

Detecting lightning infrasound using a high-altitude balloon

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Key Points:

- First lightning infrasound detected using free-flying balloon at stratospheric altitudes over Tasman Sea in May 2016.
- Infrasonic signals matched with a few lightning strokes within 100 km range of balloon as it flew over at least two thunderclouds.
- Only a fraction of the expected infrasound signals were detected, and the cause of this remains unclear.

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Abstract

Acoustic waves with a wide range of frequencies are generated by lightning strokes during thunderstorms, including infrasonic waves (0.1 to 20 Hz). The source mechanism for these low frequency acoustic waves is still debated and studies have so far been limited to ground-based instruments. Here we report the first confirmed detection of lightning generated infrasound with acoustic instruments suspended at stratospheric altitudes using a free-flying balloon. We observe high-amplitude signals generated by lightning strokes located within 100 km of the balloon as it flew over the Tasman Sea on 17 May 2016. The signals share many characteristics with waveforms recorded previously by ground-based instruments near thunderstorms. The ability to measure lightning activity with high-altitude infrasound instruments has demonstrated the potential for using these platforms to image the full acoustic wavefield in the atmosphere. Furthermore, it validates the use of these platforms for recording and characterizing infrasonic sources located beyond the detection range of ground-based instruments.

Plain-language summary

Lightning generates sound waves across a wide range of frequencies, including below the threshold for human hearing at 20 Hz. How these waves at less than 20 Hz, also known as infrasound waves, are generated during a lightning stroke is currently an area for debate. So far, measurements of lightning infrasound waves have been limited to microphones fixed to the ground and models have shown that only a small section of sound waves actually reach the ground. Here we show lightning infrasound that has been detected using microphones suspended over a thunderstorm using a balloon flying at 32 km height. This opens up the possibility of using balloons in future studies to make better measurements of infrasound waves generated by lightning activity and in turn, give a better idea of how they are generated. It also shows how balloons can be used to record infrasound waves far away from land and therefore beyond the detection limit of ground-based microphones.

1 Introduction

Acoustic signals with frequencies between 0.02 to 20 Hz are classified as infrasound and are not audible to humans. A wide variety of sources have been found to generate infrasound, including: volcanoes, earthquakes, avalanches, tsunamis, meteors, aurora, thunderstorms, wind-mountain interactions, supersonic aircraft, rockets, and chemical and nuclear explosions [*Campus and Christie*, 2010]. Infrasonic signals can travel hundreds to thousands of kilometers through the atmosphere, sampling areas from the Earth's surface up to the thermosphere. A variety of institutions maintain arrays for monitoring purposes, such as volcano observatories [*Fee and Matoza*, 2013] or the International Monitoring System [*Christie and Campus*, 2010]. The vast majority of infrasound studies currently use ground-based instrument arrays and networks. Infrasound instruments deployed on the ground may be subject to high-levels of background noise which may obscure the signals of interest, or the signals may arrive distorted due to topographic or atmospheric propagation effects [e.g. *Lacanna and Ripepe*, 2013; *Kim and Lees*, 2014]. Furthermore, the intensity of the acoustic wavefield may be as much as 50% greater directly above the source compared that at a similar horizontal distance [*Blackstock*, 2000].

A series of studies have recently taken place to explore how to fill this gap in our ability to monitor the atmospheric acoustic wavefield. These experiments have tested the use of microphones suspended underneath free-floating balloons to record infrasound at high-altitude [e.g. *Bowman and Lees*, 2015, 2017; *Bowman and Albert*, 2018]. Balloon deployments conducted as part of the NASA High-Altitude Student Platform (HASP) program described evidence of the ocean microbarom as well as other signals of unknown provenance [*Bowman and Lees*, 2015]. A follow-up experiment showed that the ocean

64 microbarom was often detectable in the stratosphere but not at ground-level, either due
 65 to low noise at the balloons, an elevated acoustic duct, or both [Bowman and Lees, 2017].
 66 In 2017, four microphone-bearing solar balloons were launched concurrently and success-
 67 fully detected and located a chemical explosion on the ground [Bowman and Albert, 2018].
 68 So far, none of these experiments have confirmed the detection of other infrasonic sig-
 69 nals from sources such as lightning storms.

70 Acoustic emissions from lightning, described as thunder, can produce a broadband
 71 range of frequencies. The audible component of thunder (20-20,000 Hz), as well as en-
 72 ergy in the infrasonic range, is understood to come mostly from shock waves produced
 73 by the rapid expansion of a lightning channel due to current flow and heating [Few *et al.*,
 74 1967]. In addition, numerous studies have observed infrasound generated by cloud-to-
 75 ground (CG) and intracloud (IC) lightning flashes [e.g. Balachandran, 1983; Assink *et al.*,
 76 2008; Farges and Blanc, 2010; Arechiga *et al.*, 2014]. The lightning signal is often a dis-
 77 crete pulse characterized by an initial compression followed by a rarefaction with maxi-
 78 mum amplitudes in the range of 0.05 to 5 Pa and a spectral peak in the range of 0.2 to
 79 2 Hz [Dessler, 1973; Bohannon *et al.*, 1977; Assink *et al.*, 2008; Campus and Christie,
 80 2010]. Multiple production mechanisms have been postulated for the infrasonic acous-
 81 tics detected during lightning storms. This includes rapid intensification of the electric
 82 field just prior to the flash, ohmic heating of the air by charge flowing into the channel,
 83 and interaction between the positive and negative charge layers in the storm cloud [Dessler,
 84 1973; Bohannon *et al.*, 1977; Pasko, 2009]. Furthermore, the acoustic wavefield gener-
 85 ated by this mechanism was predicted to be orientated vertically, restricting the hori-
 86 zontal detection range of lightning infrasound [Dessler, 1973; Pasko, 2009]. Validation
 87 of the production mechanism and acoustic wavefield has been confounded by the diffi-
 88 culty in locating the charge layers in the storm cloud, as well as characterizing the struc-
 89 ture of the parent lightning flash. Advancements in location algorithms and instruments
 90 deployments such as the Lightning Mapping Array (LMA) have produced observations
 91 that refute the previously proposed production mechanisms [Arechiga *et al.*, 2014]. In-
 92 stead, the observations suggest that the infrasonic signals from lightning flashes may be
 93 produced by electrostatic interaction of charge deposited in the streamer zone of a light-
 94 ning channel; that is, acoustic compression waves may be generated by electrostatic forces
 95 causing air within the streamer zone to expand [Arechiga *et al.*, 2014]. Lightning infra-
 96 sound has also been detected at ranges of up to 150 km from the source, contrary to pre-
 97 vious predictions of a vertically orientated acoustic wavefield [e.g. Farges and Blanc, 2010].

