

Searching for lasing threshold in a thresholdless laser

SAND2014-16570PE

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Christopher Gies and Frank Jahnke, Bremen University**

Nanolasers and thresholdless lasing

Criterial for lasing

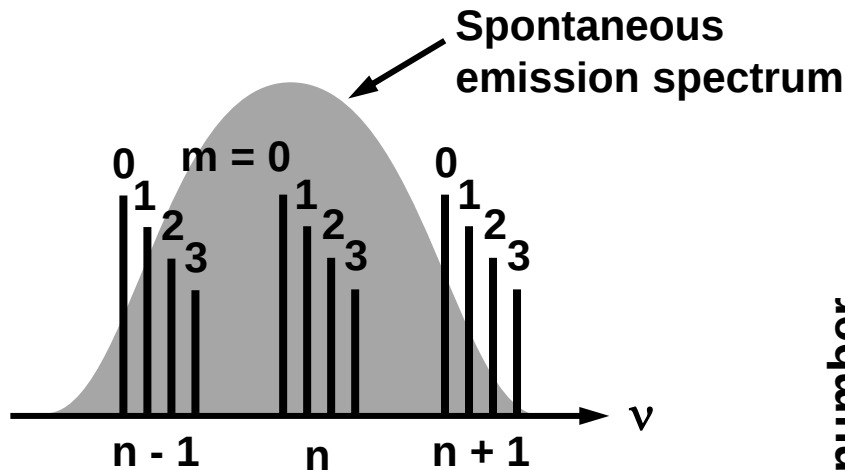
$\beta \ll 1$ versus $\beta = 1$

Thanks to:

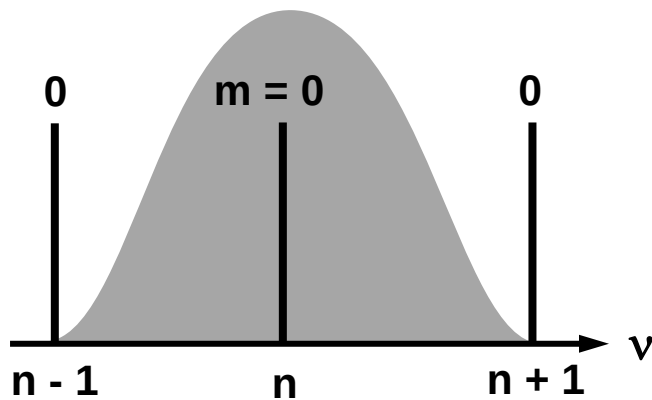
Sandia's Laboratory Directed Research & Development (LDRD) Program

Motivation: Emission entirely into single mode

Most lasers $\beta \ll 1$

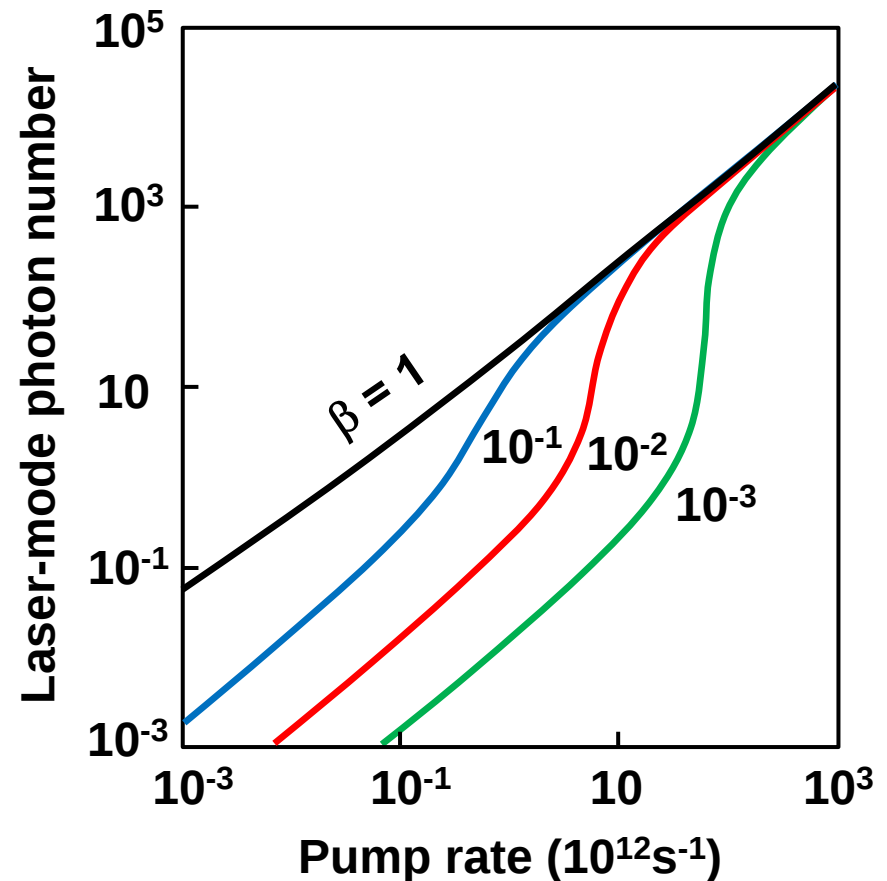


Nanolasers $\beta = 1$

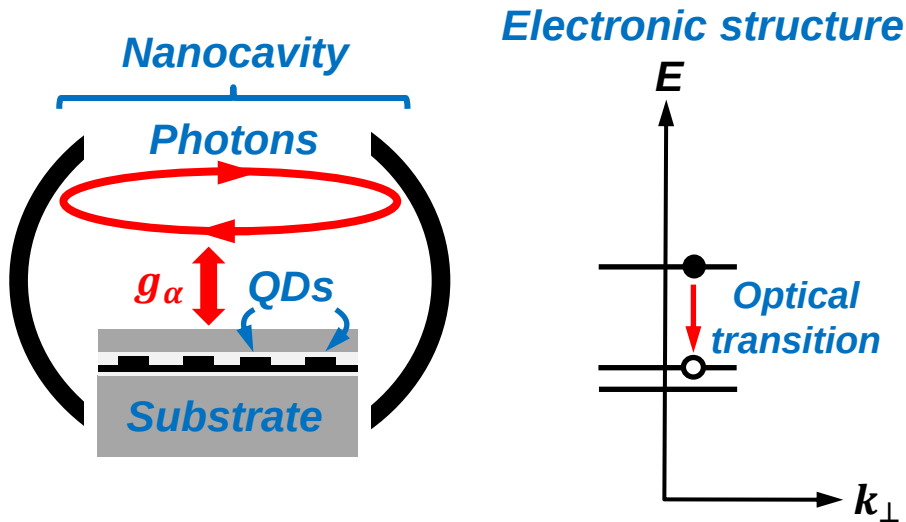


Spontaneous emission factor

$$\beta = \frac{\gamma_l}{\gamma_{sp}}$$



Nanolaser model



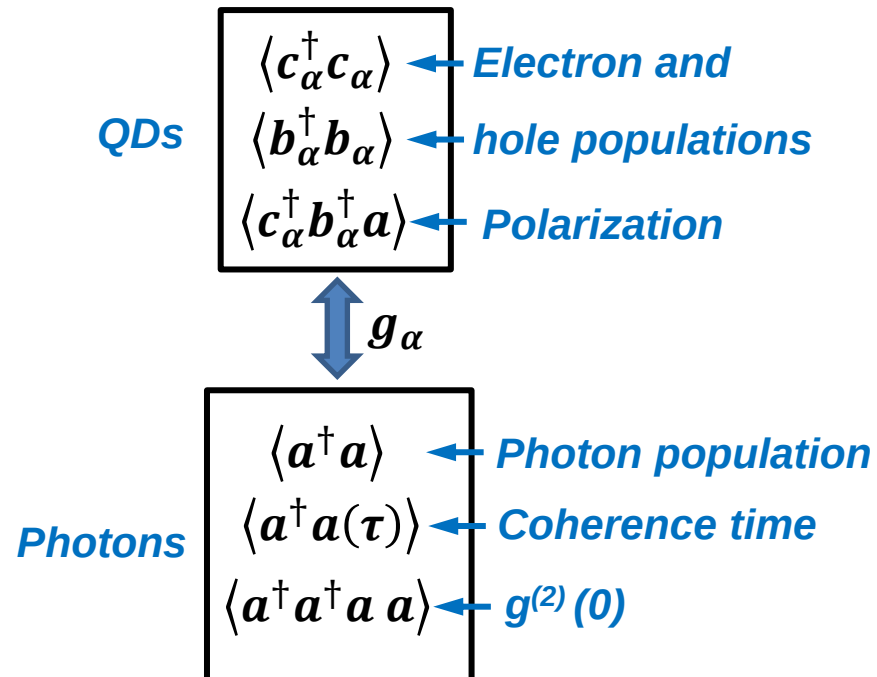
Hamiltonian

$$H = H_0 \quad \text{Single-particle energies}$$

$$\begin{aligned}
 & -\hbar \sum_{\alpha} \left(g_{\alpha} b_{\alpha}^{\dagger} c_{\alpha}^{\dagger} a - g_{\alpha}^* a^{\dagger} c_{\alpha} b_{\alpha} \right) \quad \text{Light-carrier} \\
 & + \hbar \sum_{\alpha\beta q} G_q (c_{\alpha}^{\dagger} c_{\beta} + b_{\alpha}^{\dagger} b_{\beta}) (d_q + d_q^{\dagger}) \quad \text{Carrier-phonon} \\
 & + \frac{1}{2} \sum_{\alpha\beta\sigma\eta} W_{\sigma\eta}^{\alpha\beta} c_{\alpha}^{\dagger} c_{\beta}^{\dagger} c_{\eta} c_{\sigma} + \frac{1}{2} \sum_{\alpha\beta\sigma\eta} W_{\sigma\eta}^{\alpha\beta} b_{\alpha}^{\dagger} b_{\beta}^{\dagger} b_{\eta} b_{\sigma} \\
 & - \sum_{\alpha\beta\sigma\eta} W_{\sigma\eta}^{\alpha\beta} b_{\alpha}^{\dagger} c_{\beta}^{\dagger} c_{\eta} b_{\sigma}
 \end{aligned}$$

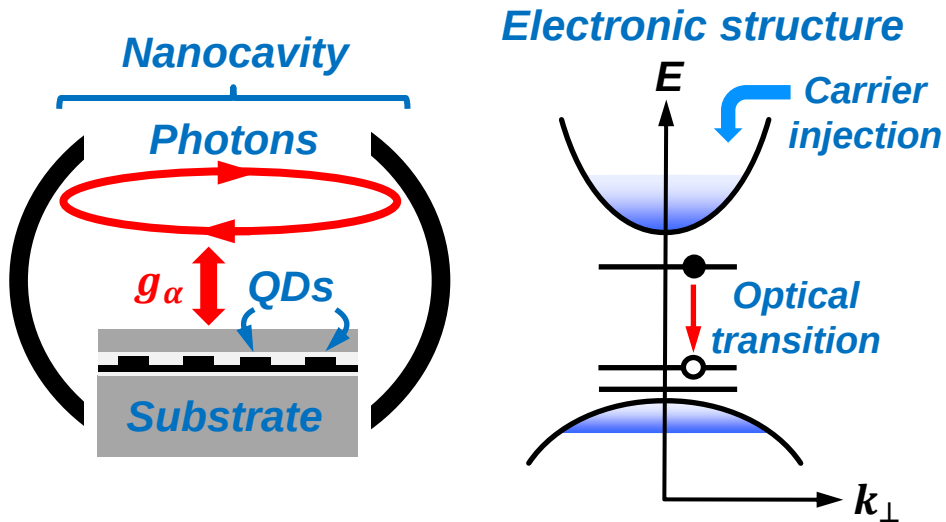
$\propto \sqrt{v/\hbar\epsilon_b V} W(R_{QD}) \sum_n C_{\alpha}(R_n) V_{\alpha}(R_n)$
 $\downarrow$ Matrix element of $\frac{e^2}{4\pi\epsilon_b |r - r'|}$

