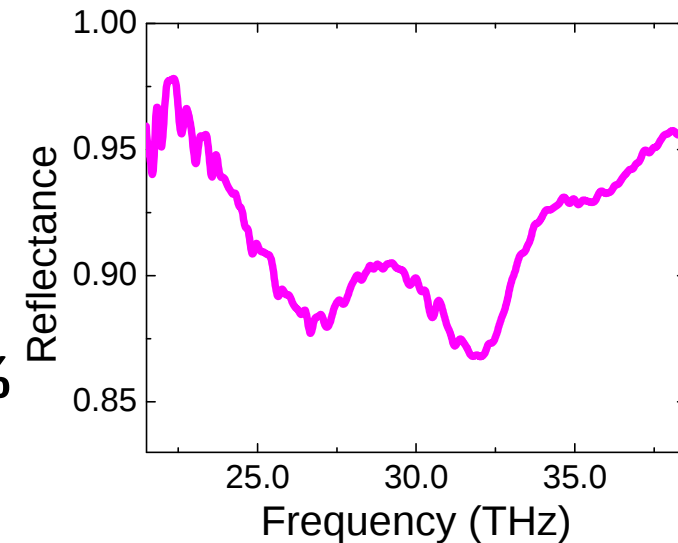
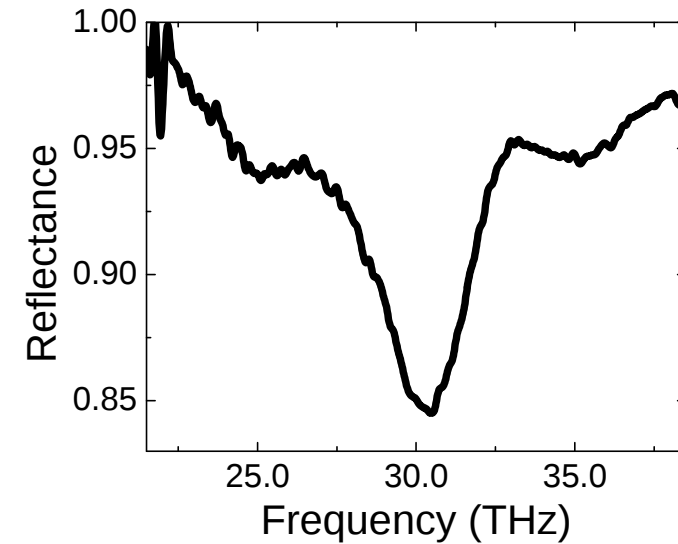
SPARES



Electrically Tunable Strong Coupling

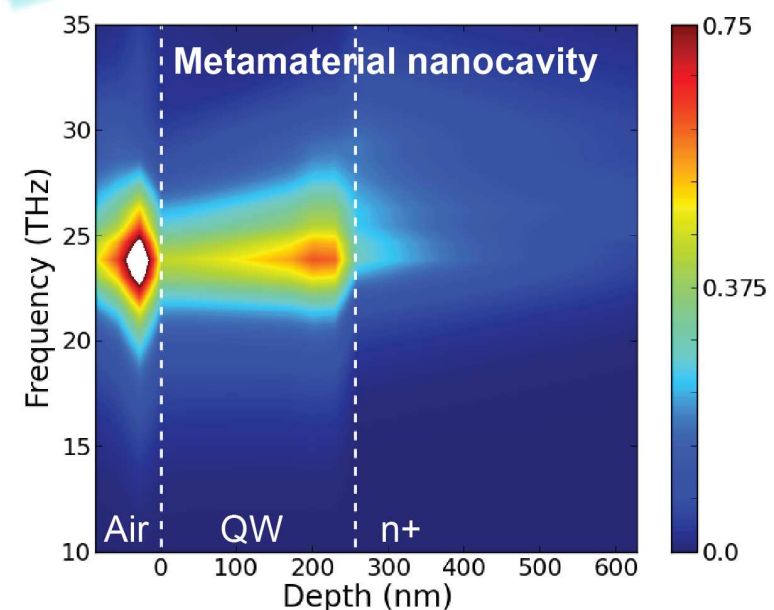
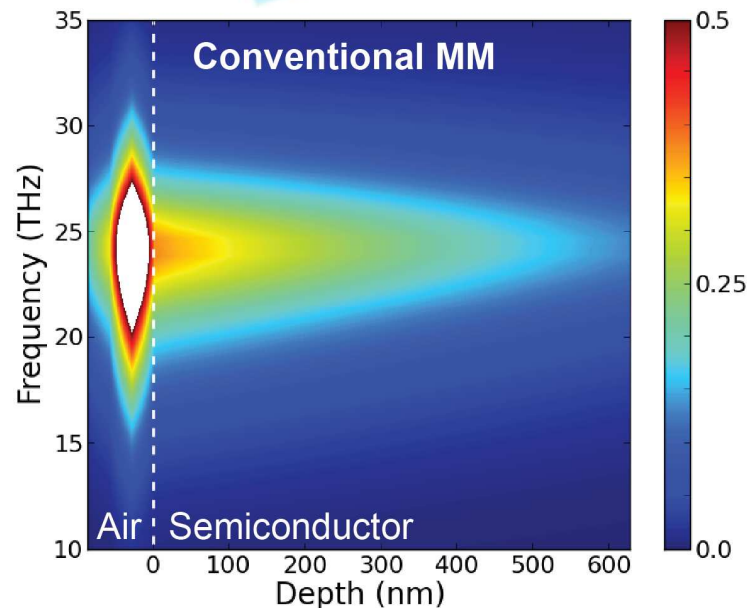
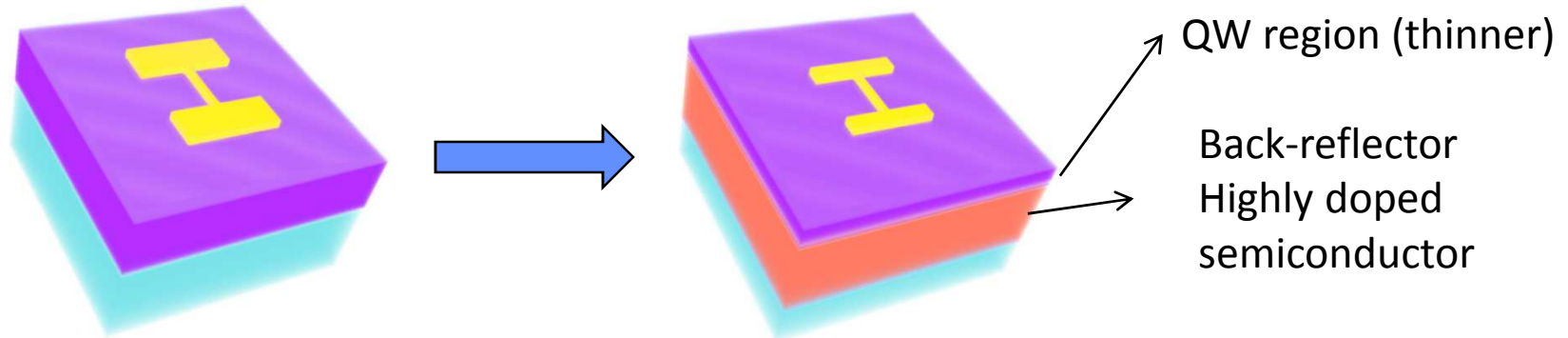


- Metamaterial spectral response tuned by voltage
- Upper polariton shifted by 2.5 THz (=8% of center frequency or one full intersubband linewidth)





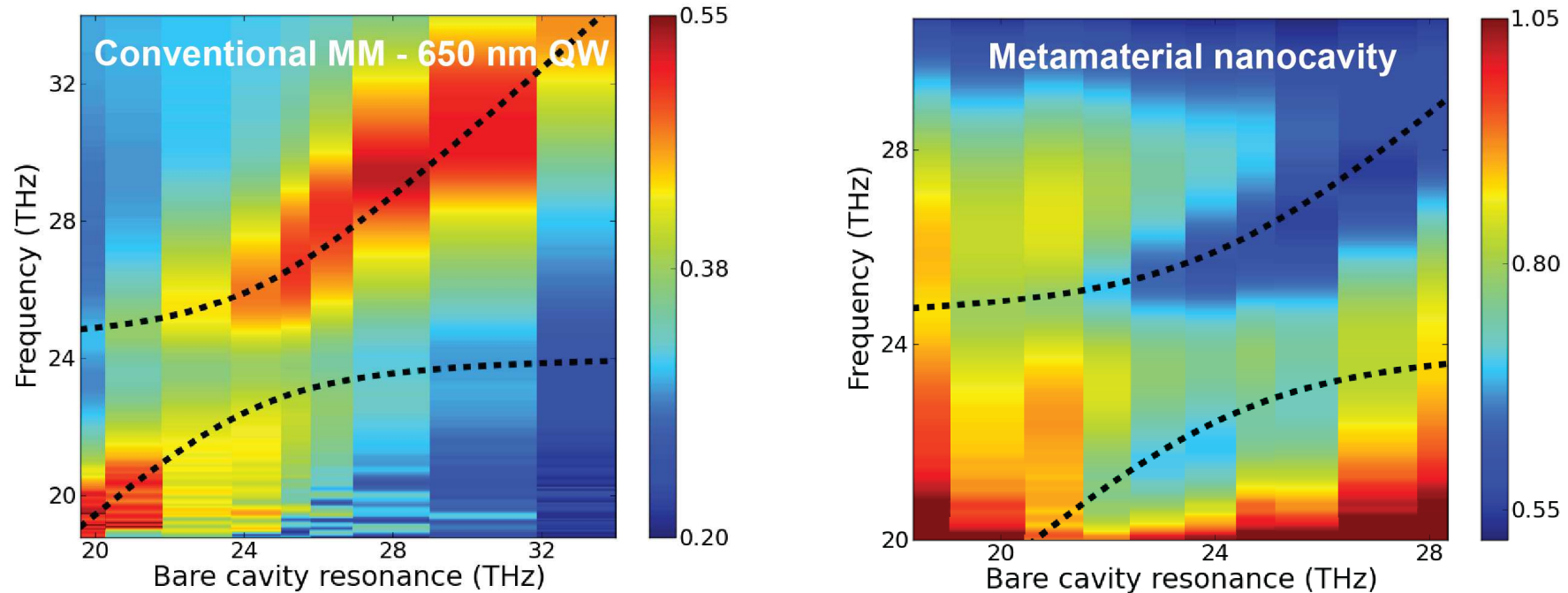
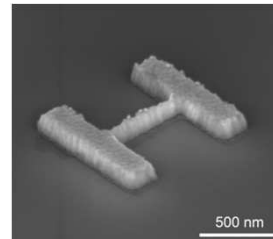
How to Increase the Rabi Frequency: Metamaterial Nanocavity + Ground-plane



- Ground plane limits field penetration into the semiconductor
- Mode volume reduced from: $2.49 \times 10^{-3} (\lambda/n)^3$ to $1.34 \times 10^{-3} (\lambda/n)^3$