98 Here we report on detections of lightning infrasound recorded by a high-altitude
 99 balloon flying over the Tasman Sea on 17 May 2016. We present evidence for signals recorded
 100 from at least two groups of lightning flashes during a 6 hour period. Measurements sug-
 101 gest that the detection of lightning infrasound was limited by the distance and atmo-
 102 spheric conditions between source and receiver. A few example signals from lightning
 103 strokes are isolated and their waveform characteristics are briefly discussed. These ob-
 104 servations, the first of their kind reported, suggest that microphones deployed on high-
 105 altitude balloons can offer additional insights into the production mechanisms of light-
 106 ning infrasound.

107 2 Data

108 An acoustic sensor package was included as a piggyback payload on the NASA Ul-
 109 tra Long Duration Balloon (ULDB) flight launched from Wanaka, New Zealand on 16
 110 May 2016. The ULDB landed in Peru on 2 July 2016 for a total flight duration of 46 days.
 111 The ULDB balloon position and height was recorded using an onboard GPS unit, and
 112 records show the full flight included a full circumnavigation of the southern hemisphere
 113 [Bowman *et al.*, 2017]. The acoustic sensor package recorded data for the first 20 days
 114 of the flight, and the ocean microbarom was recorded throughout as well as other sig-
 115 nals of unknown provenance [Bowman and Lees, 2018].

116 The sensor package contained three InfraBSU microphones [*Marcillo et al.*, 2012]:
 117 one control and a pair with reversed polarities. The reversed polarity sensor was achieved
 118 by placing the mechanical sensor on the opposite port. The reversed polarity microphone
 119 pair were combined into a single channel via:

$$M = \frac{M_+ - M_-}{2} \quad (1)$$

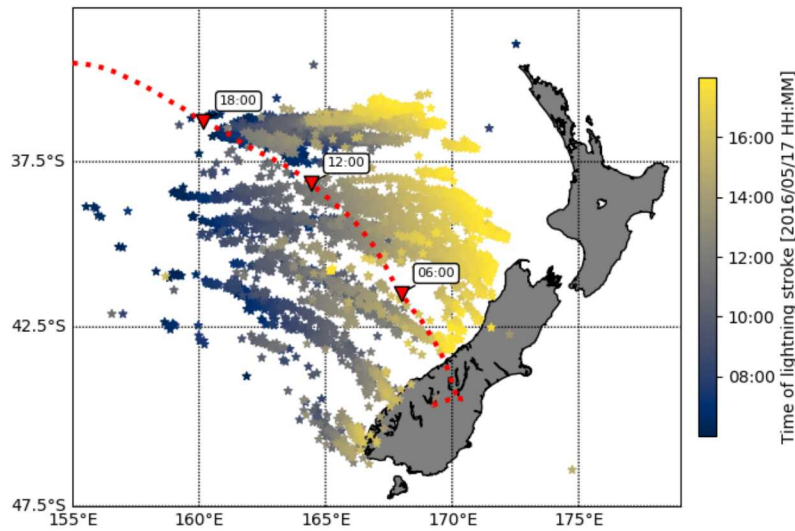
120 where M is the data analyzed in this article, and M_+ and M_- are the data from the mi-
 121 crophones with positive and negative polarities, respectively. The control sensor was a
 122 microphone that was disabled by removing the mechanical filter entirely. This acoustic
 123 sensor trio was designed to robustly distinguish between true pressure fluctuations and
 124 spurious signals, such as electronic interference [*Bowman et al.*, 2017]. Data was recorded
 125 at 200 samples per second at 64x gain using an Omnirecs Datacube digitizer. The mi-
 126 crophones were not calibrated to the pressure and temperature conditions experienced
 127 during the flight, but their primary effect should be to lower the corner period of the sen-
 128 sors [*Bowman et al.*, 2017]. The acoustic waveforms presented here are high-pass filtered
 129 at 0.6 Hz in order to remove high-amplitude signals contributed from the ocean micro-
 130 barom [*Bowman and Lees*, 2018], atmospheric gravity waves generated by thunder cloud
 131 convection [*Blanc et al.*, 2010], and balloon oscillations [*Anderson and Taback*, 1991].
 132 (Unfiltered signals recorded by the acoustic package can be seen in Figure S1 in supple-
 133 mentary information.) The microphones and digitizer were each powered by separate Lithium
 134 battery packs, and contained within high density styrofoam shipping boxes for thermal
 135 insulation. Internal temperatures within the digitizer ranged from -26 to 7 °C during the
 136 flight.

137 The lightning stroke detections and location data used in this article were detected
 138 and recorded by the World Wide Lightning Location Network (WWLLN). The WWLLN
 139 is an instrument network capable of locating and timing lightning strokes at long range
 140 (thousands of kilometers) to within <10 km and <10 μ s [*Hutchins et al.*, 2012]. The net-
 141 work uses very-low-frequency radio wave (3-30 kHz) receivers distributed around the globe
 142 to identify the time of group arrival for individual lightning waveforms, or sferics. The
 143 network is capable of detecting both CG and IC discharges, but the latter are typically
 144 underrepresented in detection databases as they produce weaker electromagnetic pulses
 145 [*Behnke and McNutt*, 2014]. As of 2010, the estimated detection efficiency for the net-
 146 work was \sim 11% for all strokes and >30% for more powerful strokes [*Hutchins et al.*, 2012].
 147 These values may seem low, but the WWLLN was not designed to detect all lightning
 148 strokes but instead to provide a global overview of lightning activity [*Dowden et al.*, 2008].
 149 Other lightning stroke datasets may exist from other global detection networks, but were
 150 not available for the analysis presented herein.

151 3 Observations

152 The ULDB was launched from Wanaka, New Zealand just before 0000 UTC on 17
 153 May 2016 started flying east as it ascended. Once the craft approached and breached
 154 30 km altitude, it turned to the west and flew out over the Tasman Sea and towards Aus-
 155 tralia (Fig. 1). During this period, the WWLLN detected intense lightning activity from
 156 multiple thunderstorms approaching New Zealand from the west (Fig. 1, Movie S1 in
 157 supporting information). From 0800 to 1400 UTC, 2994 strokes were detected and lo-
 158 cated by the WWLLN across the Tasman Sea, of which 2554 were located within 500
 159 km of the ULDB (Fig. 2a). At approximately 0945 and 1200 UTC the ULDB passed di-
 160 rectly over or near lightning activity which correlates with an increase in acoustic activ-
 161 ity recorded at the ULDB (Fig. 2a, b). Acoustic signals are recorded with peak-to-peak
 162 amplitudes of up to 0.05 Pa and a broadband range of frequencies from 0.6 to 20 Hz (Fig.
 163 2b, c).

