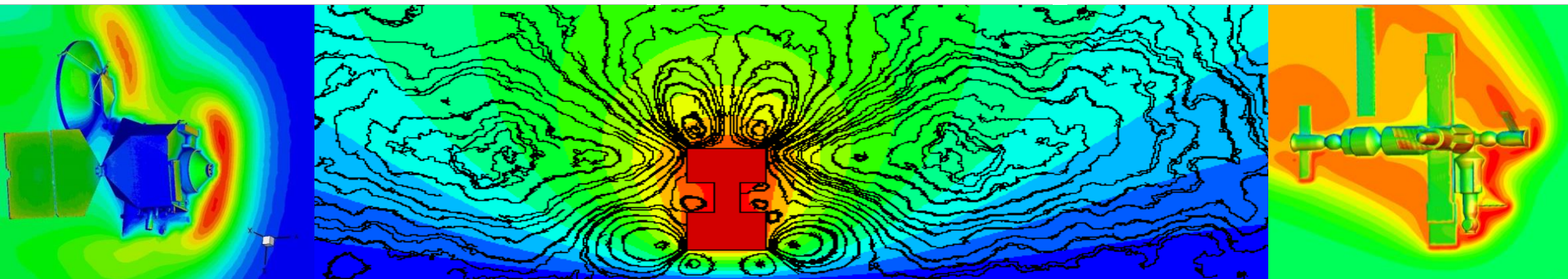


Exceptional service in the national interest



DSMC: The Quest for Speed

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Collaborators



John Torczynski



Dan Rader



Tim Koehler



Steve Plimpton

Boltzmann Equation and the Direct Simulation Monte Carlo Method



**Ludwig
Boltzmann**

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{x}} + \frac{\mathbf{F}}{m} \cdot \frac{\partial f}{\partial \mathbf{v}} = \int_{-\infty}^{\infty} \int_0^{4\pi} \left(f^* f_1^* - f f_1 \right) |\mathbf{v} - \mathbf{v}_1| S d\mathbf{W} d\mathbf{v}_1$$

molecular motion and
force-induced acceleration

pairwise molecular collisions
(molecular chaos)

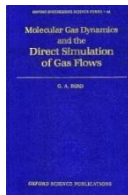
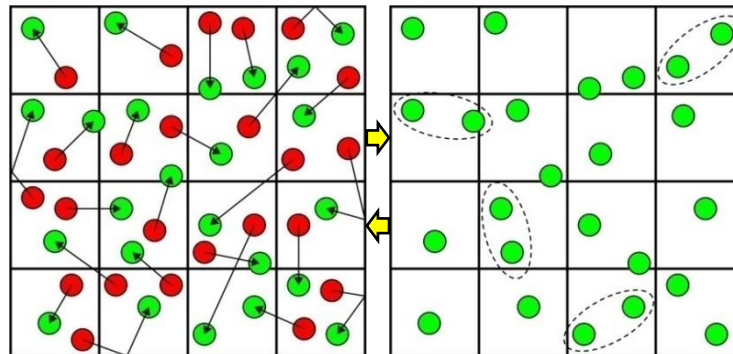


**James Clerk
Maxwell**

The objective of DSMC is to simulate complicated gas flows using only collision mechanics of simulated molecules

molecules move

molecules collide



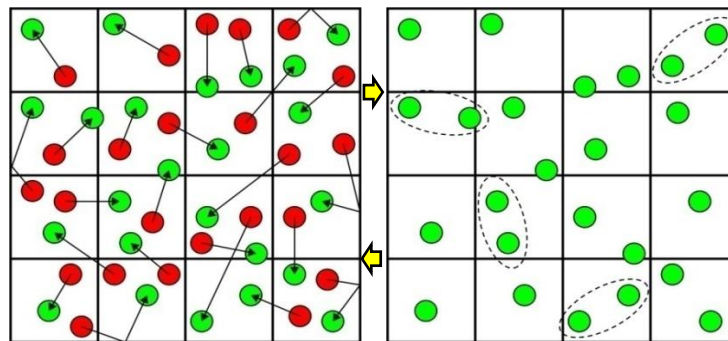
**Graeme Bird
(1963, 1994)**

DSMC vs. Boltzmann Equation

- Instead of solving Newton's laws of motion (Molecular Dynamics), DSMC replaces explicit intermolecular forces with stochastic collisions
- It has been shown that DSMC is **equivalent** to solving the Boltzmann equation (Nambu 1980, Babovsky 1989, Wagner 1992)
- DSMC has been shown to reproduce **exact** known solutions (Chapman-Enskog, Moment Hierarchy) of the Boltzmann equation (Gallis et al. 2004, 2006) for **non-equilibrium** flows
- In fact, DSMC is **superior** to solving the Boltzmann equation
 - DSMC can **model complicated processes** (e.g., polyatomic molecules, chemically reacting flows, ionized flows) for which **Boltzmann-type transport equations are not even known** (Struchtrup 2005)
 - DSMC **includes fluctuations**, which have been shown to be physically realistic (Garcia 1990) but which are **absent from the Boltzmann equation**
- **DSMC has tremendous potential**
- **1000x speedup** is desired to address problems of interest

Evolution of DSMC Procedures


- Since the original DSMC algorithm (1963), there have been no major changes to it
 - Advanced molecular interactions, physical phenomena have been added
- The DSMC algorithm has been criticized (**unfairly**) as being computationally inefficient
- Numerous alternative algorithms or procedures have been proposed that try to address this concern
- Alternative methods come with **undeniable virtues** but **accompanied vices**.
- There is no standardized test suite to allow comparative evaluations of new methods, which has led to a proliferation of *new methods*
- More than 50 years since its introduction, DSMC is still the **predominant** algorithm



Alternative Procedures

- Plume flows
 - Hybrid schemes (NS-DSMC)
- Non-reacting flows
 - Discrete velocity schemes, moment methods, higher-order CFD
- Low signal/noise flows (MEMS)
 - PDE-like, numerical solutions of the Boltzmann equation
 - Noise reduction schemes
 - Simplified molecular interactions (BGK/ES-BGK)
- High-density flows
 - Near-neighbor collisions
- Free-molecular flows
 - Analytical solutions

How can we evaluate new techniques & procedures?

- There is a very large number of alternative procedures with an **ever-growing number of variants**
 - Although there are particular applications where some variants may be preferable, it is unclear how to evaluate
 - Generality
 - Accuracy
 - Efficiency
- 
- **Accuracy** → Comparison with analytical solutions
 - **Efficiency** → Comparison with DSMC convergence rate

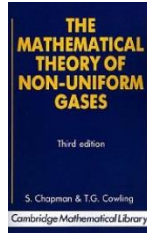
Chapman-Enskog (CE) Theory



Sydney
Chapman



David
Enskog



$$\begin{aligned}
 f &= f^{(0)}(1 + F^{(1)} + Y^{(1)}) & f^{(0)} &= (n/\rho^{3/2}c_m^3)\exp[-\tilde{c}^2] \\
 c_m &= \sqrt{2k_B T/m} & \tilde{\mathbf{c}} &= \mathbf{c}/c_m \quad \mathbf{c} = \mathbf{v} - \mathbf{u} \\
 F^{(1)} &= -\left(8/5\right)\tilde{A}[\tilde{\mathbf{c}}]\tilde{\mathbf{c}}\times\tilde{\mathbf{q}} & Y^{(1)} &= -2\tilde{B}[\tilde{\mathbf{c}}](\tilde{\mathbf{c}}\circ\tilde{\mathbf{c}}:\tilde{\mathbf{t}}) \\
 \mathbf{K} &= -(5/4)k_B c_m^2 a_1 & \mathbf{m} &= (1/2)mc_m^2 b_1 \\
 \tilde{A}[\tilde{\mathbf{c}}] &= \sum_{k=1}^{\infty} (a_k/a_1) S_{3/2}^{(k)}[\tilde{c}^2] & \tilde{B}[\tilde{\mathbf{c}}] &= \sum_{k=1}^{\infty} (b_k/b_1) S_{5/2}^{(k-1)}[\tilde{c}^2] \\
 C_p &= (5/2)(k_B/m) & \text{Pr} &= (2/3)(m_{\infty}/m_1)(K_1/K_{\infty})
 \end{aligned}$$

- Chapman and Enskog analyzed Boltzmann collision term
 - Perturbation expansion using Sonine polynomials
 - Near equilibrium, appropriate in continuum limit
- Determined velocity distribution and transport properties
 - Thermal conductivity \mathbf{K} , viscosity μ , mass self-diffusivity \mathbf{D}
 - Prandtl number Pr from “infinite-to-first” ratios K_{∞}/K_1 , μ_{∞}/μ_1
 - Distribution “shape”: Sonine polynomial coeffs. a_k/a_1 , b_k/b_1
 - Values for all Inverse-Power-Law (IPL) interactions
 - Maxwell and hard-sphere are special cases

Extracting CE Parameters from DSMC

$$\begin{aligned}
 q &= K_{\text{eff}} \left(\frac{\partial T}{\partial x} \right) & \frac{a_k}{a_1} &= \sum_{i=1}^k \left(\frac{(-1)^{i-1} k! (5/2)!}{(k-i)! i! (i + (3/2))!} \right) \left(\frac{\langle \tilde{c}^{2i} \tilde{c}_x \rangle}{\langle \tilde{c}^2 \tilde{c}_x \rangle} \right) & \tilde{\mathbf{c}} &= \frac{\mathbf{v} - \mathbf{V}}{c_m} \\
 \tau &= \mu_{\text{eff}} \left(\frac{\partial V}{\partial x} \right) & \frac{b_k}{b_1} &= \sum_{i=1}^k \left(\frac{(-1)^{i-1} (k-1)! (5/2)!}{(k-i)! (i-1)! (i + (3/2))!} \right) \left(\frac{\langle \tilde{c}^{2(i-1)} \tilde{c}_x \tilde{c}_y \rangle}{\langle \tilde{c}_x \tilde{c}_y \rangle} \right) & c_m &= \sqrt{\frac{2k_B T}{m}}
 \end{aligned}$$

DSMC moments of velocity distribution function

- Temperature T , velocity V
- Heat flux q , shear stress τ
- Higher-order moments

DSMC values for VSS molecules (variable-soft-sphere)

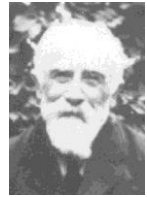
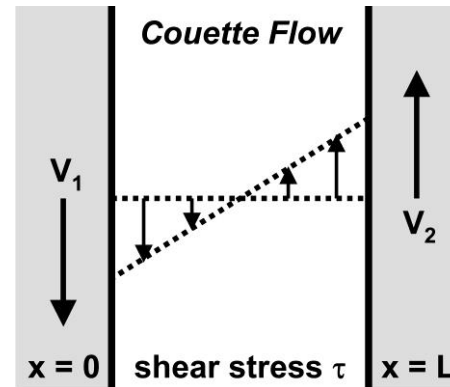
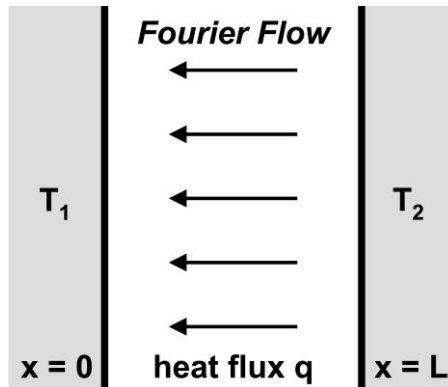
- Thermal conductivity and viscosity: K_{eff} and μ_{eff}
- Sonine-polynomial coefficients: a_k/a_1 and b_k/b_1
- Applicable for arbitrary Kn_L , Kn_q , Kn_τ

Fourier and Couette Flow



Joseph
Fourier

$$q = -K \frac{\partial T}{\partial x}$$



Maurice
Couette

$$\tau = \mu \frac{\partial v}{\partial x}$$

Investigate transport in gas between parallel plates

- Fourier flow: heat conduction in stationary gas
- Couette flow: momentum transport in isothermal shear flow

Apply DSMC to Fourier flow and Couette flow

- Heat flux, shear stress: one-dimensional, steady

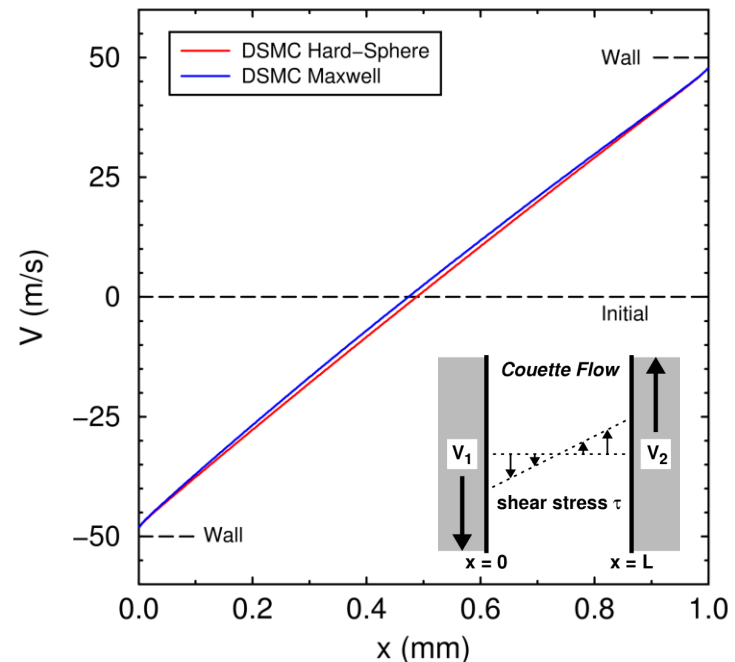
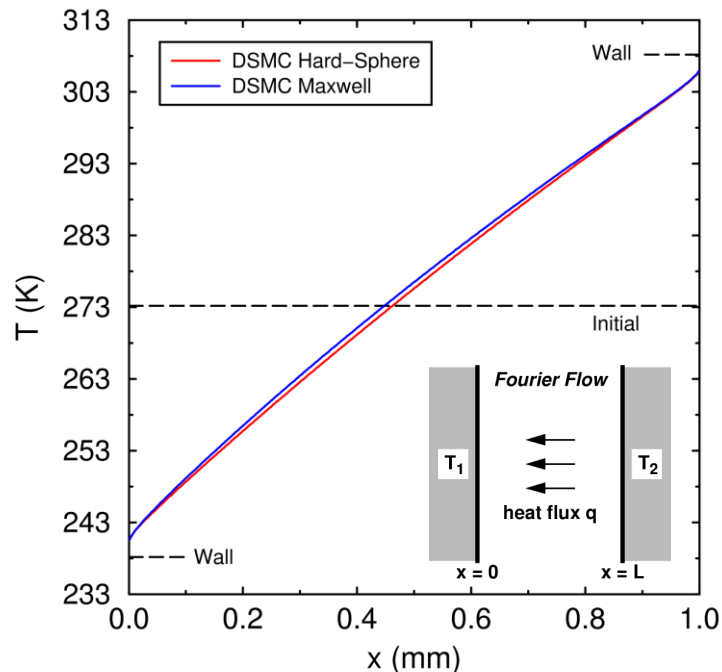
Compare DSMC to analytical “normal solutions”

- Normal: outside Knudsen layers
- Solutions: Chapman-Enskog (CE), Moment-Hierarchy (MH)

