

# Simulations of planar non-thermal plasma assisted ignition at atmospheric pressure



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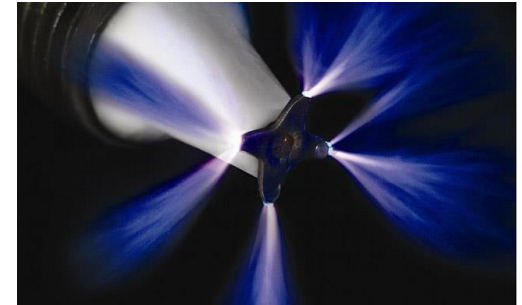


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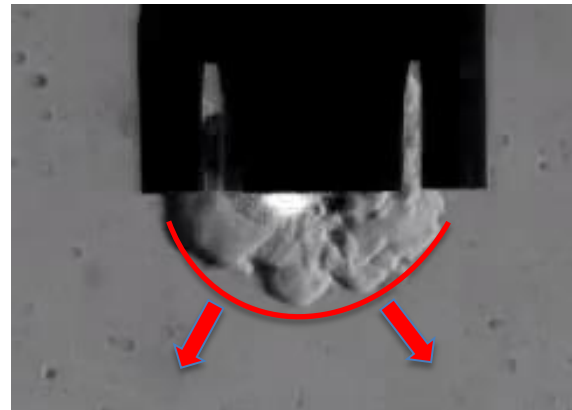


# Electric field effects: background

- Flame zones are essentially weakly ionized plasmas, generate weak self-electric fields due to local charge separation
- As such, externally applied electric fields have the ability to augment flames “non-intrusively”
- Applied electric fields have been observed to affect reacting flows by:
  - enhancing burning velocities
  - extending flammability limits (lean)
  - supporting ignition



***Corona Discharge  
Ignition System***



***Pulsed plasma  
breakdown  
ignition device***

[1] D. I. Pineda, et al., "Nanosecond Pulsed Discharge in a Lean Methane-Air Mixture," in *Laser Ignition Conference*, OSA Technical Digest (online) (Optical Society of America, 2015), paper T5A.2.



# Conceptual motivation

- thermal effects: energy transfer from accelerated electrons to neutral molecules results in bulk fluid heating
- chemical effects: energetic electron collisions with abundant neutral molecules (e.g.  $N_2$ ,  $O_2$  in air) generate excited states, ions, with increased reactivity, cause decomposition
- electro-hydrodynamic force: acceleration of ions and electrons by electric fields, resulting in increased momentum of the bulk fluid



# Electron fluid

- Strong electric fields acting on charged species can cause bulk fluid approximations to break down
- Must derive individual fluid equations for possible non-equilibrium components (electron fluid), e.g. using moments of the Boltzmann kinetic equation

$$\frac{\partial m_e n_e}{\partial t} + \nabla \cdot (m_e n_e u_e) = S_e \quad \text{electron fluid mass conservation}$$

$$\frac{\partial m_e n_e u_e}{\partial t} + \nabla \cdot (m_e n_e u_e u_e) = -\nabla p_e - \underbrace{e n_e E}_{\text{species body force}} - \underbrace{m_e n_e \nu_{en} (u_e - u_n)}_{\text{collisional drag force (Lorentz gas model)}} \quad \text{electron fluid momentum conservation}$$

*Zero-inertia approximation for electron fluid momentum*

$$n_e u_e = -\frac{e}{m_e \nu_{en}} n_e E - \frac{1}{m_e \nu_{en}} \nabla (n_e k_B T_e)$$



# Electron fluid

**This defines an electron fluid transport coefficient:**

- electron mobility

$$\mu_e = \frac{e}{m_e \nu_{en}}$$

- electron momentum eq. is now algebraic

$$n_e u_e = -\mu_e n_e E - D_e \nabla n_e$$

*“drift”*

*“diffusion”*

- electron flux is closed with specification of the collision frequency, which can be estimated from kinetic theory and knowledge of the energy/speed distribution function



# Governing equations

## 1 Electron fluid

$$\frac{\partial n_e}{\partial t} + \nabla \cdot (n_e \mathbf{u}_e) = S_e \quad \mathbf{u}_e = -\frac{e}{m_e \nu_e} \mathbf{E} - \frac{1}{n_e m_e \nu_e} \nabla (n_e k_B T_e)$$

## 2 Bulk fluid

$$\frac{\partial \rho Y_i}{\partial t} + \nabla \cdot (\rho Y_i [\mathbf{u} + \mathbf{V}_i]) = m_i S_i$$

Electron Drag

$$\frac{\partial \rho \mathbf{u}}{\partial t} + \nabla \cdot (\rho \mathbf{u} \mathbf{u}) = -\nabla p + \nabla \cdot \boldsymbol{\tau} + \rho_q \mathbf{E} - \sum_{s \neq e} n_s m_{s,e} \nu_{s,e} (\mathbf{u}_s - \mathbf{u}_e)$$

$$\frac{\partial \rho e_T}{\partial t} + \nabla \cdot (\rho \mathbf{u} e_T) = -\nabla \cdot \mathbf{Q} + \nabla \cdot (\underline{\underline{\sigma}} \cdot \mathbf{u}) - \sum_{s \neq e} 3n_e \frac{m_s^-}{m_e} \nu_{s,e} k_B (T_s - T_e) + (\epsilon_{cons} S_{e,cons} - \epsilon_{prod} S_{e,prod})$$

Electron elastic energy  
transfer model

Electron inelastic  
energy exchanges

## Gauss' Law

$$\epsilon \nabla \cdot \mathbf{E} = \sum_{s \neq e} \left( \frac{\rho q_s Y_s}{m_s} \right) - e n_e$$

s: species (electron or bulk component)

