

Exceptional service in the national interest



Hugoniot Measurements on Sylgard and HNAB via Ultrafast Time Domain Interferometry

Brook Jilek

15th International Detonation Symposium
July 13-18, 2014 San Francisco, CA



This work was performed under the auspices of the U.S. Department of Energy jointly by Lawrence Livermore National Laboratory under Contract DE-AC52-07NA27344.

Sandia National Laboratories is a multi-program laboratory managed and operated by Sandia Corporation, a wholly owned subsidiary of Lockheed Martin Corporation, for the U.S. Department of Energy's National Nuclear Security Administration under contract DE-AC04-94AL85000. SAND NO. 2011-XXXXP

Contributors to this work

From Lawrence Livermore Laboratory:

Mike R. Armstrong, Jonathan C. Crowhurst, James Lewicki, and Joseph M. Zaug.

From United States Naval Academy: Kayla M. Coleman

From Sandia National Laboratories:

Brook A. Jilek, Ian T. Kohl, Darcie A. Farrow, Junji Urayama, Robert Knepper, Gregg Radtke, Sean P. Kearney

Acknowledgements

From Sandia National Laboratories:

Michael Marquez and Alex Tappan for sample preparation.

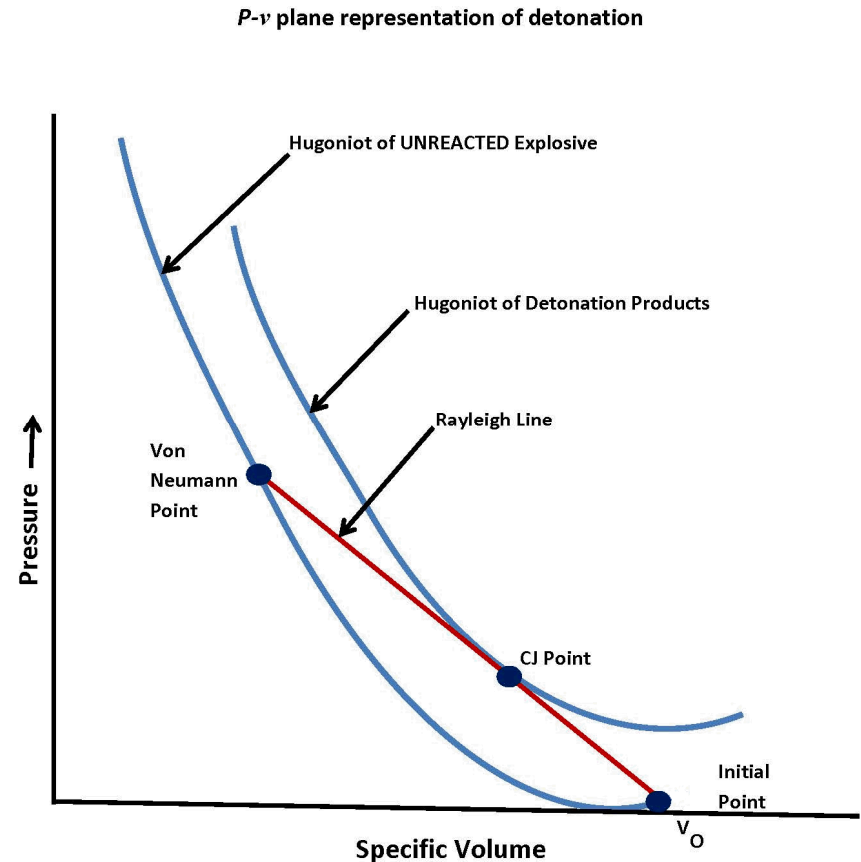
From Los Alamos National Laboratory:

Cindy Bolme, Shawn McGrane, and Katie Brown for lots of useful discussions about ultrafast shock diagnostics


Motivation and Goals

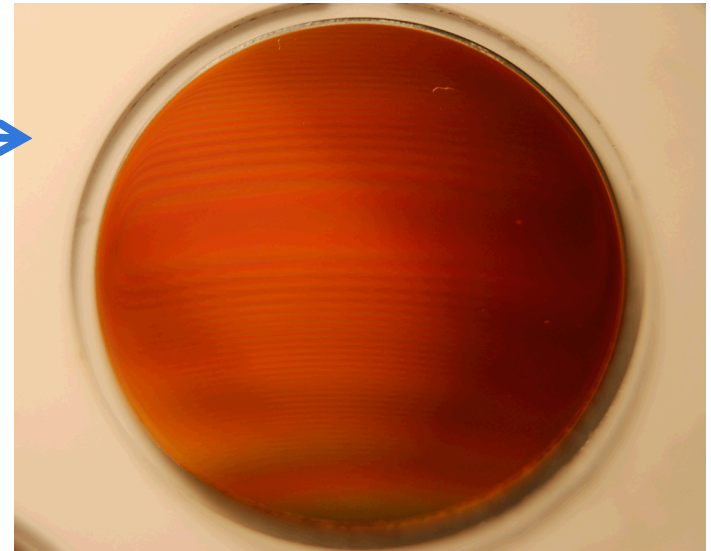
Experimentally-verified Equations of State (EOS) for energetic materials are typically limited to pressures well below detonation pressure.