Verify DSMC accuracy at arbitrary heat flux, shear stress

- Thermal conductivity, viscosity; velocity distribution

Temperature and Velocity Profiles

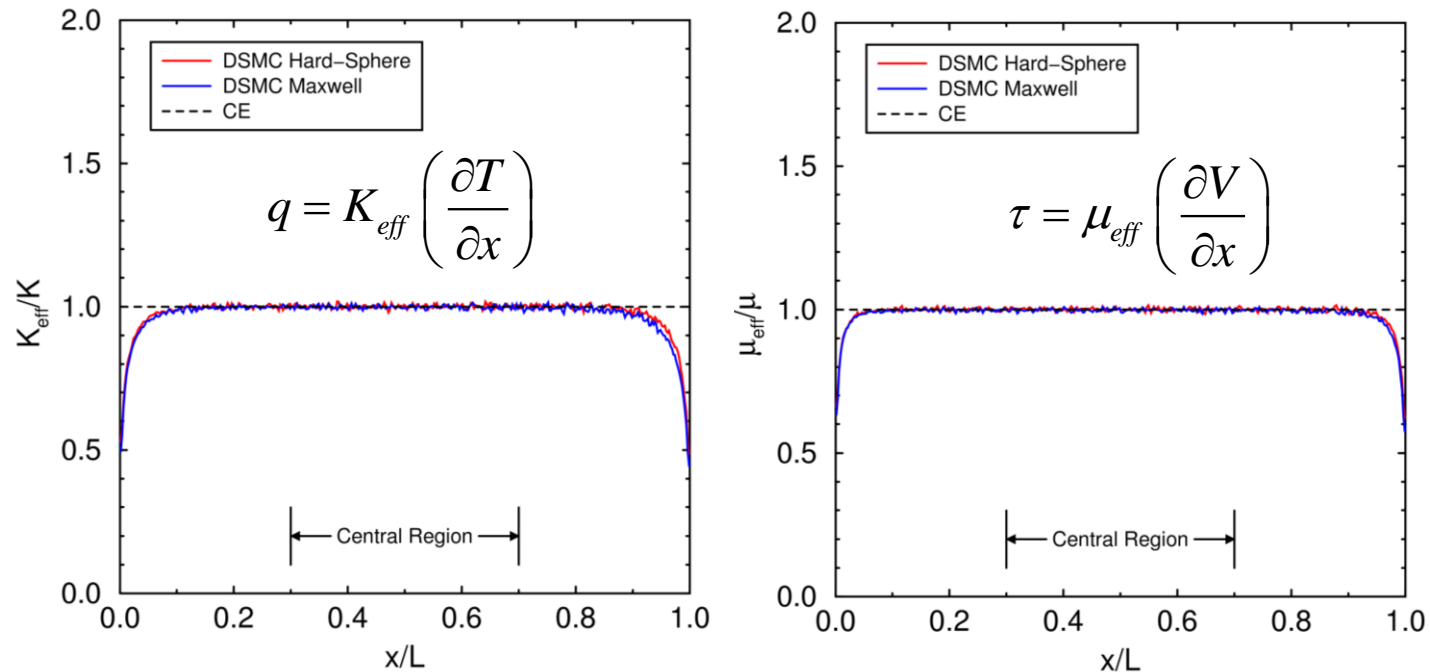


Low heat flux and shear stress: $Kn_q = 0.006$, $Kn_\tau = 0.003$

- Argon-like: initial $T = 273.15$ K, $p = 266.644$ Pa, $\lambda = 24$ μ m
- Walls: $L = 1$ mm $= 42\lambda$, $\Delta T = 70$ K, $\Delta V = 100$ m/s
- $N_c = 120$, $\Delta t = 7$ ns, $\Delta x = 2.5$ μ m, $\sim 10^9$ samples/cell, 32 runs

Small velocity slips, temperature jumps, Knudsen layers

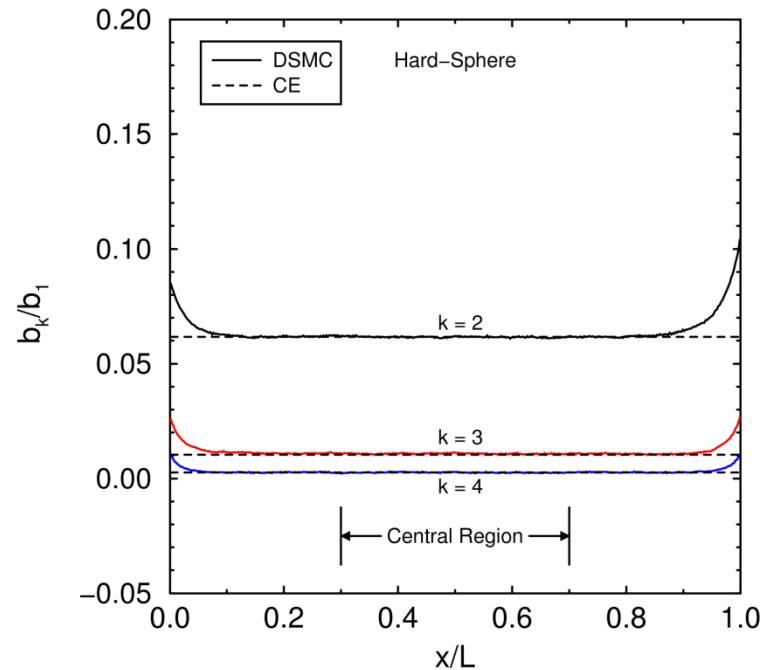
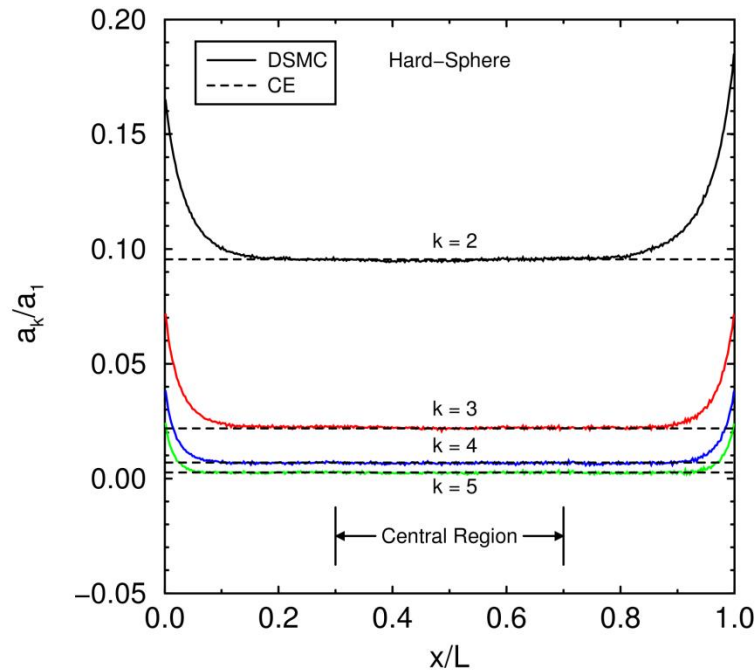
DSMC Reproduces Infinite-Approximation Chapman-Enskog Transport Coefficients



Thermal conductivity (left) and viscosity (right) away from walls

- Maxwell and hard-sphere results bound most gases
- Agreement with Chapman-Enskog theory verifies DSMC

DSMC Reproduces Infinite-Approximation Chapman-Enskog Velocity Distribution

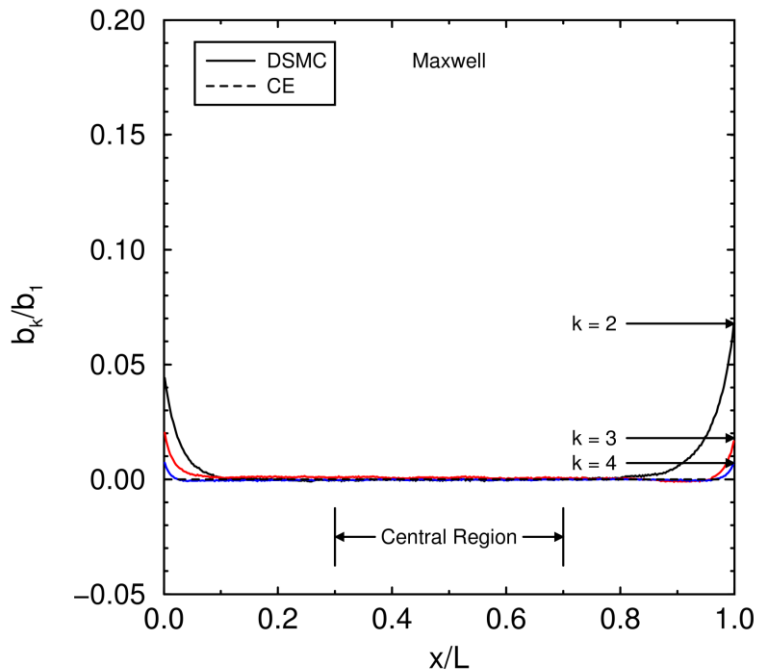
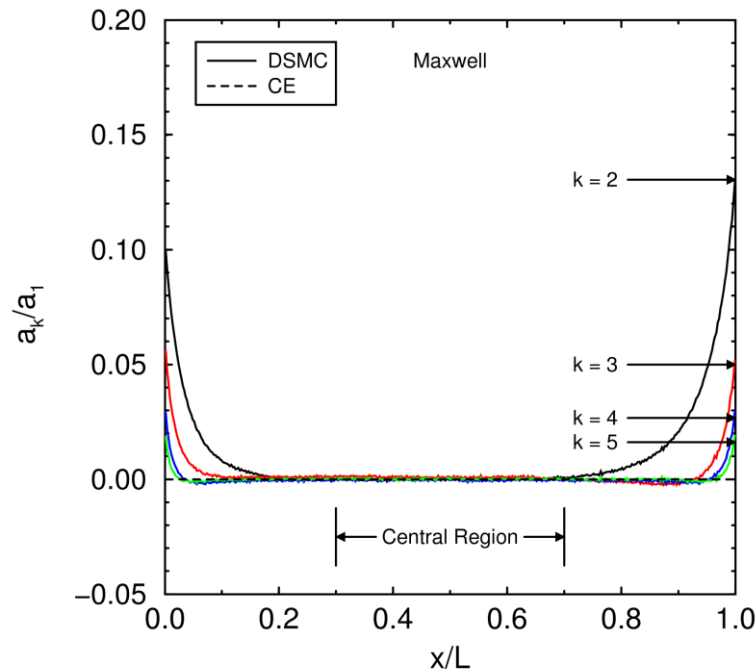


Sonine polynomial coefficients for temperature (left) & velocity (right) gradients

- Hard-sphere values are shown, other interactions have similar agreement
- Higher-order ($k > 5$) coefficients (not shown) also have similar agreement

Gallis M. A., Torczynski J. R., Rader D. J., "Molecular Gas Dynamics Observations of Chapman-Enskog Behavior and Departures Therefrom in Nonequilibrium Gases", *Physical Review E*, 69, 042201, 2004.

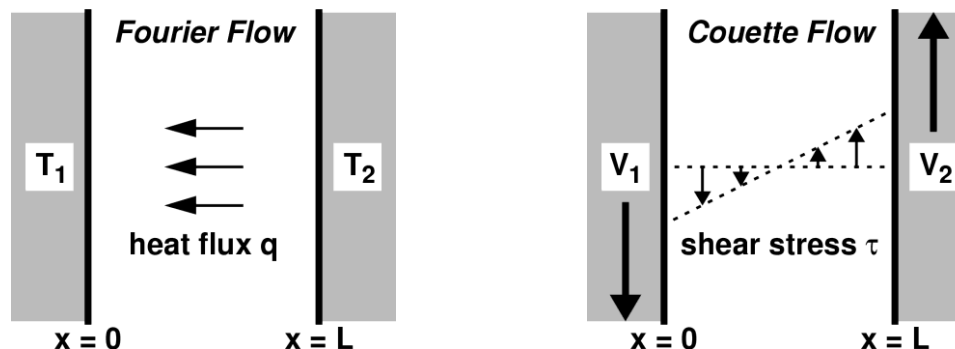
Maxwell Sonine-Coefficient Profiles



DSMC and CE Maxwell coefficients a_k/a_1 and b_k/b_1

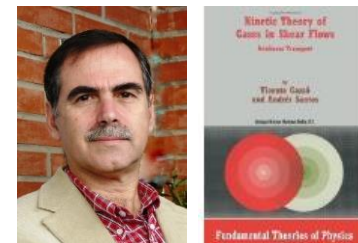
- Low heat flux, low shear stress: $Kn_q = 0.006$, $Kn_\tau = 0.003$
- Good agreement in central region: normal solution
- Knudsen layers easily observed: $\sim 10\%$ of domain

Moment-Hierarchy Method



Moment-Hierarchy (MH) normal solution

- MH solution extends CE solution to finite Kn_q and Kn_t
- Maxwell molecules: collision term quadratic in moments



Andres Santos

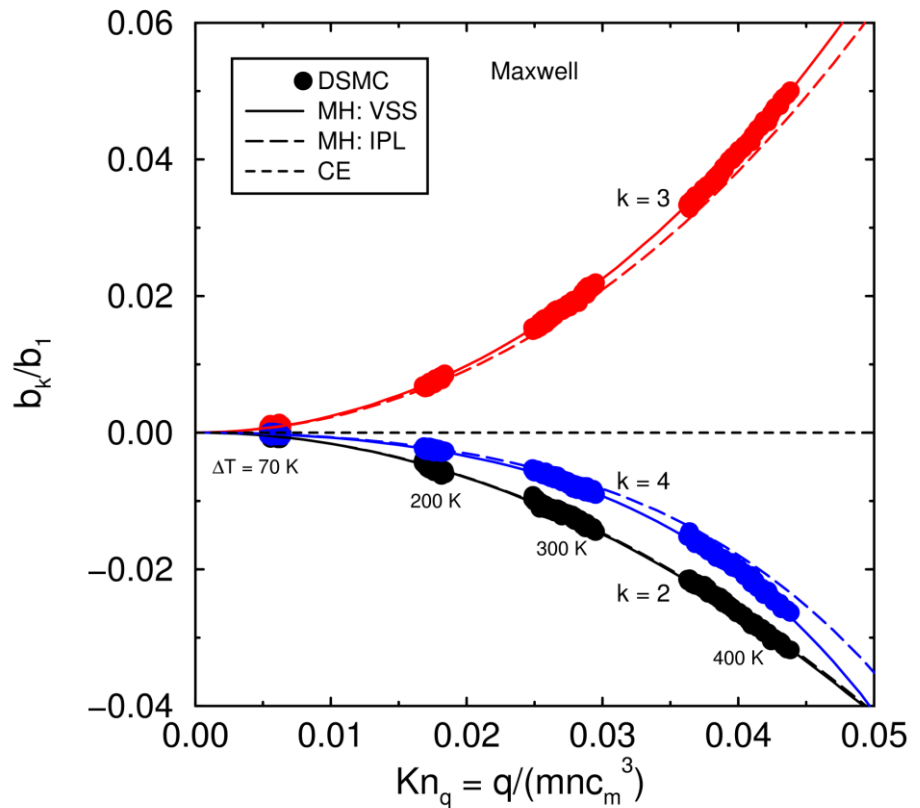
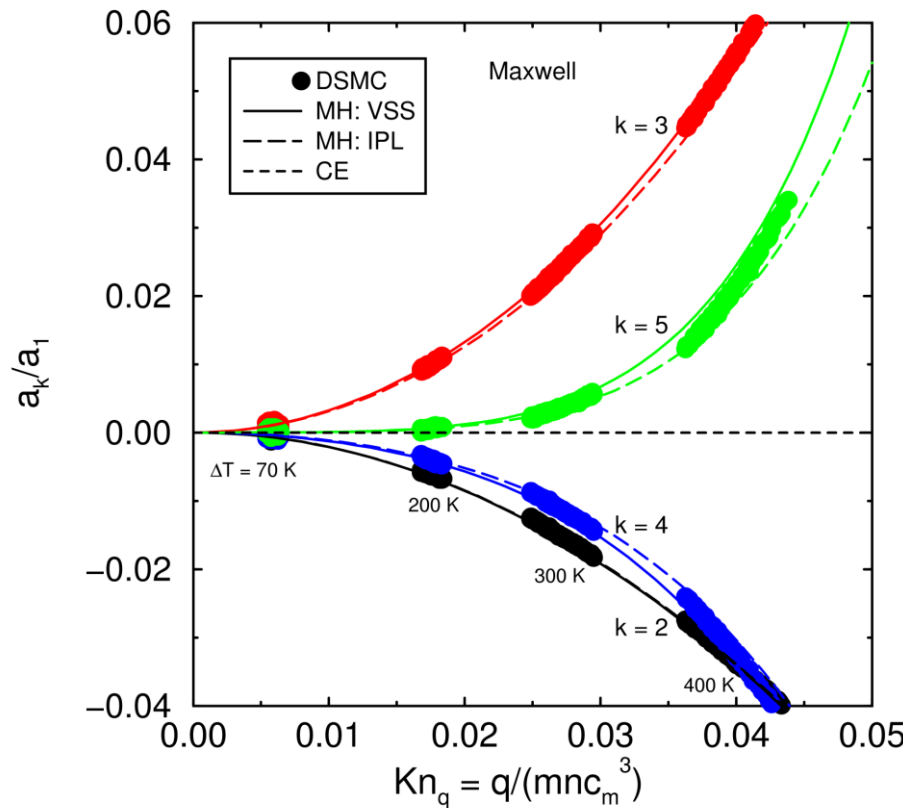
Compare DSMC to MH for Maxwell molecules

- Dependence of Sonine coefficients on Kn_q known

Apply DSMC for Maxwell molecules

Gallis M. A., Torczynski J. R., Rader D. J., Tij M., Santos A., "Normal Solutions of the Boltzmann Equation for Highly Nonequilibrium Fourier and Couette Flow", *Phys. Fluids*, 18, 017104, 2006.

