

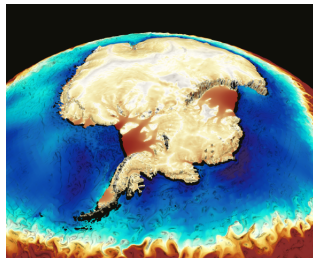
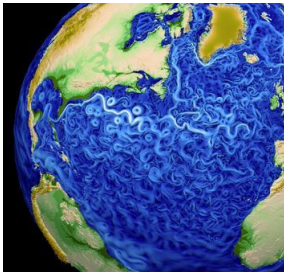


# A nonhydrostatic model for atmospheric motion in ACME-HOMME

Andrew J. Steyer, Sandia National Laboratories (SNL)

Work in collaboration with Mark Taylor (SNL) and Oksana Guba (SNL)

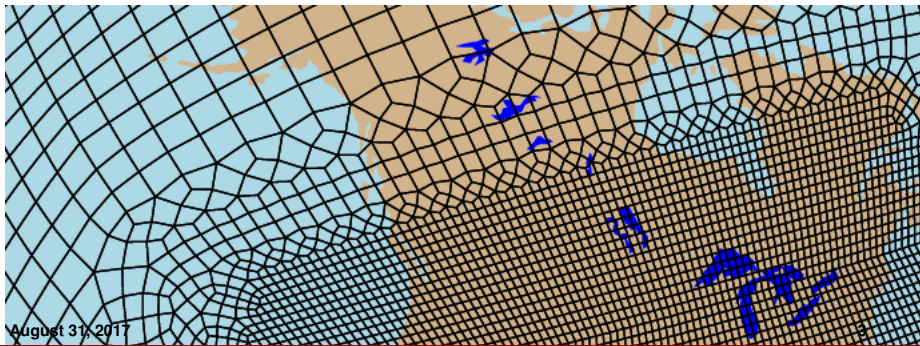
# Introduction



- ACME is an earth system model intended for scientific research and also regional prediction (temperature extremes, precipitation, etc).
- Focus is on the water cycle, biogeochemistry, and the cryosphere.
- Goal: produce a high-resolution earth system model running at 5 SYPD on DOE next-gen computers.
- ACME-HOMME is the atmospheric component of the ACME project.

# Why non-hydrostatic?

- The ACME model is the first Earth system model with variable resolution in all components.
- Variable resolution allows for regional refinement down into the nonhydrostatic regime.
- Allows for accurate regional climate modeling and cloud resolving simulations for process studies.



- Nonhydrostatic models must have  $w$  as a prognostic variable  $Dw/Dt \neq 0$ .
- We also add geopotential  $\phi = gz$  as a prognostic variable.
- We use a scaled potential temperature density  $\Theta = c_p^* \frac{\partial \pi}{\partial s} \theta$  rather than the real temperature  $T$  since  $D\theta/Dt \equiv 0$ .
- Exner pressure:  $T = \theta \left( \frac{p}{p_0} \right)^\kappa = \theta \Pi$
- Supports hydrostatic or nonhydrostatic simulation runs.
- Uses HOMME operators (2D mimetic-grad, curl, etc)
- This results in a formulation quite similar to that in (Laprise (1992)) and used in the GEM4.1 nonhydrostatic model.

# Theta-NH model

$$\left\{ \begin{array}{l} \frac{\partial \mathbf{u}}{\partial t} + (\nabla_s \times \mathbf{u} + 2\Omega) \times \mathbf{u} + \frac{1}{2} \nabla_s \mathbf{u}^2 + \dot{s} \frac{\partial \mathbf{u}}{\partial s} + \\ c_p^* \theta (\nabla_s \Pi - \frac{\partial \Pi}{\partial \kappa} \nabla_s \kappa) + \left( \frac{\partial p}{\partial s} / \frac{\partial \pi}{\partial s} \right) \nabla_s \phi = 0 \\ \frac{\partial w}{\partial t} + \mathbf{u} \cdot \nabla_s w + \dot{s} \frac{\partial w}{\partial s} + g - g \left( \frac{\partial p}{\partial s} / \frac{\partial \pi}{\partial s} \right) = 0 \\ \frac{\partial \phi}{\partial t} + \mathbf{u} \cdot \nabla_s \phi + \dot{s} \frac{\partial \phi}{\partial s} - gw = 0 \\ \frac{\partial \Theta}{\partial t} + \nabla_s \cdot (\Theta \mathbf{u}) + \frac{\partial}{\partial s} (\Theta \dot{s}) = 0 \\ \frac{\partial}{\partial t} \left( \frac{\partial \pi}{\partial s} \right) + \nabla_s \cdot (\pi_s \mathbf{u}) + \frac{\partial}{\partial s} (\pi_s \dot{s}) = 0 \\ \rho = - \frac{\partial \pi}{\partial s} / \frac{\partial \phi}{\partial s} \end{array} \right.$$