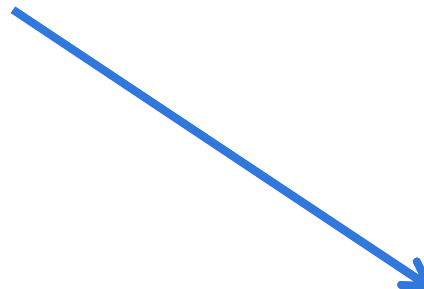
We need sub-ns time resolution to measure shock and particle velocities of energetic materials before they react.



Why HNAB?

HNAB can be vapor-deposited as a fully-dense amorphous thin film. 

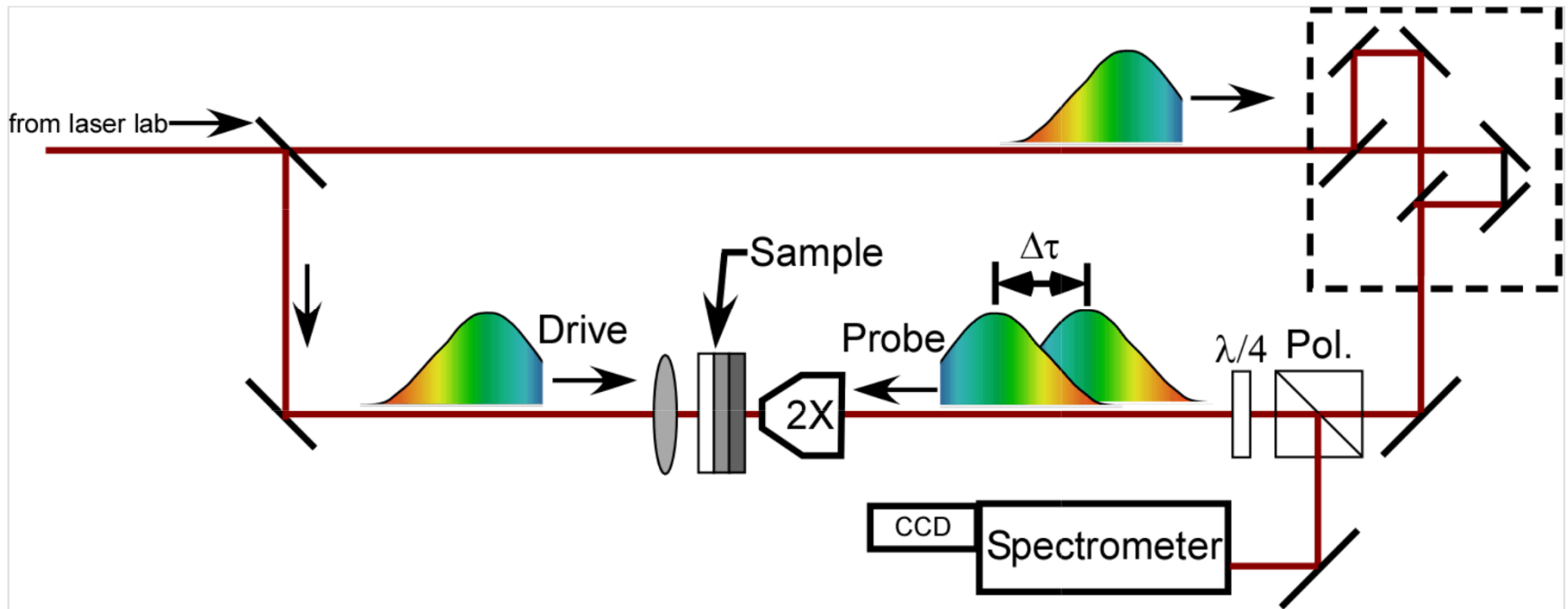


Over a matter of weeks, it crystallizes into several different forms. 