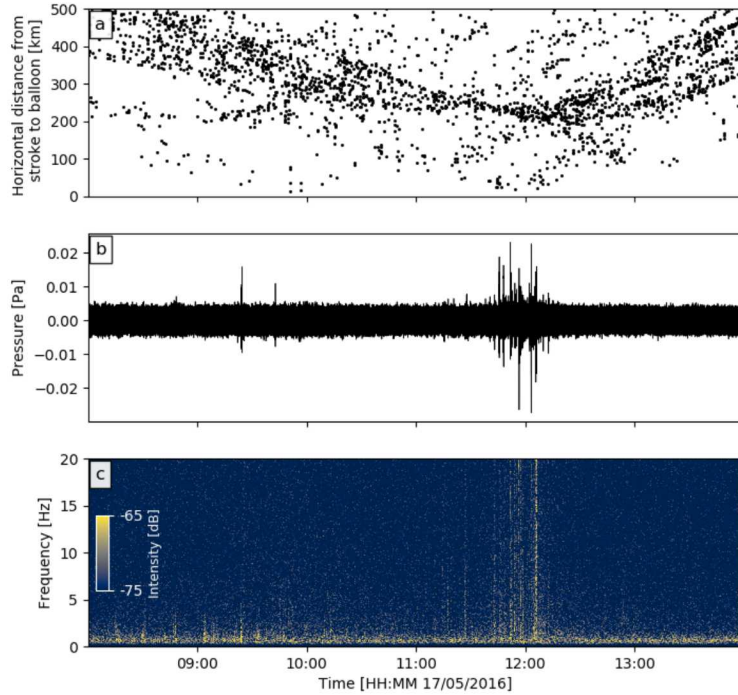
173 As the acoustic sensor package on the ULDB was fundamentally a single element
 174 station, back-azimuths and slowness vectors cannot be calculated to locate sources of de-



164 **Figure 1.** Map of the Tasman Sea with the locations of lightning detected by the WWLLN
 165 from 0600 to 1800 UTC on 17 May 2016, where color represents the progression of time (see col-
 166 orbar). Also plotted is the path of the ULDB balloon after it was launched from Wanaka, New
 167 Zealand (red dotted line), and its location at 0600, 1200 and 1800 UTC on 17 May 2016 (red
 168 triangles). (For an animated version of this figure, see Movie S1 in supporting information.)

175 tected signals. To estimate detection ranges, ray tracing was used to model infrasonic
 176 propagation paths between lightning and the ULDB. Ray tracing was performed using
 177 classical geometric acoustics techniques and a plane wave assumption, calculated within
 178 the open source GeoAc ray tracing software [Blom and Waxler, 2012]. Rays were launched
 179 at intervals of 1° from a point source at a height of 4 km. This source height was based
 180 on previously used heights for modeling lightning infrasound [Pasko, 2009; Farges and
 181 Blanc, 2010]. Atmospheric profiles were derived from the 12z Global Forecast System
 182 (GFS) analysis model run, located at the latitude/longitude coordinates for the ULDB
 183 at 1200 UTC on 17 May 2016 (Fig. S2 in supplementary information). For a source at
 184 4 km height and a receiver at 32 km height, direct arrivals from the source should only
 185 be expected <110 km horizontal distance in all directions (Fig. 3). At this distance, 34
 186 lightning strokes were recorded when the ULDB flew near a storm at approximately 1200
 187 UTC (Fig. 4a). 10 strokes were recorded within 100 km during the earlier storm at 0945
 188 UTC (Fig. S3 in supplementary information).

193 To match infrasonic signal peaks and specific lightning strokes, we compute the time
 194 needed for waveforms from each stroke within a limited distance to arrive at the ULDB.
 195 We take a simplified approach and assume that the atmosphere can be approximated
 196 with a bulk acoustic wave speed of 300 ms^{-1} . Furthermore, all acoustic waveforms ar-
 197 riving at the ULDB platform are assumed to be direct arrivals from the source. Out of
 198 34 lightning strokes within 100 km of the balloon between 1130 and 1230 UTC, multi-
 199 ple events appear to match directly with peaks in recorded acoustic amplitudes (Fig. 4b).
 200 Here we present three matches which occur at 1145, 1148, and 1152 UTC (Fig. 4c, d,
 201 and e). Other possible matches occur at 1143, 1156 and 1205 UTC but source-signal pairs
 202 cannot be distinguished due to multiple closely spaced source strokes or infrasonic ar-



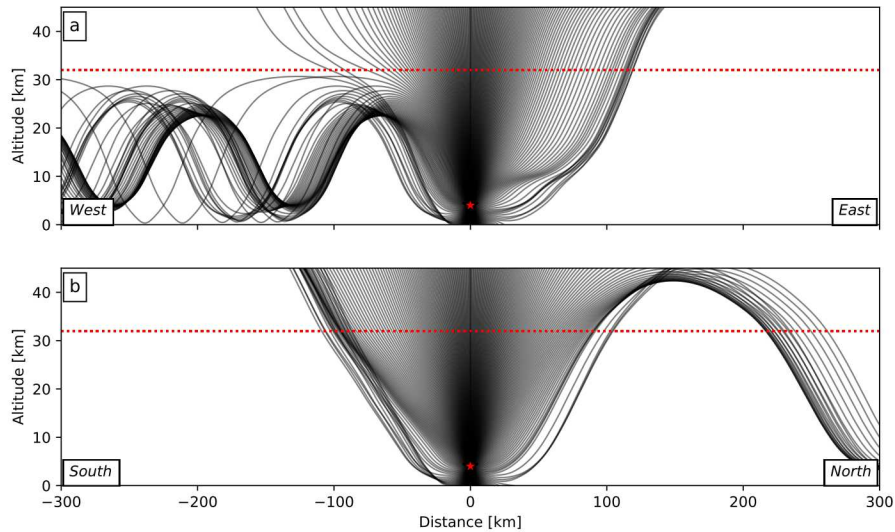
169 **Figure 2.** (a) Horizontal distances for each stroke within 500 km of the balloon’s location at
 170 the time of the stroke, from 0800 to 1400 UTC on 17 May 2016. (b) High-pass filtered (0.6 Hz)
 171 infrasound over the same time period as recorded at the ULDB. (c) Frequency spectrogram of the
 172 waveform plotted in (b).

203 rivals. None of the 10 detected strokes directly match with infrasonic signals during the
 204 earlier storm at 0945 UTC (Fig. S3 in supplementary information).

