



SAND2005-7220C

LAB 101

LOAD AND TARGET DESIGN  
DIODE AND PROBE ASSEMBLY  
WIRE ARRAY/LOAD FABRICATION



# Measurement and modeling of the implosion of wire arrays with seeded instabilities

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In collaboration with:

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and T. A. Mehlhorn**

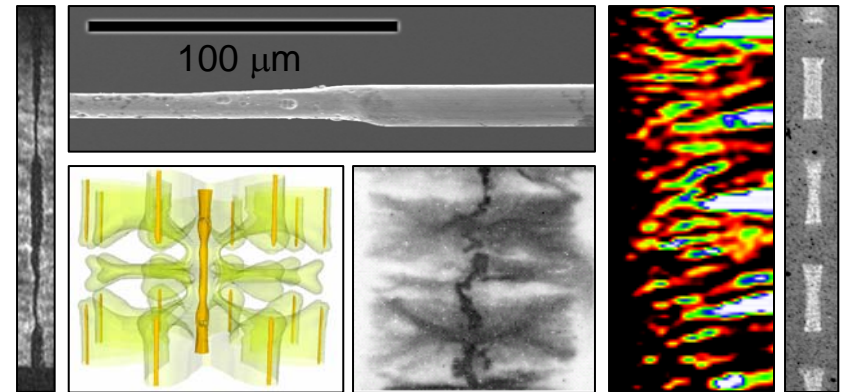
*Sandia National Laboratories*

**S. N. Bland, S. V. Lebedev,  
J. P. Chittenden, and S. C. Bott**

*Imperial College*

**B. V. Oliver**

*Mission Research Corp.*



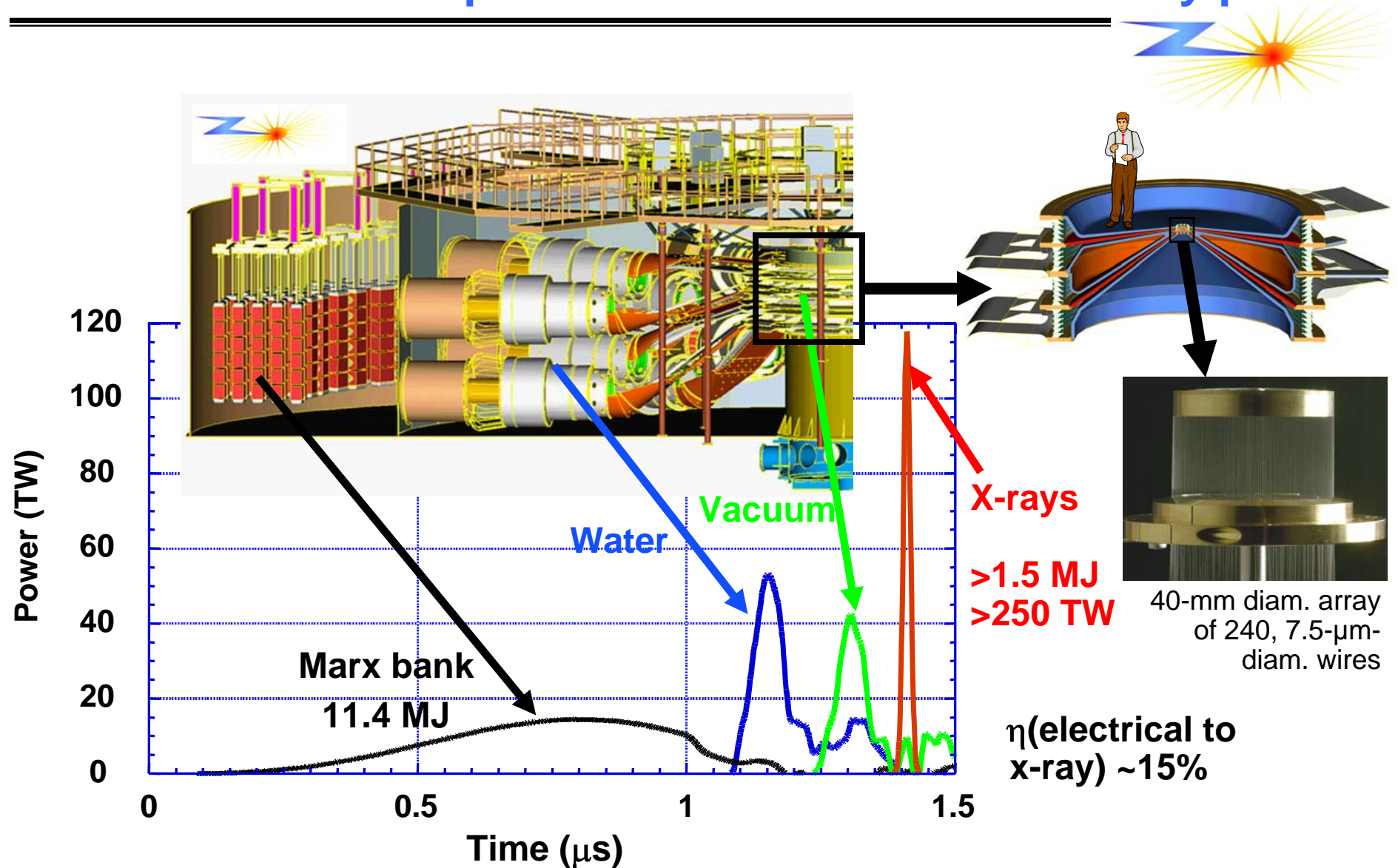
**47th Annual Meeting of the APS Division of Plasma Physics**  
**October 27, 2005**



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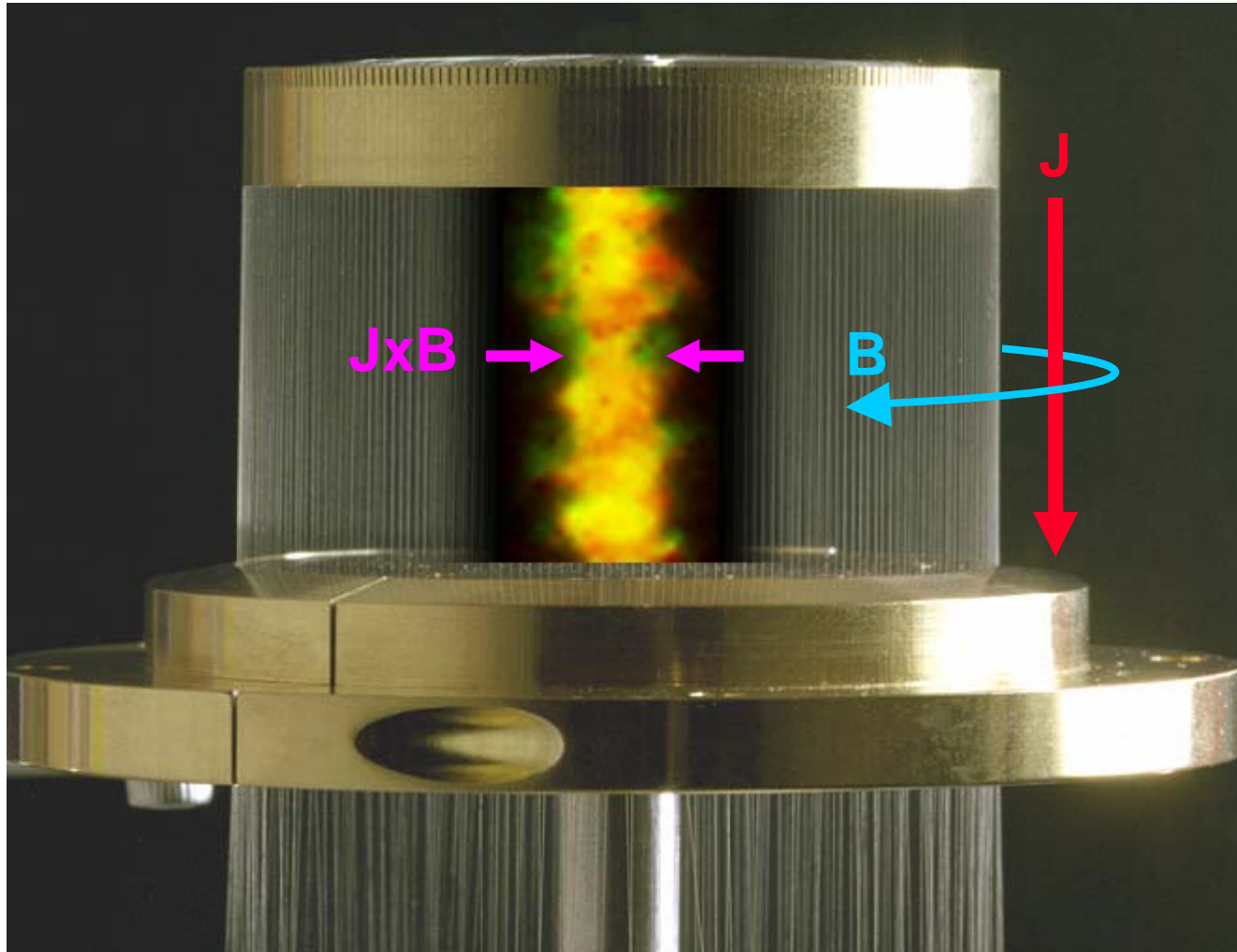
# Sandia's Z machine produces world-record soft x-ray powers



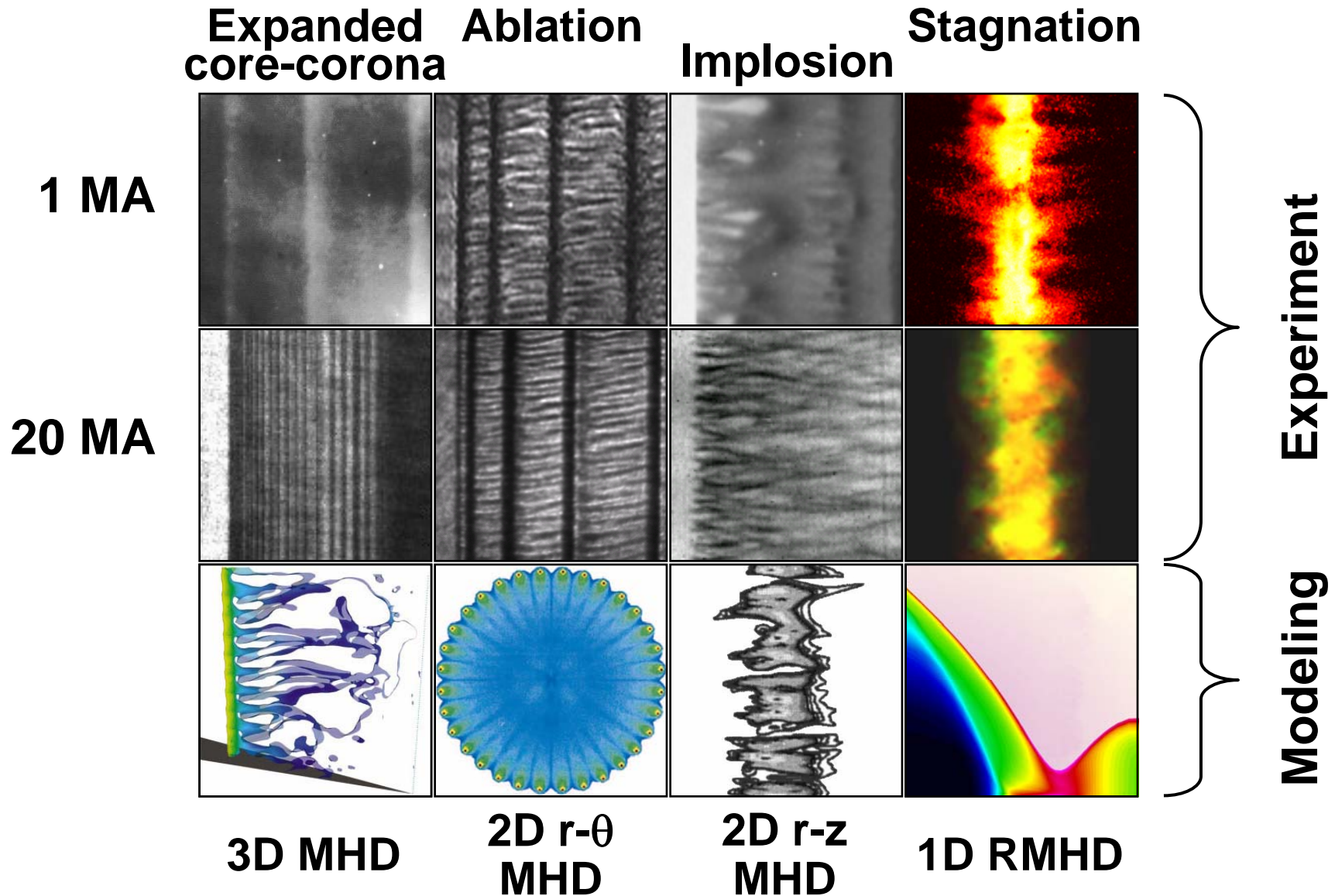
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# J x B pinches wire array into a dense, radiating plasma