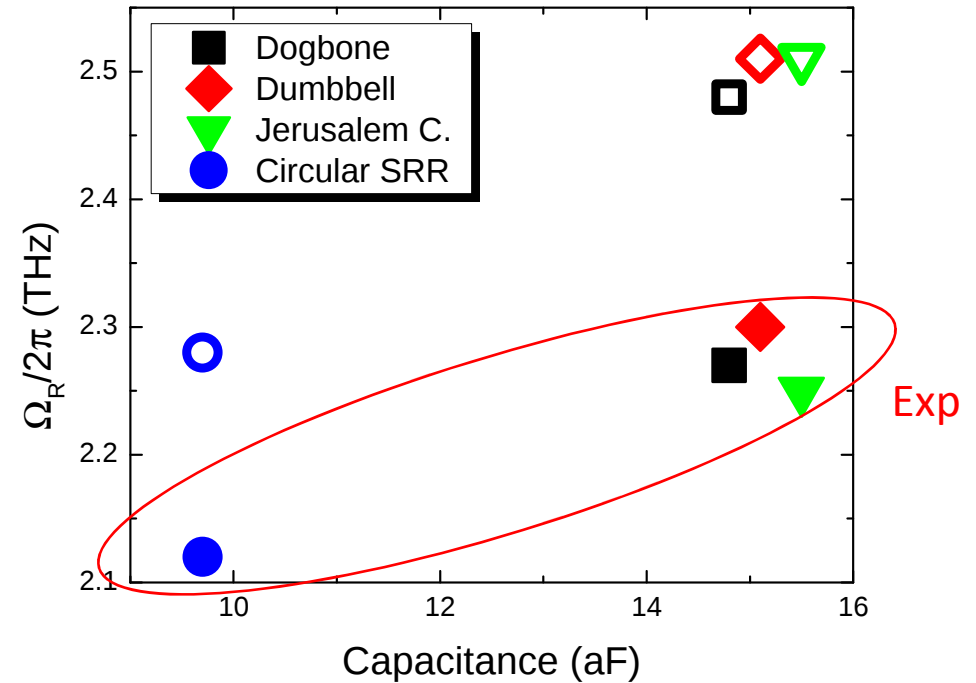
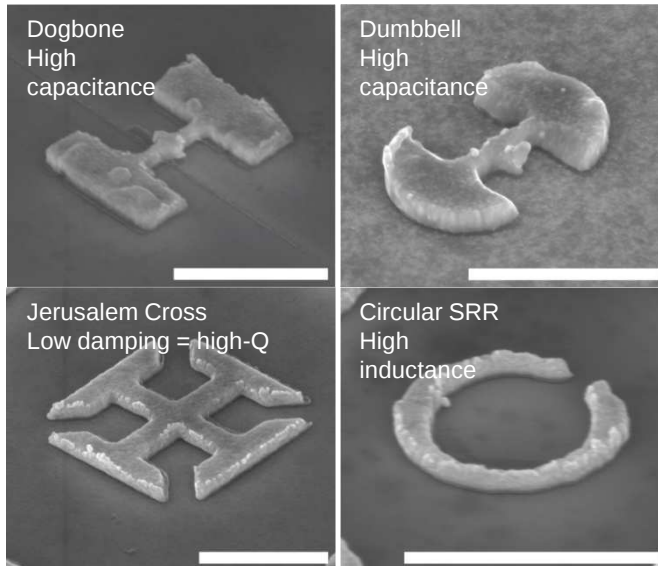
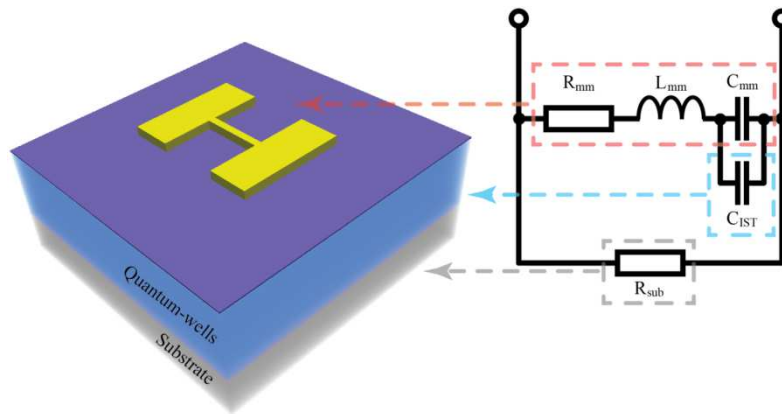
Reduced Mode Volume: Experiments



- Rabi frequency remains almost identical despite reduced QW-stack thickness (factor 3.3) or reduced mode volume (factor 1.9)
- Conventional sample: 2.1 THz vs. Nanocavity: 2.5 THz



Rabi Frequency vs. Geometry



Physical Review **B89**, 165133 (2014)

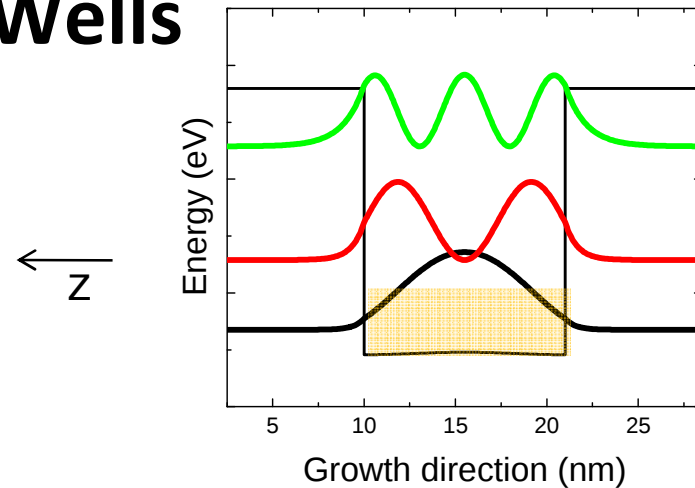
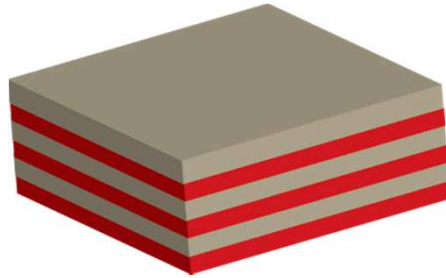
- Larger capacitance leads to larger Rabi splitting



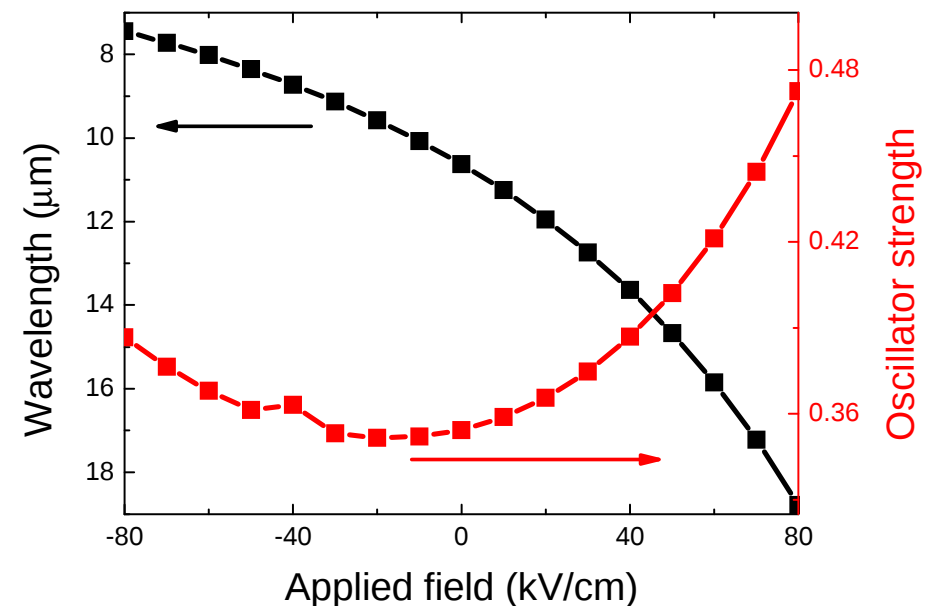
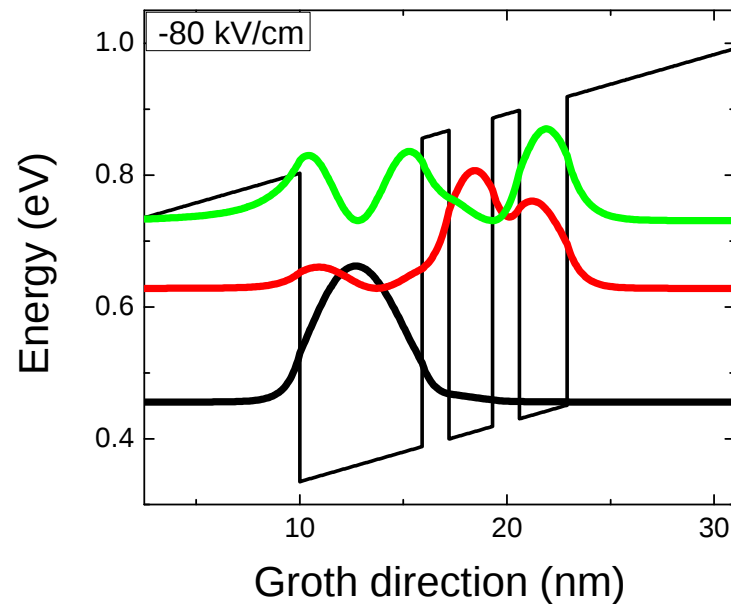
Strong Coupling to Inter-subband Transitions in Quantum Wells

- Scalable (far IR to near IR), Mature, Versatile

Opt. Express 20, 6584 (2012),
APL 98, 203103 (2011)

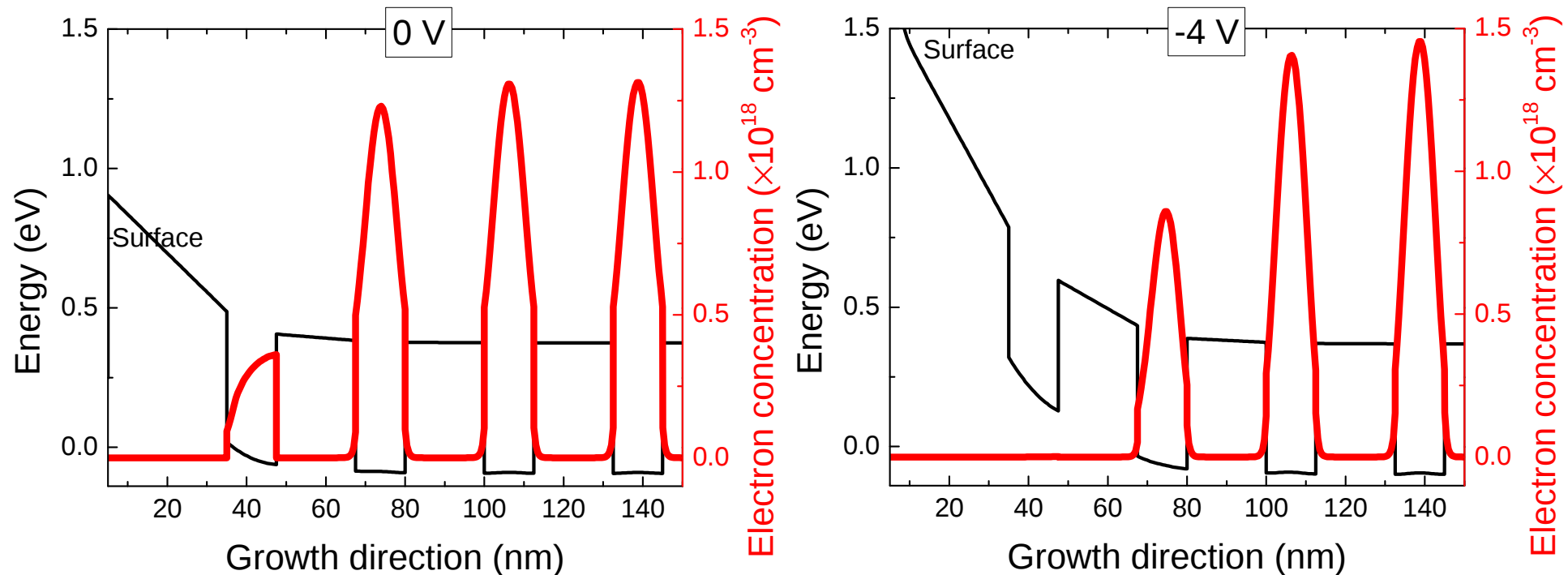


Change levels using an electric field





Reverse Biasing Highly Doped QWs: Problem!