Carrier-carrier



Emphasis now is on correlations involving light-matter interaction instead of Coulomb interaction

Nanolaser model



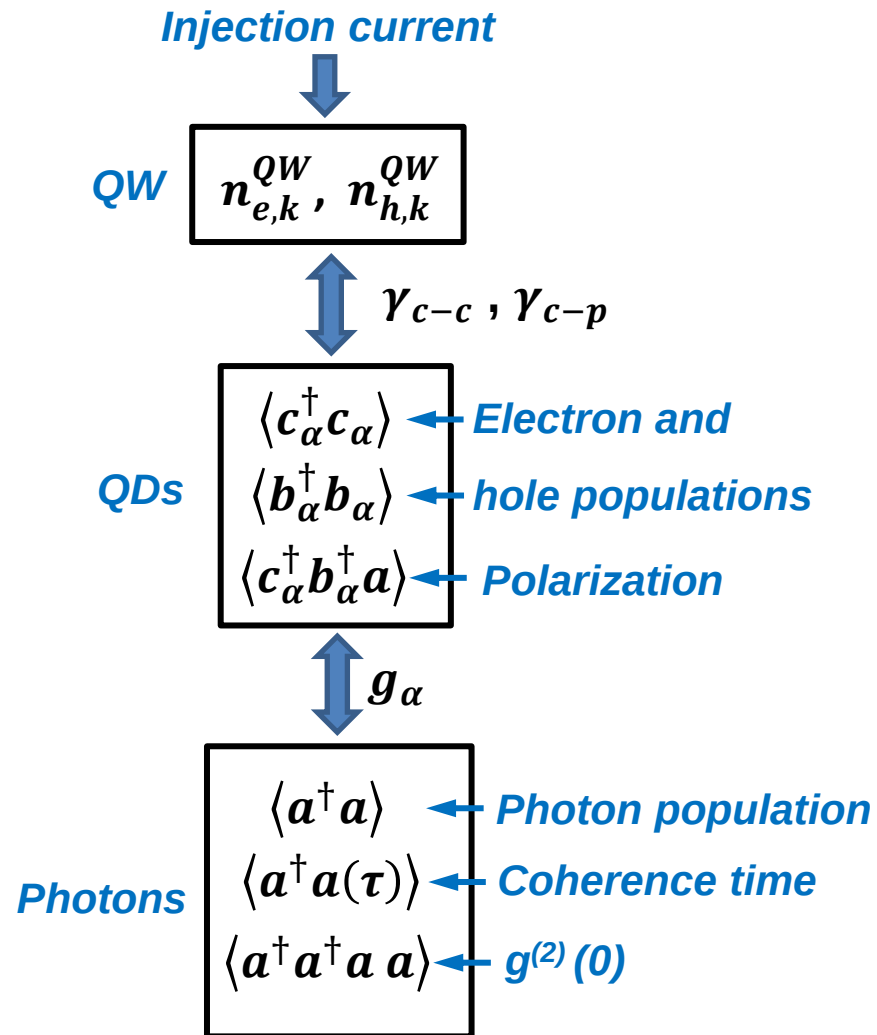
Hamiltonian

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 & + \frac{1}{2} \sum_{\alpha\beta\sigma\eta} W_{\sigma\eta}^{\alpha\beta} c_{\alpha}^{\dagger} c_{\beta}^{\dagger} c_{\eta} c_{\sigma} + \frac{1}{2} \sum_{\alpha\beta\sigma\eta} W_{\sigma\eta}^{\alpha\beta} b_{\alpha}^{\dagger} b_{\beta}^{\dagger} b_{\eta} b_{\sigma} \\
 & - \sum_{\alpha\beta\sigma\eta} W_{\sigma\eta}^{\alpha\beta} b_{\alpha}^{\dagger} c_{\beta}^{\dagger} c_{\eta} b_{\sigma}
 \end{aligned}$$

$\leftarrow \frac{e^2}{4\pi\epsilon_b|r-r'|}$ } Carrier-carrier

$\leftarrow \frac{1}{\rho\sqrt{v/\hbar\epsilon_b V}} W(R_{QD}) \sum_n C_{\alpha}(R_n) V_{\alpha}(R_n)$



Emphasis now is on correlations involving light-matter interaction instead of Coulomb interaction

Correlations

2nd order photon correlation

$$\langle a^\dagger a^\dagger a a \rangle = 2\langle a^\dagger a \rangle + \underbrace{\delta\langle a^\dagger a^\dagger a a \rangle}$$

photon-polarization

$$\delta\langle b_\alpha^\dagger c_\alpha^\dagger a^\dagger a a \rangle$$

$$\delta\langle c_\alpha b_\alpha a^\dagger a^\dagger a \rangle$$

photon-carrier

$$\delta\langle c_\alpha^\dagger c_\alpha a^\dagger a \rangle$$

$$\delta\langle b_\alpha^\dagger b_\alpha a^\dagger a \rangle$$

polarization-polarization

$$\delta\langle b_\alpha^\dagger b_\beta^\dagger c_\beta^\dagger c_\alpha^\dagger a a \rangle$$

$$\delta\langle b_\alpha^\dagger c_\alpha^\dagger c_\beta b_\beta a^\dagger a \rangle$$

$$\delta\langle c_\alpha c_\beta b_\beta b_\alpha a^\dagger a^\dagger \rangle$$

carrier-polarization

$$\delta\langle b_\alpha^\dagger c_\alpha^\dagger c_\beta^\dagger c_\beta a \rangle$$

$$\delta\langle c_\beta^\dagger c_\beta c_\alpha b_\alpha a^\dagger \rangle$$

$$\delta\langle b_\beta^\dagger b_\alpha^\dagger c_\alpha^\dagger b_\beta a \rangle$$

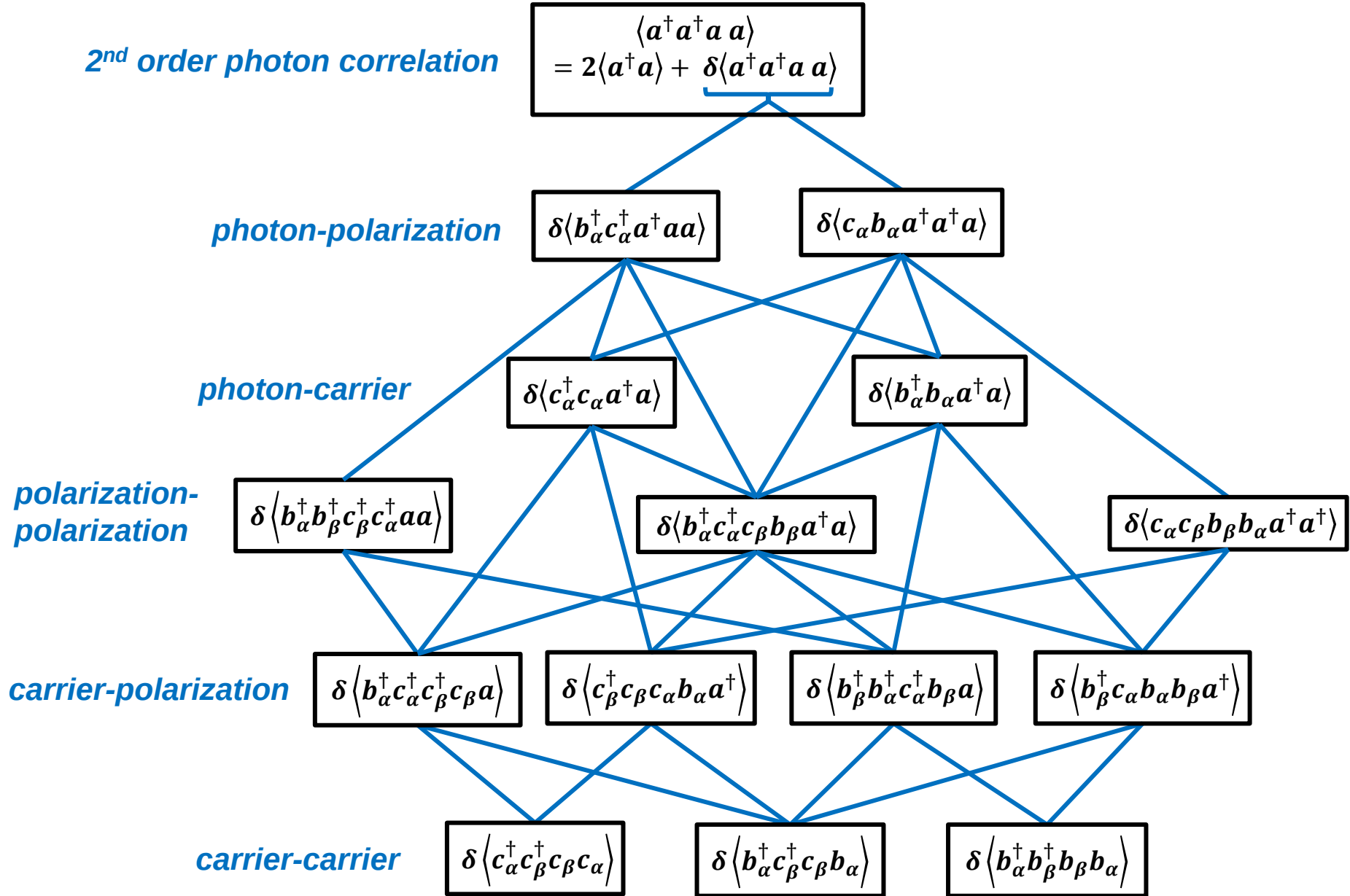
$$\delta\langle b_\beta^\dagger c_\alpha b_\alpha b_\beta a^\dagger \rangle$$

carrier-carrier

$$\delta\langle c_\alpha^\dagger c_\beta^\dagger c_\beta c_\alpha \rangle$$

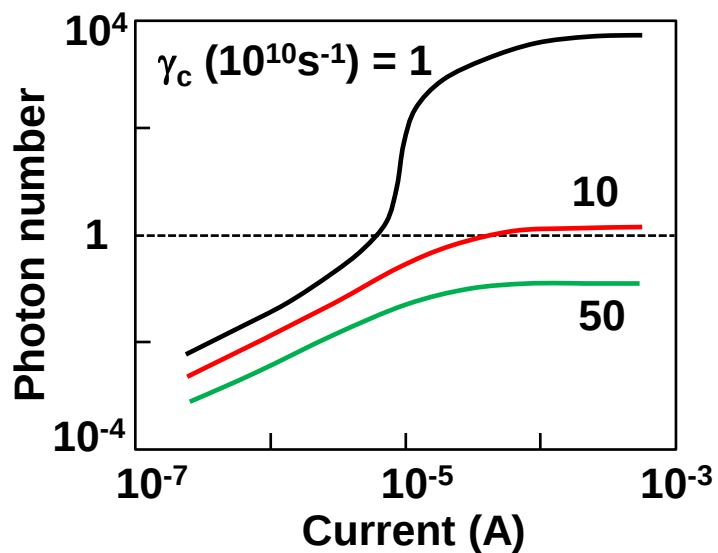
$$\delta\langle b_\alpha^\dagger c_\beta^\dagger c_\beta b_\alpha \rangle$$

