

Energy-Stable Galerkin Reduced Order Models for Nonlinear Compressible Flow

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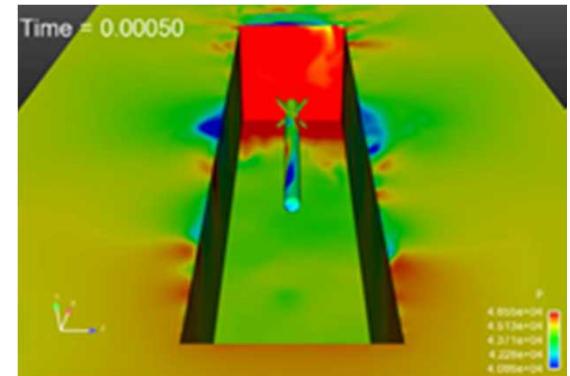
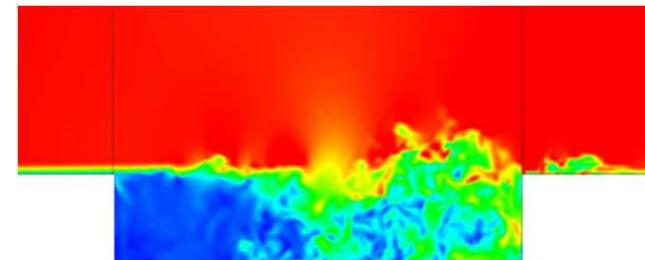
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Motivation

- Despite improved algorithms and powerful supercomputers, **“high-fidelity” models** are often too expensive for use in a design or analysis setting.
- ***Targeted application area in which this situation arises:*** compressible cavity flow problem.
 - **Large Eddy Simulations (LES)** with very fine meshes and long times are required to predict accurately dynamic pressure loads in cavity.

These simulations take **weeks** even when run in parallel on state-of-the-art supercomputers!



Proper Orthogonal Decomposition (POD)/ Galerkin Method to Model Reduction

High-Fidelity
Simulations:

Snapshot 1
Snapshot 2
⋮
Snapshot K

Step 1

Modal
Decomposition
(POD):

$$\mathbf{x}(t) \approx \Phi_M \mathbf{x}_M(t)$$

Step 2

Galerkin Projection
of LTI FOM:

$$\Phi_M^T [\dot{\mathbf{x}}(t) = \mathbf{A}\mathbf{x}(t) + \mathbf{B}\mathbf{u}(t)]$$

- **Snapshot matrix:** $\mathbf{X} = (x^1, \dots, x^K) \in \mathbb{R}^{N \times K}$
- **SVD:** $\mathbf{X} = \mathbf{U}\Sigma\mathbf{V}^T$
- **Truncation:** $\Phi_M = (\phi_1, \dots, \phi_M) = \mathbf{U}(:, 1:M)$

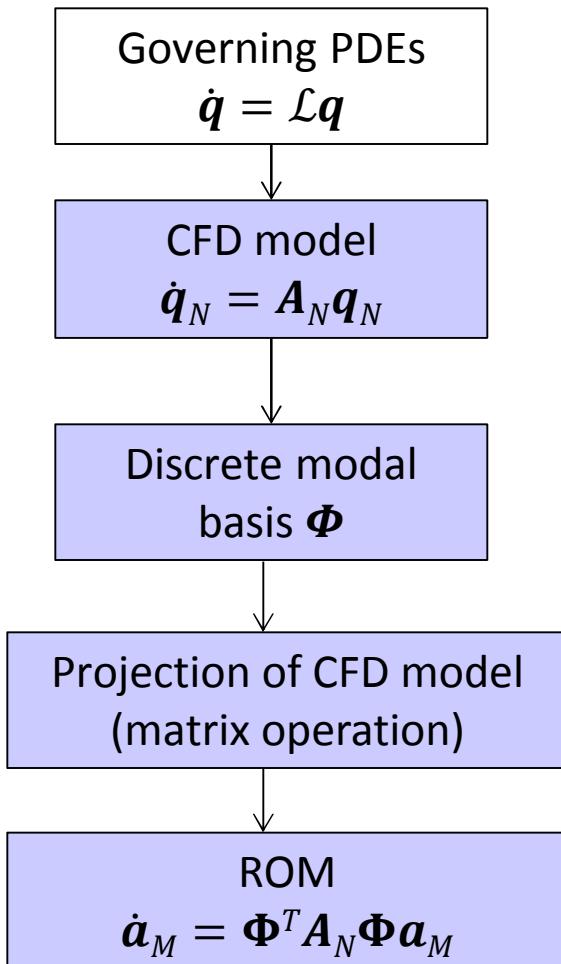
N = # of dofs in high-fidelity simulation
 K = # of snapshots
 M = # of dofs in ROM
($M \ll N, M \ll K$)