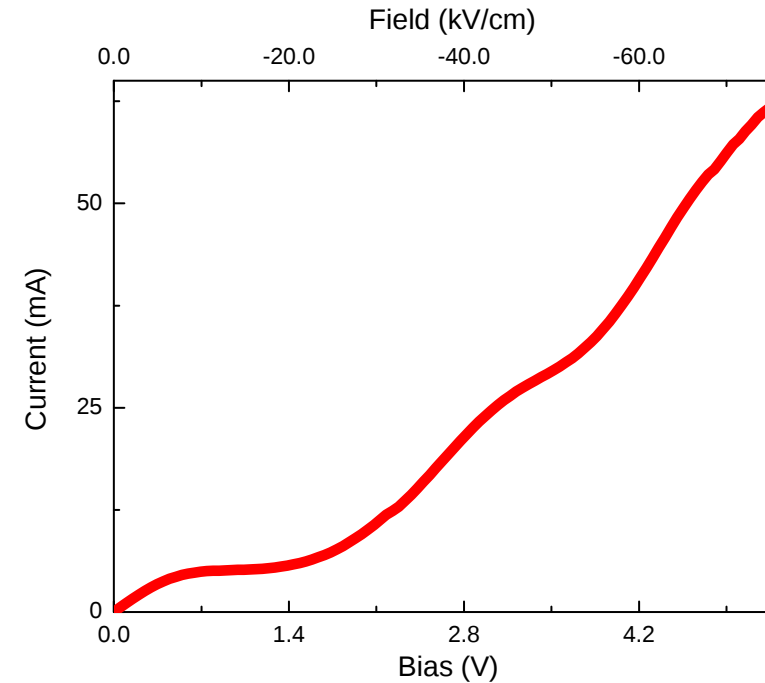
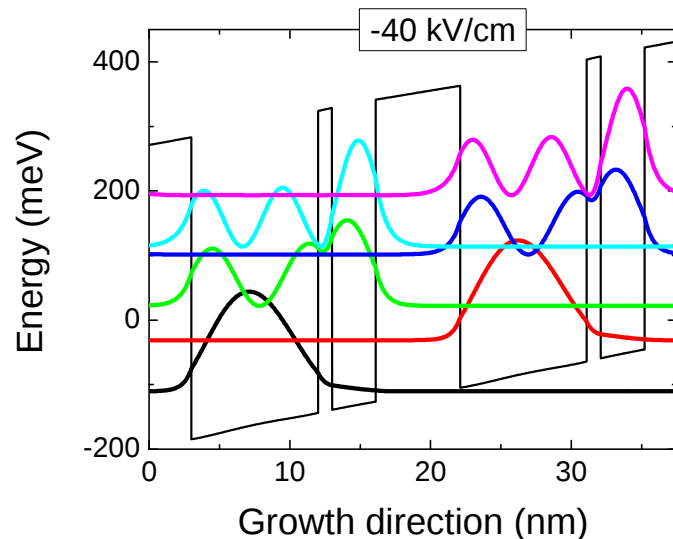
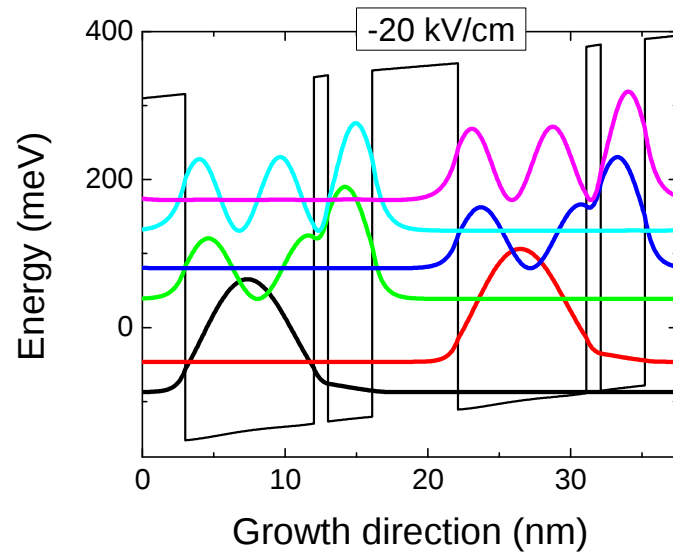


- Applying reverse bias across QW-stack depletes well by well => no continuous tuning
- High doping concentration limits effect to first few wells only



Stabilizing Internal Electric Field

Allow a small forward bias current

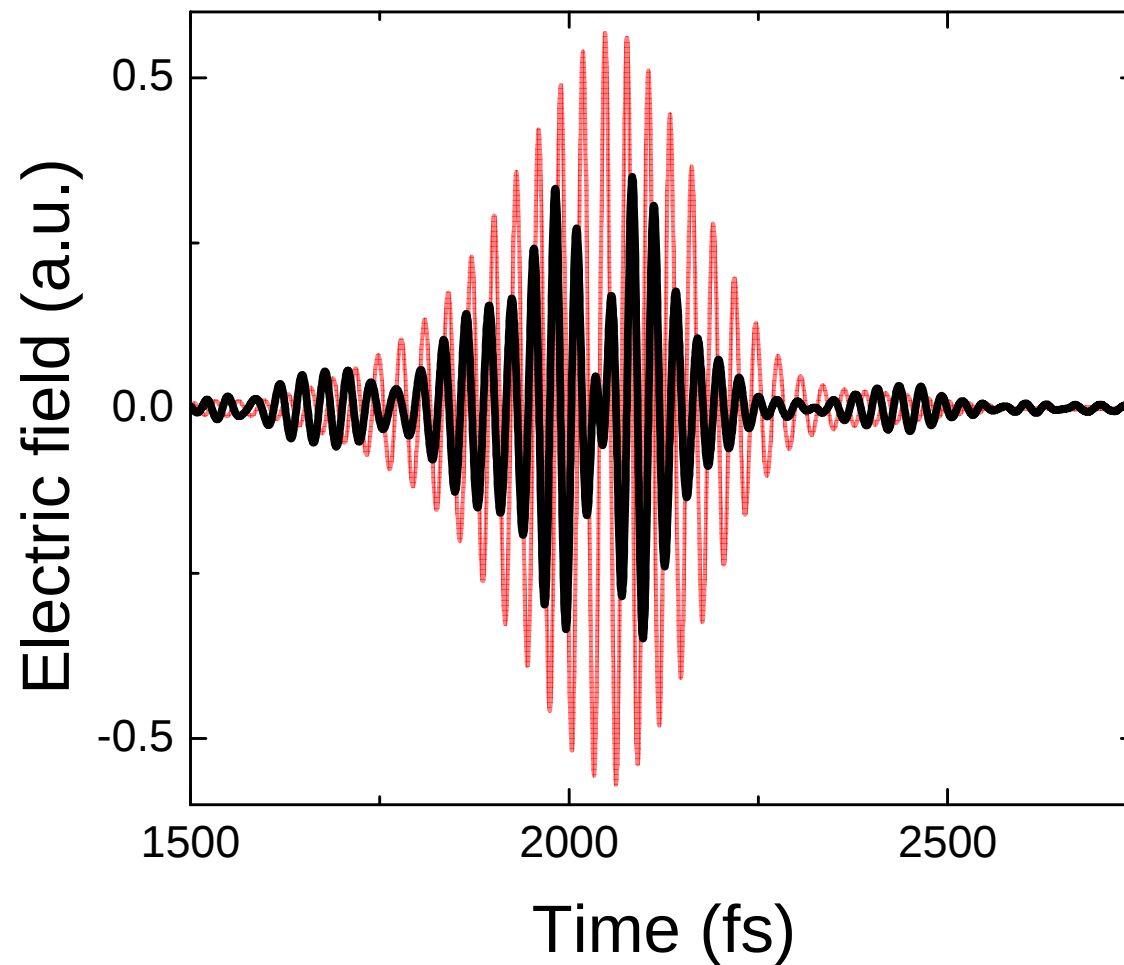


- Current flow through QW-stack stabilizes internal electric field
- Alignment and miss-alignment of subbands show up in IV curve



Experimental Rabi Oscillations

(using a mid-IR Time-domain fsec system)

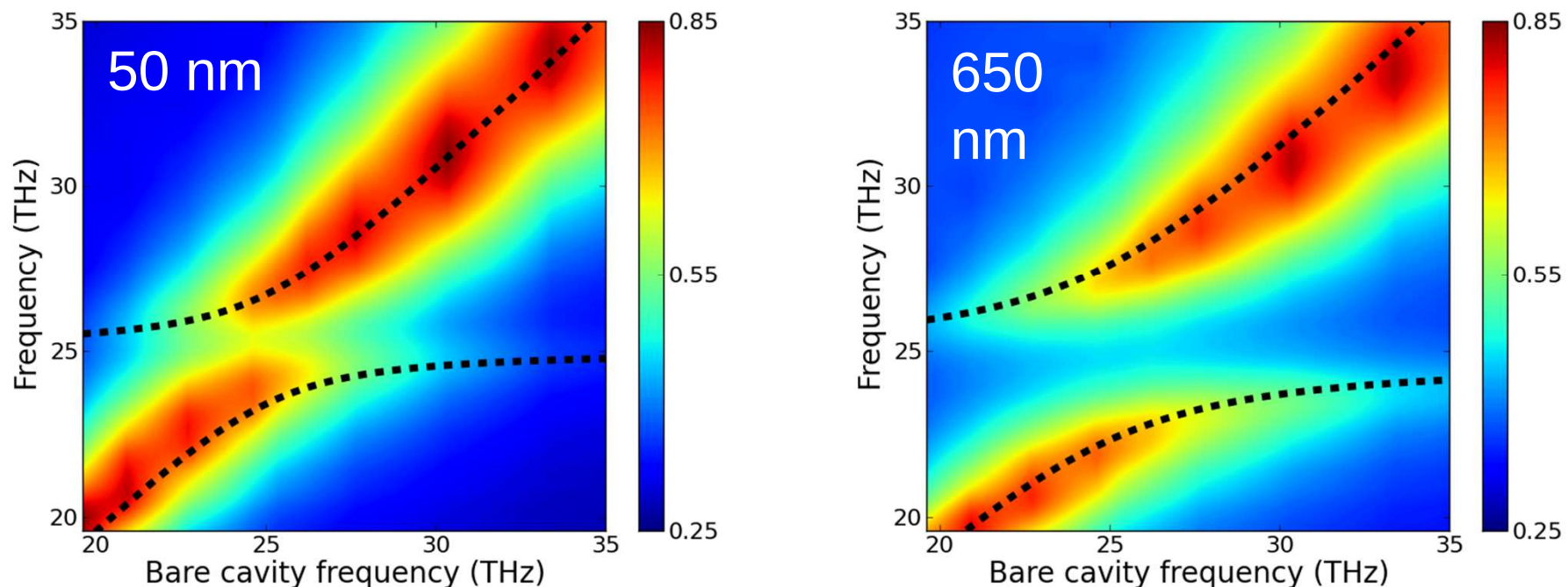


- **Energy exchange probed in time**
 - 33 fs oscillation
 - 480 fs beating
- **System strongly coupled**
- **Splitting of 4.2 THz measured**
 - 15 % of ω_{12}