$$\delta\langle b_\alpha^\dagger b_\beta^\dagger b_\beta b_\alpha \rangle$$

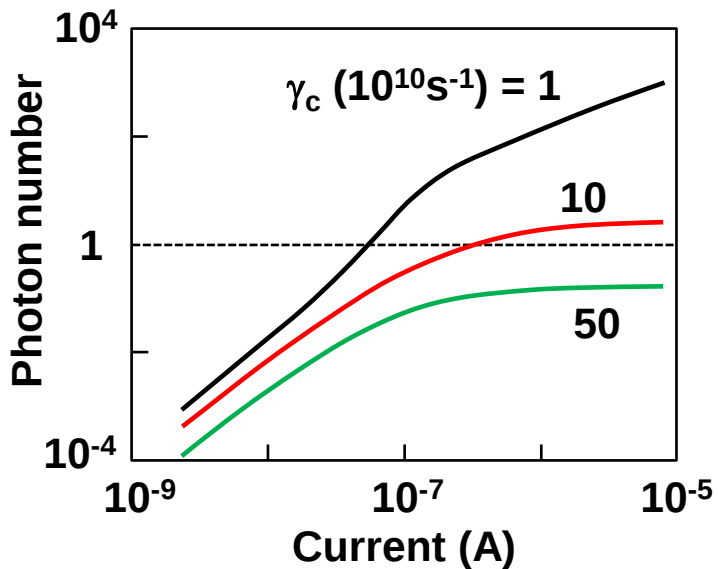


Criterial for lasing: Input/Output

$\beta = 0.01$



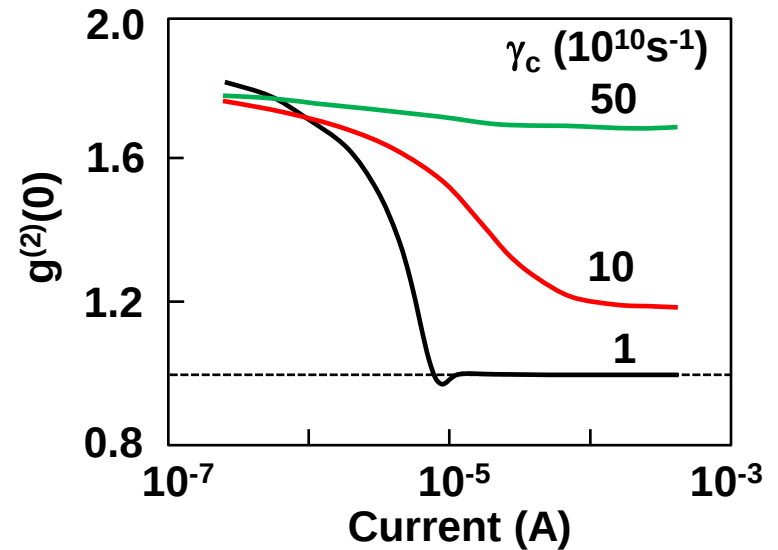
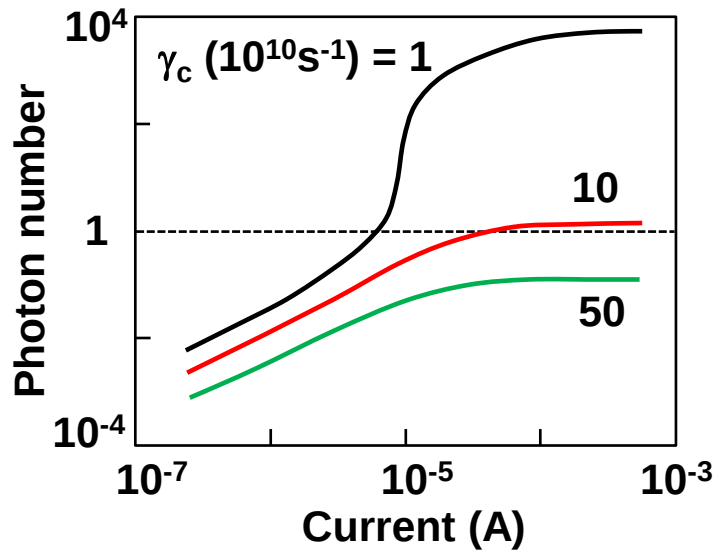
$\beta = 1$



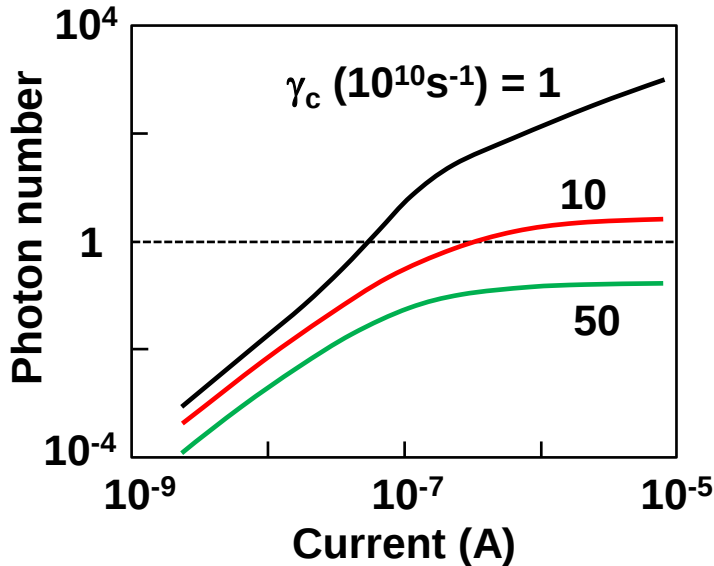
$$N_{\text{QD}} = 50, \Delta_{\text{inh}} = 20 \text{meV}$$

Criterial for lasing: $g^{(2)}(0)$

$\beta = 0.01$



$\beta = 1$

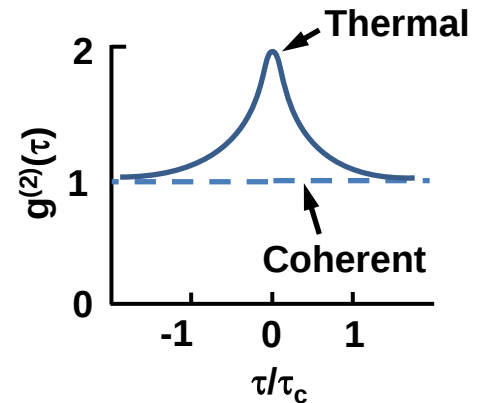
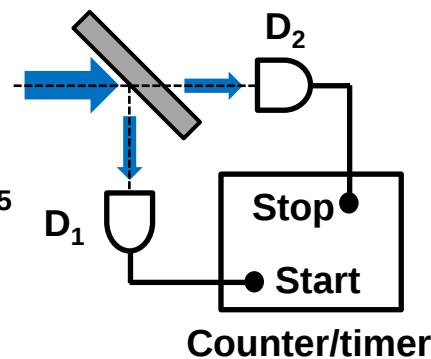


$N_{\text{QD}} = 50, \Delta_{\text{inh}} = 20\text{meV}$

Second-order photon correlation function

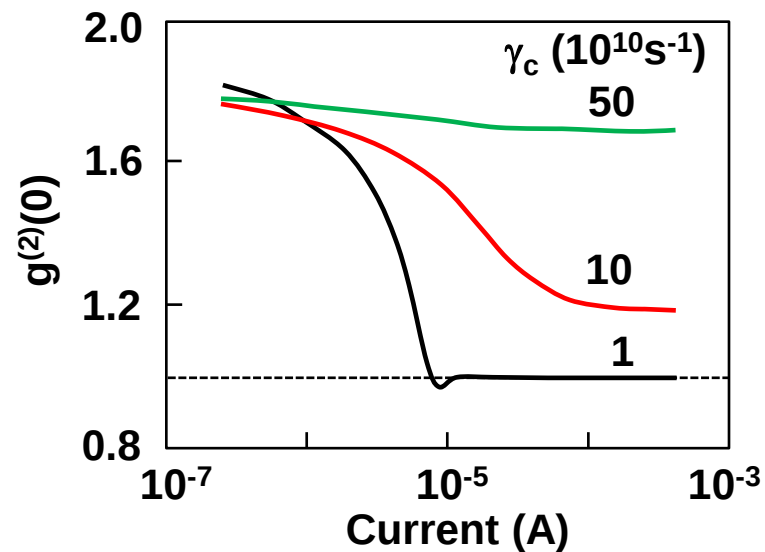
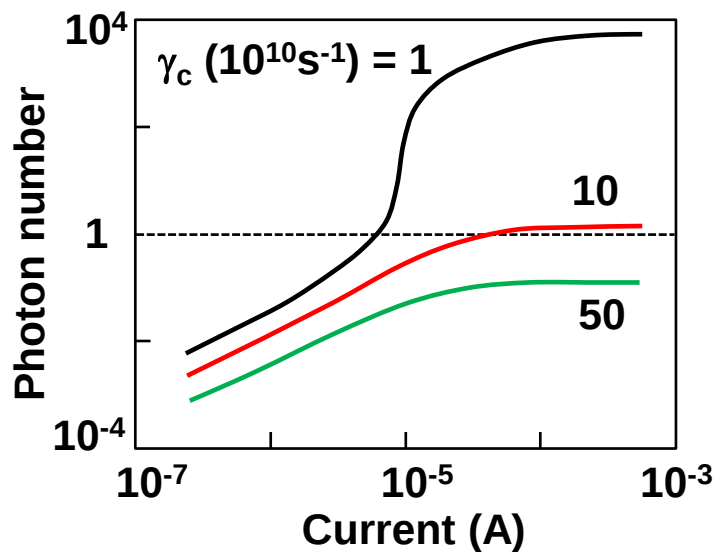
$$g^{(2)}(\tau) = \frac{\langle I(t)I(t+\tau) \rangle}{\langle I(t) \rangle^2}$$

Hanbury-Brown-Twiss experiment

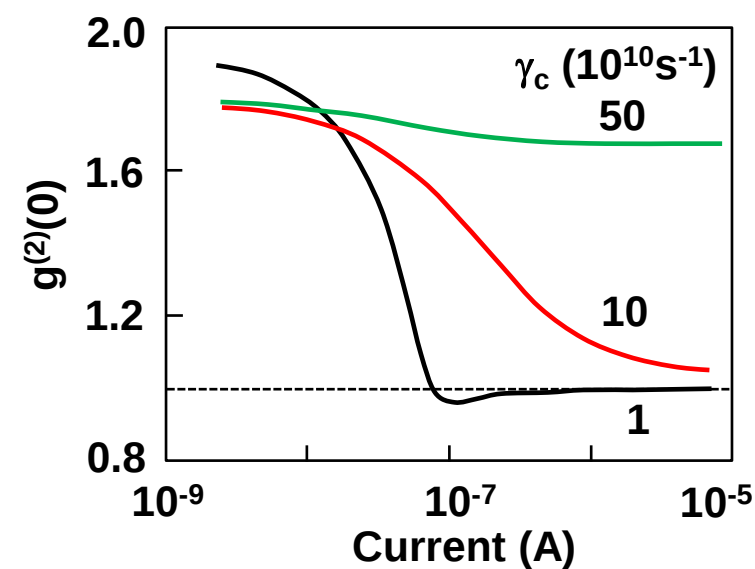
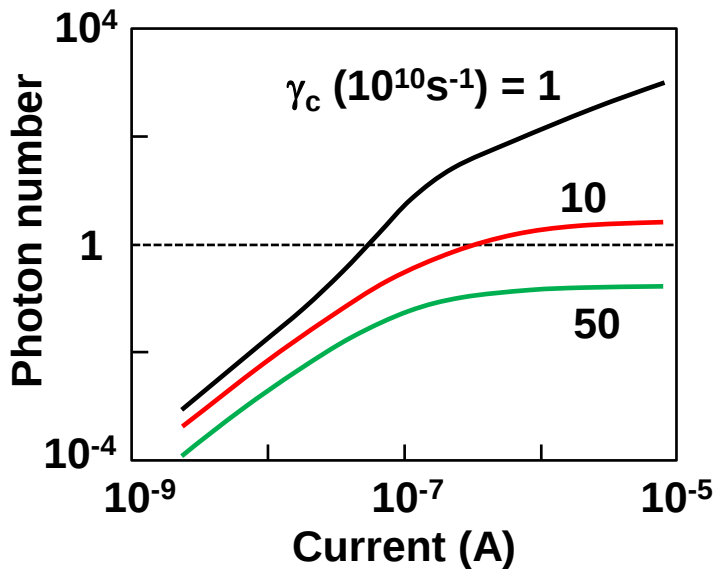


