

Ultrafast Laser Diagnostics to Interrogate High Pressure, Highly Collisional Plasma Environments

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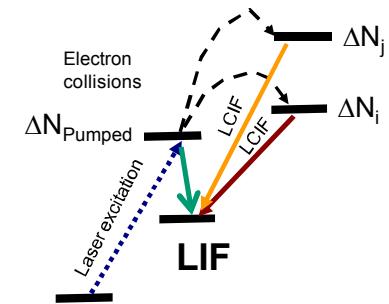
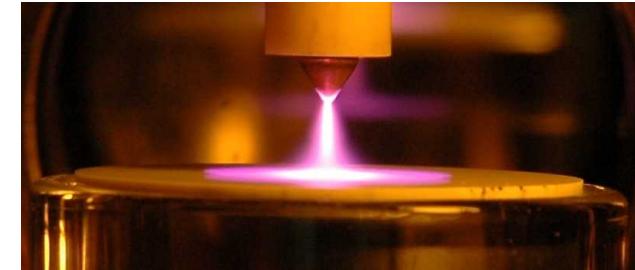
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INTRODUCTION AND MOTIVATION: CHALLENGES AND OBJECTIVES

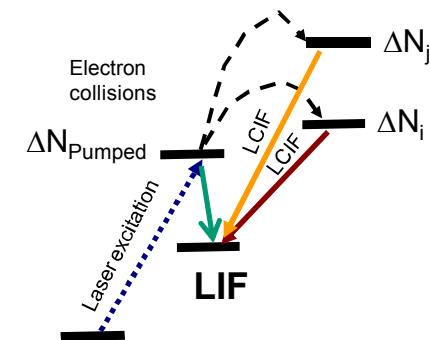
- High pressure (~ 1 ATM and beyond) plasma is challenging environment.
 - Higher densities.
 - Chemically complex environments
 - Smaller length scales.
 - Shorter lifetimes.
- Investigate diagnostics methods to access this challenging environment
 - Extend laser-collision induced fluorescence (LCIF) to atmospheric pressure helium discharge.
 - Examine suitability of ultrafast-short pulse lasers for use to these environments.



Outline challenges and demonstrate implementation of LCIF to highly collisional environments.

OVERVIEW OF CHALLENGES OF EXTENDING LCIF TO HIGH PRESSURE

- Observed LCIF is superposition of several complex processes.
 - In general need a good model to describe the redistribution.



"Electron mixing"

$$\frac{dN_j}{dt} = \left[\sum_{i \neq j} K_{ij}^e N_i - \sum_{i \neq j} K_{ji}^e N_j \right] n_e + \sum_k \left[\sum_{i \neq j} K_{ikj}^a N_i - \sum_{i \neq j} K_{jki}^a N_j \right] N_k + \left[\sum_{i > j} A_{ij} N_i - \sum_{i < j} A_{ji}^j N_j \right]$$

"Neutral mixing"

$$\frac{dN_j}{dt} = \left[\sum_{i \neq j} K_{ij}^e N_i - \sum_{i \neq j} K_{ji}^e N_j \right] n_e + \sum_k \left[\sum_{i \neq j} K_{ikj}^a N_i - \sum_{i \neq j} K_{jki}^a N_j \right] N_k + \left[\sum_{i > j} A_{ij} N_i - \sum_{i < j} A_{ji}^j N_j \right]$$

"Photon mixing"

$$\frac{dN_j}{dt} = \left[\sum_{i \neq j} K_{ij}^e N_i - \sum_{i \neq j} K_{ji}^e N_j \right] n_e + \sum_k \left[\sum_{i \neq j} K_{ikj}^a N_i - \sum_{i \neq j} K_{jki}^a N_j \right] N_k + \left[\sum_{i > j} A_{ij} N_i - \sum_{i < j} A_{ji}^j N_j \right]$$



$$\Delta N_j \sim K_{ij}^e n_e \times \Delta N_i \times \Delta t$$

$$\text{Photons emitted} \sim A_{jk} \times \Delta N_j \times \Delta t$$

Simplifications likely not to be so forthcoming at higher pressures

NEUTRAL INTERACTIONS KEY SOURCE OF INCREASED COMPLEXITY

- Neutral-impact redistribution can play dominant role at higher pressures
 - More-types and evolving nature of neutrals (dimers).
 - “Book keeping” can require sophisticated models.
 - Uncertainties in species and cross-sections limit accuracy.

Neutral mixing

$$\frac{dN_j}{dt} = \sum_k \left[\sum_{i \neq j} K_{ikj}^a N_i - \sum_{i \neq j} K_{jki}^a N_j \right] N_k$$

Bounds on detection

$$\frac{\Delta N_{Electrons}}{\Delta N_{Neutrals}} \sim \frac{K^e n_e}{K^N n_0} \sim 1$$