How to Increase the Rabi Frequency

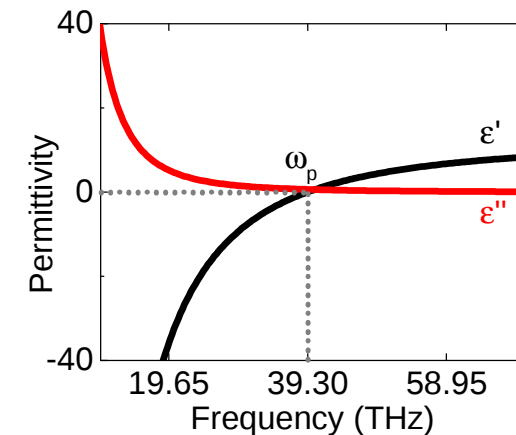
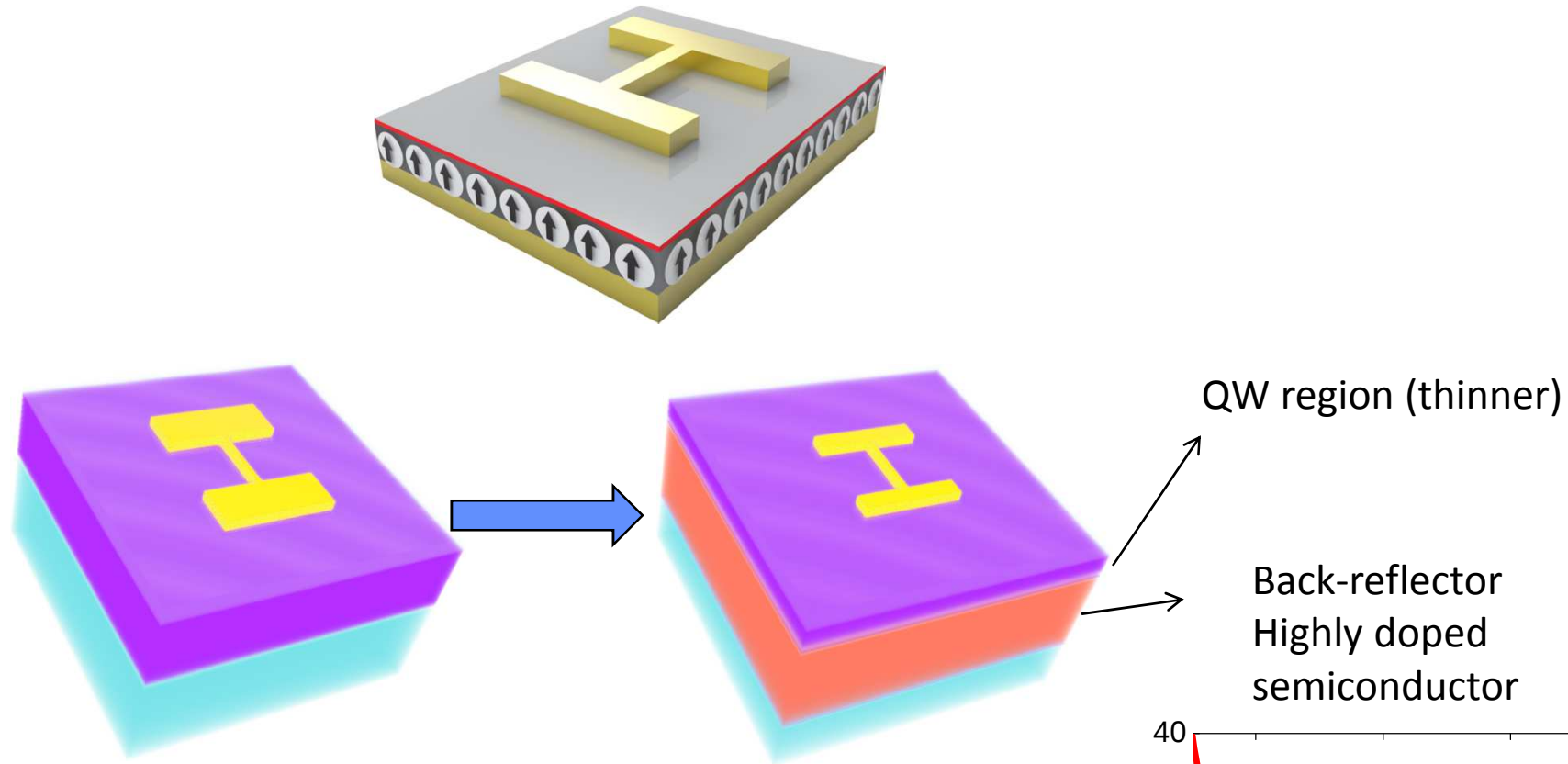
Effect of QW-stack thickness



- Thickness of the QW-stack influences the Rabi frequency
- Increased from 1.65 to 2.64 THz

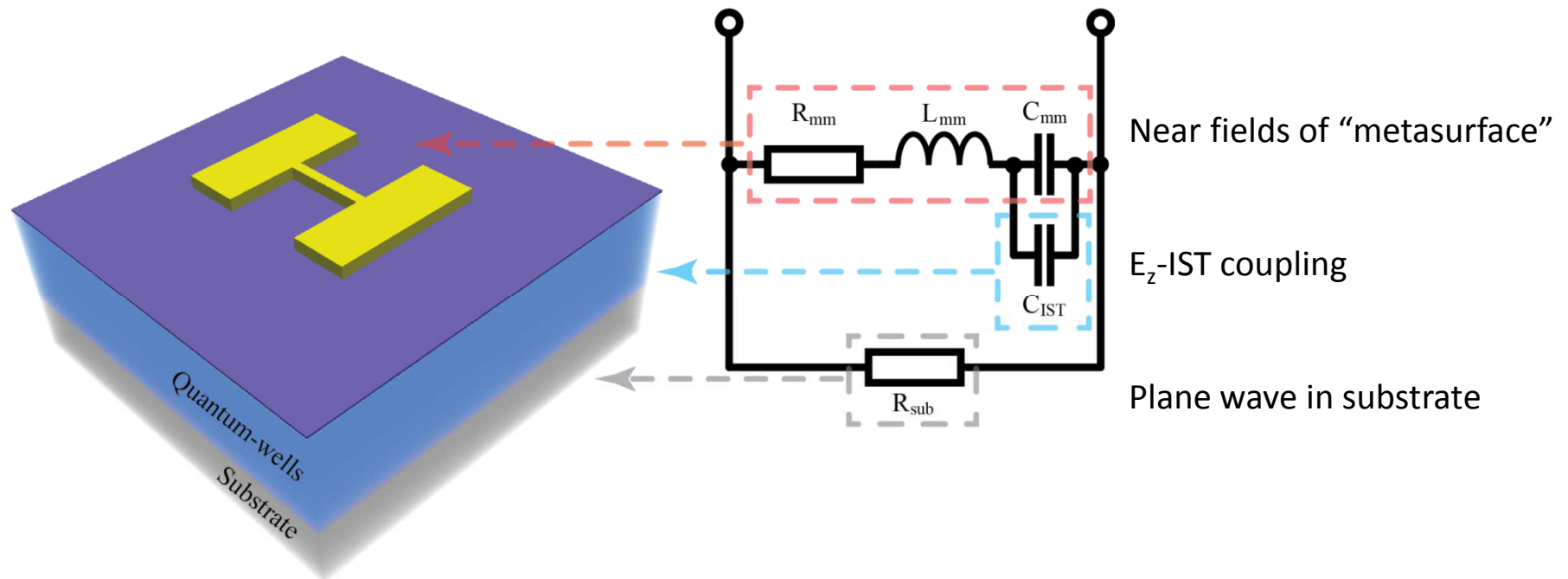


How to Increase the Rabi Frequency: Metamaterial Nanocavity + Ground-plane





Circuit Model for Strong Coupling



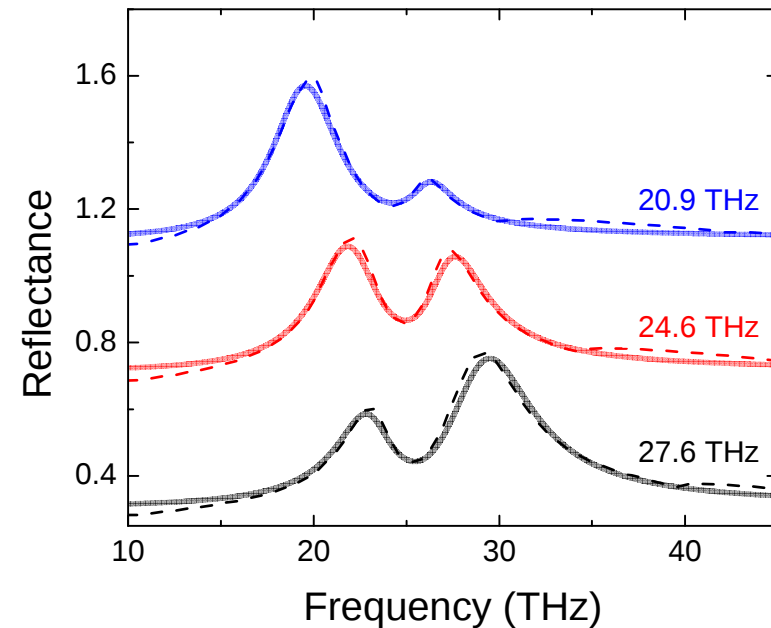
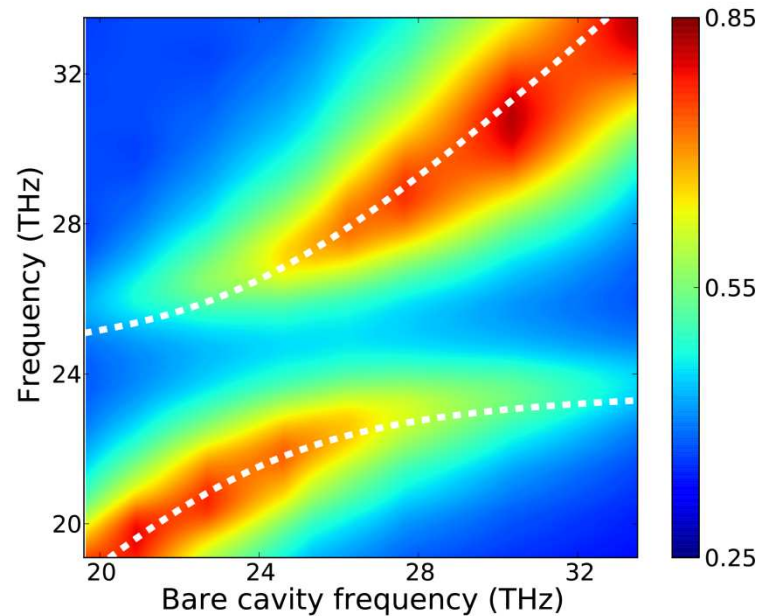
- Metamaterial modelled as simple RLC harmonic oscillator
- Strong coupling represented by dispersive capacitor
- No free fitting parameters for dispersive capacitor

$$\epsilon_{IST}(\omega) = \epsilon_{back} + \frac{\Delta N_{12} e^2}{\epsilon_0 m^*} \frac{f_{12}}{\omega_{12}^2 - \omega^2 + 2i\omega\delta}$$

$$C_{IST}(\omega) = \frac{\sqrt{\epsilon_{back}\epsilon_{IST}(\omega)} - \epsilon_{back}}{\epsilon_{back} + 1} C_{mm}$$