Criterial for lasing: $g^{(2)}(0)$

$\beta = 0.01$



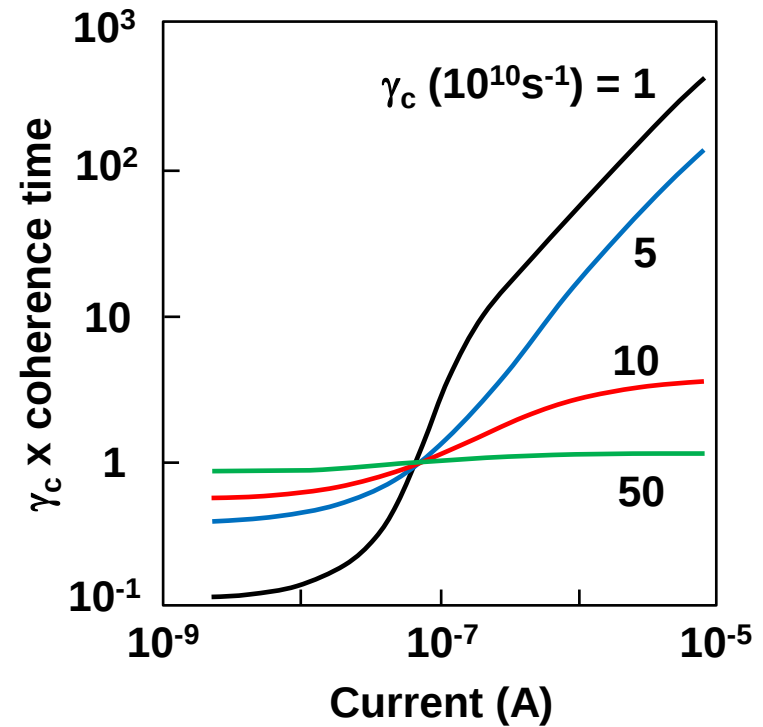
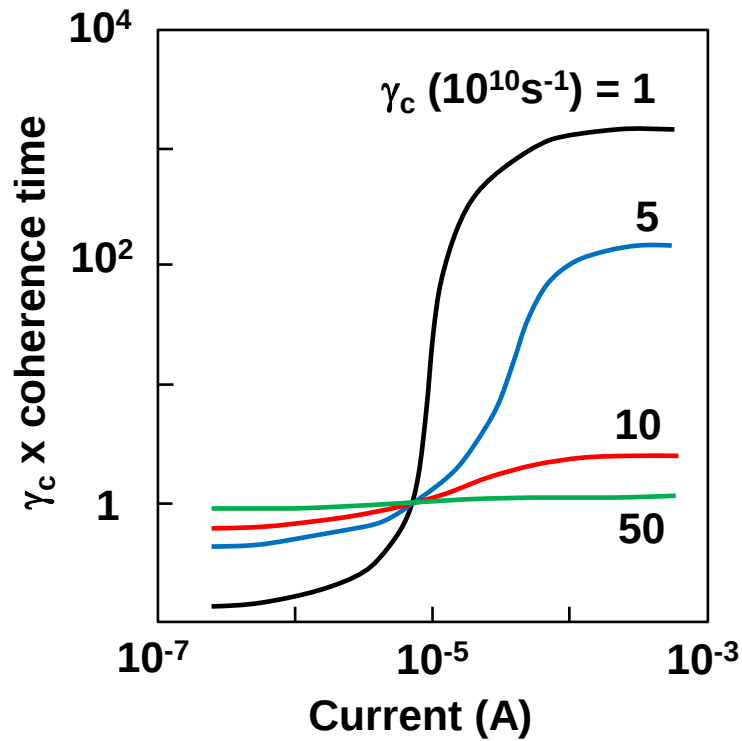
$\beta = 1$



$N_{\text{QD}} = 50, \Delta_{\text{inh}} = 20 \text{ meV}$

Criterial for lasing: coherence time

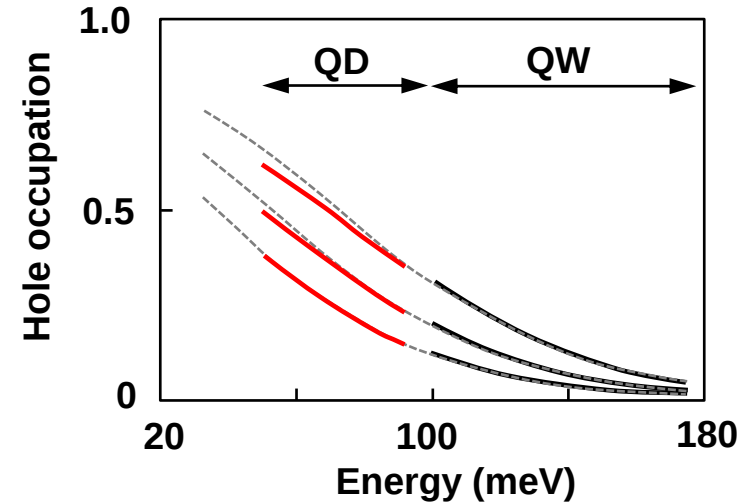
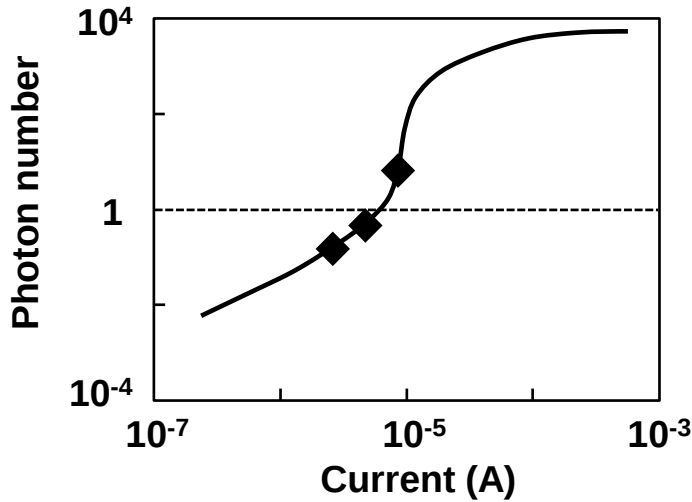
$$\tau_c = 2 \int_{-\infty}^{\infty} d\tau \left| \frac{\langle a^\dagger a(\tau) \rangle_{ss}}{\langle a^\dagger a \rangle_{ss}} \right|^2$$



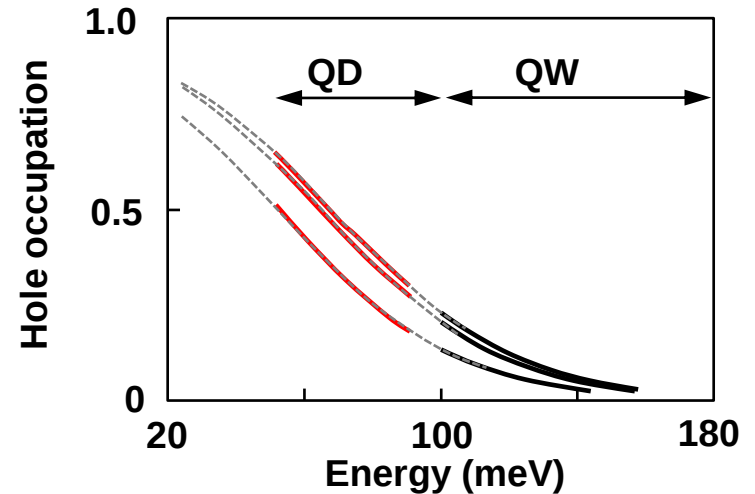
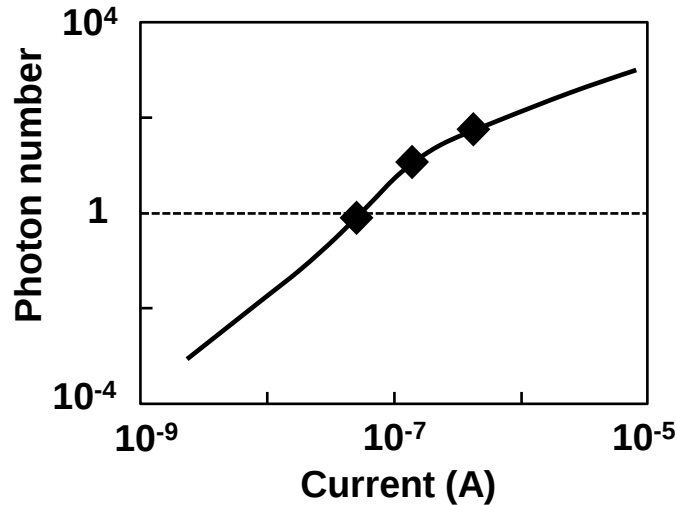
$$N_{\text{QD}} = 50, \Delta_{\text{inh}} = 20 \text{ meV}$$

Criterial for lasing: Population clamping and hole burning

$\beta = 0.01$



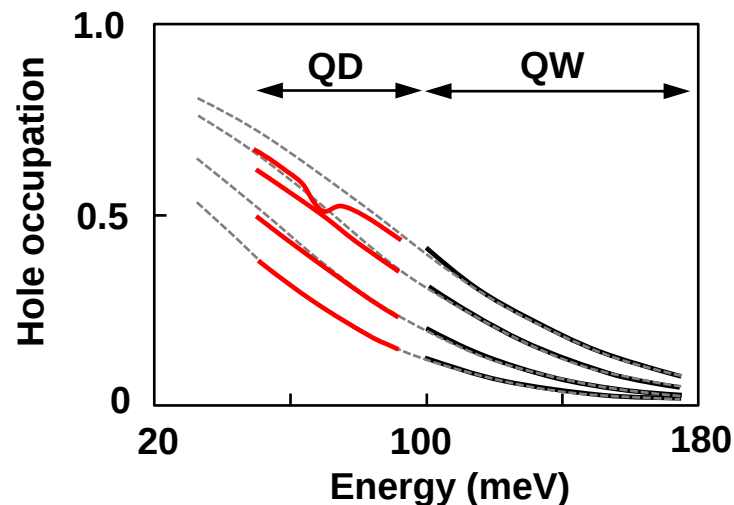
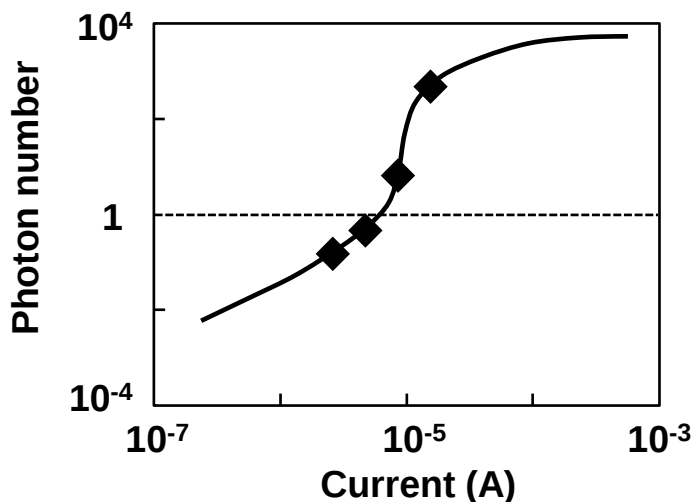
$\beta = 1$