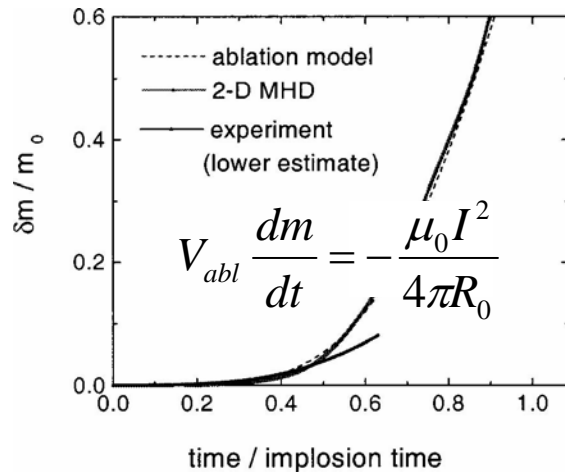


# Wire array z-pinches do not implode as a simple shell

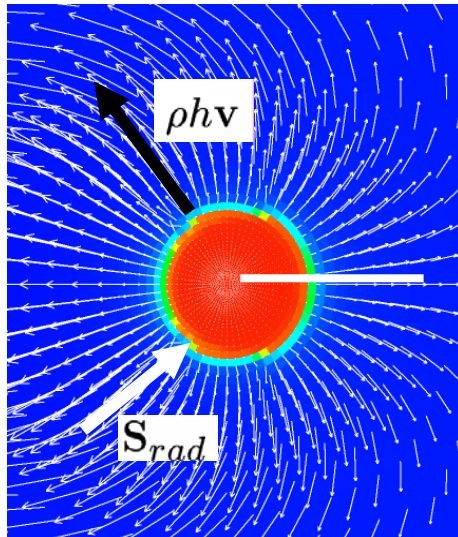




# Instabilities appear during the wire ablation phase

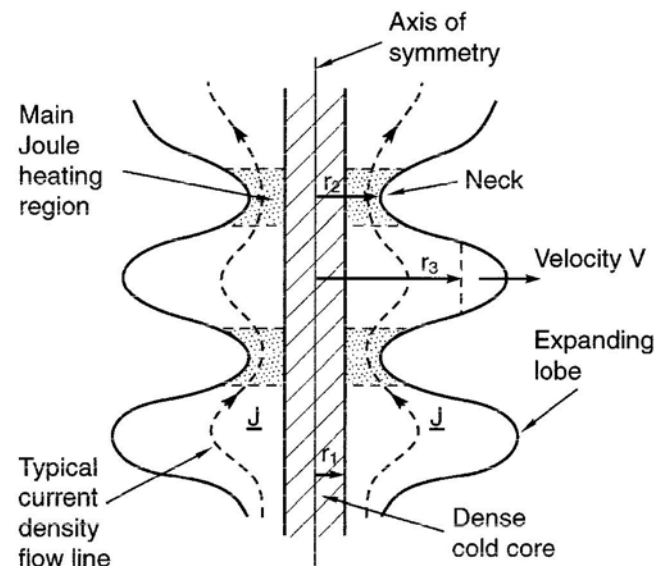


S. V. Lebedev *et al.*, Phys. Plasmas **8**, 3734 (2001).



E. P. Yu *et al.*, DPP05 C03.6

- Ablation rate is understood through an analytic rocket model and 2D  $r$ - $\theta$  MHD
- $m=0$  coronal instability and  $\mathbf{j} \times \mathbf{B}_{global}$  results in ablating flares—what determines the characteristic  $\lambda$ ?
  - MHD favored mode with  $ka \sim 1$
  - Electrothermal instability (Haines)



M. G. Haines, IEEE Trans. Plasma Sci. **30**, 588 (2002).



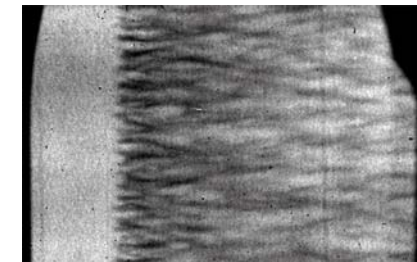
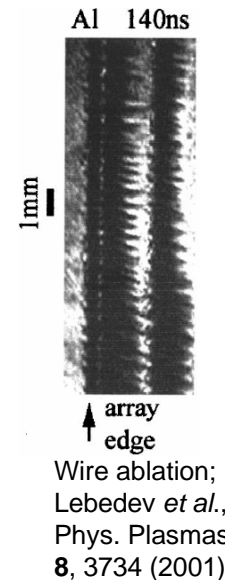
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# 3D structure impacts x-ray production in wire array z-pinches

- Instabilities in wire plasma lead to 3D structure at implosion

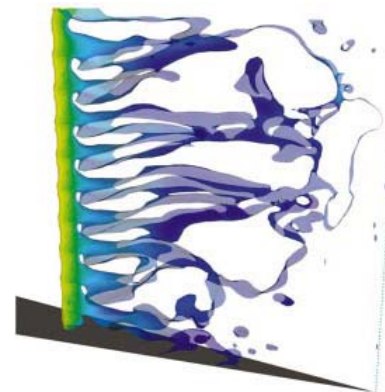
- Non-uniform wire ablation early in time leads to radial flare formation
- Magnetic Rayleigh-Taylor type implosion instability seeded by flares
- Trailing mass may impact current path
- 3D structure in stagnated pinch likely impacts x-ray generation and transport



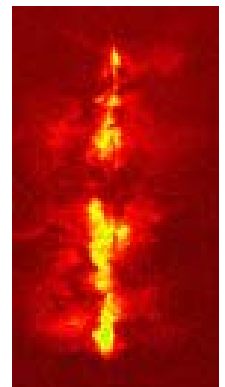
Stainless steel wire array on Z at 13.9 MA; D. B. Sinars *et al.*, Phys. Plasmas **12**, 056303 (2005).

- *Ad hoc* perturbations are seeded in 3D MHD simulations to reproduce phenomenological structure

- C. J. Garasi *et al.*, Phys. Plasmas **11**, 2729 (2004).
- J. P. Chittenden *et al.*, Plasma Phys. Control. Fusion **46**, B457 (2004).

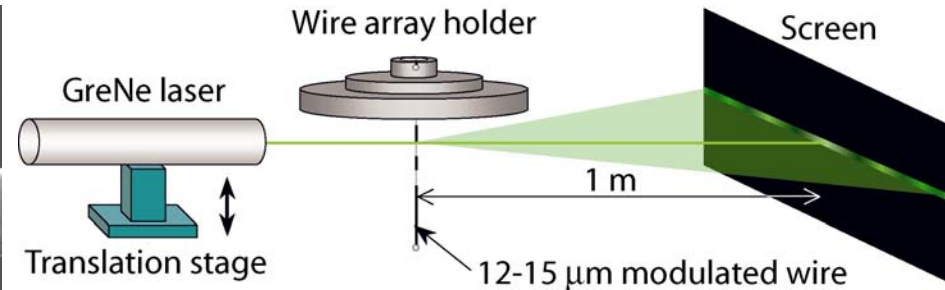
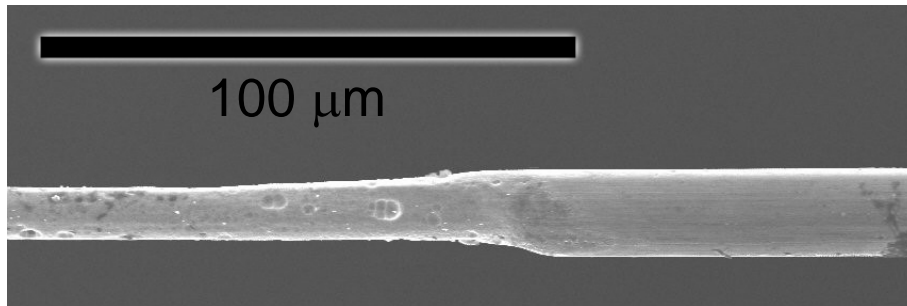


Al/Ni/Ti wire array on Z near peak x-ray power



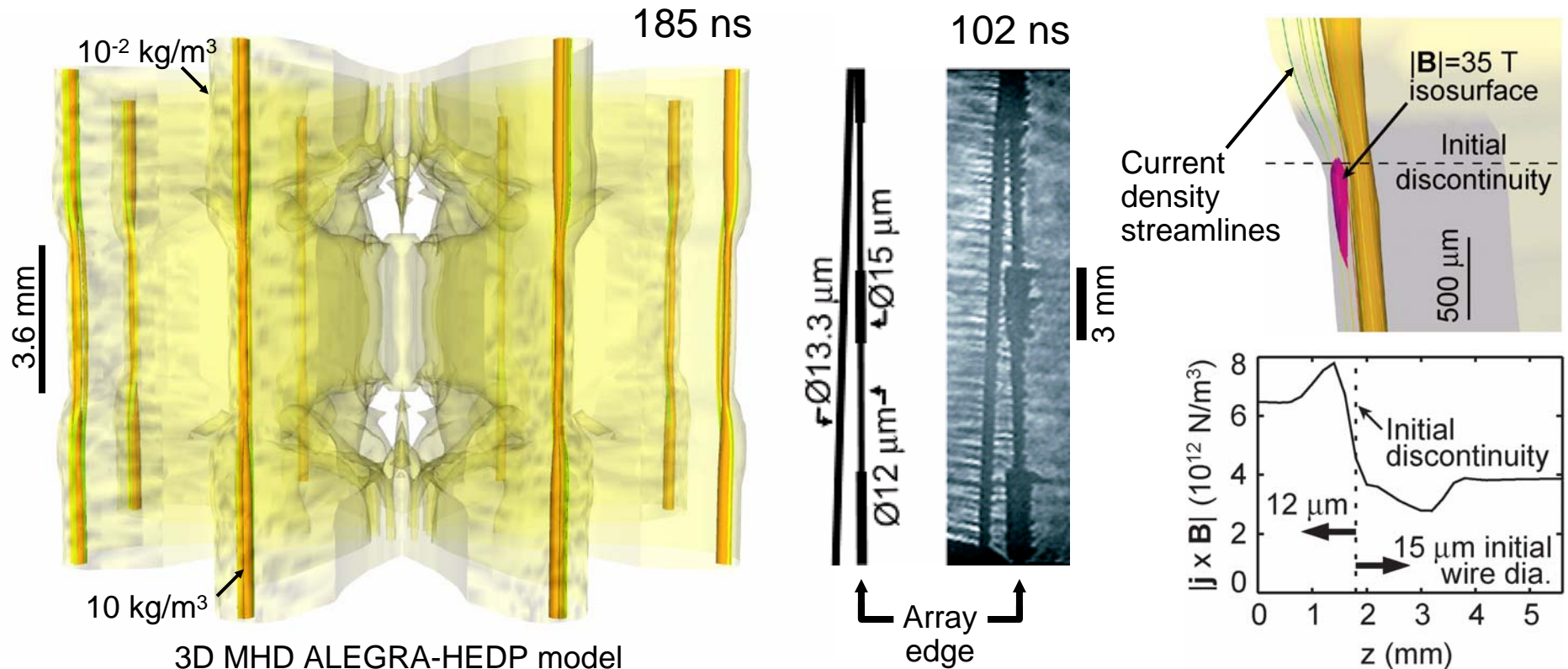
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# Mass perturbations are seeded by modulating wire radius



- Extruded  $\text{Ø}15\text{ }\mu\text{m}$  Al 5056 wires are chemically etched at MPCL/Sandia to have 20% modulation in radius with controlled axial wavelength
  - B. Jones *et al.*, Rev. Sci. Instrum. **75**, 5030 (2004).
- Aligned in wire arrays to  $\pm 100\text{ }\mu\text{m}$  with laser diffraction
- Allows seeded perturbations in z-pinch mass/length
  - Investigate radial flare formation
  - Implosion RT-type instability growth studies
  - X-ray pulse shaping by tailoring the profile of the imploding mass
- Seeded perturbations have been employed in laser plasmas
  - e.g. H. F. Robey *et al.*, Phys. Plasmas **8**, 2446 (2001).

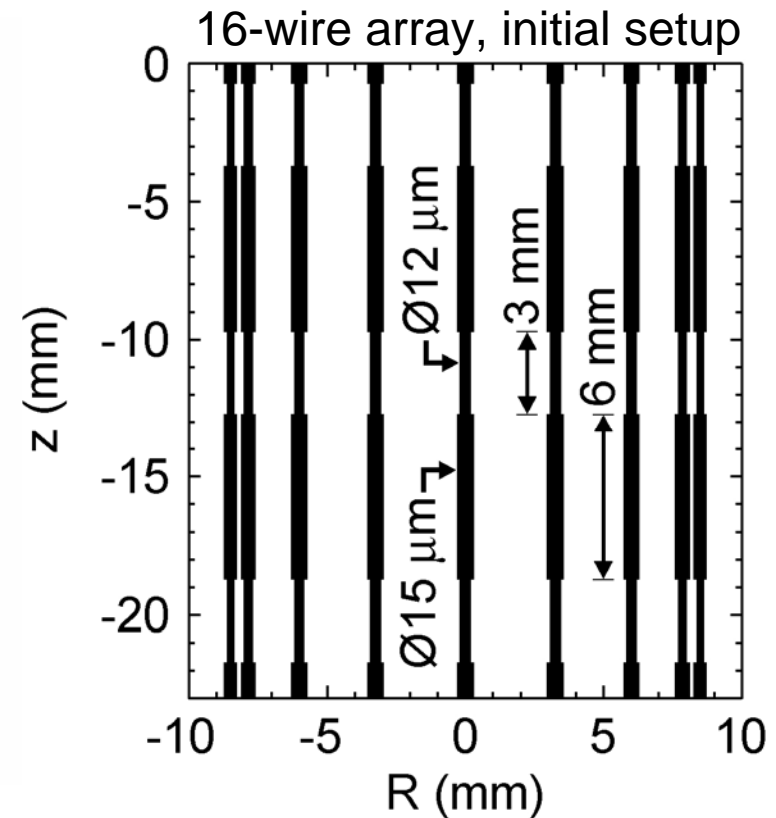
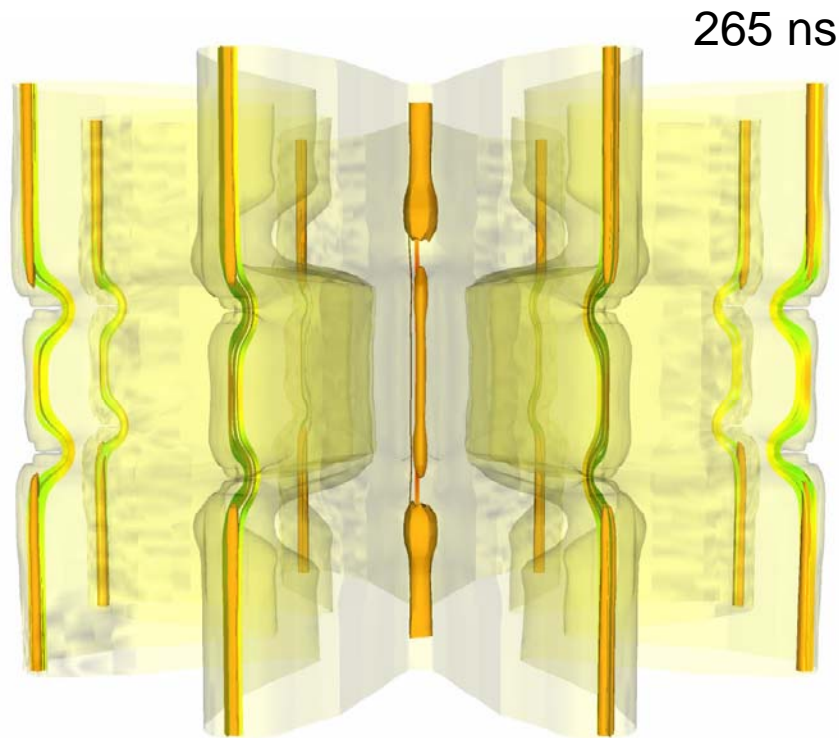
# Coronal modulation enhances $|j \times B|$ at wire discontinuities



- Enhanced magnetic field in thinner core regions constricts coronal plasma, leaving imprint of modulated structure
- Non-uniform current path near discontinuities in coronal plasma radius causes local magnetic field enhancement