Maxwell Normalized Sonine Coefficients



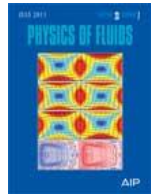
DSMC and MH Maxwell normal solutions for a_k/a_1 and b_k/b_1

- Four DSMC simulations: $\Delta T = 70, 200, 300, 400$ K
- MH: VSS-Maxwell (solid) and IPL-Maxwell (dashed) differ
- DSMC and MH VSS-Maxwell normal solutions agree

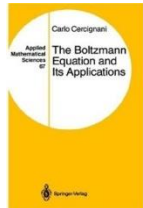
ES-BGK Kinetic Model



Lowell
Holway



Carlo
Cercignani



$$\left[\frac{df}{dt} \right]_{\text{ES}} = \nu_{\text{ES}} (f_{\text{ES}} - f)$$

$$\tilde{\mathbf{c}} = \mathbf{S}\tilde{\mathbf{C}} \quad \mathbf{P} = \int \mathbf{c}\mathbf{c}^T f d\mathbf{v}$$

$$K_{\text{ES}} = nk_B T C_p / \nu_{\text{ES}}$$

$$C_p = (5/2)(k_B/m)$$

$$f_{\text{ES}} = \frac{n}{\pi^{3/2} c_m^3 \det[\mathbf{S}]} \exp[-\tilde{\mathbf{C}}^T \tilde{\mathbf{C}}]$$

$$\mathbf{S}\mathbf{S} = \mathbf{I} - \left(\frac{1}{\text{Pr}} - 1 \right) \left(\frac{3\mathbf{P}}{\text{trace}[\mathbf{P}]} - \mathbf{I} \right)$$

$$\mu_{\text{ES}} = nk_B T \text{Pr} / \nu_{\text{ES}}$$

$$\text{Pr}_{\text{ES}} = \text{Pr}$$

Ellipsoidal-Statistical Bhatnagar-Gross-Krook
Holway (1963, 1965, 1966); Cercignani (1967, 1988)

ES-BGK collision term is generalization of BGK collision term

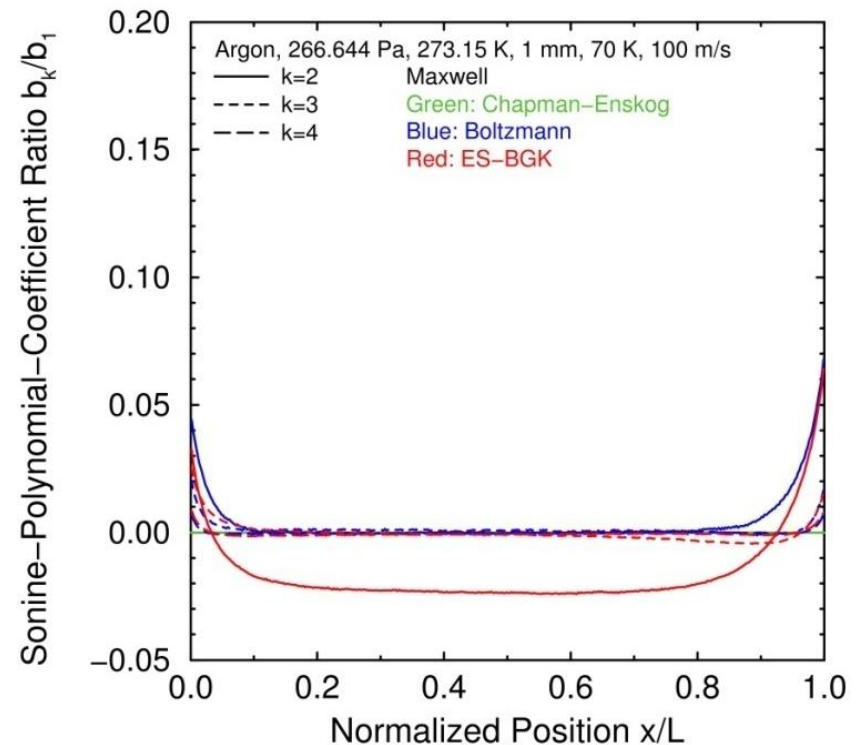
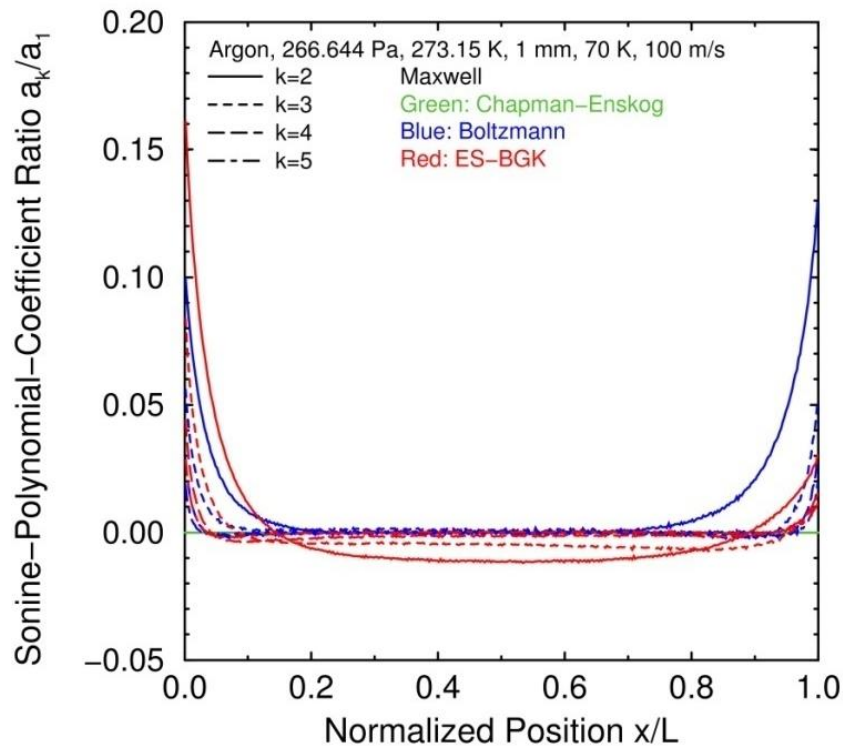
- Replaces Maxwellian with Ellipsoidal-Statistical distribution
- Prandtl number is specified independently (say, $\sim 2/3$)
- Drives distribution to equilibrium (Andries et al., 2000)

ES-BGK can presumably simulate a broad range of flows

- Can match thermal conductivity and viscosity simultaneously

Gallis M. A., Torczynski J. R., "Investigation of the ellipsoidal-statistical Bhatnagar-Gross-Krook kinetic model applied to gas-phase transport of heat and tangential momentum between parallel walls", Phys. Fluids, 23(3), 030601, 2011.

Distribution Shape: Maxwell

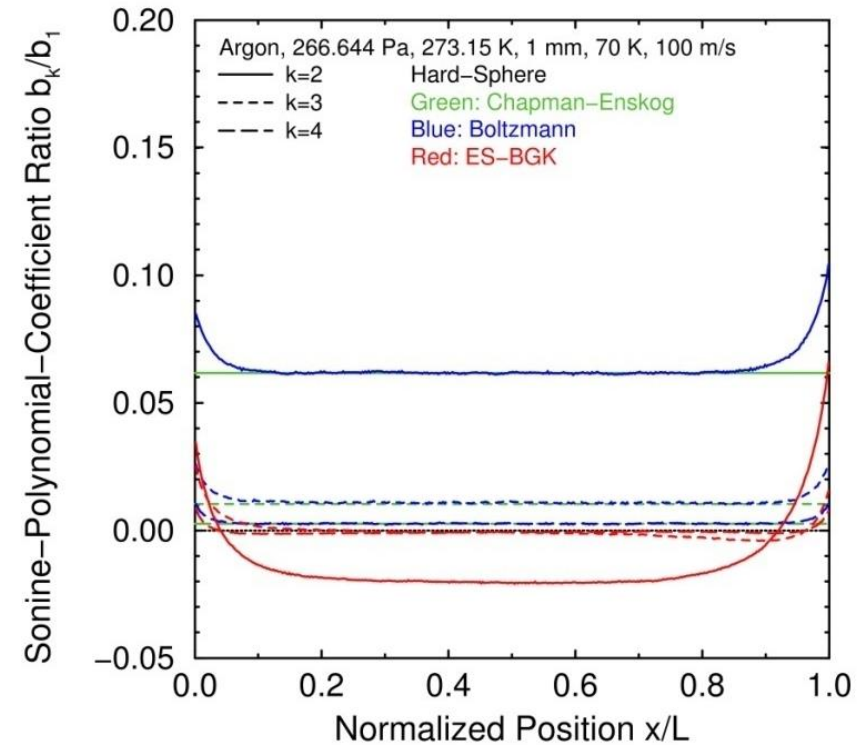
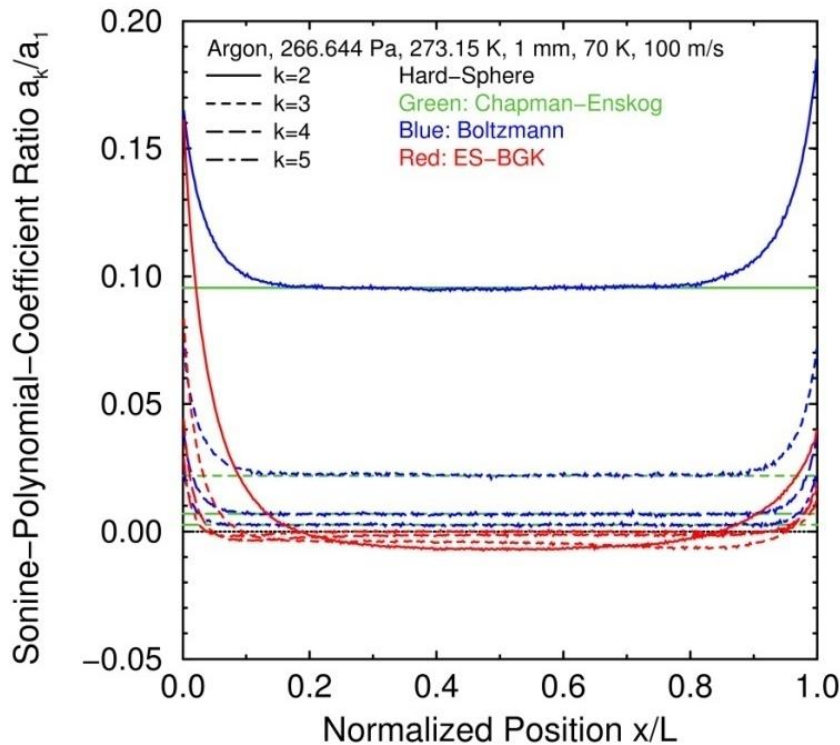


Sonine polynomial coefficients for Maxwell interaction

- Chapman-Enskog (continuum) values for $k \geq 2$ are all **zero**
- Boltzmann closely matches CE values away from walls
 - Systematic differences in Knudsen layers (transition regime)
- ES-BGK differs strongly from CE values away from walls
 - Discrepancy is largest for $k = 2$

ES-BGK distribution shape is not accurate in continuum and transition regimes

Distribution Shape: Hard-Sphere

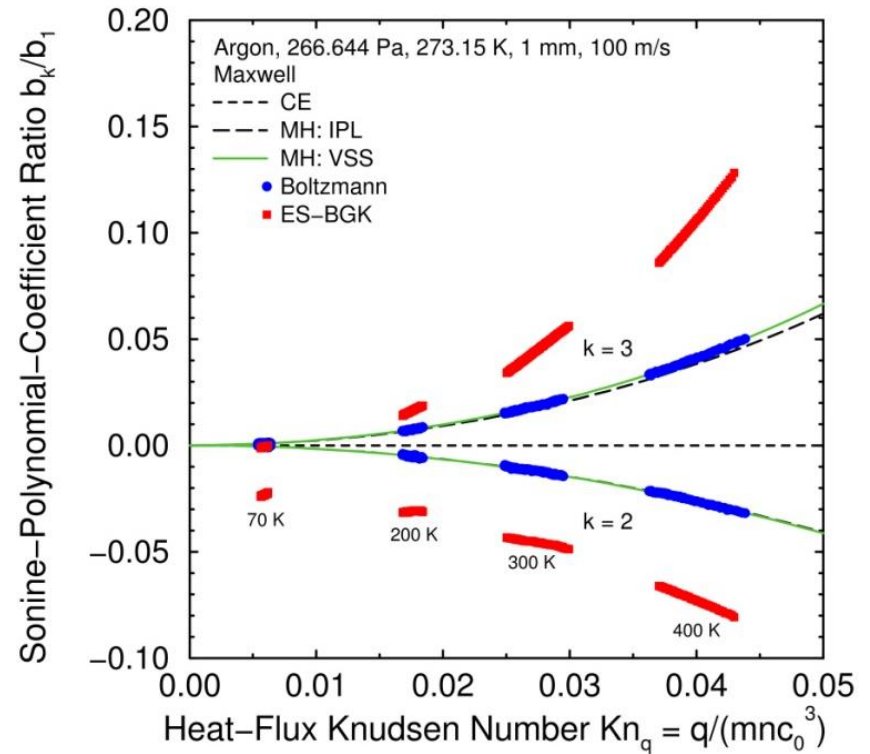
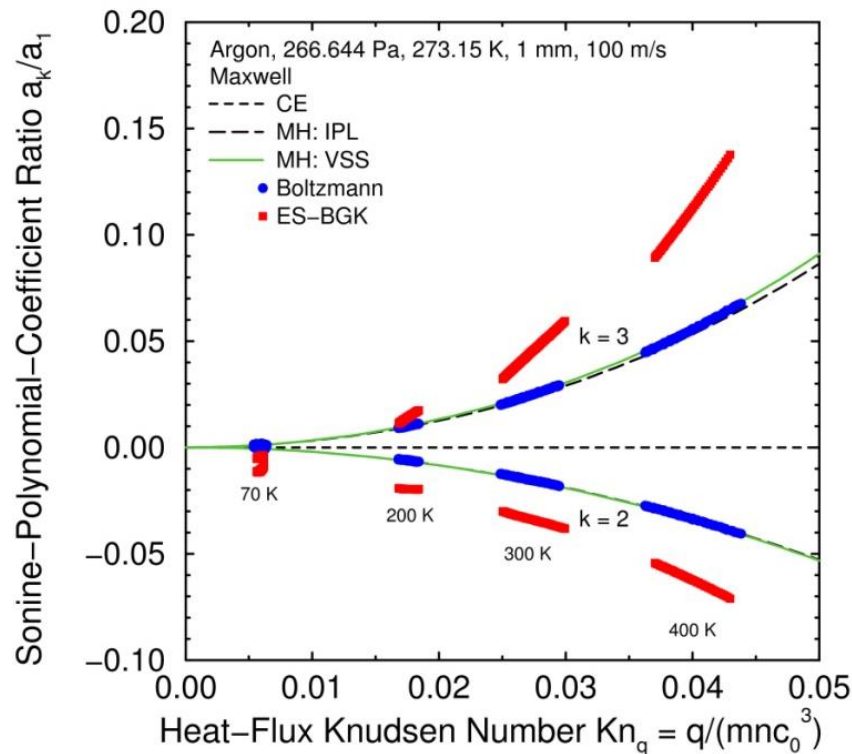


Sonine polynomial coefficients for hard-sphere interaction

- Chapman-Enskog (continuum) values for $k \geq 2$ are **nonzero**
 - Systematic differences in Knudsen layers (transition regime)
- ES-BGK differs strongly from CE values away from walls
 - ES-BGK hard-sphere & Maxwell values hardly differ but should

Variation of ES-BGK distribution shape with molecular interaction is not accurate

Distribution Shape: Maxwell

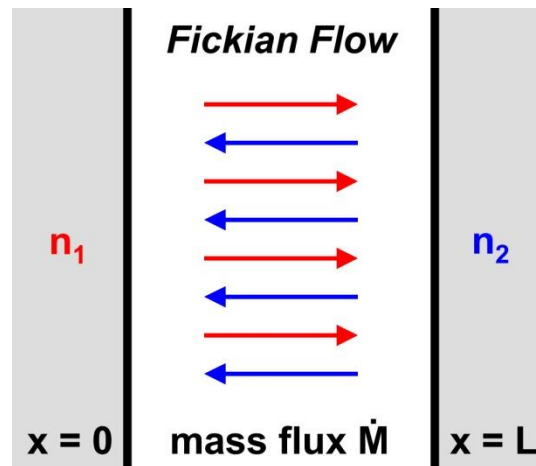


- ES-BGK shape is inaccurate at large flux or stress
- Chapman-Enskog (CE) values for $k \geq 2$ are all zero
 - Appropriate only in limit of zero heat flux & shear stress
- Moment-Hierarchy (MH) values for $k \geq 2$ are nonzero
 - Appropriate for arbitrary nonzero heat flux & shear stress