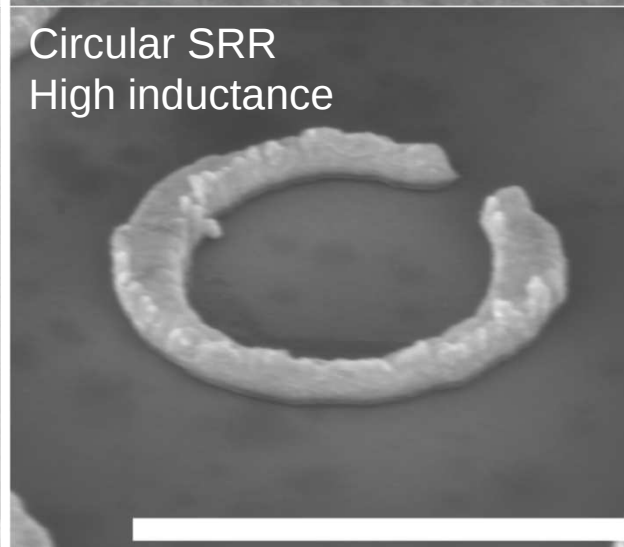
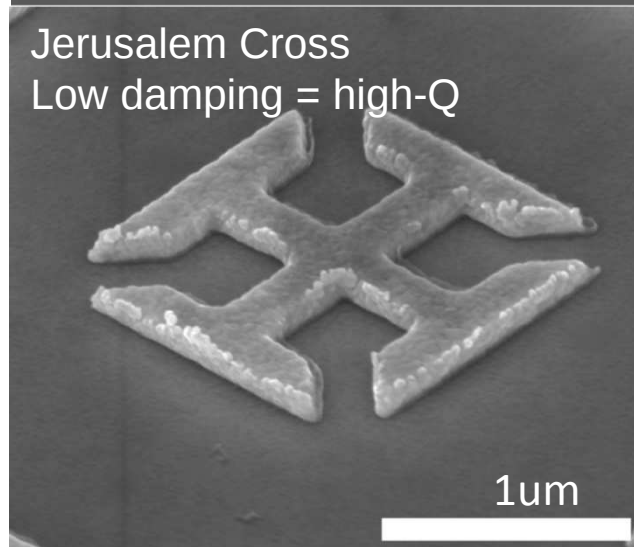
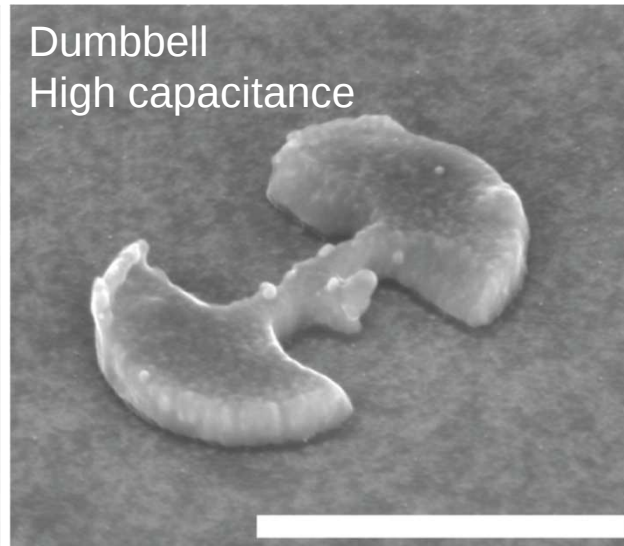
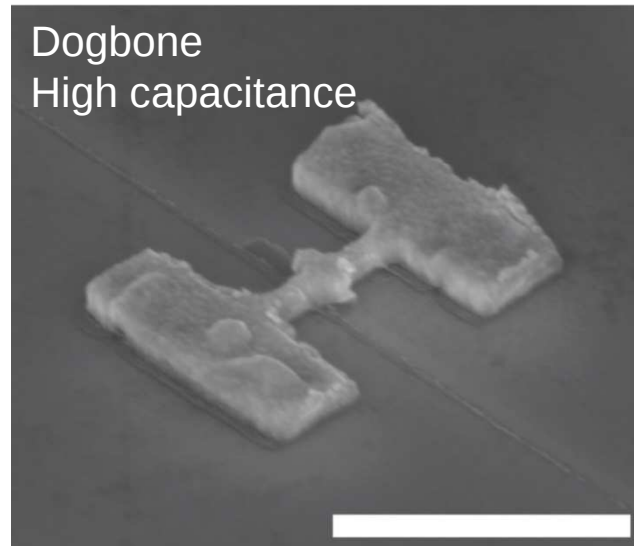
Accuracy of circuit model



- FDTD simulations show clear anti-crossing as a function bare cavity frequency
- Circuit model can recover full lineshape of reflectance (or transmittance) spectra



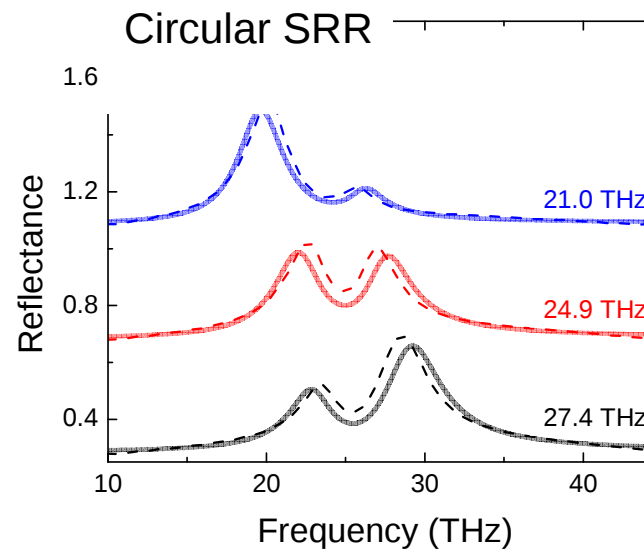
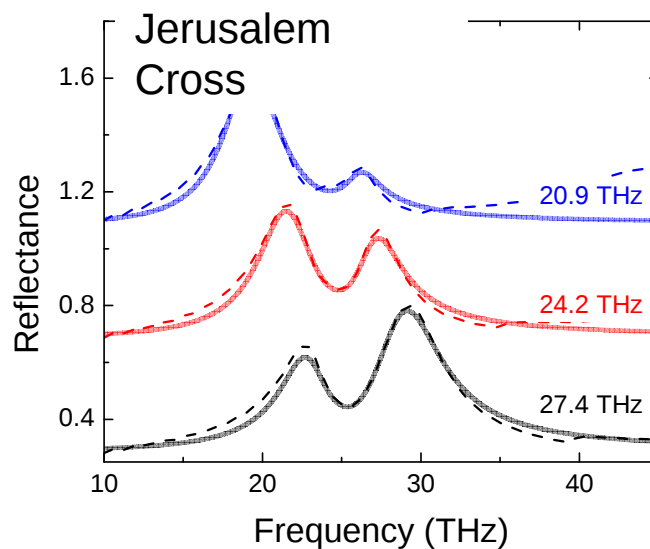
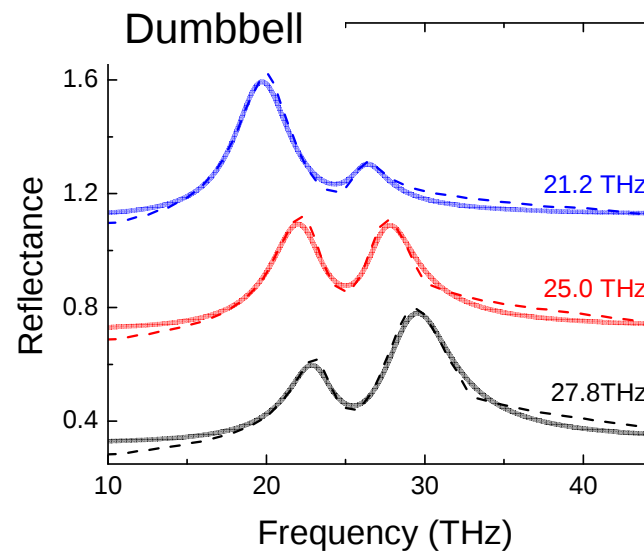
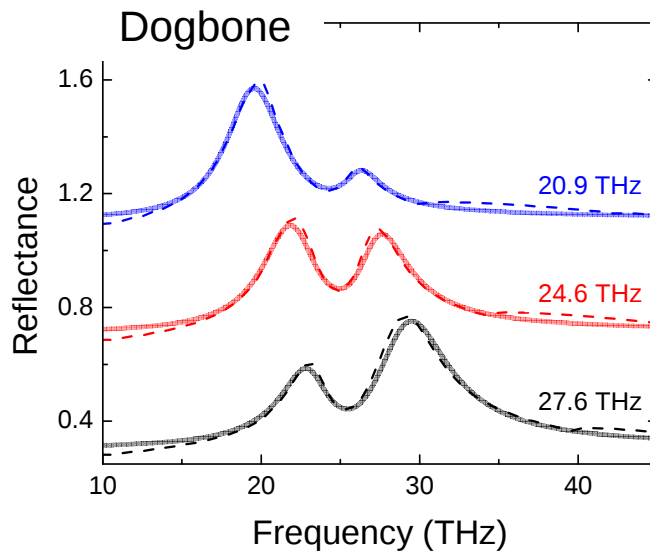
Resonator Geometries Analyzed



- All metamaterial have identical bare cavity resonances but different RLC values
- Metamaterial defined by electron-beam lithography
- Same QW-stack used for all samples
- Study only geometry effect



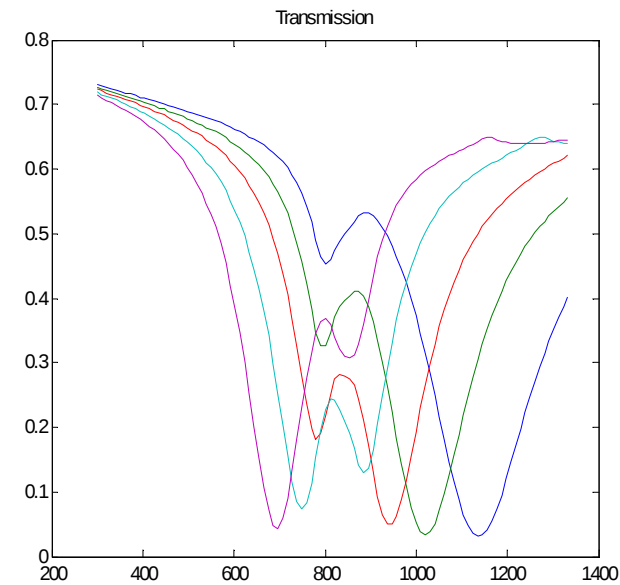
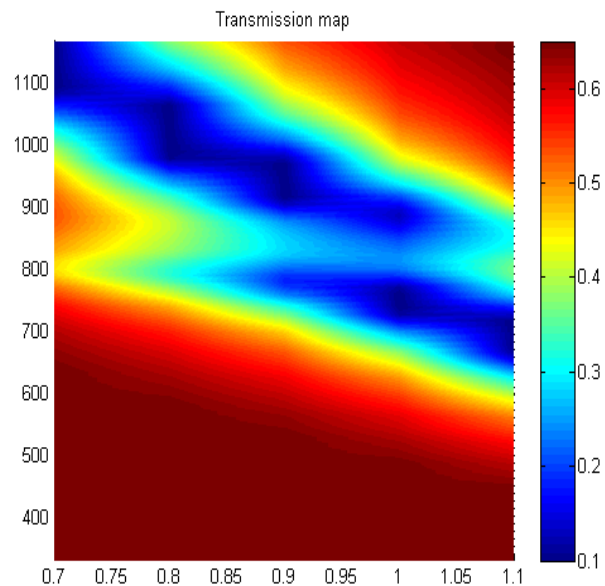
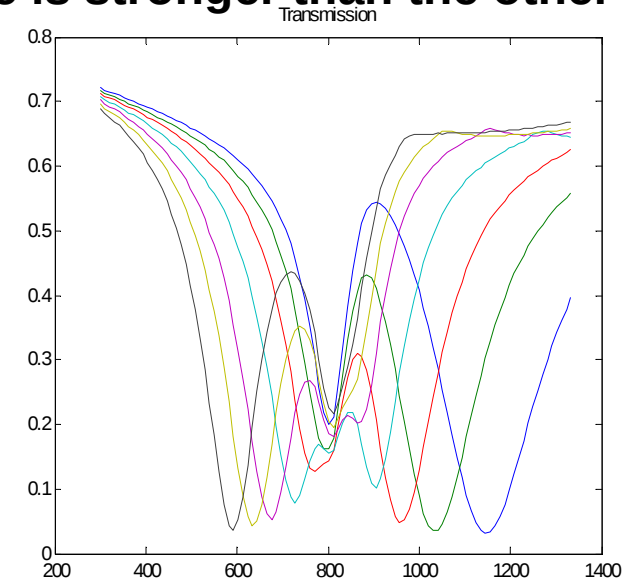
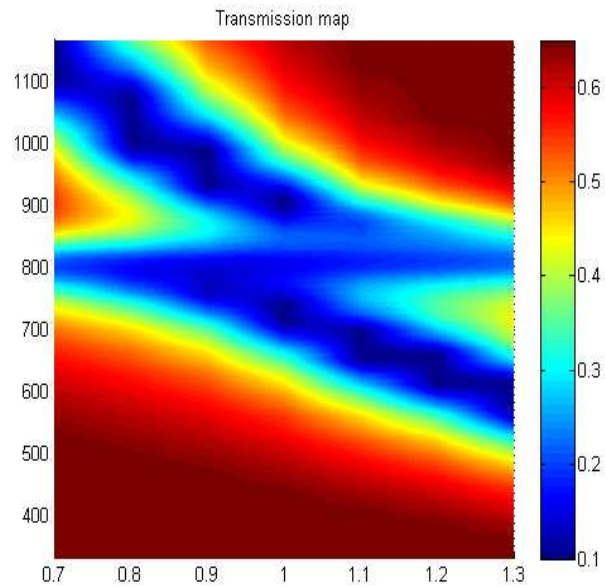
Circuit model accuracy