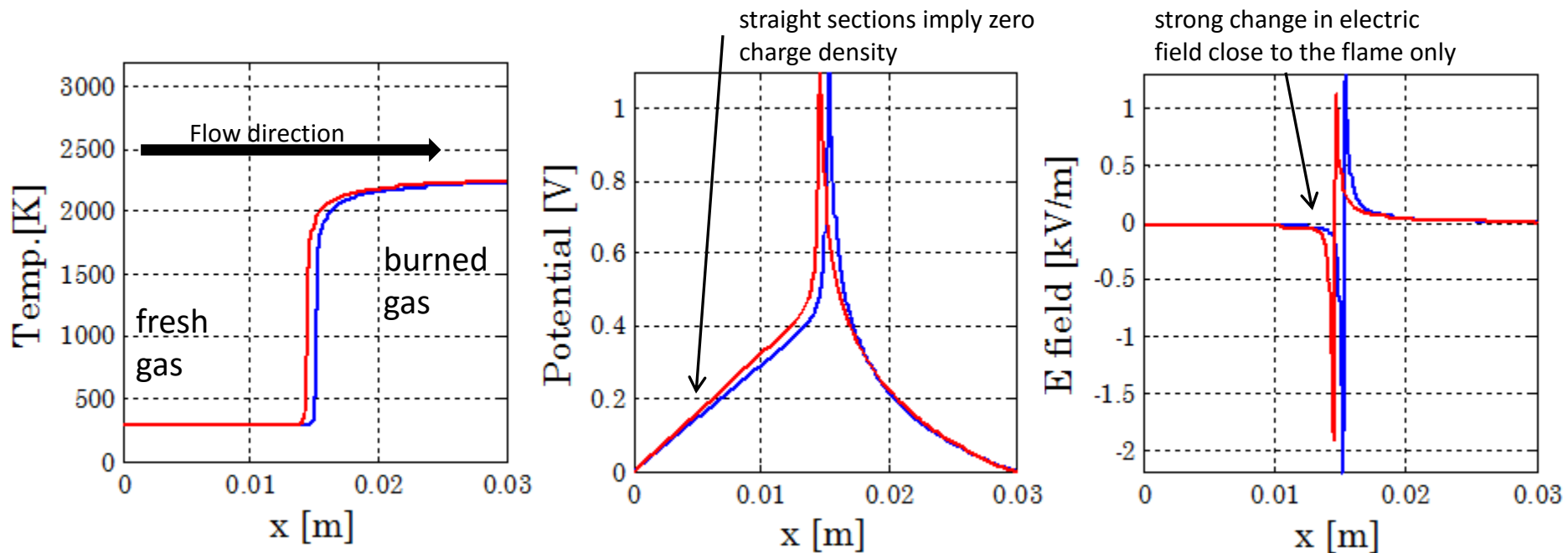
e: electron species

i: bulk fluid components



# Example: Laminar premixed methane-air flame

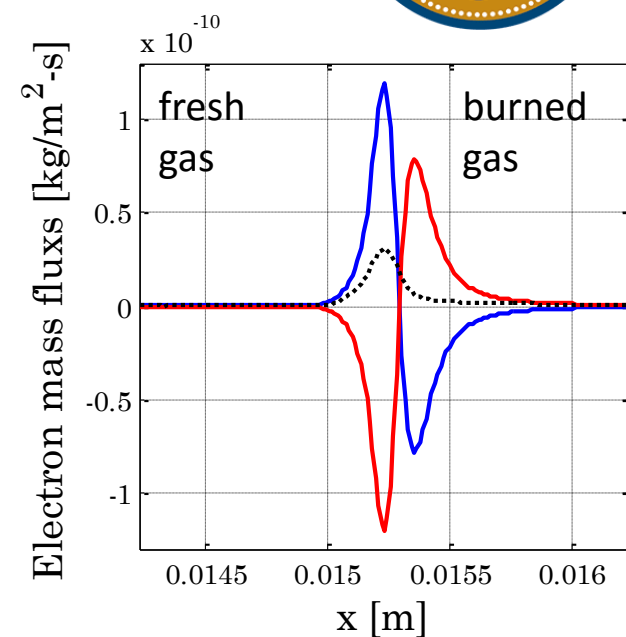
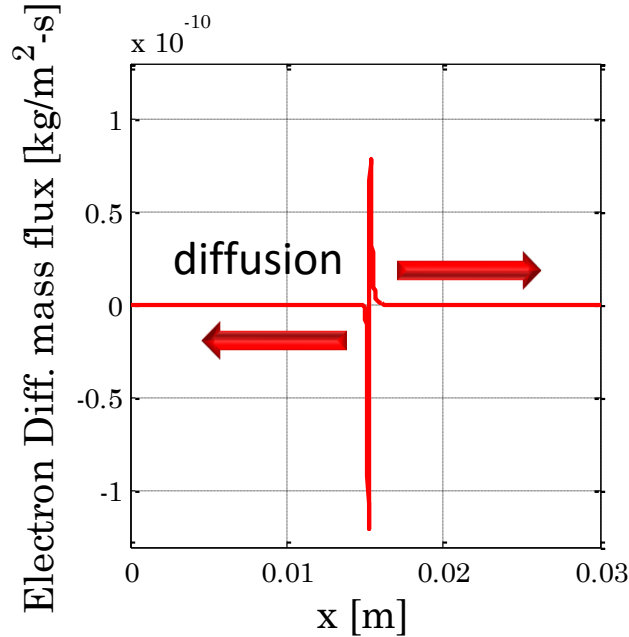
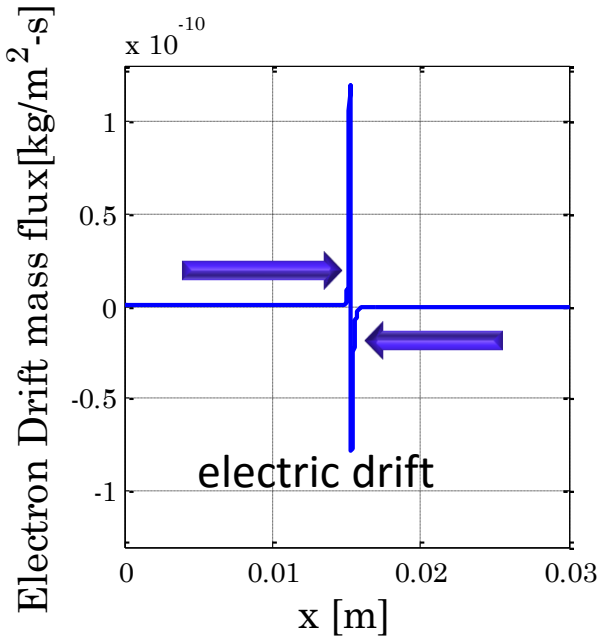
- Stoichiometric mixture, 300K fresh gas at 1atm, 0V boundaries on 3cm domain, left end inlet, right end open



Blue curves are **transient** results (S3D)  
Red curves are **steady-state** results  
using an equilibrium solver (PREMIX)



# Laminar premixed flames



- Self-field dynamic equilibrium: mass fluxes are in balance, a pool of electrons sits at the flame