“Small” ROM
LTI System:

$$\begin{aligned} \dot{\mathbf{x}}_M(t) &= \Phi_M^T \mathbf{A} \Phi_M \mathbf{x}_M(t) + \Phi_M^T \mathbf{B} \mathbf{u}(t) \\ \mathbf{y}_M(t) &= \mathbf{C} \Phi_M \mathbf{x}_M(t) \end{aligned}$$

Discrete vs. Continuous Galerkin Projection

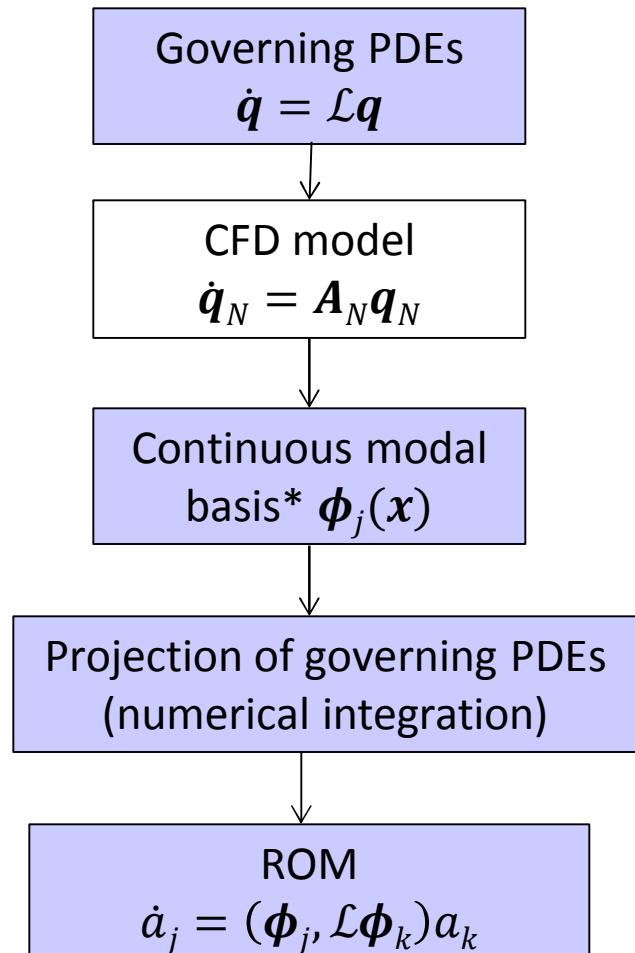
Discrete Projection



This talk
focuses on

If PDEs are linear or have polynomial non-linearities, projection can be calculated in **offline stage** of MOR.

Continuous Projection



* Continuous functions space is defined using finite elements.

Stability Issues of POD/Galerkin ROMs

Full Order Model (FOM)

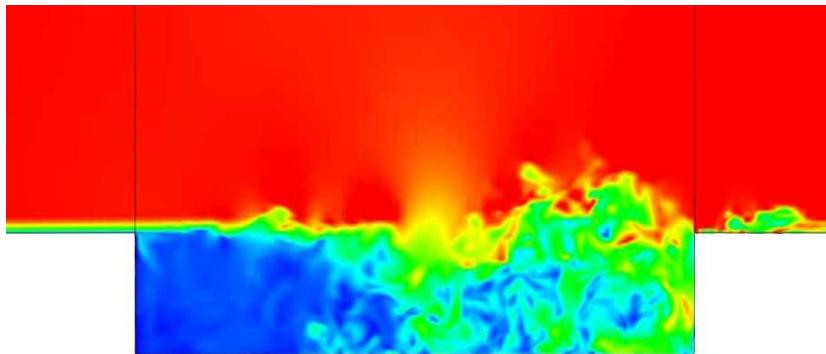
$$\dot{\mathbf{q}}(t) = \mathcal{L}\mathbf{q}(t) + \mathcal{N}(\mathbf{q}(t))$$

Reduced Order Model (ROM)

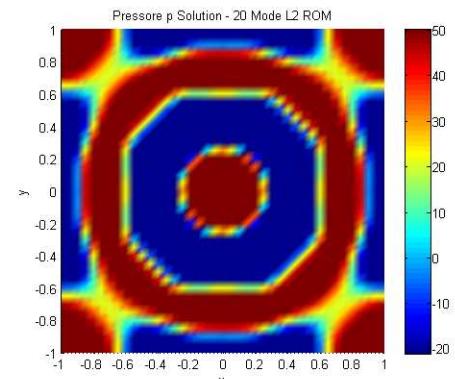
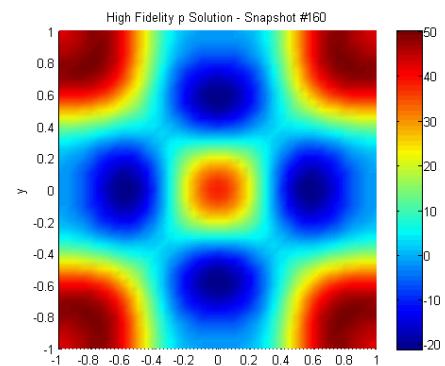
$$\dot{\mathbf{q}}_M(t) = \mathbf{A}_M \mathbf{q}_M(t) + \mathbf{N}_M(\mathbf{q}_M(t))$$

Problem: FOM stable $\not\Rightarrow$ ROM stable!