The amorphous and crystallized films are chemically identical and have almost identical densities. Microstructure is the major difference, **making it a model system for studying the effects of microstructure on performance and sensitivity.**

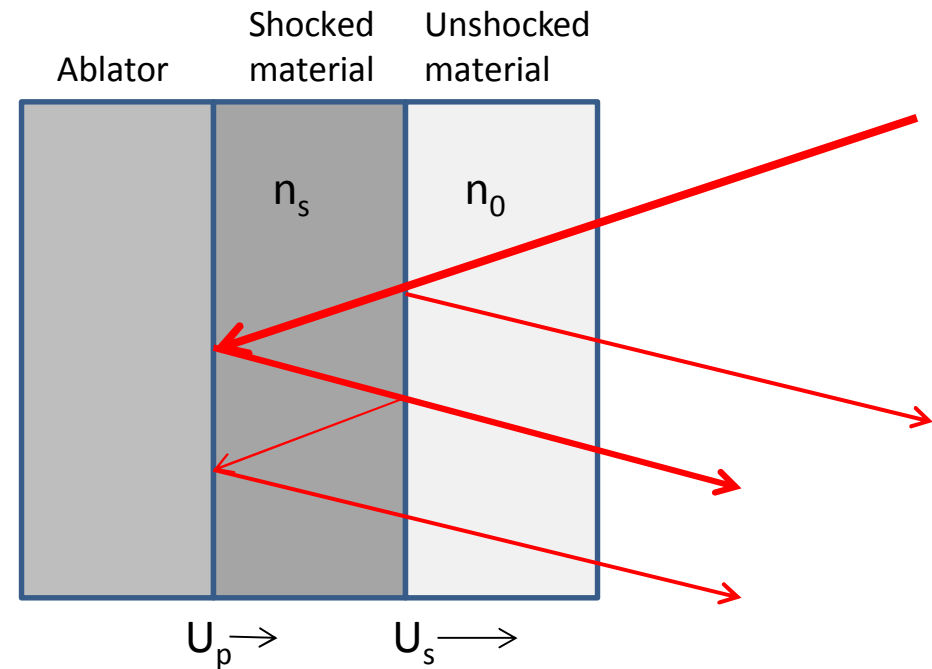
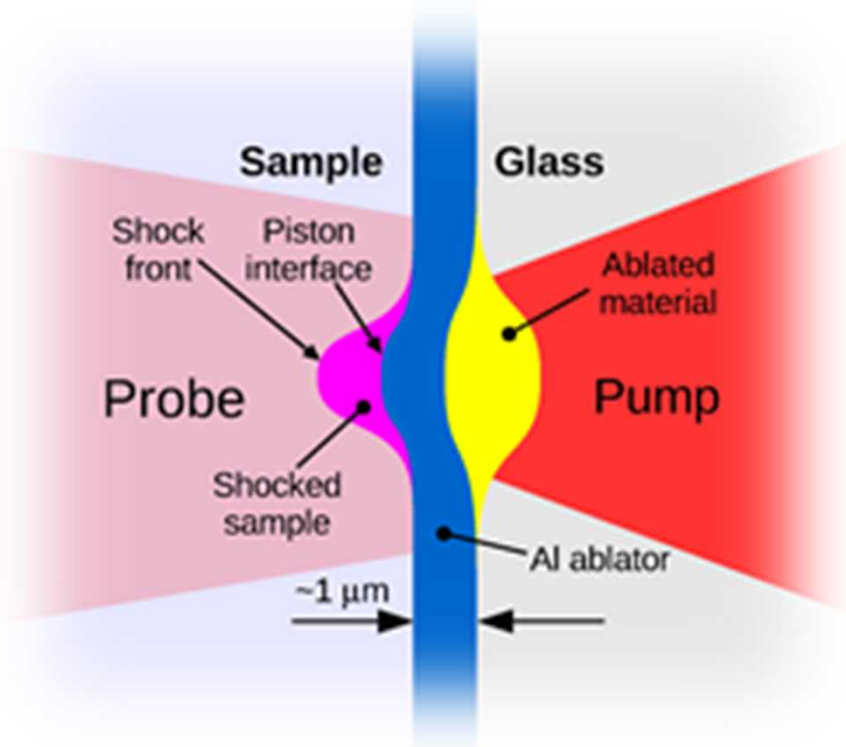
Ultrafast Time Domain Interferometry



We achieve picosecond time resolution by encoding time onto wavelength and measuring wavelength with a slow spectrometer. Pulse are $\sim 300\text{-}400$ ps long with separation of ~ 10 ps.

Timing jitter is minimized by using the same laser pulse to drive the shock and probe U_P and U_S from the other side of the target.