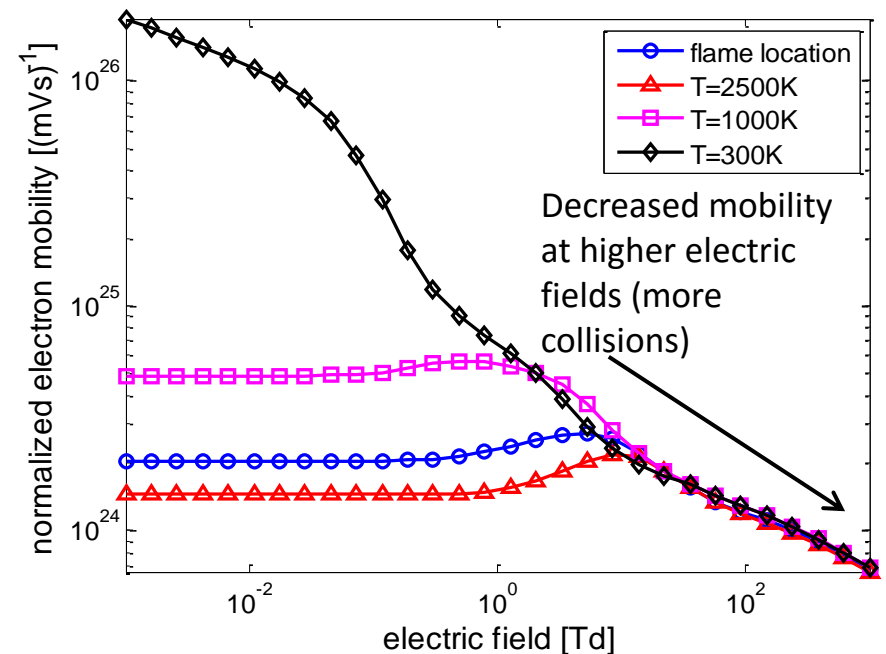
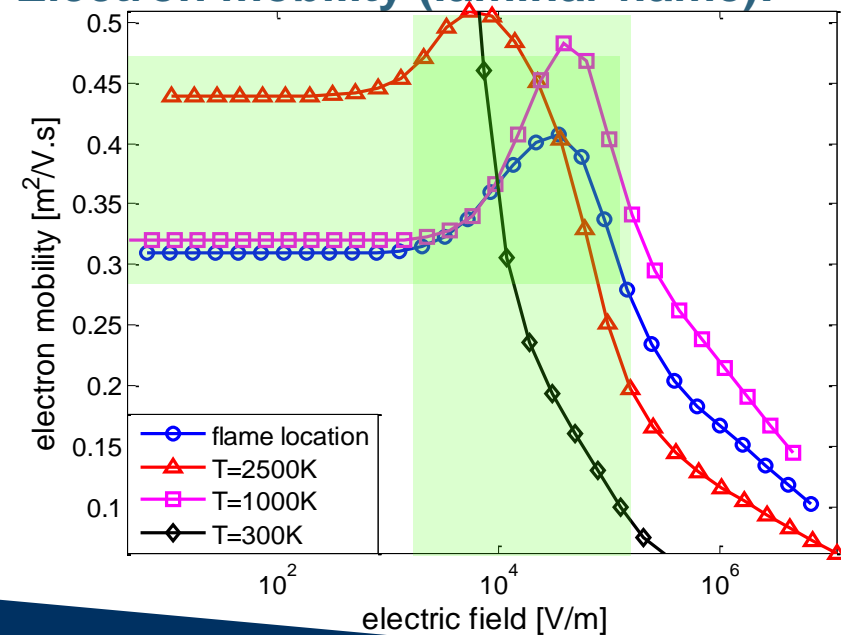


# Electron properties



- Electron Boltzmann eq. solver: BOLSIG+  
<http://www.bolsig.laplace.univ-tlse.fr/>
- Two term approximation for phase space electron density
- Problem is fully posed by specification of the neutral species densities (chemistry), gas temperature (energy loss), and electric field

## Electron mobility (laminar flame):





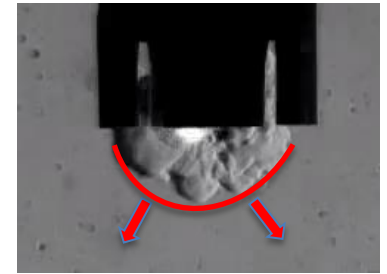
# Numerics

- Reacting Navier-Stokes equations: fully compressible transient reacting flow simulations in 1D, conservative finite difference discretization (S3D, MPI using 32 threads)
- 4<sup>th</sup> order low-storage Runge-Kutta explicit time integration, time step at typically 1ns (thermal plasma), 6<sup>th</sup> order centered spatial differencing scheme, 8 $\mu$ m resolution
- 45 species methane-air chemical model with two-temperature parameterization, non-thermal electron impact reactions using CHEMKIN. Factor of 2 speed up by pre-computing rate constants.
- Poisson equation for Gauss' law solved using geometric multigrid iterative methods
- transport properties and electron energies pre-computed using the BOLSIG+ solver, interpolated at runtime



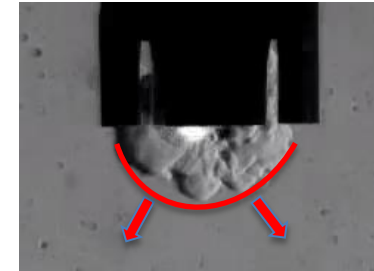
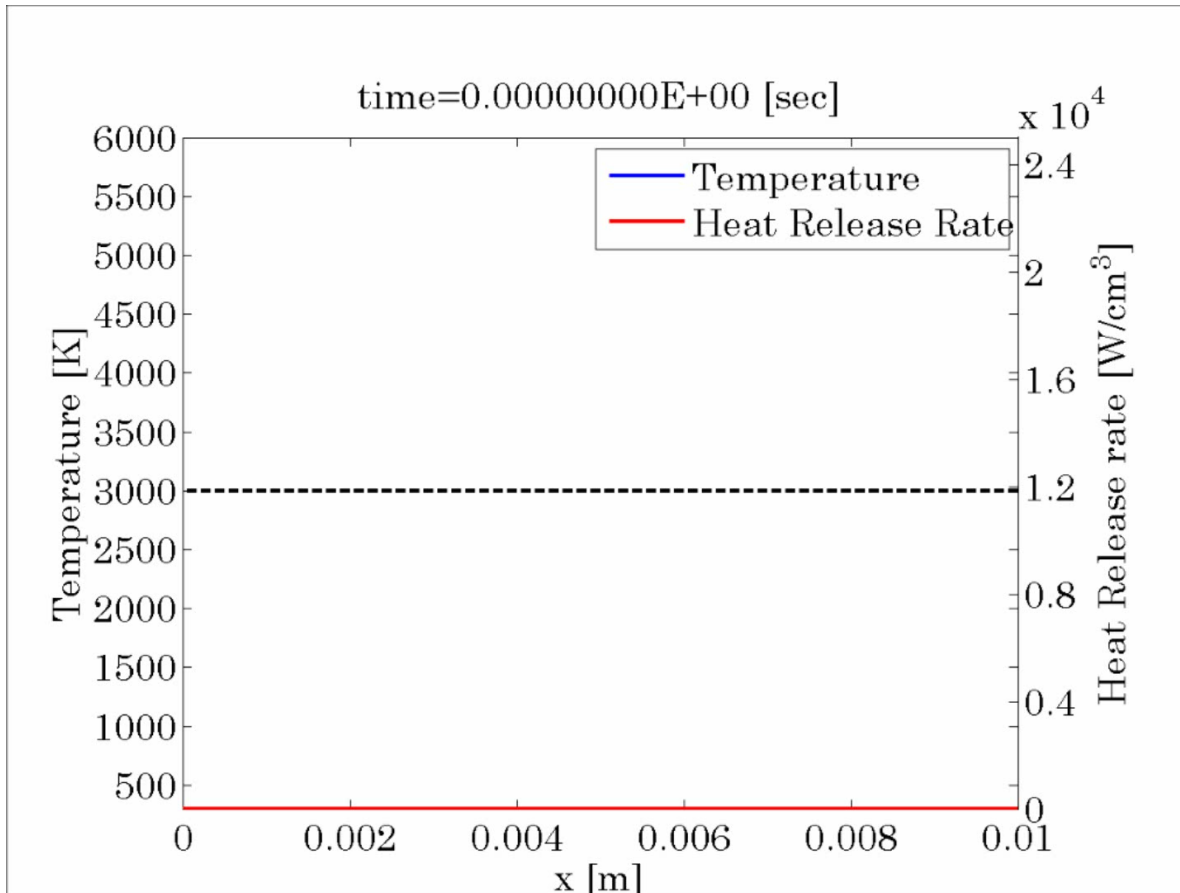
# Plasma assisted ignition