Fickian Mass Diffusion



Adolf
Fick



$$D = \frac{\mu}{\rho} \frac{D_\infty/D_1}{\mu_\infty/\mu_1} \frac{3(3\nu-5)}{5(\nu-1)} \frac{A_2[\nu]}{A_1[\nu]}$$

$$\nu = 1 + \frac{4}{2\omega-1} \quad n_1 u_1 + n_2 u_2 = 0$$

$$n_1 u_1 = -D \frac{\partial n_1}{\partial x} \quad n_2 u_2 = -D \frac{\partial n_2}{\partial x}$$

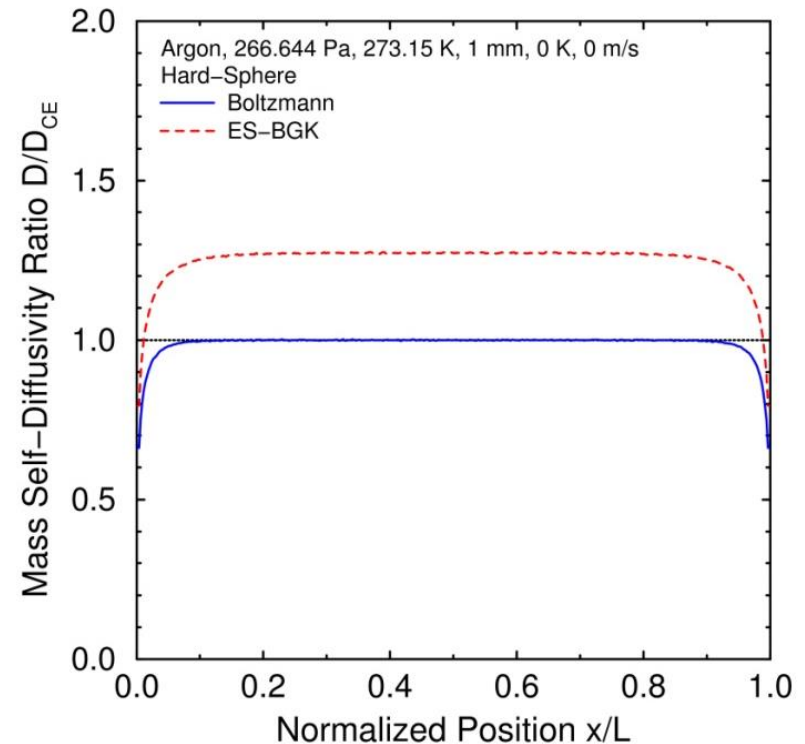
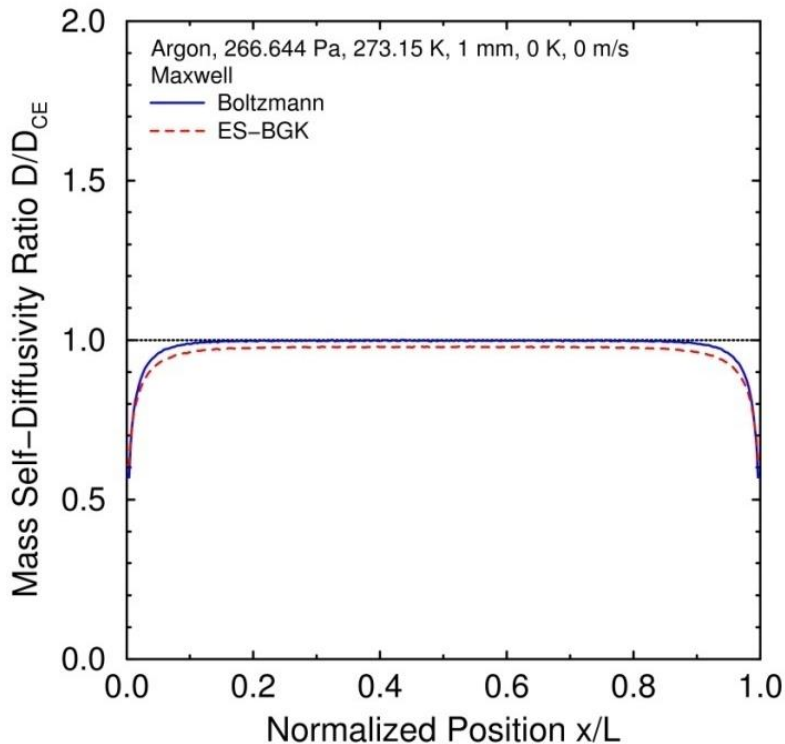
Investigate transport in gas between parallel plates

- Fickian flow: mass diffusion in stationary isothermal gas
 - Identical molecules are passively tagged either “1” or “2”
 - Reflection sets tag to “1” at left wall and “2” at right wall

Compare Boltzmann, ES-BGK results

- Number density, velocity, and mass self-diffusivity profiles
- Maxwell and hard-sphere molecular interactions
- Continuum cases and Chapman-Enskog theory

Mass Self-Diffusivity



ES-BGK is not accurate for mass diffusion

- Normalize simulation values by Chapman-Enskog values
- Boltzmann agrees well with CE values away from walls
- ES-BGK differs from CE values even away from walls
 - Slightly low for Maxwell, quite high for hard-sphere

Evaluating Computational Efficiency

- Computational efficiency is more than CPU time.
 - CPU time to achieve a particular level of accuracy
 - Convergence rate
 - Parallel efficiency (hardly ever addressed)
- Comparing low-order moments (density, temperature) is usually an **insensitive** measure of comparison.



DSMC Numerical Error

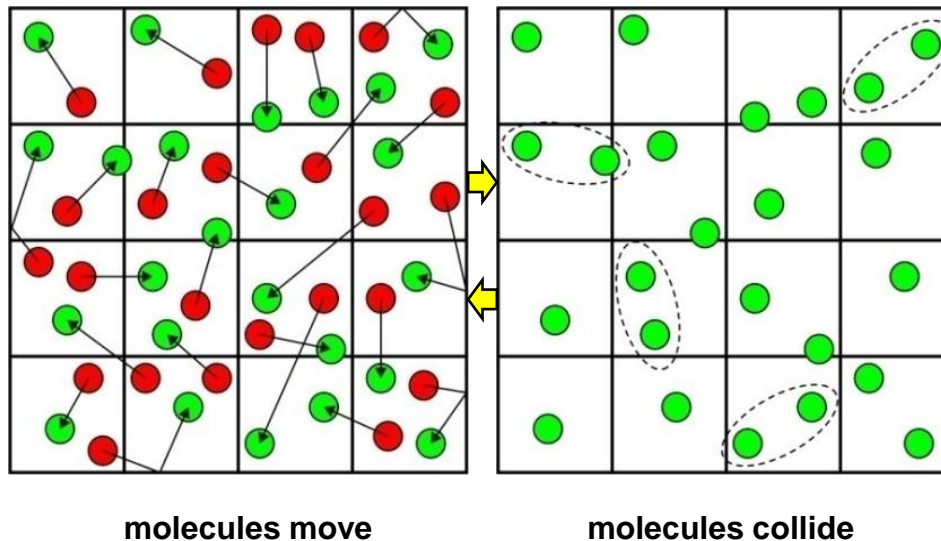
Traditional DSMC rule-of-thumb guidelines:

- Take enough samples to drive statistical error down to “acceptable” level
- Keep time step smaller than $\sim 1/4$ mean collision time
- Keep cell size smaller than $\sim 1/3$ mean free path
- Use a minimum of ~ 20 particles per cell

These guidelines give 2% error, which is similar to the uncertainty in measured transport properties for most gases

- DSMC is subject to the same constraints as other numerical methods.
- DSMC is correct to the limit of vanishing discretization.

DSMC Numerical Error



Four parameters control DSMC error:

Statistical error (1)

Samples per cell (S_c)

Discretization error (3)

- Particles per cell (N_c)
- Cell size (Δx)
- Time step (Δt)

Statistical and Particle-Number Errors

Error related to sample size

- Statistical error
- Cell sample size $S_c = N_c \times N_t$
- N_c = particles per cell; N_t = time steps

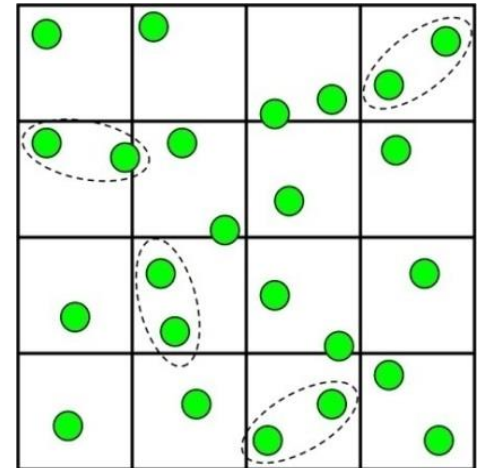
Strategies for overcoming statistical error

- Use large number of samples
- For steady flows, use time and/or ensemble averaging
- Computational expense $\sim S_c$

Error related to local number of particles

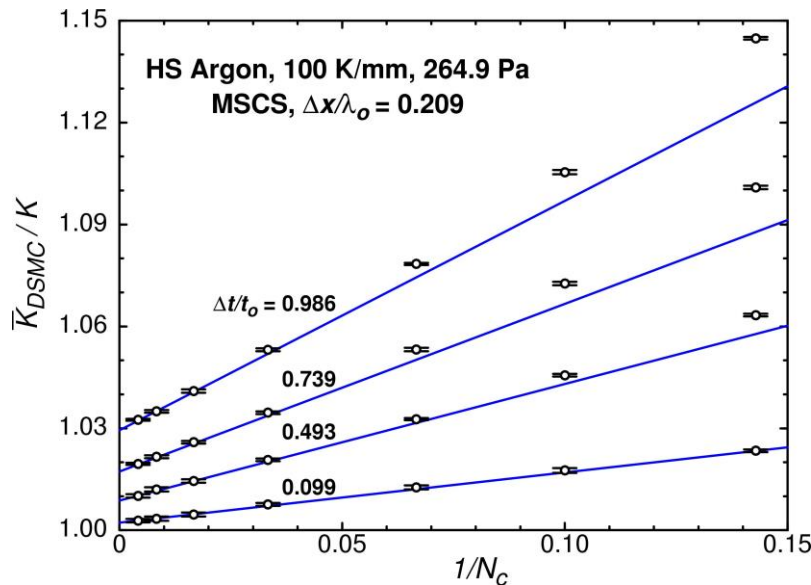
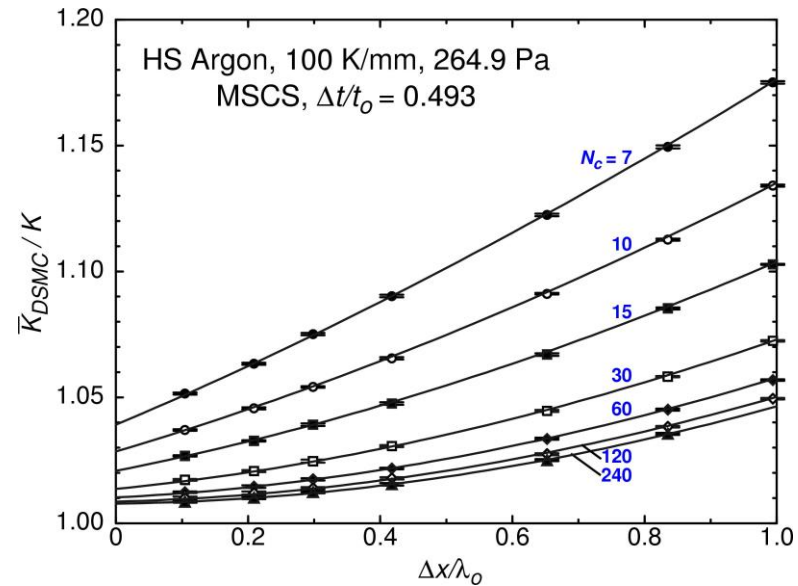
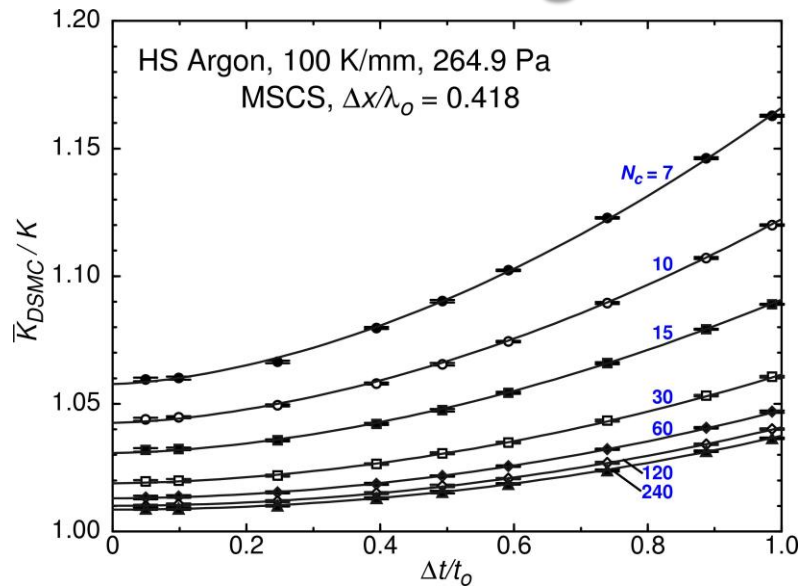
- Error $\propto 1/N_c$
- Systematic – persists even as $S_c \rightarrow \infty$

Limited number of samples per time step



Not enough particles to capture physics

DSMC Convergence



- Curves are best fits
- Error bars represent 95% confidence intervals
- Quadratic convergence for Δx , Δt
- **First-order convergence** $O(1/N_c)$, as $N_c \rightarrow \infty$
- Higher-order for long time steps
- For $N_c = 7$ and $\Delta t/t_o = 0.493$, convergence rate **appears linear** in $\Delta x/\lambda_o$

Functional Form of Error

Functional form that represents DSMC data

- Ad hoc series expansion in Δx , Δt , and $1/N_c$
- Perform least-squares fitting of entire data set

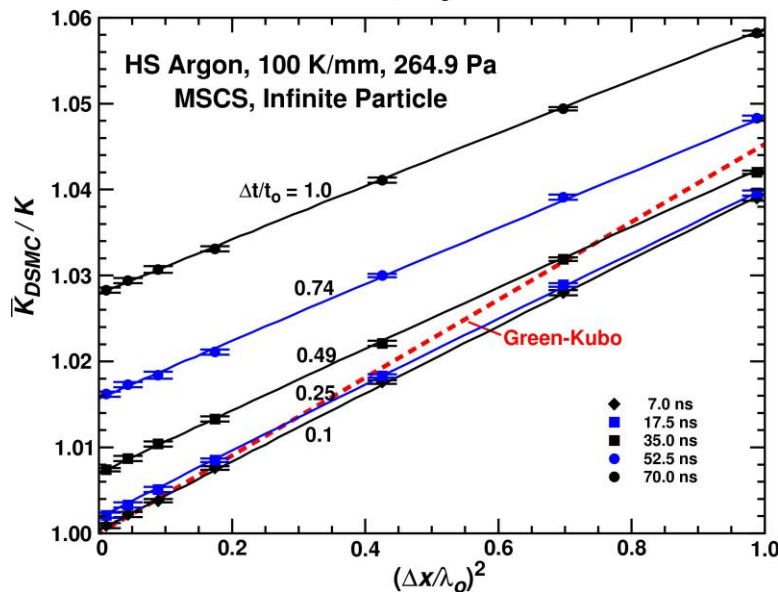
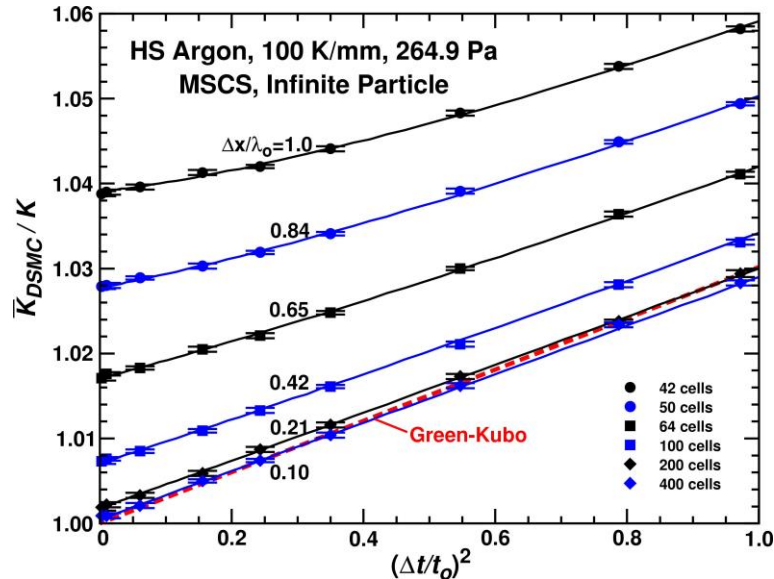
$$\frac{K_{DSMC}}{K} = 1.0000 + 0.0286\Delta\tilde{t}^2 + 0.0411\Delta\tilde{x}^2 - 0.0016\Delta\tilde{x}^3 - 0.023\Delta\tilde{t}^2\Delta\tilde{x}^2 +$$

$$-\frac{0.111}{N_c} + \frac{1}{N_c} \left[1.22\Delta\tilde{x} - 0.26\Delta\tilde{x}^2 + 0.97\Delta\tilde{t}^{3/2} + \dots \right] + 0.95\frac{\Delta\tilde{t}^2}{N_c^2} + \dots$$

Cross terms show convergence behavior is complex

Rader D. J., Gallis M. A., Torczynski J. R., Wagner W., "DSMC Convergence Behavior of the Hard-Sphere-Gas Thermal Conductivity for Fourier Heat Flow", *Phys. Fluids*, 18, 077102, 2006.