UTDI: Material interaction model

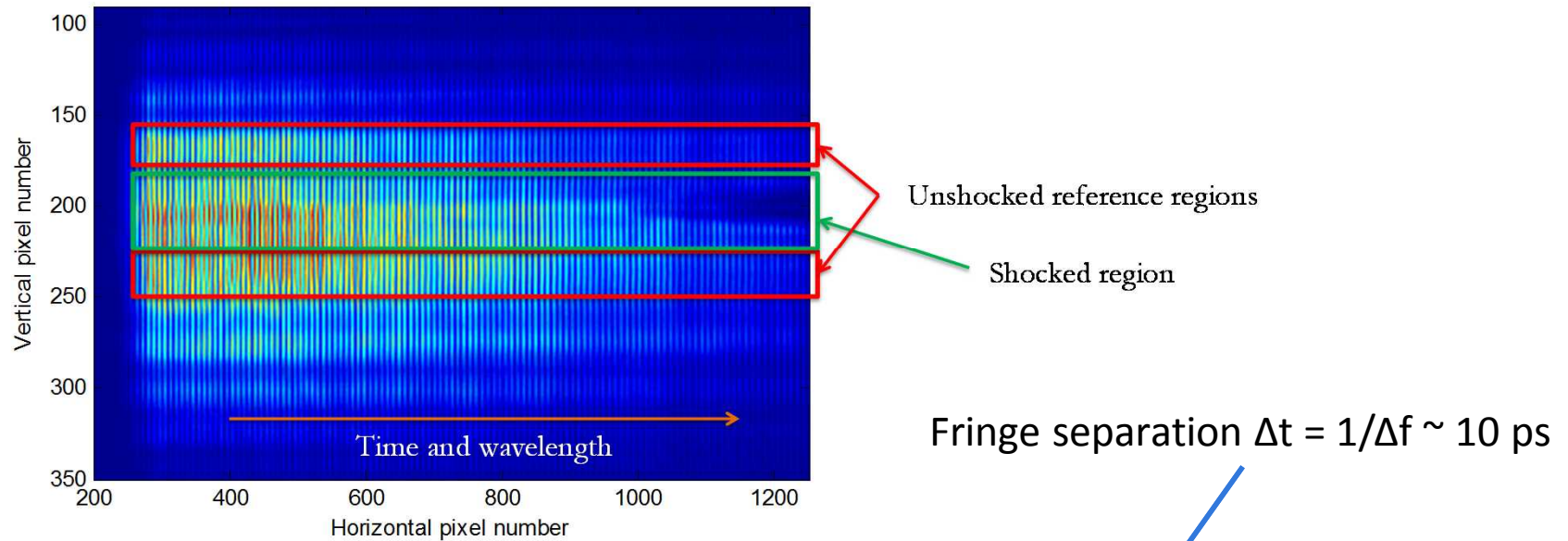


Displacement of ablator and shock are greatly exaggerated for clarity.

Shocks are nominally 1D. 20-30 μm spot size, $<1 \mu\text{m}$ surface displacement.

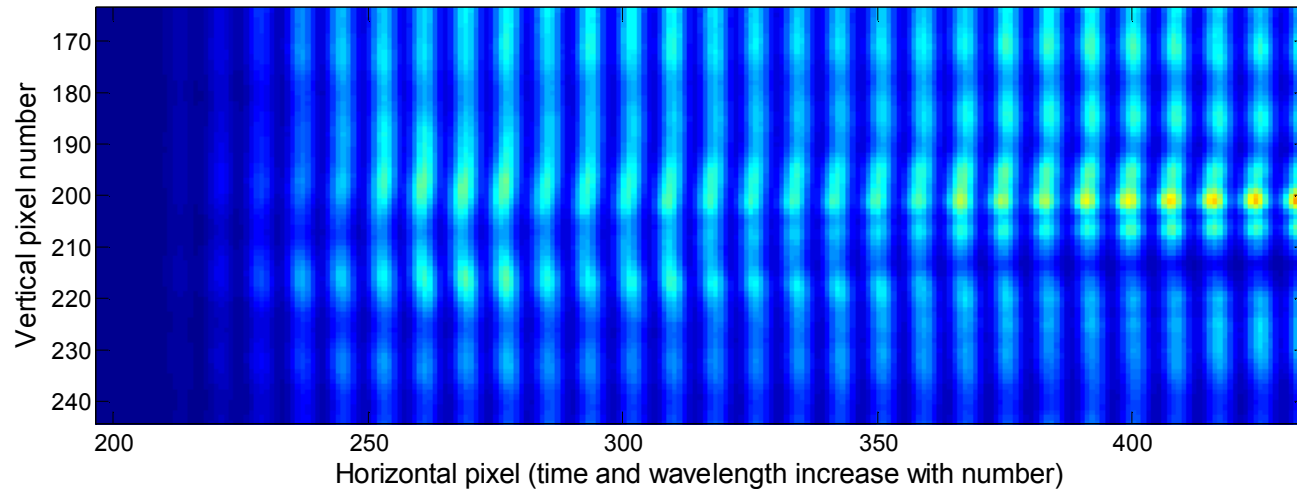
Because $U_s > U_p$, the shock front and ablator form an expanding etalon and the interference of reflections changes with time!