- Standard ignition devices (spark) use discharge arcs as energy sources for ignition. But increasing spark energy can result in decreased device life due to ablation
- Instead propose to use applied electric fields to support weak ignition kernels
- Emulate pulsed plasma assisted ignition using a single square wave pulse of 10kV magnitude and ~ns duration applied to a developing ignition kernel
- Interested in atmospheric pressure for device and experimental conditions



***nanosecond pulsed plasma  
ignition device***

# Planar ignition: CH<sub>4</sub>-Air (unsupported)

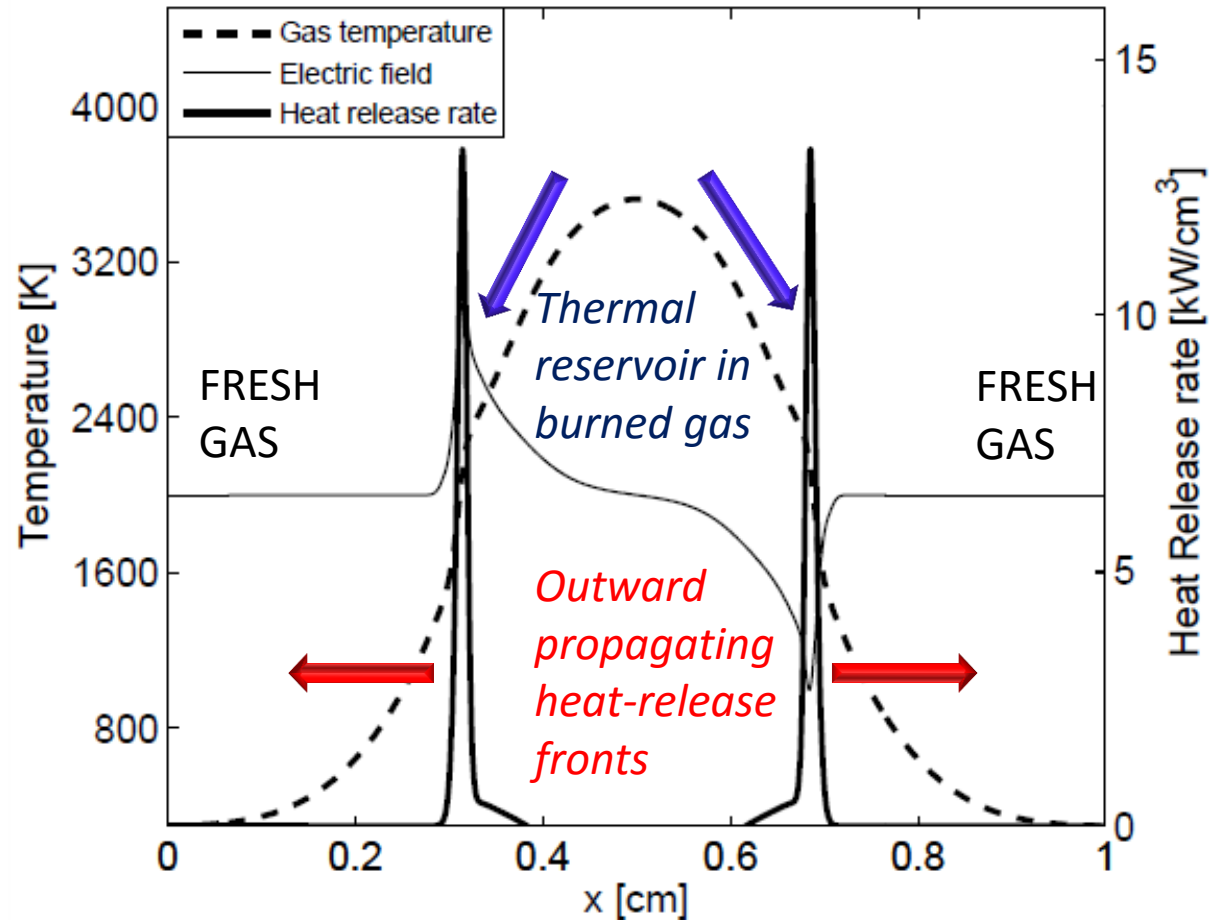


- 1cm domain
- fluid ignited by Gaussian shaped ignition kernel active for 80 $\mu$ s, 85mJ total energy
- boundaries are non-obstructive to fluid (mesh electrodes)
- ends grounded



# Planar ignition: CH<sub>4</sub>-Air

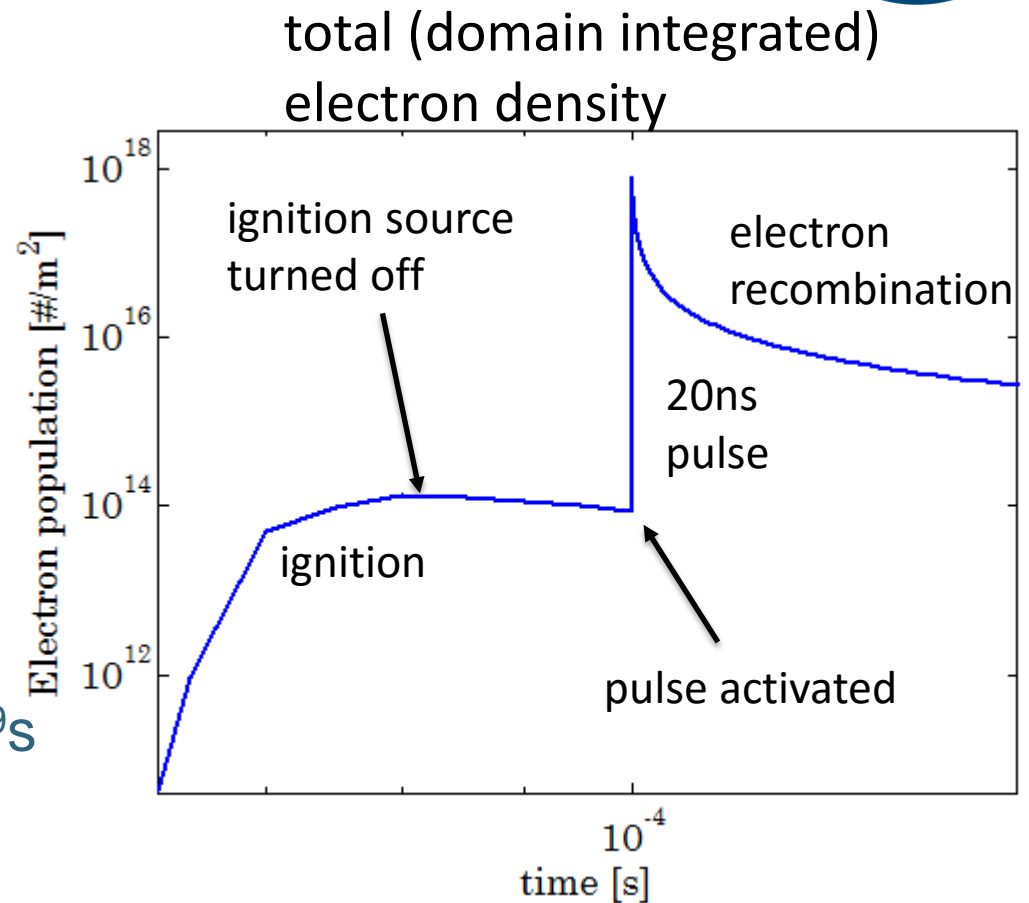
- Flame fronts propagate outwards rapidly
- Begin to decelerate as they depart from the hottest regions of burned gas and back-support diminishes
- What happens when a strong electric field is applied? (ns pulse)