$$n_e \sim \frac{K^N}{K^e} n_0 \sim \frac{10^{-11}}{10^{-5}} n_0$$

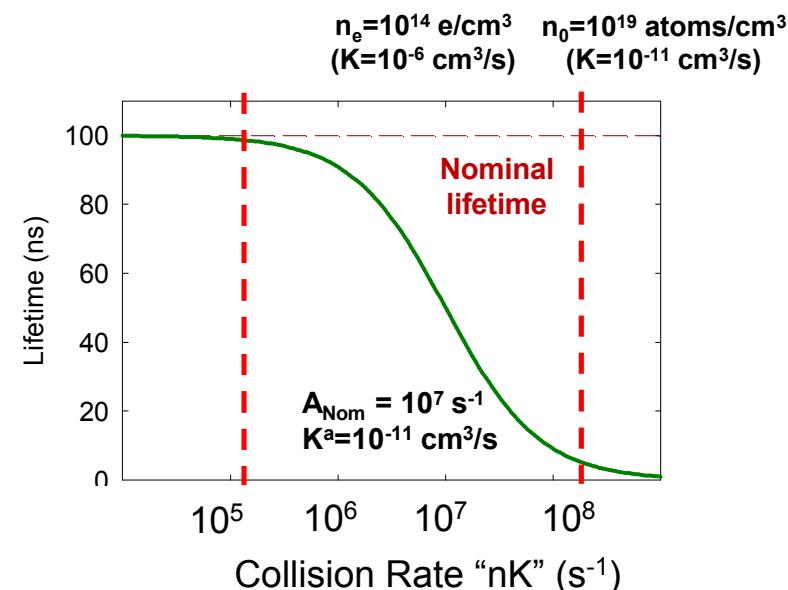
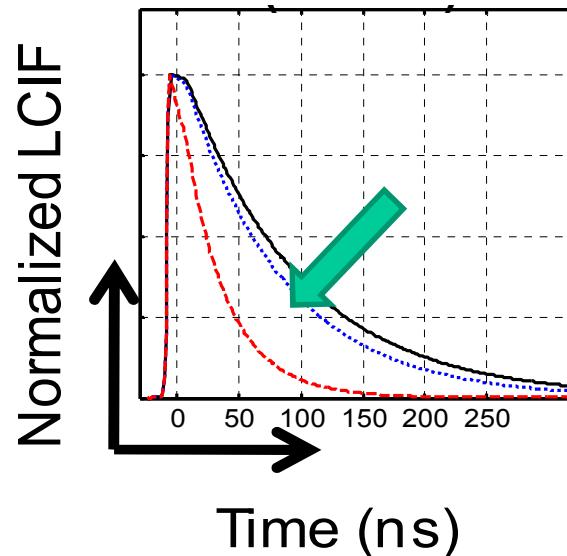
Limit $\sim 10^{13}$ e/cm³ at 1000 Torr

Neutrals are anticipated place limits on lower bound of electron detection.

LIFETIMES OF EXCITED STATES BECOME VERY SHORT AT HIGHER DENSITIES

- Physics of electron-impact redistribution is not expected to change at higher pressures.
 - Sheer number of electrons increase probability of redistribution.
 - Effective Lifetimes become reduced because of redistribution.

$$\frac{dN_j}{dt} \sim \left[\sum_{i \neq j} K_{ij}^e N_i \right] n_e$$

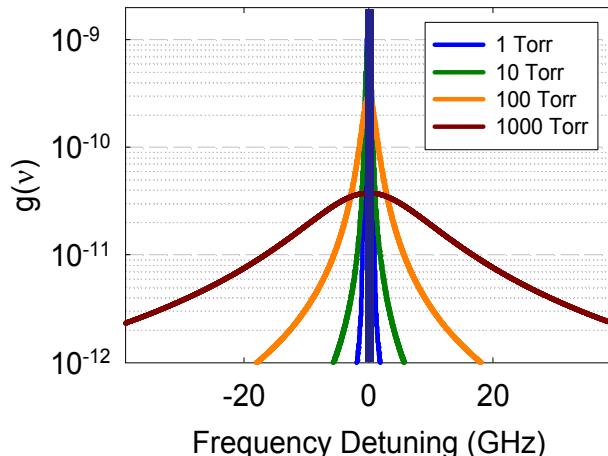


Lifetime of excited states are quite short (~5 ns) at target conditions.

ULTRASHORT-PULSE LASER IS USED FOR INITIATION LCIF EVENT

- Ti:Sapphire, regenerative laser used to generate excitation pulse.
 - Tuned amplifier to 780 nm – doubled in BBO for ~ 390 nm.
 - ~ 100 fs pulse with 10 nm bandwidth (~ 100 cm $^{-1}$).
- Short-pulse laser well suited to interrogate short lifetimes (< 10 ns) and broad absorption profiles (~ 1 nm) associated with high pressure.
 - Still realize “step-like” populating process.
 - Sample most or all of the probed states.

Anticipated absorption profiles



Estimates of linewidths

ns laser ~ 0.5 GHz or 0.1 cm $^{-1}$

fs laser ~ 500 GHz or 100 cm $^{-1}$

Pressure Broadening ~ 0.01 GHz/Torr (He)

Short pulse enables access to all of the pressure-broadened line.