Infinite-Particle Convergence



- Finite-particle error removed: values “extrapolated” to $N_c \rightarrow \infty$
- 63 extrapolated data points
- Error bars: fitting uncertainty
- Quadratic convergence in time step and cell size
- Qualitative agreement with *Green-Kubo* theory, but slopes are different
- Lines are best fits of data

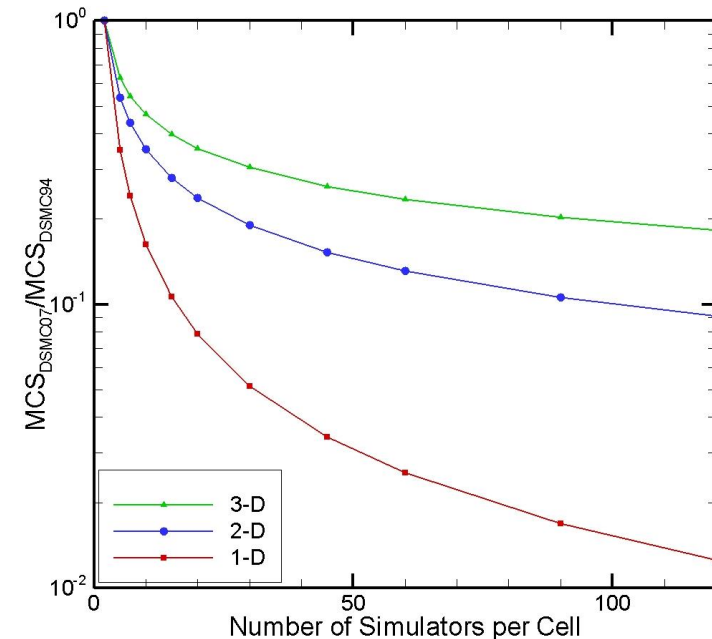
Improving DSMC

Nearest-Neighbor Procedures

- Transient Adaptive Sub-Cells (Bird 2000)
 - **Stochastically** determines the nearest neighbor
 - Creates Cartesian sub-cell structure for each collision phase
 - Number of simulators in each sub-cell is $\sim 1-2$
 - Relatively inexpensive for large N : $O(N)$
- Virtual Sub-Cells (LeBeau 2003)
 - **Deterministically** determines the nearest neighbor
 - More accurate*
 - More expensive than TASC for large N : $O(N^2)$

DSMC07 = Sophisticated DSMC

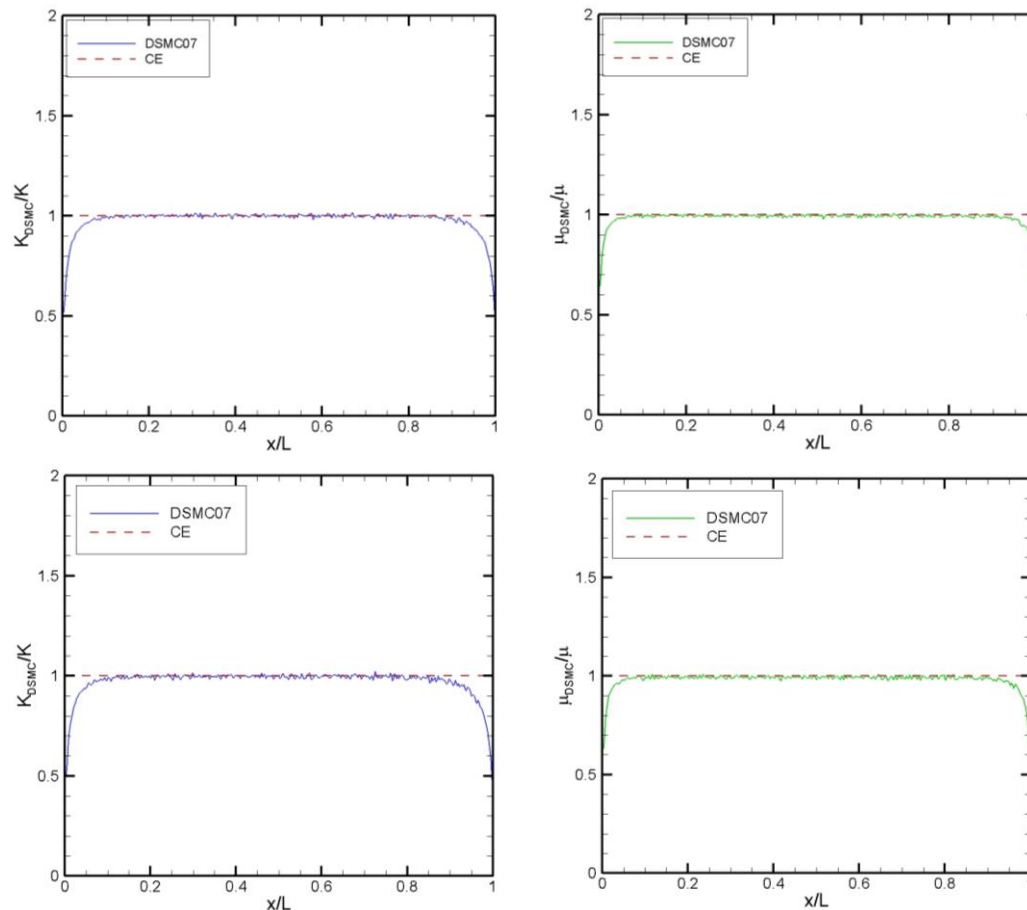
DSMC94 = Established DSMC



*A stress-tensor anisotropy is introduced in reduced-dimensionality simulations

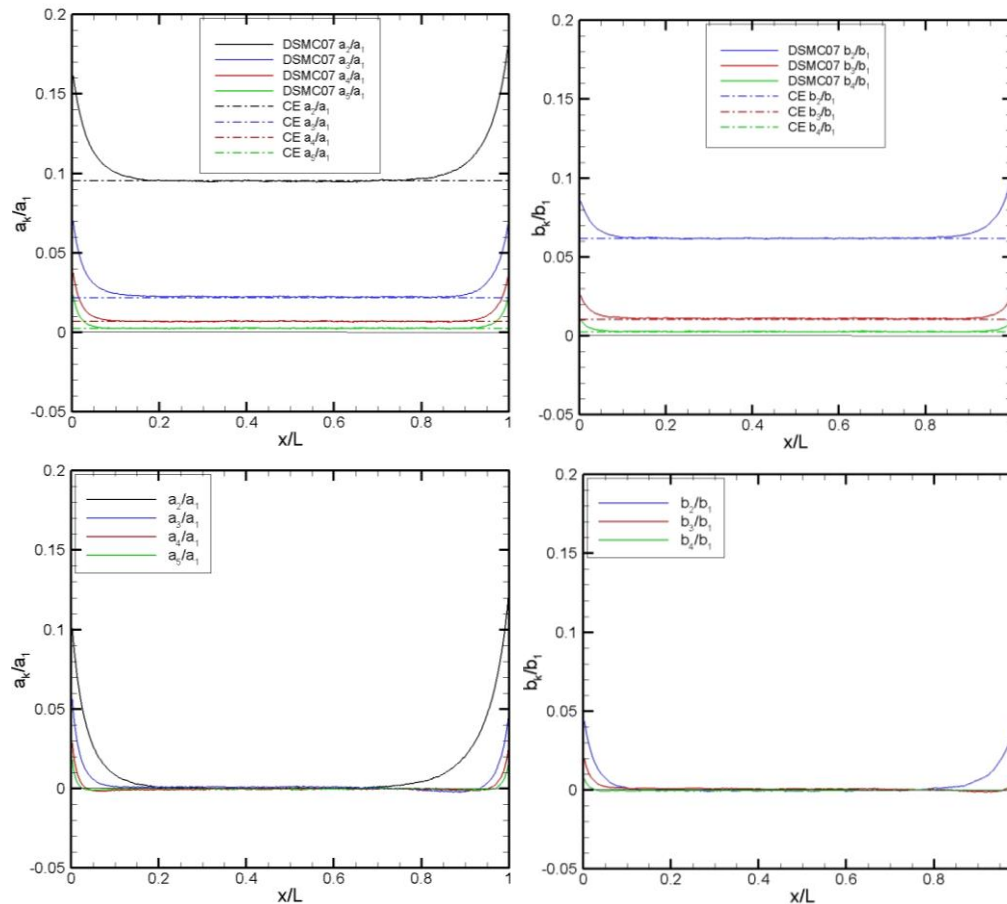
Bird G. A., Gallis M. A., Torczynski J. R., Rader D. J., "Accuracy and efficiency of the sophisticated Direct Simulation Monte Carlo algorithm", *Phys. Fluids*, 21, 017103, 2009.

Hard-Sphere & Maxwell Transport Coefficient Profiles



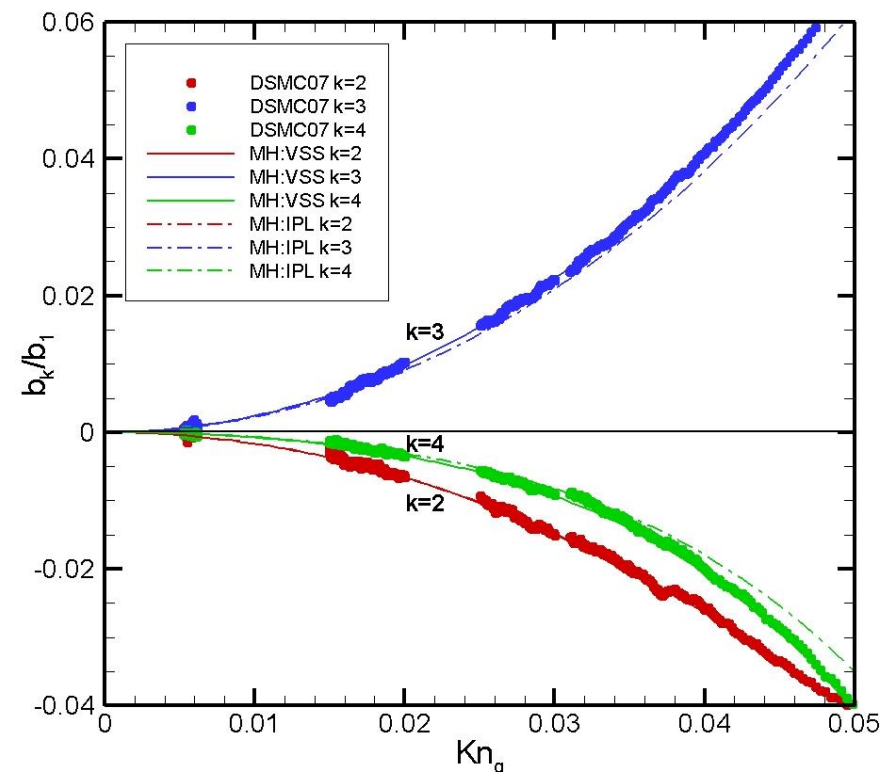
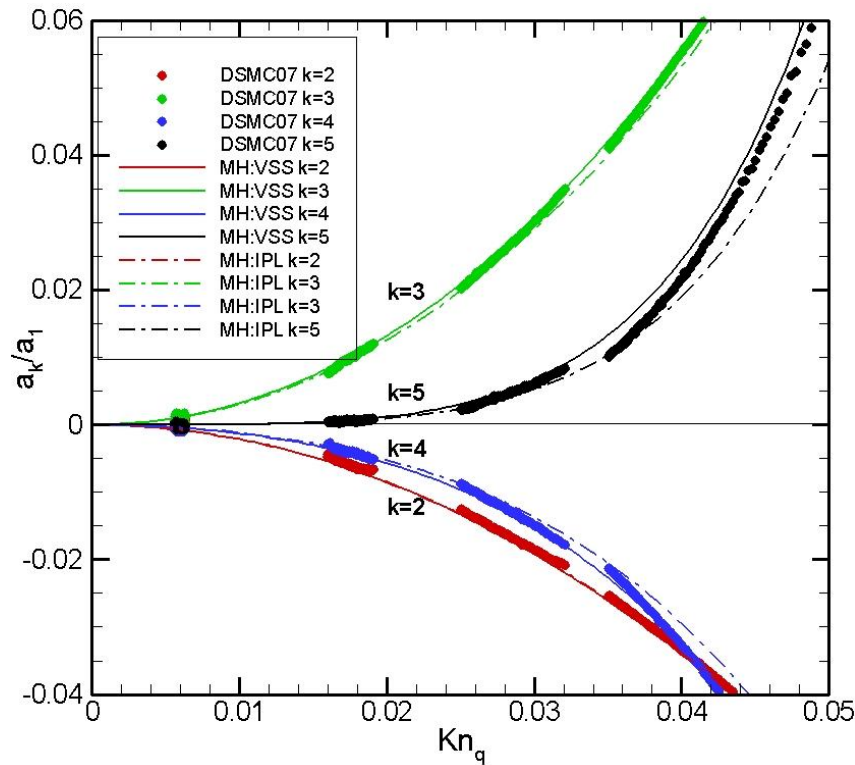
- DSMC07 and CE thermal conductivity and viscosity
 - Low heat flux, low shear stress: $Kn_q = 0.006$, $Kn_\tau = 0.003$
 - Agreement in central region: normal solution

Hard-Sphere & Maxwell Sonine-Coefficient Profiles



- **DSMC07 and CE hard-sphere Sonine polynomial coefficients a_k/a_1 and b_k/b_1**
 - Good agreement in central region: normal solution
 - Demonstrates accuracy of molecular velocity distribution

Maxwell Normal Sonine Coefficients



- DSMC07 and MH Maxwell normal solutions for a_k/a_1 and b_k/b_1
 - Four DSMC07 simulations: $\Delta T = 70, 200, 300, 400$ K
 - MH: VSS-Maxwell (solid) and IPL-Maxwell (dashed) differ
 - DSMC07 and MH VSS-Maxwell normal solutions agree

DSMC-VSC Functional Form of Error

Best-fit correlation function for sophisticated DSMC with $\Delta t = \Delta t_o$