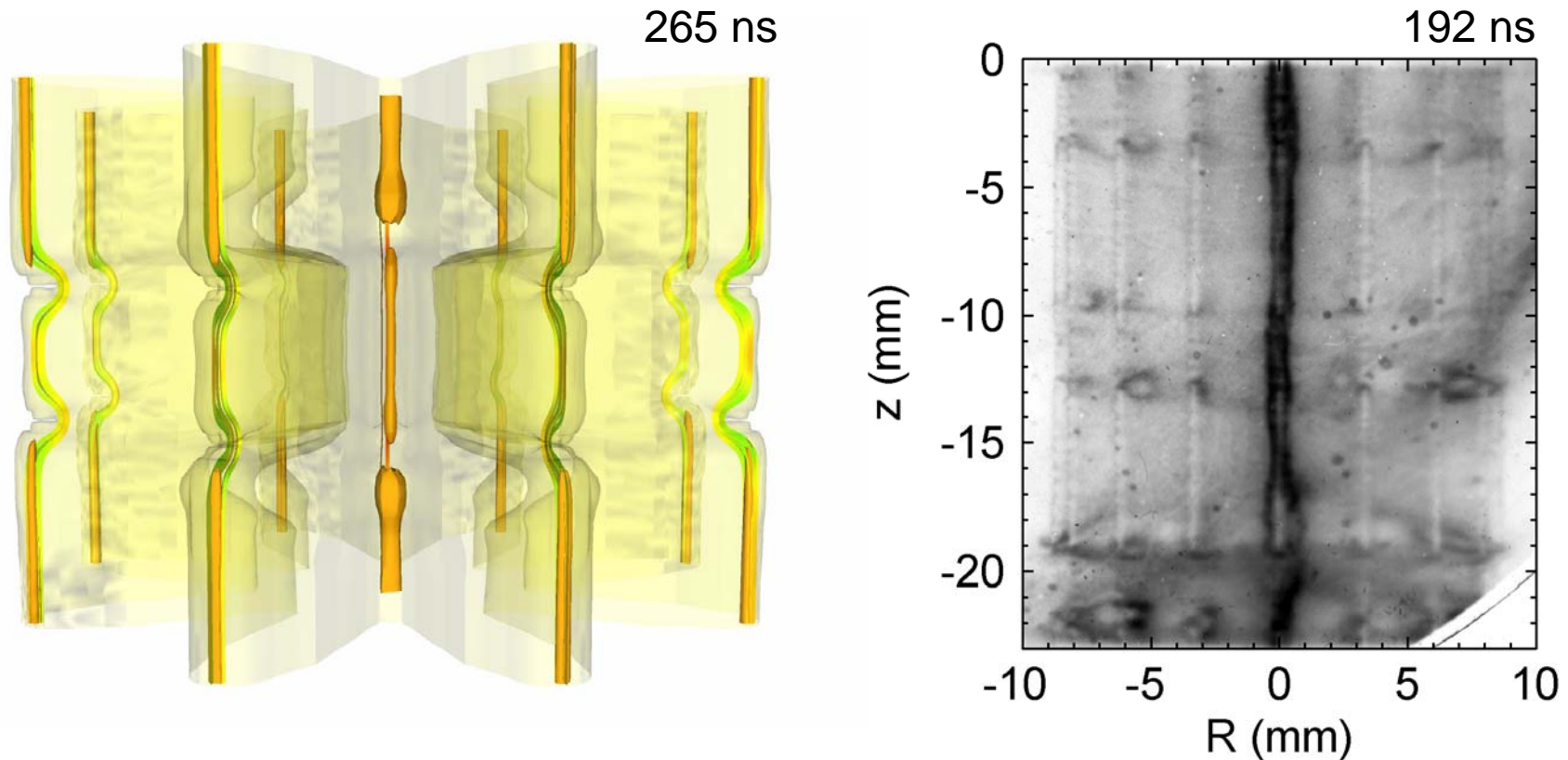


# Magnetic bubbles form at the discontinuities in wire radius



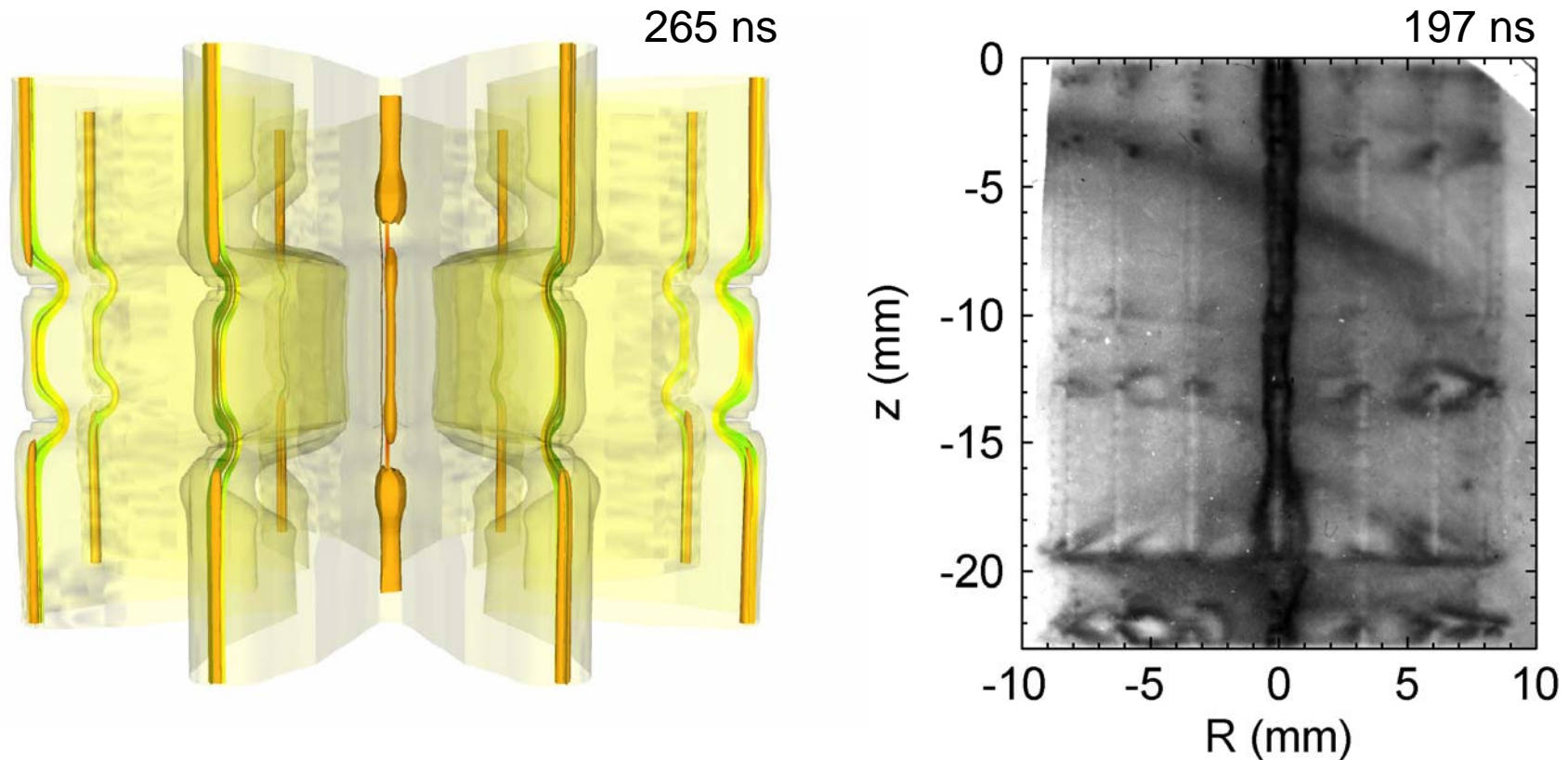
- Local  $|\mathbf{j} \times \mathbf{B}|$  enhancement at wire radius discontinuities enhances ablation, breaks wire cores, and initiates bubble implosion at discontinuities
- Bubble formation seen in both MHD model and experiment

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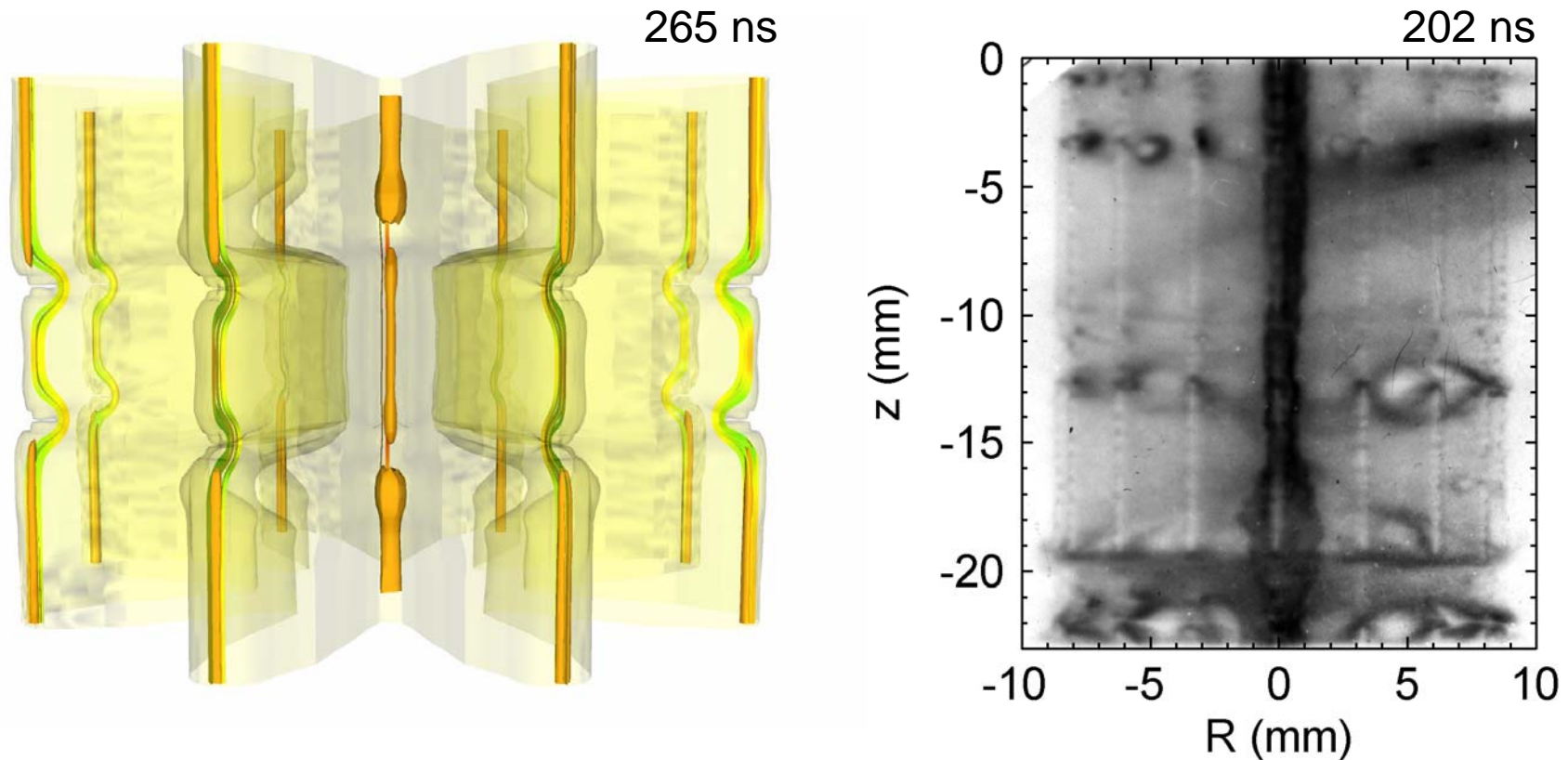
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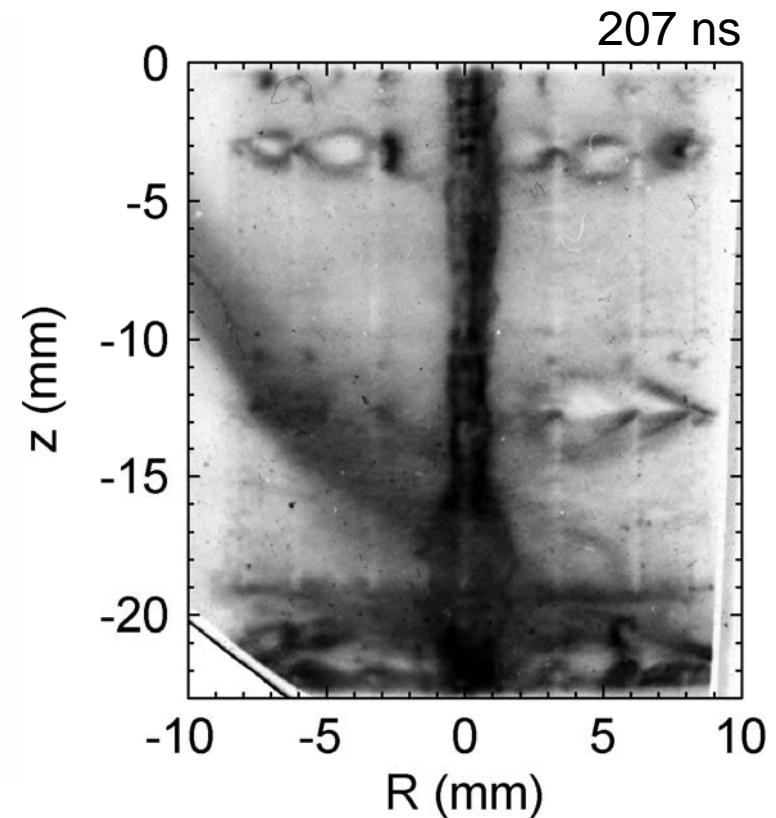
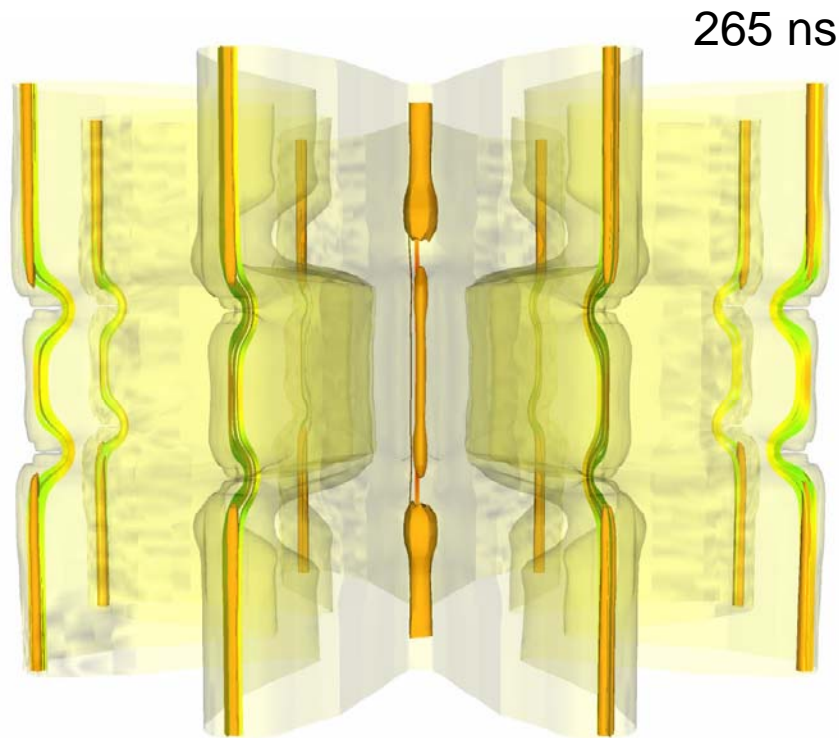
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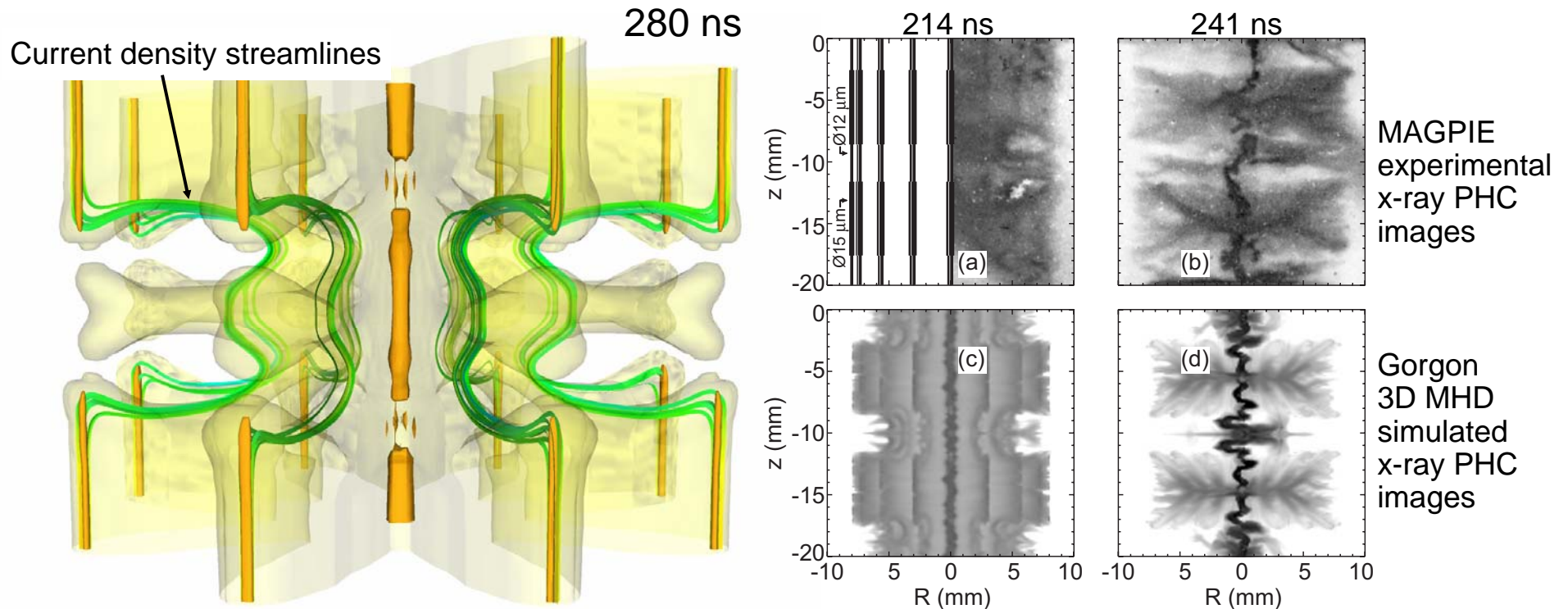