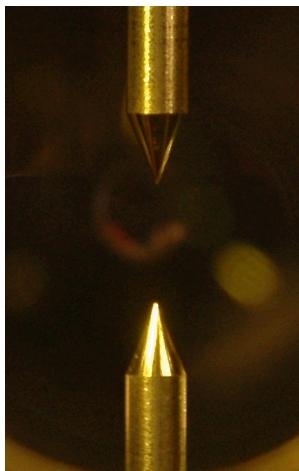


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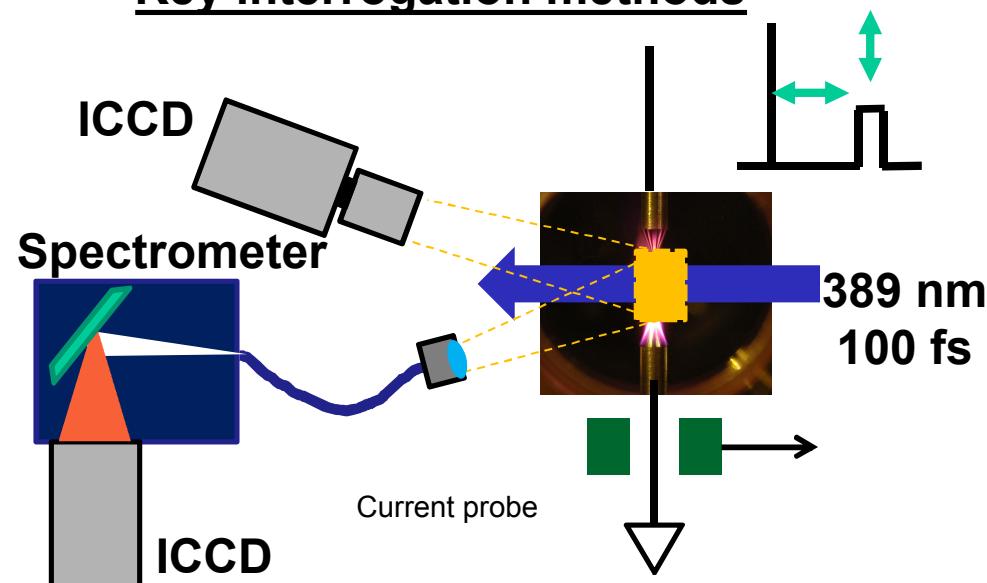
KEY CHALLENGE: GENERATING AND MANIPULATING WELL CHARACTERIZED PLASMA

- Need plasma with well controlled n_e , E/N to calibrate LCIF.
 - Plasma generation in 640 Torr He.
 - Double pulse method to separate generation and interrogation.
 - Spectrometer to identify, camera to image.

Discharge configuration



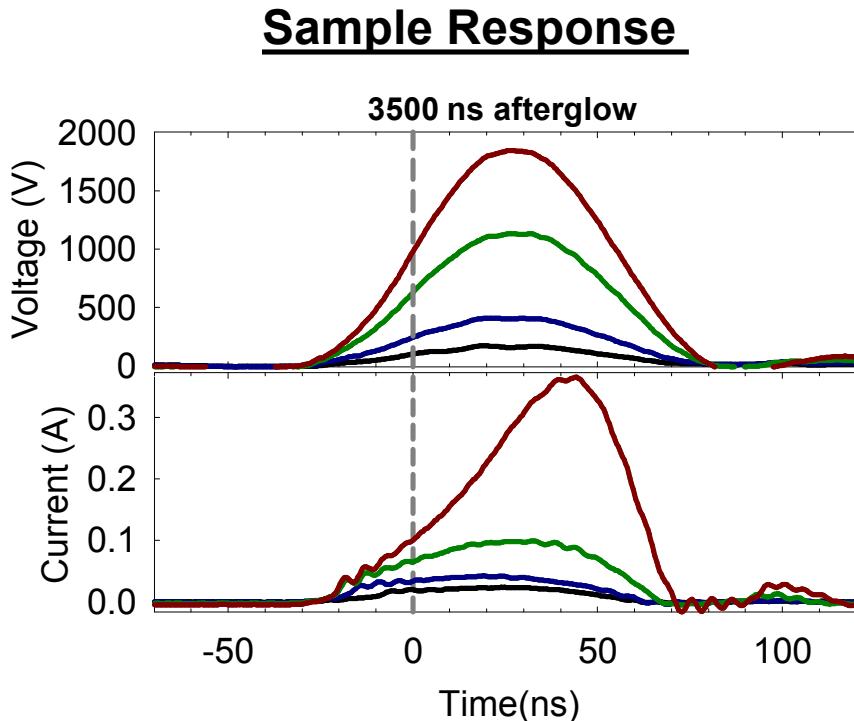
Key interrogation methods



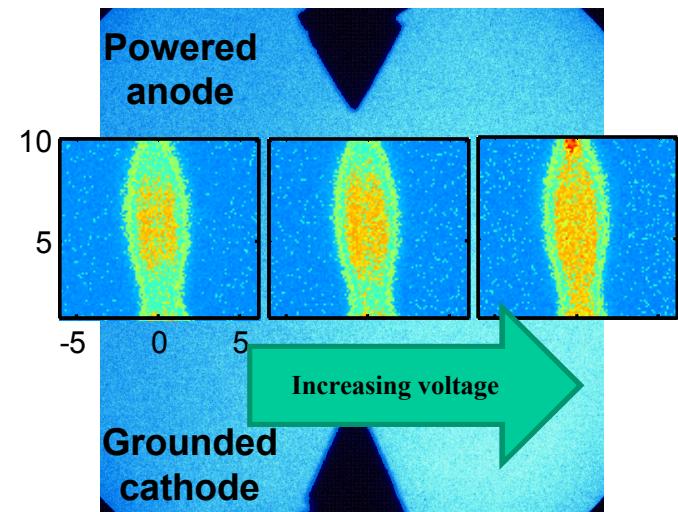
Setup enables good access and control of the plasma for LCIF calibration.

DOUBLE PULSE PLASMA SEPERATES GENERATION FROM MANIPULATION

- Need plasma with well controlled n_e , E/N to calibrate LCIF.
 - First voltage pulse dictates plasma density and distribution.
 - Second pulse drives current through afterglow.



Observed Filament



Structure of the filament does not change significantly as it is “heated” – at least initially.