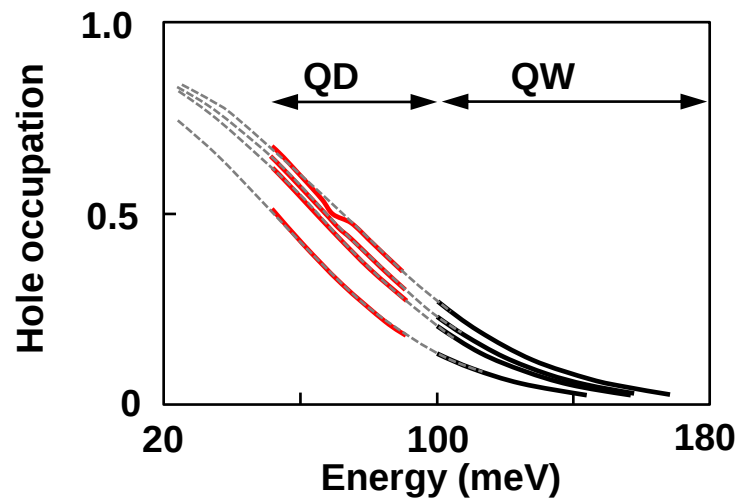
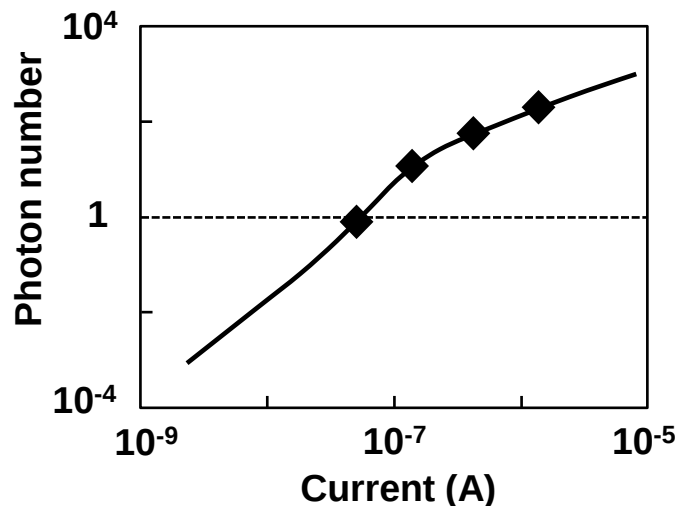
$$N_{\text{QD}} = 50, \Delta_{\text{inh}} = 20\text{meV}, \gamma_c = 10^{10}\text{s}^{-1}$$

Criterial for lasing: Population clamping and hole burning

$\beta = 0.01$



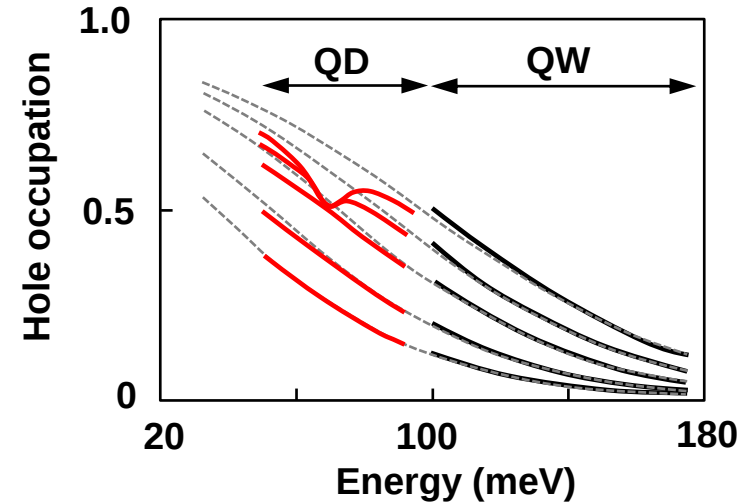
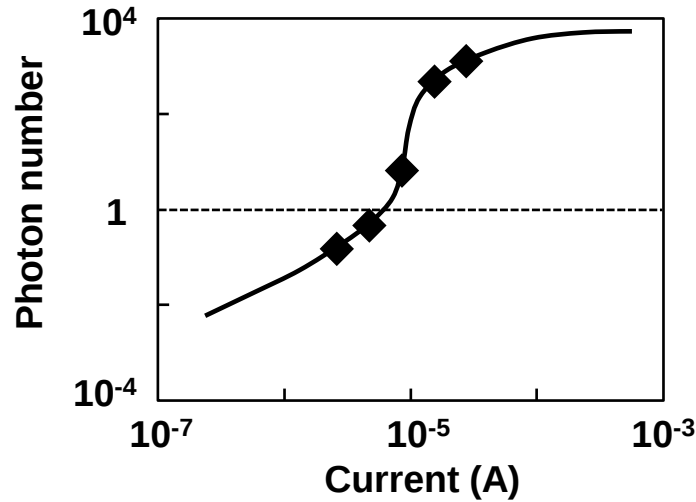
$\beta = 1$



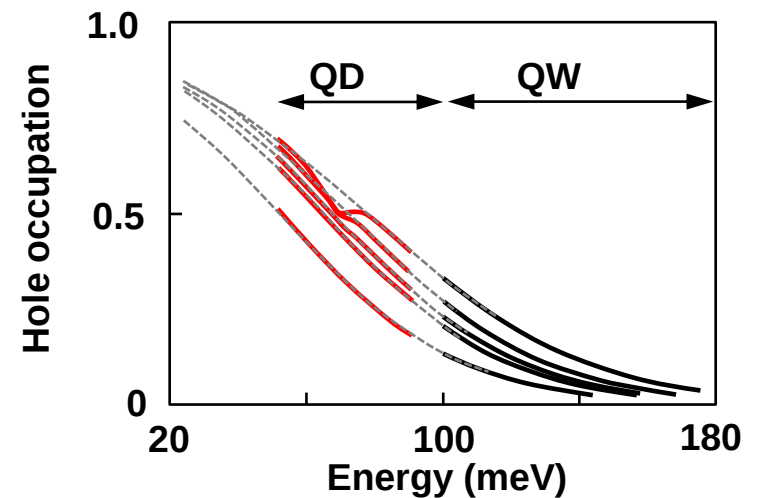
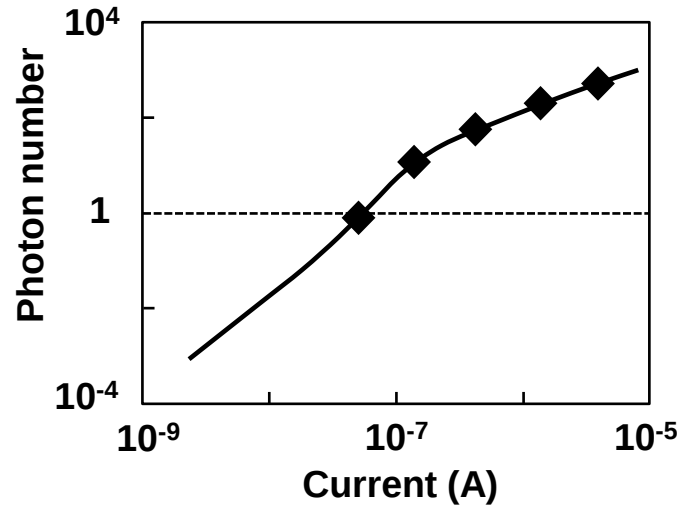
$$N_{\text{QD}} = 50, \Delta_{\text{inh}} = 20\text{meV}, \gamma_{\text{c}} = 10^{10}\text{s}^{-1}$$

Criterial for lasing: Population clamping and hole burning

$\beta = 0.01$



$\beta = 1$



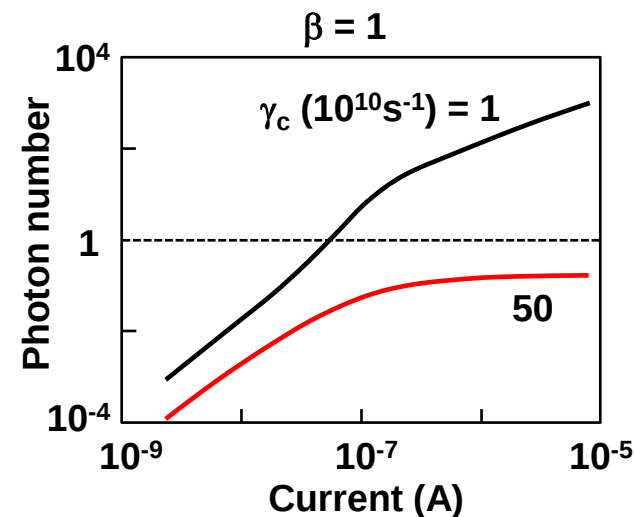
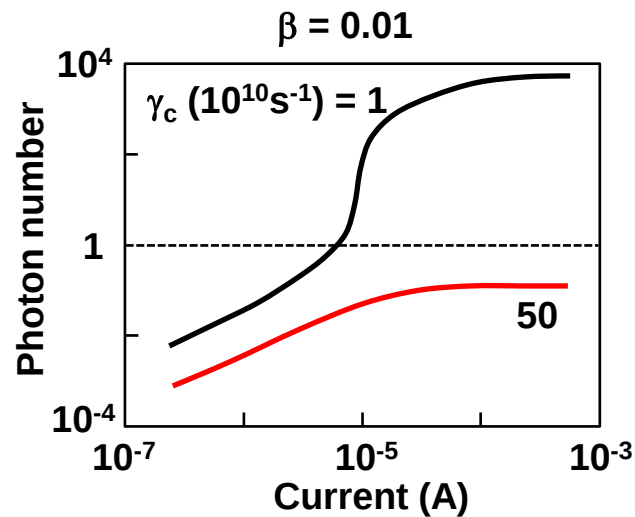
$$N_{\text{QD}} = 50, \Delta_{\text{inh}} = 20\text{meV}, \gamma_c = 10^{10}\text{s}^{-1}$$

Criterial for lasing: Stimulated emission

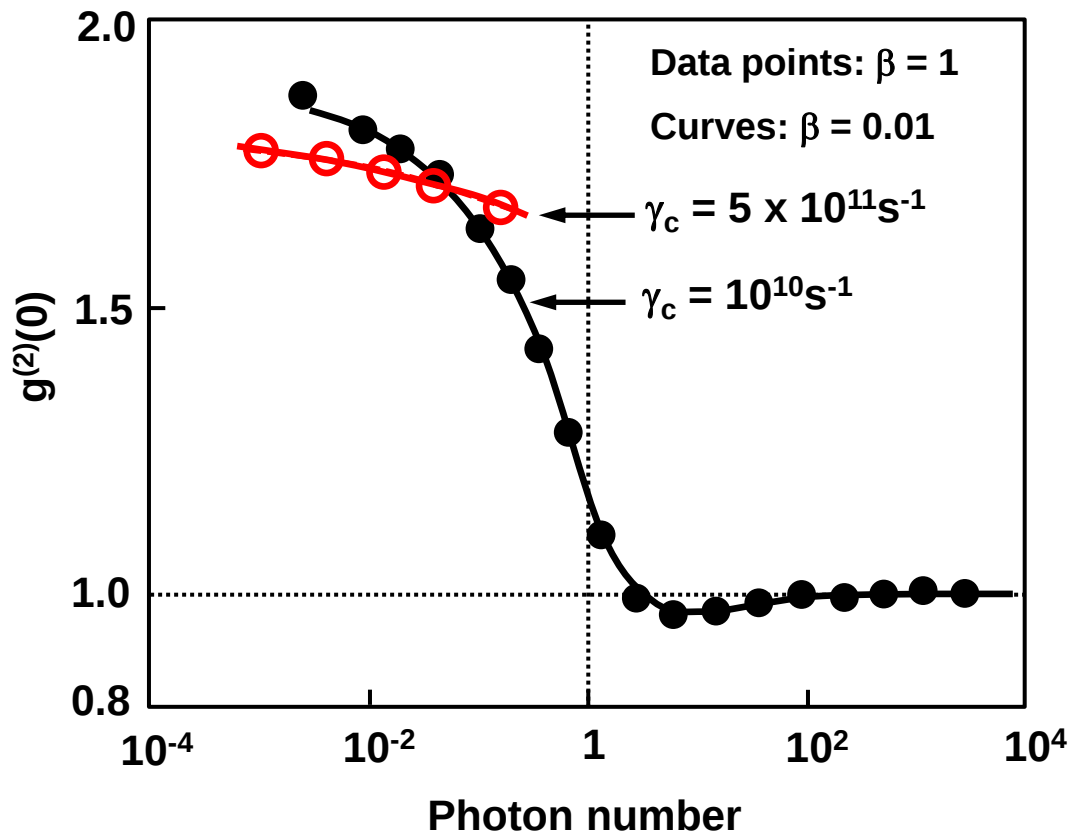
$$\frac{dn}{dt} = -2\gamma_c n + \gamma_l(n+1)N$$