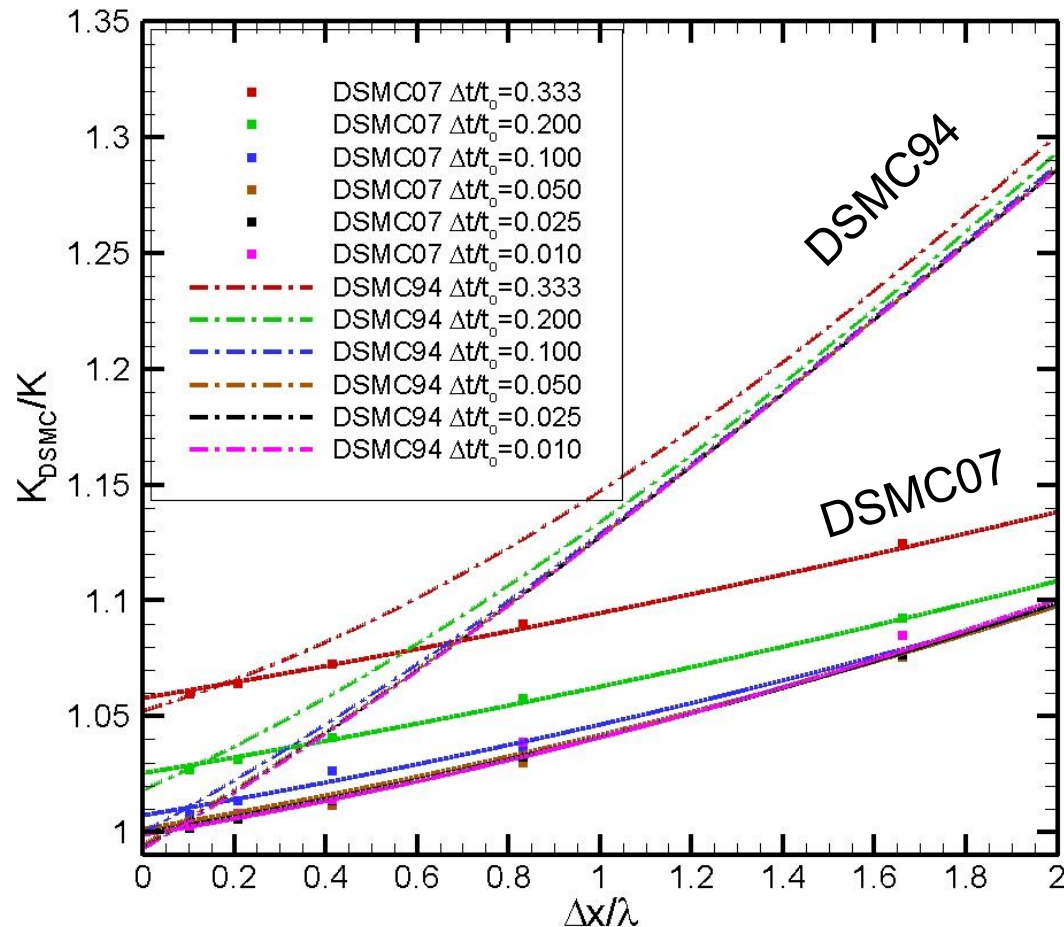
$$\frac{K_{DSMC}}{K} = 0.9953 + 0.07773 \left(\frac{2 \cdot Dt}{t_o} \right) + 0.02795 \left(\frac{2 \cdot Dt}{t_o} \right)^2 - 0.0066 \left(\frac{Dx}{l} \right)^2 \\ - 0.0234 \left(\frac{2 \cdot Dt}{t_o} \right)^2 \left(\frac{Dx}{l} \right)^2 + 0.60375 \frac{1}{N_c} + \frac{1}{N_c} F \left[\frac{2 \cdot Dt}{t_o}, \frac{Dx}{l}, \left(\frac{2 \cdot Dt}{t_o} \right)^2 \right]$$

DSMC limiting convergence differs from Green-Kubo (GK) behavior

- **Weak quadratic** convergence in cell size ($\Delta t/t_o \rightarrow 0, N_c \rightarrow \infty$)
- Linear convergence in $1/N_c$ for $N_c \geq 30$ simulators/cell
- **Linear convergence in time step** ($\Delta x/\lambda \rightarrow 0, N_c \rightarrow \infty$)

Convergence Behavior for $N_c = 10$

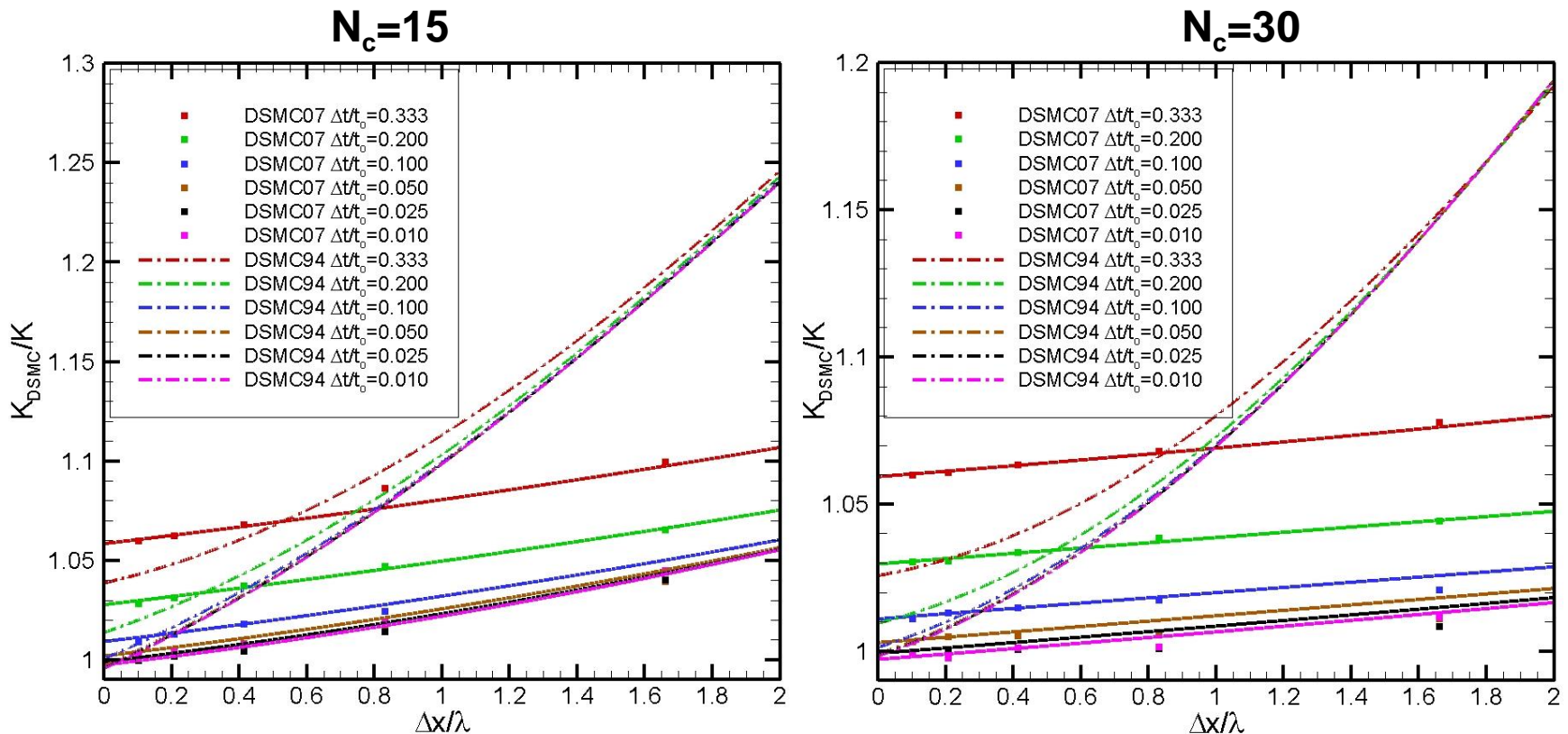
Effect of Cell Size



- For finite number of simulators, the algorithm
 - Is insensitive to spatial resolution
 - Has error below 2% for all cases as long as $\Delta t/t_0 \leq 0.2$

Convergence Behavior for $N_c = 15$ & 30

Effect of Cell Size



For a finite number of simulators, the algorithm

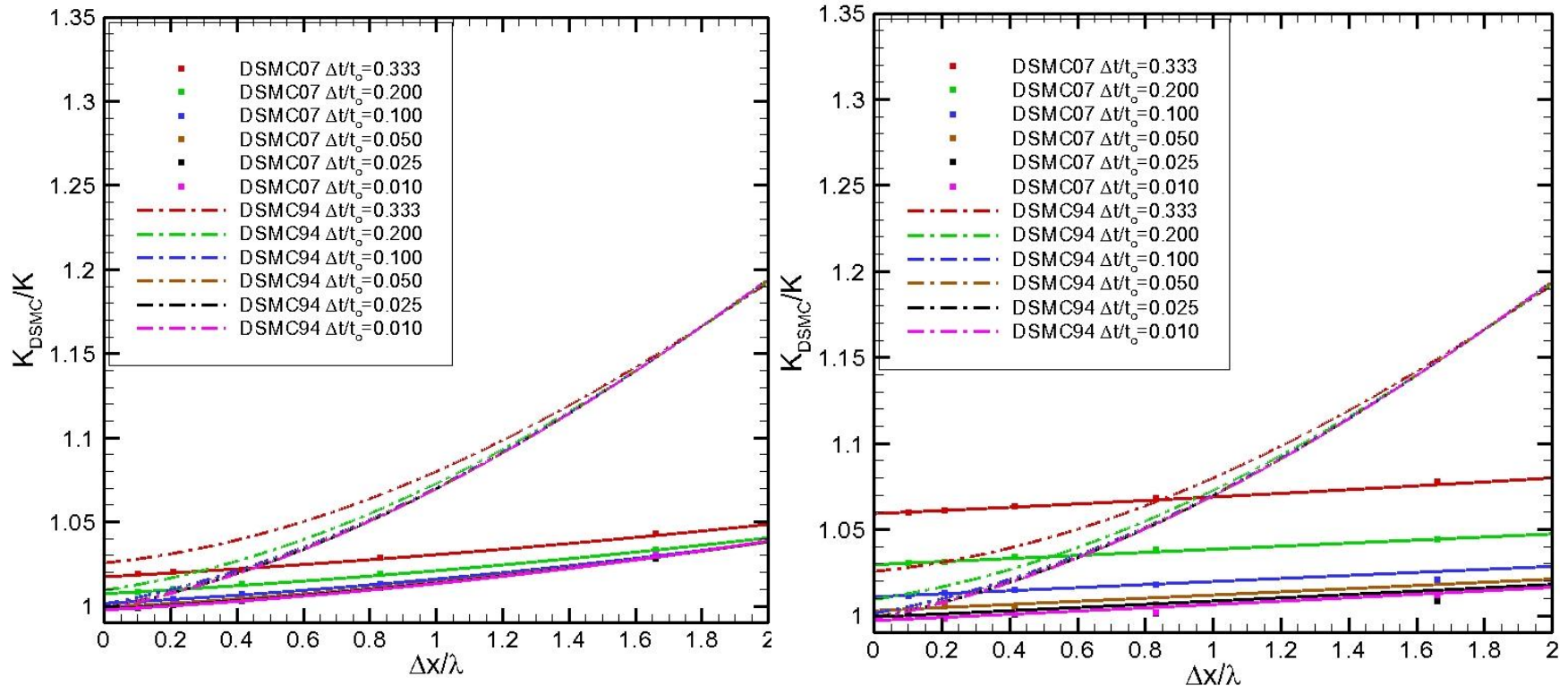
- Is insensitive to spatial resolution.

DSMC07 error is lower than DSMC94 **only** when

- $\Delta t/t_0 \leq 0.1$ ($N_c=15$), $\Delta t/t_0 \leq 0.05$ ($N_c=30$).

Convergence Behavior for $N_c = 30$

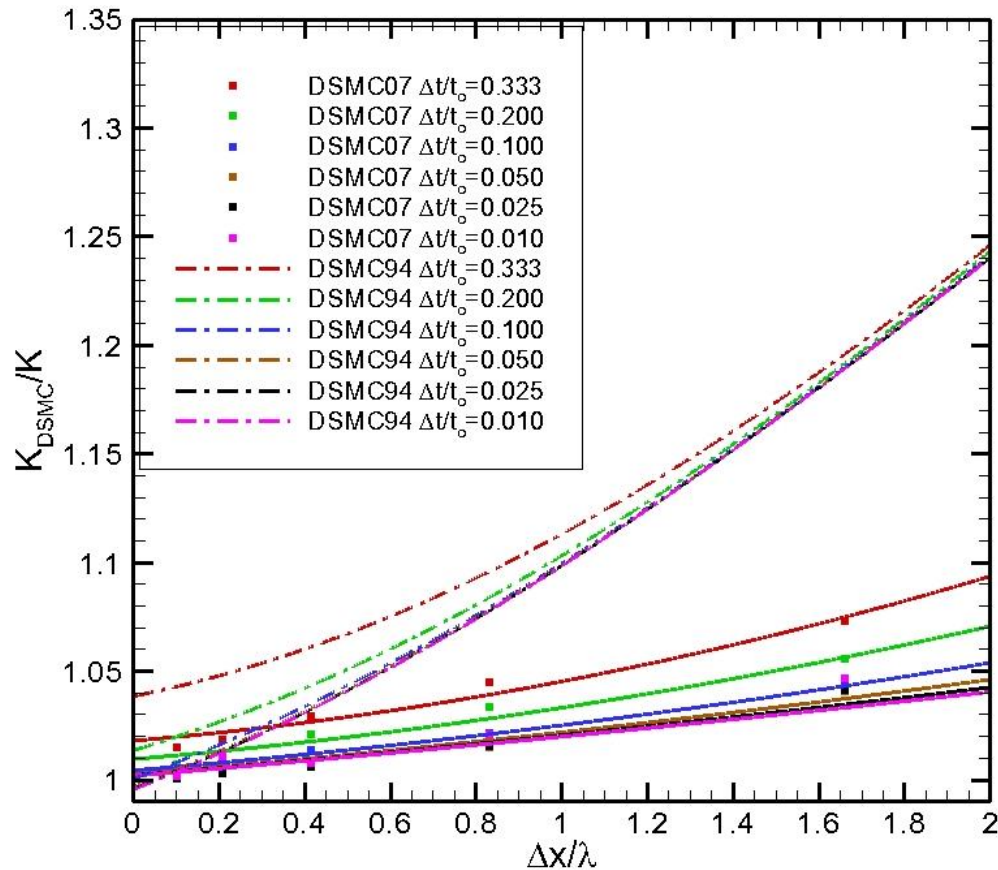
Selection Limit 10 & No Limit



- Significant performance improvement with smaller computational cost can be achieved when a random subset of simulators is used.
- A **limit of 30** should be adequate for most applications, as suggested by Bird.

Convergence Behavior for $N_c = 15$

Trajectory dependent selection



Simulators with high velocities have more simulators available as collision partners, whereas simulators with low velocities have to collide with their nearest neighbor.

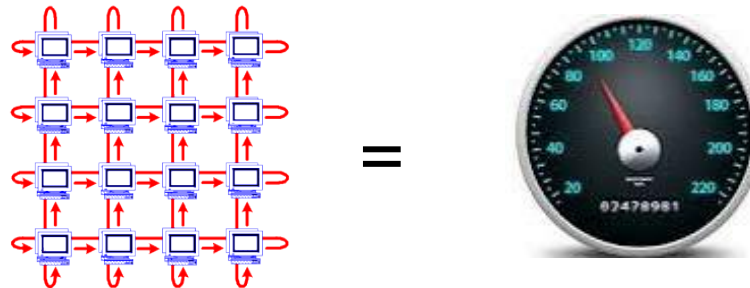
To minimize the time-step error in the VSC scheme, a **near-neighbor** collision partner is selected from within a sphere centered on the simulator with a radius proportional to the **distance traveled** by the simulator during the latest advection phase.

Gallis, M. A. Torczynski J. R., "Effect of collision-partner selection schemes on the accuracy and efficiency of the direct simulation Monte Carlo method", International Journal for Numerical Methods in Fluids, Vol. 67, No. 8, pp.1057-1072, 2011.

Parallel Efficiency: The Unfair Advantage

- The advantages of DSMC come at a cost
- DSMC is **computationally efficient** but **computationally intense**
- Its successful application to real problems depends heavily on its parallel performance
- **1000x speedup** required for some problems of interest
- Monte Carlo methods usually have good parallel performance
 - The workload depends mainly on the simulators within a cell
 - Relatively less need to communicate information between cells
 - Trivial to parallelize in velocity space

The necessary speedup can be achieved without any loss of accuracy or convergence characteristics through parallel computing



Top 5 Supercomputers (2014)

Rank	Site	System	Cores	Rmax (TFlop/s)	Rpeak (TFlop/s)
1	National Super Computer Center in Guangzhou	Tianhe-2 (MilkyWay-2) - TH-IVB-FEP Cluster, Intel Xeon E5-2692 12C 2.200GHz, TH Express-2, Intel Xeon Phi 31S1P	3,120,000	33,862.7	54,902.4
2	DOE/SC/Oak Ridge National Laboratory	Titan - Cray XK7 , Opteron 6274 16C 2.200GHz, Cray Gemini interconnect, NVIDIA K20x	560,640	17,590.0	27,112.5
3	DOE/NNSA/LLNL	Sequoia - BlueGene/Q, Power BQC 16C 1.60 GHz, Custom	1,572,864	17,173.2	20,132.7
4	RIKEN Advanced Institute for Computational Science (AICS)	K computer , SPARC64 VIIIfx 2.0GHz, Tofu interconnect	705,024	10,510.0	11,280.4
5	DOE/SC/Argonne National Laboratory	Mira - BlueGene/Q, Power BQC 16C 1.60GHz, Custom	786,432	8,586.6	10,066.3

Programming for Next Generation and Exascale Machines

- Millions of nodes likely
- Reduced memory per node
- Parallelism within node:
 - Multi-core: 16 and growing
 - Many-core: Intel Xeon Phi, 240 threads
 - GPUs: NVIDIA/AMD, 1000 warps
- Example:
 - LLNL BG/Q: 96K nodes, 16 cores/node + 4 MPI tasks/core

Programming model: MPI + X

- Goal is to decouple the science code from the hardware details

Necessary elements

- Adaptive gridding
- In-situ visualization
- Efficient communications
- Load balancing



Aiming for MPI+X via Kokkos

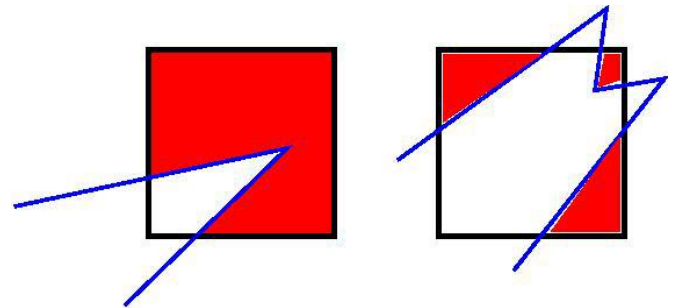
- What is Kokkos:
 - Programming model in development at Sandia
 - C++ template library
 - Open-source
 - Stand-alone
- Goal: write application kernels only once, and run them efficiently on a wide variety of hardware platforms
- Two major components:
 - Data access abstraction via Kokkos arrays optimal layout & access pattern for each device: GPU, Xeon Phi, etc.
 - Parallel dispatch of small chunks of work auto-mapped onto back-end languages: CUDA, OpenMP, etc.