UTDI: A Homodyne Interferometer like VISAR



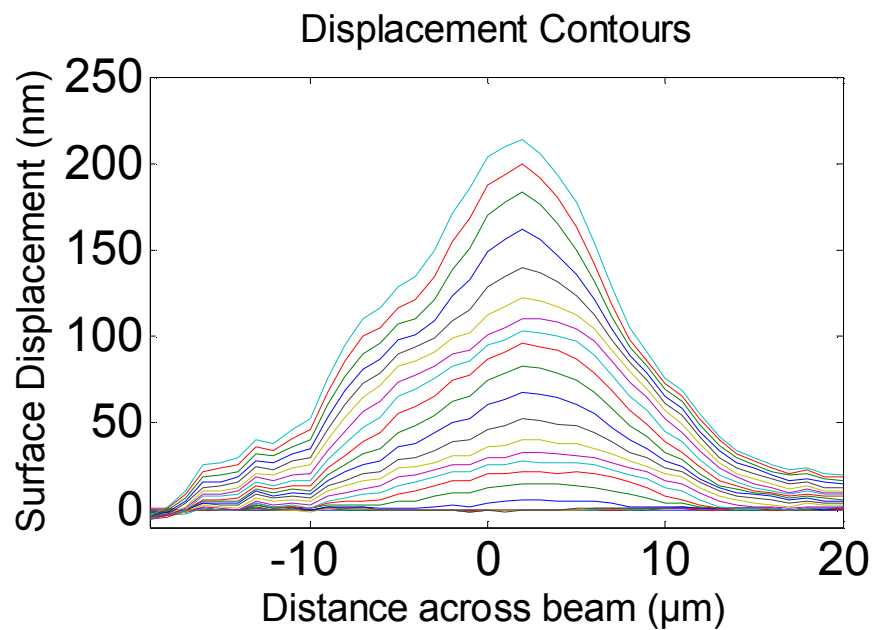
Fringe separation $\Delta t = 1/\Delta f \sim 10$ ps