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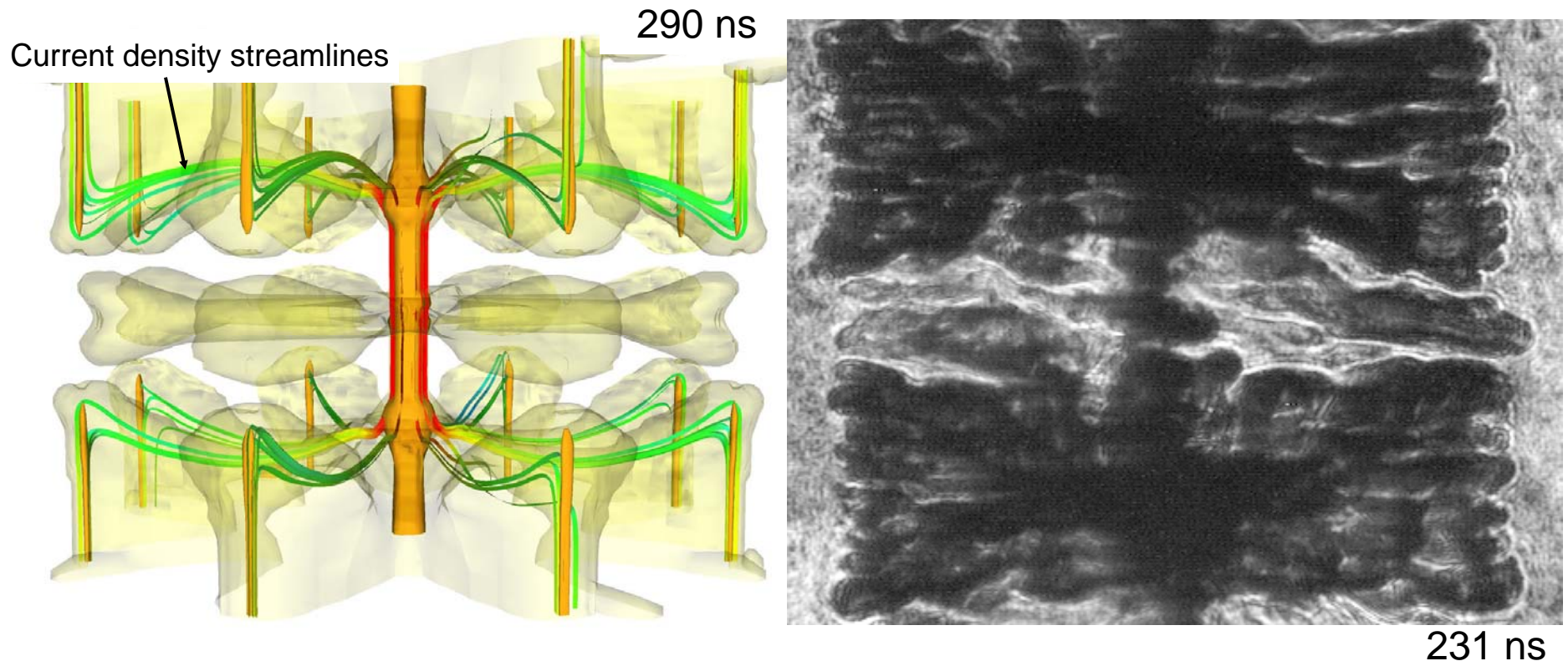
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- Bubble formation seen in both MHD model and experiment

# Correlated gaps form as magnetic bubbles implode



- Gorgon MHD model includes short-wavelength perturbations
- Same physics may explain trailing spikes in standard array
- Measurement of magnetic field in the gaps is needed to determine the actual current path

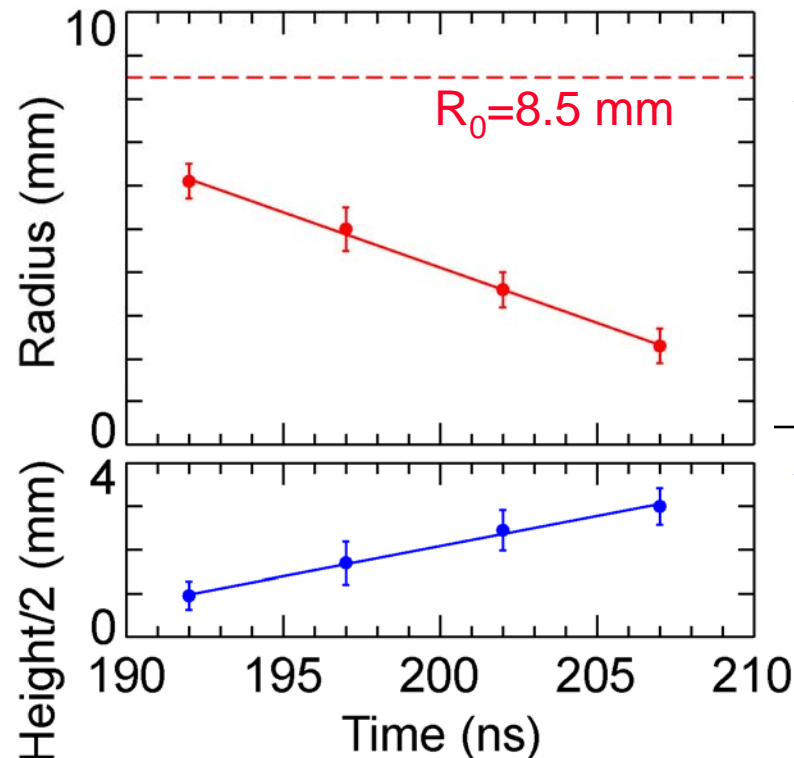
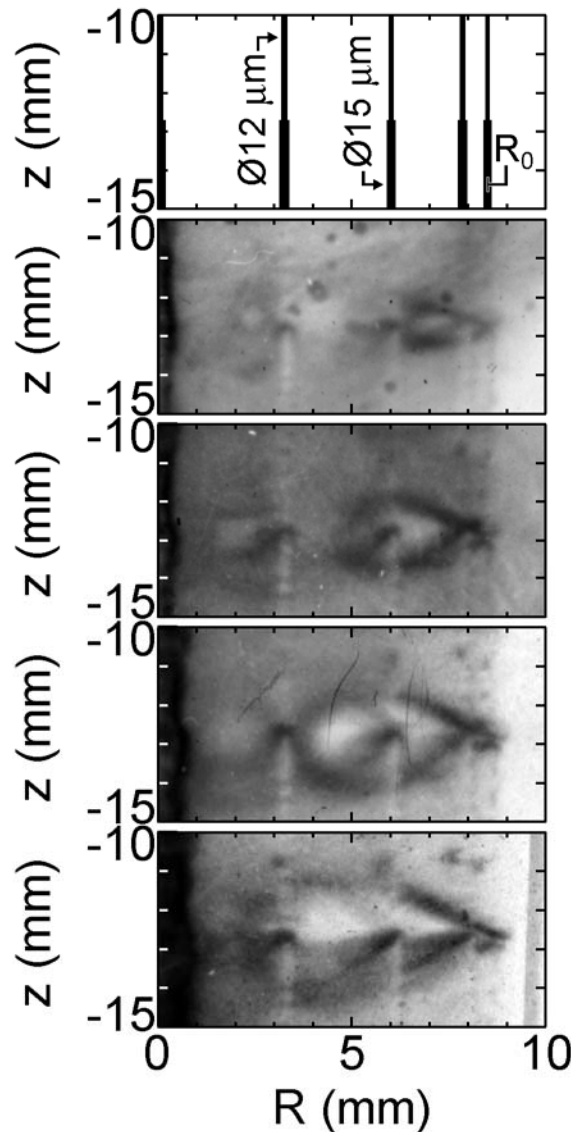
# Correlated gaps form as magnetic bubbles implode



- Same physics likely explains trailing spikes in standard array
  - Standard wire array shows similar ~2 mm bubble size on MAGPIE
- Measurement of magnetic field in the gaps is needed to determine the actual current path



# Bubbles are observed to implode rapidly—driven by $\mathbf{j} \times \mathbf{B}$ ?



Fit results –

$$v = 25.5 \pm 3.6 \text{ cm}/\mu\text{s}$$

$$R = R_0 \text{ at } t = 182.8 \pm 1.9 \text{ ns}$$

$$R = 1 \text{ mm at } t = 212.2 \pm 3.4 \text{ ns}$$

$$v = 13.9 \pm 3.4 \text{ cm}/\mu\text{s}$$

$$H = 0 \text{ at } t = 185.0 \pm 2.7 \text{ ns}$$

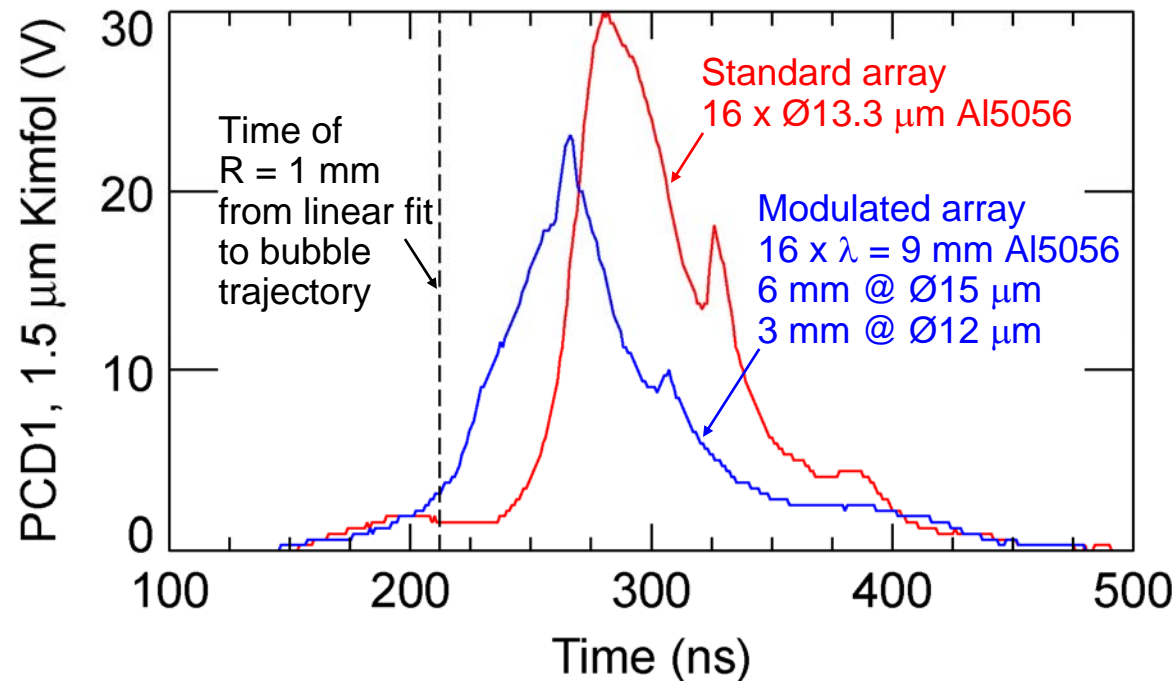
- 20-25  $\text{cm}/\mu\text{s}$  bubble implosion velocity
- Bubble blown from a point, rather than expanding spherically like a blast wave
- $\mathbf{j} \times \mathbf{B}$  would accelerate bubble, but snowplow of prefill decelerates it