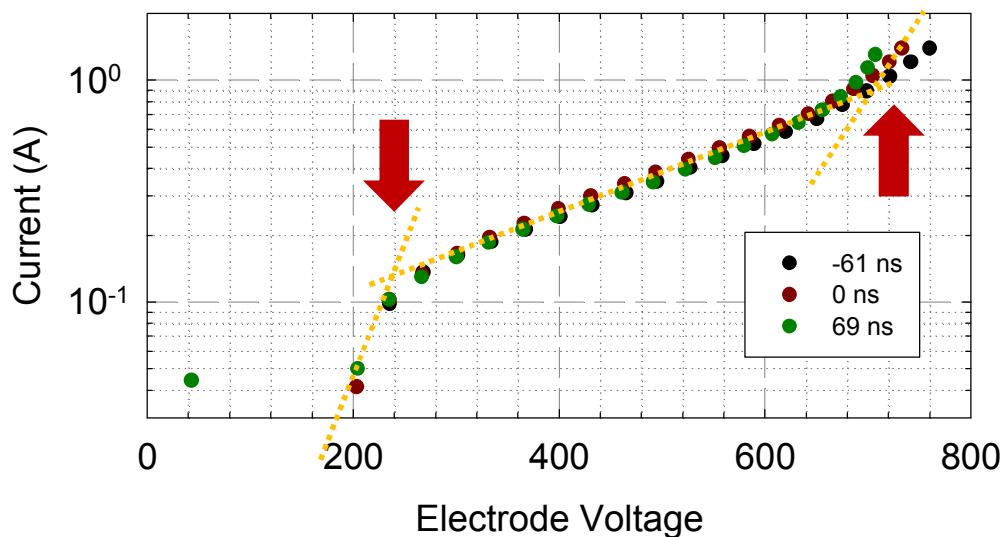


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UTILIZE CURRENT-VOLTAGE TRENDS TO BOUND PLASMA PARAMETERS

- Published drift data (Phelps) is used to bound E/N with heating voltage.
 - Knees in current correspond to knees in drift velocities.
 - Electron density remains roughly constant at lower E/N values.

Extracted Current



Anticipated plasma parameters

$$n_e = \frac{m_e}{e^2} \frac{v_m}{N} \frac{N}{E} \frac{I}{A}$$

V	I	J	E/N	Vd	ne
V	A	A/cm ²	Td	cm/s	e/cm ³
225	0.1	3.18	0.1	1.50E+05	1.32E+14
700	0.6	19.10	5	1.00E+06	1.19E+14

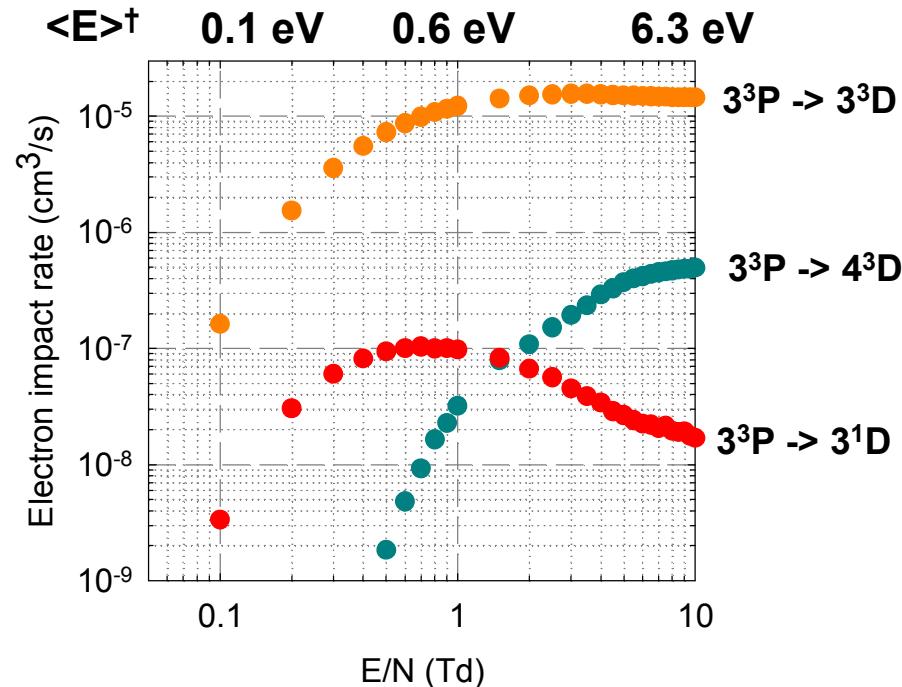
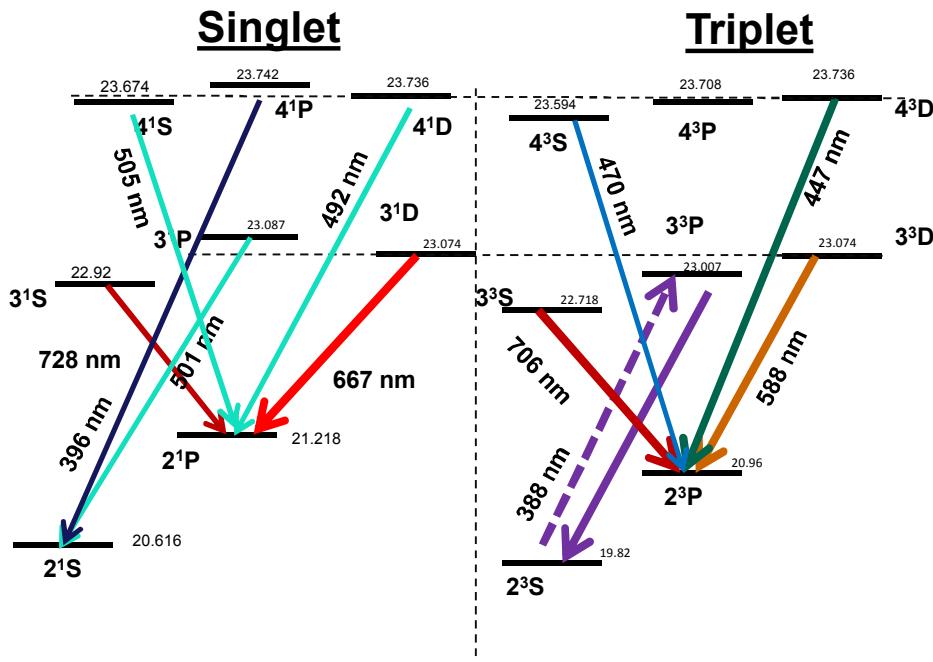
*Delay between first pulse and second pulse - Density
Magnitude of second voltage pulse – E/N*