- RLC model describes strong coupling accurately despite the exact geometry
- Only cold cavity requires fitting



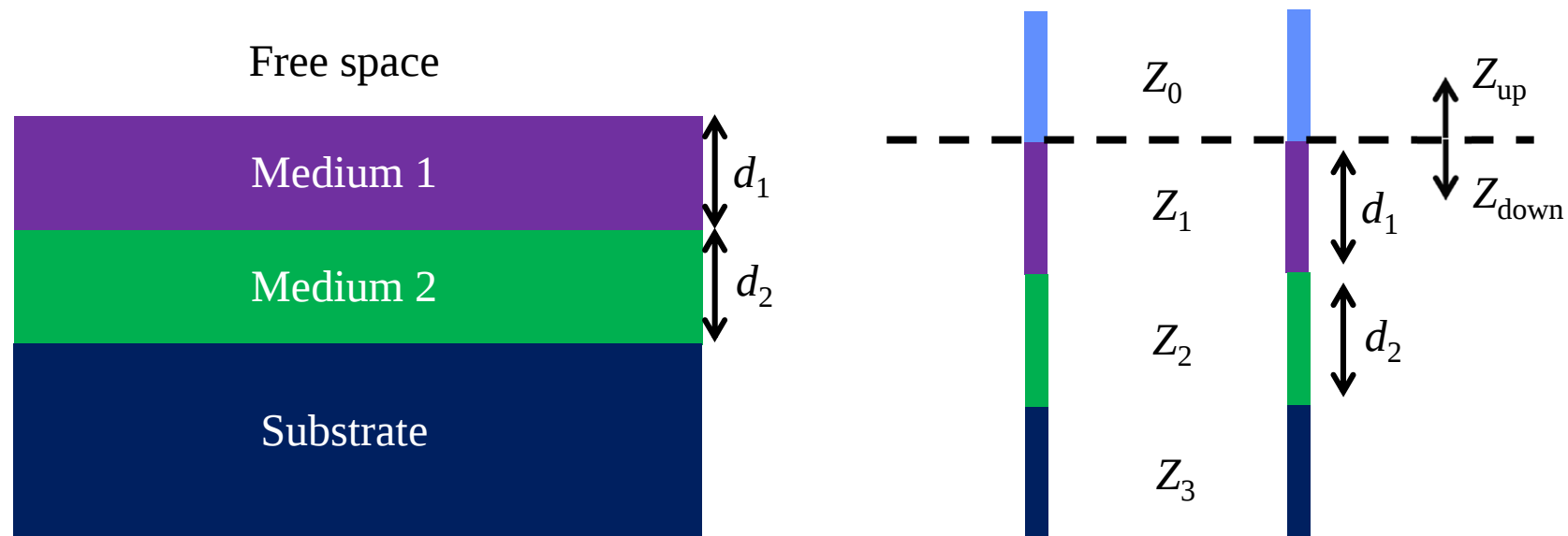
Top: Two isotropic Lorentzians; one is stronger than the other



Bottom: Two anisotropic Lorentzians; one is stronger than the other

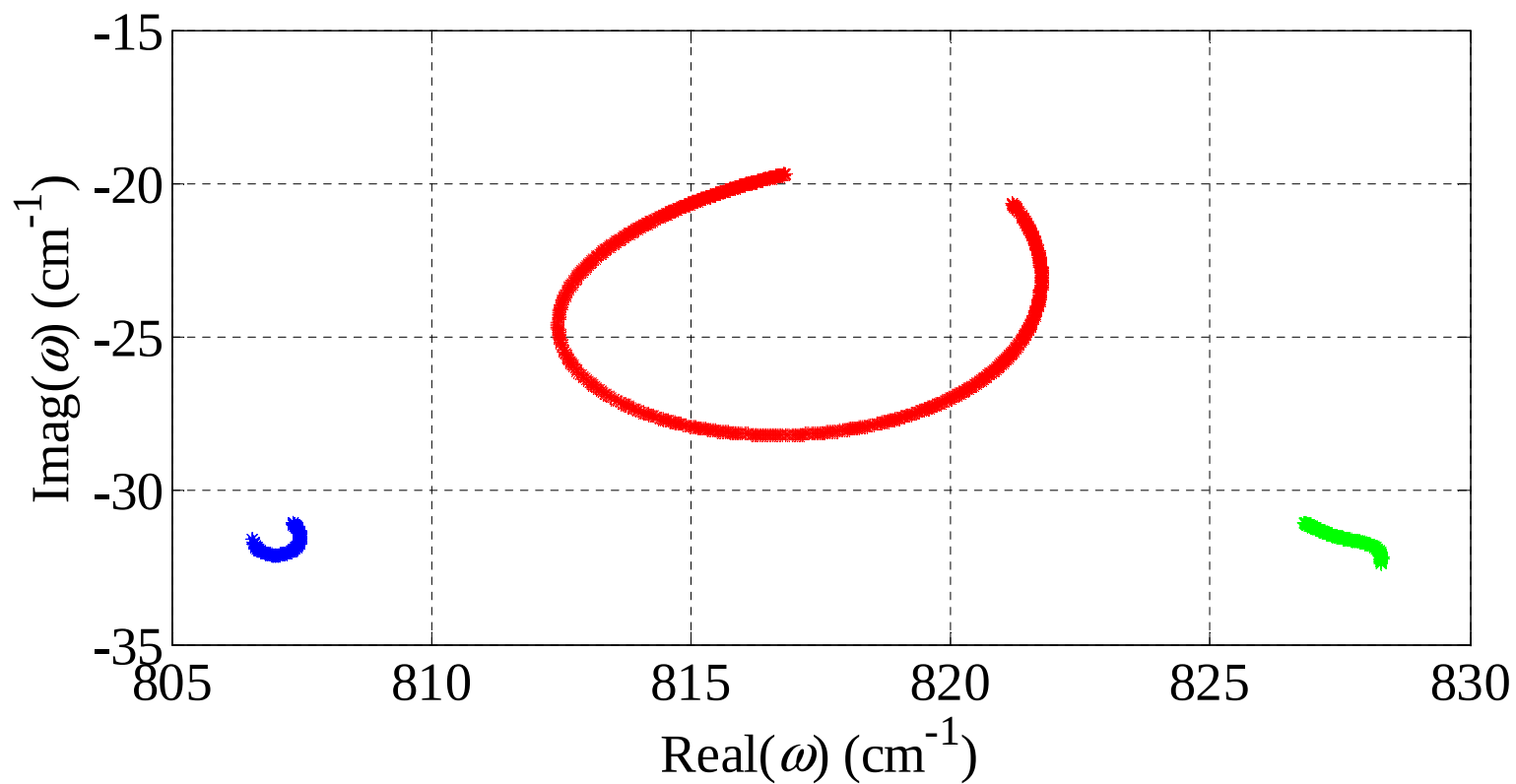
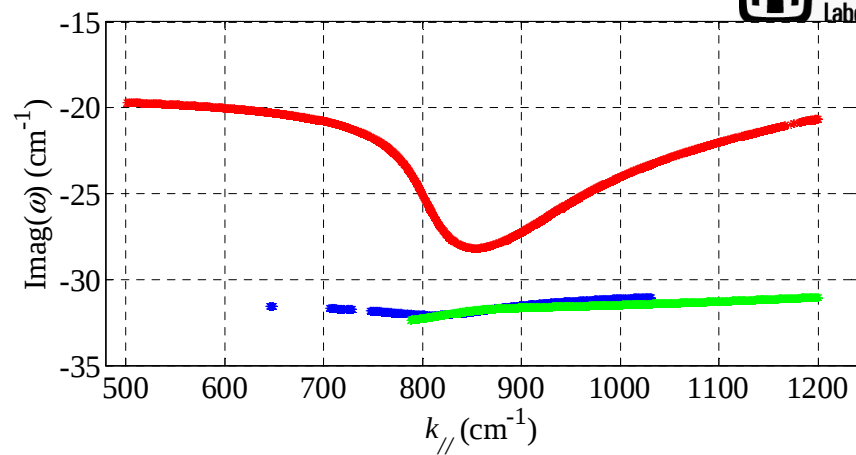
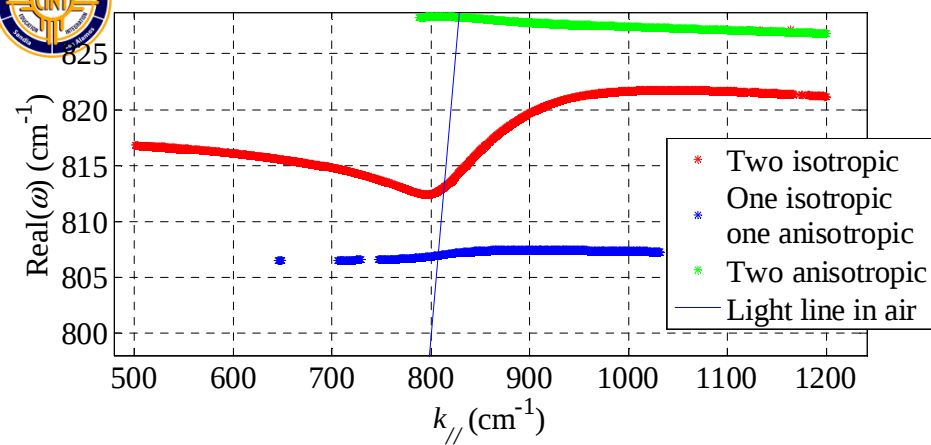
Modal analysis