Stimulated
emission

Spontaneous
emission

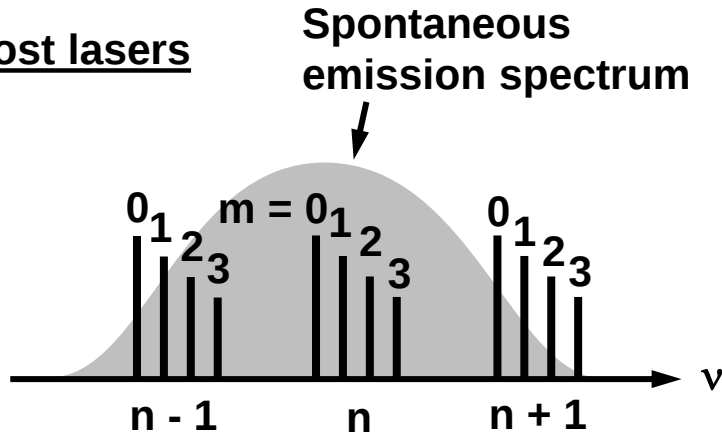


$$N_{\text{QD}} = 50, \Delta_{\text{inh}} = 20 \text{meV}$$

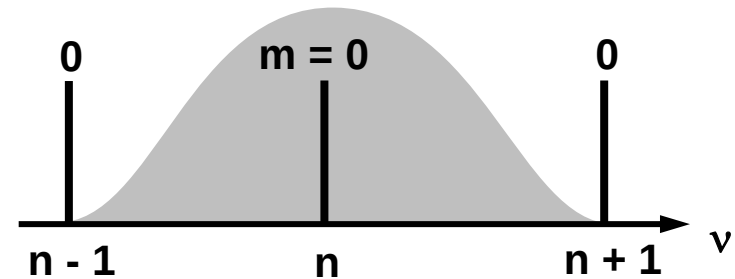


Why work on nanolasers?

Most lasers

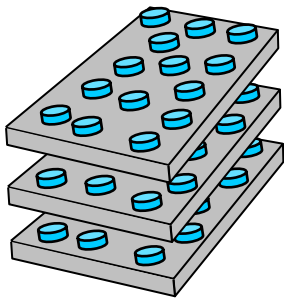


Nanolasers

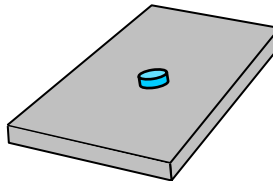


All emission into single resonator mode

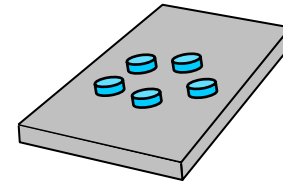
Typical QD-laser active region



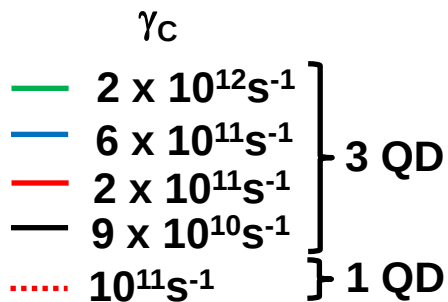
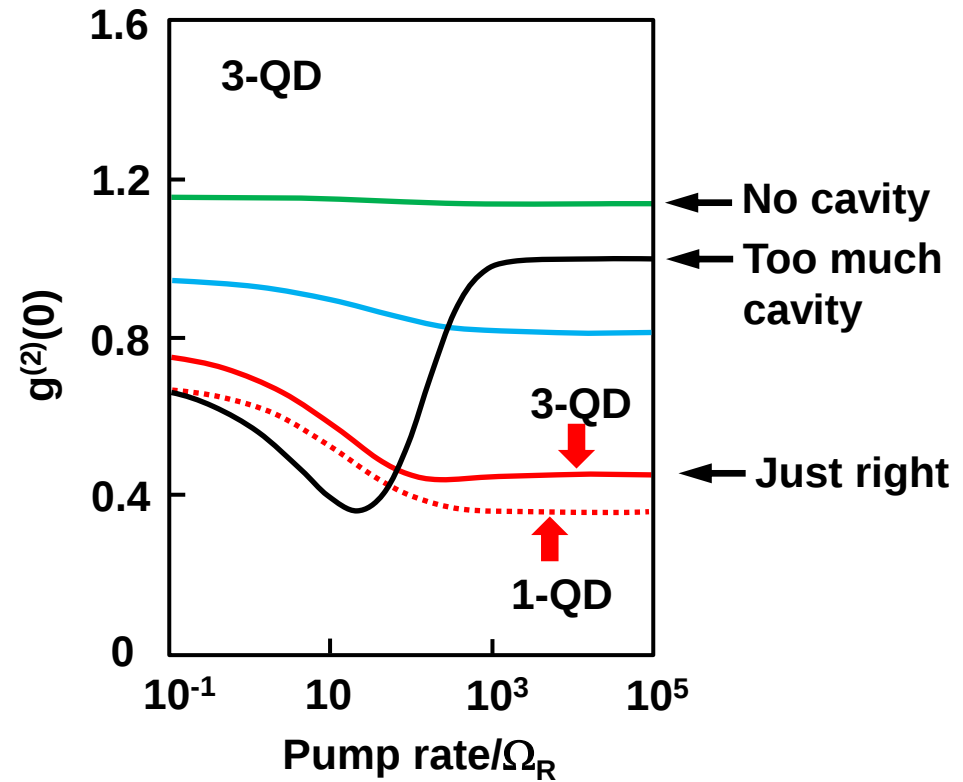
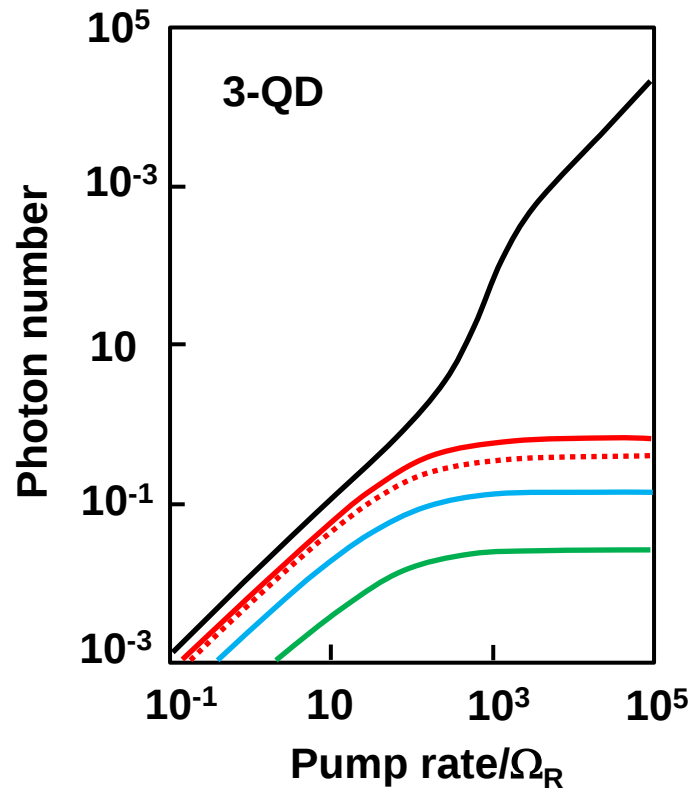
Single-photon source active region



Few-QD active region

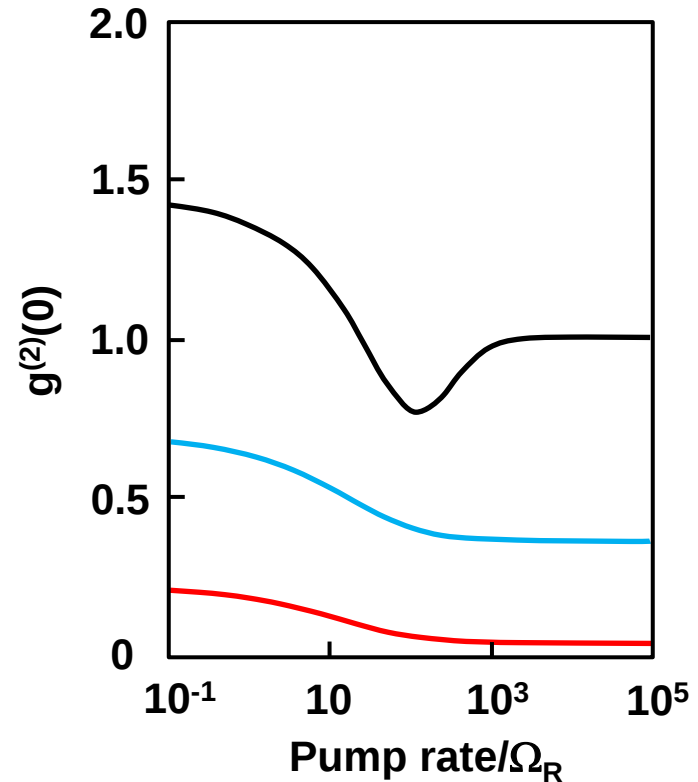
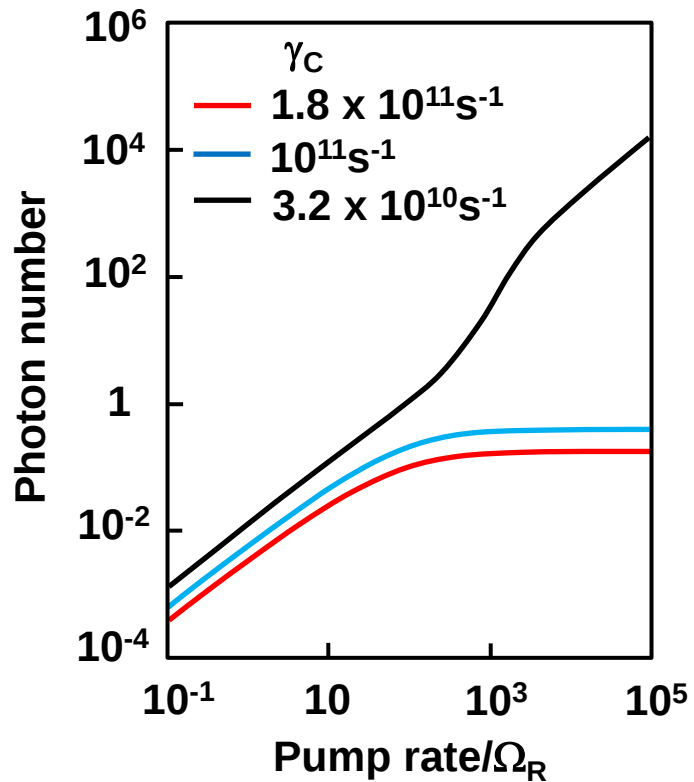