LCIF EVOLVES THROUGH SEVERAL SPECTROSCOPIC PATHWAYS

- Transitions from the laser excited 3^3P state have different E/N scaling
 - Capitalize on unique characteristics to assess $\langle E \rangle$ and E/N.

Pathway and Rates



Identify pathways that yield good signals and are “easy to” evaluate.



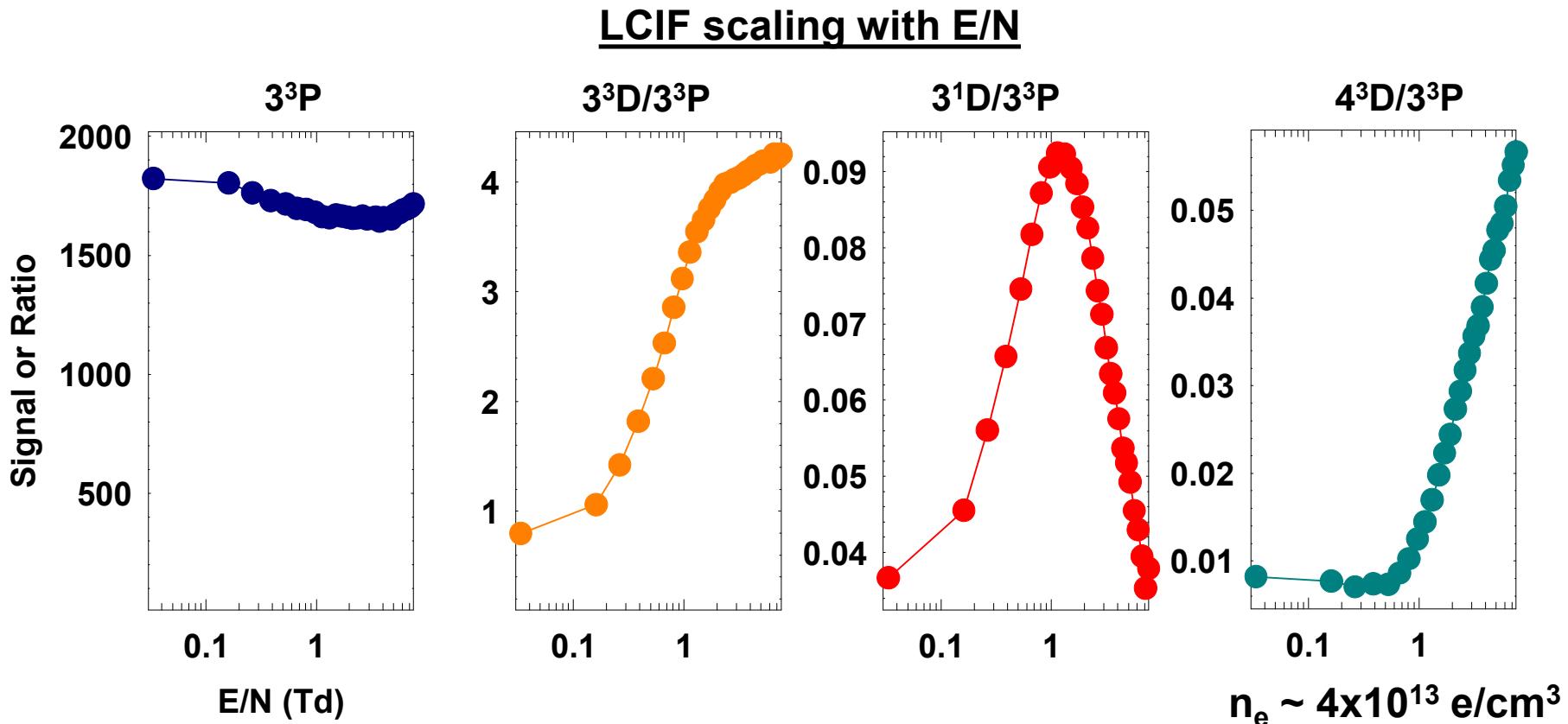
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* Cross sections: Yu. Ralchenko, Atomic Data and Nuclear Data Tables **94**, 603 (2008)

[†]Characteristic energy: Bolsig and LXCat.

RATIO OF LCIF LINES ARE UTILIZED TO IDENTIFY SCALING TRENDS

- E/N is tuned and LCIF is measured early during applied pulse.
 - Ratio of LCIF to laser excited 3^3P LIF is used to normalize signals.
 - LCIF acquired for ~ 10 ns, immediately after laser pulse.