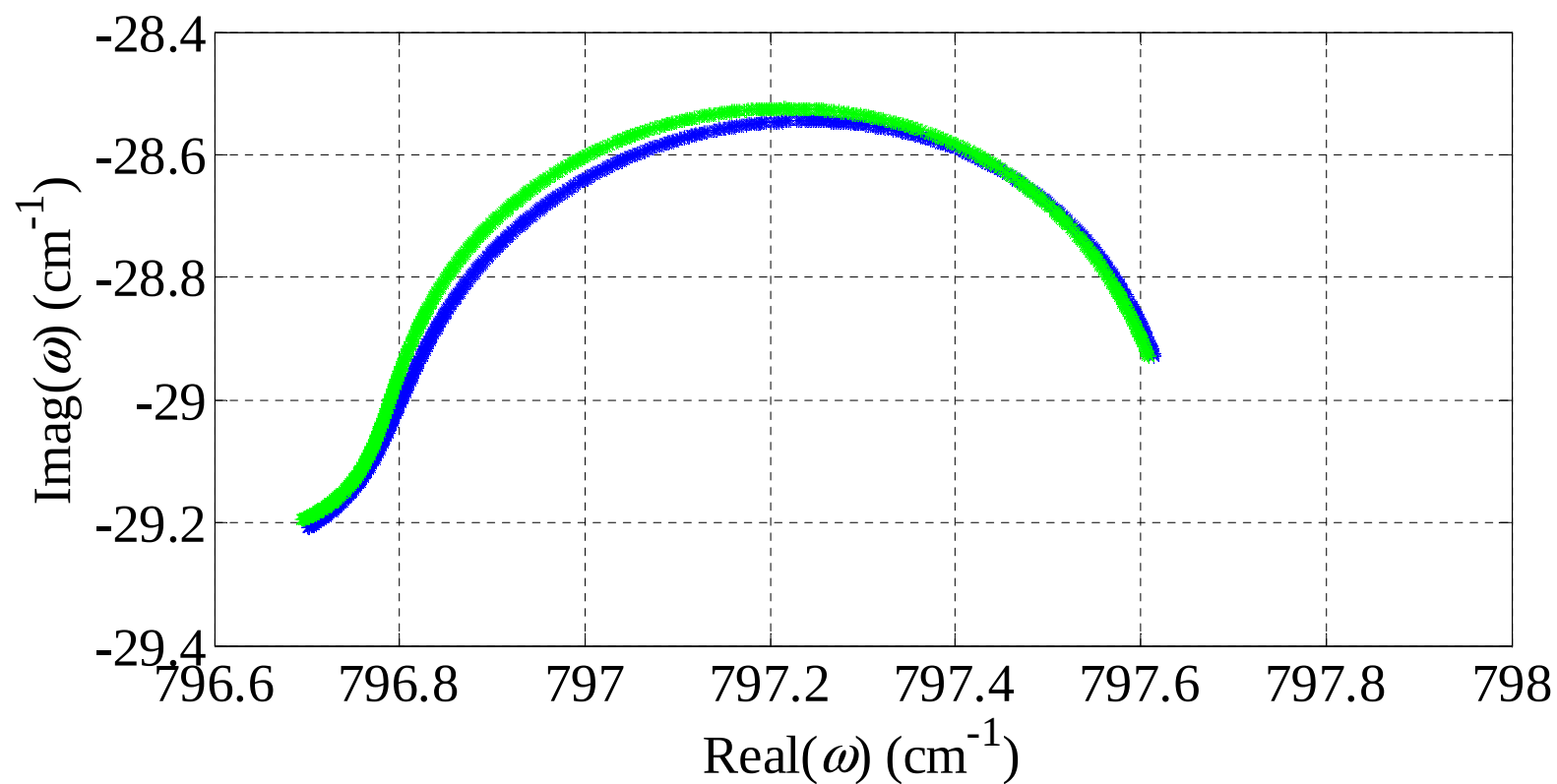
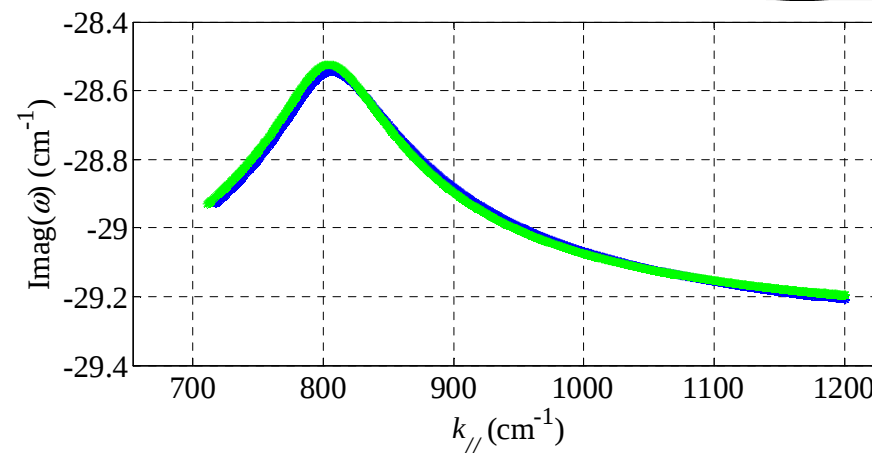
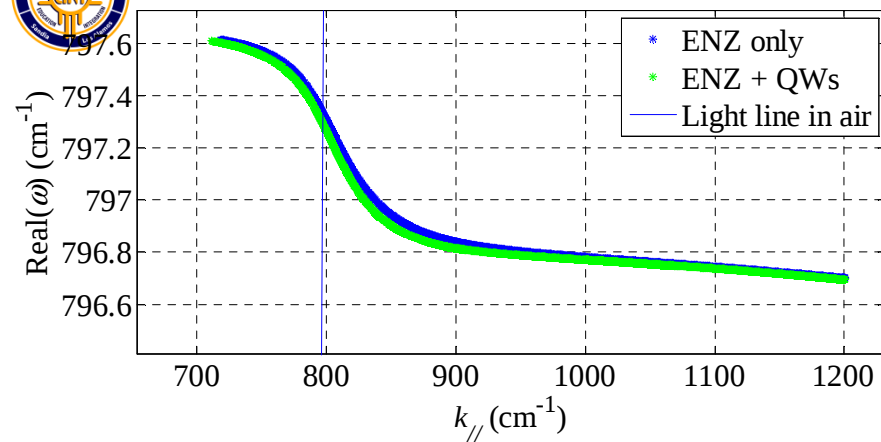
The layered medium can be modeled as a set of cascaded transmission lines



Select a reference plane, then compute the impedances looking up and looking down accounting for the dispersion in each medium

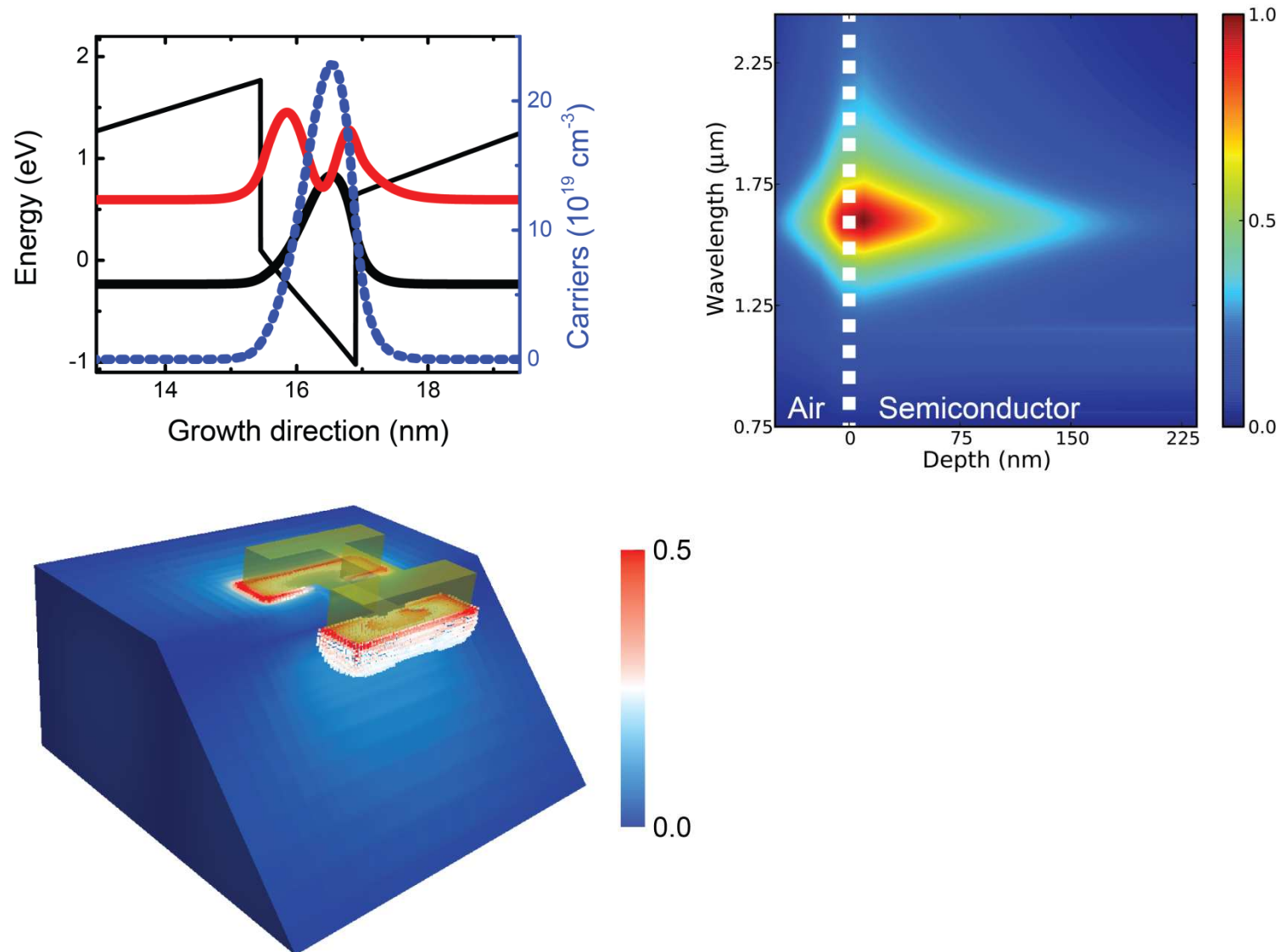
Compute the (complex) modes as the zeroes of $Z_{up} + Z_{down} = 0$





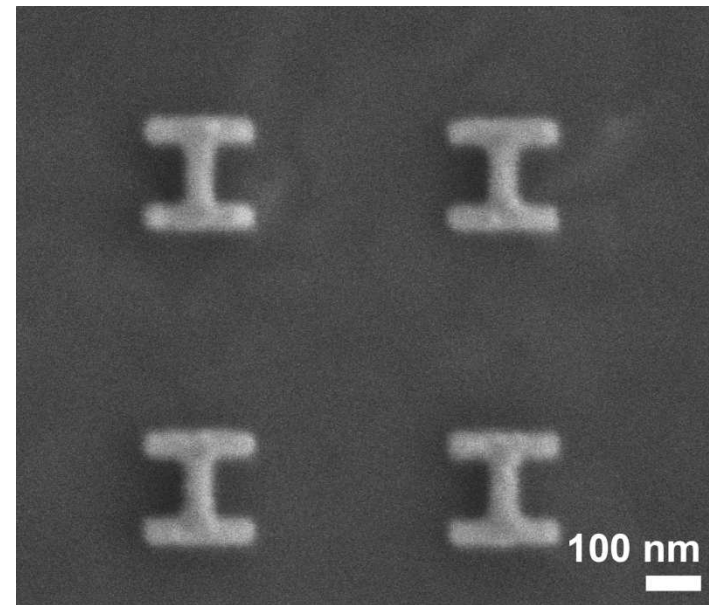
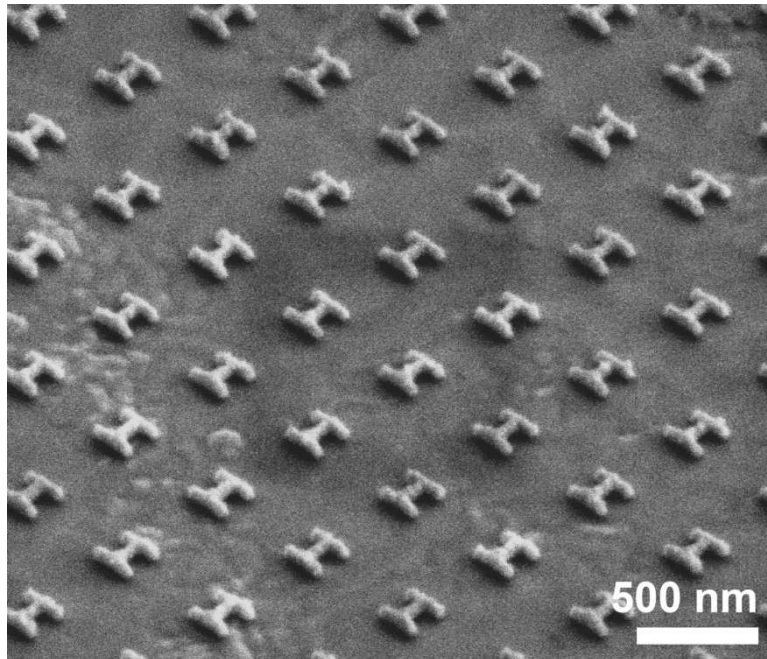