212 The energy density of an expanding acoustic shock wave from a lightning stroke
 213 can be estimated from the peak frequency of the recorded waveform [Few, 1969]. For a
 214 given acoustic waveform with peak frequency, f_p :

$$f_p = 0.63c_0\sqrt{P_0/E} \quad (2)$$

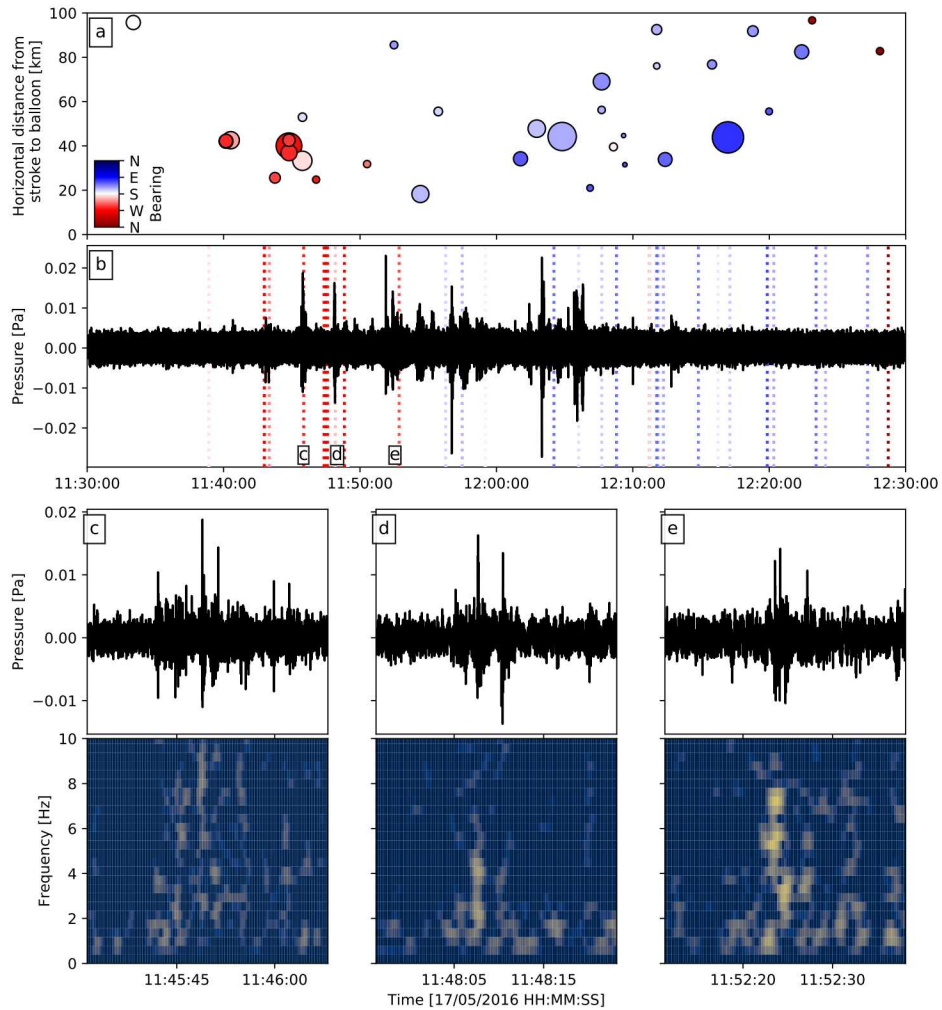
215 where c_0 is the local speed of sound (300 ms^{-1}), P_0 is the atmospheric pressure
 216 (60 kPa for a source at 4 km altitude), and E is the energy per unit length [Few, 1969].
 217 For each matched waveform plotted in Fig. 4c, d and e, we find peak frequencies of 2.65,
 218 1.27, and 5.27 Hz, respectively (Fig. S4a in supplementary info). Using equation (2) we
 219 find a positive linear relationship between the energy densities calculated from the acous-
 220 tic waveforms and the energies detected for the lightning strokes by the WWLLN (Fig.
 221 S4b in supplementary info). We also find no relationship between the calculated energy
 222 densities and the stroke-balloon distance (Fig. S4c in supplementary info).



189 **Figure 3.** Ray-tracing propagation for an acoustic source (red star) at 4 km height along
 190 East-West (a) and North-South (b) profiles using realistic atmospheric conditions derived from
 191 the Global Forecast System. The red dotted line at 32 km indicates the approximate height of
 192 the ULDB balloon on 17 May 2016.

223 4 Discussion and Conclusions

224 Here we have presented evidence that lightning infrasound was observable by acoustic
 225 instruments suspended at stratospheric altitudes by free-flying balloons. Several wave-
 226 forms were matched with detected lightning strokes through a simple time delay approach
 227 (Fig. 4c, d, e). Our matches are supported by a positive linear relationship between the
 228 WWLLN energy estimation for the matched lightning strokes and the energy densities
 229 calculated from the acoustic waveforms (Fig. S4 in supplementary info). Here, we have
 230 assumed that the signals represent direct arrivals between the source and receiver. To
 231 test this assumption, we searched for eigenray solutions using the GeoAc software pack-
 232 age. Direct arrivals for all three waveforms are found for the distances and azimuths to
 233 their associated lightning strokes (Fig. S5 in supplementary info). However, the arrival
 234 times calculated using the eigenray paths described earlier do not readily match with the
 235 recorded arrival times of signals (Fig. S4 in supplementary information). It is worth not-
 236 ing that there were a number of signals recorded that do not readily match with any light-
 237 ning strokes detected by the WWLLN, and *vice versa* (Fig. 4a, b and Fig. S3 in sup-
 238plementary info). If the estimated detection rates of the WWLLN are correct [11-30%;
 239 *Hutchins et al., 2012*], then there may have been as many as 100-300 lightning strokes
 240 within 100 km of the ULDB. This number of high-amplitude signals was not recorded
 241 at the ULDB (Fig. 4b), therefore the total number signals recorded at the ULDB un-
 242 derrepresents the true total of lightning strokes that occurred within range of the bal-
 243 loon. This is similar to detection rates of lightning infrasound by ground-based instru-
 244 ments [e.g. *Farges and Blanc, 2010*]. Complex atmospheric conditions in thunderclouds
 245 likely refract the generated acoustic waves away from the receiver [*Jones and Bedard,*
 246 *2015*]. Additionally, it is possible that not all lightning strokes generate measurable in-
 247 frasound. This may be attributed to very low signal-to-noise ratios, especially for smaller
 248 lightning strokes or those located further from the instrument than the rest of the thun-



205 **Figure 4.** (a) Horizontal distances for each stroke within 100 km of the balloon from 1130 to
 206 1230 UTC on 17 May 2016. each stroke is sized by the stroke energy, and colored by the bear-
 207 ing from the balloon to the stroke location. (b) High-pass (0.6 Hz) filtered acoustic waveform as
 208 recorded at the ULDB over the same time period. Vertical dotted lines indicate calculated time
 209 of arrivals for strokes in panel (a), colored by the bearing. Letters on bottom indicate locations
 210 of plotted example waveforms in panels c, d and e. (c,d,e) Example acoustic signals (top) from
 211 lightning strokes and their respective spectrograms (below).