The energy conserved (in the continuum) is given by

$$H = \int_A \int_s \underbrace{\frac{1}{2}\pi_s \mathbf{u}^2 + \frac{1}{2}\pi_s w^2}_{\text{kinetic energy}} + \underbrace{\Theta\Pi + \phi_s p + p_{\text{top}}\phi_{\text{top}}}_{\text{internal energy}} + \underbrace{\pi_s \phi}_{\text{potential energy}} dA ds$$

In a mass or pressure based vertical coordinate  $s$  in non-hydrostatic mode this energy is actually the Hamiltonian and we can express

$$\dot{\mathbf{q}} = J(\mathbf{q}) \frac{\delta H}{\delta \mathbf{q}}, \quad \mathbf{q} = (\mathbf{u}^T, w, \phi, \Theta, \pi_s)^T$$

where  $J$  is skew-symmetric operator. Skew-symmetry (but not conservation) is lost in hydrostatic mode.

# Discrete energy non-conservation

- Homme operators are defined in 2 dimensions and our discretization gives us horizontal and vertical discrete integration by parts.
- Can show: energy is conserved if

$$\frac{\partial \pi}{\partial s} w \frac{\partial w}{\partial t} - \frac{w^2}{2} \nabla \cdot \left( \frac{\partial \pi}{\partial s} \mathbf{u} \right) + \left( g - g \left( \frac{\partial p}{\partial s} / \frac{\partial \pi}{\partial s} \right) \right) w = 0$$

Compare to:  $\frac{\partial w}{\partial t} + \mathbf{u} \cdot \nabla_s w + \dot{s} \frac{\partial w}{\partial s} + g - g \left( \frac{\partial p}{\partial s} / \frac{\partial \pi}{\partial s} \right) = 0.$

- Using 3D vector invariant form, discrete integration by parts is sufficient for energy conservation Dubos, T. and Tort, M. (2014).

# Theta-NH model (Hydrostatic mode)

Hydrostatic approximation implies  $\left(\frac{\partial p}{\partial s} / \frac{\partial \pi}{\partial s}\right) \equiv 1$

$$\left\{ \begin{array}{l} \frac{\partial w}{\partial t} + \mathbf{u} \cdot \nabla_s w + \dot{s} \frac{\partial w}{\partial s} + g - g \left( \frac{\partial p}{\partial s} / \frac{\partial \pi}{\partial s} \right) = 0 \\ \frac{\partial \phi}{\partial t} + \mathbf{u} \cdot \nabla_s \phi + \dot{s} \frac{\partial \phi}{\partial s} - gw = 0 \end{array} \right.$$

$$dw/dt \equiv 0, \quad d\phi/dt \equiv g \cdot \text{constant}$$

Number of non-trivial prognostic equations decreases by two and discrete energy is conserved.



- Due to dissipation and vertical remap energy will never be conserved exactly by the model.

$$\frac{1}{2}w^2\nabla \cdot \left( \mathbf{u} \frac{\partial \pi}{\partial s} \right) + \frac{\partial \pi}{\partial s} w \mathbf{u} \cdot \nabla w \neq 0$$

- To prevent spurious growth/decay the energy gained/lost is tracked and added back as heat.
- All errors including time-truncation become internal energy in the model.