$$F = \frac{d}{dt}(mv) = \dot{m}v + m\dot{v} = \frac{\mu_0 I^2}{4\pi NR}$$



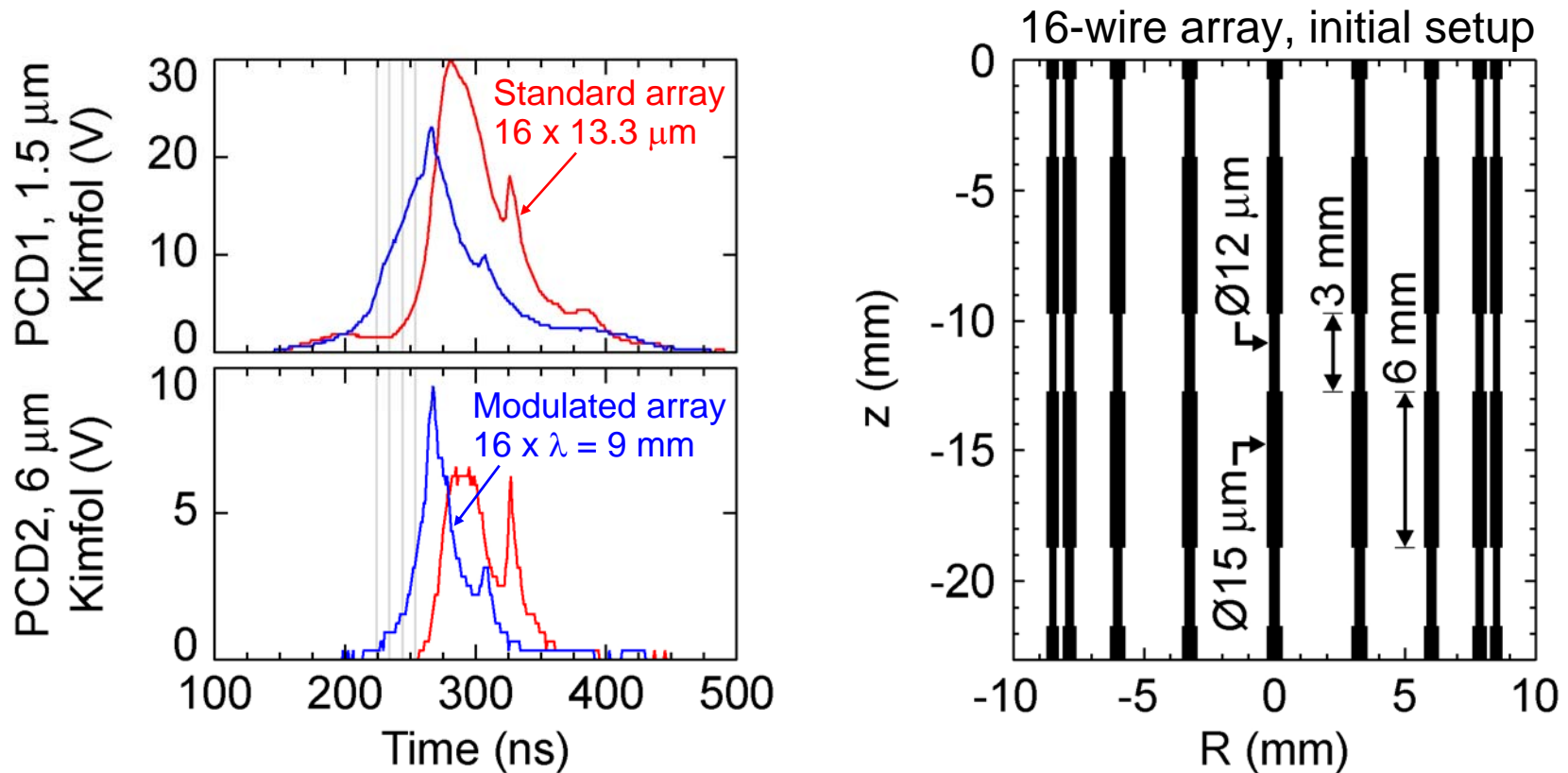


# Bubble collision on axis determines start of x-ray rise



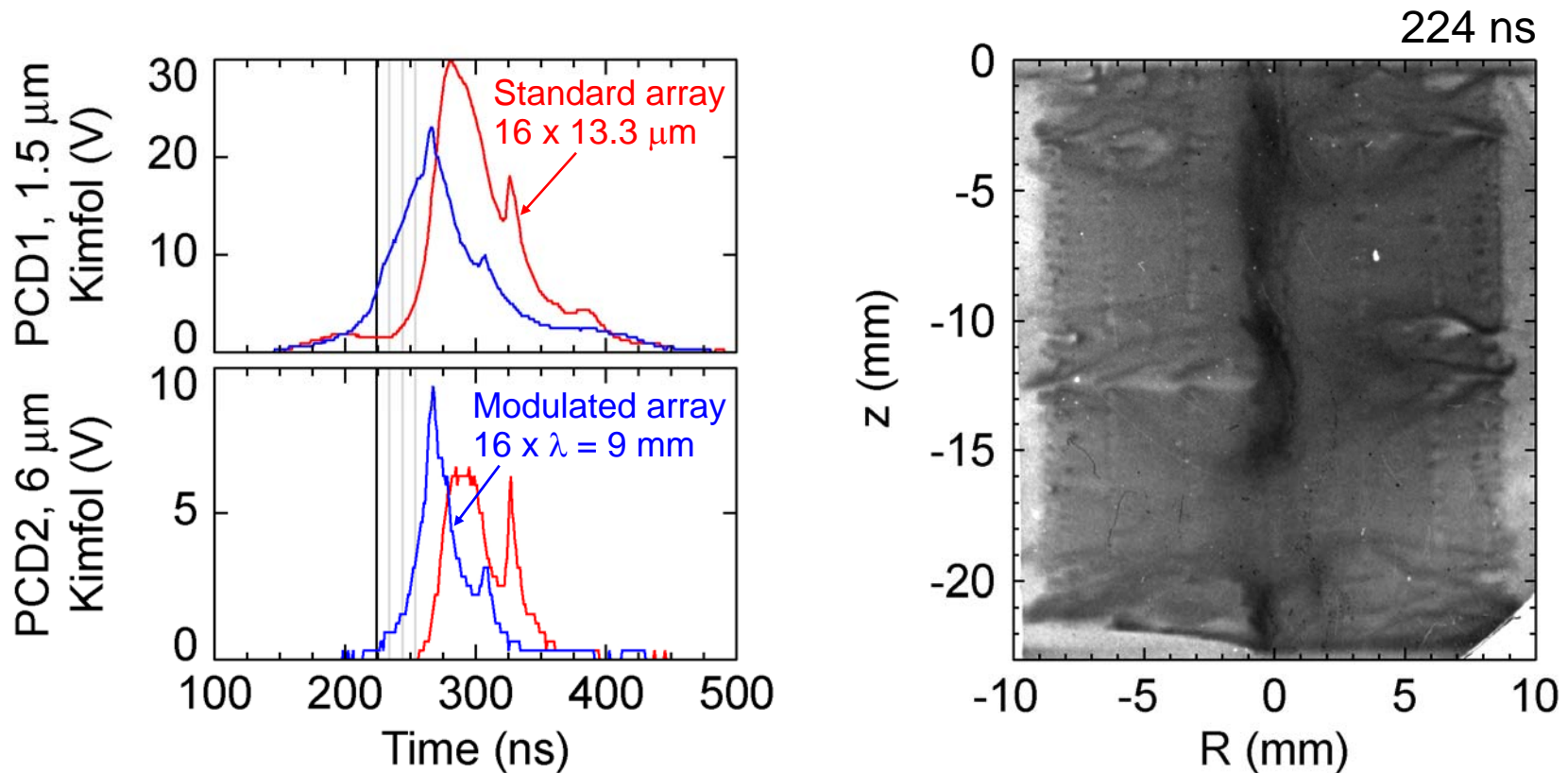
- Comparison of modulated with standard array of same mass
- Similar initial rise of precursor radiation
- Start of main x-ray rise corresponds to collision of imploding bubbles with the precursor (see also V. V. Ivanov, BP1.121)
- X-ray pulse shaping by tailoring the arrival of mass on axis

# X-ray pulse shape determined by assembly of mass on axis



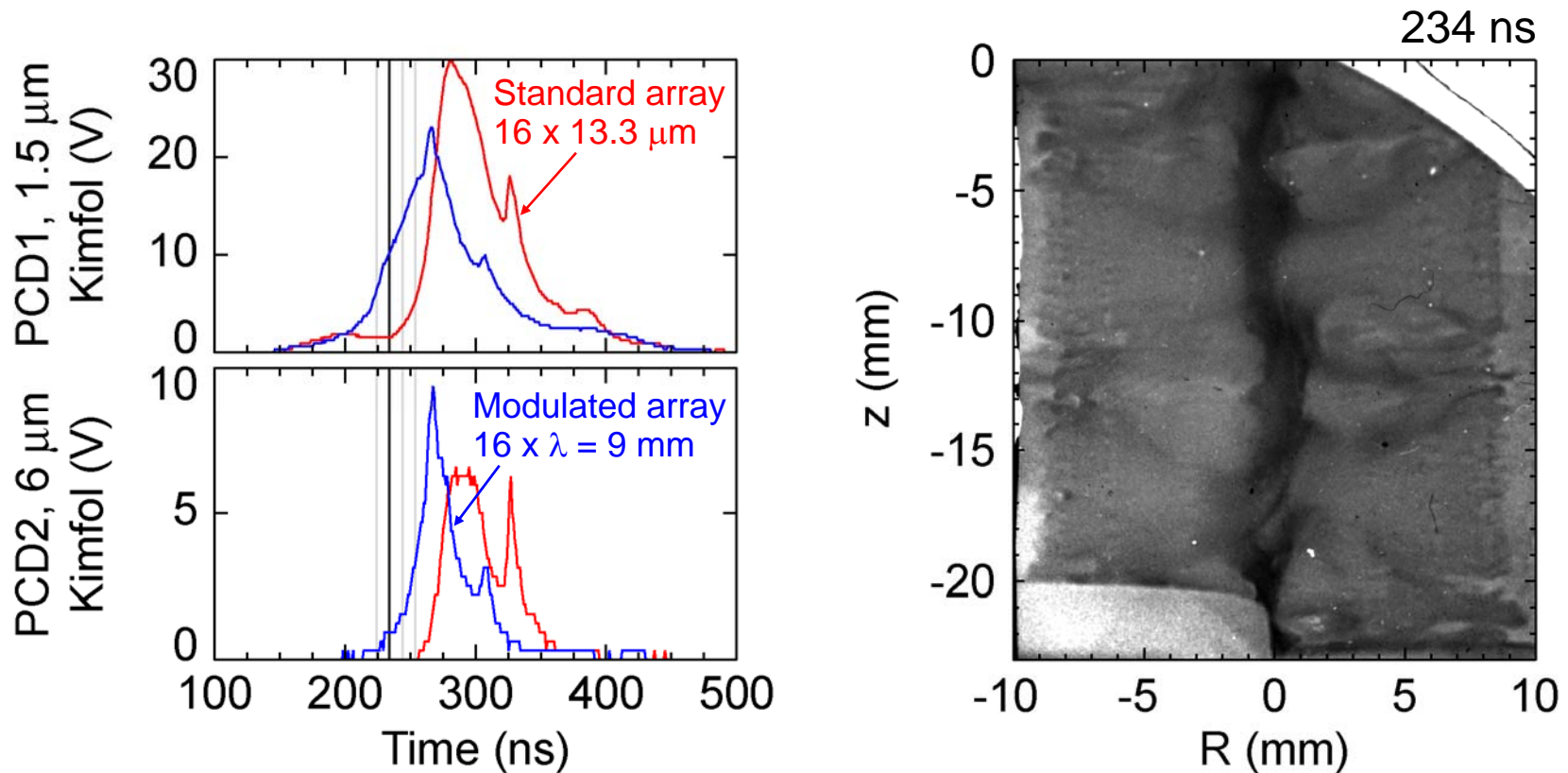
- Peak power is similar to standard wire array
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- Bubbles from discontinuities zipper on axis into a dense pinch—likely similar mechanism in standard array