249 derstorm cloud [Farges and Blanc, 2010]. This latter process is illustrated by a higher
 250 number of signals matches occurring during the storm approach towards the ULDB in-
 251 stead of during their divergence (Fig. 4b). A relatively low-frequency signal is recorded
 252 during an earlier storm at 0945 UTC (Fig. S3 in supplementary info). The low-frequency
 253 characteristic of this signal suggests an alternative source to lightning [e.g. meteors, tran-
 254 sient luminous events; Farges and Blanc, 2010; Edwards, 2010], or that the original light-
 255 ning signal was altered by absorption and/or directivity between the source and receiver.

256 For the ray propagation and eigenray modeling we have assumed a point source
 257 for the lightning infrasound at 4 km altitude. Infrasound sources from lightning have been

258 mapped up to 12 km altitude in the thundercloud [Anderson et al., 2014; Arechiga et al.,
 259 2014]. Furthermore, the mapped current flow within lightning strokes suggests the source
 260 geometry can resemble complex, dendritic structures [Anderson et al., 2014]. To test whether
 261 our source shape and height assumption was viable, we have repeated the ray tracing
 262 modeling but with a source at 32 km height instead, the ULDB flying altitude. In re-
 263 verse, this can be seen as all possible locations for sources whose acoustic waveforms will
 264 be recorded at the receiver. Results suggest that a receiver at 32 km should record sig-
 265 nals from heights up to 12 km and take-off angles between 45-90° within a 100 km hor-
 266 izontal range in all directions (Fig. S6 in supplementary information). Therefore, the source
 267 configuration used for the raypath modeling was reasonable for the lightning infrasound
 268 described here.

269 A key assumption here was that the atmospheric conditions during the generation
 270 of the lightning infrasound was relatively simple and stratified. Thunderstorms require
 271 unstable air to form and persist, and often include significant vertical wind shear due
 272 to thermal plumes. The wind and temperature profiles used for ray path modeling here
 273 was likely a highly simplified representation of reality. Ray tracing from sources directly
 274 below thermal plumes such as those found within thunderstorm clouds have shown how
 275 they may act as vertical waveguides and thus can greatly distort the acoustic wavefield
 276 [Jones and Bedard, 2015]. Ray paths are found to either converge or diverge at the top
 277 of the thermal plume dependent on the height and width of the plume [Jones and Be-
 278 dard, 2015]. Therefore the ray-paths and eigenrays presented here are likely an oversim-
 279 plified representation of the true paths taken by the acoustic waves before their detec-
 280 tion at the ULDB. Future studies of acoustic wavefields generated by lightning infrasound
 281 must take into account the complex refraction patterns induced by vertical columns of
 282 wind shear within thunderclouds.

283 The measured waveforms do not display the compression-rarefaction-compression
 284 shape that had been modeled as generated by a rapid intensification and discharge of
 285 the electric field in the thundercloud [Pasko, 2009]. It must be noted that the waveforms
 286 modeled by Pasko [2009] are of 0.1-1 Hz frequency and the acoustic wavefield was strongly
 287 oriented in the vertical direction [Dessler, 1973]. The detection of lightning infrasound
 288 from strokes at more than several tens of kilometers from the ULDB does not support
 289 the theory that acoustic wavefields generated by lightning infrasound are strictly oriented
 290 vertically. Instead, the waveforms here share amplitude, frequency and range detection
 291 characteristics with previously recorded lightning infrasound signals which were attributed
 292 to charge deposition in the lightning channels [e.g. Farges and Blanc, 2010; Arechiga et al.,
 293 2014]. The electrostatic forces caused by the charge deposition may cause the air in the
 294 streamer zone to expand, producing an acoustic compression wave whose period corre-
 295 sponds to the size of the streamer zone [Arechiga et al., 2014]. This is in addition to the
 296 rapid air expansion generated by extreme heating of the lightning channel that produces
 297 audible and infrasonic acoustic waves [Few et al., 1967].

298 The observations here of lightning infrasound recorded by a high-altitude balloon
 299 over the Tasman Sea in May 2016 are fortuitous. Yet, they also demonstrate the poten-
 300 tial for using these platforms to expand our understanding of how infrasound may be gen-
 301 erated by lightning strokes. While the observations presented here support the charge
 302 deposition mechanism postulated by Arechiga et al. [2014], we cannot yet rule out other
 303 possible mechanisms. Previous studies of these processes have thus far used only ground-
 304 based instruments which only offer an approximately two-dimensional view of the acous-
 305 tic wavefield. Future studies which incorporate both balloon- and ground-based instru-
 306 ments will have an opportunity to capture a three-dimensional view of the infrasonic waves
 307 generated during lightning storms. Furthermore, simultaneous deployments of multiple
 308 instrument-bearing balloons are able to calculate back-azimuths to each infrasound source
 309 [e.g. Bowman and Albert, 2018], as well as provide the elevation angle of incoming sources
 310 [see Supplemental information in Bowman and Lees, 2018]. The lightning infrasound sig-

311 nals presented here are derived from strokes occurring >100 km from any significant land-
 312 mass where ground-based instruments would be located, meaning these signals may never
 313 have been recorded otherwise. Therefore, balloon-based instruments offer a way to record
 314 signals from events which may not otherwise have been recorded by established infra-
 315 sound instrument networks. As well as oceanic thunderstorms, this may include Tran-
 316 sient Luminous Events [e.g. sprites; *Farges and Blanc*, 2010], meteors [e.g. *Edwards*, 2010],
 317 and supersonic auroral arcs [e.g. *Pasko*, 2012].