# Planar ignition support – DC pulse

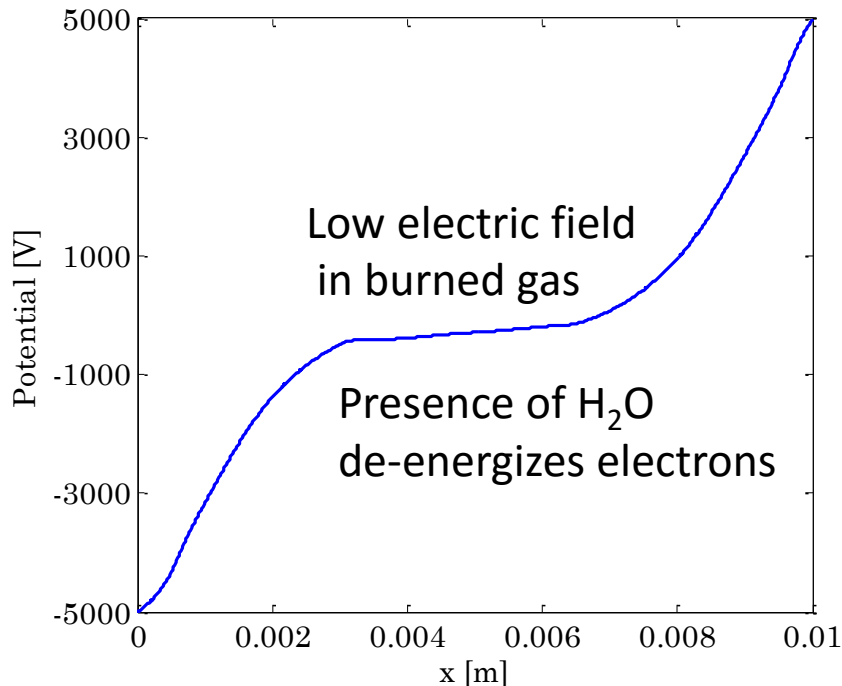
- DC pulses, 20ns-35ns duration, 10kV (-5kV left, +5kV right)
- 300K mixture, 1atm,  $\phi=0.5$
- Electrons rapidly become non-thermal due to Joule heating, activate impact chemistry
- Simulation  $\Delta t = 10^{-13}$ s to resolve electron impact kinetic rates during pulse, relax dynamically to  $\Delta t = 10^{-9}$ s post pulse



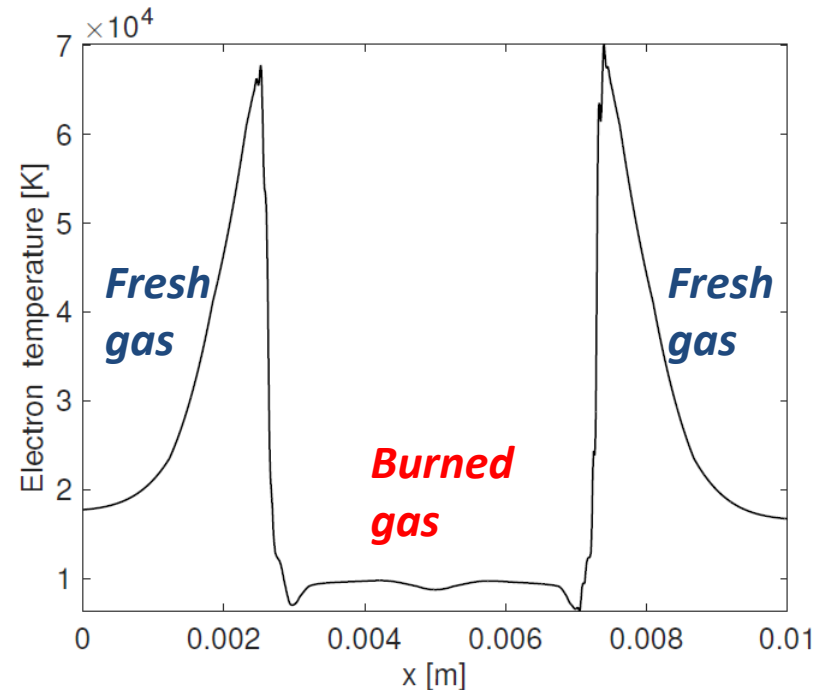


# Instantaneous ignition profiles during pulse

instantaneous potential field



instantaneous electron temperature

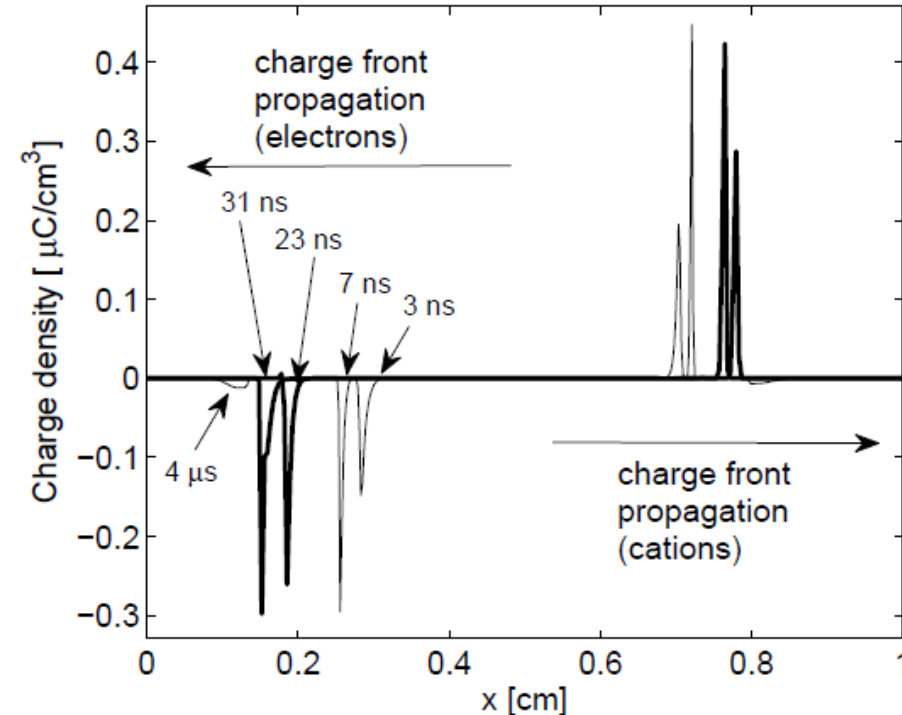


- electron temperature  $\sim 20,000$ - $70,000$  K in fresh gas
- electron temperature is conjugate to electric field strength