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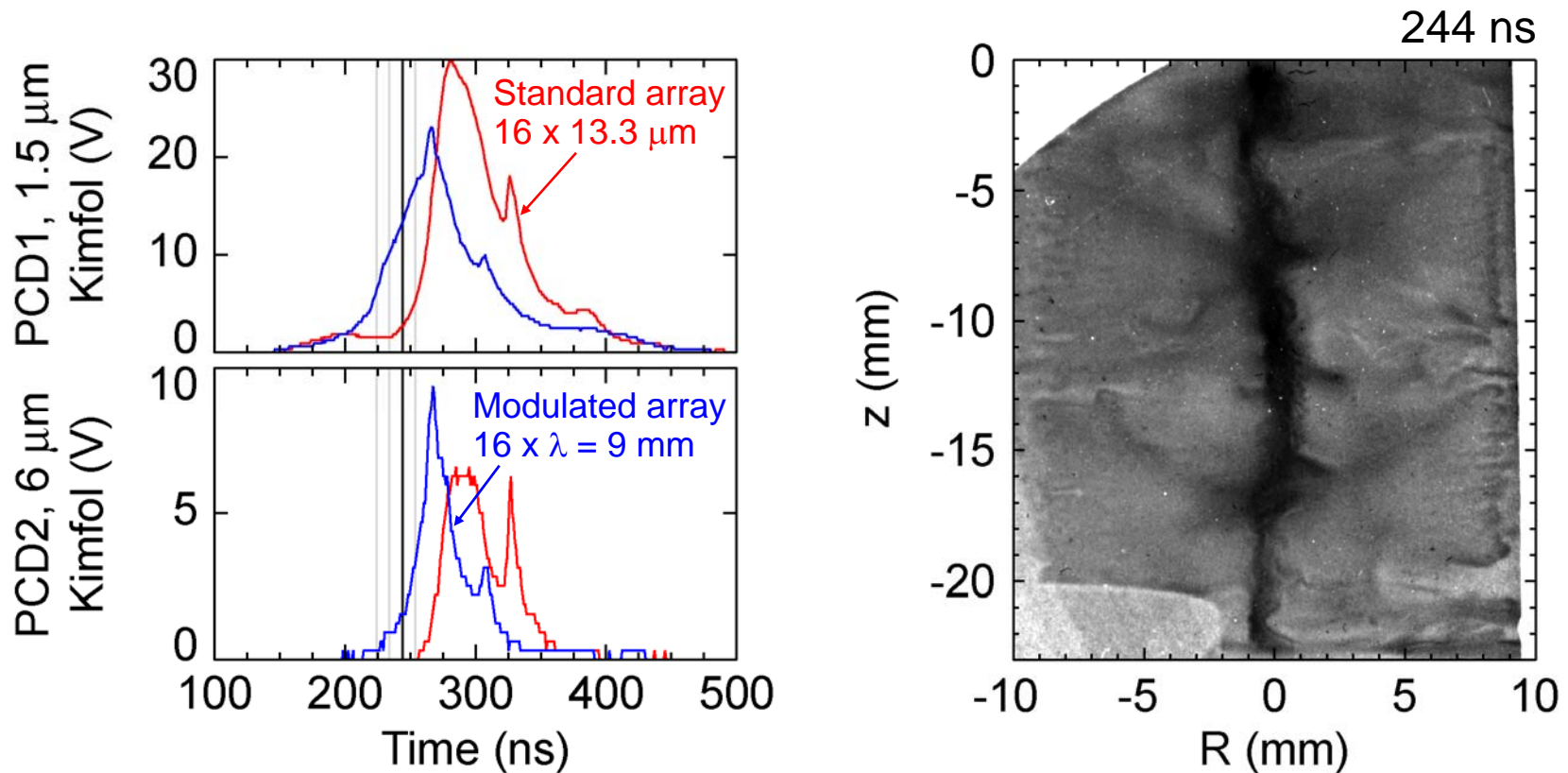
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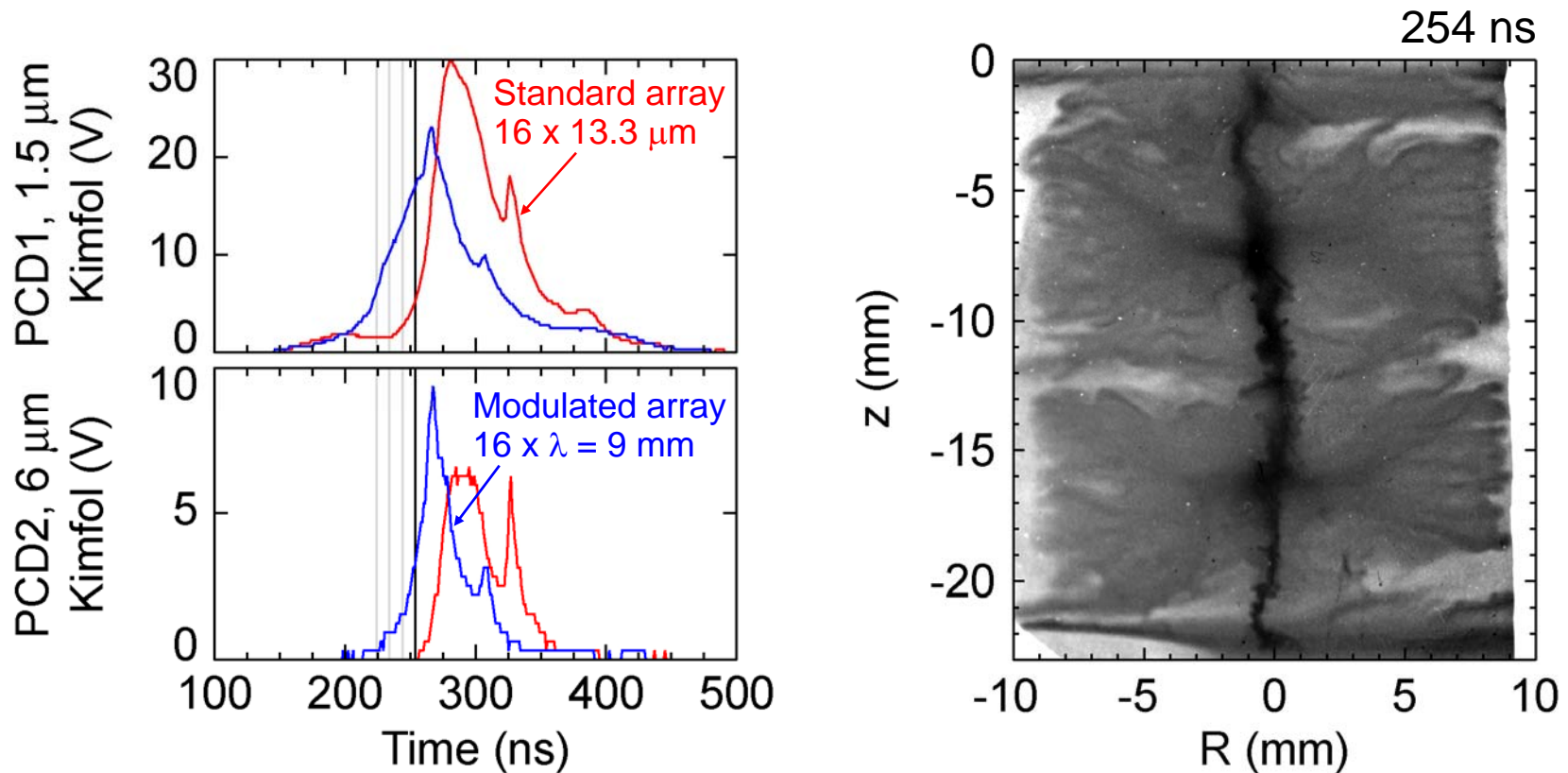


# X-ray pulse shape determined by assembly of mass on axis



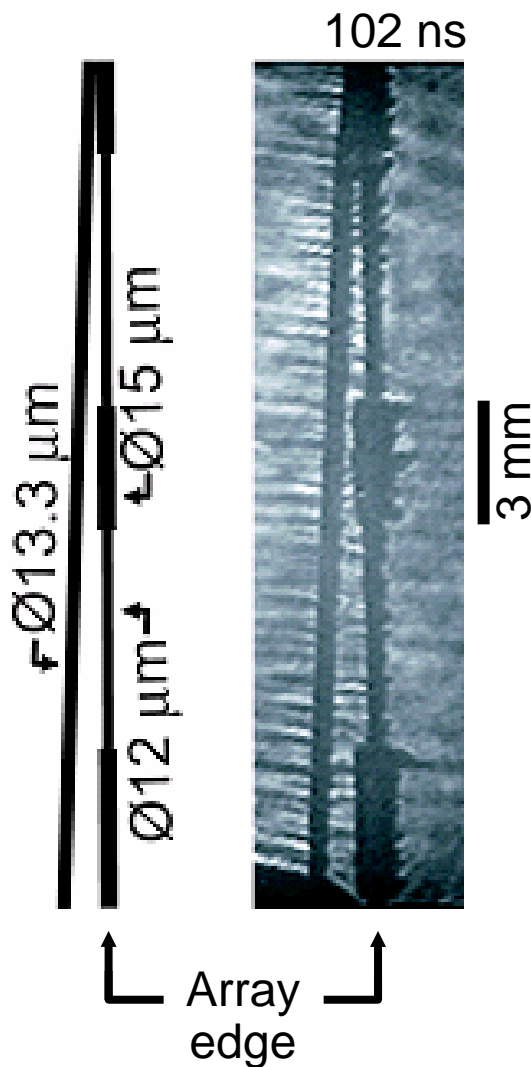
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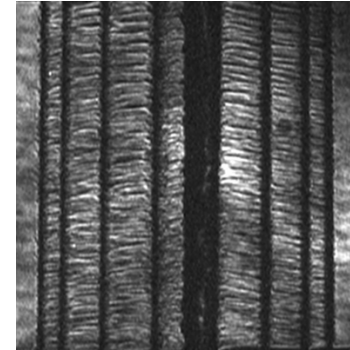
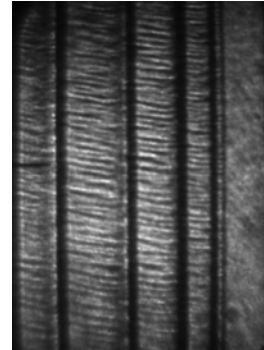


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# Natural short- $\lambda$ mode is superimposed on long- $\lambda$ perturbation



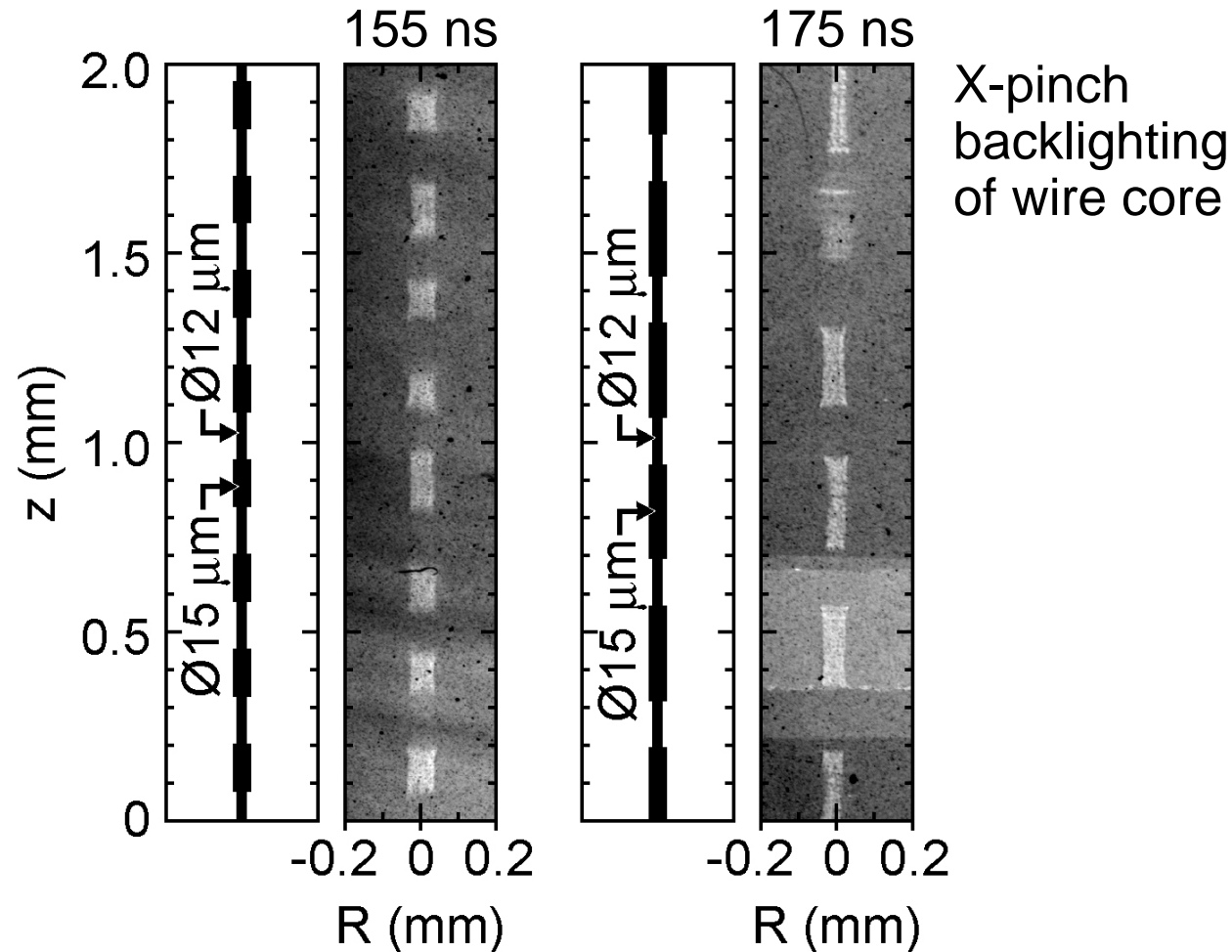
Stainless steel wire array on the 20 MA Z machine—G. Sarkisov, D. Bliss *et al.*



Aluminum wire array on the 1 MA MAGPIE driver

- Radial flares of  $\sim 0.5 \text{ mm}$  wavelength (in Al) are seen early in time on all machines due to axially non-uniform wire ablation
- In experiments with longer wavelength seeded perturbations, the  $\sim 0.5 \text{ mm}$  natural mode is seen superimposed
- What will happen if we seed  $\lambda < 0.5 \text{ mm}$ ? Will seeded and natural modes compete?
  - M. R. Douglas, C. Deeney, and N. F. Roderick, Phys. Plasmas **5**, 4183 (1998).
  - D. Ryutov and A. Toor, Phys. Plasmas **5**, 22 (1998).

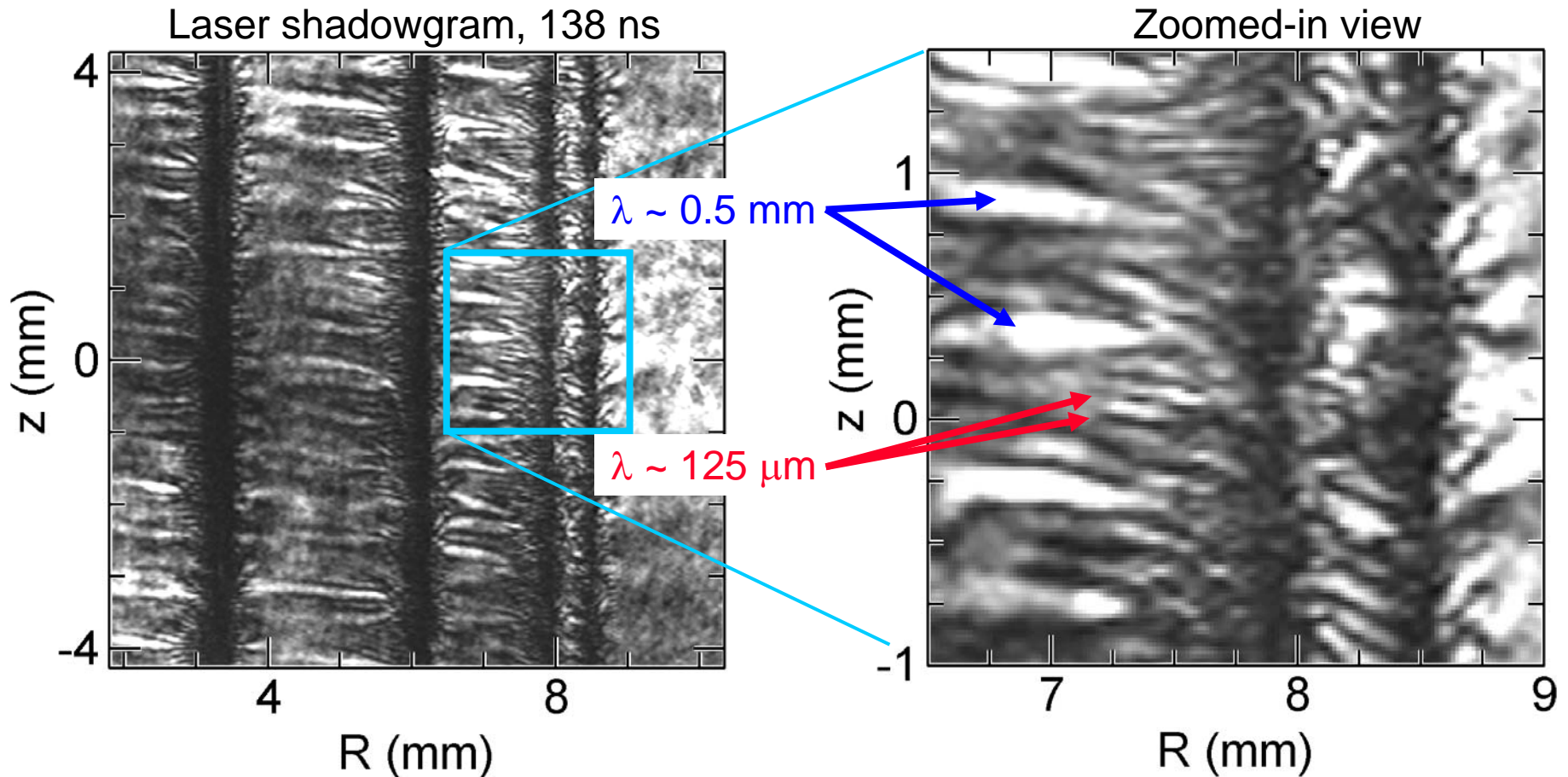
# Wire core evolution is dominated by seeded perturbation



- Mass ablated from thick segments  $<$  initial mass of thin segments—faster ablation rate near thinner wire core?