Developing an Exascale DSMC Code

SPARTA = Stochastic PArallel Rarefied-gas Time-accurate Analyzer

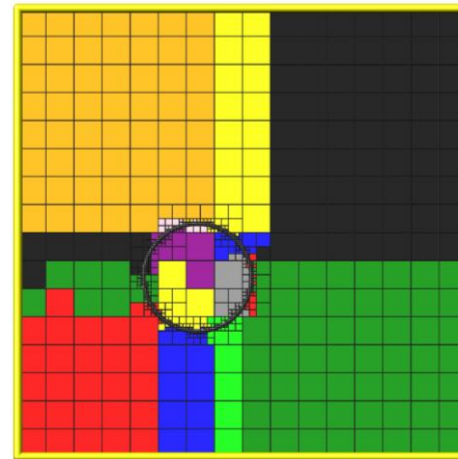
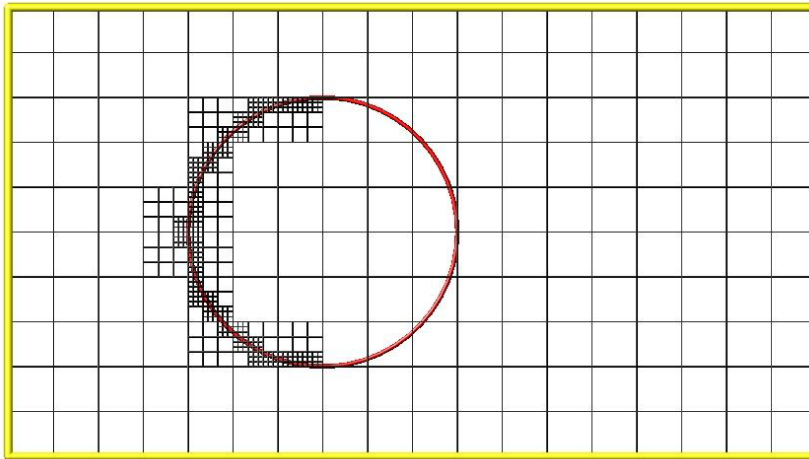
General features

- 2D or 3D, serial or parallel
- Cartesian, hierarchical grid
 - Oct-tree (up to 16 levels in 64-bit cell ID)
 - Multilevel, general NxMxL instead of 2x2x2
- Triangulated surfaces cut/split the grid cells
 - 3D via Schwartzentruber algorithm
 - 2D via Weiler/Atherton algorithm
 - Formulated so can use as kernel in 3D algorithm
- C++, but really object-oriented C
 - Designed to be easy to extend
 - New collision/chemistry models, boundary conditions, etc.



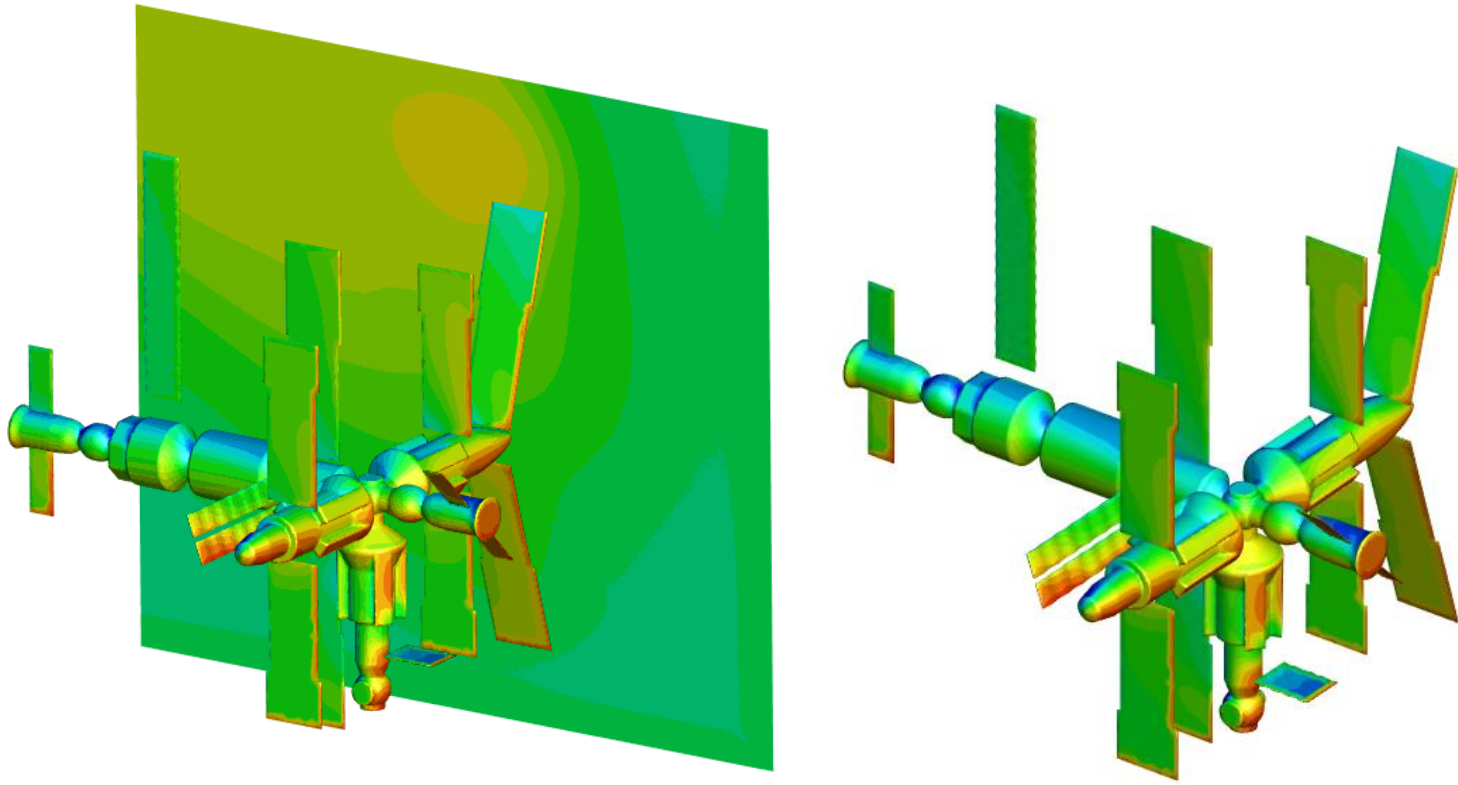
Adaptive Gridding

- Create/adapt grid in situ, rather than pre-process & read in
- Examples: Generate around surface to user-specified resolution, adapt grid based on flow properties
- Algorithms should be efficient if they require only local communications



- Another setup task: label cells as outside/inside
- Simple if pre-processing, in situ easier for large problems

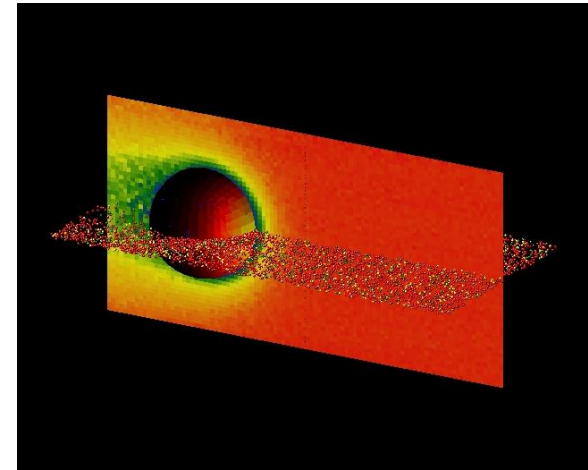
Simulation of Complicated Shapes



Grid generation (10^7 cells) completed in 0.3 seconds on 16 processors
Geometry comprises multiple “water-tight” bodies

In-Situ Visualization

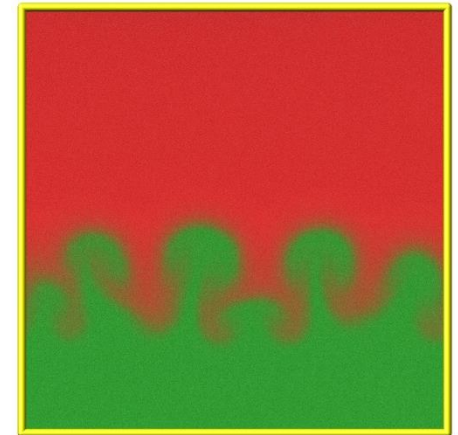
Not a replacement for interactive viz, but ...
Quite useful for **debugging** & quick analysis
At end of simulation (or during), instant movie



Render a JPG snapshot every N time steps:

- Each processor starts with blank image (1024x1024)
- Processor draws its cells/surfaces/molecules with depth-per-pixel
- Merge pairs of images, keep the pixel in front, recurse
- Draw is parallel, merge is logarithmic (like MPI Allreduce)

Images are ray-traced quality



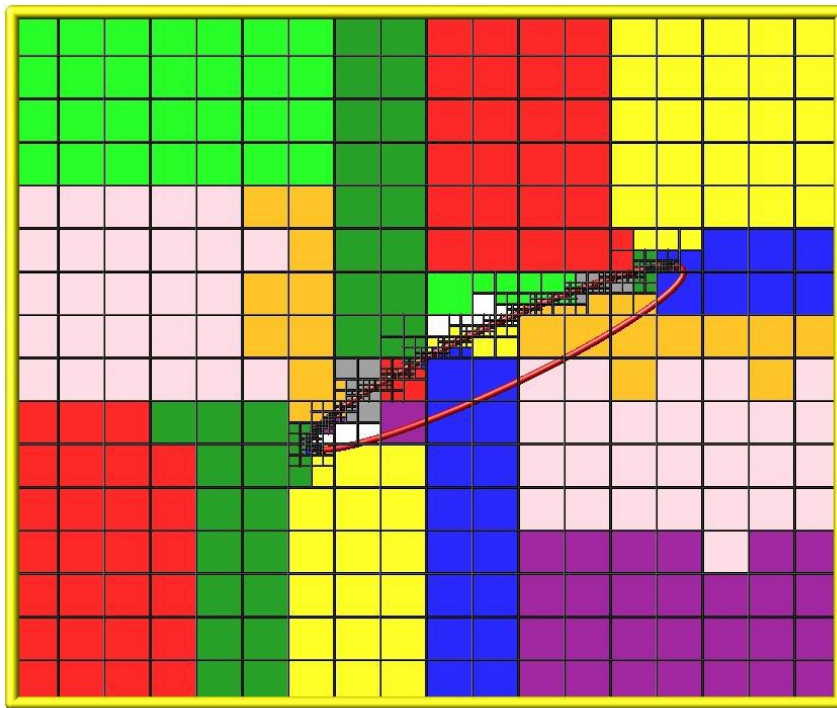
Load Balancing

Balance across processors, **static or dynamic**

Granularity = grid cell with its molecules

Geometric method: recursive coordinate bisection (RCB)

Weighted by cell count or molecules or CPU



RCB is fast

Bigger cost is **data move**

Example:

1B cells on 1024 BG/Q nodes

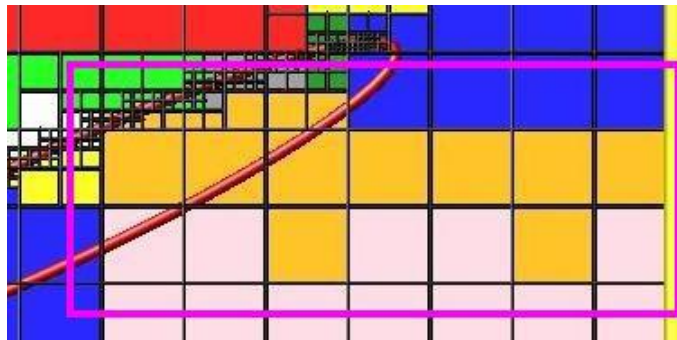
Worst case: move all cells

Balance time = 15 s:

(RCB=2, move=12, ghosts=1)

Efficient Communication

- One processor = compact clump of cells via load balancing
 - Ghost region = nearby cells within **user-defined cutoff**
 - Store surface information for ghost cells to complete move



- Efficiently distributes grid information across processors
 - With sufficient cutoff, only **one communication per step**
 - Multiple passes if needed (or can bound molecule move)
- Communication with **modest count of neighbor processors**

SPARTA Benchmarking

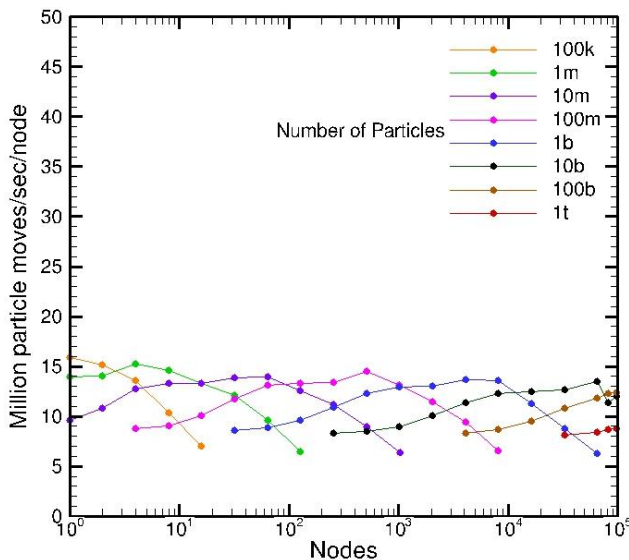
2 test cases:

- **Free-molecular**
 - Stress test for communication
 - 3D regular grid, 10^4 - 10^{11} (**0.1 trillion**) grid cells
 - 10 molecules/cell, 10^5 - 10^{12} (**1 trillion**) molecules
- **Collisional**
 - About 2x slower (sorting, collisions)
 - Same grid cell & molecule counts
- Effect of threading
 - **4 threads/core = 2x speed**