- There is no *a priori* stability guarantee for POD/Galerkin ROMs.
- Stability of a ROM is commonly evaluated *a posteriori* – **RISKY!**
- Instability of POD/Galerkin ROMs is a **real** problem in some applications...



Top right: FOM
Bottom right: ROM



...e.g., compressible flows, high-Reynolds number flows.

Stability Preserving ROM Approaches: Literature Review

Approaches for building stability-preserving POD/Galerkin ROMs found in the literature fall into **two categories**:

This talk.

1. ROMs which derive *a priori* a stability-preserving model reduction framework (usually specific to an equation set).

Can have an intrusive implementation

- ROMs based on projection in special ‘energy-based’ (not L^2) inner products, e.g., Rowley *et al.* (2004), Barone & Kalashnikova *et al.* (2009), Serre *et al.* (2012).

2. ROMs which stabilize an unstable ROM through an *a posteriori* post-processing stabilization step applied to the algebraic ROM system.

Can have inconsistencies between ROM and FOM physics

- Approaches in which an optimization problem that stabilizes an unstable ROM is formulated and solved, e.g., Amsallem *et al.* (2012), Bond *et al.* (2008), Kalashnikova *et al.* (2014).
- ROMs with increased numerical stability due to inclusion of ‘stabilizing’ terms in the ROM equations, e.g., Wang *et al.* (2012).

Energy-Stability

- **Practical Definition:** Numerical solution does not “blow up” in finite time.
- **More Precise Definition:** Numerical discretization does not introduce any spurious instabilities inconsistent with natural instability modes supported by the governing continuous PDEs.

Numerical solutions must maintain proper **energy balance**.

- Stability of ROM is intimately tied to choice of **inner product** for the Galerkin projection.
- Stability-preserving inner product derived using the **energy method**:
 - Bounds numerical solution energy in a physical way.
 - Borrowed from spectral methods community.
 - Analysis is straightforward for ROMs constructed via **continuous projection**.

Can show: if a Galerkin ROM is constructed in an **energy inner product**, the ROM system energy will be bounded in a way that is consistent with the behavior of the exact solution to the PDEs, i.e., the ROM will be **energy-stable**.

Linearized Compressible Flow Equations

Energy-Stability for Linearized PDEs:

FOM linearly stable \Rightarrow ROM built in energy inner product linearly stable ($Re(\lambda) < 0$)
(WCCM X talk and paper: Kalashnikova & Arunajatesan, 2012).

Linearized compressible Euler/Navier-Stokes equations are appropriate when a compressible fluid system can be described by small-amplitude perturbations about a steady-state mean flow.

- Linearization of full compressible Euler/Navier-Stokes equations obtained as follows:
 1. Decompose fluid field as **steady mean** plus **unsteady fluctuation**

$$\mathbf{q}(\mathbf{x}, t) = \bar{\mathbf{q}}(\mathbf{x}) + \mathbf{q}'(\mathbf{x}, t)$$

2. Linearize full nonlinear compressible Navier-Stokes equations around steady mean to yield **linear hyperbolic/incompletely parabolic** system

$$\dot{\mathbf{q}}' + \mathbf{A}_i(\bar{\mathbf{q}}) \frac{\partial \mathbf{q}'}{\partial x_i} + \frac{\partial}{\partial x_j} \left[\mathbf{K}_{ij}(\bar{\mathbf{q}}) \frac{\partial \mathbf{q}'}{\partial x_i} \right] = \mathbf{0}$$

Energy-Stable ROMs for Linearized Compressible Flow

Linearized compressible Euler/Navier-Stokes equations are **symmetrizable** (Barone & Kalashnikova, 2009; Kalashnikova & Arunajatesan, 2012).