318 Acknowledgments

319 We would like to thank the NASA Balloon Program Office and the NASA Columbia Sci-
 320 entific Ballooning Facility for hosting our instrumentation on the 2016 Ultra Long Du-
 321 ration Balloon Mission. The authors also wish to thank the World Wide Lightning Lo-
 322 cation Network (<http://wwlln.net>), a collaboration including over 50 institutions, for pro-
 323 viding the lightning data used in this paper. The acoustic data presented here is acces-
 324 sible at datadryad.org under the following doi: 10.5061/dryad.40877s4. This work was
 325 supported by National Science Foundation Grant AGS-1551999. Sandia National Labo-
 326 ratories is a multi-mission laboratory managed and operated by National Technology
 327 and Engineering Solutions of Sandia, LLC., a wholly owned subsidiary of Honeywell In-
 328 ternational, Inc., for the U.S. Department of Energy’s National Nuclear Security Admin-
 329 istration under contract DE-NA0003525. The views expressed here do not necessarily
 330 reflect the views of the United States Government, the United States Department of En-
 331 ergy, or Sandia National Laboratories. We would like to thank Jelle Assink and the anony-
 332 mous reviewers whose comments and suggestions helped improve this manuscript.

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Supporting Information for “Detecting lightning infrasound using a high-altitude balloon”

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Contents

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Additional Supporting Information (Files uploaded separately)

1. Caption for Movie S1

Introduction

Supplementary information provides Figures S1 to S6, and a caption for Movie S1.

Movie S1.

Animated version of Figure 1 in the main article. Each frame covers a 10 min interval as the ULDB (Red Triangle) flies over the Tasman Sea from 0600 to 1800 UTC on 17 May 2016. Locations of lightning strokes detected by the WWLLN are also plotted, and coloured by age for up to an hour.

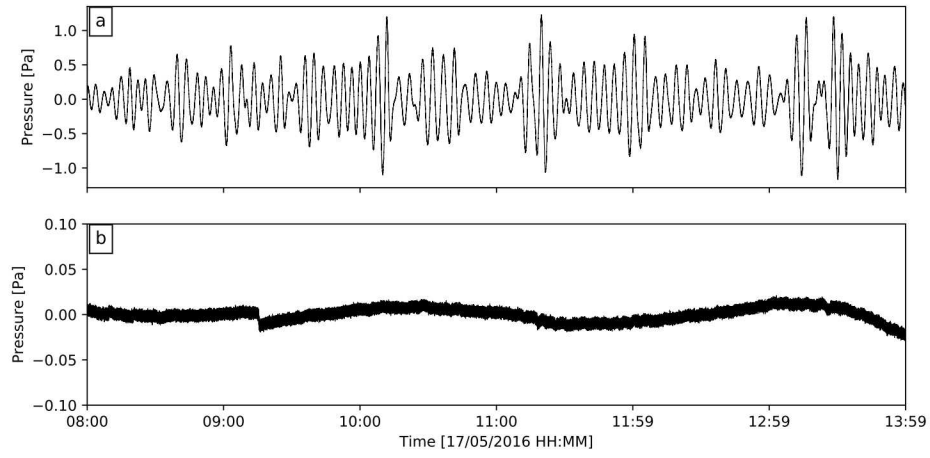
Figures S1 to S6

Figure S1 : Unfiltered acoustic waves as recorded by the acoustic package on the ULDB. (a) The combined acoustic wave as recorded by the reversed polarity pair of instruments. (b) ‘Acoustic’ waveform as recorded by the disabled ‘control’ instrument in the acoustic package. The source of the jump at 0915 UTC is unknown but has no apparent effect on the acoustics recorded by the reversed polarity pair.

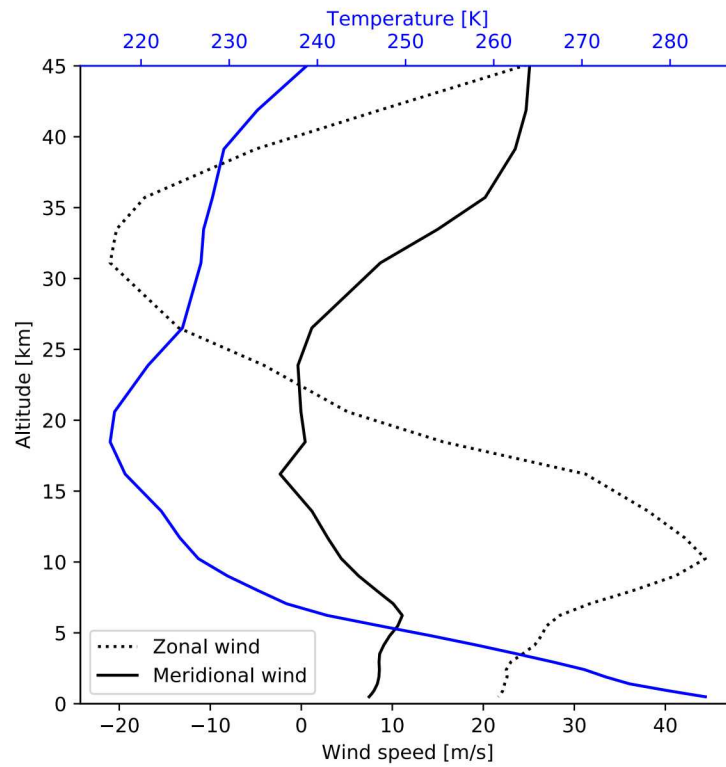


Figure S2 : Atmospheric profile above the Tasman Sea used for ray path propagation modeling, derived from the Global Forecast System model from NOMADS at 12z on 17 May 2016.