$$\begin{cases} \frac{\partial w}{\partial t} + \mathbf{u} \cdot \nabla_s w + \dot{s} \frac{\partial w}{\partial s} + g - g \left( \frac{\partial p}{\partial s} / \frac{\partial \pi}{\partial s} \right) = 0 \\ \frac{\partial \phi}{\partial t} + \mathbf{u} \cdot \nabla_s \phi + \dot{s} \frac{\partial \phi}{\partial s} - gw = 0 \end{cases}$$

- For the JW Baroclinic test case (Jablonowski, C. and Williamson, D. (2006)) the maximum usable step size drops from around  $\Delta t = 600$  seconds in the hydrostatic model to  $\Delta t = 2$  seconds in the non-hydrostatic model.
- The time-step restriction is a result of the non-hydrostatic system supporting vertically propagating acoustic waves and the vertical scale being much smaller than the horizontal scale.

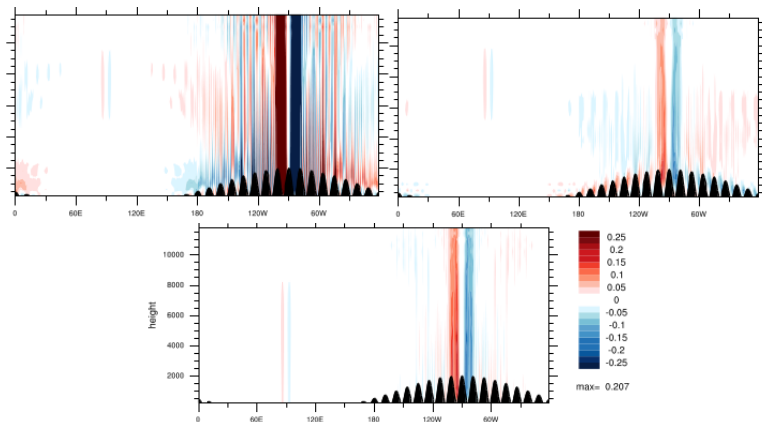
# HEVI splitting

Horizontally explicit vertically implicit splitting:

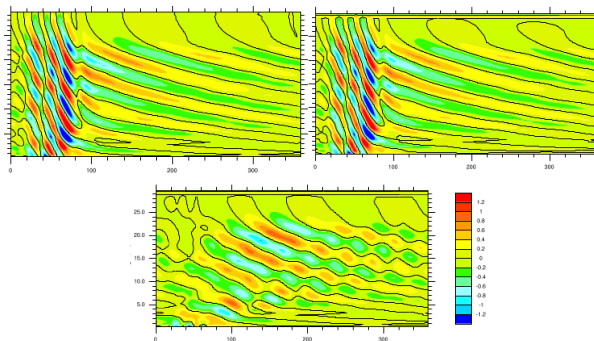
$$\left\{ \begin{array}{l} \frac{\partial \mathbf{u}}{\partial t} + (\nabla_s \times \mathbf{u} + 2\Omega) \times \mathbf{u} + \frac{1}{2} \nabla_s \mathbf{u}^2 + \dot{s} \frac{\partial \mathbf{u}}{\partial s} + \\ c_p^* \theta (\nabla_s \Pi - \frac{\partial \Pi}{\partial \kappa} \nabla_s \kappa) + \left( \frac{\partial p}{\partial s} / \frac{\partial \pi}{\partial s} \right) \nabla_s \phi = 0 \\ \frac{\partial w}{\partial t} + \mathbf{u} \cdot \nabla_s w + \dot{s} \frac{\partial w}{\partial s} + g - g \left( \frac{\partial p}{\partial s} / \frac{\partial \pi}{\partial s} \right) = 0 \\ \frac{\partial \phi}{\partial t} + \mathbf{u} \cdot \nabla_s \phi + \dot{s} \frac{\partial \phi}{\partial s} - gw = 0 \\ \frac{\partial \Theta}{\partial t} + \nabla_s \cdot (\Theta \mathbf{u}) + \frac{\partial}{\partial s} (\Theta \dot{s}) = 0 \\ \frac{\partial}{\partial t} \left( \frac{\partial \pi}{\partial s} \right) + \nabla_s \cdot (\pi_s \mathbf{u}) + \frac{\partial}{\partial s} (\pi_s \dot{s}) = 0 \\ \rho = - \frac{\partial \pi}{\partial s} / \frac{\partial \phi}{\partial s} \end{array} \right.$$

- DCMIP (Dynamical Core Model Intercomparison Project) 2012 established new non-hydrostatic test cases for atmospheric dycores.
- If interested: check out the Test Case Document by Ullrich et al on the DCMIP project website.
- Compare 3 models: PREQX (old HOMME dycore), THETA-H (theta model in hydrostatic mode), and THETA-NH (theta model in non-hydrostatic mode).
- Make use of "small planets" where non-hydrostatic effects are more apparent.