Interferogram of shocked aluminum ablator

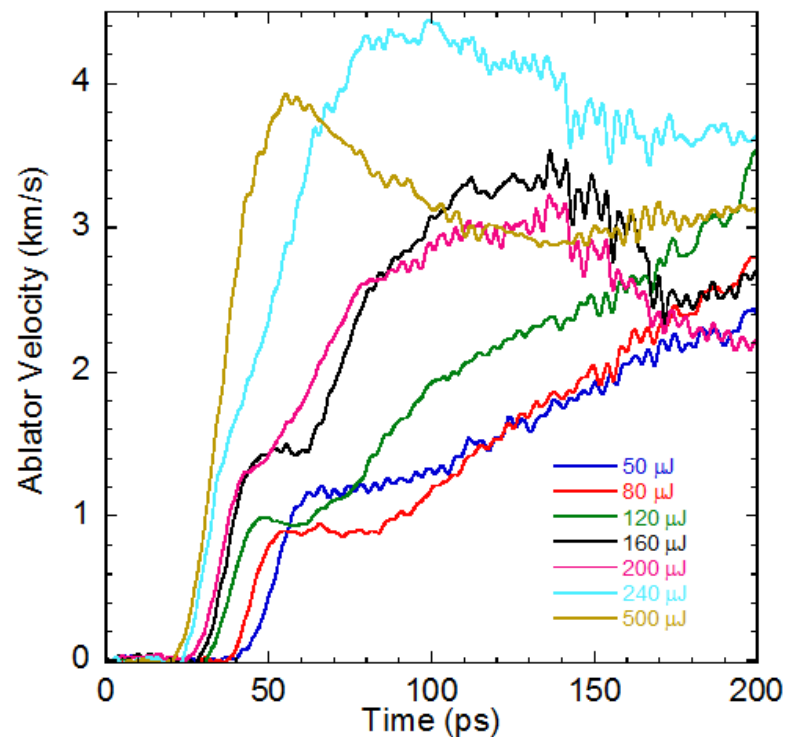


Data from SNL

Bare aluminum ablator data

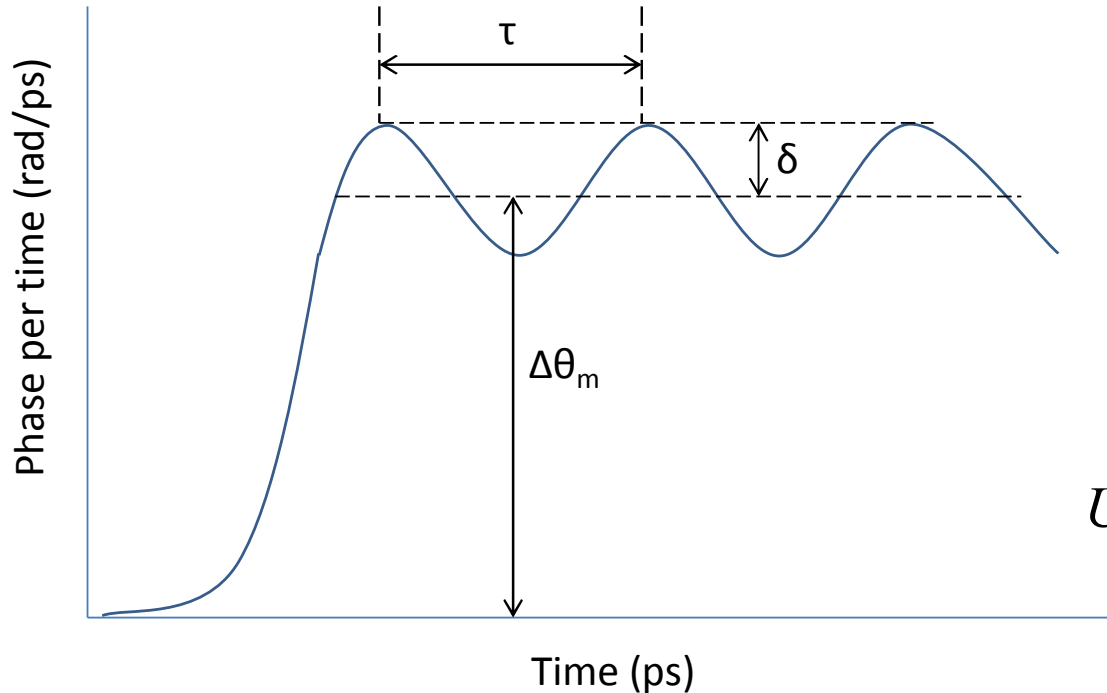


Contours generated by integrating velocity over time for various time steps. Each contour is ~ 8 ps apart.



Free-surface ablator velocity profile changes markedly with drive energy.

Canonical UTDI signal



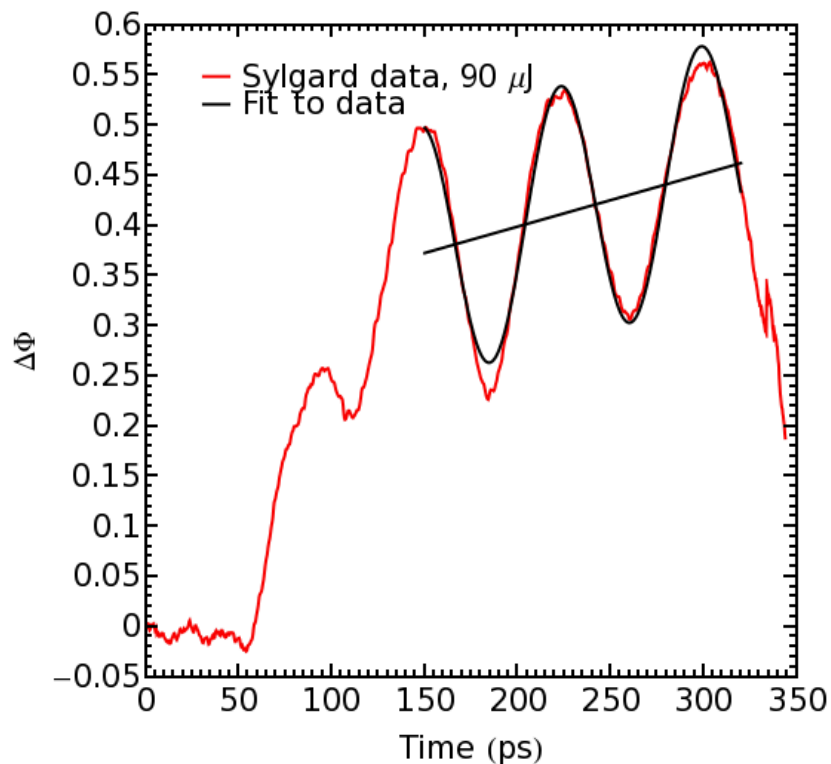
$$U_s = \frac{\lambda}{4\pi n_0} \left[\frac{\Delta\theta_m}{\Delta t} + \frac{2\pi}{\tau} \right]$$