- There exists a symmetric positive definite matrix $\mathbf{H} \equiv \mathbf{H}(\bar{\mathbf{q}})$ (system “**symmetrizer**”) s.t.:
 - The convective flux matrices $\mathbf{H}\mathbf{A}_i$ are symmetric
 - The following augmented viscosity matrix is symmetric positive semi-definite

$$\mathbf{K}^s = \begin{pmatrix} \mathbf{H}\mathbf{K}_{11} & \mathbf{H}\mathbf{K}_{12} & \mathbf{H}\mathbf{K}_{13} \\ \mathbf{H}\mathbf{K}_{21} & \mathbf{H}\mathbf{K}_{22} & \mathbf{H}\mathbf{K}_{23} \\ \mathbf{H}\mathbf{K}_{31} & \mathbf{H}\mathbf{K}_{32} & \mathbf{H}\mathbf{K}_{33} \end{pmatrix}$$

Symmetry Inner Product (weighted L^2 inner product):

$$(\mathbf{q}_1, \mathbf{q}_2)_H = \int_{\Omega} \mathbf{q}_1 \mathbf{H} \mathbf{q}_2 d\Omega$$

- If ROM is built in **symmetry inner product**, Galerkin approximation will satisfy the same energy expression as continuous PDEs:

$$\|\mathbf{q}'_M(\mathbf{x}, t)\|_H \leq e^{\beta t} \|\mathbf{q}'_M(\mathbf{x}, 0)\|_H \quad (\Rightarrow \frac{dE_M}{dt} \leq 0 \text{ for uniform base flow})$$

Symmetrizers for Several Hyperbolic/ Incompletely Parabolic Systems

- **Wave equation:** $\ddot{u} = a^2 \frac{\partial^2 u}{\partial x^2}$ or $\dot{q} = \mathbf{A} \frac{\partial q}{\partial x}$ where $q = \left(\dot{u}, \frac{\partial u}{\partial x} \right)$ $\Rightarrow \mathbf{H} = \begin{pmatrix} 1 & 0 \\ 0 & a^2 \end{pmatrix}$

- **Shallow water equations:** $\dot{q}' + \mathbf{A}_i(\bar{q}) \frac{\partial q'}{\partial x_i} = \mathbf{0}$ $\Rightarrow \mathbf{H} = \begin{pmatrix} \bar{\phi} & 0 & 0 \\ 0 & \bar{\phi} & 0 \\ 0 & 0 & 1 \end{pmatrix}$

- **Linearized compressible Euler:** $\dot{q}' + \mathbf{A}_i(\bar{q}) \frac{\partial q'}{\partial x_i} = \mathbf{0}$ $\Rightarrow \mathbf{H} = \begin{pmatrix} \bar{\rho} & 0 & 0 \\ 0 & \alpha^2 \gamma \bar{\rho}^2 \bar{p} & \bar{\rho} \alpha^2 \\ 0 & 0 & \frac{(1+\alpha^2)}{\gamma \bar{p}} \end{pmatrix}$

- **Linearized compressible Navier-Stokes:** $\dot{q}' + \mathbf{A}_i(\bar{q}) \frac{\partial q'}{\partial x_i} + \frac{\partial}{\partial x_j} \left[\mathbf{K}_{ij}(\bar{q}) \frac{\partial q'}{\partial x_i} \right] = \mathbf{0}$

$$\Rightarrow \mathbf{H} = \begin{pmatrix} \bar{\rho} & 0 & 0 \\ 0 & \frac{\bar{\rho} R}{\bar{T}(\gamma - 1)} & 0 \\ 0 & 0 & \frac{R \bar{T}}{\bar{\rho}} \end{pmatrix}$$

- Barone & Kalashnikova, *JCP*, 2009.
- Kalashnikova & Arunajatesan, *WCCM X*, 2012.
- Kalashnikova *et al.*, *SAND report*, 2014.

Continuous Projection Implementation: “Spirit” Code

“Spirit” ROM Code = 3D parallel C++ POD/Galerkin test-bed ROM code that uses data-structures and eigensolvers from **Trilinos** to build energy-stable ROMs for compressible flow problems
 → stand-alone code that can be synchronized with any high-fidelity code!

- POD modes defined using piecewise smooth finite elements.
- Gauss quadrature rules of sufficient accuracy are used to compute exactly inner products with the help of the `libmesh` library.
- Physics in Spirit:
 - ***Linearized compressible Euler*** (L^2 , energy inner product).
 - ***Linearized compressible Navier-Stokes*** (L^2 , energy inner product).
 - ***Nonlinear isentropic compressible Navier-Stokes*** (L^2 , stagnation energy, stagnation enthalpy inner product).
 - ***Nonlinear compressible Navier-Stokes*** (L^2 , energy inner product).