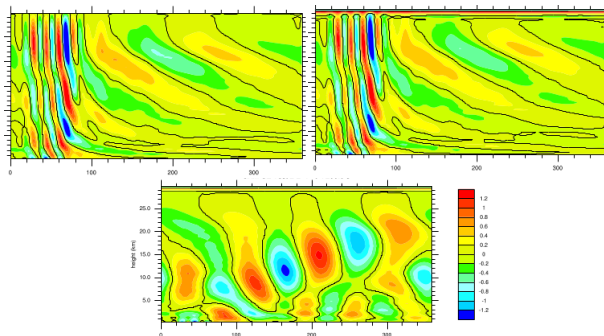
**Figure:** Atmosphere at rest, non-rotating planet, Schaefer-like mountain (Girard, C. et al. (2002)), horizontal velocity at the equator after 6 days with  $1^\circ$ , 30 vertical levels. PREQX (t. left), THETA-H (t. right), THETA-NA (bottom).



**Figure:** Temperature perturbation after 2 hours for flow over a Schaefer-like mountain without wind shear at the equator,  $1.5^\circ$  horizontal grid, 60 vertical levels, small planet  $\times 500$ . PREQX (t. left), THETA-H (t. right), THETA-NH (bottom)

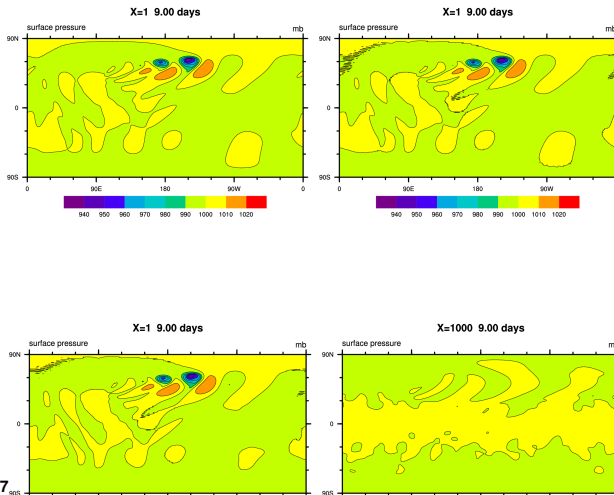


**Figure:** Temperature perturbation after 2 hours for flow over a Schaefer-like mountain with wind shear at the equator,  $1.5^\circ$  horizontal grid, 60 vertical levels, small planet  $\times 500$ . PREQX (t. left), THETA-H (t. right), THETA-NH (bottom)



# DCMIP 4 Baroclinic instability

**Figure:** Surface pressure at 1 degree. PREQX (t. left), THETA-H (t. right), THETA-NH (b. left), THETA-NH small planet x1000 (b. right)





# Fin

Questions?

# References I

- Dubos, T. and Tort, M. (2014). Equations of motion in non-Eulerian vertical coordinates: vector invariant form and quasi-Hamiltonian formulation. *Monthly Weather Review*, 142:3860–3880.
- Girard, C., Fuhrer, O., Lüthi, D., Leuenberger, D., and Schär, C. (2002). A new terrain-following vertical coordinate formulation for atmospheric prediction models. *Monthly Weather Review*, 130:2459–2480.
- Jablonowski, C. and Williamson, D. (2006). A baroclinic instability test case for atmospheric model dynamical cores. *Quarterly Journal of the Royal Meteorological Society*, 132:2943–2975.
- Laprise, R. (1992). The Euler equations of motion with hydrostatic pressure as an independent variable. *Monthly Weather Review*, 120:197–207.