# Dynamics of charge fronts

- On pulse activation, charge originating in the flame zones is ejected into the fresh gases
- As the charge front advances, ionization processes increase the charge density
- Post pulse, the charge fronts collapse due to recombination



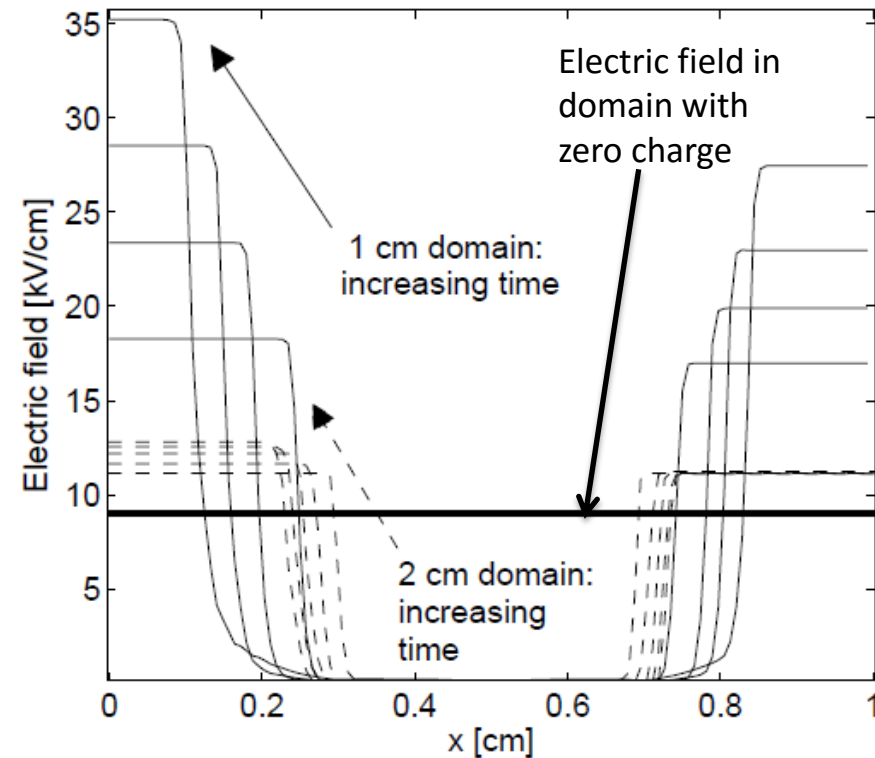
Charge front propagation for 35ns pulse. Electrons are swept to the left, cations are swept to the right



# Electric field compression



- Charge fronts are very pronounced, changes in electric field are restricted to their vicinity
- this results in a “floating electrode” effect as the potential transitions from linear to a constant value
- the advancing charge fronts steepen the slope of the potential, increasing the electric field strength



Increase in electric field strength  
between end electrodes and charge  
fronts

# Production and consumption



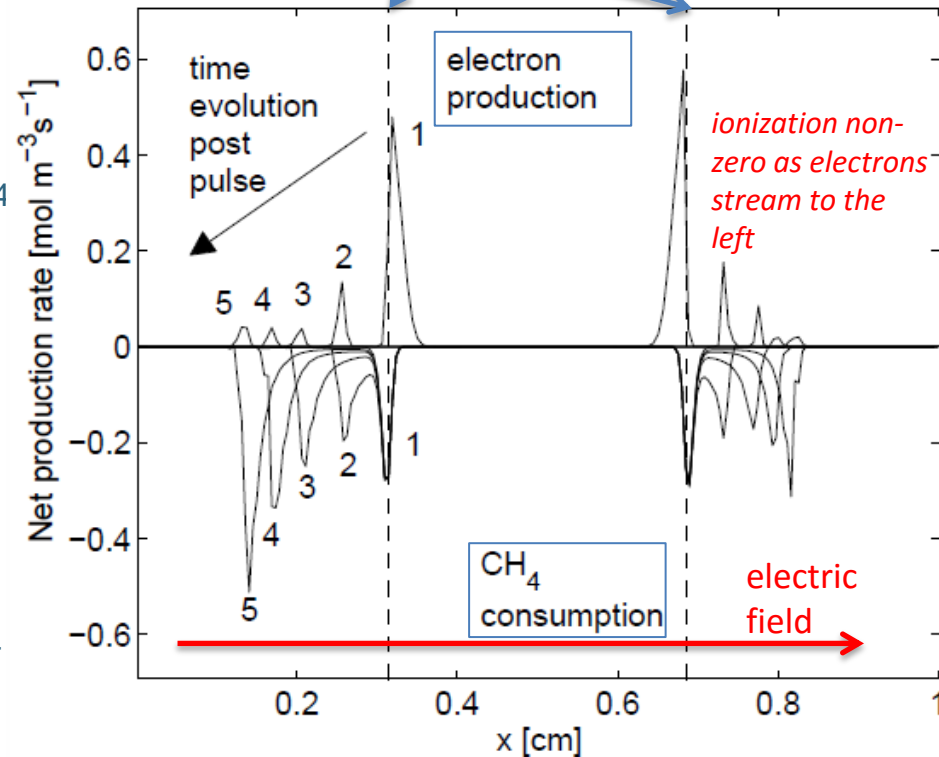
Electrons streaming into the fresh gas result in  $\text{CH}_4$  decomposition and further electron production

1: electron production is high, a large packet of thermal electrons in the reaction zone.  $\text{CH}_4$  consumption increases

2: electrons are ejected from reaction zone cause impact ionization (producing more electrons),  $\text{CH}_4$  impact decomposition occurs

3-5: ionization processes not increasing,  $\text{CH}_4$  impact decomposition and de-excitation of electronically excited  $\text{N}_2$  (also increasing electron temperature)