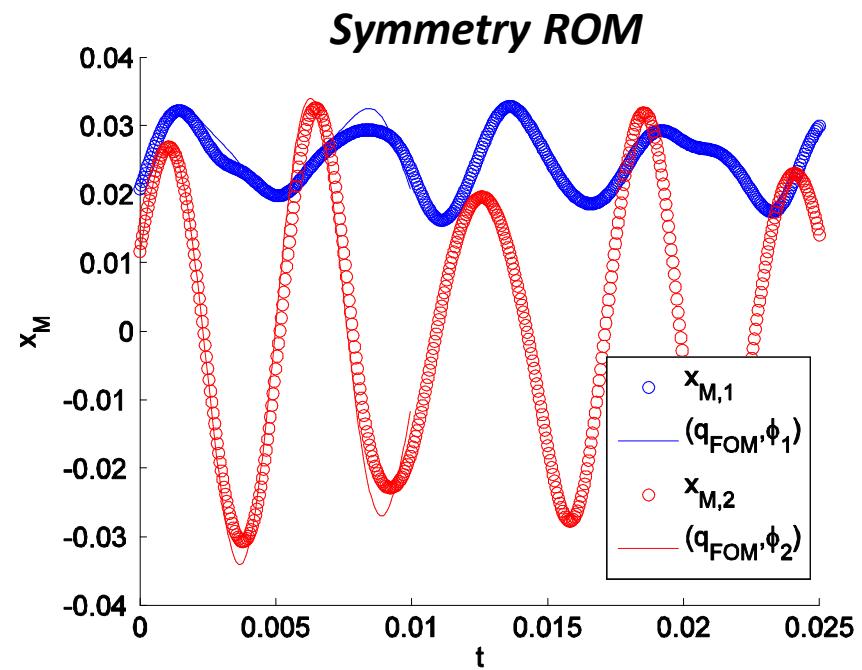
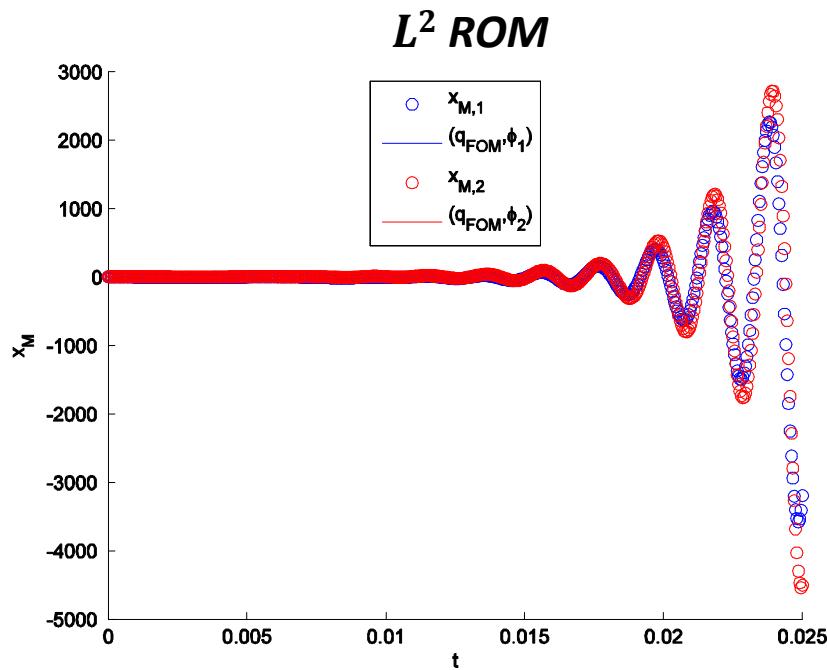
First, testing of ROMs for these physics

“SIGMA CFD” High-Fidelity Code = Sandia in-house finite volume flow solver derived from LESLIE3D (Genin & Menon, 2010), an LES flow solver originally developed in the Computational Combustion Laboratory at Georgia Tech.

Numerical Experiment: 2D Inviscid Pressure Pulse

- Inviscid pulse in a uniform base flow (linear dynamics).
- High-fidelity simulation run on mesh with 3362 nodes, up to time $t = 0.01$ seconds.
- 200 snapshots of solution used to construct $M = 20$ mode ROM in L^2 and symmetry inner products.

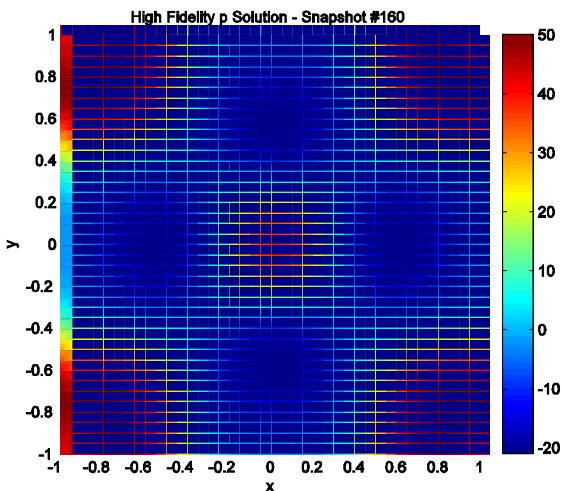
$x_{M,i}(t)$ vs. $(\mathbf{q}'_{CFD}, \phi_i)$ for $i = 1, 2$