Several LCIF transitions are observable!

RATIO OF LCIF LINES ARE UTILIZED TO IDENTIFY SCALING TRENDS

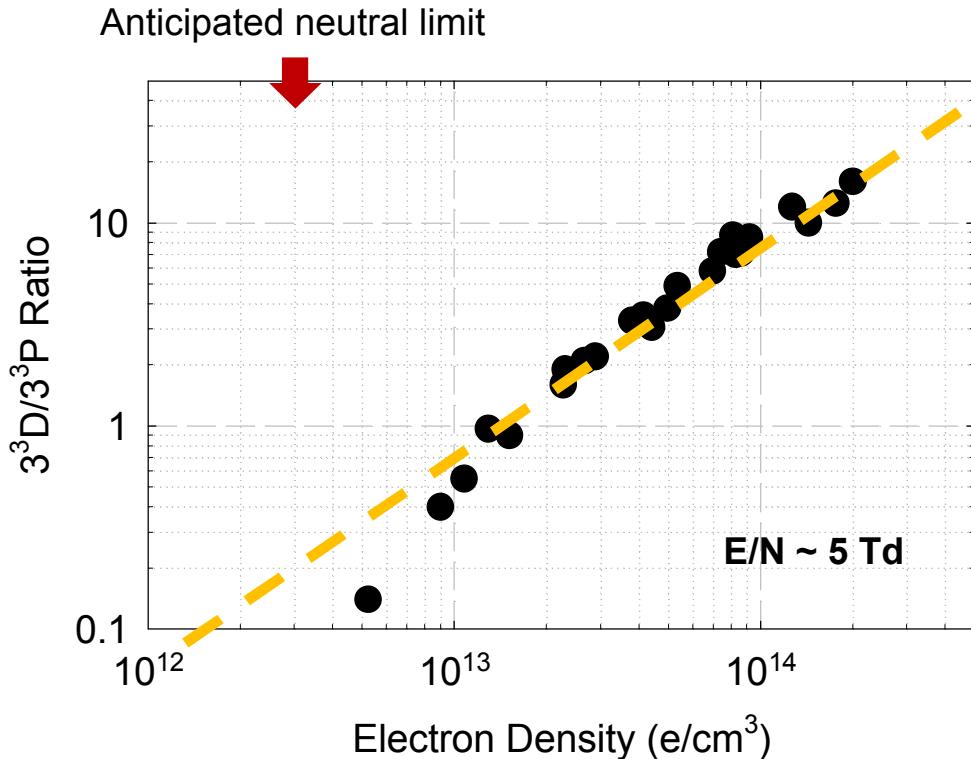
- Benchmark scaling of density dependence of 3^3D LCIF.
 - Utilize the E/N invariance of 3^3D LCIF above ~ 2 Td

Simplified Scaling

$$\Delta N_{3^3D} \sim K_{ij}^e n_e \times \Delta N_{3^3P}$$

$$\frac{\Delta N_{3^3D}}{N_{3^3P}} \sim K_{ij}^e n_e$$

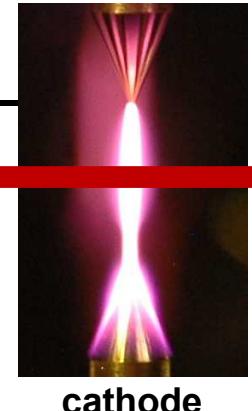
Observed Scaling of LCIF



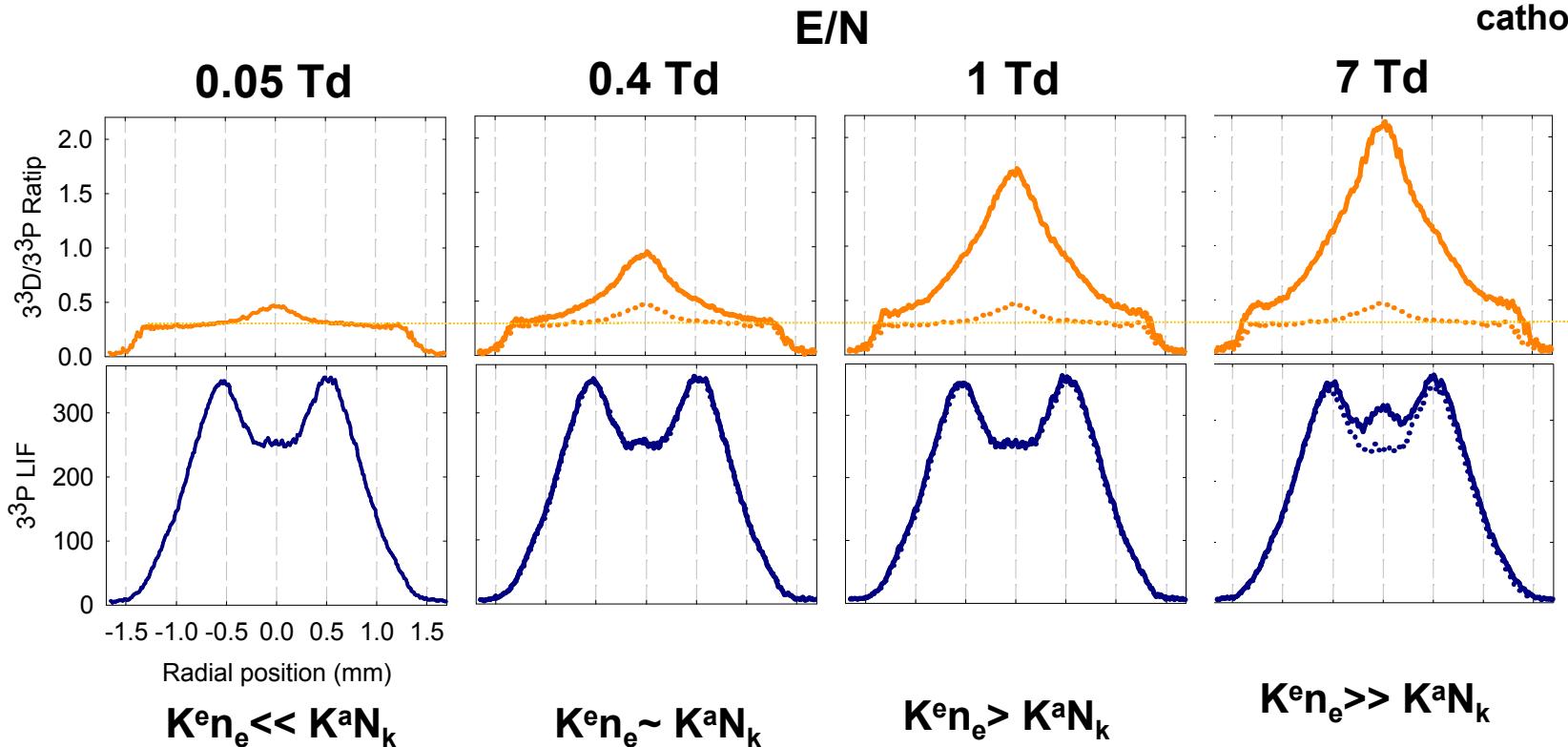
LCIF from 3³D scales linearly with electron density