Strong Coupling, Near IR



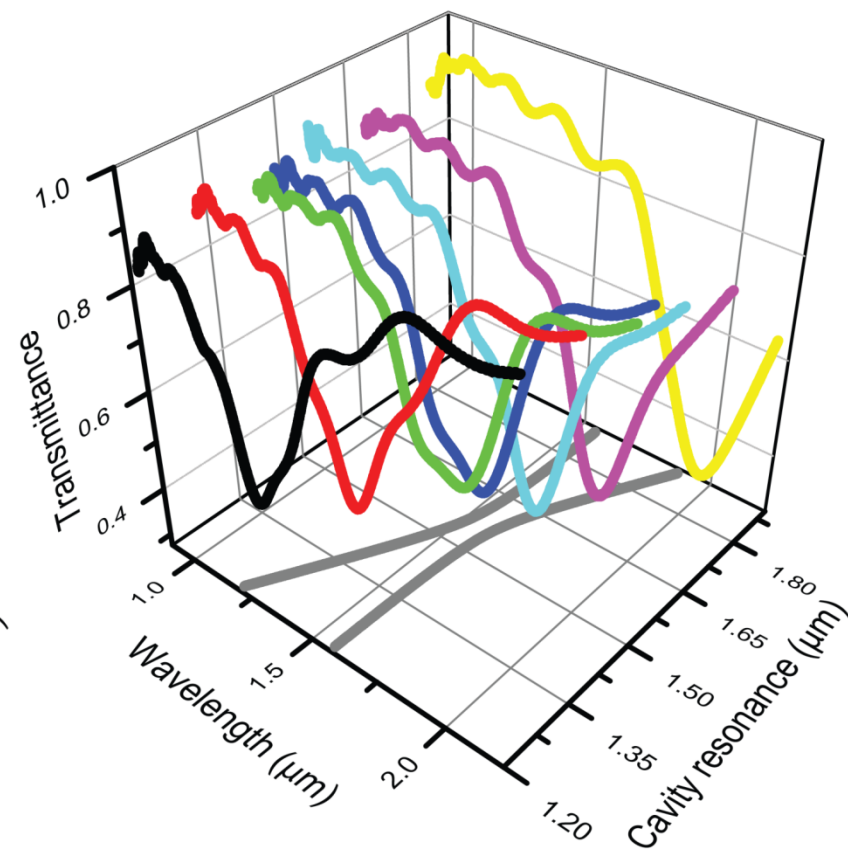
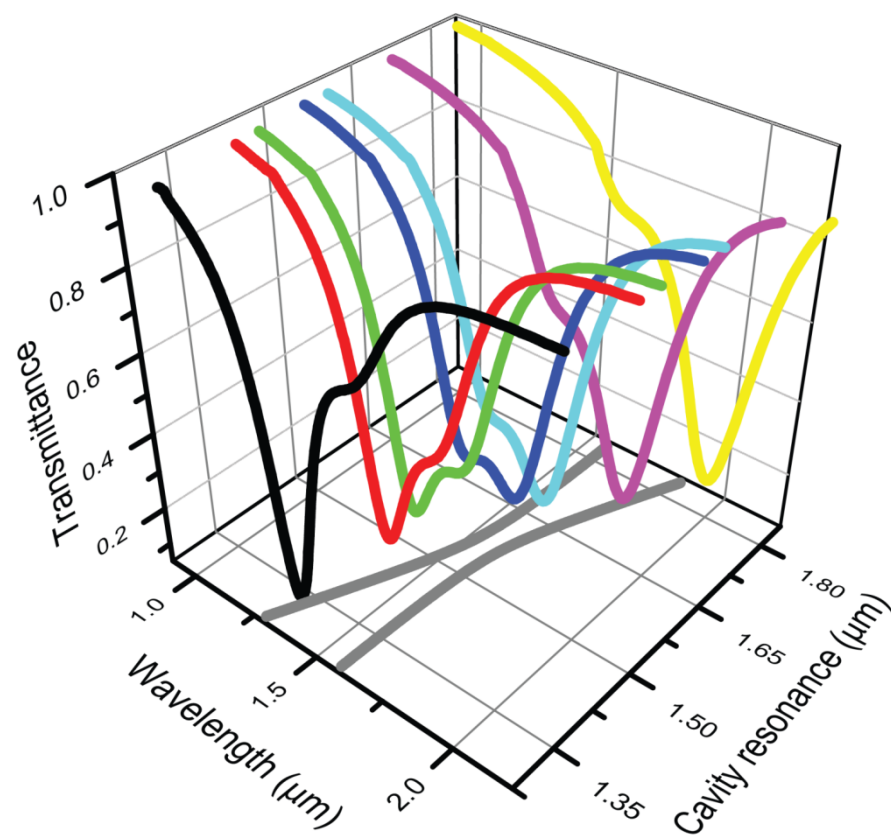


Dogbone Resonators, Near IR



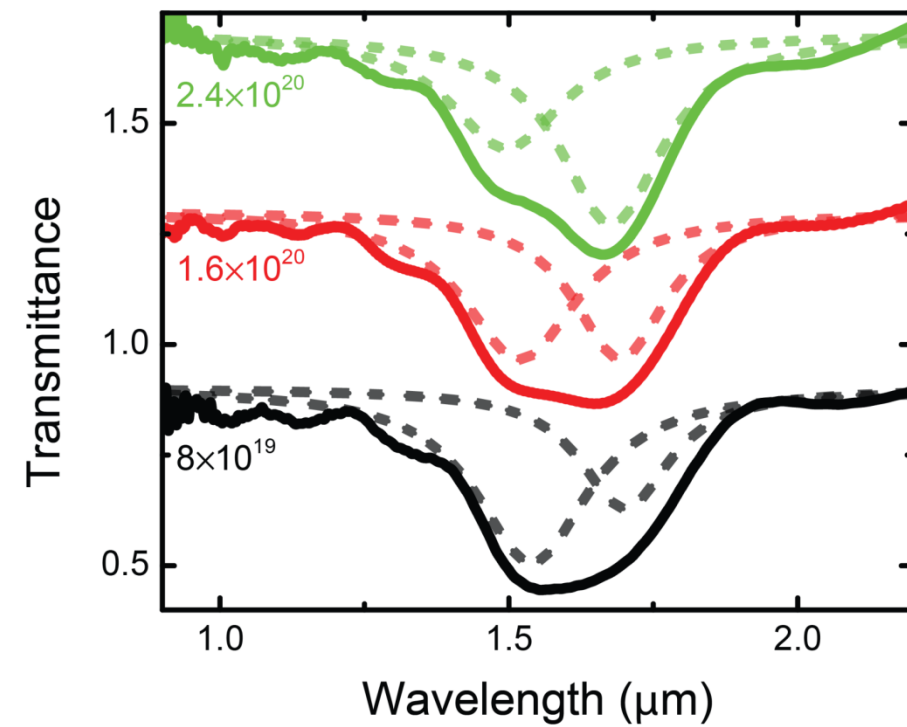
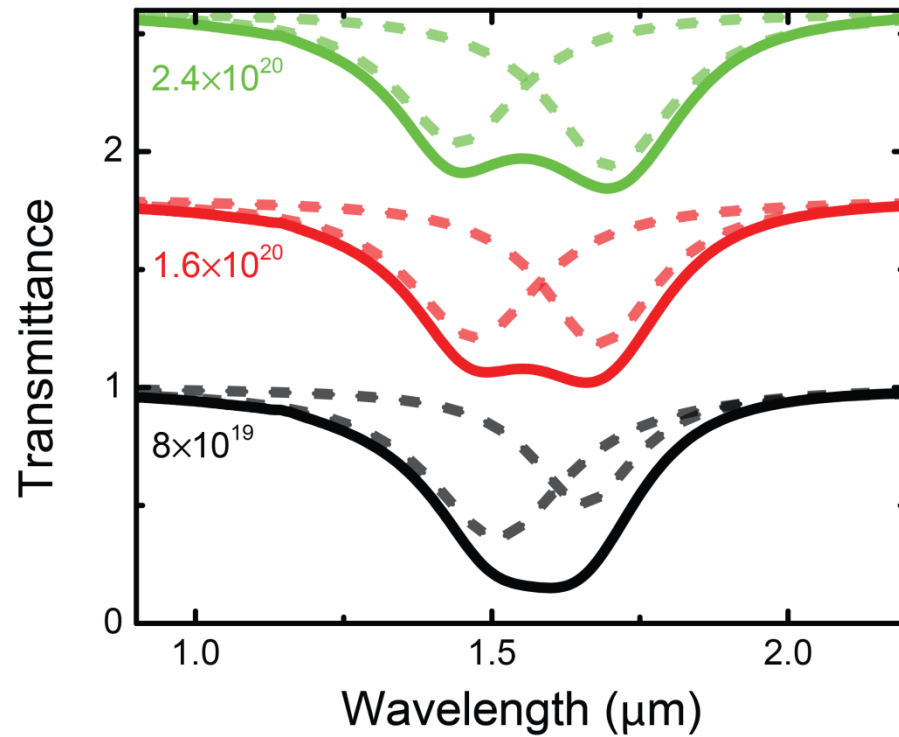


Strong Coupling, Near IR



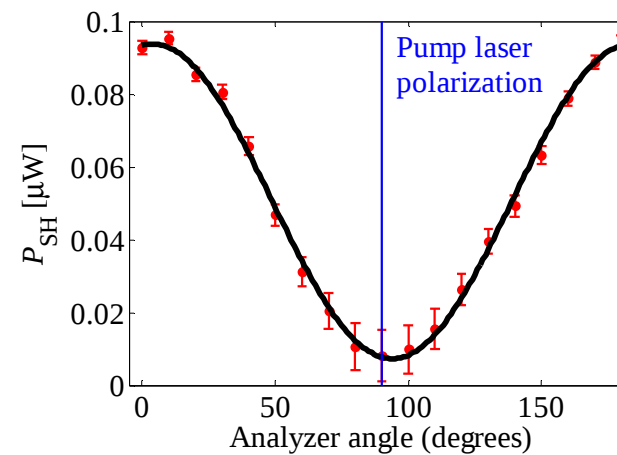
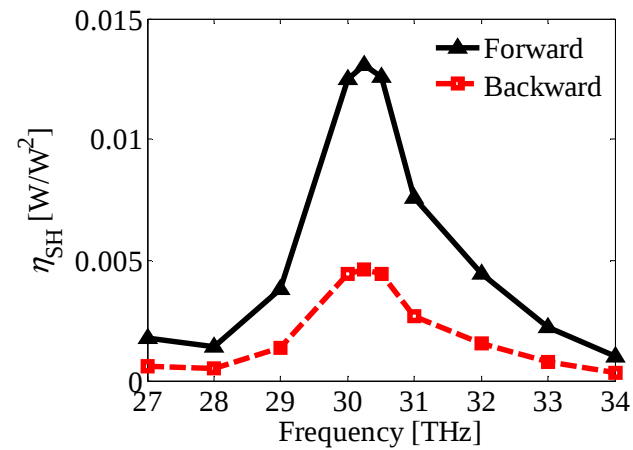
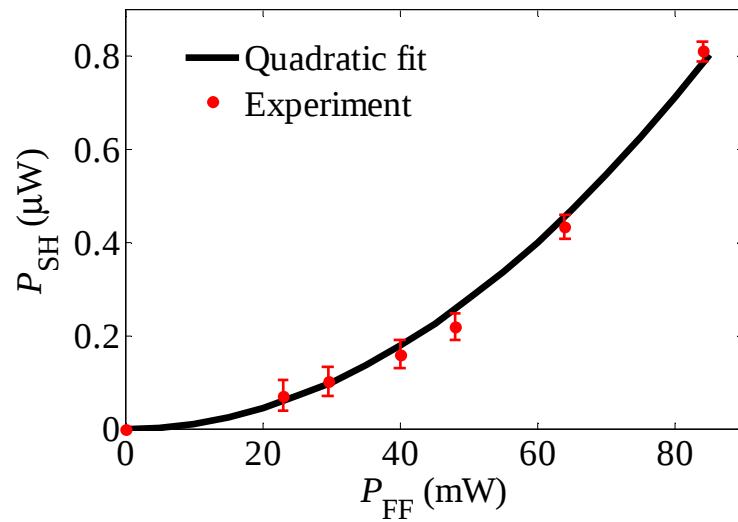


Strong Coupling, Near IR





SHG



Appl. Phys. Lett. **104**, 131104 (2014)