$$n_s = n_0 + \frac{\tau\delta n_0}{2\pi\Delta t - \tau\delta}$$

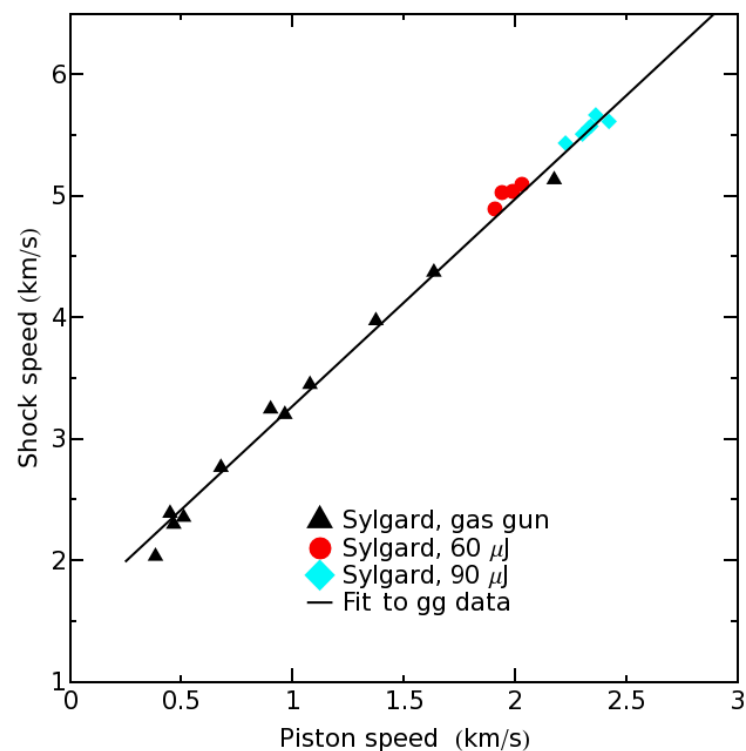
$$U_P = \frac{\lambda}{4\pi n_0} \left[\frac{\Delta\theta_m}{\Delta t} + \frac{2\pi}{\tau} \left(1 - \frac{n_0}{n_s} \right) \right]$$

Expanding etalon between the ablator and shock front causes oscillation in phase of reflected light. The offset, amplitude, and period of oscillation are related to U_P , U_S , and n_s . Besides the fitted parameters, we must also know the pulse separation, Δt , and the unshocked index of refraction, n_0 .

Sylgard Hugoniot data

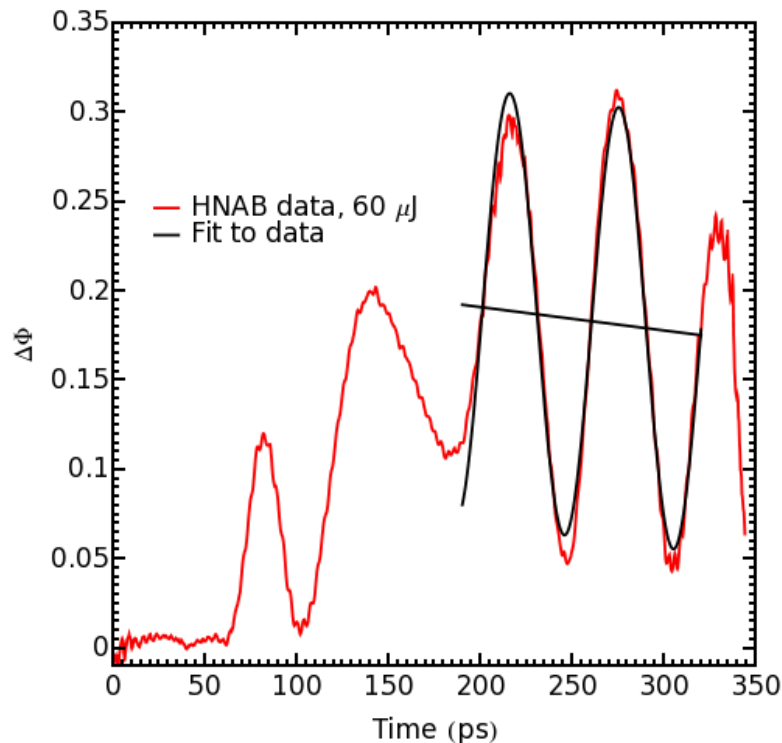


Phase data with region fit to model.