Location of “flame zones”, frozen during the ns pulse

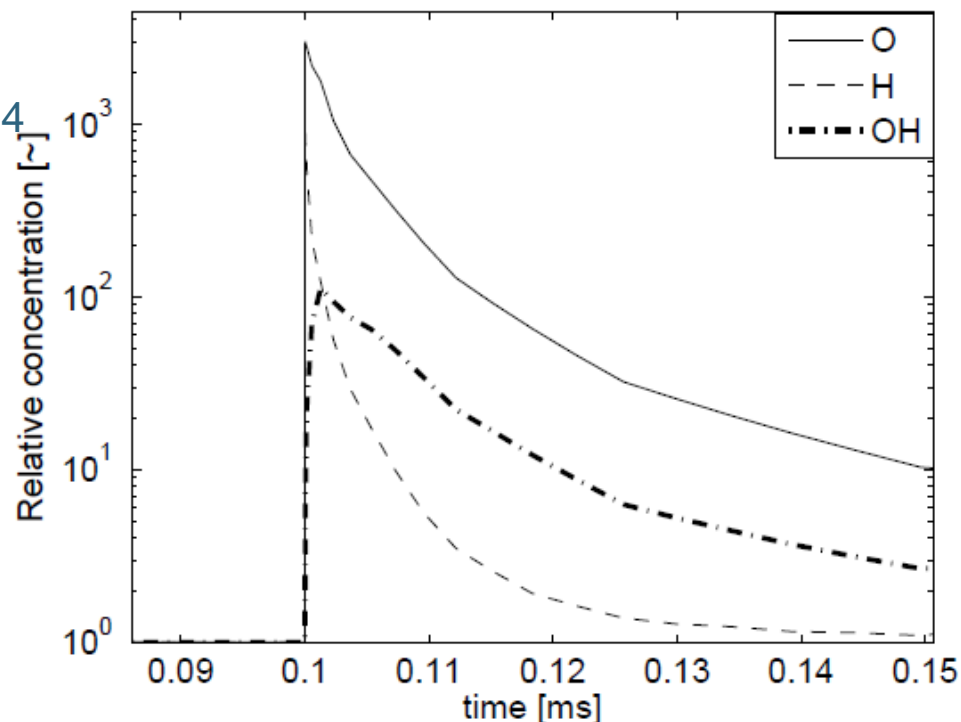


Production rate of electrons (top half plane) and consumption of  $\text{CH}_4$  (bottom half plane) during the pulse at five time instants corresponding to 1, 4, 8, 12, and 16 ns labeled “1” to “5”.



# Radical production

- Radical populations build up rapidly during the pulse as  $\text{CH}_4$  and  $\text{O}_2$  are fragmented by electron impact, by factors of 100 to 1000
- OH is not formed during the pulse, but begins to build up post pulse due to radical recombination

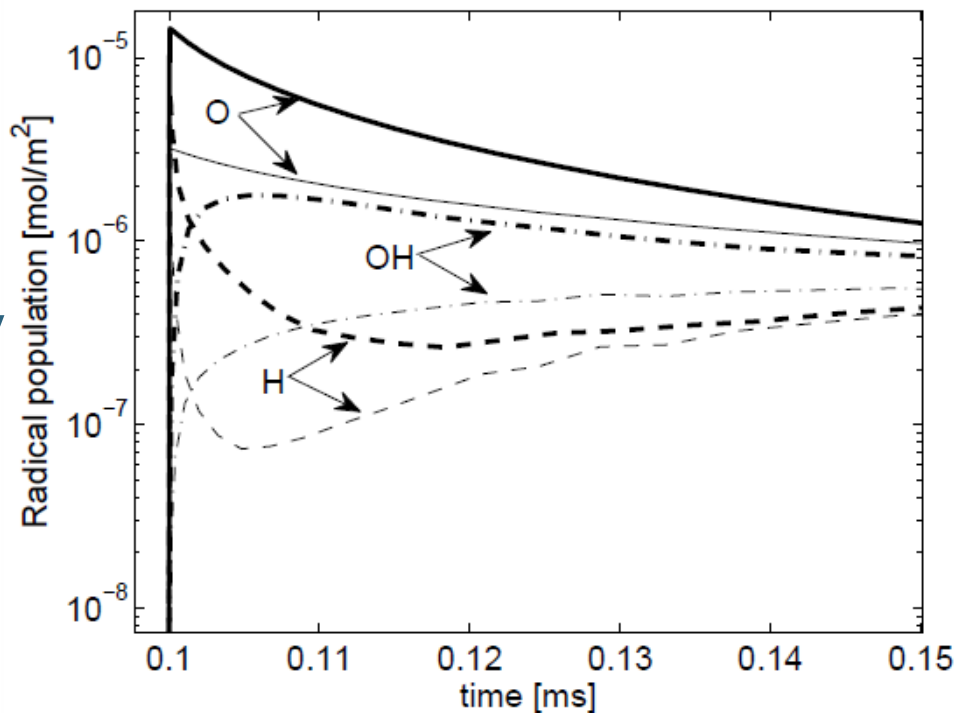


Evolution of domain integrated major radical species concentrations for a 35 ns pulse normalized by the evolution in the unsupported case



# Effect of pulse length

- Longer pulse widths increase radical populations. 25% decrease in pulse width (35ns to 26ns) reduces population by factor of 5

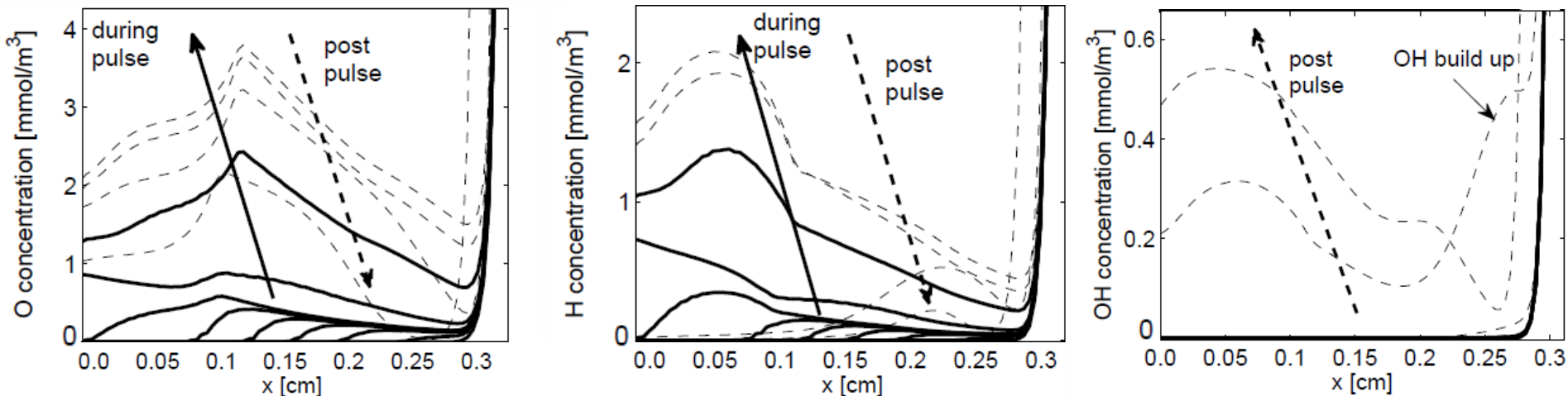


Evolution of total populations of major radical species for a 35 ns (solid lines) and 26 ns (dashed lines) pulse.



# Fresh gas radicals

Fresh gas radical spatial profiles (near left electrode), profiles are separated in time by 4 ns during the pulse, with the post pulse profiles taken at 1 ns, 56 ns, 1.7  $\mu$ s, and 8.2  $\mu$ s after the pulse has ended.