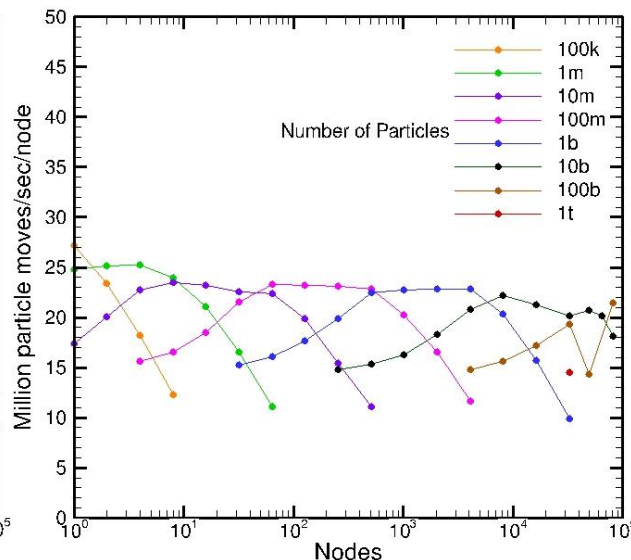


SPARTA Benchmarking

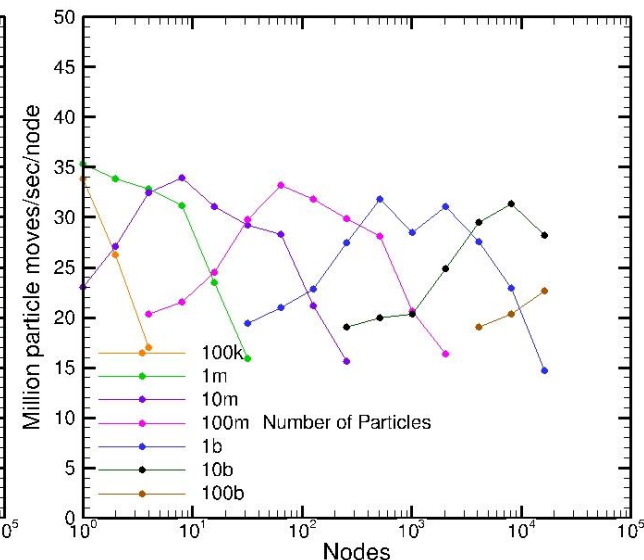
16 cores/node
1 task/core



16 cores/node
2 tasks/core



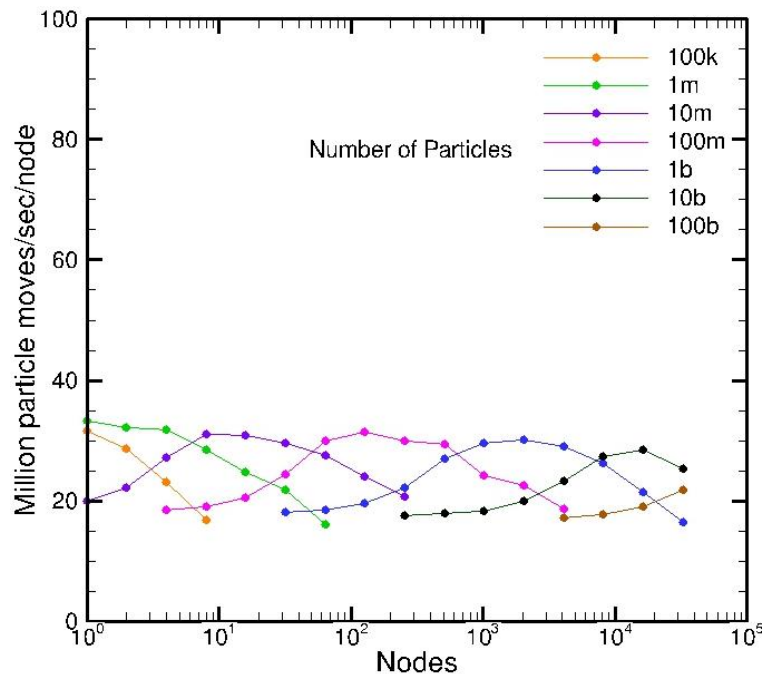
16 cores/node
4 tasks/core



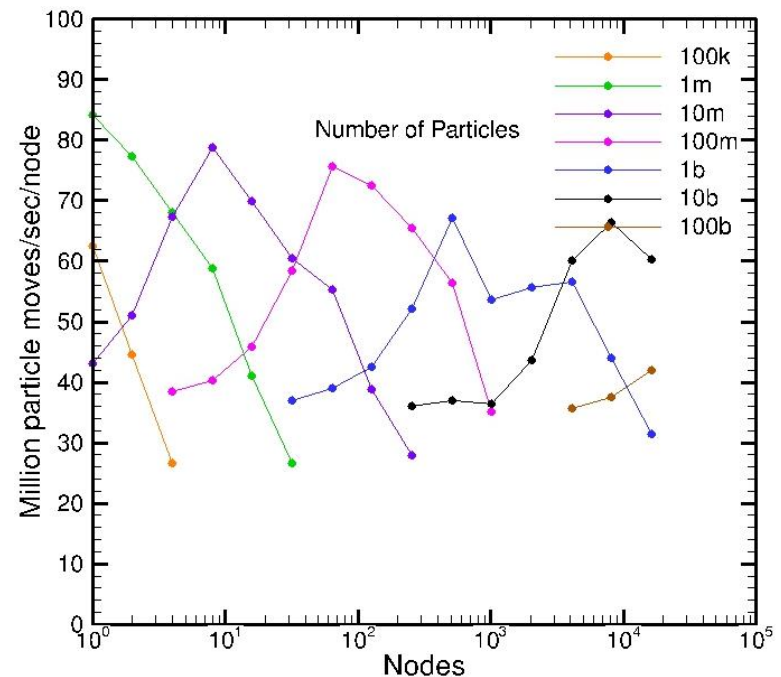
- Weak scaling indicates, 10% peak performance reduction from 1 to 10^6 cores
- 2 tasks/core gives 1.5x speedup, 4 tasks/core gives 2x speedup
- A total of **1 trillion simulators** can be simulated on **one third** of the BG/Q
- Maximum number of tasks is 2.6 million

SPARTA Benchmarking (FM)

16 cores/node, 1 task/core

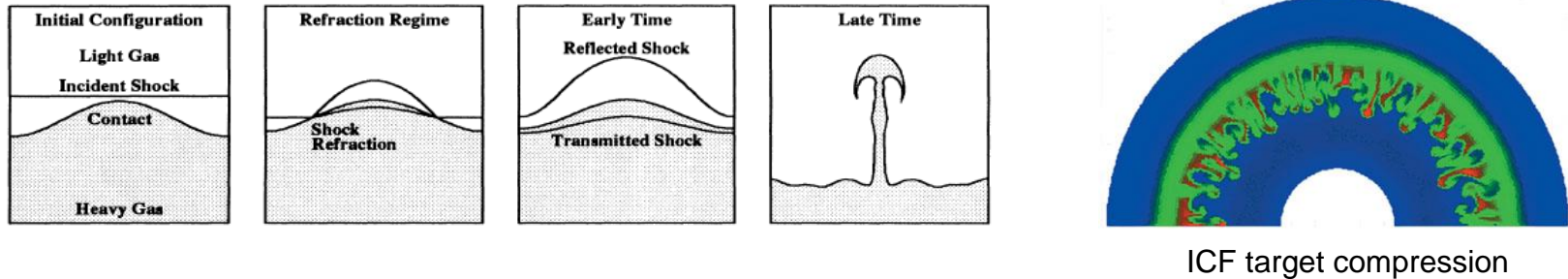


16 cores/node, 4 tasks/core



- Free-molecular (FM) calculations stress communications
- 2x speedup compared to collisional

Richtmyer-Meshkov Instability (RMI)

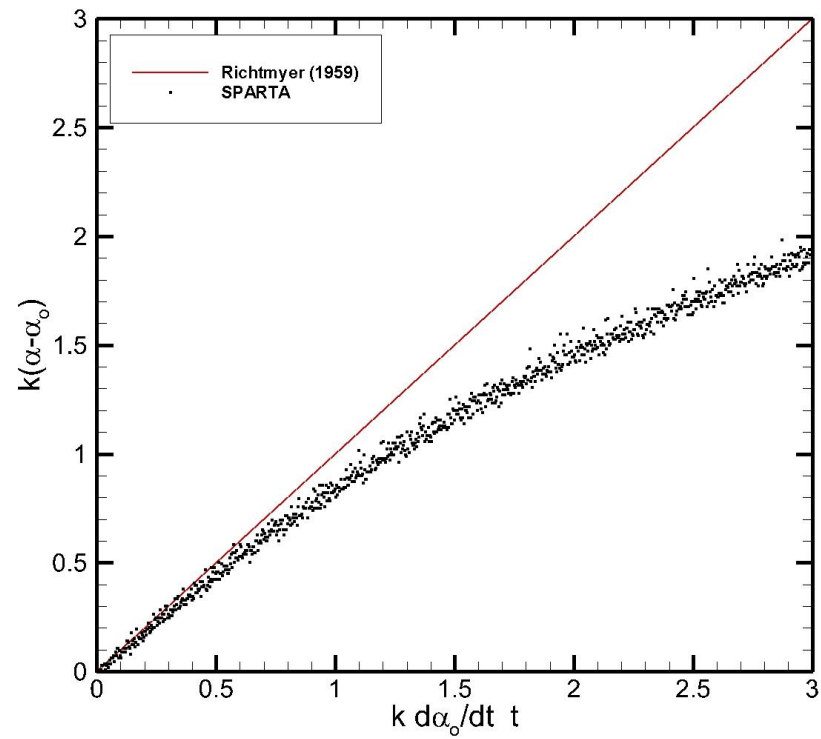


Applications include Inertial Confinement Fusion (ICF), stellar evolution models, interaction of shocks with flames

RMI combines multiple compressible phenomena

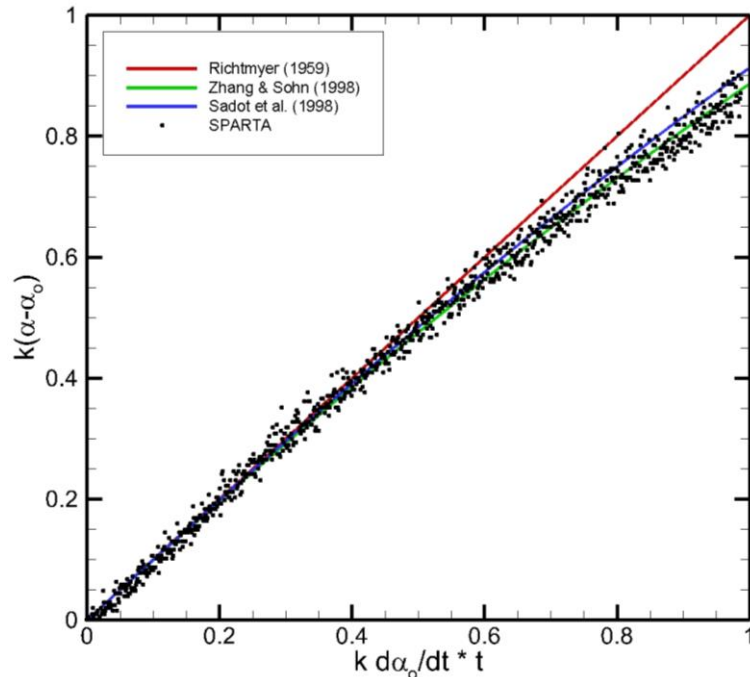
- Shock interaction, refraction, reflection, transmission
- Hydrodynamic instability, including:
- Nonlinear growth
- Subsequent transition to turbulence
- Range of Mach numbers
- Chemical reactions (combustion)

RMI in He/Ar Mixture: Mach 1.2 Shock

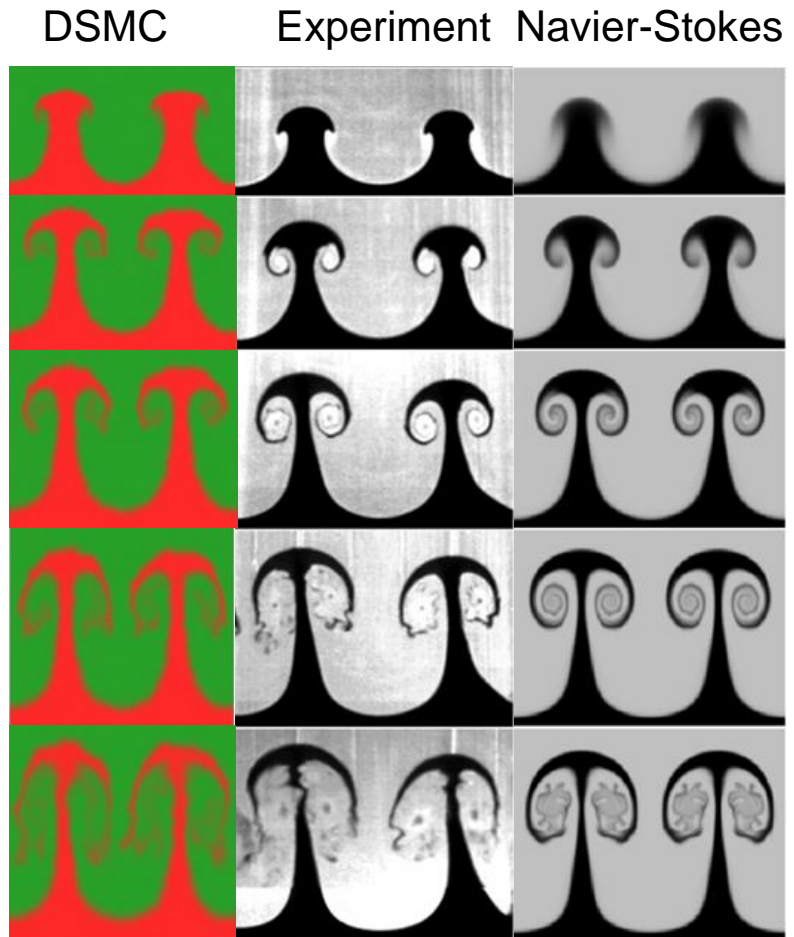


Non-dimensional amplitude for an initially small amplitude perturbation compared to Richtmyer's model for early time evolution

RMI in Air-SF₆ Mixture: Mach = 1.2 Shock



Non-dimensional amplitude for an initially small amplitude perturbation compared to theoretical/empirical models



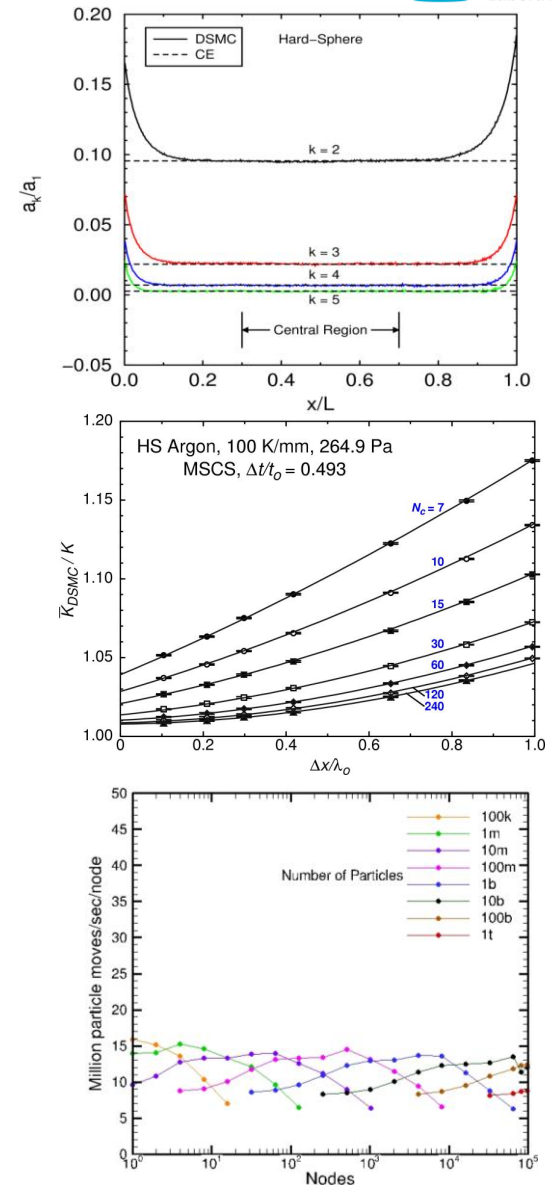
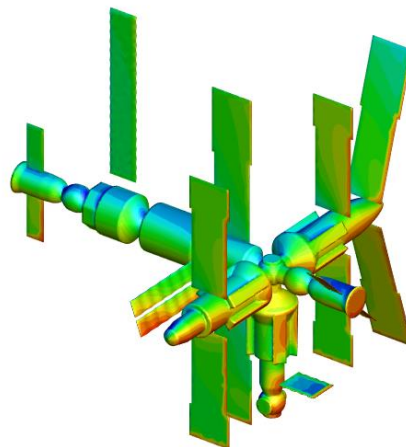
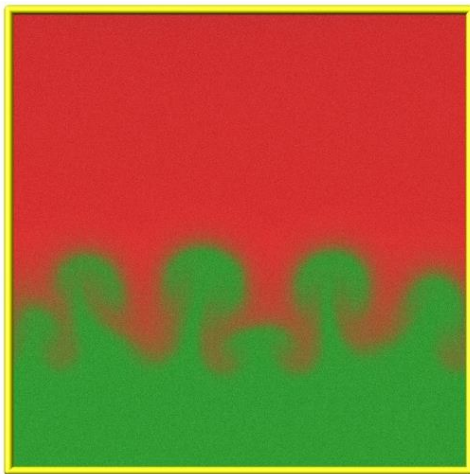
Conclusions

DSMC yields exquisite agreement with analytical results, where available

- Chapman-Enskog, Moment-Hierarchy theory
- Discretization & sampling errors understood

DSMC scales extremely well & can take full advantage of massively parallel platforms

- Can simulate unprecedented flow regimes
- Hydrodynamic instabilities, lower altitudes



Thank you!



Graeme Bird



DSMC15 Workshop, Hawaii, September 2015

*Be there.
Aloha!*