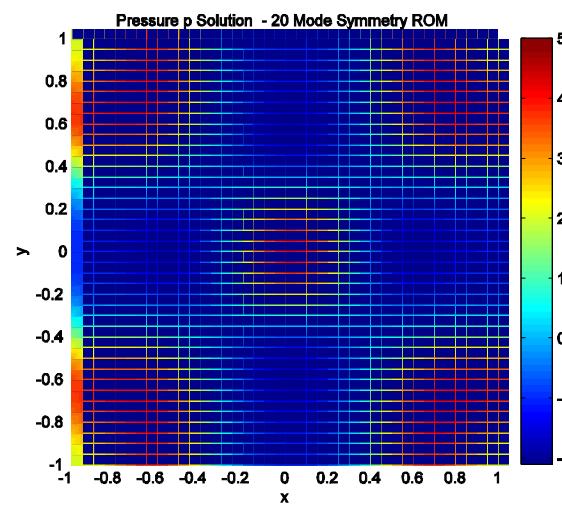
Numerical Experiment: 2D Inviscid Pressure Pulse (cont'd)

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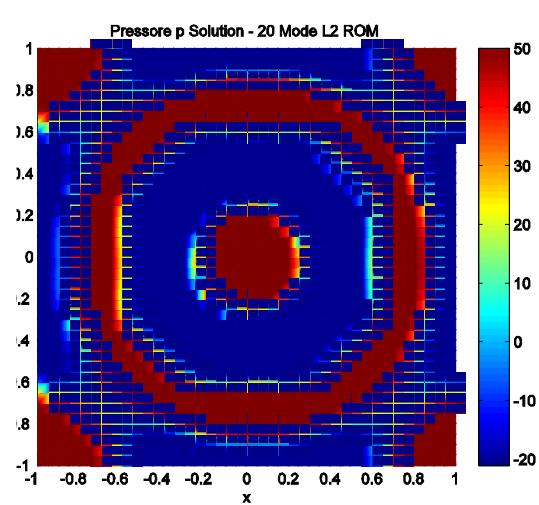
p' : High-fidelity



p' : Symmetry ROM



p' : L^2 ROM



time of snapshot 160

Nonlinear Compressible Flow Equations

Energy-Stability for Nonlinear PDEs:

ROM built in energy inner product will preserve stability of an equilibrium point at 0 for the governing nonlinear system of PDEs (Rowley, 2004; Kalashnikova *et al.*, 2014).

- Compressible isentropic Navier-Stokes equations (cold flows, moderate Mach #):

$$\begin{aligned}\frac{Dh}{Dt} + (\gamma - 1)h\nabla \cdot \mathbf{u} &= 0 \\ \frac{D\mathbf{u}}{Dt} + \nabla h - \frac{1}{Re}\Delta\mathbf{u} &= \mathbf{0}\end{aligned}$$

h = enthalpy
 \mathbf{u} = velocity vector
 ρ = density
 T = temperature
 $\boldsymbol{\tau}$ = viscous stress tensor

- Full compressible Navier-Stokes equations:

$$\begin{aligned}\rho \frac{D\mathbf{u}}{Dt} + \frac{1}{\gamma M^2} \nabla(\rho T) - \frac{1}{Re} \nabla \cdot \boldsymbol{\tau} &= \mathbf{0} \\ \frac{D\rho}{Dt} + \rho \nabla \cdot \mathbf{u} &= 0 \\ \rho \frac{DT}{Dt} + (\gamma - 1)\rho T \nabla \cdot \mathbf{u} - \frac{\gamma}{PrRe} \nabla \cdot (\kappa \nabla T) - \left(\frac{\gamma(\gamma - 1)M^2}{Re} \right) \nabla \mathbf{u} \cdot \boldsymbol{\tau} &= 0\end{aligned}$$

Energy-Stable ROMs for Nonlinear Compressible Flow (Isentropic NS)

In (Rowley, 2004), Rowley *et al.* showed that energy inner product for the compressible isentropic Navier-Stokes equations can be defined following a transformation of these equations.

- Transformed compressible isentropic Navier-Stokes equations:

$$\frac{Dc}{Dt} + \frac{\gamma - 1}{2} c \nabla \cdot \mathbf{u} = 0$$

$$\frac{D\mathbf{u}}{Dt} + \frac{2}{\gamma - 1} c \nabla c - \frac{1}{Re} \Delta \mathbf{u} = \mathbf{0}$$

c = speed of sound
 $(c^2 = (\gamma - 1)h)$
 \mathbf{u} = velocity

- Family of inner products:

$$(\mathbf{q}_1, \mathbf{q}_2)_\alpha = \int_{\Omega} \left(\mathbf{u}_1 \cdot \mathbf{u}_2 + \frac{2\alpha}{\gamma - 1} c_1 c_2 \right) d\Omega$$

$$\alpha = \begin{cases} 1 \Rightarrow \|\mathbf{q}\|_\alpha = \text{stagnation enthalpy} \\ \frac{1}{\gamma} \Rightarrow \|\mathbf{q}\|_\alpha = \text{stagnation energy} \end{cases}$$

If Galerkin projection step of model reduction is performed in α inner product, then the Galerkin projection will **preserve the stability of an equilibrium point at the origin** (Rowley, 2004).