DEMONSTRATION OF SPATIAL RESOLUTION PROVIDED BY LCIF

anode



- Radial structure observed for various E/N.
 - Initial peak electron density of $2 \times 10^{13} \text{ e/cm}^3$.
 - Measured $\sim 2 \mu\text{s}$ in afterglow of first voltage pulse.
 - Measured in first 20 ns of applied heating pulse.



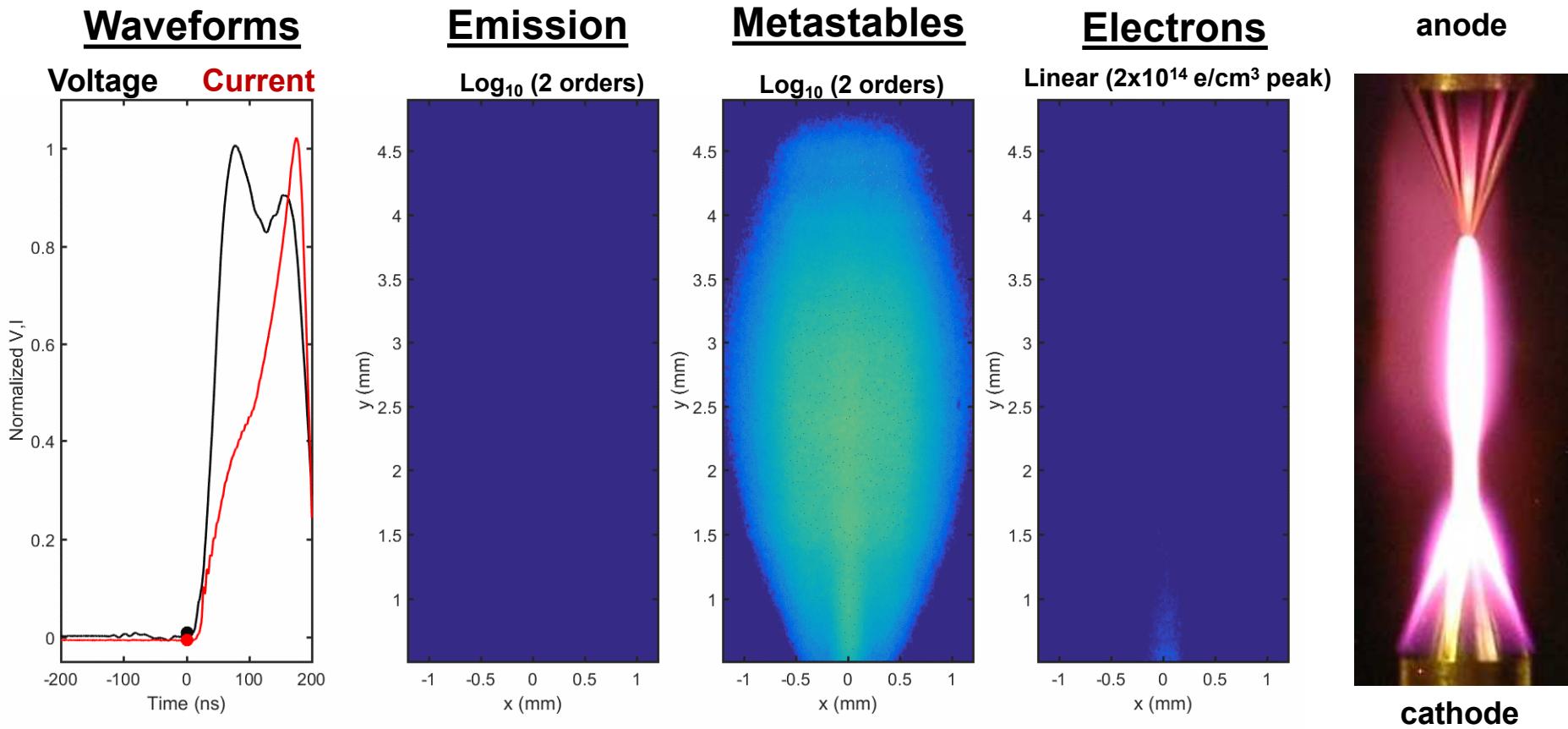
Observations consistent with initial assertions made