Higher single-photon production rate with few-QD active region



Single-quantum-dot active region

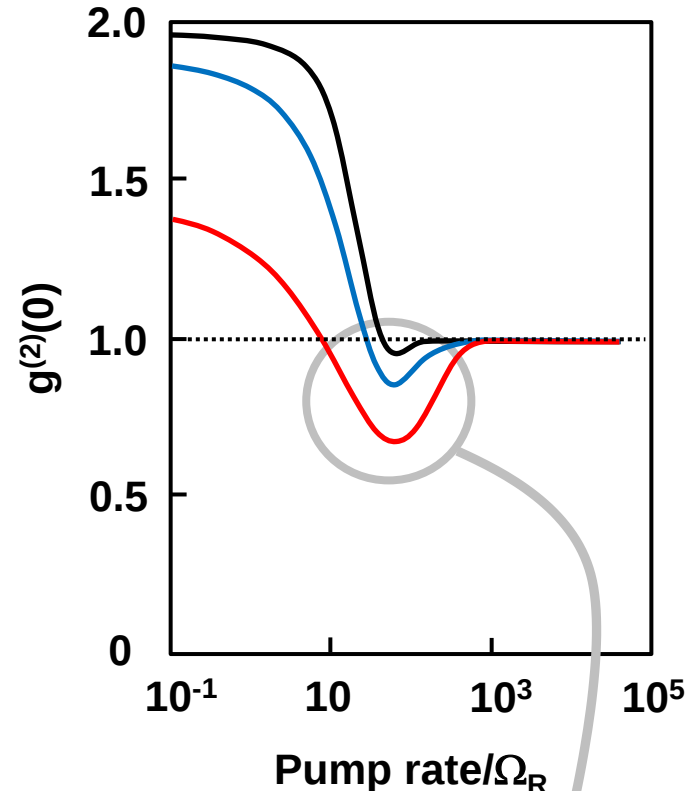
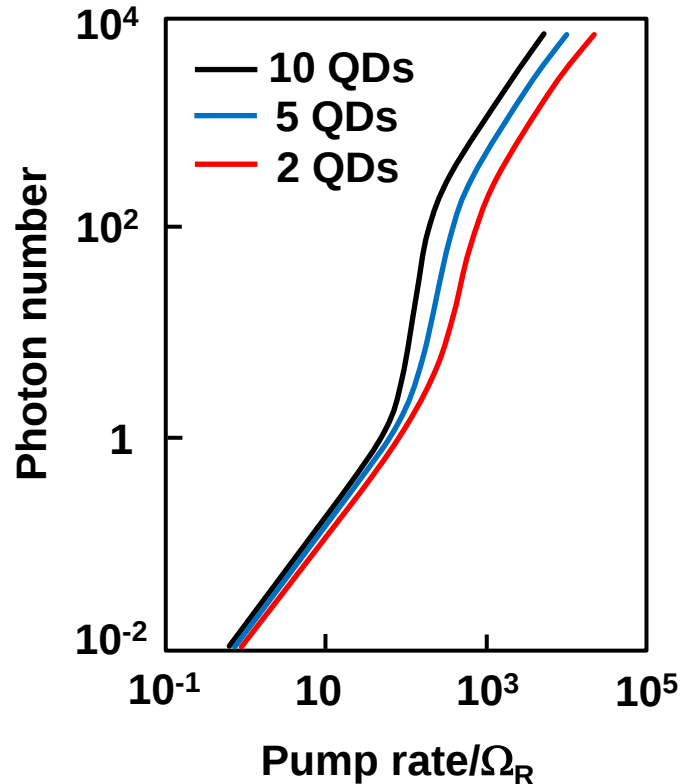
Single-photon source $\rightarrow$ Nonclassical-light source $\rightarrow$ Laser



Vacuum Rabi frequency: $\Omega_R = 2.7 \times 10^{11} \text{ s}^{-1}$

Few-quantum-dot active region

Question: Can we increase single-photon production rate?



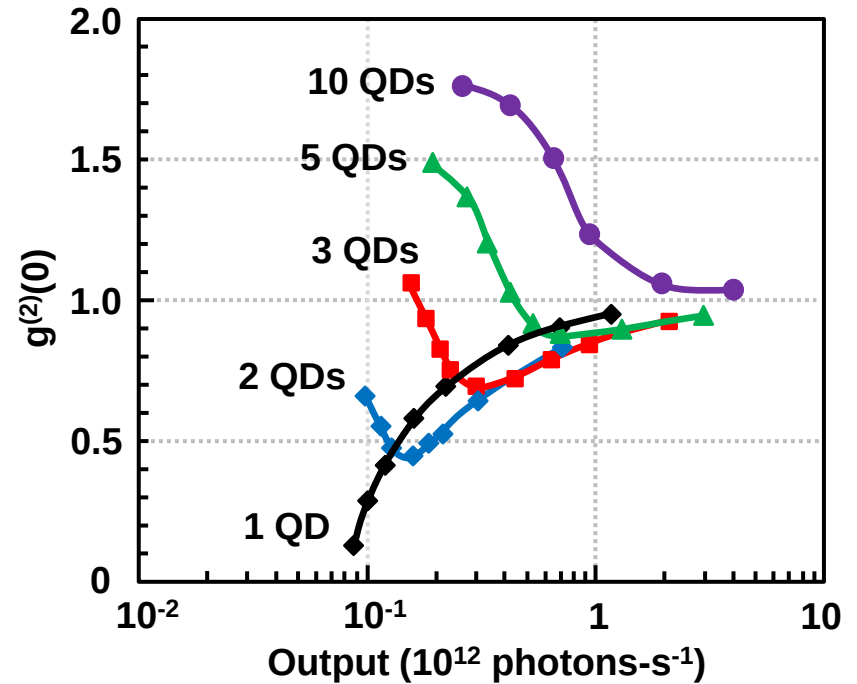
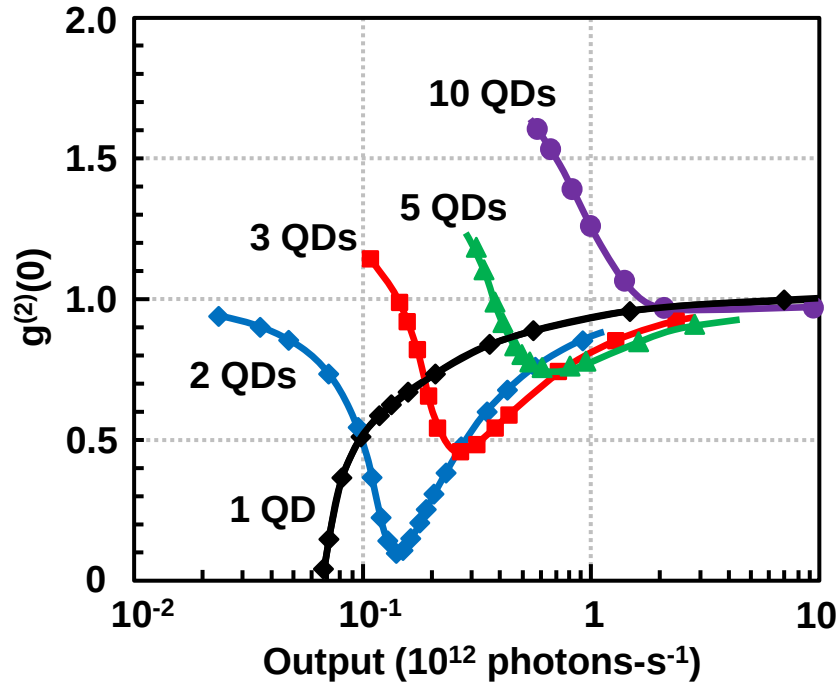
QD	γ_c
10	$2 \times 10^{11} \text{ s}^{-1}$
5	10^{11} s^{-1}
2	$6 \times 10^{10} \text{ s}^{-1}$

Light-carrier correlation

Leading terms:

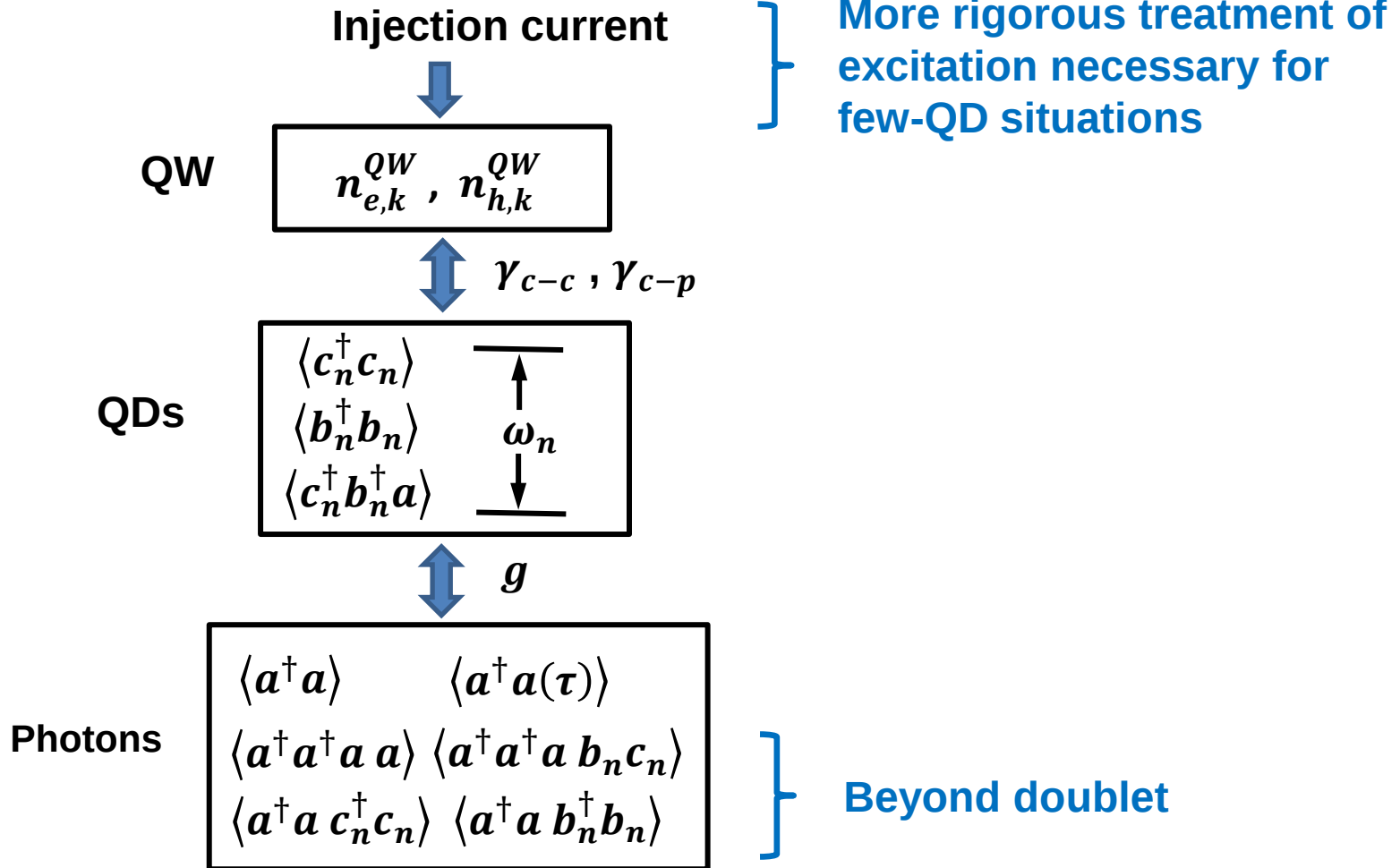
$$\delta\langle c^\dagger c a^\dagger a \rangle, \delta\langle b^\dagger b a^\dagger a \rangle, \\ \delta\langle b^\dagger c^\dagger a^\dagger a a \rangle$$

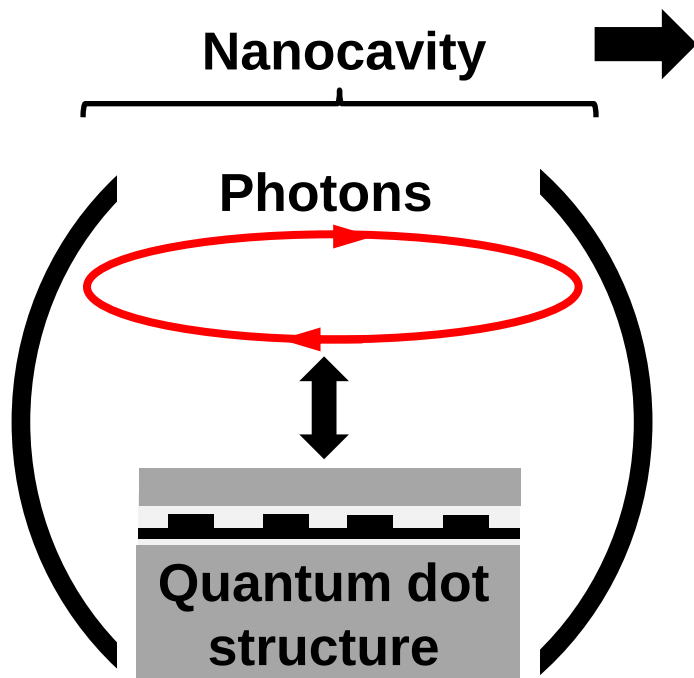
Increasing single-photon production rate with few-emitter active region