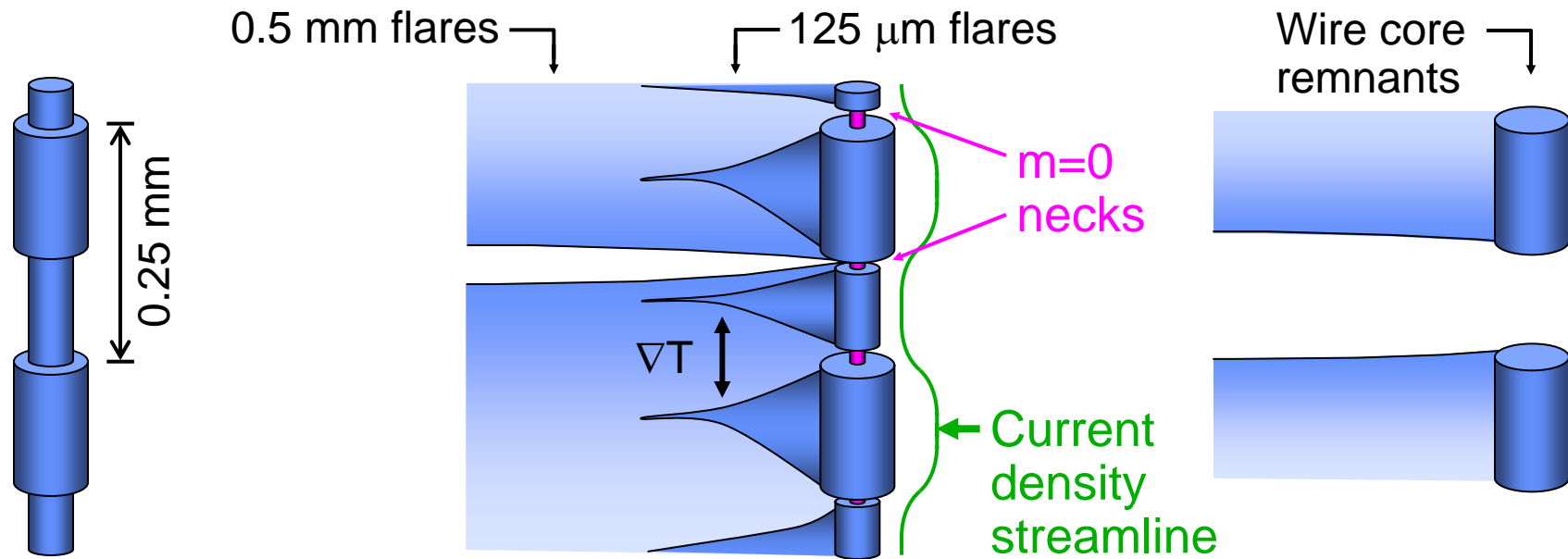


# Short wavelength perturbations compete with natural mode



- Corona is impacted by seeded features, but signature of natural mode also appears superimposed
- Natural  $\sim 0.5$  mm flare mode determined by coronal physics?

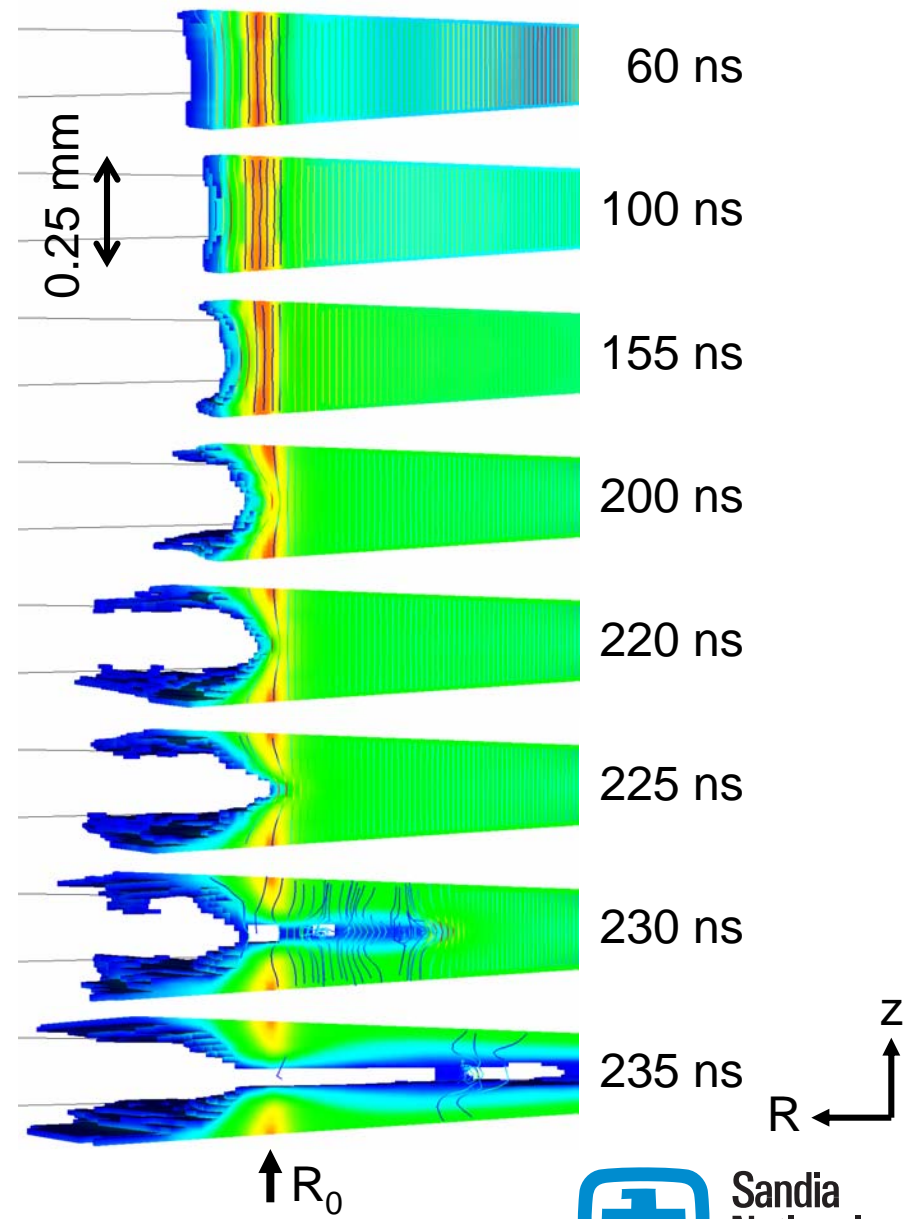
# Possible mechanism for observed behavior with short- $\lambda$ seed



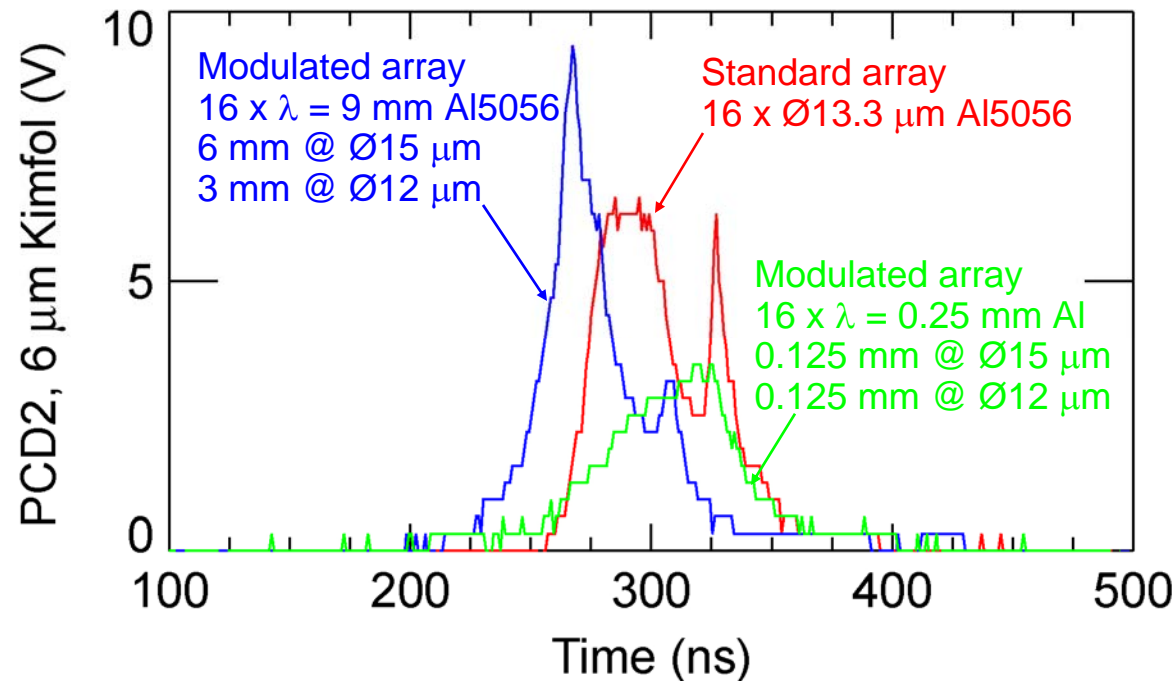
- Per the longer wavelength seeded features, discontinuities in the wire lead to enhanced  $\mathbf{j} \times \mathbf{B}$  and local necking regions
- Material flows axially out of the pinched necks, is pushed in radially by global  $\mathbf{B}$ , forms radial flares with 125  $\mu\text{m}$  period
- Thermal conduction in the ablated plasma shorts out the small- $\lambda$  seeded mode and produces  $\lambda \sim 0.5$  mm jets
  - M. G. Haines, Phys. Rev. Lett. **47**, 917 (1981).

## 3D MHD Gorgon model simulates wire core break-up

- Modulation of corona in response to seeded perturbation is seen again
- Thinner core section ablates/implodes earlier, leaving thicker wire core sections at late time
- Current stream lines press against thinner core—faster ablation rate?
- $125\text{ }\mu\text{m}$  wavelength flares are not reproduced—higher resolution needed?



# Short- $\lambda$ perturbations dramatically compromise x-ray pulse



- In short-wavelength seeded instability case, current may be shunted out of wire cores, leading to a late implosion
  - Remaining wire core sections were not expanded at normal levels
  - $m=1$  instability seen at early time in precursor column on axis
- Smaller perturbation amplitude needed to test the concept of competing modes to reduce RT radial extent



# Single wire experiments show opposite modulation trend

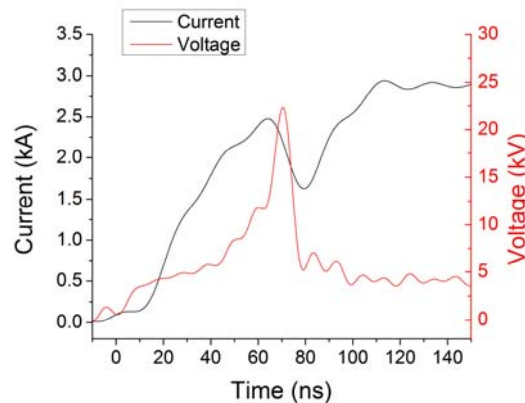
## Shot specifications

$$V_{\text{peak}} = 26.0 \text{ kV}$$

$$I_{\text{peak}} = \sim 8 \text{ kA}$$

$$t_{\text{peak}} = 72.0 \text{ ns}$$

$$dI/dt = 49.3 \text{ A/ns}$$



Initial  
configuration  
Al 5056 wire

Ø50  $\mu\text{m}$  →

Ø37.5  $\mu\text{m}$  →

Laser shadowgraph  
217 ns after start of  
the current rise

← 960 cm/ $\mu\text{s}$

← 50 cm/ $\mu\text{s}$

←  $\lambda \sim 100 \mu\text{m}$

- Thin wire sections have large corona, opposite of wire arrays
- Different breakdown mechanism? Role of global B field?