Energy-Stable ROMs for Nonlinear Compressible Flow (Full NS)

Present work extends ideas in (Rowley, 2004) to **full compressible Navier-Stokes equations**.

- First, full compressible Navier-Stokes equations are **transformed** into the following variables:

$$a = \sqrt{\rho}, \quad \mathbf{b} = a\mathbf{u}, \quad d = ae$$

e =internal energy

- Next, the following “**total energy**” inner product is defined:

$$(\mathbf{q}_1, \mathbf{q}_2)_{TE} = \int_{\Omega} (\mathbf{b}_1 \cdot \mathbf{b}_2 + a_1 d_2 + a_2 d_1) d\Omega$$

→ Norm induced by total energy inner product is the total energy of the fluid system:

$$\|\mathbf{q}\|_{TE} = \int_{\Omega} \left(\rho e + \frac{1}{2} \rho u_i u_i \right) d\Omega$$

If Galerkin projection step of model reduction is performed in total energy inner product, then the Galerkin projection will **preserve the stability of an equilibrium point at the origin** (Kalashnikova *et al.*, 2014)

- ☺ Transformed equations have only **polynomial non-linearities** (projection of which can be computed in offline stage of MOR and stored).
- ☺ Transformation introduces **higher order polynomial** non-linearities.
- ☺ Efficiency of online stage of MOR can be recovered using **interpolation** (e.g., DEIM, gappy POD).

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Numerical Experiment: Viscous Laminar Cavity

- Viscous cavity problem at $M = 0.6$, $Re = 1500$ (laminar regime).
- High-fidelity simulation:** DNS based on full nonlinear compressible Navier-Stokes equations with $\approx 117,000$ nodes (right).
- 500 snapshots collected, every $\Delta t_{snap} = 1 \times 10^{-4}$ seconds.
- Snapshots used to construct $M = 15$ mode ROM for nonlinear compressible Navier-Stokes equations in **L^2 and total energy inner products**.
- $M = 15$ mode POD bases capture $\approx 99\%$ of snapshot energy.