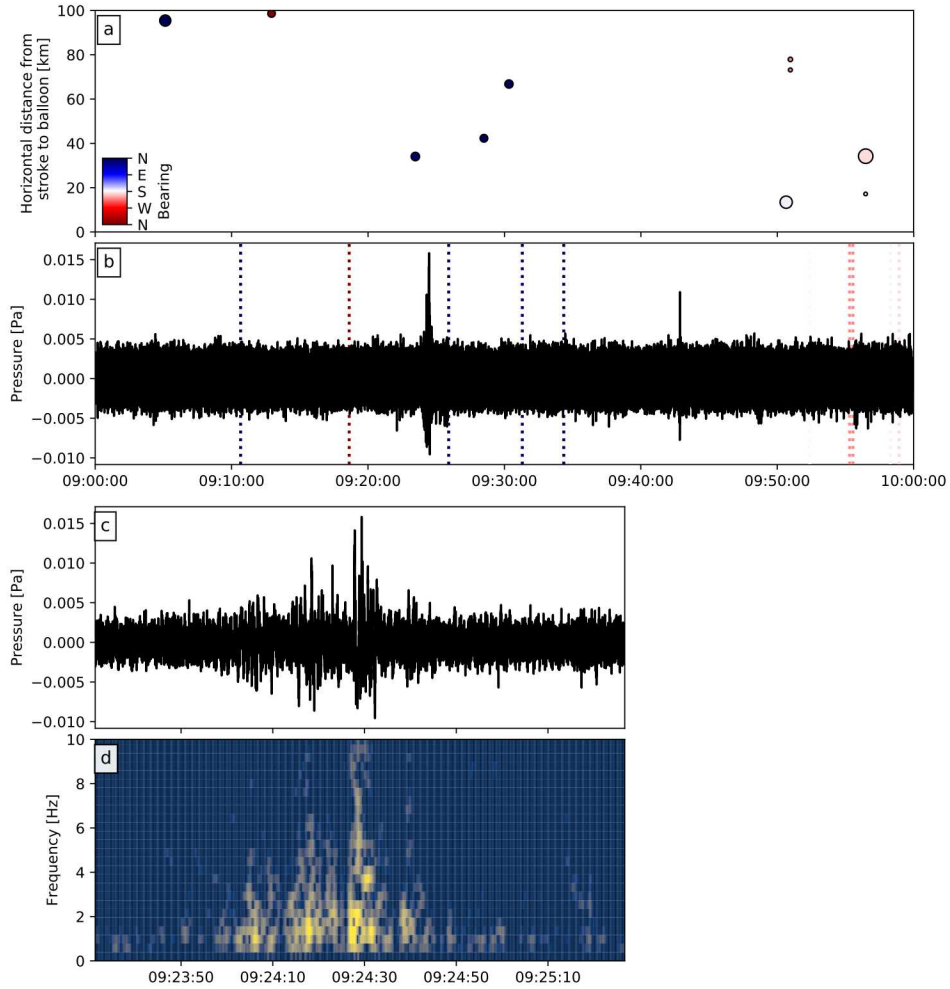


Figure S3 : (a) Horizontal distances for each stroke within 100 km of the balloon from 0900 to 1000 UTC on 17 May 2016. each stroke is sized by the stroke energy, and colored by the bearing from the balloon to the stroke location. (b) High-pass filtered acoustic waveform as recorded at the ULDB over the same time period. Vertical dotted lines indicate calculated time of arrivals for strokes in panel (a), colored by the bearing. (c) Possible example acoustic signal from lightning stroke and its frequency spectrogram (d).

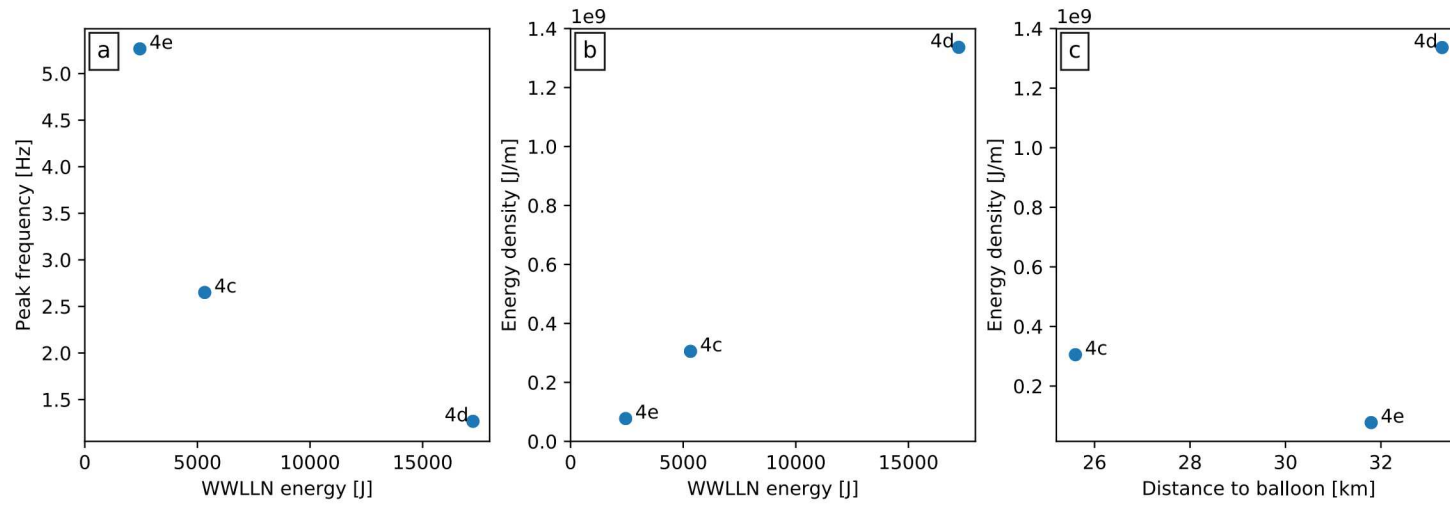


Figure S4 : Results of calculating the energy density of the lightning strokes from the acoustic waves plotted in Fig 4c, d, and e. (a) Calculated peak frequency of each wave plotted against the detected energy by the WWLLN for matching lightning strokes. (b) Calculating energy densities for each waveform versus the detected energy by the WWLLN for matching lightning strokes. (c) Calculated energy densities for each waveform versus the distance between the matching lightning stroke and balloon.