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DEMONSTRATION OF SPATIAL AND TEMPORAL RESOLUTION PROVIDED BY LCIF

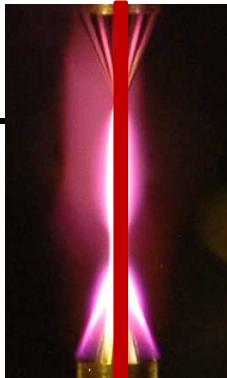
- Measured evolution of plasma formation during second pulse.
 - Initial peak electron density of $2 \times 10^{13} \text{ e/cm}^3$ (2 μs afterglow)



*LCIF provides two-dimensional
plasma structure*

UTILIZATION OF LCIF TO UNDERSTAND DISCHARGE EVOLUTION

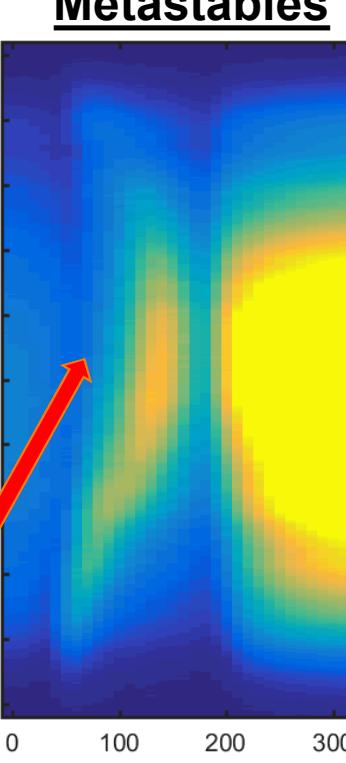
anode



cathode

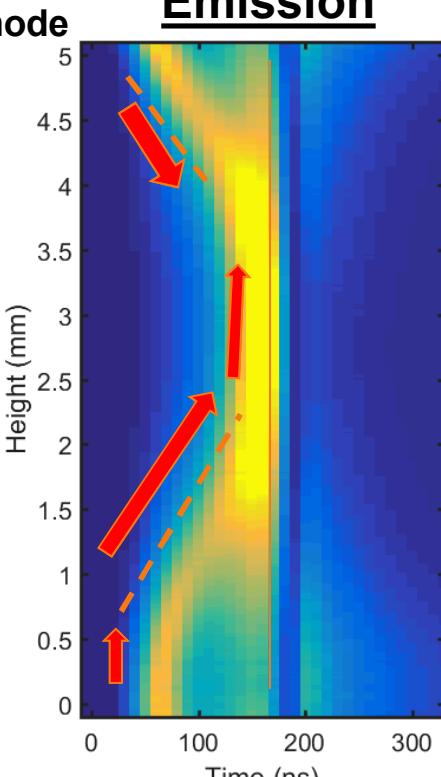
Electrons

Metastables

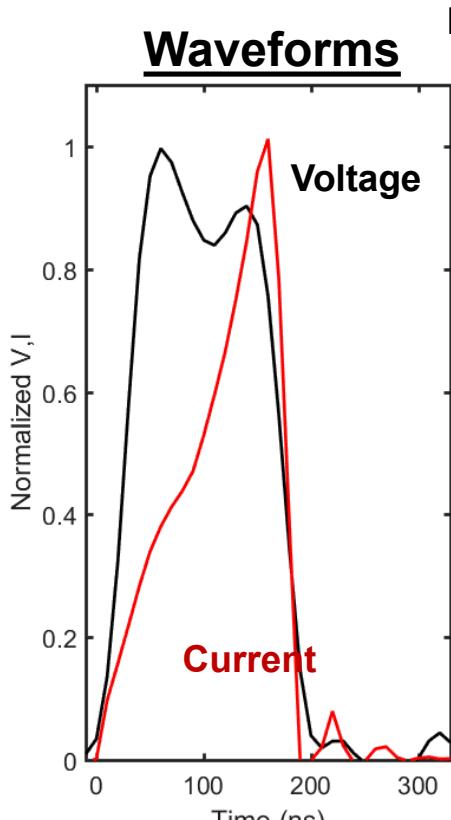


Emission

Powered anode



Waveforms



Current

Time (ns)

Grounded cathode

LCIF captures spatial and temporal evolution of plasma formation.



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CONCLUSIONS

- Ultra-fast LCIF shows promise for interrogating high pressure plasma systems.
 - Outlined pitfalls that might be encountered at higher-pressure systems.
 - Can be extended to other systems of interest (Ar, N,...).
- Preliminary LCIF results were published earlier this year
 - E.V. Barnat and A. Fierro, J.Phys.D:Appl. Phys. 50 (2017) 14LT01

Thank you for your attention!

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Center for
Predictive Control
of Plasma Kinetics:
Multi-Phase and
Bounded Systems



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UTILIZE CURRENT-VOLTAGE TRENDS TO BOUND PLASMA PARAMETERS

Published drift data

J. Pack et al. JOAP, 71 (11) p5363, 1992