Concerns and next refinements

Present model





Constitutive equation

$$D = \epsilon E = [\epsilon_r(\mathbf{r}) + i\epsilon_i(\mathbf{r})]E + \mathbf{P}$$

from QDs

Passive cavity:

$$\nabla \times E = -\frac{\partial B}{\partial t}$$

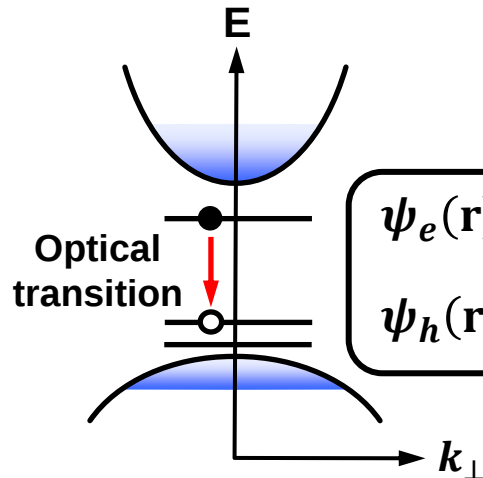
$$\nabla \times B = -\mu_0 \epsilon_r(\mathbf{r}) \frac{\partial E}{\partial t}$$

Solution

$$E \propto W(\mathbf{r})e^{-i\mathbf{v}t}$$

$$E(\mathbf{r}) = \hat{\epsilon} \sqrt{\frac{\hbar \mathbf{v}}{2\epsilon_b V}} W(\mathbf{r}) (a + a^\dagger)$$

Electronic structure



$$\psi_e(\mathbf{r}) = C(\mathbf{r}) \langle \mathbf{r} | \frac{1}{2}, s_z \rangle c_e$$

$$\psi_h(\mathbf{r}) = V(\mathbf{r}) \langle \mathbf{r} | m \rangle c_h$$

+ Adjoint

Nano-emitter model

$H = H_0$ *Single-particle energies*

$$-\hbar \sum_{\alpha} (g_{\alpha} b_{\alpha}^{\dagger} c_{\alpha}^{\dagger} a - g_{\alpha}^* a^{\dagger} c_{\alpha} b_{\alpha}) \quad \text{Light-carrier}$$

$\nearrow \wp \sqrt{v/\hbar \epsilon_b V} W(R_{QD}) \sum_n C_{\alpha}(R_n) V_{\alpha}(R_n)$

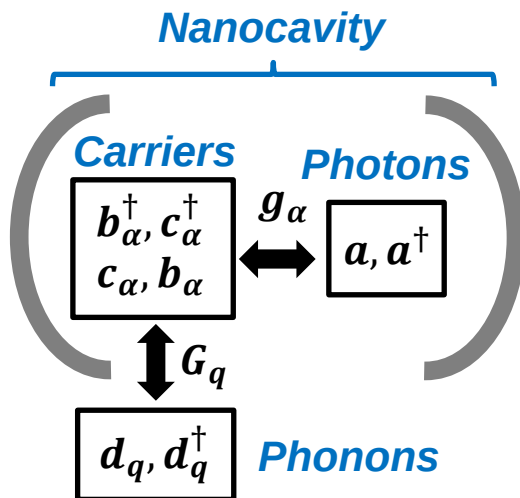
$$+\hbar \sum_{\alpha\beta q} G_q (c_{\alpha}^{\dagger} c_{\beta} + b_{\alpha}^{\dagger} b_{\beta}) (d_q + d_q^{\dagger}) \quad \text{Carrier-phonon}$$

$$+\frac{1}{2} \sum_{\alpha\beta\sigma\eta} W_{\sigma\eta}^{\alpha\beta} c_{\alpha}^{\dagger} c_{\beta}^{\dagger} c_{\eta} c_{\sigma} + \frac{1}{2} \sum_{\alpha\beta\sigma\eta} W_{\sigma\eta}^{\alpha\beta} b_{\alpha}^{\dagger} b_{\beta}^{\dagger} b_{\eta} b_{\sigma}$$

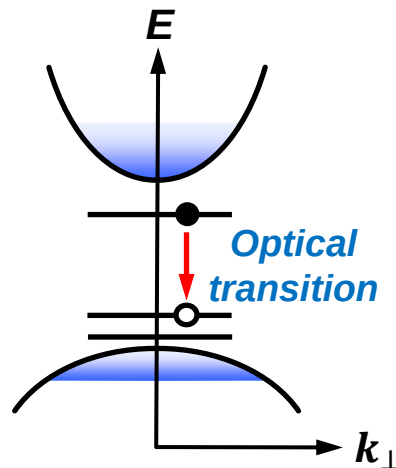
$\downarrow$ Matrix element of $\frac{e^2}{4\pi\epsilon_b |r - r'|}$

$$-\sum_{\alpha\beta\sigma\eta} W_{\sigma\eta}^{\alpha\beta} b_{\alpha}^{\dagger} c_{\beta}^{\dagger} c_{\eta} b_{\sigma}$$

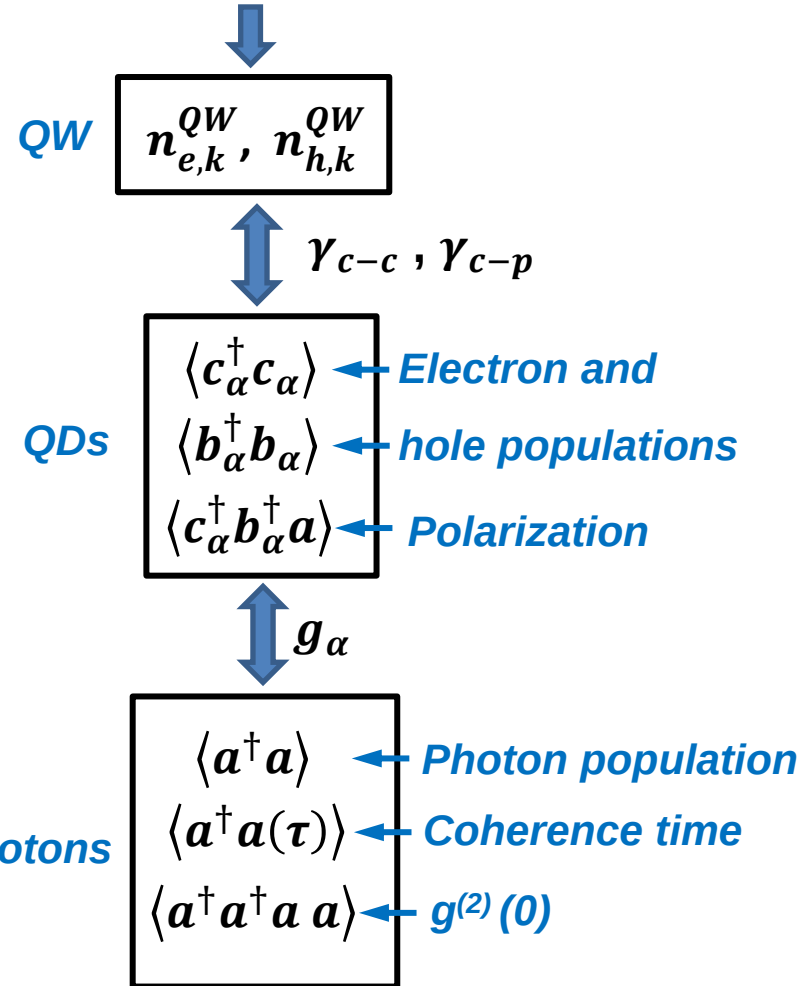
$\left. \begin{array}{l} \text{Carrier-carrier} \end{array} \right\}$



Electronic structure



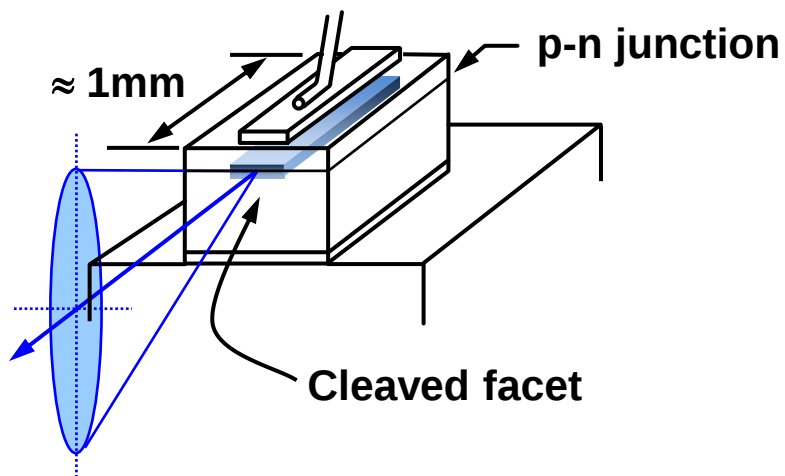
Injection current



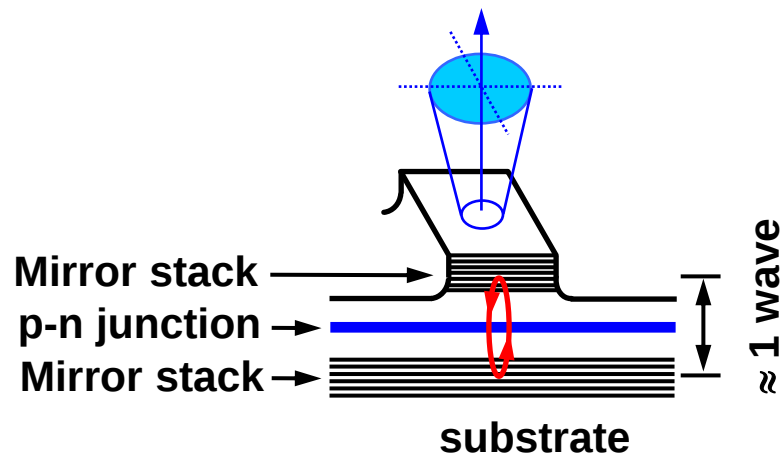
Emphasis now is on correlations involving light-matter interaction instead of Coulomb interaction

Towards smaller and smaller lasers

Edge - Emitting Laser



Vertical-Cavity Surface-Emitting Laser (VCSEL)



Nanolasers

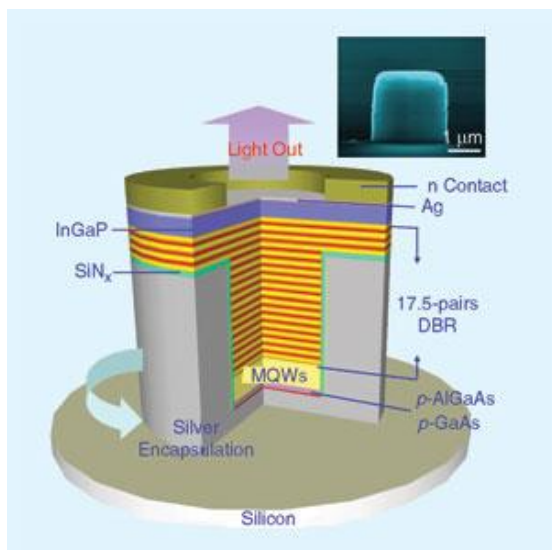


Figure 6. Combining semiconductors and metals enables nanolasers having a footprint roughly a factor 100 smaller than that of a classical VCSEL. Adapted from a figure originally created by C. Y. Lu et al., UIUC.

Taken from D. Bimberg, IEEE Photonics Society Research Highlights