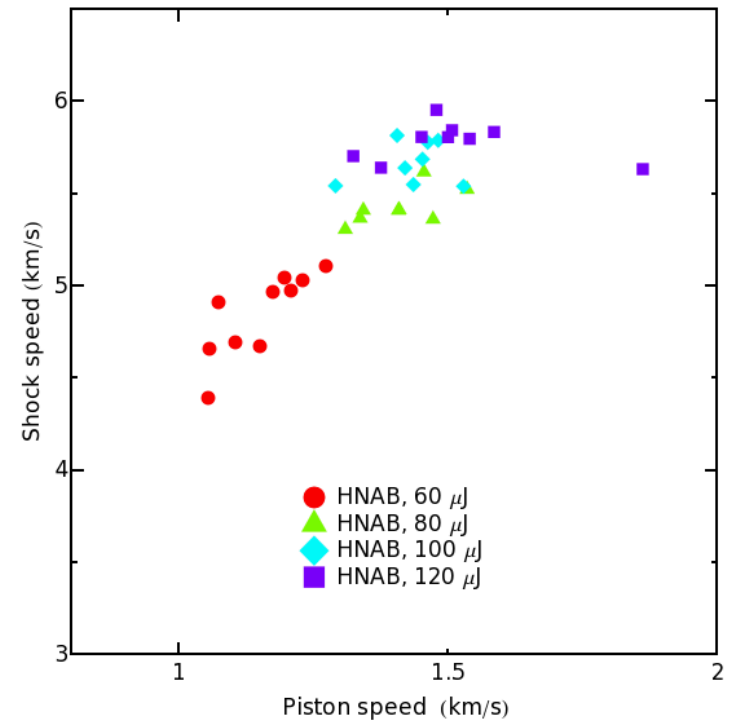


Sylgard Hugoniot compared to gas gun data from Marsh.

HNAB Hugoniot data



Phase data with region fit to model.

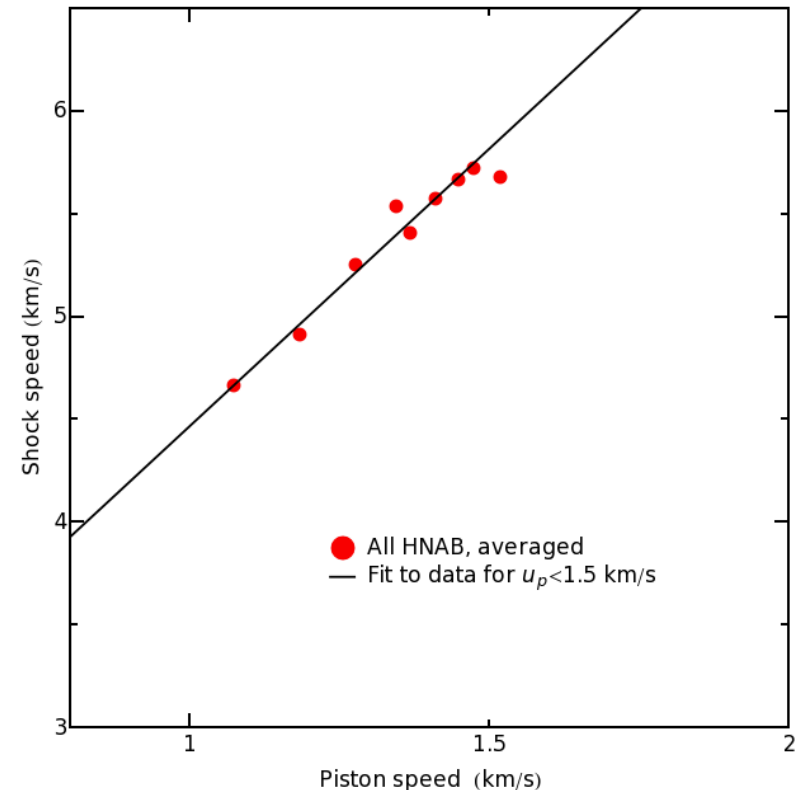


Hugoniot data for crystallized HNAB II assuming an unshocked index of refraction of 1.8.

HNAB Hugoniot details

HNAB II crystals have an index of refraction of 1.5, 1.7, or 1.8, depending on the orientation of the crystal with respect to the light. In our analysis, we assume 1.8.

HNAB II is polycrystalline with crystals oriented randomly. Which index do you pick?



HNAB data from LLNL after binning into groups of 5, averaging, and eliminating one outlier.

Best fit line: $U_s = 2.70U_p + 1.77$ [km/s]

McCrone, W. C. "Crystallographic Study of Hexanitroazobenzene (HNAB)" SAND75-7067, 1967.

Future Work

Characterize unshocked index of refraction of HNAB in crystallized and amorphous states to get more accurate Hugoniot.

Develop pulse-shaping techniques to drive sustained, steady shock in materials.

Perhaps develop a new model for analysis that relaxes requirements on steady-state shocks.