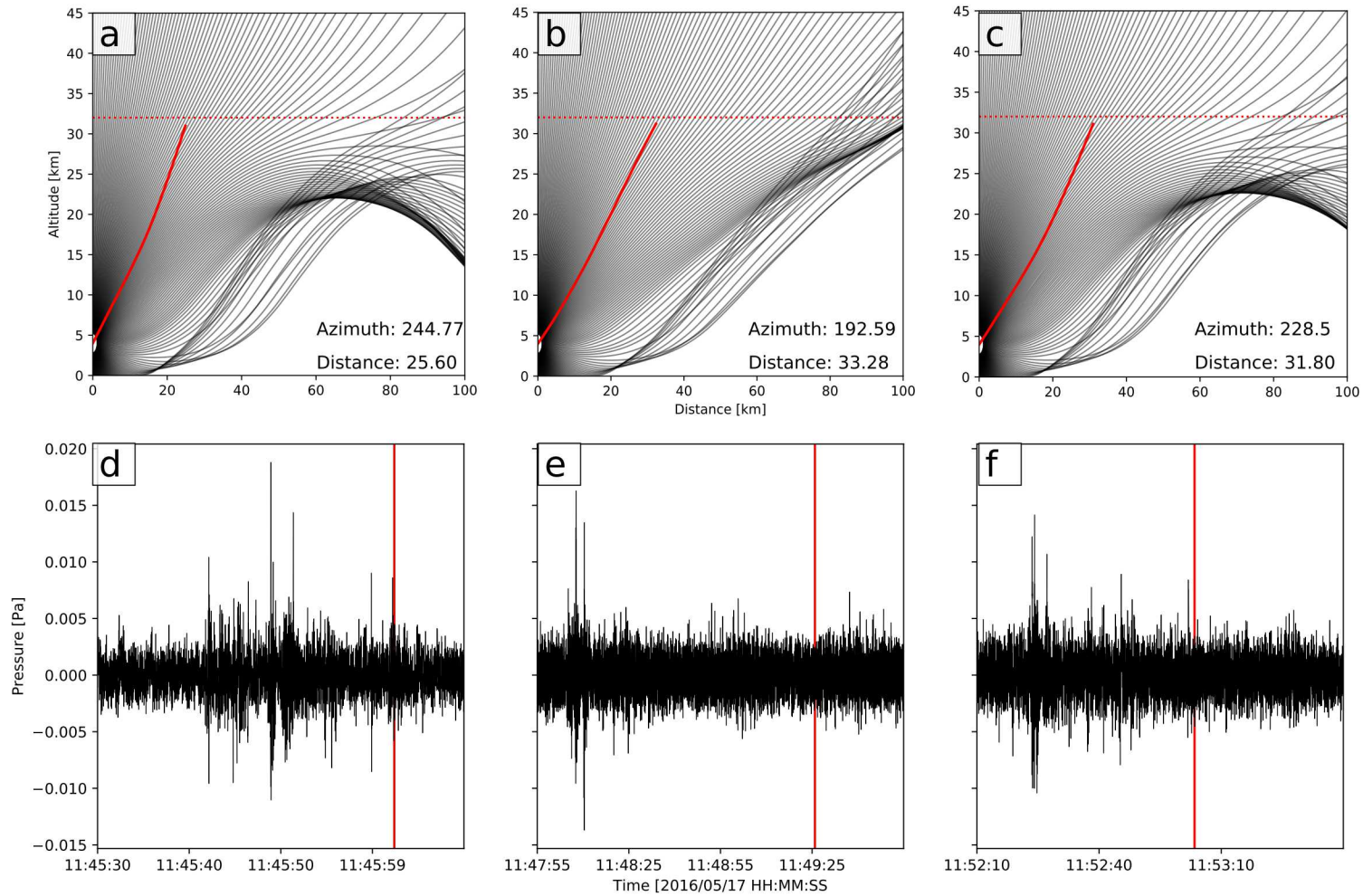


Figure S5 : Ray path propagation results (black lines) and eigenray solutions between a source and the receiver at the ULDB, for signals detailed in Figure 4c, d, and e (a, b and c, respectively). Azimuth and horizontal distances between source and receivers are detailed within each panel. Panels d, e, and f plot the waveforms with calculated arrival times from the eigenray solutions (red lines).

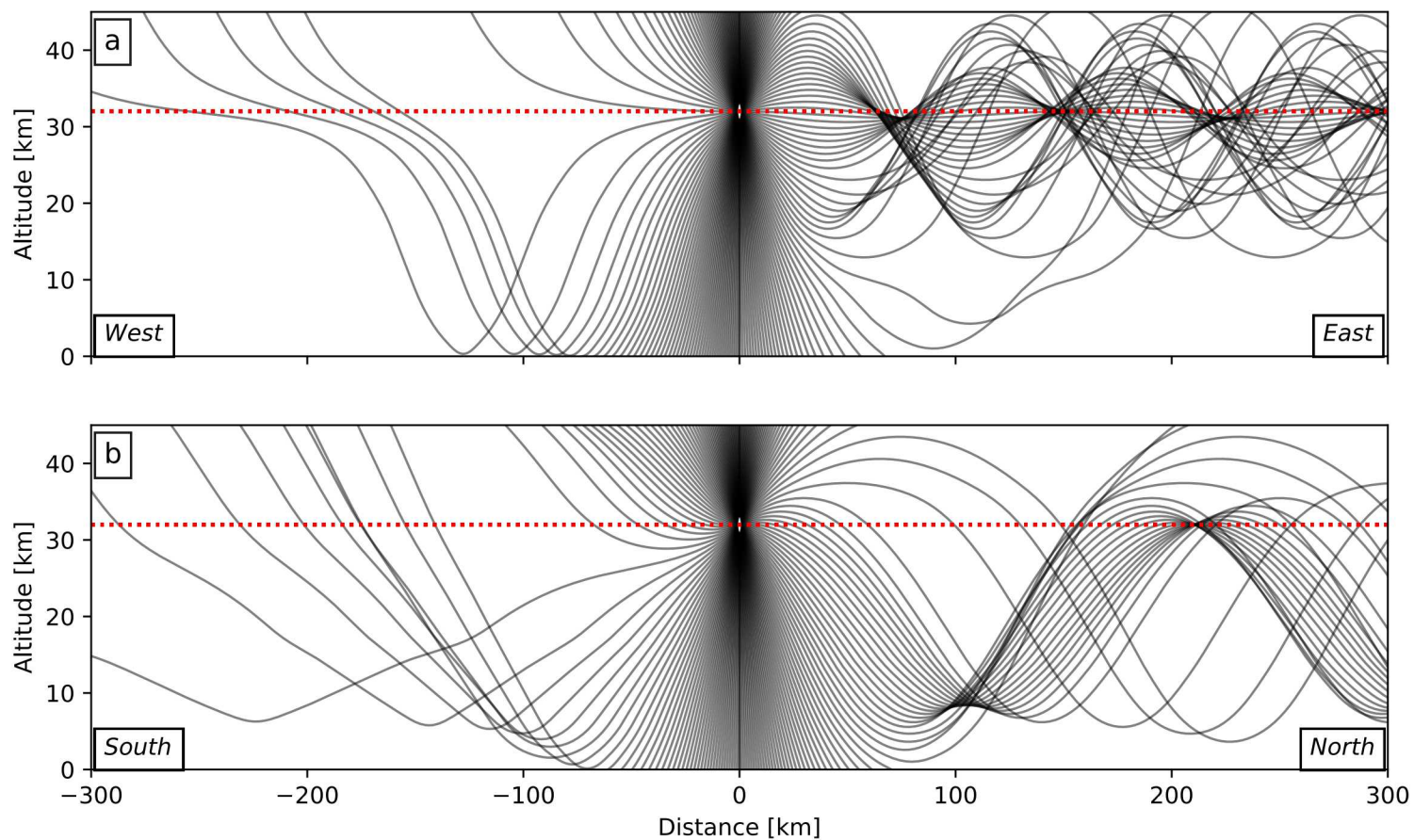


Figure S6 : Ray-tracing propagation for an acoustic source at 32 km height along East-West (a) and North-South (b) profiles using realistic atmospheric conditions derived from the Global Forecast System. The red dotted line at 32 km indicates the average height of the ULDB balloon on 17 May 2016.