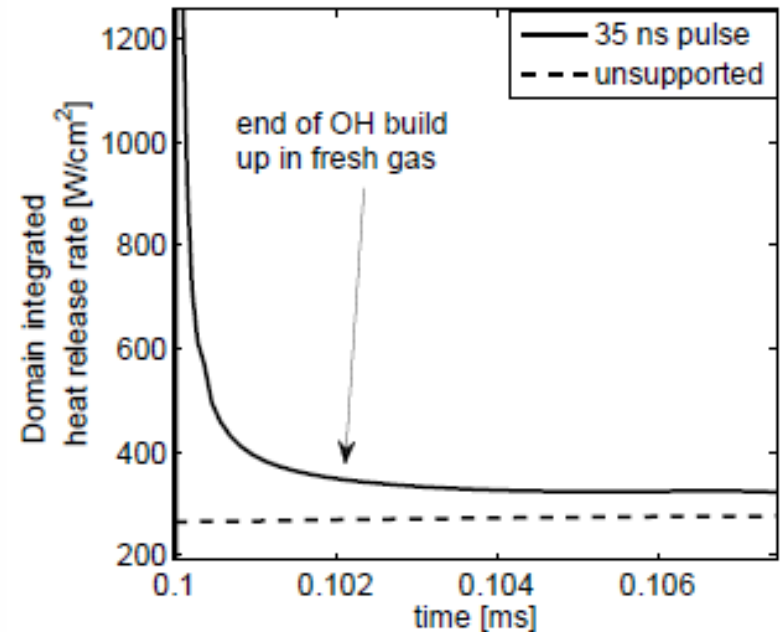


- populations slowly build up (solid lines) before reaching maxima in the vicinity of the boundary electrode (at electron temperature peak)
- OH build up post pulse in the vicinity of the flame zones



# Heat release rate

- Post-pulse, the heat-release rate decays rapidly but begins to plateau as the OH build up reaches its maximum
- This sustained heat release accelerates the flames forward into the fresh gases
- After a few ms the enhancement effect dissipates

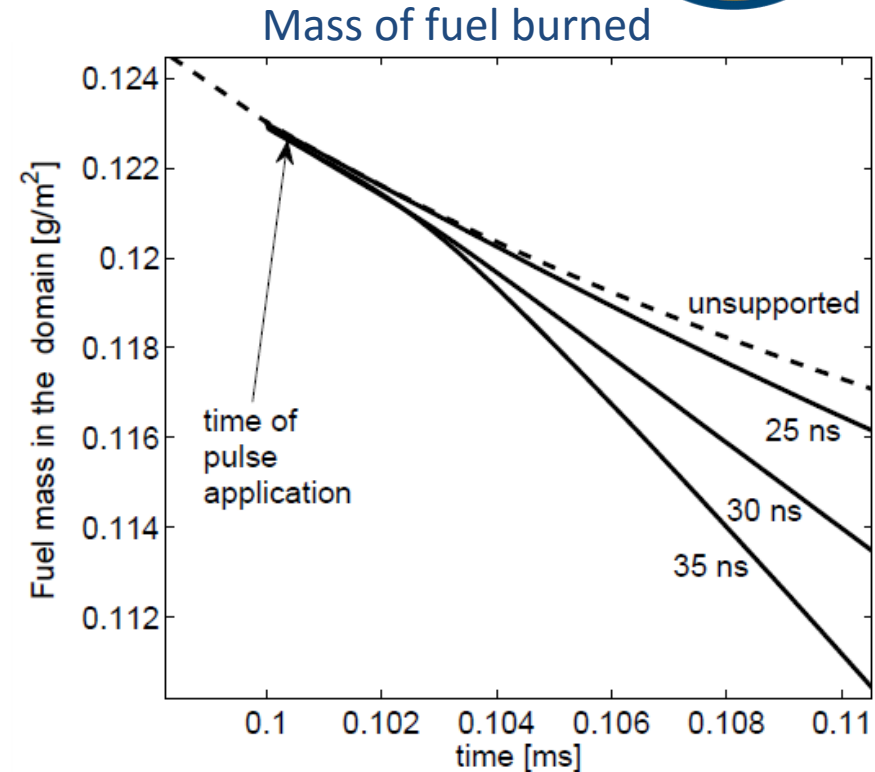
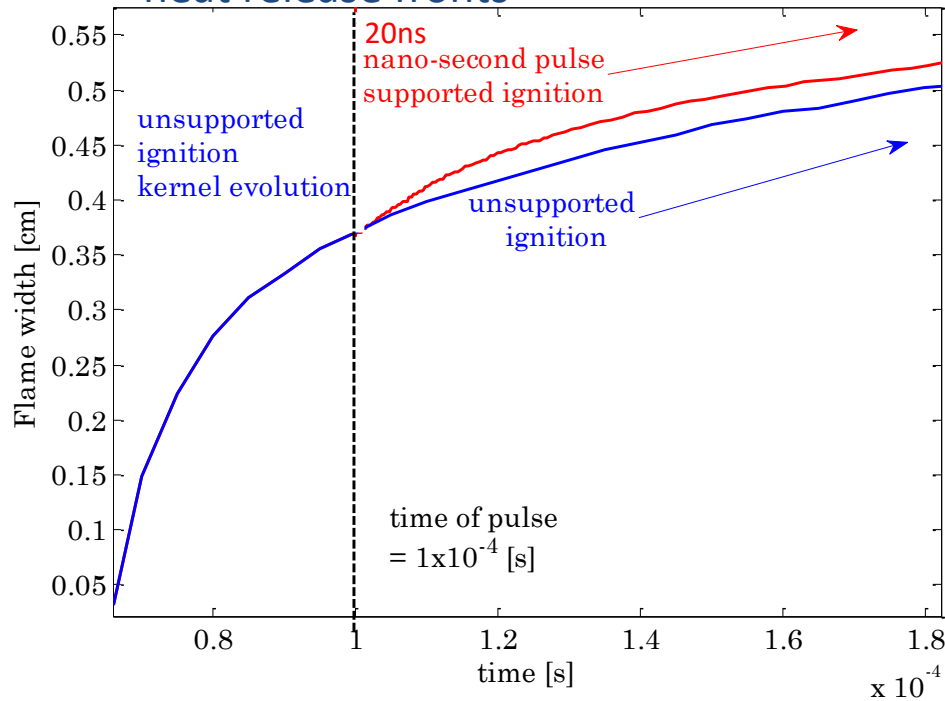


domain integrated heat release rate in the post-pulse period for a 35 ns pulse (solid line) and the unsupported case (dashed line).

# Results: global effect on ignition enhancement



flame kernel “size” – separation of heat release fronts



- Supported ignition kernel grows spatially and consumes fuel more rapidly



# Summary

- At equivalent times, up 4% increase in flame kernel size, 6% increase in mass of fuel burned as a result of DC pulse
- With respect to kernel development time, up 12% reduction in time to achieve equivalent flame kernel sizes, 33% reduction in time to achieve equivalent fuel burn
- At longer times, enhancement effect appears to diminish as the supporting effect of burned gas thermal diffusion on the flames decreases (competing effects)
- electron processes in the hot burned-gases are essentially insignificant, due to low electric field and deactivation of energetic electrons by collisions with  $\text{H}_2\text{O}$





Thanks!

Questions?