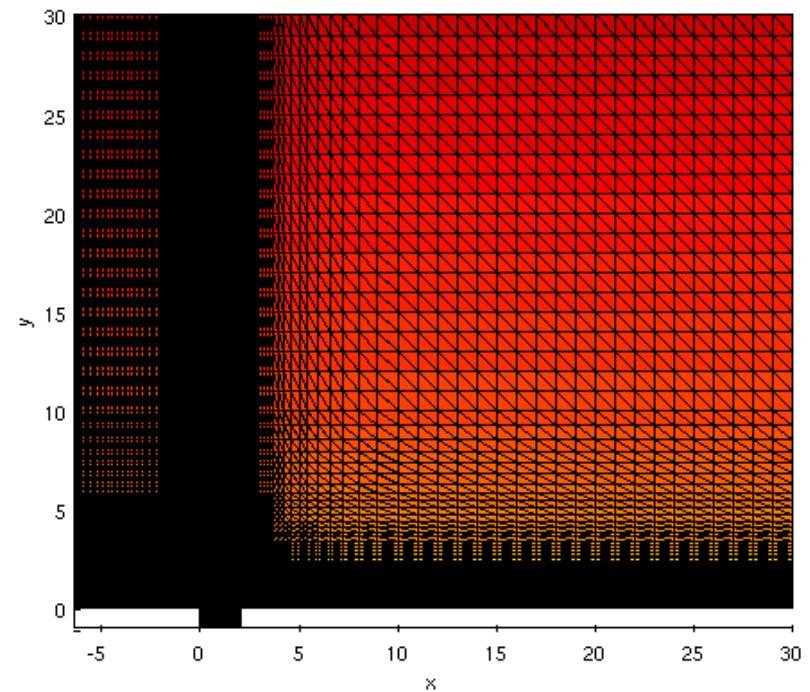
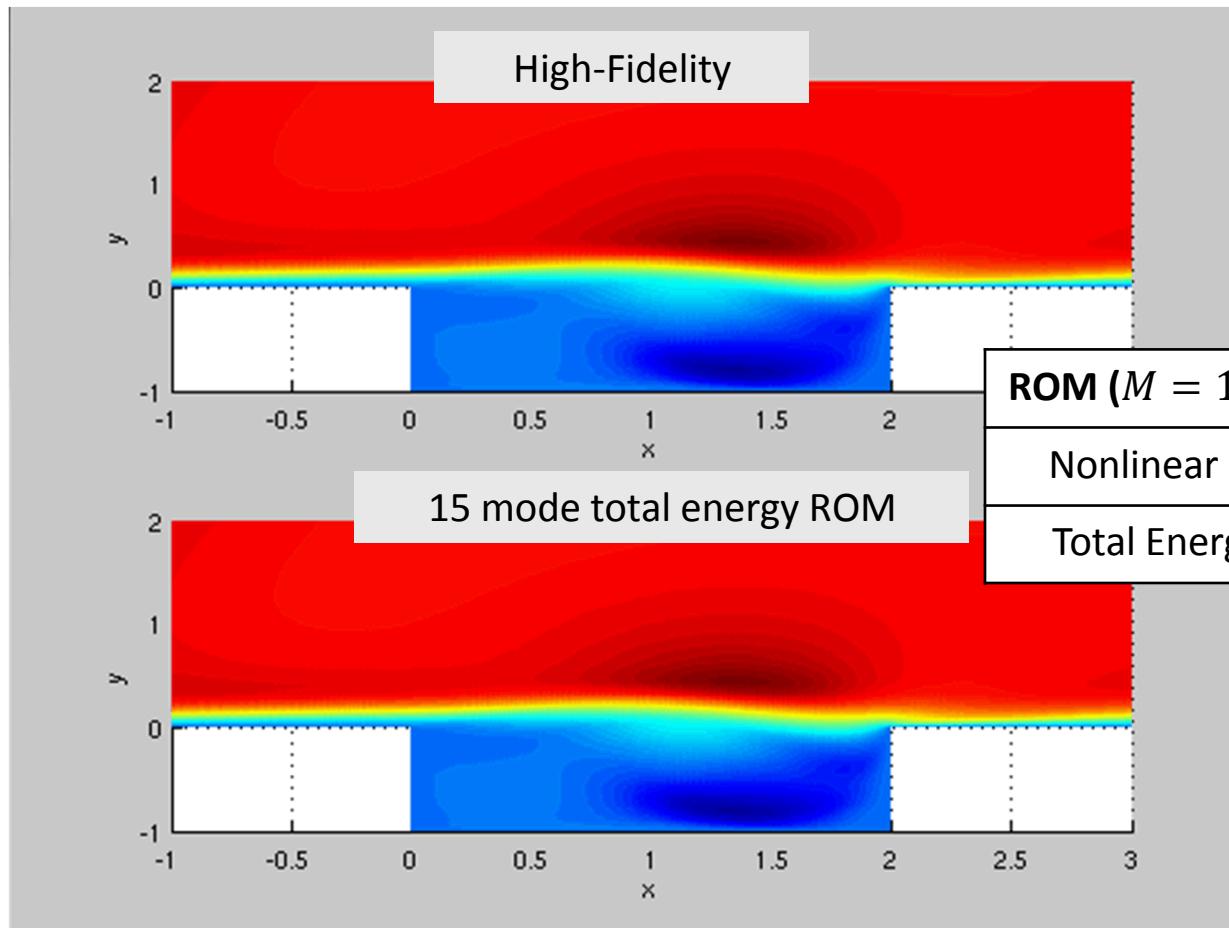


Figure above: viscous laminar cavity problem domain/mesh.

Numerical Experiment: Viscous Laminar Cavity (cont'd)



- L^2 ROM exhibited instability for for $M > 5$ modes.

- In contrast, total energy ROM remained stable and agreed well with high-fidelity solution!

Figure above: u -component of velocity as a function of time t

Summary & Future Work

- A Galerkin model reduction approach in which the ***continuous*** equations are projected onto the basis modes in a continuous inner product is proposed.
- It is shown that the choice of inner product for the Galerkin projection step is crucial to stability of the ROM.
 - For ***linearized compressible flow***, Galerkin projection in the “symmetry” inner product leads to a ROM that is energy-stable for any choice of basis.
 - For ***nonlinear compressible flow***, an inner product that induces the total energy of the fluid system is developed. A ROM constructed in this inner product will preserve the stability of an equilibrium point at 0 for the system.
- Results are promising for a nonlinear problem involving ***compressible viscous laminar flow*** over an open cavity: a total energy ROM remains stable whereas an L^2 ROM exhibits an instability.

Ongoing/Future Work

- Improve efficiency of nonlinear ROMs through interpolation (e.g., DEIM, gappy POD)
- Studies of predictive capabilities of ROMs (robustness w.r.t. parameter changes).

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Thank You! Questions?

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Some references on these ideas:

- **I. Kalashnikova**, S. Arunajatesan, M.F. Barone, B.G. van Bloemen Waanders, J.A. Fike. Reduced Order Modeling for Prediction and Control of Large-Scale Systems. *Sandia National Laboratories Report, SAND No. 2014-4693* (2014).
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