Courtesy of David Chalenski,  
Harold Barnard, Bruce Kusse



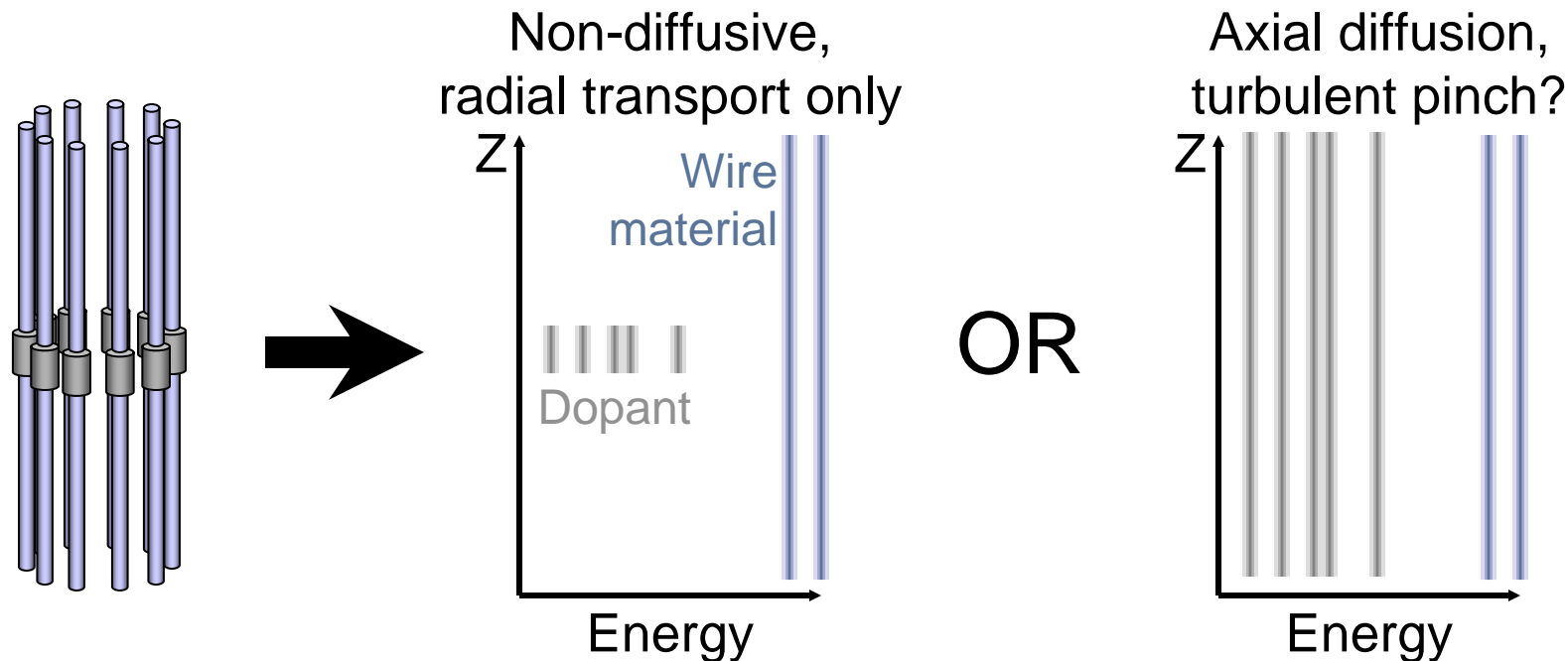
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# Localized dopant bands for particle transport studies



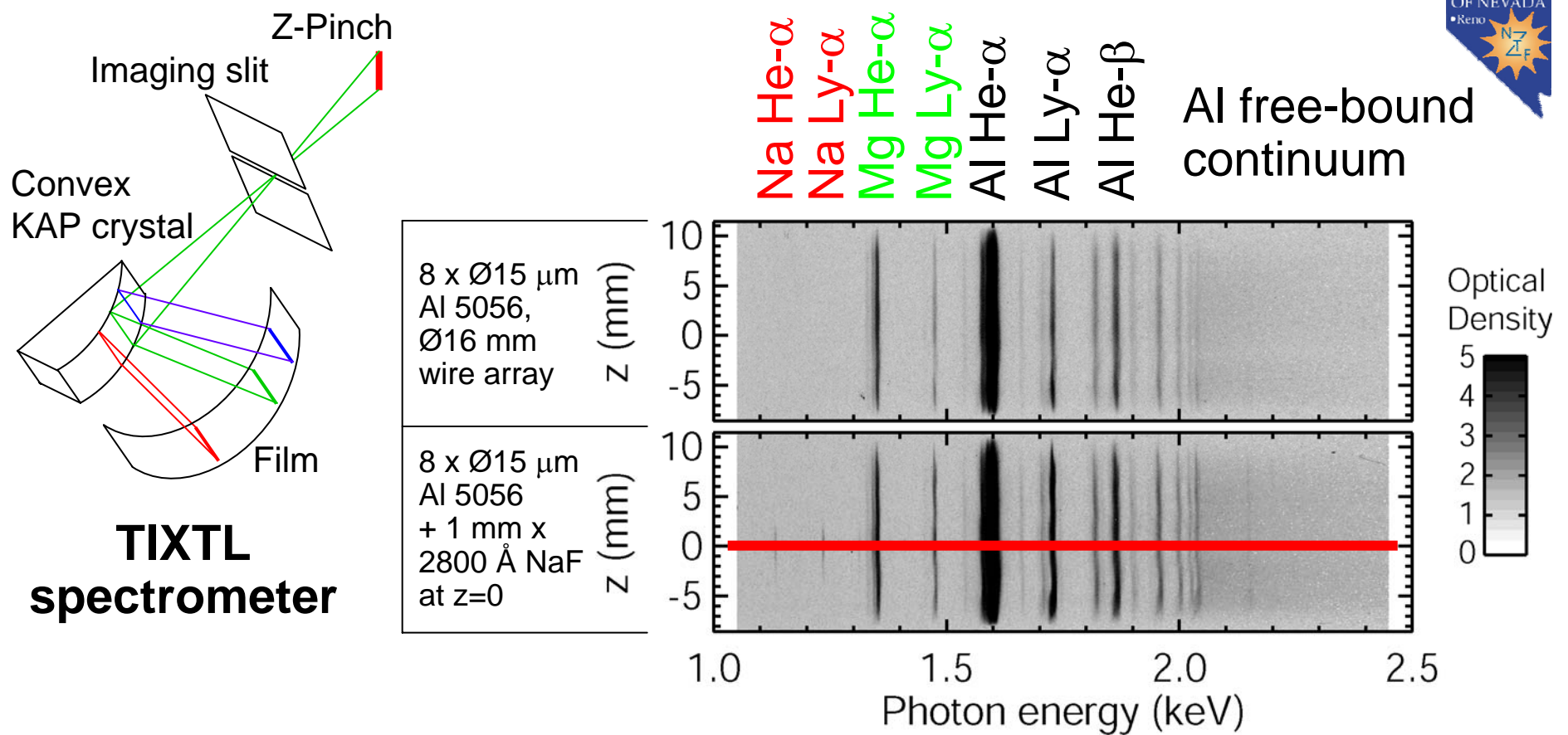
- Particle transport may be a quantitative signature of turbulence affecting z-pinch energy balance
- Localized dopants coated on wires at MPCL/Sandia
- Wire array experiments at UNR/NTF-Zebra 1 MA generator
- Localized NaF dopants previously used with laser plasmas

- Y. Al-Hadithi *et al.*, Phys. Plasmas **1**, 1279 (1994).
- K. B. Fournier *et al.*, JQSRT **71**, 339 (2001).



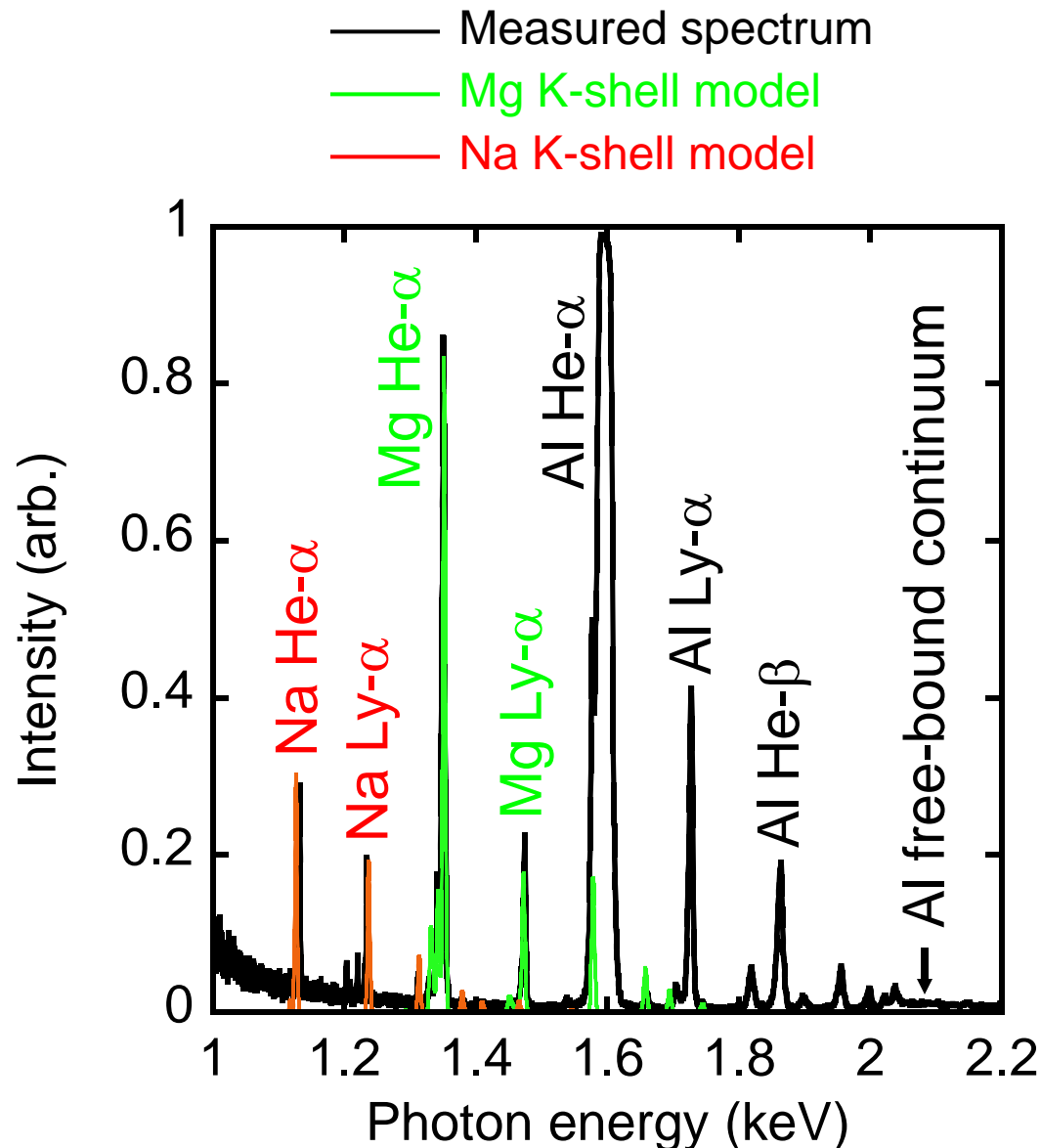
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# Localized Na K-shell emission observed on 1 MA Zebra driver



- Dopant coatings previously increased K-shell yield by mitigating opacity effects in wire array z-pinch
  - C. Deeney *et al.*, Phys. Rev. E **51**, 4823 (1995).

# Atomic physics modeling provides plasma $n_e$ , $T_e$ diagnosis



- K-shell CRE models for Mg, Na have been developed at UNR
- $n_e \sim 2 \times 10^{20} \text{ cm}^{-3}$   
 $T_e \sim 280 \text{ eV}$   
inferred from comparison with both Mg and Na lines
- Low-Z doping of tungsten wires may allow spectroscopic measurement of plasma parameters

In collaboration with Alla Safronova and Victor Kantsyrev (UNR)

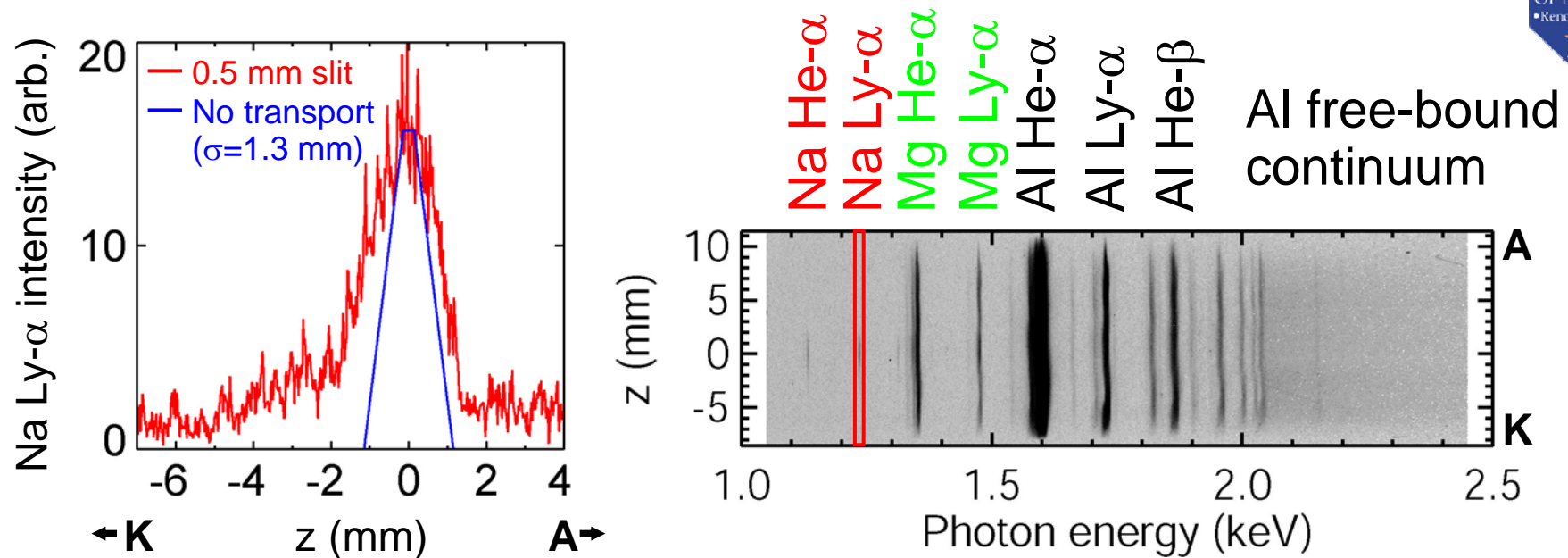


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# Na dopant is transported axially in the z-pinch plasma



- DEF film data converted to exposure using Henke tables
- Axial lineout taken over full width of the Na Ly- $\alpha$  line
- Emission is peaked near initial NaF band location
- Material has spread in axial direction (TIXTL can resolve this)
- An additional tail of material is transported in the direction of the cathode?

# Summary

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- Seeded perturbations offer a technique for exploring the role of 3D structure in wire array z-pinches
- MHD simulations with well defined initial perturbations allow study of instability growth without *ad hoc* assumptions
- Modulations in wire radius have linked current path discontinuities with seeding of RT bubbles
- Main x-ray pulse rise is due to collision of bubbles on axis assembling to form the stagnated z-pinch
- Many interesting research possibilities
  - Short-wavelength modulations and interaction with radial flares
  - Physics of magnetic bubbles
  - X-ray pulse shaping for ICF applications
  - Initiation studies with modulated single wires

